

# Photothermal spectroscopic characterization in tellurite glasses codoped with rare-earth ions

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## ABSTRACT

Thermal Lens (TL) and spectroscopic characterizations were performed in 70TeO<sub>2</sub>-19WO<sub>3</sub>-7Na<sub>2</sub>O-4Nb<sub>2</sub>O<sub>5</sub> (mol%) tellurite glasses. TL measurements were accomplished in Er<sup>3+</sup>/Tm<sup>3+</sup> co-doped tellurite glasses in function of the Tm<sub>2</sub>O<sub>3</sub> concentration (0.4 -1.6 x10<sup>20</sup> ions/cm<sup>3</sup>). Fluorescence spectra at 488 nm showed that Er<sup>3+</sup>/Tm<sup>3+</sup> co-doped tellurite glasses present several emission bands between (500-1800) nm. However, the more intense emission bands correspond to the Er<sup>3+</sup> and Tm<sup>3+</sup> transitions (<sup>4</sup>I<sub>13/2</sub> → <sup>4</sup>I<sub>15/2</sub> and <sup>3</sup>F<sub>4</sub> → <sup>3</sup>H<sub>6</sub>), respectively. The absolute nonradiative quantum efficiency (φ) was determined by TL method. Higher values of φ were obtained with the increase of Tm<sub>2</sub>O<sub>3</sub> concentration inside of the Er<sup>3+</sup>/Tm<sup>3+</sup> co-doped tellurite glasses. These results are corroborated by the Judd-Ofelt calculations.

Keyword: Er<sup>3+</sup>/Tm<sup>3+</sup> co-doped tellurite glasses, Thermal Lens, Spectroscopic properties, Judd-Ofelt

## 1. INTRODUCTION

Tellurite glasses are materials that present high rare earth ions solubility, large amplification bandwidth and higher refractive indices than the silicates and fluoride glasses, enhancing the radiative rate and emission cross-section<sup>[1-5]</sup>. In this way, these glasses have generated an increasing interest, becoming promising for applications as: optical amplifiers, optical recording, laser active media and infrared-to-visible converters<sup>[1-9]</sup>. Recently, energy transfers processes have been studied in rare earth doped and co-doped glasses using fluorescence spectroscopy<sup>[4-10]</sup>. The work is motivated by the interest in basic research phenomena and development of new solid state short-wavelength laser materials. In this way, it is very important the characterization of nonlinear optical and thermal-optical properties of the glasses, for optimization of the application of these materials.

Recently, several thermal studies have been made in different kind of glass host structures using Thermal Lens (TL) technique<sup>[11-16]</sup>. However, only Er<sup>3+</sup>, Tm<sup>3+</sup> or Yb<sup>3+</sup> were embedded in fluoroindate, silica calcium aluminosilicate, chalcogenide and phosphate glasses. In this work, we studied the spectroscopic and thermooptics properties of Er<sup>3+</sup>/Tm<sup>3+</sup> co-doped tellurite glasses. The measurements were made in function of the Tm<sub>2</sub>O<sub>3</sub> concentration (0.4 -1.6 x10<sup>20</sup> ions/cm<sup>3</sup>), with the purpose of producing multiphonon relaxation and determining the values of absolute nonradiative quantum efficiency (φ). Fluorescence spectra were obtained and the measures are in agreement with the results of Thermal Lens (TL) technique and Judd-Ofelt (JO) calculations.

