

# Model for Cameron band emission in comets: A case for EPOXI mission target comet 103P/Hartley 2

Anil Bhardwaj <sup>\*</sup> and Susarla Raghuram <sup>†</sup>

*Space Physics Laboratory, Vikram Sarabhai Space Center, Trivandrum, 695022, India*

## ABSTRACT

The CO<sub>2</sub> production rate has been derived in comets using the Cameron band ( $a^3\Pi - X^1\Sigma$ ) emission of CO molecule assuming that photodissociative excitation of CO<sub>2</sub> is the main production mechanism of CO in  $a^3\Pi$  metastable state. We have developed a model for the production and loss of CO( $a^3\Pi$ ) which has been applied to comet 103P/Hartley 2: the target of EPOXI mission. Our model calculations show that photoelectron impact excitation of CO and dissociative excitation of CO<sub>2</sub> can together contribute about 60-90% to the Cameron band emission. The modeled brightness of (0-0) Cameron band emission on comet Hartley 2 is consistent with Hubble Space Telescope observations for 3-5% CO<sub>2</sub> (depending on model input solar flux) and 0.5% CO relative to water, where photoelectron impact contribution is about 50-75%. We suggest that estimation of CO<sub>2</sub> abundances on comets using Cameron band emission may be reconsidered. We predict the height integrated column brightness of Cameron band of  $\sim 1300$  R during EPOXI mission encounter period.

**Key words:** comets:general-comets:103P/Hartley 2—molecular processes-astrochemistry

## 1 INTRODUCTION

In the exploration of the solar system, comets have been targeted by various space missions. After successful encounter of comet 9P/Tempel 1 on 4 July 2005, the NASA's Deep impact mission, also called EPOXI mission, under its extended investigation program DIXI will encounter comet 103P/Hartley 2 on 4 November 2010, with closest approach around 700 km from the nucleus. This comet has been observed by several space telescopes in different spectral regions (Weaver et al. 1994; Crovisier et al. 1999; Colangeli et al. 1999; Groussin et al. 2004; Lisse et al. 2009; Snodgrass et al. 2008, 2010).

The first clear detection of the Cameron band ( $a^3\Pi - X^1\Sigma$ ) of CO was reported by Weaver et al. (1994) in HST/FOS spectrum of comet 103 P/Hartley 2. Since Cameron band emission is a forbidden transition, resonance fluorescence is not an effective excitation mechanism. The upper state of this emission ( $a^3\Pi$ ) is a metastable state with lifetime of 3 ms (Gilijamse et al. 2007), which is quite small. Thus, Cameron band emissions are treated as “prompt emissions” and can be used to probe distribution of parent species if this emission is produced in dissociative excitation of a molecular species. Photodissociative excitation of CO<sub>2</sub> is considered as the major production mechanism of CO Cameron band and has been used to trace the distribution and abundance of CO<sub>2</sub> on comets (Feldman et al. 1997; Weaver et al. 1994).

Besides photodissociative excitation of CO<sub>2</sub>, there are other

channels of excitation of CO molecule in the  $a^3\Pi$  state. It has been shown that photoelectrons generated by solar EUV radiation also play an important role in excitation, dissociation, and ionization processes leading to emission and chemistry in cometary comae (e.g., Ip 1986; Boice et al. 1986; Körösmezey et al. 1987; Bhardwaj et al. 1990, 1996; Bhardwaj 1999, 2003; Weaver et al. 1994; Haider & Bhardwaj 2005; Capria et al. 2008). Recently, Campbell & Brunger (2009) demonstrated the importance of photoelectron impact excitation in comets, and showed that electron impact on CO gives 40% contribution to the total CO Fourth positive emission. The presence of photoelectron excitation in cometary coma is clearly publicised by detection of OI 1356 Å emission in comets (e.g., Sahnou et al. 1993; McPhate et al. 1999), since this emission being a spin-forbidden transition cannot be produced by solar fluorescence. There are other significant evidences for an important role of photoelectron excitation in cometary coma (e.g., Tozzi et al. 1998; Bhardwaj 1999; Feldman et al. 2009). In addition to photon and electron impact reactions, dissociative electron recombination reactions of CO<sup>+</sup>-bearing ions can also produce CO in the  $a^3\Pi$  excited state.

