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**Distributed optimal control of
a nonstandard nonlocal phase field system
with double obstacle potential**

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ABSTRACT

This paper is concerned with a distributed optimal control problem for a non-local phase field model of Cahn–Hilliard type, which is a nonlocal version of a model for two-species phase segregation on an atomic lattice under the presence of diffusion. The *local* model has been investigated in a series of papers by P. Podio-Guidugli and the present authors the *nonlocal* model studied here consists of a highly nonlinear parabolic equation coupled to an ordinary differential inclusion of subdifferential type. The inclusion originates from a free energy containing the indicator function of the interval in which the order parameter of the phase segregation attains its values. It also contains a nonlocal term modeling long-range interactions. Due to the strong nonlinear couplings between the state variables (which even involve products with time derivatives), the analysis of the state system is difficult. In addition, the presence of the differential inclusion is the reason that standard arguments of optimal control theory cannot be applied to guarantee the existence of Lagrange multipliers. In this paper, we employ recent results proved for smooth logarithmic potentials and perform a so-called ‘deep quench’ approximation to establish existence and first-order necessary optimality conditions for the nonsmooth case of the double obstacle potential.

1 Introduction

Let $\Omega \subset \mathbb{R}^3$ denote an open and bounded domain whose smooth boundary Γ has the outward unit normal \mathbf{n} , let $T > 0$ be a given final time, and let $Q := \Omega \times (0, T)$ and $\Sigma := \Gamma \times (0, T)$. We study in this paper distributed optimal control problems of the following form:

(**P**₀) Minimize the cost functional

$$\begin{aligned} J((\mu, \rho), u) &= \frac{\beta_1}{2} \int_0^T \int_{\Omega} |\rho - \rho_Q|^2 dx dt + \frac{\beta_2}{2} \int_0^T \int_{\Omega} |\mu - \mu_Q|^2 dx dt \\ &\quad + \frac{\beta_3}{2} \int_0^T \int_{\Omega} |u|^2 dx dt \end{aligned} \quad (1.1)$$

subject to the state system

$$(1 + 2g(\rho)) \partial_t \mu + \mu g'(\rho) \partial_t \rho - \Delta \mu = u \quad \text{a. e. in } Q, \quad (1.2)$$

$$\partial_t \rho + \xi + F'(\rho) + \mathcal{B}[\rho] = \mu g'(\rho) \quad \text{a. e. in } Q, \quad (1.3)$$

$$\xi \in \partial I_{[0,1]}(\rho) \quad \text{a. e. in } Q, \quad (1.4)$$

$$\partial_{\mathbf{n}} \mu = 0 \quad \text{a. e. on } \Sigma, \quad (1.5)$$

$$\rho(\cdot, 0) = \rho_0, \quad \mu(\cdot, 0) = \mu_0, \quad \text{a. e. in } \Omega, \quad (1.6)$$

and to the control constraints

$$\begin{aligned} u \in \mathcal{U}_{\text{ad}} &:= \{u \in H^1(0, T; L^2(\Omega)) : 0 \leq u \leq u_{\max} \text{ a. e. in } Q \\ &\quad \text{and } \|u\|_{H^1(0, T; L^2(\Omega))} \leq R\}. \end{aligned} \quad (1.7)$$

Here, $\beta_i \geq 0$, $i = 1, 2, 3$, and $R > 0$, are given constants such that $\beta_1 + \beta_2 + \beta_3 > 0$, and the threshold function $u_{\max} \in L^\infty(Q)$ is nonnegative. Moreover, $\rho_Q, \mu_Q \in L^2(Q)$ represent prescribed target functions of the tracking-type functional J . Although more general cost functionals could be admitted for large parts of the subsequent analysis, we restrict ourselves to the above situation for the sake of a simpler exposition.

The state system (1.2)–(1.6) constitutes a *nonlocal* version of a phase field model of Cahn–Hilliard type describing phase segregation of two species (atoms and vacancies, say) on a lattice, which was recently studied in [21]. In the (simpler) original *local* model, which was introduced in [30] (see also [10] and [15]), the nonlocal term $\mathcal{B}[\rho]$ is replaced by the diffusive term $-\Delta\rho$. The local model has been intensively discussed in the past years (cf. [7–11, 13–15, 18]). In particular, in [12] the analogue of the control problem (\mathbf{P}_0) for the local case was investigated for $g(\rho) = \rho$; for this special case, also the optimal boundary control problem was studied (see [17]).

The state variables of the model are the *order parameter* ρ , interpreted as a volumetric density, and the *chemical potential* μ ; for physical reasons, we must have $0 \leq \rho \leq 1$ and $\mu > 0$ almost everywhere in Q . The control function u on the right-hand side of (1.2) plays the role of a *microenergy source*. We remark at this place that the requirement encoded in the definition of \mathcal{U}_{ad} , namely that u be nonnegative, is indispensable for the forthcoming analysis; it is needed to guarantee the nonnegativity of the chemical potential μ .

The nonlinearity F is assumed smooth, while $I_{[0,1]}$ is the indicator function of the interval $[0, 1]$, so that the specific local free energy $F_{\text{loc}} := I_{[0,1]} + F$ is typically a (nonsmooth) double obstacle potential. In this connection, the subdifferential $\partial I_{[0,1]}$ of $I_{[0,1]}$ is defined by

$$\xi \in \partial I_{[0,1]}(\rho) \iff \xi \begin{cases} \leq 0 & \text{if } \rho = 0 \\ = 0 & \text{if } 0 < \rho < 1 \\ \geq 0 & \text{if } \rho = 1 \end{cases} .$$

The presence of the nonlocal term $\mathcal{B}[\rho]$ in (1.3) constitutes the main difference to the local model. Simple examples are given by integral operators of the form

$$\mathcal{B}[\rho](x, t) = \int_0^t \int_{\Omega} k(t, s, x, y) \rho(y, s) ds dy \quad (1.8)$$

and purely spatial convolutions like

$$\mathcal{B}[\rho](x, t) = \int_{\Omega} k(|y - x|) \rho(y, t) dy, \quad (1.9)$$

with sufficiently regular kernels.

The mathematical literature on control problems for phase field systems involving equations of viscous or nonviscous Cahn–Hilliard type is still scarce and quite recent. We refer in this connection to the works [4, 5, 19, 20, 26, 33]. Control problems for convective Cahn–Hilliard systems were studied in [31, 34, 35], and a few analytical contributions were made to the coupled Cahn–Hilliard/Navier–Stokes system (cf. [23–25, 27, 28]). The recent contribution [16] deals with the optimal control of a Cahn–Hilliard type system arising in the modeling of solid tumor growth. On the other hand, let us quote [2] for the analysis of the optimal control problem for a phase field system that couples an energy balance

equation with an ordinary differential inclusion characterized by the presence of the graph $\partial I_{[0,1]}$ acting on the phase variable.

The state system (1.2)–(1.6) is highly nonlinear, with nonstandard couplings involving time derivatives in (1.2). Moreover, (1.3), (1.4) define a variational inequality, and the nonlocal term $\mathcal{B}[\rho]$ is another source for possible analytical difficulties, since the absence of the Laplacian in (1.3) may entail a low regularity of the order parameter ρ . A general analysis of the state system (1.2)–(1.6) was given in [21] for the case $u = 0$ (no control), and in [22] we have investigated the control problem (\mathbf{P}_0) for the case of smooth (but singular) nonlinearities.

In this paper, we aim to employ the results established in [22] to treat the nondifferentiable double obstacle case when ξ satisfies the inclusions (1.4). Recall that it is well known that in this case all of the classical constraint qualifications fail, so that the existence of suitable Lagrange multipliers cannot be guaranteed using standard methods of optimal control. Instead, our approach is guided by a strategy employed in [4] for viscous Cahn–Hilliard systems (see also [6] for the simpler case of the Allen–Cahn equation): in [4], necessary optimality conditions for the double obstacle case could be established by performing a so-called ‘deep quench limit’ in a family of optimal control problems with differentiable nonlinearities of a form that had been previously treated in [20]. The general idea is briefly explained as follows: we replace the inclusion (1.4) by

$$\xi = \varphi(\alpha) h'(\rho), \quad (1.10)$$

where h is defined as the logarithmic potential

$$h(\rho) = \rho \ln(\rho) + (1 - \rho) \ln(1 - \rho), \quad 0 < \rho < 1, \quad h(0) = h(1) = 0, \quad (1.11)$$

and $\varphi \in C^0((0, 1])$ is a positive function satisfying

$$\lim_{\alpha \searrow 0} \varphi(\alpha) = 0. \quad (1.12)$$

We remark that we could simply choose $\varphi(\alpha) = \alpha^p$ for some $p > 0$; however, there might be situations (e. g., in the numerical approximation) in which it is advantageous to let φ have a different behavior as $\alpha \searrow 0$.

Now observe that $h'(\rho) = \ln\left(\frac{\rho}{1-\rho}\right)$ and $h''(\rho) = \frac{1}{\rho(1-\rho)} > 0$ for $\rho \in (0, 1)$. Hence, in particular, we have that

$$\begin{aligned} \lim_{\alpha \searrow 0} \varphi(\alpha) h'(\rho) &= 0 \quad \text{for } 0 < \rho < 1, \\ \lim_{\alpha \searrow 0} \left(\varphi(\alpha) \lim_{\rho \searrow 0} h'(\rho) \right) &= -\infty, \quad \lim_{\alpha \searrow 0} \left(\varphi(\alpha) \lim_{\rho \nearrow 1} h'(\rho) \right) = +\infty. \end{aligned} \quad (1.13)$$

Hence, we may regard the graphs of the functions $\varphi(\alpha) h'$ as approximations to the graph of the subdifferential $\partial I_{[0,1]}$.

We then consider, for any $\alpha \in (0, 1]$, the optimal control problem (later to be denoted by (\mathbf{P}_α)), which results if in (\mathbf{P}_0) the relation (1.4) is replaced by (1.10). For this type of problem, in [20] the existence of optimal controls $u^\alpha \in \mathcal{U}_{\text{ad}}$ as well as first-order

necessary optimality conditions have been derived. Proving a priori estimates (uniform in $\alpha \in (0, 1]$), and employing compactness and monotonicity arguments, we will be able to show the following existence and approximation result: whenever $\{u^{\alpha_n}\} \subset \mathcal{U}_{\text{ad}}$ is a sequence of optimal controls for (\mathbf{P}_{α_n}) , where $\alpha_n \searrow 0$ as $n \rightarrow \infty$, then there exist a subsequence of $\{\alpha_n\}$, which is again indexed by n , and an optimal control $\bar{u} \in \mathcal{U}_{\text{ad}}$ of (\mathbf{P}_0) such that

$$u^{\alpha_n} \rightarrow \bar{u} \quad \text{weakly-star in } \mathcal{X} \text{ as } n \rightarrow \infty, \quad (1.14)$$

where, here and in the following,

$$\mathcal{X} := H^1(0, T; L^2(\Omega)) \cap L^\infty(Q) \quad (1.15)$$

will always denote the control space. In other words, optimal controls for (\mathbf{P}_α) are for small $\alpha > 0$ likely to be ‘close’ to optimal controls for (\mathbf{P}_0) . It is natural to ask if the reverse holds, i. e., whether every optimal control for (\mathbf{P}_0) can be approximated by a sequence $\{u^{\alpha_n}\}$ of optimal controls for (\mathbf{P}_{α_n}) , for some sequence $\alpha_n \searrow 0$.

