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# Linear Driving Force Model for Adsorption onto Activated Carbon Monolith

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**Abstract**. A mathematical model on carbon coated monolith using the linear driving force model is develop. The computer program is writted in MATLAB and simulation using data from [4] for cell density 200 cpsi was used to studies the effect of different variables on breakthrough profiles. The result showed that the breakthrough curve of the monolith is very sharp. Because of its an open structure and lower pressure drop, monolith is an attractive alternative internals for separation. Copyright © 2006 Teknik Kimia UNSYIAH

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# **INTRODUCTION**

Porous carbon materials or activated carbons have been in use for many years as solid adsorbent in fields of adsorption because of their large adsorptive capacity and low cost. These activated carbons can be produced from natural source or polymer. Different types of activated carbon provide different characteristics and capacities. Activated carbons from natural source have non-uniform pore size distribution. In order to produce activated carbons which provide controlled pore size distribution, carbon derived from polymer is usually used. Applied forms of activated carbon used commonly are in the shape of powder or granule.

The development of controlled porous carbon derived from polymeric materials such as poly(furfuryl alcohol), furfuryl alcohol, furan, etc. is targeted to produce

carbon with the desired pore structure. Poly (ethylene glycol) as pore former added to poly(furfuryl alcohol) prior to pyrolysis was to create mesoporosity activated carbon (Lafiatis at al., 1991). The application of activated carbon products, in general, is implemented in packed bed column. activated carbon applied in packed bed has disadvantages, *i.e.*, high pressure drop associated with the flow through packed bed media, channeling, hot spot, and run away behavior for extremely exothermic reaction. Another problem associated with packed bed is the entrainment of particles in the flow stream. The monolithic column is an attractive alternative system to eliminate the problems mentioned previously for packed column.

The activated carbon monolith structure is known to have very high geometric surface area to volume ratio, and has very open structure. It can be located in a vertical and horizontal position or in mobile system without losing shape and is easier to be scaled up due to its simple design and uniform flow distribution [1]. Other advantages are the accessible surface area of the activated carbon monolith is larger than that of the packed column. Monolithic catalyst structures have been widely applied in single-phase flow such as automotive emission control, selective catalytic reduction of NOx [3-5], adsorption of gas [2-6], and selective oxidation. More recently, the opportunities in multiphase catalytic processes have been recognized, leading to increase research interest in this area. These processes include hydrogenation, oxidation, and decomposition reactions [3,4,7]. The production of H<sub>2</sub>O<sub>2</sub> [8] is an industrial application example.

The objective of this work is to develop an linear driving force adsorption model to simulate the effect of the total lengths of the monoliths and velocities of feeds on the breakthrough performance for linear isotherm.

## Modeling

The mathematical model on monolithic column is developed on the basis of one single square channel with a flat carbon coating on the wall. The sketch of monolithic structure is illustrated in Figure 1, where  $R_1$  and  $R_2$  are the radius of channel's wall and the radius of the channel, respectively.

The following assumptions are applied in the mathematical model development:

- The monolith channel is a square by applying diameter equivalent with inside radius R<sub>2</sub>
- No variation of the mass transfer parameters and physicochemical properties along the monolithic column.
- Instantaneous equilibrium on the wall pore surface.
- Isothermal operation.



Fig. 1. The sketch of monolithic structure

In this work, an axially dispersed plug flow model through the channel is developed. Adsorption equilibrium is included in the model by using linear isotherm.

# Mass balance in bulk phase in monolith column

Mass balance in the monolithic channel  $(t > 0; 0 \ge r \le R_2)$  can be calculated as

$$\frac{\partial C}{\partial t} = -V_z \frac{\partial C}{\partial z} + D_{ax} \frac{\partial^2 C}{\partial z^2} - \frac{(1 - \varepsilon_b)}{\varepsilon_b} \frac{1}{K} \frac{\partial \overline{C}_p}{\partial t} \quad (1)$$

where *C* is the liquid bulk phase concentration,  $\overline{C}_p$  is average pore concentration in the skeleton wall, *z* is the axial coordinate, *t* is time, and  $D_{ax}$  is the axial dispersion coefficient [4, 10, 11]. The initial and boundary conditions for Eq. (1) are as following:

Initial condition

$$t = 0; \qquad C = 0 \qquad \text{for all } z \qquad (2)$$

Boundary conditions

$$z = 0; \qquad C = C_F \qquad t > 0 \tag{3}$$

$$z = L; \qquad \frac{dC}{dz} = 0 \qquad t > 0 \tag{4}$$

#### Mass balance in channel wall

The mass balance in channel wall is given as

$$D_{p}\frac{\partial}{r\partial r}\left(r\frac{\partial C_{p}}{\partial r}\right) = \varepsilon_{p}\frac{\partial C_{p}}{\partial t} + \frac{\partial \overline{q}}{\partial t}$$
(5)

