Weierstraß-Institut für Angewandte Analysis und Sto
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im Fors
hungsverbund Berlin e.V.

Preprint ISSN 0946 - 8633

GEOMS: A software package for the numerical integration of general model equations of multibody systems

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submitted: September 25, ²⁰⁰⁷

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> No. 1259 Berlin 2007

²⁰⁰⁰ Mathematics Subject Classification. 70E55, 65L80.

This work has been supported by DFG resear
h grant no. Me790/13 during the stay of the author at TU Berlin, Fa
hgebiet Numeris
he Mathematik. The software pa
kage GEOMS is part of the PhD thesis [38] of the author developed at the TU Berlin.

Key words and phrases. differential-algebraic equations, equations of motion, multibody system, numeri
al integration, simulation.

Edited by Weierstraÿ-Institut für Angewandte Analysis und Sto
hastik (WIAS) Mohrenstraÿe 39 10117 Berlin Germany

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Abstract

In this paper we present the new numeri
al algorithm GEOMS for the numeri
al integration of the most general form of the equations of motion of multibody systems, in
luding nonholonomi onstraints and possible redundan
ies in the onstraints, as they may appear in industrial appli
ations. Besides the numeri
al integration it oers some additional features like stabilization of the model equations, use of dierent de
omposition strategies, or he
king and orre
tion of the initial values with respe
t to their onsisten
y. Furthermore, GEOMS preserves hidden onstraints and (possibly) existing solution invariants if they are provided as equations.

We will also demonstrate the performan
e and the appli
ability of GEOMS for two me
hani
al examples of dierent degrees of omplexity.

1 Introduction

The multibody system (MBS) approach is frequently used in industrial simulation pa
kages in roboti
s, vehi
le system dynami
s, and biome
hani
s. A multibody system model consists of a finite number of rigid or elastic bodies and their interconnections like, e.g., joints, springs, dampers, and actuators. The *equations of* motion may be generated in a systematic way by multibody formalisms that are based on the principles of classical mechanics [35].

The efficient and robust numerical integration of these equations is a challenging problem in the development of simulation pa
kages, sin
e dynami
al simulation is frequently used and one of the most time onsuming analysis methods for MBS models. The equations of motion with nonredundant onstraints form a nonlinear system of differential-algebraic equations $(DAEs)$ of differentiation index $(d$ -index) 3, see $[7, 10, 13, 18]$. It is well known that the numerical treatment of DAEs of high index or higher index, i.e., d-index 2 or larger than 2, respectively, is nontrivial in general. Effects arising in the numerical treatment are, for example, drift, instabilities, convergence problems, or inconsistencies. These difficulties in the numerical solution of such high index problems are discussed in $[4, 12, 14, 15, 17, 18, 24, 29, 30]$. However, the equations of motion are DAEs with a very special structure that should be exploited in the numerical solution $[7, 18]$.

In this report we will present the new software package **GEOMS** for the numerical integration of general equations of motion of multibody systems in des
riptor form. In contrast to standard textbook presentations like [18], we do not restrict ourselves to lassi
al onstrained me
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al systems but onsider the more omplex model equations that are actually used in state-of-the-art MBS simulation packages $[7, 34]$.

The software pa
kage GEOMS is suited for general equations of motion involving dynamical force elements, contact conditions, and (possibly) redundant holonomic as well as nonholonomic constraints. Furthermore, the package takes into account possibly existing information concerning solution invariants, e.g., energy conservation. The code is based on residual evaluations, i.e., the system need not be given completely in explicit form. It is sufficient that the right-hand side of the equations of motion and the mass matrix are specified.

Although, the pa
kage GEOMS is able to treat also redundant onstraints, in this paper we will restri
t our onsiderations to regular equations of motion, i.e., the onstraints are assumed to be nonredundant. For more details on equations of motion with redundant constraints we refer to $[26, 38]$.

As base of the integration method GEOMS we will propose a remodeling of the equations of motions. The aim of this remodeling is to determine an equivalent formulation, the so called *projected-strangeness free form*, which has d-index 1 but has the same set of solutions as the original equations of motion. Because of the reduced d-index, the numeri
al treatment of the proje
ted-strangeness free form by use of implicit ODE methods is not affected by instabilities arising from the higher index. Furthermore, all (hidden) constraints are preserved such that no drift-off effects arise in the numerical treatment. The proposed remodeling can be seen as regularization of the equations of motion. For more details on the regularization of equations of motion we refer to [38]. The integration method implemented in GEOMS combines an implicit Runge-Kutta-Method of order 5 with this regularization technique.

The report is organized as follows. In Section 2 we introduce the equations of motion which we want to treat numerically and we discuss the remodeling to the projectedstrangeness-free form which will be used for the discretization in GEOMS. In Section 3 we introduce the code GEOMS and we discuss its features and its applicability in detail. In Section 4 we demonstrate the properties of the software package GEOMS by two numeri
al examples. For the usage and implementation of GEOMS the manual is presented in Appendix A.

2 The Equations of Motion and their Remodeling

Here and in the following we will use the following notation.

Notation 2.1 Let f be a differentiable function $f : \mathbb{X} \to \mathbb{R}^m$, $\mathbb{X} \subset \mathbb{R}^n$, and let x be a differentiable function $x : \mathbb{I} \to \mathbb{X}$, where \mathbb{I} is an open interval in \mathbb{R} . The ith (total) derivative of $x(t)$ with respect to t is denoted by $x^{(i)}(t) = d^i x(t)/dt^i$ for $i \in \mathbb{N}_0$. Note the convention $x^{(0)}(t) = x(t), x^{(1)}(t) = \dot{x}(t)$, and $x^{(2)}(t) = \ddot{x}(t)$. The (partial) derivative of $f(x)$ with respect to x is denoted by $f_{,x}(x) = \frac{\partial}{\partial x} f(x)$. The same notation is used for differentiable vector and matrix functions. The set of *l*-times continuously differentiable functions from X to Y is denoted by $\mathcal{C}^l(X, Y)$.

In the following we investigate a spatial multibody system with holonomi as well as nonholonomic constraints [19, 33]. More precisely we consider the following initial value problem on the domain $\mathbb{I} = [t_0, t_f]$ consisting of the equations of motion in the form

$$
\dot{p} = Z(p)v,\tag{1a}
$$

$$
M(p,t)\dot{v} = f(p,v,r,w,s,\lambda,\mu,t) - Z^{T}(p)G^{T}(p,s,t)\lambda - Z^{T}(p)H^{T}(p,s,t)\mu,
$$
 (1b)

$$
\dot{r} = b(p, v, r, w, s, \lambda, \mu, t), \qquad (1c)
$$

$$
0 = d(p, v, r, w, s, \lambda, \mu, t), \qquad (1d)
$$

$$
0 = c(p, s, t), \tag{1e}
$$

$$
0 = H(p, s, t)Z(p)v + h(p, s, t) = \check{h}(p, v, s, t)),
$$
\n(1f)

$$
0 = g(p, s, t) \tag{1g}
$$

with the initial values

$$
p(t_0) = p_0 \in \mathbb{R}^{n_p}, \ v(t_0) = v_0 \in \mathbb{R}^{n_v}, \ r(t_0) = r_0 \in \mathbb{R}^{n_r}, \ w(t_0) = w_0 \in \mathbb{R}^{n_w}, s(t_0) = s_0 \in \mathbb{R}^{n_s}, \ \lambda(t_0) = \lambda_0 \in \mathbb{R}^{n_\lambda}, \ \mu(t_0) = \mu_0 \in \mathbb{R}^{n_\mu}.
$$
 (2)

Here, the *position vector* p contains arbitrary position coordinates of the multibody system. The Euler-Lagrange formalism for modeling multibody systems yields the equations of motion in se
ond order form. In order to transform the se
ond order system to an equivalent first order system we introduce a velocity vector v and get the relation (1a) between the *generalized velocities* \dot{p} and the velocities v with a matrix $Z(p)$, that determines the angular velocities. The equations (1a) are called kinematic equations. The transformation matrix $Z(p)$ occurs only if there are rotations in three dimensional spa
e, it may be determined by Poisson's kinemati
al equations [1, 7]. In the two dimensional case we have $Z(p) = I$, $\dot{p} = v$.

The equations (1b) are called *dynamic equations of motion*. They follow from the equilibrium of forces and momenta and include the mass matrix $M(p)$, the vector of the applied and gyroscopic forces $f(p, v, r, w, s, \lambda, \mu, t)$, the constraint matrices $G(p, s, t)$ and $H(p, s, t)$ of the holonomic and nonholonomic constraints, respectively, which contain the inaccessible directions of motion column-wise, the associated constraint forces $G^T(p,s,t)\lambda$ and $H^T(p,s,t)\mu$, and the Lagrange multipliers λ and μ . The holonomic constraint matrix is defined as $G(p, s, t) = \frac{d}{dp}g(p, s(p, t), t)$. The mass matrix $M(p)$ is positive semi-definite, since the kinetic energy is a nonnegative quadrati form, and in
ludes the inertia properties of the multibody system.

In a real multibody system, there are often *dynamic force elements* which are described by the vector r and determined by equations $(1c)$, see [7].

Furthermore, not all constraints of a multibody system are directly described by the position variables p or the velocity variables v , but depend on certain *contact* points with coordinates s on the surface of some bodies. The relationship between these contact point coordinates s and the position variables p are given by (1e). Furthermore, the equations of motion are affected by the n_{λ} holonomic constraints (1g) and n_{μ} nonholonomic constraints (1f). These constraints are also called the holonomic constraints on position level and the nonholonomic constraints on velocity level, respectively. Sometimes, force laws and constraints may be formulated more conveniently using *auxiliary variables* w that are implicitly defined by the n_w

possibly nonlinear equation (1d).

Here, $n = n_p + n_v + n_r + n_w + n_s + n_\lambda + n_\mu$ denotes the number of unknown variables. Furthermore, many motions of me
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al systems have known solution invariants, i.e., relations which are satisfied along any motion of the mechanical system, like the invariance of the total energy, momentum, or impulse. Let us denote the m_e equations des
ribing su
h solution invariants by

$$
0 = e(p, v, s, t). \tag{3}
$$

In particular, conservative multibody systems are energy conserving. In this case the total energy is onstant along every motion of the system. For more details on solution invariants we refer to [38].

