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Solving Stefan problem through C-NEM and level-set approach



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Final goal of the study

Develop numerical time domain approach able to simulate thermo-mechanical phenomena in Finite Transformations:

- Cutting/blanking processes in 3D
 - Matter splitting encountered in forming processes
- Laser drilling/cutting
 - Multi-phases problem with moving interfaces across the matter
- Research tool in order to be able to test new approaches and thermomechanical models





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The approach must handle:

- Large strains
- Contact
- Interfaces and discontinuities





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- FEM: induced mesh distorsions are conducting to frequent re-meshing and fields projections
 - lacktriangleright need a very efficient mesher ightarrow lack of robustness in 3D
- Mesh Free: only the distribution and number of nodes are to be managed
 - OK but need to simply take into account boundary conditions





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Mesh Free

We have choose to use a Natural Neighbor interpolant based mesh free approach \Rightarrow nodal interpolation

Existing methods:

- ${f a}$ α -NEM 1 : no geometrical description of the boundaries **but** boundaries must be quite regular
- C-Nem^{2,3}: a geometrical model is needed for the boundaries but domain can be highly non convex



¹ E. Cueto, Int. J. Numer. Meth. in Engng, 2000

² J. Yvonnet, Int. J. Numer. Meth. in Engng, 2004
³ L. Illoul, Comp. and Struc., 2011

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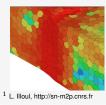
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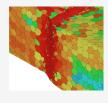
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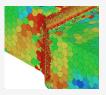
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Example of blanking process: (C-Nem simulation¹)











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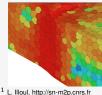
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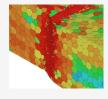
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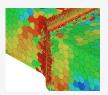
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Example of blanking process: (C-Nem simulation¹)







L. IIIoui, nttp://sn-m2p.cnrs.r

Interfaces modeling

- Full geometrical model: in 3D the shape evolution of the interface need a complex (and robust) surface mesher
 - Discontinuities: direct with duplication of the variables on nodes belonging to the interface²
- Level-set: easy but the description is linked to the nodes distribution
 - Discontinuities: X-FEM framework → X-NEM³





² J. Yvonnet, Int. J. Therm., 2005

 $^{^{3}\,}$ N. Sukumar, U.S. National Congress on Comp. Mech., 2001

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Present a numerical method to solve problems involving discontinuities on moving internal boundaries with:

- a C-Nem approach for the interpolation (based on the natural neighbours interpolation)
- a level-set technique to represent the interface
- a local enrichement through the partition of unity concept

First results in 2D for the Stefan problem are presented





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Few words on the C-Nem

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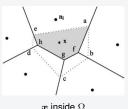
$$oldsymbol{u}^h(oldsymbol{x}) = \sum_{i=1}^n N_i(oldsymbol{x}) oldsymbol{u}_i, ~~ orall oldsymbol{x} \in \Omega$$

where $N_i({m x})$ are Natural Neighbour (NN) shape functions: one shape function per node i

Natural Neighbour shape function

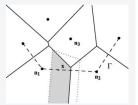
Based on :

- Voronoï diagram ⇔ Delaunay tessellation
- Systematic geometric constructions (for a given set of nodes)



Sibson shape function

$$N_i(m{x}) = rac{ ext{Area(afghe)}}{ ext{Area(abcde)}}$$



ide Ω $oldsymbol{x}$ on the boundary of Ω





Few words on the C-Nem

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Voronoï diagram with NN

Constrained Voronoï diagram with NN The constrained Voronov diagram (Delaunay tessellation) use a visibility criterion.

The Delaunay tessellation is **constrained** to respect the tessellation of $\partial\Omega$



NN supports



Constrained NN supports



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Voronoï diagram with NN

Constrained Voronoï diagram with NN The constrained Voronov diagram (Delaunay tessellation) use a visibility criterion.

