Sigma-function solution to the general Somos-6 recurrence via hyperelliptic Prym varieties

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Abstract

We construct the explicit solution of the initial value problem for sequences generated by the general Somos-6 recurrence relation, in terms of the Kleinian sigma-function of genus two. For each sequence there is an associated genus two curve X, such that iteration of the recurrence corresponds to translation by a fixed vector in the Jacobian of X. The construction is based on a Lax pair with a spectral curve S of genus four admitting an involution σ with two fixed points, and the Jacobian of X arises as the Prym variety $\mathrm{Prym}(S, \sigma)$.

1 Introduction

Somos sequences are integer sequences generated by quadratic recurrence relations, which can be regarded as nonlinear analogues of the Fibonacci numbers. They are also known as Gale-Robinson sequences, and as well as arising from reductions of bilinear partial difference equations in the theory of discrete integrable systems, they appear in number theory, statistical mechanics, string theory and algebraic combinatorics [4, 10, 12, 15].

This article is concerned with the general form of the sixth-order recurrence

$$\tau_{n+6}\tau_n = \alpha\tau_{n+5}\tau_{n+1} + \beta\tau_{n+4}\tau_{n+2} + \gamma\tau_{n+3}^2, \tag{1.1}$$

with three arbitrary coefficients α, β, γ. It was an empirical observation of Somos [29] that in the case $\alpha = \beta = \gamma = 1$ the initial values $\tau_0 = \cdots = \tau_5 = 1$ generate a sequence of integers (A006722 in Sloane's Online Encyclopedia of Integer Sequences), which begins

$$1, 1, 1, 1, 1, 1, 3, 5, 9, 23, 75, 421, 1103, 5047, 41783, 281527, \ldots. \tag{1.2}$$

Consequently, the relation (1.1) with generic coefficients is referred to as the Somos-6 recurrence, and the corresponding sequence (τ_n) as a Somos-6 sequence.

The first proof that the original Somos-6 sequence (1.2) consists entirely of integers was an unpublished result of Hickerson (see [16]); it relied on showing that the Somos-6

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recurrence has the Laurent property, meaning that the iterates are Laurent polynomials in the initial data with integer coefficients. To be precise, in the general case the iterates satisfy

$$\tau_n \in \mathbb{Z}[\tau_0^{\pm 1}, \ldots, \tau_5^{\pm 1}, \alpha, \beta, \gamma] \qquad \forall n \in \mathbb{Z},$$

which was proved by Fomin and Zelevinsky as an offshoot of their development of cluster algebras [15]. The latter proof made essential use of the fact that (1.1) is a reduction of the cube recurrence, a partial difference equation which is better known in the theory of integrable systems as Miwa's equation, or the bilinear form of the discrete BKP equation (see [9], for instance). In the general case $\alpha\beta\gamma \neq 0$, (1.1) does not arise from mutations in a cluster algebra, although it does appear in the broader framework of Laurent phenomenon algebras [23].

As was found independently by several people (see e.g. [18, 19, 27, 28] and references), the sequences generated by general bilinear recurrences of order 4 or 5 are associated with sequences of points on elliptic curves, and can be written in terms of the corresponding Weierstrass sigma-function. It was shown in [20] that sequences (τ_n) produced by (1.1) are the first ones which go beyond genus one: in general, they are parametrized by a sigma-function in two variables. To be precise, given a genus 2 algebraic curve X defined by the affine model

$$z^2 = \sum_{j=0}^{5} \bar{c}_j \, s^j \qquad \text{with} \quad \bar{c}_5 = 4 \tag{1.3}$$

in the (s, z) plane, let $\sigma(\mathbf{u})$ denote the associated Kleinian sigma-function with $\mathbf{u} = (u_1, u_2) \in \mathbb{C}^2$, as described in [2] (see also [6, 7]). It gives rise to the Kleinian hyperelliptic functions $\wp_{jk}(\mathbf{u}) = -\partial_j\partial_k \log \sigma(\mathbf{u})$, which are meromorphic on the Jacobian variety $\text{Jac}(X)$ and generalize the Weierstrass elliptic $\wp$ function.

Theorem 1 ([20]). *For arbitrary* $A, B, C \in \mathbb{C}^*$, $\mathbf{v}_0 \in \mathbb{C}^2$, *the sequence with nth term*

$$\tau_n = AB^n C^{n^2-1} \frac{\sigma(\mathbf{v}_0 + n\mathbf{v})}{\sigma(\mathbf{v})^{n^2}} \tag{1.4}$$

satisfies the recurrence (1.1) *with coefficients*

$$\alpha = \frac{\sigma^2(3\mathbf{v})C^{10}}{\sigma^2(2\mathbf{v})\sigma^{10}(\mathbf{v})} \, \hat{\alpha}, \qquad \beta = \frac{\sigma^2(3\mathbf{v})C^{16}}{\sigma^{18}(\mathbf{v})} \, \hat{\beta},$$

$$\gamma = \frac{\sigma^2(3\mathbf{v})C^{18}}{\sigma^{18}(\mathbf{v})} \left(\wp_{11}(3\mathbf{v}) - \hat{\alpha}\wp_{11}(2\mathbf{v}) - \hat{\beta}\wp_{11}(\mathbf{v}) \right), \tag{1.5}$$

where

$$\hat{\alpha} = \frac{\wp_{22}(3\mathbf{v}) - \wp_{22}(\mathbf{v})}{\wp_{22}(2\mathbf{v}) - \wp_{22}(\mathbf{v})}, \quad \hat{\beta} = \frac{\wp_{22}(2\mathbf{v}) - \wp_{22}(3\mathbf{v})}{\wp_{22}(2\mathbf{v}) - \wp_{22}(\mathbf{v})} = 1 - \hat{\alpha}, \tag{1.6}$$

provided that $\mathbf{v} \in \text{Jac}(X)$ *satisfies the constraint*

$$\det \begin{pmatrix} 1 & 1 & 1 \\ \wp_{12}(\mathbf{v}) & \wp_{12}(2\mathbf{v}) & \wp_{12}(3\mathbf{v}) \\ \wp_{22}(\mathbf{v}) & \wp_{22}(2\mathbf{v}) & \wp_{22}(3\mathbf{v}) \end{pmatrix} = 0. \tag{1.7}$$

The preceding statement differs slightly from that of Theorem 1.1 in [20], in that we have used an alternative (but equivalent) expression for $\hat{\alpha}$ in (1.6), and have included an additional parameter C which is needed in what follows. Now while the above result means that the expression (1.4) is a solution of (1.1) with suitable coefficients, it does not guarantee that it is the general solution, in the sense that the sequence (τ_n) can always be written in this way, for a generic choice of initial data and coefficients. The ultimate purpose of this paper is to show that this is indeed the case. Our main result is the solution of the initial value problem by explicit reconstruction of the parameters appearing in (1.4), which yields the following.

Theorem 2. *For a sequence of complex numbers (τ_n) generated by the recurrence (1.1) with generic values of the initial data $\tau_0, \ldots, \ldots \tau_5$ and coefficients α, β, γ, there exists a genus 2 curve X with affine model (1.3) and period lattice Λ, points $\mathbf{v}_0, \mathbf{v} \in \mathrm{Jac}(X) \simeq \mathbb{C}^2 \bmod \Lambda$ with $\mathbf{v}$ satisfying (1.7), and constants $\mathrm{A}, \mathrm{B}, \mathrm{C} \in \mathbb{C}^*$ such that the terms and coefficients are parametrized by the corresponding Kleinian functions according to (1.4) and (1.5), respectively.*

In order to solve the reconstruction problem, it will be convenient to work with a reduced version of the Somos-6 recurrence. The parameters A, B in (1.4) correspond to the group of scaling symmetries $\tau_n \to \mathrm{AB}^n \tau_n$, which maps solutions to solutions, and considering invariance under this symmetry leads to certain quantities x_n, as described in the next paragraph. The parameter C corresponds to covariance under the further scaling $\tau_n \to \mathrm{C}^{n^2} \tau_n$, which maps solutions of (1.1) to solutions of the same recurrence with rescaled coefficients; in due course we will consider quantities that are also invariant with respect to this additional symmetry.

The reduced Somos-6 map. Sequences generated by iteration of the Somos-6 recurrence (1.1) are equivalent to the orbits of the birational map

$$\varphi \, : \, (\tau_0, \tau_1, \ldots, \tau_5) \mapsto (\tau_1, \tau_2, \ldots, \tau_6), \quad \tau_6 = \frac{1}{\tau_0} \left(\alpha \tau_5 \tau_1 + \beta \tau_4 \tau_2 + \gamma \tau_3^2 \right) .$$

As was observed in [20], this map is Poisson with respect to the log-canonical bracket $\{\tau_m, \tau_n\} = (m - n)\tau_m \tau_n$, which has four independent Casimir functions

$$\mathrm{x}_j = \frac{\tau_j \, \tau_{j+2}}{\tau_{j+1}^2}, \qquad j = 0, \ldots, 3; \tag{1.8}$$

these quantities are also invariant under the scaling transformation $\tau_n \to \mathrm{AB}^n \tau_n$. The map φ induces a recurrence of order 4 for a corresponding sequence (x_n), that is

$$\mathrm{x}_{n+4} \, \mathrm{x}_{n+3}^2 \, \mathrm{x}_{n+2}^3 \, \mathrm{x}_{n+1}^2 \, \mathrm{x}_n = \alpha \, \mathrm{x}_{n+3} \, \mathrm{x}_{n+2}^2 \, \mathrm{x}_{n+1} + \beta \, \mathrm{x}_{n+2} + \gamma, \tag{1.9}$$

which is equivalent to iteration of a birational map $\hat{\varphi}$ in $\mathbb{C}^4$ with coordinates $\mathbf{x} = (\mathrm{x}_0, \ldots, \mathrm{x}_3)$. We will refer to $\hat{\varphi}$ as the reduced Somos 6 map.

The map $\hat{\varphi}$ defined by (1.9) preserves the meromorphic volume form

$$\hat{V} = \frac{1}{\mathrm{x}_0 \mathrm{x}_1 \mathrm{x}_2 \mathrm{x}_3} \, \mathrm{dx}_0 \wedge \mathrm{dx}_2 \wedge \mathrm{dx}_2 \wedge \mathrm{dx}_3$$

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for arbitrary values of α, β, γ, and has two independent rational first integrals, here denoted $K_1(\mathbf{x}), K_2(\mathbf{x})$, which are presented explicitly in section 2 below. According to [20], the map $\hat{\varphi}$ is also integrable in the Liouville–Arnold sense [25], at least in the case $\alpha\beta\gamma = 0$. In this paper are concerned with the general case $\alpha\beta\gamma \neq 0$, where a symplectic structure for the map $\hat{\varphi}$ is not known.

On the other hand, a genus 2 curve (1.3) and the corresponding sigma-function solution (1.4), (1.5) of the Somos-6 map φ imply that the solution of (1.9) is

$$\mathrm{x}_n = \frac{\mathrm{C}^2 \sigma(\mathbf{v}_0 + n\mathbf{v})\, \sigma(\mathbf{v}_0 + (n+2)\mathbf{v})}{\sigma^2(\mathbf{v}_0 + (n+1)\mathbf{v})\, \sigma^2(\mathbf{v})}. \tag{1.10}$$

In view of the addition formula for the genus 2 sigma-function [2], the right hand side of (1.10) can be written in terms of Kleinian $\wp$ functions as

$$\mathrm{x}_n = \mathrm{C}^2\Big(\wp_{22}(\mathbf{u})\wp_{12}(\mathbf{v}) - \wp_{12}(\mathbf{u})\wp_{22}(\mathbf{v}) + \wp_{11}(\mathbf{v}) - \wp_{11}(\mathbf{u})\Big) =: \mathcal{F}(\mathbf{u}), \tag{1.11}$$

where $\mathbf{u} = \mathbf{v}_0 + (n+1)\mathbf{v}$. Note that, when $\mathcal{F}$ is considered as a function on the Jacobian, $\mathcal{F}(\mathbf{u})$ is singular if and only if $\mathbf{u} \in (\sigma)$, the theta divisor in $\mathrm{Jac}(X)$ (using the notation in [6]). Then, upon setting $n = 0$, we have the map

$$\psi: \quad \mathbf{u} \mapsto \Big(\mathcal{F}(\mathbf{u}), \mathcal{F}(\mathbf{u}+\mathbf{v}), \mathcal{F}(\mathbf{u}+2\mathbf{v}), \mathcal{F}(\mathbf{u}+3\mathbf{v})\Big) = (\mathrm{x}_0, \mathrm{x}_1, \mathrm{x}_2, \mathrm{x}_3),$$

which is a meromorphic embedding $\psi: \mathrm{Jac}(X) \setminus (\sigma_{0123}) \to \mathbb{C}^4$, where (σ_{0123}) denotes the theta divisor (σ) together with its translates by $\mathbf{v}, 2\mathbf{v}$ and $3\mathbf{v}$. Once Theorem 2 is proved (see section 6), we are able to recover C, $\mathbf{v}$ and $\mathbf{u} = \mathbf{v}_0 + \mathbf{v} \in \mathrm{Jac}(X)$ from the coefficients and initial data of the map, so that we arrive at

Theorem 3. *Generic complex invariant manifolds $\mathcal{I}_K = \{K_1(\mathbf{x}) = k_1, K_2(\mathbf{x}) = k_2\}$ of the map $\hat{\varphi}$ are isomorphic to open subsets of $\mathrm{Jac}(X)$.*

For the purposes of our discussion, it will be more convenient to describe the reduced Somos-6 map in an alternative set of coordinates. We introduce the quantities

$$P_n = -\frac{\delta_1\beta}{\mathrm{x}_n\mathrm{x}_{n+1}}, \qquad R_n = \frac{\delta_1\gamma}{\mathrm{x}_n\mathrm{x}_{n+1}\mathrm{x}_{n+2}}, \qquad \text{with} \quad \delta_1 = \sqrt{-\frac{\alpha}{\beta\gamma}}, \tag{1.12}$$

so that $\mathrm{x}_n = -\gamma P_{n-2}/(\beta R_{n-2})$, and P_0, P_1, R_0, R_1 are birationally related to $\mathrm{x}_0, \mathrm{x}_1, \mathrm{x}_2, \mathrm{x}_3$. Thus, after conjugating $\hat{\varphi}: \mathbb{C}^4 \to \mathbb{C}^4$ by a birational change of variables, we can rewrite it in the form $(P_0, P_1, R_0, R_1) \mapsto (\tilde{P}_0, \tilde{P}_1, \tilde{R}_0, \tilde{R}_1)$, where

$$\tilde{P}_0 = P_1, \quad \tilde{P}_1 = \frac{\mu R_0 R_1}{P_0 P_1}, \quad \tilde{R}_0 = R_1, \quad \tilde{R}_1 = (P_0 + \lambda R_0 R_1 - \lambda P_0 R_1)^{-1}, \tag{1.13}$$

with the coefficients

$$\lambda = \frac{\delta_1\beta^2}{\alpha^2}, \qquad \mu = -\frac{\delta_1\beta^3}{\gamma^2}. \tag{1.14}$$

Observe that, from the analytic formulae (1.5) and (1.10), the quantities P_n, R_n and the coefficients λ, μ are independent of the parameter C.

