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Antenna Measurement

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1. Introduction

The antenna is an important element of radiocommunication, remote sensing and radiolocalisation systems. The measurement of the antenna radiation pattern characteristics allows one to verify the conformity of the antenna. The simplest measurement method consists in the direct far-field measurement. For large antennas, the necessary measurement distance raises a problem which was resolved by the introduction of compact ranges and near-field techniques. Compact range consists in a focusing system, as a reflector, which can create a plane wave at short distance. The principle of near-field techniques is to measure the field radiated by an antenna at a short distance on a given surface surrounding the antenna, then to calculate the far-field starting from the measured near-field. These last techniques also make it possible to have an excellent precision as required by the satellite antennas for example. The near-field techniques also make it possible to carry out the diagnosis of the antennas, i.e. to find defects on the antenna. The duration of the measurement of the large antennas, which claims a large number of measurement points, was reduced considerably by the use of rapid near-field assessment system, for which the mechanical displacement of the probe was replaced by the electronic scanning of a probes array.

2. Direct far-field measurement

2.1 Antenna pattern measurement

There are four different regions for the electromagnetic field radiated by an antenna: three near-field regions and one far-field region. The nearest region is the reactive field region which extends until a distance of one wavelength λ from the antenna surface. The second near-field region is the Rayleigh region which extends from the reactive region until a distance D²/(2 λ) from the antenna surface, a relation in which D is the tested antenna diameter. The third near-field region is the Fresnel region which extends from D²/(2 λ) until 2D²/ λ from the antenna surface. The last region is the far-field Fraunhofer region which starts at a distance of 2D²/ λ . The space variations of the electromagnetic field differ in these four regions. It is thus necessary to be at a distance sufficient (>2D²/ λ) to carry out direct far-field measurements.



Fig. 1. Different field regions for a large antenna

Direct far-field measurement can be realized either in indoor or outdoor range (Kummer & Gillepsie, 1978). Indoor range consists in an anechoic chamber with one source antenna and the tested antenna placed on a positioner. This positioner allows one to vary the tested antenna attitude with respect to the direction of the wave incidence of the wave for far-field pattern measurement. Outdoor range consists in a tower bearing a source antenna and the tested antenna placed on a positioner. The distance between the two antennas can be larger for outdoor range, i.e. the capacity of outdoor range in terms of tested antenna dimensions is higher. On the other hand the outdoor range is sensitive to parasitic signals and ground reflections.



2.2 Measurement of the characteristics of the pattern

Gain measurement

Absolute gain measurements use power budget between antennas. The Friis transmission formula gives a relation between the gains G_A and G_B of two antennas A and B:

$$P_r = P_i \ G_A \ G_B \ (\lambda/4\pi r)^2 \tag{1}$$

 P_r is the power received at a matched load connected to the receiving antenna, P_i the power accepted by the transmitting antenna, λ is the wavelength and r the distance between the two antennas. The use of this formula requires that the two antennas A and B are

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polarization matched and that the separation distance between the antennas corresponds to far-field conditions.



The first method uses the tested antenna and a reference antenna with a known gain. The measurements of P_r , P_i , r, and the knowledge of λ give the value of the tested antenna gain by means of the Friis formula. The second method used two identical tested antennas with the same gain. The Friis formula gives the value of the tested gain antenna starting from the same measurements as the preceding method. The third method used three antennas A, B and C, generally the tested antenna and two other antennas. Three different power budgets are carried out, with the three possible antennas pairs leading to the three following equations:

$$G_A G_B = (P_r/P_i)_{AB} (4\pi r/\lambda)^2$$
⁽²⁾

$$G_B G_C = (P_r / P_i)_{BC} (4\pi r / \lambda)^2$$
(3)

$$G_{\rm C} G_{\rm A} = (P_{\rm r} / P_{\rm i})_{\rm CA} (4\pi r / \lambda)^2$$
(4)

These three formulas provide the determination of the three different gains G_A , G_B and G_C .

