

THE VMC SURVEY XXII. HIERARCHICAL STAR FORMATION IN THE 30 DORADUS-N158-N159-N160 STAR-FORMING COMPLEX

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ABSTRACT

We study the hierarchical star formation in a ~ 1.5 deg² area covering the 30 Doradus-N158-N159-N160 star-forming complex with the VISTA Survey of the Magellanic Clouds. Based on the young upper main-sequence stars, we find that the surface densities cover a wide range of values, from $\log(\Sigma \cdot \text{pc}^2) \lesssim -2.0$ to $\log(\Sigma \cdot \text{pc}^2) \gtrsim 0.0$. Their distributions are highly non-uniform, showing groups that frequently have sub-groups inside. The sizes of the stellar groups do not exhibit characteristic values, and range continuously from several parsecs to more than 100 pc; the cumulative size distribution can be well described by a single power law, with the power-law index indicating a projected fractal dimension $D_2 = 1.6 \pm 0.3$. We suggest that the phenomena revealed here support a scenario of hierarchical star formation. Comparisons with other star-forming regions and galaxies are also discussed.

Subject headings: galaxies: star clusters – infrared: stars – Magellanic Clouds – stars: formation

1. INTRODUCTION

It has been suggested that star formation is hierarchical as revealed by stellar structures spanning a wide range of scales, from large stellar complexes and aggregates to small associations and clusters (Elmegreen et al. 2000). Hierarchical stellar structures have been investigated for a few famous Galactic star-forming regions, such as Taurus, Orion, Ophiuchus, etc. (Gomez et al. 1993; Larson 1995; Simon 1997); beyond the regime of binary and multiple systems, the young stars in these regions are found to be hierarchically clustered up to several parsecs. Schmeja et al. (2009) and Gouliermis et al. (2014) studied the star-forming region, NGC 345, in the Small Magellanic Cloud (SMC); specifically, bimodal clustering is found, with both a centrally-concentrated component and a self-similar, hierarchical distribution. On larger scales, a number of galaxies are shown to have hierarchically distributed young stars from several parsecs to kiloparsecs (e.g. Bastian et al. 2007, 2009; Gieles et al. 2008; Elmegreen & Elmegreen 2001; Elmegreen et al. 2006, 2014; Gouliermis et al. 2010, 2015). These studies show that hierarchical stellar structures, significant for a high degree of substructure and an absence of characteristic scales, exhibit self-similar and fractal properties. These features are also found for

the structures formed by the interstellar medium (ISM), such as clumps and filaments on various scales (e.g. Elmegreen & Falgarone 1996; Heyer et al. 2001; Roman-Duval et al. 2010). This similarity suggests that the hierarchical structures of the young stars may originate from those of their parental ISM, which are in turn related to turbulence (Larson 1981; Elmegreen 2008; Federrath et al. 2009; Girichidis et al. 2012), agglomeration with fragmentation (Carlberg & Pudritz 1990; McLaughlin & Pudritz 1996), and self-gravity (de Vega et al. 1996; see also the review of Elmegreen et al. 2000). After birth, the hierarchical stellar structures evolve toward uniform distributions, with the substructures becoming eliminated in roughly one crossing time (Gieles et al. 2008; Bastian et al. 2009; Gouliermis et al. 2015).

30 Doradus (30 Dor hereafter, also known as the Tarantula Nebula) in the Large Magellanic Cloud (LMC), with R.A. (J2000) = 05^h38^m38^s, Dec. (J2000) = $-69^{\circ}05'42''$, is among the most famous astronomical objects. It is by far the most luminous and massive star-forming region in the Local Group (Kennicutt & Hodge 1986) and hosts at its center the star cluster NGC 2070, which in turn is a collective of dense subclusters (Walborn & Blades 1997; Sabbi et al. 2012). On the other hand, three active star-forming regions, N158, N159, and N160, are

found to the south of 30 Dor (Nakajima et al. 2005; Galametz et al. 2013). They are located in the northern part of the LMC’s remarkable “molecular ridge”, which stretches ~ 2 kpc in a nearly north-south straight line and contains almost one third of the LMC’s molecular gas traced by CO (Cohen et al. 1988; Kutner et al. 1997; Johansson et al. 1998; Fukui et al. 2008; Ott et al. 2008). **Detailed studies of the hierarchical stellar structures in these star-forming regions have not been reported. Although Bastian et al. (2009) have investigated this issue for the entire LMC, their stellar sample did not reach the necessary resolution to reveal the substructures of these regions. This paper thus aims to carry out such a study for the 30 Dor-N158-N159-N160 star-forming complex based on young upper main-sequence (UMS) stars, which have been carefully selected from the VISTA Survey of the Magellanic Clouds (VMC; Cioni et al. 2011). We use dendrograms to illustrate the “parent–child” relations of star groups and subgroups identified on various density levels; and additionally the group size distribution is explored. This will demonstrate the hierarchical structures formed by the young UMS stars, which may indicate a scenario of hierarchical star formation in this complex.**

This paper is organized as follows. Section 2 describes the VMC survey and the data used in this work. Selection of UMS stars is outlined in Section 3, while in Section 4 we show their spatial distributions. We identify groups of UMS stars and reveal their “parent–child” relations using dendrograms in Section 5, and the group sizes are investigated in Section 6. We finally close this paper with a summary and conclusions.

2. THE VMC SURVEY AND DATA

Data used in this work come from the VMC survey (Cioni et al. 2011), which is carried out with the 4 m Visible and Infrared Survey Telescope for Astronomy (VISTA; Sutherland et al. 2015). The VMC survey is a multi-epoch, uniform, and homogeneous photometric survey of the Magellanic system performed in the near-infrared Y , J , and K_s filters (centered at $1.02 \mu\text{m}$, $1.25 \mu\text{m}$, and $2.15 \mu\text{m}$, respectively). On completion, the survey is expected to cover $\sim 184 \text{ deg}^2$ of the Magellanic system. A sequence of six offset positions, each one called a *paw-print*, is used to fill in the gaps between the 16 detectors of the VISTA infrared camera (VIRCAM; Dalton et al. 2006); the combined image, or a VISTA *tile*, covers an area of $\sim 1.5 \text{ deg}^2$. Each tile in the VMC survey is designed to be observed at 3 epochs in the Y and J bands, and 12 epochs in the K_s band. The total exposure times are 2400 s in the Y and J bands, and 9000 s in the K_s band for most regions of each tile, except for the tile edges and some areas of extra overlap among the detectors. The detectors usually have saturation limits of $Y = 12.9 \text{ mag}$, $J = 12.7 \text{ mag}$, and $K_s = 11.4 \text{ mag}$, and the stacked images from all epochs can provide sources with typical 5σ limiting magnitudes of $Y = 21.9 \text{ mag}$, $J = 22.0 \text{ mag}$, and $K_s = 21.5 \text{ mag}$; but the saturation limits and photometric depths also depend on sky conditions and crowding (Cioni et al. 2011; Tatton et al. 2013). This depth reaches the oldest main-sequence (MS) turn-offs in both the LMC and SMC (Kerber et al. 2009).

