

AFDELING NUMERIEKE WISKUNDE
NW 151/83
APRIL
(DEPARTMENT OF NUMERICAL MATHEMATICS)
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LINEAR MULTISTEP METHODS FOR VOLTERRA INTEGRAL AND INTEGRO-DIFFERENTIAL EQUATIONS

Preprint


Printed at the Mathematical Centre, Kruislaan 413, Amsterdam, The Netherlands.
The Mathematical Centre, founded 11 February 1946, is a non-profit institution for the promotion of pure and applied mathematics and computer science. It is sponsored by the Netherlands Government through the Netherlands Organization for the Advancement of Pure Research (Z.W.O.).

1980 Mathematics subject classification: 65R20

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Linear multistep methods for Volterra integral and integro-differential equations *)
by

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## ABS TRACT

A general class of linear multistep methods is presented for numerically solving first and second kind Volterra integral equations, and Volterra integro-differential equations. These so-called VLM methods, which include the well-known direct quadrature methods, allow for a unified treatment of the problems of consistency and convergence, and have a pendant in linear multistep methods for ODEs, as treated in any textbook on computational methods in ordinary differential equations.

General consistency and convergence results are presented (and proved in an Appendix), together with results of numerical experiments which support the theory.

KEY WORDS \& PHRASES: numerical analysis; Volterra integral and integro-differential equations; linear multistep methods; consistency; convergence

[^0]
## 1. INTRODUCTION

We consider Volterra integral equations of the form

$$
\begin{equation*}
\theta y(t)=g(t)+\int_{t_{0}}^{t} K(t, \tau, y(\tau)) d \tau, t \in I:=\left[t_{0}, T\right], \theta=0,1 \tag{1.1}
\end{equation*}
$$

This equation is called of the first kind if $\theta=0$ and of the second kind if $\theta=1$. Furthermore, we consider VoZterra integro-differential equations

$$
\begin{equation*}
\frac{d y}{d t}=f(t, y(t), z(t)), z(t)=g(t)+\int_{t_{0}}^{t} K(t, \tau, y(\tau)) d \tau, t \in I \tag{1.2}
\end{equation*}
$$

where $y\left(t_{0}\right)=y_{0}$. In these equations $y(t)$ is the unknown function and $g$, $K$ and $f$ are given, nonsingular functions on $I, S \times \mathbb{R}$ and $I \times \mathbb{R} \times \mathbb{R}$, respectively, where $S:=\left\{(t, \tau), t_{0} \leq \tau \leq t \leq T\right\}$. In order to ensure the existence of a unique, continuous solution of (1.1) and (1.2), the following conditions should be satisfied, respectively:

Conditions for the existence of a unique solution $y(t) \in C(I)$ of (1.1) with $\theta=1$
$-K(t, \tau, y)$ is continuous with respect to $t$ and $\tau$, for all ( $t, \tau) \in S$;

- K satisfies a (uniform) Lipschitz condition with respect to y, i.e., $|K(t, \tau, y)-K(t, \tau, z)| \leq L_{1}|y-z|$, for all $(t, \tau) \in S$, for all finite $y, z \in \mathbb{R} ;$
$-g(t) \in C(I)$.
Conditions for the existence of a unique solution $y(t) \in C(I)$ of (1.1) with $\theta=0$
$-K(t, \tau, y) \in C^{1}(S \times \mathbb{R}) ;$
- for $t=\tau$ the derivative $\partial K / \partial y$ is bounded away from zero: $|\partial K(t, t, y) / \partial y| \geq r_{0}>0$ for all $t \in I, y \in \mathbb{R}$;
- $K(t, \tau, y)$ satisfies a (uniform) Lipschitz condition with respect to $y$ on $S \times \mathbb{R}$;
$-g(t) \in C^{1}(I)$ with $g\left(t_{0}\right)=0$.
Conditions for the existence of a unique solution $y(t) \in C^{1}(I)$ of (1.2), for given initial value $\mathrm{y}\left(\mathrm{t}_{0}\right)=\mathrm{y}_{0}$

The following three (uniform) Lipschitz conditions:

$$
\begin{aligned}
& -\left|f\left(t, y_{1}, z\right)-f\left(t, y_{2}, z\right)\right| \leq L_{1}\left|y_{1}-y_{2}\right| \text {, for all } t \in I \text {, for all finite } \\
& \quad z, y_{1}, y_{2} \in \mathbb{R} ; \\
& -\left|f\left(t, y, z_{1}\right)-f\left(t, y, z_{2}\right)\right| \leq L_{2}\left|z_{1}-z_{2}\right| \text {, for all } t \in I \text {, for all finite }
\end{aligned}
$$ $y, z_{1}, z_{2} \in \mathbb{R}$;

$-\left|K\left(t, \tau, y_{1}\right)-K\left(t, \tau, y_{2}\right)\right| \leq L_{3}\left|y_{1}-y_{2}\right|$, for all ( $\left.t, \tau\right) \in S$, for all finite $y_{1}, y_{2} \in \mathbb{R}$.
A common, simple way of solving (1.1) numerically is obtained by writing these equations down in a sequence of equidistant points

$$
\begin{equation*}
\mathrm{t}_{\mathrm{n}}:=\mathrm{t}_{0}+\mathrm{nh}, \quad \mathrm{n}=0(1) \mathrm{N}, \mathrm{~h} \text { fixed and } \bar{E}_{\mathrm{N}}=\mathrm{T}, \tag{1.3}
\end{equation*}
$$

approximating the integral term by some suitably chosen quadrature formula, and solving the resulting equation for $y\left(t_{n}\right)$, successively for $n=n_{0}(1) N$, where $n_{0}$ is some suitable starting index. Equation (1.2) is commonly solved by integrating the differential equation in the points (1.3) (say), using an LM formula for ODEs, thereby approximating $z\left(t_{j}\right)$ with some suitably chosen quadrature formula. All these methods are called linear multistep (LM) direct quadrature (DQ) methods. DQ methods may give satisfactory results, but sometimes the results with $D Q$ are completely worthless as was demonstrated for first kind equations by $\operatorname{LINZ}$ [9, p. 67], where he applied a fourth order Gregory quadrature method to the very simple integral equation

$$
\begin{array}{r}
0=-\operatorname{sint}+\int_{0}^{t} \cos (t-\tau) y(\tau) d \tau, I=[0,2], \text { with exact solution }  \tag{1.4}\\
y(t)=1
\end{array}
$$

The "approximate" values obtained for $y(2)$ were 8.4 and $1.5 \times 10^{7}$ for $h=0.1$ and $h=0.05$, respectively. For second kind equations too the Gregory rules will fail if large Lipschitz constants for the kernel function with respect to $y$ are involved.

In this paper we present a general class of linear multistep methods for (1.1) and (1.2) which includes the DQ methods. (It should be remarked that such methods for second kind Volterra integral equations were already introduced in [6] and results were presented without proof). A characteristic
feature of this class is that it involves linear combinations, not only of $y$ - and $K$ - values, but also of values of the auxiliary function (called the Zag term)

$$
\begin{equation*}
Y(t, s):=g(t)+\int_{t_{0}}^{s} K(t, \tau, y(\tau)) d \tau \tag{1.5}
\end{equation*}
$$

for $(t, s) \in S$. Note that we may write (1.1) as $\theta y(t)=Y(t, t)$.
This general class will be called Volterra Zinear multistep (VLM) methods. VLM methods allow for a uniform treatment of the problems of consistency and convergence, and have a pendant in linear multistep methods for ordinary differential equations, as treated, e.g., by LAMBERT in [8].

In Section 2 of this paper we treat VLM methods for Volterra integral equations of the second and of the first kind jointly. Numerical experiments with several examples of VLM methods are reported. In a similar way as is done in Section 2, Section 3 treats VLM methods for integro-differential equations. It turns out that several results of Section 2 for second kind Volterra integral equations can be used in Section 3. The proofs of the theorems presented in Sections 2 and 3 are given, as far as they are nontrival, in an Appendix to this paper.
2. VLM METHODS FOR VOLTERRA INTEGRAL EQUATIONS

### 2.1. The general VLM method

In order to state our general VLM method for (1.1) we introduce numerical approximations $y_{n}$ to $y\left(t_{n}\right)$ and $Y_{n}(t)$ to $Y\left(t, t_{n}\right)$, and we let

$$
\begin{equation*}
K_{n}(t):=K\left(t, t_{n}, y_{n}\right), \quad n \geq 0 \tag{2.1.1}
\end{equation*}
$$

We assume that $Y_{n}(t), t \geq t_{n}$, will be computed by a quadrature formula of the form

$$
\begin{equation*}
Y_{n}(t):=g(t)+h \sum_{j=0}^{n} w_{n j} K_{j}(t), \quad n \geq n_{1} \tag{2.1.2}
\end{equation*}
$$

where the $w_{n j}$ are given weights and $n_{1}$ is sufficiently large in order to ensure the required order of accuracy. When we say that the order of this quadrature formula is $r$, we mean that for any $t \geq t_{n}$

$$
\begin{equation*}
E_{n}(h ; t):=\int_{t_{0}}^{t_{n}} K(t, \tau, y(t)) d \tau-h \sum_{j=0}^{n} w_{n j} K\left(t, t_{j}, y\left(t_{j}\right)\right)=O\left(h^{r}\right) \tag{2.1.3}
\end{equation*}
$$

as $h \rightarrow 0, n \rightarrow \infty$, with $t_{n}=t_{0}+n h$ fixed. An important class of quadrature formulas, which includes the well-known Gregory formulas, are the so-called $(\rho, \sigma)$ - reducible quadrature formulas [14].

Our general VLM method for (1.1) consists of
(i) the VLM formuZa
(2.1.4) $\theta \sum_{i=0}^{k} \alpha_{i} y_{n-i}+\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j} Y_{n-i}\left(t_{n+j}\right)=h \sum_{i=0}^{k} \sum_{j=-k}^{k} \gamma_{i j} K_{n-i}\left(t_{n+j}\right)$, $\mathrm{n}=\mathrm{k}^{*}(1) \mathrm{N}$,
( $k^{*}$ fixed) where $\alpha_{i}, \beta_{i j}$ and $\gamma_{i j}, i=0(1) k, j=-k(1) k$, are to be prescribed, and
(ii) the quadrature formula (2.1.2) for the computation of $Y_{n-i}\left(t_{n+j}\right)$. In the VLM method the quantities $y_{1}, \ldots, y_{k^{*}-1}$ with $k^{*}=k+n_{1}$ are assumed to be precomputed by some starting method. Then $y_{k^{*}}, \ldots, y_{N}$ can be succesively computed using (2.1.4). Since the kernel $K(t, \tau, y)$ is not necessarily defined outside $S$, we require $\beta_{i j}=\gamma_{i j}=0$ for $j<-i$. Furthermore, if $\beta_{i j}, \gamma_{i j} \neq 0$ for $j=1(1) k$ we assume that the domain of definition of $K$ can be extended to points ( $t, \tau$ ) with $t \leq T+k h, \tau \leq T$. It is convenient to characterize (2.1.4) by the matrices

$$
\begin{equation*}
A=\left(\alpha_{i}\right), \quad B=\left(\beta_{i j}\right), \quad C=\left(\gamma_{i j}\right) \tag{2.1.5}
\end{equation*}
$$

where the row index $i$ assumes the values $0(1) k$ and the column index $j$ the values $-k(1) k$. (Note that for $\theta=0$ the values of the coefficients $\alpha_{i}$ in (2.1.4) are irrelevant.)

We now describe four subclasses of (2.1.4) used as illustrating examples in this paper.

Subclass 1 Direct quadrature methods
Direct quadrature methods for (1.1) are characterized by the ( $1 \times 1$ ) matrices $A=I, B=-I, C=0$, for which (2.1.4) reduces to the simple scheme $\theta y_{n}=Y_{n}\left(t_{n}\right), n=n_{1}(1) N$.

Subclass 2 Indirect Iinear multistep methods
In [6] methods for (1.1) (with $\theta=1$ ) were considered in which the VLM formula (2.1.4) is generated by the matrices
(2.1.6)

$$
\begin{aligned}
& A=\left[\begin{array}{c}
a_{0} \\
a_{1} \\
\cdot \\
\cdot \\
\cdot \\
a_{k}
\end{array}\right], B=\left[\begin{array}{cccc}
0 & b_{0} \delta_{0} & b_{0} \delta_{1} & \cdots \\
b_{1} \delta_{0} & b_{1} \delta_{1} & \cdots & b_{0} \delta_{k} \\
\cdot & \cdot & & b_{1} \delta_{k} \\
\cdot & \cdot & & \cdot \\
\cdot & \cdot & & \cdot \\
b_{k} \delta_{0} & b_{k} \delta_{1} & \cdots, & b_{k} \delta_{k}
\end{array}\right] \text {, } \\
& C=\left[\begin{array}{cc:c}
\bigcirc & & b_{0}
\end{array}\right]
\end{aligned}
$$

where the $a_{i}$ and $b_{i}$, $i=0(1) k$, are the coefficients of some LM method for ODEs and the $\delta_{i}$, $i=0(1) k$, are the coefficients of ( $k+1$ )-point forward differentiation formula (Table 1 of the Appendix lists these for $k=1(1) 5)$. This VLM formula forms, together with (2.1.2), a so-called indirect linear multistep (ILM) method for (1.1), not only for $\theta=1$, but also for $\theta=0$. When the $a_{i}$ and $b_{i}$ are the coefficients of $a$ backward differentiation method (for $k=1$ (1) 5 these are listed in Table 2 of the Appendix), (2.1.4) represents the so-called IBD (indirect backward differentiation) method, analysed in [5] (in the case $\theta=1$ ). For this method, we have $b_{0}=1, b_{<0}=0$ and $b_{0} \delta_{j}=a_{j}, j=0(1) k$.

