

MIXED NORM ESTIMATES FOR THE CESÀRO MEANS ASSOCIATED WITH DUNKL–HERMITE EXPANSIONS

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ABSTRACT. Our main goal in this article is to study mixed norm estimates for the Cesàro means associated with Dunkl–Hermite expansions on \mathbb{R}^d . These expansions arise when one considers the Dunkl–Hermite operator (or Dunkl harmonic oscillator) $H_\kappa := -\Delta_\kappa + |x|^2$, where Δ_κ stands for the Dunkl–Laplacian. It is shown that the desired mixed norm estimates are equivalent to vector-valued inequalities for a sequence of Cesàro means for Laguerre expansions with shifted parameter. In order to obtain such vector-valued inequalities, we develop an argument to extend these Laguerre operators for complex values of the parameters involved and apply a version of three lines lemma.

1. INTRODUCTION AND MAIN RESULTS

The Dunkl operators were introduced by C. F. Dunkl in [9], where he built a framework for a theory of special functions and integral transforms in several variables related to reflection groups. Such operators are relevant in physics, namely for the analysis of quantum many body systems of Calogero–Moser–Sutherland type (see [8, 14]). From the mathematical analysis point of view, the importance of Dunkl operators lies on the fact that they generalize the theory of symmetric spaces of Euclidean type. There is a vast literature related to Dunkl transform and Dunkl Laplacian, see for instance [1, 5, 7, 12, 18, 23, 24].

In [17] M. Rösler studied the Dunkl–harmonic oscillator (which we will also call Dunkl–Hermite operator)

$$H_\kappa := -\Delta_\kappa + |x|^2,$$

where Δ_κ stands for the Dunkl–Laplacian (2.2), and introduced the *Dunkl–Hermite functions* $\Phi_{\mu,\kappa}$, $\mu \in \mathbb{N}^d$, as eigenfunctions of H_κ . When κ , which is called a multiplicity function (see (2.1)), is the null function, the situation is reduced to the standard Hermite operator H and $\Phi_{\mu,\kappa}$ become the usual Hermite functions Φ_μ , see [17, page 521]. The set $\{\Phi_{\mu,\kappa}\}_{\mu \in \mathbb{N}^d}$ forms an orthonormal basis for $L^2(\mathbb{R}^d, h_\kappa^2 dx)$ where h_κ^2 is a suitable weight function defined in terms of the corresponding reflection group and κ , see precise definitions in Section 2. Thus we have the orthogonal expansion

$$f = \sum_{\mu \in \mathbb{N}^d} (f, \Phi_{\mu,\kappa}) \Phi_{\mu,\kappa}$$

which converges to f in $L^2(\mathbb{R}^d, h_\kappa^2 dx)$. Here (\cdot, \cdot) is the product in $L^2(\mathbb{R}^d, h_\kappa^2 dx)$. In short we can rewrite this as

$$(1.1) \quad f = \sum_{j=0}^{\infty} P_{j,\kappa} f, \quad \text{where } P_{j,\kappa} f = \sum_{|\mu|=j} (f, \Phi_{\mu,\kappa}) \Phi_{\mu,\kappa}.$$

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It is known that the Dunkl–Hermite functions are of Schwartz class and hence the projections $P_{j,\kappa}f$ in (1.1) make sense for any $f \in L^p(\mathbb{R}^d, h_\kappa^2 dx)$. However, when $p \neq 2$, we do not have any convergence results for $\sum_{j=0}^\infty P_{j,\kappa}f$. In the case of $\kappa \equiv 0$, where the above series reduces to standard Hermite expansions, it is also well known that the Hermite expansions fail to converge in $L^p(\mathbb{R}^d, dx)$ for $p \neq 2$, unless $d = 1$. Even when $d = 1$, the series converges in $L^p(\mathbb{R}, dx)$ (or, equivalently, the corresponding partial sum operators are uniformly bounded) if and only if $\frac{4}{3} < p < 4$ according to a theorem by R. Askey and S. Wainger [3].

In the absence of convergence results for the partial sums, we are led to consider other summability methods such as Cesàro means or Bochner–Riesz means. For $N \in \mathbb{N}$ and $\delta > 0$ we define the Cesàro means of order δ associated with the Dunkl–Hermite expansions by

$$(1.2) \quad \sigma_{N,\kappa}^\delta f(x) := \frac{1}{A_N^\delta} \sum_{j=0}^N A_{N-j}^\delta P_{j,\kappa} f(x)$$

where

$$A_j^\delta := \binom{j+\delta}{j} = \frac{\Gamma(j+\delta+1)}{\Gamma(j+1)\Gamma(\delta+1)}$$

are the binomial coefficients. When $\kappa \equiv 0$ and $d \geq 2$, the operators $\sigma_{N,0}^\delta f$ (that we will simply denote by $\sigma_N^\delta f$) are nothing but the Cesàro means for the standard Hermite expansions. They converge to f in $L^p(\mathbb{R}^d, dx)$, $1 \leq p < \infty$, whenever $\delta > \frac{d-1}{2}$, see [21, 22]. Actually, more precise results are known, giving critical indices of summability for any given p , $1 \leq p < \infty$.

In the one dimensional case there is only one reflection group, viz. \mathbb{Z}_2 , and the generalized Hermite expansion in this case has been studied by Ó. Ciaurri and J. L. Varona [4]. Therein, the authors studied weighted norm inequalities for the Cesàro means. However, for the higher dimensional case we are not aware of any work dealing with Cesàro or Riesz means associated with Dunkl–Hermite expansions, which is the main concern of the present work.

The techniques used to study Cesàro means $\sigma_N^\delta f$ for the standard Hermite expansions are not available in the case of Dunkl–Hermite expansions. This is mainly due to the lack of explicit formulas for $\Phi_{\mu,\kappa}$ and their asymptotic properties.

However, the situation changes if we express the Cesàro means in a convenient way, by separating variables in polar coordinates. Indeed, the basic fundamental idea is to write down the Cesàro means $\sigma_{N,\kappa}^\delta f$ in terms of *spherical h -harmonics*. The spherical h -harmonics (or simply *h -harmonics*) are the restrictions of solid h -harmonics to \mathbf{S}^{d-1} , where \mathbf{S}^{d-1} is the unit sphere in \mathbb{R}^d , $d \geq 2$ (a good reference for h -harmonics is [10, Chapter 5]). By solid h -harmonics we mean homogeneous polynomials $P(x)$ satisfying $\Delta_\kappa P(x) = 0$. Let \mathcal{H}_m be the space of all h -harmonics of degree m . Then the space $L^2(\mathbf{S}^{d-1}, h_\kappa^2 d\sigma)$ is the orthogonal direct sum of the finite dimensional spaces \mathcal{H}_m over $m = 0, 1, 2, \dots$. Thus there is an orthonormal basis $\{Y_{m,j} : j = 1, 2, \dots, d(m), m = 0, 1, 2, \dots\}$, where

$$(1.3) \quad d(m) = \dim(\mathcal{H}_m),$$

for $L^2(\mathbf{S}^{d-1}, h_\kappa^2 d\sigma)$ so that for each m , $\{Y_{m,j} : j = 1, 2, \dots, d(m)\}$ is an orthonormal basis of h -harmonics of degree j for \mathcal{H}_m . Let $x \in \mathbb{R}^d$. Then, if we take $x = rx'$, $r \in (0, \infty)$, $x' \in \mathbf{S}^{d-1}$, the h -harmonic expansion of a function f is given by

$$f(rx') = \sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} f_{m,j}(r) Y_{m,j}(x'),$$

where the h -harmonic coefficients are

$$(1.4) \quad f_{m,j}(r) = \int_{\mathbf{S}^{d-1}} f(rx') Y_{m,j}(x') h_\kappa^2(x') d\sigma(x').$$

Therefore, due to the orthogonality of the h -harmonics, it is natural to introduce the mixed norm spaces for Dunkl–Hermite expansions $L^{p,2}(\mathbb{R}^d, h_\kappa^2 dx)$, namely, the space of all functions on \mathbb{R}^d for which

$$(1.5) \quad \|f\|_{(p,2)} := \left(\int_0^\infty \left(\int_{\mathbf{S}^{d-1}} |f(rx')|^2 h_\kappa^2(x') d\sigma(x') \right)^{\frac{p}{2}} r^{d+2\gamma-1} dr \right)^{\frac{1}{p}}$$

are finite. Here γ is a positive number intrinsic to the underlying multiplicity function κ , see Section 2 for the definition. Our main result is the following.

Theorem 1.1. *Let $d \geq 2$ and $\delta > \frac{d+2\gamma-1}{2}$ where γ is as in (2.3). Then, for any $1 < p < \infty$, we have the uniform estimates*

$$\|\sigma_{N,\kappa}^\delta f\|_{(p,2)} \leq C \|f\|_{(p,2)}.$$

Consequently, $\sigma_{N,\kappa}^\delta f$ converges to f in $L^{p,2}(\mathbb{R}^d, h_\kappa^2 dx)$ as $N \rightarrow \infty$.

We will sketch the outline of the proof. The first step is to write down the Cesàro means $\sigma_{N,\kappa}^\delta f$ evaluated at $x = rx'$ in terms of spherical harmonics. Then, a Funk–Hecke formula is used to rewrite the projection of the kernel associated with $\sigma_{N,\kappa}^\delta f$ in such a basis as a sum of shifted Cesàro kernels for the usual Hermite expansions. Indeed, we will prove that

$$\sigma_{N,\kappa}^\delta f(rx') = \sum_{m=0}^N \sum_{j=1}^{d(m)} T_{N,m}^{\delta,\gamma} f_{m,j}(r) Y_{m,j}(x').$$

In the expression above, $T_{N,m}^{\delta,\gamma}$ are the linear operators defined by

$$(1.6) \quad T_{N,m}^{\delta,\gamma} f_{m,j}(r) := \int_0^\infty K_{N,m}^{\delta,\gamma}(r,s) f_{m,j}(s) s^{d+2\gamma-1} ds,$$

where $f_{m,j}(r)$ are the h -harmonic coefficients of f in (1.4). The kernels $K_{N,m}^{\delta,\gamma}$ can be written in terms of the kernels of either the d -dimensional or $d+1$ -dimensional Cesàro means for the standard Hermite expansions with shifted parameters. See Corollary 4.2. But now, the standard Hermite polynomials evaluated at x are connected with the Laguerre polynomials evaluated at $|x|^2$. So when trying to prove the $(p,2)$ boundedness of the Cesàro operator, we end up with a vector-valued extension for a sequence of operators associated to Laguerre expansions.

Theorem 1.2 (Vector-valued inequalities for $T_{N,m}^{\delta,\gamma}$). *Let $0 \leq \gamma \leq \frac{1}{2}$, $\delta > \frac{d+2\gamma-1}{2}$ and $T_{N,m}^{\delta,\gamma}$ be the operators defined in (1.6). Then, for any $1 < p < \infty$, there is a constant C independent of N such that*

$$\left\| \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} |T_{N,m}^{\delta,\gamma} f_{m,j}(r)|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^+, r^{d+2\gamma-1} dr)} \leq C \left\| \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} |f_{m,j}(r)|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^+, r^{d+2\gamma-1} dr)},$$

where $d(m)$ is as in (1.3), for any sequence of functions $f_{m,j} \in L^p(\mathbb{R}^+, r^{d+2\gamma-1} dr)$ for which the right hand side is finite.

Then, it will be seen that Theorem 1.1 is a direct consequence of Theorem 1.2.

In order to prove Theorem 1.2, we need to use a sophisticated version of the three lines lemma. This will require to extend the Cesàro means for the usual Hermite expansions for complex values of δ and also complexify the type of the Laguerre functions involved.

