

Evidence for orbital motion of CW Leonis from ground-based astrometry

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Accepted 2017 May 24. Received 2017 May 24; in original form 2017 March 3

ABSTRACT

Recent Atacama Large Millimeter/submillimeter Array (ALMA) observations indicate that CW Leo, the closest carbon-rich asymptotic giant branch star to Sun, might have a low-mass stellar companion. We present archival ground-based astrometric measurements of CW Leo obtained within the context of the Torino Parallax Program and with >6 yr (1995–2001) of time baseline. The residuals to a single-star solution show significant curvature, and they are strongly correlated with the well-known *I*-band photometric variations due to stellar pulsations. We describe successfully the astrometry of CW Leo with a variability-induced motion (VIM) + acceleration model. We obtain proper motion and parallax of the centre-of-mass of the binary, the former in fair agreement with recent estimates, the latter at the near end of the range of inferred distances based on indirect methods. The VIM + acceleration model results allow us to derive a companion mass in agreement with that inferred by ALMA, they point towards a somewhat longer period than implied by ALMA, but are not compatible with much longer period estimates. These data will constitute a fundamental contribution towards the full understanding of the orbital architecture of the system when combined with *Gaia* astrometry, providing an ~ 25 yr time baseline.

Key words: astrometry – stars: AGB and post-AGB – binaries: general – circumstellar matter – stars: individual: CW Leo – stars: late-type.

1 INTRODUCTION

The brightest extrasolar object in the sky at $5\ \mu\text{m}$, CW Leo (also known as the infrared source IRC+10216) is the closest carbon-rich, long-period pulsating variable, asymptotic giant branch (AGB) star to Sun, with an estimated distance between 120 and 150 pc (e.g. Groenewegen, van der Veen & Matthews 1998, 2012; Menšchikov et al. 2001 and references therein). Embedded in an expanding cocoon of material originating from the star itself due to the mass-loss process characteristic of the AGB evolutionary phase, this remarkable object has been the subject of extensive studies across a wide range of wavelengths, aimed at understanding its circumstellar environment via detection of new molecules and the characterization of the kinematic, dynamical, chemical, structural and dust properties of the material (e.g. see references in Menten et al. 2012; Decin et al. 2015, D15 hereafter; Kim et al. 2015; Stewart et al. 2016).

One peculiar aspect of the environment surrounding CW Leo is that while the overall shape of the circumstellar mass-loss envelope (CSE) sculpted by CW Leo’s winds appears fairly spherical, high spatial resolution observations have unveiled a variety of structures at different spatial scales and in different wavelength regions: These include spiral shells and dust clumps at subarcsecond scale,

indications of bipolar structure at arcsecond scales and the presence of multiple non-concentric shells in the outer wind. The most recent observational evidence of the complex morphology and kinematics of the inner regions of CW Leo’s CSE has provided circumstantial evidence for the presence of a binary companion. Indirect inferences on the companion properties (mass, orbital separation, eccentricity) are made through comparison of the observations with models of the effects a binary companion has on the wind envelope structure. Such inferences indicate the companion likely being a low-mass K or early M dwarf, while the orbital period is relatively uncertain (from tens to several hundreds of years), and its orbital geometry is even less constrained.

Direct observational evidence for binarity in AGB stars such as CW Leo cannot be readily obtained. In particular, the very high luminosity ($\sim 10^4 L_{\odot}$) of CW Leo and its complex, dynamically evolving dusty envelope make it very difficult to directly detect a low-mass main-sequence companion (Kim et al. 2015; Stewart et al. 2016). It is a challenge to utilize photometric variability monitoring and radial-velocity measurements, because the strong variability intrinsic to CW Leo’s pulsating atmosphere can potentially mask the corresponding variability due to a companion. In this Letter, we present ground-based relative astrometry of CW Leo gathered within the context of the Torino Observatory Parallax Program (Smart et al. 2003, hereafter TOPP) covering a time baseline of over 6 yr. Positional measurements of CW Leo allow

