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Evolving surface finite element method for the Cahn-Hilliard equation

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Abstract We use the evolving surface finite element method to solve a Cahn-Hilliard equation on an evolving surface with prescribed velocity. We start by deriving the equation using a conservation law and appropriate transport formulae and provide the necessary functional analytic setting. The finite element method relies on evolving an initial triangulation by moving the nodes according to the prescribed velocity. We go on to show a rigorous well-posedness result for the continuous equations by showing convergence, along a subsequence, of the finite element scheme. We conclude the paper by deriving error estimates and present various numerical examples.

Keywords Evolving surface finite element method · Cahn-Hilliard equation · triangulated surfaces · error analysis

1 Introduction

In this paper, we will study a Cahn-Hilliard equation posed on an evolving surface with prescribed velocity. The key methodology is to discretise the equations using the evolving surface finite element method [8] originally proposed for a surface heat equation. The idea is to take a triangulation of the initial

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surface and evolve the nodes along the velocity field. This leads to a family of discrete surfaces on which we can pose a variational form of the Cahn-Hilliard equation.

There are two key results in this paper: first, we show well posedness of the continuous scheme and, second, we show convergence of a finite element scheme. The well posedness result is proven by rigorously showing convergence, along a subsequence, of the discrete scheme. In contrast to the planar setting, there are extra difficulties in this work since the classical Bochner space set-up is unavailable to us. The finite element method is analysed under the assumption of higher regularity of the solution and shown to converge to the true solution quadratically with respect to the mesh size in an L^2 norm. The paper concludes with some numerical examples to show various properties of the methodology.

1.1 The Cahn-Hilliard equation

We assume we are given an evolving surface $\{\Gamma(t)\}$, for $t \in [0, T]$, which evolves according to a given underlying velocity field v which can be decomposed into normal (v_ν) and tangential components (v_τ) so that $v = v_\nu + v_\tau$. We seek a solution u of

$$\partial^\bullet u + u \nabla_\Gamma \cdot v = \Delta_\Gamma \left(-\varepsilon \Delta_\Gamma u + \frac{1}{\varepsilon} \psi'(u) \right) \quad \text{on } \bigcup_{t \in (0, T)} \Gamma(t) \times \{t\} \quad (1.1)$$

subject to the initial condition

$$u(\cdot, 0) = u_0 \quad \text{on } \Gamma(0) = \Gamma_0. \quad (1.2)$$

Here $\partial^\bullet u$ denotes the material derivative of u and $\Delta_\Gamma u$ the Laplace-Beltrami operator of u . The function ψ is a double well potential, which we will take to be given by

$$\psi(z) = \frac{1}{4}(z^2 - 1)^2. \quad (1.3)$$

The behaviour of the Cahn-Hilliard equation in the planar case is well studied [15]. Extra effects such as spatial or concentration dependent mobilities or more physically realistic potentials could also be solved with similar methods to those suggested in this paper. Such considerations are left for future work.

This Cahn-Hilliard equation is a simplification of the model for surface dissolution set out in [14, 21] arising from a conservation law. The model [27] takes a different approach and considers a gradient flow for an energy consisting of the sum of the Ginzburg-Landau functional and a Helfrich energy on a stationary surface. One could alternatively couple the evolution of the surface to the surface field u and recover a gradient flow of the Ginzburg-Landau functional [17, 18].

The results in this work can be seen as a generalisation of the work of [6] to evolving surfaces. That work considers a fully discrete approximation of

a Cahn-Hilliard equation posed on a two-dimensional stationary surface with boundary (with a zero Dirichlet boundary condition) under the assumption $u_0 \in H_0^1(\Gamma) \cap H^2(\Gamma)$ and $\Delta_\Gamma u_0 \in H_0^1(\Gamma) \cap W^{1,2+\gamma}(\Gamma)$ for $\gamma \in (0, 1)$. Their method uses a triangulated surface for the spatial discretisation and a Crank-Nicolson scheme in time. They show an error estimate of the form

$$\max_m \|u_h^m - u^{-\ell}(t_m)\|_{L^2(\Gamma_h)} \leq c(h^2 + \tau^2),$$

where $0 = t_0 < t_1 < \dots < t_m < \dots < t_M = T$ is a partition of time with fixed time step τ and $u^{-\ell}$ is the inverse lift (3.21) of the continuous solution u .

1.2 Outline of paper

The paper is laid out as follows. In Section two, we will derive a Cahn-Hilliard equation on an evolving surface using a local conservation law. We introduce the notation for partial differential equations on evolving surfaces taken from [5, 11] and state any assumptions on the smoothness of the surfaces and its evolution we require. The third section introduces a finite element discretisation of the continuous equations. We describe the process of triangulating an evolving surface and how we formulate the space discrete-time continuous problem as a system of ordinary differential equations. This section is completed by showing some domain perturbation results relating geometric quantities on the discrete and smooth surfaces. Well posedness of the continuous equations is addressed in the fourth section. An existence result is achieved by showing convergence, along a subsequence, of the discrete solutions as the mesh size tends to zero. In Section five, we analyse the errors introduced by our finite element scheme and go on to show an optimal order error estimate. Some numerical experiments are shown in the sixth section backing up the analytical results.

We will use a Gronwall inequality as a standard tool in the analysis which leads to exponential dependence on ε in most bounds. We are not interested in taking $\varepsilon \rightarrow 0$ in this work so will simply write c_ε for a generic constant which depends on ε .

2 Derivation of continuous equations

In this section, we will derive a Cahn-Hilliard equation on an evolving surface as a conservative advection-diffusion equation. We will also introduce functional analytic setting and definition of solution that will be used.

2.1 Assumptions on the evolving surface

Given a final time $T > 0$, for each time $t \in [0, T]$, we write $\Gamma(t)$ for a compact, smooth, connected n -dimensional hypersurface in \mathbb{R}^{n+1} for $n = 1, 2$ or 3 and $\Gamma_0 = \Gamma(0)$. We assume that $\Gamma(t)$ is the boundary of an open, bounded domain

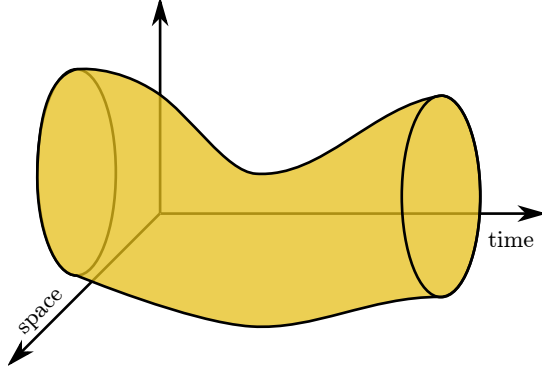


Fig. 2.1: A sketch of the space-time domain \mathcal{G}_T .

$\Omega(t)$. It follows that $\Gamma(t)$ admits a description as the zero level set of a signed distance function $d(\cdot, t): \mathbb{R}^{n+1} \rightarrow \mathbb{R}$ so that $d(\cdot, t) < 0$ in $\Omega(t)$ and $d(\cdot, t) > 0$ in $\Omega(t)^c$. We denote by \mathcal{G}_T for the space-time domain given by

$$\mathcal{G}_T = \bigcup_{t \in [0, T]} \Gamma(t) \times \{t\}. \quad (2.1)$$

For our analysis, it is sufficient to consider $d(\cdot, t)$ locally to $\Gamma(t)$. We restrict our considerations to $\mathcal{N}(t)$, an open neighbourhood of $\Gamma(t)$. We choose $\mathcal{N}(t)$ so that $|\nabla d(x, t)| \neq 0$ for $x \in \mathcal{N}(t)$ and assume that

$$d, d_t, d_{x_i}, d_{x_i x_j} \in C^2(\mathcal{N}_T) \quad \text{for } i, j = 1, \dots, n+1;$$

here $\mathcal{N}_T = \bigcup_{t \in [0, T]} \mathcal{N}(t) \times \{t\}$. The orientation of $\Gamma(t)$ is fixed by choosing ν as the outward pointing normal, so that $\nu(x, t) = \nabla d(x, t)$. For $(x, t) \in \mathcal{G}_T$, we denote $P = P(x, t)$ the projection operator onto the tangent space $T_x \Gamma(t)$, given by $P_{ij}(x, t) = \delta_{ij} - \nu_i(x, t) \nu_j(x, t)$ and by $\mathcal{H} = \mathcal{H}(x, t)$ the (extended) Weingarten map (or shape operator),

$$\mathcal{H}_{ij}(x, t) = (\nu_i(x, t))_{x_j} = d_{x_i x_j}(x, t).$$

We will use the fact that $P\mathcal{H} = \mathcal{H}P = \mathcal{H}$. Finally, we denote by $H = H(x, t)$ the mean curvature of $\Gamma(t)$

$$H(x, t) = \text{trace } \mathcal{H}(x, t) = \sum_{i=1}^{n+1} \mathcal{H}_{ii}(x, t).$$

For a function $\eta: \Gamma(t) \rightarrow \mathbb{R}$, we define its tangential gradient $\nabla_\Gamma \eta$ by

$$\nabla_\Gamma \eta = \nabla \tilde{\eta} - \nabla \tilde{\eta} \cdot \nu \nu = P \nabla \tilde{\eta},$$

where $\tilde{\eta}$ is a smooth extension of η away from $\Gamma(t)$. It can be shown that this definition is independent of the choice of extension. We denote the $n+1$ components of $\nabla_{\Gamma}\eta$ by

$$\nabla_{\Gamma}\eta = (\underline{D}_1\eta, \dots, \underline{D}_{n+1}\eta).$$

The Laplace-Beltrami operator is given by

$$\Delta_{\Gamma}\eta = \nabla_{\Gamma} \cdot \nabla_{\Gamma}\eta = \sum_{j=1}^{n+1} \underline{D}_j \underline{D}_j \eta.$$

We will denote by $d\sigma$ the surface measure on $\Gamma(t)$ which admits the following formula for partial integration for a portion $\mathcal{M}(t) \subseteq \Gamma(t)$ [11, Theorem 2.10]:

$$\int_{\mathcal{M}(t)} \nabla_{\Gamma}\eta \, d\sigma = \int_{\mathcal{M}(t)} \eta H \nu \, d\sigma + \int_{\partial\mathcal{M}(t)} \eta \mu \, d\sigma, \quad (2.2)$$

where μ is the co-normal to $\partial\mathcal{M}(t)$ which is normal to $\partial\mathcal{M}(t)$ but tangent to $\Gamma(t)$. If $\mathcal{M}(t) = \Gamma(t)$ and has no boundary, the boundary term vanishes. Furthermore, we have a Green's formula on $\Gamma(t)$ [11, Theorem 2.14]:

$$\int_{\Gamma(t)} \nabla_{\Gamma}\eta \cdot \nabla_{\Gamma}\varphi \, d\sigma = - \int_{\Gamma(t)} \varphi \Delta_{\Gamma}\eta \, d\sigma. \quad (2.3)$$

These formulae allow the definition of weak derivatives and Sobolev spaces. We define the space $W^{1,q}(\Gamma(t))$ by

$$W^{1,q}(\Gamma(t)) := \{ \eta \in L^q(\Gamma(t)) : \underline{D}_j \eta \in L^q(\Gamma(t)) \text{ for } j = 1, \dots, n+1 \},$$

with norm

$$\|\eta\|_{W^{1,q}(\Gamma(t))} = \left(\|\eta\|_{L^q(\Gamma(t))}^q + \|\nabla_{\Gamma}\eta\|_{L^q(\Gamma(t))}^q \right)^{\frac{1}{q}}.$$

This can be easily extended to higher order spaces. See [11] for details. We will use the notation $H^k(\Gamma(t))$ for $W^{k,2}(\Gamma(t))$.

We will make use of the following Sobolev embeddings:

Lemma 2.1 ([25, Theorems 2.5 and 2.6]) *For $\Gamma(t)$ as above, we have*

$$W^{1,q}(\Gamma(t)) \subset \begin{cases} L^{nq/(n-q)}(\Gamma(t)) & \text{for } q < n \\ C^0(\Gamma(t)) & \text{for } q > n. \end{cases} \quad (2.4)$$

Furthermore there exists a constant $c = c(n, q)$, independent of t , such that for any $\eta \in W^{1,q}(\Gamma(t))$,

$$\|\eta\|_{L^{nq/(n-q)}(\Gamma(t))} \leq c \|\eta\|_{W^{1,q}(\Gamma(t))} \quad \text{for } q < n \quad (2.5a)$$

$$\|\eta\|_{L^\infty(\Gamma(t))} \leq c \|\eta\|_{W^{1,q}(\Gamma(t))} \quad \text{for } q > n. \quad (2.5b)$$

In particular, this allows us to embed $H^1(\Gamma(t))$ in $L^6(\Gamma(t))$ for all dimensions ($n = 1, 2, 3$) so that $\|\psi'(\eta)\|_{L^2(\Gamma(t))} \leq c(\|\eta\|_{H^1(\Gamma(t))}^3 + \|\eta\|_{H^1(\Gamma(t))})$.

Further, we assume that for each $(x, t) \in \mathcal{N}_T$ there exists a unique $p = p(x, t) \in \Gamma(t)$, such that

$$x = p(x, t) + d(x, t)\nu(p(x, t), t). \quad (2.6)$$

See [23, Chapter 14] for a proof. We extend ν, P and \mathcal{H} to functions on \mathcal{N}_T by setting

$$\nu(x, t) = \nu(p(x, t), t) = \nabla d(x, t),$$

and similarly $P(x, t) = P(p(x, t), t) = \text{Id} - \nu(x, t) \otimes \nu(x, t)$ and $\mathcal{H}(x, t) = \nabla^2 d(x, t)$ for $(x, t) \in \mathcal{N}_T$.

Although it is sufficient to describe the evolution of the surface through a normal velocity, we wish to consider material surfaces for which a material particle, at $X(t)$ on $\Gamma(t)$, has a material velocity $\dot{X}(t)$ not necessarily only in the normal direction. The normal velocity of the surface can be calculated to be $v_\nu = -d_t \nu$. We say v_τ is a tangential velocity field if $v_\tau \cdot \nu = 0$ in \mathcal{N}_T . Given a tangential velocity field v_τ , we call

$$v := v_\tau + v_\nu$$

a material velocity field. We assume that we are given a global velocity field v so that points $X(t)$ evolve with the velocity $\dot{X}(t) = v(X(t), t)$. We will assume that $v \in C^2(\mathcal{N}_T)$.

2.2 Material derivative and transport formulae

Given a family of surfaces $\{\Gamma(t)\}$ evolving in time with normal velocity field v_ν , we define the normal time derivative ∂° of a function $\eta: \mathcal{G}_T \rightarrow \mathbb{R}$ by

$$\partial^\circ \eta := \frac{\partial \tilde{\eta}}{\partial t} + v_\nu \cdot \nabla \tilde{\eta}. \quad (2.7)$$

Here, $\tilde{\eta}$ denotes a smooth extension of η to \mathcal{N}_T . This derivative describes how a quantity η evolves in time with respect to the evolution of $\Gamma(t)$. It can be shown that this definition is an intrinsic surface derivative, independent of the choice of extension.

Given a tangential vector field v_τ , we define the material derivative of a scalar function $\eta: \mathcal{G}_T \rightarrow \mathbb{R}$, by

$$\partial^\bullet \eta := \partial^\circ \eta + v_\tau \cdot \nabla_\Gamma \eta = \frac{\partial \tilde{\eta}}{\partial t} + v \cdot \nabla \tilde{\eta}.$$

The following formula shows the significance of the material derivative. The result is a generalisation of the classical Reynolds' Transport Formula to curved domains.

Lemma 2.2 (Transport formula [12, Lemma 2.1]) *Let $\mathcal{M}(t)$ be an evolving surface with normal velocity v_ν . Let v_τ be a tangential velocity field on $\mathcal{M}(t)$. Let the boundary $\partial\mathcal{M}(t)$ evolve with velocity $v = v_\nu + v_\tau$. Assume that η, φ are functions such that all the following quantities exist. Then, we obtain the identity*

$$\frac{d}{dt} \int_{\mathcal{M}(t)} \eta \, d\sigma = \int_{\mathcal{M}(t)} \partial^\bullet \eta + \eta \nabla_\Gamma \cdot v \, d\sigma. \quad (2.8)$$

Furthermore, we have

$$\frac{d}{dt} \int_{\mathcal{M}(t)} \eta \varphi \, d\sigma = \int_{\mathcal{M}(t)} \partial^\bullet \eta \varphi + \eta \tilde{\partial}^\bullet \varphi + \eta \varphi \nabla_\Gamma \cdot v \, d\sigma. \quad (2.9)$$

Let $\mathcal{A} = \mathcal{A}(x, t)$ be a matrix which is positive definite on the tangent space to $\Gamma(t)$. Denote by $D(v)$ the rate of deformation tensor given by

$$D(v)_{ij} = \frac{1}{2} \sum_{k=1}^{n+1} (\mathcal{A}_{ik} \underline{D}_k v_j + \mathcal{A}_{jk} \underline{D}_k v_i) \quad \text{for } i, j = 1, \dots, n+1, \quad (2.10)$$

and by $\mathcal{B}(v)$ the tensor

$$\mathcal{B}(v) := \partial^\bullet \mathcal{A} + \nabla_\Gamma \cdot v \mathcal{A} - 2D(v). \quad (2.11)$$

Then we have the formula

$$\begin{aligned} \frac{d}{dt} \int_{\mathcal{M}(t)} \mathcal{A} \nabla_\Gamma \eta \cdot \nabla_\Gamma \varphi \, d\sigma &= \int_{\mathcal{M}(t)} \mathcal{A} \nabla_\Gamma \partial^\bullet \eta \cdot \nabla_\Gamma \varphi + \mathcal{A} \nabla_\Gamma \eta \cdot \nabla_\Gamma \partial^\bullet \varphi \, d\sigma \\ &\quad + \int_{\mathcal{M}(t)} \mathcal{B}(v) \nabla_\Gamma \eta \cdot \nabla_\Gamma \varphi \, d\sigma. \end{aligned} \quad (2.12)$$

We conclude this subsection with a result allowing us to extend functions defined on one surface to the whole space-time domain.

Lemma 2.3 *Fix $t \in [0, T]$ and let $\eta \in H^1(\Gamma(t))$, respectively $C^1(\Gamma(t))$. Then there exists an extension $\tilde{\eta}: \mathcal{G}_T \rightarrow \mathbb{R}$ such that $\tilde{\eta}|_t = \eta$ and $\tilde{\eta} \in H^1(\Gamma(s))$, resp. $C^1(\Gamma(s))$, for all times $s \in [0, T]$ and $\partial^\bullet \tilde{\eta} = 0$.*

Proof The ordinary differential equation:

$$\frac{d}{ds} X(s) = v(X(s), s) \quad \text{for } s \in [0, T], \quad X(t) = x,$$

determines a flow $\phi_s(x)$ on \mathcal{G}_T for $x \in \Gamma(t)$ such that

$$\phi_s(x) \in \Gamma(s) \quad \text{for all } s \in [0, T] \quad \text{and } \phi_t(x) = x.$$

Our assumptions on v imply that $\phi_s: \Gamma(t) \rightarrow \Gamma(s)$ and $(\phi_s)^{-1}: \Gamma(s) \rightarrow \Gamma(t)$ are both C^1 mappings [24, Theorem 3.1].