The natural question that arises immediately concerns the convergence of Cesàro means in $L^p(\mathbb{R}^d, h_\kappa^2 dx)$, when $\kappa \neq 0$. Nothing is known about this, neither the techniques that could be used to solve the problem.

The paper is organized as follows. In Section 2 we recall basics about the general Dunkl context and the Funk–Hecke identity. Section 3 is devoted to the study of Cesàro means for the standard

Hermite expansions, and a vector-valued inequality for an operator related to these Cesàro means is proved. In Section 4 we express the Cesàro kernels for the Dunkl–Hermite expansions in terms of the same for the standard Hermite expansions, and we bring out the connection with Laguerre expansions, which allows us to introduce an analytic family of operators. With this, we prove the main results in Section 5. For the sake of reading, we move the proofs of certain propositions in Section 4 to Section 6. Section 7 contains also some technical results and integral formulas for Bessel functions which involve ultraspherical polynomials with real and complex parameters.

Notation. Throughout the paper, we will use the following notation. In general, all the operators, kernels and functions are defined in the ambient space \mathbb{R}^d , so in these cases we will mostly omit the dimensional parameter d . For instance, we will write just Φ_k or σ_N^δ . Nevertheless, whenever the dependence of the kernels on the dimension is explicitly needed we will add a superindex d , namely we will write $\Phi_k^{(d)}$ or $\sigma_N^{\delta,d}$. Moreover, when referring to these objects in \mathbb{R}^{d+1} , we will always explicitly add a superindex $d+1$, like $\Phi_k^{(d+1)}$ or $\sigma_N^{\delta,d+1}$, to denote the Hermite functions, Cesàro means and kernels in \mathbb{R}^{d+1} . The number of operators and parameters appearing in this work is quite large, and we tried to reduce the notation to a minimum. Because of this, we present in Section 5 a table with a summary of the most important operators and kernels involved directly in the proof of the main theorems. For $1 \leq p \leq \infty$, p' will denote its conjugate, $1/p + 1/p' = 1$. Moreover, we shall write C to denote positive constants independent of significant quantities the meaning of which can change from one occurrence to another.

2. THE GENERAL DUNKL SETTING

For completeness, in this section we collect several facts concerning the general Dunkl setting and the Dunkl harmonic oscillator. For a more detailed exposition on these topics, we refer the reader to [9, 10, 17].

We use the notation $\langle \cdot, \cdot \rangle$ for the standard inner product on \mathbb{R}^d . For $\nu \in \mathbb{R}^d \setminus \{0\}$, we denote by σ_ν the orthogonal reflection in the hyperplane perpendicular to ν , i.e.,

$$\sigma_\nu(x) = x - 2 \frac{\langle \nu, x \rangle}{|\nu|^2} \nu.$$

A finite subset $R \subset \mathbb{R}^d \setminus \{0\}$ is a root system if $\sigma_\nu(R) = R$, for all $\nu \in R$. Each root system can be written as a disjoint union $R = R_+ \cup (-R_+)$, where R_+ and $-R_+$ are separated by a hyperplane through the origin. Such R_+ is called the set of all positive roots in R . The group G generated by the reflections $\{\sigma_\nu : \nu \in R\}$ is called the reflection group or Coxeter group associated with R . A function

$$(2.1) \quad \kappa : R \rightarrow [0, \infty)$$

which is invariant under the action of G on the root system R is called a multiplicity function. Let T_j , $j = 1, 2, \dots, d$, be the difference-differential operators defined by

$$T_j f(x) = \frac{\partial f}{\partial x_j}(x) + \sum_{\nu \in R_+} \kappa(\nu) \nu_j \frac{f(x) - f(\sigma_\nu x)}{\langle \nu, x \rangle}.$$

These operators, known as Dunkl operators, form a family of commuting operators. The Dunkl Laplacian Δ_κ is then defined to be the operator

$$(2.2) \quad \Delta_\kappa = \sum_{j=1}^d T_j^2$$

which can be explicitly calculated, see [10, Theorem 4.4.9]. It is known that the operators T_j have a joint eigenfunction $E_\kappa(x, y)$ satisfying

$$T_j E_\kappa(x, y) = y_j E_\kappa(x, y), \quad j = 1, \dots, d.$$

The function $(x, y) \mapsto E_\kappa(x, y)$ is called the *Dunkl kernel* or the *generalized exponential kernel* on $\mathbb{R}^d \times \mathbb{R}^d$, which is the generalization of the exponential function $e^{\langle x, y \rangle}$. Associated with the root system R and the multiplicity function κ , the weight function $h_\kappa^2(x)$ is defined by

$$h_\kappa^2(x) := \prod_{\nu \in R_+} |\langle x, \nu \rangle|^{2\kappa(\nu)}.$$

The nonnegative real number

$$(2.3) \quad \gamma = \sum_{\nu \in R_+} \kappa(\nu)$$

defined in terms of the multiplicity function $\kappa(\nu)$ plays an important role in Dunkl theory. Note that $h_\kappa^2(x)$ is homogeneous of degree 2γ , which motivates the definition of mixed norm spaces as in (1.5).

In the Dunkl setting we have a *Funk–Hecke formula for h -harmonics*. To state such formula, we need to recall the intertwining operator in the Dunkl setting. It is known that there is an operator V_κ satisfying $T_j V_\kappa = V_\kappa \frac{\partial}{\partial x_j}$. However, the explicit form of V_κ is not known, except in a couple of simple cases, but it is a useful operator. In particular, the Dunkl kernel is given by $E_\kappa(x, y) = V_\kappa e^{\langle \cdot, y \rangle}(x)$. The Funk–Hecke formula for h -harmonics is as follows (see [6, Theorem 7.2.7] or [10, Theorem 5.3.4]).

Theorem 2.1 (Funk–Hecke for h -harmonics). *Let f be a continuous function defined on $[-1, 1]$ and $\lambda = \frac{d}{2} + \gamma - 1$. Then for every $Y_{m,j} \in \mathcal{H}_m$,*

$$\int_{\mathbf{S}^{d-1}} V_\kappa f(\langle x', \cdot \rangle)(y') Y_{m,j}(y') h_\kappa^2(y') d\sigma(y') = \Lambda_m^\kappa(f) Y_{m,j}(x')$$

where $\Lambda_m^\kappa(f)$ is a constant defined by

$$\Lambda_m^\kappa(f) = \frac{\omega_d^\kappa \Gamma(\lambda + 1)}{\sqrt{\pi} \Gamma(\lambda + 1/2)} \int_{-1}^1 f(u) P_m^\lambda(u) (1 - u^2)^{\lambda - \frac{1}{2}} du$$

with

$$(2.4) \quad \omega_d^\kappa := \int_{\mathbf{S}^{d-1}} h_\kappa^2(\omega) d\sigma(\omega).$$

By applying Theorem 2.1 to the function $f(t) = e^{rst}$, $r, s \geq 0$, and using the fact $V_\kappa f(\langle x', y' \rangle) = E_\kappa(rx', sy')$, we immediately obtain the following.

Corollary 2.2 (Funk–Hecke for the Dunkl kernel). *Let $\lambda = \frac{d}{2} + \gamma - 1$. Then for every $Y_{m,j} \in \mathcal{H}_m$,*

$$\begin{aligned} & \int_{\mathbf{S}^{d-1}} E_\kappa(rx', sy') Y_{m,j}(y') h_\kappa^2(y') d\sigma(y') \\ &= \frac{\omega_d^\kappa \Gamma(\frac{d}{2} + \gamma)}{\sqrt{\pi} \Gamma(\frac{d-1}{2} + \gamma)} \left(\int_{-1}^1 e^{rsu} P_m^\lambda(u) (1 - u^2)^{\lambda - \frac{1}{2}} du \right) Y_{m,j}(x'), \end{aligned}$$

where ω_d^κ is as in (2.4).

3. CESÀRO MEANS FOR THE STANDARD HERMITE AND EXTENSION TO COMPLEX PARAMETERS

One of the key points to prove Theorem 1.1 is to express Cesàro kernels for the Dunkl–Hermite expansions in terms of Cesàro kernels for the standard Hermite expansions and then extend these operators for complex values of the parameters involved. In this section we recall some basic results concerning the L^p boundedness of Cesàro means for the standard Hermite expansions and prove others concerning the extended operators.

3.1. Cesàro and Bochner–Riesz means for the standard Hermite expansions. As explained in the introduction, when $\kappa \equiv 0$ the Dunkl–Hermite functions reduce to the standard Hermite functions Φ_μ on \mathbb{R}^d . Cesàro means of order $\delta \geq 0$ associated with the Hermite expansions (or just standard Cesàro means) are then defined by

$$\sigma_N^\delta f = \frac{1}{A_N^\delta} \sum_{j=0}^N A_{N-j}^\delta P_j f,$$

$P_j f$ being the corresponding projections. The operators σ_N^δ can be described as integrals operators with a kernel $\sigma_N^\delta(x, y)$, which is explicitly given by

$$\sigma_N^\delta(x, y) = \frac{1}{A_N^\delta} \sum_{j=0}^N A_{N-j}^\delta \Phi_j(x, y),$$

where $\Phi_j(x, y)$ is the kernel of the j th projection associated with the Hermite operator (see [21, page 6]). For $|w| < 1$, Mehler’s formula for $\Phi_j(x, y)$ reads as

$$(3.1) \quad \sum_{j=0}^{\infty} \Phi_j(x, y) w^j = \pi^{-\frac{d}{2}} (1 - w^2)^{-\frac{d}{2}} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (|x|^2 + |y|^2) + \frac{2w}{1-w^2} x \cdot y}.$$

From Mehler’s formula it follows that $\Phi_j(rx', sy')$, and consequently $\sigma_N^\delta(rx', sy')$, is a function of r, s and $u := x' \cdot y'$. Hence, sometimes we will write $\Phi_j(r, s; u)$ instead of $\Phi_j(rx', sy')$ and $\sigma_N^\delta(r, s; u)$ instead of $\sigma_N^\delta(rx', sy')$.

We also introduce the Bochner–Riesz means associated with the Hermite expansions as

$$S_R^\delta f = \sum_{j=0}^{\infty} \left(1 - \frac{2j+d}{R} \right)_+^\delta P_j f,$$

where $R > 0$ and $(1 - s^2)_+ = \max\{1 - s^2, 0\}$. In the literature, the boundedness of both Cesàro and Bochner–Riesz means have been studied. Their behaviour are similar in the sense that is possible to express the Cesàro means $\sigma_N^\delta f$ in terms of $S_R^\delta f$ and vice-versa. Indeed, we have the following theorem due to J. J. Gergen [11], adapted to our context.

Theorem 3.1 (Gergen). *Let m be the integral part of δ . Then there exist two functions U and V , $U(x) = O(x^{-2})$ as $x \rightarrow \infty$, $U(x) = O(x^{m-\delta+1})$ as $x \rightarrow 0$, and $V(x) = O(x^{-2})$ as $x \rightarrow \infty$, $V(x) = O(x^\delta)$ as $x \rightarrow 0$, such that*

$$S_R^\delta(x, y) = R^{-\delta} \sum_{k \leq R} V(R - k) A_k^\delta \sigma_k^\delta(x, y)$$

and

$$\sigma_N^\delta(x, y) = \frac{1}{A_N^\delta} \int_0^{N+1} U(N + 1 - t) t^\delta S_t^\delta(x, y) dt, \quad \text{for } N = 0, 1, \dots$$

In view of Gergen’s theorem, we can readily prove a version of [21, Theorem 3.3.3] (that states pointwise estimates for Bochner–Riesz) for Cesàro means.

