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SPINFUL ALGORITHMIZATION OF HIGH ENERGY DIFFRACTION

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DOCTORAL DISSERTATION

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SPINFUL ALGORITHMIZATION OF HIGH ENERGY DIFFRACTION

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Abstract

High energy diffraction probes fundamental interactions, the vacuum, and quantum mechanically coherent matter waves at asymptotic energies. In this work, we algorithmize our abstract ideas and develop a set of rigid rules for diffraction. To get spin under control, we construct a new Monte Carlo simulation engine, GRANITTI. It is the first event generator with custom spin-dependent scattering amplitudes for the glueball domain semi-exclusive diffraction, driven by fully multithreaded importance sampling and written in C++. Our simulations provide new computational evidence that the enigmatic glueball filter observable is a spin polarization filter for tensor resonances. For algorithmic spin studies, we automate the classic Laplace spherical harmonics inverse expansion, carefully define the geometric acceptance related phase space issues and study the harmonic mixing properties systematically in different Lorentz frames.

To improve the big picture, we generalize the standard soft diffraction observables and definitions by developing a high dimensional probabilistic framework based on incidence algebras, COMBINATORIAL SUPERSTATISTICS, and solve also a new superposition inverse problem using the Möbius inversion theorem. For inverting stochastic autoconvolution integral equations or ‘inverting the proton’, we develop a novel recursive inverse algorithm based on the Fast Fourier Transform and relative entropy minimization. The first algorithmic inverse results of the proton double multiplicity structure and multiparton interaction rates are obtained using the published LHC data, in agreement with standard phenomenology. For optimal inversion of the detector efficiency response, we build the first Deep Learning based solution working in higher phase space dimensions, DEEP EFFICIENCY, which inverts the detector response on an event-by-event basis and minimizes the event generator dependence.

Using the ALICE experiment proton-proton data at the LHC at $\sqrt{s} = 13$ TeV, we obtain the first unfolded fiducial measurement of the multidimensional combinatorial partial cross sections, the first multidimensional maximum likelihood fit of the effective soft pomeron intercept and the first multidimensional maximum likelihood fit of the single, double and non-diffractive component cross sections. Great care is taken with the fiducial and non-fiducial definitions. The second topic of measurements centers on semi-exclusive central diffractive production of hadron pairs, which we study with the ALICE data. We measure and fit the resonance spectra of identified pion and kaon pairs, which is crucial on the road towards solving the mysteries of glueballs, the proton structure fluctuations, and the pomeron.

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London, January 2020

Work of the Author

All computational machinery and code of this work are available as Open Source under the MIT license at github.com/mieskolainen for maximal scientific impact, scholarship and reproducibility. *Every* single figure or number in this thesis is produced by this code together with external libraries.

Diffraction phenomenology and computational theory

1. GRANIITTI: A new Monte Carlo event generator and algorithmic engine written in fully multithreaded C++17 with 35k lines of code. Currently the most advanced event generator for the low mass central exclusive diffraction at the LHC and beyond.

Motivation: Explicit control of a Monte Carlo event generator capable of simulating spin dependent scattering processes in the low mass central diffraction with forward proton excitation kinematics.

2. COMBINATORIAL SUPERSTATISTICS: A new statistical incidence algebra approach, Möbius inversion and abstract construction designed as a higher dimensional definition of diffractive event topologies and substructures.

Motivation: Abstract theory beyond standard tools and a self-consistent measurement framework beyond the traditional large rapidity gap counting.

Advanced algorithms

1. DEEPEFFICIENCY: The first Deep Learning based detector efficiency inversion algorithm for maximally unbiased fiducial measurements in higher dimensions.

Motivation: Maximally Monte Carlo event generator independent detector efficiency corrections of single or multidimensional observables.

2. KISU: A new Shannon entropy and the Fast Fourier Transform based algorithm for non-linear inversion of stochastic autoconvolution integral equations. The first application with public ALICE data.

Motivation: Algorithmic way to invert statistically multiparton interactions, multipomeron interactions or pileup.

3. DREM-PID: A new two-step probabilistic particle identification algorithm based on Expectation Maximization (EM) iteration.

Motivation: Mathematically optimal final state identification.

4. *F**-PROJECTION: Overcomplete MC basis projection technique relying on the combinatorial superstatistics framework.
Motivation: Multidimensional data projections with limited detector information.
5. *S*-HARMONICS: A new formulation of the classic Laplace spherical harmonics decay distribution decomposition with novel detail in defining different phase spaces: fiducial versus flat non-fiducial and their invertability and harmonic mixing properties in different Lorentz rest frames.
Motivation: Mathematically rigorous angular distribution and spin polarization measurements.

ALICE and beyond

1. *The first* measurement of unfolded and fiducial *combinatorial cross sections* at the LHC, done with the ALICE proton-proton data at the center-of-mass energy $\sqrt{s} = 13$ TeV. The mathematical construction and analysis technology built up in this thesis.
Motivation: Maximally model independent, multidimensional fiducial measurement of diffraction and soft QCD.
2. *The first* fiducial multidimensional maximum likelihood extraction of single, double and non-diffractive cross sections. A factorized definition of fiducial cross sections and then extrapolation to total inclusive cross sections.
Motivation: Mathematically nearly optimal extraction of the model based cross sections, separate fiducial and extrapolated total cross sections.
3. *The first* multidimensional maximum likelihood extraction of the effective pomeron intercept and the rapidity gap distributions with incomplete forward detector information.
Motivation: High precision diffraction phenomenology with incomplete data.
4. Resonance spectrum measurements of semi-exclusive diffraction at the LHC with the ALICE proton-proton data at the center-of-mass energy $\sqrt{s} = 7$ TeV. We introduce also a new analysis and generic data preservation strategy called *F2X*.
Motivation: Where are the glueballs and what happens with the low mass forward proton dissociation?

List of publications

1. M. Mieskolainen
GRANIITTI: A Monte Carlo Event Generator for High Energy Diffraction
[arXiv:1910.06300 \[hep-ph\]](#)
2. M. Mieskolainen
Combinatorial Superstatistics for Soft QCD
[arXiv:1910.06279 \[hep-ph\]](#)
3. M. Mieskolainen
On the Inversion of High Energy Proton
[arXiv:1905.12585 \[hep-ph\]](#)
4. M. Mieskolainen
DeepEfficiency – optimal efficiency inversion in higher dimensions at the LHC
[arXiv:1809.06101 \[physics.data-an\]](#)
5. M. Mieskolainen
Algorithmics of Diffraction
Acta Physica Polonica B Proceedings of Diffraction and Low-x 2018
[arXiv:1811.01730 \[hep-ph\]](#)
6. M. Mieskolainen, R. Orava
Observables of QCD Diffraction
AIP Proceedings of Diffraction 2016
[arXiv:1612.00980 \[hep-ph\]](#)
7. The MoEDAL collaboration
Magnetic Monopole Search with the Full MoEDAL Trapping Detector in 13 TeV pp Collisions Interpreted in Photon-Fusion and Drell-Yan Production
Physical Review Letters 123, 021802 (2019)
[arXiv:1903.08491 \[hep-ex\]](#)
8. K. Akiba et al.
LHC Forward Physics
Journal of Physics G: Nuclear and Particle Physics 43 110201 (2016)
[arXiv:1611.05079 \[hep-ph\]](#)

Contributions in papers with multiple authors

- * In 6: I wrote the manuscript, came up with the incidence algebras and did the analytical and numerical work.
- * In 7: I contributed by re-implementing independently the photon fusion based generator level simulation of spin-1/2 monopole pairs and thus re-checked the numbers in the paper before the publication. This implementation with velocity β -dependent and pure Dirac couplings, driven by collinear EPA, k_t -EPA and inclusive luxQED-pdf photon fluxes, is publicly available in GRANITTI event generator.
- * In 8: I contributed in the AD detector project and in the ALICE diffraction analysis group in which I participated in detector wrappings, in a beam test, low-level signal analysis, high-level physics analysis, a double rapidity gap trigger design planning and verified detector simulations.

Contents

1	Introduction	1
1.1	Strong and even stronger interactions	4
1.2	Fundamental open problems	8
1.3	Outline	9
2	Waves and Paths	10
2.1	Wave and Helmholtz equation	11
2.2	Helmholtz-Kirchoff integral theorem	12
2.3	Fresnel-Kirchoff and Rayleigh-Sommerfeld formulations	13
2.4	Fresnel and Fraunhofer approximations	15
2.5	Feynman path integral	17
3	ALICE and <i>combinatorial measurement</i>	19
3.1	Experimental setup	20
3.2	Experimental selections and corrections	24
3.3	Unfolded fiducial partial cross sections	29
3.4	Fiducial and total inelastic cross section	36
3.5	Diffraction cross sections and the soft pomeron	38
3.6	F^* -projected observables	42
3.7	Conclusions	46
4	ALICE and <i>glueballs</i>	47
4.1	DREM-PID: Double Recursive Expectation Maximization Particle Identification	49
4.2	F2X: A Faster Analysis of Cross Sections	52
4.3	Experimental tracking and PID setup	55
4.4	Semi-exclusive event selection	57
4.5	System invariant mass spectrum	59
4.6	System transverse momentum	64

4.7	Summary	67
5	GRANIITTI: A Monte Carlo Event Generator for High Energy Diffraction	68
5.1	Introduction	69
5.2	Dynamics	72
5.3	Kinematics and Monte Carlo sampling	111
5.4	Analysis engine	119
5.5	Technology	135
5.6	Discussion and conclusions	136
6	Combinatorial Superstatistics for Soft QCD	139
6.1	Introduction	140
6.2	Binary vector spaces and Diffraction	142
6.3	Posets, incidence algebra and Möbius inversion	154
6.4	Measurements	161
6.5	Combinatorial supercompound Poisson process	167
6.6	Invertibility simulations	180
6.7	Summary	185
7	On the Inversion of High Energy Proton	186
7.1	Introduction	187
7.2	Direct problem	189
7.3	Inverse problem	195
7.4	Algorithm	200
7.5	Simulations	205
7.6	LHC data inversion	210
7.7	Conclusions and prospects	215
8	DeepEfficiency	217
9	Conclusions	222
	Bibliography	224
	Appendix A ALICE measurements	242
A.1	Detector level marginal distributions	242
A.2	Detector response simulations	257
A.3	LHC optics	261

A.4	Trigger statistics	262
Appendix B	GRANIITTI and kinematics	263
B.1	Alternative models of pomeron	263
B.2	Kinematics of $2 \rightarrow 3$	266
B.3	The slope parameter	271
B.4	Harmonic acceptance decompositions	273
Appendix C	Combinatorics and pileup	277
C.1	Vector space subspaces over finite fields	277
C.2	Pileup combinatorics	278
C.3	Exclusive efficiency	282
C.4	Poisson pileup problem	283
C.5	Luminosity and total inelastic cross section	285
C.6	F^* -projection technique	287
C.7	Diffraction fit algorithms	288
Appendix D	GRANIITTI Code (2×2)	290
D.1	Makefile	291
D.2	MADGRAPH amplitude to GRANIITTI converter	294
D.3	C++ header files	296
D.4	C++ source files	333

1 Introduction

Algorithms and high energy physics are dual topics. The first modern computer architecture was introduced by John von Neumann in 1945 [1], the same man behind axiomatic quantum mechanics. The Manhattan project resulted in the first Monte Carlo sampling methods by Stanislaw Ulam and von Neumann. The best known physics algorithms are due to Richard Feynman [2] – the path integral description and the diagrammatic representation of the perturbation series, also the most visual pictures of fundamental interactions. The first large scale computer algebra software SCHOONSCHIP [3] was written by Martinus J.G. Veltman in 1964 in studies towards what later resulted in the renormalizability proof of Yang-Mills [4] theories by Gerard 't Hooft and Veltman. The discretized integral transform by Paul Hough in 1962 [5] was the first computer vision AI-algorithm utilized first in the bubble chamber track fitting, nowadays used by self-driving cars to keep the car between the highway lines or to geometrically align your favorite urban Instagram pictures.

Stephen Wolfram, also from high energy physics, industrialized computer algebra with Mathematica [6] in the 1980's and the World Wide Web was invented at CERN by Tim Berners Lee [7] to organize the experimental information chaos. LHC experiments produce more data faster than any other scientific experiment so far. Also, by physics standards, the most heterogeneous data ever. The ROOT technology [8] developed at CERN during 1990's was during its launch the fastest fully generic object data to disk serializer in the world, perhaps still is. Software like MadGraph have automated computations that have seemed impossible. Nowadays much-hyped quantum computers and their true supremacy are seemingly the distant future of computation, discussed first by Feynman in 1982 [9]. However, the future may rely instead on synthetic biology. Artificial neural networks are currently the mainstream target of large scale computer algebra, more precisely, the network optimization relies on the so-called automatic differentiation techniques, which we have also utilized in this thesis. What is not always known that the crucial reverse-mode automatic differentiation of the current AI industry, later re-invented as the backpropagation, was first invented by a Finnish mathematician Seppo Linnainmaa from Helsinki.

This was in 1971 during his MSc thesis, published later in [10]. Of course, also pure mathematics gets its part from high energy physics, in terms of string theory and scattering amplitude techniques, such as generalized polylogarithms. To summarize, we have listed some down to earth examples to motivate the studies of high energy interactions.

According to the de Broglie particle-wave duality, all matter must exhibit wavelike properties with the wavelength $\lambda = h/p$ inversely proportional to the momentum p , with the bounding constant of fundamental resolution being the Planck constant h . This is natural given the uncertainty principle behind non-commuting observables or the Fourier transform, on the other hand. The wavelike quantum properties of matter span all scales of physics. Experimentally the non-classical behavior has been recently observed even with a chain of 15 amino acid biomolecules [11]. However, we do not know what happens with quantum gravity or is the space-time discretized, for example.

The first documented diffraction observations are from 1665 by Francesco Maria Grimaldi [12]. Perhaps the most famous quantum mechanics experiment is the electron diffraction through the double slit, which demonstrates the probabilistic Born rule and wavelike properties of elementary particles in a highly controlled way. The double helix structure of DNA was discovered by using X-ray diffraction crystallography. Everything of diffraction in terms of quantum electrodynamics (QED) is well understood. However, the Born rule, for example, cannot be currently derived but is an axiom, which fits the data. Many algebraic arguments have been developed to support it. It works remarkably well and there are no known deviations of it. Experiments to test deviations of Born rule are often based on multi-slit diffraction. According to the path integral picture, particles are free to go several loops around the slits in space-time. However, the weight of these highly non-classical paths is usually vanishing in the full probability amplitude.

The topic of this work – hadronic high energy diffraction, is far from well understood. This is due to the non-perturbative strong nature of quantum chromodynamics (QCD), N -body parton problems and the peculiar nature of confinement. That is, there are no currently known powerful *ab initio* methods to simulate soft QCD diffraction, such as lattice methods or small parameter expansions. Phenomenological Monte Carlo modeling is the only way to get proper observables simulated for the comparison with data. However, there are some deep principles behind this modeling, in the context of Regge theory, relativistic kinematics and spin algebra, for example. It is not just parametrizations of data. In general, one needs to remember that we cannot currently even calculate the proton structure function or parton density input at the starting scale for any kind of events, not only diffractive. However,

CHAPTER 1. INTRODUCTION

once the input is fitted from data, the integro-differential evolution schemes such as DGLAP work quite well together with hard matrix elements and parton showers, and reliable predictions can be made for the LHC and elsewhere to hunt for new particles. This is based on a factorization between the soft and hard scales, which is not always exact, but extremely useful nevertheless.

1.1 Strong and even stronger interactions

The primary tool to probe strong interactions at ever-increasing energies is a high energy collider and a general purpose detector built around the interaction point. More and more powerful microscopes are being built, in essence. The initial state may be either leptons, hadrons or heavy ions. The choice of the particular initial state is dictated by the physics goals. High energy collider experiments have a long history at CERN, Brookhaven, SLAC, Berkeley and Protvino. The colliders highly relevant regarding this work are the ISR [13] at CERN, the $Spp\bar{p}S$ [14] at CERN, the LEP [15] at CERN, the HERA [16] at DESY, the Tevatron [17] at Fermilab and nowadays the RHIC [18] at Brookhaven and the LHC at CERN [19, 20]. The future is unknown, but large electron-ion colliders and massive proton-proton colliders may be expected.

The non-abelian $SU(3)$ gauge theory of strong interactions [21–24], quantum chromodynamics, is very complicated. Ideally, we would like to understand it at a similar level of detail as the quantum electrodynamics is understood. Yet understanding QED in detail does not automatically mean we understand collective N -body phenomena, such as chemistry, but only in principle. In the numerous corners of strong interactions, there are many new techniques, but the soft Regge domain has basically always defeated any other methods than the Regge-like power law scaling asymptotics or calculus based on it. There simply seems not to exist any truly powerful small parameter to expand against, other than already found decades ago. The original Reggeon field theory by Gribov was left unfinished by the master. Naturally, many papers have been written to complete it in several plausible directions. However, one should not limit oneself to this massive obstacle posed by field theories – measuring, probing and modeling QCD diffraction is possible in various other ways.

Our number one strategy in this work is to design rigorous mathematical algorithms and then using them, implement novel higher dimensional measurements. Of course, along the way we want to see in explicit detail, what is possible in terms of models, how they compare with the LHC data and introduce an advanced simulation framework, GRANITTI, with special emphasis on spin. For the mathematical formulation of inclusive diffraction observables, our approach is to use incidence algebras. These algebras were formulated by the father of modern combinatorics, mathematician Gian-Carlo Rota at MIT in the 1960’s. That is, we use a simple combinatorial counting of final states as our starting point. Simple is good because already the experimental issues are very complex. Basically, what we formulate is a new definition of diffraction. Funnily enough, this picture is visually closer to the classic diffraction

SOFT QCD	Low- $p_t \sim$ Asymptotic energy dependence, elastic scattering, inelastic diffraction: single, double and central diffraction, multiple large rapidity gaps, spin-parity selection rules in central diffraction, absorption and screening, multiplicity, transverse momentum: from exponential to power laws, soft multiparton (multipomeron) interaction observables, transition to hard scattering, long range rapidity correlations (ridge structure) and high multiplicity without jets, input for fragmentation (hadronization) models and functions
HARD QCD	High- $p_t \sim$ Deep(ly) inelastic scattering, differential jet and system kinematics, event shapes, jet composition, substructures, multijets, radiation patterns and color flow effects, quark/gluon separation, IR/CL-safe jet algorithms (hard, weighted), parton density $f(x, Q^2)$ fitting, hard (HERA) diffraction = hard system + rapidity gap, hard multiparton interactions, search for new massive BSM-resonances, running coupling α_s measurements
COSMIC RAYS	Extended Air Showers (EAS), very forward ‘fragmentation region’ measurements of energy and particle composition at colliders, Monte Carlo tuning
HEAVY IONS	Probing the unknown QCD phase diagram at different densities and temperatures, chiral symmetry recovery, search and understanding for the solid observables of Quark-Gluon Plasma (QGP): photon and lepton rates, strangeness, quarkonia, jet quenching, plasma screening effects, spherical and elliptic flow, fluctuations, Bose-Einstein correlations, phase transition temperature fits, search for GLASMA and other hypothetical (amorphous) states of matter
SPECTROSCOPY	Light mesons, baryons, glueballs and their mixing, multi-quark states
QUARKONIA	$J/\Psi(c\bar{c}), \Upsilon(b\bar{b}), \dots$ and their spin-excited states
SPIN PHYSICS	Quarkonia polarization, polarization in photoproduction Gamma-Pomeron ($\gamma - gg$) processes, proton pdf ‘spin crisis’

Table 1.1: Strong interactions topics classified by experimental observables.

CHAPTER 1. INTRODUCTION

NUCLEAR	Yukawa model (30's), Chiral perturbation theory (60's), Effective theories
REGGE DOMAIN	Regge theory (60-70's), Hard domain (Lipatov et al. since late 70's), Durham QCD (KMR) type models (00's), Stochastic calculus (70's)
LOW- x	Saturation (80's), Color Glass Condensate (90's), Classic Yang-Mills
INTEGRO-DIFFERENTIAL	DGLAP, BFKL, BK ... (70's - 90's)
PARTON DENSITIES	Bjorken scaling, Feynman-Gribov parton model (late 60's), QCD integrated pdfs, unintegrated (generalized) pdfs (80-90's)
COLLIDER QCD, JETS	Fixed order pQCD (late 70's), multileg/multiloop 'NLO revolution' (late 00's), analytic resummation, parton showers, Monte Carlo event generators (late 70's), effective collinear field theories (00's)
HIGH DENSITY	Cold and hot nuclear matter, the equation of state (EOS), Neutron stars, thermal-pQCD + Lattice QCD (70's)
HADRON SPECTROSCOPY	Regge trajectories (60's), Gell-Mann/Zweig quark/ace model (60's), Lattice QCD, Holography (90's), Supersymmetric meson-baryon spectrum (00's)
QUARKONIA	Non-Relativistic NR-QCD (80's)
HADRONIZATION	Lund strings (late 70's), Pre-confinement (Veneziano), Webber clusterization (80's)
VACUUM PROPERTIES	Instanton calculus and 't Hooft, Wilson lattice QFT, lattice QCD (late 70's)
HYDRODYNAMICS	Lattice, transport coefficients (80's)
SCATTERING AMPLITUDES	S -matrix unitarity and analyticity (60's), Spinor-Helicity / Parke-Taylor (80's), Generalized unitarity (90's), Color factorized amplitudes and string theory methods (00's), geometric 'Amplituhedron' methods (10's)

Table 1.2: Strong interactions topics classified in a theory driven way, with the approximate time of origin indicated.

pattern experiments than the de facto ‘large rapidity gap’ counting. What we also see is that it should be one of the best ways to probe the algebraic properties of the Abramovski-Gribov-Kancheli (AGK) scattering amplitude cutting rules [25]. That is, how does one see those rules from data? The interference structure contained in this calculus should be reflected in the multiplicity densities per rapidity interval. Those rules are, after all, of highly combinatorial nature.

To this end, we may mention that increasingly many collider measurements are collected under the automated RIVET [26] platform and their comparison with numerous Monte Carlo models is algorithmized under mcplots.cern.ch [27]. A technical requirement for this to work is that the measurement is a strictly fiducial one. Our novel measurements are also fiducial measurements and thus directly comparable with event generators using only the final state information. We see that highly automated and fiducial approaches should be the case for all fields of physics and science in general, not just limited to collider physics. We shall illustrate for the reader the topics of strong interactions in Tables 1.1 and 1.2. Obviously, we ignore the (effective) strong interactions in condensed matter and elsewhere and stay only in the high energy physics context.

1.2 Fundamental open problems

To expand the mind of the reader regarding where are our topic fits in the large spectrum of modern physics, we shall first list the following to span the space of outstanding problems and questions, in no specific strong order. Our topic belongs to the first one.

- ◇ **Non-perturbative strong interactions and scattering amplitudes**
- ◇ Detectors operating at the single quantum limit at different energy scales
- ◇ Early universe, big bang, inflation, monopoles and topological defects
- ◇ Origin of mass, hierarchy, flavor, matter versus antimatter and details of (the) Higgs boson
- ◇ Unification of dynamics, superstrings, extra dimensions and holography
- ◇ The cosmological constant, dark energy and vacuum(s)
- ◇ Dark matter or misunderstood gravity at large scales
- ◇ Quantum gravity, black holes and information, wormholes and entanglement
- ◇ Stochastic gravitational waves in the (early) universe
- ◇ Quantum computers; unitary port logic driven versus ‘adiabatic’ realizations
- ◇ Algorithms and the synthesis of biology, the arrow of time and entropy
- ◇ Artificial intelligence and physics of neuroscience; classical versus quantum
- ◇ Mathematical physics: number theory and physics, generalized particle statistics, topics of string theory
- ◇ Highly geometric theory of quantum mechanics and space-time; twistors, amplituhedrons, emergent unitarity, explanations for the origin of gauge symmetries
- ◇ High temperature superconductivity, exotic forms of condensed matter and quantum chemistry

1.3 Outline

The structure of this thesis goes as follows. In Chapter 2 we go through the classic picture of diffraction in terms of waves and paths and illustrate these visually brilliant topics through simulations. In Chapter 3 we described the main measurement and multidimensional diffractive cross section extractions of this thesis. In Chapter 4 we discuss new algorithms for semi-exclusive diffraction and glueball hunting and go through case studies with data. In Chapter 5 we have a computational tour of high energy diffraction, where we describe our new GRANITTI Monte Carlo event generator and advanced spin analysis tools. In Chapter 6 we describe in detail the mathematics of our combinatorial measurement framework and the related inverse problems. In Chapter 7 we study the ‘inversion of proton’, by first developing a new inverse algorithm for stochastic autoconvolution integral equations and then apply it to data. In Chapter 8 we describe the first Deep Learning based high dimensional detector efficiency inversion algorithm, DEPEFFICIENCY. Finally in Chapter 9 we end with overall conclusions.

2 Waves and Paths

We shall remind ourselves of the elegant formulations of classic diffraction. There are no radical new results in this chapter, only a compact summary of the essentials.

Huygens (1678)–Fresnel (1818) principle Every space-time point of a propagating wavefront is a source of secondary spherical wavefronts (recursion).

Babinet's principle (~1800) A geometric aperture with a hole gives the same diffraction pattern in the *far field* as the complement aperture (geometry).

2.1 Wave and Helmholtz equation

From the Maxwell equations [28], we can derive the vector valued wave equation [29]

$$\nabla \times (\nabla \times \mathbf{E}(\mathbf{x}, t)) + \frac{1}{v^2} \frac{\partial^2}{\partial t^2} \mathbf{E}(\mathbf{x}, t) = 0, \quad (2.1)$$

which is a second order partial differential equation (PDE) in space and time with the speed of propagation v . Now, the corresponding scalar wave equation reads

$$\left(\nabla^2 - \frac{1}{v^2} \frac{\partial^2}{\partial t^2} \right) U(\mathbf{x}, t) = 0, \quad (2.2)$$

where the scalar function $U(\mathbf{x}, t)$ can be taken one of the spatial components of $\mathbf{E}(\mathbf{x}, t)$. By using the scalar equation, we lose all the polarization (vector) dependent phenomena, which is just fine in the case of acoustic fields, for example.

Then substitution of a single monochromatic wave $U(\mathbf{x}, t) = u(\mathbf{x})e^{-i\omega t}$ at angular frequency ω gives us the linear Helmholtz equation with no time dependence [30]

$$(\nabla^2 + k^2) u(\mathbf{x}) = 0, \quad (2.3)$$

where the wavenumber is $k^2 = \omega^2/v^2$, $k = 2\pi/\lambda$. Usually, one uses Dirichlet and von Neumann boundary conditions. The former sets u on the boundary and the latter defines $\partial u/\partial \mathbf{n}$ on the boundary. A practical way to solve classic diffraction or field configurations in arbitrary geometries and materials is to simulate the scalar fields or Maxwell field equations numerically. An often used formulation in electrodynamics is the Finite Difference Time Domain (FDTD) method invented by Yee in 1966 [31]. However, we shall go through certain analytic classic scalar diffraction theory results. To this end, we may dream about the future where high energy QCD quantized non-linear field equations can be solved on a computer like Maxwell equations.

2.2 Helmholtz-Kirchoff integral theorem

Coordinates: Let our aperture plane and the coordinate system xy -plane coincide with a positive z -axis taken towards the detector. Let \mathbf{x} be the position vector where we evaluate the diffraction integrals at the aperture and let \mathbf{y} be the position vector in the outgoing space at the virtual detector.

We use the Green's [32] equation for a point source

$$(\nabla^2 + k^2) G(\mathbf{x}, \mathbf{y}) = -4\pi\delta(\mathbf{x} - \mathbf{y}) \quad (2.4)$$

with $G(\mathbf{x}, \mathbf{y})$ being the Green function kernel

$$G(\mathbf{x}, \mathbf{y}) = \frac{e^{ik|\mathbf{x}-\mathbf{y}|}}{|\mathbf{x} - \mathbf{y}|}. \quad (2.5)$$

In this case we have translation invariance $G(\mathbf{x} - \mathbf{y}) \equiv G(\mathbf{x}, \mathbf{y})$ by homogeneous medium, such as free space or dielectric. Thus, the solution can be written as a convolution

$$u(\mathbf{y}) = \frac{1}{4\pi} \int_V d\mathbf{x} G(\mathbf{x} - \mathbf{y}) s(\mathbf{x}), \text{ if } \mathbf{y} \in V \text{ and } u(\mathbf{y}) = 0 \text{ for } \mathbf{y} \text{ outside } V, \quad (2.6)$$

for the equation $(\nabla^2 + k^2) u(\mathbf{x}) = -s(\mathbf{x})$ within volume V .

The Helmholtz-Kirchoff (HK) integral theorem [33] states that we can express the scalar field at a point \mathbf{y} as a function of the field values at the volume boundary ∂V

$$\text{HK THEOREM: } u(\mathbf{y}) = \frac{1}{4\pi} \int_{\partial V} dS(\mathbf{x}) \left(u(\mathbf{x}) \frac{\partial G}{\partial \mathbf{n}} \Big|_{\mathbf{x}-\mathbf{y}} - G(\mathbf{x} - \mathbf{y}) \frac{\partial u}{\partial \mathbf{n}} \Big|_{\mathbf{x}} \right), \quad (2.7)$$

where \mathbf{n} is a unit normal vector oriented to inside the volume V with normal derivative $\partial u / \partial \mathbf{n} \equiv \nabla u(\mathbf{x}) \cdot \mathbf{n}$ and $S(\mathbf{x})$ is a surface element. In addition, we need boundary conditions [29]

$$\mathcal{A} : \text{Kirchoff aperture} : u(\mathbf{x}) = u_I(\mathbf{x}) \wedge \frac{\partial u}{\partial \mathbf{n}} = \frac{\partial u_I}{\partial \mathbf{n}} \quad (2.8)$$

$$\mathcal{S} : \text{Kirchoff screen} : u(\mathbf{x}) = 0 \wedge \frac{\partial u}{\partial \mathbf{n}} = 0 \quad (2.9)$$

$$\mathcal{R} : \text{Sommerfeld radiation} : \lim_{|\mathbf{x}| \rightarrow \infty} |\mathbf{x}| \left(\frac{\partial u}{\partial \mathbf{n}} - iku(\mathbf{x}) \right) = 0, \quad (2.10)$$

where we denote the aperture (hole + obstacle) with \mathcal{A} , the opaque screen or detector with \mathcal{S} and the outgoing far field radiation field half-sphere boundary with \mathcal{R} . The incoming field is denoted with $u_I(\mathbf{x})$.

2.3 Fresnel-Kirchoff and Rayleigh-Sommerfeld formula-tions

Now we use a spherical point source

$$u(\mathbf{x}) = A \frac{e^{ik|\mathbf{x}-\mathbf{x}_s|}}{|\mathbf{x}-\mathbf{x}_s|} \equiv A \frac{e^{iks}}{s} \quad (2.11)$$

with the source position \mathbf{x}_s and the amplitude A . For the Green function, we get the gradient as

$$\frac{\partial G}{\partial \mathbf{n}}|_{(\mathbf{x}-\mathbf{y})} = - \left(ik - \frac{1}{|\mathbf{x}-\mathbf{y}|} \right) G(\mathbf{x}-\mathbf{y}) \cos \delta(\mathbf{x}, \mathbf{y}) \quad (2.12)$$

$$\simeq -ikG(\mathbf{x}-\mathbf{y}) \cos \delta(\mathbf{x}, \mathbf{y}), \quad (2.13)$$

where the approximation holds when $|\mathbf{x}-\mathbf{y}| \gg \lambda$. Above the so-called inclination factor is

$$\cos \delta(\mathbf{x}, \mathbf{y}) = \mathbf{n}_x \cdot \left(\frac{\mathbf{x}-\mathbf{y}}{|\mathbf{x}-\mathbf{y}|} \right), \quad (2.14)$$

where δ is the angle between the vector $\mathbf{x}-\mathbf{y}$ and the normal vector \mathbf{n}_x (z -axis). Then using the HK theorem, we get the Fresnel-Kirchoff diffraction equation [34] as

$$\text{FK EQUATION: } u(\mathbf{y}) = -\frac{1}{4\pi} \int_{\mathcal{A}} dS(\mathbf{x}) \left(ik u(\mathbf{x}) \cos \delta(\mathbf{x}, \mathbf{y}) + \frac{\partial u}{\partial \mathbf{n}}|_{\mathbf{x}} \right) G(\mathbf{x}-\mathbf{y}). \quad (2.15)$$

However, this equation is ill-posed by inconsistent boundary conditions, presumably first noted by Sommerfeld. We would like to get rid of the normal derivative $\partial u / \partial \mathbf{n}$ term at the aperture.

Getting rid of the normal term is done by imposing the Sommerfeld radiation condition of Eq. 2.10 to the FK equation and evaluating Eq. 2.12 at the point \mathbf{x} instead of $\mathbf{x}-\mathbf{y}$, which results in a factor of 2 difference. These modifications yield the Rayleigh-Sommerfeld [35] diffraction equation

$$\text{RS EQUATION: } u(\mathbf{y}) = -i \frac{k}{2\pi} \int_{\mathcal{A}} dS(\mathbf{x}) u(\mathbf{x}) \cos \delta(\mathbf{x}, \mathbf{y}) G(\mathbf{x}-\mathbf{y}). \quad (2.16)$$

Note that sometimes in the literature, this is called the Fresnel-Kirchoff equation. This integral is a superposition of source waves (Huygen's principle) with the phase shift at the aperture by $\pi/2$ dictated by $-i$ factor. We see that the boundary conditions restricted the Helmholtz-Kirchoff integral to be non-zero only on the aperture

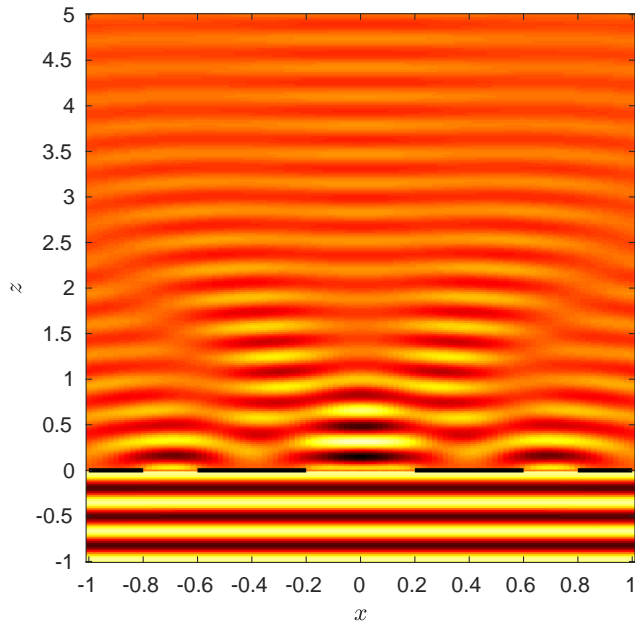


Figure 2.1: A numerical Rayleigh-Sommerfeld integral based simulation of the field values $u(\mathbf{y})$ with three square holes in \mathcal{A} , the source being isotropic at $z \rightarrow -\infty$.

boundary \mathcal{A} . That is, the field value at \mathbf{y} depends only on that. One could call this the first *holographic principle*, later made explicit by Gabor by the invention of optical holography in 1947 and then extended to extra space-time dimensions by 't Hooft, Susskind, Maldacena and others.

2.4 Fresnel and Fraunhofer approximations

By translation invariance and some approximations, we can re-write the RS diffraction equation as a convolution integral

$$u(\mathbf{y}) = \int_{\mathcal{A}} dS(\mathbf{x}) u(\mathbf{x}) h(\mathbf{x} - \mathbf{y}), \quad (2.17)$$

using the kernel function $h(\mathbf{x})$. Now different classic analytic approximation schemes can be obtained by expanding the kernel in various ways and truncating, for example in the so-called Fresnel approximation

$$\text{FRESNEL REGIME: } \frac{b^2}{z\lambda} \sim 1, \quad (2.18)$$

where b^2 is approximately the aperture dimension squared, $z \simeq |\mathbf{x} - \mathbf{y}|$ is the distance from the aperture to the detector and λ is the wavelength of the incoming radiation. Actually, by reciprocity of the FK- or RS-equation, the condition assumes also that z can be replaced by $s = |\mathbf{x} - \mathbf{x}_s|$, which is the distance from the source to the aperture. The corresponding approximated Green's function kernel is

$$h(\mathbf{x} - \mathbf{y}) = \frac{1}{i\lambda z} e^{ikz} e^{\frac{ik}{2z}|\mathbf{x}-\mathbf{y}|^2} \simeq -i \frac{k}{2\pi} \cos \delta(\mathbf{x}, \mathbf{y}) G(\mathbf{x} - \mathbf{y}) \quad (2.19)$$

obtained by Taylor expanding

$$|\mathbf{x} - \mathbf{y}| \simeq z \left(1 + \frac{1}{2} \left[\frac{|\mathbf{x} - \mathbf{y}|}{z} \right]^2 \right) \quad (2.20)$$

inside the exponential function to the first non-trivial order and outside exponential simply with $|\mathbf{x} - \mathbf{y}| \simeq z$ and substituting these in the Green's function of Eq. 2.5. The inclination factor was approximated with 1 in the forward (paraxial) limit $\delta \rightarrow 0$, such that the vector \mathbf{y} has much larger z component than transverse ones. The Fresnel diffraction integral equation is then

$$\begin{aligned} \text{FRESNEL EQUATION: } u(\mathbf{y}) &= \frac{1}{i\lambda z} e^{ikz} \int_{\mathcal{A}} dS(\mathbf{x}) u(\mathbf{x}) e^{\frac{ik}{2z}|\mathbf{x}-\mathbf{y}|^2} \\ &= \frac{1}{i\lambda z} e^{ikz} e^{\frac{ik}{2z}|\mathbf{y}|^2} \int_{\mathcal{A}} dS(\mathbf{x}) u(\mathbf{x}) e^{\frac{ik}{2z}|\mathbf{x}|^2} e^{-\frac{ik}{z}\mathbf{x}\cdot\mathbf{y}}. \end{aligned} \quad (2.21)$$

If the source and the detector are in the *far field* from the aperture or we model the outgoing field from a lens which is focusive (positive), we can model the diffraction

of the waves as plane waves and use the Fraunhofer approximation which results in linear dependence on the integration variable on the aperture, that is, a 2D-Fourier transform. The key point here is that the phase of the field is the same at each point at the aperture, due to incoming plane waves. The approximation condition can be written as

$$\text{FRAUNHOFER LIMIT: } \frac{b^2}{z\lambda} \ll 1, \quad (2.22)$$

where z can again be replaced with s . The Fraunhofer diffraction integral equation is a simplified version of the Fresnel equation

$$\text{FRAUNHOFER EQUATION: } u(\mathbf{y}) = \frac{1}{i\lambda z} e^{ikz} e^{\frac{ik}{2z}|\mathbf{y}|^2} \int_{\mathcal{A}} dS(\mathbf{x}) u(\mathbf{x}) e^{-\frac{ik}{z}\mathbf{x}\cdot\mathbf{y}}, \quad (2.23)$$

obtained by neglecting the quadratic terms from the expansion, to obtain linear dependence on the integration variables. Classic analytic solutions can be obtained by Fourier transform for rectangular ($\rightarrow \text{sinc}^2$), circular ($\rightarrow \text{Airy}$) and Gaussian ($\rightarrow \text{Gaussian}$) density slits.

2.5 Feynman path integral

The Feynman path integral based [2] complex transition amplitude or propagator for a particle, to start from (x_i, t_i) and end up in (x_f, t_f) , in a non-relativistic formulation is

$$K(x_f, t_f | x_i, t_i) = \langle x_f, t_f | x_i, t_i \rangle = \int_{\text{all paths}} D[x(t)] e^{iS[x(t)]/\hbar}, \quad (2.24)$$

where all different paths of the integral represent different quantum *phases*. The transition probability is obtained as the amplitude squared, as usual, by the Born rule. The Born rule cannot be derived, currently. In the transition amplitude, the action functional is a time integral over the Lagrangian, which encapsulates the dynamics of our physics

$$S[x(t)] = \int_{t_i}^{t_f} L[x(t), \dot{x}(t), t] dt \quad (2.25)$$

with a classical Lagrangian, here

$$L = \frac{1}{2} \dot{x}(t)^2 - V[x(t), t], \quad (2.26)$$

where V is an external potential term and the dotted variable is the time derivative. The Feynman path measure $D[x(t)]$ needs to be understood as the following discretization

$$D[x(t)] \equiv \lim_{n \rightarrow \infty} \frac{1}{(2\pi i \hbar \Delta / m)^{d/2}} \int_{\mathbb{R}^d} \prod_{k=1}^{n-1} \frac{dx_k}{(2\pi i \hbar \Delta / m)^{d/2}}, \quad (2.27)$$

with the phase oscillating exponentiated action discretized as

$$\exp(iS[x(t)]/\hbar) \rightarrow \exp\left(\frac{i}{\hbar} \sum_{z=1}^n \left[\frac{1}{2} m \left(\frac{x_z - x_{z-1}}{\Delta} \right)^2 - V(x_z) \right]\right), \quad (2.28)$$

where d denotes the number of spatial dimensions, $x_0 \equiv x_i$, $x_n \equiv x_f$ and $\Delta = (t_f - t_i)/n$. For a mathematician, this path integral measure causes some headache, especially in the case of quantum field theories. The stochastic path integral with the Wiener measure, on the other hand, is well defined. This is probably just lack of suitable mathematics. Physically, this picture is both intuitive and elegant. It also provides the path to lattice quantum field theories.

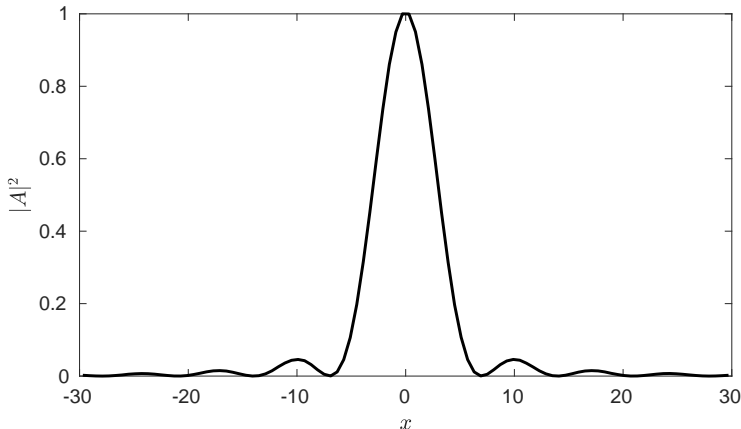


Figure 2.2: A toy Monte Carlo complex path integral simulation of diffraction.

Now a simple single or double slit diffraction can be trivially calculated in 1+1 dimensions for a free particle without any potential term. Instead of an analytical calculation, by curiosity, we did a numerical Monte Carlo simulation, where instead of doing any Wick rotation to imaginary time as is often plausible to do to avoid the highly oscillating exponential term, we brute force evaluated directly the complex path integral in real time by sampling and accumulating discretized paths without any importance sampling. The resulting diffraction pattern is illustrated in Figure 2.2. For proper simulations, one needs to remember fine enough spatial sampling in the propagation and at the detector by the Shannon-Nyquist-Whittaker-Kotelnikov sampling theorem, in order not to produce aliasing. Extreme discretization instability was observed with complex path sampling, a major problem in different quantum simulation scenarios known in general as the *sign problem*. It is an NP-hard problem with no known generic solutions [36]. Basically, this is currently *one* of the fundamental limitations to any ab initio simulations of high energy diffraction and suggests the need for quantum computers, not feasible yet.

On the other hand, deep learning based approximation techniques are already capable of producing extremely high dimensional generative sampling of photorealistic high resolution images [37]. This could be an interesting target for the future research regarding the lattice simulations. This is because image matrices can be understood analogous to extremely complicated lattice field configurations.

3 ALICE and *combinatorial measurement*

We describe the proton-proton combinatorial fiducial cross section measurement at the center-of-mass energy of $\sqrt{s} = 13$ TeV in the ALICE experiment, implemented during the thesis. The full analysis code and grid computing code, including all experimental treatments, is available online on [github](#). It is directly compatible with the ALICE experiment, but all fundamental algorithms and methods are experiment independent, by design. The measurement is the first of its kind, no similar multidimensional unfolded measurement has been attempted before. However, we shall point out that the TOTEM double diffractive measurement [38] is technically a subset of this measurement. Also already the measurements by UA5 [39] were towards this direction, however, philosophically different. Here our main goal is *not* actually to extract diffractive cross sections, but go beyond and implement a fiducial measurement of the combinatorial subspaces or measure the Grasmannian, a mathematician would say. From another perspective, these cross sections are in a sense *coded diffraction aperture* measurements – a high energy analog to the classic case of controlled coded aperture diffraction, which can be used for the phase field recovery with modern algorithms. Here we use them to recover high energy observables and model parameters.

A detailed description of the developed methodology is given in Chapter 6. In the analysis, all basic low level detector distributions were compared with the GEANT simulations driven by Monte Carlo event generators, low level detector signal quality cuts were part of the basic procedures and explicit cut flows were inspected. We shall concentrate here on the physics results but also go through all the basic steps. This and our code should inspire also others to implement a radically new type of measurements.

3.1 Experimental setup

The detector setup consisted of the AD, VZERO and SPD subdetectors, with maximal pseudorapidity coverage. The minimum bias trigger required was algebraic ‘global OR’, where each of the subdetectors had their independent signal decision criteria. The ALICE experiment uses a naming convention of C- and A-sides, where the C-side is on negative rapidities and the A-side on positive. For the generic detector performance of the ALICE experiment, see [40].

The VZERO detector [41] is a forward scintillating plastic counter made with eight cells in four radial rings giving 32 channels per rapidity side. Scintillator detectors are based on a physical mechanism where incoming charged particles induce molecular excitations in the plastic, which is de-excited by the emission of visible wavelength photons. Hamamatsu photomultipliers (PMT) are coupled directly to the scintillating plastics on the A-side and on the C-side through optic fibers. The digitized output gives signal hit time and accumulated charge information. The geometric acceptances over pseudorapidity are $\eta \in [-3.7, -1.7]$ and $[2.8, 5.1]$. The minimum bias trigger decision was based on time-domain signal arrival time cuts which filter out most of the beam induced background such as beam-gas collisions. The PMT high voltage gain setup and signal detection thresholds were adjusted near the single minimum ionizing particle (MIP) limit as described in [41], with parameters adjusted run-by-run for different beam background and noise conditions. The offline calibration was done using a beam test data and cosmic muons when the LHC beam was not active. The detector simulations were done with run anchored setups. In the offline analysis, background and noise sensitivity were studied by requiring also a minimum charge threshold for the signal decision, in addition to more stringent time windows.

The AD detector [42] is a forward shower counter built on similar technology as the VZERO detector but with different geometry and optical coupling. The geometric acceptances over pseudorapidity are $\eta \in [-7.0, -4.9]$ and $[4.8, 6.3]$. It consists of two longitudinal (z -axis) layers of scintillating plastic with four radial transverse quadrants around the beam pipe, giving 8 channels per rapidity side. The scintillation light is steered through a wavelength shifter, to optimize the light transport which is propagated through one meter of optical fibers to Hamamatsu photomultipliers. The front-end electronics is the same as in the VZERO case and the digitized output gives hit time and accumulated charge information. The detector PMT gain and signal threshold setup was similar as with the VZERO and a signal coincidence between adjacent layers was required for the trigger. In the same way as with the VZERO, in the offline analysis, we studied the noise and background by

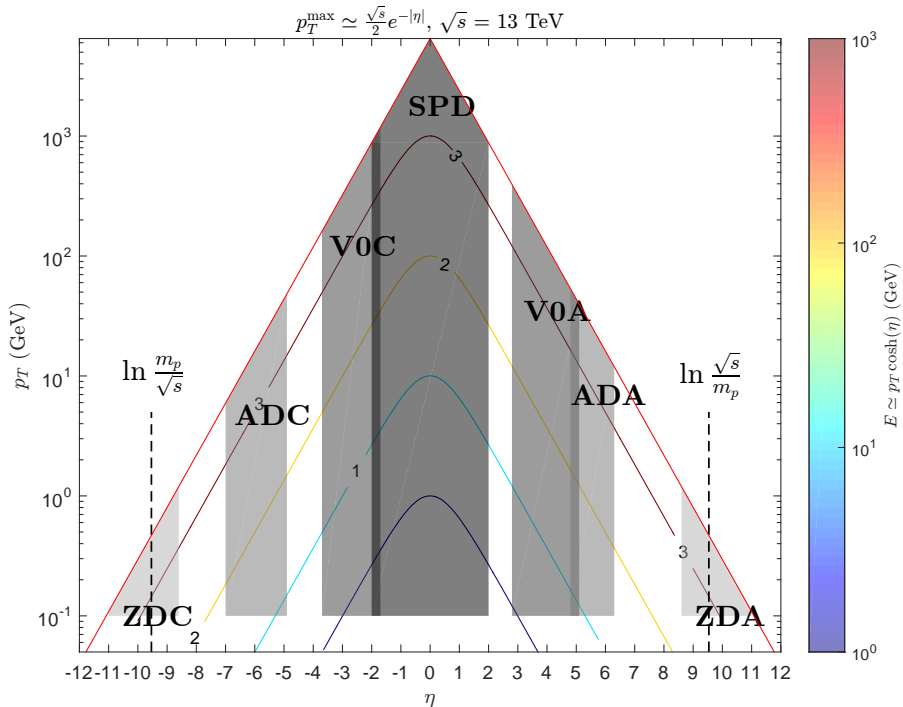


Figure 3.1: The geometric acceptance of ALICE sub-detectors. The triangle plot style adapted and extended from R. Orava.

more stringent accumulated charge and time window cuts.

The SPD (Silicon Pixel Detector) [43] of the inner tracking system (ITS) consists of two layers of pn-type silicon pixel diodes next to the interaction point with the detector cylinder radius 3.9 cm and 7.6 cm and geometric acceptance $|\eta| < 2$ and $|\eta| < 1.5$. The material thickness of each pixel is $200 \mu\text{m}$. The inner layer has 40 half-stave modules and the outer layer 80, each having 10 readout chips with 8192 pixels per chip. This gives in total 1200 readout chips with 9.83 million pixels with $50 \times 425 \mu\text{m}^2$ pitch size. The radiation damaged or otherwise faulty, inactive and noisy outlier chips were explicitly inspected from the data and bit masked also from the simulations. The low level online trigger implemented by the SPD is based on a Fast-OR decision where in each chip, the discriminator output of pixel cells is collected and algebraic OR is taken. In the offline analysis, at least one Fast-OR in the inner or outer SPD layer was required and the noise sensitivity was inspected by requiring more stringent criteria, such as hits in both layers. This, however, effectively also reduces

the acceptance to $|\eta| < 1.5$. Moving further in the transverse direction, there are also two layers of the silicon strip detector (SSD) and two outermost layers of the silicon drift detector (SDD) with acceptance $|\eta| < 0.9$. These layers we did not use due to their more limited acceptance. In the analysis, we used the fired chip level information. In the LHC run III, the inner tracking system will be replaced with a new silicon tracker with 12.5 gigapixels. In addition, we studied the ZDC (zero degree calorimeter) data, which however is not used in the analysis due to lacking calibration, but is used only as ‘spectator’ detector for general data quality cross checks. We illustrate the detector phase-space cartoon coverage in Figure 3.1.

In this analysis, three special low luminosity runs at $\sqrt{s} = 13$ TeV with identification numbers 274593, 274594 and 274595 were chosen with approximately 487k, 418k, 775k triggered minimum bias beam-beam events per run with results given here for the last one, two other used for the irreducible run-by-run uncertainties encapsulating unknown systematic uncertainties. Using the standard coding, the LHC filling scheme was

$$\begin{aligned} &\langle \text{SPACING} \rangle - \langle \text{NBb} \rangle - \langle \text{IP1/5} \rangle - \langle \text{IP2} \rangle - \langle \text{IP8} \rangle - \langle \text{code} \rangle \\ &= \text{Multi-57b-56b-25-20-24-4bpi-15inj}, \end{aligned} \quad (3.1)$$

meaning total 57 and 56 bunches circulating in the clockwise and counterclockwise, out of which 20 colliding bunch pairs in the ALICE IP2 interaction point. This type of beam configuration removes any off-time pileup. The instantaneous luminosity with these runs was very low, with Poisson $\mu \sim 10^{-3}$, given by the normalized global OR trigger rate R

$$\mu = -\ln(1 - R), \quad \text{where } R = \frac{L0b}{N_b f_{rev}} = \frac{L0b}{LMb} \in [0, 1), \quad (3.2)$$

where the LHC revolution frequency is $f_{rev} = 11.245$ kHz, N_b is the number of interacting bunch pairs and L0b (LMb) are the global OR and bunch cross trigger frequencies. The numbers in these runs give vanishing pileup corrections, which however were part of our correction algorithms through our Möbius inversion technique [44]. For the trigger details and rates, see Appendix A.4.

A van der Meer scan using the VZERO detector was used to calibrate the absolute luminosity, details and uncertainties of this are given in [45]. The basic idea is that the $VZERO_C$ and $VZERO_A$, together with their boolean AND trigger combination, provide a visible inelastic cross section

$$\sigma_{V0-AND} = 57.8 \pm 1.2 \text{ mb} \quad (3.3)$$

which is used to directly scale the event counts to visible cross sections because VZERO was part of our detector setup. The *visible* inelastic cross section seen by the VZERO is not a *fiducial* cross section, or neither a *total* inelastic cross section, to emphasize. These three different types of definitions were taken into account by the unfolding and extrapolation procedures.

3.2 Experimental selections and corrections

The offline signal selections are listed in Table 3.1. No high level reconstructed observables were used or are available in ALICE beyond $|\eta| < 0.9$, which would naturally be the optimal case with hypothetical large forward acceptance tracking and calorimetry.

Detector	Signal Cut
AD_C	$\Delta t \in [63, 69]$ ns
$V0_C$	$\Delta t \in [0, 6]$ ns
SPD_C	$F \geq 2$
SPD_A	$F \geq 2$
$V0_A$	$\Delta t \in [7, 14]$ ns
AD_A	$\Delta t \in [54, 60]$ ns

Table 3.1: Offline signal selection quality cuts in terms of the signal arrival time Δt and the number of fired chips F .

After the low-level signal decision criteria, the data was corrected for the irreducible residual beam gas, satellite collisions and noisy events. That is, diffractive events, especially single diffractive, are experimentally very similar to a fixed target like beam gas collisions passing through the time window and charge threshold cuts. To correct this, we used beam-beam, beam-empty, empty-beam and empty-empty trigger masks which were based utilizing the normal and special LHC bunch bucket sequences. The data was corrected with

$$N_j \leftarrow N_j^{(B)} - \alpha^{(A)} N_j^{(A)} - \alpha^{(C)} N_j^{(C)} + 2\alpha^{(E)} N_j^{(E)}, \quad (3.4)$$

where $N_j^{(k)}$ represent the number of events in the j -th combination with the k running over the aforementioned special bunch trigger collected event statistics. The last factor with a positive sign and a factor of two takes into account the double counting. The correction scale factors $\alpha^{(k)}$ taking into account the different bunch sequence luminosity differences and the deterministic and random trigger downscaling were obtained from the trigger statistics with

$$\alpha^{(k)} = \left[\frac{\text{LMb}(B)}{\text{LMb}(k)} \right] \left[\frac{\text{L0a}(B)/\text{L0b}(B)}{\text{L0a}(k)/\text{L0b}(k)} \right], \quad (3.5)$$

where $k = A$ is a beam from the A-side and nothing from the C-side, $k = C$ is a beam from the C-side and nothing from the A-side and $k = E$ is nothing from both

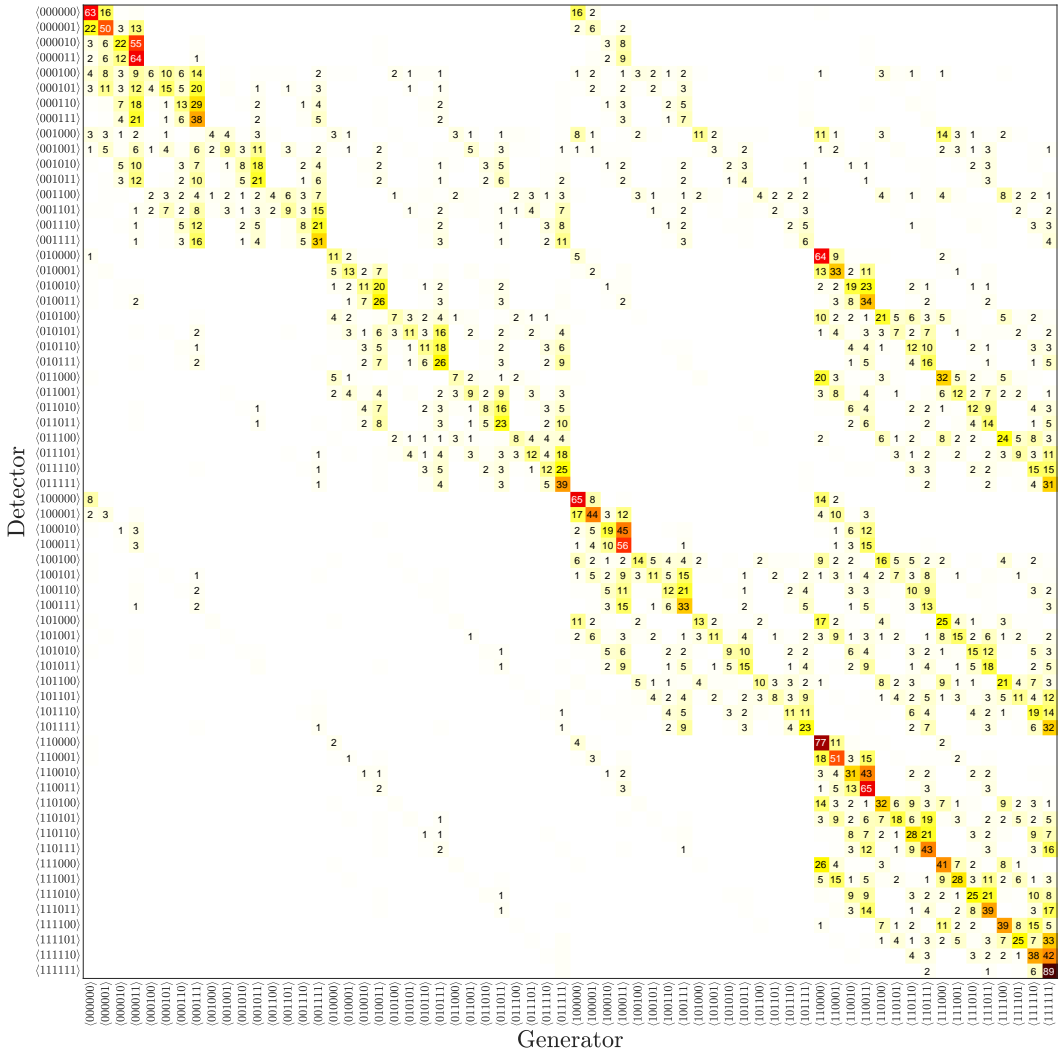


Figure 3.2: ALICE simulation at $\sqrt{s} = 13$ TeV: The 6-dimensional detector folding matrix constructed with Pythia 6 MC + GEANT. The matrix elements are probabilities (%) with each column normalized to unity. Elements $< 1\%$ are not shown.

sides. The $k = E$ case is negligible whenever detectors are operating with low noise levels. The trigger statistics are given in Tables A.3, A.4 and A.5. The detector level marginal distributions of each detector are given in Figures A.1 to A.13.

CHAPTER 3. ALICE AND *combinatorial measurement*

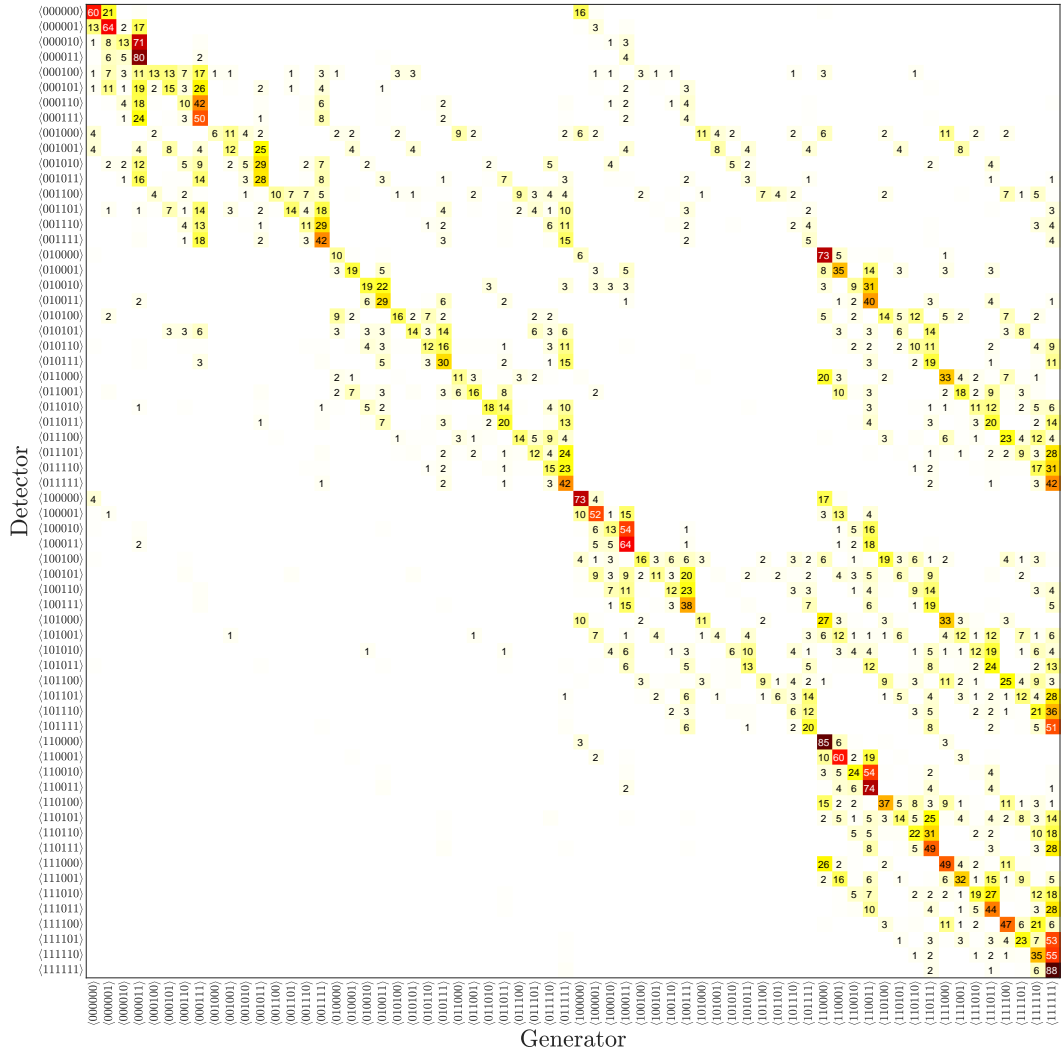


Figure 3.3: ALICE simulation at $\sqrt{s} = 13$ TeV: The 6-dimensional detector folding matrix constructed with Phojet MC + GEANT, otherwise the same as Figure 3.2.

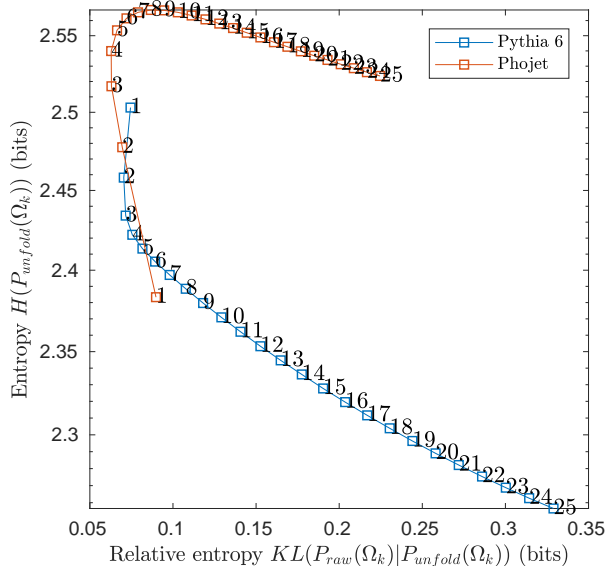


Figure 3.4: ALICE data at $\sqrt{s} = 13$ TeV: The unfolded data Shannon entropy (vertical) versus relative entropy of raw data vs the unfolded (horizontal). The evolution is shown using two MC models used in the detector response simulation. A suitable number of unfolding algorithm iterations (1...25) is near the maximum curvature.

Now speaking of 6-dimensional unfolding of data, we used the standard iterative Bayes method (D’Agostini) of the RooUnfold package [46]. However, we compared the implementation against our own implementation of the same algorithm and also Tikhonov regularized matrix inversion. To be precise, this iterative Bayes unfolding is a frequentist maximum likelihood method using the Expectation Maximization algorithm, which uses the Bayes formula, but is not a Bayesian method. The only Bayesian part is actually the number of iterations, which is an implicit regulator or prior. The unfolding here inverts the inefficiencies and the smearing flow of event topologies from one combinatorial category to another one, thus formally easily described by a folding matrix in the same way as any single dimensional histogram unfolded measurement. These combinatorial folding matrices are shown in Figures 3.2 and 3.3, which demonstrate the major role of detector material rescattering and efficiency corrections.

In the forward domain, even half of the accumulated charge seen by the detector

is from particles propagating from the material interactions, such as $\gamma \rightarrow e^+e^-$ conversions. This is demonstrated in Figure A.17. Fiducial (generator) level event topologies are easily transformed into another topology at the detector level. The simplest acceptance characterizations, here for the visualization purposes, are the diffractive mass efficiency \times acceptances shown in Figure A.16 and A.15, where we see that for the SPD the pure fiducial acceptance and the efficiency \times acceptance are very close. However, for the AD we see a large difference between these two. This is due to the material distributions which actually yield larger effective acceptance than what is geometrically expected. The difference between Pythia 6 and Phojet is due to differences in the fragmentation of the diffractive systems. Phojet has a hard (perturbative) p_t -tail, Pythia 6 generates diffractive events only with soft p_t . Also, the multiplicity distributions are different, as is visible in Figure A.14 where the detector level MC/Data ratios are shown. For a generator level comparisons together with Pythia 8, see the work in [47].

Figure 3.4 demonstrates the unfolding regularization parameter (number of iterations) abstract behavior in terms of Shannon entropy and relative entropy (Kullback-Leibler divergence) between the raw with no unfolding and the unfolded measurement. Interestingly, the two different MC models have approximately the same ‘sweet spot’ but opposite trajectories in this abstract entropy space. We found out that this technique gives a purely data driven handle to control the regularization strength. The point where the regularization should be optimal is the point of maximum curvature, by highly often used L -curve heuristics in inverse problems [48]. Or it could be also the point where trajectories between different models cross. Zero iterations gives no unfolding, whereas too many iterations will push the results towards Monte Carlo estimates, speaking in general. Based on this and simulations, we chose the point of 4 ± 2 iterations in the unfolding algorithm. Accidentally, the default values of the LHC unfolding package parameters are usually around 4 iterations. We should study this interesting topic further in the future, because it has been basically neglected in all existing rapidity gap based measurements, even unfolded ones, because they integrate out the multidimensional information.

We found out that the unfolding has here the largest experimental systematic uncertainty besides the forward detector geometry and material budget simulation, the SPD chip noise and the V0 and AD signal response modeling uncertainties. Understanding these systematic uncertainties in extreme detail would be required for higher precision measurement. Also the beam-gas correction can be in some runs very large, but our combinatorics has a self-consistent way to correct it. Still, run-by-run differences remain much larger than statistical uncertainties in some combinations. A forward tracking and calorimetry with very low material budget would be beneficial.

3.3 Unfolded fiducial partial cross sections

The 6D-measurement fiducial definition we use is charged particles within

$$\begin{aligned} \eta \in [-7.0, -4.9], [-3.7, -1.7], [-2.0, 0], [0, 2.0], [2.8, 5.1], [4.8, 6.3] \\ \wedge p_t > 0.05 \text{ GeV}. \end{aligned} \tag{3.6}$$

This corresponds to the nominal acceptance of the AD, VZERO and SPD detectors. In principle, one could have fine-tuned and optimized the fiducial definition of the measurement beyond the geometric acceptances by minimizing the off-diagonality of the folding matrix. That is, by shifting the simulated fiducial box boundaries per acceptance slice and then re-running the matrix construction, checking the off-diagonality and re-iterating. In the same way, one could also decide if only charged or both charged + neutral particles should be taken as the fiducial definition. However, here we took simply the geometric η -boundaries with a ‘nominal’ p_t -cutoff for charged particles. To denote the 6D-acceptance domains, we use binary coding with six bits.

In Figures 3.5 and 3.6 we have the main result, the unfolded combinatorial partial fiducial cross sections, compared with different minimum bias Monte Carlo event generators. The pink uncertainty bands represent total systematic uncertainty obtained by summing the individual sources in quadrature. The systematic uncertainty sources include the van der Meer scan luminosity calibration, the detector unfolding uncertainties and run-by-run variations reflecting detector calibration and simulation mismatches. The numerical values are given in Tables 3.2, 3.3, 3.4 and 3.5. Now when comparing with Monte Carlo models, the biggest differences are obtained with very forward low mass diffractive like combinations, such as $\langle 010000 \rangle$, which have also the largest uncertainty. Numerous low level cut variations in terms of the accumulated charge and time domain response were inspected, without any major difference. We see that this measurement should provide a highly systematic way to improve the event generator modeling and tunes, together with all the other existing soft QCD measurements, naturally. We emphasize that our definition of diffraction here is well-posed, we do not need to know any internal details about the event generators, only the final state information is needed for the comparison. Thus, this type of measurement is fully compatible with the increasingly popular RIVET [26] measurement-theory comparison framework requirements.

CHAPTER 3. ALICE AND *combinatorial* measurement

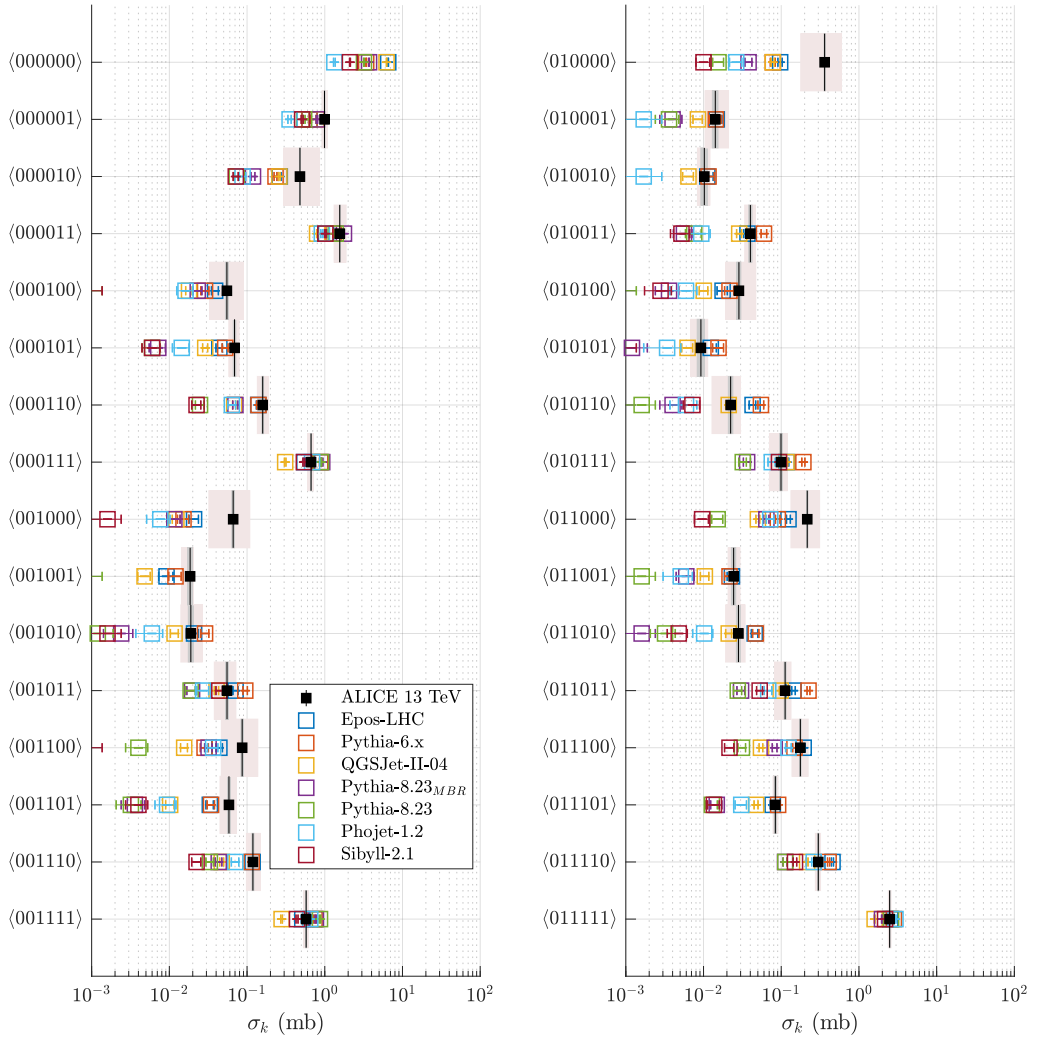


Figure 3.5: ALICE data at $\sqrt{s} = 13$ TeV and MC event generators: The fiducial partial cross sections [0-15] and [16-31].

CHAPTER 3. ALICE AND *combinatorial measurement*

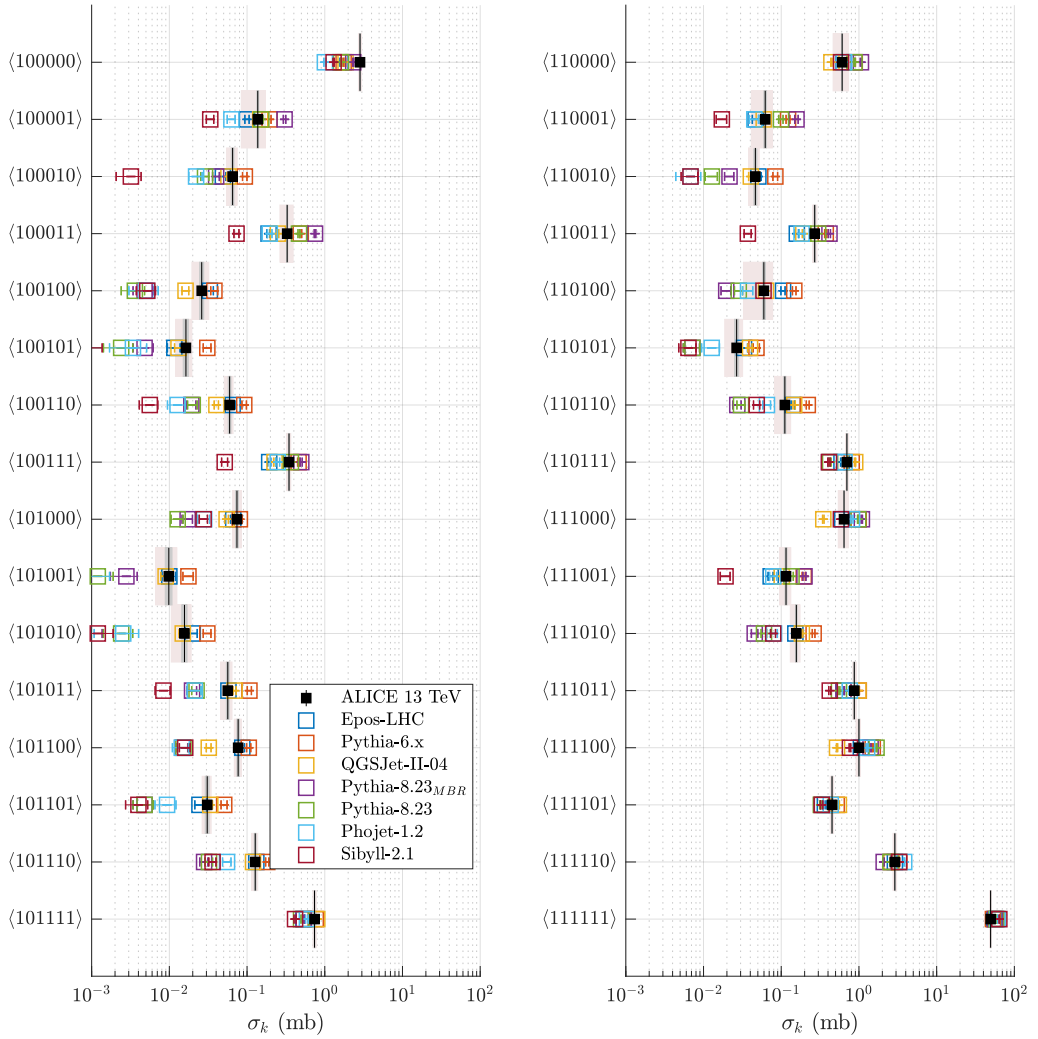


Figure 3.6: ALICE data at $\sqrt{s} = 13$ TeV and MC event generators: The fiducial partial cross sections [32-47] and [48-63].

	X-SECTION	value (mb)	stat	tot.syst	lumi	unfold	run-by-run
0	$\langle 000000 \rangle$	0.000	± 0.000	$+0.000$ -0.000	± 0.000	$+0.000$ -0.000	$+0.000$ -0.000
1	$\langle 000001 \rangle$	0.992	± 0.010	$+0.092$ -0.075	± 0.021	$+0.056$ -0.036	$+0.070$ -0.062
2	$\langle 000010 \rangle$	0.479	± 0.007	$+0.375$ -0.181	± 0.010	$+0.095$ -0.090	$+0.363$ -0.157
3	$\langle 000011 \rangle$	1.558	± 0.012	$+0.325$ -0.239	± 0.032	$+0.047$ -0.021	$+0.320$ -0.236
4	$\langle 000100 \rangle$	0.055	± 0.002	$+0.034$ -0.022	± 0.001	$+0.014$ -0.015	$+0.031$ -0.016
5	$\langle 000101 \rangle$	0.069	± 0.003	$+0.010$ -0.011	± 0.001	$+0.004$ -0.007	$+0.009$ -0.008
6	$\langle 000110 \rangle$	0.159	± 0.004	$+0.030$ -0.023	± 0.003	$+0.012$ -0.009	$+0.027$ -0.021
7	$\langle 000111 \rangle$	0.666	± 0.008	$+0.047$ -0.066	± 0.014	$+0.044$ -0.064	$+0.010$ -0.010
8	$\langle 001000 \rangle$	0.066	± 0.003	$+0.042$ -0.034	± 0.001	$+0.035$ -0.031	$+0.023$ -0.014
9	$\langle 001001 \rangle$	0.019	± 0.001	$+0.002$ -0.004	± 0.000	$+0.002$ -0.004	$+0.000$ -0.000
10	$\langle 001010 \rangle$	0.019	± 0.001	$+0.008$ -0.005	± 0.000	$+0.002$ -0.002	$+0.007$ -0.004
11	$\langle 001011 \rangle$	0.055	± 0.002	$+0.017$ -0.017	± 0.001	$+0.011$ -0.015	$+0.013$ -0.009
12	$\langle 001100 \rangle$	0.087	± 0.003	$+0.051$ -0.040	± 0.002	$+0.038$ -0.034	$+0.034$ -0.020
13	$\langle 001101 \rangle$	0.058	± 0.002	$+0.014$ -0.014	± 0.001	$+0.004$ -0.009	$+0.014$ -0.010
14	$\langle 001110 \rangle$	0.119	± 0.003	$+0.029$ -0.021	± 0.002	$+0.006$ -0.005	$+0.028$ -0.020
15	$\langle 001111 \rangle$	0.576	± 0.008	$+0.037$ -0.043	± 0.012	$+0.019$ -0.032	$+0.029$ -0.027

 Table 3.2: ALICE data at $\sqrt{s} = 13$ TeV: The fiducial partial cross sections [0,15].

CHAPTER 3. ALICE AND *combinatorial measurement*

	X-SECTION	value (mb)	stat	tot.syst	lumi	unfold	run-by-run
16	$\langle 010000 \rangle$	0.362	± 0.006	$+0.234$ -0.184	± 0.008	$+0.168$ -0.159	$+0.163$ -0.092
17	$\langle 010001 \rangle$	0.014	± 0.001	$+0.007$ -0.004	± 0.000	$+0.000$ -0.000	$+0.007$ -0.004
18	$\langle 010010 \rangle$	0.010	± 0.001	$+0.002$ -0.002	± 0.000	$+0.001$ -0.001	$+0.002$ -0.001
19	$\langle 010011 \rangle$	0.040	± 0.002	$+0.006$ -0.006	± 0.001	$+0.002$ -0.004	$+0.006$ -0.005
20	$\langle 010100 \rangle$	0.029	± 0.002	$+0.018$ -0.009	± 0.001	$+0.001$ -0.003	$+0.018$ -0.009
21	$\langle 010101 \rangle$	0.009	± 0.001	$+0.002$ -0.002	± 0.000	$+0.002$ -0.002	$+0.001$ -0.001
22	$\langle 010110 \rangle$	0.022	± 0.001	$+0.007$ -0.009	± 0.000	$+0.006$ -0.009	$+0.005$ -0.004
23	$\langle 010111 \rangle$	0.099	± 0.003	$+0.020$ -0.029	± 0.002	$+0.019$ -0.028	$+0.007$ -0.006
24	$\langle 011000 \rangle$	0.216	± 0.005	$+0.095$ -0.083	± 0.004	$+0.067$ -0.070	$+0.068$ -0.044
25	$\langle 011001 \rangle$	0.024	± 0.002	$+0.005$ -0.004	± 0.001	$+0.002$ -0.001	$+0.005$ -0.004
26	$\langle 011010 \rangle$	0.028	± 0.002	$+0.006$ -0.009	± 0.001	$+0.006$ -0.009	$+0.000$ -0.000
27	$\langle 011011 \rangle$	0.112	± 0.003	$+0.021$ -0.030	± 0.002	$+0.021$ -0.030	$+0.005$ -0.005
28	$\langle 011100 \rangle$	0.176	± 0.004	$+0.046$ -0.039	± 0.004	$+0.024$ -0.026	$+0.039$ -0.028
29	$\langle 011101 \rangle$	0.084	± 0.003	$+0.007$ -0.006	± 0.002	$+0.004$ -0.003	$+0.005$ -0.005
30	$\langle 011110 \rangle$	0.300	± 0.005	$+0.021$ -0.025	± 0.006	$+0.017$ -0.022	$+0.010$ -0.010
31	$\langle 011111 \rangle$	2.484	± 0.016	$+0.102$ -0.095	± 0.052	$+0.024$ -0.004	$+0.084$ -0.080

Table 3.3: ALICE data at $\sqrt{s} = 13$ TeV: The fiducial partial cross sections [16,31].

	X-SECTION	value (mb)	stat	tot.syst	lumi	unfold	run-by-run
32	$\langle 100000 \rangle$	2.842	± 0.017	$+0.133$ -0.126	± 0.059	$+0.024$ -0.026	$+0.117$ -0.109
33	$\langle 100001 \rangle$	0.137	± 0.004	$+0.035$ -0.053	± 0.003	$+0.033$ -0.052	$+0.012$ -0.011
34	$\langle 100010 \rangle$	0.065	± 0.003	$+0.010$ -0.011	± 0.001	$+0.002$ -0.008	$+0.009$ -0.007
35	$\langle 100011 \rangle$	0.328	± 0.006	$+0.064$ -0.063	± 0.007	$+0.015$ -0.042	$+0.062$ -0.046
36	$\langle 100100 \rangle$	0.026	± 0.002	$+0.006$ -0.007	± 0.001	$+0.004$ -0.006	$+0.004$ -0.003
37	$\langle 100101 \rangle$	0.016	± 0.001	$+0.003$ -0.004	± 0.000	$+0.003$ -0.004	$+0.001$ -0.001
38	$\langle 100110 \rangle$	0.060	± 0.002	$+0.006$ -0.009	± 0.001	$+0.004$ -0.008	$+0.005$ -0.004
39	$\langle 100111 \rangle$	0.347	± 0.006	$+0.018$ -0.021	± 0.007	$+0.003$ -0.012	$+0.017$ -0.015
40	$\langle 101000 \rangle$	0.074	± 0.003	$+0.010$ -0.009	± 0.002	$+0.009$ -0.007	$+0.005$ -0.005
41	$\langle 101001 \rangle$	0.010	± 0.001	$+0.003$ -0.003	± 0.000	$+0.002$ -0.003	$+0.002$ -0.001
42	$\langle 101010 \rangle$	0.016	± 0.001	$+0.003$ -0.005	± 0.000	$+0.003$ -0.005	$+0.001$ -0.001
43	$\langle 101011 \rangle$	0.057	± 0.002	$+0.007$ -0.011	± 0.001	$+0.007$ -0.011	$+0.002$ -0.002
44	$\langle 101100 \rangle$	0.077	± 0.003	$+0.008$ -0.008	± 0.002	$+0.007$ -0.007	$+0.003$ -0.003
45	$\langle 101101 \rangle$	0.031	± 0.002	$+0.004$ -0.004	± 0.001	$+0.003$ -0.004	$+0.002$ -0.002
46	$\langle 101110 \rangle$	0.127	± 0.004	$+0.010$ -0.012	± 0.003	$+0.009$ -0.012	$+0.003$ -0.003
47	$\langle 101111 \rangle$	0.738	± 0.009	$+0.040$ -0.035	± 0.015	$+0.030$ -0.023	$+0.022$ -0.021

 Table 3.4: ALICE data at $\sqrt{s} = 13$ TeV: The fiducial partial cross sections [32,47].

CHAPTER 3. ALICE AND *combinatorial measurement*

	X-SECTION	value (mb)	stat	tot.syst	lumi	unfold	run-by-run
48	$\langle 110000 \rangle$	0.608	± 0.008	$+0.126$ -0.143	± 0.013	$+0.119$ -0.138	$+0.040$ -0.036
49	$\langle 110001 \rangle$	0.062	± 0.002	$+0.015$ -0.021	± 0.001	$+0.012$ -0.020	$+0.008$ -0.007
50	$\langle 110010 \rangle$	0.046	± 0.002	$+0.006$ -0.009	± 0.001	$+0.004$ -0.008	$+0.004$ -0.003
51	$\langle 110011 \rangle$	0.270	± 0.005	$+0.018$ -0.016	± 0.006	$+0.009$ -0.006	$+0.015$ -0.014
52	$\langle 110100 \rangle$	0.060	± 0.002	$+0.018$ -0.027	± 0.001	$+0.017$ -0.027	$+0.006$ -0.005
53	$\langle 110101 \rangle$	0.027	± 0.002	$+0.005$ -0.008	± 0.001	$+0.005$ -0.008	$+0.001$ -0.001
54	$\langle 110110 \rangle$	0.111	± 0.003	$+0.021$ -0.029	± 0.002	$+0.019$ -0.028	$+0.008$ -0.007
55	$\langle 110111 \rangle$	0.700	± 0.008	$+0.037$ -0.026	± 0.015	$+0.029$ -0.015	$+0.017$ -0.016
56	$\langle 111000 \rangle$	0.643	± 0.008	$+0.087$ -0.101	± 0.013	$+0.063$ -0.086	$+0.058$ -0.050
57	$\langle 111001 \rangle$	0.115	± 0.003	$+0.018$ -0.020	± 0.002	$+0.008$ -0.015	$+0.015$ -0.012
58	$\langle 111010 \rangle$	0.156	± 0.004	$+0.018$ -0.026	± 0.003	$+0.017$ -0.025	$+0.005$ -0.004
59	$\langle 111011 \rangle$	0.871	± 0.009	$+0.053$ -0.039	± 0.018	$+0.047$ -0.030	$+0.017$ -0.017
60	$\langle 111100 \rangle$	1.000	± 0.010	$+0.046$ -0.082	± 0.021	$+0.032$ -0.076	$+0.024$ -0.023
61	$\langle 111101 \rangle$	0.451	± 0.007	$+0.022$ -0.023	± 0.009	$+0.001$ -0.010	$+0.019$ -0.018
62	$\langle 111110 \rangle$	2.897	± 0.017	$+0.155$ -0.153	± 0.060	$+0.041$ -0.062	$+0.137$ -0.126
63	$\langle 111111 \rangle$	49.444	± 0.070	$+1.272$ -1.319	± 1.027	$+0.108$ -0.404	$+0.743$ -0.724

Table 3.5: ALICE data at $\sqrt{s} = 13$ TeV: The fiducial partial cross sections [48,63].

3.4 Fiducial and total inelastic cross section

The fiducial inelastic measurement and the corresponding extrapolated full phase space total inelastic cross section are given numerically in Table 3.6 and visually in Figure 3.7. Figure 3.8 shows in detail the Monte Carlo model based total inelastic cross section extrapolations based on the measured partial cross sections. EPOS-LHC and QGSJet-II-04 based extrapolations give values with a closest match with respect to the total inelastic measurement done by the TOTEM experiment at $\sqrt{s} = 13$ TeV giving 79.5 ± 1.8 mb, which is based on the optical theorem and the elastic scattering with Mandelstam $t \rightarrow 0$ extrapolation [49]. That measurement has the smallest experimental uncertainty at the LHC. Alternatively, one should instrument the experiments with extremely forward shower counters to observe the total inelastic cross section directly.

Cross Section	value \pm (stat+syst)
σ_{inel}^{fid}	71.4 ± 1.8 mb
σ_{inel}^{tot} (ext.)	77.3 ± 2.1 mb

Table 3.6: ALICE data at $\sqrt{s} = 13$ TeV: The inelastic cross section in the fiducial domain and the total extrapolated one obtained using re-fitted Pythia 6.

The EPOS-LHC, QGSJet-II-04 and Pythia 6 generators have here the best overall tune, in terms of χ^2 . The χ^2 -value is simply calculated against the total inelastic re-scaled Monte Carlo estimate of the fiducial partial cross sections and the measured ones, that is, one number is being varied in the fit. We took care of the technical scale bias easily introduced in χ^2 type fits. All of the event generators have some obvious tension given that the number of degrees of freedom is here $2^N - 2 = 62$. This global tension is easily explained by certain specific deviating partial cross sections. In Figure 3.7 we have also our own re-tune of Pythia 6 based extrapolation indicated, based on the extracted diffractive cross sections and the extracted soft Pomeron intercept value. Our tune is significantly closer to the EPOS-LHC and QGSJet-II-04 than the original Pythia 6 in the CRMC package.

The overall inelastic fiducial acceptances $\sigma_{inel}^{fid}/\sigma_{inel}^{tot}$ predicted by different generators are: EPOS-LHC (0.92), QGSJet-II-04 (0.92), Pythia 8.23-MBR (0.95), Pythia 6 (0.96), Pythia 8.23 (0.96), Sibyll 2.1 (0.97) and Phojet 1.2 (0.98). This indicates that all other event generators than EPOS and QGSJet-II-04 clearly underestimate the cross section of the invisible low mass diffraction. This conclusion is compatible with our other results, such as the F^* -projected gap distributions in Section 3.6.

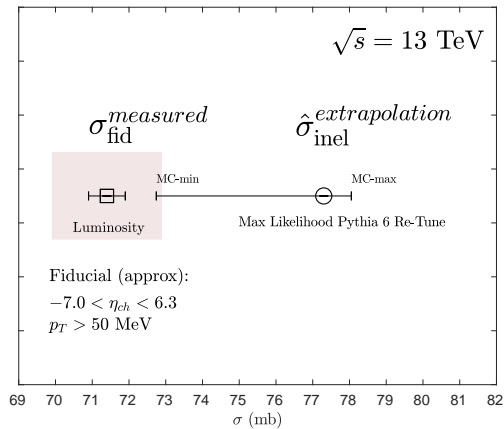


Figure 3.7: ALICE data at $\sqrt{s} = 13$ TeV: The fiducial inelastic proton-proton cross section on the left, and the extrapolated total inelastic cross section using different MC models on the right. The luminosity uncertainty is denoted with a pink box.

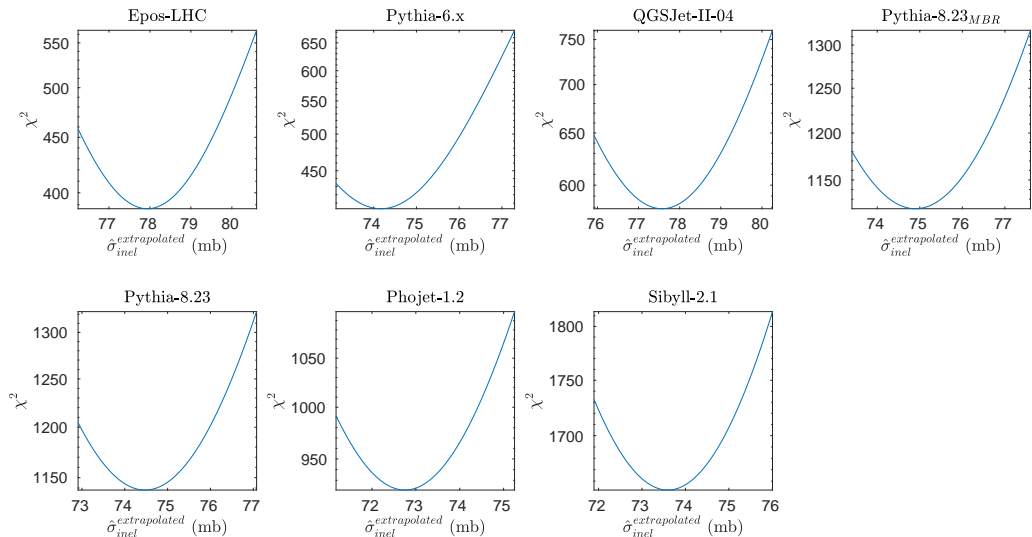


Figure 3.8: ALICE data at $\sqrt{s} = 13$ TeV: The total inelastic cross section extrapolation fit χ^2 -values with different MC models using the fiducial unfolded combinatorial cross sections as the fit input.

3.5 Diffractive cross sections and the soft pomeron

Extracting the diffractive cross section is not experimentally a uniquely posed task, but it can be mathematically clearly formulated as a multidimensional probabilistic fitting or classification problem as first done in a work with CDF data by me [50], or later by Weltri with TOTEM+CMS data in [51]. Also, in a work by Matthews [52] the trigger combinations were used to extract the diffractive cross sections using ALICE data. Here, we improve the methodology by a maximum likelihood formulation, which allows also simultaneously to vary the minimum average rapidity gap definition and the effective Pomeron intercept. But most importantly, here we implemented for the first time a proper multidimensional fiducial measurement which then allows to implement numerous fits with any future Monte Carlo event generator, using the unfolded partial cross sections from Tables 3.2, 3.3, 3.4 and 3.5.

In Figure 3.9 we have the diffractive cross section extractions using the multidimensional maximum likelihood 6D-fit machinery and their extrapolations to the full phase space with uncertainties defined as before. SDL(R) denotes the diffractive system to negative (positive) rapidities, DD denotes double diffractive and ND incoherent non-diffractive processes. The relative uncertainties are the largest with single diffraction, due to run-by-run variations (beam background). We have done the extractions as a function of the average kinematic rapidity gap, which is simply done by calculating the invariant mass of the diffractive system at the generator level and then calculating the equivalent gap sizes

$$\text{SD: } \langle \Delta Y \rangle_{\min} = -\ln \left(\frac{M^2}{s} \right) \quad (3.7)$$

$$\text{DD: } \langle \Delta Y \rangle_{\min} = -\ln \left(\frac{M_1^2 M_2^2}{s s_0} \right), \quad s_0 = 1 \text{ GeV}^2. \quad (3.8)$$

Then we excluded the simulated events with smaller gap values. Thus, one may take the traditional integration bound of $\langle \Delta Y \rangle_{\min} = 3$ and compare with previous measurements in [53]. Our definition is a sharp one, because this simply means that we have a cutoff in the diffractive model before it is fragmented. The ALICE detector has an asymmetric rapidity acceptance, thus, the fiducial single diffractive cross sections are asymmetric between forward and backward directions. Their extrapolated values are symmetric, as they should, if the detector simulation together with data is well understood. We remark here that this is by no means automatically guaranteed, given that the observables with forward shower counters are very low level, that is, no tracks and no transverse momentum being measured so the calibration of the data is non-trivial.

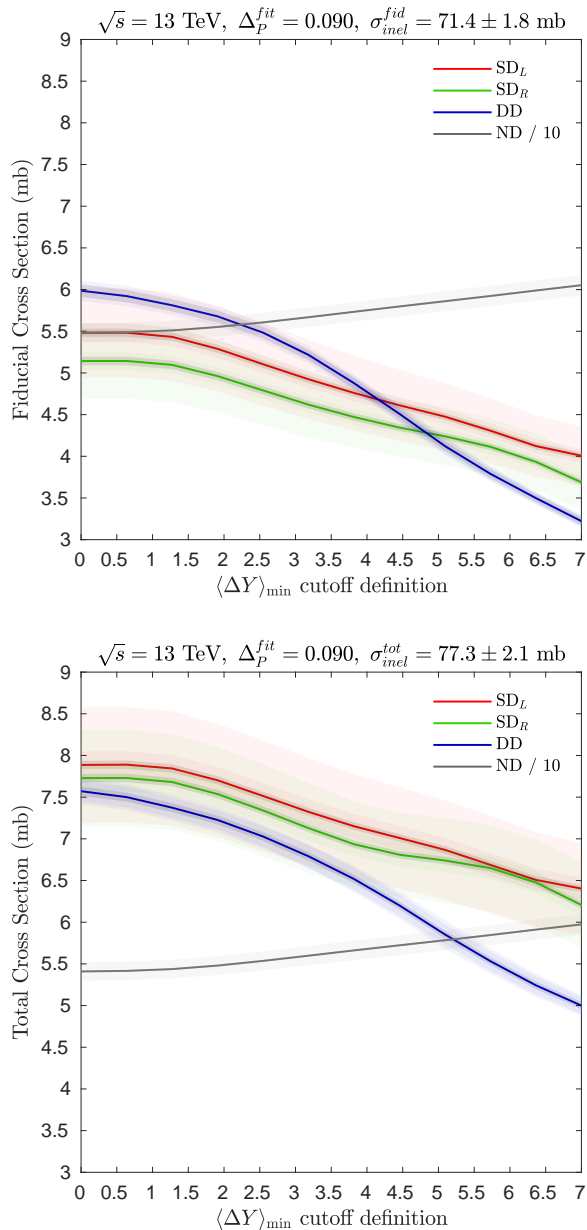


Figure 3.9: ALICE data at $\sqrt{s} = 13$ TeV: The unfolded fiducial cross sections (top) and the full phase space extrapolated (bottom) as a function of the minimum average rapidity gap cutoff based on 6D-re-fitting against Pythia 6 MC templates.

In Figure 3.10 we illustrate the typical fit cost of the soft pomeron intercept Δ_P , which is a fundamental free parameter of the Regge theory. Note that minimizing the Kullback-Leibler divergence gives a mathematically equivalent solution to the maximum likelihood, well known in the field of machine learning. Exact numerical values are given in Table 3.7 yielding a total estimate

$$\Delta_P = 0.094 \pm 0.01 \text{ (stat+syst)}, \quad (3.9)$$

where we take a simple median and standard deviation of the individual extractions. This extraction is the first maximum likelihood multidimensional fit of this parameter at the LHC, as far as we know, also the most precise one. For comparison, a recent measurement by the ATLAS+ALFA experiment obtained a more uncertain value of $\Delta_P = 0.07 \pm 0.09$ in single diffraction with forward proton measurements at $\sqrt{s} = 8$ TeV [54].

The extracted intercept value is very close to the one obtained in classic fits by Donnachie and Landshoff giving the soft pomeron $\alpha(t) = 1 + \Delta_P + \alpha' t \simeq 1.08 + 0.25t$, based on the total cross section values [55]. Actually, in later fits by DL, the intercept is a bit higher. Here, our fit is based on the differential diffractive mass distribution which we indirectly re-fit using the combinatorial cross sections as the observables and Monte Carlo model based template distributions. The pomeron intercept refit has a larger MC based diffractive mass distribution uncertainty with Pythia 6 than Phojet, because Pythia 6 includes ‘fudge factors’ at low and high masses, whereas Phojet has pure power law dependence. However, we note that this fit can be easily re-done with the future Monte Carlo models using the measured unfolded partial cross section values, with any fast (or brute force) fit techniques available. The soft pomeron slope α' can be extracted for example in the exclusive central diffraction, from the system transverse momentum dependence or forward proton measurements. We shall remark here that the Pomeron intercept and slope are model dependent quantities, in a more fundamental Regge calculus case one may talk about the ‘bare’ intercept and slope which are then dressed a bit like in field theory. These values are in principle calculable from the field theory of QCD, but in practice there no known non-perturbative techniques yet available. For more about this, see the discussion of the eikonal pomeron in Chapter 5.

MC model	Run	Δ_P pre-unfold	Δ_P post-unfold
Pythia 6	274593	0.088	0.102
Pythia 6	274594	0.094	0.106
Pythia 6	274595	0.085	0.099
Phojet	274593	0.094	0.091
Phojet	274594	0.100	0.097
Phojet	274595	0.089	0.087

Table 3.7: ALICE data at $\sqrt{s} = 13$ TeV: Run-by-run maximum likelihood 6D-fit extractions of the soft Pomeron intercept under two different MC models and before and after unfolding. Statistical uncertainties are negligible.

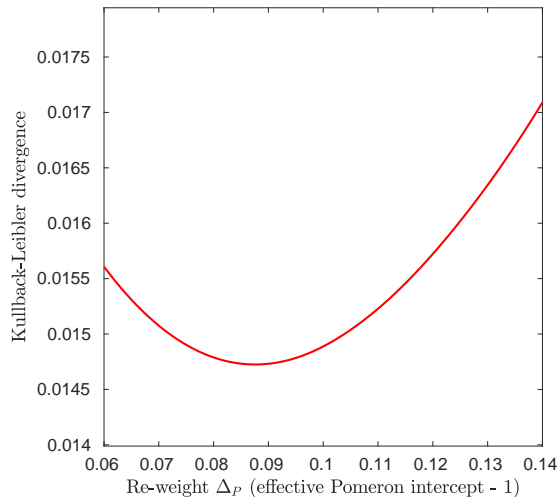


Figure 3.10: ALICE data at $\sqrt{s} = 13$ TeV: Typical behavior of the soft pomeron intercept Δ_P in the maximum likelihood 6D-fit against Pythia 6 MC templates.

3.6 F^* -projected observables

In Figures 3.12, 3.13 and 3.11 we have the unfolded pseudorapidity gap distributions projected using our F^* -projection algorithm, described in [44] and Appendix C.6. The type I boundary conditions calculate gaps with respect to the detector edge and type II between any two particle within the detector. The ALICE detector has no strict detector capabilities for direct measurement of these such as large calorimetry or tracking with p_t measurement, thus our novel MC projection algorithm based on the combinatorial cross sections, represent a state-of-the-art approach. It is not a pure measurement, but an overcomplete ‘basis’ projection.

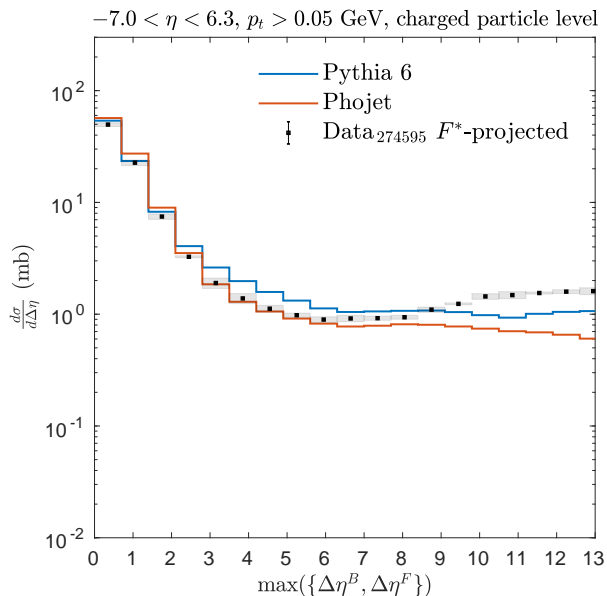


Figure 3.11: ALICE data at $\sqrt{s} = 13$ TeV and MC event generators: Type I pseudorapidity gap size distribution measured with respect to the detector edge, and the maximum of the backward *or* the forward direction gap is chosen per event. The projection uncertainties denote MC model variations.

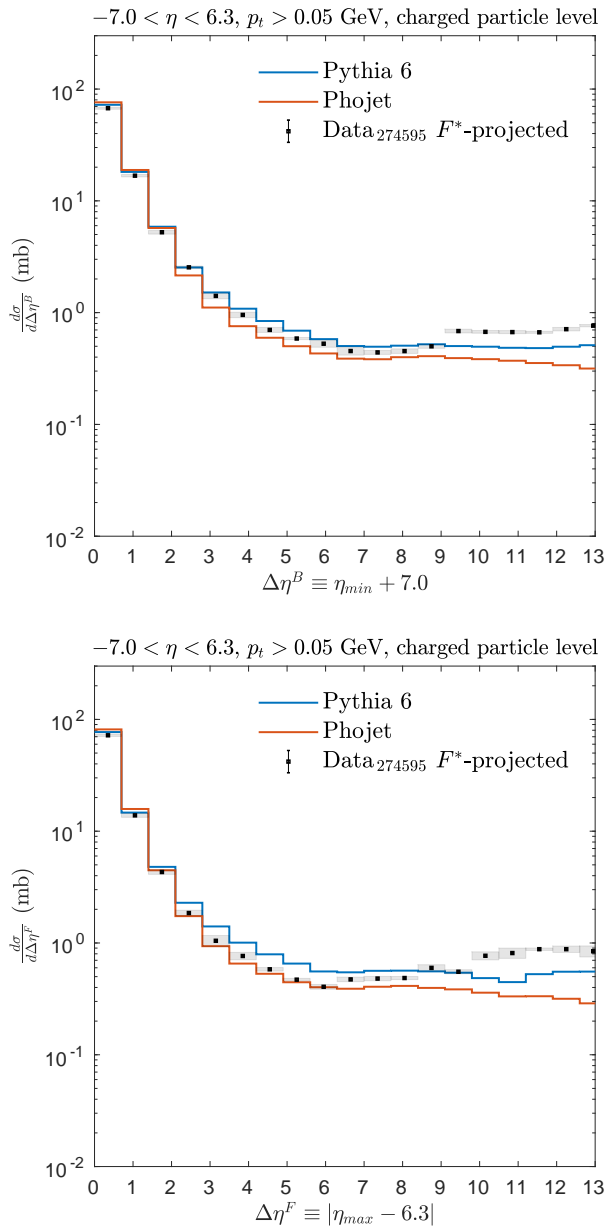


Figure 3.12: ALICE data at $\sqrt{s} = 13$ TeV and MC event generators: Type I pseudorapidity gap size distribution to the backward (top) and the forward (bottom) direction measured with respect to the detector edge. The projection uncertainties denote MC model variations.

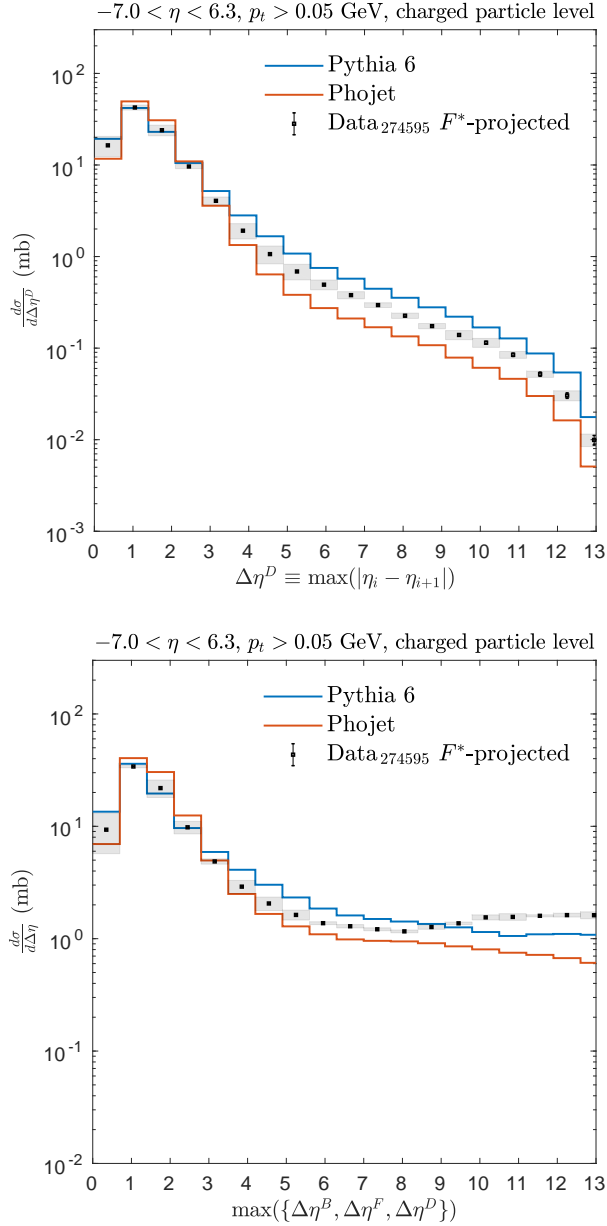


Figure 3.13: ALICE data at $\sqrt{s} = 13$ TeV and MC event generators: Type II pseudorapidity gap size distribution (top), and the maximum of type I and type II gaps per event (bottom). The projection uncertainties denote MC model variations.

One can see that the low mass diffraction, which dominates at large pseudorapidity gap sizes, is more prominent in data compared with Pythia 6 and Phojet event generators. The magnitude of low mass diffraction cannot be calculated from the Regge power law asymptotics controlled by the Regge trajectories, thus its proportion in event generators is purely model and tune dependent. At low masses, the proton structure fluctuations and baryonic resonances dominate. The values obtained here are in the line with the ATLAS measurement at $\sqrt{s} = 7$ TeV, which however has data only up to $\Delta\eta \approx 8$ [56]. They have an explicit control on the transverse energy cutoff, although the calorimeter noise starts to dominate over the signal at very low- p_t .

3.7 Conclusions

The main new results of the measurements presented in this section with the pp data at the cms energy $\sqrt{s} = 13$ TeV are as follows.

- ◇ The first unfolded multidimensional fiducial measurement of combinatorial partial cross sections.
- ◇ The first multidimensional maximum likelihood extraction of the single, double and non-diffractive cross sections. Re-extraction is possible a posteriori with future Monte Carlo models based on the measured unfolded partial cross sections.
- ◇ The first multidimensional extraction of the effective soft Pomeron intercept with $\Delta_P = 0.094 \pm 0.01$ (stat+syst).
- ◇ Low mass diffraction enhancement with respect to Monte Carlo event generator tunes is clearly visible in the forward pseudorapidity gap distributions up to $\Delta\eta \simeq 13$ units, projected using the newly developed F^* -algorithm.
- ◇ Our new measurements can be used to get more understanding of the AGK rule type calculus, which is only partially understood theoretically. These are used implicitly in the construction of the EPOS LHC and QGSJet-II-04 event generators, which give one of the best descriptions of the measured partial cross sections. However, the quality of MC parameter ‘tunes’ plays always a big factor. The detector level distributions in Appendix A.1 are also a step forward on this topic. A future measurement using \mathbb{F}_q^N with $q = 3$ or more, with three different multiplicity bins per rapidity slice, could make this topic more transparent. This higher order measurement requires more statistics at least by a factor of $3^N/2^N$.

4 ALICE and *glueballs*

We introduce new algorithms for the studies of exclusive and semi-exclusive central diffraction (CED/CEP) and describe our measurements with ALICE data. Potentially the most interesting on this topic is the future identification of glueball resonances, understanding the pomeron spin structure and the contribution of proton dissociative events. Glueballs, the resonant bound states of non-abelian gauge fields, are very difficult objects because they mix quantum mechanically with quark bound states. A curious mathematical fact was pointed out by Coleman already in 1977, that *classical* non-abelian gauge fields do *not* contain glueball like field configurations [57]. Some predictions for the quantum glueballs can be calculated using lattice QCD [58, 59] for pure glue states within the ‘quenched approximation’, or using holographic approximations within Witten-Sakai-Sugimoto model based descriptions [60, 61]. Different mixing matrix scenarios have been considered also purely phenomenologically [62].

The experimental identification of glueballs is an extremely laborious task, requiring simultaneous measurements in multiply decay channels such as pion pairs, kaon pairs but also photon pairs. We know that gluons couple only to quarks or other gluons, so a (pure) glueball decay to photons should be suppressed via quark loops. Central exclusive production is in general considered to be the best glueball production channel at the LHC [63], pomerons being presumably objects with a high gluon content. Technically speaking all f -mesons are glueball candidates at some level, some more than others because of their position on the pomeron Regge trajectory such as $f_2(1950)$ or because of lattice QCD estimates, which is the case for example with $f_0(1710)$. This way we see that it is just natural to say

Central Production Spectrum \simeq Glueball Candidate Spectrum.

Also, we point out that there is no such strict spin-parity rule as $J^{PC} = 0^{++}, 2^{++}, \dots$ in central production per se, quite often erroneously claimed. It is just that the *decay channels* which are usually measured, provide no other options by conservation laws. One may well produce resonance states with other quantum numbers too such as

axial vectors $J^P = 1^+$, as shown by the WA102 data [64]. In any case, the non-perturbative production couplings are currently theoretically incalculable.

The first evidence for the existence of double pomeron exchange type processes $pp \rightarrow p + X + p$ was measured at the CERN ISR in the low-energy range $30 \leq \sqrt{s} \leq 62$ GeV [65]. For a review of the pre-LHC era measurements see [66] and for recent measurements see Table 5.1. Our measurements do not include the forward protons but we work rather with the Babinet's complement: a forward veto type event selection. The event signature is a striking one for a high energy hadron collider: we select events with two reconstructed charge opposite tracks and require otherwise an empty detector. This semi-exclusive event selection will include also events with one or both of the forward protons excited into a low-mass state.

4.1 DREM-PID: Double Recursive Expectation Maximization Particle Identification

For the particle identification, we developed a new fast maximum likelihood algorithm that we dubbed DREM-PID. It is very simple and thus, similar to the previous algorithms such as described in [67]. However, ours is based on a double recursive (iterative) utilization of the classic Expectation Maximization (EM) algorithm, first for single particle tracks and then for particle combinations such as final state pairs. The obtained probabilities per event can be used as weights or as cuts with a suitable type I and type II frequentist error working point, also known as the ROC (receiver operating characteristics) working point. The weighted approach is the most streamlined regarding the PID related efficiency-purity corrections, in principle, no misclassification corrections are needed as long as all relevant particle pair combinations are included in the scheme, the fiducial momentum range is limited enough and the detector likelihood functions are high precision enough. In practise, some residual corrections may be necessary. Different strategies should be cross-compared using pure Monte Carlo input. We shall again remark that the EM-algorithm uses the Bayes formula, but is not a Bayesian method.

Phase I: In the first phase, one calculates the single track probabilities by EM-iteration as a function of the track 3-momentum \vec{p} , typically results binned over $|\vec{p}|$, for each three hadronic species π, K, p mass hypothesis. The detector information needs to be described in terms of signal likelihood functions, in the first approximation independently with

$$L_{tot}(\mathbf{x}, \vec{p}|h) = \prod_{i \in \text{detectors}} L_i(\mathbf{x}, \vec{p}|h), \quad (4.1)$$

where $h \in \{\pi, K, p\}$ and the product is over different detectors, such as the ionization dE/dx likelihood from the tracking detectors and the time-of-flight likelihood. If one of the detectors provides no signal, it is simply described with unit likelihood. These likelihoods are based on detector simulations for each particle hypothesis and the \mathbf{x} denotes the set of low-level signal measurements. The ionization obeys a stochastic process with Landau distribution, but a sum of multiple ionization signals obeys Gaussian approximations by central limit theorem. The algorithm can handle different likelihood functions, also multidimensional beyond the product likelihood. The EM iteration obtains the particle fractions with

$$P(h|\mathbf{x}, \vec{p}) \leftarrow \frac{L_{tot}(\mathbf{x}, \vec{p}|h)P(h|\mathbf{x}, \vec{p})}{\sum_k L_{tot}(\mathbf{x}, \vec{p}|k)P(k|\mathbf{x}, \vec{p})}, \quad (4.2)$$

which is iteratively updated starting with uniform particle fractions until convergence. This is repeated independent for each kinematic hyperbin depending on \vec{p} .

Phase II: In the second phase for the charge identified and \pm ordered tracks, one tabulates all 3^N hadronic combinations and uses the EM-iteration to extract the relative abundances of these by iterating

$$P(\{h\}_c) \leftarrow \frac{\prod_{i=1}^N P(h_c^{(i)} | \mathbf{x}_i, \vec{p}_i) P(\{h\}_c)}{\sum_{k \in \text{combs}} \prod_{i=1}^N P(h_k^{(i)} | \mathbf{x}_i, \vec{p}_i) P(\{h\}_k)}, \quad (4.3)$$

which yields the event-by-event probabilities for each of the final states.

A concrete example with $3^2 = 9$ combinations for $+ -$ ordered pairs yields

$$\{\{\pi^+ \pi^-\}, \{K^+ K^-\}, \{p \bar{p}\}, \{\pi^+ K^-\}, \{\pi^+ \bar{p}\}, \{K^+ \pi^-\}, \{K^+ \bar{p}\}, \{p \pi^-\}, \{p K^-\}\}. \quad (4.4)$$

Only the first three are physical for fully exclusive central production, unless strangeness or baryon number is violated. However, in a semi-exclusive process all are easily possible. The relative abundances of these channels yields also a handle on the exclusivity of the data sample. Four particles gives already $3^4 = 81$ combinations, which is no problem when the combinations are generated algorithmically. When the track has large $|\vec{p}|$, the particle identification based on dE/dx will fail with any algorithm, only a good resolution time of flight measurement or for example Cherenkov radiation based measurement can then extend the identification momentum range. It is $|\vec{p}| \approx 0.6$ GeV when $\langle dE/dX \rangle$ of the TPC does not yield anymore strong discrimination between pions and kaons. For protons, this goes up to $|\vec{p}| \approx 0.9$ GeV. When the TOF signal is reconstructed, the kaon discrimination will extend approximately up to 2 GeV and for protons up to 4 GeV.

Basics of binary decision theory

We shall here point out a couple of crucial things about type I and type II errors, not always explained in a concrete manner. Let us first define these frequentist errors in Table 4.1. The *power* of the statistical test is given by $1 - \beta$ whereas the statistical *significance* level is set by α .

Let us measure an observable x . The frequentist framework works purely with the likelihood functions $\mathcal{L}(x; H_i)$ of the signal and background and the optimal frequentist decision is given by the likelihood ratio via the Neyman-Pearson lemma [68]

$$R(x) = \frac{\mathcal{L}(x; H_0)}{\mathcal{L}(x; H_1)}, \quad \text{s.t. } \alpha = P(R(x) \leq c), \quad (4.5)$$

	Background H_0	Signal H_1
Accept	Type I error: $P = \alpha$ (false pos.)	$P = 1 - \beta$ (true pos.)
Reject	$P = 1 - \alpha$ (true neg.)	Type II error: $P = \beta$ (false neg.)

 Table 4.1: Frequentist decision theory probabilities P .

which yields the most powerful test with maximum true positive probability $1 - \beta$, that one can construct at the fixed significance level α . Algorithm: First one chooses the significance level of interest such as $\alpha = 0.05$, then one finds out the cut parameter c such that significance level holds. This can be done numerically with simulated likelihoods. The resulting decision domain can be very complicated already in one dimension $x \in \mathbb{R}$, if the likelihoods are complicated functions. The likelihood functions can be analytic or estimated by neural techniques, for example.

The likelihood ratio is a binary decision optimality with no prior information. The Bayes factor ratio of posteriori probabilities is the corresponding Bayesian version

$$\frac{P(H_0|x)}{P(H_1|x)} = \frac{\mathcal{L}(x; H_0) P(H_0)}{\mathcal{L}(x; H_1) P(H_1)}. \quad (4.6)$$

By using the Maximum a Posteriori (MAP) decision, one reaches the so-called Bayes error which is the irreducible probability bound of misclassifications. A fully Bayesian approach will integrate over prior distributions, usually done with Monte Carlo sampling. One must understand that the PID classification efficiency and PID background (misclassification) corrections are different and have a different type of systematics with the frequentist and Bayesian philosophies. Then, in contrast, our fully weighted approach (which can be either frequentist or Bayesian) does not strictly speaking belong into the framework of hard binary decision theory but should be considered as a complete regression (fit) algorithm.

4.2 F2X: A Faster Analysis of Cross Sections

For the central production, but not limited to, we designed a new radically simple analysis chain driven by our algorithms. This minimal approach tries to mitigate the ever increasing abstraction and complexity of high energy physics analyses. The chain is as follows.

F2X in a nutshell

1. Tracking, fiducial cuts and vetoes
2. DREM-PID final state identification weights
3. DEEPEFFICIENCY efficiency inversion weights
4. Luminosity weight
5. Data(!) events in HepMC3 containers
6. Event by event weighted analysis using only fiducial final state information

The idea is very simple: all experimental procedures are done event by event with probabilities. Final state identifications are executed with the DREM-PID algorithm and efficiency corrections with the DEEPEFFICIENCY algorithm of Chapter 8. Finally, events together with the weights are stored in HepMC3 containers which allow to treat the data and Monte Carlo events through the same analysis and plotting machinery. In practice one may work with weighted histograms of observables. Also, only final state information is used so any internal Monte Carlo information about intermediate states or propagators, parton showers etc. is explicitly forbidden within this philosophy. To keep the measurement as a measurement. This philosophy is the same as with RIVET, but we go beyond simple histogram containers.

In principle, also negative events weights could be utilized e.g. for background corrections. Corrections that cannot be implemented with event-by-event weights such as background corrections can be done a posteriori with histogram subtractions, as usual. This type of corrections may be ‘feed-down’ of higher multiplicity final states to lower ones, for example, by detector efficiency losses. One must be careful in classifying the algebraic type of corrections either into a multiplicative or additive class. This type of analysis also works as a repository container of the data and its corrections.

CHAPTER 4. ALICE AND *glueballs*

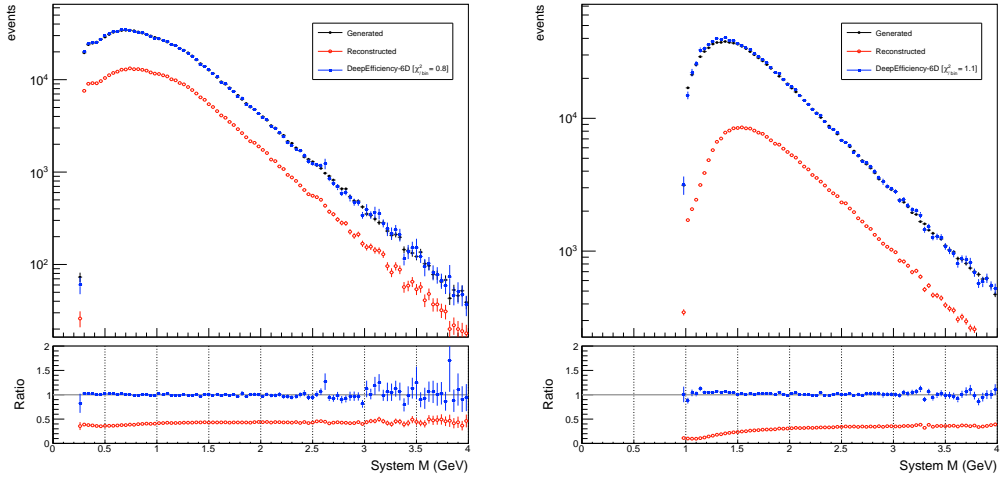


Figure 4.1: ALICE simulation: DEEPEFFICIENCY-6D based efficiency inversion of $\pi^+\pi^-$ and K^+K^- pairs. The reconstructed 1D-observable is the pair invariant mass.

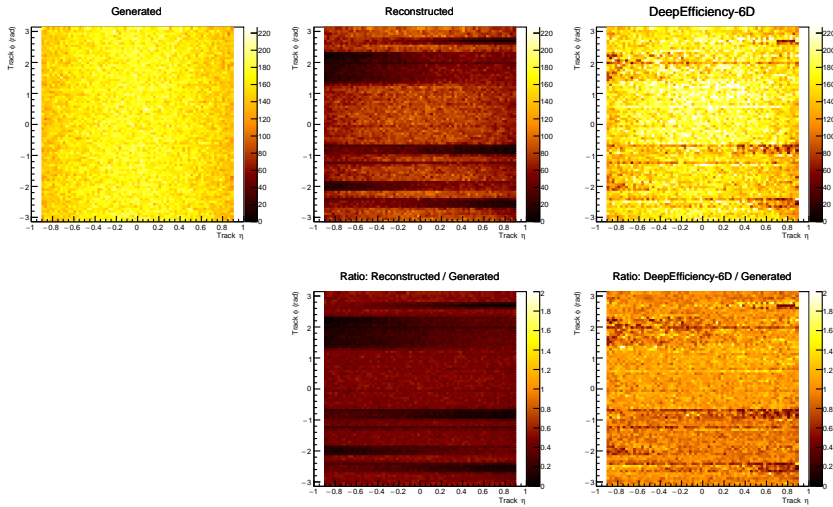


Figure 4.2: ALICE simulation: DEEPEFFICIENCY-6D based efficiency inversion of K^+K^- pairs. The reconstructed 2D-observable is the positively charged track (η, ϕ).

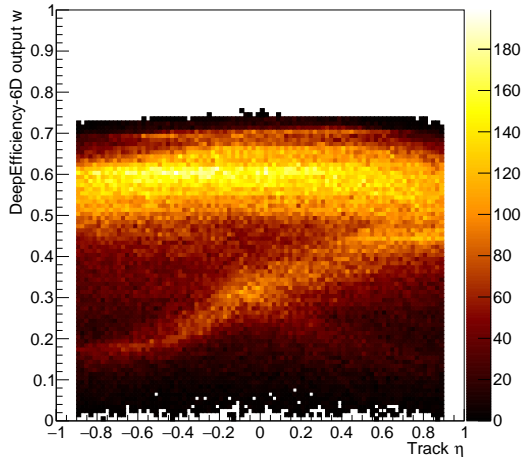


Figure 4.3: ALICE simulation: DEEPEFFICIENCY-6D event-by-event weight distribution projected for a single track pseudorapidity of the track pair.

We illustrate the DEEPEFFICIENCY algorithm efficiency inversion in Figures 4.1, 4.2 and 4.3 trained with approximately 5 million full GEANT simulated events. The 6D-input is the 3-momentum of two tracks which contains thus all correlations. The test samples are different than the training samples both statistically and systematically, by using slightly different MC generator level distributions. The inversion yields reliable results and the algorithm learns the very sharp discontinuities caused by the large inactive regimes in the SPD layers already with the modest MC training sample sizes. In these simulations we use the fiducial domain definition of $|\eta| < 0.9$ and $p_t > 0.15$ GeV, which yields nearly flat performance for the pion pairs and slightly worse for the low-mass kaon pairs. We estimate that with larger samples and the network architecture design optimizations the results will reach asymptotically near perfect inversion. However, in practice, one needs to study the performance through extensive simulations for the bias and variance. DEEPEFFICIENCY provides simultaneously a parametrization for fast detector efficiency simulations beyond the inversion application. It also provides efficiency corrections *simultaneously in all* observables. By using the EM-iteration, we could extend the algorithm to high dimensional unfolding context. However, a direct event-to-event inversion is not possible, because the unfolding mapping is not one-to-one. Thus, one needs to Monte Carlo sample densities, which poses new technical challenges due to new high dimensional density estimation techniques needed such as ‘normalizing flows’ [69].

4.3 Experimental tracking and PID setup

Tracking is based on the inner tracking silicon (ITS) detector, described earlier and the time projection chamber (TPC) [70] with a geometric acceptance of $|\eta| < 0.9$. The TPC is presumably the largest ionization gas chamber in the world currently with multi-wire proportional chambers (MWPC) end plate read out providing the cathode planes with 560k channels, which are in run III replaced with faster gas electron multipliers (GEM). The original run I gas mixture was 90% Ne and 10% CO₂ with a volume of 88 m³, with some gas mixture modifications done at the run II. The gas is under a 400 V/cm electric field, yielding a drift time of 88 ms, which clearly limits any high collision rate use of the TPC. The electronic charge read out is coupled through low input impedance charge sensitive operating amplifiers and analog pulse shapers (PASA) and then sampled by 10 bit analog-to-digital converters at 5 to 12 MHz sampling frequency. The detector gas volume is calibrated through a built-in UV laser system, to compensate for any non-uniformity in the drift field and electron paths. These corrections are a standard part of the AliROOT event reconstruction chain executed in the computing grid.

Cut	Criterion
Reconstruction	ITS + TPC
TPC cluster per track	> 70
Track χ^2/ndf	< 4
DCA _T	< 7 σ
DCA _z	< 0.5 cm
Vertex position	$ z < 10$ cm

Table 4.2: Tracking low-level cuts, where DCA means ‘distance of closest approach’.

The simultaneous track recognition and fitting algorithm are based on the Kalman filter, which is a recursive linear estimator for sequential measurements. It is an optimal algorithm in the case when the noise is Gaussian and the system is linear. Both are not exactly the case with tracking, however, it is still one of the best algorithms given its relative simplicity and well understood behavior. The alignment of the tracking is based on both cosmic rays and proton-proton collisions using different global and local fitting approaches, for more information see [71]. The warm solenoid magnetic field was operating under its default value of $B = 0.5$ Tesla. With ITS and TPC in use, the mass and transverse momentum resolution in the low mass 2-track measurements was simulated to be better than 0.5%. No unfolding of the

measurement is necessary with such a good resolution. We summarize the tracking cut parameters in Table 4.2.

ALICE has by construction numerous particle identification capabilities. The detectors which provided PID information in this study were the ITS ionization dE/dx and the TPC higher resolution dE/dx and the time-of-flight (TOF) detector for the track velocity $\beta = v/c$ measurements [72]. The TOF is built using multigap resistive plate chamber technology (MRPCs) with a total number of read out channels being 157k. The MRPCs are built using a stack of glass plates working as dielectric resistive material layers with conductive pickup cathodes on top and bottom of the stack, the anode plane set in the middle of the glassburger. In total there are 5 gaps of gas with the width of a gap being 250 μm , on top and bottom of the central anode. The incoming charged particle ionizes the gas in the stack and the charge is amplified using the high electric field resulting in an avalanche. The electrons are not free flowing through highly resistive glass layers, instead, the signal coupling is capacitive. The idea behind multiple layers is to be able to use a very high electric field of 100 kV/cm, resulting in a high signal efficiency together with a good time resolution.

The TOF time resolution in-situ is approximately 80 ps under optimal event conditions with intrinsic resolution of the MRPC elements measured to be approximately 50 ps. The TOF signal ‘matching flow’ efficiency with the ITS + TPC tracks starts to operate at $p_t \approx 0.3$ GeV, reaching approximately 50% efficiency at 0.5 GeV and the efficiency reaches plateau at 60 - 65 % around 1 GeV. To point out, this weak efficiency at low- p_t is naturally handled by probabilistic methods such as our DREM-PID algorithm. When no TOF information is available, the product likelihood simply replaces the TOF term with unity and the overall tracking efficiency is left intact. The TOF measurement requires the interaction start time t_0 , which is usually measured by the dedicated T0 detector of ALICE but in the case of low multiplicity events, it must be provided usually by the TOF itself yielding approximately 10 ps additional smearing on the time resolution compared to the T0 driven case. The third option is to use the LHC clock timing, providing t_0 with the resolution fundamentally limited by the LHC bunch spreading over the z -axis as $\propto \sigma_z/c$. It is an easy exercise to calculate how the (Gaussian) time resolution gives bounds to distinguish between particles of different mass as a function of the momentum $|\vec{p}|$. In the analysis, GEANT simulated signal templates for different particle hypothesis were used, used internally by the AliROOT PID libraries.

4.4 Semi-exclusive event selection

The event sample collected at the cms energy of $\sqrt{s} = 7$ TeV is based on a minimum bias trigger with VZERO + SPD and offline ‘double rapidity gap veto’. Also, a much larger sample at the cms energy of $\sqrt{s} = 13$ TeV was collected with a special double rapidity gap trigger with a track candidate trigger requirements in the SPD and TOF detectors. These trigger requirements effectively give a much reduced tracking efficiency at low- p_t with complicated trigger systematics to be included in the future detector simulations, not available at this stage. Thus, we study here only the $\sqrt{s} = 7$ TeV sample with readily excellent efficiency, which is essentially flat over the system mass and transverse momentum in the fiducial domain given in Table 4.3.

