

## Special linear group and generating functions

by

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### 1. Introduction

In this paper we aim at obtaining certain generating functions by using Lie theory method. The process in short involves introducing linear differential operators which form a 3-dimensional Lie algebra isomorphic to the Lie algebra  $sl(2)$  [2, p. 8]. Based on these operators local multiplier representation  $[T(g)f](x, y)g = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2)$ , is determined. Thereafter, by choosing  $f(x, y)$  in certain ways this multiplier representation leads to generating functions for Laguerre functions.

### 2. Differential operators and multiplier representation

$$(2.1) \quad u(x) = L_{m+n}^{(\gamma)}(x)$$

is a solution of [3; p. 204]

$$(2.2) \quad x \frac{d^2 u}{dx^2} + (1 + \gamma - x) \frac{du}{dx} + (m + n)u = 0.$$

By substituting  $y(\partial/\partial y)$  for  $n$  we construct the following partial differential equation.

$$(2.3) \quad \left[ x \frac{\partial^2}{\partial x^2} + (1 + \gamma - x) \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + y \right] f(x, y) = 0$$

$f(x, y) = y^n u(x)$  is a solution of (2.3). Now, we introduce the first order partial differential operators

$$(2.4) \quad \begin{aligned} J^3 &= y \frac{\partial}{\partial y} + m + \frac{\gamma + 1}{2}, \\ J^+ &= xy \frac{\partial}{\partial x} + y^2 \frac{\partial}{\partial y} + (m + \gamma + 1 - x)y, \\ J^- &= xy^{-1} \frac{\partial}{\partial x} - \frac{\partial}{\partial y} - my^{-1}, \end{aligned}$$

obeying the commutation relations

$$(2.5) \quad [J^3, J^\pm] = \pm J^\pm, \quad [J^+, J^-] = 2J^3.$$

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These  $J$ -operators from the basis of a Lie algebra isomorphic to the Lie algebra  $sl(2)$  [2].

The Casimir operator

$$(2.6) \quad C = J^+ J^- + J^3 J^3 - J^3 \\ = x^2 \frac{\partial^2}{\partial x^2} + (1 + \gamma - x) \frac{\partial}{\partial x} + xy \frac{\partial}{\partial y} + mx + \frac{\gamma^2 - 1}{4}$$

commutes with  $J^3$ ,  $J^+$  and  $J^-$ . (2.3) may be rewritten as

$$(2.7) \quad Cf(x, y) = \left( \frac{\gamma^2 - 1}{4} \right) f(x, y).$$

To determine the multiplier representation induced by  $J$  operators, we need to compute the expression [2; p. 3]  $e^{a'J^-} e^{b'J^+} e^{c'J^3}$

$$(2.8) \quad e^{a'J^-} e^{b'J^+} e^{c'J^3} f(x, y) = \exp \left\{ \left( \frac{m + \gamma + 1}{2} \right) c' \right\} \\ \times \exp \left\{ \frac{-b'xy}{(1 + a'b' - b'y)} \right\} (1 + a'b' - b'y)^{-(\gamma+1+m)} \left( 1 - \frac{a'}{y} \right)^m \\ \times f \left( \frac{xy}{(y - a')(1 + a'b' - b'y)}, \frac{y - a'}{(1 + a'b' - b'y)} e^{c'} \right).$$

The complex parameters  $a'$ ,  $b'$  and  $c'$  are related to  $g \in SL(2)$ ,

$$g = \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \quad ab - bc = 1 \quad \text{by [2; p. 8]} \\ e^{c'/2} = a, \quad a' = -c/a, \quad b' = -ab.$$

Therefore, for  $g$  in a sufficiently small neighbourhood of the identity element  $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \in SL(2)$

$$(2.9) \quad [T(g)f](x, y) = \exp \left( \frac{bxy}{d + by} \right) (d + by)^{-(\gamma+1+m)} \left( a + \frac{c}{y} \right)^m \\ \times f \left( \frac{xy}{(c + ay)(d + by)}, \frac{c + ay}{d + by} \right),$$

$$|by/d| < 1, \quad -\pi < \arg a, \quad \arg d < \pi, \quad ad - bc = 1.$$

### 3. Generating functions

A. We choose  $f(x, y)$  to be a common eigenfunction of the operators  $C$  and  $J^3 J^3 + (\gamma' - \gamma - 2m - 1)J^3 - J^+$ .

