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A Guided Wave Inspection Technique for Wedge Features

Joseph Corcoran[®], Eli Leinov[®], Alejandro Jeketo, and Michael J. S. Lowe

Abstract—Numerous engineering components feature prismatic wedge-like structures that require nondestructive 2 evaluation (NDE) in order to ensure functionality or safety. з This article focuses on the inspection of the wedge-like seal fins of a jet engine drum, though the capabilities presented 5 will be generic. It is proposed that antisymmetric flexural edge modes, feature-guided waves localized to the wedge tips, may be used for defect detection. Although analytical solutions exist that characterize the ultrasonic behavior of ideal wedges, in practice, real wedges will be irregular 10 (containing, for example, truncated tips that are built 11 12 onto an associated structure or have nonstraight edges), and therefore, generic methodologies are required to 13 characterize wave behavior in nonideal wedges. This article 14 uses a semianalytical finite-element (SAFE) methodology 15 to characterize the guided waves in wedge-like features 16 with irregular cross sections to assess their suitability for 17 NDE inspection and compare them with edge modes in 18 ideal wedges. The science and methodologies required in 19 this article are necessary to select an appropriate operating 20 frequency for the particular application at hand. In addition, 21 this article addresses the practical challenge of excitation 22 and detection of flexural edge modes by presenting 23 piezoelectric-based dry-coupled transducer system 24 а suitable for pulse-echo operation. This article, therefore, 25 presents the scientific basis required for industrial exploita-26 tion, together with the practical tools that facilitate use. The 27 study concludes with the experimental demonstration of the 28 edge wave-based inspection of a seal fin, achieving a signal-29 to-noise ratio of 28 dB from a 0.75-mm radial tip defect. 30

31 Index Terms—XXXXX.

I. INTRODUCTION

WEDGE-LIKE features are built on a range of engineering components forming, for example, blades, cutting disks, and seals. To maintain the safety and proper function of these features, it may be necessary to undertake nondestructive evaluation (NDE) to screen for defects or damage, especially at the wedge tips. In this article, we focus on the application of the seal fins that separate pressure stages of a jet engine

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Fig. 1. Photograph of engine drum with inset CAD render of a feature composed of four wedge-like seal fins, the upper two are 2 mm in height and the lower two are 4 mm.

(see Fig. 1) as an example, but the outcomes are widely 40 applicable. Aero engines comprise a number of stages between 41 which temperature and pressure are changed to ultimately 42 generate thrust. The thermodynamic performance of the engine 43 is improved by seal fins-wedge shaped structures a few 44 millimeters in height that run circumferentially between stages 45 and reduce interstage leakage. The inspection requirement in 46 this example is to detect 0.75-mm radial tip defects; this is 47 a crack that is the full width of the wedge, measured from 48 the tip of the wedge to the base and orientated normal to 49 the axial direction of the wedge. The presence of an abrasive 50 coating and restricted access means that scanning along the 51 length of the seal-fin feature is not possible and so a method 52 for screening along the length of the feature from a single 53 location is sought. 54

Antisymmetric flexural edge modes ("edge modes") can 55 propagate in wedge features, as shown in Fig. 2. Multiple 56 orders of edge modes may exist in a wedge feature, with 57 subsequent orders having increasing numbers of nodes, 58 as shown in Fig. 2(b) and (c); throughout this article, the first-59 order wedge mode will be of primary interest, but higher 60 order modes will be discussed where relevant. Wedge modes 61 were discovered by Lagasse [1] and have since been studied 62

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Fig. 2. (a) Illustration of the antisymmetric flexural edge mode, following [4]. Included in red is an illustration of the transducer motion required for excitation and detection of the flexural edge mode. Typically, it is assumed that the transducer footprint should be less than one-fifth of the wavelength. (b) Illustration of the first-order flexural edge mode. (c) Illustration of the second-order flexural edge mode.