## 2. THEORY

The TL effect is created when the excitation laser beam passes through the sample and the absorbed energy is converted into heat, changing the optical path length (s) and producing a lenslike optical element at the sample. The

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propagation of a probe laser beam (He-Ne laser) through the TL will result in either a spreading ( $ds/dT < 0$ ) or a focusing ( $ds/dT > 0$ ) of the beam, depending mainly on the sample temperature coefficients of electronic polarizability, stress and thermal expansion. TL effect can be treated through the calculation of the temporal evolution of the sample temperature profile  $\Delta T(r,t)$  caused by a Gaussian intensity distribution of the excitation beam. The variation of probe beam on-axis intensity,  $I(t)$ , can be calculated in *cw* excitation regime, in the form <sup>[11,17,18]</sup>:

$$I(t) = I(0) \left[ 1 - \frac{\theta}{2} \tan^{-1} \left( \frac{2mV}{[(1+2m)^2 + V^2] \tau_c / 2t + 1 + 2m + V^2} \right) \right]^2 \quad (1)$$

where  $m = (\omega_p/\omega_e)^2$ ,  $\omega_p$  and  $\omega_e$  are respectively the probe and excitation beam radius at the sample;  $V = z_1/z_0$ ,  $z_1$  is the distance between the sample and probe beam waist and  $z_0$  is the probe beam Rayleigh range;  $I(0)$  is the on-axis intensity when  $t$ , or  $\theta$ , is zero. In the dual beam (excitation and probe beams), mode-mismatched configuration, the transient signal amplitude,  $\theta$ , is approximately the phase difference of the probe beam at  $r = 0$  and  $r = \sqrt{2} \omega_e$  induced by the pump beam, given by:  $\theta = \Theta P_e \alpha L_{\text{eff}}$ . In this case,  $P_e(W)$  is the excitation power,  $\alpha$  ( $\text{cm}^{-1}$ ) the optical absorption coefficient at the excitation wavelength ( $\lambda_e$ ),  $L_{\text{eff}} = (1 - e^{-\alpha L})/\alpha$  is the effective length and  $L(\text{cm})$  is the sample thickness. The normalized phase shift,  $\Theta$ , can be expressed as <sup>[17]</sup>:

$$\Theta = - \frac{1}{K \lambda_p} \frac{ds}{dT} \left[ 1 - \eta \frac{\lambda_e}{\langle \lambda_{\text{em}} \rangle} \right] \quad (2)$$

where  $K$  is the thermal conductivity ( $\text{W/cmK}$ ),  $\lambda_p$  is the probe beam wavelength and  $ds/dT$  is the optical path temperature coefficient normalized by sample lengths ( $\text{K}^{-1}$ ) and  $\langle \lambda_{\text{em}} \rangle$  the average fluorescence wavelength. The term in parenthesis on the right-hand side of Eq. (2) is the fraction of absorbed energy converted into heat ( $\phi$ ) or absolute nonradiative quantum efficiency (the complementary part is converted into fluorescence). For direct excitation,  $\eta$  is defined as the ratio between the measured lifetime ( $\tau_{\text{exp}}$ ) and the radiative lifetime ( $\tau_{\text{rad}}$ ),  $\eta = \tau_{\text{exp}}/\tau_{\text{rad}}$ .

For other side,  $\tau_c = \omega_e^2/4D$  is the characteristic heat diffusion time, where  $D = K/\rho C$  is the thermal diffusivity ( $\text{cm}^2/\text{s}$ ),  $\rho$  is the density ( $\text{g/cm}^3$ ) and  $C$  is the specific heat ( $\text{J/gK}$ ).

### 3. EXPERIMENTAL

Samples studied were prepared in the proportion  $70\text{TeO}_2\text{-}19\text{WO}_3\text{-}7\text{Na}_2\text{O-}4\text{Nb}_2\text{O}_5$  (% mol). For  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses were added  $1.2 \times 10^{20}$  ions/ $\text{cm}^3$  of  $\text{Er}_2\text{O}_3$  and the concentration of  $\text{Tm}_2\text{O}_3$  was varied between  $(0.4\text{-}1.6) \times 10^{20}$  ions/ $\text{cm}^3$ . This raw material was melted and homogenized in platinum crucible at  $800^\circ\text{C}$  during 30 minutes using a RF induction furnace. All absorbance spectra measurements were carried out with a Perkin Elmer Lambda 9 spectrophotometer at room temperature, with resolution of 1 nm. The fluorescence spectra were measured by a monochromator Spectral Energy connected to either a Silicon PIN or an InAs detector for measurements from (500-1100) nm or (1000-1800) nm, respectively. The samples were excited with an  $\text{Ar}^+$  laser at  $\lambda_e = 488$  nm.

