

Dynamic characteristic analysis of two-stage quasi-zero stiffness vibration isolation system

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Abstract. A novel two-stage quasi-zero stiffness (QZS) vibration isolator was proposed for the purpose of low-frequency vibration isolation. Firstly, the dynamic model of the vibration isolation system was established; furthermore, the force transmissibility of the system under harmonic force excitation was derived by the averaging method; finally, the effects on the vibration isolation performance caused by excitation amplitude, mass ratio and damping ratio were discussed. Results show that, compared with the corresponding two-stage linear system, two-stage QZS system not only has better isolation performance, but also possesses a wider range of isolation frequency provided that the excitation amplitude, mass ratio and damping ratio is appropriate.

Keywords: two-stage quasi-zero stiffness, the averaging method, force transmissibility, vibration isolation system.

1. Introduction

Quasi-zero stiffness (QZS) isolator can obtain zero stiffness at the static equilibrium position by connecting a positive stiffness element in parallel with a negative stiffness element [1]. By reasonably selecting the geometry and stiffness parameters of the negative stiffness institutions, the vibration isolation system can support a large load statically while have low-frequency vibration isolation performance. The combination of positive and negative stiffness device possesses the characteristics of high-static-low-dynamic stiffness and can ensure that the system natural frequency is very low at small deformation [2]. Therefore, scholars have conducted in-depth research on isolator principle [3], structure design [4], properties analysis [5] and engineering application [6] of QZS, which is widely used in precision instrument isolation, bridge and house anti-earthquake, high-speed vehicle vibration reduction, marine machinery noise cancellation and other fields.

With the development of science and technology, the request of stability for the work environment of precision equipment is more and more strictly. Conventional single-stage linear or QZS vibration isolation system cannot meet requirements of industry [7, 8]. When the excitation frequency is greater than $\sqrt{2}$ times the natural frequency, vibration decay ratio is proportional to ω^{-2} in the single-stage system, while that is proportional to ω^{-4} in the two-stage system. So rigid two-stage isolation system can replace soft single-stage isolation system, which not only has good isolation effect, but also takes into account the load-bearing capacity and stability. Although the two-stage isolator is better isolation performance than the single-stage isolator, but the former has one more resonance peak and longer resonance time. To play the single QZS advantages, also take large vibration decay ratio of two-stage isolation system into consideration. Two-stage QZS vibration isolator was proposed, whose effects on the vibration isolation performance caused by excitation amplitude, mass ratio and damping ratio were studied.

2. Mechanical model of QZS vibration isolation system

Fig. 1 shows a typical QZS system. The system consists of a suspended mass m with a vertical spring and two identical oblique springs.

The force-deflection in the vertical direction is given by:

$$F_v = k_v x + 2k_h \left(1 - \frac{l_0}{\sqrt{x^2 + l^2}} \right) x, \quad (1)$$

where l_0 is the free length of horizontal springs and l is length in the horizontal position. x is the displacement deviation from equilibrium position. The stiffness of oblique spring is k_h and the stiffness of horizontal spring is k_h .

Eq. (1) can be written in non-dimensional form as:

$$\hat{f} = \frac{F_v}{k_v x_0} = \hat{x} + 2\hat{k} \left(1 - \frac{1}{\sqrt{(\hat{x}^2(1 - \hat{l}^2) + \hat{l}^2)}} \right) \hat{x}, \quad (2)$$

where $\hat{x} = x/x_0$, $\hat{l} = l/l_0$ and $\hat{k} = k_h/k_v$. Using the Maclaurin-series expansion to the third order for small \hat{x} , then Eq. (2) approximates to:

$$\hat{f} = k_a \hat{x} + k_b \hat{x}^3, \quad (3)$$

where $k_a = 1 - 2\hat{k}(1 - \hat{l})/\hat{l}$ and $k_b = \hat{k}(1 - \hat{l}^2)/\hat{l}^3$.

