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# Diagonalizable Extended Backward Differentiation Formulas 

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#### Abstract

We generalize the extended backward differentiation formulas (EBDFs) introduced by Cash and by Psihoyios and Cash such that the system matrix in the modified Newton process can be block-diagonalized. This enables an efficient parallel implementation. We construct methods which are L-stable up to order $p=6$ with the same computational complexity per processor as the conventional BDF methods. Numerical experiments with the order 6 method show that a speedup factor between 2 and 4 on four processors can be expected.


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## 1. Introduction

In [4] we discussed the parallel implementation of the extended backward differentiation formulas (EBDFs) introduced by Cash in [1] and [2] for the numerical solution of initial value problems for stiff differential equations of the form

$$
\begin{equation*}
\frac{d \mathbf{y}}{d t}=\mathbf{f}(t, \mathbf{y}), \mathbf{y}, \mathbf{f} \in \mathbb{R}^{d}, t \geq t_{0} \tag{1.1}
\end{equation*}
$$

The parallel approach described in [4] is based on block-diagonalization of the system matrix in the modified Newton process used for solving the implicit EBDF relations. The system matrix is of the form $I-(A \otimes h J)$, where $h$ is the stepsize, $I$ is the identity matrix, the matrix $A$ is determined by the EBDF method coefficients, and $J$ is an approximation to the Jacobian matrix $\partial \mathbf{f} / \partial \mathbf{y}$. Since exact block-diagonalization is not possible due to defectiveness of the matrix $A$, we applied approximate block-diagonalization. Although the rate of convergence is less than that of true modified Newton, the experiments in [4] show a speedup on a three-processor configuration of between 2 and 3 .

The same parallel approach can be applied to the more general EBDF methods which have recently been proposed by Psihoyios and Cash [12]. These more general EBDF methods also lead to defective coefficient matrices $A$ in the modified Newton process, but have the property that they can be made L-stable up to order $p=6$ (the original EBDFs are L-stable up to order $p=4$ ). However, approximate block diagonalization is now much less accurate than in the case of the
original EBDF methods. The aim of this paper is to construct methods which are L-stable up to order $p=6$ with a nondefective matrix $A$, so that exact block-diagonalization is possible.

In Section 2, we define a family of EBDF-type methods which generalizes the Cash and PsihoyiosCash methods. The order conditions, the global error for the Prothero-Robinson test equation, and stability conditions are derived. Section 3 discusses the sequential and parallel implementation of these methods and in Section 4 we derive L-stable, nondefective EBDF methods of order up to $p=6$. Per processor, the computational complexity of these methods is comparable to that of the conventional BDF methods. Finally, Section 5 reports numerical experiments for the sixth-order method. These experiments indicate that a speedup factor in the range of 2 to 4 on four processors can be expected.

## 2. EBDF-TYPE METHODS

The generalizations of the EBDF methods to be discussed in this paper are of the form

$$
\begin{align*}
& (B \otimes I) \mathbf{Y}_{n+1}-h(C \otimes I) \mathbf{F}\left(\mathbf{e} t_{n}+\mathbf{c} h, \mathbf{Y}_{n+1}\right)=(E \otimes I) \mathbf{V}_{n} \\
& \mathbf{V}_{n}=\left(\mathbf{y}_{n-s+1}^{T}, \ldots, \mathbf{y}_{n}^{T}\right)^{T} \tag{2.1}
\end{align*}
$$

Here, $\otimes$ denotes the Kronecker product, $h$ is the stepsize $t_{n+1}-t_{n}$, e and $\mathbf{c}$ are $r$-dimensional vectors, $\mathbf{e}=(1, \ldots, 1)^{T}, \mathbf{c}=\left(c_{1}, \ldots, c_{r}\right)^{T}$ with $c_{r}=1 . I$ is the $d$ by $d$ identity matrix, $B$ and $C$ are $r$ by $r$ lower triangular matrices and $E$ is an $r$ by $s$ matrix. The unknown stage vector $\mathbf{Y}_{n+1}$ contains $r$ stages $\mathbf{y}_{n+c_{i}}$ of dimension $d$, representing numerical approximations at the points $t_{n}+c_{i} h$, and $\mathbf{F}\left(\mathbf{e} t_{n}+\mathbf{c h}, \mathbf{Y}_{n+1}\right)$ contains the $r$ righthand side values $\mathbf{f}\left(t_{n}+c_{i} h, y_{n+c_{i}}\right)$. Since $B$ and $C$ are lower triangular, the first $r-1$ stage equations may be considered to be predictor formulas providing the internal stage values $\mathbf{y}_{n+c_{i}}, i=1, \ldots, r-1$, needed in the last stage equation. This last stage equation will be referred to as the corrector equation defining the output or step point value $y_{n+c_{r}}=y_{n+1}$.

We shall call (2.1) an EBDF-type method, because it can be viewed as a generalization of the original three-stage EBDF and MEBDF methods of Cash and the four-stage version recently discussed by Psihoyios and Cash. Note that the one-stage versions with $c_{1}=1$ assume the form of the conventional BDF methods.

### 2.1 Nonstiff order of accuracy

Given the abscissa vector $\mathbf{c}=\left(c_{i}\right)$, the matrices $B, C$ and $E$ can be determined such that the $i$ th stage equation in (2.1) is consistent of order $p_{i}$ provided that $p_{i}+1$ free coefficients are available for that equation. To formulate the consistency conditions, we first write (1.1) in autonomous form by adding the equation $d y_{d+1} / d t=1$, so that (2.1) also becomes autonomous. Next, we introduce the abscissa vector for the back-values $\mathbf{b}:=(1-s, 2-s, \ldots, 0)^{T}$, and the component-wise notation $g(\mathbf{v})$ associated with a scalar function $g: \mathbb{R} \rightarrow \mathbb{R}$ to denote the vector with components $g\left(v_{i}\right)$, where $\mathbf{v}=\left(v_{i}\right)$. Upon substitution of the exact solution into (2.1), that is, we set $\mathbf{Y}_{n+1}=$ $\left.\exp (\mathbf{c} h d / d t) \otimes \mathbf{y}(t)\right|_{t_{n}}$ and $\mathbf{V}_{n}=\left.\exp (\mathbf{b} h d / d t) \otimes \mathbf{y}(t)\right|_{t_{n}}$, it is easily seen that the $i$ th stage equation in (2.1) is consistent of order $p_{i}$ if

$$
\begin{equation*}
\mathbf{e}_{i}^{T}((B-h C) \exp (\mathbf{c} h)-E \exp (\mathbf{b} h))=O\left(h^{p_{i}+1}\right), i=1, \ldots, r, \tag{2.2}
\end{equation*}
$$

where $\mathbf{e}_{i}$ is the $i$ th unit vector. The conditions (2.2) lead to the order equations

$$
\begin{equation*}
\mathbf{e}_{i}^{T} E \mathbf{b}^{j}=\mathbf{e}_{i}^{T}\left(B \mathbf{c}^{j}-j C \mathbf{c}^{j-1}\right), j=0, \ldots, p_{i}, i=1, \ldots, r \tag{2.3}
\end{equation*}
$$

where we define $0^{0}=1$. If (2.3) is satisfied, then the stage order of (2.1) is defined by $\bar{p}:=\min \left\{p_{i}\right\}$. In general, the output value $\mathbf{y}_{n+c_{r}}=\mathbf{y}_{n+1}$ has nonstiff order of accuracy $p=\bar{p}$. However, if the first $r-1$ entries of the last row of the matrix $B$ in (2.1) vanish (as will be the case for the methods of Section 4) and if $p_{r}=\bar{p}+1$ (as will henceforth be assumed), then the nonstiff order of accuracy is equal to $\bar{p}+1$. The stiff order of accuracy is discussed in the following section.

