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### INFLUENCE OF CAPILLARY HETEROGENEITY ON LEAKAGE OF CO2 FROM A

### BOREHOLE

A Thesis presented in partial fulfillment of requirements for the degree of Masters of Engineering Science in the Department of Geology and Geological Engineering The University of Mississippi

by

### FRANK G. ROECKER

May 2012

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#### ABSTRACT

We used a modified invasion percolation (MIP) model to examine the effect of capillary heterogeneity, buoyancy forces, and viscous forces on the surface area and saturation of a  $CO_2$ plume leaking into a shallow aquifer. The purpose of this study is to provide a better understanding of how CO<sub>2</sub> migrates from a borehole, which is essential in implementing effective simulation and monitoring regimes to accurately detect CO<sub>2</sub> leakage from sequestration sites. The MIP model approach will simulate invasion of a light non-wetting fluid (e.g., CO<sub>2</sub>) into a medium initially saturated with a dense wetting fluid (water). The style of capillary heterogeneity, the strength of buoyancy and viscous forces, and the size of the  $CO_2$  source were systematically varied yielding 168 different simulation scenarios. We find that the interplay between capillary heterogeneity, buoyancy forces, and viscous forces controls the surface area and saturation of gaseous CO<sub>2</sub> leaking into an aquifer system. In unstructured systems with the absence of buoyancy and viscous forces, the CO<sub>2</sub> surface area and saturation are relatively large. In most cases, the CO<sub>2</sub> surface area decreases in weakly stratified systems, and as stratification increases, the CO<sub>2</sub> surface area increases and the CO<sub>2</sub> saturation decrease. Buoyancy forces stretch the invading CO<sub>2</sub> into a narrower structure, resulting in a higher surface area and a lower saturation. Our model implements weak viscous forces which cause CO<sub>2</sub> to pool around the leak source until, at a radial distance, either buoyancy or capillary forces begin to dominate  $CO_2$ invasion. The dissolution rate of gaseous CO<sub>2</sub> into groundwater is proportional to the surface area of the CO<sub>2</sub> phase. Our study shows that variations in the style of capillary heterogeneity and

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the strength of buoyancy and viscous forces can lead to large differences in  $CO_2$  dissolution rates.

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#### I. INTRODUCTION

Carbon capture and sequestration in deep geologic formations is emerging as a possible way to reduce greenhouse gas concentrations. Proposed sequestration sites are commonly located in depleted oil and gas fields, saline aquifers, or deep un-mineable coal seams (*Ide et al.*, 2006; *Folger*, 2009, *Birkholzer et al.*, 2009). On closure, most CO<sub>2</sub> sequestration sites will contain numerous well bores that penetrate cap-rocks of reservoirs and CO<sub>2</sub> sequestration horizons. Degradation of well-bore casing, plugs, and cements can create pathways for buoyant CO<sub>2</sub> to migrate towards the surface (*Ide et al.*, 2006). The consequences of CO<sub>2</sub> leakage from a storage reservoir include the contamination of shallow potable water sources and the direct or delayed release of CO<sub>2</sub> gas back into the atmosphere, negating the long term effects of CO<sub>2</sub> sequestration.

Many researchers (e.g., *Ide et al.*, 2006; *Bachu and Celia*, 2009; *Kopp et al.*, 2010; *Tao et al.*, 2011; *Zhang and Bachu.*, 2011) have identified leakage along well bores as the most likely pathway for CO<sub>2</sub> leakage from a storage reservoir. These workers have developed models that explore CO<sub>2</sub> leakage from abandoned wells [*Kopp et al.*, 2010; *Nordbotten et al.*, 2005; *Humez et al.*, 2011], along the annular cement [*Bachu and Bennion*, 2009; *Viswanathan et al.*, 2008], and through multiple geological layers and multiple leaky wells [*Nordbotten et al.*, 2009]. Others have focused on how the well-bore cement reacts in a CO<sub>2</sub> saturated environment [*Carey et al.*, 2007; *Carey et al.*, 2010; *Huerta et al.*, 2011; *Geloni et al.*, 2011]. Upon leakage from a well bore, escaped CO<sub>2</sub> may come into contact with shallow aquifers and migrate through the porous medium away from the well bore (e.g., *Oldenburg et al.*, 2010). CO<sub>2</sub> dissolution into