Unfortunately, we are not able to prove such a ‘global’ result that applies to all optimal controls of (\mathbf{P}_0) . However, a ‘local’ result can be established. To this end, let $\bar{u} \in \mathcal{U}_{\text{ad}}$ be any optimal control for (\mathbf{P}_0) . We introduce the ‘adapted’ cost functional

$$\tilde{J}((\rho, \mu), u) := J((\rho, \mu), u) + \frac{1}{2} \|u - \bar{u}\|_{L^2(Q)}^2 \quad (1.16)$$

and consider for every $\alpha \in (0, 1]$ the *adapted control problem* of minimizing \tilde{J} subject to $u \in \mathcal{U}_{\text{ad}}$ and to the constraint that (ρ, μ) solves the approximating system (1.2), (1.3), (1.10), (1.5), (1.6). It will then turn out that the following is true:

(i) There are some sequence $\alpha_n \searrow 0$ and minimizers $\bar{u}^{\alpha_n} \in \mathcal{U}_{\text{ad}}$ of the adapted control problem associated with α_n , $n \in \mathbb{N}$, such that

$$\bar{u}^{\alpha_n} \rightarrow \bar{u} \quad \text{strongly in } L^2(Q) \text{ as } n \rightarrow \infty. \quad (1.17)$$

(ii) It is possible to pass to the limit as $\alpha \searrow 0$ in the first-order necessary optimality conditions corresponding to the adapted control problems associated with $\alpha \in (0, 1]$ in order to derive first-order necessary optimality conditions for problem (\mathbf{P}_0) .

The paper is organized as follows: in Section 2, we give a precise statement of the problem under investigation, and we derive some results concerning the state system (1.2)–(1.6) and its α -approximation which is obtained if in (\mathbf{P}_0) the relation (1.4) is replaced by the relation (1.10). In Section 3, we then prove the existence of optimal controls and the approximation result formulated above in (i). The final Section 4 is devoted to the derivation of the first-order necessary optimality conditions, where the strategy outlined in (ii) is employed.

Throughout this paper, we will use the following notation: for a (real) Banach space X , we denote by $\|\cdot\|_X$ its norm and the norm of $X \times X \times X$, by X' its dual space, and by $\langle \cdot, \cdot \rangle_X$ the dual pairing between X' and X . If X is an inner product space, then the inner product is denoted by $(\cdot, \cdot)_X$. The only exception from this convention is given by

the L^p spaces, $1 \leq p \leq \infty$, for which we use the abbreviating notation $\|\cdot\|_p$ for the norms in $L^p(\Omega)$. Furthermore, we put

$$H := L^2(\Omega), \quad V := H^1(\Omega), \quad W := \{w \in H^2(\Omega) : \partial_{\mathbf{n}}w = 0 \text{ a.e. on } \Gamma\}.$$

We have the dense and continuous embeddings $W \subset V \subset H \cong H' \subset V' \subset W'$, where $\langle u, v \rangle_V = (u, v)_H$ and $\langle u, w \rangle_W = (u, w)_H$ for all $u \in H$, $v \in V$, and $w \in W$. We will make repeated use of the Young inequalities

$$ab \leq \frac{1}{p}|a|^p + \frac{1}{q}|b|^q \quad \text{and} \quad ab \leq \gamma a^2 + \frac{1}{4\gamma}b^2, \quad (1.18)$$

which are valid for all $a, b \in \mathbb{R}$, $\gamma > 0$, and $p, q \in (1, +\infty)$ with $\frac{1}{p} + \frac{1}{q} = 1$, as well as of the fact that for three dimensions of space and smooth domains the embeddings $V \subset L^p(\Omega)$, $1 \leq p \leq 6$, and $H^2(\Omega) \subset C^0(\bar{\Omega})$ are continuous and (in the first case only for $1 \leq p < 6$) compact. In particular, there are positive constants \tilde{K}_i , $i = 1, 2, 3$, which depend only on the domain Ω , such that

$$\|v\|_6 \leq \tilde{K}_1 \|v\|_V \quad \forall v \in V, \quad (1.19)$$

$$\|vw\|_H \leq \|v\|_6 \|w\|_3 \leq \tilde{K}_2 \|v\|_V \|w\|_3 \quad \forall v, w \in V, \quad (1.20)$$

$$\|v\|_\infty \leq \tilde{K}_3 \|v\|_{H^2(\Omega)} \quad \forall v \in H^2(\Omega). \quad (1.21)$$

For convenience, we also put, for $t \in [0, T]$,

$$Q_t := \Omega \times (0, t) \quad \text{and} \quad Q^t := \Omega \times (t, T). \quad (1.22)$$

Please note the difference in the position of t . About time derivatives of a time-dependent function v , we warn the reader that we will use both the notations $\partial_t v$, $\partial_t^2 v$ and the shorter ones v_t , v_{tt} .

2 Problem statement and results for the state system

Consider the optimal control problem (1.1)–(1.7). We make the following assumptions on the data:

(A1) $F \in C^3[0, 1]$, and $g \in C^3[0, 1]$ satisfies $g(\rho) \geq 0$ and $g''(\rho) \leq 0$ for all $\rho \in [0, 1]$.

(A2) $\rho_0 \in V$, $\mu_0 \in W$, where $\mu_0 \geq 0$ a.e. in Ω , and

$$\inf \{\rho_0(x) : x \in \Omega\} > 0, \quad \sup \{\rho_0(x) : x \in \Omega\} < 1. \quad (2.1)$$

(A3) The nonlocal operator $\mathcal{B}: L^1(Q) \rightarrow L^1(Q)$ satisfies the following conditions: **(i)**

For every $t \in (0, T]$, we have

$$\mathcal{B}[v]|_{Q_t} = \mathcal{B}[w]|_{Q_t} \quad \text{whenever} \quad v|_{Q_t} = w|_{Q_t}. \quad (2.2)$$

(ii) For all $p \in [2, +\infty]$, we have $\mathcal{B}(L^p(Q_t)) \subset L^p(Q_t)$ and

$$\|\mathcal{B}[v]\|_{L^p(Q_t)} \leq C_{\mathcal{B},p} (1 + \|v\|_{L^p(Q_t)}) \quad (2.3)$$

for every $v \in L^p(Q)$ and $t \in (0, T]$.

(iii) For every $v, w \in L^1(0, T; H)$ and $t \in (0, T]$, it holds that

$$\int_0^t \|\mathcal{B}[v](s) - \mathcal{B}[w](s)\|_6 ds \leq C_{\mathcal{B}} \int_0^t \|v(s) - w(s)\|_H ds. \quad (2.4)$$

(iv) It holds, for every $v \in L^2(0, T; V)$ and $t \in (0, T]$, that

$$\|\nabla \mathcal{B}[v]\|_{L^2(0,t;H)} \leq C_{\mathcal{B}} (1 + \|v\|_{L^2(0,t;V)}). \quad (2.5)$$

(v) For every $v \in H^1(0, T; H)$, we have $\partial_t \mathcal{B}[v] \in L^2(Q)$ and

$$\|\partial_t \mathcal{B}[v]\|_{L^2(Q)} \leq C_{\mathcal{B}} (1 + \|\partial_t v\|_{L^2(Q)}). \quad (2.6)$$

(vi) \mathcal{B} is continuously Fréchet differentiable as a mapping from $L^2(Q)$ into $L^2(Q)$, and the Fréchet derivative $D\mathcal{B}[\bar{v}] \in \mathcal{L}(L^2(Q), L^2(Q))$ of \mathcal{B} at \bar{v} has for every $\bar{v} \in L^2(Q)$ and $t \in (0, T]$ the following properties:

$$\|D\mathcal{B}[\bar{v}](w)\|_{L^p(Q_t)} \leq C_{\mathcal{B}} \|w\|_{L^p(Q_t)} \quad \forall w \in L^p(Q), \quad \forall p \in [2, 6], \quad (2.7)$$

$$\|\nabla(D\mathcal{B}[\bar{v}](w))\|_{L^2(Q_t)} \leq C_{\mathcal{B}} \|w\|_{L^2(0,t;V)} \quad \forall w \in L^2(0, T; V). \quad (2.8)$$

In the above formulas, $C_{\mathcal{B},p}$ and $C_{\mathcal{B}}$ denote given positive structural constants. We also notice that (2.7) implicitly requires that $D\mathcal{B}[\bar{v}](w)|_{Q_t}$ depends only on $w|_{Q_t}$; this is, however, a consequence of (2.2).

The statements related to the control problem (\mathbf{P}_0) depend on the assumptions made in the introduction. We recall them here:

(A4) J and \mathcal{U}_{ad} are defined by (1.1) and (1.7), respectively, where

$$\beta_i \geq 0, \quad i = 1, 2, 3, \quad \beta_1 + \beta_2 + \beta_3 > 0, \quad R > 0, \quad (2.9)$$

$$\rho_Q, \mu_Q \in L^2(Q), \quad \text{and } u_{\max} \in L^\infty(Q) \text{ is such that} \\ u_{\max} \geq 0 \text{ a. e. in } Q \text{ and } \mathcal{U}_{\text{ad}} \neq \emptyset. \quad (2.10)$$

REMARK 1: In [22] it was shown that the integral operator (1.9) satisfies all of the conditions (2.2)–(2.8) provided that the integral kernel satisfies $k \in C^1(0, +\infty)$ and fulfills, with suitable constants $C_1 > 0$, $C_2 > 0$, $0 < \alpha < \frac{3}{2}$, $0 < \beta < \frac{5}{2}$, the growth conditions

$$|k(r)| \leq C_1 r^{-\alpha}, \quad |k'(r)| \leq C_2 r^{-\beta}, \quad \forall r > 0.$$

Notice that these growth conditions are met by, e. g., the three-dimensional Newtonian potential, where $k(r) = \frac{c}{r}$ with some $c \neq 0$.

The following well-posedness result for the state system is a direct consequence of [21, Thms. 2.1 and 2.2].

