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A NEW PRE-LOADED MEMBRANE GEOMETRIC STIFFNESS MATRIX WITH FULL RIGID BODY CAPABILITIES

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NOTATION

a	membrane element length
b	membrane element width
[DFC]	directed force correction matrix
[DFC Total]	matrix of all corrections $([DFC] + [MCOR])$
E	Young's modulus (Modulus of elasticity)
h	element thickness
[Ke]	elastic stiffness matrix
[Kg]	geometric stiffness matrix
K	i, j element of stiffness matrix
[<i>M</i>]	mass matrix
[MCOR]	fixed end moment correction
N	distributed tensile load $(N_x = N \text{ for example},$
	where $N_y = N_{xy} = 0.$)
Nx	distributed tensile load in x-direction
N.	distributed tensile load in y-direction
$egin{aligned} & N_y \ & N_{xy} \ & lpha_i \end{aligned}$	distributed shear load on membrane
α,	element constant
β_i	element constant
ρ	mass density
v	Poisson's ratio

INTRODUCTION

Due to economic constraints associated with payload cost, space structures must consist of light-weight, and subsequently relatively flexible, components. The total stiffness of a member, in many cases, includes both the elastic stiffness, due to the material properties and member configuration, as well as geometric or initial stress stiffness due to pre-loads. Accurate prediction of the natural frequencies and mode shapes is essential for determining the adequacy of components, and for designing a controls system.

The finite element method is normally used to perform this analysis. However, a phenomenon known as "grounding" or false stiffening, occurs when the geometric stiffness matrix is used. When a pre-loaded model with free/free boundary conditions is analyzed, it behaves as if it is restrained internally. Thus, it is unable to rotate as a rigid body. Instead, erroneous "pseudo-forces" develop during rigid body motion. Free vibration analysis also yields an erroneous (non-zero) natural frequency associated with rigid body rotation. Bosela et al. [1-3] have examined the grounding phenomenon for a pre-loaded beam, and determined that it is caused by a force imbalance during rigid body rotation, and is typical of beam geometric stiffness matrices formulated by others, even those which include higher order effects. By utilizing a directed force premise, and extending to a global model the natural mode approach used by Argyris and Symeonidis [4] for developing load correction factors for non-conservative forces, a preloaded beam geometric stiffness matrix with full rigid body capabilities was developed and tested [3, 5].

This paper has the following objectives: (1) Examine the grounding phenomenon as it relates

to pre-loaded membrane elements.

(2) Develop a pre-loaded membrane element with full rigid body capabilities by following the directed force approach used successfully in developing the pre-loaded beam element.

MEMBRANE ELEMENTS

Although various membrane elements are used in commercial finite element programs, most are not readily available in explicit form. One which is stated explicitly by Yang [6] is the four-node rectangular element in Fig. 1, which has either 12 or 16 degrees of freedom (DOF). The degrees of freedom numbered as 13, 14, 15 and 16, correspond to second-order twist terms such as $\partial^2 w / \partial x \partial y$. This element is obtained by combining the shape functions for the elementary beam element in both the x and y directions.

The geometric, or initial stress stiffness matrix [Kg] for this plate element when the distributed load is in the x-direction only is

(1)

