

SUPERMODEL ANALYSIS OF A1246 AND J255: ON THE EVOLUTION OF GALAXY CLUSTERS FROM HIGH TO LOW ENTROPY STATES

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ABSTRACT

We present an analysis of high-quality X-ray data out to the virial radius for the two galaxy clusters A1246 and GMBCG J255.34805+64.23661 (J255) by means of our entropy-based SuperModel. For A1246 we find that the spherically averaged entropy profile of the intracluster medium (ICM) progressively flattens outward, and that a nonthermal pressure component amounting to $\approx 20\%$ of the total is required to support hydrostatic equilibrium in the outskirts; there we also estimate a modest value $C \approx 1.6$ of the ICM clumping factor. These findings agree with previous analyses on other cool-core, relaxed clusters, and lend further support to the picture by Lapi et al. that relates the entropy flattening, the development of the nonthermal pressure component, and the azimuthal variation of ICM properties to weakening boundary shocks. In this scenario clusters are born in a high-entropy state throughout, and are expected to develop on similar timescales a low-entropy state both at the center due to cooling, and in the outskirts due to weakening shocks. However, the analysis of J255 testifies how such a typical evolutionary course can be interrupted or even reversed by merging especially at intermediate redshift, as predicted by Cavaliere et al. In fact, a merger has rejuvenated the ICM of this cluster at $z \approx 0.45$ by reestablishing a high-entropy state in the outskirts, while leaving intact or erasing only partially the low-entropy, cool core at the center.

Key words: galaxies: clusters: general – galaxies: clusters: individual (A1246, GMBCG J255.34805+64.23661) – X-rays: galaxies: clusters

1. INTRODUCTION

The low density $n \lesssim 10^{-4} \text{ cm}^{-3}$ of the intracluster medium (ICM) in the outskirts of galaxy clusters has severely limited, until recently, the capability of investigating those regions via the thermal bremsstrahlung emission $L_X \propto n^2$ in X rays. On the other hand, cluster outskirts are extremely interesting since: they provide the connection between the ICM and the filamentary structures of the cosmic web; they are the regions where most of the baryons and of the gravitationally dominant dark matter mass reside; and they constitute the sites of several physical processes and events affecting the ICM thermodynamic properties (see Kravtsov & Borgani 2012; Cavaliere & Lapi 2013; Reiprich et al. 2013). Thus nowadays the study of cluster outskirts is a very hot topic, embracing both astrophysics and cosmology.

A breakthrough in this field has been recently obtained with the advent of the *Suzaku* X-ray observatory, thanks to its low and stable particle background. The main, somewhat unexpected findings from the *Suzaku* data taken so far can be summarized as follows.

1. The ICM temperature rapidly declines outward by a factor of three in the region $r \sim 0.3\text{--}1 r_{200}$ ⁶ and slightly beyond (see Akamatsu et al. 2011; Reiprich et al. 2013); the temperature profiles are rather similar for relaxed and disturbed galaxy clusters (Kawaharada et al. 2010; Ichikawa et al. 2013; Simionescu et al. 2013; Sato et al. 2014).
2. The entropy profile $k(r)$ flattens at $r \gtrsim 0.5 r_{200}$ (see Walker et al. 2012, 2013) relative to the shape $k \propto r^{1.1}$

expected from strong-shocked infall of external gas under pure gravitational infall (see Tozzi & Norman 2001; Lapi et al. 2005; Voit 2005).

3. The thermodynamic properties of the ICM are subject to significant azimuthal variations (see Kawaharada et al. 2010; Ichikawa et al. 2013; Sato et al. 2014). For some clusters it has been possible to ascertain that hot regions are adjacent to filamentary structures, while cold regions are in contact with low-density, void-like environments; this indicates a more efficient thermalization in the overdense infall regions, and highlights how the environment surrounding the cluster affects the physical processes in the outskirts.
4. In some relaxed clusters the mass profile derived from X-ray observations under the assumption of thermal hydrostatic equilibrium features an unphysical, decreasing behavior at large radii (e.g., Kawaharada et al. 2010; Walker et al. 2012; Ichikawa et al. 2013; Okabe et al. 2014; Sato et al. 2014); this is a consequence of the rapid temperature decline, and can be explained in terms of an ICM far from thermal equilibrium, owing to the presence of a nonthermal pressure support (see Fusco-Femiano & Lapi 2013, 2014). An example is the cluster 1835 observed by *Chandra* (Bonamente et al. 2013) and by *Suzaku* (Ichikawa et al. 2013), where the underestimate of the hydrostatic mass implies a gas mass fraction f_{gas} higher than the cosmic value at the virial radius, as also reported for several other clusters (Simionescu et al. 2011; Fusco-Femiano & Lapi 2013; 2014).

The actual existence of the entropy flattening revealed by the *Suzaku* observations was challenged by Eckert et al. (2013); these authors estimated the entropy via the relation $k \propto p/n^{5/3}$ by combining average *ROSAT* gas density profiles (Eckert et al.

⁶ Here r_Δ is the radius within which the mean density is Δ times the critical density, while R is the virial radius of the cluster. Frequently used values read $r_{500} \approx R/2$ and $r_{200} \approx 3R/4$.

2012) with average electron pressure profiles $p(r)$ from stacked Sunyaev–Zel’dovich observations of 62 clusters by the *Planck* satellite (see *Planck* Collaboration et al. 2013). However, Fusco-Femiano & Lapi (2014) showed that such a procedure is not suited to discriminate between a steady power-law increase and a flattening of the entropy; this is because in hydrostatic equilibrium the dependence of $p(r)$ on $k(r)$ is much weaker than that of the temperature $T(r)$, so that the use of X-ray temperature data is mandatory for precise determination of the entropy profile in the outskirts.