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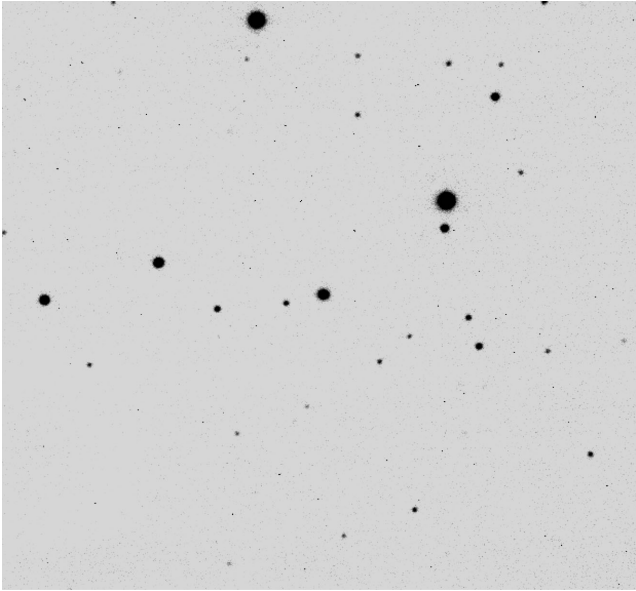


Figure 1. The base CCD image from 1997.26 centred on CW Leo. The field of view is 10×9 arcmin, north is orientated up and east is left.

us to derive a direct distance determination at the near end of the range of published indirect estimates. The astrometric residuals exhibit significant curvature, which we interpret as evidence of orbital motion due to a companion. The TOPP astrometric measurements also allow for detection of variability-induced motion (VIM) that we use to put additional constraints on the orbital architecture of the binary system.

2 OBSERVATIONS AND DATA REDUCTION

CW Leo was observed as part of TOPP from 1995 January to 2001 April. Observations were all carried out on the 1.05-m reflecting telescope that is a scaled-down version of the 1.55-m Kaj Strand Astrometric Reflector at the USNO Flagstaff Station (Strand 1966). The detector used was an EEV CCD05-30 $1296 \times 1152 @ 15 \mu\text{m pix}^{-1}$ constructed by the Astromed company that provided a pixel scale of 0.47 arcsec and a field of view of $10 \text{ arcmin} \times 9 \text{ arcmin}$. All parallax observations were carried out in the Cousins *I* filter. An image of the CW Leo field is presented in Fig. 1.

The full procedures, treatment and reduction for the TOPP programme are described in Smart et al. (1999, 2003). Here, we just briefly list the main steps. All images are flat fielded using nightly sky flats and bias-corrected using image overscan regions. All objects in the field are found and centroided using in-house Gaussian profile fitting software. We then choose a good frame in the middle of the sequence and using anonymous stars in common with other frames, we astrometrically adjust the other frames to this frame. Selection of the reference stars and usable frames is then carried out by eliminating stars with high across season errors, or eliminating frames that have a small number of common stars, or positional residuals larger than three times the average frame residual. Both TOPP astrometry and photometry of CW Leo utilized in the analysis are provided in machine-readable form in the online Journal.

A given sequence is iterated to obtain proper motions and parallaxes for all objects and the above criteria are applied until the estimated parallax of the target changes by less than 1 per cent. We then apply a correction from the calculated relative parallax to an astrophysically useful absolute parallax using estimates of the field

Table 1. Log of observations and single-star solution for CW Leo.