$$[K_g] = \begin{bmatrix} A_1 & A_2 & A_3 & A_4 & A_5 & A_6 & A_7 & A_8 & A_9 & A_{10} & A_{11} & A_{12} & A_{13} & A_{14} & A_{15} & A_{16} \\ B_1 & B_2 & B_3 & B_4 & B_5 & B_6 & B_7 & B_8 & B_9 & B_{10} & B_{11} & B_{12} & B_{13} & B_{14} & B_{15} & B_{16} \\ C_1 & C_2 & C_3 & C_4 & C_5 & C_6 & C_7 & C_8 & C_9 & C_{10} & C_{11} & C_{12} & C_{13} & C_{14} & C_{15} & C_{16} \\ D_1 & D_2 & D_3 & D_4 & D_5 & D_6 & D_7 & D_8 & D_9 & D_{10} & D_{11} & D_{12} & D_{13} & D_{14} & D_{15} & D_{16} \\ E_1 & E_2 & E_3 & E_4 & E_5 & E_6 & E_7 & E_8 & E_9 & E_{10} & E_{11} & E_{12} & E_{13} & E_{14} & E_{15} & E_{16} \\ F_1 & F_2 & F_3 & F_4 & F_5 & F_6 & F_7 & F_8 & F_9 & F_{10} & F_{11} & F_{12} & F_{13} & F_{14} & F_{15} & F_{16} \\ G_1 & G_2 & G_3 & G_4 & G_5 & G_6 & G_7 & G_8 & G_9 & G_{10} & G_{11} & G_{12} & G_{13} & G_{14} & G_{15} & G_{16} \\ H_1 & H_2 & H_3 & H_4 & H_5 & H_6 & H_7 & H_8 & H_9 & H_{10} & H_{11} & H_{12} & H_{13} & H_{14} & H_{15} & H_{16} \\ I_1 & I_2 & I_3 & I_4 & I_5 & I_6 & I_7 & I_8 & I_9 & I_{10} & I_{11} & I_{12} & I_{13} & I_{14} & I_{15} & I_{16} \\ J_1 & J_2 & J_3 & J_4 & J_5 & J_6 & J_7 & J_8 & J_9 & J_{10} & J_{11} & J_{12} & J_{13} & J_{14} & J_{15} & J_{16} \\ K_1 & K_2 & K_3 & K_4 & K_5 & K_6 & K_7 & K_8 & K_9 & K_{10} & K_{11} & K_{12} & K_{13} & K_{14} & K_{15} & K_{16} \\ I_1 & I_2 & I_3 & I_4 & I_5 & I_6 & I_7 & I_8 & I_9 & I_{10} & J_{11} & J_{12} & J_{13} & J_{14} & J_{15} & J_{16} \\ K_1 & M_2 & M_3 & M_4 & M_5 & M_6 & M_7 & M_8 & M_9 & M_{10} & M_{11} & M_{12} & M_{13} & M_{14} & M_{15} & M_{16} \\ N_1 & N_2 & N_3 & N_4 & N_5 & N_6 & N_7 & N_8 & N_9 & N_{10} & N_{11} & N_{12} & N_{13} & N_{14} & N_{15} & N_{16} \\ P_1 & P_2 & P_3 & P_4 & P_5 & P_6 & P_7 & P_8 & P_9 & P_{10} & P_{11} & P_{12} & P_{13} & P_{14} & P_{15} & P_{16} \\ Q_1 & Q_2 & Q_3 & Q_4 & Q_5 & Q_6 & Q_7 & Q_8 & Q_9 & Q_{10} & Q_{11} & Q_{12} & Q_{13} & Q_{14} & Q_{15} & Q_{15} \end{bmatrix}$$

 $A_{1} = B_{2} = C_{3} = D_{4} = \frac{78N_{x}b}{175a} \qquad A_{2} = B_{1} = C_{4} = D_{3} = -\frac{78N_{x}b}{175a}$ $A_{4} = B_{3} = C_{2} = D_{1} = \frac{27N_{x}b}{175a} \qquad A_{3} = B_{4} = C_{1} = D_{2} = -\frac{27N_{x}b}{175a}$ $A_{5} = A_{6} = D_{7} = D_{8} = \frac{13N_{x}b}{350} \qquad B_{5} = B_{6} = C_{7} = C_{8} = -\frac{13N_{x}b}{350}$ $A_{7} = A_{8} = D_{5} = D_{6} = \frac{9N_{x}b}{700} \qquad B_{7} = B_{8} = C_{5} = C_{6} = -\frac{9N_{x}b}{700}$ $A_{9} = B_{10} = C_{12} = D_{11} = \frac{11N_{x}b^{2}}{175a} \qquad A_{10} = B_{9} = C_{11} = D_{12} = -\frac{11N_{x}b^{2}}{175a}$ $A_{11} = B_{12} = C_{10} = D_{9} = \frac{13N_{x}b^{2}}{350a} \qquad A_{12} = B_{11} = C_{9} = D_{10} = -\frac{13N_{x}b^{2}}{350a}$ $A_{13} = A_{14} = C_{15} = C_{16} = \frac{11N_{x}b^{2}}{2100} \qquad B_{13} = B_{14} = D_{15} = D_{16} = -\frac{11N_{x}b^{2}}{2100}$ $B_{15} = B_{16} = D_{13} = D_{14} = \frac{13N_{x}b^{2}}{4200} \qquad A_{15} = A_{16} = C_{13} = C_{14} = -\frac{13N_{x}b^{2}}{4200}$ $E_{1} = F_{1} = G_{4} = H_{4} = \frac{13N_{x}b}{350} \qquad E_{2} = F_{2} = G_{3} = H_{3} = -\frac{13N_{x}b}{350}$