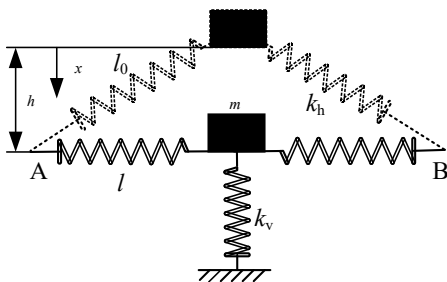


Fig. 1. Schematic diagram of the QZS vibration isolation system

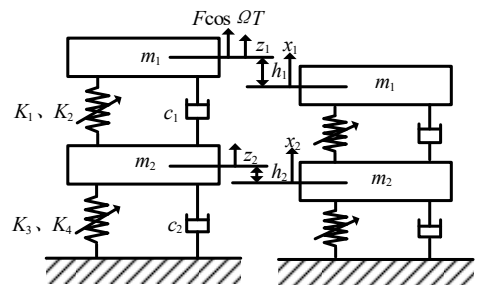


Fig. 2. Schematic of a two-stage QZS system

3. Two-stage QZS vibration isolation system

The two-stage QZS system is shown in Fig. 2.

It is subject to external harmonic force excitation $F \cos \Omega T$. The upper QZS isolator stiffness coefficient is K_1, K_2 and damping coefficient is c_1 . The lower QZS isolator stiffness coefficient is K_3, K_4 and damping coefficient is c_2 . z_1 and z_2 are the displacement of the vibration object and the middle inertia block when the respective spring in a natural state. m_1 and m_2 are the mass of vibration object and middle inertia block respectively.

Provided that only consider the vertical direction of movement, dynamic equation of two-stage QZS system is established based on Newton's second law:

$$\begin{cases} m_1 \ddot{z}_1 + c_1(\dot{z}_1 - \dot{z}_2) + K_1(z_1 - z_2) + K_2(z_1 - z_2)^3 = F \cos(\Omega T) - m_1 g, \\ m_2 \ddot{z}_2 - c_1(\dot{z}_1 - \dot{z}_2) - K_1(z_1 - z_2) - K_2(z_1 - z_2)^3 + c_2 \dot{z}_2 + K_3 z_2 + K_4 z_2^3 = -m_2 g. \end{cases} \quad (4)$$

For clarity of analysis, the parameters $\Omega_0 = \sqrt{K_1/m_1}$, $z_1 = X_1 \sqrt{K_1/K_2}$, $z_2 = X_2 \sqrt{K_1/K_2}$, $T = \Omega_0 t$ are introduced. Eq. (4) can be written in non-dimensional form as:

$$\begin{cases} \ddot{X}_1 - \xi_1(\dot{X}_2 - \dot{X}_1) - (X_2 - X_1) - (X_2 - X_1)^3 = f \cos(\omega t) - G, \\ w \ddot{X}_2 + \xi_1(\dot{X}_2 - \dot{X}_1) + (X_2 - X_1) + (X_2 - X_1)^3 + \xi_2 \dot{X}_2 + k_1 X_2 + k_2 X_2^3 = -wG, \end{cases} \quad (5)$$

where $\xi_1 = c_1/\sqrt{K_1 m_1}$, $f = F/K_1 \sqrt{K_2/K_1}$, $G = m_1 g/K_1 \sqrt{K_2/K_1}$, $k_1 = K_3/K_1$, $k_2 = K_4/K_2$, $k_3 = K_2/K_1$, $\xi_2 = c_2/\sqrt{K_1 m_1}$, $w = m_2/m_1$, $\omega = \Omega_0 \Omega$.

In order to analyze conveniently, the first and second equation of Eq. (5) are added, which can make stiffness coupling change into inertia coupling. The gravity term in Eq. (5) can be eliminated by using coordinate transform. Introducing $X_1 = Z_1 - h_1$, $X_2 = Z_2 - h_2$, $H = h_1 - h_2$, $x_1 = Z_2 - Z_1$, $x_2 = Z_2$, Eq. (5) can be transferred as:

$$\begin{cases} \ddot{x}_2 - \ddot{x}_1 - \xi_1 \dot{x}_1 - (1 + 3H^2)x_1 - 3Hx_1^2 - x_1^3 = f \cos(\omega t), \\ (1 + w)\ddot{x}_2 - \ddot{x}_1 + \xi_2 \dot{x}_2 + (k_1 + 3k_2 h_2^2)x_2 - 3k_2 h_2 x_2^2 + k_2 x_2^3 = f \cos(\omega t), \end{cases} \quad (6)$$

where $H + H^3 = G$, $k_1 h_2 + k_2 h_2^3 = (w + 1)G$.