In principle, the observed entropy flattening can be explained by gas clumping; this causes an overestimate of the ICM density and hence an underestimate of the entropy in the outskirts. On the other hand, hydrodynamical simulations show that the clumping factor $C \equiv \langle n^2 \rangle / \langle n \rangle^2 \lesssim 2$ of the ICM in the outskirts is rather limited and actually insufficient to explain the observed entropy flattening (Mathiesen et al. 1999; Nagai & Lau 2011; Vazza et al. 2013; Roncarelli et al. 2013; Zhuravleva et al. 2013; Battaglia et al. 2014; Morandi & Cui 2014), in agreement with the values for the clusters analyzed so far (Fusco-Femiano & Lapi 2013, 2014).

The presence of low-entropy gas in the outskirts can also be explained by considering that ions and electrons thermalize downstream the boundary shock on different timescales by Coulomb collisions (see Hoshino et al. 2010; Akamatsu et al. 2011). However, as noted by Okabe et al. (2014), the thermalization of the electrons can actually occur on the much shorter timescale of wave–particle interactions via plasma kinetic instabilities. A more recent proposal by Fujita et al. (2013) envisages that the entropy generation at the boundary shocks is not complete because part of the infall energy of the external gas is used to accelerate cosmic rays. We note that both these effects would be stronger in dynamically active clusters like Coma, where instead the entropy shows no clear evidences of flattening in undisturbed regions, and is much higher (even when scaled by the different mass) than in many relaxed clusters (Simionescu et al. 2013).

An alternative explanation of entropy flattening, steep temperature decline, nonthermal motions, and azimuthal variations, advocates the weakening of the boundary shocks (Lapi et al. 2010; Cavaliere et al. 2011b). In fact, weaker shocks produce less entropy, are less efficient in thermalizing the infall energy of the external gas, while allowing the residual one to seep inside and originate nonthermal motions in the form of turbulence. The shock weakening is in turn due to a reduced inflow, that can mainly occur under two circumstances: (1) either in the late evolution of a cluster, when the gas is accreted from the wings of the initial perturbations, and especially so at low redshift when the cosmic acceleration sets in, (2) or in a particular sector of a cluster facing an underdense, void-like region.

In this scenario clusters are born in a high-entropy state throughout by the strong shocks occurring at the time of formation. Then they are expected to develop synchronously a low-entropy state both at the center due to cooling (possibly balanced by gains from bubbling or rekindled active galactic nuclei in the central member galaxies, see Fabian 2012), and in the outskirts due to weakening shocks.

In Cavaliere et al. (2011a) we have predicted that the typical evolutionary course of galaxy clusters from high- to low-entropy states may be temporarily interrupted or even definitely reversed by a major merger; this can rejuvenate the ICM outskirts by adding entropy to reestablish the power-law behavior expected from strong shocks, and may penetrate deep in the cluster to

partially erase the cool core. We shall see that the analysis of the cluster GMBCG J255.34805+64.23661 (hereafter J255) demonstrates such an instance to actually occur in nature.

As mentioned above, the rapid decline of the gas temperature in the outskirts may cause the mass estimates based on X-ray observables and thermal hydrostatic equilibrium to be biased low by a systematic $\approx 10\%$ – 20% relative to the determinations via strong and weak lensing measurements (see Arnaud et al. 2007; Mahdavi et al. 2008, 2013; Lau et al. 2009; Battaglia et al. 2013). Recently, Okabe et al. (2014) have conducted a multi-wavelength analysis of four relaxed clusters reporting an average hydrostatic-to-lensing total mass ratio that decreases from $\approx 70\%$ to $\approx 40\%$ going from r_{500} to the virial radius. These values appear to be at variance with numerical simulations that report a hydrostatic-to-true total mass ratio of $\approx 80\%$ – 90% at the virial radius (Lau et al. 2013; Nelson et al. 2014).

We stress that the comparison between X-ray and weak lensing masses is fundamental for understanding the ICM physical state. In particular, the difference in the mass values obtained with these independent methods probes the level of nonthermal pressure support needed to sustain hydrostatic equilibrium. This additional pressure component may be due to turbulence originated by several processes such as mergers, supersonic motions of galaxies through the ICM, or infall of gas into the cluster from the surrounding environment (see simulations by Nagai et al. 2007; Shaw et al. 2010; Burns et al. 2010; Vazza et al. 2011; Rasia et al. 2012). We remark that such a nonthermal pressure support must be taken into account to improve both our astrophysical understanding of cluster outskirts and the accuracy in cluster mass determination for cluster cosmology (see Vikhlinin et al. 2009b).

The paper is organized as follows. In Section 2 we recall the formalism of our entropy-based SuperModel (SM; Cavaliere et al. 2009), that allows a self-consistent analysis of the X-ray observables in presence of nonthermal pressure support. In Section 3 we exploit the SM to analyze A1246 at $z = 0.19$ and J255 at $z = 0.45$ basing on the X-ray observations by *Suzaku* and *Chandra*, respectively. In Section 4 we discuss the results and draw our conclusions. Throughout the paper we adopt the standard flat cosmology (Hinshaw et al. 2013; *Planck* Collaboration et al. 2014) with parameters in round numbers: $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_\Lambda = 0.73$, $\Omega_M = 0.27$. Then 1 arcmin corresponds to 191 kpc for A1246, and to 349 kpc for J255.

2. SUPERMODEL WITH TURBULENCE

We briefly recall the entropy-based SM formalism in presence of a nonthermal component. We write the total pressure $p_{\text{tot}}(r) = p_{\text{th}}(r) + p_{\text{nth}}(r) = p_{\text{th}}(r)[1 + \delta(r)]$ in terms of the nonthermal to thermal ratio $\delta(r) \equiv p_{\text{nth}}/p_{\text{th}}$. Expressing the density in terms of temperature and entropy via $n \propto (T/k)^{3/2}$ and using the equation of hydrostatic equilibrium yields the temperature profile in the form

$$\frac{T(r)}{T_R} = \left[\frac{k(r)}{k_R} \right]^{3/5} \left[\frac{1 + \delta_R}{1 + \delta(r)} \right]^{2/5} \left\{ 1 + \frac{2}{5} \frac{b_R}{1 + \delta_R} \times \int_r^R \frac{dx}{x} \frac{v_c^2(x)}{v_R^2} \left[\frac{k_R}{k(x)} \right]^{3/5} \left[\frac{1 + \delta_R}{1 + \delta(x)} \right]^{3/5} \right\}; \quad (1)$$

here $v_c^2(r) \equiv G M(< r)/r$ is the squared circular velocity (v_R^2 is the value at the virial radius R), and b_R is the ratio of v_R^2 to the squared sound speed at R (Cavaliere et al. 2009, 2011b).

