

A Novel Physical Model and Computational Method for Non-Isentropic, Compressible Two-Fluid Flow

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Summary. This paper describes a five-equation model for compressible two-fluid flow, based on physical flow equations only. The model is conservative and pressure-oscillation free. Equations for continuous flow and jump conditions for discontinuities are given, as well as a discretisation of the equations and an adaptation of the HLL Riemann solver to two-fluid flow. Numerical tests in 1D and 2D show the accuracy of the method.

1 Introduction

Interface-capturing methods for compressible two-fluid flows are based on mixture-fluid models. The interface between the fluids appears as a numerically smeared transition from fluid 1 to fluid 2. Many conservative formulations of such models produce large pressure errors. This problem can be solved by using locally non-conservative methods [1] or by solving the full two-phase flow model [2].

Here, an intermediate approach is presented: a two-fluid method that is fully conservative and pressure-oscillation free. It is an extension of the work by Van Brummelen and Koren [3], it will be described in detail in a future paper. A similar method is derived in a different way by Kapila et al. [4].

The present method has two major advantages. First, the conservative formulation gives good capturing of shocks and interfaces, also for problems with very strong shocks. And second, the model strongly resembles a single-fluid model: it does not require a complex interface-tracking algorithm. It can thus be solved with existing techniques, even on complex, irregular grids.

2 Flow model

The physical model used here for two-fluid flow is based on a mixture model. However, the fluids are not fully mixed: the ‘mixture’ may be thought to consist of very small bits of the two pure fluids, arranged in an arbitrary pattern. Each fluid has its own pure-fluid equation of state and the fluids interact only by exerting forces on each other. In the model, the pressure and the velocity of the fluids are equal, but each fluid has its own density. The

volume fraction of fluid 1, α , is used to denote the relative amounts of the two fluids. Thus, in 1D, we have five independent state variables (p , u , ρ_1 , ρ_2 and α), so we need five differential equations to solve the flow.

The two-fluid bulk flow satisfies the standard Euler equations:

$$(\rho)_t + (\rho u)_x = 0, \quad (1a)$$

$$(\rho u)_t + (\rho u^2 + p)_x = 0, \quad (1b)$$

$$(\rho E)_t + (\rho E u + p u)_x = 0. \quad (1c)$$

In these equations, the bulk quantities are $\rho = \alpha\rho_1 + (1 - \alpha)\rho_2$ and $\rho E = \alpha\rho_1 E_1 + (1 - \alpha)\rho_2 E_2$, with the total energy for each fluid $j = 1, 2$ defined as $E_j = e_j + \frac{1}{2}u^2$. Here e_j is the internal energy of fluid j .

Two more flow equations are needed to close the system. The first one is the conservation of mass for fluid 1:

$$(\rho_1 \alpha)_t + (\rho_1 u \alpha)_x = 0. \quad (2a)$$

Together with equation (1a), this equation gives mass conservation for both fluids. For the last equation, the energy balance of fluid 1 is used. As the fluids exert forces on each other, they exchange energy, which appears as a source term in the equation:

$$(\rho_1 E_1 \alpha)_t + (\rho_1 E_1 u \alpha + p u \alpha)_x = S. \quad (2b)$$

An expression for this source term is derived in the next section.

To close the system, equations of state (EOS) are needed for the two fluids. A possible EOS is the ideal gas law,

$$p = (\gamma_1 - 1)\rho_1 e_1 = (\gamma_2 - 1)\rho_2 e_2, \quad (3)$$

with constant γ 's. For this equation, it is easy to compute the primitive variables p and α from the total energies.

3 The source term

3.1 Derivation of the source term

The source term S in equation (2b) models the exchange of energy between fluid 2 and fluid 1. Euler flow has no heat conduction, so the only energy exchanged is the work done by the force between the fluids. This force is found from a momentum analysis.

Consider a fluid element in a smooth 1D flow (Fig. 1). The element contains fluid 1 and fluid 2 (the interface is drawn schematically). The force on the entire fluid element is $p(x) - p(x + dx)$ and its bulk mass is ρdx . The force on fluid 1 in the element is $(p\alpha)(x) - (p\alpha)(x + dx) + S_M dx$. Its mass is

$\rho_1 \alpha dx$ and its acceleration is *equal* to the acceleration of the entire element (because both fluids have the same velocity). Therefore

$$\frac{p(x) - p(x + dx)}{\rho dx} = \frac{(p\alpha)(x) - (p\alpha)(x + dx) + S_M dx}{\rho_1 \alpha dx} .$$

The force S_M follows from this expression (using the mass fraction $\beta = \frac{\rho_1 \alpha}{\rho}$):

$$S_M = p\alpha_x + (\alpha - \beta) p_x . \quad (4)$$

The energy source term S is the work done by the force S_M :

$$S = u S_M = p u \alpha_x + (\alpha - \beta) u p_x . \quad (5)$$

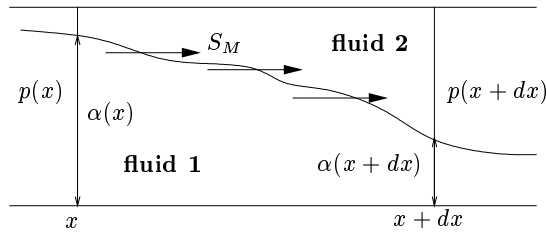


Fig. 1. Two-fluid element in smooth 1D flow.

