Accepted Manuscript

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PII:	S0377-0427(19)30120-7
DOI:	https://doi.org/10.1016/j.cam.2019.03.001
Reference:	CAM 12172
To appear in:	Journal of Computational and Applied Mathematics
Received date :	11 September 2018
Revised date :	31 January 2019



Please cite this article as: S.A. Dar and R.B. Paris, A (p, v)-extension of the Appell function $F_1(\cdot)$ and its properties, *Journal of Computational and Applied Mathematics* (2019), https://doi.org/10.1016/j.cam.2019.03.001

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A (p, ν) -extension of the Appell function $F(\cdot)$ and its properties

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Abstract

In this paper, we obtain a (p, v)-extension of the Appell hypergeometric function $F_1(\cdot)$, together with its integral representation. By using the extended Beta function $B_{p,v}(x, y)$ introduced in [9]. Also, we get some of its main properties, namely the Mellin transform, a differential formula, records of formulas and a bounded inequality. In addition, some new integral represe defines of the extended Appell function $F_{1,p,v}(\cdot)$ involving Meijer's *G*-function are obtaine.

MSC: 33C60, 33C65, 33B15, 3? .45, `3C10

Keywords: Appell's hypergeometric functions; Beta and Gamma functions; Eulerian integrals; Bessel function; *M* sijer's *G*-anction

1. Introduction and Preliminaries

In the present paper we enology the following notations:

 $\mathbb{N} = \{ 2, 2, ... \}, \ \mathbb{N}_0 := \mathbb{N} \cup \{ 0 \}, \ \mathbb{Z}_0^- := \mathbb{Z}^- \cup \{ 0 \},$

where the symbols \mathbb{N} and \mathbb{Z} denote the set of integer and natural numbers. Also, for the sake of concisen ss, we use the following notation [2, p.33]

$$\mathbb{C}_{>} := \{ p \in \mathbb{C} : \Re(p) > 0 \}.$$

where, \square summarial, the symbol $\mathbb C$ denotes the set of complex numbers

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A (p, ν) extension of the Appell function

In the available literature on hypergeometric series, the hypergeometric s ries and its generalizations appear in various branches of mathematics associated with polications. This type of series appears very naturally in quantum field theory, in particulal in the computation of analytic expressions for Feynman integrals. Such integral can be obtained and computed in different ways which may lead to identities for Appel' s ries; see [11]. On the other hand, the application of known relations for Appell pries pay lead to simplifications, help to solve problems or lead to greater insight in general field theory.

The Appell hypergeometric series are a natural two-variable extension of hypergeometric series, which are treated in detail in Erdélyi *et al.* [5]. In this paper, we highlight some of the most important properties and relations satisfied by these series. In the following, we follow to a great extent the expositions from the classical texts of F ailey [1], and Slater [14] (both contain a great amount of material on hypergeometric contexts). There are four types of Appell functions denoted by F_1, F_2, F_3, F_4 ; in the present study we shall only be concerned with the first Appell function F_1 given by [6, (10.13.1)]

$$F_{1}(b_{1}, b_{2}, b_{3}; c_{1}; x, y) = \sum_{m,n=0}^{\infty} \frac{(b_{1})_{m+n} (b_{2})_{m} (b_{3})_{n}}{(c_{1})_{m+1}} \frac{x}{r} \frac{y^{n}}{n!}$$

$$= \sum_{n,m=0}^{\infty} \frac{(b_{2})_{m} (b_{3})}{b_{1}} \frac{p'(r)_{n+1}}{b_{1}} \frac{m+n, c_{1}-b_{1}}{m!} \frac{x^{m}}{m!} \frac{y^{n}}{n!}, \quad (1.1)$$

where |x| < 1, |y| < 1. Here $(\lambda)_{\upsilon}$ $(\lambda, \upsilon \in \mathbb{C})$ C inters the Pochhammer symbol (or the shifted factorial, since $(1)_n = n!$) defined by

$$(\lambda)_{\upsilon} := \frac{\Gamma(\lambda + \upsilon)}{\Gamma(\lambda)} = \begin{cases} 1, & (\upsilon = 0, \ \lambda \in \mathbb{C} \setminus \{0\}) \\ \lambda(\lambda + 1) \dots (\lambda + n - 1), & (\upsilon = n \in \mathbb{N}, \ \lambda \in \mathbb{C}) \end{cases}$$

and $B(\alpha, \beta)$ denotes the classical B is included by [6, (5.12.1)]

$$B(\alpha,\beta) = \begin{cases} \int_{t}^{1} i^{\alpha-1} (1-t)^{\beta-1} dt, & (\Re(\alpha) > 0, \Re(\beta) > 0) \\ \frac{\Gamma(\alpha) \Gamma(t)}{\Gamma(\alpha+\beta)}, & ((\alpha,\beta) \notin \mathbb{Z}_{0}^{-}). \end{cases}$$
(1.2)

An integral representation of $F_1(\cdot)$ is given by [6, (16.15.1)]

$$F_1(b_1, b_2, b_3; c_1; x, t) = \frac{\Gamma(c_1)}{\Gamma(b_1)\Gamma(c_1 - b_1)} \int_0^1 t^{b_1 - 1} (1 - t)^{c_1 - b_1 - 1} (1 - xt)^{-b_2} (1 - yt)^{-b_3} dt,$$
(1.3)

where $\Re(c_1) > \Re(c_1) > |\arg(1-x)| < \pi$ and $|\arg(1-y)| < \pi$.