For the spherically averaged entropy profile $k(r)$ we consider the basic pattern $k(r) = k_c + (k_R - k_c)(r/R)^a$ including a central floor $k_c \sim 10\text{--}100 \text{ keV cm}^2$ (e.g., Pratt et al. 2010; or even $< 10 \text{ keV cm}^2$ for low- z cool core clusters, see Panagoulia et al. 2014) going into an outer power-law rise with slope $a \sim 1$ out to the virial value $k_R \sim \text{some } 10^3 \text{ keV cm}^2$ (see Voit 2005; Lapi et al. 2005). However, to model a possible flattening in the outer region, we modify the profile beyond a break radius r_b by allowing the slope a to change (Lapi et al. 2010); for the sake of simplicity, the entropy slope is taken to decline linearly with a gradient $a' \equiv (a - a_R)/(R/r_b - 1)$. The quantities k_c , a , r_b and a' are free parameters to be determined with their uncertainties from fitting the projected, spherically averaged X-ray observables, taking into account errorbars on both axes (i.e., radius and density/temperature); to this purpose, we exploit the physics analysis tool of the function minimization algorithm MINUIT.⁷

As to the profile of $\delta(r)$ we follow the prescription by Cavaliere et al. (2011b) based on the classic theory of turbulence generation (see Kolmogorov 1941; Monin & Yaglom 1965; Inogamov & Sunyaev 2003; Petrosian & East 2008; Brunetti & Lazarian 2011) and on indications from hydrodynamical simulations (Lau et al. 2009; Vazza et al. 2011). We adopt a shape $\delta(r) = \delta_R e^{-(R-r)^2/\ell^2}$ decaying on the scale ℓ inward of a round maximum δ_R at the boundary shock where the nonthermal component originates. The quantities δ_R and ℓ are held fixed during the fitting procedure.

The traditional equation to estimate the total gravitational mass $M(< r)$ within r must be modified to take into account the nonthermal pressure component; the outcome reads (Fusco-Femiano & Lapi 2013)

$$M(< r) = -\frac{k_B\{T(r)[1 + \delta(r)]\} r^2}{\mu m_p G} \left[\frac{1}{n_e(r)} \frac{dn_e(r)}{dr} + \frac{1}{T(r)} \frac{dT(r)}{dr} + \frac{\delta(r)}{1 + \delta(r)} \frac{2}{\ell^2} (R - r) \right]. \quad (2)$$

The ICM mass writes $M_{\text{gas}} = 4\pi\mu_e m_p \int dr n_e(r) r^2$ where $\mu_e \approx 1.16$ is the mean molecular weight of the electrons.

3. SUPERMODEL ANALYSIS FOR A1246 AND J255

In this section we present the SM analysis of A1246 and J255, by fitting the spherically averaged electron density and temperature profiles measured by *Suzaku* (Sato et al. 2014) and *Chandra* (Wang & Walker 2014), respectively.

3.1. A1246

A1246 is a cluster of galaxies at redshift $z = 0.1902$ (NASA/IPAC Extragalactic Database) that features a regular ICM distribution, as reported by *Suzaku* observations in the 0.5–5.0 keV energy range (Sato et al. 2014). The thermal emission is significantly detected out to r_{200} and the temperature at this radius is a factor ≈ 3 lower than at the peak (see Figure 1). We assume a virial radius $R = 2r_{500}$ with $r_{500} \approx 6.1$ as derived by Sato et al. (2014) from the mass value obtained under the assumption of thermal hydrostatic equilibrium. Note that at r_{500} thermal hydrostatic equilibrium holds to a very good accuracy; this is also shown by Okabe et al. (2014), that in their joint X-ray/weak-lensing study of four relaxed clusters