Theorem 3.2. *Let $d \geq 2$ and $\delta > \frac{d-1}{2}$. Then for any $q \geq 2$ and $f \in L^q(\mathbb{R}^d)$ we have the pointwise inequality*

$$\sup_N |\sigma_N^\delta f(x)| \leq CM_q f(x)$$

where $M_q f(x) = (M|f|^q)^{\frac{1}{q}}(x)$, M being the Hardy–Littlewood maximal function.

3.2. Extension of standard Cesàro means to complex parameters. In order to prove Theorem 1.1 we need to consider Cesàro means σ_N^δ when δ is complex. Let us define

$$(3.2) \quad \delta(\zeta) := \frac{d-1}{2} + \zeta \quad \text{for } \zeta \in \mathbb{C}$$

and consider $\sigma_N^{\delta(\zeta)} f$ with $\operatorname{Re}(\delta(\zeta)) \geq 0$. Let us also recall the definition of a function of admissible growth: we say that a function $F(y)$, $y \in \mathbb{R}$, is of admissible growth if there exist constants $a < \pi$ and $b > 0$ such that $|F(y)| \leq e^{be^{a|y|}}$. As a consequence of Theorem 3.2 we get the following result.

Theorem 3.3. *Let $d \geq 2$. Then for any $q \geq 2$ and $f \in L^q(\mathbb{R}^d)$ we have the pointwise inequality*

$$(3.3) \quad \sup_N |\sigma_N^{\delta(i\beta+\varepsilon)} f(x)| \leq C_\varepsilon(\beta) M_q f(x)$$

for a fixed $\varepsilon > 0$ and $\beta \in \mathbb{R}$. Here, the function $C_\varepsilon(\beta)$ is of admissible growth. Moreover, the operator $\sup_N |\sigma_N^{\delta(i\beta+\varepsilon)} f|$ is bounded on $L^p(\mathbb{R}^d)$ for any $p > 2$.

Proof. From the definition of σ_N^δ it follows that

$$(3.4) \quad \sigma_N^{\delta(i\beta+\varepsilon)} f(x) = \frac{1}{A_N^{\delta(i\beta+\varepsilon)}} \sum_{j=0}^N A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1} A_j^{\delta(\varepsilon/2)} \sigma_j^{\delta(\varepsilon/2)} f(x).$$

Indeed, the right hand side in the identity above is

$$\frac{1}{A_N^{\delta(i\beta+\varepsilon)}} \sum_{j=0}^N A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1} \sum_{k=0}^j A_{j-k}^{\delta(\varepsilon/2)} P_k f(x) = \frac{1}{A_N^{\delta(i\beta+\varepsilon)}} \sum_{k=0}^N P_k f(x) \sum_{j=k}^N A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1} A_{j-k}^{\delta(\varepsilon/2)},$$

so (3.4) will be proved once we check that

$$A_{N-k}^{\delta(i\beta+\varepsilon)} = \sum_{j=k}^N A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1} A_{j-k}^{\delta(\varepsilon/2)},$$

or equivalently

$$(3.5) \quad A_N^{\delta(i\beta+\varepsilon)} = \sum_{j=0}^N A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1} A_j^{\delta(\varepsilon/2)}.$$

Recall the following basic facts concerning generating functions. Since $(1-w)^{-r-1} = \sum_{k \geq 0} \binom{r+k}{k} w^k$ and $(1-w)^{-s-1} = \sum_{k \geq 0} \binom{s+k}{k} w^k$, if we multiply these together, we get $(1-w)^{-(r+s+1)-1} = (1-w)^{-r-1}(1-w)^{-s-1}$. Equating coefficients gives us

$$(3.6) \quad \binom{r+s+1+n}{n} = \sum_{k=0}^n \binom{r+n-k}{n-k} \binom{s+k}{k}.$$

With this, we see that (3.5) is true.

Also, by Lemma 7.1, we have that

$$\sum_{j=0}^N |A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1}| |A_j^{\delta(\varepsilon/2)}| \leq C_\varepsilon(\beta) |A_N^{\delta(i\beta+\varepsilon)}|.$$

Now, since $\delta(\varepsilon/2) > \frac{d-1}{2}$, it follows from Theorem 3.2 that $\sup_j |\sigma_j^{\delta(\varepsilon/2)} f(x)| \leq CM_q f(x)$, so (3.3) is proved.

Finally, the last statement of Theorem 3.3 follows from the fact that the maximal function M_q is bounded on $L^p(\mathbb{R}^d)$ for any $p > q \geq 2$. \square

3.3. A vector-valued inequality for an operator connected to standard Cesàro means.

Further, we introduce one more operator \mathbb{S}_N^δ related to σ_N^δ as follows. For any $f \in L^p(\mathbb{R}^+, dr)$ we define

$$(3.7) \quad \mathbb{S}_N^\delta f(r) := r^{\frac{d-1}{p}} \int_0^\infty s^{\frac{d-1}{p'}} \left(\int_{-1}^1 |\sigma_N^\delta(r, s; u)| (1-u^2)^{\frac{d-3}{2}} du \right) f(s) ds.$$

We will require a vector-valued inequality for the maximal function associated with \mathbb{S}_N^δ . In order to get this, recall the following result about vector-valued extensions for general bounded operators by J. L. Krivine (see [13] or [16, Thm. 1.f.14]).

Theorem 3.4 (Krivine). *Let X and Y be two Banach lattices and let $T : X \mapsto Y$ be a bounded linear operator. Then, for every choice of $\{x_i\}_{i=1}^n$ in X , we have*

$$\left\| \left(\sum_{i=1}^n |Tx_i|^2 \right)^{1/2} \right\| \leq K_G \|T\| \left\| \left(\sum_{i=1}^n |x_i|^2 \right)^{1/2} \right\|,$$

where K_G is the universal Grothendieck constant.

Theorem 3.5. *Let $d \geq 2$ and for $\zeta \in \mathbb{C}$, let $\delta(\zeta)$ be defined as in (3.2). Then for any $p > 2$ we have the vector-valued inequality*

$$\left(\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} \left(\sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon)} f_{m,j}(r)| \right)^2 \right)^{\frac{p}{2}} dr \right)^{\frac{1}{p}} \leq C(\beta) \left(\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |f_{m,j}(r)|^2 \right)^{\frac{p}{2}} dr \right)^{\frac{1}{p}}$$

for any sequence of functions $f_{m,j} \in L^p(\mathbb{R}^+, dr)$ for which the right hand side is finite. Moreover, $C(\beta)$ is of admissible growth.

Proof. Observe that we can consider $\sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon)} f|$ as a linear operator mapping $L^p(\mathbb{R}^+)$ into $L^p(\mathbb{R}^+, l^\infty(\mathbb{N}))$. Then, appealing to Theorem 3.4, it is enough to show that this maximal operator is bounded on $L^p(\mathbb{R}^+, dr)$.

Consider the radial function F defined by $F(x) = |x|^{-\frac{d-1}{p}} f(|x|)$ which belongs to $L^p(\mathbb{R}^d)$ as $f \in L^p(\mathbb{R}^+, dr)$. In terms of F we can consider $\mathbb{S}_N^{\delta(i\beta+\varepsilon)} f$ as a radial function on \mathbb{R}^d given by

$$|x|^{\frac{d-1}{p}} \int_{\mathbb{R}^d} |\sigma_N^{\delta(i\beta+\varepsilon)}(x, y)| F(y) dy.$$

Thus the boundedness of $\sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon)} f|$ on $L^p(\mathbb{R}^+, dr)$ follows from the fact that the maximal operator

$$\sup_N \left(\int_{\mathbb{R}^d} |\sigma_N^{\delta(i\beta+\varepsilon)}(x, y)| |F(y)| dy \right)$$

is bounded on $L^p(\mathbb{R}^d)$, which is true by Theorem 3.3. The theorem is proved. \square

4. CESÀRO MEANS FOR THE DUNKL–HERMITE AND EXTENSION TO COMPLEX PARAMETERS

In this section we establish various results in several steps:

- (1) First, we recall the Mehler’s formula in the Dunkl–Hermite setting, analogous to the Mehler’s formula for the standard Hermite functions (3.1).
- (2) In the second step, we are able to write the Cesàro kernels for Dunkl–Hermite expansions in terms of Cesàro kernels for Hermite expansions with shifted parameters. As a result, by using polar coordinates, we express the Dunkl–Cesàro means as an expansion in h -harmonics of linear operators defined in terms of standard Cesàro means, see Corollary 4.2.
- (3) In turn, we show in Proposition 4.4 that the standard shifted Cesàro kernels are connected to Cesàro kernels for Laguerre expansions.
- (4) Finally, we need to extend this connection to complex values of the parameters in Proposition 4.5.

These facts are shown in this order in the subsequent subsections. They will lead us to the proof of Theorem 1.2 and, as a consequence, of Theorem 1.1. Since the proofs of Propositions 4.4 and 4.5 are rather technical, we will state the results here, and the proofs are moved to a separate Section 6.

4.1. Mehler’s formula for Dunkl–Hermite functions. As stated in the introduction, generalized Hermite polynomials and generalized Hermite functions associated with Coxeter groups were studied in [17], where the precise definitions can be found. For our purposes, the most important result is the generating function identity or the *Mehler’s formula for the Dunkl–Hermite functions*. For $|w| < 1$, one has (see [17, Theorem 3.12])

$$(4.1) \quad \sum_{\mu \in \mathbb{N}^d} \Phi_{\mu, \kappa}(x) \Phi_{\mu, \kappa}(y) w^{|\mu|} = \frac{2}{\omega_d^\kappa \Gamma(\frac{d}{2} + \gamma)} (1 - w^2)^{-\frac{d}{2} - \gamma} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (|x|^2 + |y|^2)} E_\kappa \left(\frac{2wx}{1-w^2}, y \right).$$

4.2. Cesàro kernels for Dunkl–Hermite in terms of Cesàro kernels for Hermite. In this subsection we obtain an expression for the Cesàro kernel for the Dunkl–Hermite expansions in terms of the Cesàro kernel for the standard Hermite expansions, via the Funk–Hecke formula stated in Corollary 2.2.

Let $P_{j, \kappa}$ be the orthogonal projection described in (1.1). Then $P_{j, \kappa}$ is given by the kernel operator

$$(4.2) \quad P_{j, \kappa} f(x) = \int_{\mathbb{R}^d} \Phi_{j, \kappa}(x, y) f(y) h_\kappa^2(y) dy$$

where

$$\Phi_{j, \kappa}(x, y) = \sum_{|\mu|=j} \Phi_{\mu, \kappa}(x) \Phi_{\mu, \kappa}(y).$$

For $\delta \geq 0$, it is clear from (1.2) and (4.2) that the kernel of the Cesàro means $\sigma_{N, \kappa}^\delta$ is given by

$$(4.3) \quad \sigma_{N, \kappa}^\delta(x, y) := \frac{1}{A_N^\delta} \sum_{j=0}^N A_{N-j}^\delta \Phi_{j, \kappa}(x, y).$$

Proposition 4.1 below contains expressions for the h -harmonic coefficients of Dunkl–Cesàro kernels in terms of d -dimensional and $(d+1)$ -dimensional standard Cesàro kernels. As remarked earlier, in view of the Mehler’s formula (3.1) it follows that $\Phi_j(rx', sy')$ is a function of r, s and $u := x' \cdot y'$. The same is true for the Cesàro kernels and so again we will sometimes write $\Phi_j(r, s; u)$ instead of $\Phi_j(rx', sy')$ and $\sigma_N^\delta(r, s; u)$ instead of $\sigma_N^\delta(rx', sy')$. Let P_m^λ stand for the normalized ultraspherical polynomials of type $\lambda > -\frac{1}{2}$ and degree m . We have the following.