The 30 Dor-N158-N159-N160 star-forming complex is analyzed based on tile LMC 6.6 (see Cioni et al. 2011 for tile definitions). We retrieved the data of tile LMC 6.6 as part of VMC Data Release 3 from the VISTA Science Archive (VSA). The details of the VSA and the VISTA data flow pipeline can be found in Cross et al. (2012) and Irwin et al. (2004), respectively. We have used point-spread-function (PSF) photometry from Rubele et al. (2012) instead of aperture photometry to reduce the influence of source crowding (Tatton et al. 2013). Rubele et al. (2012) generated deep, PSF-homogenized tiles by stacking images from different epochs, each one formed by the combination of six pawprints. They performed PSF photometry on the deep tiles and estimated the photometric errors and local completeness using artificial star tests. Specifically, the local completeness is estimated in a ring of radius 0.025 deg around each source and in bins of $\pm 0.05 \text{ mag}$ separately in the Y , J , and K_s bands. We use their local completeness estimates to assign weights to stars in our sample (see Sections 3 and 4). A detailed description of the PSF homogeneity and photometry can also be found in Rubele et al. (2015).

The top half of VISTA’s detector 16 has varying quantum efficiency on short timescales, making accurate flat fielding impossible. This leads to worse signal-to-noise ratios and unreliable data in regions covered by that detector (Rubele et al. 2012; Tatton et al. 2013). For tile LMC 6.6, the region affected is the southwest corner with R.A. (J2000) $< 05^{\text{h}}33^{\text{m}}55^{\text{s}}$ and Dec. (J2000) $< -69^{\circ}43'48''$. This region is excluded from our analysis.

3. THE UPPER MAIN SEQUENCE SAMPLE

The left-hand panel of Fig. 1 shows the $(J - K_s, K_s)$ CMD of sources in tile LMC 6.6 without correcting for extinction. The MS, red-giant branch (RGB), and red clump (RC) populations are all resolved. We use the $(J - K_s)$ color instead of $(Y - K_s)$ because the former suffers less from interstellar extinction. Still, effects of interstellar extinction varying across the tile are present, evidenced by the width of the MS and the RGB, and the elongated tail associated with the RC toward fainter magnitudes and redder colors.

We estimate the extinction by taking advantage of the extinction values of $\sim 1.5 \times 10^5$ RC stars in this tile provided by Tatton et al. (2013). We bin both all stars and their RC stars into spatial grids with a grid size of $1' \times 1'$; the median extinction value of RC stars is taken as the extinction estimate for all stars in each bin. We use the extinction coefficients, $A_J = 0.283 \times A_V$ and $A_{K_s} = 0.114 \times A_V$, computed from the Cardelli et al. (1989) extinction curve with $R_V = 3.1$ (Girardi et al. 2008). **De Marchi et al. (2016) reported a new extinction law for 30 Dor, characterized by $R_V \sim 4.5$. However, using this law would not lead to significant changes in the corrections for the near-infrared wavelengths considered in this paper (Gordon et al. 2003).** Contamination by RGB stars is not important, since their intrinsic color exhibits only a small difference of $\sim 0.05 \pm 0.02 \text{ mag}$ from that of the RC stars (Tatton et al. 2013), thus introducing very minor, if any, errors in correcting the CMD. The middle panel of Fig. 1 shows the CMD after correcting for extinction. The MS and RGB become tighter and the

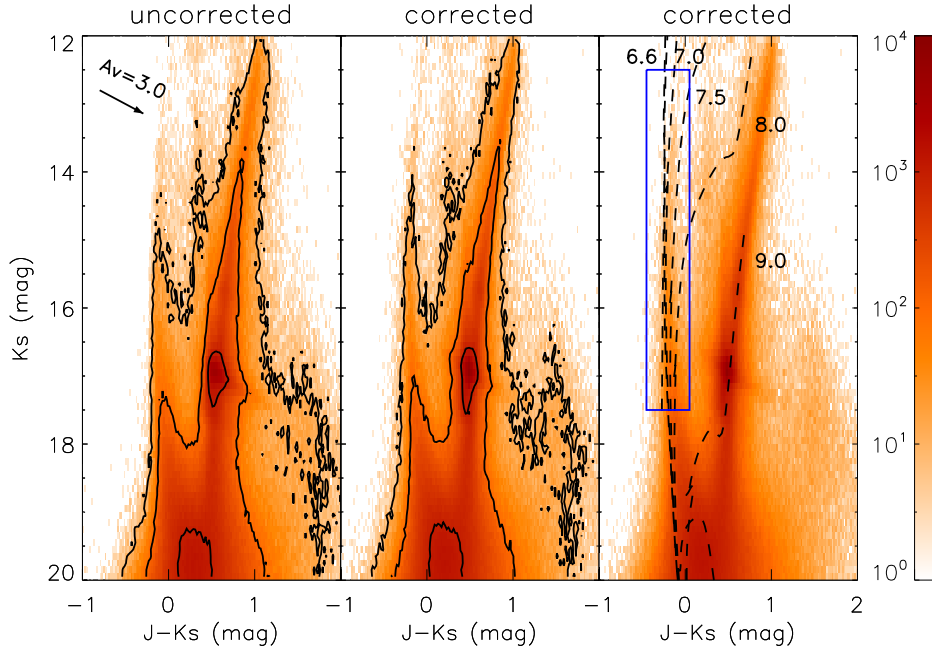


FIG. 1.— $(J - K_s, K_s)$ color-magnitude Hess diagrams of stars in tile LMC 6.6. In the left-hand panel, interstellar extinction has not been corrected for, while in the middle and right-hand panels, interstellar extinction has been corrected for according to Tattton et al. (2013). The color scales in all three panels show the number of stars in each color-magnitude bin; the bin size is 0.02 mag in color and 0.10 mag in magnitude. The contours in the left-hand and middle panels correspond to levels of 10^1 , 10^2 , and 10^3 stars per bin from the outside to the inside. The arrow in the left-hand panel shows the reddening vector corresponding to $A_V = 3.0$ mag. In the right-hand panel, the blue solid box shows the selection criterion of UMS stars and dashed lines are the MS to RGB sections of PARSEC (version 1.2S) isochrones of metallicity $Z = 0.008$ and ages $\log(\tau/\text{yr}) = 6.6, 7.0, 7.5, 8.0,$ and 9.0 , shifted by a distance modulus of $(m - M)_0 = 18.49$ mag. Offsets of 0.026 mag in J and 0.003 mag in K_s , given by Rubele et al. (2015), have been subtracted from the isochrone magnitudes to correct for the small differences between the VISTA system and the model Vega system. The color scale in the right-hand panel is the same as that in the middle panel; we overplot the contours, sample selection criterion, and isochrones in the different panels simply for clarity.

elongated tail associated with the RC appears less prominent.

The extinction in the LMC is subject to a population dependence, which is, however, difficult to model (Zaritsky et al. 2004; Cignoni et al. 2015). Compared with intermediate-age RC stars, the young stars form in gas-rich regions with high extinction, but their ionizing fluxes and stellar winds tend to reduce extinction by evacuating the surrounding material. The problem is further compounded by their different spatial distributions with respect to the dust. Along the line of sight, dust may lie in the foreground and/or background of a star, and only the foreground dust contributes to its extinction. Perpendicular to the line of sight, dust in star-forming regions also exhibits a high degree of spatial variability (e.g. Lombardi et al. 2010, 2011; Cambr esy et al. 2013). **This is also demonstrated by the *Herschel* dust emission map across the SMC star-forming region NGC 345 (also known as N66), which is highly variable on very small scales and rich in structures of knots, arcs and filaments, etc. (Hony et al. 2015).** It is very challenging and beyond the scope of this work to take into account all these considerations by obtaining a population-dependent, three-dimensional extinction map with high angular resolution. Still, our treatment of extinction is reasonable as the features become more well-defined in the extinction-corrected CMD.