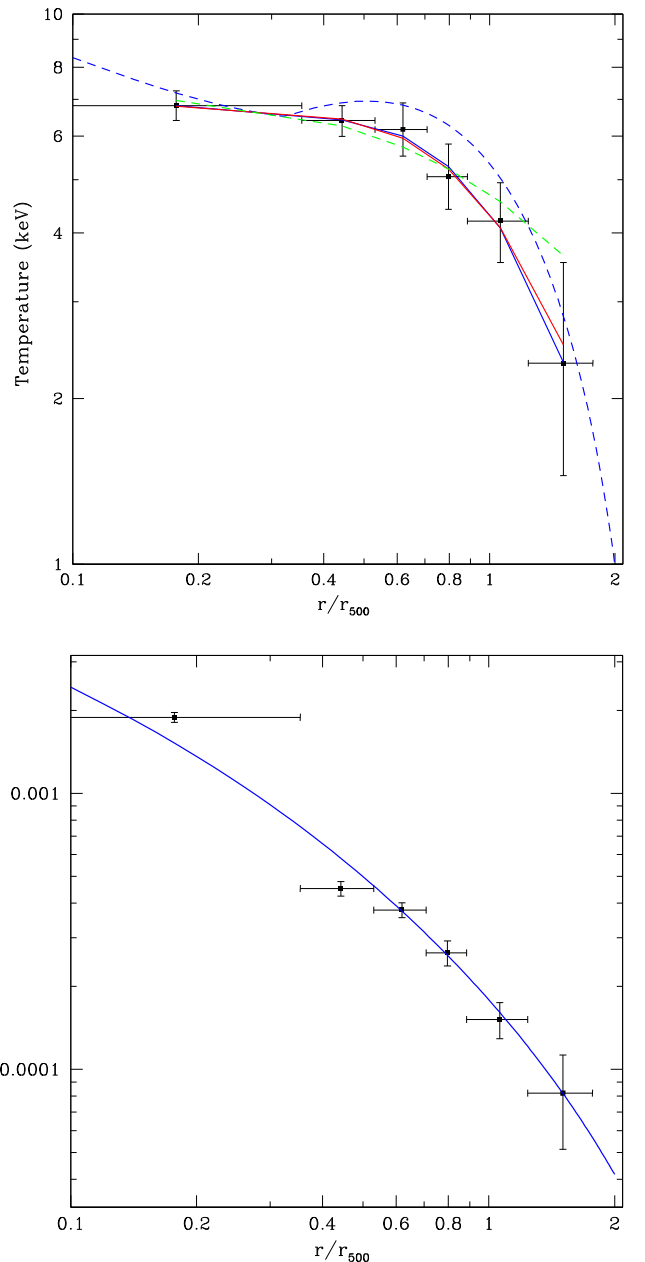


Figure 1. A1246—Top panel: projected temperature profiles. Data points from *Suzaku* (Sato et al. 2014). Blue and red lines illustrate the SM fits without ($\delta_R = 0$) and with nonthermal pressure component ($\delta_R = 0.3$, $\ell = 0.5$; see Section 2 for details), respectively. Both lines are obtained with a flattening entropy profile, while the dashed green line is obtained with a simple power-law shape; the dashed blue line shows the deprojected temperature profile corresponding to the blue line. Bottom panel: the blue line illustrates the SM fit to the electron density data points from *Suzaku* (Sato et al. 2014). We assume $r_{500} = 6.1$ as estimated by Sato et al. (2014) under the assumption of thermal hydrostatic equilibrium (see the discussion in Section 3.1).

report $dM/dr < 0$ at $r \gtrsim 1.3 r_{500}$, in agreement with previous analyses (e.g., Kawaharada et al. 2010; Ichikawa et al. 2013). A posteriori, we have also checked that our derived virial mass is consistent with the adopted value of R in yielding a mean density ≈ 100 times the critical one.

The azimuthal analysis by Sato et al. (2014) in four directions reports a slightly lower temperature in the southeast sector. Assuming spherical symmetry and a constant temperature in each annular region it is found that the de-projected electron density of the northeast and southwest sectors tends to be

⁷ See <http://seal.web.cern.ch/seal/snapshot/work-packages/mathlibs/minuit/>.

lower than that of the southeast and northwest sectors in the radial range $r \approx (0.9\text{--}1.8)r_{500}$. The derived gravitational mass, azimuthally averaged, starts flattening and then decreasing beyond r_{500} , thus indicating a break in the assumption of thermal hydrostatic equilibrium. In particular, Sato et al. (2014) found a mass of $(4.3 \pm 0.4) \times 10^{14} M_{\odot}$ within r_{500} , consistent with that of $(3.9 \pm 0.1) \times 10^{14} M_{\odot}$ obtained from *Chandra* observations by Vikhlinin et al. (2009a) through the $M_{500} - T_X$ scaling relation. In Sato et al. (2014) the X-ray mass profile is compared with that derived from the stacked weak lensing analysis by Okabe et al. (2010). The weak lensing mass amounts to ≈ 5 and $\approx 7.8 \times 10^{14} M_{\odot}$ within r_{500} and r_{200} , respectively. The X-ray and weak lensing masses are consistent within r_{500} but not at greater radii where the former is appreciably lower than the latter. Relatedly, Sato et al. (2014) found a gas mass fraction consistent with the cosmic baryon fraction at r_{500} but not at r_{200} ; on the other hand, the weak lensing mass at r_{200} yields instead a baryon fraction in agreement with the cosmic value.

The entropy profile of A1246 is similar to that reported for all the clusters observed out to the virial radius by *Suzaku* (Bautz et al. 2009; Kawaharada et al. 2010; Hoshino et al. 2010; Akamatsu et al. 2011; Walker et al. 2012; Ichikawa et al. 2013). The entropy increases with radius up to $r \approx r_{500}$ and then flattens outward. As reported also for A1835 and A1689 (see Ichikawa et al. 2013), in A1246 the entropy flattening is more pronounced in some cluster sectors; in particular, it is more evident in the southeast sector rather than in the northwest one, that appears to face a filament. Besides, in this latter sector the temperature is higher than those measured in other regions (as it happens in A1689, see Kawaharada et al. 2010), whereas the electron densities are consistent.

We fit the spherically averaged X-ray temperature profile of A1246 with the SM finding that, within the measurement uncertainties, a flattening entropy profile performs better than a simple power-law (see Figure 1). The corresponding SM fit to the gas density $n_e(r)$ is reported in the bottom panel of Figure 1. We obtain a central entropy value $k_c = 101 \pm 15$ keV cm² rather high for a relaxed cluster; this is likely due to the low spatial resolution of *Suzaku*, which is insufficient to resolve the expected presence of a cool core in the relaxed cluster A1246. For the other parameters we find $a = 0.95^{+0.22}$, $r_b = (0.53 \pm 0.07)r_{500}$ and $a' = -(1.22 \pm 0.21)$.