Let  $f(x, y)$  satisfy the simultaneous equations

$$(3.1) \quad Cf(x, y) = \left( \frac{\gamma^2 - 1}{4} \right) f(x, y),$$

$$[J^3 J^3 + (\gamma' - \gamma - 2m - 1)J^3 - J^+] f(x, y) = \left( \frac{m + \gamma + 1}{2} \right) \left( \gamma' - m - \frac{\gamma + 1}{1} \right) f(x, y)$$

which may be rewritten as

$$(3.2) \quad \left[ x \frac{\partial^2}{\partial x^2} + (\gamma + 1 - x) \frac{\partial}{\partial x} - y \frac{\partial}{\partial y} - (m + \gamma + 1) \right] e^x f(-x, y) = 0,$$

$$\left[ y \frac{\partial^2}{\partial y^2} + (\gamma' + 1 - y) \frac{\partial}{\partial y} - x \frac{\partial}{\partial x} - (m + \gamma + 1) \right] e^x f(-x, y) = 0.$$

These equations have a solution [1; 234]

$$(3.3) \quad e^x f(-x, y) = \psi_2(m + \gamma + 1; \gamma + 1, \gamma' + 1; x, y)$$

where

$$\psi_2(\alpha; \gamma, \gamma'; x, y) = \sum_{m, n=0}^{\infty} \frac{(\alpha)_{m+n} x^m y^n}{m! n! (\gamma)_m (\gamma')_n}.$$

We rewrite (3.3) as

$$(3.4) \quad f(x, y) = e^x \psi_2(m + \gamma + 1; \gamma + 1, \gamma + 1; -x, y).$$

Therefore

$$(3.5) \quad [T(g)f](x, y) = \exp\left(\frac{axy}{c+ay}\right) (d+by)^{-(\gamma+1+m)}$$

$$\times \left(a + \frac{c}{y}\right)^m \psi_2\left[m + \gamma + 1; \gamma + 1, \gamma' + 1; \frac{-xy}{(c+ay)(d+by)}, \frac{c+ay}{(d+by)}\right],$$

$$\left|\frac{by}{d}\right| < 1, \quad -\pi < \arg a, \quad \arg d < \pi, \quad ad - bc = 1.$$

$[T(g)f](x, y)$  satisfies

$$(3.6) \quad C[T(g)f](x, y) = \left(\frac{\gamma^2 - 1}{4}\right) [T(g)f](x, y).$$

(3.5) has an expansion of the form

$$(3.7) \quad [T(g)f](x, y) = \sum_{n=-\infty}^{\infty} k_n(g) L_{m+n}^{(\gamma)}(x) y^n.$$

Putting  $x=0$ , this gives

$$(3.8) \quad k_n(g) = a^m (-b)^n d^{-(\gamma+m+n+1)} \frac{\Gamma(1+m+n)}{(1+\gamma)_m}$$

$$\times \sum_{p=0}^{\infty} \frac{(-m)_p (1+\gamma+m+n)_p (bc)^p}{p! \Gamma(1+n+p) (ad)^p}$$

$$\times {}_2F_2 \left[ \begin{matrix} -n-p, & 1+m \\ \gamma'+1, & 1+m-p \end{matrix}; \frac{a}{b} \right].$$

