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Frazho, A. E.; Ran, A. C.M.

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# A note on inner-outer factorization of wide matrix-valued functions 

A.E. Frazho and A.C.M. Ran

Dedicated to our friend and mentor Rien Kaashoek on the occasion of his eightieth birthday, with gratitude for inspiring and motivating us to work on many interesting problems.


#### Abstract

In this paper we expand some of the results of [8, 9, 10]. In fact, using the techniques of $[8,9,10]$, we provide formulas for the full rank inner-outer factorization of a wide matrix-valued rational function $G$ with $H^{\infty}$ entries, that is, functions $G$ with more columns than rows. State space formulas are derived for the inner and outer factor of $G$.


Mathematics Subject Classification (2010). Primary 47B35, 47A68; Secondary 30J99 .
Keywords. Inner-outer factorization, matrix-valued function, Toeplitz operators, state space representation.

## 1. Introduction

In this note, $\mathcal{E}, \mathcal{U}$ and $\mathcal{Y}$ are finite-dimensional complex vector spaces and $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. We will present a method to compute the inner-outer factorization for certain matrix-valued rational functions $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$, defined on the closure of the unit disc. Computing inner-outer factorizations for the case when $\operatorname{dim} \mathcal{U} \leq \operatorname{dim} \mathcal{Y}$ is well developed and presented in $[4,5,13]$ and elsewhere.

Recall that a function $G_{i}$ is inner if $G_{i}$ is a function in $H^{\infty}(\mathcal{E}, \mathcal{Y})$ and $G_{i}\left(e^{i \omega}\right)$ is almost everywhere an isometry. (In particular, $\operatorname{dim} \mathcal{E} \leq \operatorname{dim} \mathcal{Y}$.) Equivalently (see, e.g., [5,13]), $G_{i}$ in $H^{\infty}(\mathcal{E}, \mathcal{Y})$ is an inner function if and only if the Toeplitz operator $T_{G_{i}}$ mapping $\ell_{+}^{2}(\mathcal{E})$ into $\ell_{+}^{2}(\mathcal{Y})$ is an isometry. A function $G_{o}$ is outer if $G_{o}$ is a function in $H^{\infty}(\mathcal{U}, \mathcal{E})$ and the range of the Toeplitz operator $T_{G_{o}}$ is dense in $\ell_{+}^{2}(\mathcal{E})$.

Let $G$ be a function in $H^{\infty}(\mathcal{U}, \mathcal{Y})$. Then $G$ admits a unique inner-outer factorization of the form $G(\lambda)=G_{i}(\lambda) G_{o}(\lambda)$ where $G_{i}$ is an inner function in $H^{\infty}(\mathcal{E}, \mathcal{Y})$ and $G_{o}$ is an outer function in $H^{\infty}(\mathcal{U}, \mathcal{E})$ for some intermediate space $\mathcal{E}$. Because $G_{i}\left(e^{i \omega}\right)$ is almost everywhere an isometry, $\operatorname{dim} \mathcal{E} \leq \operatorname{dim} \mathcal{Y}$.

Since $G_{o}$ is outer, $G_{o}\left(e^{i \omega}\right)$ is almost everywhere onto $\mathcal{E}$, and thus, $\operatorname{dim} \mathcal{E} \leq$ $\operatorname{dim} \mathcal{U}$. By unique we mean that if $G(\lambda)=F_{i}(\lambda) F_{o}(\lambda)$ is another inner-outer factorization of $G$ where $F_{i}$ is an inner function in $H^{\infty}(\mathcal{L}, \mathcal{Y})$ and $F_{o}$ is an outer function in $H^{\infty}(\mathcal{U}, \mathcal{L})$, then there exists an constant unitary operator $\Omega$ mapping $\mathcal{E}$ onto $\mathcal{L}$ such that $G_{i}=F_{i} \Omega$ and $\Omega G_{o}=F_{o}$; see $[1,5,6,13,14,15]$ for further details.

Throughout we assume that $\mathcal{U}, \mathcal{E}$ and $\mathcal{Y}$ are all finite dimensional. We say that $G_{i}$ in $H^{\infty}(\mathcal{E}, \mathcal{Y})$ is a square inner function if $G_{i}$ is an inner function and $\mathcal{E}$ and $\mathcal{Y}$ have the same dimension, that is, $G_{i}\left(e^{i \omega}\right)$ is almost everywhere a unitary operator, or equivalently, $G_{i}$ is a two-sided inner function. So if $G_{i} G_{o}$ is an inner-outer factorization of $G$ where $G_{i}$ is square, then without loss of generality we can assume that $\mathcal{E}=\mathcal{Y}$.

We say that the inner-outer factorization $G=G_{i} G_{o}$ is full rank if $G_{i}$ is a square inner function in $H^{\infty}(\mathcal{Y}, \mathcal{Y})$ and the range of $T_{G_{o}}$ equals $\ell_{+}^{2}(\mathcal{Y})$. An inner-outer factorization $G=G_{i} G_{o}$ is full rank if and only if $G_{i}$ is a square inner function and the range of $T_{G}$ is closed. If $G$ is a rational function, then $G$ admits a full rank inner-outer factorization if and only if

$$
\begin{equation*}
G\left(e^{i \omega}\right) G\left(e^{i \omega}\right)^{*} \geq \epsilon I \quad(\text { for all } \omega \in[0,2 \pi] \text { and some } \epsilon>0) \tag{1.1}
\end{equation*}
$$

see Lemma 3.1 below. Finally, if $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ admits a full rank inner-outer factorization, then $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$.

Here we are interested in computing the inner-outer factorization for full rank rational functions $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$. So throughout we assume that $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. Computing inner-outer factorizations when $G$ does not admit a full rank factorization is numerically sensitive. (In this case, our algebraic Riccati equation may not have a stabilizing solution.) Moreover, if $G$ does not admit a full rank inner-outer factorization, then a small $H^{\infty}$ perturbation of $G$ does admit such a factorization. (If $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$, does not satisfy (1.1), then a "small random" rational $H^{\infty}$ perturbation of $G$ will satisfy (1.1).) First we will present necessary and sufficient conditions to determine when $G$ admits a full rank inner-outer factorization. Then we will give a state space algorithm to compute $G_{i}$ and then $G_{o}$. Finally, it is emphasized that this note is devoted to finding inner-outer factorizations for wide rational functions $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ when $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. Finding inner-outer factorizations when $\operatorname{dim} \mathcal{U} \leq \operatorname{dim} \mathcal{Y}$ is well developed and presented in $[4,5]$ and elsewhere.