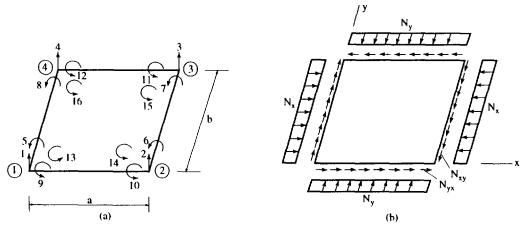


Fig. 1.(a) Four-node rectangular plate element with 16 DOF. (b) In-plane normal and shear forces on quadrilateral membrane element.

$$\begin{split} E_4 &= F_4 = G_1 = H_1 = \frac{9N_x b}{700} \qquad E_3 = F_3 = G_2 = H_2 = -\frac{9N_x b}{700} \\ E_5 &= F_6 = G_7 = H_8 = \frac{26N_x ab}{525} \qquad E_6 = F_5 = G_8 = H_7 = -\frac{13N_x ab}{1050} \\ E_7 &= F_8 = G_5 = H_6 = -\frac{3N_x ab}{700} \qquad E_8 = F_7 = G_6 = H_5 = \frac{3N_x ab}{175} \\ E_9 &= F_9 = G_{11} = H_{11} = \frac{11N_x b^2}{2100} \qquad E_{10} = F_{10} = G_{12} = H_{12} = -\frac{11N_x b^2}{2100} \\ E_{11} &= F_{11} = G_9 = H_9 = \frac{13N_x b^2}{4200} \qquad E_{12} = F_{12} = G_{10} = H_{10} = -\frac{13N_x b^2}{4200} \\ E_{13} &= F_{14} = \frac{11N_x ab^2}{1575} \qquad G_{15} = H_{16} = -\frac{11N_x ab^2}{1575} \\ G_{16} &= H_{15} = \frac{11N_x ab^2}{6300} \qquad E_{14} = F_{13} = -\frac{11N_x ab^2}{6300} \\ E_{15} &= F_{16} = \frac{13N_x ab^2}{3150} \qquad G_{13} = H_{14} = -\frac{13N_x ab^2}{12600} \\ G_{14} &= H_{13} = \frac{13N_x ab^2}{3150} \qquad E_{16} = F_{15} = -\frac{13N_x ab^2}{3150} \\ I_1 &= J_2 = K_4 = L_3 = \frac{11N_x b^2}{175a} \qquad I_2 = J_1 = K_3 = L_4 = -\frac{11N_x b^2}{350a} \\ I_4 &= J_3 = K_1 = L_2 = \frac{13N_x b^2}{350a} \qquad J_5 = J_6 = L_7 = L_8 = -\frac{11N_x b^2}{2100} \end{split}$$