The theoreti
al basis of the ode GEOMS is based on the following assumptions.

Assumption 2.2 Consider the equations of motion (1). Then the matrices

$$
d_{,w},\tag{4a}
$$

$$
b) \t\t c_{,s}, \t\t (4b)
$$

$$
c) \t\t M \t\t (4c)
$$

$$
d) \left[\begin{array}{cc} GZM^{-1}G_{\lambda} & GZM^{-1}H_{\mu} \\ HZM^{-1}G_{\lambda} & HZM^{-1}H_{\mu} \end{array} \right] \tag{4d}
$$

are assumed to be nonsingular with a bounded inverse for all $(p, v, r, w, s, \lambda, \mu, t) \in$ M , see (10), where

$$
G_{\lambda} = Z^T G^T - f_{,\lambda} + f_{,w} d_{,w}^{-1} d_{,\lambda}, \qquad (5)
$$

$$
H_{\mu} = Z^T H^T - f_{,\mu} + f_{,w} d_{,w}^{-1} d_{,\mu}.
$$
\n(6)

Furthermore, it is assumed that

$$
d \in C^1(\mathbb{M}, \mathbb{R}^{n_w}), c \in C^1(\mathbb{M}, \mathbb{R}^{n_s}), \check{h} \in C^2(\mathbb{M}, \mathbb{R}^{n_\mu}), g \in C^3(\mathbb{M}, \mathbb{R}^{n_\lambda}).
$$

Remark 2.3 a) The nonsingularity of the mass matrix M is assumed only for reasons of simplicity. It is not necessary for the successful numerical integration with GEOMS.

b) Furthermore, note that in Assumption 2.2 redundant onstraints are ex
luded. Redundant onstraints may result in a nonuniqueness of the Lagrange multipliers. Nevertheless, GEOMS is able to deal with ertain types of redundant onstraints. For more details on redundant constraints see [26, 38]. \triangleleft

Using the equations of motion (1) , the first and second derivatives with respect to t of the holonomic constraints $(1g)$ are given by

$$
0 = g^{I}(p, v, s, t) = \frac{d}{dt}g(p, s, t)
$$
\n(7a)

$$
= GZv + g_{,t} - g_{,s}c_{,s}^{-1}c_{,t}
$$
 (7b)

and

$$
0 = g^I(p, v, r, w, s, \lambda, \mu, t) = \frac{d^2}{dt^2} g(p, s, t)
$$
\n(8a)

$$
= (g_{,p}^I - g_{,s}^I c_{,s}^{-1} c_{,p}) Zv + GZM^{-1}(f - Z^T G^T \lambda - Z^T H^T \mu) + g_{,t}^I - g_{,s}^I c_{,s}^{-1} c_{,t}.
$$
 (8b)

They are called *holonomic constraints on velocity level* (7a) and *holonomic con*straints on acceleration level (8a), respectively. The first derivative with respect to t of the nonholonomic constraints (1f) is given by

$$
0 = h^{I}(p, v, r, w, s, \lambda, \mu, t) = \frac{d}{dt}(H(p, s, t)Z(p)v + h(p, s, t))
$$
\n(9a)

$$
= (\breve{h}_{,p} - \breve{h}_{,s}c_{,s}^{-1}c_{,p})Zv + HZM^{-1}(f - Z^{T}G^{T}\lambda - Z^{T}H^{T}\mu) + (\breve{h}_{,t} - \breve{h}_{,s}c_{,s}^{-1}c_{,t})
$$
(9b)

which are called *nonholonomic constraints on acceleration level* (9a).

The holonomic constraints on velocity level and on acceleration level in form (7b) and (8b), respectively, as well as the nonholonomic constraints on acceleration level in form (9b) turn out to be the hidden onstraints of the equations of motion, see [38]. The choice of values $(p, v, r, w, s, \lambda, \mu, t) \in \mathbb{R}^n \times \mathbb{I}$ is restricted by all constraints including the hidden constraints, i.e., (1d)-(1g), (7b), (8b), and (9b). Values which satisfy all of these constraints are called *consistent* and we get the set of consistency

$$
\mathbb{M} = \{ (p, v, r, w, s, \lambda, \mu, t) \in \mathbb{R}^n \times \mathbb{I} : 0 = d(p, v, r, w, s, \lambda, \mu, t),
$$
\n
$$
0 = c(p, s, t),
$$
\n
$$
0 = H(p, s, t)Z(p)v + h(p, s, t),
$$
\n
$$
0 = g(p, s, t),
$$
\n
$$
0 = h^I(p, v, r, w, s, \lambda, \mu, t),
$$
\n
$$
0 = g^I(p, v, s, t),
$$
\n
$$
0 = g^I(p, v, r, w, s, \lambda, \mu, t) \}.
$$
\n(10)

Theorem 2.4 Let the equations of motion (1) satisfy Assumptions 2.2. Then there exist matrix functions $S_p \in C^0(\mathbb{M}, \mathbb{R}^{n_{f_p}, n_p})$ and $S_v \in C^0(\mathbb{M}, \mathbb{R}^{n_{f_v}, n_v})$ with $n_{f_p} =$ $n_p - n_\lambda$ and $n_{f_v} = n_v - n_\lambda - n_\mu$ such that the matrix functions

$$
\begin{bmatrix} S_p(p,t) \\ G(p,t) \end{bmatrix} and \begin{bmatrix} S_v(p,t)M(p,t) \\ G(p,t)Z(p) \\ H(p,t)Z(p) \end{bmatrix} are nonsingular
$$
 (11)

for all $(p, v, r, w, s, \lambda, \mu, t) \in \mathbb{M}$. Then the differential-algebraic system

$$
S_p(p,t)\dot{p} = S_p(p,t)Z(p)v,\tag{12a}
$$

$$
S_v(p,t)M(p,t)\dot{v} = S_v(p,t)f(p,v,r,w,s,\lambda,\mu,t)
$$
\n(12b)

$$
-S_v(p,t)Z^T(p)G^T(p,s,t)\lambda - S_v(p,t)Z^T(p)H^T(p,s,t)\mu,
$$

$$
\dot{r} = b(p, v, r, w, s, \lambda, \mu, t), \qquad (12c)
$$

$$
0 = d(p, v, r, w, s, \lambda, \mu, t), \qquad (12d)
$$

$$
0 = c(p, s, t), \tag{12e}
$$

$$
0 = H(p, s, t)Z(p)v + h(p, s, t),
$$
\n(12f)

$$
0 = g(p, s, t), \tag{12g}
$$

$$
0 = hI(p, v, r, w, s, \lambda, \mu, t), \qquad (12h)
$$

$$
0 = g^{I}(p, v, s, t), \tag{12i}
$$

$$
0 = g^I(p, v, r, w, s, \lambda, \mu, t) \tag{12j}
$$

has d-index 1 and the same set of solutions as the equations of motion (1).

Proof. The proof can be found in [38]. \Box

Remark 2.5 a) The matrix functions S_p and S_v are called *kinematic selector* and dynamic selector, respectively.

b) We will call the DAE (12) the projected-strangeness-free formulation of the equations of motion. In $[21, 22, 23, 24]$ the strangeness-concept is introduced as tool for the classification of general nonlinear DAEs including over- and underdetermined DAEs. In particular, so called *strangeness-free* DAEs are introduced. Apart from the over- or underdeterminedness strangeness-free DAEs behave like DAEs with d-index 1 while nonstrangeness-free DAEs behave like DAEs with d-index 2 or larger. Strangeness-free DAEs do not contain hidden constraints. In particular, in $[21, 22, 23, 24]$ it is pointed out that strangeness-free DAEs and, therefore, the proje
ted-strangeness-free formulation of the equations of motion (12), are suited and preferable for the numerical treatment using stiff ODE solvers like implicit Runge-Kutta-Methods or BDF methods.

) The algorithm GEOMS is based on a proje
ted-strangeness-free from (12) of the equations of motion but it is not necessary that this form is provided by the user, i.e., the user does not have to perform the regularization to the projected-strangeness-free form. It is sufficient, if the user provides the constraints on velocity level (7b) and on acceleration level (8b) and (9b) in addition to the original equations of motion (1) and, if available, (3). By use of so alled order-n-formalisms for the evaluation of the equations of motion the constraints on velocity level and on acceleration level are computed automatically, see $[8, 34]$.

With these preparations we have presented all the tools to perform the consistency preserving index redu
tion of the equations of motion (1) as follows.

Algorithm 2.6 (Consistency preserving index reduction)

The equations of motion (1) are assumed to satisfy Assumptions 2.2. Furthermore, let $M \in \mathcal{C}^0(\mathbb{M}_p, \mathbb{R}^{n_v, n_v})$ and $Z \in \mathcal{C}^0(\mathbb{M}_p, \mathbb{R}^{n_p, n_v})$, where $\mathbb{M}_p = \mathbb{M} \cap (\mathbb{R}^{n_p} \times \mathbb{I})$ is the set of consistent (p, t) .

Then the regularization via consistency preserving index reduction is done by choosing a selector $S_p \in C^0(\mathbb{M}_p, \mathbb{R}^{n_{fp}, n_p})$ and a selector $S_v \in C^0(\mathbb{M}_p, \mathbb{R}^{n_{f_v}, n_v})$ depending on (p, u) with $n_{f_p} = n_p - n_\lambda$ and $n_{f_v} = n_v - n_\lambda - n_\mu$, in the following way.

1. Determination of selector S_p

- (a) Determine $K_p \in C^0(\mathbb{M}_p, \mathbb{R}^{n_p, n_{f_p}})$ depending on (p, t) such that the columns of $K_p(p, t)$ span ker $(G(p, s(p, t), t))$ for all $(p, t) \in M_p$.
- (b) Determine the selector $S_p \in C^0(\mathbb{M}_p, \mathbb{R}^{n_{fp}, n_p})$ depending on (p, t) such that $S_p(p,t)K_p(p,t)$ is nonsingular for all $(p,t) \in M_p$.