We define the extension $\tilde{\eta}$ by

$$\tilde{\eta}(x, s) := \eta((\phi_s)^{-1}(x)) \quad \text{for } (x, s) \in \mathcal{G}_T.$$

It is clear that since $(\phi_s)^{-1} \in C^1(\Gamma(t); \Gamma(s))$, we have $\tilde{\eta} \in H^1(\Gamma(s))$ (resp. $C^1(\Gamma(s))$) for all times $s \in [0, T]$.

Finally, we can calculate for $y = (\phi_s)^{-1}(x)$,

$$\partial^\bullet \tilde{\eta}(x, s) = \frac{d}{ds} \tilde{\eta}(\phi_s(y), s) = \frac{d}{ds} \eta(y) = 0 \quad \text{for } (x, s) \in \mathcal{G}_T,$$

which shows the result. \square

2.3 Derivation of Cahn-Hilliard equations

We will consider a conservation law on an evolving surface with a diffusive flux driven by a chemical potential. This is the approach taken by [21]. In general, the Ginzburg-Landau functional on $\Gamma(t)$ will not decrease along the trajectory of solutions.

Let u represent a density of a scalar quantity on $\Gamma(t)$. Following [12], we arrive at the pointwise conservation law

$$\partial^\circ u + u \nabla_\Gamma \cdot v_\nu + \nabla_\Gamma \cdot q = 0. \quad (2.13)$$

Here q represents the tangential flux of u on $\{\Gamma(t)\}$.

We will assume that the flux q is the sum of a diffusive flux q_d and an advective flux q_a :

$$q_d = -\nabla_\Gamma w \quad \text{and} \quad q_a = uv_\tau.$$

The diffusive flux is driven by the gradient of chemical potential w gives us the split system [16]

$$\partial^\bullet u + u \nabla_\Gamma \cdot v - \Delta_\Gamma w = 0 \quad (2.14a)$$

$$-\varepsilon \Delta_\Gamma u + \frac{1}{\varepsilon} \psi'(u) - w = 0. \quad (2.14b)$$

This leads to the fourth order Cahn-Hilliard equation on \mathcal{G}_T :

$$\partial^\bullet u + u \nabla_\Gamma \cdot v = \Delta_\Gamma \left(-\varepsilon \Delta_\Gamma u + \frac{1}{\varepsilon} \psi'(u) \right). \quad (2.15)$$

We close the system with the initial condition

$$u(\cdot, 0) = u_0 \quad \text{on } \Gamma_0. \quad (2.16)$$

There are no boundary conditions since the boundary of $\Gamma(t)$ is empty.

Remark 2.1 One can derive the Cahn-Hilliard equations posed in a Cartesian domain as an H^{-1} gradient flow of the Ginzburg-Landau functional. To obtain a gradient flow on an evolving surface, there would need to be a model for v and which would lead to a coupled system for u and v . In terms of modelling, we feel these extra terms are geometric terms determining an evolution equation for the surface, which we assume is given. Therefore, we do not consider such terms in this work.

2.4 Solution spaces

In standard parabolic theory one looks for solutions in Bochner spaces. Considering our Cahn-Hilliard equation on a Cartesian domain Ω [15], one would expect solutions to live in the spaces

$$u \in L^\infty(0, T; H^1(\Omega)), u' \in L^2(0, T; H^{-1}(\Omega)), w \in L^2(0, T; H^1(\Omega)).$$

These spaces are constructed by considering u as a function from $(0, T)$ into the Hilbert space $H^1(\Omega)$. We would like to extend this definition so that $u(t)$ is in the now time-dependent Hilbert space $H^1(\Gamma(t))$. We consider Sobolev spaces over the space-time domain \mathcal{G}_T . We will write $\nabla_{\mathcal{G}_T}$ for the space-time gradient and $d\sigma_T$ for the space-time measure on \mathcal{G}_T . This approach is similar to the Eulerian formulation of [28]. We contrast our approach with that of [33], who proposed using an equivalent formulation using a reference domain.

We start by presenting the space-time domains $L^2(\mathcal{G}_T)$ and $H^1(\mathcal{G}_T)$ defined by

$$\begin{aligned} L^2(\mathcal{G}_T) &:= \left\{ \eta \in L^1_{\text{loc}}(\mathcal{G}_T) : \int_{\mathcal{G}_T} \eta^2 d\sigma_T < +\infty \right\} \\ H^1(\mathcal{G}_T) &:= \left\{ \eta \in L^2(\mathcal{G}_T) : \nabla_{\mathcal{G}_T} \eta \in L^2(\mathcal{G}_T) \right\}. \end{aligned}$$

with norms

$$\begin{aligned} \|\eta\|_{L^2(\mathcal{G}_T)} &:= \left(\int_{\mathcal{G}_T} \eta^2 d\sigma_T \right)^{\frac{1}{2}} \\ \|\eta\|_{H^1(\mathcal{G}_T)} &:= \left(\|\eta\|_{L^2(\mathcal{G}_T)}^2 + \|\nabla_{\mathcal{G}_T} \eta\|_{L^2(\mathcal{G}_T)}^2 \right)^{\frac{1}{2}}. \end{aligned}$$

Proposition 2.1 ([25, Theorem 2.9]) *The space $H^1(\mathcal{G}_T)$ is compactly embedded into $L^2(\mathcal{G}_T)$.*

Using the identities,

$$\int_0^T \int_{\Gamma(t)} \eta d\sigma dt = \int_{\mathcal{G}_T} \frac{\eta}{\sqrt{1 + |v_\nu|^2}} d\sigma_T,$$

and

$$\nabla_{\mathcal{G}_T} \eta = \left(\nabla_\Gamma \eta + \frac{\partial^\circ \eta v_\nu}{1 + |v_\nu|^2}, \frac{\partial^\circ \eta}{1 + |v_\nu|^2} \right),$$

our assumptions on v imply that the space-time norms can be replaced with the equivalent norms

$$\begin{aligned} \|\eta\|'_{L^2(\mathcal{G}_T)} &:= \left(\int_0^T \int_{\Gamma(t)} \eta^2 d\sigma dt \right)^{\frac{1}{2}} \\ \|\eta\|'_{H^1(\mathcal{G}_T)} &:= \left(\int_0^T \int_{\Gamma(t)} \eta^2 + |\nabla_\Gamma \eta|^2 + (\partial^\circ \eta)^2 d\sigma dt \right)^{\frac{1}{2}}. \end{aligned}$$

We will use the equivalent primed norms (dropping the prime) on $L^2(\mathcal{G}_T)$ and $H^1(\mathcal{G}_T)$ in the following.

We define the space $L^2_{L^2}$ by

$$L^2_{L^2} := \left\{ \eta \in L^1_{\text{loc}}(\mathcal{G}_T) : \int_0^T \int_{\Gamma(t)} \eta \, d\sigma \, dt < +\infty \right\},$$

with the inner product

$$(\eta, \xi)_{L^2_{L^2}} := \int_0^T \int_{\Gamma(t)} \eta \xi \, d\sigma \, dt.$$

It is clear that $L^2_{L^2}$ is equivalent to $L^2(\mathcal{G}_T)$ and hence is a Hilbert space.

Next, we define the space $L^2_{H^1}$ as

$$L^2_{H^1} := \left\{ \eta : \mathcal{G}_T \rightarrow \mathbb{R} : \eta \in L^2_{L^2} \text{ and } \nabla_{\Gamma} \eta \in (L^2_{L^2})^{n+1} \right\},$$

with the inner product

$$(\eta, \xi)_{L^2_{H^1}} := \int_0^T \int_{\Gamma(t)} \nabla_{\Gamma} \eta \cdot \nabla_{\Gamma} \xi + \eta \xi \, d\sigma \, dt,$$

where $\nabla_{\Gamma} \eta$ should be interpreted in the weak sense. Notice that elements of this space are weakly differentiable at almost every time.

Lemma 2.4 *The space $L^2_{H^1}$ is a Hilbert space.*

Proof It is clear that $L^2_{H^1}$ is an inner product space and we are left to show completeness. Let η_k be a Cauchy sequence in $L^2_{H^1}$. This implies that η_k and $\nabla_{\Gamma} \eta_k$ are Cauchy sequences in $L^2(\mathcal{G}_T)$ and $(L^2(\mathcal{G}_T))^{n+1}$. This means that there exists $\eta \in L^2(\mathcal{G}_T)$, $\xi \in (L^2(\mathcal{G}_T))^{n+1}$ such that

$$\|\eta_k - \eta\|_{L^2(\mathcal{G}_T)} + \|\nabla_{\Gamma} \eta_k - \xi\|_{L^2(\mathcal{G}_T)} \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

Fix $t^* \in (0, T)$ and let $\varphi \in C^1(\Gamma(t^*))$ and $\alpha \in C(0, T)$. Using [Lemma 2.3](#), we can construct $\tilde{\varphi} : \mathcal{G}_T \rightarrow \mathbb{R}$ such that $\tilde{\varphi}(\cdot, t) = \varphi$ and $\tilde{\varphi} \in C^1(\Gamma(t))$ for each time $t \in (0, T)$. Then, for $j = 1, \dots, n+1$, we obtain

$$\begin{aligned} & \int_0^T \int_{\Gamma(t)} \eta \underline{D}_j(\alpha \tilde{\varphi}) + \xi_j(\alpha \tilde{\varphi}) \, d\sigma \, dt \\ &= \int_0^T \int_{\Gamma(t)} (\eta - \eta_k) \underline{D}_j(\alpha \tilde{\varphi}) + (\eta_k \underline{D}_j(\alpha \tilde{\varphi}) + \xi_j(\alpha \tilde{\varphi})) \, d\sigma \, dt \\ &= \int_0^T \int_{\Gamma(t)} (\eta - \eta_k) \underline{D}_j(\alpha \tilde{\varphi}) + (-\underline{D}_j \eta_k + \xi_j)(\alpha \tilde{\varphi}) \, d\sigma \, dt, \end{aligned}$$

where we have used the fact that η_k is weakly differentiable at almost every time. Taking the limit $k \rightarrow \infty$, we infer

$$\int_0^T \alpha \left(\int_{\Gamma(t)} \eta \underline{D}_j \tilde{\varphi} + \xi_j \tilde{\varphi} \, d\sigma \right) dt = 0.$$

Since this holds for all $\alpha \in C(0, T)$, by the Fundamental Lemma of the Calculus of Variations, at $t = t^*$, we have

$$\int_{\Gamma(t^*)} \eta \underline{D}_j \varphi + \xi_j \varphi \, d\sigma = 0 \quad \text{for all } \varphi \in C^1(\Gamma(t^*)).$$

Since the choice of t^* was arbitrary, we infer that ξ is the weak gradient of η for almost every time $t \in (0, T)$ and the proof is complete. \square

The equivalence of norms implies that $\eta \in L^2_{H^1}$ with $\partial^\bullet \eta \in L^2_{L^2}$ if, and only if, $\eta \in H^1(\mathcal{G}_T)$.

For $1 \leq q \leq \infty$, we will define the space $L^q_{H^1}$ by

$$L^q_{H^1} := \left\{ \eta \in L^q(\mathcal{G}_T) : \|\eta\|_{L^q_{H^1}} < +\infty \right\},$$

with norm

$$\|\eta\|_{L^q_{H^1}} := \begin{cases} \left(\int_0^T \|\eta\|_{H^1(\Gamma(t))}^q \, dt \right)^{\frac{1}{q}} & \text{for } q < \infty, \\ \operatorname{ess\,sup}_{t \in (0, T)} \|\eta\|_{H^1(\Gamma(t))} & \text{for } q = \infty. \end{cases}$$

It is clear that $L^\infty_{H^1} \subset L^2_{H^1}$ and that

$$\|\eta\|_{L^2_{H^1}} \leq \sqrt{T} \|\eta\|_{L^\infty_{H^1}} \quad \text{for all } \eta \in L^\infty_{H^1}.$$

Finally, we define $L^\infty_{H^2}$ and $L^2_{H^2}$ by

$$L^\infty_{H^2} := \left\{ \eta \in L^2(\mathcal{G}_T) : \operatorname{ess\,sup}_{t \in (0, T)} \|\eta\|_{H^2(\Gamma(t))} < +\infty \right\}$$

$$L^2_{H^2} := \left\{ \eta \in L^2(\mathcal{G}_T) : \int_0^T \|\eta\|_{H^2(\Gamma(t))}^2 \, dt < +\infty \right\}.$$

Remark 2.2 As a restriction on our analysis we will only consider $\partial^\bullet u$ as a function in $L^2_{L^2}$ since we do not wish to consider a weak material derivative. Such considerations are left to future work.

We conclude this section with a result which will take an integral in time equality into an almost everywhere in time equality. The proof is the generalisation of a similar result given in [30, Lemma 7.4] for planar domains.

Lemma 2.5 *Let $\eta \in L^2_{H^1}$ with*

$$\int_0^T \int_{\Gamma(t)} \nabla_\Gamma \eta \cdot \nabla_\Gamma \xi + \eta \xi \, d\sigma \, dt = 0 \quad \text{for all } \xi \in L^2_{H^1}. \quad (2.17)$$

Then for almost all times $t \in (0, T)$,

$$\int_{\Gamma(t)} \nabla_\Gamma \eta \cdot \nabla_\Gamma \varphi + \eta \varphi \, d\sigma = 0 \quad \text{for all } \varphi \in L^2_{H^1}. \quad (2.18)$$

Proof Fix $\varphi \in L^2_{H^1}$ and $\alpha \in C([0, T])$, then choosing $\xi = \alpha\varphi \in L^2_{H^1}$ and

$$0 = \int_0^T \int_{\Gamma(t)} \nabla_{\Gamma} \eta \cdot \nabla_{\Gamma} \xi + \eta \xi \, d\sigma \, dt = \int_0^T \alpha \left(\int_{\Gamma(t)} \nabla_{\Gamma} \eta \cdot \nabla_{\Gamma} \varphi + \eta \varphi \, d\sigma \right) dt.$$

Since the choice of α was arbitrary, the Fundamental Lemma of the Calculus of Variations implies the result. \square

2.5 Weak and variational form

We start by multiplying (2.14a, 2.14b) by a test function φ and apply integration by parts to the Laplacian terms to give the weak form. This will be the definition of solution used throughout this paper. Existence and uniqueness of solutions will be shown Section 4.

Definition 2.1 (Weak solution) We say that the pair $(u, w): \mathcal{G}_T \rightarrow \mathbb{R}^2$, with $u \in L^\infty_{H^1} \cap H^1(\mathcal{G}_T)$ and $w \in L^2_{H^1}$, are a weak solution of the Cahn-Hilliard equation (2.15) if, for almost every time $t \in (0, T)$,

$$\int_{\Gamma(t)} \partial^\bullet u \varphi + u \varphi \nabla_{\Gamma} \cdot v + \nabla_{\Gamma} w \cdot \nabla_{\Gamma} \varphi \, d\sigma = 0 \quad (2.19a)$$

$$\int_{\Gamma(t)} \varepsilon \nabla_{\Gamma} u \cdot \nabla_{\Gamma} \varphi + \frac{1}{\varepsilon} \psi'(u) \varphi - w \varphi \, d\sigma = 0, \quad (2.19b)$$

for all $\varphi \in L^2_{H^1}$,

and $u(\cdot, 0) = u_0$ pointwise almost everywhere in Γ_0 .

Restricting our thoughts to $\varphi \in H^1(\mathcal{G}_T)$, applying the transport formula to the first two terms in (2.19a) gives the variational formulation:

$$\frac{d}{dt} \left(\int_{\Gamma(t)} u \varphi \, d\sigma \right) + \int_{\Gamma(t)} \nabla_{\Gamma} w \cdot \nabla_{\Gamma} \varphi \, d\sigma = \int_{\Gamma(t)} u \partial^\bullet \varphi \, d\sigma \quad (2.20a)$$

$$\int_{\Gamma(t)} \varepsilon \nabla_{\Gamma} u \cdot \nabla_{\Gamma} \varphi + \frac{1}{\varepsilon} \psi'(u) \varphi \, d\sigma = \int_{\Gamma(t)} w \varphi \, d\sigma. \quad (2.20b)$$

We remark that this formulation has no explicit mention of the velocity field v and will be the basis of our finite element calculations.

It will be useful to write these equations using abstract bilinear forms. We define the following three to describe the above equations for $\eta, \varphi \in H^1(\Gamma(t))$:

$$\begin{aligned} m(\eta, \varphi) &= \int_{\Gamma(t)} \eta \varphi \, d\sigma & a(\eta, \varphi) &= \int_{\Gamma(t)} \nabla_{\Gamma} \eta \cdot \nabla_{\Gamma} \varphi \, d\sigma \\ g(v; \eta, \varphi) &= \int_{\Gamma(t)} \eta \varphi \nabla_{\Gamma} \cdot v \, d\sigma. \end{aligned}$$

This lets us write (2.19) as

$$\begin{aligned} m(\partial^\bullet u, \varphi) + g(v; u, \varphi) + a(w, \varphi) &= 0 \\ \varepsilon a(u, \varphi) + \frac{1}{\varepsilon} m(\psi'(u), \varphi) - m(w, \varphi) &= 0, \end{aligned} \quad (2.21)$$

and (2.20) as

$$\begin{aligned} \frac{d}{dt} m(u, \varphi) + a(w, \varphi) &= m(u, \partial^\bullet \varphi) \\ \varepsilon a(u, \varphi) + \frac{1}{\varepsilon} m(\psi'(u), \varphi) &= m(w, \varphi). \end{aligned} \quad (2.22)$$

We may also write the results of Lemma 2.2 in this form:

$$\begin{aligned} \frac{d}{dt} m(\eta, \varphi) &= m(\partial^\bullet \eta, \varphi) + m(\eta, \partial^\bullet \varphi) + g(v; \eta, \varphi) \\ \frac{d}{dt} a(\eta, \varphi) &= a(\partial^\bullet \eta, \varphi) + a(\eta, \partial^\bullet \varphi) + b(v; \eta, \varphi), \end{aligned}$$

with the addition of

$$b(v; \eta, \varphi) = \int_{\Gamma(t)} \mathcal{B}(v) \nabla_{\Gamma} \eta \cdot \nabla_{\Gamma} \varphi \, d\sigma,$$

using $\mathcal{A} = \text{Id}$ in the definition of $\mathcal{B}(v)$.

3 Finite element approximation

In this section, we propose a finite element method for approximating solutions of the Cahn-Hilliard equation (2.15) based on the evolving surface finite element method [8].

3.1 Evolving triangulation and discrete material derivative

Let $\Gamma_{h,0}$ be a polyhedral approximation of the initial surface Γ_0 with the restriction that the nodes $\{X_j^0\}_{j=1}^N$ of $\Gamma_{h,0}$ lie on Γ_0 . We evolve the nodes $\{X_j(t)\}_{j=1}^N$ by the smooth surface velocity:

$$\dot{X}_j(t) = v(X_j(t), t), \quad X_j(0) = X_j^0, \quad \text{for } j = 1, \dots, N.$$

Linearly interpolating between these nodes defines a family of discrete surfaces $\{\Gamma_h(t)\}$. At each time, we assume that we have a triangulation $\mathcal{T}_h(t)$ of $\Gamma_h(t)$, with h the maximum diameter of elements in $\mathcal{T}_h(t)$ uniformly in time:

$$h := \sup_{t \in (0, T)} \max_{E(t) \in \mathcal{T}_h(t)} \text{diam } E(t). \quad (3.1)$$

We assume this triangulation is quasi-uniform [2] uniformly in time.

