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A maximum entropy method for solving the boundary value problem of second order ordinary differential equations

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We propose a new method to solve the boundary value problem for a class of second order linear ordinary differential equations, which has a non-negative solution. The method applies the maximum entropy principle to approximating the solution numerically. The theoretical analysis and numerical examples show that our method is convergent. *Published by AIP Publishing.* <https://doi.org/10.1063/1.5029856>

I. INTRODUCTION

The maximum entropy principle was proposed by Jaynes in 1957,¹ which has been widely used in statistical physics, image processing, networks, etc.^{2–6} The maximum entropy method has been applied to computing absolutely continuous invariant measures of discrete dynamical systems^{7–9} and non-negative solutions of Fredholm integral equations.^{10,11} Theory and numerical experiments have shown that the maximum entropy method is efficient when splines are used as moment functions.¹²

There are a lot of numerical methods for solving differential equations, such as finite element methods, finite difference methods, finite volume methods, spectral methods, etc. These methods are not universal. There may be effective special techniques for special differential equations, for example, the adaptive method and the multiscale method. Numerically solving two-point boundary value problems for second order ordinary differential equations is often quite simple to formulate, but sometimes still needs particular techniques to get better results. For example, in Refs. 13 and 14, an initial value technique for singularly perturbed two-point boundary value problems using an exponentially fitted finite difference scheme was proposed. In Refs. 15 and 16, a method based on B-splines for solving a self-adjoint singularly perturbed two-point boundary value problem was given. In Ref. 17, a method was provided for the numerical solution of second order linear ordinary differential equations in the high-frequency regime. In Ref. 18, an exponentially fitted finite difference scheme was proposed for singularly perturbed two-point boundary value problems for second order ordinary differential equations with two small parameters multiplying the derivatives. In Ref. 19, a new method based on a single layer Legendre neural network model was developed to solve initial and boundary value problems.

The maximum entropy method was used to solve linear differential equations²⁰ and the Fokker-Planck equation.²¹ This method gave approximate solutions evaluated at some given points. The employed moment functions there are trigonometric functions. Another type of moment functions were used to solve some time dependent partial differential equations.^{22,23} In Ref. 24, higher order moments with respect to $\{x, x^2, x^3, \dots, x^N\}$ were evaluated to solve the Fokker-Planck equation. Such moment functions are all globally defined on the whole domain with single algebraic expressions, like all the classic maximum entropy methods, and consequently the instability issue is present due to the fact of ill-conditioning for the classic approach. Moreover, the maximum entropy principle behind

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all such methods was applied to a discrete system of algebraic equations, instead of the continuous system directly out of the Boltzmann entropy, for the Lagrange multipliers, so another error occurs in addition to the one from the use of finitely many moments. In our new method, we not only employ piecewise cubic functions to eliminate the ill-conditioning but also apply the maximum entropy principle to the whole involved integrals without choosing the evaluation points in the domain of integration.

When the solution of the boundary value problem is guaranteed to be non-negative, the maximum entropy principle, which is mainly used for the numerical recovery of unknown density functions, can be effectively explored, at least in theory. Here we try this kind of exploration by using the maximum entropy method to numerically solve the two-point boundary value problem for second order ordinary differential equations that have non-negative solutions. We shall employ cubic splines to form moment functions used in our method since such functions are locally nonzero and globally twice continuously differentiable. Our scheme provides a new way to think of the numerical solution of the special class of differential equations. In Sec. II, the maximum entropy method is proposed for the solution of the boundary value problem of a second order linear differential equation, and its convergence will be established in Sec. III. We provide an approximate maximum entropy method for nonlinear boundary value problems in Sec. IV. In Sec. V, some numerical experiments are given to show the efficiency of the method. The conclusion is given in Sec. VI.

II. A MAXIMUM ENTROPY METHOD FOR BOUNDARY VALUE PROBLEMS OF SECOND ORDER ORDINARY DIFFERENTIAL EQUATIONS

We consider the two-point boundary value problem for a second order linear ordinary differential equation

$$\begin{cases} Ly(x) \equiv -\frac{d}{dx}\left(p(x)\frac{dy}{dx}(x)\right) + q(x)y(x) = f(x), & a < x < b, \\ y(a) = \beta_1, y(b) = \beta_2, \end{cases} \quad (1)$$

where the given function p is differentiable and q and f are continuous functions. Suppose the two-point boundary value problem has a non-negative solution, and our purpose is to calculate it numerically.