in the literature theoretically and experimentally for a few 63 decades [2]–[4]. The desirable characteristics of such modes 64 for NDE are noted in the literature [4], [5], and the waves 65 are guided along the length of the feature with the energy 66 localized to the apex of the wedge; this makes the modes 67 well suited for screening wedge tips in long features. Despite 68 the numerous positive attributes of the edge wave, to date, 69 there has been little reported use in industrial applications. 70 This study aims to address some of the issues limiting uptake, 71 namely, determining dispersion characteristics and localization 72 of ultrasonic energy in irregular wedge-like features, and 73 developing a means of excitation and detection. While recent 74 literature on wedge modes' propagation has been mainly 75 focused on the damping characteristics of wedges as a means 76 to reduce structural vibration [6], [7], their application in 77 NDT has been limited mainly due to the complexity of their 78 defect-scattering phenomena [8], [9]. 79

For effective implementation of an inspection procedure, 80 the dispersion characteristics will need to be evaluated. In an 81 ideal wedge cross section (infinite length, nontruncated, sym-82 metric, and straight-edged), the edge modes are nondispersive 83 with the geometry described by a single nondimensional 84 parameter, the wedge angle θ , as shown in Fig. 2(a) [4]. 85 However, in all real components, the wedge features will be 86 nonideal as they will be of finite length and may be built 87 into a larger structure, and their tip will be truncated to some 88 extent. In nonideal wedge features, the finite geometry makes 89 the edge mode dispersive at the low frequencies that are 90 desirable for practical application; though analytical models 91 describing the influence of some of the irregular wedge-like 92 structures have been proposed [4], [10], it may be difficult to 93 apply them broadly to practical applications. Semianalytical 94 finite-element (SAFE) analysis will be used throughout this 95 study to characterize the ultrasonic behavior in the wedge-like 96 structures. SAFE is an established method for evaluating wave 97 behavior in waveguides of arbitrarily shaped cross section 98 (for example, railway lines [11], [12], beams [12], [13], and 99 stiffeners [14]-[16]) and in feature-guided waves (such as 100 bends [17], [18] and welds [14], [19]). This study extends 101 the current capabilities by specifically addressing waves in 102 irregular wedge-like features and at low frequencies where 103 dispersion may be significant. SAFE will be used to assess 104 the dispersion characteristics, the leakage of the waves into the 105



Fig. 3. Illustration of the hypothetical component with four different wedge-like features. A: Most ideal of the wedges. B: Shortened length. C: Truncated apex. D: Nonstraight sided. All dimensions in millimeters and degrees.

TABLE I ASSUMED MATERIAL PROPERTIES OF IN718 USED THROUGHOUT THIS STUDY

Material Property	Value
Density	8.22 Mg/m ³
Poisson's Ratio	0.294
Young's Modulus	208 GPa

structure the wedges that are built into, and likely sensitivity to defects by assessing the localization of the ultrasonic energy. The tools are intended to provide the scientific basis for utilizing edge modes in real structures.

Beyond establishing a theoretical appreciation of the capa-110 bilities of edge-modes, there is also a standing practical 111 challenge in their excitation and detection. In cases where the 112 wedge has a free end, shear transducers may be used for exci-113 tation and detection [20], but in closed-loop wedges such as 114 the seal fin, there is no free end for transduction. In such cases, 115 in order to excite and detect the edge mode, it is necessary to 116 laterally flex the wedge using a "footprint" significantly less 117 than a wavelength; in practice, this requires excitation and 118 detection through a submillimeter diameter footprint. Exist-119 ing studies rely on excitation by laser [5], [21] or millimeter-120 scale piezoelectric transducers [22] and detection by 121 laser [5], [21], [23] or electromagnetic acoustic transducers 122 (EMATs) [22]. A piezoelectric-based dry-coupled transducer 123



Fig. 4. Dispersion curves of the first 14 edge modes of the body shown in Fig. 3 with material properties shown in Table I. Solutions calculated using an SAFE solver programmed in MATLAB.

system has been developed that is capable of pulse-echo
excitation and detection of edge waves; this development will
enable use of edge waves in industrial applications.