Remark 3.1 In practical situations, assuming a uniformly regular mesh may not be feasible. Large surface deformations can lead to poor quality triangulations with deformed elements. In such cases, re-meshing may be required [4, 14]. Alternatively, one may use an arbitrary Lagrangian-Eulerian formulation by allowing extra tangential mesh motions [19, 20].

We define ν_h element-wise as the unit outward pointing normal to $\Gamma_h(t)$ and denote by ∇_{Γ_h} the tangential gradient on $\Gamma_h(t)$ defined element-wise by

$$\nabla_{\Gamma_h} \eta_h := \nabla \tilde{\eta}_h - (\nabla \tilde{\eta}_h \cdot \nu_h) \nu_h = (\text{Id} - \nu_h \otimes \nu_h) \nabla \tilde{\eta}_h =: P_h \nabla \tilde{\eta}_h.$$

This is a vector-valued quantity and we will denote its components by

$$\nabla_{\Gamma_h} \eta_h = (\underline{D}_{h,1} \eta_h, \dots, \underline{D}_{h,n+1} \eta_h)$$

We define the finite element space of piecewise linear functions on $\Gamma_h(t)$ by

$$S_h(t) := \{\phi_h \in C(\Gamma_h(t)) : \phi_h|_{E(t)} \text{ is affine linear, for each } E(t) \in \mathcal{T}_h(t)\}. \quad (3.2)$$

We will write $\{\phi_j^N(\cdot, t)\}_{j=1}^N$ for the nodal basis of $S_h(t)$ given by $\phi_j^N(X_i(t), t) = \delta_{ij}$.

The definition of a basis of $S_h(t)$ allows us to characterise the velocity of the surface $\{\Gamma_h(t)\}$. An arbitrary point $X(t)$ on $\Gamma_h(t)$ evolves according to the discrete velocity V_h given by

$$\dot{X}(t) = V_h(X(t), t) := \sum_{j=1}^N \dot{X}_j(t) \phi_j^N(X(t), t) = \sum_{j=1}^N v(X_j(t), t) \phi_j^N(X(t), t). \quad (3.3)$$

We will write $\mathcal{G}_{h,T}$ as the discrete equivalent to \mathcal{G}_T :

$$\mathcal{G}_{h,T} := \bigcup_{t \in (0,T)} \Gamma_h(t) \times \{t\}. \quad (3.4)$$

The discrete velocity V_h induces a discrete material derivative. For a scalar quantity η_h on $\mathcal{G}_{h,T}$, we define the discrete material derivative $\partial_h^\bullet \eta_h$ by

$$\partial_h^\bullet \eta_h := \partial_t \tilde{\eta}_h + \nabla \tilde{\eta}_h \cdot V_h, \quad (3.5)$$

where $\tilde{\eta}_h$ is an arbitrary extension of η_h to \mathcal{N}_T . This leads to the remarkable transport property of the basis functions $\{\phi_j^N\}$.

Lemma 3.1 (Transport of basis functions [8, Proposition 5.4]) *Let $\phi_j^N : \mathcal{G}_{h,T} \rightarrow \mathbb{R}$ be a nodal basis function as described above, then*

$$\partial_h^\bullet \phi_j^N = 0. \quad (3.6)$$

From a practical view point, a key advantage of this methodology is that, since basis functions have zero discrete material velocity, there is no mention of the velocity or curvature in the resulting finite element scheme.

These discrete quantities also satisfy a variant of the transport formula from [Lemma 2.2](#). We label the surface measure on $\Gamma_h(t)$ as $d\sigma_h$.

Lemma 3.2 (Transport lemma for triangulated surfaces [12, Lemma 4.2])

Let $\{\Gamma_h(t)\}$ be a discrete family of triangulated surfaces evolving with velocity V_h . Let η_h, ϕ_h be time-dependent finite element functions such that the following quantities exist. Then, we have

$$\frac{d}{dt} \int_{\Gamma_h(t)} \eta_h \, d\sigma_h = \int_{\Gamma_h(t)} \partial_h^\bullet \eta_h + \eta_h \nabla_{\Gamma_h} \cdot V_h \, d\sigma_h. \quad (3.7)$$

In particular, for the L^2 inner product this means that

$$\frac{d}{dt} \int_{\Gamma_h(t)} \eta_h \phi_h \, d\sigma_h = \int_{\Gamma_h(t)} (\partial_h^\bullet \eta_h) \phi_h + \eta_h (\partial_h^\bullet \phi_h) + \eta_h \phi_h \nabla_{\Gamma_h} \cdot V_h \, d\sigma_h, \quad (3.8)$$

and for the Dirichlet inner product, we obtain

$$\begin{aligned} & \frac{d}{dt} \int_{\Gamma_h(t)} \nabla_{\Gamma_h} \eta_h \cdot \nabla_{\Gamma_h} \phi_h \, d\sigma_h \\ &= \int_{\Gamma_h(t)} \nabla_{\Gamma_h} (\partial_h^\bullet \eta_h) \cdot \nabla_{\Gamma_h} \phi_h + \nabla_{\Gamma_h} \eta_h \cdot \nabla_{\Gamma_h} (\partial_h^\bullet \phi_h) \, d\sigma_h \\ &+ \sum_{E(t) \in \mathcal{E}_h(t)} \int_{E(t)} \mathcal{B}_h(V_h) \nabla_{\Gamma_h} \eta_h \cdot \nabla_{\Gamma_h} \phi_h \, d\sigma_h, \end{aligned} \quad (3.9)$$

where

$$\mathcal{B}_h(V_h) = \frac{1}{2} (\nabla_{\Gamma_h} \cdot V_h) \text{Id} - D_h(V_h) \quad \text{and} \quad D_h(V_h)_{ij} = \frac{1}{2} (\underline{D}_{h,i} V_{h,j} + \underline{D}_{h,j} V_{h,i}).$$

Lemma 3.3 Under our assumptions on $\{\Gamma_h(t)\}$, we have that

$$\sup_{t \in [0, T]} \left(\|\nabla_{\Gamma_h} \cdot V_h\|_{L^\infty(\Gamma_h(t))} + \|\mathcal{B}_h(V_h)\|_{L^\infty(\Gamma_h(t))} \right) \leq c \sup_{t \in [0, T]} \|v\|_{C^2(\mathcal{N}_T)}. \quad (3.10)$$

Proof The result follows from applying the geometric estimates (3.23) and (3.42) along with our assumption that $v \in C^2(\mathcal{N}_T)$. \square

3.2 Finite element scheme

We will assume that there exists a mesh size $h_0 > 0$ such that $\|U_0\|_{H^1(\Gamma_{h,0})}$ is bounded independently of h for $h < h_0$. This implies that there exists $C > 0$ such that for all $h < h_0$, we have

$$\mathcal{E}_0^h := \int_{\Gamma_{h,0}} \frac{\varepsilon}{2} |\nabla_{\Gamma_h} U_0|^2 + \frac{1}{\varepsilon} \psi(U_0) \, d\sigma_h < C. \quad (3.11)$$

Remark 3.2 One particular choice of initial condition will be to take U_0 as a suitable approximation of u_0 (for example, $\Pi_h u_0$ defined in (3.45)) for $u_0 \in H^2(\Gamma_0)$.

Our solution spaces will be

$$\begin{aligned}\tilde{S}_h^T &:= \{\phi_h \in C(\mathcal{G}_{h,T}) : \phi_h(\cdot, t) \in S_h(t) \text{ for all } t \in [0, T]\} \\ S_h^T &:= \{\phi_h \in \tilde{S}_h^T : \partial_h^\bullet \phi_h \in C(\mathcal{G}_{h,T})\}.\end{aligned}\quad (3.12)$$

The finite element scheme is: Given U_0 , find $U_h \in S_h^T$ and $W_h \in \tilde{S}_h^T$ such that for almost every time $t \in (0, T)$

$$\frac{d}{dt} \left(\int_{\Gamma_h(t)} U_h \phi_h \, d\sigma_h \right) + \int_{\Gamma_h(t)} \nabla_{\Gamma_h} W_h \cdot \nabla_{\Gamma_h} \phi_h \, d\sigma_h = \int_{\Gamma_h(t)} U_h \partial_h^\bullet \phi_h \, d\sigma_h \quad (3.13a)$$

$$\int_{\Gamma_h(t)} \varepsilon \nabla_{\Gamma_h} U_h \cdot \nabla_{\Gamma_h} \phi_h + \frac{1}{\varepsilon} \psi'(U_h) \phi_h \, d\sigma_h = \int_{\Gamma_h(t)} W_h \phi_h \, d\sigma_h \quad (3.13b)$$

for all $\phi_h \in S_h(t)$,

subject to the initial condition

$$U_h(\cdot, 0) = U_0. \quad (3.14)$$

The transport formula (3.8) implies that, for $\phi_h \in S_h^T$, (3.13a) is equivalent to

$$\int_{\Gamma_h(t)} \partial_h^\bullet U_h \phi_h + U_h \phi_h \nabla_{\Gamma_h} \cdot V_h + \nabla_{\Gamma_h} W_h \cdot \nabla_{\Gamma_h} \phi_h \, d\sigma_h = 0. \quad (3.15)$$

We can write these equations in matrix form. First, we will introduce vectors $\alpha(t), \beta(t) \in \mathbb{R}^N$ for the nodal values of U_h and W_h by

$$U_h(x, t) = \sum_{j=1}^N \alpha_j(t) \phi_j^N(x, t), \quad W_h(x, t) = \sum_{j=1}^N \beta_j(t) \phi_j^N(x, t) \quad \text{for } (x, t) \in \mathcal{G}_{h,T}.$$

In place of the bilinear forms, we have the mass matrix $\mathcal{M}(t)$ and stiffness matrix $\mathcal{S}(t)$:

$$\mathcal{M}(t)_{ij} = \int_{\Gamma_h(t)} \phi_i^N \phi_j^N \, d\sigma_h \quad \mathcal{S}(t)_{ij} = \int_{\Gamma_h(t)} \nabla_{\Gamma_h} \phi_i^N \cdot \nabla_{\Gamma_h} \phi_j^N \, d\sigma_h,$$

and in place of the non-linear term, we will write

$$\mathcal{F}(\alpha(t))_j = \int_{\Gamma_h(t)} \psi' \left(\sum_{i=1}^N \alpha_i(t) \phi_i^N \right) \phi_j^N \, d\sigma_h.$$

Using the transport of basis property (Lemma 3.1), we can write (3.13) as

$$\frac{d}{dt} (\mathcal{M}(t)\alpha(t)) + \mathcal{S}(t)\beta(t) = 0 \quad (3.16a)$$

$$\varepsilon \mathcal{S}(t)\alpha(t) + \frac{1}{\varepsilon} \mathcal{F}(\alpha(t)) - \mathcal{M}(t)\beta(t) = 0. \quad (3.16b)$$

Alternatively, eliminating $\beta(t)$, this can be written as

$$\frac{d}{dt} (\mathcal{M}(t)\alpha(t)) + \mathcal{S}(t)\mathcal{M}(t)^{-1} \left(\varepsilon \mathcal{S}(t)\alpha(t) + \frac{1}{\varepsilon} \mathcal{F}(\alpha(t)) \right) = 0. \quad (3.17)$$

One could also use lumped mass integration [32, Chapter 15] instead of the full mass matrix.

Notice that this is the same structure as a finite element discretisation of a Cahn-Hilliard equation posed on a planar domain. We now have time dependent matrices which need to be assembled on each time step. Various time stepping schemes have been considered for second-order parabolic problems on evolving surfaces [10, 13, 26].

Next, we introduce abstract notation which permit a more compact writing of the analysis that follows:

$$\begin{aligned} m_h(\eta_h, \phi_h) &= \int_{\Gamma_h(t)} \eta_h \phi_h \, d\sigma_h & a_h(\eta_h, \phi_h) &= \int_{\Gamma_h(t)} \nabla_{\Gamma_h} \eta_h \cdot \nabla_{\Gamma_h} \phi_h \, d\sigma_h \\ g_h(V_h; \eta_h, \phi_h) &= \int_{\Gamma_h(t)} \eta_h \phi_h \nabla_{\Gamma_h} \cdot V_h \, d\sigma_h. \end{aligned}$$

This lets us write (3.13) as

$$\begin{aligned} \frac{d}{dt} m_h(U_h, \phi_h) + a_h(W_h, \phi_h) &= m_h(U_h, \partial_h^\bullet \phi_h) \\ \varepsilon a_h(U_h, \phi_h) + \frac{1}{\varepsilon} m_h(\psi'(U_h), \phi_h) &= m_h(W_h, \phi_h), \end{aligned}$$

and (3.15) as

$$m_h(\partial_h^\bullet U_h, \phi_h) + g_h(V_h; U_h, \phi_h) + a_h(W_h, \phi_h) = 0.$$

The transport laws from Lemma 3.2 transfer to the abstract setting also:

$$\begin{aligned} \frac{d}{dt} m_h(\eta_h, \phi_h) &= m_h(\partial_h^\bullet \eta_h, \phi_h) + m_h(\eta_h, \partial_h^\bullet \phi_h) + g_h(V_h; \eta_h, \phi_h) \\ \frac{d}{dt} a_h(\eta_h, \phi_h) &= a_h(\partial_h^\bullet \eta_h, \phi_h) + a_h(\eta_h, \partial_h^\bullet \phi_h) + b_h(V_h; \eta_h, \phi_h), \end{aligned}$$

where

$$b_h(V_h; \eta_h, \phi_h) = \sum_{E(t) \in \mathcal{T}_h(t)} \int_{E(t)} \mathcal{B}_h(V_h) \nabla_{\Gamma_h} \eta_h \cdot \nabla_{\Gamma_h} \phi_h \, d\sigma_h.$$

Under the above assumptions, the following estimates are possible.

Theorem 3.1 (Well-posedness of the finite element scheme (3.13))

Under the above assumptions on U_0 and $\{\Gamma_h(t)\}$, there exists a unique solution pair $(U_h, W_h) \in S_h^T \times \tilde{S}_h^T$, both with C^1 in time nodal values, to the finite element scheme (3.13) and $\int_{\Gamma_h(t)} U_h \, d\sigma_h$ is conserved:

$$\int_{\Gamma_h(t)} U_h \, d\sigma_h = \int_{\Gamma_{h,0}} U_0 \, d\sigma_h \quad \text{for all } t \in (0, T). \quad (3.18)$$

Furthermore, there exists h_1 , $0 < h_1 < h_0$, and $C_0 > 0$, which depend on the final time T and the $H^1(\Gamma_{h,0})$ -norm of the initial condition U_0 , such that for all $h < h_1$ the following bound is satisfied:

$$\sup_{t \in (0, T)} \int_{\Gamma_h(t)} \frac{\varepsilon}{2} |\nabla_{\Gamma_h} U_h|^2 + \frac{1}{\varepsilon} \psi(U_h) d\sigma_h + \frac{1}{2} \int_0^T \|\nabla_{\Gamma_h} W_h\|_{L^2(\Gamma_h(t))}^2 dt \leq C_0. \quad (3.19)$$

The proof will be shown after we have proven some intermediate results.

3.3 Lifted finite elements

The following analysis will rely on lift operators defined using a time dependent closest point operator p (2.6). This lifting process will also be applied to the surface triangulation. This will induce a further discrete material velocity v_h which will describe how the lifts of triangles on $\{\Gamma(t)\}$ evolve.

First, for a function $\eta_h : \mathcal{G}_{h,T} \rightarrow \mathbb{R}$, we define its lift, $\eta_h^\ell : \mathcal{G}_T \rightarrow \mathbb{R}$, implicitly, by:

$$\eta_h^\ell(p(x, t), t) = \eta_h(x, t), \quad (3.20)$$

and, for a function $\eta : \mathcal{G}_T \rightarrow \mathbb{R}$, we define its inverse lift, $\eta^{-\ell} : \mathcal{G}_{h,T} \rightarrow \mathbb{R}$ by

$$\eta^{-\ell}(x, t) := \eta(p(x, t), t). \quad (3.21)$$

It is clear that these operations are inverses of each other

$$(\eta^{-\ell})^\ell = \eta \quad \text{and} \quad (\eta_h^\ell)^{-\ell} = \eta_h.$$

Furthermore, (2.6) allows us to define a lifted triangulation $\mathcal{T}_h^\ell(t)$ of $\Gamma(t)$ by

$$\mathcal{T}_h^\ell = \{e(t) = E^\ell(t) : E(t) \in \mathcal{T}_h(t)\}, \quad E^\ell(t) := \{p(x, t) : x \in E(t)\}. \quad (3.22)$$

This defines an exact triangulation of $\Gamma(t)$.

Lemma 3.4 (Stability of lift [8, for $q = 2$]) *Let $\eta_h : \mathcal{G}_{h,T} \rightarrow \mathbb{R}$, with lift $\eta_h^\ell : \mathcal{G}_T \rightarrow \mathbb{R}$, be such that the following quantities exist. For $1 \leq q \leq +\infty$, there exists $c_1, c_2 > 0$, independent of h , but depending on q , such that for each time $t \in [0, T]$ and each element $E(t) \in \mathcal{T}_h(t)$ with associated lifted element $e(t) \in \mathcal{T}_h^\ell(t)$, the following hold:*

$$c_1 \|\eta_h^\ell\|_{L^q(e(t))} \leq \|\eta_h\|_{L^q(E(t))} \leq c_2 \|\eta_h^\ell\|_{L^q(e(t))} \quad (3.23a)$$

$$c_1 \|\nabla_{\Gamma} \eta_h^\ell\|_{L^q(e(t))} \leq \|\nabla_{\Gamma_h} \eta_h\|_{L^q(E(t))} \leq c_2 \|\nabla_{\Gamma} \eta_h^\ell\|_{L^q(e(t))} \quad (3.23b)$$

$$\|\nabla_{\Gamma_h}^2 \eta_h\|_{L^2(E(t))} \leq c \left(\|\nabla_{\Gamma}^2 \eta_h^\ell\|_{L^2(e(t))} + h \|\nabla_{\Gamma} \eta_h^\ell\|_{L^2(e(t))} \right). \quad (3.23c)$$

This result allows us to give Sobolev embeddings for discrete surfaces:

Lemma 3.5 For $\Gamma_h(t)$ as above,

$$W^{1,q}(\Gamma_h(t)) \subset \begin{cases} L^{nq/(n-q)}(\Gamma_h(t)) & \text{for } q < n \\ L^\infty(\Gamma_h(t)) & \text{for } q > n. \end{cases} \quad (3.24)$$

Furthermore there exists a constant $c = c(n, q)$, independent of h , such that for any $\eta_h \in W^{1,q}(\Gamma_h(t))$

$$\|\eta_h\|_{L^{nq/(n-q)}(\Gamma_h(t))} \leq c \|\eta_h\|_{W^{1,q}(\Gamma_h(t))} \quad \text{for } q < n \quad (3.25a)$$

$$\|\eta_h\|_{L^\infty(\Gamma_h(t))} \leq c \|\eta_h\|_{W^{1,q}(\Gamma_h(t))} \quad \text{for } q > n. \quad (3.25b)$$

Proof To see the embedding result, we apply [Lemma 2.1](#). The bounds then follow using the stability of the lift ([Lemma 3.4](#)). \square

We will write $S_h^\ell(t)$ for the space of lifted finite element functions:

$$S_h^\ell(t) = \{\varphi_h = \phi_h^\ell : \phi_h \in S_h(t)\}.$$

This space comes with the standard approximation property:

Proposition 3.1 (Approximation property) *The Lagrangian interpolation operator $I_h : C(\Gamma(t)) \rightarrow S_h^\ell(t)$ is well defined and, for $z \in H^2(\Gamma(t))$, satisfies the bound*

$$\|z - I_h z\|_{L^2(\Gamma(t))} + h \|\nabla_\Gamma(z - I_h z)\|_{L^2(\Gamma(t))} \leq ch^2 \|z\|_{H^2(\Gamma(t))}. \quad (3.26)$$

Let $1 \leq q \leq \infty$ be such that $H^1(\Gamma(t))$ embeds into $L^q(\Gamma(t))$, then

$$\|\nabla_\Gamma(z - I_h z)\|_{L^q(\Gamma(t))} \leq ch^{1+\min(0, n/q - n/2)} \|z\|_{H^2(\Gamma(t))}. \quad (3.27)$$

Proof The proof is given in [7] for the case $q = 2$ and can be easily extended using standard interpolation theory [3, Theorem 3.1.6] to the case $q \neq 2$. \square

Remark 3.3 For the remainder of the paper, we will write lower case letters for the lift finite element functions with capital letters (i.e. $U_h^\ell = u_h$ and $W_h^\ell = w_h$) and φ_h for the lift of ϕ_h .