$$\begin{split} I_{1} &= I_{8} = K_{5} = K_{6} = \frac{13N_{*}b^{2}}{4200} \qquad J_{7} = J_{8} = L_{5} = L_{6} = -\frac{13N_{*}b^{2}}{4200} \\ I_{8} &= J_{10} = K_{11} = L_{12} = \frac{2N_{*}b^{3}}{175a} \qquad I_{10} = K_{12} = L_{11} = J_{9} = -\frac{2N_{*}b^{3}}{175a} \\ I_{11} &= J_{12} = K_{9} = L_{10} = \frac{3N_{*}b^{3}}{350a} \qquad I_{12} = J_{11} = K_{10} = L_{9} = -\frac{3N_{*}b^{3}}{350a} \\ I_{13} &= I_{14} = L_{15} = L_{16} = M_{9} = N_{9} = Q_{12} = P_{12} = \frac{N_{*}b^{3}}{1050} \\ J_{13} &= J_{14} = K_{15} = K_{16} = M_{10} = P_{11} = Q_{11} = -\frac{N_{*}b^{3}}{1050} \\ J_{15} &= J_{16} = K_{13} = K_{14} = M_{11} = N_{11} = P_{10} = Q_{10} = \frac{N_{*}b^{3}}{1400} \\ I_{15} &= I_{16} = L_{13} = L_{14} = M_{12} = N_{12} = P_{9} = Q_{9} = -\frac{N_{*}b^{3}}{1400} \\ M_{1} &= N_{1} = P_{3} = Q_{2} = \frac{11N_{*}b^{2}}{2100} \qquad M_{2} = N_{2} = P_{4} = Q_{4} = -\frac{11N_{*}b^{2}}{2100} \\ M_{4} &= N_{4} = P_{2} = Q_{2} = \frac{13N_{*}b^{2}}{4200} \qquad M_{3} = N_{3} = P_{1} = Q_{1} = -\frac{13N_{*}b^{2}}{4200} \\ M_{5} &= N_{6} = \frac{11N_{*}ab^{2}}{1575} \qquad M_{6} = N_{5} = -\frac{11N_{*}ab^{2}}{12600} \\ M_{13} &= N_{14} = P_{15} = Q_{16} = \frac{2N_{*}ab^{3}}{1575} \qquad M_{14} = N_{13} = P_{16} = Q_{13} = -\frac{N_{*}ab^{3}}{3150} \\ M_{15} &= N_{16} = P_{13} = Q_{14} = \frac{N_{*}ab^{3}}{4200} \qquad M_{16} = N_{15} = P_{14} = Q_{13} = -\frac{N_{*}ab^{3}}{3150} \\ M_{15} &= N_{16} = P_{13} = Q_{14} = \frac{N_{*}ab^{3}}{1050} \qquad M_{16} = N_{15} = P_{14} = Q_{13} = -\frac{N_{*}ab^{3}}{3150} \\ M_{15} &= N_{16} = P_{13} = Q_{14} = \frac{N_{*}ab^{3}}{4200} \qquad M_{16} = N_{15} = P_{14} = Q_{13} = -\frac{N_{*}ab^{3}}{3150} \\ P_{5} &= Q_{6} = \frac{13N_{*}ab^{2}}{12600} \qquad N_{10} = -\frac{N_{*}b^{3}}{1050} \qquad P_{5} = Q_{6} = \frac{11N_{*}ab^{2}}{3150} \\ P_{5} &= Q_{7} = \frac{11N_{*}ab^{2}}{6300} \qquad P_{7} = Q_{8} = -\frac{11N_{*}ab^{2}}{1575}. \end{cases}$$

In order to solve for the natural frequencies, the corresponding four-node quadrilateral mass matrix developed by Yang was also used. The coefficients of the mass matrix are determined by using the

equation

$$\mathbf{M}_{ij} = \rho h \int_0^b \int_0^a f_i(x, y) f_j(x, y) \, \partial x \, \partial y.$$
 (3)

(2)