This article starts with a demonstration of the SAFE 127 methodology for establishing the dispersion characteristics of 128 various arbitrarily shaped wedge-like structures with finite 129 geometry. A comparison is made between the edge modes in 130 ideal wedges, and wedge with irregular cross section is made 131 throughout. A further demonstration is then given to illustrate 132 the frequency-dependent localization of ultrasonic energy; this 133 is used to evaluate both leakage and sensitivity to tip defects. 134 The piezoelectric-based dry-coupled transducer system will 135 then be described and demonstrated. This article concludes 136 with the experimental demonstration of the inspection method-137 ology on an example jet engine seal fin with simulated defects. 138

II. SAFE ANALYSIS OF FEATURE-GUIDED WAVES IN WEDGE-LIKE FEATURES

141 A. SAFE Background

The SAFE methodology has been outlined in numerous previous publications [11], [12], [24]–[26] but will be briefly outlined here for completeness. The methodology allows the calculation of the properties of waves in waveguides of arbitrarily shaped cross section. The SAFE method requires the geometry of the cross section of the waveguide to be represented by a 2-D FE discretization and assumes a prismatic geometry, so there is no geometric variation along its axis and the behavior in the propagation direction can be written in an analytical form

$$u_j(x, y, z, t) = U_j(x, y)e^{i(kz - \omega t)}$$
 (1) 152

in which k is the wavenumber, $\omega = 2\pi f$ is the angular frequency, f is the frequency, t is the time variable, and the subscript j = 1, 2, 3. i is the imaginary number $(-1)^{1/2}$. The function U_j represents the behavior in the cross section of the waveguide. For general anisotropic media, the equation of dynamic equilibrium is written in the following form of an eigenvalue problem:

$$c_{iqjl}\frac{\partial^2 U_j}{\partial x_q \partial x_l} + I(c_{i3jl} + c_{iqj3})\frac{\partial (kU_j)}{\partial x_q}$$

$$-kc_{i3j3}(kU_j) + \rho \omega^2 \delta_{ij}U_j = 0 \quad (2) \quad 161$$



Fig. 5. Dispersion curves of the first-order flexural edge modes for the four-wedge structures shown in Fig. 3 with material properties shown in Table I. Solutions calculated using an SAFE solver programmed in MATLAB. Also shown are the mode shapes (arbitrary displacement) of the 2-D discretization at 300 kHz and 1.5 MHz.

with summation over the indices j = 1, 2, 3 and q, l = 1, 2. The coefficients c_{ikjl} are the stiffness moduli, ρ is the density, and δ_{ii} is the Kronecker symbol.

For chosen values of angular frequency ω , eigenvalues of complex wavenumber *k* are found, in which the real part describes the harmonic wave propagation, while its imaginary part presents the attenuation. Many cases of study address waveguides with zero attenuation, which requires only the solution of real eigenvalues, but the possibility of attenuation must be included in the formulation (using the complex *k* representation) for waveguides where there might the loss of energy by material damping or by leakage into the loss of energy by materials. Each solution at a chosen that frequency reveals the wavenumbers of all possible modes at that frequency; then, the full dispersion curve spectrum is that frequency; then, the full dispersion curve spectrum is the desired range of frequencies. The solution of the loss of the solution of the loss of the loss

An in house MATLAB (The Mathworks, Inc., 2014) 179 script was used for the real wavenumber solutions. This 180 is then followed by a separate study using COMSOL 181



Fig. 6. Illustration of the model of an example wedge feature built into a semicircular body to investigate localization of ultrasonic energy. The wedge is elastic, while the surrounding region is absorbing.

(Comsol Inc., 2012), which includes an absorbing boundary
layer in order to investigate possible attenuation by leakage [27].