Outline of the paper. In the next section, we describe the first of our main tools, namely the 3×3 Lax pair for the map $\hat{\varphi}$, which (as announced in [20]) is obtained from the associated Lax representation for the discrete BKP equation. The corresponding spectral curve S yields the first integrals K_1, K_2. However, S is not the required genus 2 curve X, but rather it is trigonal of genus 4, having an involution σ with two fixed points. Then it turns out that the 2-dimensional Jacobian of X, which is the complex invariant manifold of the map $\hat{\varphi}$ according to Theorem 3, can be identified with the Prym subvariety $\mathrm{Prym}(S, \sigma)$ of $\mathrm{Jac}(S)$. (An analogous situation was described recently for an integrable Hénon-Heiles system [11].)

To obtain an explicit algebraic description of $\mathrm{Prym}(S, \sigma)$ and, therefore, of the curve X, we make use of recent work by Levin [24] on the general case of double covers of hyperelliptic curves with two branch points. All relevant details are given in section 3.

In section 4 it is shown how the discrete Lax pair allows a description of the map $\hat{\varphi}$ as a translation on $\mathrm{Prym}(S, \sigma) \subset \mathrm{Jac}(S)$ by a certain vector. This translation is subsequently identified with a specific degree zero divisor on X representing the required vector $\mathbf{v} \in \mathrm{Jac}(X)$, and in section 5 we also explicitly find degree zero divisors on X representing the vectors $2\mathbf{v}, 3\mathbf{v}$. This enables us to rewrite the determinantal constraint (1.7) in terms of the above three divisors, and then observe that it is trivially satisfied.

In section 6, all of the required ingredients are ready to present the reconstruction of the sigma-function solution (1.4) from the initial data and coefficients, which proves Theorem 2. We also provide a couple of explicit examples, including the original Somos-6 sequence (1.2). The paper ends with some conclusions, followed by an appendix which includes the derivation of the Lax pair and another technical result.

2 The Lax pair, its spectral curve, and related Jacobian varieties

The key to the solution of the initial value problem for the Somos-6 recurrence is the Lax representation of the map $\hat{\varphi}$.

Theorem 4. *The mapping $\hat{\varphi} : \mathbb{C}^4 \to \mathbb{C}^4$ is equivalent to the discrete Lax equation*

$$\tilde{\mathbf{L}}\mathbf{M} = \mathbf{M}\mathbf{L}, \tag{2.1}$$

with

$$\mathbf{L}(x) = \begin{pmatrix} \dfrac{A_2 x^2 + A_1 x}{x + \lambda} & \dfrac{A_2' x^2 + A_1' x}{x + \lambda} & \dfrac{A_1'' x + A_0''}{x + \lambda} \\ B_2\, x^2 + B_1 x & B_1' x & B_1'' x + B_0'' \\ C_2 x^2 + C_1 x & C_2' x^2 + C_1' x & C_1'' x + C_0'' \end{pmatrix}, \tag{2.2}$$

$$\mathbf{M}(x) = \frac{1}{R_0} \begin{pmatrix} -1 & 1 & 0 \\ -\frac{x}{\lambda} - 1 & 1 & \frac{1}{\lambda} \\ 0 & (\lambda P_0 R_1 R_2 + 1)x & -P_0 R_2 \end{pmatrix}, \tag{2.3}$$

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where $R_2 = \tilde{R}_1$ as in (1.13), and

$$
\begin{aligned}
A_1 = \ & P_0\left(P_1 + \frac{1}{R_1} - \frac{P_1}{R_0 R_1}\right) + \mu\left(\frac{(R_0 P_1 + 1)R_1}{P_0 P_1} - R_0 - R_1\right) \\
& + \lambda\left(R_0 + R_1 - P_0 - P_1 + P_1 R_0 R_1 - P_0 P_1 R_1 + \frac{P_0 P_1 - 1}{R_0}\right) \\
& + \lambda\mu\left(R_0 R_1 - \frac{(P_1 R_0 + 1)R_0 R_1}{P_0 P_1}\right) + \frac{\mu}{\lambda},
\end{aligned}
$$

$$
\begin{aligned}
A_2 = \ & P_0 + P_1 + \lambda(R_0 - P_0)R_1 + \frac{1}{R_0 R_1}\left(R_0 - P_1 - \frac{1}{\lambda}\right) \\
& + \mu\frac{(R_1 - P_0 - P_1)R_0 - P_0 R_1}{P_0 P_1} + \frac{\mu}{\lambda}\left(\frac{1}{P_0} + \frac{1}{P_1}\right) + \lambda\mu\frac{(P_0 - R_0)R_0 R_1}{P_0 P_1},
\end{aligned}
$$

$$
\begin{aligned}
A_1' = \ & \lambda\left(P_0 - R_1 - P_0 P_1 R_1 + \frac{1 - P_0 P_1}{R_0}\right) + \mu\left(1 - \frac{1}{P_0 P_1}\right)R_1 \\
& + \lambda\mu\left(\frac{1}{P_0 P_1} - 1\right)R_0 R_1, \qquad A_2' = \lambda P_0 R_1 + \mu\frac{R_1}{P_1} - \lambda\mu\frac{R_0 R_1}{P_1},
\end{aligned}
$$

$$
A_0'' = \frac{P_0 P_1}{R_0 R_1} - P_0 P_1 - \frac{P_0}{R_1} + \mu R_0 - \frac{\mu}{\lambda},
$$

$$
A_1'' = \frac{P_1}{R_0 R_1} - \frac{1}{R_1} - P_0 - P_1 + \mu\left(R_0 - \frac{1}{\lambda}\right)\left(\frac{1}{P_0} + \frac{1}{P_1}\right) + \frac{1}{\lambda R_0 R_1},
$$

$$
B_1 = R_0 - P_0 - P_1 + \frac{P_0 P_1 - 1}{R_0} + \frac{1}{\lambda}\left(\frac{P_0(R_0 - P_1)}{R_0 R_1}\right), \qquad B_2 = -\frac{P_1}{\lambda R_0 R_1},
$$

$$
B_1' = P_0 + \frac{1 - P_0 P_1}{R_0}, \qquad B_0'' = \frac{P_0(P_1 - R_0)}{\lambda R_0 R_1}, \qquad B_1'' = \frac{P_1}{\lambda R_0 R_1},
$$

$$
C_1 = \mu\left(\frac{R_0 P_1 + 1}{P_0 P_1} - 1\right)R_1 + \frac{\mu}{\lambda},
$$

$$
C_2 = P_0 + \lambda(R_0 - P_0)R_1 + \mu\frac{(R_0 - P_0)R_1}{P_0 P_1} - \frac{1}{\lambda R_0 R_1} + \frac{\mu}{\lambda}\left(\frac{1}{P_0} + \frac{1}{P_1}\right),
$$

$$
C_1' = \mu\left(1 - \frac{1}{P_0 P_1}\right)R_1, \qquad C_2' = \lambda P_0 R_1 + \mu\frac{R_1}{P_1},
$$

$$
C_0'' = -\frac{\mu}{\lambda}, \qquad C_1'' = -P_0 + \frac{1}{\lambda R_0 R_1} - \frac{\mu}{\lambda}\left(\frac{1}{P_0} + \frac{1}{P_1}\right).
$$

Proof: The equation (2.1) can be checked directly with computer algebra. For the rather more straightforward origin of this complicated-looking Lax pair, see the appendix. $\square$

The characteristic equation $\det(\mathbf{L}(x) - y\mathbf{1})$ defines the spectral curve $S \subset \mathbb{C}^2(x, y)$, which, after elimination of the common factor $1/(x + \lambda)$, is given by

$$
f(x, y) := (x+\lambda)\,y^3 + (x\,K_1 + \mu + x^2\,K_2)\,y^2 - (\mu\,x^4 + K_1\,x^3 + x^2\,K_2)\,y - \lambda\,x^4 - x^3 = 0, \quad (2.4)
$$

where K_1, K_2 are independent first integrals, namely

$$K_1 = \frac{\hat{K}_1(P_0, P_1, R_0, R_1)}{P_0 P_1 R_0 R_1}, \quad K_2 = \frac{\hat{K}_2(P_0, P_1, R_0, R_1)}{P_0 P_1 R_0 R_1}, \tag{2.5}$$

with
$$\begin{aligned}
\hat{K}_1 ={}& \lambda \mu R_0^2 R_1^2 (R_0 \, P_1 - P_0 \, P_1 + 1) \\
&+ \lambda R_0 \, R_1 \, P_0 \, P_1 \, (R_1 \, P_0 \, P_1 + P_0 - P_1 R_0 R_1 - R_1 + P_1 - R_0) \\
&- \mu R_0 R_1 (R_1 - P_0 - R_1 \, P_0 \, P_1 - R_0 \, P_0 \, P_1 - P_1 + P_1 \, R_0 \, R_1) \\
&- P_0 \, P_1 \, (P_0 \, P_1 \, R_0 \, R_1 - P_0 \, P_1 + 1 + P_0 \, R_0), \\
\hat{K}_2 ={}& \lambda \, \mu {R_0}^2 \, {R_1}^2 \, (R_0 - P_0) - \lambda R_0 R_1^2 P_0 P_1 \, (R_0 - P_0) \\
&+ \mu \, R_0 \, R_1 \, (P_0 R_1 - R_1 \, R_0 + P_0 R_0 + R_0 \, P_1) \\
&- P_0 \, P_1 \, (-R_1 \, P_0 \, P_1 + P_0 \, R_0 \, R_1 + P_1 R_0 R_1 + R_1 - P_1 + R_0). \tag{2.6}
\end{aligned}$$

Remark 5. *Replacing the variables P_0, P_1, R_0, R_1 by the expressions (1.12) yields the first integrals $K_1(\mathbf{x}), K_2(\mathbf{x})$ of the reduced map $\hat{\varphi}$ in the original variables $x_0, \ldots, x_3$. These are seen to be rescaled versions of the quantities H_1, H_2 derived in [20] from higher order bilinear relations, according to*

$$K_1 = \frac{\beta H_1}{\alpha \gamma^2}, \qquad K_2 = \frac{\delta_1 \beta H_2}{\alpha \gamma}. \tag{2.7}$$

The curve S is trigonal of genus 4 and has an interesting involution $\sigma : (x, y) \to (1/x, 1/y)$ with two fixed points, namely $(1, 1)$ and $(-1, 1)$.

We compactify S by embedding it in $\mathbb{P}^2$ with homogeneous coordinates $(X : Y : Z)$, where $x = X/Y, y = Y/Z$. The compact curve has a singularity at $(0 : 1 : 0)$. After regularization, this point gives two points at infinity: the first one is $\bar{\mathcal{O}} = (x = \infty, y = \infty)$ with the Laurent expansion

$$x = \frac{1}{\tau^2} + O(\tau^{-1}), \quad y = \frac{1}{\sqrt{\mu}\,\tau^3} + O(\tau^{-2})$$

with respect to a local parameter τ near $\bar{\mathcal{O}}$; and the second is $\bar{\mathcal{O}}_1 = (x = -\lambda, y = \infty)$, with the Laurent expansion
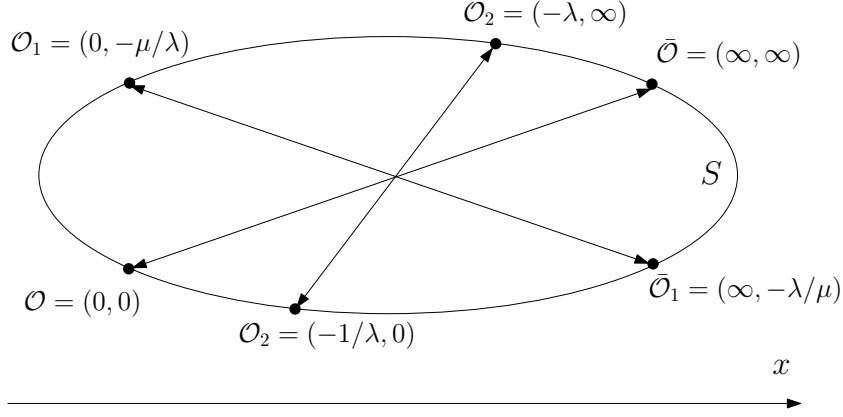
$$x = -\lambda + O(\tau), \quad y = -\frac{\lambda}{\tau} + O(1).$$

The third point at infinity $\bar{\mathcal{O}}_2 = (x = \infty, y = -\lambda/\mu)$ comes from $(1 : 0 : 0)$ and has the expansion

$$x = \frac{1}{\tau} + O(1), \quad y = -\frac{\lambda}{\mu} + O(\tau).$$

Under the action of σ, these points are in involution with the following three finite points:

$$\mathcal{O} = (x = 0, y = 0) \quad \text{with} \quad x = \tau^2, \quad y = \frac{\tau^3}{\sqrt{\mu}} + O(\tau^4);$$

$$\mathcal{O}_1 = (x = 0, y = -\mu/\lambda) \quad \text{with} \quad x = \tau;$$
$$\mathcal{O}_2 = (x = -1/\lambda, y = 0) \quad \text{with} \quad y = \tau + O(\tau^2).$$

The above three pairs of points on S will play an important role, so we depict them on the diagram above, with arrows denoting the involution σ. The curve S can be viewed as a 3-fold cover of $\mathbb{P}^1$ with affine coordinate x. As follows from the above description, the points $\mathcal{O}, \bar{\mathcal{O}}$ are ordinary branch points of the covering, and there is no branching at the points $\mathcal{O}_1, \bar{\mathcal{O}}_1, \mathcal{O}_2, \bar{\mathcal{O}}_2$. It follows that the divisors of zeros and poles of the coordinates x, y on S are

$$(x) = 2\mathcal{O} + \mathcal{O}_1 - 2\bar{\mathcal{O}} - \bar{\mathcal{O}}_1, \quad (y) = 3\mathcal{O} + \mathcal{O}_2 - 3\bar{\mathcal{O}} - \bar{\mathcal{O}}_2. \tag{2.8}$$

Observe that a generic complex 2-dimensional invariant manifold $\mathcal{I}_K$ of the reduced Somos-6 map $\hat{\varphi}$ cannot be the Jacobian of S, as the latter has genus 4. The curve S is a 2-fold covering of a curve $G = S/\sigma$ whose genus is 2, by the Riemann–Hurwitz formula. The involution σ extends to $\mathrm{Jac}(S)$ which then contains two Abelian subvarieties: the Jacobian of G, which is invariant under σ, and the 2-dimensional Prym variety, denoted $\mathrm{Prym}(S, \sigma)$, which is anti-invariant with respect to σ. It will play a key role in the description of the complex invariant manifolds of the map $\hat{\varphi}$. For this purpose it is convenient to recall some properties of Prym varieties corresponding to our case.

3   Hyperelliptic Prym varieties

Generic double cover of a hyperelliptic curve with two branch points.   Consider a genus g hyperelliptic curve C: $v^2 = f(u)$, where $f(u)$ is a polynomial of degree $2g + 1$ with simple roots. As was shown in [24], any double cover of C ramified at two finite points $P = (u_P, y_P), Q = (u_Q, y_Q) \in C$ (which are not related to each other by the hyperelliptic involution on C, i.e., $u_P \neq u_Q$) can be written as a space curve of the form

$$\widetilde{C} : \mathrm{z}^2 = v + h(u), \quad v^2 = f(u), \tag{3.1}$$

where $h(u)$ is a polynomial of degree $g + 1$ such that

$$h^2(u) - f(u) = (u - u_P)(u - u_Q)\rho^2(u)$$

with $\rho(u)$ being a polynomial of degree g. (Here u_P or u_Q may or may not coincide with roots of $\rho(u)$.) Thus $\widetilde{C}$ admits the involution $\sigma : (u, v, \mathrm{z}) \mapsto (u, v, -\mathrm{z})$, with fixed points

$(u_P, y_P, 0), (u_Q, y_Q, 0) \in \widetilde{C}$. Then the genus of $\widetilde{C}$ is $2g$, and it was shown by Mumford [26] and Dalaljan [8] that

- $\mathrm{Jac}(\widetilde{C})$ contains two g-dimensional Abelian subvarieties: $\mathrm{Jac}(C)$ and the Prym subvariety $\mathrm{Prym}(\widetilde{C}, \sigma)$, with the former invariant under the extension of σ to $\mathrm{Jac}(\widetilde{C})$, and the latter anti-invariant;

- $\mathrm{Prym}(\widetilde{C}, \sigma)$ is principally polarized and is the Jacobian of a *hyperelliptic* curve C'.