Directivity measurement

There are two methods which make it possible to measure the directivity Dir of an antenna. The first method is based on the definition of the directivity. The knowledge of the relative far-field radiated by the antenna in all the directions is sufficient to know the directivity of this antenna. This method is called the pattern integration method because it uses the integration of the radiated power density (dP_e/dS) on all the directions to calculate the total power radiated by the antenna P_e . The power density is related to the electric field E at a distance r by the relation:

$$(dP_e/dS) = E^2/Z_0$$
(5)

 $Z_0 = 120\pi \Omega$ is the free space impedance of wave.

$$P_{e} = \int_{\varphi=0}^{2\pi} \int_{\theta=0}^{\pi} (dP_{e}/dS) r \sin\theta \, d\phi \, d\theta$$
(6)

Dir =
$$(dP_e/dS) / [P_e/(4\pi r^2)]$$
 (7)

 $\langle \alpha \rangle$

The second method uses the relation between the directivity Dir and the gain G of a given antenna:

$$Dir = G/\eta \tag{8}$$

in which η is the efficiency of the antenna. The measurements of the gain and the efficiency of the antenna result in the knowledge of the directivity.

The first method is used more for directive antennas while the second is rather used more for omnidirectional antennas.

Efficiency

There are essentially three different methods to measure the efficiency of an antenna. The first method consists in the measurement of the gain G and the directivity D of the antenna, and then the relation between gain, directivity and efficiency η results in to obtain the efficiency:

$$\eta = G/Dir \tag{9}$$

The second method uses the measurement of the antenna input impedance and is called the Wheeler Cap Method. For certain antennas, microstrip patches for example, the losses can be modeled by a series resistance or a parallel resistance R_1 with the antenna input resistance R_r . Two different antenna input impedance measurements are performed. The first measurement is done with the antenna in free space and the second measurement with the antenna inside a metallic hemisphere, and the input resistance of the antenna takes the values R_1 and R_2 respectively for these two measurements. If the loss resistance occurs in series, then it is short-circuited by the cap, and if the loss resistance occurs in parallel, then it is open-circuited by the cap. It is possible to have the relation between the efficiency, the radiated power and the losses power. This relation allows one to obtain the efficiency in function of R_r and R_1 and then in function of the measured input resistance R_1 and R_2 . For the series modelling the efficiency is:

$$\eta = R_r / (R_r + R_l) = (R_1 - R_2) / R_1$$
(10)

And for the parallel modelling:

$$\eta = R_{l} / (R_{r} + R_{l}) = (R_{2} - R_{1}) / R_{2}$$
(11)

The third method is called radiometric method. It consists in the measurement of the available noise power at the output of the antenna by means of a radiometer (Ashkenazy et al, 1985). This power U is related to the effective temperature T_e of the antenna and the equivalent temperature T_n of the radiometer by:

$$U = C(T_e + T_n) \tag{12}$$

in which C is a constant. The equivalent temperature T_n is derivated from the noise figure of the radiometer. The effective temperature T_e of the antenna, at physical temperature T_a , is related to the efficiency η of the antenna by:



Fig. 4. The Wheeler Cap Method

 T_t is the temperature target aimed by the antenna. The measurement of the noise power U for two different target temperatures, respectively the cold temperature T_c and the warm temperature T_w , allows the determination of the efficiency η . The cold target temperature is a clear sky and the warm one is an extended absorber at room temperature T_a .

$$\eta = \left[(T_n + T_a)(1 - \varepsilon) \right] / \left[\varepsilon \left(T_c + T_a \right) \right]$$
(14)

$$\varepsilon = U_c / U_w \tag{15}$$

In fact it is also better to measure another antenna with high efficiency, such as a horn, for which the efficiency is equal to 1, with the same cold and warm temperatures as the first antenna. For this second antenna:

$$U_{c}/U_{w} = \delta = (T_{n} + T_{a})/(T_{n} + T_{c})$$
 (16)

The efficiency η of the first antenna is:

$$\eta = [\delta (\varepsilon - 1)] / [\varepsilon (\delta - 1)]$$
(17)
Clear sky
Absorber material
Tempearture Ta
$$\bigcup_{U_c}$$
U_c

Fig. 5. The radiometric method

To achieve high sensitivity measurements the radiometer should have a low internal noise temperature and the temperature difference between the two targets should be as high as possible. The gain and directivity method and the radiometric method are well suited for directive antennas and the Wheeler Cap method for small antennas. A comparison study of these three methods for microstrip antennas has an accuracy of about 2% for the Wheeler Cap method, 10% for the radiometric method and 20% for the gain and directivity method (Pozar & Kaufman, 1988).