### 2.2 Stiff order of accuracy

We study the global error of the EBDF-type method (2.1) when applied to the Prothero-Robinson equation $d y(t) / d t=\lambda y(t)+\phi(t)$, where $\phi$ is a given function. By means of this test equation we can obtain insight into the behavior of the error components in the integration of the general ODE system (1.1) by interpreting $\lambda$ as an eigenvalue of the matrix $J$, where $J$ denotes the Jacobian of the ODE system. For general linear methods, Hundsdorfer [9] has derived an upper bound for the global error, so that by rewriting (2.1) as a general linear method, we can use his results. However, the rather special form of (2.1) makes it easier to derive such error bounds directly.

Applying (2.1) to the Prothero-Robinson equation yields the EBDF solution

$$
\begin{equation*}
y_{n+1}=\mathbf{e}_{r}^{T} \mathbf{Y}_{n+1}=\mathbf{e}_{r}^{T}(B-z C)^{-1}\left(h C \phi\left(\mathbf{e} t_{n}+\mathbf{c} h\right)+E \mathbf{V}_{n}\right), z:=h \lambda \tag{2.4}
\end{equation*}
$$

Furthermore, upon substitution of the exact solution $y(t)$ into (2.4) we define the local error $\delta_{n+1}$ by the relation

$$
\begin{align*}
& y\left(t_{n+1}\right)=\mathbf{e}_{r}^{T}(B-z C)^{-1}\left(h C \phi\left(\mathbf{e} t_{n}+\mathbf{c} h\right)+E \tilde{\mathbf{V}}_{n}\right)+\delta_{n+1}  \tag{2.5}\\
& \tilde{\mathbf{V}}_{n}=\left(y\left(t_{n-s+1}\right), \ldots, y\left(t_{n}\right)\right)^{T}
\end{align*}
$$

By subtracting (2.4) from (2.5) and defining the global errors $\varepsilon_{n}:=y\left(t_{n}\right)-y_{n}$, we obtain

$$
\begin{equation*}
\varepsilon_{n+1}=\mathbf{e}_{r}^{T}(B-z C)^{-1} E\left(\varepsilon_{n-s+1}, \ldots, \varepsilon_{n}\right)^{T}+\delta_{n+1} . \tag{2.6}
\end{equation*}
$$

From this global error recursion we derive the following result:

Theorem 2.1 Let $\bar{p}$ be the stage order of the EBDF-type method (2.1). Then, the global error of the Prothero-Robinson equation behaves according to $\varepsilon_{n+1}=O\left(z^{-1} h^{\bar{p}+1}\right)$ as $h \rightarrow 0$ and $z=h \lambda \rightarrow$ $\infty$.

Proof. First an explicit expression for the global error $\varepsilon_{n+1}$ in terms of the local errors $\delta_{i}$ is derived. In this derivation, it is convenient to rewrite the multistep difference equation (2.6) in one-step form. Let us define the $s$-dimensional vector function $\mathbf{u}(z)$, the $s$-dimensional local error vector $\boldsymbol{\delta}_{n+1}$, and the $s$-dimensional global error vector $\boldsymbol{\varepsilon}_{n+1}$ by

$$
\begin{equation*}
\mathbf{u}^{T}(z):=\mathbf{e}_{r}^{T}(B-z C)^{-1} E, \quad \boldsymbol{\delta}_{n+1}:=\delta_{n+1} \mathbf{e}_{s}, \quad \varepsilon_{n+1}:=\left(\varepsilon_{n-s+2}, \ldots, \varepsilon_{n+1}\right)^{T} . \tag{2.7}
\end{equation*}
$$

Then, assuming that $\varepsilon_{0}=0$, we obtain

$$
\begin{aligned}
& \varepsilon_{n+1}=R \varepsilon_{n}+\boldsymbol{\delta}_{n+1}=\sum_{i=0}^{n} R^{i} \boldsymbol{\delta}_{n+1-i}, \\
& R=R(z):=\left(\begin{array}{ccccc}
1 & & & \\
& & 1 & & \\
& & & \ddots & \\
u_{1}(z) & u_{2}(z) & u_{3}(z) & \cdots & u_{s}(z)
\end{array}\right) .
\end{aligned}
$$

Applying partial summation (see [6], p. 242), we arrive at the expression

$$
\begin{equation*}
\varepsilon_{n+1}=\mathbf{e}_{s}^{T}(I-R)^{-1}\left(I-R^{n+1}\right) \mathbf{e}_{s} \delta_{1}+\mathbf{e}_{s}^{T} \sum_{i=1}^{n}(I-R)^{-1}\left(I-R^{n+1-i}\right) \mathbf{e}_{s}\left(\delta_{i+1}-\delta_{i}\right) \tag{2.8}
\end{equation*}
$$

provided that $I-R$ is nonsingular.

Next, we express $\varepsilon_{n+1}$ in terms of the derivatives of the exact solution $y(t)$. Using the relation $\phi(t)=y^{\prime}(t)-\lambda y(t)$, it follows from (2.5) that $\delta_{n+1}$ has the Taylor expansion

$$
\begin{align*}
\delta_{n+1} & =\sum_{j=0}^{\infty} \frac{1}{j!} \gamma_{j}(z) h^{j} y^{(j)}\left(t_{n}\right)  \tag{2.9}\\
\gamma_{0}(z) & :=1-\mathbf{e}_{r}^{T}(B-z C)^{-1}(E \mathbf{e}-z C \mathbf{e}) \\
\gamma_{j}(z) & :=1-\mathbf{e}_{r}^{T}(B-z C)^{-1}\left(E \mathbf{b}^{j}+j C \mathbf{c}^{j-1}-z C \mathbf{c}^{j}\right), j \geq 1
\end{align*}
$$