It was further shown recently by Levin [24] that the second curve C' can be written explicitly as

$$w^2 = h(u) + Z, \quad Z^2 = h^2(u) - f(u) \equiv (u - u_P)(u - u_Q)\rho^2(u), \tag{3.2}$$

which is equivalent to the plane curve $w^4 - 2h(u)w^2 + f(u) = 0$. The latter can be transformed explicitly to a hyperelliptic form by an algorithm given in [24].

In order to apply the above results to obtain an explicit description of the Prym variety $\mathrm{Prym}(S, \sigma) \subset \mathrm{Jac}(S)$ in our case, we will need

Proposition 6. *1) The quotient of S by the involution σ is the genus 2 curve $G \subset \mathbb{C}^2(T, Y)$ given by the equation*

$$G(T, Y) := AY^3 + B(T)Y^2 + [C(T) - 3A]Y + D(T) = 0, \tag{3.3}$$

where $A = -\mu^2$, $B(T) = \mu(-K_2 T^2 + T K_1 + 2\lambda + 2K_2 - T^3 + 3T)$,

$$\begin{aligned}
C(T) &= \lambda T^4 \mu + (K_1 \mu + \lambda) T^3 + (\lambda K_2 - 4\mu\lambda - K_2\mu + K_1) T^2 \\
&\quad + (K_1 K_2 - 3\lambda - K_2 - \lambda K_1 - \mu - 3K_1\mu) T + 2\mu\lambda - \lambda^2 - (K_1{}^2 + K_2{}^2 + 1) \\
&\quad - 2(K_1 + \lambda K_2 - \mu K_2), \\
D(T) &= (\mu^2 + \lambda^2) T^4 + 2\lambda T^3 K_1 + (K_1{}^2 - 4\lambda^2 - 4\mu^2 - 2\lambda K_2 + 1) T^2 \\
&\quad + (-2K_1 K_2 - 6\lambda K_1 + 2K_2 - 2\lambda)T + 2(\lambda^2 + \mu^2) \\
&\quad - 2(K_1{}^2 - K_2{}^2 + 1) - 4(K_1 - \lambda K_2).
\end{aligned}$$

The double cover $\pi : S \to G$ is described by the relations

$$T = \frac{x}{y} + \frac{y}{x}, \quad Y = y + \frac{1}{y}, \tag{3.4}$$

and the images of the branch points $P = (x = 1, y = 1), Q = (x = -1, y = 1) \in S$ on G are $P = (T = 2, Y = 2), Q = (T = -2, Y = 2)$.

2) The curve G is equivalent to the following curve $C \subset \mathbb{C}^2(u, v)$ in hyperelliptic form:

$$\begin{aligned}
C : \quad v^2 &= \mathrm{P}_6(u), \tag{3.5} \\
\mathrm{P}_6(u) &= 1 + (4\mu - 2K_2) u + (4\mu\lambda + K_2{}^2 - 2K_1) u^2 \\
&\quad + (2\lambda - 10\mu + 2K_1 K_2) u^3 + (-8\mu\lambda + K_1{}^2 - 2\lambda K_2 + 2\mu K_2) u^4 \\
&\quad + (4\mu - 2\lambda K_1 + 2\mu K_1) u^5 + (\mu + \lambda)^2 u^6.
\end{aligned}$$

The birational transformation between G and C is described by the relations

$$T = -\frac{1}{2} \frac{(\lambda + \mu)u^3 + (2 + K_1)u^2 + (K_2 + 2\lambda)u + 1 - v}{u\,(1 + \lambda\,u)}, \tag{3.6}$$

$$Y = -\frac{1}{2\mu\,u^3(1 + \lambda\,u)} \Big[(\lambda - \mu)^2 u^4 - (\lambda - \mu)(K_1 - 1)u^3$$

$$+ ((\mu - \lambda)K_2 + 4\lambda\mu - K_1)u^2 + (\lambda + 3\mu - K_2)u + 1 - (1 + \mu u + \lambda u)v \Big]. \tag{3.7}$$

The branch points $P = (T = 2, Y = 2), Q = (T = -2, Y = 2) \in G$ on C are, respectively, $P = (u_P, v_P), Q = (u_Q, v_Q)$ with

$$u_P = -\frac{F_2}{F_1} := -\frac{2 + \lambda - \mu + K_2}{2\,\lambda + 1 + 2\,\mu + K_1} \qquad u_Q = \frac{\bar{F}_2}{\bar{F}_1} := \frac{-2 + \lambda - \mu + K_2}{2\,\lambda - 1 + 2\mu - K_1}, \tag{3.8}$$

$$v_P = \frac{1}{F_1^3}\left(F_1^3 - (2 + \lambda + \mu)F_1^2 F_2 + (2\lambda - 2\mu + 1)F_1 F_2^2 - (\lambda + \mu)F_2^3 \right),$$

$$v_Q = \frac{1}{\bar{F}_1^3}\left(\bar{F}_1^3 + (-2 + \lambda + \mu)\bar{F}_1^2 \bar{F}_2 + (-2\lambda + 2\mu + 1)\bar{F}_1 \bar{F}_2^2 - (\lambda + \mu)\bar{F}_2^3 \right). \tag{3.9}$$

Proof: Applying the substitution (3.4) to the polynomial $G(T, Y)$ we get the product $f(x, y)\tilde{f}(x, y)/(x^4\,y^2)$, where

$$\tilde{f} = (x^3\,y^2 - x\,y^5)\,K_1 + (x^2\,y^4 - x^2 y^3)K_2 + \lambda x^4\,y + \mu\,x^4 - x\,y^4 + y^3\,x^3 - y^7\,\mu - \lambda\,y^6,$$

and this product is zero due to the equation (2.4). Hence T, Y satisfy $G(T, Y) = 0$. The proof of the other items is a direct calculation (which we made with Maple). $\qquad \square$

Prym(S, σ) is isomorphic to the Jacobian of a second genus 2 curve C', and in order to find its equation by applying the algorithm of [24] described above, it is convenient to represent the curve S in a form similar to (3.1).

Proposition 7. *The spectral curve S is equivalent to the space curve*

$$\tilde{C} : \ v^2 = \mathrm{P}_6(u), \quad w^2 = 4\mu u^6(1 + \lambda u)^2\,(Y(u, v) - 2)(Y(u, v) + 2) \equiv h(u) + g(u)v, \tag{3.10}$$

where $\mathrm{P}_6(u)$ is given by (3.5), $Y(u, v)$ by (3.7), and $g(u), h(u)$ are polynomials of degree 5 and 8 respectively, obtained by replacing v^2 in the right hand side of the second equation by $\mathrm{P}_6(u)$. On $\tilde{C}$ the involution σ is given by $(u, v, w) \to (u, v, -w)$, and its fixed points are $P = (u_P, v_P, 0), Q = (u_Q, v_Q, 0)$.

Explicit expressions for $h(u), g(u)$ are relatively long and are not shown here. Observe that $h(u), g(u), \mathrm{P}_6(u)$ do not have common roots and that

$$h^2(u) - g^2(u)\mathrm{P}_6(u) = 16\,\mu^2\,u^6\,(1 + \lambda\,u)^2(\bar{F}_2 - u\bar{F}_1)(F_2 + uF_1) \cdot Q^2(u), \tag{3.11}$$

$$Q(u) = 2(\lambda - \mu)\,u^2 + (1 - K_1)u - K_2 - \lambda + \mu,$$

which, in view of the expressions (3.8) for u_P, u_Q, yields

$$h^2(u) - g^2(u)\mathrm{P}_6(u) = F_1\bar{F}_1\,(u - u_P)(u - u_Q)\,u^6\,(1 + \lambda\,u)^2\,Q^2(u). \tag{3.12}$$

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Thus, the function w^2 in (3.10) is meromorphic on the hyperelliptic curve C and has simple zeros only at P, Q, and even order zeros elsewhere. It also has only even order poles at the two points at infinity on C. Hence, $\tilde{C}$ is a double cover of C ramified at P, Q only, as expected.

The above description of the curves and coverings can be briefly summarized in the diagram below.

$$
\begin{array}{ccc}
S & = & \tilde{C} \\
{\scriptstyle 2:1}\Big\downarrow & & \Big\downarrow{\scriptstyle 2:1} \\
G & = & C
\end{array}
$$

(The horizontal equals signs denote birational equivalence.)

Proof of Proposition 7: In view of relations (3.4), (3.7), the curve S can be written as

$$
v^2 = \mathrm{P}_6(u), \quad \frac{1}{y} + y = Y(u, v) \iff 2y = Y(u, v) + \sqrt{Y^2(u, v) - 4}.
$$

Under the birational transformation $(u, v, y) \to (u, v, \hat{w} = 2y - Y(u, v))$, the above reads

$$
v^2 = \mathrm{P}_6(u), \quad \hat{w}^2 = Y^2(u, v) - 4.
$$

Then by the substitution $\hat{w} = w/(2\mu u^3(1 + \lambda u))$, the latter equations take the form (3.10). Next, since σ leaves the curve C invariant, it does not change the coordinates u, v, so it only flips the sign of w. Finally, since $Y(P) = Y(Q) = 2$, from (3.10) we get $w(P) = w(Q) = 0$. $\qquad\square$

Observe that the equation (3.10) of $\tilde{C}$ does not have the same structure as that of the model curve (3.1): the degrees of the corresponding polynomials do not match. Hence, the formula (3.2) for the second hyperelliptic curve C' is not directly applicable to (3.10). For this reason, below we adapt the derivation of (3.2) to our situation.

Tower of curves and Jacobians. Following the approach of [8], consider the tower of curves

$$
C \xleftarrow{\ \pi\ } \tilde{C} \xleftarrow{\ \tilde{\pi}\ } \widetilde{\tilde{C}} \xrightarrow{\ \pi_1\ } C' \tag{3.13}
$$

where $\tilde{C}$ is given by (3.10) and $\widetilde{\tilde{C}}$ is a double cover of $\tilde{C}$ given by

$$
\widetilde{\tilde{C}} : v^2 = \mathrm{P}_6(u), \quad w^2 = h(u) + g(u)v, \quad \bar{w}^2 = h(u) - g(u)v. \tag{3.14}
$$

The covering $\widetilde{\tilde{C}} \to \tilde{C}$ is ramified at the points on $\tilde{C}$ where the function $\bar{w}^2$ has *simple* zeros or poles. As shown above, the function $w^2 = h(u) + g(u)v$ has precisely 2 simple zeros $P = (u_P, v_P), Q = (u_Q, v_Q)$ and no simple poles on C. Hence, $h(u) - g(u)v$ has only two simple zeros $\bar{P} = (u_P, -v_P), \bar{Q} = (u_Q, -v_Q)$ on C. Since $\tilde{C}$ is a double cover of C and $\bar{P}, \bar{Q}$ are not branch points of $\tilde{C} \to C$, the function $\bar{w}^2$ has four simple zeros on $\tilde{C}$, namely $\bar{P}, \bar{Q}, \sigma(\bar{P}), \sigma(\bar{Q})$. Hence $\widetilde{\tilde{C}} \to \tilde{C}$ is ramified at the latter four points, and so, by the Riemann–Hurwitz formula, the genus of $\widetilde{\tilde{C}}$ equals 9.

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The "big" curve $\widetilde{\widetilde{C}}$ has various involutions, one of which is

$$\sigma_1 \; : \; (u, v, w, \bar{w}) \to (u, -v, \bar{w}, w),$$

and the last curve in the tower (3.13) is the genus 2 quotient curve $C' = \widetilde{\widetilde{C}}/\sigma_1$. The corresponding projection $\widetilde{\widetilde{C}} \to C'$ is denoted π_1. The projections $\tilde{\pi}$ and π_1 are explicitly given by

$$\tilde{\pi}(u, v, w, \bar{w}) = (u, v, w), \quad \pi_1(u, v, w, \bar{w}) = (u, W = (w + \bar{w})/\sqrt{2}, Z = w \cdot \bar{w}). \quad (3.15)$$

The tower (3.13) is a part of a tree of curves introduced in [8] for the general case of a genus g hyperelliptic curve C. As was shown in [8], the tree of curves implies relations between the corresponding Jacobian varieties described by the following diagram,

$$
\begin{array}{ccc}
\mathrm{Jac}(C') & \xrightarrow{\;\pi_1^*\;} & \mathrm{Jac}(\widetilde{\widetilde{C}}) \\[4pt]
\parallel & & \tilde{\pi}^* \big\uparrow \\[4pt]
\mathrm{Prym}(\tilde{C}, \sigma) & \longrightarrow & \mathrm{Jac}(\tilde{C}) \\[4pt]
& & \pi^* \big\uparrow \\[4pt]
& & \mathrm{Jac}(C)
\end{array}
$$

where arrows denote inclusions. The diagram indicates that $\mathrm{Jac}(C')$ is isomorphic to $\mathrm{Prym}(\tilde{C}, \sigma)$.

Following [24], the curve C' can be written in terms of u and the symmetric functions $w + \bar{w}$, $w\bar{w}$. In view of (3.14), we have

$$(w + \bar{w})^2 = 2h(u) + 2w\,\bar{w}, \quad (w\,\bar{w})^2 = h^2(u) - v^2 g^2(u).$$

Setting here $W = (w+\bar{w})/\sqrt{2}, Z = w\,\bar{w}$, one obtains equations defining C' in $\mathbb{C}^3(u, Z, W)$:

$$C' \; : \; W^2 = h(u) + Z, \quad Z^2 = h^2(u) - g^2(u)\mathrm{P}_6(u); \quad (3.16)$$

this leads to the single equation $C' \; : \; W^4 - 2h(u)W^2 + g^2(u)\mathrm{P}_6(u) = 0$.

Remark 8. *In the special case that the polynomial $h^2(u) - g^2(u)\mathrm{P}_6(u)$ is a perfect square $f^2(u)$, the latter equation admits the factorization*

$$(W^2 - h(u) - f(u))(W^2 - h(u) + f(u)) = 0, \quad (3.17)$$

and C' is a union of two elliptic curves (possibly singular ones). This situation will be considered in detail elsewhere.