For some boundary value problem, the maximum principle²⁵ can guarantee that it has a non-negative solution, as the following theorem shows.

Theorem 2.1. *If $p(x) \geq 0$, $f(x) \geq 0$, and $q(x) > 0$ for $x \in (a, b)$ and $\beta_1 \geq 0$ and $\beta_2 \geq 0$, then the solution y_* of the boundary value problem (1) satisfies that $y_*(x) \geq 0$ for all $x \in [a, b]$.*

Proof. Let \hat{x} be the minimum point of y_* in (a, b) and $y_*(\hat{x}) < 0$. Then $y'_*(\hat{x}) = 0$ and $y''_*(\hat{x}) \geq 0$, and so

$$f(\hat{x}) = -\frac{d}{dx}\left(p\frac{dy_*}{dx}\right)(\hat{x}) + q(\hat{x})y_*(\hat{x}) = -p(\hat{x})\frac{d^2y_*}{dx^2}(\hat{x}) + q(\hat{x})y_*(\hat{x}) < 0,$$

which contradicts to the fact that $f(\hat{x}) \geq 0$. □

For the purpose of using a piecewise polynomial maximum entropy method to numerically recover the non-negative solution y_* from above, we first divide the interval $[a, b]$ into n equal subintervals of length $h = (b - a)/n$. The nodes of the partition are $x_i = a + ih$, $i = -1, 0, 1, \dots, n, n + 1$, where x_{-1} and x_{n+1} are the two artificial ones to help dealing with the boundary condition.

We shall use the cubic B-spline functions associated with the above partition of $[a, b]$ to develop our maximum entropy method. They are constructed via translation and scaling from the “mother” cubic B-spline function defined on $(-\infty, \infty)$,²⁶ which can be directly calculated out as follows by the

fundamental property of its being second order continuously differentiable:

$$B(x) = \begin{cases} \frac{1}{6}(x^3 + 6x^2 + 12x + 8), & -2 \leq x \leq -1, \\ \frac{1}{6}(-3x^3 - 6x^2 + 4), & -1 \leq x \leq 0, \\ \frac{1}{6}(3x^3 - 6x^2 + 4), & 0 \leq x \leq 1, \\ \frac{1}{6}(-x^3 + 6x^2 - 12x + 8), & 1 \leq x \leq 2, \\ 0, & x \notin [-2, 2]. \end{cases}$$

Then the corresponding cubic B-spline functions are

$$B_i(x) = B\left(\frac{x - x_i}{h}\right), \quad i = -1, 0, \dots, n+1.$$

Now we begin to develop the maximum entropy method with the cubic B-spline functions. Since y_* solves Eq. (1), there is the equality

$$-\frac{d}{dx}\left(p\frac{dy_*}{dx}\right) + qy_* = f.$$

For $i = -1, 0, \dots, n+1$, multiplying B_i to both sides of the above equality and integrating from a to b , we get

$$\int_a^b B_i(x) \left[-\frac{d}{dx}\left(p\frac{dy_*}{dx}\right)(x) + q(x)y_*(x) \right] dx = \int_a^b B_i(x)f(x)dx. \quad (2)$$

Using integration by parts two times gives

$$\begin{aligned} & -\int_a^b B_i(x) \frac{d}{dx}\left(p\frac{dy_*}{dx}\right)(x) dx \\ &= -B_i(x)p(x)\frac{dy_*}{dx}(x)\Big|_a^b + \int_a^b p(x)y'_*(x)B'_i(x)dx \\ &= -B_i(x)p(x)\frac{dy_*}{dx}(x)\Big|_a^b + p(x)y_*(x)B'_i(x)\Big|_a^b \\ & \quad - \int_a^b y_*(x)[p'(x)B'_i(x) + p(x)B''_i(x)]dx. \end{aligned}$$

For each $i = -1, 0, \dots, n+1$, let

$$\tilde{g}_i(x) = -[p'(x)B'_i(x) + p(x)B''_i(x)] + q(x)B_i(x) \quad (3)$$

and

$$\tilde{m}_i = \int_a^b B_i(x)f(x)dx.$$

Then equality (2) can be written as

$$\int_a^b y_*(x)\tilde{g}_i(x)dx + \left[-B_i(x)p(x)\frac{dy_*}{dx}(x)\Big|_a^b + p(x)y_*(x)B'_i(x)\Big|_a^b \right] = \tilde{m}_i. \quad (4)$$