Polarization measurement

For harmonic mode, the electromagnetic field radiated by an antenna is polarized. Generally speaking, in one time period, the electric and magnetic fields, observed at a given point, describe a plane curve which is an ellipse. In the case of linear polarization the ellipse is reduced to its major axis. For circular polarization the modulus of the field remains constant. The knowledge of the polarization is equivalent to the knowledge of the ellipse: its axial ratio, the slope of its major axis relative to a reference direction and the sense of the displacement along the ellipse.

The experimental determination of the characteristics of the polarization ellipse can be carried out in several ways. It is possible to use amplitude-only measurements, or two amplitude and phase measurements with two different antennas with independent polarizations. The simplest method consists in using two antennas with linear polarization orthogonal one with the other. The same antenna can be used with two different orientations with orthogonal linear polarizations. Measurements provide the amplitude and the phase of the two field orthogonal components from which it is possible to calculate the characteristics of the polarization ellipse.



Fig. 6. Characteristics of the polarization ellipse

3. Compact antenna test ranges

3.1 Classical compact antenna test ranges

The compact ranges allow the direct far-field pattern measurement of antennas at near-field distance. It is then possible to carry out indoor antenna far field pattern measurements. The principle of compact range is to produce a plane wave at short distance, by means of a focusing system, like a reflector or a lens, illuminated by a primary feed (Johnson et al, 1969). There are two principal kinds of compact ranges. The first kind uses an offset parabolic reflector, with a possible Cassegrain geometry. The second kind uses two orthogonal cylindrical parabolic reflectors. This second geometry is equivalent to the first one with a significantly reduced cost due to the simpler reflectors geometry. The choice of an offset feed prevents aperture blockage and reduces the diffracted energy from the feed structure in the test region.

The main causes of errors on the field in the quiet zone are direct radiation from the primary feed, diffraction from the feed support, diffraction from the edges of the reflector, reflector surface deviations from the theoretical surface and room reflections. For surface deviations two parameters have to be taken into account: the value of the deviation and its extent over the surface. Calculations (Johnson et al, 1973) show that a 0.5dB error on the field corresponds to a 0.007 wavelength surface deviation. The use of serrations or rolled edges reduces the effect of the diffraction from the reflector edges. The positions and the length of the serrations are empirically adjusted so that the energy is diffracted in directions away from the test region for a broad band of frequencies. Room reflections and diffraction from the feed support are minimized by the use of absorbing material. The wise choice of the primary feed position compared to the tested antenna reduces the coupling due to the direct radiation of the feed. Time-gating of the measurement signal improve the performances of the compact range in relation to these last three errors causes.

The dimensions of the quiet zone are about the third of the reflector dimensions. Broadband compact ranges are available: 0.7GHz-100GHz with 3.6m cubic quiet zone. For such ranges,

several feed antennas are used, each antenna covering a half octave frequency band. The different feed antennas are automatically positioned at the reflector focus point and connected to the instrumentation, according to the required frequency. Typical amplitude variations of 1dB and phase variations of 10°, in the quiet zone, can be achieved. The cross polarization is better than 30dB.





Fig. 8. Compact range reflector with serrations or rolled edges

3.2 Hologram compact antenna test range

It is difficult to test large antennas at frequencies above 100GHz. The use of near-field antenna measurements requires deformable coaxial cables or rotary joints with high performances, which is impossible at frequencies that are too high. Conventional reflector-

type compact antenna test ranges require one or more reflectors for which the surface accuracy needs to be better than about 0.01 wavelength, that is to say, for example, $15\mu m$ at 200GHz.

