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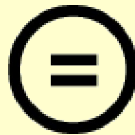
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理學博士 學位論文

**Dye-Sensitized Solar Cells based on
Titanium Oxide(IV) Nanotube Arrays**

티타니아 나노튜브 층이 도입된
염료감응형 태양전지

2017年 8月

서울대학교 大學院

化學部

千命煥

**Dye-Sensitized Solar Cells based on Titanium
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티타니아 나노튜브 층이 도입된 염료감응형 태양전지

指導教授 徐 正 雙 博士

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서울大學校 大學院

化學部

千 命 煥

천명환의 이학박사 학위논문을 인준함

2017년 8월

위 원 장 정 연 준 (인)

부위원장 서 정 쌍 (인)

위 원 홍 종 인 (인)

위 원 정 대 흥 (인)

위 원 윤 천 호 (인)

Abstract

We prepared dye-sensitized solar cells (DSSCs) with enhanced energy conversion efficiency using open-ended TiO₂ nanotube arrays with a TiO₂ scattering layer. As compared to closed-ended TiO₂ nanotube arrays, the energy conversion efficiency of the open-ended TiO₂ nanotube arrays was increased from 5.63% to 5.92%, which is an enhancement of 5.15%. With the TiO₂ scattering layer, the energy conversion efficiency was increased from 5.92% to 6.53%, which is an enhancement of 10.30%. After treating the open-ended TiO₂ nanotube arrays with TiCl₄, the energy conversion efficiency was increased from 5.92% to 6.89%, a 16.38% enhancement, which is attributed to improved light harvesting and increased dye adsorption.

Dye-sensitized solar cells (DSSCs) were fabricated using open-ended freestanding TiO₂ nanotube arrays functionalized with Ag nanoparticles (NPs) in the channel to create a plasmonic effect, and then coated with large TiO₂ NPs to create a scattering effect in order to improve energy conversion efficiency. Compared to closed-ended freestanding TiO₂ nanotube array-based DSSCs without Ag or large TiO₂ NPs, the energy conversion efficiency of closed-ended DSSCs improved by 9.21%(actual efficiency, from 5.86% to 6.40%) with Ag NPs, 6.48%(actual efficiency, from 5.86% to 6.24%) with TiO₂ NPs, and

14.50%(actual efficiency, from 5.86% to 6.71%) with both Ag NPs and TiO₂ NPs. By introducing Ag NPs and/or large TiO₂ NPs to open-ended freestanding TiO₂ nanotube array-based DSSCs, the energy conversion efficiency was improved by 9.15%(actual efficiency, from 6.12% to 6.68%) with Ag NPs and 8.17%(actual efficiency, from 6.12% to 6.62%) with TiO₂ NPs, and by 15.20%(actual efficiency, from 6.12% to 7.05%) with both Ag NPs and TiO₂ NPs. Moreover, compared to closed-ended freestanding TiO₂ nanotube arrays, the energy conversion efficiency of open-ended freestanding TiO₂ nanotube arrays increased from 6.71% to 7.05%. We demonstrate that each component—Ag NPs, TiO₂ NPs and open-ended freestanding TiO₂ nanotube arrays—enhanced the energy conversion efficiency, and the use of a combination of all components in DSSCs resulted in the highest energy conversion efficiency.

Keywords: open-ended freestanding TiO₂ nanotube arrays; dye-sensitized solar cells; plasmonic; scattering; anodization

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Chapter 1 Introduction

1. 1. Anodization

Anodization of metal oxide has applications to prevent corrosion, to form capacitor dielectrics [1], to template nanomaterials [2-7], and to apply to catalysis, optics, and electronics [8-10]. Anodic aluminum oxide (AAO), known as porous anodic oxide, was first reported [11-14] and is now commercially available because its pores can be used as template for preparing nanostructures. Additionally, anodic metal oxide has been achieved on surface of many other metals such as tungsten [15], zirconium [16], titanium [17-19], tantalum [20], hafnium [21], niobium [22], and vanadium [23].

The formation of anodic metal oxides is based on two continuous processes, one is oxide dissolution at the electrolyte/oxide interface and the other is oxidation of metal at the oxide/metal interface. To optimize the anodization condition for pore ordering and pore size, the electrochemical process and chemical reagents are more investigated [24, 25].

1. 1. 1. Titanium (Ti) anodization - Anodic titanium oxide (ATO)

Anodic titanium oxide, TiO₂ nanotubes, has better properties compared to many other forms of TiO₂ for application in photocatalysis [26, 27],

gas sensors [28-31], photoelectrolysis [10, 32, 33], and photovoltaics [34, 36]. Since Zwilling reported the TiO₂ nanotubes in chromic acid and hydrofluoric acid in 1999 [37], TiO₂ nanotubes have been made in the fabrication, characterization, application, and formation mechanism [38].

Gong reported TiO₂ nanotubes in a 0.5 wt% HF aqueous solution at room temperature using different anodizing voltages as shown in Figure 1. 1. [17].

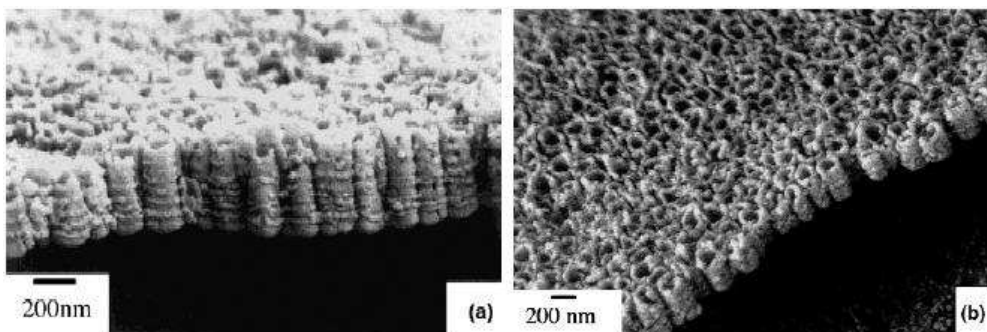


Figure 1. 1. FE-SEM cross-sectional images of titanium oxide nanotubes. The sample was anodized in 0.5 wt% HF solution at 20 V for 20 min [17].

The film thickness could not be increased further from 400-500 nm using HF-based electrolyte. Fluoride solution dissolves the TiO₂ by forming TiF₆²⁻ anions. Too strong acidity of HF electrolyte results in a too fast dissolution of the TiO₂. To overcome this problem, Mor

reported that addition of acetic acid to a 0.5 wt% HF electrolyte in a 1:7 ratio was resulted in more mechanically robust TiO₂ nanotubes without changing their shape and size [29, 39]. The surface morphology of TiO₂ nanotubes are anodized in an electrolyte containing 2.5 % HNO₃ and 1 % HF at 20 V for 4 h showed a uniform, clean, regular TiO₂ nanotube structure with a length about 400 nm that was reported by Ruan. An electrolyte of 0.5 M H₃BO₃, 2.5 % HNO₃, and 1 % HF in anodization at 20 V for 4 h led to irregular pore of TiO₂ nanotubes but long length of TiO₂ nanotubes about 560 nm [40]. To increase the length of TiO₂ nanotubes, Ti was anodized with KF or NaF in electrolyte [41]. The acidity of electrolyte might be controlled by adding HF, H₂SO₄ or Na₂SO₄ to adjust the balance of dissociation of TiO₂ at the electrolyte/oxide interface and oxidation of TiO₂ at the oxide/metal interface [42, 43].

Grimes and co-workers obtained TiO₂ nanotubes up to approximately 1000 nm using a variety of organic electrolytes such as dimethyl sulfoxide, formamide, ethylene glycol, and N-methylformamide as shown in Figure 1. 2. [19, 44, 45].

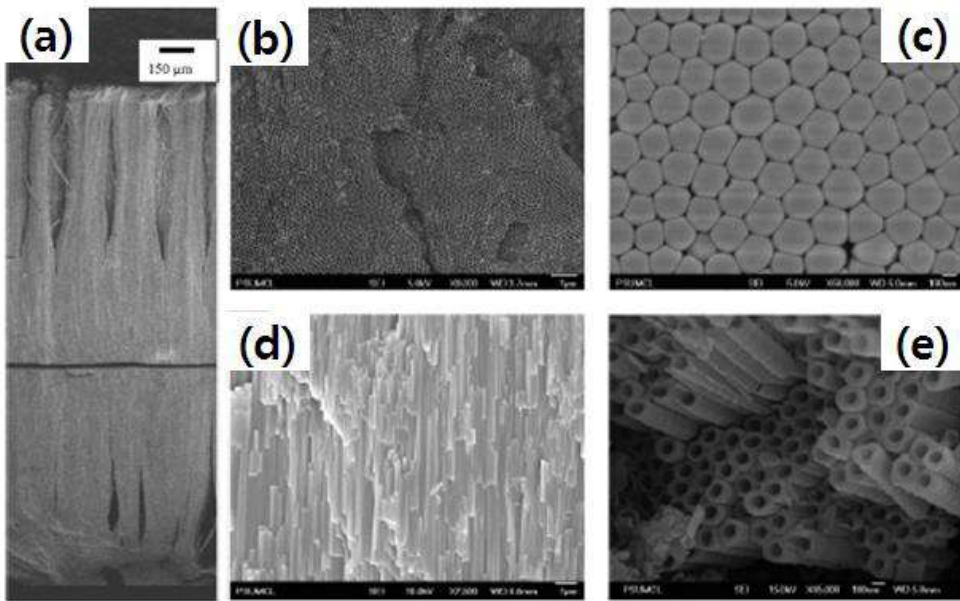


Figure 1. 2. TiO_2 nanotubes over 2 mm in thickness. (a) Cross-sectional view of self-standing TiO_2 nanotubes, (b) top side of nanotube array film, (c) back side or barrier layer side of nanotube array film, and (d, e) cross sectional image of mechanically fractured nanotube array film showing its tubular nature [45].

The key of long length of TiO_2 nanotubes is to minimize water content in the anodization bath to less than 5 %. In organic electrolytes, a little water content reduces the dissociation of the oxide in the fluorine containing electrolytes.

Tahya reported that water content affect the morphology of titanium

oxide array films. [46]

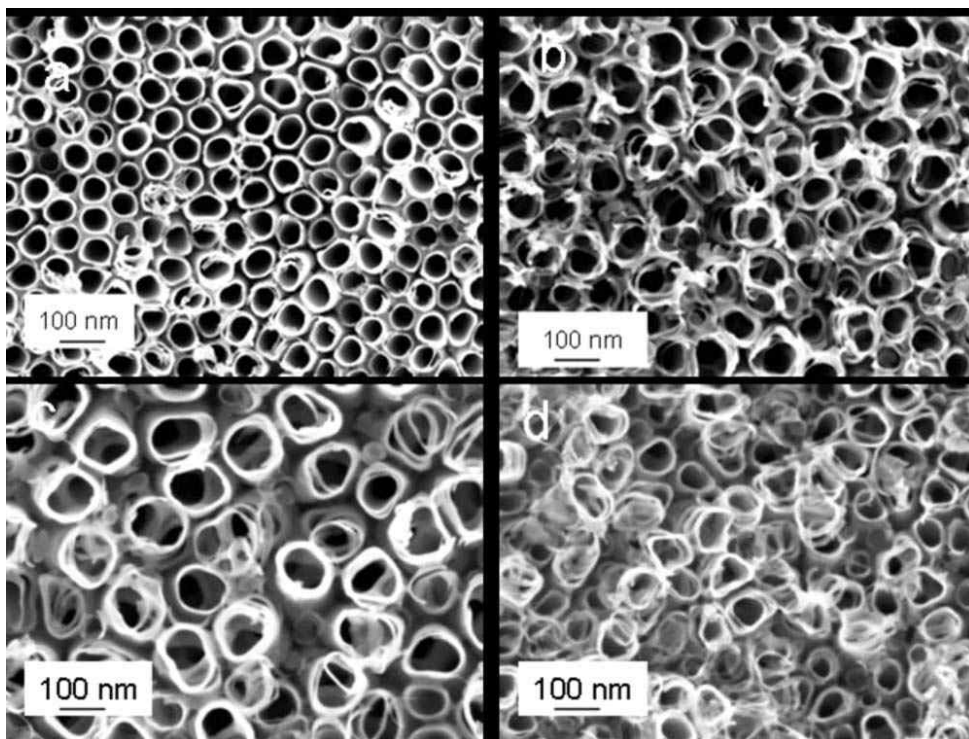


Figure 1. 3. SEM images of TiO_2 NTs grown by electrochemical anodization in $\text{NH}_4\text{F} + \text{H}_2\text{O} + \text{glycerol}$ electrolyte at different H_2O concentrations: (a) 0% H_2O , (b) 20%, (c) 50%, and (d) 70%. The same NH_4F concentration 0.5% was used in these experiments. [46]

1. 1. 2. Chemical reaction of Titanium anodization

The chemical reaction of titanium anodization should be the same as that of aluminum anodization [19, 38, 47, 48].

The anodization process begins with an initial oxide layer formed by interaction of the surface Ti^{4+} ions with oxygen ions in the electrolyte. In the initial stages of the anodization process, field-assisted dissolution dominates the chemical dissolution due to the relatively large electric field across the thin oxide layer. Further oxide growth is controlled by field-assisted ion transport (O^{2-} and Ti^{4+} ions) through the growing oxide. The anode oxidation of titanium is shown in Figure 1. 4.

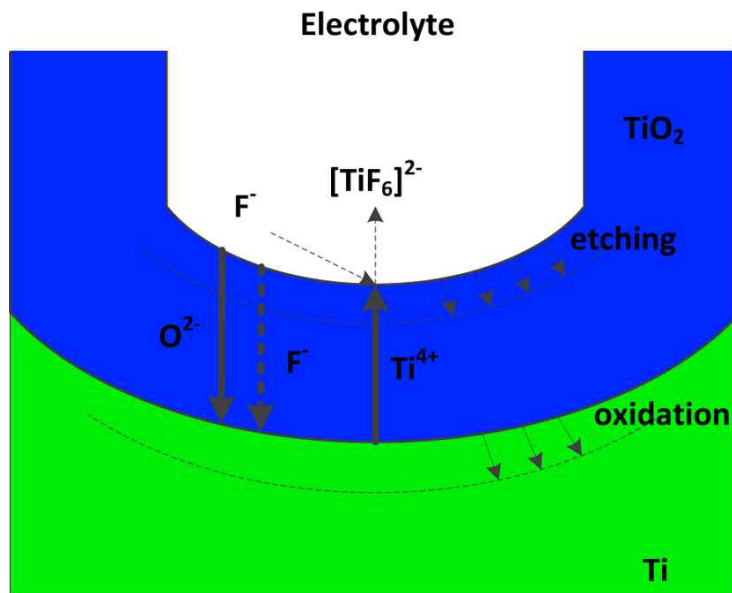
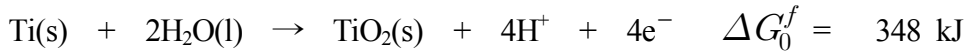


Figure 1. 4. Schematic drawing of adonization process. [49]

Anode:

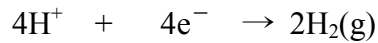


The overall anode reaction for anodic oxidation of titanium is represented as

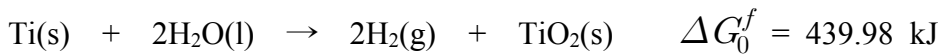


Meanwhile, hydrogen gas is released at the cathode that causes the flotation. The electrode reaction is described as follows:

Cathode:

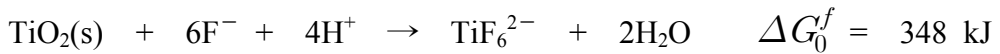


So the overall electrode reaction at this stage can be described as



After the formation of the oxide layer, the $\text{TiO}_2\text{--Ti}$ interface is locally activated, and the chemical and physical diffusions occur in the meantime. Small pores are formed first due to the localized dissolution of the oxide. Then, these pits are converted into bigger pores, and the pore density increases as governed by both electrochemical etching and chemical dissolution. Fluoride ions in the electrolyte have the ability to form water-soluble TiF_6^{2-} , and their small ionic radius makes them

suitable to enter the growing TiO₂ lattice and to be transported through the oxide by the applied electric field. After a while, the individual pores will interfere with each other, balance the available current, and finally reach a steady state condition, which results in a growth of the self-ordered nanotube structure. The nanotubes will grow steadily when the rate of TiO₂ formation equals the rate of TiF₆²⁻ formation and dissolution. The chemical dissolution process of the oxide is described as a chemical reaction



The basic premise of electrochemical anodization is the competition between the formation of the compact TiO₂ layer and the formation/dissolution of Ti⁴⁺ in the nanotube TiO₂ layer. After the formation of the compact TiO₂ oxide layer, Ti⁴⁺ ions at the metal–oxide interface will move toward the oxide–electrolyte interface under the applied electric field, as illustrated in Figure 1. 5. [50]

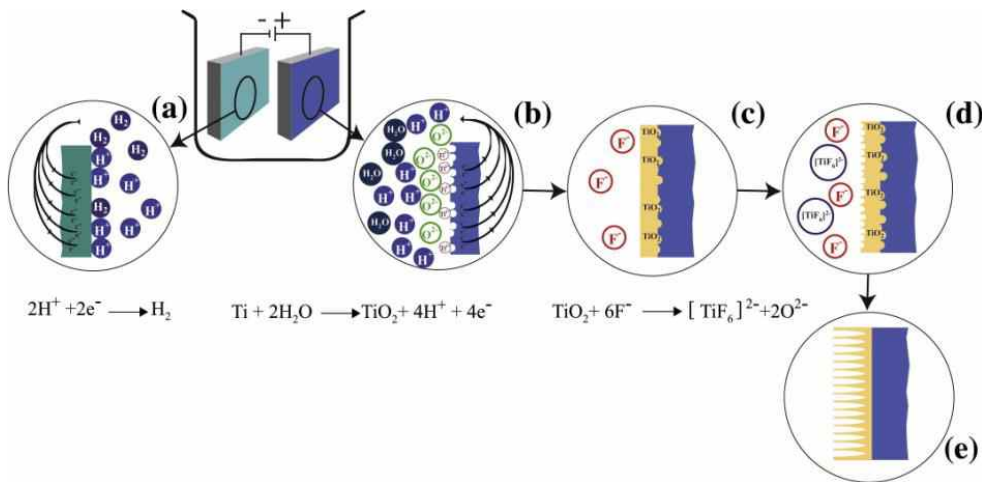


Figure 1. 5. Growth of regular TiO₂ nanotubes: (a) cathodic reaction, (b) anodic reaction, (c) transition state of TiO₂ layer, (d) starting of nanotube formation and (e) TiO₂ nanotubes [51].

1. 1. 3. Formation of TiO₂ nanotubes

The pores are developed from pits on the Ti plate surface and a schematic diagram for the equifield strength model as shown in Figure 1. 6.