Cut	Criterion
Track fiducial	2 tracks with $ \eta < 0.9$, $p_t > 0.15$ GeV and charge $\sum_i Q_i = 0$
Exclusivity veto	Veto in $\eta \in [-3.7, -0.9]$ and $\eta \in [0.9, 5.1]$ (VZERO & FMD)
Track PID	No unassociated tracklets in SPD within $\eta \in [-2.0, 2.0]$ ITS + TPC + TOF MAP decision for both tracks using pp -minbias Bayesian priors

Table 4.3: Event selection cuts.

The low-level event selection was developed during the years by numerous people with codes available under github.com/alispw/Aliphysics/tree/master/PWGUD. The high-level analysis codes of this work are available on my [github](#). In this study, we require Maximum a Posteriori probability PID decision for both tracks (MAP \otimes 2), using kinematics dependent but otherwise fixed pp -minimum bias priors for the particle fractions. With larger statistics one could try to EM-fit the fractions from data, as described earlier. We compared the results with the fully weighted PID, with nearly identical mass distributions shapes and event counts. This is an expected result given the good performance of the PID for low-momentum tracks which dominate our sample.

The purity of the sample was studied counting the number of events using the exclusive selection, but requiring the charge sum $Q = \pm 2$. This was measured to be $< 1.5\%$ of the full sample after veto and track cuts, indicating good tracking and veto efficiency within the central domain, but is not a fully conclusive statement. Based on this, we may gain an estimate of the ‘feed-down’ background of $Q = 0$ from high multiplicity central production. For example, a system X with four charged pions with two lost tracks due to inefficiency will yield four possible sum $Q = 0$

combinations and $Q = \pm 2$. The four $Q = 0$ combinations will, unfortunately, pass the selection. With larger statistics, one could subtract this using the $Q = \pm 2$ spectrum multiplied by two. One could also in a data-driven way try to mathematically invert the ‘complete ladder’ of feed-down, by using the measured 2,4,6,...-track events with different final states. This requires solving a large system of possibly non-linear equations with some coefficients needed from Monte Carlo.

Final state	MAP \otimes 2	(%)	Weighted	(%)
$\pi^+\pi^-$	32608	(90.234)	32352.8	(88.776)
K^+K^-	958	(2.651)	1003.6	(2.754)
$p\bar{p}$	119	(0.329)	119.4	(0.328)
π^+K^-	1027	(2.842)	1233.1	(3.384)
$\pi^+\bar{p}$	163	(0.451)	190.5	(0.523)
$K^+\pi^-$	1097	(3.036)	1336.4	(3.667)
$K^+\bar{p}$	14	(0.039)	19.9	(0.055)
$p\pi^-$	136	(0.376)	166.9	(0.458)
pK^-	15	(0.042)	20.3	(0.056)
Σ	36137		36443.0	

Table 4.4: ALICE data at $\sqrt{s} = 7$ TeV: Integrated final state event counts with two different PID selection strategies.

The total number of $Q = 0$ events are shown in Table 4.4. The number of seemingly strangeness conservation ‘violating’ pairs of π^+K^- or π^-K^+ is thus approximately the same as K^+K^- pairs. These lonely kaons may be for example due to efficiency losses (feed-down) or forward excitation, which can leak them from the forward system fragmentation. Also, these may come from random minimum bias processes generating large rapidity gaps, even if with exponentially suppressed probability. The states $K^*(892)$ and $K^*(1430)$ which decay dominantly into pion-kaon pairs, were observed in the pion-kaon samples. Interestingly, PDG states that the (glueball candidate) $f_2(1950)$ is seen to branch into $K^*(892)\bar{K}^*(892)$. These decay channels should be measured to construct a clear hierarchy of f -mesons and glueballs. We remark that we see some asymmetry in Table 4.4 at the level of a few standard deviations between the charge conjugate channels, possibly due to material interaction effects and a calibration mismatch. We conclude that extensive MC background studies are needed for the future cross section measurements, with complete efficiency and background corrections, which goes beyond what is shown here. We shall now concentrate on the pion and kaon pair channels spectroscopy.

4.5 System invariant mass spectrum

To study the invariant mass spectrum in detail we implemented a simple, but complete spectrum fit. We point out that theoretically better fits can be done with our GRANITTI MC, where individual phenomenological scattering amplitudes are summed together at amplitude level and evaluated as a function of the Lorentz scalars. The corresponding fit code is part of the package. However, for simplicity we shall fit the spectrum this way. The total spectrum parametrization is

$$f(m) = \left| \sum_k A_k e^{i\phi_k} g(m; M_0^{(k)}, \Gamma_0^{(k)}) + \psi(m) \right|^2, \quad (4.7)$$

where the sum is over resonances with complex weights $A_k e^{i\phi_k}$. We take the sum over four resonances in the case of $\pi^+\pi^-$ pairs and over two resonances in the case of K^+K^- pairs. For the underlying continuum, we take an ansatz

$$\psi(m) = iA \left[\frac{m - D}{\gamma} e^{-\frac{m-D}{\gamma}} \right]^{1/2}, \quad (4.8)$$

where A is the weight, γ is the scale and D is the shift parameter. The resonances are parametrized with relativistic Breit-Wigner poles

$$g(m; M_0, \Gamma_0) = \frac{\sqrt{2m}}{m^2 - M_0^2 + iM_0\Gamma(m)}, \quad (4.9)$$

where M_0 is the pole mass and $\Gamma(M)$ is the momentum dependent width [73]

$$\Gamma(m) = \Gamma_0 \frac{M_0}{m} \left[\frac{m^2 - 4m_f^2}{M_0^2 - 4m_f^2} \right]^{3/2}, \quad (4.10)$$

where m_f is the decay daughter mass. The denominator factor $\sqrt{2m}$ originates from the change of variables $|G(m^2)|^2 dm^2 \mapsto |g(m)|^2 dm$ of the absolute squared density of the relativistic propagator expression $G(m^2) = 1/[m^2 - M_0^2 + iM_0\Gamma]$. This needs to be taken into account, because our expressions are as a function of m .

We remark that the fit exercises with Breit-Wigners are complicated for composite (non-elementary) resonances with varying spins, phase space and production channel effects. No simple parametrization is expected to hold perfectly.

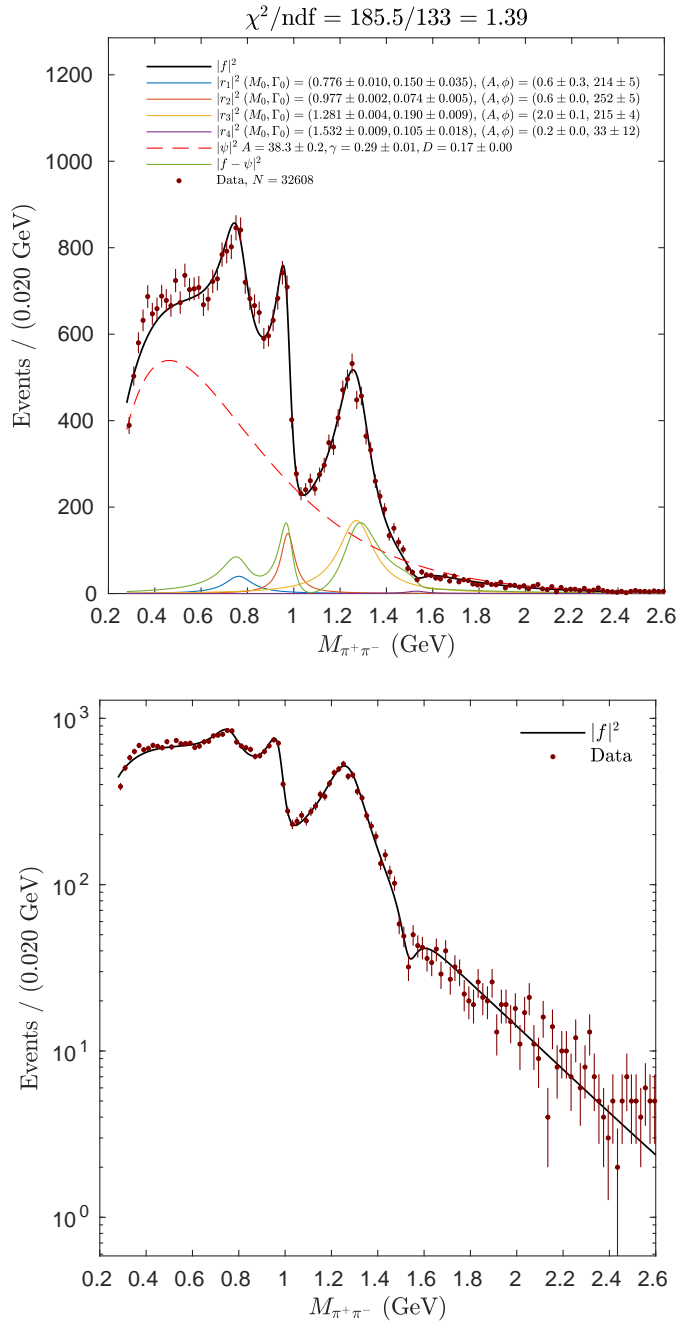


Figure 4.4: ALICE data at $\sqrt{s} = 7$ TeV: The semi-exclusive $\pi^+\pi^-$ spectrum with $|\eta(\pi)| < 0.9$ and $p_t(\pi) > 0.15$ GeV.

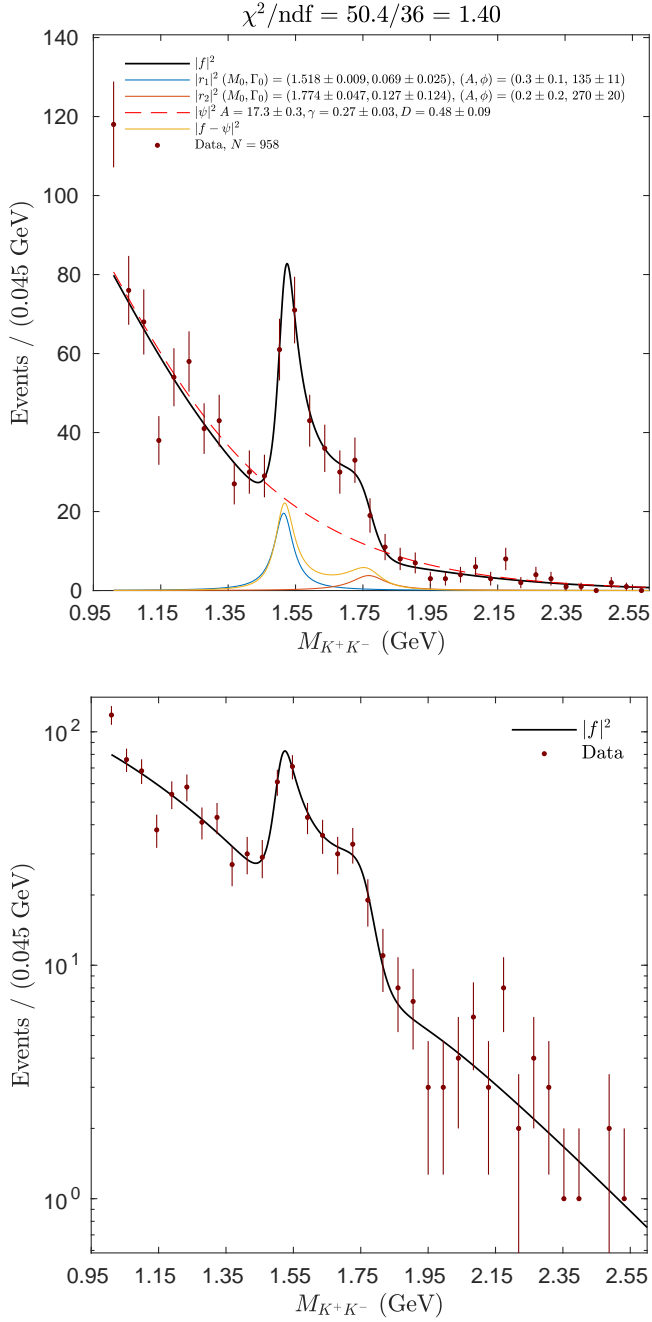


Figure 4.5: ALICE data at $\sqrt{s} = 7$ TeV: The semi-exclusive K^+K^- spectrum with $|\eta(K)| < 0.9$ and $p_t(K) > 0.15$ GeV.

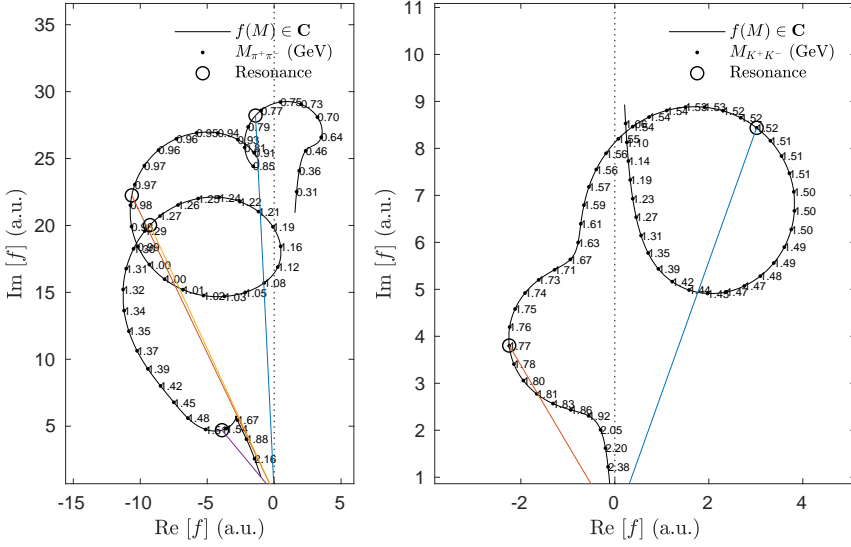


Figure 4.6: ALICE data at $\sqrt{s} = 7$ TeV: The complex plane fit curves in $\pi^+\pi^-$ on the left and in K^+K^- on the right. Colored lines are for visualizing the resonance pole positions.

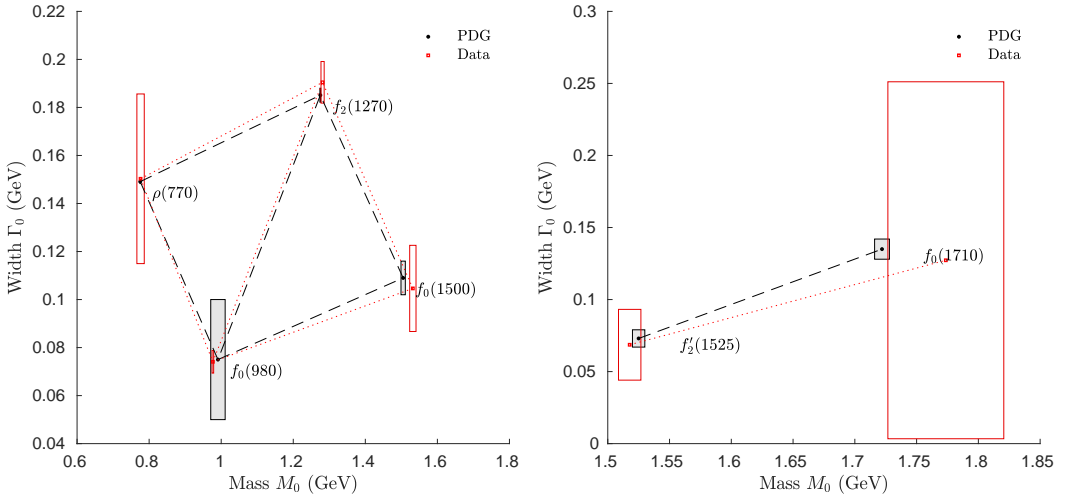


Figure 4.7: ALICE data at $\sqrt{s} = 7$ TeV: The fit extracted resonance parameters with statistical 1σ uncertainties and comparisons with the established PDG resonance candidates in $\pi^+\pi^-$ on the left and in K^+K^- on the right.

The results for the $\pi^+\pi^-$ and K^+K^- spectrum and the corresponding parameter values with statistical uncertainties (phases in degrees) are shown in Figures 4.4 and 4.5, and the extracted resonance mass and width parameters are compared with the closest matching PDG resonance values in Figure 4.7. The fits have near unity χ^2/ndf values. The amplitude for the $f_2(1270)$ candidate is significantly larger than for the other resonances, which can be also seen by fitting the spectra by using GRANITTI scattering amplitudes, so it is not just an effect from this particular fit parametrization. This is interesting given that it is a spin-2 state and not a spin-0 (ground state). The obtained statistical fit uncertainties for the mass and width of the $f_0(980)$ candidate are very competitive, better than the PDG global average. We see a sharp drop in the invariant mass spectrum near $M_{\pi^+\pi^-} \simeq 1$ GeV, which may have something to do with the K^+K^- channel opening and unitarity. Fitting the mass spectrum around that point relies on complex interference. The total fit functions in the complex plane are shown in Figure 4.6.

We observe that the continuum slope γ has the same value for both of the final states within statistical uncertainties, yielding knowledge about the non-resonant scattering dynamics. This could be easily compared with GRANITTI based simulations, by adjusting there the continuum amplitude off-shell form factor parametrizations. It can shed light on the question is the continuum amplitude dominated by a sub- t channel exchange or not. At very low masses, the enigmatic $f_0(500)$ resonance was not manifestly present given that the continuum parametrization alone was able to describe the data. Again, this is something to be studied in more detail with Monte Carlo and spherical harmonic angular decompositions. The identification of possible glueball states from the spectrum requires further systematic studies. Because glueballs are expected to mix with ordinary meson states, to decompose the mixing experimentally especially without forward protons, one needs do simultaneous high statistics measurements of pion, kaon and gamma pair final states. Also the four body decays should be measured, which can be used to probe resonances decaying to intermediate vectors such as $\rho^0(770)$ or $\phi(1020)$ pairs. Our DEEPEFFICIENCY algorithm solves the problem of unknown angular distributions affecting the efficiency corrections, which can be hard to implement otherwise with soft 4-body final states.

4.6 System transverse momentum

The system transverse momentum is a crucial observable, related to the transverse plane Gribov diffusion and the proton incoherent structure fluctuations. In this section we describe a fast Monte Carlo template fit to the system transverse momentum spectrum. As a stochastic distribution driving the fit templates, we use for the elastic coherent proton p_x, p_y -components the Gaussian distribution and for the excited incoherent proton p_t^2 , we use the Gamma distribution ansatz. The forward-backward system transverse angle separation $\Delta\varphi$ is sampled from the uniform distribution, which is the Occam's razor choice without forward proton measurements. Then we get the central system transverse momentum by momentum conservation, and we also use the vertex factorization assumption for different legs in the process. The corresponding distribution parameters together with the process fit fractions f_i are shown in the legend of Figure 4.9, where b is the coherent forward proton exponential p_t^2 -slope related to the Gaussian components by $b = 1/(2\sigma^2)$, as derived in Appendix B.3. The parameter k is the shape of the incoherent Gamma distribution with the scale parameter set as $\theta = 1/b$, and the mean is given by $k\theta$. Effectively, the Gamma distribution includes both exponential and power law (hard tail) behavior.

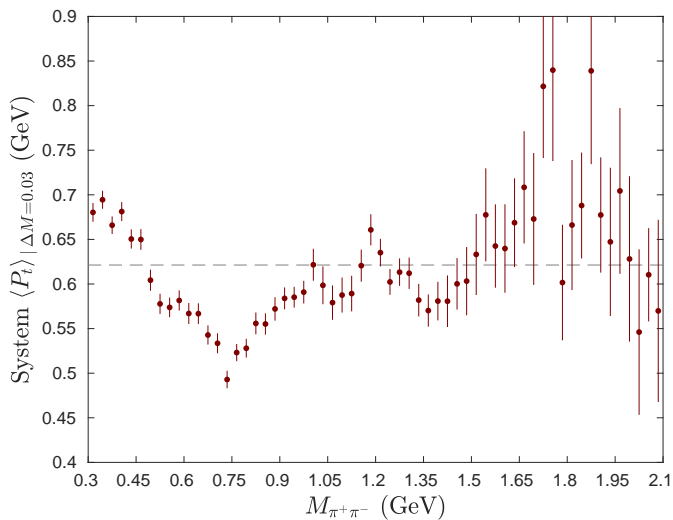


Figure 4.8: ALICE data at $\sqrt{s} = 7$ TeV: The $\pi^+\pi^-$ system mean transverse momentum as a function of the invariant mass with $|\eta(\pi)| < 0.9$ and $p_t(\pi) > 0.15$ GeV. The dashed line is the mean transverse momentum over the full mass range.

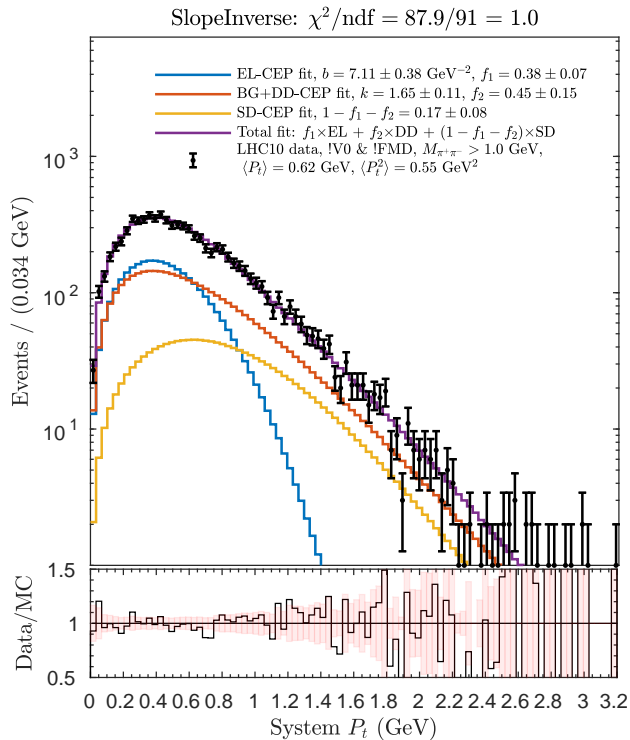


Figure 4.9: ALICE data at $\sqrt{s} = 7 \text{ TeV}$: The $\pi^+\pi^-$ system transverse momentum fit with $M_{\pi^+\pi^-} > 1 \text{ GeV}$, $|\eta(\pi)| < 0.9$ and $p_t(\pi) > 0.15 \text{ GeV}$.

We use a minimum mass cutoff at 1 GeV, to remove the contribution from the photoproduction of ρ^0 . The mean transverse momentum $\langle P_t \rangle$ as a function of the invariant mass is shown in Figure 4.8. We observe a clear dip a bit below the ρ^0 pole mass, which is a strong indication for the photoproduction origin. The figure shows also up and down variation near other resonances, which is expected. Now based on the fit shown in Figure 4.9, we obtain the fraction of events with elastic forward protons

$$f_{el} = 0.38 \pm 0.07 \text{ (stat)} \quad (4.11)$$

using the event selection of Table 4.3. The rest are inelastic events with one or two protons dissociated, which still pass through the forward veto. This ratio is compatible with GRANITTI based simulations given in Chapter 5, which are based on a more sophisticated phenomenology, with MC parameters fixed by completely

different measurements. We found out that the ratio between single and double dissociation is strongly dependent on the model details, such as the vertex factorization and the incoherent stochastic distribution ansatz. It can be studied reliably only with larger event statistics. This fraction extraction is presumably done for the first time at the LHC, at least with the ALICE data. The fraction depends naturally on the forward veto acceptance of the experiment, thus a careful fiducial calibration should be done.

The obtained exponential elastic slope is

$$b = 7.11 \pm 0.38 \text{ (stat) GeV}^{-2}. \quad (4.12)$$

The value is realistic and gives confidence to the elastic fraction estimate of the sample. Systematic errors in these are essentially a fit model driven, something which needs to be studied and simulated with full GRANITTI based fits in the future, which allows one also to study e.g. the effects of screening corrections. For a comparison, the ATLAS+ALFA experiment obtained a value of $b = 7.65 \pm 0.34 \text{ GeV}^{-2}$ in single diffraction measurements at $\sqrt{s} = 8 \text{ TeV}$ [54].

4.7 Summary

The main results of the studies with semi-exclusive pp data with identified pion and kaon pair final states at the cms energy $\sqrt{s} = 7$ are as follows.

- ◇ The $\pi^+\pi^-$ continuum and the enigmatic $f_0(500)$ are ambiguous, that is, the simple spectrum fit cannot uniquely separate these. There is a sharp fit interference at $M_{\pi^+\pi^-} \simeq 1$ GeV, also visible by eye in the sharp drop of the invariant mass spectrum. Perhaps related to the K^+K^- channel opening and unitarity.
- ◇ The single dimensional complex spectrum fit is easily unstable in way that a heavy functional dependence was observed with the continuum parametrization ansatz. However, within statistical uncertainties, the same continuum parametrization slope γ was extracted for the pion and kaon channels, interesting for the phenomenology.
- ◇ The $f_0(980)$ mass and width extraction seems to yield significantly smaller (statistical) uncertainties than the PDG global average.
- ◇ The fit amplitude A_k of $f_2(1270)$ is a factor of a few larger than for the other resonances, a similar observation can be made with GRANITTI MC based fits. This is somewhat puzzling.
- ◇ Photoproduction of $\rho(770)$ is visible, which was verified also qualitatively observing a dip in the system $\langle P_t \rangle$ distribution, also $\phi(1020)$ seems to be visible given the good low- p_t efficiency.
- ◇ Glueball candidates such as $f_0(1710)$ require more statistics to make systematic branching ratio comparisons feasible. Technically, all f -mesons are at least mixed candidates for glueballs.
- ◇ The visible fraction of events with elastic vs inelastic central production and the corresponding elastic proton b -slope in the process were extracted through fast Monte Carlo template fits, with values compatible with other measurements and phenomenology.

5 GRANIITTI: A Monte Carlo Event Generator for High Energy Diffraction

We describe the physics and computational power of GRANIITTI Monte Carlo event generator, a new fully multithreaded engine designed for high energy diffraction, written in modern C++. The emphasis is especially on the low-mass domain of central exclusive processes of the S -matrix, where exotic QCD phenomena such as glueballs are expected and holographic dualities with gravity may be tested. The generator includes photon-photon, photon-pomeron, pomeron-pomeron, Durham QCD model and Tensor pomeron model type scattering amplitudes and advanced spin analysis tools. The generator combines a full parametric resonance spectrum with continuum interference and forward + central spin correlations together with event-by-event eikonal screening loop and forward proton excitation kinematics. We demonstrate the state-of-the-art capabilities and show how the enigmatic ‘glueball filter’ observable is driven by spin polarization.

Chapter in: [arXiv:1910.06300](https://arxiv.org/abs/1910.06300) [hep-ph]

5.1 Introduction

High energy diffraction probes ‘asymptotic’ strong interactions in the kinematic domain of small invariant momentum transfer $-t$ and large energy squared $s \rightarrow \infty$ or small x . The main processes of interest here are

$$pp \rightarrow p^{(*)} + X + p^{(*)}, \quad (5.1)$$

where the system X is produced via color and charge singlet exchanges between two protons, via photon or pomeron fusion. These enigmatic objects known as pomerons were constructed originally in the complex angular momentum theory of scattering amplitudes $\mathcal{A}(s, j)$, a method devised by Regge first in a non-relativistic scattering context [74] and then extended to the relativistic domain by Gribov, who also built the most intense but incomplete interacting pomeron calculus [75]. It is clear that one cannot currently consider these asymptotics as being properly explained by current QCD in the soft domain. But it is the future theory together with novel measurements which should eventually provide the detailed ‘space-time picture’ with fluctuating proton structure.

This ‘Pomeranchuk’ singularity can be a simple soft pole or a mathematically almost arbitrary complex set of cuts, such as described in the pioneering work by Mandelstam [76]. Later, a hard pomeron as a fixed branch point singularity was found in the field theory by Balitsky, Fadin, Kuraev and Lipatov [77] (BFKL). In general, this property of field theories is known as ‘Reggeization’ of the strongly interacting amplitudes in the Regge asymptotic. In gauge/string duality, the soft and hard pomeron have been unified [78]. Well known is that the Regge amplitudes behaving like $\propto s^{\alpha(t)}$ have a Gribov diffusion picture, obtained via a Fourier transform to the impact parameter space. They are directly related to the quantum mechanical stochastic branching random walk. An interesting discussion of the parton model, strings and black holes by Susskind is fundamental in this context [79], after all, we are talking about asymptotics. The elastic scattering asymptotics of the proton are fascinating in terms of the evolution of the ‘blackness’ of the disc and the transverse size of this disc, noting that we are now talking about average properties.

In our main process illustrated in Figure 5.1, the forward protons may stay coherent which is known as the central exclusive process or fluctuate into a dissociated state, denoted with p^* , giving semi-exclusive processes experimentally. The dissociation will result in a higher transverse momentum for the system X , experimentally one of the only signs of low mass dissociation bypassing the forward veto detectors. Other indirect observables may be related to the central system spin polarization and

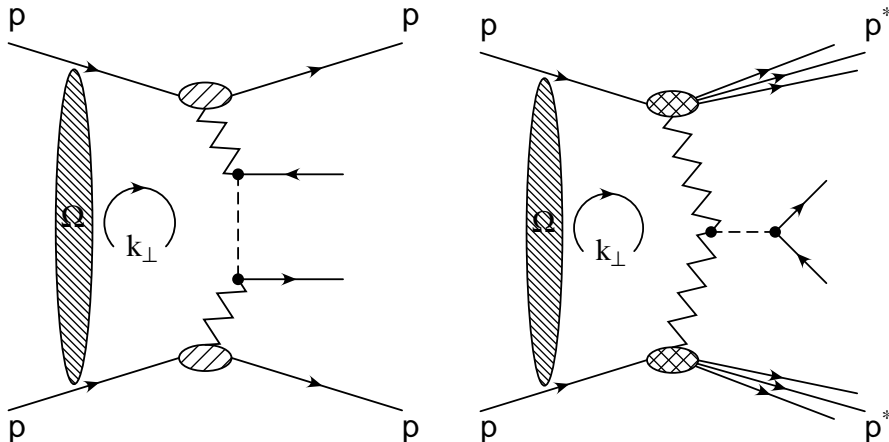


Figure 5.1: A pomeron-pomeron continuum amplitude on the left and a pomeron-pomeron resonance production process with eikonal Ω screening and double dissociation of the forward protons on the right.

relative magnitude of resonances. The emphasis here is in the non-perturbative domain of strong interactions, however, also certain perturbative QCD, QED and EW processes are included. The design of the engine is based on a synthesis \leftrightarrow analysis philosophy, where the analysis side is accelerated via automated phenomenology of fiducial observables and spin polarization decomposition algorithms. The purpose of the GRANIITI Monte Carlo event generator is to provide the necessary tools for the LHC experiments and beyond.

Some of the currently unique features of GRANIITI, to our knowledge, are: central production of arbitrary spin $J = 0, 1, 2, 4, \dots$ resonances and their spin polarization driven helicity amplitudes together with a non-resonant continuum, a full effective Lagrangian Tensor pomeron model implementation and forward proton dissociation combined with a screening (absorption) loop integral. The full potential of the generator is obtained when interfaced together with forward fragmentation such as generated by PYTHIA [80], especially interesting and useful in the case of ALICE and LHCb experiments without forward protons. Experiments with forward protons such as CMS+TOTEM, ATLAS+ALFA or STAR will be highly interesting due to the helicity driven ‘glueball filter’ type forward-central correlations. Thus both types of measurements will be crucial, their cross sections bootstrapping the theory together like a generalized Babinet’s principle.

CHAPTER 5. GRANIITI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

Outline This article is organized as follows. We first describe dynamics involved, then kinematics, analysis algorithms and briefly the technology behind the generator. Finally, we give some conclusions.

5.2 Dynamics

The scattering amplitudes currently included span from QED, EW, Regge theory and QCD. Advanced users can add more amplitudes e.g. from MadGraph, the normalization conventions between the phase space and matrix element squared obey the standard field theory conventions.

Eikonal pomeron

The soft pomeron is an enigmatic object in terms of QCD, however, it is the best effective description of diffraction or asymptotically slowly growing total cross sections so far. The simple eikonal pomeron model we use is based on the model from [81], however our eikonalization procedure is slightly different. The single pomeron exchange elastic pp -scattering amplitude is the basic building block

$$\mathcal{A}_{el}(s, t) = [g_{ppP}F(t)]^2 \eta(\alpha_P(t), +) \left(\frac{s}{s_0} \right)^{\alpha_P(t)}, \quad (5.2)$$

with the proton-proton-pomeron coupling g_{ppP} (GeV^{-1}) and the proton form factor parametrization

$$F(t) = \frac{1}{1 - t/a_1} \frac{1}{1 - t/a_2}, \quad (5.3)$$

with two free parameters a_1, a_2 (GeV^2). The coupling and form factor represent together a factorized real valued residue function at the vertex. The pomeron propagator is $(s/s_0)^{\alpha(t)}$ with the trajectory $\alpha_P(t)$ and the (even +) signature phase factor for the pomeron

$$\eta(\alpha(t), \pm) = -\frac{1 \pm \exp(-i\pi\alpha(t))}{\sin(\pi\alpha(t))}, \quad (5.4)$$

which comes from the Sommerfeld-Watson transform [82]. It is an analytical integral continuation of the integer spin partial wave series – the mathematical basis behind the Regge pole exchanges with $\alpha(t) \sim J$. This expression has trivial singularities (poles) at integer values of $\pi\alpha(t)$, thus $t \rightarrow 0$ limit is sometimes taken. The elastic amplitude gives us, by the optical theorem, the total cross section behavior $\sigma_{tot} \sim s^{\alpha_P(t)-1}$. Experimentally at high energies, the intercept $\alpha_P(0)$ needs to be larger than 1. We use a non-linear parametrization of the pomeron trajectory with a pion loop, described in detail in [81].

We note here that it is the signature factor which gives the ratio between real and imaginary part of the amplitude. In order to proceed with the eikonal model, the

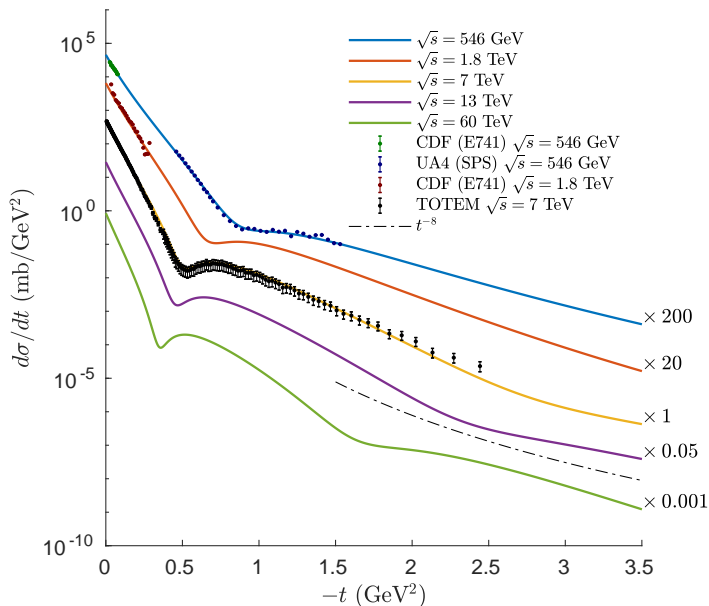


Figure 5.2: Elastic proton-proton scattering with the eikonal pomeron. Data from [83–85].

complex density in the impact parameter b_t -space is obtained by numerical (inverse) Fourier 2D-transform with proper normalization

$$\Omega(s, b_t) = \frac{1}{(2\pi)^2 s} \int \mathcal{A}_{el}(s, -\mathbf{k}_t^2) e^{i\mathbf{b}_t \cdot \mathbf{k}_t} d^2 \mathbf{k}_t \quad (5.5)$$

$$= \frac{1}{2\pi s} \int_0^\infty \mathcal{A}_{el}(s, -\mathbf{k}_t^2) J_0(bk_t) k_t dk_t, \quad (5.6)$$

where J_0 is the 0-th Bessel function of the first kind – that is, we used the Fourier-Bessel transform because there is no azimuthal dependence. Now under the eikonal approximation, the multiple re-scattering amplitude which is an approximation to the more complicated cut singularities than just a single pomeron pole exchange, is given by exponentiation

$$\mathcal{A}_{el}^{eik}(s, b_t) = i(1 - \exp(i\Omega(s, b_t)/2)). \quad (5.7)$$

Then finally, we get the eikonalized amplitude back in momentum space with a

(forward) Fourier 2D-transform

$$\mathcal{A}_{el}^{eik}(s, t = -\mathbf{k}_t^2) = 2s \int \mathcal{A}_{el}^{eik}(s, b_t) e^{-i\mathbf{b}_t \cdot \mathbf{k}_t} d^2\mathbf{b}_t \quad (5.8)$$

$$= 4\pi s \int_0^\infty \mathcal{A}_{el}^{eik}(s, b_t) J_0(bk_t) b_t db_t. \quad (5.9)$$

We calculate these integrals automatically for each cms energy \sqrt{s} , generate lookup tables and interpolate the values event-by-event. The eikonalized amplitude is used to generate elastic scattering events and in the screening loop for the central production.

Now using unitarity (optical theorem), we get total, elastic and inelastic cross sections as

$$\sigma_{tot}(s) = 2 \int \text{Im} \mathcal{A}_{el}^{eik}(s, b_t) d^2\mathbf{b}_t \quad (5.10)$$

$$\sigma_{el}(s) = \int |\mathcal{A}_{el}^{eik}(s, b_t)|^2 d^2\mathbf{b}_t \quad (5.11)$$

$$\sigma_{inel}(s) = \int 2\text{Im} \mathcal{A}_{el}^{eik}(s, b_t) - |\mathcal{A}_{el}^{eik}(s, b_t)|^2 d^2\mathbf{b}_t. \quad (5.12)$$

The integrals are taken numerically using $\int \dots d^2\mathbf{b}_t \rightarrow 2\pi \int \dots b_t db_t$. We point out here that the ‘optimal’ pomeron trajectory parameters are significantly different within the eikonalized amplitude and non-eikonalized single exchange. This means that one simply cannot take the eikonalized pomeron and use that naively with central production amplitudes, for which we use the effective linear pomeron $\alpha_P(t) \simeq 1.08 + 0.25t$.

The characteristic shrinking of the ‘diffractive cone’ is shown in Figure 5.2, which comes out naturally and was originally predicted by Gribov. The dip structure of data is reproduced. We remark here that at high $-t$, a secondary dip develops due to the eikonalization procedure, which is a well known feature with a single Pomeron amplitude. Depending on the non-linear pomeron trajectory parameters, the proton form factor and the fit dataset extension to higher values in $-t$, it can be tuned to flatten out at LHC energies and to appear only at larger energies. Thus using it as a test statistic or model separation tool is not completely conclusive. From perturbative QCD at large values of $-t$ one expects a power law t^{-8} dependence. This amplitude is driven by 3-gluon exchange mechanism, with gluons between three different quarks [86, 87], which is a $C = -1$ exchange. See [87] for discussions about the expected dip behavior in pp and $p\bar{p}$.

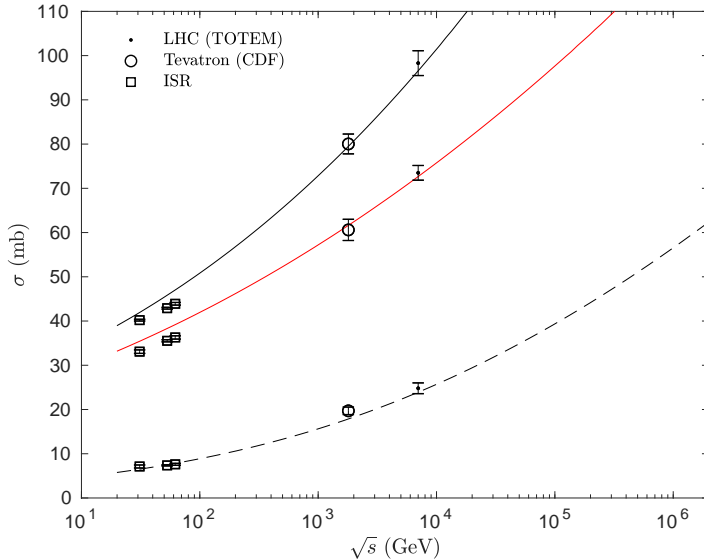


Figure 5.3: The eikonal pomeron based calculations of the total (black), inelastic (red) and elastic (dashed) proton-proton cross sections. Data from [88].

The behavior of total, inelastic and elastic cross section is shown in Figure 5.3. We see that at ISR energies the inelastic cross section, and thus the total, are slightly higher than what is being calculated after parameter tuning, indicating that additional degrees of freedom are important. This is well known phenomenologically, given that Reggeon trajectories have $\alpha_R(t) \simeq 0.5 + 0.9t$, a vanishing role asymptotically due to intercept $\alpha_R(0) < 1$, but a significant at lower energies. Our engine includes automated fit machinery to improve, simplify or extend the eikonal pomeron description and fit more extensive datasets.

Unitarity via screening loop

The screening loop can be turned on for all processes. The basic idea is that the screening loop will induce ‘cut contributions’, via the eikonal pomeron, to the central production which makes the total scattering amplitude unitary and in effect significantly suppresses the cross section. The screening integral is calculated event-by-event with 4-momentum conservation along vertices. The bare amplitude + loop

amplitude give the total screened amplitude

$$\mathcal{A}^S(s, \mathbf{p}_{t,1}, \mathbf{p}_{t,2}) \equiv \mathcal{A}^0(s, \mathbf{p}_{t,1}, \mathbf{p}_{t,2}) + \mathcal{A}^{loop}(s, \mathbf{p}_{t,1}, \mathbf{p}_{t,2}), \quad (5.13)$$

where for notation purposes, we leave out the dependence on the other kinematic variables, only implicitly affected by the loop integral. We do reconstruct the full kinematics for each discrete loop evaluation point.

As a high energy limit, the loop integral is done numerically in the 2D-transverse momentum \mathbf{k}_t -space

$$\mathcal{A}^{loop}(s, \mathbf{p}_{t,1}, \mathbf{p}_{t,2}) = \frac{i}{s} \int \frac{d^2\mathbf{k}_t}{8\pi^2} \mathcal{A}_{el}^{eik}(s, -\mathbf{k}_t^2) \mathcal{A}^0(s, \mathbf{p}_{t,1}, \mathbf{p}_{t,2}), \quad (5.14)$$

where $-\mathbf{p}_{t,1} + \mathbf{k}_t = \mathbf{q}_{t,1}$, $-\mathbf{p}_{t,2} - \mathbf{k}_t = \mathbf{q}_{t,2}$ with $\mathbf{q}_{t,1,2}$ being the transverse momentum transferred to the central system and $\mathbf{p}_{t,1,2}$ the transverse momentum of forward systems (protons). The screening has the largest differential impact to the forward observables, as illustrated in Figure 5.4.

Now relatively straightforward but by no means a complete generalization here would be to consider a ‘multichannel’ eikonal amplitude, responsible for the proton structure and its excitations. That is, the incoming proton is treated as a coherent superposition of eigenstates $|p\rangle$, $|p^*\rangle$, $|p^{**}\rangle \dots$ and the eikonal screening is calculated over the set of pair of states and the corresponding eikonal densities. This approach for the screening loop is in use in DIME [89] and SUPERCHIC [90] event generators. This shall be investigated in the future with more precise data to fit the parameters. Technically, our code includes already some prototype constructions. Perhaps, one could use new approaches for describing the Good-Walker excitations, such as lattice Yang-Mills simulation driven [91], a bit like its application in photoproduction in [92]. A classic approach that does not consider details of the structure is described in [81]. Our minimal description philosophy requires that each new free parameter for the dissociation part should be well motivated by data.

In Table 5.1 we show how the screening loop suppresses the cross sections for processes with a different type of scattering amplitudes and compare with experimental measurements. Clearly, photon-photon processes producing lepton pairs l^+l^- or W^+W^- pairs are more peripheral in \mathbf{b}_t -space with smaller average transverse momentum whereas pomeron-pomeron production of pion pairs results in larger average absorption

$$\langle S^2 \rangle \equiv \frac{\int d\Pi_N |\mathcal{A}^0(\Pi_N) + \mathcal{A}^{loop}(\Pi_N)|^2}{\int d\Pi_N |\mathcal{A}^0(\Pi_N)|^2}, \quad (5.15)$$

which is the phase space integrated cross section ratio between screened and unscreened processes. We denote the full abstract set of 4-momentum variables by Π_N

CHAPTER 5. GRANIITTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

		MEASUREMENT		GRANIITTI		
	\sqrt{s}	PHASE SPACE CUTS	value \pm stat \pm syst	σ_S	σ_0	$\langle S^2 \rangle$
gg	13	$ y < 2.5, p_t > 20$ GeV		0.872	14.1 nb	0.06
$\pi^+\pi^-_{EL}$	7	$ \eta < 0.9, p_t > 0.15$ GeV		3.84	19.6 μ b	0.20
$\pi^+\pi^-_{SD}$	7	$ \eta < 0.9, p_t > 0.15, M_1 < 5$ GeV		1.25	9.34 μ b	0.13
$\pi^+\pi^-_{DD}$	7	$ \eta < 0.9, p_t > 0.15, M_{1,2} < 5$ GeV		0.269	3.13 μ b	0.09
$\pi^+\pi^-$	7	$ Y_X < 0.9$		9.27	45 μ b	0.21
$\pi^+\pi^-$	0.2	$ \eta < 0.7, p_t > 0.2$ GeV		1.85	6.5 μ b	0.28
W^+W^-	7	Full 4π		29.6	39 fb	0.76
e^+e^-	7	ATLAS [93]	$0.428 \pm 0.035 \pm 0.018$	0.462	0.494 pb	0.94
$\mu^+\mu^-$	7	ATLAS [93]	$0.628 \pm 0.032 \pm 0.021$	0.738	0.789 pb	0.94
$\pi^+\pi^-$	13	ATLAS (Thesis) [94]	$18.75 \pm 0.048 \pm 0.770$	20	106 μ b	0.19
$\mu^+\mu^-$	13	ATLAS [95]	$3.12 \pm 0.07 \pm 0.14$	3.35	3.49 pb	0.96
e^+e^-	1.96	CDF [96]	$1.6 \pm 0.5 \pm 0.3$	1.65	1.74 pb	0.95
e^+e^-	1.96	CDF [97]	$2.88 \pm 0.57 \pm 0.63$	3.24	3.37 pb	0.96
$\pi^+\pi^-$	1.96	CDF [98]	$\dagger \pm \dagger \pm \dagger$	1.93	8.96 μ b	0.22
$\mu^+\mu^-$	7	CMS [99]	$3.38 \pm 0.58 \pm 0.21$	3.88	4.09 pb	0.95
$\pi^+\pi^-_{EL}$	7	CMS [100]	$26.5 \pm 0.3 \pm 5.12$	11.5	57 μ b	0.20
$\pi^+\pi^-_{SD}$	7	$ y < 2, p_t > 0.2, M_1 < 5$ GeV		3.77	28.6 μ b	0.13
$\pi^+\pi^-_{DD}$	7	$ y < 2, p_t > 0.2, M_{1,2} < 5$ GeV		0.851	9.4 μ b	0.09
$\pi^+\pi^-_{EL}$	13	CMS [101]	$19.0 \pm 0.6 \pm 3.2$	14.6	80.2 μ b	0.18
$\pi^+\pi^-_{SD}$	13	$ \eta < 2.4, p_t > 0.2, M_1 < 5$ GeV		4.22	34.2 μ b	0.12
$\pi^+\pi^-_{DD}$	13	$ \eta < 2.4, p_t > 0.2, M_{1,2} < 5$ GeV		0.903	11 μ b	0.08

Table 5.1: LHC and Tevatron integrated fiducial cross section measurements compared with GRANIITTI using the screening loop σ_S and without σ_0 . Processes are with elastic forward protons, unless otherwise stated. See the papers for the corresponding fiducial cuts. Note that data contains also non-exclusive background, which experiments have estimated and subtracted to obtain fully exclusive cross section in the case of lepton pairs. † The paper [98] includes a differential cross section in invariant mass, but no integrated cross section.

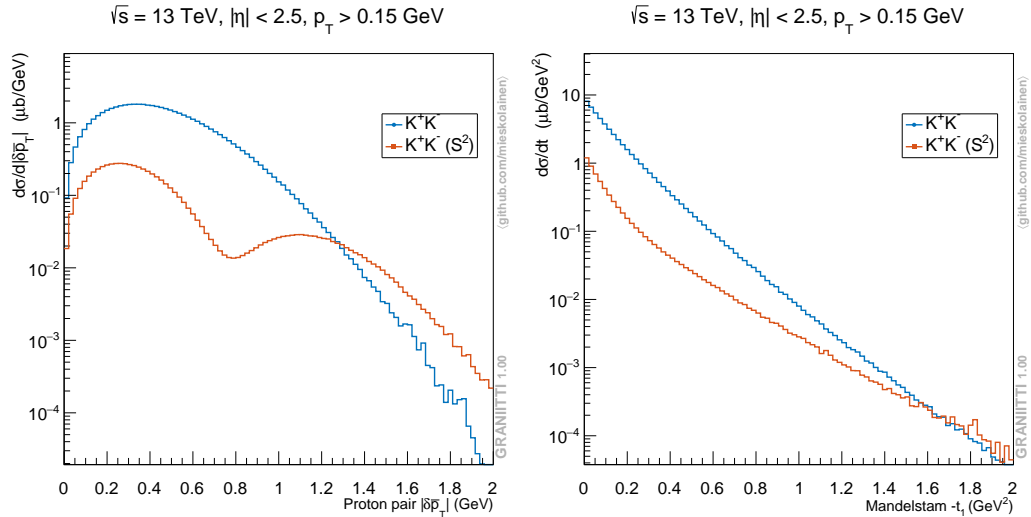


Figure 5.4: The screening loop effect on the K^+K^- continuum forward proton ‘glueball filter’ observable on the left and the Mandelstam $t_{1(2)}$ distribution on the right.

for the $2 \rightarrow N$ process with the corresponding measure $d\Pi_N$. For details, see Section 5.3. The largest absorption is obtained in Durham QCD model based production of a gluon pair gg . We see that lepton pair measurements have a systematically slightly lower cross section than the simulations. The simulations were done with k_t -EPA amplitudes, but this discrepancy is not explained by using the full tree-level $2 \rightarrow 4$ QED amplitude either, which is available in GRANIETTI. First one must remember that all photon-photon measurements in the table are non-exclusive and the exclusive cross section has been obtained by the collaborations via subtraction procedures. Regarding the semi-exclusive pion pair production, we see that our simulations are reasonably close to the experimental values obtained by CMS.

Forward dissociation

Central production together with single or double forward dissociation is implemented with exact skeleton kinematics, for details see Section 5.3. For dynamics the basic idea is that we can use the well known triple-pomeron limit $m_p^2/M^2 \ll 1$ and $M^2/s \ll 1$ of the forward system dissociation, which can be derived under the Mueller’s generalized 3-body optical theorem [82]. This neglects low mass baryonic resonance structure and all internal structure fluctuations of the proton. In princi-

ple, one could think in terms of stochastic processes of partons. A coherent proton is controlled by a diffusive Gaussian process where only small changes at a given time interval are present. This results in a limited exponential type momentum transfer distribution. The dissociation represents an incoherent jump process, perhaps of Cauchy type, where a radical change may take over any finite time interval. Unifying these together under a relativistic QCD framework is a challenge.

Single and double dissociation

The triple leg unitarity diagrams give at the cross section level the well known results [82]

$$|\mathcal{A}_{SD}|^2 \sim [F(t)g_{ppP}]^2 g_{ppP} g_{PPP}(t) \left(\frac{s}{M_1^2}\right)^{\alpha_P(t_i)+\alpha_P(t_j)} \left(\frac{M_1^2}{s_0}\right)^{\alpha_P(0)} \quad (5.16)$$

$$|\mathcal{A}_{DD}|^2 \sim [g_{ppP} g_{PPP}(t)]^2 \left(\frac{s_0 s}{M_1^2 M_2^2}\right)^{\alpha_P(t_i)+\alpha_P(t_j)} \left(\frac{M_1^2}{s_0}\right)^{\alpha_P(0)} \left(\frac{M_2^2}{s_0}\right)^{\alpha_P(0)}, \quad (5.17)$$

where we replace $\alpha_P(t_i) + \alpha_P(t_j) = 2\alpha_P(t)$ and take the normalization scale $s_0 = 1$ GeV². In general, one could easily include here a sum over all $2^3 = 8$ different pomeron and reggeon combinations in the triple vertex with corresponding g_{ijk} couplings and trajectories. The i, j and k indicate the Pomeron positions in the diagrams, where k is the one with the inelastic cut. For simplicity, we however use only one triple pomeron coupling term with no t -dependence. The essential feature is that with double dissociation, there is no proton form factor dependent term and the sensitivity to the crucial triple pomeron coupling is squared. In the limit $t \rightarrow 0$, we obtain the most essential triple pomeron scaling law

$$\frac{d\sigma}{dM_1^2} \sim \frac{1}{(M_1^2)^{\alpha(0)}}. \quad (5.18)$$

Different reggeon-pomeron triple terms give different scaling laws driven by different trajectory combinations, respectively. The power law mass scaling gives clues in the direction of critical phenomena. The proton structure function parametrizations for the forward legs usually incorporate these Regge type scaling laws. The triple pomeron coupling has a fundamental role in studies of the Reggeon field theory in more fundamental terms, weak versus strong coupling and higher orders, parton branching picture and QCD. With eikonal screening we obtain approximately a value of $g_{PPP}/g_{ppP} \simeq 0.15$ with $g_{ppP} \simeq 8$ GeV⁻² when no other triple leg terms are included. The value is small enough perhaps to consider that higher orders from

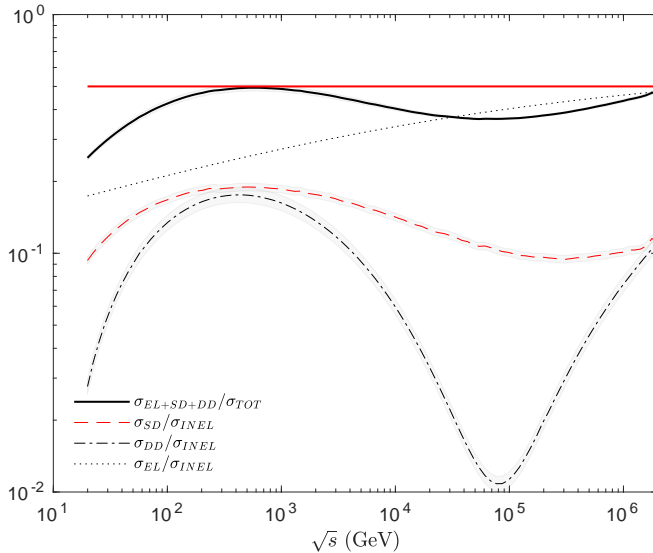


Figure 5.5: Inclusive cross section ratios and the Pumplin bound (solid red).

the full interacting Reggeon theory, which is an incomplete theory currently, can be neglected. It is well known that Eqs. 5.16 and 5.17 are not unitary, but scale with $s^{2\alpha_P(0)}$ and this is not compensated by the phase space which is for $2 \rightarrow 2$ given by $d\sigma/dt \sim 1/(16\pi s^2)$, combined together giving $\sim s^{2(\alpha_P(0)-1)}$. This is faster than the total cross section which grows like $s^{\alpha_P(0)-1}$ with a single pomeron exchange at large s .

To implement the eikonal screening for $2 \rightarrow 2$ single (SD) and double dissociation (DD) which restores unitarity, we use the same loop machinery as with central production, but here we need to recover the approximate equivalent bare amplitude by taking the square root and insert the pomeron signature factor phase – a crude first order inversion procedure which needs to be verified a posteriori. The integrated cross sections as a function of CMS energy are shown in Figure 5.6 with phase space integration cuts as indicated, to be able to compare with experimental data. In general, it is unknown both experimentally and theoretically what are the maximum forward mass or ‘coherence’ limits other than given by 4-momentum conservation, also the fragmentation gives irreducible smearing of the boundaries which we do not consider here.

An interesting feature is the energy evolution of the double dissociation cross

CHAPTER 5. GRANIITI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

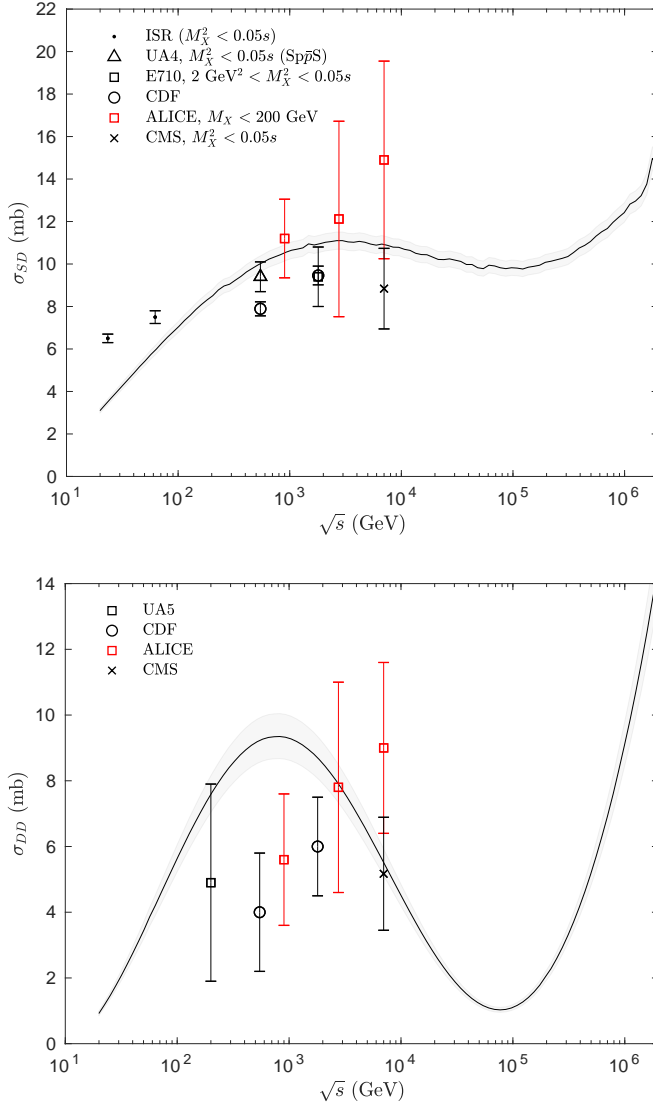


Figure 5.6: Eikonal screened single (forward + backward) and double diffraction integrated cross sections within phase space $|t| < 6 \text{ GeV}^2$, $M_X^2 < 0.05s$ and $\Delta Y_{DD} > 3$, respectively. Data from [102–107].

section, which develops a strong ‘bouncy’ oscillation in our screening computation. Because of large uncertainties, higher precision data and collider energies up to $\sqrt{s} = 100$ TeV are needed to verify the behavior experimentally. The cross section ratios are shown in Figure 5.5, where the Pumplin unitarity bound [108] $\sigma_{EL+SD+DD}/\sigma_{TOT} < 1/2$ is not violated. The $\sigma_{EL}/\sigma_{INEL}$ ratio is seen to have near power law scaling towards value 1/2 on the energy interval shown. What is known as the black disc (fully absorptive) limit corresponds to $\sigma_{EL} = \sigma_{INEL}$. Naturally, near the Planck scale or beyond we may expect some unexpected behavior.

To point out for clarity, we are not using the triple pomeron amplitudes in the central production dissociation, but pick up only the leading invariant mass dependence. We shall describe this next.

Central production with dissociation

The dissociative legs in the pomeron exchange central production are implemented via replacing the elastic vertex at amplitude level with an inelastic ansatz

$$g_{ppP}F(t) \mapsto [g_{ppP}g_{PPP}]^{\frac{1}{2}}\mathcal{F}(t, M^2), \quad (5.19)$$

where we take a coupling ansatz $[g_{ppP}g_{PPP}]^{\frac{1}{2}}$ motivated by the triple pomeron expressions. For the inelastic function \mathcal{F} we use a minimal parametrization similar to the Donnachie-Landshoff proton structure function parametrization [109]

$$\mathcal{F}(t, M^2) = \left[\frac{s_0}{M^2} \frac{|t|}{|t| + a} \right]^{\frac{1}{2}\alpha_P(0)}, \quad (5.20)$$

where $a = 0.56 \text{ GeV}^2$ and $s_0 = 1 \text{ GeV}^2$, obeying the triple pomeron scaling. Once more measurements are available, this vertex and its normalization may be re-defined. In photon exchange processes, the inelastic vertex is described in Section 5.2.

We illustrate the fully elastic and dissociation cross sections obtained for the pion pair continuum production in Figure 5.7. The results depend on the process, so we may take the continuum amplitude as the most simple example. Before screening, we obtain the average ratios over a large range of energies

$$\frac{\sigma(\pi^+\pi_{SD}^-)}{\sigma(\pi^+\pi_{EL}^-)} \simeq 0.44, \quad \frac{\sigma(\pi^+\pi_{DD}^-)}{\sigma(\pi^+\pi_{SD}^-)} \simeq 0.39. \quad (5.21)$$

Interesting is to compare with the ratio 1/2 of HERA results on dissociative photo-production of vector mesons V with $\sigma_{\gamma p \rightarrow VX}(W_{\gamma p}) \simeq \frac{1}{2}\sigma_{\gamma p \rightarrow Vp}(W_{\gamma p})$. The HERA

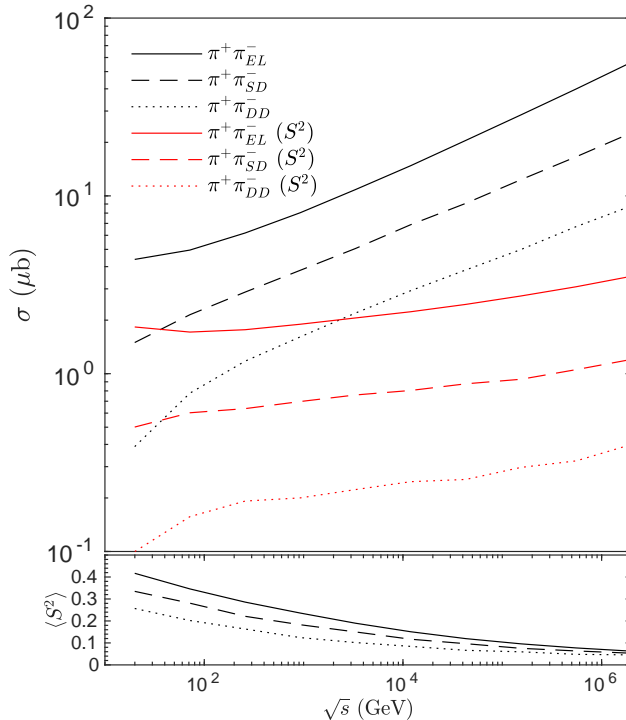


Figure 5.7: Continuum $\pi^+\pi^-$ pair production with elastic forward protons (EL), single dissociative (SD) and double dissociative (DD), without and with (S^2) the screening loop. In the lower figure the integrated screening ratio $\langle S^2 \rangle$ is shown. Fiducial cuts: $|\pi^\pm| < 0.9$ and $p_t(\pi^\pm) > 0.15$ GeV with forward system $M^2 < 0.05s$ and $\Delta Y_{DD} > 3$.

case is free from eikonal screening, thus it represents an interesting reference. After turning on the screening loop, we obtain the ratios

$$\frac{\sigma_S(\pi^+\pi_{SD}^-)}{\sigma_S(\pi^+\pi_{EL}^-)} \simeq 0.34, \quad \frac{\sigma_S(\pi^+\pi_{DD}^-)}{\sigma_S(\pi^+\pi_{SD}^-)} \simeq 0.29. \quad (5.22)$$

In Table 5.1 we have simulation results of semi-exclusive production for the LHC experiments. The simulation forward mass $M_{1,2}$ cuts given in the table of the dissociative production are approximations to the veto capabilities of the experiments. For a given experiment, there is a smooth forward mass (veto) efficiency function with asymptotics: $\epsilon(M) \rightarrow 0$ when $M \rightarrow 0$ and $\epsilon(M) \rightarrow \infty$ when $M \rightarrow \infty$. However, due to lack of instrumentation, the forward mass is not a direct experimental

observable at the LHC. Thus, this efficiency curve is obtained only through simulations and its verification is highly non-trivial, also because of significant detector material re-scattering effects. More precise effective cut estimates would require forward fragmentation by PYTHIA together with reliable effective fiducial pseudorapidity and transverse momentum cuts of the forward domain detectors obtained via GEANT simulation. Anyway, we observe that there should be a significant dissociative contribution in the data after the forward vetoes. This contribution can be experimentally disentangled on a statistical basis by inspecting the transverse momentum spectrum.

Pomeron-Pomeron interactions

The double pomeron fusion production of two, four and six central final state continuum is provided, generalizing the $pp \rightarrow pM\bar{M}p$ meson pair amplitude described in [89, 110]. The simplest case is the two body continuum with the sub- t and sub- u channel amplitudes, here for a pion pair production

$$\mathcal{A}_{pp \rightarrow p\pi^+\pi^-p}^{\hat{t}} = F(t_1)g_{ppP}\Delta_P(s_{13}, t_1)g_{\pi\pi P}F_M(\hat{t})\Delta_M(\hat{t}) \times F_M(\hat{t})g_{\pi\pi P}\Delta_P(s_{24}, t_2)g_{ppP}F(t_2), \quad (5.23)$$

where the pomeron-pion coupling $g_{\pi\pi P}$ is obtainable from the total cross section fits. The pomeron propagator part with the signature factor is

$$\Delta_P(s, t) = \eta(\alpha_P(t), +) \left(\frac{s}{s_0} \right)^{\alpha_P(t)}, \quad (5.24)$$

with the normalization scale typically set $s_0 = 1 \text{ GeV}^2$. The non-reggeized off-shell meson propagator is

$$\Delta_M(\hat{t}) = \frac{1}{\hat{t} - M_0^2}, \quad (5.25)$$

where M_0 is the pion mass. Here we use a typical linear pomeron trajectory parametrization and evaluate the signature factor with $\eta(t) \simeq \eta(0)e^{-i\frac{\pi}{2}\alpha't}$ limit to avoid signature poles encountered with the linear trajectory at high $|t|$. The sub- u channel amplitude is obtained by crossing $3 \leftrightarrow 4$ and thus $\hat{t} \leftrightarrow \hat{u}$.

The kinematic invariants are $\hat{t} = (q_1 - p_3)^2$, $\hat{u} = (q_1 - p_4)^2$ and $s_{ij} = (p_i + p_j)^2$, where $i = 1, 2$ denote the forward systems. The 4-momentum transfer squared are $t_1 = (p_A - p_1)^2$, $t_2 = (p_B - p_2)^2$ and the total CMS energy squared is $s = (p_A + p_B) = (p_1 + p_2 + p_3 + p_4)^2$. For the four and six body central states we need a set of

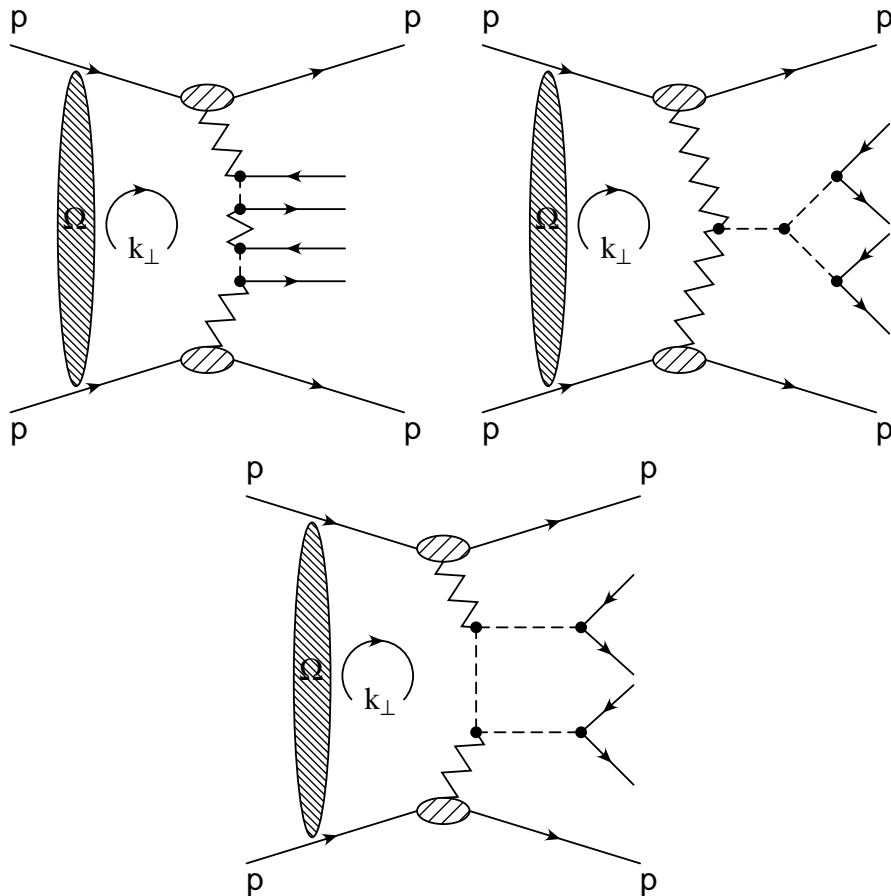


Figure 5.8: A pomeron-pomeron 4-body continuum amplitude on the left, resonance production with intermediate states on the right and 2-body continuum with intermediate states at the bottom, all with eikonal screening.

generalized invariants. For the straightforward details, we refer to the code which calculates these and the amplitudes algorithmically. Enumerating all permutations of charged particle final state legs gives 2, 16 and 288 sub-amplitudes growing like Cayley-Menger determinant absolute coefficients. Some of the amplitude topologies are illustrated in Figure 5.8.

For the off-shell meson sub-form factor $F_M(\hat{t})$, not to be mixed with the proton form factor, there are several options already familiar from [89], such as an exponen-

tial, Orear-like and a power law

$$F_M^{\text{exp}}(\hat{t}; \beta) = \exp(-\beta|t'|) \quad (5.26)$$

$$F_M^{\text{Orear}}(\hat{t}; \kappa_1, \kappa_2) = \exp\left(-\kappa_2\sqrt{|t'| + \kappa_1^2} + \kappa_1\kappa_2\right) \quad (5.27)$$

$$F_M^{\text{pow}}(\hat{t}; \kappa) = \frac{1}{1 - t'/\kappa}, \quad \text{with } t' \equiv \hat{t} - M_0^2. \quad (5.28)$$

By default, we use the power law which seems to work quite well. These form factors *do* contribute significant differences, effectively encapsulating unknown non-perturbative effects in QCD, which in the perturbative limit typically has power law behavior. We emphasize that the continuum amplitudes here have no explicit active helicity degrees of freedom, but the produced angular distributions are still far from isotropic $J = 0$. This is due to the amplitude structure. For a scenario where also the continuum amplitudes incorporate helicities, see the Tensor pomeron model amplitudes in Section 5.2. For the spin analysis purposes, we provide a generation mode of the crucial pure $J = 0$ continuum with exponential or limited $t_{1,2}$ -distributions for any number of final states produced by the isotropic decay.

The production of resonances interfering with the continuum is described with simple relativistic Breit-Wigner poles

$$\Delta_{BW}^R(m^2) = \frac{1}{m_R^2 - m^2 - im_R\Gamma}, \quad (5.29)$$

with the fixed or running width Γ and the pole mass m_R . The amplitude in the pomeron-pomeron fusion for the production of a resonance R decaying to particles 3 and 4 is

$$\begin{aligned} \mathcal{A}_{pp \rightarrow pR(\rightarrow 34)p}^{\hat{s}} &= F(t_1)\Delta_P(s_1, t_1)g_{ppP}\Delta_{BW}^R(m^2) \times \\ &\quad \mathcal{V}_{PP}^{R \rightarrow 34}(q_1, q_2, p_3, p_4)\Delta_P(s_2, t_2)g_{ppP}F(t_2), \end{aligned}$$

The kinematic variables are standard Lorentz scalars, discussed in Appendix B.2. The pomeron-resonance-pomeron vertex has an unknown effective functional structure, which we write as

$$\begin{aligned} &\mathcal{V}_{PP}^{R \rightarrow 34}(q_1, q_2, p_3, p_4) \\ &= \left[\frac{s_0}{(q_1 + q_2)^2} \right]^\omega \mathcal{A}_{PP \rightarrow R}(q_1, q_2)\mathcal{A}_{R \rightarrow 34}(p_3, p_4)F_R(q_1, q_2). \end{aligned} \quad (5.30)$$

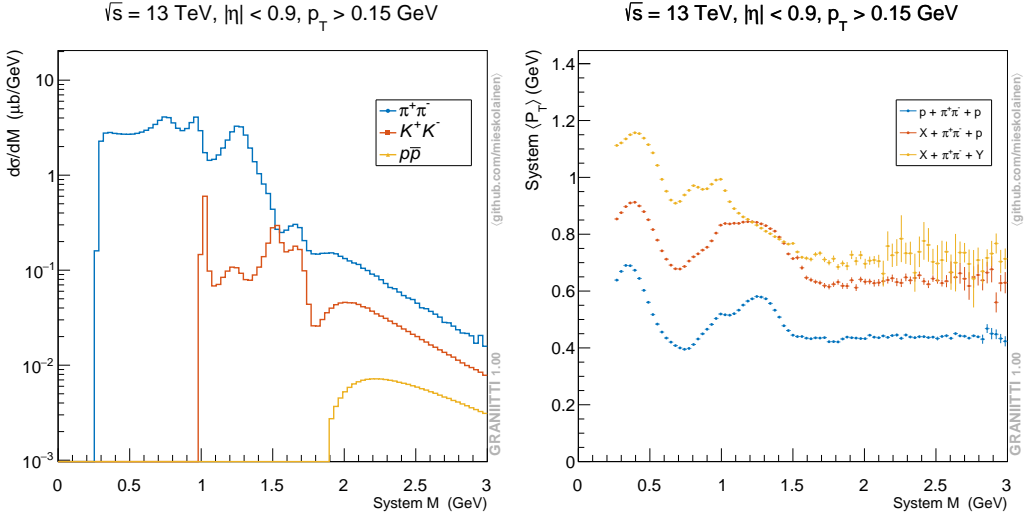


Figure 5.9: The invariant mass spectrum of $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pair on the left and the mean transverse momentum of the $\pi^+\pi^-$ spectrum with elastic, single and double dissociative production on the right. A constant $\langle S^2 \rangle \equiv 0.15$ applied.

The terms $\mathcal{A}_{PP \rightarrow R}$ and $\mathcal{A}_{R \rightarrow 34}$ encapsulate the production and the decay couplings and helicity dependence, respectively. The resonance production couplings $g_{PP \rightarrow R}$ embedded in $\mathcal{A}_{PP \rightarrow R}$ are free parameters (one per resonance) which we allow to be complex for maximal flexibility, to set up the relative complex phase between the continuum and resonance amplitudes. The rest of the production parameters for non-scalar resonances belong in our description to the resonance spin polarization density matrix. The vertex dependence in terms of the invariant mass squared $(q_1 + q_2)^2$ is parametric and the value $\omega = 1/2$ seems suitable. Technically, this is related to the resonance form factor

$$F_R(q_1, q_2) = \exp\left(-[(q_1 + q_2)^2 - m_R^2]/\Lambda_0^4\right), \quad \Lambda_0 \approx 1.0 \text{ GeV} \quad (5.31)$$

which is responsible for the finite size of the resonance. These both together modify the spectral shape beyond the effect of rising low-mass phase space, the Breit-Wigner propagator, Pomeron propagators and interference with the continuum amplitude. Again, alternative form factors are clearly possible.

Now we see very clearly that it is non-trivial to say which are the ‘true’ mass and width parameters of the composite resonances here because they depend effectively

on the production process, a well known feature. Figure 5.9 shows the invariant mass spectrum and average transverse momentum for $\pi^+\pi^-$, K^+K^- and $p\bar{p}$ pairs. We emphasize that the system transverse momentum allows one to gain control of the dissociative contribution, assuming that the elastic central production t -distribution is properly understood and controlled by forward proton measurements. The elastic case is driven by the proton form factors which are here fixed by the eikonal pomeron fits and the rest of the t -dependence comes from the pomeron slope α' dependent parts of Eq. 5.24.

Minimal spin pomeron

The polarized decay part $\mathcal{A}_{R\rightarrow 34}$ for non-scalar resonances is driven by the spin polarization density matrix of the resonance and Jacob-Wick helicity $1 \rightarrow 2$ amplitudes, details of this formalism are given in Section 5.2. For the production part $\mathcal{A}_{PP\rightarrow R}$, we first assume a *minimal effective spin* for the fusing pomerons compatible with the basic conservation laws. We take the pomeron spin to be $J = 0$ for the production of scalar or tensor meson resonances, also in photoproduction, and take $J = 1$ for the production of pseudoscalars. We have also other implementations in the code such as the sliding helicity trajectories, as we discuss in Appendix B.1. Our minimal choice introduces in the most resonance spin-parity cases no additional free parameters.

After fixing the spin of pomerons, we use the helicity amplitudes in $2 \rightarrow 1$ direction for the two pomeron fusion. This generates the forward proton azimuthal distribution $\Delta\varphi$ modulation which is dependent on the spin polarization matrix elements. The feasibility of this computational trick depends on the chosen Lorentz rest frame because these explicitly *non-covariant* amplitudes are parametrized by a frame-dependent spin polarization density matrix. Some rotated frames where pomerons have fixed $(\cos\theta, \varphi)$ event-by-event, such as the Gottfried-Jackson with pomerons back-to-back on the z -axis, are not directly suitable, neither is the helicity frame resulting (seemingly) in a too strong system rapidity dependence. We found out that a suitable generation frame is the Collins-Soper frame, which gives results in a good agreement with the preliminary ATLAS+ALFA data with $\Delta\varphi < \pi/2$ and $\Delta\varphi > \pi/2$ cuts [94]. This simulation is shown in Figure 5.10, also directly relevant to the upcoming CMS+TOTEM measurements. We see also that the t -acceptance cut results in a strong suppression of the photoproduced ρ^0 , as expected. When the pomeron spin is not zero, one should couple the proton legs with pomerons due to the pomeron polarization information. Our machinery is capable of this, but we encounter it only in the case of pseudoscalar resonances because of our minimal description. The pseudoscalar case is calculated in a faster way by using the results

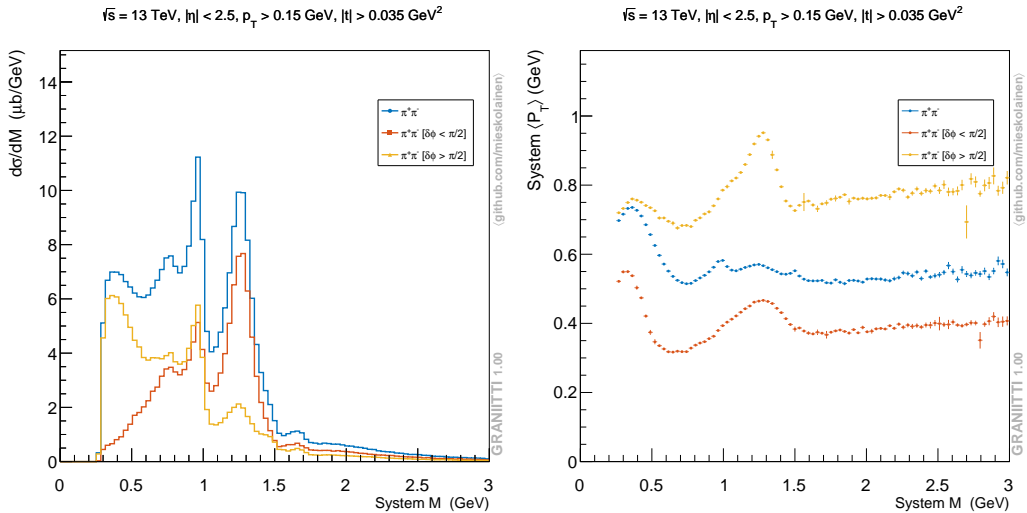


Figure 5.10: The exclusive $\pi^+\pi^-$ pair production without and with forward proton transverse plane angle cuts on the left. The mean transverse momentum with the same transverse plane cuts on the right. A constant $\langle S^2 \rangle \equiv 0.15$ applied.

discussed in [111]. People interested in *covariant* $2 \rightarrow N$ spin-dependent amplitudes may generate events using the tensor pomeron model described in Section 5.2.

The parameters of the spin polarization density can be changed to arbitrary values, but currently we take a diagonal ansatz of $|J_z| = \pm 1$ in photoproduction and $|J_z| = \pm 2$ for the tensor resonances such as $f_2(1270)$. This discussion is intimately related to the glueball filter introduced in [64]. Finally, the total amplitude is a coherent sum over the continuum and resonances

$$\mathcal{A}^{\text{tot}} \equiv \mathcal{A}^{\hat{t}} + \mathcal{A}^{\hat{u}} + \sum_R \mathcal{A}_R^{\hat{s}}. \quad (5.32)$$

The highly modular structure of our code and the input engine allows ‘experimenting’ with many of the phenomenological aspects of the scattering amplitudes, or plugging in completely new ones. Interesting would be to see if models such as ‘dual amplitudes’ of early strings would provide useful input here for the resonance coupling systematics [112]. A full theory would take into account also the analyticity and unitarity of the S -matrix and derive all the resonance couplings, something which is currently much beyond even the most radical amplitude technology.

Photoproduction

We have a (k_t -EPA)-pomeron based photoproduction of vector mesons, with a purpose of working as a suitable semi-reference process between the measurements done with and without forward proton tagging. In future we may implement dipole picture based models, such as variants of Golec-Biernat and Wüsthoff [113]. Due to the typical experimental t -acceptance at the LHC, the photoproduced vector mesons should disappear from the spectrum, but be well visible without the proton tagging. This seems to be the case in data as seen in [101] and [94]. The amplitude for the photoproduction of a vector meson V with a spin polarization dependent decay is

$$\mathcal{A}_{pp \rightarrow pV(\rightarrow 34)p} = \mathcal{A}_{\gamma P \rightarrow V}^{V \rightarrow 34} + \mathcal{A}_{P\gamma \rightarrow V}^{V \rightarrow 34}, \quad (5.33)$$

where

$$\mathcal{A}_{\gamma P \rightarrow V}^{V \rightarrow 34} = \left[\frac{1}{\xi_1} f_\gamma^{EL}(\xi_1, \mathbf{q}_{t,1}) \right]^{\frac{1}{2}} \Delta_{BW}^V(m^2) \mathcal{V}_{yP}^{V \rightarrow 34}(q_1, q_2, p_3, p_4) \mathcal{A}_P^V(s_2, t_2). \quad (5.34)$$

The amplitude $\mathcal{A}_{P\gamma}$ is obtained by permuting the variables with $1 \leftrightarrow 2$ and the central vertex \mathcal{V} is

$$\mathcal{V}_{yP}^{V \rightarrow 34}(q_1, q_2, p_3, p_4) = \mathcal{A}_{\gamma P \rightarrow V}(q_1, q_2) \mathcal{A}_{V \rightarrow 34}(p_3, p_4) F_R(q_1, q_2). \quad (5.35)$$

For more details about the photon flux f_{EL} and kinematic variables see Section 5.2. The relative sign of Eq. 5.33 would be negative in the case of proton-antiproton collisions, due to the anti-symmetric initial state. The pomeron side factor combining the s_2 and t_2 -dependence is

$$\mathcal{A}_P^V(s_2, t_2) = \Delta_P(s_2, t_2) \exp(b_0 t_2/2). \quad (5.36)$$

See also that $s_2 \equiv W_{\gamma p}^2$ is the typical notation in the deep inelastic scattering context. The normalization scale $s_0 = W_0^2 = 90^2 \text{ GeV}^2$ in the pomeron propagator of Eq. 5.24 is the most typical scale to fix the parameters. Note that the energy dependence of the subprocess is experimentally at larger energies approximately [114]

$$\sigma(\gamma p \rightarrow \rho^0 p) \propto W_{\gamma p}^{0.22} \quad (5.37)$$

$$\sigma(\gamma p \rightarrow J/\psi p) \propto W_{\gamma p}^{0.65}, \quad (5.38)$$

which motivates in the literature the large intercept ~ 0.3 hard pomeron (BFKL ladder) type production interpretation of J/ψ and Υ , where as for low mass vectors the soft pomeron seems suitable.

The gamma-pomeron-vector meson coupling and the pomeron side exponential b_0 -slope have been fixed with HERA data. The slope parameters b_0 are typically $\sim 11 \text{ GeV}^{-2}$ for ρ^0 production down to $\sim 4 - 5 \text{ GeV}^{-2}$ for J/ψ , reflecting the intrinsic transverse size of the vector meson $q\bar{q}$ dipole. These experimental fit values combine the effect of both the vector meson and the proton form factor. The photon side has a much steeper t -dependence than the pomeron side, which is included in our description. The main expected difference between photoproduction and speculative odderon-pomeron fusion is thus in the transverse momentum dependence, but possibly also in the polarization of the produced vector mesons, which is easily changed in the simulation. Also, we have simple odderon-pomeron amplitudes within a simplified description. We take the odderon simply as odd signature pomeron with unknown couplings to be fixed.

Decay couplings

The effective decay coupling for two-body decays $M_R \rightarrow m_1 + m_2$ is calculated according to the branching ratios $\text{BR}_i \equiv \Gamma_i/\Gamma$ imported from the PDG and the standard partial decay width formula which factorizes, in the $1 \rightarrow 2$ case, between the phase space and the decay matrix element squared. The partial decay width is the well known

$$\Gamma_i = \frac{1}{S} \frac{\sqrt{\lambda(M_R^2, m_1^2, m_2^2)}}{16\pi M_R^3} |\mathcal{M}_D|^2, \quad (5.39)$$

where the $\sqrt{\lambda}$ is the standard Källén triangle function and S is the statistical symmetry factor. That is, we simply invert the relation to find out the effective decay coupling $|\mathcal{M}_D|^2 \sim |g_D|^2$ given the measured branching ratios. The full decay matrix element is non-perturbative and unknown here for the f -mesons and glueballs but may be estimated under certain frameworks such as holography. For non-scalar decays, one needs to take into account the spin related normalization factors in the machinery in Section 5.2. To obtain higher precision, we could calculate this using our generic phase space sampling functions which integrate also phase space volumes, more suitable for large width resonances and near threshold behavior. The machinery allows one to add arbitrary many new resonances and their decays. We provide arbitrary deep decay chains according to the phase space but also with (cascaded) spin correlations initiated by the resonance amplitude. The cascaded spin correlations require the intermediate decay ls -couplings as an input. A flexible decay chain machinery is highly important for many experimental analyses, which need to test different hypotheses and evaluate the significance of various ‘feed-down’ contribu-

tions, for example.

Photon-Photon interactions

Two photon interactions are generated according the k_t -EPA Equivalent Photon Approximation formalism for the photon fluxes from protons, with both elastic and inelastic fluxes at the cross section level [115, 116]

$$\begin{aligned}
 f_\gamma^{EL}(\xi, \mathbf{q}_t) &= \frac{\alpha}{\pi} \left[(1 - \xi) \Delta^2 \frac{4m_p^2 G_E^2(Q^2) + Q^2 G_M^2(Q^2)}{4m_p^2 + Q^2} + \frac{\xi^2}{4} \Delta G_M^2(Q^2) \right] \\
 f_\gamma^{IN}(\xi, \mathbf{q}_t, M_X^2) &= \frac{\alpha}{\pi} \left[(1 - \xi) \Delta^2 \frac{F_2(x_{Bj}, Q^2)}{Q^2 + M_X^2 - m_p^2} + \frac{\xi^2}{4x_{Bj}^2} \Delta \frac{2x_{Bj} F_1(x_{Bj}, Q^2)}{Q^2 + M_X^2 - m_p^2} \right] \\
 \text{with } \Delta &\equiv \frac{\mathbf{q}_t^2}{\mathbf{q}_t^2 + \xi(M_X^2 - m_p^2) + \xi^2 m_p^2}, \tag{5.40}
 \end{aligned}$$

where $\xi = 1 - p'_z/p_z$ is the longitudinal momentum loss of the proton, \mathbf{q}_t the photon transverse momentum, $Q^2 \equiv -t$ the 4-momentum transfer squared, $x_{Bj} = Q^2/(Q^2 + M_X^2 - m_p^2)$ is the Bjorken- x identically one for the elastic case and M_X^2 is the forward system invariant mass. These fluxes are then matched with the exact $2 \rightarrow N$ phase space construction taking into account the kinematic factors by a transform at the cross section level

$$f_\gamma \mapsto 16\pi^2 [\xi \mathbf{q}_t^2]^{-1} f_\gamma, \tag{5.41}$$

which we have verified against the full QED $pp \rightarrow pl^+l^-p$ tree level amplitude.

The coherent proton electromagnetic form factors in Sachs form [117] are G_E and G_M . By construction, the linear relation between Sachs and F_1 (Dirac) and F_2 (Pauli) form factors can be written as

$$\begin{pmatrix} G_E \\ G_M \end{pmatrix} = \begin{pmatrix} 1 & -\tau \\ 1 & 1 \end{pmatrix} \begin{pmatrix} F_1 \\ F_2 \end{pmatrix} \Leftrightarrow \begin{pmatrix} F_1 \\ F_2 \end{pmatrix} = \begin{pmatrix} \frac{1}{\tau+1} & \frac{\tau}{\tau+1} \\ -\frac{1}{\tau+1} & \frac{1}{\tau+1} \end{pmatrix} \begin{pmatrix} G_E \\ G_M \end{pmatrix}, \tag{5.42}$$

where $\tau = Q^2/(4m_p^2)$. A simple dipole parametrization is used

$$G_E(Q^2) = G_D(Q^2) = \frac{1}{(1 + Q^2/\lambda^2)^2} \tag{5.43}$$

$$G_M(Q^2) = \mu_p G_D(Q^2), \tag{5.44}$$

where the proton magnetic moment $\mu_p = 2.7928$ in units of the nuclear magneton $\mu_N = e\hbar/(2m_p)$ and $\lambda^2 = 0.71 \text{ GeV}^2$. The normalization here gives $F_1(0) = 1$

and $F_2(0) = 1.7928$. In addition, we have included more complex but still simple parametrization from [118].

In the inelastic case, the proton structure functions are $F_2(x, Q^2)$ and $F_1(x, Q^2)$, for which we use some classic parametrizations, to be plug-in replaced easily by more relevant up-to-date external libraries or implemented by the user. In the parton model the Callan-Gross relation holds

$$2xF_1(x) = F_2(x), \quad (5.45)$$

due to spin-1/2 quarks, but in QCD with gluons a longitudinal structure function component is present

$$F_L(x, Q^2) \equiv \left(1 + \frac{4x^2m_p^2}{Q^2}\right)F_2(x, Q^2) - 2xF_1(x, Q^2), \quad (5.46)$$

which may be extracted from the scaling violations. The longitudinal component is especially relevant at small values of Bjorken- x where gluon density rises steeply. The terminology between transverse and longitudinal stems from the deep inelastic scattering and the corresponding virtual photon polarization component cross sections. With high enough Q^2 , for the F_2 here one could use the QCD evolved parton density description

$$F_2(x, Q^2) = x \left(\sum_{i \in u_v, d_v} e_i^2 f_i(x) + \sum_{i \in u_s, d_s, s_s} e_i^2 [f_i(x) + \bar{f}_i(x)] \right), \quad (5.47)$$

with $e_i^2 = 4/9$ ($1/9$) for up (down) type quarks, where the parton densities are readily available through the LHAPDF 6 library interface [119]. The longitudinal structure function F_L can be related to gluon densities within pQCD, an interesting topic also algorithmically, for a recent work see [120].

To cross check the k_t -EPA implementation differentially, we have implemented the $pp \rightarrow p\ell\bar{\ell}p$ tree level $2 \rightarrow 4$ full QED process with the standard covariant current

$$J^\mu = ie\bar{v}(p', \lambda') [\gamma_\mu F_1(Q^2) + \frac{i\sigma_{\mu\nu}}{2m_p} (p - p')^\nu F_2(Q^2)] v(p, \lambda), \quad (5.48)$$

where the term with F_1 is the helicity λ conserving part and the term with F_2 is the helicity flip (non-conserving) part and $\sigma_{\mu\nu} = \frac{i}{2}(\gamma_\mu\gamma_\nu - \gamma_\nu\gamma_\mu)$. This implementation also provides proper distributions e.g. for very low invariant masses of the lepton pair system, where the k_t -EPA + on-shell $\gamma\gamma \rightarrow X$ matrix elements are on the edge of their validity.

Currently, we have included on-shell matrix elements for the lepton pair and W^+W^- production with helicity amplitudes imported from MadGraph5 C++ export [121]. Because our kinematics construction is exact $2 \rightarrow N$, the photon kinematics have always small but finite $q_i^2 < 0$, but MadGraph $\gamma\gamma \rightarrow X$ amplitudes assume $q_i^2 = 0$. We correct this by transforming the initial state photons to the closest point at light cone which is found by Lagrange multipliers and iterate the final state kinematics so that energy-momentum is conserved and amplitudes can be safely evaluated. This procedure might be optimized in the future.

In addition, we have also $j = 1/2$ monopodium pair production and monopodium bound state $J = 0$ resonance production with pure Dirac and velocity-dependent couplings, as simple scenarios of fundamental magnetic monopoles. As magnetic monopoles are strongly coupled, intrinsically non-perturbative and currently lacking rigorous field theory framework, the QED matrix element replacement represents a somewhat uncontrolled approximation, but still useful. The bound state modeling is based on a simple Schrödinger equation type solution.

Durham QCD model

For the Durham QCD (KMR) model [122], we include the numerical gluon loop with spin-parity projection and a generalized gluon pdf transformation which includes the Shuvaev transform and Sudakov radiation suppression. The main interest for us here is the transition region from the low mass Regge domain to this QCD domain. We go now through the formulation, for more details see [123, 124]. The formulation starts with the amplitude at parton level for $qq \rightarrow q + X + q$ with a gluon loop, which is most easily derived under the high energy eikonal forward quark vertex limit

$$-i\bar{u}(p', \lambda') ig_s T_{ij}^a \gamma^\mu u(p, \lambda) \rightarrow -2g_s p^\mu T_{ij}^a \delta_{\lambda, \lambda'}, \quad (5.49)$$

where $T_{ij}^a = \lambda_{ij}^a/2$ is the $SU(3)$ generator matrix element $\langle i|T^a|j \rangle$ and color indices run $a = 1, \dots, 8$ and $i, j = 1, 2, 3$. The amplitude will be dominantly imaginary in the forward limit and the imaginary part of the loop amplitude can be obtained most easily with Cutkosky cutting rules [125], which replace propagators by delta functions. The Durham model is illustrated in Figure 5.11, where the Cutkosky line goes vertically through the gluon loop.

We denote the quark momentum with color indices in the parenthesis

$$\text{Left side of the cut: } p_1(i) + p_2(j) \rightarrow p_1(m) + p_2(n) \quad (5.50)$$

$$\text{Right side of the cut: } p_1(m) + p_2(n) \rightarrow p'_1(i) + p'_2(j), \quad (5.51)$$

where the color is oriented along the quark lines so that the system X is color singlet and the amplitude will be compatible at the proton level. The screening gluon carries momentum Q with color c on the left side of the cut and the fusing gluons $q_1(a), q_2(b)$ on the right side. Now writing these down gives [123]

$$\text{Im } \mathcal{A} = \underbrace{\frac{1}{2}}_{\text{cut-rule}} \times \underbrace{2}_{\# \text{diagrams}} \times \int d\Pi_2 \frac{2g_s p_1^\rho \cdot 2g_s p_{2\rho}}{Q^2} \frac{2g_s p_1^\mu}{q_1^2} \frac{2g_s p_2^\nu}{q_2^2} \times \quad (5.52)$$

$$V_{\mu\nu}^{ab} T_{mi}^c T_{im}^a T_{nj}^c T_{jn}^b \delta((p_1 - Q)^2) \delta((p_2 + Q)^2). \quad (5.53)$$

At this point we take the integral over transverse momentum space in the eikonal limit

$$\int d\Pi_2 \equiv \frac{1}{2s} \int \frac{d^2 \mathbf{Q}_t}{(2\pi)^2}, \quad (5.54)$$

which can be overall justified by Sudakov decomposition, and we will change the momentum variables to transverse variables. The average and sum over colors at amplitude level gives

$$\frac{1}{N_C^2} T_{mi}^c T_{im}^a T_{nj}^c T_{jn}^b = \frac{1}{N_C^2} \text{Tr}[T^c T^a] \text{Tr}[T^c T^b] = \frac{1}{4N_C^2} \delta^{ca} \delta^{cb} = \frac{\delta^{ab}}{4N_C^2} \rightarrow \frac{2}{9}, \quad (5.55)$$

where the last step is obtained, when the sub-amplitude gives equal $N_C^2 - 1$ contributions such as the SM Higgs production. The normalization color factor is $T_R = 1/2$ from $\text{Tr}[T^a T^b] = T_{ij}^a T_{ij}^b = T_R \delta^{ab}$ associated with the gluon splitting into quark pair, $C_F = (N_C^2 - 1)/(2N_C) = 4/3$ is the color factor with gluon emission from quarks $\sum_a T_{ik}^a T_{kj}^a = \sum_a (T^a T^a)_{ij} \equiv C_F \delta_{ij}$ and the gluon splitting into gluon color factor $f_{acd} f_{bcd} = C_A \delta_{ab}$ with $C_A \equiv N_C = 3$. The generators are traceless $\text{Tr}[T^a] = \sum_{i=1}^{N_C} T_{ii}^a = 0$.

Kinematics and coupling factors give us

$$\frac{1}{2} \times 2 \times \frac{1}{2s} \frac{1}{(2\pi)^2} \times \frac{s}{2} \times (2\sqrt{4\pi\alpha_s})^4 = 16\alpha_s^2. \quad (5.56)$$

Combined together with color factors we get

$$\frac{\delta^{ab}}{4N_C^2} \times 16\alpha_s^2 \xrightarrow{V_{\mu\nu}^{ab} \sim \delta^{ab}} \frac{N^2 - 1}{N^2} 4\alpha_s^2, \quad (5.57)$$

where the last step is in the case of SM Higgs like amplitude.

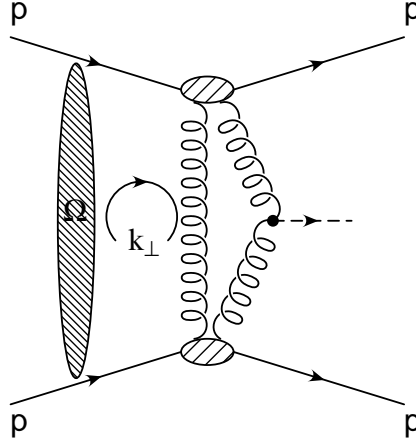


Figure 5.11: Durham QCD model with eikonal screening.

The vertex contraction is manipulated into fusing gluon form by setting $q_i = x_i p_i + q_{t,i}$ and $2q_1 \cdot q_2 \simeq M_X^2$

$$p_1^\mu p_2^\nu V_{\mu\nu}^{ab} \simeq \frac{q_{t,1}^\mu}{x_1} \frac{q_{t,2}^\nu}{x_2} V_{\mu\nu}^{ab} \xrightarrow{s x_1 x_2 \simeq M_X^2} \frac{s}{M_X^2} q_{t,1}^\mu q_{t,2}^\nu V_{\mu\nu}^{ab}, \quad (5.58)$$

where one could use gauge invariance $q_1^\mu V_{\mu\nu}^{ab} = q_2^\nu V_{\mu\nu}^{ab} = 0$. The contraction with the vertex is as with external polarization vectors $\epsilon \sim q$, which means a transverse polarization with $\epsilon_1 = -\epsilon_2$ from $Q_t = -q_{1,t} = q_{2,t}$ in the pure forward limit $p'_{t,1,2} \rightarrow 0$ of outgoing quarks giving the $J_z = 0$ selection rule.

Now we denote the color averaged sub-amplitude as [126]

$$\langle \mathcal{M} \rangle \equiv \frac{2}{M_X^2} \frac{1}{N_C^2 - 1} \sum_{a,b=1}^{N_C^2 - 1} \delta^{ab} q_{1,t}^\mu q_{2,t}^\nu V_{\mu\nu}^{ab}. \quad (5.59)$$

The total result matching the derivation above can be written as

$$\frac{\text{Im } \mathcal{A}}{s} \simeq C_F^2 \alpha_s^2 \int \frac{d^2 \mathbf{Q}_t}{\mathbf{Q}_t^2 \mathbf{q}_1^2 \mathbf{q}_2^2} \langle \mathcal{M} \rangle. \quad (5.60)$$

What one notices, is the infrared divergence in the loop Q_t^2 which is to be tamed by the Sudakov resummation.

Now, the description above was at the parton level. At the hadron level, we need the generalized gluon ‘Durham flux’ from the proton [127]

$$f_g(x, x', Q_t^2, \mu^2) = \frac{\partial}{\partial \ln Q_t^2} \left[\sqrt{T(Q_t^2, \mu^2)} G\left(\frac{x}{2}, \frac{x}{2}, Q_t^2\right) \right], \quad (5.61)$$

where the generalized (skewed) gluon pdf is obtained by a ‘Shuvaev transform’ of the standard integrated gluon pdf $g(x, Q^2)$ with an integral method [128]

$$G\left(\frac{x}{2}, \frac{x}{2}, Q_t^2\right) = \frac{4x}{\pi} \int_{\frac{x}{4}}^1 dy g\left(\frac{x}{4y}, Q_t^2\right) (1-y)^{1/2} y^{1/2}, \quad (5.62)$$

and the Sudakov suppression vetoing real radiation is

$$T(Q_t^2, \mu^2) = \exp \left(- \int_{Q_t^2}^{\mu^2} \frac{dk_t^2}{k_t^2} \frac{\alpha_s(k_t^2)}{2\pi} \int_0^{1-\Delta} dz \left[z P_{gg}(z) + \sum_{q \in \text{flavors}} P_{qg}(z) \right] \right), \quad (5.63)$$

where $P_{gg}(z)$ and $P_{qg}(z)$ are leading order DGLAP splitting kernels. In the integral the upper bound $\Delta = k_t/\mu$ is taken as described in detail in [124]. For the discussion of single and double logarithmic terms and the origin of the derivative in Eq. 5.61, see [127]. For completeness, the splitting kernels are for gluon $P_{g \leftarrow g}$ and quark $P_{q \leftarrow g}$ emissions

$$P_{gg}(z) = 2C_A \left[\frac{z}{(1-z)_+} + \frac{1-z}{z} + z(1-z) \right] + \delta(1-z) \frac{11C_A - 4n_f T_R}{6} \quad (5.64)$$

$$P_{qg}(z) = T_R [z^2 + (1-z)^2], \quad (5.65)$$

where the ‘plus prescription’ is

$$\int_0^1 dz \frac{f(z)}{(1-z)_+} = \int_0^1 dz \frac{f(z) - f(1)}{1-z}, \quad \text{with } (1-z)_+ = 1-z, \quad \text{for } z < 1, \quad (5.66)$$

so formally the soft gluon divergence at $z = 1$ cancellation relies on $f(1)$. The integration and differentiation are done numerically, the results being cached into look-up tables and interpolated. The Sudakov suppression is illustrated in Figure 5.12.

Finally at the proton level, the scattering amplitude for the total process is given by

$$\mathcal{T} = \frac{\text{Im } \mathcal{A}}{s} = \pi^2 \int \frac{d^2 \mathbf{Q}_t}{\mathbf{Q}_t^2 (\mathbf{Q}_t - \mathbf{p}_{t,1})^2 (\mathbf{Q}_t + \mathbf{p}_{t,2})^2} \times f_g(x_1, x'_1, Q_1^2, \mu^2) F(t_1) f_g(x_2, x'_2, Q_2^2, \mu^2) F(t_2) \langle \mathcal{M}_{gg \rightarrow X} \rangle \quad (5.67)$$

CHAPTER 5. GRANIETTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

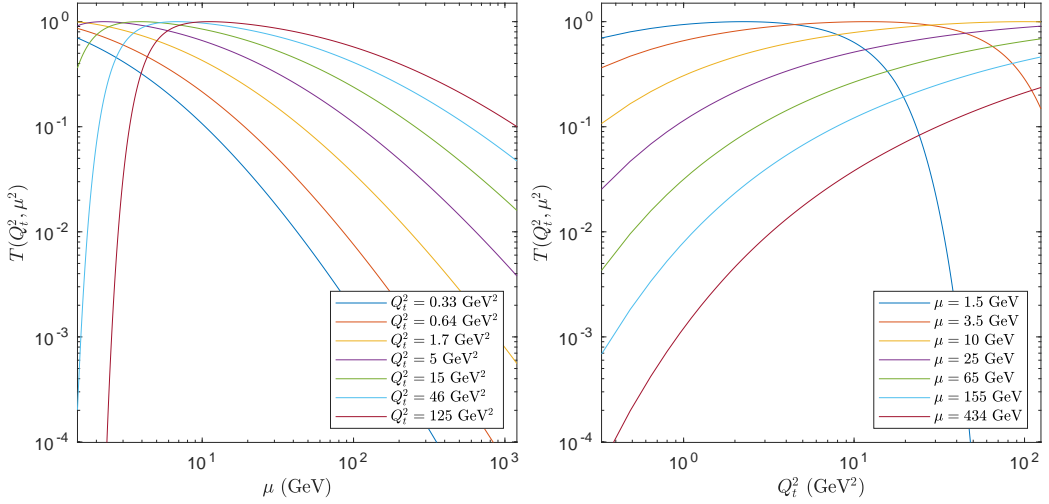


Figure 5.12: The Sudakov suppression factor $T(Q^2, \mu^2)$ evolution.

where the π^2 factor comes from the parton to proton replacement $\alpha_s C_F / \pi \rightarrow f_g$. Here we simply assume that the *coherent* proton form factors factorize with the gluon fluxes. The fusing gluon vectors in the loop are $\mathbf{q}_1 = \mathbf{Q}_t - \mathbf{p}_{t,1}$ and $\mathbf{q}_2 = \mathbf{Q}_t + \mathbf{p}_{t,2}$, with the outgoing proton transverse momentum $\mathbf{p}_{t,1,2}$. Here, one sees that the process cross section has a quadratic dependence on the gluon pdfs. This 2D-loop integral is calculated numerically event-by-event and the default scales are chosen as $\mu^2 = M_X^2$ and $Q_i^2 = \min(|\mathbf{Q}_t|^2, |\mathbf{q}_i|^2)$, but these can be varied easily with the program input.

For the sub-amplitude, it is useful to use the helicity basis and use the decomposition [126] of $q_1^i q_2^j \mathcal{M}_{ij}$ which gives

$$J_z^P = 0^+ : -\frac{1}{2}(\mathbf{q}_1 \cdot \mathbf{q}_2) [\mathcal{M}_{++} + \mathcal{M}_{--}] \quad (5.68)$$

$$J_z^P = 0^- : -\frac{i}{2}(\mathbf{q}_1 \times \mathbf{q}_2) [\mathcal{M}_{++} - \mathcal{M}_{--}] \quad (5.69)$$

$$J_z^P = +2^+ : +\frac{1}{2}([\mathbf{q}_1 \ominus \mathbf{q}_2] + i[\mathbf{q}_1 \oplus \mathbf{q}_2]) \mathcal{M}_{+-} \quad (5.70)$$

$$J_z^P = -2^+ : +\frac{1}{2}([\mathbf{q}_1 \ominus \mathbf{q}_2] - i[\mathbf{q}_1 \oplus \mathbf{q}_2]) \mathcal{M}_{+-} \quad (5.71)$$

with $\mathbf{q}_1 \oplus \mathbf{q}_2 \equiv q_{1,x}q_{2,y} + q_{1,y}q_{2,x}$ and $\mathbf{q}_1 \ominus \mathbf{q}_2 \equiv q_{1,x}q_{2,x} - q_{1,y}q_{2,y}$. For each outgoing helicity combination, let us point out that the sum over incoming helicity states is here *coherent*. The current first implementation includes $gg \rightarrow gg$ continuum,

$gg \rightarrow \chi_{c0}$ resonance and $gg \rightarrow m\bar{m}$ pseudoscalar meson pair amplitudes. The gluon pdfs and running couplings are provided by the LHAPDF6 library [119]. We encourage users to implement their own processes and pay attention to the normalization conventions. The phase space normalization factors are chosen such that the process is compatible with $2 \rightarrow N$ kinematics. Interfacing with automated matrix element generators should be the target for the future. For more processes readily available within Durham model we refer the reader to SUPERCHIC 3 MC [90].

Tensor pomeron model

The tensor pomeron model [129] implementation includes central exclusive processes of a two body continuum production of pseudoscalar pairs, vector meson pairs and baryon pairs. The resonance processes implemented here include the production of scalar resonances f_0 , pseudoscalar resonances η, η' , vector mesons $\rho^0(770)$ and $\varphi(1020)$ via photoproduction and f_2 tensor mesons interfering with the continuum at the amplitude level. The model takes the ansatz that the pomeron should carry a definite Lorentz structure, namely a rank-2 graviton like current, coupling thus in a symmetric way between matter and antimatter. In this picture, it is the vector odderon which provides the anti-symmetric coupling.

A tensor like pomeron has been also recently studied in a holographic AdS/QCD context in [130] constructing a duality between the triple-graviton vertex and the double pomeron fusion production of glueballs. In the classic Gribov Regge theory, the pomeron carries ‘sliding spin’ as an infinite sum and, depending on the interaction, may or may not coincide with definite Lorentz structures.

We consider the model implemented here as a practical one due to its explicit enough computational rules, however, see also the aesthetics behind other descriptions. An interesting problem is to see, how uniquely can the upcoming LHC or RHIC measurements with forward protons constrain these structures. Based on this and ansatz structures for couplings, one can write down a diverse set of interactions with mesons and baryons using the corresponding Feynman rules [129]. As an example, the tensor pomeron propagator ansatz is

$$i\Delta_{\mu\nu,\kappa\lambda}(s, t) = \frac{1}{4s}(g_{\mu\kappa}g_{\nu\lambda} + g_{\mu\lambda}g_{\nu\kappa} - \frac{1}{2}g_{\mu\nu}g_{\kappa\lambda})(-is\alpha'_P)^{\alpha_P(t)-1}, \quad (5.72)$$

obeying permutation symmetries and identities

$$\Delta_{\mu\nu,\kappa\lambda}(s, t) = \Delta_{\mu\nu,\lambda\kappa}(s, t) = \Delta_{\nu\mu,\kappa\lambda}(s, t) = \Delta_{\kappa\lambda,\mu\nu}(s, t) \quad (5.73)$$

$$g^{\kappa\lambda}\Delta_{\mu\nu,\kappa\lambda}(s, t) = 0, \quad g^{\mu\nu}\Delta_{\mu\nu,\kappa\lambda}(s, t) = 0. \quad (5.74)$$

The tensorial coupling with the proton or antiproton is

$$i\Gamma_{\mu\nu}(p', p) = -i3\beta F_1(t) \left[\frac{1}{2}[\gamma_\mu(p' + p)_\nu + \gamma_\nu(p' + p)_\mu] - \frac{1}{4}g_{\mu\nu}(\not{p}' + \not{p}) \right], \quad (5.75)$$

where $t = (p - p')^2$ and $\beta = 1.87 \text{ GeV}^{-1}$. We may illustrate the central vertex by the double pomeron production of a pseudoscalar resonance. It involves two amplitude structures

$$i\Gamma_{\mu\nu,\kappa\lambda}^{(1)}(q_1, q_2) = i(g_{\mu\kappa}\epsilon_{\nu\lambda\rho\sigma} + g_{\nu\kappa}\epsilon_{\mu\lambda\rho\sigma} + g_{\mu\lambda}\epsilon_{\nu\kappa\rho\sigma} + g_{\nu\lambda}\epsilon_{\mu\kappa\rho\sigma}) \times (q_1 - q_2)^\rho (q_1 + q_2)^\sigma \quad (5.76)$$

$$i\Gamma_{\mu\nu,\kappa\lambda}^{(2)}(q_1, q_2) = i\{\epsilon_{\nu\lambda\rho\sigma}[q_{1\kappa}q_{2\mu} - q_1 \cdot q_2 q_{\mu\kappa}] + \epsilon_{\mu\lambda\rho\sigma}[q_{1\kappa}q_{2\nu} - q_1 \cdot q_2 q_{\nu\kappa}] + \epsilon_{\nu\kappa\rho\sigma}[q_{1\lambda}q_{2\mu} - q_1 \cdot q_2 q_{\mu\lambda}] + \epsilon_{\mu\kappa\rho\sigma}[q_{1\lambda}q_{2\nu} - q_1 \cdot q_2 q_{\nu\lambda}]\} \times (q_1 - q_2)^\rho (q_1 + q_2)^\sigma, \quad (5.77)$$

which correspond to ls -couplings $(1, 1)$ and $(3, 3)$, respectively [131]. In addition, one adds couplings and form factors in the vertex. For the rest of somewhat lengthy building blocks, we refer to the original papers [129, 131], or directly to our code written in high level C++, which follows closely the algebraic notation. The effective decay couplings of resonances, when possible, are computed according to the model spin structure and PDG branching ratios, which in the scalar case matches directly Eq. 5.39.

In the code, we have implemented the Dirac and Lorentz algebra including gamma matrices, Dirac spinors, massive and massless spin-1 polarization vectors and spin-2 tensors, photon, pomeron and fermion propagators and resonance coupling tensors and evaluate the scattering amplitudes helicity combination by combination with explicit component representations. Continuum and resonance amplitudes are summed together at the amplitude level. Computationally speaking, the most complex sub-amplitudes are the rank-6 Lorentz structures for $J = 2$ resonances. Some of the resulting observables are shown in Figure 5.13. All basic identities are implemented as code unit tests, such as completeness relations and normalizations. Performance wise, this slightly ‘brute force’ approach could be improved in the future by a more streamlined spinor-helicity formalism. The Lorentz index contractions are accelerated with the FTensor tensor algebra C++ expression template library [132] originally designed for numerical general relativity.

We may discuss a bit the Lagrangian structures. As an example: the negative parity of η -meson production is ‘compensated’, to conserve parity, by the anti-symmetric

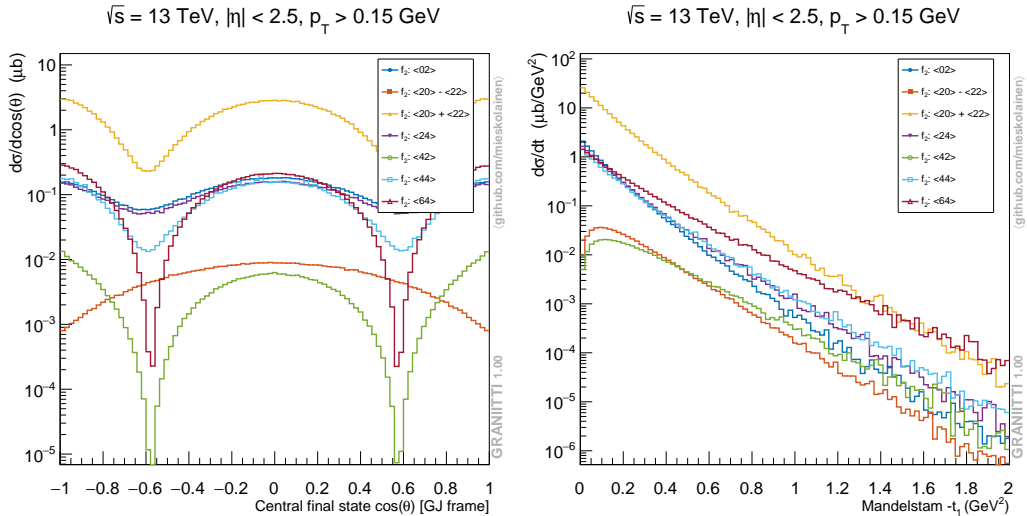


Figure 5.13: PPf_2 -vertex rank-6 tensor structures and the resulting angular distributions of final state pions in the system Gottfried-Jackson rest frame and Mandelstam $t_{1(2)}$ distributions. A constant $\langle S^2 \rangle \equiv 0.15$ applied.

epsilon tensor structure which produces characteristic $\sin^2 \Delta\varphi$ cross section dependence motivated experimentally by WA102 data – a well known example. We can take an analogous case from the two photon decays of pseudoscalar and scalar mesons or Higgs. At the effective Lagrangian level, for a pseudoscalar this is a term like $F^{\mu\nu} F^{\kappa\rho} \epsilon_{\mu\nu\kappa\rho} \eta$ whereas for a scalar it would be $F_{\mu\nu} F^{\mu\nu} f_0$, where η, f_0 are the scalar fields and $F^{\mu\nu}$ is the photon vector field. The corresponding amplitudes behave like $e^{\mu\nu\kappa\rho} \epsilon_\mu^1 q_\nu^1 \epsilon_\kappa^2 q_\rho^2$ and $(\epsilon_\mu^1 q_\nu^1 - \epsilon_\nu^1 q_\mu^1)(\epsilon^{2\mu} q^{2\nu} - \epsilon^{2\nu} q^{2\mu})$. Scalar versus pseudoscalar cases can be discriminated if photons are virtual and decay to e^+e^- , a classic example of the π^0 parity determination. This involves analyzing the angular distributions of fermions, which then yields the conclusion that the virtual photon polarizations were orthogonal. Thus, here the forward protons work in an analogous role with different central resonances, which then bootstrapped together with the central system decay products, yield the maximal information about the scattering dynamics and the spin of pomeron. Thus, the full information is in the multidimensional angular distributions.

The main open problems regarding this model parameters are related to the resonance-by-resonance couplings, which require full spectrum simulation compar-

isons with data simultaneously in several differential observables, which is possible. In principle at the LHC energies, one needs to take into account the screening loop (absorption) effects, which is also supported by our code. We return in the discussion to some of the properties of the model.

Jacob and Wick helicity amplitudes

To be able to generate maximally model-independent two-body angular distributions for arbitrary resonances and daughter spin-parities, we have encoded in the relativistic ‘wave function free’ Jacob and Wick helicity amplitude formalism [133]. Wave function free means that it involves no spinors, polarization vectors and so on. The formalism is abstract, it does not consider the underlying dynamic microscopic details. Because it is algorithmically very easy to rotate the frame or spin density matrix event-by-event, we provide several different spin quantization z -axis (rest frame) options for the event generation, for the definitions, see Section 5.4. We use the Collins-Soper frame (CS) as the default frame for calculating helicity amplitudes. Using this frame the two basic requirements hold: 1. Parity symmetry is manifest, 2. φ -angle dependence stays flat in the laboratory or CM rest frame, as required by rotational invariance.

Canonical and helicity states

In this formalism, there are two different state representations which are intimately connected. The two particle *canonical states* $|JJ_z ls\rangle$ and the *helicity states* $|JJ_z \lambda_1 \lambda_2\rangle$ are constructed from the single particle canonical states $|\vec{p}, jm\rangle_i$ and helicity states $|\vec{p}, j\lambda\rangle_i$ for $i = 1, 2$, respectively. Their relation is

$$|\vec{p}, j\lambda\rangle = D_{m\lambda}^{(j)}(\theta_R, \varphi_R)|\vec{p}, jm\rangle, \quad (5.78)$$

where the helicity rotation is defined by (θ_R, φ_R) in terms of Wigner D function. The explicit constructive definitions can be found in [134, 135]. The two particle helicity state is

$$\begin{aligned} |JJ_z \lambda_1 \lambda_2\rangle &= \mathcal{N}_J \sum_{m_1 m_2} \int d\Omega D_{J_z \lambda}^{*(J)}(\theta, \varphi) D_{m_1 \lambda_1}^{(s_1)}(\theta, \varphi) D_{m_2 - \lambda_2}^{(s_2)}(\theta, \varphi) |\theta \varphi m_1 m_2\rangle \quad (5.79) \\ &= \mathcal{N}_J \int d\Omega D_{J_z \lambda}^{*(J)}(\theta, \varphi) |\theta \varphi \lambda_1 \lambda_2\rangle, \quad (5.80) \end{aligned}$$

where $\Omega = (\cos \theta, \varphi)$ denotes the direction of the particle 1 and the normalization which gives simple unit completeness relations is

$$\mathcal{N}_J = \left(\frac{2J+1}{4\pi} \right)^{1/2}. \quad (5.81)$$

Now the explicit relation between the two particle states in the canonical and helicity basis are [135]

$$|JJ_z \lambda_1 \lambda_2\rangle = \sum_{ls} \langle \langle l, s, \lambda_1, \lambda_2 | J, s_1, s_2 \rangle \rangle |JJ_z ls\rangle \quad (5.82)$$

$$|JJ_z ls\rangle = \sum_{\lambda_1 \lambda_2} \langle \langle l, s, \lambda_1, \lambda_2 | J, s_1, s_2 \rangle \rangle |JJ_z \lambda_1 \lambda_2\rangle, \quad (5.83)$$

where the re-coupling coefficient double products are denoted with

$$\langle \langle l, s, \lambda_1, \lambda_2 | J, s_1, s_2 \rangle \rangle \equiv \mathcal{N}_J^l \langle J\lambda | ls0\lambda \rangle \langle s\lambda | s_1 s_2 \lambda_1, -\lambda_2 \rangle, \quad \boxed{\lambda = \lambda_1 - \lambda_2} \quad (5.84)$$

where minus sign between λ_1 and λ_2 comes naturally because individual helicities contribute in opposite directions. The normalization is

$$\mathcal{N}_J^l = \left(\frac{2l+1}{2J+1} \right)^{1/2}. \quad (5.85)$$

The two inner products denote Clebsch-Gordan $SU(2)$ decomposition coefficients, which we evaluate algorithmically via Wigner $3j$ symbols via Racah formula taking into account also the algebraic selection rules. The re-coupling coefficients simply connect the two basis states together

$$\langle J' J'_z l s | JJ_z \lambda_1 \lambda_2 \rangle = \langle \langle l, s, \lambda_1, \lambda_2 | J, s_1, s_2 \rangle \rangle \delta_{JJ'} \delta_{J_z J'_z}. \quad (5.86)$$

To make the notation clear, we have the following variables for the process with spins $J \rightarrow s_1 + s_2$:

$$\text{Angular momentum projection : } -J \leq J_z \leq J \text{ with } \mathbf{J} \equiv \mathbf{1} + \mathbf{s} \quad (5.87)$$

$$\text{Helicity : } \lambda_J \equiv \mathbf{J} \cdot \mathbf{p} / |\mathbf{p}| \quad (5.88)$$

$$\text{Daughter helicities : } -s_1 \leq \lambda_1 \leq s_1, -s_2 \leq \lambda_2 \leq s_2. \quad (5.89)$$

Note that in our notation $J_z \equiv M$, which we use to emphasize the physical meaning. In the helicity frame, one has $J_z = \lambda_J$. We emphasize that the quantum numbers l

and s are rotational invariants of the canonical basis, just like helicities are rotational invariants of the helicity basis. They are defined by equations for the total spin and fixed orbital angular momentum [135]

$$|\theta\varphi sm_s\rangle = \sum_{m_1 m_2} \langle sm_s | s_1 m_1 s_2 m_2 \rangle |\theta\varphi m_1 m_2\rangle \quad (5.90)$$

$$|lm sm_s\rangle = \int d\Omega Y_l^m(\theta, \phi) |\theta\varphi sm_s\rangle, \quad (5.91)$$

which makes the relations clear.

Density matrix

The resonance state spin polarization is encoded in a fixed spin density matrix ρ_i , which is a $(2J + 1)$ hermitian matrix obeying the standard von Neumann density matrix properties and expectation values of operators $\langle A \rangle_\rho = \text{Tr}[\rho A]$, where A can be a spin operator such one of the Pauli matrices for the spin-1/2 case or their generalization. The following properties of the density matrix hold always: A. $\text{Tr} \rho = 1$, B. $\rho_{ij}^* = \rho_{ji}$, C. $\rho_{ii} \geq 0$, D. $\text{Tr} \rho^2 \leq 1$, E. Positive semi-definite \leftrightarrow non-negative eigenvalues. In general, the matrix can be described by using $(2J + 1)^2 - 1$ real parameters. If the density matrix is very complicated, more convenient ways have been developed, see [134].

For a statistical mixture of *pure states*, giving only diagonal entries, the density matrix is

$$\rho = \sum_i p_i |J, J_z\rangle_i \langle J, J_z|_i \quad \text{s.t.} \quad \sum_i p_i = 1. \quad (5.92)$$

As an example a transverse only polarization for $J = 1$ is

$$\rho = \frac{1}{2} |1, -1\rangle \langle 1, -1| + \frac{1}{2} |1, 1\rangle \langle 1, 1| = \begin{bmatrix} 1/2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1/2 \end{bmatrix}. \quad (5.93)$$

Thus, the representation is

$$|1, -1\rangle \equiv [1, 0, 0]^T, \quad |1, 0\rangle \equiv [0, 1, 0]^T, \quad |1, 1\rangle \equiv [0, 0, 1]^T, \quad (5.94)$$

which is easily continued for higher spins. Off-diagonal elements in the density matrix are constrained by hermiticity and parity, and they generate the φ -angle dependence. In general, the off-diagonal elements are responsible for the quantum superposition (coherence). A user can freely parametrize the matrices or generate

completely random ones from Gaussian random matrix ensembles. Clearly, the spin density matrix is not covariant but its elements depend on the chosen Lorentz rest frame. An *Unpolarized* process has an equal probability for every helicity state, that is, a diagonal density matrix with elements $1/(2J + 1)$ which results in a uniform angular distribution in $(\cos\theta, \varphi)_{\text{r.f.}}$. In the most general case taking into account the parity conservation and hermiticity, $J = 1$ requires 4 and $J = 2$ requires 12 independent parameters.

A natural frame The most natural Lorentz frame for the given process, which always exists but might be non-trivial to know a priori, is the one which gives the most simple, diagonal spin density structure without azimuthal (off-diagonal) dependence. It can be argued to exist, in a mathematical sense, because the density matrix is a hermitian matrix which can be always diagonalized by a suitable rotation. By reciprocity, a process analyzed in an *unnatural frame* will have spurious off-diagonal dependence and mixing of moments. We believe that the CS frame is probably the most natural one, or very close, for central exclusive processes by analog with the Drell-Yan process.

Amplitudes

The decay dynamics is encoded in the helicity amplitude matrix, which is given by the linear combination [136]

$$T_{\lambda_1\lambda_2}^{(J)} = \sum_{ls} \alpha_{ls}^{(J)} \langle\langle l, s, \lambda_1, \lambda_2 | J, s_1, s_2 \rangle\rangle, \quad (5.95)$$

where $0 \leq s \leq s_1 + s_2$ and $0 \leq l \leq J + s$ with dimensions $(2s_1 + 1)(2s_2 + 1)$. This relation can be easily inverted, to obtain the canonical ls -representation coefficients in terms of the helicity amplitudes. The unknown coupling weights are denoted with α_{ls} , which are left as a user input. These are normalized such that

$$\sum_{\lambda_1\lambda_2} |T_{\lambda_1\lambda_2}^{(J)}|^2 = \sum_{ls} |\alpha_{ls}^{(J)}|^2 = 1. \quad (5.96)$$

The algorithm takes the sum over all allowed values of ls given the angular momentum conservation, parity and spin-statistics and gives user a list of the required $\alpha_{ls}^{(J)}$ coupling input. Here we remark that some combinations of quantum numbers have very simple α_{ls} coupling structure giving only $T = 1$, and thus only the spin polarization density matrix is unknown. This is the case for example with a $J^P = 2^+$ state into a pseudoscalar pair.

Now we have all the necessary input and the decay transition amplitude matrix element is given by [136]

$$f_{\lambda_1 \lambda_2, J_z}(\theta, \varphi) = \mathcal{D}_{\lambda J_z}^{(J)}(\theta, \varphi) T_{\lambda_1 \lambda_2}^{(J)}, \quad (5.97)$$

with dimensions $(2s_1 + 1)(2s_2 + 1) \times (2J + 1)$. Above, the Wigner spin rotation matrix in the spin space is

$$D_{mm'}^{(J)}(\theta, \varphi) = e^{im'\varphi} d_{mm'}^{(J)}(\theta), \quad (5.98)$$

written using the Wigner small d symbol matrix which we calculate algorithmically. The phase convention is fixed by $e^{im'\varphi}$.

Finally, the decay weight or decay amplitude squared of the event is given by the standard expectation value trace

$$|A|^2 = \mathcal{N}_J \text{Tr}[\rho_f] = \mathcal{N}_J \text{Tr}[f \rho_i f^\dagger], \quad \mathcal{N}_J = 2J + 1 \quad (5.99)$$

where the operator product maps the initial state ρ_i spin density matrix to the final state spin density matrix ρ_f , also known as the ‘Krauss operator’ map in quantum mechanics. With the normalization in use, we get $|A|^2 = 1$ for unpolarized decays. In addition, we can directly also calculate coherent spin correlated decay chains by tensor products. For example, $X \rightarrow A \rightarrow \{A_1 + A_2\} + B \rightarrow \{B_1 + B_2\}$ gives

$$f_{tot} = [f_A \otimes f_B] f_X \quad (5.100)$$

which are supported. The individual transition amplitudes f are evaluated in the corresponding rest frames of each decay.

To choose different quantization axes, the rotation along direction (θ_R, φ_R) of the quantization coordinate system is obtained via Wigner rotation matrices

$$|J, J_z\rangle = \sum_{J_z=-J}^J \mathcal{D}_{J_z J_z}^{(J)}(\theta_R, \varphi_R) |J, J_z\rangle. \quad (5.101)$$

The change of basis for the initial density matrix is obtained via a similarity transform

$$\rho'_i = \mathcal{D}^{\dagger(J)}(\theta_R, \varphi_R) \rho_i \mathcal{D}^{(J)}(\theta_R, \varphi_R), \quad (5.102)$$

which is used to change the reference frame of the spin polarization density. The rotation of the density matrix keeps its eigenvalues unchanged, thus also its von

CHAPTER 5. GRANIITI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

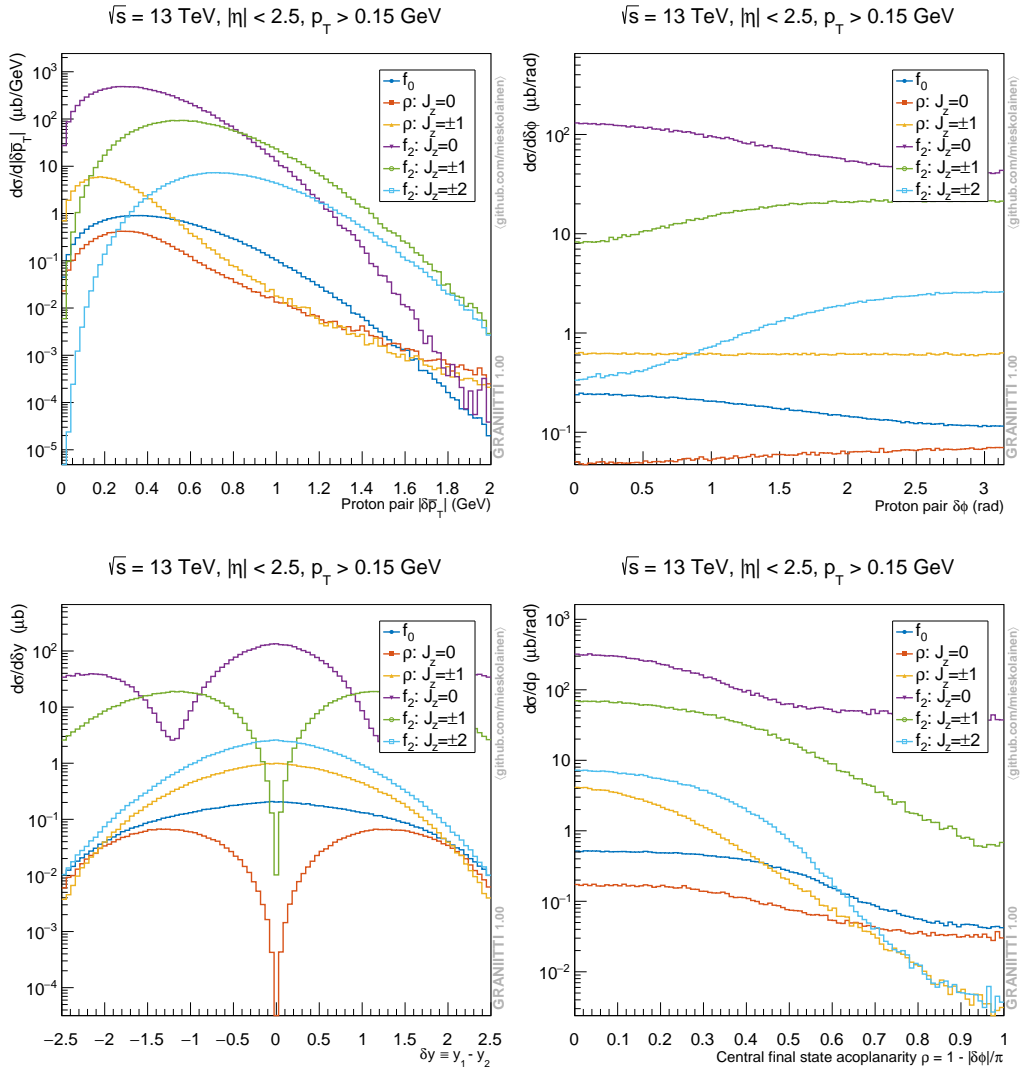


Figure 5.14: The glueball filter forward observable (top left), the forward proton pair transverse angle separation (top right), the central pion pair rapidity separation (bottom left) and the central pion pair acoplanarity (bottom right). A constant $\langle S^2 \rangle \equiv 0.15$ and 0.7 (photoproduction) applied.