Thus the generating function (3.7) becomes

$$\begin{aligned}
(3.9) \quad & \exp\left(\frac{axy}{c+ay}\right)\left(1+\frac{by}{d}\right)^{-(\gamma+1+m)}\left(1+\frac{c}{ay}\right)^m \\
& \times \psi_2\left[m+\gamma+1; \gamma+1, \gamma'+1; \frac{-xy}{(c+ay)(d+by)}, \frac{c+ay}{d+by}\right] \\
& = \sum_{n=-\infty}^{\infty} \frac{\Gamma(1+m+n)(-b/d)^n y^n}{(1+\gamma)_m} L_{m+n}^{(\gamma)}(x) \sum_{p=0}^{\infty} \frac{(-m)_p (1+\gamma+m+n)_p}{p! \Gamma(1+n+p)} \\
& \times {}_2F_2\left[\begin{matrix} -n-p, & 1+m \\ \gamma'+1, & 1+m+p \end{matrix}; \frac{a}{b}\right], \\
& \left|\frac{c}{ay}\right| < 1, \left|\frac{by}{d}\right| < 1, ad-bc=1,
\end{aligned}$$

where the terms corresponding to  $n = -1, -2, \dots$  are well defined because of the relation

$$\begin{aligned}
(3.10) \quad & L_{\mu \rightarrow k} \sum_{p=0}^{\infty} \frac{(1+\gamma+m+\mu)_p (-m)_p (bc)^p}{p! \Gamma(1+\mu+p)} \left(\frac{bc}{ad}\right)^p {}_2F_2\left[\begin{matrix} -\mu-p, & -1-\gamma-m \\ \gamma'+1, & 1+m-p \end{matrix}; \frac{a}{b}\right] \\
& = \frac{(1+\gamma+m-k)_k (-m)_k (bc)^k}{k!} \sum_{p=0}^{\infty} \frac{(1+\gamma+m)(-m+k)_p (bc)^p}{p! (1+k)_p} \left(\frac{bc}{ad}\right)^p \\
& \times {}_2F_2\left[\begin{matrix} -p, & 1+m \\ \gamma'+1, & 1+m-p-k \end{matrix}; \frac{a}{b}\right], \quad k=1, 2, 3, \dots
\end{aligned}$$

*Special Cases:*

(3.9) gives the following special cases:

$$\begin{aligned}
(3.11) \quad & \frac{-xy}{e^{w-y}} \left(1-\frac{w}{y}\right)^m \psi_2\left(m+\gamma+1; \gamma+1, \gamma'+1; \frac{xy}{w-y}, w-y\right) \\
& = \sum_{n=-\infty}^{\infty} \frac{(1+n)_m}{(1+\gamma)_m (1+\gamma')_n} L_{m+n}^{(\gamma)}(x) {}_2F_2\left[\begin{matrix} 1+\gamma+m+n, & 1+\lambda-\gamma+n \\ 1+n, & \gamma'+1+n \end{matrix}; w\right]
\end{aligned}$$

and where the terms corresponding to  $n = -1, -2, \dots$  are well defined because of the relation of the type (3.10);

$$\begin{aligned}
3.12 \quad & e^x (1-y)^{-(\gamma+1+m)} \psi_2\left(m+\gamma+1; \gamma+1, \gamma'+1; \frac{-x}{1-y}, \frac{-wy}{1-y}\right) \\
& = \sum_{n=0}^{\infty} \frac{\Gamma(1+n+m)}{(\gamma+1)_m (\gamma'+1)_n} L_{m+n}^{(\gamma)}(x) L_n^{(\gamma')}(w) y^n
\end{aligned}$$

**B.** We choose  $f(x, y)$  to be a common eigenfunction of the operator  $C$  and  $J^- J^3 - (m + (\gamma + 1)/-2) J^- J^3$ .

Let  $f(x, y)$  satisfy the simultaneous equations

$$\begin{aligned}
(3.13) \quad & Cf(x, y) = \left(\frac{\gamma^2-1}{4}\right) f(x, y), \\
& \left[J^- J^3 - \left(m + \frac{\gamma+1}{2}\right) J^- - J^3\right] f(x, y) = 0.
\end{aligned}$$

which may be rewritten as

$$(3.14) \quad \begin{aligned} & \left[ x \frac{\partial^2}{\partial x^2} + (1 + \gamma - x) \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + m \right] f(x, y) = 0, \\ & \left[ -y \frac{\partial^2}{\partial y^2} + x \frac{\partial^2}{\partial x \partial y} - (1 + m + y) \frac{\partial}{\partial y} - \left( m + \frac{\gamma + 1}{2} \right) \right] f(x, y) = 0 \end{aligned}$$

