## Research Note

# CVBEM for a system of second-order elliptic partial differential equations 

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#### Abstract

A boundary element method based on the Cauchy's integral formulae and the theory of complex hypersingular integrals is devised for the numerical solution of boundary value problems governed by a system of second-order elliptic partial differential equations. The elliptic system has applications in physical problems involving anisotropic media. © 1998 Elsevier Science Ltd. All rights reserved


Key words: complex variable boundary element method, elliptic partial differential equations, anisotropic media.

## 1 INTRODUCTION

Consider the system of second-order elliptic partial differential equations given by

$$
\begin{equation*}
\sum_{j=1}^{2} \sum_{p=1}^{2} \sum_{k=1}^{N} a_{i j k p} \frac{\partial^{2} \phi_{k}}{\partial x_{j} \partial x_{p}}=0(i=1,2, \cdots, N), \tag{1}
\end{equation*}
$$

where $\phi_{k}(k=1,2, \cdots, N)$ are functions of $x_{1}$ and $x_{2}$ and $a_{i j k p}(j, p=1,2$ and $i, k=1,2, \cdots$,
$N$ ) are real constant coefficients which satisfy the symmetry conditions $a_{i j k p}=a_{k p i j}$ and are such that

$$
\begin{equation*}
\sum_{j=1}^{2} \sum_{p=1}^{2} \sum_{i=1}^{N} \sum_{k=1}^{N} a_{i j k p} \lambda_{i j} \lambda_{k p}>0 \tag{2}
\end{equation*}
$$

for every non-zero $N \times 2$ real matrix $\left[\lambda_{i j}\right]$.
We are interested in solving eqn (1) in a region $\mathcal{R}$ bounded by a simple closed curve $C$ (on the $0 x_{1} x_{2}$ plane) subject to

$$
\left.\begin{array}{lll}
\phi_{k}\left(x_{1}, x_{2}\right)=\mu_{k}\left(x_{1}, x_{2}\right) & \text { for } \quad\left(x_{1}, x_{2}\right) \in C_{1}  \tag{3}\\
P_{i}\left(x_{1}, x_{2}\right)=Q_{i}\left(x_{1}, x_{2}\right) & \text { for } \quad\left(x_{1}, x_{2}\right) \in C_{2}
\end{array}\right\}
$$

where $\mu_{k}$ and $Q_{i}$ are suitably prescribed functions of $x_{1}$ and $x_{2}, C_{1}$ and $C_{2}$ are non-intersecting curves such that $C=C_{1} \cup$ $\mathcal{C}_{2}$ and

$$
\begin{equation*}
P_{i}=\sum_{j=1}^{2} \sum_{p=1}^{2} \sum_{k=1}^{N} a_{i j k p} \frac{\partial \phi_{k}}{\partial x_{p}} n_{j}(i=1,2, \cdots, N) \tag{4}
\end{equation*}
$$

[^0]with $n_{j}(j=1,2)$ being components of the unit outer normal vector to $R$ on $C$.

The boundary value problem defined by eqns (1) and (3) has important applications in engineering. As an example, the steady-state temperature distribution in a flat plate which is thermally anisotropic and homogeneous obeys eqn (1) with $N=1$. The temperature and heat flux are given by $\phi_{1}$ and ( $P_{1}, P_{2}$ ) respectively, and $a_{1 j 1 p}$ are the heat conduction coefficients.

The plane static deformation of a homogeneous anisotropic elastic solid is governed by eqn (1) with $N=2$ and $x_{1}$ and $x_{2}$ as the Cartesian coordinates. The Cartesian displacement and traction are given by ( $\phi_{1}, \phi_{2}$ ) and ( $P_{1}, P_{2}$ ) respectively. The coefficients $a_{i j k p}$ are the elastic moduli of the material occupying the solid. For a specific case, the elastostatic behaviour of a transversely isotropic material which has transverse planes perpendicular to the $0 x_{1} x_{2}$ plane and which undergoes plane deformation is governed by

$$
\begin{align*}
& C \frac{\partial^{2} \phi_{1}}{\partial x_{1}^{2}}+L \frac{\partial^{2} \phi_{1}}{\partial x_{2}^{2}}+(F+L) \frac{\partial^{2} \phi_{2}}{\partial x_{1} \partial x_{2}}=0,  \tag{5}\\
& C \frac{\partial^{2} \phi_{2}}{\partial x_{2}^{2}}+L \frac{\partial^{2} \phi_{2}}{\partial x_{1}^{2}}+(F+L) \frac{\partial^{2} \phi_{1}}{\partial x_{1} \partial x_{2}}=0,
\end{align*}
$$

which is a special case that can be recovered from eqn (1) if we let $N=2$ and $a_{2222}=A, a_{1111}=C, a_{1122}=a_{2211}=F$, $a_{1212}=a_{2121}=a_{1221}=a_{2112}=L$ and the remaining $a_{i j k l}$ be zero. [The constants $A, F, C$ and $L$ are independent elastic coefficients.]


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