CHAPTER 5. GRANIITTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

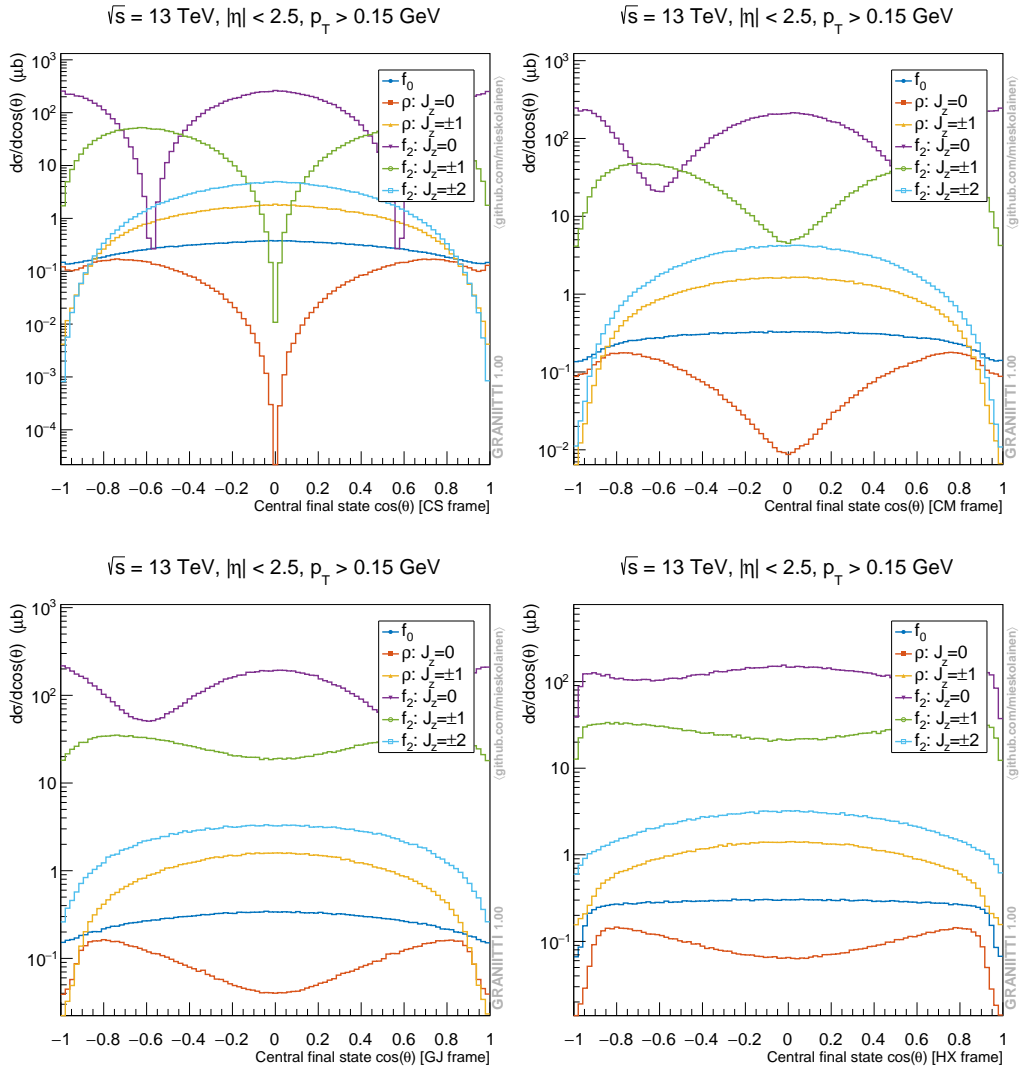


Figure 5.15: Central final state pion $\cos\theta$ measured in different rest frames. A constant $\langle S^2 \rangle \equiv 0.15$ and 0.7 (photoproduction) applied.

Neumann entropy. With parity conservation, the elements of the density matrix obey

$$\rho_{mm'} = (-1)^{m-m'} \rho_{-m-m'}. \quad (5.103)$$

So we need to remember that if we unsuitably rotate the density matrix, the parity conservation may not be manifest anymore, which is clear given that parity is a spatial symmetry.

In Figure 5.14 we demonstrate how the glueball filter $|\Delta\mathbf{p}_t|$, the forward proton transverse angle separation $\Delta\varphi_{pp}$, the central pion rapidity separation Δy and the central acoplanarity $1 - |\Delta\phi|/\pi$ are being driven by the spin polarization of $J = 1$ and $J = 2$ resonances. We see that the forward observables are strongly correlated with central observables, however, the analysis will be more difficult and ambiguous without forward protons. The $J = 2$ states with longitudinal and transverse polarization modes have opposite behavior in terms of forward azimuthal angle, also the glueball filter observable peaks in a different domain. Perhaps the glueball filter should be called a non-perturbative helicity amplitude filter. We can postulate a hypothesis that $J = 2$ glueballs may be produced dominantly with $J_z = 0$ and quark states with $|J_z| = 2$ polarization, but naturally they can have quantum mechanical mixing. In any case, we cannot say that the picture is complete at this point, especially without new data and models working truly at a non-perturbative parton level. Figure 5.15 shows the decay daughter $\cos\theta$ in different rest frames, which demonstrates clearly how different frames smear the distributions due to different rotations. Similarly, different frame rotations will induce non-flat φ -angle dependencies. We also point out that the tensor pomeron model cannot produce $|J_z| = 1$ polarization modes for $J = 2$ resonances, by angular momentum conservation.

Symmetries

We check algorithmically the required symmetries of the amplitudes to obtain the allowed subset of ls values:

Spin statistics The Bose-Einstein statistics requires $l - s$ to be even for identical boson pairs. The Fermi-Dirac statistics requires $l + s$ to be even for identical fermion pairs. These come simply from the symmetric wavefunction requirement for bosons and anti-symmetric for fermions.

Spatial parity [P] The parity operator or spatial inversion operator \hat{P} with $P|0\rangle = P^{-1}|0\rangle = |0\rangle$, which exchanges a left handed field to a right handed one, operates on the helicity and canonical states as [135]

$$\hat{P}|JJ_z\lambda_1\lambda_2\rangle = P_1P_2(-1)^{J-s_1-s_2}|JJ_z-\lambda_1-\lambda_2\rangle \quad (5.104)$$

$$\hat{P}|JJ_zls\rangle = P_1P_2(-1)^l|JJ_zls\rangle, \quad (5.105)$$

by flipping the sign of helicities λ_1, λ_2 , but not the angular momentum projection J_z , because the angular momentum is an axial-vector. Thus, with parity conservation, the helicity amplitudes obey a selection rule [135]

$$T_{\lambda_1\lambda_2} = PP_1P_2(-1)^{J-s_1-s_2}T_{-\lambda_1-\lambda_2}, \quad (5.106)$$

where P, P_1, P_2 are the parity ± 1 of the resonance and daughters. Also, the spherical harmonics are the eigenfunctions of parity

$$\hat{P}Y_l^m(\theta, \varphi) = Y_l^m(\pi - \theta, \varphi + \pi) = (-1)^l Y_l^m(\theta, \varphi). \quad (5.107)$$

Thus, the parity associates with the orbital angular momentum l .

Time reversal [T] The anti-unitary time operator \hat{T} operates to the helicity and canonical states as [135]

$$\hat{T}|JJ_z\lambda_1\lambda_2\rangle = P_1P_2(-1)^{J-J_z}|J-J_z\lambda_1\lambda_2\rangle \quad (5.108)$$

$$\hat{T}|JJ_zls\rangle = P_1P_2(-1)^l|J-J_zls\rangle. \quad (5.109)$$

by flipping the sign of J_z but not the helicity of decay daughters.

Charge conjugation [C] The charge parity operator \hat{C} operates by changing the sign of internal quantum numbers. This gives for boson and fermion pairs

$$\hat{C}|\pi^+\pi^-\rangle = (-1)^l(-1)^s|\pi^+\pi^-\rangle = (-1)^{l+s}|\pi^+\pi^-\rangle \quad (5.110)$$

$$\hat{C}|f\bar{f}\rangle = (-1)^l(-1)^{s+1}(-1)|f\bar{f}\rangle = (-1)^{l+s}|f\bar{f}\rangle, \quad (5.111)$$

where the charge conjugation action operates on the orbital part like parity and the third factor in the Fermi case is the particle statistics requirement.

5.3 Kinematics and Monte Carlo sampling

We follow along the lines of the exact 4-body phase space construction suitable for diffraction used in [110, 137], but extend it to include forward proton excitation and generalize it from $N = 4$ process to the case N in two ways: using the exact phase space factorization and a ladder-type direct $2 \rightarrow N$ construction.

Skeleton kinematics

The standard QFT cross section in terms of the phase space and amplitude squared is

$$\sigma = \frac{1}{F} \frac{1}{S} \int \prod_{i=1}^N \frac{d^3 p_i}{(2\pi)^3 2E_i} (2\pi)^4 \delta^{(4)} \left(p_A + p_B - \sum_{f=1}^N p_f \right) |\mathcal{A}_{2 \rightarrow N}|^2, \quad (5.112)$$

where the Möller flux is $F = 4\sqrt{(p_A \cdot p_B)^2 + m_A^2 m_B^2} \simeq 2s$ and S is the statistical QFT symmetry factor to take into account identical final states. Now using the relation from Cartesian to collider variables

$$E \frac{d^3 \sigma}{d^3 p} = \frac{d^3 \sigma}{d\varphi dy p_t dp_t} \leftrightarrow \frac{d^3 p}{2E} = \frac{1}{2} d\varphi dy p_t dp_t, \quad (5.113)$$

we can turn the sampling over 3-momentum into rapidity, transverse momentum and azimuthal angle. Above, one uses the identity $dy/dp_z = 1/E$.

The total number of Lorentz scalars or non scalar variables needed for $2 \rightarrow N$ process is $3N - 4$ for $N \geq 2$. Thus a $2 \rightarrow 3$ process needs 5 variables, which we use as our starting point. We can eliminate redundant variables using the energy-momentum conservation. By Lorentz invariance of the expressions, let us work in the frame where $\sum_{i=1}^3 \vec{p}_i = \vec{0}$ and write

$$d\sigma = \frac{1}{F} \frac{1}{S} (2\pi)^4 \delta(E_1 + E_2 + E_3 - \sqrt{s}) \times \delta^{(2)}(\vec{p}_{t,1} + \vec{p}_{t,2} + \vec{p}_{t,3}) \delta(p_{z,1} + p_{z,2} + p_{z,3}) \prod_{i=1}^3 \frac{d^3 p_i}{(2\pi)^3 2E_i}. \quad (5.114)$$

1. Eliminate $d^2 p_{t,3}$ by $\vec{p}_{t,3} = -(\vec{p}_{t,1} + \vec{p}_{t,2})$ dependence using an implicit integral over the delta function

$$d\sigma = \frac{1}{F} \frac{1}{S} \frac{(2\pi)^4}{(2\pi)^9} \delta(E_1 + E_2 + E_3 - \sqrt{s}) \delta(p_{z,1} + p_{z,2} + p_{z,3}) \frac{d^3 p_1}{2E_1} \frac{d^3 p_2}{2E_2} \frac{dp_{z,3}}{2E_3}. \quad (5.115)$$

CHAPTER 5. GRANIETTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

2. Eliminate $dp_{z,2}$ by $p_{z,2} = -(p_{z,1} + p_{z,3})$ dependence using $d^3p = d^2p_t d_z = d\varphi p_t dp_t d_z$ and an implicit integral over the delta function

$$d\sigma = \frac{1}{F} \frac{1}{S} \frac{1}{(2\pi)^5} \delta(E_1 + E_2 + E_3 - \sqrt{s}) \frac{d^3p_1}{2E_1} \frac{d\varphi_2 p_{t,2} dp_{t,2}}{2E_2} \frac{dp_{z,3}}{2E_3}. \quad (5.116)$$

3. Then treat the last energy conservation delta function

$$d\sigma = \frac{1}{F} \frac{1}{S} \frac{1}{(2\pi)^5} \delta(f(p_{z,1})) [d\varphi_1 \frac{p_{t,1}}{2E_1} dp_{t,1} dp_{z,1}] [d\varphi_2 \frac{p_{t,2}}{2E_2} dp_{t,2}] [\frac{dp_{z,3}}{2E_3}], \quad (5.117)$$

where we denote

$$f(p_{z,1}) = \sqrt{M_{1,t}^2 + p_{z,1}^2} + \sqrt{M_{2,t}^2 + p_{z,2}^2} + E_3 - \sqrt{s} \quad (5.118)$$

$$p_{z,2}(p_{z,1}) = -(p_{z,1} + p_{z,3}) \quad (5.119)$$

with the transverse mass $M_t^2 \equiv M^2 + p_t^2 = E^2 - p_z^2$. Then, we obtain a factor

$$|\Delta| = \left| \frac{df}{dp_{z,1}} \right| = \left| \frac{p_{z,1}}{\sqrt{M_{t,1}^2 + p_{z,1}^2}} - \frac{p_{z,2}}{\sqrt{M_{t,2}^2 + p_{z,2}^2}} \right| = \left| \frac{p_{z,1}}{E_1} - \frac{p_{z,2}}{E_2} \right|, \quad (5.120)$$

which is in most kinematic cases approximately 2. The need for this is based on the relation

$$\delta(f(x)) = \sum_{x_i: f(x_i)=0} \delta(x - x_i) \left| \frac{df(x_i)}{dx_i} \right|^{-1}, \quad (5.121)$$

where the sum runs over solutions of $f(x) = 0$ (roots). We leave the sum implicit in the notation, because we will use only one root, as we will see.

4. Finally, the change of a variable with $dy_3 = dp_{z,3}/E_3$ gives

$$d\sigma = \frac{1}{F} \frac{1}{S} \frac{1}{(2\pi)^5} |\Delta|^{-1} [d\varphi_1 \frac{p_{t,1}}{2E_1} dp_{t,1}] [d\varphi_2 \frac{p_{t,2}}{2E_2} dp_{t,2}] [\frac{1}{2} dy_3]. \quad (5.122)$$

The variables which are thus left to be sampled are the forward system φ_1 , φ_2 , $p_{t,1}$, $p_{t,2}$ and the central system rapidity y_3 . Then, to be able to include variable invariant masses for the forward and central legs, we sample over M_1^2, M_2^2 of the forward systems and over M_3^2 of the central system, the squared masses. The overall Monte Carlo event phase space weight of the main skeleton kinematics is

$$W_{2 \rightarrow 3} = V_{2 \rightarrow 3} \frac{1}{F} \frac{1}{S} \frac{1}{2} \frac{1}{(2\pi)^5} \frac{p_{t,1}}{2E_1} \frac{p_{t,2}}{2E_2} |\Delta|^{-1}. \quad (5.123)$$

The sampling volume is

$$V_{2 \rightarrow 3} = [p_{t,1}] \times 2\pi \times [p_{t,2}] \times 2\pi \times [y_3] \times [M_3^2], \quad (5.124)$$

where $[x] \equiv |x_{\max} - x_{\min}|$ denotes the sampling interval. In addition, one includes the sampling volumes of M_1^2 and M_2^2 , if excitation is included. Also, we need the phase space factor related with the central system phase space, which we will calculate in two ways.

Kinematic polynomials

The variables $p_{z,1}, p_{z,2}, E_1, E_2$ of the event skeleton kinematics are found in a closed form by solving the non-linear system of equations

$$0 = p_{z,1} + p_{z,2} + p_{z,3} \quad (5.125)$$

$$s^{1/2} = E_1 + E_2 + E_3 \quad (5.126)$$

$$E_1^2 = M_1^2 + p_{z,1}^2 + p_{t,1}^2 \quad (5.127)$$

$$E_2^2 = M_2^2 + p_{z,2}^2 + p_{t,2}^2. \quad (5.128)$$

The resulting expressions are very lengthy due to the forward legs and can be found in the code, together with the symbolic machine algebra code solution of the non-linear system, which we used to generate the corresponding C++ code. The polynomial root branch which results in a non-flip of the forward-backward momentum, is chosen.

1. Using the chosen polynomial branch, we calculate

$$p_{z,1} = \text{sol}(M_1, M_2, p_{t,1}, p_{t,2}, p_{z,3}, E_3), \quad (5.129)$$

where the central system energy and longitudinal momentum are

$$\{E_3, p_{z,3}\} = \{M_{t,3} \cosh(y_3), M_{t,3} \sinh(y_3)\}, \quad (5.130)$$

which are obtained in terms of the transverse mass, by solving the central system transverse momentum from the forward system transverse variables, by momentum conservation.

2. Then we get by momentum conservation $p_{z,2} = -(p_{z,1} + p_{z,3})$. The variables E_1 and E_2 are then obtained directly by substitution.

Factorized phase space

To be able to include the central system phase space, we use the exact factorization relation of the phase space

$$\begin{aligned} d^N \Pi(s; p_1, p_2, \dots, p_N) \\ = \frac{1}{2\pi} dM_X^2 d^3 \Pi(s; p_1, p_2, p_X) d^{N-2} \Pi(M_X^2; p_3, p_4, \dots, p_N), \end{aligned} \quad (5.131)$$

where p_1, p_2 are the outgoing forward legs, p_X the central system 4-momentum with $M_X^2 \equiv p_X^2$ and $d^N \Pi$ abstracts the corresponding Lorentz invariant phase space measure. The integral over M_X^2 represents the integral over the central system mass squared. The central system flat $1 \rightarrow N - 2$ phase space is constructed recursively following the classic algorithm of Kopylov and Raubold-Lynch described by James in [138], which calculates also the exact phase space volume weight $W_{1 \rightarrow N-2}$. The basic idea behind the algorithm is to split the phase space into $N - 2$ sequential $1 \rightarrow 2$ decays with intermediate masses, for the explicit details of this well known algorithm, we refer reader to the program code. The two body phase space is nearly trivial and thus works as the building block. The total weight of the event is now

$$W_{2 \rightarrow N} = W_{2 \rightarrow 3} \frac{1}{2\pi} W_{1 \rightarrow N-2}. \quad (5.132)$$

The classic algorithm can be plug-in replaced easily with alternative algorithms, such as variants of RAMBO [139]. Also, we have a simple ‘chain recursive’ phase space implemented which can be useful for long decay chains with intermediate propagators. However, using it requires some care due to intermediate-mass squared sampling. If the matrix element of the process contains all information about the intermediate states, then the sampling should be done with flat masses squared within reasonable ranges, given that the final state leg permutations for different sub-amplitudes may probe different mass regimes in the phase space. Alternatively, a $1 \rightarrow N - 2$ central phase space can always be constructed, which is safe but with low efficiency if N is large. Finally, if the matrix element does not contain full decay chain information, then a relativistic Breit-Wigner sampling in mass squared is applied as a simple dynamic propagator model.

Direct phase space

As an alternative formulation of the phase space, instead of the factorized phase space, we have constructed a ‘direct’ $2 \rightarrow N$ kinematics based on solving a certain linear system of equations. Let us denote the number of central states with $K =$

$N - 2$. For the transverse degrees of freedom, we need $K - 1$ transverse momentum k_t variables, $K - 1$ azimuthal variables φ . For the longitudinal degrees of freedom, we need K rapidity y variables as our Monte Carlo sampling variables, in addition to the 4 forward system variables as before. Thus, the total number of variables is $3N - 4$, which is the minimum possible.

Let us have so-called difference momentum transverse vectors

$$\vec{k}_{t,j} = (k_{t,j} \cos \varphi_j, k_{t,j} \sin \varphi_j), \quad j = 1, \dots, K - 1, \quad (5.133)$$

which we use as the basis of the construction. Let us then write a system of equations

$$\mathbf{b} = A\mathbf{p}, \quad (5.134)$$

where \mathbf{p} denotes the system vector of central final states. This is solved separately for x and y components, which we leave implicit in the notation below. We construct the system vector of difference momentum

$$b_j = \begin{cases} \vec{k}_{t,j} - \vec{w}, & \text{when } j = 1 \\ -\vec{k}_{t,j-1} - \vec{w}, & \text{when } j = 2, \dots, K, \end{cases} \quad (5.135)$$

where the forward system transverse vector sum is $\vec{w} = \vec{p}_1 + \vec{p}_2$. Then, we can solve the central final state transverse momentum components by

$$\mathbf{p} = A^{-1}\mathbf{b}. \quad (5.136)$$

The system matrices are full rank with components

$$A_2 = \begin{pmatrix} 2 & 0 \\ 0 & 2 \end{pmatrix}, A_3 = \begin{pmatrix} 2 & 0 & 1 \\ 0 & 2 & 1 \\ 1 & 0 & 2 \end{pmatrix}, A_4 = \begin{pmatrix} 2 & 0 & 1 & 1 \\ 0 & 2 & 1 & 1 \\ 1 & 0 & 2 & 1 \\ 1 & 1 & 0 & 2 \end{pmatrix} \dots, \quad (5.137)$$

which can be constructed with a simple algorithm up to any K and the inverses are taken by symbolic machine algebra and saved. For example

$$A_2^{-1} = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}, A_3^{-1} = \begin{pmatrix} \frac{2}{3} & 0 & -\frac{1}{3} \\ \frac{1}{6} & \frac{1}{2} & -\frac{1}{3} \\ -\frac{1}{3} & 0 & \frac{2}{3} \end{pmatrix}, A_4^{-1} = \begin{pmatrix} \frac{7}{8} & \frac{1}{8} & -\frac{1}{2} & -\frac{1}{4} \\ \frac{1}{8} & \frac{1}{8} & -\frac{1}{2} & -\frac{1}{4} \\ -\frac{1}{8} & \frac{1}{8} & -\frac{1}{2} & -\frac{1}{4} \\ -\frac{1}{8} & -\frac{1}{8} & -\frac{1}{2} & \frac{3}{4} \end{pmatrix}, \quad (5.138)$$

which demonstrate the algebraic structures. Finally, the longitudinal momentum and energy for the central final states are obtained with

$$p_{z,j} = m_{t,j} \sinh y_j \quad (5.139)$$

$$E_j = m_{t,j} \cosh y_j, \quad \text{for } j = 1, \dots, K, \quad (5.140)$$

where $m_{t,j}^2 = m_j^2 + p_{t,j}^2$. We sample K rapidity variables y_j independently, which then fix the central system rapidity and we can proceed with the skeleton kinematics polynomials.

Then proceeding in a same way as in the $2 \rightarrow 3$ case, the total Monte Carlo phase space weight is

$$W_{2 \rightarrow N} = V_{2 \rightarrow N} \frac{1}{F} \frac{1}{S} \frac{1}{2^{2(N-2)}} \frac{1}{(2\pi)^{3N-4}} \frac{p_{t,1}}{2E_1} \frac{p_{t,2}}{2E_2} |\Delta|^{-1} \prod_{j=1}^{K-1} k_{t,j}, \quad \text{for } N \geq 4, \quad (5.141)$$

where the sampling volume is

$$V_{2 \rightarrow N} = [p_{t,1}] \times 2\pi \times [p_{t,2}] \times 2\pi \times [y_j]^K \times ([k_{t,j}] \times 2\pi)^{K-1}. \quad (5.142)$$

Again, if the forward excitation is included, the sampling volumes of M_1^2 and M_2^2 are inserted. Let us point out that we have found this space construction to be easily unstable with VEGAS with high leg count $N \geq 4$ Regge like amplitudes, due to the complicated integration boundaries and a non-trivial structure of the high dimensional Lorentz manifold. Typically, a larger number of integrand evaluations has somewhat stabilized the behavior. In practice we recommend the factorized phase space as the default stable option. However, this phase space construction provides a cross check and may turn out to be of good use with alternative importance sampling techniques. In general, we check all kinematic algorithms against the well known exact volume formulas for the massive two body case and the massless N -body case.

Also, as the simplest possible construction, we include a collinear phase space option $2 \rightarrow N$, where N is the number of final states excluding proton remnants. This phase space is suitable for simple amplitudes convoluted with parton densities or collinear EPA fluxes.

Monte Carlo sampling

The basic idea behind Monte Carlo integration with importance sampling in n -dimensional space is

$$I = \int_{\Omega} d^n \mathbf{x} f(\mathbf{x}) = \int_{\Omega} d^n \mathbf{x} \frac{f(\mathbf{x})}{q(\mathbf{x})} q(\mathbf{x}) \simeq \frac{V_{\Omega}}{N} \sum_{i=1}^N \frac{f(\mathbf{x}_i)}{q(\mathbf{x}_i)}, \quad \text{when } N \rightarrow \infty, \quad (5.143)$$

where $V_\Omega = \int_\Omega d^n \mathbf{x}$ is the integration boundary volume, typically a box volume and N is the number of samples. So instead of sampling \mathbf{x} uniformly from $[0, 1]^n$, we sample \mathbf{x} according to $q(\mathbf{x})$ and then compensate for this in the integral by weighting events with $1/q(\mathbf{x})$. When $q(\mathbf{x}) \rightarrow 0$, no samples are generated, thus there is no division by zero. However, unless $f(\mathbf{x})$ also goes to zero simultaneously, the integral will be biased. Also the normalization

$$\int d\mathbf{x} q(\mathbf{x}) = 1 \quad (5.144)$$

needs to hold, which usually needs to be estimated simultaneously. We want to find out the optimal importance sampling pdf

$$q_{\text{opt}}(\mathbf{x}) \equiv \frac{|f(\mathbf{x})|}{\int d^n \mathbf{x} |f(\mathbf{x})|}. \quad (5.145)$$

This would give a vanishing variance for the integral estimate, however, the problem of adaptive learning of $q(\mathbf{x})$ is non-trivial.

For sampling the phase space and integrating cross sections, the engine includes a fully multithreaded implementation of classic VEGAS adaptive importance sampling [140], where multithreading is implemented using C++17 standard library threading support by distributing the integrand samples over a fixed number of threads. We have tested the scalability up to thousands of threads. The correctness of VEGAS implementation can be cross-checked with a naive flat sampling mode. In a standard Monte Carlo way, we operate over a unit hypercube $[0, 1]^n$ and scale and shift each dimension with $x_i \rightarrow a_i + (b_i - a_i)x_i$ to the interval $[a_i, b_i]$. Because describing the correlations in high dimensional phase space is in general highly non-trivial and requires neural networks or similar techniques, VEGAS takes a factorized simplification

$$q(\mathbf{x}) = \prod_{i=1}^n q_i(x_i). \quad (5.146)$$

Clearly, neglecting correlations gives bad efficiency if the process kinematics \times matrix element squared does not ‘align’ along the dimensions. The complexity scales as $\mathcal{O}(nB)$, where B is the number of bins per dimension. In contrast, full n -dimensional histogram representation would scale exponentially fast $\mathcal{O}(B^n)$. Other classic alternatives or extensions are mixture model importance densities, also known as multi-channeling in high energy physics.

VEGAS histogram grids for each dimension need to be initialized over a few number of iterations. The number of iterations R and the number of samples per

iteration N_k in the initialization burn-in phase and in the integration phase are free parameters of the algorithm. In the integration phase, we use also a maximum relative error $\sigma_I/\langle I \rangle$ target as a criterion.

For each k -th iteration, the integral estimate and its variance are

$$\langle I_k \rangle = \frac{1}{N_k} \sum_{i=1}^{N_k} \frac{f(\mathbf{x})}{q^{(k)}(\mathbf{x})}, \quad \sigma_{I_k}^2 = \frac{1}{N_k - 1} \sum_{i=1}^{N_k} \left[\left(\frac{f(\mathbf{x}_i)}{q^{(k)}(\mathbf{x}_i)} \right)^2 - \langle I_k \rangle^2 \right], \quad (5.147)$$

where one accumulates the sum of values $f(\mathbf{x})/q(\mathbf{x})$ and their squares during the operation. We evaluate the quality of the set of integral estimates using the χ^2 test

$$\chi^2/\text{ndf} = \frac{1}{N_k - 1} \sum_{k=1}^R \frac{(\langle I_k \rangle - \langle I \rangle)^2}{\sigma_{I_k}}, \quad (5.148)$$

with values close to unity being an indicator of solid results. Above, the global estimate of the integral is the weighted sum

$$\langle I \rangle = \left[\sum_{k=1}^R \frac{N_k}{\sigma_{I_k}} \right]^{-1} \sum_{k=1}^R \frac{N_k \langle I_k \rangle}{\sigma_{I_k}}. \quad (5.149)$$

The binning algorithm operates with a fixed number of bins per dimension and shifts the bin boundaries during the operations according to the original description [140]. The stability of the re-binning is controlled with an additional regularization parameter λ . There are also variants of VEGAS, where stratified sampling is combined with importance sampling, but we did not find them effective enough to compensate for the additional complexity. The unweighted event generation is based on a standard hit-and-miss, where estimate for the crucial maximum weight $\max [f(\mathbf{x})/q(\mathbf{x})]$ is obtained during the pre-event generation phase. After the event generation, the user is given the statistics of events overshooting the maximum weight, thus indicating a need for a longer pre-event generation phase.

In addition to VEGAS importance sampling, variables related to steeply falling spectra such as the forward system invariant masses are MC sampled in log space together with the corresponding Jacobian inside the integrand, often with much improved behavior. The rest of the standard kinematics not described here is based typically on heavy use of the Källén triangle function. To point out, we have experimented with deep learning techniques, similar to [141], which could provide superior scaling in higher dimensions and with difficult scattering amplitudes. The results are promising and we may expect these techniques for learning the distribution $q(\mathbf{x})$ to be included in the future versions.

5.4 Analysis engine

The analysis engine includes highly efficient plotting machinery to gain quick understanding of fiducial observables for different processes and theoretical constructions, such as eikonal densities. Naturally, these are also part of the automated test suite to control the quality of the code at a high level.

Lorentz rest frames

There is an infinite number of different Lorentz rest frames for the system X , obtained by fixed or event-by-event kinematics dependent $SO(3)$ rotations. However, certain Lorentz rest frames have more special properties than others. These different frames have been originally designed to be more ‘natural’ either for s -channel or t -channel dominated processes or mitigate the effects of the system transverse momentum, which is the case with Collins-Soper frame. In practice, it is trivial and highly recommended to repeat the analysis in multiple frames. Different frames give different projections of the angular distributions and spherical moments, by definition.

Let us have beam protons in the lab frame p_1^{lab} and p_2^{lab} and their boosted versions p_1^X, p_2^X in the system X rest frame. We now define a set of frames X' , which are related to the frame X by a rotation. The definitions of the z -axis are as follows

$$\text{CM: } \mathbf{z} = [0, 0, 1]^T \quad (5.150)$$

$$\text{HX: } \mathbf{z} = u(-(\mathbf{p}_1^X + \mathbf{p}_2^X)) \quad (5.151)$$

$$\text{CS: } \mathbf{z} = u(u(\mathbf{p}_1^X) - u(\mathbf{p}_2^X)) \quad (5.152)$$

$$\text{AH: } \mathbf{z} = u(u(\mathbf{p}_1^X) + u(\mathbf{p}_2^X)) \quad (5.153)$$

$$\text{PG: } \mathbf{z} = u(\mathbf{p}_1^X) \text{ or } \mathbf{z} = u(\mathbf{p}_2^X), \quad (5.154)$$

where $u(\mathbf{x}) \equiv \mathbf{x}/\|\mathbf{x}\|$ returns a unit vector.

Center of Momentum (CM) The quantization z -axis in the center of momentum X' is the same as in the X frame. No rotation is involved in the transformation to this frame, only a boost from the colliding proton-proton beam frame (lab).

Helicity (HX) The quantization z -axis is defined by the system X momentum vector direction in the lab frame, or equivalently, by the negative direction of the sum of the initial state proton \mathbf{p}_1^X and \mathbf{p}_2^X momentum in the system X rest frame [142]. This is a common analysis frame in quarkonium studies.

Collins-Soper (CS) The quantization z -axis is defined by the bisector vector between the initial state proton \mathbf{p}_1^X and $-\mathbf{p}_2^X$ (negative) momentum in the system X rest frame. This originated from the context of Drell-Yan process [143], but is not limited to it.

Anti-Helicity (AH) The quantization z -axis is defined by the bisector vector between the initial state \mathbf{p}_1^X and \mathbf{p}_2^X (positive) momentum in the system X rest frame. This frame may be interesting for pure symmetry reasons, because it is perpendicular to the CS frame.

Pseudo-Gottfried-Jackson (PG) The quantization axis defined by the initial state proton \mathbf{p}_1^X (or \mathbf{p}_2^X) momentum vector in the system X rest frame [144]. Note that sometimes, this is known directly as the Gottfried-Jackson frame.

For all frames except CM which has $\mathbf{y} = [0, 1, 0]^T$, we define the y -axis as the normal vector from the plane spanned by the initial states

$$\mathbf{y} = u(u(\mathbf{p}_1^X) \times u(\mathbf{p}_2^X)). \quad (5.155)$$

Finally, the x -axis is obtained by taking the cross product

$$\mathbf{x} = \mathbf{y} \times \mathbf{z}, \quad (5.156)$$

with the axes being orthonormal. These give us a rotation matrix

$$R = [\mathbf{x}, \mathbf{y}, \mathbf{z}]^T, \quad (5.157)$$

which we apply to all boosted final states in the system X

$$\mathbf{p}_i^{X'} \leftarrow R\mathbf{p}_i^X \quad (5.158)$$

to transform them to the new frame X' .

Forward proton dependent frames The quantization z -axis may be defined by the momentum transfer vector $q_1^{lab} = p_1^{lab} - p_1'^{lab}$ (or $q_2^{lab} = p_2^{lab} - p_2'^{lab}$) momentum boosted to the system X rest frame, also known as the Gottfried-Jackson (GJ) frame. In the lab frame the production plane is spanned by \mathbf{q}_1^{lab} and \mathbf{q}_2^{lab} . In the GJ rest frame the momentum transfer vectors are back-to-back $\mathbf{q}_1^{GJ} = -\mathbf{q}_2^{GJ}$ along the z -axis. The transform to this frame from the lab can be obtained by two rotations after the boost, which fix also the x - and y -axes. We note here that the boosts and rotations do not commute, so the order counts.

Density matrix in terms of spherical tensors

It is often useful to represent the density matrix in terms of generalized spin- J operators, also known as spherical tensors [134]

$$[T_l^m]_{J_z J'_z} \equiv \langle J J_z | \hat{T}_l^m | J J'_z \rangle \equiv \langle J J_z | J J'_z; l m \rangle, \quad 0 \leq l \leq 2J, \quad -l \leq m \leq l, \quad (5.159)$$

defined in terms of Clebsch-Gordan coefficients. However, we note here that there can be many other representations too, this one being one of the most common. The rank of the tensor operator is denoted with l . Now the matrix valued multipole expansion is

$$\rho = \frac{1}{2J+1} \sum_{l,m} (2l+1) t_l^{*m} T_l^m \quad (5.160)$$

with expansion coefficients (multipole parameters)

$$t_l^{*m} = \sum_{J_z, J'_z} \langle J J_z | J J'_z; l m \rangle \rho_{J_z J'_z} \quad \text{with} \quad t_{-l}^m = (-1)^m t_l^{*m}. \quad (5.161)$$

These are normalized with $t_0^0 = \text{Tr} \rho = 1$, such that the expectation value is $\text{Tr}(\rho T_l^m) = t_l^m$ or $\text{Tr}(\rho T_l^{\dagger m}) = t_l^{*m}$. Next we describe how to extract the expansion coefficients from the data.

Spherical harmonics inverse expansion

The angular distribution expansion engine is based on a complete spherical harmonics expansion described in [145]. We shall re-derive it, add some extra rigor regarding the phase spaces and inversion algorithms and clarify some aspects relevant at high energies. Experimentally, with forward protons and Mandelstam t_1, t_2 measured, one is interested in general in the multidifferential cross section

$$\frac{d^5 N}{dt_1 dt_2 dM^2 dY d\Omega} \sim \mathcal{I}(\Omega; t_1, t_2, M^2, Y). \quad (5.162)$$

Alternatively, without forward protons and also integrating over typically flat rapidity Y dependence of the process in the central domain, one has

$$\frac{d^3 N}{dM^2 dP_t d\Omega} \sim \mathcal{I}(\Omega; M^2, P_t), \quad (5.163)$$

where $\Omega \equiv (\cos \theta, \varphi)$ is the decay daughter momentum direction in the chosen rest frame of the system X with invariant mass squared M^2 and transverse momentum P_t . For every single crucial kinematic variable which is integrated (summed) over, a Monte Carlo event generator dependent bias is being induced through the acceptance expansion – unless the event generator is one-to-one with reality. This event generator dependence is even higher with naive single dimensional histogram based efficiency corrections, naturally. That is, fully MC generator independent results can be obtained *only* through fully differential hyperbinning, or by fully differential continuum techniques such as our *DeepEfficiency* which is based on inverting the high dimensional efficiency response via Deep Learning [146].

Phase spaces

We need three different spaces: a detector space Ω^{det} , a fiducial phase space Ω^{fid} and an angular flat phase space Ω^{flat} . Our definition of the detector space contains the reconstructed and selected events *with the fiducial cuts applied*, thus this space includes any finite efficiency effects. Naturally, it includes also possible biases (offset) and variance (resolution) effects of the track momentum measurements, which can be corrected a posteriori after the spherical expansion, if needed, by unfolding the spectrum histograms. We remark here that in general, there can be also measured events outside the fiducial domain, such as events with very low transverse momentum tracks, but still reconstructed. Thus, one needs careful definitions. The pure fiducial phase space contains the geometric η -acceptance of the detector equipped with a minimum p_t -cutoff, basically the geometric ‘ideal’ of the detector space and minimally extrapolating. The angular flat phase space is the space that leaves scalar decays uniform over the solid angle, for example, a limited range on the system rapidity – any cut applied only on the system will leave scalar decays uniform, by Lorentz invariance. One could take also the full 4π solid angle space as the flat phase space definition, but that usually means massive extrapolation from the detector (or fiducial) phase space at high energy experiments, easily a factor of 10. An example of a practical definition is

$$\text{Angular flat phase space } \Omega^{\text{flat}} = \{|Y_X| < 2.5\}. \quad (5.164)$$

$$\text{Fiducial phase space } \Omega^{\text{fid}} = \{|\eta| < 2.5 \wedge p_t > 0.15 \text{ GeV}\} \quad (5.165)$$

$$\text{Detector phase space } \Omega^{\text{det}} = \Omega^{\text{fid}} \otimes \text{detector response function}. \quad (5.166)$$

That is, the system’s rapidity limit is taken the same as the pseudorapidity of the decay daughters, which is fine given that always monotonically $|\eta| \geq |y|$ for single

final states, which bounds the system rapidity. That is, we must have

$$\Omega^{\text{det}} \subseteq \Omega^{\text{fid}} \subseteq \Omega^{\text{flat}} \subseteq \Omega^{4\pi}, \quad (5.167)$$

so that all events in the detector space belong to the fiducial phase space, and all in the fiducial phase space belong to the flat phase space. In Eq. 5.167, the first relation dictates the efficiency corrections, the second relation dictates the amount of geometric-kinematic inversion (extrapolation) done by the spherical harmonics expansion and the last relation is obvious. The ratio $|\Omega^{\text{fid}}|/|\Omega^{\text{flat}}|$ is order of 0.4 - 0.7 at the LHC for our processes with the type of definition given by Eqs. 5.165 and 5.164. The generator level ‘ground truth’ sample needs to contain cuts of the flat phase space. Our code includes examples that illustrate this straightforward but crucial detail. The crucial thing to understand is that the spherical moments are mixed in the pure fiducial phase space, unless the detector acceptance is extremely good. This is the ultimate reason to use the angular flat phase as the inversion target, even in principle the fiducial measurements should be always the primary target, because they minimize data extrapolations, by construction.

Spherical harmonics

The ‘intensity’ function \mathcal{I} describes the angular distribution. This function is expanded in terms of spherical harmonic moment expansion

$$\mathcal{I}(\Omega) = \sum_{l=0}^{l_{\text{max}}} \sum_{m=-l}^l t_{lm} Y_{lm}(\Omega), \quad (5.168)$$

where we use a *real valued representation* of the Laplace spherical harmonics, which is an alternative to the complex representation. We illustrate these functions in Figure 5.16. Both the real and complex representation provide a full orthonormal basis for square-integrable functions. The conversion from the complex representation is

$$Y_{lm} \equiv \text{Real}[Y_l^m] \equiv \begin{cases} \sqrt{2}(-1)^m \text{Im}[Y_l^m] & \text{if } m < 0 \\ Y_l^0 & \text{if } m = 0 \\ \sqrt{2}(-1)^m \text{Re}[Y_l^m] & \text{if } m > 0. \end{cases} \quad (5.169)$$

The conversion in the inverse direction is

$$Y_l^m \equiv \text{Complex}[Y_{lm}] \equiv \begin{cases} \frac{1}{\sqrt{2}}(Y_{l|m|} - iY_{l,-|m|}) & \text{if } m < 0 \\ Y_{l0} & \text{if } m = 0 \\ \frac{(-1)^m}{\sqrt{2}}(Y_{l|m|} + iY_{l,-|m|}) & \text{if } m > 0, \end{cases} \quad (5.170)$$

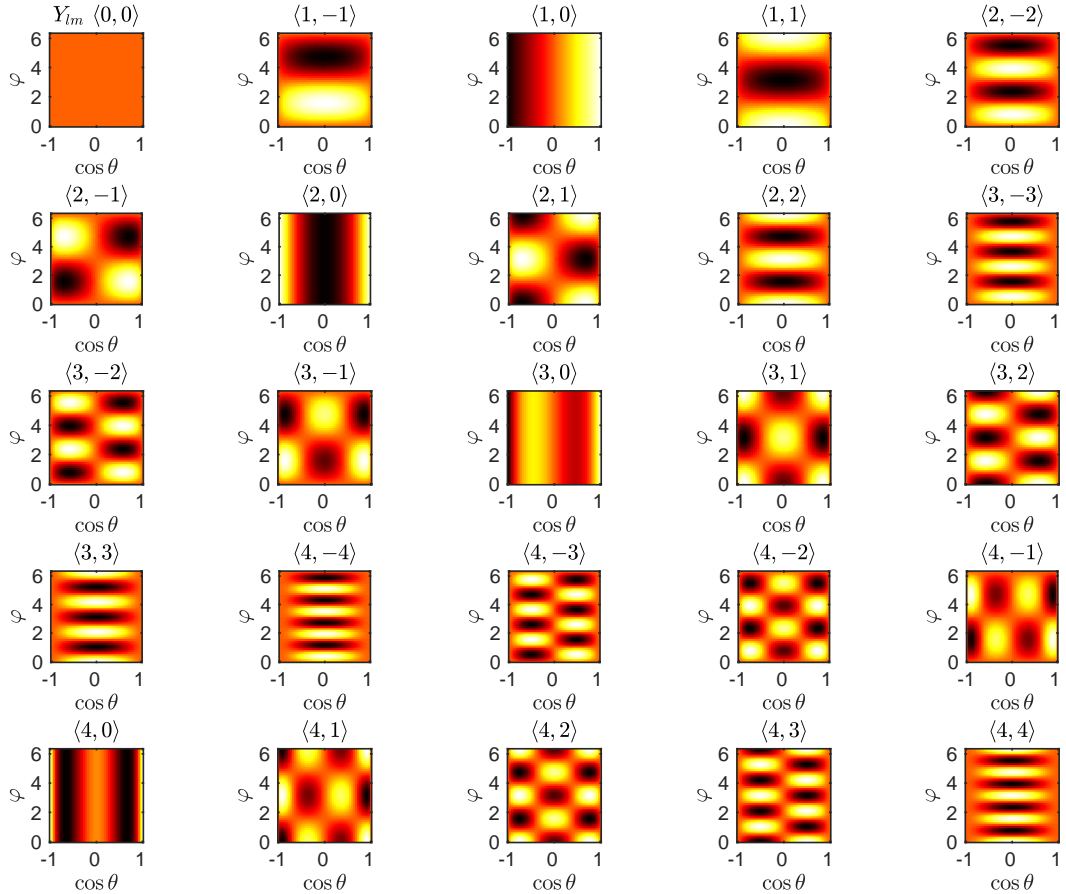


Figure 5.16: Real spherical harmonics Y_{lm} basis functions up to $l_{\max} = 4$.

which obey the basic symmetry relation

$$Y_l^{m*}(\theta, \varphi) = (-1)^m Y_l^{-m}(\theta, \varphi). \quad (5.171)$$

The harmonic functions are calculated using standard numerical methods up to any lm -value using the quantum mechanics normalization conventions.

The expansion coefficients (moments) t_{lm} we are interested in are defined by the inner product integral

$$t_{lm} \equiv \int d\Omega \mathcal{I}(\Omega) Y_{lm}(\Omega) \quad (5.172)$$

with normalization

$$t_{00} = \int d\Omega \mathcal{I}(\Omega), \quad (5.173)$$

returning the number of events, typically. The inner product works because our basis functions are orthonormal.

Parity inversion $(\theta, \varphi) \mapsto (\pi - \theta, \pi + \varphi)$ gives the relation

$$Y_{lm}(-\mathbf{r}) = (-1)^l Y_{lm}(\mathbf{r}), \quad (5.174)$$

where \mathbf{r} is a unit vector and $(-1)^l \equiv$ parity. With parity conservation in the process and the chosen rest frame, only *even* values of l are needed, because then the values of t_{lm} will integrate to zero for odd l . Now if the processes are also rotation symmetric with respect to φ , we see that it is enough to write the expansion in Eq. 5.168 only over non-negative m . This is seen also in Figure 5.16, where for negative m the basis functions are odd under translations of φ over $0 \mapsto 2\pi$, thus giving a vanishing t_{lm} integral for φ -symmetric distributions. However, if the detector response induces major asymmetries which should be also included in the detector simulation, one may need to use the full expansion.

Inverse expansion

The crucial acceptance \times efficiency function moments are

$$\mathcal{E}_{lm} = \int d\Omega A(\Omega) Y_{lm}, \quad (5.175)$$

where the acceptance \times efficiency function $A(\Omega)$ is known only indirectly through detector simulated samples. In its simplest form, it corresponds to geometric fiducial cuts of the detector. The rest of the detector efficiency effects are then variations on this manifold.

The acceptance \times efficiency mixing matrix is

$$\mathcal{E}_{lm,l'm'} = \int d\Omega Y_{lm}(\Omega) Y_{l'm'}(\Omega) A(\Omega), \quad (5.176)$$

which describes how the experimental acceptance cuts will ‘mix’ different spherical moments. In the limit $A(\Omega) \rightarrow 1$ we get $\mathcal{E}_{lm,l'm'} = \delta_{ll'} \delta_{mm'}$, by orthonormality. We calculate the Singular Value Decomposition (SVD) based condition number of the matrix \mathcal{E} as

$$\kappa(\mathcal{E}) = \frac{\sigma_{\max}(\mathcal{E})}{\sigma_{\min}(\mathcal{E})}, \quad (5.177)$$

which we use to characterize the geometric ill-posedness of the problem characterizing the limited detector geometric acceptance, but also the simulation Monte Carlo sample size sufficiency given the chosen expansion truncation l_{max} . The maximum and minimum singular values of the matrix are denoted with $\sigma_{max}, \sigma_{min}$. The identity matrix has a condition number 1, whereas an ill-posed problem has a very large value of κ .

The *observed* moments of data are

$$t_{lm}^{obs} = \frac{\int d\Omega Y_{lm}(\Omega) \mathcal{I}(\Omega) A(\Omega)}{\int d\Omega \mathcal{I}(\Omega) A(\Omega)}, \quad (5.178)$$

which are calculated through a finite sum over the hyperbin sample. One needs to pay attention to normalization, that is, we use the standard quantum mechanics normalization of spherical harmonics $\int |Y_l^m|^2 d\Omega = 1$ which conserves probability. To conserve number of events, we multiply each finite sum with $\sqrt{4\pi}$.

The inverse estimate in the flat phase space is given by the direct algebraic inverse

$$\hat{t}_{lm}^{flat} = \sum_{l'm'} [\mathcal{E}_{lm,l'm'}]^{-1} t_{lm}^{obs}, \quad (5.179)$$

which we take through SVD with possible regularization. That is, one does not necessarily need to do Maximum Likelihood optimization. Note here that once we have \hat{t}_{lm} in the flat phase space, we can *push forward* it to the fiducial phase space using the fiducial acceptance map \mathcal{F} which is calculated analogously to the acceptance \times efficiency matrix, leading to

$$\hat{t}_{lm}^{fid} = \sum_{l'm'} \mathcal{F}_{lm,l'm'} \hat{t}_{lm}^{flat}. \quad (5.180)$$

With low event count statistics, usually the most optimal approach to find \hat{t}_{lm}^{fid} is a $d\Omega$ -unbinned extended Maximum Likelihood formulation with Poissonian event fluctuations in the hyperbin. The likelihood functional is

$$\mathcal{L} = \frac{\langle n \rangle^n}{n!} e^{-\langle n \rangle} \prod_{i=1}^n \frac{\mathcal{I}(\Omega_i)}{\int \mathcal{I}(\Omega) A(\Omega) d\Omega}, \quad (5.181)$$

$$\text{where } \langle n \rangle \equiv \int \mathcal{I} d\Omega (\Omega) A(\Omega) = \sum_{lm} \mathcal{E}_{lm} \hat{t}_{lm}^{fid}, \quad (5.182)$$

where the expected number of *observed* events is $\langle n \rangle$, thus the denominator and Poisson term cancel partially. Now taking the negative logarithm for the minimization

gives

$$-\ln \mathcal{L} = -\sum_{i=1}^n \ln \sum_{lm} \hat{t}_{lm}^{\text{flat}} Y_{lm}(\Omega_i) + \sum_{lm} \mathcal{E}_{lm} \hat{t}_{lm}^{\text{flat}}, \quad (5.183)$$

where we dropped the constant terms depending only on n . This non-convex optimization problem is minimized numerically via MINUIT routines, more specifically by Davidon-Fletcher-Powell quasi-Newton algorithm (MIGRAD), initial estimate given by the algebraic inverse. Once the set of moments $\{\hat{t}_{lm}^{\text{flat}}\}$ is extracted, one may construct other observables based on this set, such as density matrices. We point out here that for *each hyperbin independently*, one repeats the whole chain of calculations including the detector expansion matrices. Technical extensions of this could include interpolation between hyperbins, to suppress statistical fluctuations of the simulation and data samples.

We demonstrate the spherical harmonics expansion in Figure 5.17, 5.18, 5.19, 5.20, 5.21 and 5.22 with cuts suitable for the ATLAS and CMS experiments, where we have neglected the forward proton cuts, thus also the photoproduced ρ^0 is well visible. We emulated the detector p_t -efficiency transfer function with a smooth hyperbolic tangent function for illustration and used flat η -efficiency within the acceptance cube. The corresponding acceptance decompositions are shown in Figures B.1, B.2 and B.3 in Appendix. The somewhat peculiar acceptance bowl at low masses is due to the interplay between the peripheral kinematics and phase space definitions.

The spectra are normalized to one, but naturally, one may use event counts or cross sections. The figures contain results in both the fiducial phase space and in the inverted flat phase space. The pure $J = 0$ process verifies the correctness of the inversion algorithm giving no visible moments other than $lm = \langle 0, 0 \rangle$ in the flat phase space, as should be the case. All other frames than the CM frame have non-zero coefficients with $m \neq 0$. The case $m = 0$ reduces always to the ordinary Legendre polynomials

$$Y_{l0}(\theta, \varphi) = \sqrt{\frac{2l+1}{4\pi}} P_l(\cos \theta), \quad (5.184)$$

with no φ dependence. However, one needs to remember that also in this case the flat phase space inversion machinery is crucial for easy interpretations, otherwise we could just use directly the ordinary Legendre polynomials without any algorithmic machinery.

The reference $J = 0$ process used for the acceptance expansion is on purpose here taken, for the realism, with a slightly different effective b -slope than the full spectrum process. The result is that we obtain slightly varying inversion results in

CHAPTER 5. GRANIITI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

different frames which is visible especially at low masses. This is just the Monte Carlo model dependence which is propagated differently in different frames. This dependence is minimized by using a b -slope value closely matching the data. If event statistics allow, binning over the system transverse momentum solves the problem.

CHAPTER 5. GRANIETTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

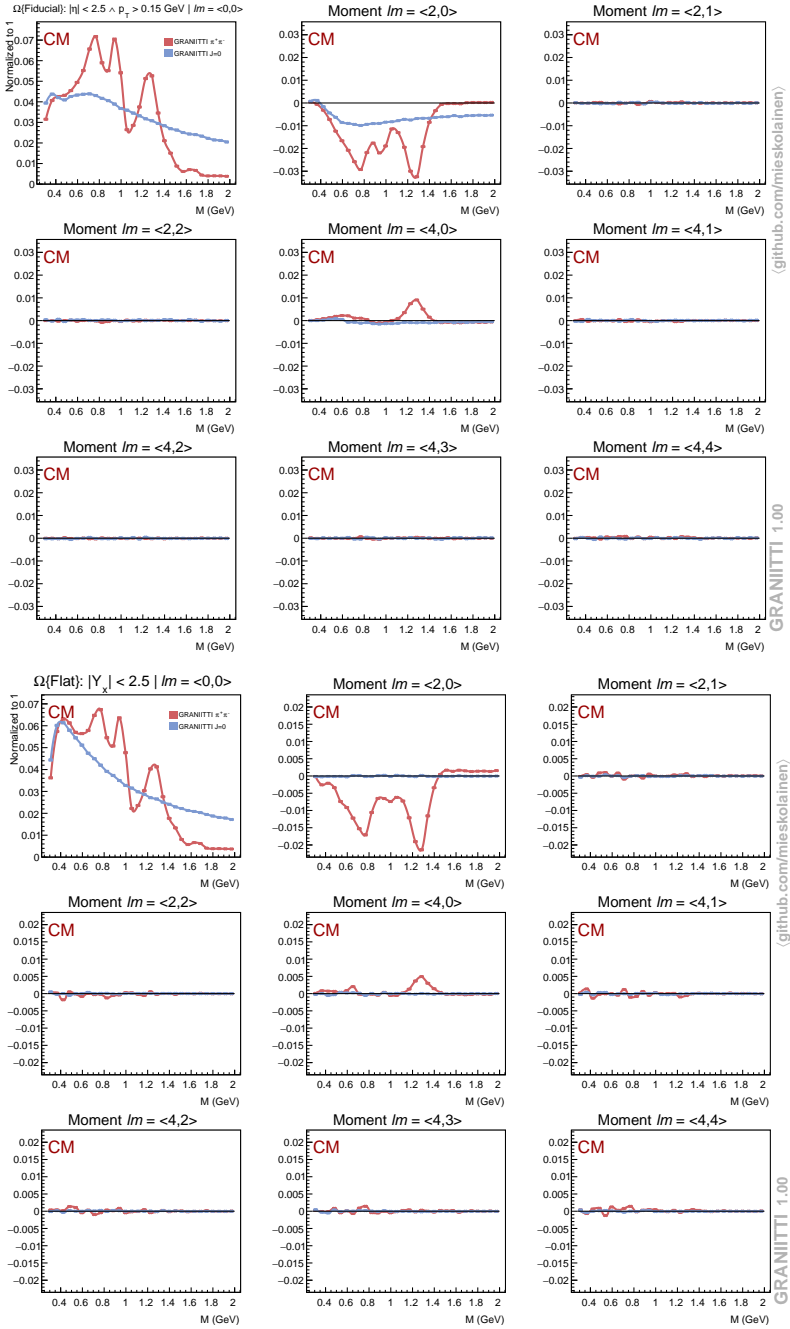


Figure 5.17: CM frame: Harmonic moments in the fiducial phase space (rows 1-3) and in the flat phase space (rows 4-6).

CHAPTER 5. GRANITTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

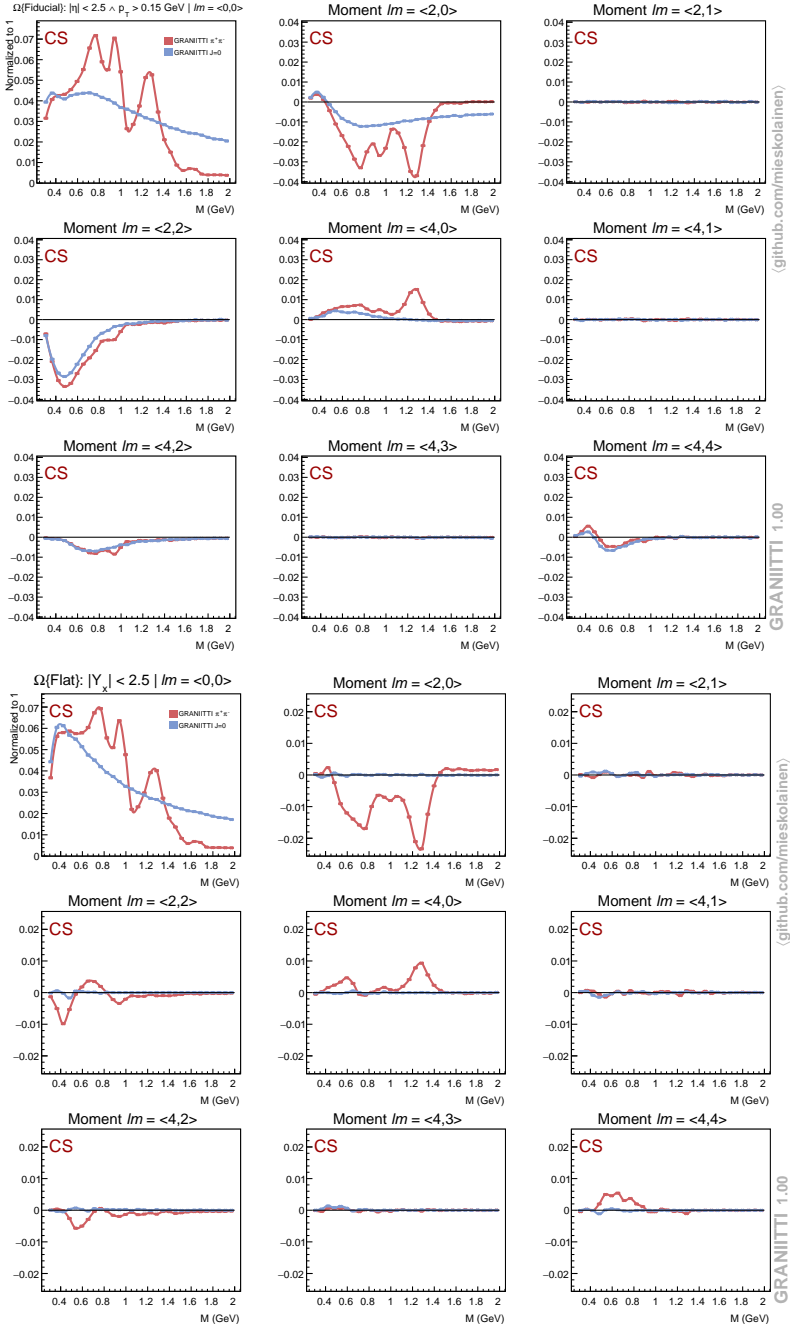


Figure 5.18: CS frame: Harmonic moments in the fiducial phase space (rows 1-3) and in the flat phase space (rows 4-6).

CHAPTER 5. GRANITTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

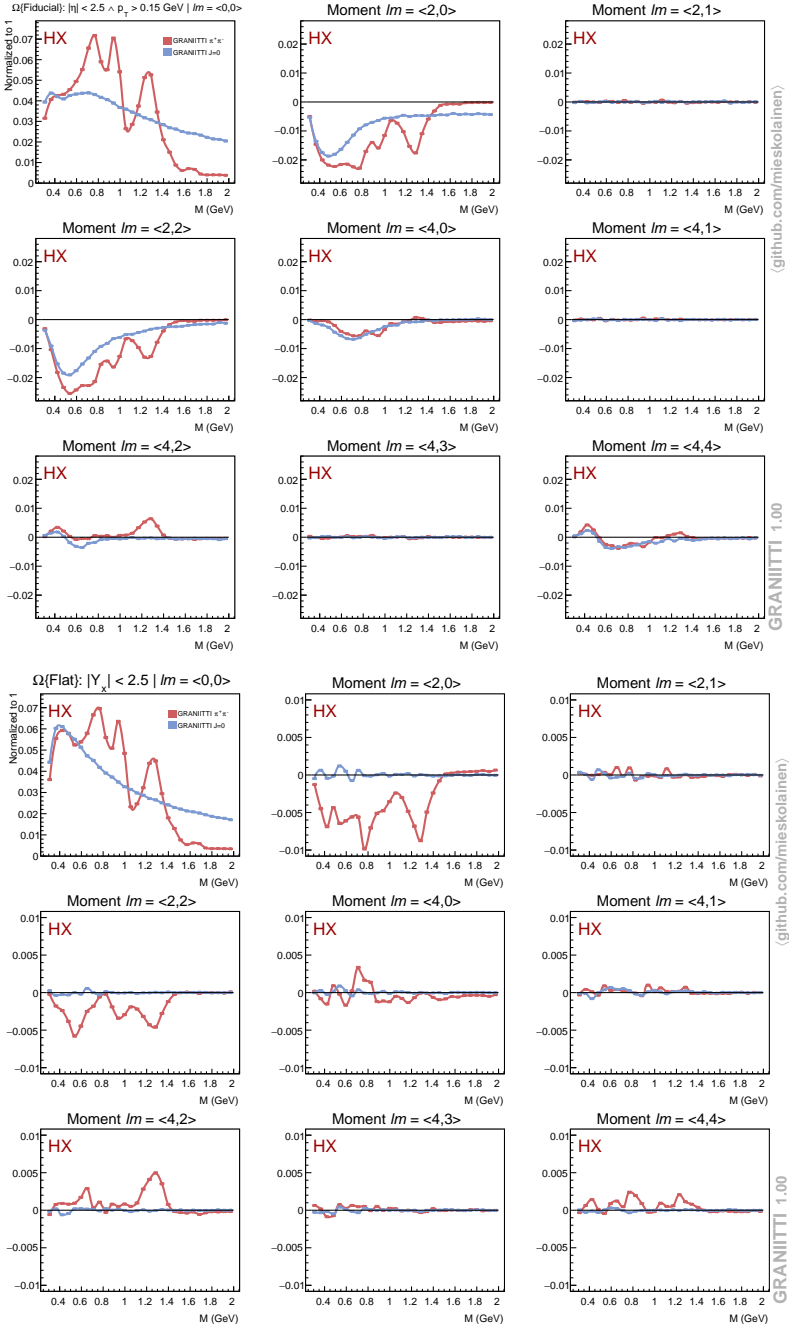


Figure 5.19: HX frame: Harmonic moments in the fiducial phase space (rows 1-3) and in the flat phase space (rows 4-6).

CHAPTER 5. GRANIETTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

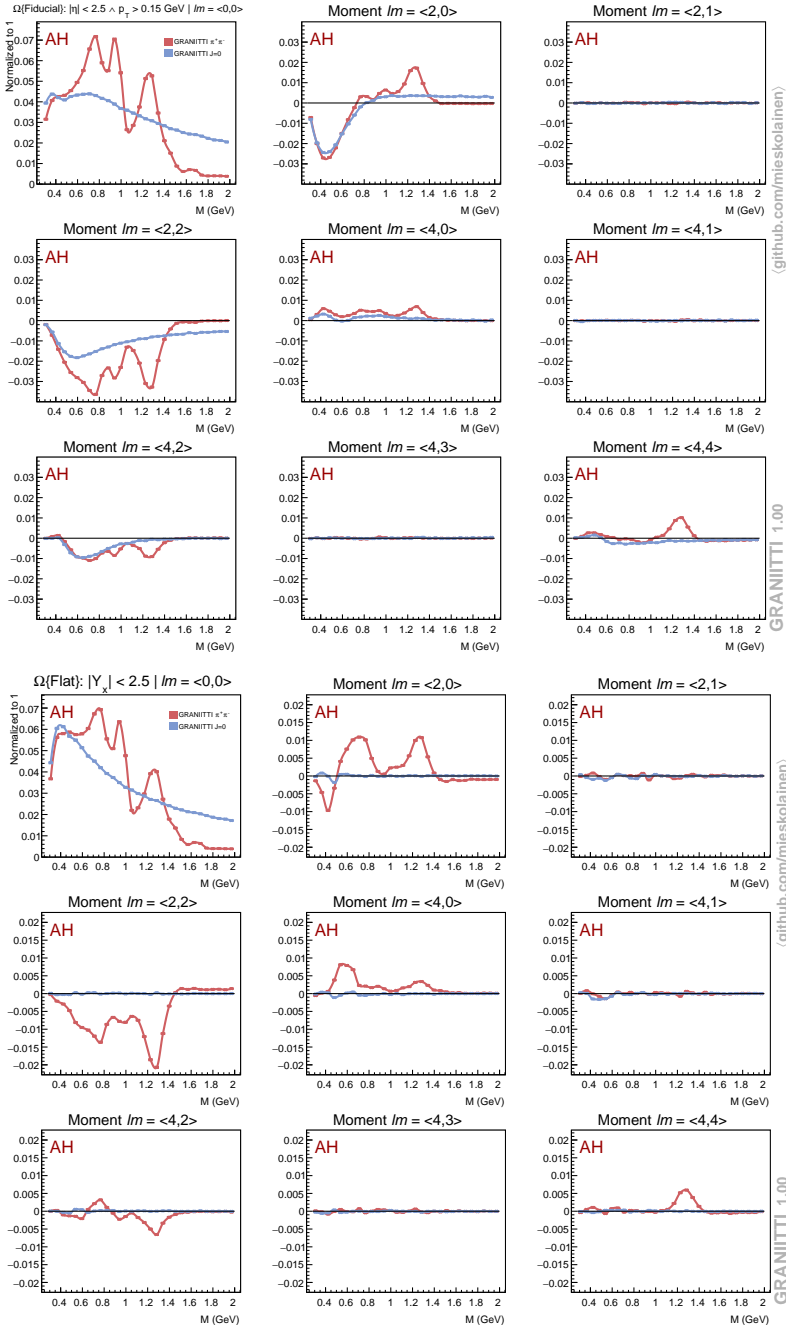


Figure 5.20: AH frame: Harmonic moments in the fiducial phase space (rows 1-3) and in the flat phase space (rows 4-6).

CHAPTER 5. GRANITTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

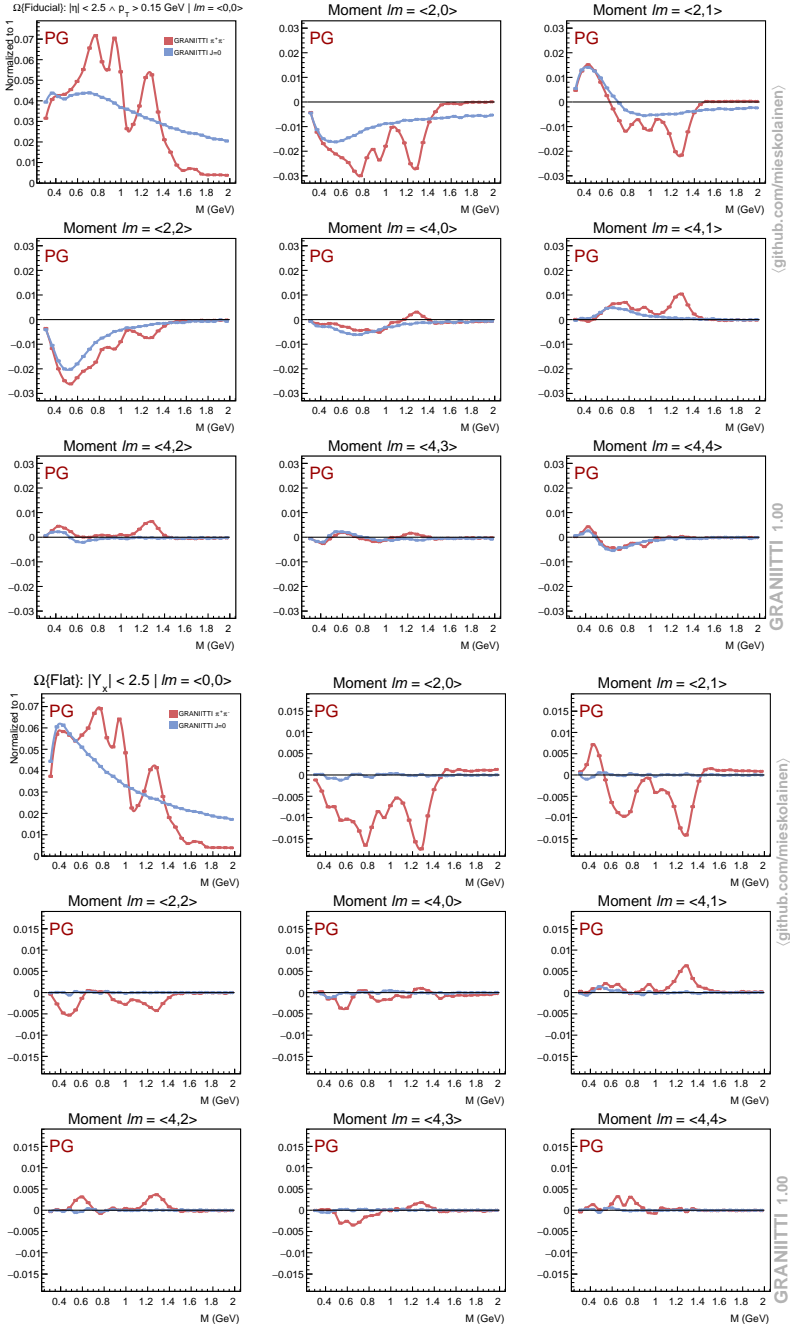


Figure 5.21: PG frame: Harmonic moments in the fiducial phase space (rows 1-3) and in the flat phase space (rows 4-6).

CHAPTER 5. GRANITTI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

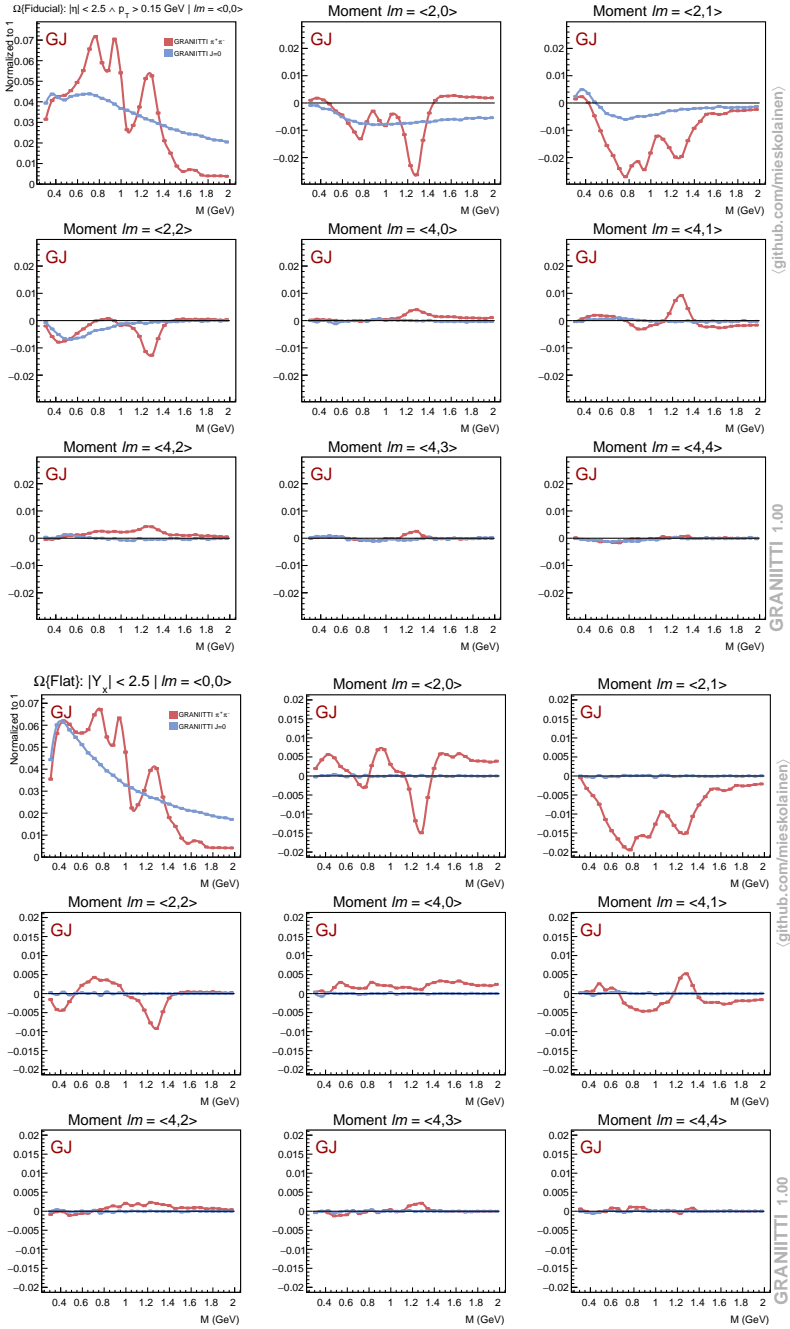


Figure 5.22: GJ frame: Harmonic moments in the fiducial phase space (rows 1-3) and in the flat phase space (rows 4-6).

5.5 Technology

The event generator is implemented in modern C++, utilizing features from the C++17 language standard with full parallel processing using multithreading, basically with no limit in the number of threads. As an example of the current state-of-the-art hardware, Intel Xeon Platinum CPU provides hardware support up to 112 threads with 56 cores. The default random number generator in use is 48-bit RANLUX [147] with easy command line seeding for the distributed grid computing use. The inline commented and physics literature referenced codebase is currently $\mathcal{O}(35\text{k})$ lines.

The generator steering is done using JSON5 style input control cards and a command line interface. To ease out using the software correctly, we check the input and use extensively exception throwing for failure situations. Free parameters of the models are easily changed by modifying the system JSON cards which are grouped under the same folder which allows creating easily different ‘tunes’. Event output is provided in HepMC3 (default), HepMC2 and HEPEVT formats and a converter to LHE format is provided. The ROOT 6 library is used for the analysis algorithms, fitting and plotting, but the generator side is a standalone code and should compile on any modern Linux platform. Compiling the source code is fully automated with standard MAKE tools and a GCC7+ or Clang5+ compiler is needed.

The code is constructed with modern quality standards, driven by an extensive set of fully automated and semi-automated custom test cases and Catch2 library, a software development aspect which we have found highly crucial for reducing the likelihood of the code ‘regression’ and other types of problems.

5.6 Discussion and conclusions

We have seen that GRANIITTI provides a new computational edge on the topic of central diffraction, accelerating the progress on the road towards understanding the topic from the first principles, whatever they are in the future. This concrete code, hopefully, also demystifies several aspects of the expert literature. We simulated how the glueball filter observable is being driven by the spin polarization components of the resonance spin density matrix, thus providing an ansatz that $J = 2$ glueballs produced in central production could be produced with different polarization compared with reasonably established tensor mesons such as $f_2(1270)$ produced most probably in pure transverse $|J_z| = 2$ polarization, which is our hypothesis – of course a mixture of pure states is possible, also with off-diagonal density matrix elements.

For the higher multiplicity final states, we point out here that the ‘parallel processes’ of simultaneous multiple pomeron-pomeron fusion are interesting. In principle, one could assume these to be distributed according to a Poisson distribution in the first approximation. We see that the ‘serial process’ of peripheral ladder exchange, which we have implemented, generates a cross section orders of magnitude lower than what is measured for the $\pi^+\pi^-\pi^+\pi^-$ in the preliminary ATLAS measurement [94] or seen in ALICE data [148]. Within the ALICE fiducial phase space at $\sqrt{s} = 7$ TeV, the cross section for the two body, four body and six body central states with double rapidity gap veto scales approximately $\propto 3^{-N}$, with $N = 2, 4, 6, \dots$. The peripheral ladder exchange simulation has more like $\propto 10^{-N}$ scaling in this phase space, unless the ladder has some unknown non-perturbative mechanism enhancing the couplings or modifying the form factors within ladder vertices. There is always the possibility that the production of intermediate enigmatic $f_0(500)$ mesons and their sequential decays, which is readily available within the generator, would provide the explanation. Discriminating experimentally between the parallel, serial and sequential production is non-trivial, but possible based on the final state kinematics and higher dimensional statistical techniques. We emphasize the need for rigorous (multidimensional) fiducial measurements: cut and count and efficiency correct with possible generalized unfolding. Model-based interpretations or fits of data are always of limited use for the theory development, such will be any attempts to remove any physics background in the soft domain. Such subtractions can be done as a bonus exercise.

The tensor pomeron model predicts a distinctive ‘dip’ in the rapidity separation of a central proton-antiproton production, due to the fermion spin-1/2 structure. In essence, \hat{t} - and \hat{u} -sub-amplitudes have a negative relative sign in this case. Unfortunately, we see that this dip can be destroyed to become unobservable by the screening

loop but also by certain modifications of the unknown internal form factors. The destructive effect of the screening loop can be readily simulated with GRANIITTI. In general, it will be interesting to understand what is the detailed mechanism behind the baryon pair production in central production, what are the effective degrees of freedom and whether models such as Lund strings or other QCD motivated pictures can provide better descriptions.

Regarding spin, GRANIITTI is currently the only event generator that can generate arbitrary spin-dependent scattering amplitudes in low-mass central production. For the spin analyses, our engine puts the fiducial acceptance inversion on a rigorous footing. We demonstrated the ‘implicit complexity’ of the Lorentz rest frames measured by non-zero coefficients of the acceptance decomposition. This is shown in Figures B.1, B.2 and B.3. Unsurprisingly, the most simple one by this first-order measure is the direct Center of Momentum (CM) rest frame without any rotation from the lab and the most complex one seems to be the Pseudo-Gottfried-Jackson frame (PG). We recommend to implement and publish the experimental analysis in several different frames, which makes the measurements more future proof. The full dataset together with detector simulations should be made eventually available at portals such as opendata.cern.ch, because this allows re-studying the higher dimensional kinematic correlations and arbitrary observables. After all, even the spherical harmonic decomposition is only a projection to a certain basis and frame. We probed also dualities between central and forward observables. This is especially interesting for the ALICE and LHCb cases which do not measure forward protons.

We shall here shortly mention the related open inverse problems. What is the optimal tensor basis, not necessarily the spherical tensor basis, which would allow reconstructing the density matrix ρ given the measured set of moments t_{lm} ? The formulation of this problem was proposed originally by Pauli for spin-1/2, which has a particularly simple solution in terms of Pauli matrices. A related question is what is the ‘best way’ to solve amplitudes in the partial wave basis from the measured spherical moments, a problem which is known to have polynomial ambiguities in numerous solutions [135]. Finding out a single physical solution is ill-posed for the case of overlapping resonances and continuum processes, which sum coherently at the amplitude level. The partial wave problem may be, however, actually more tractable for the full $2 \rightarrow 4$ process. To characterize the intrinsic invertibility of the detector acceptance through spherical harmonic expansion, we proposed calculating the condition number of the detector acceptance moment mixing matrix through singular value decomposition. The measured and inverted moments in multiple Lorentz frames, as we showed, give directly useful information for the theory development, independent of the underlying models. Also, GRANIITTI is fully equipped to fit the

CHAPTER 5. GRANIITI: A MONTE CARLO EVENT GENERATOR FOR HIGH ENERGY DIFFRACTION

parameters of the Tensor pomeron model against spherical moment expansions or any other observables – this requires only some CPU time. Finally, it would be interesting to construct maximally model independent and fully covariant spin measures from the frame dependent moment set.

In the future, one could include more scattering amplitudes, lattice simulations fused together with generative machine learning techniques for event-by-event proton structure fluctuations which would be a truly novel case, higher dimensional spin analysis algorithms and deep learning driven ultra efficient Monte Carlo importance sampling.

6 Combinatorial Superstatistics for Soft QCD

High energy diffraction and soft QCD span exciting final state topologies and fluctuations which have not yet been measured or characterized in a fully exhaustive way. In this work, we go beyond the standard measures and formulate a new framework to generalize rapidity gap counting with an emphasis on abstract structure, fiducial observables, experimental limitations and factorizing model dependence. The construction is based on a higher dimensional Bernoulli vector statistics with a combinatorial incidence algebra structure, independent of the underlying field theory. This is the first experimentally feasible framework designed to probe the seemingly arcane sign alternating scattering amplitude cutting rules of AGK-style, which are of interest in Regge theory, perturbative QCD and even in stringy black hole calculus. As an additional novel show case, we pose, construct and solve a highly related combinatorial stochastic superposition Poisson inverse problem using the Möbius inversion theorem.

Chapter in: [arXiv:1910.06279](https://arxiv.org/abs/1910.06279) [hep-ph]

6.1 Introduction

The physics signatures used in searching and analysing interactions at high energy colliders are often based on specific final state ‘topologies’. Examples of these span from the soft diffractive QCD processes, where different event classes are identified by detector responses in pseudorapidity space to new physics searches spanning complex event topologies with jets, leptons and missing transverse energy. In this work, we represent the final state topologies using binary vector spaces. This allows us to use powerful mathematical tools of combinatorics, namely, the incidence algebras introduced by G.C. Rota [149].

The combinatorial problems in high energy physics are numerous. To name a few, they span from combinatorial symmetry factors and Feynman diagram enumeration in QFT, Wick’s theorem (Isserlis’ theorem), Gribov-Glauber cancellations in heavy ion physics, Abramovski-Gribov-Kancheli cutting rules in Regge and field theory, numerous experimental combinatorial puzzles in track fitting and primary vertex association and in the resonance search analysis of final states with combinatorial background. Here we take a slightly different angle on this topic than what has been done previously. From the experimental point of view, the framework here can be utilized for the extraction of diffractive cross sections and completely novel fiducial measurements, and can be considered as a complete generalization of the double diffractive measurement done by the TOTEM experiment [38]. The approach developed in this work is the first explicit utilization of the Möbius inversion theorem in high energy physics, as far as we know. Perhaps the most well known Möbius inversion paper in physics is by Chen [150]. Kowalski, Polyzou and Redish used these tools in the context of partition combinatorics of a non-relativistic multiparticle scattering operator formalism [151]. The connection to supersymmetric mathematical physics models was first developed in papers by Spector [152] and Julia [153], and later by Onofri, Veneziano and Wosiek, which they called ‘supercombinatorics’ [154].

In soft QCD, different minimum bias processes have very different final state configurations over rapidity. However, when the luminosity grows high, nearly all capabilities to study diffractive processes are lost due to the multiple proton-proton interactions per bunch crossing – the pileup. This can be both on-time and off-time pileup, that is, it can propagate from the same or previous bunch cross due to the short 25 ns spacing at the LHC and long detector integration time windows. As an illustrative example, when the average number of inelastic interactions is $\mu \sim 10$, the so-called large rapidity gap events vanish almost completely when we talk about low- p_T events, and we are left only with uniform event signatures as our observables over pseudorapidity. Thus, there are naturally some fundamental limitations in how

large the luminosity can be even without considering all instrumentation and detector level complications. The optimal running setup depends heavily on the physics signal and background. The pile-up rates, being an order of 30-50 in ATLAS and CMS, even higher in the future, pose a major challenge for physics analysis, cross section measurements and searches for new physics phenomena. However, some high- p_T channels are still experimentally manageable up to $\mu \sim \mathcal{O}(10^2)$ or more. Thus we point out that one practical application of our work is the inversion of (minimum bias) trigger combinations under pileup, which was treated in work by ATLAS [155] and ALICE [156] collaborations with less powerful tools. In our other paper [157], we solve a related inverse problem of multiplicity distributions undergoing a compound Poisson autoconvolution process, by fusing techniques of characteristic functions and algorithms.

Section 6.2 starts with multivariate Bernoulli observables together with vector spaces over the finite Galois field $\text{GF}(2)$ and sets them in the context of high energy diffraction. In Section 6.3 we go through the partially ordered sets, incidence algebras and the Möbius inversion and Section 6.4 is devoted for concrete measurements of observables. The other half of the paper deals with the combinatorial superposition problem in a compound Poisson scenario. Section 6.5 constructs the compound Poisson problem and a mathematical model is built, which is then solved using the combinatorial techniques from the first half of the paper. Finally in Section 6.6, we study the inverse problem numerically and summarize in Section 6.7. The appendix provides supplementary information.

6.2 Binary vector spaces and Diffraction

For the algebraic construction of Bernoulli observables, we need a vector space \mathbb{F}_2^N over the finite Galois field $\text{GF}(2)$ with addition as a component wise Boolean OR (\vee) and multiplication as a component wise Boolean AND (\wedge). As the inner product we can use $\langle \mathbf{a} | \mathbf{b} \rangle = \bigvee_{i=1}^N a_i \wedge b_i \in \mathbb{F}_2$, where $\mathbf{a}, \mathbf{b} \in \mathbb{F}_2^N$. The vector space is spanned by N standard basis vectors as $\text{span}(\mathbb{F}_2^N) = \{\mathbf{e}_k\}_{k=1}^N$ and the number of elements in an N -dimensional binary vector space is $|\mathbb{F}_2^N| = 2^N$, illustrated in Figure 6.2. The total number of subspaces is given by a sum over q -binomial coefficients, given in Appendix C.1.

The most useful property concerning the physics involved here is that the subspaces form an orthomodular poset (partially ordered set). Another aspect is that Boolean operators which are diagonal in one basis are diagonal in all bases. Projection operators are always diagonal, which is semi-intuitive if one thinks about the N -dimensional unit-hypercube structure. In [158] it was pointed out that the multivariate Bernoulli can be considered as a certain generalization of the Ising model. However, here we apply these to high energy physics.

Bernoulli random variables

Every component observable of vectors in \mathbb{F}_2^N is treated as a Bernoulli random variable B_i with $1 \leq i \leq N$. The random variables B_i follow the distribution

$$\text{Ber}(B_i | m_i) = P_B(B_i = b | m_i) = m_i^b (1 - m_i)^{1-b}, \quad (6.1)$$

where $b \in \{0, 1\}$. The mean is m_i and variance $\text{Var}[B_i] = m_i(1 - m_i)$. Together, these can be encapsulated with a multivariate Bernoulli distribution $P(B_1 = b_1, B_2 = b_2, \dots, B_N = b_N | \theta)$. The full distribution is either described directly with $2^N - 1$ parameters (-1 from normalization) associated with every 2^N elements of the binary vector space, in what we may call the *fundamental* or *natural* representation, or with parameters describing the expectation values and multipoint correlations between different Bernoulli variables, what we may call the *correlation* representations. We use the notation and representations constructed by Teugels [159].

Example $N = 3$: A centralized correlation representation is

$$m_i = \langle B_i \rangle, \quad i = 1, 2, 3 \quad (6.2)$$

$$\theta_{ij} = \text{Cov}[B_i, B_j] = \langle (B_i - m_i)(B_j - m_j) \rangle \quad (6.3)$$

$$= \langle B_i B_j \rangle - \langle B_i \rangle \langle B_j \rangle, \quad (i, j) \in \{(1, 2), (1, 3), (2, 3)\} \quad (6.4)$$

$$\theta_{123} = \langle (B_1 - m_1)(B_2 - m_2)(B_3 - m_3) \rangle, \quad (6.5)$$

where the moments or correlation functions are centralized, that is, the mean is subtracted. The structure is only slightly more complicated for larger values of N , and is described in the following chapters. One interesting property is that for the multivariate Bernoulli distribution component variables which are pairwise uncorrelated $\text{Cov}[B_i, B_j] = 0$ are also always independent, that is, the density factorizes $P(B_i, B_j) = P(B_i)P(B_j)$ [158]. This does not hold in general for multivariate probability distributions.

Probabilities and cross sections

In the fundamental representation, a probability vector \mathbf{p} encapsulates the probabilities of *mutually exclusive* combinations of observables in the binary vector space

$$\mathbf{p} = (P_{\mathbf{b}_0}, P_{\mathbf{b}_1}, \dots, P_{\mathbf{b}_{2^N-1}})^T, \quad \sum_c p_c = 1 \quad (6.6)$$

where indices go through all the elements of the binary vector space. The one-to-one correspondence is given uniquely by the binary expansion¹

$$c = \sum_{i=1}^N b_i 2^{i-1} \Leftrightarrow \mathbf{b}^c = (b_1, b_2, \dots, b_N), \quad (6.7)$$

between the index $0 \leq c \leq 2^N - 1$ and the binary vector $\mathbf{b}^c \in \mathbb{F}_2^N$.

Now turning into physical observables, the differential cross section element is given by the usual

$$d\sigma_{n'} = \frac{1}{F} \frac{1}{S_{n'}} d\Pi_{n'} |\mathcal{M}_{i \rightarrow n'}|^2, \quad (6.8)$$

where the matrix element squared $|\mathcal{M}_{i \rightarrow n'}|^2$ is the probability amplitude for the transition from the initial state to the n' -particle final state, encapsulating all dynamics. The $S_{n'}$ is a QFT combinatorial symmetry factor for identical final states, $F = 4\sqrt{(p_1 \cdot p_2)^2 - m_1^2 m_2^2} \rightarrow 2s$ (high energy limit) is the incoming invariant Møller flux and the standard Lorentz invariant phase space element is

$$d\Pi_{n'} = (2\pi)^4 \prod_{j=1}^{n'} \frac{d^3 k_j}{(2\pi)^3 2k_j^0}. \quad (6.9)$$

¹Here we remark on the difference between left and right ordered bit strings. We use the convention of most significant ‘bit’ (MSB) on *right*. In the most significant ‘bit’ on *left* representation, this expansion would be $c = \sum_{i=1}^N b_i 2^{d-i}$.

We now divide the fiducial phase space seen by the detector into N elements. Using the combinatorial incidence algebra notation, then the so-called 2^N partial cross sections can be expressed with

$$\begin{aligned} \sigma = \frac{1}{F} \sum_{n'} \frac{1}{S_{n'}} \int d\Pi_{n'} \delta^{(4)} \left(p_1 + p_2 - \sum_{n'} k_j \right) |\mathcal{M}_{2 \rightarrow n'}|^2 \begin{pmatrix} 1 & -1 \\ 0 & 1 \end{pmatrix}^{\otimes N} \\ \left(A_{\{\{k_j\}; \Xi_N\}} \right) \otimes \\ \left(A_{\{\{k_j\}; \Xi_{N-1}\}} \right) \otimes \cdots \otimes \left(A_{\{\{k_j\}; \Xi_1\}} \right). \end{aligned} \quad (6.10)$$

The expression above is a 2^N -vector, where the component σ_0 denotes the null part outside the total fiducial phase space. The acceptance (indicator) functions $A : \{k_j\} \rightarrow \{0, 1\}$ encoding the Bernoulli trials are

$$A_i(\{k_j\}; \Xi_i) = \begin{cases} 1, & \text{if } \exists k_j \in \Xi_i \\ 0, & \text{otherwise.} \end{cases} \quad (6.11)$$

The acceptance function returns 1 if any of the final state particles was inside the phase-space domain Ξ_i defined for B_i and returns 0, if none were inside the domain. This function can be constructed with simple phase space cuts suitable for analytical calculations. In practice, they are most easily calculated with soft QCD Monte Carlo event generators, where arbitrary fiducial phase space cuts are possible. We abstract out the detailed singularity and theory structure connections with the phase space cuts out at this point, they are theory specific, but certainly collinear and infrared singularities may be used to design a specific structure for the acceptance functions.

Experimental efficiency and resolution functions work as a mixing map $\sigma_c \mapsto \sigma_{c'}$, which is a highly complicated stochastic non-linear operator modulating the pure acceptance indicator function. For the expectation values, it can be represented with a finite linear *folding* operator acting on the final states partial cross section vector as $\mathcal{F} : \sigma \rightarrow \tilde{\sigma}$, where σ denotes the partial cross sections at the fiducial phase space particle level and $\tilde{\sigma}$ denotes the visible partial cross sections at the detector level. Both can have almost arbitrary definitions, depending on the problem. The hierarchy goes as follow

$$d\mathcal{O} \xrightarrow[\forall k]{\sum_{n'} \int_{\Omega_{n'}}} \sigma \xrightarrow[-\sigma_E, +\sigma_L]{\mathcal{F}} \tilde{\sigma}, \quad (6.12)$$

where the first map is the ‘theory’ or ‘generator’ and the second one is the ‘detector’, to illustrate. The zero-th component σ_0 of σ denotes the invisible fiducial cross section in the fiducial phase space division (geometrically invisible), only accessible theoretically or via indirect measurements. The inverse operator \mathcal{F}^{-1} is the so-called *unfolding* process, which can be naively solved with a $2^N - 1$ matrix inversion, obtained via detector simulation. The σ_E denotes the integrated efficiency loss, which is not the same as pure acceptance loss σ_0 . Integrated leakage from outside the fiducial to inside the visible is denoted with σ_L , which may happen due to e.g. material re-scattering – this can be a non-negligible effect with very forward detectors. We emphasize here that the unfolding is for efficiency and smearing induced corrections, not to extrapolate the detector geometry.

As an example of the indicator functions one could consider a simple division of the (pseudo)rapidity into N intervals with suitable p_t thresholds per interval. Experimentally, the intervals may overlap over pseudorapidity such that $[\eta_{i,\min}, \eta_{i,\max}] \cup [\eta_{j,\min}, \eta_{j,\max}] \neq \{0\}$ for a detector pair (ij) , and thus making them trivially correlated. It is possible, in principle, via suitable non-overlapping fiducial definition to take this into account in the unfolding procedure which de-correlates the data. Another option is simply to take care that no overlap happens physically. Also, the detectors may have gaps over η such that $[\eta_{i,\min}, \eta_{i,\max}] \cup [\eta_{j,\min}, \eta_{j,\max}] = \{0\}$. However, it is worth pointing out that the formalism presented here does not assume statistical independence of any kind, either implicit or explicit.

Abramovski-Gribov-Kancheli cutting rules

We collect here some illustrating results of the AGK cutting rule calculus [25]. By no means this is a complete description, but rather of an introductory nature. These rules are somewhat similar to more well known Cutkosky-Landau cutting rules based on the S -matrix unitarity, the rules which treat scattering amplitude discontinuities. The AGK rules were introduced in the Regge theory context, but they are actually flexible in terms of the underlying theory. What we need here is the reggeization, which happens in field theories under Regge limit but also in string theory. The AGK calculus results here are collected from a stringy black hole context [160] and from QCD papers [161, 162]. We assume the usual unitary \mathcal{S} -matrix and \mathcal{T} matrix related with

$$\mathcal{S} = 1 + i\mathcal{T}. \tag{6.13}$$

Let us have the elastic scattering amplitude $\mathcal{T}_{AB}(s, t)$ with a Sommerfield-Watson integral transform representation in the complex angular momentum space as

$$\mathcal{T}_{AB}(s, t) = \int \frac{d\omega}{2i} \xi(\omega) s^{1+\omega} \mathcal{F}(\omega, t) \quad \text{with} \quad \xi(\omega) = \frac{\tau - e^{-i\pi\omega}}{\sin \pi\omega}, \quad (6.14)$$

where $\mathcal{F}(\omega, t)$ is the corresponding partial wave, $\omega = J - 1 \in \mathbb{C}$ being the conjugate variable to s together with the signature factor ξ for the positive or negative signature $\tau = \pm 1$. The positive signature case, such as Pomeron exchange, can be written more compactly as $\xi(\omega) = i + \tan(\frac{\pi}{2}\omega)$. Now assume an exchange of n Reggeons, that is, exchange of Regge trajectories.

The n -cut discontinuity over ω of $\mathcal{F}(\omega, t)$ can be written in terms of vertex functions $\mathcal{V}_n^a, \mathcal{V}_n^b$ and integrating over transverse momentum phase space as [162]

$$\text{DISC}_\omega^{(n)}[\mathcal{F}(\omega, t)] = 2\pi i \int \frac{d\Omega_n}{n!} \delta(\omega - \sum_j \beta_j) \Pi(\{\beta_j\}) \mathcal{V}_n^a(\{\mathbf{k}_j, \omega\}) \mathcal{V}_n^b(\{\mathbf{k}_j, \omega\}) \quad (6.15)$$

$$d\Omega_n \equiv (2\pi)^2 \delta^{(2)}\left(\mathbf{q} - \sum_{j=1}^n \mathbf{k}_j\right) \prod_{j=1}^n \frac{d^2\mathbf{k}_j}{(2\pi)^2} \quad (6.16)$$

$$\Pi(\{\beta_j\}) \equiv \text{Im} \left[-i \prod_j i\xi_j \right] = (-1)^{n-1} \frac{\cos \left[\frac{\pi}{2} \sum_j (\beta_j + \frac{1+\tau_j}{2}) \right]}{\prod_j \cos \left[\frac{\pi}{2} (\beta_j + \frac{1-\tau_j}{2}) \right]}, \quad (6.17)$$

where $\beta(-\mathbf{k}^2) \equiv \alpha(-\mathbf{k}^2) - 1$ with $\alpha(t_j = -\mathbf{k}_j^2)$ being the Regge trajectory function, for the soft Pomeron often used parametrization is $\alpha(t) \simeq 1.08 + 0.25t$. The factor $n!$ comes from all possible exchange diagram planar and crossed non-planar orderings.

Now the amplitude for the n -cut exchange can be written using the Sommerfield-Watson transform and the cut discontinuity as [161]

$$\mathcal{T}_{AB}^{n\text{-cut}}(s, t) = \int \frac{d\omega}{2i} \xi(\omega) s^{1+\omega} \text{DISC}_\omega^{(n)}[\mathcal{F}(\omega, t)] \quad (6.18)$$

$$= \pi \int \frac{d\Omega_n}{n!} \xi(\sum_j \beta_j) s^{1+\sum_j \beta_j} \Pi(\{\beta_j\}) \mathcal{V}_n^a(\{\mathbf{k}_j, \sum_i \beta_j\}) \mathcal{V}_n^b(\{\mathbf{k}_j, \sum_j \beta_j\}). \quad (6.19)$$

Next we write down the Abramovski-Gribov-Kancheli sign alternating combinatorial factors [162]

$$G_k^n = \begin{cases} 2^{n-1}(-1)^n + \Pi(\{\beta_j\}), & \text{for } k = 0 \text{ cut Pomerons,} \\ 2^{n-1}(-1)^{n-k} \frac{n!}{(n-k)!k!}, & \text{for } 0 < k \leq n, \end{cases} \quad (6.20)$$

n, k	0	1	2	3	4	5	6
1	0	1	0	0	0	0	0
2	1	-4	2	0	0	0	0
3	-3	12	-12	4	0	0	0
4	7	-32	48	-32	8	0	0
5	-15	80	-160	160	-80	16	0
6	31	-192	480	-640	480	-192	32

Table 6.1: AGK factors G_k^n for the even signature (Pomeron) case with n exchanges and k of them cut.

which we tabulate for illustration in Table 6.1.

Let us now point some integer sequence properties of the AGK series, which might be illuminating in our hyperspace context. The series $|G_{k=1}^n|$ is the number of edges in an n -hypercube. The series $|G_{k=2}^n|$ is the number of diagonals in an n -hypercube of length $\sqrt{2}$. The series $|G_{k=3}^n|$ is the number of 3-cycles in the halved n -cube graph. These suggest to us directly that in order to be (experimentally) sensitive to n Pomeron exchanges, the number of rapidity slices N should be at least $n + 1$. Now the use of AGK cutting factors is that they allow us to represent the total discontinuity over s -space scattering amplitude as a weighted sum over all k -cut Pomeron contributions [162]

$$\mathcal{A}^n(s, t) \equiv \text{DISC}_s[\mathcal{T}_{AB}^{n-\text{cut}}(s, t)] = \sum_{k=0}^n a_k^n(s, t), \quad \text{where}$$

$$a_k^n = 2\pi i G_k^n \int \frac{d\Omega_n}{n!} s^{1+\sum_i \beta_i} \mathcal{V}_n^a(\{\mathbf{k}_j, \sum_i \beta_i\}) \mathcal{V}_n^b(\{\mathbf{k}_j, \sum_i \beta_i\}). \quad (6.21)$$

The last step in our short discussion on this topic is the impact parameter b space picture and eikonalization. The eikonal opacity function in the impact parameter space is obtained using the (inverse) Fourier transform of a single discontinuity integrated over ω [162]

$$\Omega(s, b) \equiv \frac{1}{\pi i} \int \frac{d^2 \mathbf{q}}{(2\pi)^2} e^{-i\mathbf{b} \cdot \mathbf{q}} \int d\omega \text{DISC}_\omega[\mathcal{F}(\omega, -\mathbf{q}^2)] s^\omega \quad (6.22)$$

with the amplitude recovered in the s -space by (forward) Fourier transform of the n -times exponentiated amplitude, by *assuming factorization* of each exchange as

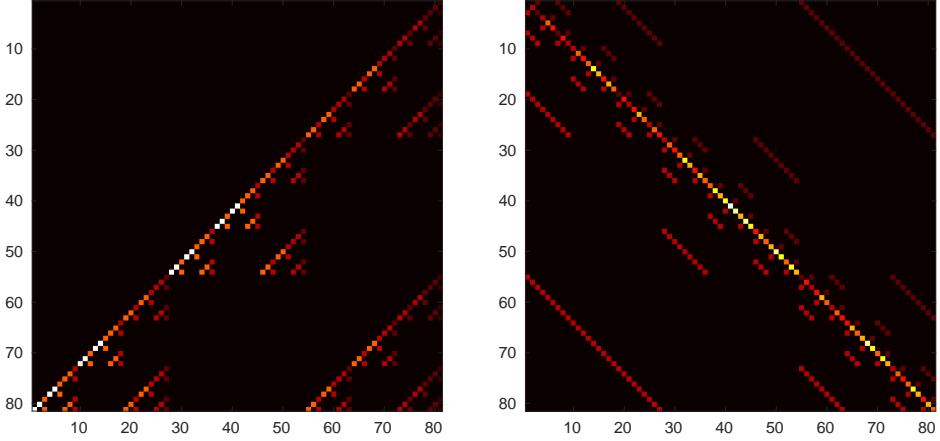


Figure 6.1: A tensor product chain $M^{\otimes 4}$ on the left and a tensor sum chain $\frac{1}{4}M^{\oplus 4}$ on the right. Black to red to white denote values in $[0,1]$.

$\mathcal{V}_n^a = [\mathcal{V}_1]^n$, giving

$$\mathcal{A}_k^n(s, t) = 4is \frac{(-1)^{n-k}}{k!(n-k)!} \int d^2\mathbf{b} e^{i\mathbf{b}\cdot\mathbf{q}} [\Omega(s, b)]^n \quad (6.23)$$

with summation $n \geq k$ resulting in

$$\mathcal{A}_k(s, t) = 4is \int d^2\mathbf{b} e^{i\mathbf{b}\cdot\mathbf{q}} P(k, s, b). \quad (6.24)$$

Where the probability is a sum over the \mathcal{S} -matrix elements [162]

$$P(k, s, b) \equiv \sum_{n \geq k} [\mathcal{S}(s, b) \mathcal{S}(s, b)^\dagger]_k^n = \frac{[\Omega(s, b)]^k}{k!} e^{-\Omega(s, b)}, \quad (6.25)$$

which gives a Poisson distribution with $\langle k(s, b) \rangle = \text{Var}[k(s, b)] = \Omega(s, b)$. That is, the number of k -cut Pomerons at given energy squared s and impact parameter b is Poisson distributed in this picture. Now, in typical Monte Carlo implementations, one needs also the parton densities to pick the longitudinal momentum fractions from vertex factors $\mathcal{V}_1^{a,b}$ for each exchanged Pomeron and at least triple-Pomeron interactions with diffractive cuts, in addition (hard) QCD matrix elements, parton shower and fragmentation.

As an example, we shall refer to the description of QGSJet-II event generator for details [163], which implements advanced Reggeon calculus for soft and hard interactions. The scope of this paper is to provide new algebraic methods for measuring the observables. We use the Poissonian AGK distribution as our motivation for the later construction. We shall also note here that a left stochastic matrix devised by Ryskin and Bartels [161], describes the ‘non-diagonal’ cut transitions necessary for rapidity gap configurations

$$M = \begin{pmatrix} 0 & 0 & \frac{1}{2} \\ 0 & 1 & 0 \\ 1 & 0 & \frac{1}{2} \end{pmatrix}, \quad (6.26)$$

which has $\mathbf{x} = [1, -4, 2]^T$ as its eigenvector, the first non-trivial row from Table 6.1. That is, for every rapidity interval, these ratios hold, and the transition matrix describes the stochastic probability of transition from different configurations per chain. In Figure 6.1 we illustrate the nature of this matrix under the Kronecker product $A \otimes B$ and the Kronecker sum defined as $A \oplus B = A \otimes I_A + I_B \otimes B$, where I is an identity matrix. The product matrix belongs to the class of inclusion-exclusion matrices whereas the sum matrix resembles a hypercube adjacency (connectivity) matrix. Later we learn how these concepts appear in the incidence algebra.

Rapidity gaps as a spacing distribution

It is typical that ‘non-diffractive events’ have an exponentially decreasing maximum rapidity gap distribution $\sim \exp(-\ell_c \Delta y)$, where the correlation length $\ell_c \sim 1$ depends on the average multiplicity density. This can be derived assuming n independent final state emissions over rapidity. Diffractive events are those which have ‘anomalously’ large gaps. The rapidity gap itself is a kinematic consequence given the excited (dissociative) proton system mass is small enough and the initial state boost large enough, however, the M^2 distribution itself is driven dynamically. The standard triple Pomeron limit distributions obtained via generalized multibody unitarity together with Regge asymptotics give $d\sigma/dM^2 \sim 1/(M^2)^{\alpha_P(0)}$, where the $\alpha_P(0)$ is the effective Pomeron intercept $\alpha_P(0) \equiv 1 + \Delta_P \sim 1.08$. Then by standard kinematic² change of a variable $\phi^{-1}(\Delta y) \equiv s \exp(-\Delta y) = M^2$, we obtain the single diffractive

²Beam rapidity is $y_{\pm} = \pm \ln(\sqrt{s}/m_p)$ and the diffractive system spans $y_{\text{span}} = \ln(M^2/m_p^2)$.

rapidity gap distribution

$$\begin{aligned}
 P_D(\Delta y) &\sim \frac{d\sigma}{d\Delta y} = \frac{d\sigma}{dM^2}(\phi^{-1}(\Delta y)) \left| \frac{d\phi^{-1}}{d\Delta y} \right| \\
 &= \frac{1}{(s \exp(-\Delta y))^{\alpha_P(0)}} \left| \frac{-s}{e^{\Delta y}} \right| = e^{\Delta y(\alpha_P(0)-1)}. \quad (6.27)
 \end{aligned}$$

Above, we neglected Mandelstam t -dependence of the triple Regge expressions by taking $t \rightarrow 0$, but this is not relevant for this discussion. Now Equation 6.27 is a slowly rising exponential, instead of decreasing such as for the random non-diffractive processes. In the case of a double diffractive process, the equivalent derivation goes via the Jacobian determinant $\det J(\Delta y, y_0) = sm_p^2 e^{-\Delta y}$, from $\phi_i^{-1}(\Delta y, y_0) = (sm_p^2 e^{-\Delta y \pm 2y_0})^{1/2} = M_i^2$ with $i = 1, 2$ for the two systems and $\Delta_y^{DD} = -\ln(M_1^2 M_2^2 / (sm_p^2))$ and $y_0 = \frac{1}{2} \ln(M_1^2 / M_2^2)$. This gives us

$$\begin{aligned}
 P_{DD}(\Delta y) &\sim \frac{d\sigma}{d\Delta y dy_0} = \frac{d^2\sigma}{dM_1^2 dM_2^2}(\phi_{1,2}^{-1}(\Delta y, y_0)) |\det(J)| \\
 &= \prod_i \left[\frac{1}{(sm_p^2 \exp(-\Delta y \pm 2y_0))^{\alpha_P(0)/2}} \right] \frac{sm_p^2}{e^{\Delta y}} \quad (6.28) \\
 &\sim e^{\Delta y(\alpha_P(0)-1)}, \quad (6.29)
 \end{aligned}$$

where the dependence on y_0 vanishes and the functional dependence is the same as in the single diffractive case. The boundary conditions are different, because the both edges of the gap are moving in the case of double diffractive scattering. In general, the gap may be taken with respect to the detector geometric boundary or with respect to the minimum or maximum of the rapidity interval (type I), or we can count distances between particles (type II). One must remember that any unitarization scheme may modify these results. That is, the triple Pomeron expressions do not satisfy unitarity under $s \rightarrow \infty$, which is a long standing problem, with numerous partially plausible solutions (eikonalization, triple Pomeron decoupling, ...). Currently, there is no detailed data enough to gain new insights on this problem neither the non-perturbative QCD techniques are advanced enough to solve it. However, new LHC data may change the situation.

At this point, we refer the reader to mcplots.cern.ch for the measured rapidity gap distributions at the LHC with different minimum p_t thresholds and Monte Carlo event generator comparisons. The large (diffractive) gaps spanning several units of rapidity originate from charge-neutral exchanges, such as exchange of color neutral gluon systems (pomeron) or photons. The interaction rate of gap events may be

screened by additional elastic or inelastic interactions, this is known as the gap survival discussion. An important property is that measuring the final states down to an arbitrary infrared cutoff, not experimentally accessible, modifies the observed rates of rapidity gap events. This corresponds e.g. to extra emission of soft gluons. Thus, the rapidity gap events are not clearly ‘infrared safe’ without an explicit p_t threshold definitions. Later in this work, in Figure 6.6, we show with a toy Monte Carlo simulation how the combinatorial gap rates run as a function of p_t thresholds for four different particle densities per rapidity slice. We see that it is possible, given different p_t thresholds and particle densities, to run the relative event rates of different final state topologies to wildly different values.

For the non-diffractive events, now let Y_0, Y_1, \dots, Y_n be $n + 1$ rapidity intervals spanned by an n particle final state proceed e.g. via ‘string breaking process’. To calculate statistics of these, we may use the Darling’s contour integral technique from mathematical statistics [164] via the Laplace transform

$$\mathbb{E}[f_0(Y_0)f_1(Y_1)\dots f_n(Y_n)] = \frac{n!}{2\pi i} \int_{c-i\infty}^{c+i\infty} e^z \prod_{j=0}^n \int_0^\infty dz dr_j e^{-r_j z} f_j(r_j). \quad (6.30)$$

The path of the integral is along the straight line $\text{Re}\{z\} = c$ and $f_j(x)$ are arbitrary real valued functions. Using this expectation integral, one can derive numerous results regarding the spacing statistics. For example:

MAXGAP: The probability that all rapidity gaps are smaller than β is

$$P(Y_j < \beta \forall j) = \sum_{0 \leq j < 1/\beta} \binom{n+1}{j} (-1)^j (1 - \beta j)^n. \quad (6.31)$$

MINGAP: The probability that all rapidity gaps are larger than α is

$$P(Y_j > \alpha \forall j) = (1 - (n+1)\alpha)^n \quad \text{s.t.} \quad \alpha < 1/(n+1). \quad (6.32)$$

Both of these are derived in Darling’s paper [164] under the uniform distribution assumptions. We see that given uniformly distributed rapidities, the number of final states gives physical boundaries for the probabilities, visible in the inequalities in Equation 6.31 and 6.32. Now clearly, one can consider higher-order spacing statistics such as distributions of the largest and the second largest gap and so on. Clearly, our incidence algebra construction is an integrating embedding of these taking also into account the translation (boost) position on the rapidity axis.

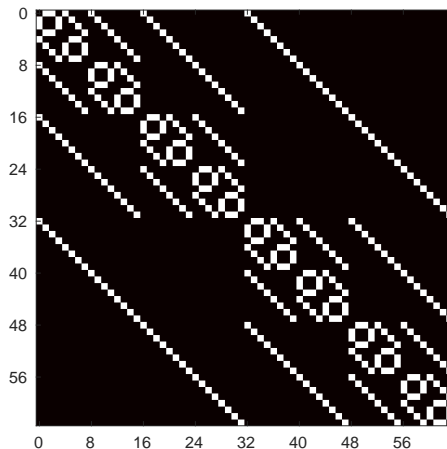


Figure 6.2: The binary vector space illustrated as an $N = 6$ hypercube graph with its corresponding adjacency matrix (black is 0, white is 1).

Our construction includes, in the limit $N \rightarrow \infty$, exactly the ‘traditional’ rapidity gap description. That is, one can recover the one dimensional rapidity gap distribution from the N -dimensional combinatorial distribution. Moreover, in practice, N needs not to be infinite due to the random correlation length over rapidity which smears out useful information at very small gap values.

Hypergraph and topology

In Figure 6.2 we demonstrate the vector space as an equivalent hypercube $G(E, V)$. Each vertex in set E corresponds to one particular final state configuration and partial cross section component. The number of edges in the graph is $N2^{N-1}$, also counting the number of non-zero phase space elements (bits). The total number of subspaces is 2825, for $N = 6$. That is, it is a radically richer space than a single dimensional projection of rapidity gap distributions. In the following sections, a complete algebraic construction for any finite N will be given.

We point out that our definition of ‘event topology’ is reasonably well defined even if slightly abstract at first sight. Regarding the topological invariants of the hypercube, the Euler characteristic gives us $V - E + F = 2 - 2\gamma(G_N)$ where F is the number of faces and the minimum genus of N -cube for $N \geq 2$ is [165]

$$\gamma(G_N) = (N - 4)2^{N-3} + 1. \quad (6.33)$$

That is, if $N \leq 3$, then $\gamma(G_N) = 0$ and we can embed the planar graph without edge crossings on a sphere, $N = 4$ requires a one hole torus and then exponentially more complicated topologies are generated. The connection between final state geometry (kinematics) and topology is now well defined – more complicated ‘event topologies’ require larger N . An interesting construction would be the limit $N \rightarrow \infty$, that is, from finite dimensional hypercubes to an infinite dimensional construction. In practice, approximation schemes such as neural networks could be useful proxies for that.

6.3 Posets, incidence algebra and Möbius inversion

To have the necessary mathematical preliminaries to attack our problems further, we need to first define here a few basic concepts well known in combinatorics. For more information, we refer the reader to the original papers on incidence algebras by G.C. Rota [149] and a textbook on combinatorics by R. Stanley [166].

Let P be a set. A binary relation \leq between elements $x, y, z \in P$ satisfying

1. Reflexivity: $x \leq x$ for all $x \in P$ (6.34)

2. Antisymmetry: if $x \leq y$ and $y \leq x$, then $x = y$ (6.35)

3. Transitivity: if $x \leq y$ and $y \leq z$, then $x \leq z$ (6.36)

is called a *partial order*. Thus, a set equipped with a partial order is called a *poset*.

Now (P, \leq) is a poset. An *incidence algebra* is constructed by first defining an interval $[x, y] = \{z \in P \mid x \leq z \leq y\}$. Then the incidence algebra on P is

$$I(P) = \{f : P \times P \rightarrow \mathbb{R} \mid f(x, y) = 0 \text{ unless } x \leq y\}. \quad (6.37)$$

Thus, it gives us the precise definition of the inclusion operation. The incidence algebra has two operations of addition and multiplication defined as

1. Addition : $(f + g)(x, y) = f(x, y) + g(x, y)$ (6.38)

2. Multiplication : $(f * g)(x, y) = \sum_{x \leq z \leq y} f(x, z)g(z, y)$ (6.39)

so the addition is point-wise and the ‘multiplication’ is a convolution sum. Thus, division can be understood as a deconvolution operation. Common examples of locally finite posets P are: the positive integers ordered by divisibility, an inclusive order of the subspaces of a vector space over a finite field (projective spaces) or all the inclusively ordered subsets of a given set.

In order to proceed, we need a ζ -function³ defined as

$$\zeta(x, y) = \begin{cases} 1, & \text{if } x \leq y. \\ 0, & \text{otherwise.} \end{cases} \quad (6.40)$$

and its inverse ζ^{-1} is the so-called Möbius function μ (as in number theory) defined

³This has a relation with the Riemann ζ -function, which explains the name.

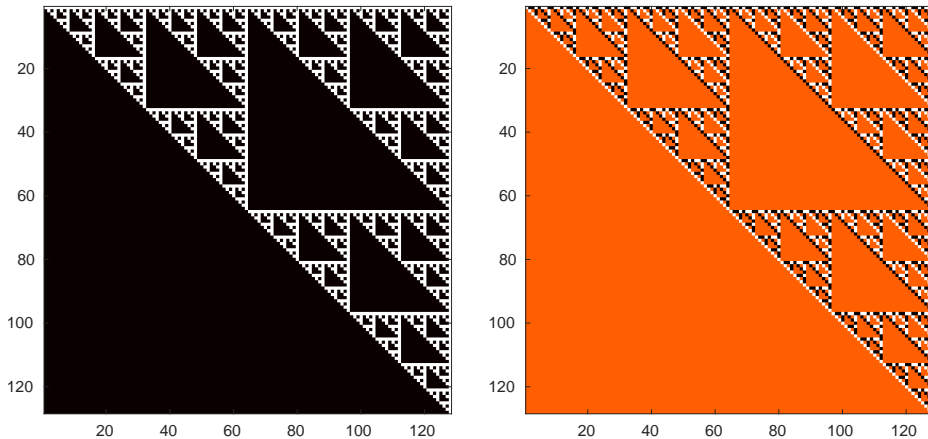


Figure 6.3: Matrices $\zeta_{128 \times 128}$ and $\zeta_{128 \times 128}^{-1}$ for $N = 7$. On the left black is 0 and white is 1. On the right, orange is 0, black is -1 and white is 1.

with a recursive relation

$$u(x, y) = \begin{cases} 1, & \text{if } x = y \\ -\sum_{x \leq z \leq y} u(x, z) & \text{for } x < y \\ 0, & \text{otherwise.} \end{cases} \quad (6.41)$$

The identity or delta-function works as expected in $I(P)$: $\delta(x, y) = 1$ if $x = y$, and 0 otherwise. In what follows, the ζ -function can be represented with an upper triangular matrix of size $2^N \times 2^N$ with Boolean elements, unit diagonal and determinant 1. Thus, the Möbius matrix is found easily by inverting the ζ -matrix which results in an upper triangular matrix with unit diagonal and alternating signed elements.

The Möbius inversion theorem

The *Möbius inversion theorem* goes as follows [166]. Let (P, \leq) be a finite poset and f and g elements of the incidence algebra $I(P)$. The theorem states two equivalent

formulas in direct $g \mapsto f$ and inverse $f \mapsto g$ directions

$$f(x, y) = \sum_{x \leq z \leq y} g(x, z) \quad \Leftrightarrow \quad f = g * \zeta \quad (6.42)$$

$$g(x, y) = \sum_{x \leq z \leq y} f(x, z)u(z, y) \quad \Leftrightarrow \quad g = f * u. \quad (6.43)$$

Intuitively, multiplying by ζ can be understood as integration and by u as a differentiation operation. The same convolution products are easily written as

$$\mathbf{f} = \zeta \mathbf{g} \Leftrightarrow \mathbf{g} = \zeta^{-1} \mathbf{f}, \quad (6.44)$$

using a representation with column vectors \mathbf{f}, \mathbf{g} and triangular square matrices ζ, ζ^{-1} . When increasing the dimension N , the recursive Sierpinski triangle fractal structure in these matrices will manifest itself. This is shown in Figure 6.3.

The principle of inclusion-exclusion

The Principle of Inclusion Exclusion (PIE) is the Möbius inversion for subsets [166]. The basic ‘sieving’ principle of PIE itself has been known for long. Now let different scattering cross section domains for each Bernoulli observable be represented with subsets $D_1, D_2, \dots, D_N \subseteq D$. By defining two functions f and g on the poset (P, \subseteq) having $|P| = 2^N$ and $\mathbb{F}_2^N \cong P$ as

$$f(\chi) = \left| \bigcap_{i \in \chi} D_i \right| \quad (6.45)$$

$$g(\chi) = |\{x \in D \mid x \in D_i \text{ for all } i \in \chi \text{ holds } x \notin D_j \text{ for all } j \notin \chi\}|, \quad (6.46)$$

where $\chi \subseteq P$. This can be written in the forward direction with

$$f(\Upsilon) = \sum_{\chi \subseteq \Upsilon \subseteq P} g(\chi). \quad (6.47)$$

Now the inverse, the PIE, follows directly as

$$g(\chi) = \sum_{\chi \subseteq \Upsilon \subseteq P} (-1)^{|\Upsilon - \chi|} f(\Upsilon) \quad (6.48)$$

with the earlier introduced Möbius function

$$u(\chi, \Upsilon) = \begin{cases} (-1)^{|\Upsilon - \chi|}, & \text{if } \chi \subseteq \Upsilon \\ 0, & \text{otherwise.} \end{cases} \quad (6.49)$$

Next we apply these tools.

Algebraic representations

The vector valued Bernoulli distributions can be now written using

$$\mathbf{p} = \begin{pmatrix} 1 & -1 \\ 0 & 1 \end{pmatrix}^{\otimes N} \mathbf{m} \Leftrightarrow \mathbf{m} = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}^{\otimes N} \mathbf{p}, \quad (6.50)$$

where we immediately recognize on left ζ^{-1} and on right the ζ -matrix obtained via recursive use of the Kronecker tensor product. This is the form given by Teugels in [159], but he did not point out the connection to the incidence algebras or Möbius functions. The form is very similar to ones often used in quantum mechanics with finite degrees of freedom, for example the Hilbert space decompositions into subsystems. Using the Kronecker tensor products, the probability vector \mathbf{p} is

$$p_c = \left\langle \left(\begin{pmatrix} 1 & -1 \\ 0 & 1 \end{pmatrix}^{\otimes N} \begin{pmatrix} 1 \\ B_N \end{pmatrix} \otimes \begin{pmatrix} 1 \\ B_{N-1} \end{pmatrix} \otimes \cdots \otimes \begin{pmatrix} 1 \\ B_1 \end{pmatrix} \right) \right\rangle_c. \quad (6.51)$$

The polynomial vector of *ordinary moments* \mathbf{m} is

$$m_c = \left\langle \prod_{i=1}^N B_i^{b_i^c} \right\rangle = \left\langle \begin{pmatrix} 1 \\ B_N \end{pmatrix} \otimes \begin{pmatrix} 1 \\ B_{N-1} \end{pmatrix} \otimes \cdots \otimes \begin{pmatrix} 1 \\ B_1 \end{pmatrix} \right\rangle_c. \quad (6.52)$$

The polynomial vector of *central moments* $\boldsymbol{\delta}$ is

$$\begin{aligned} \delta_c &= \left\langle \prod_{i=1}^N (B_i - \langle B_i \rangle)^{b_i^c} \right\rangle \\ &= \left\langle \begin{pmatrix} 1 \\ B_N - \langle B_N \rangle \end{pmatrix} \otimes \begin{pmatrix} 1 \\ B_{N-1} - \langle B_{N-1} \rangle \end{pmatrix} \otimes \cdots \otimes \begin{pmatrix} 1 \\ B_1 - \langle B_1 \rangle \end{pmatrix} \right\rangle_c, \end{aligned} \quad (6.53)$$

where c is given by the binary expansion, $0 \leq c \leq 2^N - 1$ and $b_i^c \in \{0, 1\}$ of \mathbf{b}^c . The central moments describe the multipoint correlations ($\# 2^N - N - 1$) between any 2 or more subspaces in phase-space (such as rapidity intervals) and B_i are the corresponding Bernoulli random variables. The correlation functions $\langle B_i \dots B_j \rangle$ factorize to a product of $\langle B_i \rangle \dots \langle B_j \rangle$ if no correlations are present, as usual. This property can be used to test factorization of physics, such as factorization properties of (soft) QCD amplitudes.

Subspace decompositions

Now we drop the zero vector case $c = 0$, which corresponds to the case $B_1 = B_2 = \dots B_N = 0$. This gives no contribution being the additive identity of the algebra, physically outside the fiducial phase space of the definition. Thus the vector \mathbf{p} will be then $2^N - 1$ dimensional, and $\sum p_c = 1$ still by definition. By a slight abuse of notation, we will continue with this convention for the rest of the paper unless otherwise mentioned. Now analogously to Equation 6.50, we write

$$\Lambda \mathbf{p} = \mathbf{r}, \quad (6.54)$$

where \mathbf{r} is an auxiliary vector which collects the additive subspace probabilities of different mutually exclusive final states and Λ is a full rank (invertible) square Boolean matrix of size $2^N - 1$. The number of different subspaces of varying dimensionality is given by q -binomials of Equation C.1 in Appendix C.1.

Example $N = 2$: Denote the Bernoulli probability of the observable B_1 with P_{B_1} and of B_2 with P_{B_2} . The probability of events obeying fiducial binary combinations $\mathbf{b}_1 = (1, 0) \equiv [B_1] \wedge [\neg B_2]$, $\mathbf{b}_2 = (0, 1) \equiv [\neg B_1] \wedge [B_2]$ and $\mathbf{b}_3 = (1, 1) \equiv [B_1] \wedge [B_2]$ can be obtained by inverting a system

$$\begin{pmatrix} 1 & 0 & 1 \\ 0 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} P_{1,0} \\ P_{0,1} \\ P_{1,1} \end{pmatrix} = \begin{pmatrix} P_{B_1} \\ P_{B_2} \\ P_{B_1 \vee B_2} = P_{B_1} + P_{B_2} - P_{B_1 \wedge B_2} \equiv 1 \end{pmatrix}, \quad (6.55)$$

$$\Lambda_{3 \times 3}^{-1} = \begin{pmatrix} 0 & -1 & 1 \\ -1 & 0 & 1 \\ 1 & 1 & -1 \end{pmatrix}. \quad (6.56)$$

Example $N = 3$: Observables are now B_1, B_2, B_3 . We get in a completely analogous way as in the $N = 2$ case

$$\begin{pmatrix} 1 & 0 & 1 & 0 & 1 & 0 & 1 \\ 0 & 1 & 1 & 0 & 0 & 1 & 1 \\ 0 & 0 & 0 & 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 0 & 1 & 1 & 1 \\ 1 & 0 & 1 & 1 & 1 & 1 & 1 \\ 0 & 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 & 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} P_{1,0,0} \\ P_{0,1,0} \\ P_{1,1,0} \\ P_{0,0,1} \\ P_{1,0,1} \\ P_{0,1,1} \\ P_{1,1,1} \end{pmatrix} = \begin{pmatrix} P_{B_1} \\ P_{B_2} \\ P_{B_3} \\ P_{B_1 \vee B_2} = P_{B_1} + P_{B_2} - P_{B_1 \wedge B_2} \\ P_{B_1 \vee B_3} = P_{B_1} + P_{B_3} - P_{B_1 \wedge B_3} \\ P_{B_3 \vee B_2} = P_{B_2} + P_{B_3} - P_{B_2 \wedge B_3} \\ P_{B_1 \vee B_2 \vee B_3} = P_{B_1} + P_{B_2} + P_{B_3} - Z \equiv 1 \end{pmatrix}, \quad (6.57)$$

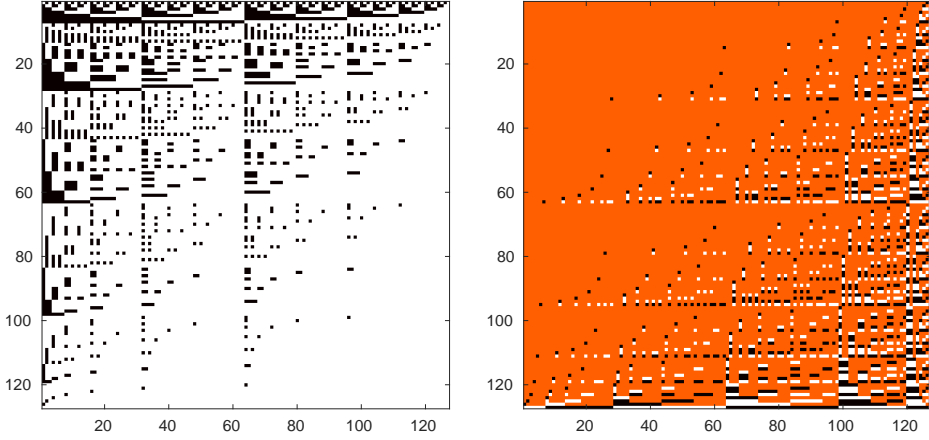


Figure 6.4: Matrices $\Lambda_{127 \times 127}$ and $\Lambda_{127 \times 127}^{-1}$ for $N = 7$. On the left black is 0 and white is 1. On the right, orange is 0, black is -1 and white is 1.

$$\Lambda_{7 \times 7}^{-1} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & -1 & 1 \\ 0 & 0 & 0 & 0 & -1 & 0 & 1 \\ 0 & 0 & -1 & 0 & 1 & 1 & -1 \\ 0 & 0 & 0 & -1 & 0 & 0 & 1 \\ 0 & -1 & 0 & 1 & 0 & 1 & -1 \\ -1 & 0 & 0 & 1 & 1 & 0 & -1 \\ 1 & 1 & 1 & -1 & -1 & -1 & 1 \end{pmatrix}, \quad (6.58)$$

where $Z = P_{B_1 \wedge B_2} + P_{B_1 \wedge B_3} + P_{B_2 \wedge B_3} - P_{B_1 \wedge B_2 \wedge B_3}$.

We can proceed in a same way up to an arbitrary number N of Bernoulli observables, which span a number of $2^N - 1$ non-zero combinations. To summarize: the idea is to generate a system of all linear combinations of event rates described by a full rank (invertible) matrix Λ . This matrix and its inverse will be used directly in the solution to the compound inverse problem. The matrix Λ , or more precisely the inverse of it, can be generated recursively utilizing the principle of inclusion-exclusion

$$P\left(\bigcup_{i=1}^N D_i\right) = \sum_{i=1}^N \left((-1)^{i-1} \sum_{I \subset \{1, \dots, N\}, |I|=i} P(D_I) \right). \quad (6.59)$$

The inner sum runs over all subsets I constructed by the indices $1, \dots, N$, with

cardinality of I identically i elements. Also above we have defined

$$D_I := \bigcap_{i \in I} D_i. \quad (6.60)$$

Figure 6.4 shows an algorithmically generated example where the fractal structure is again evident. Finally the relation between the inverses of ζ -matrix and Λ matrix can be written as

$$[\Lambda^{-1}(\Lambda^{-1})^T]_{i,j} = \begin{cases} [(\zeta^{-1})^T \zeta^{-1}]_{i+1,j+1}, & \text{when } 1 \leq i, j \leq 2^N - 1 \\ [(\zeta^{-1})^T \zeta^{-1}]_{i+1,j+1} - 1, & \text{if } i = j = 2^N - 1. \end{cases}, \quad (6.61)$$

which is evident by careful inspection of the definitions. Because our space \mathbb{F}_2^N is linear, all subspaces \mathbb{F}_2^M with $1 \leq M < N$ are naturally included in it as $\mathbb{F}_2^M \subseteq \mathbb{F}_2^N$, and we used that property. An additional funny remark is that the probability to have an $n \times n$ matrix over \mathbb{F}_2 with determinant $\neq 0$ is $1/4$. This is easily proven.

6.4 Measurements

Regarding proper observables and measurements, we could start to write down suitable theory formulations from which one could perhaps derive semi-analytical distributions for the Bernoulli observables, at least with toy models. However, here our abstraction level is different. This construction is essentially a new model-independent definition of high energy diffraction: different correlations over space-time observables, in terms of rapidity and transverse momentum (and multiplicity), are embedded in the construction and predictions can be obtained directly with any Monte Carlo event generator. We rely only on the observable fiducial final state information. Incidentally, our construction has some similarities with the coherence definition by Glauber in quantum optics [167].

Partial cross sections

To be concrete, a practical measurement procedure of the combinatorial partial cross sections, which may be called also vector fiducial cross sections, goes as follows.