The thermo-optic properties of tellurite glasses were investigated by TL method <sup>[11,17]</sup>. TL transient measurements were performed using the mode-mismatched dual-beam (excitation and probe) configuration. A He-Ne laser ( $\lambda_p = 632.8$  nm) was used as the probe beam and either an  $\text{Ar}^+$  laser ( $\lambda_e = 488$  nm) or a Ti-sapphire laser ( $\lambda_e = 785$  nm) was used as the excitation beam. The absorption of the relatively intense excitation beam generates the TL heat profile and the induced phase shift, which is proportional to  $\theta$ . For other side,  $\theta$  is measured by the weak probe beam that counter propagates nearly collinear with the excitation beam. Details of the experimental setup can be found elsewhere <sup>[11,17,18]</sup>. The optical absorption coefficients ( $\alpha$ ) were determined applying the same experimental configuration used for realization of TL measurements. The reflections at the sample surface were considered in  $\alpha$  calculus.

Lifetime ( $\tau_{\text{exp}}$ ) measurements of the  ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$  transition were done using a diode laser at  $\lambda_e = 980$  nm ( $P_e \approx 120$  mW). The signal detected by a InGaAs fast response detector, after passing through a bandwidth centred at 1550 nm band pass filter, was acquired using a digital oscilloscope.

## 4. RESULTS AND DISCUSSIONS

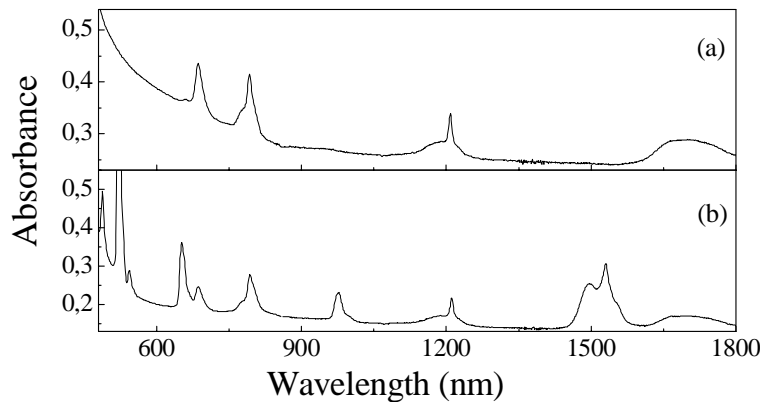
### 4.1. Spectroscopic properties

The tellurite glasses studied in this paper are presented at Table 1. Absorbance spectra of  $\text{Tm}^{3+}$  ( $0.78 \times 10^{20}$  ions/cm<sup>3</sup>) doped tellurite and  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite ( $S_{1.2,0.4}$ ) glasses are showed in the Fig. 1 (a-b). The absorption bands of  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses (Fig. 1(b)) that do not appear in Fig. 1 (a) are characteristics of the  $\text{Er}^{3+}$  ion.

Tellurite	$\text{Er}^{3+}$ ( $10^{20}$ ions/cm <sup>3</sup> )	$\text{Tm}^{3+}$ ( $10^{20}$ ions/cm <sup>3</sup> )	$\xi$	$\eta_{\text{JO}}$
$S_{0,0}$	0	0	-	-
$S_{1,2,0}$	1.2	0	220	0.93
$S_{1,2,0.4}$	1.2	0.39	13	0.34
$S_{1,2,0.8}$	1.2	0.78	5	0.28
$S_{1,2,1.2}$	1.2	1.2	4	0.18
$S_{1,2,1.6}$	1.2	1.6	3	0.15

**Table I-** Composition of the analyzed tellurite glasses.  $\xi$  is the ratio between the fluorescence intensities of the  $T_{\text{Er}}$  and  $T_{\text{Th}}$  transitions.  $\eta_{\text{JO}}$  is the radiative quantum efficiency obtained by Judd-Ofelt method.