r		- (
	$a^{2}b^{2}$	C7011
	$\frac{a^2b^2}{11025} - \frac{a^2b^2}{-a^2b^2}$	14/00
	$\frac{a^2b^2}{11025} - \frac{a^2b^2}{14700}$	00061
	$\frac{a^2b^2}{11025} - \frac{a^2b^2}{14700} - \frac{a^2b^2}{18600} - \frac{a^2b^2}{14700}$	
	$\frac{13b^3}{3675}$ $\frac{-11ab^2}{58800}$ $\frac{-13ab^2}{58800}$ $\frac{-13ab^2}{11ab^2}$	00077
SYMMETRIC	$\frac{13b^2}{3675}$ $\frac{3b^3}{3675}$ $\frac{3b^3}{2450}$ $\frac{3b^3}{2450}$ $\frac{3b^3}{29400}$ $\frac{111ab^2}{22050}$ $\frac{229400}{13ab^2}$	
MINAS	$\frac{13b^2}{3675} - \frac{13b^2}{4900} - \frac{13b^2}{4900} - \frac{13b^2}{7800} - \frac{1400}{11ab^2} - \frac{11ab^2}{22050} - \frac{11ab^2}{22400} - \frac{11ab^2}{22400} - \frac{13ab^2}{5800} - \frac{13ab^2}{5800} + \frac{11ab^2}{5800} - \frac{13ab^2}{5800} + \frac{11ab^2}{5800} + \frac{11ab^2}{500} $	00000
	$\frac{13b^2}{3675} \frac{13b^2}{3675} \frac{3675}{2450} \frac{3b^2}{2800} \frac{3b^2}{9800} \frac{-9b}{9800} \frac{-9b}{4400} \frac{4900}{11ab^2} \frac{-13b^3}{58800} \frac{44100}{58800} \frac{11ab^2}{58800} \frac{50400}{50400} \frac{222050}{50400} \frac{222050}{5000} \frac{222050}{50400} \frac{222050}{5000} 222050$	
	$\frac{13a^{3}}{3675}$ $\frac{3675}{3675}$ $\frac{3675}{88200}$ $\frac{88200}{176400}$ $\frac{176400}{176400}$ $\frac{176400}{44100}$ $\frac{143ab}{44100}$ $\frac{44100}{13a^{2}b}$ $\frac{13a^{2}b}{29400}$ $\frac{13a^{2}b}{29400}$ $\frac{11a^{2}b}{29400}$ $\frac{29400}{2050}$	00077
$\frac{13a^2}{3675}$	$\begin{array}{r} 3675 \\ -13a^{3} \\ -1900 \\ 176400 \\ 176400 \\ 176400 \\ 176400 \\ 176400 \\ 188200 \\ 188200 \\ 88200 \\ 188200 \\ 88200 \\ 13a^{2}b \\ 13a^{2}b \\ 13a^{2}b \\ 13a^{2}b \\ 22050 \\ 11a^{2}b \\ 22050 \\ 11a^{2}b \end{array}$	00467
$\frac{13a^2}{3675}$ $\frac{3a^2}{2450}$	$\begin{array}{r} 2450 \\ - 9a^{3} \\ \overline{9800} \\ 98200 \\ 88200 \\ 88200 \\ 143ab \\ 44100 \\ 15820 \\ 169ab \\ 176400 \\ 169ab \\ 176400 \\ 11a^{2}b \\ 229400 \\ 11a^{2}b \\ 22050 \\ - 13a^{2}b \\ 44100 \\ 13a^{2}$	00000
$\frac{13a^2}{3675} - \frac{13a^2}{4900} - \frac{9a^2}{2}$	$\begin{array}{r} 9800\\ 3a^{3}\\ 2450\\ 121ab\\ 44100\\ 188200\\ 88200\\ 176400\\ 176400\\ 1143ab\\ 88200\\ 11a^{2}b\\ 88200\\ 11a^{2}b\\ 22050\\ 13a^{2}b\\ 58800\\ 58800\\ -11a^{2}b\\ 58800\\ -11a^{2}b\\ 58800\\ 13a^{2}b\\ 58800\\ -11a^{2}b\\ -11a^{2}b\\$	8 F
169 1225 33 <i>a</i> 4900 9800 - 169 <i>a</i> 14700	$\begin{array}{r} 14700\\ 14700\\ \hline 7350\\ 169b\\ 169b\\ \hline 169b\\ \hline 9800\\ \hline 9800\\ \hline 9800\\ \hline 98200\\ \hline 7350\\ \hline 143ab\\ \hline 88200\\ \hline 88200\\ \hline 88200\\ \hline 88200\\ \hline 88200\\ \hline 143ab\\ \hline 88200\\ \hline 88200\\ \hline 88200\\ \hline 44100\\ \hline \end{array}$	8
	7350 169a 14700 9800 14700 14700 14700 14700 14700 14700 169b 14700 169b 14700 17350 14700 17350 17350 14700 173500 173500 173500 173500 173500 173500 173500 173500 1735000 173500 173500 1735000000000000000000000000000000000000	
	4900 9800 9800 1436 1436 14700 9800 9800 9800 14100 143 <i>ab</i> 9800 98200 88200 1143 <i>ab</i> 143 <i>ab</i> 98200 143 <i>ab</i> 143 <i>bb</i> 143 <i>bb</i> 14	
	9800 33 <i>a</i> 4900 143 <i>b</i> 4900 - 39 <i>b</i> 9800 - 143 <i>ab</i> 121 <i>ab</i> 14700 121 <i>ab</i> 14700 121 <i>ab</i> 14700 14700 169 <i>ab</i> 173 <i>b</i> 88200 169 <i>ab</i> 173 <i>b</i> 88200 169 <i>ab</i> 173 <i>b</i> 88200 169 <i>ab</i> 173 <i>b</i> 88200 169 <i>ab</i> 173 <i>b</i> 88200 169 <i>ab</i> 173 <i>b</i> 88200 88200	