THEOREM 2.1. *Suppose that **(A1)**–**(A4)** are satisfied. Then the state system (1.2)–(1.6) has for every $u \in \mathcal{U}_{\text{ad}}$ a unique solution (ρ, μ, ξ) such that*

$$\mu \in H^1(0, T; H) \cap L^\infty(0, T; V) \cap L^2(0, T; W^{2,3/2}(\Omega)) \cap L^\infty(Q), \quad (2.11)$$

$$\rho \in L^\infty(0, T; V), \quad \rho_t \in L^6(Q), \quad \xi \in L^6(Q), \quad (2.12)$$

$$\mu \geq 0, \quad 0 \leq \rho \leq 1, \quad \text{a. e. in } Q. \quad (2.13)$$

It follows from Theorem 2.1 that under the given assumptions the control-to-state operator for problem (\mathbf{P}_0) is well defined on \mathcal{U}_{ad} . In the following, we denote by \mathcal{S}_0 the operator

$$\mathcal{S}_0 : \mathcal{U}_{\text{ad}} \ni u \mapsto \mathcal{S}_0(u) := (\mu, \rho). \quad (2.14)$$

We now turn our interest to the α -approximating state system, which results when (1.4) is replaced by (1.10). For later convenience, we write this system explicitly:

$$(1 + 2g(\rho^\alpha)) \partial_t \mu^\alpha + \mu^\alpha g'(\rho^\alpha) \partial_t \rho^\alpha - \Delta \mu^\alpha = u \quad \text{a. e. in } Q, \quad (2.15)$$

$$\partial_t \rho^\alpha + \varphi(\alpha) h'(\rho^\alpha) + F'(\rho^\alpha) + \mathcal{B}[\rho^\alpha] = \mu^\alpha g'(\rho^\alpha) \quad \text{a. e. in } Q, \quad (2.16)$$

$$\partial_{\mathbf{n}} \mu^\alpha = 0 \quad \text{a. e. on } \Sigma, \quad (2.17)$$

$$\rho^\alpha(\cdot, 0) = \rho_0, \quad \mu^\alpha(\cdot, 0) = \mu_0, \quad \text{a. e. in } \Omega. \quad (2.18)$$

For this system, we have the following regularity result (cf. [22, Thm. 2.1]):

THEOREM 2.2. *Suppose that **(A1)**–**(A4)** are satisfied. Then the system (2.15)–(2.18) has for every $u \in \mathcal{U}_{\text{ad}}$ and $\alpha \in (0, 1]$ a unique solution $(\mu^\alpha, \rho^\alpha)$ such that*

$$\mu^\alpha \in W^{1,\infty}(0, T; H) \cap H^1(0, T; V) \cap L^\infty(0, T; W) \cap C^0(\bar{Q}), \quad (2.19)$$

$$\rho^\alpha \in H^2(0, T; H) \cap W^{1,\infty}(0, T; L^\infty(\Omega)) \cap H^1(0, T; V). \quad (2.20)$$

Moreover, there are constants $0 < \rho_*(\alpha) < \rho^*(\alpha) < 1$, $\mu^*(\alpha) > 0$, which depend only on the given data and α , such that for every $u \in \mathcal{U}_{\text{ad}}$ and every $\alpha \in (0, 1]$ the corresponding solution $(\mu^\alpha, \rho^\alpha)$ satisfies

$$0 < \rho_*(\alpha) \leq \rho^\alpha \leq \rho^*(\alpha) < 1, \quad 0 \leq \mu^\alpha \leq \mu^*(\alpha), \quad \text{a. e. in } Q. \quad (2.21)$$

From Theorem 2.2, we infer that the α -control-to-state operator

$$\mathcal{S}_\alpha : \mathcal{U}_{\text{ad}} \ni u \mapsto \mathcal{S}_\alpha(u) := (\mu^\alpha, \rho^\alpha) \quad (2.22)$$

is well defined on \mathcal{U}_{ad} for every $\alpha \in (0, 1]$. We now establish some a priori estimates for $(\mu^\alpha, \rho^\alpha)$, which are uniform in α . At first, we conclude from (2.21), (2.3), and **(A1)**, that there is some constant $K_1^* > 0$, which depends only on the data of the problem, such that

$$\max_{0 \leq i \leq 3} \left(\|F^{(i)}(\rho^\alpha)\|_{L^\infty(Q)} + \|g^{(i)}(\rho^\alpha)\|_{L^\infty(Q)} \right) + \|\mathcal{B}[\rho^\alpha]\|_{L^\infty(Q)} \leq K_1^*, \quad (2.23)$$

whenever $(\mu^\alpha, \rho^\alpha) = \mathcal{S}_\alpha(u)$ for some $u \in \mathcal{U}_{\text{ad}}$ and some $\alpha \in (0, 1]$.

Unfortunately, we cannot exclude the possibility that $\rho_*(\alpha) \searrow 0$ and/or $\rho^*(\alpha) \nearrow 1$ as $\alpha \searrow 0$. This is the main reason why the following a priori estimate, which will however be fundamental for our analysis, is comparatively weak.

THEOREM 2.3. *Suppose that (A1)–(A4) are satisfied. Then there is some constant $K_2^* > 0$, which depends only on the given data, such that for every $u \in \mathcal{U}_{\text{ad}}$ and every $\alpha \in (0, 1]$ the corresponding solution $(\mu^\alpha, \rho^\alpha)$ to (2.15)–(2.18) satisfies*

$$\begin{aligned} & \|\mu^\alpha\|_{H^1(0,T;H) \cap L^\infty(0,T;V) \cap L^2(0,T;H^2(\Omega)) \cap L^\infty(Q)} \\ & + \|\rho^\alpha\|_{W^{1,6}(0,T;L^6(\Omega)) \cap L^\infty(0,T;V)} + \|\varphi(\alpha) h'(\rho^\alpha)\|_{L^6(Q)} \\ & + \int_0^T \int_\Omega \varphi(\alpha) \frac{1}{\rho^\alpha(1-\rho^\alpha)} |\nabla \rho^\alpha|^2 dx dt \leq K_2^*. \end{aligned} \quad (2.24)$$

PROOF: Let $u \in \mathcal{U}_{\text{ad}}$ be fixed and $(\mu^\alpha, \rho^\alpha) = \mathcal{S}_\alpha(u)$ for $\alpha \in (0, 1]$. We establish the validity of (2.24) in a series of steps. In what follows, we denote by C_i , $i \in \mathbb{N}$, positive constants that may depend on the data of our problem but not on $\alpha \in (0, 1]$. For the sake of a better readability, we will omit the arguments of functions if there is no danger of confusion, as well as the superscript $\alpha \in (0, 1]$, which will only be written at the end of each step of estimation.

STEP 1: First observe that $\partial_t \left(\left(\frac{1}{2} + g(\rho) \right) \mu^2 \right) = (1 + 2g(\rho)) \mu \mu_t + g'(\rho) \rho_t \mu^2$. Hence, multiplication of (2.15) by μ and integration over Q_t , for $0 < t \leq T$, yields the identity

$$\begin{aligned} & \int_\Omega \left(\frac{1}{2} + g(\rho(t)) \right) \mu^2(t) dx + \int_0^t \int_\Omega |\nabla \mu|^2 dx ds \\ & = \int_\Omega \left(\frac{1}{2} + g(\rho_0) \right) \mu_0^2 dx + \int_0^t \int_\Omega u \mu dx ds, \end{aligned}$$

whence, using (A1), Young's inequality, and Gronwall's lemma, we easily deduce that

$$\|\mu^\alpha\|_{L^\infty(0,T;H) \cap L^2(0,T;V)} \leq C_1 \quad \forall \alpha \in (0, 1]. \quad (2.25)$$

Moreover, since the embedding $(L^\infty(0, T; H) \cap L^2(0, T; V)) \subset L^{10/3}(Q)$ is continuous, we have (by possibly choosing a larger constant C_1) that

$$\|\mu^\alpha\|_{L^{10/3}(Q)} \leq C_1 \quad \forall \alpha \in (0, 1]. \quad (2.26)$$

STEP 2: Next, we multiply (2.16) by ρ_t and integrate over Q_t , where $0 < t \leq T$. It follows

$$\begin{aligned} & \int_0^t \int_\Omega |\rho_t|^2 dx ds + \int_\Omega \varphi(\alpha) h(\rho(t)) dx - \int_\Omega \varphi(\alpha) h(\rho_0) dx \\ & = \int_0^t \int_\Omega (-F'(\rho) - \mathcal{B}[\rho] + \mu g'(\rho)) \rho_t dx ds. \end{aligned}$$

Thus, employing (2.1), (2.23), and (2.25), we readily obtain from Young's inequality that

$$\|\rho^\alpha\|_{H^1(0,T;H)} \leq C_2 \quad \forall \alpha \in (0, 1]. \quad (2.27)$$

STEP 3: Now observe that we have the identity

$$\nabla \rho_t + \varphi(\alpha) h''(\rho) \nabla \rho = -F''(\rho) \nabla \rho - \nabla \mathcal{B}[\rho] + g'(\rho) \nabla \mu + \mu g''(\rho) \nabla \rho.$$

Taking the euclidean scalar product of this identity with $\nabla \rho$, and integrating the result over Q_t , where $0 < t \leq T$, we find that

$$\begin{aligned} & \frac{1}{2} \|\nabla \rho(t)\|_H^2 - \frac{1}{2} \|\nabla \rho_0\|_H^2 + \int_0^t \int_\Omega \varphi(\alpha) h''(\rho) |\nabla \rho|^2 dx ds \\ &= \int_0^t \int_\Omega (-F''(\rho) \nabla \rho - \nabla \mathcal{B}[\rho] + g'(\rho) \nabla \mu) \cdot \nabla \rho dx ds \\ & \quad + \int_0^t \int_\Omega \mu g''(\rho) |\nabla \rho|^2 dx ds, \end{aligned}$$

where, since $h''(\rho) \geq 0$, the integral on the left-hand side is nonnegative, while, thanks to the fact that $g''(\rho) \leq 0$ and $\mu \geq 0$, the second integral on the right-hand side is nonpositive. Moreover, by using (2.5), (2.23), (2.25), and Young's inequality, we easily find that the first integral on the right-hand side is bounded by an expression of the form $C_3 (1 + \|\rho\|_{L^2(0,t;V)}^2)$. Hence, we can infer from Gronwall's lemma and (2.27) that

$$\|\rho^\alpha\|_{L^\infty(0,T;V)} + \int_0^T \int_\Omega \varphi(\alpha) h''(\rho^\alpha) |\nabla \rho^\alpha|^2 dx ds \leq C_4 \quad \forall \alpha \in (0, 1]. \quad (2.28)$$

In particular, this yields the desired bound for the last term on the right-hand side of (2.24).