2. Determination of selector S_n

(a) Determine $K_v \in C^0(\mathbb{M}_p, \mathbb{R}^{n_v, n_{f_v}})$ depending on (p, t) such that the columns of $K_n(p,t)$ span

$$
\ker\left(\begin{bmatrix} G(p, s(p, t), t)Z(p) \\ H(p, s(p, t), t)Z(p) \end{bmatrix}\right)
$$

for all $(p, t) \in \mathbb{M}_p$.

(b) Determine the selector $S_v \in C^0(\mathbb{M}_p, \mathbb{R}^{n_{f_v}, n_v})$ depending on (p, t) such that $S_v(p, t)M(p, t)K_v(p, t)$ is nonsingular for all $(p, t) \in M_p$.

3. Pro je
ted strangeness-free form of the equations of motion

By appending the constraints on velocity level (7b) and the constraints on acceleration level (8b) and (9b), the projected-strangeness-free form of the equations of motion is given by (12).

With this algorithm we are able to determine a projected-strangeness-free form (12) of the equations of motion which contains all information of the set of consistency (10) . The projected-strangeness-free form (12) that is created in this way is analyti
ally equivalent to the original equations of motion in the sense that both have the same set of solutions. Therefore and because of Remark 2.5, the projectedstrangeness-free form (12) can be seen as a regularization technique. In particular, the semi-impli
it form of the proje
ted-strangeness-free form (12) is of great advantage, sin
e all onstraints are stated as purely algebrai equations, and there are no redundancies among the algebraic constraints and the differential equations.

Remark 2.7 Note that Selectors S_p and S_v satisfying the rank conditions (11) are not uniquely determined. Rather it is possible to choose the selectors in a piecewise onstant fashion. In prin
iple, the sele
tors may be kept onstant as long as the Newton iteration matrix $\mathfrak N$ (see Page 15) remains nonsingular. But the choice of the selectors influences the conditioning of the projected-strangeness-free formulation. Therefore, with respe
t to the onditioning of the linear systems whi
h have to be solved during the Newton iteration, the sele
tors should be re
omputed early enough and not just shortly before rea
hing a state, where the Newton iteration matrix be
omes singular. This fa
t is treated in GEOMS by the re
omputation of the selectors if the column pivoting with respect to the algebraic constraints changes or convergence problems of the Newton iteration occur. This is demonstrated in two simulation scenarios which are depicted in Tables 3 and 4.

Note that the pie
ewise onstant hoi
e of the sele
tors is of great advantage and importance for the numerical integration, because it offers the possibility to reduce the amount of computational work for the computation of the selectors. In particular, this means, that the condition number of the Newton iteration matrix \mathfrak{N} depends directly on the choice of the selectors. △△

Example 2.8 The mathematical pendulum: Let us consider a mathematical pendulum, of length $L > 0$ which represents a point mass moving without friction along a vertical circle of radius L under gravity denoted by the gravity acceleration q. For the description of the configuration of the pendulum we choose Cartesian coordinates $p = \left[\begin{array}{cc} x & y\end{array}\right]^T$ denoting the position of the mass m in the two dimensional space \mathbb{R}^2 . The equations of motion of first order have the form

$$
\left[\begin{array}{c}\n\dot{p}_1 \\
\dot{p}_2\n\end{array}\right] = \left[\begin{array}{c}\n v_1 \\
 v_2\n\end{array}\right],\n\tag{13a}
$$

$$
\begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} \dot{v}_1 \\ \dot{v}_2 \end{bmatrix} = \begin{bmatrix} 0 \\ -mg \end{bmatrix} - \begin{bmatrix} 2p_1 \\ 2p_2 \end{bmatrix} \begin{bmatrix} \lambda_1 \end{bmatrix},
$$
(13b)

$$
0 = [p_1^2 + p_2^2 - L^2]. \tag{13c}
$$

The holonomic constraints on velocity level and on acceleration level are given by

$$
0 = [2p_1v_1 + 2p_2v_2], \qquad (13d)
$$

$$
0 = \left[2v_1^2 + 2v_2^2 - 2p_2g - \frac{4}{m}(p_1^2 + p_2^2)\lambda_1 \right],
$$
 (13e)

respectively. Following Algorithm 2.6 we have to consider $G = \left[\begin{array}{cc} 2p_1 & 2p_2 \end{array}\right]$. The matrix function K_p can be determined as

$$
K_p = \left[\begin{array}{c} -p_2 \\ p_1 \end{array} \right]
$$

and, therefore, the selector S_p can be chosen as

$$
S_p = \left[\begin{array}{cc} -p_2 & p_1 \end{array} \right]
$$

such that

$$
S_p K_p = \begin{bmatrix} -p_2 & p_1 \end{bmatrix} \begin{bmatrix} -p_2 \\ p_1 \end{bmatrix} = \begin{bmatrix} p_2^2 + p_1^2 \end{bmatrix} = \begin{bmatrix} L^2 \end{bmatrix},
$$

see the constraints (13c). Since the mass matrix is given by $M = mI$, we can use $S_v = S_p$ and we get the projected-strangeness-free formulation

$$
-p_2\dot{p}_1 + p_1\dot{p}_2 = -p_2v_1 + p_1v_2, \tag{14a}
$$

$$
-mp_2\dot{v}_1 + mp_1\dot{v}_2 = -mgp_1, \qquad (14b)
$$

$$
0 = p_1^2 + p_2^2 - L^2, \tag{14c}
$$

$$
0 = 2p_1v_1 + 2p_2v_2, \t\t(14d)
$$

$$
0 = 2v_1^2 + 2v_2^2 - 2p_2g - \frac{4}{m}(p_1^2 + p_2^2)\lambda_1.
$$
 (14e)

As mentioned in Remark 2.7, the selectors S_p and S_v are not uniquely determined by the conditions (11) or the Algorithm 2.6. In particular, the selectors can be chosen to be pie
ewise onstant.

Let us consider this fact for the pendulum with the initial state $p_1 = 0$ and $p_2 = -L$. i.e., the pendulum is hanging downwards. In this position the selectors can be determined as

$$
S_p(p, u) = S_v(p, u) = \begin{bmatrix} L & 0 \end{bmatrix}.
$$
 (15)

Keeping these selectors constant, the leading matrix of the left-hand side of the underlying ordinary differential equations, (obtained by substituting the algebraic equations in (14) by their derivatives with respect to t) is

$$
\begin{bmatrix} L & 0 & 0 & 0 & 0 \\ 0 & 0 & mL & 0 & 0 \\ 2p_1 & 2p_2 & 0 & 0 & 0 \\ \times & \times & 2p_1 & 2p_2 & 0 \\ \times & \times & \times & \times & \frac{4}{m}(p_1^2 + p_2^2) \end{bmatrix}.
$$
 (16)

Obviously, the rank conditions (11) are fulfilled and the leading matrix (16) is nonsingular, as long as p_2 does not become zero. In particular, this means that as long as the pendulum does not reach one of the horizontal positions, i.e., $p_1 = \pm L$ and $p_2 = 0$, the selectors can be chosen constant as in (15). Otherwise, if the pendulum reaches or passes the horizontal position, the matrix (16) becomes singular and the first and third as well as the second and fourth equations are redundant such that the solution is not uniquely defined. Furthermore, the condition number of matrix (16) goes to infinity as p_2 goes to zero.

For these reasons, in the neighborhood of the horizontal position of the pendulum new selectors have to be determined. See also the Example 4.1 for numerical results. ⊳

3 GEOMS

The ode GEOMS is implemented in FORTRAN77 and furthermore, there exists a MATLAB [20] interface via MEX files for the direct usage of GEOMS in MATLAB.

However, in the following we only dis
uss the in FORTRAN77 implementation of GEOMS.

In GEOMS the 3-stage implicit Runge-Kutta Method Radau IIa of order 5, see [18]. as dis
retization of the proje
ted-strangeness-free formulation (12) of the equations of motion is implemented. Although, GEOMS bases on the presented stabilization technique developed in [38] and presented in Theorem 2.4, i.e., GEOMS uses the proje
ted-strangeness-free formulation (12) for the dis
retization, the user does not have to provide the projected-strangeness-free formulation. Instead the user has to provide all ne
essary information, i.e., in parti
ular, the hidden onstraints in addition to the original equations of motion (1) and , if available, (3).

The Runge-Kutta matrix A, the weight vector b , and the node vector c are given by the But
her tableau

$$
\begin{array}{c|c}\nC \uparrow A \\
\hline\nB \uparrow\n\end{array} \iff\n\begin{array}{c|c}\n\frac{4-\sqrt{6}}{10} & \frac{88-7\sqrt{6}}{360} & \frac{296-169\sqrt{6}}{1800} & \frac{-2+3\sqrt{6}}{225} \\
\frac{4+\sqrt{6}}{10} & \frac{296+169\sqrt{6}}{1800} & \frac{88+7\sqrt{6}}{36} & \frac{-2-3\sqrt{6}}{225} \\
\hline\n\frac{16-\sqrt{6}}{36} & \frac{16+\sqrt{6}}{36} & \frac{1}{9} \\
\hline\n\frac{16-\sqrt{6}}{36} & \frac{16+\sqrt{6}}{36} & \frac{1}{9}\n\end{array} \tag{17}
$$

see $[17, 18]$. The algorithm GEOMS is designed to handle equations of motion of the form (1) with possible redundant onstraints as well as with possibly known solution invariants (3) which may be provided as additional equations. If the mass matrix M is nonsingular and the onstraints are nonredundant then the equations of motion have to satisfy Assumption 2.2. If this is not the case some further rank assumptions have to be satisfied. For more details see [38].

Here and in the following we will use the typewriter style for obje
ts whi
h are part of the source codes of the implemented numerical algorithms. In particular, this involves names of subroutines like GEOMS, GEERREST, IVCOND, and variables like T, X, NWTMAT, CALSEL.

In the following we will discuss the features of GEOMS in detail. For the use and implementation of GEOMS see the manual in Appendix A.

The information of the equations of motion needed from the integration algorithm has to be provided in the following form.

The vector of unknown variables has to be in the form

$$
x^T = \mathbf{X}^T = \begin{bmatrix} w^T & \lambda^T & \mu^T & r^T & v^T & s^T & p^T \end{bmatrix}
$$

and the right-hand side in (1) and (3) of the hidden constraints has to be specified in a user-supplied subroutine with a name given by the user. The different parts

have to be given in the following order provided they occur.