A natural extension of the Gauss hypergeometric series $_2F_1$ is the generalized hypergeometric series F_q [1., p.42, Eq.(1)] (see also [5]) with p numerator parameters $\alpha_1, ..., \alpha_p$ and q denominator parameters $\beta_1, ..., \beta_q$ defined by

$${}_{p}F_{q}\left(\begin{array}{cc}\alpha_{1},...,\alpha_{p};\\\beta_{1},...,\beta_{q};\end{array}\right) = \sum_{n=0}^{\infty} \frac{(\alpha_{1})_{n}...(\alpha_{p})_{n}}{(\beta_{1})_{n}...(\beta_{q})_{n}} \frac{z^{n}}{n!},\qquad(1.4)$$

where $\alpha_j \in \mathbb{C}$ $(1 \leq j \leq p)$ and $\beta_j \in \mathbb{C} \setminus \mathbb{Z}_0^ (1 \leq j \leq q)$ and $p, q \in \mathbb{N}_0$. The series ${}_pF_q(\cdot)$ is convergent for $|z| < \infty$ if $p \leq q$, and for |z| < 1 if p = q + 1. Furthermore, if the set

$$\omega = \sum_{j=1}^{q} \beta_j - \sum_{j=1}^{p} \alpha_j,$$

it is known that the ${}_{p}F_{q}$ series, with p = q + 1, is absolutely convergent on the unit circle |z| = 1 if $\Re(\omega) > 0$, and conditionally convergent on the unit circle $|z| = 1, z \neq 1$, if $-1 < \Re(\omega) \le 0$.

The modified Bessel function of the second kind $K_{\nu}(z)$ of rare ν (also known as the Macdonald function) is defined by (see [6, p. 251], [10])

$$K_{\nu}(z) = \sqrt{\pi} (2z)^{\nu} e^{-z} U(\nu + \frac{1}{2}, 2\nu + 1, 2z)$$
(1.5)

where U(a, b, z) is the confluent hypergeometric function [6, p. ??']. The Meijer *G*-function is defined by means of the Mellin-Barnes contour integral [6, (16.17.1)]

$$G_{p,q}^{m,n}\left(z \mid \alpha_{1},...,\alpha_{n};\alpha_{n+1},...,\alpha_{p}\right)$$

$$= \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} \frac{\prod_{j=1}^{m} \Gamma(\gamma_{j},-\zeta_{j}) \prod_{j=1}^{n} \Gamma(1-\alpha_{j}+\zeta)}{\prod_{j=n,-1}^{q} \Gamma(1-\beta_{j}+\zeta) \prod_{j=n+1}^{p} \Gamma(\alpha_{j}-\zeta)} z^{\zeta} d\zeta, \qquad (1.6)$$

where $z \neq 0$, and m, n, p, q are non-negative integers such that $1 \leq m \leq q$, $0 \leq n \leq p$ and $p \leq q$. The integral (1.6) converges in the sector $|\arg z| < \pi \kappa$, where $\kappa = m + n - \frac{1}{2}(p+q)$ and it is supposed that $\kappa > 0$.

The G-function is important in a_{μ_1} is d mathematics and formulas developed for the G-function become master or key formulas from which a very large number of relations can be deduced for Bessel function s, their combinations and many other related functions. Thus the following list of some particular cases of Meijer's G-function associated with the Bessel function $K_{\nu}(z)$ has been betained mainly from several papers by C. S. Meijer (see also [5, pp. 219-220, Eq. (4) 50)], [12, p.48, Eq.(12)])

$$\kappa_{\nu}(z) = \sqrt{\pi} \ e^{z} \ G_{1,2}^{2,0} \left(2z \left| \begin{array}{c} \frac{1}{2} \\ \nu \end{array} \right|, -\nu \right), \tag{1.7}$$

$$= \frac{\cos(\pi\nu)}{\sqrt{\pi}} e^{-z} G_{1,2}^{2,1} \left(2z \middle| \begin{array}{c} \frac{1}{2} \\ \nu \\ \end{array}, -\nu \right), \qquad (1.8)$$

$$= z^{-\mu} 2^{\mu-1} G_{0,2}^{2,0} \left(\frac{z^2}{4} \bigg|_{\frac{\mu+\nu}{2}}, \frac{\mu-\nu}{2} \right), \qquad (1.9)$$

$$= \cos(\pi\nu) \frac{(2z)^{-\mu} e^z}{\sqrt{\pi}} G_{1,2}^{2,1} \left(2z \middle| \begin{array}{c} \mu + \frac{1}{2} \\ \mu + \nu, \ \mu - \nu \end{array} \right), \tag{1.10}$$

$$= \frac{z^{-\mu}4^{\mu-1}}{\pi} G_{0,4}^{4,0} \left(\frac{z^4}{256} \right|_{\frac{\mu+\nu}{4}}, \frac{2+\mu+\nu}{4}, \frac{\mu-\nu}{4}, \frac{2+\mu-\nu}{4} \right), \tag{1.11}$$

where μ is a free parameter and in all these expressions we have $z \neq 0$.