From (2.3) we see that $E \mathbf{b}^{j}=B \mathbf{c}^{j}-j C \mathbf{c}^{j-1}$ for $j=0, \ldots, \bar{p}, \bar{p}$ being the stage order, so that the first $\bar{p}+1$ terms in the Taylor expansion vanish. Since (2.7) implies $\mathbf{u}^{T}(z)=-z^{-1} \mathbf{e}_{r}^{T} C^{-1} E+$ $O\left(z^{-2}\right)$, it follows from the structure of the matrix $R(z)$ that $R(z)^{n}$ vanishes as $z \rightarrow \infty$ for $n \geq s$. Hence, we conclude from (2.8) and (2.9) that

$$
\begin{aligned}
\varepsilon_{n+1} & =\mathbf{e}_{s}^{T}(I-R(z))^{-1} \mathbf{e}_{s} \delta_{1}(z)+O\left(h^{\bar{p}+2}\right) \\
& =\frac{1}{(\bar{p}+1)!} \mathbf{e}_{s}^{T}(I-R(z))^{-1} \gamma_{\bar{p}+1}(z) \mathbf{e}_{s} h^{\bar{p}+1} y^{(\bar{p}+1)}\left(t_{n}\right)+O\left(h^{\bar{p}+2}\right) \text { as } z \rightarrow \infty, n \geq s .
\end{aligned}
$$

The theorem now follows from the fact that $(I-R(z))^{-1}$ is bounded as $z \rightarrow \infty$ and that $\gamma_{\bar{p}+1}(z)=$ $1-\mathbf{e}_{r}^{T} \mathbf{c}^{\bar{p}+1}+O\left(z^{-1}\right)=O\left(z^{-1}\right)$.
If the stiff order of accuracy is defined by the order of $\varepsilon_{n+1}$ in $h$ as $z \rightarrow \infty$, then we conclude from this theorem that the stiff order of EBDF-type methods is $\bar{p}+1$. This favorable property of EBDF-type methods perhaps explains the impressive performance of the MEBDF methods observed in [11].

### 2.3 Stability

From the linear difference equation (2.4) it follows that, with respect to the stability test equation $y^{\prime}=\lambda y$ (i.e. $\phi=0$ ), EBDF-type methods are stable if the characteristic equation associated with (2.4) has roots only on the unit disk. Using the identity

$$
\mathbf{p}^{T} P^{-1} \mathbf{q}=\frac{\operatorname{det}\left(P+\mathbf{q} \mathbf{p}^{T}\right)}{\operatorname{det}(P)}-1
$$

which holds for any nonsingular $m$ by $m$ matrix $P$ and any two $m$-dimensional vectors $\mathbf{p}$ and $\mathbf{q}$, we find that (2.4) can be written as

$$
y_{n+1}=\mathbf{e}_{r}^{T}(B-z C)^{-1} E \mathbf{V}_{n}=\frac{\operatorname{det}\left(B-z C+E \mathbf{V}_{n} \mathbf{e}_{r}^{T}\right)}{\operatorname{det}(B-z C)}-1
$$

Hence, the characteristic equation is given by

$$
\begin{equation*}
\zeta^{s}=\frac{\operatorname{det}\left(B-z C+E \boldsymbol{\Gamma}(\zeta) \mathbf{e}_{r}^{T}\right)}{\operatorname{det}(B-z C)}-1, \boldsymbol{\Gamma}(\zeta):=\left(1, \zeta, \ldots, \zeta^{s-2}, \zeta^{s-1}\right)^{T} \tag{2.10}
\end{equation*}
$$

First of all, we require that (2.1) is zero-stable; that is, we require that for $z=0$ the characteristic equation (2.10) has one simple root at 1 and $s-1$ roots on the unit disk with only simple roots on the unit circle.

Theorem 2.2 Let $B$ be nonsingular and let the row vectors of the matrix $B^{-1} E$ be denoted by $\mathbf{w}_{i}^{T}$. The EBDF-type method (2.1) is zero-stable if the equation $\zeta^{s}=\mathbf{w}_{r}^{T} \boldsymbol{\Gamma}(\zeta)$ has one simple root $\zeta_{1}=1$ and $s-1$ roots $\zeta_{i}, i=2, \ldots, s-1$, on the unit disk with only simple roots on the unit circle.

Proof. For $z=0$, the characteristic equation (2.10) simplifies to

$$
\zeta^{s}=\frac{\operatorname{det}\left(B+E \boldsymbol{\Gamma}(\zeta) \mathbf{e}_{r}^{T}\right)}{\operatorname{det} B}-1=\operatorname{det}\left(I+B^{-1} E \boldsymbol{\Gamma}(\zeta) \mathbf{e}_{r}^{T}\right)-1
$$

The matrix $B^{-1} E \boldsymbol{\Gamma}(\zeta) \mathbf{e}_{r}^{T}$ has zero columns, except for its last column which has entries $\mathbf{w}_{i}^{T} \boldsymbol{\Gamma}(\zeta)$, $i=1, \ldots, r$. Hence, $\left(2.4^{\prime}\right)$ reduces to $\zeta^{s}=\mathbf{w}_{r}^{T} \boldsymbol{\Gamma}(\zeta)$, which proves the assertion of the theorem.

Note that this theorem holds for any general linear method of the form (2.1) such that the output (step point) value is given by one of the stages, regardless of the structures of the matrices $B, C$ and $E$.

Secondly, the stability region of (2.1) is defined by the points in the $z$-plane where the zeros of (2.10) are on the unit disk. Setting $\zeta=\exp (i \theta)$, the boundary of this region is defined by the boundary locus equation

$$
\begin{equation*}
\left(\mathrm{e}^{i s \theta}+1\right) \operatorname{det}(B-z C)-\operatorname{det}\left(B-z C+E \boldsymbol{\Gamma}\left(\mathrm{e}^{i \theta}\right) \mathbf{e}_{r}^{T}\right)=0,0 \leq \theta<2 \pi \tag{2.11}
\end{equation*}
$$

This equation can be used for plotting stability regions.
Finally, we remark that an $\mathrm{A}(\alpha)$-stable method is automatically $\mathrm{L}(\alpha)$-stable, because the characteristic equation (2.10) reduces to $\zeta^{s}=0$ as $z \rightarrow \infty$.

## 3. SEQUENTIAL AND PARALLEL iteration

The solution of (2.1) can be obtained by successively solving $r$ subsystems, each of dimension $d$ (recall that $B$ and $C$ are assumed to be lower triangular). If a (modified) Newton method is applied, then the iteration scheme for the $i$ th stage $\mathbf{y}_{n+c_{i}}$ of $\mathbf{Y}_{n+1}$ assumes the form

$$
\begin{align*}
\left(I-h \tilde{C}_{i i} J\right)\left(\mathbf{y}_{n+c_{i}}^{(j)}-\mathbf{y}_{n+c_{i}}^{(j-1)}\right) & =-\mathbf{y}_{n+c_{i}}^{(j-1)}+h \tilde{C}_{i i} \mathbf{f}\left(t_{n+c_{i}}, \mathbf{y}_{n+c_{i}}^{(j-1)}\right) \\
& +h \sum_{k=1}^{i-1} \tilde{C}_{i k} \mathbf{f}\left(t_{n+c_{k}}, y_{n+c_{k}}\right)+\sum_{k=1}^{s} \tilde{E}_{i k} \mathbf{y}_{n-s+k}, j=1, \ldots, m_{i}, \tag{3.1}
\end{align*}
$$

where $\tilde{C}_{i k}$ and $\tilde{E}_{i k}$ denote the entries of the matrices $B^{-1} C$ and $B^{-1} E$, respectively, $J$ is an approximation to the Jacobian matrix of the righthand side function in (1.1) at $t_{n+1}$, and $\mathbf{y}_{n+c_{i}}^{(0)}$ is an initial approximation to $\mathbf{y}_{n+c_{i}}$. This amounts to the solution of $\bar{m} r$ linear systems per step, where $\bar{m}$ denotes the (average) number of Newton iterations needed in the $r$ subsystems. This approach will be called sequential iteration.