The motion of the edges of the simplices in the triangulation $\{\mathcal{T}_h^\ell(t)\}$ defines a discrete material velocity for the surface $\{\Gamma(t)\}$. Let $X(t)$ be the trajectory of a point on $\{\Gamma_h(t)\}$ with velocity $V_h(X(t), t)$. We set $Y(t) = p(X(t), t)$ then define v_h by

$$v_h(Y(t), t) := \dot{Y}(t) = \frac{\partial p}{\partial t}(X(t), t) + \nabla p(X(t), t) \cdot V_h(X(t), t), \quad (3.28)$$

so that for $x \in \Gamma_h(t)$, using (2.6), we have

$$v_h(p(x, t), t) = (P(x, t) - d(x, t)\mathcal{H}(x, t))V_h(x, t) - d_t(x, t)\nu(x) - d(x, t)\nu_t(x, t).$$

This defines another discrete material derivative for functions $\varphi_h(\cdot, t) \in S_h^\ell(t)$. We define the discrete material derivative on \mathcal{G}_T element-wise by

$$\partial_h^\bullet \varphi_h := \partial_t \varphi_h + v_h \cdot \nabla \varphi_h. \quad (3.29)$$

A quick calculation [12] shows that for all $\phi_h \in S_h(t)$, with lift $\varphi_h \in S_h^\ell(t)$,

$$\partial_h^\bullet \varphi_h = (\partial_h^\bullet \phi_h)^\ell. \quad (3.30)$$

It can be shown, similarly to (3.6), that $\partial_h^\bullet (\phi_j^N)^\ell = 0$. We will write $S_h^{\ell, T}$ and $\tilde{S}_h^{\ell, T}$ for the lifts of the spaces S_h^T and \tilde{S}_h^T defined by (3.12). It is clear that from Lemma 3.4 that

$$S_h^{\ell, T} \subset H^1(\mathcal{G}_T) \quad \text{and} \quad \tilde{S}_h^{\ell, T} \subset L_{H^1}^2.$$

We remark that the continuous and discrete material velocities on $\{\Gamma(t)\}$ only differ in the tangential direction. This implies that the difference between the two material derivatives on $\{\Gamma(t)\}$ only depends on the tangential gradient of the original function and not on any time derivatives.

These definitions also permit transport formulae:

Lemma 3.6 (Transport lemma for smooth triangulated surfaces [12, Lemma 4.2]) *Let $\{\Gamma(t)\}$ be an evolving surface decomposed at each time into a family curved elements $\{\mathcal{T}_h^\ell(t)\}$ whose edges evolve with velocity v_h . Then the following relations hold for functions $\eta_h, \varphi_h: \mathcal{G}_T \rightarrow \mathbb{R}$ such that the following quantities exist:*

$$\frac{d}{dt} \int_{\Gamma(t)} \eta_h \, d\sigma = \int_{\Gamma(t)} \partial_h^\bullet \eta_h + \eta_h \nabla_{\Gamma} \cdot v_h \, d\sigma, \quad (3.31)$$

and

$$\frac{d}{dt} m(\eta_h, \varphi_h) = m(\partial_h^\bullet \eta_h, \varphi_h) + m(\eta_h, \partial_h^\bullet \varphi_h) + g(v_h; \eta_h, \varphi_h) \quad (3.32)$$

$$\frac{d}{dt} a(\eta_h, \varphi_h) = a(\partial_h^\bullet \eta_h, \varphi_h) + a(\eta_h, \partial_h^\bullet \varphi_h) + b(v_h; \eta_h, \varphi_h). \quad (3.33)$$

3.4 Proof of finite element scheme well-posedness

Before showing stability of the finite element scheme, we will show a generalised Gronwall inequality:

Lemma 3.7 *Let $y_h(t), z_h(t) \geq 0$ and satisfy the following differential inequality for $\tilde{C} \geq 0$*

$$\begin{aligned} \frac{d}{dt} y_h(t) + z_h(t) &\leq c(y_h(t) + h y_h(t)^2 + h^2 y_h(t)^3) + \tilde{C} \quad \text{for } 0 \leq t \leq T, \\ y_h(0) &= y_0. \end{aligned} \quad (3.34)$$

Let h be sufficiently small so that $1 - h(y_0 + \tilde{C})^2(e^{2(1+h)ct} - 1) > 0$, then y_h, z_h satisfy the bound

$$y_h(t) + \int_0^t z_h(s) \, ds \leq \frac{e^{(1+h)ct}(y_0 + \tilde{C})}{\sqrt{1 - h(y_0 + \tilde{C})^2(e^{2(1+h)ct} - 1)}}. \quad (3.35)$$

Proof Let $\eta_h(t) = y_h(t) + \int_0^t z_h(s) \, ds + \tilde{C}$ and $\eta_0 = y_0 + \tilde{C}$. We note that $y_h(t)^q \leq \eta_h(t)^q$ ($q = 1, 2, 3$) and $\eta_h(t) \geq 0$. Then η_h satisfies

$$\frac{d}{dt}\eta_h(t) \leq c(\eta_h(t) + h\eta_h(t)^2 + h^2\eta_h(t)^3) \leq c((1+h)\eta_h(t) + (h+h^2)\eta_h(t)^3).$$

This implies

$$\frac{1}{2+2h} \frac{d}{dt} \log \left(\frac{\eta_h(t)^2}{1+h\eta_h(t)^2} \right) = \frac{\frac{d}{dt}\eta_h(t)}{(1+h)\eta_h(t) + (h+h^2)\eta_h(t)^3} \leq c.$$

Assuming that h is sufficiently small so that $1 - h\eta_0^2(e^{2(1+h)ct} - 1) > 0$, integrating this inequality in time implies

$$\eta_h(t)^2 \leq \frac{e^{2(1+h)ct}\eta_0^2}{1 - h\eta_0^2(e^{2(1+h)ct} - 1)}.$$

Rearranging this inequality gives the desired result. \square

We can now show the stability result in Theorem 3.1.

Proof (Proof of Theorem 3.1) Considering (3.17), since $\mathcal{M}(t)$ is positive definite, $\mathcal{S}(t)$ positive semi-definite and \mathcal{F} is locally Lipschitz, standard theory of ordinary differential equations gives a unique short-time solution $\alpha \in C^1([0, T_0]; \mathbb{R}^N)$ for some $T_0 < T$. From (3.7), we know $\mathcal{S}(t)$ and $\mathcal{M}(t)$ are C^1 in time, and $\mathcal{M}(t)^{-1} \in C^1$ by the Inverse Function Theorem. Thus, we infer

$$\beta(t) = \mathcal{M}(t)^{-1}\mathcal{S}(t)\mathcal{M}(t)^{-1}(\varepsilon\mathcal{S}(t)\alpha(t) + \frac{1}{\varepsilon}\mathcal{F}(\alpha(t))) \in C^1([0, T_0]; \mathbb{R}^N).$$

This is easily translated into solutions U_h, W_h in the appropriate spaces.

Since $\phi_h = 1$ is an admissible test function in (3.13a), it is clear that $\int_{\Gamma_h(t)} U_h \, d\sigma_h$ is conserved.

To extend to the long-term solution, we construct an energy bound. We start by testing (3.13a) with W_h and (3.13b) with $\partial_h^\bullet U_h$ and sum to see

$$\begin{aligned} & \varepsilon a_h(U_h, \partial_h^\bullet U_h) + \frac{1}{\varepsilon} m_h(\psi'(U_h), \partial_h^\bullet U_h) + a_h(W_h, W_h) \\ &= -\frac{d}{dt} m_h(U_h, W_h) + m_h(U_h, \partial_h^\bullet W_h) + m_h(\partial_h^\bullet U_h, W_h). \end{aligned}$$

Applying the transport formulae from [Lemma 3.2](#), we obtain

$$\begin{aligned} & \frac{d}{dt} \left(\varepsilon a_h(U_h, U_h) + \frac{1}{\varepsilon} m_h(\psi(U_h), 1) \right) + a_h(W_h, W_h) \\ &= \frac{\varepsilon}{2} b_h(V_h; U_h, U_h) + \frac{1}{\varepsilon} g_h(V_h; \psi(U_h), 1) - g_h(V_h; U_h, W_h). \end{aligned}$$

Next, we introduce the $L^2(\Gamma_h(t))$ projection $\Lambda_h: L^2(\Gamma_h(t)) \rightarrow S_h(t)$. For $z \in L^2(\Gamma_h(t))$, we define $\Lambda_h z$ as the unique solution of

$$m_h(\Lambda_h z, \phi_h) = m_h(z, \phi_h) \quad \text{for all } \phi_h \in S_h(t). \quad (3.36)$$

For $z \in H^1(\Gamma_h(t))$, we will make use of the following bounds:

$$\|\Lambda_h z\|_{H^1(\Gamma_h(t))} \leq c \|z\|_{H^1(\Gamma_h(t))}, \quad \|z - \Lambda_h z\|_{L^2(\Gamma_h(t))} \leq ch \|z\|_{H^1(\Gamma_h(t))}. \quad (3.37)$$

These bounds follow since our triangulation is quasi-uniform.

We first note that from our assumptions on v , we have $U_h(\nabla_\Gamma \cdot v)^{-\ell} \in H^1(\Gamma_h(t))$ with $\|U_h(\nabla_\Gamma \cdot v)^{-\ell}\|_{H^1(\Gamma_h(t))} \leq c \|U_h\|_{H^1(\Gamma_h(t))}$. Next, we test [\(3.13b\)](#) with $\Lambda_h(U_h(\nabla_\Gamma \cdot v)^{-\ell})$ and using [\(3.37\)](#) and the Sobolev embedding ([Lemma 3.5](#)), we see that

$$\begin{aligned} & \left| \int_{\Gamma_h(t)} W_h \Lambda_h(U_h(\nabla_\Gamma \cdot v)^{-\ell}) \, d\sigma_h \right| \\ & \leq \varepsilon |a_h(U_h, \Lambda_h(U_h(\nabla_\Gamma \cdot v)^{-\ell}))| + \frac{1}{\varepsilon} |m_h(\psi'(U_h), U_h(\nabla_\Gamma \cdot v)^\ell)| \\ & \quad + \frac{1}{\varepsilon} |m_h(\psi'(U_h), \Lambda_h(U_h(\nabla_\Gamma \cdot v)^{-\ell}) - U_h(\nabla_\Gamma \cdot v)^\ell)| \\ & \leq c \left(\varepsilon a_h(U_h, U_h) + \frac{1}{\varepsilon} m_h(\psi(U_h), 1) \right) + \frac{ch}{\varepsilon} \|U_h\|_{H^1(\Gamma_h(t))}^4. \end{aligned}$$

Similarly, testing [\(3.13b\)](#) with W_h leads to

$$m_h(W_h, W_h) \leq c\varepsilon a_h(U_h, U_h) + \frac{1}{2} a_h(W_h, W_h) + \frac{c}{\varepsilon^2} \|U_h\|_{H^1(\Gamma_h(t))}^6.$$

Applying the geometric bound [\(3.42\)](#), the two previous bounds, a Poincaré inequality and the fact that the mass of U_h is conserved, we infer that

$$\begin{aligned} & |g(V_h; U_h, W_h)| \\ & \leq |m_h(W_h, U_h(\nabla_{\Gamma_h} \cdot V_h - (\nabla_\Gamma \cdot v)^{-\ell}))| + |m_h(W_h, \Lambda_h(U_h(\nabla_\Gamma \cdot v)^{-\ell}))| \\ & \leq cm_h(U_h, U_h) + ch^2 m_h(W_h, W_h) + |m_h(W_h, \Lambda_h(U_h(\nabla_\Gamma \cdot v)^{-\ell}))| \\ & \leq c \left(\varepsilon a_h(U_h, U_h) + \frac{1}{\varepsilon} m_h(\psi(U_h), 1) \right) + \frac{1}{2} a_h(W_h, W_h) \\ & \quad + c_\varepsilon (h a_h(U_h, U_h)^2 + h^2 a_h(U_h, U_h)^3) + \tilde{C}_0(U_0), \end{aligned}$$

where $\tilde{C}_0(U_0)$ is a constant which only depends on the integral of U_0 on $\Gamma_{h,0}$. This leads to the estimate

$$\begin{aligned} & \frac{d}{dt} \left(\varepsilon a_h(U_h, U_h) + \frac{1}{\varepsilon} m_h(\psi(U_h), 1) \right) + a_h(W_h, W_h) \\ & \leq c_\varepsilon \left(\varepsilon a_h(U_h, U_h) + \frac{1}{\varepsilon} m_h(\psi(U_h), 1) + h a_h(U_h, U_h)^2 + h^2 a_h(U_h, U_h)^3 \right) \\ & \quad + \tilde{C}_0(U_0). \end{aligned}$$

We will use the generalised Gronwall inequality from (3.35) with $y_h = \varepsilon a_h(U_h, U_h) + \frac{1}{\varepsilon} m_h(\psi(U_h), 1)$, $z_h = a_h(W_h, W_h)$ and $\tilde{C} = \tilde{C}_0(U_0)$. Given T , there exists $h_1, 0 < h_1 < h_0$, such that for $h < h_1$, we have

$$1 - h(\mathcal{E}_0^h + \tilde{C}_0(U_0))^2(e^{2(1+h)ct} - 1) > 0.$$

This gives the energy bound in (3.19) with C_0 given by

$$C_0 := \frac{e^{(1+h)cT}(\mathcal{E}_0^h + \tilde{C}_0(U_0))}{\sqrt{1 - h(\mathcal{E}_0^h + \tilde{C}_0(U_0))^2(e^{2(1+h)ct} - 1)}}$$

This implies, that if $h < h_1$, we have an energy bound on $(0, T)$ and hence can turn the short-time existence result in to existence over $(0, T)$ where T is arbitrary. \square

3.5 Geometric estimates

In this section, we will simply state the following geometric estimates without proof. Details can be found in [12, Section 5] except for (3.39c) and (3.39d) which can be found in [29, Lemma 3.3.14].

Lemma 3.8 *Let μ_h denote the quotient of surface measures $d\sigma$ on $\Gamma(t)$ and $d\sigma_h$ on $\Gamma_h(t)$ such that $\mu_h d\sigma_h = d\sigma$; then*

$$\sup_{t \in [0, T]} \sup_{\Gamma_h(t)} |1 - \mu_h| \leq ch^2 \quad (3.38a)$$

$$\sup_{t \in [0, T]} \|\partial_h^\bullet \mu_h\|_{L^\infty(\Gamma_h(t))} \leq ch^2. \quad (3.38b)$$

Lemma 3.9 *Let $Z_h, \phi_h \in S_h(t)$ with lifts $z_h, \varphi_h \in S_h^\ell(t)$. Then the following estimates hold for the given bilinear forms:*

$$|m_h(Z_h, \phi_h) - m(z_h, \varphi_h)| \leq ch^2 \|Z_h\|_{L^2(\Gamma_h(t))} \|\phi_h\|_{L^2(\Gamma_h(t))} \quad (3.39a)$$

$$|a_h(Z_h, \phi_h) - a(z_h, \varphi_h)| \leq ch^2 \|\nabla_{\Gamma_h} Z_h\|_{L^2(\Gamma_h(t))} \|\nabla_{\Gamma_h} \phi_h\|_{L^2(\Gamma_h(t))} \quad (3.39b)$$

$$|g_h(V_h; Z_h, \phi_h) - g(v_h; z_h, \varphi_h)| \leq ch^2 \|Z_h\|_{L^2(\Gamma_h(t))} \|\phi_h\|_{L^2(\Gamma_h(t))} \quad (3.39c)$$

$$|b_h(V_h; Z_h, \phi_h) - b(v_h; z_h, \varphi_h)| \leq ch^2 \|\nabla_{\Gamma_h} Z_h\|_{L^2(\Gamma_h(t))} \|\nabla_{\Gamma_h} \phi_h\|_{L^2(\Gamma_h(t))}. \quad (3.39d)$$

Using the same reasoning, it is also clear that

$$|m_h(\psi'(Z_h), \phi_h) - m(\psi'(z_h), \varphi_h)| \leq ch^2 \|\psi'(Z_h)\|_{L^2(\Gamma_h(t))} \|\phi_h\|_{L^2(\Gamma_h(t))}. \quad (3.40)$$

Similar results apply if the first argument is the material derivative of a finite element function:

Lemma 3.10 *For $Z_h \in S_h^T$, $\phi_h \in \tilde{S}_h^T$ with lifts $z_h, \varphi_h \in S_h^\ell(t)$ for each time, we have*

$$|m_h(\partial_h^\bullet Z_h, \varphi_h) - m(\partial_h^\bullet z_h, \phi_h)| \leq ch^2 \|\partial_h^\bullet Z_h\|_{L^2(\Gamma_h(t))} \|\phi_h\|_{L^2(\Gamma_h(t))} \quad (3.41a)$$

$$|a_h(\partial_h^\bullet Z_h, \varphi_h) - a(\partial_h^\bullet z_h, \phi_h)| \leq ch^2 \|\nabla_{\Gamma_h}(\partial_h^\bullet Z_h)\|_{L^2(\Gamma_h(t))} \|\nabla_{\Gamma_h} \phi_h\|_{L^2(\Gamma_h(t))}. \quad (3.41b)$$

The next lemma bounds errors from the approximation of v by v_h :

Lemma 3.11 *The difference between the continuous velocity v and the discrete velocity v_h on $\Gamma(t)$ can be estimated by*

$$|v - v_h| + h |\nabla_{\Gamma}(v - v_h)| \leq ch^2 \|v\|_{C^2(\mathcal{N}_T)} < ch^2. \quad (3.42)$$

This allows us to bound the error between the material derivatives on $\Gamma(t)$:

Corollary 3.1 *Suppose that $\eta: \mathcal{G}_T \rightarrow \mathbb{R}$ and $\partial^\bullet \eta$ and $\partial_h^\bullet \eta$ exist. For $\eta \in H^1(\Gamma(t))$, we have the estimate*

$$\|\partial^\bullet \eta - \partial_h^\bullet \eta\|_{L^2(\Gamma(t))} \leq ch^2 \|\nabla_{\Gamma} \eta\|_{L^2(\Gamma(t))}, \quad (3.43)$$

and for $\eta \in H^2(\Gamma(t))$, we obtain

$$\|\nabla_{\Gamma}(\partial^\bullet \eta - \partial_h^\bullet \eta)\|_{L^2(\Gamma(t))} \leq ch^2 \|\eta\|_{H^2(\Gamma(t))}. \quad (3.44)$$

3.6 Ritz projection

We conclude this section by constructing a discrete projection operator, similar to an interpolation operator. We define the Ritz projection operator, $\Pi_h z \in S_h(t)$, of $z \in H^1(\Gamma(t))$ as the unique solution of

$$a_h(\Pi_h z, \phi_h) = a(z, \varphi_h) \quad \text{for all } \phi_h \in S_h(t), \text{ with lift } \varphi_h \in S_h^\ell(t) \quad (3.45)$$

and

$$\int_{\Gamma_h(t)} \Pi_h z \, d\sigma_h = \int_{\Gamma(t)} z \, d\sigma.$$

We will write $\pi_h z = (\Pi_h z)^\ell$ for the lift of the Ritz projection.