185 B. SAFE Analysis of Dispersion Characteristics

In this study, a hypothetical component with four wedges 186 built into it has been devised for demonstrative purposes, 187 as shown in Fig. 3. The geometries have been selected to 188 illustrate the forms of geometric characteristic that arises in 189 waveguides of interest such as the example case of engine 190 drum seal fins where multiple wedges share a common base 191 structure. The wedges, labeled A-D, all taper with an angle 192 of 20° but are each different according to: A) the most ideal 193 of the wedges, albeit of finite length and built into a larger 194 structure; B) shortened length; C) truncated apex; and D) 195 nonstraight sided. The engine drum seal fin of interest in 196 this study is similar to Tip C, being of both finite length 197 and containing a truncated tip. The dispersion of the edge 198 waves is influenced by the solid base, but the detail of it is 199 unimportant for the purpose of this study. In principle, there 200 may be coupling due to the common base, but it will be shown 201 that this is insignificant due to the localized nature of the 202 edge modes. Using a common base was chosen for ease of 203 computation as all wedges could be evaluated using a single 204 model. For this analysis, a MATLAB code was used to find 205 the real k solutions. The solver identified the first 14 modes 206 and their dispersion characteristics up to 2.5 MHz. Like the 207 seal fins of interest, the hypothetical component is assumed 208 to be composed of IN718 material, with properties given 209

in Table I. The wavenumber and phase velocity are plotted against frequency in Fig. 4. The solver allowed identification of the different modes, as labeled. 212

The key to understanding the dispersion characteristics 213 of the edge modes is the frequency-dependent localization. 214 Fig. 5 shows only the first-order flexural edge modes for each of the four wedges for improved clarity, together with the corresponding displacement fields. It is shown that the modes 217 become increasingly localized at higher frequency. 218

The first consequence of the increased localization is that 219 the three nontruncated wedges (A, B, and D) converge to the 220 same high-frequency solutions. At low frequencies, the modes 221 interact with the nonwedge-like structures making the mode 222 dispersive. At higher frequencies, the modes become localized 223 only to the wedge tips (which are identical) and so the 224 dispersion characteristics must converge. This convergence 225 appears at the highest frequency for Wedge B as it requires the 226 greatest localization before it becomes insensitive to the built-227 in base; the mode shape and dispersion curves shown in Fig. 5 228 indicate its occurrence at around 300 kHz. Equivalent behavior 229 can be observed in the higher order edge modes. Successively 230 higher order modes are less localized, and therefore, higher 231 frequencies are required to constrain the mode to within the 232 wedge tip and for the dispersion curves to converge. 233

The second consequence of the localization is that the nontruncated wedges converge to be nondispersive. At sufficiently high frequency, the edge modes become wholly localized to the wedge tip and so are insensitive to the remaining structure that causes dispersion. The nontruncated wedge tips may be considered ideal and therefore are 239





Fig. 7. Results of the SAFE study illustrated in Fig. 6. (a) Real part of wavenumber. (b) Imaginary part of wavenumber. (c) Wavelength.

nondispersive. The high-frequency asymptotic behavior 240 is dictated by the nondimensional wedge angle and the 241 nondispersive phase velocity may be calculated analytically 242 from the analysis presented in [4]. For the material properties 243 in Table I, the shear velocity is 3127 m/s, and for a 20° tip, 244 the phase velocity for the first three wedge modes are 1103, 245 2159, and 3226 m/s, respectively. This is in good agreement 246 with the high-frequency asymptotic values shown in Fig. 4. 247

Tip C has a truncated tip, which is of interest due to the
impossibility of fabricating a perfectly sharp wedge. The truncation creates a length scale that in turn makes it dispersive.
McKenna *et al.* [10] calculated that the dispersion measure
that is defined as

 $C \equiv \frac{\omega}{v_g} \frac{dv_g}{d\omega} \tag{3}$

where v_g is the group velocity and a function of the truncation width, h_0 (0.3 mm in Wedge C), and the angular frequency ω according to

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$$C = 6 \left(\frac{2\pi}{\lambda} \frac{h_0}{\operatorname{Tan}\frac{\theta}{2}}\right)^2 \tag{4}$$

Fig. 8. Power flow density of the first-order flexural edge mode at 200 kHz, 500 kHz, and 1 MHz for the feature shown in Fig. 6. The color indicates the axial power flow (for unit power flow, note different color scales at different frequencies). (a) 200 kHz. (b) 500 kHz. (c) 1 MHz.

where $\beta = 2\pi/\lambda$. The dispersion is reduced when the truncation width is small compared with the localization, and this may be achieved by having a small truncation width or reducing the frequency (a longer wavelength).