These equations have a solution [1; 235]

$$(3.15) \quad f(x, y) = H_4 \left[ -m, m + \frac{\gamma + 1}{2}; 1 + \gamma; x, y \right]$$

where

$$H_4[\alpha, \gamma; \delta; x, y] = \sum_{m, n=0}^{\infty} \frac{(\alpha)_{m-n} (\gamma)_n}{m! n! (\delta)_m} x^m y^n.$$

Therefore

$$(3.16) \quad \begin{aligned} [T(g)f](x, y) &= \exp\left(\frac{bxy}{d+by}\right) (d-by)^{-(\gamma+1+m)} \left(a + \frac{c}{y}\right)^m \\ &\times H_4 \left[ -m, m + \frac{\gamma + 1}{2}; 1 + \gamma; \frac{xy}{(c+ay)(d+by)}, \frac{c+ay}{d+by} \right] \\ &\left| \frac{by}{d} \right| < 1, \quad -\pi < \arg a, \arg d < \pi, \quad ad - bc = 1. \end{aligned}$$

By the similar analysis as we did in sec. A. we obtain the generating function

$$(3.17) \quad \begin{aligned} & e^{(bxy/d+by)} \left(1 + \frac{by}{d}\right)^{-(\gamma+1+m)} \left(1 + \frac{c}{ay}\right)^m \\ & \times H_4 \left[ -m, m + \frac{\gamma + 1}{2}; 1 + \gamma; \frac{xy}{(c+ay)(d+by)}, \frac{c+ay}{d+by} \right] \\ & = \sum_{n=-\infty}^{\infty} \frac{\Gamma(1+m+n)}{(1+\gamma)_m} \left(\frac{-by}{d}\right)^n L_{m+n}^{(\gamma)}(x) \\ & \times \sum_{p=0}^{\infty} \frac{(-m)_p (1+m+n+\gamma)_p}{p! \Gamma(1+n+p)} \left(\frac{bc}{ad}\right)^p \\ & \times {}_2F_2 \left[ \begin{matrix} -n-p, m + \frac{\gamma + 1}{2} \\ m+1-p, 1+\gamma+m \end{matrix}; -\frac{a}{b} \right], \\ & \left| \frac{c}{ay} \right| < 1, \quad \left| \frac{by}{d} \right| < 1, \quad ad - bc = 1, \end{aligned}$$

where the terms corresponding to  $n = -1, -2, \dots$  are well defined because of the relation (3.10).

*Special Cases:*

$$(3.18) \quad e^{(-xy/1-y)}(1-y)^{-(1+\gamma+m)} H_4 \left[ -m, m + \frac{\gamma+1}{2}; 1+\gamma; \frac{x}{1-y}, \frac{-wy}{1-y} \right]$$

$$= \sum_{n=-\infty}^{\infty} \frac{(1+n)_m}{(1+\gamma)_m} L_{m+n}^{(\gamma)}(x) {}_2F_2 \left[ \begin{matrix} -n, m + \frac{\gamma+1}{2} \\ m+1, 1+m+\gamma \end{matrix}; w \right] y^n$$

where the terms corresponding to  $n=1, -2, \dots$  are well defined because of the relation of type 3.10.;

$$(3.19) \quad \left(1 - \frac{w}{y}\right)^m H_4 \left[ -m, m + \frac{\gamma+1}{2}; 1+\gamma; \frac{xy}{w-y}, w-y \right] e^w$$

$$= \sum_{n=0}^{\infty} \frac{\Gamma(1+m)}{(1+\gamma)_m} L_{m+n}^{(\gamma)}(x) L_{m+(r-1/2)}^{(n)}(w) y^n$$

C. We choose  $f(x, y)$  to be a common eigenfunction of the operator C and  $(-J^-J^3 + ((m+\gamma+1)/2)J^- + J^3)$ .

Let  $f(x, y)$  satisfy the simultaneous equations

$$(3.20) \quad Cf(x, y) = \left(\frac{\gamma^2-1}{4}\right)f(x, y),$$

$$\left(-J^-J^3 + \left(\frac{m+\gamma+1}{2}\right)J^- + J^3\right)f(x, y) = \left(\frac{m+\gamma-1}{2}\right)f(x, y).$$