groundwater can lower pH, altering the aquifer geochemistry [*Wang and Jaffe.*, 2004; *Zheng et al.*, 2009; *Kharaka et al.*, 2010; *Keating et. al.*, 2010; *Peter et al.*, 2011] and possibly triggering movement of trace metals [*Apps et al.*, 2010; *Peter et al.*, 2011].

Direct detection of CO<sub>2</sub> leakage in the shallow subsurface can be difficult due to large variations in natural background CO<sub>2</sub> fluxes (*Lewicki et al.*, 2007). Other workers [*Romanak et al*, 2012] have suggested that well-bore leaks into aquifers may be identified through long-term monitoring of groundwater. Measuring accurate amounts of dissolved CO<sub>2</sub> in groundwater is difficult, (*Newell et al.*, 2008), but solute plumes containing newly dissolved minerals may be more readily identified (*Benson and Cole*, 2008). The character of solute plumes resulting from CO<sub>2</sub>-driven alteration of aquifer materials will vary with the local CO<sub>2</sub> dissolution rate and the total mass of CO<sub>2</sub> within the aquifer, which control the evolution of aquifer pH near leaking boreholes. The dissolution rate of CO<sub>2</sub> into groundwater is limited by the surface area of the CO<sub>2</sub>-groundwater interface, and the mass of CO<sub>2</sub> within the aquifer is a function of the CO<sub>2</sub>-phase saturation. CO<sub>2</sub>-phase saturation and surface area are directly influenced by the capillary heterogeneity within the aquifer, the magnitude of local buoyancy forces, and the influence of viscous forces arising from the displacement of groundwater by CO<sub>2</sub>.

We examine the effect of capillary heterogeneity, buoyancy forces, and viscous forces on the surface area and saturation of a CO<sub>2</sub> plume in a shallow aquifer, resulting from a slow wellbore leak. The displacement of the denser wetting phase (e.g., water) by a lighter non-wetting phase (e.g., CO<sub>2</sub> gas) is simulated using a quasi-3D, Modified Invasion Percolation (MIP) model [e.g., *Ioannidis et al.*, 1996; *Glass et al.*, 2001; *Holt*, 2003] that includes capillary forces, buoyancy forces, and a first-order approximation of viscous forces. Over 8,400 simulations are

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employed in a Monte Carlo approach used to quantify the influence of and uncertainty due to variations in capillary heterogeneity, buoyancy forces, and viscous forces on  $CO_2$  surface area and saturation. The style of capillary heterogeneity is systematically varied from unstructured to stratified to mimic the range of capillary heterogeneity found in nature. Similarly, the strength of buoyancy and viscous forces are also varied, and the influence of the leak source size is evaluated. All simulations are non-dimensional and can be scaled using the gas-entry pressure of a porous medium.

We find that the interplay between capillary heterogeneity, buoyancy forces, and viscous forces controls the surface area and saturation of gaseous  $CO_2$  derived from well-bore leaks in aquifer systems. In stratified systems, the  $CO_2$  surface area increases and the  $CO_2$  saturation decreases with increasing stratification and buoyancy forces, both of which focuses  $CO_2$  invasion into narrow pathways.  $CO_2$  surface area decreases and  $CO_2$  saturation increases in the presence of viscous forces, which leads to pooling of  $CO_2$  near the source. As the leak source size increases,  $CO_2$  invades along more pathways near the source, increasing the  $CO_2$  surface area and saturation.