α, δ (J2000)	9:47:57.3, +13:16:43.5
Mean epoch (yr)	1997.2644
π_{abs} (mas)	10.79 ± 4.60
μ_{α^*} (mas yr $^{-1}$)	37.50 ± 1.54
μ_{δ} (mas yr $^{-1}$)	30.22 ± 2.02
Rel. to abs. correction (mas)	1.72
Obs. start, time span (yr)	02/01/1995, 6.30
No. of reference stars, obs.	31, 139

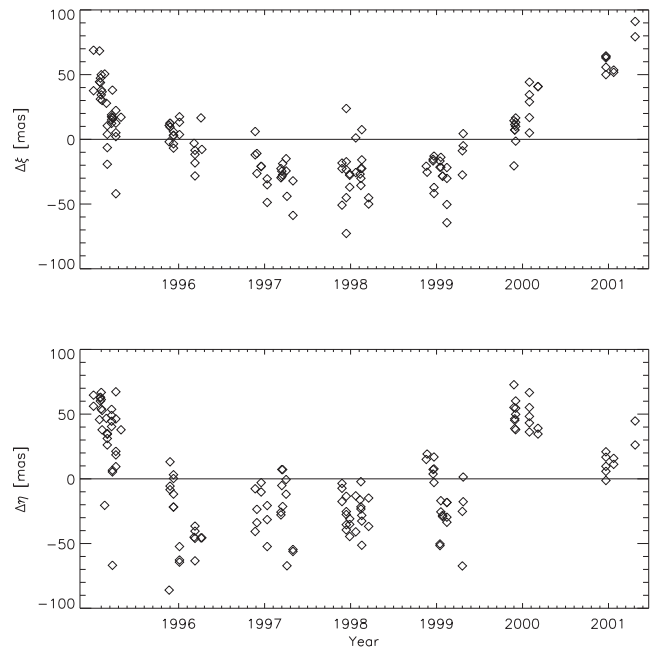


Figure 2. Residuals in standard coordinates ξ, η to a simple single-star solution for all observations not rejected by the standard pipeline rejection criteria.

star distances from the Mendez & van Altena (1998) Galaxy model. The error of the final parallax is found by adding in quadrature to the formal error 33 per cent of the relative to absolute correction.

By default this program assumes the target is a single star, we do not attempt to model at this stage any possible binary nature. In Table 1, we list the parameters of the CW Leo sequence using this single-star astrometric solution. The solution has a reduced chi-square $\chi_r^2 \sim 42$. The excess residuals indicate the presence of unmodelled effects that are not captured by the single-star model.

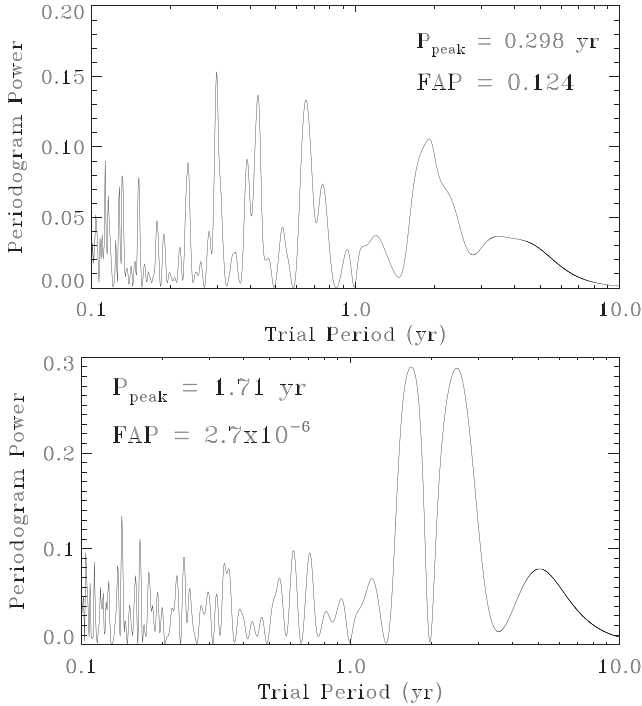
In Fig. 2, we plot the residuals of the CW Leo fit over the observational sequence. The residuals exhibit a distinct curvature, hinting at a period exceeding the ~ 6 yr of the observational campaign. Prompted by the recent analysis of D15, we further analysed the data to see if a binary solution could explain the observed residuals.