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The seal fins of interest closely resemble Wedge C. The first-262 order flexural edge wave is dispersive over the full frequency 263 range of interest: at low frequency due to the finite length 264 of the fin and interaction with the base structure and at high 265 frequency due to the increasing significance of the influence of 266 the truncation. Despite this, it is assessed that the dispersion is 267 sufficiently low for practical use above approximately 500 kHz 268 and therefore may be suitable for NDE inspection. 269

C. SAFE Analysis of Leakage

Utilizing edge modes for NDE requires that both the trans-271 mitted and reflected waves can propagate effectively along 272 the wedge over sufficient distances to provide component 273 coverage from a reasonable number of measurement points. 274 From the previous analysis, it was shown that at high frequen-275 cies, the edge modes are localized to the wedge tips, but at 276 lower frequencies, in real components, they may interact with 277 nonwedge features, and this interaction causes energy from 278







Fig. 10. Photograph showing inspection of a seal fin feature of a jet engine drum.

the guided wave to travel out of the feature of interest into the
surrounding media, a process known as leakage [24]. The wave
attenuation due to leakage has to be assessed to ensure that the
mode excited could propagate along the structure effectively.

The phenomenon of wave mode leakage could be quantified using SAFE formalism, and in this case, the SAFE analysis was implemented using Comsol finite element software. Considering the four wedges presented in Section II-B, Wedge B is likely to have the greatest leakage as it is the shortest and therefore requires greater localization for the energy to be contained wholly within the wedge; we therefore chose to study this wedge in greater detail. 290

The implementation in Comsol software comprised the 291 wedge built into a semicircular base, as shown in Fig. 6. 292 The wedge and inner portion of the semicircle are modeled as 293 perfectly elastic, while the outer portion of the semicircle is 294 absorbing modeled by an absorbing material, having the same 295 stiffness properties but also damping properties. The damping 296 properties are introduced as a growing function with radius, 297 following the approach that has been documented in [24]. 298

Any attenuation will, therefore, be a result of wave modes 299 leaking into the absorbing region. The solutions to the wave 300 equation for the modeled cross section will be complex, 301 the real part providing the real wavenumber and the imaginary 302 part indicating the effective attenuation or leakage. 303

The complex wavenumber and wavelength are shown in 304 Fig. 7 for the first- and second-order edge modes. In both 305 cases, the magnitude of the imaginary part of the wavenumber 306 increases with decreasing frequency, and this is expected as the 307 mode becomes less localized and interacts with the attenuating 308 absorbing region. The second-order mode is less localized, and 309 therefore, leaking will occur up to higher frequencies. The root 310 of the first-order mode becomes entirely real above around 311 220 kHz, while the root of the second-order mode becomes 312 entirely real above 500 kHz, indicating that the edge mode is 313 entirely localized to the elastic region above these frequencies 314 and effective propagation can be expected. 315

For inspection purposes, the presence of higher order modes 316 with faster velocities (see Fig. 4) may interfere with the 317 interpretation of inspection results, which utilize the first-order 318 mode. In this regard, leakage-related attenuation of the higher 319 order modes is beneficial. The results of Fig. 7 show that 320 for this particular geometry at 500 kHz, the imaginary part 321 of the first-order mode is zero, while the second-order mode 322 -5.5 m^{-1} . This indicates that if excited, the second-order is 323 mode will be attenuated by 48 dB/m and therefore fall below 324 the noise level within a meter of propagation, i.e., no influence 325 on the detection of flaws is expected from higher order modes. 326 It is worth noting that this particular geometry is only 2 mm 327 high and taller wedges will have less leakage. 328