#### **II. MODIFIED INVASION PERCOLATION MODEL**

An MIP approach is used to simulate invasion of a light non-wetting fluid (gas) into a medium initially saturated with a dense wetting fluid (water) under the combined influence of capillary, buoyancy, and viscous forces. Compressibility and dissolution of CO<sub>2</sub> phase, trapping, and other dynamic effects, such as fragmentation and pulsation [e.g., *Wagner et al.*, 1997; *Vedvik et al.*, 1998; *Glass and Yarrington*, 2003], are neglected here, as our focus is on the broad factors affecting CO<sub>2</sub> dissolution rates and concentrations in aquifer systems. The current MIP differs from Invasion Percolation (IP) [e.g., *Wilkinson and Willemsen*, 1983] through the definition of macro, near-pore scale capillarity. Here, individual pore throats and necks are not considered. Instead, a near-pore scale block is defined and characterized by a local threshold spanning pressure (a local block-scale breakthrough pressure) that represents the behavior of the subscale network (*Holt et al.*, 2003). The model domain is discretized into an array of grid blocks with assigned spanning pressures. The invasion pressure for each block is determined by the sum of the spanning pressure, gravity forces, and viscous forces:

$$P_I^* = P_S^* + \Delta \rho g Z^* + \Delta P_v^* \tag{1}$$

where  $P_S^*$  is the capillary spanning pressure (ML<sup>-1</sup>T<sup>-1</sup>),  $\Delta \rho = \rho_{defender} - \rho_{invader}$  is the difference in densities of the defending phase and invading phase (ML<sup>-3</sup>), *g* is the acceleration due to gravity in the Z-direction (LT<sup>-2</sup>), *Z*<sup>\*</sup> is the coordinate in the Z-direction (L), and  $\Delta P_v^* = \left(P_v^{defender^*} - P_v^{invader^*}\right)$  is the difference in viscous pressure (ML<sup>-1</sup>T<sup>-1</sup>). Dividing equation (1) by the average spanning pressure  $\langle P_S^* \rangle$ , our dimensionless invasion pressure becomes:

$$P_{I} = \frac{P_{S}^{*}}{\left\langle P_{S}^{*} \right\rangle} + B_{O}Z + \frac{\left(P_{v}^{d^{*}} - P_{v}^{i^{*}}\right)}{\left\langle P_{S}^{*} \right\rangle}$$
(2)

where  $\langle P_S^* \rangle$  is the average spanning pressure (ML<sup>-1</sup>T<sup>-1</sup>), the dimensionless Bond Number is defined as:

$$B_{O} = \frac{\Delta \rho g \Delta Z^{*}}{\left\langle P_{S}^{*} \right\rangle} \tag{3}$$

 $\Delta Z^*$  is the grid block length in the Z-direction (L). The dimensionless Z-coordinate is defined as:

$$Z = \frac{Z^*}{\Delta Z^*} \tag{4}$$

The Bond Number reflects the strength of the buoyancy forces. For the simulations presented here, the Bond Number was systematically varied between 0, 0.1, and 1 to represent increasing buoyancy forces.

The viscous pressure in the defending fluid was set at zero for all simulations, to simulate static, non-flowing groundwater conditions. The viscous pressure in the invading fluid is calculated by a first-order approximation for steady flow from a point source in spherical coordinates with the reduced equation being:

$$\frac{P_{v}^{i^{*}}}{\left\langle P_{S}^{*}\right\rangle} = VSN\left(\frac{1}{r} - \frac{1}{R}\right)$$
(5)

where *VSN* is the dimensionless viscous scaling number, *r* is the dimensionless grid-block radial coordinate (-), and *R* is the dimensionless radius of influence (-). An *R* value of 1000 was kept constant for all simulations. We selected *VSN* values of 0 and 250; a *VSN* value of 250 is equivalent to a CO<sub>2</sub> volumetric discharge rate of 0.03 m<sup>3</sup>/d in a typical sandstone and  $3 \times 10^{-4}$  m<sup>3</sup>/d in a typical mudrock or shale.