## 2. Preliminaries

Let $R=\sum_{-\infty}^{\infty} e^{i \omega n} R_{n}$ be the Fourier series expansion for a function $R$ in $L^{\infty}(\mathcal{Y}, \mathcal{Y})$. Then $T_{R}$ is the Toeplitz operator on $\ell_{+}^{2}(\mathcal{Y})$ defined by

$$
T_{R}=\left[\begin{array}{cccc}
R_{0} & R_{-1} & R_{-2} & \cdots  \tag{2.2}\\
R_{1} & R_{0} & R_{-1} & \cdots \\
R_{2} & R_{1} & R_{0} & \cdots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right] \text { on } \ell_{+}^{2}(\mathcal{Y})
$$

The function $R$ is called the symbol for $T_{R}$. Recall that the Toeplitz operator $T_{R}$ is strictly positive if and only if there exists an $\epsilon>0$ such that $R\left(e^{i \omega}\right) \geq \epsilon I$ almost everywhere. The Toeplitz operator $T_{G}$ with symbol $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$, is given by

$$
T_{G}=\left[\begin{array}{cccc}
G_{0} & 0 & 0 & \cdots  \tag{2.3}\\
G_{1} & G_{0} & 0 & \cdots \\
G_{2} & G_{1} & G_{0} & \cdots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right]: \ell_{+}^{2}(\mathcal{U}) \rightarrow \ell_{+}^{2}(\mathcal{Y})
$$

where $G(\lambda)=\sum_{0}^{\infty} \lambda^{n} G_{n}$ is the Taylor series expansion for $G$ about the origin. Moreover, if $G$ is in $H^{\infty}(\mathcal{U}, \mathcal{Y})$, then the Hankel operator $H_{G}$ mapping $\ell_{+}^{2}(\mathcal{U})$ into $\ell_{+}^{2}(\mathcal{Y})$ is defined by

$$
H_{G}=\left[\begin{array}{cccc}
G_{1} & G_{2} & G_{3} & \cdots  \tag{2.4}\\
G_{2} & G_{3} & G_{4} & \cdots \\
G_{3} & G_{4} & G_{5} & \cdots \\
\vdots & \vdots & \vdots & \vdots
\end{array}\right]: \ell_{+}^{2}(\mathcal{U}) \rightarrow \ell_{+}^{2}(\mathcal{Y})
$$

Finally, for $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ it is well know and easy to verify that

$$
\begin{equation*}
T_{G} T_{G}^{*}=T_{G G^{*}}-H_{G} H_{G}^{*} . \tag{2.5}
\end{equation*}
$$

## 3. Inner-outer factorization

First a characterization of the existence of a full rank inner-outer factorization is presented.

Lemma 3.1. Let $G$ be a rational function in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ where $\mathcal{U}$ and $\mathcal{Y}$ are finite-dimensional spaces satisfying $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. Then $G$ admits a full rank inner-outer factorization if and only if

$$
\begin{equation*}
G\left(e^{i \omega}\right) G\left(e^{i \omega}\right)^{*} \geq \epsilon I \quad(\text { for all } \omega \in[0,2 \pi] \text { and some } \epsilon>0) \tag{3.6}
\end{equation*}
$$

or equivalently, the Toeplitz operator $T_{G G^{*}}$ is strictly positive.
Proof. Let $G=G_{i} G_{o}$ be the inner-outer factorization for $G$ where $G_{i}$ is an inner function in $H^{\infty}(\mathcal{E}, \mathcal{Y})$ and $G_{o}$ is an outer function in $H^{\infty}(\mathcal{U}, \mathcal{E})$. Clearly,

$$
\begin{equation*}
G\left(e^{i \omega}\right) G\left(e^{i \omega}\right)^{*}=G_{i}\left(e^{i \omega}\right) G_{o}\left(e^{i \omega}\right) G_{o}\left(e^{i \omega}\right)^{*} G_{i}\left(e^{i \omega}\right)^{*} \tag{3.7}
\end{equation*}
$$

Because $G_{i}$ is an inner function, $G\left(e^{i \omega}\right) G\left(e^{i \omega}\right)^{*}$ and $G_{o}\left(e^{i \omega}\right) G_{o}\left(e^{i \omega}\right)^{*}$ have the same nonzero spectrum and rank almost everywhere. The range of $T_{G_{o}}$ equals $\ell_{+}^{2}(\mathcal{E})$ if and only if the operator $T_{G_{o}} T_{G_{o}}^{*}$ is strictly positive. If $T_{G_{o}} T_{G_{o}}^{*}$ is strictly positive, then $T_{G_{o} G_{o}^{*}}=T_{G_{o}} T_{G_{o}}^{*}+H_{G_{o}} H_{G_{o}}^{*}$ implies that $T_{G_{o} G_{o}^{*}}$ is also strictly positive. So if the range of $T_{G_{o}}$ equals $\ell_{+}^{2}(\mathcal{E})$, then $G_{o}\left(e^{i \omega}\right) G_{o}\left(e^{i \omega}\right)^{*} \geq$ $\epsilon I_{\mathcal{E}}$ for some $\epsilon>0$.

In addition, if $G=G_{i} G_{o}$ is a full rank inner-outer factorization, then $G_{i}\left(e^{i \omega}\right)$ is a unitary operator. In this case, equation (3.7) shows that (3.6) holds.

On the other hand, assume that (3.6) holds, or equivalently, the Toeplitz operator $T_{G G^{*}}$ is strictly positive. Because $G$ is rational, the range of $H_{G}$ is finite dimensional. Using $T_{G} T_{G}^{*}=T_{G G^{*}}-H_{G} H_{G}^{*}$, we see that $T_{G} T_{G}^{*}$ equals a strictly positive operator $T_{G G^{*}}$ minus a finite rank positive operator $H_{G} H_{G}^{*}$. Clearly, $T_{G G^{*}}$ is a Fredholm operator with index zero. Since $T_{G} T_{G}^{*}$ is a finite rank perturbation of $T_{G G^{*}}$, it follows that $T_{G} T_{G}^{*}$ is also a Fredholm operator with index zero. In particular, the range of $T_{G}$ is closed. Hence the range of $T_{G_{o}}$ is also closed. Because $G\left(e^{i \omega}\right) G\left(e^{i \omega}\right)^{*}$ and $G_{o}\left(e^{i \omega}\right) G_{o}\left(e^{i \omega}\right)^{*}$ have the same rank and $\operatorname{dim} \mathcal{E} \leq \operatorname{dim} \mathcal{Y}$, equation (3.7) with (3.6) shows that $\mathcal{E}$ and $\mathcal{Y}$ are of the same dimension. In particular, $G_{i}$ is a square inner function. Therefore the inner-outer factorization $G=G_{i} G_{o}$ is of full rank.