If, however, a parallel computer system is available, then one may change to a more efficient parallel approach. In [4] we developed for the original EBDF and MEBDF methods of Cash a highly parallel iterative method for solving the implicit relations in (2.1). This parallel approach can also be applied to methods of the form (2.1) with more general matrices $B, C$ and $E$. It is based on the approximate block-diagonalization of the modified Newton method applied to the full (block) system (2.1). Let us define the residue function

$$
\begin{equation*}
\mathbf{R}_{n}(\mathbf{Y}):=\mathbf{Y}-h\left(B^{-1} C \otimes I\right) \mathbf{F}\left(\mathbf{e} t_{n}+\mathbf{c} h, \mathbf{Y}\right)-\left(B^{-1} E \otimes I\right) \mathbf{V}_{n} \tag{3.2}
\end{equation*}
$$

Then, solving (2.1) by $m$ modified Newton iterations amounts to

$$
\begin{equation*}
\left(I-B^{-1} C \otimes h J\right)\left(\mathbf{Y}_{n+1}^{(j)}-\mathbf{Y}_{n+1}^{(j-1)}\right)=-\mathbf{R}_{n}\left(\mathbf{Y}_{n+1}^{(j-1)}\right), j=1, \ldots, m \tag{3.3}
\end{equation*}
$$

Unfortunately, if we use an abscissa vector of the form $\mathbf{c}=(1,2, \ldots, r-1,1)^{T}$ and assume the same zero structure of the matrices $B, C$ and $E$ as in the original EBDF and MEBDF methods, then the matrix $B^{-1} C$ is defective, so that we cannot directly diagonalize (3.1) by applying a similarity transformation. One option is to replace the matrix $B^{-1} C$ in (3.3) by a diagonalizable approximation $A^{*}$, for example, by $A^{*}=\operatorname{diag}\left(B^{-1} C\right)$. The rate of convergence will be less than that of the modified Newton method, however. In the case of the three-stage EBDF and MEBDF methods, the loss in rate of convergence is modest (see the experiments in [4]) because the diagonalizable approximation is quite accurate. In fact, even with the simple choice $A^{*}=\operatorname{diag}\left(B^{-1} C\right)$, we obtained surprisingly fast convergence. However, for higher-stage methods, where diagonalizable approximations are less accurate, the rate of convergence is expected to decrease significantly.

### 3.1 Nondefective methods

Rather than applying approximate block-diagonalization, we follow an alternative approach in which the abscissa vector is changed to the form $\mathbf{c}=\left(c_{1}, 2,3, \ldots, r-1,1\right)^{T}$ and in which we choose $c_{1} \neq 1$ such that $B^{-1} C$ is no longer a defective matrix (except for the degenerate case $s<r-2$ ). We shall call such EBDF methods nondefective EBDF methods. Nondefective EBDF methods can directly be diagonalized by the transformation $\tilde{\mathbf{Y}}^{(j)}=\left(Q^{-1} \otimes I\right) \mathbf{Y}^{(j)}$, where $Q$ is such that $Q^{-1}\left(B^{-1} C\right) Q=D$ with $D$ diagonal. This yields the transformed iteration method

$$
\begin{align*}
(I-D \otimes h J)\left(\tilde{\mathbf{Y}}^{(j)}-\tilde{\mathbf{Y}}^{(j-1)}\right) & =-\left(Q^{-1} \otimes I\right) \mathbf{R}_{n}\left((Q \otimes I) \tilde{\mathbf{Y}}^{(j-1)}\right), j=1, \ldots, m  \tag{3.4}\\
\mathbf{Y}_{n+1} & =(Q \otimes I) \tilde{\mathbf{Y}}^{(m)}
\end{align*}
$$

and will be called parallel iteration. We emphasize that (3.4) is algebraically equivalent to (3.3).
The introduction of the free parameter $c_{1}$ in the abscissa vector $\mathbf{c}=\left(c_{1}, 2, \ldots, r-1,1\right)^{T}$ preserves the attractive property that all stage values, except for the first one, can be reused in the initial approximation $\mathbf{Y}^{(0)}$ needed in the succeeding time step. In fact, setting $c_{1}=1$ has no additional advantages, because it 'duplicates' the output value at $t_{n+1}$.

### 3.2 Convergence condition

Defining the iteration error $\boldsymbol{\varepsilon}^{(j)}:=\mathbf{Y}^{(j)}-\mathbf{Y}_{n+1}$, we derive for (3.3) the error recursion

$$
\begin{aligned}
\varepsilon^{(j)} & =h K \Phi\left(\varepsilon^{(j-1)}\right), j=1, \ldots, m \\
K & :=\left(I-B^{-1} C \otimes h J\right)^{-1}\left(B^{-1} C \otimes I\right) \\
\Phi(\varepsilon) & :=\mathbf{F}\left(\mathbf{e} t_{n}+\mathbf{c} h, \mathbf{Y}_{n+1}+\boldsymbol{\varepsilon}\right)-\mathbf{F}\left(\mathbf{e} t_{n}+\mathbf{c h}, \mathbf{Y}_{n+1}\right)-(I \otimes J) \varepsilon
\end{aligned}
$$

Let $\Phi(\varepsilon)$ have at $\varepsilon=0$ a Lipschitz constant $L_{\Phi}$ with respect to the Euclidean norm and let the problem be dissipative, i.e. $\mu_{2}[J] \leq 0$, where $\mu_{2}[\cdot]$ denotes the logarithmic norm associated with the Euclidean norm. Then, by applying the matrix version of von Neumann's theorem (see [6], p. 356), we conclude that for dissipative problems

$$
\begin{equation*}
\left\|\varepsilon^{(j)}\right\|_{2} \leq h L_{\Phi} L_{K}\left\|\varepsilon^{(j-1)}\right\|_{2}, L_{K}=\max \left\{\left\|\left(I-z B^{-1} C\right)^{-1} B^{-1} C\right\|_{2}: \operatorname{Re}(z) \leq 0\right\} \tag{3.5}
\end{equation*}
$$

Hence, for dissipative problems, a sufficient condition for convergence is

$$
h \leq \frac{1}{L_{\Phi} L_{K}}
$$

Thus, difference in convergence of two EBDF-type methods is mainly determined by differences in the upper bound $L_{K}$.

### 3.3 Analysis of computational expenses

Finally, the computational expenses of (3.1) when implemented on one processor (sequential iteration of the subsystems) are compared with those of (3.4) implemented on $r$ processors (parallel iteration). In (3.1) we define $\bar{m}:=r^{-1}\left(m_{1}+\cdots+m_{r}\right)$ and we denote the number of distinct diagonal entries of $C$ by $r_{0}$. Table 3.1 lists the numbers of floating point operations to advance the solution one time step using a fixed stepsize. In this table, $C_{f}$ and $C_{J}$ respectively denote the average numbers of operations needed to compute a component of $f$ and an entry of $J$.