We compute the total gravitational mass within r (see the blue line of Figure 2) from Equation (2) by assuming thermal hydrostatic equilibrium ($\delta = 0$). In agreement with the analysis of Sato et al. (2014) the mass profile starts to flatten at $r \approx r_{500}$ and then declines downward at $r \gtrsim 1.6r_{500}$, yielding a gas mass fraction greater than the cosmic value at the virial radius (blue line of Figure 2). The decreasing mass profile clearly indicates that the outskirts of A1246 are not in thermal hydrostatic equilibrium. We evaluate the level and radial shape of the nonthermal component by requiring the X-ray mass to be in agreement (as suggested by Sato et al. 2014) with the stacked weak lensing mass profile from Okabe et al. (2013, updating Okabe et al. 2010). We find that the agreement is recovered for $\delta_R = 0.3$ and $\ell = 0.5$, implying a nonthermal pressure component $\approx 20\%$ of the total pressure at the virial radius. We note that our SM analysis gives a mass value consistent at r_{500} with the revised estimate of Okabe et al. (2013), and with the value obtained by Sato et al. (2014); this cross-checks the consistency in our adopted value of the virial radius. Despite the presence of a nonthermal pressure component that gives a higher total mass consistent with the weak lensing determination, the

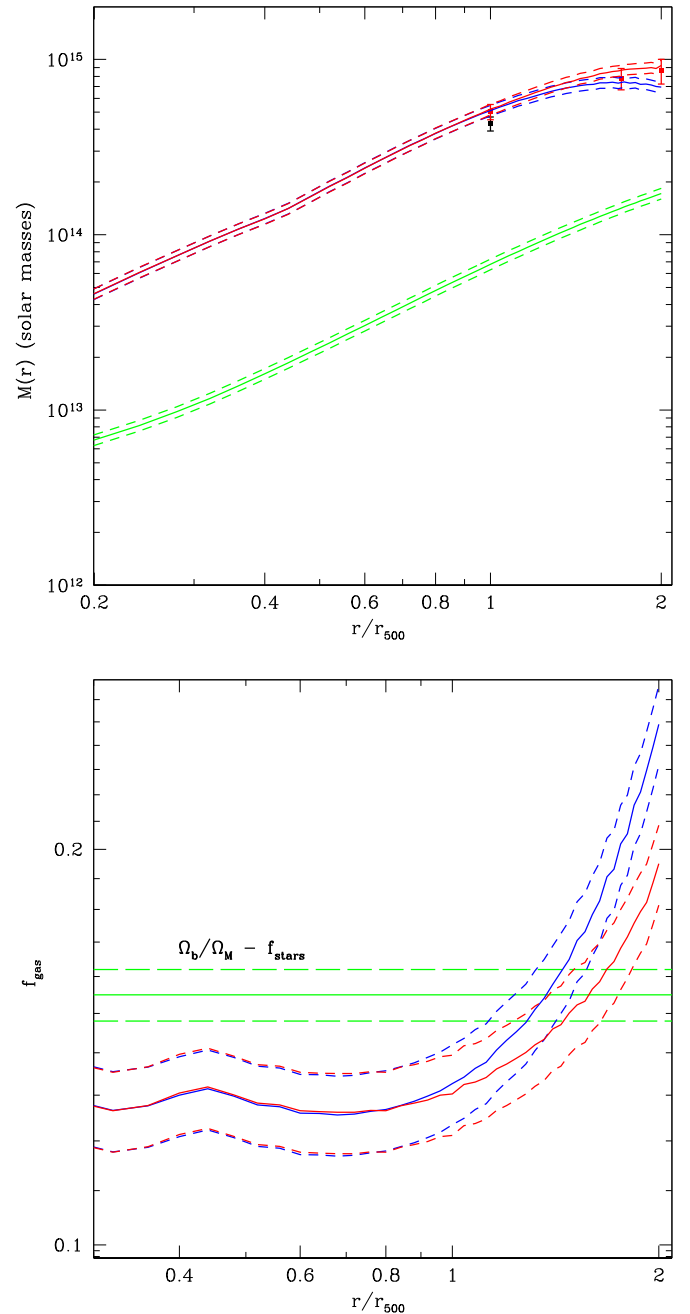


Figure 2. A1246—Top panel: mass profiles. The blue line illustrates the X-ray cluster mass obtained with $\delta_R = 0$, while the red line is obtained with $\delta_R = 0.3$ and $\ell = 0.5$; the green line illustrates the gas mass obtained by the gas density from *Suzaku* (Sato et al. 2014). The dashed lines mark the 1σ uncertainty region from the SM fit. The black point shows the mass derived at r_{500} by Sato et al. (2014), while the red points refer to the weak lensing mass values from the stacking analysis by Okabe et al. (2013, updating Okabe et al. 2010). Bottom panel: gas mass fraction derived from the above mass profiles. Blue line is with $\delta_R = 0$; red line is with the above values of δ_R and ℓ ; green lines show the measured difference of the cosmic baryon fraction and the fraction of baryons in stars and galaxies, $\Omega_b/\Omega_M - f_{\text{stars}} = 0.155 \pm 0.007$ (Komatsu et al. 2011; Gonzalez et al. 2007).

gas mass fraction remains higher than the cosmic value (see red line of Figure 2), highlighting the presence of gas clumping with a factor $C \approx 1.6$ at the virial radius.

Figure 3 shows the entropy profiles obtained by the SM analysis, confronted with that observed by *Suzaku*. We again find that the flattening profile matches better the one derived by Sato et al. (2014). With the above value of clumping, the entropy

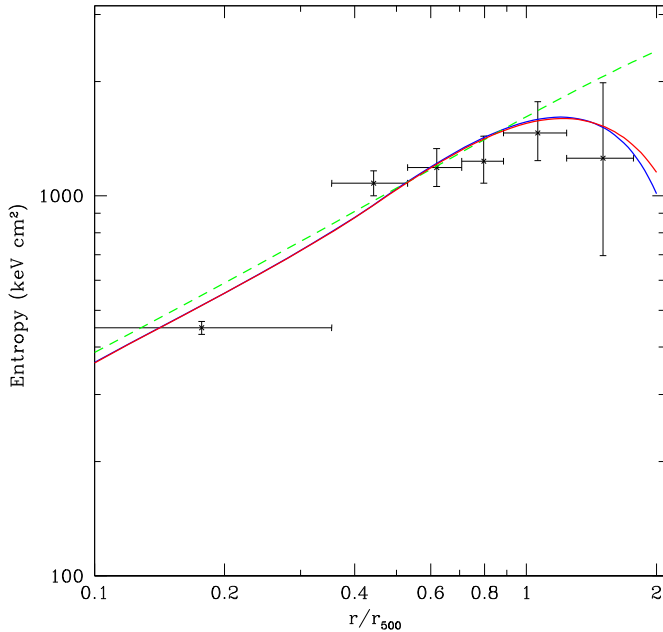


Figure 3. A1246—Entropy profiles. The blue line illustrates the entropy profile obtained with the deprojected temperature profile (dashed blue line) and the gas density profile in Figure 1. The red line shows the entropy profile obtained with the deprojected temperature profile relative to $\delta_R = 0.3$ and $\ell = 0.5$ (see the red line of Figure 1). The dashed green line is obtained by the fit to the temperature data with a power-law increase of the entropy. Data points show the entropy values obtained from the analysis by Sato et al. (2014).