The use of a planar hologram constitutes another solution (Hirvonen et al, 1997). The surface accuracy requirements for an amplitude hologram are less stringent than those for a reflector, and its planar geometry simplifies its realization. The hologram compact antenna test range is a low-cost and easy-to-fabricate structure.

The principle of the hologram is to change the spherical wave front radiated by a source antenna (a horn for example) in a plane wave by means of the transmission through the hologram. It is possible to numerically calculate the structure required to change the known input field into the desired output field. The fabrication of the hologram is simplified by binary amplitude quantization: the local transmittance of the hologram is either 1 or 0. This is obtained by the use of a copper-plated Kapton film using an etching procedure. A hologram of 3 meters diameter has been used for the test of a 1.5 meter diameter reflector antenna at 322GHz. The envelopes of the measured and simulated far-field patterns are similar, but there are relatively important differences between the two far-field patterns.



Fig. 9. Hologram and hologram compact antenna test range

4. Near-field techniques and applications

4.1 Near-field techniques

The principle of near-field techniques is to measure the field radiated by an antenna at a short distance on a given surface surrounding the antenna, then to calculate the far-field starting from the measured near-field (Yaghjian, 1986). Of the several formulations for these techniques, the two principal formulations are the Huygens principle and the modal expansion of the field.

Huygens principle

The tangential components of the electric and magnetic fields E_t and H_t are measured on an arbitrary surface S enclosing the tested antenna. These components allow one to calculate the equivalent electric and magnetic currents J_s and M_s . Then the electric and magnetic

fields can be evaluated everywhere out of the surface S starting from the equivalent currents. This method uses simple calculations, but for large antenna of diameter D the computer time varies like $(D/\lambda)^3$ and can become very long. Moreover the method requires calibrated and ideal probes and generally the measurement of the four field components. The electric and magnetic far-field **E** and **H** are given by the relations:



Fig. 10. The Huygens principle

Modal expansion of the field

In free space the electric and magnetic fields verify the propagation equation. This equation has elementary solutions or modes and a given field is a linear combination of these modes. The knowledge of the field of an antenna is equivalent to the knowledge of the coefficients of the linear combination. The expression of the modes is known for the different systems of orthogonal coordinates: cartesian, cylindrical and spherical. The coefficients of the linear combination are obtained by means of the two tangential field components measurement on a reference surface of the used coordinates system, then using an orthogonality integration.

The case of the planar scanning is simple (Slater, 1991). The measurement of the two tangential components of the field, the electric field $E_t(x,y,z)$ for example, is realized on a plane z=0 following a two dimensional regular grid (axis x and y). The antenna is located at z<0. The tangential components of the plane wave spectrum are obtained from the measured field of the orthogonality integration:

$$A_{t}(k_{x},k_{y},z) = 1/(2\pi) \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} E_{t}(x,y,z) e^{j(k_{x}x+k_{y}y)} dx dy$$
(21)

It is then possible to calculate the electric field in any point thanks to:

$$E(x,y,z) = 1/(2\pi) \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} A(k_x,k_y) e^{-j(k_x x + k_y y + k_z z)} dk_x dk_y$$
(22)

$$k^2 = \omega^2 \varepsilon_0 \mu_0$$
 $k^2 = k_x^2 + k_y^2 + k_z^2$ (23)

The normal component $A_z(k_x,k_y)$ of vector $\mathbf{A}(k_x,k_y)$ is obtained from the local Gauss equation:

$$\mathbf{k} \mathbf{A}(\mathbf{k}_{x}, \mathbf{k}_{y}) = 0 \qquad \mathbf{k} = \mathbf{k}_{x} \mathbf{e}_{x} + \mathbf{k}_{y} \mathbf{e}_{y} + \mathbf{k}_{z} \mathbf{e}_{z}$$
(24)

It is then possible to obtain the near-field of the antenna everywhere from the measurement of the near-field on a given plane. The electric far-field in the direction θ , ϕ and at a distance r is given by the relation:

$$\mathbf{E}(\mathbf{r},\theta,\phi) = \mathbf{j} \ \mathbf{k} \ \cos\theta \ e^{-\mathbf{j}\mathbf{k}\mathbf{r}}/\mathbf{r} \ \mathbf{A}(\mathbf{k}\sin\theta\cos\phi,\mathbf{k}\sin\theta\sin\phi) \qquad \mathbf{k}^{2} = \omega^{2}\varepsilon_{0}\mu_{0}$$
(25)

It would be possible to obtain the magnetic field from the Maxwell-Faraday equation with the knowledge of the electric field.