Our aim in this paper is to study various production and loss mechanisms of CO( $a^3\Pi$ ) and to estimate the contribution of photoelectron impact excitation of CO and CO<sub>2</sub> in the production of Cameron band for different relative abundances of CO<sub>2</sub> on comet 103 P/Hartley 2: the target of EPOXI mission. Since model calculations depend on input solar flux, we have estimated its sensitivity on the calculated intensity of Cameron band emission. We show that photoelectron impact on CO and CO<sub>2</sub> are dominant processes ( $\sim 60$ – $90$ % contribution) in producing CO molecule in ( $a^3\Pi$ ) state.

<sup>\*</sup> E-mail: anil.bhardwaj@vssc.gov.in; bhardwaj\_spl@yahoo.com

<sup>†</sup> E-mail: raghuramsusarla@gmail.com

## 2 Anil Bhardwaj and Susarla Raghuram

Around the EPOXI encounter epoch predictions are made for the brightness of Cameron band on for comet 103P/Hartley 2, which will be observed by several space-based telescopes including HST.

### 2 MODEL

We have developed a model for the production of Cameron band emission on comets, which uses the basic coupled chemistry model described in detail in our earlier papers (Bhardwaj et al. 1996; Bhardwaj 1999; Haider & Bhardwaj 2005). Various sources of production and loss of CO( $a^3\Pi$ ) are summarized in Table 1. The total water production rate is taken as  $6.3 \times 10^{28} \text{ s}^{-1}$  for comet 103P/Hartley 2 (Weaver et al. 1994).

To evaluate the effect of solar EUV flux on model calculations, we have considered 2 solar flux models: EUVAC model of Richards et al. (1994) and SOLAR 2000 (S2K) model of Tobiska et al. (2000). The degradation of the solar UV-EUV radiation and solar EUV-generated photoelectrons in the coma is modeled using the method of Bhardwaj et al. (1990, 1996) and further developed by Bhardwaj (1999, 2003). The electron impact production rates are calculated using the Analytical Yield Spectrum (AYS) approach, which is based on the Monte Carlo method. Details of AYS approach are given in several of the previous papers (Bhardwaj et al. 1990, 1996; Bhardwaj 1999; Bhardwaj & Michael 1999a,b; Bhardwaj & Jain 2009). The current model takes into account the most recently published cross sections for the photon impact and electron impact dissociation, ionization, and excitation processes for the gases in the coma. The cross section for photodissociative excitation of CO<sub>2</sub> producing CO in  $a^3\Pi$  state is calculated using absorption cross sections of CO<sub>2</sub> and yield of Cameron band measured by Lawrence (1972). The cross section for electron impact excitation of CO( $a^3\Pi$ ) from CO is taken from Jackman et al. (1977) and for dissociative excitation of CO<sub>2</sub> is taken from Bhardwaj & Jain (2009). The electron temperature profile required for dissociative recombination reactions is taken from Körösmezey et al. (1987) and is assumed to be same as on comet Halley. Calculations are made for the comet 103P/Hartley 2 at heliocentric distance of 0.96 AU.

### 3 RESULTS AND DISCUSSION

The photodissociation of CO<sub>2</sub> producing CO in  $a^3\Pi$  state is determined by solar flux mainly in the wavelength region 550 to 1050 Å. Table 2 presents the calculated photon production frequencies of CO( $a^3\Pi$ ) for two different solar flux models. The CO( $a^3\Pi$ ) production frequencies calculated for photoelectron impact on CO<sub>2</sub> and CO are also shown in the same table for the corresponding solar flux models. Our calculated photodissociation frequencies are about 50% to a factor of 2 lower than those reported by Huebner et al. (1992).