The mass matrix is as follows: $[M] = \rho abh$

GROUNDING OF MEMBRANE ELEMENTS

Suppose one considers a free/free model with two 12 DOF elements, and a uniform tensile load applied in the x-direction only, as shown in Fig. 2. That element should possess five rigid body modes: (1) translation in the y-direction; (2) rigid body rotation about the x-axis; (3) rotation of the edge line between nodes 1, 2 and 5; (4) rotation of the edge lines between nodes 4, 3 and 6; (5) rigid body rotation about the y-axis. It should be noted that (2) and (4) only occur because for the example problem, there is no tension in the y-direction, and we are neglecting elastic stiffness effects.

In matrix form, assembling [Kg] and multiplying a matrix of these rigid body modes yields

to rigid body rotation about the y-axis, was not present. Instead, pseudo-forces developed in degree of freedoms 1, 4, 9, 12, 13, 14, 17 and 18. The pseudo-forces in degrees of freedom 1, 4, 13 and 14 correspond to the forces required to keep the rotated model in equilibrium, which is consistent with the behavior of a pre-loaded beam element [2]. Hence, the procedure used successfully for the pre-loaded beam element [3] will be adapted.

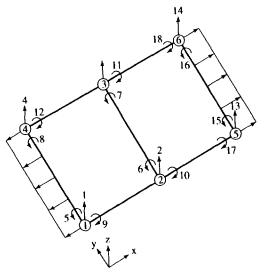
The pseudo-moments developed in degrees of freedom 9, 12, 17 and 18 did not occur with rotation of the pre-loaded beam element. It should be noted, however, that the form resembles a fixed-end moment caused by a distributed load.

	Γ 1	-1	0	0	– ר ו							
	1	-1	0	0	0		Г о	0	0	0	\underline{Nb}	
	1	1	0	0	0)				$\frac{1}{2a}$	
	1	1	0	0	-1		0	0	0	0	0	
	0	0	0	0	1		0	0	0	0	0	
		0	0	0	a 1		0	0	0	0	$-\frac{Nb}{2a}$	
	0	0	0	0	$\frac{1}{a}$		0	0	0	0	0	
	0	0	0	0	$\frac{1}{2}$		0	0	0	0	0	
					a 1		0	0	0	0	0	
	0	0	0	0	$\frac{1}{a}$		0	0	0	0	0	
	0	$\frac{2}{b}$	1	0	0		0	0	0	0	$-\frac{Nb^2}{12a}$	
[<i>Kg</i>] ·		2				*	0	0	0	0	0	. (5)
,	0	$\frac{2}{b}$	1	0	0		0	0	0	0	0	
	0	$\frac{2}{b}$	0	1	0		0	0	0	0	$\frac{Nb^2}{12a}$	
	0	$\frac{2}{\overline{b}}$	0	1	0		0	0	0	0	$\frac{Nb}{2a}$	
	1	-1 1	0 0	0 0	1		0	0	0	0	$\frac{Nb}{2a}$	
			•		1		0	0	0	0	0	
	0	0	0	0	$\frac{1}{a}$		0	0	0	0	0	
	0	0	0	0	$\frac{1}{a}$		0	0	0	0	$\frac{Nb^2}{12a}$	
	0	$\frac{2}{b}$	1	0	0		0	0	0	0	$-\frac{Nb^2}{12a}$	
	0	$\frac{2}{b}$ $\frac{2}{b}$	0	1	0		L				_	