Correlated spanning-pressure fields are generated using an FFT approach (e.g., Robin et al., 1993). First, a two-dimensional, correlated random field (512 by 512) with a standardnormal distribution is generated, then the value at each nodal location is mapped to an associated probability, and finally the nodal spanning pressure is determined from the nodal probability (P) using the van Genuchten (1980) pressure-saturation model:

$$P_{S} = \frac{1}{\alpha} \left\{ \left[ \left( \frac{1}{P} \right)^{\frac{1}{1-\frac{1}{N}}} \right] - 1 \right\}^{\frac{1}{N}}$$

$$\tag{8}$$

where  $P_s$  is the spanning pressure,  $\alpha$  is a model parameter related to the air-entry pressure, N is a model parameter associated with the pore-size distribution. Alpha ( $\alpha$ ) values of 1 and N values of 1.5 were kept constant for all simulations. The style of capillary heterogeneity in each field is defined by a fixed correlation ratio, the ratio of the horizontal to vertical correlation lengths. Correlation ratios of 1, 10, 20, 30, 50, 100, 500, and 1000 were used to reflect increasing horizontal stratification. Values of 1, 10, and 50 are used as the dimensionless vertical correlation lengths (Figure 1). As the vertical correlation length increases, the horizontal correlation lengths are increased to maintain specified correlation ratios.

During simulations, non-wetting phase invasion proceeds from a source on the left edge of the domain towards the top or right edge domain boundary according to an IP algorithm that implements equation (2). We apply a network connectivity of 8 (communication is allowed with all surrounding blocks) to approximate three-dimensional behavior within a two-dimensional network (e.g., Holt et al., 2003). The IP algorithm sorts all the invadable nodes along the wetted surface and invades the block with the lowest invasion pressure. This process continues until the blocks along the active interface reach either the top or right edge boundaries, the percolation threshold. In the absence of buoyancy and viscous forces ( $B_0 = 0$  and  $\Delta P_v^* = 0$ ), standard percolation will exactly reproduce the drainage pressure-saturation curve given by equation (8) (*Holt et al.*, 2003). However, when  $\Delta P_v^* \neq 0$  and  $B_o \neq 0$ , this is not the case.

Our model simulates leakage from a point source and a line source along the left edge of the domain. The model defines an initial wetted surface of  $CO_2$  that acts as the leak source. The wetted surface is defined as all the possible nodes that are in contact with the growing cluster. The initial wetted surface for a point source allows invasion of  $CO_2$  only into the bottom-most node on the left-hand side of the domain. This small source size is analogous to  $CO_2$  escaping through a fracture in the well-bore casing or cement. For the larger leak source, the bottom 102 nodes (10% of the domain size) along the left-hand side of the domain are considered invadable at the start of the simulation. This arrangement simulates  $CO_2$  leaking from a well bore through degraded annular cement. Studies have shown that well-bore cement can degrade after 200 years to allow potential fluid flow analogues to that through a silty sand (*EPA*, 2009).

Simulations including viscous forces were only conducted for a point source in random spanning pressure fields with a vertical correlation length of 1 grid block, as more complicated source sizes and styles of heterogeneity would lead to greater errors in our first-order approximation.

In the following discussion, the term "surface area" is defined as the surface area per unit volume of invaded.

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#### **III. RESULTS**

We focus on the final  $CO_2$  surface area (ratio of the  $CO_2$ -water surface area to  $CO_2$ volume) and the final  $CO_2$  saturation at percolation to either the right or top boundary. 8,400 simulations were run to evaluate the influence of capillary heterogeneity, buoyancy forces, and viscous forces on the  $CO_2$  surface area and saturation. 168 different combinations were created by varying the correlation ratio, Bond Number, strength of first-order viscous forces, the Zcorrelation length, and leak source size (Table 1). 50 simulations were run for each combination.  $CO_2$  surface area and saturation results are summarized using the average and standard deviation for each 50-simulation set.