In the right-hand panel of Fig. 1, we overplot PARSEC isochrones (version 1.2S; Bressan et al. 2012) of

metallicity $Z = 0.008$ and ages $\log(\tau/\text{yr}) = 6.6, 7.0, 7.5, 8.0,$ and 9.0 , shifted by an LMC distance modulus of $(m - M)_0 = 18.49$ mag (Pietrzyński et al. 2013; de Grijs et al. 2014) and assuming zero extinction. We have subtracted offsets of 0.026 mag in J and 0.003 mag in K_s , given by Rubele et al. (2015), from the isochrones to correct for the small differences between the VISTA system and the model Vega system. The differences are caused by the unaccounted for non-linearities in the color-color relation used to calibrate VISTA’s photometric zeropoint with 2MASS stars (for details, see Rubele et al. 2012, 2015). It is clear that the upper part of the MS corresponds to a very young population, since it coincides well with isochrones of $\log(\tau/\text{yr}) = 6.6$ and 7.0 . We thus select a sample of UMS stars using color and magnitude cuts, $-0.45 \leq J - K_s \leq 0.05$ mag and $12.5 \leq K_s \leq 17.5$ mag, which are shown as the blue solid-line box in the right-hand panel of Fig. 1. This sample contains $\sim 1.5 \times 10^4$ stars.

We confirm that the UMS stars in our sample comprise indeed a young population by comparing their K_s -band luminosity function with synthetic ones (Fig. 2). For the observed UMS sample, we take into account the photometric completeness by assigning to each star a weight $w = 1/f$, where f is the local completeness; because we are using the J and K_s bands, f is estimated from the lowest value of the J - and K_s -band local completenesses, f_J and f_{K_s} , as estimated by Rubele et al. (2012, 2015; Section 2). On the other hand, we cre-

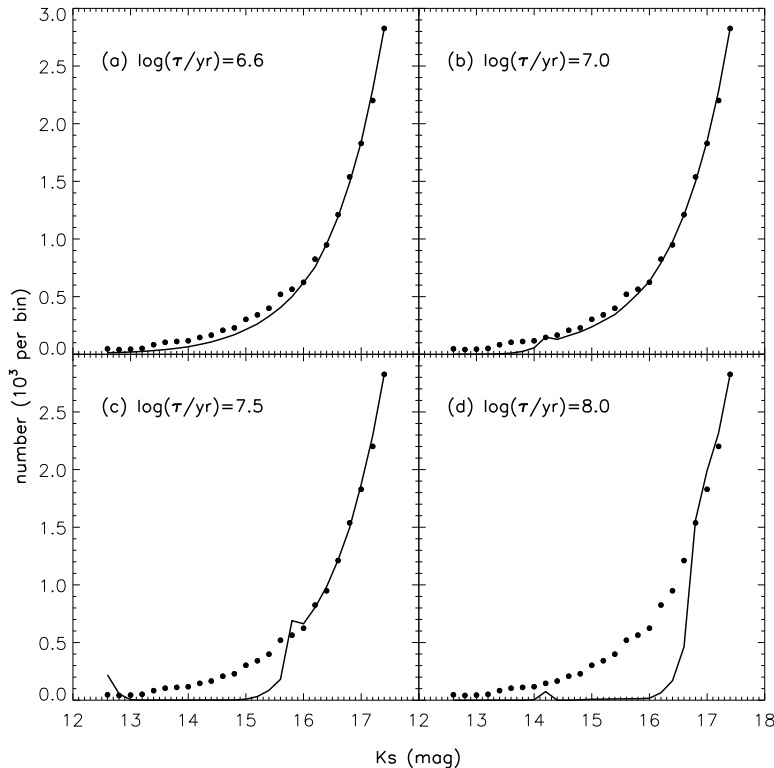


FIG. 2.— K_s -band luminosity functions of the observed UMS sample (dots) and synthetic ones (lines) with ages $\log(\tau/\text{yr}) = 6.6, 7.0, 7.5,$ and 8.0 (see the text for details). The observed UMS sample has been corrected for interstellar extinction and local completeness. The luminosity functions of the synthetic samples are normalized so that they have equal numbers of stars in the faintest magnitude bin as the observed sample. The two younger synthetic UMS samples have luminosity functions quite similar to the observed one, while the two older synthetic UMS samples have significantly fewer bright stars because they have evolved away from the MS. The Poissonian error bars of the observed luminosity function are smaller than or comparable to the symbol size.

ate four synthetic populations of metallicity $Z = 0.008$ and ages $\log(\tau/\text{yr}) = 6.6, 7.0, 7.5,$ and 8.0 . We assume a Chabrier (2001) stellar initial mass function (IMF), a 30% binary fraction, and that the binaries are non-interacting systems with primary/secondary mass ratios evenly distributed from 0.7 to 1.0 (Kerber et al. 2009; Rubele et al. 2012, 2015). In each synthetic population, sufficiently large numbers of stars are created and placed at the LMC’s distance modulus of 18.49 mag (Pietrzyński et al. 2013; de Grijs et al. 2014). Synthetic UMS samples are obtained by applying the same selection criterion to the synthetic populations as to the real data; their luminosity functions are then calculated.

Figure 2 shows that the luminosity functions of the $\log(\tau/\text{yr}) = 6.6$ synthetic UMS sample and the observed one are consistent, although the latter seems to contain slightly more stars with $K_s < 16.0$ mag (Fig. 2a); the $\log(\tau/\text{yr}) = 7.0$ synthetic sample also has a very similar luminosity function, except for a slight lack of stars brighter than $K_s = 14.0$ mag (Fig. 2b). Synthetic samples with $\log(\tau/\text{yr}) = 7.5$ and 8.0 deviate from the observed sample even more significantly, in that they have many fewer stars brighter than $K_s = 15.7$ mag or 16.8 mag, respectively (Fig. 2c and 2d). The comparison leads us to suggest that the observed UMS sample is indeed composed of very young stars with ages of a few million years.

The VLT-FLAMES Tarantula Survey (VFTS; Evans et al. 2011) has collected fiber spectroscopy of about 800 early-type stars in a $25'$ diameter

field of the 30 Dor region, and spectroscopic classification of 352 O–B0 and 438 B-type stars has been done by Walborn et al. (2014) and Evans et al. (2015), respectively. We compare our UMS sample with their OB-type stars; 279 O-type and 313 B-type stars observed by VFTS are covered by our UMS sample, and their spectral types range from O2 (earliest) to B5 (latest). This comparison again indicates the youth of the UMS sample.

Contamination by foreground Galactic stars and background galaxies is minimal. The former reside at typically $(J - K_s) = 0.7$ mag, and background galaxies have colors even redder than this value. Thus, they do not overlap with the UMS stars in the CMD and are not included by the color–magnitude selection criterion.