3 RESULTS

The evidence of a long-term trend in the residuals to the standard five-parameter solution is statistically solid, based on the Generalized Lomb-Scargle (GLS) periodogram analysis (Zechmeister & Kürster 2009), with a bootstrap-based false-alarm probability FAP $\sim 10^{-40}$. We then attempted to describe the data in terms of a seven-parameter model that includes derivatives of the proper motion. The

Table 2. The acceleration solution for CW Leo. Uncertainties are derived via bootstrap method.

Offset in α (mas)	21.05 ± 1.31
Offset in δ (mas)	30.69 ± 1.34
μ_{α^*} (mas yr $^{-1}$)	30.41 ± 1.47
μ_{δ} (mas yr $^{-1}$)	23.73 ± 1.45
π_{abs} (mas)	12.33 ± 1.90
$\dot{\mu}_{\alpha^*}$ (mas yr $^{-2}$)	8.34 ± 1.24
$\dot{\mu}_{\delta}$ (mas yr $^{-2}$)	8.57 ± 1.23


Figure 3. Top: GLS periodogram of the ξ -axis residuals to the acceleration solution. Bottom: the same, for the η -axis.

best-fitting solution is presented in Table 2. The reduced chi-square of the acceleration solution ($\chi_r^2 \sim 19$) is improved with respect to the five-parameter model value. We quantify the likelihood that acceleration in the astrometric data of CW Leo is detected based on an F -test (in essence, a likelihood-ratio test) that evaluates the significance of the decrease of the χ^2 resulting from the addition of two parameters. The F -test gave a probability of $\sim 10^{-26}$ that the single-star solution is a better description of the data, thus the accelerated motion of CW Leo is considered detected with high statistical confidence.

The addition of statistically significant acceleration terms, however, does not bring χ_r^2 close to unity. One possibility for the excess residuals is that our formal per-measurement uncertainties in the astrometry, having a median of ~ 5 mas, might be underestimated. We nevertheless investigated the residuals of the acceleration solution to see if any additional correlations could be identified that might point to other sources of variability in our positional measurements. Fig. 3 shows the results of the GLS periodogram analysis run on the post-acceleration fit residuals. While no significant power at any period is measured on the ξ -axis, a low-FAP (as determined via a standard bootstrap method) signal at 1.71 ± 0.07 yr is clearly evident in the η -axis residuals (with a second peak of similar power at approximately 2.5 yr). Quite interestingly, this value is very close to the 1.77 ± 0.02 yr period clearly seen in the TOPP photometry

Table 3. The VIM solution with accelerated motion for CW Leo. Uncertainties are derived via bootstrap method.

Offset in α_B (mas)	41.13 ± 2.59
Offset in δ_B (mas)	16.45 ± 2.61
$\mu_{\alpha^*,B}$ (mas yr $^{-1}$)	33.38 ± 1.60
$\mu_{\delta,B}$ (mas yr $^{-1}$)	25.43 ± 1.69
$\pi_{\text{abs},B}$ (mas)	10.56 ± 2.02
offset in α_1 (mas)	-33.39 ± 2.41
offset in δ_1 (mas)	27.55 ± 2.42
$\mu_{\alpha^*,1}$ (mas yr $^{-1}$)	-2.98 ± 1.66
$\mu_{\delta,1}$ (mas yr $^{-1}$)	-2.49 ± 1.75
$\dot{\mu}_{\alpha^*,1}$ (mas yr $^{-2}$)	24.24 ± 1.91
$\dot{\mu}_{\delta,1}$ (mas yr $^{-2}$)	21.69 ± 1.89
$m_{\text{c.o.m.}}$ (mag)	11.24 ± 0.16

that corresponds to the well-known pulsation period of CW Leo (in the range 1.72–1.78 yr; see e.g. Le Bertre 1992; Kim et al. 2015). Indeed, the η -axis residuals to the single-star (resp. acceleration) fit and the photometric time series show a rank-correlation coefficient of 0.53 (resp. 0.49), while no correlation is apparent between the photometry and the ξ -axis residuals.