Next, we recall some results on the inner-outer factorization in terms of a stable finite-dimensional realization for a rational function $G$. To this end, let $\{A$ on $\mathcal{X}, B, C, D\}$ be a stable realization for $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$, that is,

$$
\begin{equation*}
G(\lambda)=D+\lambda C(I-\lambda A)^{-1} B \tag{3.8}
\end{equation*}
$$

Here $A$ is a stable operator on a finite-dimensional space $\mathcal{X}$ and $B$ maps $\mathcal{U}$ into $\mathcal{X}$ while $C$ maps $\mathcal{X}$ into $\mathcal{Y}$ and $D$ maps $\mathcal{U}$ into $\mathcal{Y}$. By stable we mean that all the eigenvalues for $A$ are inside the open unit disc. Note that $\{A, B, C, D\}$ is a realization for $G$ if and only if

$$
\begin{equation*}
G_{0}=D \quad \text { and } \quad G_{n}=C A^{n-1} B \quad(\text { for } n \geq 1) \tag{3.9}
\end{equation*}
$$

where $G(\lambda)=\sum_{0}^{\infty} \lambda^{n} G_{n}$ is the Taylor series expansion for $G$. Let $W_{o}$ be the observability operator mapping $\mathcal{X}$ into $\ell_{+}^{2}(\mathcal{Y})$ and $W_{c}$ the controllability operator mapping $\ell_{+}^{2}(\mathcal{U})$ into $\mathcal{X}$ defined by

$$
\begin{align*}
W_{o} & =\left[\begin{array}{c}
C \\
C A \\
C A^{2} \\
\vdots
\end{array}\right]: \mathcal{X} \rightarrow \ell_{+}^{2}(\mathcal{Y}), \\
W_{c} & =\left[\begin{array}{lll}
B & A B & A^{2} B \cdots
\end{array}\right]: \ell_{+}^{2}(\mathcal{U}) \rightarrow \mathcal{X} . \tag{3.10}
\end{align*}
$$

Let $P=W_{c} W_{c}^{*}=\sum_{0}^{\infty} A^{n} B B^{*} A^{* n}$ be the controllability Gramian for the pair $\{A, B\}$. Then $P$ is the solution to the following Stein equation

$$
\begin{equation*}
P=A P A^{*}+B B^{*} . \tag{3.11}
\end{equation*}
$$

Using (3.9), we see that the Hankel operator $H_{G}$ is equal to

$$
\begin{equation*}
H_{G}=W_{o} W_{c} \tag{3.12}
\end{equation*}
$$

In particular, it follows that the Hankel operator $H_{G}$ admits a factorization of the form $H_{G}=W_{o} W_{c}$ where $W_{o}$ is an operator mapping $\mathcal{X}$ into $\ell_{+}^{2}(\mathcal{Y})$ and $W_{c}$ is an operator mapping $\ell_{+}^{2}(\mathcal{U})$ into $\mathcal{X}$. Using $P=W_{c} W_{c}^{*}$ with (2.5), we obtain

$$
\begin{equation*}
H_{G} H_{G}^{*}=W_{o} P W_{o}^{*} \quad \text { and } \quad T_{G} T_{G}^{*}=T_{G G^{*}}-W_{o} P W_{o}^{*} . \tag{3.13}
\end{equation*}
$$

Consider the algebraic Riccati equation

$$
\begin{align*}
Q & =A^{*} Q A+\left(C-\Gamma^{*} Q A\right)^{*}\left(R_{0}-\Gamma^{*} Q \Gamma\right)^{-1}\left(C-\Gamma^{*} Q A\right) \\
\Gamma & =B D^{*}+A P C^{*} \quad \text { and } \quad R_{0}=D D^{*}+C P C^{*} \tag{3.14}
\end{align*}
$$

We say that $Q$ is a stabilizing solution to this algebraic Riccati equation if $Q$ is positive, $R_{0}-\Gamma^{*} Q \Gamma$ is strictly positive, and the following operator $A_{o}$ on $\mathcal{X}$ is stable:

$$
\begin{equation*}
A_{o}=A-\Gamma\left(R_{0}-\Gamma^{*} Q \Gamma\right)^{-1}\left(C-\Gamma^{*} Q A\right) \tag{3.15}
\end{equation*}
$$

Moreover, if the algebraic Riccati equation (3.14) admits a stabilizing solution $Q$, then the stabilizing solution $Q$ can be computed by

$$
\begin{align*}
Q & =\lim _{n \rightarrow \infty} Q_{n}  \tag{3.16}\\
Q_{n+1} & =A^{*} Q_{n} A+\left(C-\Gamma^{*} Q_{n} A\right)^{*}\left(R_{0}-\Gamma^{*} Q_{n} \Gamma\right)^{-1}\left(C-\Gamma^{*} Q_{n} A\right)
\end{align*}
$$

subject to the initial condition $Q_{0}=0$. In particular, if the limit in (3.16) does not exist or $A_{o}$ is not stable, then the algebraic Riccati equation (3.14) does not have a stabilizing solution; see $[8,9]$ for further details.

If $\Theta$ is an inner function in $H^{\infty}(\mathcal{E}, \mathcal{Y})$, then $\mathfrak{H}(\Theta)$ is the subspace of $\ell_{+}^{2}(\mathcal{Y})$ defined by

$$
\begin{equation*}
\mathfrak{H}(\Theta)=\ell_{+}^{2}(\mathcal{Y}) \ominus T_{\Theta} \ell_{+}^{2}(\mathcal{E})=\operatorname{ker} T_{\Theta}^{*} . \tag{3.17}
\end{equation*}
$$

Because $T_{\Theta}$ is an isometry, $I-T_{\Theta} T_{\Theta}^{*}$ is the orthogonal projection onto $\mathfrak{H}(\Theta)$. It is noted that $\mathfrak{H}(\Theta)$ is an invariant subspace for the backward shift $S_{\mathcal{Y}}^{*}$ on $\ell_{+}^{2}(\mathcal{Y})$. According to the Beurling-Lax-Halmos Theorem if $\mathfrak{H}$ is any invariant subspace for the backward shift, then there exists a unique inner function $\Theta$ in $H^{\infty}(\mathcal{E}, \mathcal{Y})$ such that $\mathfrak{H}=\mathfrak{H}(\Theta)$. By unique we mean that if $\mathfrak{H}=\mathfrak{H}(\Psi)$ where $\Psi$ is an inner function in $H^{\infty}(\mathcal{L}, \mathcal{Y})$, then there exists a constant unitary operator $\Omega$ from $\mathcal{E}$ onto $\mathcal{L}$ such that $\Theta=\Psi \Omega$; see $[5,11,12,13,14,15]$ for further details. By combining Lemma 3.1 with the results in [9], we obtain the following result. (For part (v) compare also Lemma 4.1 below.)