In particular, the right-hand side has to be ordered such that the *algebraic part* i.e., the upper part (18A), contains the algebraic constraints ordered with respect to their dependencies, i.e., first $(18a)$, the constraints which restrict the additional variables w as well as the Lagrange multipliers λ and μ , second (18b), the constraints on velocity level and the information concerning solution invariants which restrict the velocities v , and third (18c), the constraints on position level, which restrict the position p and the contact variables s . The specified order leads to a Jacobian of the algebraic part with respect to x which has already block upper triangular structure that will be exploited in GEOMS.

The differential part, i.e., the second part (18D), contains the right-hand side of the differential equations also ordered in the same way as the algebraic part. We first (18d) have the equations that describe the behavior of the dynamical force elements followed by (18e) the dynamical equations of motion and, finally, (18f) the kinemati
al equations of motion.

In some cases the constraints of acceleration level (8) and (9), i.e., $0 = g^I$ and $0 = h^I$, are not explicitly available or difficult to evaluate. In this case GEOMS is also applicable. But one should note that only if all algebraic information, including $0 = g^I$, $0 = g^I$, and $0 = h^I$ are provided, instabilities and drift can be avoided by GEOMS. It is preferable to provide as mu
h information as possible. In the ase that the constraints on acceleration level are missing, the provided information is similar to a DAE that behaves like a DAE with d-index 2.

This fact has to be communicated by the user to the code GEOMS with help of the option IOPT(5)=FORM. If IOPT(5)=0 then the projected-strangeness-free form (12) of the equations of motion will be expe
ted as basis for the dis
retization. Thus, the user has to specify all information of the hidden constraints, i.e., up to acceleration level. If $IOPT(5)=1$, then the discretization will be done without specifying the constraints on acceleration level $0 = g^I$ and $0 = h^I$. In the latter case the used formulation of the equations of motion behaves like a system of d-index 2, i.e., it is not strangeness-free. Be
ause of the fa
t that the used formulation is not strangeness-free, the success of the numerical integration depends highly sensitively on the problem and on the consistency of the initial values, in particular, on the

Option	Name	Feature	Page
		preserving invariant solutions	4
		preserving hidden constraints	5
		preserving nonholonomic constraints	$\overline{2}$
		taking into account of redundancies in the con-	19
		straints	
IOPT(2)	LUN	optional output for integration information	
IOPT(3)	NIT	maximal number of Newton iterations	17
IOPT(4)	STARTN	starting values for the internal stages in the New-	15
		ton iteration	
IOPT(5)	FORM	incomplete regularization	11
IOPT(6)	NMAX	maximal number of integration steps	19
IOPT(8)	PRFD	step size control	18
IOPT(9)	NWTMAT	approximation of the Newton matrix at x_0 or one	15
		of the extrapolated stages possible	
IOPT(10)	NWTUPD	update of the Newton matrix	17
IOPT(11)	DECOMPC	LU, QR, or SV decomposition for the algebraic	17
		part	
IOPT(12)	DECOMPD	LU or QR decomposition for the differential part	17
IOPT(13)	SELCOMP	selector control	18
IOPT(14)	AUTONOM	exploitation of autonomous equations of motion	19
IOPT(15)	MASSTRCT	exploitation of the structure of the mass matrix	19
IOPT(17)	IVCNSST	check and correction of the initial values with	14
		respect to its consistency	

Table 1: Options and features of GEOMS

consistency of the Lagrange multipliers λ and μ .

An overview over the features of GEOMS is given in Table 1. Furthermore, in Table 2 the subroutines belonging to GEOMS and their task are listed.

The initial values are of great importan
e for the existen
e and the uniqueness of the solution. For the existen
e of a ontinuous solution the onsisten
y of the initial values is necessary. In particular, admissible initial values are restricted by the (hidden) constraints. On the other hand consistent initial values, in particular, onsistent initial Lagrange multipliers, are not automati
ally given by the modeling process and their determination by solving a system of nonlinear algebraic equations is difficult for complex multibody systems with a large number of constraints. Therefore, the algorithm GEOMS provides the possibility to determine onsistent initial values.

In addition to the algebraic equations determining the set of consistency M, see (10) , the user has to define in a subroutine IVCOND additional conditions to determine consistent initial values. Such conditions offer the possibility to determine some of the freely hoosable variables or to give further relations whi
h allows a unique determination of onsistent initial values.

Subroutines contained in the code GEOMS				
GEBSUBST	backward substitution of the algebraic part			
GECORE	core routine			
GEDECCLU	decomposition of the algebraic part with LU decomposition			
GEDECCOR	decomposition of the algebraic part with QR decomposition			
GEDECCSV	decomposition of the algebraic part with SV decomposition			
GEDECDLU	LU decomposition of the differential part			
GEELIMFXQ	elimination in the differential part according to QR decomposition			
	of the algebraic part			
GEELIMFXS	elimination in the differential part according to SV decomposition			
	of the algebraic part			
GEELIMMIQ	elimination in the mass matrix and the identity of the kinemati-			
	cal equations of motion according to QR decomposition of the			
	algebraic part			
GEELIMMIS	elimination in the mass matrix and the identity of the kinemat-			
	ical equations of motion according to SV decomposition of the			
	algebraic part			
GEERREST	error estimation, see Page 18			
GEFXNUM	numerical approximation of the Jacobian of the right-hand side of			
	the equations of motion			
GEGREPEQ	picking relevant columns of the differential part according to QR			
GEGREPES	decomposition picking relevant columns of the differential part according to LU			
	and SV decomposition			
GEINIVAL	determination of consistent initial values, see Page 14			
GEOMS	main routine			
GESOLDLU	solving the differential part by use of LU decomposition			
GESOLDQR	solving the differential part by use of QR decomposition			
GETRFRHSC	transformation of the right-hand side according to the algebraic			
	part			
User-supplied subroutines				
EOM	provides the reduced derivative array RDA (18)			
IVCOND	provides additional initial conditions needed for the consistent ini-			
	tialization, see Page 12			
JAC	provides the Jacobian of the reduced derivative array			
MAS	provides the mass matrix			
SOLOUT	output of the numerical solution and additional information during			
	integration			

Table 2: Subroutines of GEOMS

Example 3.1 The mathematical pendulum: In Example 2.8 we have introduced the mathematical pendulum. The position variables p are restricted to the circle with radius L, i.e., the constraint on position level is given by $0 = p_1^2 + p_2^2 - L^2$. If one of the position variables is given, the other is uniquely determined up to the

sign.

By defining additional conditions via the subroutine IVCOND the user can force the pendulum into a deviation of $\pi/4$ by setting $p_1 = L/\sqrt{2}$ or by $0 = p_1 + p_2$, for instance. Furthermore, a certain angular velocity ω can be prescribed by 0 = $\sqrt{v_1^2 + v_2^2}$ $\frac{2}{2}/L - \omega$.

The determination of onsistent initial values is done in the subroutine GEINIVAL and is based on the collection of all algebraic constraints $(1d)-(1g)$ and (7) , (8) , and (9) together with the onditions dened in the subroutine IVCOND.

The user has to decide if the given initial values are assumed to be consistent or not. By setting $IOPT(17) = IVCNSST = 1$, the initial values are assumed to be consistent and no check of consistency or correction of the initial values is done during the run of GEOMS. Note that non
onsistent initial values ould lead to onvergen
e problems in the integration pro
ess whi
h leads to an abort of the run of GEOMS. Otherwise, by setting $IOPT(17)=0$, the initial values are considered to be possibly inconsistent. Thus, consistency will be checked and the initial values will be corrected during the run of GEOMS, if necessary. If the user does not provide sufficiently many additional conditions, only the consistency is checked. If the initial values are consistent, then the integration will be ontinued, otherwise the run of GEOMS will be stopped. If the user provides more additional onditions than ne
essary, then the orre
tion (if necessary) is done regarding the overdetermined nonlinear system. If all conditions together are non
ontradi
tory, then onsistent initial values will be determined. Otherwise, the Newton iteration used in this pro
ess will diverge and the run of GEOMS will be stopped.

The solution of the nonlinear system of equations is obtained via a simplified Newton method with the possibility of a ertain number of updates of the iteration matrix, as des
ribed at Page 17. The stopping riterion is the same as that for the simplied Newton method during the integration pro
ess des
ribed at Page 16.

Remark 3.2 Note the fact that the conditions provided to IVCOND by the user dominate the given initial guess, i.e., if the given initial guess is consistent but does not satisfy the (possibly wrong) onditions provided by IVCOND, then the initial guess will be corrected in such a way that both, the constraints $(1d)-(1g)$ and (7) . (8), and (9) and the initial conditions provided to **IVCOND** are satisfied.

In case of an initial guess which is consistent to the constraints, the option IOPT(17) an be set to one to avoid su
h a orre
tion. Otherwise, the onditions provided to IVCOND should be adapted. △

If there is only interest in the omputation of onsistent initial values, the user has to set $T=TRND$ and $IOPT(17)=0$. Then the code GEOMS determines consistent initial values, will call the user-supplied subroutine SOLOUT, and finally will return to the alling subroutine.

In the following we will discuss the approach which is used in the algorithm GEOMS for the numeri
al integration of the equations of motion (1) and, if available, (3) by use of the three stage Runge-Kutta method of type Radau IIa of order 5. Let $s = 3$ denote the number of stages.

As mentioned above, the ode GEOMS ombines the dis
retization method with the regularization technique presented in Theorem 2.4. Therefore, the algorithm uses the proje
ted-strangeness-free form (12) as basis for the dis
retization. For more details on the discretization we refer to [38]. This discretization leads to a nonlinear stage equation for the determination of the three stages $X_{ki} \in \mathbb{R}^n$, $i = 1, 2, 3$ on the current integration interval $[t_k, t_{k+1}]$ with $t_{k+1} = t_k + h_k$. Here h_k denotes the current step size. The stages $X_{ki} \in \mathbb{R}^n$, $i = 1, 2, 3$ approximate the solution at the points $t_{ki} = t_k + c_i h_k$. The nonlinear stage equation has to be solved by use of a (simplified) Newton method.