It should be remarked that the ILM methods require the extension of the domain of definition $S$ with the points $\left\{(t, \tau) \mid T<t \leq T+k h, t_{0} \leq \tau \leq T\right\}$.

In this connection we observe that if $S$ can also be extended to points with $t<\tau$ we may use backward instead of forward differentiation coefficients $\delta_{i}$ in the IBD method, i.e., the matrix $B$ is replaced by the matrix ( $\beta_{i j}$ ) all elements of which vanish, except for those in the first row, which are given by $\left(b_{0} \delta_{k}, b_{0} \delta_{k-1}, \ldots, b_{0} \delta_{0}, 0, \ldots, 0\right)$.

## Subclass 3 Multilag methods

In $[16,17]$ we find methods for (1.1) with $\theta=1$ which can be characterized by the matrices
(2.1.7)

$$
A=\left[\begin{array}{c}
a_{0} \\
0 \\
\cdot \\
\cdot \\
\cdot \\
0
\end{array}\right], \quad B=\left[\begin{array}{c}
0 \\
a_{1} \\
\cdot \\
. \\
\cdot \\
a_{k}
\end{array}\right], \quad C=\left[\begin{array}{c}
b_{0} \\
-b_{1} \\
\cdot \\
. \\
\cdot \\
b_{k}
\end{array}\right]
$$

These methods were called multilag methods (ML) for (1.1) with $\theta=1$. Here, the $a_{i}$ and $b_{i}$, $i=0(1) k$ may be the coefficients of any LM method for ODEs. Wolkenfelt has pointed out that in the case that the lag term $Y_{n}(t)$ is computed by using a quadrature rule which is reducible to an LM-method for ODEs with the same coefficients $a_{i}$ and $b_{i}$, then the resulting method is, in fact, equivalent to a $D Q$ method based on the same quadrature rule (provided, of course, that identical starting values are used).

## Subclass 4 Modified multilag methods

In [16] Wolkenfelt introduced a modification of the ML methods, viz., the so-called modified multilag (MML) methods for (1.1) characterized by the matrices
(2.1.8)

$$
A=\left[\begin{array}{c}
a_{0} \\
a_{1} \\
\cdot \\
\cdot \\
\cdot \\
a_{k}
\end{array}\right], \quad B=\left[\begin{array}{cc} 
& 0 \\
\bigcirc & -a_{1} \\
a_{1} \\
\cdot & \cdot \\
\cdot & \cdot \\
\cdot & \cdot \\
-a_{k} & a_{k}
\end{array}\right], \quad C=\left[\begin{array}{c}
b_{0} \\
b_{1} \\
\cdot \\
\cdot \\
\cdot \\
b_{k}
\end{array}\right]
$$

The $a_{i}$ and $b_{i}$ are, again, the coefficients of any LM method for ODEs. A common choice are the Adams-MouZton formulas (listed in Table 3 in the Appendix, for $k=1(1) 5)$. As for $M \mathbb{L}$ methods, the MML method is algebraically equivalent with the $D Q$ method if in both methods the lag term formula is based on the LM formula $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$ and if the starting values would be identical, where $y_{n}=Y_{n}\left(t_{n}\right)$ for $n_{1} \leq n \leq n_{1}+k$.

### 2.2. Consistency of VLM formulas for Volterra integral equations

Let $C_{2}^{1}(S)$ denote the space of continuous functions $Y(t, s)$, differentiable with respect to $s$ for $\operatorname{all}(t, s) \in S$. Along the lines of the theory developed for LM formulas for ODEs, we associate with the VLM formula (2.1.4) the linear difference - differential operator ${\underset{\sim}{n}}_{n}: C_{2}^{1}(S) \rightarrow \mathbb{R}$, defined by

$$
\begin{equation*}
{\underset{\sim}{n}}_{n}[Y]:=\sum_{i=0}^{k}\left\{\theta \alpha_{i} Y\left(t_{n-i}, t_{n-i}\right)+\sum_{j=-k}^{k}\left[\beta_{i j}-\gamma_{i j} h \frac{\partial}{\partial s}\right] Y\left(t_{n+j}, t_{n-i}\right)\right\} \tag{2.2.1}
\end{equation*}
$$

As is usual in the consistency analysis of numerical schemes for functional equations, we now substitute the exact solution $y(t)$ of (1.1) into (2.1.4), and analyze the resulting residue. With the relations $\theta y(t)=Y(t, t)$ and $\partial Y(t, s) / \partial s=K(t, s, y(s))$, and (2.1.3) we obtain the equation

$$
\begin{align*}
& \sum_{i=0}^{k}\left\{\left(\theta_{i} y\left(t_{n-i}\right)+\sum_{j=-k}^{k}\left(\beta _ { i j } \left[g\left(t_{n+j}\right)+h \sum_{\ell=0}^{n-i} w_{n-i}, \ell^{K\left(t_{n+j}, t_{\ell}, y\left(t_{\ell}\right)\right]}\right.\right.\right.\right.  \tag{2.2.2}\\
& \left.\left.-h \gamma_{i j} K\left(t_{n+j}, t_{n-i}, y\left(t_{n-i}\right)\right)\right)\right\}={\underset{\sim}{n}}_{n}[Y]-\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j} E_{n-i}\left(h ; t_{n+j}\right) .
\end{align*}
$$

The second term in the residual originates from the quadrature formula $\{(2.1 .2)-(2.1 .3)\}$, and is $O\left(h^{r}\right)$ if $E_{n}(h ; t)$ is $O\left(h^{r}\right)$. The first term originates from the VLM formula (2.1.4) and will be called the local truncation error of the VLM formula. In order to analyze it, we use the following

Definition 2.2.1. The operator (2.2.1) and the associated VLM formula (2.1.4) are said to be consistent of order $p$ with equation (1.1) if for all $\mathrm{X} \in \mathrm{C}^{\mathrm{p}+1}(\mathrm{~S})$, which in the case of first kind equations $(\theta=0)$ vanish on the Zine $t=\tau$, we have ${\underset{\sim}{n}}^{L}[X]=O\left(h^{p+1}\right)$ as $h \rightarrow 0$ with nonvanishing error constant.

The following two theorems express the p-th order consistency conditions in terms of the parameters occurring in (2.1.4) for (1.1) in the cases $\theta=1$ and $\theta=0$, respectively.

THEOREM 2.2.1. The operator ${\underset{\sim}{n}}^{n}$ and the associated VLM formula (2.1.4) are consistent of order p with $\{(1.1), \theta=1\}$ if $\mathrm{C}_{\mathrm{q} \ell}=0$ for $\mathrm{q}=0(1) \mathrm{p}, \ell=0(1) \mathrm{q}$, where

$$
\begin{equation*}
C_{q \ell}:=\frac{1}{(q-\ell)!\ell!} \sum_{i=0}^{k}\left[(-i)^{q^{\prime}} \alpha_{i}-\sum_{j=-k}^{k} j^{q-\ell}(-i)^{\ell-1}\left(i \beta_{i j}+\ell \gamma_{i j}\right)\right], \tag{2.2.3}
\end{equation*}
$$

with the convention that $(-i)^{\ell-1} \ell=0$ if $i=\ell=0$.

THEOREM 2.2.2. The operator ${\underset{\sim}{n}}$ and the associated VLM formula (2.1.4) are consistent of order $p$ with $\{(1.1), \theta=0\}$ if $B_{q l}=0$ for $q=1(1) p, l=1(1) q$, where

$$
\begin{array}{r}
B_{q \ell}:=\frac{1}{(q-\ell)!\ell!} \sum_{i=0}^{k} \sum_{j=-k}^{k}(j-i)^{q-\ell-1}(j+i)^{\ell-1}\left[\beta_{i j}\left(j^{2}-i^{2}\right)-\right.  \tag{2.2.4}\\
\left.-\gamma_{i j}(q j+q i-2 \ell j)\right] .
\end{array}
$$

The various orders of consistency of the illustrating subclasses introduced in Section 2.1 can now be found by substituting the relevant values of the parameters $\alpha_{i}, \beta_{i j}$ and $\gamma_{i j}$ into the above two theorems. The results are summarized in the following corollary.

Corollary 2.2.1. Let $\tilde{p}$ be the order of consistency of the LM formula for ODEs defining the coefficients $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$ in Subclasses 2, 3 and 4. Then the order of consistency $p$ of the operator ${\underset{\sim}{n}}^{n}$ and of the associated VLM formula (2.1.4) with equation (1.1) is given by

$$
p=\left\{\begin{array}{l}
\infty \text { for Subclass } 1 \text { (DQ); } \\
\min \{k, \tilde{p} / \theta\} \text { for Subclass } 2 \text { (ILM); } \\
\tilde{p} \text { both for Subclasses } 3 \text { and } 4 \text { (ML resp. MML). }
\end{array}\right.
$$

Note that for $\theta=0$ in the ILM case this corollary gives $p=k$, independent of the order of consistency $\tilde{p}$, and, in fact, in this case the $\left\{a_{i}, b_{i}\right\}$ need
not represent an LM method for ODEs at all.
Let us assume in the rest of this section that the VLM formula (2.1.4) is consistent of order $p$ with equation (1.1). From the proofs of Theorems 2.2 .1 resp. 2.2 .2 it follows that the local truncation errror ${\underset{\sim}{n}}_{\mathrm{L}}^{\mathrm{L}}[\mathrm{Y}]$ can be expressed in terms of the constants defined in (2.2.3) resp. (2.2.4) as follows:

$$
\begin{equation*}
\mathrm{L}_{\mathrm{n}}[\mathrm{Y}]=\mathrm{h}^{\mathrm{p}+1} \sum_{\ell=0}^{\mathrm{p}+1} C_{\mathrm{p}+1}, \ell^{Y^{(\mathrm{p}+1-\ell, \ell)}}+O\left(\mathrm{~h}^{\mathrm{p}+2}\right), \text { as } \mathrm{h} \rightarrow 0 \text {, } \tag{2.2.5}
\end{equation*}
$$

resp.
where

$$
\begin{aligned}
& Y^{(i, j)}:=\left.(\partial / \partial t)^{i}(\partial / \partial s)^{j} y_{Y}(t, s)\right|_{t=s=t_{n}}, \\
& Z^{(i, j)}:=\left.(\partial / \partial u)^{i}(\partial / \partial v)^{j} Z_{(u, v)}\right|_{u=2 t_{n}, v=0}, \\
& Z(u, v):=Y\left(\frac{u+v}{2}, \frac{u-v}{2}\right) .
\end{aligned}
$$

In order to compare the values of the error constants $C_{p+1, \ell}$ and $B_{p+1, \ell}$ for the various subclasses introduced in Section 2.1 we have evaluated and simplified the expressions for these constants as much as possible, and obtained the following results.

For $\theta=1$ Corollary 2.2.1 gives for the ILM formula: $p=k$ provided that $\tilde{p} \geq k$ (which is a reasonable assumption, true, e.g., when the $L M$ formula for ODEs is a Backward Differentiation formula ( $\tilde{p}=k$ ) or an Adams-Moulton formula ( $\tilde{p}=k+1)$ ). Hence,

$$
c_{p+1, \ell}=\frac{(-1)^{p-1}}{(p+1-\ell)!\ell!} \sum_{i=0}^{k}\left\{i^{\left.\left.p_{\left[i a_{i}\right.}+(p+1) b_{i}\right]-R\right\}, p=k, ~}\right.
$$

with $R=k!b_{i}$ if $\ell=0$ and $R=0$ if $\ell=1,2, \ldots, p^{+1}$. For the (M)ML formula, Corollary 2.2 .1 gives $p=\tilde{p}$ and we found

$$
c_{p+1, \ell}=\left\{\begin{array}{ll}
0 & \text { if } \ell \leq p, \\
\frac{(-1)^{p-1}}{(p+1-\ell)!\ell!} & \sum_{i=0}^{k} i_{\left[i a_{i}+(p+1) b_{i}\right]} \text { if } \ell=p+1,
\end{array} \quad \text { = } \tilde{p} .\right.
$$

In Table 2.2.1 the numerical values of the constants $C_{p+1, \ell}$ are explicitly given for two popular choices of the coefficients $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$, viz., the $B D$ formulas and the $A M$ formulas.