1. **INPUT:** Use minimum bias or zero-bias (bunch cross-over) triggered data.
2. **BEAM BACKGROUND FILTER:** Filter event by event the beam-gas and satellite interactions via detector time-domain cuts, if possible. Tune the beam background cuts using a combination of Monte Carlo (pure beam-beam) and data in a way that beam-beam interaction efficiency is preserved.
3. **FIDUCIAL CUT DEFINITION:** Apply fiducial cuts for pseudorapidity η and transverse momentum p_t . A single fixed set of cuts gives $2^N - 1$ non-zero observable binary combinations. A sliding threshold cuts give a ‘trajectory of observables’.
4. **RESIDUAL BEAM BACKGROUND CORRECTION:** Correct residual beam-gas background at statistical level using beam-empty, empty-beam, empty-empty trigger masks: $N \leftarrow N - \sum_j \alpha_j N_j$, where weight factors α_j are obtained from the j -th trigger mask statistics with pre-determined and/or random trigger downscaling and N is the number of events.
5. **PILE-UP:** Correct pile-up via statistical Möbius inversion using Eq. 6.79, if the run was with high instantaneous luminosity.
6. **LUMINOSITY:** Map event counts to visible cross section units via integrated luminosity, calibrated via van der Meer scans, see Appendix C.5.
7. **MULTIDIMENSIONAL UNFOLDING:** Unfold visible partial cross sections to fiducial partial cross sections. One may use ‘standard’ unfolding algorithms.
8. **OUTPUT:** A vector σ (set) of fiducial cross sections.

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

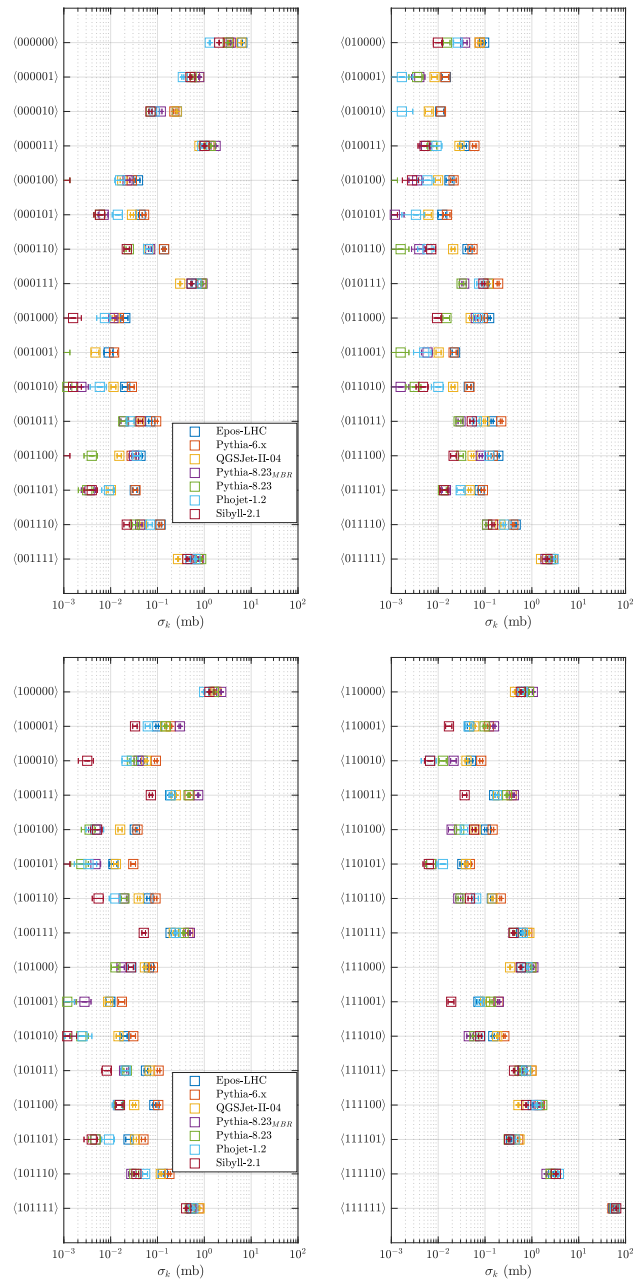


Figure 6.5: Partial cross sections evaluated with proton-proton minimum bias Monte Carlo event generators at $\sqrt{s} = 13$ TeV. Statistical uncertainties are indicated with error bars. The total inelastic cross section for all models is set to 80 mb.

In addition, one may now use the unfolded partial cross sections to obtain highly constraining multidimensional fit extractions of diffractive cross sections and the Pomeron parameters or the maximum diffractive system mass limit, sometimes called the ‘coherence limit’ [168]. Also, re-projections of other observables are possible. For highly efficient algorithms, see Appendix C.6 and C.7.

In Figure 6.5 we have several Monte Carlo model [163, 169–174] fiducial proton-proton cross sections at $\sqrt{s} = 13$ TeV with cuts suitable for the ALICE experiment. The fiducial acceptance domains for charged particles with $N = 6$ are

$$\eta \in [-7.0, -4.9], [-3.7, -1.7], [-2.0, 0], [0, 2.0], [2.8, 5.1], [4.8, 6.3] \wedge p_t > 0.05 \text{ GeV}, \quad (6.62)$$

motivated by the nominal acceptances of AD, V0 and SPD detectors of the ALICE experiment. A large variation between models is seen, representing different tunes and design choices. In certain sense, the vector fiducial cross sections are closer to classic diffraction measurements (aperture + screen), than the typical rapidity gap counting. Regarding the multidimensional unfolding, we recommend using several reference MC models together with the detector simulation. One should also characterize the unfolding regularization parameter strength in a data-driven way, for example, using the Shannon information entropy

$$S = - \sum_i P_i \ln P_i, \quad (6.63)$$

where $P_i \equiv \sigma_i / \sum_j \sigma_j$, by calculating the entropy of the unfolded spectrum and the relative entropy between the unfolded and raw combination rates

$$D = \sum_i P_i \ln \left(\frac{P_i}{\bar{P}_i} \right). \quad (6.64)$$

Naturally, pure simulation studies should be done also. The non-trivial systematics behind efficiency corrections are made here very explicit by the multidimensional unfolding, which accounts for both efficiency losses and the flow of event topologies from one combination to another. This flow is not just because of efficiency losses but also due to major material re-scattering and other detector effects, often hardly understood with single dimensional rapidity gap counting measurements.

To this end, we point out that in principle one could extend the fiducial measurement vector space from the GF(2) finite field \mathbb{F}_2^N to \mathbb{F}_q^N , where q is an integer power of a prime number. The prime power is a finite field requirement. We see that this could be useful regarding different rapidity slices containing different extreme (low or high) multiplicity densities in the same event, not necessarily hard scale or jets driven, for example.

Fractal marginal distributions

We proceed as before, but now measure all single dimensional distributions available per binary combination, such as the multiplicity $P(N_{ch})$ and transverse momentum $P(p_t)$ distributions, once the slicing is done over (pseudo)rapidity. For example, assume $N = 3$ slices over pseudorapidity. For the combination $\langle 0, 0, 1 \rangle$ we have one non-zero component $P(N_{ch})_{\langle 0,0,[1] \rangle}$, for $\langle 0, 1, 1 \rangle$ we get $P(N_{ch})_{\langle 0,1,[1] \rangle}$, $P(N_{ch})_{\langle 0,[1],1 \rangle}$ and similarly for transverse momentum. As an illustrative example, we tabulate those in Tables 6.2 and 6.3 with $N = 3$. The number $a(n)$ of non-zero distributions per vector combination is enumerated with $a(0) = 0$, $a(2n) = a(n)$, $a(2n + 1) = a(n) + 1$, producing a fractal sequence $0, 1, 1, 2, 1, 2, 2, 3, \dots$. That is simply the Hamming weight. The total number of distributions is $N2^{N-1}$, the number of edges in an N -hypercube.

1	\emptyset	\emptyset	$f(\mathcal{O}_x)_{\langle 0,0,[1] \rangle}$
2	\emptyset	$f(\mathcal{O}_x)_{\langle 0,[1],0 \rangle}$	\emptyset
3	\emptyset	$f(\mathcal{O}_x)_{\langle 0,[1],1 \rangle}$	$f(\mathcal{O}_x)_{\langle 0,1,[1] \rangle}$
\vdots		\vdots	
8	$f(\mathcal{O}_x)_{\langle [1],1,1 \rangle}$	$f(\mathcal{O}_x)_{\langle 1,[1],1 \rangle}$	$f(\mathcal{O}_x)_{\langle 1,1,[1] \rangle}$
$\Sigma \downarrow$	$\underbrace{f(\mathcal{O}_x)_{\langle [1],-,- \rangle}}_{\int_{\Delta\eta_1} d\eta}$	$\underbrace{f(\mathcal{O}_x)_{\langle -,[1],- \rangle}}_{\int_{\Delta\eta_2} d\eta}$	$\underbrace{f(\mathcal{O}_x)_{\langle -,-,[1] \rangle}}_{\int_{\Delta\eta_3} d\eta}$

Table 6.2: Combinatorially embedded marginal 1D-distributions of the observable \mathcal{O}_x with $N = 3$ slices over pseudorapidity.

1	\emptyset	\emptyset	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle 0,0,[1] \rangle}$
2	\emptyset	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle 0,[1],0 \rangle}$	\emptyset
3	\emptyset	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle 0,[1],1 \rangle}$	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle 0,1,[1] \rangle}$
\vdots		\vdots	
8	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle [1],1,1 \rangle}$	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle 1,[1],1 \rangle}$	$f(\mathcal{O}_x, \mathcal{O}_y)_{\langle 1,1,[1] \rangle}$
$\Sigma \downarrow$	$\underbrace{f(\mathcal{O}_x, \mathcal{O}_y)_{\langle [1],-,- \rangle}}_{\int_{\Delta\eta_1} d\eta}$	$\underbrace{f(\mathcal{O}_x, \mathcal{O}_y)_{\langle -,[1],- \rangle}}_{\int_{\Delta\eta_2} d\eta}$	$\underbrace{f(\mathcal{O}_x, \mathcal{O}_y)_{\langle -,-,[1] \rangle}}_{\int_{\Delta\eta_3} d\eta}$

Table 6.3: Combinatorially embedded marginal 2D-distributions of observables $(\mathcal{O}_x, \mathcal{O}_y)$ with $N = 3$ slices over pseudorapidity.

We emphasize here that this is just cut-and-count with bookkeeping, a well defined fiducial measurement. The problem of unfolding from detector level to particle level, however, can be non-trivial. We believe that these fractal distributions together with combinatorial partial cross sections can be key handles to test the AGK rules and (topological) fluctuations, perhaps driven by anomalous QCD color field configurations. A potentially related theory discussion can be found in [175, 176].

Threshold gapflow trajectories

The idea here is a straightforward one. One measures $2^N - 1$ fiducial vector cross sections over rapidity as a function of a threshold condition. The threshold condition can be global over rapidity, or local. For example, a minimum transverse momentum threshold, a minimum multiplicity threshold or a minimum flavor content threshold (e.g. strangeness).

The measurement is illustrated in Figure 6.6, where the transverse momentum threshold is changed for four different average particle densities, particles drawn from a simple Poisson distribution with the given mean value. Here we used a simple soft exponential p_t^2 distribution with particles generated uniformly over rapidity with $N = 3$ slices. The exponential in p_t^2 results in a Rayleigh type p_t -distribution with Gaussian distributed p_x and p_y components. Now it is very easy to see that rapidity gaps are not invariant under transverse momentum cutoffs, a well known fact but sometimes a forgotten truth. Depending on the particle density and the cutoff, different event topology rates may have a cross over. The transition to pQCD with power law tails will be interesting.

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

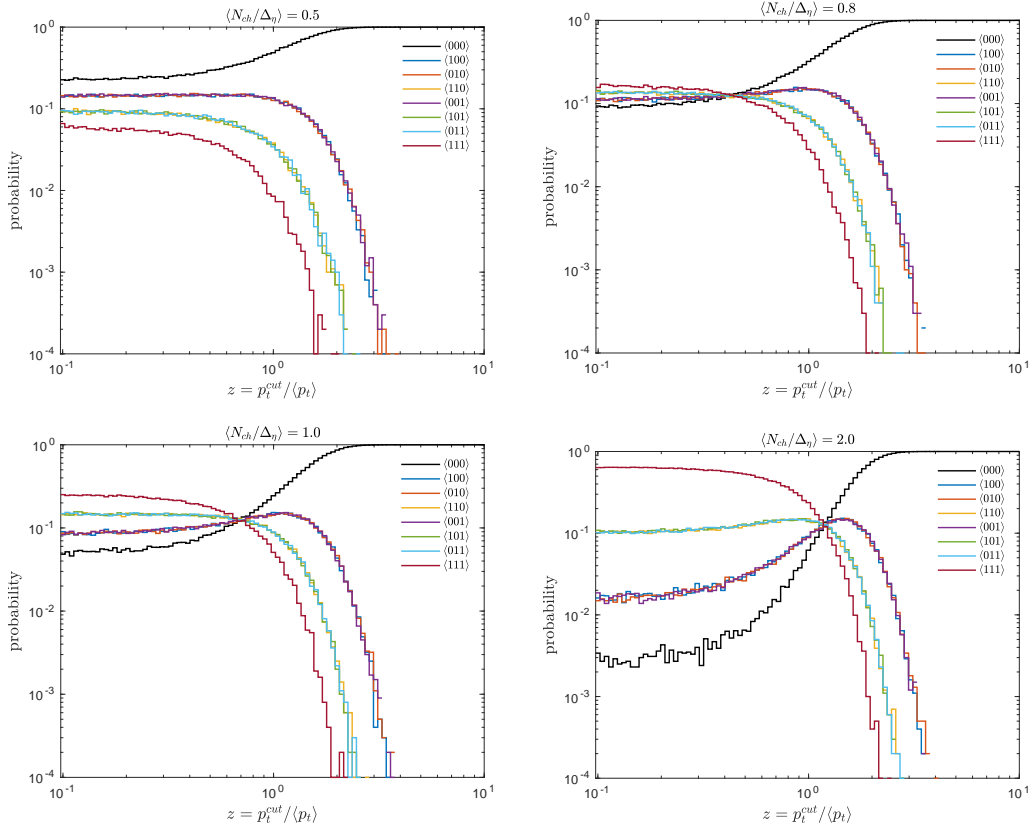


Figure 6.6: A toy simulation of ‘gapflow trajectories’ as a function of the normalized transverse momentum cutoff $z = p_t^{cut} / \langle p_t \rangle$ with four different particle density.

6.5 Combinatorial supercompound Poisson process

By using the tools introduced earlier, we formulate a Poisson superposition problem and its inverse. This is essentially a minimal non-trivial mathematical model needed to understand the problematics involved, before trying to deconstruct more complicated scenarios. This model has no interfering amplitude structures, *a priori*, whereas the AGK type compound interactions do contain interfering amplitudes, by construction. We will see later that certain boundary conditions will force the solutions to this problem towards more complicated scenarios.

We start by assuming a Poisson fluctuating random variable $K \sim \frac{\mu^k}{k!} e^{-\mu}$ with mean μ and argument $k \in \mathbb{N}$ denoting the number of simultaneous interactions, such as multiparton interactions or multiple independent proton-proton interactions. Both cases are treated analogously, modulo the energy-momentum conservation differences. The probability generating function for Poisson is $G(z) = \mathbb{E}(z^K) = \exp(\mu(z - 1))$ which is useful for deriving the compound moments. In general, the expectation values and variances of a random variable $C_{|K>0} = \sum_{j=1}^K X_j$ after compounding stochastics are obtained from the compound relations

$$\langle C \rangle = \mathbb{E}_{K>0}[\mathbb{E}_{C|K>0}(C)] = \langle K > 0 \rangle \langle X \rangle \quad (6.65)$$

$$\begin{aligned} \text{Var}[C] &= \mathbb{E}_{K>0}[\text{Var}_{C|K>0}(C)] + \text{Var}_{K>0}[\mathbb{E}_{C|K>0}(C)] \\ &= \langle K > 0 \rangle \text{Var}[X] + \text{Var}[K > 0] (\langle X \rangle)^2. \end{aligned} \quad (6.66)$$

The zero-truncated Poisson mean and variance of K are given in Appendix C.4. We need the zero truncated distribution, because the 0 case is not a physical observable. However, all the N -point moments necessary to describe the multidimensional case, i.e. the case $N > 1$, are *not* handled with these simple formulas but only via the full machinery which follows.

For preliminaries, the ordinary binomial coefficient is

$$C_k^n = \binom{n}{k} = \frac{n!}{k!(n-k)!}, \quad 0 \leq k \leq n \quad (6.67)$$

which describes how many ways k numbers can be chosen from a set of n numbers. Counting Bernoulli random variables in a binomial setup, one obtains the Binomial distribution. By counting random variables in a multivariate Bernoulli setup, we obtain analogously the multivariate Binomial distribution. However, for the sake of practicality we use now the multinomial distribution with random variables $X_1, X_2, \dots, X_{2N-1}$, which is as many as we have non-zero fiducial vectors in

the binary space. The multinomial probability distribution is

$$P_m(X_1 = x_1, \dots, X_{2^N-1} = x_{2^N-1} | \mathbf{p}) = \frac{k!}{\prod_{c=1}^{2^N-1} x_c!} \prod_{c=1}^{2^N-1} p_c^{x_c} \quad (6.68)$$

$$\text{with } \sum_{c=1}^{2^N-1} x_c = k, \quad x_c \in \mathbb{N} \text{ and } \sum_c p_c = 1, \quad (6.69)$$

which represents sampling or occupying a different number x_i of different fiducial vectors per mixed event, the number of vectors in superposition per event is the number k . The first term is the multinomial coefficient easily derived using the binomial coefficient, also known in Maxwell-Boltzmann statistics. The probability generating function is here given by $G(\mathbf{z}) = \left(\sum_{c=1}^{2^N-1} p_c z_c \right)^k$, $\mathbf{z} \in \mathbb{C}^n$ which is useful, when expanded, in interpreting multinomial distribution as polynomial coefficients p_c with unit sum.

Here *distinguishability* means that we can separate between different final state vectors but not between a vector of the same type. This is a sampling with replacement scheme. Sampling without replacement gives, for example, a multivariate hypergeometric distribution instead of the multinomial. Similar non-multinomial sampling scenarios are encountered with Fermi-Dirac and Bose-Einstein statistics, for example.

Direct model

We define a generative direct model $f : \Theta \rightarrow Y$, where Θ denotes the space of parameters and Y the space of observables. The unknown parameters are in $\mathbf{p} = [p_1, p_2, \dots, p_n]^T \in \Theta$ and the compounding process diluted or enhanced combination rate measurements are in $\mathbf{y} = [y_1, y_2, \dots, y_n]^T \in Y$. The model estimates for the observed rates y_c are generated, as a sum over product of combinatorially selected multinomial distribution terms W_{ck} and Poisson probabilities $P_p(k = j)$ order by order as

$$\begin{pmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{pmatrix} = \frac{1}{1 - e^{-\mu}} \begin{pmatrix} W_{11} & W_{12} & W_{13} \dots W_{1\infty} \\ W_{21} & W_{22} & W_{23} \dots W_{2\infty} \\ & \vdots & \ddots \\ W_{n1} & W_{n2} & W_{n3} \dots W_{n\infty} \end{pmatrix} \begin{pmatrix} P_p(k = 1; \mu) \\ P_p(k = 2; \mu) \\ P_p(k = 3; \mu) \\ \vdots \\ P_p(k = \infty; \mu) \end{pmatrix}, \quad (6.70)$$

where we have an infinite series. The (re)-normalization factor $P_p(0; \mu) = 1 - e^{-\mu}$ is needed, because we exclude (truncate) the null case $k = 0$ of the Poisson density. Solving the inverse of Eq. 6.70 is a non-linear problem, because the matrix elements of W depend on the unknown parameters. Now more explicitly written, the *compound Poisson* model equation for each probability component is

$$\begin{aligned} y_c &= \frac{1}{1 - e^{-\mu}} \sum_{k=1}^{\infty} \frac{\mu^k}{k!} e^{-\mu} W_{ck} \\ &= \frac{e^{-\mu}}{1 - e^{-\mu}} \sum_{k=1}^{\infty} \frac{\mu^k}{k!} \left\{ \sum_{\{x_j\} \subset \Omega_{ck}} \frac{k!}{\prod_{j=1}^n x_j!} \prod_{j=1}^n p_j^{x_j} \right\}, \end{aligned} \quad (6.71)$$

with $\sum_c y_c = 1$. The multinomial term and its values of x_j are evaluated over all valid combinations or partitions for y_c from the set of n -tuples Ω_{ck} , that is, those which are allowed by pileup poset combinatorics. Formally, this set of compositions can be enumerated with

$$\Omega_{ck} = \left\{ (x_1, \dots, x_j, \dots, x_n) \mid \bigvee_j x_j \mathbf{b}_j = \mathbf{b}_i \text{ and } \sum_j x_j = k \right\}, \quad (6.72)$$

where the boolean \bigvee operator takes care of ‘summing’ the binary vectors \mathbf{b}_j of multiplicity x_j and thus evaluating the pileup compositions. The p_j are the probabilities we want to find out by inverting the compound Poisson process, that is, the probabilities of each $n = 2^N - 1$ fiducial final state.

Example $N = 2$: a (heavily) truncated Poisson series at $\mathcal{O}(\mu^2)$ is

$$W_{3 \times 2} = \begin{pmatrix} p_1 & P_m(2, 0, 0) \\ p_2 & P_m(0, 2, 0) \\ p_3 & P_m(1, 1, 0) + P_m(1, 0, 1) + P_m(0, 1, 1) + P_m(0, 0, 2) \end{pmatrix}. \quad (6.73)$$

Now a brute force inverse $f^{-1} : Y \rightarrow \Theta$ would be to use Eq. 6.71 and then estimate parameters by minimizing a non-linear χ^2 -function between the generated \hat{y}_c and measured y_c values $\hat{\mathbf{p}} = \arg \min_{\mathbf{p}} \chi^2$, where $\chi^2 = (\hat{\mathbf{y}} - \mathbf{y})^T \Sigma_{\mathbf{y}}^{-1} (\hat{\mathbf{y}} - \mathbf{y})$ and usually $\Sigma_{\mathbf{y}}^{-1} = \text{diag}(1/\delta y_c^2)$. This is a typical fit approach, which is the statistical Maximum Likelihood (ML) solution when the measurement noise (counting fluctuations) is approximated with Gaussian distribution. A ML solution taking into account the multinomial counting fluctuations can be formulated also. However, if the direct model is constructed just order by order in k , this approach is computationally

extremely expensive at high values of k (and N) due to factorial growth. This can be seen in Table C.2 in Appendix C.2. Instead of using this, we will formulate an ‘all-order’ solution based on the Möbius inversion.

Solution for $N \leq 2$

The first case $N = 1$ is trivial, because we have just one Bernoulli observable, no combinations involved and because of our normalization condition $\sum_c p_c = 1$. Thus, the most interesting quantity is directly the Poisson distribution μ -value and probabilities for $k = 0, 1, 2, \dots$.

The case $N = 2$ is straightforward. An explicit formula for y_1 is obtained by noticing that the set Ω_{1k} allows only ‘autocompound’ for all k . Thus $B_1 \equiv k$ and $B_2 \equiv B_3 \equiv 0$ for all values of k , which gives

$$\begin{aligned} y_1 &= \frac{e^{-\mu}}{1 - e^{-\mu}} \sum_{k=1}^{\infty} \frac{\mu^k}{k!} \left(\frac{k!}{k!0!0!} p_1^k p_2^0 p_3^0 \right) \\ &= \frac{e^{-\mu}}{1 - e^{-\mu}} \sum_{k=1}^{\infty} \frac{\mu^k}{k!} p_1^k \\ &= \frac{e^{\mu p_1} - 1}{e^{\mu} - 1}. \end{aligned} \quad (6.74)$$

Above we used the definition of Taylor series for e^x . By calculating similarly for y_2 by replacing above p_1 with p_2 , then we get y_3 using the conservation of probability

$$y_3 = 1 - \sum_{c=1}^2 \frac{e^{\mu p_c} - 1}{e^{\mu} - 1}. \quad (6.75)$$

Now the inverse is given by

$$\hat{p}_c = \frac{\ln((e^{\mu} - 1)y_c + 1)}{\mu}, \quad \text{for } c = 1, 2 \quad (6.76)$$

and again by conservation of probability

$$\hat{p}_3 = 1 - \sum_{c=1}^2 \frac{\ln((e^{\mu} - 1)y_c + 1)}{\mu}, \quad (6.77)$$

which completes the solution.

Solution for $N \geq 1$

The arbitrary N case follows recursively from the $N = 2$ case and can be done in practice by constructing the matrix Λ and its inverse according to the incidence algebra described earlier. The main result follows. The combinatorial statistics is

$$\mathbf{y}(\mathbf{p}; \mu) = \Lambda^{-1} \frac{\exp[-\mu\Lambda\mathbf{p}] - 1}{e^{-\mu} - 1} \quad (6.78)$$

which has also the inverse we are looking for

$$\hat{\mathbf{p}}(\mathbf{y}; \mu) = \Lambda^{-1} \frac{\ln[(e^{-\mu} - 1)\Lambda\mathbf{y} + 1]}{-\mu}. \quad (6.79)$$

The exponential and logarithm are taken vector entry wise, such that the Taylor series representation is using the Hadamard power: $\exp[\mathbf{v}] = \sum_{k=0}^{\infty} \frac{1}{k!} \mathbf{v}^{\odot k}$. Equations 6.78, 6.79 are relying on the matrix Λ constructing linear combinations of the probabilities of elements in \mathbb{F}_2^N and the property that multinomial probability distribution is factorized as a product of different terms. The multinomial distribution obeys the so-called grouping property: re-partitioning conserves the multinomial distribution. Similarly Λ^{-1} deconstructs the linear combinations, and in between there is a non-linear transformation according to the compound Poisson statistics. One can understand this also in terms of convolution theorem which turns the convolution into a usual product in the transform (Fourier) domain or vice verse. Here the Möbius inversion has the same role as Fourier decomposition.

The example with $N = 3$ gives a vector

$$\hat{\mathbf{p}} = \frac{1}{\mu} \begin{pmatrix} \ln(e^{\mu} y_1 + 1) \\ \ln(e^{\mu} y_2 + 1) \\ - \sum_{c=1,2} \ln(e^{\mu} y_c + 1) + \ln(1 + \sum_{c=1,2,3} e^{\mu} y_c) \\ \ln(e^{\mu} y_4 + 1) \\ - \sum_{c=1,4} \ln(e^{\mu} y_c + 1) + \ln(1 + \sum_{c=1,4,5} e^{\mu} y_c) \\ - \sum_{c=2,4} \ln(e^{\mu} y_c + 1) + \ln(1 + \sum_{c=2,4,6} e^{\mu} y_c) \\ \mu + \sum_{c=1,2,4} \ln(e^{\mu} y_c + 1) - \ln(1 + \sum_{c=1,2,3} e^{\mu} y_c) - \ln(1 + \sum_{c=1,4,5} e^{\mu} y_c) - \ln(1 + \sum_{c=2,4,6} e^{\mu} y_c) \end{pmatrix},$$

where by conservation of probability we chose to fix $y_7 = 1 - \sum_{c=1}^6 y_c$ and for compacting the notation we set $e^{\mu}_- \equiv e^{\mu} - 1$. The fractal structure of the problem is visible again. Components $c = 1, 2, 4$ only undergo autocompound, thus no mixing

between other vectors. These are the unit basis vectors $[1, 0, \dots, 0]^T$, $[0, 1, \dots, 0]^T$ and $[0, 0, \dots, 1]^T$ and cannot be thus constructed as a linear combination of other vectors. A practical remark when working with non-linear equations with exponential and logarithmic functions: numerical evaluations are prone to floating point accuracy problems at high μ values.

Maximum Input Entropy

The constraints used in the construction required that both \mathbf{p} and \mathbf{y} are probability distributions and that the problem satisfies the combinatorial incidence algebra, together with the multinomial statistics embedded in the Poisson process. Next, we treat certain special solutions.

One interesting case to study is the maximum input entropy case $\mathbf{p} \equiv \mathbf{1}/n$, that is, all different combinations are equally probably. We go through the case $N = 3$ here explicitly. First we solve

$$\begin{aligned} \ln(e_-^\mu y_1 + 1)/\mu &= \frac{1}{n} = \frac{1}{7} \\ y_1 &= \frac{e^{\mu/7} - 1}{e_-^\mu} \equiv y_2 \equiv y_4. \end{aligned} \quad (6.80)$$

Then get we the third component by first substituting y_1

$$\begin{aligned} [-2\frac{\mu}{7} + \ln(1 + 2(e^{\mu/7} - 1) + e_-^\mu y_3)]/\mu &= \frac{1}{7} \\ y_3 &= \frac{-2e^{\mu/7} + e^{3\mu/7} + 1}{e_-^\mu} \equiv y_5 \equiv y_6. \end{aligned} \quad (6.81)$$

The last component is obtained by conservation of probability

$$y_7 = 1 - y_{1:6} = \frac{3e^{\mu/7}(1 - e^{2\mu/7})}{e_-^\mu} + 1, \quad (6.82)$$

where $y_{1:6}$ denotes a sum over all the given indices. It is interesting also to take a look at

$$f(\mu) = \sum_{c=3,5,6} y_c - \sum_{c=1,2,4} y_c = \frac{3(2 - 3e^{\mu/7} + e^{3\mu/7})}{e_-^\mu}. \quad (6.83)$$

Maximum input entropy solutions are represented in Figure 6.7 as a function of μ . We see that there is a fixed maximum output entropy point when $\mathbf{y} = \mathbf{1}/n$ with the critical μ -value depending on N .

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

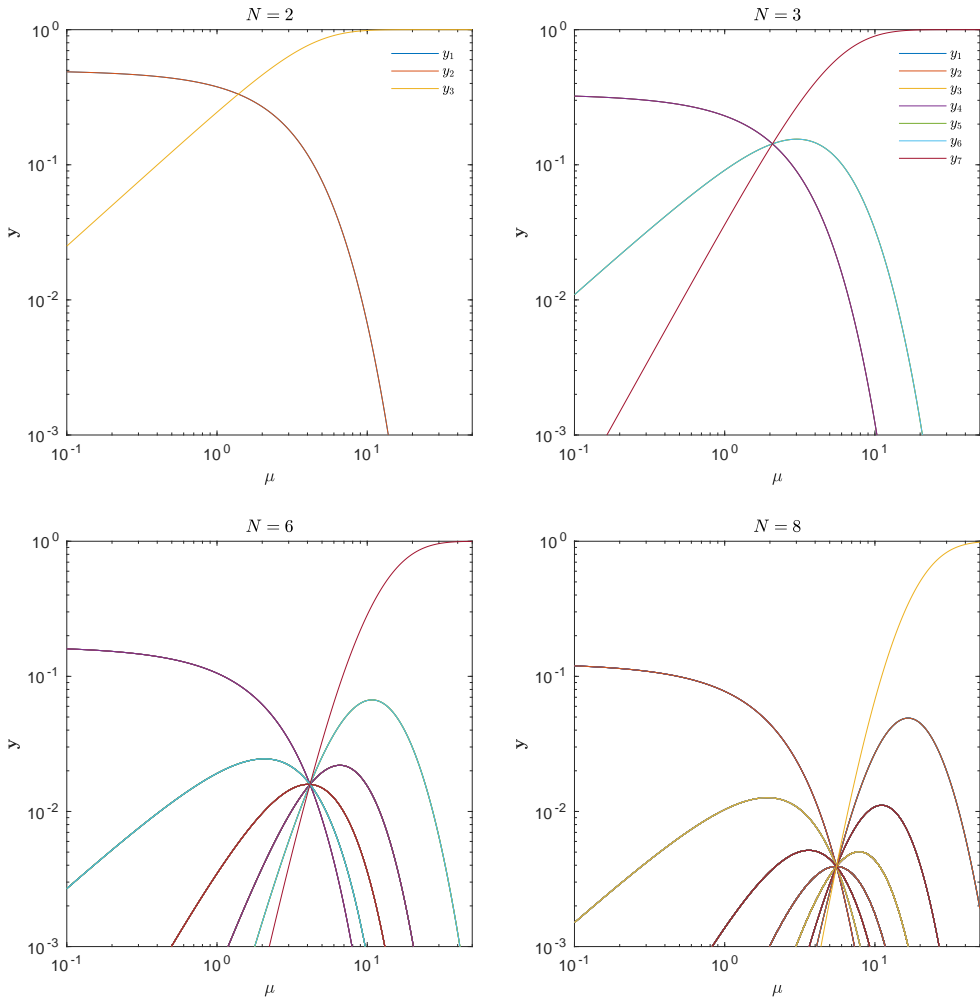


Figure 6.7: Trajectories of \mathbf{y} with the maximum input entropy initial condition $\mathbf{p} \equiv \mathbf{1}/n$.

Maximum output entropy

Here we construct solutions with maximum output entropy. Examples of functional behavior are given in Figure 6.8. We see that it is not possible to obey the constraints of all probabilities being non-negative, thus we obtain \mathbf{p} with negative elements when μ is high enough. An interesting oscillating phenomenon appears with nodes at higher μ due to the polynomial structure of our construction.

Negative values are seemingly unphysical unless one extends the domain of the problem to include also ‘negative probabilities’ – sometimes called *quasi-probability distributions*, utilized first by Wigner [177]. In Wigner’s construction, these are fundamentally non-classical distributions of quantum states. The construction here has certain analogous behavior, essentially because it is the vector \mathbf{y} which are observables and the vector \mathbf{p} contains the parameters of the states, which may have non-classical distributions.

Non-classical distributions of \mathbf{p} here may thus be useful for quantum interference, absorption, shadowing and diffraction. A natural extension is complex vector distributions. Negative probabilities as a useful intermediate step of calculations were also discussed in a well-known article by Feynman [178] in the context of two non-commuting observables.

Master Equation

The first order derivative of \mathbf{y} with respect to μ results in

$$\frac{\partial \mathbf{y}}{\partial \mu} = \Lambda^{-1} \left(\frac{\exp[-\mu \Lambda \mathbf{p}] - 1}{e^\mu (e^{-\mu} - 1)^2} - \frac{\Lambda \mathbf{p} \odot \exp[-\mu \Lambda \mathbf{p}]}{e^{-\mu} - 1} \right). \quad (6.84)$$

The higher order derivatives follow similarly, with lengthy expressions which we evaluated with symbolic computer algebra. It is this expression which can be called the *master equation* of describing the evolution of probabilities in terms of the μ -value. Visualized, we obtain wavelet like behavior shown in Figure 6.9. Now obvious future direction would be to probe a set of stochastic equations which would use these basic building blocks. In Figure 6.9 we show also the ‘analytically continued’ case with negative μ -values for completeness and leave the possible applications out of discussion here. The values of \mathbf{y} are positive with positive μ values, but obtain oscillating positive and negative values at negative μ . The behavior is actually opposite, if we invert the case and study values of \mathbf{p} at negative μ when $\mathbf{y} \equiv \mathbf{1}/n$.

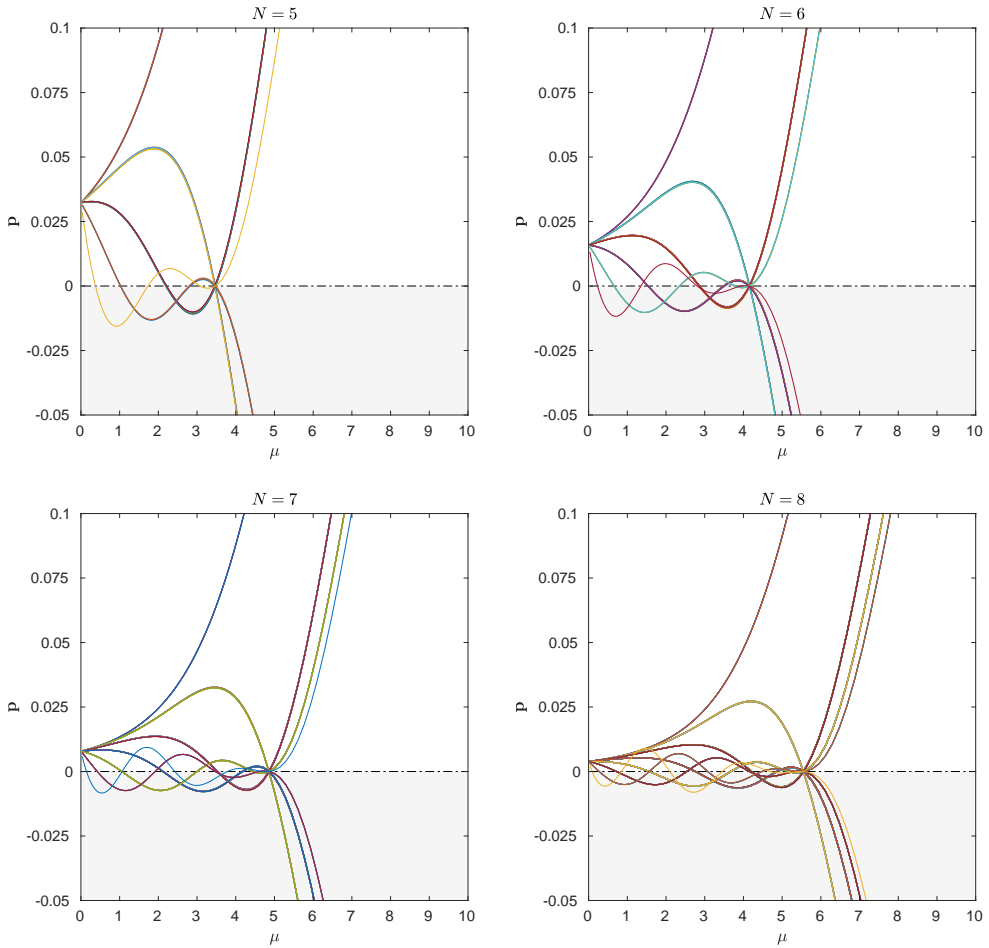


Figure 6.8: Inverted trajectories of \mathbf{p} with the maximum output entropy $\mathbf{y} \equiv \mathbf{1}/n$ boundary condition. The gray area shows the non-positive definite domain.

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

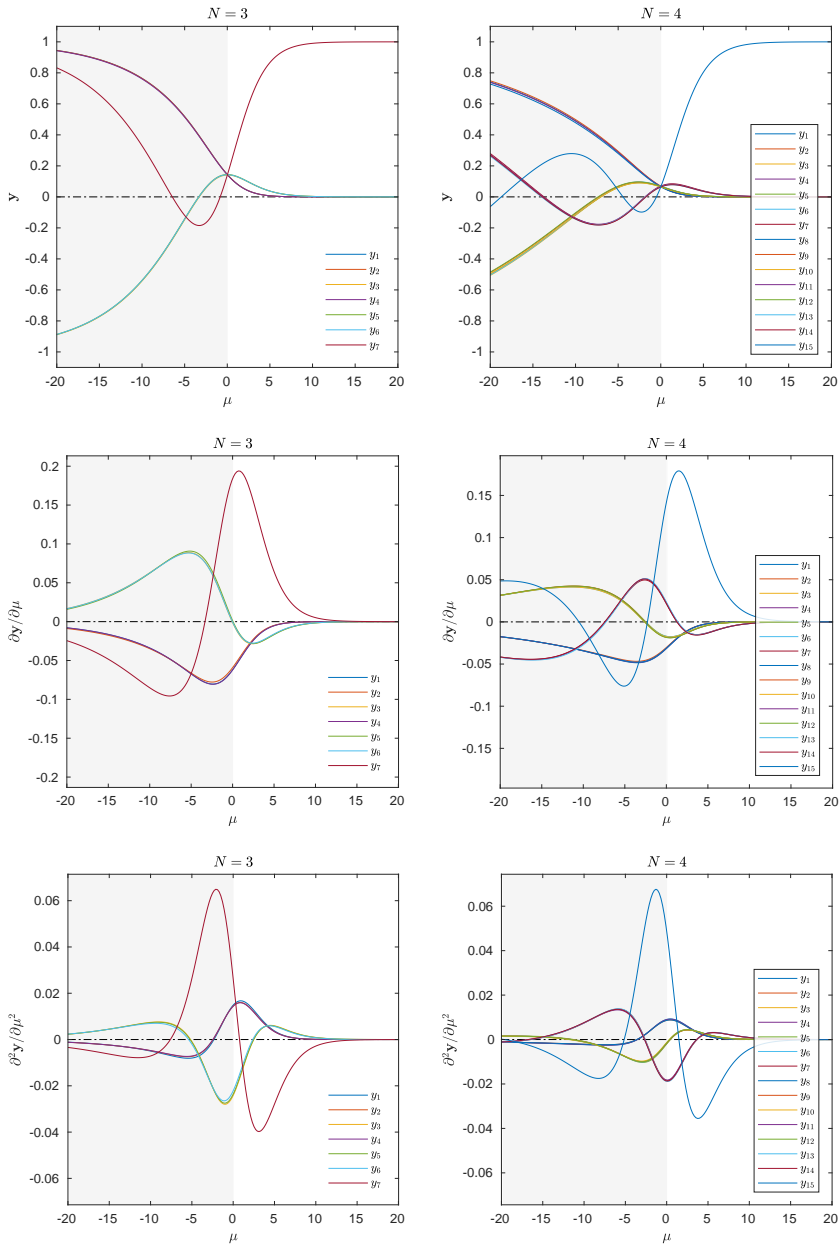


Figure 6.9: Trajectories (1st row) and first and second order derivatives (2nd and 3rd rows) with the $\mathbf{p} \equiv \mathbf{1}/n$ input, for $N = 3$ and $N = 4$. We have added a tiny perturbation to the input for visualization (parity degeneracy).

Semi-Bose-Einstein construction

The construction above was the most relaxed one, in a sense that we allow for example $[0, 1] + [0, 1] \rightarrow [0, 1]$ type vector combinations. However, one can construct a modified statistics by forbidding same the vector combination appearing more than once, but not limiting occupation number per vector dimension. This will limit a large number of possible combinations meaning all ‘autocompound’ is forbidden. In even more strict construction, on the other hand, we would allow only one occupation per vector dimension. We recall that different canonical particle statistics are most easily derived using the combinatorial method presented in any advanced book on statistical mechanics.

The case $N = 2$ is

$$\begin{aligned}
 y_1 &= P_{K=1}P_m(1, 0, 0) \\
 y_2 &= P_{K=1}P_m(0, 1, 0) \\
 y_3 &= P_{K=1}P_m(0, 0, 1) \\
 &\quad + P_{K=2}[P_m(1, 1, 0) + P_m(1, 0, 1) + P_m(0, 1, 1)] \\
 &\quad + P_{K=3}P_m(1, 1, 1),
 \end{aligned} \tag{6.85}$$

where the notation is the same as before.

The case $N = 3$ is

$$\begin{aligned}
 y_1 &= P_{K=1}P_m(1, 0, 0, 0, 0, 0, 0) \\
 y_2 &= P_{K=1}P_m(0, 1, 0, 0, 0, 0, 0) \\
 y_3 &= P_{K=1}P_m(0, 0, 1, 0, 0, 0, 0) \\
 &\quad + P_{K=2}[P_m(1, 1, 0, 0, 0, 0, 0) + P_m(1, 0, 1, 0, 0, 0, 0) + P_m(0, 1, 1, 0, 0, 0, 0)] \\
 &\quad + P_{K=3}P_m(1, 1, 1, 0, 0, 0, 0) \\
 y_4 &= P_{K=1}P_m(0, 0, 0, 1, 0, 0, 0) \\
 y_5 &= P_{K=1}P_m(0, 0, 0, 0, 1, 0, 0) \\
 &\quad + P_{K=2}[P_m(1, 0, 0, 1, 0, 0, 0) + P_m(1, 0, 0, 0, 1, 0, 0) + P_m(0, 0, 0, 1, 1, 0, 0)] \\
 &\quad + P_{K=3}P_m(1, 0, 0, 1, 1, 0, 0) \\
 y_6 &= P_{K=1}P_m(0, 0, 0, 0, 0, 1, 0) \\
 &\quad + P_{K=2}[P_m(0, 1, 0, 1, 0, 0, 0) + P_m(0, 0, 0, 1, 1, 0, 0) + P_m(0, 1, 0, 0, 1, 0, 0)] \\
 &\quad + P_{K=3}P_m(0, 1, 0, 1, 0, 1, 0) \\
 y_7 &= 1 - y_{1:6},
 \end{aligned} \tag{6.86}$$

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

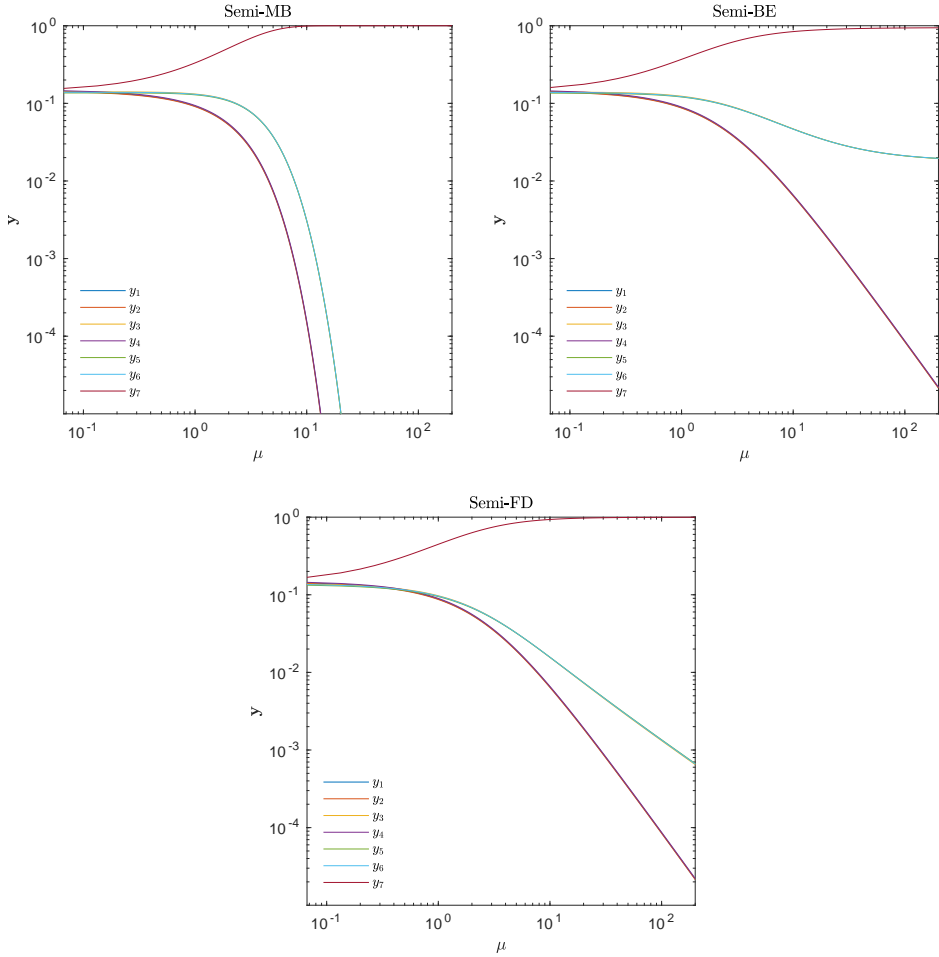


Figure 6.10: Semi-Maxwell-Boltzmann, Semi-Bose-Einstein and Semi-Fermi-Dirac boundary conditions with the maximum input entropy $\mathbf{p} = \mathbf{1}/n$ initial condition.

where the last component is obtained simply by conservation of probability. In this ‘Semi-Bosonic’ case, we did not find such a simple closed-form formula as the matrix equation 6.79, but solutions for any N can be found simply algorithmically using symbolic computer algebra.

Semi-Fermi-Dirac construction

In the Semi-Fermi-Dirac construction, which is the most strict one we consider here, we allow only one occupation per vector element.

The case $N = 2$ is

$$\begin{aligned}
 y_1 &= P_{K=1}P_m(1, 0, 0) \\
 y_2 &= P_{K=1}P_m(0, 1, 0) \\
 y_3 &= P_{K=1}P_m(0, 0, 1) + P_{K=2}P_m(1, 1, 0).
 \end{aligned} \tag{6.87}$$

The case $N = 3$ is

$$\begin{aligned}
 y_1 &= P_{K=1}P_m(1, 0, 0, 0, 0, 0, 0) \\
 y_2 &= P_{K=1}P_m(0, 1, 0, 0, 0, 0, 0) \\
 y_3 &= P_{K=1}P_m(0, 0, 1, 0, 0, 0, 0) + P_{K=2}P_m(1, 1, 0, 0, 0, 0, 0) \\
 y_4 &= P_{K=1}P_m(0, 0, 0, 1, 0, 0, 0) \\
 y_5 &= P_{K=1}P_m(0, 0, 0, 0, 1, 0, 0) + P_{K=2}P_m(1, 0, 0, 1, 0, 0, 0) \\
 y_6 &= P_{K=1}P_m(0, 0, 0, 0, 0, 1, 0) + P_{K=2}P_m(0, 1, 0, 1, 0, 0, 0) \\
 y_7 &= 1 - y_{1:6}.
 \end{aligned} \tag{6.88}$$

Again, the any N case can be constructed algorithmically.

The results are shown in Figure 6.10, where different constraints on the combinations result in different asymptotic behavior in the tails.

6.6 Invertibility simulations

To evaluate the finite sample statistics invertibility from \mathbf{y} into \mathbf{p} using Equation 6.79, simulations are necessary. Standard numerical uncertainty methods such as bootstrapping may be used to quantify the uncertainties in-situ. We do the simulation in a fully controlled way and use the multidimensional Dirichlet distribution

$$f_D(x_1, \dots, x_n; \alpha_1, \dots, \alpha_n) = \frac{1}{B(\boldsymbol{\alpha})} \prod_{i=1}^n x_i^{\alpha_i-1}, \text{ where } B(\boldsymbol{\alpha}) = \frac{\prod_{i=1}^n \Gamma(\alpha_i)}{\Gamma(\sum_{i=1}^n \alpha_i)} \quad (6.89)$$

as a random distribution for probabilities of \mathbf{p} with a constraint $\sum_c p_c = 1$. This constraint is by construction included in the Dirichlet distribution. The normalization factor is the multivariate Beta function, seemingly closely related to the famous Veneziano amplitudes of early string theory. We illustrate the Dirichlet distribution drawn realizations in Figure 6.13, where we observe how the Shannon entropy $S(\mathbf{y})$ of the superposition result is running as a function of the compounding Poisson process μ -value. Interestingly, the entropy has often a non-monotonic, somewhat universal behavior. Picking up random distributions allows us to sample the spectrum of different probability distributions of \mathbf{p} maximally. This means a wide range of possible applications can be effectively probed, in a model independent way. We set the parameter vector of the Dirichlet distribution $\boldsymbol{\alpha} = \alpha \mathbf{1}$ with a concentration parameter $\alpha \in \{0.5, 1\}$. The case $\alpha = 1$ corresponds to the maximally uniform case in the multidimensional space of probabilities.

As a measure of agreement between the estimated probabilities and the true ones, we use the Kolmogorov-Smirnov type test statistic

$$D = \max |F(i) - \hat{F}(i)| \in [0, 1], \quad i = 1, \dots, n \quad (6.90)$$

where $F(i) = \sum_{j=1}^i p_j$ is the cumulative distribution function (CDF) of the true components of \mathbf{p} and $\hat{F}(i)$ is the CDF of the estimate $\hat{\mathbf{p}}$. We are now neglecting all the technicalities related to how rigorously extend the KS test to discrete distributions. We calculate only the test statistic itself, and do not calculate the significance level or p -value. In principle this could be done using the Kolmogorov distribution.

The inversion results of several simulations are shown in Figures 6.11 and 6.12. The dashed lines correspond to uncorrected distributions and the solid ones for inverted results and error bars are 1 sigma (68 CL) uncertainties around the median value, obtained via Monte Carlo runs. The μ -value is varied between 10^{-3} and 30. To demonstrate the $\sqrt{N_E}$ scaling or asymptotic efficiency of the algorithm, we vary the number of simulated events N_E . We observe a clear saturation in the inversion

process at high μ -values, because at that point all event signatures are almost identical. That is, all binary observables are $B_i \equiv 1$ for all $i = 1, \dots, N$. The distributions for the relative components errors

$$\Delta_c \equiv (\hat{p}_c - p_c)/p_c, \tag{6.91}$$

are shown in Figure 6.12 sampled over different μ -value runs. In the histogram legend, the mean square error $\langle (\hat{p}_c - p_c)^2 \rangle$ of the estimates are shown. The estimates are seen to be numerically unbiased, as expected from the construction, thus the mean square error of the estimates is equal to the variance of the estimator. In Figure 6.12 we see also that the relative component error distribution variance is seen to grow as a function of dimension N , demonstrating growing complexity in the combinatorial mixing process. The uncorrected residuals are significantly biased, by the nature of the problem which results in extreme mixing between event topologies. These simplified simulations give also new orthogonal insight into the hardness of understanding soft inclusive QCD processes in detail where mixing may happen at more complicated amplitude level.

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

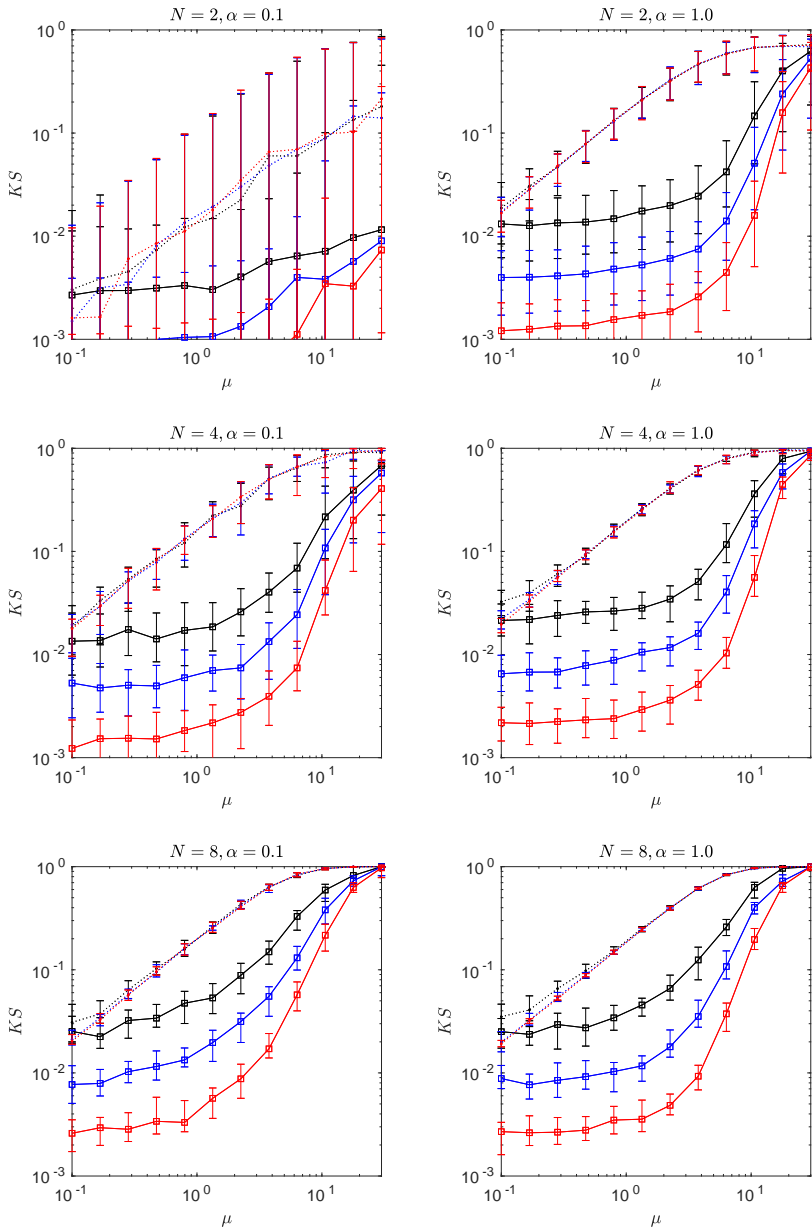


Figure 6.11: Kolmogorov-Smirnov error on the vertical axis. Median values and $\pm 1\sigma$ intervals as a function of μ in different dimensions N . The simulated number of events are $N_E = 10^3, 10^4, 10^5$ counts, denoted with black, blue and red markers. Solid lines are after, and dashed lines before the inversion.

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

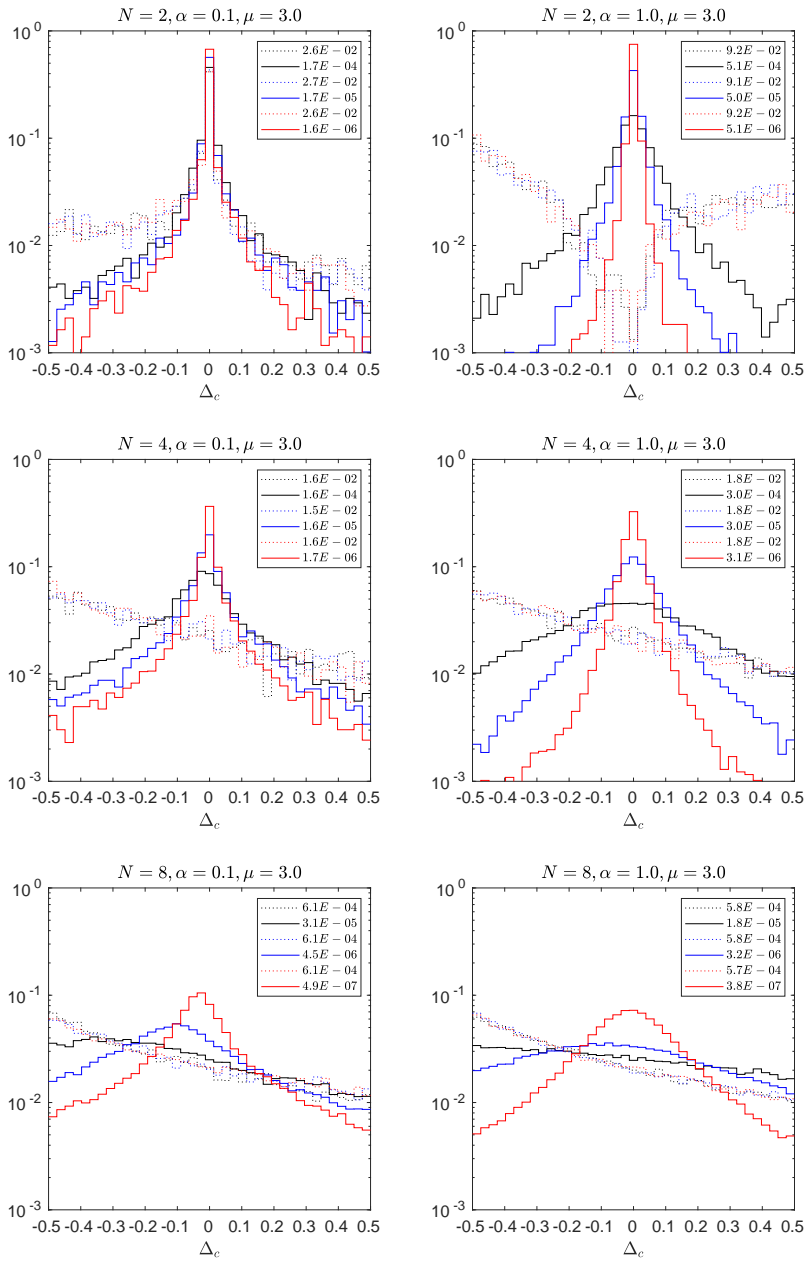


Figure 6.12: Relative component error Δ_c distributions in different dimensions N . The simulated number of events are $N_E = 10^3, 10^4, 10^5$ counts, denoted with black, blue and red markers. Solid lines are after, and dashed lines before the inversion.

CHAPTER 6. COMBINATORIAL SUPERSTATISTICS FOR SOFT QCD

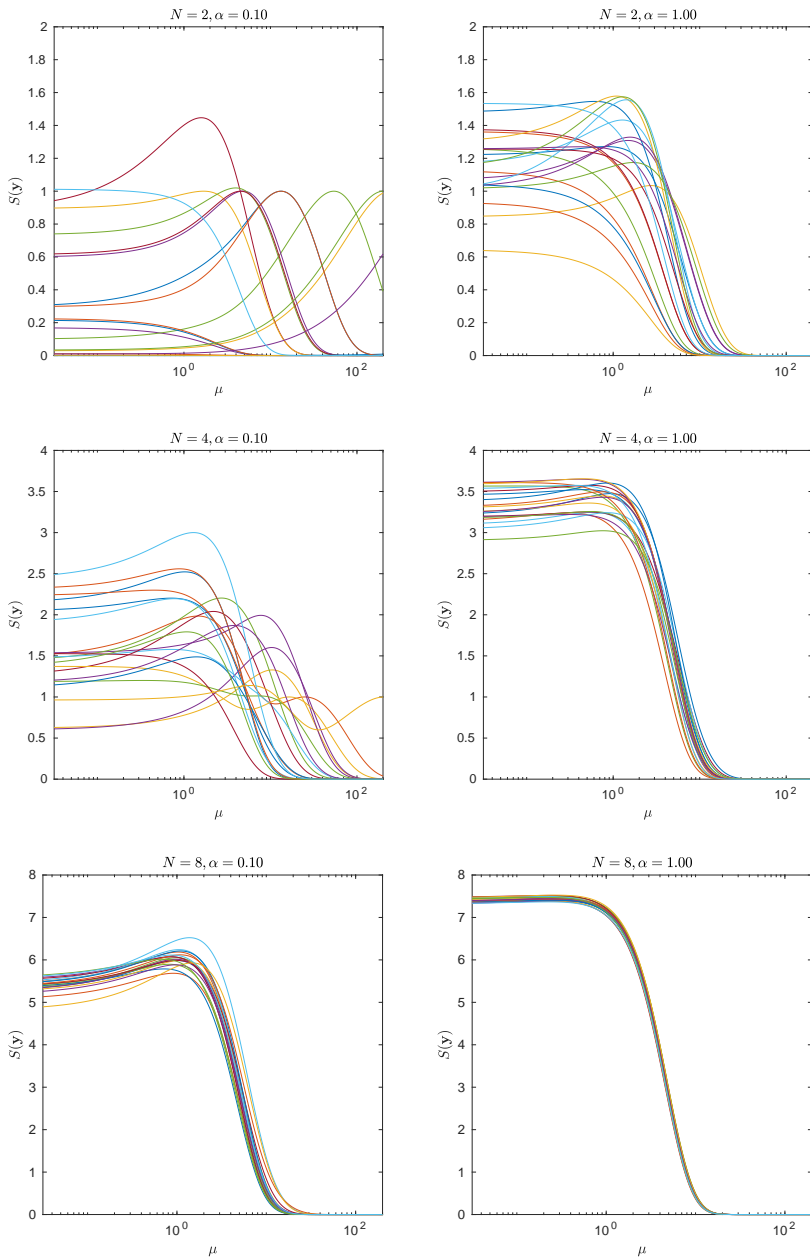


Figure 6.13: Shannon entropy $S(\mathbf{y})$ of the Dirichlet simulation realizations in different dimensions N running as a function of μ , where individual random realizations are visualized with different colors.

6.7 Summary

We built explicitly the following new constructions:

- ◇ A novel algebraic formulation for N -point correlations and complex event ‘topologies’ over rapidity, transverse momentum and multiplicity designed as new observables of high energy soft QCD diffraction. The framework includes vector fiducial partial cross sections, fractal marginal distributions, transverse momentum or multiplicity dependent gap trajectories and multidimensional diffraction fit and projection algorithms.
- ◇ A novel description of combinatorial compound processes based on the incidence algebras and inversion of the corresponding superposition Poisson problem based on the Möbius inversion theorem. The invertibility with finite statistics was simulated.

As a related technical remark regarding the superposition problem, the current event by event pileup correction methods at the LHC are essentially non-linear filtering operations (such as the median filter) applied to the observables event by event. A combination of the statistical approaches developed here and the event by event methods of [179–182] could be possible.

7 On the Inversion of High Energy Proton

Inversion of the K -fold stochastic autoconvolution integral equation is an elementary nonlinear problem, yet there are no de facto methods to solve it with finite statistics. To fix this problem, we introduce a novel inverse algorithm based on a combination of minimization of relative entropy, the Fast Fourier Transform and a recursive version of Efron's bootstrap. This gives us power to obtain new perspectives on non-perturbative high energy QCD, such as probing the ab initio principles underlying the approximately negative binomial distributions of observed charged particle final state multiplicities, related to multiparton interactions, the fluctuating structure and profile of proton and diffraction. As a proof-of-concept, we apply the algorithm to ALICE proton-proton charged particle multiplicity measurements done at different center-of-mass energies and fiducial pseudorapidity intervals at the LHC, available on HEPData. A strong double peak structure emerges from the inversion, barely visible without it.

Chapter in: [arXiv:1905.12585](https://arxiv.org/abs/1905.12585) [hep-ph]

7.1 Introduction

Differential distributions of final state particles in high energy collisions, such as multiplicity distributions, are notoriously difficult to model precisely from the first principles of QCD, due to the nature of confinement and non-perturbative infrared physics. The topic has been studied since the early days of the Hagedorn fireball model [183], Feynman scaling [184] and Kobe-Nielsen-Olesen/Polyakov self-similarity (fractal) scaling [185, 186]. Hadron production has been treated as resulting from a breakdown of classic strings [187], the Feynman-Field model [188] and Lund strings of PYTHIA [189, 190], cluster hadronization of HERWIG [191], the topological expansion of cylindrical pomeron particle production model or ‘quark-gluon strings’ [192], to name some well known concepts. From the perturbative side, the local QCD parton-hadron duality [193] is perhaps the most appealing, with much support already from LEP data. Basically, it states that the hadronization happens at a low virtuality scale independent of the hard scattering scale. However, then the case of pure soft hadron-hadron scatterings without any hard scale involved is very interesting for collectivity or global color coherence reasons. These concepts are in one way or another implemented algorithmically in the Monte Carlo event generators, with a varying number of free parameters to be fixed by data, with goals of factorizing universal, the initial state and center-of-mass energy (in)-dependent properties.

In proton-nucleus (pA) and nucleus-nucleus (AA) physics a whole new class of phenomenological topics appear, from the space-time evolution of relativistic hydrodynamics and medium expansion to quark-gluon plasma signatures and the separation of different stages and collective phases of the process. A very interesting topic is the question which of these are due to heavy ion physics and which of them already happen in the elementary proton-proton interactions. To probe the ab initio physics behind many phenomenological models, and to factorize effects between heavy ions and elementary pp , solid mathematical approaches are crucial as analysis tools.

In this work, we introduce a novel inverse algorithm to invert a certain simple stochastic integral equation, the K -fold autoconvolution. By autoconvolution we mean convolving the distribution with itself and K -fold means repeating the operation recursively K -times. This is a nonlinear operation. In addition, we take the K to be a random number. The ‘direct problem’ and closely related problems are well studied in probability theory and implicitly ubiquitous in high energy physics, given the much studied pQCD logarithmic evolution integro-differential resummation schemes such as DGLAP in $\ln(Q^2)$ and BFKL in $\ln(1/x)$, which require as input the non-perturbative parton densities. The ‘inverse problem’ has been studied much less, considering the variety of methods for standard deconvolution and more generally,

methods to invert approximately linear and nonlinear integral equations with varying kernels. In experimental high energy physics, the unfolding of detector responses is more well known problem [194–197]. It corresponds to the usual deconvolution in the special case where the detector response matrix is of a convolution kernel type. One reason why inverse methods have not been studied extensively is, perhaps, the Monte Carlo event generators, which are extensively used as direct models. The demanding inverse problem, in that case, is the multidimensional parameter estimation or tuning of the event generator itself.

The problem of inverting the pileup effects in inclusive soft QCD multiplicity measurements by inverting a K -fold autoconvolution was treated in [52]. However, the subtractive iterative solution there was proposed without a statistical treatment, explicit regularization or uncertainty estimation and only for a small number K . We address these issues here by developing an all-order inverse. As far as we know, our approach which combines statistical modeling, Fast Fourier Transform (FFT) and Efron’s bootstrap, is the first of its kind in high energy physics, but also including other fields. However, we mention that related problems have been studied also in a different context in laser optics [198], inverse problems and mathematics [199–203], and the problem is also closely related to the inverse Born series problem [204].

In Section 7.2 we introduce the mathematical direct problem and related formalism together with some simple examples and in Section 7.3 and 7.4 we go through the corresponding inverse problem, its discretization, regularization and the algorithmic solution. Section 7.5 is devoted to simulations and finally in Section 7.6 we do a proof-of-concept study of LHC-ALICE proton-proton multiplicity data. In Section 7.7 we discuss further applications and research directions.

7.2 Direct problem

Let the underlying continuous or discrete differential distribution be f_X , with normalization $\int f_X(x) dx = 1$ and the corresponding random variable be X . The distribution can be discrete, for example, in a case of charged particle multiplicity spectrum measurements. The K -fold autoconvolution means that the measurement g_Y is a sum of mutually independent identically distributed (i.i.d) random variables $Y = X_1 + X_2 + \dots + X_K$. If we take $K \sim \text{Poisson}(\mu)$ and all $X_i \sim f_X$, this is known as a *compound Poisson* distribution. The convolution arises here naturally, because the sum of random variables is equivalent to a convolution between the probability distributions of the corresponding random variables.

The autoconvolution operation is defined with

$$[f_X \circledast f_X](y) \equiv \int_{-\infty}^{\infty} f_X(y-x)f_X(x) dx. \quad (7.1)$$

If we take the support of f_X on $[0, \infty)$, because physical observables are usually non-negative, then the autoconvolution integral range lower bound is zero. As a remark, in a different setup, a translation dependent Green's function $G(x, y)$ would be used instead of $f(x-y)$. That would give a Fredholm's integral equation instead of a convolution integral, leading us to the 'inverse Schrödinger equation' type of problems, as they are often called in the field of inverse problems.

For now, let us assume that the measurement follows a compound Poisson. However, in our algorithmic formulation we allow also any discrete distribution for K , not just Poisson distribution. The distribution of random variable Y is now given, by construction, as an autoconvolution series weighted with Poisson probabilities

$$\begin{aligned} g_Y(y) &= P_1 f_X(y) + P_2 [f_X \circledast f_X](y) + P_3 [[f_X \circledast f_X] \circledast f_X](y) + \dots \\ &= \frac{1}{1 - e^{-\mu}} \sum_{n=1}^{\infty} \frac{\mu^n}{n!} e^{-\mu} f_X^{\circledast n}(y), \end{aligned} \quad (7.2)$$

where the *convolution power* \circledast^n is defined recursively as $f^{\circledast n} = f^{\circledast(n-1)} \circledast f$ and $f^{\circledast 1} = f$. We have removed the unobservable case $n = 0$ which gives $Y = 0$ and renormalized the remaining Poisson probabilities $P_n, n = 1, 2, 3, \dots$ to sum to one. The zero suppressed mean of the variable K is

$$\mathbb{E}[K > 0] = \mu / (1 - \exp(-\mu)) > 1, \quad (7.3)$$

where as the mean and variance (second central moment) of the compound distribution are given by useful formulas

$$\mathbb{E}[Y] = \mathbb{E}[K]\mathbb{E}[X] \tag{7.4}$$

$$\text{Var}[Y] = \mathbb{E}[K]\text{Var}[X] + \text{Var}[K](\mathbb{E}[X])^2. \tag{7.5}$$

See [205] for all higher order moments and central moments of generic compound distributions.

In the Fourier spectral domain, the characteristic function (CHF) φ_X is defined as

$$\varphi_X(t) = \mathbb{E}[e^{itX}] = \int_{\mathbb{R}} e^{itX} dF_X(x) = \int_{\mathbb{R}} e^{itX} f_X(x) dx = \int_0^1 e^{itQ_X(p)} dp, \tag{7.6}$$

where $F_X(x)$ and $Q_X(p)$ are the cumulative and the inverse cumulative distribution functions of the random variable $X \in \mathbb{R}$. A probability generating function (PGF) G_X corresponding to a discrete probability mass function $p \sim X \in \mathbb{N}$ is a Laurent series around zero defined as

$$G_X(z) = \mathbb{E}[z^X] = \sum_{n=-\infty}^{\infty} p(n)z^n. \tag{7.7}$$

For the rest of the work we assume that the characteristic and generating functions are defined in the complex plane within their domain of convergence of the corresponding integral and sum, respectively, and leave the extensive algorithmic treatment of possible singularities for future work.

Using these tools, the well-known characteristic function for the compound Poisson is obtained with

$$\begin{aligned} \varphi_Y(t) &= \mathbb{E}[e^{itY}] = \mathbb{E}_K \left[(\mathbb{E}[e^{itX}])^K \right] = \mathbb{E}_K \left[(\varphi_X(t))^K \right] \\ &= \sum_{n=0}^{\infty} \varphi_X(t)^n P(K = n) \\ &= \sum_{n=0}^{\infty} \varphi_X(t)^n \frac{\mu^n}{n!} e^{-\mu} \\ &= e^{\mu(\varphi_X(t)-1)} \end{aligned} \tag{7.8}$$

and solving similarly for the case $K > 0$ ($n = 1, 2, \dots$) gives

$$\varphi_g(t) \equiv \varphi_{Y|K>0}(t) = \frac{e^{-\mu} (e^{\mu\varphi_f(t)} - 1)}{1 - e^{-\mu}} = \frac{1}{e^{\mu} - 1} (e^{\mu\varphi_f(t)} - 1), \tag{7.9}$$

where we denoted the characteristic function of $f_X(t)$ with $\varphi_f(t)$. A direct attempt to invert this would be to take the inverse Fourier transform, but we would encounter the problem of multivalued complex logarithm. Defining the complex logarithm in a suitable way was proposed together with a kernel estimator, for ‘decompounding’ in [202]. Here we take a different approach, where we *avoid* altogether using complex logarithms or other multivalued complex operations such as complex roots.

We will obtain a practical formal notation for the algorithm by representing the problem with an operator $\mathcal{F} : \Omega_X \rightarrow \Omega_Y$, where Ω_X and Ω_Y are the original and the smeared domain

$$\mathcal{F}(f) + \delta g = g. \quad (7.10)$$

Our main goal is to invert this nonlinear mapping taking also into account the statistical fluctuations δg . The nonlinearity comes directly from the fact that the auto-convolution operator \mathcal{F} is constructed from f .

As a side remark, the existence of K -fold convolution is guaranteed by the *infinitely divisibility* of a probability distribution, which is the Lévy-Khintchine theorem. That is, for a given distribution f , there exists for any n , the corresponding n -th convolution power. All members of the stable family - for example Gaussian, Poisson, Negative binomial, Gamma, Cauchy (non-relativistic Breit-Wigner), Lévy and Landau - are infinitely divisible. They are also closed under convolution.

Examples

An interesting phenomenological fact, already observed in the CERN proton-antiproton UA5 data at $\sqrt{s} = 540$ GeV [206, 207], is that the charged particle multiplicity distributions tend to follow approximately a negative binomial distribution (NBD) or two of them, as the classic two component models may suggest. The same observation is still approximately valid at the LHC [208]. Mathematically speaking the NBD follows from a compound Poisson distribution of number of K *logarithmic series* $f(n; p) = p^n / (-n \ln(1-p))$, $n \in \mathbb{N}_+$ i.i.d variables each with shape parameter $p \in (0, 1)$ while K being Poisson distributed with parameter $\mu \in \mathbb{R}_+$. Thus we obtain a negative binomial distribution on \mathbb{N}_+ as the compound distribution with parameters $1-p \in (0, 1)$ and $-\mu / \ln(1-p) \in \mathbb{R}_+$. This mathematical picture may be useful to keep in mind while thinking the possible dynamical explanations.

Monte Carlo models reproduce NBD like distributions by multipomeron cuts or multiparton interaction modeling – the low multiplicity shape of the distribution is usually *very sensitive* to the non-perturbative proton profile or eikonal density driving the number K with impact parameter dependent distributions, and also to the diffraction contribution. Kinematically, a system with the cms energy W will

span a longitudinal rapidity range $Y \sim \ln(W^2/W_0^2)$, where W_0 is at the order of the proton mass. In Lund model like scenarios, also the average number of particles scales like $\langle N \rangle \sim \ln(W^2/W_0^2)$. A multistring system total multiplicity consisting of K -strings will behave as $\langle N \rangle \sim \sum_{i=0}^{K-1} \ln(s_{i,i+1}/W_0^2)$ with $s_{i,i+1} = (k_i + k_{i+1})^2$ being sub-Mandelstam invariants constructed from the parton 4-momenta $\{k_i\}$ spanning the strings and the variance $\text{Var}[N]$ behaves proportionally to $\langle N \rangle$ [209], which is a Poisson like behavior. So if the sub- W^2 distributions behave typically like falling power laws and the number of K -simultaneous strings are also modeled as Poisson like with impact parameter dependence or not, one will generate negative binomial like distributions. Let us emphasize at this point, the soft transverse momentum dependence with multiplicity is not understood from first principles.

Now let us think instead of the experimentally visible superposition of ‘pileup’ proton-proton interactions at the LHC for a given Poisson μ . At the LHC Poisson μ is measured for low values of $\mu \sim 0.05$ (ALICE) using the minimum bias trigger occupancy¹ $\langle R \rangle = 1 - P_0(\mu) = 1 - e^{-\mu} \Leftrightarrow \mu = -\ln(1 - \langle R \rangle)$, where $R \in [0, 1]$. For large values of $\mu \sim 50$ when the occupancy is completely saturated (ATLAS, CMS), the scaling of the track multiplicity or the number of primary vertices is used to determine μ . A combination of both can be useful for intermediate values of $\mu \sim 5$ (LHCb), for example. The pileup is illustrated in Figure 7.1. The convolution of a negative binomial with a negative binomial is yet again a negative binomial – this is the closure under convolution. It is worth pointing out that the estimation of μ based solely on the measured distribution $g(x)$ itself cannot be done in general. This is because if the distribution is from the family of distributions which are closed under convolutions, then we have no capability of identifying if the measured distribution is pileup free, or, one with a larger scale and position parameter than the expected one, or the autoconvoluted version of one with smaller scale and position parameter. Thus, μ must be a parameter of the inversion and must be inferred from other measurements or theoretical constructions.

Clearly, our discussion is now very close to the KNO / Polyakov scaling [186], which states that there is one elegant scaling function ψ giving the probability of multiplicity n as a function of the scaling variable $n/\langle n(s) \rangle$

$$P(n; s) = \frac{1}{\langle n(s) \rangle} \psi \left(\frac{n}{\langle n(s) \rangle} \right), \quad (7.11)$$

where s is the center of mass energy squared Mandelstam variable, a Lorentz scalar. However, this scaling is known to be violated at some level since the ISR times, and

¹ $\langle R \rangle_{t, N_B} = \frac{\frac{1}{\Delta t} \int_{\Delta t} f_R(t) dt}{N_B f_O}$, where N_B is the number of colliding bunch pairs, f_O the LHC revolution frequency (Hz) and $f_R(t)$ the instantaneous trigger rate (Hz).

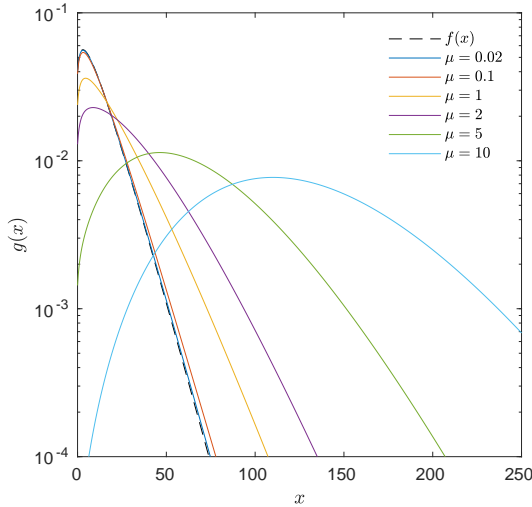


Figure 7.1: K -fold compound Poisson autoconvolution of the negative binomial distribution with varying μ .

at the LHC it holds merely in limited regions of the full phase space, in a narrow central rapidity window. In general, when inspecting scaling or its violation one needs to take into account the collision initial state, e.g. pp versus e^+e^- , the fiducial phase space of the measurement and contribution of diffraction. At the model level, multiparton or multipomeron exchanges running as a function of energy are usually creating significant deviations from the KNO scaling. Additional effects arise from soft versus hard scale interactions. See the books [210, 211] for a discussion on a variety of perturbative QCD and other scenarios exhibiting scaling.

By the central limit theorem, the K -fold convolution will map almost² every distribution eventually into a Gaussian distribution when $\mathbb{E}[K] \rightarrow \infty$. This is clearly visible in Figure 7.1, where we used the direct mapping of Equation 7.2 with different Poisson μ -values to convolve the negative binomial distribution

$$P_{\text{NBD}}(n; \langle n \rangle, k) = \frac{\Gamma(k+n)}{\Gamma(k)\Gamma(n+1)} \left[\frac{\langle n \rangle}{k + \langle n \rangle} \right]^n \left[\frac{k}{k + \langle n \rangle} \right]^k \quad (7.12)$$

with a fit to the LHC charged particle multiplicity data in proton-proton at $\sqrt{s} = 7$ TeV for pseudorapidity range $\eta \in [-0.9, 0.9]$ obtained from [208]. The fit parameters

²For example Cauchy / non-relativistic Breit-Wigner type distributions do not turn into Gaussian, also moments for these distributions are undefined.

central values are $\langle n \rangle = 12.5$ and $k = 1.4$, where $\langle n \rangle$ is the average multiplicity of NBD and k is related to the variance $D^2 = \langle n^2 \rangle - \langle n \rangle$ by $D^2/\langle n \rangle^2 = 1/\langle n \rangle + 1/k$.

In an explicit experimental realization, to take into account the detector response and its effects on the resolution and efficiency, an unfolding algorithm with uncertainty estimation and other corrections should be processed first – if necessary. The output of this can be then propagated to the algorithm described here. For the rest of the work, we assume this to be the case.

7.3 Inverse problem

Solving the K -fold autodeconvolution means here a statistical inverse of Equation 7.10. Unfortunately, the measurement noise and modeling errors will usually always forbid us executing a direct inversion in practice, even if we would have a closed-form expression to do so. This ill-posedness requires usually either implicit or explicit regularization which turns the problem into a more well-posed one.

We calculate the convolution by pointwise multiplication in the Fourier domain using the convolution theorem $f \circledast g = \mathfrak{F}[f]\mathfrak{F}[g]$, which works also the other way around $\mathfrak{F}[f] \circledast \mathfrak{F}[g] = fg$, which is also often called ambiguously convolution or folding of distributions of f and g . Now, a fixed $K \in \mathbb{N}$ autoconvolution is given by

$$\mathfrak{F}[f^{\circledast K}] = (\mathfrak{F}[f])^K = \mathfrak{F}[g], \quad (7.13)$$

where the power is applied as $z^K = (|z|e^{i(\varphi+2\pi n)})^K = |z|^K e^{iK\varphi} e^{2\pi inK}$, $n \in \mathbb{Z}$. Thus, in the spectral domain the stochastic autoconvolution results in a random scaling of amplitudes and a rotation of phases of the characteristic function. Now when K is not an integer this results in a multi-valued operation, because $e^{2\pi inK}$ does not map a full rotation around the complex plane for different values of n . As a reminder, if we shift the probability density $f(x-x_0)$, this results in a phase shift in the frequency domain by $\exp(-i\omega x_0)\mathfrak{F}[f]$.

A formal ‘naive’ inverse is obtained by taking the complex K -th root in the Fourier domain, and then proceeding with an inverse Fourier transform

$$(\mathfrak{F}[f]^K)^{1/K} = (\mathfrak{F}[g])^{1/K} \quad (7.14)$$

$$f = \mathfrak{F}^{-1} \left[(\mathfrak{F}[g])^{1/K} \right], \quad (7.15)$$

analogously to the standard naive Fourier deconvolution. However, complex root is a multivalued operation $|z|^{1/K} e^{i(\text{Arg}(z)+2\pi n)/K}$ with $n = 0, 1, \dots, K-1$ roots with suitable branch cuts to be specified algorithmically, in principle. Now we see that the only strictly well posed case is the case when the characteristic function is real and non-negative \leftrightarrow a mirror symmetric (even) probability density around zero by the properties of the Fourier transform, a case which is clearly too restrictive to be of interest here. Also, this ‘solution’ does not take into account that K is a random variable, neither it takes into account the finite statistics of g and mismatches between the measurement and modeling causing instabilities to be regulated. For some related discussion, see for example [212].

Discretization

To solve the finite sample version of the problem, we need a (discrete) representation of the problem. For simplicity, computational efficiency and due to the high energy physics analysis convention – as a practical format of the measurement, we use histograms as the density estimators of f and g , even if simple kernel (Parzen) estimators and even more advanced density estimation techniques could be used in principle. Also, we consider only a 1D-version of the problem. However, extensions to higher dimensions are formally straightforward.

The histograms are defined in terms of D, KD bins with corresponding fixed bin width $\Delta x \equiv \Delta y$. In the context of algorithms histograms are represented with finite non-negative column vectors, which we denote with bold symbols

$$\mathbf{f}[n] = |\{f_i \in [x_n, x_n + \Delta x)\}|, \quad n \in \mathbb{Z}_{\Omega_X} := \{0, 1, \dots, D - 1\} \quad (7.16)$$

and similarly for a K -fold autoconvoluted

$$\mathbf{g}[n] = |\{g_i \in [y_n, y_n + \Delta y)\}|, \quad n \in \mathbb{Z}_{\Omega_Y} := \{0, 1, \dots, K(D - 1)\}, \quad (7.17)$$

where $f_i, g_i \in \mathbb{R}$ denote sample elements and x_n, y_n denote the bin lower edges. We extended the domain of autoconvoluted by K -times and because K is a random variable, in practice we take some large enough integer which avoids any truncation problems in the upper tail. Only values of $\{g_i\}$ are directly observable and measured. Now the discrete autoconvolution is defined as

$$[\mathbf{f} \circledast \mathbf{f}][n] = \sum_{m=0}^{2(D-1)} \mathbf{f}[m] \mathbf{f}[n - m], \quad (7.18)$$

and the K -fold version follows recursively. In the algorithmic implementations, when necessary to extend the domain of vectors due to convolution, we do it explicitly by zero padding.

In principle, the binning can be also non-uniform, such as logarithmic. However, that must be taken explicitly into account in the evaluation of the discrete convolution, which is defined using uniform sampling. One classic solution could be interpolation and resampling schemes with B-splines, for example. Here, however, we use only fixed binning. The bin width Δx needs to be chosen such that the spectrum resolution criteria and event count statistics are taken into account. This in order not the create a ‘noisy spectrum’ with low bin counts, or on the other hand, induce a loss of resolution. Thus, it has also a role as an implicit regulator.

Relative entropy minimization

The measured histogram \mathbf{g} is assumed to be bin-by-bin Poisson distributed, a usual assumption, given also the Poisson superposition property $\sum_i \text{Poisson}(\mu_i) \sim \text{Poisson}(\sum_i \mu_i)$. The Poisson likelihood with a known constant background histogram \mathbf{b} is then written as a product over the histogram bins

$$p(\mathbf{g}|\mathbf{f}) = \prod_i \frac{[\mathcal{F}\mathbf{f} + \mathbf{b}]_i^{g_i} e^{-[\mathcal{F}\mathbf{f} + \mathbf{b}]_i}}{g_i!} \quad (7.19)$$

and its negative log-likelihood is now our fit quality term

$$J(\mathbf{f}|\mathbf{g}) = -\log p(\mathbf{g}|\mathbf{f}) = \sum_i -g_i \log([\mathcal{F}\mathbf{f} + \mathbf{b}]_i) + [\mathcal{F}\mathbf{f}]_i + b_i + \log(g_i!) \quad (7.20)$$

and by neglecting terms which do not affect the solution and rearranging gives

$$J(\mathbf{f}|\mathbf{g}) = \sum_i g_i \frac{g_i}{\log([\mathcal{F}\mathbf{f} + \mathbf{b}]_i)} + [\mathcal{F}\mathbf{f}]_i - g_i. \quad (7.21)$$

This often used formulation is the minimum entropy solution, and is equivalent to minimizing generalized Kullback-Leibler divergence (relative entropy) [213], known also as the Csiszár I-divergence [214]. The gradient of this functional is

$$\nabla J(\mathbf{f}) = \frac{\partial}{\partial \mathbf{f}} (-\mathbf{g}^T \log(\mathcal{F}\mathbf{f} + \mathbf{b}) + \mathbf{1}^T \mathcal{F}\mathbf{f}) = \mathcal{F}^T \mathbf{1} - \mathcal{F}^T \frac{\mathbf{g}}{\mathcal{F}\mathbf{f} + \mathbf{b}} = \mathcal{F}^T \left(\mathbf{1} - \frac{\mathbf{g}}{\mathcal{F}\mathbf{f} + \mathbf{b}} \right), \quad (7.22)$$

where divisions are understood component wise. If distributions are normalized to unity, then $\mathcal{F}^T \mathbf{1} \equiv \mathbf{1}$ holds. Here, we use event counts.

The gradient based iterative update or a fixed point iteration is now obtained from the steepest descent update $\mathbf{f}_{k+1} = \mathbf{f}_k - \gamma D_k \nabla J(\mathbf{f}_k)$, where D_k is a gradient scaling matrix. If we set $D_{\mathbf{f}_k} = \text{diag}(\mathbf{f}_k)$ and $\gamma = 1$ which is well known to correspond in this case to the expectation maximization (EM) [215] formulation of a frequentist maximum likelihood solution under the Poisson likelihood, then we obtain a Richardson-Lucy [216, 217] type multiplicative deconvolution formula well known in optics and astrophysics

$$\mathbf{u}_k = (\mathbf{f}_k \oslash (\mathcal{F}^T \mathbf{1})) \odot (\mathcal{F}^T \mathbf{g} \oslash (\mathcal{F}\mathbf{f}_k + \mathbf{b})), \quad (7.23)$$

where \odot and \oslash are Hadamard's vector component wise product and division, respectively. This was also re-invented by D'Agostini [197] in high energy physics

unfolding context. For the rest of the paper, we set the background $\mathbf{b} = 0$. The RL-type update conserves non-negativity of the solution and the total number of events. Another way to derive the RL type update is to set gradient Equation 7.22 to zero, move terms on both sides, multiply by \mathbf{f} and do fixed point iteration.

Regularization

The problem of choosing the spectral bandwidth cut-off or regularization strength in the non-transformed domain are the most common but also the most difficult topics of inverse problems. Motivated by the fact that the distributions of observables for us are usually smooth and non-discontinuous, we use a traditional variational regularization with a Tikhonov ℓ_2 -norm type regularity functional

$$J_R(f) = \lambda \int_{\Omega} \|\nabla f(x)\|^2 dx, \quad (7.24)$$

where $\Omega \subset \mathbb{R}$. The minimum of this is given by

$$\begin{aligned} \nabla J_R(f) &= \lambda \nabla \left(\int_{\Omega} \|\nabla f(x)\|^2 dx \right) \\ &= \lambda \nabla^* \nabla f(x) = -\lambda \operatorname{div}(\nabla f(x)) = -\lambda \nabla^2 f(x) = 0, \end{aligned} \quad (7.25)$$

where we applied to the integral von Neumann boundary conditions and Gauss divergence-theorem, the standard vector analysis identities and ∇^2 is the usual Laplacian. That is, the gradient descent direction is given by the negative Laplacian.

Instead of using the abstract functional gradient, the regularization term is discretized using a finite difference matrix and the corresponding finite Laplacian

$$\nabla_M \equiv \begin{pmatrix} 1 & -1 & 0 & \cdots & 0 \\ 0 & 1 & -1 & & \vdots \\ \vdots & & & \ddots & 0 \\ 0 & \cdots & 0 & 1 & -1 \end{pmatrix}, \quad \nabla_M^2 \equiv \nabla_M^T \nabla_M = \begin{pmatrix} -1 & 2 & -1 & \cdots & 0 \\ 0 & -1 & 2 & -1 & \vdots \\ \vdots & & & \ddots & 0 \\ 0 & \cdots & -1 & 2 & -1 \end{pmatrix}. \quad (7.26)$$

It is naturally also possible to modify the boundary conditions which affect these matrices, based on a priori assumed or known properties of the distribution $f(x)$.

Regularization parameter λ depends on the number of observed events count N , the number of histogram bins D ($=$ problem discretization) and the shape of the distribution of interest. The average pileup μ and the functional smoothness of the distribution under inversion are very important factors affecting the nonlinear direct

operator and thus the necessary regularization strength. In the limit $N_E \rightarrow \infty$ and $\mu \rightarrow 0$, we can take asymptotically $\lambda \rightarrow 0$. There are several semi-heuristic ‘data-driven’ methods developed for selecting the regularization. The well-known principles are the L-curve, generalized cross validation (GCV), Morozov’s discrepancy principle and various covariance and residual spectrum analysis methods [218]. Monte Carlo event generator driven or a toy model based analysis together with data driven approaches is also a reasonable choice in high energy physics. We use a simple variational equilibrium approach to choose the regularization parameter. This is described in Section 7.5.

The full solution taking together fidelity + regularity usually means directly minimizing in the direction of combined gradient of the full cost

$$\min_{\mathbf{f}} J(\mathbf{f}|\mathbf{g}) + \lambda J_R(\mathbf{f}) \quad (7.27)$$

as implemented for example in [219] in the context of image deconvolution. However, this we observed to give very unstable results in this problem. Thus, we implemented the regularization step through forward-backward slitting type gradient update, similar in fashion to proximal optimization algorithms. It means that we update the result after the EM-update step as

$$\mathbf{f}_{k+1} = \mathcal{P}_+(\mathbf{u}_k - \lambda \nabla_M^2 \mathbf{u}_k), \quad (7.28)$$

where $\mathcal{P}_+(\cdot)$ is the positive solution projector operator $\mathcal{P}_+ : \mathbb{R}^n \rightarrow \mathbb{R}_+^n$ such that $\mathcal{P}_+(\mathbf{x}) = \mathbf{y}$ gives $y_i = x_i$ if $x_i > 0$, else, $y_i = 0$, $\forall i = 1, \dots, n$. Separating the regularization as a separate step resulted in a much improved behavior. However, the approach chosen here is the best performing we found from the vast phase space of optimization algorithms and regularization approaches. One must remember that we are not dealing here with an inverse problem with a fixed kernel, thus this problem has its own peculiarities.

7.4 Algorithm

For every iteration, we construct the autoconvolution operator by taking the discrete Fourier transform using FFT as $F = \mathfrak{F}[\mathbf{f}]$ and then construct the empirical characteristic function of the projected measurement with a nonlinear complex map $G = \frac{1}{e^{\mu}-1} (e^{\mu F} - 1)$ using Equation 7.9, where complex exponential is the single valued function $\exp(z) = \sum_{n \geq 0} z^n / n!$, $z \in \mathbb{C}$. The autoconvolution ‘kernel’ in the spectral domain is then obtained by taking a point wise division $H = G \oslash F$, inverse transforming $\mathbf{h} = \mathfrak{F}^{-1}[H]$ and finally representing it as a Toeplitz convolution matrix to obtain $\mathcal{F} = \mathcal{T}[\mathbf{h}]$. The operator \mathcal{T} represents here a simple standard way to construct a Toeplitz matrix. We assume F to be non-zero for all of its components, which might be violated in pathological cases. These are intrinsically non-invertible without extra a priori information.

We remark here that this nonlinear map and division in the spectral domain has a different goal than in the usual Fourier domain linear deconvolution such as in Wiener filters [220], where the noise amplification is regularized in the spectral domain. In Wiener filtering, the division gives a solution to the problem non-recursively. Here, on the other hand, we use the spectral domain as a way to obtain the Toeplitz matrix representation in the probability domain extremely efficiently via FFT. The full solution is recursive and regularized in the probability domain.

Discretization of the operator \mathcal{F} in the case of arbitrary (non-Poisson) compound distribution is done with a finite number of convolution Toeplitz matrices $F^{(n)}$, $n = 1, \dots, \tilde{K}$ with $F^{(1)} = I$ up to a numerically suitable finite order \tilde{K} . The maximum order is basically limited only by the available computing resources. These Toeplitz matrices are then added together with weights P_n . That is, we evaluate the discretized and truncated version of Equation 7.2. This numerical finite order implementation was cross checked against the exact all order Fourier domain method in the Poisson case, as described above, and they agreed within floating point accuracy.

Now summarizing: the multiplicative algorithm describe above corresponds to a non-negative solution with Poisson bin fluctuations likelihood with the generalized Kullback-Leibler divergence being the fit criteria. On the other hand, subtractive algorithms are usually encountered with iterative solutions to unconstrained optimization under Gaussian noise and ℓ_2 -minimum norm minimizing the least squares cost $J_{\text{LS}}(\mathbf{f}) = \frac{1}{2} \|\mathcal{F}\mathbf{f} - \mathbf{g}\|^2$ with an iterative solution

$$\mathbf{f}_{k+1} = \mathbf{f}_k - \gamma \nabla J_{\text{LS}} = \mathbf{f}_k - \gamma \mathcal{F}^T (\mathcal{F}\mathbf{f}_k - \mathbf{g}) \quad (7.29)$$

which is known as the Landweber iteration scheme [221]. Non-negativity of the solution should be enforced simply with the projector $\mathcal{P}_+(\cdot)$, because it is not a

Algorithm 1 KISU: K-fold Inverse of Stochastic Autoconvolution

INPUT: Smearred histogram \mathbf{g} with N events, regularization strength $\lambda \in \mathbb{R}_+$, Poisson mean $\mu \in \mathbb{R}_+$ **OR** discrete distribution P_n with $\sum_{n=1}^{\tilde{K}} P_n = 1$, number of iterations $R \in \mathbb{N}_+$, background histogram \mathbf{b} or otherwise $\mathbf{b} = 0$.

If μ given, use option **§I** below, **else**, use input P_n and option **§II** below

Initialize $\mathbf{f}_1 \leftarrow \mathcal{Z}_{\tilde{K}}(\mathbf{g} - \mathbf{b})$ with \tilde{K} -fold zero padding $\mathcal{Z}_K(\cdot)$

for all $k = 1, \dots, R$ **do**

Step 1. Construction of the direct operator:

§I: Poisson case; ‘all-order’ map:

$$F \leftarrow \mathfrak{F}[\mathbf{f}_k], G \leftarrow \frac{1}{e^{\mu} - 1} (e^{\mu F} - 1), H \leftarrow G \otimes F, \mathbf{h} \leftarrow \mathfrak{F}^{-1}[H], \mathcal{F}_k \leftarrow \mathcal{T}[\mathbf{h}]$$

§II: Generic P_n case; evaluated ‘order-by-order’:

for all $n = 2, \dots, \tilde{K}$ **do**

$$F_k^{(n)} \leftarrow \mathcal{T}[\mathbf{f}_k^{\otimes n-1}] \quad \# \text{ Construct convolution Toeplitz matrix}$$

$$\mathbf{f}_k^{\otimes n} \leftarrow \mathbf{f}_k^{\otimes n-1} \otimes \mathbf{f}_k \quad \# \text{ Calculate convolution power}$$

end for

$$\mathcal{F}_k \leftarrow \sum_{n=1}^{\tilde{K}} P_n F_k^{(n)}, \text{ where } F_k^{(1)} = I \quad \# \text{ Weighted sum of convolution matrices}$$

Step 2. EM-step + regularization/smoothing with positivity constraint:

$$\mathbf{u}_k \leftarrow (\mathbf{f}_k \otimes (\mathcal{F}^T \mathbf{1})) \odot (\mathcal{F}^T \mathbf{g} \otimes (\mathcal{F} \mathbf{f}_k + \mathbf{b}))$$

$$\mathbf{f}_{k+1} \leftarrow \mathcal{P}_+(\mathbf{u}_k - \lambda \nabla_M^2 \mathbf{u}_k)$$

end for

OUTPUT: Deconvolution estimate $\hat{\mathbf{f}} \leftarrow \mathcal{Z}_{\tilde{K}}^{-1}(\mathbf{f}_R)$ with the same support as \mathbf{g} .

built-in constraint of the standard least squares. The γ is a relaxation parameter that controls the convergence and can be optimized by several different line search methods or by fixing it to a constant ~ 1 , which gives (much) slower convergence. This Gaussian version of the algorithm could be utilized with large event samples. With low event count histograms, the multiplicative solution was observed to be superior in terms of the stability of the solutions.

Uncertainty estimation

Point estimates are obtained with Algorithm 1, which we call KISU. The algorithm first solves the direct problem and then one EM-step plus regularization of the inverse, and recursively alternate between these two until convergence. The regularization is done explicitly, because it is well known that the EM-type maximum likelihood solution alone will amplify the high frequencies (\sim Poisson noise) when the number of iterations $k \rightarrow \infty$. That is, iteration first fits the low frequency Fourier components and later starts to fit the high frequency noise to the solution. A semi-explicit regularization strategy would correspond to an early iteration stop. However, with an explicit regularization scheme as here, the iteration can be allowed to continue till convergence within some numerical threshold.

The iterative bias estimation and uncertainty estimation are done using numerical Monte Carlo bootstrap with Algorithms 2 and 3. We call these the daughter and mother bootstrap, respectively. The algorithm is started by calling the mother bootstrap, which calls the daughter bootstrap, which calls KISU. The bootstrap is selected here due to the nonlinear and recursive nature of the problem which makes the usual linearized Taylor expansion error propagation and analytical (Gaussian) approximations unreliable or semi-impossible. The bootstrap was introduced in the seminal paper by Efron in 1979 [222] and the recursive bias correction was first discussed in 1986 [223]. In high energy physics unfolding context, bootstrap iterated bias correction and related confidence intervals were utilized only quite recently in [195] and we follow a similar strategy. The underlying inverse problem is nonlinear here, in contrast. The bootstrap is a relatively easy method to implement and variants of it have already long been used in high energy physics dubbed often under the large umbrella of ‘toy Monte Carlo’. A practical problem is the high computational cost, fortunately, bootstrap sampling is trivially parallelizable. However Algorithm 1 and iterative bias correction loop of Algorithm 2 are recursive, and thus cannot be made fully parallel.

The bias of a parameter estimator is by definition $\text{Bias}[\hat{\theta}] = \mathbb{E}[\hat{\theta}] - \theta$. An iterative bootstrap estimate of the bias replaces the unknown expectation with a

Algorithm 2 ‘Daughter bootstrap’ – Iterative Bias Correction

INPUT: Smearred histogram \mathbf{g} , number of bias correction iterations $M \in \mathbb{N}_+$, bootstrap sample size per iteration $B \in \mathbb{N}_+$ and Algorithm **A1** parameters.

Run **A1** using \mathbf{g} as input, obtain $\hat{\mathbf{f}}_0$
 Set $\hat{\mathbf{f}}_\star \leftarrow \hat{\mathbf{f}}_0$, $\hat{\mathbf{g}}_\star \leftarrow \mathbf{g}$
for all $i = 1, \dots, M$ **do**
 Set $S_i \leftarrow \emptyset$
 for all $j = 1, \dots, B$ **do**
 Draw $\mathbf{g}_j^* \sim \hat{\mathbf{g}}_\star$ with replacement using toy Monte Carlo
 Run **A1** using \mathbf{g}_j^* as input, obtain \mathbf{f}_j^* , update $S_i \leftarrow S_i \cup \{\mathbf{f}_j^*\}$
 end for
 Step 1. Iterated bias estimate using bootstrap sample S_i mean:
 $\widehat{\text{Bias}}[\hat{\mathbf{f}}] \leftarrow \text{Mean}[S_i] - \hat{\mathbf{f}}_\star$
 Step 2. Point estimate with bias correction, and non-negativity re-enforced:
 $\hat{\mathbf{f}}_\star \leftarrow \mathcal{P}_+(\hat{\mathbf{f}}_0 - \widehat{\text{Bias}}[\hat{\mathbf{f}}])$
 Step 3. Using direct operator estimation **Step 1.** of **A1**, generate:
 $\hat{\mathbf{g}}_\star \leftarrow \mathcal{F}_\star(\hat{\mathbf{f}}_\star)$
end for

OUTPUT: Bias corrected estimate $\hat{\mathbf{f}}_\star$, first estimate $\hat{\mathbf{f}}_0$ and the sample S_B for the standard bootstrap confidence interval evaluation.

Algorithm 3 ‘Mother bootstrap’

INPUT: Smearred histogram \mathbf{g} , bootstrap sample size $Q \in \mathbb{N}_+$, Algorithm **A2** and **A1** parameters.

Set $S_\star \leftarrow \emptyset$, $S_0 \leftarrow \emptyset$
for all $q = 1, \dots, Q$ **do**
 Draw $\mathbf{g}_q^* \sim \mathbf{g}$ with replacement using toy Monte Carlo
 Run **A1** using \mathbf{g}_q^* as input, obtain $\hat{\mathbf{f}}_{\star q}^*$ and $\hat{\mathbf{f}}_{0q}^*$, update $S_\star \leftarrow S_\star \cup \{\hat{\mathbf{f}}_{\star q}^*\}$, $S_0 \leftarrow S_0 \cup \{\hat{\mathbf{f}}_{0q}^*\}$
end for

OUTPUT: Samples S_\star and S_0 for calculating the confidence intervals, median etc., for both bias corrected (\star) and uncorrected (0) estimates.

bootstrap sample mean and the true parameter value with the current estimate of the parameter. This is described in algorithm 2. However, the variance of the estimator $\text{var}[\hat{\theta}] = \mathbb{E}[(\hat{\theta}) - E[\hat{\theta}]]^2$ can grow as a result of the bias correction. That is, thinking in terms of classic but non-robust mean squared error $\text{MSE}[\hat{\theta}] = \mathbb{E}[(\hat{\theta} - \theta)^2] = \text{var}(\hat{\theta}) + (\text{Bias}[\hat{\theta}])^2$, it is clear that a good estimator is a combination of both qualities. The classic goal has usually been to find out the minimum variance unbiased estimator (MVUE), which clearly minimizes MSE among the unbiased estimators. For more information about the iterative bootstrap, we refer the reader to the book by Hall [224].

We calculate the uncertainty estimates by ordering the sample obtained from the mother algorithm and calculate the vector component (pointwise) $1 - 2\alpha$ *percentile intervals* at level α as $[\hat{\mathbf{f}}_{\star/0,\alpha}^*, \hat{\mathbf{f}}_{\star/0,1-\alpha}^*]$, where $\star/0$ means we consider both bias corrected (\star) and uncorrected (0) estimates and their intervals. As another option, the so-called *basic bootstrap intervals* would be obtained with $[2\hat{\mathbf{f}} - \hat{\mathbf{f}}_{\star/0,1-\alpha}^*, 2\hat{\mathbf{f}} - \hat{\mathbf{f}}_{\star/0,\alpha}^*]$, where $\hat{\mathbf{f}}$ is the estimate obtained directly running KISU. Basic bootstrap intervals have clearly a possibility of flipping the intervals upside down, which we observed also in the simulations. One could also go further and estimate the spectrum global uncertainty coverage, not just point by point local one. For more information, see [225].

7.5 Simulations

The first scenario to study is the simplest fixed K autoconvolution inverse of the negative exponential distribution without any uncertainty estimation. We illustrate this in Figure 7.2 and compare with the naive Fourier domain inverse with spectral cut-off regularization, that is, we set high frequencies above the cut-off Λ_ω to zero, with $\omega \in [0, 1]$. The spectral cut-off is tuned to give the best performance for this scenario. We observe nearly perfect reconstruction with KISU and very large oscillations with the naive Fourier domain algorithm. These oscillations seemed to appear immediately after a small amount of counting fluctuations in the observed spectrum g , which we simply added here as Gaussian noise. For a typical iteration trajectories, see Figure 7.3. We see that the variational regularity cost is kept as almost constant, and the error to the ground truth and the re-projection error behave in the same way, which is always called for behavior.

We tried to two different ‘data-driven’ ways of choosing the regularization parameter, see Figure 7.4. First, the regularization parameter λ was selected as the equilibrium point $\min(\alpha + \beta)$, where the re-projection error is $\alpha = \chi_{\hat{g}}^2 = \|(\mathbf{g} - \hat{\mathbf{g}})/\sqrt{\mathbf{g}}\|_{\ell_2}^2$ and the regularity cost is $\beta = \chi_{\nabla^2 \hat{f}}^2 = \|\nabla^2 \hat{\mathbf{f}}/\sqrt{\hat{\mathbf{f}}}\|_{\ell_2}^2$, with operations taken element wise. Re-projection is simply defined as $\hat{g} = \hat{\mathcal{F}}(\hat{f})$, available after the inversion. The

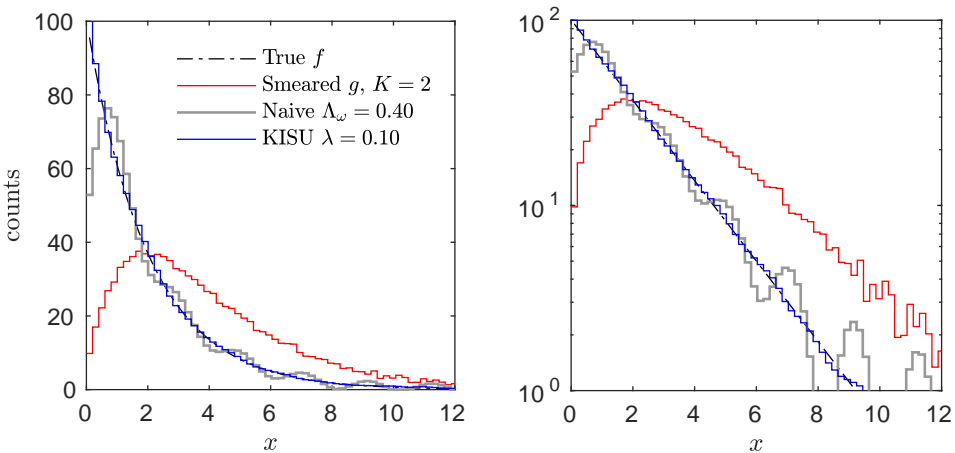


Figure 7.2: Simplified fixed K autoconvolution inverse using the naive Fourier domain inverse with spectral cut-off and KISU based inverse. The true distribution is $f(x) = 1/\alpha \exp(-x/\alpha)$ with $\alpha = 2$.

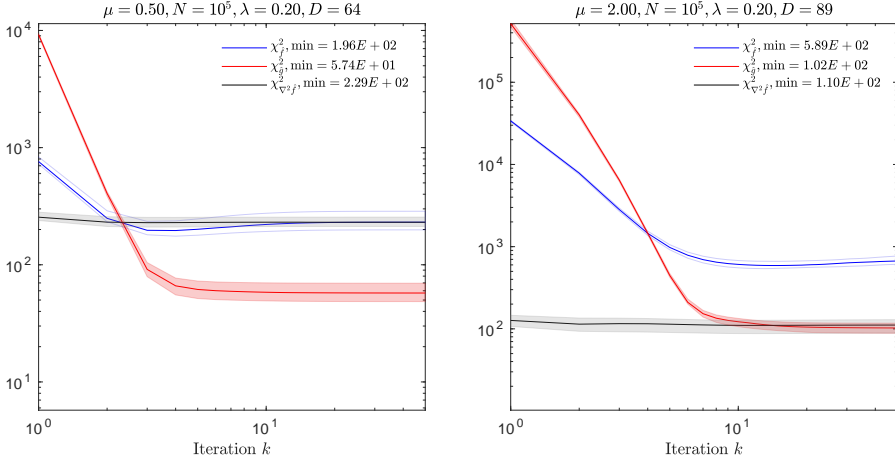


Figure 7.3: Typical KISU algorithm iteration trajectories. Red lines denote the re-projection error, blue lines are the error with respect to the true distribution (available only in simulation) and black lines denote the variational regularity (smoothness).

values for different λ values are obtained via brute force loop. This was repeated for each simulation scenario separately. The second criteria was to use the Morozov like discrepancy principle point where the re-projection error α is approximately the same as the number of histogram bins D . For the rest of the simulations, we use the equilibrium solution. When $\lambda \rightarrow 1$, the solution becomes over smooth which results in a loss of high frequency details and when $\lambda \rightarrow 0$, we start to fit the noise in high frequencies. In general, slightly larger λ values were always obtained as the equilibrium solution compared to the discrepancy principle.

In more demanding simulations we have a scenario that corresponds approximately to inverting minimum bias pileup at the LHC. We chose fully analytical distributions as the ground truth f and draw samples via von Neumann acceptance-rejection Monte Carlo sampling. The distribution is a negative binomial with parameters given in Figure 7.5 and exponential in Figure 7.6. The number of events N and the Poisson μ were varied, with the histogram bin size Δx kept approximately fixed. One must remember that the overall statistics scales also with the μ -value, so there must be an optimal point with respect to the pileup deterioration and the collected number of events.

The convergence of the iterative algorithms was obtained in every scenario, and the fundamental limitations were hit when the smeared distribution g was almost

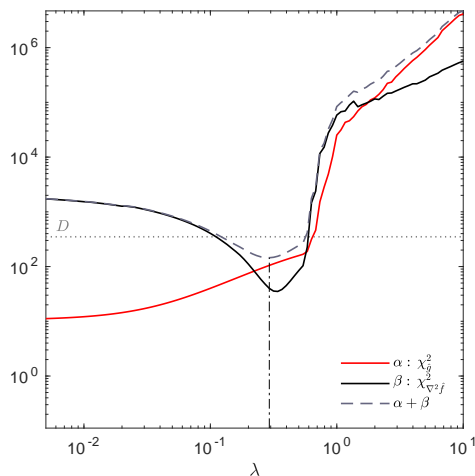


Figure 7.4: Data driven equilibrium between the re-projection cost α and the variational regularity cost β . The horizontal line denotes the number of bins D in the measured histogram \sim number of degrees of freedom.

purely Gaussian, which was the case when $\mu \sim 10$ or more. In that case, the inverted distribution was lacking any fine structure. Frequentist pointwise uncertainty bands ($\alpha = 0.05$) \sim 95 CL were obtained from the bootstrap procedures. The coverage is heuristically reasonable when comparing with respect to the truth, however, the nonlinear oscillation is not completely covered by the bootstrap. A 3σ signal to uncertainty ratio, shown in each figure, seems to be a reasonable safe threshold cut.

In Figure 7.5, we observe some oscillation or ‘nucleation’ of certain eigenmodes in the solution, which is not an unusual phenomenon in inverse problems. We see that the bias correction seems to increase the variance of the estimates slightly, which is visible in bin-by-bin fluctuations. This is expected due to the bias-variance tradeoff. The optimal number of bias iterations should be studied further, because it is a free parameter unless taken till convergence. In the NBD simulation case, the bias correction seems to have amplified slightly the oscillation via recursion. Here, we used three iterations. Thus, the bias iteration can behave as an additional regulator of the problem. The size of the bootstrap samples for the daughter and mother algorithms should be taken as large as computational resources allow. As a practical guideline, extensive simulations should be always used to investigate the problem-specific stability, regularization, bias-variance and uniqueness.

CHAPTER 7. ON THE INVERSION OF HIGH ENERGY PROTON

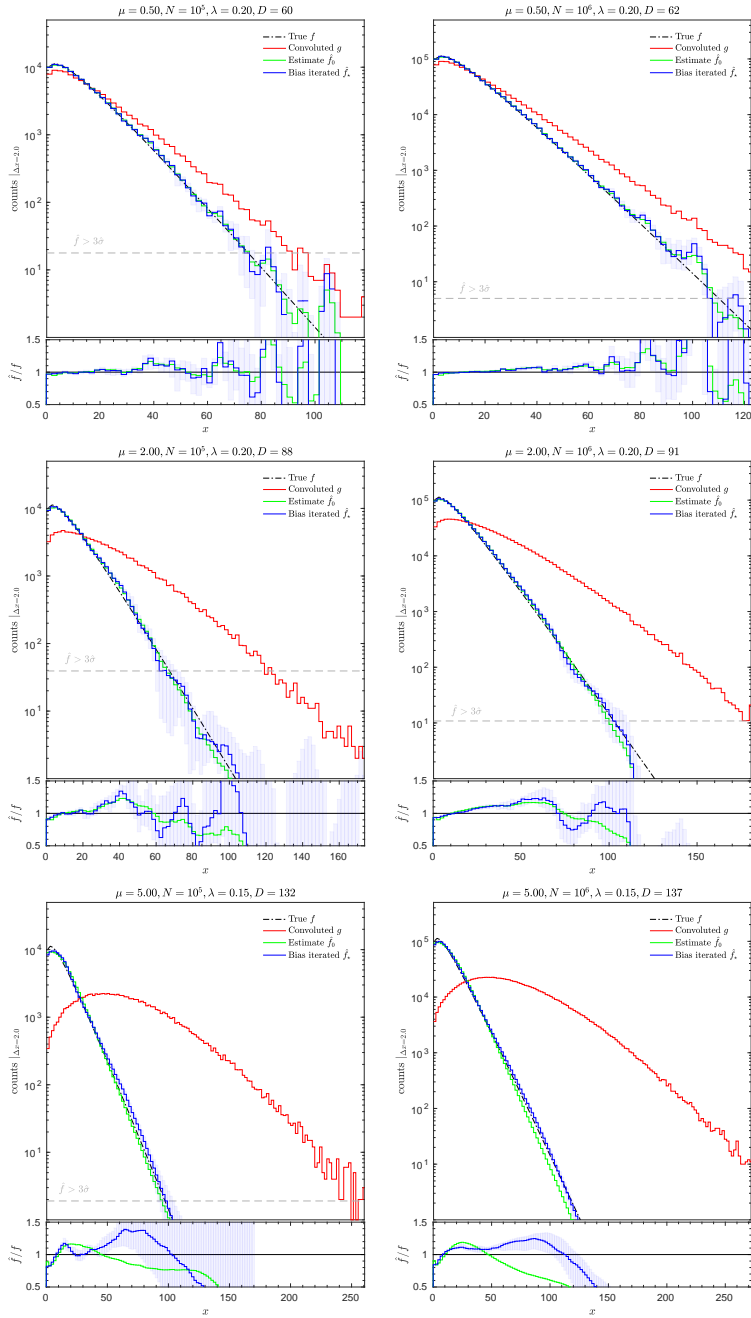


Figure 7.5: Simulations with $f(x) = P_{\text{NBD}}(x; \langle n \rangle = 12.5, k = 1.4)$. The bootstrap based uncertainties are shown in blue.

CHAPTER 7. ON THE INVERSION OF HIGH ENERGY PROTON

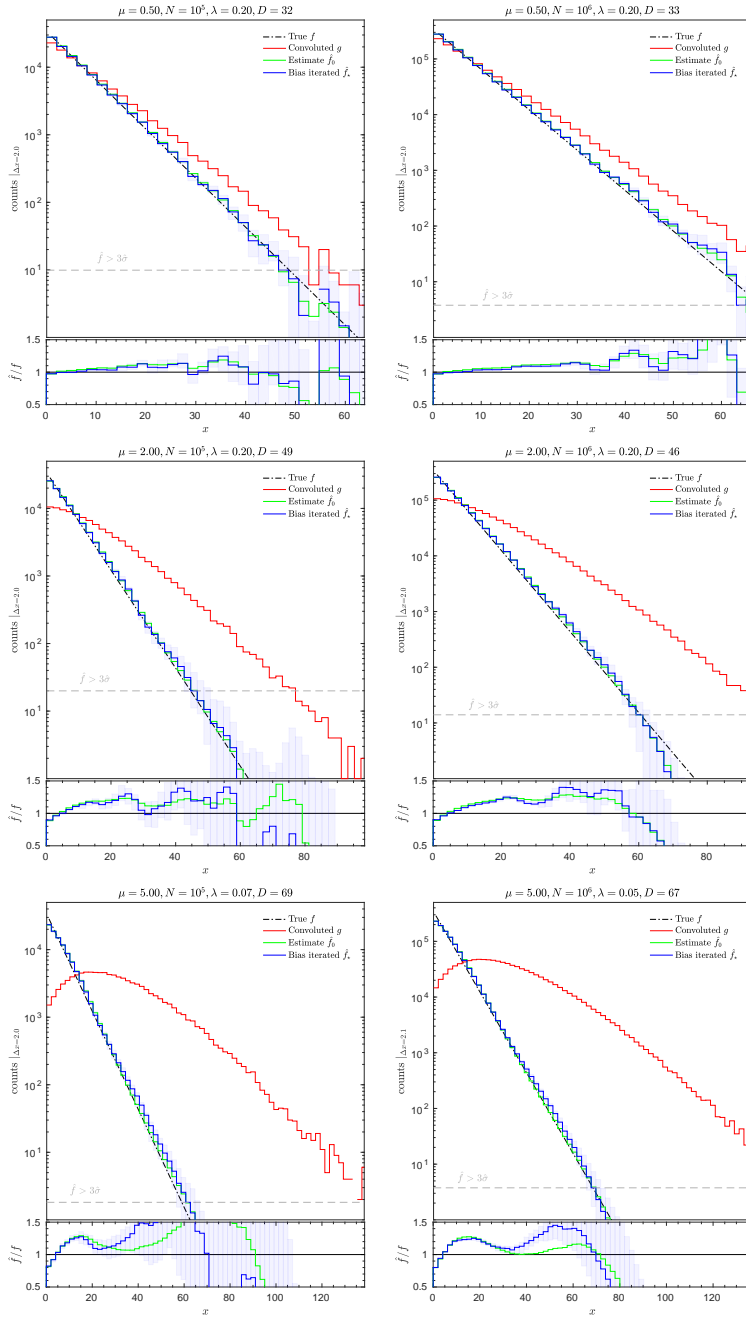


Figure 7.6: Simulations with $f(x) = 1/\alpha \exp(-x/\alpha)$ with $\alpha = 6$. The bootstrap based uncertainties are shown in blue.

7.6 LHC data inversion

For this proof-of-concept study, we use ALICE proton-proton charged particle multiplicity spectra measured at $\sqrt{s} = 0.9$ and 7 TeV from [HEPData](#) [226]. The publication includes also data at $\sqrt{s} = 8$ TeV, which we leave out here because of very similarity with $\sqrt{s} = 7$ TeV, but we include it in our analysis code available online. Phenomenologically, the average central multiplicity density $dN_{ch}/d\eta$ in proton-proton collisions follows Regge like power law scaling $\propto s^{\alpha(0)-1}$ with the effective pomeron intercept $\alpha(0) \simeq 1.1$. The event selection definitions are: the INEL (minimum bias inelastic) which in this case means minimal activity in any of the trigger subdetectors and the NSD (non-single-diffractive), which is event rapidity topology based selection. The idea behind the NSD is to suppress *qualitatively* single diffraction events within forward system mass range which result in a large enough pseudorapidity gap on forward or backward pseudorapidity side of the detector. The single diffractive suppression is done simply by requiring event activity on both forward and backward triggers. In addition, the INEL events may have either $N_{ch} \geq 0$ or ≥ 1 particles (tracklets) required in the central region $|\eta| < 1$. We use the $N_{ch} \geq 0$ class, because diffractive events can be with $N_{ch} = 0$ at central but trigger forward detectors.

We executed KISU inversion simply based on the Poisson compounding hypothesis with different μ values and the regularization being fixed to small $\lambda = 0.03$ in order not to bias the distribution shapes. Varying the regularization strength results in principle in a systematic uncertainty, but we observed it mainly affects the small scale structure. The ALICE data uncertainties are a combination of systematic factors with technically unknown distribution coverage, for example, soft QCD Monte Carlo modeling affecting estimated trigger efficiencies and detector unfolding, and statistical counting fluctuations. For the systematic shape variations we did a minimum and maximum global shape shift and evaluated statistical bin-by-bin bootstrap Poisson re-sampling on top of that, with the given event sample sizes $\mathcal{O}(7 \cdot 10^6)$ at 0.9 TeV and $\mathcal{O}(6 \cdot 10^7)$ at 7 TeV from [226]. The uncertainties in Figures 7.7 and 7.8 represent the 95CL values of this procedure. In principle, one could have done also bin-by-bin bootstrap for the systematic shapes, but because that procedure would neglect all bin-to-bin continuity correlations, it would give way too large high frequency fluctuations not reflecting typical systematic spectrum distortion or bias variations.

As a surprise from the inversion, a strong hidden secondary NBD like peak structure appears at certain μ hypothesis, most strongly with Poisson $\mu \approx 3 \dots 4$ number of simultaneous ‘mini-collisions’, with the zero-suppressed average given by Equation

CHAPTER 7. ON THE INVERSION OF HIGH ENERGY PROTON

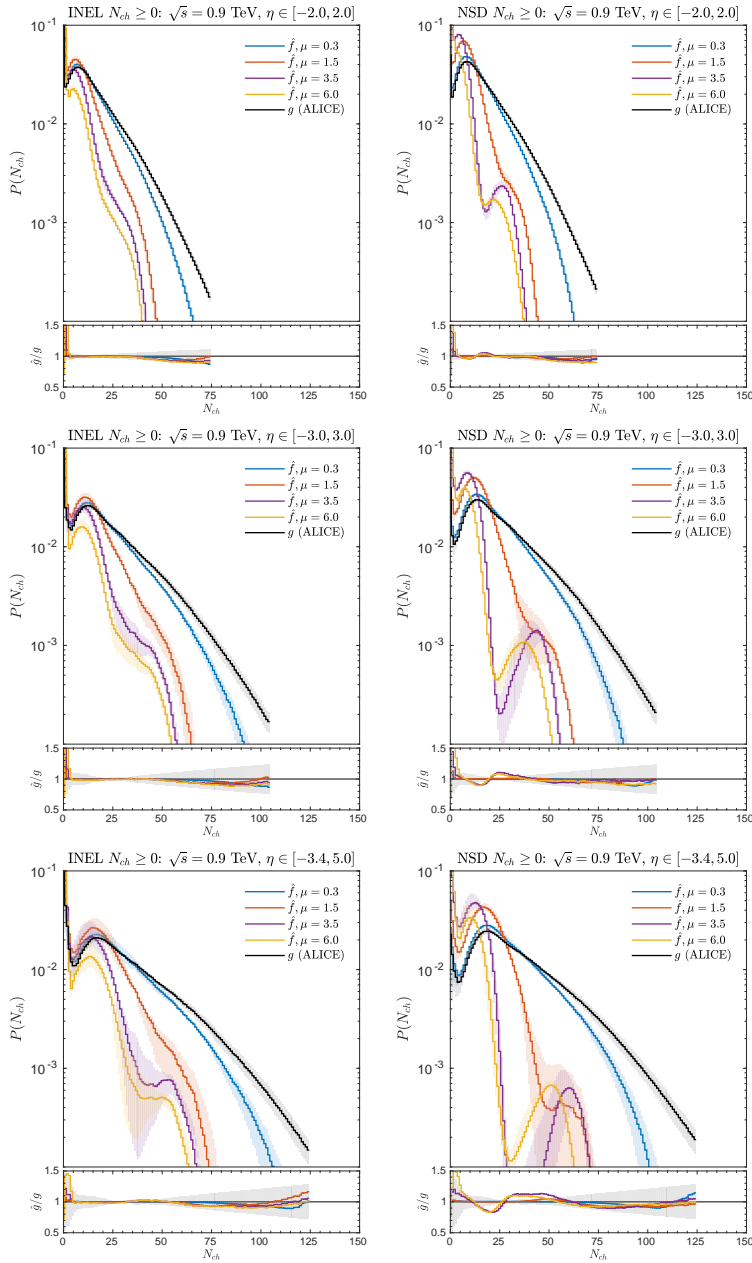


Figure 7.7: ALICE INEL $\sqrt{s} = 0.9$ TeV data on the left and NSD data on the right [226], for three different pseudorapidity intervals and the KISU inversion results under different μ -hypothesis.

CHAPTER 7. ON THE INVERSION OF HIGH ENERGY PROTON

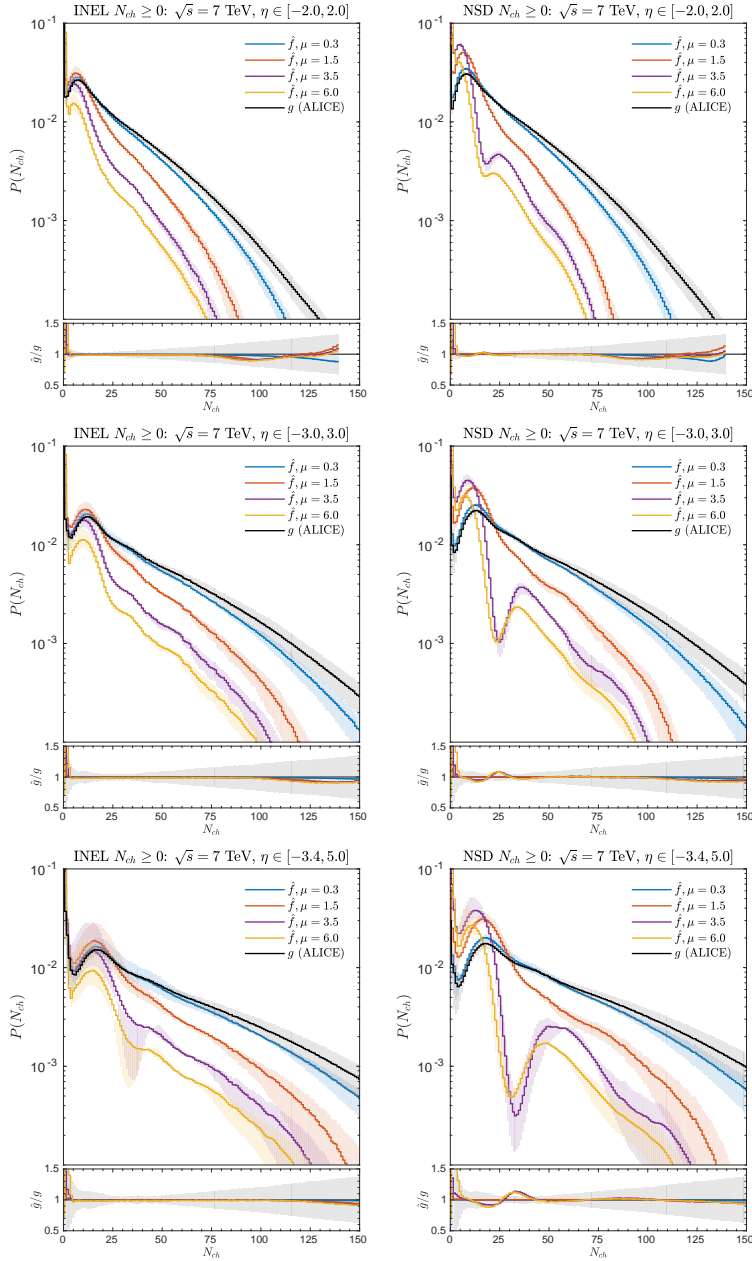


Figure 7.8: ALICE INEL $\sqrt{s} = 7$ TeV data on the left and NSD data on the right [226], for three different pseudorapidity intervals and the KISU inversion results under different μ -hypothesis.

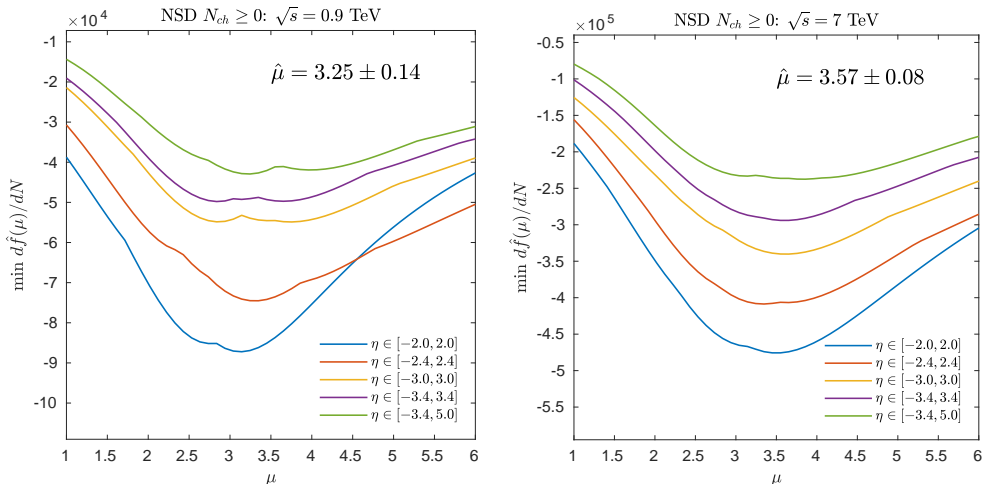


Figure 7.9: ALICE NSD data parameter μ -scans for two different energies, at $\sqrt{s} = 0.9$ TeV on the left and at 7 TeV on the right.

