

Unburned carbon measurement in fly ash using laser-induced breakdown spectroscopy with short nanosecond pulse width laser

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[Abstract]

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Abstract

The unburned carbon in fly ash is one of the important factors for the boiler combustion condition. Controlling the unburned carbon in fly ash is beneficial for fly ash recycle and to improve the combustion efficiency of the coal. Laser-induced breakdown spectroscopy (LIBS) technology has been applied to measure the fly ash contents due to its merits of non-contact, fast response, high sensitivity, and real-time measurement. In this study, experimental measurements have been adopted for fly ash flows with the surrounding gases of N₂ and CO₂, while the CO₂ concentration varied to evaluate the CO₂ effect on the unburned carbon signal from fly ash powder. Two kinds of pulse width lasers, 6ns and 1ns, were separately adopted to compare the influence of laser pulse width. Results showed that compared with 6ns pulse width laser, plasma temperature was lower and had less dependence on delay time when using 1ns pulse width laser, and spectra had more stable background. By using 1ns pulse width laser, the emission signal from surrounding CO₂ also decreased because of the less surrounding gas breakdown. The solid powder breakdown signals also became more stable when using 1ns pulse width laser. So it is demonstrated that 1ns pulse width laser has the merits for fly ash flow measurement using LIBS.

Keywords: laser-induced breakdown spectroscopy; fly ash powder flow; short nanosecond pulse width laser; solid powder breakdown; surrounding gas breakdown.

1. Introduction

The unburned carbon in fly ash is one of the important factors to evaluate the combustion efficiency of the boiler [1-3]. A fast-response measurement method of unburned carbon in fly ash can improve the control of boiler combustion. The improvement of the boiler combustion efficiency is beneficial for the fuels saving and the environment protection. Recycling the fly ash as the source of concrete is also important for its useful utilization. The carbon content in the fly ash is not conducive to its recycling. Therefore, the carbon content should be controlled. The online measurement of fly ash compositions is helpful to control the carbon content in time, because of the continuous coal combustion. The variation range of the unburned carbon in fly ash is very different for different boilers and combustion conditions. For large capacity boilers, the ash carbon content is usually about 2-15% [4]. For circulating fluidized bed boilers, it is about 10-20% [5]. When the coals species and load change in the plant, the fly ash content is also different.

A variety of fly ash detection methods have been proposed and applied in fly ash measurement. The traditional methods, such as thermal gravimetric analysis method, microwave attenuation method, and electrostatic capacitance measurement method and so on [6, 7], are widely applied. X-ray diffraction (XRD) powder analysis or SEM methods [1, 8-11] have also been applied to measure the compositions and microstructure of the fly ash. For online measurement, however, these methods meet the difficulties due to the sample preparation. Therefore, laser-induced breakdown spectroscopy (LIBS) is proposed for fly ash measurement due to its merits of non-contact, fast-response, multi-element detection and online measurement. Up to now, LIBS technology has been applied in different fields such as power plants, steel making processes, environment, marine, food safety, and so on [12-14].

LIBS measurements of fly ash have also attracted much attention. The measurement properties and quantitative analysis method of fly ash were extensively studied [3, 15-23]. The influence of ambient pressure and temperature, powder sizes, laser energy and so on were discussed [3, 15-17]. Some calibration methods were proposed [18-21] to improve the quantitative measurement, such as internal standard method, plasma temperature correction method, multivariate calibration method, wavelet neural method. In some studies [22, 23], the fly ash samples were prepared as pellet to detect the elements of Ca, Mg, Fe and so on.

78 For the online measurement of fly ash flow in the boiler, LIBS measurement of fly ash has its own problem.
79 LIBS measurement of powder flow should consider both the solid powder breakdown and the surrounding
80 gas breakdown. Besides, the components of flue gas are very complex, such as NO_x , SO_x , CO_2 and so on.
81 The concentration of CO_2 is usually up to the level of around 10% [24]. The induced CO_2 plasma emits the
82 carbon signal which introduces measurement error to the unburned carbon content of fly ash powder [25].
83 To reduce the CO_2 effect, the measurement systems were designed and applied to reduce CO_2 concentration
84 [26-28], which were able to measure the fly ash flow continuously. For example, a cyclone system was
85 applied to separate the fly ash powder with the flue gas [26]. Another system was designed as that the fly ash
86 was dropped onto the conveying belt. The measurement position was covered by a chamber when
87 introducing air to reduce the CO_2 concentration [27, 28]. Reduction of CO_2 concentration by sampling and
88 measurement systems was a valid method to reduce the CO_2 influence on unburned carbon signal for online
89 measurement of fly ash. However, the detailed influence of CO_2 on unburned carbon signal was not
90 discussed entirely in these studies. Because the CO_2 effect is rather sensitive to the concentration in
91 surrounding gas and less than 1% CO_2 causes the spurious C signal in LIBS spectra of fly ash [26].

92 The control of the solid powder breakdown and surrounding gas breakdown is also a way to decrease the
93 surrounding gas effects. The laser pulse width is one important factor for laser-induced plasma, and many
94 researches have testified the pulse width effects on LIBS measurement [29-35] for both gas-phase target and
95 solid measurement. For gas-phase target measurement, the shorter pulse width laser induced a better target
96 element measurement. For solid measurement, when using the shorter pulse width laser, more laser beam
97 energy was concentrated on the solid surface, and craters on solid surface became more regular, such as
98 picosecond or femtosecond laser. So the LIBS detection ability can be improved, such as the higher signal to
99 noise ratio, lower background emission. Therefore, the laser-induced plasma processes of solid powder and
100 surrounding gas can be changed when using different pulse width lasers. In this study, the CO_2 effect on the
101 unburned carbon signal was studied according to the laser-induced plasma processes. Two different
102 nanosecond pulse width lasers were employed to discuss the laser-induced plasma processes of fly ash flow.
103

104 2. Theory

105
106 In the LIBS measurement of powder flow, the sample is in the multi-phase condition including solid
107 powder and surrounding gas. The laser-induced plasma formation processes of the solid powder and its
108 surrounding gas can be summarized as follow [35, 36]: When the laser beam is focused on the sample, the
109 solid powders firstly absorb energy and evaporate as the material vapour. Plasma is firstly generated within
110 the material vapour. Because of the high pressure and temperature, the material vapour and plasma expand
111 and mix with the surrounding gas and the residual part of the powder, which are also broken down through
112 the electron collision, which can be enhanced through plasma reheating by laser pulse. The formation and
113 development processes of plasma are illustrated in Fig.1. Fig.1 (a) is the plasma structure concept of the
114 solid powder with the surrounding gas. When the laser energy is small, surrounding gas cannot be broken
115 down directly due to its high breakdown threshold [19]. So the solid powder plasma is firstly generated, then
116 expands and induces the breakdown of surrounding gas. Fig.1 (b) shows the diagram of different laser pulse
117 width, and Fig.1 (c) illustrates the concept of differences of plasma evolution processes using different pulse
118 width lasers. As mentioned in plasma formation processes, there is a time gap between the beginning time of
119 solid breakdown and gas breakdown, while the end time of breakdown is near to the end time of laser pulse.
120 When the laser pulse width becomes shorter, the time of plasma reheating by laser pulse will be shorter [30],
121 and more plasma is formed within the solid material vapour. What's more, when the laser pulse energy is
122 same, the peak power is higher when the pulse width is shorter, as shown in Fig.1 (b), and more laser pulse
123 energy is allocated to solid breakdown. In this way, the solid powder breakdown and surrounding gas
124 breakdown can be controlled by the laser pulse width. Previously, the pulse width influence was discussed
125 by comparing the nanosecond laser with picosecond or femtosecond lasers. The shorter pulse width laser can
126 better concentrate energy for the target breakdown, but the complex laser system challenges their on-line
127 applications. It has been mentioned [36] that the time of solid evaporation is usually within 1ns; while up to
128 around 5ns, the material vapour expands to induce the surrounding gas breakdown. Therefore, even if 1ns
129 pulse width, the plasma formation processes may also have obvious difference. Besides, 1ns nanosecond
130 laser can keep its relative structure simplicity, which is good for online measurement.