Table 2.2.1. Error constants $C_{p+1, \ell}=C_{p+1, \ell}^{*} \frac{1}{\ell!(p+1-\ell)!}, \ell=0(1) p+1$

| VLM[LM] | p |  | $\mathrm{k}=1$ | $\mathrm{k}=2$ | $\mathrm{k}=3$ | $\mathrm{k}=4$ | $\mathrm{k}=5$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ILM[BD] | k | [ $\mathrm{C}^{*}$ |  | 0 | $-72$ | 0 | $-\frac{14400}{137}$ |
|  |  | $\left\{\begin{array}{l}\text { c } \\ \text { +1, }\end{array}\right.$ |  | 0 | $-\frac{11}{11}$ | 0 | 137 |
|  |  | $\left\{c_{p+1, ~}^{*}>0\right.$ | $=-1$ | $-\frac{4}{3}$ | $-\frac{36}{11}$ | $-\frac{288}{25}$ | $-\frac{7200}{137}$ |
| . ILM[AM] | k | $\left\{C^{*}{ }^{*}+1,0\right.$ | $=-1$ | 2 | -6 | 24 | -120 |
|  |  | $\left\{c_{p+1, ~}^{*}>0\right.$ | = 0 | 0 | 0 | 0 | 0 |
| (M)ML[BD] | k | ${ }_{\mathrm{p}+1}^{*}, \ell<\mathrm{p}+1$ | $=0$ | 0 | 0 | 0 | 0 |
|  |  |  | $=-1$ | $-\frac{4}{3}$ | $-\frac{36}{11}$ | $-\frac{288}{25}$ | $-\frac{7200}{137}$ |
| (M)ML [AM] | k+1 | $\int c_{p+1, \ell<p+1}^{*}$ | $=0$ | 0 | 0 | 0 | 0 |
|  |  | $\left\{\begin{array}{l}c^{*} \\ p^{*}+1, p+1\end{array}\right.$ | $=-\frac{1}{2}$ | -1 | $-\frac{19}{6}$ | $-\frac{27}{2}$ | $-\frac{863}{12}$ |

For $\theta=0$ and in the case of the ILM formula, Corollary 2.2.1 gives $p=k$, while the coefficients $\left\{b_{i}\right\}$ can still be chosen freely. For the error constants we found

$$
B_{p+1, \ell}=\frac{(-1)^{k} k!}{(p+1-\ell)!\ell!} \sum_{i=0}^{k} b_{i}, \quad 1 \leq \ell \leq p+1, \quad p=k .
$$

For the MML formula, Corollary 2.2 .1 gives $p=\widetilde{p}$, and for the error
constants we found

$$
\left.B_{p+1, \ell}=\frac{(-1)^{p+1-\ell}}{(p+1-\ell)!\ell!} \sum_{i=0}^{k} i^{p^{p}\left[i a_{i}\right.}+(p+1) b_{i}\right], \quad 1 \leq \ell \leq p+1
$$

In Table 2.2.2 the $B_{p+1, \ell}$ are given for the $B D$ formulas (cf. Table 2.2.1)。

Table 2.2.2 Error constants $B_{p+1, \ell}=B_{p+1, \ell}^{*} \frac{1}{\ell!(p+1-\ell)!}$,

$$
-\ell=1(1) \mathrm{p}+1
$$

| VLM[LM] | $p$ | $k=1$ | $k=2$ | $k=3$ | $k=4$ | $k=5$ |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| ILM[BD] | $k$ | $B_{p+1, \ell}^{*}$ | $=-1$ | $\frac{4}{3}$ | $-\frac{36}{11}$ | $\frac{288}{25}$ | $-\frac{7200}{137}$ |
| $M M L[B D]$ | $k$ | $(-1)^{p+1-\ell_{B}^{*}}$ |  |  |  |  |  |

### 2.3. Convergence

We first give a definition of convergence.

Definition 2.3.1. A VLM method is said to yield a convergent solution for (1.1) if $y_{n} \rightarrow y\left(t_{n}\right)$ as $h \rightarrow 0$, with $t_{n}=t$ fixed, holds for all convergent starting values $\left\{y_{i}, Y_{i}\left(t_{j}\right)\right\}, i=1, \ldots, k-1, j=-i,-i+1, \ldots, k-i$.

Before considering the convergence of VLM methods for (1.1) we answer the question to what equation the numerical scheme (2.1.4) converges if we substitute a sufficiently differentiable function $y(t)$ (not neccessarily the exact solution) and if we then let $h$ tend to zero in a fixed point. To that end, we define the polynomial

$$
\begin{equation*}
\alpha(z):=\sum_{i=0}^{k} \alpha_{i} z^{k-i} \tag{2.3.1}
\end{equation*}
$$

and the quantities

$$
\begin{equation*}
A_{q}:=\sum_{i=0}^{k}(-i)^{q_{\alpha_{i}}}, \quad q=0,1, \ldots \tag{2.3.2}
\end{equation*}
$$

We observe that $A_{0}=\alpha(1), A_{1}=\alpha^{\prime}(1)-k \alpha(1), \ldots$.
THEOREM 2.3.1. If sufficiently differentiable functions $y(t), g(t)$ and $K(t, \tau, y)$ are substituted into the VLM method, then the method converges to the equation

$$
\begin{align*}
\sum_{q=0}^{m} h^{q}\left\{\frac{\theta A}{q!} \frac{d^{q}}{d t}{ }^{q} y(t)\right. & \left.+\sum_{\ell=0}^{q}\left[C_{q \ell}-\frac{A_{q}}{\ell!(q-\ell)!}\right]\left(\frac{\partial}{\partial t}\right)^{q-\ell}\left(\frac{\partial}{\partial s}\right)^{\ell} Y(t, t)\right\}  \tag{2.3.3}\\
& =0\left(h^{r}+h^{m+1}\right) \text { as } h \rightarrow 0
\end{align*}
$$

where $r$ is defined in (2.1.3), $C_{q l}$ in (2.2.3), and $m$ is some integer $\geq 0$ determined by the differentiability of $\mathrm{y}, \mathrm{g}$ and k .

Examples. In the case of the $D Q$ method for $\{(1.1), \theta=1\}$ we have $A_{0}=1$, so that we infer from Corollary 2.2.1 and Theorem 2.3.1 that the numerical scheme converges to the equation $y(t)-Y(t, t)=0$ as $h \rightarrow 0$, which is the original equation (1.1). In the case of the $D Q$ method for $\{(1.1), \theta=0\}$ we have $C_{00}=1$, so that it easily follows that the numerical scheme converges as $h \rightarrow 0$ to the equation $Y(t, t)=0$, also the original equation (1.1). In the case of the ILM method for (1.1) it is not difficult to show that if the coefficients $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$ in (2.1.6) correspond to a convergent LM formula for ODEs, then the numerical scheme converges, as $h \rightarrow 0$, to the differentiated Volterra equation $\theta y^{\prime}(t)=K(t, t, y(t))+Y_{t}(t, t)$.

In order to present convergence theorems for VLM methods, we need the following concepts and definitions: A polynomial is called simple von Neumann if its zeros lie on the unit disk, those on the unit circle being simple. A polynomial is called Schur if its zeros lie within the
unit circle. Besides $\alpha(z)$ defined above, we define
(2.3.4a) $\quad \beta(z):=\sum_{i=0}^{k} \beta_{i} z^{k-i}$ where $\beta_{i}:=\sum_{j=-i}^{k-i} \beta_{i j}$;
(2.3.4b)

$$
\gamma(z):=\sum_{i=0}^{k} \gamma_{i} z^{k-i} \text { where } \gamma_{i}:=\sum_{j=-i}^{k-i} \gamma_{i j} .
$$

Furthermore, we need

$$
\begin{equation*}
b:=\max _{i, j}\left|\beta_{i j}\right|, c:=\max _{i, j}\left|\gamma_{i j}\right|, w:=\max _{i, j}\left|w_{i j}\right| ; \tag{2.3.5}
\end{equation*}
$$

(2.3.6a) $\left\{\begin{array}{l}\Delta K\left(t, s, y, y^{*}\right):=K(t, s, y)-K\left(t, s, y^{*}\right), \\ \Delta E(h):=\max _{i \leq j \leq N}\left|E_{i}\left(h ; t_{j}\right)-E_{i}\left(h ; t_{j+\ell}\right)\right|,\end{array}\right.$

where $y(t)$ is the exact solution of (1.1) resp. (1.2) and $Y(t, s), Y_{n}(t)$ are the corresponding functions defined in (1.5) resp. (2.1.2). E(h) is the maximal error arising in the approximation of the lag terms $Y_{n}(t)$ during the integration process until $t=T, T(h)$ may be considered as the maximal local truncation error of the VLM formula (2.1.4) until $t=T$, and $\delta(h)$ is the maximal starting error.

### 2.3.1. Second kind equations

We are now in a position to state a general convergence theorem for VLM methods in the case of second kind equations ( $\theta=1$ ), which provides an estimate for the global error

$$
(2.3 .7) \quad \varepsilon_{n}:=y\left(t_{n}\right)-y_{n}
$$

THEOREM 2.3.2. Let the conditions for the existence of a unique solution $\mathrm{y} \in \mathrm{C}(\mathrm{I})$ of $\{(1,1), \theta=1\}$ be satisfied (see Section 1).
(a) If $\alpha(z) \equiv \alpha_{0} z^{k}, \alpha_{0} \neq 0$, then there exists a constant $C>0$, independent of $h$, such that for $h$ sufficiently small

$$
\left|\varepsilon_{n}\right| \leq C[h \delta(h)+E(h)+T(h)], \quad n=k^{*}, \ldots, N .
$$

(b) If $\alpha(z)$ is simple von Neumann, if $\beta(z) \equiv 0$, and if $\Delta K$ satisfies the (uniform) Lipschitz condition

$$
\left|\Delta K\left(t, s, y, y^{*}\right)-\Delta K\left(t^{*}, s, y, y^{*}\right)\right| \leq L\left|t-t^{*}\right|\left|y-y^{*}\right|,
$$

for all $(\mathrm{t}, \mathrm{s}, \mathrm{y}),\left(\mathrm{t}, \mathrm{s}, \mathrm{y}^{*}\right),\left(\mathrm{t}^{*}, \mathrm{~s}, \mathrm{y}\right),\left(\mathrm{t}^{*}, \mathrm{~s}, \mathrm{y}^{*}\right) \in \mathrm{S} \times\{|\mathrm{y}|<\infty\}$, where the Lipschitz constant $L$ is independent of $t, t^{*}, y$ and $y^{*}$, then there exists $a$ constant C > 0 , independent of $h$, such that for $h$ sufficiently small

$$
\left|\varepsilon_{n}\right| \leq C h^{-1}[h \delta(h)+\Delta E(h)+T(h)], \quad n=k^{*}, \ldots, N
$$

Using Theorem 2.3.2 it is easy to derive the orders of convergence of the subclasses introduced in Section 2.1. The results are given in the following

Corollary 2.3.1. Let $\tilde{p}$ be the order of consistency of the LM formula for ODEs defining the coefficients $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$ employed in the ILM and (M)ML methods, let $\delta(\mathrm{h})=O\left(\mathrm{~h}^{\mathrm{s}}\right), \mathrm{E}(\mathrm{h})=O\left(\mathrm{~h}^{\mathrm{r}}\right)$ and $\Delta \mathrm{E}(\mathrm{h})=O\left(\mathrm{~h}^{\mathrm{r}+1}\right)$ as $\mathrm{h} \rightarrow 0$. Then the order of convergence $p$ is given by

$$
p=\left\{\begin{array}{ll}
\min \{s+1, r\} & \text { for the DQ method } \\
\min \{s, r, \tilde{p}, k\} & \text { for the ILM method } \\
\min \{s+1, r, \tilde{p}+1\} & \text { for the ML method } \\
\min \{s, r, \tilde{p}\} & \text { for the MML method }
\end{array} .\right.
$$

The convergence analysis of the DQ methods goes back to KOBAYASI [7], LINZ [9] and NOBLE [13]. The (M)ML methods were proved to be convergent in WOLKENFELT [16].

It is known (cf. $[6,16]$ ) that VLM methods which have $\beta(z) \equiv 0$ behave more stable then DQ methods if large Lipschitz constants for $\partial K / \partial y$ are involved. The maximal attainable order of convergence of these VLM methods is expressed in the following

Coro1lary 2.3.2. Let $\delta(\mathrm{h})=O\left(\mathrm{~h}^{\mathrm{s}}\right), \Delta \mathrm{E}(\mathrm{h})=O\left(\mathrm{~h}^{\mathrm{r}+1}\right)$ as $\mathrm{h} \rightarrow 0$, let $\alpha(\mathrm{z})$ be simple von Neumann and let $\beta(z) \equiv 0$. Then the order of convergence $p$ of the k-step VLM method \{(2.1.4); (2.1.2)\} satisfies

$$
p \leq \begin{cases}\min (s, r, k+1) & \text { for } k \text { odd } \\ \min (s, r, k+2) & \text { for } k \text { even } .\end{cases}
$$

From this corollary it follows that the MML methods are of maximal attainable order of convergence if we choose the generating LM formula $\{\rho, \sigma\}$ to be optimal, that is, of order $k+1$ when $k$ is odd and of order $\mathrm{k}+2$ when k is even. We note that the restriction $\mathrm{p} \leq \mathrm{k}$ in the ILM methods is due to the use of a k-step forward differentiation formula $\left\{\delta_{i}\right\}$ in the generating matrix $B$ (see (2.1.6)).