Remark 3.4 This operator is the Ritz projection used by [6], but different to that used in other surface finite element analyses such as [9, 12], which use the operator $\mathcal{R}_h: H^2(\Gamma(t)) \rightarrow S_h^\ell(t)$ given as the unique solution of

$$a(\mathcal{R}_h z, \varphi_h) = a(z, \varphi_h) \quad \text{for all } \varphi_h \in S_h^\ell(t) \quad \text{and} \quad \int_{\Gamma(t)} \mathcal{R}_h z \, d\sigma = 0.$$

The following bounds are immediate:

Theorem 3.2 For $z \in H^1(\Gamma(t))$,

$$\|\pi_h z\|_{H^1(\Gamma(t))} \leq c \|z\|_{H^1(\Gamma(t))}, \quad \|\pi_h z - z\|_{L^2(\Gamma(t))} \leq ch \|z\|_{H^1(\Gamma(t))}. \quad (3.46)$$

For $z \in H^2(\Gamma(t))$,

$$\|\pi_h z - z\|_{L^2(\Gamma(t))} + h \|\nabla_\Gamma(\pi_h z - z)\|_{L^2(\Gamma(t))} \leq ch^2 \|z\|_{H^2(\Gamma(t))}, \quad (3.47)$$

and for $1 \leq q \leq \infty$, such that $H^1(\Gamma(t))$ embeds into $L^q(\Gamma(t))$,

$$\|\nabla_\Gamma(\pi_h z - z)\|_{L^q(\Gamma(t))} \leq ch^{1+\min(0, n/q-n/2)} \|z\|_{H^2(\Gamma(t))}. \quad (3.48)$$

Proof The H^1 stability result is clear and the L^2 error bound for a H^1 function follows from an Aubin-Nitsche trick. The L^2 results for $z \in H^2(\Gamma(t))$ follow from standard error estimates for the surface finite element [7]. The L^q result follows from the same splitting argument along with an inverse inequality. \square

Corollary 3.2 The Ritz projection is bounded in L^∞ and we have the bound

$$\|I_h z\|_{L^\infty(\Gamma_h(t))} \leq \|\pi_h z\|_{L^\infty(\Gamma_h(t))} \leq c \|z\|_{H^2(\Gamma(t))}. \quad (3.49)$$

Proof Let $1 < q < \infty$ be such that $H^1(\Gamma(t))$ embeds into $L^q(\Gamma(t))$ and such that $W^{1,q}(\Gamma(t))$ embeds into $L^\infty(\Gamma(t))$. The previous result, a Poincaré inequality and (3.38a) imply that

$$\begin{aligned} \|\pi_h z - z\|_{W^{1,q}(\Gamma(t))} &\leq c \|\nabla_\Gamma(\pi_h z - z)\|_{L^q(\Gamma(t))} + \left| \int_{\Gamma(t)} \pi_h z - z \, d\sigma \right| \\ &\leq c \|z\|_{H^2(\Gamma(t))}, \end{aligned}$$

for h sufficiently small. We use a Sobolev embedding (Lemma 2.1), to see

$$\|\pi_h z\|_{L^\infty(\Gamma(t))} \leq c \|\pi_h z\|_{W^{1,q}(\Gamma(t))} \leq c \|z\|_{H^2(\Gamma(t))}.$$

It is clear that

$$\|I_h z\|_{L^\infty(\Gamma_h(t))} = \|\pi_h z\|_{L^\infty(\Gamma(t))},$$

which completes the proof. \square

Since $\partial_h^\bullet \Pi_h z \neq \Pi_h \partial_h^\bullet z$, we also wish to have a bound on the discrete material derivative of this error for a function. We will assume that $z \in H^2(\Gamma(t))$ and $\partial^\bullet z \in H^2(\Gamma(t))$ for each t . Under this assumption, we may take a time derivative of (3.45), so that for all $\phi_h \in S_h^T$ with lift $\varphi_h \in S_h^{\ell, T}$,

$$a_h(\partial_h^\bullet \Pi_h z, \phi_h) = a(\partial_h^\bullet z, \varphi_h) + (b(v_h; z, \varphi_h) - b_h(V_h; \Pi_h z, \phi_h)). \quad (3.50)$$

In fact using similar arguments to Lemma 2.3, we can construct a similar extension of a finite element function $\phi_h \in S_h(t)$ to a function $\tilde{\phi}_h \in S_h^T$ by

$$\tilde{\phi}_h(x, s) = \sum_{j=1}^N \gamma_j \phi_j^N(x, s) \quad \text{for } (x, s) \in \mathcal{G}_{h, T} \quad \text{where } \phi_h(x) = \sum_{j=1}^N \gamma_j \phi_j^N(x, t).$$

Hence, we deduce that (3.50) applies at each time $t \in (0, T)$ for $\phi_h \in S_h(t)$.

We start by proving two technical lemmas:

Lemma 3.12 *Given $z: \mathcal{G}_T \rightarrow \mathbb{R}$ with $z \in H^2(\Gamma(t))$ and $\partial^\bullet z \in H^2(\Gamma(t))$ for almost every time $t \in (0, T)$, then $\partial_h^\bullet \Pi_h z$ exists and we have the bound*

$$\|\nabla_{\Gamma_h}(\partial_h^\bullet \Pi_h z)\|_{L^2(\Gamma_h(t))} \leq c(\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))}). \quad (3.51)$$

Proof To show the bound, we start from (3.50), using a Young's inequality, (3.47) and (3.39d) gives

$$a_h(\partial_h^\bullet \Pi_h z, \phi_h) \leq c(\|\nabla_{\Gamma} \partial^\bullet z\|_{L^2(\Gamma(t))}^2 + \|z\|_{H^2(\Gamma(t))}^2) + \frac{1}{2} \|\nabla_{\Gamma_h} \phi_h\|_{L^2(\Gamma_h(t))}^2.$$

Applying this bound with $\phi_h = \partial_h^\bullet \Pi_h z$ gives the estimate (3.51). \square

Lemma 3.13 *Define the function T_h on $S_h^\ell(t)$ by*

$$T_h(\varphi_h) := a(\partial_h^\bullet(\pi_h z - z), \varphi_h). \quad (3.52)$$

Then we have the bound

$$|T_h(\varphi_h)| \leq ch \left(\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))} \right) \|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))}. \quad (3.53)$$

Furthermore, for any $\eta \in H^2(\Gamma(t))$, we have that

$$\begin{aligned} |T_h(\varphi_h)| &\leq ch \|z\|_{H^2(\Gamma(t))} \|\nabla_{\Gamma}(\varphi_h - \eta)\|_{L^2(\Gamma(t))} + ch^2 \|z\|_{H^2(\Gamma(t))} \|\eta\|_{H^2(\Gamma(t))} \\ &\quad + ch^2 \left(\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))} \right) \|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))}. \end{aligned} \quad (3.54)$$

Proof Using (3.45) and (3.50), we see for $\phi_h \in S_h(t)$, with lift $\varphi_h \in S_h^\ell(t)$,

$$\begin{aligned} T_h(\varphi_h) &= a(\partial_h^\bullet \pi_h z, \varphi_h) - a(\partial_h^\bullet z, \varphi_h) \\ &= b(v_h; z - \pi_h z, \varphi_h) + (a(\partial_h^\bullet \pi_h z, \varphi_h) - a_h(\partial_h^\bullet \Pi_h z, \phi_h)) \\ &\quad + (b(v_h; \pi_h z, \varphi_h) - b_h(V_h; \Pi_h z, \phi_h)). \end{aligned}$$

Using our bound on the Ritz projection (3.47), and two geometric estimates (3.39d) and (3.41b), we have that

$$\begin{aligned} |T_h(\varphi_h)| &\leq ch \|z\|_{H^2(\Gamma(t))} \|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))} \\ &\quad + ch^2 \left(\|\nabla_{\Gamma} \partial_h^{\bullet} \pi_h z\|_{L^2(\Gamma(t))} + \|\nabla_{\Gamma} \pi_h z\|_{L^2(\Gamma(t))} \right) \|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))} \\ &\leq ch \left(\|z\|_{H^2(\Gamma(t))} + \|\partial^{\bullet} z\|_{H^2(\Gamma(t))} \right) \|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))}. \end{aligned}$$

We can improve this estimate by comparing v_h to the smooth velocity v and introducing a smooth function $\eta \in H^2(\Gamma(t))$. Then, we split the first term in $T_h(\varphi_h)$ into

$$\begin{aligned} b(v_h; \pi_h z - z, \varphi_h) \\ = b(v_h - v; \pi_h z - z, \varphi_h) + b(v; \pi_h z - z, \varphi_h - \eta) + b(v; \pi_h z - z, \eta). \end{aligned}$$

Using the smoothness of η , the final term, $b(v; \pi_h z - z, \eta)$, is bounded using an integration by parts argument given by [12, p. 21]:

$$b(v; \varphi, \eta) = \int_{\Gamma(t)} \sum_{i,j=1}^{n+1} H\nu_j \mathcal{B}(v)_{ij} \varphi \underline{D}_i \eta \, d\sigma - \int_{\Gamma(t)} \varphi \sum_{i,j=1}^{n+1} \underline{D}_j (\mathcal{B}(v)_{ij} \underline{D}_i \eta) \, d\sigma.$$

Hence, we obtain

$$|b(v; \varphi, \eta)| \leq c \|\varphi\|_{L^2(\Gamma(t))} \|\eta\|_{H^2(\Gamma(t))}.$$

Combining these calculations with (3.42) and (3.47), we get

$$\begin{aligned} |b(v_h; \pi_h z - z, \varphi_h)| &\leq ch \|z\|_{H^2(\Gamma(t))} \|\nabla_{\Gamma}(\varphi_h - \eta)\|_{L^2(\Gamma(t))} \\ &\quad + ch^2 \|z\|_{H^2(\Gamma(t))} \left(\|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))} + \|\eta\|_{H^2(\Gamma(t))} \right). \end{aligned}$$

Hence, we have

$$\begin{aligned} |T_h(\varphi_h)| &\leq ch \|z\|_{H^2(\Gamma(t))} \|\nabla_{\Gamma}(\varphi_h - \eta)\|_{L^2(\Gamma(t))} + ch^2 \|z\|_{H^2(\Gamma(t))} \|\eta\|_{H^2(\Gamma(t))} \\ &\quad + ch^2 \left(\|z\|_{H^2(\Gamma(t))} + \|\partial^{\bullet} z\|_{H^2(\Gamma(t))} \right) \|\nabla_{\Gamma} \varphi_h\|_{L^2(\Gamma(t))}, \end{aligned}$$

which is the second estimate. \square

These results allow us to show an estimate for the difference between the material derivative of a function and its Ritz projection.

Lemma 3.14 *For $z: \mathcal{G}_T \rightarrow \mathbb{R}$ with $z, \partial^{\bullet} z \in H^2(\Gamma(t))$, we have*

$$\begin{aligned} \|\partial_h^{\bullet}(\pi_h z - z)\|_{L^2(\Gamma(t))} + h \|\nabla_{\Gamma} \partial_h^{\bullet}(\pi_h z - z)\|_{L^2(\Gamma(t))} \\ \leq ch^2 \left(\|z\|_{H^2(\Gamma(t))} + \|\partial^{\bullet} z\|_{H^2(\Gamma(t))} \right). \end{aligned} \tag{3.55}$$

Proof We start by rewriting the error as

$$\begin{aligned} & a(\partial_h^\bullet(\pi_h z - z), \partial_h^\bullet(\pi_h z - z)) \\ &= a(\partial_h^\bullet(\pi_h z - z), \partial_h^\bullet \pi_h z - I_h(\partial^\bullet z)) + a(\partial_h^\bullet(\pi_h z - z), I_h(\partial^\bullet z) - \partial^\bullet z) \\ & \quad + a(\partial_h^\bullet(\pi_h z - z), \partial^\bullet z - \partial_h^\bullet z). \end{aligned} \tag{3.56}$$

We can bound the first term on the right-hand side using (3.53) by

$$\begin{aligned} & |a(\partial_h^\bullet(\pi_h z - z), \partial_h^\bullet \pi_h z - I_h(\partial^\bullet z))| = |T_h(\partial_h^\bullet \pi_h z - I_h(\partial^\bullet z))| \\ & \leq ch(\|\partial^\bullet z\|_{H^2(\Gamma(t))} + \|z\|_{H^2(\Gamma(t))}) \|\nabla_\Gamma(\partial_h^\bullet \pi_h z - I_h(\partial^\bullet z))\|_{L^2(\Gamma(t))} \\ & \leq ch(\|\partial^\bullet z\|_{H^2(\Gamma(t))} + \|z\|_{H^2(\Gamma(t))}) \|\nabla_\Gamma \partial_h^\bullet(\pi_h z - z)\|_{L^2(\Gamma(t))} \\ & \quad + ch \|\nabla_\Gamma \partial_h^\bullet(\pi_h z - z)\|_{L^2(\Gamma(t))}^2. \end{aligned}$$

The second term is bounded using the approximation property (3.26):

$$|a(\partial_h^\bullet(\pi_h z - z), I_h(\partial^\bullet z) - \partial^\bullet z)| \leq ch \|\nabla_\Gamma \partial_h^\bullet(\pi_h z - z)\|_{L^2(\Gamma(t))} \|\partial^\bullet z\|_{H^2(\Gamma(t))}.$$

Finally, we use our estimate of the difference of material derivatives (3.43) to bound the third term:

$$|a(\partial_h^\bullet(\pi_h z - z), \partial^\bullet z - \partial_h^\bullet z)| \leq ch^2 \|\nabla_\Gamma \partial_h^\bullet(\pi_h z - z)\|_{L^2(\Gamma(t))} \|z\|_{H^2(\Gamma(t))}.$$

Combining these three bounds in (3.56), we get the desired gradient norm bound for h sufficiently small.

To show the L^2 bound, we use the Aubin-Nitsche trick. We start by writing $e = \partial_h^\bullet(\pi_h z - z)$, then e is in L^2 so can be set as the right-hand side for the dual problem: Find $\zeta \in H^1(\Gamma(t))$ such that

$$a(\varphi, \zeta) = m(e - c_0, \varphi) \quad \text{for all } \varphi \in H^1(\Gamma(t)), \text{ and } \int_{\Gamma(t)} \zeta \, d\sigma = 0, \tag{3.57}$$

where $c_0 = \frac{1}{|\Gamma(t)|} \int_{\Gamma(t)} e \, d\sigma$. We know [1] that (3.57) has a unique solution and satisfies the regularity result

$$\|\zeta\|_{H^2(\Gamma(t))} \leq c \|e\|_{L^2(\Gamma(t))}. \tag{3.58}$$

We note that from $\int_{\Gamma_h(t)} \Pi_h z \, d\sigma_h = \int_{\Gamma(t)} z \, d\sigma$, that

$$\begin{aligned} |\Gamma(t)| |c_0| &= \int_{\Gamma(t)} \partial_h^\bullet(\pi_h z - z) \, d\sigma \\ &= \frac{d}{dt} \int_{\Gamma(t)} \pi_h z - z \, d\sigma - \int_{\Gamma(t)} (\pi_h z - z) \nabla_\Gamma \cdot v_h \, d\sigma. \end{aligned}$$

We remark that from (3.38a) and (3.38b), we have

$$\begin{aligned} \frac{d}{dt} \int_{\Gamma(t)} \pi_h z - z \, d\sigma &= \frac{d}{dt} \left(\int_{\Gamma(t)} \pi_h z \left(1 - \frac{1}{\mu_h}\right) \, d\sigma \right) \\ &\leq ch^2 (\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))}), \end{aligned}$$

and using (3.47), we infer

$$\int_{\Gamma(t)} (\pi_h z - z) \nabla_{\Gamma} \cdot v_h \, d\sigma \leq c \|\pi_h z - z\|_{L^2(\Gamma(t))} \leq ch^2 \|z\|_{H^2(\Gamma(t))}.$$

This implies

$$|c_0| \leq ch^2 (\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))}).$$

These calculations lead to

$$m(e, e) - |\Gamma(t)|^2 c_0^2 = a(\zeta, e) = a(\zeta - I_h \zeta, e) + T_h(I_h \zeta). \quad (3.59)$$

The first term on the right-hand side is bounded using the approximation property (3.26) and the gradient norm bound on e , together with the dual regularity result (3.58):

$$|a(\zeta - I_h, e)| \leq ch^2 \|e\|_{L^2(\Gamma(t))} (\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))}).$$

The second term is estimated using the improved bound (3.54) on $T_h(I_h \zeta)$ with $\eta = \zeta$. Applying the approximation (3.26) we see

$$|T_h(I_h \zeta)| \leq ch^2 (\|z\|_{H^2(\Gamma(t))} + \|\partial^\bullet z\|_{H^2(\Gamma(t))}) \|e\|_{L^2(\Gamma(t))}.$$

Applying these two bounds in (3.59) gives the desired result. \square

4 Well-posedness of the continuous problem

We use this section to show some properties of the continuous scheme based on the energy estimates coming from Theorem 3.1 along with further some estimates. We will use these properties in later sections but they are also important results in their own right.