Proposition 4.1 (Funk–Hecke for Cesàro–Dunkl–Hermite). *For γ defined as in (2.3), let*

$$c_{d,\gamma} = \frac{2\pi^{(d-1)/2}}{\Gamma(\frac{d-1}{2} + \gamma)}$$

and $\lambda = \frac{d}{2} + \gamma - 1$. Then, for any spherical h -harmonic $Y_{m,\ell}$ of degree m , we have the identity

$$(4.4) \quad \int_{\mathbf{S}^{d-1}} \sigma_{N,\kappa}^{\delta,d}(rx', sy') Y_{m,\ell}(y') h_\kappa^2(y') d\sigma(y') \\ = \frac{c_{d,\gamma}}{A_N^\delta} \sum_{j=0}^{[N/2]} A_j^{\gamma-1} A_{N-2j}^\delta \left(\int_{-1}^1 \sigma_{N-2j}^{\delta,d}(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du \right) Y_{m,\ell}(x').$$

Moreover, we also have

$$(4.5) \quad \int_{\mathbf{S}^{d-1}} \sigma_{N,\kappa}^{\delta,d}(rx', sy') Y_{m,\ell}(y') h_\kappa^2(y') d\sigma(y') \\ = \frac{c_{d+1,\gamma-\frac{1}{2}}}{A_N^\delta} \sum_{j=0}^{[N/2]} A_j^{\gamma-\frac{3}{2}} A_{N-2j}^\delta \left(\int_{-1}^1 \sigma_{N-2j}^{\delta,d+1}(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du \right) Y_{m,\ell}(x').$$

Proof. By integrating Mehler’s formula for the Dunkl–Hermite functions (4.1) against $Y_{m,\ell}(y')$ and using the Funk–Hecke formula in Corollary 2.2 we get

$$(4.6) \quad \sum_{N=0}^{\infty} w^N \int_{\mathbf{S}^{d-1}} \Phi_{N,\kappa}^{(d)}(rx', sy') Y_{m,\ell}(y') h_\kappa^2(y') d\sigma(y') \\ = \frac{2(1-w^2)^{-d/2-\gamma}}{\sqrt{\pi} \Gamma(\frac{d-1}{2} + \gamma)} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (r^2+s^2)} \left(\int_{-1}^1 e^{\frac{2rsu}{1-w^2}} P_m^\lambda(u) (1-u^2)^{\lambda-\frac{1}{2}} du \right) Y_{m,\ell}(x').$$

Now, comparing the right hand side of the above with the Mehler’s formula for the Hermite expansions (3.1), we conclude that

$$(4.7) \quad \sum_{N=0}^{\infty} w^N \int_{\mathbf{S}^{d-1}} \Phi_{N,\kappa}^{(d)}(rx', sy') Y_{m,\ell}(y') h_\kappa^2(y') d\sigma(y') \\ = \frac{2\pi^{(d-1)/2} (1-w^2)^{-\gamma}}{\Gamma(\frac{d-1}{2} + \gamma)} \sum_{N=0}^{\infty} \left(\int_{-1}^1 \Phi_N^{(d)}(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-\frac{1}{2}} du \right) w^N Y_{m,\ell}(x').$$

By multiplying the identity (4.7) by $(1-w)^{-\delta-1}$ and using the expansion $(1-w)^{-\delta-1} = \sum_{N=0}^{\infty} A_N^\delta w^N$, we obtain

$$\sum_{N=0}^{\infty} w^N \sum_{j=0}^N A_{N-j}^\delta \left(\int_{\mathbf{S}^{d-1}} \Phi_{j,\kappa}^{(d)}(rx', sy') Y_{m,\ell}(y') h_\kappa^2(y') d\sigma(y') \right) \\ = c_{d,\gamma} (1-w^2)^{-\gamma} \sum_{N=0}^{\infty} w^N \sum_{j=0}^N A_{N-j}^\delta \left(\int_{-1}^1 \Phi_j^{(d)}(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-\frac{1}{2}} du \right) Y_{m,\ell}(x') \\ = c_{d,\gamma} (1-w^2)^{-\gamma} \sum_{N=0}^{\infty} w^N A_N^\delta \left(\int_{-1}^1 \sigma_N^{\delta,d}(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du \right) Y_{m,\ell}(x').$$

Since $(1 - w^2)^{-\gamma} = \sum_{N=0}^{\infty} A_N^{\gamma-1} w^{2N}$, the above can be written as

$$\begin{aligned} & \sum_{N=0}^{\infty} w^N \sum_{j=0}^N A_{N-j}^{\delta} \left(\int_{\mathbf{S}^{d-1}} \Phi_{j,\kappa}^{(d)}(rx', sy') Y_{m,\ell}(y') h_{\kappa}^2(y') d\sigma(y') \right) \\ &= c_{d,\gamma} \sum_{N=0}^{\infty} w^N \sum_{j=0}^{[N/2]} A_j^{\gamma-1} A_{N-2j}^{\delta} \left(\int_{-1}^1 \sigma_{N-2j}^{\delta,d}(r, s; u) P_m^{\lambda}(u) (1 - u^2)^{\lambda-1/2} du \right) Y_{m,\ell}(x'). \end{aligned}$$

Comparing the coefficients of w^N on both sides in the above and in view of (4.3) we obtain (4.4).

Observe that the right hand side of (4.6) remains the same if we replace d by $d + 1$ and γ by $\gamma - \frac{1}{2}$. Consequently, in view of Mehler's formula for the Hermite expansions (3.1), the right hand side of (4.6) is also equal to

$$\frac{2\pi^{d/2}(1 - w^2)^{-\gamma+\frac{1}{2}}}{\Gamma\left(\frac{d-1}{2} + \gamma\right)} \sum_{N=0}^{\infty} \left(\int_{-1}^1 \Phi_N^{(d+1)}(r, s; u) P_m^{\lambda}(u) (1 - u^2)^{\lambda-\frac{1}{2}} du \right) w^N Y_{m,\ell}(y').$$

Therefore, by repeating an analogous reasoning, we also arrive at (4.5).

This completes the proof of the proposition. \square

By expanding $\sigma_{N,\kappa}^{\delta} f(x)$ in terms of h -harmonics and using Proposition 4.1 we can easily deduce the following.

Corollary 4.2. *For any Schwartz class function f we have*

$$\sigma_{N,\kappa}^{\delta} f(rx') = \sum_{m=0}^N \sum_{j=1}^{d(m)} T_{N,m}^{\delta,\gamma} f_{m,j}(r) Y_{m,j}(x'),$$

where $T_{N,m}^{\delta,\gamma}$ are the operators defined in (1.6) with kernels $K_{N,m}^{\delta,\gamma}$ that can be expressed either as

$$(4.8) \quad K_{N,m}^{\delta,\gamma}(r, s) := \frac{c_{d,\gamma}}{A_N^{\delta}} \sum_{j=0}^{[N/2]} A_j^{\gamma-1} A_{N-2j}^{\delta} \int_{-1}^1 \sigma_{N-2j}^{\delta,d}(rx', sy') P_m^{\lambda}(u) (1 - u^2)^{\lambda-1/2} du$$

or

$$(4.9) \quad K_{N,m}^{\delta,\gamma}(r, s) = \frac{c_{d+1,\gamma-\frac{1}{2}}}{A_N^{\delta}} \sum_{j=0}^{[N/2]} A_j^{\gamma-\frac{3}{2}} A_{N-2j}^{\delta} \int_{-1}^1 \sigma_{N-2j}^{\delta,d+1}(r, s; u) P_m^{\lambda}(u) (1 - u^2)^{\lambda-1/2} du.$$

Remark 4.3. Therefore, we can express the kernels $K_{N,m}^{\delta,\gamma}(r, s)$ and consequently the operators $T_{N,m}^{\delta,\gamma} f$ in terms of both d -dimensional and $(d+1)$ -dimensional Hermite Cesàro kernels. Indeed, the expression (4.8) comes from the use of identity (4.4) in Proposition 4.1, and the expression (4.9) is obtained in view of identity (4.5) in Proposition 4.1.

4.3. The Laguerre connection. We are going to express the kernel $K_{N,m}^{\delta,\gamma}(r, s)$ in (4.8) in terms of Laguerre functions. This is contained in Proposition 4.4 below. This result will not be enough for our purposes; an improved version with complex parameters is required, see Proposition 4.5. Nevertheless, we first show the non-complex version in detail since it is interesting to understand that, philosophically, there is an underlying phenomenon of transplantation.

Let us recall some basic facts about Laguerre functions. For $\alpha > -1$, let ψ_k^{α} be the normalized Laguerre functions given by

$$\psi_k^{\alpha}(r) = \left(\frac{2\Gamma(k+1)}{\Gamma(k+\alpha+1)} \right)^{\frac{1}{2}} L_k^{\alpha}(r^2) e^{-\frac{1}{2}r^2}, \quad k = 0, 1, \dots,$$

where L_k^α are the Laguerre polynomials of order α , see [15, page 76].

We have the following generating function identity for Laguerre functions (see [21, (1.1.47)]). For $|w| < 1$,

$$\sum_{k=0}^{\infty} \psi_k^\alpha(r) \psi_k^\alpha(s) w^{2k} = 2(1-w^2)^{-1} (rsw)^{-\alpha} e^{-\frac{1}{2}\left(\frac{1+w^2}{1-w^2}\right)(r^2+s^2)} I_\alpha\left(\frac{2rsw}{1-w^2}\right).$$

Here, I_α is the modified Bessel function, see Subsection 7.2 for the definition and further results concerning these functions. From this identity, we easily deduce that

$$(4.10) \quad \sum_{k=0}^{\infty} (rs)^m \psi_k^{\alpha+m}(r) \psi_k^{\alpha+m}(s) w^{2k+m} = 2(1-w^2)^{-1} e^{-\frac{1}{2}\left(\frac{1+w^2}{1-w^2}\right)(r^2+s^2)} (rsw)^{-\alpha} I_{\alpha+m}\left(\frac{2rsw}{1-w^2}\right).$$

Proposition 4.4. *For $d \geq 2$, let $\lambda = \frac{d}{2} + \gamma - 1$. Then the kernel $K_{N,m}^{\delta,\gamma}$ expressed either in (4.8) or in (4.9) can be written as*

$$K_{N,m}^{\delta,\gamma}(r, s) = \frac{1}{A_N^\delta} \sum_{j=0}^{[(N-m)/2]} A_{N-m-2j}^\delta (rs)^m \psi_j^{\lambda+m}(r) \psi_j^{\lambda+m}(s)$$

provided $N \geq m$. For other values of N , $K_{N,m}^{\delta,\gamma}(r, s) = 0$.

4.4. Extension to complex parameters. In order to prove Theorem 1.2 we need to extend our operators for complex parameters and get proper expressions for them.

First, observe that the kernels $K_{N,m}^{\delta,\gamma}$ in (4.8), initially expressed in terms of standard Cesàro kernels, are also expressible in terms of Laguerre functions $\psi_k^{\lambda+m}$ by Proposition 4.4. Then, since Laguerre functions of complex parameters are well defined, the kernels $K_{N,m}^{\delta,\gamma}$ make sense even if δ and γ are complex. For $\zeta \in \mathbb{C}$ let us define

$$(4.11) \quad \lambda(\zeta) := \frac{d}{2} + \zeta - 1.$$

Fix $\varepsilon > 0$ and consider the sequence of operators

$$(4.12) \quad \mathcal{T}_{N,m}^{\varepsilon,\zeta} f(r) := \int_0^\infty \mathcal{K}_{N,m}^{\varepsilon,\zeta}(r, s) f(s) ds$$

defined for $f \in L^p(\mathbb{R}^+, dr)$ with kernel

$$(4.13) \quad \mathcal{K}_{N,m}^{\varepsilon,\zeta}(r, s) := r^{\frac{2\delta(\zeta)}{p}} s^{\frac{2\delta(\zeta)}{p'}} K_{N,m}^{\delta(\zeta+\varepsilon),\zeta}(r, s)$$

where $p > 2$. Then, we can go back and prove that the kernels $\mathcal{K}_{N,m}^{\varepsilon,\zeta}(r, s)$, that are now defined in terms of Laguerre functions of complex parameters, can be expressed also in terms of the d -dimensional and $(d+1)$ -dimensional kernels of Cesàro means involving parameters with certain values of ζ . This result is contained in Proposition 4.5.