Eq. (6) in matrix form is:

$$\mathbf{M}\ddot{\mathbf{X}} + \mathbf{C}\dot{\mathbf{X}} = \mathbf{F}, \quad (7)$$

where:

$$\mathbf{M} = \begin{bmatrix} -1 & 1 \\ -1 & 1 + w \end{bmatrix}, \quad \mathbf{C} = \begin{bmatrix} -\xi_1 & 0 \\ 0 & \xi_2 \end{bmatrix}, \quad \mathbf{X} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix},$$

$$\mathbf{F} = \begin{bmatrix} f \cos(\omega t) + (1 + 3H^2)x_1 + 3Hx_1^2 + x_1^3 \\ f \cos(\omega t) - (k_1 + 3k_2 h_2^2)x_2 + 3k_2 h_2 x_2^2 - k_2 x_2^3 \end{bmatrix}.$$

By the averaging method, suppose steady-state response solution of the above system as:

$$\begin{cases} \mathbf{X} = \mathbf{U} \cos(\omega t) + \mathbf{V} \sin(\omega t), \\ \mathbf{X}' = -\omega \mathbf{U} \sin(\omega t) + \omega \mathbf{V} \cos(\omega t), \end{cases} \quad (8)$$

where $\mathbf{U} = [U_1, U_2]^T$ and $\mathbf{V} = [V_1, V_2]^T$ change slowly with time. Differentiate the first equation of Eq. (8) with respect to t and eliminate the second equation of Eq. (8):

$$\mathbf{U}' \cos(\omega t) + \mathbf{V}' \sin(\omega t) = 0. \quad (9)$$

Differentiate the second equation of Eq. (8) with respect to t and take it into Eq. (7):

$$(\omega \mathbf{M}\mathbf{V}' - \omega^2 \mathbf{M}\mathbf{U} + \omega \mathbf{C}\mathbf{V}) \cos(\omega t) - (\omega \mathbf{M}\mathbf{U}' + \omega^2 \mathbf{M}\mathbf{V} + \omega \mathbf{C}\mathbf{U}) \sin(\omega t) = \mathbf{F}. \quad (10)$$

Based on Eqs. (9) and (10), it can be inferred as:

$$\begin{cases} \omega \mathbf{M}\mathbf{U}' = -\mathbf{F} \sin(\omega t) + (-\omega^2 \mathbf{M}\mathbf{U} + \omega \mathbf{C}\mathbf{V}) \cos(\omega t) \sin(\omega t) - (\omega^2 \mathbf{M}\mathbf{V} + \omega \mathbf{C}\mathbf{U}) \sin^2(\omega t), \\ \omega \mathbf{M}\mathbf{V}' = \mathbf{F} \cos(\omega t) + (\omega^2 \mathbf{M}\mathbf{V} + \omega \mathbf{C}\mathbf{U}) \sin(\omega t) \cos(\omega t) + (\omega^2 \mathbf{M}\mathbf{U} - \omega \mathbf{C}\mathbf{V}) \cos^2(\omega t). \end{cases} \quad (11)$$

Note that, \mathbf{U} and \mathbf{V} are the slowly changing function of time. The right side of Eq. (11) can be approximately represented with the average value of (ωt) in a period. Provided that \mathbf{U} and \mathbf{V} remain unchangeable in a period, average equation can be obtained as:

$$\begin{cases} \omega \mathbf{M}\mathbf{U}' = \int_0^{2\pi} \left[-\mathbf{F} \sin(\omega t) + (-\omega^2 \mathbf{M}\mathbf{U} + \omega \mathbf{C}\mathbf{V}) \cos(\omega t) \sin(\omega t) - (\omega^2 \mathbf{M}\mathbf{V} + \omega \mathbf{C}\mathbf{U}) \sin^2(\omega t) \right] / 2\pi d\omega t, \\ \omega \mathbf{M}\mathbf{V}' = \int_0^{2\pi} \left[\mathbf{F} \cos(\omega t) + (\omega^2 \mathbf{M}\mathbf{V} + \omega \mathbf{C}\mathbf{U}) \sin(\omega t) \cos(\omega t) + (\omega^2 \mathbf{M}\mathbf{U} - \omega \mathbf{C}\mathbf{V}) \cos^2(\omega t) \right] / 2\pi d\omega t. \end{cases} \quad (12)$$