Using EUVAC solar flux, the calculated radial profile of volume production rate for the various sources of CO( $a^3\Pi$ ) at the relative abundance of 4% CO<sub>2</sub> and 0.5% CO are shown in Figure 1. At 100 km cometocentric distance, the dominant source of production of CO( $a^3\Pi$ ) is electron impact of CO<sub>2</sub> (~50%) followed by electron impact of CO (~25%), and photodissociation of CO<sub>2</sub> (~15%). The contributions from dissociative recombination reactions are quite small ( $\leq 5\%$ ) at lower cometocentric distances, but the recombination of CO<sub>2</sub><sup>+</sup> is a significant (<30%) source at 1000

km and beyond. Figure 2 shows radial profile of various loss processes of CO( $a^3\Pi$ ) for the same relative composition of CO<sub>2</sub> and CO. Since lifetime of CO in excited state ( $a^3\Pi$ ) is very short (~3 ms; Gilijamse et al. 2007), the radiative decay is the dominant loss process. Collisional quenching of CO( $a^3\Pi$ ) by cometary neutral species is negligible since 103P/Hartley 2 is a low production rate comet. But in the case of large production rate comets, like Hale-Bopp, quenching by water would be a dominant loss process in the innermost part of the coma.

Figure 3 shows the modeled limb brightness profiles of Cameron band emission for different production processes of CO( $a^3\Pi$ ). The cometary coma is assumed to be spherically symmetric. The production rates are integrated up to 10<sup>5</sup> km along the line of sight at a given projected distances from the cometary nucleus, and converted into brightness. The brightness profiles are averaged over the projected area of slit (2870 × 954 km) corresponding to the HST observation (Weaver et al. 1994). The volume emission rate for 3 transitions (0-0, 1-0, 0-1) of Cameron band emission are calculated using the following formula

$$V_{\nu'\nu''}(r) = q_{\nu\nu'}(A_{\nu'\nu''} / \sum_{\nu''} A_{\nu'\nu''}) V(r) \exp(-\tau) \quad (1)$$

where V(r) is total volume excitation rate of CO( $a^3\Pi$ ) at cometocentric distance r,  $q_{\nu\nu'}$  is the Franck-Condon factor for transition,  $A_{\nu'\nu''}$  is Einstein transition probability from upper state  $\nu'$  to lower state  $\nu''$ , and  $\tau$  is optical depth. Since resonance fluorescence is not an effective excitation mechanism for the Cameron band and the total gas production rate is only  $6.3 \times 10^{28} \text{ s}^{-1}$ , the cometary coma can be safely assumed to be optically thin. The Franck-Condon factors are taken from Nicholls (1962) and branching ratios from Conway (1981). The relative contributions of (1-0), (0-0), (0-1) transitions to the total Cameron band are 13.9%, 10.4%, and 14.7%, respectively.

Table 3 presents the model calculated slit-averaged brightness of (1-0), (0-0), (0-1) transitions of Cameron band, as well as total Cameron band brightness and height-integrated column brightness for different relative abundances of CO and CO<sub>2</sub> corresponding to the HST observation of comet 103P/Hartley 2 on September 18-19, 1991. Due to the absence of CO Fourth positive emission in this comet (Weaver et al. 1994), the abundance of CO is constrained to 0.5%. However, we do consider a case of 1% of CO to evaluate its implications on the results. This table also depicts fractional contribution of photodissociation of CO<sub>2</sub>, photoelectron impact of CO and CO<sub>2</sub>, and dissociative recombination of CO<sub>2</sub><sup>+</sup> to the total calculated brightness at 3 projected distances (10<sup>2</sup>, 10<sup>3</sup>, and 10<sup>4</sup> km) from the nucleus. Since the production rates of photodissociative excitation of CO<sub>2</sub>, and photoelectron impact of CO and CO<sub>2</sub>, are dependent on input solar flux model, results are presented for the EUVAC and S2K solar fluxes relevant to the date of comet observation which was in solar maximum condition.