The excellent match between the periodicity seen in the astrometric residuals and the photometry of CW Leo calls for interpretation in terms of VIM. This is an observable effect in astrometric measurements caused by brightness variations in one of the components of an unresolved double source or blended image, which manifest themselves as a strongly correlated shift of the optical photocentre. VIM effects, suggesting the presence of a binary companion, have been detected in *Hipparcos* observations of variable stars, especially long-period Miras (e.g. Pourbaix et al. 2003, and references therein) and, more recently, in extensive re-analyses of *Kepler* photometry and astrometry (Makarov & Goldin 2016).

We used the Wielen (1996) modelling approach described in Wielen (1996) to gain further insight on the possible architecture of the putative binary companion to CW Leo. As the VIM shows clearly only as a higher-order effect, we first modelled the residuals to the acceleration solution assuming a binary with a fixed relative geometry, described by a linear model with two parameters. This allows us to derive the position angle Θ and a lower limit to the angular distance ϱ between the components (for details, see Wielen 1996). We obtain $\Theta = -12.7 \pm 2.1$ (with Θ measured from north through east) and $\varrho \gtrsim 44$ mas.

We then used the Wielen (1996) formalism in the case of binaries with accelerated motion. In this case, we model simultaneously astrometry and photometry to derive position offsets (in rectangular coordinates), proper motion ($\mu_{\alpha^*,B}$ and $\mu_{\delta,B}$) and parallax $\pi_{\text{abs},B}$ of the centre of mass of the binary (five linear parameters), position offsets (in rectangular coordinates as measured from the barycentre), orbital proper motion ($\mu_{\alpha^*,1}$ and $\mu_{\delta,1}$) in rectangular coordinates moving with the barycentre) and the acceleration ($\dot{\mu}_{\alpha^*,1}$ and $\dot{\mu}_{\delta,1}$) in the coordinate system of the barycentre of the variable component (six linear parameters), plus the total magnitude $m_{\text{c.o.m.}}$ of the system (one non-linear parameter) when the photocentre coincides with the centre-of-mass of the pair. We searched for the minimum of χ_r^2 for the VIM model with accelerated motion as a function of a dense grid of input values for $m_{\text{c.o.m.}}$. Table 3 reports the best-fitting solution ($\chi_r^2 \simeq 15.5$) obtained using a local minimization procedure (Levenberg-Marquardt). Based on an F -test, the VIM model with accelerated motion appears to be superior to the acceleration-only model (probability that the acceleration model better describes the data of 3.0×10^{-9}). The best-fitting VIM+acceleration model is shown superposed to the data in Fig. 4. Attempts at fitting a full orbital VIM model based on a dense grid of trial periods exceeding

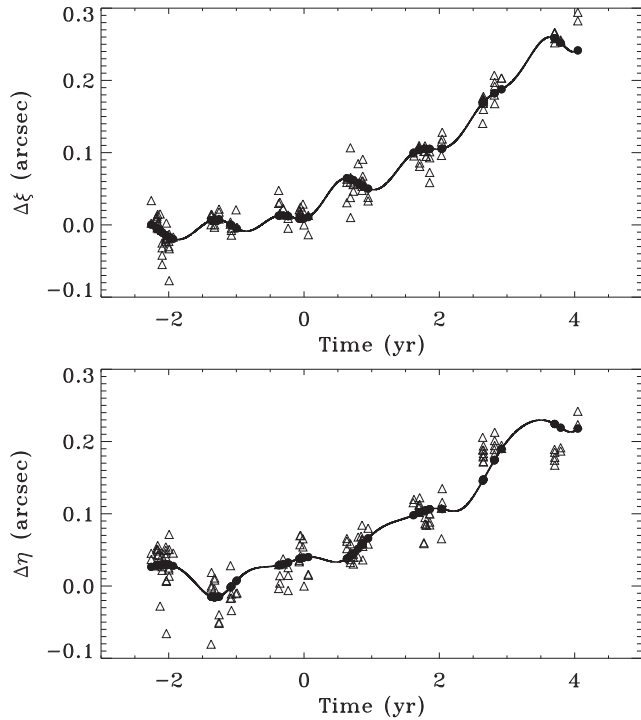


Figure 4. Top: ξ coordinate versus time. The best-fitting VIM+acceleration solution is superposed. Time is relative to the mean epoch (see Table 1). Bottom: the same, for the η coordinate.