7.3. Also interesting is that this peak is not strongly visible in the INEL class but appears only once single diffractive events have been suppressed using the NSD class. It seems that the first NBD like peak is similar at both energies, but the second runs with energy. The INEL with $N_{ch} \geq 1$ event class (figures available using our [code](#)) gives results roughly similar to the NSD class, this requirement effectively also suppress diffraction. Also, in general the secondary peak increases in relative amplitude when the pseudorapidity window is enlarged and also the peak position shifts. A thorough interpretation of this observation requires further studies. A typical interpretation of the negative binomial double peak structures is that there are two separate mechanisms for the particle production initiator, soft confining and hard point like. The high multiplicity tail is also interesting, usually difficult to describe perfectly by many Monte Carlo models. For comparisons, see [226].

In Figure 7.9, we scan numerically over μ for the minimum of $\frac{df(\mu)}{dN}$ as the criteria for the steepest dip. The results for $\hat{\mu}$ show the mean and its standard error over different pseudorapidity intervals, value growing with energy, which is expected in the multiparton interaction picture with increasing particle densities. At lower energy, there seems to be larger variation with the solutions. We found out that using larger regularization λ values pushes the minimum towards lower values of μ , perhaps simply (over)-smoothing the solutions. Finally, we see in Figure 7.7 and 7.8 ratio plots

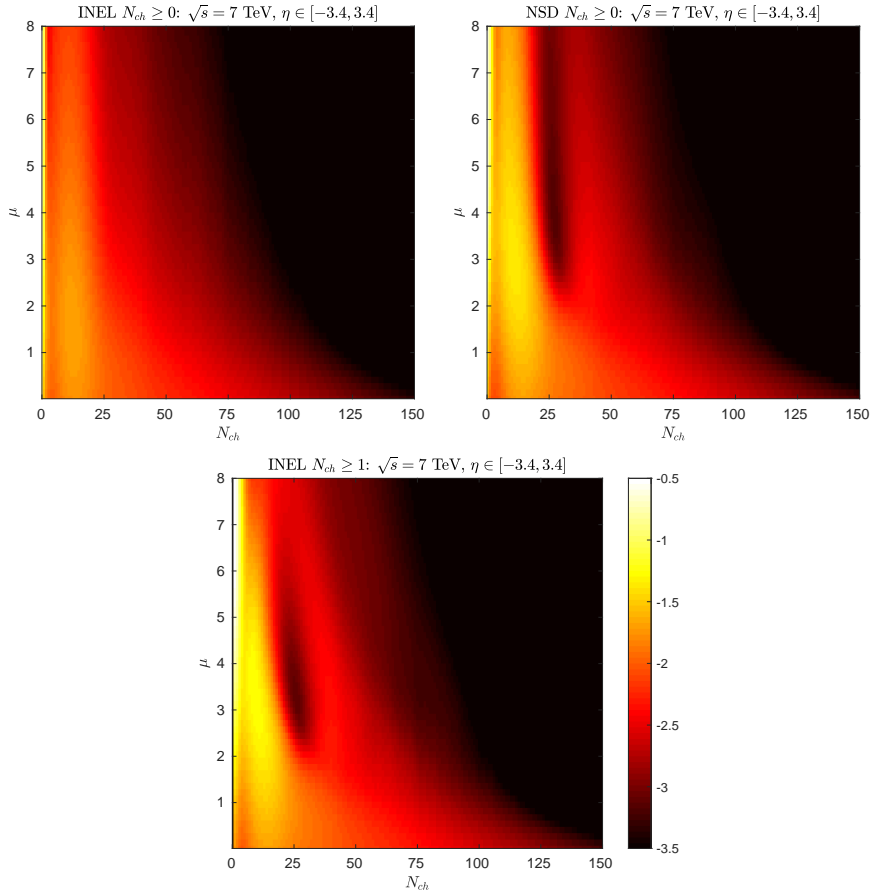


Figure 7.10: Inversion results of \hat{f} in the (N_{ch}, μ) plane at $\sqrt{s} = 7$ TeV for $\eta \in [-3.4, 3.4]$. Colors denote values of $\log_{10}[P(N_{ch})]$.

that the re-projection $\hat{g} = \hat{\mathcal{F}}(\hat{f})$ versus the measured g has the largest tension at very low multiplicities $N_{ch} < 5$, in the domain dominated by diffraction with qualitatively different behavior than the NBD peaks. The second interesting domain is around $N_{ch} \sim 25$, where oscillations are manifest. Figure 7.10 demonstrates the ‘topological’ differences between the inversion results as a function of the μ parameter. The distribution splits and develops the secondary peak at a certain μ -value in the case of the NSD $N_{ch} \geq 0$ event class. In the case of the INEL $N_{ch} \geq 1$ splitting also happens, and the re-fusion of two peaks happens faster as a function of the growing μ -value. For the INEL $N_{ch} \geq 0$ class, on the other hand, no splitting is visible.

7.7 Conclusions and prospects

We studied the inversion of a stochastic autoconvolution integral equation and showed that useful inverse solutions can be obtained with a novel algorithm. Special effort was invested in the algorithmic construction from bottom-up and in statistical and systematic uncertainty estimation. A new class of semi-model free phenomenological interpretations and statistical tests of inclusive distributions in soft QCD are possible with the algorithm. It is possibly also useful in jet substructure studies when used together with jet algorithms. The algorithm can be also utilized directly as an experimental pileup inversion algorithm in minimum bias studies. There are also possibilities to combine the statistical approach derived here with event-by-event pileup correction algorithms such as [179, 227]. Future directions may include more detailed analytic treatment in terms of both phenomenology of soft QCD, but also technical development of the algorithm for other types of statistical integral equations. The ultimate goal could be higher dimensional tomography of high energy proton, which may be realistic with the help of deep neural network techniques.

Inverting the final state multiplicity spectrum caused by independent multiparton interactions (MPI) is an obvious research target. We did a simple one free parameter inversion using ALICE multiplicity measurements, which exposed an interesting secondary peak structure most prominent with the NSD event selection class. We are not aware of a similar study, typically superposition fits with multiple negative binomial distributions are done, or Monte Carlo model-based studies. One could do a further study where different cms-energies and rapidity intervals are matched together using the autoconvolution picture in a direct or inverse direction. Speaking in the Regge theory language, an approximately similar concept is the multiple independent pomeron exchanges with ‘AGK cuts’ [25]. Given a model of exchange probabilities of multipomeron exchange with inelastic cuts, such as the inelastic eikonal approximation, in principle we can from the measured observables reconstruct a single pomeron exchange observables without a Monte Carlo model. Deviation from expectations indicates multipomeron (enhanced) interactions which are nonlinear effects not fitting in the K -fold autoconvolution picture. Thus, the inverse algorithm result can be seen as a statistical null-hypothesis generator. Probing the fuzzy interface between multiplicities driven by a semi-hard scale interaction versus soft cut pomerons, could be interesting. Also, the nonlinear gluon saturation conjecture at low Bjorken- x and the robust experimental observables of it are open problems.

Combining the approach here with other mathematical tools, the shadow of diffraction can be studied in a new light. That is, it would be interesting to see

if data can *directly* show evidence for the AGK cuts like phenomena – certain specific additive algebraic rules, multiplicity density, rapidity gap topologies. So far, no water-proof observables with one-to-one mapping with AGK have been measured, as far as we know. We also see that methods here could give new insight to the ‘centrality’ determination in heavy ion collisions when combined with the Gribov-Glauber model. Inverting the measured transverse energy distribution could lead to a more precise understanding of pure ‘binary scaling’ (independent nucleon-nucleon collisions) versus strong nuclear effects. The mathematical description here is, essentially, what is approximately meant by binary scaling.

8 DeepEfficiency

We introduce a new high dimensional algorithm for efficiency corrected, maximally Monte Carlo event generator independent fiducial measurements at the LHC and beyond. The approach is driven probabilistically using a Deep Neural Network on an event-by-event basis, trained using detector simulation and even only pure phase space distributed events. This approach gives also a glimpse into the future of high energy physics, where experiments publish new type of measurements in a radically multidimensional way.

Chapter in: [arXiv:1809.06101](https://arxiv.org/abs/1809.06101) [[physics.data-an](#)]

High energy physics data needs to be corrected for experimental effects induced by the kinematically non-uniform response of the detector and reconstruction algorithms. These corrections are usually dubbed efficiency times acceptance corrections, implemented often by dividing the measured number of background corrected events in a histogram bin with the efficiency value obtained by simulations. In addition, unfolding of histograms of observables is often done to account for distortions of resolution (variance) and absolute value (bias) induced by the detector. Unfolding is usually treated with a stochastic smearing matrix, obtained via simulations. It is the art of unfolding and optimization algorithms, with suitable regularization and possible additional physical boundary conditions, which then “inverts” this response matrix in an algebraic or probabilistic way. An important part is the uncertainty estimation, both in purely statistical terms but also in systematic ways. Well known is that many of the naive bin-by-bin methods are severely biased towards “Monte Carlo truth” [228].

A crucial step of the measurement is the proper definition of the fiducial phase space. Ideally, this should be defined in terms of proper final state observables, such as transverse momentum and pseudorapidity of charged particles, within geometrically visible and electrically active detector volume. A fiducial measurement, by its very definition, minimizes extrapolations of data and thus also minimizes the model dependence. In this work, we go further and propose a *maximally* Monte Carlo event generator independent way to implement fiducial efficiency corrections. For the best of our knowledge, this philosophy and approach what we call *DeepEfficiency*, based on Deep Neural Networks (DNN), is a new one. However, in other contexts the networks are well known in high energy physics, especially in signal-background separation. For a recent review see [229], or for the pioneering studies [230].

For concreteness, let us define our observable degrees of freedom at the level of individual final state particles. Our final state consist of N charged particles (tracks) with 3-momentum \vec{p} of each being measured. This final state spans a real valued vector space \mathbb{R}^{3N} , from which one can construct a set of Lorentz scalars and other non-scalar observable distributions, such as transverse momenta or angular ones. The detector has a finite probability \mathbb{P} of seeing the N -track event, given the trigger and tracking efficiencies. This expected probability is represented with an efficiency mapping $\mathcal{E} : \mathbb{R}^{3N} \rightarrow \mathbb{P}$. Clearly, by using space-time symmetries of interest, one can often do dimensional reduction. However, treating maximally the abstract detector space, forbids us from doing so. That is, we want to take into account all the correlations that the detector may induce in terms of efficiency losses.

A simplification, dimensional reduction but also a loss of information would occur, if one would for example factorize all tracks and get the full event detection efficiency

as $\mathcal{E}_{tot} = \mathcal{E}_1 \mathcal{E}_2 \cdots \mathcal{E}_N$. However, for very high dimensional problems, one may do this. We point out here that the efficiency function \mathcal{E} can be used also within the context of Matrix Element Method (MEM) type likelihood analyses, in addition to the fiducial measurements which are of our interest here. For typical approximations within MEM applications, see [231].

What we simply now do is learn this high dimensional detector efficiency function \mathcal{E} using a fully connected, multilayer DNN in a regression mode. That is, we simply use the neural network as a high dimensional function approximator and interpolator - as a mathematical hammer. The exact architecture, the cost function and its regularization and the gradient descent methods are art forms each of their own, we return to these in extended descriptions. Because we use the fully differential observable kinematic information of the final states, we can even use a pure phase space Monte Carlo event generator as an input to the detector simulation. Perhaps the most intuitive proof of the vanishing dependence on the event generator is based on a multidimensional histogram division picture with hyperbin volumes approaching zero.

A practical requirement for the generator is that a large enough event statistics is produced continuously for all corners of the phase space of interest, naturally. For final states containing jets, a pure phase space Monte Carlo is not feasible. Thus in practice one proceeds as usual with a realistic MC generator of interest. One may do this also in other cases, naturally. Here we do not touch the highly interesting but complex topic of matching optimally parton level, hadron level and detector level objects and sub-structures, which is basically always non-trivial, sometimes ill-posed, with hard or soft QCD processes.

In optimizing (training) the network parameter set Θ , we minimize the so-called cross entropy cost function suitable for our probabilistic inference problem

$$L(\Theta, \mathcal{S}) = -\frac{1}{|\mathcal{S}|} \sum_{i \in \mathcal{S}} [R_i \ln \mathcal{E}_i + (1 - R_i) \ln(1 - \mathcal{E}_i)], \quad (8.1)$$

where the sum is over all events in the simulation sample \mathcal{S} . The known response R is 0 for efficiency lost events, and 1 for selected events, both within the fiducial phase space. That is, one must not include events outside the fiducial phase space in the sample, by construction. A cutoff regularization is used for the rare $\mathcal{E} \rightarrow 0$ singularity. The word entropy comes from the functional form of the cost, related to the Kullback-Leibler [213] divergence which is a difference between cross entropy and entropy, and the KL-divergence itself is related to the Maximum Likelihood principle. In addition, one can add typical regularization schemes such as ℓ_1 - (sparsity, Laplace) or ℓ_2 -norm (smoothness, Gaussian) based, in order to avoid overfitting. This is avoided also

effectively by choosing the minimal network architecture which results in efficiency corrections with minimal bias and variance on simulated samples. Also the training sample size needs to be high enough, given that the deep network architectures can contain $\mathcal{O}(10^4 - 10^6)$ free parameters.

After training, the efficiency inversion of data and arbitrary observables of interest are obtained with

$$\frac{dN}{d\mathcal{O}} = h_{\mathcal{O}}(\{\vec{p}\}) \odot [\mathcal{E}(\{\vec{p}\})]^{-1}, \quad (8.2)$$

where $h_{\mathcal{O}}$ is a probability distribution estimator operator, typically a bin width $\Delta\mathcal{O}$ normalized histogram. That is, the discretization of observables only enters at this point. The point-wise operator \odot is defined as an integral (sum) over the event sample and the weight $\mathcal{E}(\{\vec{p}\})$ is obtained from the neural network. Simply put, event-by-event, one constructs the observable of interest and calls a weighted histogram fill with a weight \mathcal{E}^{-1} . Thus, in an extraordinary smooth way, one can efficiency correct *simultaneously* arbitrary number of single or multidimensional histograms of observables using the same weight. For a related mathematical discussion, see the Horvitz-Thompson unbiased estimator [232]. Differential cross sections are obtained in a standard way normalizing by integrated luminosity. A straightforward way to obtain statistical uncertainties is via Efron’s bootstrap re-sampling. The possible algorithmic (network) bias should be empirically tested using simulations, observable by observable, which is easily automated.

The performance was numerically evaluated using the full ALICE detector simulation with input being low-mass, low- p_T central exclusive QCD diffraction, which is a prominent “glueball production process” at the LHC, decaying in this case to charged meson final states ($\pi^+\pi^-$, K^+K^-). We obtained solid results with a 5-layer network with ~ 100 neurons per layer with a 6-dimensional input, using hyperbolic tangent activation functions and a final layer with one sigmoidal output function. The inversion performance in terms of χ^2/bin was close to unity for a variate of one and two dimensional differential distributions, also in the steep tails. TensorFlow 1.9 [233] was used as the computational framework. The event generator independence lemma was observed to hold with Monte Carlo samples driven by different scattering amplitudes, such as photoproduction of $\rho^0 \rightarrow \pi^+\pi^-$ versus $\pi^+\pi^-$ production via double Pomeron exchange. A typical efficiency correction bias coming from unknown angular dependence or spin polarization densities, due to uncertain non-perturbative scattering amplitudes for both continuum and resonance production - is a solved problem now. Also the efficiency correction bias, propagating from a priori unknown but parametrized non-perturbative proton elastic form factors and inelastic structure functions driving the system soft transverse momentum distributions, vanishes.

Summarizing, we introduced a formally optimal and generic efficiency inversion algorithm for fiducial measurements. As a reminder, fiducial measurements are measurements minimizing the kinematic-geometrical acceptance extrapolations. However, in certain cases non-fiducial extrapolations are needed, for example the case of a fully angular flat phase space definition requirement for spin polarization studies. Another interesting topic is the high dimensional, continuous unfolding of distributions, which we did not discuss here. However, one can fuse single dimensional matrix unfolding algorithms with the efficiency inversion of DeepEfficiency.

9 Conclusions

The topic of this thesis was high energy diffraction. We obtained numerous *new results*, the most important are as follows.

1. We developed a novel algebraic-probabilistic framework which we called COMBINATORIAL SUPERSTATISTICS, as a new observable definition of soft QCD inclusive diffraction. In addition, we posed and solved a new combinatorial Poisson superposition problem with the Möbius inversion theorem. Then we applied the developed tools to the proton-proton data measured in ALICE at the cms energy $\sqrt{s} = 13$ TeV. This resulted in the first explicitly high dimensional unfolded fiducial measurement of soft QCD, as far as we know. One could call this a measurement of the Grassmannian manifold encapsulating the partial cross sections. All different finite field \mathbb{F}_2^N vector space subspaces are living on this manifold, and thus they hold very rich information to be understood in detail. Also, we implemented the first simultaneous multidimensional maximum likelihood extraction of diffractive cross sections and the effective Pomeron intercept. Finally, we demonstrated our novel F^* projection algorithm by reconstructing pseudorapidity gap distributions from incomplete data. The results showed an enhancement for the low mass diffraction in comparison with typical minimum bias event generator tunes.
2. For the exclusive or semi-exclusive diffraction, we developed a new Monte Carlo event generator, GRANIITI, written in modern C++. It is currently the leading event generator on this topic regarding custom spin-dependent low mass diffraction scattering amplitudes and computational performance. Using our simulations in comparison with the preliminary ATLAS+ALFA data at the cms energy $\sqrt{s} = 13$ TeV, we conjectured that tensor mesons and tensor glueballs could be produced with different spin polarization mixtures of transverse and longitudinal components. Thus, we claim that this explains non-perturbatively the enigmatic ‘glueball filter’, at least in the case of tensor resonances. A significant computational future problem is, how to calculate ab initio the resonance

couplings for the low mass resonance spectrum, which are currently free parameters. In addition, the proton dissociation probability and proton structure fluctuations are significant open problems, perhaps to be solved by future advances in lattice field theories. Experimentally, we did pioneering studies with the ALICE semi-exclusive central diffraction data regarding the invariant mass and transverse momentum composition and provided highly automated spherical harmonic decomposition code systematically in different Lorentz rest frames for advanced spin analyses, available in our GRANIETTI engine.

3. For studying multiparton interactions from a new perspective, we developed a novel inversion algorithm KISU, for K -fold stochastic autoconvolution integral equations. Then we applied the algorithm to the published ALICE charged particle multiplicity data measured with different cms energies and final state pseudorapidity ranges. We obtained strong evidence that the average number of soft interactions per collision is between 3 to 4 and that two different physical mechanisms are driving the multiplicity spectra because the inversion results in a clear separation between two-component distributions. Both results match the phenomenology often implemented in Monte Carlo event generators. The two-component model hypothesis is a classic one, such as a soft confining QCD and hard point like scattering. Our algorithm provides a new approach to attack this topic.
4. For precision fiducial measurements, we introduced the first fully differential detector efficiency inversion algorithm working event-by-event in higher phase space dimensions. This was based on deep neural networks and we called it DEEP EFFICIENCY. Effectively, we relied on the universal function approximation lemma of deep multilayer neural networks. In classic mathematics, the corresponding statement is the Stone-Weierstrass theorem which states that every continuous function can be represented with polynomials on a closed interval. A more generic extension of the efficiency inversion problem would be the higher dimensional unfolding of distributions. A general mathematical theory is still lacking for deep neural networks, which we see as the fundamental future problem to be solved. That would give theoretical tools to design and understand the networks beyond the current style of ‘empirical mathematics’. Also this algorithm still relies on a reliable detector simulation, which we see as an open problem to be solved in the future. Higher dimensional variants of the data-driven ‘tag-and-probe’ methods to estimate efficiencies could be possible.

CHAPTER 9. CONCLUSIONS

In terms of the big picture, tools from algebraic geometry and topology would be interesting to be studied in the context of high energy diffraction. From the experimental side, the future forward instrumentation should be improved significantly for precision measurements and discoveries in the forward domain, perhaps with the help of novel material innovations from the condensed matter physics.

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A ALICE measurements

A.1 Detector level marginal distributions

The detectors in Figures [A.1-A.13](#) are from left to right: ZDN/P_C, AD_C, V0_C, SPD_C, SPD_A, V0_A, AD_A, ZDN/P_A.

For each signal combination, the detector level charge Q (analog-to-digital conversion units) is shown in the upper row and the signal time t (ns) in the lower row. An exception to this is ZDN/P, where the upper row shows the ZDN charge and the lower row is for the ZDP charge. The general color coding is: white (signal accepted), gray (signal not accepted) and pink (spectator, not used in the analysis). The simulations are run through the full ALICE GEANT simulation with a run anchored calibration.

APPENDIX A. ALICE MEASUREMENTS

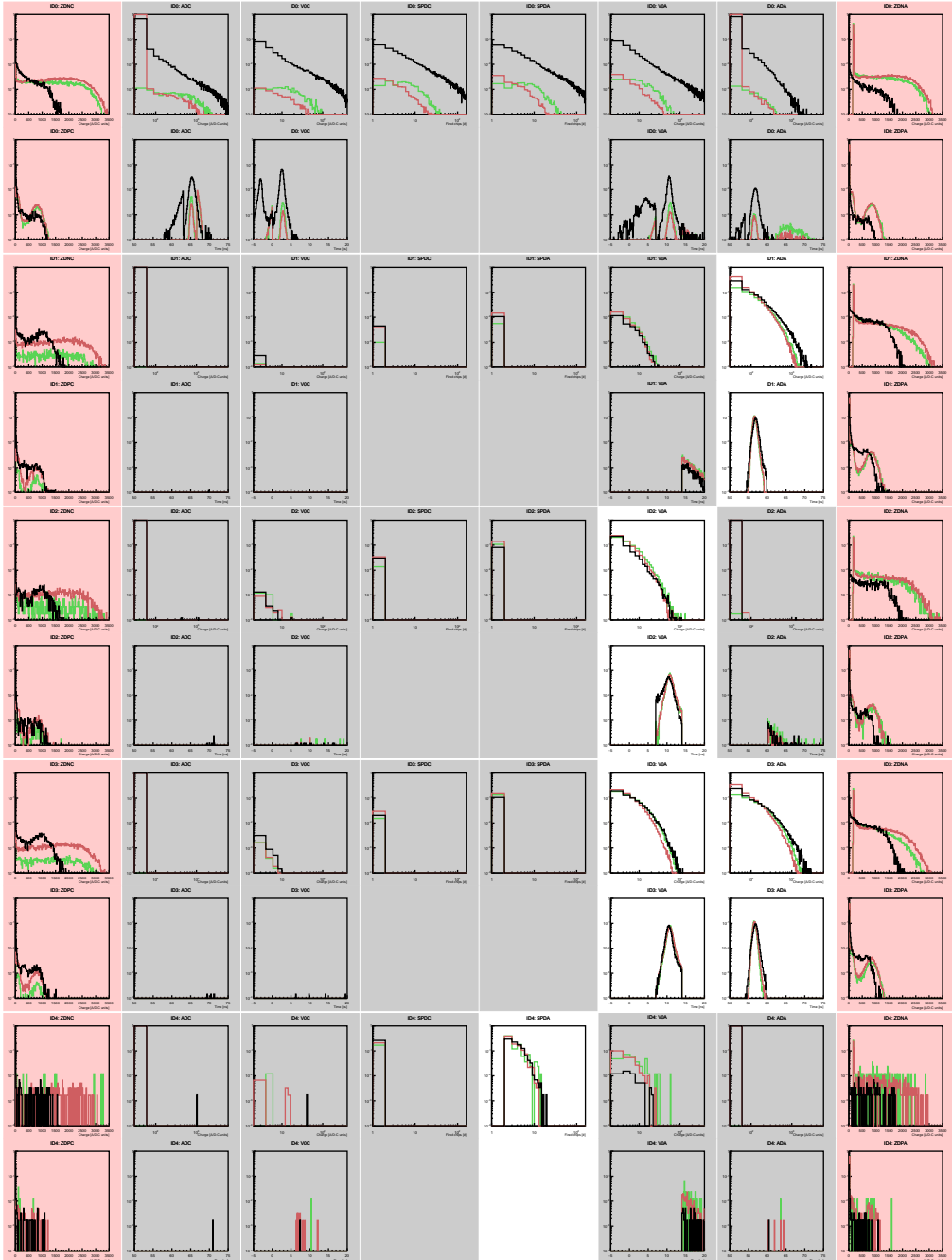


Figure A.1: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [0-4]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

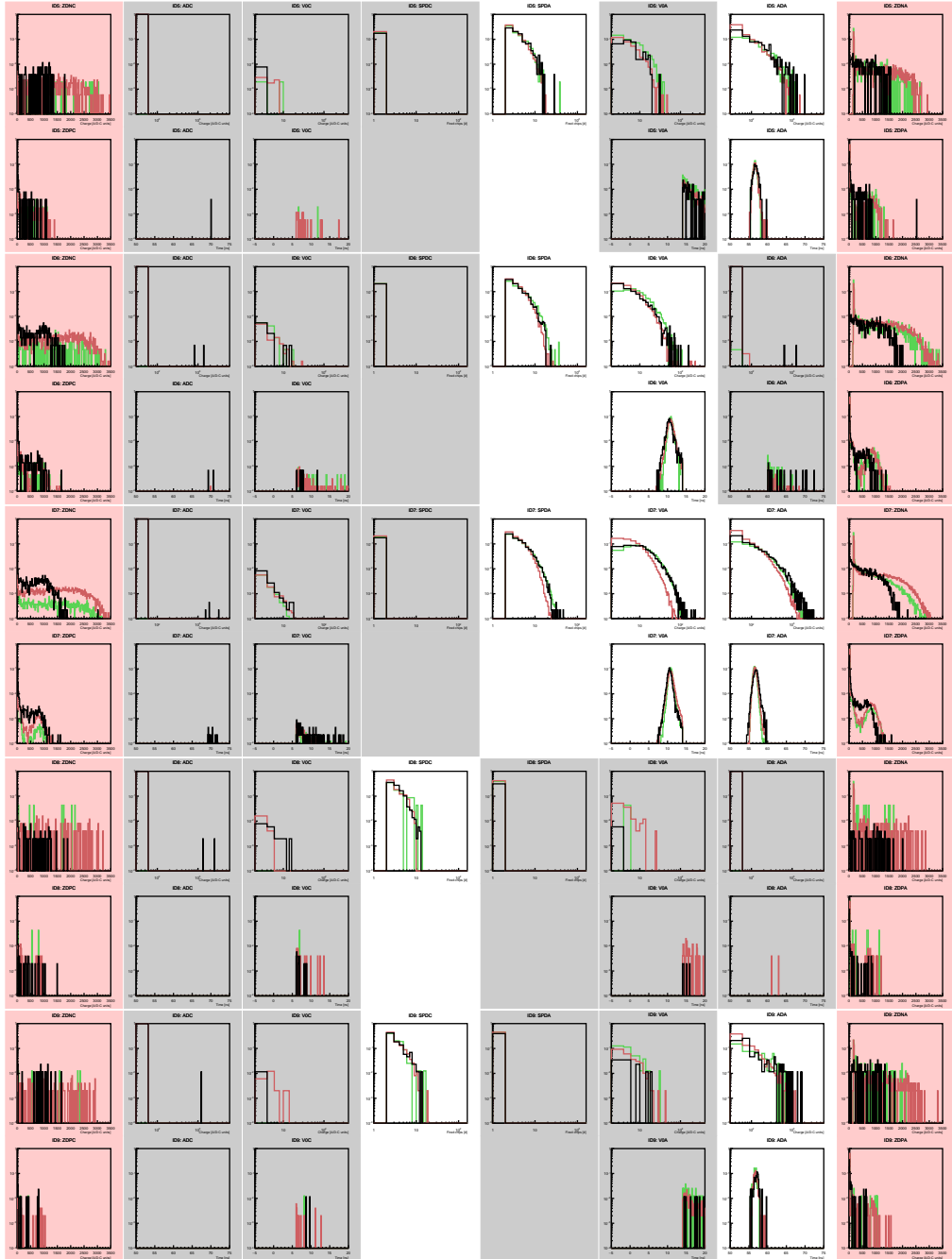


Figure A.2: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [5-9]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

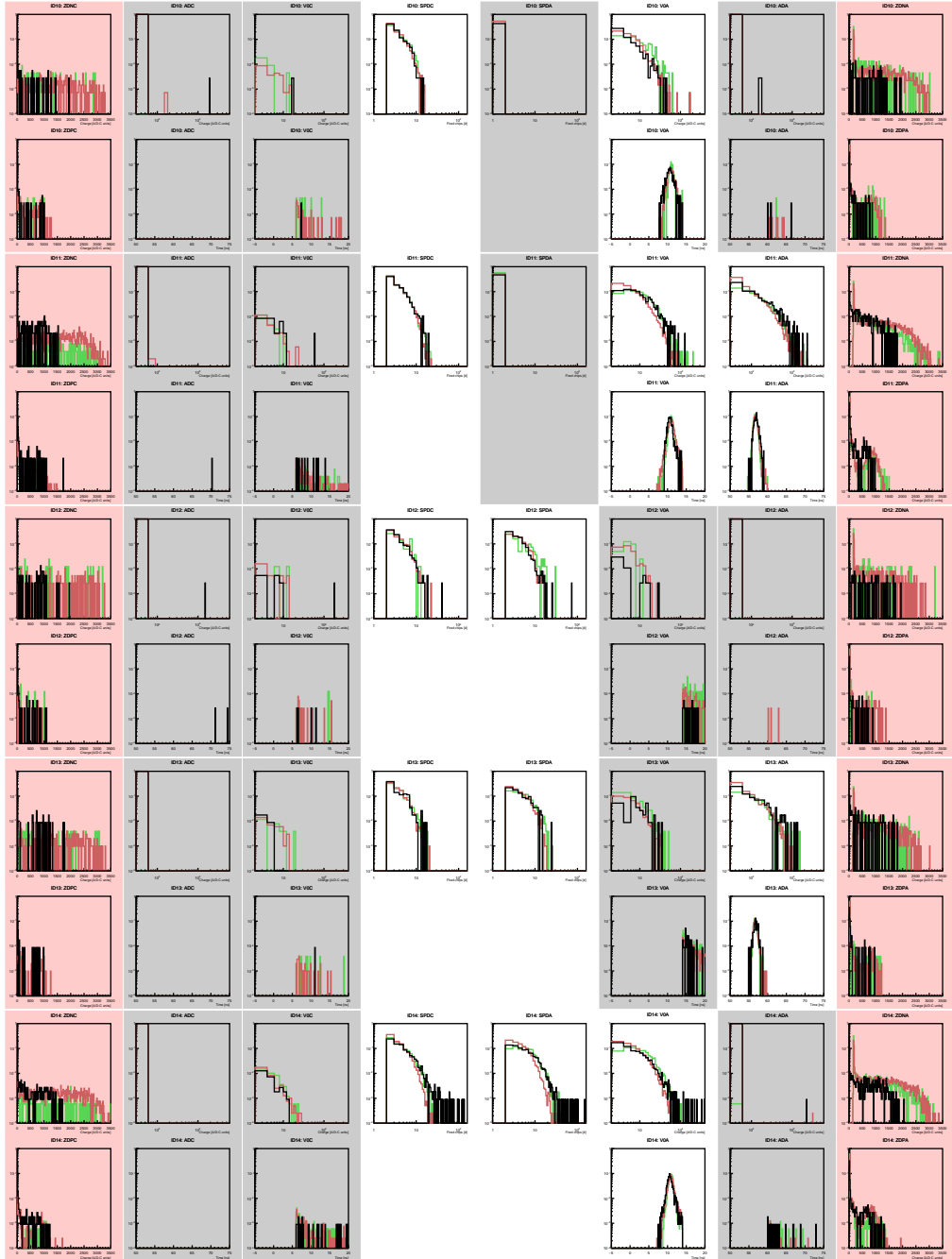


Figure A.3: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [10-14]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

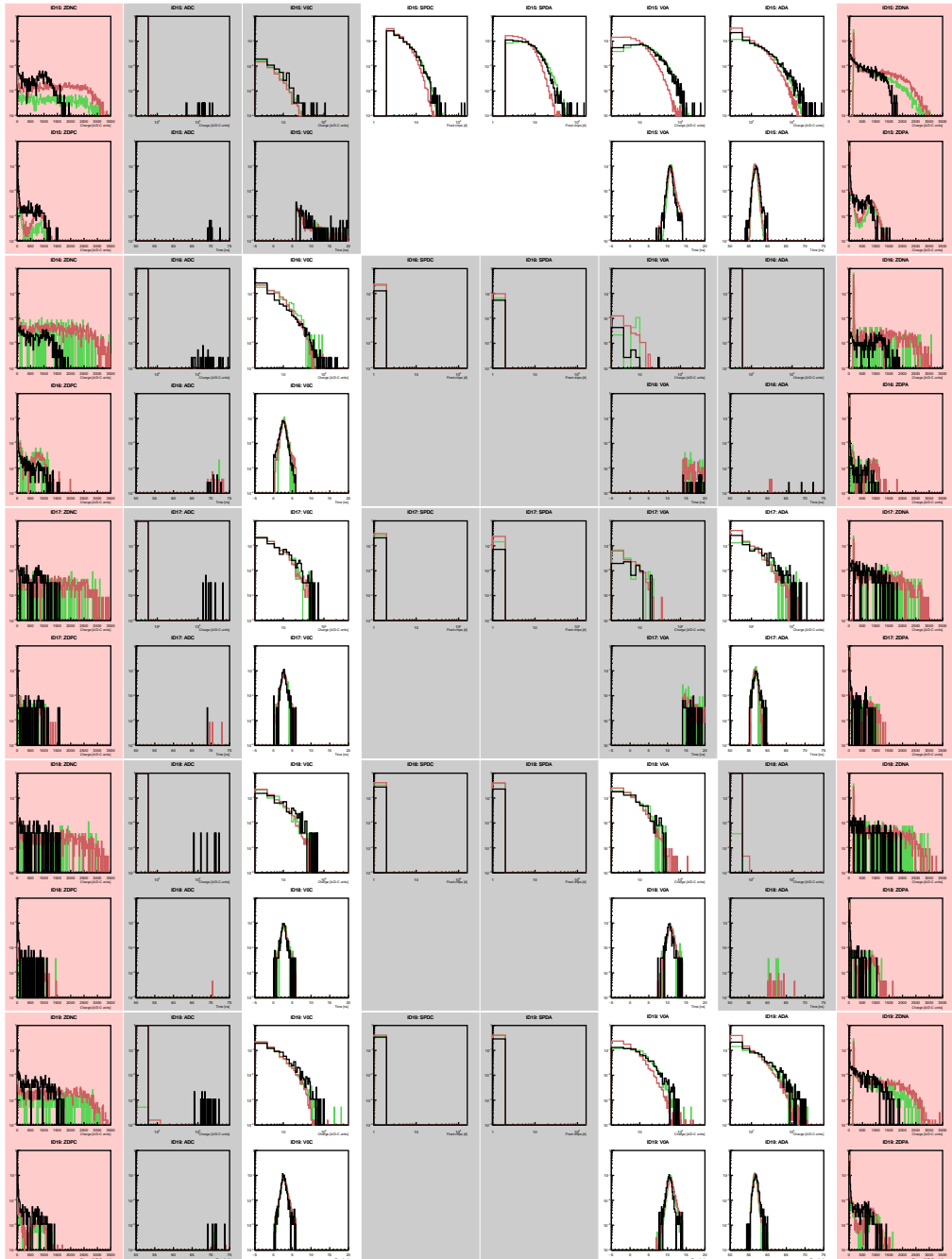


Figure A.4: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [15-19]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

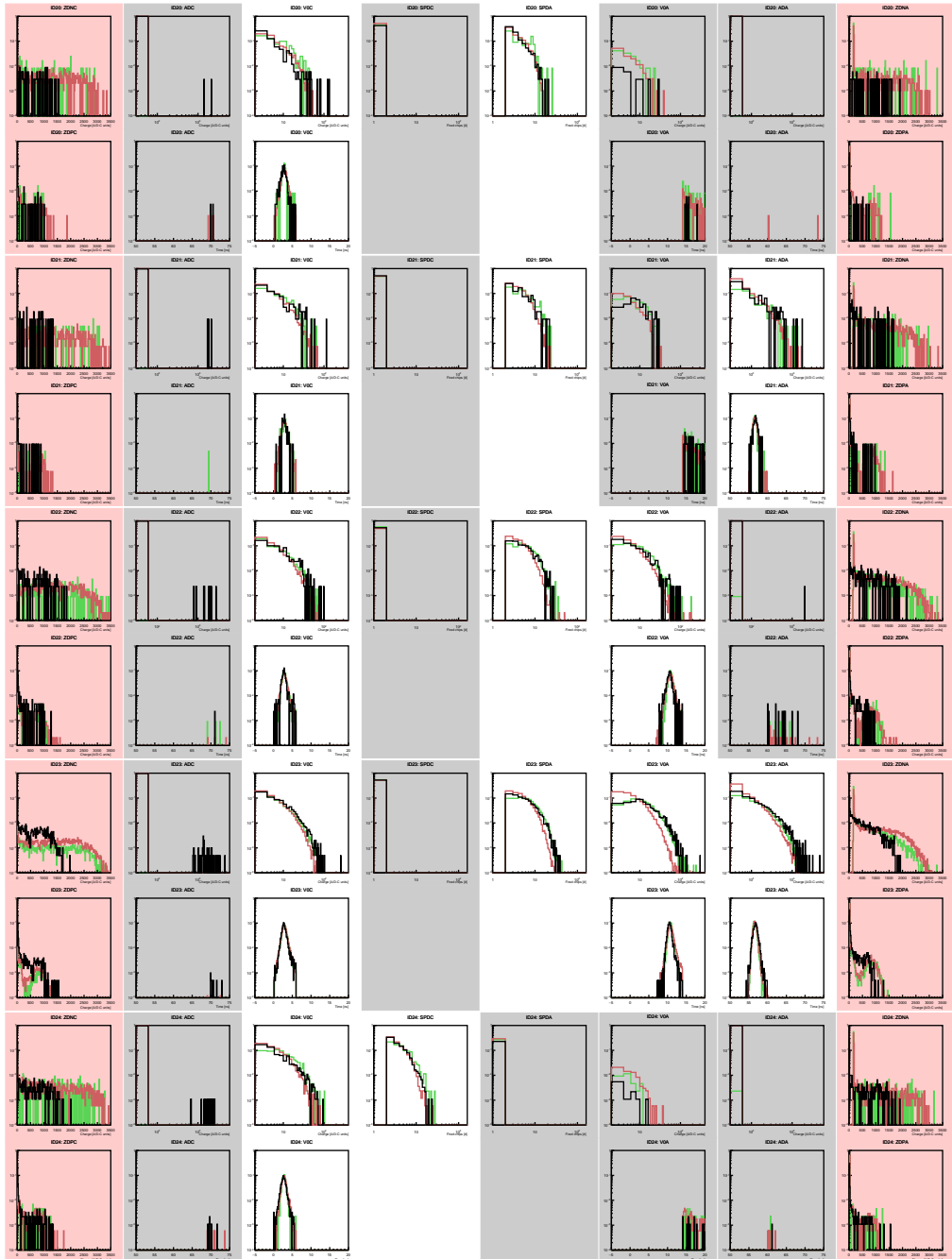


Figure A.5: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [20-24]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

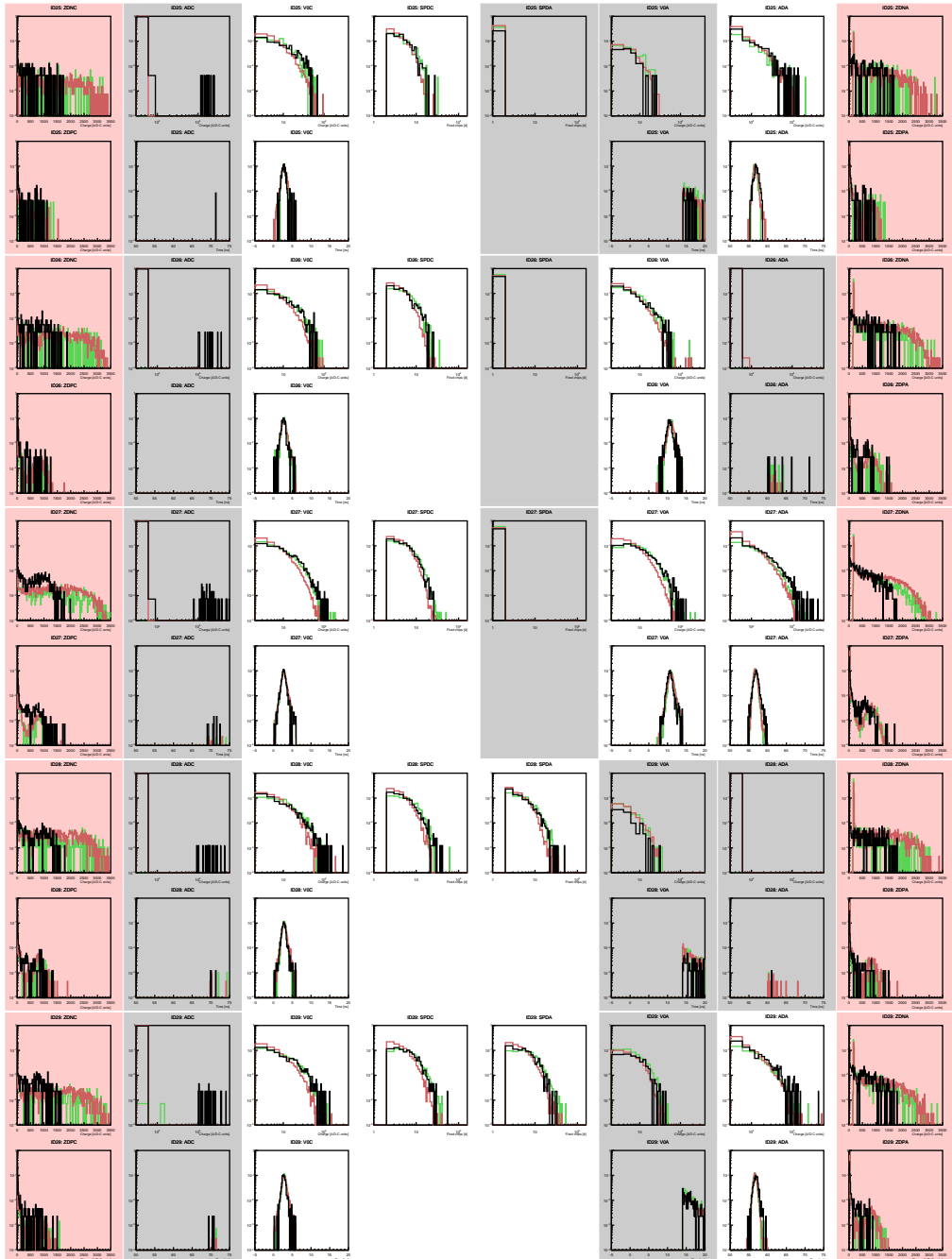


Figure A.6: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [25-29]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

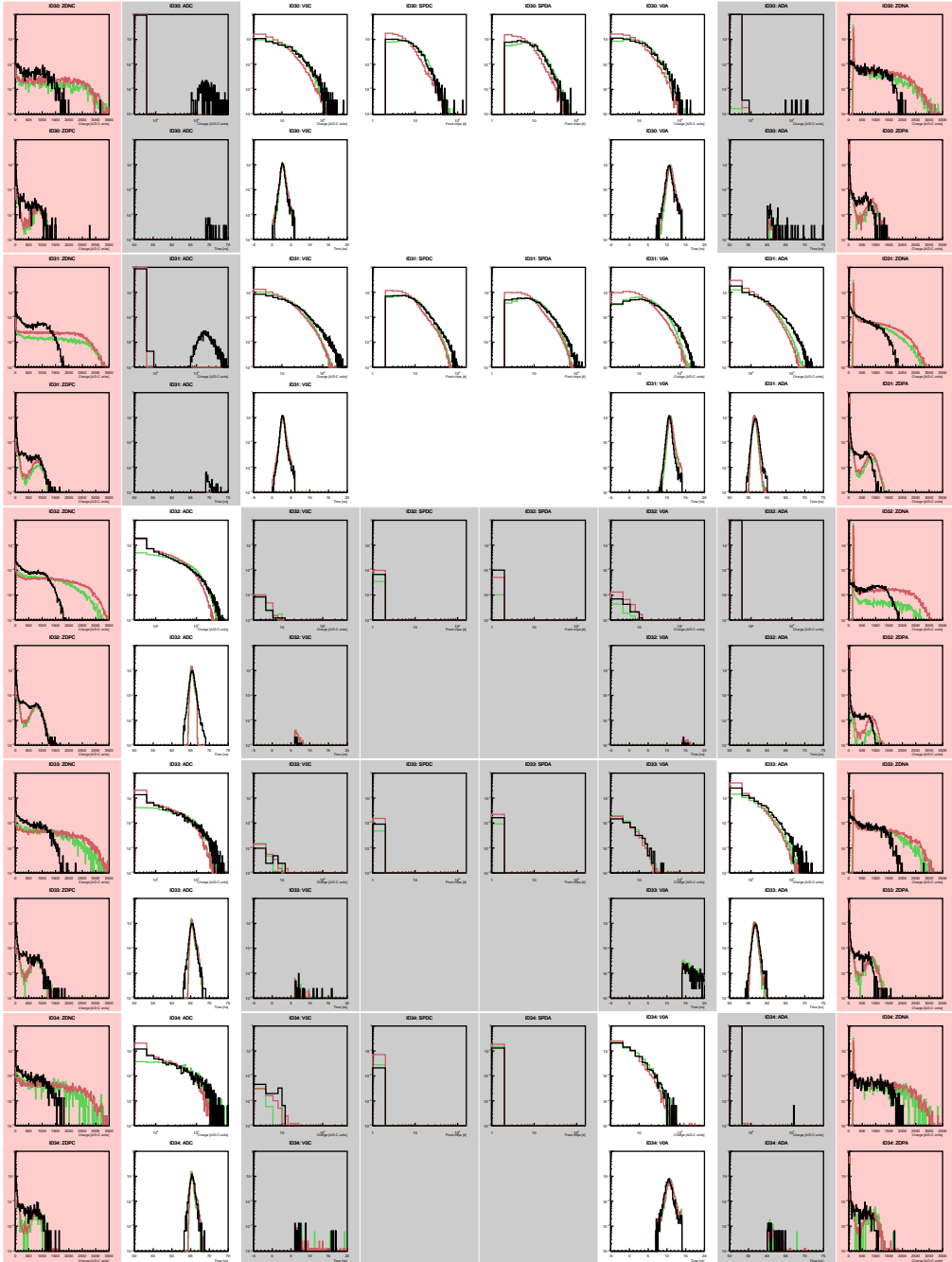


Figure A.7: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [30-34]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

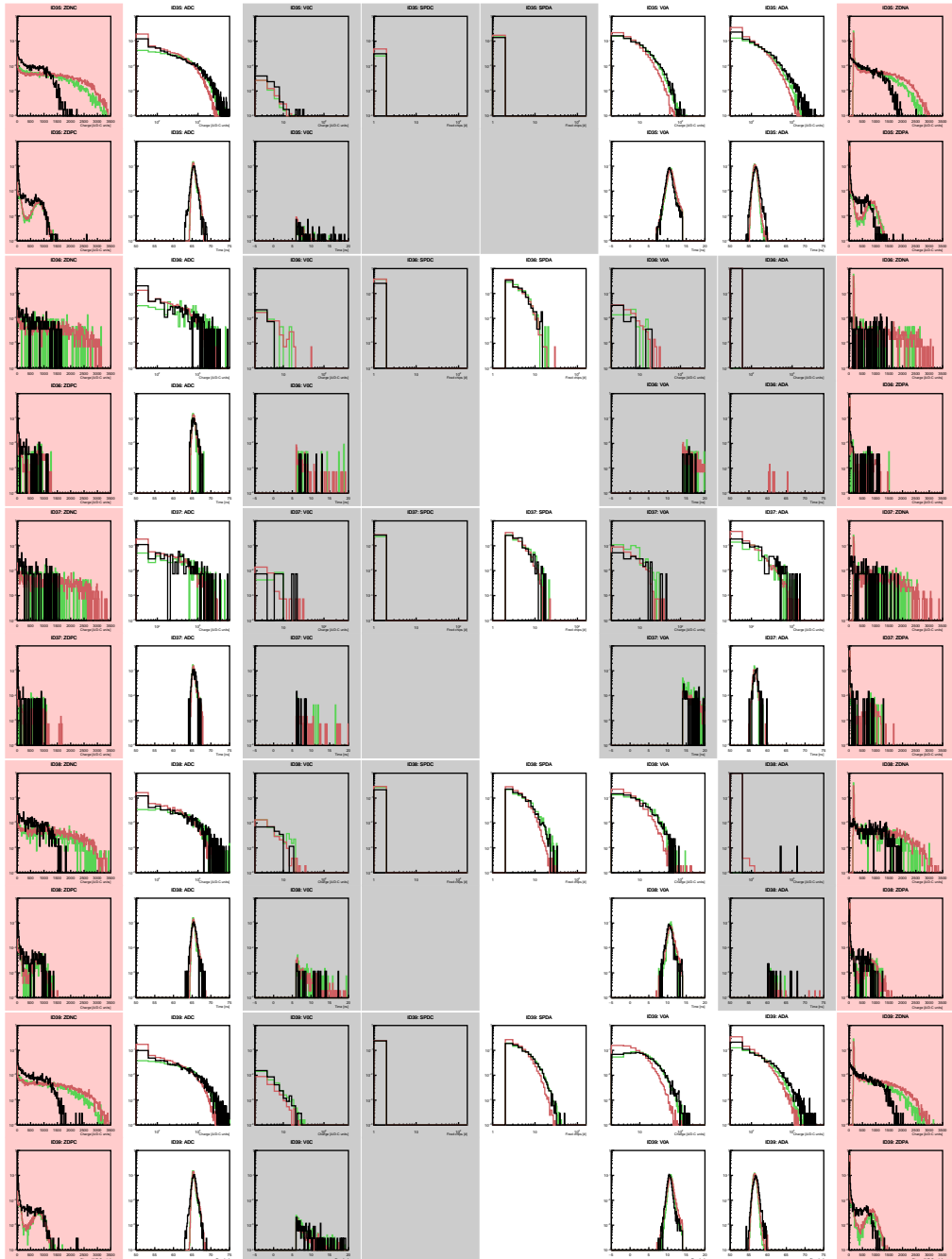


Figure A.8: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [35-39]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

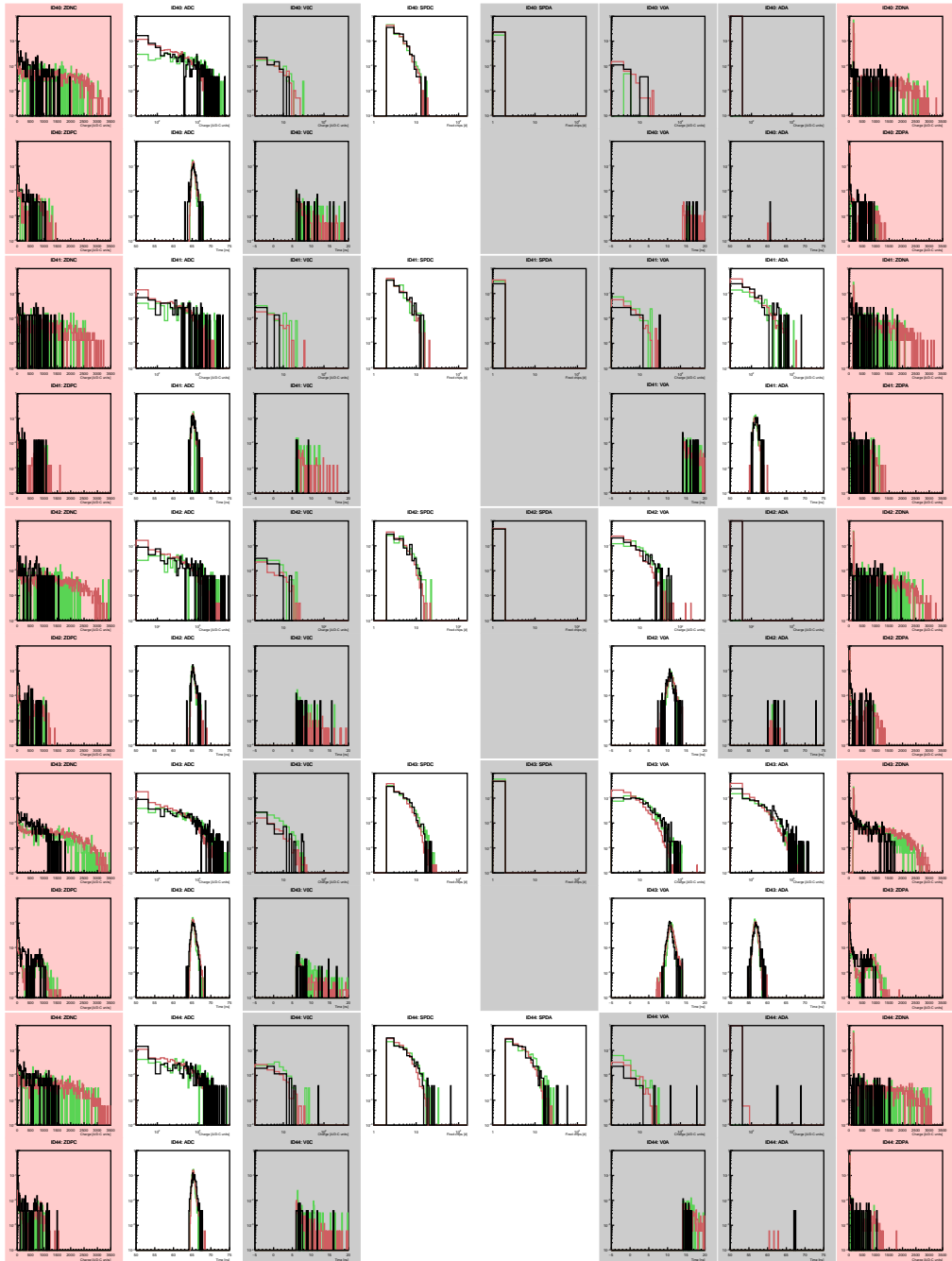


Figure A.9: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [40-44]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

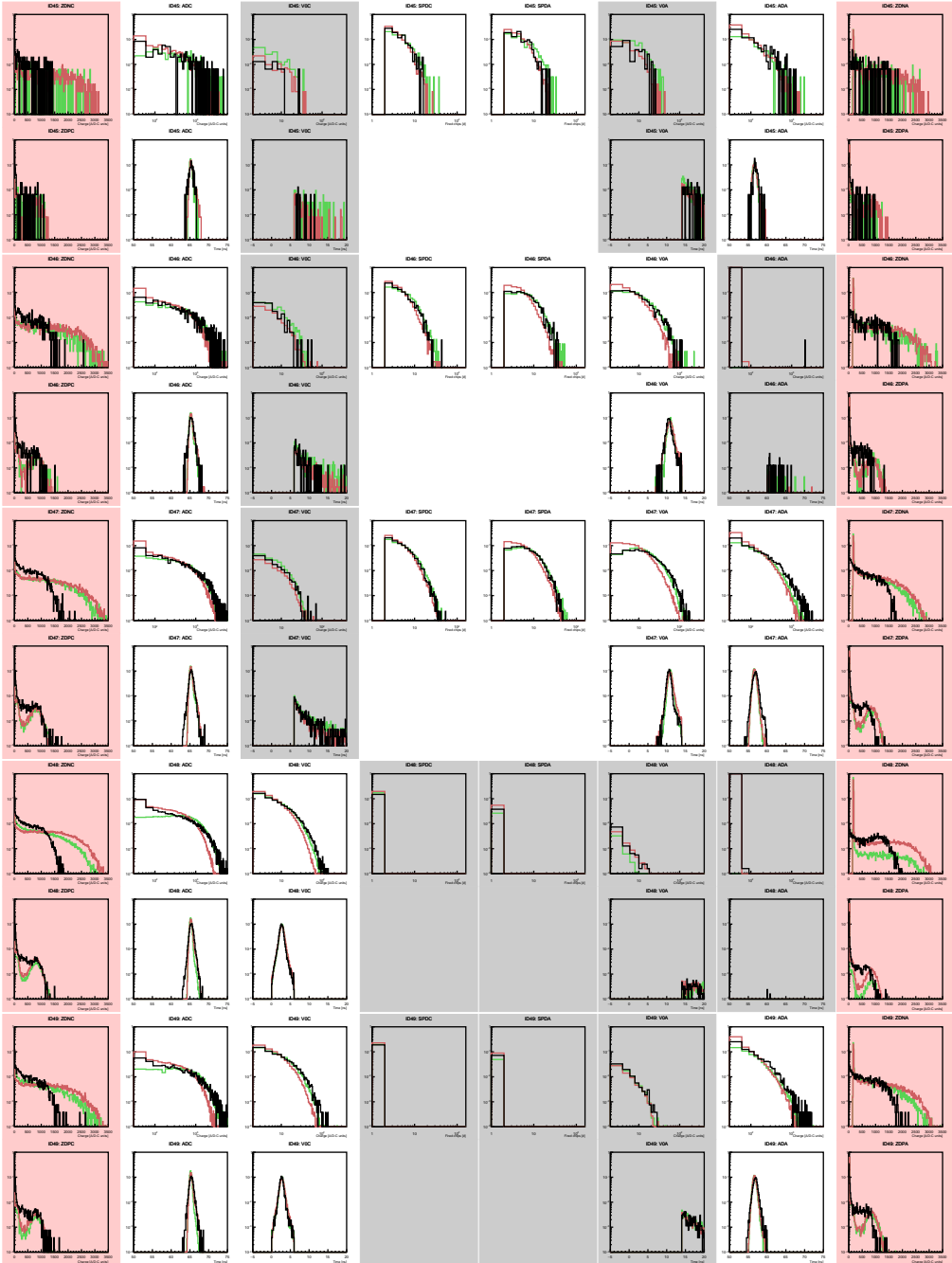


Figure A.10: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [45-49]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

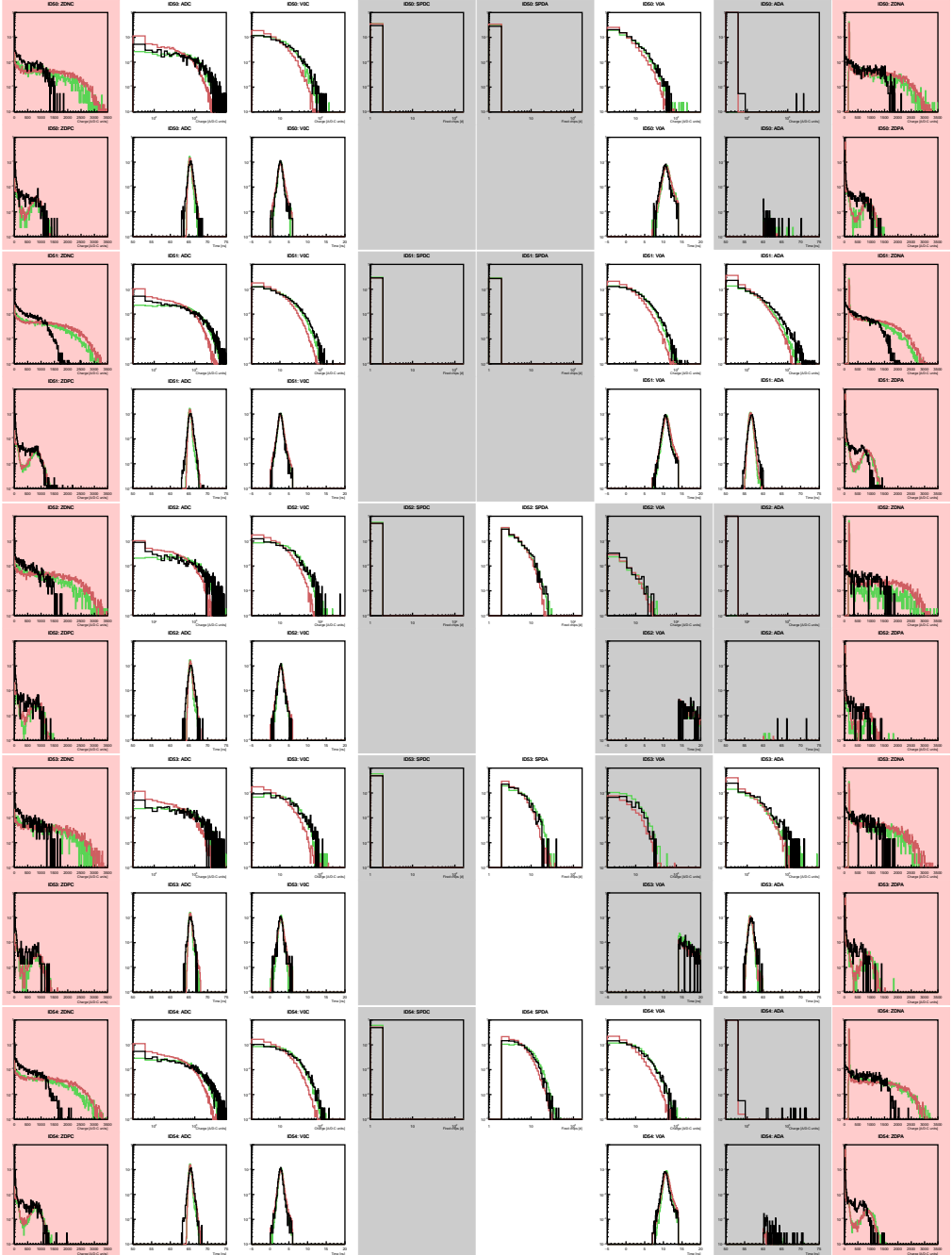


Figure A.11: ALICE data at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [50-54]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

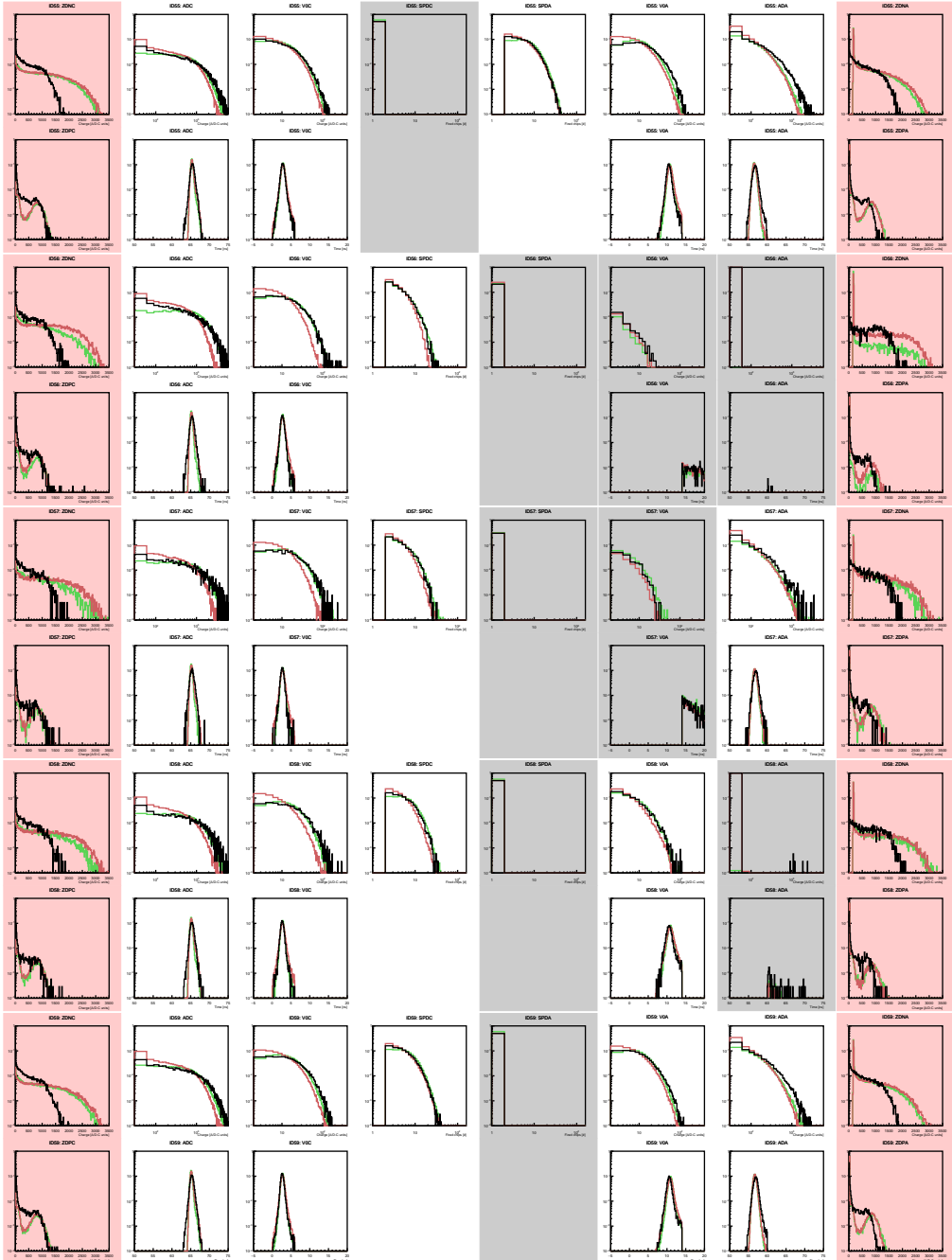


Figure A.12: ALICE data at $\sqrt{s} = 13$ TeV (black) and Pythia 6 (red), Phojet (green) simulations for combinations [55-59]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

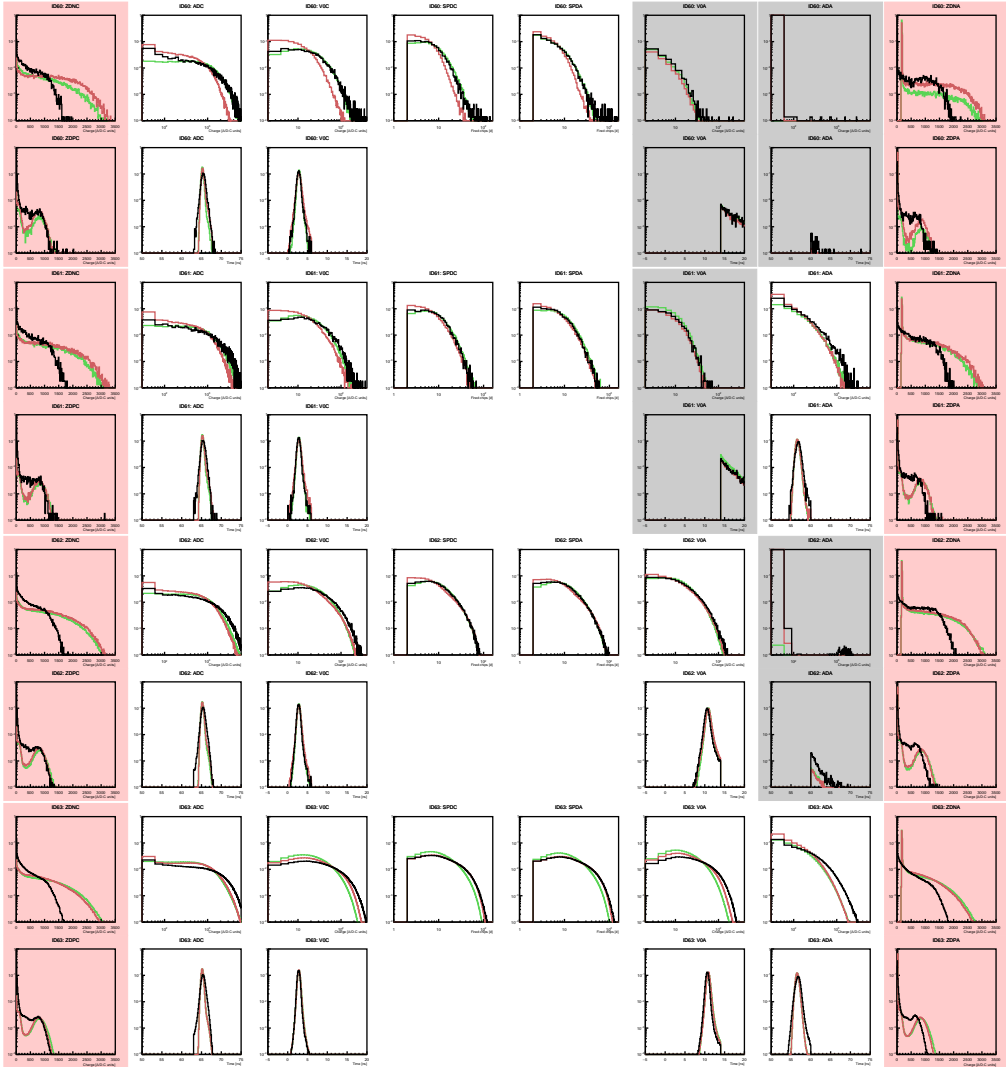


Figure A.13: ALICE **data** at $\sqrt{s} = 13$ TeV (black) and **Pythia 6** (red), **Phojet** (green) simulations for combinations [60-63]: Detector level charge and time distributions.

APPENDIX A. ALICE MEASUREMENTS

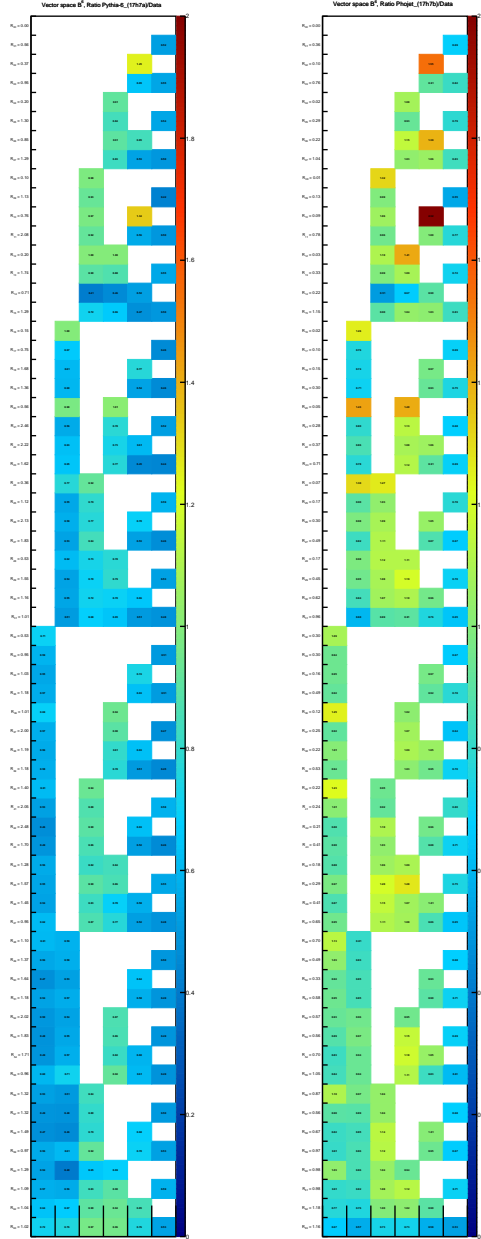


Figure A.14: ALICE data at $\sqrt{s} = 13$ TeV: Detector level visible cross section ratios for Pythia 6/Data (left) and Phojet/Data (right). Colored cells are mean signal ratios for: $\langle Q \rangle_{\text{ADC}}$, $\langle Q \rangle_{\text{V0C}}$, $\langle F \rangle_{\text{SPDC}}$, $\langle F \rangle_{\text{SPDA}}$, $\langle Q \rangle_{\text{V0A}}$, $\langle Q \rangle_{\text{ADA}}$, where Q is the charge and F is the number of fired chips.

A.2 Detector response simulations

APPENDIX A. ALICE MEASUREMENTS

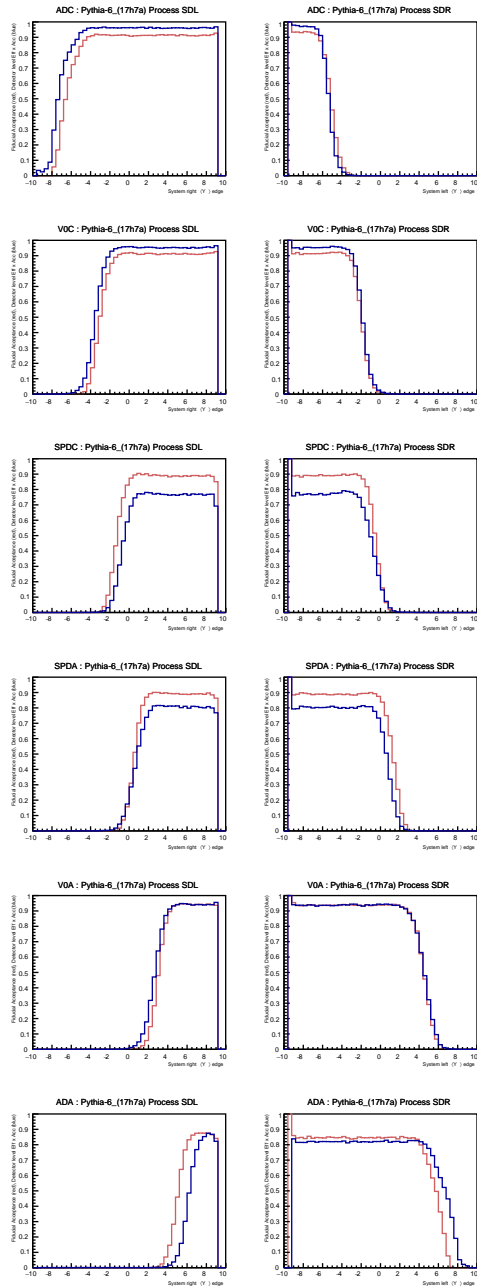


Figure A.15: ALICE simulation at $\sqrt{s} = 13$ TeV: Pythia 6 MC + GEANT based acceptance of diffractive systems (SDL left column, SDR right column) for different sub-detectors (as rows) in terms of the kinematic rapidity gap edge. Red is at the fiducial acceptance and blue at the detector efficiency \times acceptance \times smearing level.

APPENDIX A. ALICE MEASUREMENTS

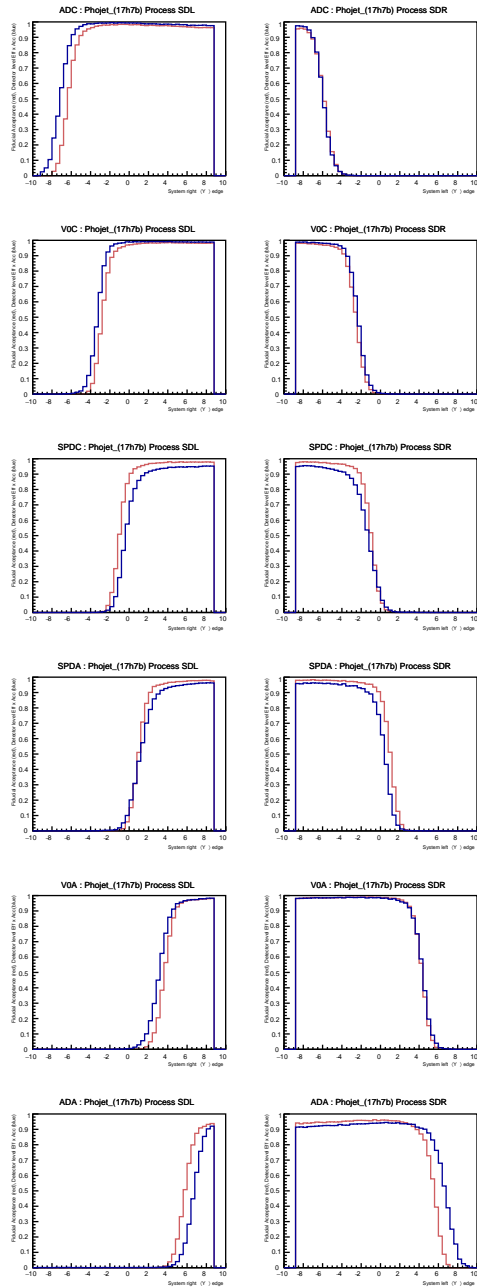


Figure A.16: ALICE simulation at $\sqrt{s} = 13$ TeV: Phojet MC + GEANT based acceptance of diffractive systems (SDL left column, SDR right column) for different sub-detectors (as rows) in terms of the kinematic rapidity gap edge. **Red** is at the fiducial acceptance and **blue** at the detector efficiency \times acceptance \times smearing level.

APPENDIX A. ALICE MEASUREMENTS

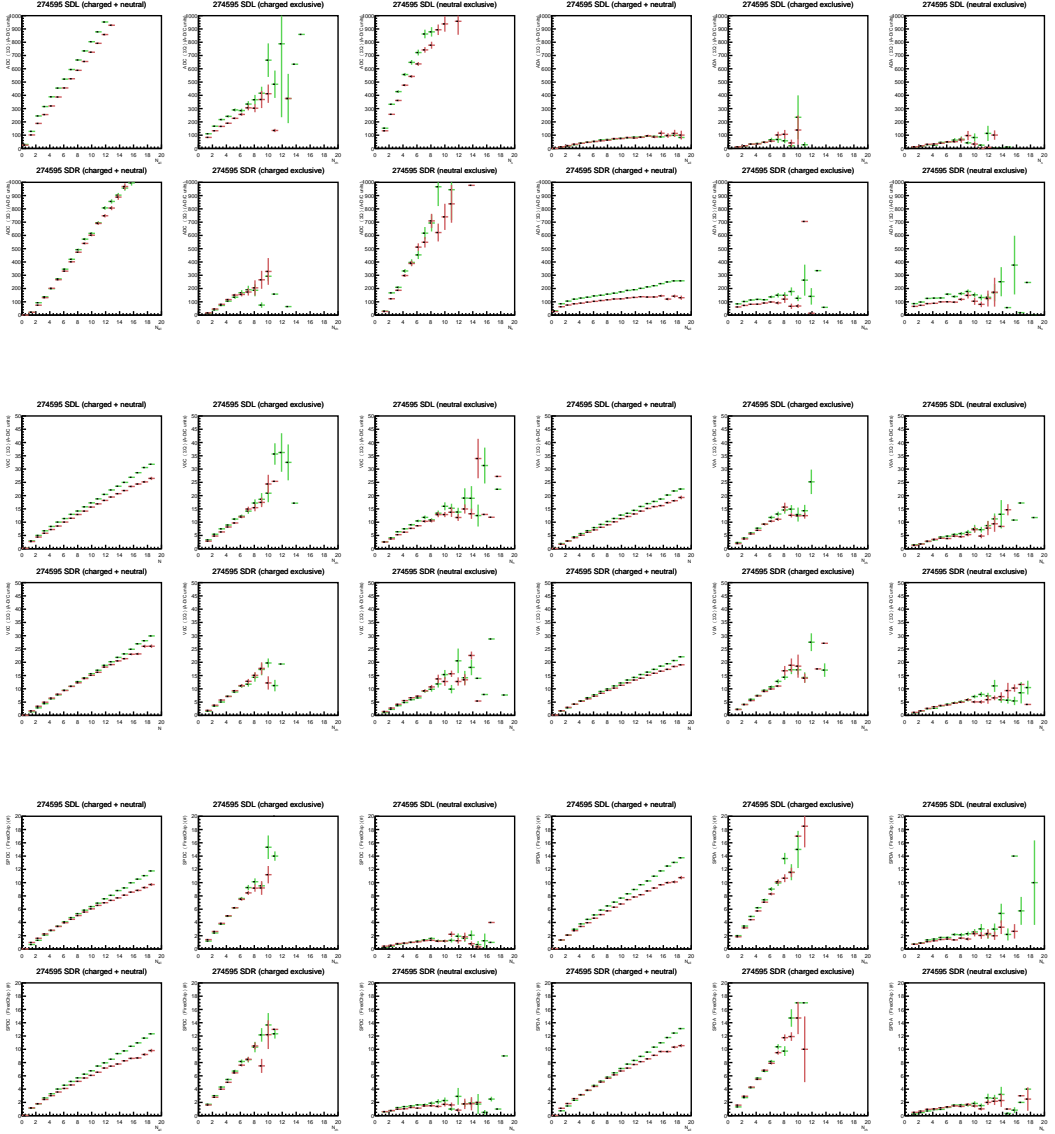


Figure A.17: ALICE simulation at $\sqrt{s} = 13$ TeV: **Pythia** (red) and **Phojet** (green) + GEANT simulated detector signal response as a function of the fiducial acceptance domain multiplicity containing charged, neutral or both type of particles for $AD_{C,A}$ in the upper row, $VZERO_{C,A}$ in the middle row and $SPD_{C,A}$ in the bottom row.

A.3 LHC optics

The nominal LHC optics parameters at the ALICE IP2 are given in Table A.1, which show the larger β^* , transverse plane separation and crossing angle than other LHC experiments. These are used to reduce the luminosity by three (+two) orders of magnitude compared to the ATLAS and CMS (+LHCb) cases.

Plane of crossing	-	Vertical
Bunch population	N	1.25×10^{11}
Revolution frequency	f_{rev}	11245 Hz
Beam energy	E	6500 GeV
Normalized transverse emittance	ϵ	2.3×10^{-6} m
Beta*	β^*	10 m
Full crossing angle (xz)	α	480 μ rad
Transverse plane separation	d	220 μ m
Bunch length	σ_z	0.0787 m (1σ)
RF frequency	f_{RF}	400 MHz ($\Delta_t = 2.5$ ns)

Table A.1: LHC IP2 nominal run setup from, but here with an (ad-hoc) adjusted transverse plane separation (from 357 μ m in lpc.web.cern.ch).

The transverse beam size is given by

$$\sigma_x = \sigma_y = [\epsilon\beta^*]^{1/2} = 57.61 \mu\text{m}. \quad (\text{A.1})$$

The crossing angle affects the effective areas by

$$\Sigma_x = [2\sigma_x^2 \cos^2(\alpha/2) + 2\sigma_z^2 \sin^2(\alpha/2)]^{1/2} = 85.74 \mu\text{m} \quad (\text{A.2})$$

$$\Sigma_y = [2\sigma_y^2]^{1/2} = 81.74 \mu\text{m}. \quad (\text{A.3})$$

These values yield instantaneous luminosity per bunch pair

$$L_{bb} = \frac{f_{rev} N^2 \cos^2(\alpha/2) F}{2\pi \Sigma_x \Sigma_y} = 1.046 \times 10^{28} \text{ Hz/cm}^2, \quad (\text{A.4})$$

where the effective separation factor is $F = \exp\left(-\frac{d^2}{2\Sigma_y^2}\right)$. The instantaneous luminosity integrated over all colliding bunch pairs n_{bb} is $L = n_{bb} L_{bb}$. The numbers here are from the LHC beam division luminosity calculator lpc.web.cern.ch.

APPENDIX A. ALICE MEASUREMENTS

A.4 Trigger statistics

The trigger acronyms are explained in Table A.2. The trigger system ‘vetoes’ or downtime can be due to the CTP (central trigger processor) being busy, the trigger cluster being busy or the trigger system having dead time for other countable reasons in the digital logic or due to a pre-determined trigger downscaling factor. The trigger system takes these into account with before (b) and after (a) counters. For more information about the ALICE trigger system, see [40].

Name	Meaning
ID	Trigger list identifier
BCMask	Bunch cross mask identifier (A,B,C,I,E)
DS	Deterministic down scaling factor (percent)
LMb(a)	Level minus-trigger before(after) trigger ‘vetoes’
L0b(a)	Level 0-trigger before(after) trigger ‘vetoes’
L1b(a)	Level 1-trigger before(after) trigger ‘vetoes’
L2b(a)	Level 2-trigger before(after) trigger ‘vetoes’

Table A.2: Trigger terminology.

ID	Name	BCMask	DS	Duration	LMb	LMa	L0b	L0a	L1b	L1a	L2b	L2a	LMb/Hz	LMa/Hz	L0b/Hz	L0a/Hz	L1b/Hz	L1a/Hz	L2b/Hz	L2a/Hz	LMa/LMb	L0a/L0b	L1a/L1b	L2a/L2b
34	CINT114NOFFCENTNOTRD	B		002431	46484908	46484908	744883	489433	489433	489433	489433	489433	22447	22447	358	236.3	236.3	236.3	236.3	236.3	1	0.66	1	1
35	CINT142NOFFCENTNOTRD	I		002431	46484919	46484919	69300	45437	45437	45437	45437	45437	22447	22447	30.3	22.1	22.1	22.1	22.1	22.1	1	0.67	1	1
36	CINT14ASOFFCENTNOTRD	A	30	002431	83607352	83607352	124765	21729	21729	21729	21729	21729	415264	415264	60.2	10.5	10.5	10.5	10.5	10.5	1	0.17	1	1
37	CINT14CNOFFCENTNOTRD	C	30	002431	83678740	83678740	92587	17574	17574	17574	17574	17574	404040	404040	44.7	8.5	8.5	8.5	8.5	8.5	1	0.19	1	1
38	CINT14ENOFFCENTNOTRD	E		002431	43484892	43484892	1336	663	663	663	663	663	22447	22447	0.63	0.32	0.32	0.32	0.32	0.32	1	0.5	1	1

Table A.3: Trigger statistics for the run 274593.

ID	Name	BCMask	DS	Duration	LMb	LMa	L0b	L0a	L1b	L1a	L2b	L2a	LMb/Hz	LMa/Hz	L0b/Hz	L0a/Hz	L1b/Hz	L1a/Hz	L2b/Hz	L2a/Hz	LMa/LMb	L0a/L0b	L1a/L1b	L2a/L2b
34	CINT114NOFFCENTNOTRD	B		002558	34983252	34983252	60896	31922	31922	31922	31922	31922	22447	22447	44.7	239.4	239.4	239.4	239.4	239.4	1	0.61	1	1
35	CINT142NOFFCENTNOTRD	I		002558	34983253	34983253	62926	38450	38450	38450	38450	38450	22447	22447	40.3	24.7	24.7	24.7	24.7	24.7	1	0.61	1	1
36	CINT14ASOFFCENTNOTRD	A	30	002558	64718012	64718012	127846	21336	21336	21336	21336	21336	415265	415265	82	13.7	13.7	13.7	13.7	13.7	1	0.17	1	1
37	CINT14CNOFFCENTNOTRD	C	30	002558	62969543	62969543	102327	39388	39388	39388	39388	39388	404041	404041	65.7	12.4	12.4	12.4	12.4	12.4	1	0.19	1	1
38	CINT14ENOFFCENTNOTRD	E		002558	34983253	34983253	2933	1557	1557	1557	1557	1557	22447	22447	1.9	1	1	1	1	1	1	0.53	1	1

Table A.4: Trigger statistics for the run 274594.

ID	Name	BCMask	DS	Duration	LMb	LMa	L0b	L0a	L1b	L1a	L2b	L2a	LMb/Hz	LMa/Hz	L0b/Hz	L0a/Hz	L1b/Hz	L1a/Hz	L2b/Hz	L2a/Hz	LMa/LMb	L0a/L0b	L1a/L1b	L2a/L2b
34	CINT114NOFFCENTNOTRD	B		002950	40188874	40188874	1623287	778911	778911	778911	778911	778911	22447	22447	906.7	435	435	435	435	435	1	0.48	1	1
35	CINT142NOFFCENTNOTRD	I		002950	40188874	40188874	145953	69786	69786	69786	69786	69786	22447	22447	81.5	39	39	39	39	39	1	0.48	1	1
36	CINT14ASOFFCENTNOTRD	A	30	002950	73209470	73209470	136000	13056	13056	13056	13056	13056	415265	415265	84.8	7.7	7.7	7.7	7.7	7.7	1	0.22	1	1
37	CINT14CNOFFCENTNOTRD	C	30	002950	723399732	723399732	80662	11244	11244	11244	11244	11244	404041	404041	48.1	6.3	6.3	6.3	6.3	6.3	1	0.33	1	1
38	CINT14ENOFFCENTNOTRD	E		002950	40188874	40188874	1237	457	457	457	457	457	22447	22447	0.69	0.26	0.26	0.26	0.26	0.26	1	0.37	1	1

Table A.5: Trigger statistics for the run 274595.

B GRANIITI and kinematics

B.1 Alternative models of pomeron

The original pomeron

The original ‘sliding helicity’ Gribov pomeron has been implemented in the code partially – this is only partial because the calculations involve unknown functions at the central vertex. However, we shall explore it a bit. In the following, we shall strip off all the form factors and concentrate only on the helicity dependent terms. The helicity amplitude is written as [111]

$$\begin{aligned}
 & T_{\lambda_A \lambda_B}^{\lambda_1 \lambda_3 \lambda_2}(s_1, s_2, t_1, t_2, \varphi) \\
 &= \sum_{ab} g_{\lambda_A \lambda_1}^a(t_1) \Delta_a(s_1, t_1) g_{ab}^{\lambda_3}(t_1, t_2, \varphi) \Delta_b(s_2, t_2) g_{\lambda_B \lambda_2}^b(t_2), \quad (\text{B.1})
 \end{aligned}$$

where the sum runs over exchanged Reggeon with propagators Δ_i as defined earlier, to first order at high energies, over only two pomerons. The particle-Reggeon-particle vertex helicity structure at small $-t$ is

$$g_{\lambda_A \lambda_1}^a(t) \simeq (-t)^{|\lambda_A - \lambda_1|/2}. \quad (\text{B.2})$$

The spin structure of the crucial central vertex is [111]

$$g_{ab}^{\lambda_3}(t_1, t_2, \varphi) = \sum_{m_1, m_2 = -\infty}^{\infty} e^{im_1 \varphi} \gamma_{m_1 m_2}^{\lambda_3}(t_1, t_2), \quad \text{subject to } \lambda_3 = m_1 - m_2, \quad (\text{B.3})$$

where

$$\gamma_{m_1 m_2}^{\lambda_3}(t_1, t_2) \simeq (-t_1)^{|m_1|/2} (-t_2)^{|m_2|/2}, \quad (\text{B.4})$$

a limit which holds for small values of $-t_1$ and $-t_2$. The projection of the angular momentum $j_{a(b)}$ of the Reggeon $a(b)$ is denoted with $m_{1(2)}$, which is an analytical continuation over all values over the sliding Regge trajectory, thus technically an infinite sum. However, the sum over $m_{1(2)}$ should be truncated to a first few, given that

$-t_{1(2)}$ are small and there is no full information about the central vertex available. We make remark that one can see easily that in order the parity conservation to hold for all spin-parities, the vertex Eq. B.4 cannot be always symmetric e.g. under $(m_1, m_2) \leftrightarrow (-m_1, -m_2)$, but should change sign – thus this equation is a slightly formal one. Now we have hopefully emphasized enough that the original pomeron has no fixed spin structure but is an infinite sum. It is essentially the spin-parity of the central state which ‘forces’ the pomeron Lorentz structure to look like a non-conserved or conserved vector (or tensor) current, from a Regge theory point of view, as discussed in [111].

Vector current pomeron

This is a model from [234], which illustrates many points and is described here shortly for completeness. In some sense the Tensor pomeron model is a superset of this with complete Feynman rules and higher rank spin structure. The spin density matrix of the i -th pomeron in this conserved vector current model is obtained as

$$\rho_i^{\lambda\lambda'} = (-1)^{\lambda+\lambda'} \epsilon^\mu(q_i, \lambda) \epsilon^\nu(q_i, \lambda') \rho^{\mu\nu}, \quad (\text{B.5})$$

where $\lambda = 0, \pm 1$ and the space-like $q^2 < 0$ spin-1 polarization vectors are denoted with ϵ . On the right hand side, the covariant density matrix is hermitian $\rho^{\mu\nu*} = \rho^{\nu\mu}$ and the spin-1 polarization vectors obey

$$\epsilon^{\mu*}(\pm 1) = -\epsilon^\mu(\mp 1) \quad (\text{B.6})$$

$$\epsilon^{\mu*}(0) = -\epsilon^\mu(0). \quad (\text{B.7})$$

We denote the 4 independent elements of the density matrix by ρ^{++}, ρ^{00} , which are on the diagonal, and the off-diagonal terms by $|\rho^{+0}|e^{i\tilde{\varphi}_i}$ and $|\rho^{+-}|e^{i\tilde{\varphi}_i}$, where $\tilde{\varphi}_i$ is the forward proton azimuthal angle in the pomeron-pomeron CM frame. This angle is not exactly the same, clearly, as the forward angle φ in the proton-proton CM frame or LHC lab frame.

The unnormalized covariant density matrix for the pomeron emission from the i -th proton leg is

$$\rho_i^{\mu\nu} = -\frac{1}{q_i^2} \sum_{\text{helicities}} J_\mu J_\nu^*, \quad (\text{B.8})$$

where J_μ is a Dirac like current of the proton-pomeron-proton vertex as the one in QED given by Eq. 5.48. However, pomeron may couple with different strengths than photon to the electric and magnetic terms of the proton form factors. The differential

APPENDIX B. GRANIETTI AND KINEMATICS

cross section is then obtained by coupling $\rho_1^{\lambda_1\lambda'_1}$ and $\rho_2^{\lambda_2\lambda'_2}$ elements together at the central vertex [234]

$$\begin{aligned}
 d\sigma \sim & 2\rho_1^{++}\rho_2^{++} \times [W(++++) + W(+-,+-)] \\
 & + 2\rho_1^{++}\rho_2^{00}W(+0,+0) + 2\rho_1^{00}\rho_2^{++}W(0+,0+) \\
 & + \rho_1^{00}\rho_2^{00}W(00,00) \\
 & + 2|\rho_1^{+-}\rho_2^{+-}|W(++,-) \cos 2\tilde{\varphi} \\
 & - 4|\rho_1^{+0}\rho_2^{+0}| \times [W(+++,00) + W(0+,-)] \cos \tilde{\varphi},
 \end{aligned} \tag{B.9}$$

which depends on eight helicity structure functions $W(\lambda_1\lambda_2, \lambda'_1\lambda'_2)$, where the helicity conservation requires $\lambda_1 - \lambda_2 = J_z = \lambda'_1 - \lambda'_2$, otherwise $W = 0$. The azimuthal angle is $\tilde{\varphi} = \tilde{\varphi}_1 + \tilde{\varphi}_2$. Six of these structure function can be measured with unpolarized initial state protons. Now, let us write this in terms of the helicity amplitudes

$$W(\lambda_1\lambda_2, \lambda'_1\lambda'_2) \sim A_{\lambda_1\lambda_2}(t_1, t_2)A_{\lambda'_1\lambda'_2}^*(t_1, t_2)\delta((q_1 + q_2)^2 - M^2). \tag{B.10}$$

Parity conservation requires that $A_{-\lambda_1-\lambda_2} = \eta_1\eta_2\eta_M A_{\lambda_1\lambda_2}$, where $\eta_{1,2}$ denote the naturality of pomerons and η_M the naturality of the boson, which is $+1$ if $P = (-1)^J$ and -1 if $P = (-1)^{J-1}$. Bose-Einstein symmetry (statistics) requires $A_{\lambda_1\lambda_2}(t_1, t_2) = (-1)^J A_{\lambda_2\lambda_1}(t_2, t_1)$, and time invariance requires invariance under under $\lambda_1\lambda_2 \leftrightarrow \lambda'_1\lambda'_2$, which gives somewhat more complicated relations [234]. This model has some definite properties regarding the observables. Most of them are direct consequences of the parity conservation and vector currents.

B.2 Kinematics of $2 \rightarrow 3$

Lorentz invariants

Let us have the $2 \rightarrow 3$ process

$$p_A + p_B \rightarrow p_1 + P_3 + p_2, \quad (\text{B.11})$$

where p_1 and p_2 are the forward systems. Kinematically this is fully characterized by 5 linearly independent Lorentz scalars. The most typical set is $s = (p_A + p_B)^2 = (p_1 + P_3 + p_2)^2$, $t_1 = (p_A - p_1)^2$, $t_2 = (p_B - p_2)^2$, $s_1 = (p_1 + P_3)^2$, $s_2 = (p_2 + P_3)^2$.

Colliding beam frame

Now, we need (at least) 5 variables. For practical reasons, we will use 6. The redundancy is between forward system transverse angles φ_1 and φ_2 which we can remove by the rotational invariance and use only the difference $\Delta\varphi \equiv \varphi_1 - \varphi_2$, whenever necessary. We define 6 variables, which fully characterize the exact kinematics and separate between the longitudinal and transverse degrees of freedom

$$\xi_{1(2)} \equiv 1 - \frac{p_{z,1(2)}}{p_{z,A(B)}} \in [0, 1] \quad (\text{B.12})$$

$$p_{t,1(2)} \equiv \left(p_{x,1(2)}^2 + p_{y,1(2)}^2 \right)^{1/2} \in \mathbb{R}_+ \quad (\text{B.13})$$

$$\varphi_{1(2)} \equiv \tan^{-1} \left(\frac{p_{y,1(2)}}{p_{x,1(2)}} \right) \in (-\pi, \pi]. \quad (\text{B.14})$$

Due to the 4-momentum conservation, not all arbitrary combinations of the values of these variables are physically valid.

4-momentum components

The momentum $p^\mu = [E, \vec{p}]$ of forward systems are now given directly by

$$p_{1(2)}^\mu = \left[\left(M_{1(2)}^2 + p_{t,1(2)}^2 + (-m_p^2 + \frac{s}{4})(\xi_{1(2)} - 1)^2 \right)^{1/2}, \right. \\ \left. p_{t,1(2)} \cos(\varphi_{1(2)}), p_{t,1(2)} \sin(\varphi_{1(2)}), \mp \frac{(\xi_{1(2)} - 1)(-4m_p^2 + s)^{1/2}}{2} \right]. \quad (\text{B.15})$$

We allow also excitation of forward protons with masses M_1 and M_2 . For the elastic forward protons, we set $M_{1(2)} \equiv m_p$. The central system 4-momentum is

$$\begin{aligned}
 P_3^\mu &= q_1^\mu + q_2^\mu = \\
 &[s^{1/2} - \left(M_1^2 + p_{t,1}^2 + (-m_p^2 + \frac{s}{4})(\xi_1 - 1)^2\right)^{1/2} \\
 &- \left(M_2^2 + p_{t,2}^2 + (-m_p^2 + \frac{s}{4})(\xi_2 - 1)^2\right)^{1/2}, \\
 &- p_{t,1} \cos(\varphi_1) - p_{t,2} \cos(\varphi_2), -p_{t,1} \sin(\varphi_1) - p_{t,2} \sin(\varphi_2), \frac{(\xi_1 - \xi_2)(-4m_p^2 + s)^{1/2}}{2}].
 \end{aligned} \tag{B.16}$$

4-momentum transfer squared

The Lorentz scalars $t_{1(2)} < 0$, which are the 4-momentum transfer squared, are written as

$$\begin{aligned}
 t_{1(2)} &= \frac{\left(\left((\xi_{1(2)} - 1)^2(-4m_p^2 + s) + 4M_{1(2)}^2 + 4p_{t,1(2)}^2\right)^{1/2} - s^{1/2}\right)^2}{4} \\
 &- \frac{\xi_{1(2)}^2(-4m_p^2 + s)}{4} - p_{t,1(2)}^2.
 \end{aligned} \tag{B.17}$$

In the limit $\xi_1, \xi_2 \rightarrow 0$ and $s \rightarrow \infty$, we get the familiar

$$\boxed{t_{1(2)} \simeq -p_{t,1(2)}^2}. \tag{B.18}$$

This approximation is very good. When $p_{t,1(2)}^2 \lesssim 1 \text{ GeV}^2$, the relative error we make is proportional to the scale of ξ and for low mass CEP at the LHC, this scale is $\xi \sim 10^{-4}$.

Central system rapidity

The boost (rapidity) definition $Y_X = \frac{1}{2} \ln \left(\frac{E+p_z}{E-p_z} \right)$ along the beam line gives

$$\begin{aligned}
 Y_X = \frac{1}{2} \ln [& \left((M_1^2 + p_{t,1}^2 + (-m_p^2 + \frac{s}{4})(\xi_1 - 1)^2)^{1/2} + \right. \\
 & \left. \left(M_2^2 + p_{t,2}^2 + (-m_p^2 + \frac{s}{4})(\xi_2 - 1)^2 \right)^{1/2} - \frac{(\xi_1 - \xi_2)(-4m_p^2 + s)^{1/2}}{2} - s^{1/2} \right] \\
 & / [\left((M_1^2 + p_{t,1}^2 + (-m_p^2 + \frac{s}{4})(\xi_1 - 1)^2)^{1/2} + \right. \\
 & \left. \left(M_2^2 + p_{t,2}^2 + (-m_p^2 + \frac{s}{4})(\xi_2 - 1)^2 \right)^{1/2} + \frac{(\xi_1 - \xi_2)(-4m_p^2 + s)^{1/2}}{2} - s^{1/2} \right],
 \end{aligned}$$

where taking massless limits and collinear $p_{t,1(2)} \rightarrow 0$ gives

$$\boxed{Y_X \simeq \frac{1}{2} \ln \left(\frac{\xi_1}{\xi_2} \right)}. \quad (\text{B.19})$$

This is all around a very good approximation, relative error $10^{-3} \dots 10^{-6}$.

Subsystem energy invariants

The scalars $s_1 = (p_1 + P_3)^2 > 0$ and $s_2 = (p_2 + P_3)^2 > 0$ are

$$s_1 = s - 2s^{1/2} \left(M_2^2 + p_{t,2}^2 + (-m_p^2 + \frac{s}{4})(\xi_2 - 1)^2 \right)^{1/2} + M_2^2 \quad (\text{B.20})$$

$$s_2 = s - 2s^{1/2} \left(M_1^2 + p_{t,1}^2 + (-m_p^2 + \frac{s}{4})(\xi_1 - 1)^2 \right)^{1/2} + M_1^2. \quad (\text{B.21})$$

The obvious massless and $p_{t,1(2)} \rightarrow 0$ limits of these are

$$\boxed{s_1 \simeq s\xi_2, \quad s_2 \simeq s\xi_1}. \quad (\text{B.22})$$

Relations above are sometimes also called as ‘Regge domain criteria’, equivalent with $s \gg |t_{1(2)}|, s \rightarrow \infty$. Note how terms are crossed $1 \leftrightarrow 2$, due to the momentum flow. The approximations are very good for high ξ values and even at $\xi \sim 10^{-6}$, the relative error is order of 10^{-3} . In terms of the central system invariant mass and rapidity, we can also write

$$\boxed{s_1 \simeq M_X s^{1/2} \exp(-Y_X), \quad s_2 \simeq M_X s^{1/2} \exp(Y_X)}, \quad (\text{B.23})$$

often used in the case of photoproduction. These can be derived using using the collinear relations

$$\boxed{\xi_1 \simeq \frac{M_X}{s^{1/2}} \exp(-Y_X), \quad \xi_2 \simeq \frac{M_X}{s^{1/2}} \exp(Y_X)} \quad (\text{B.24})$$

in the limit $p_{t,1(2)} \rightarrow 0$.

Central system invariant mass

The central system invariant mass squared $M_X^2 \equiv P_3^2 = (q_1 + q_2)^2$, is exactly written as

$$\begin{aligned} M_X^2 = & \left[\left(M_1^2 + p_{t,1}^2 + \left(-m_p^2 + \frac{s}{4} \right) (\xi_1 - 1)^2 \right)^{1/2} \right. \\ & + \left(M_2^2 + p_{t,2}^2 + \left(-m_p^2 + \frac{s}{4} \right) (\xi_2 - 1)^2 \right)^{1/2} - s^{1/2} \left. \right]^2 \\ & - (p_{t,1} \cos(\varphi_1) + p_{t,2} \cos(\varphi_2))^2 - (p_{t,1} \sin(\varphi_1) + p_{t,2} \sin(\varphi_2))^2 \\ & - \frac{(\xi_1 - \xi_2)^2 (-4m_p^2 + s)}{4}. \end{aligned} \quad (\text{B.25})$$

Now if we take limit $p_{t,1(2)} \rightarrow 0$ and also $m_p, M_1, M_2 \rightarrow 0$, that is, $s \gg m_p^2, M_1^2, M_2^2$, we get

$$\boxed{M_X^2 \simeq \xi_1 \xi_2 s.} \quad (\text{B.26})$$

This approximation is *not* especially good for the low mass CEP, because the $p_{t,1,2}$ scale is close to the central system mass scale. The relative error can be order of one ($\sim 100\%$). Note that the complement scalar is

$$s_{12} = (p_1 + p_2)^2 \simeq (\xi_1 - 1)(\xi_2 - 1)s, \quad (\text{B.27})$$

which is not in common use.

Spanning set of variables

We show here a generic method for checking if the constructed set of Lorentz scalars is a spanning set. In a row order s, t_1, t_2, s_1, s_2 , we construct a matrix

$$A = \begin{bmatrix} 1 & 1 & 0 & 0 & 0 \\ 1 & 0 & -1 & 0 & 0 \\ 0 & 1 & 0 & -1 & 0 \\ 0 & 0 & 1 & 0 & 1 \\ 0 & 0 & 0 & 1 & 1 \end{bmatrix}, \quad (\text{B.28})$$

APPENDIX B. GRANIITI AND KINEMATICS

where each row/column correspond to the corresponding 4-momenta indices with order p_A, p_B, p_1, p_2, P_3 and the matrix element is $A_{ij} = \pm$ when $(p_i \pm p_j)^2$. This matrix should have full rank \Leftrightarrow a non-zero determinant. Computer algebra gives full rank(A) = 5 and determinant -2, which fine.

How about the 4-momentum representations? For this, we can construct the Gram matrix between all 4-dot products $G_{ij} = \langle p_i, p_j \rangle$. The determinant of this matrix is identically zero, because we have 5 vectors in 4-dimensional space. Thus, the check goes for example by choosing (only) the final states (p_3, p_4, p_5) , and calculating the 3×3 Gram matrix, and its determinant. The determinant obtained is an extremely lengthy expression, and gives zero for *linearly dependent final state configurations*, which is fine. An example of a linearly dependent case: $(\xi_1 \equiv \xi_2, p_{t,1} \equiv p_{t,2}, \Delta\varphi \equiv \varphi_1 - \varphi_2 \equiv \pi) \rightarrow$ system produced at rest gives $\det(G) = 0$.

B.3 The slope parameter

Experimentally in soft processes $d\sigma/dt \sim \exp(-b|t|)$ holds for small values of $|t|$. For example, in photoproduction the pomeron exchange t -distribution slope is often written as

$$b = b_0 + 2\alpha' \ln \left(\frac{W^2}{W_0^2} \right), \quad (\text{B.29})$$

where b_0 is a process (e.g. vector meson mass) dependent constant, W^2 is the γ^*p subsystem invariant such as s_1 or s_2 here and W_0^2 is the (fit) normalization scale squared, often $W_0 \simeq 90$ GeV in HERA fits. The second term with the pomeron slope $\alpha' \simeq 0.05 \dots 0.25$ GeV $^{-2}$ originates from the Regge phenomenology, where as the first term is simply based on an exponential ansatz. That is, this exponential ansatz may encapsulate both the proton form factor and the vector meson form factor or transverse size, however, a model dependent factorization ansatz between them is also possible.

The photon side t -distribution is different and is driven by QED photon singularity $1/q^2$ and proton EM-form factors. The slope parameter has dimensions (GeV $^{-2}$) which is the same as the cross section and may be interpreted as the average size of the transverse interaction region, which naively could do logarithmic grow infinitely, perhaps without confinement. So if α' is the ‘inverse string tension’ coupled with the logarithmically growing Gribov diffusion term, what is the constant b_0 term then? Perhaps it is the Fermi $\langle k_t^2 \rangle \sim 1/b$ of the ‘vacuum noise’, by definition.

Slope inference fit

We shall make a remark that an exponential $-t \simeq p_t^2$ -distribution results in a Rayleigh distributed p_t -distribution, which then means that p_x and p_y components are following a Gaussian stochastic process with a variance σ^2 . The relation between the exponential and the Rayleigh can be shown directly by the change of a variable theorem

$$f(p_t^2) = b \exp(-bp_t^2) \Leftrightarrow g(p_t) = \left| \frac{dp_t^2}{dp_t} \right| f(p_t^2) = 2p_t b \exp(-bp_t^2). \quad (\text{B.30})$$

Using this, the relation between the parameter b and the forward system average transverse momentum $\langle p_t \rangle$ with the used distribution assumptions is

$$b = \frac{\pi}{4\langle p_t \rangle^2} = \frac{1}{2\sigma^2}. \quad (\text{B.31})$$

APPENDIX B. GRANITTI AND KINEMATICS

We note that the square here is of the average, not the other way around.

How to obtain an estimate of b without forward proton measurements? First assume that the forward proton separation $\Delta\varphi$ follows a specific distribution, such as the uniform, generate forward proton p_t^2 values according to the chosen b , construct the forward proton transverse momentum vectors $\vec{p}_{t,1(2)}$, take the sum $|\vec{P}_t| = |\vec{p}_{t,1} + \vec{p}_{t,2}|$ and compare this sample with the measured central system transverse momentum $|\vec{P}_t|$ -distribution in a fit loop. Extension of this is to take into account the forward proton dissociation, which results qualitatively in a different distribution presumably with a hard (point like) power law tail, perhaps compatible with a Lévy flight like stochastic process in the transverse plane. Clearly, one needs to take into account in the fit the mixture of elastic-elastic, elastic-inelastic, inelastic-inelastic events. The most extensive fits are most easily done by generating samples using the full GRANITTI machinery.

B.4 Harmonic acceptance decompositions

APPENDIX B. GRANITTI AND KINEMATICS

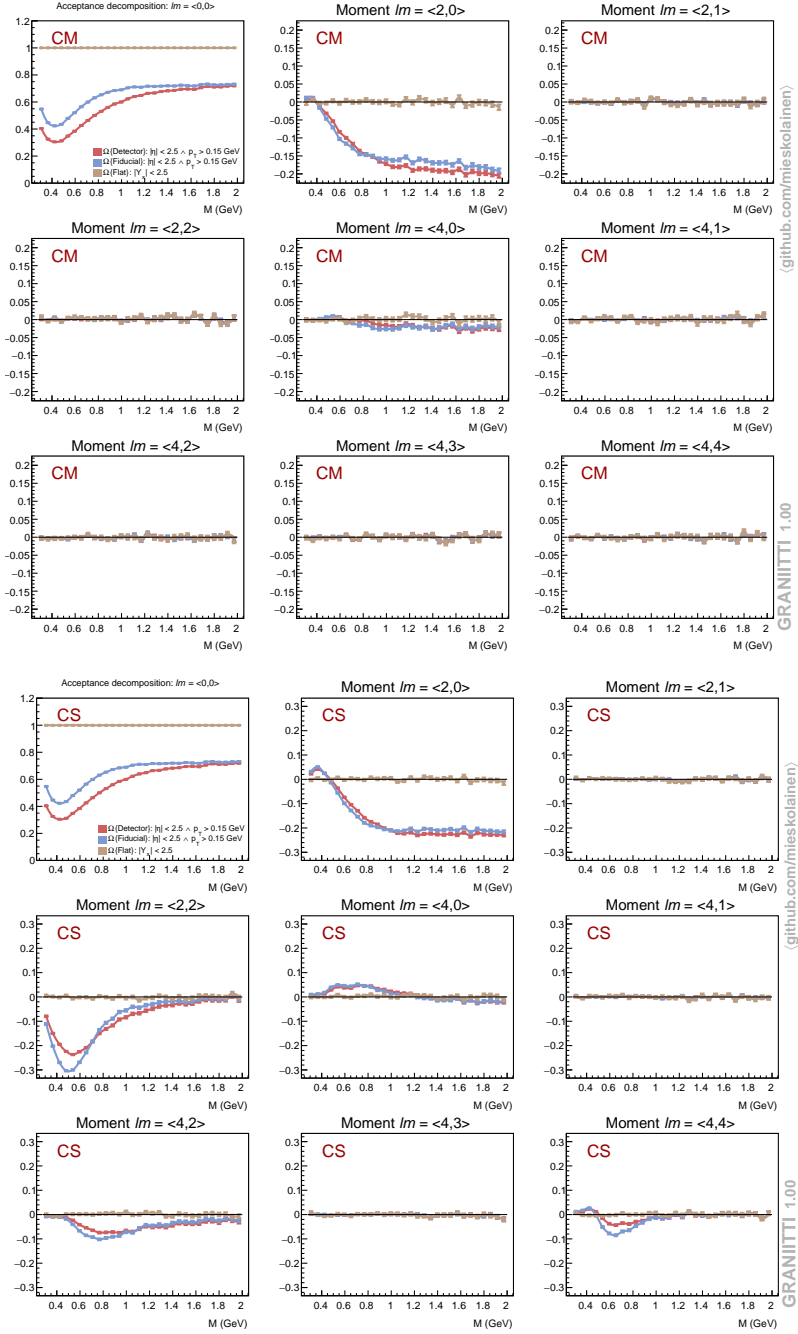


Figure B.1: CM frame (up) and CS frame (down): Acceptance decomposition.

APPENDIX B. GRANITTI AND KINEMATICS

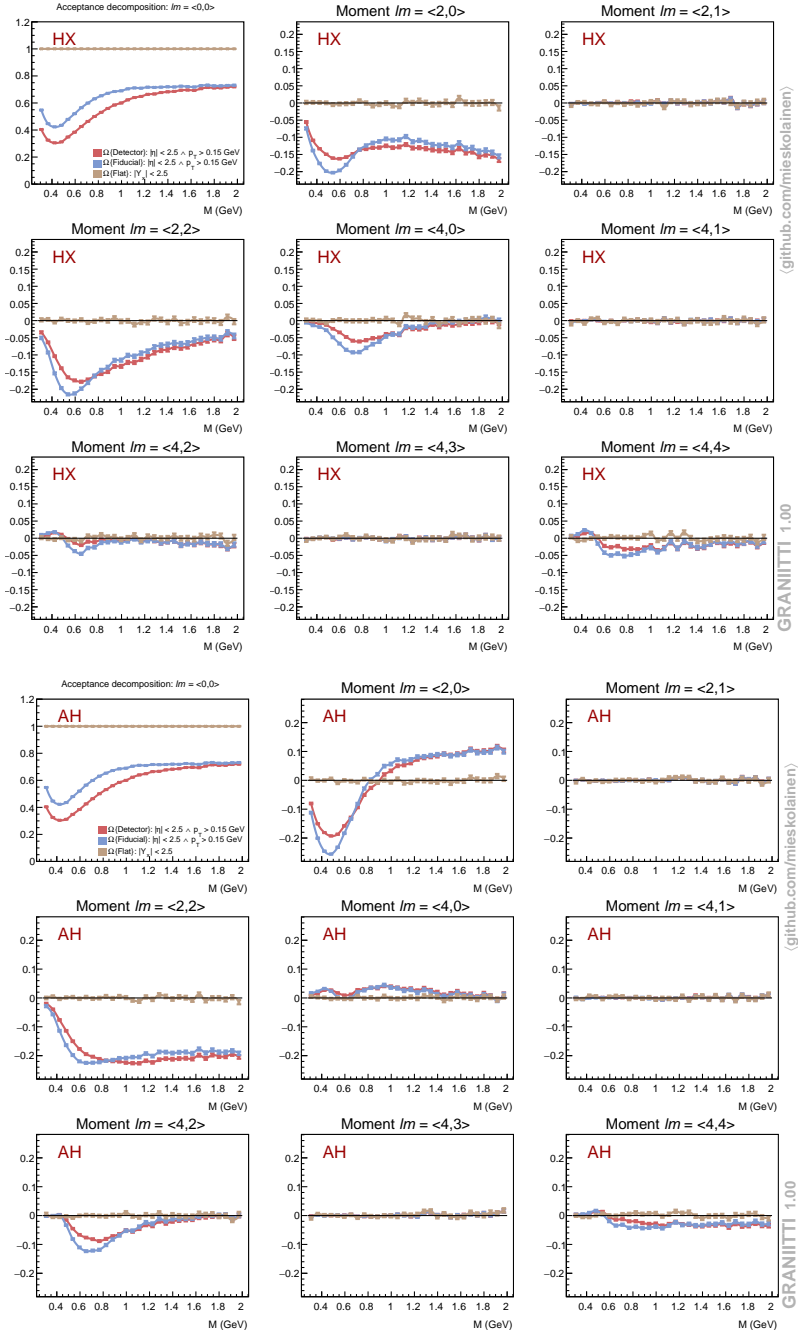


Figure B.2: HX frame (up) and AH frame (down): Acceptance decomposition.

APPENDIX B. GRANITTI AND KINEMATICS

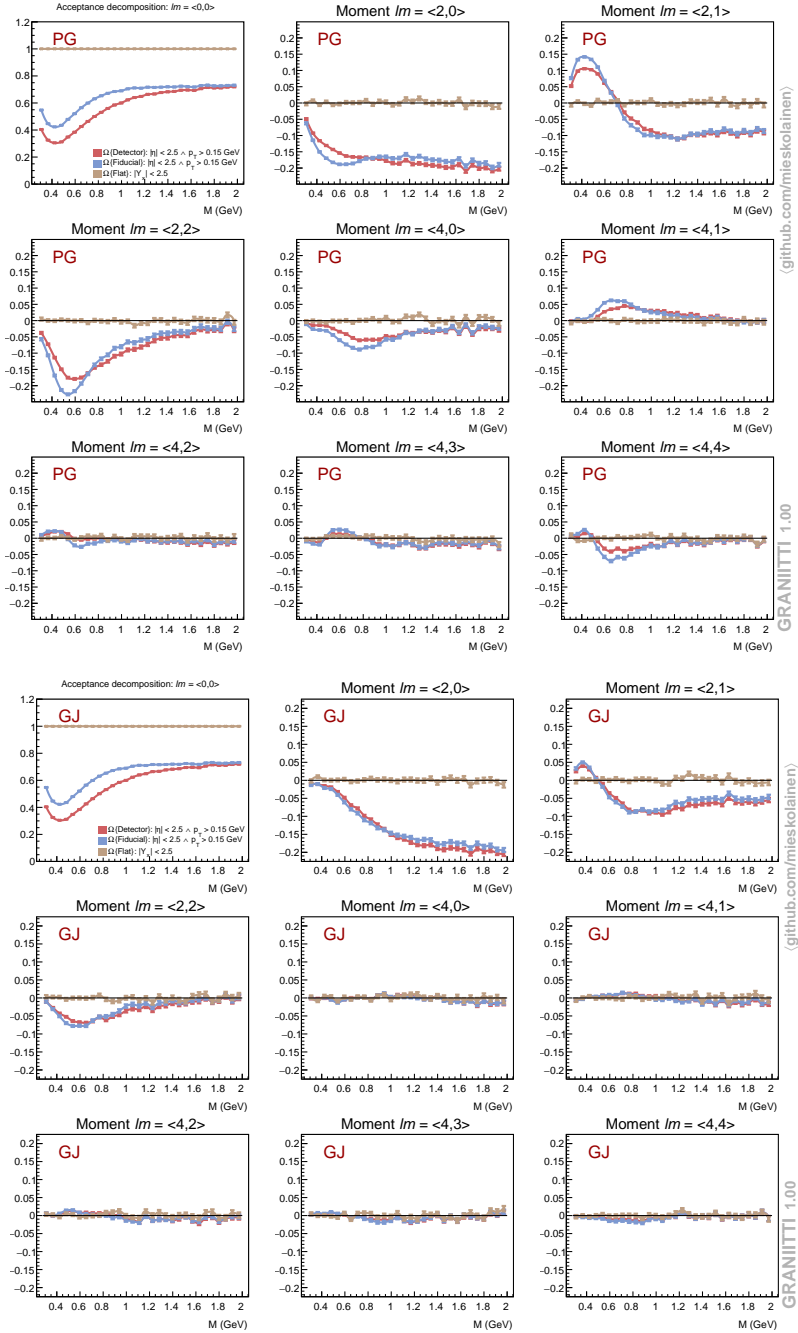


Figure B.3: PG frame (up) and GJ frame (down): Acceptance decomposition.

C Combinatorics and pileup

C.1 Vector space subspaces over finite fields

The q -binomial coefficient counts the number of subspaces of dimension r in a vector space of dimension N , when the vector space is over a finite field with a prime power q , is

$$\begin{bmatrix} N \\ r \end{bmatrix}_q = \begin{cases} \frac{(1-q^N)(1-q^{N-1})\dots(1-q^{N-r+1})}{(1-q)(1-q^2)\dots(1-q^r)}, & \text{if } r \leq N. \\ 0, & r > N. \end{cases} \quad (\text{C.1})$$

These subspaces are encapsulated in the Grassmannian manifold $\text{Gr}(r, N, \mathbb{F}_q)$ with $\dim(\text{Gr}) = r(N - r)$.

$N \setminus r$	0	1	2	3	4	5	6	7	8	Σ
1	1	1	0	0	0	0	0	0	0	2
2	1	3	1	0	0	0	0	0	0	5
3	1	7	7	1	0	0	0	0	0	16
4	1	15	35	15	1	0	0	0	0	67
5	1	31	155	155	31	1	0	0	0	374
6	1	63	651	1395	651	63	1	0	0	2825
7	1	127	2667	11811	11811	2667	127	1	0	29212
8	1	255	10795	97155	200787	97155	10795	255	1	417199

Table C.1: Total number of r -dimensional subspaces of \mathbb{F}_2^N .

The total number of subspaces is obtained with $\sum_{r=0}^N \begin{bmatrix} N \\ r \end{bmatrix}_q$. This count gives also the so-called simplicity of the N -cube in a minimal corner-cut triangulation of the cube, which is geometrically intuitive. For $q = 2$ and $N = 1, \dots, 8$, these are listed in Table C.1. Note that r starts from zero.

C.2 Pileup combinatorics

Let $M : \Theta \rightarrow Y$ be a probabilistic or stochastic mixing matrix

$$\mathbf{y} = M(\mathbf{p}, \mu)\mathbf{p}, \quad \sum_i M_{ij} = 1 \quad \forall j \tag{C.2}$$

which has each column with unit sum and size $2^N - 1$. Each matrix element

$$M_{ij} \equiv P(c = i | c = j) \tag{C.3}$$

gives the conditional probability of an event vector originating from the j -th final state class c propagating to the i -th class, due to pileup. The matrix is thus interpreted vertically, which makes it a left stochastic matrix. This matrix is illustrated in Figure C.1.

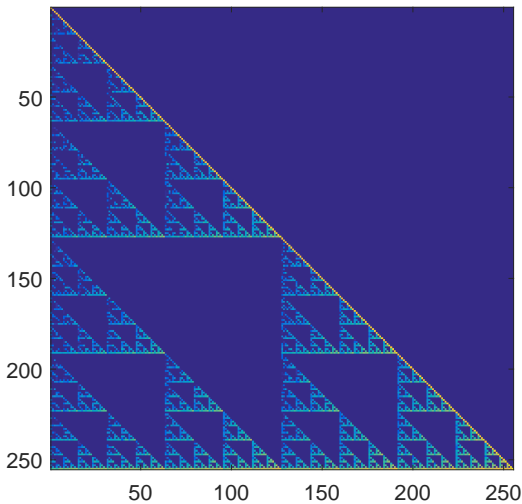


Figure C.1: Simulated pileup mixing matrix $M_{255 \times 255}$ for $N = 8$ and $\mu = 1.5$.

The matrix was obtained by a toy Monte Carlo pileup simulation using Poisson random numbers with a mean μ and by generating different final states according to uniform probabilities $p_c \sim 1/n$. The mixing matrix was obtained by counting the mixed final states event by event. Due to the permutation ambiguity, the originating class c was selected always to be the first of the random list, thus randomly. Because

APPENDIX C. COMBINATORICS AND PILEUP

the matrix M here is a function of the (unknown) \mathbf{p} , and also μ , it is not directly useful as an algorithmic inverse solution. Its function is to demonstrate algebraic properties. In general, the combinatorial enumerations can be approached by the ‘twelve-fold way’, which is a systematic approach developed by G.C. Rota. Here we list some combinatorial aspects of this problem to make it more transparent.

The number of non-zero elements in the pile-up mixing matrix per row are given by

$$a(i) = \sum_{r=1}^i \binom{i}{r} \bmod 2, \quad (\text{C.4})$$

where mod denotes here the remainder after division and i is the index of the row. The sequence for the first 15 rows is

$$1, 1, 3, 1, 3, 3, 7, 1, 3, 3, 7, 3, 7, 7, 15, \quad i = 1, \dots, 15. \quad (\text{C.5})$$

Similarly, the number of 1 in the binary vector representation, or the so-called Hamming weight, is given by adding one and taking base-2 logarithm of Equation C.4

$$a(i) = \log_2 \left(\sum_{r=1}^i \binom{i}{r} \bmod 2 + 1 \right). \quad (\text{C.6})$$

As an example

$$1, 1, 2, 1, 2, 2, 3, 1, 2, 2, 3, 2, 3, 3, 4, \quad i = 1, \dots, 15. \quad (\text{C.7})$$

This sequence, namely the position of 1 in this sequence, give us directly the binary vectors which undergo only autocompound. That is, their signature cannot be imitated by a linear combination of other vectors. An example of this is given in Table C.3, for example when $N = 3$, these are $c = 1, 2$ and 4. These are the unit basis vectors of \mathbb{F}_2^N . We also see the ‘fractal’ structure, that is, increasing the dimension of binary space keeps the multiplicity structure from lower dimensions which is evident from Table C.3. This is clear given the Hamming weight sequence.

The number of different terms for $c = 3$ (or 5, 6, 9, ...) follows from a sequence

$$a(k) = \binom{k+1}{2} + k - 1 = \left\{ \begin{matrix} k+1 \\ k \end{matrix} \right\} + k - 1, \quad (\text{C.8})$$

which generates a multiplicity sequence 1, 4, 8, 13, 19, 26, ... and the curly brackets on right give the Stirling partition number (Stirling number of the second kind). This is the base of all higher dimensional binary vector space sequences. All the higher order

APPENDIX C. COMBINATORICS AND PILEUP

$\#_{(N,k)}$	1	2	3	4	5	6
2	3	6	10	15	21	28
3	7	28	84	210	462	924
4	15	120	680	3060	11628	38760
5	31	496	5456	46376	324632	1947792
6	63	2016	43680	720720	9657648	109453344
7	127	8128	349504	11358880	297602656	6547258432
8	255	32640	2796160	180352320	9342250176	404830840960
9	511	130816	22369536	2874485376	296071993728	25462191460608
10	1023	523776	178956800	45902419200	9428356903680	1615391816163840

Table C.2: Combinatorial total multiplicities for $N \in [2, 10]$ and $k \in [1, 6]$.

multiplicity matrices are self repeating results of this sequence or identity sequence $1, 1, \dots$, which follows the algebraic binary expansion ordering. An example of this is given in Table C.3.

The total combinatorial or multinomial multiplicity (multiset multiplicity) is

$$\#_{(N,k)} = \binom{(2^N - 1) + k - 1}{(2^N - 1) - 1} \tag{C.9}$$

which is tabulated in Table C.2 for a finite number of N and k values. The factorial growth is obvious for high Poisson orders and binary dimensions.

APPENDIX C. COMBINATORICS AND PILEUP

$N = 2$	k					
c	1	2	3	4	5	6
1	1	1	1	1	1	1
2	1	1	1	1	1	1
3	1	4	8	13	19	26
Σ	3	6	10	15	21	28
$N = 3$	k					
c	1	2	3	4	5	6
1	1	1	1	1	1	1
2	1	1	1	1	1	1
3	1	4	8	13	19	26
4	1	1	1	1	1	1
5	1	4	8	13	19	26
6	1	4	8	13	19	26
7	1	13	57	168	402	843
Σ	7	28	84	210	462	924
$N = 4$	k					
c	1	2	3	4	5	6
1	1	1	1	1	1	1
2	1	1	1	1	1	1
3	1	4	8	13	19	26
4	1	1	1	1	1	1
5	1	4	8	13	19	26
6	1	4	8	13	19	26
7	1	13	57	168	402	843
8	1	1	1	1	1	1
9	1	4	8	13	19	26
10	1	4	8	13	19	26
11	1	13	57	168	402	843
12	1	4	8	13	19	26
13	1	13	57	168	402	843
14	1	13	57	168	402	843
15	1	40	400	2306	9902	35228
Σ	15	120	680	3060	11628	38760

Table C.3: Combinatorial multiplicities for each individual binary vector (unit-hypercube vertex) with $N \in [2, 4]$ and $k \in [1, 6]$.

C.3 Exclusive efficiency

A term *exclusive efficiency* is often used in the context of pileup and large pseudo-rapidity gap veto based measurements of central diffraction. The probability for an event to be exclusive in a bunch cross is simply obtained from the zero-truncated Poisson distribution

$$P_P^{>0}(k; \mu) = P(X = k | X > 0) = \frac{P_P(k = 1; \mu)}{1 - P_P(k = 0; \mu)} = \frac{\mu^k}{(e^\mu - 1)k!}. \quad (\text{C.10})$$

Then setting $k = 1$, which is the probability of one visible interaction, gives us the value of interest $P_P^{>0}(k = 1; \mu) = \mu/(e^\mu - 1)$. This runs asymptotically to zero when $\mu \rightarrow \infty$. The exclusive efficiency should not be mixed with the fundamental probability of rapidity gap survival, such as random re-scattering or production of secondaries filling the rapidity gap in the elementary pp -interaction. Neither it should be mixed with signal selection efficiency or direct veto *inefficiency*, these are either simulation based or control sample based corrections. The signal selection efficiency is driven by tracking and trigger efficiency whereas veto inefficiency is driven by efficiency of (forward) detectors to see all soft particles. A signal efficiency loss is simply of type multiplicative efficiency correction, the veto inefficiency, on the other hand, can be either multiplicative or additive (subtractive background) correction.

One ‘obvious’ idea would be to extrapolate the visible distributions of veto detectors down to a zero multiplicity or zero energy deposit by assuming a certain shape for the distributions (such as negative binomial distribution or similar), to estimate the veto inefficiency from data. However, this approach relies on several assumptions.

C.4 Poisson pileup problem

Let us assume for now that the Poissonian fluctuations are fluctuations in the number of simultaneous pp -interactions. That is, driven purely by instantaneous luminosity of the accelerator. The parameter μ describes the mean of visible simultaneous pp -interactions. We highlight the word visible because $k = 0$ includes not just ‘empty’ or ‘non-interacting’ bunch crossings but also experimentally non-visible inelastic interactions, the so-called low mass diffraction processes. These are the non-perturbative Regge domain QCD processes. Also, the elastic interactions belong to this category with cross section $\sim 25\%$ of the total pp -cross section, which is $\sigma_{tot}^{pp} \sim \mathcal{O}(100)$ mb at the LHC. The Poisson distribution comes from the law of *rare events* and is an approximation of the binomial distribution $\text{Bin}(n, p)$, when the number of trials $n \rightarrow \infty$ and $np = \mu$, then $\text{Bin}(n, p) \rightarrow \text{Poi}(\mu)$. For proton bunches circulating the LHC, n approximately $\sim 10^{11}$ – the number of protons per bunch, and p is order of $\sim 1/n$.

This gives us a finite probability of the triggered bunch cross event to originate from different Poisson terms as

$$P_P(k \geq 1; \mu) = 1 - P_P(0; \mu) = P_P(1; \mu) + P_P(2; \mu) + P_P(3; \mu) + \dots \quad (\text{C.11})$$

That is the sum of probabilities to have one visible pp -interaction, two pp -interactions (first order pileup) and so on, with the ‘zero-suppressed’ or truncated mean value: $\mathbb{E}[K|K > 0] = \mu/(1 - e^{-\mu}) \geq 1$ and the variance $\text{Var}[K|K > 0] = (\mu + \mu^2)/(1 - e^{-\mu}) - \mu^2/(1 - e^{-\mu})^2$, which both converge to μ at high values of μ . Experimentally, the Poisson distribution μ -value can be obtained from the mean trigger normalized rate $R \in [0, 1)$ by

$$\boxed{R \equiv 1 - P_P(0; \mu) = 1 - e^{-\mu} \Leftrightarrow \mu = -\ln(1 - R)}. \quad (\text{C.12})$$

Clearly when $R \ll 1$, then $\mu \simeq R$. The mean normalized trigger rate R itself is obtained from the measured trigger frequency f_R (Hz) as a function of time t . This we get by using the accelerator orbital frequency $f_O = 11.2455$ kHz at the LHC and the number of colliding bunch pairs N_B circulating

$$\langle R \rangle_{t, N_B} = \frac{\frac{1}{\Delta t} \int_{\Delta t} f_R(t) dt}{N_B f_O} \quad \text{or} \quad \langle R \rangle_{N_B} = \frac{N_E}{N_{BC}}. \quad (\text{C.13})$$

The interval Δt corresponds to the given data run span over time. The alternative version on right is just a reformulation using the number of triggered events N_E and the number of bunch crosses N_{BC} . For the formalism it is important to emphasize that the trigger corresponds to Boolean OR (\vee) operation between different

(sub)detectors, which is the minimum bias trigger. The reason is that this operator is the most inclusive Boolean operator.