Figures 2 and 3 show the influence of varying styles of stratification on the evolution of the CO<sub>2</sub> phase in the absence of buoyancy and viscous forces for the case of a point source and a line source, respectively. Both source sizes show similar results. For unstructured media (Z-correlation length of 1 and correlation ratio of 1), the CO<sub>2</sub> surface area is fairly large (~1) as the CO<sub>2</sub> phase distribution is controlled by the random distribution of spanning pressures. The addition of structure, in the form of either vertical or horizontal correlation, decreases the surface area (~0.6), as invasion fills adjacent grid blocks, which have similar spanning pressures. As the degree of stratification increases, the CO<sub>2</sub> surface area increases, reflecting less vertical overlap in the CO<sub>2</sub> phase. As the correlation ratio becomes large, invasion tends to proceed similarly regardless of the vertical correlation length, and the average surface areas for all vertical correlation lengths become similar. The average surface area from a line source (Figure 3) (~1.5) is greater than that of a point source (Figure 2) (~1.2) at large correlation ratios, as the i

nvading  $CO_2$  takes multiple pathways from the source. Average  $CO_2$  saturation increases with increasing Z-correlation length, as regions of similar spanning pressure become completely filled with the invading  $CO_2$ . The average  $CO_2$  saturation decreases with increasing correlation ratio, because  $CO_2$  invasion becomes focused to fewer strata and the percolation threshold is reached in fewer invasion steps.

Figures 4 and 5 display how intermediate buoyancy forces ( $B_0 = 0.1$ ) affect CO<sub>2</sub> phase invasion from a point and a line force respectively. Both source sizes show similar results. For unstructured media (Z-correlation lengths of 1 and a correlation ratio of 1), the CO<sub>2</sub> surface area is fairly large (~1.3) and CO<sub>2</sub> saturation is low (~0.02), as the CO<sub>2</sub> phase is being stretched vertically into a narrow structure by buoyancy forces. When the vertical correlation ratio increases from 1 to 10 in un-stratified media, CO<sub>2</sub> fills adjacent grid blocks with similar spanning pressures, decreasing CO<sub>2</sub> surface area (~0.6) and slightly increasing CO<sub>2</sub> saturation (~0.03). When the vertical correlation length reaches 50, however, stronger vertical correlation of spanning pressures allows buoyancy forces to dominate the local capillary heterogeneity, leading to slightly higher surface areas.

For both leak source sizes (Figures 4 and 5), buoyancy forces become less influential over the  $CO_2$  plume as the medium becomes more stratified, and the percolation threshold is reached on the right boundary. Instead of a narrow, vertical structure, the final  $CO_2$  phase distribution reflects trapping beneath horizontal capillary barriers of higher spanning pressure. When the vertical correlation length is 1, buoyancy forces are strongest, leading to vertical stacking of  $CO_2$ -invaded layers, lower  $CO_2$  surface areas, and higher  $CO_2$  saturations. Weakly stratified systems (correlation ratio of 10) show lower surface areas (~0.7) and higher saturations (~0.05) than the unstructured systems, as  $CO_2$  fills more vertically adjacent gridblocks. As

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stratification increases,  $CO_2$  invasion proceeds along an increasingly horizontal pathway, increasing  $CO_2$  surface area and decreasing  $CO_2$  saturation. At large correlation ratios, the average surface area from a line source (Figure 5) (~1.5) is greater than that from a point source (Figure 4) (~1.2), because  $CO_2$  can take multiple pathways from the source.

Figures 6 and 7 show how strong buoyancy forces ( $B_0 = 1$ ) affect CO<sub>2</sub> invasion from a point source and from a line source, respectively. Both source sizes show similar results. In unstructured media (Z-correlation lengths of 1 and a correlation ratio of 1), the CO<sub>2</sub> surface area is very large (~2.2) and the CO<sub>2</sub> saturation is very low (~0.004), because the CO<sub>2</sub> phase is being stretched vertically into a very narrow structure. In weakly stratified systems, buoyancy forces remain strong relative to capillary forces, and vertical stacking of CO<sub>2</sub>-invaded layers causes lower surface areas (~0.8 for Z = 1) and slightly higher saturations (~0.025 for Z = 1). This process is particularly important for simulations run with a vertical correlation length of 1, where strong buoyancy forces lead to percolation at the upper boundary and increased vertical stacking to correlation ratios of 30 and 50. As the correlation ratios increase, horizontal capillary forces dominate the system, vertical stacking is reduced, and the percolation threshold is quickly reached at the right boundary; under these conditions, CO<sub>2</sub> surface areas increase, and CO<sub>2</sub> saturations decrease. The CO<sub>2</sub> surface areas for line sources are generally slightly larger than those observed from point sources.