STEP 4: Next, we multiply (2.16) by $|\varphi(\alpha) h'(\rho)|^{7/3} \text{sign}(\varphi(\alpha) h'(\rho))$ (where we put $\text{sign}(0) := 0$) and integrate over Q_t , where $0 < t \leq T$. We obtain the inequality

$$\begin{aligned} & \int_\Omega (\Phi(x, t) - \Phi(x, 0)) dx + \|\varphi(\alpha) h'(\rho)\|_{L^{10/3}(Q_t)}^{10/3} \\ & \leq \int_0^t \int_\Omega |-F'(\rho) - \mathcal{B}[\rho] + \mu g'(\rho)| |\varphi(\alpha) h'(\rho)|^{7/3} dx ds, \end{aligned} \quad (2.29)$$

where, by the definition of h , the function

$$\Phi(x, t) := \int_{1/2}^{\rho(x,t)} |\varphi(\alpha) h'(r)|^{7/3} \text{sign}(\varphi(\alpha) h'(r)) dr$$

is nonnegative almost everywhere in Q , and where $\Phi(\cdot, 0)$ is bounded in Ω independently of $\alpha \in (0, 1]$ by (2.1). Moreover, by virtue of (2.23) and (2.26), we have

$$\|-F'(\rho) - \mathcal{B}[\rho] + \mu g'(\rho)\|_{L^{10/3}(Q)} \leq C_5.$$

Employing Young's inequality with $q = 10/7$, we thus can infer from (2.29) that

$$\|\varphi(\alpha) h'(\rho)\|_{L^{10/3}(Q_t)}^{10/3} \leq C_6 + \frac{7}{10} \|\varphi(\alpha) h'(\rho)\|_{L^{10/3}(Q_t)}^{10/3},$$

whence we conclude that

$$\|\varphi(\alpha) h'(\rho^\alpha)\|_{L^{10/3}(Q)} \leq C_7 \quad \forall \alpha \in (0, 1].$$

Comparison in (2.16) then shows that also

$$\|\rho_t^\alpha\|_{L^{10/3}(Q)} \leq C_8 \quad \forall \alpha \in (0, 1]. \quad (2.30)$$

STEP 5: Now that (2.30) is established, we are in a position to test (2.15) by μ_t . Integrating over Q_t , where $0 < t \leq T$, and using the fact that g is nonnegative, we obtain the estimate

$$\begin{aligned} & \int_0^t \int_\Omega |\mu_t|^2 dx ds + \frac{1}{2} \|\nabla \mu(t)\|_H^2 - \frac{1}{2} \|\nabla \mu_0\|_H^2 \\ & \leq \int_0^t \int_\Omega |u| |\mu_t| dx ds + \int_0^t \int_\Omega |g'(\rho)| |\rho_t| |\mu| |\mu_t| dx ds =: I_1 + I_2. \end{aligned} \quad (2.31)$$

Owing to Young's inequality, we have that

$$I_1 \leq \frac{1}{4} \int_0^t \int_\Omega |\mu_t|^2 dx ds + C_9,$$

and invoking (2.23) and (1.19), as well as Hölder's and Young's inequalities, we find that

$$\begin{aligned} I_2 & \leq C_{10} \int_0^t \|\mu_t(s)\|_2 \|\rho_t(s)\|_3 \|\mu(s)\|_6 ds \\ & \leq \frac{1}{4} \int_0^t \int_\Omega |\mu_t|^2 dx ds + C_{11} \int_0^t \|\rho_t(s)\|_3^2 \|\mu(s)\|_V^2 ds. \end{aligned}$$

Now observe that $\frac{10}{3} > 3 > 2$, and therefore Hölder's and Young's inequalities can be employed to bound the $L^1(0, T)$ norm of the function $s \mapsto \|\rho_t(s)\|_3^2$ as follows:

$$\int_0^T \|\rho_t(s)\|_3^2 ds \leq C_{12} \int_0^T \|\rho_t(s)\|_{10/3}^2 ds \leq C_{13} \left(1 + \int_0^T \|\rho_t(s)\|_{10/3}^{10/3} ds \right),$$

where the last expression is by (2.30) bounded independently of $\alpha \in (0, 1]$. In conclusion, we may employ Gronwall's lemma to conclude from (2.31) the estimate

$$\|\mu^\alpha\|_{H^1(0,T;H) \cap L^\infty(0,T;V)} \leq C_{14}. \quad (2.32)$$

STEP 6: Next, recalling the continuity of the embedding $H^1(\Omega) \subset L^6(\Omega)$ and invoking (2.23) and (2.32), we find that

$$\| -F'(\rho) - \mathcal{B}[\rho] + \mu g'(\rho) \|_{L^6(Q)} \leq C_{15}.$$

Therefore, if we multiply (2.16) by $(\varphi(\alpha) h'(\rho))^5$ and integrate over Q_t , where $0 < t \leq T$, then we can conclude from Young's inequality with $q = 6/5$ the inequality

$$\begin{aligned} & \int_{\Omega} (\Psi(x, t) - \Psi(x, 0)) dx + \|\varphi(\alpha) h'(\rho)\|_{L^6(Q_t)}^6 \\ & \leq C_{16} + \frac{5}{6} \|\varphi(\alpha) h'(\rho)\|_{L^6(Q_t)}^6, \end{aligned} \quad (2.33)$$

where the function

$$\Psi(x, t) := \int_{1/2}^{\rho(x, t)} (\varphi(\alpha) h'(r))^5 dr$$

is obviously nonnegative, and where $\Psi(\cdot, 0)$ is by (2.1) bounded, uniformly in $\alpha \in (0, 1]$. We thus conclude from (2.33) that

$$\|\varphi(\alpha) h'(\rho^\alpha)\|_{L^6(Q)} \leq C_{17} \quad \forall \alpha \in (0, 1], \quad (2.34)$$

whence, by comparison in (2.16), also

$$\|\rho_t^\alpha\|_{L^6(Q)} \leq C_{18} \quad \forall \alpha \in (0, 1]. \quad (2.35)$$

STEP 7: From (2.35), we have, in particular, that

$$\|\rho_t^\alpha\|_{L^{7/3}(0, T; L^{14/3}(\Omega))} \leq C_{19} \quad \forall \alpha \in (0, 1]. \quad (2.36)$$

We may therefore account for the conditions $\mu_0 \in W$ and $\|u\|_{L^\infty(Q)} \leq R$ in order to repeat the argument developed in the proof of [10, Thm. 2.3], which is based on the above summability of ρ_t^α . We should remark that the quoted proof is performed for $g(\rho) = \rho$ and $u \equiv 0$, but only minor changes are needed to arrive at the same conclusion in the present situation (see also the proof of the analogous [15, Thm. 3.7] in an even more complicated case). We thus can infer the bound

$$\|\mu^\alpha\|_{L^\infty(Q)} \leq C_{20} \quad \forall \alpha \in (0, 1]. \quad (2.37)$$

Finally, we use the previously shown estimates to conclude that the expression

$$(1 + 2g(\rho^\alpha)) \mu_t^\alpha + \mu^\alpha g'(\rho^\alpha) \rho_t^\alpha - u \quad \text{is bounded in } L^2(Q)$$

independently of $\alpha \in (0, T]$, which then also holds true for $\Delta\mu^\alpha$, by comparison in (2.15). We thus have

$$\|\mu^\alpha\|_{L^2(0, T; H^2(\Omega))} \leq C_{21} \quad \forall \alpha \in (0, 1]. \quad (2.38)$$

This concludes the proof of the assertion. ■

3 Existence and approximation of optimal controls

Our first aim is to show the existence result stated below. For its proof we do not use the direct method since this would force us to find bounds for the states corresponding to the

chosen minimizing sequence of controls. On the contrary, Theorem 2.3 already provides a number of estimates. For that reason, we pass through the related approximating control problem.

THEOREM 3.1. *Suppose that the conditions (A1)–(A4) are fulfilled. Then the optimal control problem (P₀) admits a solution.*

PROOF: Recall the properties (1.11) and (1.12) of the functions h and φ . We consider for $\alpha \in (0, 1]$ the approximating optimal control problem

(P_α) Minimize $J((\mu^\alpha, \rho^\alpha), u)$ subject to $u \in \mathcal{U}_{\text{ad}}$ and $(\mu^\alpha, \rho^\alpha) = \mathcal{S}_\alpha(u)$, where \mathcal{S}_α is the solution operator of (2.15)–(2.18) defined in (2.22).

According to [22, Thm. 4.1], (P_α) enjoys for every $\alpha \in (0, 1]$ at least one solution u^α . Now let $\{\alpha_n\} \subset (0, 1]$ be any sequence such that $\alpha_n \searrow 0$ as $n \rightarrow \infty$. By virtue of Theorem 2.3, there exists a subsequence of α_n (still denoted by α_n) such that

$$u^{\alpha_n} \rightarrow u \quad \text{weakly-star in } L^\infty(Q) \text{ and weakly in } H^1(0, T; H), \quad (3.1)$$

$$\begin{aligned} \mu^{\alpha_n} &\rightarrow \mu \quad \text{weakly-star in } L^\infty(Q) \cap L^\infty(0, T; V) \\ &\quad \text{and weakly in } H^1(0, T; H) \cap L^2(0, T; W), \end{aligned} \quad (3.2)$$

$$\begin{aligned} \rho^{\alpha_n} &\rightarrow \rho \quad \text{weakly-star in } L^\infty(Q) \cap L^\infty(0, T; V) \\ &\quad \text{and weakly in } H^1(0, T; H), \end{aligned} \quad (3.3)$$

$$\rho_t^{\alpha_n} \rightarrow \rho_t \quad \text{weakly in } L^6(Q), \quad (3.4)$$

$$\varphi(\alpha_n)h'(\rho^{\alpha_n}) \rightarrow \xi \quad \text{weakly in } L^6(Q). \quad (3.5)$$

At first, we note that the set \mathcal{U}_{ad} is bounded, convex, and closed (and thus weakly sequentially closed) in $H^1(0, T; H)$, whence it follows that $u \in \mathcal{U}_{\text{ad}}$.

We now claim that (μ, ρ, ξ) coincides with the unique solution to the state system (1.2)–(1.6) established in Theorem 2.1. At first, $\mu \in L^2(0, T; W)$ satisfies the boundary condition (1.5), and we have that $\Delta\mu^{\alpha_n} \rightarrow \Delta\mu$ weakly in $L^2(Q)$. Moreover, we can infer from standard compactness results (cf. [32, Sect. 8, Cor. 4]) that, for every $q \in [1, 6)$, it holds that

$$\rho^{\alpha_n} \rightarrow \rho \quad \text{strongly in } C^0([0, T]; L^q(\Omega)), \quad (3.6)$$

$$\mu^{\alpha_n} \rightarrow \mu \quad \text{strongly in } C^0([0, T]; L^q(\Omega)), \quad (3.7)$$

and both pointwise a. e. in Q , at least for a subsequence of α_n . Consequently, μ and ρ fulfill the initial conditions from (1.6), and (2.13) holds true. In addition, we conclude from (A1) that, for $\Phi \in \{g, g', F'\}$ and $1 \leq q < 6$,

$$\Phi(\rho^{\alpha_n}) \rightarrow \Phi(\rho) \quad \text{strongly in } C^0([0, T]; L^q(\Omega)) \quad (3.8)$$

and pointwise a. e. in Q . From (3.2)–(3.8), it readily follows that we have the convergence given by

$$\begin{aligned} g(\rho^{\alpha_n})\mu_t^{\alpha_n} &\rightarrow g(\rho)\mu_t, \quad \mu^{\alpha_n}g'(\rho^{\alpha_n})\rho_t^{\alpha_n} \rightarrow \mu g'(\rho)\rho_t, \quad \mu^{\alpha_n}g'(\rho^{\alpha_n}) \rightarrow \mu g'(\rho), \\ &\quad \text{all weakly in } L^1(Q), \text{ as } n \rightarrow \infty. \end{aligned} \quad (3.9)$$

Moreover, assumption **(A3,vi)** (see, in particular, condition (2.7)) implies that \mathcal{B} is Lipschitz continuous on $L^2(Q)$, which entails that

$$\mathcal{B}[\rho^{\alpha_n}] \rightarrow \mathcal{B}[\rho] \quad \text{strongly in } L^2(Q). \quad (3.10)$$

In conclusion, the triple (μ, ρ, ξ) also fulfills (1.2) and (1.3). It remains to show that (1.4) holds true as well. Once this is shown, we can infer that $(\mu, \rho) = \mathcal{S}_0(u)$, i. e., that the pair $((\mu, \rho, \xi), u)$ is admissible for (\mathbf{P}_0) .