The sampling spacing on the measurement surface is $\lambda/2$ following rectilinear axis (planar and cylindrical scanning) and $\lambda/2(R+\lambda)$ for angular variable (cylindrical and spherical scanning), R is the radius of the minimal sphere, i.e. the sphere whose centre is on the rotation axis, which contains the whole of the antenna and whose radius is minimal.



Fig. 11. Sampling spacing for the different scanning geometries: $\Delta x = \Delta y = \Delta z = \lambda/2$, $\Delta \phi = \Delta \theta = \lambda/2$ (R+ λ)

Probe correction

In practice, the probe is not an ideal electric or magnetic dipole which measures the nearfield in a point. The far-field pattern of the probe differs appreciably from the far-field of an elementary electric and magnetic dipole. For the accurate determination of electric and magnetic fields from near-field measurements, it is necessary to correct the nonideal receiving response of the probe. The probe remains oriented in the same direction with planar scanning, and the sidelobe field is sampled at an angle off the boresight direction of

the probe. Thus it is necessary to apply probe correction to planar near-field measurements. The problem is the same with cylindrical scanning for the rectilinear axis, and probe correction is also necessary in this case. For spherical scanning, the probe always points toward the test antenna and probe correction is not necessary if the measurement radius is large enough.

The formulation of probe correction is simple for planar scanning. The plane wave spectrum of the measurement $A_{m\nu}$ as definite previously, is the scalar product of the plane wave spectra of the tested antenna A_a and the probe A_{pr} .

$$\mathbf{A}_{\mathbf{m}} = \mathbf{A}_{\mathbf{a}} \mathbf{A}_{\mathbf{p}} = \mathbf{A}_{\mathbf{a}\mathbf{x}} \mathbf{A}_{\mathbf{p}\mathbf{x}} + \mathbf{A}_{\mathbf{a}\mathbf{y}} \mathbf{A}_{\mathbf{p}\mathbf{y}}$$

(26)

The measurement is repeated twice, for two orthogonal orientations between them, of the probe. This results in two equations on A_{ax} and A_{ay} and it is enough to invert this linear system of equations to obtain A_{ax} and A_{ay} .

Different coordinates systems comparison

In the case of planar cartesian and cylindrical coordinates systems, the measurement surface is truncated because the length of a rectilinear axis is limited. In practice, the measurement surface is a rectangle for planar exploration and a cylinder with a finite height for cylindrical exploration. Thus to minimize the effect of the measurement surface truncation, planar nearfield systems are devoted to two-dimensional directive antennas and the cylindrical system requires antennas with directive pattern in at least one plane. Spherical near-field systems are convenient for omni-directional and directive antennas.

Phaseless method

The use of near-field techniques at frequencies above 100GHz is very difficult. This is due to the phase errors induced by coaxial cables or rotary joints whose performances are degraded at these frequencies. In counterpart, it is possible to measure the amplitude of the near-field until very high frequencies. This is why the phaseless methods appeared. These methods consist in the measurement of the near-field on two different surfaces, two parallel planes in front of the antenna for example, and to try to find the phase using an iterative process (Isernia & Leone, 1994). This iterative process consists in passing alternatively from one surface to the other by a near-field to near-field transformation. At the beginning, the distribution of the near-field phase on a surface is arbitrarily selected, a constant phase for example. Then when the near-field is calculated on the other surface, the calculated phase is preserved, and one associates it with the measured near-field amplitude. Then the near-field is calculated on the first surface and one starts again the process again. The process is stopped when the difference between the amplitudes of the computed and measured fields is lower than a given value.