3.2 Characteristic analysis of the system for ideal gas

The source term (5) is valid for any EOS. Substitution of the ideal gas law (3) allows a characteristic analysis of the flow equations. This results in five wave speeds,

$$\lambda_1 = u - c, \lambda_{2,3,4} = u, \lambda_5 = u + c, \quad \text{with } c = \sqrt{(\alpha\gamma_1 + (1 - \alpha)\gamma_2)p/\rho} . \quad (6)$$

This combination of wave speeds is physically correct. It can be proved that (5) is the only possible source term that gives such a combination.

3.3 Source term in discontinuities

To allow weak solutions with discontinuities of the two-fluid flow equations, we need a proper definition of the flow across a discontinuity. The first four equations, (1a) – (1c) and (2a), satisfy the Rankine-Hugoniot condition $[\mathbf{f}] = c_s[\mathbf{q}]$, with c_s the speed of the discontinuity. For the fifth equation, (2b), this condition becomes

$$[\mathbf{f}] = c_s[\mathbf{q}] + \int_{x_L}^{x_R} S dx , \quad (7)$$

where x_L and x_R denote the coordinates of the left and right side of the discontinuity. The integral must be evaluated across the discontinuity, which is impossible. However, if we assume that the discontinuity is the inviscid limit of a viscous layer and thus has a continuous internal structure (the precise shape is unimportant), then we can write the state variables as continuous functions of p and integrate the source term:

$$\int_{x_L}^{x_R} S dx = [pu\alpha] + \frac{1}{2}\beta_L\rho_L(u_L - c_s)[u]^2 + \frac{1}{\rho_L(u_L - c_s)} \int_{p_L}^{p_R} p\alpha dp. \quad (8)$$

A derivation of this expression will be given in a future paper. The last integral can be evaluated exactly, but it requires an EOS. So there is a unique jump condition for the present two-fluid model, but, unlike the single-fluid jump condition, it depends on the material properties of the fluids.

4 Numerical method

4.1 Second-order accurate discretisation

The flow equations are discretised with a second-order accurate finite-volume scheme. Fluxes are computed with an improved version of Linde's three-wave HLL approximate Riemann solver [5], combined with a limited reconstruction of the cell interface states. The limiter is applied to the primitive variables ρ , u , p , α and β . Time stepping is done with a two-step scheme (see [7]).

4.2 Numerical source term

A discretisation of the source term is needed in two places. First, an approximation of the source term in a discontinuity is needed to compute the HLL flux. The HLL solver models a Riemann problem with three discontinuous waves. The easiest way to incorporate the source term in these waves is to compute only one approximate solution of (8), using the left and right cell interface state, and to divide this source term proportionally over the three waves. This procedure causes some small inaccuracies, but it is fast and straightforward.

Secondly, the source term for the time integration is computed. It consists of two parts:

- i) sources in the discontinuities at the cell faces, that are summed over all HLL waves on interfaces $i - \frac{1}{2}$ and $i + \frac{1}{2}$, that actually run into cell i ,
- ii) sources in the continuous flow in the cell, which are integrated over the piecewise linear approximations to the primitive variables, that follow from the use of the limiter.

These two sources are summed per cell.

5 Numerical results

5.1 Shock-tube test

The method is tested first on a 1D Riemann problem for ideal gases: a two-fluid variant of Sod's problem, with a ten times higher left pressure and density, giving it a pressure ratio of 100:1. Figure 2 shows that the discontinuities (shock and two-fluid interface) are in the proper locations. The pressure is constant over the contact discontinuity and the volume fraction is constant over both the shock and the expansion fan. A convergence study for this problem shows that the L^1 -errors in ρ , u and p converge approximately with the power 0.96 of the mesh width. The volume fraction converges with the power 0.78 of the mesh width. This rate of convergence is comparable to that for similar single-fluid solutions.