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NN supports



Constrained NN supports

C-Nem: Constrained Natural Element Method

C-NEM use the constrained NN shape functions





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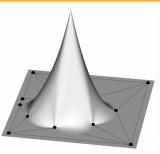
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Properties of (constrained) NN interpolant

- Delta Kronecker:
 - $N_i(\boldsymbol{x}_j) = \delta_{ij}$
- Positivity: $0 \le N_i(\boldsymbol{x}) \le 0$
- Partition of unity:
- $\sum_{i=1}^{n} N_i(\boldsymbol{x}) = 1$
- Local coordinate property:
 - $\boldsymbol{x} = \sum_{i=1}^{n} N_i(\boldsymbol{x}) \boldsymbol{x}_i$
 - \Rightarrow exact interpolation of linear fields
 - \Rightarrow reproduction of large solid motions

A Sibson shape function



Continuity

Natural neighbor shape functions are C^{∞} at any point except :

- \blacksquare at the nodes: C^0
- lacksquare on the boundary of the Delaunay circles (spheres in 3D): C^1





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Coupling C-Nem with a level-set approach

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Main aspects

We propose to use a X-FEM like strategy by enriching the C-Nem approximation space through the **partition of unity technique**.

As for the X-FEM, the **location of the discontinuity interface** is defined by a **level-set function**. This latter being defined by the nodal values of the level-set function with the **C-Nem approximation**.

We need to define:

- the adequate enrichment function (depending on the discontinuity), based on the level-set (distance) function
- the selection of the nodes subjected to enrichment (near the interface)
- the quadrature rules for the weak forms





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Local enrichment of the Constrained NN interpolant

$$T^{h}(\mathbf{x},t) = \sum_{i} N_{i}(\mathbf{x}) a_{i}(t) + \sum_{j \in I(t)} \underbrace{N_{j}(\mathbf{x}) \psi(\mathbf{x},t)}_{M_{j}(\mathbf{x},t)} b_{j}(t)$$

- $\psi(\mathbf{x},t)$ is the enrichment function depending on the interface position
- \blacksquare I(t) is the set of the nodes subjected to enrichment
- $lackbox{N}_i(oldsymbol{x})$ are the Constrained NN shape function verifying the partition of unity. If the geometry of the domain do not evolve, these shape functions do not depend on time.

Selection of the nodes subjected to enrichment

In order to define the set I(t) we use the constrained Delaunay tessellation.



Constrain Delaunay Tesselation (black)
+ Discontinuity interface (green)



Enriched nodes selection (in red)

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Enrichment function ψ

Here we have chosen a enrichment function, proposed by Moes et al¹, in order handle weak discontinuity (gradient discontinuity):

$$\psi(\mathbf{x},t) = \sum_{j \in I(t)} N_j(\mathbf{x}) \left| \Phi(\mathbf{x}_j,t) \right| - \left| \sum_j N_j(\mathbf{x}) \Phi(\mathbf{x}_j,t) \right|$$

where $\Phi(\mathbf{x}_j, t)$ is the level-set (distance) function

 $^{1}\,$ N. Moes, Comput. Methods Appl. Mech. Engrg, 2003

Representation of the level-set and enrichment functions



 $\Phi(\mathbf{x})$: Level-set contours



 $\Phi(\mathbf{x})$: Level-set contours in 3D



 $\psi(\mathbf{x})$: Schematic representation



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Quadrature

- For the integration of the weak forms we use the constrained Delaunay tessellation
- As for the X-Fem, the triangles (tetrahedrons in 3D) intersecting the interface, are re-meshed in order to be compatible with the interface and to improve the quadrature.

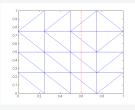




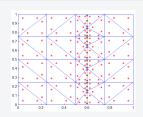


- ◆ Intersecting points
- Surface (volume in 3D) quadrature points
- Line (surface in 3D) quadrature points

Example of quadrature points distribution



Initial Delaunay mesh



Refined Delaunay mesh - Quadrature points



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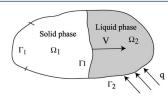
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Strong form

Heat equation :

$$\begin{split} \rho & \frac{\partial}{\partial t} \left(c_1 T \right) = \nabla \cdot \left(k_1 \nabla T \right) & \text{ in } \Omega_1(t); \quad \rho \frac{\partial}{\partial t} \left(c_2 T \right) = \nabla \cdot \left(k_2 \nabla T \right) & \text{ in } \Omega_2(t) \\ c_i, k_i : \text{ heat capacities, thermal conductivities} & \rho = \rho_1 = \rho_2 : \text{ density} \end{split}$$