the time span of the observations did not produce an improvement in χ_r^2 . The rms of the residuals to the VIM+accelerated motion model (~ 20 mas) still significantly exceeds the typical formal errors in the astrometry. However, a GLS analysis of the residuals to the VIM+acceleration fit did not indicate any significant periodicity, and no correlation was found between the residuals and the photometry. The secondary peak at ~ 2.5 yr initially seen in the η -axis residuals of the acceleration fit is thus to be interpreted as an alias of the main pulsation period. A typical long-term astrometric accuracy at the 15–20 mas level might be more representative of the quality of our astrometric measurements and would bring χ_r^2 close to unity. We also note that the improved fit to the data reduces the parallax error by more than a factor of 2 with respect to the single-star model, bringing it in line with the typical precision of parallax measurements (2–3 mas) as determined for other TOPP targets in the past (e.g. Smart et al. 2003).

4 DISCUSSION

The analysis of >6 yr of unpublished ground-based astrometric measurements from the TOPP program allows us to measure the trigonometric parallax of CW Leo (IRC+10216), the closest carbon-rich star to Sun nearing the end of its AGB lifetime and characterized by extreme mass-loss. The direct distance estimate sits at the lower end of the range of indirect determinations reported in the recent literature (e.g. Groenewegen et al. 2012). Most importantly, we find convincing evidence for the presence of a binary companion to CW Leo, based on the detection of significant curvature in the residuals to a single-star model and on the identification of VIM effects that are successfully modelled simultaneously in astrometry and photometry. Any inferences that can be drawn on the orbital architecture (position angle, separation, orbital period)

and mass of the companion ought to be seen in the context of the growing evidence for the existence of such an object.

To our knowledge, the first claim of binarity for CW Leo dates back to the work of Guélin, Lucas & Cernicharo (1993), who proposed the possible existence of a $1 M_\odot$ companion with a period of ~ 800 yr based on arcsecond-level displacement of a molecular shell from the expected position of CW Leo at mm wavelengths. The tentative identification of a low-mass stellar companion with inferred period in the range 200–800 yr, similar to that proposed by Guélin et al. (1993), was recently announced by (Kim et al. 2015) using the *HST* archival data. However, the detection might be spurious (Stewart et al. 2016). The innermost regions (a few arcsec) of the environment surrounding CW Leo have recently been probed with Atacama Large Millimeter/submillimeter Array (ALMA). Hints of rotating spiral structures in the inner wind envelope have been detected by D15, Cernicharo et al. (2015) and Quintana-Lacaci et al. (2016). Such structures can be interpreted in terms of the presence of a binary companion, but the exact geometry of the spirals is still uncertain. This has important repercussions on the possibility to constrain orbital elements (position angle, separation, eccentricity) and mass of the putative companion. For example, edge-on spiral structures induced by a K- or M-type companion star might place the secondary at ~ 20 – 25 au from CW Leo (depending on the adopted mass for the primary), with an approximate period of 55 yr (D15), but if the spiral structure is seen closer to face-on, the inferred separation might be much larger (~ 65 au) for a similar companion mass (Cernicharo et al. 2015; Quintana-Lacaci et al. 2016).

The results presented here lend further support to the existence of a stellar companion to CW Leo and help in constraining the range of possible orbital architectures and mass ratio of the system. The VIM analysis allows us to infer a small negative value (assuming fixed geometry) for the position angle ($\Theta = -12^\circ$) some 15 yr before the ALMA observations presented by D15, who argue for $\Theta \sim 20^\circ$. The direction of the accelerated motion of the primary agrees with the companion having moved eastwards. The combined effect of the presence of the companion might also help to explain why the μ_{α^*} value derived in this work agrees with recent estimates obtained at mm wavelengths (Menten et al. 2012), while a statistically significant difference in μ_δ with the Menten et al. (2012) value is observed. The discrepancy might be due to confusion between the VIM (contributing mostly in declination) and the proper motion. It then remains to be seen what inferences can be made on the mass and period of the companion based on the magnitudes of the perturbation and VIM effect.