In 1997, Chaudhry *et al.* [3, Eq.(1.7)] gave a *p*-extension of the Bet_{*i*} function B(x, y) in the form

$$B(x,y;p) = \int_0^1 t^{x-1} (1-t)^{y-1} \exp\left[\frac{-p}{t(1-t)}\right] dt \qquad (\mathcal{Y}(p) > 0),$$

and they proved that this extension has connections with the M \ldots lonal, error and Whittaker functions. Also, Chaudhry *et al.* [4] extended the Gau sian hypergeometric series ${}_{2}F_{1}(\cdot)$ and its integral representations. Recently, Parmar *et al.* [9]. Jave given a further extension of the extended Beta function B(x, y; p) by addir g one more parameter ν , which we denote and define by

$$B_{p,\nu}(x,y) = \sqrt{\frac{2p}{\pi}} \int_0^1 t^{x-\frac{3}{2}} (1-t)^{y-\frac{3}{2}} \sum_{t+\frac{1}{2}} \left(\frac{p}{1-t} \right) dt, \qquad (1.12)$$

where $\Re(p) > 0, \nu \ge 0$ and $K_{\nu+\frac{1}{2}}(\cdot)$ is the modified essel function of order $\nu + \frac{1}{2}$. When $\nu = 0$, (1.12) reduces to B(x, y; p), since $K_{\frac{1}{2}}(z) = (\pi/(2z))^{1/2}e^{-z}$. A different generalization of the Beta function has been g. en m_1 8].

Motivated by some of the above-mentioned elements of special functions, many authors have studied integral representations $e^+ e^+ F_1(\cdot)$ function. Our aim in this paper is to introduce a (p, v)-extension of the Appell bypergeometric function in (1.1) based on the extended Beta function in (1.12), which we denote by $F_{1,p,\nu}(\cdot)$, and to systematically investigate some properties of this extended function. We consider the Mellin transform, a differential formula, recursion form (as all a bounded inequality satisfied by this function. Also, we obtain some integral representations for $F_{1,p,\nu}$ containing Meijer's *G*-function.

The plan of this paper as ollows. The extended Appell function $F_{1,p,\nu}(\cdot)$ and its integral representation are deduced in Section 2. Some new integral representations for $F_{1,p,\nu}(\cdot)$ involving the Meijer $G_{-1,\nu}$ et on are given. The main properties $F_{1,p,\nu}(\cdot)$, namely its Mellin transform, a differential formula, recurrence relation and a bounded inequality are established in Sections 3–6. The concluding remarks are made in Section 7.

2 The $_{(\nu)}$, ν)-extended Appell function $F_{1,p,\nu}(\cdot)$

In [7], Özarslan *et al.* Sive an extension of Appell's hypergeometric function $F_1(\cdot)$ together with its integrate representation. Here we consider the following (p, v)-extension of the Appell hypergeometric function, which we denote by $F_{1,p,\nu}(\cdot)$, based on the extended beta function $B_{T,\nu}(x, y)$ defined in (1.12). This is given by

$$F_{1,n,\nu}(b_1, \dot{z}_2, a_3; c_1; x, y) = \sum_{n,m=0}^{\infty} \frac{(b_2)_m (b_3)_n \ B_{p,\nu}(b_1 + m + n, c_1 - b_1)}{B(b_1, c_1 - b_1)} \frac{x^m}{m!} \frac{y^n}{n!}, \qquad (2.1)$$

where |x| < 1, |y| < 1 and $b_1, b_2, b_3 \in \mathbb{C}$ and $c_1 \in \mathbb{C} \setminus \mathbb{Z}_0^-$. This definition clearly reduces to the original function when $\nu = 0$ and p = 0.

An integral representation for the function $F_{1,p,\nu}(\cdot)$ is given by

-

$$F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) = \frac{\Gamma(c_1)}{\Gamma(b_1)\Gamma(c_1 - b_1)}$$

$$\times \sqrt{\frac{2p}{\pi}} \int_0^1 t^{b_1 - \frac{3}{2}} (1 - t)^{c_1 - b_1 - \frac{3}{2}} (1 - xt)^{-b_2} (1 - yt)^{-b_3} K_{\nu + \frac{1}{2}} \left(\frac{p}{\sqrt{1 - t}}\right) dt, \qquad (2.2)$$

where $\Re(p) > 0$, $\nu \ge 0$, $|\arg(1-x)| < \pi$ and $|\arg(1-y)| < \pi$ as a we impose the condition $\Re(c_1) > \Re(b_1) > 0$ for the multiplicative factor $1/B(b_1, c_1 - b_1)$ to be finite. That this representation yields (2.1) can be shown by binomially expanding the factors $(1 - xt)^{-b_2}$ and $(1 - yt)^{-b_3}$ when |x|, |y| < 1, reversing the order of summation and integration and evaluating the resulting integral by (1.12).