Observe that the operators $\mathcal{T}_{N,m}^{\varepsilon,\zeta}$ can be understood as “extensions”, for complex parameters, of the operators $T_{N,m}^{\delta,\gamma}$. In turn, the kernels $\mathcal{K}_{N,m}^{\varepsilon,\zeta}$ can be understood as “extensions”, for complex parameters, of the kernels $K_{N,m}^{\delta,\gamma}$.

Proposition 4.5. Let $C_d = \frac{4\pi^{(d-1)/2}}{\Gamma((d-1)/2)}$. For any $\beta \in \mathbb{R}$ we have

$$(4.14) \quad \mathcal{K}_{N,m}^{\varepsilon,i\beta}(r,s) = \frac{C_d}{A_N^{\delta(i\beta+\varepsilon)}} \sum_{k=0}^N \binom{\varepsilon/2}{N-k} \\ \times \sum_{j=0}^{\lfloor k/2 \rfloor} A_j^{\frac{\varepsilon}{2}+i\beta-1} A_{k-2j}^{\delta(i\beta+\frac{\varepsilon}{2})} r^{\frac{2\delta(i\beta)}{p}} s^{\frac{2\delta(i\beta)}{p'}} \int_{-1}^1 \sigma_{k-2j}^{\delta(i\beta+\frac{\varepsilon}{2}),d}(r,s;u) Q_m^{\lambda(i\beta)}(u) du$$

and

$$(4.15) \quad \mathcal{K}_{N,m}^{\varepsilon,\frac{1}{2}+i\beta}(r,s) = \frac{C_{d+1}}{A_N^{\delta(\frac{1}{2}+i\beta+\varepsilon)}} \sum_{k=0}^N \binom{\varepsilon/2}{N-k} \\ \times \sum_{j=0}^{\lfloor k/2 \rfloor} A_j^{\frac{\varepsilon}{2}+i\beta-1} A_{k-2j}^{\delta(\frac{1}{2}+i\beta+\frac{\varepsilon}{2})} r^{\frac{2\delta(\frac{1}{2}+i\beta)}{p}} s^{\frac{2\delta(\frac{1}{2}+i\beta)}{p'}} \int_{-1}^1 \sigma_{k-2j}^{\delta(\frac{1}{2}+i\beta+\frac{\varepsilon}{2}),d+1}(r,s;u) Q_m^{\lambda(\frac{1}{2}+i\beta)}(u) du.$$

Here, the functions $Q_m^{\lambda(\zeta)}(u)$ are the ones defined in (7.2).

5. MIXED NORM ESTIMATES FOR THE CESÀRO MEANS: PROOFS OF THEOREMS 1.1 AND 1.2

In this Section we prove the main theorems. Theorem 1.1 easily follows from Theorem 1.2. In its turn, the proof of Theorem 1.2 needs the ingredients from the previous sections, and the different operators and their extensions with complex parameters play a role. Since the quantity of operators and the corresponding kernels is large and the notation cumbersome, we present the Table 1 with a summary of notation and references that will make the reading of this section more comfortable.

TABLE 1. Summary of operators and kernels

$T_{N,m}^{\delta,\gamma}$	operator defined in (1.6)
$K_{N,m}^{\delta,\gamma}$	kernel of $T_{N,m}^{\delta,\gamma}$
$\mathcal{T}_{N,m}^{\varepsilon,\zeta}$	operator defined in (4.12)
$\mathcal{K}_{N,m}^{\varepsilon,\zeta}$	kernel of $\mathcal{T}_{N,m}^{\varepsilon,\zeta}$, related to $K_{N,m}^{\delta,\gamma}$ by (4.13)
S_N^δ	operator defined in (3.7) in terms of the standard Cesàro means σ_N^δ

Proof of Theorem 1.1. In view of Corollary 4.2 we have that

$$\begin{aligned} \|\sigma_{N,\kappa}^\delta f\|_{(p,2)} &= \left(\int_0^\infty \left(\int_{\mathbf{S}^{d-1}} |\sigma_{N,\kappa}^\delta f(rx')|^2 h_\kappa^2(x') d\sigma(x') \right)^{\frac{p}{2}} r^{d+2\gamma-1} dr \right)^{\frac{1}{p}} \\ &= \left(\int_0^\infty \left(\int_{\mathbf{S}^{d-1}} \left| \sum_{m=0}^\infty \sum_{j=1}^{d(m)} |T_{N,m}^{\delta,\gamma} f_{m,j}(r) Y_{m,j}(x')|^2 h_\kappa^2(x') d\sigma(x') \right)^{\frac{p}{2}} r^{d+2\gamma-1} dr \right)^{\frac{1}{p}} \\ &= \left\| \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |T_{N,m}^{\delta,\gamma} f_{m,j}(r)|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^+, r^{d+2\gamma-1} dr)}. \end{aligned}$$

Then, with Theorem 1.2, we can conclude. \square

Therefore, it remains to prove Theorem 1.2. Given sequences of functions $(f_{m,j})$ and $(g_{m,j})$ such that

$$\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |f_{m,j}(r)|^2 \right)^{\frac{p}{2}} r^{d+2\gamma-1} dr = \int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |g_{m,j}(r)|^2 \right)^{\frac{p'}{2}} r^{d+2\gamma-1} dr = 1$$

we define $\tilde{f}_{m,j}(r) = f_{m,j}(r)r^{\frac{2\delta(\gamma)}{p}}$ and $\tilde{g}_{m,j}(r) = g_{m,j}(r)r^{\frac{2\delta(\gamma)}{p'}}$, where $\delta(\gamma)$ is as in (3.2). Then it follows that

$$\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\tilde{f}_{m,j}(r)|^2 \right)^{\frac{p}{2}} dr = \int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\tilde{g}_{m,j}(r)|^2 \right)^{\frac{p'}{2}} dr = 1.$$

We consider the function $F_N(\zeta)$ on the strip $0 \leq \operatorname{Re}(\zeta) \leq \frac{1}{2}$ defined by

$$F_N(\zeta) := \int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} \mathcal{T}_{N,m}^{\varepsilon,\zeta} \tilde{f}_{m,j}(r) \overline{\tilde{g}_{m,j}(r)} \right) dr.$$

It is clear that when $f_{m,j}$ are compactly supported on \mathbb{R}^+ , $F_N(\zeta)$ is a holomorphic function in the interior of the strip $0 \leq \operatorname{Re}(\zeta) \leq \frac{1}{2}$ and continuous up to the boundary. We will prove that $F_N(\operatorname{Re}(\zeta))$ is bounded in the strip. In order to do that, we need the following variant of three lines lemma proved in [20, Ch. V, Lemma 4.2].

Lemma 5.1 (Stein-Weiss). *Suppose F is a function defined and continuous on the unit strip $\mathcal{S} = \{z \in \mathbb{C} : 0 \leq \operatorname{Re}(z) \leq 1\}$ that is analytic in the interior of \mathcal{S} satisfying*

$$\sup_{x+iy \in \mathcal{S}} e^{-a|y|} \log |F(x+iy)| < \infty$$

for some $a < \pi$. Then

$$(5.1) \quad \log |F(x)| \leq \frac{1}{2} \sin \pi x \int_{-\infty}^\infty \left(\frac{\log |F(iy)|}{\cosh \pi y - \cos \pi x} + \frac{\log |F(1+iy)|}{\cosh \pi y + \cos \pi x} \right) dy$$

whenever $0 < x < 1$.

Proposition 5.2. *Let $\zeta := \gamma + i\beta$, with $0 \leq \gamma \leq 1/2$ and $\beta \in \mathbb{R}$. We have*

$$|F_N(i\beta)| \leq C_0(\beta) \quad \text{and} \quad |F_N(\tfrac{1}{2} + i\beta)| \leq C_1(\beta)$$

where $C_0(\beta)$ and $C_1(\beta)$ are independent of N and of admissible growth. Moreover,

$$|F_N(\gamma)| \leq C, \quad \text{for } 0 \leq \gamma \leq 1/2.$$

Proof. In view of (4.14) in Proposition 4.5, the operator $\mathcal{T}_{N,m}^{\varepsilon,i\beta}$ is given by

$$\begin{aligned} \mathcal{T}_{N,m}^{\varepsilon,i\beta} f(r) &= \frac{C_d}{A_N^{\delta(i\beta+\varepsilon)}} \sum_{k=0}^N \binom{\varepsilon/2}{N-k} \sum_{j=0}^{[k/2]} A_j^{\frac{\varepsilon}{2}+i\beta-1} A_{k-2j}^{\delta(i\beta+\frac{\varepsilon}{2})} \\ &\quad \times \int_0^\infty r^{\frac{2\delta(i\beta)}{p}} s^{\frac{2\delta(i\beta)}{p'}} \int_{-1}^1 \sigma_{k-2j}^{\delta(i\beta+\frac{\varepsilon}{2}),d}(r,s;u) Q_m^{\lambda(i\beta)}(u) du f(s) ds. \end{aligned}$$

The operators $\mathcal{T}_{N,m}^{\varepsilon,i\beta}$ can be bounded in terms of the operators $\mathbb{S}_k^{\delta(i\beta+\varepsilon/2),d}$. Indeed, taking into account the fact that $\sum_{j=0}^\infty \binom{\varepsilon/2}{j} < \infty$, we have

$$|\mathcal{T}_{N,m}^{\varepsilon,i\beta} \tilde{f}_{m,j}(r)| \leq C_0(\beta) \sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon/2),d} \tilde{f}_{m,j}(r)|$$

provided we have the estimate

$$\frac{C_d}{|A_N^{\delta(i\beta+\varepsilon)}|} \sum_{j=0}^{[k/2]} |A_j^{\varepsilon/2+i\beta-1} A_{k-2j}^{\delta(i\beta+\varepsilon/2)}| \leq C_0(\beta)$$

for all $0 \leq k \leq N$, $N = 1, 2, 3, \dots$. But this can be proved similarly as in the case of Lemma 7.1, and hence we leave the details to the reader.

With the above observations, we obtain, for $p > 2$,

$$\begin{aligned} |F_N(i\beta)| &\leq \int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\mathcal{T}_{N,m}^{\varepsilon,i\beta} \tilde{f}_{m,j}(r)|^2 \right)^{1/2} \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\tilde{g}_{m,j}(r)|^2 \right)^{1/2} dr \\ &\leq C_0(\beta) \int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} \left(\sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon/2),d} \tilde{f}_{m,j}(r)| \right)^2 \right)^{1/2} \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\tilde{g}_{m,j}(r)|^2 \right)^{1/2} dr \\ &\leq C_0(\beta) \left(\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} \left(\sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon/2),d} \tilde{f}_{m,j}(r)| \right)^2 \right)^{p/2} dr \right)^{1/p} \\ &\quad \times \left(\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\tilde{g}_{m,j}(r)|^2 \right)^{p'/2} dr \right)^{1/p'} \\ &\leq C_0(\beta) \left(\int_0^\infty \left(\sum_{m=0}^\infty \sum_{j=1}^{d(m)} |\tilde{f}_{m,j}(r)|^2 \right)^{p/2} dr \right)^{1/p} = C_0(\beta), \end{aligned}$$

where in the last inequality we applied Theorem 3.5. Similarly, by (4.15) in Proposition 4.5, we get the estimate

$$|\mathcal{T}_{N,m}^{\varepsilon,1/2+i\beta} \tilde{f}_{m,j}(r)| \leq C_1(\beta) \sup_N |\mathbb{S}_N^{\delta(i\beta+\varepsilon/2),d+1} \tilde{f}_{m,j}(r)|.$$

Then, an analogous reasoning with proper modifications, leads us to $|F_N(\frac{1}{2} + i\beta)| \leq C_1(\beta)$.