To obtain an accurate reconstructed phase, it is necessary that the near-fields on the two planes are sufficiently different, i.e. the two planes are separated by a sufficient distance. A study shows good results for a low sidelobe shaped reflector antenna with an elliptical aperture with axes 155cm x 52cm at 9GHz (Isernia & Leone, 1995). The two planes are at a distance respectively of 4.2cm and 17.7cm from the antenna. The far-field pattern obtained

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from the near to far-field transformation with phaseless method shows agreement with the reference far-field pattern, up to a -25dB level approximately.



Fig. 16. Phaseless method with two parallel planes configuration

Near-field measurement errors analysis

One of the difficulties related to the use of the near-field techniques is the evaluation of the effect on the far field, of the measurement errors intervening on the near-field. A study allows the identification of the error sources, an evaluation of their level and the value of the induced uncertainties on the far-field, in the case of planar near-field measurements (Newell, 1988). About twenty different error sources are identified as probe relative pattern, gain, polarization, or multiple reflections between probe and tested antenna, measurement area truncation, temperature drift... The main error sources on the maximum gains are the multiple reflections between probe and tested antenna, the power measurement, for a global induced error of 0.23dB. For sidelobe measurement, the main error sources are the multiple reflections between probe and tested antenna, the phase errors, the probe position errors and the probe alignment for a global induced error of 0.53dB on a -30dB sidelobe level. A comparison of the results obtained with four different near-field European ranges shows agreement on the copolar far-field pattern and directivity of a contoured beam antenna (Lemanczyk, 1988).

4.2 Near field applications

Electromagnetic antenna diagnosis

Antenna diagnosis consists in the detection of defects on an antenna. There are essentially two different electromagnetic diagnosis: reflector antenna diagnosis and array antenna diagnosis.



Cylindrical near-field range Fig. 12. Near-field ranges at Supélec.

Spherical near-field range

Reflector antenna diagnosis

For reflector antennas, the diagnosis consists mainly in checking the reflector surface. It is possible to use an optical method to measure the reflector surface. This is a photogrammetric triangulation method (Kenefick, 1971). This method utilizes two or more long-focal length cameras that take overlapping photographs of the surface. This surface is uniformly covered with self-adhesive photographic targets whose images appear on the photographic record. The two-dimensional measurements of the image of the targets are processed with a least squares triangulation to provide the three-dimensional coordinates of each target. The accuracy of this method is of the order of one part in 100000 of the reflector diameter.

It is also possible to perform electromagnetic diagnosis of reflector antenna (Rahmat Samii, 1985). For this method, the knowledge of the amplitude and phase far-field pattern is required. This far-field can be obtained by means of near-field, compact range or direct far-field measurement. The relation between the two-dimensional amplitude and phase far-field and the electric current on the reflector surface is known. This relation can take the form of a two-dimensional Fourier transform at the cost of some approximations, and can then be inverted easily. Finally, the phase of the currents can be interpreted like a deformation starting from the theoretical geometry of the reflector. A study of this method using spherical near-field measurements on a large reflector antenna give good results: small deformations of about one λ diameter and a $\lambda/10$ thickness are detected (Rahmat Samii, 1988).

Array antenna diagnosis

The electromagnetic diagnosis of array antennas consists in detecting defective or badly fed elements on the antenna. To obtain this detection, it is sufficient to rebuild the feeding law of the antenna elements. There are two methods of array antenna diagnosis that primarily exist. The first method uses backward transform from the measurement plane to the antenna surface and is called the spectral method (Lee et al, 1988). The measurement of the radiated near field is performed on a plane parallel to the antenna surface. Then the measured near field is processed to obtain the near field at the location of each element of the array. This processing contains element and probe patterns correction. The feeding of each element is then considered as being proportional to the near field at the location of the element. The second method uses the linear relation between the feeding of each element and the measured near field and is called the matrix method (Wegrowicz & Pokuls, 1991), (Picard et al, 1996), (Picard et al, 1998). The near field is also measured on a plane parallel to the antenna surface. The number of space points is higher than or equal to the number of elements in the array. The linear equation system is numerically inverted. The advantage of the matrix method, compared to the spectral method, is that it uses a number of measurement points significantly weaker. The accuracy of these methods on the reconstructed feeding law is of the order of a few degrees and a few tenth of dB.