for t > 0;  $R_2 \ge r \le R_1$ 

The initial and Boundary conditions are: Initial condition

$$t = 0; \qquad C_p = 0 \quad \text{for all } r \tag{6}$$

Boundary conditions

$$r = R_2; \qquad C_p \Big|_{R_2} = C \Big|_{R_2} \tag{7}$$

$$r = R_1; \qquad \frac{\partial C_p}{\partial r}\Big|_{R_1} = 0 \tag{8}$$

where  $C_p$  is the liquid concentration in the wall pores and  $\overline{q}$  is the average wall adsorbed phase concentration.  $D_p$  is the wall pore diffusivity. If transport within the macro-pores occurs only by molecular diffusion, the  $D_p$  is given by

$$D_p = \frac{D_m}{\tau} \tag{9}$$

where  $D_m$  is the molecular diffusivity and  $\tau$  is the wall pore tortuosity. For straight, randomly oriented, cylindrical pores,  $\tau$  can be taken as 3 [12].

The adsorption equilibrium on the wall pore surface is described by linear adsorption isotherm:

$$q = KC_p \tag{10}$$

Averaging the pore concentration over the monolith wall and assuming the parabolic pore concentration profile through the wall, Eq. (5) is transformed in linear driving force (LDF) equation form:

$$(\varepsilon_p + K)\frac{\partial \overline{C}_p}{\partial t} = k_{LDF}(C|_{R2} - \overline{C}_p) \quad (11)$$

where  $k_{LDF}$  is the internal mass transfer coefficient:

$$k_{LDF} = \frac{8D_p R_2^2}{(R_1^2 - R_2^2)^2}$$
(12)

## Numerical simulation

The model set of partial differential equations (PDEs) was solved by using MATLAB 6.5. The PDEs were discretised in the spatial domain with the method of orthogonal collocation (OC). Detailed information on the method of OC can be found in [13-15]. The resulting set of ordinary differential equations (ODEs) was solved by using MATLAB subroutine ODE15s.

#### Validation of the discretisation method

The OC method is first used to simulate a case of heat conduction for which analytical equation can be obtained. It uses a polynomial of order 50. The equation for the transient heat conduction is given as:

$$\frac{\partial T}{\partial t} = \alpha_h \frac{\partial^2 T}{\partial z^2} \tag{13}$$

where  $\alpha_h$  is the thermal diffusivity (m<sup>2</sup> s<sup>-1</sup>), *T* is the temperature inside the slab (°C), *t* is time (s) and z is the axial co-ordinate (m).

# Heat conduction in a slab of thickness 2L

A slab with thickness  $2L = 3.0 \times 10^{-2}$  m,  $\alpha_h = 5.0 \times 10^{-6}$  m<sup>2</sup> s<sup>-1</sup> and initial temperature  $T_{initial} = 1000$  °C is suddenly immersed into an ice bath of 0 °C. The temperature profile inside the slab in this case is symmetric. The boundary conditions are

$$T\Big|_{z=0} = 0;$$
  $\frac{\partial T}{\partial z}\Big|_{z=L} = 0$  (14)

The analytical solution is given as [9]:

$$T = \frac{4T_{initial}}{\pi} \sum_{n=0}^{\infty} \frac{1}{(2n+1)} \exp\left(\frac{-\alpha_h (2n+1)^2 \pi^2 t}{(2L)^2}\right) \sin\frac{(2n+1)\pi z}{2L}$$
(15)

Figure 2 shows T as a function of dimensionless length (z/L), calculated using OC. Excellent agreement is obtained.



Fig. 2. Comparison of the numerical calculations and analytical solution of heat conduction in a slab of thickness 2L

#### Validation of column model

The simulation programs are validated by simplifying a case of packed bed adsorber to compare the numerical calculation with analytical solution. The equations describing the packed bed adsorber are:

$$V_0 \frac{\partial C}{\partial z} + \varepsilon_b \frac{\partial C}{\partial t} + (1 - \varepsilon_b) \frac{\partial q}{\partial t} = 0 \quad (16)$$

where

$$\frac{\partial \overline{q}}{\partial t} = k_{LDF} \left( q^* - \overline{q} \right) \tag{17}$$

 $q^*$  is the adsorption equilibrium on solid phase. It is equivalent to q in equation (10).

Initial conditions and boundary conditions are

$$C(0,\xi) = 0 \tag{18}$$

$$C_p(0,\xi) = 0$$
 (19)

$$C\Big|_{\xi=0} = C_F \tag{20}$$

$$\frac{\partial C}{\partial \xi}\Big|_{\xi=1} = 0 \tag{21}$$

The analytical solution for equation (16) - (18) are given by Rice and Do (1995) as

$$\frac{C}{C_0} = J(\zeta, \eta) \tag{22}$$

where  $J(\varsigma, \eta)$  is defined as

$$J(\zeta,\eta) = 1 - e^{-\tau} \int_0^{\zeta} e^{-t} I_0 \left( 2\sqrt{\eta \cdot t} \right) dt \qquad (23)$$

where 
$$\zeta = \left(\frac{k_c a}{\varepsilon_b}\right) \cdot \frac{z}{V_z}$$
 (24)

$$\eta = \left[\frac{k_c a}{K(1 - \varepsilon_b)}\right] \cdot \left(t - \frac{z}{V_z}\right)$$
(25)

where  $k_c a$  is the mass transfer constant  $(s^{-1})$ ,  $\varepsilon_b$  is bed porosity, z is bed length (m), and t is time (s), and  $V_z$  is interstitial velocity (ms<sup>-1</sup>). Input parameters for validation of column are given in Table 1. Time as function dimensionless concentration  $(C/C_0)$  is shown in Figure 3. Excellent agreement between the simulation and analytical solution is obtained as well.