To this end, we recall that h is convex and bounded in $[0, 1]$ and that φ is nonnegative. We thus have, for every $n \in \mathbb{N}$,

$$\begin{aligned} & \varphi(\alpha_n) \operatorname{Min}_{[0,1]} h \cdot \operatorname{meas} Q + \int_0^T \int_{\Omega} \varphi(\alpha_n) h'(\rho^{\alpha_n}) (z - \rho^{\alpha_n}) dx dt \\ & \leq \int_0^T \int_{\Omega} \varphi(\alpha_n) h(\rho^{\alpha_n}) dx dt + \int_0^T \int_{\Omega} \varphi(\alpha_n) h'(\rho^{\alpha_n}) (z - \rho^{\alpha_n}) dx dt \\ & \leq \int_0^T \int_{\Omega} \varphi(\alpha_n) h(z) dx dt \\ & \text{for all } z \in \mathcal{K} = \{v \in L^2(Q) : 0 \leq v \leq 1 \text{ a. e. in } Q\}. \end{aligned} \quad (3.11)$$

Now, we notice that the first term and the last one of this chain tend to zero as $n \rightarrow \infty$ by (1.12). Hence, by ignoring the middle line of (3.11) and invoking (3.5) and (3.6), passage to the limit as $n \rightarrow \infty$ yields

$$\int_0^T \int_{\Omega} \xi (\rho - z) dx dt \geq 0 \quad \forall z \in \mathcal{K}.$$

This entails that ξ is an element of the subdifferential of the extension \mathcal{I} of $I_{[0,1]}$ to $L^2(Q)$, which means that $\xi \in \partial \mathcal{I}(\rho)$ or, equivalently (cf. [3, Ex. 2.3.3., p. 25]), that $\xi \in \partial I_{[0,1]}(\rho)$ a. e. in Q .

It remains to show that $((\mu, \rho, \xi), u)$ is in fact optimal for (\mathbf{P}_0) . To this end, let $v \in \mathcal{U}_{\text{ad}}$ be arbitrary. In view of the convergence properties (3.1)–(3.7), and using the weak sequential lower semicontinuity properties of the cost functional, we obtain

$$\begin{aligned} J((\mu, \rho), u) &= J(\mathcal{S}_0(u), u) \leq \liminf_{n \rightarrow \infty} J(\mathcal{S}_{\alpha_n}(u^{\alpha_n}), u^{\alpha_n}) \\ &\leq \liminf_{n \rightarrow \infty} J(\mathcal{S}_{\alpha_n}(v), v) = \lim_{n \rightarrow \infty} J(\mathcal{S}_{\alpha_n}(v), v) = J(\mathcal{S}_0(v), v), \end{aligned} \quad (3.12)$$

where for the last equality the continuity in $L^2(Q) \times L^2(Q)$ of the cost functional with respect to (μ, ρ) was used (see however the next statement). With this, the assertion is completely proved. \blacksquare

COROLLARY 3.2. *Let the general assumptions **(A1)**–**(A4)** satisfied, and let sequences $\{\alpha_n\} \subset (0, 1]$ and $\{u^{\alpha_n}\} \subset \mathcal{U}$ be given such that, as $n \rightarrow \infty$, $\alpha_n \searrow 0$ and $u^{\alpha_n} \rightarrow u$ weakly-star in \mathcal{X} . Then, with $(\mu^{\alpha_n}, \rho^{\alpha_n}) = \mathcal{S}_{\alpha_n}(u^{\alpha_n})$, $n \in \mathbb{N}$, and $(\mu, \rho) = \mathcal{S}_0(u)$, (3.2)–(3.10) hold true, where $\xi \in L^6(Q)$ satisfies $\xi \in \partial I_{[0,1]}(\rho)$ almost everywhere in Q . Moreover, we have*

$$\lim_{n \rightarrow \infty} J(\mathcal{S}_{\alpha_n}(v), v) = J(\mathcal{S}_0(v), v) \quad \forall v \in \mathcal{U}_{\text{ad}}. \quad (3.13)$$

PROOF: By the same arguments as in the first part of the proof of Theorem 3.1, we can conclude that (3.2)–(3.10) hold true at least for some subsequence. But, as we have just seen, the limit is given by the unique solution triple to the state system (1.2)–(1.6). Hence, the limit is the same for all convergent subsequences, and thus (3.2)–(3.10) are true for the entire sequence, as claimed.

Now, let $v \in \mathcal{U}_{\text{ad}}$ be arbitrary. Then, owing to (3.6), (3.7), $\mathcal{S}_{\alpha_n}(v)$ converges strongly to $\mathcal{S}_0(v)$ in $L^2(Q) \times L^2(Q)$. The validity of (3.13) is then a consequence of the fact that J is continuous in $L^2(Q) \times L^2(Q)$ with respect to (μ, ρ) . \blacksquare

Theorem 3.1 does not yield any information on whether every solution to the optimal control problem (\mathbf{P}_0) can be approximated by a sequence of solutions to the problems (\mathbf{P}_α) . As already announced in the introduction, we are not able to prove such a general ‘global’ result. Instead, we give a ‘local’ answer for every individual optimizer of (\mathbf{P}_0) . For this purpose, we employ a trick due to Barbu [1]. Now let $\bar{u} \in \mathcal{U}_{\text{ad}}$ be an arbitrary optimal control for (\mathbf{P}_0) , and let $(\bar{\mu}, \bar{\rho}, \bar{\xi})$ be the associated solution triple to the state system (1.2)–(1.6) in the sense of Theorem 2.1. In particular, $(\bar{\mu}, \bar{\rho}) = \mathcal{S}_0(\bar{u})$. We associate with this optimal control the ‘adapted cost functional’

$$\tilde{J}((\mu, \rho), u) := J((\mu, \rho), u) + \frac{1}{2} \|u - \bar{u}\|_{L^2(Q)}^2 \quad (3.14)$$

and a corresponding ‘adapted optimal control problem’

$$(\tilde{\mathbf{P}}_\alpha) \quad \text{Minimize } \tilde{J}((\mu^\alpha, \rho^\alpha), u) \quad \text{subject to } u \in \mathcal{U}_{\text{ad}} \text{ and (2.15)–(2.18).}$$

With a standard direct argument that needs no repetition here, we can show the following result.

LEMMA 3.3. *Suppose that the assumptions (A1)–(A4) are fulfilled. Then the optimal control problem $(\tilde{\mathbf{P}}_\alpha)$ has for every $\alpha \in (0, 1]$ a solution.*

We are now in the position to give a partial answer to the question raised above. We have the following result.

THEOREM 3.4. *Let the general assumptions (A1)–(A4) be fulfilled, suppose that $\bar{u} \in \mathcal{U}_{\text{ad}}$ is an arbitrary optimal control of (\mathbf{P}_0) with associated state triple $(\bar{\mu}, \bar{\rho}, \bar{\xi})$, and let $\{\alpha_n\} \subset (0, 1]$ be any sequence such that $\alpha_n \searrow 0$ as $n \rightarrow \infty$. Then there exist a subsequence $\{\alpha_{n_k}\}$ of $\{\alpha_n\}$ and, for every $k \in \mathbb{N}$, an optimal control $u^{\alpha_{n_k}} \in \mathcal{U}_{\text{ad}}$ of the adapted problem $(\tilde{\mathbf{P}}_{\alpha_{n_k}})$ with associated state $(\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}})$ such that, as $k \rightarrow \infty$,*

$$u^{\alpha_{n_k}} \rightarrow \bar{u} \quad \text{strongly in } L^2(Q), \quad (3.15)$$

$$(3.2)–(3.10) \text{ hold true, where } (\mu, \rho, \xi) \text{ is replaced} \\ \text{by } (\bar{\mu}, \bar{\rho}, \bar{\xi}) \text{ and the index } n \text{ is replaced by } n_k, \quad (3.16)$$

$$\tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}) \rightarrow J((\bar{\mu}, \bar{\rho}), \bar{u}). \quad (3.17)$$

PROOF: Let $\alpha_n \searrow 0$ as $n \rightarrow \infty$. For any $n \in \mathbb{N}$, we pick an optimal control $u^{\alpha_n} \in \mathcal{U}_{\text{ad}}$ for the adapted problem $(\tilde{\mathbf{P}}_{\alpha_n})$ and denote by $(\mu^{\alpha_n}, \rho^{\alpha_n}) = \mathcal{S}_{\alpha_n}(u^{\alpha_n})$ the associated solution to problem (2.15)–(2.18), where the right-hand side in (2.15) is given by u^{α_n} . Then, by virtue of Theorems 2.2 and 2.3, (2.19)–(2.21) and (2.23)–(2.24) are satisfied. Arguing as in the proof of Corollary 3.2, and using the fact that \mathcal{U}_{ad} is a bounded subset of \mathcal{X} (see (1.15)), we can infer that there is a subsequence $\{n_k\}_{k \in \mathbb{N}}$ of \mathbb{N} such that

$$u^{\alpha_{n_k}} \rightarrow u \quad \text{weakly-star in } \mathcal{X} \quad \text{as } k \rightarrow \infty, \quad (3.18)$$

with some $u \in \mathcal{U}_{\text{ad}}$, and such that (3.16) holds true with $(\mu, \rho) := \mathcal{S}_0(u)$ and a suitable $\xi \in \partial I_{[0,1]}(\rho)$. In particular, the pair $((\mu, \rho, \xi), u)$ is admissible for (\mathbf{P}_0) .