The influence of viscous forces on  $CO_2$  invasion with increasing levels of stratification and increasing buoyancy forces from a point source is show in Figures 8. Here, results for VSN = 250 and a vertical correlation length of 1 are presented. Viscous forces lead to pooling of  $CO_2$ around the source, which decreases  $CO_2$  surface area and increases  $CO_2$  saturation when compared with the results for simulations where viscous forces are absent (e.g., Figures 2, 3, and 6). The effect of  $CO_2$  pooling is most apparent in strongly stratified systems, where the  $CO_2$  surface area is greatly reduced (~0.7) when compared with simulations without viscous forces. For example, in Figure 2, where  $CO_2$  invasion is only controlled by spanning pressures,  $CO_2$  surface areas were around 1.2 in strongly stratified systems (correlation ratio of 1000). As buoyancy forces were increased, the  $CO_2$  surface area increases to ~1.2 and ~1.5 at Bond Numbers of 0.1 (Figure 4) and 1 (Figure 6), respectively.

The overall trends observed in the surface area and saturation plots in Figure 8 are similar to those reported above. Here, the average CO<sub>2</sub> surface area in unstructured media increases as buoyancy forces are increased, while the average CO<sub>2</sub> saturation decreases. Increasing stratification initially leads to vertical stacking of CO<sub>2</sub> saturated layers, reducing the surface area. As stratification is increased further, the CO<sub>2</sub> surface area increases slightly, but does not reach the high values observed in the absence of viscous forces. When buoyancy forces are absent, CO<sub>2</sub> saturation is very large (~0.2), and then decreases as stratification increases. In the presence of buoyancy forces, saturations peak (~0.08) in weakly stratified media as CO<sub>2</sub> is vertically stacked in overlapping layers. As stratification increases, CO<sub>2</sub> invasion becomes more strongly focused into horizontal layers, reducing saturations (~0.02).

The coefficient of variation ( $C_v$ ) of the CO<sub>2</sub> surface area (Figure 9) reflects the relative uncertainty in the average CO<sub>2</sub> surface area. The pattern of the surface area  $C_v$  versus correlation ratio or stratification strength is varies mainly with the vertical correlation length. When the vertical correlation length is 1 (Figure 9a), the  $C_v$  generally increases as stratification becomes stronger. Here, the unstructured systems are more closely ergodic, and the systems become less ergodic (more variability in invasion pathways) as the correlation ratio increases. When the vertical correlation length is 10 (Figure 9b), the surface area  $C_v$  initially increases and then decreases, as stratification becomes stronger, reflecting a shift from mainly upward to dominantly horizontal invasion patterns. When the vertical correlation length is 50 (Figure 9c), the surface area  $C_v$  generally decreases with increasing stratification, as CO<sub>2</sub> invasion proceeds along increasingly narrow, horizontal pathways. The surface area  $C_v$  during CO<sub>2</sub> invasion with viscous forces into media with a vertical correlation length of 1 is similar to that observed in Figure 9a, indicating that the addition of viscous forces does not lead to significantly different CO<sub>2</sub> invasion patterns between simulations.

In general, the uncertainty increases when  $B_o$  increases from 0 to 0.1, as CO<sub>2</sub> invasion pathways are pulled upward by buoyancy forces. A further increase in  $B_o$  from 0.1 to 1, generally leads to a decrease in the uncertainty, as invasion is focused into narrower pathways.