A good choice of *starting values* X_{ki}^0 , $i = 1, 2, 3$ is very important for the convergence of the Newton iteration. In the code GEOMS two different possibilities for the determination of starting values for the integration step from t_k to t_{k+1} are implemented. The user has to define in advance which of both shall be used during the integration pro
ess.

By setting $IOPT(4) = STARTN=1$ the starting values for the internal stages are chosen by $X_{ki}^0 = x_k$, $i = 1, 2, 3$, where x_k denotes the already known value which approximates the solution at the point t_k . This x_k corresponds either to the initial value in the first integration step, i.e., $k = 0$, or it corresponds to the value determined at the end of the pre
eding integration step.

On the other hand setting $IOPT(4)=0$ (which is the default) the starting values X_{ki}^0 , $i = 1, 2, 3$ for the Newton iteration are obtained by evaluating the interpolation polynomial $q(t)$ of degree s over the already passed integration interval $[t_{k-1}, t_k]$ with $t_{k-1} = t_k - h_{k-1}$ and with $q(t_{k-1}) = x_{k-1}$, $q(t_{k-1} + c_i h_{k-1}) = X_{k-1i}$, $i = 1, 2, 3$. In this way we obtain the starting values for the Newton iteration as $X_{ki}^0 = q(t_k + c_i h_k)$, $i =$ 1, 2, 3, where x_{k-1} denotes the numerical solution at the point t_{k-1} . In particular, this means that the new starting values in the integration step from t_k to t_{k+1} are obtained by extrapolation to the points $t_k + c_i h_k$, $i = 1, 2, 3$ based on the internal stages of the earlier integration step from t_{k-1} to t_k . Of course, this is not possible in the first step. For more details see [18].

In GEOMS a simplified Newton method is implemented. For more details on Newton methods we refer to $[6]$. In particular, this means that a constant Newton iteration matrix $\mathfrak N$ is used during the whole or several parts of the Newton iteration inside the current integration step $[t_k, t_{k+1}]$. We use the simplified Newton method, since a constant Newton iteration matrix reduces the amount of computation because of the saved evaluation of Ja
obians and saved de
ompositions of the Newton iteration matrix in every except the first Newton iteration step. But the particular choice of the Newton iteration matrix influences the convergence of the Newton iteration. For this reason, the code GEOMS offers the possibility to choose between several reference points (X^*,t^*) for the determination of the Newton matrix. The choice has to be determined by the user by setting the option $IOPT(9)$ =NWTMAT. The range of possible choices is related to the stages during the integration step. As discussed previously,

there are two possibilities for the hoi
e of initial values for the Newton iteration for the determination of the internal stages. In case of $IOPT(4)=0$ the initial values are obtained by extrapolation of the solution computed so far in the points $t_k + c_i h_k$, $i = 1, 2, 3$. This offers the possibility to approximate the Newton iteration matrix at four different reference points $(X^*, t^*) = (X_{ki}^0, t_k + c_i h_k)$ for $i = 0, ..., 3$, where $c_0 = 0$ and c_i , $i = 1, 2, 3$ correspond to the node vector of the Runge-Kutta method, see Table 17. Furthermore, X_{ki}^0 corresponds to the extrapolated starting values for the internal stages at the times $t_k + c_i h_k$, $i = 0, ..., 3$, and, in particular, $X_{k0}^0 = x_k$ orresponds to the initial state of the urrent integration interval. Note that this possibility is only given if the initial values for the Newton iteration are extrapolated, i.e., if $\texttt{IOPT(4)}=0$. In the case of initial values chosen such that $X_{ki}^0=x_k$ for all $i = 1, 2, 3$ this possibility is not given and the Newton iteration matrix will be approximated at the initial point (x_k, t_k) with the initial state of the current integration step $[t_k, t_{k+1}]$.

Several numerical experiments have shown that the convergence of the Newton iteration can be improved by use of extrapolated initial values, i.e., $IOPT(4)=0$ in onne
tion with an approximation of the Newton iteration matrix at the se
ond internal stage, i.e., $(X^*, U^*) = (X_{k2}^0, t_k + c_2h_k)$ with IOPT(9)=2. But, if the Newton iteration dete
ts onvergen
e problems, and the integration step has to be repeated with a smaller step size, then the Newton iteration matrix has to be recomputed such that the overall computation time may increase if the number of times a convergence problems is detected is large. This number is reflected in the counter NCRJCT=IWORK(11) which corresponds to the number of step rejections caused by onvergen
e test failures.

It should be noted that the choice of different Newton iteration matrices within the Newton iteration is not available in the ode RADAU5. Furthermore, the ode GEOMS offers the possibility of a certain number of updates of the Newton iteration matrix during the Newton iteration inside of one integration step, see the following.

The convergence rate of the simplified Newton method is investigated in detail in [6], see also [17, 27]. One important question in the use of an iterative method for solving nonlinear systems inside an integration pro
ess is when to stop the iteration such that the obtained accuracy of the computed solution of the nonlinear system is within the pres
ribed toleran
e without performing too many Newton iteration steps.

The convergence estimation and the stopping criterion implemented in GEOMS is described in [18] and adopted from the code RADAU5 [17, 18]. The estimation of the convergence is based on the *weighted root square norm* $|| \cdot ||_{sc}$ which is defined for a $\zeta \in \mathbb{R}^n$ by

$$
||\zeta||_{sc} = \sqrt{\frac{1}{n} \sum_{i=1}^{n} \left(\frac{\zeta_i}{s c_i}\right)^2}
$$
(19)

with $sc_i = \texttt{ATOL}(i) + \max(|x_{ki}|, |x_{k+1i}|)$ RTOL (i) , see [18]. This norm allows to prescribe that some solution components have to be more precisely approximated than

other. This can be specified in the vectors ATOL and RTOL prescribing the absolute and relative toleran
e, respe
tively. For more details on the error estimation and the stopping criterion of the Newton iteration we refer to $|18, 38|$.

In the case of a very slow convergence or, in particular, in the case of divergence, the number of Newton iteration steps has to be restri
ted by a maximal number k_{max} = NIT = IOPT(3). Thus, the Newton iteration will stop unsuccessfully if a) the stopping criterion is not satisfied within the maximal number k_{max} of allowed Newton iteration steps, or if b) the iteration diverges.

In case a) the user has to decide whether the whole integration step has to be reje
ted be
ause of onvergen
e failures and to be repeated with a redu
ed step size, or if the Newton iteration should be ontinued with an updated Newton iteration matrix. In GEOMS this decision is made by defining the maximal number of updates in the option $IOPT(10) = NWTUPD$. However, several numerical results suggest that the number of allowed updates should not ex
eed 1.

It should be noted that the possibility of an update of the Newton iteration matrix within the Newton iteration is not available in the ode RADAU5.

During the Newton iteration a linear system has to be solved in each step. This has to be done in an efficient but stable way. The code GEOMS offers the possibility to decompose the differential part and the algebraic part via different decomposition methods. The user has to specify in the option $IDPT(11)=DECOMPC$ if the algebraic part, i.e., the Ja
obian of the onstraints, should be de
omposed by use of the LU decomposition with full pivoting ($IOPT(11)=1$), by a QR decomposition with pivoting (IOPT(11)=2), or by a SV decomposition (IOPT(11)=3). Heuristically seen, the LU decomposition with (partial) pivoting is a good compromise concerning efficiency and stability. Therefore, it is the default in GEOMS, although, the SV decomposition offers excellent stability properties but is more expensive.

Furthermore, with the option $IOPT(12) = DECOMPD$, the user can specify how to decompose the differential part. By setting $IOPT(12)=0$ the LU decomposition with partial pivoting is used and by setting $IOPT(12)=1$ the QR decomposition is used.

Remark 3.3 a) The separate de
omposition implemented in GEOMS has the advantage that the decomposition of the algebraic part can be done independently of the step size h. Only the de
omposition of the dierential part has to be done separately depending on h . In particular, if the Newton iteration has convergence problems and the algorithm interrupts the Newton for to redu
e the step size, then the information with respect to the algebraic part may be recycled which saves computational work. b) For the linear algebra computations like QR decompositions and SV decompositions we use BLAS⁺ (Basic Linear Algebra Subprograms) [25] and LAPACK² (Linear Algebra PACKage) [2] subroutines. △ △

For strangeness-free differential-algebraic systems in semi-implicit form like the projected-strangeness-free form (12) the scaling of the algebraic constraints with $1/h$ is

¹ BLAS - http://www.netlib.org/blas/

⁻LAPACK - http://www.netiib.org/lapack/

recommended in [32], where h is the current step size. Since the numerical integration of the equations of motion in GEOMS is based on the projected-strangeness-free formulation of the equations of motion, the constraints are scaled by $1/h$.

The *step size control* of the integration process is a very sensitive topic in the implementation of numeri
al algorithms for the integration of ODEs as well as for DAEs. An overview over several step size control strategies is given in [37], see also $[4, 5, 11, 18]$. The code GEOMS works with two different step size control strategies as used in the ode RADAU5, but adapted to the stru
ture of the equations of motion (1). The basis for a step size ontrol me
hanism is a lo
al error estimation. For more details we refer to $[18]$. The error estimation is implemented in the subroutine GEERREST. For the choice of a new step size for the next integration step or a repeated integration step two possibilities are implemented in GEOMS whi
h have to be selected by use of the option $IOPT(8)$ =PRED. With $IOPT(8)=2$ the *classical step* size controller developed in [11] is used and with $IDPT(8)=1$ the predictive step size $controler$, developed by Gustafsson in [16], is used. The predictive step size control is not possible in the first step, so, the classical step size controller will be used instead. The predictive step size controller needs slightly more work and storage than the classical step size controller but is more flexible in adaptating the step size. By use of the predictive step size controller a faster reduction of the step size without step rejections is possible than by use of the classical step size controller. This leads to a possible reduction of the overall amount of computation by use of the predictive step size controller. Experiments suggest that the predictive step size ontroller seems to produ
e safer results for simple problems. On the other hand, the choice of the classical controller often produces slightly faster runs, see also [18]. The predictive step size controller will be used in **GEOMS** by default.