Theorem 3.2. Let $\{A$ on $\mathcal{X}, B, C, D\}$ be a minimal realization for a rational function $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ where $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. Let $R$ be the function in $L^{\infty}(\mathcal{Y}, \mathcal{Y})$ defined by $R\left(e^{i \omega}\right)=G\left(e^{i \omega}\right) G\left(e^{i \omega}\right)^{*}$. Let $P$ the unique solution to the Stein equation $P=A P A^{*}+B B^{*}$. Then the following statements are equivalent.
(i) The function $G$ admits a full rank inner-outer factorization;
(ii) the Toeplitz operator $T_{R}$ is invertible;
(iii) there exists a stabilizing solution $Q$ to the algebraic Riccati equation (3.14).

In this case, $Q=W_{o}^{*} T_{R}^{-1} W_{o}$ and the following holds.
(iv) The eigenvalues of $Q P$ are real numbers contained in the interval $[0,1]$.
(v) If $G_{i}$ is the inner factor of $G$, then the dimension of $\mathfrak{H}\left(G_{i}\right)$ is given by

$$
\begin{equation*}
\operatorname{dim} \mathfrak{H}\left(G_{i}\right)=\operatorname{dim} \operatorname{ker} T_{G_{i}}^{*}=\operatorname{dim} \operatorname{ker} T_{G}^{*}=\operatorname{dim} \operatorname{ker}(I-Q P) . \tag{3.18}
\end{equation*}
$$

(vi) The McMillan degree of $G_{i}$ is given by

$$
\begin{equation*}
\delta\left(G_{i}\right)=\operatorname{dim} \mathfrak{H}\left(G_{i}\right)=\operatorname{dim} \operatorname{ker}(I-Q P) \tag{3.19}
\end{equation*}
$$

In particular, the McMillan degree of $G_{i}$ is less than or equal to the McMillan degree of $G$.
(vii) The operator $T_{R}^{-1} W_{o}$ is given by

$$
\begin{align*}
T_{R}^{-1} W_{o} & =\left[\begin{array}{c}
C_{o} \\
C_{o} A_{o} \\
C_{o} A_{o}^{2} \\
\vdots
\end{array}\right]: \mathcal{X} \rightarrow \ell_{+}^{2}(\mathcal{Y}), \\
C_{o} & =\left(R_{0}-\Gamma^{*} Q \Gamma\right)^{-1}\left(C-\Gamma^{*} Q A\right): \mathcal{X} \rightarrow \mathcal{Y} \tag{3.20}
\end{align*}
$$

Finally, because $\{C, A\}$ is observable, $T_{R}^{-1} W_{o}$ is one-to-one and $\left\{C_{o}, A_{o}\right\}$ is a stable observable pair.

Let us present the following classical result; see Theorem 7.1 in [7], Sections 4.2 and 4.3 in [5] and Section XXVIII. 7 in [11].

Lemma 3.3. Let $\Theta$ be an inner function in $H^{\infty}(\mathcal{Y}, \mathcal{Y})$ where $\mathcal{Y}$ is finite dimensional. Then the Hankel operator $H_{\Theta}$ is a partial isometry whose range equals $\mathfrak{H}(\Theta)$, that is,

$$
\begin{equation*}
P_{\mathfrak{H}(\Theta)}=H_{\Theta} H_{\Theta}^{*} \tag{3.21}
\end{equation*}
$$

where $P_{\mathfrak{H}(\Theta)}$ denotes the orthogonal projection onto $\mathfrak{H}(\Theta)$. Furthermore, the following holds.
(i) The subspace $\mathfrak{H}(\Theta)$ is finite dimensional if and only if $\Theta$ is rational.
(ii) The dimension of $\mathfrak{H}(\Theta)$ equals the McMillan degree of $\Theta$.

Proof. For completeness a proof is given. By replacing $G$ by $\Theta$ in (2.5), we see that

$$
T_{\Theta} T_{\Theta}^{*}=T_{\Theta \Theta^{*}}-H_{\Theta} H_{\Theta}^{*}
$$

Because $\Theta$ is a square inner function, $\Theta\left(e^{i \omega}\right) \Theta\left(e^{i \omega}\right)^{*}=I$ almost everywhere on the unit circle. Hence $T_{\Theta \Theta^{*}}=I$. This readily implies that

$$
H_{\Theta} H_{\Theta}^{*}=I-T_{\Theta} T_{\Theta}^{*}=P_{\mathfrak{H}(\Theta)}
$$

Therefore (3.21) holds and $H_{\Theta}$ is a partial isometry whose range equals $\mathfrak{H}(\Theta)$.
It is well know that the range of a Hankel operator $H_{F}$ is finite dimensional if and only if its symbol $F$ is rational. Moreover, the dimension of the range of the Hankel operator $H_{F}$ equals the McMillan degree of $F$. Therefore parts (i) and (ii) follow from the fact that $\mathfrak{H}(\Theta)=\operatorname{ran} H_{\Theta}$.

Let $\left\{A_{i}\right.$ on $\left.\mathcal{X}_{i}, B_{i}, C_{i}, D_{i}\right\}$ be a minimal state space realization for a rational function $\Theta$ in $H^{\infty}(\mathcal{Y}, \mathcal{Y})$. It is well known (see, e.g., [7], Section III.7) that $\Theta$ is a square inner function if and only if

$$
\left[\begin{array}{cc}
A_{i}^{*} & C_{i}^{*}  \tag{3.22}\\
B_{i}^{*} & D_{i}^{*}
\end{array}\right]\left[\begin{array}{cc}
Q_{i} & 0 \\
0 & I
\end{array}\right]\left[\begin{array}{cc}
A_{i} & B_{i} \\
C_{i} & D_{i}
\end{array}\right]=\left[\begin{array}{cc}
Q_{i} & 0 \\
0 & I
\end{array}\right]
$$

where $Q_{i}=A_{i}^{*} Q_{i} A_{i}+C_{i}^{*} C_{i}$. Moreover, in this case,

$$
\mathfrak{H}(\Theta)=\operatorname{ran} H_{\Theta}=\operatorname{ran} W_{i}
$$

where $W_{i}$ is the observability operator for $\left\{C_{i}, A_{i}\right\}$ defined by

$$
W_{i}=\left[\begin{array}{c}
C_{i}  \tag{3.23}\\
C_{i} A_{i} \\
C_{i} A_{i}^{2} \\
\vdots
\end{array}\right]: \mathcal{X}_{i} \rightarrow \ell_{+}^{2}(\mathcal{Y})
$$

It is noted that $S_{\mathcal{Y}}^{*} W_{i}=W_{i} A_{i}$. So the range of $W_{i}$ is a finite-dimensional invariant subspace for the backward shift $S_{\mathcal{Y}}^{*}$. Finally, $Q_{i}=W_{i}^{*} W_{i}$.