Figure 2 and 3 show the fluorescence spectra of the  $\text{Er}^{3+}$  and  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses. The measurements were made under low power laser  $P_e \approx 10$  mW. The matrix of tellurite glass without rare earth ions presented no emission band with intensity comparable to those that measurements were performed.  $\text{Tm}^{3+}$  doped tellurite glass presented emission bands in  $\sim 804$ ,  $1470$  and  $1780$  nm. For  $\text{Er}^{3+}$  doped tellurite glasses, the more intense emission band is centered at  $1528$  nm (Fig. 3 (a)).  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped glasses presented suitable emission bands with those measured with the tellurite glass doped either with erbium or thulium (Fig. 2 e 3 (b-c)). In this case, it is possible to observe that are more intense, the emission bands correspondent to the transition of  $\text{Er}^{3+}$ :  $^4I_{13/2} \rightarrow ^4I_{15/2}$  ( $T_{\text{Er}}$ ) and of  $\text{Tm}^{3+}$ :  $^3F_4 \rightarrow ^3H_6$  ( $T_{\text{Th}}$ ), presented in the Fig. 3. The ratio of the fluorescence intensities ( $\xi$ ) of the  $T_{\text{Er}}$  and  $T_{\text{Th}}$  transitions, are presented in the Table I. For  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glass with highest concentration of thulium ( $S_{1.2,1.6}$ ) presents the emission intensity  $T_{\text{Er}}$  just 3 times larger than the emission of  $T_{\text{Th}}$  transition. In general form, for  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses, the emission bands between (500-800) nm and (850-1400) nm are respectively  $\sim 100$  and 5 times less intense than the  $T_{\text{Er}}$  and  $T_{\text{Th}}$  emissions. In this form, the emission bands at wavelength  $< 1400$  nm could be neglectful as a good approach.



**Figure 1-** Absorbance ( $-\ln(I/I_0)$ ) for: (a)  $\text{Tm}^{3+}$  ( $0.78 \times 10^{20}$  ions/cm<sup>3</sup>) doped tellurite glass ( $L = (1.25 \pm 0.03)$  mm) and (b)  $S_{1.2,0.4}$  glass ( $L = (1.75 \pm 0.03)$  mm).

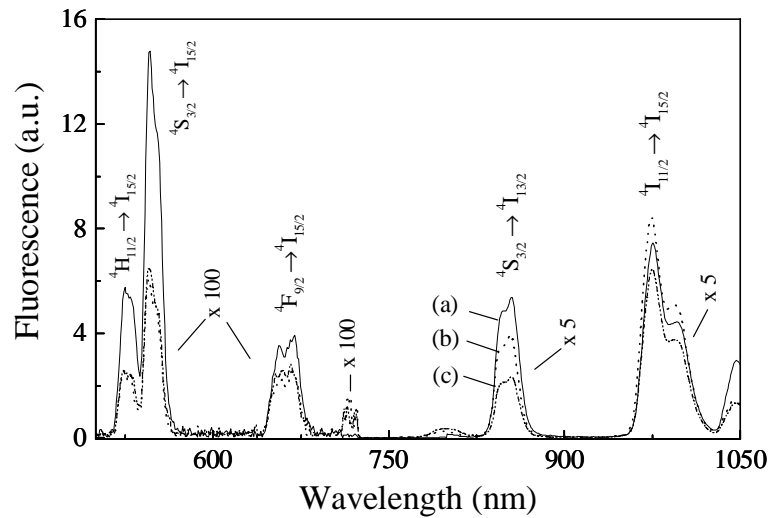


Figure 2- Fluorescence spectra between (500-1050) nm for (a)  $S_{1,2,0}$ ; (b)  $S_{1,2,0,4}$  and (c)  $S_{1,2,1,6}$  at  $\lambda_e=488$  nm ( $P_e \approx 10$  mW).