According to the orthogonality of trigonometric function, Eq. (12) can be simplified as:

$$\begin{cases} \omega \mathbf{M}\mathbf{U}' = -\frac{1}{2}(\omega^2 \mathbf{M}\mathbf{V} + \omega \mathbf{C}\mathbf{U}) - \frac{1}{2} \begin{pmatrix} Q_1 \\ Q_2 \end{pmatrix}, \\ \omega \mathbf{M}\mathbf{V}' = \frac{1}{2}(\omega^2 \mathbf{M}\mathbf{U} - \omega \mathbf{C}\mathbf{V}) + \frac{1}{2} \begin{pmatrix} Q_3 + f \\ Q_4 + f \end{pmatrix}, \end{cases} \quad (13)$$

$$\begin{aligned} Q_1 &= \left(1 + 3H^2 + \frac{3U_1^2}{4}\right)V_1 + \frac{3V_1^3}{4}, & Q_2 &= \left(-k_1 - 3k_2h_2^2 - \frac{3k_2U_2^2}{4}\right)V_2 - \frac{3V_2^3}{4}, \\ Q_3 &= \left(1 + 3H^2 + \frac{3V_1^2}{4}\right)U_1 + \frac{3U_1^3}{4}, & Q_4 &= \left(-k_1 - 3k_2h_2^2 - \frac{3k_2V_2^2}{4}\right)U_2 - \frac{3U_2^3}{4}. \end{aligned} \quad (14)$$

Then, the response of system is the solution of following equation:

$$\begin{cases} -\omega^2 V_1 + \omega^2 V_2 - \xi_1 \omega U_1 + Q_1 = 0, \\ -\omega^2 V_1 + (1 + w)\omega^2 V_2 + \xi_2 \omega U_2 + Q_2 = 0, \\ -\omega^2 U_1 + \omega^2 U_2 + \xi_1 \omega V_1 + Q_3 + f = 0, \\ -\omega^2 U_1 + (1 + w)\omega^2 U_2 + \xi_2 \omega V_2 + Q_4 + f = 0. \end{cases} \quad (15)$$

The force transmitted to the base includes elastic restoring force and damping force of the lower vibration isolator, which can be expressed as:

$$f_a = \xi_2 \dot{x}_2 + (k_1 + 3k_2h_2^2)x_2 - 3k_2h_2x_2^2 + k_2x_2^3. \quad (16)$$

Substitute Eq. (6) in Eq. (14) and neglect higher harmonics:

$$f_a = f_b \cos(\omega t) + f_c \sin(\omega t) + f_d, \quad (17)$$

where:

$$\begin{aligned} f_b &= (k_1 + 3k_2h_2^2)U_2 + \omega \xi_2 V_2 + \frac{3k_2U_2(U_2^2 + V_2^2)}{4}, \\ f_c &= (k_1 + 3k_2h_2^2)V_2 - \omega \xi_2 U_2 + \frac{3k_2V_2(V_2^3 + U_2^2)}{4}, & f_d &= -\frac{3k_2h_2(U_2^2 + V_2^2)}{2}. \end{aligned}$$

Force transmissibility of two-stage QZS system is calculated as:

$$T_f = \left(\sqrt{f_b^2 + f_c^2} + |f_d| \right) / f. \quad (18)$$

As a comparison, the two oblique springs are removed.

By successively applying the same procedure as for the model of equivalent two-stage linear system, the response of the system is the solution of Eq. (15), where

$$Q_1 = 3V_1/4, \quad Q_2 = k_1V_2, \quad Q_3 = U_1, \quad Q_4 = k_1U_2. \quad (19)$$

Similarly, the force transmissibility of two-stage linear system can be founded:

$$T_{fl} = \sqrt{(4w^2\xi_2^2V_2^2 + k_1^2)(U_2^2 + V_2^2)} / f. \quad (20)$$

4. Influence of the system parameters on force transmissibility

As discussed above, force transmissibility is closely related to f , w and ξ . Next, the effects of

different system parameters on force transmissibility of two-stage QZS system are investigated by controlling variables method. The Force transmissibility of two-stage linear system at the same condition is also plotted together to compare the isolation performance of two systems. All the force transmissibility results are plotted in dB, i.e. as $20\log_{10}T$.

Fig. 3 shows the force transmissibility of system under different f . Two-stage linear system are not influenced by the f . However, for the two-stage QZS system, two order resonant frequencies and transmissibility peaks increase with the rising of f . The less the f is, the better the isolation performance of two-stage QZS system compared with the two-stage linear system. Fig. 4 shows the force transmissibility of system under different w , for the two-stage linear system, two order resonant frequencies increase with the rising of w , the transmissibility of the first order resonant frequency decreases, while the transmissibility of the second order resonant frequency increases. For the two-stage QZS system, the first order resonant frequency transmissibility decreases with the rising of w , while the second resonant order frequency and transmissibility increases with the rising of w . Which indicates that the isolation bandwidth broadens and isolation performance near the second order resonant frequency weakens with the rising of w .