On the other hand, if $\left\{C_{i}, A_{i}\right.$ on $\left.\mathcal{X}_{i}\right\}$ is a stable observable pair where $\mathcal{X}_{i}$ is finite dimensional, then there exists operators $B_{i}$ mapping $\mathcal{Y}$ into $\mathcal{X}_{i}$ and $D_{i}$ on $\mathcal{Y}$ such that

$$
\begin{equation*}
\Theta(\lambda)=D_{i}+\lambda C_{i}\left(I-\lambda A_{i}\right)^{-1} B_{i} \tag{3.24}
\end{equation*}
$$

is a square inner function in $H^{\infty}(\mathcal{Y}, \mathcal{Y})$. Moreover, $\mathfrak{H}(\Theta)=\operatorname{ran} W_{i}$ and (3.22) holds. The Beurling-Lax-Halmos Theorem guarantees that the inner function $\Theta$ is unique up to a unitary constant on the right. The operators $B_{i}$ and $D_{i}$ are called the complementary operators for the pair $\left\{C_{i}, A_{i}\right\}$. To compute the complementary operators $B_{i}$ and $D_{i}$ explicitly, let

$$
\left[\begin{array}{l}
E_{1}  \tag{3.25}\\
E_{2}
\end{array}\right]: \mathcal{Y} \rightarrow\left[\begin{array}{l}
\mathcal{X}_{i} \\
\mathcal{Y}
\end{array}\right]
$$

be an isometry from $\mathcal{Y}$ onto the kernel of $\left[\begin{array}{ll}A_{i}^{*} Q_{i}^{\frac{1}{2}} & C_{i}^{*}\end{array}\right]$. Then set

$$
\begin{equation*}
B_{i}=Q_{i}^{-\frac{1}{2}} E_{1} \quad \text { and } \quad D_{i}=E_{2} \tag{3.26}
\end{equation*}
$$

Because the pair $\left\{C_{i}, A_{i}\right\}$ is observable, the operator $W_{i}$ defined in (3.23) is one to one, and the complementary operators $B_{i}$ and $D_{i}$ together with $A_{i}$ and $C_{i}$ form a minimal realization $\left\{A_{i}, B_{i}, C_{i}, D_{i}\right\}$ for a square inner function $\Theta$ such that $\operatorname{ran} W_{i}=\mathfrak{H}(\Theta)$. For further details see Lemma XXVIII7.7 in [11] and Sections 4.2 and 4.3 in [5].

We are now in a position to present our main result. The proof is given in Section 5.

Let $\{A, B, C, D\}$ be a minimal realization for a rational function $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ where $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. To compute a full rank inner-outer factorization $G=G_{i} G_{o}$ for $G$, let $P$ be the controllability Gramian for the pair $\{A, B\}$ (see (3.11)) and $Q$ the stabilizing solution to the algebraic Riccati equation (3.14). If this algebraic Riccati equation does not admit a stabilizing solution, then $G$ does not have a full rank inner-outer factorization.

Theorem 3.4. Let $\{A, B, C, D\}$ be a minimal realization for a rational function $G$ in $H^{\infty}(\mathcal{U}, \mathcal{Y})$ where $\operatorname{dim} \mathcal{Y} \leq \operatorname{dim} \mathcal{U}$. Assume there exists a stabilizing solution $Q$ to the algebraic Riccati equation (3.14).

Let $\mathcal{X}_{i}$ be any space isomorphic to the kernel of $I-Q P$. Let $U$ be any isometry from $\mathcal{X}_{i}$ onto the kernel of $I-Q P$. In particular, $U=Q P U$. Let $A_{i}$ on $\mathcal{X}_{i}$ and $C_{i}$ mapping $\mathcal{X}_{i}$ into $\mathcal{Y}$ be the operators computed by

$$
\begin{equation*}
A_{i}=U^{*} Q A_{o} P U \quad \text { and } \quad C_{i}=C_{o} P U \tag{3.27}
\end{equation*}
$$

Then $\left\{C_{i}, A_{i}\right\}$ is a stable observable pair. Let $B_{i}$ and $D_{i}$ be the complementary operators for $\left\{C_{i}, A_{i}\right\}$ as constructed in (3.25) and (3.26). Then the square inner factor $G_{i}$ for $G$ is given by

$$
\begin{equation*}
G_{i}(\lambda)=D_{i}+\lambda C_{i}\left(I-\lambda A_{i}\right)^{-1} B_{i} . \tag{3.28}
\end{equation*}
$$

The outer factor $G_{o}$ for $G$ is given by

$$
\begin{equation*}
G_{o}(\lambda)=D_{i}^{*} D+B_{i}^{*} U^{*} B+\lambda\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right)(I-\lambda A)^{-1} B . \tag{3.29}
\end{equation*}
$$

## 4. An auxiliary lemma

To prove that the inner-outer factorization of $G=G_{i} G_{o}$ is indeed given by (3.28) and (3.29), let us begin with the following auxiliary result.

Lemma 4.1. Let $T$ be a strictly positive operator on $\mathcal{H}$ and $P$ a strictly positive operator on $\mathcal{X}$. Let $W$ be an operator mapping $\mathcal{X}$ into $\mathcal{H}$ and set $Q=W^{*} T^{-1} W$. Then the following two assertions hold.
(i) Let $\mathfrak{X}$ and $\mathfrak{H}$ be the subspaces defined by

$$
\begin{equation*}
\mathfrak{X}=\operatorname{ker}(I-Q P) \quad \text { and } \quad \mathfrak{H}=\operatorname{ker}\left(T-W P W^{*}\right) . \tag{4.1}
\end{equation*}
$$

Then the operators

$$
\begin{equation*}
\Lambda_{1}=W^{*} \mid \mathfrak{H}: \mathfrak{H} \rightarrow \mathfrak{X} \quad \text { and } \quad \Lambda_{2}=T^{-1} W P \mid \mathfrak{X}: \mathfrak{X} \rightarrow \mathfrak{H} \tag{4.2}
\end{equation*}
$$

are both well defined and invertible. Moreover, $\Lambda_{1}^{-1}=\Lambda_{2}$.
(ii) The operator $T-W P W^{*}$ is positive if and only if $P^{-1}-Q$ is positive, or equivalently, $P^{\frac{1}{2}} Q P^{\frac{1}{2}}$ is a contraction. In this case, the spectrum of $Q P$ is contained in $[0,1]$. In particular, if $\mathcal{X}$ is finite dimensional, then the eigenvalues for $Q P$ are contained in $[0,1]$.