The analysis of LIBS is usually based on the signal intensity of characteristic lines. The theoretical equation between the signal intensity and the concentration of measured species with the local thermal equilibrium (LTE) assumption has been proposed. The detailed explanation is shown in elsewhere [12]. The signal can be calculated as the ratio of $I_{\text{target}}/I_{\text{reference}}$ [16, 20, 27] to discard the absolute intensity fluctuation influence of the spectra. Plasma temperature is one of the main influencing factors for signal variation. To evaluate plasma temperature, similar as the Saha-Boltzmann multi-line graph theory [37], the plasma temperature indicator is proposed [3]. Plasma temperature indicator is defined as the ratio of same element from different upper energy levels: $I_{i,j1}/I_{i,j2}$. i means the emission element, $j1$, $j2$ mean different upper energy levels. According to the relation between $I_{\text{target}}/I_{\text{reference}}$ and $I_{i,j1}/I_{i,j2}$, the plasma temperature correction factor can be determined [5, 19, 26]. Plasma temperature correction method is applied to reduce the influence of plasma temperature fluctuation on the signal intensity. In this study, I_{Mg1}/I_{Mg2} was defined as the plasma temperature indicator. The functional relation between I_C/I_{Si1} and I_{Mg1}/I_{Mg2} was fitted into some exponential relation and the exponential term of I_{Mg1}/I_{Mg2} was the temperature correction factor, which can be calculated theoretically but actually determined according to the experiment due to the complex laser-induced plasma processes [3, 16].

3. Experimental

The experimental diagram of fly ash measurement using LIBS is presented in Fig.2. It contained two lasers, a spectrometer, an ICCD camera and auxiliary equipment. These two lasers were both nanosecond lasers with different pulse widths operating at 1064nm. The laser marked as Laser 1 was a laser with the pulse width of 6ns (LOTIS TII 2132-UTF). The laser marked as Laser 2 was a laser with the pulse width of 1ns (Hamamatsu Photonics, Microchip Laser L12968-01). During the experiment, one laser was operated and the optical path of laser beam was altered by the up-and-down adjustment of Mirror 2. Both lasers were operated at the frequency of 10Hz. The Focus Lens 1 with the focal length of 200mm was employed to focus the laser beams onto the fly ash powder flow. In order to compare the influence of laser pulse width on fly ash measurement, the conditions of these two laser beams including laser energy, beam shape, beam

diameter and beam axis in front of Focus Lens 1 at the same point were adjusted to the same by a series of optical components. The laser energy was 7mJ/pulse and the beam diameter was 8mm. The low laser energy was set to induce solid powder plasma and to avoid the direct breakdown of surrounding gas. Because of the transmittance of the gas, its breakdown threshold is usually higher than that of solid [29, 38, 39]. In this study it was testified that 7mJ/pulse is lower than the gas breakdown threshold and it cannot induce the gas breakdown directly. The plasma emission was detected at the coaxial direction with the laser beam. The plasma emission was reflected by a filter that reflects the light below 400nm, and then focused onto the fiber by the Focus Lens 2 with the focal length of 220mm.

The emission signals from the plasma were detected by the combination of a spectrometer (SOL, NP-250-2), an ICCD camera (Andor, iStar DH334T-18U-03), and auxiliary equipment. The spectrometer with two channels was employed, which can simultaneously detect the spectra with different resolution when using the different gratings. In this study, the resolution of Channel 1 was 0.076nm/pixel with the grating of 600l/mm. The wavelength region was 240~320nm. The resolution of Channel 2 was 0.012nm/pixel with the grating of 3600l/mm. The wavelength region was 244~256nm. The gate width was set as 500ns. The delay time was different when using different pulse width lasers. When Laser 1 was used, the delay time was 300, 500, 800, 1000, 1500ns. When Laser 2 was used, the delay time was 5, 10, 25, 50, 100ns. The accumulation was 100 times of laser shot and the experimental result under the same condition was measured 3 times.

In this study, the employed fly ash sample with unburned carbon of 24.9% was sampled from a furnace in the lab. The compositions of fly ash were checked by the conventional chemical analysis methods (Japanese Industrial Standards JIS-M-8801, JIS-M-8815) [16]. The main components of fly ash sample are listed in Table 1. The fly ash was introduced to the measurement area by a feeder (Nisshin Engineering Inc. Feedcon- μ Mtype). The feeding speed of the feeder was 200mg/min. A chamber covered the exit of the feeder to mix the fly ash with the surrounding gas of N₂ and CO₂ mixture. The N₂ and CO₂ were pre-mixed before introducing into the chamber. The pre-mixed surrounding gas was divided into two flows to totally cover and mix the fly ash powder under the N₂ and CO₂ mixture condition during the measurement process. One surrounding gas flow was introduced from the top of the chamber with the flow rate of 1L/min. The fly ash dropped out from the outlet pipe with the diameter of 3mm. Another gas was introduced to the underside of

the chamber to surround the outlet of pipe, with the flow rate of 4L/min. In order to clarify the CO₂ effect in this study, the CO₂ volume percentage was set as 0%, 1.30%, 3.80%, and 6.18% in the measurement system. The laser beam was focused closely to this outlet pipe to ensure the measured fly ash mixed with the surrounding gas of N₂ and CO₂ mixture.

4. Results and Discussion

The spectrometer employed in this study with two channels can measure two wavelength ranges of the spectra simultaneously. The main elements in fly ash, such as C, Si, Al, Fe, Ca, Mg, were measured using Channel 1 with the wide wavelength range from 240nm to 320nm, with the resolution of 0.076nm/pixel. According to the previous studies, the C emission line at 247.86nm can be interfered by the neighbor lines such as Fe line at 247.98nm [11, 40, 41]. In order to clarify the unburned carbon in fly ash, the clear C emission line was detected using Channel 2 with the wavelength range from 244nm to 256nm, with the resolution of 0.012nm/pixel. The C content of fly ash sample employed in this study was 24.9% to discuss the surrounding gas effect on the solid powder measurement. It was verified that Fe line showed less interference on C signal in this study.

The typical measured spectra of fly ash without CO₂ using Laser 1 with the pulse width of 6ns is shown in Fig.3. The measured species and their specific wavelengths are listed in Table 2 according to the NIST database [42]. The emission intensity of the spectra in this paper was normalized with the corresponding maximum signal in each measured result.