In general, care is required in determining the μ -value experimentally because it is effectively a factor of both on-time and off-time pileup due to the detector time integration windows which can span over several bunch crosses. It must be emphasized that the formalism here requires only an *effective* μ -value of the Poisson distribution. For the rest of the work, we assume that the effective modeling and measurement of the μ -value is under control. The μ value determination at large μ is problematic directly from the $\mu = -\ln(1 - R)$ due to $R \simeq 1$. However, under high luminosity, the μ value is measured with less saturated signals, by using for example track or primary vertex counting.

We assume that the pileup is linear concerning the detector responses. Thus, we can take a linear incoherent superposition of the final states propagating from different simultaneous pp -interactions. Also to point out explicitly: when we ‘sum’ binary vectors, we use the component-wise OR, for example: $[1, 0] \vee [1, 1] = [1, 1] \in \mathbb{F}_2^2$.

C.5 Luminosity and total inelastic cross section

The μ -value can be understood also naturally in the context of instantaneous luminosity L ($\text{cm}^{-2} \cdot \text{s}^{-1}$) measurement by using a simplified definition as in [155]

$$\mu_{abs} \equiv \frac{\sigma_{inel} L}{N_B f_O}, \quad (\text{C.14})$$

where σ_{inel} is the total inelastic cross section (~ 80 mb in pp at $\sqrt{s} = 13$ TeV), dictated by non-perturbative strong interactions, N_B is the number of colliding bunch pairs and f_O is the LHC orbital frequency (Hz). The μ_{abs} is by definition the absolute mean number of inelastic interactions per colliding bunch cross.

However, because the complete σ_{inel} is essentially unknown due to limited low-mass diffraction acceptance at the LHC, also the μ_{abs} is unknown. The absolute mean number of inelastic interactions per bunch crossing μ_{abs} , is reduced to the measured μ by the total integrated detector multiplicative acceptance \times efficiency factor $0 \leq \epsilon \leq 1$ defining¹ the visible inelastic cross section as

$$\sigma_{vis} \equiv \epsilon \sigma_{inel}. \quad (\text{C.15})$$

Then using linearity of Eq. C.14 with respect to the cross section, we can write

$$\frac{\mu}{\mu_{abs}} = \frac{\sigma_{vis}}{\sigma_{inel}} \quad (\text{C.16})$$

and reformulating

$$L = \mu_{abs} \frac{N_B f_O}{\sigma_{inel}} = \mu \frac{N_B f_O}{\sigma_{vis}}. \quad (\text{C.17})$$

Now L or equivalently σ_{vis} , can be measured by doing a van der Meer (vdM) scan, which is an absolute luminosity measurement calibration technique. However, the ‘total inelastic measurement’ is based on an extrapolation inverting Eq. C.15. The efficiency \times acceptance factor ϵ here has an uncertainty at least order of $\delta\epsilon/\epsilon \sim 0.1$. Also, a crucial thing is to factorize the acceptance and efficiency, but experimentally one cannot always distinguish between these two and this creates the problem of defining the fiducial acceptance domain. This is because not all particles and their momentum of multibody decays of systems are measured, thus, the efficiency and acceptance have intrinsic detector simulation and minimum bias soft QCD Monte Carlo model dependence.

¹Technically, we may have also ‘leakage’ of events from outside the geometric acceptance, which do not exactly obey this definition, but nevertheless the effective ϵ is never larger than 1.

APPENDIX C. COMBINATORICS AND PILEUP

It is a significant challenge for experiments to carefully define the final state observables in the forward domain. Fiducial definitions may be in terms of single particle acceptance, described by (η, p_t) or $|\vec{p}|$, or in terms of the (diffractive) system invariant mass. Intuitively, the fiducial definition in terms of single particle kinematics should be less model-dependent than definitions relying on the invariant mass, at least if momentum or energy measurements are available. A special case to the discussion here is an indirect inference based on the elastic scattering, *extrapolation* of $d\sigma_{el}/dt$ down to $t \rightarrow 0$ and using the optical theorem (unitarity) relating the total cross section and the imaginary part of the elastic forward amplitude.

C.6 F^* -projection technique

We shall here shortly outline a data-MC hybrid interpolation and projection technique, what we may call the F^* -projection. The idea is simple: observables which cannot be directly reconstructed with the detector, may be multidimensionally interpolated using the Monte Carlo event generator based simulations, given the measured partial cross sections. This can be done at the generator level once the measured partial cross sections have been unfolded. This is closely related to non-orthogonal and overcomplete basis projection techniques known in the Wavelet and Compressed Sensing literature, for those see [235]. In Monte Carlo, for each $j = 1, \dots, n = 2^N - 1$ combinatorially selected sub-sample, we construct the probability distribution $f^{\text{MC}}(\mathcal{O})|_j$ of the observable \mathcal{O} . We may use normalized histograms to represent the distributions, for example. This gives us a set

$$\{f^{\text{MC}}(\mathcal{O})|_j\}_{j=1}^n. \quad (\text{C.18})$$

Then we measure the set of combinatorial partial cross sections in data

$$\{\sigma_j^{\text{DATA}}\}_{j=1}^n. \quad (\text{C.19})$$

The F^* -projected estimate of the differential cross section is obtained with

$$\frac{d\sigma^{F^*}}{d\mathcal{O}} = \sum_j \sigma_j^{\text{DATA}} f^{\text{MC}}(\mathcal{O})|_j, \quad (\text{C.20})$$

which shows clearly that the result is a sum over Monte Carlo driven ‘dictionary functions’ weighted with data driven coefficients. It is possible to show that for certain observables, when $N \rightarrow \infty$, the dependence on Monte Carlo is minimized. In practice, one wants to make the projection using several different Monte Carlo samples to obtain estimate of the model dependence. Natural distributions to reconstruct using this technique are rapidity gap size distributions $d\sigma/d\Delta y$ with varying boundary conditions, for example.

C.7 Diffraction fit algorithms

Let us have \mathcal{C} different scattering processes for which we have a Monte Carlo event generator based predictions of the multidimensional combinatorial partial cross sections, normalized to probability densities. We write these down as rows in a model likelihood matrix \mathcal{L} with size $\mathcal{C} \times 2^N - 1$. Now given the measured partial cross sections, we want to obtain the Maximum Marginal Likelihood solution: the best fit of the measurement as a weighted incoherent sum of the Monte Carlo model based ‘multidimensional template distributions’. The solution to this is obtained via Expectation Maximization iteration given in Algorithm 4. If in addition, we want to re-weight the event generator distributions to simultaneously fit e.g. the effective Pomeron intercept, we proceed with Algorithm 5. Statistical fit uncertainties can be obtained via bootstrap re-sampling. These algorithms can be used either at unfolded fiducial cross section or at visible (detector) level.

Algorithm 4 Maximum Marginal Likelihood estimator via EM-iteration.

INPUT: (Unfolded) measurement vector σ ($2^N - 1 \times 1$), model likelihood matrix \mathcal{L} ($\mathcal{C} \times 2^N - 1$) with $\forall i : \sum_j [\mathcal{L}]_{i,j} = 1$, initial estimate Θ ($\mathcal{C} \times 1$), otherwise $\Theta = \mathbf{1}_{\mathcal{C}}/|\mathcal{C}|$.

repeat

A. Inverse step using Bayes theorem:

Diagonal priors matrix \times Likelihood matrix:

$\mathcal{P} \leftarrow \text{diag}(\Theta_k) \mathcal{L}$

Normalize each column to obtain the probabilistic mixed density operator:

for all $j = 1, \dots, 2^N - 1$ **do**

$[\mathcal{P}]_{:,j} \leftarrow [\mathcal{P}]_{:,j} / \sum_{i=1}^{\mathcal{C}} [\mathcal{P}]_{i,j}$

end for

B. Forward step:

Operate on data with the mixed density operator:

$\Theta_{k+1} \leftarrow \mathcal{P} \sigma$

until Convergence $\|\Theta_{k+1} - \Theta_k\|_2 < \delta$

OUTPUT: The process cross section estimates $\hat{\Theta} = (\sigma^1, \sigma^2, \dots, \sigma^{\mathcal{C}})^T$.

Algorithm 5 Event-by-event MC re-weighting via M parameter brute force grid scan.

INPUT:

for all M parameters in $\Phi = [\phi_a^1, \phi_b^1] \times [\phi_a^2, \phi_b^2] \times \dots \times [\phi_a^M, \phi_b^M]$ **do**
for all MC events **do**

 Construct the MC observables functionally dependent on the parameters

 Re-weight the event according to $(\phi^1, \phi^2, \dots, \phi^M)$

end for

A. Construct a re-weighted MC process Likelihood matrix with weighted event selection

B. Estimate the process cross sections (mixing weights) using Algorithm A4

$\hat{\Theta} \leftarrow \text{A4}(\tilde{\sigma}, \mathcal{L}_{RW})$

C. Construct the new Monte Carlo estimate of $2^N - 1$ partial cross sections

$\tilde{\sigma}^{MC} \leftarrow \mathcal{L}_{RW}^T \hat{\Theta}$

D. Calculate Kullback-Leibler divergence $D(\text{data}|\text{model})$ with

$D \leftarrow D(\tilde{\sigma}|\sigma^{MC})$

end for

OUTPUT: Kullback-Leibler divergences for the parameter (hyper)grid $\Phi \subseteq \mathbb{R}^M$. Use these to infer optimal values and parameter sensitivity/uncertainty.

D GRANIITTI Code (2×2)

APPENDIX D. GRANIETTI CODE (2×2)

D.1 Makefile

Makefile

1/6

```
1: # GRAMITTT Makefile
2: #
3: # (c) 2017-2020 Mikael Hieskolainen.
4: # Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: #
6: #
7: # COMPILING:
8: # make -j4
9: #
10: # CLEANING:
11: # make clean (objects)
12: # make superclean (objects + binaries)
13: #
14: # To compile with clang: make CXX=clang
15: # -|- unit tests: make TEST=TRUE
16: #
17: # To compile with ROOT using -std=c++11: make ROOT1=TRUE
18: #
19: # To compile with profiling: make VALGRIND=TRUE
20: #
21: #-----
22: #
23: # EXTERNAL LIBRARIES SETUP:
24: #
25: # [HEPMC3] and [LHAFDP]: See ./install folder
26: # [ROOT] Example: export ROOTSYS=/home/user/sw/ROOT6
27: #
28: # GENERAL:
29: #
30: # If you just installed e.g. HEPMC3 libraries
31: # $ sudo ldconfig
32: #
33: #-----
34: #
35: # If you see symbol errors related to HepMC such that:
36: # "undefined symbol: _ZN5HepMC10FilterBase9init_has_endVertexEv"
37: # check that you do not have a collision of HEPMC3 with an existing
38: # HepMC2 shared library installation!
39: #
40: # Check below:
41: #
42: # For tracking the libraries of the program:
43: # $ ld -l /executable
44: #
45: # Check if the shared object (.so) file contains the missing symbol:
46: # $ nm -D libHepMC.so
47: #
48: #
49: # Creating shared libraries, use -fPIC flag, then:
50: # g++ -shared MH1.0 -o libMH1.so
51: #
52: #
53: # USE [TABS] for indentation while modifying this file!
54: #-----
55: #
56: #
57: # Detect compiler version if using by default g++
58: #
59: #
60: ifeq ($(CXX),g++)
61:
62: # KERNEL_GCC_VERSION
63: # := $(shell cat /proc/version | sed -n 's/^.*gcc version \([0-9]*\.[0-9]*\.[0-9]*\) .*$/\1/p' | head -n 1)
64: # := $(info Info KERNEL_GCC_VERSION = $(KERNEL_GCC_VERSION))
65:
66: # CXX_VERSION = $(shell g++ -dumpversion)
67: # $Info Found CXX_VERSION = $(CXX_VERSION)
68: # CXX_REQUIRED = 7.0
69:
70: # Use bc command for double comparison, a common utility on Unix platforms
71: # FOO=$(shell echo "$((CXX_VERSION) < $(CXX_REQUIRED))" | bc)
72:
73: ifeq ($(FOO),1)
74: #error Your CXX_VERSION is too old, need at least CXX_VERSION = $(CXX_REQUIRED), update your GCC chain
75: endif
76:
77: #
78: #
79: #
80: #
81: # PATH setup
82: #
83: # HEPMC3 installation
```

Makefile

3/6

```
167: INCLUDES += -I$(HEPMC3SYS)/include
168: INCLUDES += -I$(LHAFDPFYS)/include
169:
170: # C++
171: INCLUDES += -I/user/include
172: INCLUDES += -I/user/local/include
173:
174: # Obj
175: INCLUDES += -I
176: INCLUDES += -Iinclude
177:
178: # Internal libraries
179: INCLUDES += -Ilibs
180:
181: # Ftnesor
182: INCLUDES += -Ilibs/Ftnesor
183:
184: # Eigen
185: INCLUDES += -Ilibs/Eigen/unsupported/
186:
187: # Pytorch
188: #INCLUDES += -Ilibs/libtorch/include/
189: #INCLUDES += -Ilibs/libtorch/include/torch/csrc/api/include
190:
191: #
192: #-----
193: # Compiler and its options
194: #
195: CXX = g++
196:
197: CXXVER = -std=c++17
198: CXXVER_OLD = -std=c++17
199:
200: # Use this for alternative ROOT installations
201: ifeq ($(ROOT1),TRUE)
202: CXXVER_OLD = -std=c++14
203: endif
204:
205: OPTIM = -O3 -DNDEBUG -fno-vectorize -fno-signed-zeros
206: CXXFLAGS = -Wall -fPIC -pipe $(OPTIM)
207:
208: # Profiling & debug
209: ifeq ($(VALGRIND),TRUE)
210: OPTIM += -pg
211: endif
212:
213:
214: # Needed by PyTorch if using pre-compiled (ABI = Application Binary Interface)
215: # gcc < 5.1 is 0, later versions use 1 by default
216: # CXXFLAGS += -D_GLIBCXX_USE_CXX11_ABI=0
217:
218: # Automatic dependencies on
219: CXXFLAGS += -MD -MP
220:
221: # -Wall, compiler warnings full on
222: # -fno-vectorize, Autovectorization on
223: # -fno-signed-zeros Optimization (floating point) which ignores the signedness of zero
224: # -fPIC, Position independent code (PIC) for shared libraries
225: # -O2, -O3 Optimization level
226: # -DNDEBUG Release flag, no debug
227: # -pipe, Faster compilation using pipes
228: # -O0 -pg, Profiling and debugging (HIGH PERFORMANCE HIT)
229: # -rpath-link Needed for /so which links to another /so
230: #
231: #
232: # Check your CPU instruction set with: cat /proc/cpuinfo
233: #
234: #
235: # DANGEROUS:
236: #
237: # -march=native, CPU specific instruction set usage
238: # (gives unknown flops problem, factor 1/4 wrong results on i5-4570 with g++7.4, do not use!)
239: # -ffast-math, Heavy floating point optimization
240: # (fast but breaks IEEE flops standards, do not use!)
241: #
242: #-----
243: # Sources, objects and dependency files
244: #
245: OBJ_DIR = obj
246:
247: #
248: #
249: #
```

Makefile

2/6

```
84: #HEPMC3SYS = $(shell HepMC-config --prefix)
85:
86: # LHAFDP Installation
87: #LHAFDPFYS = $(shell lhafdp-config --prefix)
88:
89: # ROOT installation
90: #ROOTSYS = $(shell root-config --prefix)
91:
92:
93: ifeq ($(HEPMC3SYS),)
94: #error Please set HEPMC3SYS environment variable and paths [perhaps "source ./install/setenv.sh"]
95: else
96: $Info Found HEPMC3SYS = $(HEPMC3SYS)
97: endif
98:
99: ifeq ($(LHAFDPFYS),)
100: #error Please set LHAFDPFYS environment variable and paths [perhaps "source ./install/setenv.sh"]
101: else
102: $Info Found LHAFDPFYS = $(LHAFDPFYS)
103: endif
104:
105: # If ROOT seems to be installed, tag ROOT dependent compilation on
106: ifeq ($(ROOTSYS),)
107: $Info Did not find ROOTSYS - compilation without ROOT <root.cern.ch> dependent tools)
108: ROOT=FALSE
109: else
110: $Info Found ROOTSYS = $(ROOTSYS)
111: ROOT=TRUE
112:
113: # Info messages
114: $(Info )
115: $(Info *)
116: $(Info **)
117: $(Info *** If compilation with ROOT libraries fails, try with: make -j4 ROOT1=TRUE **)
118: $(Info ****)
119: $(Info *****)
120: $(Info)
121: endif
122:
123:
124: #-----
125: # External libraries to be linked
126:
127: # HEPMC3 (lib64 needed on some systems)
128: HEPMC3lib = -L$(HEPMC3SYS)/lib -l$(HEPMC3SYS)/lib64 -lHepMC3 -lHepMC3search
129:
130: # LHAFDPFYS (lib64 needed on some systems)
131: LHAFDPFlib = -L$(LHAFDPFYS)/lib -l$(LHAFDPFYS)/lib64 -lLHAFDPF
132:
133: # PyTorch
134: # Note -Wl,-rpath-link= handles the recursive dependency (for linker)
135: #PYTORCHlib = -L./libs/libtorch/lib -ltorch -ltorch -lqomp -lcaffe2 \
136: # -Wl,-rpath-link=./libs/libtorch/lib
137:
138: # ROOT
139: ifeq ($(ROOT),TRUE)
140: ROOTlib = -L$(ROOTSYS)/lib -L$(ROOTSYS)/lib/64bit -lCore -lRIO -lNet \
141: -lHist -lICarf -lICorah3d -lIquad -lTree -lRint \
142: -lPostscript -lMatrix -lPhysics -lMathCore \
143: -lThread -lGui -lRoofit -lMinuit
144: endif
145:
146: # C++ standard
147: STANDARDlib = -std=c++11 -pthread -lrt
148: # -ldl -rdynamic
149:
150: # GSL, on ubuntu run: sudo apt-get install libgsl-dev
151: #GSLlib = -lgsl -lgslcblas
152:
153:
154: # External libraries (THESE TWO FIRST for priority!!)
155: LDLIBS = $(HEPMC3lib)
156: LDLIBS += $(LHAFDPFlib)
157:
158: # The rest
159: LDLIBS += $(STANDARDlib)
160: LDLIBS += $(PYTORCHlib)
161:
162: #
163: #
164: # Header files
165:
166: # External libraries (THESE TWO FIRST for priority!!)
167:
168:
169:
170:
171:
172:
173:
174:
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179:
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235:
236:
237:
238:
239:
240:
241:
242:
243:
244:
245:
246:
247:
248:
249:
250: # LIBRARY
251:
252: #-----
253: SRC_DIR_0 = src
254: SRC_0 = $(wildcard $(SRC_DIR_0)/*.cc)
255: OBJ_0 = $(SRC_0:$(SRC_DIR_0)/%.cc=$(OBJ_DIR)/%.o)
256: DEP_0 = $(OBJ_0:$(OBJ_DIR)/%.o=.d)
257: #
258:
259: #-----
260: SRC_DIR_1 = src/Amplitude
261: SRC_1 = $(wildcard $(SRC_DIR_1)/*.cc)
262: OBJ_1 = $(SRC_1:$(SRC_DIR_1)/%.cc=$(OBJ_DIR)/%.o)
263: DEP_1 = $(OBJ_1:$(OBJ_DIR)/%.o=.d)
264: #
265:
266: # Create collection of all library objects
267: OBJ = $(OBJ_0) $(OBJ_1)
268:
269: #-----
270: ifeq ($(ROOT),TRUE)
271: SRC_DIR_ROOT = src/Analysis
272: SRC_2 = $(wildcard $(SRC_DIR_ROOT)/*.cc)
273: OBJ_2 = $(SRC_2:$(SRC_DIR_ROOT)/%.cc=$(OBJ_DIR)/%.o)
274: DEP_2 = $(OBJ_2:$(OBJ_DIR)/%.o=.d)
275: endif
276:
277: #-----
278: #
279: ifeq ($(TEST),TRUE)
280: SRC_DIR_3 = tests/catchlibrary
281: SRC_3 = $(wildcard $(SRC_DIR_3)/*.cc)
282: OBJ_3 = $(SRC_3:$(SRC_DIR_3)/%.cc=$(OBJ_DIR)/%.o)
283: DEP_3 = $(OBJ_3:$(OBJ_DIR)/%.o=.d)
284: endif
285: #
286:
287: #-----
288: # PROGRAM compiled against the library
289:
290: #-----
291: SRC_DIR_PROGRAM = src/Program
292: SRC_PROGRAM = $(wildcard $(SRC_DIR_PROGRAM)/*.cc)
293: OBJ_PROGRAM = $(SRC_PROGRAM:$(SRC_DIR_PROGRAM)/%.cc=$(OBJ_DIR)/$(BIN_DIR)/%.o)
294: DEP_PROGRAM = $(OBJ_PROGRAM:$(OBJ_DIR)/$(BIN_DIR)/%.o=.d)
295: #
296: #-----
297: SRC_DIR_TEST = tests
298: SRC_TEST = $(wildcard $(SRC_DIR_TEST)/*.cc)
299: OBJ_TEST = $(SRC_TEST:$(SRC_DIR_TEST)/%.cc=$(OBJ_DIR)/$(BIN_DIR)/%.o)
300: DEP_TEST = $(OBJ_TEST:$(OBJ_DIR)/$(BIN_DIR)/%.o=.d)
301: #
302: #-----
303: ifeq ($(ROOT),TRUE)
304: SRC_DIR_PROGRAM_ROOT = src/ProgramAnalysis
305: SRC_PROGRAM_ROOT = $(wildcard $(SRC_DIR_PROGRAM_ROOT)/*.cc)
306: OBJ_PROGRAM_ROOT = $(SRC_PROGRAM_ROOT:$(SRC_DIR_PROGRAM_ROOT)/%.cc=$(OBJ_DIR)/$(BIN_DIR)/%.o)
307: DEP_PROGRAM_ROOT = $(OBJ_PROGRAM_ROOT:$(OBJ_DIR)/$(BIN_DIR)/%.o=.d)
308: #
309: #-----
310: SRC_DIR_PROGRAM_TEST = tests
311: SRC_PROGRAM_TEST = $(wildcard $(SRC_DIR_PROGRAM_TEST)/*.cc)
312: OBJ_PROGRAM_TEST = $(SRC_PROGRAM_TEST:$(SRC_DIR_PROGRAM_TEST)/%.cc=$(OBJ_DIR)/$(BIN_DIR)/%.o)
313: DEP_PROGRAM_TEST = $(OBJ_PROGRAM_TEST:$(OBJ_DIR)/$(BIN_DIR)/%.o=.d)
314: #
315: #-----
316: #
317: #
318: #
319: # PROGRAM
320: #
321: #
322: # Directory
323: BIN_DIR = bin
324:
325: # SUFFIXES: .o .cc
326:
327: # Normal
328: EXE_NAMES = qr scan minbias hepctlab data2hepctl pathmark pedbench sommerfeld ot
329: PROGRAM = $(EXE_NAMES:=$(BIN_DIR)/%)
330:
331: ifeq ($(ROOT),TRUE)
332: EXE_ROOT_NAMES = analyze fitsort fitocentral fitharmonic
```

Makefile

5/6

```

333: PROGRAM_ROOT = $(EXE_ROOT_NAMES:%=$(BIN_DIR)/%)
334: endif
335:
336: PROGRAM_TEST =
337: ifeq ($(TEST),TRUE)
338: EXE_TEST_NAME = testbench0 testbench1 testbench2 testbench3
339: PROGRAM_TEST = $(EXE_TEST_NAMES:%=$(BIN_DIR)/%)
340: endif
341:
342: # Multicore
343: export MAKEFLAGS="-$@" $(grep -c "processor /proc/cpuinfo")
344:
345: # -----
346: #
347: # RULES for linking
348:
349: all: $(PROGRAM) $(PROGRAM_ROOT) $(PROGRAM_TEST)
350:     @echo " "
351:     @echo "PROGRAM:" $(PROGRAM) $(PROGRAM_ROOT) $(PROGRAM_TEST)
352:     @echo " "
353:     @echo "Compilation of '$@' done."
354:
355: $(PROGRAM) : $(OBJ) $(OBJ_PROGRAM)
356:     $(CXX) $(OBJ_DIR)/%.o $(OBJ) -o $@ $(CXXFLAGS) $(LDLIBS)
357:
358: $(PROGRAM_ROOT) : $(OBJ) $(OBJ_2) $(OBJ_PROGRAM_ROOT)
359:     $(CXX) $(OBJ_DIR)/%.o $(OBJ) $(OBJ_2) -o $@ $(CXXFLAGS) $(LDLIBS) $(ROOTLIB)
360:
361: # Unit tests (note, we use catchmain.o from $(OBJ_2) for linking with catch2)
362: $(PROGRAM_TEST) : $(OBJ) $(OBJ_2) $(OBJ_PROGRAM_TEST)
363:     $(CXX) $(OBJ_DIR)/%.o $(OBJ) $(OBJ_2) -o $@ $(CXXFLAGS) $(LDLIBS)
364:
365: # -----
366: #
367: # RULES to generate library dependencies and compile
368:
369: # LIBRARY objects
370: #
371: # Note that for different ROOT installations, we need both:
372: # -I$(ROOTSYS)/include
373: # -I$(ROOTSYS)/include/root
374:
375: # -----
376: $(OBJ_DIR)/%.o: $(SRC_DIR)/%.cc
377:     @echo " "
378:     @echo "Generating dependencies and compiling $<..."
379:     $(CXX) $(CXXOVER) -c $< -o $@ $(CXXFLAGS) $(INCLUDES)
380: # -----
381: #
382: #
383: $(OBJ_DIR)/%.o: $(SRC_DIR_1)/%.cc
384:     @echo " "
385:     @echo "Generating dependencies and compiling $<..."
386:     $(CXX) $(CXXOVER) -c $< -o $@ $(CXXFLAGS) $(INCLUDES)
387: # -----
388: #
389: $(OBJ_DIR)/%.o: $(SRC_DIR_2)/%.cc
390:     @echo " "
391:     @echo "Generating dependencies and compiling $<..."
392:     $(CXX) $(CXXOVER_OLD) -c $< -o $@ $(CXXFLAGS) $(INCLUDES) \
393:         -I$(ROOTSYS)/include -I$(ROOTSYS)/include/root
394: # -----
395: #
396: #
397: $(OBJ_DIR)/%.o: $(SRC_DIR_3)/%.cc
398:     @echo " "
399:     @echo "Generating dependencies and compiling $<..."
400:     $(CXX) $(CXXOVER) -c $< -o $@ $(CXXFLAGS) $(INCLUDES)
401: # -----
402: #
403: #
404: # PROGRAM objects
405:
406: # -----
407: $(OBJ_DIR)/$(BIN_DIR)/%.o: $(SRC_DIR_PROGRAM)/%.cc
408:     @echo " "
409:     @echo "Generating dependencies and compiling $<..."
410:     $(CXX) $(CXXOVER) -c $< -o $@ $(CXXFLAGS) $(INCLUDES)
411: # -----
412: #
413: #
414: $(OBJ_DIR)/$(BIN_DIR)/%.o: $(SRC_DIR_PROGRAM_ROOT)/%.cc
415:     @echo " "

```

Makefile

6/6

```

416:     @echo "Generating dependencies and compiling $<..."
417:     $(CXX) $(CXXOVER_OLD) -c $< -o $@ $(CXXFLAGS) $(INCLUDES) \
418:         -I$(ROOTSYS)/include -I$(ROOTSYS)/include/root
419: # -----
420: #
421: #
422: $(OBJ_DIR)/$(BIN_DIR)/%.o: $(SRC_DIR_PROGRAM_TEST)/%.cc
423:     @echo " "
424:     @echo "Generating dependencies and compiling $<..."
425:     $(CXX) $(CXXOVER) -c $< -o $@ $(CXXFLAGS) $(INCLUDES)
426: # -----
427: #
428: #
429: # Clean up of object files
430: # - ignores return code error
431: # @ is silent
432: .PHONY: clean
433: clean:
434:     @echo "Cleaning objects"
435:     -rm $(OBJ_DIR)/*.o 2>/dev/null || true
436:     -rm $(OBJ_DIR)/*.d 2>/dev/null || true
437:     -rm $(OBJ_DIR)/$(BIN_DIR)/*.o 2>/dev/null || true
438:     -rm $(OBJ_DIR)/$(BIN_DIR)/*.d 2>/dev/null || true
439:
440: .PHONY: superclean
441: superclean: clean
442:     @echo "Cleaning binaries"
443:     -rm $(BIN_DIR)/* 2>/dev/null || true
444:     -rm $(BIN_DIR)/* 2>/dev/null || true
445:
446: # -----
447: # Dependencies listed here
448: -include $(DEP_0)
449: -include $(DEP_1)
450: ifeq ($(ROOT),TRUE)
451: -include $(DEP_2)
452: endif
453: ifeq ($(TEST),TRUE)
454: -include $(DEP_3)
455: endif
456:
457: # DEBUG printing
458:
459: #print_vars:
460:     @ echo P00 = $(OBJ)
461:

```

APPENDIX D. GRANIETTI CODE (2×2)

D.2 MADGRAPH amplitude to GRANIETTI converter

D.3 C++ header files


```

./include/Granitti/MQuasiElastic.h 1/1
1: // QuasiElastic (EL,SD,DD) and soft MD class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MQUASIELASTIC_H
7: #define MQUASIELASTIC_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // HepMC3
15: #include "HepMC3/GenEvent.h"
16:
17: // Own
18: #include "Granitti/M4Vec.h"
19: #include "Granitti/MBox.h"
20: #include "Granitti/MMatrix.h"
21: #include "Granitti/MProcess.h"
22: #include "Granitti/MSpin.h"
23:
24: namespace gra {
25: // Matrix element dimension: "GeV" << -(2*external_legs - 8)
26: class MQuasiElastic : public MProcess {
27: public:
28: MQuasiElastic(const string process, const std::vector<aux::OneCMD> &syntax);
29: virtual ~MQuasiElastic();
30: void post_Constructor();
31:
32: //
33: // double operator()(const std::vector<double> &rvecvec, AuxInData &aux) {
34: //     return EventWeight(rvecvec, aux);
35: // }
36: //
37: // double EventWeight(const std::vector<double> &rvecvec, AuxInData &aux);
38: // bool EventRecord(HepMC3::GenEvent &evt);
39: // void PrintInit(bool silent) const;
40:
41: private:
42: void Initialize();
43: bool LoopKinematics(const std::vector<double> &ip, const std::vector<double> &ip2);
44: bool FiducialCuts(const std::vector<double> &cuts);
45:
46: // Non-Diffractive
47: std::complex<double> PolySoft(const std::vector<double> &rvecvec);
48:
49: // 2/3-Dim phase space, 2>2 QuasiElastic
50: bool B3BuildKin(const std::vector<double> &rvecvec);
51: bool B3BuildKin(double s3, double s4, double t);
52: bool B3GetLorentzScalars();
53: double B3IntegrateVolume() const;
54: double B3PhaseSpaceWeight() const;
55:
56: // Event by event integration boundaries
57: double t_max = 0.0;
58: double t_min = 0.0;
59: double DD_M2_max = 0.0;
60: double log_DD_M2_max = 0.0;
61:
62: // Multipomeron weight table
63: std::vector<double> MMAPPMW;
64:
65: //
66: } // namespace gra
67:
68: #endif

```

```

./include/Granitti/MVEGAS.h 2/3
84: }
85:
86: // VEGAS rebinning function (algorithm adapted from Numerical Recipes)
87: void Rebin(double ac, unsigned int j, const VEGASPARAM &param) {
88:     unsigned int k = 0;
89:     double dz = 0.0;
90:     double zn = 0.0;
91:     double zo = 0.0;
92:
93:     for (std::size_t i = 0; i < param.BINS - 1; ++i) {
94:         while (ac > dz) { dz += rvec[i+1] - 1; }
95:         if (k > 1) { zo = xmat[k - 2][j]; }
96:         zn = xmat[k - 1][j];
97:         dz = zn - zo;
98:         kxache[i] = zn - (zn - zo) * dz / rvec[k - 1];
99:     }
100:
101:     for (std::size_t i = 0; i < param.BINS - 1; ++i) { xmat[i][j] = kxache[i]; }
102:
103:     xmat[param.BINS - 1][j] = 1.0;
104: }
105:
106: // VEGAS grid optimizing function (algorithm adapted from Numerical Recipes)
107: void OptimizeGrid(const VEGASPARAM &param) {
108:     double zo = 0;
109:     double zn = 0;
110:     double ac = 0;
111:
112:     for (std::size_t j = 0; j < FDIM; ++j) {
113:         zo = f2mat[0][j];
114:         zn = f2mat[1][j];
115:         f2mat[0][j] = (zo + zn) / 2.0;
116:         dcache[j] = f2mat[0][j];
117:
118:         for (std::size_t i = 2; i < param.BINS; ++i) {
119:             ac = zo + zn;
120:             zo = zn;
121:             zn = f2mat[i][j];
122:             f2mat[i - 1][j] = (ac + zn) / 3.0;
123:             dcache[i] += f2mat[i - 1][j];
124:         }
125:         f2mat[param.BINS - 1][j] = (zo + zn) / 2.0;
126:         dcache[j] += f2mat[param.BINS - 1][j];
127:     }
128:
129:     for (std::size_t j = 0; j < FDIM; ++j) {
130:         ac = 0.0;
131:         for (std::size_t i = 0; i < param.BINS; ++i) {
132:             f2mat[i][j] = (f2mat[i][j] < param.EPS ? param.EPS : f2mat[i][j]);
133:             rvec[i] = std::pow(1.0 - f2mat[i][j], dcache[j]);
134:             rvec[i] = (std::ilog(dcache[j]) - std::ilog(f2mat[i][j]) + param.EPS) /
135:                     param.LAMBDA);
136:
137:             // Floating point precision (integrand close to 0)
138:             if (std::isnan(rvec[i]) || std::isinf(rvec[i]) || rvec[i] == param.EPS) {
139:                 ac += rvec[i];
140:             }
141:             Rebin(ac / param.BINS, j, param);
142:         }
143:
144: // Initialize sampling region [0,1] x [0,1] x ... x [0,1]
145: void InitRegion(unsigned int fdim) {
146:     FDIM = fdim;
147:     region.resize(2 * FDIM, 0.0);
148:
149: // Lower bound [zero]
150: for (std::size_t i = 0; i < FDIM; ++i) { region[i] = 0.0; }
151: // Upper bound [one]
152: for (std::size_t i = FDIM; i < 2 * FDIM; ++i) { region[i] = 1.0; }
153: }
154:
155: // Clear integrated data
156: void ClearIntegral() {
157:     sumdata = 0.0;
158:     sumch2 = 0.0;
159:     sweight = 0.0;
160: }
161:
162: // Full Init
163: void ClearAll(const VEGASPARAM &param) {
164: // Vectors [MAXFDIM]
165: dcache = std::vector<double>(param.MAXFDIM, 0.0);
166: dxvec = std::vector<double>(param.MAXFDIM, 0.0);

```

```

./include/Granitti/MVEGAS.h 1/3
1: // VEGAS integrator class grid routines
2: //
3: // Future step:
4: // Factorize all VEGAS functions out from MGranitti to here (use function pointers etc.)
5: //
6: // (c) 2017-2020 Mikael Mieskolainen
7: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
8:
9: #ifndef MVEGAS_H
10: #define MVEGAS_H
11:
12: #include <vector>
13:
14: namespace gra {
15:
16: // VEGAS MC default parameters
17: struct VEGASPARAM {
18:     unsigned int BINS = 128; // Maximum number of bins per dimension (EVEN NUMBER!)
19:     double LAMBDA = 1.5; // Regularization parameter
20:
21: // Initialization
22:     unsigned int NCALL = 20000; // Number of calls per iteration
23:     unsigned int ITER = 15; // Number of iterations
24:     double CH2MAX = 10.0; // Maximum chi2 in initialization
25:     double PRECISION = 0.01; // Maximum relative error of cross section integral
26:     int DEBUG = -1; // Debug mode
27:
28: // User cannot set these
29:     unsigned int MAXFDIM = 100; // Maximum integral dimension
30:     double EPS = 1.0e-30; // Epsilon parameter
31:
32: //
33: //
34: // VEGAS MC adaptation data
35: struct VEGASData {
36: // VEGAS initialization function
37: void Init(unsigned int init, const VEGASPARAM &param) {
38: // First initialization: Create the grid and initial data
39: if (init == 0) {
40: ClearAll(param);
41: InitGridDependent(param);
42: // Use the previous grid but NOT integral data
43: else if (init == 1) {
44: ClearIntegral();
45: InitGridDependent(param);
46: // Use the previous grid and its integral data
47: else if (init == 2) {
48: InitGridDependent(param);
49: //
50: //
51: InitGridDependent(param);
52: } else {
53: throw std::invalid_argument("VEGASData::Init: Unknown init parameter = " +
54:                               std::to_string(init));
55: }
56: }
57:
58: // Create number of calls per thread, they need to sum to calls
59: std::vector<unsigned int> GetLocalCalls(int calls, int N_threads) {
60: std::vector<unsigned int> localcalls(N_threads, 0);
61: int sum = 0;
62: for (int k = 0; k < N_threads; ++k) {
63: localcalls[k] = std::floor(calls / N_threads);
64: sum += localcalls[k];
65: }
66: // Add remainder to the thread number 0
67: localcalls[0] += calls - sum;
68: return localcalls;
69: }
70:
71: // Initialize
72: void InitGridDependent(const VEGASPARAM &param) {
73: // Create grid spacing
74: for (std::size_t j = 0; j < FDIM; ++j) { dxvec[j] = region[j] + FDIM - region[j]; }
75:
76: // If binning parameter changed from previous call
77: if (param.BINS != BINS_prev) {
78: for (std::size_t i = 0; i < std::max(param.BINS, BINS_prev); ++i) { rvec[i] = 1.0; }
79: for (std::size_t j = 0; j < FDIM; ++j) {
80: Rebin(BINS_prev / static_cast<double>(param.BINS), j, param);
81: }
82: BINS_prev = param.BINS;
83: }

```

```

./include/Granitti/MVEGAS.h 3/3
167: // Vectors [BINS]
168: dxvec = std::vector<double>(param.BINS, 0.0);
169: kxache = std::vector<double>(param.BINS, 0.0);
170:
171: // Matrices [BINS x MAXFDIM]
172: fmat = std::vector<std::vector<double>>(param.BINS, std::vector<double>(param.MAXFDIM, 0.0));
173: f2mat = std::vector<std::vector<double>>(param.BINS, std::vector<double>(param.MAXFDIM, 0.0));
174: xmat = std::vector<std::vector<double>>(param.BINS, std::vector<double>(param.MAXFDIM, 0.0));
175:
176: // Init with 1!
177: for (std::size_t j = 0; j < FDIM; ++j) { xmat[0][j] = 1.0; }
178:
179: // VEGAS integrals data
180: sumdata = 0.0;
181: sumch2 = 0.0;
182: sweight = 0.0;
183:
184: // Previous binning
185: BINS_prev = 1;
186:
187: }
188:
189: // VEGAS scalars
190: unsigned int BINS_prev = 0;
191: unsigned int FDIM = 0;
192:
193: double fsum = 0.0;
194: double f2sum = 0.0;
195:
196: double sumdata = 0.0;
197: double sumch2 = 0.0;
198: double sweight = 0.0;
199:
200: // Vectors
201: std::vector<double> region;
202:
203: // Vectors
204: std::vector<double> dcache;
205: std::vector<double> dxvec;
206:
207: // Vectors
208: std::vector<std::vector<double>> fmat;
209: std::vector<std::vector<double>> f2mat;
210: std::vector<std::vector<double>> xmat;
211:
212: // Matrices
213: std::vector<std::vector<double>> fmat;
214: std::vector<std::vector<double>> f2mat;
215: std::vector<std::vector<double>> xmat;
216:
217: // struct VEGASData
218: } // namespace gra
219:
220: #endif

```

```

./include/Granitti/MFlux.h 1/1
1: // Photon and other fluxes
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MFLUX_H
7: #define MFLUX_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/MForm.h"
16: #include "Granitti/MGlobale.h"
17: #include "Granitti/MKinematics.h"
18:
19: namespace gra {
20: namespace flux {
21:
22:
23: double ApplyKEPFfluxes(double amp2, gra::LORENTZSCALAR& lts);
24: double ApplyKEFluxes(double amp2, gra::LORENTZSCALAR& lts);
25: double ApplyMFluxes(double amp2, gra::LORENTZSCALAR& lts);
26:
27:
28: } // namespace flux
29: } // namespace gra
30:
31:
32: #endif

```

```

./include/Granitti/MCW.h 1/2
1: // Monte Carlo Weight Objects (HEADER ONLY file)
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MCW_H
7: #define MCW_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <valarray>
13: #include <vector>
14:
15: #include "Granitti/MMath.h"
16:
17:
18: namespace gra {
19: namespace kinematics {
20:
21: // A simple Monte Carlo weight (container)
22: class MCW {
23: public:
24: MCW() {
25: W = 0.0;
26: W2 = 0.0;
27: N = 0.0;
28: }
29: MCW(double w, double w2, double n) {
30: W = w;
31: W2 = w2;
32: N = n;
33: }
34: // operator
35: MCW operator+(const MCW &obj) {
36: MCW res;
37: res.W = W + obj.W;
38: res.W2 = W2 + obj.W2;
39: res.N = N + obj.N;
40: return res;
41: }
42: // operator
43: MCW operator+=(const MCW &obj) {
44: this->W += obj.W;
45: this->W2 += obj.W2;
46: this->N += obj.N;
47: return *this; // return by reference
48: }
49: // * operator (does scale W and W^2 but not N)
50: MCW operator*(double scale) { return MCW(GetW() * scale, GetW2() * scale * scale, GetN()); }
51:
52: // Push new weight
53: void Push(double w) {
54: W += w;
55: W2 += w * w;
56: N += 1.0;
57: }
58: // MC estimate of the integral
59: double Integral() const { return (N > 0.0) ? W / N : 0.0; }
60: // MC estimate of the integral error squared (standard error of the mean)
61: double IntegralError2() const {
62: if (N > 0.0) {
63: return (W2 / N - gra::math::pow2(W / N)) / N;
64: } else {
65: return 0.0;
66: }
67: }
68: double IntegralError() const { return gra::math::msqrt(IntegralError2()); }
69: double GetW() const { return W; }
70: double GetW2() const { return W2; }
71: double GetN() const { return N; }
72: void SetW(double w) { W = w; }
73: void SetW2(double w2) { W2 = w2; }
74: void SetN(double n) { N = n; }
75:
76: private:
77: double W; // sum of weights
78: double W2; // sum of weights^2
79: double N; // trials
80: };
81:
82:
83: // A weighted sum of MCW container integral values (use with VEGAS, for example)

```

```

./include/Granitti/MCW.h 2/2
84: class MCWSum {
85: public:
86: MCWSum() {}
87:
88: void Add(const MCW &w, double weight) {
89: wsum += weight;
90: wintsum += weight * x.Integral();
91: error2sum += (weight * weight) * x.IntegralError2();
92: }
93: double Integral() const {
94: if (wsum > 0.0) {
95: return wintsum / wsum;
96: } else {
97: return 0.0;
98: }
99: }
100: double IntegralError2() const {
101: if (wsum > 0.0) {
102: return error2sum / gra::math::pow2(wsum);
103: } else {
104: return 0.0;
105: }
106: }
107: double IntegralError() const { return gra::math::msqrt(IntegralError2()); }
108:
109: private:
110: // sum_i weight_i
111: double wsum = 0.0;
112:
113: // sum_i weight_i * integral_i
114: double wintsum = 0.0;
115:
116: // sum_i weight_i^2 * error_i^2
117: double error2sum = 0.0;
118: };
119:
120: } // namespace kinematics
121: } // namespace gra
122:
123: #endif

```

```

./include/Granitti/MDurham.h 1/2
1: // "Durham QCD" Processes and Amplitudes
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MDURHAM_H
7: #define MDURHAM_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/MVVec.h"
16: #include "Granitti/MAmplitudes.h"
17: #include "Granitti/MAux.h"
18: #include "Granitti/MGlobale.h"
19: #include "Granitti/MKinematics.h"
20: #include "Granitti/MSudakov.h"
21:
22:
23: namespace gra {
24:
25: // Durham loop integral discretization technical parameters
26: struct MDurhamParam {
27: unsigned int N_qt = 0; // (> 30)
28: unsigned int N_phi = 0; // (> 10)
29:
30: double qt2_MIN = 0; // Loop momentum qt^2 minimum (GeV^2)
31: double qt2_MAX = 0; // Loop momentum qt^2 maximum (GeV^2)
32:
33: std::string PDF_scale = "MIN"; // Scheme
34: double alpha_scale = 4.0; // PDF Factorization scale
35: double MAXCOS = 0.9; // Neutron amplitude |cos(theta*)| < MAXCOS
36:
37: // THESE ARE CALCULATED FROM ABOVE
38: double qt_MIN = 0;
39: double qt_MAX = 0;
40: double qt_STEP = 0;
41: double phi_STEP = 0;
42: bool initialized = false;
43:
44: // Read parameters from file
45: void ReadParameters(const std::string &modelfile) {
46: using json = nlohmann::json;
47: const std::string data = gra::aux::GetInputData(modelfile);
48: json j;
49:
50: try {
51: j = json::parse(data);
52:
53: // JSON block identifier
54: const std::string XID = "PARAM_DURHAMQCD";
55: N_qt = j.at(XID).at("N_qt");
56: N_phi = j.at(XID).at("N_phi");
57: qt2_MIN = j.at(XID).at("qt2_MIN");
58: qt2_MAX = j.at(XID).at("qt2_MAX");
59: PDF_scale = j.at(XID).at("PDF_scale");
60: alpha_scale = j.at(XID).at("alpha_scale");
61: MAXCOS = j.at(XID).at("MAXCOS");
62:
63: // How calculate sum
64: qt_MIN = math::msqrt(qt2_MIN);
65: qt_MAX = math::msqrt(qt2_MAX);
66:
67: qt_STEP = (qt_MIN - qt_MAX) / N_qt;
68: phi_STEP = (2.0 * math::PI) / N_phi;
69:
70: initialized = true;
71: } catch (...) {}
72: std::string str = "MDurhamParam:ReadParameters: Error parsing " + modelfile +
73: " (Check for extra/missing commas)";
74: throw std::invalid_argument(str);
75: }
76:
77: };
78:
79: class MDurham : public MAmplitudes {
80: public:
81: MDurham(gra::LORENTZSCALAR &lts, const std::string &modelfile);
82: MDurham() {}
83:

```

```

./include/Granitti/MDurham.h      2/2
84: double DurhamQCD(gra:LORENTZSCALAR &Its, const std::string &process);
85: double DQLoop(gra:LORENTZSCALAR &Its, std::vector<std::vector<double>>> &Amp);
86: inline void DScaleChoice(double q2_1, double q1_2, double q2_2, double q1_2_scale,
87: double q2_2_scale) const;
88:
89: inline void Dq2zhic0(const gra:LORENTZSCALAR & Its,
90: std::vector<std::vector<std::complex<double>>> &Amp,
91: const std::vector<double> &q2) const;
92:
93: inline void DHelicity(const std::vector<double> &q1, const std::vector<double> &q2,
94: std::vector<std::complex<double>> &IzF) const;
95:
96: inline std::complex<double> DReIProj(const std::vector<std::complex<double>> &A,
97: const std::vector<std::complex<double>> &IzF) const;
98:
99: void Dq2zg(const gra:LORENTZSCALAR &Its, std::vector<std::vector<std::complex<double>>> &Amp);
100:
101: void Dq2zqbar(const gra:LORENTZSCALAR & Its,
102: std::vector<std::vector<std::complex<double>>> &Amp);
103:
104: void Dq2zMHbar(const gra:LORENTZSCALAR & Its,
105: std::vector<std::vector<std::complex<double>>> &Amp);
106: double phi_CZ(double x, double P0) const;
107: std::vector<double> EvalPhi(int N, int p0) const;
108:
109: double Asum = 0.0;
110: double Nsum = 0.0;
111:
112: private:
113: // Parameters
114: MDurhamParam Param;
115: };
116:
117: // namespace gra
118:
119: #endif

```

```

./include/Granitti/MProcess.h    2/5
84: virtual double EventWeight(const std::vector<double> &randvec, AuxIntData &aux) = 0;
85: virtual bool EventRecord(HepMC3::GenEvent &evt) = 0;
86:
87: // Set central system decay structure
88: void SetDecayMode(const std::string &str);
89: void SetupBranching();
90: void ProcessHelicityTree(MDecayBranch &branch);
91: HELMatrix HelicityDecay(const MParticle &p, const std::vector<MParticle> &daughter) const;
92:
93: // Set initial state
94: void SetInitialState(const std::vector<std::string> &beam, const std::vector<double> &energy);
95:
96: // Set beam energies
97: void SetBeamEnergies(double E1, double E2);
98:
99: // ISOLATE phase space in C++ class processes
100: void SetISOLATE(bool in) { Its.PS_active = in; }
101: bool GetISOLATE() { return Its.PS_active; }
102: void SetFLATMSS2(bool in) {
103: aux:PrintNotice();
104: std::cout << "rang:fg:ired << "MProcess: Set Flat in mass? "
105: << "MProcess: SetFLATMSS2: Set Flat in mass? "
106: "sampling in decay trees. "
107: << "rang:fg:reset << std::endl;
108: FLATMSS2 = in;
109: FLATMSS2_user = true; // user has tagged it
110: }
111: bool GetOFFSHELL(bool in) { return FLATMSS2; }
112: void SetOFFSHELL(double in) {
113: aux:PrintNotice();
114: std::cout << "rang:fg:ired << "MProcess: Set number of "
115: "decay widths in decay trees. "
116: << "rang:fg:reset << std::endl;
117: OFFSHELL = in;
118: OFFSHELL_user = true; // user has tagged it
119: }
120: }
121: void SetOFFSHELL() { return OFFSHELL; }
122:
123: // Set generation spin correlations
124: void SetSPINGEN(const bool SPINGEN) {
125: std::cout << "rang:fg:ired << "MProcess: SetSPINGEN: Set generation 2-1 spin correlations: "
126: << "SPINDEC ? "true" : "false" << "rang:fg:reset << std::endl;
127: for (const auto &x : Its.RESONANCES) { Its.RESONANCES[x.first].SPINGEN = SPINGEN; }
128: }
129: // Set decay spin correlations
130: void SetSPINDEC(const bool SPINDEC) {
131: std::cout << "rang:fg:ired << "MProcess: SetSPINDEC: Set decay 1-2 spin correlations: "
132: << "SPINDEC ? "true" : "false" << "rang:fg:reset << std::endl;
133: for (const auto &x : Its.RESONANCES) { Its.RESONANCES[x.first].SPINDEC = SPINDEC; }
134: }
135: // Set common Lorentz frame for all resonances
136: void SetFRAME(const std::string &FRAME) {
137: std::cout << "rang:fg:ired << "MProcess: SetFRAME: Set common Lorentz "
138: << "FRAME for the resonance amplitudes: "
139: "FRAME << "rang:fg:reset << std::endl;
140: for (const auto &x : Its.RESONANCES) { Its.RESONANCES[x.first].FRAME = FRAME; }
141: }
142: }
143: // Set maximum sliding pomeron helicity for all resonances
144: void SetJMAX(const int &JMAX) {
145: std::cout << "rang:fg:ired << "MProcess: SetJMAX: Set common maximum Pomeron helicity "
146: << "JMAX << "rang:fg:reset << std::endl;
147: for (const auto &x : Its.RESONANCES) { Its.RESONANCES[x.first].JMAX = JMAX; }
148: }
149: // Get initial state
150: std::vector<gra:MParticle> GetInitialState() {
151: std::vector<gra:MParticle> beams = {Its.beam1, Its.beam2};
152: return beams;
153: }
154: // Phase space dimension
155: unsigned int GetLIPEDim() { return ProcPtr.LIPEDim; }
156:
157: // Pomeron loop screening
158: void SetScreening(bool value) { SCREENING = value; }
159: bool GetScreening() { return SCREENING; }
160:
161: // Get/Get input eikonal
162: void SetEikonal(const MEikonal &in) { Eikonal = in; }

```

```

./include/Granitti/MProcess.h    1/5
1: // Abstract process class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MPROCESS_H
7: #define MPROCESS_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: #include <HepMC3>
15: #include "HepMC3/FourVector.h"
16: #include "HepMC3/GenEvent.h"
17: #include "HepMC3/GenParticle.h"
18:
19: // Own
20: #include "Granitti/MVAec.h"
21: #include "Granitti/MBox.h"
22: #include "Granitti/MEikonal.h"
23: #include "Granitti/MGLobals.h"
24: #include "Granitti/MH2.h"
25: #include "Granitti/MH2.h"
26: #include "Granitti/MHEMatrix.h"
27: #include "Granitti/MHMemoria.h"
28: #include "Granitti/MFDD.h"
29: #include "Granitti/MRandom.h"
30: #include "Granitti/MSubProc.h"
31: #include "Granitti/MSubProc.h"
32: #include "Granitti/MSubProc.h"
33:
34: namespace gra {
35:
36: // Event-by-event auxiliary data for integration
37: struct AuxIntData {
38: // Aux weight
39: double vegasweight = 1.0;
40: bool Burn_in_mode = false;
41:
42: // Event-by-event assertions (init all with true!)
43: bool amplitude_ok = true;
44: bool kinematics_ok = true;
45: bool fidcuts_ok = true;
46: bool vetocuts_ok = true;
47:
48: // Forced acceptance of the event
49: bool forced_accept = false;
50:
51: bool Valid() const { return kinematics_ok && fidcuts_ok && vetocuts_ok; }
52: };
53:
54: // Multipomeron kinematics
55: struct MPI {
56: MVAec p1;
57: MVAec p2;
58:
59: MVAec q1;
60: MVAec q2;
61:
62: MVAec k;
63:
64: MVAec p1f;
65: MVAec p2f;
66:
67: MVAec p3;
68: MVAec p4;
69: };
70:
71: // Abstract process class
72: class MProcess: public MUserHistograms {
73: public:
74: // or polymorphic type
75: virtual ~MProcess() {} // MUST HAVE IT HERE as virtual
76:
77: // Pure virtual functions, without definitions here
78: // WITHOUT =0, undefined reference to vtable for MProcess will occur in
79: // compilation
80: virtual void post_Constructor() = 0;
81: virtual void PrintInit(bool silent) const = 0;
82:
83: virtual double operator()(const std::vector<double> &randvec, AuxIntData &aux) = 0;

```

```

./include/Granitti/MProcess.h    3/5
167: MEikonal GetEikonal() const { return Eikonal; }
168:
169: // Set LMAPDET name
170: void SetLMAPDET(const std::string &in) {
171: std::cout << "MProcess: SetLMAPDET: " << in << std::endl;
172: Its.LMAPDET = in;
173: }
174:
175: // Set cuts
176: void SetGenCuts(const gra:GENCUT &in) { gcuts = in; }
177: void SetFIDCuts(const gra:FIDCUT &in) { fcuts = in; }
178: void SetUserCuts(int in) { USERCUTS = in; }
179: void SetVetoCuts(const gra:VETOCUT &in) { vetocuts = in; }
180:
181: double GetMandelstam_s() const { return Its.s; }
182:
183: // Set proton excitation to low-mass N*
184: void SetExcitation(int in) {
185: if (in > 2) in = 0;
186: std::string str =
187: "MProcess: SetExcitation: Not valid input "
188: "0,1,2,3,4
189: std::to_string(in) + " ";
190: throw std::invalid_argument(str);
191: }
192: EXCITATION = in;
193:
194: if (EXCITATION > 0) {
195: aux:PrintWarning();
196: std::cout << "rang:fg:ired << "MProcess: SetExcitation: Proton "
197: << "MProcess: SetExcitation: Proton "
198: "excitation is under construction / some processes contain only kinematic part: "
199: << in << "rang:fg:reset << std::endl;
200: }
201: }
202:
203: // Set flat matrix element mode
204: void SetFLATAMP(int in) {
205: if (in > 0) {
206: aux:PrintNotice();
207: std::cout << "rang:fg:ired << "MProcess: SetFLATAMP: Flat matrix element FLATAMP: " << in
208: << "rang:fg:reset << std::endl;
209: FLATAMP = in;
210: }
211: }
212:
213: // pp invariant Holler flux (high energy limit)
214: double HOLLERFLUX(int &in) { return 2.0 * Its.s; }
215:
216: // Flat amplitudes (for HENUS)
217: double GetFlatAmp2(const gra:LORENTZSCALAR &Its) const;
218:
219: // Set input resonances
220: void SetResonances(const std::map<std::string, gra:PARAM_RES &in> (Its.RESONANCES = in);
221:
222: // Get input resonances
223: std::map<std::string, gra:PARAM_RES> GetResonances() const { return Its.RESONANCES; }
224:
225: // Eikonal (screening) functions
226: MEikonal Eikonal;
227:
228: // Lets keep these public for easy access
229: gra:LORENTZSCALAR Its; // Lorentz scalars and others for kinematics
230:
231: // Cut structures
232: gra:GENCUT gcuts; // Generator sampling cuts (phase space boundaries)
233: gra:FIDCUT fcuts; // Fiducial cuts (phase space boundaries)
234: gra:VETOCUT vetocuts; // Veto cuts
235:
236: void PrintDecayTree(const gra:MDecayBranch &branch) const;
237:
238: void CalculatePhaseSpace(const gra:MDecayBranch &branch, double &product, double &product2p,
239: double &volume, int &N_final) const;
240: void PrintPhaseSpace(const gra:MDecayBranch &branch, double &product, double &product2p,
241: int &N_final) const;
242:
243: // Random numbers (keep it public for testing)
244: MRandom random;
245:
246: // Subprocess (amplitudes)
247: MSubProc ProcPtr;
248:
249: protected:

```

```

./include/Granitti/MProcess.h      4/5
250: // Copy and assignment made private
251: // MProcess(const MProcess other);
252: // MProcess operator=(const MProcess rhs);
253:
254: // Internal virtual functions, without definitions here
255: virtual void Initialize() = 0;
256: virtual bool LoopKinematics(const std::vector<double> &kp, const std::vector<double> &kp2) = 0;
257: virtual bool FiducialCuts() const = 0;
258:
259: // Cascade phase-space factor
260: double CascadeP() const;
261:
262: // Amplitude squared
263: double GetAmp2();
264:
265: // Eikonal screening loop
266: double S3ScreenedAmp();
267:
268: // First print
269: void PrintSetup() const;
270:
271: // Setup process
272: void SetProcess(std::string sProcess, const std::vector<aux::iOneCMD> &syntax);
273:
274: // QFT symmetry factor
275: void CalculateSymmetryFactor();
276:
277: // Recursive function to treat decay trees
278:
279: void SaveBranch(HepMC3::GenEvent &evt, const gra::HDecayBranch &branch,
280:                const HepMC3::GenParticlePtr &pp);
281:
282: bool CommonRecord(HepMC3::GenEvent &evt);
283: bool VetOCuts() const;
284: bool CommonCuts() const;
285: void FindDecayCuts(const gra::HDecayBranch &branch, bool &ok) const;
286: void FindVetOCuts(const gra::HDecayBranch &branch, bool &ok) const;
287: bool ConstructDecayKinematics(gra::HDecayBranch &branch);
288: void WriteDecayKinematics(const gra::HDecayBranch &branch, const HepMC3::GenParticlePtr &another,
289:                            HepMC3::GenEvent &evt);
290: void PrintFiducialCuts() const;
291:
292: void GetOffShellMass(const gra::HDecayBranch &branch, double &mass);
293: void SetTechnicalBoundaries(gra::iGENCUT &cuts, unsigned int EXCITATION);
294: double ForwardVolume() const;
295:
296: // Lorentz scalars
297: bool GetLorentzScalars(unsigned int Nf);
298:
299: void SampleForwardMasses(std::vector<double> &mvct, const std::vector<double> &randvec);
300:
301: // -----
302: // System fragmentation
303:
304: bool ExciteNstar(const MVec &nstar, gra::HDecayBranch &forward, const HParticle &pbarn);
305: bool ExciteContinuum(const MVec &nstar, gra::HDecayBranch &forward, double Q2scale, int B_sum,
306:                      int Q_sum, const std::string &pt_distribution = "powexp");
307: void BranchForwardSystem(const std::vector<MVec> &ps, const std::vector<HParticle> &sp,
308:                          const MVec &nstar, gra::HDecayBranch &forward);
309: bool CEPForwardFragment();
310:
311: // -----
312: void ParseCMD(const std::string &str, std::string &first, std::string &second,
313:              std::string &third) const;
314:
315: // Check std::nan/std::inf
316: bool CheckInNaN(double W) {
317:     if (std::isnan(W)) {
318:         ++N_nan;
319:         W = 0;
320:         return false;
321:     } else if (std::isinf(W)) {
322:         ++N_inf;
323:         W = 0;
324:         return false;
325:     }
326: }
327: if (N_nan > 50) {
328:     throw std::invalid_argument(
329:         "MProcess::CheckInNaN: Too many NaN weights "
330:         "  - Check model parameters "
331:         "  - and cuts!");
332: }

```

```

./include/Granitti/MUserCuts.h      1/1
1: // Custom user defined cuts
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MUSERCUTS_H
7: #define MUSERCUTS_H
8:
9: // C++
10: #include <vector>
11:
12: // Own
13: #include "Granitti/MKinematics.h"
14:
15: namespace gra {
16: // User cuts (return false for events not passing the cuts)
17: bool UserCut(int id, const gra::tLorentzScalar &ts);
18:
19: // namespace gra
20:
21: #endif

```

```

./include/Granitti/MProcess.h      5/5
333: if (N_inf > 50) {
334:     throw std::invalid_argument(
335:         "MProcess::CheckInNaN: Too many Inf weights "
336:         "  - Check model parameters "
337:         "  - and cuts!");
338: }
339: return true;
340: }
341: unsigned int N_inf = 0;
342: unsigned int N_nan = 0;
343:
344: // Cross section statistical 1/S symmetry factor (for identical final states)
345: double S_factor = 0.0;
346:
347: // -----
348: // Steering parameters
349:
350: std::string PROCESS; // Process identifier string
351: std::string CID; // phase space sampler identifier such as "T" or "C"
352: std::string DECAYMODE; // Decaymode identifier string
353: bool SCREENING = false; // Pomeron loop on/off
354: int EXCITATION = 0; // Forward proton excitation (0 = off, 1 = single, 2 = double)
355: int USERCUTS = 0; // User custom cuts identifier
356: int FLATAMP = 0; // Flat matrix element mode
357:
358: // -----
359: // Phase-space control
360:
361: bool FLATM2S2 = false; // Flat in M^2 instead of Breit-Wigner sampling
362: bool FLATM2S2_user = false;
363: bool OFFSHELL = 5; // How many full widths to sample particles in cascades
364: bool OFFSHELL_user = false;
365:
366: // Forward excitation minimum/maximum M^2 boundaries
367: double M2_f_min = 0.0;
368: double M2_f_max = 0.0;
369: double log_M2_f_min = 0.0;
370: double log_M2_f_max = 0.0;
371:
372: static constexpr double ZERO_EPS = 1e-12; // To use with log(0+ZERO_EPS)
373: // -----
374: // Non-Diffraction
375: std::vector<CMD> etree;
376: double bc = 0.0;
377:
378: };
379:
380: // namespace gra
381:
382: #endif

```

```

./include/Granitti/MPDG.h           1/2
1: // PDG Class
2: //
3: // (c) 2017-2019 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MPDG_H
7: #define MPDG_H
8:
9: // C++
10: #include <complex>
11: #include <fstream>
12: #include <map>
13: #include <random>
14: #include <vector>
15:
16: // Own
17: #include "Granitti/MAux.h"
18: #include "Granitti/MParticle.h"
19:
20: namespace gra {
21: namespace PDG {
22:
23: // Common PDG ids, http://pdg.lbl.gov/2007/reviews/montecarlopp.pdf
24: // For MC internal, use 01-000
25: constexpr int PDG_p = 2212; // Proton
26: constexpr int PDG_n = 2212; // ...
27: constexpr int PDG_pip = 2112;
28: constexpr int PDG_pim = -2112;
29: constexpr int PDG_pi0 = 111;
30: constexpr int PDG_kp = 321;
31: constexpr int PDG_ks = -321;
32: constexpr int PDG_gamma = 22;
33: constexpr int PDG_gluon = 21;
34: constexpr int PDG_muon = 11;
35: constexpr int PDG_neutrino = 113;
36:
37: constexpr int PDG_pomeron = 990;
38: constexpr int PDG_regegon = 110;
39: constexpr int PDG_monopole = 992; // spin-1/2 monopole
40: constexpr int PDG_NSTAR = 90000; // Wexon resonance
41: constexpr int PDG_NSTAR = 90210; // Baryon resonance
42: constexpr int PDG_system = 90; // "Contains" / "Not" / "system"
43: constexpr int PDG_fragment = 91; // Proton fragment
44: constexpr int PDG_propagator = 99; // Generic propagator
45:
46: // HepMC STABLE definitions conventions
47: constexpr int PDG_BEAM = 4; // Beam particles
48: constexpr int PDG_STABLE = 1; // Final state "stable" particles
49: constexpr int PDG_DECAY = 2; // Such as pi0
50: constexpr int PDG_INTERMEDIATE = 81; // Such as gamma, pomeron etc.
51:
52:
53: // For CPU efficiency [GeV]
54: constexpr double mp = 0.938272081;
55: constexpr double mn = 0.93827;
56: constexpr double mpi = 0.13957018;
57: constexpr double mpi0 = 0.13497660;
58:
59: // [hbar] = [M][L]^2[T]^-1 and [c] = [L][T]^-1
60: // set [hbar] = c = 1
61: // then [M] = [L]^2 = T^-1
62:
63: // http://pdg.lbl.gov/2018/reviews/rpp2018-rev-physics-constants.pdf
64: // Basic constants
65: constexpr double c = 2.99792458E8; // c [m/s] (EXACT/DEFINITION)
66: constexpr double hbar = 6.58211954E-25; // hbar = [GeV*s]
67: constexpr double eV = 1.6021766208E-19; // e = [Joule], e*0.303 in nat.u.
68:
69: // Basic definition
70: constexpr double barn2m = 1E-28; // 1 barn to [m^2]
71:
72: // Standard conversions
73: constexpr double GeV2m = hbar * c; // 1 GeV^-1 to [m]
74: constexpr double GeV2fm = GeV2m * 1E15; // 1 GeV^-1 to [fermi] (1 fm)
75: constexpr double GeV2m2 = GeV2m * GeV2m; // 1 GeV^-2 to [m^2]
76: constexpr double GeV2a = hbar; // 1 GeV^-1 to [a]
77:
78: constexpr double GeV2barn = GeV2m2 / barn2m; // 1 GeV^-2 to barns
79: constexpr double GeV2mb = GeV2barn * 1E3; // 1 GeV^-2 to mbarns
80: // constexpr double GeV2ub = GeV2barn * 1E6; // 1 GeV^-2 to microbarn
81: // constexpr double GeV2nb = GeV2barn * 1E9; // 1 GeV^-2 to nanobarn
82: // constexpr double GeV2pb = GeV2barn * 1E12; // 1 GeV^-2 to picobarns
83: // constexpr double mb2GeV = 1.0 / GeV2mb; // 1 mb to GeV^-2

```

```

./include/Granitti/MPDG.h          2/2
84:
85: // More conversions
86: //constexpr double GeV2J  = eV * 1e9; // 1 GeV [1] to [kg*m^2/s^2]=Joule
87: //constexpr double GeV2kg = GeV2J / GeV2m * pow(GeV2s); // 1 GeV [1] to [kg]
88: //constexpr double GeV2N  = GeV2J / GeV2m; // 1 GeV [2] to [N]=kg*m/s^2
89: // (force)
90: //constexpr double GeV2mom = GeV2J / GeV2m * GeV2a; // 1 GeV [1] to [kg*m/s]
91: // (Momentum)
92:
93: // cross section: [value] x [GeV^-2] = [value] x [hbar c]^2
94: // decay rates: [value] x [GeV] = [value] / hbar
95: // length: [value] x [GeV^-1] = [value] x [hbar x c]
96:
97: // Mason decay constants [PDG] [p10, p1+, K+, K0]
98: static const std::map<int, double> DM_meson{
99:     {111, 0.1300}, {211, 0.1307}, {321, 0.1598}, {311, 0.1598}};
100:
101: | // namespace PDG
102:
103: class MPDG {
104: public:
105:     MPDG();
106:     ~MPDG();
107:
108:     void ReadParticleData(const std::string &filepath);
109:     void TokenizeProcess(const std::string &str, int depth);
110:     const gra::MParticle &MDecayBranch(const std::string &branchname) const;
111:     bool IsDecay(const std::string &str) const;
112:
113:     void PrintPDGTable() const;
114:     const gra::MParticle &FindByPDG(int pdgcode) const;
115:     const gra::MParticle &FindByPDGname(const std::string &pdgname) const;
116:
117:     // PDG tables
118:     std::map<int, gra::MParticle> PDG_Table;
119: };
120:
121: | // namespace gra
122:
123: #endif

```

```

./include/Granitti/MEikonal.h      1/4
1: // Eikonal proton density and Screening amplitude class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MEIKONAL_H
7: #define MEIKONAL_H
8:
9: // C++
10: #include <complex>
11: #include <vector>
12:
13: // Own
14: #include "Granitti/M4Vec.h"
15: #include "Granitti/MFom.h"
16: #include "Granitti/MMatrix.h"
17: #include "Granitti/MMatrix.h"
18: #include "Granitti/MPDG.h"
19:
20: // External
21: #include "rang.hpp"
22:
23: namespace gra {
24:
25: // Numerical integration parameters
26: struct MEikonalNumerics {
27:     static constexpr double MinKt2 = 1E-6;
28:     static constexpr double MaxKt2 = 25.0;
29:     unsigned int N = 0;
30:     bool logKt = false;
31:
32:     static constexpr double MinM = 1E-6;
33:     static constexpr double MaxM = 10.0 / PDG::GeV2m;
34:     unsigned int Npart = 0;
35:     bool logM = false;
36:
37:     static constexpr double FIntegralMinKt = 1E-9;
38:     static constexpr double FIntegralMaxKt = 30.0;
39:     static constexpr unsigned int FIntegralN = 10000;
40:     static constexpr double MinLoopKt = 1E-4;
41:
42:     double MaxLoopKt = 1.75;
43:
44:     unsigned int NpartLoopKt = 15; // Number of kt steps (default minimum)
45:     unsigned int NpartLoopM = 12; // Number of m steps (default minimum)
46:
47:     // User setup (ND can be negative, to get below the default)
48:     void SetLoopIntegration(int ND) {
49:         NpartLoopKt = std::max(3 * ND + (int)NpartLoopKt);
50:         NpartLoopM = std::max(3 * ND + (int)NpartLoopM);
51:
52:         if (ND < 0) { std::cout << "rang::if::iread }";
53:         std::cout << "MEikonalNumerics::SetLoopIntegration: ND = " << ND << " << endl;
54:         std::cout << "NpartLoopKt = " << NpartLoopKt << " << endl;
55:         std::cout << "NpartLoopM = " << NpartLoopM << " << endl;
56:         std::cout << "rang::if::iread }";
57:     }
58:
59:     // Unique hash
60:     std::string GetHashString() {
61:         std::string str = std::ito_string(MinKt2) + std::ito_string(MaxKt2) + std::ito_string(NpartLoopKt) +
62:             std::ito_string(logKt) + std::ito_string(MinM) + std::ito_string(MaxM) +
63:             std::ito_string(NpartLoopKt) + std::ito_string(logM) +
64:             std::ito_string(FIntegralMinKt) + std::ito_string(FIntegralMaxKt) +
65:             std::ito_string(FIntegralN);
66:         return str;
67:     }
68:
69:     // Read parameters from file
70:     void ReadParameters() {
71:         // Read and parse
72:         using json = nlohmann::json;
73:         const std::string inputfile = gra::aux::GetBasePath(2) + "/modeldata/" + "NUMERICS.json";
74:         const std::string data = gra::aux::GetInputData(inputfile);
75:         json j;
76:
77:         try {
78:             j = json::parse(data);
79:
80:             // JSON block identifier
81:             const std::string XID = "NUMERICS_EIKONAL";
82:
83:

```

```

./include/Granitti/MRandom.h      1/1
1: // Random number class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MRANDOM_H
7: #define MRANDOM_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: namespace gra {
15: class MRandom {
16: public:
17:     // Calling constructors of gaussian functions
18:     MRandom(); flat(0, 1), gaussian(0, 1) {
19:         rng.seed(); // default initialization
20:     }
21:     ~MRandom() {}
22:
23:     // Set random number engine seed
24:     void SetSeed(int seed) {
25:         const int SEEDMAX = 2147483647;
26:         if (seed > SEEDMAX) {
27:             std::string str = "MRandom::SetSeed: Invalid input seed: " + std::ito_string(seed) +
28:                 " > SEEDMAX = " + std::ito_string(SEEDMAX);
29:             throw std::invalid_argument(str);
30:         }
31:         rng.seed(seed);
32:         RNSEED = seed;
33:     }
34:
35:     // Return current random seed
36:     unsigned int GetSeed() const { return RNSEED; }
37:
38:     // Random sampling functions
39:     double U(double a, double b);
40:     double G(double mu, double sigma);
41:     double PowerRandom(double a, double b, double alpha);
42:     double RelativisticGammaRandom(double m0, double Gamma, double LIMIT = 5.0, double M_MIN = 0.0);
43:     double CauchyRandom(double m0, double Gamma, double LIMIT = 5.0, double M_MIN = 0.0);
44:     int NBRRandom(double avgM, double k, int maxvalue);
45:     int PoissonRandom(double lambda);
46:     double ExpRandom(double lambda);
47:     int LogRandom(double p, int maxvalue);
48:     void D1Random(const std::vector<double> &alpha, std::vector<double> &gamma);
49:
50:     double NBRRandf(int n, double avgM, double k);
51:     double Logpdff(int k, double p);
52:
53:     // For generators, check: https://nullprogram.com/blog/2017/09/21/
54:
55:     // 44-bit Mersenne Twister by Matsumoto and Nishimura, fast, basic
56:     // For some (possible) source problems, see:
57:     // REFERENCE: Harase, https://arxiv.org/abs/1708.06018
58:     // std::mt19937_64 rng;
59:
60:     // 48-bit RANDUX (a bit slower)
61:     // REFERENCE: Luecher, https://arxiv.org/abs/hep-lat/9309020
62:     std::ranlux48 rng;
63:
64:     // Distribution engines
65:     std::uniform_real_distribution<double> flat;
66:     std::normal_distribution<double> gaussian;
67:     unsigned int RNSEED = 0; // Random seed set
68:
69: private:
70:     // Nothing
71:
72:
73: } // namespace gra
74:
75: #endif

```

```

./include/Granitti/MEikonal.h      2/4
84:     NumberKt2 = j.at(XID).at("NumberKt2");
85:     logKt2 = j.at(XID).at("logKt2");
86:     NumberM = j.at(XID).at("NumberM");
87:     logM = j.at(XID).at("logM");
88:
89:     } catch (...) {
90:         std::string str =
91:             "MEikonalNumerics: Error parsing " + inputfile + " (Check for extra/missing commas)";
92:         throw std::invalid_argument(str);
93:     }
94:
95: };
96:
97:
98: // Interpolation container
99: class IArrayID {
100: public:
101:     IArrayID();
102:     ~IArrayID();
103:
104:     std::string name;
105:     double MIN = 0;
106:     double MAX = 0;
107:     unsigned int N = 0;
108:     double STEP = 0;
109:     bool islog = false;
110:
111:     // Setup discretization
112:     void Setf(std::string &name, double_min, double_max, double_N, double_logarithmic) {
113:         // AT least two intervals
114:         if (N < 2) {
115:             throw std::invalid_argument("IArrayID::Set: Error: N = " + std::ito_string(N) + " < 2");
116:         }
117:
118:         // Check sanity
119:         if (MIN >= MAX) {
120:             throw std::invalid_argument("IArrayID::Set: Error: Variable " + name + " MIN = " +
121:                 std::ito_string(MIN) + " >= MAX = " + std::ito_string(MAX));
122:         }
123:
124:         // Check sanity
125:         if (logarithmic && min < 1e-9) {
126:             throw std::invalid_argument("
127:             "IArrayID::Set: Error: Variable " + name +
128:             " is using logarithmic stepping with boundary MIN = " + std::ito_string(MIN));
129:         }
130:
131:         islog = logarithmic;
132:
133:         // Logarithm taken here!
134:         MIN = islog ? std::log(MIN) : MIN;
135:         MAX = islog ? std::log(MAX) : MAX;
136:         N = N - 1;
137:         STEP = (MAX - MIN) / N;
138:         name = &name;
139:     }
140:
141:
142:     // Call this last
143:     void InitArray() {
144:         // Note: F is
145:         F = MMatrix<std::complex<double>>(N + 1, 2, 0.0);
146:
147:         // N: size
148:         MMatrix<std::complex<double>> F;
149:
150:         // Use energy (needed for boundary conditions)
151:         double sqrts = 0.0;
152:
153:         std::string GetHashString() const {
154:             std::string str = std::ito_string(islog) + std::ito_string(MIN) + std::ito_string(MAX) +
155:                 std::ito_string(N) + std::ito_string(sqrts);
156:             return str;
157:         }
158:
159:         bool WriteArray(const std::string &filename, bool overwrite) const;
160:         bool ReadArray(const std::string &filename);
161:         std::complex<double> InterpolateID(double a) const;
162:
163:         static const unsigned int X = 0;
164:         static const unsigned int Y = 1;
165:
166: };

```

```

./include/Granitti/MEikonal.h      3/4
167:
168: class MEikonal {
169: public:
170:     MEikonal();
171:     ~MEikonal();
172:
173:     // Construct outside
174:     void S3Constructor(double s_in, const std::vector<grai::MParticle> &initialState_in,
175:                       bool onlyDensity = false, int NumberRT2 = 0);
176:
177:     // Get ekonal and amplitude values
178:     // std::complex<double> S3DensityInterpolator(double bt);
179:     // std::complex<double> S3ScreeningInterpolator(double kt2);
180:
181:     // Initialization already done
182:     bool TIsInitialized() const {
183:         if (S3INIT == true) { return true; }
184:         return false;
185:     }
186:
187:     // Get total cross sections
188:     void GetTotXS(double ktot, double sel, double iin) const;
189:
190:     static std::complex<double> SingleAmpLastic(double s, double t, int type);
191:     static std::complex<double> S3Density(double bt) const;
192:     static std::complex<double> S3Screening(double kt2) const;
193:
194:     // Get random number of cut Pomerons
195:     template <typename T>
196:     void S3GetRandomCutsBT(unsigned int em, double sht, T krng) {
197:         const double STEP = (Numerics.MaxBT - Numerics.MinBT) / Numerics.NumberBT;
198:
199:         // Numerical integral loop over impact parameter (b,t) space
200:         // C++11, thread_local is also static
201:         thread_local std::uniform_real_distribution<double> flat(0, 1);
202:         thread_local std::uniform_int_distribution<unsigned int> randint(0, Numerics.NumberBT);
203:         thread_local std::uniform_int_distribution<unsigned int> randint1(MCUT - 1); // 1,2,...
204:
205:         // Acceptance-Rejection
206:         unsigned int n = 0;
207:         while (true) {
208:             // Draw random impact parameter flat
209:             n = randint(rng);
210:
211:             const int value = randint(rng);
212:             if (flat(rng) < P_Array[value][n]) {
213:                 n = value; // m cut Pomerons
214:                 break;
215:             }
216:         }
217:         bt = Numerics.MinBT + n * STEP;
218:     }
219:
220:     // Get random number of cut Pomerons
221:     template <typename T>
222:     unsigned int S3GetRandomCuts(T krng) {
223:         // Random integer from [1, Nbins-1]
224:         // C++11, thread_local is also static
225:         thread_local std::uniform_real_distribution<double> flat(0, 1);
226:         thread_local std::uniform_int_distribution<unsigned int> RANDI(1, P_cut.size() - 1);
227:
228:         // Acceptance-Rejection
229:         while (true) {
230:             const unsigned int m = RANDI(rng);
231:             if (flat(rng) < P_cut[m]) {
232:                 return m; // m cut Pomerons
233:             }
234:         }
235:     }
236:
237:     // Arrays
238:     IArrayD MBT;
239:     IArrayD MSA;
240:
241:     // Parameters
242:     MEikonalNumerics Numerics;
243:
244: private:
245:     static const unsigned int MCUT = 25; // Maximum number of cut Pomerons
246:
247:     // S3 sett survival initialized
248:     bool S3INIT = false;
249:

```

```

./include/Granitti/MEikonal.h      4/4
250:     void S3CalculateArray(IArrayD kadd, std::complex<double> &(MEikonal::**f)(double) const);
251:     void S3CalcXS();
252:     void S3InitCutPomerons();
253:
254:     // Handwritten s
255:     double s = 0.0;
256:
257:     // Initial state
258:     std::vector<grai::MParticle> INITIALSTATE;
259:
260:     // Integrated ekonal based total cross sections
261:     double sigma_tot = 0.0;
262:     double sigma_el = 0.0;
263:     double sigma_inel = 0.0;
264:
265:     // Cut Pomerons probabilities
266:     std::vector<double> P_cut;
267:     std::vector<std::vector<double>> P_Array;
268:
269:     // Two-Channel ekonal based cross sections
270:     double sigma_diff[4] = {0.0};
271:
272:     static const unsigned int X = 0;
273:     static const unsigned int Y = 1;
274: };
275:
276: // namespace gra
277:
278: #endif

```

```

./include/Granitti/MTensor.h      1/2
1: // Minimal tensor class
2:
3: // Example: rank-6 tensor with dim=4 per dimension
4:
5: // MTensor<double> tensor = MTensor({4,4,4,4,4,4}, 0.0);
6: // tensor({0,3,0,1,2}) = 1.0;
7:
8:
9: // (c) 2017-2020 Mikael Mieskolainen
10: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
11:
12: #ifndef MTENSOR_H
13: #define MTENSOR_H
14:
15: #include <initializer_list>
16: #include <iomanip>
17: #include <iostream>
18:
19: namespace gra
20: {
21:     template <typename T>
22:     class MTensor {
23:     public:
24:         MTensor() { data = nullptr; }
25:         MTensor(const std::vector<std::size_t> &newdim) {
26:             dim = newdim;
27:             data = new T[Prod(dim)]; // No initialization
28:         }
29:         MTensor(const std::vector<std::size_t> &newdim, T value) {
30:             dim = newdim;
31:             data = new T[Prod(dim)];
32:             std::fill(data, data + Prod(dim), T(value)); // Initialization
33:         }
34:         MTensor() {
35:             // Delete dynamically allocated memory
36:             delete[] data;
37:         }
38:         // Copy constructor
39:         MTensor(const MTensor &a) {
40:             dim = a.dim;
41:             data = new T[Prod(dim)];
42:         }
43:         // Copy all elements
44:         Copy(a);
45:     };
46:     // Assignment operator
47:     MTensor &operator=(const MTensor &rhs) {
48:         if (data != rhs.data && rhs.data != nullptr) {
49:             ReSize(rhs.dim);
50:             Copy(rhs);
51:         }
52:         return *this;
53:     }
54:     T &operator()(const std::vector<std::size_t> &ind) {
55:         const std::size_t i = Index(ind);
56:         return data[i];
57:     }
58:     T &operator()(const std::vector<std::size_t> &ind) const {
59:         const std::size_t i = Index(ind);
60:         return data[i];
61:     }
62:     // Size operators
63:     std::size_t size(const std::size_t ind) const { return dim[ind]; }
64:
65:     /*
66:     // Print full tensor
67:     void Print() {
68:
69:         X(, , 1, 1, 1) =
70:         std::vector<std::size_t> state(dim.size()-2, 0);
71:
72:         while (true) {
73:             for (std::size_t k = 2; k < dim.size(); ++k) {
74:                 }
75:
76:             for (std::size_t i = 0; i < dim[0]; ++i) {
77:                 for (std::size_t j = 0; j < dim[1]; ++j) {
78:                     const std::size_t ind = Index({i,j});
79:                     std::cout <<
80:
81:                 }
82:             }
83:

```

```

./include/Granitti/MTensor.h      2/2
84:         std::cout << data[i]
85:         ,
86:         //
87:         );
88:     private:
89:         // Re-allocate memory
90:         // Row-major order
91:         std::size_t Index(const std::vector<std::size_t> &ind) const {
92:             if (ind.size() != dim.size()) {
93:                 throw std::invalid_argument(
94:                     "MTensor: Error: Index vector with rank = " + std::to_string(ind.size()) +
95:                     " c.f. Tensor has rank = " + std::to_string(dim.size()));
96:             }
97:             std::size_t sum = 0;
98:             for (std::size_t i = 0; i < ind.size(); ++i) {
99:                 if (ind[i] == dim[i]) {
100:                     throw std::invalid_argument("MTensor: Error: Input index " + std::to_string(i) +
101:                                                 " over bounds: " + std::to_string(ind[i]) +
102:                                                 " = " + std::to_string(dim[i]));
103:                 }
104:                 std::size_t product = 1;
105:                 for (std::size_t j = i + 1; j < ind.size(); ++j) { product *= dim[j]; }
106:                 sum += product * ind[i];
107:             }
108:             return sum;
109:         }
110:         // Copy data from a to *this (after ReSize)
111:         void Copy(const MTensor &a) {
112:             T *p = data + Prod(dim);
113:             T *q = a.data + Prod(dim);
114:             while (p > data) { *p = *q; }
115:         }
116:         // Re-allocate
117:         void ReSize(const std::vector<std::size_t> &newdim) {
118:             if (data != nullptr) { delete[] data; }
119:             dim = newdim;
120:             data = new T[Prod(newdim)];
121:         }
122:         // Product to get array volume (memory)
123:         std::size_t Prod(const std::vector<std::size_t> &x) const {
124:             std::size_t product = 1;
125:             for (std::size_t i = 0; i < x.size(); ++i) { product *= x[i]; }
126:             return product;
127:         }
128:
129:         // Dimensions
130:         std::vector<std::size_t> dim;
131:
132:         T *data;
133:
134:     };
135: } // namespace gra
136:
137: #endif

```

```

./include/Granitti/MParticle.h      1/2
1: // Particle and branch objects [HEADER Only file]
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MPARTICLE_H
7: #define MPARTICLE_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <valarray>
13: #include <vector>
14:
15: #include "Granitti/MAux.h"
16: #include "Granitti/MCW.h"
17: #include "Granitti/MHELMMatrix.h"
18: #include "Granitti/MMatrix.h"
19:
20: namespace gra {
21:
22:
23: // Particle class
24: struct MParticle {
25: // These values are read from the PDG-file
26: std::string name;
27: int pdg = 0; // PDG code
28: int chargeQ3 = 0; // 0 x 3
29: int spinX2 = 0; // 0 x 2
30: int color = 0; // QCD color code
31:
32: double mass = 0.0;
33: double width = 0.0;
34: double tau = 0.0; // hbar / width
35:
36: // J^PC
37: int P = 1; // Default even P-parity
38: int C = 0; // Default zero C-parity (e.g. fermions do not have)
39: unsigned int L = 0; // For Mesons/Baryons
40: bool glue = false; // Glueball state
41:
42: void setPCL(int _P, int _C, unsigned int _L) {
43: P = _P;
44: C = _C;
45: L = _L;
46: }
47:
48: // Width cut (in Breit-Wigner sampling)
49: double wcut = 0.0;
50:
51: void print() const {
52: std::cout << " NAME: " << name << std::endl;
53: std::cout << " ID: " << pdg << std::endl;
54: std::cout << " M: " << mass << std::endl;
55: std::cout << " W: " << width << std::endl;
56: std::cout << " J^PC: " << aux:SpinToString(spinX2) << " " << aux:ParityToString(P)
57: << aux:ParityToString(C) << std::endl
58: << std::endl;
59: }
60: };
61:
62:
63: // Recursive decay tree branch
64: struct MDecayBranch {
65: MDecayBranch() {
66: f = MMatrix(std::complex<double>)(1, 1, 1.0); // Unit element
67: }
68:
69: // Offshell mass picked event by event
70: double m_offshell = 0.0;
71:
72: MParticle pj; // PDG particle
73: M4Vec p4; // 4-momentum
74: std::vector<MDecayBranch> legs; // Daughters
75: M4Vec decay_position; // Decay 4-position
76: gra:HELMMatrix hel; // Decay helicity information
77:
78: // Used with helicity amplitudes
79: MMatrix(std::complex<double>) f;
80:
81: // MC weight container
82: gra:kinematics:MCW w;
83:

```

```

./include/Granitti/MParticle.h      2/2
84: // Active in the factorized phase space product (by default, no)
85: // This is controlled by the specific amplitudes
86: bool P2_active = false;
87:
88: // Decay tree current level
89: int depth = 0;
90: };
91:
92: } // namespace gra
93:
94: #endif

```

```

./include/Granitti/MContinuum.h     1/1
1: // Continuum 2->N phase space class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MCONTINUUM_H
7: #define MCONTINUUM_H
8:
9: // C++
10: #include <complex>
11: #include <vector>
12:
13: // Own
14: #include "Granitti/M4Vec.h"
15: #include "Granitti/MAux.h"
16: #include "Granitti/MMatrix.h"
17: #include "Granitti/MProcess.h"
18: #include "Granitti/MSpin.h"
19:
20: // HepMC3
21: #include "HepMC3/GenEvent.h"
22:
23: namespace gra {
24: class MContinuum : public MProcess {
25: public:
26: MContinuum();
27: MContinuum(std::string process, const std::vector<aux:OneCMD> &syntax);
28: virtual ~MContinuum();
29:
30: void post_Constructor();
31:
32: double operator()(const std::vector<double> &randvec, AuxInData &aux) {
33: return EventWeight(randvec, aux);
34: }
35:
36: double EventWeight(const std::vector<double> &randvec, AuxInData &aux);
37: bool EventRecord(HepMC3:GenEvent &evt);
38: void PrintInit(bool silent) const;
39: private:
40: void Initialize();
41: bool LoopKinematics(const std::vector<double> &ip1, const std::vector<double> &ip2p);
42: bool FiducialCut() const;
43:
44: // 3^N-4 dimensional phase space, 2->N
45: bool BNRandomIn(assigned int Nf, const std::vector<double> &randvec);
46: bool BNBuildIn(assigned int Nr, double pt1, double pt2, double phi1, double phi2,
47: const std::vector<double> &kt, const std::vector<double> &phi,
48: const std::vector<double> &v, double m1, double m2);
49:
50: void BilinearSystem(const std::vector<M4Vec> &p, const std::vector<M4Vec> &q, const M4Vec &ip1,
51: const M4Vec &ip2) const;
52:
53: double BNIntegralVolume() const;
54: double BNPhaseSpaceWeight() const;
55:
56: // Auxiliary (kt) vectors
57: std::vector<M4Vec> pkt_;
58: };
59:
60: } // namespace gra
61:
62: #endif

```

```

./include/Granitti/MResonance.h     1/2
1: // Resonance container class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MRESONANCE_H
7: #define MRESONANCE_H
8:
9: // C++
10: #include <complex>
11: #include <map>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/M4Vec.h"
16: #include "Granitti/MHELMMatrix.h"
17: #include "Granitti/MParticle.h"
18:
19: namespace gra {
20:
21: // Resonance parameters
22: class PARAM_RES {
23: public:
24: PARAM_RES() {}
25: void PrintParam(double sqrts) const {
26: std::cout << rangi::style::bold << "Custom resonance parameters:" << rangi::style::reset
27: << std::endl;
28: }
29:
30: print{"~ PDG ID: " <v n", p.pdg};
31: print{"~ Mass M0: " <v 0.5F GeV n", p.mass};
32: print{"~ Width W: " <v 0.5F GeV n", p.width};
33: print{"~ Spin Jx2: " <v n", p.spinX2};
34: print{"~ Parity P: " <v n", p.P};
35: std::cout << std::endl;
36:
37: print{"~ Production W"};
38: print{"~ Ineffective vertex constant g_I: " <v 0.1E x exp(1 x 0.1F) n", std:riabs(g), std:riarg(g)};
39: print{"~ Form factor parameter: " <v 0.2F n", g_FF};
40:
41: std::cout << std::endl;
42: print{"~ Decay W"};
43: print{"~ BR: " <v 0.3En", hel.BR};
44: print{"~ Ineffective vertex constant g_F: " <v 0.3E", hel.g_decay};
45: std::cout << std::endl << std::endl;
46:
47: if (p.spinX2 != 0) {
48: std::cout << " ~ Polarization Lorentz frame " << FRAME << std::endl;
49: std::cout << std::endl;
50: std::cout << " ~ Spin polarization density matrix [rho]" << std::endl << std::endl;
51:
52: // Print elements
53: for (std::size_t i = 0; i < rho.size_row(); ++i) {
54: for (std::size_t j = 0; j < rho.size_col(); ++j) {
55: std::string delim = (j < rho.size_col() - 1) ? ", " : "";
56: print{"~> Re[rho[" <v i, j, ", std:riabs(rho[i][j]), std:riarg(rho[i][j]), delim, c_str(0)});
57: }
58: std::cout << std::endl;
59: }
60: std::cout << std::endl << std::endl;
61: // print{"Density matrix von Neumann entropy S = 0.2F
62: // "\n", MSpin::VonNeumannEntropy(rho)};
63: }
64: }
65:
66: // Particle class
67: MParticle p;
68:
69: // -----
70: // Tensor Pomeron couplings
71: std::vector<double> g_Tensor; // Production couplings
72: // -----
73:
74: // (Complex) production coupling constant
75: std::complex<double> g = 0.0;
76:
77: // Form factor parameter
78: double g_FF = 0.0;
79:
80: // Breit-Wigner type
81: int BW = 0;
82:
83: // 2->1 (generation spin correlation), 1->2 (decay spin correlations)

```

```


```

```

./include/Granitti/MResonance.h      2/2
84: bool SPINGEN = true;
85: bool SPINDEC = true;
86:
87: // Lorentz frame of the spin distribution
88: std::string FRAME = "null";
89:
90: // Maximum sliding pomeron helicity
91: int JMAX = 0;
92:
93: // Spin-density matrix (constant)
94: Matrix<std::complex<double>> rho;
95:
96: // -----
97: // Helicity and decay amplitude information
98: HELMatrix hel;
99: // -----
100: };
101:
102: } // namespace gra
103:
104: #endif

```

```

./include/Granitti/MTimer.h          1/1
1: // Timer class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MTIMER_H
7: #define MTIMER_H
8:
9: // C++
10: #include <chrono>
11:
12: namespace gra {
13: class MTimer {
14: public:
15: // Declare explicit, means here no conversion to bool allowed
16: explicit MTimer(bool reset = true) {
17: if (reset) { Reset(); }
18: }
19: // << operator
20: template <typename T, typename Traits>
21: friend std::basic_ostream<T, Traits> &operator<<(std::basic_ostream<T, Traits> &out,
22: const MTimer & timer) {
23: return out << timer.Elapsed().count();
24: }
25: // Time in sec
26: double ElapsedSec() const { return Elapsed().count() / 1000.0; }
27: // Time in msec
28: std::chrono::milliseconds Elapsed() const {
29: return std::chrono::duration_cast<std::chrono::milliseconds>(
30: std::chrono::high_resolution_clock::now() - start);
31: }
32: // Reset timer
33: void Reset() { start = std::chrono::high_resolution_clock::now(); }
34:
35: private:
36: std::chrono::high_resolution_clock::time_point start;
37: };
38:
39: } // namespace gra
40:
41: #endif

```

```

./include/Granitti/MGlobals.h        1/1
1: // GRANITTI Monte Carlo all global variables collected here
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MGLOBAL_H
7: #define MGLOBAL_H
8:
9: // Own
10: #include "Granitti/MStdakov.h"
11:
12: // LHAFDP
13: #include "LHAFDP/LHAFDP.h"
14:
15: namespace gra {
16: // =====
17: // These variables are initialized by MGranitti.cc
18: // =====
19:
20: // Model tune
21: extern std::string MODELPARAM;
22:
23: // Multithreading lock
24: extern std::mutex g_mutex;
25:
26: // For multithreaded VEGAS, to handle the exceptions from forked threads
27: extern std::exception_ptr g_globalExcept_ptr;
28:
29: // =====
30:
31: } // namespace gra
32:
33: #endif

```

```

./include/Granitti/MForm.h          1/2
1: // Form factors, structure functions, Regge trajectories etc. parametrizations
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MFORM_H
7: #define MFORM_H
8:
9: // C++
10: #include <complex>
11: #include <map>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/MVee.h"
16: #include "Granitti/MWw.h"
17: #include "Granitti/MKinematics.h"
18: #include "Granitti/MParticle.h"
19: #include "Granitti/MRandom.h"
20: #include "Granitti/MResonance.h"
21:
22: // Libraries
23: #include "json.hpp"
24:
25: namespace gra {
26:
27: // Model parameters
28: namespace PARAM_SOFT {
29: // Pomeron trajectory
30: extern double DELTA_P;
31: extern double ALPHA_P;
32:
33: // Couplings
34: extern double gN_P;
35: extern double gL_O;
36:
37: extern double gP_P;
38: extern double gamma;
39:
40: // Proton form factor
41: extern double fC1;
42: extern double fC2;
43: extern double fC3;
44:
45: extern bool COOPERON_ON;
46:
47: void PrintParam();
48: std::string GetHashStrin();
49:
50: void ReadParameters(const std::string& modelfile);
51: extern bool initialized;
52: } // namespace PARAM_SOFT
53:
54: namespace PARAM_STRUCTURE {
55: extern std::string F2;
56: extern std::string EM;
57: extern std::string QED_alpha;
58:
59: void ReadParameters(const std::string& modelfile);
60: extern bool initialized;
61: } // namespace PARAM_STRUCTURE
62:
63: // Flat (SERBUS) amplitude parameters
64: namespace PARAM_FLAT {
65: extern double b;
66:
67: void ReadParameters(const std::string& modelfile);
68: extern bool initialized;
69: } // namespace PARAM_FLAT
70:
71: // Forward proton excitation
72: namespace PARAM_NSTAR {
73: extern std::string Fragment;
74: extern std::vector<double> rcj;
75:
76: void ReadParameters(const std::string& modelfile);
77: extern bool initialized;
78: } // namespace PARAM_NSTAR
79:
80: namespace form {
81:
82: // Read resonance
83: gra::PARAM_RES ReadResonance(const std::string& resparam_str, MRandom& rng)

```



```

./include/Granitti/MForm.h          2/2
84:
85: // Regge signature
86: std::complex<double> ReggeEta(double alpha_t, double sigma);
87: std::complex<double> ReggeEtaLinear(double t, double alpha_t0, double ap, double sigma);
88:
89: // Proton form factor, elastic and inelastic
90: double SF(double t);
91: double SFINEI(double t, double M2);
92:
93: // Proton structure functions
94: double F2Q2(double x1, double Q2);
95: double F1Q2(double x1, double Q2);
96:
97: // Pomeron trajectory
98: double S3FomAlpha(double t);
99:
100: // Pion loop insert
101: double S3PI(double tau, double t);
102:
103: // t-integrated collinear EPA flux
104: double DDFlux(double x);
105:
106: // Coherent gamma flux
107: double F1(double Q2);
108: double F2(double Q2);
109: double G_M(double Q2);
110: double G_E(double Q2);
111:
112: double G_M_DIPOLE(double Q2);
113: double G_E_DIPOLE(double Q2);
114: double G_M_RELLY(double Q2);
115: double G_E_RELLY(double Q2);
116:
117: double CobFlux(double x, double t, double pt);
118: double IncohFlux(double x, double t, double pt, double M2);
119:
120:
121: double mu_ratio();
122:
123: // Breit-Wigner functions
124: double deltaWksee(double shat, double M0, double Gamma);
125: double deltaWkmp(double shat, double M0, double Gamma);
126:
127: std::complex<double> CBW(const gra:LORENTZSCALAR& lra, const gra:PARAM_RES& resonance);
128: std::complex<double> CBW_F1(double m2, double M0, double Gamma);
129: std::complex<double> CBW_F2(double m2, double M0, double Gamma);
130: std::complex<double> CBW_B1(double m2, double M0, double Gamma, int l, double mA, double mB);
131: std::complex<double> CBW_B2(double m2, double M0, double Gamma, double s);
132:
133: } // namespace form
134: } // namespace gra
135:
136: #endif

```

```

./include/Granitti/MH2.h           2/2
84: void SetAutoBufferSize(int n) { AUTOBUFSIZE = n; }
85: void FlushBuffer();
86:
87: // Symmetric bounds
88: void SetAutoSymmetry(const std::vector<bool>& tin) {
89:     if (tin.size() != 2)
90:         throw std::invalid_argument("MH2::SetAutoSymmetry: Input should be size 2 boolean vector");
91:     AUTOSYMMETRY = tin;
92: }
93:
94:
95: void GetBounds(int xbins, double xmin, double xmax, int ybins, double ymin,
96:                double ymax) const {
97:     xbins = XMIN;
98:     xmin = XMIN;
99:     xmax = XMAX;
100:    ybins = YMIN;
101:    ymin = YMIN;
102:    ymax = YMAX;
103: }
104:
105: void FuseBuffer(const MH2 &rhs) {
106:    buff_values.insert(buff_values.end(), rhs.buff_values.begin(), rhs.buff_values.end());
107:    buff_weights.insert(buff_weights.end(), rhs.buff_weights.begin(), rhs.buff_weights.end());
108: }
109:
110: // Keep it public for buffer fusion
111: std::vector<std::vector<double>> buff_values;
112: std::vector<double> buff_weights;
113:
114: private:
115:     std::string name; // Histogram name
116:
117:     //-----
118:     // For autorange
119:     bool FILLBUFF = false;
120:     int AUTOBUFSIZE = 10000; // Default AUTOBUFSIZE
121:     std::vector<bool> AUTOSYMMETRY = {false, false};
122:     //-----
123:
124:     // Boundary conditions
125:     double XMIN = 0.0;
126:     double XMAX = 0.0;
127:     int XBINS = 0;
128:
129:     double YMIN = 0.0;
130:     double YMAX = 0.0;
131:     int YBINS = 0;
132:
133:     // Number of underflow and overflow counts
134:     long long int fills = 0;
135:     std::vector<long long int> overflow = {0, 0};
136:     std::vector<long long int> underflow = {0, 0};
137:     long long int nanflow = 0;
138:
139:     // Logarithmic binning
140:     bool LOGX = false;
141:     bool LOGY = false;
142:     bool ValidBin(int xbin, int ybin) const;
143:     int GetBin(double value, double minval, double maxval, int nbins, bool logbins) const;
144:
145:     // Weights (in unweighted case weights = counts)
146:     MMatrix<double> weights;
147:     MMatrix<double> weights2;
148:
149:     // Counts
150:     MMatrix<long long int> counts;
151:
152: } // namespace gra
153:
154: } // namespace gra
155:
156: #endif

```

```

./include/Granitti/MH2.h           1/2
1: // 2D histogram class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MH2_H
7: #define MH2_H
8:
9: // C++
10: #include <complex>
11: #include <vector>
12:
13: // Own
14: #include "Granitti/MMatrix.h"
15:
16: namespace gra {
17: class MH2 {
18: public:
19:     MH2(int xbins, double xmin, double xmax, int ybins, double ymin, double ymax,
20:         std::string namestr = "noname");
21:
22:     MH2(int xbins, int ybins, std::string namestr = "noname");
23:     MH2();
24:     ~MH2();
25:
26:     void ResetBounds(int xbins, double xmin, double xmax, int ybins, double ymin, double ymax);
27:
28:     // Get full histogram data
29:     MMatrix<double> GetWeights() const { return weights; }
30:     MMatrix<double> GetWeights2() const { return weights2; }
31:     MMatrix<long long int> GetCounts() const { return counts; }
32:
33:     void Fill(double xvalue, double yvalue);
34:     void Fill2(double xvalue, double yvalue, double weight);
35:     void Clear();
36:     std::pair<double, double> WeightMeanAndError() const;
37:
38:     double GetMeanX(int power) const;
39:     double GetMeanY(int power) const;
40:
41:     double SumWeights() const;
42:     double SumWeights2() const;
43:     long long int SumBinCounts() const;
44:     long long int FillCount() const { return fills; }
45:
46:     double GetBinWeight(int xbin, int ybin) const;
47:     long long int GetBinCount(int xbin, int ybin) const;
48:     double GetBinWeight2(int xbin, double yvalue, double ymax, int kxbin, int kybin) const;
49:     double GetMaxWeight() const;
50:     double GetMinWeight() const;
51:     void Print() const;
52:     double ShannonEntropy() const;
53:
54:     // Set logarithmic binning
55:     // User needs to take care that XMIN and XMAX > 0
56:     void SetLogX() { LOGX = true; }
57:     void SetLogY() { LOGY = true; }
58:     void SetLogXY() {
59:         LOGX = true;
60:         LOGY = true;
61:     }
62:
63:     // Overload + operator to add two histograms
64:     MH2 operator+(const MH2 &rhs) {
65:         if ((this->XBINS != rhs.XBINS) || (this->YBINS != rhs.YBINS)) {
66:             throw std::domain_error("MH2 + operator: Histograms with different number of bins");
67:         }
68:         MH2 h(this->XBINS, this->XMIN, this->XMAX, this->YBINS, this->YMIN, this->YMAX, this->name);
69:
70:         h.fills = this->fills + rhs.fills;
71:         h.underflow = (this->underflow[0] + rhs.underflow[0], this->underflow[1] + rhs.underflow[1]);
72:         h.overflow = (this->overflow[0] + rhs.overflow[0], this->overflow[1] + rhs.overflow[1]);
73:         h.nanflow = this->nanflow + rhs.nanflow;
74:
75:         // DATA
76:         h.weights = this->weights + rhs.weights;
77:         h.weights2 = this->weights2 + rhs.weights2;
78:         h.counts = this->counts + rhs.counts;
79:
80:         return h;
81:     }
82:
83:     // AUTOBUFSIZE for autorange buffer size

```

```

./include/Granitti/MGamma.h       1/1
1: // Gamma amplitudes
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MGAMMA_H
7: #define MGAMMA_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/MHVec.h"
16: #include "Granitti/Mplitudes.h"
17: #include "Granitti/MKinematics.h"
18:
19: namespace gra {
20:
21: // Monopole wave functions and parameters
22: namespace PARAM_MONOPOLE {
23:
24:     extern bool printed; // Printing called
25:     extern bool initialized; // For lazy initialization
26:
27:     extern int Enj // Bound state energy level
28:     extern double M0; // Monopole mass
29:     extern double Gamma; // Monopoleium width
30:     extern std::string coupling; // Coupling scenarios
31:     extern int gn; // Monopole charge n = 1,2,3,...
32:
33:     double EnergyMP(double n);
34:     double GammaMP(double n, double alpha_g);
35:     double F1MP(double n);
36:
37:     void PrintParameters(double sqrts);
38:     void ReadParameters(const std::string &modelFile);
39:
40: } // namespace PARAM_MONOPOLE
41:
42:
43: class MGamma : public MAmplitudes {
44: public:
45:     MGamma(gra:LORENTZSCALAR &its, const std::string &modelFile);
46:     ~MGamma() {}
47:
48:     // yy->resonance X
49:     double yyX(gra:LORENTZSCALAR &its, gra:PARAM_RES &resonance) const;
50:
51:     // yy->lepton pair, or monopole antimonopole amplitude
52:     double yyfbar(gra:LORENTZSCALAR &its);
53:
54:     // yy->SM Higgs
55:     double yyHiggs(gra:LORENTZSCALAR &its) const;
56:
57:     // yy->monopoleium
58:     double yyMP(gra:LORENTZSCALAR &its) const;
59:
60: protected:
61:
62:
63: } // namespace gra
64:
65: #endif

```

./include/Granitti/MSudakov.h

1/3

```
1: // Shuvaev PDF and Sudakov suppression class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MSUDAKOV_H
7: #define MSUDAKOV_H
8:
9: // C++
10: #include <complex>
11: #include <memory>
12: #include <vector>
13:
14: // LHAPOF
15: #include "LHAPOF/LHAPOF.h"
16:
17: // Own
18: #include "Granitti/Muon.h"
19: #include "Granitti/MSuon.h"
20: #include "Granitti/MTimer.h"
21:
22: // Libraries
23: #include "json.hpp"
24:
25: namespace gra {
26:
27: struct MSudakovNumerics {
28:     double q2_MIN = 0.0;
29:     double q2_MAX = 0.0;
30:     double M_MIN = 0.0;
31:     double M_MAX = 0.0;
32:     double x_MIN = 0.0;
33:     const double x_MAX = 1.0 - 1E-9;
34:
35:     // Numerical integral discretization [SET HERE]
36:     const unsigned int SudakovIntegralN = 1000;
37:     const unsigned int ShuvaevIntegralN = 400;
38:
39:     // Logarithmic stepping true/false
40:     std::vector<bool> SHUV_log_ON;
41:     std::vector<bool> SHUV_log_OFF;
42:
43:     // Number of discrete intervals
44:     std::vector<unsigned int> SUDA_N;
45:     std::vector<unsigned int> SHUV_N;
46:
47:     // CMS energy
48:     double sqrts = 0.0;
49:
50:     bool DEBUG = false;
51:
52:     void ReadParameters() {
53:         // Read and parse
54:         using json = nlohmann::json;
55:         const std::string inputfile = gra::aux::GetBasePath(2) + "/moddata/" + "NUMERICS.json";
56:         const std::string data = gra::aux::GetInputData(inputfile);
57:         json j;
58:
59:         try {
60:             j = json::parse(data);
61:
62:             // JSON block identifiers
63:             const std::string MID = "NUMERICS_SUDAKOV";
64:
65:             q2_MIN = j.at(XID).at("q2_MIN");
66:             q2_MAX = j.at(XID).at("q2_MAX");
67:             M_MIN = j.at(XID).at("M_MIN");
68:             M_MAX = j.at(XID).at("M_MAX"); // Post-Setup
69:             x_MIN = j.at(XID).at("x_MIN"); // Post-Setup
70:             x_MAX = j.at(XID).at("x_MAX"); // Set above
71:
72:             // Logarithmic stepping true/false (assign needed for <cast>)
73:             SUDA_log_ON = xx = j.at(XID).at("SUDA").at("log_ON");
74:             SUDA_log_OFF = xxx;
75:             SHUV_log_ON = yy = j.at(XID).at("SHUV").at("log_ON");
76:             SHUV_log_OFF = yyy;
77:
78:             // Number of node points (assign needed for <cast>)
79:             std::vector<unsigned int> nn = j.at(XID).at("SUDA").at("N");
80:             SUDA_N = nn;
81:             std::vector<unsigned int> mm = j.at(XID).at("SHUV").at("N");
82:             SHUV_N = mm;
83:
84:         }
85:     }
86:
87:     // Discretization
88:     // ShuvaevIntegralN = j.at(XID).at("ShuvaevIntegralN");
89:     // SudakovIntegralN = j.at(XID).at("SudakovIntegralN");
90:     // DEBUG = j.at(XID).at("DEBUG");
91:     catch (...) {
92:         std::string str = "MSudakovNumerics: ReadParameters: Error parsing " + inputfile +
93:             " (Check for extra/missing commas)";
94:         throw std::invalid_argument(str);
95:     }
96: };
97:
98: // Interpolation container
99: class TArray2D {
100: public:
101:     TArray2D(int) {}
102:     TArray2D(int, int) {}
103:
104:     std::string name[2];
105:     double MIN[2] = {0};
106:     double MAX[2] = {0};
107:     unsigned int N[2] = {0};
108:     double STEP[2] = {0};
109:     bool islog[2] = {false};
110:
111:     // Setup discretization
112:     void Set(unsigned int VAR, std::string_name, double_min, double_max, double_N,
113:             double_logarithmic) {
114:         // Out of index
115:         if (VAR > 1) {
116:             throw std::invalid_argument("TArray2D:Set: Error: VAR = " + std::ito_string(VAR) + " > 1");
117:         }
118:
119:         // At least two intervals
120:         if (N[VAR] < 2) {
121:             throw std::invalid_argument("TArray2D:Set: Error: N = " + std::ito_string(N[VAR]) + " < 2");
122:         }
123:
124:         // Sanity
125:         if (MIN[VAR] >= MAX[VAR]) {
126:             throw std::invalid_argument("TArray2D:Set: Error: Variable " + name[VAR] + " MIN = " +
127:                 std::ito_string(MIN[VAR]) + " >= MAX = " + std::ito_string(MAX[VAR]));
128:         }
129:
130:         // Sanity
131:         if (!logarithmic || MIN[VAR] <= 1e-9) {
132:             throw std::invalid_argument("TArray2D:Set: Error: Variable " + name[VAR] +
133:                 " is using logarithmic stepping with boundary MIN = " + std::ito_string(MIN[VAR]) +
134:                 " <= 1e-9");
135:         }
136:         islog[VAR] = _logarithmic;
137:
138:         // Logarithm taken here!
139:         MIN[VAR] = islog[VAR] ? std::log(MIN[VAR]) : MIN[VAR];
140:         MAX[VAR] = islog[VAR] ? std::log(MAX[VAR]) : MAX[VAR];
141:         N[VAR] = _N[VAR];
142:
143:         STEP[VAR] = (MAX[VAR] - MIN[VAR]) / N[VAR];
144:         name[VAR] = _name[VAR];
145:     }
146:
147:     // Call this last
148:     void InitArray() {
149:         // Note Wrt
150:         F = std::vector<std::vector<std::vector<double>>>(N[1] + 1, std::vector<double>(4, 0.0));
151:     }
152:
153:     // Wrt per dimension!
154:     std::vector<std::vector<std::vector<double>>>>;
155:     std::vector<std::vector<std::vector<double>>>>;
156:
157:     // CMS energy (needed for boundary conditions)
158:     double sqrts = 0.0;
159:
160:     std::string GetHashString() const {
161:         std::string str = std::ito_string(islog[0]) + std::ito_string(islog[1]) + std::ito_string(MIN[0]) +
162:             std::ito_string(MIN[1]) + std::ito_string(MAX[0]) + std::ito_string(MAX[1]) +
163:             std::ito_string(N[0]) + std::ito_string(N[1]) + std::ito_string(sqrts);
164:         return str;
165:     }
166: };
167:
168: }
```

./include/Granitti/MSudakov.h

3/3

```
167: bool WriteArray(const std::string filename, bool overwrite) const;
168: bool ReadArray(const std::string filename);
169: std::pair<double, double> Interpolate2D(double A, double B) const;
170: };
171:
172: // Sudakov suppression and skewed pdf
173: //
174: // Take care when copying this class - there is a pointer to LHAPOF
175: class MSudakov {
176: public:
177:     MSudakov();
178:     ~MSudakov();
179:
180:     void Init(double sqrts_in, const std::string &PDFSET, bool init_arrays = true);
181:     void InitLHAPOF(const std::string &PDFSET);
182:     std::pair<double, double> ShuvaevJ(double q2, double x);
183:     std::pair<double, double> SudakovJ(double q2, double M);
184:
185:     double fg_XQM(double x, double q2, double M) const;
186:     double Alpha2_Q2(double q2) const;
187:     double RunFlavor(double q2) const;
188:     double sq_Q2(double x, double q2) const;
189:     void TestPDF() const;
190:
191:     bool initialized = false;
192:
193: private:
194:     int init_trials = 0;
195:
196:     void InitArrays();
197:     double diff_sq_Q2_wrt_Q2(double x, double q2) const;
198:     double AP_sq(double delta, double q2) const;
199:     double AP_gg(double delta, double q2) const;
200:     void CalculateArray(TArray2D &arr, std::pair<double, double> (MSudakov::*f)(double, double));
201:
202:     std::string PDFSETNAME;
203:     LHAPOF::PDF *PDFset = nullptr;
204:
205:     TArray2D veto; // Sudakov veto
206:     TArray2D sqd; // Shuvaev pdf
207:
208:     MSudakovNumerics Numerics; // Numerics
209: };
210:
211: // namespace gra
212:
213: #endif
```

./include/Granitti/MSudakov.h

2/3

```
84: // Discretization
85: // ShuvaevIntegralN = j.at(XID).at("ShuvaevIntegralN");
86: // SudakovIntegralN = j.at(XID).at("SudakovIntegralN");
87: // DEBUG = j.at(XID).at("DEBUG");
88: catch (...) {
89:     std::string str = "MSudakovNumerics: ReadParameters: Error parsing " + inputfile +
90:         " (Check for extra/missing commas)";
91:     throw std::invalid_argument(str);
92: }
93:
94: };
95:
96:
97:
98: // Interpolation container
99: class Tarray2D {
100: public:
101:     Tarray2D(int) {}
102:     Tarray2D(int, int) {}
103:
104:     std::string name[2];
105:     double MIN[2] = {0};
106:     double MAX[2] = {0};
107:     unsigned int N[2] = {0};
108:     double STEP[2] = {0};
109:     bool islog[2] = {false};
110:
111:     // Setup discretization
112:     void Set(unsigned int VAR, std::string_name, double_min, double_max, double_N,
113:             double_logarithmic) {
114:         // Out of index
115:         if (VAR > 1) {
116:             throw std::invalid_argument("TArray2D:Set: Error: VAR = " + std::ito_string(VAR) + " > 1");
117:         }
118:
119:         // At least two intervals
120:         if (N[VAR] < 2) {
121:             throw std::invalid_argument("TArray2D:Set: Error: N = " + std::ito_string(N[VAR]) + " < 2");
122:         }
123:
124:         // Sanity
125:         if (MIN[VAR] >= MAX[VAR]) {
126:             throw std::invalid_argument("TArray2D:Set: Error: Variable " + name[VAR] + " MIN = " +
127:                 std::ito_string(MIN[VAR]) + " >= MAX = " + std::ito_string(MAX[VAR]));
128:         }
129:
130:         // Sanity
131:         if (!logarithmic || MIN[VAR] <= 1e-9) {
132:             throw std::invalid_argument("TArray2D:Set: Error: Variable " + name[VAR] +
133:                 " is using logarithmic stepping with boundary MIN = " + std::ito_string(MIN[VAR]) +
134:                 " <= 1e-9");
135:         }
136:         islog[VAR] = _logarithmic;
137:
138:         // Logarithm taken here!
139:         MIN[VAR] = islog[VAR] ? std::log(MIN[VAR]) : MIN[VAR];
140:         MAX[VAR] = islog[VAR] ? std::log(MAX[VAR]) : MAX[VAR];
141:         N[VAR] = _N[VAR];
142:
143:         STEP[VAR] = (MAX[VAR] - MIN[VAR]) / N[VAR];
144:         name[VAR] = _name[VAR];
145:     }
146:
147:     // Call this last
148:     void InitArray() {
149:         // Note Wrt
150:         F = std::vector<std::vector<std::vector<double>>>(N[1] + 1, std::vector<double>(4, 0.0));
151:     }
152:
153:     // Wrt per dimension!
154:     std::vector<std::vector<std::vector<double>>>>;
155:     std::vector<std::vector<std::vector<double>>>>;
156:
157:     // CMS energy (needed for boundary conditions)
158:     double sqrts = 0.0;
159:
160:     std::string GetHashString() const {
161:         std::string str = std::ito_string(islog[0]) + std::ito_string(islog[1]) + std::ito_string(MIN[0]) +
162:             std::ito_string(MIN[1]) + std::ito_string(MAX[0]) + std::ito_string(MAX[1]) +
163:             std::ito_string(N[0]) + std::ito_string(N[1]) + std::ito_string(sqrts);
164:         return str;
165:     }
166: };
167:
168: }
```

./include/Granitti/MParton.h

1/1

```
1: // Simple parton model 2->2 phase space class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MPARTON_H
7: #define MPARTON_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/Muon.h"
16: #include "Granitti/MSuon.h"
17: #include "Granitti/MMatrix.h"
18: #include "Granitti/MProcess.h"
19: #include "Granitti/MSuon.h"
20:
21: // HepMC3
22: #include "HepMC3/GenEvent.h"
23:
24: namespace gra {
25: class MParton : public MProcess {
26: public:
27:     MParton();
28:     MParton(std::string process, const std::vector<aux::OneCMD &syntax>);
29:     virtual ~MParton();
30:
31:     void post_Constructor();
32:
33:     double operator()(const std::vector<double> &randvec, AuxIntData &aux) {
34:         return EventWeight(randvec, aux);
35:     }
36:
37:     double EventWeight(const std::vector<double> &randvec, AuxIntData &aux);
38:     void PrintInfo(bool silent) const;
39:
40: private:
41:     void Initialize();
42:     bool LogProbabilities(const std::vector<double> &pip, const std::vector<double> &ppp);
43:     bool PhysicalCut() const;
44:
45:     // 2->2 dim phase space
46:     bool B2RandomKin(const std::vector<double> &randvec);
47:     bool B2BuildKin(double xb1, double xb2);
48:     bool B2RecordEvent(HepMC3::GenEvent &ev);
49:
50:     double B2IntegralVolume() const;
51:     double B2PhaseSpaceWeight() const;
52:
53:     void DecayWidthHPS(double &exact) const;
54: };
55:
56: } // namespace gra
57:
58: #endif
```

```

./include/Granitti/MRegge.h          1/2
1: // Regge Amplitudes
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MREGGE_H
7: #define MREGGE_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/M4Vec.h"
16: #include "Granitti/MRelativistics.h"
17:
18: namespace gra {
19:
20: // Regge amplitude parameters
21: namespace PARAM_REGGE {
22:
23: extern bool initialized;
24:
25: extern std::vector<double> a0;
26: extern std::vector<double> ap;
27: extern std::vector<double> sgn;
28:
29: extern double s0;
30:
31: extern int offshellFF;
32: extern double b_EXP;
33: extern double a_OREAR;
34: extern double b_OREAR;
35: extern double b_POW;
36:
37: extern bool reggeize;
38:
39: extern double omega;
40:
41: // Meson/Baryon couplings and Pomeron, Reggeon, Reggeon exchanges
42: extern std::vector<double> c; // coupling
43: extern std::vector<bool> n; // on/off
44:
45: void PrintParam();
46: void ReadParameters(int PDG, const std::string &modelFile);
47:
48: std::complex<double> JPC_CR_coupling(const gra::LORENTZSCALAR &its,
49:                                     const gra::PARAM_RES &resonance);
50:
51: // Amplitude form factors
52: double Proton_FF(double tprime, double b);
53: double Meson_FF(double that, double M2);
54: double Baryon_FF(double that, double M2);
55: double ResonanceFormFactor(double that, double M2, double S0);
56:
57: // Propagators
58: double Meson_prop(double that, double M2);
59: double Baryon_prop(double that, double M2);
60: std::complex<double> FEI_prop(double that, double M2);
61: // namespace PARAM_REGGE
62:
63: // Matrix element REGGE * GvV * <<- 2*external_legs - 8)
64: class MRegge {
65: public:
66:     MRegge(gra::LORENTZSCALAR &its, const std::string &modelFile);
67:     MRegge() {}
68:
69: // Regge amplitudes
70: std::complex<double> ME3(gra::LORENTZSCALAR &its) const;
71: std::complex<double> ME6(gra::LORENTZSCALAR &its) const;
72: std::complex<double> ME4(gra::LORENTZSCALAR &its, double sign) const;
73: std::complex<double> ME2(gra::LORENTZSCALAR &its, int mode) const;
74: std::complex<double> ME3(gra::LORENTZSCALAR &its, gra::PARAM_RES &resonance) const;
75: std::complex<double> ME3DD(gra::LORENTZSCALAR &its, gra::PARAM_RES &resonance) const;
76:
77: std::complex<double> ME3HEL(gra::LORENTZSCALAR &its, gra::PARAM_RES &resonance) const;
78: std::complex<double> PhotoME3(gra::LORENTZSCALAR &its, gra::PARAM_RES &resonance) const;
79:
80: void PomPomProtonVertex(const gra::LORENTZSCALAR &its, double &FF_A, double &FF_B) const;
81:
82: // Constructor amplitude leg permutations
83:

```

```

./include/Granitti/MRegge.h          2/2
84: void ConstructPerm(int type);
85:
86: // Regge propagators
87: std::complex<double> PropOnly(double s, double t) const;
88: std::complex<double> OdderonProp(double s, double t) const;
89: std::complex<double> PhotonProp(double s, double t, double s2, bool excite,
90:                                 double M2_forward) const;
91:
92: // Helicity functions
93: double g_VerTex(double t, double lambda_d, double lambda_f) const;
94: std::complex<double> g_H_VerTex(double t1, double t2, double sphi, int lambda_d, int P,
95:                                 int JMAA) const;
96: double gammaLambda(double t1, double t2, double s1, double s2) const;
97: int
98:   k3(int J, int P, int t_1, int sigma_1, int P_k, int sigma_k) const;
99:
100: private:
101: // Amplitude permutations
102: std::vector<std::vector<int>> permutations4;
103: std::vector<std::vector<int>> permutations;
104:
105: // namespace gra
106:
107: #endif

```

```

./include/Granitti/M4Vec.h          1/4
1: // 4-vectors (HEADER ONLY class)
2: // with standard metric (+,-,-,-) and MC initialization convention (px,py,pz,0)
3: //
4: // However, note that ^, operators index in "normal textbook convention",
5: // to be more streamlined with Lorentz index contractions of amplitudes
6: //
7: // p^0 = E
8: // p^1 = px
9: // p^2 = py
10: // p^3 = pz
11: //
12: // p_0 = E
13: // p_1 = -px
14: // p_2 = -py
15: // p_3 = -pz
16: //
17: // (c) 2017-2020 Mikael Mieskolainen
18: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
19:
20: #ifndef M4VEC_H
21: #define M4VEC_H
22:
23: #include <complex>
24: #include <iostream>
25: #include <vector>
26:
27: namespace gra {
28: // C++17 allows class template argument default deduction without brackets
29: //
30: // M4Vec s;
31: // M4Vec<double> b;
32: // M4Vec<double> c;
33: //
34: // a,b,c are all valid.
35: //
36: // template <typename T = double>
37: //
38: // Complex 4-vectors could be implemented this way, TBD.
39: //
40: //
41: class M4Vec {
42: public:
43: // All default to zero
44: M4Vec() { k = {0.0, 0.0, 0.0, 0.0}; }
45: // Initialize
46: M4Vec(double x, double y, double z, double t) { k = {t, x, y, z}; }
47: // Copy constructor
48: M4Vec(const M4Vec &rhs) { k = rhs.k; }
49:
50: // SET methods
51: void SetP3(double v) { k[X_] = v; }
52: void SetPy(double v) { k[Y_] = v; }
53: void SetPz(double v) { k[Z_] = v; }
54: void SetE(double v) { k[T_] = v; }
55:
56: void SetX(double v) { k[X_] = v; }
57: void SetY(double v) { k[Y_] = v; }
58: void SetZ(double v) { k[Z_] = v; }
59: void SetT(double v) { k[T_] = v; }
60:
61: void Set(double x, double y, double z, double t) {
62:     k[X_] = x;
63:     k[Y_] = y;
64:     k[Z_] = z;
65:     k[T_] = t;
66: }
67:
68: void SetP3PzPz(double x, double y, double z, double m) {
69:     k[X_] = x;
70:     k[Y_] = y;
71:     k[Z_] = z;
72:     k[R_] = msqrt(P3mod2() + m * m);
73: }
74:
75: void SetP3PyPz(double x, double y, double z) {
76:     k[X_] = x;
77:     k[Y_] = y;
78:     k[Z_] = z;
79: }
80:
81: void SetP3Y(double x, double y) {
82:     k[X_] = x;
83:

```

```

./include/Granitti/M4Vec.h          2/4
84:     k[Y_] = y;
85:
86: void SetPzE(double z, double e) {
87:     k[Z_] = z;
88:     k[T_] = e;
89: }
90:
91: void SetP3(const std::vector<double> &vec) {
92:     if (vec.size() != 3) { throw std::invalid_argument("M4Vec::SetP3: input should be 3-vector"); }
93:     k[X_] = vec[0];
94:     k[Y_] = vec[1];
95:     k[Z_] = vec[2];
96: }
97:
98: // Apply metric tensor: eta_{lmunu} k^lunu = k_lunu
99: void Flip3() {
100:     k[X_] = -k[X_];
101:     k[Y_] = -k[Y_];
102:     k[Z_] = -k[Z_];
103: }
104:
105: // GET methods
106: double Px() const { return k[X_]; }
107: double Py() const { return k[Y_]; }
108: double Pz() const { return k[Z_]; }
109: double E() const { return k[T_]; }
110:
111: double X() const { return k[X_]; }
112: double Y() const { return k[Y_]; }
113: double Z() const { return k[Z_]; }
114: double T() const { return k[T_]; }
115:
116: // Return 3-vector
117: std::vector<double> P3() const { return {k[X_], k[Y_], k[Z_]}; }
118:
119: // ALGEBRA methods
120:
121: // Space-time invariants
122: double Invariant0() const { return E() * E() - P3mod(); }
123: double M2() const { return Invariant(); }
124: double M() const { return (M2() > 0.0) ? msqrt(M2()) : -msqrt(-M2()); }
125:
126: // gamma = E/m = 1/sqrt(1-v^2/c^2) = 1/sqrt(1-beta^2)
127: double Gamma() const { return E() / M(); }
128: double Beta() const { return P3mod() / E(); }
129:
130: // Transverse 2-vector norm and norm^2
131: double Perp() const { return Pz(); }
132: double PerpP() const { return Pz(); }
133: double Pt() const { return msqrt(Pz()); }
134: double Pt2() const { return k[X_] * k[X_] + k[Y_] * k[Y_]; }
135:
136: // Total 3-vector norm and norm^2
137: double P3mod() const { return msqrt(P3mod()); }
138: double P3mod2() const { return k[X_] * k[X_] + k[Y_] * k[Y_] + k[Z_] * k[Z_]; }
139:
140: // Transverse mass (invariant under boost in z-direction)
141: double Mt() const { return msqrt(M2()); }
142: double Mt2() const { return M2(); }
143:
144: // Coincides with transverse mass for single particle
145: double Et() const { return M(); }
146: double Et2() const { return M2(); }
147:
148: // Angles
149: double Phi() const {
150:     return (Px() == 0.0 && Py() == 0.0) ? 0.0 : std::atan2(Py(), Px());
151: }
152: // ( / / / , , , range [0,PI]
153: double Theta() const {
154:     return (Px() == 0.0 && Pz() == 0.0 && Pt() == 0.0) ? 0.0 : std::atan2(Pt(), Pz());
155: }
156: // [PI, / 2
157: double CosTheta() const { return std::cos(Theta()); }
158:
159: // Pseudorapidity and rapidity (boost) in z-direction
160: double Eta() const { return 0.5 * std::log(P3mod() + Pz()) / (P3mod() - Pz()); }
161: double Rap() const { return 0.5 * std::log(E() + Pz()) / (E() - Pz()); }
162:
163: // Lightcone variables: k_+ = E + p_z
164: double LightconePos() const { return E() + Pz(); }
165: // Lightcone variables: k_- = E - p_z
166: double LightconeNeg() const { return E() - Pz(); }

```

```

./include/Granitti/M4Vec.h          3/4
167:
168: // SPINOR-HELICITY CO-VARIABLES
169:
170: std::complex<double> ComplexPz() const { return Px() + std::complex<double>(0, 1) * Py(); }
171: std::complex<double> ExpCPz() const {
172:     return ComplexPz() / sqrt(LightconePz() * LightconeQz());
173: }
174:
175: // 2-BODY ALGEBRA
176:
177: // Minkowski 4-product
178: double DotM(const M4Vec &rhs) const {
179:     return E() * rhs.E() - (Px() * rhs.Px() + Py() * rhs.Py() + Pz() * rhs.Pz());
180: }
181:
182: // 3-vector dot product
183: double Dot3(const M4Vec &rhs) const {
184:     return Px() * rhs.Px() + Py() * rhs.Py() + Pz() * rhs.Pz();
185: }
186:
187: // Transverse 2-vector dot product
188: double DeltaP(const M4Vec &rhs) const { return Px() * rhs.Px() + Py() * rhs.Py(); }
189:
190: // 3-vector cross product (return vector with 0 energy/time)
191: M4Vec Cross3(const M4Vec &rhs) const {
192:     M4Vec a(Py() * rhs.Pz() - Pz() * rhs.Py(), Pz() * rhs.Px() - Px() * rhs.Pz(),
193:           Px() * rhs.Py() - Py() * rhs.Px(), 0.0);
194:     return a;
195: }
196:
197: // Azimuth angle difference between [-Pz,Pz]
198: double DeltaPhi(const M4Vec &v) const {
199:     double D = Phi() - v.Phi();
200:     while (D >= PI) { D -= 2.0 * PI; }
201:     while (D < -PI) { D += 2.0 * PI; }
202:     return D;
203: }
204:
205: // Azimuth angle between [D,PI]
206: double DeltaPhiAbs(const M4Vec &v) const { return std::abs(DeltaPhi(v)); }
207:
208: // OPERATORS
209: double operator[](size_t mu) const { return k[mu]; } // for reading only
210: double operator|(size_t mu) const { return k[mu]; } // for substituting
211:
212: // Access operator in normal contravariant (upper index) indexing
213: double operator*(size_t mu) const { return k[mu]; }
214:
215: // Access operator with simultaneous lowering with metric (covariant index)
216: double operator|(size_t mu) const {
217:     if (mu == E_) {
218:         return k[E_];
219:     } else {
220:         return -k[mu];
221:     }
222: }
223:
224: #ifdef M4VEC_BOUNDS
225: #define throw std::out_of_range("M4Vec: operator %%% index out of bounds!");
226: #else
227: #define throw std::out_of_range("M4Vec: operator %%% index out of bounds!");
228: #endif
229:
230: // Minkowski scalar product
231: double operator*(const M4Vec &rhs) const { return DotM(rhs); }
232:
233: // 4-vector + 4-vector
234: M4Vec operator+(const M4Vec &rhs) const {
235:     return M4Vec(k[X_] + rhs.k[X_], k[Y_] + rhs.k[Y_], k[Z_] + rhs.k[Z_], k[R_] + rhs.k[R_]);
236: }
237:
238: M4Vec operator-(const M4Vec &rhs) const {
239:     return M4Vec(k[X_] - rhs.k[X_], k[Y_] - rhs.k[Y_], k[Z_] - rhs.k[Z_], k[R_] - rhs.k[R_]);
240: }
241:
242: // 4-vector * scalar
243: M4Vec operator*(const double rhs) const {
244:     return M4Vec(k[X_] * rhs, k[Y_] * rhs, k[Z_] * rhs, k[R_] * rhs);
245: }
246:
247: // Comparison
248: bool operator==(const M4Vec &rhs) const {
249:     const double EPS = 1e-10;