The HST observation of 0-0 transition of Cameron band is 35 Rayleigh (Weaver et al. 1994), which is consistent with model calculated brightness for the relative abundance of 4 to 5% of CO<sub>2</sub> and 0.5% CO when EUVAC solar flux is used. In this case, at 100 km, the photoelectron impact of CO<sub>2</sub> (50%) and CO (25%) contribute around 75%, while photodissociative excitation of CO<sub>2</sub> is <15%. At 1000 km and beyond, the contribution due to electron impact of CO<sub>2</sub> and CO is about 60–70% while those of dissociative recombination of CO<sub>2</sub><sup>+</sup> is ~15–30% and of photodissociative excitation of CO<sub>2</sub> ~10% only. On an average, the photoelectron impact of CO<sub>2</sub> and CO contribute about 60–75% to the production of Cameron band emission, while photodissociative excitation of

CO<sub>2</sub> contribute about 10-15% only. In the case of S2K solar flux model, the CO<sub>2</sub> abundance of 3 to 4% is required to match HST-observed Cameron band 0-0 transition brightness. Here the contribution of photodissociative excitation of CO<sub>2</sub> is ~20%, while the electron impact of CO<sub>2</sub> and CO together contribute ~65%, to the total Cameron band emission. When the CO abundance is doubled to 1% of water the relative contribution due to electron impact on CO increases resulting in the reduction (by ~1%) in the requirement for CO<sub>2</sub> abundance to match the HST-observation brightness. However, there is no major change in the percentage contributions due to photodissociation and photoelectron impact excitation of CO and CO<sub>2</sub>. Hence, we conclude that the photodissociative excitation of CO<sub>2</sub> is not the dominant source for the production of Cameron band in comets.

#### 4 SUMMARY

The summary of Table 3 and the above results is that whatever be the relative abundances of CO<sub>2</sub> and CO, and the solar input flux, the photodissociation contribution is in the range of 5–20%, while photoelectron impact excitation contribution is about 50–80%. We conclude that photoelectron impact excitation of Cameron band in comets is much more important than was thought previously. Further, the contribution due to photoelectron impact excitation of CO is significant and vary between 20-50%. We have also shown that dissociative recombination of CO<sub>2</sub><sup>+</sup> ion is also a significant source of Cameron band at cometocentric distances  $\geq 1000$  km with contribution of as high as 30%. Hence, the derivation of CO<sub>2</sub> abundances based on the assumption that photodissociation of CO<sub>2</sub> is major mechanism should be revisited. The results presented in this paper suggest that Cameron band emission may not be used to probe cometary CO<sub>2</sub> abundance uniquely.

#### 5 PREDICTION

The EPOXI mission will encounter the comet 103P/Hartley 2 on 4 November 2010 when comet will be at heliocentric distance of 1.06 AU and geocentric distance 0.15 AU. In the current solar minimum period, using our model we have estimated the intensity of Cameron band around mission encounter date by taking solar flux on 1 sept. 2010 for both EUVAC and S2K solar flux models. The calculated Cameron band intensity is tabulated in table 3 for the relative abundance of 4% of CO<sub>2</sub> and 0.5% of CO. For EUVAC solar flux, the predicted height-integrated column intensity of Cameron band is 1365 R, in which 1-0, 0-0, and 0-1 band transitions contribute 189, 142, 200 R, respectively. Figure 3 presents the total Cameron band predicted brightness profile as a function of projected distances from the nucleus. For the same relative abundances of CO<sub>2</sub> and CO and the solar S2K flux, the predicted intensity is 1258 R. The results presented in the paper will help in understanding and interpretation of the extensive data to be obtained during EPOXI encounter period.