Now we evaluate the expression in the brackets above for various indices i . Clearly $B_i(a) = B_i(b) = 0$ and $B'_i(a) = B'_i(b) = 0$ for $i = 2, 3, \dots, n-2$. So for such indices i , (4) is reduced to

$$\int_a^b y_*(x)\tilde{g}_i(x)dx = \tilde{m}_i.$$

Since $B_i(b) = B'_i(b) = 0$ for $i = -1, 0, 1$ and

$$\begin{aligned} B_{-1}(a) &= \frac{1}{6}, & B_0(a) &= \frac{2}{3}, & B_1(a) &= \frac{1}{6}, \\ B'_{-1}(a) &= -\frac{1}{2h}, & B'_0(a) &= 0, & B'_1(a) &= \frac{1}{2h}, \end{aligned}$$

the equalities in (4) with $i = -1, 0, 1$ become

$$\begin{cases} \int_a^b y_*(x) \tilde{g}_{-1}(x) dx + \frac{1}{6} p(a) \frac{dy_*}{dx}(a) + \frac{1}{2h} p(a) y_*(a) = \tilde{m}_{-1}, \\ \int_a^b y_*(x) \tilde{g}_0(x) dx + \frac{2}{3} p(a) \frac{dy_*}{dx}(a) = \tilde{m}_0, \\ \int_a^b y_*(x) \tilde{g}_1(x) dx + \frac{1}{6} p(a) \frac{dy_*}{dx}(a) - \frac{1}{2h} p(a) y_*(a) = \tilde{m}_1. \end{cases}$$

In the above equalities, multiplying the first and last by 4 and subtracting the second lead to

$$\int_a^b y_*(x) [4\tilde{g}_{-1}(x) - \tilde{g}_0(x)] dx = 4\tilde{m}_{-1} - \tilde{m}_0 - \frac{2}{h} p(a) y_*(a)$$

and

$$\int_a^b y_*(x) [4\tilde{g}_1(x) - \tilde{g}_0(x)] dx = 4\tilde{m}_1 - \tilde{m}_0 + \frac{2}{h} p(a) y_*(a).$$

Similarly, from $B_i(a) = B'_i(a) = 0$ for $i = n-1, n, n+1$ and

$$\begin{aligned} B_{n-1}(b) &= \frac{1}{6}, & B_n(b) &= \frac{2}{3}, & B_{n+1}(b) &= \frac{1}{6}, \\ B'_{n-1}(b) &= -\frac{1}{2h}, & B'_n(b) &= 0, & B'_{n+1}(b) &= \frac{1}{2h}, \end{aligned}$$

the equalities (4) with $i = n-1, n, n+1$ are

$$\begin{cases} \int_a^b y_*(x) \tilde{g}_{n-1}(x) dx - \frac{1}{6} p(b) \frac{dy_*}{dx}(b) - \frac{1}{2h} p(b) y_*(b) = \tilde{m}_{n-1}, \\ \int_a^b y_*(x) \tilde{g}_n(x) dx - \frac{2}{3} p(b) \frac{dy_*}{dx}(b) = \tilde{m}_n, \\ \int_a^b y_*(x) \tilde{g}_{n+1}(x) dx - \frac{1}{6} p(b) \frac{dy_*}{dx}(b) + \frac{1}{2h} p(b) y_*(b) = \tilde{m}_{n+1}, \end{cases}$$

from which it follows that

$$\int_a^b y_*(x) [4\tilde{g}_{n-1}(x) - \tilde{g}_n(x)] dx = 4\tilde{m}_{n-1} - \tilde{m}_n + \frac{2}{h} p(b) y_*(b)$$

and

$$\int_a^b y_*(x) [4\tilde{g}_{n+1}(x) - \tilde{g}_n(x)] dx = 4\tilde{m}_{n+1} - \tilde{m}_n - \frac{2}{h} p(b) y_*(b).$$