A hyperelliptic form of C'. We return to the general case when $h^2(u) - g^2(u)\mathrm{P}_6(u)$ is not a perfect square. Using the factorization (3.12), in (3.16) we can write

$$Z^2 = 16\,\mu^2\,u^6\,(1 + \lambda\,u)^2\,t^2\,(u - u_P)^2 F_1^2\,Q^2(u), \quad (3.18)$$

where we set

$$t^2 := -\frac{\bar{F}_1}{F_1}\left(\frac{u - u_Q}{u - u_P}\right) = \frac{-\bar{F}_1 u + \bar{F}_2}{F_1 u + F_2} \tag{3.19}$$

and $Q(u)$ is specified in (3.12). Solving the last equation with respect to u, we get

$$u = \frac{\bar{F}_2 - t^2 F_2}{t^2 F_1 + \bar{F}_1}. \tag{3.20}$$

Then equations (3.16) read

$$W^2 = h(u) + 4\mu\, t\, u^3\, (1 + \lambda\, u)\, Q(u)\, (F_1 u + F_2). \tag{3.21}$$

Replacing u here by (3.20), we obtain

$$W^2 = 2H\frac{\mathcal{P}^2(t)}{(t^2 F_1 + \bar{F}_1)^8}(t+1)^2\, \mathrm{R}_6(t), \tag{3.22}$$

$$H = (\lambda + \mu)K_2 + \lambda^2 - \mu^2 - K_1 - 1 = (\lambda + \mu)F_2 - F_1, \tag{3.23}$$

where $\mathcal{P}(t)$ is a polynomial in t of degree 4, and $\mathrm{R}_6(t)$ is specified below. We write

$$\mathcal{P}(t) = \frac{8H^2}{(F_1 u + F_2)^2}\left(2(\mu - \lambda)u^2 + (K_1 - 1)u + F_2 - 2\right) - 2Ht(t^2 F_1 + \bar{F}_1),$$

in a concise form, where u should be replaced by (3.20).

Removing perfect squares from the right hand side of (3.22), i.e. introducing a new variable $\mathcal{W}$ such that

$$W = \frac{\sqrt{2H}\,\mathcal{P}(t)(t+1)}{(t^2 F_1 + \bar{F}_1)^4}\,\mathcal{W}, \tag{3.24}$$

we finally obtain C' in hyperelliptic form as

$$\mathcal{W}^2 = \mathrm{R}_6(t) := c_6 t^6 + c_5 t^5 + \cdots + c_1 t + c_0, \tag{3.25}$$

where
$$c_6 = (\lambda + \mu)F_2^3 - (2\lambda - 2\mu + 1)F_1 F_2^2 + (\lambda + \mu + 2)F_1^2 F_2 - F_1^3,$$
$$c_0 = (\lambda + \mu)\bar{F}_2^3 - (2\lambda - 2\mu - 1)\bar{F}_1 \bar{F}_2^2 + (\lambda + \mu - 2)\bar{F}_1^2 \bar{F}_2 + \bar{F}_1^3,$$
$$c_4 = F_1^3 - (\mu + \lambda + 6)F_1^2 F_2 - 8(\mu + \lambda - 1)F_1^2 + (6\lambda - 6\mu + 1)F_1 F_2^2$$
$$\qquad + 8(\mu^2 + 2\mu\lambda + \lambda^2 + 4\mu - 1)F_2 F_1 - (\mu + \lambda)F_2^3 + 8(\mu^2 - \lambda^2 + \mu + \lambda)F_2^2,$$
$$c_2 = -\bar{F}_1^3 - (\mu + \lambda - 6)\bar{F}_2\bar{F}_1^2 + 8(\mu + \lambda + 1)\bar{F}_1^2 + (6\lambda - 6\mu - 1)\bar{F}_1\bar{F}_2^2$$
$$\qquad + 8(\mu^2 + 2\mu\lambda + \lambda^2 - 4\mu - 1)\bar{F}_1\bar{F}_2 - (\mu + \lambda)\bar{F}_2^3 + 8(\mu^2 - \lambda^2 - \mu - \lambda)\bar{F}_2^2,$$
$$c_5 = 2H(F_1^2 - F_2^2), \quad c_3 = 4H(F_1\bar{F}_1 + F_2\bar{F}_2), \quad c_1 = 2H(\bar{F}_1^2 - \bar{F}_2^2).$$

Equivalently, the equation (3.25) can be written compactly by using both variables t and u, related by (3.20):

$$\frac{(F_2 + F_1 u)^3 \mathcal{W}^2}{32H^3} = (K_1 + 1)u^3 + (3\mu - 3\lambda + K_2)u^2 + (K_1 - 3)u \tag{3.26}$$
$$\qquad\qquad + \lambda - \mu + K_2 - t(F_2 + F_1 u)(u^2 - 1).$$

We can summarize the results of this section with

Theorem 9. *The Jacobian of the spectral curve S in (2.4) contains a 2-dimensional Prym variety, isomorphic to the Jacobian of the genus 2 curve C' given by (3.25) or (3.26).*

It is also worth mentioning the following relation between the roots of $R_6(t)$ and of the polynomial $P_6(u)$ defining the first genus 2 curve $C = S/\sigma$.

Proposition 10. *If $t = \hat{t}$ is a root of $R_6(t)$ then (3.20) gives a root of $P_6(u)$.*

Proof: For $t = \hat{t}$ we have $\mathcal{W} = 0$, which, in view of (3.24), implies $W = 0$ (provided that the denominator in (3.24) does not vanish for $t = \hat{t}$, and this condition always holds). In view of the definition $W = (w + \bar{w})/\sqrt{2}$, in this case $w = -\bar{w}$, which by (3.14) gives $g(u)\,v = 0$. Next, since $g(u), P_6(u)$ do not have common roots and deg $g(u) = 5$, the last equation defines 11 values of u, which, via (3.19), correspond to 11 zeros of the right hand side of (3.24). Further calculations show that $\mathcal{P}(t)\,(t + 1) = 0$ implies $g(u) = 0$. As a result, the 6 zeros of $\mathcal{W}$ correspond to the 6 zeros of v, i.e., the roots of $P_6(u)$. $\square$

4   Translation on $\mathrm{Prym}(S, \sigma)$ and on $\mathrm{Jac}(C')$

Below we represent the reduced Somos 6 map $\hat{\varphi} : \mathbb{C}^4 \to \mathbb{C}^4$ as a translation on the Jacobian of the spectral curve S, given by a divisor $\mathcal{V}$, and show that it belongs to $\mathrm{Prym}(S, \sigma)$. Then the translation will be described in terms of degree zero divisors on the curve C'.

First, recall that the Jacobian variety of an algebraic curve X can be defined as the additive group of degree zero divisors on X considered modulo divisors of meromorphic functions on X. Equivalence of divisors $\mathcal{D}_1, \mathcal{D}_2$ will be denoted as $\mathcal{D}_1 \equiv \mathcal{D}_2$.

Let $\mathcal{J}_K$ be the isospectral manifold, the set of the matrices $\mathbf{L}(x)$ of the form (2.3) having the same spectral curve S. Consider the *eigenvector map*

$$\mathcal{E} : \mathcal{J}_K \longrightarrow \mathrm{Jac}(S),$$

defined as follows: a matrix $\mathbf{L}(x) \in \mathcal{J}_K$ induces the eigenvector bundle $\mathbb{P}^2 \to S$; for any point $p = (x, y) \in S$

$$p \longrightarrow \psi(p) = (\psi_1(p), \psi_2(p), \psi_3(p))^T \quad \text{such that} \quad \mathbf{L}(x)\psi(p) = y\,\psi(p).$$

We assume that the eigenvector $\psi(p)$ is normalized: $\langle \alpha, \psi(p) \rangle = 1$, for a certain non-zero $\alpha \in \mathbb{C}^3$. This defines the divisor $\mathcal{D}$ of poles of $\psi(p)$ on S. For any choice of normalization, such divisors form an equivalence class $\{\mathcal{D}\}$. The latter defines a point $\mathcal{D} - Np_0 \in \mathrm{Jac}(S)$ with a certain base point $p_0 \in S$. Here $N = \mathrm{degree}(\mathcal{D})$, and for the case at hand $N = 6$. Then the composition $\mathbf{L}(x) \longrightarrow \mathcal{D} - 6p_0$ defines the above map $\mathcal{E}$. As was shown in e.g., [1, 3], the eigenvector map is injective.

Theorem 11. *Under the map $\mathcal{E}$, the transformation $\mathbf{L}(x) \to \tilde{\mathbf{L}}(x)$ defined by the discrete Lax representation (2.3) is the translation on $\mathrm{Jac}(S)$ given by the degree zero divisor*

$$\mathcal{V} = \bar{\mathcal{O}} - \mathcal{O},$$

where, as above, $\mathcal{O} = (x = 0, y = 0)$, $\bar{\mathcal{O}} = (x = \infty, y = \infty) \in S$.

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Proof: As follows from the intertwining relation (2.1), if $\psi(p)$ is a normalized eigenvector of $\mathbf{L}(x)$, then $\widehat{\psi}(p) = \mathbf{M}(x)\psi(p)$ is an eigenvector of $\tilde{\mathbf{L}}(x)$ with the same eigenvalue. Note that, in contrast to $\psi(p)$, $\widehat{\psi}(p)$ is not normalized, so we consider its normalization $\tilde{\psi}(p) = f^{-1}(p)\widehat{\psi}(p)$, $f = \langle \alpha, \widehat{\psi}(p) \rangle$ for a generic non-zero normalization vector $\alpha \in \mathbb{C}^3$.

Compare the divisors $\mathcal{D}, \tilde{\mathcal{D}}$ of poles of $\psi(p), \tilde{\psi}(p)$. Using (1.13) (or (7.11) in the appendix) we have

$$\det \mathbf{M}(x) = \frac{R_1 R_2}{R_0^2}\, x,$$

which implies that $\mathbf{M}(x), \mathbf{M}^{-1}(x)$ are non-degenerate apart from the points of S over $x = 0$ and $x = \infty$. Then $\mathcal{D}, \tilde{\mathcal{D}}$ can differ only by the points $\mathcal{O} = (0,0), \mathcal{O}_1 = (0, -\mu/\lambda)$ and $\bar{\mathcal{O}} = (\infty, \infty), \bar{\mathcal{O}}_1 = (\infty, -\lambda/\mu)$ or their multiples.

According to the structure of the matrix in (2.3), $\mathbf{M}(0)$ has a eigenvalue 0 with multiplicity 2, with a 1-dimensional eigenspace spanned by the vector $(1, 1, 0)^T$, and $\psi(\mathcal{O})$ is parallel to $(1, 1, 0)^T$, whereas $\psi(\mathcal{O}_1)$ is not. That is,

$$\widehat{\psi}(\mathcal{O}) = \mathbf{M}(0)\psi(\mathcal{O}) = 0, \quad \widehat{\psi}(\mathcal{O}_1) = \mathbf{M}(0)\psi(\mathcal{O}_1) \neq 0.$$

Further, for $\tau = \sqrt{x}$ being a local parameter on S near $\mathcal{O}$, we have the expansion $\psi(p) = (1 + O(\tau), 1 + O(\tau), O(\tau))^T$. Then near $\mathcal{O}$, $\widehat{\psi}(p) = (O(\tau), O(\tau), O(\tau))^T$, hence the normalizing factor $f = \langle \alpha, \widehat{\psi}(p) \rangle$ has a *simple* zero at $\mathcal{O}$ and does not vanish at $\mathcal{O}_1$.

Similarly, by considering the expansions of $\psi(p)$ near the points $\bar{\mathcal{O}}, \bar{\mathcal{O}}_1$ at infinity, one observes that $f(p)$ has a simple pole at $\bar{\mathcal{O}}$ and no poles at $\bar{\mathcal{O}}_1$. As a result,

$$(f) = \mathcal{O} + \mathcal{U} - \mathcal{D} - \bar{\mathcal{O}},$$

for a certain effective divisor $\mathcal{U}$. Then the divisor of poles of $\tilde{\psi}(p) = f^{-1}(p)\widehat{\psi}(p)$ equals $\mathcal{U}$. Indeed, the zeros of $\widehat{\psi}(p)$ and $f(p)$ at $\mathcal{O}$, as well as their poles at $\mathcal{D} + \bar{\mathcal{O}}$ cancel each other. Since $f(p)$ is meromorphic on S, we conclude that $\tilde{\mathcal{D}}$ is equivalent to $\mathcal{D} + \bar{\mathcal{O}} - \mathcal{O}$. Thus the images of $\mathcal{D}, \tilde{\mathcal{D}}$ in $\mathrm{Jac}(S)$ differ by the translation $\mathcal{V}$. □

Clearly $\sigma(\mathcal{V}) = -\mathcal{V}$, hence the divisor $\mathcal{V}$ represents a vector on $\mathrm{Prym}(S, \sigma)$. This leads us to the following

Theorem 12. *A generic complex invariant manifold $\mathcal{I}_K$ of the map $\hat{\varphi}$ is isomorphic to an open subset of the Jacobian of the genus 2 curve C' given by* (3.25) *or* (3.26).

Since Theorem 3 states that $\mathcal{I}_K$ is also an open subset of the Jacobian of the curve X in (1.3), we can conclude that the curves C' and X must be birationally equivalent.

Proof of Theorem 12: First note that in our case the manifolds $\mathcal{I}_K$ and $\mathcal{J}_K$ coincide. Then, since the map $\mathcal{E}$ is injective, $\mathcal{I}_K$ must be a subset of $\mathrm{Jac}(S)$. Theorem 11 specifies that $\mathcal{I}_K$ must be a subset of $\mathrm{Prym}(S, \sigma)$, which is isomorphic to $\mathrm{Jac}(C')$. As the dimensions of $\mathcal{I}_K$ and $\mathrm{Jac}(C')$ are equal, we arrive at the statement of the theorem. □

Now observe that for the special points $\mathcal{O}_1, \bar{\mathcal{O}}_1, \mathcal{O}_2, \bar{\mathcal{O}}_2 \in S$ described in Section 2, the degree zero divisors $\mathcal{V}_1 = \mathcal{O} - \mathcal{O}_1, \mathcal{V}_2 = \mathcal{O}_2 - \bar{\mathcal{O}}_2$ are also antisymmetric with respect to the involution σ, hence they represent vectors in $\mathrm{Prym}(S, \sigma)$ and, therefore, in $\mathrm{Jac}(C')$. Our next objective is to describe $\mathcal{V}, \mathcal{V}_1, \mathcal{V}_2$ in terms of degree zero divisors on C'.