#### REFERENCES

- Avakyan S. V., Avakian S. V., Il'in R. N., et al., eds, 1998, Collision Processes and Excitation of UV Emission from Planetary Atmospheric Gases. Gordon and Breach Science Publishers
- Bhardwaj A., 1999, *J. Geophys. Res.*, 104, 1929
- Bhardwaj A., 2003, *Geophys. Res. Lett.*, 30
- Bhardwaj A., Haider S. A., Singhal R. P., 1990, *Icarus*, 85, 216
- Bhardwaj A., Haider S. A., Singhal R. P., 1996, *Icarus*, 120, 412
- Bhardwaj A., Jain S. K., 2009, *J. Geophys. Res.*, 114
- Bhardwaj A., Michael M., 1999a, *J. Geophys. Res.*, 104, 713
- Bhardwaj A., Michael M., 1999b, *Geophys. Res. Lett.*, 26, 393
- Boice D. C., Huebner W. F., Keady J. J., et al., 1986, *Geophys. Res. Lett.*, 13, 381
- Campbell L., Brunger M. J., 2009, *Geophys. Res. Lett.*, 36
- Capria M. T., Cremonese G., Bhardwaj A., et al., 2008, *A&A.*, 479, 257
- Colangeli L., Epifani E., Brucato J. R., et al., 1999, *A&A.*, 343
- Conway R. R., 1981, *J. Geophys. Res.*, 86, 4767
- Crovisier J., Encrenaz T., Lellouch E., et al., 1999, Proceedings of the conference "The Universe as seen by ISO", 121, 161
- Feldman P. D., Festou M. C., Tozzi G. P., et al., 1997, *ApJ.*, 475, 829
- Feldman P. D., Lupu R. E., McCandliss S. R., Weaver H. A., 2009, *ApJ.*, 699
- Gilijamse J. J., Hoekstra S., Meek S. A., et al., 2007, *J. Chem. Phys.*, 127, 221102
- Groussin O., Lamy P., Jorda L., Toth I., 2004, *A&A.*, 413, 1163
- Haider S. A., Bhardwaj A., 2005, *Icarus*, 177, 196
- Huebner W. F., Keady J. J., Lyon S. P., 1992, *Ap&SS.*, 1, 195
- Ip W. H., 1986, *Adv. Space Res.*, 5, 47
- Jackman C. H., Garvey R. H., Green A. E. S., 1977, *J. Geophys. Res.*, 82, 5081
- Körösmezey A., Cravens T. E., Gombosi T. I., et al., 1987, *J. Geophys. Res.*, 92, 7331
- Lawrence G. M., 1972, *J. Chem. Phys.*, 56, 3435
- Lisse C., Fernandez Y., Reach W., et al., 2009, *PASP*, 121, 968
- McPhate J. B., Feldman P. D., McCandliss S. R., Burgh E. B., 1999, *Ap. J.*, 521, 920
- Nicholls R. W., 1962, *J. Quant. Spectr. Radiat. Transfer*, 2, 433
- Richards P. G., Fennelly J. A., Torr D. G., 1994, *J. Geophys. Res.*, 99, 8981
- Rosati R. E., Johnsen R., Golde M. F., 2003, *J. Chem. Phys.*, 119, 11630
- Rosati R. E., Skrzypkowski M. P., Johnsen R., Golde M. F., 2007, *J. Chem. Phys.*, 126, 154302
- Sahnow D. J., Feldman P. D., McCandliss S. R., Marinez M. E., 1993, *Icarus*, 101, 71
- Schmidt H. U., Wegmann R., Huebner W. F., Boice D. C., 1988, *Comp. Phy. Comm.*, 49, 17
- Seiersen K., Al-Khalili A., Heber O., et al., 2003, *Phys. Rev. A*, 68, 022708
- Shunk R. W., Nagy A. F., 2009, *Ionospheres - physics, plasma physics, and chemistry*. Cambridge university press
- Skrzypkowski M. P., Gougousi T., Johnsen R., Golde M. F., 1998, *J. Chem. Phys.*, 108, 8400
- Snodgrass C., Lowry S. C., Fitzsimmons A., 2008, *MNRAS*, 385, 737
- Snodgrass C., Meech K., Hainaut O., 2010, *A&A.*, 516, L9
- Tobiska W. K., Woods T., Eparvier F., et al., 2000, *J. Atmos. Solar-Terr. Phys.*, 62, 1233
- Tozzi G. P., Feldman P. D., Festou M. C., 1998, *A&A.*, 330, 753