Theorem 1. Each of the following integral representations of $F_{1,p,\nu}(\cdot)$ associated with Meijer's G-function holds for $p \in \mathbb{C}_{>}$.

$$F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) = \frac{\Gamma(c_1)\sqrt{2p}}{\Gamma(b_1)\Gamma(c_1 - b_1)} \int_0^1 t^{b - \frac{3}{2}} (1 - t)^{c_1 - b_1 - \frac{3}{2}} (1 - xt)^{-b_2} (1 - yt)^{-b_3} \\ \times e^{\frac{p}{t(1 - t)}} G_{1,2}^{2,0} \left(\frac{2p}{t(1 - t)} \Big|_{\nu - t_1}^2 , -\nu - \frac{1}{2} \right) dt$$

$$\Gamma(c_1)_{1/2p} = \cos \pi (\nu + \frac{1}{2}) \int_0^1 t^{-3} dt \qquad (2.3)$$

$$= \frac{\Gamma(c_1)\sqrt{2p}}{\Gamma(b_1)\Gamma(c_1 - b_1)} \frac{\cos \pi(\nu + \frac{1}{2})}{\pi} \int_0^1 \frac{1}{2} \frac{b_1 - \frac{3}{2}}{(1 - t)^{c_1 - b_1 - \frac{3}{2}}} (1 - xt)^{-b_2} (1 - yt)^{-b_3} \\ \times e^{\frac{-p}{t(1 - t)}} G_{1,2}^{2,1} \left(\frac{2p}{t(1 - t)}\right) \left| \begin{array}{c} \frac{1}{2} \\ \nu + \frac{1}{2} \\ \nu + \frac{1}{2} \\ -\nu - \frac{1}{2} \end{array} \right) dt$$
(2.4)

$$= \frac{\Gamma(c_1)2^{\mu-\frac{1}{2}}p^{-\mu+\frac{1}{2}}}{\Gamma(b_1)\Gamma(c_1-b_1)} \int_0^t t^{b_1+\mu-\frac{\gamma}{2}} (1-t)^{c_1-b_1+\mu-\frac{3}{2}} (1-xt)^{-b_2} (1-yt)^{-b_3} \times G_{0,2}^{2,\ell} \left(\frac{r}{4t^2(1-t)^2} \bigg| \frac{2\mu+2\nu+1}{4}, \frac{2\mu-2\nu-1}{4} \bigg) dt$$
(2.5)

$$= \frac{\Gamma(c_1)(2p)^{-\mu+\frac{1}{2}}}{\Gamma(b_1)\Gamma(c_1-b_1)} \left(\frac{2s_7}{\pi} \frac{(\nu+\frac{1}{2})}{\pi} \int_0^1 t^{b_1+\mu-\frac{3}{2}} (1-t)^{c_1-b_1+\mu-\frac{3}{2}} (1-xt)^{-b_2} (1-yt)^{-b_3} \right) \\ \times \left| \frac{t^{-p}}{t^{(1-s)}} G_{1,2}^{2,1} \left(\frac{2p}{t(1-t)} \right) \frac{\frac{2\mu+1}{2}}{\frac{2\mu+2\nu+1}{2}} , \frac{2\mu-2\nu-1}{2} \right) dt$$
(2.6)

$$= \frac{\Gamma(c_1)p^{-\mu-\frac{1}{2}}2^{2\mu-\frac{3}{2}}}{\Gamma(b_1,\Gamma(c_1-b_1)\pi^{\frac{3}{2}}} \int_0^1 t^{b_1+\mu-\frac{3}{2}}(1-t)^{c_1-b_1+\mu-\frac{3}{2}}(1-xt)^{-b_2}(1-yt)^{-b_3} \times G_{0,4}^{4,0}\left(\frac{p^4}{(4t)^4(1-t)^4}\right| \frac{2\mu+2\nu+1}{8}, \frac{2\mu+2\nu+5}{8}, \frac{2\mu-2\nu-1}{8}, \frac{2\mu-2\nu+3}{8}\right) dt, \qquad (2.7)$$

where $5, c_{1}, \gamma^{(0)}(b_1) > 0, |\arg(1-x)| < \pi, |\arg(1-y)| < \pi$ and μ is a free parameter.

Proof: The above integral representations (2.3)-(2.7) are obtained by using (1.7)-(1.11) in the expression of the extended Appell function in (2.2). Similarly, other , we integral representations of $F_{1,p,v}(\cdot)$ associated with the confluent hypergeometric function, can be obtained using (1.5) in (2.2).

The following transformation formula can be derived from the integ. 1^{-1} epresentation (2.2) for $F_{1,p,\nu}(\cdot)$.

Theorem 2. The following transformation formula holds:

$$F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) = (1-x)^{-b_2} (1-y)^{-b_3} F_{1,p,\nu}\left(c_1 - b_1, \frac{1}{2}, b_3; c_1, \frac{x}{x-1}, \frac{y}{y-1}\right).$$
(2.8)

Proof: Put $t = 1 - \zeta$ in (2.2) to obtain

$$F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) = \frac{\Gamma(c_1)(1 - x)^{-c_2}(1 - y)^{-b_3}}{\Gamma(b_1)\Gamma(c_1 - b_1)}$$

$$\times \sqrt{\frac{2p}{\pi}} \int_0^1 \zeta^{c_1 - b_1 - \frac{3}{2}} (1 - \zeta)^{b_1 - \frac{3}{2}} \left(1 - \frac{x}{x - 1}\zeta\right)^{-b_2} \left(1 - \frac{z}{y - 1}\zeta\right)^{-b_3} K_{\nu + \frac{1}{2}} \left(\frac{p}{\zeta(1 - \zeta)}\right) d\zeta.$$

Identification of the above integral as a $F_{1,p,\nu}(\cdot)$ function then yields (2.8).