For a linear problem, only one Newton iteration is needed. Hence, assuming that the costs of building and factoring the Jacobian are negligible, it follows from Table 3.1 that the parallel speedup can be estimated by

$$
S=r \frac{\left(2-r^{-1}\right) C_{f}+2 d+2 s+6-2 r^{-1}}{C_{f}+2 d+2 s+3 r+3} .
$$

Table 3.1: Operation costs per processor to advance the solution one time step.

|  | Sequential iteration | Parallel iteration |
| :--- | :---: | :---: |
| Once per Jacobian update |  |  |
| Jacobian evaluation | $C_{J} d^{2}$ | $\frac{1}{r} C_{J} d^{2}$ |
| System matrix | $r_{0} d$ | $d$ |
| LUD of system matrix | $\frac{2}{3} r_{0} d^{3}$ | $\frac{2}{3} d^{3}$ |
| Once per time step | $\left(C_{f}+2 s+1\right) r d-\left(C_{f}+2\right) d$ | $(2 s-1) d$ |
| Righthand side |  |  |
| Per Newton iteration | $2 r d^{2}$ | $2 d^{2}$ |
| Forward/backward | $r\left(C_{f}+5\right) d$ | $\left(C_{f}+r+4\right) d$ |
| Updates | - | $2 r d$ |
| Transformations |  |  |

At the other extreme, assume a very stiff nonlinear problem such that the Jacobian must be evaluated once per step. Then, we obtain

$$
S=r \frac{\bar{m}\left(2 d+C_{f}+5\right)+\left(C_{f}+2 s+1\right)+r^{-1}\left(C_{J} d-C_{f}-2\right)+r_{0} r^{-1}\left(1+\frac{2}{3} d^{2}\right)}{m\left(2 d+C_{f}+3 r+4\right)+(2 s-1)+r^{-1} C_{J} d+\left(1+\frac{2}{3} d^{2}\right)},
$$

from which the following observations can be made:

- If the evaluation of the Jacobian dominates the computation, then $S \approx r$.
- If factoring the Jacobian dominates the computation, then $S \approx r_{0}$.
- If the iterations dominate the computation, then $S \approx r \bar{m} m^{-1}$.


## 4. Construction of nondefective EBDF methods

We shall construct nondefective versions of the original three-stage and four-stage EBDF-type methods given in [1] and [12].

### 4.1 Three-stage methods

We consider methods of the form (2.1) with $r=3$ and

$$
\begin{align*}
& c=\left(\begin{array}{c}
c_{1} \\
2 \\
1
\end{array}\right), B=\left(\begin{array}{ccc}
1 & 0 & 0 \\
B_{21} & 1 & 0 \\
0 & 0 & 1
\end{array}\right), C=\left(\begin{array}{ccc}
C_{11} & 0 & 0 \\
0 & C_{22} & 0 \\
C_{31} & C_{32} & C_{33}
\end{array}\right)  \tag{4.1}\\
& E=\left(\begin{array}{cccc}
E_{11} & E_{12} & \cdots & E_{1 s} \\
0 & E_{22} & \cdots & E_{2 s} \\
E_{31} & E_{32} & \cdots & E_{3 s}
\end{array}\right) .
\end{align*}
$$

Given the abscissa $c_{1}$ and one of the parameters $C_{3 j}$, the remaining entries in the arrays in (4.1) can be computed by means of the order conditions (2.3) such that $p_{1}=p_{2}=s$ and $p_{3}=s+1$. Hence, the order of accuracy (both stiff and nonstiff) is $p=s+1$. The cases $\left\{c_{1}=1, C_{31}=0\right\}$ and $\left\{c_{1}=1, C_{33}=C_{11}=C_{22}\right\}$ respectively define the original EBDF and MEBDF methods. For future reference, Table 4.1 lists for $p=3, \ldots, 6$ the MEBDF values of the angle of unconditional stability $\alpha$; the parameters $D_{1}$ and $D_{2}$ determining the rectangle $\left\{z:-D_{1} \leq \Re(z) \leq 0,-D_{2} \leq\right.$ $\left.\Im(z) \leq D_{2}\right\}$ containing the region of instability in the left half-plane; and the maximal modulus of the characteristic roots $\zeta$ in this region of instability. For larger values of $p$, the angle $\alpha$ quickly decreases, so that the resulting integration methods are less useful for solving general stiff problems.

As we already observed, the MEBDF methods of Table 4.1 are defective, so that direct diagonalization is not possible. Therefore, we used the two free parameters $c_{1}$ and $C_{31}$ to construct a

Table 4.1: Three-stage MEBDF methods of Cash with $c_{1}=1, C_{33}=C_{11}=C_{22}$.

| $p$ | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: |
| $\alpha$ | $90^{\circ}$ | $90^{\circ}$ | $88.4^{\circ}$ | $83.1^{\circ}$ |
| $\left(D_{1}, D_{2}\right)$ | $(0,0)$ | $(0,0)$ | $(0.040,1.8)$ | $(0.246,2.6)$ |
| $\|\zeta\|_{\max }$ | 1 | 1 | 1.029 | 1.121 |

nondefective, zero-stable and $\mathrm{L}(\alpha)$-stable EBDF method with (i) a relatively large $\alpha$ and (ii) a well-conditioned transformation matrix $Q$. Requiring that $Q$ be lower triangular with unit diagonal entries, we found by a straightforward numerical search the results listed in Table 4.2 (for the L-stable third- and fourth- order methods, the generating matrices $B^{-1} C, B^{-1} E, D$ and $Q$ needed in (3.4) are given in the Appendix to this paper). We mention only that there is a lot of freedom in choosing the parameters $c_{1}$ and $C_{31}$ to determine L-stable 3-stage methods satisfying (i) and (ii). For the 4-stage methods of the next section, the L-stable parameter space is much more restricted.