Observe that with the estimates just proven, we apply Lemma 5.1 to our function F_N , so that right hand side of (5.1) is finite, therefore ensuring the boundedness of $F_N(\gamma)$, for $0 < \gamma < 1/2$. \square

We recall the following result that corresponds with an idea of J. L. Rubio de Francia tacitly contained in [19, Remark (a)]. Here we stated it as a theorem.

Theorem 5.3 (Rubio de Francia). *Let $T : L^p(\mathbb{R}^d) \rightarrow L^p(\mathbb{R}^d)$ be a bounded linear operator which is rotation invariant, i.e. it commutes with the action of the rotation group $SO(d)$ on $L^p(\mathbb{R}^d)$. Then, T is also bounded in $L^p(\ell^2)$.*

We are ready to prove Theorem 1.2.

Proof of Theorem 1.2. We first observe that when $\gamma = 0$ the Dunkl–Hermite expansion reduces to the standard Hermite expansion and hence Theorem 1.2 is true. Indeed, we have $L^p(\mathbb{R}^d)$ boundedness of σ_N^δ for any $\delta > \frac{d-1}{2}$, $1 \leq p < \infty$ (see [22]). On the other hand, it is easy to check that the Cesàro kernels associated to standard Hermite expansions are rotation invariant. Then, from Theorem 5.3, we immediately get the desired vector-valued extension. When $\gamma = \frac{1}{2}$, the same reasoning holds by considering Cesàro means $\sigma_N^{\delta,d+1}$. Hence we can restrict ourselves to the case $0 < \gamma < \frac{1}{2}$.

Given $\delta > \frac{d+2\gamma-1}{2}$ we can choose $\varepsilon > 0$ so that $\delta = \frac{d+2\gamma+2\varepsilon-1}{2} = \delta(\gamma + \varepsilon)$, which is defined in (3.2). It suffices to prove that, for $p > 2$,

$$\left\| \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} |T_{N,m}^{\delta(\gamma+\varepsilon),\gamma} f_{m,j}(r)|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^+, r^{d+2\gamma-1} dr)} \leq C \left\| \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} |f_{m,j}(r)|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^+, r^{d+2\gamma-1} dr)}.$$

We can write

$$\begin{aligned} & \int_0^{\infty} \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} T_{N,m}^{\delta(\gamma+\varepsilon),\gamma} f_{m,j}(r) \overline{g_{m,j}(r)} \right) r^{d+2\gamma-1} dr \\ &= \int_0^{\infty} \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} \int_0^{\infty} K_{N,m}^{\delta(\gamma+\varepsilon),\gamma}(r,s) f_{m,j}(s) s^{d+2\gamma-1} ds \overline{g_{m,j}(r)} \right) r^{d+2\gamma-1} dr \\ &= \int_0^{\infty} \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} \int_0^{\infty} r^{\frac{2\delta(\gamma)}{p}} s^{\frac{2\delta(\gamma)}{p'}} K_{N,m}^{\delta(\gamma+\varepsilon),\gamma}(r,s) f_{m,j}(s) s^{\frac{2\delta(\gamma)}{p}} ds \overline{g_{m,j}(r)} \right) r^{\frac{2\delta(\gamma)}{p'}} dr \\ &= \int_0^{\infty} \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} \int_0^{\infty} \mathcal{K}_{N,m}^{\varepsilon,\gamma}(r,s) \tilde{f}_{m,j}(s) ds \overline{\tilde{g}_{m,j}(r)} \right) dr \\ &= \int_0^{\infty} \left(\sum_{m=0}^{\infty} \sum_{j=1}^{d(m)} T_{N,m}^{\varepsilon,\gamma} \tilde{f}_{m,j}(r) \overline{\tilde{g}_{m,j}(r)} \right) dr = F_N(\gamma), \end{aligned}$$

and the last one is bounded, by Proposition 5.2. By an argument of duality, the estimate is valid for all $1 < p < \infty$. \square

6. PROOFS OF PROPOSITIONS 4.4 AND 4.5

Proof of Proposition 4.4. We integrate Mehler's formula for $\Phi_N(x, y)$ in (3.1) against the function $P_m^\lambda(u)(1-u^2)^{\lambda-1/2}$ over $(-1, 1)$, and apply Lemma 7.2 with $z = \frac{2w}{1-w^2}rs$, so that we get

$$\begin{aligned} & \sum_{N=0}^{\infty} w^N \int_{-1}^1 \Phi_N(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du \\ &= \pi^{-\frac{d}{2}} (1-w^2)^{-\frac{d}{2}} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (r^2+s^2)} \int_{-1}^1 P_m^\lambda(u) (1-u^2)^{\lambda-1/2} e^{zu} du \\ &= \pi^{-(d-1)/2} \Gamma(\lambda+1/2) (1-w^2)^{-\frac{d}{2}} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (r^2+s^2)} \frac{I_{\lambda+m}(z)}{(z/2)^\lambda}. \end{aligned}$$

Comparing this with the generating function identity for the Laguerre functions (4.10) we obtain

$$(6.1) \quad \sum_{N=0}^{\infty} w^N \int_{-1}^1 \Phi_N(r, s; u) P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du = \frac{\Gamma(\lambda+1/2)}{2\pi^{(d-1)/2}} (1-w^2)^\gamma (wrs)^m \sum_{N=0}^{\infty} \psi_N^{\lambda+m}(r) \psi_N^{\lambda+m}(s) w^{2N}.$$

On the one hand, multiplying $\sum_{N=0}^{\infty} w^N \Phi_N(r, s; u)$ by $(1-w)^{-\delta-1}$, we have

$$(6.2) \quad (1-w)^{-\delta-1} \left(\sum_{N=0}^{\infty} w^N \Phi_N(r, s; u) \right) = \sum_{N=0}^{\infty} \left(\sum_{j=0}^N A_{N-j}^\delta \Phi_j(r, s; u) \right) w^N = \sum_{N=0}^{\infty} A_N^\delta \sigma_N^\delta(r, s; u) w^N.$$

On the other hand, we also have

$$(6.3) \quad (1-w)^{-\delta-1} \left(\sum_{N=0}^{\infty} \psi_N^{\lambda+m}(r) \psi_N^{\lambda+m}(s) w^{2N} \right) = \sum_{N=0}^{\infty} \left(\sum_{j=0}^{[N/2]} A_{N-2j}^{\delta} \psi_j^{\lambda+m}(r) \psi_j^{\lambda+m}(s) \right) w^N.$$

Therefore, in view of (6.2) and (6.3), from (6.1) we get

$$\begin{aligned} & \sum_{N=0}^{\infty} A_N^{\delta} \left(\int_{-1}^1 \sigma_N^{\delta}(r, s; u) P_m^{\lambda}(u) (1-u^2)^{\lambda-1/2} du \right) w^N \\ &= \frac{\Gamma(\lambda+1/2)}{2\pi^{(d-1)/2}} (1-w^2)^{\gamma} (wrs)^m \sum_{N=0}^{\infty} \left(\sum_{j=0}^{[N/2]} A_{N-2j}^{\delta} \psi_j^{\lambda+m}(r) \psi_j^{\lambda+m}(s) \right) w^N. \end{aligned}$$

Now, we multiply both sides by $(1-w^2)^{-\gamma}$ and rearrange to get

$$\begin{aligned} & \sum_{N=0}^{\infty} \left(\sum_{j=0}^{[N/2]} A_j^{\gamma-1} A_{N-2j}^{\delta} \int_{-1}^1 \sigma_{N-2j}^{\delta}(r, s; u) P_m^{\lambda}(u) (1-u^2)^{\lambda-1/2} du \right) w^N \\ &= \frac{\Gamma(\lambda+1/2)}{2\pi^{(d-1)/2}} \sum_{N=0}^{\infty} \left(\sum_{j=0}^{[N/2]} A_{N-2j}^{\delta} (rs)^m \psi_j^{\lambda+m}(r) \psi_j^{\lambda+m}(s) \right) w^{N+m} \\ &= \frac{\Gamma(\lambda+1/2)}{2\pi^{(d-1)/2}} \sum_{N=m}^{\infty} \left(\sum_{j=0}^{[(N-m)/2]} A_{N-m-2j}^{\delta} (rs)^m \psi_j^{\lambda+m}(r) \psi_j^{\lambda+m}(s) \right) w^N. \end{aligned}$$

Equating the coefficients of w^N on both sides, by (4.8) we see that

$$K_{N,m}^{\delta,\gamma}(r, s) = \frac{1}{A_N^{\delta}} \sum_{j=0}^{[(N-m)/2]} A_{N-m-2j}^{\delta} (rs)^m \psi_j^{\lambda+m}(r) \psi_j^{\lambda+m}(s)$$

for $N \geq m$. It is clear that $K_{N,m}^{\delta,\gamma}(r, s) = 0$ for $N < m$. This completes the proof of the proposition.

Repeating the same procedure, but starting with Mehler's formula (3.1) for $\Phi_N^{(d+1)}(x, y)$, $x, y \in \mathbb{R}^{d+1}$, we obtain the proposition for the kernel $K_{N,m}^{\delta,\gamma}(r, s)$ expressed as in (4.9). \square

Proof of Proposition 4.5. Let $Q_m^{\lambda(\zeta)}(u)$ be as in Lemma 7.3, where $\lambda(\zeta)$ is taken as in (4.11). We integrate Mehler's formula for $\Phi_N(x, y)$ in (3.1) against $Q_m^{\lambda(i\beta)}(u)$, and use (7.9) with $z = \frac{2w}{1-w^2}rs$, so that

$$\begin{aligned} & \sum_{N=0}^{\infty} w^N \int_{-1}^1 \Phi_N(r, s; u) Q_m^{\lambda(i\beta)}(u) du \\ &= \pi^{-\frac{d}{2}} (1-w^2)^{-\frac{d}{2}} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (r^2+s^2)} \int_{-1}^1 Q_m^{\lambda(i\beta)}(u) e^{zu} du \\ &= \frac{\Gamma((d-1)/2)}{2\pi^{(d-1)/2}} (1-w^2)^{-\frac{d}{2}} e^{-\frac{1}{2} \left(\frac{1+w^2}{1-w^2} \right) (r^2+s^2)} \frac{I_{\lambda(i\beta)+m}(z)}{(z/2)^{\lambda(i\beta)}}. \end{aligned}$$

From this and (4.10) we have

$$(6.4) \quad \sum_{N=0}^{\infty} w^N \int_{-1}^1 \Phi_N(r, s; u) Q_m^{\lambda(i\beta)}(u) du \\ = \frac{\Gamma((d-1)/2)}{4\pi^{(d-1)/2}} (1-w^2)^{i\beta} \sum_{N=0}^{\infty} (rs)^m \psi_N^{\lambda(i\beta)+m}(r) \psi_N^{\lambda(i\beta)+m}(s) w^{2N+m}.$$