$I_{\text{target}}/I_{\text{reference}}$ and $I_{i,j1}/I_{i,j2}$ were defined according to the ratios of the characteristic emission lines, which were selected from the same channel. For the $I_{\text{target}}/I_{\text{reference}}$, Si was selected as the reference element, because Si was a major component in fly ash. For unburned carbon measurement, the ratio of C (247.86nm)/Si1 (251.61nm) in Channel 2 was used to evaluate the carbon content [3, 16]. Other elemental signals should also be considered [3], such as the ratio of Fe/Si2, Ca1/Si2, and Al/Si2 in Channel 1. The intensity of Al was the sum of Al1 and Al2. For the intensity ratio of the same element from different upper energy levels $I_{i,j1}/I_{i,j2}$,

Mg1 (ion, 279.55nm, 280.27nm)/Mg2 (atom, 285.21nm) was defined as the plasma temperature indicator, which showed the relation with the plasma temperature [16].

4.1 Pulse width effect on plasma temperature

Plasma temperature indicator I_{Mg1}/I_{Mg2} was evaluated in different delay time, which means the plasma evolution process. The dependence of I_{Mg1}/I_{Mg2} on delay time using different lasers under different CO₂ concentration conditions is shown in Fig.4. From Fig.4, the plasma temperature was not affected by the CO₂ concentration obviously. However, the plasma temperature decreased when delay time increased. Comparing Fig.4 (a) and Fig.4 (b), the plasma states using two nanosecond lasers had obvious differences. When delay time was 300ns using Laser 1, I_{Mg1}/I_{Mg2} was averagely 7.65 under different CO₂ concentration conditions. While using Laser 2, I_{Mg1}/I_{Mg2} was averagely only 3.03 when delay time was 5ns. The setting of delay time with Laser 1 cannot be shorter than 300ns, because the plasma temperature would be very high and the emission intensity would be very strong, saturating the ICCD detector. But it can be understood that the plasma temperature would be much higher if in delay time of 5ns using Laser 1. Considering that the delay time of 5ns was just after the plasma formation, it means at the formation time, plasma temperature using Laser 1 was much higher than that using Laser 2. Along the evolution of plasma, using Laser 1, I_{Mg1}/I_{Mg2} varied as 7.65, 5.66, 4.51, 3.50, 2.00 from 300ns to 1500ns, with obvious dependence on delay time. While using Laser 2, I_{Mg1}/I_{Mg2} varied as 3.03, 3.02, 2.96, 2.75, 2.44 when delay time increased from 5ns to 100ns, with less dependence on delay time.

The high plasma temperature can introduce some problems for the LIBS measurement. The measured spectra without CO₂ in delay time of 300ns using Laser 1 are shown in Fig.5. Comparing Fig.5 in delay time of 300ns and Fig.3 (a) in delay time of 800ns, the background signal was not horizontal and some emission lines were immersed in the background signals at higher plasma temperature in Fig.5. The Al line background in Fig.3 (a) was 0.048, and in Fig.5 was 0.185. At high plasma temperature, some effects influence the background, such as the strong bremsstrahlung, wide line broadenings, or self-absorption. What's more, when the plasma temperature is high, the absolute emission intensity is strong, which is easy to saturate for the detector. The spectra behaviors change with the plasma formation and emission processes

[36]. The measured spectra using Laser 2 in delay time of 10ns and 100ns are shown in Fig.6, the backgrounds of Al line were 0.037 and 0.033, which didn't show the obvious variation of the spectra in different delay time, so the signal background could be clearly judged.

According to the theory, when using short 1ns pulse width laser, more energy is allocated to solid breakdown. While using long 6ns pulse width laser, plasma can be more reheated during gas breakdown. From Fig.4, it has testified that the plasma formation processes have obvious differences when using 6ns pulse width laser and 1ns pulse width laser, and plasma reheating processes have more obvious influence to increase plasma temperature. When using 1ns, the plasma temperature at the formation time was lower and had less dependence on delay time.

The obvious dependence on delay time of spectra will induce the variation of the signal in the practical measurement. Because plasma formation process is very fast and not fixed, obvious dependence on delay time means plasma states have obvious differences with different delay time. In this way, the setting of delay time should carefully consider plasma evolution process. However, for Laser 2, the setting of delay time will be more robust. When the LIBS spectra are stabilized, it is benefit for the LIBS measurement of the powder flow.

4.2 Pulse width effect on C signal

4.2.1 Surrounding CO₂ effects using different pulse width lasers

The laser pulse width not only affects the plasma states, but also the breakdown of surrounding gas. The surrounding gas can also be broken down to generate the plasma. For example, the breakdown of surrounding CO₂ in fly ash emits the carbon signal, which results in the spurious carbon signal to increase the calculated unburned carbon content in fly ash.

The comparison of the spectra using two different lasers should be at similar plasma states. When delay time was 1000ns using Laser 1, I_{Mg1}/I_{Mg2} showed the similar level as that using Laser 2. As for Laser 2, when delay time was 5ns and 10ns, even though I_{Mg1}/I_{Mg2} showed the similar level as Laser 1, there would be other effects influencing signals just after plasma formation, so they were not suitable for the comparison. While after delay time of 50ns, I_{Mg1}/I_{Mg2} was relatively small. Therefore, the measured spectra in delay time

of 25ns using Laser 2 were compared with that in delay time of 1000ns using Laser 1 because of their similar plasma temperature.

The measured spectra of fly ash without CO₂ and with 3.80% CO₂ are shown in Fig.7 using Laser 1 in delay time of 1000ns in Channel 2. Comparing Fig.7 (a) and Fig.7 (b), when CO₂ concentration increased from 0% to 3.80%, the C emission line intensity increased from 0.298 to 0.439. The C signal from the CO₂ breakdown resulted in the inaccurately quantitative analysis of the unburned carbon content in fly ash. The measured spectra of fly ash with different CO₂ concentration using Laser 2 in delay time of 25ns in Channel 2 are shown in Fig.8. Comparing Fig.8 (a) without CO₂ and Fig.8 (b) with 3.80% CO₂, the C emission line intensity changed from 0.257 to 0.259, which didn't show the obvious change compared to that using Laser 1.

Therefore, the breakdown rate of surrounding gas was different from different lasers due to the pulse width effect. When the laser pulse width was shorter, such as 1ns, the breakdown rate of the surrounding gas decreased. The C signal from the CO₂ breakdown reduced.

According to the measured spectra, the dependence of I_C/I_{Si1} on CO₂ concentration is shown in Fig.9. Compared with I_C/I_{Si1} under 0% CO₂ condition, when the CO₂ concentration increased, the maximum enhancement of I_C/I_{Si1} using Laser 1 was 41.2%, while the maximum enhancement of I_C/I_{Si1} using Laser 2 was 26.7%. The intensity enhancement of I_C/I_{Si1} using Laser 1 was higher than that using Laser 2. It indicated that the CO₂ effect on unburned carbon measurement of solid fly ash powder could be reduced when using Laser 2. The increase of I_C/I_{Si1} with CO₂ in Fig.9 didn't follow the linear relation due to the non-uniformity of the fly ash powder and the mixture of powder and surrounding gas [19]. The measured carbon signals contained both the powder carbon signal and gas carbon signal. The fluctuation of the powder flow caused the fluctuation and inhomogeneity of the plasma. Therefore, one of the merits of fly ash measurement using Laser 2 is the reduction of gas breakdown to diminish the CO₂ effect on unburned carbon measurement because of the effect on the solid powder and gas plasma processes, which can be applicable for other powder conditions. Another merit is that Laser 2 is a microchip laser with high durability and long lifetime. Laser 2 employed here is more available for online measurement of LIBS system due to its flexibility and durability.