Table 4.2: Three-stage, nondefective EBDF methods of the form (4.1).

| $p$ | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: |
| $c_{1}$ | $5 / 4$ | $5 / 4$ | $5 / 4$ | $5 / 4$ |
| $C_{31}$ | 0 | 0 | $2 / 7$ | $3 / 13$ |
| $\\|Q\\|_{\infty}$ | 6.1 | 7.9 | 6.6 | 8.3 |
| $\alpha$ | $90^{\circ}$ | $90^{\circ}$ | $88.5^{\circ}$ | $83.9^{\circ}$ |
| $\left(D_{1}, D_{2}\right)$ | $(0,0)$ | $(0,0)$ | $(0.04,2.1)$ | $(0.24,3.9)$ |
| $\|\zeta\|_{\max }$ | 1 | 1 | 1.029 | 1.121 |

### 4.2 Higher-stage methods

The original EBDF and MEBDF methods have $c=(1,2,1)^{T}$, so that there is one 'future point' at $t_{n}+2 h$. This method can be interpreted as the successive application of the $s$-step BDF formula at $t_{n}+h$ and $t_{n}+2 h$ for predicting the future point value at $t_{n}+2 h$ needed in the $(s+1)$-step (M)EBDF corrector formula. More generally, we may introduce further future points by using $c=(1,2,3, \ldots, r-1,1)^{T}$. Considering only the stability of the corrector formula (last stage equation), we verified that up to order 18 the maximal order of L-stable formulas increases by 2 and the maximal order of $\mathrm{L}(\alpha)$-stable formulas increases by 3 with each additional future point. Of course, the use of BDF predictors will reduce the stability of the overall method, but we may still hope for improvement: Psihoyios and Cash [12] have shown that there exist L-stable 4 -stage methods of order 6 . However, just as in the case of the three-stage EBDF, choosing $c=(1,2,3, \ldots, r-1,1)^{T}$ yields defective matrices $B^{-1} C$. Therefore, we shall consider abscissae vectors of the form

$$
\begin{equation*}
c=\left(c_{1}, 2,3, \ldots, r-1,1\right)^{T} \tag{4.2a}
\end{equation*}
$$

where $c_{1}$ is a free parameter. According to the structure of the original (M)EBDF methods, we impose the following sparsity pattern on the matrices $B, C$ and $E$ :

$$
B:=\left(\begin{array}{ccccc}
1 & & & &  \tag{4.2b}\\
* & 1 & & & \\
\vdots & \vdots & \ddots & & \\
* & * & \cdots & 1 & \\
0 & 0 & \cdots & 0 & 1
\end{array}\right), C:=\left(\begin{array}{ccccc}
* & & & & \\
& * & & & \\
& & \ddots & & \\
& & & * & \\
* & * & \cdots & * & *
\end{array}\right), E:=\left(\begin{array}{cccccc}
* & * & \cdots & * & \cdots & * \\
& * & \cdots & * & \cdots & * \\
& & \ddots & & \vdots & \\
& & & * & \cdots & * \\
* & * & \cdots & * & \cdots & *
\end{array}\right),
$$

The entries in the matrices $B, C$ and $E$ can be determined such that the first $r-1$ stage equations in (4.2b) are consistent of order $s$. The last stage equation contains $r+s$ free parameters, so that
it can be made consistent of order $r+s-1$. Since the order of accuracy of (4.2a, 4.2 b$)$ cannot exceed $s+1$, we shall choose the entries in the corrector equation such that it is consistent of order $s+1$, leaving $r-2$ free parameters. Together with the free parameter $c_{1}$, we obtain an ( $r-1$ )-parameter family of EBDF-type methods with stage order $\bar{p}=s$ and order of accuracy $p=s+1$. From this family, we want nondefective, L-stable methods, again under the condition of zero-stability and a well-conditioned transformation matrix $Q$.
Let us consider the case of four stages $(r=4)$ with three free parameters. As already mentioned, Psihoyios and Cash have considered the defective case $c_{1}=1$ and shown that L-stable, sixthorder methods exist for a particular choice of the remaining two free parameters. For example, they verified that the parameters $C_{41}=1 / 10$ and $C_{43}=1 / 20$ generate an L-stable method with $p=s+1=6$. This motivated us to search for nondefective, L- and zero-stable methods by choosing $c_{1} \neq 1$. A numerical search produced for $p=s+1=5$ the values $c_{1}=3 / 2, C_{41}=3 / 10$, $C_{43}=7 / 50$ giving $\|Q\|_{\infty} \approx 31.5$ and for $p=s+1=6$ the values $c_{1}=6 / 5, C_{41}=11 / 100$, $C_{43}=1 / 20$, giving $\|Q\|_{\infty} \approx 167.5$. The corresponding generating matrices $B^{-1} C, B^{-1} E, D$ and $Q$ needed in (3.4) are given in the Appendix.
Together with the conventional BDF methods of order $p=1$ and $p=2$, and the three-stage nondefective EBDF methods of order $p=3$ and $p=4$ derived in the preceding section, we now have L-stable methods up to order six, all having a comparable effective computational complexity per step, provided that we employ three processors for $p=3,4$ and four processors for $p=5,6$.

## 5. Numerical experiments

Preliminary numerical experiments have been conducted using Matlab. We compare two methods from the three-parameter family of four-stage, 6th-order EBDF-type methods of the form (4.2a, 4.2 b ) with free parameters $c_{1}, C_{41}$ and $C_{43}$. The first method is due to Psihoyios and Cash and is defined by $c_{1}=1, C_{41}=0.10, C_{43}=0.05$. It is L-stable, but defective, so that sequential iteration has to be applied (see Section 2.4). The second method is defined by $c_{1}=1.2, C_{41}=0.11$, $C_{43}=0.05$. It also is L-stable, but nondefective, so that the parallel iteration method (3.4) can be applied. In the following, we call these methods the Defective and Nondefective EBDF methods, respectively. The values of the parameter $L_{K}$ in (3.5) are $L_{K} \approx 1.88$ for the Defective EBDF method and $L_{K} \approx 1.68$ for the Nondefective EBDF method, so we would expect the methods to have similar convergence behaviors. In addition to the these methods, we reproduced the results from [4] obtained for the original three-stage, 6th-order EBDF method of Cash when iterated by the diagonal iteration method (3.2) with $A^{*}=\operatorname{diag}\left(B^{-1} C\right)$, to be referred to as Diagonal EBDF. By mutual comparison of the three methods we can see what we have gained by the introduction of nondefective EBDF methods.
Following [4] the initial iterates for the iteration processes are obtained by taking the most recent approximation available or, if not yet available (in the case of the future value at $t_{n+r-1}$ and at $t_{n+c_{1}}$ ), by 6 -point extrapolation of already computed approximations. The Jacobian matrix $J$ is evaluated in each step using the future-point-approximation to $y_{n+1}$ from the preceding step. The starting values were obtained either from the exact solution (if available) or by applying the 5 th-order Radau IIA method with a 5 times smaller stepsize and using 10 Newton iterations per step.
Three of the test problems are the same as in [4], viz. the Kaps problem [10]

$$
\begin{align*}
& \frac{d y_{1}}{d t}=-1002 y_{1}+1000 y_{2}^{2} \\
& \frac{d y_{2}}{d t}=y_{1}-y_{2}\left(1+y_{2}\right)  \tag{5.1}\\
& y_{1}(0)=y_{2}(0)=1,0 \leq t \leq 5
\end{align*}
$$

the eight-dimensional 'High Irradiance RESponse' problem given in ([6], p. 157):
HIRES on the interval [5, 321.8122],
where the initial conditions at $t=5$ were obtained by applying the RADAU 5 code $[7]$ on $[0,5]$; and the non-autonomous Robertson problem

$$
\begin{array}{ll}
y_{1}^{\prime}=-0.04 y_{1}+10^{4} y_{2} y_{3}-0.96 \mathrm{e}^{-t}, & y_{1}(0)=1, \\
y_{2}^{\prime}=0.04 y_{1}-10^{4} y_{2} y_{3}-10^{7}\left(y_{2}\right)^{2}-0.04 \mathrm{e}^{-t}, & y_{2}(0)=0, \quad 0 \leq t \leq 1  \tag{5.3}\\
y_{3}^{\prime}=3 \times 10^{7}\left(y_{2}\right)^{2}+\mathrm{e}^{-t}, & y_{3}(0)=0
\end{array}
$$