4.1 Improved bounds on the finite element scheme

In order to derive some improved bounds on $\partial_h^\bullet U_h$ and W_h , we will assume that $U_{h,0} = \Pi_h u_0$ with $u_0 \in H^2(\Gamma_0)$. It is clear that assumption (3.11) still holds in this case. In fact, we will make use of the bound

$$\mathcal{E}_0^h + \tilde{C}_0(U_0) \leq c_\varepsilon (1 + \|u_0\|_{H^2(\Gamma_0)}^2 + \|u_0\|_{H^2(\Gamma_0)}^4) + \tilde{\tilde{C}}_1(u_0) =: \tilde{C}_1(u_0), \quad (4.1)$$

where

$$\tilde{\tilde{C}}_1(u_0) := c_\varepsilon \left(\left| \int_{\Gamma_0} u_0 \, d\sigma \right| + \left| \int_{\Gamma_0} u_0 \, d\sigma \right|^2 + \left| \int_{\Gamma_0} u_0 \, d\sigma \right|^3 \right).$$

This implies the constant $C_0(U_0)$ from Theorem 3.1 can be bounded by

$$C_0(U_0) \leq \max \left\{ \frac{\exp(cT)\tilde{C}_1(u_0)}{\sqrt{1 - (\exp(2cT) - 1)(h\tilde{C}_1(u_0) + h^2\tilde{C}_1(u_0)^2)}}, \tilde{C}_1(u_0) \right\} \\ =: C_1(u_0).$$

This is not essential for well-posedness of the finite element method but will be used for the well-posedness results for the continuous problem.

First, we need a bound on $W_h|_{t=0}$:

Lemma 4.1 *Under the assumption that $u_0 \in H^2(\Gamma_0)$, the following bound holds for $W_h|_{t=0}$:*

$$\|W_h(0, \cdot)\|_{L^2(\Gamma_h(0))} \leq c_\varepsilon (\|u_0\|_{H^2(\Gamma_0)} + \|u_0\|_{H^2(\Gamma_0)}^3). \quad (4.2)$$

Proof Since α, β are $C^1([0, T]; \mathbb{R}^N)$ in time (Theorem 3.1), we know that (3.13b) holds at time $t = 0$. We see that from the choice $U_{h,0} = \Pi_h u_0$, using Green's formula (2.3), we have

$$a_h(U_{h,0}, W_h(0, \cdot)) = a(u_0, w_h(0, \cdot)) = -m(\Delta_\Gamma u_0, w_h(0, \cdot)).$$

This implies that

$$m_h(W_h(0, \cdot), W_h(0, \cdot)) = \varepsilon a_h(U_{h,0}, W_h(0, \cdot)) + \frac{1}{\varepsilon} m_h(\psi'(U_{h,0}), W_h(0, \cdot)) \\ \leq c_\varepsilon (\|u_0\|_{H^2(\Gamma_0)} + \|u_0\|_{H^2(\Gamma_0)}^3) \|W_h(0, \cdot)\|_{L^2(\Gamma_h(0))}.$$

In the last line we have used (3.23) and the Sobolev embedding of $H^1(\Gamma(t)) \hookrightarrow L^6(\Gamma(t))$ (Lemma 2.1). \square

From Theorem 3.1, we see that $\beta \in C^1([0, T], \mathbb{R}^N)$ so $\partial_h^\bullet W_h$ exists. Hence, we may take the time derivative of (3.13b) to see, for $\phi_h \in S_h^T$,

$$\varepsilon (a_h(\partial_h^\bullet U_h, \phi_h) + b_h(V_h; U_h, \phi_h)) \\ + \frac{1}{\varepsilon} (m_h(\psi''(U_h)\partial_h^\bullet U_h, \phi_h) + g_h(V_h; \psi'(U_h), \phi_h)) \\ - (m_h(\partial_h^\bullet W_h, \phi_h) + g_h(V_h; W_h, \phi_h)) = 0. \quad (4.3)$$

Lemma 4.2 *Under the assumption that $u_0 \in H^2(\Gamma_0)$, we have the bound*

$$\varepsilon \int_0^T \|\partial_h^\bullet U_h\|_{L^2(\Gamma_h(t))}^2 dt + \sup_{t \in (0, T)} \|W_h\|_{L^2(\Gamma_h(t))}^2 \leq C_2(u_0). \quad (4.4)$$

with $C_2(u_0)$ given by

$$C_2(u_0) := c_\varepsilon (\|u_0\|_{H^2(\Gamma_0)} + \|u_0\|_{H^2(\Gamma_0)}^3 + C_1(u_0) + C_1(u_0)^2).$$

Proof We start by subtracting (4.3) tested with W_h from (3.15) tested with $\varepsilon \partial_h^\bullet U_h$ and use the transport formula (3.8) to arrive at

$$\begin{aligned} & \varepsilon m_h(\partial_h^\bullet U_h, \partial_h^\bullet U_h) + \frac{1}{2} \frac{d}{dt} m_h(W_h, W_h) \\ &= -\varepsilon (g_h(V_h; \partial_h^\bullet U_h, U_h) + b_h(V_h; U_h, W_h)) \\ & \quad + \frac{1}{\varepsilon} (m_h(\psi''(U_h) \partial_h^\bullet U_h, W_h) + g_h(V_h; \psi'(U_h), W_h)) - \frac{1}{2} g_h(V_h; W_h, W_h). \end{aligned} \quad (4.5)$$

Note that using a Hölder inequality, Young's inequality with ε , and the Sobolev embedding (Lemma 3.5) we have

$$\begin{aligned} & |m_h(\psi''(U_h) \partial_h^\bullet U_h, W_h)| \\ &= \left| \int_{\Gamma_h(t)} \psi''(U_h) \partial_h^\bullet U_h W_h \, d\sigma_h \right| = \left| \int_{\Gamma_h(t)} (3U_h^2 - 1) \partial_h^\bullet U_h W_h \, d\sigma_h \right| \\ &\leq \frac{\varepsilon}{4} \|\partial_h^\bullet U_h\|_{L^2(\Gamma_h(t))}^2 + c_\varepsilon (\|U_h\|_{H^1(\Gamma_h(t))}^2 \|W_h\|_{H^1(\Gamma_h(t))}^2 + \|W_h\|_{L^2(\Gamma_h(t))}^2). \end{aligned}$$

Applying this estimate in (4.5), we have

$$\begin{aligned} & \varepsilon m_h(\partial_h^\bullet U_h, \partial_h^\bullet U_h) + \frac{d}{dt} m_h(W_h, W_h) \\ &\leq c_\varepsilon \left(\|U_h\|_{H^1(\Gamma_h(t))}^2 + \|W_h\|_{H^1(\Gamma_h(t))}^2 + \|U_h\|_{H^1(\Gamma_h(t))} \|W_h\|_{H^1(\Gamma_h(t))}^2 \right). \end{aligned}$$

Integrating in time using a Gronwall inequality gives us

$$\begin{aligned} & \varepsilon \int_0^T \|\partial_h^\bullet U_h\|_{L^2(\Gamma_h(t))}^2 \, dt + \sup_{t \in (0, T)} \|W_h\|_{L^2(\Gamma_h(t))}^2 \\ &\leq \|W_h(\cdot, 0)\|_{L^2(\Gamma_h(t))}^2 + c_\varepsilon \int_0^T \left(\|U_h\|_{H^1(\Gamma_h(t))}^2 + \|W_h\|_{H^1(\Gamma_h(t))}^2 \right) \, dt \\ & \quad + c_\varepsilon \sup_{t \in (0, T)} \|U_h\|_{H^1(\Gamma_h(t))}^2 \int_0^T \|W_h\|_{H^1(\Gamma_h(t))}^2 \, dt. \end{aligned}$$

Applying the bounds from Theorem 3.1, Lemma 4.1 and (4.1) completes the proof. \square

4.2 Existence

The idea of the existence proof is to show that the lift of the solutions to finite element scheme (3.13) converges, along a subsequence, to a solution of the continuous equations.

We suppose that $u_0 \in H^2(\Gamma_0)$ is a given function. In this section, we will take $U_{h,0} = \Pi_h u_0$ with Π_h the Ritz projection defined in (3.45). Since the Ritz projection is stable in H^1 , the stability bound in Theorem 3.1 holds independently of h . Furthermore, the stability bounds from Lemma 3.4 imply

we may transform this bound to $\{\Gamma(t)\}$ and bound the lifts $u_h = U_h^\ell$ and $w_h = W_h^\ell$ by

$$\sup_{t \in (0, T)} \int_{\Gamma(t)} \frac{\varepsilon}{2} |\nabla_\Gamma u_h|^2 + \frac{1}{\varepsilon} \psi(u_h) \, d\sigma + \int_0^T \|w_h\|_{H^1(\Gamma(t))}^2 \, dt \leq C_1(u_0).$$

Our assumption that $u_0 \in H^2(\Gamma_0)$ allows the use of the improved bounds in [Lemma 4.2](#). Using similar lifting arguments we have

$$\int_0^T \|\partial^\bullet u_h\|_{L^2(\Gamma(t))}^2 \, dt \leq C_2(u_0).$$

These bounds, along with the conservation of mass property [\(3.18\)](#), imply that u_h is uniformly bounded in $L_{H^1}^\infty \cap H^1(\mathcal{G}_T)$ and w_h in $L_{H^1}^2$. Hence, we may extract subsequences (for which we will still use the subscript h), and functions \bar{u} and \bar{w} with $\bar{u} \in L_{H^1}^\infty \cap H^1(\mathcal{G}_T)$, and $\bar{w} \in L_{H^1}^2$ such that

$$u_h \rightharpoonup \bar{u} \quad \text{weakly in } H^1(\mathcal{G}_T) \quad w_h \rightharpoonup \bar{w} \quad \text{weakly in } L_{H^1}^2. \quad (4.6)$$

We remark that these results imply $\partial^\bullet \bar{u} \in L_{L^2}^2$ and $\partial^\bullet u_h \rightharpoonup \partial^\bullet \bar{u}$ weakly in $L_{L^2}^2$. Furthermore, from the compactness result ([Proposition 2.1](#)) we infer that we may take a further subsequence (still denoted u_h) such that

$$u_h \rightarrow \bar{u} \quad \text{almost everywhere in } \mathcal{G}_T.$$

Using a Dominated Convergence Theorem-type argument [[30](#), Lemma 8.3], since $\|\psi'(u_h)\|_{L^2(\mathcal{G}_T)} \leq c \|u_h\|_{H^1(\mathcal{G}_T)}^3$ is bounded independently of h , we infer that

$$\psi'(u_h) \rightharpoonup \psi'(\bar{u}) \quad \text{weakly in } L_{L^2}^2. \quad (4.7)$$

We will show that \bar{u} and \bar{w} satisfy [\(2.19\)](#). For $\varphi \in L_{H^1}^2$, we write $\phi_h = \Pi_h \varphi$, where Π_h is the Ritz-projection [\(3.45\)](#), and $\varphi_h = \phi_h^\ell = \pi_h \varphi$.

Using [\(3.13\)](#), we have

$$\begin{aligned} & m(\partial^\bullet \bar{u}, \varphi) + g(v; \bar{u}, \varphi) + a(w, \varphi) \\ &= (m(\partial^\bullet \bar{u}, \varphi) - m_h(\partial_h^\bullet U_h, \phi_h)) + (g(v; \bar{u}, \varphi) - g_h(V_h; U_h, \phi_h)) \\ & \quad + (a(w, \varphi) - a_h(W_h, \phi_h)). \end{aligned} \quad (4.8)$$

and

$$\begin{aligned} & \varepsilon a(\bar{u}, \varphi) + \frac{1}{\varepsilon} m(\psi'(\bar{u}), \varphi) - m(\bar{w}, \varphi) \\ &= \varepsilon (a(\bar{u}, \varphi) - a_h(U_h, \phi_h)) + \frac{1}{\varepsilon} (m(\psi'(\bar{u}), \varphi) - m_h(\psi'(U_h), \phi_h)) \\ & \quad - (m(\bar{w}, \varphi) - m_h(W_h, \phi_h)). \end{aligned} \quad (4.9)$$

We may use the geometric estimates shown in [Section 3.5](#) and the bounds on the Ritz projection from [Section 3.6](#) to see to bound the terms on the right-hand sides of these equations. We will denote by $c(h)$ a generic constant

depending on h , which may also depend on ε , such that $c(h) \rightarrow 0$ as $h \rightarrow 0$. Integrating in time, this implies

$$\begin{aligned} & \int_0^T m(\partial^\bullet \bar{u}, \varphi) + g(v; \bar{u}, \varphi) + a(\bar{w}, \varphi) dt \\ & \leq \int_0^T m(\partial^\bullet \bar{u} - \partial^\bullet u_h, \varphi) + g(v; \bar{u} - u_h, \varphi) + a(\bar{w} - w_h, \varphi) dt \\ & \quad + c(h) \int_0^T (\|\partial_h^\bullet u_h\|_{L^2(\Gamma(t))} + \|u_h\|_{H^1(\Gamma(t))} + \|\nabla_{\Gamma} w_h\|_{L^2(\Gamma(t))}) \|\varphi\|_{H^1(\Gamma(t))} dt, \end{aligned}$$

and

$$\begin{aligned} & \int_0^T \varepsilon a(\bar{u}, \varphi) + \frac{1}{\varepsilon} m(\psi'(\bar{u}), \varphi) - m(\bar{w}, \varphi) dt \\ & \leq \int_0^T \varepsilon a(\bar{u} - u_h, \varphi) + \frac{1}{\varepsilon} m(\psi'(\bar{u}) - \psi'(u_h), \varphi) - m(\bar{w} - w_h, \varphi) dt \\ & \quad + c(h) \int_0^T (\|u_h\|_{H^1(\Gamma(t))} + \|w_h\|_{L^2(\Gamma(t))}) \|\varphi\|_{H^1(\Gamma(t))} dt. \end{aligned}$$

We may send $h \rightarrow 0$ in the right-hand sides of both previous equations, and use the convergence results (4.6) and (4.7), so that for all $\varphi \in L^2_{H^1}$ we arrive at

$$\begin{aligned} & \int_0^T m(\partial^\bullet \bar{u}, \varphi) + g(v; \bar{u}, \varphi) + a(\bar{w}, \varphi) dt = 0 \\ & \int_0^T \varepsilon a(\bar{u}, \varphi) + \frac{1}{\varepsilon} m(\psi'(\bar{u}), \varphi) - m(\bar{w}, \varphi) dt = 0. \end{aligned}$$

Finally, we use Lemma 2.5 to transform this equality into a almost everywhere in time equality so that the pair \bar{u}, \bar{w} satisfy (2.19).

To show that \bar{u} achieves the initial condition, we start by choosing $\varphi \in C^2(\bar{\mathcal{G}}_T)$ and continue with the notation $\varphi_h = \pi_h \varphi$. Using the discrete transport formula (3.32), the lift of the finite element solution u_h satisfies

$$\int_0^T m(u_h, \varphi_h) \dot{\alpha} dt = - \int_0^T (m(\partial_h^\bullet u_h, \varphi_h) + m(u_h, \partial_h^\bullet \varphi_h) + g(v_h; u_h, \varphi_h)) \alpha dt,$$

for all $\alpha \in C_c^\infty(0, T)$. Using similar limiting arguments as above, with the addition of (3.55), we obtain the identity

$$\int_0^T m(\bar{u}, \varphi) \dot{\alpha} dt = - \int_0^T (m(\partial^\bullet \bar{u}, \varphi) + m(\bar{u}, \partial^\bullet \varphi) + g(v; \bar{u}, \varphi)) \alpha dt.$$

In fact, by density of $C^2(\mathcal{G}_T)$ functions in $H^1(\mathcal{G}_T)$ [25, Theorem 2.4], we see that this equality holds for all $\varphi \in H^1(\mathcal{G}_T)$. This implies that $m(\bar{u}, \varphi)$ is weakly differentiable as a function on $(0, T)$ with weak derivative $m(\partial^\bullet \bar{u}, \varphi) + m(\bar{u}, \partial^\bullet \varphi) + g(v; \bar{u}, \varphi)$. Since $\bar{u}, \varphi \in H^1(\mathcal{G}_T)$, this weak derivative is a function

in $L^1(0, T)$, and hence we infer that $m(\bar{u}, \varphi)$ is absolutely continuous on $[0, T]$ [22, Section 4.9, Theorem 1]. In particular, $\|\bar{u}\|_{L^2(\Gamma(t))}$ is absolutely continuous, which means that we can interpret $\bar{u}(\cdot, 0)$ as an $L^2(\Gamma_0)$ function. The absolute continuity of $m(\bar{u}, \varphi)$ for $\varphi \in C^2(\mathcal{G}_T)$ also implies that

$$\begin{aligned} & m(\bar{u}(\cdot, t), \varphi(\cdot, t)) - m(\bar{u}(\cdot, 0), \varphi(\cdot, 0)) \\ &= \int_0^t m(\partial^\bullet \bar{u}, \varphi) + g(v; \bar{u}, \varphi) + m(\bar{u}, \partial^\bullet \varphi) \, ds. \end{aligned} \quad (4.10)$$

Next, we choose $\varphi \in C^2(\mathcal{G}_T)$ with $\varphi(\cdot, T) = 0$. It is clear that $\varphi \in L^2_{H^1}$, hence we can use the limiting equation and (4.10) to see that

$$\int_0^T -m(\bar{u}, \partial^\bullet \varphi) + a(\bar{w}, \varphi) \, dt = m(\bar{u}(\cdot, 0), \varphi(\cdot, 0)).$$

We can do the same in the finite element scheme for $\phi_h = \Pi_h \varphi$, using the transport formula (3.8):

$$\int_0^T -m_h(U_h, \partial_h^\bullet \phi_h) + a_h(W_h, \phi_h) \, dt = m_h(\Pi_h u_0, \phi_h(\cdot, 0)).$$

The above calculations show that we are able to take the limit $h \rightarrow 0$ (in the appropriate sense) to see that

$$\int_0^T -m(\bar{u}, \partial^\bullet \varphi) + a(\bar{w}, \varphi) \, dt = m(u_0, \varphi(\cdot, 0)).$$

Therefore, by comparing terms, we have shown that $\bar{u}(\cdot, 0) = u_0$ almost everywhere in Γ_0 by the Fundamental Lemma of the Calculus of Variations.

Hence we have shown the following result:

Theorem 4.1 *Given $u_0 \in H^2(\Gamma_0)$ there exists a weak solution pair (u, w) of the Cahn-Hilliard equation in the sense of Definition 2.1. Furthermore the solution satisfies the energy bound*

$$\sup_{t \in (0, T)} \int_{\Gamma(t)} \frac{\varepsilon}{2} |\nabla_\Gamma u|^2 + \frac{1}{\varepsilon} \psi(u) \, d\sigma + \int_0^T \|w\|_{H^1(\Gamma(t))}^2 \, dt \leq C_2(u_0). \quad (4.11)$$

4.3 Uniqueness

To show the uniqueness result, we require an inverse Laplacian on $\Gamma(t)$. For $z \in L^2(\Gamma(t))$ with $\int_{\Gamma(t)} z \, d\sigma = 0$, we define $\mathcal{G}z$ the inverse Laplacian of z as the unique solution of

$$a(\mathcal{G}z, \varphi) = m(z, \varphi) \quad \text{for all } \varphi \in H^1(\Gamma(t)), \text{ and } \int_{\Gamma(t)} \mathcal{G}z \, d\sigma = 0. \quad (4.12)$$

We will write

$$\|z\|_{-1} := \|\nabla_{\Gamma} \mathcal{G}z\|_{L^2(\Gamma(t))} = a(\mathcal{G}z, \mathcal{G}z)^{\frac{1}{2}}.$$

and remark that

$$\|z\|_{-1}^2 = m(\mathcal{G}z, z).$$

It is clear that if $z \in L^2(\Gamma(t))$ then $\mathcal{G}z \in H^1(\Gamma(t))$. We also have a similar result for the material derivative of $\mathcal{G}z$.