3. The Mellin t. s. sfo. m of $F_{1,p,\nu}(\cdot)$

The Mellin transform of a locally integrate function f(x) on $(0,\infty)$ is defined by

$$\mathcal{M}\{f(x)\}(s) = \int_0^\infty x^{s-1} f(x) \, dx$$

when the integral converges.

Theorem 3. The following Meu. \uparrow tr nsform of the extended Appell hypergeometric function $F_{1,p,v}(\cdot)$ holds true:

$$\mathcal{M}\left\{F_{1,p,v}(b_{1}, b_{2}, b_{3}, c_{1}; x, y)\right\}(s) = \int_{0}^{\infty} p^{s-1} F_{1,p,v}(b_{1}, b_{2}, b_{3}; c_{1}; x, y) dp$$
$$= \frac{2^{s-1}}{\sqrt{2}} \Gamma\left(\frac{s-v}{2}\right) \Gamma\left(\frac{s+v+1}{2}\right) F_{1}\left(b_{1}+s, b_{2}, b_{3}; c_{1}+s; x, y\right),$$
(3.1)

where $\Re(s-v) > 0$, $\Re(z-v) > -1$, $\Re(s) > 0$ and $c_1 + s \in \mathbb{C}/\mathbb{Z}_0^-$.

Proof: Substituting the extended Appell function (2.1) into the left-hand side of (3.1) and changing the under of integration (by the uniform convergence of the integral), we obtain

$$\Lambda: \{ \overline{c}, \dots, (b_1, b_2, b_3; c_1; x, y) \} (s) = \frac{\Gamma(c_1)}{\Gamma(b_1)\Gamma(c_1 - b_1)} \sqrt{\frac{2}{\pi}} \int_0^1 t^{b_1 - \frac{3}{2}} (1 - t)^{c_1 - b_1 - \frac{3}{2}} dt^{b_1 - \frac{3}{2}} (1 - t)^{c_1 - b_1 - \frac{3}{2}} dt^{b_1 - \frac{3}{2}} (1 - t)^{c_1 - b_1 - \frac{3}{2}} dt^{b_1 - \frac{3}{2}} (1 - t)^{c_1 - b_1 - \frac{3}{2}} dt^{b_1 - \frac{3}{2}} dt^{b$$

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$$\times (1 - xt)^{-b_2} (1 - yt)^{-b_3} \left\{ \int_0^\infty p^{s - \frac{1}{2}} K_{v + \frac{1}{2}} \left(\frac{p}{t(1 - t)} \right) dp \right\} dt.$$
(3.2)

Application of the result [6, (10.43.19)]

$$\int_0^\infty w^{s-\frac{1}{2}} K_{\alpha+\frac{1}{2}}(w) dw = 2^{s-\frac{3}{2}} \Gamma\left(\frac{s-\alpha}{2}\right) \Gamma\left(\frac{s+\alpha+1}{2}\right) \qquad (|\Re(\alpha_j| < \Re(s))| + |\Re(\alpha_j| < \Re(s))| + |\Re(s)| + |$$

followed by the substitution $w = p/\{t(1-t)\}$ in (3.2) then yields

$$\varphi(s) \equiv \mathcal{M}\left\{F_{1,p,v}(b_1, b_2, b_3; c_1; x, y)\right\}(s) = \frac{2^{s-1}}{\sqrt{\pi}} \Gamma\left(\frac{s-v}{2}\right) \Gamma\left(\frac{s+v+1}{2}\right)$$
$$\times \frac{\Gamma(c_1)}{\Gamma(b_1)\Gamma(c_1-b_1)} \int_0^1 t^{b_1+s-1} (1-t)^{c_1+s-b_1-1} (1-xt)^{-b_3} (1-yt)^{-b_2} dt.$$
(3.3)

Finally, using the definition of the Appell function $F_1(\cdot) \stackrel{\cdot}{\to} (1.2)$ we obtain the right-hand side of (3.1).

Corollary: The following inverse Mellin formula to. $F_{1,p,v}(\cdot)$ holds:

$$F_{1,p,v}(b_1, b_2, b_3; c_1; x, y) = \mathcal{M}^{-1} \{\varphi(s)\}$$

$$= \frac{1}{4\pi i \sqrt{\pi}} \int_{-i\infty}^{+i\infty} \left(\frac{2}{p}\right)^s \Gamma\left(\frac{s-v}{2}\right) \Gamma\left(\frac{s+v-1}{2}\right) F_1\left(b_1+s, b_2, b_3; c_1+s; x, y\right) \, ds, \quad (3.4)$$

where $c > \nu$.