### 2.3.2 First kind equations

Now we shall give convergence theorems for VLM methods for Volterra first kind equations $(\theta=0)$. We restrict our attention here to Zinear equations, i.e., we assume in (1.1) that

$$
\begin{equation*}
K(t, \tau, y(\tau))=K(t, \tau) y(\tau) . \tag{2.3.8}
\end{equation*}
$$

We first give the following convergence theorem of WOLKENFELT [14] for ( $\rho, \sigma$ )-reducible DQ methods.

THEOREM 2.3.3. Let the conditions for the existence of a unique solution $y(t) \in C(I)$ of $\{(1.1), \theta=0\}$ be satisfied (see Section 1), where $K(t, \tau, y)$ is of the form given in (2.3.8). Let $\mathrm{A}=0, \mathrm{~B}=1, \mathrm{C}=0$ in (2.1.5) (DQ formula) and let the weights in (2.1.2) be given by a $(\rho, \sigma)$-reducible quadrature sormula of order $\mathrm{r} \geq 1$, where $\sigma$ is simple von Neumann. Then there exists a constant $C>0$, independent of $h$, such that for $h$ sufficiently
smaZ2

$$
\left|\varepsilon_{n}\right| \leq C[\delta(h)+E(h)], \quad n=k, \ldots, N .
$$

PROOF. See [14].
THEOREM 2.3.4. Let the conditions for the existence of a unique solution $y \in C(I)$ of $\{(1.1), \theta=0\}$ be satisfied. Let $\beta(z) \equiv 0$ and let $\gamma(z)$ be Schur. Then there exists a constant $\mathrm{C}>0$, independent of h such that for h sufficiently small

$$
\left|\varepsilon_{n}\right| \leq C h^{-1}[h \delta(h)+\Delta E(h)+T(h)], n=k^{\bar{\Psi}}, \ldots, N .
$$

Observe that this convergence result is identical to that obtained for VLM methods for $\{(1.1), \theta=1\}$ with $\beta(z) \equiv 0$ (Theorem 2.3.2(b)). Now it is easy to derive from Theorems 2.3.3 and 2.3.4 the orders of convergence of the $D Q$, ILM and MML methods for $\{(1.1), \theta=0\}$. The results are summarized in the following

Corollary 2.3.3. Let $\bar{p}$ be the order of consistency of the LM formula $\{\rho, \sigma\}$ employed in the $D Q$ lag term formula (2.1.2) and let $\tilde{p}$ be the order of consistency of the LM formula for ODEs employed in the MML formula (2.1.8) (with $A \equiv 0$ ). Furthermore, let $\delta(h)=O\left(h^{s}\right), E(h)=O\left(h^{r}\right)$ and $\Delta E(h)=O\left(h^{r+1}\right)$ as $\mathrm{h} \rightarrow 0$. Then the order of convergence p is given by

$$
\mathrm{p}= \begin{cases}\min \{\mathrm{s}, \overline{\mathrm{p}}\} & \text { for the } D Q \text { method with } \sigma \text { being simple von Neumann, } \\ \min \{\mathrm{s}, \mathrm{r}, \mathrm{k}\} & \text { for the ILM method with } \gamma \text { being Schur, } \\ \min \{\mathrm{s}, \mathrm{r}, \tilde{p}\} & \text { for the MML method with } \gamma \text { being Schur. }\end{cases}
$$

WOLKENFELT [15] has also given a convergence theorem for MML methods for nonlinear equations $\{(1.1), \theta=0\}$, with the following restrictions on the coefficients $\beta_{i j}$ and $\gamma_{i j}: \beta(z) \equiv 0, \gamma_{00} \neq 0$ and all other $\gamma_{i j}$ vanish.

### 2.4 Numerical experiments

In this section we present the results of numerical experiments in
order to support and illustrate the convergence behaviour of VLM methods for (1.1) as predicted by Corollaries 2.3.1 and 2.3.3, by applying various DQ , ILM and (M)ML methods to a number of problems.

The required starting values for $y_{i}, 0 \leq i<n_{1}+k$, are taken from the exact solution $y\left(t_{i}\right)$ (so that $s=\infty$ in Corollaries 2.3.1 and 2.3.3), and values of the lag term $Y_{n}(t)$ required in (2.1.4) for $n \geq n_{1}$ are computed with a Gregory quadrature rule in (2.1.2) of the proper order. The coefficients $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$ in the ILM and (M)ML formulas are taken from Tables 2 and 3 of the Appendix. The values of $r, \tilde{p}$ and $k$ in Corollary 2.3 .1 and of $r, \bar{p}, \tilde{p}$ and k in Corollary 2.3.3 are chosen as small as is allowed by the theoretical order to be tested.

In the tables of results we present the number of correct significant digits in the end point $T$, i.e., the value of
(2.1.4) $\quad \operatorname{sd}(h):=-\log _{10}\left(\left|y(T)-y_{N}\right| /|y(T)|\right), \quad T=t_{N}=N h$.

Moreover, we list the effective order of the method, viz., the value of $(s d(h)-s d(2 h)) / \log _{10} 2$. This value should tend to the asymptotic order of convergence as $h \rightarrow 0$ and will tell us therefore
(i) whether the asymptotic, theoretical order of the numerical scheme is correct, and
(ii) how fast the asymptotic order is reached.
2.1.4. Second kind equations ((1.1) with $\theta=1$ )

Example 2.1.4 (GAREY [2], adapted)

$$
\begin{aligned}
K(t, \tau, y) & =-\lambda \cdot \ln (1+t-\tau) y, \\
g(t) & =1-t+\lambda\left[\frac{1}{2}\left(1-t^{2}\right) \ln (1+t)+\frac{3}{4} t^{2}-\frac{1}{2} t\right] \\
y(t) & =1-t, \\
{\left[t_{0}, T\right] } & =[0,4] .
\end{aligned}
$$

For $\lambda=4$ and $\lambda=100$ Tables 2.4 .1 and 2.4 .2 give the results obtained with $D Q$, IIM and (M)ML methods of asymptotic order 5, where for the coefficients $\left\{a_{i}, b_{i}\right\}$ employed in the IIM and (M)ML methods we used the coefficients of the Adams-Moulton formula of the proper order. $G_{r}$ means that for the lag term we used a Gregory rule of order $r$ and $A M \sim$ means that a $\tilde{p}$-th order Adams-Moulton formula was used.

Table 2.4.1 Examp1e 2.4.1 with $\lambda=4$

| h | $\mathrm{DQ}\left(\mathrm{G}_{5}\right)$ | $\operatorname{ILM}\left(\mathrm{G}_{5}-\mathrm{AM}_{6}\right)$ | $\mathrm{ML}\left(\mathrm{G}_{5}-\mathrm{AM}_{4}\right)$ | $\operatorname{MML}\left(\mathrm{G}_{5}-\mathrm{AM}_{5}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
| 1/4 | 4.6 | $3.4$ | $4.3$ | $6.1$ |
|  | $>4.6$ | $>3.8$ | $>4.8$ | $>3.9$ |
| $1 / 8$ | 6.0 | 4.5 | 5.7 | 7.3 |
|  | > 5.0 | > 4.5 | >4.6 | >3.2 |
| 1/16 | 7.5 | 5.9 | 7.1. | 8.2 |
|  | > 5.0 | > 4.7 | >5.0 | >4.0 |
| 1/32 | 9.0 | 7.3 | 8.6 | 9.4 |
|  | $>5.0$ | >4:8 | >5.0 | >4.6 |
| 1/64 | 10.5 | 8.8 | 10.1 | 10.8 |

Table 2.4.2 Example 2.4.1 with $\lambda=100$

| h | $\mathrm{DQ}\left(\mathrm{G}_{5}\right)$ | $\operatorname{ILM}\left(\mathrm{G}_{5}-\mathrm{AM}_{6}\right)$ | $\mathrm{ML}\left(\mathrm{G}_{5}-\mathrm{AM}_{4}\right)$ | $\operatorname{MML}\left(\mathrm{G}_{5}-\mathrm{AM}_{5}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
| $1 / 4$ | $-6.5$ | 1.8 | -3.7 | -2.4 |
|  | >30 | >8.9 | >25 | >22 |
| 1/8 | 2.3 | 4.5 | 3.7 | 4.2 |
|  | >13 | >4.1 | >8.2 | $>16$ |
| 1/16 | 6.3 | 5.8 | 6.2 | 9.0 |
|  | >5.8 | >4.4 | >4.7 | >2.3 |
| 1/32 | 8.1 | 7.1 | 7.6 | 9.7 |
|  | >6.8 | >6.4 | >5.8 | >2.5 |
| 1/64 | 10.1 | 9.0 | 9.3 | 10.4 |

The example with $\lambda=4$ shows that the correct asymptotic order $p=5$ is reached by all methods for not too small integration steps. For relatively large values of $h$ the MML method shows the most accurate results. The ILM method shows an accuracy about 1-2 digits less than the other methods, due to larger error constants (cf. Table 2.2.1). The example with $\lambda=100$ shows that the ILM method is stable for "realistic" values of $h$ (in view of the behaviour of the exact solution, integration steps $h=1 / 4$ or $h=1 / 8$ should be small enough for representing the function $y(t)=1-t$ ), whereas the other methods develop instabilities. If $h$ is decreased the ( $M$ ) ML methods become more accurate than the ILM method. Except for the ML method this experiment does not yet demonstrate the asymptotic order $\mathrm{p}=5$.

Thus, we conclude that the (M)ML methods are superior for nonstiff problems ( $h|\partial K / \partial y| s m a l l$ ), and the ILM methods superior for stiff problems (h $|\partial K / \partial y|$ large).

### 2.4.2 First kind equations ( $(1.1)$ with $\theta=0)$

Example 2.4.2 (GLADWIN [3])

$$
\begin{aligned}
K(t, \tau, y) & =\cos (t-\tau) y \\
g(t) & =-\exp (t)-\sin (t)+\cos (t), \\
y(t) & =\exp (t), \\
{\left[t_{0}, T\right] } & =[0,4] .
\end{aligned}
$$

Table 2.4.3 gives the results obtained with $D Q$, ILM and MML methods of asymptotic order 4 and 5 , where for the coefficients $\left\{a_{i}, b_{i}\right\}$ in the ILM and MML methods we used the coefficients of the backward differentiation formulas of the proper order. $\mathrm{BD}_{\mathrm{k}}$ means that a k-step (k-th order) $B D$ formula was used.

Table 2.4.3 Example 2.4.2

| h | $\mathrm{DQ}\left(\mathrm{G}_{4}\right)$ | $\operatorname{ILM}\left(G_{4}-\mathrm{BD}_{4}\right)$ | $\operatorname{MML}\left(\mathrm{G}_{4}-\mathrm{BD}_{4}\right)$ | DQ (G) | $\operatorname{ILM}\left(\mathrm{G}_{5}-\mathrm{BD}_{5}\right)$ | $\mathrm{MML}\left(\mathrm{G}_{5}-\mathrm{BD}_{5}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1/10 | -7.6 | 4.3 | 3.9 | -11 | 5.6 | 4.9 |
|  |  | >4.1 | >3.8 |  | >4.9 | >4.9 |
| 1/20 | -21 | 5.5 | 5.1 | \| -29 | 7.0 | 6.4 |
|  |  | >4.0 | >3.9 | 1 | >5.0 | >5.0 |
| 1/40 | -50 | 6.7 | 6.3 | 1-65 | 8.5 | 7.9 |
|  |  | >4.0 | >4.0 |  | >5.2 | $>5.0$ |
| 1/80 | -109 | 7.9 | 7.5 | , -140 | 10.1 | 9.4 |
|  |  |  |  | 1 |  |  |

For the ILM and the MML methods the results show that the correct asymptotic order is reached already for relatively large values of $h$. The apparent unstable behaviour of the $D Q\left(G_{4}\right)$ and $D Q\left(G_{5}\right)$ methods is explained by the fact that the Gregory quadrature formulas of order $\geq 3$ are ( $\rho, \sigma$ )-quadrature formulas for which the $\sigma$-polynomial is not simple von Neumann (cf. Corollary 2.3.3). Unlike its performance for second kind equations, the ILM method is here more accurate than the MML method.
3. VLM METHODS FOR VOLTERRA INTEGRO-DIFFERENTIAL EQUATIONS

### 3.1. The general VLM method

In analogy to the VLM formula (2.1.4) for Volterra integral equations we formally define the VLM formula for Volterra integro-differential equations (1.2) as follows:

$$
\begin{equation*}
\sum_{i=0}^{k} \alpha_{i}{ }^{*} y_{n-i}=h \sum_{i=0}^{k} \gamma_{i}^{*} f_{n-i}, \quad f_{n}:=f\left(t_{n}, y_{n}, z_{n}\right), \tag{3.1.1a}
\end{equation*}
$$

$$
\begin{array}{r}
\sum_{i=0}^{k} \alpha_{i} z_{n-i}+\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j} Y_{n-i}\left(t_{n+j}\right)=h \sum_{i=0}^{k} \sum_{j=-k}^{k} \gamma_{i j} j_{n-i}\left(t_{n+j}\right)  \tag{3.1.1b}\\
n=k^{*}(1) N
\end{array}
$$

where $Y_{n}(t)$ is defined as in (2.1.2) and $\left\{\alpha_{i}{ }^{*}, \gamma_{i}{ }^{*}\right\}_{i=0}^{k}$ are the coefficients of some LM method for ODEs. These formulas (3.1.1), combined with (2.1.2), will be called a VLM method for integro-differential equations. Formula
(3.1.1b) can be characterized by the parameter matrices $A=\left(\alpha_{i}\right), B=\left(\beta_{i j}\right)$, $C=\left(\gamma_{i j}\right)$.