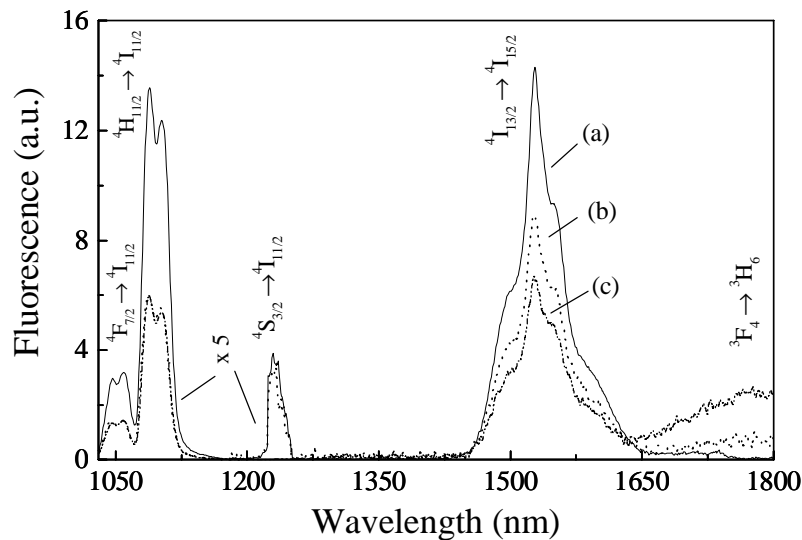


Figure 3- Fluorescence spectra between (1050-1800) nm for (a)  $S_{1,2,0}$ ; (b)  $S_{1,2,0,4}$  and (c)  $S_{1,2,1,6}$  at  $\lambda_e=488$  nm ( $P_e \approx 10$  mW).

#### 4.2. Thermal lens measurements

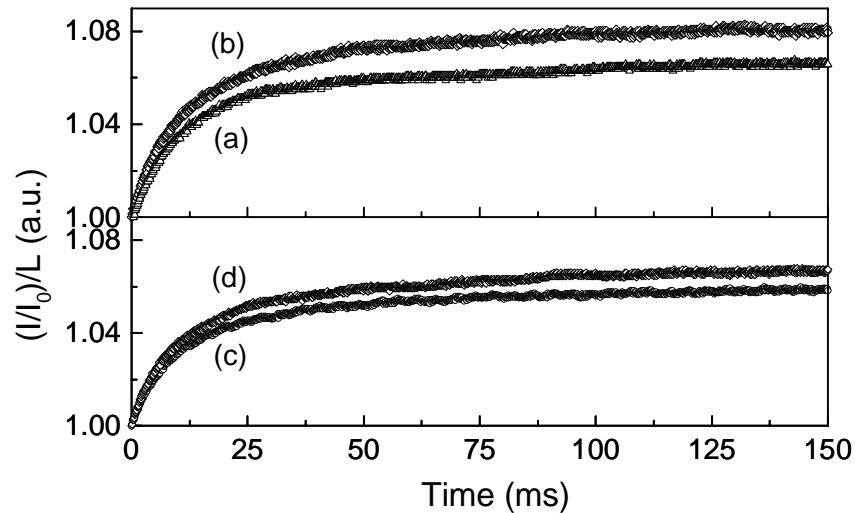
Figure 4 shows typical TL transient signals for the  $\text{Er}^{3+}$  doped,  $\text{Tm}^{3+}$  doped and  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses. The curve behaviors of the Fig. 4 indicate that  $ds/dT$  is positive, i.e., the created TL focalize the probe beam in the far field. Fitting the experimental data of Fig. 4 by Eq. (1),  $\theta$  and  $\tau_c$  were obtained. From  $D = w_e^2/4\tau_c$  and using the measured value  $w_e = 2.6 \times 10^{-3}$  cm, the thermal diffusivity was determined. The pattern sample, whose thermal properties are well known, has been used to test the calibration of our optical system.