Proof. The proof is based on some ideas involving Schur complements; see [2] and Section 2.2 in [3]. Consider the operator matrix

$$
\begin{aligned}
M=\left[\begin{array}{cc}
T & W \\
W^{*} & P^{-1}
\end{array}\right] & =\left[\begin{array}{cc}
I & W P \\
0 & I
\end{array}\right]\left[\begin{array}{cc}
T-W P W^{*} & 0 \\
0 & P^{-1}
\end{array}\right]\left[\begin{array}{cc}
I & 0 \\
P W^{*} & I
\end{array}\right] \\
& =\left[\begin{array}{cc}
I & 0 \\
W^{*} T^{-1} & I
\end{array}\right]\left[\begin{array}{cc}
T & 0 \\
0 & P^{-1}-Q
\end{array}\right]\left[\begin{array}{cc}
I & T^{-1} W \\
0 & I
\end{array}\right] .
\end{aligned}
$$

From this we conclude several things: first, by the fact that both $T$ and $P$ are strictly positive, the congruences above imply that $T-W^{*} P W$ is positive if and only if $P^{-1}-Q$ is positive, or equivalently, $P^{\frac{1}{2}} Q P^{\frac{1}{2}}$ is a contraction. In particular, if $T-W^{*} P W$ is positive, then the spectrum of $Q P$ is contained in the interval $[0,1]$. This proves part (ii) of the lemma.

To prove part (i) observe that we can describe ker $M$ in two different ways. Based on the first factorization we have

$$
\operatorname{ker} M=\left\{\left.\left[\begin{array}{c}
h \\
-P W^{*} h
\end{array}\right] \right\rvert\, h \in \operatorname{ker}\left(T-W P W^{*}\right)=\mathfrak{H}\right\} .
$$

The second factorization yields

$$
\begin{aligned}
& \operatorname{ker} M=\left\{\left.\left[\begin{array}{c}
-T^{-1} W y \\
y
\end{array}\right] \right\rvert\, y \in \operatorname{ker}\left(P^{-1}-Q\right)=P \mathfrak{X}\right\} \\
& \operatorname{ker} M=\left\{\left.\left[\begin{array}{c}
T^{-1} W P x \\
-P x
\end{array}\right] \right\rvert\, x \in \operatorname{ker}(I-Q P)=\mathfrak{X}\right\}
\end{aligned}
$$

Together these equalities prove the first assertion in Lemma 4.1. Indeed,

$$
\Phi_{1}=\left[\begin{array}{c}
I \\
-P W^{*}
\end{array}\right]: \mathfrak{H} \rightarrow \operatorname{ker} M
$$

is a one-to-one operator from $\mathfrak{H}$ onto ker $M$. Likewise,

$$
\Phi_{2}=\left[\begin{array}{c}
T^{-1} W P \\
-P
\end{array}\right]: \mathfrak{X} \rightarrow \operatorname{ker} M
$$

is a one-to-one operator from $\mathfrak{X}$ onto ker $M$. Because the first component of $\Phi_{1}$ is the identity operator on $\mathfrak{H}$, we see that $T^{-1} W P$ maps $\mathfrak{X}$ onto $\mathfrak{H}$. Since the second component of $\Phi_{2}$ is $-P$ and $P$ is invertible, $W^{*}$ maps $\mathfrak{H}$ onto $\mathfrak{X}$. Therefore the operators $\Lambda_{1}$ and $\Lambda_{2}$ in (4.2) are well defined.

If $x$ is in $\mathfrak{X}$, then $\Phi_{2} x=\Phi_{1} h$ for some unique $h$ in $\mathfrak{H}$, that is,

$$
\left[\begin{array}{c}
\Lambda_{2} x \\
-P x
\end{array}\right]=\left[\begin{array}{c}
T^{-1} W P x \\
-P x
\end{array}\right]=\Phi_{2} x=\Phi_{1} h=\left[\begin{array}{c}
h \\
-P W^{*} h
\end{array}\right]=\left[\begin{array}{c}
\Lambda_{2} x \\
-P W^{*} \Lambda_{2} x
\end{array}\right] .
$$

The last equality follows from the fact that $h=\Lambda_{2} x$. The second component of the previous equation shows that $x=W^{*} \Lambda_{2} x$, and thus, $\Lambda_{1}=W^{*} \mid \mathfrak{H}$ is the left inverse of $\Lambda_{2}$. On the other hand, if $h$ is in $\mathfrak{H}$, then $\Phi_{1} h=\Phi_{2} x$ for some unique $x$ in $\mathfrak{X}$, that is,

$$
\left[\begin{array}{c}
h \\
-P W^{*} h
\end{array}\right]=\Phi_{1} h=\Phi_{2} x=\left[\begin{array}{c}
T^{-1} W P x \\
-P x
\end{array}\right]=\left[\begin{array}{c}
\Lambda_{2} x \\
-P x
\end{array}\right] .
$$

By consulting the second component, we have $\Lambda_{1} h=W^{*} h=x$. Substituting $x=\Lambda_{1} h$ into the first component, yields $h=\Lambda_{2} \Lambda_{1} h$. Therefore $\Lambda_{1}$ is the right inverse of $\Lambda_{2}$ and $\Lambda_{1}^{-1}=\Lambda_{2}$.

## 5. Proof of the inner-outer factorization

Proof. Assume that the algebraic Riccati equation (3.14) admits a stabilizing solution $Q$. In other words, assume that $T_{R}$ is strictly positive, or equivalently, $G$ admits a full rank inner-outer factorization $G=G_{i} G_{o}$. Using $P=W_{c} W_{c}^{*}$ with $H_{G}=W_{o} W_{c}$, we have

$$
T_{G} T_{G}^{*}=T_{R}-H_{G} H_{G}^{*}=T_{R}-W_{o} P W_{o}^{*}
$$

Recall that the subspace $\mathfrak{H}\left(G_{i}\right)=\ell_{+}^{2}(\mathcal{Y}) \ominus T_{G_{i}} \ell_{+}^{2}(\mathcal{Y})$. Then

$$
\mathfrak{H}\left(G_{i}\right)=\operatorname{ker} T_{G_{i}}^{*}=\operatorname{ker} T_{G}^{*}=\operatorname{ker}\left(T_{R}-W_{o} P W_{o}^{*}\right) .
$$