4.2.2 Plasma temperature correction of C signal

According to the previous study, the carbon signal intensity was affected by the plasma temperature [3, 16]. As shown in Fig.4, plasma temperature showed the obvious delay time dependence using Laser 1. However, the plasma temperature showed a little change using Laser 2. Plasma temperature correction is necessary to reduce the plasma temperature effect on signals, especially for the measurement using Laser 1. The temperature correction factor for carbon signal was determined according to the relation between I_C/I_{Si1} and I_{Mg1}/I_{Mg2} . The variation of plasma temperature with delay time was different using different lasers. Therefore, the plasma temperature correction factor for each laser was determined respectively including all the CO_2 concentration and delay time conditions.

After applying the plasma temperature correction, the relation between the corrected I_C/I_{Si1} and CO_2 concentration is shown in Fig.10. In Fig.10, the corrected I_C/I_{Si1} in each CO_2 concentration condition was the averaged result in all delay time. After plasma temperature correction, the maximum enhancement of I_C/I_{Si1} using Laser 1 was 30.4%, and the maximum enhancement of I_C/I_{Si1} using Laser 2 was 19.7%. Compared with Fig.9, the enhancement of I_C/I_{Si1} decreased after the plasma temperature correction. As for the standard deviation, the largest standard deviation in Fig.9 was 0.0712 of Laser 1 and 0.0646 of Laser 2, while the largest standard deviation in Fig.10 was 0.0649 of Laser 1 and 0.0624 of Laser 2. Plasma temperature correction didn't make obvious difference on standard deviation. In a word, plasma temperature correction can partially decrease the influence of surrounding CO_2 breakdown.

4.3 Pulse width effect on solid breakdown signals

According to the C quantitative calculation equation, solid powder breakdown signals of I_{Ca1}/I_{Si2} , I_{Fe}/I_{Si2} , I_{Al}/I_{Si2} should also be concerned [16]. So the pulse width effect on these solid powder breakdown signals was also discussed here. As mentioned in Theory, for 6ns pulse width laser, more energy is introduced to plasma reheating process, while for 1ns pulse width laser, more energy is allocated to solid breakdown. In order to study the pulse width effect on these solid powder breakdown signals, I_{Ca1}/I_{Si2} was discussed here in detail under different experimental conditions. I_{Fe}/I_{Si2} , I_{Al}/I_{Si2} showed the consistent results [3, 16]. In the

case of I_C/I_{Si1} analysis which was affected by the surrounding CO_2 gas, it has been discussed above in detail.

The CO_2 concentration dependencies of I_{Ca1}/I_{Si2} in different delay time when using Laser 1 and Laser 2 are shown in Fig.11. I_{Ca1}/I_{Si2} didn't show the obvious change with CO_2 concentration. When using Laser 1, the averaged I_{Ca1}/I_{Si2} under different CO_2 concentration conditions was as high as 0.575 in delay time of 300ns, and it decreased as delay time increased, varying as 0.536, 0.496, 0.380, 0.225, due to the higher upper energy level of Ca1 than that of Si2. However, the averaged I_{Ca1}/I_{Si2} using Laser 2 had less variation, varying as 0.342, 0.349, 0.347, 0.308, and 0.297 along the increase of delay time. Using Laser 2, I_{Ca1}/I_{Si2} also had less dependence on delay time compared with that using Laser 1. The values of I_{Ca1}/I_{Si2} using Laser 2 in different delay time were similar as that using Laser 1 in delay time of 1000ns. The results in Fig.11 corresponded to the plasma temperature in Fig.4. Therefore, it indicated that plasma temperature has crucial effects on the solid breakdown signals. The signal stability under different lasers was also different, and the large relative standard deviation indicated the fluctuation of signals. For Laser 1, 8 results of total 20 results in Fig.11 (a) had the relative standard deviation more than 10%, and 3 of them were more than 20%. For Laser 2, 6 results in Fig.11 (b) had the relative standard deviation more than 10%, while only 1 of them was more than 20%. So the signals using Laser 2 had less fluctuations, which is beneficial for the measurement.

5. Conclusions

The unburned carbon in fly ash is an important factor to evaluate combustion efficiency. The low carbon content in fly ash means the good combustion of the boiler. Besides, the recycling of the fly ash also requires the low carbon content in fly ash. Online measurement of fly ash compositions helps to control the carbon content of fly ash, and LIBS technology can realize the online measurement compared to other methods. The LIBS measurement of fly ash powder flow should consider both the solid powder breakdown and surrounding gas breakdown, while the surrounding gas breakdown will introduce the carbon signal from CO_2 gas to the unburned carbon measurement of fly ash. Experiments were designed using 6ns pulse width laser and 1ns pulse width laser to measure fly ash separately to compare their properties. The results are concluded as follow.

346 1. At the formation time of plasma, plasma temperature using 6ns pulse width laser was much higher than
347 that using 1ns pulse width laser. High plasma temperature affected the background of spectra and
348 measurement accuracy. Using 6ns pulse width laser, the plasma temperature had obvious dependence on
349 delay time, the parameters setting should consider the plasma states. It has testified that when the laser pulse
350 width changes from 6ns to 1ns, the plasma formation process had obvious change, which showed the lower
351 plasma temperature at the formation time, more stable background, and less dependence on delay time.

352 2. The I_C/I_{Si1} increased obviously with the increase of the CO_2 concentration when using 6ns pulse width
353 laser. The CO_2 concentration dependence of I_C/I_{Si1} was reduced when using 1ns pulse width laser under all
354 the CO_2 concentration conditions. The CO_2 effect on unburned carbon measurement could be reduced when
355 using 1ns pulse width laser. When applying the plasma temperature correction method, the increase of I_C/I_{Si1}
356 with the CO_2 concentration became smaller.

357 3. The laser pulse width also influenced the solid breakdown signals stability. The powder breakdown
358 signals using 1ns pulse width laser showed less delay time dependence and generally smaller relative
359 standard deviations than that using 6ns pulse width laser.

360 In summary, the CO_2 effect on the unburned carbon signal can be reduced when using the laser with pulse
361 width of 1ns, which also induces less variation of plasma temperature and solid powder breakdown signals
362 like I_{Ca1}/I_{Si2} , I_{Fe}/I_{Si2} , I_{Al}/I_{Si2} . All these are beneficial for fly ash flow measurement. At the same time, the laser
363 with pulse width of 1ns is a microchip laser with high durability and long lifetime, which is more suitable
364 for the LIBS system integration for the industrial online application. In a word, 1ns pulse width laser is more
365 suitable for the fly ash powder flow measurement using LIBS.

367 **Acknowledgements**

368
369 This work was supported by National Natural Science Foundation of China (No. 51506171), the National
370 Key Basic Research Development Plan (No. 2015CB251504).