Define the following numbers, which are called the *moments*, as

$$m_i = \begin{cases} 4\tilde{m}_{-1} - \tilde{m}_0 - \frac{2}{h} p(a) y_*(a), & i = 0, \\ 4\tilde{m}_1 - \tilde{m}_0 + \frac{2}{h} p(a) y_*(a), & i = 1, \\ \tilde{m}_i, & i = 2, \dots, n-2, \\ 4\tilde{m}_{n-1} - \tilde{m}_n + \frac{2}{h} p(b) y_*(b), & i = n-1, \\ 4\tilde{m}_{n+1} - \tilde{m}_n - \frac{2}{h} p(b) y_*(b), & i = n \end{cases}$$

and the corresponding *moment functions* as

$$g_i(x) = \begin{cases} 4\tilde{g}_{-1}(x) - \tilde{g}_0(x), & i = 0, \\ 4\tilde{g}_1(x) - \tilde{g}_0(x), & i = 1, \\ \tilde{g}_i(x), & i = 2, \dots, n-2, \\ 4\tilde{g}_{n-1}(x) - \tilde{g}_n(x), & i = n-1, \\ 4\tilde{g}_{n+1}(x) - \tilde{g}_n(x), & i = n. \end{cases} \quad (5)$$

Then the following is valid:

$$\int_a^b y_*(x) g_i(x) dx = m_i, \text{ for } i = 0, 1, \dots, n.$$

The above equalities provide a framework for numerically solving the boundary value problem (1) by the maximum entropy method. That is, among all the non-negative solutions y of the equality constraints

$$\int_a^b y(x)g_i(x)dx = m_i, \text{ for } i = 0, 1, \dots, n, \quad (6)$$

we find the one, denoted as y_n , which maximizes the corresponding *generalized differential entropy* defined by

$$H(y) = - \int_a^b y(x) \ln y(x) dx + \int_a^b y(x) dx, \quad \forall y \in L^1(a, b), y \geq 0. \quad (7)$$

Namely, we solve the following constrained optimization problem:¹¹

$$\max \left\{ H(y) : y \geq 0, \int_a^b y(x)g_i(x)dx = m_i, i = 0, 1, \dots, n \right\}. \quad (8)$$

To get the expression of the solution of (8), we need the following lemma.

Lemma 2.1. *Given any two non-negative functions $f, g \in L^1(a, b)$,*

$$\int_a^b f(x)dx - \int_a^b f(x) \ln f(x)dx \leq \int_a^b g(x)dx - \int_a^b f(x) \ln g(x)dx.$$

Proof. In the Gibbs inequality²⁷

$$u - u \ln u \leq v - u \ln v, \quad \forall u, v \geq 0,$$

letting $u = f(x)$ and $v = g(x)$ and then integrating both sides, we obtain the required integral inequality. \square

Theorem 2.2. *The non-negative function in $L^1(a, b)$ that maximizes the entropy functional (7) under the constraints (6) is*

$$y_n(x) = \exp \left(\sum_{i=0}^n \lambda_i g_i(x) \right)$$

if the constants $\lambda_0, \dots, \lambda_n$, which are called the Lagrange multipliers, satisfy the given constraints

$$\int_X g_i(x) \exp \left(\sum_{j=0}^n \lambda_j g_j(x) \right) dx = m_i, \quad i = 0, \dots, n. \quad (9)$$

Proof. Let $y \in L^1(a, b)$ be non-negative and satisfy the moment constraints

$$\int_a^b y(x)g_j(x)dx = m_j, \quad j = 0, \dots, n.$$

Then, using Lemma 2.1, we have

$$\begin{aligned} H(y) &= - \int_a^b y(x) \ln y(x) dx + \int_a^b y(x) dx \\ &\leq \int_a^b y_n(x) dx - \int_a^b y(x) \ln y_n(x) dx \\ &= \int_a^b y_n(x) dx - \sum_{i=0}^n \lambda_i \int_a^b y(x) g_i(x) dx \\ &= \int_a^b y_n(x) dx - \sum_{i=0}^n \lambda_i m_i \\ &= \int_a^b y_n(x) dx - \int_a^b y_n(x) \left[\sum_{i=0}^n \lambda_i g_i(x) \right] dx \\ &= \int_a^b y_n(x) dx - \int_a^b y_n(x) \ln y_n(x) dx = H(y_n). \end{aligned}$$

This shows that y_n is an optimal solution. Furthermore, since H is strictly concave,²⁸ the optimal solution is unique. \square

The nonlinear equation (9) can be solved by Newton's method as in other maximum entropy algorithms. Since the support of each $B_i(x)$ is local, from (3) to (5), the Jacobian matrix of the nonlinear equation is banded, which can be shown to be positive definite with the same technique as in Ref. 8. For more details of the maximum entropy method, see Refs. 7, 8, and 11.