The force transmissibility curves of system under different ξ are illustrated in Fig. 5 and Fig. 6 respectively. Two order resonant frequencies of two-stage linear system are not influenced by the ξ_1 , while the peaks of corresponding force transmissibility reduce. Force transmissibility of system under different ξ_1 is shown in Fig. 5. The resonance branch of the two-stage QZS system shortens and the peaks of corresponding force transmissibility reduce with the rising of ξ_1 . But the isolation performance will decrease when the excessive damping ratio. Force transmissibility of system under different ξ_1 is shown in Fig. 6. Two order resonant frequencies increase with the rising of ξ_2 . The transmissibility of the first order resonant frequency decreases, while the transmissibility of the second order resonant frequency increases, which indicates that the isolation performance near the second order resonant frequency weakens with the rising of ξ_2 .

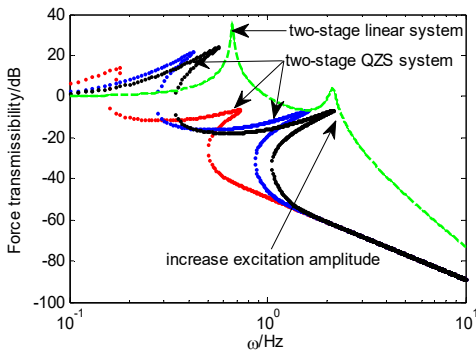


Fig. 3. Force transmissibility under different f

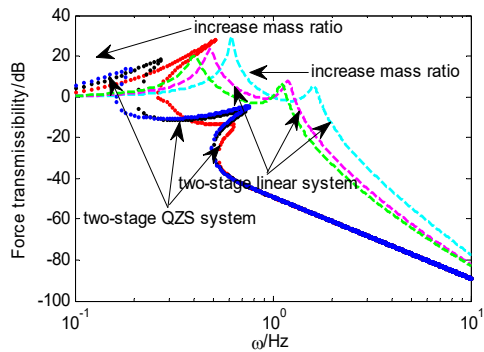


Fig. 4. Force transmissibility under different w

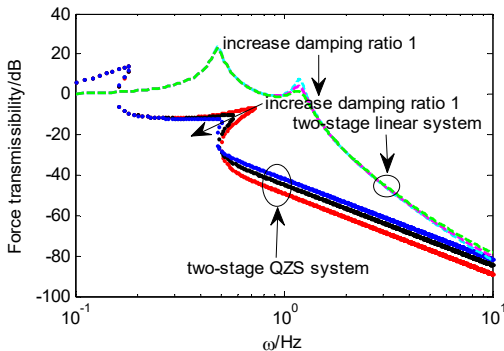


Fig. 5. Force transmissibility under different ξ_1

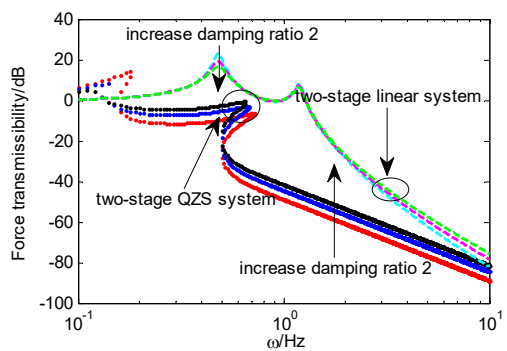


Fig. 6. Force transmissibility under different ξ_2

5. Conclusions

In this study, a novel two-stage quasi-zero stiffness vibration isolator was presented. The conclusions were summarized as follows:

1) The dynamic models of two-stage QZS and linear vibration isolation system were established. The force transmissibility under harmonic force excitation was derived by using the averaging method.

2) Decrease the excitation amplitude and increase the mass ratio as well as damping ratio properly, which can broaden the isolation bandwidth, increase the vibration decay ratio and enhance the isolation performance of two-stage QZS system.

3) Compared to the corresponding two-stage linear system, the two-stage QZS system not only has smaller initial isolation frequency, wider isolation band and better isolation performance, but also possesses excellent load-bearing capacity and stability.

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