**Table 1.** Input parameter for validation ofpacked bed adsorber model [14].

Input Parameter	Value	Unit
Column height, L	1.0	т
Bulk porosity, $\mathcal{E}_b$	0.4	-
Linear adsorption constant, K	2.0	-
Mass transfer constant, $k_c a$	0.1	$s^{-1}$
Superficial velocity, V <sub>0</sub>	0.04	$m \cdot s^{-1}$
Linear driving force constant, $k_{LDF} = [k_c a / K(1 - \varepsilon_b)]$	0.083	s <sup>-1</sup>
Interstitial velocity, $V_{z} = V_0 / \varepsilon_b$	0.1	$m \cdot s^{-1}$



**Fig. 3**. Comparison of the numerical calculations using *OC* method and analytical solution of a packed bed adsorber

#### **RESULT AND DISCUSSION**

Figure 4 presents the simulation result of the breakthrough profiles for plug flow velocity profile obtained using data from [4] with 200 cpsi monoliths having total length of 5 cm, 10 cm, 15 cm are plotted. The longer the total length of the monolith, the higher a displacement of breakthrough time value. The adsorption performance however remains identical. It can also be explained that short diffusion length occurs in the additional monolith. An advantage of carbon coated activated on monolith compare to carbon packed beds is the low pressure drop that monoliths produce in the system [3]. The pressure drop associated with the monolithic system was estimated by using the Hagen-Poiseulle equation. The value in this operation conditions is 2.8 Pa/m [3], while the pressure drop for a carbon packed bed calculated by the dimensionless friction factor using Chilton-Colburn correlations [12] gave a value of around 150 Pa/m.

The breakthrough curves for different velocity 1, 3, and 5 cm/s are presented in Figure 5. The slope of breakthrough curve for the velocity 1 cm/s decreases gradually. It is clear that there is no constant pattern behavior for the monolithic column. The effect of velocity which contributes to axial mixing is lumped together into the axial dispersion coefficient.



**Fig. 4**. Breakthrough profiles for monolith having 200 *cpsi* at *v* 1 cm/s ( $\diamond$ =5 cm;  $\Box$ =10 cm;  $\Delta$ =15 cm)



**Fig. 5**. Breakthrough profiles for monolith having 200 cpsi at length 10 cm and different velocities ( $\Delta$ =1 cm/s;  $\Diamond$ =3 cm/s;  $\Box$ =5 cm/s).

#### **CONCLUSION**

The breakthrough curves shown by monoliths are very sharp and it gives extremely low pressure drop. It also shows that dynamic adsorptive capacity behavior of monolith column was very good. Monoliths are a promising alternative for the adsorption system. The width of the breakthrough profile scales with respect to the length of monolith takes place, if a laminar flow is over the channel. This case can be reduced by increasing the cell density.

# **NOMENCLATURE**

- С bulk liquid phase concentration  $(kg/m^3)$
- liquid concentration in the skeleton  $C_{n}$ pores  $(kg/m^3)$
- $\overline{C}_{n}$ average concentration in the skeleton pores  $(kg/m^3)$
- initial (feed) concentration  $(kg/m^3)$  $C_{F}$
- axial dispersion coefficient  $(m^2/s)$  $D_{ax}$
- molecular diffusion  $(m^2/s)$  $D_m$
- pore diffusion  $(m^2/s)$  $D_{p}$
- linear driving force mass transfer  $k_{LDF}$ coefficient (1/s)
- K linear adsorption constants
- mass transfer constant  $(s^{-1})$ k<sub>c</sub>a
- length of the monolithic column (m) L
- equilibrium adsorbed phase qconcentration in the monolithic skeleton (kg/m<sup>3</sup>)
- equilibrium  $q^*$ adsorbed phase concentration in packed bed  $(kg/m^3)$ radial variable (m) r
- radius of the channel's skeleton  $R_1$ wall (m)
- radius of the channel (m)  $R_2$
- time variable (s) t
- Т temperature (K)
- superficial velocity (ms<sup>-1</sup>)  $V_0$
- interstitial velocity (ms<sup>-1</sup>)  $V_{z}$
- axial variable (m) Ζ.

### Greek letters

- thermal diffusivity  $(m^2 s^{-1})$  $\alpha_h$
- bulk porosity  $\mathcal{E}_{b}$
- internal porosity  $rac{arepsilon}{ au}_p$
- skeleton pore tortuosity

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