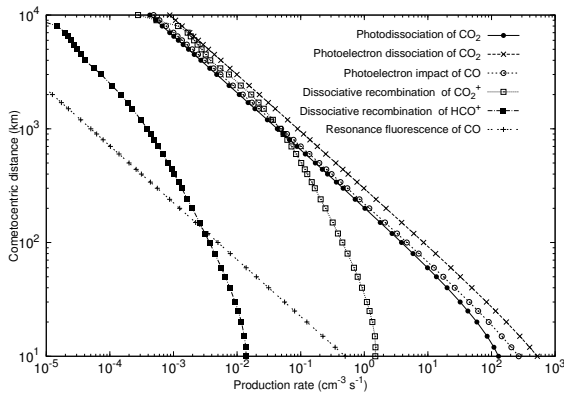
**Table 1.** Reactions for the production and loss of CO(*a*<sup>3</sup>Π)

Reaction	Rate(cm <sup>-3</sup> s <sup>-1</sup> or s <sup>-1</sup> )	Reference
CO <sub>2</sub> + hν → CO( <i>a</i> <sup>3</sup> Π) + O( <sup>3</sup> P)	Model	Shunk & Nagy (2009); Lawrence (1972)
CO + hν → CO( <i>a</i> <sup>3</sup> Π)	1.69 × 10 <sup>-9</sup>	Weaver et al. (1994)
CO <sub>2</sub> + e <sub>ph</sub> <sup>-</sup> → CO( <i>a</i> <sup>3</sup> Π) + O + e <sup>-</sup>	Model	<i>Present work</i>
CO + e <sub>ph</sub> <sup>-</sup> → CO( <i>a</i> <sup>3</sup> Π) + e <sup>-</sup>	Model	<i>Present work</i>
CO <sub>2</sub> <sup>+</sup> + e <sup>-</sup> → CO( <i>a</i> <sup>3</sup> Π) + O	6.5 × 10 <sup>-7</sup> (300/Te) <sup>0.8</sup> × 0.87 <sup>†</sup> × 0.29 <sup>‡</sup>	Seiersen et al. (2003); Rosati et al. (2003),
HCO <sup>+</sup> + e <sup>-</sup> → CO( <i>a</i> <sup>3</sup> Π) + H	0.23 <sup>§</sup> × 2.4 × 10 <sup>-7</sup> (300/Te) <sup>0.7</sup>	Rosati et al. (2007); Schmidt et al. (1988)
CO( <i>a</i> <sup>3</sup> Π) + hν → C + O	7.2 × 10 <sup>-5</sup>	Huebner et al. (1992)
CO( <i>a</i> <sup>3</sup> Π) + hν → CO <sup>+</sup> + e <sup>-</sup>	8.58 × 10 <sup>-6</sup>	Huebner et al. (1992)
CO( <i>a</i> <sup>3</sup> Π) + hν → O + C <sup>+</sup> + e <sup>-</sup>	2.45 × 10 <sup>-8</sup>	Huebner et al. (1992)
CO( <i>a</i> <sup>3</sup> Π) + hν → C + O <sup>+</sup> + e <sup>-</sup>	2.06 × 10 <sup>-8</sup>	Huebner et al. (1992)
CO( <i>a</i> <sup>3</sup> Π) + H <sub>2</sub> O → CO + H <sub>2</sub> O	3.3 × 10 <sup>-10</sup>	Wysong (2000)
CO( <i>a</i> <sup>3</sup> Π) + CO <sub>2</sub> → CO + CO <sub>2</sub>	1.0 × 10 <sup>-11</sup>	Skrzypkowski et al. (1998)
CO( <i>a</i> <sup>3</sup> Π) + CO → CO + CO	5.7 × 10 <sup>-11</sup>	Wysong (2000)
CO( <i>a</i> <sup>3</sup> Π) + e <sub>ph</sub> <sup>-</sup> → CO <sup>+</sup> + 2e <sup>-</sup>	Model	<i>Present work</i>
CO( <i>a</i> <sup>3</sup> Π) → CO + hν	1.6 × 10 <sup>-9</sup>	Weaver et al. (1994)