4. A differentiation formula for $F_{1,p,\nu}(\cdot)$

Theorem 4. The following different tration formula for $F_{1,p,v}(\cdot)$ holds:

$$\frac{\partial^{M+N}}{\partial x^M \partial y^N} F_{1,p,\nu}(b_1, b_2, b_3; c_1; z, y)$$

$$= \frac{(b_1)_{M+N}(b_2)_M(b_3)_I}{(c_1)_{M+N}} F_{1,p,\nu}(c_1 + M + N, b_2 + M, b_3 + N; c_1 + M + N; x, y), \quad (4.1)$$
where $M, N \in \mathbb{N}_0$.

Proof: If we differ ntize the series for $F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y)$ in (2.1) with respect to x, we obtain

$$\frac{\partial}{\partial x}F_{1,p,\nu}(b_1,k_2,b_3;c_1;x,y) = \sum_{m=1}^{\infty}\sum_{n=0}^{\infty}\frac{(b_2)_m(b_3)_n}{B(b_1,c_1-b_1)}\frac{B(b_1,b_2-b_1)}{B(b_1,c_1-b_1)}\frac{x^{m-1}}{(m-1)!}\frac{y^n}{n!}$$

Making us of the fact that

$$B(b_1, c_1 - b_1) = \frac{c_1}{b_1} B(b_1 + 1, c_1 - b_1)$$
(4.2)

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and $(b_2)_{m+1} = b_2(b_2 + 1)_m$, we have upon setting $m \to m + 1$

$$\frac{\partial}{\partial x}F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) = \frac{b_1 b_2}{c_1} \sum_{m,n=0}^{\infty} \frac{(b_2 + 1)_m (b_3)_n B_{p,\nu}(b_1 + 1 + m + n, c_1 - b_1)}{B(b_1 + 1, c_1 - b_1)} \frac{x^m}{m!} \frac{y^n}{n!} \\
= \frac{b_1 b_2}{c_1} F_{1,p,\nu}(b_1 + 1, b_2 + 1, b_3; c_1 + 1; x, y).$$
(4.3)

Repeated application of (4.3) then yields for M = 1, 2, ...

$$\frac{\partial^M}{\partial x^M}F_{1,p,\nu}(b_1,b_2,b_3;c_1;x,y) = \frac{(b_1)_M(b_2)_M}{(c_1)_M}F_{1,p,\nu}(b_1+M,t_2+M,t_3;c_1+M;x,y).$$

A similar reasoning shows that

$$\frac{\partial^{M+1}}{\partial x^{M} \partial y} F_{1,p,\nu}(b_{1}, b_{2}, b_{3}; c_{1}; x, y)$$

$$= \frac{(b_{1})_{M}(b_{2})_{M}}{(c_{1})_{M}} \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} \frac{(b_{2} + M)_{m}(b_{3})_{n} B_{p,\nu}(b_{1} + M + ... + n, c_{1} - b_{1})}{B(b_{1} + M - c_{1} - \cdot}) \frac{x^{m}}{m!} \frac{y^{n-1}}{(n-1)!}$$

$$= \frac{(b_{1})_{M+1}(b_{2})_{M}b_{3}}{(c_{1})_{M+1}} F_{1,p,\nu}(b_{1} + M + 1; c_{2} + M, b_{3} + 1; c_{1} + M + 1; x, y). \quad (4.4)$$

Repeated differentiation of (4.4) with respert to y then readily produces the result stated in (4.1). The result (4.1) has been derived assuming that |x| < 1, |y| < 1 but can be extended to all values of x and y satisfying $|x| < (1 - x)| < \pi$, $|\arg(1 - y)| < \pi$ by appeal to analytic continuation.

5. An $\mathbf{u}_{\mathbf{F},\mathbf{P}\mathbf{e}'}$ bound for $F_{1,p,\nu}(\cdot)$

Theorem 5. Let the parameters j_1 , j_2 , b_3 , c_1 and the variables x, y be real. Then the following bound for $F_{1,p,\nu}(\cdot)$ holds:

$$|F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y)| \leq \frac{2^{\nu} |p|^{\nu+1} \Gamma(\nu + \frac{1}{2})}{\sqrt{\pi}(\Re(p))^{\nu} \nu + 1} \frac{\beta(b_1 + \nu, c_1 - b_1 + \nu)}{B(b_1, c_1 - b_1)} F_1(b_1 + \nu, b_2, b_3; c_1 + 2\nu; x, y),$$
(5.1)

where $\Re(p) > 0$.