When we compare (3.1.1) with (2.1.4), it is clear that aZl methods defined for second kind Volterra integral equations by specifying the matrices $A, B$ and $C$ in (2.1.5) and the quadrature weights $w_{n j}$ in (2.1.2), can be extended to methods for Volterra integro-differential equations by specifying the coefficients $\left\{\alpha_{i}{ }^{*}, \beta_{i}{ }^{*}\right\}$ of some LM method for ODEs. In this way, we define $D Q$, ILM, ML and MML methods for (1.2) where the matrices A, B and C are specified in Section 3.1, Subclasses 1, 2, 3 and 4, respectively. For example, any $D Q$ method for (1.2) is specified by $A=1, B=-1$, $C=0$, which gives $z_{n}=Y_{n}\left(t_{n}\right)$ for (3.1.1b)。

An alternative way to arrive at the VLM formulas (3.1.1) is obtained as follows. We first integrate (1.2) formally, which results in the system of Volterra integral equations of the second kind

$$
\left\{\begin{array}{l}
y(t)=y\left(t_{0}\right)+\int_{t_{0}}^{t} f(\tau, y(\tau), z(\tau)) d \tau  \tag{3.1.2}\\
z(t)=g(t)+\int_{t_{0}}^{t_{0}} K(t, \tau, y(\tau)) d \tau
\end{array}\right.
$$

Next, we apply the VLM formula (2.1.4) to this system with parameter matrices $\left(A^{*}, B^{*}, C^{*}\right)$ and ( $A, B, C$ ) for the respective components, i.e.,

$$
\sum_{i=0}^{k} \alpha_{i}{ }^{*} y_{n-i}+\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j}^{*} y_{n-i}^{*}\left(t_{n+j}\right)=h \sum_{i=0}^{k} \sum_{j=-k}^{k} \gamma_{i j}^{*} f\left(t_{n-i} y_{n-i}, z_{n-i}\right)
$$

(3.1.1')

$$
\sum_{i=0}^{k} \alpha_{i} z_{n-i}+\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j} Y_{n-i}\left(t_{n+j}\right)=h \sum_{i=0}^{k} \sum_{j=-k}^{k} \gamma_{i j} K\left(t_{n+j}, t_{n-i}, y_{n-i}\right)
$$

where $Y_{n}^{*}\left(t_{j}\right)$ is an approximation to

$$
Y^{*}(t, s):=y\left(t_{0}\right)+\int_{t_{0}}^{s} f(\tau, y(\tau), z(\tau)) d \tau
$$

at $t=t_{j}, s=t_{n}$. Here, however, both $Y^{*}$ and $f$ do not depend on $t$ so that, by putting $\beta_{i}^{*}:=\sum_{j=-k}^{k} \beta_{i j}^{*}=0$ and writing $\gamma_{i}^{*}:=\Sigma_{j=-k}^{k} \gamma_{i j}^{*}$ we have reduced (3.1.1') to (3.1.1).

### 3.2. Consistency of VLM formulas for integro-differential equations

With the numerical schemes (3.1.1) we associate the linear differencedifferential operators ${\underset{\sim}{\sim}}^{*}: C^{1}(I) \rightarrow \mathbb{R}$ and $\underset{\sim}{\sim}{ }_{n}: C_{2}^{1}(S) \rightarrow \mathbb{R}$, defined by (3.2.1a) $\quad \underset{\sim}{L_{n}^{*}}[y]:=\sum_{i=0}^{k}\left[\alpha_{i}^{*}-\gamma_{i}^{*} h \frac{d}{d t}\right] y\left(t_{n-i}\right)$
and
(3.2.1b) $\quad \underset{\sim}{L}{ }_{n}[Y]:=\sum_{i=0}^{k}\left\{\alpha_{i} Y\left(t_{n-i}, t_{n-i}\right)+\sum_{j=-k}^{k}\left[\beta_{i j}-\gamma_{i j} h \frac{\partial}{\partial s}\right] Y\left(t_{n+j}, t_{n-i}\right)\right\}$,
where $y$ and $Y$ are arbitrary functions from $C^{1}(I)$ and $C_{2}^{1}(S)$, respectively. Now we substitute the exact solution $y(t)$ and $z(t)$ of (1.2) into (3.2.1) and obtain (cf. (2.2.2))

$$
\begin{aligned}
& \sum_{i=0}^{k}\left[\alpha_{i}^{*} y\left(t_{n-i}\right)-\gamma_{i}^{*} h f\left(t_{n-i}, y\left(t_{n-i}\right), z\left(t_{n-i}\right)\right)\right]=L_{\sim}^{*}[y] \\
& \begin{aligned}
\sum_{i=0}^{k}\left\{\alpha_{i} z\left(t_{n-i}\right)\right. & \left.+\sum_{j=-k}^{k}\left[\beta_{i j} \tilde{Y}_{n-i}\left(t_{n+j}\right)-\gamma_{i j} h \tilde{K}_{n-i}\left(t_{n+j}\right)\right]\right\}= \\
& ={\underset{\sim}{n}}_{n}^{L}[Y]-\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j} E_{n-i}\left(h ; t_{n+j}\right) \\
& ={\underset{\sim}{n}}_{n}^{L}[Y]+O\left(h^{r}\right) \text { as } h \rightarrow 0,
\end{aligned}
\end{aligned}
$$

where $\tilde{Y}_{n-i}\left(t_{n+j}\right)$ and $\tilde{K}_{n-i}\left(t_{n+j}\right)$ are defined by (2.1.1) and (2.1.2) with $y_{n}$ replaced by $y\left(t_{n}\right)$ and when $r$ is the order of the quadrature error $E_{n}$. This shows the connection of the operators (3.2.1) with the VLM formula (3.1.1). The quantities ${\underset{\sim}{\sim}}^{*}[y]$ and ${\underset{\sim}{n}}^{L}[Y]$, with $y$ and $Y$ corresponding to the exact solution of (1.3), are called the local truncation errors of the VLM formulas (3.1.1). In analogy to Section 2.2 we use the following

Definition 3.2.1. The operators (3.2.1) and the associated VLM formulas (3.1.1) are said to be consistent of order $p^{*}$ and $p$ with the equations (1.2) if for all $\mathrm{y} \in \mathrm{C}^{\mathrm{P}^{\star+1}}(\mathrm{I})$ and for all $\mathrm{Y} \in \mathrm{C}^{\mathrm{P}+1}(\mathrm{~S})$, we have $\mathrm{I}^{\star}[\mathrm{y}]=0\left(\mathrm{~h}^{\mathrm{P}^{\star}+1}\right)$ and $L_{n}[Y]=O\left(h^{+1}\right)$ as $h \rightarrow 0$, with non-vanishing error constants.

Since ${\underset{\sim}{n}}_{n}^{*}$ is of the same form as the linear operator ${\underset{\sim}{n}}^{n}$ occurring in ODE theory (compare LAMBERT [8, p. 23]) the consistency conditions for ${\underset{\sim}{n}}^{*}$ are also of the same form. (It should be remarked that in the derivation of the consistency conditions we expand $y\left(t_{n-i}\right)$ and $y^{\prime}\left(t_{n-i}\right)$ as Taylor series about $t_{n}$, whereas Lambert expands about $t_{n-k}$.) Similarly, since ${\underset{\sim}{n}}_{n}$ defined in (3.2.1b) is identical to the operator defined in (2.2.1), the consistency conditions for ${\underset{\sim}{n}}^{n}$ are also known already. Therefore, the following consistency theorem is immediate.

THEOREM 3.2.1. The operators ${\underset{\sim}{~}}_{\mathrm{n}}^{*}$ and ${\underset{\sim}{\sim}}_{\mathrm{n}}$ and the associated VLM formulas (3.1.1) are consistent of order $p^{*}$ and $p$ with (1.2) if $c_{q}=0$ for $\mathrm{q}=0,1, \ldots, \mathrm{p}^{*}$ and $\mathrm{C}_{\mathrm{q} \ell}=0$ for $\mathrm{q}=0,1, \ldots, \mathrm{p}, \ell=0,1, \ldots, \mathrm{q}$, where

$$
c_{q}:=\frac{(-1)^{q}}{q!} \sum_{i=0}^{k} i^{q-1}\left[i \alpha_{i}^{*}+q \gamma_{i}^{*}\right]
$$

and where $\mathrm{C}_{\mathrm{q}}$ is defined in (2.2.3).
Evidently, $p^{*}$ equals the order of consistency of the LM method for ODEs with coefficients $\left\{\alpha_{i}^{*}, \gamma_{i}^{*}\right\}$. Furthermore, in the case of the DQ, ILM, $M L$ and MML formulas, $p$ is determined by the expressions as derived for the operator ${\underset{\sim}{n}}^{\mathrm{n}}$ for second kind Volterra integral equations in Corollary 2.2.1.

The values of the error constants $C_{p^{*}+1}$ and $C_{p+1, \ell}, 0 \leq \ell \leq p+1$, follow easily from those given in Table 2.2 .1 (for a number of popular methods for second kind Volterra integral equations).

### 3.3. Convergence

As we did for the Volterra equations of first and second kind we first consider the continuous problem to which the VLM method \{(3.1.1); (2.1.2)\} converges as $h \rightarrow 0$. We assume that the $L M$ formulas in (3.1.1) are consistent and that $A_{1}=\alpha^{\prime}(1)-k \alpha(1) \neq 0$ (see Section 2.3). Then, for sufficiently smooth functions $g, K$ and $f$ the VLM method converges to the equations

$$
\begin{equation*}
y^{\prime}(t)=f(t, y(t), z(t)) \tag{3.3.1}
\end{equation*}
$$

$$
\alpha(1)[z(t)-Y(t, t)]+\left(\alpha^{\prime}(1)-k \alpha(1)\right) h\left[z^{\prime}(t)-Y_{t}(t, t)-Y_{S}(t, t)\right]=0
$$

as $h \rightarrow 0$ (see the proof of Theorem 2.3.1). Thus, if $\alpha(1) \neq 0$ (DQ and $M L$ method), then the vLM method is a direct discretization of (1.3). If $\alpha(1)=0$ (ILM and MML) (and $\alpha^{\prime}(1) \neq 0$ by assumption), then the linear method converges to the system

$$
y^{\prime}(t)=f(t, y(t), z(t))
$$

$$
\begin{equation*}
z^{\prime}(t)=K(t, t, y(t))+g^{\prime}(t)+\int_{t_{0}}^{t} K_{t}(t, \tau, y(\tau)) d \tau \tag{3.3.2}
\end{equation*}
$$

that is the system (1.3) where the expression for $z(t)$ is differentiated with respect to $t$.

Next we present a general convergence theorem. In the proof it is convenient to use, in addition to the notation introduced in Section 2.3 , the notations

$$
\begin{align*}
& \eta_{n}=z\left(t_{n}\right)-z_{n}, \\
& \Delta f_{n}:=f\left(t_{n}, y\left(t_{n}\right), z\left(t_{n}\right)\right)-f\left(t_{n}, y_{n}, z_{n}\right), \\
& T^{*}(h):=\max _{i \leq N}\left|{\underset{\sim}{i}}_{i}^{*}[y]\right|, \delta^{*}(h):=\max _{j \leq k-1}\left|z\left(t_{j}\right)-z_{j}\right|,  \tag{3.3.3}\\
& \alpha^{*}(z):=\sum_{i=0}^{k} \alpha_{i}^{*} z^{k-i},
\end{align*}
$$

where $z$ and $y$ correspond to the exact solution.

THEOREM 3.3.1. Let the conditions for the existence of a unique solution $y \in C^{1}(I)$ of (1.2) be satisfied (see Section 1). Let $\alpha(z)$ and $\alpha^{*}(z)$ be simple von Neumann.
(a) If $\alpha(z)=\alpha_{0} z^{k}, \alpha_{0} \neq 0$, then there exists a constant $C>0$, independent of $h$ such that for $h$ sufficiently small

$$
\left|\varepsilon_{n}\right| \leq C\left[\delta(h)+h \delta^{*}(h)+E(h)+T(h)+h^{-1} T^{*}(h)\right], n=k^{*}, \ldots, N .
$$

(b) If $\beta(z) \equiv 0$, then there exists a constant $C>0$, independent of $h$ such that for $h$ sufficiently small

$$
\begin{array}{r}
\left|\varepsilon_{n}\right| \leq C\left[\delta(h)+\delta^{*}(h)+h^{-1} \Delta E(h)+h^{-1} T(h)+h^{-1} T^{*}(h)\right], \\
n=k^{*}, \ldots, N .
\end{array}
$$

Using this theorem, it is easy to derive the orders of convergence of the various examples of VLM methods for Volterra integro-differential equations described in Section 3.1. The results are given in the following Corollary 3.3.1. Let $\mathrm{p}^{*}$ and $\tilde{p}$ be the order of consistency of the VLM formula $\left\{\alpha_{i}^{*}, \gamma_{i}^{*}\right\}$ employed in (3.1.1a) and of the LM formula $\left\{\mathrm{a}_{\mathrm{i}}, \mathrm{b}_{\mathrm{i}}\right\}$ employed in the (M)ML and ILM methods, respectively; let $\delta(h)=O\left(h^{s}\right), \delta^{*}(h)=0\left(h^{*}\right)$, $E(h)=O\left(h^{r}\right), \Delta E(h)=O\left(h^{r+1}\right)$ as $h \rightarrow 0$. Then the order of convergence $p$ of the VLM method $\{(3.1 .1)$; (2.1.2) \} is given by

$$
p= \begin{cases}\min \left(s, s^{*}+1, r, p^{*}\right) & \text { for the } D Q \text { method } \\ \min \left(s, s^{*}+1, r, p^{*}, \tilde{\left.p^{+}+1\right)}\right. & \text { for the ML method } \\ \min \left(s, s^{*}, r, p^{*}, \tilde{p}^{2}\right) & \text { for the MML method } \\ \min \left(s, s^{*}, r, p^{*}, \tilde{p}, k\right) & \text { for the ILM method. }\end{cases}
$$

The convergence of the conventional DQ method has already been studied by LINZ [10] and MOCARSKY [12]. The convergence results for the (M)ML methods has already been given in WOLKENFELT [16].