372 **References**

- 373 [1]. Kazemian H, Naghdali Z, Kashani T G, et al, Conversion of high silicon fly ash to Na-P1 zeolite:
374 Alkaline fusion followed by hydrothermal crystallization, *Advanced Powder Technology*, 21 (2009) 279-
375 283.
- 376 [2]. Noda N, Makino H, Influence of operating temperature on performance of electrostatic precipitator for
377 pulverized coal combustion boiler, *Advanced Powder Technology*, 21 (2010) 495-499.
- 378 [3]. Noda M, Deguchi Y, Iwasaki S, et al, Detection of carbon content in a high-temperature and high-
379 pressure environment using laser-induced breakdown spectroscopy, *Spectrochimica Acta Part B Atomic*
380 *Spectroscopy*, 57 (2002) 701-709.
- 381 [4]. A.K. Ouazzane, J.L. Castagner, A.R. Jones, et al, Design of an optical instrument to measure the carbon
382 content of fly ash, *Fuel*, 81 (2002) 1907-1911.
- 383 [5]. Y. Qiu, M.C. Tian, L. Cheng, et al, The Theory of the Reduction of Fly-ash Carbon Content in a
384 Circulating Fluidized Boiler and Its Applied Research, 20 (2005) 369-373.
- 385 [6]. Xie H L, Zhang, The Necessity of Monitoring Carbon Content in Fly-ash and the Development of
386 Measuring Technique, *Science Technology & Engineering*, 16 (2016) 115-121.
- 387 [7]. Cheng Q M, Xiao-Qing H U, Wang Y F, et al, Summary of Measurement Methods of Carbon Content in
388 Fly Ash, *Journal of Shanghai University of Electric Power*, 27 (2011) 519-524.
- 389 [8]. Marrero J, Polla G, Rebagliati R J, et al, Characterization and determination of 28 elements in fly ashes
390 collected in a thermal power plant in Argentina using different instrumental techniques, *Spectrochimica*
391 *Acta Part B Atomic Spectroscopy*, 62 (2007) 101-108.
- 392 [9]. Tanaka H, Fujii A, Fujimoto S, et al, Microwave-Assisted Two-Step Process for the Synthesis of a
393 Single-Phase Na-A Zeolite from Coal Fly Ash, *Advanced Powder Technology*, 19 (2008) 83-94.
- 394 [10]. Li Q, Pang J, Wang B, et al, Preparation, characterization and microwave absorption properties of
395 barium-ferrite-coated fly-ash cenospheres, *Advanced Powder Technology*, 24 (2013) 288-294.
- 396 [11]. Chindaprasirt P, Role of microwave radiation in curing the fly ash geopolymer, *Advanced Powder*
397 *Technology*, 24 (2013) 703-707.
- 398 [12]. Y. Deguchi, *Industrial applications of laser diagnostics*, CRC Press, 2011.

- 399 [13]. Wang Z Z, Deguchi Y, Zhang Z Z, et al, Laser-induced breakdown spectroscopy in Asia, *Frontiers of*
400 *Physics*, 11 (2016) 231-255.
- 401 [14]. Wang Z, Yuan T B, Hou Z Y, et al, Laser-induced breakdown spectroscopy in China. *Frontiers of*
402 *Physics*, 9 (2014) 419-438.
- 403 [15]. Zhang L, Ma W G, Yan X J, et al, Research on parameters optimization of laser-induced breakdown
404 spectroscopy based experimental device, *Spectroscopy & Spectral Analysis*, 31 (2011) 2355-2360.
- 405 [16]. Wang Z Z, Deguchi Y, Kuwahara M, et al, Quantitative elemental detection of size-segregated
406 particles using laser-induced breakdown spectroscopy, *Spectrochimica Acta Part B Atomic Spectroscopy*,
407 87 (2013) 130-138.
- 408 [17]. Bai K J, Tian H C, Yao S C, et al, Influence of laser energy on measurement of unburned carbon in fly
409 Ash particle flow, *Spectroscopy & Spectral Analysis*, 34 (2014) 1407.
- 410 [18]. S. C. Yao, J. D. Lu, C. L. Xie, et al, Quantitative analysis of laser induced carbon plasma by intensity
411 ratio calibration, *High Power Laser and Particle Beams*, 20 (2008) 1089-1092.
- 412 [19]. Wang Z Z, Deguchi Y, Watanabe H, et al, Improvement on Quantitative Measurement of Fly Ash
413 Contents Using Laser-Induced Breakdown Spectroscopy, *Journal of Flow Control Measurement &*
414 *Visualization*, 03 (2015) 10-21.
- 415 [20]. Yao S, Lu J, Zheng J, et al, Analyzing unburned carbon in fly ash using laser-induced breakdown
416 spectroscopy with multivariate calibration method, *Journal of Analytical Atomic Spectrometry*, 27 (2012)
417 473-478.
- 418 [21]. Zhang T, Yan C, Qi J, et al, Classification and discrimination of coal ash by laser-induced breakdown
419 spectroscopy (LIBS) coupled with advanced chemometric methods, *Journal of Analytical Atomic*
420 *Spectrometry*, 32 (2017) 1960-1965.
- 421 [22]. Stankova A, Gilon N, Dutruch L, et al, A simple LIBS method for fast quantitative analysis of fly
422 ashes, *Fuel*, 89 (2010) 3468-3474.
- 423 [23]. Ctvrtnickova T, Mateo M P, Yañez A, et al, Characterization of coal fly ash components by laser-
424 induced breakdown spectroscopy, *Spectrochimica Acta Part B Atomic Spectroscopy*, 64 (2009) 1093-1097.

- 425 [24]. Zhao Q X, Zhang Z X, Wen-Zhi D U, et al, Experimental Study on Sulfuric Acid Dewpoint Corrosion
426 in Simulated Atmospheric Conditions, *Journal of Chinese Society of Power Engineering*, 32 (2012) 420-424.
- 427 [25]. J. T. Bian, K. J. Yin, S. C. Yao, et al, Quantitative Analysis of Unburned Carbon in Fly Ash by Laser-
428 Induced Breakdown Spectroscopy in Different Atmosphere, *Laser & Optoelectronics Process*, 4 (2016) 234-
429 239.
- 430 [26]. Kurihara M, Ikeda K, Izawa Y, et al, Optimal boiler control through real-time monitoring of unburned
431 carbon in fly ash by laser-induced breakdown spectroscopy, *Appl Opt*, 42 (2003) 6159-6165.
- 432 [27]. Zhang L, Jia S T, Hu Z Y, et al, Recent progress on laser-induced breakdown spectroscopy for the
433 monitoring of coal quality and unburned carbon in fly ash, *Frontiers of Physics*, 7 (2012) 690-700.
- 434 [28]. Zhang L, Ma W, Dong L, et al, Development of an apparatus for online analysis of unburned carbon in
435 fly ash using laser-induced breakdown spectroscopy (LIBS), *Applied Spectroscopy*, 65 (2011) 790-796.
- 436 [29]. Wang Z Z, Deguchi Y, Kuwahara M, et al, Enhancement of Laser-Induced Breakdown Spectroscopy
437 (LIBS) Detection Limit Using a Low-Pressure and Short-Pulse Laser-Induced Plasma Process, *Applied*
438 *Spectroscopy*, 67 (2013) 1242-1251.
- 439 [30]. Eland K L, Stratis D N, Lai T, et al, Some Comparisons of LIBS Measurements using Nanosecond and
440 Picosecond Laser Pulses, *Applied Spectroscopy*, 55 (2001) 279-285.
- 441 [31]. Rieger G W, Taschuk M, Tsui Y Y, et al, Comparative study of laser-induced plasma emission from
442 microjoule picosecond and nanosecond KrF-laser pulses, *Spectrochimica Acta Part B Atomic Spectroscopy*,
443 58 (2003) 497-510.
- 444 [32]. Rohwetter P, Yu J, Méjean G, et al. Remote LIBS with ultrashort pulses: characteristics in picosecond
445 and femtosecond regimes[J]. *Journal of Analytical Atomic Spectrometry*, 19 (2004) 437-444.
- 446 [33]. Abdel-Salam Z A, Nanjing Z, Anglos D, et al, Effect of experimental conditions on surface hardness
447 measurements of calcified tissues via LIBS, *Applied Physics B*, 94 (2009) 141-147.
- 448 [34]. S. V. Rao, S. Sreedhar, M. A. Kumar, P. P. Kiran, S. P. Tewari and G. M. Kumar, Laser Induced
449 Breakdown Spectroscopy of high energy materials using nanosecond, picosecond, and femtosecond pulses:
450 Challenges and opportunities, *Proc. SPIE*, 2011, 8173.