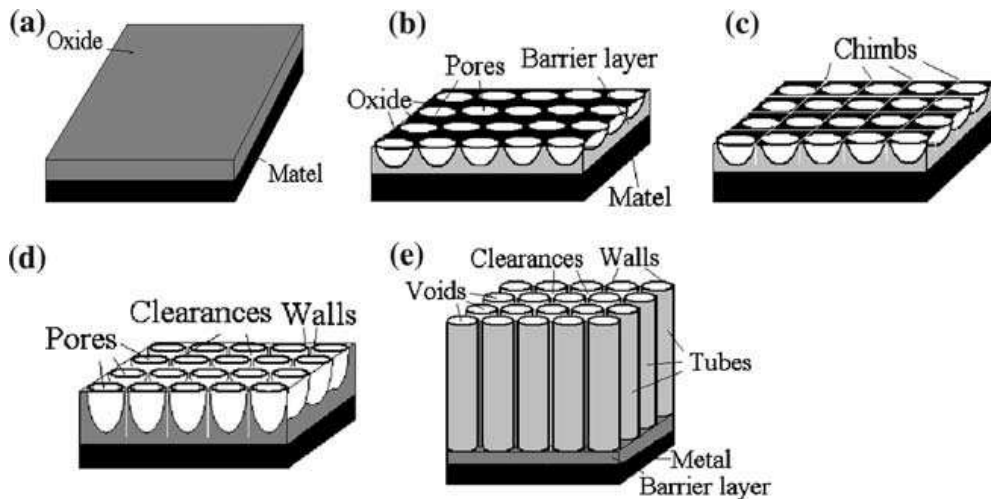


Figure 1. 6. Schematic diagram of the evolution of titania nanotubes in anodization as follows: (a) oxide layer formation; (b) pore formation on the oxide layer; (c) chimbs formation between pores; (d) growth of the pores and the chimbs; (e) fully developed of the titania nanotube arrays. [52]

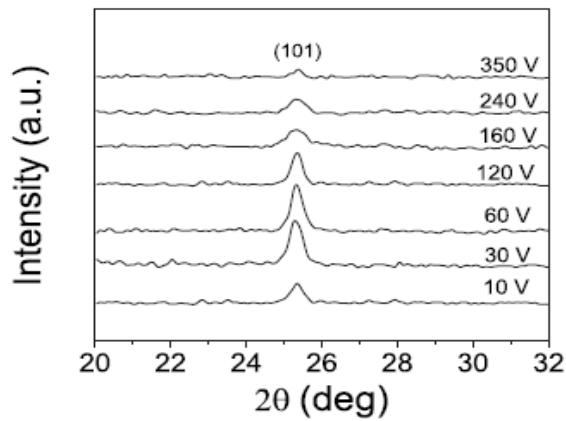
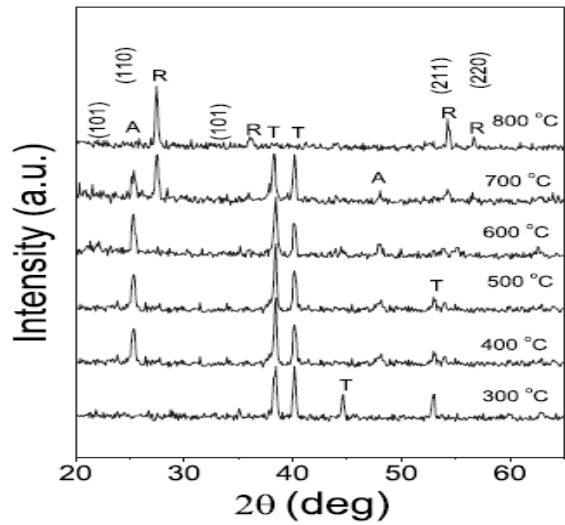
The initiation and growth of pores are associated with accelerated dissolution of TiO_2 with an influence of electric field. Many defects like impurities, dislocation, grain boundaries, or nonmetallic inclusions could cause a faster dissolution rate and lead to a pit growth [53, 54]. When Ti^{4+} cations are gotten out from the oxide surface that are caused by an applied field, Ti^{4+} cation vacancies can arise and accumulate to form high density voids in the oxide layer, which can help the propagation of pits [55, 56]. The pores are developed from

pits after dissolution, oxide at the electrolyte/oxide interface, and oxidation, titanium metal at the oxide/metal interface, are stable.

1. 1. 4 Anodization condition and annealing

For good appearance of titanium oxide array, anodization voltage is somewhat important. In Yahya's report, titanium oxide suitably evolve to anatase at 500°C in regular Voltage of 60s.

After anodization period, array was dealt in high temperature annealing. Through annealing process, titanium oxide expand its crystallography and strength. In Figure 1. 7, recommended anodization condition is 60V anodization and about 500°C annealing temperature.



(a)

(b)

Figure 1. 7. (a) GAXRD spectra of TiO_2 NTs grown at 0.7% acid concentration as a function of annealing temperature. Samples were annealed in temperature range of 300–800 °C with 100 °C step. A, R, and T refer to anatase, rutile, and titanium, respectively. (b) (101) anatase XRD diffraction peak of TiO_2 NTs as a function of anodization voltage. [46]

1. 2. Dye Sensitized Solar Cell

1.2.1. Overview

The first embodiment of modern day Dye-sensitized Solar Cell (DSC) dates back to late 1980s. [57] However, not until the fundamental work of Grätzel and O'Regan in 1991, [58] it was proven that DSCs can be a feasible alternative energy source. They built the first DSSCs which showed 7.1% energy conversion efficiency. [58] After that work, numerous studies had been reported. [60-68]

Although the power conversion efficiency of DSSCs is not as good as compared to other inorganic 1st and 2nd generation solar cells, it has an edge over them at some points.

In the normal operating temperature range of 25–65°C, DSSCs efficiency is nearly temperature-independent. In diffuse sunlight or cloudy conditions, DSSCs shows even better efficiency than polycrystalline Si solar cell. Performance is less sensitive to the incident angle of the light radiation. Although a mass commercial production of DSSCs is still not available, it can be expected that it has a cost advantage over all thin film devices. Only low cost and abundantly available materials are needed. DSSCs materials are biocompatible and abundantly available.

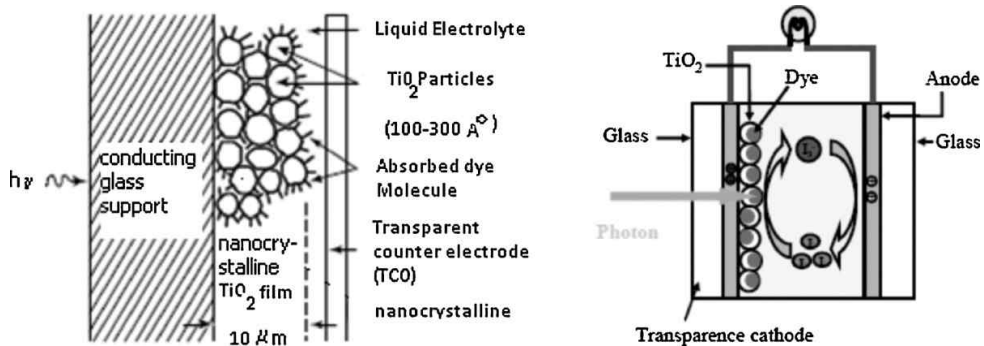


Figure 1. 8. Cross section of dye-sensitized solar cell [62, 67]

1. 2. 2. Components of DSSCs

DSSCs has four components, semiconductor, sensitizer(dye), electrolyte, counter electrode.

1. 2. 2. 1 Semiconductor

The semiconductor as a photoanode in DSSCs provides multi-function as supporter of sensitizer(dye) to load, and transporter of excited electrons from sensitizer to external circuit.

Among potential electron acceptors for DSSCs, TiO_2 is the most versatile. It delivers the highest efficiencies, is chemically stable, non-toxic, and available in large quantities. Generally, TiO_2 nanoparticles (NPs) have several crystal phase, such as rutile, anatase, and brookite. Its anatase nanocrystalline form gives the most efficiency

with optimization of surface area, porosity, pore diameter, transparency and film thickness. TiO_2 is the semiconductor of choice due to its wide bandgap, nanoporous structure with large surface area and, suitable HOMO and LUMO alignment with the electrolyte and dye respectively. The rutile phase has a bandgap of 3 eV, while 3.2 eV for anatase; corresponding to photon absorption edge of 413 and 388 nm. [60, 69, 70]

1. 2. 2. 2 Dye

The photo sensitizer dye is the heart of operation of DSSCs. The dye molecule commonly used in DSSCs are polypyridyl-type ruthenium (Ru) complexes such as namely N719 and N3 dye. [59, 60, 71, 72, 73] Upon sensitization of the semiconductor film, the dyes form a monomolecular coating, with the stoichiometry of the anchored photosensitizer as high as $\sim 1:100$ (Ru per Ti). [74] Stringent engineering of the dye is necessary for efficient light absorption, charge injection and collection of the cell. It is well-known that the Ru complexes absorb the visible light due to a metal to ligand charge transfer (MLCT) process. Furthermore, it has a proper excited and ground state energy levels, relatively long excited-state lifetime, and high chemical stability. [71]

1. 2. 2. 3 Electrolyte

In DSSCs, electrolyte act as reductor to regenerate of the oxidized dye. [75] The iodide/triiodide electrolyte system is commonly used. It's advantage is easy preparation, high conductivity, good permeability, and fast oxidation of I^- at the photoanode/electrolyte interface. [76, 78]

Incorporation of 4-tertbutylpyridine in the electrolyte increases the open-circuit voltage and fill factor by decreasing the dark current at the semiconductor-electrolyte junction.

The following relation holds for regenerative photo-electrochemical systems,

$$V_{oc} = \left(\frac{kT}{e} \right) \ln \left(\frac{I_{inj}}{n_{cb} K_{et} [I_3^-]} \right)$$

Where I_{inj} is the charge flux from sensitizer injection, K_{et} is the rate constant for triiodide reduction and n_{cb} is the electron concentration on the semiconductor surface. [60] Although the TiO_2 surface is covered by a dye monolayer, the reduction of triiodide by conduction band electrons causes the dark current. Due to its relatively small size the triiodide ions either cross the monolayer or have entrance to nanometer-sized pores into which the dye molecules cannot enter. In

the latter case, the surface of TiO_2 is exposed to redox mediator. 4-tert-butylpyridine effectively decreases the rate of the reduction of triiodide, increasing the open-circuit voltage of the cell. [75-83]

1. 2. 2. 4 Counter Electrode

Platinum coated FTO glass substrate is used to as the counter electrode for DSSCs. The counter electrode has a multiple role as a catalyst to regenerate the electrolyte and electron collector from external circuit. [84] Thus, the counter electrode should have advantages of high catalytic activity and electrical conductivity. [85] Pt counter electrode fulfilled this purpose.

RF sputtered Pt layer showed a good performance, however, for a nanoporous titania film fully optimized for efficient light absorption, chemically deposited platinum from a solution of 0.05 M hexachloroplatinic acid can perform alike.

The electrode substrate TCO affects the cell performance in 2 ways. The sheet resistance of the FTO influences the series resistance of the cell and the transmittance control the light absorption. Considering sheet resistance, the lower the better, but lower FTO sheet resistance require thicker FTO layer causing reduced light transmittance and efficiency. [95, 98-100]

1. 2. 3 Basic Operating Principle

The DSSCs is the useful photovoltaic device that utilizes separate mediums for light absorption/carrier generation (dye) and carrier transport (TiO_2 nano-particles).

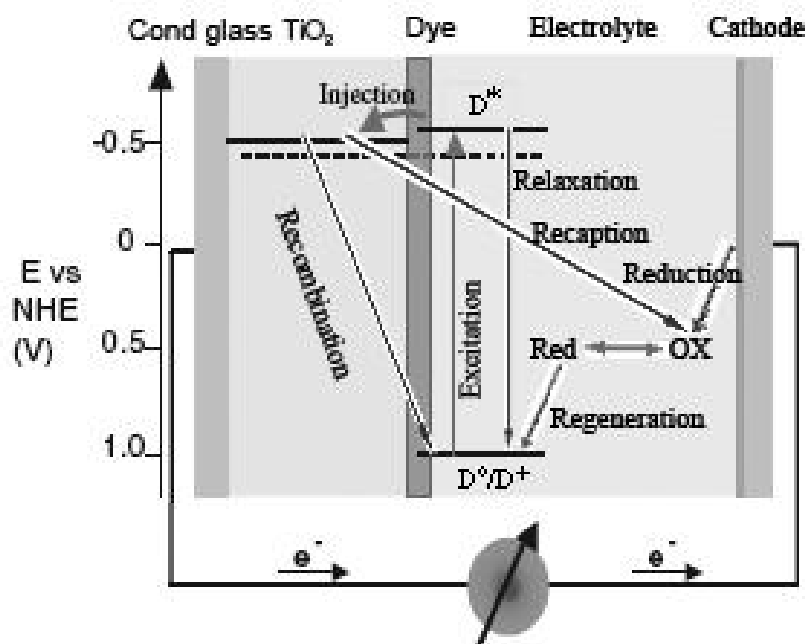


Fig. 1. 9. Principle of operation of DSSCs. [86]

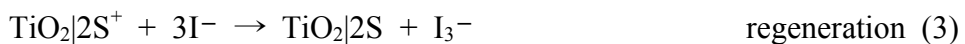
DSSCs include a substrate of fluorine-doped SnO_2 conducting glass (FTO), a porous nanocrystalline semiconductor oxide (the most employed is TiO_2) film sensitized by a dye (typically bipyridine ruthenium complexes) for absorbing visible light, a redox electrolyte (usually an organic solvent containing a redox system, such as iodide/triiodide couple) layer for deoxidizing oxidized dye, and a

platinized cathode to collect electrons and catalyze the redox couple regeneration reaction [58, 62, 65].

The light-to-electricity conversion in a DSSC is based on the injection of electron from the photoexcited state of the sensitized dye into the conduction band of TiO₂. The dye is regenerated by electron donation from iodide in the electrolyte. The iodide is restored, in turn, by the reduction of triiodide at the cathode, with the circuit being completed via electron migration through the external load.

The voltage generated under illumination corresponds to the difference between the Fermi level of the electron in the TiO₂ and the redox potential of the electrolyte. Overall, the device generates electric power from light without suffering any permanent chemical transformation [58, 62, 87, 88].

The photoelectric chemical process in DSSC can be expressed as equations. (1)–(6).





The photoexcited electron injects into the conduction band of TiO_2 . [71, 89-91] The dark reaction equation (5) and (6) also occur during the light-to-electricity conversion, but do not play a remarkable negative effect on photovoltaic performance of DSSCs owing to their slow reaction speed compared with that of equation (2) [92-94].

The operation steps are the following.

1. 2. 3. 1 Excitation

The light is absorbed by a sensitizer dye molecule, it goes over an electronic state change from the ground (S) to the excited state (S*). The lifetime of the excited state is in the order of nanoseconds.

1. 2. 3. 2 Injection

The sensitizing dye molecules are adsorbed on the surface of a wide band gap semiconductor (typically TiO_2). Upon absorption of a photon (excitation), the dye gains the ability to transfer an electron to the conduction band of the semiconductor. The internal electric field of the nanoparticles causes the electron extraction and the dye becomes

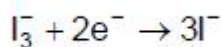
oxidized (S^+). The injection rate constant is in the femtosecond range for singlet state.

1. 2. 3. 3 Diffusion in TiO_2

The nonporous TiO_2 film consists of spherical anatase particles of diameter ~ 20 nm. The presence of oxygen vacancies in the lattice makes it a weakly n-doped material. [71, 72] As the TiO_2 particle diameter is too small for electric field to build up, the dominant electron transport mechanism is diffusion via trapping and de-trapping.

1. 2. 3. 4 Iodine Reduction

The electron travels through the outer circuit performing work, reaches the back FTO electrode, and reduces the iodine in the electrolyte. The platinum layer on the FTO acts as a catalyst for the reduction. The dark cathode reaction:

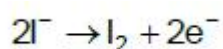


The iodine reduction can also occur at the excited dye molecules causing recombination of the photo-generated electrons. For efficient charge transfer, the rate of iodine reduction at the counter electrode has to be orders of magnitude faster than the recombination at the TiO_2 /electrolyte interface.

1. 2. 3. 5 Dye Regeneration

The reduced iodide ion replenishes the highest occupied molecular orbital (HOMO) of the dye - regenerating its original form, and makes it ready for electron generation again.

The photoanode reaction:



This prevents buildup of S^+ , which could lead to the conduction band electrons going back to the dye molecules. The maximum output voltage equals to the difference between the Fermi level of the semiconductor and the redox potential of the mediator. [72]

Thus, the device is can produce electricity from light without undergoing any permanent physical and chemical change.

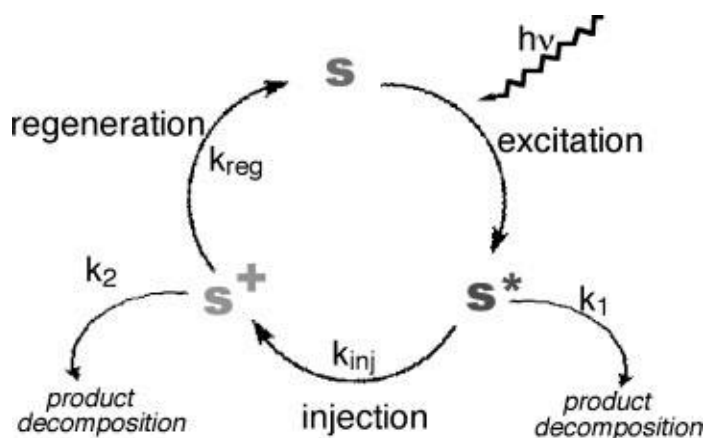


Figure 1. 10. The overall catalytic cycle of the sensitizer during DSSCs operation [68]

1. 2. 4 Solar Cell Terminologies

1. 2. 4. 1 Equivalent circuit of solar cells

A current source in parallel with a forward biased diode expresses the equivalent circuit of an ideal solar cell. Series and parallel resistances are added to account for various loss mechanisms.

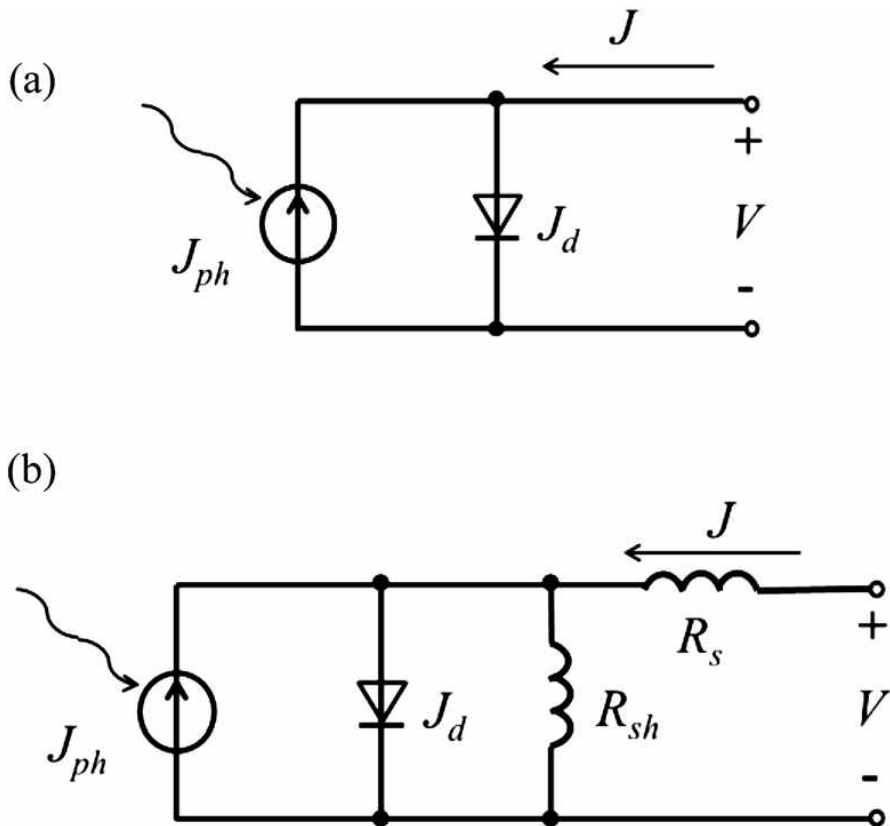


Figure 1. 11. Equivalent circuit of (a) an ideal solar cell and (b) a practical solar cell. [95]

Under illumination a solar cell works like a current source, providing a current density of J_{ph} . A proportion of this current counteracts the junction current of the diode, and the residual current flows to the load. According to Shockley's theory, the junction current density of the diode is presented by:

$$J_d = J_0[\exp(qV/k_B T) - 1]$$

(T : temperature, k_B : Boltzmann constant, J_0 : the reverse saturation current density (in a p-n junction solar cell, J_0 represents the current density of minority carriers, which is a summation of the hole current in the n region and the electron current in the p region), V : the output voltage)

Series resistance R_s and shunt resistance R_{sh} have to be taken into account in a practical solar cell, whose equivalent circuit is shown in Figure 1. 11(b). R_s is composed of the resistance of the bulk active layer, the resistance of electrodes and the contact resistance between the active layer and the electrodes, etc.

R_{sh} originates from various kinds of current leakage, such as current leakage in the p-n junction, current leakage from the edge of the cell, current leakage induced by impurities in the cell, etc.