We now aim to prove that $u = \bar{u}$. Once this is shown, the uniqueness result of Theorem 2.1 yields that also $(\mu, \rho, \xi) = (\bar{\mu}, \bar{\rho}, \bar{\xi})$, which then implies that (3.16) holds true. Indeed, we have, owing to the weak sequential lower semicontinuity of \tilde{J} , and in view of the optimality property of $((\bar{\mu}, \bar{\rho}), \bar{u})$ for problem (\mathbf{P}_0) ,

$$\begin{aligned} \liminf_{k \rightarrow \infty} \tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}) &\geq J((\mu, \rho), u) + \frac{1}{2} \|u - \bar{u}\|_{L^2(Q)}^2 \\ &\geq J((\bar{\mu}, \bar{\rho}), \bar{u}) + \frac{1}{2} \|u - \bar{u}\|_{L^2(Q)}^2. \end{aligned} \quad (3.19)$$

On the other hand, the optimality property of $((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}})$ for problem $(\tilde{\mathbf{P}}_{\alpha_{n_k}})$ yields that for any $k \in \mathbb{N}$ we have

$$\tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}) = \tilde{J}(\mathcal{S}_{\alpha_{n_k}}(u^{\alpha_{n_k}}), u^{\alpha_{n_k}}) \leq \tilde{J}(\mathcal{S}_{\alpha_{n_k}}(\bar{u}), \bar{u}), \quad (3.20)$$

whence, taking the limit superior as $k \rightarrow \infty$ on both sides and invoking (3.13) in Corollary 3.2,

$$\begin{aligned} \limsup_{k \rightarrow \infty} \tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}) &\leq \tilde{J}(\mathcal{S}_0(\bar{u}), \bar{u}) \\ &= \tilde{J}((\bar{\mu}, \bar{\rho}), \bar{u}) = J((\bar{\mu}, \bar{\rho}), \bar{u}). \end{aligned} \quad (3.21)$$

Combining (3.19) with (3.21), we have thus shown that $\frac{1}{2} \|u - \bar{u}\|_{L^2(Q)}^2 = 0$, so that in fact $u = \bar{u}$ and thus also $(\mu, \rho, \xi) = (\bar{\mu}, \bar{\rho}, \bar{\xi})$. Moreover, (3.19) and (3.21) also imply that

$$\begin{aligned} J((\bar{\mu}, \bar{\rho}), \bar{u}) &= \tilde{J}((\bar{\mu}, \bar{\rho}), \bar{u}) = \liminf_{k \rightarrow \infty} \tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}) \\ &= \limsup_{k \rightarrow \infty} \tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}) = \lim_{k \rightarrow \infty} \tilde{J}((\mu^{\alpha_{n_k}}, \rho^{\alpha_{n_k}}), u^{\alpha_{n_k}}), \end{aligned} \quad (3.22)$$

which proves (3.17) and, at the same time, also (3.15). The assertion is thus completely proved. \blacksquare

4 The optimality system

In this section, we aim to establish first-order necessary optimality conditions for the optimal control problem (\mathcal{P}_0) . This will be achieved by a passage to the limit as $\alpha \searrow 0$

in the first-order necessary optimality conditions for the adapted optimal control problems $(\tilde{\mathbf{P}}_\alpha)$ that can be derived by arguing as in [22] with minor changes. This procedure will yield certain generalized first-order necessary conditions of optimality in the limit. In this entire section, we assume that $\bar{u} \in \mathcal{U}_{\text{ad}}$ is a fixed optimal control of problem (\mathbf{P}_0) and that $(\bar{\mu}, \bar{\rho}, \bar{\xi})$ is the associated solution to the state system (1.2)–(1.6) established in Theorem 2.1, that is, we have $(\bar{\mu}, \bar{\rho}) = \mathcal{S}_0(\bar{u})$ and $\bar{\xi} \in \partial I_{[0,1]}(\bar{\rho})$ almost everywhere in Q . We begin our analysis by formulating for arbitrary $\alpha \in (0, 1]$ the adjoint state system for the adapted control problem $(\tilde{\mathbf{P}}_\alpha)$ corresponding to \bar{u} . We assume that $\bar{u}^\alpha \in \mathcal{U}_{\text{ad}}$ is an arbitrary optimal control for $(\tilde{\mathbf{P}}_\alpha)$ and that $(\bar{\mu}^\alpha, \bar{\rho}^\alpha) = \mathcal{S}_\alpha(\bar{u}^\alpha)$ is the corresponding solution to the associated state system (2.15)–(2.18), which then enjoys the regularity properties (2.19), (2.20) and fulfills the boundedness conditions (2.21), (2.23) and (2.24). It then follows (see [22, Eqs. (4.3)–(4.6)]) that the corresponding adjoint system has the form

$$\begin{aligned} & - (1 + 2g(\bar{\rho}^\alpha)) p_t^\alpha - g'(\bar{\rho}^\alpha) \bar{\rho}_t^\alpha p^\alpha - \Delta p^\alpha - g'(\bar{\rho}^\alpha) q^\alpha \\ & = \beta_2(\bar{\mu}^\alpha - \mu_Q) \quad \text{in } Q, \end{aligned} \quad (4.1)$$

$$\begin{aligned} & - q_t^\alpha + \varphi(\alpha) h''(\rho^\alpha) q^\alpha + F''(\bar{\rho}^\alpha) q^\alpha - \bar{\mu}^\alpha g''(\bar{\rho}^\alpha) q^\alpha \\ & = -g'(\bar{\rho}^\alpha)(\bar{\mu}_t^\alpha p^\alpha - \bar{\mu}^\alpha p_t^\alpha) - D\mathcal{B}[\bar{\rho}^\alpha]^*(q^\alpha) + \beta_1(\bar{\rho}^\alpha - \rho_Q) \quad \text{in } Q, \end{aligned} \quad (4.2)$$

$$\partial_{\mathbf{n}} p^\alpha = 0 \quad \text{on } \Sigma, \quad (4.3)$$

$$p^\alpha(T) = q^\alpha(T) = 0 \quad \text{in } \Omega, \quad (4.4)$$

where $D\mathcal{B}[\bar{\rho}^\alpha]^* \in \mathcal{L}(L^2(Q), L^2(Q))$ denotes the adjoint operator associated with the operator $D\mathcal{B}[\bar{\rho}^\alpha] \in \mathcal{L}(L^2(Q), L^2(Q))$, which is defined by the identity

$$\int_0^T \int_\Omega D\mathcal{B}[\bar{\rho}^\alpha]^*(v) w \, dx \, dt = \int_0^T \int_\Omega v D\mathcal{B}[\bar{\rho}^\alpha](w) \, dx \, dt \quad \forall v, w \in L^2(Q). \quad (4.5)$$

According to [22, Thm. 4.2], the system (4.1)–(4.3) has for every $\alpha \in (0, 1]$ a unique solution pair (p^α, q^α) such that

$$p^\alpha \in H^1(0, T; H) \cap L^\infty(0, T; V) \cap L^2(0, T; W), \quad q^\alpha \in H^1(0, T; H). \quad (4.6)$$

Moreover, as in [22, Cor. 4.3] it follows that the following variational inequality is satisfied:

$$\int_0^T \int_\Omega (p^\alpha + \beta_3 \bar{u}^\alpha + (\bar{u}^\alpha - \bar{u}))(v - \bar{u}^\alpha) \, dx \, dt \geq 0 \quad \forall v \in \mathcal{U}_{\text{ad}}. \quad (4.7)$$

We now prove an a priori estimate that will be fundamental for the derivation of the optimality conditions for (\mathbf{P}_0) . To this end, we set, for all $\alpha \in (0, 1]$,

$$\lambda^\alpha := \varphi(\alpha) h''(\bar{\rho}^\alpha) q^\alpha, \quad (4.8)$$

and we introduce the function space

$$Y := \{v \in H^1(0, T; H) : v(0) = 0\}, \quad (4.9)$$

which is a Hilbert space when equipped with the standard inner product of $H^1(0, T; H)$. As $H^1(0, T; H)$ is continuously embedded in $C^0([0, T]; H)$, the initial condition encoded

in the definition of Y is meaningful. We denote by Y' the dual space of Y and by $\langle \cdot, \cdot \rangle_Y$ the duality pairing between Y' and Y . Observe that $Y \subset L^2(Q) \subset Y'$ with dense and continuous embeddings, where it is understood that for all $w \in Y$ and $v \in L^2(Q)$ it holds that

$$\langle v, w \rangle_Y = \int_0^T \int_{\Omega} v w \, dx \, dt. \quad (4.10)$$

We have the following result.

LEMMA 4.1. *Suppose that the conditions **(A1)**–**(A4)** are fulfilled. Then there is a constant $K_3^* > 0$, which depends only on the data of problem **(P₀)**, such that, for every $\alpha \in (0, 1]$,*

$$\begin{aligned} & \|p^\alpha\|_{H^1(0,T;H) \cap L^\infty(0,T;V) \cap L^2(0,T;H^2(\Omega))} + \|q^\alpha\|_{L^\infty(0,T;H)} \\ & + \|\partial_t q^\alpha\|_{Y'} + \|\lambda^\alpha\|_{Y'} \leq K_3^*. \end{aligned} \quad (4.11)$$

PROOF: In the following, C_i , $i \in \mathbb{N}$, denote positive constants, which are independent of $\alpha \in (0, 1]$. For the sake of a better readability, we will in the following steps omit the superscript α , writing it only at the final estimate of each step. We also will make repeated use of the global bounds (2.21), (2.23), (2.24) without further reference.

STEP 1: We first add p on both sides of (4.1), multiply the result by $-p_t$, and integrate over Q^t (recall (1.22)), where $0 \leq t < T$. Using the fact that g is nonnegative, we then obtain that

$$\int_t^T \int_{\Omega} |p_t|^2 \, dx \, ds + \frac{1}{2} \|p(t)\|_V^2 \leq \sum_{j=1}^4 I_j, \quad (4.12)$$

where the quantities I_j , $1 \leq j \leq 4$, will be specified and estimated below. At first, we obtain from Young's inequality the estimates

$$I_1 := - \int_t^T \int_{\Omega} p p_t \, dx \, ds \leq \frac{1}{8} \int_t^T \int_{\Omega} |p_t|^2 \, dx \, ds + C_1 \int_t^T \int_{\Omega} |p|^2 \, dx \, ds, \quad (4.13)$$

$$\begin{aligned} I_3 &:= - \int_t^T \int_{\Omega} g'(\bar{\rho}) q p_t \, dx \, ds \\ &\leq \frac{1}{8} \int_t^T \int_{\Omega} |p_t|^2 \, dx \, ds + C_2 \int_t^T \int_{\Omega} |q|^2 \, dx \, ds, \end{aligned} \quad (4.14)$$

$$I_4 := - \int_t^T \int_{\Omega} \beta_2(\bar{\mu} - \mu_Q) p_t \, dx \, ds \leq \frac{1}{8} \int_t^T \int_{\Omega} |p_t|^2 \, dx \, ds + C_3. \quad (4.15)$$

Moreover, owing to Hölder's and Young's inequality, and using the continuity of the embedding $V \subset L^6(\Omega)$,

$$\begin{aligned} I_2 &:= - \int_t^T \int_{\Omega} g'(\bar{\rho}) \bar{\rho}_t p p_t \, dx \, ds \leq C_4 \int_t^T \|p_t(s)\|_2 \|\bar{\rho}_t(s)\|_3 \|p(s)\|_6 \, ds \\ &\leq \frac{1}{8} \int_t^T \int_{\Omega} |p_t|^2 \, dx \, ds + C_5 \int_t^T \|\bar{\rho}_t(s)\|_3^2 \|p(s)\|_V^2 \, ds. \end{aligned} \quad (4.16)$$

Combining the estimates (4.12)–(4.16), we have thus shown that

$$\begin{aligned} & \int_t^T \int_{\Omega} |p_t^\alpha|^2 dx ds + \|p^\alpha(t)\|_V^2 \\ & \leq C_6 \left(1 + \int_t^T \int_{\Omega} (|p^\alpha|^2 + |q^\alpha|^2) dx ds + \int_t^T \|\bar{\rho}_t^\alpha(s)\|_3^2 \|p^\alpha(s)\|_V^2 ds \right). \end{aligned} \quad (4.17)$$

