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MANY FERMION GREEN FUNCTIONS AND DYNAMICAL ALGEBRA

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ABSTRACT

The mean-field Hamiltonian for a many-fermion system is an element in a Dynamical Algebra - a classical Lie Algebra. The thermal Matsubara and T=0 Green Functions are sums of products of factors determined by certain structure constants of the algebra. These describe the automorphisms of the Algebra.

In the mean-field approximation the Hamiltonian of an interacting many-fermion system is a sum of bilinear products of fermion creation/ destruction operators with complex coefficients. This applies for systems such as a singlet BCS superconductor; and a coexisting charge-density-wave superconductor; and coexisting charge - and spin-density-wave superconductor. The Hamiltonian can be taken as

$$H = \sum_k H_k$$

where  $(k)$  is a wave-vector index. We define a basis  $(A_i)$  by relabelling the standard Fermi operators as  $(A_1, A_2, \dots) = (a_{k\uparrow}, a_{-k\downarrow}^\dagger, \dots)$ . Then  $H_k = \sum_{ij} \lambda_{ij} X_{ij}$  with  $X_{ij} = A_i^\dagger A_j$ . The matrix  $\lambda$  with  $(\lambda)_{ij} = \lambda_{ij}$  represents the complex coupling coefficients. The set of all pairs  $X_{ij}$  closes under Lie Bracket:  $[X_{ij}, X_{\ell n}] = X_{im} \delta_{\ell j} - X_{\ell j} \delta_{im}$ . If  $(i, j) = 1 \dots N$  the algebra is an  $SO(4N)$ , (or D series) algebra. The set of all singles  $(A_i)$  and  $(A_i^\dagger)$  plus all pairs  $(X_{ij})$  generates an  $SO(4N+1)$  (or

B series) algebra under Lie Bracket. Evidently the Hamiltonian  $H_k$  in the mean-field approximation is an element in the  $SO(4N)$  algebra. If we pass to a Cartan-Weyl realization of the basis, then  $H_k$  can be re-written as

$$H_k = \sum_j \mu_j h_j + \sum_{\alpha^\pm} \mu_\alpha e_\alpha \quad (1)$$

where the  $(h_j, e_\alpha)$  are the generators in C-W form and  $(\mu_j, \mu_\alpha)$  are some combinations of  $\lambda_{ij}$ .

The  $T=0$  Green Function can be written, (suppressing spin) as

$$G(k,t) = -i \langle \phi' | T_t a_k(t) a_k^\dagger | \phi' \rangle \quad (2)$$

here  $|\phi'\rangle$  is the (mean-field) ground state of  $H_k$ , and

$$a_k(t) = (\exp i H_k t) a_k (\exp -i H_k t) \quad (3)$$

Using the  $(A_i)$  notation we find

$$a_k(t) = A_r(t) = \sum_j (e^{-it\lambda})_{rj} A_j \quad (4)$$

so,

$$a_k(t) a_k^\dagger \equiv A_r(t) A_r^\dagger = \sum_j (e^{-it\lambda})_{rj} X_{rj}^\dagger \quad (5)$$

Examining the above we can identify the matrix elements  $(e^{-it\lambda})_{rj}$  as "structure constants" of an "Heisenberg" automorphism  $\phi$ :

$A_r \xrightarrow{\phi} \{A_r\}$  of single fermion operators onto themselves in the B algebra.

The  $H_k$  can be rotated to "diagonal" form:  $U H_k U^{-1} = \sum_j \gamma_j h_j \equiv H'_k$ , by means of a unitary transformation  $U = \exp i (\sum_{\alpha} \theta_{\alpha} e_{\alpha}^j)$ , with coefficients  $\theta_{\alpha}$ . We can identify the coefficients  $\gamma_j$  as a second set of "structure constants" of an automorphism  $\phi'$  of pairs of fermion operators onto themselves:  $(h_j, e_{\alpha}) \xrightarrow{\phi'} (h_j)$  produced by the mapping  $U$ .

After rotation  $H_k$  is a sum of generators  $(h_j)$  of the Cartan sub-algebra of D. It is natural to label the kets  $|\{\lambda\}\rangle$  of  $H'_k$  by eigenvalues of the  $h_j$ :  $h_j |\{\lambda'\}\rangle = \lambda'_j |\{\lambda'\}\rangle$  where  $(\lambda')$  are the lowest eigenvalues. In terms of the eigenfunction  $|\{\lambda'\}\rangle$  the state  $|\phi'\rangle$  is  $|\phi'\rangle = U^{-1} |\{\lambda'\}\rangle$ .

The Green Function  $G(k,t)$  can be written for  $t>0$  as

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$$-i \sum_j (e^{-it\lambda})_{rj} \langle \{\lambda'\} | U X_{rj}^\dagger U^{-1} | \{\lambda'\} \rangle \quad (6)$$

But  $U X_{rj}^\dagger U^{-1} \Rightarrow U(h_\ell, e_\beta) U^{-1}$  which is again an automorphism in  $D$   
 INDEPT  $\{h_\ell, e_\beta\} \phi' \{h_\ell, e_\beta\}$  with structure constants  $(\gamma)$

$$U(h_\ell, e_\beta) U^{-1} = \sum_j \gamma_j h_j + \sum_{\alpha \pm} \gamma_\alpha e_\alpha \quad (7)$$

Taking the diagonal matrix element of (7) only the terms  $\sum_j \gamma_j h_j$  contribute. Then for  $t > 0$

$$G(k, t) = (-i) \sum_j (e^{-it\lambda})_{rj} \sum_\ell \gamma_\ell \lambda'_\ell \quad (8)$$

with a similar expression for  $t < 0$ . This result gives the factorization of  $G(k, t)$  into a sum of products of factors determined by the Dynamical Algebras (D and B), namely:

- (1) Structure constants  $(e^{-it\lambda})$  of  $\phi$ ;
- (2) Structure constants  $\gamma_\ell$  of  $\phi'$ ;
- (3) Eigenvalues  $\{\lambda'\}$  of the operators  $h_\ell$  of the Cartan sub-algebra.

A similar factorization applies for the  $T=0$  Gor'kov (anomalous) function  $F(k, t) = \langle \phi' | T_t a_k(t) a_{-k} | \phi' \rangle$

The Thermal Matsubara Green Function ( $T \neq 0$ ) is defined as

$$G(k, \tau) \equiv - \langle T_\tau a_k(\tau) a_k^\dagger \rangle \quad (9)$$

with  $(\tau = it)$  and in this case the thermal average is defined as

$$\langle \hat{O} \rangle \equiv \text{Tr} (e^{-\beta H \hat{O}}) / \text{Tr} e^{-\beta H} \quad (10)$$

A similar analysis as in the  $T=0$  case gives  $(\tau > 0)$ .

$$G(k, \tau) = \sum_j (e^{-\tau\lambda})_{rj} \left( \sum_{\{\lambda_p\}} \left( \sum_p e^{-\beta E_p \lambda_p} \gamma_p \lambda_p \right) / \sum_{p\{\lambda_p\}} e^{-\beta E_p \lambda_p} \right) \quad (11)$$

with a similar expression for  $\tau < 0$ . Again,  $G(k, \tau)$  is constructed as a sum of terms each of which is composed of factors which are symmetry

determined from structure constants of automorphisms of the Dynamical Algebra and eigenvalues of the generators ( $h_\ell$ ). Mutatis mutandis the thermal Gor'kov anomalous function has a similar structure.

Details of the analysis and illustrations for simple BCS-like theory of superconductivity are given elsewhere<sup>1), 2)</sup>.

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