The integrat representation of the extension $F_{1,p,\nu}(\cdot)$ in (2.2) is associated with the modified Besser function of the second kind, for which we have the following expression [6, (10.32.8)]

$$\mathcal{K}_{1}(z) = \frac{\sqrt{\pi} \left(\frac{1}{2}z\right)^{\nu + \frac{1}{2}}}{\Gamma(\nu + 1)} \int_{1}^{\infty} e^{-zt} (t^{2} - 1)^{\nu} dt, \qquad (\nu > -1, \ \Re(z) > 0).$$

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In our problem we have $\nu > 0$, $\Re(z) > 0$. Further, we let $x = \Re(z)$, so that

$$\begin{aligned} |K_{\nu+\frac{1}{2}}(z)| &\leq \frac{\sqrt{\pi} \left(\frac{1}{2}|z|\right)^{\nu+\frac{1}{2}}}{\Gamma(\nu+1)} \left| \int_{1}^{\infty} e^{-zt} (t^{2}-1)^{\nu} dt \right| < \frac{\sqrt{\pi} \left(\frac{1}{2}|z|\right)^{\nu+\frac{1}{2}}}{\Gamma(\nu+1)} \int_{\gamma}^{\infty} t^{2\nu} e^{-x_{\nu}} dt \\ &= \frac{\sqrt{\pi} \left(\frac{1}{2}|z|\right)^{\nu+\frac{1}{2}}}{\Gamma(\nu+1)} \frac{\Gamma(2\nu+1,x)}{x^{2\nu+1}}, \end{aligned}$$
(5.2)

where $\Gamma(a, z)$ is the upper incomplete gamma function [6, (8.2.2)]. Alone, gh this bound is numerically found to be quite sharp when z is real, it involves the incomplete gamma function which would make the integral for $F_{1,p,\nu}(b_1, b_2, b_3; c_1; \cdot, y)$ difficult to bound. We can simplify (5.2) by making use of the simple inequality $\Gamma(2\nu + 1, x) < \Gamma(2\nu + 1)$ to find

$$|K_{\nu+\frac{1}{2}}(z)| < \frac{\sqrt{\pi} \left(\frac{1}{2}|z|\right)^{\nu+\frac{1}{2}}}{\Gamma(\nu+1)} \frac{\Gamma(2\nu+1)}{x^{2\nu+1}} = \frac{1}{2} \left(\frac{\Gamma_{\nu}|z|}{x} \right)^{\nu+\frac{1}{2}} \Gamma(\nu+\frac{1}{2}), \tag{5.3}$$

upon use of the duplication formula for the gamma funct. n. The bound (5.3) is less sharp than (5.2) but has the advantage of being evier to handle in the integral for $F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y).$

Proof: Setting z = p/(t(1-t)), where $t \in (0,1)$ a.¹ $\Re(p) > 0$ in (5.3), we obtain

$$\left| K_{\nu+\frac{1}{2}} \left(\frac{p}{t(1-t)} \right) \right| < \frac{1}{2} \left(\frac{2|\mu|^{t}(1-\iota)}{(\Re_{\lfloor \iota \rceil}))^{2}} \right)^{\nu+\frac{1}{2}} \Gamma(\nu+\frac{1}{2}).$$

For ease of presentation we shall assume that the parameters b_1 , b_2 , b_3 and c_1 are real; the extension to complex parameters is strai, but and. In addition, we shall consider only real values of the variables x and y. Then, from (2.2),

$$\begin{aligned} |F_{1,p,\nu}(b_{1}b_{2},b_{3};c_{1};x,y)| \\ &\leq \frac{\sqrt{2|p|/\pi}}{B(b_{1},c_{1}-b_{1})} \int_{0}^{1} \left| t^{b_{1}-\frac{3}{2}} (1-t)^{c_{1}-c_{-}-\frac{3}{2}} (1-xt)^{-b_{2}} (1-yt)^{-b_{3}} K_{\nu+\frac{1}{2}} \left(\frac{p}{t(1-t)}\right) \right| dt \\ &< \frac{2^{\nu}|p|^{\nu+1}}{\sqrt{\pi}(\Re(p))^{2\nu+1}} \frac{\Gamma(\nu+\frac{1}{2})}{B(b_{1},c_{1}-b_{1})} \int_{0}^{t} t^{b_{1}+\nu-1} (1-t)^{c_{1}-b_{1}+\nu-1} (1-xt)^{-b_{2}} (1-yt)^{-b_{3}} dt \\ &< \frac{2^{\nu}|p|^{\nu+1}\Gamma(\nu+\frac{1}{2})}{\sqrt{\pi}(\Re(p))^{2\nu-1}} \int_{0}^{t} \frac{B(b_{1}+\nu,c_{1}-b_{1}+\nu)}{B(b_{1},c_{1}-b_{1})} F_{1}(b_{1}+\nu,b_{2},b_{3};c_{1}+2\nu;x,y), \end{aligned}$$
(5.4)

which is the result s' ated n. (5.1).

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If we have x < 0, y < 0 (resp. x > 0, y > 0) and suppose further that $b_2 > 0$, $b_3 > 0$ (resp. $b_2 < 0, b_3 < 0$) her we obtain the simpler bound

$$|F_{1,\nu}(b_1, b^r, b_3; c_1; x, y)| < \frac{2^{\nu} |p|^{\nu+1} \Gamma(\nu + \frac{1}{2})}{\sqrt{\pi} (\Re(p))^{2\nu+1}} \frac{B(b_1 + \nu, c_1 - b_1 + \nu)}{B(b_1, c_1 - b_1)}$$

In Table 1 we predent some values of the bound (5.1) compared with those of $F_{1,p,\nu}(\cdot)$ for several values of the parameters p and ν .