Sin
e the ode GEOMS is based on the ombination of dis
retization and regularization to the projected-strangeness-free formulation of the equations of motion which is influenced by the choice of the selectors S_p and S_v , see Theorem 2.4, an efficient computation of these selectors is also important and will be discussed in the following.

In general, it is not necessary to recompute selectors in every integration step, see Remark 2.7.

If the LU decomposition is used for the differential part then it is possible to decide whether the determination of the selectors is done in each integration step $(IOPT(13) = SELCOMP=1)$ or the selectors are kept constant for those integration steps where the pivoting in the algebraic part does not change $(IOPT(13)=0)$. The latter ase is the default.

The code GEOMS offers the possibility to integrate the equations of motion of form (1) with possibly redundant constraints. As discussed in the literature [26], see also Remark 2.3, the solution may not be unique in this case, but under certain conditions the nonuniqueness is only restricted to the Lagrange multipliers λ , μ , and w. For more details see [38].

Very important for the integration of equations of motion with redundant onstraints is the dete
tion of the degree of redundan
y, i.e., the determination of the rank of the Jacobian associated with the constraints. The reliable numerical determination of the rank of a matrix is a deli
ate task and the SV de
omposition is a ommonly used tool for doing this. Therefore, the numeri
al integration of equations of motion with redundant onstraints is only allowed via the SV de
omposition for the constraints, i.e., $IOPT(11)=DECOMPC=3$.

The rank of the onstraints will be determined in every integration step. If it is detected in the first step that the constraints are redundant, a reliable numerical integration requires the use of the SV de
omposition at least for the de
omposition of the onstraints. Furthermore, if a possibly hange of the rank from one step to another is dete
ted, then the integration possibly has rea
hed a singular point and will be stopped with an error message.

If the equations of motion have solution invariants (3), then it is often desirable to preserve these solution invariants explicitly. GEOMS is able to preserve solution invariants if they are provided by the user as equations (3) in the RDA (18). See the Example 4.1.

The user may restrict the maximal number of allowed integration steps by setting the option $IOPT(6) = NMAX$. The default value of NMAX is 100000. Furthermore, the user may force the code to exploit some special structures of the problem. If the problem is autonomous the amount of computational work for the numerical integration may be reduced. By setting $IOPT(14)=AUTONOM=1$ the user tells the code that the problem is autonomous and the code GEOMS exploits this in the integration process. The default is $IOPT(14)=0$, i.e., the problem is not autonomous. In particular, if the mass matrix is constant and/or diagonal a large amount of computational work can be saved. Therefore, the user can specify by use of $IOPT(15) =$ MASSTRKT if the mass matrix is diagonal and constant $IDPT(15)=4$, full and constant $IDPT(15)=3$, diagonal and time and/or state dependent $IDPT(15)=2$, or full and time and/or state dependent $IOPT(15)=1$ (default).

4 Numeri
al experiments

In the following we will demonstrate the applicability and the performance of the new solver GEOMS. The integration with GEOMS will be performed for three different formulations of the regularized equations of motion. First, the numerical results obtained with GEOMS using the proje
ted-strangeness-free form (12) of the equations of motion will be abbreviated by **GEOMS**($psfEoM$). Second, the numerical results obtained with GEOMS without providing the constraints on acceleration level, see option IOPT(5) on Page 11, will be abbreviated by GEOMS(pEoM1). Furthermore, if the solution of the considered example satisfies some solution invariants we will use the projected-strangeness-free form of the equations of motion with explicit forcing of the solution invariants in addition to the two formulations above, see Page 4. The

numerical results in this case are denoted by **GEOMS**($psfEoM+1$).

The numerical integrations are done on an AMD Athlon XP 1800+, 1533 MHz.

Let us note that we will abstain from the use of physical units like meters or seconds.

Example 4.1 The mathematical pendulum: In Example 2.8 we introduced the equations of motion of the mathemati
al pendulum and we did regularize them to the proje
ted-strangeness-free form (14) whi
h is used for the numeri
al integration via GEOMS.

For the numerical simulations of the movement we used the mass $m = 1$, the length $L = 1$, and the gravitational acceleration $q = 13.75$. Let us note that we did modify the gravitational acceleration to approximately $q = 13.75$ such that the exact solution has a period of 2 which allows the comparison of the accuracy every period.

Figure 1: Mathemati
al Pendulum: Conservation of the total energy by the numerical solutions for prescribed RTOL=ATOL= 10^{-7} on the time domain $\mathbb{I} = [0, 1000]$

The mathematical pendulum modeled as in 13 represents a mechanical system which onserves the total energy. This total energy is given by

$$
E(p, v) = \frac{1}{2}m(v_1^2 + v_2^2) + mgp_2
$$
\n(20)

and is onserved su
h that

$$
0 = E(p(t), v(t)) - E_0 = e(p, v) \text{ with } E_0 = E(p_0, v_0)
$$
\n(21)

for $t \in \mathbb{I}$ and every solution of the equations of motion (13).

Let us consider the holonomic constraints $(13c)$ and their derivatives, which restrict

the motion of the pendulum in a nonredundant way, in omparison to the onservation of the total energy (21). We have

$$
0 = p_1^2 + p_2^2 - L^2, \tag{22a}
$$

$$
0 = 2p_1v_1 + 2p_2v_2, \t\t(22b)
$$

$$
0 = 2v_1^2 + 2v_2^2 - 2p_2g - \frac{4}{m}(p_1^2 + p_2^2)\lambda_1, \tag{22c}
$$

$$
0 = \frac{1}{2}m(v_1^2 + v_2^2) + mgp_2 - E_0.
$$
 (22d)

The constraints (22) are nonredundant for all p, v, and λ satisfying (22). In particular, in addition to the holonomic constraints and their derivatives the energy onservation restri
ts the solution as well. The dimension of the solution manifold with the energy conservation is therefore smaller than without the energy conservation.

For comparison, in Figure 1 the total energy in the numerical solution is depicted. In addition to GEOMS, the numerical solution is computed with RADAU5 $[17, 18]$ for different formulations, i.e., (EoM) the equations of motion (1) of d-index 3, (EoM2) the dindex 2 formulation (using the constraints on velocity level instead of the holonomic constraints), $(EoM1)$ the d-index 1 formulation (using the constraints on acceleration level instead of the holonomic constraints), and (GGL) the *Gear-Gupta-Leimkuhler formulation*, see [13]. Furthermore, the solution is computed with **ODASSL** [9, 10]. DASSL [4, 31], MEXAX [28], and HEDOP5 [3]. Expecting GEOMS($psfEoM+1$) the numeri
ally omputed total energy is far from being onstant. This an be expe
ted be
ause the energy onservation is ontained as an equation in the used formulation and is therefore explicitely forced during the numerical integration. However, even the other numeri
al results obtained with GEOMS satisfy the onservation of total energy very accurately. The preserving of the total energy yields a stabilization of the solution.

In the Figures 2 and 3 the efficiency is depicted, i.e., the relation between the obtained accuracy and the consumed computation time of the different used formulations. Obviously, the integration with use of GEOMS based on the projectedstrangeness-free formulation (14) plus solution invariants GEOMS($psfEoM+1$) offers the best performan
e for this example. Note that the approximation of the Lagrange multipliers by GEOMS(psfEOM+I) is mu
h better than of the other results.

A very important fact for the numerical integration and the stability of the numerical algorithms regarding the integration of DAEs is the satisfa
tion of the onstraints, in
luding the hidden onstraints. In Figure 4 the residual of the onstraints of position level, of velocity level, and of acceleration level depending on the simulation time is depicted. As one can see, **GEOMS** satisfies all constraints well.

Above we discussed the strategy for the determination of appropriate selectors, concerning IOPT(13), see Page 18. Furthermore, the projected-strangeness-free formulation (14) of the pendulum has been developed in Example 2.8 and the choice of the selectors S_p and S_v has been considered. We have stated above that, in principle the sele
tors may be kept onstant as long as the deviation of the pendulum does

Figure 2: Mathematical Pendulum: Efficiency of the solvers based on residual evaluations. Simulations are done on the time domain $\mathbb{I} = [0, 1000]$.

Figure 3: Mathematical Pendulum: Efficiency of the solvers based on residual evaluations. Simulations are done on the time domain $\mathbb{I} = [0, 1000]$.

Figure 4: Mathematical Pendulum: Residual of the constraints depending on $t \in$ [0, 1000] for prescribed $RTOL=ATOL=10^{-7}$

Figure 5: Mathemati
al Pendulum: Residual of the onstraints depending on the prescribed tolerance. Simulations are done on the time domain $\mathbb{I} = [0, 1000]$ with solvers based on residual evaluations.

```
Example 01_SimpPend
Integration with GEOMS(psfEoM)
TSTART = 0.00 TEND = 5.00 H0 = 0.100E-01
                         \equivTOLMIN = 1.0D- 7 TOLMAX = 1.0D- 9
Initial velocity 2.80 rad
[SimpPend] GEOMS(psfEoM) starts with IDID= 0 H= 0.100E-01 TOL = 0.100E-06
[SimpPend] GEOMS(psfEoM) finished with IDID= 0 H= 0.254E-01 at T= 0.500E+01
  NACCPT = 187 | NEOM = 2167 | NPDEC =187
                                                187
  NERJCT = 16 | NJAC = 187 | NEDEC = 204NCRJCT = 1 | NMAS= 1 | NBSUB = 660
                                                  \mathcal{L}\overline{\phantom{a}}CPUTIME= 0.060s | 2.060s | NSEL = 2.060s | 2.06
[SimpleEnd] GEOMS(psfEoM) starts with IDID= 0 H= 0.100E-01 TOL = 0.100E-07
[SimpPend] GEOMS(psfEoM) finished with IDID=
                                              0 H= 0.176E-01 at T= 0.500E+01
  NACCPT =270
                     \vert NEOM= 2961 | NPDEC =
                                                270
  NERJCT = 13 | NJAC= 270 | NEDEC = 284
  NCRJCT = 1 | NMAS = 1 | NBSUB =897
  CPUTIME = 0.060s | NSEL
                                           \overline{a}\Box\mathcal{D}[SimpPend] GEOMS(psfEoM) starts with IDID= \qquad 0 H= 0.100E-01 TOL = 0.100E-08
[Simplend] GEOMS(psfEoM) finished with IDID= 0 H= 0.118E-01 at T= 0.500E+01
  NACCPT = 391 | NEOM= 4141 | NPDEC = 391
  NERJCT = 9 | NJAC= 391 | NEDEC = 401
  NCRJCT = 1 | NMAS = 1 | NBSUB = 1250\BoxCPUTIME= 0.080s | | NSEL = 2
```
Table 3: Mathemati
al Pendulum: Statisti
al results for the numeri
al simulation with GEOMS using the psfEoM with initial velocity $v_{10} = 2.8$

not reach 90 degrees with respect to the initial state. The strategy for choosing the selectors is demonstrated in two simulation scenarios which are depicted in Tables 3 and 4.