Practically, to provide full coverage for the inspection of 329 a circumferential seal fin and to exclude interference from 330 higher order modes, measurements are taken from at least two 331 locations. As the interpretation of any signal is done on the 332 assumption of wave speed of the first-order mode, any noise 333 from an unwanted mode will indicate different locations for 334 each test and so can be ignored. 335

To conclude this analysis, as the seal fins of the jet 336 engine are longer than the 2 mm in this example, we can 337 conclude that down to at least 220 kHz, the localization of 338 the first-order mode is sufficient that leakage will be negli-339 gible, leading to effective propagation. It should, therefore, 340 be possible to screen long lengths of the seal fin structure 341 from a single inspection location by choosing to inspect 342 at frequencies greater than around 200 kHz. Higher order 343 modes will also propagate without leakage-related attenua-344 tion at higher frequencies; this provides a motivation to use 345 lower frequencies that must be balanced against the increased 346 dispersion. 347

D. SAFE Analysis of Anticipated Reflection From 348 Defects 349

Given the very localized nature of the edge modes, it is 350 anticipated that the sensitivity to defects in the wedge tips is 351 very high. This has recently been shown for an ideal wedge 352 using FE simulations and experimentally by Chen et al. [5], 353 obtaining reflection coefficients (defined as the reflected 354

energy divided by incident energy) of >0.5 from 0.3-mm 355 defects in 40° wedges at 2 MHz.

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The SAFE analysis described in Section II-A may also be 357 used to approximate the likely response to candidate defects. 358

Fig. 8 shows the power flow density into the cross section 359 of the fin calculated from the stress and particle velocity 360 obtained by the model. The power flow fields have been 361 normalized to the peak power flow density of the cross section 362 to allow comparison across the different frequencies. This 363 corresponds to the wave energy moving in the propagating 364 direction (perpendicular to the cross section) per unit time and 365 area. The presence of a defect at the wedge will disrupt the 366 power flow; the reflection from a defect will increase with its 367 size and the transmission will decrease. 368

As anticipated, it can be seen that the power flow becomes 369 more localized to the fin tip as frequency increases. The 370 reflection coefficient of a through crack will be roughly 371 proportional to the fraction of the power flow that it interrupts; 372 assuming that the defect is situated at the wedge tip, then the 373 reflection coefficient will increase with frequency and defect 374 length (measured from the tip of the wedge to the base). 375

Fig. 8 shows that a defect of just 0.25 mm would interrupt 376 the vast majority of energy of the first-order flexural edge mode at 500 kHz, indicating that the reflection coefficient is expected to be very high. Further quantitative analysis on reflection and transmission coefficients is possible, as shown by Chen et al. [5], and SAFE analysis may complement this effort.

The inspection requirement for the seal fin inspection prob-383 lem is to detect a 0.75-mm radial tip defect. The analysis 384 presented here indicates that when using frequencies of the 385 order of 100 kHz and above, the reflection coefficient is likely 386 to be very high, making the first-order flexural edge mode 387 well suited to the inspection problem. It should be noted 388 that there may potentially be a negative side of the high 389 sensitivity to tip defects; if the mode is highly localized to 390 the tip, then transmission may be impeded by scattering from 391 roughness along the length of the wedge tip. It may, therefore, 392 be necessary to find a compromise between sensitivity and 393 sought defects while suppressing the effects of benign tip 394 roughness. 395

III. EXCITATION AND DETECTION

In order to excite the edge waves, the transducer must 397 laterally displace the wedge tip, as shown in Fig. 2. It must 398 do so with a transducer footprint diameter of less than half a 399 wavelength, as material more than half a wavelength away 400 should be free to be displaced in the opposite direction. 401 Furthermore, practical studies by Fromme et al. [28] have 402 indicated that smaller is better; in practice, we use a "rule-403 of-thumb" of one-fifth, but interested readers could follow 404 up on the work of Fromme or the earlier work by Wilcox 405 cited therein. Fig. 7 shows the calculated wavelengths of 406 the edge mode for the example wedge feature used in 407 Sections II-A and II-B. It can be seen that the transducer 408 footprints of submillimeter are required. 409