Lemma 4.3 *If $z \in H^1(\mathcal{G}_T)$, with $\int_{\Gamma(t)} z \, d\sigma = 0$, then $\mathcal{G}z \in H^1(\mathcal{G}_T)$.*

Proof It is clear that $\mathcal{G}z \in L^2_{H^1}$ for $z \in L^2_{H^1}$. It is left to show $\partial^\bullet \mathcal{G}z \in L^2_{L^2}$. We start by taking a time derivative of (4.12) so that for $\xi \in H^1(\mathcal{G}_T)$:

$$a(\partial^\bullet \mathcal{G}z, \xi) + a(\mathcal{G}z, \partial^\bullet \xi) + b(v; \mathcal{G}z, \xi) = m(\partial^\bullet z, \xi) + m(z, \partial^\bullet \xi) + g(v; z, \xi).$$

From Lemma 2.3, given $\varphi \in H^1(\Gamma(t^*))$, we can construct $\tilde{\varphi}: \mathcal{G}_T \rightarrow \mathbb{R}$, with $\tilde{\varphi} \in H^1(\Gamma(t))$ for all $t \in [0, T]$ and $\partial^\bullet \tilde{\varphi} = 0$. Thus, we have that

$$a(\partial^\bullet \mathcal{G}z, \tilde{\varphi}) + b(v; \mathcal{G}z, \tilde{\varphi}) = m(\partial^\bullet z, \tilde{\varphi}) + g(v; z, \tilde{\varphi}) \quad \text{for } t \in (0, T),$$

and, in particular, at $t = t^*$,

$$a(\partial^\bullet \mathcal{G}z, \varphi) + b(v; \mathcal{G}z, \varphi) = m(\partial^\bullet z, \varphi) + g(v; z, \varphi).$$

Also, we have that

$$m(\partial^\bullet z, 1) + g(v; z, 1) - b(v; \mathcal{G}z, 1) = \frac{d}{dt} \int_{\Gamma(t)} z \, d\sigma = 0.$$

These calculations imply that $\partial^\bullet \mathcal{G}z$ solves the elliptic problem:

$$a(\partial^\bullet \mathcal{G}z, \varphi) = m(\partial^\bullet z, \varphi) + g(v; z, \varphi) - b(v; \mathcal{G}z, \varphi) \quad \text{for all } \varphi \in H^1(\Gamma(t^*)).$$

This implies that $\partial^\bullet \mathcal{G}z \in H^1(\Gamma(t^*))$ with the bound

$$\|\partial^\bullet \mathcal{G}z\|_{H^1(\Gamma(t^*))} \leq c \left(\|\partial^\bullet z\|_{L^2(\Gamma(t^*))} + \|z\|_{L^2(\Gamma(t^*))} + \|z\|_{-1} \right).$$

Integrating in time gives the desired result. \square

Theorem 4.2 *There is at most one solution to (2.19).*

Proof We suppose that (u_1, w_1) and (u_2, w_2) are solutions to (2.19). We will write $\eta^u = u_1 - u_2$ and $\eta^w = w_1 - w_2$. For $\varphi \in L^2_{H^1}$, we know that

$$m(\partial^\bullet \eta^u, \varphi) + g(v; \eta^u, \varphi) + a(\eta^w, \varphi) = 0 \quad (4.13a)$$

$$\varepsilon a(\eta^u, \varphi) + \frac{1}{\varepsilon} (\psi'(u_1) - \psi'(u_2), \varphi) - m(\eta^w, \varphi) = 0. \quad (4.13b)$$

Testing (4.13a) with $\varphi = 1$ tells us that

$$\int_{\Gamma(t)} \eta^u \, d\sigma = \int_{\Gamma_0} \eta^u \, d\sigma = 0,$$

Hence, since $\mathcal{G}\eta^u$ is well defined and $\mathcal{G}\eta^u \in H^1(\mathcal{G}_T)$, we may test the first equation with $\mathcal{G}\eta^u$, and apply (2.9), to obtain

$$\frac{d}{dt} \|\eta^u\|_{-1}^2 + m(\eta^w, \eta^u) = m(\eta^u, \partial^\bullet \mathcal{G}\eta^u). \quad (4.14)$$

Next, using the monotonicity of $z \mapsto z^3$, testing the second equation with η^u gives

$$\varepsilon a(\eta^u, \eta^u) - \frac{1}{\varepsilon} m(\eta^u, \eta^u) \leq m(\eta^w, \eta^u). \quad (4.15)$$

Taking the sum of (4.15) and (4.14), we obtain

$$\frac{d}{dt} \|\eta^u\|_{-1}^2 + \varepsilon \|\nabla_{\Gamma} \eta^u\|_{L^2(\Gamma(t))}^2 \leq \frac{1}{\varepsilon} m(\eta^u, \eta^u) + m(\eta^u, \partial^\bullet \mathcal{G}\eta^u).$$

For the first term on the right-hand side, we see that

$$\frac{1}{\varepsilon} m(\eta^u, \eta^u) = \frac{1}{\varepsilon} a(\eta^u, \mathcal{G}\eta^u) \leq \frac{\varepsilon}{2} \|\nabla_{\Gamma} \eta^u\|_{L^2(\Gamma(t))}^2 + c_\varepsilon \|\eta^u\|_{-1}^2,$$

and for the second, we have

$$\begin{aligned} m(\eta^u, \partial^\bullet \mathcal{G}\eta^u) &= a(\mathcal{G}\eta^u, \partial^\bullet \mathcal{G}\eta^u) = \frac{1}{2} \frac{d}{dt} a(\mathcal{G}\eta^u, \mathcal{G}\eta^u) - \frac{1}{2} b(v; \mathcal{G}\eta^u, \mathcal{G}\eta^u) \\ &\leq \frac{1}{2} \frac{d}{dt} \|\eta^u\|_{-1}^2 + c \|\eta^u\|_{-1}^2. \end{aligned}$$

Combining these terms, we obtain the estimate

$$\frac{d}{dt} \|\eta^u\|_{-1}^2 + \varepsilon \|\nabla_{\Gamma} \eta^u\|_{L^2(\Gamma(t))}^2 \leq c_\varepsilon \|\eta^u\|_{-1}^2.$$

We next use a Gronwall inequality and integration in time to see

$$\sup_{t \in (0, T)} \|\eta^u\|_{-1}^2 + \varepsilon \int_0^T \|\nabla_{\Gamma} \eta^u\|_{L^2(\Gamma(t))}^2 \leq c_\varepsilon \|\eta^u|_{t=0}\|_{-1}^2 = 0.$$

Since $\int_{\Gamma(t)} \eta^u \, d\sigma = 0$, we apply a Poincaré inequality to arrive at

$$\int_0^T \|\eta^u\|_{L^2(\Gamma(t))}^2 \, dt \leq \int_0^T \|\nabla_{\Gamma} \eta^u\|_{L^2(\Gamma(t))}^2 \, dt = 0.$$

This shows that $u_1 = u_2$.

Now, we know that $\eta^u = 0$ and thus testing (4.13a) with η^w gives

$$m(\eta^w, \eta^w) = \varepsilon a(\eta^u, \eta^w) + \frac{1}{\varepsilon} m(\psi'(u_1) - \psi'(u_2), \eta^w) = 0.$$

This shows that $w_1 = w_2$. □

4.4 Regularity

In this section, we show that the solution enjoys H^2 regularity.

Theorem 4.3 (Regularity) *Let $u_0 \in H^2(\Gamma_0)$ and (u, w) be the solution pair of (2.19), then $u \in L^\infty_{H^2}$ and $w \in L^2_{H^2}$, with the bounds*

$$\varepsilon \sup_{t \in (0, T)} \|u\|_{H^2(\Gamma(t))}^2 + \int_0^T \|w\|_{H^2(\Gamma(t))}^2 \leq C_2(u_0). \quad (4.16)$$

Proof Using the improved estimates from Lemma 4.2, we have that

$$\varepsilon \int_0^T \|\partial^\bullet u\|_{L^2(\Gamma(t))}^2 + \sup_{t \in (0, T)} \|w\|_{L^2(\Gamma(t))}^2 \leq C_2(u_0). \quad (4.17)$$

Now, we can translate the fact that (u, w) are solutions of (2.19) into

$$\begin{aligned} \varepsilon a(u, \varphi) &= m(f_1, \varphi) \\ a(w, \varphi) &= m(f_2, \varphi) \quad \text{for all } \varphi \in H^1(\Gamma(t)), \end{aligned}$$

for $f_1 = w - \frac{1}{\varepsilon} \psi'(u)$ and $f_2 = \partial^\bullet u + u \nabla_\Gamma \cdot v$. Notice that

$$\int_{\Gamma(t)} f_1 \, d\sigma = \int_{\Gamma(t)} f_2 \, d\sigma = 0.$$

The above improved bounds combined with the bounds in Theorem 4.1 gives $f_1 \in L^\infty_{L^2}$ and $f_2 \in L^2_{L^2}$. Standard theory of elliptic partial differential equations [1] gives $u \in L^\infty_{H^2}$ and $w \in L^2_{H^2}$. The proof is completed by using the bounds in (4.11) and (4.17) on f_1 and f_2 . \square

5 Error analysis of finite element scheme

In this section, we show an error bound for the surface finite element method described in Section 3. The proof relies on decomposing the errors into errors between the smooth solution and Ritz projection and between the Ritz projection and discrete solution. In contrast to previous studies of partial differential equations on surfaces [7, 8, 12], we show an error bound on $\Gamma_h(t)$ instead of $\Gamma(t)$. This allows an easier treatment of the non-linear terms.

We will assume that u_0, u and w are bounded in the following norms

$$\|u_0\|_{H^2(\Gamma_0)}^2 + \sup_{t \in (0, T)} \|u\|_{H^2(\Gamma(t))}^2 + \int_0^T \|w\|_{H^2(\Gamma(t))}^2 + \|\partial^\bullet u\|_{H^2(\Gamma(t))}^2 \, dt < +\infty. \quad (5.1)$$

Section 4.4 shows how to bound some of these terms. Again, we will assume that the initial condition of the finite element scheme is given by the Ritz projection:

$$U_{h,0} = I_h u_0. \quad (5.2)$$

The error bound we will show is stated as follows:

Theorem 5.1 *Let u, w solve (2.15) and satisfy (5.1). Let U_h, W_h solve (3.13) with initial condition (5.2). We have that*

$$\varepsilon \sup_{t \in (0, T)} \|u^{-\ell} - U_h\|_{L^2(\Gamma_h(t))}^2 + \int_0^T \|w^{-\ell} - W_h\|_{L^2(\Gamma_h(t))}^2 \leq Ch^4, \quad (5.3)$$

and

$$\varepsilon \sup_{t \in (0, T)} \|\nabla_{\Gamma_h}(u^{-\ell} - U_h)\|_{L^2(\Gamma_h(t))}^2 + \int_0^T \|\nabla_{\Gamma_h}(w^{-\ell} - W_h)\|_{L^2(\Gamma_h(t))}^2 \leq Ch^2, \quad (5.4)$$

with C given by

$$C = c_\varepsilon \sup_{t \in [0, T]} \|u\|_{H^2(\Gamma(t))}^2 + c_\varepsilon \int_0^T (\|\partial^\bullet u\|_{H^2(\Gamma(t))}^2 + \|w\|_{H^2(\Gamma(t))}^2) dt.$$

5.1 Pointwise bound on the discrete solution

In the following error analysis, a pointwise bound on the discrete solution uniformly in space and time will be extremely useful. This will allow us to convert the local Lipschitz property of ψ and ψ' into global results.

Theorem 5.2 *The discrete solution U_h is bounded uniformly in space and time, independently of h , and we have the bound*

$$\sup_{t \in (0, T)} \|U_h\|_{L^\infty(\Gamma_h(t))}^2 \leq C_2(u_0). \quad (5.5)$$

Proof Let $F_h = W_h - \frac{1}{\varepsilon}\psi'(U_h)$, then $F_h \in L^\infty(0, T; L^2(\Gamma_h(t)))$ with the estimate

$$\sup_{t \in (0, T)} \|F_h\|_{L^2(\Gamma_h(t))}^2 \leq C_2(u_0). \quad (5.6)$$

This follows immediately from [Theorem 3.1](#) and [Lemma 4.2](#) combined with a Sobolev inequality ([Lemma 3.5](#)) and (4.1). Furthermore, since $\phi_h = 1$ is an admissible test function in (3.13b), the mean value of F_h is zero:

$$\int_{\Gamma_h(t)} F_h d\sigma_h = 0. \quad (5.7)$$

We define $\tilde{F}_h = F_h^\ell / \mu_h^\ell$, so that

$$\int_{\Gamma(t)} \tilde{F}_h d\sigma = \int_{\Gamma(t)} F_h^\ell \frac{1}{\mu_h^\ell} d\sigma = \int_{\Gamma_h(t)} F_h d\sigma_h = 0.$$

Let $\bar{u}: \mathcal{G}_T \rightarrow \mathbb{R}$ solve

$$-\varepsilon \Delta_\Gamma \bar{u} = \tilde{F}_h \quad \text{on } \Gamma(t), \text{ and } \int_{\Gamma(t)} \bar{u} d\sigma = \int_{\Gamma_h(t)} U_h d\sigma_h \text{ for each } t \in (0, T).$$

Then it is clear that $\Pi_h \bar{u} = U_h$. Standard elliptic theory [1] and the L^∞ bound on Π_h (3.49) gives that

$$\begin{aligned} \|U_h\|_{L^\infty(\Gamma_h(t))} &= \|\Pi_h \bar{u}\|_{L^\infty(\Gamma_h(t))} \leq c \|\bar{u}\|_{H^2(\Gamma(t))} \\ &\leq c \left\| \tilde{F}_h \right\|_{L^2(\Gamma(t))} \leq c \|F_h\|_{L^2(\Gamma_h(t))}. \end{aligned}$$

We apply this inequality uniformly in time, with (5.6), to give the desired estimate. \square

5.2 Splitting the error

We split the error into two parts using the Ritz projection Π_h from Section 3.6:

$$\begin{aligned} u^{-\ell} - U_h &= (u^{-\ell} - \Pi_h u) + (\Pi_h u - U_h) = \rho^u + \theta^u \\ w^{-\ell} - W_h &= (w^{-\ell} - \Pi_h w) + (\Pi_h w - W_h) = \rho^w + \theta^w. \end{aligned}$$

We note that from Theorem 3.2, we already have estimates for ρ^u and ρ^w and it is left to bound θ^u and θ^w . Notice that, the assumptions in (5.1) imply that $\theta^u \in S_h^T$ and $\theta^w \in \tilde{S}_h^T$.

To derive equations for θ^u and θ^w , we start by rewriting (3.13a) using the definition of Π_h and (2.14a) to obtain for $\phi_h \in S_h^T$ with lift $\varphi_h \in S_h^{\ell,T}$ that

$$\begin{aligned} &\frac{d}{dt} m_h(\theta^u, \phi_h) + a_h(\theta^w, \phi_h) - m_h(\theta^u, \partial_h^\bullet \phi_h) \\ &= (m_h(\partial_h^\bullet \Pi_h u, \phi_h) - m(\partial_h^\bullet u, \varphi_h)) + (g_h(V_h; \Pi_h u, \phi_h) - g(v_h; u, \varphi_h)) \\ &\quad + m(u, \partial_h^\bullet \varphi_h - \partial_h^\bullet \varphi_h) \\ &=: E_1(\phi_h) + E_2(\phi_h) + E_3(\phi_h). \end{aligned} \tag{5.8}$$

Next, we rewrite (3.13b) using (2.14b) this time to see for $\phi_h \in \tilde{S}_h^T$ with lift $\varphi_h \in \tilde{S}_h^{\ell,T}$ that

$$\begin{aligned} &\varepsilon a_h(\theta^u, \phi_h) + \frac{1}{\varepsilon} m_h(\psi'(\Pi_h u) - \psi'(U_h), \phi_h) - m_h(\theta^w, \phi_h) \\ &= \frac{1}{\varepsilon} (m_h(\psi'(\Pi_h u), \phi_h) - m(\psi'(u), \varphi_h)) - (m_h(\Pi_h w, \phi_h) - m(w, \varphi_h)) \\ &=: E_4(\phi_h) + E_5(\phi_h). \end{aligned} \tag{5.9}$$

The quantities $E_j(\phi_h)$, for $j = 1, \dots, 5$, are consistency terms involving the approximation properties of the finite element spaces and the geometric perturbation.

Lemma 5.1 For $\phi_h \in S_h^T$ we have

$$|E_1(\phi_h)| \leq ch^2 (\|\partial_h^\bullet u\|_{H^2(\Gamma(t))} + \|u\|_{H^2(\Gamma(t))}) \|\phi_h\|_{L^2(\Gamma_h(t))} \tag{5.10}$$

$$|E_2(\phi_h)| \leq ch^2 \|u\|_{H^2(\Gamma(t))} \|\phi_h\|_{L^2(\Gamma_h(t))} \tag{5.11}$$

$$|E_3(\phi_h)| \leq ch^2 \|u\|_{L^2(\Gamma(t))} \|\phi_h\|_{H^1(\Gamma_h(t))}, \tag{5.12}$$

and for $\phi_h \in \tilde{S}_h^T$:

$$|E_4(\phi_h)| \leq c \frac{h^2}{\varepsilon} \|u\|_{H^2(\Gamma(t))} \|\phi_h\|_{L^2(\Gamma_h(t))} \quad (5.13)$$

$$|E_5(\phi_h)| \leq ch^2 \|w\|_{H^2(\Gamma(t))} \|\phi_h\|_{L^2(\Gamma_h(t))}. \quad (5.14)$$

Proof The proof is a combination of the geometric bounds from [Section 3.5](#) and the bounds of Π_h from [Theorem 3.2](#) and [Lemma 3.14](#). \square

5.3 Error bounds

In this section, we derive bounds on θ^u and θ^w based on the error equations derived in the previous section and natural energy methods for the partial differential equation system go on to show the final error estimate.