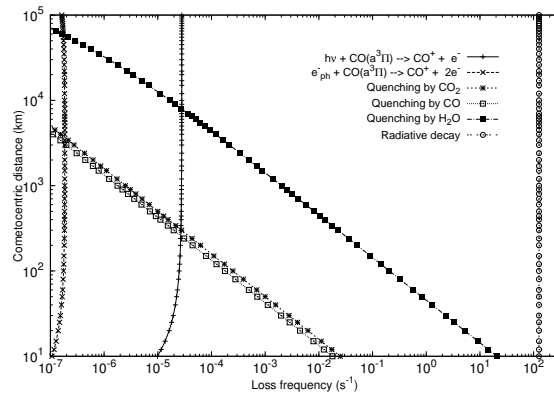
<sup>†</sup>0.87 is yield of dissociative recombination of CO<sub>2</sub><sup>+</sup> producing CO; <sup>‡</sup>0.29 is yield of CO(*a*<sup>3</sup>Π) produced from CO; <sup>§</sup>0.23 is yield of dissociative recombination of HCO<sup>+</sup> producing CO(*a*<sup>3</sup>Π); e<sub>ph</sub><sup>-</sup> = photoelectron.

**Table 2.** Production frequency (s<sup>-1</sup>) of CO(*a*<sup>3</sup>Π) for three different processes at 1 AU

Solar flux	Photodissociation of CO <sub>2</sub>		Photoelectron impact of CO <sub>2</sub>		Photoelectron impact of CO	
	Solar min (2010 Sep 1)	Solar max (1991 Sep 18)	Solar min (2010 Sep 1)	Solar max (1991 Sep 18)	Solar min (2010 Sep 1)	Solar max (1991 Sep 18)
S2K	1.1 × 10 <sup>-7</sup>	5.6 × 10 <sup>-7</sup>	3.4 × 10 <sup>-7</sup>	9.7 × 10 <sup>-7</sup>	1.5 × 10 <sup>-6</sup>	4.3 × 10 <sup>-6</sup>
EUVAC	1.7 × 10 <sup>-7</sup>	2.6 × 10 <sup>-7</sup>	3.5 × 10 <sup>-7</sup>	8.9 × 10 <sup>-7</sup>	1.5 × 10 <sup>-6</sup>	3.7 × 10 <sup>-6</sup>



**Figure 1.** Radial profiles of the production processes of CO(*a*<sup>3</sup>Π) in comet 103P/Hartley 2 for 0.5% CO and 4% CO<sub>2</sub> using EUVAC solar flux on 18-19 sept 1991.

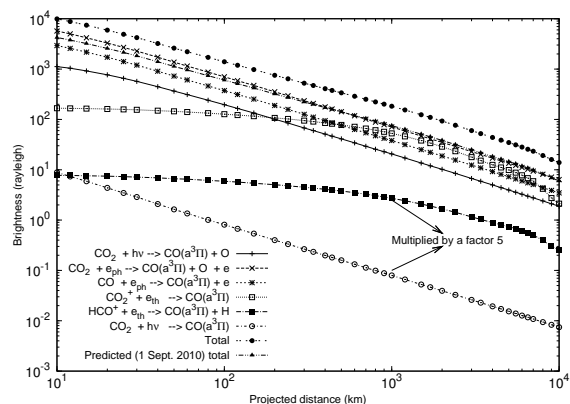


**Figure 2.** Radial profiles of the loss frequencies of CO(*a*<sup>3</sup>Π) in comet 103P/Hartley 2 for 0.5% CO and 4% CO<sub>2</sub> using EUVAC solar flux on 18-19 sept 1991.