The medium value of  $D$  obtained for  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite samples in function of thulium concentration, at  $\lambda_e = 488$  nm and 785 nm, is  $D = (3.1 \pm 0.2) \times 10^{-3} \text{ cm}^2/\text{s}$ . This result of  $D$  is similar to the values obtained for us to  $\text{Er}^{3+}$  doped tellurite glass ( $D = (3.1 \pm 0.2) \times 10^{-3} \text{ cm}^2/\text{s}$ ),  $\text{Tm}^{3+}$  doped tellurite glass ( $D = (3.2 \pm 0.4) \times 10^{-3} \text{ cm}^2/\text{s}$ ), and for others glasses obtained in literature, as: ZBLAN, YABC, PGIZCa, ISZn and InSBZnGdN [11,19].

In order to calculate  $\phi$ , we supposed a negligible fluorescence to the tellurite glass without earth rare ion ( $S_{0,0}$ ), where all absorbed energy is converted into heat by the sample, i.e.,  $\phi_{S_{0,0}} = 1$  ( $\eta_{S_{0,0}} = 0$ ).  $\Theta_{S_{0,0}}$  value was determined through transient thermal lens measurements. However, normalizing the thermal parameters of erbium (E) and thulium (T) co-doped tellurite samples ( $S_{E,T}$ ), i.e.  $\Theta_{S_{E,T}}$  by  $\Theta_{S_{0,0}}$  value for the undoped tellurite matrix, we obtain:

$$\phi = 1 - \eta \left( \frac{\lambda_e}{\langle \lambda_{em} \rangle} \right) = \frac{\Theta_{S_{E,T}}}{\Theta_{S_{0,0}}} \quad (3)$$

where  $\Theta_{S_{E,T}} = \theta_{S_{E,T}} / P_e \alpha L_{\text{eff}}$  and  $\theta_{S_{E,T}}$  is the amplitude of TL signal for  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glass. In this case,  $\Theta_{S_{E,T}}$  was changed by  $\Theta_{S_{1,2,0}}$ ,  $\Theta_{S_{1,2,0.4}}$ ,  $\Theta_{S_{1,2,0.8}}$ ,  $\Theta_{S_{1,2,1.2}}$  or  $\Theta_{S_{1,2,1.6}}$ , that correspond the samples  $S_{1,2,0}$ ,  $S_{1,2,0.4}$ ,  $S_{1,2,0.8}$ ,  $S_{1,2,1.2}$  and  $S_{1,2,1.6}$ , respectively (Table I).  $\phi$  results obtained by TL method are presented in Fig. 5.