- 451 [35]. Labutin T A, Lednev V N, Ilyin A A, et al, Femtosecond Laser-Induced Breakdown Spectroscopy,
452 Journal of Analytical Atomic Spectrometry, 31 (2015) 90-118.
- 453 [36]. Noll R. Laser-induced breakdown spectroscopy: fundamentals and applications, Springer, 2011.
- 454 [37]. Bye C A, Scheeline A., Saha-Boltzmann Statistics for Determination of Electron Temperature and
455 Density in Spark Discharges Using an Echelle/CCD System, Applied Spectroscopy, 47 (1993) 2022-2030.
- 456 [38]. Naeem T M, Matsuta H, Wagatsuma K, Effect of plasma gas for spectrometric analysis of tin and zinc
457 using low-pressure laser-induced plasma, Spectrochimica Acta. Part B: Atomic Spectroscopy, 58 (2003)
458 891-899.
- 459 [39]. Cristoforetti G, Legnaioli S, Palleschi V, et al, Characterization of a collinear double pulse laser-
460 induced plasma at several ambient gas pressures by spectrally- and time-resolved imaging, Applied Physics
461 B: Lasers and Optics, 80 (2005) 559-568.
- 462 [40]. Yao S C, Chen J C, Lu J D, et al, Influence of C-Fe Lines Interference Correction on Laser-Induced
463 Breakdown Spectroscopy Measurement of Unburned Carbon in Fly Ash, Spectroscopy & Spectral Analysis,
464 35 (2015) 1719-1723.
- 465 [41]. Bai K, Yao S, Lu J, et al, Correction of C-Fe line interference for the measurement of unburned carbon
466 in fly ash by LIBS[J]. Journal of Analytical Atomic Spectrometry, 31 (2016) 3418-2426.
- 467 [42]. Y. Ralchenko, NIST atomic spectra database. Laboratory Space Science Workshop, 1998.
468 https://physics.nist.gov/PhysRefData/ASD/lines_form.html
- 469
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Figure Captions

Fig.1 Plasma development processes during LIBS.

Fig.2 Fly ash measurement system of LIBS using different lasers.

Fig.3 Measured spectra of fly ash without CO₂ using Laser 1. Conditions: pulse width: 6ns, delay time: 800ns.

Fig.4 Delay time dependence of I_{Mg1}/I_{Mg2} using different lasers under different CO₂ concentration conditions.

Fig.5 Measured spectra without CO₂ using Laser 1 in Channel 1. Conditions: pulse width: 6ns, delay time: 300ns, resolution of 0.076nm/pixel.

Fig.6 Measured spectra without CO₂ using Laser 2 in Channel 1 in different delay time. Conditions: pulse width: 1ns, resolution of 0.076nm/pixel.

Fig.7 Comparison of measured spectra of fly ash with different CO₂ concentration using Laser 1 in Channel 2. Conditions: pulse width: 6ns, delay time: 1000ns, resolution: 0.012nm/pixel.

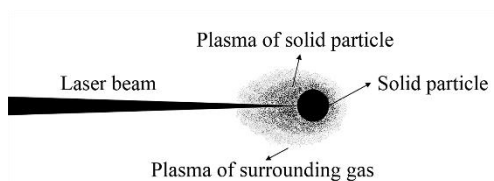
Fig.8 Comparison of measured spectra of fly ash with different CO₂ concentration using Laser 2 in Channel 2. Conditions: pulse width: 1ns, delay time: 25ns, resolution: 0.012nm/pixel.

Fig.9 Dependence of I_C/I_{Si1} on CO₂ concentration using different lasers in Channel 2. Conditions of Laser 1: pulse width: 6ns, delay time: 1000ns. Conditions of Laser 2: pulse width: 1ns, delay time: 25ns.

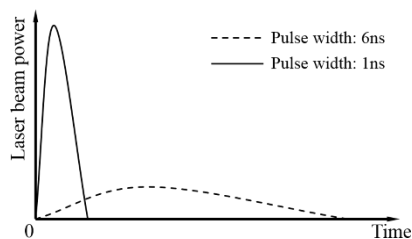
Fig.10 Dependence of averaged I_C/I_{Si1} on CO₂ concentration after plasma temperature correction in Channel 2. Conditions: pulse width of Laser 1: 6ns, pulse width of Laser 2: 1ns.

Fig.11 Dependence of I_{Ca1}/I_{Si2} on CO₂ concentration in different delay time using Laser 1 and Laser 2. Conditions: Laser 1, pulse width: 6ns, delay time: 300ns, 500ns, 800ns, 1000ns, 1500ns. Laser 2, pulse width: 1ns, delay time: 5ns, 10ns, 25ns, 50ns, 100ns.

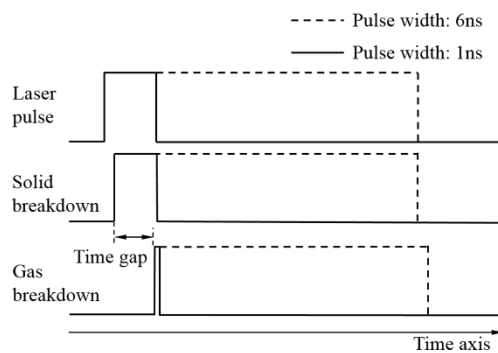
Fig.1



(a) Plasma structure concept of powder with surrounding gas



(b) Diagram of the laser pulse width



(c) Concept of time evolution process during laser-induced plasma

Fig.1 Plasma development processes during LIBS.

Fig.2

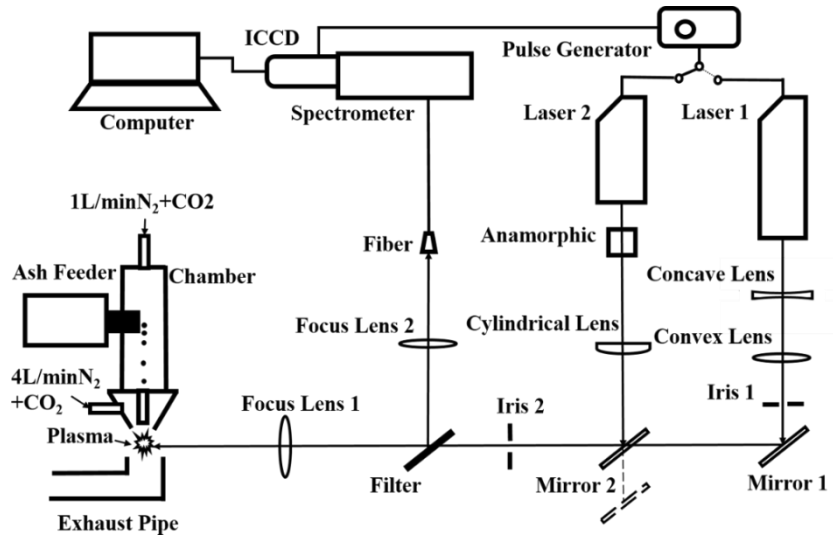
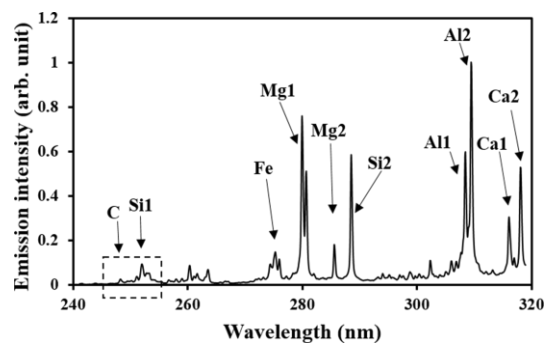
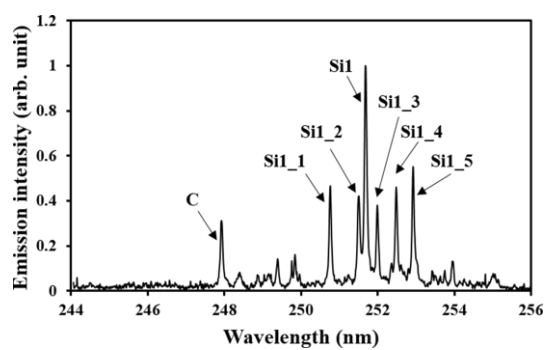