This may be rewritten as

$$(3.21) \quad \left[x \frac{\partial^2}{\partial x^2} + (1+\gamma-x) \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + m\right]f(x, y) = 0,$$

$$\left[y \frac{\partial^2}{\partial y^2} - x \frac{\partial^2}{\partial x \partial y} + (1+m+y) \frac{\partial}{\partial y} + 1\right]f(x, y) = 0.$$

These equations have a solution [1; 235]

$$(3.22) \quad f(x, y) = H_5[-m; 1+\gamma; x, y]$$

where

$$H_5[\alpha; \delta; x, y] = \sum_{m,n=0}^{\infty} \frac{(\alpha)_{m-n}}{m! n! (\delta)_m} x^m y^n.$$

The generating function that we obtain now is

$$(3.23) \quad \exp\left(\frac{bxy}{d+by}\right) \left(1 + \frac{by}{d}\right)^{-(r+1+m)} \left(1 + \frac{c}{ay}\right)^m$$

$$\times H_5 \left[ -m; 1+\gamma; \frac{xy}{(c+ay)(d+by)}, \frac{c+ay}{(d+by)} \right]$$

$$\begin{aligned}
 &= \sum_{n=-\infty}^{\infty} \frac{\Gamma(1+m+n)}{(1+\gamma)_{m+n}} \left(\frac{-by}{d}\right)^n L_{m+n}^{(\gamma)}(x) \sum_{p=0}^{\infty} \frac{(-m)_p}{p!} \left(\frac{bc}{ad}\right)^p \\
 &\quad \times L_{n+p}^{(\gamma+m)}\left(\frac{-a}{b}\right).
 \end{aligned}$$

where the terms corresponding to  $n = -1, -2, \dots$  are well defined because of the relation of the type (3.10).

*Special Cases:*

$$\begin{aligned}
 (3.24) \quad &\exp\left(\frac{-xy}{1-y}\right)(1-y)^{-(r+1+m)} H_5\left[-m; 1+\gamma; \frac{x}{1-y}, \frac{-wy}{1-y}\right] \\
 &= \sum_{p=0}^{\infty} \frac{\Gamma(1+m+n)}{(1+\gamma)_{m+n}} L_{m+n}^{(\gamma)}(x) L_n^{(\gamma+m)}(w) y^n
 \end{aligned}$$

and

$$\begin{aligned}
 (3.25) \quad &\left(1-\frac{w}{y}\right)^m H_5\left[-m; 1+\gamma; \frac{-xy}{w-y}, w-y\right] \\
 &= \sum_{n=-\infty}^{\infty} \frac{\Gamma(1+m)}{\Gamma(1+n)(1+\gamma)_m} L_{m+n}^{(\gamma)}(x) {}_0F_1[-m; 1+n; -w] y^n.
 \end{aligned}$$

**D.** Choose  $f(x, y)$  to be a common eigenfunction of the operator  $C$  and  $[J^3 J^3 - J^- J^3 + (m + (\gamma + 1)/2) J^- - (m + 1) J^3]$ .

Let  $f(x, y)$  satisfy the simultaneous equations

$$\begin{aligned}
 (3.26) \quad &Cf(x, y) = \left(\frac{\gamma^2 - 1}{4}\right) f(x, y), \\
 &\left[J^3 J^3 - J^- J^3 + \left(m + \frac{\gamma + 1}{2}\right) J^- - (m + 1) J^3\right] f(x, y) = 0
 \end{aligned}$$

which may be rewritten as

$$\begin{aligned}
 (3.27) \quad &\left[x \frac{\partial^2}{\partial x^2} + (1 + \gamma - x) \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} + m\right] f(x, y) = 0, \\
 &\left[y(y+1) \frac{\partial^2}{\partial y^2} - x \frac{\partial^2}{\partial x \partial y} + \{1 + m + (m + \gamma + 1)y\} \frac{\partial}{\partial y} \right. \\
 &\quad \left. + \left(\frac{\gamma - 1}{2}\right) \left(m + \frac{\gamma + 1}{2}\right)\right] f(x, y) = 0.
 \end{aligned}$$