III. CONVERGENCE OF THE METHOD

The weak convergence of the classic maximum entropy method was first obtained by Mead and Papanicolaou in their pioneering paper,²⁹ and Borwein and Lewis²⁸ gave the first rigorous convergence analysis for the moment problem under the L^1 -norm, based on a more general convergence theory in the context of topological spaces, as developed in Ref. 30. The norm convergence of our cubic spline maximum entropy method for the boundary value problem (1) can be established via the general convergence theory from Ref. 28. To put our algorithm in the framework for the convergence analysis in Ref. 28, we divide $[a, b]$ into 2^r equal subintervals in succession with $r = 2, 3, \dots$. And we use y_r instead of y_{2^r} to denote the approximate maximum entropy solution of (1) with respect to the corresponding partition of $[a, b]$ into 2^r equal sub-intervals. In other words, we calculate the function y_r among all non-negative functions of $L^1(a, b)$ that satisfy the $2^r + 1$ moment constraints (6).

Denote by D_r the feasible set of the maximum entropy problem, that is, the collection of all non-negative solutions of (6). Then $D_r \supset D_{r+1}$ for all r . Clearly y_r solves

$$\max \left\{ H(y) : y \in D_r = \bigcap_{k=2}^r D_k \right\}.$$

On the other hand, it is easy to see that the unique solution y_* of the original boundary value problem (1) is the unique solution of

$$\max \left\{ H(y) : y \in \bigcap_{r=2}^{\infty} D_r \right\}.$$

Using the same idea for the convergence analysis in Ref. 11, we can prove the following theorem.

Theorem 3.1. *Suppose that $H(y_*) > -\infty$. Let y_r be the maximum entropy solution of (8) with respect to the equal partition of $[a, b]$ into 2^r subintervals. Then $\lim_{r \rightarrow \infty} \|y_r - y_*\|_1 = 0$.*

As for the convergence rate of our algorithm, from Theorem 4.7 of Ref. 28,

$$\|y_r - y_*\|_1 \leq d_r e^{\frac{d_r}{2}} = O(d_r),$$

where

$$d_r = \inf \left\{ \left\| \ln y_* - \sum_{i=0}^{2^r} \alpha_i g_i \right\| : \forall \alpha_0, \alpha_1, \dots, \alpha_{2^r} \right\}$$

is the minimal distance of the function $\ln y_*$ to the subspace spanned by all g_i 's. From the expressions (5) of the moment functions g_i 's, which involve the differential operator L as well as the cubic B-spline functions B_i 's, we see that in general, it is not an easy task to estimate d_r theoretically. However, our numerical experiments in Sec. V show that the convergence rate of the method is $\|y_n - y_*\|_1 = O(1/n^2)$, where y_n is the approximate solution of the boundary value problem when $[a, b]$ is partitioned into n equal sub-intervals.

IV. A CLASS OF NONLINEAR BOUNDARY VALUE PROBLEMS

The idea of the above maximum entropy method for linear boundary value problems may be extended to solving nonlinear boundary value problems numerically, resulting in an approximate

maximum entropy method for a class of boundary value problems of nonlinear ordinary differential equations.

Let the right-hand side of the boundary value problem (1) be $f(x, y)$ that is a general nonlinear differentiable function of its variables. Motivated by the explicit expression of maximum entropy solutions, we choose an initial vector $(\lambda_0^0, \dots, \lambda_n^0)$ for the Lagrange multipliers and form the initial approximate maximum entropy solution

$$y_n^0(x) = \exp\left(\sum_{i=0}^n \lambda_i^0 g_i(x)\right)$$

for the boundary value problem. Next, we substitute y_n^0 into the right side of the differential equation for the variable y in the expression of $f(x, y)$ so that we have obtained the corresponding linear boundary value problem (1) as

$$\begin{cases} Ly(x) \equiv -\frac{d}{dx}\left(p(x)\frac{dy}{dx}(x)\right) + q(x)y(x) = f(x, y_n^0(x)), & a < x < b, \\ y(a) = \beta_1, y(b) = \beta_2. \end{cases} \quad (10)$$

We can apply the maximum entropy method developed in Secs. II and III to solve (10). Then we get new values $\lambda_0^1, \dots, \lambda_n^1$ of the Lagrange multipliers, and a new approximate maximum entropy solution of the nonlinear boundary value problem is

$$y_n^1(x) = \exp\left(\sum_{i=0}^n \lambda_i^1 g_i(x)\right).$$

By the same token, we substitute y_n^1 into the right-hand side of (10), solve the updated linear boundary value problem by the maximum entropy method, and get the next approximate maximum entropy solution y_n^2 . From this iterative process, we can get a sequence of approximate maximum entropy solutions $\{y_n^0, y_n^1, \dots\}$ by our maximum entropy method. The fourth example of Sec. V will illustrate this approach for a famous boundary value problem.