Fig.2 Fly ash measurement system of LIBS using different lasers.

Fig.3



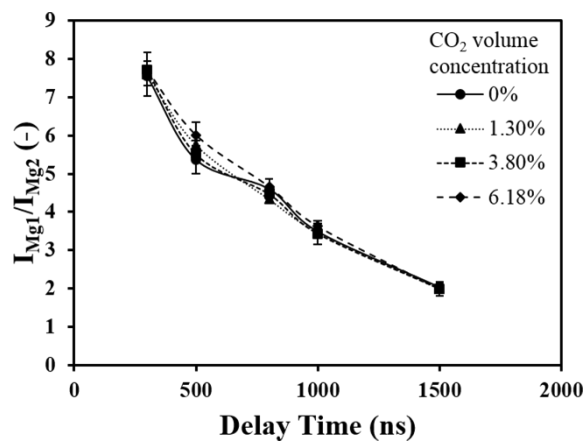
(a) Channel 1 with resolution of 0.076nm/pixel



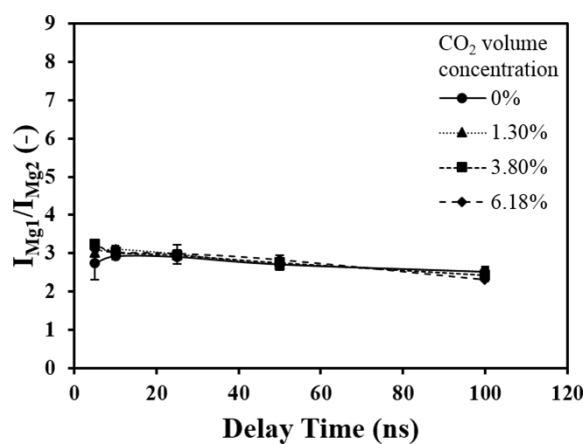
(b) Channel 2 with resolution of 0.012nm/pixel

Fig.3 Measured spectra of fly ash without CO₂ using Laser 1. Conditions: pulse width: 6ns, delay time: 800ns.

Fig.4



(a) Laser 1 with pulse width of 6ns



(b) Laser 2 with pulse width of 1ns

Fig.4 Delay time dependence of I_{Mg1}/I_{Mg2} using different lasers under different CO₂ concentration conditions.

Fig.5

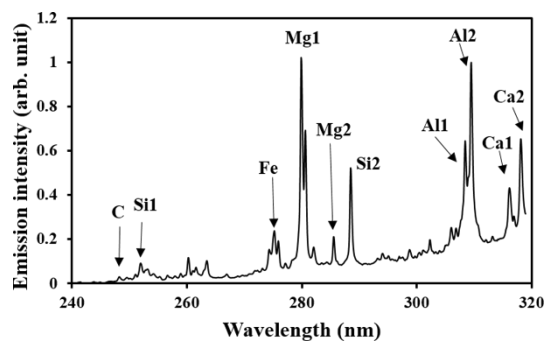
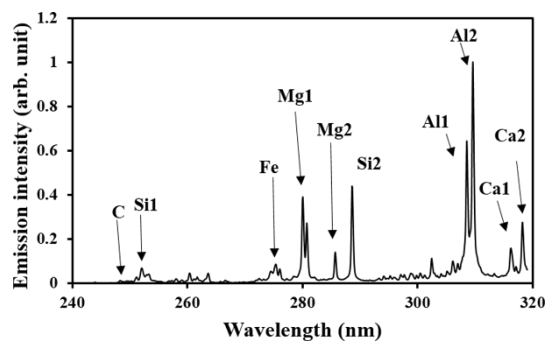
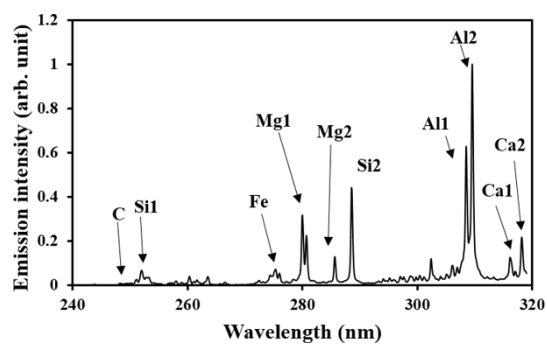


Fig.5 Measured spectra without CO₂ using Laser 1 in Channel 1. Conditions: pulse width: 6ns, delay time: 300ns, resolution of 0.076nm/pixel.

Fig.6



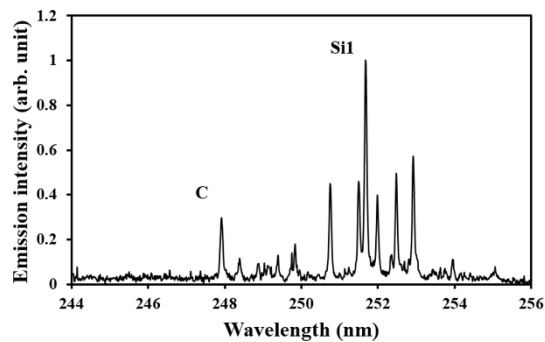
(a) Delay time of 10ns



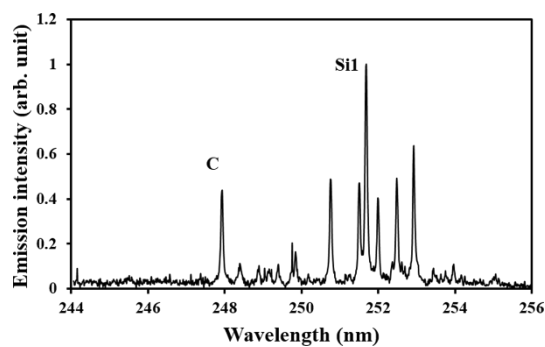
(b) Delay time of 100ns

Fig.6 Measured spectra without CO₂ using Laser 2 in Channel 1 in different delay time. Conditions: pulse width: 1ns, resolution of 0.076nm/pixel.