Both s
enarios simulate the motion of the pendulum starting with the downward hanging initial position $p_0 = \begin{bmatrix} 0 & -1 \end{bmatrix}^T$ and an initial velocity $v_0 = \begin{bmatrix} v_{10} & 0 \end{bmatrix}^T$ over the time domain $\mathbb{I} = [0, 5]$. In Table 3 the simulation starts with an initial velocity of $v_{10} = 2.8$. This initial velocity leads to the highest deviation of $p = \left[\begin{array}{cc} \pm 0.699 & -0.715 \end{array} \right]$ which does not reach the deviation of 45 degrees. Because the constraint matrix has the form $G = \left[\begin{array}{cc} 2p_1 & 2p_2 \end{array} \right]$ and because of $|2p_1| < |2p_2|$ for all $t \in \mathbb{I}$, a change of the pivoting is not necessary such that a (re-)computation of the selector is only necessary at the beginning of the integration process and after every detected convergence failure. Therefore, the number NSEL of (re-)computations of the sele
tor equals the number NCRJCT of reje
tions be
ause of onvergen
e failures plus one initial omputation. The situation hanges ompletely if the pendulum passes the deviation of 45 degrees with respe
t to the initial state. This happens if the initial velocity is increased to $v_{10} = 2.9$. The numerical results are depicted in Table 4. Obviously, the (re-)
omputations of the sele
tor NSEL happened 13 times

```
Example 01_SimpPend
Integration with GEOMS(psfEoM)
TSTART = 0.00 TEND = 5.00 H0 = 0.100E-01
                       \alpha = 1TOLMIN = 1.0D- 7 TOLMAX = 1.0D- 9
Initial velocity 2.90 rad
[SimpPend] GEOMS(psfEoM) starts with IDID= 0 H= 0.100E-01 TOL = 0.100E-06
[SimpPend] GEOMS(psfEoM) finished with IDID= 0 H= 0.236E-01 at T= 0.500E+01
  NACCPT = 207 | NEOM= 2730 | NPDEC = 207
             207
  NERJCT = 7 | NJAC = 207 | NEDEC = 227NCRJCT = 13 | NMAS= 1 | NBSUB = 841
  CPUTIME= 0.060s | 2600s | NSEL = 2600s | 2700s | 2700s | 2700s
[SimpleEnd] GEOMS(psfEoM) starts with IDID= 0 H= 0.100E-01 TOL = 0.100E-07
[SimpPend] GEOMS(psfEoM) finished with IDID=
                                             0 H= 0.121E-02 at T= 0.500E+01
  NACCPT =303
                    | NEOM= 3729 | NPDEC =
                                               303
  NERJCT = 12 | NJAC = 303 | NEDEC = 32112
  NCRJCT = 6 | NMAS = 1 | NBSUB = 1142CPUTIME=
             CPUTIME= 0.080s | | NSEL = 19
                    \blacksquare\overline{a}[SimpPend] GEOMS(psfEoM) starts with IDID= \qquad 0 H= 0.100E-01 TOL = 0.100E-08
[SimpleEnd] GEOMS(psfEoM) finished with IDID= 0 H= 0.109E-01 at T= 0.500E+01
  NACCPT = 441 | NEOM = 5208 | NPDEC = 441NERJCT = 12 | NJAC = 441 | NEDEC = 459NCRJCT = 6 | NMAS = 1 | NBSUB = 1589CPUTIME= 0.090s | | NSEL = 19
```
Table 4: Mathemati
al Pendulum: Statisti
al results for the numeri
al simulation with GEOMS using the psfEoM with initial velocity $v_{10} = 2.9$

more often than convergence problems NCRJCT are detected. In Figure 6 the motion of the pendulum is depi
ted. One an see that the altitude of the pendulum passes 12 times the altitude of a deviation of 45 degrees. Therefore, the number NSEL of (re-)
omputations of the sele
tors ex
eeds the number NCRJCT of onvergen
e problems by 13, i.e., 12 plus one initial computations of the selectors. \triangleleft

Example 4.2 The truck model: In [36] a planar nonlinear model of a truck is introduced as benchmark example. In Figure 7 the topology as well as the coordinates, bodies, joints, and for
e elements are depi
ted. The model onsists of eleven coordinates $p_i, i = 1, ..., 11$ describing the motion of seven rigid bodies and one Lagrange multiplier λ_1 , see Table 5.

We omit to specify the equations of motion in detail and refer to [36] instead. Note that the equations of motion of the tru
k model are badly s
aled, sin
e the solution of the Lagrange multiplier λ_1 is of magnitude 10^4 but the solution of the other independent variables p and v are of magnitude 10^{-2} .

Figure 6: Mathematical Pendulum: Solution p_2 for initial velocity $v_{10} = 2.8$ and $v_{10} = 2.9$ on the time domain $\mathbb{I} = [0, 5]$.

Figure 7: Topology of the tru
k

Body			Coordinate
$\mathbf{1}$	rear wheel	p_1	vertical motion
2	front wheel	p ₂	vertical motion
3	truck chassis	p_3	vertical motion
		p_4	rotation about y -axis
4	engine	p_5	vertical motion
		p_6	rotation about y -axis
5	driver cabin	p ₇	vertical motion
		p_8	rotation about y -axis
6	driver seat	p_9	vertical motion
7	loading area	p_{10}	vertical motion
		p_{11}	rotation about y -axis
		λ_1	Lagrange multiplier with respect to the joint
			between loading area and truck chassis

Table 5: Nonlinear tru
k model

Figure 8: Truck: Efficiency of the solvers based on residual evaluations. Simulations are done on the time domain $\mathbb{I} = [0, 20]$.

The obtained accuracy of the numerical solutions is compared with the numerical solution RADAU5(GGL) obtained with a prescribed tolerance RT0L=AT0L= 10^{-15} . The precision of all results obtained by a prescribed tolerance are of similar accuracy but the consumed computation time differs, as seen in the Figure 8. By the use of the code GEOMS no problem occurred in the numerical integrations for any prescribed

toleran
es RTOL=ATOL≥ 10−¹⁵

$\overline{5}$ **Summary**

In this paper we have presented the new numerical algorithm GEOMS for the numeri
al integration of general equations of motion.

In particular, the algorithm GEOMS has been developed to carry out the numerical integration of the most general form of the equations of motion, including nonholonomic constraints and possible redundancies in the constraints, as they may appear in industrial applications. Besides the numerical integration it offers some additional features like preservation of invariant solutions, preservation of hidden constraints, use of different decomposition strategies, use of an incomplete regularization, and also checking and correction of the initial values with respect to their consistency. Subsequently, we have demonstrated the performance and the applicability of the algorithm for two mechanical examples of different degrees of complexity. The experien
e with these numeri
al examples and several other numeri
al tests suggest that the code GEOMS is an efficient and robust method for the numerical integration of the equations of motion.