Our results show that the interplay between capillary heterogeneity, buoyancy forces, and viscous force controls the dissolution rate and aqueous concentrations of CO<sub>2</sub> derived from wellbore leaks in aquifer systems. The dissolution rate of gaseous CO<sub>2</sub> into groundwater is proportional to the surface area of the CO<sub>2</sub> phase. Our study shows that variations in the style of capillary heterogeneity and the strength of buoyancy and viscous forces can lead to differences in CO<sub>2</sub> dissolution rates of well over an order of magnitude (factor of ~17 in this study). Depending on the scenario, the CO<sub>2</sub> surface area increases by a factor between 3 and 9, as the degree of stratification increases. As buoyancy forces increase (from  $B_o = 0$  to  $B_o = 1$ ), the CO<sub>2</sub> surface area typically increases by a factor of ~2. The addition of weak viscous forces lowers the CO<sub>2</sub> surface area by a factor of 1.1 in unstructured media and a factor of 2 in well stratified media. Changing from a point to a line source can increase CO<sub>2</sub> surface areas by a factor of 1.3.

The  $CO_2$  saturation is a metric for the dissolved  $CO_2$  concentration in groundwater, assuming all gaseous  $CO_2$  is dissolved into groundwater, and the geochemical impact of the  $CO_2$  well-bore leak, as high CO<sub>2</sub> saturations will eventually lead to high dissolved CO<sub>2</sub> concentrations. Increasing stratification decreases CO<sub>2</sub> saturation by a factor of between 8 and 60, depending upon the scenario. The addition of strong buoyancy forces ( $B_o = 1$ ) decreases CO<sub>2</sub> saturation by a factor of 23 to 35 in unstructured media. Strong buoyancy forces reduce saturations by a factor of 1.2 to 3.6 in weakly-stratified media, and increase saturations by a factor of 1.1 to 2.5 in well-stratified media. Saturations increase by factors between 1.2 and 6, when weak viscous forces are added. The presence of a line source, rather than a point source, can increase CO<sub>2</sub> saturations by up to a factor of 4.

#### **IV. SUMMARY AND DISCUSSION**

We used a quasi-3D, MIP model to examine the effect of capillary heterogeneity, buoyancy forces, and first-order viscous forces on the surface area (per unit volume) and saturation of a  $CO_2$  plume in a shallow aquifer, resulting from a slow well-bore leak. Compressibility and dissolution of  $CO_2$  phase, trapping, and other dynamic effects, such as fragmentation and pulsation, were neglected. The style of capillary heterogeneity, the strength of buoyancy and viscous forces, and the size of the  $CO_2$  source were systematically varied yielding 168 different scenarios. Each scenario was then simulated 50 times and the  $CO_2$  surface area and saturation was determined. The style of capillary heterogeneity was systematically varied from unstructured to stratified to mimic the range of capillary heterogeneity found in nature.

We find that the  $CO_2$  surface area and saturation is governed by the interplay between capillary heterogeneity, buoyancy forces, and viscous forces. In stratified systems, the  $CO_2$ surface area increases and the  $CO_2$  saturation decreases with increasing stratification and buoyancy forces, both of which focuses  $CO_2$  invasion into narrow pathways. Viscous forces, which lead to pooling of the  $CO_2$  near the source, lead to lower  $CO_2$  surface area and higher  $CO_2$ saturation. As the leak source size increases,  $CO_2$  invades along more pathways near the source, increasing the  $CO_2$  surface area and saturation. Variations in the style of capillary heterogeneity and the strength of buoyancy and viscous forces can lead to differences in  $CO_2$  dissolution rates of well over an order of magnitude, and  $CO_2$  saturations can vary by nearly two orders of magnitude. Our results provide critical insight into the processes affecting  $CO_2$  dissolution rates and aqueous  $CO_2$  concentrations resulting from well-bore leaks into groundwater aquifers and can inform continuum-based modeling of well-bore leak processes and geochemical reactions within the aquifer. Knowledge of the style of capillary heterogeneity and the strength of buoyancy and viscous forces is essential for determining apparent diffusion coefficients and dissolution rates for  $CO_2$  gas into groundwater. The capillary pressure-saturation relationships used in most continuum-based models should be modified to include the impact of stratification, the directionality of the invasion process, and the strength of local buoyancy and viscous forces. Prior to modeling  $CO_2$  dissolution processes in aquifer materials, site-specific geologic data (core descriptions, samples, etc.) are required to identify and characterize the extent of capillary heterogeneity.