```

```

./include/Granitti/M4Mplitudes.h  1/1
1: // All external amplitudes (matrix elements) from MadGraph etc. collected here
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef M4MPLITUDES_H
7: #define M4MPLITUDES_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/M4Vec.h"
16: #include "Granitti/MHELematics.h"
17:
18: // MadGraph
19: #include "Granitti/Amplitude/AMP_MG3_gg_gg.h"
20: #include "Granitti/Amplitude/AMP_MG3_gg_ggg.h"
21: #include "Granitti/Amplitude/AMP_MG3_yy_1l.h"
22: #include "Granitti/Amplitude/AMP_MG3_yy_1l.h"
23: #include "Granitti/Amplitude/AMP_MG3_yy_wv.h"
24:
25: namespace gra {
26:
27:     // Matrix element dimension: * GaV* ** <- (2*external_legs - 8)
28:     class M4Mplitudes {
29:     public:
30:         M4Mplitudes() {}
31:         M4Mplitudes(int l) {}
32:         M4Mplitudes(int l) {}
33:
34:         AMP_MG3_yy_wv AmpMG3_yy_wv;
35:         AMP_MG3_yy_1l AmpMG3_yy_1l;
36:         AMP_MG3_yy_wvew AmpMG3_yy_wvew;
37:         AMP_MG3_gg_gg AmpMG3_gg_gg;
38:         AMP_MG3_gg_ggg AmpMG3_gg_ggg;
39:
40:     protected:
41:     };
42:
43:     // namespace gra
44: };
45: #endif

```

```

./include/Granitti/M4Vec.h          4/4
250:     return std::abs(k[R_] - rhs.k[R_]) < EPS && std::abs(k[X_] - rhs.k[X_]) < EPS &&
251:           std::abs(k[Y_] - rhs.k[Y_]) < EPS && std::abs(k[Z_] - rhs.k[Z_]) < EPS;
252: }
253:
254: bool operator!=(const M4Vec &rhs) const { return !(this == rhs); }
255:
256: // 4-vector +- 4-vector
257: void operator+=(const M4Vec &rhs) {
258:     k[R_] += rhs.k[R_];
259:     k[X_] += rhs.k[X_];
260:     k[Y_] += rhs.k[Y_];
261:     k[Z_] += rhs.k[Z_];
262: }
263:
264: void operator-=(const M4Vec &rhs) {
265:     k[R_] -= rhs.k[R_];
266:     k[X_] -= rhs.k[X_];
267:     k[Y_] -= rhs.k[Y_];
268:     k[Z_] -= rhs.k[Z_];
269: }
270:
271: // 4-vector */= scalar
272: void operator*=(const double rhs) {
273:     k[R_] *= rhs;
274:     k[X_] *= rhs;
275:     k[Y_] *= rhs;
276:     k[Z_] *= rhs;
277: }
278:
279: void operator/=(const double rhs) {
280:     k[R_] /= rhs;
281:     k[X_] /= rhs;
282:     k[Y_] /= rhs;
283:     k[Z_] /= rhs;
284: }
285:
286: void Print(const std::string name = "") const {
287:     std::cout << "M4Vec: " << name << " Pz (E): " << Pz() << ", Py (Y): " << Py()
288:           << ", Pa (Z): " << Pz() << ", E (E): " << E() << ", M (S): " << M()
289:           << ", theta: " << Theta() << ", phi: " << Phi() << std::endl;
290: }
291:
292: // Particle 4-position starting starting propagation from (0,0,0,0)
293: // - p is the particle 4-momentum in the lab frame
294: // - tau is the particle flight time in its rest frame
295: M4Vec PropagatePosition(double tau, double scale) const {
296:     const double gamma = Gamma();
297:     const double tau_ = gamma * tau;
298:     // Velocity
299:     const double beta = Beta();
300:     // End point 4-position
301:     return M4Vec(tau_ * c * beta * Pz() / P3mod() * scale, tau_ * c * beta * Py() / P3mod() * scale,
302:                 tau_ * c * beta * Px() / P3mod() * scale, tau_ * c * scale);
303: }
304:
305: private:
306:     static constexpr const double PI =
307:         3.141592653589793238462643383279502884197169399375105820974944L;
308:
309:     // speed of light, c = [m/s] (EXACT/DEFINITION)
310:     static constexpr const double c = 2.99792458E8;
311:
312: // Indices
313: static const int E_ = 0;
314: static const int X_ = 1;
315: static const int Y_ = 2;
316: static const int Z_ = 3;
317:
318: // Safe sqrt
319: double sqrt(double x) const { return std::sqrt(std::max(0.0, x)); }
320:
321: // 4-vector * vector
322: std::vector<double> k;
323:
324: // namespace gra
325:
326: #endif

```

```

./include/Granitti/MHELMatrix.h     1/1
1: // Kinematic objects and helicity coupling structures (HEADER ONLY file)
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MHELMATRIX_H
7: #define MHELMATRIX_H
8:
9: // C++
10: #include <complex>
11: #include <vector>
12:
13: #include "Granitti/M4Vec.h"
14:
15: namespace gra {
16:
17:     // Decay helicity amplitude information
18:     struct HELMatrix {
19:     void InitAlphaToZero() {
20:         const unsigned int N = 20;
21:         alpha_ = HELMatrix(std::complex<double>(N, N, 0.0));
22:         alpha_set_ = HELMatrix<bool>(N, N, true);
23:     }
24:
25:     // Decay amplitude matrix (calculated by SU(2) decomposition routines)
26:     gqa: HELMatrix<std::complex<double>> T;
27:
28:     // Helicity amplitude decay l-couplings (constant)
29:     gqa: HELMatrix<std::complex<double>> alpha;
30:     gqa: HELMatrix<bool> alpha_set;
31:
32:     // Parity conservation
33:     bool P_conservation = true;
34:
35:     // -----
36:     // Tensor Pomeron couplings
37:     std::vector<double> g_decay_tensor; // Decay couplings
38:
39:     // Branching Ratio to a particular decay
40:     double BR = 1.0; // Equivalent decay coupling
41:
42:     // namespace gra
43: };
44:
45: #endif

```

```

./include/Granitti/MH1.h 1/4
1: // Templated real valued or complex ID-histogram class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MH1_H
7: #define MH1_H
8:
9: // C++
10: #include <complex>
11: #include <vector>
12:
13: namespace grn {
14: template <class T>
15: class MH1 {
16: public:
17: MH1(int xbins, double xmin, double xmax, std::string namestr = "noname");
18: MH1(int xbins, std::string namestr = "noname");
19: MH1();
20: ~MH1();
21:
22: void Fill(double xvalue);
23: void Fill(double xvalue, T weight);
24: void Clear();
25:
26: void RawOutput() const;
27:
28: long long int FillCount() const { return fills; }
29: long long int SumBinCounts() const;
30: double GetMean(int power) const;
31: std::pair<double, double> WeightMeanAndError() const;
32:
33: // Probability density
34: std::vector<double> GetProbDensity() const {
35: double sum = 0;
36: for (std::size_t i = 0; i < weights.size(); ++i) { sum += GetPositiveDefinite(i); }
37: std::vector<double> p(weights.size(), 0.0);
38: if (sum > 0)
39: for (std::size_t i = 0; i < weights.size(); ++i) { p[i] = GetPositiveDefinite(i) / sum; }
40: }
41: return p;
42: }
43:
44: // Get full histogram data
45: std::vector<T> GetWeights() const { return weights; }
46: std::vector<T> GetWeights2() const { return weights2; }
47: std::vector<long long int> GetCounts() const { return counts; }
48:
49: void GetXPositiveDefinite(std::valarray<double> &x, std::valarray<double> &y) const;
50:
51: T SumWeights() const;
52: T SumWeights2() const;
53:
54: T GetBinWeight(int idx) const;
55: T GetBinWeights2(int idx) const;
56: double GetBinError(int idx) const;
57:
58: double GetMaxWeight() const;
59: double GetMinWeight() const;
60:
61:
62: long long int GetBinCount(int idx) const;
63: double GetPositiveDefinite(int i) const;
64:
65: double GetBinNval(int idx, int boundary = 0) const;
66: void GetBinInfo(double xvalue, int side);
67: void Print(double width = 1.25) const; // default argument
68:
69: void ResetBounds(int xbins);
70: void ResetBounds(int xbins, double xmin, double xmax);
71:
72: // Set logarithmic binning
73: void SetLogX() {
74: if (OMIN < 1e-9) {
75: throw std::invalid_argument(
76: "MH1::SetLogX: Error: Minimum boundary XMIN = " + std::to_string(OMIN) + " < 0");
77: }
78: LOGX = true;
79: }
80:
81: // Number of bins
82: std::size_t XBins() const { return XBINs; }
83:

```

```

./include/Granitti/MH1.h 3/4
167:
168: MH1<T> h(this->XBINS, this->XMIN, this->XMAX, this->name);
169:
170: h.fills = this->fills;
171: h.underflow = this->underflow;
172: h.overflow = this->overflow;
173: h.nanflow = this->nanflow;
174:
175: // DATA
176: h.weights = this->weights;
177: h.weights2 = this->weights2;
178: h.counts = this->counts;
179: for (std::size_t i = 0; i < h.weights.size(); ++i) {
180: h.weights[i] = (std::abs(h.weights[i]) > 0 ? h.weights[i] / h.weights[i] : 0;
181: h.weights2[i] = (std::abs(h.weights2[i]) > 0 ? h.weights2[i] / h.weights2[i] : 0;
182: h.counts[i] = (h.counts[i] > 0 ? h.counts[i] / h.counts[i] : 0;
183: }
184: return h;
185: }
186:
187: // AUTOBUFSIZE for autorange buffer size
188: void SetAutoBufferSize(int n) { AUTOBUFSIZE = n; }
189: void FlushBuffer();
190:
191: // Symmetric bounds
192: void SetAutoSymmetry(bool in) { AUTOSYMMETRY = in; }
193:
194: // Get histogram bounds
195: void GetBounds(int &xbins, double &xmin, double &ymax) const {
196: xbins = XBINs;
197: xmin = XMIN;
198: xmax = XMAX;
199: }
200:
201: void FuseBuffer(const MH1<T> &rhs) {
202: buff_values.insert(buff_values.end(), rhs.buff_values.begin(), rhs.buff_values.end());
203: buff_weights.insert(buff_weights.end(), rhs.buff_weights.begin(), rhs.buff_weights.end());
204: }
205:
206: // Keep it public for buffer fusion
207: std::vector<double> buff_values;
208: std::vector<T> buff_weights;
209:
210: private:
211: std::string name; // Histogram name
212:
213: // For autorange
214:
215: bool FILLBUFF = false;
216: int AUTOBUFSIZE = 10000; // Default AUTOBUFSIZE
217: bool AUTOSYMMETRY = false;
218:
219: // Boundary conditions
220:
221: double XMIN = 0.0;
222: double XMAX = 0.0;
223: int XBINs = 0;
224:
225: // Number of fills, overflow and underflow counts
226: long long int fills = 0;
227: long long int overflow = 0;
228: long long int underflow = 0;
229: long long int nanflow = 0;
230:
231: // weights (in unweighted case weights = counts)
232: std::vector<T> weights;
233: std::vector<T> weights2;
234: std::vector<long long int> counts; // counts
235:
236: // Logarithmic binning
237: bool LOGX = false;
238:
239: bool ValidBin(int idx) const;
240: int GetIdx(double value, double minval, double maxval, int nbins, bool logbins) const;
241:
242: // Comparing do we have double or std::complex<double>
243: constexpr bool IsReal() const { return std::is_same<T, double>::value; }
244: constexpr bool IsComplex() const { return std::is_same<T, complex<double>>::value; }
245:
246:
247:
248: // Include the implementation
249: #include "MH1.tcc"

```

```

./include/Granitti/MH1.h 2/4
84: // Copy constructor (DEFAULT is fine)
85: MH1<T> MH1(const MH1<T> &obj) {}
86:
87: // Overload + operator to add two histograms
88: MH1<T> operator+(const MH1<T> &rhs) {
89: if (this->XBINS != rhs.XBINs) {
90: throw std::domain_error("MH1<T> + operator: Histograms with different number of bins");
91: }
92:
93: MH1<T> h(this->XBINS, this->XMIN, this->XMAX, this->name);
94:
95: h.fills = this->fills + rhs.fills;
96: h.underflow = this->underflow + rhs.underflow;
97: h.overflow = this->overflow + rhs.overflow;
98: h.nanflow = this->nanflow + rhs.nanflow;
99:
100: // DATA
101: h.weights = this->weights;
102: h.weights2 = this->weights2;
103: h.counts = this->counts;
104: for (std::size_t i = 0; i < h.weights.size(); ++i) {
105: h.weights[i] += rhs.weights[i];
106: h.weights2[i] += rhs.weights2[i];
107: h.counts[i] += rhs.counts[i];
108: }
109: return h;
110: }
111:
112: // Overload - operator to subtract two histograms
113: MH1<T> operator-(const MH1<T> &rhs) {
114: if (this->XBINS != rhs.XBINs) {
115: throw std::domain_error("MH1<T> - operator: Histograms with different number of bins");
116: }
117:
118: MH1<T> h(this->XBINS, this->XMIN, this->XMAX, this->name);
119:
120: h.fills = this->fills - rhs.fills;
121: h.underflow = this->underflow - rhs.underflow;
122: h.overflow = this->overflow - rhs.overflow;
123: h.nanflow = this->nanflow - rhs.nanflow;
124:
125: // DATA
126: h.weights = this->weights;
127: h.weights2 = this->weights2;
128: h.counts = this->counts;
129: for (std::size_t i = 0; i < h.weights.size(); ++i) {
130: h.weights[i] -= rhs.weights[i];
131: h.weights2[i] -= rhs.weights2[i];
132: h.counts[i] -= rhs.counts[i];
133: }
134: return h;
135: }
136:
137: // Overload * operator to multiply two histograms
138: MH1<T> operator*(const MH1<T> &rhs) {
139: if (this->XBINS != rhs.XBINs) {
140: throw std::domain_error("MH1<T> * operator: Histograms with different number of bins");
141: }
142:
143: MH1<T> h(this->XBINS, this->XMIN, this->XMAX, this->name);
144:
145: h.fills = this->fills;
146: h.underflow = this->underflow;
147: h.overflow = this->overflow;
148: h.nanflow = this->nanflow;
149:
150: // DATA
151: h.weights = this->weights;
152: h.weights2 = this->weights2;
153: h.counts = this->counts;
154: for (std::size_t i = 0; i < h.weights.size(); ++i) {
155: h.weights[i] *= rhs.weights[i];
156: h.weights2[i] *= rhs.weights2[i];
157: h.counts[i] *= rhs.counts[i];
158: }
159: return h;
160: }
161:
162: // Overload / operator to divide two histograms
163: MH1<T> operator/(const MH1<T> &rhs) {
164: if (this->XBINS != rhs.XBINs) {
165: throw std::domain_error("MH1<T> / operator: Histograms with different number of bins");
166: }

```

```

./include/Granitti/MH1.h 4/4
250:
251: // namespace grn
252:
253: #endif

```



```

./include/Granitti/MSubProc.h      3/8
167: PROC_3() : MProc("yy", "CON", ("Continuum 1+-, qgsar, W+W-, monopolepair", "kt-EPA")) {}
168: PROC_3() {}
169: virtual double Amp2(gra:LorentzScalar& lts) {
170:     InitGamma(lts);
171:     double amp = GammaGammaCON(lts);
172:     return lts.Regge->ME3(lts, lts);
173: }
174: };
175: class PROC_4 : public MProc {
176: public:
177:     PROC_4() : MProc("yy", "QED", ("Continuum 1+-, qgsar", "FULL QED")) {}
178:     PROC_4() {}
179:     virtual double Amp2(gra:LorentzScalar& lts) {
180:         InitTensor(lts);
181:         if (!AssertN(2, lts.decaytree.size())) {
182:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2-body final state);
183:         }
184:         return Tensor->ME4(lts);
185:     }
186: };
187:
188:
189:
190: // -----
191: // Inclusive processes
192: // -----
193:
194: class PROC_5 : public MProc {
195: public:
196:     PROC_5() : MProc("X", "EL", ("Elastic", "Eikonal Pomeron", "Use with screening loop on")) {}
197:     PROC_5() {}
198:     virtual double Amp2(gra:LorentzScalar& lts) {
199:         InitRegge(lts);
200:         return abs2(Regge->ME2(lts, l));
201:     }
202: };
203:
204: class PROC_6 : public MProc {
205: public:
206:     PROC_6() : MProc("X", "SD", ("Single Diffractive", "Triple Pomeron", "With TOY fragmentation")) {}
207:     PROC_6() {}
208:     virtual double Amp2(gra:LorentzScalar& lts) {
209:         InitRegge(lts);
210:         return abs2(Regge->ME2(lts, 2));
211:     }
212: };
213:
214: class PROC_7 : public MProc {
215: public:
216:     PROC_7() : MProc("X", "DD", ("Double Diffractive", "Triple Pomeron", "With TOY fragmentation")) {}
217:     PROC_7() {}
218:     virtual double Amp2(gra:LorentzScalar& lts) {
219:         InitRegge(lts);
220:         return abs2(Regge->ME2(lts, 3));
221:     }
222: };
223:
224: class PROC_8 : public MProc {
225: public:
226:     PROC_8() : MProc("X", "ND", ("Non-Diffractive", "M-cut soft Pomerons", "With TOY fragmentation")) {}
227:     PROC_8() {}
228:     virtual double Amp2(gra:LorentzScalar& lts) { return 1.0; }
229: };
230:
231: // -----
232: // Pomeron-Pomeron processes
233: // -----
234:
235: class PROC_9 : public MProc {
236: public:
237:     PROC_9() : MProc("PP", "RES", ("Parametric resonance", "Pomeron")) {}
238:     PROC_9() {}
239:     virtual double Amp2(gra:LorentzScalar& lts) {
240:         InitRegge(lts);
241:         std::complex<double> A = 0.0;
242:         // Coherent sum of Resonances (loop over)
243:         for (auto& x : lts.RESONANCES) {
244:             const int J = static_cast<int>(x.second.p.spinX2 / 2.0);
245:             // Gamma-Pomeron for vectors (could be Pomeron-Odderon too)
246:             if (J == 1 && x.second.p.P == -1) {
247:                 A += Regge->PhotoME3(lts, x.second);
248:             }
249:             // Pomeron-Pomeron, J = 0,1,2,... all ok

```

```

./include/Granitti/MSubProc.h      5/8
333:     Tensor->ME4(lts);
334:
335:     // Add to temp vector (init with zero)
336:     std::vector<std::complex<double>> tempsum(lts.hamp.size(), 0.0);
337:     for (const auto& i : aux::indices(lts.hamp)) { tempsum[i] += lts.hamp[i]; }
338:
339:     // 2. Evaluate resonance matrix elements -> helicity amplitudes to lts.hamp
340:     if (lts.RESONANCES.size() != 0) {
341:         // No loop over resonances inside ME3
342:         Tensor->ME3(lts);
343:         // Add to temp vector
344:         for (const auto& i : aux::indices(lts.hamp)) { tempsum[i] += lts.hamp[i]; }
345:     }
346:
347:     // Set helicity amplitudes for the screening loop
348:     lts.hamp = tempsum;
349:     // -----
350:     // Get total amplitude squared 1/4 |sum_u [A_u]|^2
351:     double amp2 = 0.0;
352:     for (const auto& i : aux::indices(lts.hamp)) { amp2 += gra.math::abs2(lts.hamp[i]); }
353:     amp2 /= 4; // initial state helicity average
354:
355:     return amp2;
356: }
357:
358: };
359:
360:
361: class PROC_15 : public MProc {
362: public:
363:     PROC_15() : MProc("PP", "CON", ("Hadron continuum 2/4/6-body", "Pomeron")) {}
364:     PROC_15() {}
365:     virtual double Amp2(gra:LorentzScalar& lts) {
366:         InitRegge(lts);
367:         std::complex<double> A = 0.0;
368:         if (!AssertN(2, lts.decaytree.size())) {
369:             A = Regge->ME4(lts, l);
370:         } else if (!AssertN(4, lts.decaytree.size())) {
371:             A = Regge->ME4(lts, l);
372:         } else if (!AssertN(6, lts.decaytree.size())) {
373:             A = Regge->ME8(lts);
374:         } else {
375:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2, 4 or 6-body final state);
376:         }
377:         return abs2(A);
378:     }
379: };
380: };
381:
382:
383: class PROC_16 : public MProc {
384: public:
385:     PROC_16() : MProc("PP", "CON", ("Hadron continuum 2-body with [t-u] amp.", "Pomeron")) {}
386:     PROC_16() {}
387:     virtual double Amp2(gra:LorentzScalar& lts) {
388:         InitRegge(lts);
389:         if (!AssertN(2, lts.decaytree.size())) {
390:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2-body final state);
391:         }
392:         std::complex<double> A = Regge->ME4(lts, -1);
393:         return abs2(A);
394:     }
395: };
396: };
397:
398:
399: class PROC_17 : public MProc {
400: public:
401:     PROC_17() : MProc("PP", "RES+CON", ("Hadron resonances + continuum 2-body", "Pomeron / yP")) {}
402:     PROC_17() {}
403:     virtual double Amp2(gra:LorentzScalar& lts) {
404:         InitRegge(lts);
405:         std::complex<double> A = 0.0;
406:         if (!AssertN(2, lts.decaytree.size())) {
407:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2-body final state);
408:         }
409:         // 1. Continuum matrix element
410:         A = Regge->ME4(lts, l);
411:         // 2. Coherent sum of Resonances (loop over)
412:         for (auto& x : lts.RESONANCES) {

```

```

./include/Granitti/MSubProc.h      4/8
250:     else {
251:         A += Regge->ME3(lts, x.second);
252:     }
253: }
254:
255: // -----
256: // Set for the screening loop
257: lts.hamp = A;
258: // -----
259:
260: return abs2(A);
261: }
262: };
263:
264: class PROC_10 : public MProc {
265: public:
266:     PROC_10()
267:     : MProc("PP", "RESHEL", ("Elastic pomeron helicity amplitudes", "Pomeron", "DEVELOPER ONLY PROCESS!")) {}
268:     PROC_10() {}
269:     virtual double Amp2(gra:LorentzScalar& lts) {
270:         InitRegge(lts);
271:         std::complex<double> A = 0.0;
272:         A = Regge->ME3HEL(lts, lts.RESONANCES.begin()->second);
273:         return abs2(A);
274:     }
275: };
276:
277:
278: class PROC_11 : public MProc {
279: public:
280:     PROC_11() : MProc("PP", "RESTENSOR", ("Parametric resonance", "Tensor Pomeron")) {}
281:     PROC_11() {}
282:     virtual double Amp2(gra:LorentzScalar& lts) {
283:         InitTensor(lts);
284:         return Tensor->ME3(lts);
285:     }
286: };
287:
288: class PROC_12 : public MProc {
289: public:
290:     PROC_12() : MProc("PP", "CONTENSOR", ("Hadron continuum 2-body", "Tensor Pomeron")) {}
291:     PROC_12() {}
292:     virtual double Amp2(gra:LorentzScalar& lts) {
293:         InitTensor(lts);
294:         if (!AssertN(2, lts.decaytree.size())) {
295:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2-body final state);
296:         }
297:         return Tensor->ME4(lts);
298:     }
299: };
300: };
301:
302: class PROC_13 : public MProc {
303: public:
304:     PROC_13()
305:     : MProc("PP", "CONTENSOR24", ("Hadron resonances + continuum 2-body > 4-body", "Tensor Pomeron", "DEVELOPER ONLY PROCESS!")) {}
306:     PROC_13() {}
307:     virtual double Amp2(gra:LorentzScalar& lts) {
308:         InitTensor(lts);
309:         if (!AssertN(2, lts.decaytree.size()) && !AssertN(4, lts.decaytree.size())) {
310:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2 (2s) or 4-body final state);
311:         }
312:         return Tensor->ME6(lts);
313:     }
314: };
315: };
316:
317: };
318:
319: class PROC_14 : public MProc {
320: public:
321:     PROC_14()
322:     : MProc("PP", "RES+CONTENSOR", ("Hadron resonances + continuum 2-body", "Tensor Pomeron / yP")) {}
323:     PROC_14() {}
324:     virtual double Amp2(gra:LorentzScalar& lts) {
325:         InitTensor(lts);
326:         if (!AssertN(2, lts.decaytree.size())) {
327:             throw std::invalid_argument(ISTATE + "[+ CHANNEL +]" requires 2-body final state);
328:         }
329:         // 1. Evaluate continuum matrix element -> helicity amplitudes to lts.hamp

```

```

./include/Granitti/MSubProc.h      6/8
416:     const int J = static_cast<int>(x.second.p.spinX2 / 2.0);
417:     // Gamma-Pomeron for vectors
418:     if (J == 1 && x.second.p.P == -1) {
419:         A += Regge->PhotoME3(lts, x.second);
420:     }
421:     // Pomeron-Pomeron, J = 0,1,2,... all ok
422:     else {
423:         A += Regge->ME3(lts, x.second);
424:     }
425: }
426:
427: // -----
428: // Set helicity amplitudes for the screening loop
429: lts.hamp = A;
430: // -----
431:
432: return abs2(A);
433: }
434: };
435:
436: // -----
437: // Gamma (Odderon)-Pomeron processes
438: // -----
439:
440: class PROC_18 : public MProc {
441: public:
442:     PROC_18() : MProc("yP", "RES", ("Photoproduced resonance", "kt-EPA x Pomeron")) {}
443:     PROC_18() {}
444:     virtual double Amp2(gra:LorentzScalar& lts) {
445:         InitRegge(lts);
446:         std::complex<double> A = 0.0;
447:         // Coherent sum of Resonances (loop over)
448:         for (auto& x : lts.RESONANCES) {
449:             const int J = static_cast<int>(x.second.p.spinX2 / 2.0);
450:             // Vectors only
451:             if (J == 1 && x.second.p.P == -1) {
452:                 A += Regge->PhotoME1(lts, lts.RESONANCES.begin()->second);
453:             }
454:             // amplitude squared
455:         }
456:         return abs2(A);
457:     }
458: };
459: };
460:
461: class PROC_19 : public MProc {
462: public:
463:     PROC_19() : MProc("yP", "RESTENSOR", ("Photoproduced resonance", "QED x Tensor Pomeron")) {}
464:     PROC_19() {}
465:     virtual double Amp2(gra:LorentzScalar& lts) {
466:         InitTensor(lts);
467:         return Tensor->ME3(lts);
468:     }
469: };
470: };
471:
472: class PROC_20 : public MProc {
473: public:
474:     PROC_20() : MProc("OP", "RES", ("Oddproduced resonance", "Odderon x Pomeron")) {}
475:     PROC_20() {}
476:     virtual double Amp2(gra:LorentzScalar& lts) {
477:         InitRegge(lts);
478:         std::complex<double> A = 0.0;
479:         // Coherent sum of Resonances (loop over)
480:         for (auto& x : lts.RESONANCES) {
481:             const int J = static_cast<int>(x.second.p.spinX2 / 2.0);
482:             // Vectors only
483:             if (J == 1 && x.second.p.P == -1) { A += Regge->ME3ODD(lts, x.second); }
484:         }
485:         return abs2(A); // amplitude squared
486:     }
487: };
488: };
489:
490: // -----
491: // Durham QCD processes
492: // -----
493:
494: class PROC_21 : public MProc {
495: public:
496:     PROC_21() : MProc("gg", "chic(0)", ("QCD resonance chic(0)", "Durham QCD")) {}
497:     PROC_21() {}
498:     virtual double Amp2(gra:LorentzScalar& lts) {

```

```
./include/Granitti/MSubProc.h 7/8
```

```
499:   InitDurham(lts);
500:   double amp2 = 0.0;
501:   amp2 = Durham->DurhamQCD(lts, CHANNEL);
502:   return amp2;
503: }
504: }
505:
506: class PROC_22 : public MFProc {
507: public:
508:   PROC_22()
509:   : MFProc("gg", "COM",
510:           {"QCD continuum gg, 2 pseudoscalar", "Durham QCD", "ONDER VALIDATION!"}) {}
511:   PROC_22(l) {"QCD continuum gg, 2 pseudoscalar", "Durham QCD", "ONDER VALIDATION!"} {}
512:   virtual double Amp2(gra::LORENTZSCALAR& lts) {
513:     InitDurham(lts);
514:
515:     double amp2 = 2.0;
516:     if (AssertN(121, 21, PDLlist(lts)) {
517:       amp2 = Durham->DurhamQCD(lts, "gg");
518:     } else if (AssertN(2, lts.decaytree.size())) {
519:       amp2 = Durham->DurhamQCD(lts, "HBBbar");
520:     } else {
521:       ThrowUnknownFinalState();
522:     }
523:     return amp2;
524: }
525: }
526:
527: class PROC_23 : public MFProc {
528: public:
529:   PROC_23()
530:   : MFProc("gg", "FLUX", {"Durham flux with |A|^2 = 1", "Durham QCD", "SYSTEM TEST PROCESS!"}) {}
531:   PROC_23(l) {}
532:   virtual double Amp2(gra::LORENTZSCALAR& lts) {
533:     InitDurham(lts);
534:
535:     return Durham->DurhamQCD(lts, "FLUX");
536: }
537: }
538:
539:
540: // -----
541: // Gamma-Gamma processes
542: // -----
543:
544: class PROC_24 : public MFProc {
545: public:
546:   PROC_24()
547:   : MFProc("yy_LUX", "COM", {"Continuum l+l, qqbar, WW-, monopolepair", "Collinear LUX-PDF"}) {}
548:   PROC_24(l) {}
549:   virtual double Amp2(gra::LORENTZSCALAR& lts) {
550:     InitGamma(lts);
551:     double amp2 = GammaGammaCOM(lts);
552:     return Flux::ApplyLUXfluxes(amp2, lts);
553: }
554: }
555:
556: class PROC_25 : public MFProc {
557: public:
558:   PROC_25()
559:   : MFProc("yy_DS", "COM",
560:           {"Continuum l+l, qqbar, WW-, monopolepair", "Collinear Drees-Zeppenfeld EPA"}) {}
561:   PROC_25(l) {"Continuum l+l, qqbar, WW-, monopolepair", "Collinear Drees-Zeppenfeld EPA"} {}
562:   virtual double Amp2(gra::LORENTZSCALAR& lts) {
563:     InitGamma(lts);
564:     double amp2 = GammaGammaCOM(lts);
565:     return Flux::ApplyDZfluxes(amp2, lts);
566: }
567: }
568:
569: class PROC_26 : public MFProc {
570: public:
571:   PROC_26()
572:   : MFProc("yy", "FLUX", {"kt-EPA flux with |A|^2 = 1", "kt-EPA", "SYSTEM TEST PROCESS!"}) {}
573:   PROC_26(l) {}
574:   virtual double Amp2(gra::LORENTZSCALAR& lts) {
575:     InitGamma(lts);
576:     double amp2 = 1.0;
577:     return Flux::ApplyktEPAfluxes(amp2, lts);
578: }
579: }
580:
581: class PROC_27 : public MFProc {
```

```
./include/Granitti/MFactorized.h 1/1
```

```
1: // Factorized 2->3 phase space class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MFACTORIZED_H
7: #define MFACTORIZED_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/M4Vec.h"
16: #include "Granitti/MBox.h"
17: #include "Granitti/MMatrix.h"
18: #include "Granitti/MFProcess.h"
19: #include "Granitti/MHepMC.h"
20:
21: // HepMC3
22: #include "HepMC3/GenEvent.h"
23:
24: namespace gra {
25: class MFactorized : public MFProcess {
26: public:
27:   MFactorized();
28:   MFactorized(std::string process, const std::vector<aux::OneCMD>& syntax);
29:   virtual ~MFactorized();
30:
31:   void post_Constructor();
32:
33:   double operator()(const std::vector<double>& randvec, AuxInData &aux) {
34:     return EventWeight(randvec, aux);
35: }
36:
37:   double EventWeight(const std::vector<double>& randvec, AuxInData &aux);
38:   bool EventReactor(HepMC3::GenEvent &evt);
39:   void FPrintInit(bool silent);
40: private:
41:   void Initialize();
42:   bool LoopKinematics(const std::vector<double>& k1p, const std::vector<double>& k2p);
43:   bool FiducialCut4() const;
44:
45:   // s+1-Dim phase space, 2->3
46:   bool B51RandomKin(const std::vector<double>& randvec);
47:   bool B51BuildKin(double p1, double p2, double phi1, double phi2, double yX, double m2X,
48:                   double m1, double m3);
49:   void B51RecordEvent(HepMC3::GenEvent &evt);
50:
51:   double B51IntegralVolume() const;
52:   double B51PhaseSpaceWeight() const;
53:
54:   void DecayWithPS(double &exact) const;
55:
56:   // Dynamic sampling boundaries based on resonance position and width
57:   double M_MIN = 0.0;
58:   double M_MAX = 0.0;
59: };
60:
61: // namespace gra
62: }
63: #endif
```

```
./include/Granitti/MSubProc.h 8/8
```

```
582: public:
583:   PROC_27()
584:   : MFProc(
585:     "yy_DS", "FLUX",
586:     {"T02 flux with |A|^2 = 1", "Collinear Drees-Zeppenfeld EPA", "SYSTEM TEST PROCESS!"}) {}
587:   PROC_27(l) {"T02 flux with |A|^2 = 1", "Collinear Drees-Zeppenfeld EPA", "SYSTEM TEST PROCESS!"} {}
588:   virtual double Amp2(gra::LORENTZSCALAR& lts) {
589:     InitGamma(lts);
590:     double amp2 = 1.0;
591:     return Flux::ApplyDZfluxes(amp2, lts);
592: }
593: }
594:
595: // Umbrella class
596: class MSubProc {
597: public:
598:   MSubProc(const std::vector<std::string>& istate, const std::string& mc);
599:   void Initialize(const std::string& istate, const std::string& channel);
600:
601:   MSubProc(l) {}
602:   MSubProc(f) {}
603:
604:   double GetBareAmplitude2(gra::LORENTZSCALAR& lts);
605:
606:   std::string ISTATE; // "gg", "yy", "gg" etc.
607:   std::string CHANNEL; // "COM", "RES" etc.
608:   unsigned int LIPSDIM = 0; // Lorentz Invariant Phase Space Dimension
609:
610:   // Process descriptions
611:   std::map<std::string, std::vector<std::string>> Processes;
612:
613:   bool ProcessExist(std::string process) const;
614:   std::vector<std::string> PrintProcesses() const;
615:   std::vector<std::string> GetProcessDescription(std::string process) const;
616:
617: private:
618:   void ConstructDescriptions(const std::string& istate, const std::string& mc);
619:   void CreateProcesses();
620:   void DeleteProcesses();
621:   void ActivateProcess();
622:
623:   std::vector<std::string> prj;
624:   std::vector<std::string> prj;
625: };
626:
627: // namespace gra
628: }
629: #endif
```

```
./include/Granitti/MKinematics.h 1/20
```

```
1: // Standard relativistic kinematics and related structures [HEADER ONLY file]
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MKINEMATICS_H
7: #define MKINEMATICS_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13: #include <valarray>
14: #include <vector>
15:
16: // Own
17: #include "Granitti/M4Vec.h"
18: #include "Granitti/MBox.h"
19: #include "Granitti/MBox.h"
20: #include "Granitti/MBox.h"
21: #include "Granitti/MBox.h"
22: #include "Granitti/MBox.h"
23: #include "Granitti/MBox.h"
24: #include "Granitti/MBox.h"
25: #include "Granitti/MBox.h"
26: #include "Granitti/MBox.h"
27:
28: namespace gra {
29: namespace kinematics {
30:
31: // Case m3 = m4
32: inline double SolveP3_3(double m3, double pt3, double pt4, double pz3, double pz4, double E3, double E4) {
33:   const double t2 = E3 * E3;
34:   const double t3 = pt3 * pt3;
35:   const double t4 = pt4 * pt4;
36:   const double t5 = pz3 * pz3;
37:   const double t6 = m3 * m3;
38:   const double t7 = std::sqrt(s);
39:   const double t8 = s * t2 * t2;
40:   const double t9 = std::pow(s, 3.0 / 2.0);
41:   const double t10 = t2 * t2;
42:   const double t11 = t3 * t3;
43:   const double t12 = t4 * t4;
44:   const double t13 = t5 * t5;
45:   const double t14 = s * s;
46:   const double t15 = t5 * t6 * t6;
47:   const double t16 = t3 * t3 * 2.0;
48:   const double t17 = t4 * t5 * 2.0;
49:   const double t18 = E3 * t6 * t7 * t7;
50:   const double t19 = E3 * t3 * t7 * t7;
51:   const double t20 = E3 * t4 * t7 * t7;
52:   const double t21 = E3 * t5 * t7 * t7;
53:   const double t22 = t8 * t10 + t11 * t12 + t13 + t14 + t15 + t16 + t17 + t18 + t19 + t20 + t21 -
54:     E3 * t9 + t2 * t3 + 2.0 * t2 * t3 - s * t4 * 2.0 - s * t5 * 2.0 - s * t6 * 4.0 -
55:     t2 * t3 * 2.0 - t2 * t4 * 2.0 - t2 * t5 * 2.0 - t3 * t4 * 2.0 - t2 * t6 * 4.0 -
56:     E3 * t2 * t7 * t7;
57:
58:   const double t23 = std::sqrt(t22);
59:   const double t0 = pz3 * (-1.0 / 2.0) - (E3 * t23 * (1.0 / 2.0) - t7 * t23 * (1.0 / 2.0) +
60:     pz4 * (t3 + (1.0 / 2.0) - t4 * (1.0 / 2.0)) /
61:     (s * t2 - t3 - E3 * t7 * 2.0));
62:
63:   return t0;
64: }
65:
66: //
67: // Old version for reference
68: inline double SolveP3_3(double m1, double m2, double pt1, double pt2, double
69:   pz, double pz, double
70:   s) {
71:
72:   using gra::math::msqrt;
73:   using gra::math::pow2;
74:   using gra::math::pow3;
75:   using gra::math::pow4;
76:   using gra::math::pow5;
77:
78:   const double sgt_s = msqrt(s);
79:   const double p12 =
80:     (-pow2(m1) * pow2(p2) * pz + pow2(m2) * pow2(p2) * pz -
81:     pow(p2) * pz - pow2(p2) * pow2(p1) * pz +
82:     pow(p2) * pow2(p1) * pz + pow2(m1) * pow2(p2) -
83:     pow(m2) * pow3(p2) + 2.0 * pow2(p2) * pow3(p2) +
```



```

./include/Granitti/MKinematics.h      6/20
416: double costheta, sintheta, phi;
417: FlatIsotropic(costheta, sintheta, phi, rng);
418:
419:
420: // Jacobian of spherical coordinates
421: const std::valarray<double> k = {pnorm * sintheta * std::cos(phi),
422:   pnorm * sintheta * std::sin(phi), pnorm * costheta};
423:
424: // Energies by on-shell condition
425: const std::valarray<double> e = {msqr(pow2(m1) + pow2(pnorm)), msqr(pow2(m2) + pow2(pnorm))};
426:
427: // Back-to-back
428: p1 = T1(k[0], k[1], k[2], e[0]);
429: p2 = T1(-k[0], -k[1], -k[2], e[1]);
430:
431:
432: // 2-Body isotropic phase space decay in the rest frame
433: // ----- m0 -----
434: // ----- m1 -----
435: // ----- m2 -----
436: // For floating point / efficiency reasons, the mother mass needs to be provided
437: // outside
438:
439: // ----- m1 -----
440: // ----- m2 -----
441:
442: template <typename T1, typename T2>
443: inline MCW NBodyPhaseSpace(const T1 smother, double M0, const std::vector<double> km,
444:   double M1, double M2, const std::vector<T1> sp, T2 kring) {
445:   // Two particles
446:   p.resize(2);
447:
448:   // First isotropic decay
449:   const double pnorm = DecayMomentum(M0, m[0], m[1]);
450:   Isotropic(pnorm, p[0], p[1], m[0], m[1], rng);
451:
452:   // Boost daughters to the lab frame
453:   const int sign = 1;
454:   for (const auto k1 : {0, 1}) { LorentzBoost(mother, M0, p[k1], sign); }
455:
456:   // return phase space weight
457:   const double weight = dPhi2(M0, pnorm);
458:   return MCW(weight, gra:math:pow2(weight), 1);
459: }
460:
461: // 3-Body isotropic (Dalitz) phase space decay in the rest frame;
462: // by factorizing the 3-body phase space into two 2-body
463: // ----- m0 ----- m[1] -----
464: // ----- m1 ----- m[2] -----
465: // ----- m2 ----- m[3] -----
466: // ----- m12 ----- m2 -----
467: // ----- m12 ----- m2 -----
468:
469: // dPhi_3(pM, p0, p1, p2) ~ dm_12^2 dphi_1_2 (pM, p0, p12) dphi_1_3 (p12, p1, p2)
470: // ----- m12 ----- m2 -----
471: // Returns: phase space weight for a valid fragmentation and -1.0 for a
472: // kinematically impossible.
473: // When unweight == true, this function returns also number of trials for proper
474: // MC error treatment
475: // (MCW)
476:
477: template <typename T1, typename T2>
478: inline MCW ThreeBodyPhaseSpace(const T1 smother, double M0, const std::vector<double> km,
479:   double M1, double M2, const std::vector<T1> sp, bool unweight, T2 kring) {
480:   // Three final states
481:   T1 p[2];
482:
483:   // Phase space boundaries [min,max]
484:   const std::vector<double> m12bound = {m[1] + m[2], M0 - m[0]};
485:   double w_max = DecayMomentum(M0, m[0], m12bound[0]) * DecayMomentum(m12bound[1], m[1], m[2]);
486:   if (unweight == false) { w_max = 0; }
487:
488:   // Acceptance-Rejection
489:   const unsigned int MAXTRIAL = 1e8;
490:   std::vector<double> pnorm = {0.0, 0.0};
491:   double m12 = 0;
492:   MCW
493:   do {
494:     m12 = rng.U(m12bound[0], m12bound[1]); // Flat mass (in GeV, not GeV^2)
495:     pnorm[0] = DecayMomentum(M0, m[0], m12);
496:     pnorm[1] = DecayMomentum(m12, m[1], m[2]);
497:
498:     const double w = (m12 / gra:math:PI) * dPhi2(M0, pnorm[0]) * dPhi2(m12, pnorm[1]);

```

```

./include/Granitti/MKinematics.h      8/20
582: MCW
583: double
584: do {
585:   // 1. Ordered random numbers
586:   std::for_each(randvec.begin(), randvec.end(), fillrandom);
587:   std::sort(randvec.begin(), randvec.end()); // Ascending
588:   // std::sort(randvec.begin(), randvec.end(), [](const double a, const double
589:   // b) { return a > b; }); // Descending
590:   // return a > b; // Descending
591:
592:   // 2. Calculate effective masses
593:   calcmass();
594:
595:   // 3. Calculate momentum norm values
596:   w = calcmomentum();
597:
598:   x.Push(w);
599:   if (x.GetN() > MAXTRIAL) { return MCW(-1, 0, 0); } // Impossible kinematics
600:   while (x < rng.U(0.0, w_max))
601:
602:   // Normalize the phase space integral
603:   const double volume = 1.0 / (2.0 * std::pow(2.0 * gra:math:PI, 2 * N - 3)) *
604:     std::pow(DELTA, N - 2) / gra:math:factorial(N - 2) / M0;
605:   x = x * volume; // operator overloaded
606:
607:   // Return phase space weight
608:   return x;
609: }
610:
611: // Flat N-body phase space by implementing a number of (N-2) recursive 2-body
612: // decays.
613: // Input: sm [mother 4-vector, vector of daughter masses, vector of final
614: // 4-vectors]
615:
616: // [REFERENCE: F. James, https://cds.cern.ch/record/23743/files/CERN-68-15.pdf]
617:
618: // ----- m_n -----
619: // ----- m_{n-1} -----
620: // ----- m_{n-2} -----
621: // ----- m_{n-3} -----
622: // ----- m_{n-4} -----
623: // ----- m_{n-5} -----
624: // ----- m_{n-6} -----
625: // ----- m_{n-7} -----
626: // ----- m_{n-8} -----
627: // ----- m_{n-9} -----
628: // ----- m_{n-10} -----
629:
630: // For faster alternatives, investigate:
631: // M.H. Block, Monte Carlo phase space evaluation, Comp. Phys. Commun.
632: // 69, 459 (1992)
633: // S. Platzer, RAMBO on diet, https://arxiv.org/abs/1308.2922
634:
635: // Return: weight for a valid fragmentation and -1.0 for a kinematically
636: // impossible
637: // When unweight == true, this function returns also number of trials for proper
638: // MC error treatment
639: // (MCW obj)
640:
641:
642: template <typename T1, typename T2>
643: inline MCW NBodyPhaseSpace(const T1 smother, double M0, const std::vector<double> km,
644:   double M1, double M2, const std::vector<T1> sp, bool unweight, T2 kring) {
645:   using gra:math:msqr;
646:   using gra:math:pow2;
647:
648:   const int N = m.size();
649:
650:   // Generate effective masses and decay momentum
651:   std::vector<double> M_eff(N, 0.0); // Effective masses
652:   std::vector<double> pnorm(N, 0.0); // Decay momentum weights
653:   const MCW
654:   if (x.GetN() < 0) { return x; } // Impossible kinematics
655:
656:   // Setup decay daughters size
657:   p.resize(0);
658:
659:   // Recursively go down from N-body to 1-particle
660:   T1 k = smother;
661:   for (std::size_t i = N - 1; i >= 1; --i) {
662:     double costheta, sintheta, phi;
663:     FlatIsotropic(costheta, sintheta, phi, rng);
664:     const double pt = pnorm[i] * sintheta;

```

```

./include/Granitti/MKinematics.h      7/20
499:   x.Push(w);
500:   if (x.GetN() > MAXTRIAL) { return MCW(-1, 0, 0); } // Impossible kinematics
501:   while (pnorm[0] + pnorm[1] < rng.U(0.0, w_max))
502:
503:   // p0 -> p0 + p12 in the pM r.f. and p12 -> p1 + p2 in the p12 r.f.
504:   Isotropic(pnorm[0], p[0], p12, m[0], m[2], rng);
505:   Isotropic(pnorm[1], p[1], p2, m[1], m[2], rng);
506:
507:   // Boost p[1] and p[2] out from the pM rest frame to the pM rest frame
508:   const int sign = 1; // Boost sign
509:   for (const auto k1 : {1, 2}) { LorentzBoost(p12, m2, p[k1], sign); }
510:
511:   // Boost p[0], p[1] and p[2] to the lab frame
512:   for (const auto k1 : {0, 1, 2}) { LorentzBoost(mother, M0, p[k1], sign); }
513:
514:   // Phasespace weights [note w versus w^2 Jacobian, taken into account]
515:   // [m12 volume (GeV) x dm_12^2 x dphi1_2 x dphi2]
516:   const double volume = (m12bound[1] - m12bound[0]);
517:   x = x * volume; // operator overloaded
518:
519:   // Return phase space weight
520:   return x;
521: }
522:
523: // Decay setup for NBodyPhaseSpace
524: // Uses the "sorting algorithm", see description in F. James, CERN/68
525: // 01.1, 1968
526:
527: // Returns: W for a valid and -1.0 for a kinematically impossible
528:
529: template <typename T1, typename T2>
530: inline MCW NBodySetup(const T1 smother, double M0, const std::vector<double> km,
531:   double M1, double M2, const std::vector<double> km_eff, std::vector<double> ipnorm, bool unweight,
532:   T2 kring) {
533:   // Decay multiplicity (i -> N decay)
534:   const unsigned int N = m.size();
535:   std::vector<double> randvec(N - 2, 0.0);
536:
537:   // Random variables for volume
538:   auto fillrandvec = [&](double value) -> void { value = rng.U(0, 1); };
539:
540:   // Effective (intermediate) masses
541:   M_eff[0] = m[0]; // First daughter
542:   M_eff[N - 1] = M0; // Mother
543:
544:   // Cumulative sum of daughter masses
545:   std::vector<double> sumvec;
546:   gra:math::CumulativeSum(sumvec);
547:
548:   // Sampling mass interval
549:   const double DELTA = M_eff[N - 1] - sumvec[N - 1];
550:
551:   // Generate effective masses functor, where sumvec[i] defines the minimum
552:   auto calcmass = [&](i) { double {
553:     for (std::size_t i = 1; i < N - 1; ++i) { M_eff[i] = sumvec[i] + randvec[i - 1] * DELTA; }
554:   } };
555:
556:   // Recursively find maximum weight for Acceptance-Rejection
557:   double w_max = 1.0;
558:   std::vector<double> m_minmax = {0.0, DELTA * m[N - 1]};
559:   for (std::size_t i = 1; i < N; ++i) {
560:     m_minmax[i] = m[i + 1];
561:     m_minmax[i] += m[i];
562:     w_max = DecayMomentum(m_minmax[i], m_minmax[0], m[i]);
563:   }
564:
565:   // Functor to calculate decay momentum to pnorm[i], and return product
566:   auto calcmomentum = [&](i) { double {
567:     double w = 1.0;
568:     for (std::size_t i = 1; i < N; ++i) {
569:       pnorm[i] = DecayMomentum(M_eff[i], M_eff[i - 1], m[i]);
570:       w *= pnorm[i];
571:     }
572:     return w;
573:   } };
574:
575:   // ----- m12 -----
576:   // Acceptance-Rejection
577:   // ----- m12 -----
578:   if (unweight == false) { w_max = 0; }
579:
580:   const unsigned int MAXTRIAL = 1e8;

```

```

./include/Granitti/MKinematics.h      9/20
665:   p[i] = T1(pt * std::cos(phi), pt * std::sin(phi), pnorm[i] * costheta,
666:     msqr(pow2(m11) + pow2(pnorm[i])));
667:
668:   // Boost to the lab frame
669:   const int sign = 1; // positive
670:   LorentzBoost(k, k.M(), p[i], sign);
671:
672:   // Subtract the momentum
673:   k = p[i];
674:
675:   // Finally, we are left with only one daughter
676:   p[0] = k;
677:
678:   return x; // Phase-space weight
679: }
680:
681:
682: // SO(3) Rotations
683:
684: // Rotation matrix from 3-vector a to 3-vector b (both need to be unit vectors with ||x|| = 1)
685: // ----- m12 -----
686: inline MMatrix<double> RotFromOne(const std::vector<double> ka, const std::vector<double> kb) {
687:   if (a.size() != 3 || b.size() != 3) {
688:     throw std::invalid_argument("Kinematics::RotFromOne: Input should be 3-vectors");
689:   }
690:
691:   // Dot product
692:   const double c = mtoper:VecVecMultiply(a, b);
693:
694:   // Cross product
695:   const std::vector<double> w = mtoper:Cross(a, b);
696:
697:   // Skew symmetric cross product matrix
698:   const MMatrix<double> Vx = {{0, -v[2], v[1]}, {v[2], 0, -v[0]}, {-v[1], v[0], 0}};
699:   // Identity
700:   const MMatrix<double> I(3, 3, "eye");
701:
702:   // Rotation matrix, with singularity at c = -1 (and a b back-to-back)
703:   const MMatrix<double> R = I + Vx * (Vx * Vx) * (1.0 / (1.0 + c));
704:
705:   return R;
706: }
707:
708: // Rotate 4-vector spatial part by (theta, phi)
709:
710: template <typename T>
711: inline void Rotate4(T sp, double theta, double phi) {
712:   const double c1 = std::cos(theta);
713:   const double s1 = std::sin(theta);
714:   const double c2 = std::cos(phi);
715:   const double s2 = std::sin(phi);
716:
717:   // 3x3 Rotation matrix
718:   const MMatrix<double> R = {{c1 * c2, -s2, s1 * c2}, {c1 * s2, c2, s1 * s2}, {-s1, 0.0, c1}};
719:
720:   // Rotate
721:   const std::vector<double> p3 = {p.x(), p.y(), p.z()};
722:   const std::vector<double> p3new = R * p3;
723:
724:   p = T(p3new[0], p3new[1], p3new[2], p.E());
725:
726: // Active SO3-rotation counter-clockwise around x-axis
727:
728: template <typename T>
729: inline void RotateX(T sp, double angle) {
730:   const double cosh = std::cos(angle);
731:   const double sinh = std::sin(angle);
732:
733:   // Rotation matrix applied
734:   const double px = p.x();
735:   const double py = cosh * p.y() - sinh * p.z();
736:   const double pz = sinh * p.y() + cosh * p.z();
737:
738:   p = T(px, py, pz, p.E());
739:
740: // Active SO3-rotation counter-clockwise around y-axis
741:
742: template <typename T>
743: inline void RotateY(T sp, double angle) {
744:   const double cosh = std::cos(angle);
745:   const double sinh = std::sin(angle);
746:
747:   // Rotation matrix applied
748:   const double px = cosh * p.x() + sinh * p.z();
749:   const double py = p.y();

```

./include/Granitti/MKinematics.h 10/20

```
748: const double pz = -sinh * p.Fx() + cosh * p.Fz();
749:
750: p = T(pw, py, pz, p.E());
751:
752:
753: // Active SO3-rotation counterclockwise around z-axis
754: template <typename T>
755: inline void RotateT(T & p, double angle) {
756:   const double cosh = std::cosh(angle);
757:   const double sinh = std::sinh(angle);
758:
759:   // Rotation matrix applied
760:   const double px = cosh * p.Fx() - sinh * p.Fy();
761:   const double py = sinh * p.Fx() + cosh * p.Fy();
762:   const double pz = p.Fz();
763:
764:   p = T(pw, py, pz, p.E());
765: }
766:
767:
768: // Find the closest 4-vector on lightcone
769: //
770: // Solution by Lagrange multipliers Wabla f = lambda Wabla g
771: //
772: // eq(1) = 2*(px-px0) = -2*lambda*px
773: // eq(2) = 2*(py-py0) = -2*lambda*py
774: // eq(3) = 2*(pz-pz0) = -2*lambda*pz
775: // eq(4) = 2*(E-E0) = 2*lambda*E
776: // eq(5) = E^2 = (px^2 + py^2 + pz^2)
777: //
778: // S = solve(eq, {E, px, py, pz, lambda})
779: //
780: inline M4Vec LagrangeLightCone(const M4Vec sp0) {
781:   // Spacelike (q^2 < 0) or timelike (q^2 > 0) input
782:   const double sign = p0.M0() <= 0 ? 1.0 : -1.0;
783:
784:   const double p3mod2 = p0.F3Mod2();
785:   const double p3mod = math::msqrt(p3mod2);
786:
787:   const double E = p0.E() / 2 + sign * p3mod / 2;
788:   const double alpha = (p3mod2 + sign * p0.E() / p3mod) / (2 * p3mod);
789:   return M4Vec(alpha * p0.Fx(), alpha * p0.Fy(), alpha * p0.Fz(), E);
790: }
791:
792:
793: // Kinematic transform in process: p1 + p2 -> (p), where
794: //
795: // p1, p2 are massless spacelike q^2 < 0 -> transformed to lightlike q^2 = 0
796: // (p) massive/massless final states with q^2 = m^2
797: //
798: //
799: inline void OffShellLightCone(M4Vec sp1, M4Vec sp2, std::vector<M4Vec> & p) {
800:   const int MAXITER = 15;
801:   const double STOPEPS = 1e-10;
802:
803:   // 4-momentum sum
804:   auto psumfunc = [&] () {
805:     M4Vec sum(D, D, D, 0);
806:     for (const auto & i : aux::indices(p)) { sum += p[i]; }
807:     return sum;
808:   };
809:
810:   // Energy sum
811:   auto Esumfunc = [&] () {
812:     double sum = 0.0;
813:     for (const auto & i : aux::indices(p)) { sum += p[i].E(); }
814:     return sum;
815:   };
816:
817:   //
818:   //
819:   // Fractions
820:   auto Efrac = [&] () {
821:     const double Esum = Esumfunc();
822:     std::vector<double> f(p.size());
823:     for (const auto & i : aux::indices(p)) { f[i] = p[i].E() / Esum; }
824:     return f;
825:   };
826:
827:   //
828:   // Set initial state spacelike (q^2 < 0) particles to lightcone by
829:   // conserving 3-momentum and z-increasing energy
830: }
```

./include/Granitti/MKinematics.h 12/20

```
914: // "HE" : Helicity
915: // "PG" : Pseudo-Gottfried-Jackson
916: //
917: // Collins-Soper: Quantization z-axis defined by the bi-vector between
918: // initial state p1 and (-p2) (NEGATIVE) directions in the (resonance) system rest frame,
919: // where p1 and p2 are the initial state proton 3-momentum.
920: //
921: // Anti-Helicity: Quantization z-axis defined by the bisector vector between
922: // initial state p1 and (p2) (POSITIVE) directions in the (resonance) system rest frame,
923: // where p1 and p2 are the initial state proton 3-momentum.
924: //
925: // Helicity: Quantization axis defined by the resonance system
926: // 3-momentum vector in the colliding beams frame (lab frame).
927: //
928: // Pseudo-Gottfried-Jackson: Quantization axis defined by the initial state
929: // proton p1 (or p2) 3-momentum vector in the (resonance) system rest frame.
930: //
931: //
932: template <typename T>
933: inline void LorentzFramePrepare(const std::vector<T> & p, const T & X, const T & pbeam1,
934:                               const T & pbeam2, T & pboost, T & pboost, std::vector<T> & pboost) {
935:   const double M = X.M();
936:
937:   // 2. Boost each particle to the system rest frame
938:   pboost = p;
939:   for (const auto & k : gra::aux::indices(p)) {
940:     gra::kinematics::LorentzBoost(X, M, pboost[k], -1); // note minus sign
941:   }
942:
943:   // 3. Boost initial state protons
944:   pboost = pbeam1;
945:   pboost = pbeam2;
946:   gra::kinematics::LorentzBoost(X, M, pboost, -1); // note minus sign
947:   gra::kinematics::LorentzBoost(X, M, pboost, -1); // note minus sign
948: }
949:
950:
951: template <typename T>
952: inline void LorentzFrame(std::vector<T> & pfour, const T & pboost, const T & pboost,
953:                         int direction) {
954:   const std::vector<double> pboost3 = pboost.F3();
955:   const std::vector<double> pboost3 = pboost3.F3();
956:
957:   // Frame rotation x-y-z-axes
958:   std::vector<double> xaxis;
959:   std::vector<double> yaxis;
960:   std::vector<double> zaxis;
961:   std::vector<double> xaxis;
962:
963:   // NON-ROTATED FRAME AXIS DEFINITION ##
964:   if (frameType == "CM") {
965:     xaxis = {0, 0, 1};
966:     yaxis = {0, 1, 0};
967:   }
968:
969:   // ## COLLINS-SOPER FRAME POLARIZATION AXIS DEFINITION ##
970:   else if (frameType == "CS" || frameType == "CS") {
971:     xaxis = gra::matoper::Unit(
972:       gra::matoper::Minus(gra::matoper::Unit(pboost3)), gra::matoper::Unit(pboost3));
973:   }
974:   // ## ANTI-HELICITY FRAME POLARIZATION AXIS DEFINITION ##
975:   else if (frameType == "AH" || frameType == "CS") {
976:     xaxis = gra::matoper::Unit(
977:       gra::matoper::Plus(gra::matoper::Unit(pboost3)), gra::matoper::Unit(pboost3));
978:   }
979:   // ## HELICITY FRAME POLARIZATION AXIS DEFINITION ##
980:   else if (frameType == "HE" || frameType == "CS") {
981:     xaxis = gra::matoper::Unit(gra::matoper::Negat(gra::matoper::Plus(pboost3, pboost3)));
982:   }
983:   // ## PSEUDO-GOTTFRIED-JACKSON AXIS DEFINITION: |1) or |2) ##
984:   else if (frameType == "PJ" || frameType == "GJ") {
985:     if (direction == -1) {
986:       xaxis = gra::matoper::Unit(pboost3);
987:     } else if (direction == 1) {
988:       xaxis = gra::matoper::Unit(pboost3);
989:     }
990:     throw std::invalid_argument("gra::kinematics::LorentzFrame: Invalid direction <+>");
991:   }
992:   else {
993:     throw std::invalid_argument("gra::kinematics::LorentzFrame: Unknown frame <+> + frameType + >");
994:   }
995: }
996: // y-axis
```

./include/Granitti/MKinematics.h 11/20

```
831: p1.SetE(p1.F3Mod());
832: p2.SetE(p2.F3Mod());
833:
834:
835: // Get sum
836: const M4Vec q = p1 + p2;
837:
838: // Normalize vvec[q] to unit length
839: const std::vector<double> qvec = gra::matoper::Unit(q.F3());
840:
841: // Final state masses
842: std::vector<double> m(p.size());
843: for (const auto & i : aux::indices(p)) { m[i] = p[i].M(); }
844:
845: int iter = 0;
846:
847: // for (const auto i : aux::indices(p)) { p[i].Print(); }
848: // std::cout << "iter " * iter << " * " * q - psumfunc().M2() << " <std::endl;
849:
850: while (true) {
851:   // 3-Momentum scaling and energy subtraction step
852:   //
853:   // New energy difference
854:   const double dE = Esumfunc() - q.E();
855:
856:   // Current energy fractions
857:   std::vector<double> f = Efrac();
858:
859:   for (const auto & i : aux::indices(p)) {
860:     const double E = p[i].E() - dE / N; // Scaling distributed by simple 1/N
861:     // gra::matoper::ScaleVector(D3_this, f[i]); // Scaling distributed by fractions
862:     const double a = gra::math::msqrt(E * E - m[i] * m[i]) / p[i].F3Mod();
863:     p[i].SetE(a * p[i].F3(), a * p[i].Fy(), a * p[i].Fz(), E);
864:   }
865:
866:   //
867:   // 3-Momentum subtraction step
868:   //
869:   // New 3-momentum difference
870:   const std::vector<double> D3 = (psumfunc() - q).F3();
871:
872:   // Current energy fractions
873:   f = Efrac();
874:
875:   for (const auto & i : aux::indices(p)) {
876:     std::vector<double> D3_this = D3;
877:     gra::matoper::ScaleVector(D3_this, 1.0 / N); // Scaling distributed by simple 1/N
878:     // gra::matoper::ScaleVector(D3_this, f[i]); // Scaling distributed by fractions
879:     p[i].SetE3(gra::matoper::Minus(p[i].F3(), D3_this));
880:     p[i].SetE(gra::math::msqrt(p[i].F3Mod() * m[i] * m[i]));
881:   }
882:
883:   //
884:   // 3-Momentum rotation step
885:   //
886:   // Add and normalize to unit length, find rotation
887:   const std::vector<double> avc = gra::matoper::Unit(psumfunc().F3());
888:   const gra::Matrix<double> R = gra::kinematics::Rotemove(avc, qvec);
889:
890:   // Rotate
891:   for (const auto & i : aux::indices(p)) { p[i].SetE3(R * p[i].F3()); }
892:
893:   //
894:   //
895:   //
896:   //
897:   //
898:   //
899:   const double IVM = std::abs(q - (psumfunc())).M2();
900:   if (iter > MAXITER || IVM < STOPEPS) break;
901:
902:   // std::cout << "iter " * iter << " * " * IVM << " <std::endl;
903:
904:   // for (const auto & i : aux::indices(p)) { p[i].Print(); }
905:   // std::cout << " * " * iter << " * " * IVM << " <std::endl;
906:
907:   //
908:   //
909:   // A unified Lorentz Transform function to 'frameType' give below:
910:   //
911:   // "CS" : Collins-Soper
912:   // "AH" : Anti-Helicity (Anti-CS)
913:
914:
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./include/Granitti/MKinematics.h      14/20
1080: }
1081: //
1082: //
1083: // From Lab to Collins-Soper frame
1084: // Quantization z-axis is defined as the bisector of two beams, in the rest
1085: // frame of the resonance
1086: //
1087: // Input:  p = Set of 4-momentum to be transformed
1088: //         X = System 4-momentum
1089: //         p1 = Beam 1 4-momentum
1090: //         p2 = Beam 2 4-momentum
1091: //
1092: template <typename T>
1093: inline void CFrame(std::vector<T> &p, const T &X, const T &p1, const T &p2, bool DEBUG = false) {
1094: // *****
1095: if (DEBUG) {
1096:   printf("\n\n :COLLINS-SOPER FRAME: \n\n");
1097:   printf("CFrame: Daughters in LAB FRAME: \n");
1098:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1099: }
1100: // *****
1101: // Boost particles to the central system rest frame
1102: for (const auto &i : gra::aux::indices(p)) {
1103:   gra::kinematics::LorentzBoost(X, X.M(), p[i], -1); // Note the minus sign
1104: }
1105: //
1106: //
1107: T pib = p1;
1108: T p2b = p2;
1109: //
1110: // Boost the beam particles
1111: gra::kinematics::LorentzBoost(X, X.M(), pib, -1); // Note the minus sign
1112: gra::kinematics::LorentzBoost(X, X.M(), p2b, -1); // Note the minus sign
1113: //
1114: // Now get the 3-momenta
1115: const std::vector<double> pibP3 = pib.P3();
1116: const std::vector<double> p2bP3 = p2b.P3();
1117: //
1118: // Frame rotation x-y-z axes
1119: std::vector<double> xaxis;
1120: std::vector<double> yaxis;
1121: std::vector<double> zaxis;
1122: //
1123: // Collins-Soper bisector vector
1124: xaxis = gra::matoper::Unit(
1125:   gra::matoper::Minus(gra::matoper::Unit(pibP3), gra::matoper::Unit(p2bP3)));
1126: // * ALTERNATIVE WAY, but gives random reflection (rotation around Z by PI)
1127: //
1128: MVec b1jactor;
1129: b1jactor.SetP3(xaxis);
1130: //
1131: // Now get the rotation angle, note the minus
1132: const double Z_angle = -b1jactor.Phi();
1133: const double Y_angle = -b1jactor.Theta();
1134: //
1135: // Rotate final states
1136: for (const auto &i : gra::aux::indices(p)) {
1137:   gra::kinematics::RotateZ(p[i], Z_angle);
1138:   gra::kinematics::RotateY(p[i], Y_angle);
1139:   gra::kinematics::RotateX(p[i], math::PI);
1140: }
1141: //
1142: //
1143: //
1144: xaxis = gra::matoper::Unit(
1145:   gra::matoper::Cross(gra::matoper::Unit(pibP3), gra::matoper::Unit(p2bP3)));
1146: //
1147: xaxis = gra::matoper::Unit(gra::matoper::Cross(yaxis, zaxis)); // x = y [cross product] z
1148: //
1149: // Create SO(3) rotation matrix for the new coordinate axes
1150: const Matrix<double> R = {xaxis, yaxis, zaxis}; // Axes as rows
1151: //
1152: // Rotate all vectors
1153: for (const auto &k : gra::aux::indices(p)) {
1154:   const std::vector<double> pNew = R * p[k].P3(); // Spatial part rotation Rp -> p'
1155:   p[k] = T(pNew[0], pNew[1], pNew[2], p[k].E()); // Full 4-momentum [px; py; pz; E]
1156: }
1157: //
1158: // *****
1159: if (DEBUG) {
1160:   printf("CFrame: Daughters in CS FRAME: \n");
1161:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1162: }

```

```

./include/Granitti/MKinematics.h      16/20
1246: // Input:  p = Set of 4-momentum to be transformed
1247: //         X = System 4-momentum
1248: //         direction = -1 or 1
1249: //         p_beam_plus = 4-momentum of beam1
1250: //         p_beam_minus = 4-momentum of beam2
1251: //
1252: template <typename T>
1253: inline void PFrame(std::vector<T> &p, const T &X, const T &dir, const T &p_beam_plus,
1254:   const T &p_beam_minus, const T &XNEW, bool DEBUG = false) {
1255: // *****
1256: printf("\n\n :PSEUDO-GOTTFRIED-JACKSON FRAME: \n\n");
1257: printf("PFrame: Daughters in LAB FRAME: \n");
1258: printf("PFrame: Daughters in LAB FRAME: \n");
1259: for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1260: }
1261: // *****
1262: // Boost particles to the central system rest frame
1263: for (const auto &i : gra::aux::indices(p)) {
1264:   gra::kinematics::LorentzBoost(X, X.M(), p[i], -1); // Note the minus sign
1265: }
1266: // Boost initial state protons to the central system rest frame
1267: T proton_p = p_beam_plus;
1268: T proton_m = p_beam_minus;
1269: //
1270: gra::kinematics::LorentzBoost(X, X.M(), proton_p, -1); // Note the minus sign
1271: gra::kinematics::LorentzBoost(X, X.M(), proton_m, -1); // Note the minus sign
1272: //
1273: // Now get the rotation angles, note the minus
1274: double Z_angle = 0;
1275: double Y_angle = 0;
1276: //
1277: if (direction == -1) {
1278:   Z_angle = -proton_p.Phi();
1279:   Y_angle = -proton_p.Theta();
1280: } else if (direction == 1) {
1281:   Z_angle = -proton_m.Phi();
1282:   Y_angle = -proton_m.Theta();
1283: } else {
1284:   printf("PFrame: Input direction %d not valid (1,-1)", direction);
1285:   return;
1286: }
1287: //
1288: // Rotation function
1289: const auto rotate = [&](T &p) {
1290:   gra::kinematics::RotateZ(p, Z_angle);
1291:   gra::kinematics::RotateY(p, Y_angle);
1292:   // gra::kinematics::RotateX(p, math::PI); // Reflection
1293: };
1294: //
1295: // Rotate final states
1296: for (const auto &i : gra::aux::indices(p)) { rotate(p[i]); }
1297: //
1298: // *****
1299: if (DEBUG) {
1300:   printf("PFrame: Daughters in Pseudo-Gottfried-Jackson FRAME: \n");
1301:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1302: }
1303: //
1304: // Rotate proton
1305: rotate(proton_p);
1306: //
1307: // Rotate other proton
1308: rotate(proton_m);
1309: //
1310: printf("Protons are on (sx)-plane: \n");
1311: printf("USER chosen direction along %d beam direction \n", direction);
1312: printf("sx = %s Proton in Pseudo-Gottfried-Jackson FRAME: \n");
1313: proton_p.Print();
1314: //
1315: // Rotate other proton
1316: rotate(proton_m);
1317: printf("sx = %s Proton in Pseudo-Gottfried-Jackson FRAME: \n");
1318: proton_m.Print();
1319: //
1320: // Mandelstam invariant
1321: printf("Mandelstam s = %0.5f \n", (proton_p + proton_m).M());
1322: // *****
1323: //
1324: //
1325: //
1326: // From Lab to the Helicity frame
1327: // Quantization z-axis as the direction of the resonance in the Lab Frame

```

```

./include/Granitti/MKinematics.h      15/20
1163: // *****
1164: //
1165: //
1166: // From Lab to Gottfried-Jackson frame
1167: // Quantization z-axis spanned by the propagator (momentum transfer vector)
1168: // momentum in the rest frame of the central system
1169: //
1170: // Input:  p = Set of 4-momentum to be transformed
1171: //         X = System 4-momentum
1172: //         direction = -1 or 1
1173: //         q1 = 4-momentum transfer vector 1
1174: //         q2 = 4-momentum transfer vector 2
1175: //
1176: template <typename T>
1177: inline void GJFrame(std::vector<T> &p, const T &X, int direction, const T &q1, const T &q2,
1178:   bool DEBUG = false) {
1179: // *****
1180: if (DEBUG) {
1181:   printf("\n\n :GOTTFRIED-JACKSON FRAME: \n\n");
1182:   printf("GJFrame: Daughters in LAB FRAME: \n");
1183:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1184: }
1185: // *****
1186: // Boost particles to the central system rest frame
1187: for (const auto &i : gra::aux::indices(p)) {
1188:   gra::kinematics::LorentzBoost(X, X.M(), p[i], -1); // Note the minus sign
1189: }
1190: //
1191: // Boost the propagators
1192: T qboost = q1;
1193: T qboost = q2;
1194: gra::kinematics::LorentzBoost(X, X.M(), qboost, -1); // Note the minus sign
1195: gra::kinematics::LorentzBoost(X, X.M(), qboost, -1); // Note the minus sign
1196: //
1197: // Now get the rotation angle, note the minus
1198: double Z_angle = 0;
1199: double Y_angle = 0;
1200: if (direction == -1) {
1201:   Z_angle = -qboost.Phi();
1202:   Y_angle = -qboost.Theta();
1203: } else if (direction == 1) {
1204:   Z_angle = -qboost.Phi();
1205:   Y_angle = -qboost.Theta();
1206: } else {
1207:   printf("GJFrame: direction not -1 or 1");
1208: }
1209: //
1210: // these std::invalid_argument("GJFrame: direction not -1 or 1");
1211: //
1212: //
1213: const auto rotate = [&](T &p) {
1214:   gra::kinematics::RotateZ(p, Z_angle);
1215:   gra::kinematics::RotateY(p, Y_angle);
1216:   // gra::kinematics::RotateX(p, math::PI); // Reflection
1217: };
1218: //
1219: // Rotate final states
1220: for (const auto &i : gra::aux::indices(p)) { rotate(p[i]); }
1221: // *****
1222: if (DEBUG) {
1223:   printf("GJFrame: Daughters in Gottfried-Jackson FRAME: \n");
1224:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1225: }
1226: //
1227: // Rotate propagator
1228: rotate(qboost);
1229: rotate(qboost);
1230: //
1231: printf("GJFrame: Propagators are along z-axis \n");
1232: printf("GJFrame: Propagator 1 in Gottfried-Jackson FRAME: \n");
1233: qboost.Print();
1234: //
1235: printf("GJFrame: Propagator 2 in Gottfried-Jackson FRAME: \n");
1236: qboost.Print();
1237: // *****
1238: //
1239: //
1240: //
1241: // From Lab to Pseudo-Gottfried-Jackson frame
1242: // Quantization z-axis spanned by the beam proton +z (or -z) momentum
1243: // in the rest frame of the central system
1244: //
1245: //

```

```

./include/Granitti/MKinematics.h      17/20
1329: // Input:  p = Set of 4-momentum to be transformed
1330: //         X = System 4-momentum
1331: //         direction = -1 or 1
1332: //
1333: template <typename T>
1334: inline void HXFrame(std::vector<T> &p, const T &X, bool DEBUG = false) {
1335: // *****
1336: printf("\n\n :HELICITY FRAME: \n\n");
1337: printf("HXFrame: Daughters in LAB FRAME: \n");
1338: printf("HXFrame: Daughters in LAB FRAME: \n");
1339: for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1340: }
1341: // *****
1342: // Rotation angles, note the minus
1343: const double Z_angle = -X.Phi();
1344: const double Y_angle = -X.Theta();
1345: //
1346: const auto rotate = [&](T &p) {
1347:   gra::kinematics::RotateZ(p, Z_angle);
1348:   gra::kinematics::RotateY(p, Y_angle);
1349:   gra::kinematics::RotateX(p, math::PI); // Reflection
1350: };
1351: //
1352: for (const auto &i : gra::aux::indices(p)) { rotate(p[i]); }
1353: //
1354: // *****
1355: if (DEBUG) {
1356:   printf("HXFrame: Daughters in ROTATED LAB FRAME: \n");
1357:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1358: }
1359: //
1360: T ex(1, 0, 0, 0);
1361: T ey(0, 1, 0, 0);
1362: T ez(0, 0, 1, 0);
1363: //
1364: // x -> x', y -> y', z -> z'
1365: rotate(ex);
1366: rotate(ey);
1367: rotate(ez);
1368: //
1369: printf("HXFrame: AXIS vectors after rotation: \n");
1370: //
1371: ex.Print();
1372: ey.Print();
1373: ez.Print();
1374: //
1375: //
1376: //
1377: // Construct the central system 4-momentum in ROTATED FRAME
1378: T XNEW = X;
1379: rotate(XNEW);
1380: //
1381: // Boost particles to the central system rest frame
1382: // -> Helicity frame obtained
1383: for (const auto &i : gra::aux::indices(p)) {
1384:   gra::kinematics::LorentzBoost(XNEW, XNEW.M(), p[i], -1); // Note the minus sign
1385: }
1386: // *****
1387: if (DEBUG) {
1388:   printf("HXFrame: Daughters after boost in HELICITY FRAME: \n");
1389:   for (const auto &i : gra::aux::indices(p)) { p[i].Print(); }
1390: }
1391: //
1392: printf("HXFrame: Central system in LAB FRAME: \n");
1393: XNEW.Print();
1394: //
1395: printf("HXFrame: Central system in ROTATED LAB FRAME: \n");
1396: XNEW.Print();
1397: //
1398: printf("\n");
1399: // *****
1400: //
1401: //
1402: //
1403: //
1404: // "Rest frame" (no boost, beam axis = z-axis / spin quantization axis)
1405: //
1406: // Input:  p = Set of 4-momentum to be transformed
1407: //         X = System 4-momentum
1408: //
1409: template <typename T>
1410: inline void CRFrame(std::vector<T> &p, const T &X, bool DEBUG = false) {
1411: // *****

```

```

./include/Granitti/MKinematics.h      18/20
1412: if (DEBUG) {
1413:   printf("\n\n : REST FRAME: \n");
1414: }
1415:   printf("CMFrame: Daughters in LAB FRAME: \n");
1416:   for (const auto &i : gra::aux::indices(p) | p[i].Print(); )
1417:   }
1418: }
1419: // -----
1420: // Boost particles to the central system rest frame
1421: for (const auto &i : gra::aux::indices(p) |
1422:   gra::kinematics::LorentzBoost(X, X.M(), p[i], -1) // Note the minus sign
1423:   );
1424: }
1425: // -----
1426: if (DEBUG) {
1427:   printf("CMFrame: Daughters in REST FRAME: \n");
1428:   for (const auto &i : gra::aux::indices(p) | p[i].Print(); )
1429:   }
1430: }
1431: }
1432: // -----
1433: } // namespace kinematics
1434: }
1435: // -----
1436: // Lorentz scalars and other common kinematic variables
1437: class LORENTZSCALAR {
1438: public:
1439:   LORENTZSCALAR() {}
1440:   ~LORENTZSCALAR() {}
1441:   delete GlobalSudakovPtr;
1442:   delete GlobalPdfPtr;
1443: }
1444: // -----
1445: // Particle Database
1446: MFDG PDG;
1447: // -----
1448: // -----
1449: // PDF access
1450: // Sudakov/pdf routines
1451: MSudakov *GlobalSudakovPtr = nullptr;
1452: // Normal pdfs
1453: std::string LHAPDFSET = "null";
1454: LHAPDF::PDF *GlobalPdfPtr = nullptr;
1455: int pdf_trials = 0;
1456: // -----
1457: // Cascade sampling forced control initiated by corresponding amplitude functions
1458: bool FORCE_FLATBASE2 = false;
1459: double FORCE_OFFSHELL = -1;
1460: // -----
1461: // Resonances read from JHUP input
1462: std::map<std::string, PARAM_REP> RESONANCES;
1463: // Helicity amplitudes returned by amplitude functions,
1464: // used in screening loop etc.
1465: std::vector<std::complex<double>> hamp;
1466: // Central system decaytree
1467: std::vector<MDecayBranch> decaytree;
1468: // Forward system decaytree
1469: MDecayBranch decayforward1;
1470: MDecayBranch decayforward2;
1471: // Integral weight container (central system phase space)
1472: gra::kinematics::MCW MW;
1473: // Active in the factorized phase space product (by default, true)
1474: bool FS_active = true;
1475: // Sum containers
1476: gra::kinematics::MCMSUM DM_sum;
1477: gra::kinematics::MCMSUM DM_sum_exact;
1478: // Initial states
1479: gra::MParticle beam1;
1480: gra::MParticle beam2;
1481: }

```

```

./include/Granitti/MKinematics.h      20/20
1578: // Fiducial cuts (default parameters set here)
1579: struct FIDUCIAL {
1580:   bool active = false;
1581: };
1582: // "Central particles"
1583: double eta_min = -30.0;
1584: double eta_max = 30.0;
1585: // Rapidity
1586: double rap_min = -30.0;
1587: double rap_max = 30.0;
1588: // Transverse momentum
1589: double pt_min = 0.0;
1590: double pt_max = 1000000.0;
1591: // Energy
1592: double Et_min = 0.0;
1593: double Et_max = 1000000.0;
1594: // "Central system"
1595: double M_min = 0.0;
1596: double M_max = 1000000.0;
1597: // Rapidity
1598: double Y_min = -30.0;
1599: double Y_max = 30.0;
1600: // Transverse momentum
1601: double Pt_min = 0.0;
1602: double Pt_max = 1000000.0;
1603: // "Forward system"
1604: double forward_M_min = 0.0;
1605: double forward_M_max = 1000000.0;
1606: double forward_Y_min = -30.0;
1607: double forward_Y_max = 30.0;
1608: double forward_Pt_min = 0.0;
1609: double forward_Pt_max = 1000000.0;
1610: }
1611: // For veto
1612: struct VETODOMAIN {
1613:   double eta_min = -30.0;
1614:   double eta_max = 30.0;
1615:   double pt_min = 0.0;
1616:   double pt_max = 1000000.0;
1617: };
1618: struct VETOCUT {
1619:   bool active = false;
1620:   std::vector<VETODOMAIN> cuts;
1621: };
1622: } // namespace gra
1623: }
1624: #endif

```

```

./include/Granitti/MKinematics.h      19/20
1495: // Four momenta of initial and final states
1496: M4Vec pbeam1;
1497: M4Vec pbeam2;
1498: std::vector<M4Vec> pfinal;
1499: std::vector<M4Vec> pfinal_orig;
1500: // -----
1501: // Basic Lorentz scalars
1502: double s = 0.0;
1503: double t = 0.0;
1504: double u = 0.0;
1505: double spt_s = 0.0;
1506: // Sub-energies^2
1507: double s1 = 0.0;
1508: double s2 = 0.0;
1509: // 4-momentum transfer squared
1510: double t1 = 0.0;
1511: double t2 = 0.0;
1512: // Sub-Mandelstam variables
1513: double s_hat = 0.0;
1514: double t_hat = 0.0;
1515: double u_hat = 0.0;
1516: // Central system
1517: double m2 = 0.0;
1518: double Y = 0.0;
1519: double Pt = 0.0;
1520: // -----
1521: // Maximum 10 particles in the final state
1522: // (3 protons + 8 direct central system)
1523: double sz[10][10] = {{0,0}};
1524: double tt_i[10] = {0,0};
1525: double tt_j[10] = {0,0};
1526: double tt_xy[10][10] = {{0,0}};
1527: // Longitudinal momentum losses
1528: double x1 = 0.0;
1529: double x2 = 0.0;
1530: // Shower-x
1531: double xb1 = 0.0;
1532: double xb2 = 0.0;
1533: // Forward proton excitation tag
1534: bool excitel = false;
1535: bool excite2 = false;
1536: // Propagators from proton1 (up) and proton2 (down)
1537: M4Vec q1;
1538: M4Vec q2;
1539: // Propagator pt
1540: double qT1 = 0.0;
1541: double qT2 = 0.0;
1542: // -----
1543: // Generator cut (default parameters set here)
1544: struct GENCUT {
1545:   // Continuum phase space <C>
1546:   double rap_min = -9.0;
1547:   double rap_max = 9.0;
1548:   double kt_min = 0.0;
1549:   double kt_max = -1.0; // Keep at -1 for user setup trigger
1550:   // Factorized phase space <F>
1551:   double Y_min = -9.0;
1552:   double Y_max = 9.0;
1553:   double M_min = 0.0;
1554:   double M_max = 0.0;
1555:   // Both <C> and <F> class forward legs
1556:   double forward_pt_min = -1.0; // Keep at -1 for user setup trigger
1557:   double forward_pt_max = -1.0; // Keep at -1 for user setup trigger
1558: };
1559: // -----
1560: // Quasi-Elastic phase space <Q> or forward excitation
1561: double XI_min = 0.0;
1562: double XI_max = 1.0;
1563: };

```

```

./include/Granitti/MFragment.h        1/1
1: // Toy fragmentation class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: #include "MFragmentation.h"
6: #define MFRAGMENTATION_H
7: // C++
8: #include <complex>
9: #include <map>
10: #include <random>
11: #include <vector>
12: // Own
13: #include "Granitti/MFDG.h"
14: #include "Granitti/MRandom.h"
15: namespace gra {
16: class MFragment {
17: public:
18:   MFragment() {}
19:   ~MFragment() {}
20:   static double TubeFragment(const M4Vec &mother, double M0, const std::vector<double> &lm,
21:     std::vector<M4Vec> &lp, double q, double T, double lambda, double maxpt,
22:     MRandom &rng, const std::string &distri);
23:   static bool SolveAlpha(double alpha, double M0, const std::vector<double> &lm,
24:     const std::valarray<double> &int, const std::valarray<double> &iy);
25:   static void ExpPowMND(double q, double T, double maxpt, const std::vector<double> &smas,
26:     std::vector<double> &ix, MRandom &rng);
27:   static void GetDecayStatus(const std::vector<int> &spdgcode, std::vector<bool> &isstable);
28:   static void GetForwardMass(double smas1, double smas2, bool excitel1, bool excitel2,
29:     unsigned int excite, MRandom &random);
30:   static void GetSingleForwardMass(double smas, MRandom &random);
31:   static void HitAreaTable(int Q, double M0, std::vector<int> &spdgcode, MRandom &rng);
32:   static bool PickParticles(double M, unsigned int N, int S, int B, int Q, int Q,
33:     std::vector<double> &smas, std::vector<int> &spdgcode, const MFDG &PDG,
34:     MRandom &rng);
35: };
36: } // namespace gra
37: #endif

```

```

./include/Granitti/MDimArray.h 1/1
1: // Multidimensional fixed size arrays via template alias technique
2: //
3: //
4: // Usage:
5: // MultidimArray<int, 3, 2> arr {1, 2, 3, 4, 5, 6};
6: // assert(arr[1][1] == 4);
7: //
8: // (c) 2017-2020 Mikael Mieskolainen
9: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
10:
11: #ifndef MDIMARRAY_H
12: #define MDIMARRAY_H
13:
14: #include <array>
15: #include <complex>
16:
17: namespace gra {
18: // Note the index order below, in order to get the row major format
19: //
20: template <typename T, size_t D1, size_t D2, size_t... DN>
21: struct GetArray {
22: using type = std::array<typename GetArray<T, D2, DN...>::type, D1>;
23: };
24:
25: template <typename T, size_t D1, size_t D2>
26: struct GetArray<T, D1, D2> {
27: using type = std::array<std::array<T, D2>, D1>;
28: };
29:
30: template <typename T, size_t D1, size_t D2, size_t... DN>
31: using MDimArray = typename GetArray<T, D1, D2, DN...>::type;
32:
33: } // namespace gra
34:
35: #endif

```

```

./include/Granitti/MUserHistograms.h 1/1
1: // Container class for different type of histograms
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: //
6: #ifndef MUSERHISTOGRAMS_H
7: #define MUSERHISTOGRAMS_H
8:
9: // C++
10: #include <complex>
11: #include <vector>
12:
13: // Own
14: #include "Granitti/MBL.h"
15: #include "Granitti/MBR.h"
16: #include "Granitti/MBKinematics.h"
17:
18: namespace gra {
19: class MUserHistograms {
20: public:
21: // Constructors, destructor
22: MUserHistograms() {}
23: ~MUserHistograms() {}
24:
25: void InitHistograms();
26: void FillHistograms(double totalweight, const gra::LORENTZSCALAR &scalar);
27: void PrintHistograms();
28: void SetHistograms(unsigned int in) { HIST = in; }
29: void FillCosThetaPhi(double totalweight, const gra::LORENTZSCALAR &scalar);
30:
31: // Histograms indexed by name std::string
32: std::map<std::string, MBR<double>> h1;
33: std::map<std::string, MBR2> h2;
34:
35: unsigned int HIST = 0; // Histogramming level
36: };
37:
38: } // namespace gra
39:
40: #endif

```

```

./include/Granitti/MSpin.h 1/1
1: // Spin polarization and correlation functions
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: //
6: #ifndef MSPIN_H
7: #define MSPIN_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "Granitti/MBKinematics.h"
16: #include "Granitti/MRandom.h"
17:
18: namespace gra {
19: namespace spin {
20:
21: MMatrix<std::complex<double>> CalculateMMatrix(const MDecayBranch &branch);
22: void TreeRecursion(MDecayBranch &branch);
23: void TensorTree(const MDecayBranch &branch, MMatrix<std::complex<double>> &out);
24:
25: std::complex<double> ProdAmp(const gra::LORENTZSCALAR &ltts, gra::PARAM_RES &res);
26: std::complex<double> DecayAmp(gra::LORENTZSCALAR &ltts, gra::PARAM_RES &res);
27: void GetRotation(const gra::LORENTZSCALAR &ltts, const std::string &FRAME, double &theta_R,
28: double &phi_R);
29:
30: MMatrix<std::complex<double>> DMatrix(double J, double theta_mother, double phi_mother);
31: MMatrix<std::complex<double>> fMatrix(const MMatrix<std::complex<double>> &T, double J, double s1,
32: double s2, double theta, double phi);
33:
34: std::vector<double> SpinProjections(double J);
35: void InitMatrix(gra::HELMATRIX &hc, const gra::MParticle &p, const gra::MParticle &pl,
36: const gra::MParticle &pd);
37:
38: // Spin-Statistics
39: bool BoseFermi(int l, int s);
40: bool FermiBose(int l, int s);
41:
42: // Spin algebra functions
43: double ClebschJordan(double j1, double j2, double m1, double m2, double j, double m);
44: double CGrules(double j1, double j2, double m1, double m2, double j, double m);
45: bool ChairInt(double x);
46: bool iequal(double x, double y);
47: bool WJ(double j1, double j2, double j3, double m1, double m2, double m3);
48: bool WJrules(double j1, double j2, double j3, double m1, double m2, double m3);
49: double TriangleCoeff(double j1, double j2, double j3);
50: double WignerD(double theta, double phi, double m, double mp, double J);
51: double WignerSmall(double theta, double m, double mp, double J);
52:
53: // Density matrix functions
54: bool Positivity(const MMatrix<std::complex<double>> &rho, unsigned int J);
55: MMatrix<std::complex<double>> RandomBo(unsigned int J, bool parity, MRandom &rng);
56: double VonNeumannEntropy(const MMatrix<std::complex<double>> &rho);
57:
58: // namespace spin
59: } // namespace gra
60:
61: #endif

```

```

./include/Granitti/MMatrix.h 1/5
1: // Minimal templated matrix class [HEADER ONLY class]
2: //
3: //
4: // (c) 2017-2020 Mikael Mieskolainen
5: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
6: //
7: #ifndef MMATRIX_H
8: #define MMATRIX_H
9:
10: #include <initializer_list>
11: #include <memory>
12: #include <iostream>
13:
14: namespace gra {
15: // Constant size matrix
16: template <typename T>
17: class MMatrix {
18: public:
19: MMatrix(): rows(0), cols(0) { data = nullptr; }
20:
21: MMatrix(std::size_t r, std::size_t c) {
22: rows = r;
23: cols = c;
24: data = new T[rows * cols];
25: std::fill(data, data + rows * cols, T(0)); // No initialization
26: }
27:
28: MMatrix(std::size_t r, std::size_t c, T value) {
29: rows = r;
30: cols = c;
31: data = new T[rows * cols];
32: std::fill(data, data + rows * cols, T(value)); // Initialization
33: }
34: MMatrix(std::size_t r, std::size_t c, const std::string &special) {
35: rows = r;
36: cols = c;
37: data = new T[rows * cols];
38:
39: if (special == "eye") {
40: std::fill(data, data + rows * cols, T(0.0));
41: Identity();
42: } else if (special == "minkowski") {
43: std::fill(data, data + rows * cols, T(0.0));
44: Minkowski();
45: } else {
46: throw std::invalid_argument("MMatrix: Unknown initialization string: " + special);
47: }
48: }
49: // Set matrix to Identity (diagonal = 1, otherwise 0)
50: void Identity() {
51: for (std::size_t i = 0; i < rows; ++i)
52: for (std::size_t j = 0; j < cols; ++j) (this->operator()(i, j) = (i == j) ? 1.0 : 0.0; )
53: }
54:
55: // Set matrix to Minkowski metric
56: void Minkowski() {
57: for (std::size_t i = 0; i < rows; ++i) {
58: for (std::size_t j = 0; j < cols; ++j) {
59: double sign = (i > 0 && j > 0) ? -1.0 : 1.0;
60: this->operator()(i, j) = (i == j) ? sign * 1.0 : 0.0;
61: }
62: }
63: }
64:
65: // For initializing with a = (row-vector, row-vector, ...)
66: // where each row is std::vector<T>
67:
68: // Row-major order
69: MMatrix(std::initializer_list<std::vector<T>> list) {
70: rows = list.size();
71: cols = list.begin()->size();
72: data = new T[rows * cols];
73:
74: std::size_t i = 0;
75: for (const auto &v : list) {
76: for (std::size_t j = 0; j < cols; ++j) { data[cols * i + j] = v[j]; }
77: ++i;
78: }
79: }
80:
81: // For initializing with a = { (i, j), (i, j) }
82: // Row-major order!

```

```

./include/Granitti/MMatrix.h          2/5
84:  MMatrix(std::initializer_list<std::initializer_list<>> list) {
85:  rows = list.size();
86:  cols = list.begin()->size();
87:  data = new T[rows * cols];
88:
89:  std::size_t i = 0;
90:  for (const auto &v : list) {
91:  std::size_t j = 0;
92:  for (const auto &w : v) {
93:  data[cols * i + j] = w;
94:  ++j;
95:  }
96:  ++i;
97:  }
98:  }
99:
100: // Copy constructor
101: MMatrix(const MMatrix &a) {
102: rows = a.rows;
103: cols = a.cols;
104: data = new T[rows * cols];
105:
106: // Copy all elements
107: Copy(a);
108: }
109:
110: MMatrix() {
111: // delete dynamically allocated memory
112: delete[] data;
113: }
114:
115: // Assignment operator
116: MMatrix operator=(const MMatrix &rhs) {
117: if (data != rhs.data && rhs.data != nullptr) {
118: Resize(rhs.rows, rhs.cols);
119: Copy(rhs);
120: }
121: return *this;
122: }
123: // For indexing with [i][j]
124: // Row-major order!
125: operator[](const std::size_t row) { return data + cols * row; }
126: const T *operator[](const std::size_t row) const { return data + cols * row; }
127:
128: // Row-major order!
129: T operator()(std::size_t i, std::size_t j) {
130: if (i >= rows || j >= cols) {
131: throw std::out_of_range("MMatrix: Index over matrix dimensions!");
132: }
133: return data[cols * i + j];
134: }
135: const T operator()(std::size_t i, std::size_t j) const {
136: if (i >= rows || j >= cols) {
137: throw std::out_of_range("MMatrix: Index over matrix dimensions!");
138: }
139: return data[cols * i + j];
140: }
141:
142: // -----
143: // Add to the left.
144: MMatrix operator+(const MMatrix &rhs) {
145: for (std::size_t i = 0; i < rows; ++i) {
146: for (std::size_t j = 0; j < cols; ++j) { this->operator()(i, j) += rhs[i][j]; }
147: }
148: return *this;
149: }
150: // Subtract to the left
151: MMatrix operator-(const MMatrix &rhs) {
152: for (std::size_t i = 0; i < rows; ++i) {
153: for (std::size_t j = 0; j < cols; ++j) { this->operator()(i, j) -= rhs[i][j]; }
154: }
155: return *this;
156: }
157:
158: // -----
159: // Return negated matrix
160: MMatrix operator-() const {
161: MMatrix out(rows, cols);
162: for (std::size_t i = 0; i < rows; ++i) {
163: for (std::size_t j = 0; j < cols; ++j) { out[i][j] = -this->operator()(i, j); }
164: }
165: return out;
166: }
167:
168: // -----
169: // Frobenius norm
250: double FrobNorm() const { return std::sqrt(FrobNorm2()); }
251:
252: // Frobenius norm squared
253: double FrobNorm2() const {
254: double sum = 0.0;
255: for (std::size_t i = 0; i < rows; ++i) {
256: for (std::size_t j = 0; j < cols; ++j) {
257: const double value = std::abs(this->operator()(i, j));
258: sum += value * value; // |aij|2
259: }
260: }
261: return sum;
262: }
263:
264: // Get transposed matrix
265: MMatrix Transpose() const {
266: MMatrix out(cols, rows);
267: for (std::size_t i = 0; i < rows; ++i) {
268: for (std::size_t j = 0; j < cols; ++j) { out[j][i] = this->operator()(i, j); }
269: }
270: return out;
271: }
272:
273: // Get conjugate transposed matrix (dagger)
274: MMatrix ConjTranspose() const {
275: MMatrix out(cols, rows);
276: for (std::size_t i = 0; i < rows; ++i) {
277: for (std::size_t j = 0; j < cols; ++j) { out[j][i] = std::conj(this->operator()(i, j)); }
278: }
279: return out;
280: }
281:
282: // Get diagonal vector
283: std::vector<T> GetDiag() const {
284: if (rows != cols) {
285: throw std::invalid_argument("MMatrix: GetDiag: Only defined for square matrices, rows = " +
286: std::to_string(rows) + ", cols = " + std::to_string(cols));
287: }
288: std::vector<T> d(rows);
289: for (std::size_t i = 0; i < rows; ++i) { d[i] = this->operator()(i, i); }
290: return d;
291: }
292:
293: // Get trace
294: T Trace() const {
295: if (rows != cols) {
296: throw std::invalid_argument("MMatrix: Trace: Only defined for square matrices, rows = " +
297: std::to_string(rows) + ", cols = " + std::to_string(cols));
298: }
299: T sum = 0.0;
300: for (std::size_t i = 0; i < rows; ++i) { sum += this->operator()(i, i); }
301: return sum;
302: }
303:
304: // Print
305: void Print(const std::string &name = "") const {
306: std::cout << "MMatrix: Print: " << name << " [" << rows << " x " << cols << "]" << std::endl;
307: std::cout << std::setprecision(4);
308: for (std::size_t i = 0; i < rows; ++i) {
309: for (std::size_t j = 0; j < cols; ++j) { std::cout << this->operator()(i, j) << "\t"; }
310: std::cout << std::endl;
311: }
312: }
313:
314: // Size operators
315:

```

```

./include/Granitti/MMatrix.h          3/5
167:
168: // Add two matrices
169: MMatrix operator+(const MMatrix &rhs) const {
170: MMatrix out(rows, cols);
171: if (rows != rhs.size_row() || cols != rhs.size_col()) {
172: throw std::invalid_argument("MMatrix: Matrix + Matrix with invalid dimensions");
173: }
174:
175: for (std::size_t i = 0; i < rows; ++i) {
176: for (std::size_t j = 0; j < cols; ++j) { out[i][j] = this->operator()(i, j) + rhs[i][j]; }
177: }
178: return out;
179: }
180:
181: // Subtract two matrices
182: MMatrix operator-(const MMatrix &rhs) const {
183: MMatrix out(rows, cols);
184: if (rows != rhs.size_row() || cols != rhs.size_col()) {
185: throw std::invalid_argument("MMatrix: Matrix - Matrix with invalid dimensions");
186: }
187:
188: for (std::size_t i = 0; i < rows; ++i) {
189: for (std::size_t j = 0; j < cols; ++j) { out[i][j] = this->operator()(i, j) - rhs[i][j]; }
190: }
191: return out;
192: }
193:
194: // -----
195: // We do not allow addition or subtraction by scalar, for dimensional
196: // safety reasons (can lead to abstracted results), thus only multiplication
197: // operators
198:
199: // Multiply by a scalar
200: MMatrix operator*(const T &rhs) const {
201: MMatrix out(rows, cols);
202: for (std::size_t i = 0; i < rows; ++i) {
203: for (std::size_t j = 0; j < cols; ++j) { out[i][j] = this->operator()(i, j) * rhs; }
204: }
205: return out;
206: }
207:
208: // Divide by a scalar
209: MMatrix operator/(const T &rhs) const {
210: MMatrix out(rows, cols);
211: for (std::size_t i = 0; i < rows; ++i) {
212: for (std::size_t j = 0; j < cols; ++j) { out[i][j] = this->operator()(i, j) / rhs; }
213: }
214: return out;
215: }
216:
217: // -----
218: // Matrix [this] * Vector [rhs] multiplication
219: std::vector<T> operator*(const std::vector<T> &rhs) const {
220: if (this->size() != cols) {
221: throw std::invalid_argument(
222: "MMatrix: matrix * vector product with invalid dimension on rhs");
223: }
224:
225: std::vector<T> out(rows, 0.0); // Init!
226: for (std::size_t i = 0; i < rows; ++i) {
227: for (std::size_t j = 0; j < cols; ++j) { out[i] += this->operator()(i, j) * rhs[j]; }
228: }
229: return out;
230: }
231:
232: // Matrix [this] * Matrix [rhs] multiplication
233: MMatrix operator*(const MMatrix &rhs) const {
234: const std::size_t m = rows;
235: const std::size_t n = cols;
236: const std::size_t p = rhs.size_col();
237:
238: MMatrix<T> C(m, p, 0.0); // Note initialization to 0.0!
239:
240: if (cols != rhs.size_row()) {
241: throw std::invalid_argument("MMatrix: matrix * matrix product with invalid dimensions");
242: }
243:
244: // Over [i,j] of C
245: for (std::size_t i = 0; i < m; ++i) {
246: for (std::size_t j = 0; j < p; ++j) {
247: for (std::size_t k = 0; k < n; ++k) {
248: C[i][j] += this->operator()(i, k) * rhs[k][j]; // notice plus
249: }
250: }
251: }
252: }
253:
254: // Copy data from a to *this (after Resize)
255: void Copy(const MMatrix &a) {
256: T *p = data + rows * cols;
257: T *q = a.data + rows * cols;
258: while (p > data) { *p = *q; }
259: }
260:
261: // Re-Allocate
262: void ReSize(std::size_t r, std::size_t c) {
263: if (data != nullptr) { delete[] data; }
264: rows = r;
265: cols = c;
266: data = new T[rows * cols];
267: }
268:
269: std::size_t rows;
270: std::size_t cols;
271: T * data;
272:
273: // namespace gra
274:
275: #endif

```

```

./include/Granitti/MMatrix.h          4/5
250:
251: }
252: }
253: return C;
254: }
255:
256: // -----
257: // Sum all elements
258: T Sum() const {
259: T sum(0.0);
260: for (std::size_t i = 0; i < rows; ++i) {
261: for (std::size_t j = 0; j < cols; ++j) { sum += this->operator()(i, j); }
262: }
263: return sum;
264: }
265:
266: // Frobenius norm
267: double FrobNorm() const { return std::sqrt(FrobNorm2()); }
268:
269: // Frobenius norm squared
270: double FrobNorm2() const {
271: double sum = 0.0;
272: for (std::size_t i = 0; i < rows; ++i) {
273: for (std::size_t j = 0; j < cols; ++j) {
274: const double value = std::abs(this->operator()(i, j));
275: sum += value * value; // |aij|2
276: }
277: }
278: return sum;
279: }
280:
281: // Get transposed matrix
282: MMatrix Transpose() const {
283: MMatrix out(cols, rows);
284: for (std::size_t i = 0; i < rows; ++i) {
285: for (std::size_t j = 0; j < cols; ++j) { out[j][i] = this->operator()(i, j); }
286: }
287: return out;
288: }
289:
290: // Get conjugate transposed matrix (dagger)
291: MMatrix ConjTranspose() const {
292: MMatrix out(cols, rows);
293: for (std::size_t i = 0; i < rows; ++i) {
294: for (std::size_t j = 0; j < cols; ++j) { out[j][i] = std::conj(this->operator()(i, j)); }
295: }
296: return out;
297: }
298:
299: // Get diagonal vector
300: std::vector<T> GetDiag() const {
301: if (rows != cols) {
302: throw std::invalid_argument("MMatrix: GetDiag: Only defined for square matrices, rows = " +
303: std::to_string(rows) + ", cols = " + std::to_string(cols));
304: }
305: std::vector<T> d(rows);
306: for (std::size_t i = 0; i < rows; ++i) { d[i] = this->operator()(i, i); }
307: return d;
308: }
309:
310: // Get trace
311: T Trace() const {
312: if (rows != cols) {
313: throw std::invalid_argument("MMatrix: Trace: Only defined for square matrices, rows = " +
314: std::to_string(rows) + ", cols = " + std::to_string(cols));
315: }
316: T sum = 0.0;
317: for (std::size_t i = 0; i < rows; ++i) { sum += this->operator()(i, i); }
318: return sum;
319: }
320:
321: // Print
322: void Print(const std::string &name = "") const {
323: std::cout << "MMatrix: Print: " << name << " [" << rows << " x " << cols << "]" << std::endl;
324: std::cout << std::setprecision(4);
325: for (std::size_t i = 0; i < rows; ++i) {
326: for (std::size_t j = 0; j < cols; ++j) { std::cout << this->operator()(i, j) << "\t"; }
327: std::cout << std::endl;
328: }
329: }
330:
331: // Size operators
332:

```

```

./include/Granitti/MMatrix.h          5/5
333: std::size_t size_row() const { return rows; }
334: std::size_t size_col() const { return cols; }
335:
336: private:
337: void Copy(const MMatrix &a) {
338: T *p = data + rows * cols;
339: T *q = a.data + rows * cols;
340: while (p > data) { *p = *q; }
341: }
342:
343: // Re-Allocate
344: void ReSize(std::size_t r, std::size_t c) {
345: if (data != nullptr) { delete[] data; }
346: rows = r;
347: cols = c;
348: data = new T[rows * cols];
349: }
350:
351: std::size_t rows;
352: std::size_t cols;
353: T * data;
354:
355: // namespace gra
356:
357: #endif

```



```
./include/Granitti/MTransport.h 1/1
1: // Optimal Transport Class
2: //
3: //
4: // (c) 2017-2020 Mikael Mieskolainen
5: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
6:
7: #ifndef MTRANSPORT_H
8: #define MTRANSPORT_H
9:
10: // C++
11: #include <algorithm>
12: #include <complex>
13: #include <random>
14: #include <stdexcept>
15: #include <vector>
16:
17: // Eigen
18: #include <Eigen/Dense>
19:
20: // Own
21: #include "Granitti/MMatrix.h"
22: #include "Granitti/MFence.h"
23:
24: namespace gra {
25: namespace opt {
26:
27: void ConvKernel(std::size_t n, std::size_t m, double lambda, MMatrix<double>& K);
28: void GibbsKernel(double lambda, const MMatrix<double>& C, MMatrix<double>& K);
29:
30: double SinkHorn(MMatrix<double>& P, const MMatrix<double>& K, std::vector<double>& p,
31:               std::vector<double>& q, std::size_t iter);
32:
33: } // namespace opt
34: } // namespace gra
35:
36: #endif
```

```
./include/Granitti/MNeuroJacobian.h 2/6
84: // Sigmoid
85: inline dual sigmoid(dual z) { return 1.0 / (1.0 + exp(-z)); }
86: // Modified Sigmoid
87: inline dual logexp(dual z) {
88:   const double alpha = 25.0;
89:   return 1.0 / alpha * log(1.0 + exp(alpha * z)) / (1.0 + exp(alpha * (z - 1.0)));
90: }
91:
92: // Gaussian
93: dual gaussprob(const VectorXd& z, double mu, double sigma) {
94:   dual prod = 1.0;
95:   for (int i = 0; i < z.size(); ++i) {
96:     prod *= (1.0 / sqrt(2 * M_PI * sigma * sigma)) * exp(-0.5 * (z[i] - mu) * (z[i] - mu) / (2.0 * sigma * sigma));
97:   }
98:   return prod;
99: }
100:
101: // Feedforward Network Layer
102: class Layer {
103: public:
104:   Layer(int M, int N) {
105:     w = MatrixXd(M, N);
106:     b = VectorXd(M);
107:   }
108:   Layer() {}
109:   MatrixXd w;
110:   MatrixXd b;
111:   VectorXd x;
112: };
113:
114: // Parameter struct
115: struct NetParams {
116:   std::vector<Layer> L;
117:
118:   // Return number of parameters
119:   int size_param() {
120:     int k = 0;
121:     for (std::size_t l = 0; l < L.size(); ++l) { k += L[l].w.rows() * L[l].w.cols(); }
122:     for (std::size_t l = 0; l < L.size(); ++l) { k += L[l].b.size(); }
123:     return k;
124:   }
125:
126:   int D = 0; // Integral dimension
127:   dual alpha;
128: };
129:
130: // Global parameters
131: NetParams par;
132:
133: // Neural network function G: D^N -> D^M, D = [0 ... 1]
134: //
135: // This depends on global parameter struct p
136: //
137: //
138: VectorXd G_net(const VectorXd& x) {
139:   const double beta = 0.5;
140:
141:   // Network input
142:   VectorXd in(x.size());
143:   for (std::size_t i = 0; i < (unsigned int)x.size(); ++i) { in[i] = x[i]; }
144:
145:   // Network layers
146:   for (std::size_t l = 0; l < par.L.size() - 1; ++l) {
147:     VectorXd out(par.L[l].w.rows());
148:     for (std::size_t i = 0; i < (unsigned int)par.L[l].w.rows(); ++i) {
149:       for (std::size_t j = 0; j < (unsigned int)par.L[l].w.cols(); ++j) {
150:         out[i] += par.L[l].w(i, j) * in[j];
151:       }
152:       out[i] = hypertanh(out[i] + par.L[l].b[i]) * beta + (1.0 - beta) * out[i];
153:     }
154:     in = out; // Next layer input is this layer output
155:   }
156:
157:   // Output layer compression to [0,1] range
158:   const unsigned int last = par.L.size() - 1;
159:   VectorXd out(par.L[last].w.rows());
160:   for (std::size_t i = 0; i < (unsigned int)par.L[last].w.rows(); ++i) {
161:     for (std::size_t j = 0; j < (unsigned int)par.L[last].w.cols(); ++j) {
162:       out[i] += par.L[last].w(i, j) * in[j];
163:     }
164:     out[i] = logexp(out[i] + par.L[last].b[i]);
165:   }
166:   return out;
```

```
./include/Granitti/MNeuroJacobian.h 1/6
1: // GRANITTI - Monte Carlo event generator for high energy diffraction
2: // https://github.com/mieskolainen/granitti
3: //
4: // <NeuroJacobian neural net prototype>
5: //
6: //
7: // (c) 2017-2020 Mikael Mieskolainen
8: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
9:
10: #ifndef MNEUROJACOBIAN_H
11: #define MNEUROJACOBIAN_H
12:
13: // C++ includes
14: #include <future>
15: #include <iostream>
16:
17: // Own
18: #include <Granitti/MAux.h>
19: #include <Granitti/MRNG.h>
20: #include <Granitti/MMath.h>
21: #include <Granitti/MProcess.h>
22: #include <Granitti/MRandom.h>
23:
24: // Eigen includes
25: #include <Eigen/Core>
26:
27: // L-BFGS algorithms
28: #include <LBFGS/include/LBFGS.h>
29:
30: // autodiff library
31: #include <autodiff/forward.hpp>
32: #include <autodiff/forward/eigen.hpp>
33: #include <autodiff/reverse.hpp>
34: #include <autodiff/reverse/eigen.hpp>
35:
36: using namespace gra;
37: using gra::aux::indices;
38: using gra::math::PI;
39: using gra::math::pow2;
40:
41: using Eigen::VectorXd;
42: using namespace Eigen;
43: using namespace autodiff;
44: using namespace LBFGSpp;
45:
46: namespace gra {
47: namespace neurojac {
48: MRandom rand;
49: unsigned int BATCHSIZE = 0;
50: bool FLAT_TARGET = false;
51:
52: // declare a pointer to member function
53: MProcess *pprocess;
54: double (MProcess::*ptfptr)(const std::vector<double>& randvec,
55:                          AuxInData & aux) = &MProcess::EventWeight;
56:
57: // Function to be integrated
58: double Func(std::vector<double>& x) {
59:   if (FLAT_TARGET) { return 1.0; }
60:
61:   gra::AuxInData aux;
62:   aux.vegasweight = 1.0;
63:   aux.burn_in_mode = false;
64:   double y = (pprocess->neurojac)(*graineurojac::ptfptr)(x, aux);
65:
66:   // Numerical protection
67:   if (std::isnan(y) || std::isinf(y) || y < 0.0) {
68:     return y;
69:   }
70:
71:   // Exponential linear unit
72:   inline dual elu(dual z) {
73:     if (z >= 0.0) {
74:       return z;
75:     } else {
76:       return exp(z) - 1.0;
77:     }
78:   }
79:
80:   // Hyperbolic Functions
81:   inline dual hyperbolic(dual z) { return (exp(z) - exp(-z)) / 2; }
82:   inline dual hyperbolic(dual z) { return (exp(2.0 * z) - 1.0) / (exp(2.0 * z) + 1.0); }
83:
167: }
```

```
./include/Granitti/MNeuroJacobian.h 3/6
167: }
168:
169: class MNeuroJacobian {
170: public:
171:   MNeuroJacobian() {
172:     MNeuroJacobian() {}
173:
174:   // Vector format to parameters
175:   static void VecPar(const std::vector<double>& w, NetParams par) {
176:     std::size_t k = 0;
177:     for (std::size_t l = 0; l < par.L.size(); ++l) {
178:       for (std::size_t i = 0; i < (unsigned int)par.L[l].w.rows(); ++i) {
179:         for (std::size_t j = 0; j < (unsigned int)par.L[l].w.cols(); ++j) {
180:           par.L[l].w(i, j) = w[k];
181:           ++k;
182:         }
183:       }
184:     }
185:
186:     for (std::size_t l = 0; l < par.L.size(); ++l) {
187:       for (std::size_t i = 0; i < (unsigned int)par.L[l].b.size(); ++i) {
188:         par.L[l].b[i] = w[k];
189:         ++k;
190:       }
191:     }
192:   }
193:
194:   // Random init for parameter vector
195:   std::vector<double> RandInit(NetParams &par, double sigma) {
196:     std::vector<double> w(par.size_param());
197:     unsigned int k = 0;
198:
199:     // Connection matrices with random initialization
200:     for (std::size_t l = 0; l < par.L.size(); ++l) {
201:       for (std::size_t i = 0; i < (unsigned int)par.L[l].w.rows(); ++i) {
202:         for (std::size_t j = 0; j < (unsigned int)par.L[l].w.cols(); ++j) {
203:           w[k] = randn(0.0, sigma) * sqrt(2.0 / par.L[l].w.cols()); // Heur all style
204:           ++k;
205:         }
206:       }
207:     }
208:
209:     // Bias initialized to zero
210:     for (std::size_t l = 0; l < par.L.size(); ++l) {
211:       for (std::size_t i = 0; i < (unsigned int)par.L[l].b.size(); ++i) {
212:         w[k] = 0.0;
213:         ++k;
214:       }
215:     }
216:     return w;
217:   }
218:
219:   // Kullback-Leibler divergence loss function
220:
221:   // For KL-divergence and algorithms with Jacobians, see:
222:   // [REFERENCE: Leow et al., Statistical Properties of Jacobian Maps ..., 2007]
223:   // http://www.sop.inria.fr/asclopos/cours/MVA/Module2/Papers/Leow_TMI07_LogUnbiased.pdf
224:
225:   // For neural networks and Jacobians, see:
226:   // [REFERENCE: Dillon et al., Tensorflow Bijector API,
227:   // https://arxiv.org/abs/1711.10365v1]
228:
229:   // Matrix derivative results, see:
230:   // [REFERENCE: M. Giles, http://eprints.maths.ox.ac.uk/1079/1/NA-08-01.pdf]
231:
232:   // MC Integration:
233:   // [REFERENCE: J. Bendavid, https://arxiv.org/pdf/1707.00028.pdf]
234:
235:   static double singleLoss(VectorXd& z, const NetParams &par) {
236:     // Input likelihood value
237:     double p = val(gaussprob(z, 0, 1));
238:
239:     // Evaluate network
240:     VectorXd u = G_net(z);
241:
242:     // Jacobian matrix du/dz (AUTODIFF)
243:     MatrixXd J = jacobian(G_net, u, z);
244:
245:     // Absolute Jacobian determinant
246:     double absDet = abs(val(J.determinant()));
247:
248:     // =====
249: }
```



```

./include/Granitti/MNeuroJacobian.h      4/6
250: // Evaluate integrand function
251: std::vector<double> u(u.size());
252: for (std::size_t i = 0; i < (unsigned int)u.size(); ++i) { u[i] = val(u[i]); }
253: double fG = func(u);
254: // =====
255: // Prior term
256: double KL = 0.0;
257: if (p > 0) { KL = log(p); }
258: // Jacobian determinant term
259: if (absDetJ > 0) { KL = KL - log(absDetJ); }
260: // Target function fidelity term
261: if (fG > 0) { KL = KL - log(fG); }
262: return KL;
263: }
264: static dual loss(VectorXd& z, const NetParams& par) {
265:   VectorXd x(par.D);
266:   dual sumloss = 0.0;
267:   // Loop over all samples
268:   int k = 0;
269:   while (true) {
270:     // Pick prior p(z) distribution samples
271:     for (std::size_t i = 0; i < (unsigned int)par.D; ++i) {
272:       z[i] = Z[k];
273:       ++k;
274:     }
275:     sumloss = sumloss + singleloss(z, par);
276:     if (k == Z.size()) break;
277:   }
278:   return sumloss;
279: }
280: // Evaluate loss and gradient vector
281: static double CostGrad(const std::vector<double>& w, std::vector<double>& gradv) {
282:   // Generate new virtual batch from noise distribution (prior)
283:   VectorXd Z(BATCHSIZE * par.D);
284:   for (std::size_t i = 0; i < (unsigned int)Z.size(); ++i) { Z[i] = randx.G(0, 1); }
285:   // Update vector to global parameters
286:   VecPar w(par);
287:   // =====
288:   // Evaluate gradient vector components numerically (not very optimal..)
289:   std::size_t k = 0;
290:   gradv.resize(w.size());
291:   // Numerical derivative epsilon
292:   const double h = 1e-4;
293:   for (std::size_t l = 0; l < par.L.size(); ++l) {
294:     for (std::size_t j = 0; j < (unsigned int)par.L[l].w.rows(); ++j) {
295:       for (std::size_t i = 0; i < (unsigned int)par.L[l].w.cols(); ++i) {
296:         par.L[l].w(i, j) = w[k] + h;
297:         const double f_pos = val(loss(z, par));
298:         par.L[l].w(i, j) = w[k] - h;
299:         const double f_neg = val(loss(z, par));
300:         gradv[k] = (f_pos - f_neg) / (2 * h);
301:         par.L[l].w(i, j) = w[k]; // back to previous value
302:         ++k;
303:       }
304:     }
305:   }
306:   for (std::size_t l = 0; l < par.L.size(); ++l) {
307:     for (std::size_t j = 0; j < (unsigned int)par.L[l].b.size(); ++j) {
308:       par.L[l].b(j) = w[k] + h;
309:       const double f_pos = val(loss(z, par));
310:       par.L[l].b(j) = w[k] - h;
311:       const double f_neg = val(loss(z, par));
312:       gradv[k] = (f_pos - f_neg) / (2 * h);
313:       par.L[l].b(j) = w[k]; // back to previous value
314:       ++k;
315:     }
316:   }
317: }
318: }
319: }
320: }
321: }
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323: }
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329: }
330: }
331: }
332: }

```

```

./include/Granitti/MNeuroJacobian.h      6/6
416: std::cout << "z(x) = " << fx << std::endl;
417: }
418: }
419: private:
420: class Neurocost {
421: public:
422:   Neurocost(int n) { n = n; }
423: }
424: double fx = 0.0;
425: double operator()(const VectorXd& x, VectorXd& gradv) {
426:   // Map to vectors
427:   std::vector<double> w(x.size(), 0.0);
428:   std::vector<double> gradv(w.size(), 0.0);
429:   for (std::size_t i = 0; i < (unsigned int)x.size(); ++i) { w[i] = x[i]; }
430:   fx = CostGrad(w, gradv);
431:   for (std::size_t i = 0; i < (unsigned int)x.size(); ++i) { gradv[i] = gradv[i]; }
432:   return fx;
433: }
434: }
435: }
436: }
437: }
438: }
439: private:
440: int n;
441: };
442: };
443: }; // MNeuroJacobian
444: }; // Namespace neurojac
445: }; // Namespace gra
446: }; // Namespace gra
447: };
448: #endif

```

```

./include/Granitti/MNeuroJacobian.h      5/6
333: const double cost = val(loss(z, par));
334: std::cout << "loss = " << cost << std::endl;
335: return cost;
336: }
337: }
338: // Naive gradient descent
339: void NaiveGradDescent(std::vector<double>& w, unsigned int MAXITER, double rate) {
340: // =====
341: // Gradient Descent
342: // w_n+1 = w_n - gamma*GradF(w_n)
343: unsigned int iter = 0;
344: // gradient vector
345: std::vector<double> gradv(w.size(), 0.0);
346: while (true) {
347:   CostGrad(w, gradv);
348:   // Naive gradient descent update
349:   for (std::size_t k = 0; k < gradv.size(); ++k) {
350:     w[k] = w[k] - rate * gradv[k];
351:     printf("w[%3d] = %6.3f <gradv[%3d] = %10.3E> \n", k, w[k], k, gradv[k]);
352:   }
353:   ++iter;
354: }
355: }
356: // if (iter > MAXITER) break;
357: }
358: }
359: }
360: }
361: // Test network dimensions
362: void TestNetworkDimensions() {
363:   for (std::size_t l = 0; l < par.L.size() - 1; ++l) {
364:     if (par.L[l + 1].w.cols() != par.L[l].w.rows()) {
365:       throw std::invalid_argument("Network layer dimension mismatch!");
366:     }
367:   }
368: }
369: }
370: // Optimize network
371: void Optimize() {
372: // =====
373: // Random init (too high sigma -> may get trapped to a bad local minima)
374: double sigma = 1e-4;
375: std::vector<double> w = RandomInit(par, sigma);
376: // =====
377: // By Jacobi formula: d/dt ln det A(t) = tr(A(t)^{-1} d/dt A)
378: // =====
379: }
380: }
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```

```

./include/Granitti/MAux.h                 1/5
1: // I/O aux functions
2: }
3: // (c) 2017-2020 Mikael Hieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: }
6: #ifndef MAUX_H
7: #define MAUX_H
8: }
9: // C++
10: #include <limits>
11: #include <unistd.h>
12: #include <complex>
13: #include <mutex>
14: #include <random>
15: #include <regex>
16: #include <sstream>
17: #include <stdexcept>
18: #include <vector>
19: // #include <experimental/filesystem>
20: }
21: // Own
22: #include "Granitti/M4Vec.h"
23: #include "Granitti/MMatrix.h"
24: }
25: // LAPDP
26: #include "LAPDP/LAPDPF.h"
27: }
28: // HepMC3
29: #include "HepMC3/FourVector.h"
30: }
31: // Eigen
32: #include <Eigen/Dense>
33: }
34: // Libraries
35: #include "rang.hpp"
36: }
37: namespace gra {
38: namespace aux {
39: // Example:
40: //
41: // @SOMECOMMAND: value
42: // @PFG[992][M300, W:0]
43: //
44: // Format options:
45: //
46: // @id: value is represented with map<"_SINGLETON_", value>
47: //
48: // @id: value
49: // @id: arg
50: // @id: [arg]
51: // @id: [target1, target2, ...] [arg]
52: }
53: struct OneCMD {
54:   std::string id;
55:   std::vector<std::string> target;
56:   std::map<std::string, std::string> arg;
57: };
58: void Print() {
59:   std::cout << "id : " << id << std::endl;
60: }
61: std::cout << "[target] : [";
62: for (std::size_t i = 0; i < target.size(); ++i) {
63:   std::cout << target[i];
64:   if (i < target.size() - 1) { std::cout << ", "; }
65: }
66: std::cout << "]" << std::endl;
67: }
68: std::cout << "arg : " << std::endl;
69: for (const auto& arg) { std::cout << x.first << " : " << x.second << std::endl; }
70: std::cout << std::endl;
71: }
72: }
73: }
74: // M4Vec to HepMC3::FourVector
75: inline HepMC3::FourVector M4Vec2HepMC3(const M4Vec& v) {
76:   return HepMC3::FourVector(v.X(), v.Y(), v.Z(), v.E());
77: }
78: // HepMC3::FourVector to M4Vec
79: inline M4Vec HepMC3M4Vec(const HepMC3::FourVector& v) { return M4Vec(v.X(), v.Y(), v.Z(), v.E()); }
80: }
81: // Eigen to std::vector
82: inline std::vector<double> EigenVector(const Eigen::VectorXd& x) {
83:   std::vector<double> y(x.size());

```

```

./include/Granitti/MAux.h          2/5
84:  for (int i = 0; i < x.size(); ++i) { y[i] = x[i]; }
85:  return y;
86:  }
87:  }
88:  // std::vector to Eigen
89:  template <typename T>
90:  inline Eigen::VectorX<T> Eigen(const std::vector<T> &x) {
91:  Eigen::VectorX<T> y(x.size());
92:  for (std::size_t i = 0; i < x.size(); ++i) { y[i] = x[i]; }
93:  return y;
94:  }
95:  // Matrix to Eigen matrix
96:  template <typename T>
97:  inline Eigen::MatrixX<T> Eigen(const Matrix<T> &m) {
98:  Eigen::MatrixX<T> mEigen(m.size_row(), m.size_col());
99:  for (std::size_t i = 0; i < m.size_row(); ++i) {
100:  for (std::size_t j = 0; j < m.size_col(); ++j) { mEigen(i, j) = M[i][j]; }
101:  }
102:  return mEigen;
103:  }
104:  }
105:  // System information
106:  void PrintKey(int argc, char *argv[]);
107:  void AutoDownloadLDAP(const std::string pdfname);
108:  std::string ExecCommand(const std::string &cmd);
109:  std::string GetExecutablePath();
110:  std::string GetBasePath(std::size_t level);
111:  //
112:  //
113:  //
114:  std::string GetCurrentPath();
115:  bool FileExists(const std::experimental::filesystem::path &p,
116:  std::experimental::filesystem::file_status &s =
117:  std::experimental::filesystem::file_status());
118:  //
119:  //
120:  std::uintmax_t GetFileSize(const std::string filename);
121:  void GetProcessMemory(double &peak_use, double &resident_use);
122:  void GetDiskUsage(const std::string &path, int64_t &size, int64_t &free, int64_t &used);
123:  unsigned long long TotalSystemMemory();
124:  std::string SystemName();
125:  std::string HostName();
126:  const std::string DateTime();
127:  //
128:  // Progress bar
129:  void PrintProgress(double ratio);
130:  void ClearProgress();
131:  //
132:  // djb2hash function
133:  unsigned long djb2hash(const std::string &s);
134:  //
135:  // Simple CSV reader
136:  void ReadCSV(const std::string inputfile, std::vector<std::vector<std::string>> &output);
137:  //
138:  // Input processing
139:  std::string GetInputData(const std::string inputfile);
140:  //
141:  bool IsIntegerDigits(const std::string &str);
142:  //
143:  // Trim extra spaces of a string
144:  void TrimExtraSpaces(std::string &value);
145:  void TrimLeadSpace(std::string &value);
146:  void TrimTrailingSpace(std::string &value);
147:  void TrimWhiteSpace(std::string &value);
148:  void TrimAllSpace(std::string &value);
149:  //
150:  // Number to string with formatting
151:  template <typename T>
152:  std::string ToString(const T &value, const unsigned int n = 6) {
153:  std::stringstream out;
154:  out.precision(n);
155:  out << std::fixed << value;
156:  return out.str();
157:  }
158:  //
159:  // Quantum numbers as a string
160:  std::string ParityToString(int &value);
161:  std::string ChargeXtoString(int &q);
162:  std::string SpinXtoString(int &Jz);
163:  //
164:  // String splitting
165:  std::vector<std::string> SplitStr2Str(const std::string input, const char delim = ',');

```

```

./include/Granitti/MAux.h          4/5
250:  if (cut.size() != 2) {
251:  throw std::invalid_argument("AssertCutRange: Input '" + name + "' vector size not 2");
252:  }
253:  if (!AssertCut(cut, name, dothrow)) { return false; }
254:  //
255:  if (cut[0] < bounds[0] || cut[1] > bounds[1]) {
256:  if (dothrow) {
257:  std::stringstream message = "AssertCutRange: Input '" + name + "' with [" + std::ito_string(cut[0]) +
258:  ", " + std::ito_string(cut[1]) + "]" + " invalid given bounds: [" +
259:  std::ito_string(bounds[0]) + ", " + std::ito_string(bounds[1]) + "]" ;
260:  throw std::invalid_argument(message);
261:  }
262:  return false;
263:  }
264:  return true;
265:  }
266:  //
267:  // Assert compare function, threshold 0.01 means 1 percent accuracy
268:  template <typename T>
269:  bool AssertRatio(T value, T reference, T threshold, const std::string &name = "") {
270:  bool ok = false;
271:  const T ratio = value / reference;
272:  if (ratio < (1.0 - threshold) && ratio > (1.0 + threshold)) { ok = true; }
273:  if (ok && dothrow) {
274:  throw std::invalid_argument("AssertRatio: Input '" + name + "' = " + std::ito_string(value) +
275:  " not within reference = " + std::ito_string(reference) +
276:  " under threshold = " + std::ito_string(threshold));
277:  }
278:  }
279:  return ok;
280:  }
281:  //
282:  // Assert range [a,b]
283:  template <typename T>
284:  bool AssertRange(T value, std::vector<T> range, const std::string &name = "",
285:  bool dothrow = false) {
286:  if (range.size() != 2) {
287:  throw std::invalid_argument("AssertRange: Input '" + name + "' , range vector size is not 2!");
288:  }
289:  bool ok = false;
290:  if (value >= range[0] && value <= range[1]) { ok = true; }
291:  if (ok && dothrow) {
292:  throw std::invalid_argument("AssertRange: Input '" + name + "' = " + std::ito_string(value) +
293:  " out of range [" + std::ito_string(range[0]) + ", " +
294:  std::ito_string(range[1]) + "]" );
295:  }
296:  return ok;
297:  }
298:  //
299:  // Assert value if found from a set of numbers
300:  template <typename T>
301:  bool AssertSet(T value, std::vector<T> set, const std::string &name = "", bool dothrow = false) {
302:  bool ok = false;
303:  for (const auto &i : set) {
304:  if (value == i) {
305:  ok = true;
306:  break;
307:  }
308:  }
309:  if (ok && dothrow) {
310:  throw std::invalid_argument("AssertSet: Input '" + name + "' = " + std::ito_string(value) +
311:  " not found from the input set");
312:  }
313:  return ok;
314:  }
315:  //
316:  // -----
317:  // Templates for enhanced index based looping
318:  // for (const auto &i : indices(whatever) { vector[i] = foo; ... }
319:  //
320:  template <typename T>
321:  struct index_range {
322:  struct iterator {
323:  bool operator==(iterator &x) const { return index == x.index; }
324:  iterator operator++() {
325:  ++index;
326:  return *this;
327:  }
328:  T operator*() const { return index; }
329:  T index;
330:  };
331:  };
332:  }

```