is underestimated only by a factor $C^{1/3} \approx 1.2$ at the virial radius, insufficient to explain the observed entropy flattening. This implies that also for A1246 the entropy flattening in the outskirts is mainly related to the rapid decline of the temperature, and not to an overestimate of the density caused by gas clumping; this is in agreement with the results recently reported by Fusco-Femiano & Lapi (2014) for a sample of other four clusters.

3.2. J255

To investigate the ICM in the cluster outskirts a feasible possibility is given by deep *Chandra* observations of distant clusters at redshift $\gtrsim 0.2$. In particular, diffuse emission has been observed out to the virial radius in A1835 (Bonamente et al. 2013), showing a sharp decline of the temperature consistent with the *Suzaku* observations (Ichikawa et al. 2013).

The internal regions of J255 present an elongated X-ray structure most likely due to a recent merger event, while the cluster appears relaxed at distance $r \gtrsim 1'$. This cluster shows a projection overlap with the X-ray emission of A2246 at $z = 0.23$, that appears evident only in a limited sector (Wang & Walker 2014). A cool core is also detected (see Figure 4). The conservative extent of the cluster is $\approx 1.1 r_{200}$ where $r_{200} = (4 \pm 0.2)'$ is estimated by the average temperature (5.5 ± 0.4 keV) using the scaling relations by Arnaud et al. (2005); the total mass M_{200} amounts to $5.0^{+0.8}_{-0.7} \times 10^{14} M_{\odot}$.

We perform a fit with the SM to the de-projected temperature profile derived by Wang & Walker (2014), assuming spherical symmetry and a virial radius $R = 4/3 r_{200}$. Despite the central value of the outermost bin is limited to $\sim r_{500}$, it is evident that the observed temperature profile does not show the rapid decline observed in other clusters, and is well fitted by a simple power-law increase of the entropy (see Figure 4) $k \propto r^a$ with $a = 0.98^{+0.2}$ (and $k_c = 12.5 \pm 3.1$ keV cm²), in agreement with

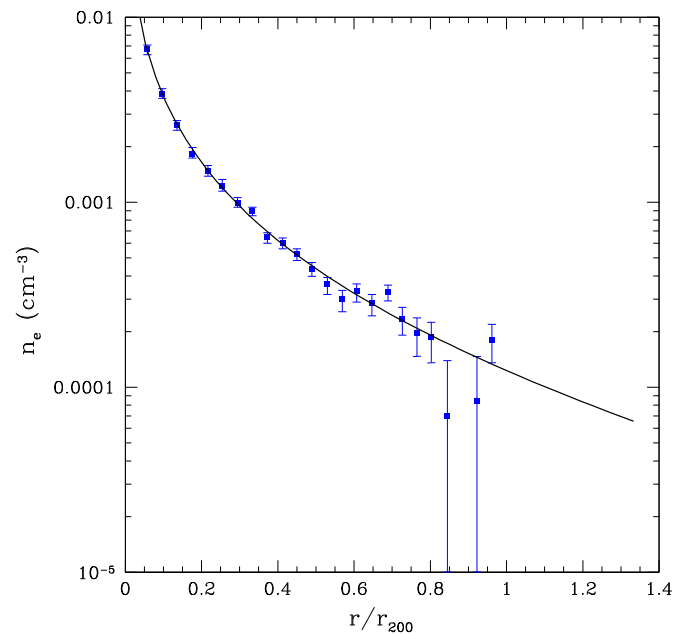
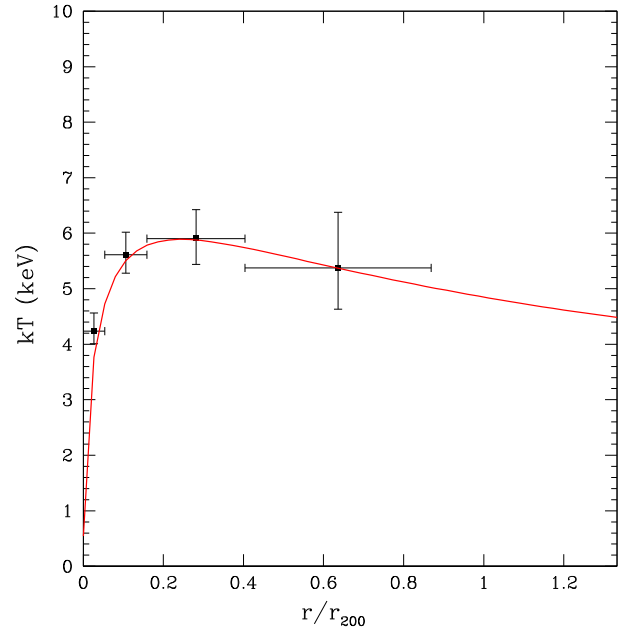


Figure 4. J255—Top panel: temperature profile. The deprojected data points are derived by *Chandra* observations (Wang & Walker 2014). The red line illustrates the fit obtained with an entropy profile following a simple power-law increase; this fit is indistinguishable from that obtained with entropy flattening. Bottom panel: the black line shows the gas density profile obtained by the SM fit to the points derived by *Chandra* observations (Wang & Walker 2014). We assume $r_{200} = 4'$ as estimated by Wang & Walker (2014) using the scaling relations by Arnaud et al. (2005).

the expected value of 1.1 from strong-shocked accretion under pure gravitational infall. The corresponding fit to the gas density profile observed by *Chandra* is reported in Figure 4.