4. SPATIAL DISTRIBUTIONS

Figure 3 shows the surface density, Σ , of the UMS sample in tile LMC 6.6. We estimate the surface density using two methods. The first method is simple star counts in spatial grids of $1' \times 1'$, with each star assigned a weight in the same way as described in Section 3. The second is the so-called n^{th} nearest-neighbor method (Casertano & Hut 1985; Megeath et al. 2016), and surface densities are calculated locally around each UMS star, Schmeja (2011) compared different algorithms in constructing stellar density maps and identifying star clusters as overden-

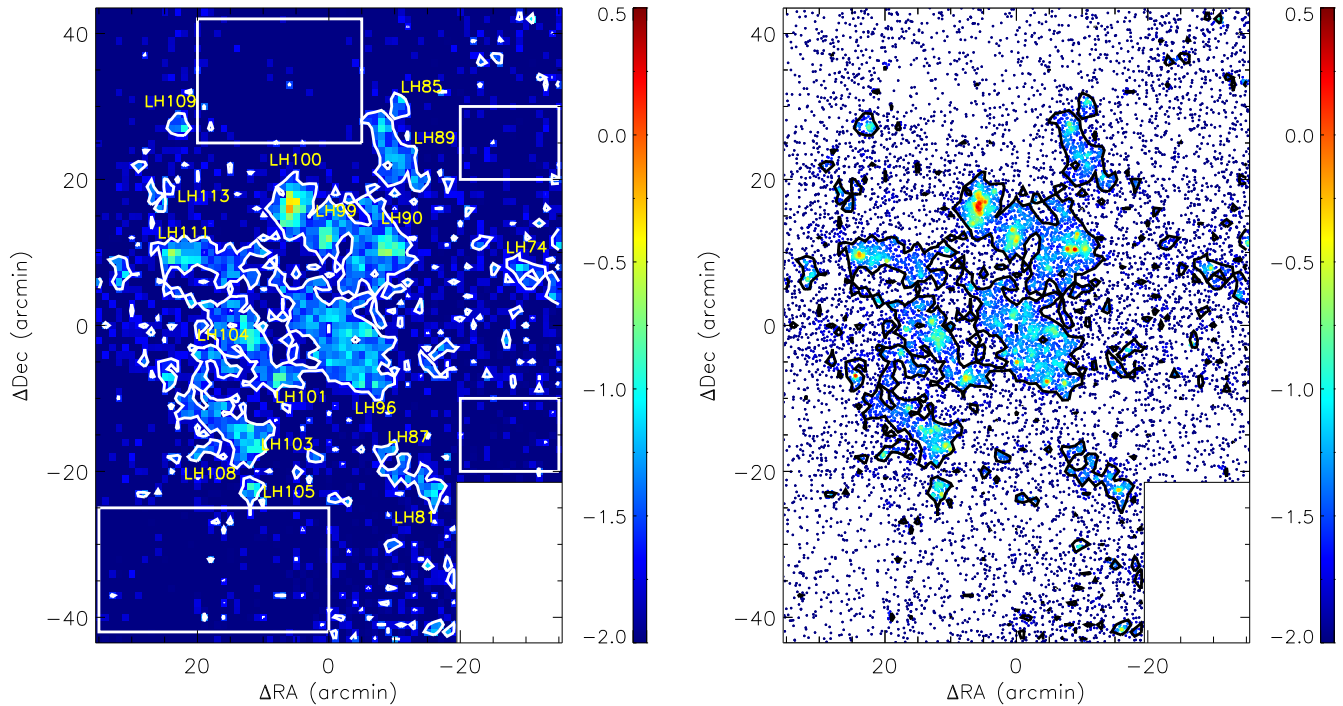


FIG. 3.— Surface-density map of the UMS sample. In the left-hand panel, surface densities are derived from simple star counts in grids of $1' \times 1'$, while in the right-hand panel, surface densities are calculated locally around each UMS star with the n^{th} nearest-neighbor method ($n = 10$; see the text for details). Surface-density values are shown as a color scale, according to the color bar in units of logarithmic number of stars pc^{-2} . The four white boxes in the left-hand panel are the regions where we estimate the LMC background, and the level corresponding to 3σ above the average of the background is overlaid as white and black contours in the left- and right-hand panels, respectively. **The contours are computed only from the surface densities derived with simple star counts (i.e., the left-hand panel), but the same contours are also overlaid on the right-hand panel to indicate the large-scale stellar distributions.** The yellow annotations in the left-hand panel label the associations catalogued by Lucke & Hodge (1970). The (0, 0) position corresponds to R.A. (J2000) = $05^{\text{h}}37^{\text{m}}40^{\text{s}}$, Dec. (J2000) = $-69^{\circ}22'18''$. The southwest blank corner corresponds to the region affected by worse signal-to-noise ratios associated with detector 16.

sities; he showed that clusters with low overdensities or highly hierarchical structure are only reliably identified by methods with inherent smoothing, such as simple star counts and the n^{th} nearest-neighbor method (Casertano & Hut 1985; Megeath et al. 2016). Thus, in this section we apply these two methods to the UMS sample to estimate their surface densities; the surface-density map derived with the n^{th} nearest-neighbor method is then used to identify stellar groups in Section 5. In the first method, we use a spatial grid of $1' \times 1'$ cells and determine the total weight of stars in each bin, with each star assigned a weight in the same way as described in Section 3. In the second method, surface densities are calculated around each star with

$$\Sigma_n = \frac{w_{\text{tot}}}{\pi r_n^2}, \quad (1)$$

where r_n is the distance to the n^{th} nearest neighbor, and w_{tot} is the sum of the weights of all $n+1$ UMS stars at a radius of $r \leq r_n$:

$$w_{\text{tot}} = (w_0 - 1) + \sum_{i=1}^{n-1} w_i + 0.5 \times (w_n - 1). \quad (2)$$

Since we measure the density inside the annulus between the central star and its n^{th} nearest neighbor, Casertano & Hut (1985) argued that the n^{th} nearest-neighbor method provides an unbiased estimate of density if neither the central star nor its n^{th} nearest neighbor are counted. In other words, we take into account the total weight inside the annulus between the central star and its n^{th} nearest neighbor. Thus, unity is subtracted from the central star, and its excess weight above unity, i.e. $w_0 - 1$, represents the over-density inside the annulus; similarly, 1 is subtracted from its n^{th} nearest neighbor, and half of its excess weight above 1, i.e. $0.5 \times (w_n - 1)$, is assumed to be inside the annulus. We use $n = 10$, which has been shown to be a reasonable choice to balance statistical noise and locality (Gouliermis et al. 2010; Megeath et al. 2016). The separations between stars and their 10^{th} nearest neighbors range from $0.1'$ in the densest area, to $3.6'$ in low-surface-density regions, with a median value of $0.9'$. **Surface-density maps derived with the two methods are shown in Fig. 3.**

Figure 3 shows that the surface densities are generally low in the outer part of the tile. The UMS stars in the low-surface-density areas could come from the LMC's general background field, which are not associated with

the star-forming complex. As we will show later in this paper, the young stars are hierarchically grouped in a way that small and compact stellar structures reside in larger and looser ones. Thus, there is also a possibility that the UMS stars in the low-surface-density areas belong to a larger stellar structure created in the star-forming process, but our sky coverage is not large enough to reveal it. In this work we do not try to reach scales larger than a single tile, thus we characterize this low surface density only from the LMC’s background contribution. This should not make significant differences to our results, considering its very small values compared with the higher surface densities of the star-forming complex. We estimate the background level using four subregions, which are devoid of surface-density enhancements and outlined with white boxes in the left-hand panel of Fig. 3. We calculate the average and standard deviation of the surface densities derived with simple star counts of the bins in the subregions, iteratively rejecting 3σ outliers. The average value is $\log(\bar{\Sigma}\cdot\text{pc}^2) = -2.30$ and the standard deviation is $\log(\sigma\Sigma\cdot\text{pc}^2) = -2.28$. The former is used as a measure of the background level, and the latter is an estimate of its fluctuations.