1. 2. 4. 2 Short-circuit Current

In closed circuit, short-circuit Current is the current obtained from the cell when short-circuited or in other words when the load resistance is zero. Solar cell current is normally represented as current density, J_{sc} :

$$J_{sc} = \frac{I_{sc}}{A} \quad (\text{mA/cm}^2)$$

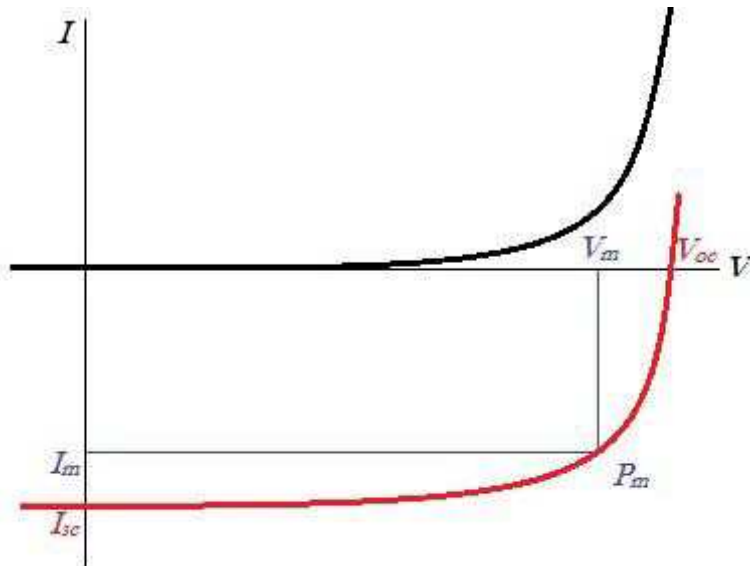
where A is the effective area of the solar cell. [69]

1. 2. 4. 3 Open-circuit Voltage

Open-circuit voltage is the maximum voltage available from a solar cell and is obtained when a load with infinite resistance is attached to its terminals. It is a function of the semiconductor bandgap and charge recombination in the cell. For DSC the V_{oc} is given by:

$$V_{oc} = \frac{E_{CB}}{q} + \frac{kT}{q} \ln\left(\frac{n}{N_{CB}}\right) - \frac{E_{redox}}{q} \quad (\text{volts})$$

where, n is the number of electron in TiO_2 conduction band and N_{CB} is the effective density of states. [96] The first two terms defines the quasi-fermi level of TiO_2 and E_{redox} is the Nernst potential of the redox mediator.



Light I-V response(below line) and dark I-V response(upper line).

Figure 1. 12. Typical current-voltage relationship of a solar cell. [95, 97]

1. 2. 4. 4 Series Resistance

Series resistance, R_s in a solar cell is the result of contact resistance and charge transfer resistance in the semiconductor material. Series resistance reduces the fill factor affecting the maximum power output, while excessively high value of R_s can also reduce the short-circuit current. The open-circuit voltage is not affected since, at V_{oc} the total

current flow through cell itself and hence through the series resistance is zero. An approximation of the series resistance can be determined from the slope of the IV curve at the open-circuit voltage point. [98, 99, 100]

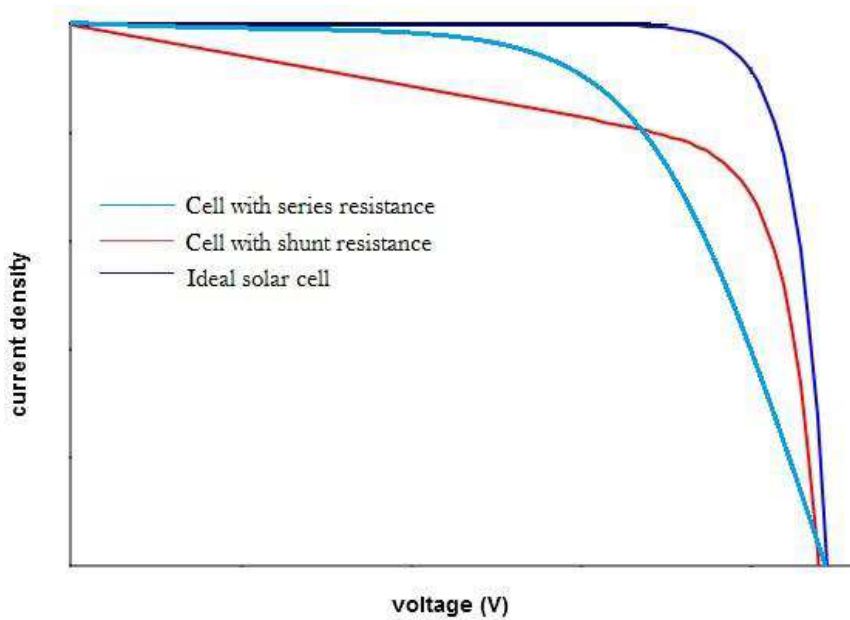


Figure 1. 13. Current-voltage response of a solar cell with series and shunt resistance.

1. 2. 4. 5 Shunt Resistance

Low shunt resistance, R_{sh} provides an alternate current path for the photo-generated current causing significant power loss. The effect of low shunt resistance is reduced fill factor and lower open-circuit voltage affecting the maximum power output. The short-circuit voltage is not affected unless for a very low value, since at J_{SC} the total current flows

through the outer path and hence through the shunt resistance is low. An approximation of the shunt resistance can be calculated from the slope of the I-V curve at the short circuit current point. [95]

1. 2. 4. 6 Fill Factor

The fill factor (FF) is a measure of the maximum power output from a solar cell. It represents the squareness of the I-V curve and is defined as the ratio of the maximum power to the product of V_{OC} and I_{SC} for the solar cell:

$$FF = \frac{V_m \times I_m}{V_{OC} \times I_{SC}}$$

Where, V_m and I_m are the voltage and current at maximum power point. Fill factor, being a ratio of the same physical parameters, has no unit. Fill factor is a function of the series and shunt resistance of the solar cell. For D_{SC} , it reflects the extent of electrical and

electrochemical losses during cell operation. To obtain higher fill factor improvement of the shunt resistance and decrement of the series resistance, with reduction of the overvoltage for diffusion and charge transfer is required. [101]

1. 2. 4. 7. Efficiency

The efficiency of a solar cell is defined as the ratio of maximum electrical energy output to the energy input from the sun. Thus the mathematical definition of Efficiency:

$$\eta = \left(\frac{V_{oc} \times I_{sc} \times FF}{P_{in}} \right)$$

where, P_{in} is the power input from the sunlight. Efficiency is generally expressed in percentage. [89]

1. 2. 4. 8. Quantum Efficiency

Quantum efficiency (QE) or ‘External Quantum Efficiency (EQE)’, sometimes also referred to as Incident Photon to Charge Carrier Efficiency (IPCE) is a measure of how efficient a solar cell is in producing photo-generated charge at a given frequency. It is defined as

the photocurrent density produced in the external circuit under illumination condition, divided by the photon flux of excitation wavelength. [71]

For DSSCs, the term is defined as:

$$\text{IPCE} = \frac{\text{photocurrent density}}{\text{wavelength} \times \text{photon flux}}$$

$$\text{IPCE}(\lambda) = \text{LHE}(\lambda) \times \Phi(\text{inj}) \times \eta(\text{coll})$$

Where, $\text{LHE}(\lambda)$ is the light-harvesting efficiency at excitation wavelength λ , $\Phi(\text{inj})$ is the electron injection quantum yield for the excited sensitizer to the semiconductor oxide conduction band and $\eta(\text{coll})$ is the efficiency for the collection of electrons. [71, 102]

1. 3 Additional strategies to enhance the efficiency of DSSCs

1. 3. 1 Blocking Layer

Charge recombination is one of the main reasons of lower current for J_{SC} . Recombination occurs at both Electrode/Electrolyte and TiO_2 /Electrolyte interface. A compact blocking layer of TiO_2 by RF sputtering [103] or spray pyrolysis [104] between the conducting electrode and the nanocrystalline TiO_2 layer can effectively prevent the recombination at electrode/electrolyte interface. Introduction of this layer prevents electrolyte from reaching the electrode (increases J_{SC}) and also enhance electron transport from nanocrystalline titania to the electrode (increases V_{OC}). This blocking layer is absolutely necessary for planar organic dyes, while ruthenium based sensitizers can perform this anode insulation themselves against recombination losses. [105, 106]

A rather cost effective and easier way of incorporating this blocking layer is through hydrolysis of $TiCl_4$. $TiCl_4$ treatment on FTO substrate is found to suppress the dark current, shifting its commencement by some hundred millivolts. This is due to a positive shift in the conduction band edge of highly doped SnO_2 by about 0.5 V which results in a higher electron density in the FTO substrate. [107]

1. 3. 2 TiCl₄ Treatment

To reduce charge recombination at TiO₂/electrolyte interface, another blocking layer of TiO₂ is implemented via TiCl₄ treatment with subsequent annealing. The average pore diameter and specific surface area decreases with the TiCl₄ treatment, because of the surface epitaxial growth of TiO₂, resulting in particle-necking and a new TiO₂ layer. The increase of the inter-particle necking will reduce the resistance in electron transport through the nanoparticles reducing the recombination probability. At the same time, it improves the surface roughness factor and light absorbance. [107] These lead to a decrease in the dark current and an increase of the photocurrent. [108]

Hydrolysis from the acidic TiCl₄ solution prevents deposition of impurities in the TiCl₄, such as Fe³⁺, because of the higher solubility of iron oxide compared to TiO₂. The P25 powder, which is the most common source of producing nanocrystalline Titania, contains as far as 100 ppm of Fe₂O₃, which is known for causing interference with electron injection from the excited state of the dye. The TiCl₄ treatment covers this relatively impure semiconductor core with a TiO₂ thin layer of ultra-high purity. This improves the injection efficiency and the blocking character of the semiconductor-electrolyte junction. [60, 83]

Theoretically, this overlayer can be formed of other metal-oxide films (examples include MgO, ZnO, Al₂O₃, and Nb₂O₅) which have conduction band edges higher than TiO₂. [109]

The performance of these oxides depends on their protonation / deprotonation capability of the TiO₂ surface, which enhances dye adsorption and a positive shift of the TiO₂ flat band. [110] For example, Al₂O₃ being basic, will tend to deprotonate the TiO₂ film, resulting in the improvement in cell performance.

1. 3. 3 Light Scattering Layer

One limitation of the sensitizing dyes is their poor performance in the near infra-red spectrum of light. A way of improving this is - by introducing an additional light scattering layer of larger titania particles. These can be mixed with or screen-printed on top of the film of 15-30nm sized TiO₂ particles. This allows the scattered photons to be contained in the film by means of multiple reflections, increasing their optical path length substantially beyond the film thickness. Consequently, the solar light absorption is enhanced, especially in the red to near-IR regions. With the use of 200-400nm sized anatase particles as light-scattering centers, an increment of the J_{SC} by 3-4

mA/cm² was observed for N719-based DSCs because of the enhanced light absorption.

Moreover, approximately 4% loss occurs because of the reflection of incident light on the glass substrate. This can be partially overcome by introducing an anti-reflecting film, which can also act as a UV cut-off filter. [107, 111-131]

1. 4 Plasmonic DSSCs

1. 4. 1. Surface plasmon resonance for solar cells

Plasmonic solar cells are photovoltaic devices that use the noble metal surface plasmons. Surface plasmons are the collective oscillation of excited free electrons.

Most of solar cells have a weak absorber. To trap the more absorption and long wavelength of light, the morphology of substrate is etched like pyramids with a size of 2-10 μm that is about wavelength-scale texture. This structure not only increases the light trapping to solar cell, but also increases the surface recombination and low material quality. [132, 133]

Another is to use the noble metal nanoparticles in the solar cells. The

light is absorbed and scattered from noble metal nanoparticles that is excited at the surface plasmon resonance. It has been based on Raman scattering that can be increased by order of magnitude when the metal nanoparticles are used [134, 135].

Surface plasmon resonance (SPRs) are coherent oscillations of free electrons at the boundaries between metal and dielectric which are often categorized into two classes: (1) propagating surface plasmons (PSPs), which generate at the metal film and (2) localized surface plasmons (LSPs), which generate at the metal NPs. [136-139]

In the case the case of PSPs, surface electromagnetic field propagates along the dielectric interface of metal in the x-/y-directions, but decays exponentially in the z-direction [137, 140, 141]

Whereas in LSPs, metallic NPs interact with light when they are much smaller than the incident wavelength, leading to a polarization at the surface of metal NPs. Consequently, induce a strong enhancement of electromagnetic field in the near-field region (resonance amplification), polarization make oscillate electrons locally at the surface of metal NPs. Therefore, light absorption can be enhanced strongly, as depicted in figure [137, 140, 141]

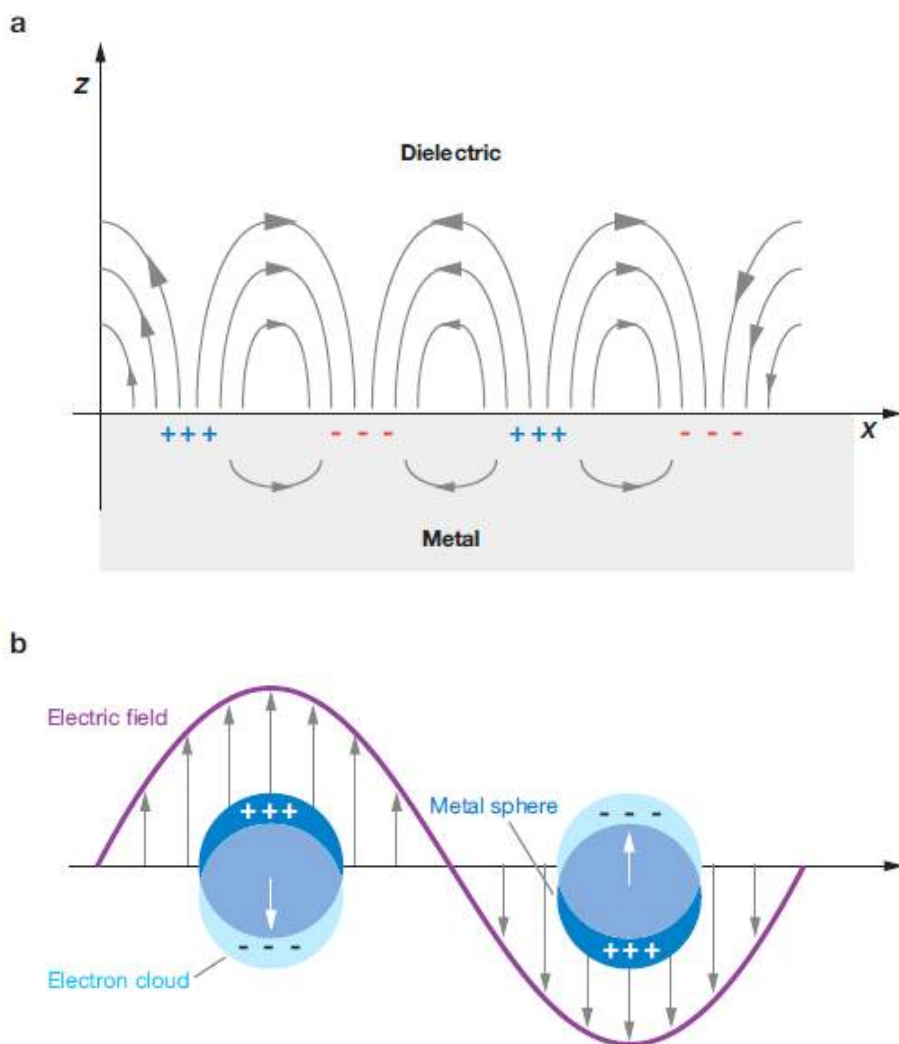


Figura 1. 14. Schematic diagrams illustrating (a) a surface plasmon polariton (or propagating plasmon) and (b) a localized surface plasmon. [140]

1. 4. 2 Localized surface plasmon resonance

Localized surface plasmons (LSPs) are associated with the collective oscillations of electrons confined locally by metal nanostructures. The most representative examples of LSPs are metal NPs. The particle plasmons are excited when the frequency of the incident photons matches the resonance frequency of the NPs. Their resonance wavelength depends on the particle shape, size, and the dielectric parameters of the surrounding environment. From the quasistatic approximation, the polarizability (P) of a spherical NP can be expressed as:

$$P = 4\pi a^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m} \quad [142]$$

Where a is the diameter of the NP and ε and ε_m are the dielectric constants of the surrounding dielectric medium and of the metal NP itself, respectively. [142] From equation, we conclude that the value of P reaches its maximum when ε_m is equal to $-2\varepsilon_m$, resulting in a resonance condition. Because the excited plasmons are localized and cannot propagate within the nanostructure, this process is known as localized surface plasmon resonance (LSPR).

1. 4. 3. Plasmonic effects in solar cells

Proposal mechanism of plasmonic solar cells is to explain photocurrent improvement by metal nanoparticles that are incorporated into or on solar cells as shown in Figure 1. 15. One is scattering from the metal nanoparticles. The role of metal nanoparticles is used as subwavelength scattering elements to couple and trap propagating plane waves from the Sun into semiconductor thin film, by folding the light into a thin absorber layer as shown in Figure 1. 15. (a). In inorganic plasmonic solar cells, the photocurrent improvement is increased by scattering from metal nanoparticles. Another is near-field enhancement from metal nanoparticles. Metal nanoparticles are used as subwavelength antennas in which the plasmonic near field is coupled to the semiconductor, increasing its effective absorption cross-section as shown in Figure 1. 15. (b). In organic plasmonic solar cells, the photocurrent improvement is increased by near-field enhancement. The other is direct generation of charge carriers in the semiconductor substrate. A corrugated metallic film on the back surface of a thin photovoltaic absorber layer can couple sunlight into surface plasmon polaritons (SPP) modes. Surface plasmon polaritons are supported at the metal/semiconductor interface as well as guided modes in the semiconductor slab. Accordingly, the light is converted to photocarriers in the semiconductor as shown in Figure 1. 15. (c). [143].

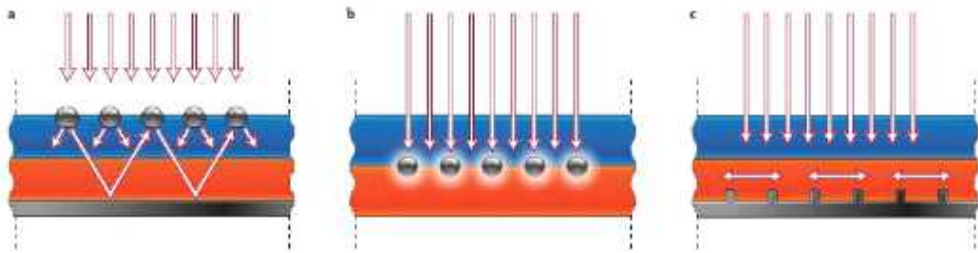


Figure 1. 15. Plasmonic light-trapping geometries for thin-film solar cells. (a) Light trapping by scattering from metal nanoparticles at the surface of the solar cell. Light is preferentially scattered and trapped into the semiconductor thin film by multiple and high-angle scattering, causing an increase in the effective optical path length in the cell. (b) Light trapping by the excitation of localized surface plasmons in metal nanoparticles is embedded in the semiconductor. The excited particles' near-field causes the creation of electron-hole pairs in the semiconductor. (c) Light trapping by the excitation of surface plasmon polaritons at the metal/semiconductor interface. A corrugated metal back surface couples light to surface plasmon polariton or photonic modes that propagate in the plane of the semiconductor layer [143].

Chapter 2 Experimental Section

2. 1. Preparation of TiO₂ nanotubes arrays

2. 1. 1. Materials

Titanium plate (Alfa aesar, 99.7% purity, 0.25 mm thickness), ammonium fluoride (NH₄F, Showa Chemical Co., 97.0%), ethylene glycol (Daejung chemical, 99%), hydrogen peroxide (H₂O₂, Daejung chemical, 30%), fluorine-doped tin oxide(FTO) glass (Pilkington, TEC-A7), Titanium diisopropoxide bis(acetylacetonate) solution (Aldrich, 75 wt% in isopropanol), n-butanol (Daejung chemical, 99%), TiO₂ paste (Ti-Nanoxide T/SP, solaronix), scattering TiO₂ paste (18NR-AO, Dyesol), silver nitrate (AgNO₃, Aldrich, 99%) titanium chloride (TiCl₄, Aldrich, 0.09 M in 20% HCl), Dye (cis-diisothiocyanato-bis(2,2'-bipyridyl-4,4'-dicarboxylato) ruthenium(II) bis(tetrabutylammonium), N719, Solaronix), chloroplatinic acid hexahydrate(H₂PtCl₆·6H₂O, Aldrich), 1-butyl-3-methyl-imidazolium iodide (BMII, Aldrich, 99%), iodine (I₂, Aldrich, 99%), guanidium thiocyanate (GSCN, Aldrich, 99%), 4-tertbutylpyridine (TBP, Aldrich, 96%), acetonitrile (CH₃CN, Aldrich, 99.8%), valeronitrile (CH₃(CH₂)₃CN, Aldrich, 99.5%)

2. 1. 2. Titanium (Ti) Anodization

TiO₂ nanotubes were fabricated by anodizing thin Ti plates (99.7 % purity, Alpha, 2.5 cm x 4.0 cm x 200 μm) in an electrolyte composed of 0.8 wt% of NH₄F and 2 vol% of H₂O in ethylene glycol at 25°C and at a constant applied voltage of 60 V DC for 2 h as shown in Figure 2. 1. [45, 144, 145].