The shape of the inferred entropy profile and the morphologically disturbed features of J255 concur to suggest that the cluster has undergone a recent merger event; on the other hand, the presence of a cool core implies that either the merger has not reached the core yet, or it has only partially destroyed the cool core (Fusco-Femiano et al. 2009; Rossetti & Molendi 2010).

We find that the total cluster mass (see Figure 5) from the SM is consistent with the M_{200} value estimated by Wang & Walker

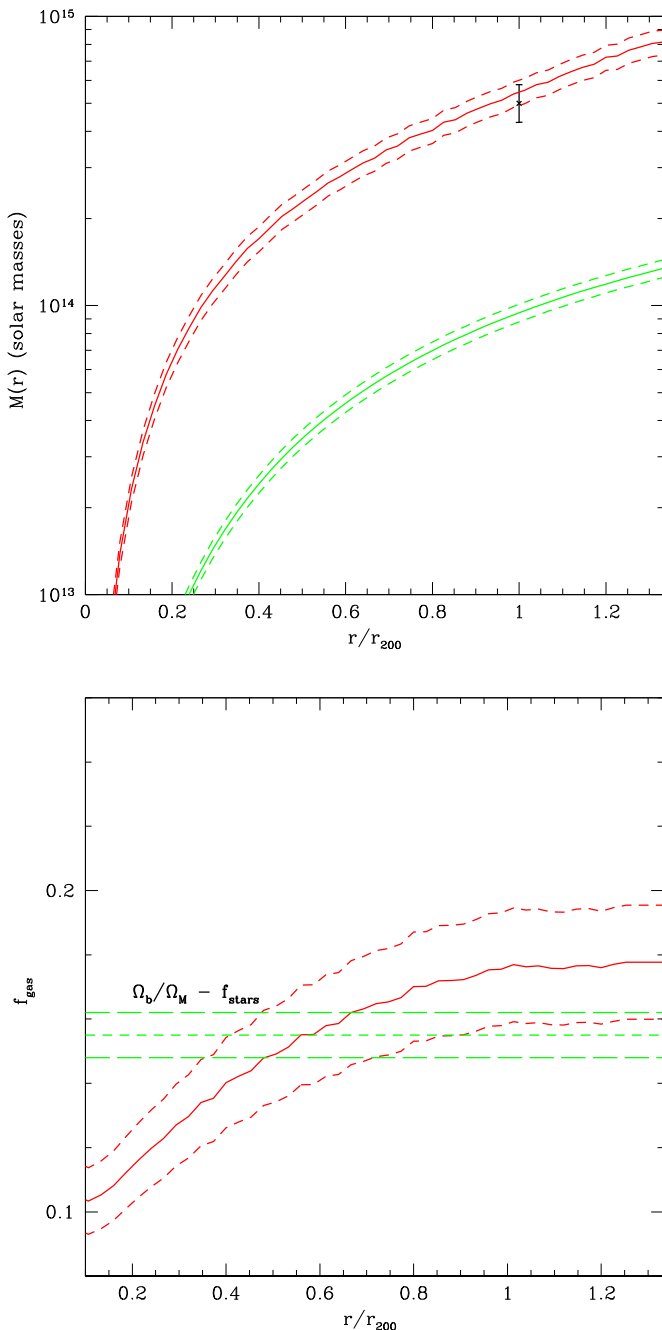


Figure 5. J255—Top panel: mass profiles. The red line illustrates the mass profile obtained with the gas temperature and density profiles of Figure 4; the data point is the value of M_{200} obtained by Wang & Walker (2014) using scaling relations with the temperature by Arnaud et al. (2005); the green line shows the gas mass profile. The dashed lines mark the 1σ uncertainty region from the SM fit. Bottom panel: the red curve is obtained with the mass profiles reported in the top panel. Green lines are the same as in Figure 2.

(2014) and with a mean density that at R is ≈ 100 times the critical one, so cross-checking the consistency in our assumed value of the virial radius; the overall mass corresponds to a gas mass fraction at R slightly above the cosmic value. This may be explained either by a gas clumping factor $C \approx 1.3$, or by a modest level of turbulence ($p_{\text{nth}} \approx 5\% p_{\text{tot}}$).

4. DISCUSSION AND CONCLUSIONS

Our SM analysis of the X-ray temperature and brightness profiles can shed light on the dynamical state of the ICM;

specifically, it enables us to determine the shape of the entropy profile throughout the cluster volume, from the inner core region out to the outskirts. In the latter, our analysis provides both a measure of the nonthermal pressure support needed to sustain hydrostatic equilibrium, and an estimate of gas clumping. In this paper we have applied the SM analysis to the two clusters A1246 at $z = 0.19$ and J255 at $z = 0.45$, that have been recently observed out to the virial radius by *Suzaku* and *Chandra*, respectively.

We have determined the entropy profile of A1246 by fitting with the SM the *Suzaku* temperature data (see Figure 1); we have found that the entropy progressively flattens outward relative to the simple power-law increase $k \propto r^{1.1}$ expected from strong-shocked accretion under pure gravitational infall; this is in accordance with the entropy analysis by Sato et al. (2014).

We have determined the X-ray mass on assuming pure thermal hydrostatic equilibrium, finding results consistent with the revised stacked cluster lensing measurements by Okabe et al. (2013) at r_{500} . On the other hand, the determination at the virial radius is biased low due to the break of thermal hydrostatic equilibrium, related to the rapid decrease of the temperature as also found by Okabe et al. (2014); quantitatively, we derive a bias of $\approx 20\%$ consistent with the X-ray to true total mass ratio within r_{200} found by hydrodynamical simulations (see Lau et al. 2013; Nelson et al. 2014), but appreciably lower than the result obtained by Okabe et al. (2014) using a joint X-ray and weak lensing analysis for a sample of four relaxed clusters.