High surface densities are found in the central part of the tile, with roughly $-20' \leq \Delta\text{R.A.} \leq 25'$, $-20' \leq \Delta\text{Dec.} \leq 30'$. The peak surface density, which occurs in the center of 30 Dor with $\Delta\text{R.A.} = 6'$ and $\Delta\text{Dec.} = 17'$, reaches $\log(\Sigma\cdot\text{pc}^2) \sim -0.20$, derived from simple star counts, or $\log(\Sigma\cdot\text{pc}^2) \sim 0.25$, obtained with the n^{th} nearest-neighbor method, where $n = 10$. The discrepancies between results obtained with the two methods arise from their different smoothing scales. The star count method has a fixed smoothing scale of $1'$; the n^{th} nearest-neighbor method, in contrast, has variable smoothing scales, which are the separations between stars and their n^{th} nearest neighbors. As previously mentioned, the separations range from $0.1'$ to $3.6'$ for $n = 10$, depending on the surface densities.

Also note that the majority of UMS stars are distributed in a number of highly fractured groups. The largest ones, with scales of several to $\sim 20'$, are consistent with the OB associations catalogued by Lucke & Hodge (1970), whose designations have been labeled in the figure. **Here we recall that it can be difficult to define clusters, associations, or larger structures, and equally difficult to make clear distinctions between them (Bressert et al. 2010; Gieles & Portegies Zwart 2011; Krumholz 2014).** Here we recall that it can be difficult to define clusters, associations, or larger structures, and equally difficult to make clear distinctions between them (Bressert et al. 2010; Krumholz 2014). Gieles & Portegies Zwart (2011) suggested that bound and unbound stellar systems can be distinguished by comparing the ages and crossing times; yet observationally it is not trivial to obtain these parameters. In order to avoid ambiguity, we will refer to any stellar structure with a surface density enhancement as a group, regardless of its size or dynamical state; the critical surface density to define an “enhancement” is not fixed but can be varied arbitrarily, and by doing this we will reveal the hierarchy of stellar groups on a range of scales in Sections 5 and 6. But here we temporarily

set the critical surface density at 3σ above the average of the background level, so that we can investigate the stellar distributions on large scales. This corresponds to $\log(\Sigma_{\text{cr}}\cdot\text{pc}^2) \sim -1.7$, which is shown as white or black contours in both panels of Fig. 3.

Most prominently in Fig. 3, the 30 Dor Nebula contains three populous groups, LH 100, LH 99, and LH 90 from its center to the west. In the vicinity of 30 Dor, there are three large, diffuse stellar groups, LH 89, LH 111, and LH 96, to its northwest, east, and south, respectively. In the southeastern part of the tile, there are another few groups with significant stellar overdensities, i.e., LH 101, LH 104, LH 103, and LH 105, which correspond to the N158, N160, and N159 star-forming regions (Bolatto et al. 2000; Nakajima et al. 2005; Galametz et al. 2013). A number of smaller groups can also be seen across the tile, e.g. LH 74 in the west, and even more that have not been catalogued by Lucke & Hodge (1970).

Many of the groups described above are, in turn, highly sub-grouped. For example, LH 96 is far from centrally concentrated and contains multiple surface density peaks. This has been noted by Lucke & Hodge (1970), who identify four smaller associations inside, i.e. LH 93, LH 94, LH 97, and LH 98. Similarly, the star-forming region, N158, contains two surface density peaks, LH 101 and LH 104, as catalogued by Lucke & Hodge (1970) and also revealed here. Although not obvious in Fig. 3, LH 100 contains two star clusters, NGC 2070 at the center of 30 Dor and Hodge 301 slightly to the northwest (Cignoni et al. 2015); the former is, in turn, a collective of smaller star clusters (Walborn & Blades 1997; Sabbi et al. 2012). A statistical analysis of the hierarchical structures is given below.

5. GROUP IDENTIFICATION AND DENDROGRAMS

In this section, we investigate the hierarchy of UMS star groups through the so-called *dendrograms* (Houllahan & Scalo 1992; Rosolowsky et al. 2008; Gouliermis et al. 2010). Dendrograms are structure trees showing the “parent-child” relations of groups with surface density enhancements over different critical values. **We vary the critical surface density from $\log(\Sigma_{\text{cr}}\cdot\text{pc}^2) = -1.7$ to -0.1 in steps of 0.2. The lower limit corresponds to 3σ above average of the background level, and the upper limit is chosen close to the peak surface density (see Section 4). We vary the critical surface density from $\log(\Sigma_{\text{cr}}\cdot\text{pc}^2) = -1.7$ to -0.1 ; the lower limit corresponds to 3σ above average of the background level, and the upper limit is chosen close to the peak surface density (see Section 4). Between them we define levels of critical surface densities with equal logarithmic steps of 0.2 dex. This step is arbitrarily chosen in order to avoid too many levels in displaying the dendrograms, considering that the surface densities span two orders of magnitude; reducing this step and increasing the number of levels would not change the conclusion reached in this section.** On each level we use a friends-of-friends group identification algorithm (Megeath et al. 2016). To identify stellar groups, we use the surface densities estimated with the n^{th} nearest-neighbor method ($n = 10$; Section 4). For any star with $\Sigma_n \geq \Sigma_{\text{cr}}$, its 10 nearest neighbors which also show $\Sigma_n \geq \Sigma_{\text{cr}}$ are its

friends; friends of its friends are also friends, and so on. To avoid spurious detections, we require additionally that each group should contain at least N_{\min}^* UMS stars. Smaller values of N_{\min}^* would lead to more spurious detections, but larger values of N_{\min}^* would miss small groups which are not massive enough to contain sufficient number of UMS stars at the high-mass end of their IMF. **In the following analysis, we use three different values of $N_{\min}^* = 5, 10, \text{ and } 15$ to assess the effect of changing N_{\min}^* . As we shall see in Section 6, this effect is significant for groups of sizes smaller than 8 pc. In this section, however, this effect does not affect the conclusions reached; thus we arbitrarily set $N_{\min} = 10$ in the following analysis.** The group identification process is repeated on each level of critical surface density, and we keep track of the connections between “parent” groups on lower levels and “child” groups on higher levels.

The groups found at the first four levels with $\log(\Sigma_{\text{cr}} \cdot \text{pc}^2) = -1.7, -1.5, -1.3, \text{ and } -1.1$ are shown in Fig. 4. **Here we only show the results obtained with $N_{\min}^* = 10$. With smaller N_{\min}^* there are more smaller-sized groups and vice versa, but the conclusion reached below is not altered.** To help distinguish different groups, each one is enclosed with its convex hull, shown as thick-lined polygons in the figure; 39 groups are found on the lowest level with $\log(\Sigma_{\text{cr}} \cdot \text{pc}^2) = -1.7$. Most prominent are the two large groups, labeled “A” and “B” in Fig. 4a, residing in the central part of the tile. Group A consists of the 30 Dor complex, LH 100-LH 99-LH 90, and the diffuse LH 96, LH 111 to its south and east, which have already been mentioned in Section 4. On the other hand, Group B is composed of LH 101-LH 104 and LH 103 in the star-forming regions N158 and N160, respectively. 48 groups are found on the second level with $\log(\Sigma_{\text{cr}} \cdot \text{pc}^2) = -1.5$, and the groups are generally smaller in size compared with those on the first level. Group A now splits into two large groups, corresponding to LH 100-LH 99-LH 90 and LH 96, a smaller LH 111 in the east, and several even smaller ones among them. Group B also splits, with LH 101, LH 104, LH 103, and many other groups gaining their independence. 45 and 40 groups are found on the third and fourth levels, respectively. Results on higher levels are not shown here, but similar processes are continuously repeated, i.e., larger groups vanish and smaller groups emerge.