A piezoelectric-driven pointed waveguide is proposed, 410 which will focus on the energy to a small, point-like, 411



Fig. 11. A-Scan from (a) defect-free seal fin and (b)–(d) seal fin with a 0.75-mm radial defect. 500 kHz, five-cycle toneburst; 30-V output; $600-\mu s$ sample length; 50-dB receiver gain; 400–600-kHz receiver bandpass filter; 16 averages. Amplitude of each signal is individually normalized to the maximum of the signal beyond 100 μs .

footprint, as shown in Fig. 9(a) and (b). The power transfer efficiency through the waveguide is not critical as the input power can simply be increased to achieve the desired output. Instead, the key performance parameter is the fidelity, and the signal transmitted through the tip should not be too prolonged or distorted by reverberations or undesired modes in the waveguide as a prolonged signal 418 will act to mask any reflected signals being received. The 419 design of waveguides to suppress the undesired modes is 420 an active field of research; the design of the waveguide 421 used in this study was inspired by the flexural damping 422 design of [29]. 423 444

It was found that the capability of the transducer system 424 was not particularly sensitive to the details of the design, but 425 for completeness, the design used in this study was as follows. 426 The transducer is a commercially available Olympus, 10-mm 427 diameter, 2.25-MHz center frequency, piezoelectric transducer 428 (part number: U8477174). The waveguide is 6 mm diameter 429 at the base, <0.5 mm diameter at the tip, and 11 mm long. 430 It is fabricated from 316 stainless steel and the curvature is 431 approximately cubic. The excitation pulse is not dispersion 432 compensated. 433

A simple experiment is conducted to demonstrate the 434 waveguide performance; two transducers are used in a 435 simple through transmission arrangement, as shown in 436 Fig. 9(a) and (b). Measurements were taken with a standard 437 flaw detector within its normal range of operation. The exci-438 tation and received signals are shown in Fig. 9(c). Only a 439 slight reverberation is evident following the initial wave that is 440 unlikely to be an issue in most applications. The reverberations 441 may be further suppressed if necessary by more advanced 442 waveguide design or damping of the free surfaces. 443

IV. EXPERIMENTAL DEMONSTRATION

An example jet engine drum is used to demonstrate the 445 use of edge waves for screening of defects, as shown in 446 Figs. 1 and 10. The drum has several seal fin features, 447 including two with diameters of approximately 420 mm. The 448 wedge feature was similar to Wedge C from Fig. 3, with a 449 0.3-mm-wide flat truncation. The heights of the wedges were 450 greater than the wedge used for the leakage analysis, and 451 therefore, minimal leakage is anticipated. The two fins were 452 otherwise identical except one that was defect free, while in 453 the other, a 0.75-mm radial EDM notch was introduced to 454 replicate a defect. 455

The piezoelectric transducer with pointed tip waveguide, 456 as shown in Fig. 9, was used for the inspections in the pulse-457 echo configuration. The transducer system was lightly pressed 458 normal to the wedge, causing a lateral flexural excitation, 459 relying on dry coupling. A standard handheld ultrasonic flaw 460 detector was used with a 30-V output, 500-kHz five-cycle 461 toneburst output. The receiver chain had 50-dB amplification 462 and a 400–600-kHz bandpass filter, ~ 100 averages were taken. 463 Fig. 11(a) shows an A-scan taken of the defect-free compo-464 nent. The first $\sim 50 \ \mu s$ show the transmitted signal and rever-465 berations as the ultrasound passes through the waveguide into 466 the component. At around 510 μ s, a wave packet arrives that 467 has traveled around the full circumference of the component, 468 a distance of approximately 1.32 m. This result shows the 469 effectiveness of the transducer system and also the excellent 470 propagation of the edge mode. 471

The noise observed in this experiment is coherent. It is believed to originate from excitation of undesired modes, interaction with the bulk of the body, and imperfect surface condition. Two or more readings from separate locations will help suppress the influence of coherent noise.