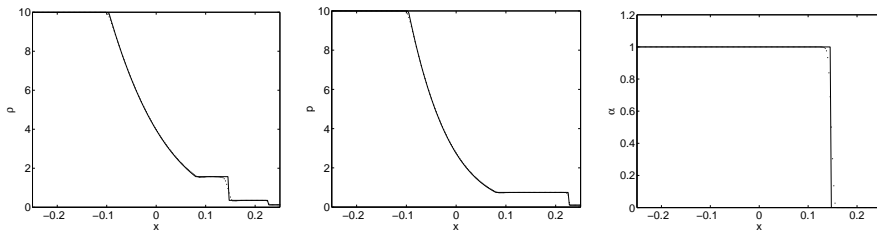


Fig. 2. High-pressure, two-fluid Sod problem. $(\rho, u, p)_L = (10, 0, 10)$, $(\rho, u, p)_R = (0.125, 0, 0.1)$, $\gamma_L = 1.4$ and $\gamma_R = 1.6$. The grid has 200 cells, 160 time steps, $\Delta t/\Delta x = 0.2$ (CFL = $|\lambda_{\max}|\Delta t/\Delta x = 0.56$). Solid lines: exact solution.

5.2 Shock hitting helium bubble in air

This 2D test case has been taken from literature [6]. It consists of a cylindrical helium bubble in air, which is hit by an incoming shock wave. The problem is solved on a grid of 200×400 cells, with $\Delta t = 1.25 \times 10^{-5}$. Figure 3 shows the solution at two times. The (half) bubble is visible between $x = -0.025$ and $x = 0.025$. The incident shock, coming from the right, can be seen in the air above the bubble, the curved shock in the bubble runs ahead of this shock. The rightmost wave is an expansion wave, reflected into the air behind the shock. At the later time, a complicated λ -shock structure has developed above the bubble. Figure 4 shows the pressure and the volume fraction for this time. Of the waves appearing in the density plot, the shock waves and expansions are visible in the pressure plot only and the interface in the volume fraction plot only, as it should be. The pressure is continuous over the interface.

The speeds of the shocks and the interface at the centerline ($y = 0$) have been compared with results from Quirk and Karni [6] (obtained on a very fine, adapted grid). The difference is between 0.7% and 2.2%.

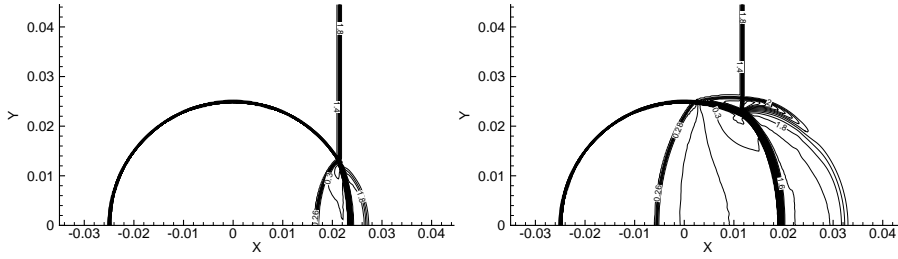


Fig. 3. Shock hitting helium bubble, density at $t = 2.74 \times 10^{-3}$ and $t = 10.74 \times 10^{-3}$.

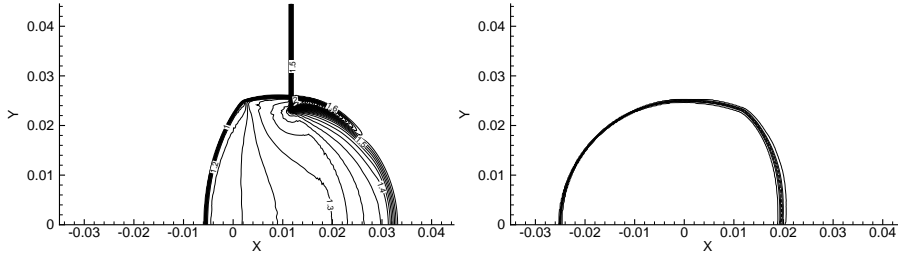


Fig. 4. Shock hitting helium bubble, pressure (left) and volume fraction (right) at $t = 10.74 \times 10^{-3}$.

6 Conclusions

A model for compressible two-fluid flow is proposed, that is conservative and pressure-oscillation free. 1D tests show that the model resolves contact discontinuities without creating pressure errors and that it accurately handles problems with strong shocks. A 2D test shows that the method properly resolves curved shocks and interfaces too.

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References

1. R. Abgrall, S. Karni: *J. Comp. Phys.* **169**, 594 (2001)
2. R. Abgrall, R. Saurel: *J. Comp. Phys.* **186**, 361 (2003)
3. E.H. van Brummelen, B. Koren: *J. Comp. Phys.* **185**, 289 (2002)
4. A.K. Kapila, J.B. Bdzil, R. Menikoff, S.F. Son, D.S. Steward: *Phys. Fluids* **13**, 3002 (2001)
5. T. Linde: *Int. J. Num. Meth. Fluids* **40**, 391 (2002)
6. J.J. Quirk, S. Karni: On the dynamics of a shock-bubble interaction. ICASE Report 94-75, NASA Langley Research Center, Hampton, VA (1994)
7. J. Wackers, B. Koren: A simple and efficient space-time adaptive grid technique for unsteady compressible flows. AIAA Paper 2003-3825 (2003)