Examining the right hand side of eqn (5) reveals that rigid body modes 1-4 are present in the global [Kg]. However, mode #5, which corresponds

DIRECTED-FORCE APPROACH

Consider Argyris's methodology [4] for the directed force problem. Let $\phi \cong (u_4 - u_{18})/2a$ or



becomes

Taking

Fig. 2. Two element 18 DOF.

$$\frac{\partial \mathbf{R}_i^{\text{DFC}}}{\partial u_i}$$

 $(u_1 - u_{17})/2a$. If one neglects the change in the axial component during a small rotation, one obtains $N(b/2)\cos\phi \cong Nb/2$. The load vector for this force

and assuming small angles yields

	$-\frac{bN}{4a}$	0	0	0	0	0	0	0	0	0	0	0	$\frac{bN}{4a}$	0	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	}	
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	0	0	0	$-\frac{bN}{4a}$	0	0	0	0	0	0	0	0	0	$\frac{bN}{4a}$	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	ļ	
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	ĺ	
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
$[K^{\rm DFC}] =$	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		(7)
[w]-	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		(7)
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	$\frac{bN}{4a}$	0	0	0	0	0	0	0	0	0	0	0	$-\frac{bN}{4a}$	0	0	0	0	0		
	0	0	0	$\frac{bN}{4a}$	0	0	0	0	0	0	0	0	0	$-\frac{bN}{4a}$	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0		
	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0]	

	Го	0	0	0	0]
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	$-\frac{Nb^2}{12a}$
$[[Kg] + [K^{DFC}] \cdot [Rigid Body Modes] =$	0	0	0	0	0
$[[Kg] + [K^{-1}] \cdot [Kigid Body Modes] =$	0	0	0	0	0
	0	0	0	0	$\frac{Nb^2}{12a}$
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	0
	0	0	0	0	$\frac{Nb^2}{12a}$
	0	0	0	0	$-\frac{Nb^2}{12a}$

(8)

Note that the pseudo-force terms associated with degrees of freedom 1, 4, 13 and 14 have disappeared, but the pseudo-moment terms associated with degrees of freedom 9, 12, 17 and 18 remain.

Once again, consider a rotation ϕ about the y-axis, and let $\phi = (u_{13} + u_{14} - u_1 - u_4)/4a$. We have considered the equivalent concentrated vertical loads at the nodes due to the rotation in the determination of $[K^{DFC}]$. However, this vertical component is really based on a distributed load. In order to replace it with concentrated loads at the node points, we must also include fixed-end moments. The load vector for the fixed end moments caused by the vertical distributed loads would be

Taking $\frac{\partial R_i^{\text{FEM}}}{\partial v_i}$, assuming small angles and enforcing symmetry yields