On the one hand, the value of $m_{c.o.m.}$ we derive is proportional to that of the (assumed) constant secondary via the quantity $(1+q)/q$, with $q = M_2/M_1$ being the binary mass ratio. An amplitude of the VIM effect of tens of mas is most likely obtained if the magnitude difference between the two components is not large ($\Delta m \lesssim 2$ mag; Wielen 1996), thus, depending on the assumed mass for the primary, a late-type dwarf companion ($q \sim 0.25$ – 0.40) at ~ 100 pc would roughly fit the scenario outlined by D15. On the other hand, at a distance of ~ 80 – 110 pc as determined from our VIM+accelerated motion solution, the expected angular semimajor axis of the orbital motion of CW Leo around the barycentre would be $\alpha \sim 70$ – 125 mas for a $1 M_\odot$ secondary at 20–25 au, assuming a primary mass of $2.5 M_\odot$, and a circular orbit (D15). The variation of the position angle of the primary over the time span of the TOPP observations (as derived from the VIM+acceleration model) is $\sim 17^\circ$, implying an orbital period $P \approx 130$ yr, which is approximately compatible with the excursion of $\sim 35^\circ$ over ~ 15 yr between the Θ value obtained with the fixed-geometry VIM model and the one argued for by D15.

The amount of curvature in the astrometric residuals to a single-star fit, with an excursion of some 50 mas over 6 yr, would suggest that a significant fraction (10–20 per cent) of the orbital period has been covered by the TOPP observations. This can be made compatible with the tentative P estimate above if the perturbation size were to be increased due to a lower primary mass. Indeed, Decin et al. (2011) argue for an envelope mass of CW Leo of $\sim 2 M_{\odot}$, thus the actual primary mass today might be around $1.5\text{--}2 M_{\odot}$. A value of $\alpha \approx 120\text{--}200$ mas would then be more compatible with the inferred value of P and would indicate an instantaneous orbital proper motion $2\pi\alpha/P \sim 5\text{--}7$ mas yr $^{-1}$, in agreement within 1σ with that derived in our VIM+acceleration solution.

In summary, our results support the D15 evidence for the existence of a low-mass stellar companion to CW Leo with an orbital period shorter than those ($\sim 200\text{--}1000$ yr) derived for other carbon-rich AGB stars, for which a binary companion is thought to be the cause of the detected spiral arm structure. However, our results would tend to indicate a longer orbital period (by a factor of 2 or so) than that tentatively reported by D15. The TOPP observations cannot directly constrain the orbital eccentricity that could very well be significant as a consequence of the mass-loss and mass exchange between the binary components (e.g. Bonačić Marinović, Glebbeek & Pols 2008; Kim et al. 2017).

Future astrometric and photometric measurements with *Gaia* will help tremendously in furthering our understanding of this remarkable system, but the interpretation of *Gaia*'s space-borne astrometry of CW Leo will be improved significantly by the combination with the observations presented here, taking advantage of a time baseline of over 25 yr.

ACKNOWLEDGEMENTS

PG gratefully acknowledges support from INAF through the 'Progetti Premiali' funding scheme of the Italian Ministry of Education, University, and Research. RLS's research was supported by a visiting professorship from the Leverhulme Trust (VP1-2015-063). We thank U. Abbas, J.-L. Halbwachs, P. Lucas and E. Poretti for useful discussion. An anonymous referee provided insightful comments that helped to improve the manuscript.

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SUPPORTING INFORMATION

Supplementary data are available at *MNRASL* online. The data will also be made available at CDS (Strasbourg, France).

Table 4. TOPP epoch astrometry and photometry of CW Leo.

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