V. NUMERICAL RESULTS

In this section, we present some numerical examples for solving boundary value problems for second-order linear ordinary differential equations using our method.

Example 1. The first example is

$$\begin{cases} y'' - y = x, & 0 < x < 1, \\ y(0) = 0, y(1) = 1, \end{cases}$$

whose exact solution is

$$y(x) = \frac{2(e^x - e^{-x})}{e - e^{-1}} - x.$$

The L^1 -norm errors $\|y_n - y_*\|_1$ and the L^∞ -norm ones $\|y_n - y_*\|_\infty$ of the numerical approximations are shown in Table I. The numerical solutions converge, and the convergence rate is approximately $O(1/n^2)$ for the L^1 -norm errors and $O(1/n)$ for the L^∞ -norm errors.

Example 2. We consider a singular perturbation two-point boundary value problem¹⁶

$$\begin{cases} -\varepsilon y'' + y = 1 + 2\sqrt{\varepsilon} \left[\exp\left(-\frac{x}{\sqrt{\varepsilon}}\right) + \exp\left(\frac{x-1}{\sqrt{\varepsilon}}\right) \right], & 0 < x < 1, \\ y(0) = 0, y(1) = 0, \end{cases}$$

and the exact solution is

$$y(x) = 1 - (1 - x) \exp\left(-\frac{x}{\sqrt{\varepsilon}}\right) - x \exp\left(\frac{x-1}{\sqrt{\varepsilon}}\right).$$

TABLE I. Errors for Example 1.

n	L^1 -norm error	L^∞ -norm error
4	5.3×10^{-3}	2.6×10^{-2}
8	1.6×10^{-3}	1.3×10^{-2}
16	4.5×10^{-4}	6.4×10^{-3}
32	1.3×10^{-4}	3.2×10^{-3}
64	3.6×10^{-5}	1.6×10^{-3}
128	9.9×10^{-6}	8.0×10^{-4}
256	2.7×10^{-6}	4.0×10^{-4}

Example 3. The third example is also from Ref. 16, which is another singularly perturbed two-point boundary value problem

$$\begin{cases} -\varepsilon y'' + y = x, & 0 < x < 1, \\ y(0) = 1, & y(1) = 1 + \exp\left(-\frac{1}{\sqrt{\varepsilon}}\right), \end{cases}$$

whose exact solution is

$$y(x) = x + \exp\left(-\frac{x}{\sqrt{\varepsilon}}\right).$$

The numerical results for Examples 2 and 3 are shown in Tables II and III, respectively. The second and third columns are the results when $\varepsilon = 1/16$ and the fourth and fifth columns are the results when $\varepsilon = 1/32$. We see that the convergence rate for Example 2 is the same as for Example 1, but both L^1 -norm and L^∞ -norm errors are roughly $O(1/n^2)$ for Example 3.

Example 4. The last example is a nonlinear two-point boundary value problem for the so-called Thomas-Fermi equation,³¹

TABLE II. Errors for Example 2.

n	L^1 -norm error	L^∞ -norm error	L^1 -norm error	L^∞ -norm error
4	4.99×10^{-2}	1.91×10^{-1}	6.62×10^{-2}	0.27×10^{-1}
8	1.45×10^{-2}	9.21×10^{-2}	1.89×10^{-2}	1.26×10^{-1}
16	4.40×10^{-3}	4.5×10^{-2}	5.45×10^{-3}	6.10×10^{-2}
32	1.31×10^{-3}	2.26×10^{-2}	1.63×10^{-3}	3.02×10^{-2}
64	3.81×10^{-4}	1.13×10^{-2}	4.80×10^{-4}	1.51×10^{-2}
128	1.09×10^{-4}	5.67×10^{-3}	1.37×10^{-4}	7.56×10^{-3}
256	3.05×10^{-5}	2.84×10^{-3}	3.89×10^{-5}	3.79×10^{-3}