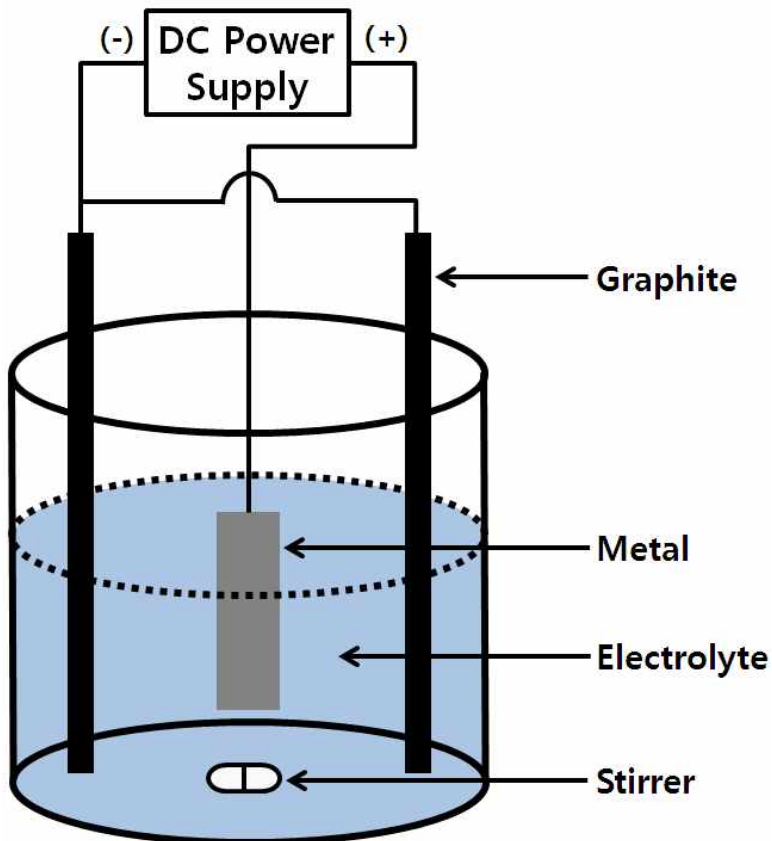


Figure 2. 1. Schematic diagram of titanium (Ti) anodization.

2. 1. 3. Preparation of free-standing crystallized TiO₂ nanotubes

The Ti plate-formed TiO₂ nanotubes arrays were annealed at 450°C for 1 h under an ambient atmosphere to improve their crystallinity. To detach the free-standing TiO₂ nanotube arrays from the Ti plate, a second anodization was done at a constant applied voltage of 30 V DC for 10 min, and then the plate was immersed in the 10 % of H₂O₂ for 24 h [146].

2. 1. 4. Preparation of free-standing crystallized TiO₂ nanotube arrays

To prepare TiO₂ nanotube arrays, the barrier layer of TiO₂ nanotubes was etched. However, free-standing crystallized TiO₂ nanotubes was very stable in chemicals. So the barrier layer was not etched by acidic or basic condition.

The barrier layer of TiO₂ nanotubes was removed by ion milling with Ar⁺ bombardment for several minutes to prepare open-ended freestanding TiO₂ nanotube arrays. [146, 147].

2. 2. Fabrication of dye-sensitized solar cells based on TiO₂ nanotubes array with scattering layer

2. 2. 1. Overview of DSSCs manufacturing

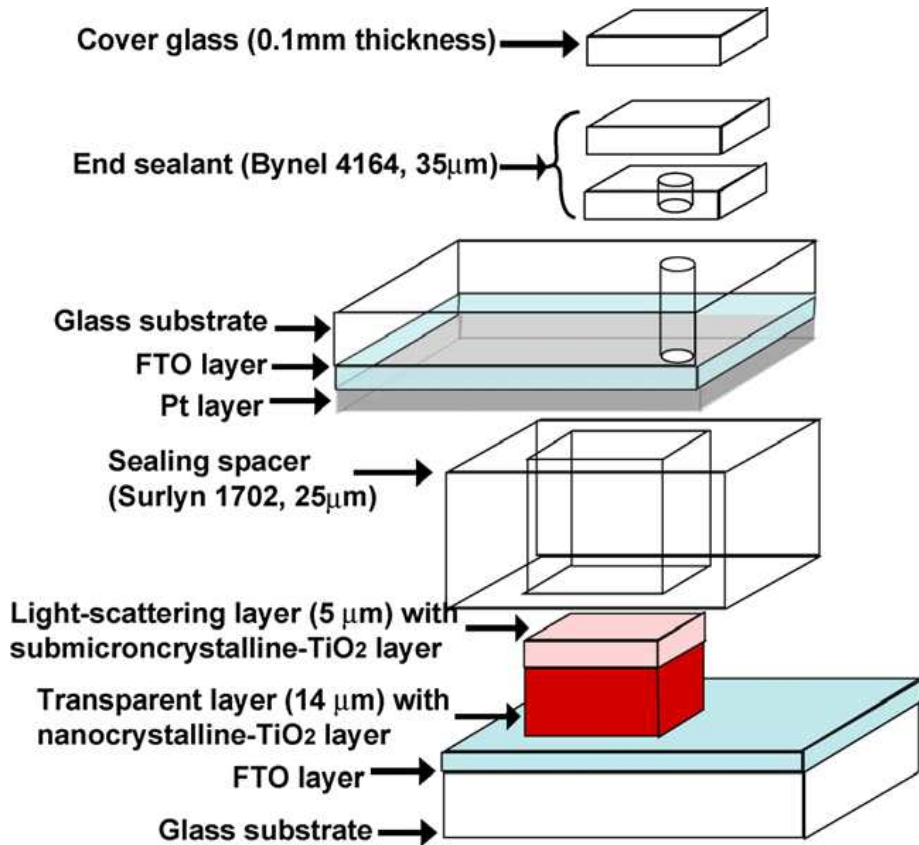


Figure 2. 2. Configuration of the dye sensitized solar cells. [107]

2. 2. 2. Preparation of blocking layer

A TiO₂ blocking layer was formed on fluorine-doped tin oxide (FTO) glass by spin-coating with 5 wt% of titanium di-isopropoxide bis(acetylacetonate) in butanol and then by heating at 450°C for 30 min under ambient conditions. [106].

2. 2. 3. Introduction to TiO₂ nanotubes arrays on FTO glass

TiO₂ nanotubes or nanomembranes were attached on the fluorine-doped thin oxide (FTO) glass formed a TiO₂ blocking layer on its surface by using TiO₂ nanoparticles viscous paste. The viscous TiO₂ paste (from Solaronix) was printed onto FTO glass by the doctor blade technique and the closed- and open-ended TiO₂ nanotube arrays were introduced on the paste. The substrate was then sintered at 450°C for 1 h under ambient conditions.

2. 2. 4 Adapting scattering layer

The TiO₂ scattering layer (~400nm-diameter particles) was coated onto the closed- and open-ended TiO₂ nanotube arrays using a doctor blade and sintered at 450°C for 30 min under ambient conditions.

2. 2. 5. TiCl₄ treatment

TiO₂ nanotubes arrays on FTO glass were treated with TiCl₄ by dipping the plate in a 0.1 M TiCl₄ aqueous solution for 30 min and then sintered at 450°C for 30 min. By the sintering process, the crystal structure of TiO₂ nanoparticles was changed to anatase.

2. 2. 6. Dye absorption

Dye molecules were absorbed by immersing the nanotube film-attached FTO glass plate in an ethanol solution of 0.5 mM cis-diisothiocyanato-bis(2,2'-bipyridyl-4,4'-dicarboxylato) ruthenium(II) bis (tetrabutylammonium), (N719, Solaronix), at 50°C over 8 h.

2. 2. 7. Electrolyte

The composition of the electrolyte was as follows: 0.7 M 1-butyl-3-methyl-imidazolium iodide (BMII), 0.03 M I₂, 0.1 M guanidium thiocyanate (GSCN), and 0.5 M 4-tertbutylpyridine (TBP) in a mixture of acetonitrile and valeronitrile (85:15 v/v).

2. 2. 8. Counter electrode

The counter-electrode was prepared by sputtering Pt on FTO glass or prepared by spin-coating the H₂PtCl₆ in isopropanol on FTO glass and then sintering at 400°C for 20 min [106].

2. 2. 9. Fabrication of DSSCs

The working electrode was further sandwiched with the Pt-coated FTO glass, separated by a 60- μ m-thick hot-melt spacer.

2. 3. Fabrication of DSSCs with freestanding TiO₂ nanotube arrays with channels containing Ag NPs

2. 3. 1. Same the former process of free standing TiO₂ nanotube arrays with scattering layer

Former process of free standing TiO₂ nanotube arrays with scattering layer, preparation of blocking layer and introduction to TiO₂ nanotube arrays on FTO glass, are same.

2. 3. 2. Ag nanoparticle formation

The substrate was dipped in 0.3 mM AgNO₃ aqueous solution and exposed to 254 nm UV irradiation for reduction of Ag cations.

2. 3. 3. Same the latter process of free standing TiO₂ nanotube arrays with scattering layer

Latter process of free standing TiO₂ nanotube arrays with scattering layer, adapting scattering layer, TiCl₄ treatment, dye absorption,

electrolyte filling, preparation of counter electrode and fabrication of DSSCs, are same.

2. 4. Characterization

2. 4. 1. FE-SEM

The morphology and thickness of the free-standing TiO_2 nanotube arrays, which were TiO_2 nanotube arrays after detachment from the Ti plate, were analyzed using a field emission scanning electron microscope (FE-SEM, JSM-6330F, JEOL Inc.).

2. 4. 2. TEM

The morphology, thickness, size, and presence of Ag NPs in the channel of freestanding TiO_2 nanotube arrays were confirmed using a field emission scanning electron microscope (FE-SEM, JSM-6330F, JEOL Inc.) and the high angular annular dark field (HAADF) technique with a scanning transmission electron microscope (TEM) (JEM-2200FS, JEOL).

2. 4. 3. Solar simulator

The current density–voltage (I – V) characteristics of the DSSCs were measured using an electrometer (KEITHLEY 2400) under AM 1.5 illumination (100 mW/cm^2) provided by a solar simulator (1 kW xenon with AM 1.5 filter, PEC-L01, Peccel Technologies).

2. 4. 4. Incident photon-to-current conversion efficiency (IPCE)

The incident photon-to-current conversion efficiency (IPCE) was measured using McScience (model K3100) with reference to the calibrated diode. A 300 W xenon lamp was used as light source for generation of a monochromatic beam. The bias light was supplied by a 150 W halogen lamp.

Chapter 3

Improved energy conversion efficiency of dye-sensitized solar cells fabricated using open-ended TiO_2 nanotube arrays with scattering layer

3. 1 Overall view of free standing TiO₂ nanotube arrays DSSCs

Dye-sensitized solar cells (DSSCs) have attracted immense interest due to their high energy conversion efficiency and low cost [58, 62, 107]. However, the energy conversion efficiency of DSSCs still needs to be improved so that it compares favorably with conventional photovoltaic devices [148]. There are several parameters that can be investigated, including the dimensionality of TiO₂ for electron transport [149, 150], light-harvesting capability [151, 152], molar absorption coefficient [153], energetically suitable HOMO–LUMO levels [154], available surface area for dyes [155], transport kinetics of the electrons [156], regeneration by a redox couple [157], and losses due to recombination and back reactions [158].

TiO₂ nanotubes can enhance electron transport and charge separation by creating direct pathways and accelerating the charge transfer between interfaces [144-146, 159-163]. These properties make them an attractive candidate for DSSC applications. TiO₂ nanotube arrays that are prepared by electrochemical anodization have a highly oriented and vertically aligned tubular structure [155, 156]. Thus, the arrays have a high degree of electron transport and minor charge recombination in comparison to TiO₂ nanoparticle films [164]. Hence, although current DSSCs fabricated using TiO₂ nanotube arrays have a low energy conversion efficiency as compared to DSSCs fabricated using TiO₂

nanoparticle films, they have immense potential. Recently, we prepared DSSCs using open-ended TiO₂ nanotube arrays and demonstrated that nanotube arrays whose barrier layers were removed by ion milling have 24% higher energy conversion efficiency as compared with DSSCs fabricated using TiO₂ nanoparticle films [163].

Introducing a scattering layer such as TiO₂, ZrO₂, or SiO₂ can increase the total energy conversion efficiency of DSSCs [165]. TiO₂ is a good material to use for a scattering layer due to its chemical stability and dye adsorption capability; hence, several DSSCs fabricated using TiO₂ nanoparticle films use a TiO₂ scattering layer on the active layer.

To the best of our knowledge, TiO₂ nanotube arrays have not been combined with scattering layers. In this paper, we report the improved energy conversion efficiency of DSSCs using open-ended TiO₂ nanotube arrays with a TiO₂ scattering layer. In this study, we compared the energy conversion efficiency of 1) closed- and open-ended TiO₂ nanotube arrays 2) with and without a TiO₂ scattering layer. In addition, we compared the energy conversion efficiency of fabricated DSSCs treated with TiCl₄ to untreated DSSCs.

3. 2. Characteristics of closed- and open-ended TiO₂ nanotube arrays

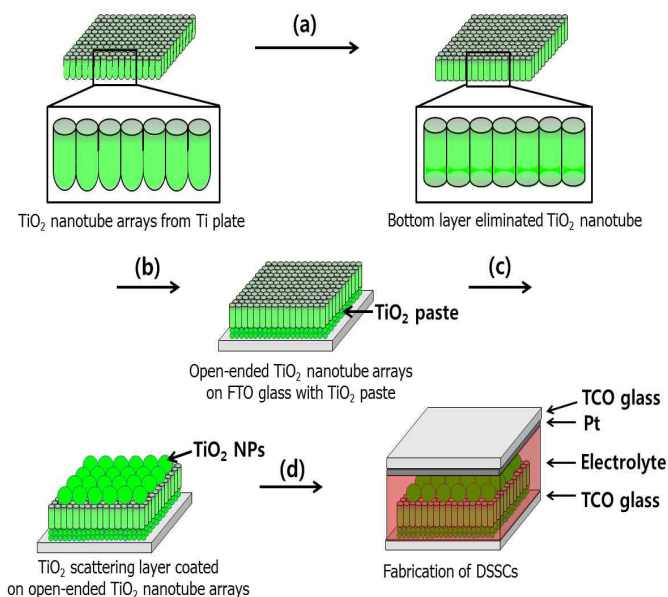


Figure 3. 1. Overall scheme of fabrication of DSSCs using the closed- or open-ended TiO₂ nanotube arrays. (a) Elimination of the bottom layer of closed-ended TiO₂ nanotube arrays by ion milling, (b) introduction of the closed- or open-ended TiO₂ nanotube arrays on FTO glass with TiO₂ paste, (c) coating the TiO₂ scattering layer on closed-ended or open-ended TiO₂ nanotube arrays by doctor blade, and (d) fabrication of DSSCs.

Figure 3. 1. shows the fabrication flow of the DSSCs using the closed- and open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer for improved energy conversion efficiency. After sintering at 450°C for 1 h under ambient conditions, the TiO₂ nanotube arrays have a crystalline form similar to anatase.

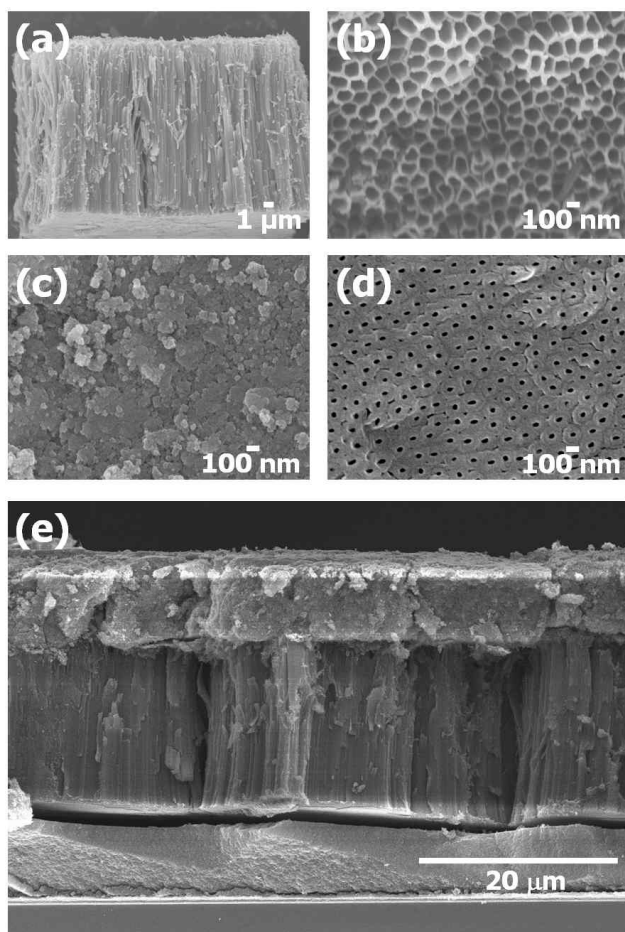


Figure 3. 2. FE-SEM images of (a) side view of TiO_2 nanotube arrays, (b) top view of TiO_2 nanotube arrays, (c) bottom view of the closed-ended TiO_2 nanotube arrays, (d) bottom view of open-ended TiO_2 nanotube arrays after ion milling, and (e) side view of TiO_2 nanotube arrays on FTO glass with TiO_2 nanoparticles and TiO_2 scattering layer.

SEM images of the side, top, and bottom of the free-standing TiO₂ nanotube arrays are shown in Figures 3. 2a, b, and c, respectively. The length of the free-standing TiO₂ nanotube arrays was approximately 18 μm, as shown in Figure 3. 2a. The upper pores were well ordered and their diameter was ca. 100 nm. The morphology of the bottom layer under the closed-ended TiO₂ nanotube array was very rough due to chemical etching with the H₂O₂ solution after secondary anodization, as shown in Figure 2c.

To prepare the open-ended TiO₂ nanotube arrays, the bottom layer of the closed-ended TiO₂ nanotube array was eliminated by ion milling to remove the barrier layer. Most of the bottom tips were opened after ion milling for 90 min and they had an approximate diameter of 20 nm, as shown in Figure 3. 2d. Figure 3. 2e shows the closed- and open-ended TiO₂ nanotube arrays on FTO glass after attachment using a TiO₂ paste and sintering at 450°C for 1 h under ambient conditions. The TiO₂ scattering layer was coated onto the closed- and open-ended TiO₂ nanotube arrays using a doctor blade and then the dye (N719) was adsorbed. DSSCs were fabricated by assembling the working electrode (the closed- and open-ended TiO₂ nanotube arrays with TiO₂ scattering layer) and the counter electrode (Pt).