Theorem 13. *Under the transformation of S to the canonical form $\tilde{C}$ given by (3.10), the points $\mathcal{O}, \bar{\mathcal{O}}$ become*

$$\mathcal{R}_1 = (u = 0, v = -1, w = 2), \quad \mathcal{R}_2 = (u = 0, v = -1, w = -2). \tag{4.1}$$

The points $\mathcal{O}_1, \bar{\mathcal{O}}_1 \in S$ become $\mathcal{R}'_1, \mathcal{R}'_2 \in \tilde{C}$, the two preimages of the infinite point $\infty_2 \in \{\infty_1, \infty_2\} \subset C$ specified by the Laurent expansions $u = 1/\tau$, $v = -(\lambda + \mu)/\tau^3 + O(\tau^{-2})$ with the local parameter $\tau = 1/u$. The expansions of w near $\mathcal{R}'_1, \mathcal{R}'_2$ are

$$w = \pm \frac{2(\lambda^2 - \mu^2)}{\tau^4} + O(\tau^{-3}),$$

respectively. Also, $\mathcal{O}_2, \bar{\mathcal{O}}_2$ become the points

$$\mathcal{R}''_{1,2} = \left(u = -\frac{1}{\lambda}, \; v = \frac{\varkappa}{\lambda^3}, \; w = \pm 2\mu \frac{\varkappa}{\lambda^4} \right), \quad \varkappa = \lambda^3 + K_2\lambda^2 - (K_1 + 1)\lambda + \mu. \tag{4.2}$$

Proof: Relations (3.4) describing the projection $\pi : S \to G$ imply that the coordinate Y has poles at $\pi(\mathcal{O}), \pi(\bar{\mathcal{O}}) \in G$. More precisely, in view of the behavior (2.8), these are triple poles. On the other hand, as follows from (3.4), (3.7), on the curve C the function $Y(u, v)$ has a triple pole only at $(u = 0, v = -1)$. (In particular, at $(u = 0, v = 1)$ this function has a pole of lower order, and it has no poles at the two points at infinity on C.) Hence, $\mathcal{O}, \bar{\mathcal{O}} \in S$ are projected to the same point $(u = 0, v = -1)$ on C. Since it is not a branch point of $\tilde{C} \to C$, the point has two preimages on $\tilde{C}$. To find their w-coordinates, we note that (3.7) implies

$$\lambda \mu u^3 (1 + \lambda u) \, Y(u, v) \Big|_{u=0, v=-1} = 2, \quad \lambda \mu u^3 (1 + \lambda u) \, Y(u, v) \Big|_{u=0, v=1} = 0, \tag{4.3}$$

and for $u = 0, v = -1$ the equation (3.10) gives $w^2 = 4$, thus we get (4.1). The proof of the rest of the theorem goes along similar lines. $\square$

Now recall [26] that for a double cover of curves $\pi : C_2 \to C_1$ with an involution $\sigma : C_2 \to C_2$, there are two natural maps between $\mathrm{Jac}(C_2)$ and $\mathrm{Jac}(C_1) \subset \mathrm{Jac}(C_2)$:

- the pullback $\pi^* : \mathrm{Jac}(C_1) \to \mathrm{Jac}(C_2)$;

  degree 0 divisor D on C_1 $\to$ degree 0 divisor $\tilde{D} = \pi^{-1}(D)$ on C_2;

- The projection (*Norm map*) $\mathrm{Nm}_{C_1} : \mathrm{Jac}(C_2) \to \mathrm{Jac}(C_1)$;

  degree 0 divisor $\tilde{D}$ on C_2 $\to$ degree 0 divisor $\pi(\tilde{D})$ on C_1.

Notice that for any degree zero divisor D on C_1, $\mathrm{Nm}_{C_1}(\pi^*(C_1)) = 2\,C_1$. This property should be understood on the level of equivalence classes of divisors, that is, for a degree zero divisor D on C_1, let $\tilde{D}'$ be any divisor on C_2 equivalent to $\tilde{D} = \pi^{-1}(D)$. Then $\pi(\tilde{D}') \equiv 2D$ on C_1. Then let $\mathcal{A}$, $\tilde{\mathcal{A}}$ be the Abel maps with images in $\mathrm{Jac}(C_1)$, $\mathrm{Jac}(C_2)$ respectively. Hence

$$\tilde{\mathcal{A}}(\tilde{D}) = \frac{1}{2}\tilde{\mathcal{A}}(\pi^* \circ \mathrm{Nm}_{C_1}(\tilde{D}')) \quad \text{and} \quad \mathcal{A}(\mathrm{Nm}_{C_1} \tilde{D}') = 2\,\mathcal{A}(D). \tag{4.4}$$

Now apply the above to the tower (3.13), introducing the map

$$\mathrm{Nm}_{C'} \; : \; \mathrm{Jac}(\widetilde{\widetilde{C}}) \to \mathrm{Jac}(C') \subset \mathrm{Jac}(\widetilde{\widetilde{C}})$$

as follows: for a degree zero divisor $\mathcal{Q}$ on $\widetilde{\widetilde{C}}$, $\mathrm{Nm}_{C'}(\mathcal{Q}) = \pi_1(\mathcal{Q})$. Next, consider the sequence of divisors

$$\mathcal{V} = \underbrace{\mathcal{R}_1 - \mathcal{R}_2}_{\in\,\mathrm{Prym}(\tilde{C},\sigma)} \xrightarrow{\;\tilde{\pi}^{-1}\;} \tilde{R}_1^+ + \tilde{R}_1^- - \tilde{R}_2^+ - \tilde{R}_2^- \xrightarrow{\;\pi_1\;} \mathcal{S} = \underbrace{S_1^+ + S_1^- - S_2^+ - S_2^-}_{\in\,\mathrm{Jac}(C')}, \quad (4.5)$$

where $\tilde{R}_j^{\pm}$ are the preimages of $\mathcal{R}_j$ on $\widetilde{\widetilde{C}}$.

Let $\mathcal{V}_0$ be a degree zero divisor on C' such that the pullbacks $\tilde{\pi}^{-1}(\mathcal{R}_1 - \mathcal{R}_2)$ and $\pi_1^{-1}(\mathcal{V}_0)$ give equivalent divisors on $\widetilde{\widetilde{C}}$. That is, the vector $\mathbf{w} = \mathcal{A}(\mathcal{V}_0) \in \mathrm{Jac}(C') \subset \mathrm{Jac}(\widetilde{\widetilde{C}})$ coincides with the Abel image of $\mathcal{R}_1 - \mathcal{R}_2$ in $\mathrm{Prym}(\tilde{C},\sigma) \subset \mathrm{Jac}(\widetilde{\widetilde{C}})$. Then, in view of the property (4.4),

$$\mathbf{w} = \frac{1}{2}\mathcal{A}(S_1^+ + S_1^- - S_2^+ - S_2^-). \qquad (4.6)$$

Proposition 14. *On the curve C' written in the hyperelliptic form (3.25) the divisor $\mathcal{V}_0$ is determined by*

$$S_1^+ = (t_*, \mathcal{W}_*^+), \quad S_1^- = (-t_*, \mathcal{W}_*^-), \quad S_2^+ = (t_*, -\mathcal{W}_*^+), \quad S_2^- = (-t_*, -\mathcal{W}_*^-),$$

$$t_* = \sqrt{\frac{\bar{F}_2}{F_2}}, \quad \mathcal{W}_*^{\pm} = 4\frac{H^{3/2}}{F_2}(1 \pm t_*) = 4\frac{H^{3/2}}{F_2^{3/2}}(\sqrt{F_2} \pm \sqrt{\bar{F}_2}).$$

The proof of the proposition is quite technical and is reserved for the appendix. Note that the squares of $\mathcal{W}_*^{\pm}$ given above coincide with the right hand side of (3.26) for $(u = 0, t = \pm t_*)$, as expected.

One now can observe that the divisor $\mathcal{S} = S_1^+ + S_1^- - S_2^+ - S_2^-$ is anti-invariant with respect to the hyperelliptic involution $\iota : (t, \mathcal{W}) \to (t, -\mathcal{W})$ on C'. Hence, modulo period vectors of $\mathrm{Jac}(C')$,

$$\mathcal{A}(S_1^+ - S_2^+) = 2\mathcal{A}(S_1^+ - P_0), \quad \mathcal{A}(S_1^- - S_2^-) = 2\mathcal{A}(S_1^- - P_0),$$

where P_0 is any Weierstrass point on C'. Then, using the relation (4.6), we arrive at

Theorem 15. *The reduced Somos 6 map $\hat{\varphi}$ is described by translation by the following vector on $\mathrm{Jac}(C')$:*

$$\mathbf{w} = \int_{P_0}^{(t_*, \mathcal{W}_*^+)} (\omega_1, \omega_2)^T + \int_{P_0}^{(-t_*, \mathcal{W}_*^-)} (\omega_1, \omega_2)^T = \int_{(-t_*, -\mathcal{W}_*^-)}^{(t_*, \mathcal{W}_*^+)} (\omega_1, \omega_2)^T , \qquad (4.7)$$

where $\omega_{1,2}$ are holomorphic differentials on C' and $t_*, \mathcal{W}_*^{\pm}$ are given in Proposition 14.

Since the curve C' is birationally equivalent to the genus 2 curve X in (1.3), the translation vector $\mathbf{v} \in \mathrm{Jac}(X)$ in the sigma-function solution (1.4) is the above vector $\mathbf{w}$ written in an appropriate basis of holomorphic differentials on C'.

It is convenient to describe also the degree zero divisors $\mathcal{V}_1 = \mathcal{O} - \mathcal{O}_1, \mathcal{V}_2 = \mathcal{O}_2 - \bar{\mathcal{O}}_2$ on S, or equivalently, $\mathcal{R}_1' - \mathcal{R}_2', \mathcal{R}_1'' - \mathcal{R}_2''$ on $\tilde{C}$, antisymmetric under σ, in terms of divisors on $\mathrm{Jac}(C')$. Namely, let $\mathcal{V}_{01}, \mathcal{V}_{02}$ be degree zero divisors on C' such that the vectors

$$\mathbf{w}_1 = \mathcal{A}(\mathcal{V}_{01}), \qquad \mathbf{w}_2 = \mathcal{A}(\mathcal{V}_{02}) \in \mathrm{Jac}(C') \subset \mathrm{Jac}(\widetilde{\tilde{C}})$$

coincide with the Abel images of $\mathcal{R}_1' - \mathcal{R}_2', \mathcal{R}_1'' - \mathcal{R}_2''$ respectively, in $\mathrm{Prym}(\tilde{C}, \sigma) \subset \mathrm{Jac}(\widetilde{\tilde{C}})$.

Theorem 16. *The divisors $\mathcal{V}_{01}, \mathcal{V}_{02}$ are equivalent to*

$$(t_*', \mathcal{W}_*^{+'}) - (-t_*', -\mathcal{W}_*^{-'}) \quad and \quad (t_*'', \mathcal{W}_*^{+''}) - (-t_*'', -\mathcal{W}_*^{-''}),$$

respectively, that is,

$$\mathbf{w}_1 = \int_{(-t_*', -\mathcal{W}_*^{-'})}^{(t_*', \mathcal{W}_*^{+'})} (\omega_1, \omega_2)^t, \quad \mathbf{w}_2 = \int_{(-t_*'', -\mathcal{W}_*^{-''})}^{(t_*'', \mathcal{W}_*^{+''})} (\omega_1, \omega_2)^t,$$

where

$$t_*' = \sqrt{\frac{-\bar{F}_1}{F_1}}, \quad \mathcal{W}_*^{\pm'} = 4\frac{H^{3/2}}{F_1}(1 \pm t_*') = 4\frac{H^{3/2}}{F_1^{3/2}}\left(\sqrt{F_1} \pm \sqrt{-\bar{F}_1}\right),$$

$$t_*'' = \sqrt{\frac{\lambda \bar{F}_2 + \bar{F}_1}{\lambda F_2 - F_1}}, \quad \mathcal{W}_*^{\pm''} = \frac{4H^{3/2}}{\lambda F_2 - F_1}\left(\lambda - 1 \pm (\lambda + 1)t_*''\right) \tag{4.8}$$

$$= \frac{4H^{3/2}\left((\lambda+1)\sqrt{\lambda F_2 - F_1} \pm (\lambda - 1)\sqrt{\lambda \bar{F}_2 + \bar{F}_1}\right)}{(\lambda F_2 - F_1)^{3/2}}.$$

The proof follows the same lines as that of Theorem 15; it uses Theorem 13 describing the coordinates of the pairs $\mathcal{R}_{1,2}', \mathcal{R}_{1,2}''$ on $\tilde{C}$. For each pair we construct a sequence of degree zero divisors analogous to (4.5), which gives rise to the divisors $\mathcal{V}_{01}, \mathcal{V}_{02}$ on C' and, in view of the relation (4.6), the vectors $\mathbf{w}_1, \mathbf{w}_2$.

5   The shift vector and the determinantal constraint

The above vectors $\mathbf{w}, \mathbf{w}_1, \mathbf{w}_2 \in \mathrm{Jac}(C')$ arising from the three special involutive pairs $\mathcal{O}, \bar{\mathcal{O}}, \mathcal{O}_1, \bar{\mathcal{O}}_1, \mathcal{O}_2, \bar{\mathcal{O}}_2$ on the spectral curve S have the following remarkable property.

Proposition 17. *The vectors in Theorem 15 and Theorem 16 are related by*

$$\mathbf{w}_1 = -2\mathbf{w}, \quad \mathbf{w}_2 = -3\mathbf{w}. \tag{5.1}$$

Proof: Following (2.8), the divisor of the meromorphic function y/x on S is

$$(y/x) = \mathcal{O} + \bar{\mathcal{O}}_1 + \mathcal{O}_2 - \bar{\mathcal{O}} - \mathcal{O}_1 - \bar{\mathcal{O}}_2 \equiv 0, \tag{5.2}$$

so it corresponds to zero in $\mathrm{Jac}(S)$, in $\mathrm{Prym}(S, \sigma) = \mathrm{Prym}(\tilde{C}, \sigma)$, and, therefore, in $\mathrm{Jac}(C')$. By Theorems 15 and 16, the degree zero divisors $\mathcal{O} - \bar{\mathcal{O}}, \mathcal{O}_1 - \bar{\mathcal{O}}_1, \mathcal{O}_2 - \bar{\mathcal{O}}_2$ with Abel image in $\mathrm{Prym}(S, \sigma)$ are represented, respectively, by the divisors $\mathcal{V}_0, \mathcal{V}_{01}, \mathcal{V}_{02}$ on C'. Then (5.2) implies $\mathcal{V}_0 - \mathcal{V}_{01} + \mathcal{V}_{02} \equiv 0$, which, under the Abel map, yields $\mathbf{w} - \mathbf{w}_1 + \mathbf{w}_2 = 0$. Similarly, we have

$$(y/x^2) = \bar{\mathcal{O}} + 2\bar{\mathcal{O}}_1 + \mathcal{O}_2 - \mathcal{O} - 2\mathcal{O}_1 - \bar{\mathcal{O}}_2 \equiv 0,$$

which implies $-\mathbf{w} - 2\mathbf{w}_1 + \mathbf{w}_2 = 0$. These two relations prove the proposition. $\qquad\square$

The divisors representing $\mathbf{w}_1, \mathbf{w}_2$ can also be derived (in a much more tedious way) by using addition formulae on $\mathrm{Jac}(C')$ described in terms of pairs of points on the hyperelliptic curve C'. These formulae can be obtained algorithmically using the Bäcklund transformation presented in [13, 21], which also allows us to calculate the divisor of the form $(t_+''', -\mathcal{W}_+''') - (t_-''', \mathcal{W}_-''')$ corresponding to the vector $4\mathbf{w}$. Since it will be needed in section 6, here we simply record that $t = t_\pm'''$ are found to be the roots of the quadratic

$$\left((\bar{F}_1 F_2 + F_1 \bar{F}_2)(\lambda+1) + 4F_1 \right) t^2 + 2\lambda\left(\bar{F}_1 F_2 + F_1 \bar{F}_2 \right) t + (\bar{F}_1 F_2 + F_1 \bar{F}_2)(\lambda-1) + 4\bar{F}_1. \tag{5.3}$$

The sigma-function, Kleinian functions, and the determinantal constraint.
Now let an appropriate Möbius transformation $t = \mathcal{T}(s)$ take the curve C' to a canonical odd order form X given by (1.3) (there are several possible transformations of this kind). Choosing the canonical basis of holomorphic differentials $\mathrm{d}s/z, s\,\mathrm{d}s/z$ on X, define the Abel map for a degree zero divisor $(s_1, z_1) + (s_2, z_2) - 2\infty$ by

$$\mathbf{u} = (u_1, u_2)^T = \int_\infty^{(s_1, z_1)} \left(\frac{ds}{z}, \frac{s\,ds}{z} \right)^T + \int_\infty^{(s_2, z_2)} \left(\frac{ds}{z}, \frac{s\,ds}{z} \right)^T \in \mathrm{Jac}(X).$$

It can be inverted by means of the the Bolza formulae

$$(s - s_1)(s - s_2) = s^2 - \wp_{22}(\mathbf{u})x - \wp_{12}(\mathbf{u}),$$
$$z_1 = \wp_{222}(\mathbf{u})s_1 + \wp_{122}(\mathbf{u}), \quad z_2 = \wp_{222}(\mathbf{u})s_2 + \wp_{122}(\mathbf{u}),$$

which involve $\wp_{jk\ell}(\mathbf{u}) = -\partial_j \partial_k \partial_\ell \log \sigma(\mathbf{u})$ in addition to the Kleinian hyperelliptic functions $\wp_{jk}(\mathbf{u})$. In particular, this yields

$$s_1 + s_2 = \wp_{22}(\mathbf{u}), \quad s_1 s_2 = -\wp_{12}(\mathbf{u}), \tag{5.4}$$

and for $\wp_{11}$ there is also Klein's formula

$$z_1 z_2 = \frac{1}{2} \sum_{k=0}^{2} (s_1 s_2)^k \left(2\bar{c}_{2k} + (s_1 + s_2)\bar{c}_{2k+1} \right) - 2(s_1 - s_2)^2 \wp_{11}(\mathbf{u}). \tag{5.5}$$

Now let $(\bar{s}_1, \bar{z}_1), (\bar{s}_2, \bar{z}_2) \in X$ be the images of the points $(t_*, \mathcal{W}_*^+), (-t_*, -\mathcal{W}_*^-) \in C'$ described in Proposition 14, and let $\{(\bar{s}_1', \bar{z}_1'), (\bar{s}_2', \bar{z}_2')\}$, $\{(\bar{s}_1'', \bar{z}_1''), (\bar{s}_2'', \bar{z}_2'')\} \in X \times X$ be the images of $\{(t_*', \mathcal{W}_*^{+'}), (-t_*', -\mathcal{W}_*^{-'})\}$, $\{(t_*'', \mathcal{W}_*^{+''}), (-t_*'', -\mathcal{W}_*^{-''})\} \in C' \times C'$, respectively, as specified in (4.8).