We have confirmed previous analysis (Sato et al. 2014) in finding that the mass profile reconstructed from X-ray observables features in the outskirts an unphysical, non-monotonic behavior. This is likely due to the break of thermal hydrostatic equilibrium in the outskirts. Thus we have exploited the capability of our SM to include a nonthermal pressure component. We have quantified such nonthermal levels by requiring the reconstructed mass to be consistent with the weak lensing determination (see Figure 2); we find a nonthermal pressure component of about 20% of the total at the virial radius.

We have also estimated the level of gas clumping by matching the resulting gas mass fraction to the cosmic value; we find a clumping factor $C \approx 1.6$ at the virial radius. Modest values of C are consistent with our previous analysis of relaxed clusters (see Fusco-Femiano & Lapi 2013, 2014) mainly based on gas density profile measured by *ROSAT*, which by its coarser resolution is less sensitive to clumping. Our findings also agree with the bounds $C \lesssim 2$ at the virial radius from numerical simulations (see Nagai & Lau 2011; Vazza et al. 2013; Battaglia et al. 2014). Moreover, the modest levels of gas clumping imply that the entropy at the virial radius may be underestimated only by a factor $C^{1/3} \approx 1.2$, in agreement with our previous analysis of relaxed clusters and with the more recent study by Okabe et al. (2014). This indicates that the entropy flattening in the outskirts is strictly related to the steeply declining temperature profile and not to an overestimate of the density because of clumping.

The shape of the entropy profile strongly suggests that A1246 is a relaxed cluster that probably had time to develop a cool core in the inner region, though it cannot be resolved because of the insufficient spatial resolution of *Suzaku*. The indication of the relaxed nature of A1246 is confirmed by the physical correlation between the outskirts entropy and the virial mass suggested by Walker et al. (2012); Sato et al. (2012) and recently investigated by Okabe et al. (2014) using the joint X-ray/weak lensing analysis of four relaxed clusters. Specifically, this correlation is between the average entropy K_{out} in the range $r_{500} - R$, and

$M_{\text{vir}} E(z)$, with $E(z) = [\Omega_{m,0}(1+z)^3 + \Omega_{\Lambda}]^{1/2}$ the Hubble expansion rate. For A1246 we found an average entropy K_{out} of $\approx 1300 \text{ keV cm}^2$ at $M_{\text{vir}} E(z) [10^{14} M_{\odot}] \approx 10$ consistent with the tight correlation found by Okabe et al. (2014) for relaxed clusters.

This implies that the actual weak lensing mass profile of A1246 is not much different from that adopted by us on the basis of Okabe et al. (2013, updating Okabe et al. 2010); thus our derived values of δ_R and C are not significantly affected. The azimuthal analysis performed by Sato et al. (2014) reveals the entropy flattening to be more pronounced in the cluster sectors facing low-density, void-like environments, while to be almost absent in the sectors facing high-density filaments of the cosmic web.

These findings agree with previous analyses on other cool-core, relaxed clusters, and lend further support to the picture by Lapi et al. (2010) and Cavaliere et al. (2011a) that relates the entropy flattening and the development of nonthermal pressure component to weakening boundary shocks; the latter produce less entropy, while allowing more bulk inflow energy to seep inside and develop nonthermal pressure in the form of turbulence. The shock weakening mainly occurs under two conditions: (1) either at late time in cluster evolution when external gas is accreted from the wings of the initial perturbation, and especially so at low redshift in an accelerating background universe; (2) or in a particular sector of a cluster facing an underdense, void-like region. In this scenario clusters are born in a high-entropy state throughout, and are expected to develop synchronously a low-entropy state both at the center due to cooling, and in the outskirts due to weakening shocks.

However, the analysis of J255 testifies how such a typical evolutionary course can be interrupted or even reversed by merging, especially at intermediate redshift. We have determined the entropy profile by fitting with SM the azimuthally average temperature profile observed by *Chandra*. We have found that the entropy steadily increase with radius as a power-law $k \propto r^{1.1}$, close to the expectation for strong shocks (see Figure 4).

We have also determined the mass profile (see Figure 5), finding a value of $M_{200} \approx 5.6 \pm 0.4 \times 10^{14} M_{\odot}$ in agreement with that derived by Wang & Walker (2014) using the mass–temperature scaling relation (Arnaud et al. 2005). The gas mass fraction f_{gas} at the virial radius is slightly greater than the cosmic value (see Figure 5). This can be easily explained by a clumping factor $C \approx 1.3$ or by a very modest level around 5% of nonthermal pressure; the latter would imply a total mass in excess of $\approx 7\%$ than the value reported in Figure 5 at R .

The value of the average outer entropy $K_{\text{out}} \approx 2140 \text{ keV cm}^2$ at $M_{\text{vir}} E(z) [10^{14} M_{\odot}] \approx 10$ is much larger than the value derived for relaxed clusters; this concurs with the power-law shape of the entropy profile in indicating that J255 as an unrelaxed cluster. This is also confirmed by the interesting morphological structure of J255. The cluster features a cool core, but with an elongated X-ray morphology in the inner regions, most likely due to a recent merger event.

So the typical evolutionary course from high- to low-entropy state expected to occur simultaneously both in the core due to cooling and in the outskirts by reduced entropy production in weakening shocks has been interrupted by a merger (see Cavaliere et al. 2011a; Cavaliere & Lapi 2013). This event may have indeed reheated locally the ICM, rejuvenating the cluster outskirts to a high-entropy state but leaving intact or destroying only partially the low-entropy cool core (Fusco-Femiano et al. 2009; Rossetti & Molendi 2010). In fact, such a behavior

is in pleasing agreement with the predictions by Cavaliere et al. (2011a).

Finally, we stress that the analysis of these two clusters have yielded a further evidence that the entropy content in the outskirts is jointly determined by the dynamical state of the cluster and by the surrounding environment.

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