Slow leaks from boreholes may be difficult to detect in well-stratified geologic materials, as high  $CO_2$  dissolution rates and low  $CO_2$  saturations may make geochemical changes in the aquifer difficult to identify. In unstructured (e.g., sedimentary media affected by bioturbation or pedogenesis) and weakly stratified materials,  $CO_2$  well-bore leaks may be more readily identifiable from geochemical changes in groundwater, as  $CO_2$  saturations will be higher.

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List of Appendices

Appendix: A

Table 1.	
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Simulation variables and values	
Viscous Scaling Number (VSN)	250
Leak Source Sizes	1, 102
Vertical Correlation Lengths	1, 10, 50
Correlation Ratios	1, 10, 20, 30, 50, 100, 500, 1000
Bond Number	0, 0.1, 1

Simulation variables and values -

Table 2.	Summary of	simulation results a	as changes to	the CO	O <sub>2</sub> surface	area and	saturation.
Values an	re factors of in	ncrease or decrease					

Increase	<u>Surface Area</u>	<u>Saturation</u>	
Stratification	1.5	8-60	
Vertical Correlation Length	3-4	1.3 - 1.6	
Source Size	1.2 – 1.7	1.2-4	
Buoyancy Forces	1.3 – 3.7	6.5 – 35	
Viscous Forces	1.1 - 1.7	1.7 – 6.2	



**Figure 1.** Random fields with (a) a horizontal correlation length of 20 and a vertical correlation length of 1, and (b) a horizontal correlation length of 100 and a vertical correlation length of 50. The horizontal correlation length will vary to maintain correlation ratios.



**Figure 2.**  $CO_2$  invasion history (I) from a point source in the absence of buoyancy forces and viscous forces in (a) unstructured media, (b) weakly stratified media, (c) more stratified media, and (d) highly stratified media. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 3.**  $CO_2$  invasion history (I) from a line source with a Bond Number = 0 in (a) unstructured media, (b) weakly stratified media, (c) more stratified media, and (d) highly stratified media. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 4.**  $CO_2$  invasion history (I) from a point source with a Bond Number = 0.1 and no viscous forces in (a) unstructured media, (b) weakly stratified media, (c) more stratified media, and (d) highly stratified media. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 5.**  $CO_2$  invasion history (I) from a line source with a Bond Number = 0.1 and no viscous forces in (a) unstructured media, (b) weakly stratified media, (c) more stratified media, and (d) highly stratified media. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 6.**  $CO_2$  invasion history (I) from a point source with a Bond Number = 1 and no viscous forces in (a) unstructured media, (b) weakly stratified media, (c) more stratified media, and (d) highly stratified media. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 7.**  $CO_2$  invasion history (I) from a line source with a Bond Number = 1 and no viscous forces in (a) unstructured media, (b) weakly stratified media, (c) more stratified media, and (d) highly stratified media. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 8.**  $CO_2$  invasion history (I) from a point source with VSN = 250 with Bon Numbers of (A) 0, (B) 0.1, (C) 1. Invasion pictures for correlation ratios of 100, 500, and 1000 are not shown because the invasion phase only populates a few layers near the bottom of the domain.  $CO_2$  invasion progress from early to late is reflected by the progression of dark to light colors. Grid blocks saturated with the wetted phase are shown in gray. Plots displaying the (II) average surface area, and (III) average saturation for theses simulations versus increasing stratification. The average saturation is plotted between -0.5 to 0.3 to clearly illustrate results close to zero. No simulations produced saturations less than zero.



**Figure 9.** Plots displaying the coefficient of variation (Cv) of the CO<sub>2</sub> surface area for correlation lengths of (a) 1, (b) 10, and (c) 50 from a line source. Plot (d) displaying Cv of the CO<sub>2</sub> surface area for a correlation of 1 with a VSN = 250 invading from a point source.

Appendix: B

The appendix for this paper is not attached. It includes a copy of the program code, user manuals, and all program files.

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