3. 3 Performance Measurements of the DSSCs which have closed-/open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer

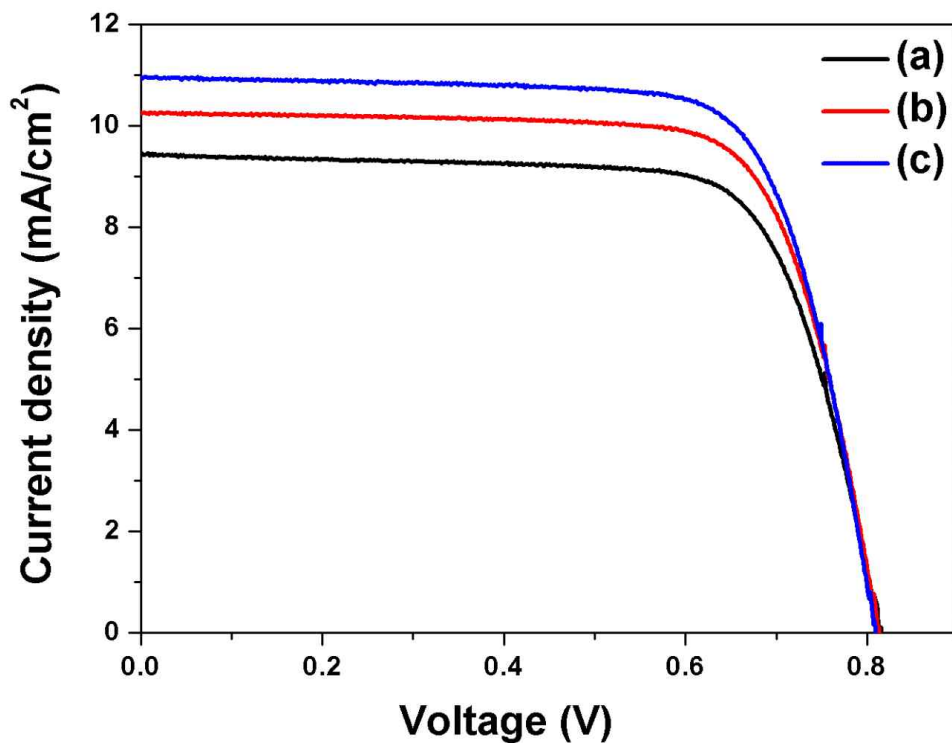


Figure 3. 3. I-V curves of DSSCs fabricated using (a) the closed-ended TiO₂ nanotube arrays, (b) the closed-ended TiO₂ nanotube arrays with the TiO₂ scattering layer, and (c) the closed-ended TiO₂ nanotube arrays with the TiO₂ scattering layer treated with TiCl₄.

Figure 3. 3. presents the current density–voltage curves of three different DSSCs fabricated using the closed-ended TiO₂ nanotube arrays attached to the FTO glass using TiO₂ paste. The measurements were taken using AM 1.5-simulated sunlight.

Table 3. 1. Photovoltaic properties of DSSCs fabricated using the closed-ended TiO₂ nanotube arrays

	J _{sc} (mA/cm ²)	V _{oc} (V)	ff	η (%)
Closed-ended TiO ₂ nanotube arrays	9.44	0.82	0.73	5.63± 0.14
Closed-ended TiO ₂ nanotube arrays with TiO ₂ scattering layer	10.24	0.81	0.74	6.17± 0.18
Closed-ended TiO ₂ nanotube arrays with TiO ₂ scattering layer treated with TiCl ₄ .	10.96	0.81	0.74	6.54± 0.20

The values of the open-circuit voltage (V_{oc}), short-circuit current (J_{sc}), fill factor (ff), and energy conversion efficiency (h) are summarized in Table 3. 1. For the DSSC fabricated using just the

closed-ended TiO₂ nanotube arrays, the energy conversion efficiency was $5.63 \pm 0.14\%$. For the DSSC fabricated using the closed-ended TiO₂ nanotube arrays and the TiO₂ scattering layer, the energy conversion efficiency was $6.17 \pm 0.18\%$.

By introducing the TiO₂ scattering layer on the closed-ended TiO₂ nanotube arrays, the energy conversion efficiency improved significantly, with a 9.59% enhancement. When the closed-ended TiO₂ nanotube arrays with the TiO₂ scattering layer were treated with TiCl₄, the energy conversion efficiency increased from $5.63 \pm 0.14\%$ to $6.54 \pm 0.20\%$, corresponding to a 16.2% enhancement due to increasing dye adsorption on the surface of the TiO₂ nanotube arrays [146].

By introducing the TiO₂ scattering layer on the closed-ended TiO₂ nanotube arrays, the energy conversion efficiency was improved due to increased light harvesting by the scattering layer.

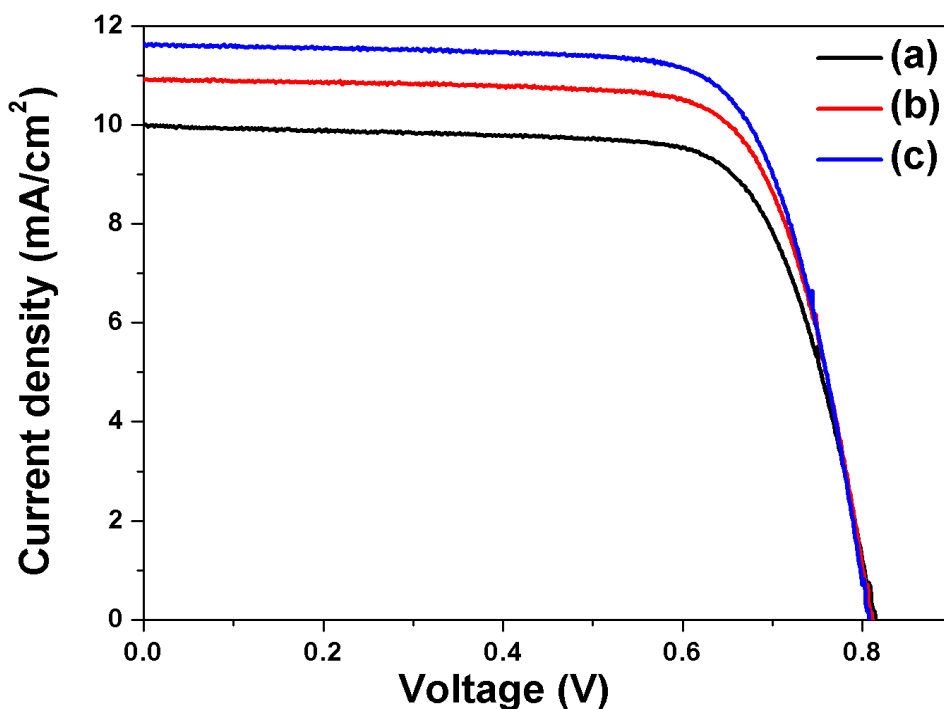


Figure 3. 4. I-V curves of DSSCs fabricated using (a) the open-ended TiO₂ nanotube arrays, (b) the open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer, and (c) the open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer treated with TiCl₄.

Figure 3. 4. presents the current density-voltage curves of three different DSSCs fabricated using the open-ended TiO₂ nanotube arrays attached to the FTO glass using TiO₂ paste.

Table 3. 2. Photovoltaic properties of DSSCs fabricated using open-ended TiO₂ nanotube arrays

	Jsc (mA/cm ²)	Voc (V)	ff	η (%)
Open-ended TiO ₂ nanotube arrays	9.99	0.82	0.73	5.92 ± 0.19
Open-ended TiO ₂ nanotube arrays with TiO ₂ scattering layer	10.92	0.81	0.74	6.53 ± 0.13
Open-ended TiO ₂ nanotube arrays with TiO ₂ scattering layer treated with TiCl ₄	11.63	0.81	0.73	6.89 ± 0.16

The values of Voc, Jsc, ff, and h are summarized in Table 3. 2. For the DSSC fabricated using just the open-ended TiO₂ nanotube arrays, the energy conversion efficiency was 5.92 ± 0.19%. For the DSSC fabricated using the open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer, the energy conversion efficiency was 6.53 ± 0.13%, a 10.30% enhancement. When the open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer were treated with TiCl₄, the energy

conversion efficiency improved from $5.92 \pm 0.19\%$ to $6.89 \pm 0.16\%$, corresponding to a 16.38% enhancement.

The energy conversion efficiency increased from $5.63 \pm 0.14\%$ for the DSSCs with the closed-ended TiO_2 nanotube arrays to $5.92 \pm 0.19\%$ for the DSSCs with open-ended TiO_2 nanotube arrays, a 5.15% enhancement. With the introduction of the scattering layer, the efficiency increased from $6.17 \pm 0.18\%$ for the DSSCs with the closed-ended TiO_2 nanotube arrays to $6.53 \pm 0.13\%$ for the DSSCs with open-ended TiO_2 nanotube arrays, an improvement of 5.83%.

Upon treatment with TiCl_4 , the enhancement was 5.35%, from $6.54 \pm 0.20\%$ for the DSSCs with the closed-ended TiO_2 nanotube arrays to $6.89 \pm 0.16\%$ for the DSSCs with open-ended TiO_2 nanotube arrays.

In previous our works, the barrier layer in the closed-ended TiO_2 nanotube arrays affected the electron transport in the DSSCs, so the barrier layer was removed by ion milling in order to prepare the open-ended TiO_2 nanotube arrays [163].

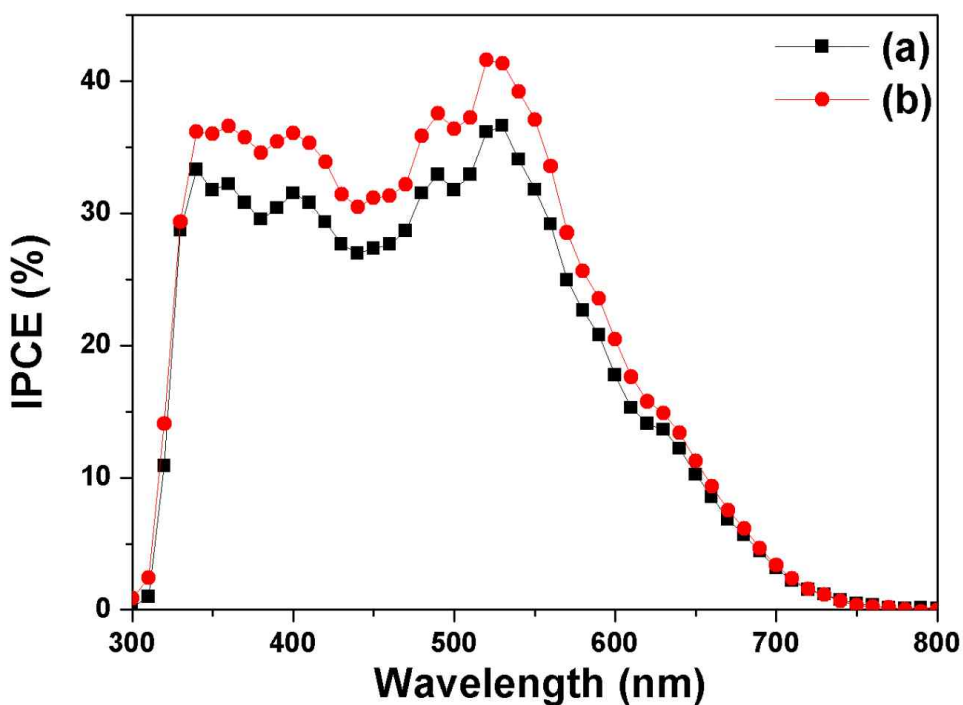


Figure 3. 5. IPCE spectra of DSSCs fabricated using (a) the open-ended TiO₂ nanotube arrays and (b) the open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer.

The IPCE spectra of the DSSCs fabricated using the open-ended TiO₂ nanotube arrays and open-ended TiO₂ nanotube arrays with the TiO₂ scattering layer are shown in Figure 3. 5. The IPCE spectra are similar but the DSSC with the TiO₂ scattering layer had a higher intensity.

3. 4. Conclusions

We prepared dye-sensitized solar cells (DSSCs) with enhanced energy conversion efficiency using open-ended TiO_2 nanotube arrays with a TiO_2 scattering layer.

As compared to closed-ended TiO_2 nanotube arrays, the energy conversion efficiency of the open-ended TiO_2 nanotube arrays was increased from 5.63% to 5.92%, which is an enhancement of 5.15%. Due to the removal of the barrier layer, which was present in the closed-ended TiO_2 nanotube arrays, causing an improvement in electron transport.

By introducing the TiO_2 scattering layer on the open-ended TiO_2 nanotube arrays, the energy conversion efficiency was increased from 5.92% to 6.53%, which is an enhancement of 10.30%.

After treating the open-ended TiO_2 nanotube arrays with TiCl_4 , the energy conversion efficiency was increased from 5.92% to 6.89%, a 16.38% enhancement, which is attributed to improved light harvesting and increased dye adsorption.

Chapter 4

Ag nanoparticle-functionalized
open-ended freestanding TiO₂
nanotube arrays with scattering
layer for improved energy
conversion efficiency in
dye-sensitized solar cells

4. 1. Overall view of Ag nanoparticle-functionalized open-ended freestanding TiO₂ nanotube arrays with scattering layer for improved energy conversion efficiency in dye-sensitized solar cells

Since the original work by O'Regan and Grätzel in 1991 [58], dye-sensitized solar cells (DSSCs) have been investigated extensively because of their high energy conversion efficiency and low cost.[168-175] Generally, mesoporous TiO₂ nanoparticle (NP) films and ruthenium sensitizers are used for DSSCs. [64, 168-181] However, the efficiency of mesoporous TiO₂ NP film-based DSSC is limited by grain boundaries, defects, and numerous trapping sites. Moreover, mesoporous TiO₂ NP films can cause charge recombination and mobility. [82, 182]

TiO₂ nanotubes, which enhance electron transport and charge separation by creating direct pathways and accelerating charge transfer between interfaces, have great potential to overcome the problems of mesoporous TiO₂ NP films. [183-186] TiO₂ nanotubes can be prepared by a hydrothermal method [187] or an electrochemical method [18], known as anodization. TiO₂ nanotube arrays prepared by anodization have a well-ordered and vertically oriented-tubular structure that facilitates a high degree of electron transport and less charge recombination than mesoporous TiO₂ NP films. [38, 149, 164] There is much room for improvement in the energy conversion efficiency of

current DSSCs based on TiO₂ nanotube arrays compared to the relatively extensively researched mesoporous TiO₂ NP film-based DSSCs. [188]

To date, several approaches for improving the energy conversion efficiency of TiO₂ nanotube array-based DSSCs have been reported. Metal NPs, which can harvest light via surface plasmon resonance (SPR), have been used to enhance the energy conversion efficiency of DSSCs by introducing Au or Ag NPs into TiO₂ nanotube arrays. [143, 189-191] Barrier layers remove TiO₂ nanotube arrays, so open-ended TiO₂ nanotube arrays, which can also be classified as arrays of columnar nanopores, have been used for DSSCs to provide increased energy conversion efficiency.[163] Moreover, the energy conversion efficiency of TiO₂ nanotube array-based DSSCs can be further increased by introducing a scattering layer to the active layer.[192]

So far, TiO₂ nanotubes that make use of a scattering layer[192] or plasmonic materials [179] have been reported, but a scattering layer with plasmonic materials has not been used in TiO₂ nanotube-based DSSCs. In this study, we report the development of freestanding TiO₂ nanotube arrays filled with Ag NPs and large TiO₂ NPs, which improve the energy conversion efficiency of DSSCs. Furthermore, we compare the effects of Ag NPs and large TiO₂ NPs in open- and close-ended freestanding TiO₂ nanotube arrays in DSSCs. The energy

conversion efficiencies of the following eight types of DSSC were compared: closed-ended freestanding TiO₂ nanotube arrays with/without Ag NPs and/or TiO₂ scattering layer and open-ended freestanding TiO₂ nanotube arrays with/without Ag NPs and/or a TiO₂ scattering layer.

4. 2. Results and discussion

4. 2. 1 Overall scheme of DSSCs with freestanding TiO_2 nanotube arrays with channels containing Ag NPs

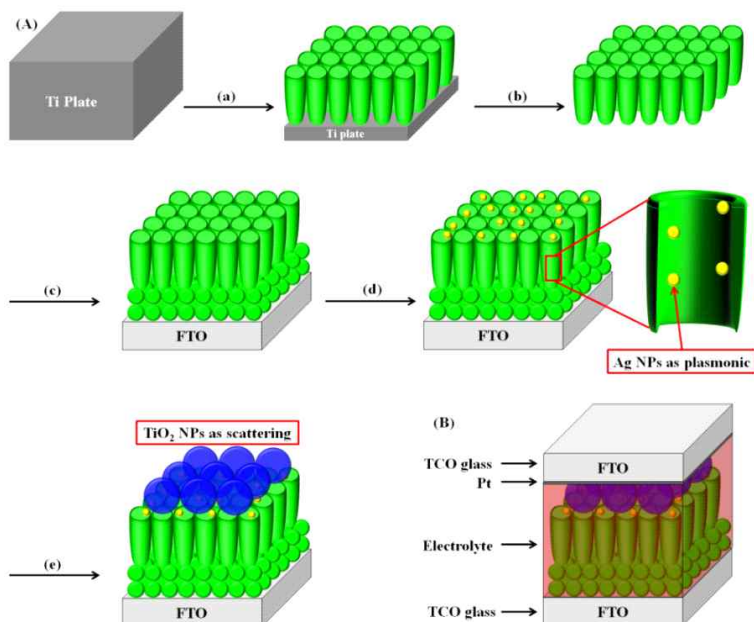


Figure 4. 1. Overall scheme of DSSCs with open-ended freestanding TiO_2 nanotube arrays with Ag NPs and large TiO_2 NPs. (A) (a) Ti anodization for TiO_2 nanotube arrays, (b) freestanding TiO_2 nanotube arrays and etching by ion milling, (c) transference of open-ended freestanding TiO_2 nanotube arrays on FTO glass, (d) formation of Ag NPs by UV irradiation, and (e) introducing large TiO_2 NPs. (B) Structure of DSSC with freestanding TiO_2 nanotube arrays and large TiO_2 NPs.

Figure 4. 1. illustrates the fabrication of DSSCs with Ag NPs and large TiO_2 NPs to enable plasmonic and scattering effect in open-ended freestanding TiO_2 nanotube array-based DSSCs. Ti plates were anodized and then annealed at 500C for 1 h to prepare anatase TiO_2 nanotube arrays. After carrying out secondary anodization, the TiO_2 nanotube arrays were easily detached from the Ti plate. TiO_2 nanotube arrays, once separated from the Ti plate, are termed “closed-ended freestanding TiO_2 nanotube arrays”. Freestanding TiO_2 nanotube arrays have a barrier layer at the bottom that disturbs electron transport and electrolyte diffusion.

This barrier layer was removed using the ion-milling method with several minutes of Ar^+ bombardment to yield “open-ended freestanding TiO_2 nanotube arrays”. The closed- and open-ended freestanding TiO_2 nanotube arrays were transferred to fluorine-doped tin oxide(FTO) glass using TiO_2 paste and annealed to enhance the adhesion between the closed- and open-ended freestanding TiO_2 nanotube arrays and the fluorine-doped tin oxide(FTO) glass.

To improve the energy conversion efficiency by the plasmonic effect, Ag NPs were embedded in the channel of freestanding TiO_2 nanotube arrays using 254 nm ultraviolet(UV) irradiation with aqueous silver nitrate. To further enhance the energy conversion efficiency, large TiO_2 NPs (400 nm) as a scattering layer were coated onto the active layer by the doctor blade method. This substrate was sandwiched with the

counter electrode and filled with electrolyte. The active area of the DSSCs was $\sim 0.25 \text{ cm}^2$.

4. 2. 2 Characterization of freestanding TiO_2 nanotube arrays with channels containing Ag NPs

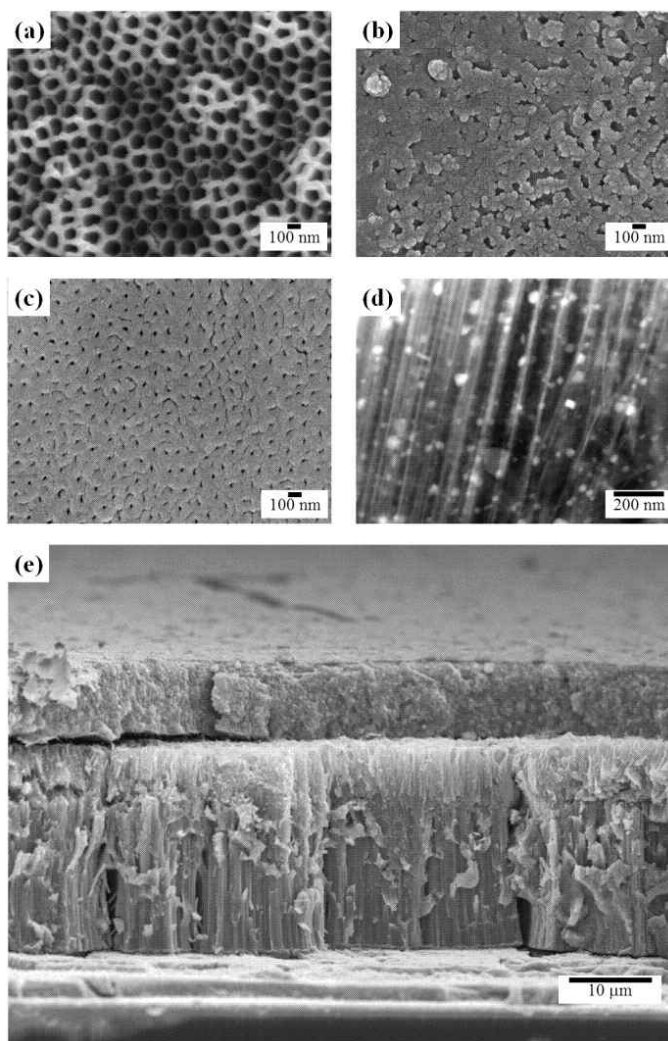


Figure 4. 2. FE-SEM images of the (a) top, (b) bottom, (c) bottom of post-ion milling freestanding TiO_2 nanotube arrays, (d) a high-angle annular dark-field (HAADF) image of Ag NPs in the channel of TiO_2 nanotube arrays, and (e) a side view of the active layer with freestanding TiO_2 nanotube arrays and the scattering layer.