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Theorem 18. *The vector*

$$\mathbf{v} = \int_{(\bar{s}_2, \bar{z}_2)}^{(\bar{s}_1, \bar{z}_1)} \left(\frac{ds}{z}, \frac{s\,ds}{z} \right)^T \in \mathrm{Jac}(X) \tag{5.6}$$

satisfies the determinantal constraint (1.7).

Corollary 19. *For given values of K_1, K_2, λ, μ, take the associated genus 2 curve C' from (3.25), transform it into the canonical form X, as in (1.3), and pick the vector $\mathbf{v} \in \mathrm{Jac}(X)$ defined by (5.6). Then for the function $\sigma(\mathbf{u})$ associated with X, and for any $\mathbf{v}_0 \in \mathrm{Jac}(X)$ and $\mathrm{C} \in \mathbb{C}^*$, the expression (1.10) produces a sequence (x_n) satisfying the reduced Somos 6 recurrence (1.9) with coefficients given by (1.5) and first integrals K_1, K_2.*

Proof of Theorem 18: Observe that the vectors $\mathbf{v}$ and

$$\mathbf{v}_1 = \int_{(\bar{s}_2', \bar{W}_2)}^{(\bar{s}_1', \bar{W}_1)} \left(\frac{ds}{W}, \frac{s\,ds}{W} \right)^T, \quad \mathbf{v}_2 = \int_{(\bar{s}_2'', \bar{W}_2)}^{(\bar{s}_1'', \bar{W}_1)} \left(\frac{ds}{W}, \frac{s\,ds}{W} \right)^T$$

are just $\mathbf{w}, \mathbf{w}_1, \mathbf{w}_2$ written in the coordinates corresponding to the canonical differentials on X. Hence, the relations (5.1) imply $\mathbf{v}_1 = -2\mathbf{v}$, $\mathbf{v}_2 = -3\mathbf{v}$. Then, in view of the Bolza expressions (5.4), the determinant constraint (1.7) can be written as

$$\det \begin{pmatrix} 1 & 1 & 1 \\ -\bar{s}_1\bar{s}_2 & -\bar{s}_1'\bar{s}_2' & -\bar{s}_1''\bar{s}_2'' \\ \bar{s}_1 + \bar{s}_2 & \bar{s}_1' + \bar{s}_2' & \bar{s}_1'' + \bar{s}_2'' \end{pmatrix} = 0. \tag{5.7}$$

Next, we apply the inverse Möbius transformation $s = \mathcal{T}^{-1}(t)$ to each of the above s-coordinates, and observe that, after dividing out common denominators from each column, the rows of the resulting matrix are linear combinations of the rows of

$$\begin{pmatrix} 1 & 1 & 1 \\ (t_*)^2 & (t_*')^2 & (t_*'')^2 \\ t_* - t_* & t_*' - t_*' & t_*'' - t_*'' \end{pmatrix}. \tag{5.8}$$

As the third row of the latter matrix is zero, the condition (5.7) is trivially satisfied. $\square$

Remark 20. *To give an exact transformation from the equation of C' to the canonical form (1.3) for X one needs to know at least one root of the degree 6 polynomial $\mathrm{R}_6(t)$ in (3.25) (or, in view of Proposition 10, at least one root of $\mathrm{P}_6(u)$ in (3.5)). Yet in general it appears that the equation $\mathrm{R}_6(t) = 0$ is not solvable in radicals.*

6   Solution of the initial value problem

Before proceeding with the proof of Theorem 2, it is worth commenting on the meaning of the word "generic" appearing in its statement. Various non-generic situations arise:

- Some of the initial data or coefficients can be zero.

- For special values of λ, μ, K_1, K_2, the spectral curve S can acquire singularities (in addition to the singularity at $(0:1:0) \in \mathbb{P}^2$ for the projective curve).

- For special values of λ, μ, K_1, K_2, the curve C' becomes a product of two elliptic curves given by the factorization (3.17).

- For special values of λ, μ, K_1, K_2, one of the multiples $2\mathbf{v}, 3\mathbf{v}, 4\mathbf{v}$ of the shift vector can lie on the theta divisor $(\sigma) \subset \mathrm{Jac}(X)$.

A set of non-zero initial data $\tau_0, \ldots, \tau_5$ and coefficients α, β, γ determine the values of λ, μ, K_1, K_2, and these in turn determine the spectral curve S, the curve C' (hence X), and the vector $\mathbf{v} \in \mathrm{Jac}(X)$. Yet the sequence (τ_n) may contain zero terms; for instance, $\tau_0 = 0$ when $\mathbf{v}_0 = \mathbf{0}$ in (1.4). Iteration of the recurrence (1.1) requires non-zero initial data, but if an isolated zero appears in the sequence, then the Laurent phenomenon can be used to pass through this apparent singularity, by evaluating suitable Laurent polynomials in order to avoid division by zero.

Another degenerate possibility is that one of α, β, γ is zero, in which case (1.1) can be obtained as a reduction of the Hirota-Miwa (discrete KP) equation, and the solutions require a separate treatment in each case. For each of these three special cases there is also an associated cluster algebra, and by results in [14, 17] this means that a log-canonical symplectic structure is available for the reduced map (1.9) (see [20] for details).

If $\mathbf{v} \in (\sigma)$, so that $\sigma(\mathbf{v}) = 0$, then the expression (1.4) does not make sense; in that case one should replace $\sigma(\mathbf{v})$ by $\sigma_2(\mathbf{v}) = \partial_2 \sigma(\mathbf{v})$ in the denominator of the formula for τ_n, and then it satisfies a Somos-8 recurrence [5]. This situation is not relevant to our construction, since it can be checked directly that for λ, μ, K_1, K_2 such that (3.25) defines a curve of genus 2, neither S_1^+ nor S_1^- in Proposition 14 can be a Weierstrass point on C', hence $\mathbf{v} \notin (\sigma)$. However, it may happen that one of $2\mathbf{v}, 3\mathbf{v}$ or $4\mathbf{v} \in (\sigma)$, and in each case the formulae in Theorem 1 and/or our method for solving the initial value problem require certain adjustments. We illustrate this below in the case that $2\mathbf{v} \in (\sigma)$, which is needed for reconstruction of the original Somos-6 sequence (1.2).

Reconstruction of the constant C **and the initial phase** $\mathbf{v}_0$. The attentive reader might wonder why the constant C in (1.4) should be necessary to represent the general solution of the Somos-6 recurrence. Indeed, making the scaling $s \to \zeta^2 s$, $z \to \zeta^5 z$ in (1.3) changes the coefficients $\bar{c}_j$ but preserves the form of the curve X, and rescales the sigma-function so that $\sigma(\mathbf{u}) \to \zeta^{-3}\sigma(\mathbf{u})$, which means that C can always be set to 1. However, while the constraint (1.7) defines a curve in $\mathrm{Jac}(X)$, the discrete Lax pair considered above yields a unique point on this curve, so it is necessary to allow different values of C in our construction.

Now if a particular set of non-zero initial data $\tau_0, \ldots, \tau_5$ and coefficients α, β, γ are given, then the associated initial values x_0, x_1, x_2, x_3 for the reduced map $\hat{\varphi}$ are found from (1.8). The latter values can be used in the formulae for H_1, H_2 in [20], and from (2.7) these produce the values of K_1, K_2; alternatively, putting the x_j into (1.12) yields P_0, P_1, R_0, R_1, and then K_1, K_2 can be obtained directly from the expressions (2.5). The values of λ and μ are specified according to (1.14), and thus by Corollary 19 only C and $\mathbf{v}_0$ are needed to reconstruct the reduced sequence (x_n).

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Supposing that $\mathbf{v}, \mathbf{v}_0$ and C have already been found for a particular initial value problem, the parameters A and B are immediately obtained in the form

$$A = \frac{C\tau_0}{\sigma(\mathbf{v}_0)}, \qquad B = \frac{\sigma(\mathbf{v})\tau_1}{A\sigma(\mathbf{v}_0 + \mathbf{v})} = \frac{\sigma(\mathbf{v})\sigma(\mathbf{v}_0)\tau_1}{C\sigma(\mathbf{v}_0 + \mathbf{v})\tau_0}. \tag{6.1}$$

Thus the only outstanding problem is the determination of C and $\mathbf{v}_0$.

As an intermediate step, we introduce the sequence $(\tilde{\phi}_n)$ defined by

$$\tilde{\phi}_n = C^{n^2-1} \frac{\sigma(n\mathbf{v})}{\sigma(\mathbf{v})^{n^2}} \tag{6.2}$$

Apart from the powers of C, the latter is the same as Kanayama's phi-function introduced in [22] in genus 2, and considered for hyperelliptic curves of arbitrary genus in [30]. The sequence $(\tilde{\phi}_n)$ is a natural companion to (τ_n): it satisfies the same Somos-6 recurrence (1.1) and produces the same values of the first integrals K_1, K_2. In fact, it turns out that for each n, $\tilde{\phi}_n$ is an algebraic function of the quantities $\alpha, \beta, \gamma, H_1, H_2$. (The proof will be presented elsewhere.) For our current purposes, it is enough to consider only the first few terms of the sequence.

Lemma 21. *The terms of the sequence (6.2) for $n = 0, \ldots, 4$ are fixed up to signs by*

$$\tilde{\phi}_0 = 0, \quad \tilde{\phi}_1 = 1, \quad \tilde{\phi}_2^2 = \frac{\hat{\alpha}\beta}{\alpha\hat{\beta}}, \quad \tilde{\phi}_3^2 = \frac{\beta}{\hat{\beta}}, \quad \tilde{\phi}_4 = \tilde{\phi}_2^{-1}(\alpha\tilde{\phi}_3 - \gamma). \tag{6.3}$$

Proof: For $n = 0, 1$ the result is immediate from the definition. The formulae for $n = 2, 3$ follow from (1.5). To obtain $\tilde{\phi}_4$, set $n = -2$ in (1.1), replace τ_j with $\tilde{\phi}_j$ throughout, and use the fact that $\tilde{\phi}_{-j} = -\tilde{\phi}_j$ for all j. $\qquad \square$

Now from the coordinates of the points on C' considered in Proposition 14 and Theorem 16, and by applying the Möbius transformation $t = \mathcal{T}(s)$ followed by the Bolza formulae together with (5.5), as in the proof of Theorem 18, the values of the hyperelliptic functions $\wp_{jk}(m\mathbf{v})$ for $m = 1, 2, 3$ are all determined algebraically in terms of λ, μ, K_1, K_2. In turn, this allows $\hat{\alpha}, \hat{\beta}$ to be found from (1.6), which means that the expressions (6.3) determine $\tilde{\phi}_2$, $\tilde{\phi}_3$ and $\tilde{\phi}_4$, up to fixing the signs of $\tilde{\phi}_2$, $\tilde{\phi}_3$. Then, upon rearranging the formula for γ in (1.5), we see that

$$C^2 = \frac{\gamma}{\tilde{\phi}_3^2}\Big(\wp_{11}(3\mathbf{v}) - \hat{\alpha}\wp_{11}(2\mathbf{v}) - \hat{\beta}\wp_{11}(\mathbf{v})\Big)^{-1}, \tag{6.4}$$

which determines C up to a choice of sign; and this sign is irrelevant, since from (6.1) the prefactor $AB^n C^{n^2-1}$ in (1.4) is seen to be invariant under sending $C \to -C$. Once C is known, the sign of $\tilde{\phi}_3$ can then be fixed from an application of Baker's addition formula, which gives the identity $\tilde{\phi}_3 = \tilde{\phi}_2^2 \mathcal{F}(2\mathbf{v})$, where $\mathcal{F}$ is as in (1.11); the sign of $\tilde{\phi}_2$ will not be needed in what follows: it corresponds to the overall freedom to send $\mathbf{v} \to -\mathbf{v}, \mathbf{v}_0 \to -\mathbf{v}_0$ in the solution, which is removed once the signs of the z coordinates of the points in (5.6) are fixed.