Manual of GEOMS \mathbf{A} A Manual of GEOMS

```
SUBROUTINE GEOMS (
      SUBROUTINE GEOMS(CONTINUES) IN THE CONTINUES OF THE 
        NP, NV, NR, NW, NS, NL, NM, NI, M, N, NIVCOND,
     ## X,T,TEND,H,RTOL,ATOL,ITOL,IOPT,ROPT,
     #IVCOND, EOM, MAS, JAC, IJAC,
     #SOLOUT, IOUT,
     # LIWORK,IWORK,LRWORK,RWORK,
     #RPAR, IPAR, IERR,
     ## IDID)
  C
C NAME : (G)eneral (E)quations (O)f (M)otion (S)olver
\mathcal{C}\mathcal{C}PURPOSE
            : This subroutine performs the numerical simulation
\mathcal{C}of a multibody system whose state is described by
C
C p - position variables c of dimension NP,
\mathcal{C}v - velocity variables of dimension NV,
C The original force element variables of dimension NR,
C w - auxiliary variables of dimension NW,
C 3 - contact point variables 5 of dimension NS,
C 1 - holonomic Lagrange multipliers of dimension NL,
C m - nonholonomi
 Lagrange multipliers of dimension NM
C
C by numerical integration of the equations of motion
C of the form
C
```


```
C VERSION : April 12, 2006
\mathcal{C}C REVISIONS : -
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her
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\mathcal{C}C REFERENCES: This 
ode is part of the PhD thesis:
\mathcal{C}A.Steinbrecher. Numerical Solution of Quasi-Linear Differential-
\mathcal{C}Algebraic Equations and Industrial Simulation of Multibody
C Systems. PhD thesis, TU Berlin, Institut fuer Mathematik, 2006
\mathcal{C}C KEYWORDS : numeri
al simulation of me
hani
al systems, equations of motion,
\mathcal{C}differential-algebraic equations, projected-strangeness-free
C formulation, proje
ted-strangeness-index-1 formulation
C
C NOTE : The (basi
) linear algebra routines are provided by the
C libraries BLAS and LAPACK
\mathcal{C}C DISCLAIMER: Warranty dis
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C limited to loss of profits or loss of data, for any reason
C whatsoever, whether such liability is asserted on the basis
C of contract, tort (including negligence or strict liability),
C or otherwise, even if any of said parties has been warned of
C the possibility of such loss or damages.
\mathcal{C}\mathcal{C}\mathcal{C}\mathcal{C}C CALL
\mathcal{C}______________________________
\mathcal{C}C SUBROUTINE GEOMS(
C # NP,NV,NR,NW,NS,NL,NM,NI,M,N,NIVCOND,
\mathcal{C}C # X,T,TEND,H,RTOL,ATOL,ITOL,IOPT,ROPT,
\mathcal{C}# IVCOND, EOM, MAS, JAC, IJAC,
```

```
C # SOLOUT,IOUT,
C # LIWORK, IWORK, LRWORK, RWORK,
\overline{C}# RPAR, IPAR, IERR,
C # IDID)
C IMPLICIT NONE
C INTEGER NP, NV, NR, NW, NS, NL, NM, NI, M, N, NIVCOND,
C # ITOL,IJAC,IOUT,LIWORK,LRWORK,IERR,IDID,
     #C # IOPT(40),IWORK(LIWORK),IPAR(*)
C DOUBLE PRECISION T,TEND,H,
C # X(N), RTOL(*), ATOL(*), ROPT(40), RWORK(LRWORK),
     #C RPAR(*)
\mathcal{C}EXTERNAL IVCOND, EOM, MAS, JAC, SOLOUT
\mathcal{C}\mathcal{C}INPUT- AND OUTPUT-ARGUMENTS
\mathcal{C}____________________________
C
C NP Input : integer
\mathcal{C}Number of position variables p.
C
C NV Input : integer
C Number of velocity variables v.
C
C NR Input : integer
C Number of dynami
al for
e element variables r.
C
C NW Input : integer
C Number of auxiliary variables w.
C
C NS Input : integer
C Mumber of contact point variables s.
C
C NL Input : integer
\mathcal{C}Number of Lagrange multipliers l=lambda for holonomic
\mathcal{C}constraints.
C
C NM Input : integer
\mathcal{C}Number of Lagrange multipliers m=mu for nonholonomic
C constraints.
C
C NI Input : integer
C Number of invariants, e.g., energy 
onservation.
C
C M Input : integer
C Total number of provided equations (M.GE.N), i.e., dimension of
C RDA, see subroutine EOM. In the 
ase of the use of the
C \bullet projected-strangeness-free formulation we have
C M=NP+NV+NR+NW+NS+3*NL+2*NM+NI,C * proje
ted-strangeness-index-1 formulation we have
\mathcal{C}C M=NP+NV+NR+NW+NS+2*NL+NM+NI.
C
C N Input : integer
C Number of unknowns (M.GE.N), i.e., dimension of X. We have
N=NP+NV+NR+NW+NS+NL+NM.
\mathcal{C}
```
C NIVCOND Input : integer Number of initial value conditions, which have to be satisfied \mathcal{C} \overline{C} in addition to the constraints obtained from the provided equa-C tions of motion. See subroutine IVCOND. C C X Input : double precision array X(N) C Initial values for X. The array X ontains the (initial) state C of the me
hani
al system in the following order C $X(1:NW)$ $=w$ C $X(W+1:NW+NL)$ $=1$ $(=1ambda)$ \mathcal{C} $X(NW+NL+1:NW+NL+NM)$ =m (=mu) $\mathbf C$ \mathcal{C} $X(NL+NM+NW+1:NL+NM+NW+NR)$ $=r$ \mathcal{L} \mathcal{C} $X(NL+NM+NW+NR+1:NL+NM+NW+NR+NV)$ \mathcal{C} C -- $X(NL+NM+NW+NR+NV+1:NL+NM+NW+NR+NV+NS)$ =s \mathcal{C} \mathcal{C} $X(NL+NM+NN+NK+NN+NS+1:NL+NM+NW+NR+NV+NS+NP) = p$ C C Output : C Numerical approximation of the solution at the last successfully C reached time T. C C T Input : double precision C Initial time. C Output : C 12 Last successfully reached time. If the whole integration was C successful then T=TEND. \overline{C} C TEND Input : double precision C Final time. \mathcal{C} \mathcal{C} $\,$ H Input : double precision \mathcal{C} Initial step size. \mathcal{C} Output : \mathcal{C} Last used step size. C RTOL Input : double pre
ision RTOL (or array RTOL(N)) C ATOL Input : double precision ATOL (or array ATOL(N)) C Relative and absolute error tolerances. They can be both C scalars or else both vectors of length N. C **In the case of a scalar the prescribed relative and absolute** C between to tolerances are valid for every component of the vector of C cannowns X. The code keeps, roughly, the local error of $X(I)$ \mathcal{C} $below$ RTOL*ABS $(X(I))$ +ATOL. C 1. In the case of a vector of dimension N the prescribed relative \mathcal{C} tolerances RTOL(I) and absolute tolerances ATOL(I) are valid \mathcal{C} for the I-th component $X(I)$ of the vector of unknowns X. \mathcal{C} The code keeps, roughly, the local error of $X(I)$ below C RTOL $(I)*ABS(Y(I))+ATOL(I)$. C C ITOL Input : integer \mathcal{C} Switch for RTOL and ATOL:

41

C IOUT Input : integer \mathcal{C} Switch for the calling of subroutine SOLOUT. \mathcal{C} IOUT=0 Subroutine is never called. C IOUT=1 Subroutine is available for output. C C LIWORK Input : integer C De
lares the length of the array IWORK. LIWORK has to be at least C 20. C C IWORK Output: integer array IWORK(LIWORK) C Statisti
al information \mathcal{C} IWORK(1) NACCPT - Number of accepted integration steps \mathcal{C} IWORK(2) NEOM - Number of evaluations of the right-hand side \mathcal{C} of the equations of motion \mathcal{C} IWORK(3) NMAS - Number of evaluations of the mass matrix C IWORK(4) NJAC - Number of evaluations of the Ja
obian of the \mathcal{C} right-hand side of the equations of motion \mathcal{C} IWORK(5) NSEL - Number of determinations of suitable selectors \mathcal{C} IWORK(6) NPDEC - Number of predecompositions, i.e., of FX, M, C and IKIN C IWORK(7) NEDEC - Number of E-de
ompositions, i.e., of E1 and E2 C **IWORK(8) NBSUB** - Number of backward substitutions C IWORK(9) NSTEP - Number of steps C IWORK(10) NERJCT - Number of reje
tions aused by error test C failures C IWORK(11) NCRJCT - Number of reje
tions aused by onvergen
e C C problems of the Newton process C C LRWORK Input : integer \overline{C} Declares the length of the array RWORK. C 6 A safe choice for all possible setting in IOPT is C LRWORK at least 5*N \mathcal{C} Depending on IOPT it is sufficient ... \mathcal{C} If $IOPT(17)=IVCNSST=0$ then LRWORK has to be at least $5*N$ \mathcal{C} If $IOPT(11)=DECOMPC=3$ then LRWORK has to be at least \mathcal{C} 5*MAX(M1, M2, M3, N1, N3, N4), see comments to subroutine JAC. \mathcal{C} If $IOPT(11)=DECOMPC=2$ then LRWORK has to be at least N C If IOPT(12)=DECOMPD=1 then LRWORK has to be at least 2*N C For good performance, LRWORK should generally be larger. C C RWORK Intern : integer array IWORK(LIWORK) C C IPAR Input/Output : integer array IPAR(*) C **Integer parameters which are only used by the user. They are** \mathcal{C} unused and unchanged by GEOMS. \mathcal{C} C RPAR Input/Output: double pre
ision array RPAR(*) \mathcal{C} Double precision parameters which are only used by the user. \mathcal{C} RPAR is unused and unchanged by GEOMS. \overline{C} C IERR Input/Output : integer C 1ndicator of success. IERR is only used by user \mathcal{C} supplied subroutines. After every call of a user supplied \mathcal{C} subroutine the status of IERR is checked. If IERR is negative

```
C the run of GEOMS will be interrupted and GEOMS returns to the
             C 
alling program. IERR is un
hanged by GEOMS.
\mathcal{C}\mathcal{C}C IDID Output : integer
C Reports success upon return. The first two digits
C indicate the subroutine which causes trouble.
C
C IDID=-10.. An error occurred in the subroutine GEOMS
C -1001 Option array IOPT or ROPT or toleran
es RTOL or ATOL
C contains wrong data
C Check the output in UNIT=IOPT(2) for more information
\mathcal{C}If the option IOPT(2) equals 0 turn on the output.
\mathcal{C}-1002 Initial IDID lower than 0
\mathcal{C}\mathcal{C}IDID=-11.. An error occurred in the subroutine GECOR
C -1101 Stop initialized by SOLOUT
\mathcal{C}-1102 Stop initialized by EOM
\mathcal{C}-1103 Stop initialized by MAS
\overline{C}-1104 Stop initialized by JAC
C -1105 Initial conditions not consistent
C -1106 Final time TEND before initial time T
C -1111 QR-Decomposition of FX1 not possible
C -1112 QR-Decomposition of FX2 not possible
C -1113 QR-Decomposition of FX3 not possible
C -1114 QR-Decomposition of E1 or E2 not possible
C -1115 Newton method repeatedly does not converge NSING.GE.5
C -1116 Newton method repeatedly does not converge NSING.GE.5
C -1117 More than NMAX steps are needed
C -1118 Step size too small
\overline{C}-1128 An error occurred during use of DORMQR
C -1129 An error occurred during use of DORMQR
C
\mathcal{C}IDID=-12.. An error occurred in the subroutine GEFXNUM
\mathcal{C}-1201 Stop initialized by EOM
\mathcal{C}\mathcal{C}IDID=-14.. An error occurred in the subroutine GEDECCQR
\mathcal{C}-1401 Constraints redundant or dd/dw singular
C (FX1 rank deficient).
C Try the integration again with IOPT(11)=3 (SVD).
C -1402 Constraints or the invariant equations are redundant
C (FX2 rank deficient). Try the integration again with
C IOPT(11)=3 (SVD).
C -1403 Constraints redundant or dc/ds singular
C (FX3 rank deficient).
C Try the integration again with IOPT(11)=3 (SVD).
\mathcal{C}\mathcal{C}IDID=-18.. An error occurred in the subroutine GETRFRHSC
\mathcal{C}-1801 Multiplication with Q1 not possible
\mathcal{C}-1802 Multiplication with Q2 not possible
\overline{C}-1803 Multiplication with Q3 not possible
C -1804 Multiplication with Q4 not possible
\mathcal{C}\mathcal{C}IDID=-20.. An error occurred in the subroutine GEERREST
\mathcal{C}-2004 Multiplication with Q4 not possible
```


Referen
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