To better display the relations between groups on different levels, we show all the convex hulls of groups found on the first four levels in Fig. 5, and the dendrograms of Group A and Group B in Fig. 6. One can, of course, construct a single dendrogram containing all groups found on all levels. Here we show only the dendrograms starting from Group A and Group B, which are its “child dendrograms”. Both Fig. 5 and Fig. 6 show that the stellar groups on lower levels very frequently split up into several smaller ones on higher levels. Thus, the stellar structures display a high degree of sub-grouping. This result agrees well with previous works, which also reveal hierarchically sub-grouped young stellar structures in individual star-forming regions (e.g. Kirk & Myers 2011) (e.g. Schmeja et al. 2009; Kirk & Myers 2011; Gouliermis et al. 2014) or over the whole galaxies (e.g. Gouliermis et al. 2010, 2015; Gusev 2014).

6. GROUP SIZES

The stellar groups identified in Section 5 have a wide range of sizes, from the largest Groups A and B, which are more than 100 pc in extent, to the smallest ones on parsec scales. Fig. 7 shows their cumulative size distribution, which includes stellar groups found on all levels of critical surface densities. The size of a group is estimated using

$$R = \frac{\sum_i |x_i - x_c| w_i}{\sum_i w_i}, \quad (3)$$

where x_i is the position vector of the i^{th} group member, w_i its weight assigned in the same way as described in Section 3 to account for the local photometric completeness, and x_c the group center, which is calculated as

$$x_c = \frac{\sum_i x_i w_i}{\sum_i w_i}. \quad (4)$$

The cumulative size distribution is approximately a power law between ~ 10 pc to > 100 pc, with significant flattening at several parsecs. **As mentioned in Section 5, stellar groups are identified on levels of critical surface densities spaced by equal logarithmic steps of 0.2 dex; with smaller steps and thus more levels, more stellar groups can be included in obtaining the cumulative size distribution. We carried out a test by reducing this step to 0.1 and 0.005 dex and repeating the group identification process; the shapes of the resultant cumulative size distributions are almost unchanged (not shown), suggesting that this does not affect the conclusion reached in this section. On the other hand,** the flattening is due to incompleteness in our group-identifying algorithm. As mentioned in the previous section, smaller groups may be missed if they contain fewer UMS stars than the required minimum number of stars, N_{\min}^* . Thus, we have calculated the cumulative size distribution multiple times for stellar groups identified with $N_{\min}^* = 5, 10, \text{ and } 15$. As expected, fewer small-sized groups are identified with increasing values of N_{\min}^* . Beyond $R = 8$ pc, however, the cumulative size distribution remains almost unaltered. Thus we suggest that the stellar groups found by our group-identifying algorithm are complete above $R = 8$ pc.

In the range not affected by incompleteness, the cumulative size distribution can be well fitted with a single power law with power-law index $\alpha = -1.6 \pm 0.3$ (fitting error). This indicates a scale-free behavior and agrees with the scenario of hierarchical star formation (Elmegreen et al. 2000). On the other hand, the cumulative size distribution for substructures inside a fractal follows

$$N(> R) \propto R^{-D}, \quad (5)$$

where N is the number of substructures with sizes larger than R , and D is the fractal dimension (Mandelbrot 1983; Elmegreen & Falgarone 1996). Thus the UMS stars have a fractal dimension of $D_2 = 1.6 \pm 0.3$, by comparing the power-law cumulative group size distribution with Equation 5. Note D_2 is the fractal dimension of the UMS stars’ *projections* on the two-dimensional plane perpendicular to the line of sight. **The UMS stars’ volume distributions have a fractal dimen-**

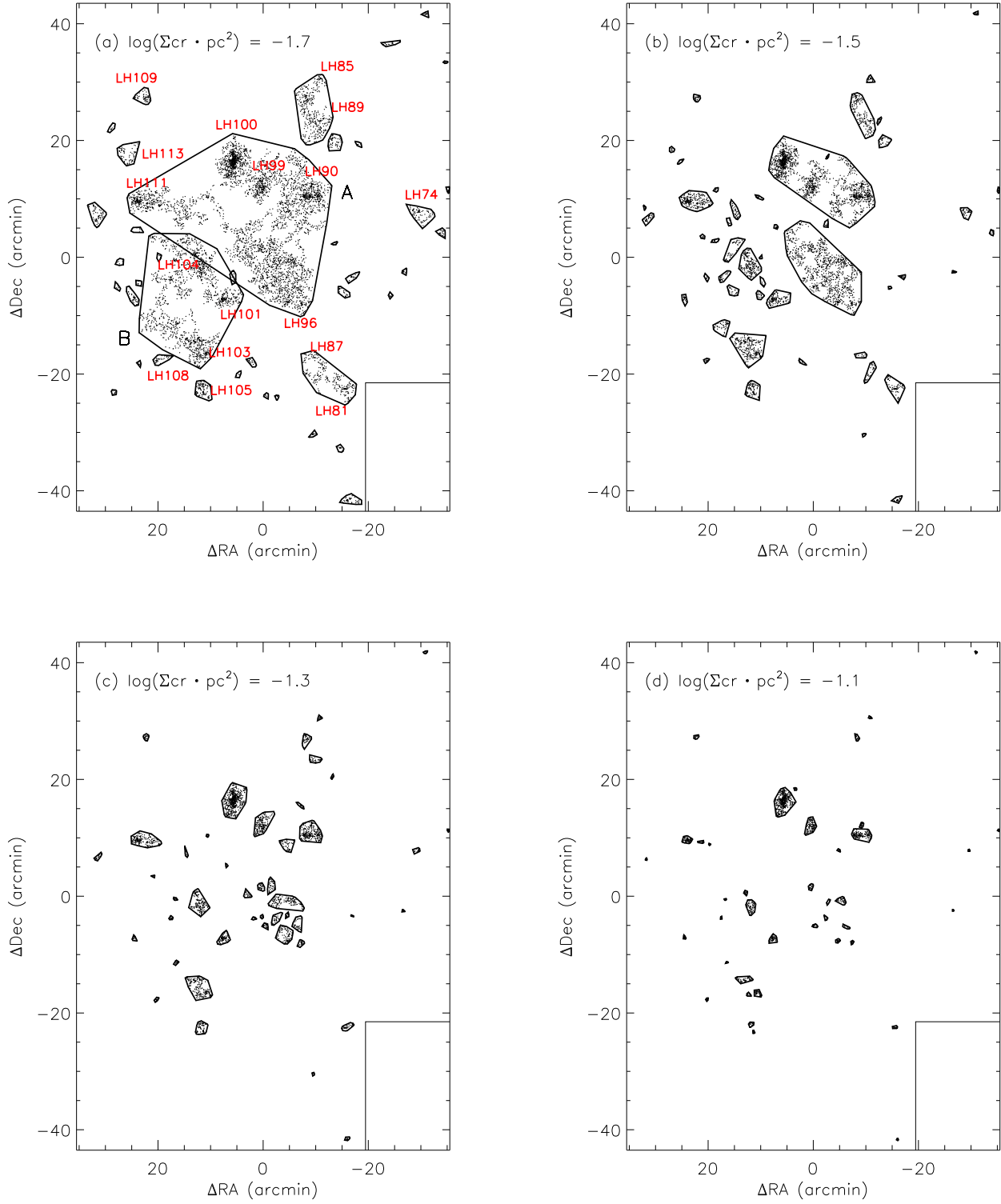


FIG. 4.— Groups of UMS stars identified with $N_{\text{min}}^* = 10$ and different critical surface densities labeled in each panel (see the text in Section 5 for details). Each group is enclosed by its convex hull to distinguish itself from others. The red annotations in the top left-hand panel label the associations catalogued by Lucke & Hodge (1970). The (0, 0) position corresponds to R.A. (J2000) = $05^{\text{h}}37^{\text{m}}40^{\text{s}}$, Dec. (J2000) = $-69^{\circ}22'18''$. The southwest blank corner corresponds to the region affected by worse signal-to-noise ratios associated with detector 16.