STEP 2: We now multiply (4.2) by q and integrate over Q^t , where $0 \leq t < T$. We obtain that

$$\frac{1}{2} \|q(t)\|_H^2 + \int_t^T \int_{\Omega} (\varphi(\alpha) h''(\bar{\rho}) - \mu g''(\bar{\rho})) |q|^2 dx ds \leq \sum_{j=1}^5 J_j, \quad (4.18)$$

where, since $\varphi(\alpha) h''(\bar{\rho}) \geq 0$ and $\mu g''(\bar{\rho}) \leq 0$, the integral on the left-hand side is non-negative and the quantities J_j , $1 \leq j \leq 5$, will be defined and estimated below. We have

$$J_1 := - \int_t^T \int_{\Omega} F''(\bar{\rho}) |q|^2 dx ds \leq C_7 \int_t^T \int_{\Omega} |q|^2 dx ds, \quad (4.19)$$

as well as, using Young's inequality,

$$\begin{aligned} J_3 & := \int_t^T \int_{\Omega} g'(\bar{\rho}) \mu p_t q dx ds \\ & \leq \frac{1}{4} \int_t^T \int_{\Omega} |p_t|^2 dx ds + C_8 \int_t^T \int_{\Omega} |q|^2 dx ds, \end{aligned} \quad (4.20)$$

$$J_5 := \int_t^T \int_{\Omega} \beta_1 (\bar{\rho} - \rho_Q) q dx ds \leq \int_t^T \int_{\Omega} |q|^2 dx ds + C_9, \quad (4.21)$$

$$\begin{aligned} J_4 & := - \int_t^T \int_{\Omega} D\mathcal{B}[\bar{\rho}]^*(q) q dx ds = - \int_t^T \int_{\Omega} D\mathcal{B}[\bar{\rho}](q) q dx ds \\ & \leq C_{10} \int_t^T \int_{\Omega} |q|^2 dx ds, \end{aligned} \quad (4.22)$$

where in the last estimate we have used (4.5) and (2.7). Finally, we take advantage of the continuity of the embedding $H^2(\Omega) \subset L^\infty(\Omega)$ and of Hölder's and Young's inequality to conclude that, for every $\gamma > 0$ (to be chosen later),

$$\begin{aligned} J_2 & := - \int_t^T \int_{\Omega} g'(\bar{\rho}) \bar{\mu}_t p q dx ds \leq C_{11} \int_t^T \|\bar{\mu}_t(s)\|_2 \|p(s)\|_\infty \|q(s)\|_2 ds \\ & \leq \gamma \int_t^T \|p(s)\|_{H^2(\Omega)}^2 ds + \frac{C_{12}}{\gamma} \int_t^T \|\bar{\mu}_t(s)\|_2^2 \|q(s)\|_2^2 ds. \end{aligned} \quad (4.23)$$

In view of the boundary condition (4.3), we can infer from classical elliptic estimates that

$$\|p(s)\|_{H^2(\Omega)} \leq C_\Omega (\|p(s)\|_H + \|\Delta p(s)\|_H) \quad \text{for a. e. } s \in (0, T), \quad (4.24)$$

where the constant $C_\Omega > 0$ depends only on the domain Ω .

Combining all the estimates (4.18)–(4.24), we have therefore shown that

$$\begin{aligned} \|q^\alpha(t)\|_H^2 &\leq C_{13} + \frac{1}{2} \int_t^T \int_\Omega |p_t^\alpha|^2 dx ds + \gamma C_{14} \int_t^T \int_\Omega (|p^\alpha|^2 + |\Delta p^\alpha|^2) dx ds \\ &\quad + C_{15} (1 + \gamma^{-1}) \int_t^T \int_\Omega (\|\bar{\mu}_t^\alpha(s)\|_2^2 + 1) |q^\alpha|^2 dx ds. \end{aligned} \quad (4.25)$$

STEP 3: We now use (4.1) to estimate Δp^α directly. Indeed, we have

$$\begin{aligned} \int_t^T \int_\Omega |\Delta p^\alpha|^2 dx ds &= \|\Delta p^\alpha\|_{L^2(Q^t)}^2 \\ &\leq C_{16} + C_{17} \left(\|p_t^\alpha\|_{L^2(Q^t)}^2 + \|q^\alpha\|_{L^2(Q^t)}^2 + \|\bar{\rho}_t^\alpha p^\alpha\|_{L^2(Q^t)}^2 \right). \end{aligned} \quad (4.26)$$

Now, owing to (1.20), it holds that

$$\|\bar{\rho}_t^\alpha p^\alpha\|_{L^2(Q^t)}^2 \leq C_{18} \int_t^T \|\bar{\rho}_t^\alpha(s)\|_3^2 \|p^\alpha(s)\|_V^2 ds. \quad (4.27)$$

In conclusion, we have shown that

$$\begin{aligned} \int_t^T \int_\Omega |\Delta p^\alpha|^2 dx ds &\leq C_{16} + C_{17} \int_t^T \int_\Omega (|p_t^\alpha|^2 + |q^\alpha|^2) dx ds \\ &\quad + C_{19} \int_t^T \|\bar{\rho}_t^\alpha(s)\|_3^2 \|p^\alpha(s)\|_V^2 ds. \end{aligned} \quad (4.28)$$

STEP 4: At this point, we are in a position to combine the estimates (4.17), (4.25) and (4.28). Choosing $\gamma > 0$ appropriately small (such that $0 < \gamma C_{14} C_{17} < \frac{1}{2}$), and noting that the global estimates (2.24) imply that the mappings $s \mapsto \|\bar{\rho}_t^\alpha(s)\|_3^2$ and $\|\bar{\mu}_t^\alpha(s)\|_2^2$ are bounded in $L^1(0, T)$ independently of $\alpha \in (0, 1]$, we conclude from Gronwall's lemma (taken backward in time) the estimate

$$\|p^\alpha\|_{H^1(0, T; H) \cap L^\infty(0, T; V) \cap L^2(0, T; H^2(\Omega))} + \|q^\alpha\|_{L^\infty(0, T; H)} \leq C_{20}. \quad (4.29)$$

Now assume that $v \in Y$ is arbitrary. As $q^\alpha(T) = v(0) = 0$, integration by parts with respect to time (which is permitted since $q^\alpha, v \in H^1(0, T; H)$) yields

$$\begin{aligned} |\langle q_t^\alpha, v \rangle_Y| &= \left| \int_0^T \int_\Omega q_t^\alpha v dx dt \right| = \left| - \int_0^T \int_\Omega q^\alpha v_t dx dt \right| \\ &\leq C_{21} \|q^\alpha\|_{L^\infty(0, T; H)} \|v\|_{H^1(0, T; H)}, \end{aligned} \quad (4.30)$$

and it follows from (4.29) that

$$\|q_t^\alpha\|_{Y'} \leq C_{22}. \quad (4.31)$$

It remains to show the bound for λ^α (see (4.8)). To this end, notice that (4.2) yields

$$\begin{aligned} \lambda^\alpha &= q_t^\alpha - g'(\bar{\rho}^\alpha) \bar{\mu}_t^\alpha p^\alpha - F''(\bar{\rho}^\alpha) q^\alpha + \bar{\mu}^\alpha g''(\bar{\rho}^\alpha) q^\alpha + \bar{\mu}^\alpha g'(\bar{\rho}^\alpha) p_t^\alpha \\ &\quad - D\mathcal{B}[\bar{\rho}^\alpha]^*(q^\alpha) + \beta_1(\bar{\rho}^\alpha - \rho_Q), \end{aligned} \quad (4.32)$$

where it easily follows from the estimates (2.23), (2.24) and (4.29) that the last five summands on the right-hand side are bounded in $L^2(Q)$ (and thus in Y') independently of $\alpha \in (0, 1]$. To estimate the remaining term, let $v \in Y$ be arbitrary and observe that (2.19), (4.6), the continuous embedding $V \subset L^6(\Omega)$ and Hölder's inequality imply

$$\bar{\mu}_t^\alpha p^\alpha \in L^\infty(0, T; L^{3/2}(\Omega)) \cap L^1(0, T; L^3(\Omega)) \subset L^2(0, T; L^2(\Omega)),$$

whence also $g'(\bar{\rho}^\alpha) \bar{\mu}_t^\alpha p^\alpha \in L^2(Q)$. Therefore, using (4.10), (2.24) and (4.29), as well as the continuity of the embeddings $H^2(\Omega) \subset L^\infty(\Omega)$ and $H^1(0, T; H) \subset C^0([0, T]; H)$, we have

$$\begin{aligned} | \langle -g'(\bar{\rho}^\alpha) \bar{\mu}_t^\alpha p^\alpha, v \rangle_{Y'} | &= \left| \int_0^T \int_\Omega g'(\bar{\rho}^\alpha) \bar{\mu}_t^\alpha p^\alpha v \, dx \, dt \right| \\ &\leq C_{23} \int_0^T \|\bar{\mu}_t^\alpha(t)\|_2 \|p^\alpha(t)\|_\infty \|v(t)\|_2 \, dt \\ &\leq C_{24} \max_{0 \leq t \leq T} \|v(t)\|_H \|\bar{\mu}_t^\alpha\|_{L^2(Q)} \|p^\alpha\|_{L^2(0, T; H^2(\Omega))} \\ &\leq C_{25} \|v\|_{H^1(0, T; H)}. \end{aligned}$$

In consequence, we have that

$$\|\lambda^\alpha\|_{Y'} \leq C_{26}, \quad (4.33)$$

which concludes the proof of the assertion. \blacksquare

We draw some consequences from the previously established results. Assume that $\{\alpha_n\} \subset (0, 1]$ satisfies $\alpha_n \searrow 0$ as $n \rightarrow \infty$. Then, thanks to Theorem 3.4, there is a subsequence, without loss of generality $\{\alpha_n\}$ itself, such that, for any $n \in \mathbb{N}$, we can find an optimal control $\bar{u}^{\alpha_n} \in \mathcal{U}_{\text{ad}}$ for $(\tilde{\mathbf{P}}_{\alpha_n})$ and an associated state $(\bar{\mu}^{\alpha_n}, \bar{\rho}^{\alpha_n})$ such that the convergence given by (3.15) and (3.17) (where one reads α_n in place of α_{n_k}) holds true. As in the proof of Theorem 3.1, we may without loss of generality assume that, for $1 \leq q < 6$ and $\Phi \in \{F''', g, g', g''\}$,

$$\Phi(\bar{\rho}^{\alpha_n}) \rightarrow \Phi(\bar{\rho}) \quad \text{strongly in } C^0([0, T]; L^q(\Omega)). \quad (4.34)$$