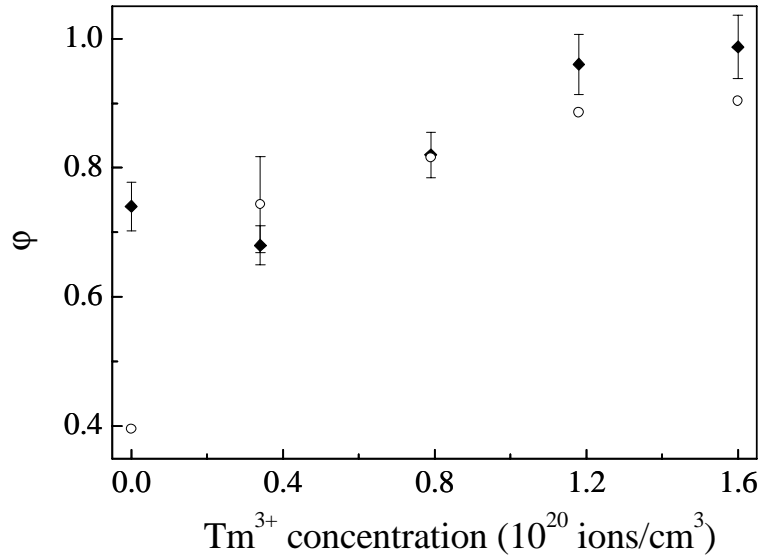


**Figure 4** TL transient signal normalized by  $L$  versus time for  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses: (a)  $S_{0,0}$ , (b)  $\text{Tm}^{3+}$  ( $0.78 \times 10^{20}$  ions/ $\text{cm}^3$ ) doped tellurite glass, (c)  $S_{1,2,0}$  and (d)  $S_{1,2,0.4}$ , respectively. The samples were excited with an  $\text{Ar}^+$  laser at 488 nm with  $P_e = 20$  mW and probe laser beam at  $\lambda_p = 632.8$  nm.

### 4.3. Judd-Ofelt (JO) results

The quantum efficiencies for  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped tellurite glasses were determined through the Judd-Ofelt (JO) model [5,20,21]. The radiative lifetime value ( $\tau_{\text{rad}}$ ) was calculated as being 5.6 ms and using the experimental measurements of lifetime  $\tau_{\text{exp}}$  at 980 nm, the fluorescence quantum efficiency was estimated by  $\eta_{\text{JO}} = \tau_{\text{exp}}/\tau_{\text{rad}}$  for  ${}^4\text{I}_{13/2} \rightarrow {}^4\text{I}_{15/2}$  (Table I). Then using the Eq. (3) considering  $\lambda_e = 980$  nm, we calculated  $\phi_{\text{JO}}$  values. The results of  $\phi_{\text{JO}}$  in function of thulium concentration for  $\text{Er}^{3+}/\text{Tm}^{3+}$  co-doped glasses are presented in the Fig. 5. There are good agreement between both TL and JO methods for  $\phi$ . For  $\text{Er}^{3+}$ -doped tellurite glass, the difference of the results should be due to resonance at 980 nm with  ${}^4\text{I}_{15/2} \rightarrow {}^4\text{I}_{13/2}$  transition and consequently a decreasing of cross-relaxation process and lose of energy by phonons.

Finally, Table I shows, by JO theory, an increase of radiative quantum efficiency  $\sim 6$  for the  $S_{1.2,0}$  glass with relation to the  $S_{1.2,1.6}$  sample. Then, the presence of thulium ( $1.6 \times 10^{20}$  ions/cm<sup>3</sup>) inside of the tellurite matrix co-doped with  $Er^{3+}/Tm^{3+}$ , corroborate with the lost of energy by phonons, producing an increase of the fraction of energy turned into heat, and consequently the quantum efficiency approaches zero.



**Figure 5-**  $\phi$  versus thulium concentration for  $Er^{3+}/Tm^{3+}$  co-doped tellurite glasses, obtained by TL and JO (open circle) methods.

## 5. CONCLUSIONS

Spectroscopic characterizations of the  $Er^{3+}/Tm^{3+}$  co-doped tellurite glasses were performed using Thermal Lens (TL) and fluorescence techniques at visible wavelength (488 and 785 nm). The behavior of the nonradiative quantum efficiency ( $\phi$ ) was determined in  $70TeO_2-19WO_3-7Na_2O-4Nb_2O_5$  (% mol) tellurite matrix. The measurements were made in function of the  $TmO_3$  concentration in  $Er^{3+}/Tm^{3+}$  co-doped tellurite glasses and  $\phi$  values are in agreement with the results obtained by Judd-Ofelt theory.

## 6. ACKNOWLEDGMENTS

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