																	-
- 0	0	0	0	0	0	0	0	$-\frac{Nb^2}{10}$	0	0	$\frac{Nb^2}{10}$	0	0	0	0	$\frac{Nb^2}{48}$	$-\frac{Nb^2}{48a}$
						_	_	48a	_	_	48 <i>a</i>			~	~	48a	
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	Nb^2	0	0	$\underline{Nb^2}$	0	0	0	0	$\frac{Nb^2}{2}$	$-\frac{Nb^2}{m}$
Ŭ	v	Ŭ	Ū	0	Ŭ	v	Ũ	48 <i>a</i>	Ũ	v	48 <i>a</i>					48 <i>a</i>	48 <i>a</i>
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
Nb ²	0	0	Nb ²	0	0	0	0	0	0	0	0	Nb^2	$\underline{Nb^2}$	0	0	0	0
$-\frac{1}{48a}$	U	0	$-\frac{1}{48a}$	U	U	U	U	U	U	U	0	48 <i>a</i>	48 <i>a</i>	U	v	Ū	Ŭ
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
$\frac{Nb^2}{48a}$	0	0	$\frac{Nb^2}{48a}$	0	0	0	0	0	0	0	0	$-\frac{Nb^2}{48a}$	$-\frac{Nb^2}{48a}$	0	0	0	0
404			400					Nb ²			Nb ²	704				Nb ²	Nb ²
0	0	0	0	0	0	0	0	$\frac{1}{48a}$	0	0	$-\frac{110}{48a}$	0	0	0	0	$-\frac{110}{48a}$	48 <i>a</i>
0	0	0	0	0	0	0	0	$\underline{Nb^2}$	0	0	$-\frac{Nb^2}{2}$	0	0	0	0	$-\frac{Nb^2}{Nb^2}$	Nb ²
	Ŭ	v				v		48 <i>a</i>			48 <i>a</i>					48 <i>a</i>	48 <i>a</i>
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
$\frac{Nb^2}{48a}$	0	0	$\frac{Nb^2}{48a}$	0	0	0	0	0	0	0	0	$-\frac{Nb^2}{48a}$	$-\frac{Nb^2}{48a}$	0	0	0	0
$-\frac{Nb^2}{48a}$	0	0	$-\frac{Nb^2}{48a}$	0	0	0	0	0	0	0	0	$\frac{Nb^2}{48a}$	$\frac{Nb^2}{48a}$	0	0	0	0
-																	(10)
																	(10)

At this point,

	_				_	_
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
$[[Kg] - [K^{DFC}] + [K^{FEM}]]$	0	0	0	0	0	
\cdot [Rigid Body Modes] =	0	0	0	0	0	· (
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	0	0	0	0	0	
	-					-

Thus, the all pseudo-forces and pseudo-moments have been eliminated.

COMPARISON OF FREQUENCIES

In order to determine the effect of the correction factors, a comparison of the first 12 natural frequencies of vibration was performed for the two element model. The results are included in Table 1.

Table 1. Comparison of frequencies of vibration for twoelement free/free model

(11)	Frequency	Yang's [<i>Kg</i>] 16 DOF	Yang's [Kg] 24 DOF	Corrected 16 DOF	Corrected 24 DOF						
	1	0	0	0	0						
	2	0	0	0	0						
	3	0	0	0	0						
	4	0	0	0	0						
	5	103.2	0	0	0.0152						
	6	103.2	0	103.2	103.2						
	7	118.0	0	118.0	103.2						
	8	118.0	0	206.4	17 4.7						
	9	206.4	0	206.4	206.4						
	10	206.4	0	206.4	206.4						
	11	206.4	0	206.4	206.4						
	12	206.4	0	268.3	206.4						

As the table indicates, the missing rigid body mode (mode #5, frequency = 0), was obtained when the total directed force correction factors were added to Yang's [Kg].

SUMMARY

Based upon this investigation, the following conclusions can be made:

(1) Pre-loaded membrane elements exhibit the same "grounding" problems associated with preloaded beam elements. In other words, an internal grounding, or false stiffening, occurs when they are subjected to a rigid body rotation, which generates erroneous pseudo-forces.

(2) In addition to the pseudo-forces, erroneous pseudo-moments also develop during rigid body rotation.

(3) An erroneous non-zero frequency develops during rigid body rotation, instead of the required zero frequency associated with a rigid body mode.

(4) Using a directed force premise, and extending the Argyris approach for developing correction factors for non-conservative forces to the global level, the pseudo-forces were eliminated, but the pseudomoments still remained.

(5) By considering the fixed-end moments associated with the distributed load, and similarly developing correction factors, the pseudo-moments were eliminated, and the missing zero frequency associated with rigid body rotation were obtained.

(6) Future work should include comparison of results for larger models and with test results.

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