Initial and boundary conditions :

$$\begin{cases} T(\boldsymbol{x},t = 0) = T_0 & \forall \boldsymbol{x} \in \Omega \\ T(\boldsymbol{x},t) = \overline{T}(\boldsymbol{x},t) & \forall \boldsymbol{x} \in \Gamma_1, \forall i \in [0,t_{\max}] \\ -k_i \nabla T(\boldsymbol{x},t) \cdot \boldsymbol{n}_{12} = \overline{q}(\boldsymbol{x},t) & \forall \boldsymbol{x} \in \Gamma_1, \forall i \in [0,t_{\max}] \end{cases}$$

Interface velocity: depends on L the latent heat of fusion

$$\boldsymbol{V}(\boldsymbol{x} \in \Gamma_I(t)) = \tfrac{[q]}{L} \boldsymbol{n}_{12}(\boldsymbol{x}) \quad \text{where } [q] = (k_1 \nabla T \big|_{\Gamma_{12}^-(t)} - k_2 \nabla T \big|_{\Gamma_{12}^+(t)}) \cdot \boldsymbol{n}_{12}$$

Constraint prescribed on the interface $\Gamma_I(t)$: $T(x,t) = T_m \quad \forall x \in \Gamma_I(t); \quad T_m$: melting temperature



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Find $T \in H^1(\Omega)$ with $T = \overline{T}$ on Γ_1 such that

$$\begin{split} \int_{\Omega} \rho c \frac{\partial T}{\partial t} \delta T \, d\Omega + \int_{\Omega} k \, \nabla T \cdot \nabla \delta T \, d\Omega &= \int_{\Gamma_I} \alpha \left(T - T_m \right) \delta T d\Gamma + \int_{\Gamma_I} \left[\!\!\left[\mathbf{q} \cdot \mathbf{n}_{12} \right]\!\!\right] \delta T d\Gamma \\ &\qquad \qquad \text{(Simplify form : } \overline{q}(t) = 0\text{)} \end{split}$$

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Time discretization using implicit scheme 1

The implicit backward Euler integration scheme between t^{n-1} and t^n gives:

$$\int_{\Omega} \rho c \frac{T^{n} - T^{n-1}}{dt} \delta T^{n} d\Omega + \int_{\Omega} k \nabla T^{n} \cdot \nabla \delta T^{n} d\Omega = \int_{\Gamma_{I}} \alpha \left(T^{n} - T_{m} \right) \delta T^{n} d\Gamma + \int_{\Gamma_{I}} \left[\left[\mathbf{q}^{n} \cdot \mathbf{n}_{12}^{n} \right] \right] \delta T^{n} d\Gamma$$



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Time discretization using implicit scheme 1

$$\int_{\Omega} \rho c \frac{T^{n} - T^{n-1}}{dt} \delta T^{n} d\Omega + \int_{\Omega} k \nabla T^{n} \cdot \nabla \delta T^{n} d\Omega = \int_{\Gamma_{I}} \alpha \left(T^{n} - T_{m} \right) \delta T^{n} d\Gamma + \int_{\Gamma_{I}} \left[\left[\mathbf{q}^{n} \cdot \mathbf{n}_{12}^{n} \right] \right] \delta T^{n} d\Gamma$$

Time discretization using implicit scheme 2

$$\int_{\Omega} \rho c T^{n} \delta T^{n} d\Omega + dt \int_{\Omega} k \nabla T^{n} \cdot \nabla \delta T^{n} d\Omega = \int_{\Omega} \rho c T^{n-1} \delta T^{n} d\Omega$$
$$+ dt \int_{\Gamma_{I}} \alpha \left(T^{n} - T_{m} \right) \delta T^{n} d\Gamma + dt \int_{\Gamma_{I}} \left(k_{1} - k_{2} \right) \left(\nabla T^{n} \cdot \mathbf{n}_{12}^{n} \right) \delta T^{n} d\Gamma$$





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Time discretization using implicit scheme 2

$$\begin{split} \int_{\Omega} \rho c T^n \delta T^n \ d\Omega + dt \ \int_{\Omega} k \, \nabla T^n \cdot \nabla \delta T^n \ d\Omega &= \int_{\Omega} \rho c T^{n-1} \delta T^n \ d\Omega \\ + dt \ \int_{\Gamma_I} \alpha \left(T^n - T_m \right) \delta T^n d\Gamma + dt \ \int_{\Gamma_I} \left(k_1 - k_2 \right) \left(\nabla T^n \cdot \mathbf{n}_{12}^n \right) \delta T^n d\Gamma \end{split}$$