Table 1: Values of $F_{1,p,\nu}(\cdot)$ and the bound (5.1) for different p and ν when $b_1 = 2/3$, $b_2 = 1$, $b_3 = 7/5$, $c_1 = 3$ and x = 1/4, y = 1/3. $c_1 = 3$ and x = 1/4, y = 1/3.

p	ν	$F_{1,p,\nu}(\cdot)$	Bound	p	ν	$F_{1,p,\nu}(\cdot)$	B Jund
0.05	0.50	1.27094	1.47965	0.25	0.50	0.27000	0.26172
0.05	1.00	3.17018	3.34772	0.25	1.00	0.39194	ີ 6695.
0.05	2.00	37.9432	38.4515	0.25	2.00	1.2064′	1.53006
0.10	0.50	0.75218	1.04627	0.50	0.50	0.07198	^۱ 46791
0.10	1.00	1.43982	1.67386	0.50	1.00	0.0^{\prime} 025	0.33477
0.10	2.00	9.15901	9.61288	0.50	2.00	0.1 `230	.38452

6. Recursion formulas for $F_{1,\nu}$, $\nu(\cdot)$

In view of the recursion formulas for the Appell function. $F_1(\cdot)$ (see [13] and [15]) we give the following recursion formulas for the extended Appe.' function $F_{1,p,\nu}(\cdot)$.

Theorem 6. The following recursion formulas for 'he extended Appell function with respect to the numerator parameters b_2 and b_3 ?

$$F_{1,p,\nu}(b_1, b_2 + n, b_3; c_1; x, y) = F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) + \frac{b_1}{c_1} \sum_{\ell=1}^n r_{1,p,\nu}(b_1 + 1, b_2 + \ell, b_3; c_1 + 1; x, y)$$
(6.1)

and

and $F_{1,p,\nu}(b_1, b_2, b_3 + n; c_1; x, y)$

$$=F_{1,p,\nu}(b_1,b_2,b_3;c_1;x,y) + \frac{{}'_1 y}{c_1} \sum_{\ell=1}^n F_{1,p,\nu}(b_1+1,b_2,b_3+\ell;c_1+1;x,y)$$
(6.2)

for positive integer n.

Proof: From (2.1) an the result $(b_2 + 1)_m = (b_2)_m (1 + m/b_2)$, we obtain

$$F_{1,p,\nu}(b_1, b_2 + 1 \ b_3; \ \gamma_1; x \ y) = \sum_{m,n=0}^{\infty} \frac{(b_2 + 1)_m (b_3)_n B_{p,\nu}(b_1 + m + n, c_1 - b_1)}{B(b_1, c_1 - b_1)} \frac{x^m}{m!} \frac{y^n}{n!}$$

$$=F_{1,p,\nu}(b_1,b_2,\cdots,c_1;z,y)+\frac{x}{b_2}\sum_{m=1}^{\infty}\sum_{n=0}^{\infty}\frac{(b_2)_m(b_3)_nB_{p,\nu}(b_1+m+n,c_1-b_1)}{B(b_1,c_1-b_1)}\frac{x^{m-1}}{(m-1)!}\frac{y^n}{n!}.$$

Setting $m \rightarrow m + 1$ and using $(b_2)_{m+1} = b_2(b_2 + 1)_m$ together with (4.2), we find

$$F_{1,p,
u}(\mathbf{1},\mathbf{0})=1,b_3;c_1;x,y)$$

$$=F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) + \frac{b_1 x}{c_1} \sum_{m,n=0}^{\infty} \frac{(b_2 + 1)_m (b_3)_n B_{p,\nu}(b_1 + 1 + m + n, c_1 - b_1)}{B(b_1 + 1, c_1 - b_1)} \frac{x^m}{m!} \frac{y^n}{m!} \frac{y^n}{n!}$$
$$=F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) + \frac{b_1 x}{c_1} F_{1,p,\nu}(b_1 + 1, b_2 + 1, b_3; c_1 + 1, x, y).$$
(6.3)

From (6.3) we obtain, upon putting $b_2 \rightarrow b_2 + 1$,

$$F_{1,p,\nu}(b_1, b_2 + 2, b_3; c_1; x, y)$$

$$= F_{1,p,\nu}(b_1, b_2 + 1, b_3; c_1; x, y) + \frac{b_1 x}{c_1} F_{1,p,\nu}(b_1 + 1, b_2 - 2, b_3; r_1 + 1; x, y)$$

= $F_{1,p,\nu}(b_1, b_2, b_3; c_1; x, y) + \frac{b_1 x}{c_1} \sum_{\ell=1}^2 F_{1,p,\nu}(b_1 + 1, b_2 + 2, r_3; c_1 + 1; x, y).$

Repeated application of the recursion (6.3) in this manner then immediately leads to the result stated in (6.1). The proof of (6.2) is obtained in the same way by interchanging b_2 and b_3 .

7. Concludi gromarks

We have introduced the extended Appell 1 pergeometric function $F_{1,p,\nu}(\cdot)$ given in (2.1) by use of the extended Beta function defined n. (1.12). Also, we have described some properties of this function, namely the Gaussian transform, a differential formula, some recurrence relations and a bounded inequality. In addition, we have also obtained some new integral representations of the extended Appell function involving Meijer's *G*-function and indicated other possible representations in terms of the confluent hypergeometric function.

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