### 3.4. Numerical experiments

In order to illustrate the convergence behaviour of VLM methods for (1.2) we have tested the DQ, ILM and (M)ML methods of orders 2, 3 and 4. For the two ODE-LM formulas involved in (3.1.1) we chose the backward differentiation formulas. As in the experiments for (1.1), the lag term $Y_{n}(t)$ was evaluated with a Gregory quadrature rule of the proper order.

Example 3.4.1. (LINZ [10], MOCARSKY [12], MAKROGLOU [11])

$$
\begin{aligned}
& f(t, y, z)=1-t \cdot \exp \left(-t^{2}\right)+y-2 z, \quad y(0)=0 \\
& K(t, \tau, y)=t \tau \exp \left(-y^{2}\right), \quad g(t)=0 \\
& y(t)=t, \quad\left[t_{0}, T\right]=[0,2]
\end{aligned}
$$

Table 3.4.1 gives the results of our experiments. The ILM method is the less accurate one, the $D Q$ and (M)ML methods exhibit a varying accuracy behaviour.

Table 3.4.1 Example 3.4.1
second order methods (with $\left\{\alpha_{i}^{*}, \gamma_{i}^{*}\right\}=\mathrm{BD}_{2}$ in (3.1.1a))

| h | DQ ( $\mathrm{G}_{2}$ ) | $\operatorname{ILM}\left(\mathrm{G}_{2}-\mathrm{BD}_{2}\right)$ | $\mathrm{ML}\left(\mathrm{G}_{2}-\mathrm{BD} 1\right)$ | $\operatorname{MML}\left(\mathrm{G}_{2}-\mathrm{BD}_{2}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
| 1/10 | 2.2 | 3.3 | 2.2 | 1.8 |
|  | >2.0 | >-2 | >2.2 | >2.0 |
| 1/20 | 2.8 | 2.6 | 2.8 | 2.4 |
|  | >2.0 | >1.3 | >2.0 | >2.0 |
| 1/40 | 3.4 | 3.0 | 3.5 | 3.0 |

third order methods (with $\mathrm{BD}_{3}$ in (3.1.1a))

| h | $\mathrm{DQ}\left(\mathrm{G}_{3}\right)$ | $\operatorname{ILM}\left(\mathrm{G}_{3}-\mathrm{BD}_{3}\right)$ | $\mathrm{ML}\left(\mathrm{G}_{3}-\mathrm{BD} 2\right)$ | $\operatorname{MML}\left(\mathrm{G}_{3}-\mathrm{BD}_{3}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
| 1/10 | 3.6 | 2.4 | 2.9 | 3.3 |
|  | >2.8 | $>2.3$ | >2.7 | >4.7 |
| 1/20 | 4.5 | 3.1 | 3.7. | 4.7 |
|  | >3.9 | >2.8 | >3.0 | >4.1 |
| 1/40 | 5.4 | 3.9 | 4.6 | 6.0 |

fourth order methods (with $\mathrm{BD}_{4}$ in (3.1.1a))

| h | $\mathrm{DQ}\left(\mathrm{G}_{4}\right)$ | $\operatorname{ILM}\left(\mathrm{G}_{4}-\mathrm{BD}_{4}\right)$ | $\mathrm{ML}\left(\mathrm{G}_{4}-\mathrm{BD}_{3}\right)$ | $\mathrm{MML}\left(\mathrm{G}_{4}-\mathrm{BD}_{4}\right)$ |
| :--- | :--- | :--- | :--- | :--- |
| $1 / 10$ | 4.0 | 3.2 | 3.6 | 3.6 |
| $1 / 20$ | $5.1^{>3.7}$ | $4.6^{>4.5}$ | 4.8 | $>4.0$ |
| $1 / 40$ | $6.3^{>3.9}$ | $6.4^{>6.2}$ | $6.1^{>4.1}$ | 4.6 |

## ACKNOWLEDGEMENT

We would like to thank Joke Blom for programming the numerical experiments.

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APPENDIX

In this appendix we present, successively,
(i) three tables of coefficients of forward differentiation formulas, and of two common LM formulas for ODEs, viz., backward differentiation formulas and Adams-Moulton formulas;
(ii) two lemmas which are needed in:
(iii) proofs of the main results of this paper, as far as they are nontrival (in the opinion of the authors).

Table 1 Coefficients of forward differentiation formulas

$$
f^{\prime}\left(t_{n}\right) \approx \frac{-1}{h} \sum_{\ell=0}^{k} \delta \ell^{f\left(t_{n+\ell}\right), \quad t_{n+\ell}=t_{n}+\ell h}
$$

| k | $\delta_{0}$ | $\delta_{1}$ | $\delta_{2}$ | $\delta_{3}$ | $\delta_{4}$ | $\delta_{5}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | -1 |  |  |  |  |
| 2 | $3 / 2$ | -2 | $1 / 2$ |  |  |  |
| 3 | $11 / 6$ | -3 | $3 / 2$ | $-1 / 3$ |  |  |
| 4 | $25 / 12$ | -4 | 3 | $-4 / 3$ | $1 / 4$ | $5 / 4$ |
| 5 | $137 / 60$ | -5 | 5 | $-10 / 3$ | $-1 / 5$ |  |

Table 2 Coefficients of the backward differentiation formulas

$$
\text { for ODEs } f^{\prime}(t)=g(t): \sum_{i=0}^{k} a_{i} f_{n-i}=b_{0} g_{n}
$$

| $k$ | $a_{0}$ | $a_{1}$ | $a_{2}$ | $a_{3}$ | $a_{4}$ | $a_{5}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | -1 |  |  | $b_{0}$ |  |
| 2 | 1 | $-4 / 3$ | $1 / 3$ |  |  | 1 |
| 3 | 1 | $-18 / 11$ | $9 / 11$ | $-2 / 11$ |  | $2 / 3$ |
| 4 | 1 | $-48 / 25$ | $36 / 25$ | $-16 / 25$ | $3 / 25$ |  |
| 5 | 1 | $-300 / 137$ | $300 / 137$ | $-200 / 137$ | $75 / 137$ | $-12 / 137$ |

Table 3 Coefficients of the Adams-Moulton formulas for ODEs $f^{\prime}(t)=g(t): f_{n}-f_{n-1}=\sum_{i=0}^{k} b_{i} g_{n-i}$

| $k$ | $b_{0}$ | $b_{1}$ | $b_{2}$ | $b_{3}$ | $b_{4}$ | $b_{5}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | $1 / 2$ | $1 / 2$ |  |  |  |  |
| 2 | $5 / 12$ | $2 / 3$ | $-1 / 12$ |  |  |  |
| 3 | $3 / 8$ | $19 / 24$ | $-5 / 24$ | $1 / 24$ |  |  |
| 4 | $251 / 720$ | $323 / 360$ | $-11 / 30$ | $53 / 360$ | $-19 / 720$ |  |
| 5 | $95 / 288$ | $1427 / 1440$ | $-133 / 240$ | $241 / 720$ | $-137 / 1440$ | $3 / 160$ |

LEMMA A.1. Let $\mathrm{z}_{\mathrm{n}} \geq 0$ for $\mathrm{n}=0,1, \ldots, \mathrm{~N}$, and suppose that

$$
z_{n} \leq h C_{1} \sum_{i=0}^{n-1} z_{i}+C_{2}, \quad n=k, k+1, \ldots, N
$$

where $\mathrm{k}>0$, $\mathrm{h}>0$ and $\mathrm{C}_{\mathrm{i}}>0(\mathrm{i}=1,2)$. Suppose, moreover, that $\mathrm{z}_{\mathrm{j}} \leq \mathrm{z} / \mathrm{k}$ for $\mathrm{j}=0,1, \ldots, \mathrm{k}-1$. Then

$$
z_{n} \leq\left(h C_{1} z+C_{2}\right)\left(1+h C_{1}\right)^{n-k}, \quad n=k, k+1, \ldots, N
$$

PROOF. See [ 4 ].

LEMMA A.2. Consider the Zinear inhomogeneous difference equation with constant coefficients $\zeta_{j}$ :

$$
\begin{equation*}
\zeta_{0} \mathrm{y}_{\mathrm{n}+\mathrm{k}}+\zeta_{1} \mathrm{y}_{\mathrm{n}+\mathrm{k}-1}+\ldots+\zeta_{\mathrm{k}} \mathrm{y}_{\mathrm{n}}=\mathrm{g}_{\mathrm{n}+\mathrm{k}}, \quad \mathrm{n} \geq 0 \tag{A.1}
\end{equation*}
$$

where $\left\{g_{\mathrm{n}}\right\}$ is a given sequence, independent of the $\mathrm{y}_{\mathrm{p}}$.
(i) Iff the characteristic polynomial $\zeta(z):=\sum_{j=0}^{k} \zeta_{j} z^{k-j}$ is simple von Neumann (cf. Section 2.3) then the solution of (A.1) satisfies the inequality

$$
\left|y_{n}\right| \leq C\left\{\max _{0 \leq j \leq k-1}\left|y_{j}\right|+\sum_{j=k}^{n}\left|g_{j}\right|\right\}, \quad n \geq k
$$

where $c$ is independent of $n$.
(ii) Iff $\zeta(z)$ is Schur (cf. Section 2.3) then the solution of (A.1) satisfies the inequality

$$
\left|y_{n}\right| \leq C\left\{\max _{0 \leq j \leq k-1}\left|y_{j}\right|+\max _{k \leq j \leq n}\left|g_{j}\right|\right\}, \quad n \geq k,
$$

where $C$ is independent of $n$.

PROOF. See [ 4 ].

PROOF OF THEOREM 2.2.1. Taylor expansion of $Y\left(t_{n+j}, t_{n-i}\right)$ around $\left(t_{n}, t_{n}\right)$ yields

$$
\begin{aligned}
& L_{n}[Y]=\sum_{i=0}^{k}\left\{\alpha_{i} \sum_{q=0}^{p} \frac{1}{q!} h^{q}\left(-i \frac{\partial}{\partial t}-i \frac{\partial}{\partial s}\right)^{q} Y(t, s)\right. \\
& \left.+\sum_{j=-k}^{k}\left[\beta_{i j}-\gamma_{i j} \frac{h}{\partial s}\right] \sum_{q=0}^{p} \frac{1}{q!} h^{q}\left(j \frac{\partial}{\partial t}-i \frac{\partial}{\partial s}\right)^{q} Y(t, s)\right\} \mid\left(t_{n}, t_{n}\right) \\
& +O\left(h^{P+1}\right) \text { as } h \rightarrow 0 .
\end{aligned}
$$

Writing this formula in the form

$$
{\underset{\sim}{n}}_{n}[Y]=\left.\sum_{q=0}^{p} \frac{1}{q!} h^{q}\left(D_{q} Y\right)\right|_{\left(t_{n}, t_{n}\right)}+O\left(h^{p+1}\right)
$$

and expanding the differential operator $\mathrm{D}_{\mathrm{q}}$ by the binomial theorem we find

$$
\begin{aligned}
& D_{q}=\sum_{i=0}^{k}\left\{\alpha_{i}\left(-i \frac{\partial}{\partial t}-i \frac{\partial}{\partial s}\right)^{q}+\sum_{j=-k}^{k}\left[j \beta_{i j} \frac{\partial}{\partial t}-\left(i \beta_{i j}+q \gamma_{i j}\right) \frac{\partial}{\partial s}\right]\left[j \frac{\partial}{\partial t}-i \frac{\partial}{\partial s}\right]^{q-1}\right\} \\
& =\sum_{\ell=0}^{q} \sum_{i=0}^{k}\left\{(-i)^{q} \alpha_{i}-\sum_{j=-k}^{k} j^{q-\ell}(-i)^{\ell-1}\left[i \beta_{i j}+\ell \gamma_{i j}\right]\right\}\left(\frac{q}{\ell}\right)\left(\frac{\partial}{\partial t}\right)^{q-\ell}\left(\frac{\partial}{\partial s}\right)^{\ell},
\end{aligned}
$$

where $(-i)^{\ell-1} \ell$ is assumed to be zero for $i=\ell=0$. Equating to zero all terms in the $\Sigma_{\ell=0}^{q}$ yields the order equations (2.2.3) and at the same time ${\underset{\sim}{n}}^{\mathrm{L}}(\mathrm{Y})=O\left(\mathrm{~h}^{\mathrm{P}+1}\right)$ as required in Definition 2.2.1.