Now we follow a similar procedure used in the proof of Proposition 4.4 but with modifications. Namely, instead of multiplying both sides by $(1-w^2)^{-\gamma}(1-w)^{-\delta-1}$ we multiply by the factor $(1-w^2)^{-i\beta}(1-w)^{-\delta(i\beta+\varepsilon)-1}$, which we rewrite as $(1-w^2)^{-\frac{\varepsilon}{2}-i\beta}(1-w)^{-\delta(i\beta+\frac{\varepsilon}{2})-1}(1+w)^{\frac{\varepsilon}{2}}$, where $\delta(\zeta)$ is as in (3.2). Then, left hand side of (6.4) delivers

$$(1-w^2)^{-\frac{\varepsilon}{2}-i\beta}(1-w)^{-\delta(i\beta+\frac{\varepsilon}{2})-1}(1+w)^{\frac{\varepsilon}{2}} \sum_{N=0}^{\infty} w^N \int_{-1}^1 \Phi_N(r, s; u) Q_m^{\lambda(i\beta)}(u) du \\ = (1+w)^{\frac{\varepsilon}{2}} \sum_{N=0}^{\infty} w^N \left(\sum_{j=0}^{[N/2]} A_j^{\frac{\varepsilon}{2}+i\beta-1} A_{N-2j}^{\delta(i\beta+\frac{\varepsilon}{2})} \int_{-1}^1 \sigma_{N-2j}^{\delta(i\beta+\frac{\varepsilon}{2})}(r, s; u) Q_m^{\lambda(i\beta)}(u) du \right)$$

On the other hand, from the right hand side of (6.4) we get

$$C_d^{-1} (1-w^2)^{-\frac{\varepsilon}{2}} (1-w)^{-\delta(i\beta+\frac{\varepsilon}{2})-1} (1+w)^{\frac{\varepsilon}{2}} \sum_{N=0}^{\infty} (rs)^m \psi_N^{\lambda(i\beta)+m}(r) \psi_N^{\lambda(i\beta)+m}(s) w^{2N+m} \\ = C_d^{-1} (1-w)^{-\delta(i\beta+\varepsilon)-1} \sum_{N=0}^{\infty} (rs)^m \psi_N^{\lambda(i\beta)+m}(r) \psi_N^{\lambda(i\beta)+m}(s) w^{2N+m} \\ = C_d^{-1} \sum_{N=m}^{\infty} \left(\sum_{j=0}^{[(N-m)/2]} A_{N-m-2j}^{\delta(i\beta+\varepsilon)} (rs)^m \psi_j^{\lambda(i\beta)+m}(r) \psi_j^{\lambda(i\beta)+m}(s) \right) w^N \\ = C_d^{-1} \sum_{N=m}^{\infty} w^N A_N^{\delta(i\beta+\varepsilon)} K_{N,m}^{\delta(i\beta+\varepsilon), i\beta}(r, s),$$

where the last equality is true in view of Proposition 4.4. Altogether, we have

$$\sum_{N=0}^{\infty} w^N \left(\sum_{j=0}^{[N/2]} A_j^{\frac{\varepsilon}{2}+i\beta-1} A_{N-2j}^{\delta(i\beta+\frac{\varepsilon}{2})} \int_{-1}^1 \sigma_{N-2j}^{\delta(i\beta+\frac{\varepsilon}{2})}(r, s; u) Q_m^{\lambda(i\beta)}(u) du \right) \\ = \frac{(1+w)^{-\frac{\varepsilon}{2}}}{C_d} \sum_{N=m}^{\infty} w^N A_N^{\delta(i\beta+\varepsilon)} K_{N,m}^{\delta(i\beta+\varepsilon), i\beta}(r, s).$$

Multiplying both sides above by $r^{\frac{2\delta(i\beta)}{p}} s^{\frac{2\delta(i\beta)}{p'}} (1+w)^{\varepsilon/2}$ and using $(1+w)^{\varepsilon/2} = \sum_{j=0}^{\infty} \binom{\varepsilon/2}{j} w^j$, we see that

$$\sum_{N=0}^{\infty} w^N \left(\sum_{k=0}^N \binom{\varepsilon/2}{N-k} \sum_{j=0}^{[k/2]} A_j^{\frac{\varepsilon}{2}+i\beta-1} A_{k-2j}^{\delta(i\beta+\frac{\varepsilon}{2})} r^{\frac{2\delta(i\beta)}{p}} s^{\frac{2\delta(i\beta)}{p'}} \int_{-1}^1 \sigma_{k-2j}^{\delta(i\beta+\frac{\varepsilon}{2})}(r, s; u) Q_m^{\lambda(i\beta)}(u) du \right) \\ = \frac{1}{C_d} \sum_{N=m}^{\infty} w^N A_N^{\delta(i\beta+\varepsilon)} \mathcal{K}_{N,m}^{\varepsilon, i\beta}(r, s).$$

Equating the coefficients of w^N on both sides, we get (4.14).

Similarly, by using the Mehler's formula for $\Phi_N^{(d+1)}(x, y)$, $x, y \in \mathbb{R}^{d+1}$, we have

$$\begin{aligned} \sum_{N=0}^{\infty} w^N \int_{-1}^1 \Phi_N^{(d+1)}(r, s; u) Q_m^{\lambda(1/2+i\beta)}(u) du \\ = \frac{\Gamma(d/2)}{4\pi^{d/2}} (1-w^2)^{i\beta} \sum_{N=0}^{\infty} (rs)^m \psi_N^{\lambda(1/2+i\beta)+m}(r) \psi_N^{\lambda(1/2+i\beta)+m}(s) w^{2N+m}. \end{aligned}$$

Proceeding as above with this identity we obtain (4.15).

This completes the proof of the proposition. \square

7. TECHNICAL RESULTS

7.1. A lemma concerning binomial coefficients with complex parameters. The estimate contained in the following lemma is used in the proof of Theorem 3.3. Moreover, a slightly different version of such lemma (whose explicit statement and proof are omitted) is also used in the proof of Proposition 5.2.

Lemma 7.1. *For $\zeta \in \mathbb{C}$, let $\delta(\zeta)$ be defined as in (3.2). For any $\varepsilon > 0$ and $\beta \in \mathbb{R}$ we have the estimate*

$$\frac{1}{|A_N^{\delta(i\beta+\varepsilon)}|} \sum_{j=0}^N |A_{N-j}^{i\beta+\frac{\varepsilon}{2}-1}| |A_j^{\delta(\varepsilon/2)}| \leq C_\varepsilon (1+|\beta|) \cosh(\pi\beta)$$

where C_ε is independent of N .

Proof. In order to prove the lemma we make use of the fact that

$$(7.1) \quad |\Gamma(\alpha + i\beta)| \leq \Gamma(\alpha),$$

for $\alpha, \beta \in \mathbb{R}$, $\alpha > 0$, which follows from the very definition of Gamma function. We also have the Beta function

$$\frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)} = \int_0^1 (1-t)^{x-1} t^{y-1} dt$$

which leads to the estimate

$$(7.2) \quad \left| \frac{\Gamma(\alpha)\Gamma(i\beta+1)}{\Gamma(\alpha+i\beta+1)} \right| \leq \frac{\Gamma(\alpha)}{\Gamma(\alpha+1)}$$

for any $\alpha > 0$. We rewrite

$$\sum_{j=0}^N A_{N-j}^{i\beta+\varepsilon/2-1} A_j^{\delta(\varepsilon/2)} = \frac{\Gamma(\varepsilon/2)}{\Gamma(\varepsilon/2+i\beta)} \sum_{j=0}^N \frac{\Gamma(N-j+\varepsilon/2+i\beta)}{\Gamma(N-j+1)\Gamma(\varepsilon/2)} A_j^{\delta(\varepsilon/2)}.$$

By (7.1) and (3.6) we have

$$\begin{aligned} \left| \sum_{j=0}^N A_{N-j}^{i\beta+\varepsilon/2-1} A_j^{\delta(\varepsilon/2)} \right| &\leq \frac{\Gamma(\varepsilon/2)}{|\Gamma(\varepsilon/2+i\beta)|} \sum_{j=0}^N \frac{\Gamma(N-j+\varepsilon/2)}{\Gamma(N-j+1)\Gamma(\varepsilon/2)} A_j^{\delta(\varepsilon/2)} \\ &= \frac{\Gamma(\varepsilon/2) A_N^{\delta(\varepsilon)}}{|\Gamma(\varepsilon/2+i\beta)|}. \end{aligned}$$

Thus, we are left with estimating

$$\left| \frac{\Gamma(\varepsilon/2)}{\Gamma(\varepsilon/2+i\beta)} \frac{A_N^{\delta(\varepsilon)}}{A_N^{\delta(\varepsilon+i\beta)}} \right| = \left| \frac{\Gamma(\varepsilon/2)}{\Gamma(\varepsilon/2+i\beta)} \frac{\Gamma(\delta(\varepsilon)+i\beta+1)}{\Gamma(\delta(\varepsilon)+1)} \frac{\Gamma(N+\delta(\varepsilon)+1)}{\Gamma(N+\delta(\varepsilon)+i\beta+1)} \right|.$$

The middle term in the right hand side is clearly bounded by (7.1). The first term can be written as

$$\frac{(\varepsilon/2 + i\beta)\Gamma(\varepsilon/2)\Gamma(i\beta + 1)}{\Gamma(i\beta + 1)\Gamma(i\beta + \varepsilon/2 + 1)}$$

which leads, by (7.2), to the estimate

$$\left| \frac{\Gamma(\varepsilon/2)}{\Gamma(\varepsilon/2 + i\beta)} \right| \leq \frac{\Gamma(\varepsilon/2)}{\Gamma(\varepsilon/2 + 1)} \frac{(1 + |\beta|)}{|\Gamma(i\beta + 1)|}.$$

Similarly, the third term gives the estimate $|\Gamma(i\beta + 1)|^{-1}$. Therefore, the expression in Lemma 7.1 is bounded by

$$\frac{\Gamma(\varepsilon/2)}{\Gamma(\varepsilon/2 + 1)} (1 + |\beta|) |\Gamma(i\beta + 1)|^{-2}.$$

Once again the expression with the Beta function

$$\frac{\Gamma(\frac{1}{2})\Gamma(\frac{1}{2} + i\beta)}{\Gamma(1 + i\beta)} = \int_0^1 (1-t)^{-\frac{1}{2}} t^{i\beta - \frac{1}{2}} dt$$

leads to the estimate $|\Gamma(i\beta + 1)|^{-1} \leq \sqrt{\pi} |\Gamma(\frac{1}{2} + i\beta)|^{-1}$. Therefore,

$$|\Gamma(i\beta + 1)|^{-2} \leq \pi |\Gamma(\frac{1}{2} + i\beta)|^{-2} = \pi (\Gamma(\frac{1}{2} + i\beta)\Gamma(\frac{1}{2} - i\beta))^{-1}.$$

In view of the identity $\Gamma(z)\Gamma(1-z) = \frac{\pi}{\sin \pi z}$ we obtain

$$|\Gamma(i\beta + 1)|^{-2} \leq \sin \pi (\frac{1}{2} + i\beta) = \cosh \pi \beta.$$

This completes the proof of the lemma. \square

7.2. Some results on ultraspherical polynomials and Bessel functions. We will prove some technical results involving ultraspherical polynomials and modified Bessel functions with complex parameters. Let us recall several facts concerning both special functions.