Also, by virtue of Lemma 4.1, we may without loss of generality assume that the corresponding adjoint state variables $(p^{\alpha_n}, q^{\alpha_n})$ satisfy

$$p^{\alpha_n} \rightharpoonup p \quad \text{weakly-star in } H^1(0, T; H) \cap L^\infty(0, T; V) \cap L^2(0, T; W), \quad (4.35)$$

$$q^{\alpha_n} \rightharpoonup q \quad \text{weakly-star in } L^\infty(0, T; H), \quad (4.36)$$

$$\lambda^{\alpha_n} \rightharpoonup \lambda \quad \text{weakly in } Y', \quad (4.37)$$

for a suitable triple (p, q, λ) . Therefore, passing to the limit as $n \rightarrow \infty$ in the variational inequality (4.7), written for α_n , $n \in \mathbb{N}$, we obtain that p satisfies

$$\int_0^T \int_\Omega (p + \beta_3 \bar{u})(v - \bar{u}) \, dx \, dt \geq 0 \quad \forall v \in \mathcal{U}_{\text{ad}}. \quad (4.38)$$

Next, we will show that in the limit as $n \rightarrow \infty$ a limiting adjoint system for (\mathcal{P}_0) is satisfied. To this end, we note that it is not difficult to check that

$$\begin{aligned} g(\bar{\rho}^{\alpha_n}) p_t^{\alpha_n} &\rightarrow g(\bar{\rho}) p_t, & g'(\bar{\rho}^{\alpha_n}) \bar{\rho}_t^{\alpha_n} p^{\alpha_n} &\rightarrow g'(\bar{\rho}) \bar{\rho}_t p, \\ g'(\bar{\rho}^{\alpha_n}) q^{\alpha_n} &\rightarrow g'(\bar{\rho}) q, & &\text{all weakly in } L^1(Q). \end{aligned} \quad (4.39)$$

Hence, passage to the limit as $n \rightarrow \infty$ in (4.1) (written for α_n) yields

$$\begin{aligned} &-(1 + 2g(\bar{\rho})) p_t - g'(\bar{\rho}) \bar{\rho}_t p - \Delta p - g'(\bar{\rho}) q \\ &= \beta_2(\bar{\mu} - \mu_Q) \quad \text{a. e. in } Q. \end{aligned} \quad (4.40)$$

We also have that (see (4.35))

$$\partial_{\mathbf{n}} p = 0 \quad \text{a. e. on } \Sigma, \quad p(T) = 0 \quad \text{a. e. in } \Omega. \quad (4.41)$$

In order to derive an equation resembling (4.2), we multiply (4.2) (written for α_n) by an arbitrary element v belonging to the space

$$Y_0 := \{v \in C^1([0, T]; C(\bar{\Omega})) : v(\cdot, 0) = 0\},$$

which is a dense subset of Y . Integrating over Q and by parts with respect to t , we then obtain the equation

$$\begin{aligned} &\int_0^T \int_{\Omega} \lambda^{\alpha_n} v \, dx \, dt + \int_0^T \int_{\Omega} q^{\alpha_n} v_t \, dx \, dt + \int_0^T \int_{\Omega} F''(\bar{\rho}^{\alpha_n}) q^{\alpha_n} v \, dx \, dt \\ &- \int_0^T \int_{\Omega} \bar{\mu}^{\alpha_n} g''(\bar{\rho}^{\alpha_n}) q^{\alpha_n} v \, dx \, dt + \int_0^T \int_{\Omega} g'(\bar{\rho}^{\alpha_n}) \bar{\mu}_t^{\alpha_n} p^{\alpha_n} v \, dx \, dt \\ &= \int_0^T \int_{\Omega} g'(\bar{\rho}^{\alpha_n}) \bar{\mu}^{\alpha_n} p_t^{\alpha_n} v \, dx \, dt - \int_0^T \int_{\Omega} D\mathcal{B}[\bar{\rho}^{\alpha_n}]^*(q^{\alpha_n}) v \, dx \, dt \\ &+ \int_0^T \int_{\Omega} \beta_1(\bar{\rho}^{\alpha_n} - \rho_Q) v \, dx \, dt. \end{aligned} \quad (4.42)$$

By virtue of the previously established convergence properties of the involved sequences, it is not difficult to show that we may pass to the limit as $n \rightarrow \infty$ in all of the summands occurring in (4.42) up to the penultimate one. We may leave the details of these straightforward calculations to the reader. For the penultimate summand, we have

$$\begin{aligned} &\int_0^T \int_{\Omega} (D\mathcal{B}[\bar{\rho}^{\alpha_n}]^*(q^{\alpha_n}) - D\mathcal{B}[\bar{\rho}]^*(q)) v \, dx \, dt \\ &= \int_0^T \int_{\Omega} q^{\alpha_n} (D\mathcal{B}[\bar{\rho}^{\alpha_n}](v) - D\mathcal{B}[\bar{\rho}](v)) \, dx \, dt \\ &+ \int_0^T \int_{\Omega} (q^{\alpha_n} - q) D\mathcal{B}[\bar{\rho}](v) \, dx \, dt. \end{aligned}$$

While the second summand tends to zero as $n \rightarrow \infty$, nothing can be said about the first one: we need an additional assumption. The following condition is obviously sufficient to guarantee that also the first summand approaches zero as $n \rightarrow \infty$:

(A5) The mapping $v \mapsto D\mathcal{B}[v]$ is continuous from $L^2(Q)$ into $\mathcal{L}(L^2(Q), L^2(Q))$.

In conclusion, if **(A5)** is valid, then the passage to the limit as $n \rightarrow \infty$ results in the following identity:

$$\begin{aligned}
& \langle \lambda, v \rangle_Y + \int_0^T \int_{\Omega} q v_t dx dt + \int_0^T \int_{\Omega} F''(\bar{\rho}) q v dx dt \\
& - \int_0^T \int_{\Omega} \bar{\mu} g''(\bar{\rho}) q v dx dt + \int_0^T \int_{\Omega} g'(\bar{\rho}) \bar{\mu}_t p v dx dt \\
& = \int_0^T \int_{\Omega} g'(\bar{\rho}) \bar{\mu} p_t v dx dt - \int_0^T \int_{\Omega} D\mathcal{B}[\bar{\rho}]^*(q) v dx dt \\
& + \int_0^T \int_{\Omega} \beta_1(\bar{\rho} - \rho_Q) v dx dt \quad \forall v \in Y_0.
\end{aligned} \tag{4.43}$$

We claim that the variational equation (4.43) holds in fact true for all $v \in Y$. To see this, we employ a standard density argument. Indeed, if $v \in Y$ is given, then there is some sequence $\{v_n\} \subset Y_0$ such that $v_n \rightarrow v$ in the norm of $H^1(0, T; H)$. Now observe that (4.43) is valid for $v = v_n$. It is now easily checked that we may pass to the limit as $n \rightarrow \infty$ in each term occurring in (4.43). As an example, we give the details for the most difficult term, which is the last one on the left-hand side. We have, denoting by C_i , $i \in \mathbb{N}$, constants that do not depend on n :

$$\begin{aligned}
& \left| \int_0^T \int_{\Omega} g'(\bar{\rho}) \bar{\mu}_t p (v - v_n) dx dt \right| \leq C_1 \int_0^T \|\bar{\mu}_t(t)\|_2 \|p(t)\|_{\infty} \|v(t) - v_n(t)\|_2 dt \\
& \leq C_2 \max_{0 \leq t \leq T} \|v(t) - v_n(t)\|_H \|\bar{\mu}_t\|_{L^2(Q)} \|p\|_{L^2(0, T; H^2(\Omega))} \\
& \leq C_3 \|v - v_n\|_{H^1(0, T; H)} \rightarrow 0 \quad \text{as } n \rightarrow \infty.
\end{aligned}$$

Next, we show that the limit λ satisfies some sort of a complementarity slackness condition. Indeed, we have

$$\liminf_{n \rightarrow \infty} \int_0^T \int_{\Omega} \lambda^{\alpha_n} q^{\alpha_n} dx dt = \liminf_{n \rightarrow \infty} \int_0^T \int_{\Omega} \varphi(\alpha_n) h''(\bar{\rho}^{\alpha_n}) |q^{\alpha_n}|^2 dx dt \geq 0. \tag{4.44}$$

Moreover, there is some indication that the limit λ should somehow be concentrated on the set where $\bar{\rho} = 0$ or $\bar{\rho} = 1$ (which, however, we cannot prove rigorously). To this end, we test λ^{α_n} by the function $\bar{\rho}^{\alpha_n}(1 - \bar{\rho}^{\alpha_n})\phi$, where ϕ is any smooth test function satisfying $\phi(0) = 0$. By recalling that $h''(r) = \frac{1}{r(1-r)}$, we then obtain that

$$\lim_{n \rightarrow \infty} \int_0^T \int_{\Omega} \lambda^{\alpha_n} \bar{\rho}^{\alpha_n} (1 - \bar{\rho}^{\alpha_n}) \phi dx dt = \lim_{n \rightarrow \infty} \int_0^T \int_{\Omega} \varphi(\alpha_n) q^{\alpha_n} \phi dx dt = 0. \tag{4.45}$$

We now collect the results established above. We have the following statement.

THEOREM 4.2: *Let the assumptions **(A1)**–**(A5)** be satisfied, and let $\bar{u} \in \mathcal{U}_{\text{ad}}$ be an optimal control for **(P₀)** with the associated solution $(\bar{\mu}, \bar{\rho}, \bar{\xi})$ to the state system (1.2)–(1.6) in the sense of Theorem 2.1. Moreover, let $\{\alpha_n\} \subset (0, 1]$ with $\alpha_n \searrow 0$ as $n \rightarrow \infty$ be*

such that there are optimal pairs $((\bar{\mu}^{\alpha_n}, \bar{\rho}^{\alpha_n}), \bar{u}^{\alpha_n})$ for the adapted problem $(\tilde{\mathbf{P}}_{\alpha_n})$ satisfying (3.15)–(3.17) (such sequences exist by virtue of Theorem 3.4) and having the associated adjoint variables $\{(p^{\alpha_n}, q^{\alpha_n})\}$. Then, for any subsequence $\{n_k\}_{k \in \mathbb{N}}$ of \mathbb{N} , there are a subsequence $\{n_{k_\ell}\}_{\ell \in \mathbb{N}}$ and some triple (p, q, λ) such that

- $p \in H^1(0, T; H) \cap L^\infty(0, T; V) \cap L^2(0, T; W)$, $q \in L^\infty(0, T; H)$, and $\lambda \in Y'$,
- the relations (4.35)–(4.37), (4.44), and (4.45) are valid (where the sequences are indexed by n_{k_ℓ} and the limits are taken for $\ell \rightarrow \infty$), and
- the variational inequality (4.38) and the adjoint system equations (4.40), (4.41) and (4.43) are satisfied, where (4.43) holds for every $v \in Y$.

REMARK: We are unable to show that the limit triple (p, q, λ) solving the adjoint problem associated with the optimal pair $((\bar{\mu}, \bar{\rho}), \bar{u})$ is uniquely determined. Therefore, it may well happen that the limiting pairs differ for different subsequences. However, it follows from the variational inequality (4.38) that, for any such limit, it holds with the orthogonal projection $\mathbb{P}_{\mathcal{U}_{\text{ad}}}$ onto \mathcal{U}_{ad} with respect to the standard inner product in $L^2(Q)$ that in the case $\beta_3 > 0$ we have $\bar{u} = \mathbb{P}_{\mathcal{U}_{\text{ad}}}(-\beta_3^{-1}p)$.

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