The fourth test problem is the 15 -dimensional circuit analysis problem due to Horneber [8] and extensively discussed in [3] and [5]. In our implementation, we used the specification given in [11]:

$$
\begin{equation*}
\text { Ring modulator on the interval }\left[0,10^{-3}\right] \text { with } C_{s}=10^{-9} . \tag{5.4}
\end{equation*}
$$

In our numerical experiments, we denoted the number of steps by $N$, the number of iterations in each iteration process by $m$, and the total number of iterations by $M$ (not including the iterations needed to compute the starting values). Note that for fixed values of $m$ and $N$, Defective EBDF requires four times more sequential righthand side evaluations and forward-backward substitutions than the Diagonal and Nondefective EBDF methods, because Defective EBDF solves four subsystems per step. Hence, for Defective EBDF the value of $M$ is four times greater. The accuracy is given by the number of significant correct digits $s c d$; that is, we write the maximal absolute end point error in the form $10^{-s c d}$. In the tables of results, we shall indicate negative $s c d$-values by *.

### 5.1 Fixed numbers of iterations

The Tables 5.1, 5.2 and 5.3 list for given values of $m$ and $N$ the resulting $s c d$-values for the first three problems (5.1)-(5.3). These results show that the three methods converge to solutions with comparible accuracy. Furthermore, the convergence rate is for Diagonal EBDF slightly less than for the other two methods.

Table 5.1: Values of $s c d$ for the Kaps problem (5.1).

| $N$ | Method | $m=1$ | $m=2$ | $\ldots$ | $m=\infty$ |
| :---: | :--- | :---: | :---: | :--- | :---: |
| 0 | Defective EBDF | 5.0 |  | $\ldots$ | 5.0 |
|  | Nondefective EBDF | 5.2 |  | $\ldots$ | 5.2 |
|  | Diagonal EBDF | $*$ | 4.7 | $\ldots$ | 4.5 |
| 20 | Defective EBDF | 6.8 |  | $\ldots$ | 6.8 |
|  | Nondefective EBDF | 6.9 |  | $\ldots$ | 6.9 |
|  | Diagonal EBDF | $*$ | 6.4 | $\ldots$ | 6.3 |
| 40 | Defective EBDF | 8.5 |  | $\ldots$ | 8.5 |
|  | Nondefective EBDF | 8.8 |  | $\ldots$ | 8.8 |
|  | Diagonal EBDF | $*$ | 8.2 | $\ldots$ | 8.1 |

Table 5.2: Values of $s c d$ for the HIRES problem (5.2).

| $N$ | Method | $m=1$ | $m=2$ | $m=3$ | $m=4$ | $\ldots$ | $m=\infty$ |
| :---: | :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| 10 | Defective EBDF | 3.1 | 3.0 | 2.0 | 3.2 | $\ldots$ | 3.1 |
|  | Nondefective EBDF | 3.4 | 3.0 | 2.0 | 2.9 | $\ldots$ | 2.8 |
|  | Diagonal EBDF | $*$ | 2.8 | 2.5 | 2.7 | $\ldots$ | 2.7 |
| 20 | Defective EBDF | 2.6 | 3.7 | 3.9 | 3.9 | $\ldots$ | 3.8 |
|  | Transformed EBDF | 3.6 | 3.7 | 3.6 | 3.6 | $\ldots$ | 3.6 |
|  | Diagonal EBDF | $*$ | 3.6 | 3.4 | 3.3 | $\ldots$ | 3.3 |
| 40 | Defective EBDF | 5.0 | 4.8 | 4.9 |  | $\ldots$ | 4.9 |
|  | Nondefective EBDF | 4.4 | 4.7 | 4.8 |  | $\ldots$ | 4.8 |
|  | Diagonal EBDF | $*$ | 4.4 | 4.3 |  | $\ldots$ | 4.3 |

Table 5.3: Values of $s c d$ for the Robertson problem (5.3).

| $N$ | Method | $m=1$ | $m=2$ | $\ldots$ | $m=\infty$ |
| :---: | :--- | :---: | :---: | :---: | :---: |
| 10 | Defective EBDF | 7.6 |  | $\ldots$ | 7.6 |
|  | Nondefective EBDF | 7.7 |  | $\ldots$ | 7.7 |
|  | Diagonal EBDF | 7.8 | 7.9 | $\ldots$ | 7.9 |
| 20 | Defective EBDF | 9.3 |  | $\ldots$ | 9.3 |
|  | Nondefective EBDF | 9.3 |  | $\ldots$ | 9.3 |
|  | Diagonal EBDF | $*$ | 9.6 | $\ldots$ | 9.6 |
| 40 | Defective EBDF | 11.0 |  | $\ldots$ | 11.0 |
|  | Nondefective EBDF | 11.0 |  | $\ldots$ | 11.0 |
|  | Diagonal EBDF | $*$ | 11.3 | $\ldots$ | 11.3 |

### 5.2 Variable number of iterations

In our dynamic iteration strategy, we used the stopping strategy described in ([6], p. 130) using for the tolerance parameter Tol an estimate of the local truncation error LTE (see [4] for details on this modification). In the Defective EBDF and Diagonal EBDF cases, the difference $\mathbf{y}_{n-1+c_{1}}-\mathbf{y}_{n-1+c_{r}}$, available from the preceding step, provides us with a free estimate of LTE. In the Nondefective EBDF case, we used $\mathbf{y}_{n-2+c_{2}}-\mathbf{y}_{n-1+c_{r}}$ (i.e. the difference between the order $p-1$ initial guess and the order $p$ converged solution at time step $n$.) All further iteration strategy parameters are the same as in [4].

For the three most difficult problems (5.2), (5.3) and (5.4), we performed experiments in which the number of steps was chosen such that a prescribed $s c d$-value was obtained. For these problems, the maximal number of Newton iterations in the subsequent iteration processes was limited to 10. The Tables 5.4, 5.5 and 5.6 list the total number of iterations $M$ needed to obtain a given $s c d$-value. From these values, we may conclude that the two parallel methods Nondefective EBDF and Diagonal EBDF need about two to four times fewer iterations then the sequential method Defective EBDF.