Fig.7



(a) Without CO₂

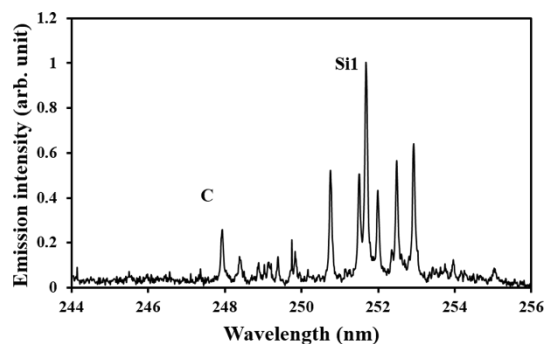


(b) With 3.80% CO₂

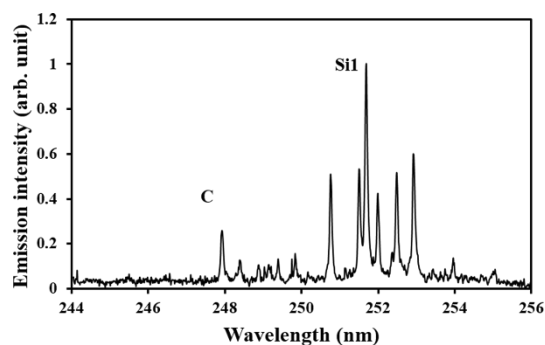
Fig.7 Comparison of measured spectra of fly ash with different CO₂ concentration using Laser 1 in Channel

2. Conditions: pulse width: 6ns, delay time: 1000ns, resolution: 0.012nm/pixel.

Fig.8



(a) Without CO₂



(b) With 3.80% CO₂

Fig.8 Comparison of measured spectra of fly ash with different CO₂ concentration using Laser 2 in Channel 2. Conditions: pulse width: 1ns, delay time: 25ns, resolution: 0.012nm/pixel.

Fig.9

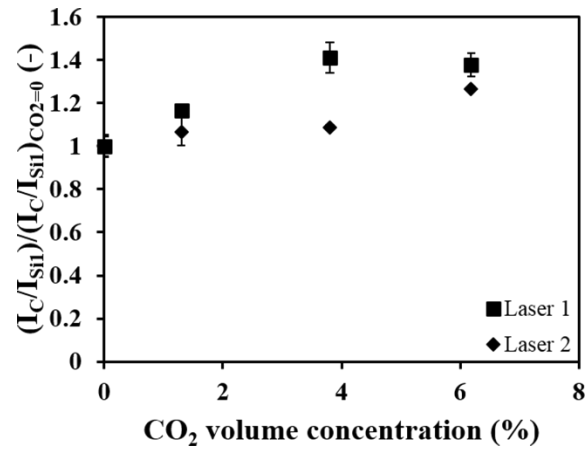


Fig.9 Dependence of I_C/I_{Si1} on CO₂ concentration using different lasers in Channel 2. Conditions of Laser 1: pulse width: 6ns, delay time: 1000ns. Conditions of Laser 2: pulse width: 1ns, delay time: 25ns.

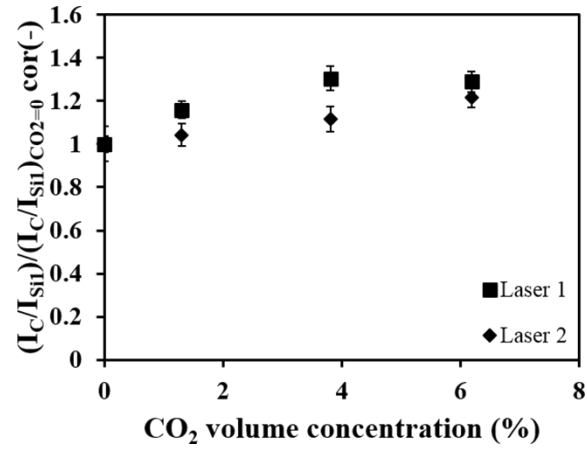
Fig.10

Fig.10 Dependence of averaged I_C/I_{Si1} on CO_2 concentration after plasma temperature correction in Channel 2. Conditions: pulse width of Laser 1: 6ns, pulse width of Laser 2: 1ns.

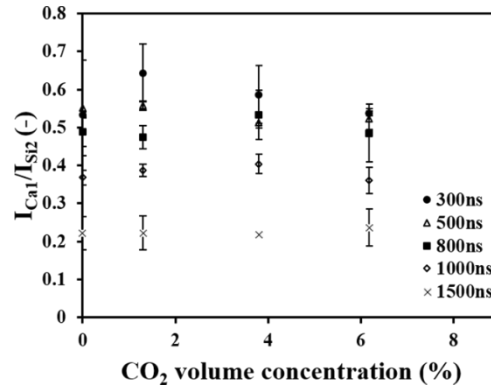
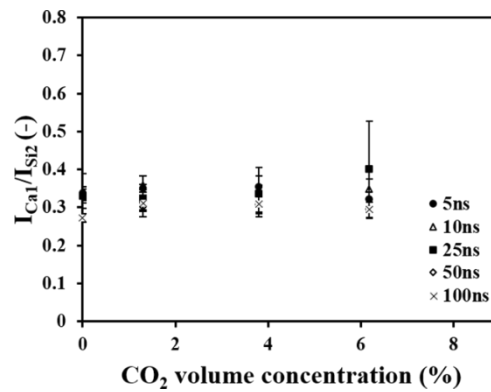
Fig.11(a) I_{Ca1}/I_{Si2} using Laser 1(b) I_{Ca1}/I_{Si2} using Laser 2

Fig.11 Dependence of I_{Ca1}/I_{Si2} on CO_2 concentration in different delay time using Laser 1 and Laser 2. Conditions: Laser 1, pulse width: 6ns, delay time: 300ns, 500ns, 800ns, 1000ns, 1500ns. Laser 2, pulse width: 1ns, delay time: 5ns, 10ns, 25ns, 50ns, 100ns.

Tables

Table 1 Main compositions of fly ash sample

Table 2 Measured species and the specific wavelengths [42]

Table 1

Table 1 Main compositions of fly ash sample

SiO ₂	Al ₂ O ₃	Fe ₂ O ₃	K ₂ O	Ash Content (wt%)					Na ₂ O	Unburned Carbon (wt%)
				CaO	TiO ₂	P ₂ O ₅	SO ₃	MgO		
37.38	25.91	6.26	0.39	1.36	0.8	0.36	0.53	0.59	0.78	24.9

Table 2

Table 2 Measured species and the specific wavelengths [42]

species	wavelength(nm)	$E_i(\text{cm}^{-1})$	$E_k(\text{cm}^{-1})$
C	247.86	21648.01	61981.82
Si1	251.61	223.16	39955.053
Si1_1	250.69	77.12	39955.053
Si1_2	251.43	0	39760.29
Si1_3	251.92	77.12	39760.29
Si1_4	252.41	77.12	39683.16
Si1_5	252.85	223.16	39760.29
Si2	288.16	6298.85	40991.884
Fe	274.20-	415.93-	36686.16-
(atom and ion)	275.63	8846.7	45289.80
Ca1(ion)	315.89	25191.51	56839.25
Ca2 (ion)	318.13	25414.4	56839.25
Al1	308.22	0	32435.45
Al2	309.27	112.06	32436.79
Mg1 (ion)	279.55	0	35760.88
Mg1 (ion)	280.27	0	35669.31
Mg2	285.21	0	35051.26