It is noted that $\mathfrak{H}\left(G_{i}\right)$ is an invariant subspace for the backward shift $S_{\mathcal{Y}}^{*}$ on $\ell_{+}^{2}(\mathcal{Y})$. Recall that $Q=W_{o}^{*} T_{R}^{-1} W_{o}$. Let $k=\operatorname{dim} \operatorname{ker}(I-Q P)$, and put $\mathcal{X}_{i}=\mathbb{C}^{k}$. Let $U$ be an isometry from $\mathcal{X}_{i}$ onto $\operatorname{ker}(I-Q P)$. According to Lemma 4.1, the operator

$$
\Lambda_{2}=T_{R}^{-1} W_{o} P U=\left[\begin{array}{c}
C_{o} \\
C_{o} A_{o} \\
C_{o} A_{o}^{2} \\
C_{o} A_{o}^{3} \\
\vdots
\end{array}\right] P U: \mathcal{X}_{i} \rightarrow \mathfrak{H}\left(G_{i}\right)
$$

is invertible, where we also use (3.20). In particular, the dimension of $\mathfrak{H}\left(G_{i}\right)$ equals $\operatorname{dim} \mathcal{X}_{i}$. Since $P$ is invertible and $U$ is an isometry, the operator $P U$ from $\mathcal{X}_{i}$ into $\mathcal{X}$ is one to one.

Because $\mathfrak{H}\left(G_{i}\right)$ is an invariant subspace for the backward shift $S_{\mathcal{Y}}^{*}$, there exists an operator $A_{i}$ on $\mathcal{X}_{i}=\mathbb{C}^{k}$ such that

$$
\begin{equation*}
S_{\mathcal{Y}}^{*} T_{R}^{-1} W_{o} P U=T_{R}^{-1} W_{o} P U A_{i} . \tag{5.1}
\end{equation*}
$$

Since $T_{R}^{-1} W_{o} P U$ is one to one and $S_{\mathcal{Y}}^{* n}$ converges to zero pointwise, $A_{i}$ is stable.

Now observe that

$$
\left[\begin{array}{c}
C_{o} \\
C_{o} A_{o} \\
C_{o} A_{o}^{2} \\
\vdots
\end{array}\right] A_{o} P U=S_{\mathcal{Y}}^{*} T_{R}^{-1} W_{o} P U=\left[\begin{array}{c}
C_{o} \\
C_{o} A_{o} \\
C_{o} A_{o}^{2} \\
\vdots
\end{array}\right] P U A_{i}
$$

Since the observability matrix for $\left\{C_{o}, A_{o}\right\}$ is one to one, $A_{o} P U=P U A_{i}$. Because $P U$ is one to one, the spectrum of $A_{i}$ is contained in the spectrum of $A_{o}$. Multiplying $A_{o} P U=P U A_{i}$ by $U^{*} Q$ on the left and using $Q P U=U$ shows that

$$
\begin{equation*}
A_{o} P U=P U A_{i} \quad \text { and } \quad A_{i}=U^{*} Q A_{o} P U \tag{5.2}
\end{equation*}
$$

Setting $C_{i}=C_{o} P U$ and using $A_{o}^{j} P U=P U A_{i}^{j}$ for all positive integers $j$, we obtain

$$
\begin{equation*}
C_{o} A_{o}^{j} P U=C_{o} P U A_{i}^{j}=C_{i} A_{i}^{j} \quad(\text { for all integers } j \geq 0) \tag{5.3}
\end{equation*}
$$

In particular,

$$
T_{R}^{-1} W_{o} P U=\left[\begin{array}{c}
C_{o} \\
C_{o} A_{o} \\
C_{o} A_{o}^{2} \\
\vdots
\end{array}\right] P U=\left[\begin{array}{c}
C_{i} \\
C_{i} A_{i} \\
C_{i} A_{i}^{2} \\
\vdots
\end{array}\right] .
$$

Since $T_{R}^{-1} W_{o} P U$ is one to one, $\left\{C_{i}, A_{i}\right\}$ is a stable observable pair. Let $B_{i}$ mapping $\mathcal{Y}$ into $\mathcal{X}_{i}=\mathbb{C}^{k}$ and $D_{i}$ on $\mathcal{Y}$ be the complementary operators for the
pair $\left\{C_{i}, A_{i}\right\}$. Since $\mathfrak{H}\left(G_{i}\right)$ equals the range of $T_{R}^{-1} W_{o} P U$, the inner function $G_{i}$ (up to a unitary constant on the right) is given by

$$
G_{i}(\lambda)=D_{i}+\lambda C_{i}\left(I-\lambda A_{i}\right)^{-1} B_{i} .
$$

To find the outer factor $G_{o}$, first notice that

$$
Q=W_{o}^{*} T_{R}^{-1} W_{o}=\left[\begin{array}{lll}
C_{o}^{*} & A_{o}^{*} C_{o}^{*} & A_{o}^{* 2} C_{o}^{*} \ldots
\end{array}\right]\left[\begin{array}{c}
C \\
C A \\
C A^{2} \\
\vdots
\end{array}\right]=\sum_{j=0}^{\infty} A_{o}^{* j} C_{o}^{*} C A^{j}
$$

The second equality follows from (3.20). In other words, $Q$ satisfies the Stein equation

$$
\begin{equation*}
Q=A_{o}^{*} Q A+C_{o}^{*} C . \tag{5.4}
\end{equation*}
$$

Now note that $U^{*} P C_{o}^{*} C=C_{i}^{*} C$, so that $C_{i}^{*} C=U^{*} P\left(Q-A_{o}^{*} Q A\right)$. Moreover, $U^{*} P A_{o}^{*}=A_{i}^{*} U^{*} P$ and $U^{*} P Q=U^{*}$. Hence

$$
\begin{equation*}
C_{i}^{*} C=U^{*}-A_{i}^{*} U^{*} P Q A=U^{*}-A_{i}^{*} U^{*} A . \tag{5.5}
\end{equation*}
$$

It follows that $U^{*}=\sum_{j=0}^{\infty} A_{i}^{* j} C_{i}^{*} C A^{j}$.
Next observe that $T_{G}=T_{G_{i} G_{o}}=T_{G_{i}} T_{G_{o}}$. Multiplying by $T_{G_{i}}^{*}$ on the left, with the fact that $T_{G_{i}}$ is an isometry, we have $T_{G_{i}}^{*} T_{G}=T_{G_{o}}$. Using this with $U^{*}=\sum_{j=0}^{\infty} A_{i}^{* j} C_{i}^{*} C A^{j}$, we see that the first column of $T_{G_{o}}$ is given by

$$
\begin{aligned}
T_{G_{i}}^{*}\left[\begin{array}{c}
D \\
C B \\
C A B \\
\vdots
\end{array}\right] & =\left[\begin{array}{ccccc}
D_{i}^{*} & B_{i}^{*} C_{i}^{*} & B_{i}^{*} A_{i}^{*} C_{i}^{*} & B_{i}^{*} A_{i}^{* 2} C_{i}^{*} & \ldots \\
0 & D_{i}^{*} & B_{i}^{*} C_{i}^{*} & B_{i}^{*} A_{i}^{*} C_{i}^{*} & \ldots \\
0 & 0 & D_{i}^{*} & B_{i}^{*} C_{i}^{*} & \ldots \\
0 & 0 & 0 & D_{i}^{*} & \ldots \\
\vdots & \vdots & \vdots & \vdots & \ldots
\end{array}\right]\left[\begin{array}{c}
D \\
C B \\
C A B \\
C A^{2} B \\
\vdots
\end{array}\right] \\
& =\left[\begin{array}{c}
D_{i} D^{*}+B_{i}^{*} U^{*} B \\
\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right) B \\
\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right) A B \\
\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right) A^{2} B \\
\vdots
\end{array}\right] .
\end{aligned}
$$