PROOF OF THEOREM 2.2.2. Taylor expansion of $Y\left(t_{n+j}, t_{n-i}\right)$ around $\left(t_{n}, t_{n}\right)$ yields

$$
\begin{aligned}
Y\left(t_{n+j}, t_{n-i}\right) & =\left.\sum_{q=0}^{p} \frac{1}{q!} h^{q}\left[j \frac{\partial}{\partial t}-i \frac{\partial}{\partial s}\right]^{q} Y(t, s)\right|_{\left(t_{n}, t_{n}\right)}+ \\
& +O\left(h^{p+1}\right) \text { as } h \rightarrow 0 .
\end{aligned}
$$

In order to exploit the fact that $Y(t, t) \equiv 0$ (see definition 2.2.1), we introduce the variables $u=t+s$ and $v=t-s$ and write

$$
Y(t, s)=Y\left(\frac{u+v}{2}, \frac{u-v}{2}\right)=: Z(u, v) .
$$

The identity $\mathrm{Y}(\mathrm{t}, \mathrm{t}) \equiv 0$ implies that Z and all its derivatives with respect to $u$ vanish for $u=2 t$ and $v=0$. In the following we use the notation

$$
z^{(n, m)}:=\frac{\partial^{n} \partial^{m} z}{\partial u^{n} \partial v^{m}}\left(2 t_{n}, 0\right)
$$

By means of the binomial theorem we have
(A.2)

$$
\begin{aligned}
& Y\left(t_{n+j}, t_{n-i}\right)=\sum_{q=0}^{p} \frac{1}{q!} h^{q}\left[(j-i) \frac{\partial}{\partial u}+\left.(j+i) \frac{\partial}{\partial v} 4 q_{Z(u, v)}\right|_{\left(2 t_{n}, 0\right)}+O\left(h^{p+1}\right)\right. \\
& =\sum_{q=0}^{p} \sum_{\ell=0}^{q} \frac{1}{q!} h^{q}\left({ }_{\ell}^{q}\right)(j-i)^{q-\ell}(j+i)^{\ell} Z(q-\ell, \ell)+O\left(h^{p+1}\right) \text { as } h \rightarrow 0
\end{aligned}
$$

and
(A3) $h Y_{s}\left(t_{n+j}, t_{n-i}\right)=\sum_{q=0}^{p} \sum_{\ell=0}^{q} \frac{1}{q!} h^{q+1}\left(l_{\ell}^{q}\right)(j-i)^{q-\ell}(j+i)^{\ell}\left[z^{(q-\ell+1, \ell)}-z^{(q-\ell, \ell+1)}\right]$

$$
\begin{aligned}
&+O\left(h^{p+1}\right) \\
&=\sum_{q=0}^{p} \sum_{l=0}^{q} \frac{1}{q!} h^{q}\left(l_{\ell}^{q}\right)(j-i)^{q-\ell-1}(j+i)^{\ell-1}[q j+q i-2 \ell j] z(q-\ell, \ell) \\
&+O\left(h^{p+1}\right) \text { as } h \rightarrow 0
\end{aligned}
$$

Substitution of (A.2) and (A.3) into ${\underset{\sim}{n}}_{n}^{[Y]}$ and using $Z^{(q, 0)}=0$ yields

$$
\cdot{\underset{\sim}{n}}^{\mathrm{L}}[\mathrm{Y}]=\sum_{\mathrm{q}=1}^{\mathrm{p}} \frac{1}{\mathrm{q}!} \mathrm{h}^{\mathrm{q}} \sum_{\ell=1}^{\mathrm{q}}\left(\mathrm{q}^{\mathrm{q}}\right) \mathrm{B}_{\mathrm{q} \ell^{Z}}^{(\mathrm{q}-\ell, \ell)}+O\left(\mathrm{~h}^{\mathrm{p}+1}\right) \text { as } \mathrm{h} \rightarrow 0
$$

where $\mathrm{B}_{\mathrm{q} \ell}$ is defined in (2.2.4). This proves the theorem.
PROOF OF THEOREM 2.3.1.
PROOF. Taylor expansion in a fixed point $t=t_{n}$ yields, respectively,

$$
\begin{aligned}
& y\left(t_{n-i}\right)=\sum_{q=0}^{m} \frac{1}{q!}\left(-i h \frac{d}{d t}\right)^{q} y\left(t_{n}\right)+O\left(h^{m+1}\right), \\
& Y_{n-i}\left(t_{n+j}\right)=Y\left(t_{n+j}, t_{n-i}\right)-E_{n-i}\left(h ; t_{n+j}\right) \\
& =\sum_{q=0}^{m} \frac{1}{q!} h^{q}\left(j \frac{\partial}{\partial t}-i \frac{\partial}{\partial s}\right) q_{Y}\left(t_{n}, t_{n}\right)+O\left(h^{r}+h^{m+1}\right) \\
& =\sum_{q=0}^{m} \frac{1}{q!} h^{q} \sum_{l=0}^{q} j^{q-\ell}(-i)^{\ell}\left(\frac{q}{l}\right) \frac{\partial^{q-\ell-\partial} l_{Y}}{\partial t^{q-\ell} \partial s}\left(t_{n}, t_{n}\right) \\
& +O\left(h^{r}+h^{m+1}\right) \\
& K_{n-i}\left(t_{n+j}\right)=K\left(t_{n+j}, t_{n-i}, y\left(t_{n-i}\right)\right)=\frac{\partial}{\partial s} Y\left(t_{n+j}, t_{n-i}\right) \\
& =\sum_{q=0}^{m} \frac{1}{q!} h^{q-1} \sum_{\ell=0}^{q} j^{q-\ell}(-i)^{\ell-1}\left(l^{q}\right) \ell \frac{\partial^{q-\ell} \partial^{\ell}{ }_{Y}}{\partial t^{q-\ell} \partial l^{\ell}}\left(t_{n}, t_{n}\right) \\
& +O\left(h^{m}\right) .
\end{aligned}
$$

From these expansions it is immediate that the VLM formula (2.1.4) satisfies the relation

$$
\begin{align*}
& \sum_{i=0}^{k}\left\{\alpha_{i} \theta y\left(t_{n-i}\right)+\sum_{j=-k}^{k}\left[\beta_{i j} Y_{n-i}\left(t_{n+j}\right)-h \gamma_{i j} K_{n-i}\left(t_{n+j}\right)\right]\right\}  \tag{A.4}\\
& =\sum_{q=0}^{m} h q\left\{\frac{\theta A}{q!} \frac{d^{q} y}{d t}\left(t_{n}\right)+\sum_{\ell=0}^{q}\left(C_{q \ell}-\frac{A q}{l!(q-\ell)!}\right) \frac{\partial q}{\partial t}{ }^{q-\ell} \frac{\partial}{\partial s} Y\left(t_{n}, t_{n}\right)\right\} \\
& +O\left(h^{r}+h^{m+1}\right)
\end{align*}
$$

where $A_{q}$ and $C_{q \ell}$ are defined by (2.3.2) and (2.2.3), respectively. Under the conditions of the theorem it is easily verified that this equation leads to (2.3.3). Furthermore, (2.3.3) is obviously the m-times differentiated form of equation (1.1).

PROOF OF THEOREM 2.3.2. Let $Y(t, s)$ be given by ( 1.6 ) where $y(t)$ is the exact solution of (1.1), then we may write for $n \geq k$

$$
\begin{aligned}
{\underset{\sim}{n}}_{n}(Y) \equiv & L_{n}(Y)-\sum_{i=0}^{k}\left[\alpha_{i} y_{n-i}+\sum_{j=-k}^{k}\left(\beta_{i j} Y_{n-i}\left(t_{n+j}\right)-h \gamma_{i j} K_{n-i}\left(t_{n+j}\right)\right)\right] \\
= & \sum_{i=0}^{k}\left\{\alpha_{i} \varepsilon_{n-i}+\sum_{j=-k}^{k}\left[\beta_{i j}\left(Y\left(t_{n+j}, t_{n-i}\right)-Y_{n-i}\left(t_{n+j}\right)\right)\right.\right. \\
& \left.\left.-h \gamma_{i j}\left(K\left(t_{n+j}, t_{n-i}, y\left(t_{n-i}\right)\right)-K_{n-i}\left(t_{n+j}\right)\right)\right]\right\}
\end{aligned}
$$

Substitution of the functions $Y(t, s)$ and $Y_{n}(t)$ and using (2.1.3) and (2.3.6b) leads to

$$
\begin{align*}
L_{n}(Y)= & \sum_{i=0}^{k}\left\{\alpha_{i} \varepsilon_{n-i}+\sum_{j=-k}^{k}\left[\beta _ { i j } \left(h \sum_{\ell=0}^{n-i} w_{n-i, l} \Delta K\left(t_{n+j}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right)\right.\right.\right.  \tag{A.5}\\
& \left.\left.\left.+E_{n-i}\left(h ; t_{n+j}\right)\right)-h \gamma_{i j} \Delta K\left(t_{n+j}, t_{n-i}, y\left(t_{n-i}\right), y y_{n-i}\right)\right]\right\} .
\end{align*}
$$

Thus, we have found for the errors $\varepsilon_{\mathrm{n}}$ the relation

$$
\begin{align*}
& \sum_{i=0}^{k} \alpha_{i} \varepsilon_{n-i}=v_{n}, \quad n \geq k^{*} \text {, where }  \tag{A.6}\\
& v_{n}={\underset{\sim}{n}}(Y)-\sum_{i=0}^{k} \sum_{j=-k}^{k}\left[h \beta_{i j} \sum_{\ell=0}^{n} w_{n-i}, \ell^{\Delta K\left(t_{n+j}\right.}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right) \\
& \left.+\beta_{i j} E_{n-i}\left(h ; t_{n+j}\right)-h \gamma_{i j} \Delta K\left(t_{n+j}, t_{n-i}, y\left(t_{n-i}\right), y_{n-i}\right)\right] .
\end{align*}
$$

We now proceed with the two cases (a) and (b) separately.
(a) $\quad \alpha(z) \equiv \alpha_{0} z^{k}, \quad \alpha_{0} \neq 0$.

We want to apply the discrete Gronwall inequality stated in Lemma A. 1 in order to derive an upper bound for the solution of this linear difference equation, and therefore we need an upper bound.for $\left|v_{n}\right|$. A straightforward calculation yields
(A.7)

$$
\begin{aligned}
\left|v_{n}\right| & \leq T(h)+\sum_{i=0}^{k} \sum_{j=-k}^{k}\left[b w L_{1} h \sum_{\ell=0}^{n}\left|\varepsilon_{\ell}\right|+c L_{1} h\left|\varepsilon_{n-i}\right|+b E(h)\right] \\
& \leq C_{0} h \sum_{\ell=0}^{n}\left|\varepsilon_{\ell}\right|+C_{1} E(h)+T(h),
\end{aligned}
$$

where $C_{0}$ and $C_{1}$ are constants independent of $h$ and $n$ (in the following all constants $C_{j}$ will be independent of $h$ and $n$ ). From (A.6) it follows that

$$
\left|\alpha_{0}\right|\left|\varepsilon_{n}\right| \leq C_{0} h \sum_{\ell=0}^{n}\left|\varepsilon_{\ell}\right|+C_{1} E(h)+T(h)
$$

so that for $h$ sufficiently small

$$
\begin{aligned}
\left|\varepsilon_{n}\right| & \leq \frac{1}{\alpha_{0} \mid-C_{0} h}\left[C_{0} h \sum_{\ell=0}^{n-1}\left|\varepsilon_{\ell}\right|+C_{1} E(h)+T(h)\right] \\
& \leq C_{2} h \sum_{\ell=0}^{n-1}\left|\varepsilon_{\ell}\right|+C_{3}[E(h)+T(h)] .
\end{aligned}
$$

Application of Lemma A. 1 (with $z=k^{*} \delta(h)$ ) yields

$$
\begin{array}{r}
\left|\varepsilon_{n}\right| \leq\left(1+C_{2} h\right)^{n-k^{*}}\left(k^{*}{ }_{h C_{2}} \delta(h)+C_{3}[E(h)+T(h)]\right), \\
n=k^{*}, \ldots, N .
\end{array}
$$

Since $n h \leq T-t_{0}$, part (a) of the theorem is immediate.
(b) $\quad \alpha(z)$ is simple von Neumann, $\beta(z) \equiv 0$.