Finally, given the six non-zero initial data, one should iterate the Somos-6 recurrence forwards/backwards to obtain additional terms, in order to calculate ratios of the form

$\tau_j \tau_{-j} / \tau_0^2$ for $j = 1, 2, 3, 4$. (By adjusting the offset of the index if necessary, a maximum of three iterations are needed to obtain 9 adjacent terms $\tau_{-4}, \tau_{-3}, \ldots, \tau_4$ with generic initial data.) Matching these ratios with the analytic formula (1.4), and using Baker's addition formula, yields four linear equations for the quantities $\wp_{jk}(\mathbf{v}_0)$, namely

$$\frac{\tau_j \tau_{-j}}{\tau_0^2} = C^2 \tilde{\phi}_j^2 \left(\wp_{12}(j\mathbf{v}) \wp_{22}(\mathbf{v}_0) - \wp_{22}(j\mathbf{v}) \wp_{12}(\mathbf{v}_0) + \wp_{11}(j\mathbf{v}) - \wp_{11}(\mathbf{v}_0) \right) \qquad (6.5)$$

for $j = 1, 2, 3, 4$, Now, observing that the first three equations are linearly dependent, due to the constraint (1.7), it is necessary to use any two of the first three together with the fourth; for instance, picking $j = 1, 3, 4$ produces the 3×3 matrix equation

$$\begin{pmatrix} \wp_{22}(\mathbf{v}) & \wp_{12}(\mathbf{v}) & 1 \\ \wp_{22}(3\mathbf{v}) & \wp_{12}(3\mathbf{v}) & 1 \\ \wp_{22}(4\mathbf{v}) & \wp_{12}(4\mathbf{v}) & 1 \end{pmatrix} \begin{pmatrix} -\wp_{12}(\mathbf{v}_0) \\ \wp_{22}(\mathbf{v}_0) \\ -\wp_{11}(\mathbf{v}_0) \end{pmatrix} = \begin{pmatrix} C^{-2}\rho_1 - \wp_{11}(\mathbf{v}) \\ C^{-2}\rho_3 - \wp_{11}(3\mathbf{v}) \\ C^{-2}\rho_4 - \wp_{11}(4\mathbf{v}) \end{pmatrix}, \qquad (6.6)$$

where we set

$$\rho_j = \tau_j \tau_{-j} / (\tilde{\phi}_j \tau_0)^2.$$

In order to make this formula effective, the values of $\wp_{jk}(4\mathbf{v})$ are required; these can be found by taking the roots of (5.3) and transforming them with $s = \mathcal{T}^{-1}(t)$ to the corresponding s-coordinates on X, or by directly applying the Bäcklund transformation for the genus 2 odd Mumford system [13, 21] to perform the addition $3\mathbf{v} + \mathbf{v} = 4\mathbf{v}$ on $\mathrm{Jac}(X)$. Upon solving (6.6), the quantities $\wp_{jk}(\mathbf{v}_0)$ are found, so that $\mathbf{v}_0 \in \mathrm{Jac}(X)$ is

$$\mathbf{v}_0 = \int_{\infty}^{(s_1^{(0)}, z_1^{(0)})} \left(\frac{\mathrm{d}s}{z}, \frac{s\,\mathrm{d}s}{z} \right)^T + \int_{\infty}^{(s_2^{(0)}, z_2^{(0)})} \left(\frac{\mathrm{d}s}{z}, \frac{s\,\mathrm{d}s}{z} \right)^T, \qquad (6.7)$$

corresponding to the Abel map for the divisor $\mathcal{D}_0 = (s_1^{(0)}, z_1^{(0)}) + (s_2^{(0)}, z_2^{(0)}) - 2\infty$ on X, where the coordinates of the points $(s_1^{(0)}, z_1^{(0)}), (s_2^{(0)}, z_2^{(0)})$ are obtained by using (5.4) and (5.5) with $\mathbf{u} = \mathbf{v}_0$. (An overall choice of sign for $z_j^{(0)}$ is left undetermined; this can be fixed by doing a single iteration, taking $\mathbf{v}_0 \to \mathbf{v}_0 + \mathbf{v}$ and checking the result.) This completes the proof of Theorem 2. $\qquad \square$

A numerical example. For illustration of the main result, we consider the following choice of initial data and coefficients:

$$(\tau_0, \ldots, \tau_5) = (1, 1, -1, 1, -3, -3), \qquad \alpha = 1, \qquad \beta = 2, \qquad \gamma = -2.$$

This produces an integer sequence which extends both backwards and forwards,

$$\ldots, 1, -1, 1, 1, 1, 1, -1, 1, -3, -3, 1, -25, 49, 1, 385, 1489, 503, 10753, -82371, \ldots, \quad (6.8)$$

so that it has the symmetry $\tau_n = \tau_{-n-1}$. The corresponding initial data for (1.9) are $(x_0, x_1, x_2, x_3) = (-1, 1, 3, -1/3)$, and so the first integrals presented in [20] take the values $H_1 = -12, H_2 = 8$. Then (fixing a choice of square root) $\delta_1 = 1/2$, which gives

$$\lambda = 2, \qquad \mu = -1, \qquad K_1 = -6, \qquad K_2 = -4. \qquad (6.9)$$

After rescaling $\mathcal{W}$ suitably, the curve C' is found from (3.25) to be

$$C': \qquad \mathcal{W}^2 = (t^2 - 1)(19t^4 + 16t^3 + 2t^2 - 80t - 37). \qquad (6.10)$$

With the Möbius transformation $t = \mathcal{T}(s) = (2s - 5)/(2s + 5)$ this is transformed to

$$X: \qquad z^2 = 4s^5 + 52s^4 + 35s^3 + 25s^2 - \frac{375}{4}s. \qquad (6.11)$$

m	$\wp_{12}(m\mathbf{v})$	$\wp_{22}(m\mathbf{v})$	$\wp_{11}(m\mathbf{v})$
1	$-25/4$	$-5/2$	$-125/8$
2	$-25/4$	$-25/2$	$-25/8$
3	$-25/4$	$15/2$	$575/8$
4	$-25/36$	$-5/18$	$475/72$

Table 6.1: Values of Kleinian functions at multiples of $\mathbf{v}$ for sequence (6.8).

Now to obtain the vector $\mathbf{v} \in \mathrm{Jac}(X)$ as in Theorem 18, we start from the points on C' given in Proposition 14, and applying the inverse Möbius transformation $s = \mathcal{T}^{-1}(t)$ to find the degree zero divisor $\mathcal{D}$ on X corresponding to $\mathbf{v}$. This produces

$$\mathcal{D} = (s_1, z_1) + (s_2, z_2) - 2\infty, \qquad \mathbf{v} = \int_\infty^{(s_1, z_1)} \left(\frac{\mathrm{d}s}{z}, \frac{s\,\mathrm{d}s}{z} \right)^T + \int_\infty^{(s_2, z_2)} \left(\frac{\mathrm{d}s}{z}, \frac{s\,\mathrm{d}s}{z} \right)^T,$$

$$\text{(6.12)}$$

$$\text{where} \qquad s_{1,2} = -\frac{5}{2} e^{\mp i\pi/3}, \qquad z_{1,2} = 25\sqrt{2} e^{\pm i\pi/3}.$$

(Note that we slightly changed notation here compared with (5.6); in particular, we dropped bars on the coordinates.) Similarly, by applying the same (inverse) Möbius transformation to the coordinates given in (4.8) and (5.3), or via the Bäcklund transformation in [13, 21], we find the divisors corresponding to $2\mathbf{v}, 3\mathbf{v}, 4\mathbf{v} \in \mathrm{Jac}(X)$, namely $\mathcal{D}' = (s_1', z_1') + (s_2', z_2') - 2\infty, \mathcal{D}'' = (s_1'', z_1'') + (s_2'', z_2'') - 2\infty, \mathcal{D}''' = (s_1''', z_1''') + (s_2''', z_2''') - 2\infty$, where

$$s_{1,2}' = -\frac{5}{4}(5 \pm \sqrt{21}), \quad z_{1,2}' = \frac{25}{2}\sqrt{2}(5 \pm \sqrt{21}), \qquad s_{1,2}'' = \frac{5}{4}(3 \pm \sqrt{5}), \quad z_{1,2}'' = \frac{25}{2}\sqrt{2}(11 \pm 5\sqrt{5}),$$

$$s_{1,2}''' = \frac{5}{36}(-1 \pm i\sqrt{35}), \qquad z_{1,2}''' = \frac{25}{486}\sqrt{2}(103 \mp 13i\sqrt{35}).$$

Using the Bolza formulae (5.4) and (5.5), this allows us to calculate the values of the Kleinian functions $\wp_{jk}(m\mathbf{v})$ for $m = 1, 2, 3, 4$, as presented in Table 6.1.

The values in the latter table, together with (1.6), yield $\hat{\alpha} = -1$, $\hat{\beta} = 2$, so from (6.3) we see that $\tilde{\phi}_2^2 = -1$, $\tilde{\phi}_3^2 = 1$ and $\tilde{\phi}_4 = \tilde{\phi}_2^{-1}(\tilde{\phi}_3 + 2)$. From (6.4), this is enough to determine that $C^2 = -1/50$, and then from $\tilde{\phi}_3 = \tilde{\phi}_2^2 \mathcal{F}(2\mathbf{v})$ we find

$$\tilde{\phi}_2 = \pm i, \qquad \tilde{\phi}_3 = 1, \qquad \tilde{\phi}_4 = \mp 3i.$$

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Then we have $\rho_1 = 1, \rho_3 = -1, \rho_4 = 1/3$, which means that the linear system (6.6) can be solved for $\wp_{jk}(\mathbf{v}_0)$, to yield

$$\mathrm{C} = \frac{i}{\sqrt{50}}, \qquad \wp_{12}(\mathbf{v}_0) = -5/4, \quad \wp_{22}(\mathbf{v}_0) = 3/2, \quad \wp_{11}(\mathbf{v}_0) = 175/8$$

(where we have recorded a particular choice of sign for C). Hence, after fixing signs of the z-coordinates appropriately, the coordinates of the points in the divisor $\mathcal{D}_0$ are

$$s_{1,2}^{(0)} = \frac{1}{4}(3 \pm i\sqrt{11}), \qquad z_{1,2}^{(0)} = \sqrt{2}(1 \mp 3i\sqrt{11}),$$

so that $\mathbf{v}_0$ is given by (6.7). Finally, with these values of $\mathbf{v}_0$ and $\mathbf{v}$, the constants A, B are found from (6.1) to be $\mathrm{A} = i/(\sqrt{50}\,\sigma(\mathbf{v}_0))$, $\mathrm{B} = -i\sqrt{50}\,\sigma(\mathbf{v}_0)\sigma(\mathbf{v})/\sigma(\mathbf{v}_0 + \mathbf{v})$.

The special case where $2\mathbf{v} \in (\sigma)$. In order to illustrate the modifications that are needed in a degenerate case, we briefly consider the situation where $2\mathbf{v}$ lies on the theta divisor. This corresponds to having

$$2\mathbf{v} = \int_\infty^{(s',z')} \left(\frac{\mathrm{d}s}{z}, \frac{s\,\mathrm{d}s}{z} \right)^T, \tag{6.13}$$

the image of a single point (s', z') under the Abel map based at infinity. The formula (1.4) still makes sense, but (since $\wp_{jk}(2\mathbf{v})$ become singular) the expressions (1.5) for the coefficients are no longer appropriate, and Theorem 1 requires a slight reformulation.

Theorem 22. *For $\mathbf{v} \in \mathrm{Jac}(X)$ such that $2\mathbf{v}$ has the form (6.13) modulo periods, with arbitrary $\mathrm{A}, \mathrm{B}, \mathrm{C} \in \mathbb{C}^*$, $\mathbf{v}_0 \in \mathbb{C}^2$, the sequence with nth term (1.4) satisfies the recurrence (1.1) with coefficients given by*

$$\alpha = \frac{\tilde{\phi}_3^2}{\hat{\phi}_2^2}\,\hat{\alpha}, \qquad \beta = \tilde{\phi}_3^2, \qquad \gamma = \tilde{\phi}_3^2\,\mathrm{C}^2\Big(\wp_{11}(3\mathbf{v}) + \hat{\alpha}(s')^2 - \wp_{11}(\mathbf{v})\Big), \tag{6.14}$$

where $\hat{\phi}_2 = \mathrm{C}^3\,\sigma_2(2\mathbf{v})/\sigma(\mathbf{v})^4$, $\hat{\alpha} = \wp_{22}(\mathbf{v}) - \wp_{22}(3\mathbf{v})$, provided that $\mathbf{v}$ satisfies the constraint

$$\det \begin{pmatrix} 1 & 0 & 1 \\ \wp_{12}(\mathbf{v}) & -s' & \wp_{12}(3\mathbf{v}) \\ \wp_{22}(\mathbf{v}) & 1 & \wp_{22}(3\mathbf{v}) \end{pmatrix} = 0.$$

The coefficients satisfy the condition

$$\alpha^2 \beta = \gamma^2. \tag{6.15}$$

Proof: The main formulae above arise from Theorem 1 by taking the limit $\sigma(2\mathbf{v}) \to 0$ with $\sigma_2(2\mathbf{v}) \neq 0$, or directly by using Baker's addition formula and its limiting case for a shift on the theta divisor [5]. For the necessary condition (6.15) note that by (1.1) for $n = -2$ with $\tau_j = \tilde{\phi}_j$ for all j and $\tilde{\phi}_2 = 0$, the identity $\alpha\tilde{\phi}_3 = \gamma$ holds; squaring both sides of the latter and comparing with β in (6.14) yields the condition. $\qquad\square$

For the purpose of the reconstruction problem, we need an additional formula, namely

$$\tilde{\phi}_4^2 = \alpha^3 + H_1. \tag{6.16}$$

Its proof is based on the fact that the companion sequence also satisfies a Somos-10 recurrence (see Proposition 2.5 in [20]), but we omit further details.

The original Somos-6 sequence. For the original sequence (1.2) considered by Somos, we choose to index the terms so that

$$(\tau_0, \ldots, \tau_5) = (1, 1, 1, 3, 5, 9), \qquad \alpha = 1, \qquad \beta = 1, \qquad \gamma = 1.$$

We have $H_1 = 19, H_2 = 14$, as noted in [20]. Then (upon fixing a sign) $\delta_1 = i$, giving

$$\lambda = i, \qquad \mu = -i, \qquad K_1 = 19, \qquad K_2 = 14i. \tag{6.17}$$

The curve C', found from (3.25), takes the form

$$C' : \qquad \mathcal{W}^2 = (t - 1) \, \mathrm{Q}(t), \tag{6.18}$$

where, after removal of a numerical prefactor, $\mathrm{Q}(t) = 159025t^5 + \ldots + 154607 + 37224i$ is a quintic polynomial with Gaussian integer coefficients whose real and imaginary parts have 5 or 6 digits. The Möbius transformation $t = \mathcal{T}(s) = (s - i)/(s + i)$ sends the root $t = 1$ to infinity, and transforms C' to the canonical quintic curve

$$X : \qquad z^2 = 4s^5 - 233s^4 + 1624s^3 - 422s^2 + 36s - 1. \tag{6.19}$$

j	$\wp_{12}(j\mathbf{v})$	$\wp_{22}(j\mathbf{v})$	$\wp_{11}(j\mathbf{v})$
1	1	-16	$51/2$
2	∞	∞	∞
3	1	-36	$11/2$
4	0	0	$49/2$

Table 6.2: Values of Kleinian functions at multiples of $\mathbf{v}$ for sequence (1.2).

As in the previous example, by rewriting the points in Proposition 14 in terms of points in X, we obtain the divisor $\mathcal{D}$ and corresponding vector $\mathbf{v} \in \mathrm{Jac}(X)$ in the form (6.12), where

$$s_{1,2} = -8 \pm \sqrt{65}, \qquad z_{1,2} = 20i(129 \mp 16\sqrt{65}).$$

The condition (6.15) clearly holds, but it is necessary, not sufficient for $2\mathbf{v} \in (\sigma)$. However, from the first formula in (4.8) we find that $(t'_*)^2 = 1$, meaning that one of the points in the divisor $\mathcal{V}_{01}$ is the Weierstrass point $(1, 0) \in C'$, and under the Möbius transformation this means that $2\mathbf{v}$ has the form (6.13) with $(s', z') = (0, -i)$, corresponding to the divisor

$$\mathcal{D}' = (0, -i) - \infty$$

on X. Application of the formula for t''_* in (4.8) leads to the divisor $\mathcal{D}'' = (s''_1, z''_1) + (s''_2, z''_2) - 2\infty$ corresponding to $3\mathbf{v}$, where

$$s''_{1,2} = -18 \pm 5\sqrt{13}, \qquad z''_{1,2} = 20i(-667 \pm 185\sqrt{13}),$$

while for $4\mathbf{v}$ we have $\mathcal{D}''' = 2\mathcal{D}'$. This means that the finite values of $\wp_{jk}(m\mathbf{v})$ for $m = 1, 3, 4$, as presented in Table 6.2, can be obtained in the usual way, except that (5.5) is no longer valid for $\mathbf{u} = 4\mathbf{v}$. Instead, in order to find $\wp_{11}(4\mathbf{v})$, we use the equation

of the Kummer surface (see e.g. [2]), which provides a quartic relation between the functions $\wp_{jk}(\mathbf{u})$.