TABLE III. Errors for Example 3.

n	L^1 -norm error	L^∞ -norm error	L^1 -norm error	L^∞ -norm error
4	4.36×10^{-3}	3.39×10^{-2}	6.23×10^{-3}	5.39×10^{-2}
8	9.74×10^{-4}	8.12×10^{-3}	1.28×10^{-3}	1.16×10^{-2}
16	2.36×10^{-4}	2.07×10^{-3}	3.06×10^{-4}	2.98×10^{-3}
32	5.86×10^{-5}	5.24×10^{-4}	7.56×10^{-5}	7.63×10^{-4}
64	1.46×10^{-5}	1.32×10^{-4}	1.88×10^{-5}	1.93×10^{-4}
128	3.65×10^{-6}	3.31×10^{-5}	4.70×10^{-6}	4.86×10^{-5}
256	9.12×10^{-7}	8.29×10^{-6}	1.18×10^{-6}	1.22×10^{-5}

TABLE IV. Errors for Example 4.

n	L^1 -norm error	L^∞ -norm error	$y'(0)$
4	1.1×10^{-3}	1.0×10^{-2}	-1.0855
8	3.1×10^{-4}	4.3×10^{-3}	-1.2032
16	8.3×10^{-5}	1.7×10^{-3}	-1.2980
32	2.2×10^{-5}	6.6×10^{-4}	-1.3725
64	5.6×10^{-6}	2.5×10^{-4}	-1.4298
128	1.4×10^{-6}	9.2×10^{-5}	-1.4729
256	4.3×10^{-7}	3.4×10^{-5}	-1.5049

$$\begin{cases} y'' = \frac{y^{\frac{3}{2}}(x)}{\sqrt{x}}, & 0 < x < +\infty, \\ y(0) = 1, \quad \lim_{x \rightarrow \infty} y(x) = 0, \end{cases}$$

whose exact solution is unknown. There are lots of papers about the above equation; see Ref. 32 and the references therein. To test our method in Sec. IV, we consider the following boundary value problem of the Thomas-Fermi equation on a finite interval

$$\begin{cases} y'' = \frac{y^{\frac{3}{2}}(x)}{\sqrt{x}}, & 0 < x < 1, \\ y(0) = 1, \quad y(1) = y_*(1), \end{cases}$$

where the boundary condition is chosen according to an approximate analytic solution in the following:³³

$$y_*(x) = \frac{(1 + 1.81061x^{\frac{1}{2}} + 0.60112x)^2}{(1 + 1.81061x^{\frac{1}{2}} + 1.39515x + 0.77112x^{\frac{3}{2}} + 0.21465x^2 + 0.04793x^{\frac{5}{2}})^2}. \quad (11)$$

In this example, the right-hand side of the differential equation contains the unknown solution y . So the moments used in the maximum entropy method cannot be defined as in Eq. (5). However, we can follow the iterative process described in Sec. IV to obtain an approximate maximum entropy solution accurate to a desired precision.

The L^1 -norm errors $\|y_n - y_*\|_1$ and the L^∞ -norm ones $\|y_n - y_*\|_\infty$ of the numerical approximations are shown in Table IV, where y_* is given by (11). The numerical solutions seem to converge, and the convergence rate is approximately $O(1/n^2)$ for the L^1 -norm errors and faster than $O(1/n)$ for the L^∞ -norm errors. We also list an estimated value of the derivative $y'(0)$ by our method, which is important in physics³² that obtained a better approximate value $-1.588071\dots$ for $y'(0)$ after a rather complicated iterative process.

VI. CONCLUSIONS

In this study, we proposed a maximum entropy method, based on cubic B-spline functions, to approximate non-negative solutions of boundary value problems of second order linear ordinary differential equations. The theoretical analysis and numerical results show the convergence of the method. Furthermore, since the main numerical work is solving a system of nonlinear equations for the Lagrange multipliers without ill-conditioning, the algorithm is stable and easy to be implemented. We also applied this method to the boundary value problem for some special second order nonlinear ordinary differential equations.

The maximum entropy method with splines as moment functions provides a serious and promising numerical approach for the recovery of non-negative function solutions of differential, integral, and operator equations. Extending our method to multi-variable differential equations and other difficult problems in stochastic analysis and random dynamical systems will be our future work when non-negative or density functions need our computational exploration.

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