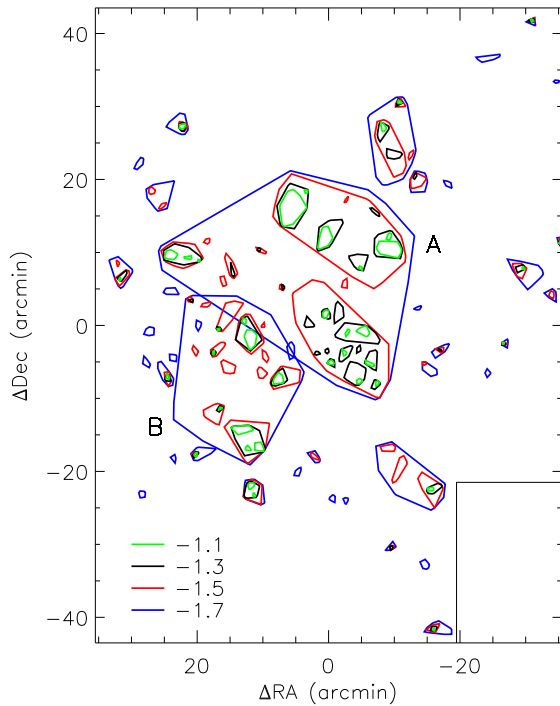


FIG. 5.— Convex hulls enclosing the groups found with $N_{\min}^* = 10$ and critical surface densities of $\log(\Sigma_{\text{cr}} \cdot \text{pc}^2) = -1.7$ (blue), -1.5 (red), -1.3 (black), and -1.1 (green). The $(0, 0)$ position corresponds to R.A. (J2000) = $05^{\text{h}}37^{\text{m}}40^{\text{s}}$, Dec. (J2000) = $-69^{\circ}22'18''$. The southwest blank corner corresponds to the region affected by worse signal-to-noise ratios associated with detector 16.

sion of $D_3 = 2.6 \pm 0.3$, according to the relation $D_3 = D_2 + 1$ given by Beech (1992). It is not easy to recover the fractal dimension of the UMS stars' volume distributions, D_3 , from D_2 . Beech (1992) suggested a relation $D_3 = D_2 + 1$; however, this relation holds only if the perimeter-area dimension of a projected three-dimensional structure is the same as that of a slice (Elmegreen & Scalo 2004), and moreover, this relation is often shown to be unrealistic for star-forming regions (e.g., Sánchez et al. 2005). Based on simulations, Gouliermis et al. (2014) proposed a more general conversion between D_3 and D_2 , which would suggest $D_3 = D_2 = 1.6$ for this case; but this may still remain uncertain, since this real star-forming complex can be different in its three-dimensional structures from, and maybe also more complicated than, their simulated fractals.

The ISM also exhibits hierarchical structures of clumps and filaments on all scales (Larson 1981; Elmegreen 2008). These structures also have power-law cumulative size distributions, with typical fractal dimensions close to $D_3 = 2.4$ (e.g., Elmegreen & Falgarone 1996; Heyer et al. 2001; Roman-Duval et al. 2010). If $D_2 = D_3 - 1 = 1.4$ holds for the ISM, this value is consistent with that of the UMS stars, within the errors. While quantitatively this may not be valid, still the structures formed by the UMS stars and ISM at least

qualitatively share very similar scale-free, power-law behaviors in the cumulative size distributions. This supports a scenario of hierarchical star formation, where the newly-born stars follow the the gas distribution (Elmegreen et al. 2000). For the origin of hierarchical structures in the ISM, the role of turbulence as the driving source is highlighted by recent studies (e.g., Federrath et al. 2009); as suggested by Elmegreen et al. (2000), however, this is also possibly accompanied by agglomeration with fragmentation (Carlberg & Pudritz 1990; McLaughlin & Pudritz 1996) as well as self-gravity (de Vega et al. 1996).

Beyond the regime of binary and multiple systems on very small scales, the young stars in a few Galactic star-forming regions are hierarchically clustered with fractal dimensions $D_2 = 1.4$, as derived by Larson (1995) for Taurus (see also Gomez et al. 1993), or $D_2 = 1.5$, as estimated by Simon (1997) for Taurus, Orion, and Ophiuchus. The projected fractal dimension we have derived agrees well with these studies. Gouliermis et al. (2014) reported bimodal stellar clustering for the SMC star-forming region, NGC 346; the auto-correlation function (ACF) of young stars is well described by a broken power law, with a shallower slope for scales smaller than ~ 5.8 pc and a steeper one for scales above this value. The slope of the latter part corresponds a projected fractal dimension of $D_2 = 1.4$, also consistent with our result. They have shown that the ACF is best explained by two distinct stellar clustering components, with a centrally concentrated one dominant in the central part of the region, and an extended, hierarchical distribution across the observed field. In the case of 30 Dor-N158-N159-N160, however, while there may be centrally concentrated, non-hierarchical components on small scales (e.g., star clusters), their existence is not revealed by the cumulative size distribution shown here.

Significant for the intense star formation, the 30 Dor-N158-N159-N160 complex has been proposed to arise from galaxy interactions. de Boer et al. (1998) suggested that the star formation in this area may have been induced by the bow-shock formed as the LMC moves through the Milky Way's halo. Alternatively, Fujimoto & Noguchi (1990) and Bekki & Chiba (2007) considered the interaction between the Magellanic Clouds. Despite these possible external influences, however, no significant difference in D_2 is found compared with the star-forming regions mentioned above, given the measurement uncertainties. It is possible that the energy injected by the external process(es), occurring on larger scales, have cascaded down to smaller scales via turbulence, clump agglomeration, and fragmentation, thus establishing the hierarchical structures in the ISM and in turn in the young stars formed within; during this process, any difference on larger scales may have been erased and lost on smaller scales, leading to very similar D_2 values for this star-forming complex.