Fig. 13. Array antenna diagnosis: measurement configuration

Antennas coupling

The coupling coefficient between two antennas can be obtained by using the fields radiated by these two antennas separately (Yaghjan, 1982). The reciprocity theorem makes it possible to show that the voltage V_{BA} induced by the radiation of an antenna A at the output of an antenna B is

$$V_{BA} = -\iint_{S} [E_a x H_b + H_a x E_b] n dS$$
(27)

S is a close surface surrounding the antenna B,

n is the normal vector to S with the outside orientation,

E_a, H_a electric and magnetic fields radiated by the antenna A,

 E_b,H_b electric and magnetic fields radiated by the antenna A for the emission mode with unit input current,



Fig. 14. The two antennas system for coupling evaluation

The advantage of this method is that it can predict, by calculation, the coupling between the two antennas for any relative position, only by means of their separate radiated near-fields measurements.

Determination of the safety perimeter of base station antennas

An application of the cylindrical near-field to near-field transformations is the determination of the base station antennas safety perimeter. The electric and magnetic near-fields level can be evaluated from the near-field measurements and from the power accepted by the antenna. The comparison of this level with the ICNIRP reference level allows the determination of the safety perimeter (Ziyyat et al, 2001), (ICNIRP, 1998). The accuracy obtained by this method is within a few percent on the calculated near-field. Rapid near-field assessment system

The near-field measurement of a large antenna requires a considerable number of measurement points. Computers' computing power has increased regularly and was multiplied by approximately 100000 between 1981 and 2006. The result is from it that the duration of the far-field calculation decreases regularly and is no longer a problem. On the other hand the duration of measurement can be very important. This is due to the slowness of mechanical displacements. The replacement of the mechanical displacement of the probe by the electronic scanning of a probes array makes it possible to accelerate considerably the measurement rate and to reduce the measurement duration (Picard et al., 1992), (Picard et al, 1998).



Fig. 15. Rapid near-field range at Supélec and principle of rapid near-field assessment systems

5. Electromagnetic field measurement method

The measurement of the radiation of the antennas is indissociable from the measurement of high frequency electromagnetic field. Primarily four different methods for high frequency electromagnetic field measurement exist. These methods differ primarily by the type of connection between the probe and the receiver, this connection could possibly be the cause of many disturbances. The first method is the simplest one. It consists in using of a small dipole probe connected to a receiver with a coaxial line. In order to limit the parasitic effects of the line on the measurement signal, a balun is placed between the line and the dipole. This method makes it possible the measurement of the local value of one component of the electric or magnetic field.



Fig. 17. Measurement of the electric and magnetic field with a dipole probe

In the case of a field whose space variations are very fast, the modulated scattering technique can be used advantageously. This second method consists in the use of a small probe loaded with a nonlinear element like a PIN diode, which is low frequency modulated (Callen & Parr, 1955), (Richmond, 1955), (Bolomey & Gardiol, 2001). The electromagnetic field scattered by this probe is collected by the emitting antenna (monostatic arrangement) or by a specific or auxiliary antenna (bistatic arrangement) called auxiliary antenna. The signal provided by the emitting antenna is proportional to the square of the field radiated at the probe location for the monostatic arrangement, and that provided by the auxiliary antenna is proportional to this field for the bistatic arrangement. These two signals are low frequency modulated like the scattered field, and this amplitude modulation allows one to retrieve this signal among parasitic signals, with coherent detection for example. The low frequency modulation of the diode may be conveyed by resistive lines or by an optical fiber in the case of the optically modulated scattering technique (Hygate & Nye, 1990) so as to limit the perturbations.