By taking the Fourier transform of the first column of $T_{G_{o}}$, we obtain the following state space formula:

$$
G_{o}(\lambda)=D_{i}^{*} D+B_{i}^{*} U^{*} B+\lambda\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right)(I-\lambda A)^{-1} B .
$$

This completes the proof.
For completeness we shall also provide a slightly different derivation of the last part of the proof, that is, the formula for $G_{o}$. The idea is similar in nature but slightly different in execution: for $|\lambda|=1$ we compute $G_{o}(\lambda)=$ $G_{i}(\lambda)^{*} G(\lambda)$ using the realization formulas for $G_{i}$ and $G$. This leads to

$$
\begin{aligned}
G_{o}(\lambda)= & G_{i}(\lambda)^{*} G(\lambda) \\
= & \left(D_{i}^{*}+\bar{\lambda} B_{i}^{*}\left(I-\bar{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*}\right)\left(D+\lambda C(I-\lambda A)^{-1} B\right) \\
= & D_{i}^{*} D+\bar{\lambda} B_{i}^{*}\left(I-\bar{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*} D+\lambda D_{i}^{*} C(I-\lambda A)^{-1} B \\
& +|\lambda|^{2} B_{i}^{*}\left(I-\bar{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*} C(I-\lambda A)^{-1} B .
\end{aligned}
$$

Since we consider $|\lambda|=1$ this is equal to

$$
\begin{aligned}
G_{o}(\lambda)= & D_{i}^{*} D+\frac{1}{\lambda} B_{i}^{*}\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*} D+\lambda D_{i}^{*} C(I-\lambda A)^{-1} B \\
& +B_{i}^{*}\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*} C(I-\lambda A)^{-1} B .
\end{aligned}
$$

Consider the Stein equation $C_{i}^{*} C=U^{*}-A_{i}^{*} U^{*} A$; see (5.5). This may be used to compute

$$
\begin{aligned}
& \left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*} C(I-\lambda A)^{-1} \\
& =\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1}\left(U^{*}-\frac{1}{\lambda} A_{i}^{*} U^{*}(\lambda A)\right)(I-\lambda A)^{-1} \\
& =\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1}\left(U^{*}(I-\lambda A)+\left(I-\frac{1}{\lambda} A_{i}^{*}\right) U^{*}(\lambda A)\right)(I-\lambda A)^{-1} \\
& =\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1} U^{*}+\lambda U^{*} A(I-\lambda A)^{-1}
\end{aligned}
$$

Inserting this in the formula for $G_{o}(\lambda)$ we obtain

$$
\begin{aligned}
G_{o}(\lambda)= & D_{i}^{*} D+\frac{1}{\lambda} B_{i}^{*}\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1} C_{i}^{*} D+\lambda D_{i}^{*} C(I-\lambda A)^{-1} B+ \\
& +B_{i}^{*}\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1} U^{*} B+\lambda B_{i}^{*} U^{*} A(I-\lambda A)^{-1} B \\
= & D_{i}^{*} D+B_{i}^{*}\left(\lambda I-A_{i}^{*}\right)^{-1} C_{i}^{*} D \\
& +B_{i}^{*}\left(I+\left(I-\frac{1}{\lambda} A_{i}^{*}\right)^{-1}-I\right) U^{*} B \\
& +\lambda\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right)(I-\lambda A)^{-1} B \\
= & D_{i}^{*} D+B_{i}^{*} U^{*} B+B_{i}^{*}\left(\lambda I-A_{i}^{*}\right)^{-1}\left(C_{i}^{*} D+A_{i}^{*} U^{*} B\right) \\
& +\lambda\left(D_{i}^{*} C+B_{i}^{*} U^{*} A\right)(I-\lambda A)^{-1} B .
\end{aligned}
$$

Because $G_{o}$ is analytic in the open unit disc, we know that the term $B_{i}^{*}\left(\lambda I-A_{i}^{*}\right)^{-1}\left(C_{i}^{*} D+A_{i}^{*} U^{*} B\right)$ must be zero. Let us give a direct proof of this fact. This turns out to be an easy consequence of formula (3.23) in [9]. Indeed, this formula states that

$$
C_{1}^{*} C_{1}=(Q-Q P Q)-A_{0}^{*}(Q-Q P Q) A_{0},
$$

where $C_{1}=D^{*} C_{0}+B^{*} Q A_{0}$. Multiplying the above formula with $P U$ on the right and $U^{*} P$ on the left, we obtain

$$
U^{*} P C_{1}^{*} C_{1} P U=U^{*} P(Q-Q P Q) P U-U^{*} P A_{0}^{*}(Q-Q P Q) A_{0} P U
$$

Since $Q P U=U$ it follows that $(Q-Q P Q) P U=0$, so the first term on the right hand side is zero. Further, since $A_{0} P U=P U A_{i}$ it follows that also the second term on the right hand side is zero. Hence $C_{1} P U=0$, which means

$$
0=\left(D^{*} C_{0}+B^{*} Q A_{0}\right) P U=D^{*} C_{i}+B^{*} Q P U A_{i}=D^{*} C_{i}+B^{*} U A_{i} .
$$

Thus the formula for $G_{o}$ can also be established by a direct computation using the realizations of $G_{i}$ and $G$.

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A.E. Frazho<br>Department of Aeronautics and Astronautics, Purdue University West Lafayette, IN 47907<br>USA<br>e-mail: frazho@ecn.purdue.edu<br>A.C.M. Ran<br>Department of Mathematics, Faculty of Science, VU Amsterdam<br>De Boelelaan 1081a, 1081 HV Amsterdam<br>The Netherlands<br>and<br>Unit for BMI, North-West University<br>Potchefstroom<br>South Africa<br>e-mail: a.c.m.ran@vu.nl


[^0]:    E-mail address:
    vuresearchportal.ub@vu.nl