FE-SEM images of freestanding TiO₂ nanotube arrays are shown in Figure 4. 2. The top side of the freestanding TiO₂ nanotube arrays had 100-nm-diameter pores, as shown in Figure 4. 2(a). The bottom layer of closed-ended freestanding TiO₂ nanotube arrays before ion milling lacked pores, as shown in Figure 4. 2(b).

However, after ion milling, 20-nm-diameter pores were evident on the bottom layer of open-ended freestanding TiO₂ nanotube arrays, as shown in Figure 4. 2(c). Open ended TNT arrays can be prepared by the chemical etching [161] and physical etching method. [162, 163] In the chemical etching method, the bottom layer of TNT arrays were easily removed by etchant. However, the surface morphology and length of TNT arrays were also dissolved in etchant and TNT arrays are easily fragile when they are attached on substrate due to their amorphous crystallinity. In the physical etching method, the bottom layer of TNT arrays were removed by the plasma or ion milling process, which is not simple. However, the surface morphology and length of TNT arrays were not damaged and they are very stable when they are attached on substrate because TNT arrays have the annealing process to change the crystallinity from the amorphous to anatase phase. After UV irradiation using a silver source, ~30 nm Ag NP were found in the channels of freestanding TiO₂ nanotubes in high-angle-annular dark-field (HAADF) images, as shown in Figure 4. 2(d). The length of TiO₂ nanotubes was ~22 μm and the length of the

scattering layer, which consisted of 400 nm TiO₂ NPs, was ~10 nm.

The UV-Vis spectrum of Ag NPs in the channels of freestanding TiO₂ nanotubes was shown in Figure 4. 3. A broad absorption peak centered at 402 nm was observed, which indicates the presence of Ag NPs.[193, 194] The absorption band of Ag NPs is matched with dye, N719 has two visible absorption bands; 390 nm 531 nm, [195] that was affected by the plasmon band. Moreover, the shell of Ag NPs was prepared by TiCl₄ to prevent the trapping electron by Ag NPs that will be better to electron transport in the channel of TiO₂ nanotube arrays.

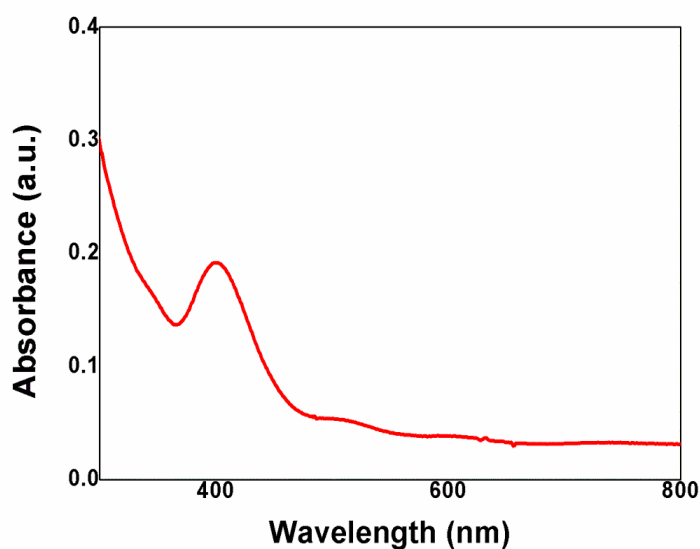


Figure 4. 3. UV-vis spectrum of Ag NP-functionalized TiO₂ nanotubes

4. 2. 3. DSSCs with closed-ended freestanding TiO₂ nanotube arrays with channels containing Ag NPs and large TiO₂ NPs

The photocurrent-voltage curves of DSSCs fabricated using closed-ended freestanding TiO₂ nanotube arrays measured under air mass 1.5 illumination (100 mW/cm²) are shown in Figure 4. 4. and Table 4. 1.

Four types of closed-ended freestanding TiO₂ nanotube array-based DSSCs were fabricated to assess the effect of each component on the energy conversion efficiency: closed-ended freestanding TiO₂ nanotube array based DSSCs without Ag or large TiO₂ NPs (a), with Ag NPs (b), with large TiO₂ NPs (c), and with Ag NPs and large TiO₂ NPs (d). The open-circuit voltage (V_{oc}), short-circuit current (J_{sc}), fill factor (ff), and energy conversion efficiency values are shown in Table 4. 1. The energy conversion efficiency of DSSCs based on closed-ended freestanding TiO₂ nanotube arrays lacking NPs was 5.86%. The energy conversion efficiency of the DSSCs based on closed-ended freestanding TiO₂ nanotube arrays with Ag NPs, with large TiO₂ NPs, and with Ag NPs and large TiO₂ NPs was 6.40%, 6.24%, and 6.71%, respectively.

Introduction of Ag NPs increased the energy conversion efficiency significantly, by 9.21% comparing to closed-ended freestanding TiO₂ nanotube array-based DSSCs without Ag and large TiO₂ NPs (actual efficiency; 6.40%), due to increased light harvesting by the plasmonic

effect.

Introduction of large TiO_2 NPs increased the energy conversion efficiency significantly, by 6.48% (actual efficiency; 6.40%), due to increased light harvesting by the scattering effect. Moreover, introduction of both Ag NPs and large TiO_2 NPs increased the energy conversion efficiency significantly, by 14.50% (actual efficiency; 6.71%), due to increased light harvesting by both the plasmonic and scattering effects.

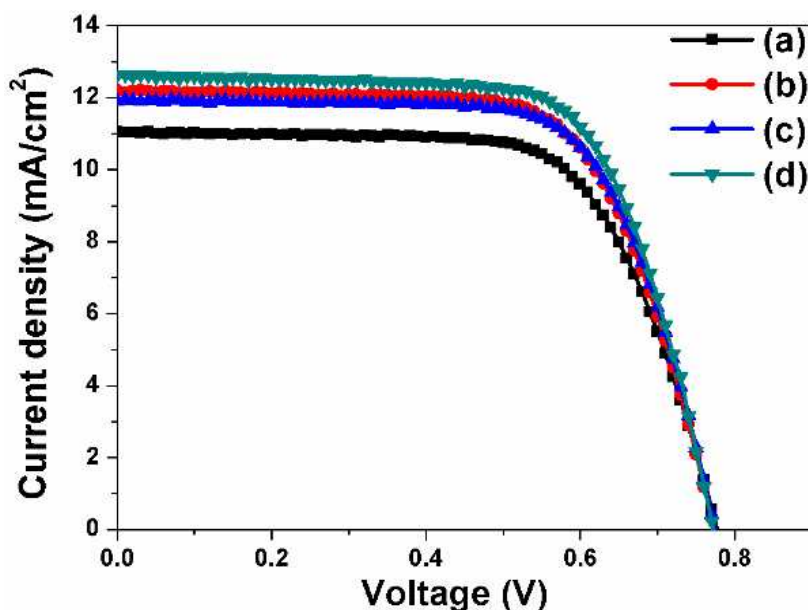


Figure 4. 4. I-V curves of DSSC-based closed-ended freestanding TiO_2 nanotube arrays fabricated without NPs (a), with Ag NPs (b), with large TiO_2 NPs (c), and with Ag NPs and large TiO_2 NPs (d).

Table 4. 1. Photovoltaic properties of DSSCs based on closed-ended freestanding TiO₂ nanotube arrays

DSSCs		Jsc (mA/cm ²)	Voc (V)	ff	η (%)
a	Closed-ended freestanding TiO ₂ nanotube arrays without any NPs	11.05	0.78	0.68	5.86
b	Closed-ended freestanding TiO ₂ nanotube arrays with Ag NPs	12.22	0.77	0.68	6.40
c	Closed-ended freestanding TiO ₂ nanotube arrays with large TiO ₂ NPs	11.90	0.76	0.69	6.24
d	Closed-ended freestanding TiO ₂ nanotube arrays with Ag NPs and large TiO ₂ NPs	12.63	0.77	0.69	6.71

4. 2. 4. DSSCs with open-ended freestanding TiO₂ nanotube arrays with channels containing Ag NPs and large TiO₂ NPs

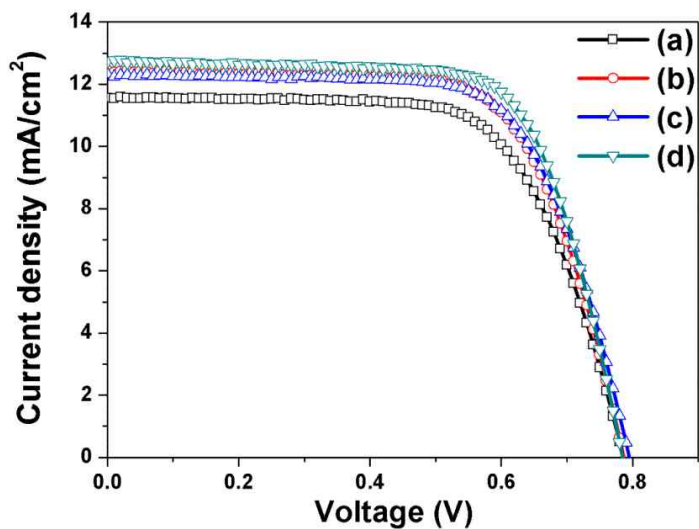


Figure 4. 5. I-V curves of DSSCs based open-ended freestanding TiO₂ nanotube arrays fabricated without NPs (a), with Ag NPs (b), with large TiO₂ NPs (c), and with Ag NPs and large TiO₂ NPs (d).

Table 4. 2. Photovoltaic properties of DSSCs based on open-ended freestanding TiO₂ nanotube arrays

DSSCs		J _{sc} (mA/cm ²)	V _{oc} (V)	ff	η (%)
a	Open-ended freestanding TiO ₂ nanotube arrays without any NPs	11.56	0.79	0.67	6.12
b	Open-ended freestanding TiO ₂ nanotube arrays with Ag NPs	12.45	0.79	0.68	6.68
c	Open-ended freestanding TiO ₂ nanotube arrays with large TiO ₂ NPs	12.33	0.79	0.68	6.62
d	Open-ended freestanding TiO ₂ nanotube arrays with Ag NPs and large TiO ₂ NPs	12.74	0.78	0.71	7.05

The photocurrent-voltage curves of DSSCs fabricated using open-ended freestanding TiO₂ nanotube arrays are shown in Figure 4. 5. and Table 4. 2. to assess the effect of each component on the energy conversion efficiency.

Four types of DSSC-based open-ended freestanding TiO₂ nanotube array were fabricated: open-ended freestanding TiO₂ nanotube array based DSSCs without Ag or large TiO₂ NPs, (a), with Ag NPs (b), with large TiO₂ NPs (c), and with Ag NPs and large TiO₂ NPs (d). The Voc, Jsc, ff, and η values are summarized in Table 4. 2.

The energy conversion efficiency of DSSCs based on open-ended freestanding TiO₂ nanotube arrays lacking NPs was 6.12%. The energy conversion efficiency of DSSCs based on open-ended freestanding TiO₂ nanotube arrays with Ag NPs, with large TiO₂ NPs, and with Ag NPs and large TiO₂ NPs was 6.68%, 6.62%, and 7.05%, respectively.

The introduction of Ag NPs and/or large TiO₂ NPs increased the energy conversion efficiency by 9.15%, 8.17%, and 15.20%, respectively. Compared to closed-ended freestanding TiO₂ nanotube arrays, the energy conversion efficiency of DSSCs based on open-ended freestanding TiO₂ nanotube arrays was 5.07% (6.71–7.05%) higher due to enhanced electron transport and electrolyte diffusion. [163]

Although TiO₂ nanotube arrays based DSSCs have great potential, as best we know, the theoretical maximum improvement by Ag NPs or TiO₂ scattering layer of TiO₂ nanotube based DSSCs were not reported yet. However, the opened-end TiO₂ nanotube-based device exhibited an increase in one-sun efficiency from 5.3% to 9.1%, corresponding to 70% which is a much higher increase compared to our results. [161] We believe that there is lots of room to improve the efficiency by

combining each parameter with an optimal condition based on theoretical studies.

4. 3. Conclusions

In this study, we compared the natural consequence of several parameters such as the plasmonic effect, the scattering effect and open-closed ended freestanding TiO₂ nanotube as basic experiments for better efficiency.

We demonstrated that the plasmonic and scattering effects enhanced the energy conversion efficiency of freestanding TiO₂ nanotube arrays in DSSCs. Ag NPs were added to the channels of TiO₂ nanotube arrays by UV irradiation to induce a plasmonic effect, and large TiO₂ NPs were introduced to TiO₂ nanotube arrays to induce a scattering effect.

The energy conversion efficiency of DSSCs with both Ag NPs and large TiO₂ NPs was higher than that of DSSCs without Ag NPs due to the plasmonic effect, and was higher than that of DSSCs without large TiO₂ NPs due to the scattering effect. Compared to closed-ended freestanding TiO₂ nanotube arrays, open-ended freestanding TiO₂

nanotube arrays exhibited enhanced energy conversion efficiency.

We demonstrate that Ag NPs, TiO₂ NPs and open-ended freestanding TiO₂ nanotube arrays, enhanced the energy conversion efficiency; furthermore, the combination of all components exhibited the highest energy conversion efficiency. Our research suggests that the energy conversion efficiency of DSSCs is improved by both the plasmonic and scattering effects and have applications in organic solar cells, hybrid solar cells, and perovskite solar cells.

Chapter 5

Dual Functionalized Freestanding TiO₂ Nanotube Arrays Coated with Ag Nanoparticles and Carbon Materials for Dye-Sensitized Solar Cells

5.1. Overall view of dual Functionalized Freestanding TiO₂ Nanotube Arrays Coated with Ag Nanoparticles and Carbon Materials for Dye-Sensitized Solar Cells

Highly ordered, freestanding TiO₂ nanotube arrays (TiO₂ NTAs) were prepared using an electrochemical method. The barrier layer was etched to open the bottom of each array, aptly named “open-ended TiO₂ NTAs”. These arrays were coated with silver nanoparticles (Ag NPs) and/or carbon materials to enhance electron generation and transport. The energy conversion efficiency of the resulting dye-sensitized solar cells (DSSCs) with open-ended freestanding TiO₂ NTAs increased from 5.32% to 6.14% (by 15%) when coated with Ag NPs due to plasmonic interactions. Meanwhile, coating the open-ended freestanding TiO₂ NTAs with carbon materials increased the energy conversion efficiency from 5.32% to 6.07% (by 14%) due to π - π conjugation. When the Ag NPs and carbon materials were simultaneously applied to the open-ended freestanding TiO₂ NTAs, the energy conversion efficiency increased to 6.91% an enhancement of 30% due to the additive effects of plasmonics and π - π conjugation.

Global research of dye-sensitized solar cells (DSSCs) has continued since their initial development in 1991 by the Grätzel group [58], due to their low cost, ease of fabrication, and high power conversion efficiency [64,197]. Titanium dioxide nanoparticles (TiO₂ NPs) are

typically used as the photoanode in DSSCs, because they have a desirable direct band gap (3.2 eV) and a large surface area for adsorbing dyes, both helping to generate electrons [198-201]. However, TiO₂ NPs are randomly networked, and the countless grain boundaries within them lead to material defects and charge recombinations that inhibit the smooth electron transport [62,202,203].

For the last few years, TiO₂ nanotube arrays (NTAs) have been explored as an alternative to TiO₂ NPs [204-206]. The TiO₂ NTAs could be fabricated using an electrochemical method (i.e., anodization) [204,207], and they have a well-ordered and vertically aligned tubular structure that serves as direct electron pathways to enhance not only electron transport, but also charge separation [208,209]. Despite their merits, however, the barrier layer at the bottom of the TiO₂ NTAs could impede charge transfer and electrolyte diffusion. To overcome this problem, we recently removed the bottom layer of TiO₂ NTAs using argon ion (Ar⁺) milling, resulting in improved electron transport and electrolyte diffusion [210].

There have been an increasing number of studies that add carbon to TiO₂ NTAs in order to improve the charge separation and transfer of electrons, due to the superior electrical properties of the π - π conjugation [211-213]. Many researchers have explored the application of carbon materials in solar cell technologies. Carbon 60 (C₆₀ or “fullerene”) and carbon nanotubes (CNTs) are well known for their

roles as electron acceptors and charge separators in organic solar cells [214,215]. When incorporated in organic solar cells, CNTs act as exciton dissociation sites and hopping centers for hole transport [216], and graphene mixed with TiO₂ NPs in DSSCs plays a role in promoting charge separation and movement [217]. As mentioned, TiO₂ NTAs were developed as alternatives to TiO₂ NPs. While it is not simple to blend carbon materials with TiO₂ NTAs, we recently reported a method for enriching freestanding TiO₂ NTAs with carbon for use in DSSCs. When a small amount of carbon was deposited on TiO₂ NTAs, the energy conversion efficiency increased by approximately 22.4% compared to those without carbon enrichment [218]. We tentatively attributed this to the π - π conjugation introduced by carbon enrichment that improved the efficiency of electron transport.

A plasmonic effect triggered by metal NPs such as silver and gold can be used to enhance photoabsorption in solar cells [219-221]. When incident photons pass by Ag NPs, they cause electron vibration and photo scattering of the nanoparticles, which facilitate more efficient photon control [222]. The metal NPs were incorporated by mixing with TiO₂ sol in the DSSCs or with precursors of the active layer in organic solar cells. In the fabrication of DSSCs based on TiO₂ NTAs, however, it is difficult to insert metal NPs into the channels of TiO₂ NTAs. We recently devised a simple method for the complete formation of Ag NPs in the channels of TiO₂ NTAs using UV

irradiation [223-225]. These NPs helped the dyes generate electrons, as demonstrated by a high current density in the DSSCs.

To date, we have confirmed the separate effects of carbon enrichment and the incorporation of Ag NPs in previous studies. However, the effects of adding both carbon and Ag NPs remain unknown. Here, we report the effects of enriching freestanding TiO₂ NTAs with one or both materials on the performance of DSSCs, in terms of enhanced electron transport and plasmonic effects. Carbon materials were synthesized by chemical vapor deposition (CVD) and deposited on the wall of TiO₂ NTAs. The Ag NPs were formed using UV irradiation within the channels of TiO₂ NTAs.

5.2. Materials and Methods

5.2.1 Preparation of closed- and open-ended TiO₂ NTAs

Titanium foils (Alfa Aesar, 99.7% purity, 2.5 cm × 4.0 cm × 320 μm) were prepared and anodized using an electrochemical method to fabricate TiO₂ NTAs. The electrolyte was composed of 0.8 wt.% NH₄F and 2 vol.% H₂O in ethylene glycol. Carbon rods served as the cathode material. A 60 V DC potential was supplied to the titanium foils at 25 °C for 2 h. Later, the anodized titanium foils were annealed in a tube furnace at 450 °C for 1 h, and a second anodization process was then conducted on the samples at 30 V for 10 min. To detach the TiO₂ NTAs from the titanium foils, the sample after the second anodization was immersed in 10% H₂O₂ for 24 h. Ion milling with argon (Ar⁺) bombardment was used to remove the bottom of the TiO₂ NTAs to make open-ended tips [218].