Table 5.4: Values of $M$ for the HIRES problem (5.2).

| Method | $s c d=4$ | $s c d=5$ | $s c d=6$ | $s c d=7$ |
| :--- | :---: | :---: | :---: | :---: |
| Defective EBDF | 96 | 168 | 573 | 928 |
| Nondefective EBDF | 73 | 102 | 195 | 343 |
| Diagonal EBDF | 83 | 133 | 189 | 241 |

Table 5.5: Values of $M$ for the Robertson problem (5.3).

| Method | $s c d=8$ | $s c d=9$ | $s c d=10$ | $s c d=11$ | $s c d=12$ | $s c d=13$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| Defective EBDF | 36 | 63 | 107 | 169 | 265 | 415 |
| Nondefective EBDF | 10 | 18 | 29 | 47 | 74 | 123 |
| Diagonal EBDF | 9 | 17 | 29 | 49 | 74 | 114 |

Table 5.6: Values of $M$ for the Ring modulator (5.4).

| Method | $s c d=6$ | $s c d=7$ | $s c d=8$ |
| :--- | :---: | :---: | :---: |
| Defective EBDF | 49900 | 72400 | 104500 |
| Nondefective EBDF | 14800 | 22300 | 33800 |
| Diagonal EBDF | 20000 | 29700 | 42800 |

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## Appendix A Coefficients of some nondefective EBDF methods

For reference we provide the coefficient matrices of the L-stable nondefective EBDF methods considered in this paper. For each method we give the matrices $B^{-1} C, B^{-1} E$ and $Q$ needed for parallel implementation of (3.4). Obviously, the diagonal matrix $D$ needed for the implementation of 3.4 is given by $D=\operatorname{diag}\left(B^{-1} C\right)$. The coefficients listed are exact, expressed in fractional form, and were determined by Maple.

The 3 -stage L-stable method of order $p=3$ is defined by $c_{1}=5 / 4, C_{31}=0$. The method coefficients and transformation matrix are given by

$$
\begin{align*}
& B^{-1} C=\left(\begin{array}{ccc}
\frac{45}{56} & 0 & 0 \\
\frac{72}{77} & \frac{6}{11} & 0 \\
0 & -\frac{4}{23} & \frac{22}{23}
\end{array}\right), B^{-1} E=\left(\begin{array}{cc}
-\frac{25}{56} & \frac{81}{56} \\
-\frac{40}{77} & \frac{117}{77} \\
-\frac{5}{23} & \frac{28}{23}
\end{array}\right),  \tag{A.1}\\
& Q=\left(\begin{array}{ccc}
1 & 0 & 0 \\
\frac{192}{53} & 1 & 0 \\
\frac{43008}{10441} & \frac{11}{26} & 1
\end{array}\right) .
\end{align*}
$$

A 3 -stage L-stable method of order $p=4$ is defined by $c_{1}=5 / 4, C_{31}=0$. The method coefficients and transformation matrix are given by

$$
\begin{align*}
& B^{-1} C=\left(\begin{array}{ccc}
\frac{585}{908} & 0 & 0 \\
\frac{192}{227} & \frac{6}{13} & 0 \\
0 & -\frac{18}{197} & \frac{150}{197}
\end{array}\right), B^{-1} E=\left(\begin{array}{ccc}
\frac{2025}{7264} & -\frac{4225}{3632} & \frac{13689}{7264} \\
\frac{1080}{2951} & -\frac{4204}{2951} & \frac{6075}{2951} \\
\frac{17}{197} & -\frac{99}{197} & \frac{279}{197}
\end{array}\right),  \tag{A.2}\\
& Q=\left(\begin{array}{ccc}
1 & 0 & 0 \\
\frac{3328}{719} & 1 & 0 \\
\frac{18130944}{5022215} & \frac{39}{128} & 1
\end{array}\right) .
\end{align*}
$$

A 4-stage L-stable method of order $p=5$ is defined by $c_{1}=3 / 2, C_{41}=3 / 10, C_{43}=7 / 50$. The method coefficients and transformation matrix are given by

$$
B^{-1} C=\left(\begin{array}{cccc}
\frac{315}{496} & 0 & 0 & 0 \\
\frac{864}{1147} & \frac{12}{37} & 0 & 0 \\
\frac{2768}{3441} & \frac{32}{37} & \frac{4}{9} & 0 \\
\frac{3}{10} & -\frac{3059487}{4001600} & \frac{7}{50} & \frac{5279163}{4001600}
\end{array}\right)
$$

$$
\begin{align*}
& B^{-1} E=\left(\begin{array}{cccc}
-\frac{1225}{3968} & \frac{6075}{3968} & -\frac{11907}{3968} & \frac{11025}{3968} \\
-\frac{420}{1147} & \frac{2043}{1147} & -\frac{3884}{1147} & \frac{3408}{1147} \\
-\frac{12110}{30969} & \frac{2118}{1147} & -\frac{3907}{1147} & \frac{91382}{30969} \\
\frac{2153579}{24009600} & -\frac{3413921}{8003200} & \frac{4631823}{8003200} & \frac{3640463}{4801920}
\end{array}\right),  \tag{A.3}\\
& Q=\left(\begin{array}{cccc}
1 & 0 & 0 & 0 \\
\frac{4608}{1901} & 1 & 0 & 0 \\
\frac{24616704}{1617751} & -\frac{36}{5} & 1 & 0 \\
-\frac{38599642812960}{45767552496101} & \frac{145802607}{81838795} & -\frac{5042016}{31506067} & 1
\end{array}\right) .
\end{align*}
$$

A 4-stage L-stable method of order $p=6$ is defined by $c_{1}=6 / 5, C_{41}=11 / 100, C_{43}=1 / 20$. The method coefficients and transformation matrix are given by

$$
\begin{aligned}
& B^{-1} C=\left(\begin{array}{ccccc}
\frac{16016}{32525} & 0 & 0 & 0 \\
\frac{40625}{49438} & \frac{15}{38} & 0 & 0 \\
\frac{39040625}{41626796} & \frac{30375}{31996} & \frac{180}{421} & 0 \\
\frac{11}{100} & -\frac{120153318}{388515625} & \frac{1}{20} & \frac{1497086157}{1554062500}
\end{array}\right) \\
& B^{-1} E=\left(\begin{array}{ccccc}
\frac{569184}{4065625} & -\frac{10469888}{12196875} & \frac{9018009}{4065625} & -\frac{12719616}{4065625} & \frac{32064032}{12196875} \\
\frac{5775}{24719} & -\frac{101768}{74157} & \frac{82350}{24719} & -\frac{105400}{24719} & \frac{227750}{74157} \\
\frac{5549775}{20813398} & -\frac{46526500}{31220097} & \frac{70906923}{20813398} & -\frac{42611025}{10406699} & \frac{90894625}{31220097} \\
-\frac{211339877}{6216250000} & \frac{939457771}{4662187500} & -\frac{168763034}{388515625} & \frac{333046763}{1554062500} & \frac{19629003023}{18648750000}
\end{array}\right),
\end{aligned}
$$

$$
Q=\left(\begin{array}{cccc}
1 & 0 & 0 & 0  \tag{A.4}\\
\frac{1015625}{120733} & 1 & 0 & 0 \\
\frac{7376452890625}{53619698494} & -\frac{405}{14} & 1 & 0 \\
-\frac{475587595010650768146875}{51052091899348840572958} & \frac{241922892409}{78349451754} & -\frac{32713015625}{350542022097} & 1
\end{array}\right) .
$$