The third method uses an electro-optic probe. This probe is a small one like a dipole, and is loaded with an electro-optic crystal like LiNbO₃. The refraction index of the crystal linearly depends on the radiofrequency electric field which is applied to it. The light of a laser is conveyed by an optical fiber and crosses the crystal. The phase variations of the light transmitted through the crystal are measured and are connected linearly to the radiofrequency electric field applied to the crystal. The calibration of the probe makes it possible to know the proportionality factor between the variation of the phase undergone by the light and the amplitude of the measured radiofrequency electric field. This method makes it possible to produce probes with very broad band performances (Loader et al, 2003). In particular, an electric dipole of this type is an excellent time-domain probe: the measurement signal is proportional to the measured time-domain electric field.

The last method is simpler and less expensive than the two preceding ones while making it possible to carry out very local measurements without the disturbances due to the connection between the probe and the receiver. This method uses detected probes (Bowman, 1973) to measure the local electric field. Such a probe is loaded with a schottky diode and detects the RF currents induced by the electric field, to obtain a continuous voltage. This voltage can be measured by a voltmeter. The lines connecting the dipole and the voltmeter are made highly resistive to reduce their parasitic effect. The main defects of this method are

its poor sensitivity and that it provides only the amplitude of the measured field. If the knowledge of the phase is necessary for the application, it must be obtained by means of phaseless methods.



Fig. 19. Electro-optic dipole probe



Fig. 20. Detected probe

6. Instrumentation

The instrumentation used for antenna measurements depends on the temporal mode used: time domain or frequency domain. Network or spectrum analyzer and frequency synthetizer are used for frequency domain measurements. Real time or sampling oscilloscope and pulse generator are used for time domain measurements. Frequency domain antenna measurements

The system of emission-reception the most used for antenna measurements is the vector network analyzer. It allows the measurement of transmission coefficients and it supplies the phase. Its intermediary frequency bandwidth can reach 1MHz, i.e. it allows very high speed measurements, and its dynamic can reach 140dB. It can have several ways of measurement so as to be able to measure simultaneously direct and cross polarizations. Its maximum frequency bandwidth of operation is 30kHz to 1000GHz (with several models). It is also possible to use a scalar network analyzer or a spectrum analyzer coupled to a frequency synthetizer when the measurement of the phase is not necessary as for far-field for example. Time domain antenna measurements

Antenna measurements in the time domain are less frequent than in the frequency domain. The measurement signal is delivered by a pulse generator. Certain characteristics of the pulse can be adjusted: the rise and fall times, the duration, the repetition rate and the amplitude. The receiver is a fast oscilloscope. The real time oscilloscope acquires the measured time response in one step, but its sensitivity is limited and it is very expensive. The sampling oscilloscope requires numerous repetitions of the measurement signal to acquire its time response, but its sensitivity is better and its price is lower than those of the real time oscilloscope. In 2009, the maximum frequency of operation for real time oscilloscope is 20GHz and 75GHz for sample oscilloscope.

Probe

Direct far-field measurements use a source antenna. The dimensions of this source antenna are limited by the distance between this antenna and the tested antenna. A large source antenna increases the measurement signal and decreases the parasitic reflections. The measured polarization is the one of the source antenna.

Near-field measurements can use several types of probe: open-ended circular or rectangular metallic waveguide and electric dipole for narrow band operation (half a octave) and ridged waveguide for broad band operation (a decade).

7. Conclusion

Currently it is possible to measure all the characteristics of an antenna with a good accuracy. Far-field ranges do not have a very good accuracy, due to parasitic reflections for the outdoor ranges and because of the limited distance between the source antenna and the tested antenna for the indoor ranges. The compact range allows one to obtain a direct far-field cut in a relatively short time. The near-field techniques are the most accurate and the most convenient for global antenna radiation testing. Their main defect is the duration of the measurement which rises from the large number of necessary space points. Rapid near-field measurement systems allow one to solve this problem, but the accuracy is less good, and the frequency bandwidth is limited. Progress is necessary in this field. Research relates to the rise in frequency, with for solutions the hologram compact antenna test range and the phaseless methods. The hologram compact antenna test range must improve their accuracy while the phaseless methods must improve their reliability. Electromagnetic diagnosis of antenna must be optimized on a case-by-case basis.

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