5.2.2 Preparation of photoanodes for DSSCs based on the TiO₂ NTAs

Fluorine-doped tin oxide (FTO) glass was washed and sonicated in ethanol and acetone to remove impurities. Titanium diisopropoxide bis(acetylacetonate) (5 wt.% in n-butanol) was spin-coated on the clean FTO glass to form a compact TiO₂ blocking layer after annealing at

450 °C for 1 h. A TiO₂ paste (Solaronix, T/SP) was applied to the FTO glass using a doctor blade method, in order to attach the closed- or open-ended TiO₂ NTAs. Finally, the samples were annealed in a furnace at 450 °C for 30 min.

5.2.3 Synthesis of Ag NPs on the TiO₂ NTAs by UV irradiation

The samples were placed in a 0.3 mM AgNO₃ aqueous solution. Ag NPs were synthesized in the channels of closed- or open-ended TiO₂ NTAs using a 254 nm UV lamp for 3 min.

5.2.4 Synthesis of carbon materials by CVD on TiO₂ NTAs

The samples were placed in a quartz tube furnace filled with nitrogen (200 standard cubic centimeter per minute [sccm]). Hydrogen gas (30 sccm) and ethylene gas (40 sccm) were flowed into the tube furnace at 450 °C for 30 s.

5.2.5 Fabrication of dye-sensitized solar cells

All DSSC samples were post-treated with 10 mM TiCl₄ solution at 50 °C for 30 min, and then annealed at 450 °C for 1 h. These steps not only enhanced the photocurrent but also prevented the Ag NPs from dissolution upon contacting the iodine-iodide electrolyte. Each

treated sample was stained using dye molecules (0.5 mM solutions of N719 ((Bu₄N)₂Ru(dobpyH)₂(NCS)₂, Solaronix) in ethanol at 50 °C for 8 h. Following this treatment, samples were washed with ethanol to eliminate physisorbed dye molecules. To fabricate counter electrodes, chloroplatinic acid (H₂PtCl₆) in ethanol was drop-casted onto clean FTO glass and annealed in a tube furnace at 400 °C for 1 h.

The electrolyte used to separate the electrodes contained 0.7 M 1-butyl-3-methyl-imidazolium iodide (BMII), 0.03 M I₂, 0.1 M guanidium thiocyanate (GSCN), and 0.5 M 4-tert-butyl pyridine (TBP) in a mixture of acetonitrile and valeronitrile (85:15 v/v). A 60- μ m thick hot-melt Surlyn spacer (Solaronix) was put between the photoanode and counter electrode. The electrolyte was injected into the space provided by the Surlyn spacer.

5.2.6 Characterization of dye-sensitized solar cells

The structures of TiO₂ NTAs on FTO glass were confirmed using a field emission scanning electron microscope (FE-SEM, JSM-6330F, JEOL Inc.). The existence of Ag NPs in the channels of TiO₂ NTAs was verified by the high-angle annular dark-field (HAADF) imaging technique using a scanning transmission electron microscope (TEM, JEM-2200FS, JEOL Inc.). Raman spectra were measured with a Raman spectrometer (LabRAM HV Evolution spectrometer, HORIBA). The

UV-Vis spectra were recorded using a UV-Vis spectrophotometer (NEOSYS-2000, SCINCO). Current density-voltage measurements were carried out using an electrometer (Keithley 2400) and a solar simulator (1 kW xenon with AM 1.5 filter, PEC-L01, Peccel Technologies). Electrochemical impedance spectroscopy (EIS) data were collected using a potentiostat (Solartron 1287) equipped with a frequency response analyzer (Solartron 1260) between 10^{-2} and 10^6 Hz under AM 1.5 light illumination, and analyzed using Z-View software (Solartron Analytical). The applied bias voltage and AC amplitude were set at the open circuit voltage (V_{oc}) of the DSSCs and 10 mV, respectively.

5.3. Results and Discussion

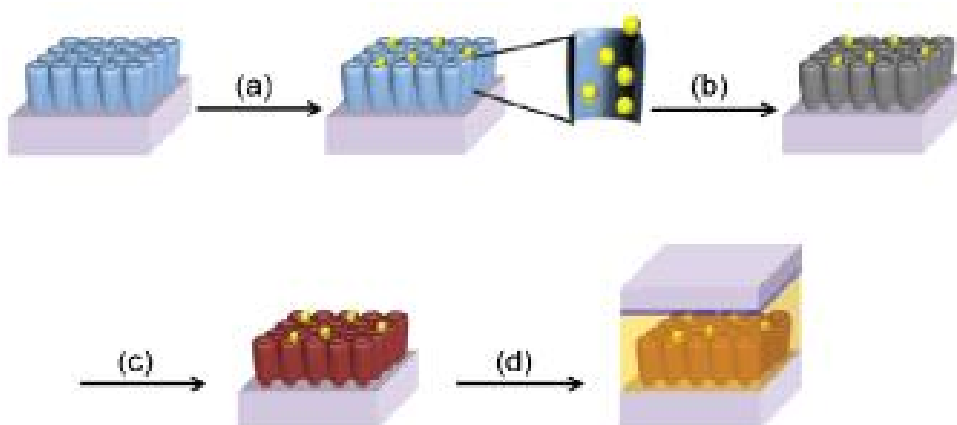


Figure 5.1. Overall scheme of the fabrication of DSSCs based on freestanding TiO₂ NTAs coated with Ag NPs and carbon materials. (a) Synthesis of Ag NPs in the channel of TiO₂ NTAs, (b) deposition of carbon materials, (c) dye adsorption, and (d) fabrication of the DSSC.

The fabrication of DSSCs based on freestanding TiO₂ NTAs is shown in Figure 5.1. The bottom layer was present in the closed-ended freestanding TiO₂ NTAs, but removed by ion milling in the open-ended ones. The DSSCs were fabricated from both types of freestanding TiO₂ NTAs to compare their energy conversion efficiency. In both cases, the freestanding TiO₂ NTAs were attached to the FTO glass with TiO₂ paste, and Ag NPs were synthesized using UV irradiation as shown in Figure 5.1(a). Carbon materials were synthesized using CVD as shown in Figure 5.1(b). By using the UV irradiation and CVD, Ag NPs and

carbon materials were deposited in the channel of highly ordered TiO₂ NTAs without any distortion. The dye (N719) was adsorbed onto both types of freestanding TiO₂ NTAs, as shown in Figure 5.1(c). Finally, DSSCs were fabricated by assembling the working electrode (freestanding TiO₂ NTAs on FTO glass) and the counter electrode (Pt on FTO glass) as shown in Figure 5.1(d).

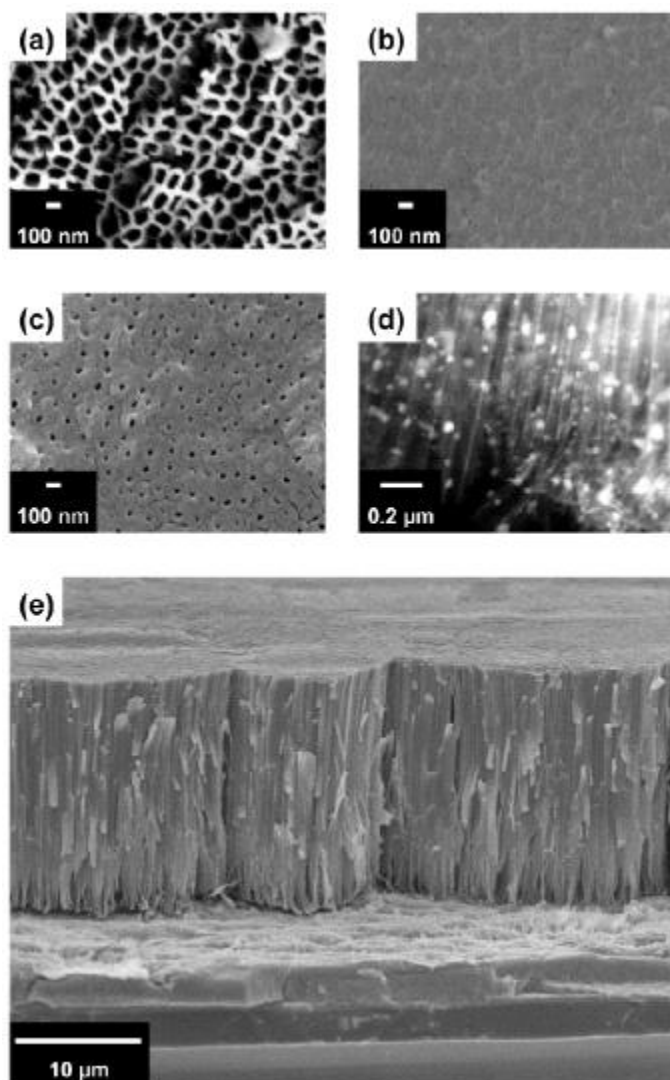


Figure 5.2. FE-SEM images of TiO₂ NTAs: (a) top view, (b) bottom view, and (c) bottom view after ion milling. (d) HAADF image of Ag NPs in the channel of TiO₂ NTAs. (e) Side view of TiO₂ NTAs on FTO glass.

FE-SEM images of TiO₂ NTAs are shown in Figure 5.2. The top view in Figure 5.2(a) shows a pore size of approximately 100 nm after applying the electrochemical method. The bottom of the TiO₂ NTAs before ion milling (the “closed-ended TiO₂ NTA”) is shown in Figure 5.2(b), with a total bottom pore size of approximately 100 nm including that of the wall thickness. However, when the bottom was removed by ion milling to produce the “open-ended TiO₂ NTAs” (Figure 5.2(c)), the bottom pore was reduced to 30 nm in size, and the wall thickness was approximately 35 nm. An HAADF image of Ag NPs in the channels of TiO₂ NTAs is shown in Figure 5.2(d), and the diameter of Ag NPs was approximately 30 nm. This allowed the Ag NPs to be successfully immobilized inside the channel of TiO₂ NTAs by UV irradiation, and the resulting plasmonic interactions may have affected all the surface areas. A side view of TiO₂ NTAs attached to the FTO glass by TiO₂ paste after being sintered at 450 °C is shown in Figure 5.2(e). The main role of the TiO₂ paste is to connect the TiO₂ NTAs with the FTO glass surface. The thickness of the TiO₂ film layer was 3 μm, and the length of TiO₂ NTAs was approximately 18 μm.

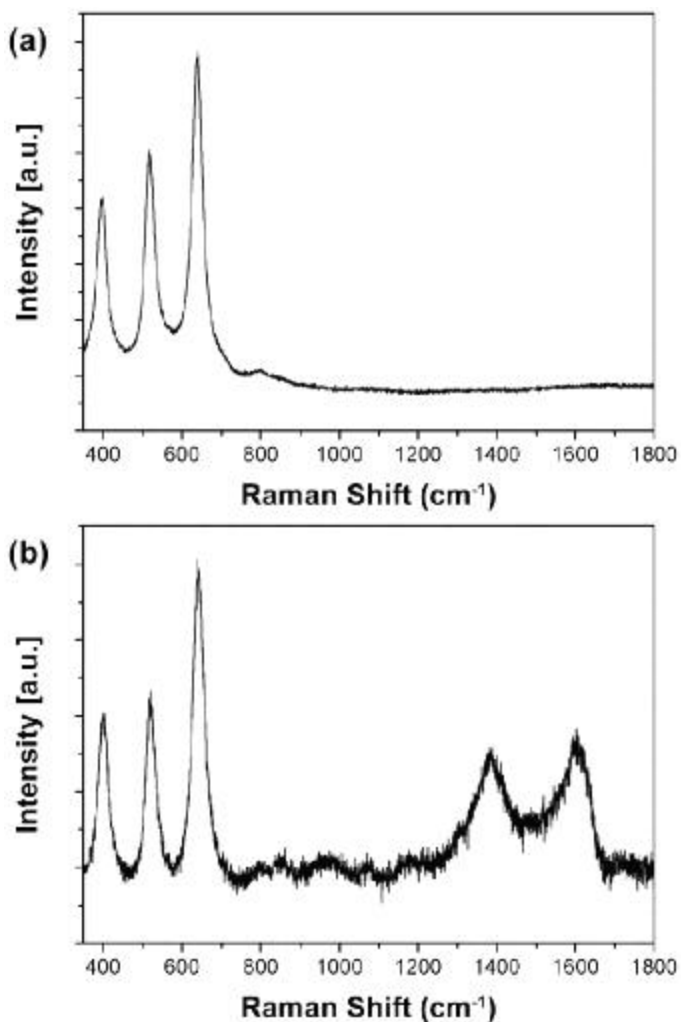


Figure 5.3. Raman spectra of TiO₂ NTAs (a) without and (b) with carbon materials.

Carbon materials on the TiO₂ NTAs were synthesized by CVD, and their structure was confirmed by Raman spectroscopy, as shown in Figure 5.3. In a previous publication, we reported the optimization of

TiO₂ NTAs for DSSCs using carbon materials [218]. The B1g (397 cm⁻¹), A1g (518 cm⁻¹), and Eg (641 cm⁻¹) peaks indicated that the TiO₂ NTAs were in the form of anatase TiO₂, as shown in Figure 5.3(a) [225]. When carbon materials were synthesized on the TiO₂ NTAs using CVD, the G band at 1600 cm⁻¹ represented graphite, while the D band at 1384 cm⁻¹ was due to the disorderly network of sp² and sp³ sites in the carbon materials, as shown in Figure 5.3(b). The sp² sites of carbon materials resulted in a π - π conjugation that improved the efficiency of electron transport across the TiO₂ NTAs.

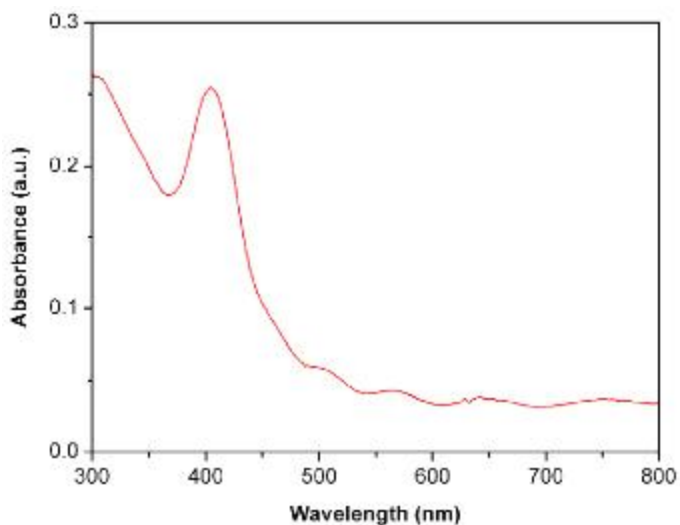


Figure 5.4. UV-Vis spectrum of Ag NPs on the TiO₂ NTAs.

Ag NPs were synthesized on the TiO₂ NTAs using UV irradiation, and this was confirmed by the ultraviolet-visible (UV-Vis) spectrum. The size of Ag NPs was confirmed as approximately 30 nm, using the HAADF image shown in Figure 5.2(d). An absorption peak centered at 405 nm was observed, as shown in Figure 5.4. Our previous paper [218] reported on the optimization of TiO₂ NTAs using Ag NPs. Other researchers have reported that Ag NPs with sizes of approximately 30 nm had UV-Vis absorption peaks at 420 nm. However, in this case, the Ag NPs were synthesized using UV irradiation (at 254 nm) without the addition of any stabilizing or reducing agents. As such, the Ag NPs were immobilized in the TiO₂ NTAs, which would affect absorption in the UV-Vis spectrum. The absorption band of Ag NPs is within the same range as that of the dye N719 (cis-diisothiocyanato-bis(2,2'-bipyridyl-4,4'-dicarboxylato) ruthenium(II) bis(tetrabutylammonium)), 390–530 nm, leading to enhanced electron generation from the dye by means of plasmonic interactions.

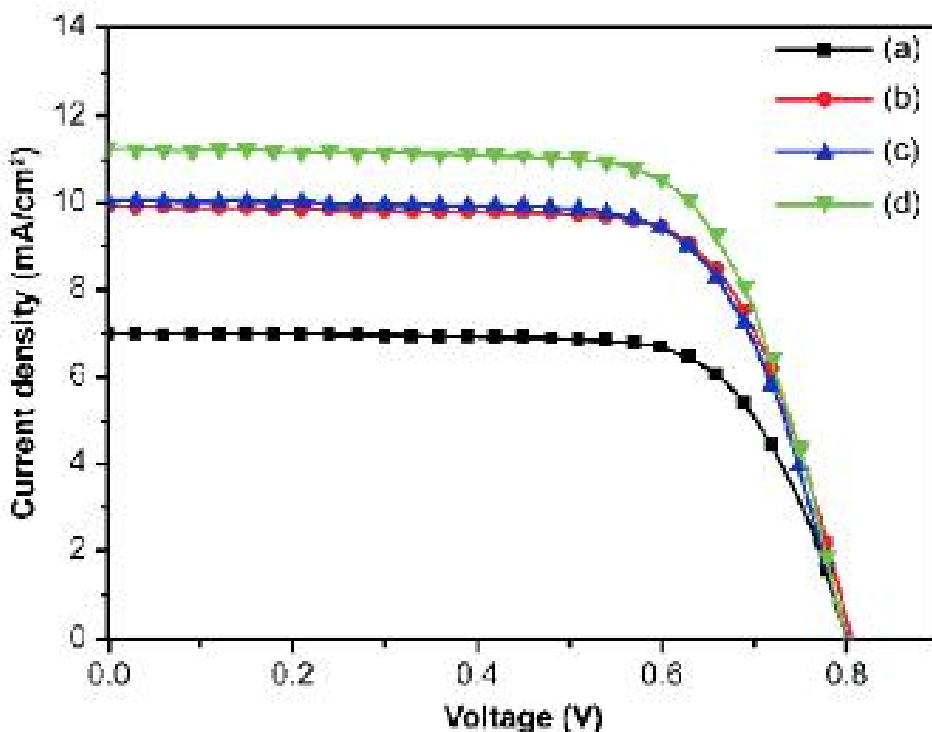


Figure 5.5. Current density-voltage curves of DSSCs based on (a) unmodified closed-ended TiO_2 NTAs, (b) with embedded Ag NPs, (c) with applied carbon materials, and (d) with both Ag NPs and carbon materials.

The current density-voltage curves of DSSCs using closed-ended TiO_2 NTAs with or without modification were measured under air-mass (AM) 1.5 sunlight, and the results are presented in Figure 5.5. The V_{oc} , short-circuit current density (J_{sc}), fill factor (ff), and energy conversion efficiency (η) of the DSSCs are summarized in Table 5.1. For the DSSCs without any treatment, the energy conversion efficiency

was 4.10%, which increased to 5.73% when Ag NPs were embedded via UV irradiation, corresponding to an overall increase of 40%. When carbon materials were added to the closed-ended TiO₂ NTAs via CVD, the energy conversion efficiency improved to 5.69%, corresponding to a 39% increase. With both Ag NPs and carbon materials, the energy conversion efficiency improved even further to 6.36%, corresponding to an overall increase of 55%. Note that when Ag NPs were treated with TiCl₄, the core-shell type Ag@TiO₂ nanoparticles were formed. Because the dye are adsorbed on Ag@TiO₂, the amount of dye loading might not be significantly reduced as shown table 5.1. As previously reported [218], a large amount of carbon doping materials could lower the conversion efficiency by decreasing of dye loading. However, in this case only a trace amount of carbon material was deposited, which did not decrease the dye loading significantly.

Table 5.1. Photovoltaic properties of DSSCs based on closed-ended TiO₂ NTAs with Ag NPs and/or carbon materials.