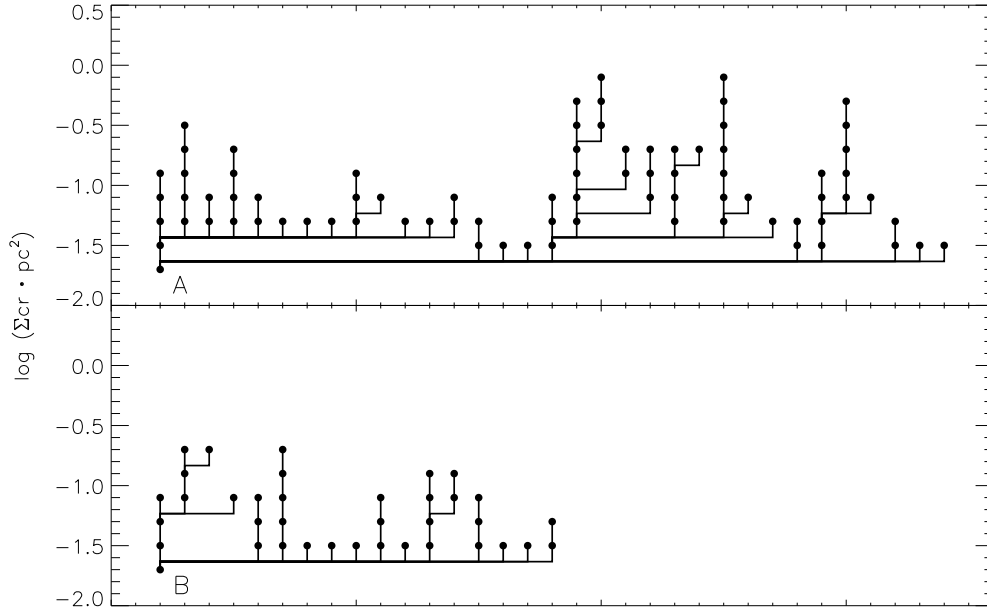


FIG. 6.— Dendrograms of Group A (top) and Group B (bottom), illustrating the groups found at different levels and their connections. This figure only shows the results obtained with $N_{\min}^* = 10$ in the group-identifying algorithm (see the text in Section 5 for details).

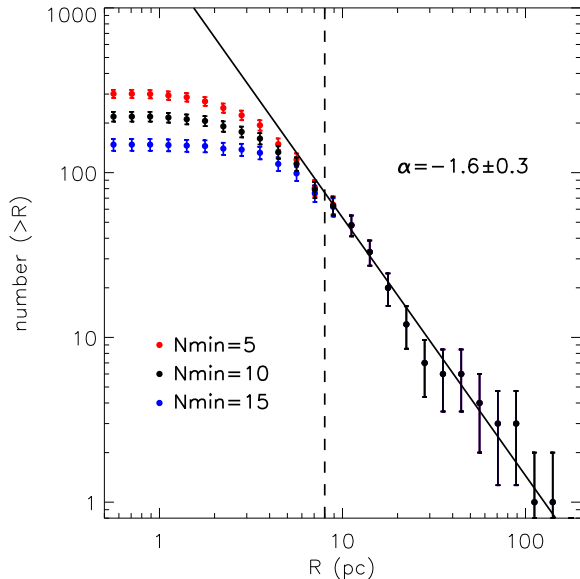


FIG. 7.— Cumulative size distribution of stellar groups with Poissonian errorbars. The red, black and blue data points correspond to different values of N_{\min}^* used in the group-identifying algorithm. **At large R , only the black points are seen because points of different N_{\min}^* overlap with each other.** The vertical dashed line shows the radius beyond which the stellar groups are complete. The solid line is the best-fitting power-law relation to the data points unaffected by incompleteness. Data points of different colors overlap at large sizes.

On larger scales, young stars also exhibit galaxy-wide hierarchical distributions. Specifically, Elmegreen & Elmegreen (2001), Elmegreen et al. (2006), Elmegreen et al. (2014) and Goulier-

mis et al. (2015) reported power-law cumulative size distributions of the young stellar groups in a sample of galaxies; in contrast, Bastian et al. (2007, 2009) suggested that the young stellar groups in M33 and the LMC have log-normal size distributions. However, if we deem the flattening on small scales in their cumulative size distributions (their Figures 5 and 1, respectively) as caused by incompleteness, the large-scale parts can also be well described by single power-laws. Note that they have required each of their identified stellar groups should contain at least 5 sources, which may miss small-sized groups; but they did not discuss the issue of incompleteness as has been done earlier in this paper.

On the other hand, the reported projected fractal dimensions of the galaxy-wide hierarchies span a range from close to 1.0 (e.g. NGC 5477; Elmegreen et al. 2014) to as large as 1.8 or more (e.g. NGC 1705; Elmegreen et al. 2014). D_2 derived in this work is close to those for NGC 628 ($D_2 = 1.5$; Elmegreen et al. 2006), NGC 6503 ($D_2 = 1.7$; Gouliermis et al. 2015), and some of the galaxies in the sample of Elmegreen et al. (2014), but it can deviate more than the measurement uncertainties compared with, e.g., the galaxies in Elmegreen & Elmegreen (2001) and Elmegreen et al. (2014). Elmegreen et al. (2014) found that compared with large spiral galaxies or low surface brightness dwarfs, the starburst dwarfs or HII galaxies have larger projected fractal dimensions if they contain one or two dominant young stellar complexes; thus, the difference in D_2 may reflect different clustering under different conditions. Secondly, older stellar populations have a smaller degree of substructures due to dynamical evolutions (Gieles et al. 2008; Bas-

tian et al. 2009; Parker et al. 2014; Gouliermis et al. 2015); this effect may also contribute to the difference in D_2 , since these studies rely on different stellar samples which may be complicated mixtures of populations of various ages. Moreover, it is worth-mentioning that the derived fractal dimensions may also depend on the adopted method (Federrath et al. 2009), but in this work we do not attempt to explore this effect.

Using the same VMC tile as in this work, Romita et al. (2016) report the discovery of abundant embedded star clusters in molecular clouds. The size distribution of the embedded clusters spans the range of 0.25–2.25 pc with a peak at ~ 1.1 pc, very different from ours. Through their visual inspections, they preferentially identified small and compact cluster candidates, which have high contrast of stellar surface density with respect to their surrounding areas. The group identification algorithm used in this work, however, does not introduce a preferential scale and reveals star groups over a continuous range of sizes. On the other hand, as mentioned above, the star groups identified with the relatively rare UMS stars are incomplete below $R = 8$ pc. Thus, Romita et al. (2016) and this work focus on different subsections of the full hierarchy of young stellar structures. Bonatto & Bica (2010) investigated the size distributions of star clusters and “non-clusters” which are essentially nebular complexes and stellar associations. The former class has a steeper distribution, and the latter has a shallower distribution than ours. Their results are different from ours owing to two reasons. First, they have treated the two classes separately, while we do not try to distinguish star clusters, associations, or larger stellar structures. Second, they have used the catalog of extended objects of Bica et al. (2008), which in turn was compiled from many previous works. Thus the effect of selection bias is hard to assess.

7. SUMMARY AND CONCLUSIONS

In this work we investigate the hierarchical stellar structures of the 30 Dor-N158-N159-N160 star-forming complex traced by young UMS stars observed with the VMC survey. We calculate the surface densities of the UMS stars, whose values cover two orders of magnitudes. We identify groups of UMS stars on different levels of critical surface densities. The larger-sized groups on lower levels often contain several smaller-sized ones on higher levels, and a high degree of sub-grouping is revealed by the dendrograms constructed to show the “parent–child” relations between the groups on different levels. The stellar groups have sizes ranging continuously from several parsecs to more than 100 pc. Without characteristic sizes, the stellar groups have a power-law cumulative size distribution, **with the power-law index indicating a projected fractal dimension $D_2 = 1.6 \pm 0.3$. We suggest that the cumulative size distribution of the UMS stars is related to their parental ISM in a scenario of hierarchical star formation, in which newly-born stars follow the gas distribution. The projected fractal dimension derived here is close to those reported for other star-forming regions as well as for several galaxies; the difference with respect to some other galaxies is also discussed.**

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