**Table 3.** Calculated brightness of Cameron band on comet 103P/Hartley 2 for different conditions

Relative abundance		Percentage contribution to total Cameron band for different process at three different projected radial distances (km)															Total Cameron band brightness (R)		
HST-slit averaged																			
CO <sub>2</sub> (%)	CO (%)	brightness (R)	Photodissociation of CO <sub>2</sub>			Photoelectron impact of CO <sub>2</sub>			Photoelectron impact of CO			Recombination of CO <sub>2</sub> <sup>+</sup>			HST-Slit Avg.	Height integrated nadir view			
		(1-0)	(0-0)	(0-1)	10 <sup>2</sup>	10 <sup>3</sup>	10 <sup>4</sup>	10 <sup>2</sup>	10 <sup>3</sup>	10 <sup>4</sup>	10 <sup>2</sup>	10 <sup>3</sup>	10 <sup>4</sup>	10 <sup>2</sup>	10 <sup>3</sup>	10 <sup>4</sup>			
EUVAC *																			
7	0.5	68	51	72	15	12	15	57	42	51	17	12	15	10	32	17	489	4769	
5	0.5	51	39	54	14	11	14	53	40	48	22	17	21	9	30	16	369	3710	
4	0.5	43	32	45	13	11	13	50	38	45	26	20	25	9	29	15	309	3172	
3	0.5	34	26	36	12	10	12	46	36	42	32	26	30	8	26	13	250	2628	
2	0.5	26	20	28	11	9	10	39	32	36	42	34	40	7	23	12	190	2078	
4	1	51	38	54	11	9	11	40	32	36	41	33	38	7	24	12	311	3948	
3	1	43	32	46	10	8	10	35	29	32	48	40	45	6	21	11	311	3420	
S2K *																			
7	0.5	86	64	91	26	21	25	48	36	44	15	11	14	10	31	16	622	5928	
5	0.5	65	48	68	25	20	24	45	35	42	20	15	19	9	29	15	467	4586	
4	0.5	54	40	57	23	19	22	43	33	40	24	19	22	8	28	14	390	3903	
3	0.5	43	32	46	22	18	21	40	31	37	30	23	28	8	26	13	313	3210	
2	0.5	33	25	35	19	16	18	35	28	32	39	32	37	8	23	11	236	2508	
3	1	53	40	56	17	15	17	31	25	29	45	37	42	6	21	10	384	4093	
2	1	43	32	45	14	12	14	25	21	23	55	47	52	5	18	8	308	3411	
Prediction EUVAC †																			
4	0.5	189 <sup>‡</sup>	142 <sup>‡</sup>	200 <sup>‡</sup>	21	18	20	46	37	42	25	20	23	7	23	12	-	1365	
S2K †																			
4	0.5	175 <sup>‡</sup>	130 <sup>‡</sup>	185 <sup>‡</sup>	16	14	16	50	41	46	28	23	25	7	23	13	-	1258	

\* Solar flux for 18 sept. 1991; † Solar flux for 1 sept. 2010 ‡ Relative contribution of band in total nadir view intensity.



**Figure 3.** Brightness profile of Cameron band on comet 103P/Hartley 2 as a function of projected distance from nucleus for 4% CO<sub>2</sub> and 0.5% of CO using EUVAC solar flux on 18-19 sept 1991. Resonance fluorescence values and dissociative recombination profiles are plotted after multiplication by a factor 5. Total brightness predicted profile for EPOXI mission period is also shown