DSSCs based on closed-ended TiO ₂ NTAs decorated	J_{sc} (mA/cm ²)	Voc (V)	ff	η (%)	Dye loading (nmol/cm ²)
without Ag NPs and carbon materials	7.02	0.81	0.72	4.10 ±0.28	144
with Ag NPs	9.92	0.81	0.72	5.73 ±0.31	142
with carbon materials	10.03	0.80	0.71	5.69 ±0.26	139
with Ag NPs and carbon materials	11.25	0.80	0.71	6.36 ±0.34	141

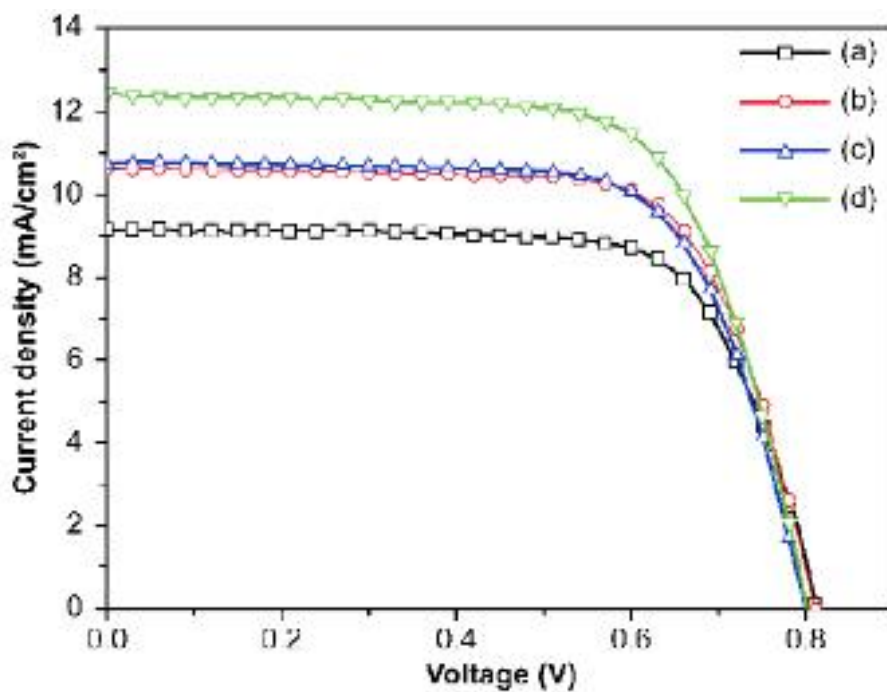


Figure 5.6. Current density-voltage curves of DSSCs based on (a) unmodified open-ended TiO_2 NTAs, (b) embedded with Ag NPs, (c) applied with carbon materials, and (d) with both Ag NPs and carbon materials.

Table 5.2. Photovoltaic properties of DSSCs based on open-ended TiO₂ NTAs with Ag NPs and/or carbon materials.

DSSCs based on open-ended TiO ₂ NTAs decorated	J_{sc} (mA/cm ²)	Voc (V)	ff	η (%)	Dye loading (nmol/cm ²)
without Ag NPs and carbon materials	9.12	0.81	0.72	5.32 ±0.36	153
with Ag NPs	10.61	0.81	0.71	6.14 ±0.46	151
with carbon materials	10.73	0.80	0.71	6.07 ±0.30	147
with Ag NPs and carbon materials	12.41	0.80	0.69	6.91 ±0.41	149

The current density-voltage curves for DSSCs based on open-ended TiO₂ NTAs with or without modification were also measured under AM 1.5 sunlight, and the results are presented in Figure 5.6. The Voc, Jsc, ff, η values of these DSSCs are summarized in Table 5.2. When using unmodified TiO₂ NTAs, DSSCs based on open-ended TiO₂ NTAs had higher energy conversion efficiency (5.32%) compare to those based on closed-ended ones (4.10%). The closed-end barrier of the TiO₂ NTA disturbs electron transport between the TiO₂ layer and the electrode [210,218].

When Ag NPs were embedded in the open-ended TiO₂ NTAs, the energy conversion efficiency improved from 5.32% to 6.14%, corresponding to a 15% enhancement. In this case, electron generation in the DSSCs was enhanced by the plasmonics from the NPs, despite the slightly diminished dye loading (from 153 to 151 nmol/cm²). When carbon materials alone were applied to TiO₂ NTAs, the energy conversion efficiency improved to 6.07%, a 14% increase. In this case, the carbon materials improved electron transport due to the π - π conjugation across the small quantity of carbon materials, in spite of a diminished dye load (153 to 147 nmol/cm², which is even less than that with Ag NPs). The reason is that the carbon materials were distributed to interact between the TiO₂ and the dye, making up for the loss of dye loading in term of the energy conversion efficiency. When Ag NPs and carbon materials were both applied to the open-ended TiO₂ NTAs, the energy conversion efficiency improved to 6.91%, corresponding to a 30% enhancement compared to the unmodified open-ended TiO₂ NTAs. In this case, the Ag NPs and carbon materials produced additive effects with their respective plasmonics and π - π conjugations, in spite of a slightly reduced dye loading of 149 nmol/cm². Comparing the performance parameters in Table 5.2, the Voc and ff decreased with treatment; the conduction band of the TiO₂ NTAs shifted down, which in turn affected the Voc and the charge recombination through electron density that suppressed

the ff. However, the J_{sc} was increased by the plasmonic activity in conjunction with π - π conjugation, which improved the energy conversion efficiency of DSSCs.

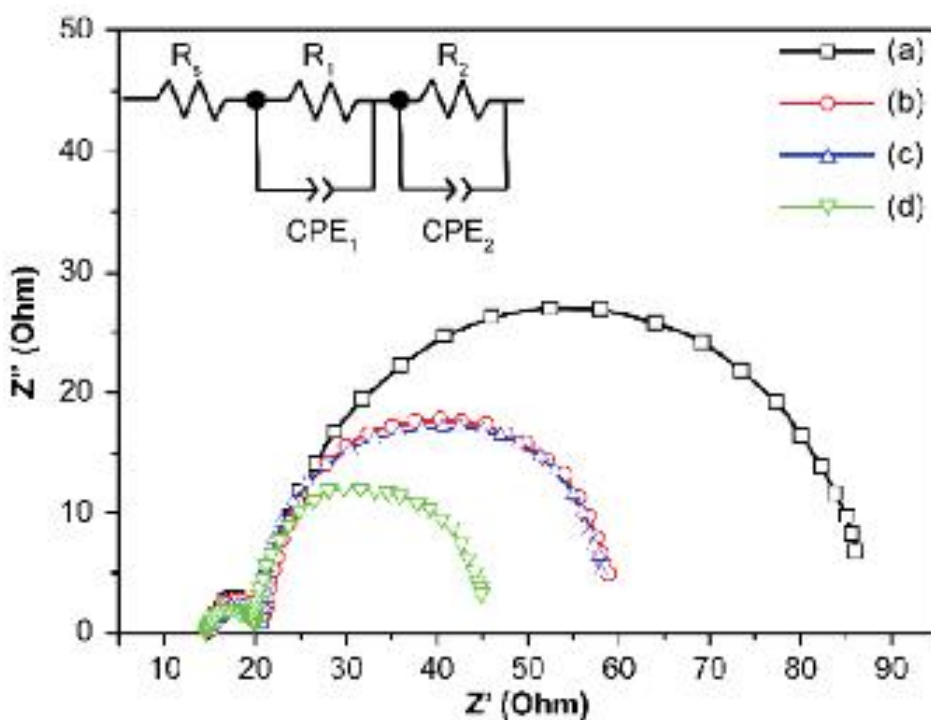


Figure 5.7. EIS data of DSSCs based on (a) unmodified open-ended TiO_2 NTAs, (b) embedded with Ag NPs, (c) applied with carbon materials, and (d) with both Ag NPs and carbon materials.

Table 5.3. EIS fitting results for DSSCs with open-ended TiO₂ NTAs.

DSSCs based on open-ended TiO ₂ NTAs decorated	R _s (Ω)	R ₁ (Ω)	CPE ₁ (F)	R ₂ (Ω)	CPE ₂ (F)
without Ag NPs and carbon materials	15.50	5.58	6.91×10 ⁻⁶	61.12	1.99×10 ⁻³
with Ag NPs	15.52	5.54	8.65×10 ⁻⁶	36.90	2.10×10 ⁻³
with carbon materials	15.56	5.07	1.62×10 ⁻⁵	36.40	2.03×10 ⁻³
with Ag NPs and carbon materials	14.99	4.88	1.16×10 ⁻⁶	24.55	2.99×10 ⁻³

The DSSCs based on open-ended TiO₂ NTAs were characterized by EIS across the frequency range from 10⁻² to 10⁶ Hz, as shown in Figure 5.7. The applied bias voltage was set at the Voc with 10 mV of AC amplitude. The data were analyzed using an equivalent circuit (inset of Figure 5.7), and the fit parameters are listed in Table 5.3. The ohmic series resistance (R_s) is due to the sheet resistance that corresponds to the x-axis value where the first semicircle begins on the left-hand side in Figure 5.7. The value of R_s is similar with or without Ag NPs or carbon materials, meaning that the additional deposits did not affect the sheet's resistance to FTO or the current collector. The R₁ value is given by the sum of the small semicircle at

high frequency that is assigned to the parallel combination of resistances, and the capacitances at the Pt-FTO/electrolyte and the FTO/TiO₂ interfaces. The R₂ value is given by the sum of the large semicircle at low frequency that is associated with the resistance, and the capacitance at the dye-adsorbed TiO₂/electrolyte interface, plus the transport resistance. The value of R₁ without and with Ag NPs was approximately 5.58 and 5.54 Ω, respectively. However, the value of R₂ with Ag NPs (36.90 Ω) is much lower than that without Ag NPs (61.12 Ω). More electrons were generated by plasmonic activities than produced at the dye-adsorbed TiO₂/electrolyte interface. As a result, the R₂ value was reduced in the presence of Ag NPs. The value of R₁ with carbon materials (5.07 Ω) was less than those without or with Ag NPs (5.58 and 5.54 Ω, respectively), whereas the value of R₂ (36.40 Ω) was less than that without Ag NPs (61.12 Ω). Electrons are better transported by π-π conjugation that is affected by the FTO/TiO₂ and dye-adsorbed TiO₂/electrolyte interfaces. Hence, the values of both R₁ and R₂ decreased in the presence of carbon materials. In the presence of both Ag NPs and carbon materials, the values of R₁ (4.88 Ω) and R₂ (24.55 Ω) were the lowest. More electrons were generated and better transported in this case, by a combination of plasmonics and the π-π conjugation that affected the FTO/TiO₂ and dye-adsorbed TiO₂/electrolyte interfaces. Therefore, the values of R₁ and R₂ were reduced.

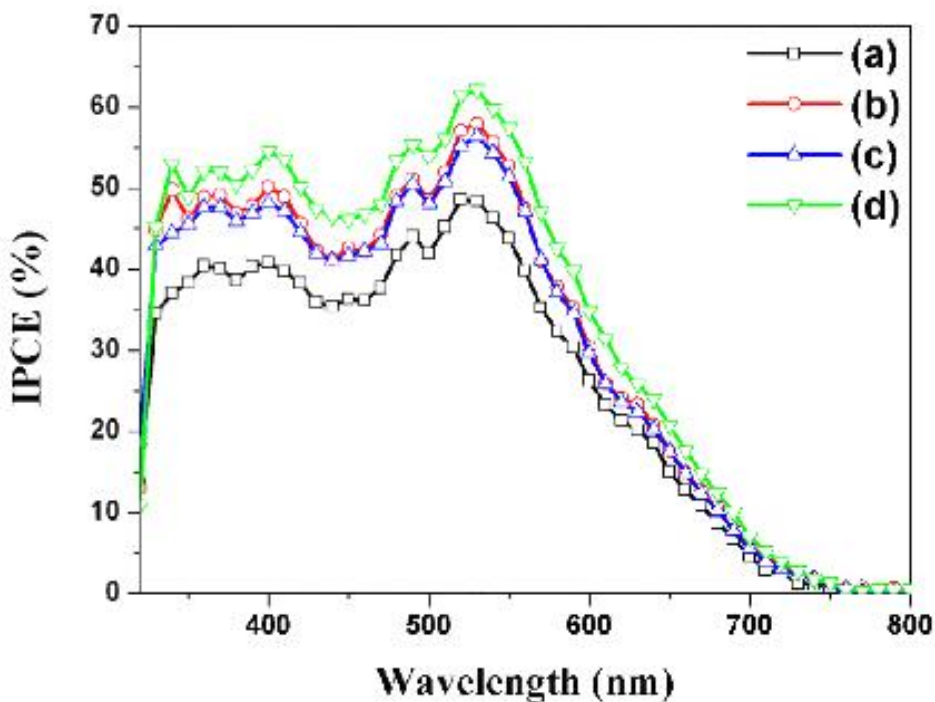


Figure 5.8. IPCE of DSSCs based on (a) unmodified open-ended TiO_2 NTAs, (b) embedded with Ag NPs, (c) applied with carbon materials, and (d) with both Ag NPs and carbon materials.

The incident photon-to-electron conversion efficiency (IPCE) of DSSCs based on the open-ended TiO_2 NTAs is shown in Figure 5.8. Plasmon is a kind of quasiparticle consisting of free electrons collectively vibrating within the metal. At the interface between a metal with negative dielectric constant and a medium having positive dielectric constant, surface plasmon resonance (SPR) combines a spreading electromagnetic wave on this interface (from visible to near-infrared

frequency) and the plasmon. This combination generates plasmon-polariton, which leads to optical absorption, and strong electric field is generated in some parts. During SPR, the light energy accumulates on the metal nanoparticle surface, and optical control is possible in the frequency range below the optical diffraction limit. Therefore, the intensity of DSSC based on open-ended TiO₂ NTAs embedded with Ag NPs is higher than that without. This may mean that more electrons were generated by the plasmonic activities, which increased the short circuit current. The current intensity in the DSSCs based on open-ended TiO₂ NTAs with carbon materials is also higher than that without. This may mean that electrons were better transported by π - π conjugation, which also increased the short circuit current. Moreover, the current intensity is the strongest in the presence of both Ag NPs and carbon materials. In this case, electrons are generated in large quantities and are better transported by plasmonic activities and π - π conjugation.

4. Conclusions

We deposited Ag NPs and carbon materials in the channels of closed- and open-ended TiO₂ NTAs using UV irradiation and CVD, respectively. These modifications improved the energy conversion efficiency of the corresponding DSSCs: the electron generation is enhanced by plasmonics from the Ag NPs, while the resistance of TiO₂ NTAs is suppressed via the π - π conjugation from the carbon materials. DSSCs made of freestanding TiO₂ NTAs coated with both Ag NPs and carbon materials had the best energy conversion efficiency, due to the combination of these two factors. Between the open-ended and closed-ended TiO₂ NTAs (both with Ag NPs and carbon materials), the energy conversion efficiency of the DSSCs was higher for the former.

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국 문 초 록

매우 규칙적인 구조를 가지는 이산화티타늄 나노튜브 수직 배열층을 전기화학적 방법으로 제조하였다. 이 수직 배열층의 바닥면은 벌집통의 바닥처럼 막혀있다. 이런 형태를 ‘끝이 닫힌 이산화티타늄 나노튜브 수직 배열층’이라고 하고, 이 배열층 바닥면을 이온밀링의 방법으로 깎아내면 위 아래가 완전히 관통된 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’이 만들어진다.

제조된 이산화티타늄 나노튜브 수직 배열층을 염료감응형 태양전지에 도입하여 기능을 강화하였다. 이어지는 연구는 이 구조에 여러 가지 변인들의 효과를 알아보는 과정이다.

첫 번째 연구는 염료감응형 태양전지에 이산화티타늄 나노튜브 수직 배열층을 도입하는 것이다. 수직으로 배열된 이산화티타늄들에 보다 많은 염료가 부착되고 나노입자일 때보다 grain boundary 손실을 줄이고 전자의 재결합을 줄여서 태양전지의 에너지 전환 효율을 증대시킬 수 있다.

이를 위해 ‘끝이 닫힌 이산화티타늄 나노튜브 수직 배열층’과 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’을 도입한 결과를 각각 비교하였다.

‘끝이 닫힌 이산화티타늄 나노튜브 수직 배열층’보다 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’을 도입할 때 에너지 변환 효율이 5.15% 더 많이 증가되었다(5.92% 대 5.92%). 여기에 비교적 입자의 크기가 큰 이산화티타늄 나노입자로 부가적인 산란층을 더 만들어준 경우 에너지 전환 효율은 10.30%가 더 증가되었다(5.92% 대 6.53%).

추가적으로 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’에 사염화티타늄 용액 처리를 해주었더니 에너지 전환 효율이 16.38% 더 증가하였다(5.92%에서 6.89%로). 이것은 염료의 빛 흡수율 증가와 빛 수확 비율이 개선된 결과이다.

두 번째 연구는 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’을 산란층으로 염료감응 태양전지에 도입하고 여기에 은 나노입자들을 수직배열층 채널들 사이사이에 생성시키는 것이다. 이 은 나노입자들은 태양전지가 작동할 때 플라즈몬 효과를 나타내게 되어 빛의 흡수율을 높이고, 아울러 추가적으로 도입된 비교적 입자의 크기가 큰 이산화티타늄 나노입자들에 의한 산란효과가 추가되어 결국 전체 에너지 전환 효율이 높아진다.

비교실험으로 먼저 ‘끝이 닫힌 이산화티타늄 나노튜브 수직 배열층’이 도입된 염료감응 태양전지를 만들어 에너지 전환 효율을 측정하였다. 그 후에 비교적 입자의 크기가 큰 이산화티타늄 나노입자를 산란층으로 추가 도입하였더니 6.48%(5.86%에서 6.24%로)의 효율 증가가 나타났다.

한편, 은 나노입자를 추가 도입하였더니 9.21%(5.86%에서 6.40%로)의 효율 증가가 나타나서 이산화티타늄 나노입자에 의한 추가적인 산란층 도입 효과보다는 은 나노입자를 생성시켜서 얻는 플라즈몬 효과가 더 큰 것을 확인하였다.

두 가지 변인 즉, 은 나노입자에 의한 플라즈몬 효과와 이산화티타늄 나노입자에 의한 추가적 산란 효과를 ‘끝이 닫힌 이산화티타늄 나노튜브 수직배열층’이 도입된 염료감응 태양전지에 함께 적용한 결과 14.50%(5.86%에서 6.71%로)의 에너지 효율증가를 확인하였다.

이번에는 ‘끝이 열린 이산화티타늄 나노튜브 수직배열층’이 도입된 염료감응 태양전지에 동일한 변인들의 효과를 조사하였다. 그 결과 비교적 큰 이산화티타늄 나노입자층을 산란층으로 추가 도입하여 8.17% (6.12%에서 6.62%로)의 에너지 전환 효율 증가를, 은 나노입자의 도입을 통한 플라즈몬 효과로 9.15%(6.12에서 6.68%로)의 에너지 전환 효율 증가를 확인하였고, 두 변인을 모두 적용하면 15.20% (6.12%에서 7.05%로)의 에너지 전환 효율 증가를 나타내었다.

이산화티타늄 나노튜브 수직 배열층의 바닥이 열려있는지 여부가 에너지 전환 효율과 얼마나 관련이 있는지를 비교하면 끝이 닫힌 수직 배열층보다 열린 경우에 더 에너지 전환 효율이 높았다(6.71% 대 7.05%).

결론으로 ‘끝이 열린 이산화티타늄 나노튜브 수직배열층’을 염료감응형 태양전지에 도입하고, 은 나노입자를 채널 속에 고르게 생성시켰으며, 여기에 더해 비교적 입자의 크기가 큰 이산화티타늄 나노입자를 추가 산란층으로 도입하였을 때, 산란효과와 플라즈몬 효과가

함께 증대되어 에너지 전환 효율이 증가된다.

세 번째 연구는 이산화티타늄 나노튜브 수직 배열층에 은 나노입자와 함께 탄소 물질을 도입하는 것이다. 먼저 선행 연구와 같이 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’을 염료감응 태양전지에 도입하고 은 나노입자들을 채널들 사이사이에 생성시키고 에너지 전환 효율을 측정하여 15%(5.32%에서 6.14%로)의 증가를 확인하였다. 이것은 플라즈몬 효과에 기인한다. 이번에는 ‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’이 도입된 염료감응 태양전지에 (고온의 탄화 수소를 산화시키는 방법으로) 탄소물질을 생성시켜 만든 태양전지는 에너지 전환 효율이 14%(5.32%에서 6.07%로) 증가하였다. 이것은 파이-파이 공액(공명, 결합)의 결과이다.

‘끝이 열린 이산화티타늄 나노튜브 수직 배열층’이 도입된 염료감응 태양전지에 은 나노입자와 탄소물질이 모두 생성된 경우, 그렇지 않은 경우에 비해 30%(5.32%에서 6.91%로)의 에너지 전환 효율 증가를 나타내었다. 이것은 플라즈몬 효과와 파이-파이 공액(공명, 결합)의 효과가 함께 작용한 결과이다.

Keywords: open-ended freestanding TiO_2 nanotube arrays, dye-sensitized solar cells, plasmonic, scattering, anodization