These equations have a solution [1; 235]

$$(3.28) \quad f(x, y) = H_{11}\left[-m, \frac{\gamma - 1}{2}, m + \frac{\gamma + 1}{2}; 1 + \gamma; x, y\right]$$

where

$$H_{11}(\alpha, \beta, \gamma; \delta; x, y) = \sum_{m, n=0}^{\infty} \frac{(\alpha)_{m-n} (\beta)_n (\gamma)_n}{m! n! (\delta)_m} x^m y^n \quad |y| < 1.$$

Thus, we arrive at a generating function

$$\begin{aligned}
 (3.29) \quad & \exp\left(\frac{bxy}{d+by}\right)\left(1+\frac{by}{d}\right)^{-(\gamma+1+m)}\left(1+\frac{c}{ay}\right)^m \\
 & \times H_{11}\left[-m, \frac{\gamma-1}{2}, m+\frac{\gamma+1}{2}; \gamma+1; \frac{xy}{(c+ay)(d+by)}, \frac{c+ay}{d+by}\right] \\
 & = \sum_{n=-\infty}^{\infty} \frac{\Gamma(1+m+n)}{(1+\gamma)_m} \left(\frac{-by}{d}\right)^n L_{m+n}^{(\gamma)}(x) \\
 & \times \sum_{p=0}^{\infty} \frac{(-m)_p (1+\gamma+m+n)_p}{p! \Gamma(1+n+p)} \left(\frac{bc}{ad}\right)^p \\
 & \times {}_3F_2\left[\frac{\gamma-1}{2}, m+\frac{\gamma+1}{2}, -n-p; \frac{a}{b}; \frac{m+1-p, 1+\gamma+m}{m+1-p, 1+\gamma+m}\right], \\
 & \left|\frac{c}{ay}\right| < 1, \left|\frac{by}{d}\right| < 1, \quad ad-bc=1.
 \end{aligned}$$

where the terms corresponding to  $n=-1, -2, \dots$  are well defined because of the relation of the type 3.10.

*Special Cases:*

$$\begin{aligned}
 (3.30) \quad & \exp\left(\frac{xy}{1-y}\right)(1-y)^{-(\gamma+m+1)} \\
 & \times H_{11}\left[-m, \frac{\gamma-1}{2}, m+\frac{\gamma+1}{2}; 1+\gamma; \frac{x}{1-y}, \frac{wy}{1-y}\right] \\
 & = \sum_{n=-\infty}^{\infty} \frac{(1+n)_m}{(1+\delta)_m} L_{m+n}^{(\gamma)}(x) {}_3F_2\left[\frac{\gamma-1}{2}, m+\frac{\gamma+1}{2}, -n; w; \frac{m+1, 1+m+\gamma}{m+1, 1+m+\gamma}\right] (y)^n.
 \end{aligned}$$

and

$$\begin{aligned}
 (3.31) \quad & \left(1-\frac{w}{y}\right)^m H_{11}\left[-m, \frac{\gamma-1}{2}, m+\frac{\gamma+1}{2}; 1+\gamma; \frac{-xy}{w-y}, w-y\right] \\
 & = \sum_{n=-\infty}^{\infty} \frac{\Gamma(1+m)\left(\frac{\gamma-1}{2}\right)_n \left(m+\frac{\gamma+1}{2}\right)_n}{\Gamma(1+n)(1+\gamma)_m} L_{m+n}^{(\gamma)}(x) \\
 & \times {}_2F_1\left[m+n+\frac{\gamma-1}{2}, +n\frac{\gamma+1}{2}; -w; y^n\right]
 \end{aligned}$$

where the terms corresponding to  $n=-1, -2, \dots$  are well defined by a relation of the type (3.10).



## References

- [1] ERDELYI, A.; *Higher Transcendental Functions* (Bateman Manuscript Project) Vol. 1  
McGraw Hill New York 1953.
- [2] MILLER, W.; *Lie Theory and Special Functions* Academic Press New York 1968.
- [3] RAINVILLE, E. D.; *Special Functions* (MacMillan New York 1960).