Fig. 11(b)–(d) shows the A-scans of the defective component, while the transducer location is moved successively further away from the defect. In each case, the reflection from the defect is very apparent, rising well above the noise floor. A typical signal-to-noise ratio (SNR) of 28 dB was calculated from the maximum amplitude between $100-200 \ \mu s$ (the signal) and $200-400 \ \mu s$ (the noise) from Fig. 11(b).

It is, of course, apparent that at least two inspection loca-484 tions are required for a full coverage of a circular component. 485 "Blind spots" are created in the near field of the transducer 486 due to breakthrough and also diametrically opposite from 487 the transducer due to the signal transmitted around the full 488 circumference arriving at the same time as a defect reflection. 489 Minimizing the edge mode dispersion and improving the 490 waveguide fidelity are the key to minimizing the extent of the 491 blind spots. It should be emphasized that efforts have not been 492 made to size the defects but only to screen for their presence. 493 It is envisaged that after they have been identified, they may 494 be sized with a more local inspection procedure. 495

V. CONCLUSION

This article provides the scientific basis to develop practical NDE tools for the inspection of wedge features, demonstrating the analytical tools required for characterization of edge modes in arbitrarily shaped wedge-like structures together with the practical tools required for inspection in industry. 501

Introducing irregular features to a wedge-like geometry pro-502 vides length scales, which may cause the edge wave propagat-503 ing along them to become dispersive. SAFE analysis provides 504 an opportunity to establish the dispersion characteristics of 505 arbitrarily shaped wedge-like features to assess their suitability 506 for guided wave inspection. The SAFE analysis shows that 507 increasing the ultrasonic frequency leads to increasing local-508 ization of the edge modes to the wedge tips. If the tip of 509 the wedge approximates an ideal wedge (straight edged and is 510 not truncated), then at sufficiently high frequencies, the energy 511 becomes sufficiently focused to the dimensionless tip that it 512 becomes nondispersive. 513

For wedges with truncated tips, which will be true of any 514 real fabricated wedge to some extent, the dispersion measure 515 is proportional to the squared ratio of the truncation width 516 and the wavelength. The truncation provides a length scale 517 that will be increasingly significant at higher frequencies; for 518 a given wedge geometry, the dispersion can be suppressed by 519 reducing the frequency (increasing the wavelength). In such 520 cases, the edge mode will be dispersive over the full frequency 521 range of interest: at low frequency due to the finite length of 522 the fin and at high frequency due to the increasing significance 523 of the truncation. An understanding of the science presented 524 in this article is therefore crucial; the frequency must be 525 chosen cognizant of the opposing requirements for sufficient 526 localization to avoid interaction with the base structure, but 527 not too localized for the influence of the truncation to become 528 too significant. 529

SAFE analysis allows the assessment of the anticipated effect of attenuation causing leakage into surrounding structures. It was shown using one example geometry that by increasing the measurement frequency, the ultrasonic energy could be sufficiently localized into the tip so that it becomes insensitive to the surrounding structure allowing good propagation. SAFE analysis also offers a tool for the assessment of

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likely sensitivity to defects by analyzing the power flow in the 537 feature. The more energy that a defect interrupts, the greater 538 the likely reflection coefficient; this could be manipulated by 539 focusing the energy by choice of measurement frequency. 540

The practical use of the short-wavelength edge waves for 541 NDE purposes requires transducers with a submillimeter point-542 like footprint. A simple piezoelectric transducer system with 543 a pointed tip waveguide has been presented for this purpose. 544 The transducer system may be dry coupled and used in a pitch-545 catch configuration. 546

For the example application of the seal fin in this article, 547 500 kHz was found to be suitable; balancing the needs of limit-548 ing dispersion, sensitivity, and wavelength limited transducer 549 footprint. The seal fins of an ex-service engine drum have 550 been used to demonstrate the use of edge modes for defect 551 screening. The demonstration shows the propagation of the 552 edge modes over distances of greater than 1 m and the clear 553 identification of defects from the reflected signals; a typical 554 SNR of 28 dB was achieved. 555

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