Instead of directly applying Lemma A. 1 to the inequality (obtained from (A.6))

$$
\sum_{i=0}^{k}\left|\alpha_{i}\right|\left|\varepsilon_{n-i}\right| \leq\left|v_{n}\right|
$$

we first apply Lemma A. 2 (i) to obtain the "sharper" inequality

$$
\begin{equation*}
\left|\varepsilon_{n}\right| \leq C_{0}\left[\delta(h)+\sum_{j=k}^{n}\left|v_{j}\right|\right], \quad n \geq k^{*} . \tag{A.8}
\end{equation*}
$$

Unfortunately, if we use the upper bound (A.7) for $\left|v_{j}\right|$ and then apply Lemma A.1, we cannot prove convergence. However, by using the property $\beta(z) \equiv 0$, that is $\beta_{i}=\sum_{j=-k}^{k} \beta_{i j}=0$, a sharper upper bound than (A.7) can be derived. To that end we write

$$
\begin{aligned}
& \left|\sum_{j=-k}^{k} \beta_{i j} \Delta K\left(t_{n+j}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right)\right|=\sum_{j=-k}^{k} \beta_{i j}\left[\Delta K\left(t_{n}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right)\right. \\
& \left.\quad+\Delta K\left(t_{n+j}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right)-\Delta K\left(t_{n}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right)\right] \mid \\
& \quad \leq b L h \sum_{j=-k}^{k}\left|j \varepsilon_{\ell}\right|
\end{aligned}
$$

and, similarly,

$$
\left|\sum_{j=-k}^{k} \beta_{i j} E_{n-i}\left(h ; t_{n+j}\right)\right| \leq b \sum_{j=-k}^{k} \Delta E(h) .
$$

In this way we obtain instead of (A.7) the upper bound

$$
\begin{align*}
\left|v_{n}\right| & \leq T_{n}(h)+\sum_{i=0}^{k} \sum_{j=-k}^{k}\left[b w L|j| h^{2} \sum_{\ell=0}^{n}\left|\varepsilon_{\ell}\right|+c L_{1} h\left|\varepsilon_{n-i}\right|+b \Delta E(h)\right]  \tag{A.9}\\
& \leq c_{1} h \sum_{i=0}^{k}\left[\left|\varepsilon_{n-i}\right|+h \sum_{\ell=0}^{n}\left|\varepsilon_{\ell}\right|\right]+C_{2} \Delta E(h)+T(h) .
\end{align*}
$$

Substitution into (A.8) yields the inequality

$$
\left|\varepsilon_{n}\right| \leq C_{3}\left\{\delta(h)+h \sum_{j=k}^{n}\left[\sum_{i=0}^{k}\left|\varepsilon_{j-i}\right|+h \sum_{\ell=0}^{j}\left|\varepsilon_{\ell}\right|+h^{-1} \Delta E(h)+h^{-1} T(h)\right]\right\}
$$

It is easily verified that

$$
\sum_{j=k}^{n} \sum_{i=0}^{k}\left|\varepsilon_{j-i}\right| \leq(k+1) \sum_{j=0}^{n}\left|\varepsilon_{j}\right| .
$$

Hence,

$$
\left|\varepsilon_{n}\right| \leq c_{4}\left\{\delta(h)+h\left[(1+n h) \sum_{l=0}^{n}\left|\varepsilon_{\ell}\right|+n h^{-1} \Delta E(h)+n h^{-1} T(h)\right]\right\} .
$$

Since $n h \leq T-t_{0}$ we find for $h$ sufficiently small

$$
\left|\varepsilon_{n}\right| \leq C_{5} h \sum_{\ell=0}^{n-1}\left|\varepsilon_{\ell}\right|+C_{6} h^{-1}[h \delta(h)+\Delta E(h)+T(h)] .
$$

Finally, by applying Lemma A.l we arrive at the estimate

$$
\left|\varepsilon_{n}\right| \leq\left(1+C_{5} h\right)^{n-k^{*}}\left(k^{*}{ }_{h C_{5}} \delta(h)+C_{6} h^{-1}[h \delta(h)+\Delta E(h)+T(h)]\right)
$$

from which part (b) of the theorem follows.

PROOF OF THEOREM 2.3.4. Following the first lines of the proof of Theorem 2.3 .2 we obtain the following relation, analogous to (A.5), where $K_{r s}:=K\left(t_{r}, t_{s}\right)$
(A10) $\sum_{i=0}^{k} \sum_{j=-k}^{k} \gamma_{i j} K_{n+j, n-i} \varepsilon_{n-i}=\sum_{i=0}^{k} \sum_{j=-k}^{k} \beta_{i j}\left[\sum_{\ell=0}^{n} w_{\left.n-i, \ell^{K}{ }_{n+j}, \ell^{\varepsilon_{j}}{ }^{n} h^{-1} E_{n-i}\left(h ; t_{n+j}\right)\right]}\right.$

$$
-h^{-1}{\underset{\sim}{n}}^{n}(Y), \quad n \geq k^{*}
$$

Now we write $K_{n+j, n-i}=K_{n n}+\left(K_{n+j, n-i}-K_{n n}\right)$ and $K_{n+j, \ell}=K_{n \ell}+\left(K_{n+j, \ell}-K_{n \ell}\right)$ and rewrite (A.10) to obtain
(A.11)

$$
\sum_{i=0}^{k} \gamma_{i} \varepsilon_{n-i}=v_{n}, \quad n \geq k^{*}
$$

where

$$
\begin{aligned}
K_{n n} v_{n} & =h \sum_{i, j} \gamma_{i j}\left(\frac{K_{n n}-K_{n+j, n-i}}{h}\right) \varepsilon_{n-i}+\sum_{i, j} \beta_{i j} \sum_{\ell} w_{n-i, \ell^{K} n \ell \ell^{\varepsilon_{l}}}+ \\
& +h \sum_{i, j} \beta_{i j} \sum_{\ell} w_{n-i, \ell}\left(\frac{\left.K_{n+j, \ell^{-K} n \ell}^{h}\right)}{n} \varepsilon_{\ell}+\right. \\
& +h^{-1} \sum_{i, j} \beta_{i j} E_{n-i}\left(h ; t_{n+j}\right)-h^{-1}{\underset{\sim}{L}}_{n}(Y) .
\end{aligned}
$$

Since $\gamma(z)$ is Schur, we may apply Lemma A. 2 (ii) to (A.11) and find

$$
\text { (A.12) } \quad\left|\varepsilon_{n}\right| \leq c\left\{\delta(h)+\max _{k \leq j \leq n}\left|v_{j}\right|\right\}, n \geq k^{*} \text {. }
$$

where $C$ (and all subsequent $C_{i}$ ) is independent of $h$ and $n$. So we have to find bounds on $\left|v_{j}\right|$. Using the conditions of the theorem, we find

$$
\begin{aligned}
&\left|v_{r}\right| \leq C_{1} h \sum_{i, j}\left|\gamma_{i j}\right|(j+i)\left|\varepsilon_{r-i}\right|+\left|\sum_{i=0}^{k} \beta_{i} \sum_{\ell=0}^{r}{ }^{w} r-i, \ell^{K} l^{\varepsilon} \ell\right| \\
&+C_{2} h w \sum_{i, j} j\left|\beta_{i j}\right| \sum_{\ell=0}^{r}\left|\varepsilon_{\ell}\right|+h^{-1}\left|\sum_{i, j} \beta_{i j}{ }^{E}{ }_{r-i}\left(h ; t_{r+j}\right)\right|+h^{-1}\left|L_{r}(Y)\right|, \\
& r \geq k^{*} .
\end{aligned}
$$

Now we use the condition $\beta(z) \equiv 0$, i.e., $\beta_{i}=0$, and (2.3.6a) to obtain (cf. the derivation of (A.9) in the proof of Theorem 2.3.2)

$$
\begin{aligned}
\left|v_{r}\right| & \leq c_{3}\left\{h \sum_{i=0}^{k}\left|\varepsilon_{r-i}\right|+h \sum_{\ell=0}^{r}\left|\varepsilon_{\ell}\right|+h^{-1} \Delta E(h)\right\}+h^{-1} T(h), \quad r \geq k^{*}, \\
& \leq c_{4}\left\{h \sum_{\ell=0}^{r}\left|\varepsilon_{\ell}\right|+h^{-1} \Delta E(h)\right\}+h^{-1} T(h) .
\end{aligned}
$$

Substituting this into (A.12) we find, for $h$ sufficiently small,

$$
\left|\varepsilon_{n}\right| \leq C_{5}\left\{\delta(h)+h^{-1} \Delta E(h)+h^{-1} T(h)+h \sum_{\ell=0}^{n-1}\left|\varepsilon_{\ell}\right|\right\}
$$

and application of Lemma A. 1 yields the result of the theorem.

PROOF OF THEOREM 3.3.1. Proceeding as in the proof of Theorem 2.3.2 we derive the relations
(A.13)

$$
{\underset{\sim}{n}}^{*}[y]=\sum_{i=0}^{k}\left[\alpha_{i}^{*} \varepsilon_{n-i}-h \gamma_{i}^{*} \Delta f_{n-i}\right]
$$

$$
\begin{aligned}
{\underset{\sim}{n}}_{n}[y] & =\sum_{i=0}^{k}\left\{\alpha_{i} \eta_{n-i}+\sum_{j=-k}^{k}\left[\beta _ { i j } \left(h \sum_{\ell=0}^{n-i} w_{n-i}, \ell \Delta K\left(t_{n+j}, t_{\ell}, y\left(t_{\ell}\right), y_{\ell}\right)\right.\right.\right. \\
& \left.\left.\left.+E_{n-i}\left(h ; t_{n+j}\right)\right)-h \gamma_{i j} \Delta K\left(t_{n+j}, t_{n-i}, y\left(t_{n-i}\right), y_{n-i}\right)\right]\right\} .
\end{aligned}
$$

The first relation is written as (cf. (A.6))
(A.14)

$$
\sum_{i=0}^{k} \alpha_{i}^{*} \varepsilon_{n-i}=v_{n}^{*}
$$

where $v_{n}^{*}$ satisfies the inequality (using (1.3') and (1.3'))

$$
\begin{aligned}
\left|v_{n}^{*}\right| & :=\left|{\underset{\sim}{n}}_{*}^{*}[y]+h \sum_{i=0}^{k} \gamma_{i}^{*} \Delta f_{n-i}\right| \\
& \leq F_{n}^{*}(h)+h \sum_{i=0}^{k}\left|\gamma_{i}^{*}\right|\left[L_{1}\left|\varepsilon_{\ddot{n}-i}\right|+L_{2}\left|\eta_{n-i}\right|\right] .
\end{aligned}
$$

Application of Lemma A. 2 (i) yields (because $\alpha^{*}(z)$ is simple von Neumann)
(A.15)

$$
\left|\varepsilon_{n}\right| \leq C_{0}\left[h \sum_{j=0}^{n}\left[\left|\varepsilon_{j}\right|+\left|\eta_{j}\right|\right]+\delta(h)+\sum_{j=k}^{n} T_{j}^{\star}(h)\right]
$$

where $C_{0}$ is some constant independent of $n$ and $h$.
For $\eta_{n}$ we derive from the second relation in (A.13)
(A.16)

$$
\sum_{i=0}^{k} \alpha_{i} \eta_{n-i}=v_{n}
$$

where $v_{n}$ is defined as in (A.6).
(a) In the case where $\alpha(z)=\alpha_{0} z^{k}$ we have from (A.7):

$$
\left|n_{n}\right| \leq C_{1}\left[E_{n}(h)+h \sum_{\ell=0}^{n}\left|\varepsilon_{\ell}\right|+T_{n}(h)\right], \quad n \geq k^{*}
$$

for some constant $C_{1}$. Substitution into (A.15) yields

$$
\begin{aligned}
\left|\varepsilon_{n}\right| \leq & C_{2}\left\{h \sum _ { j = k } ^ { n } \left[\left|\varepsilon_{j}\right|+h \sum_{\ell=0}^{j}\left|\varepsilon_{l}\right|+E_{j}(h)+\right.\right. \\
& \left.\left.+T_{j}(h)+h^{-1} T_{j}^{*}(h)\right]+\delta(h)+h \delta^{*}(h)\right\} \\
\leq & C_{3}\left\{h \sum_{j=0}^{n}\left|\varepsilon_{j}\right|+E_{n}(h)+T_{n}(h)+h^{-1} T_{n}^{*}(h)+\delta(h)+h \delta^{*}(h)\right\}
\end{aligned}
$$

where we have used that $n h \leq T-t_{0}$. From Lemma A. 1 , part (a) of the theorem easily follows.
(b) Since $\alpha(z)$ is simple von Neumann, we apply Lemma A. 2 (i) to (A.16) and use (A.9) (since $\beta(z) \equiv 0$ ) to find

$$
\begin{aligned}
\left|\eta_{n}\right| & \leq C_{4}\left\{\delta^{*}(h)+\sum_{j=k}^{n}\left[\sum_{i=0}^{k}\left(h\left|\varepsilon_{j-i}\right|+h^{2} \sum_{\ell=0}^{j}\left|\varepsilon_{\ell}\right|\right)+\Delta E_{j}(h)+T_{j}(h)\right]\right\} \\
& \leq C_{5}\left\{h \sum_{j=0}^{n}\left|\varepsilon_{j}\right|+\delta^{*}(h)+h^{-1} \Delta E_{n}(h)+h^{-1} T_{n}(h)\right\}
\end{aligned}
$$

Substitution into (A.15) and applying Lemma A. 1 leads to part (b) of the theorem.


[^0]:    *) This report will be submitted for publication elsewhere.