From the identities in Theorem 22 and its proof we see that $\hat{\alpha} = 20$ and $\tilde{\phi}_3 = 1$, giving $\hat{\phi}_2^2 = 20$, $C^2 = -1/20$ from the first and last formulae in (6.14), and also $\hat{\phi}_4^2 = 20$ by (6.16). The equation (6.5) should be modified when $j = 2$, but is valid for $j = 1, 3, 4$, which means that we can still solve (6.6) with $\rho_1 = 1, \rho_3 = 3, \rho_4 = 3/4$, and (fixing the sign of C) we find

$$C = \frac{i}{\sqrt{20}}, \qquad \wp_{12}(\mathbf{v}_0) = -1, \quad \wp_{22}(\mathbf{v}_0) = 10, \quad \wp_{11}(\mathbf{v}_0) = 79/2.$$

Hence $\mathbf{v}_0$ is given by (6.7), where the associated divisor $\mathcal{D}_0$ contains the coordinates

$$s_{1,2}^{(0)} = 5 \pm 2\sqrt{6}, \qquad z_{1,2}^{(0)} = 4i(71 \pm 29\sqrt{6}).$$

Thus, from (6.1), we see that $A = i/(\sqrt{20}\,\sigma(\mathbf{v}_0))$, $B = -i\sqrt{20}\,\sigma(\mathbf{v}_0)\sigma(\mathbf{v})/\sigma(\mathbf{v}_0 + \mathbf{v})$.

7   Conclusion

The explicit solution (1.4) is equivalent to an expression $\tau_n = A\,B^n\,C^{n^2}\,\Theta(\mathbf{z}_0 + n\mathbf{z})$ in terms of a Riemann theta function in two variables, for suitable constants A, B, C and vectors $\mathbf{z}_0, \mathbf{z} \in \mathbb{C}^2$. In fact (see http://somos.crg4.com/somos6.html), a numerical fit of (1.2) with a two-variable Fourier series was performed by Somos some time ago.

There are a number of aspects of the solution that we intend to consider in more detail elsewhere. The companion sequence (6.2) deserves more attention, since its properties should be helpful in proving that other Somos-6 sequences consist entirely of integers, in cases where the Laurent property is insufficient; for example, take

$$(\tau_0, \ldots, \tau_5) = (2, 3, 6, 18, 54, 108), \qquad \alpha = 18, \qquad \beta = 36, \qquad \gamma = 108,$$

which defines a sequence belonging to an infinite family found by Melanie de Boeck. We also propose to examine Somos-6 sequences that are parametrized by elliptic functions, including the case where the factorization (3.17) holds.

The solution of the initial value problem for (1.1) raises the further possibility of performing separation of variables for the reduced map $\hat{\varphi}$ defined by (1.9), and finding a 2×2 matrix Lax representation for it. In principle, such a representation might also be used to obtain a symplectic structure for $\hat{\varphi}$ when $\alpha\beta\gamma \neq 0$, which could shed some light on the open problem of finding compatible Poisson or (pre)symplectic structures for general Laurent phenomenon algebras (see [23]).

Appendix A. Derivation of the Lax pair

The general Somos-6 recurrence arises by reduction from the discrete BKP equation, which is a given by a bilinear relation for a tau function $T(n_1, n_2, n_3)$ that depends on three independent variables. For convenience we use indices to write $T(n_1, n_2, n_3) = T_{jk\ell}$, $(n_1, n_2, n_3) = (j, k, \ell)$, so that the discrete BKP equation takes the form

$$
\begin{aligned}
T_{j+1,k+1,\ell+1}T_{jk\ell} \;-\;& T_{j+1,k,\ell}T_{j,k+1,\ell+1} \\
+\;& T_{j,k+1,\ell}T_{j+1,k,\ell+1} - T_{j,k,\ell+1}T_{j+1,k+1,\ell} = 0.
\end{aligned}
\tag{7.1}
$$

Following [9], this equation arises as a compatibility conditon of the linear triad

$$
\begin{aligned}
\Psi_{j+1,k+1,\ell} - \Psi_{jk\ell} &= \frac{T_{j+1,k,\ell}T_{j,k+1,\ell}}{T_{j+1,k+1,\ell}T_{jk\ell}}\Big(\Psi_{j+1,k,\ell} - \Psi_{j,k+1,\ell}\Big), \\
\Psi_{j,k+1,\ell+1} - \Psi_{jk\ell} &= \frac{T_{j,k+1,\ell}T_{j,k,\ell+1}}{T_{j,k+1,\ell+1}T_{jk\ell}}\Big(\Psi_{j,k+1,\ell} - \Psi_{j,k,\ell+1}\Big), \\
\Psi_{j+1,k,\ell+1} - \Psi_{jk\ell} &= \frac{T_{j+1,k,\ell}T_{j,k,\ell+1}}{T_{j+1,k,\ell+1}T_{jk\ell}}\Big(\Psi_{j+1,k,\ell} - \Psi_{j,k,\ell+1}\Big).
\end{aligned}
\tag{7.2}
$$

Now construct a tau function which, apart from a gauge transformation by the exponential of a quadratic form, depends only on the single independent variable

$$
n = n_1 + 2n_2 + 3n_3 = j + 2k + 3\ell;
\tag{7.3}
$$

so we set

$$
T_{jk\ell} = \delta_1^{k\ell}\delta_2^{j\ell}\delta_3^{jk}\,\tau_n
\tag{7.4}
$$

for some parameters δ_j, $j = 1, 2, 3$. Substituting this into (7.1) produces the Somos-6 recurrence for τ_n with the coefficients

$$
\alpha = \frac{1}{\delta_2\delta_3}, \quad \beta = -\frac{1}{\delta_1\delta_3}, \quad \gamma = \frac{1}{\delta_1\delta_2}.
\tag{7.5}
$$

The derivation of the Lax pair for Somos-6 is somewhat more involved, but proceeds by applying an analogous reduction procedure to the system (7.2). We suppose that the wave function also depends primarily on the same dependent variable n in (7.3), apart from a gauge factor, taking the form $\Psi_{jk\ell} = x^{-k}y^{-\ell}\psi_n$, where x and y are spectral parameters. By imposing this form for the wave function, together with (7.4), we find that (7.1) gives a scalar system for the reduced wave function ψ_n, namely

$$
\begin{aligned}
\psi_{n+3} &= -P_n\,\psi_{n+2} + x\,P_n\,\psi_{n+1} + x\,\psi_n, \\
\psi_{n+5} &= -x\,Q_n\,\psi_{n+3} + y\,Q_n\,\psi_{n+2} + xy\,\psi_n, \\
\psi_{n+4} &= -R_n\,\psi_{n+3} + y\,R_n\,\psi_{n+1} + y\,\psi_n,
\end{aligned}
\tag{7.6}
$$

where we have introduced new dependent variables

$$
P_n = \frac{1}{\delta_3}\,\frac{\tau_{n+1}\tau_{n+2}}{\tau_n\tau_{n+3}}, \quad Q_n = \frac{1}{\delta_1}\,\frac{\tau_{n+2}\tau_{n+3}}{\tau_n\tau_{n+5}}, \quad R_n = \frac{1}{\delta_2}\,\frac{\tau_{n+1}\tau_{n+3}}{\tau_n\tau_{n+4}}.
\tag{7.7}
$$

After identifying the prefactors from (7.5), it is clear that the above formulae for P_n and R_n are the same as (1.12) rewritten in terms of tau functions. Using the first equation

to eliminate ψ_{n+3}, the third and second equations in (7.6) provide expressions for ψ_{n+4} and ψ_{n+5}, respectively, as linear combinations of ψ_n, ψ_{n+1} and ψ_{n+2}, that is

$$
\begin{aligned}
\psi_{n+4} &= P_n R_n \psi_{n+2} - x P_n R_n \psi_{n+1} - x R_n \psi_n + y \phi_n^{(1)}, \\
\psi_{n+5} &= x P_n Q_n \psi_{n+2} - x^2 P_n Q_n \psi_{n+1} - x^2 Q_n \psi_n + y \phi_n^{(2)},
\end{aligned}
\tag{7.8}
$$

where in each case we have isolated the coefficient of y as

$$
\phi_n^{(1)} = R_n \psi_{n+1} + \psi_n, \qquad \phi_n^{(2)} = Q_n \psi_{n+2} + x \psi_n.
$$

Now by shifting $n \to n+1$ and $n \to n+2$ in the first equation of (7.6), we obtain alternative expressions for ψ_{n+4} and ψ_{n+5} as linear combinations of ψ_n, ψ_{n+1} and ψ_{n+2}, which combine with (7.8) to yield a pair of linear equations of the form

$$
L^{(1)}(\psi_n, \psi_{n+1}, \psi_{n+2}) = y \phi_n^{(1)}, \quad L^{(2)}(\psi_n, \psi_{n+1}, \psi_{n+2}) = y \phi_n^{(2)}.
\tag{7.9}
$$

Next, setting $\phi_n^{(0)} = \psi_n$, we have $\phi_{n+1}^{(0)} = \psi_{n+1} = (\phi_n^{(1)} - \phi_n^{(0)})/R_n$, and similar equations for the shifts $\phi_{n+1}^{(1)}$ and $\phi_{n+1}^{(2)}$ and so, using a tilde to denote the shift $n \to n+1$, this produces the matrix equation

$$
\tilde{\boldsymbol{\Phi}} = \mathbf{M}(x)\,\boldsymbol{\Phi}, \qquad \boldsymbol{\Phi} = (\phi_n^{(0)}, \phi_n^{(1)}, \phi_n^{(2)})^T,
\tag{7.10}
$$

where $\mathbf{M}$ (for $n=0$) is given by (2.3), including the parameter $\lambda = Q_n/(R_n R_{n+1}) = \delta_2^2/\delta_1 = \delta_1 \beta^2/\alpha^2$, as in (1.14). To obtain the simplest-looking version of $\mathbf{M}$ we used

$$
(\lambda P_n R_{n+1} - \lambda R_n R_{n+1} - P_n) R_{n+2} + 1 = 0,
\tag{7.11}
$$

which is equivalent to the Somos-6 recurrence (1.1) for τ_n.

The system (7.9) is incomplete, because it lacks an equation for $y \phi_n^{(0)}$, but applying the shift $n \to n+1$ to the first equation of this pair, and using (7.10), we obtain the missing relation. This results in the eigenvalue equation

$$
\mathbf{L}(x)\,\boldsymbol{\Phi} = y\,\boldsymbol{\Phi},
\tag{7.12}
$$

with the Lax matrix taking the form (2.2). Upon using (7.11) and introducing the additional parameter $\mu = P_n P_{n+1} P_{n+2}/(R_n R_{n+1}) = \delta_2^3/\delta_3^3 = -\delta_1 \beta^3/\gamma^2$, as in (1.14), the coefficients $A_2, \ldots, C_0''$ can be written in terms of $P_n, P_{n+1}, R_n, R_{n+1}$ and the constants λ, μ, and for $n=0$ the resulting expressions coincide with those in Theorem 4.

Equations (7.10) and (7.12) form a linear system for $\boldsymbol{\Phi}$, whose compatibility condition is the discrete Lax equation (2.1), or equivalently $\tilde{\mathbf{L}} = \mathbf{M}\mathbf{L}\mathbf{M}^{-1}$, meaning that the shift $n \to n+1$ is an isospectral evolution. This explains the origin of Theorem 4.

Appendix B. Proof of Proposition 14

By (3.14) and (4.1), $\bar{w}^2 = h(u) - v g(u) = h(0) + g(0) = 0$ holds at $\mathcal{R}_1, \mathcal{R}_2$, hence

$$
\tilde{\pi}^{-1}(\mathcal{R}_1) = (u = 0, v = -1, w = 2, \bar{w} = 0), \quad \tilde{\pi}^{-1}(\mathcal{R}_2) = (u = 0, v = -1, w = -2, \bar{w} = 0).
$$

However, observe that on $\tilde{C}$ the function $\bar{w}^2 = h(u) - g(u)v$ has zeros of order 6 at $\mathcal{R}_1, \mathcal{R}_2$. Hence $\widetilde{\widetilde{C}}$ is singular at the above 2 points. To regularize it, observe that near $u = 0, v = -1$ the coordinate $\bar{w}$ admits two Taylor expansions

$$\bar{w}(u) = \pm\mu\,(F_2 + \bar{F}_2)\sqrt{F_2\bar{F}_2}\cdot u^3 + O(u^4)\,. \tag{7.13}$$

As a result, on the regularized $\widetilde{\widetilde{C}}$ each of these points $\tilde{\pi}^{-1}(\mathcal{R}_1), \tilde{\pi}^{-1}(\mathcal{R}_2)$ gives rise to a pair of points, which we denote as $\tilde{R}_j^-, \tilde{R}_j^+$, $j = 1, 2$, according to the sign in (7.13).

Next, in view of (3.15), we have

$$\pi_1\,(\tilde{R}_1^\pm) = (u = 0, W = (2 + 0)/\sqrt{2} = \sqrt{2}, Z = 0),$$
$$\pi_1\,(\tilde{R}_2^\pm) = (u = 0, W = (-2 + 0)/\sqrt{2} = -\sqrt{2}, Z = 0),$$

which, by (3.19), gives

$$t^2(S_1^\pm) = t^2(S_2^\pm) = \frac{\bar{F}_2}{F_2} = \frac{K_2 - 2 + \lambda - \mu}{K_2 + 2 + \lambda - \mu}, \tag{7.14}$$
$$W(S_1^\pm) = \sqrt{2}, \quad W(S_2^\pm) = -\sqrt{2}. \tag{7.15}$$

To determine signs of $t(S_1^+), \ldots, t(S_2^-)$, we use the expression

$$t = \frac{W^2 - h(u)}{4\mu u^3\,(1 + \lambda\,u)Q(u)\,(F_1 u + F_2)} \tag{7.16}$$

obtained from (3.16). However, this expression gives the indeterminate result $0/0$ for $u = 0$. To resolve it, we use the Puiseux expansions (7.13) for $\bar{w}(u)$, as well as the expansion of $w(u)$, and substitute them into $W = (w + \bar{w})/\sqrt{2}$ to get the expansions of $W^2 - h(u)$ in powers of u. Near $\tilde{R}_1^\pm$ and $\tilde{R}_1^\pm$ we get

$$W^2 - h(u) = \pm 2\mu\,(F_2 + \bar{F}_2)\sqrt{F_2\bar{F}_2}\cdot u^3 + O(u^4).$$

Putting this into (7.16) and taking the limit $u \to 0$, yields

$$t(S_{1,2}^+) = t_* = \sqrt{\bar{F}_2/F_2}, \quad t(S_{1,2}^-) = -t_*\,.$$

Finally, to find $\mathcal{W}(S_1^\pm), \mathcal{W}(S_2^\pm)$, we substitute the values of $t(S_{1,2}^\pm)$ and $W(S_{1,2}^\pm)$ from (7.15) into (3.24). After simplifications, we get

$$\mathcal{W}(S_1^\pm) = 4\frac{H^{3/2}}{F_2}(1 \pm t_*), \quad \mathcal{W}(S_2^\pm) = -4\frac{H^{3/2}}{F_2}(1 \pm t_*),$$

which completes the proof. $\qquad\qquad\square$

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