Rodrigues' formula for the normalized ultraspherical polynomials takes the form

$$(7.3) \quad (1-u^2)^{\lambda-1/2} P_m^\lambda(u) = \frac{(-1)^m}{2^m (\lambda + 1/2)_m} \frac{d^m}{du^m} ((1-u^2)^{m+\lambda-1/2}), \quad m = 0, 1, \dots,$$

where $(\lambda)_m = \frac{\Gamma(\lambda+m)}{\Gamma(\lambda)}$ is the Pochhammer symbol. The normalized ultraspherical polynomials are also given by the explicit formula $P_m^\lambda(\cos \theta) = \frac{C_m^\lambda(\cos \theta)}{C_m^\lambda(1)}$ where (see [2, page 302])

$$(7.4) \quad C_m^\lambda(\cos \theta) = \sum_{k=0}^m \frac{(\lambda)_k (\lambda)_{m-k}}{k! (m-k)!} \cos(m-2k)\theta$$

for $\theta \in [0, \pi]$. It is immediate from (7.4) that

$$(7.5) \quad |P_m^\lambda(u)| \leq 1.$$

Since the functions $\lambda \mapsto (\lambda)_k$ are holomorphic for each fixed $u \in [-1, 1]$, the ultraspherical polynomials $P_m^\lambda(u)$ can be extended as a holomorphic function of λ on the domain $\{\lambda \in \mathbb{C} : \operatorname{Re}(\lambda) > -\frac{1}{2}\}$.

Let I_λ be the modified Bessel function of the first kind and order λ , defined as

$$I_\lambda(z) = \sum_{m=0}^{\infty} \frac{1}{m! \Gamma(m + \lambda + 1)} \left(\frac{z}{2}\right)^{2m+\lambda}.$$

Note that the functions I_λ can also be defined for complex values of λ . We will use Schläfli's integral representation of Poisson type for modified Bessel function, see [15, (5.10.22)]

$$(7.6) \quad I_\lambda(z) = \frac{(z/2)^\lambda}{\sqrt{\pi} \Gamma(\lambda + 1/2)} \int_{-1}^1 e^{zu} (1-u^2)^{\lambda-1/2} du, \quad |\arg z| < \pi, \quad \lambda > -\frac{1}{2}.$$

The following lemma is useful to express the Dunkl–Cesàro mean kernel in terms of the Cesàro kernel for the Laguerre expansions in Subsection 4.3.

Lemma 7.2. *Let $z \in \mathbb{C}$ and $\lambda > -\frac{1}{2}$. Then the following holds*

$$(7.7) \quad \int_{-1}^1 e^{zu} P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du = \sqrt{\pi} \Gamma(\lambda+1/2) (z/2)^{-\lambda} I_{\lambda+m}(z), \quad m = 0, 1, \dots$$

Proof. In view of Rodrigues' formula (7.3), we get

$$\int_{-1}^1 e^{zu} P_m^\lambda(u) (1-u^2)^{\lambda-1/2} du = \frac{(-1)^m}{2^m (\lambda+1/2)_m} \int_{-1}^1 e^{zu} \frac{d^m}{du^m} ((1-u^2)^{m+\lambda-1/2}) du$$

Integrating by parts and using (7.6) we see that

$$\begin{aligned} \frac{(-1)^m}{2^m (\lambda+1/2)_m} \int_{-1}^1 e^{zu} \frac{d^m}{du^m} ((1-u^2)^{m+\lambda-1/2}) du &= \frac{(z/2)^m}{(\lambda+1/2)_m} \int_{-1}^1 e^{zu} (1-u^2)^{m+\lambda-1/2} du \\ &= \sqrt{\pi} \Gamma(\lambda+1/2) (z/2)^{-\lambda} I_{\lambda+m}(z) \end{aligned}$$

since $\Gamma(\lambda+1/2) = \frac{\Gamma(\lambda+m+1/2)}{(\lambda+1/2)_m}$. This proves the lemma. \square

In order to estimate the Cesàro kernels for Dunkl–Hermite expansions, we needed to consider a variant of Lemma 7.2 where the parameter λ has to be taken complex. In the following lemma we express $I_{\lambda+m}$ in terms of certain variants of ultraspherical polynomials defined for complex λ . Furthermore, we also obtain some good estimates for these variants.

For $u \in [-1, 1]$, define

$$Q_m^{\lambda, \varepsilon}(u) := \frac{1}{\Gamma(i\beta + \varepsilon)} \int_0^1 \chi_{[|u|, 1]}(s) P_m^\alpha\left(\frac{u}{s}\right) (s^2 - u^2)^{\alpha-1/2} (1-s)^{i\beta+\varepsilon-1} s^{m+1} ds$$

for $m = 0, 1, 2, \dots$

Lemma 7.3. *Let $\varepsilon > 0$ and $\lambda = \alpha + i\beta$ with $\alpha > \frac{1}{2}$ and $\beta \in \mathbb{R}$. Then, for $z \in \mathbb{C}$, we have the identity*

$$(7.8) \quad \int_{-1}^1 e^{zu} Q_m^{\lambda, \varepsilon}(u) du = \frac{\sqrt{\pi} \Gamma(\alpha+1/2)}{2} \cdot \frac{I_{\lambda+m+\varepsilon}(z)}{(z/2)^{\lambda+\varepsilon}}.$$

Moreover, we have the uniform estimates

$$|Q_m^{\lambda, \varepsilon}(u)| \leq C \frac{(1-u^2)^{\alpha-1/2}}{|\Gamma(i\beta+1)|},$$

for all $-1 \leq u \leq 1$. Consequently,

$$Q_m^\lambda(u) := \lim_{\varepsilon \rightarrow 0} Q_m^{\lambda, \varepsilon}(u)$$

exists, satisfies (7.3) and

$$(7.9) \quad \int_{-1}^1 e^{zu} Q_m^\lambda(u) du = \frac{\sqrt{\pi} \Gamma(\alpha+1/2)}{2} \cdot \frac{I_{\lambda+m}(z)}{(z/2)^\lambda}.$$

Proof. For $\operatorname{Re}(\nu) > 0$ and $\operatorname{Re}(\mu) > 0$, we have (see [2, Theorem 4.11.1])

$$\frac{J_{\nu+\mu}(z)}{z^{\nu+\mu}} = \frac{2^{1-\mu}}{\Gamma(\mu)} \int_0^1 \frac{J_\nu(sz)}{(sz)^\nu} (1-s^2)^{\mu-1} s^{2\nu+1} ds,$$

where J_ν is the Bessel function of order ν . Using the relation $\frac{J_\nu(iz)}{(iz/2)^\nu} = \frac{I_\nu(z)}{(z/2)^\nu}$ and taking $\nu = \alpha + m$ and $\mu = \varepsilon + i\beta$, i.e. $\nu + \mu = \lambda + m + \varepsilon$, the above can be written as

$$\begin{aligned} \frac{I_{\lambda+m+\varepsilon}(z)}{(z/2)^{\lambda+m+\varepsilon}} &= \frac{2}{\Gamma(i\beta + \varepsilon)} \int_0^1 \frac{I_{\alpha+m}(sz)}{(sz/2)^{\alpha+m}} (1-s^2)^{i\beta+\varepsilon-1} s^{2\alpha+2m+1} ds \\ &= \frac{2}{\Gamma(i\beta + \varepsilon)} \int_0^1 \frac{I_{\alpha+m}(sz)}{(sz/2)^\alpha} (1-s^2)^{i\beta+\varepsilon-1} s^{2\alpha+m+1} ds. \end{aligned}$$

Therefore, in view of this and (7.7), we get

$$\frac{I_{\lambda+m+\varepsilon}(z)}{(z/2)^{\lambda+\varepsilon}} = \frac{2}{\Gamma(i\beta + \varepsilon)\sqrt{\pi}\Gamma(\alpha + 1/2)} \int_0^1 \int_{-1}^1 e^{szu} P_m^\alpha(u) (1-u^2)^{\alpha-\frac{1}{2}} du (1-s^2)^{i\beta+\varepsilon-1} s^{2\alpha+m+1} ds.$$

By making a change of variables and then a change of the order of integration, the right hand side of the above can be written as

$$\frac{2}{\Gamma(i\beta + \varepsilon)\sqrt{\pi}\Gamma(\alpha + 1/2)} \int_{-1}^1 e^{zu} \int_0^1 \chi_{[|u|,1]}(s) P_m^\alpha\left(\frac{u}{s}\right) (s^2 - u^2)^{\alpha-\frac{1}{2}} (1-s^2)^{i\beta+\varepsilon-1} s^{m+1} ds du.$$

Thus we have the desired identity

$$\frac{I_{\lambda+m+\varepsilon}(z)}{(z/2)^{\lambda+\varepsilon}} = \frac{2}{\sqrt{\pi}\Gamma(\alpha + 1/2)} \int_{-1}^1 e^{zu} Q_m^{\lambda,\varepsilon}(u) du.$$

We now proceed to show that $\lim_{\varepsilon \rightarrow 0} Q_m^{\lambda,\varepsilon}(u)$ exists and satisfies the required estimate. In order to do so, let us write

$$f_m(s) = P_m^\alpha(s)(1-s^2)^{\alpha-\frac{1}{2}}, \quad g_m(s) = s^{2\alpha+m-1}$$

so that we can write

$$Q_m^{\lambda,\varepsilon}(u) = \frac{i\beta + \varepsilon}{\Gamma(i\beta + \varepsilon + 1)} \int_{|u|}^1 f_m\left(\frac{u}{s}\right) (1-s^2)^{i\beta+\varepsilon-1} g_m(s) s ds.$$

Integrating by parts and noting that the boundary terms vanish (since $\alpha > \frac{1}{2}$) we have

$$Q_m^{\lambda,\varepsilon}(u) = \frac{1}{2\Gamma(i\beta + \varepsilon + 1)} \int_{|u|}^1 \frac{d}{ds} \left(f_m\left(\frac{u}{s}\right) g_m(s) \right) (1-s^2)^{i\beta+\varepsilon} ds.$$

Observe that $|Q_m^{\lambda,\varepsilon}(u)| \leq \frac{1}{2|\Gamma(i\beta+1)|} \int_{|u|}^1 \left| \frac{d}{ds} \left(f_m\left(\frac{u}{s}\right) g_m(s) \right) \right| (1-s^2)^{i\beta} ds$. Besides, it is easy to check that we can now pass to the limit as $\varepsilon \rightarrow 0$ and define

$$Q_m^\lambda(u) := \lim_{\varepsilon \rightarrow 0} Q_m^{\lambda,\varepsilon}(u) = \frac{1}{2\Gamma(i\beta + 1)} \int_{|u|}^1 \frac{d}{ds} \left(f_m\left(\frac{u}{s}\right) g_m(s) \right) (1-s^2)^{i\beta} ds.$$

On one hand, the boundedness (7.5) of $P_m^\alpha(u)$ leads to the estimate

$$\int_{|u|}^1 \left| f_m\left(\frac{u}{s}\right) \right| g'_m(s) ds \leq C(g_m(1) - g_m(|u|)) \leq 2C.$$

On the other hand, from Rodrigues' formula (7.3) it is easy to see that $|f'_m(u)| \leq C(1-u^2)^{\alpha-\frac{3}{2}}$ and therefore,

$$\int_{|u|}^1 \left| f'_m\left(\frac{u}{s}\right) \right| \frac{|u|}{s^2} g_m(s) ds \leq C \int_{|u|}^1 (s^2 - u^2)^{\alpha-\frac{3}{2}} \frac{|u|}{s} s^{m+1} ds$$

which gives the estimate (as $|u| \leq s \leq 1$)

$$|Q_m^{\lambda,\varepsilon}(u)| \leq \frac{C}{|\Gamma(i\beta + 1)|} (1-u^2)^{\alpha-\frac{1}{2}}.$$

As $Q_m^\lambda(u)$ is defined as the limit of $Q_m^{\lambda,\varepsilon}(u)$ it follows that $Q_m^\lambda(u)$ also satisfies the same estimate. Passing to the limit in (7.8) we see that

$$\int_{-1}^1 e^{zu} Q_m^\lambda(u) du = \frac{\sqrt{\pi}\Gamma(\alpha + 1/2)}{2} \cdot \frac{I_{\lambda+m}(z)}{(z/2)^\lambda}.$$

This completes the proof. \square

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