To bound θ^u and θ^w we start by testing [\(5.8\)](#) with $\varepsilon\theta^u$ and [\(5.9\)](#) with θ^w and subtract to see that

$$\begin{aligned} & \varepsilon \frac{d}{dt} m_h(\theta^u, \theta^u) + m_h(\theta^w, \theta^w) \\ &= \varepsilon m_h(\theta^u, \partial_h^\bullet \theta^u) + \frac{1}{\varepsilon} m_h(\psi'(\Pi_h u) - \psi'(U_h), \theta^w) \\ & \quad + E_1(\varepsilon\theta^u) + E_2(\varepsilon\theta^u) + E_3(\varepsilon\theta^u) - E_4(\theta^w) - E_5(\theta^w). \end{aligned}$$

Applying [Lemma 5.1](#) and the transport lemma [\(3.8\)](#), with the local Lipschitz property of ψ' , this result gives that

$$\begin{aligned} & \frac{\varepsilon}{2} \frac{d}{dt} m_h(\theta^u, \theta^u) + m_h(\theta^w, \theta^w) \\ & \leq c \frac{\varepsilon}{2} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + \frac{1}{\varepsilon} \|\theta^u\|_{L^2(\Gamma_h(t))} \|\theta^w\|_{L^2(\Gamma_h(t))} \\ & \quad + c\varepsilon h^2 (\|\partial^\bullet u\|_{H^2(\Gamma(t))} + \|u\|_{H^2(\Gamma(t))}) \|\theta^u\|_{L^2(\Gamma_h(t))} \\ & \quad + \frac{ch^2}{\varepsilon} (\|u\|_{H^2(\Gamma(t))} + \|w\|_{H^2(\Gamma(t))}) \|\theta^w\|_{L^2(\Gamma_h(t))} \\ & \quad + c\varepsilon h^2 \|u\|_{H^2(\Gamma(t))} \|\nabla_{\Gamma_h} \theta^u\|_{L^2(\Gamma_h(t))}. \end{aligned} \quad (5.15)$$

We apply a Young's inequality to find that

$$\begin{aligned} & \varepsilon \frac{d}{dt} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + \|\theta^w\|_{L^2(\Gamma_h(t))}^2 \\ & \leq \frac{1}{\varepsilon^2} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + c\varepsilon \|\nabla_{\Gamma_h} \theta^u\|_{L^2(\Gamma_h(t))}^2 \\ & \quad + \frac{ch^4}{\varepsilon^2} (\|\partial^\bullet u\|_{H^2(\Gamma(t))}^2 + \|u\|_{H^2(\Gamma(t))}^2 + \|w\|_{H^2(\Gamma(t))}^2). \end{aligned} \quad (5.16)$$

Next, in order to bound the $\nabla_{\Gamma_h} \theta^u$ term in the previous equation, we test [\(5.9\)](#) with θ^u . Using [Theorem 3.2](#) and [Lemma 5.1](#) and the L^∞ bound on u and

U_h , we have for some $\delta > 0$,

$$\begin{aligned}
\varepsilon a_h(\theta^u, \theta^u) &= m(\theta^w, \theta^u) - \frac{1}{\varepsilon} m_h(\psi'(\Pi_h u) - \psi'(U_h), \theta^u) + E_4(\theta^u) + E_5(\theta^u) \\
&\leq c \frac{1}{\varepsilon} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + \|\theta^w\|_{L^2(\Gamma_h(t))} \|\theta^u\|_{L^2(\Gamma_h(t))} \\
&\quad + c \frac{h^2}{\varepsilon} (\|u\|_{H^2(\Gamma(t))} + \|w\|_{H^2(\Gamma(t))}) \|\theta^u\|_{L^2(\Gamma_h(t))} \\
&\leq c \frac{1}{\varepsilon} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + \delta \|\theta^w\|_{L^2(\Gamma(t))}^2 \\
&\quad + c \frac{h^4}{\varepsilon^2} (\|u\|_{H^2(\Gamma(t))}^2 + \|w\|_{H^2(\Gamma(t))}^2).
\end{aligned} \tag{5.17}$$

Applying this bound in the right-hand side of (5.16), we may choose δ small enough so that

$$\begin{aligned}
\varepsilon \frac{d}{dt} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + \|\theta^w\|_{L^2(\Gamma_h(t))}^2 \\
\leq c \frac{1}{\varepsilon} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + c \frac{h^4}{\varepsilon^2} (\|\partial^\bullet u\|_{H^2(\Gamma(t))}^2 + \|u\|_{H^2(\Gamma(t))}^2 + \|w\|_{H^2(\Gamma(t))}^2).
\end{aligned} \tag{5.18}$$

We recall from (5.2): $U_{h,0} = \Pi_h u_0$, hence we have that $\theta^u|_{t=0} = \Pi_h u_0 - U_{h,0} = 0$. Applying a Gronwall inequality and integrating in time gives the following bounds on θ^u and θ^w :

$$\varepsilon \sup_{t \in (0, T)} \|\theta^u\|_{L^2(\Gamma_h(t))}^2 + \int_0^T \|\theta^w\|_{L^2(\Gamma_h(t))}^2 dt \leq Ch^4, \tag{5.19}$$

with $C = C(u, w, \varepsilon, T)$ given by

$$C = c_\varepsilon \int_0^T (\|\partial^\bullet u\|_{H^2(\Gamma(t))}^2 + \|u\|_{H^2(\Gamma(t))}^2 + \|w\|_{H^2(\Gamma(t))}^2) dt.$$

Proof (Proof of Theorem 5.1) The previous bound can then be combined with the bounds on ρ^u and ρ^w from Theorem 3.2 to give the L^2 error (5.3). One can also apply an inverse inequality to derive gradient bounds on θ^u and θ^w to give the H^1 error bound (5.4). \square

6 Numerical results

The above finite element method discretised in time using semi-implicit time stepping. Given U_0 and a partition of time $0 = t_0, t_1, \dots, t_M = T$, for $k = 0, \dots, M-1$, we find (U_{k+1}, W_{k+1}) as the solution the matrix system

$$\begin{aligned}
\mathcal{M}(t_{k+1})U_{k+1} + (t_{k+1} - t_k)\mathcal{S}(t_{k+1})W_{k+1} &= \mathcal{M}(t_k)U_k \\
\varepsilon\mathcal{S}(t_{k+1})U_{k+1} - \mathcal{M}(t_{k+1})W_{k+1} &= -\frac{1}{\varepsilon}\mathcal{F}(U_k).
\end{aligned}$$

Full analysis of the fully discrete problem is left to future work. Based on ideas from [10], we expect stability subject to $\tau < \varepsilon$ and convergence rate order $\tau + h^2$ for the discrete version of the norms in [Theorem 5.1](#).

The method was implemented using the ALBERTA finite element toolbox [31] and the full block linear system solved using a direct solver.

6.1 Fourth-order linear problem

We start by showing the derived orders of convergence can be achieved for a fourth order linear problem. We calculate with $\psi \equiv 0$ and choose $\varepsilon = 0.1$. We couple $\tau \approx h^2$ to ensure we see the full order of convergence. The surface is given by $\Gamma(t) = \{x \in \mathbb{R}^3 : \Phi(x, t) = 0\}$ with

$$\Phi(x, t) = \frac{x_1^2}{a(t)} + x_2^2 + x_3^2 - 1. \quad (6.1)$$

We have chosen $a(t) = 1.0 + 0.25 \sin(10\pi t)$ and solve for $t \in (0, 0.1)$. The exact solution is given by $u(x, t) = e^{-6t} x_1 x_2$, where right hand side f is calculated from

$$f = u_t + v \cdot \nabla u + u \nabla_{\Gamma} \cdot v + \varepsilon \Delta_{\Gamma}^2 u.$$

The convergence is shown in [Table 6.1](#) for the errors in the L^2 norm. The experimental order of convergence (eoc) is calculated via the formula [\(6.2\)](#): Given an error E_i and E_{i-1} at two different mesh sizes h_i and h_{i-1} , we calculate the experimental order of convergence (eoc) by

$$(\text{eoc})_i = \frac{\log(E_i/E_{i-1})}{\log(h_i/h_{i-1})}. \quad (6.2)$$

The results for the H^1 norm are not shown here, however we observe first order convergence in h .

h	$\ u^{-\ell} - U_h\ _{L^2(\Gamma_h(T))}$	(eoc)
$5.564983 \cdot 10^{-1}$	$9.424750 \cdot 10^{-3}$	—
$2.866409 \cdot 10^{-1}$	$3.001764 \cdot 10^{-3}$	1.724571
$1.443332 \cdot 10^{-1}$	$8.068147 \cdot 10^{-4}$	1.914955
$7.229393 \cdot 10^{-2}$	$2.033971 \cdot 10^{-4}$	1.993007

h	$\ w^{-\ell} - W_h\ _{L^2(\Gamma_h(T))}$	(eoc)
$5.564983 \cdot 10^{-1}$	$4.796888 \cdot 10^{-3}$	—
$2.866409 \cdot 10^{-1}$	$1.432177 \cdot 10^{-3}$	1.821993
$1.443332 \cdot 10^{-1}$	$3.824468 \cdot 10^{-4}$	1.924429
$7.229393 \cdot 10^{-2}$	$9.651516 \cdot 10^{-5}$	1.991496

Table 6.1: Error table of the solution of a fourth-order linear problem with surface defined by [\(6.1\)](#).

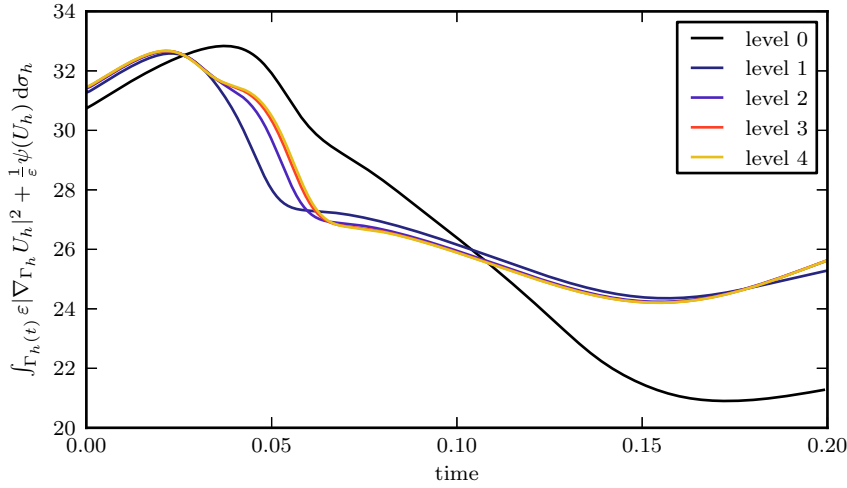


Fig. 6.1: A plot of the Ginzburg-Landau energy over five levels of refinement.

6.2 Cahn-Hilliard equation on a periodically evolving surface

In this example, we consider the same surface as above but now with the full non-linearity as considered in the above analysis over the time interval $t \in (0, 0.8)$.

The initial condition for the simulations was the interpolant of a small perturbation about zero given by

$$u_0(x, y, z) = 0.1 \cos(2\pi x) \cos(2\pi y) \cos(2\pi z).$$

We present two plots to show the behaviour of the numerical solution. First, in [Figure 6.1](#), we see that for short times we have good convergence of the solution. The second, [Figure 6.2](#), demonstrates that the energy does not decrease monotonically along solutions. Running for a longer time suggests that the solution converges to a time periodic solution. We show a plot of the solution at level 2 at different times in [Figure 6.3](#). The system is solved with a fixed time step of 10^{-4} .

6.3 An example with tangential motion

We show the flexibility of the method with an other example with larger surface deformation and tangential motion. The initial condition is taken to be a small random perturbation about zero.

We take a surface given by the level set function

$$\Phi(x, t) = x_1^2 + x_2^2 + a(t)^2 G(x_3^2/L(t)) - a(t)^2, \quad (6.3)$$

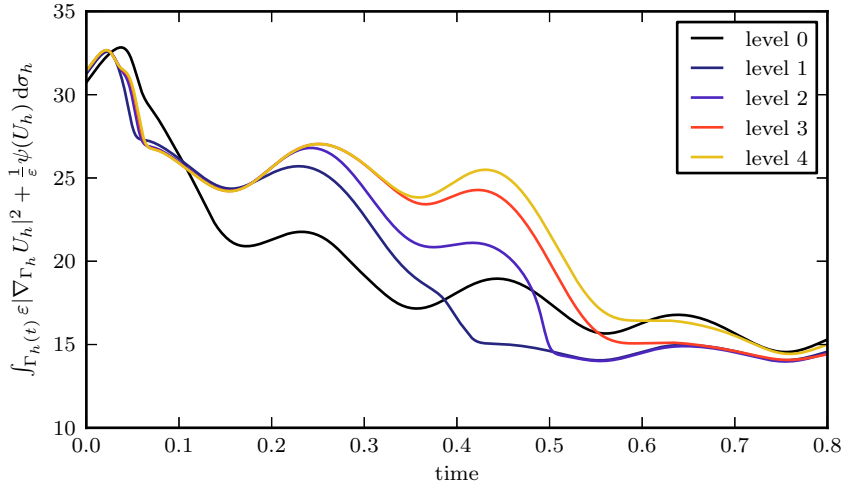


Fig. 6.2: A plot of the Ginzburg-Landau energy over five levels of refinement over a longer time interval.

where

$$\begin{aligned} G(s) &= 200s(s - 199/100) \\ a(t) &= 0.1 + 0.05 \sin(2\pi t) \\ L(t) &= 1 + 0.2 \sin(4\pi t). \end{aligned}$$

In addition, we will prescribe a tangential velocity so that we will consider points moving according to

$$X(t) = \left(X_1(0) \frac{a(t)}{a(0)}, X_2(0) \frac{a(t)}{a(0)}, X_3(0) \frac{L(t)}{L(0)} \right).$$

We plot the solution at different times in [Figure 6.4](#). In particular, we notice that under this flow the nodes remain uniformly distributed.

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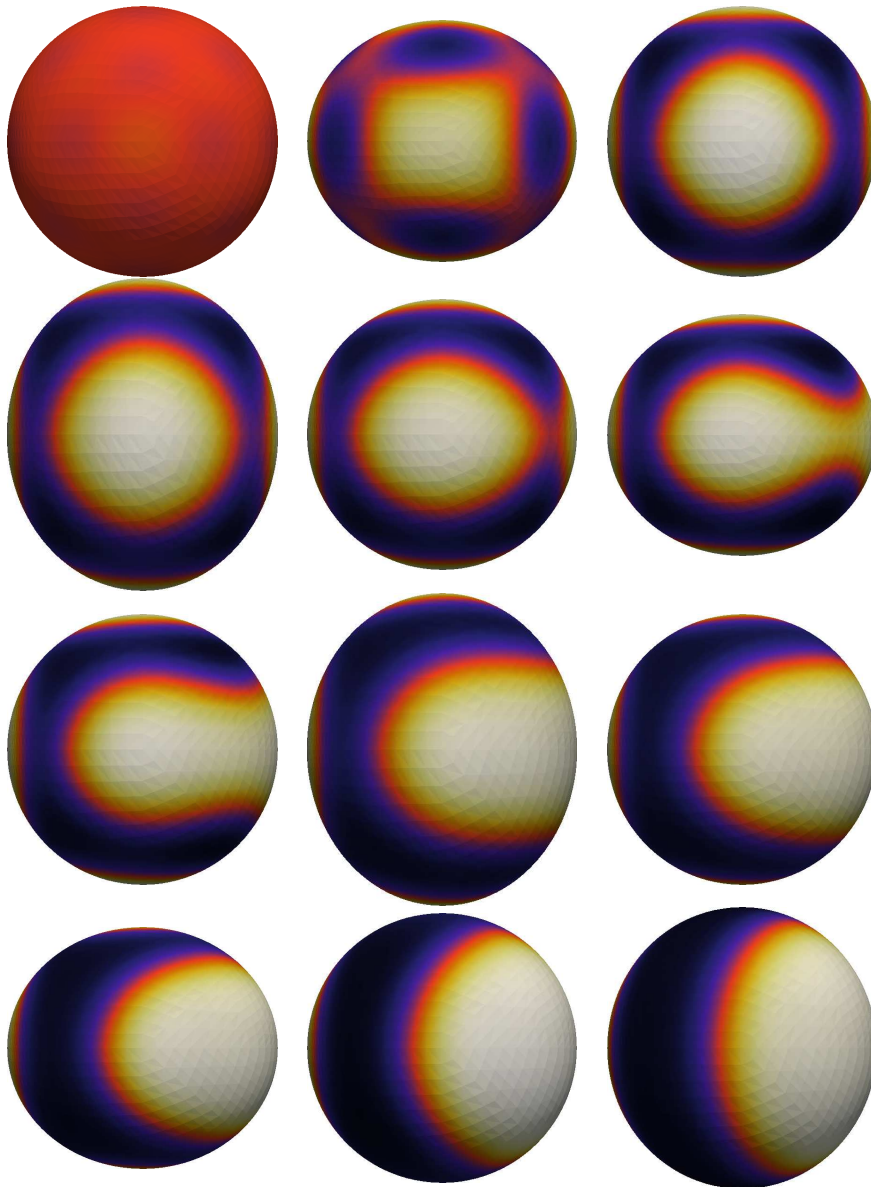


Fig. 6.3: Plot of the solution of the Cahn-Hilliard equation at level two for time $t = 0.0, 0.05, 0.1, 0.15, 0.2, 0.25, 0.3, 0.35, 0.4, 0.45, 0.5, 1.0$. The colour scheme represents values between -1 and 1 .

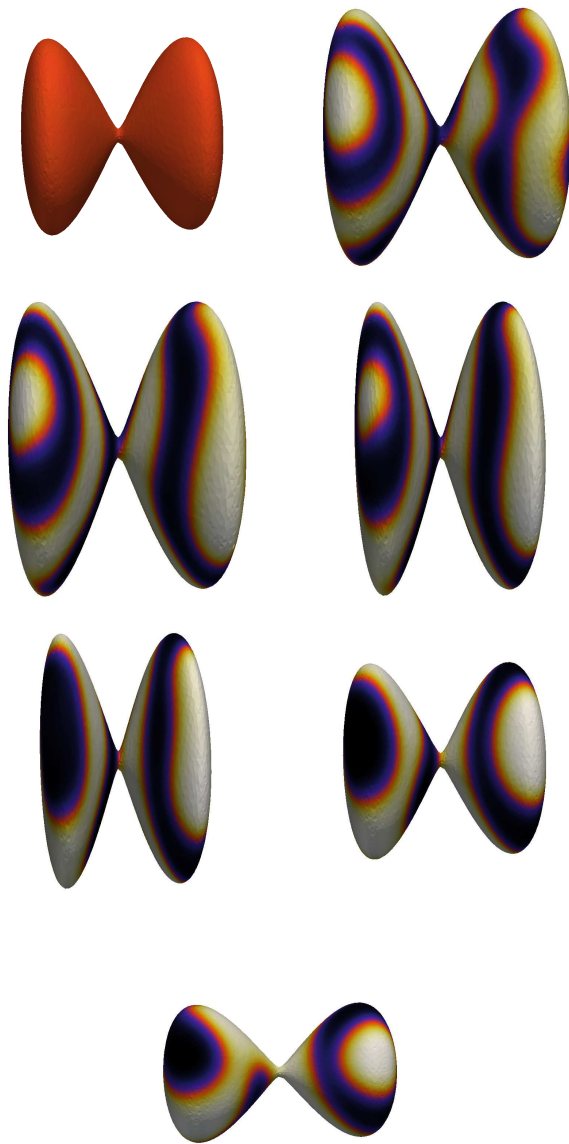


Fig. 6.4: Plot of the solution on the surface defined by (6.3) at times $t = 0, 0.1, 0.2, 0.3, 0.4, 0.5, 0.6$.

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