Matrix Form

$$(\mathbf{C} + dt\mathbf{K})\mathbf{T}^n = \mathbf{F}$$

with

$$\mathbf{C} = \int \rho c \mathbf{N}^{nT} \mathbf{N}^{n} d\Omega$$

$$\mathbf{K} = \int_{\Omega} k \, \mathbf{B}^{nT} \mathbf{B}^{n} \, d\Omega - \int_{\Gamma} \alpha \mathbf{N}^{nT} \mathbf{N}^{n} d\Gamma + \int_{\Gamma} \left(k_{2} - k_{1} \right) \mathbf{N}^{nT} \left(\mathbf{B}^{n} \cdot \mathbf{n}_{12}^{n} \right) d\Gamma$$

$$\mathbf{F} = \int_{\Omega} \rho c \mathbf{N}^{nT} \left(\mathbf{N}^{n-1} \mathbf{T}^{n-1} \right) d\Omega + dt \int_{\Gamma_{\mathbf{I}}} (\alpha T_m) \mathbf{N}^{nT} d\Gamma$$



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Interface convection¹

Velocity extension

$$sign(\Phi)\nabla F \cdot \nabla \Phi = 0$$

with $F = \mathbf{V} \cdot \mathbf{n}_{12}$ on Γ_I

Level-set updating

$$\frac{\partial \Phi}{\partial t} + \mathbf{V} \cdot \nabla \Phi = 0$$

¹ J. Chessa, Int. J. Numer. Meth. Engng, 2002

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Pseudo-code

Let \mathbf{T}^{n-1} and Φ^{n-1} be known.

- Compute the velocity of the interface \mathbf{V}^{n-1} on Γ_I $\mathbf{V}^{n-1} = \frac{[\mathbf{q}]}{2} \mathbf{n}_{i,2}^{n-1}$
- $\begin{tabular}{l} \blacksquare & \text{Extend this velocity to the whole domain } \Omega & \text{solving} \\ & \text{sign}(\Phi)\nabla F \cdot \nabla \Phi = 0 & \text{with } F = \mathbf{V}^{n-1} \cdot \mathbf{n}_{12} & \text{on } \Gamma_I \\ \end{tabular}$
- Determine Φ^n by updating the level-set function through $\frac{\partial \Phi}{\partial t} + F |\nabla \Phi| = 0$
- Localize integration points by dividing the elements cut by Γ_I into sub-elements matching Γ_I using Φ^n only
- Build matrices C & K and vector F
- Compute \mathbf{T}^n by solving the heat equation $(\mathbf{C} + dt\mathbf{K}) \mathbf{T}^n = \mathbf{F}$





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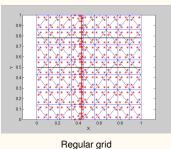
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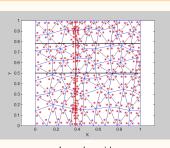
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Motion on the interface across the mesh







Irregular grid



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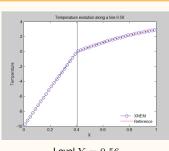
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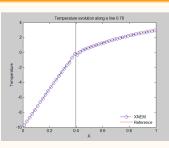
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Motion of the interface at two Y levels







Level Y = 0.78



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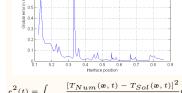
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Global and local errors on temperature





$$e^{2}(t) = \int_{\Omega(t)} \frac{\left[T_{Num}(\mathbf{x}, t) - T_{Sol}(\mathbf{x}, t)\right]^{2}}{\left[T_{Sol}(\mathbf{x}, t)\right]^{2}} dS$$



$$e(t) = \sup_{\mathbf{x} \in \Omega} \frac{|T_{Num}(\mathbf{x}, t) - T_{Sol}(\mathbf{x}, t)|}{|T_{right} - T_{left}|}$$



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Weak form

- First results are encouraging
- Partition of unity technique seems to work as well with the C-Nem than with the FEM
- It is a first approach in 2D, investigation must be done on more complex geometries and in 3D
- Main errors are observed in the "enriched zones" where partition of unity not exactly respected are observed (error $\approx 10^{-2}$)

Work still in progress ...









Thanks for your attention Any questions?

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