```

./include/Granitti/MAux.h          3/5
167:  bool trimextraspaces = true);
168:  std::vector<int> SplitStr2Int(const std::string input, const char delim = ',');
169:  std::vector<std::string> Extract(const std::string &str);
170:  //
171:  // Split string to int or double
172:  template <class T>
173:  std::vector<T> SplitStr(const std::string input, T type, const char delim = ',') {
174:  std::vector<T> output;
175:  std::stringstream ss(input);
176:  //
177:  // Get inputfiles by comma
178:  while (ss.good()) {
179:  std::string substr;
180:  std::getline(ss, substr, delim);
181:  TrimExtraSpace(substr);
182:  TrimExtraSpace(substr);
183:  //
184:  // Detect type >>
185:  // int
186:  if (std::is_same<T, int>::value) {
187:  output.push_back(std::stoi(substr));
188:  }
189:  // double
190:  else if (std::is_same<T, double>::value) {
191:  output.push_back(std::stod(substr));
192:  }
193:  }
194:  return output;
195:  }
196:  //
197:  // Check if file exists
198:  bool FileExists(const std::string &name);
199:  //
200:  void PrintStatus();
201:  void PrintWarning();
202:  void PrintGameOver();
203:  void PrintFlashScreen(rang::rg_color);
204:  void PrintVersion();
205:  //
206:  // Bar print
207:  void PrintBar(const std::string &str, unsigned int N = 74);
208:  //
209:  // Create directory
210:  void CreateDirectory(const std::string &fulpath);
211:  //
212:  // Version information
213:  void CheckUpdate();
214:  void CreateVersionJSON();
215:  //
216:  double GetVersion();
217:  std::string GetVersionType();
218:  std::string GetVersionDate();
219:  std::string GetVersionPath();
220:  std::string GetVersionString();
221:  std::string GetVersionText();
222:  std::string GetVersionDate();
223:  //
224:  // Assert functions
225:  //
226:  //
227:  // Check out, lower value needs to be smaller than upper value
228:  template <typename T>
229:  bool AssertCut(std::vector<T> cut, const std::string &name = "", bool dothrow = false) {
230:  if (cut.size() != 2) {
231:  throw std::invalid_argument("AssertCut: Input '" + name + "' vector size not 2");
232:  }
233:  if (cut[1] <= cut[0]) {
234:  if (dothrow) {
235:  std::stringstream message = "AssertCut: Input '" + name + "' with [" + std::ito_string(cut[0]) +
236:  ", " + std::ito_string(cut[1]) +
237:  "]" (maximum value smaller than minimum value);
238:  throw std::invalid_argument(message);
239:  }
240:  return false;
241:  }
242:  return true;
243:  }
244:  //
245:  // Check cut obeys given boundaries
246:  template <typename T>
247:  bool AssertOutRange(std::vector<T> cut, std::vector<T> bounds, const std::string &name = "",
248:  bool dothrow = false) {

```

```

./include/Granitti/MAux.h          5/5
333:  iterator begin() const { return 0; }
334:  iterator end() const { return (n); }
335:  //
336:  T n;
337:  //
338:  //
339:  template <typename T, typename Index = typename T::size_type>
340:  index_range<Index> indices(const T &container) {
341:  return {container.size()};
342:  }
343:  // -----
344:  //
345:  std::vector<std::string> SplitCommands(const std::string &fullstr);
346:  std::vector<std::size_t> FindOccurrences(const std::string &str, const std::string &sub);
347:  //
348:  // namespace aux
349:  // namespace gra
350:  //
351:  #endif

```

```
./include/Granitti/MFFT.h 1/1
1: // Simple in-place Fast Fourier Transform with radix-2 using std::valarray
2: //
3: // Normalization is 1 to 1, that is: fft(fft(x)) = x
4: //
5: //
6: // (c) 2017-2020 Mikael Mieskolainen
7: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
8:
9: #ifndef MFFT_H
10: #define MFFT_H
11:
12: // C++
13: #include <complex>
14: #include <valarray>
15:
16: namespace gra {
17: #ifndef M_PI
18: #define M_PI 3.141592653589793238462643383279502884197169399375105820974944L
19: #endif
20:
21: namespace MFFT {
22: // In-place Cooley's 200/123 Tukey FFT
23: template <typename T>
24: void fft(std::valarray<std::complex<T>> &x) {
25:     const std::size_t N = x.size();
26:     if (N <= 1) return; // Trivial case/recursion ends, X = x
27:     if ((N & (N - 1)) != 0) {
28:         throw std::invalid_argument("ERROR: MFFT: fft: Input x.size() = " + std::ito_string(N) +
29:             " not a power of 2!");
30:     }
31:
32:     // Radix-2 step
33:     std::valarray<std::complex<T>> E = x[std::slice(0, N / 2, 2)];
34:     std::valarray<std::complex<T>> O = x[std::slice(1, N / 2, 2)];
35:
36:     // Even and odd part via recursion
37:     MFFT::fft(E);
38:     MFFT::fft(O);
39:
40:     for (std::size_t k = 0; k < N / 2; ++k) {
41:         const std::complex<T> t = std::exp(std::complex<T>(0, -2.0 * M_PI * k / N));
42:         x[k] = E[k] + O[k] * t;
43:         x[k + N / 2] = E[k] - O[k] * t;
44:     }
45: }
46:
47: // In-place Cooley-Tukey IFFT
48: template <typename T>
49: void ifft(std::valarray<std::complex<T>> &x) {
50:     x = x.apply(std::conj); // Conjugate
51:     MFFT::fft(x); // FFT
52:     x = x.apply(std::conj); // Conjugate
53:     x /= x.size(); // Normalize
54: }
55:
56: } // namespace MFFT
57:
58: } // namespace gra
59: #endif
```

```
./include/Granitti/MGraniitti.h 2/4
84:
85: // Weight statistics FLAT MC
86: double Wsum = 0.0;
87: double W2sum = 0.0;
88: double maxW = 0.0;
89:
90: // Weight statistics VEGAS MC
91: double maxF = 0.0;
92: double chi2 = 0.0;
93: };
94:
95: class MGraniitti {
96: public:
97: // Destructor & Constructor
98: MGraniitti();
99: ~MGraniitti();
100:
101: // For cross-section for HepMC3 output
102: void ForSec(double &x) { xForceSec = x; }
103:
104: // Read parameters
105: void ReadParam(const std::string &inputfile, const std::string &cmd_PROCESS = "null");
106:
107: // Initialize memory
108: void InitProcessMemory(std::string process, unsigned int NRESEED);
109: void InitMultiMemory();
110:
111: // Set simple MC parameters
112: void SetMCParam(MCFANUM &in);
113:
114: // Set VEGAS parameters
115: void SetVegasParam(const VEGASPARAM &in);
116:
117: // Set external file handle
118: void SetHepMC3Output(std::shared_ptr<HepMC3::WriterAscii<HepMC2>> &hepmc,
119:     const std::string & OUTPUTNAME) {
120:     OUTPUT = OUTPUTNAME;
121:     FORMAT = "hepmc2";
122:     outputHepMC2 = hepmc;
123: }
124: void SetHepMC3Output(std::shared_ptr<HepMC3::WriterAscii> &hepmc, const std::string &OUTPUTNAME) {
125:     OUTPUT = OUTPUTNAME;
126:     FORMAT = "hepmc3";
127:     outputHepMC3 = hepmc;
128: }
129:
130: // Maximum weight set/get
131: void SetMaxWeight(double w);
132: double GetMaxWeight() const;
133:
134: // Get number of CPU cores
135: void SetCores(int N) {
136:     CORES = N;
137:     if (CORES == 0) { // SETUP number of threads automatically
138:         CORES = std::round(std::thread::hardware_concurrency() * 1.5);
139:     }
140:     // If autodetection fails, set 1
141:     if (CORES < 1) { CORES = 1; }
142: }
143:
144: if (CORES < 0) {
145:     std::string str = "MGraniitti::SetCORES: CORES = 0";
146:     throw std::invalid_argument(str);
147: }
148:
149: int GetCores() const { return CORES; }
150: void SetWeight(const std::string &integrator) { INTEGRATOR = integrator; }
151: void SetWeighted(bool weighted) { WEIGHED = weighted; }
152:
153: // Output file name
154: void SetOutput(const std::string &output) { OUTPUT = output; }
155: // Output file format
156: void SetFormat(const std::string &format) {
157:     if (format == "hepmc3" || format == "hepmc2" || format == "hepvt") {
158:         FORMAT = format;
159:     } else {
160:         throw std::invalid_argument("MGraniitti::SetFormat: Unknown output format: " + format +
161:             " (valid: hepmc3, hepmc2, hepvt)");
162:     }
163: }
164: // Special method for combining histograms from each thread
165: void HistogramFusion();
166:
167: // Get cross section and error
```

```
./include/Granitti/MGraniitti.h 1/4
1: // GRANITTI Monte Carlo main class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: //
6: #ifndef MGRANITTI_H
7: #define MGRANITTI_H
8:
9: // C++
10: #include <complex>
11: #include <mutex>
12: #include <random>
13: #include <thread>
14: #include <vector>
15:
16: // HepMC3
17: #include "HepMC3/GenEvent.h"
18: #include "HepMC3/WriterAscii.h"
19: #include "HepMC3/WriterAsciiHepMC2.h"
20: #include "HepMC3/WriterHEPEVT.h"
21:
22: // Own
23: #include "Granitti/MVeg.h"
24: #include "Granitti/MHus.h"
25: #include "Granitti/MContinuum.h"
26: #include "Granitti/MEikonal.h"
27: #include "Granitti/MFactized.h"
28: #include "Granitti/MEKinematics.h"
29: #include "Granitti/MMatrix.h"
30: #include "Granitti/MParton.h"
31: #include "Granitti/MProcess.h"
32: #include "Granitti/MQuasiElastic.h"
33: #include "Granitti/MBjrm.h"
34: #include "Granitti/MWim.h"
35: #include "Granitti/MWeg.h"
36:
37: namespace gra {
38: // Simple MC parameters
39: struct MCFANUM {
40:     double PRECISION = 0.05; // Integral relative precision
41:     unsigned int MIN_EVENTS = 100000; // Minimum number of events to be sampled
42: };
43:
44: // Integration statistics
45: class Stats {
46: public:
47:     void Accumulate(const gra::AuxInData &aux) {
48:         evaluations += 1.0;
49:
50:         amplitude_ok += (aux.amplitude_ok ? 1.0 : 0.0);
51:         kinematics_ok += (aux.kinematics_ok ? 1.0 : 0.0);
52:         fidcuts_ok += (aux.fidcuts_ok ? 1.0 : 0.0);
53:         vetocuts_ok += (aux.vetocuts_ok ? 1.0 : 0.0);
54:     }
55:
56:     // Calculate integrated cross section sigma and its
57:     // error for direct (simple) sampling. NOT TO BE USED WITH VEGAS!
58:     void CalculateCrossSection() {
59:         // <<
60:         sigma = Wsum / evaluations;
61:         // err^2 = <f^2> - <f>^2
62:         sigma_err2 = W2sum / evaluations - gra::math::pow(sigma);
63:         // err_trial = err^2 / sqrt(trials) (standard error of the mean)
64:         sigma_err = gra::math::msqrt(sigma_err2 / evaluations);
65:     }
66:
67:     double amplitude_ok = 0.0;
68:     double kinematics_ok = 0.0;
69:     double fidcuts_ok = 0.0;
70:     double vetocuts_ok = 0.0;
71:
72:     // Keep as double to avoid overflow of range
73:     double evaluations = 0.0; // Integrand evaluations
74:     double trials = 0.0; // Event generation trials
75:
76:     unsigned int generated = 0.0; // Event generation
77:     unsigned int n_overflow = 0.0; // Weight overflows
78:
79:     // Cross section and its error
80:     double sigma = 0.0;
81:     double sigma_err = 0.0;
82:     double sigma_err2 = 0.0;
83:
84:     double amplitude_ok = 0.0;
85:     double kinematics_ok = 0.0;
86:     double fidcuts_ok = 0.0;
87:     double vetocuts_ok = 0.0;
88:
89:     // Set number of (un)weighted events to be generated
90:     void SetNumberOfEvents(int n) {
91:         if (n <= 0) {
92:             throw std::invalid_argument(
93:                 "MGraniitti::SetNumberOfEvents(): Error: Number "
94:                 "of events < 0!");
95:         }
96:         NEVENTS = n;
97:     }
98:
99:     int GetNumberOfEvents() const { return NEVENTS; }
100:
101: // Return processes
102: std::vector<std::string> GetProcessNumbers() const;
103:
104: // Initialize generator
105: void Initialize();
106: void Initialize(const MEikonal &eikonal_in);
107:
108: // Process object pointer, public so methods can be accessed
109: MProcess *proc = nullptr;
110:
111: void PrintHistograms();
112: void ReadGeneralParam(const std::string &inputfile);
113: void ReadProcessParam(const std::string &inputfile, const std::string &cmd_PROCESS = "null");
114: void ReadIntegParam(const std::string &inputfile);
115: void ReadReadOuts(const std::string &inputfile);
116: void ReadReadIn(const std::string &inputfile);
117: void ReadVetoCuts(const std::string &inputfile);
118: void ReadModelParam(const std::string &inputfile) const;
119:
120: void Generate();
121:
122: std::string PROCESS = "null"; // Physics process identifier
123:
124: bool WEIGHTED = false; // Unweighted or weighted event generation
125: int NEVENTS = 0; // Number of events to be generated
126: int CORES = 0; // Number of CPU cores (threads) in use
127: std::string INTEGRATOR = "null"; // Integrator (VEGAS, Flat, ...)
128:
129: // SILENT OUTPUT
130: bool HILJAA = false;
131:
132: private:
133: // -----
134: // Copy and assignment disabled by making the private
135: MGraniitti(const MGraniitti &);
136: MGraniitti &operator=(const MGraniitti &);
137: // -----
138:
139: // Integration statistics
140: Stats stat;
141:
142: // Histogram fusion
143: bool hist_fusion_done = false;
144:
145: // Input string
146: std::string FULL_INPUT_STR = "null";
147:
148: // HepMC3 outputfile
149: std::string FULL_OUTPUT_STR = "null";
150: std::string OUTPUT = "null";
151: std::string FORMAT = "null"; // hepmc3 or hepmc2 or hepvt
152:
153: std::shared_ptr<HepMC3::GenRunInfo> runInfo = nullptr;
154: std::shared_ptr<HepMC3::WriterAscii> outputHepMC3 = nullptr;
155: std::shared_ptr<HepMC3::WriterAsciiHepMC2> outputHepMC2 = nullptr;
156: std::shared_ptr<HepMC3::WriterHEPEVT> outputHEPEVT = nullptr;
157:
158: // VEGAS creates copies here
159: std::vector<MProcess> *pvec;
160: MContinuum proc_C;
161: MFactized proc_F;
162: MQuasiElastic proc_Q;
163: MParton proc_P;
164:
165: // Forced cross-section for HepMC3 output
166: double xForceSec = 1;
```

```
./include/Granitti/MGraniitti.h 2/4
84:
85: // Weight statistics FLAT MC
86: double Wsum = 0.0;
87: double W2sum = 0.0;
88: double maxW = 0.0;
89:
90: // Weight statistics VEGAS MC
91: double maxF = 0.0;
92: double chi2 = 0.0;
93: };
94:
95: class MGraniitti {
96: public:
97: // Destructor & Constructor
98: MGraniitti();
99: ~MGraniitti();
100:
101: // For cross-section for HepMC3 output
102: void ForSec(double &x) { xForceSec = x; }
103:
104: // Read parameters
105: void ReadParam(const std::string &inputfile, const std::string &cmd_PROCESS = "null");
106:
107: // Initialize memory
108: void InitProcessMemory(std::string process, unsigned int NRESEED);
109: void InitMultiMemory();
110:
111: // Set simple MC parameters
112: void SetMCParam(MCFANUM &in);
113:
114: // Set VEGAS parameters
115: void SetVegasParam(const VEGASPARAM &in);
116:
117: // Set external file handle
118: void SetHepMC3Output(std::shared_ptr<HepMC3::WriterAscii<HepMC2>> &hepmc,
119:     const std::string & OUTPUTNAME) {
120:     OUTPUT = OUTPUTNAME;
121:     FORMAT = "hepmc2";
122:     outputHepMC2 = hepmc;
123: }
124: void SetHepMC3Output(std::shared_ptr<HepMC3::WriterAscii> &hepmc, const std::string &OUTPUTNAME) {
125:     OUTPUT = OUTPUTNAME;
126:     FORMAT = "hepmc3";
127:     outputHepMC3 = hepmc;
128: }
129:
130: // Maximum weight set/get
131: void SetMaxWeight(double w);
132: double GetMaxWeight() const;
133:
134: // Get number of CPU cores
135: void SetCores(int N) {
136:     CORES = N;
137:     if (CORES == 0) { // SETUP number of threads automatically
138:         CORES = std::round(std::thread::hardware_concurrency() * 1.5);
139:     }
140:     // If autodetection fails, set 1
141:     if (CORES < 1) { CORES = 1; }
142: }
143:
144: if (CORES < 0) {
145:     std::string str = "MGraniitti::SetCORES: CORES = 0";
146:     throw std::invalid_argument(str);
147: }
148:
149: int GetCores() const { return CORES; }
150: void SetWeight(const std::string &integrator) { INTEGRATOR = integrator; }
151: void SetWeighted(bool weighted) { WEIGHED = weighted; }
152:
153: // Output file name
154: void SetOutput(const std::string &output) { OUTPUT = output; }
155: // Output file format
156: void SetFormat(const std::string &format) {
157:     if (format == "hepmc3" || format == "hepmc2" || format == "hepvt") {
158:         FORMAT = format;
159:     } else {
160:         throw std::invalid_argument("MGraniitti::SetFormat: Unknown output format: " + format +
161:             " (valid: hepmc3, hepmc2, hepvt)");
162:     }
163: }
164: // Special method for combining histograms from each thread
165: void HistogramFusion();
166:
167: // Get cross section and error
```

```
./include/Granitti/MGraniitti.h 3/4
167: void GetXS(double &xs, double &xs_err) const {
168:     xs = stat.sigma;
169:     xs_err = stat.sigma_err;
170: }
171:
172: // Set number of (un)weighted events to be generated
173: void SetNumberOfEvents(int n) {
174:     if (n <= 0) {
175:         throw std::invalid_argument(
176:             "MGraniitti::SetNumberOfEvents(): Error: Number "
177:             "of events < 0!");
178:     }
179:     NEVENTS = n;
180: }
181:
182: int GetNumberOfEvents() const { return NEVENTS; }
183:
184: // Return processes
185: std::vector<std::string> GetProcessNumbers() const;
186:
187: // Initialize generator
188: void Initialize();
189: void Initialize(const MEikonal &eikonal_in);
190:
191: // Process object pointer, public so methods can be accessed
192: MProcess *proc = nullptr;
193:
194: void PrintHistograms();
195: void ReadGeneralParam(const std::string &inputfile);
196: void ReadProcessParam(const std::string &inputfile, const std::string &cmd_PROCESS = "null");
197: void ReadIntegParam(const std::string &inputfile);
198: void ReadReadOuts(const std::string &inputfile);
199: void ReadReadIn(const std::string &inputfile);
200: void ReadVetoCuts(const std::string &inputfile);
201: void ReadModelParam(const std::string &inputfile) const;
202:
203: void Generate();
204:
205: std::string PROCESS = "null"; // Physics process identifier
206:
207: bool WEIGHTED = false; // Unweighted or weighted event generation
208: int NEVENTS = 0; // Number of events to be generated
209: int CORES = 0; // Number of CPU cores (threads) in use
210: std::string INTEGRATOR = "null"; // Integrator (VEGAS, Flat, ...)
211:
212: // SILENT OUTPUT
213: bool HILJAA = false;
214:
215: private:
216: // -----
217: // Copy and assignment disabled by making the private
218: MGraniitti(const MGraniitti &);
219: MGraniitti &operator=(const MGraniitti &);
220: // -----
221:
222: // Integration statistics
223: Stats stat;
224:
225: // Histogram fusion
226: bool hist_fusion_done = false;
227:
228: // Input string
229: std::string FULL_INPUT_STR = "null";
230:
231: // HepMC3 outputfile
232: std::string FULL_OUTPUT_STR = "null";
233: std::string OUTPUT = "null";
234: std::string FORMAT = "null"; // hepmc3 or hepmc2 or hepvt
235:
236: std::shared_ptr<HepMC3::GenRunInfo> runInfo = nullptr;
237: std::shared_ptr<HepMC3::WriterAscii> outputHepMC3 = nullptr;
238: std::shared_ptr<HepMC3::WriterAsciiHepMC2> outputHepMC2 = nullptr;
239: std::shared_ptr<HepMC3::WriterHEPEVT> outputHEPEVT = nullptr;
240:
241: // VEGAS creates copies here
242: std::vector<MProcess> *pvec;
243: MContinuum proc_C;
244: MFactized proc_F;
245: MQuasiElastic proc_Q;
246: MParton proc_P;
247:
248: // Forced cross-section for HepMC3 output
249: double xForceSec = 1;
```

```

./include/Graniitti/MGraniitti.h      4/4
250: // Global timing
251: #ifimer global_timer;
252: #ifimer local_timer;
253: #ifimer atime;
254: double time_0 = 0.0;
255: double itertime = 0.0;
256:
257: // -----
258: // Process MC integration generic variables
259:
260: // Generation mode (integration = 0, event generation = 1)
261: unsigned int GMODE = 0;
262:
263: // -----
264: // FLAT MC
265: MCFPARAM mparam;
266:
267: // -----
268: // VEGAS MC
269:
270: // Parameters
271: VEGASPARAM vparam;
272:
273: // DATA
274: VEGASData VD;
275:
276: void VEGASInit(unsigned int init, unsigned int calls);
277: int VEGAS(unsigned int init, unsigned int calls, unsigned int iter, unsigned int N);
278: void VEGASMultiThread(unsigned int N, unsigned int tid, unsigned int init,
279:                        unsigned int LOCALcalls);
280:
281: // -----
282:
283: void UnifyHistogramBounds();
284:
285: // Calculate cross section
286: void CalculateCrossSection();
287:
288: // Vegas wrapper
289: double VegasWrapper(std::vector<double> &randvec, double wgt);
290:
291: // Event sampling/generation
292: void CallIntegrator(unsigned int N);
293: void SampleVegas(unsigned int N);
294: void SampleFlat(unsigned int N);
295: void SampleNeuro(unsigned int N);
296:
297: // Helper functions
298: void PrintInit() const;
299: int GetEventParameters("pr, double MAXM, const gra:auxintData aux);
300: void PrintStatus(unsigned int events, unsigned int N, #ifimer itime, double timercut);
301: void PrintStatistics(unsigned int N);
302: gra:PARAM_RES ResoFactorized(const std::string &tparam_str);
303: void InitFileOutput();
304:
305: // Interpreter commands
306: std::vector<aux::OneCMD> syntax;
307:
308: // -----
309: // Launcher functions
310: void MGraniitti(MGraniitti gen, int randomseed, int tid, int events);
311: void MGraniitti(MGraniitti agen);
312:
313: // namespace gra
314:
315: #endif

```

```

./include/Graniitti/Analysis/MROOT.h 1/4
1: // ROOT style namespace
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT license <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MROOT_H
7: #define MROOT_H
8:
9: #include <complex>
10: #include <memory>
11: #include <string>
12: #include <vector>
13:
14: // Own
15: #include "Graniitti/MMath.h"
16:
17: // ROOT
18: #include "TCanvas.h"
19: #include "TColor.h"
20: #include "TLabel.h"
21: #include "TMath.h"
22: #include "TROOT.h"
23: #include "TStyle.h"
24:
25: // Own
26: #include "Graniitti/MAux.h"
27:
28: namespace gra {
29: namespace rootstyle {
30:
31: // Returns a new colormap
32: inline void CreateColorMap(std::vector<int> & color, std::vector<std::shared_ptr<TColor>> & rootcolor,
33:                            int colorscheme = 1) {
34:     std::vector<std::vector<double>> colormap(150);
35:
36:     if (COLORSCHEME == 1) {
37:         // "Modern colormap"
38:         std::vector<std::vector<double>> cm = {{0, 0.4470, 0.7410}, {0.8500, 0.3250, 0.0980},
39:                                             {0.9290, 0.6940, 0.1230}, {0.4940, 0.1840, 0.5560},
40:                                             {0.4660, 0.6740, 0.1880}, {0.75, 0, 0.75}, {0.75, 0, 0.75}, {0.75, 0.75, 0},
41:                                             {0.6350, 0.0780, 0.1840}};
42:         colormap = cm;
43:     }
44:
45:     if (COLORSCHEME == 2) {
46:         // "Classic colormap"
47:         std::vector<std::vector<double>> cm = {{0, 0, 0.3}, {0, 0.5, 0}, {0.9, 0, 0},
48:                                             {0, 0.25, 0.75}, {0.75, 0, 0.75}, {0.75, 0.75, 0},
49:                                             {0.25, 0.25, 0.25}};
50:         colormap = cm;
51:     }
52:
53:     color = std::vector<int>(colormap.size(), 0);
54:     rootcolor = std::vector<std::shared_ptr<TColor>>(color.size(), nullptr);
55:
56:     for (const auto & i : aux::indices(color)) {
57:         color[i] = TColor::GetFreeColorIndex(i);
58:         // color[i] = 3000 + i; // some big number not used
59:     }
60:
61:     // ROOT style, we need to create some hidden memory part
62:     rootcolor[i] = TColor::GetFreeColorIndex(i);
63:     std::make_shared<TColor>(color[i], colormap[i][0], colormap[i][1], colormap[i][2]);
64: }
65:
66: // Create grid canvas with N subpads
67: inline void AutoGridCanvas(std::shared_ptr<TCanvas> & ci, unsigned int N) {
68:     unsigned int ADD = 0;
69:     while (true) { // Adjust grid size
70:         const unsigned int val = std::sqrt(N + ADD);
71:         if (val * val == (N + ADD)) { break; }
72:         ++ADD;
73:     }
74:
75:     // Calculate if we have a full empty row -> remove that
76:     unsigned int DEL = 0;
77:     if (ADD * ADD - ADD == N) { DEL = 1; }
78:
79:     const unsigned int COLS = std::sqrt(N + ADD);
80:     const unsigned int ROWS = std::sqrt(N + ADD) - DEL;
81:
82:     // Adjust aspect ratio
83:     if (ROWS == COLS) {
167: double red[NRGBs] = {0.00, 0.00, 0.87, 1.00, 0.51};
168: double green[NRGBs] = {0.00, 0.81, 1.00, 0.20, 0.00};
169: double blue[NRGBs] = {0.51, 1.00, 0.12, 0.60, 0.00};
170: TColor::CreateGradientColorTable(NRGBs, stops, red, green, blue, kCont);
171:
172: // See https://root.cern.ch/doc/master/classTColor.html
173: gStyle->SetPalette(53); // kDarkBodyRadiator
174: gStyle->SetPalette(51); // kHepData
175: gStyle->SetPalette(87); // kLightTemperature
176: gStyle->SetPalette(109); // kThermometer
177: gStyle->SetPalette(71); // kMicrocannelow
178: gStyle->SetPalette(57); // kBird
179: gStyle->SetPalette(19); // kCherry
180: gStyle->SetPalette(112); // kVividize
181:
182: // TColor::InvertPalette(); // Palette inversion
183: #else if (style == "gray") {
184:     const int NRGBs = 5;
185:
186:     double stops[NRGBs] = {0.00, 0.34, 0.61, 0.84, 1.00};
187:     double red[NRGBs] = {1.00, 0.84, 0.61, 0.34, 0.00};
188:     double green[NRGBs] = {1.00, 0.84, 0.61, 0.34, 0.00};
189:     double blue[NRGBs] = {1.00, 0.84, 0.61, 0.34, 0.00};
190:     TColor::CreateGradientColorTable(NRGBs, stops, red, green, blue, kCont);
191: #else if (style == "cubehelix") {
192:     const int
193:     const std::vector<std::vector<double>> NRGBs = 256;
194:     const std::vector<std::vector<double>> M = Cubehelix(NRGBs, 0.5, -1.5, 1.2, 1.0);
195:     double stops[NRGBs];
196:     double red[NRGBs];
197:     double green[NRGBs];
198:     double blue[NRGBs];
199:
200:     for (std::size_t i = 0; i < NRGBs; ++i) {
201:         stops[i] = M[i][1];
202:         red[i] = M[i][4];
203:         green[i] = M[i][5];
204:         blue[i] = M[i][3];
205:     }
206:
207:     TColor::CreateGradientColorTable(NRGBs, stops, red, green, blue, kCont);
208: }
209:
210: gStyle->SetNumberContours(NCont);
211:
212: gStyle->SetTitleOffset(1.6, "x"); // title offset from axis
213: gStyle->SetTitleOffset(1.0, "y"); // title offset from axis
214: gStyle->SetTitleSize(0.03, "x"); // title size
215: gStyle->SetTitleSize(0.03, "y"); // title size
216: gStyle->SetTitleSize(0.03, "z"); // title size
217: gStyle->SetLabelOffset(0.025);
218:
219: // Necessary with multiple plots per canvas
220: gStyle->SetPadBottomMargin(0.1);
221: gStyle->SetPadLeftMargin(0.1);
222: gStyle->SetPadRightMargin(0.09);
223: gStyle->SetPadTopMargin(0.1);
224:
225: // Global Style Setup
226: inline void SetROOTStyle() {
227:     gStyle->SetOptStat(0); // Statistics Box OFF [0,1]
228:
229:     gStyle->SetOptFit(1); // Fit parameters
230:
231:     gStyle->SetTitleSize(0.04, "x"); // Title with "*" (or anything else than xyz)
232:     gStyle->SetTitleSize(0.04, "y"); // Title with "*" (or anything else than xyz)
233:     gStyle->SetTitleSize(0.04, "z"); // Title with "*" (or anything else than xyz)
234:     gStyle->SetStat(1.0);
235:     gStyle->SetStat(0.15);
236:     gStyle->SetStat(0.09);
237:
238:     SetPlotStyle();
239: }
240:
241: // Before calling this, call mother TCanvas cd->()
242: inline void TransparentPad(std::shared_ptr<TPad> pad) {
243:     pad->SetLineStyle(kDashed); // "transparent pad", 0, 0, 1, 1;
244:     pad->SetFillStyle(4000);
245:     pad->Draw();
246:     pad->cd();
247: }
248:
249: // Create GRANITTI Text

```

```

./include/Graniitti/Analysis/MROOT.h 2/4
84: ci = std::make_shared<TCanvas>("ci", "ci", 700, 600); // horizontal, vertical
85:
86: #if (ROWS != COLS) {
87:     ci = std::make_shared<TCanvas>("ci", "ci", 700, 450); // horizontal, vertical
88: }
89:
90: ci->Divide(COLS, ROWS, 0.002, 0.001);
91:
92: // This is needed
93: gStyle->SetPadLeftMargin(0.15);
94:
95:
96: // CubeHelix colormap generator
97:
98: // N = number of discrete stops
99: // start color (1 = red ... 2 = green ... 3 = red)
100: // R = number of helix rotations
101: // hue = hue, with 0 gives black/white
102: // gamma = intensity correction
103:
104: // Default CubeHelix(256, 0.5, -1.5, 1.2, 1.0)
105:
106: // REFERENCE: D.A. Green, https://arxiv.org/abs/1108.3083
107: // https://www.mrao.cam.ac.uk/dag/CUBEHELIX
108:
109: inline std::vector<std::vector<double>> CubeHelix(int N, double start, double R, double hue,
110:                                                double gamma) {
111:     auto limitfunc = [](std::vector<double> & x) {
112:         for (std::size_t i = 0; i < x.size(); ++i) {
113:             if (x[i] < 0.0) { x[i] = 0.0; }
114:             if (x[i] > 1.0) { x[i] = 1.0; }
115:         }
116:     };
117:
118:     // Color matrix
119:     const std::vector<std::vector<double>> A = {
120:         {{0.14861, 1.76271}, {-0.23227, -0.305649}, {1.97294, 0}};
121:         const double PI = 3.14159265359;
122:
123:         // Steps, Red, Green, Blue
124:         std::vector<std::vector<double>> M(4, std::vector<double>(N, 0.0));
125:
126:         // Lightning
127:         const double maxlight = 1.0;
128:         const double minlight = 0.0;
129:         const double lightstep = (maxlight - minlight) / N;
130:
131:         for (int i = 1; i <= N; ++i) {
132:             // Rotation angle and shift
133:             double alpha = (i - 1.0) / (N - 1.0);
134:             const double phi = 2.0 * PI * (start + 0.34 + i.0 + R * alpha);
135:
136:             // Apply gamma-correction
137:             alpha = std::pow(alpha, gamma);
138:             const double a = hue * alpha * (1.0 - alpha) / 2.0;
139:
140:             // Affine Map
141:             const std::vector<double> x = {a * std::cos(phi), a * std::sin(phi)};
142:             std::vector<double>
143:             A[1][0] * x[0] + A[1][1] * x[1] + alpha,
144:             A[2][0] * x[0] + A[2][1] * x[1] + alpha;
145:             limitfunc(); // Limit values to [0,1]
146:
147:             // Save values
148:             M[0][i - 1] = minlight + lightstep * (i - 1);
149:             for (std::size_t j = 0; j < 3; ++j) { M[j][i - 1] = y[j]; }
150:         }
151:
152:         return M;
153:     };
154:
155:
156: // Set "nice" 2D-plot style
157: inline void SetPlotStyle() {
158:     // Set smooth color gradients
159:     const int NCont = 256;
160:
161:     const std::string style = "default";
162:
163:     if (style == "default") {
164:         const int NRGBs = 5;
165:
166:         double stops[NRGBs] = {0.00, 0.34, 0.61, 0.84, 1.00};

```

```

./include/Granitti/Analysis/MROOT.h      4/4
250: inline void MadeInFinland(std::shared_ptr<Latex> & l1, std::shared_ptr<Latex> & l2,
251:                             double xpos = 0.935, std::vector<double> ypos = {0.03, 0.58}) {
252:     if (ypos.size() != 2) {
253:         throw std::invalid_argument("MROOT::MadeInFinland: argument ypos should be of size 2");
254:     }
255: }
256:
257: l1 = std::make_shared<Latex>(xpos, ypos[0], gra::aux::GetVersionLatex().c_str());
258: l2->SetNDC(1) // Normalized coordinates
259: l1->SetTextAngle(90);
260: l1->Draw();
261: l2 = std::make_shared<Latex>(xpos, ypos[1], gra::aux::GetWebText(0).c_str());
262: l2->SetNDC(1) // Normalized coordinates
263: l2->SetTextAngle(90);
264: l2->Draw();
265: }
266:
267: } // namespace rootstyle
268: } // namespace gra
269:
270: #endif

```

```

./include/Granitti/Analysis/MAnalyzer.h  1/2
1: // Fast analysis class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MANALYZER_H
7: #define MANALYZER_H
8:
9: #include <complex>
10: #include <memory>
11: #include <vector>
12:
13: // ROOT
14: #include "TBranch.h"
15: #include "TCanvas.h"
16: #include "TColor.h"
17: #include "TFile.h"
18: #include "TFile.h"
19: #include "TH1.h"
20: #include "TH2.h"
21: #include "TLegend.h"
22: #include "TLorentzVector.h"
23: #include "TProfile.h"
24: #include "TROOT.h"
25: #include "TRandom3.h"
26: #include "TString.h"
27: #include "TStyle.h"
28: #include "TTree.h"
29:
30: // HepMC3
31: #include "HepMC3/FourVector.h"
32: #include "HepMC3/GenEvent.h"
33: #include "HepMC3/GenParticle.h"
34: #include "HepMC3/GenVertex.h"
35: #include "HepMC3/Print.h"
36: #include "HepMC3/ReaderAscii.h"
37: #include "HepMC3/Relativis.h"
38: #include "HepMC3/Selector.h"
39: #include "HepMC3/WriterAscii.h"
40:
41: // Own
42: #include "Granitti/Analysis/Multiplet.h"
43: #include "Granitti/MNux.h"
44:
45: namespace gra {
46:
47: namespace analyzer {
48:
49: // Different Lorentz frame labels
50: const std::vector<std::string> FRAMES = {"CM", "BK", "CS", "PC", "GJ", "LAB"};
51:
52: } // namespace analyzer
53:
54: class MAnalyzer {
55: public:
56:     // Constructor, destructor
57:     MAnalyzer(const std::string &ID);
58:     ~MAnalyzer();
59:
60:     // Default daughter particle name string
61:     std::string get = "daughter";
62:
63:     // -----
64:     // Forward system quantities
65:     std::shared_ptr<TH1D> hE_Pions;
66:     std::shared_ptr<TH1D> hE_Gamma;
67:     std::shared_ptr<TH1D> hE_Neutron;
68:     std::shared_ptr<TH1D> hE_GammaNeutron;
69:
70:     std::shared_ptr<TH1D> hXF_Pions;
71:     std::shared_ptr<TH1D> hXF_Gamma;
72:     std::shared_ptr<TH1D> hXF_Neutron;
73:
74:     std::shared_ptr<TH1D> hEta_Pions;
75:     std::shared_ptr<TH1D> hEta_Gamma;
76:     std::shared_ptr<TH1D> hEta_Neutron;
77:     std::shared_ptr<TH1D> hM_NHTM;
78:
79:     // -----
80:     // Angular observables
81:     std::shared_ptr<TProfile> hP1[8];
82:
83:     static constexpr unsigned int NFR = 6; // number of frames

```

```

./include/Granitti/Analysis/MAnalyzer.h  2/2
84:
85: // Correlations between frames
86: std::shared_ptr<TH2D> h2CoorThea[NFR][NFR];
87: std::shared_ptr<TH2D> h2Phi[NFR][NFR];
88: // -----
89:
90: // HepMC3 reader
91: double HepMC3_GraceFill(const std::string inputfile, unsigned int multiplicity, int finalPOG,
92:                         unsigned int MAXEVENTS,
93:                         std::map<std::string, std::shared_ptr<CMultiplet>> & h1,
94:                         std::map<std::string, std::shared_ptr<CMultiplet>> & h2,
95:                         std::map<std::string, std::shared_ptr<CProfMultiplet>> &hP,
96:                         unsigned int SID);
97:
98: // Plot out all local histograms
99: void PlotAll(const std::string scitextstr);
100:
101: double cross_section = 0;
102:
103: double CheckEnergyMomentum(HepMC3::GenEvent evt) const;
104: void FrameObservables(double W, HepMC3::GenEvent evt, const M4Vec ep_beam_plus,
105:                      const M4Vec ep_beam_minus, const M4Vec ep_final_plus,
106:                      const M4Vec ep_final_minus, const std::vector<M4Vec> &sp,
107:                      const std::vector<M4Vec> &spin);
108: void NStarObservables(double W, HepMC3::GenEvent evt);
109:
110: private:
111: double sqrts = 0.0;
112:
113: // Name of the HepMC33 input
114: std::string inputfile;
115:
116: // Proton excitation is turned on
117: bool N_STAR_ON = false;
118: };
119:
120: } // namespace gra
121:
122: #endif

```

```

./include/Granitti/Analysis/Multiplet.h  1/2
1: // Multiplet of ROOT histograms
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MULTIPLETCCLASS_H
7: #define MULTIPLETCCLASS_H
8:
9: // C++
10: #include <string>
11:
12: // ROOT
13: #include "TCanvas.h"
14: #include "TH1.h"
15: #include "TH2.h"
16: #include "TProfile.h"
17:
18: // Own
19: #include "Granitti/MNux.h"
20:
21: namespace gra {
22:
23: // Histogram boundaries
24: class hIBound {
25: public:
26:     hIBound(unsigned int N, double min, double max) {
27:         N = _N;
28:         min = _min;
29:         max = _max;
30:     }
31:     ~hIBound() {}
32:
33:     unsigned int N = 0;
34:     double min = 0.0;
35:     double max = 0.0;
36: };
37:
38: class hMultiplet {
39: public:
40:     hMultiplet(const std::string sname, const std::string llabeltext, int N, double minval,
41:               double maxval, const std::vector<std::string> &legendtext);
42:     ~hMultiplet() {}
43:     for (const auto &i : gra::aux::indices(h)) { delete h[i]; }
44:
45: void MultiFill(const std::vector<double> &x) {
46:     for (const auto &i : gra::aux::indices(h)) { h[i]->Fill(x[i]); }
47: }
48: void MultiFill(const std::vector<double> &x, const std::vector<double> &weight) {
49:     for (const auto &i : gra::aux::indices(h)) { h[i]->Fill(x[i], weight[i]); }
50: }
51: void NormalizeAll(const std::vector<double> &cross_section,
52:                  const std::vector<double> &multiplier);
53:
54: // Plot and save ID-histogram Multiplet
55: std::vector<double> SaveFig(const std::string &fullpath, bool RATIOPLOT = true) const;
56:
57: int N_;
58: double minval_;
59: double maxval_;
60: std::string legendposition_;
61: std::string name_;
62:
63: // Histogram pointers here
64: std::vector<TH1 *> h;
65: std::vector<std::string> legendtext_;
66:
67: class h2Multiplet {
68: public:
69:     h2Multiplet(const std::string sname, const std::string llabeltext, int N1, double minval1,
70:                double maxval1, int N2, double minval2, double maxval2,
71:                const std::vector<std::string> &legendtext);
72:
73:     ~h2Multiplet() {}
74:     for (const auto &i : gra::aux::indices(h)) { delete h[i]; }
75:
76: void MultiFill(const std::vector<std::vector<double>> &x) {
77:     for (const auto &i : gra::aux::indices(h)) { h[i]->Fill(x[i][0], x[i][1]); }
78: }
79: void MultiFill(const std::vector<std::vector<double>> &x, const std::vector<double> &weight) {
80:     for (const auto &i : gra::aux::indices(h)) { h[i]->Fill(x[i][0], x[i][1], weight[i]); }
81: }
82: void NormalizeAll(const std::vector<double> &cross_section,
83:                  const std::vector<double> &multiplier);

```

./include/Granitti/Analysis/MMultiplet.h 2/2

```
84:
85: // Plot and save histogram Multiplet
86: double SaveFig(const std::string fullpath, bool RATIO_PLOT = true) const;
87:
88: int N1;
89: double minval1;
90: double maxval1;
91:
92: int N2;
93: double minval2;
94: double maxval2;
95:
96: std::string name_;
97:
98: // Histogram pointers here
99: std::vector<T> * h;
100: std::vector<std::string> legendtext_;
101: };
102:
103: class HProfMultiplet {
104: public:
105:   HProfMultiplet(const std::string fname, const std::string llabeltext, int N, double minval,
106:                 double maxval1, double maxval2,
107:                 const std::vector<std::string> legendtext);
108:
109:   "HProfMultiplet()" {
110:     for (const auto &i : gra::raun::indices(h)) { delete h[i]; }
111:   }
112:   void MultiFill(const std::vector<std::vector<double>> &w) {
113:     for (const auto &i : gra::raun::indices(h)) { h[i]->Fill(x[i][0], x[i][1]); }
114:   }
115:   void MultiFill(const std::vector<std::vector<double>> &w, const std::vector<double> &weight) {
116:     for (const auto &i : gra::raun::indices(h)) { h[i]->Fill(x[i][0], x[i][1], weight[i]); }
117:   }
118:
119:   // Plot and save histogram Multiplet
120:   double SaveFig(const std::string fullpath, bool RATIO_PLOT = true) const;
121:
122:   int N_;
123:   double minval1;
124:   double maxval1;
125:
126:   double minval2;
127:   double maxval2;
128:
129:   std::string name_;
130:   std::string legendposition_;
131:
132:   // Histogram pointers here
133:   std::vector<TProfile * > h;
134:   std::vector<std::string> legendtext_;
135: };
136:
137: } // namespace gra
138:
139: #endif
```

./include/Granitti/Analysis/MHarmonic.h 2/2

```
84:         const std::string koutputpath) const;
85:
86: void PlotFigures(const std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>> &tensor,
87:                 unsigned int OBSERVABLE, const std::string STYPESTRING, int barcolor,
88:                 const std::string koutputpath) const;
89:
90: void PlotFigures2D(const std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>> &tensor,
91:                   const std::vector<int> &OBSERVABLEZ, const std::string STYPESTRING,
92:                   int barcolor, const std::string koutputpath) const;
93:
94: void PlotIDefficiency(unsigned int OBSERVABLE, const std::string koutputpath) const;
95:
96: // Parameters
97: HPARAM param;
98: int NCOORD;
99:
100: private:
101:   MMatrix<double> Y_min; // Calculate once for speed
102:   std::vector<gra::spherical::Omega> DATA_events; // Input data
103:
104: // ===== MINUIT fit/const function
105: // Needed by MINUIT fit/const function
106:
107: // Currently active
108: std::vector<std::size_t> DATA_ind; // Hypercell indices
109: std::vector<std::size_t> activecell; // Hypercell indices
110: std::vector<double> t_lm; // Fitted moments
111: std::vector<double> t_lm_error; // Their errors
112:
113: // Error and covariance matrices
114: MMatrix<double> errmat;
115: MMatrix<double> covmat;
116:
117: // =====
118:
119: // Active moments bookkeeping here
120: std::vector<bool> ACTIVE;
121: unsigned int ACTIVEIND = 0;
122:
123: // Detector expansion tensor data
124: MTensor<gra::spherical::SH_DET> det_DET;
125: MTensor<gra::spherical::SH_DET> fit_DET;
126: MTensor<gra::spherical::SH_DET> fit_DET;
127:
128: // Expanded data in a map for different data sources
129: std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>> det;
130: std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>> fit;
131: std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>> fit;
132:
133: // Discrimination
134: std::vector<std::vector<EDGE>> grid;
135:
136: // Plotting
137: const std::vector<int> colors = {46, 38, 43, 30, 25, 14, 9, 19, 29, 39, 49};
138: const std::vector<std::string> xlabel = {"M", "E", "E", "E", "E", "E", "E", "E", "E", "E", "E"};
139: };
140:
141: } // namespace gra
142:
143: #endif
```

./include/Granitti/Analysis/MHarmonic.h 1/2

```
1: // Spherical harmonic expansion and analysis class
2:
3:
4: // (c) 2017-2020 Mikael Mieskolainen
5: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
6:
7: #ifndef M_HARMONIC_H
8: #define M_HARMONIC_H
9:
10: // C++
11: #include <complex>
12: #include <random>
13: #include <vector>
14:
15: // Own
16: #include "Granitti/MBox.h"
17: #include "Granitti/MDimArray.h"
18: #include "Granitti/MMatrix.h"
19: #include "Granitti/Random.h"
20: #include "Granitti/MSpherical.h"
21: #include "Granitti/MTensor.h"
22:
23: namespace gra {
24: class MHarmonic {
25: public:
26:   // Bin edges
27:   struct EDGE {
28:     double min = 0;
29:     double max = 0;
30:   };
31:   double center() const { return (max + min) / 2.0; }
32: };
33:
34: // Parameters
35: struct HPARAM {
36:   std::vector<double> M = {0, 0, 0};
37:   std::vector<double> Y = {0, 0, 0};
38:   std::vector<double> PT = {0, 0, 0};
39:
40:   int LM_MAX = 2; // Maximum spherical harmonic truncation degree (non-negative)
41:   // Integer
42:   bool REMOVE_ZERO = true; // Fix odd moments to zero (due to lacking specific spin)
43:   // States, for example
44:   bool REMOVE_NEGATIVE = true; // Fix negative m to zero (due to parity)
45:   // Conservation, for example
46:   bool EML = false; // Use Extended Maximum Likelihood fit
47:   double SVDRG = 0.001; // SVD regularization strength in algebraic inverse (put 0
48:   // for no regularization)
49:   double LIREG = 0.001; // L1-norm regularization in EML fit (put 0 for no
50:   // regularization)
51:
52:   void Print() {
53:     std::cout << "HARMONIC EXPANSION PARAMETERS:" << std::endl << std::endl;
54:
55:     std::cout << "LM_MAX" << std::endl;
56:     std::cout << "REMOVE_ZERO" << std::endl;
57:     std::cout << "REMOVE_NEGATIVE" << std::endl;
58:     std::cout << "EML" << std::endl;
59:     std::cout << "SVDRG" << std::endl;
60:     std::cout << "LIREG" << std::endl;
61:   }
62: };
63:
64: // Constructor, destructor
65: MHarmonic();
66: ~MHarmonic();
67:
68: void Init(const HPARAM &hp);
69:
70: void HyperLoop(void (*f)(int &int, double *, double *, double *, int),
71:               const std::vector<gra::spherical::Omega> &ME,
72:               const std::vector<gra::spherical::Data> &DATA, const HPARAM &hp);
73:
74: void MomentFit(const gra::spherical::Meta &META, const std::vector<std::size_t> &cell,
75:               void (*f)(int &int, double *, double *, int));
76:
77: double PrintOutHyperCell(const gra::spherical::Meta &META, const std::vector<std::size_t> &cell);
78: void logFunc(int npar, double *gin, double *sf, double *par, int iflag) const;
79: bool PrintLoop(const std::string &outpath) const;
80: void PlotAll(const std::string &outpath) const;
81:
82: void Plot2DExpansion(const std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>> &tensor,
83:                    unsigned int OBSERVABLE, const std::string STYPESTRING, int barcolor,
```

./include/Granitti/MMatOper.h 1/4

```
1: // Simple matrix/vector operations not included in MMatrix. [HEADER ONLY FILE]
2:
3:
4: // (c) 2017-2020 Mikael Mieskolainen
5: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
6:
7: #ifndef M_MATOPER_H
8: #define M_MATOPER_H
9:
10: // C++
11: #include <complex>
12: #include <iostream>
13: #include <vector>
14:
15: // Own
16: #include "Granitti/MBox.h"
17: #include "Granitti/MMatrix.h"
18: namespace gra {
19: namespace MatOper {
20:
21: // MATIAR style meshgrid
22:
23: // X is a matrix with each row = x
24: // Y is a matrix with each column = y
25: template <typename T>
26: inline void MeshGrid(const std::vector<T> &x, const std::vector<T> &y, MMatrix<T> &X,
27:                    MMatrix<T> &Y) {
28:   X = MMatrix<T>(y.size(), x.size());
29:   Y = MMatrix<T>(y.size(), x.size());
30:
31:   for (std::size_t i = 0; i < X.size_row(); ++i) {
32:     for (std::size_t j = 0; j < X.size_col(); ++j) { X[i][j] = x[j]; }
33:   }
34:   for (std::size_t i = 0; i < Y.size_row(); ++i) {
35:     for (std::size_t j = 0; j < Y.size_col(); ++j) { Y[i][j] = y[j]; }
36:   }
37: }
38:
39: // diag(n) * A * diag(y) product
40: template <typename T>
41: inline MMatrix<T> diagdiag(const std::vector<T> &x, const MMatrix<T> &A, const std::vector<T> &y) {
42:   // Over [i,j] of A
43:   MMatrix<T> A(x.size(), x.size());
44:   const std::size_t n = A.size_row();
45:   const std::size_t m = A.size_col();
46:
47:   MMatrix<T> C(n, m);
48:   for (std::size_t i = 0; i < n; ++i) {
49:     for (std::size_t j = 0; j < m; ++j) { C[i][j] = x[i] * A[i][j] * y[j]; }
50:   }
51:   return C;
52: }
53:
54: // Return diagonal matrix constructed from a vector
55: template <typename T>
56: inline MMatrix<T> Diag(const std::vector<T> &x) {
57:   MMatrix<T> A(x.size(), x.size());
58:   for (std::size_t i = 0; i < x.size(); ++i) { A[i][i] = x[i]; }
59:   return A;
60: }
61:
62: // Return L2-norm vector ||x||
63: template <typename T>
64: inline std::vector<T> Unit(const std::vector<T> &x) {
65:   double norm = 0.0;
66:   for (std::size_t k = 0; k < x.size(); ++k) { norm += gra::math::abs2(x[k]); }
67:   norm = gra::math::msqrt(norm); // L2-norm
68:
69:   std::vector<T> y = x;
70:   for (std::size_t k = 0; k < y.size(); ++k) {
71:     y[k] /= norm; // normalize
72:   }
73:   return y;
74: }
75:
76: // Return unit sum normalized (sum_i x_i = 1)
77: template <typename T>
78: inline std::vector<T> Normalized(const std::vector<T> &x) {
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492: }
493:
494: // Return unit sum normalized (sum_i x_i = 1)
495: template <typename T>
496: inline std::vector<T> Unit(const std::vector<T> &x) {
497:   double norm = 0.0;
498:   for (std::size_t k = 0; k < x.size(); ++k) { norm += gra::math::abs2(x[k]); }
499:   norm = gra::math::msqrt(norm); // L2-norm
500:
501:   std::vector<T> y = x;
502:   for (std::size_t k = 0; k < y.size(); ++k) {
503:     y[k] /= norm; // normalize
504:   }
505:   return y;
506: }
507:
508: // Return unit sum normalized (sum_i x_i = 1)
509: template <typename T>
510: inline std::vector<T> Normalized(const std::vector<T> &x) {
511:   double sum = 0.0;
512:   for (std::size_t k = 0; k < x.size(); ++k) { sum += x[k]; }
513:   std::vector<T> y = x;
514:   if (sum > 0) {
515:     for (std::size_t k = 0; k < y.size(); ++k) {
516:       y[k] /= sum; // normalize
517:     }
518:   }
519: }
520:
521: // Return unit sum normalized (sum_i x_i = 1)
522: template <typename T>
523: inline std::vector<T> Unit(const std::vector<T> &x) {
524:   double norm = 0.0;
525:   for (std::size_t k = 0; k < x.size(); ++k) { norm += gra::math::abs2(x[k]); }
526:   norm = gra::math::msqrt(norm); // L2-norm
527:
528:   std::vector<T> y = x;
529:   for (std::size_t k = 0; k < y.size(); ++k) {
530:     y[k] /= norm; // normalize
531:   }
532:   return y;
533: }
534:
535: // Return unit sum normalized (sum_i x_i = 1)
536: template <typename T>
537: inline std::vector<T> Normalized(const std::vector<T> &x) {
538:   double sum = 0.0;
539:   for (std::size_t k = 0; k < x.size(); ++k) { sum += x[k]; }
540:   std::vector<T> y = x;
541:   if (sum > 0) {
542:     for (std::size_t k = 0; k < y.size(); ++k) {
543:       y[k] /= sum; // normalize
544:     }
545:   }
546: }
547:
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550: inline std::vector<T> Unit(const std::vector<T> &x) {
551:   double norm = 0.0;
552:   for (std::size_t k = 0; k < x.size(); ++k) { norm += gra::math::abs2(x[k]); }
553:   norm = gra::math::msqrt(norm); // L2-norm
554:
555:   std::vector<T> y = x;
556:   for (std::size_t k = 0; k < y.size(); ++k) {
557:     y[k] /= norm; // normalize
558:   }
559:   return y;
560: }
561:
562: // Return unit sum normalized (sum_i x_i = 1)
563: template <typename T>
564: inline std::vector<T> Normalized(const std::vector<T> &x) {
565:   double sum = 0.0;
566:   for (std::size_t k = 0; k < x.size(); ++k) { sum += x[k]; }
567:   std::vector<T> y = x;
568:   if (sum > 0) {
569:     for (std::size_t k = 0; k < y.size(); ++k) {
570:       y[k] /= sum; // normalize
571:     }
572:   }
573: }
574:
575: // Return unit sum normalized (sum_i x_i = 1)
576: template <typename T>
577: inline std::vector<T> Unit(const std::
```

```

./include/Granitti/MMatOper.h      2/4
84: }
85: }
86: return y;
87: }
88:
89: // Multiply vector by scalar
90: template <typename T, typename T2>
91 inline void ScaleVector(std::vector<T> &A, T2 scale) {
92   for (std::size_t i = 0; i < A.size(); ++i) { A[i] *= scale; }
93: }
94:
95: // Add a constant matrix: A + (full ones) * c
96: template <typename T, typename T2>
97 inline void AddConstant(MMatrix<T> &A, T2 c) {
98   for (std::size_t i = 0; i < A.size_row(); ++i) {
99     for (std::size_t j = 0; j < A.size_col(); ++j) { A[i][j] += c; }
100: }
101: }
102:
103: // Vector-Matrix-Vector Multiplication
104: // where a is (1 x n)
105: // B is (m x n)
106: // c is (n x 1)
107: template <typename T>
108 inline T VecMatVecMultiply(const std::vector<T> &a, const MMatrix<T> &B, const std::vector<T> &c) {
109:   const std::size_t m = A.size();
110:   const std::size_t n = C.size();
111:   T value = 0.0;
112:   for (std::size_t i = 0; i < m; ++i) {
113:     for (std::size_t j = 0; j < n; ++j) {
114:       value += A[i] * B[i][j] * C[j]; // notice plus
115:     }
116:   }
117:   return value;
118: }
119:
120: // Vector-Matrix Multiplication c = aB,
121: // where a is (1 x n)
122: // B is (m x n)
123: // which gives c (1 x n)
124: template <typename T>
125 inline std::vector<T> VecMatMultiply(const std::vector<T> &a, const MMatrix<T> &B) {
126:   const std::size_t m = A.size();
127:   const std::size_t n = B.size_col();
128:   std::vector<T> C(m, 0.0); // init with zero!
129:   // Over [j] of c
130:   for (std::size_t j = 0; j < n; ++j) {
131:     for (std::size_t k = 0; k < m; ++k) {
132:       C[j] += A[k] * B[k][j]; // notice plus
133:     }
134:   }
135:   return C;
136: }
137:
138: // Vector-Vector multiplication (strictly not a dot product, which would have
139: // complex conjugation for the other)
140: template <typename T>
141 inline T VecVecMultiply(const std::vector<T> &a, const std::vector<T> &b) {
142:   T out = 0.0;
143:   for (std::size_t i = 0; i < A.size(); ++i) {
144:     out += A[i] * B[i]; // notice plus
145:   }
146:   return out;
147: }
148:
149: // Vector Conjugate Transpose (Dagger)
150: // X is (n x 1)
151: // Y is (1 x n)
152: template <typename T>
153 inline std::vector<T> VecDagger(const std::vector<T> &x) {
154:   const unsigned int n = X.size();
155:   std::vector<T> Y(n, 0.0);
156:   // Conjugate element by element
157:   for (std::size_t i = 0; i < n; ++i) { Y[i] = std::conj(X[i]); }
158:   return Y;
159: }
160:
161: // Cross product for two 3-vectors
162:

```

```

./include/Granitti/MMatOper.h      4/4
250: inline void PrintMatrix(const MMatrix<T> &A, const std::string name = "") {
251:   // Print elements
252:   std::cout << name << " : ";
253:   for (std::size_t i = 0; i < A.size_row(); ++i) {
254:     for (std::size_t j = 0; j < A.size_col(); ++j) {
255:       const std::string delim = (j < A.size_col() - 1) ? ", " : "";
256:       printf("%6.3f%6.3f%e", std::real(A[i][j]), std::imag(A[i][j]), delim.c_str());
257:     }
258:     std::cout << std::endl;
259:   }
260:   std::cout << std::endl;
261: }
262:
263: // Print out matrix elements as two separate matrix
264: template <typename T>
265 inline void PrintMatrixSeparate(const MMatrix<T> &A) {
266:   // Print real part
267:   std::cout << "Re:" << std::endl;
268:   for (std::size_t i = 0; i < A.size_row(); ++i) {
269:     for (std::size_t j = 0; j < A.size_col(); ++j) {
270:       const std::string delim = (j < A.size_col() - 1) ? ", " : "";
271:       printf("%6.3f%e", std::real(A[i][j]), delim.c_str());
272:     }
273:     std::cout << std::endl;
274:   }
275:   std::cout << std::endl;
276: }
277: // Print imaginary part
278: std::cout << "Im:" << std::endl;
279: for (std::size_t i = 0; i < A.size_row(); ++i) {
280:   for (std::size_t j = 0; j < A.size_col(); ++j) {
281:     const std::string delim = (j < A.size_col() - 1) ? ", " : "";
282:     printf("%6.3f%e", std::imag(A[i][j]), delim.c_str());
283:   }
284:   std::cout << std::endl;
285: }
286: std::cout << std::endl;
287: }
288:
289: // namespace matorp
290: } // namespace gra
291:
292: #endif

```

```

./include/Granitti/MMatOper.h      3/4
167: template <typename T>
168: inline std::vector<T> Cross(const std::vector<T> &a, const std::vector<T> &b) {
169:   std::vector<T> c(a.size(), 0.0);
170:   a[0] * b[1] - a[1] * b[0];
171:   return c;
172: }
173:
174: // Negate vector
175: template <typename T>
176: inline std::vector<T> Negat(const std::vector<T> &a) {
177:   std::vector<T> c(a.size(), 0.0);
178:   for (std::size_t k = 0; k < a.size(); ++k) { c[k] = -a[k]; }
179:   return c;
180: }
181:
182: // Subtract two vectors
183: template <typename T>
184: inline std::vector<T> Minus(const std::vector<T> &a, const std::vector<T> &b) {
185:   std::vector<T> c(a.size(), 0.0);
186:   for (std::size_t k = 0; k < a.size(); ++k) { c[k] = a[k] - b[k]; }
187:   return c;
188: }
189:
190: // Add two vectors
191: template <typename T>
192: inline std::vector<T> Plus(const std::vector<T> &a, const std::vector<T> &b) {
193:   std::vector<T> c(a.size(), 0.0);
194:   for (std::size_t k = 0; k < a.size(); ++k) { c[k] = a[k] + b[k]; }
195:   return c;
196: }
197:
198: // Outer (tensor) product of two vectors: u (n) v = u v^T
199: template <typename T>
200: inline MMatrix<T> OuterProd(const std::vector<T> &u, const std::vector<T> &v) {
201:   const unsigned int n = u.size();
202:   const unsigned int m = v.size();
203:   MMatrix<T> out(m, n);
204:   for (std::size_t i = 0; i < m; ++i) {
205:     for (std::size_t j = 0; j < n; ++j) { out[i][j] = u[i] * v[j]; }
206:   }
207:   return out;
208: }
209:
210: // Kronecker tensor product of two matrices: AxB = C
211: // Dyadic product, e.g. 2 x R^3 vectors span 6=0 space
212: // Full R^6 matrix has 3 components (ize product state vs entanglement)
213: template <typename T>
214: inline MMatrix<T> TensorProd(const MMatrix<T> &A, const MMatrix<T> &B) {
215:   const unsigned int rowA = A.size_row();
216:   const unsigned int rowB = B.size_row();
217:   const unsigned int colA = A.size_col();
218:   const unsigned int colB = B.size_col();
219:   MMatrix<T> C(rowA * rowB, colA * colB);
220:   for (std::size_t i = 0; i < rowA; ++i) {
221:     for (std::size_t j = 0; j < rowB; ++j) {
222:       C[i * rowA + j][k] = A[i][k] * B[j][k];
223:     }
224:   }
225:   return C;
226: }
227:
228: // Print out vector elements
229: template <typename T>
230: inline void PrintVector(const std::vector<T> &A, const std::string name = "") {
231:   std::cout << name << " : ";
232:   for (std::size_t i = 0; i < A.size(); ++i) {
233:     printf("%6.3f%6.3f%e", std::real(A[i]), std::imag(A[i]));
234:   }
235:   std::cout << std::endl;
236: }
237:
238: // Print out matrix elements
239: template <typename T>
240: inline void PrintMatrix(const MMatrix<T> &A, const std::string name = "") {
241:   std::cout << name << " : ";
242:   for (std::size_t i = 0; i < A.size_row(); ++i) {
243:     for (std::size_t j = 0; j < A.size_col(); ++j) {
244:       printf("%6.3f%6.3f%e", std::real(A[i][j]), std::imag(A[i][j]));
245:     }
246:     std::cout << std::endl;
247:   }
248:   std::cout << std::endl;
249: }
250:
251: // Functional methods for Spherical Harmonic Expansions
252:
253: // (c) 2017-2020 Mikael Mieskolainen
254: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
255: #define SPHERICAL_H
256: #define SPHERICAL_H
257:
258: // C++
259: #include <complex>
260: #include <iostream>
261: #include <vector>
262:
263: // Own
264: #include "Granitti/MSMath.h"
265: #include "Granitti/MMatrix.h"
266:
267: namespace gra {
268: namespace spherical {
269: // Metadata
270: struct Meta {
271:   std::string NAME; // Input name ID
272:   std::string LEGEND; // Legend string
273:   std::string VARS; // Variable string
274:   std::string MODE; // NC or DATA
275:   bool FASTSIM = false; // Fast simulation on
276:   std::string FRAME; // Locust frame
277:   double SCALE = 1.0; // Scale/normalization value
278: };
279:
280: std::vector<std::string> TITLES; // Phase space titles
281:
282: // So we can use this in std::map
283: bool operator<(const Meta &lhs) const { return lhs.NAME < rhs.NAME; }
284:
285: void Print() const {
286:   std::cout << "NAME: " << NAME << std::endl;
287:   std::cout << "LEGEND: " << LEGEND << std::endl;
288:   std::cout << "VARS: " << VARS << std::endl;
289:   std::cout << "MODE: " << MODE << std::endl;
290:   std::cout << "FRAME: " << FRAME << std::endl;
291:   std::cout << "FASTSIM: " << (FASTSIM ? "true" : "false") << std::endl;
292:   std::cout << "SCALE: " << SCALE << std::endl;
293: }
294:
295: std::cout << std::endl;
296: for (std::size_t i = 0; i < TITLES.size(); ++i) {
297:   printf("TITLES[%i] = %s\n", i, TITLES[i].c_str());
298: }
299: }
300:
301: // Microevent structure
302: struct Omega {
303:   // Decay daughter
304:   // rest frame variables
305:   double costheta = 0.0;
306:   double phi = 0.0;
307:
308:   // Invariant / system lab frame variables
309:   double M = 0.0;
310:   double Pt = 0.0;
311:   double Y = 0.0;
312:
313:   bool fiducial = false;
314:   bool selected = false;
315: };
316:
317: // Container
318: struct Data {
319:   // Metadata
320:   spherical::Meta META;
321:
322:   // Events
323:   std::vector<spherical::Omega> EVENTS;
324: };
325:
326: // Detector data for one hypercell, e.g., in (M, Pt, Y)
327: struct SH_DET {
328:   // Moment mixing matrix
329:   MMatrix<double> MIXM;
330:
331:   // Efficiency decomposition coefficients
332:   std::vector<double> E_m;
333:   std::vector<double> E_m_error;

```

```
./include/Granitti/MSpherical.h 2/2
84: }
85:
86: // Data for one hypercell, e.g., in (M, P, Y)
87: struct SH {
88:     // Directly (algebraic) observed moments
89:     std::vector<double> t_lm_MPF;
90:     std::vector<double> t_lm_MPF_error;
91:
92:     // Extended Maximum Likelihood fitted Moments
93:     std::vector<double> t_lm_EM;
94:     std::vector<double> t_lm_EM_error;
95: };
96:
97: MMatrix<double> GetMixing(const std::vector<Omega> &events, const std::vector<std::size_t> &ind,
98:     int lMAX, const std::string &mode);
99:
100: std::pair<std::vector<double>, std::vector<double>> GetEIM(const std::vector<Omega> & MC,
101:     const std::vector<std::size_t> &ind,
102:     int lMAX, const std::string &mode);
103:
104: std::vector<double> SphericalMoments(const std::vector<Omega> & input,
105:     const std::vector<std::size_t> &ind, int lMAX,
106:     const std::string &mode);
107:
108: MMatrix<double> YLM(const std::vector<Omega> &events, int lMAX);
109:
110: double HarmonicProd(const std::vector<double> &G, const std::vector<double> &ix,
111:     const std::vector<bool> &ACTIVE, int lMAX);
112:
113: void PrintOutMoments(const std::vector<double> &ix, const std::vector<double> &ix_error,
114:     const std::vector<bool> &ACTIVE, int lMAX);
115:
116: std::vector<std::size_t> GetIndices(const std::vector<Omega> &events, const std::vector<double> &IM,
117:     const std::vector<double> &fT, const std::vector<double> &V);
118:
119: void TestSphericalIntegrals(int lMAX);
120: int LinearEIM(int l, int N);
121: double CalcError(double f2, double f, double N);
122: void PrintMatrix(FILE *fp, const std::vector<std::vector<double>> &A);
123:
124: MMatrix<double> Y_real_synthesis(const std::vector<double> &G_lm, const std::vector<bool> &ACTIVE,
125:     std::size_t N, std::vector<double> &costheta,
126:     std::vector<double> &phi, bool normalized = false);
127:
128: std::vector<double> ErrorProp(const MMatrix<double> &A, const std::vector<double> &ix);
129:
130: // namespace spherical
131: // namespace gra
132: #endif
133:
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135:
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./include/Granitti/MMath.h 4/13
250: inline double sqrt(double x) { return std::sqrt(std::max(0.0, x)); }
251:
252: // Complex amplitude squared
253: inline double abs2(const std::complex<double> &M) { return pow2(std::abs(M)); }
254:
255: // ||v||^-2
256: template <typename T>
257: inline double vpow2(const std::vector<T> &v) {
258:     double sum = 0.0;
259:     for (std::size_t i = 0; i < v.size(); ++i) { sum += abs2(v[i]); }
260:     return sum;
261: }
262:
263: // ||v||
264: template <typename T>
265: inline double vnorm(const std::vector<T> &v) {
266:     double sum = vpow2(v);
267:     return sqrt(sum);
268: }
269:
270: // Template print
271: template <template <typename T> class container_type, class value_type>
272: inline void PrintArray(container_type<value_type> &x, std::string name) {
273:     std::cout << "PrintArray: " << name << std::endl;
274:     for (unsigned int i = 0; i < x.size(); ++i) { std::cout << x[i]; }
275:     std::cout << std::endl << std::endl;
276: }
277:
278: // Degrees to radians
279: constexpr double Deg2Rad(double deg) { return (deg / 180.0) * gra::math::PI; }
280:
281: // Radians to degrees
282: constexpr double Rad2Deg(double rad) { return (rad * 180.0) / gra::math::PI; }
283:
284: // Binomial Coefficient C(n, k)
285: constexpr int Chinom(int n, int k) {
286:     if (k == 0 || k == n) { return 1; }
287: }
288:
289: // Recursive function
290: return Chinom(n - 1, k - 1) + Chinom(n - 1, k);
291: }
292:
293: // N-dim epsilon tensor e_{l1, l2, ..., lN}
294: // + 1 if even permutation of arguments
295: // - 1 if odd permutation of arguments
296: // 0 otherwise
297: inline MTensor<int> epsTensor(std::size_t N) {
298:     // Maximum range value for each for-loop
299:     const std::size_t MAX = N;
300:
301:     // These hold for-loop index for each nested for loop
302:     std::vector<std::size_t> ind(N, 0);
303:
304:     // Permutation tensor definition
305:     auto permutation = &lt;int>();
306:     int value = 1;
307:     for (std::size_t i = 0; i < N; ++i) {
308:         for (std::size_t j = i + 1; j < N; ++j) {
309:             // Even permutation +1, Odd permutation -1, Otherwise 0
310:             if (ind[i] > ind[j]) {
311:                 value = -value;
312:             } else if (ind[i] == ind[j]) {
313:                 return 0;
314:             }
315:         }
316:     }
317:     return value;
318: }
319:
320: // -----
321: const std::vector<std::size_t> dimensions(N, MAX);
322: MTensor<int> T(dimensions, ind(0));
323:
324: // Nested for-loop
325: std::size_t index = 0;
326: while (true) {
327:     // -----
328:     // Evaluate the function value
329:     T(ind) = permutation();
330:     // -----
331:     ind[0]++;
332: }
```

```
./include/Granitti/MMath.h 6/13
416: T i = 3 * hstep * f(0) + 3 * T_first + 2 * T_second + f(3 * N) / 8;
417:
418: return i;
419: }
420:
421: // -----
422: // 2D-Simpson's 1/3 rule, see functions below
423: // -----
424: // f, W with dim[N+1, N+1] where N,N are even
425: // hstepM, hstepM are discretization step sizes (b-a)/M, (d-c)/N
426: // -----
427: // Based on Fubini's theorem <https://en.wikipedia.org/wiki/Fubini%27s_theorem>
428:
429: template <typename T>
430: inline T SimpsonIntegral2D(const MMatrix<T> &f, const MMatrix<double> &W, double hstepM,
431:     double hstepN) {
432:     if (((f.size_row() - 1) & 2 != 0) || ((f.size_col() - 1) & 2 != 0)) { // Must be even
433:         std::string str =
434:             "FATAL ERROR: gra::math::SimpsonIntegral2D,M,N must be "
435:             "even!";
436:         throw std::invalid_argument(str);
437:     }
438:
439:     T i = 0.0;
440:     for (std::size_t i = 0; i < f.size_row(); ++i) {
441:         for (std::size_t j = 0; j < f.size_col(); ++j) { i += W[i][j] * f[i][j]; }
442:     }
443:     i *= (hstepM * hstepN) / 9;
444:
445:     return i;
446: }
447:
448: // -----
449: // 2D-Simpson's 3/8 rule, see functions below
450: // -----
451: // f, W with dim[N+1, N+1] where N,N are multiple of 3
452: // hstepM, hstepM are discretization step sizes (b-a)/M, (d-c)/N
453: // -----
454: // Based on Fubini's theorem <https://en.wikipedia.org/wiki/Fubini%27s_theorem>
455:
456: template <typename T>
457: inline T Simpson3Integral2D(const MMatrix<T> &f, const MMatrix<double> &W, double hstepM,
458:     double hstepN) {
459:     if (((f.size_row() - 1) & 3 != 0) || ((f.size_col() - 1) & 3 != 0)) {
460:         std::string str =
461:             "FATAL ERROR: gra::math::Simpson3Integral2D,M,N must be "
462:             "multiple of 3!";
463:         throw std::invalid_argument(str);
464:     }
465:
466:     T i = 0.0;
467:     for (std::size_t i = 0; i < f.size_row(); ++i) {
468:         for (std::size_t j = 0; j < f.size_col(); ++j) { i += W[i][j] * f[i][j]; }
469:     }
470:     i *= (hstepM * hstepN) / 64; // 3^2, 3^2
471:
472:     return i;
473: }
474:
475: // Example with N = 6, N = 8:
476: // 1 4 2 4 2 4 2 4 2
477: // 4 16 8 16 8 16 8 16 4
478: // 2 8 4 8 4 8 4 8 2
479: // 4 16 8 16 8 16 8 16 4
480: // 2 8 4 8 4 8 4 8 2
481: // 4 16 8 16 8 16 8 16 4
482: // 1 4 2 4 2 4 2 4 2
483:
484: // Simpson's 1/3 rule weight vector dim[N+1]
485: inline std::vector<double> Simpson1Weight(unsigned int N) {
486:     if (N & 2 != 0) { // Must be even
487:         std::string str =
488:             "FATAL ERROR: gra::math::SimpsonWeight,N = " + std::to_string(N) + " is not even!";
489:         throw std::invalid_argument(str);
490:     }
491:
492:     std::vector<double> W(N + 1, 2.0); // Init with 2
493:
494:     // Set j as the first and the last
495:     W[0] = 1.0;
496:     W[N] = 1.0;
497:
498:     for (std::size_t j = 1; j < N; j += 2) { W[j] = 4.0; }
499:     return W;
500: }
```

```
./include/Granitti/MMath.h 5/13
333:
334: // Carry
335: while (ind[index] == MAX) {
336:     if (index == N - 1) { return T; }
337:
338:     ind[index++] = 0;
339:     ind[indindex]++;
340: }
341: index = 0;
342: }
343: }
344:
345: // Composite trapezoidal integral:
346: // Has geometric (fast) convergence for periodic functions:
347: // - a circle in complex plane
348: // - interval on real line
349: // - Hankel contour around (-inf, 0)
350:
351: // hstep = (b-a)/N is the discretization size
352:
353: // Indexing: j = 0, 1, ..., N-1, N, => length of f = N+1, N = even
354:
355: template <typename T>
356: inline T CTrapezIntegral(const std::vector<T> &f, double hstep) {
357:     if ((f.size() - 1) & 2 != 0) { // Must be even
358:         std::string str = "FATAL ERROR: gra::math::CTrapezIntegral,N = " + std::to_string(f.size() - 1) +
359:             " is not even!";
360:         throw std::invalid_argument(str);
361:     }
362:
363:     const std::size_t N = f.size() - 1;
364:     T sum = 0.0;
365:
366:     for (std::size_t j = 1; j <= N - 1; ++j) { sum += f[j]; }
367:     T i = hstep * (f[0] + 2 * sum + f[N]) / 2;
368:
369:     return i;
370: }
371:
372: // Composite Simpson's rule integral (quadratic interpolation),
373: // - hstep=(b-a)/N the discretization size
374: // -----
375: // Indexing: j = 0, 1, ..., N-1, N, => length of f = N+1, N = even
376:
377: template <typename T>
378: inline T CSIntegral(const std::vector<T> &f, double hstep) {
379:     if ((f.size() - 1) & 2 != 0) { // Must be even
380:         std::string str =
381:             "FATAL ERROR: gra::math::CSIntegral,N = " + std::to_string(f.size() - 1) + " is not even!";
382:         throw std::invalid_argument(str);
383:     }
384:
385:     const std::size_t N = f.size() - 1;
386:     T sum = 0.0;
387:     T i_odd = 0.0;
388:
389:     for (std::size_t j = 1; j <= N / 2 - 1; ++j) { sum += f[2 * j]; }
390:     for (std::size_t j = 1; j <= N / 2; ++j) { i_odd += f[2 * j - 1]; }
391:     T i = hstep * (f[0] + 2.0 * sum + 4.0 * i_odd + f[N]) / 3.0;
392:
393:     return i;
394: }
395:
396: // Composite Simpson's 3/8 rule integral (cubic interpolation),
397: // - hstep=(b-a)/3N the discretization size
398: // -----
399: // Indexing: j = 0, 1, ..., 3N, => length of f = 3N+1
400:
401: template <typename T>
402: inline T CS3Integral(const std::vector<T> &f, double hstep) {
403:     if ((f.size() - 1) & 3 != 0) { // Must be multiple of 3
404:         std::string str = "FATAL ERROR: gra::math::CS3Integral,N = " +
405:             std::to_string(f.size() - 1) / 3 + " is not multiple of 3!";
406:         throw std::invalid_argument(str);
407:     }
408:
409:     const std::size_t N = (f.size() - 1) / 3;
410:
411:     T i_first = 0.0;
412:     T i_second = 0.0;
413:
414:     for (std::size_t j = 1; j <= N; ++j) { i_first += f(3 * j - 2) + f(3 * j - 1); }
415:     for (std::size_t j = 1; j <= N - 1; ++j) { i_second += f(3 * j); }
416: }
```

```
./include/Granitti/MMath.h 7/13
499:
500: // Simpson's 3/8 rule weight vector dim[N+1]
501: inline std::vector<double> Simpson3Weight(unsigned int N) {
502:     if (N & 3 != 0) { // Must be multiple of 3
503:         std::string str =
504:             "FATAL ERROR: gra::math::SimpsonWeight,N = " + std::to_string(N) + " is not multiple of 3!";
505:         throw std::invalid_argument(str);
506:     }
507:
508:     std::vector<double> W(N + 1, 3.0); // Init with 3
509:
510:     // Set j as the first and the last
511:     W[0] = 1.0;
512:     W[N] = 1.0;
513:
514:     for (std::size_t i = 3; i < N; i += 3) { W[i] = 2.0; }
515:
516:     return W;
517: }
518:
519: // Create a weight matrix dim[N+1, N+1] for 2D-Simpson's 1/3 rule
520: inline MMatrix<double> Simpson1Weight2D(unsigned int M, unsigned int N) {
521:     MMatrix<double> W(M + 1, N + 1, 0.0);
522:     const std::vector<double> U = Simpson1Weight(M);
523:     const std::vector<double> V = Simpson3Weight(N);
524:
525:     // Outer (tensor) product: W = U*V^T
526:     // print("u,v");
527:     for (std::size_t i = 0; i <= M; ++i) {
528:         for (std::size_t j = 0; j <= N; ++j) { W[i][j] = U[i] * V[j]; }
529:     }
530:
531:     return W;
532: }
533:
534: // Create a weight matrix dim[N+1, N+1] for 2D-Simpson's 3/8 rule
535: inline MMatrix<double> Simpson3Weight2D(unsigned int M, unsigned int N) {
536:     MMatrix<double> W(M + 1, N + 1, 0.0);
537:     const std::vector<double> U = Simpson3Weight(M);
538:     const std::vector<double> V = Simpson3Weight(N);
539:
540:     // Outer (tensor) product: W = U*V^T
541:     // print("u,v");
542:     for (std::size_t i = 0; i <= M; ++i) {
543:         for (std::size_t j = 0; j <= N; ++j) { W[i][j] = U[i] * V[j]; }
544:     }
545:
546:     return W;
547: }
548:
549: // Construct binary reflect Gray code (BRGC)
550: // which is the Hamiltonian Path on a N-dim unit hypercube (2^N Boolean vector
551: // of space)
552: // The operator ^ is Exclusive OR (XOR) and the operator >> is bit shift right
553: constexpr unsigned int BinaryGray(unsigned int number) { return number ^ (number >> 1); }
554:
555: // Convert BRGC (binary reflect Gray code) to a binary number
556: // Gray code is obtained as XOR of all more significant bits
557: constexpr unsigned int GrayBinary(unsigned int number) {
558:     unsigned int mask = number >> 1; mask = mask >> 1; (number = number ^ mask);
559:     return number;
560: }
561:
562: // Boolean vector to index representation
563: // [the normal order: 0^0, 1^0, 1^1, 0^1, 2^0, 2^1, 3^0, 3^1, ...]
564: inline int Vec2Ind(std::vector<bool> &vec) {
565:     // Swap bit direction (comment this for the reverse ordering)
566:     std::reverse(vec.begin(), vec.end());
567:
568:     int retval = 0;
569:     int i = 0;
570:
571:     for (std::vector<bool>::iterator it = vec.begin(); it != vec.end(); ++it, ++i) {
572:         if (*it) { retval |= 1 << i; }
573:     }
574:
575:     return retval;
576: }
577:
578: // Index to Boolean vector representation
579: // [the normal order: 0^0, 1^0, 1^1, 0^1, 2^0, 2^1, 3^0, 3^1, ...]
580: // Input: x = index representation
581: // d = Boolean vector space dimension
582: inline std::vector<bool> Ind2Vec(unsigned int ind, unsigned int d) {
583:     std::vector<bool> ret;
584:
585:     while (ind) {
586:         ret.push_back(ind & 1);
587:         ind >>= 1;
588:     }
589:
590:     while (ret.size() < d) { ret.push_back(0); }
591:
592:     return ret;
593: }
```



```

./include/Granitti/MMath.h 12/13
914: )
915: }
916:
917: // Manually written Spherical Harmonics to cross-check algorithmic
918: // implementations
919: inline std::complex Y_complex_basis_ref(double costheta, double phi, int l, int m) {
920:     const double theta = std::acos(costheta);
921:     922:     if (l == 0 && m == 0) return 0.5 * std::sqrt(1.0 / PI);
923:     924:     // -----
925:     if (l == 1 && m == -1)
926:         return 0.5 * std::sqrt(3.0 / (2.0 * PI)) * std::exp(-gra::math::i * phi) * std::sin(theta);
927:     928:     if (l == 1 && m == 0) return 0.5 * std::sqrt(3.0 / PI) * costheta;
929:     930:     if (l == 1 && m == 1)
931:         return -0.5 * std::sqrt(3.0 / (2.0 * PI)) * std::exp(gra::math::i * phi) * std::sin(theta);
932:     933:     // -----
934:     if (l == 2 && m == -2)
935:         return 0.25 * std::sqrt(15.0 / (2.0 * PI)) * std::exp(-2.0 * gra::math::i * phi) *
936:             std::pow(std::sin(theta), 2);
937:     938:     if (l == 2 && m == -1)
939:         return 0.5 * std::sqrt(15.0 / (2.0 * PI)) * std::exp(-gra::math::i * phi) * std::sin(theta) *
940:             costheta;
941:     942:     if (l == 2 && m == 0) return 0.25 * std::sqrt(5.0 / PI) * (3.0 * std::pow(costheta, 2) - 1.0);
943:     944:     if (l == 2 && m == 1)
945:         return -0.5 * std::sqrt(15.0 / (2.0 * PI)) * std::exp(gra::math::i * phi) * std::sin(theta) *
946:             costheta;
947:     948:     if (l == 2 && m == 2)
949:         return 0.25 * std::sqrt(15.0 / (2.0 * PI)) * std::exp(2.0 * gra::math::i * phi) *
950:             std::pow(std::sin(theta), 2);
951:     952:     // -----
953:     if (l == 3 && m == -3)
954:         return 1.0 / 8.0 * std::sqrt(35.0 / PI) * std::exp(-3.0 * gra::math::i * phi) *
955:             std::pow(std::sin(theta), 3);
956:     957:     if (l == 3 && m == -2)
958:         return 0.25 * std::sqrt(105.0 / (2.0 * PI)) * std::exp(-2.0 * gra::math::i * phi) *
959:             std::pow(std::sin(theta), 2) * costheta;
960:     961:     if (l == 3 && m == -1)
962:         return 1.0 / 8.0 * std::sqrt(21.0 / PI) * std::exp(-gra::math::i * phi) * std::sin(theta) *
963:             (5.0 * std::pow(costheta, 2) - 1.0);
964:     965:     if (l == 3 && m == 0)
966:         return 0.25 * std::sqrt(7.0 / PI) * (5.0 * std::pow(costheta, 3) - 3.0 * costheta);
967:     968:     if (l == 3 && m == 1)
969:         return -1.0 / 8.0 * std::sqrt(21.0 / PI) * std::exp(gra::math::i * phi) * std::sin(theta) *
970:             (5.0 * std::pow(costheta, 2) - 1.0);
971:     972:     if (l == 3 && m == 2)
973:         return 0.25 * std::sqrt(105.0 / (2.0 * PI)) * std::exp(2.0 * gra::math::i * phi) *
974:             std::pow(std::sin(theta), 2) * costheta;
975:     976:     if (l == 3 && m == 3)
977:         return -1.0 / 8.0 * std::sqrt(35.0 / PI) * std::exp(3.0 * gra::math::i * phi) *
978:             std::pow(std::sin(theta), 3);
979:     980:     // -----
981:     if (l == 4 && m == -4)
982:         return 3.0 / 16.0 * std::sqrt(35.0 / (2.0 * PI)) * std::exp(-4.0 * gra::math::i * phi) *
983:             std::pow(std::sin(theta), 4);
984:     985:     if (l == 4 && m == -3)
986:         return 3.0 / 8.0 * std::sqrt(35.0 / (PI)) * std::exp(-3.0 * gra::math::i * phi) *
987:             std::pow(std::sin(theta), 3) * costheta;
988:     989:     if (l == 4 && m == -2)
990:         return 3.0 / 8.0 * std::sqrt(5.0 / (2.0 * PI)) * std::exp(-2.0 * gra::math::i * phi) *
991:             std::pow(std::sin(theta), 2) * (7.0 * std::pow(costheta, 2) - 1.0);
992:     993:     if (l == 4 && m == -1)
994:         return 3.0 / 8.0 * std::sqrt(5.0 / PI) * std::exp(-gra::math::i * phi) * std::sin(theta) *
995:             (7.0 * std::pow(costheta, 2) - 1.0);

```

```

./include/Granitti/MQED.h 1/1
1: // Running QED coupling (HEADER ONLY FILE)
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6:
7: #ifndef MQED_H
8: #define MQED_H
9:
10: #include <vector>
11:
12: // Own
13: #include "Granitti/MMath.h"
14:
15: namespace gra {
16: namespace qed {
17: namespace qed {
18:
19: constexpr double alpha_ref = 0.0072973525693; // reference value of alpha at scale Q
20: constexpr double Q_ref = 0.005109989461; // reference scale (GeV)
21: constexpr double alpha_0 =
22:     Q_ref * alpha_ref; // value of alpha at Q = 0 (Zine structure constant)[-1]
23:
24: // -----
25:
26: // Get number of charged leptons at given scale
27: inline int get_n_leptons(double Q) {
28:     if (Q > 1.77686E+03) { return 3; }
29:     if (Q > 1.056583745E+01) { return 2; }
30:     return 1;
31: }
32:
33: // The coefficient of one loop beta function
34: inline double beta_QED(int N_lf) { return -1.0 / (3.0 * math::PI) * N_lf; }
35:
36: // QED running coupling
37: inline double alpha_QED(double Q, const std::string& order = "LL") {
38:     // Freeze the coupling at Q < Q_ref
39:     if (order == "ZERO" || Q < Q_ref * Q_ref) { return alpha_0; }
40:
41:     // Input
42:     const double Q = math::msqrt(Q);
43:     const int N_lf = get_n_leptons(Q);
44:
45:     // Beta-function quantities
46:     const double beta0 = beta_QED(N_lf);
47:     const double R = std::log(Q2 / (Q_ref * Q_ref));
48:
49:     if (order == "LL") { // leading log [one loop geometric series resummation]
50:         return alpha_ref / (1.0 + alpha_ref * beta0 * R);
51:     } else {
52:         throw std::invalid_argument("alpha_QED: Unknown order parameter: " + order);
53:     }
54: }
55:
56: // QED coupling at Q = 0
57: inline double alpha_QED0() { return alpha_QED(0, "ZERO"); }
58:
59: // QED: Electric charge in natural units ~ 0.3
60: inline double e_QED(double Q2) { return math::msqrt(alpha_QED(Q2) * 4.0 * math::PI); }
61: inline double e_QED0() { return math::msqrt(alpha_QED0() * 4.0 * math::PI); }
62:
63:
64: } // namespace qed
65: } // namespace gra
66:
67: #endif

```

```

./include/Granitti/MMath.h 13/13
997: (7.0 * std::pow(costheta, 3) - 3.0 * costheta);
998:
999: if (l == 4 && m == 0)
1000:     return 3.0 / 16.0 * std::sqrt(1.0 / PI) *
1001:         (35.0 * std::pow(costheta, 4) - 30.0 * std::pow(costheta, 2) + 3.0);
1002:
1003: if (l == 4 && m == 1)
1004:     return -2.0 / 8.0 * std::sqrt(5.0 / PI) * std::exp(gra::math::i * phi) * std::sin(theta) *
1005:         (7.0 * std::pow(costheta, 3) - 3.0 * costheta);
1006:
1007: if (l == 4 && m == 2)
1008:     return 3.0 / 8.0 * std::sqrt(5.0 / (2.0 * PI)) * std::exp(2.0 * gra::math::i * phi) *
1009:         std::pow(std::sin(theta), 2) * (7.0 * std::pow(costheta, 2) - 1.0);
1010:
1011: if (l == 4 && m == 3)
1012:     return -2.0 / 8.0 * std::sqrt(35.0 / (PI)) * std::exp(3.0 * gra::math::i * phi) *
1013:         std::pow(std::sin(theta), 3) * costheta;
1014:
1015: if (l == 4 && m == 4)
1016:     return 3.0 / 16.0 * std::sqrt(35.0 / (2.0 * PI)) * std::exp(4.0 * gra::math::i * phi) *
1017:         std::pow(std::sin(theta), 4);
1018:
1019: throw std::invalid_argument("Y_complex_basis_ref: Not supported l = " + std::ito_string(l) +
1020:     ", m = " + std::ito_string(m));
1021: }
1022:
1023: } // namespace math
1024: } // namespace gra
1025:
1026: #endif

```

```

./include/Granitti/MTensorPomeron.h 1/4
1: // Tensor Pomeron amplitudes
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: #ifndef MTENSORPOMERON_H
7: #define MTENSORPOMERON_H
8:
9: // C++
10: #include <complex>
11: #include <random>
12: #include <vector>
13:
14: // Tensor algebra
15: #include "TTensor.hpp"
16:
17: // Own
18: #include "Granitti/MVec.h"
19: #include "Granitti/MGSObs.h"
20: #include "Granitti/MGSObs.h"
21: #include "Granitti/MMinimizations.h"
22: #include "Granitti/MMath.h"
23: #include "Granitti/MTensor.h"
24:
25: namespace gra {
26: namespace MTensorPomeronParam {
27:
28: struct MTensorPomeronParam {
29:     // Trajectory parameters
30:     double delta_P = 0;
31:     double delta_O = 0;
32:     double delta_IR = 0;
33:     double delta_2R = 0;
34:     double delta_2R = 0;
35:
36:     double ap_P = 0; // GeV[-2]
37:     double ap_O = 0; // GeV[-2]
38:     double ap_IR = 0; // GeV[-2]
39:     double ap_2R = 0; // GeV[-2]
40:
41:     double eta_O = 1; // +- 1
42:
43:     // Scales
44:     double M_O = 0; // GeV
45:     double M_IR = 0; // GeV
46:
47:     // Couplings
48:     double gPNN = 0; // GeV[-1], Pomeron-Proton-Proton
49:     double gPpP = 0; // GeV[-1], Pomeron-Pion-Pion
50:     double gPRK = 0; // GeV[-1], Pomeron-Rhoon-Rhoon
51:     std::vector<double> gPRrho; // GeV[-3], GeV[-1], Pomeron-Rho-Rho
52:     std::vector<double> gPpphpi; // GeV[-3], GeV[-1], Pomeron-Phi-Phi
53:
54:     bool initialized = false;
55:
56:     // Read parameters from file
57:     void ReadParameters(const std::string& modelfile) {
58:         using json = nlohmann::json;
59:         const std::string data = gra::aux::GetInputData(modelfile);
60:         json j;
61:
62:         try {
63:             j = json::parse(data);
64:         }
65:
66:         // JSON block identifier
67:         const std::string XID = "PARAM_TENSORPOM";
68:
69:         delta_P = 3.at(XID).at("delta_P");
70:         delta_O = 3.at(XID).at("delta_O");
71:         delta_IR = 3.at(XID).at("delta_IR");
72:         delta_2R = 3.at(XID).at("delta_2R");
73:
74:         ap_P = 3.at(XID).at("ap_P");
75:         ap_O = 3.at(XID).at("ap_O");
76:         ap_IR = 3.at(XID).at("ap_IR");
77:         ap_2R = 3.at(XID).at("ap_2R");
78:
79:         eta_O = 3.at(XID).at("eta_O");
80:         M_O = 3.at(XID).at("M_O");
81:         M_IR = 3.at(XID).at("M_IR");
82:
83:         gPNN = 3.at(XID).at("gPNN");
84:         gPpP = 3.at(XID).at("gPpP");

```

```

./include/Granitti/MTensorPomeron.h      2/4
84:  gPFK = j.at(KID).at("gPFK");
85:
86:  std::vector<double> a = j.at(KID).at("gPrho");
87:  gPrho = a;
88:
89:  std::vector<double> b = j.at(KID).at("gPhiphi");
90:  gPhiphi = b;
91:
92:  initialized = true;
93: } catch (...) {
94:   std::string str = "MTensorPomeronParam: ReadParameters: Error parsing " + modelfile +
95:     " (Check for extra/missing commas)";
96:   throw std::invalid_argument(str);
97: }
98: }
99: }
100:
101:
102: // Matrix element dimension: * GeV^4 << -(2*external_legs - 8)
103: class MTensorPomeron : public MDIrac {
104: public:
105:   MTensorPomeron(gra:LorentzSCALAR k1ts, const std::string smodelfile);
106:   ~MTensorPomeron() {}
107:
108:   // Decay coupling (static so we can call it independently)
109:   static double GDecay(int J, double M, double Gamma, double mf, double BR);
110:
111:   // Amplitude squared
112:   double ME3(gra:LorentzSCALAR k1ts) const;
113:   double ME4(gra:LorentzSCALAR k1ts) const;
114:   double ME6(gra:LorentzSCALAR k1ts) const;
115:
116:   // Scalar, Pseudoscalar, Tensor coupling structures
117:   FTensor<std::complex<double>, 4, 4, 4, 4> iG_PFS_0(const M4Vec k1, const M4Vec k2,
118:     double g_PFS) const;
119:   FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_PFS_0(const M4Vec k1, const M4Vec k2,
120:     double g_PFS) const;
121:   FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_PFS_1(const M4Vec k1, const M4Vec k2,
122:     double g_PFS) const;
123:   FTensor<std::complex<double>, 4, 4, 4, 4> iG_PFS_1(const M4Vec k1, const M4Vec k2,
124:     double g_PFS) const;
125:   MTensor<std::complex<double>, 4, 4, 4, 4> iG_PFT_12(const M4Vec k1, const M4Vec k2, double g_PFT,
126:     int mode) const;
127:   MTensor<std::complex<double>, 4, 4, 4, 4> iG_PFT_03(const M4Vec k1, const M4Vec k2, double g_PFT) const;
128:   MTensor<std::complex<double>, 4, 4, 4, 4> iG_PFT_04(const M4Vec k1, const M4Vec k2, double g_PFT) const;
129:   MTensor<std::complex<double>, 4, 4, 4, 4> iG_PFT_05(const M4Vec k1, const M4Vec k2, double g_PFT) const;
130:   MTensor<std::complex<double>, 4, 4, 4, 4> iG_PFT_06(const M4Vec k1, const M4Vec k2, double g_PFT) const;
131:
132:   // Vertex functions
133:   FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_vv2pps(const std::vector<M4Vec> k,
134:     int PDG) const;
135:
136:   FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_PFS_Total(
137:     const M4Vec k1, const M4Vec k2, double M0, const std::string smode,
138:     const std::vector<double> g_PFS) const;
139:   MTensor<std::complex<double>, 4, 4, 4, 4> iG_PFT_Total(const M4Vec k1, const M4Vec k2, double M0,
140:     const std::vector<double> g_PFT) const;
141:
142:   FTensor<Tensor<std::complex<double>, 4, 4> iG_PppH(const M4Vec kprime, const M4Vec p) const;
143:
144:   FTensor<Tensor<std::complex<double>, 4> iG_vee(
145:     const M4Vec kprime, const M4Vec k, const std::vector<std::complex<double>> sbar,
146:     const std::vector<std::complex<double>> s) const;
147:
148:   FTensor<Tensor<std::complex<double>, 4, 4> iG_yeehbar(
149:     const std::vector<std::complex<double>> sbar, const MMatrix<std::complex<double>> sISF,
150:     const std::vector<std::complex<double>> sv) const;
151:
152:   FTensor<Tensor<std::complex<double>, 4> iG_ypp(
153:     const M4Vec kprime, const M4Vec k, const std::vector<std::complex<double>> sbar,
154:     const std::vector<std::complex<double>> s) const;
155:
156:   FTensor<Tensor<std::complex<double>, 4, 4> iG_Ppp(
157:     const M4Vec kprime, const M4Vec k, const std::vector<std::complex<double>> sbar,
158:     const std::vector<std::complex<double>> s) const;
159:
160:   FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_PpphBar(
161:     const M4Vec kprime, const std::vector<std::complex<double>> sbar, const M4Vec k,
162:     const MMatrix<std::complex<double>> sISF, const std::vector<std::complex<double>> sv,
163:     const M4Vec k) const;
164:
165:   FTensor<Tensor<std::complex<double>, 4, 4> iG_Ppps(const M4Vec kprime, const M4Vec k,
166:     double g) const;

```

```

./include/Granitti/MTensorPomeron.h      4/4
250: FTensor<Tensor<double, 4, 4, 4, 4> R_DDDU;
251: FTensor<Tensor<double, 4, 4, 4, 4> R_UUDD;
252:
253: MTensor<double> T1;
254: MTensor<double> T2;
255: MTensor<double> T3;
256:
257: // Parameters
258: MTensorPomeronParam Param;
259: }
260:
261: } // namespace gra
262:
263: #endif

```

```

./include/Granitti/MTensorPomeron.h      3/4
167: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_Pvv(const M4Vec kprime, const M4Vec k,
168:   double g1, double g2) const;
169:
170: std::complex<double> iG_f0ss(const M4Vec k1, const M4Vec k2, double g1) const;
171: FTensor<Tensor<std::complex<double>, 4, 4> iG_f0vv(const M4Vec k1, const M4Vec k2, double M0,
172:   double g1, double g2) const;
173:
174: FTensor<Tensor<std::complex<double>, 4, 4> iG_pvv(const M4Vec k1, const M4Vec k2, double M0,
175:   double g1) const;
176: FTensor<Tensor<std::complex<double>, 4> iG_ypps(const M4Vec k1, const M4Vec k2, double M0,
177:   double g1) const;
178:
179: FTensor<Tensor<std::complex<double>, 4, 4> iG_f2pps(const M4Vec k1, const M4Vec k2,
180:   double M0, double g1) const;
181: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_f2vv(const M4Vec k1, const M4Vec k2,
182:   double M0, double g1, double g2) const;
183: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iG_f2yy(const M4Vec k1, const M4Vec k2,
184:   double M0, double g1, double g2) const;
185:
186: FTensor<Tensor<std::complex<double>, 4, 4> iG_vY(double g2, int pdg) const;
187:
188: // Polarization sums
189: std::vector<FTensor<Tensor<std::complex<double>, 4, 4> MassiveSpinPolSum(
190:   const M4Vec k1, const M4Vec k2, const M4Vec k3,
191:   const M4Vec k4) const;
192: std::vector<FTensor<Tensor<std::complex<double>, 4, 4> MasslessSpinPolSum(
193:   const M4Vec k1, const M4Vec k2,
194:   const M4Vec k3) const;
195:
196: std::vector<std::complex<double>> MasslessSpinPolSum(
197:   const FTensor<Tensor<std::complex<double>, 4, 4> kM, const M4Vec k1,
198:   const M4Vec k2) const;
199: std::vector<std::complex<double>> MassiveSpinPolSum(
200:   const FTensor<Tensor<std::complex<double>, 4, 4> kM, const M4Vec k1,
201:   const M4Vec k2) const;
202:
203: // Propagators
204: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iD_F(double s, double t) const;
205: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iD_2R(double s, double t) const;
206: FTensor<Tensor<std::complex<double>, 4, 4> iD_0(double s, double t) const;
207: FTensor<Tensor<std::complex<double>, 4, 4> iD_1R(double s, double t) const;
208:
209: FTensor<Tensor<std::complex<double>, 4, 4> iD_V(const M4Vec k, double M0, double s14) const;
210: std::complex<double> iD_MES(const M4Vec k, double M0) const;
211: std::complex<double> iD_MES(const M4Vec k, double M0, double Gamma) const;
212: FTensor<Tensor<std::complex<double>, 4, 4> iD_YMES(const M4Vec k, double M0, double Gamma,
213:   bool INDEX_UP) const;
214: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> iD_TME3(const M4Vec k, double M0,
215:   double Gamma, bool INDEX_UP) const;
216:
217: // Tensor functions
218: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> Gamma0(const M4Vec k1, const M4Vec k2) const;
219: FTensor<Tensor<std::complex<double>, 4, 4, 4, 4> Gamma2(const M4Vec k1, const M4Vec k2) const;
220: void CalcRTensor();
221:
222: // Trajectories
223: double alpha_P(double t) const;
224: double alpha_0(double t) const;
225: double alpha_1R(double t) const;
226: double alpha_2R(double t) const;
227:
228: // Form factors
229: double F1(double t) const;
230: double F2(double t) const;
231:
232: double F1_1(double t) const;
233: double F2_1(double t) const;
234:
235: double GD(double t) const;
236: double FM(double g2, double LAMBDA) const;
237: double F(double g2, double M0, double LAMBDA) const;
238:
239: private:
240: // Hukawaki metric tensor
241: FTensor<Tensor<double, 4, 4> gT;
242:
243: // Kpallon tensors
244: FTensor<Tensor<double, 4, 4, 4, 4> eps_1a;
245: FTensor<Tensor<double, 4, 4, 4, 4> eps_1b;
246:
247: // Aux tensors
248: FTensor<Tensor<double, 4, 4, 4, 4> R_DDDD;
249: FTensor<Tensor<double, 4, 4, 4, 4> R_DDDU;

```

D.4 C++ source files


```

./src/MUserHistograms.cc 1/2
1: // Container class for different type of histograms
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <iostream>
9: #include <vector>
10:
11: // Own
12: #include "GrantiLL/Muon.h"
13: #include "GrantiLL/Muon.h"
14: #include "GrantiLL/Muonematics.h"
15: #include "GrantiLL/Muon.h"
16: #include "GrantiLL/Muon.h"
17: #include "GrantiLL/MUserHistograms.h"
18:
19: using gra::aux::indices;
20: using gra::math::PI;
21:
22: namespace gra {
23: void MUserHistograms::InitHistograms() {
24: // Level 1
25: unsigned int Nbins = 40;
26:
27: h1["M"] = MH1<double>(Nbins, "Central System M (GeV)");
28: h1["Rap"] = MH1<double>(Nbins, "Central System Rap");
29: h1["Rap"].SetAxisSymmetry(true);
30: h1["Pt"] = MH1<double>(Nbins, 0.0, 2.5, "Central System Pt (GeV)");
31: h1["dPhi_pp"] = MH1<double>(Nbins, 0.0, gra::math::PI, "Forward dPhi (rad)");
32: h1["pT1"] = MH1<double>(Nbins, 0.0, 2.0, "Forward Pt (GeV)");
33: h2["rapRap2"] = MH2(Nbins, Nbins, "Rapidity vs Rapidity");
34: h2["rapRap2"].SetAxisSymmetry(true, true);
35:
36: h1["FM"] = MH1<double>(Nbins, "Forward M (GeV)");
37: h1["m0"] = MH1<double>(Nbins, "Intermediate daughter M (GeV)");
38:
39: // Level 2
40: Nbins = 40;
41:
42: h2["cosTheta_CS"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [CS [Collins-Soper frame]]");
43: h2["cosTheta_CS"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [CS [Collins-Soper frame]]");
44: h2["cosTheta_BK"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [BK [Rapidity frame]]");
45: h2["cosTheta_BK"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [BK [Rapidity frame]]");
46: h2["cosTheta_AB"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [AB [Anti-Helicity frame]]");
47: h2["cosTheta_AB"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [AB [Anti-Helicity frame]]");
48: h2["cosTheta_PG"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [PG [Pseudo-GJ frame]]");
49: h2["cosTheta_PG"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [PG [Pseudo-GJ frame]]");
50: h1["cosTheta_J"] = MH1<double>(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [GJ [Gottfried-Jackson frame]]");
51: h2["cosTheta_CM"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [CM [Direct system rest frame]]");
52: h2["cosTheta_CM"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [CM [Direct system rest frame]]");
53: h2["cosTheta_LA"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [LA [Laboratory frame]]");
54: h2["cosTheta_LA"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [LA [Laboratory frame]]");
55: h2["cosTheta_LA"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [LA [Laboratory frame]]");
56: h2["cosTheta_LA"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [LA [Laboratory frame]]");
57: }
58:
59: // Input as the total event weight
60: void MUserHistograms::FillHistograms(double totalweight, const gra::LORENTZSCALAR &its) {
61: // Level 1
62: if (HIST >= 1) {
63: h1["M"].Fill(its.math::msqrt(its.m2), totalweight);
64: h1["Rap"].Fill(its.y, totalweight);
65: h1["Pt"].Fill(its.pt, totalweight);
66: h1["dPhi_pp"].Fill(its.dphi(its.pfinal[1], DeltaPhi(its.pfinal[2])), totalweight);
67: h1["pT1"].Fill(its.pfinal[1].pt(), totalweight);
68:
69: // Cascade decay
70: if (its.decaytree[0].legs.size() > 0) { h1["m0"].Fill(its.decaytree[0].p4.M(), totalweight); }
71:
72: // Dissociated proton
73: if (its.pfinal[1].M() > 1.0) { h1["FM"].Fill(its.pfinal[1].M(), totalweight); }
74: }
75:
76: // Level 2
77: if (HIST >= 2) {
78: h2["rapRap2"].Fill(its.decaytree[0].p4.Rap(), its.decaytree[1].p4.Rap(), totalweight);
79: h2["cosTheta_CS"] = MH2(Nbins, -1.0, 1.0, Nbins, -PI, PI, "(cos theta, phi) [CS [Collins-Soper frame]]");
80: }
81: }
82:
83: // (cos theta, phi) of daughter in different Lorentz frames, input as the total

```

```

./src/MSudakov.cc 1/8
1: // Shuvaev PDF and Sudakov suppression class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <fstream>
9: #include <future>
10: #include <iostream>
11: #include <memory>
12: #include <vector>
13:
14: // C file processing
15: #include <sys/types.h>
16: #include <sys/types.h>
17: #include <unistd.h>
18:
19: // Own
20: #include "GrantiLL/Muon.h"
21: #include "GrantiLL/Muon.h"
22: #include "GrantiLL/MSudakov.h"
23: #include "GrantiLL/MTimer.h"
24:
25: // LHAPDF
26: #include "LHAPDF/LHAPDF.h"
27:
28: // Libraries
29: #include "rang.hpp"
30:
31: namespace gra {
32: using aux::indices;
33: using math::msqrt;
34: using math::pow2;
35: using math::izi;
36:
37: // Constructor
38: MSudakov::MSudakov() {}
39:
40: // Destructor
41: ~MSudakov() { delete PdfFits; }
42:
43: // Init
44: void MSudakov::Init(double_sqrts, const std::string& PDFSET, bool init_arrays) {
45: // First this
46: Numerics::ReadParameters();
47:
48: // Setup energy dependent boundaries
49: Numerics::sqrts = sqrts;
50: Numerics::M_MAX = sqrts;
51: Numerics::M_MIN = sqrts;
52:
53: // Collinear case: M^2 = x1x2s, let x1 = 1.0 gives minimum x2 = M^2/s
54: Numerics::x_MIN = math::pow2(Numerics::M_MIN / Numerics::sqrts);
55:
56: // InitLHAPDF("CT1016");
57: InitLHAPDF(PDFSET);
58:
59: if (init_arrays) { InitArrays(); }
60:
61:
62: void MSudakov::InitArrays() {
63: // Init Sudakov variables
64:
65: std::cout << "Initializing <Sudakov> array:" << std::endl;
66:
67: // q2,M
68: veto_sqrts = Numerics::sqrts; // FIRST THIS
69: veto_Set(0, "q2", Numerics::q2_MIN, Numerics::q2_MAX, Numerics::SHUV_N[0],
70: Numerics::SUDA_log_ON[0]);
71: veto_Set(1, "M", Numerics::M_MIN, Numerics::M_MAX, Numerics::SUDA_M[1], Numerics::SUDA_log_ON[1]);
72: veto_InitArray(); // Initialize (call last)
73:
74: const unsigned long hash = gra::aux::djb2hash(veto.GetHashString());
75: const std::string filename = gra::aux::GetBasePath(2) + "/" + "skonal/SUDA_" +
76: gra::aux::ToStdString(veto_sqrts, 0) + "_" + PDFSETNAME + "_" +
77: std::ito_string(hash);
78:
79: // Try to read pre-calculated
80: bool ok = veto.ReadArray(filename);
81:
82: while (tok) { // Problem, re-calculate
83: // Pointer to member function: ReturnType

```

```

./src/MUserHistograms.cc 2/2
84: // event weight
85: void MUserHistograms::FillCosThetaPhi(double totalweight, const gra::LORENTZSCALAR &its) {
86: // Central Trial status
87: std::vector<M4Vec> pf;
88: pf.push_back(its.decaytree[0].p4);
89: pf.push_back(its.decaytree[1].p4);
90:
91: // System 4-momentum
92: M4Vec X;
93: for (const auto &i : indices(its.decaytree)) { X += its.decaytree[i].p4; }
94:
95: // Choose Pseudo-Gottfried-Jackson beam direction (-1,1)
96: const int direction = 1;
97:
98:
99: {
100: // ** PREPARE LORENTZ TRANSFORMATION COMMON VARIABLES **
101: M4Vec p1boost; // beam1 particle boosted
102: M4Vec p2boost; // beam2 particle boosted
103: std::vector<M4Vec> pfbost; // central particles boosted
104: gra::kinematics::LorentzFramePrepare(pf, X, its.pbeam1, its.pbeam2, p1boost, p2boost, pfbost);
105:
106: // ** TRANSFORM TO DIFFERENT LORENTZ FRAMES **
107: std::vector<std::string> frametype = {"CS", "BK", "AB", "PG", "CM"};
108: for (const auto &k : indices(frametype)) {
109: // Transform and histogram
110: std::vector<M4Vec> pfout;
111: gra::kinematics::LorentzFrame(pfout, p1boost, p2boost, pfbost, frametype[k], direction);
112: h2["cosTheta_" + frametype[k]].Fill(std::cos(pfout[0].Theta()), pfout[0].Phi(),
113: totalweight);
114: // h2["Phi_" + frametype[k]].Fill(its.math::msqrt(its.m2), pfout[0].Phi(),
115: // totalweight);
116: }
117: }
118:
119: // -----
120: // Gottfried-Jackson frame
121: std::vector<M4Vec> pf5 = pf;
122: gra::kinematics::GJFrame(pf5, its.pfinal[0], direction, its.q1, its.q2, false);
123: h2["cosTheta_GJ"].Fill(std::cos(pf5[0].Theta()), pf5[0].Phi(), totalweight);
124:
125: // Laboratory frame
126: h2["cosTheta_LA"].Fill(std::cos(pf[0].Theta()), pf[0].Phi(), totalweight);
127: }
128:
129: // Print all histograms out
130: void MUserHistograms::PrintHistograms() {
131: if (HIST == 1) {
132: for (auto const & h1 : h1(x.first).Print());
133: }
134: if (HIST == 2) {
135: for (auto const & h2 : h2(x.first).Print());
136: }
137: }
138:
139: // namespace gra
140:
141: // (CppType:*) (ParameterTypes...)
142: std::pair<double, double> (MSudakov::f)(double, double) = &MSudakov::Sudakov_T;
143:
144: veto.WriteArray(filename, true);
145: ok = veto.ReadArray(filename);
146:
147:
148: // Init Shuvaev PDF variables
149:
150: std::cout << "Initializing <Shuvaev> array:" << std::endl;
151:
152: // q2,x
153: spdf_sqrts = Numerics::sqrts; // FIRST THIS
154: spdf_Set(0, "q2", Numerics::q2_MIN, Numerics::q2_MAX, Numerics::SHUV_N[0],
155: Numerics::SHUV_log_ON[0]);
156: spdf_Set(1, "x", Numerics::x_MIN, Numerics::x_MAX, Numerics::SHUV_M[1], Numerics::SHUV_log_ON[1]);
157: spdf_InitArray(); // Initialize (call last)
158:
159: const unsigned long hash = gra::aux::djb2hash(spdf.GetHashString());
160: const std::string filename = gra::aux::GetBasePath(2) + "/" + "skonal/SHUV_" +
161: gra::aux::ToStdString(spdf_sqrts, 0) + "_" + PDFSETNAME + "_" +
162: std::ito_string(hash);
163:
164: // Try to read pre-calculated
165: bool ok = spdf.ReadArray(filename);
166:
167: while (tok) { // Problem, re-calculate
168: // Pointer to member function: ReturnType
169:
170: // (CppType:*) (ParameterTypes...)
171: std::pair<double, double> (MSudakov::f)(double, double) = &MSudakov::Shuvaev_T;
172:
173: spdf.WriteArray(filename, true);
174: ok = spdf.ReadArray(filename);
175:
176: }
177:
178: // -----
179: // Finally
180: initialized = true;
181: std::cout << std::endl;
182:
183: if (Numerics::DEBUG) { TestPDF(); }
184:
185: // Return gluon_xg(q2) from LHAPDF
186: double MSudakov::xg_q2(double x, double q2) const {
187: const int pid = 21; // gluon
188:
189: return PdfFits->xFxQ2(pid, x, q2);
190: }
191:
192: catch (...) {
193: std::string str =
194: "MSudakov::xg_q2: Problem with x = " + std::ito_string(x) + ", q2 = " + std::ito_string(q2);
195: throw std::invalid_argument(str);
196: }
197:
198: // PDF access
199:
200: // [REFERENCE: LHAPDF6, arxiv.org/abs/1412.7420]
201: void MSudakov::InitLHAPDF(const std::string& pdfname) {
202: PDFSETNAME = pdfname;
203:
204: // LHAPDF init
205:
206: try {
207: PdfFits = LHAPDF::mkPDF(pdfname, 0);
208: } catch (...) {
209: ++init_trials;
210:
211: std::string str = "MSudakov::InitLHAPDF: Trials = " + std::ito_string(init_trials) +
212: " : Problem with reading a pdfset " + pdfname + " ";
213:
214: if (init_trials >= 2) { // Too many failures
215: throw std::invalid_argument(str);
216: }
217:
218: std::cout << str << std::endl;
219:
220: aux::AutoDownloadLHAPDF(pdfname);
221: InitLHAPDF(pdfname); // try again
222: }
223:
224: // PDF test routine

```

```

./src/MSudakov.cc 2/8
84: // (CppType:*) (ParameterTypes...)
85: std::pair<double, double> (MSudakov::f)(double, double) = &MSudakov::Sudakov_T;
86:
87: veto.WriteArray(filename, true);
88: ok = veto.ReadArray(filename);
89:
90:
91: // Init Shuvaev PDF variables
92:
93: std::cout << "Initializing <Shuvaev> array:" << std::endl;
94:
95: // q2,x
96: spdf_sqrts = Numerics::sqrts; // FIRST THIS
97: spdf_Set(0, "q2", Numerics::q2_MIN, Numerics::q2_MAX, Numerics::SHUV_N[0],
98: Numerics::SHUV_log_ON[0]);
99: spdf_Set(1, "x", Numerics::x_MIN, Numerics::x_MAX, Numerics::SHUV_M[1], Numerics::SHUV_log_ON[1]);
100: spdf_InitArray(); // Initialize (call last)
101:
102: const unsigned long hash = gra::aux::djb2hash(spdf.GetHashString());
103: const std::string filename = gra::aux::GetBasePath(2) + "/" + "skonal/SHUV_" +
104: gra::aux::ToStdString(spdf_sqrts, 0) + "_" + PDFSETNAME + "_" +
105: std::ito_string(hash);
106:
107: // Try to read pre-calculated
108: bool ok = spdf.ReadArray(filename);
109:
110: while (tok) { // Problem, re-calculate
111: // Pointer to member function: ReturnType
112:
113: // (CppType:*) (ParameterTypes...)
114: std::pair<double, double> (MSudakov::f)(double, double) = &MSudakov::Shuvaev_T;
115:
116: spdf.WriteArray(filename, true);
117: ok = spdf.ReadArray(filename);
118:
119: }
120:
121: // -----
122: // Finally
123: initialized = true;
124: std::cout << std::endl;
125:
126: if (Numerics::DEBUG) { TestPDF(); }
127:
128: // Return gluon_xg(q2) from LHAPDF
129: double MSudakov::xg_q2(double x, double q2) const {
130: const int pid = 21; // gluon
131:
132: return PdfFits->xFxQ2(pid, x, q2);
133: }
134:
135: catch (...) {
136: std::string str =
137: "MSudakov::xg_q2: Problem with x = " + std::ito_string(x) + ", q2 = " + std::ito_string(q2);
138: throw std::invalid_argument(str);
139: }
140:
141: // PDF access
142:
143: // [REFERENCE: LHAPDF6, arxiv.org/abs/1412.7420]
144: void MSudakov::InitLHAPDF(const std::string& pdfname) {
145: PDFSETNAME = pdfname;
146:
147: // LHAPDF init
148:
149: try {
150: PdfFits = LHAPDF::mkPDF(pdfname, 0);
151: } catch (...) {
152: ++init_trials;
153:
154: std::string str = "MSudakov::InitLHAPDF: Trials = " + std::ito_string(init_trials) +
155: " : Problem with reading a pdfset " + pdfname + " ";
156:
157: if (init_trials >= 2) { // Too many failures
158: throw std::invalid_argument(str);
159: }
160:
161: std::cout << str << std::endl;
162:
163: aux::AutoDownloadLHAPDF(pdfname);
164: InitLHAPDF(pdfname); // try again
165: }
166:
167: // PDF test routine

```

```
./src/MSudakov.cc 3/8
167: void MSudakov::TestPDF() const {
168:     const double MINLOGX = std::log10(Numerics.x_MIN);
169:     const double MAXLOGX = std::log10(Numerics.x_MAX);
170:     const int NX = 5; // Number of points - 1
171:     const double stepX = (MAXLOGX - MINLOGX) / NX;
172:
173:     const double MINLOGQ2 = std::log10(Numerics.q2_MIN);
174:     const double MAXLOGQ2 = std::log10(Numerics.q2_MAX);
175:     const int NQ2 = 5; // Number of points - 1
176:     const double stepQ2 = (MAXLOGQ2 - MINLOGQ2) / NQ2;
177:
178:     const double MINLOGM = std::log10(Numerics.M_MIN);
179:     const double MAXLOGM = std::log10(Numerics.M_MAX);
180:     const int NM = 5; // Number of points - 1
181:     const double stepM = (MAXLOGM - MINLOGM) / NM;
182:
183:     // Test loop
184:     for (std::size_t i = 0; i < NM + 1; ++i) {
185:         const double log1M = MINLOGM + i * stepM;
186:         const double M = std::exp(log1M, log10M);
187:
188:         printf("M = %0.1F GeV: |alpha_s(Q = M GeV) = %0.3F \n\n", M, PdfTrF-alpha2(M * M));
189:
190:         for (std::size_t j = 0; j < NX + 1; ++j) {
191:             const double log1X = MINLOGX + j * stepX;
192:             const double x = std::exp(log1X, log10X);
193:
194:             printf("x = %0.3E \n", x);
195:             for (std::size_t k = 0; k < NQ2 + 1; ++k) {
196:                 const double log1Q2 = MINLOGQ2 + k * stepQ2;
197:                 const double q2 = std::exp(log1Q2, log10Q2);
198:
199:                 // Normal gluon pdf
200:                 const double xf = xg_xQ2(x, q2);
201:
202:                 // Durham flux
203:                 const double hxf = fg_xQM(x, q2, M);
204:
205:                 printf(
206:                     "M = %0.3E, q2 = %0.2E, M = %0.1E: [gluon pdf: xg(x,q2), "
207:                     "Durham flux: fg(x,q2,M)] = (%0.2F,%0.2E) \n",
208:                     x, q2, M, xf, hxf);
209:             }
210:             std::cout << std::endl;
211:         }
212:         std::cout << std::endl;
213:     }
214: }
215:
216: // Access QCD coupling alpha_s(Q^2) from LHAPDF
217: double MSudakov::Alpha_Q2(double q2) const {
218:     try {
219:         return PdfTrF-alpha2(q2);
220:     } catch (...) {
221:         std::string str = "MSudakov::Alpha_Q2: Problem with q2 = " + std::to_string(q2);
222:         throw std::invalid_argument(str);
223:     }
224: }
225:
226: // Calculate differentiation dxy/dQ2 via "Richardson's extrapolation"
227: // f'(x) = (f(x) - f_0(x)) / (x - x_0)
228: // Requires 4 evaluations of densities
229:
230: // [REFERENCE: en.wikipedia.org/wiki/Richardson_extrapolation]
231: double MSudakov::dYf_xQ2_ext_Q2(double x, double q2) const {
232:     // Do not take h less than 1E-4 with 64-bit double
233:     const double h = 1E-4;
234:     const double hX2 = 2 * h * h;
235:
236:     // Calculate central differences
237:     const double DO_A = (xg_xQ2(x, q2 + h) - xg_xQ2(x, q2 - h)) / (2 * h);
238:     const double DO_B = (xg_xQ2(x, q2 + hX2) - xg_xQ2(x, q2 - hX2)) / (2 * hX2);
239:
240:     return (4.0 * DO_A - DO_B) / 3.0;
241: }
242:
243: // Durham flux (skewed gluon pdf)
244: // Shuvev transformed gluon pdf a Sudakov suppression
245: // f(x,q2) = f_0(x,q2) * exp(-int_0^x (alpha_s(t)/t) dt)
246: // f(x,q2) = f_0(x,q2) * exp(-int_0^x (alpha_s(t)/t) dt)
247: // f(x,q2) = f_0(x,q2) * exp(-int_0^x (alpha_s(t)/t) dt)
248: // f(x,q2) = f_0(x,q2) * exp(-int_0^x (alpha_s(t)/t) dt)
249: // f(x,q2) = f_0(x,q2) * exp(-int_0^x (alpha_s(t)/t) dt)
```

```
./src/MSudakov.cc 4/8
250: double MSudakov::fg_xQ2(double x, double q2, double M) const {
251:     // Calculate Shuvev transformation
252:     std::pair<double, double> out1 = Shuvev_B(q2, x);
253:     std::pair<double, double> out2 = Interpolated2(q2, x);
254:     const double Hg = out1.first;
255:     const double dHg = out1.second;
256:
257:     // Calculate Sudakov veto
258:     std::pair<double, double> out2 = Sudakov_T(q2, M);
259:     std::pair<double, double> out2 = veto.Interpolated2(q2, M);
260:     const double Tg = out2.first;
261:     const double dTg = out2.second;
262:
263:     // Chain rule s' d/dln(q^2) [ ... ]
264:     double total = 0.0;
265:
266:     if (Tg > 1e-15) {
267:         total = dHg * math::imgstr(Tg) + Hg * dTg / (2.0 * math::imgstr(Tg));
268:     } else {
269:         // printf("MSudakov::fg_xQ2: Sudakov factor Tg = %0.2E [x = %0.1E, q2 = %0.1E] \n", Tg, x, q2);
270:         // GeV^2, M = %0.1F GeV \n", Tg, x, q2, M);
271:     }
272:
273:     // And because our functions return partial derivatives w.r.t. q^2 (not ln q^2)
274:     // We transform:
275:     // dln(q^2)/dq^2 = 1/q^2 <=> dln(q^2) = dq^2/q^2 <=>
276:     // d[...]/dln(q^2) = d[...]/dq^2 * q^2, applied below
277:
278:     // Total = q^2;
279:
280:     // Truncate negative values to give proper zero
281:     total = (total < 0.0) ? 0.0 : total;
282:
283:     return total;
284: }
285:
286: // Calculate Shuvev integral transform
287: // Identity:
288: // H_g(x, |x| -> 0) = xg(x)
289: // H_g = numerical integral transformed from standard pdf
290: // dHg = dHg/dq^2 (differentiated numerically)
291: // [REFERENCE: Harland-Lang, arxiv.org/abs/1306.6661]
292:
293: std::pair<double, double> MSudakov::Shuvev_B(double q2, double x) {
294:     const double y_MIN = x / 4.0;
295:     const double y_MAX = 1.0;
296:     const double y_STEP = (y_MAX - y_MIN) / Numerics.ShuvevIntegral;
297:
298:     double Hg = 0.0;
299:     double dHg = 0.0;
300:
301:     // Check that we are within valid domain (take into account floating points)
302:     const double EPS = 1e-5;
303:     if (x < Numerics.x_MIN * (1 - EPS) && x <= Numerics.x_MAX * (1 + EPS) &&
304:         q2 >= Numerics.q2_MIN * (1 - EPS) && q2 <= Numerics.q2_MAX * (1 + EPS)) {
305:         // We transform:
306:         std::vector<double> fA(Numerics.ShuvevIntegral + 1, 0.0);
307:         std::vector<double> fB(Numerics.ShuvevIntegral + 1, 0.0);
308:
309:         for (const auto i: indices(fA)) {
310:             const double y = y_MIN + i * y_STEP;
311:             const double argument = x / (4.0 * y);
312:
313:             // H_g(x/2, x/2, Q^2) = 4x/gp1 * ln(x/4) * y
314:             // y^(1/2) (1-y)^(1/2) / (2y(x/4y), Q^2)
315:
316:             // We take into account that LHAPDF provides xg(), not g(), gives:
317:             const double norm = math::imgstr(math::pow3(y) * (1 - y));
318:             fA[i] = factor * xg_Q2(argument, q2);
319:             fB[i] = factor * diff_xg_Q2_ext_Q2(argument, q2);
320:
321:             const double norm = 16.0 * math::PI;
322:             Hg = norm * math::CIntegral(fA, y_STEP);
323:             dHg = norm * math::CIntegral(fB, y_STEP);
324:         }
325:     } else {
326:         // Fatal error
327:     }
328: }
329:
330: // With sum over active quark flavors with mass threshold qT
331: // sum_q TR = ln(q^2 + (1-s)^2) * delta, z = 0 ... (1-delta)
332:
333: double MSudakov::fAG(double delta, double q2) const {
334:     const double TR = 0.0; // Structure constant
335:     return TR * (-2.0 * math::pow3(delta) / 3.0 + math::pow2(delta) - delta + 2.0 / 3.0 *
336:         NumFlavor(q2));
337: }
338:
339: // Return the number of quark flavors at scale q^2
340: double MSudakov::NumFlavor(double q2) const {
341:     const double n_charm = 1.275; // GeV, PDG-2018 (default definition)
342:     const double n_bottom = 4.18; // GeV
343:
344:     if (q2 < math::pow2(m_charm)) {
345:         return 3.0;
346:     } else if (q2 < math::pow2(m_bottom)) {
347:         return 4.0;
348:     } else {
349:         return 5.0;
350:     }
351: }
352:
353: // Constructs interpolation array values
354: void MSudakov::CalculateArray(Array2D arr,
355:     int& nLines, int& nCols) const {
356:     for (const auto i: indices(arr.F)) {
357:         const double a = arr.MIN[0] + i * arr.STEP[0];
358:         // Transform input to linear if log stepping, for the function
359:         const double var1 = (arr.log[0]) ? std::exp(a) : a;
360:         const double var2 = (arr.log[1]) ? std::exp(a) : a;
361:         // Do not write if file exists already
362:         // Write the array to a file
363:         bool fWriteArray(const std::string& filename, bool overwrite) const {
364:             if (std::ifstream(filename).is_open()) {
365:                 // Do not write if file exists already
366:                 return true;
367:             }
368:             // Write to file
369:             std::ofstream file;
370:             file.open(filename);
371:             if (!file.is_open()) {
372:                 std::string str = "Array2D::WriteArray: Fatal I/O-error with: " + filename;
373:                 throw std::invalid_argument(str);
374:             }
375:             std::cout << "Array2D::WriteArray: "
376:                 + "unsaved int line_number = 0"
377:                 + "\n";
378:             for (const auto i: indices(F)) {
379:                 for (const auto j: indices(F)) {
380:                     file << F[i][j][0] << ", " << F[i][j][1] << ", " << F[i][j][2] << ", "
381:                         << F[i][j][3] << std::endl;
382:                 }
383:             }
384:             ++line_number;
385:         }
386:     }
387: }
388:
389: // See standard literature, e.g.
390: // [REFERENCE: B.R. Webber, CERN lectures 09,
391: // www.hep.phy.cam.ac.uk/theory/webber/C09lec3.pdf]
392: // [REFERENCE: R.K. Ellis, W.J. Stirling, B.R. Webber, QCD and Collider Physics]
393: // [REFERENCE: Higgs/pdg-1.1-gov/2018/download/ds118.pdf]
394:
395: // [int z - gyz(z)]
396: // [int z + 1/2 * CA * int [z/(1-z)] + (1-z)/z + 2*(1-z)]
397: // [1/6*(11*CA - 4*Nf*Tr)*deltafunc(1-z) dz, z = 0 ... (1-delta) =
398: // (deltaw)
399:
400: // First part: ...
401: // Second part: int 1/6*(11*CA - 4*Nf*Tr)*DiracDelta[1-x] dx, x = 0 ... 1 - delta,
402: // this vanishes with positive delta
403: // where + is the "plus-description"
404:
405: double MSudakov::fAG(double delta) const {
406:     const double CA = 3.0; // Structure constant
407:     return 2.0 * CA *
408:         (std::log(1.0 / delta) -
409:             math::pow2(1.0 - delta) * (3.0 * math::pow2(delta) - 2.0 * delta + 11.0) / 12.0);
410: }
411:
412: // Altarelli-Parisi splitting function definite integral over z,
413: // [int sum_q P_qg(z) dz
```



```

./src/MSudakov.cc          7/8
499:         }
500:     }
501: } catch (...) {
502:     throw std::invalid_argument("TArray2D::WriteArray: Error in file " + filename + " at line " +
503:         std::ito_string(line_number));
504: }
505:
506: file.close();
507: return true;
508: }
509:
510: // Read the array from a file
511: bool TArray2D::ReadArray(const std::string& filename) {
512:     std::ifstream file;
513:     file.open(filename);
514:     if (!file.is_open()) {
515:         std::string str = "TArray2D::ReadArray: Fatal IO-error with: " + filename;
516:         return false;
517:     }
518:
519:     std::string line;
520:     unsigned int fills = 0;
521:     unsigned int line_number = 0;
522:     std::cout << "TArray2D::ReadArray: ";
523:
524:     try {
525:         for (const auto& i : indices(F)) {
526:             for (const auto& j : indices(F[i]) ) {
527:                 // Read every line from the stream
528:                 getline(file, line);
529:                 std::istringstream stream(line);
530:                 std::vector<std::string> columns;
531:                 std::string element;
532:
533:                 // Get every line element (if them) separated by separator
534:                 int k = 0;
535:                 while (getline(stream, element, ',') ) {
536:                     F[i][j][k] = std::stod(element); // string to double
537:                     ++k;
538:                 }
539:                 ++fills;
540:                 ++line_number;
541:             }
542:         }
543:     }
544: } catch (...) {
545:     throw std::invalid_argument("TArray2D::ReadArray: Error in file " + filename + " at line " +
546:         std::ito_string(line_number));
547: }
548:
549: file.close();
550:
551: if (fills != 4 * (N[0] + 1) * (N[1] + 1)) {
552:     std::string str = "Corrupted file: " + filename;
553:     std::cout << str << std::endl;
554:     return false;
555: }
556:
557: std::cout << rang::fg::green << "[DONE]" << rang::fg::reset << std::endl;
558: return true;
559: }
560:
561: // Standard 2D-bilinear interpolation f(a,b) = z, and derivative dz
562:
563:
564: // [REFERENCE: en.wikipedia.org/wiki/Bilinear_interpolation]
565: std::pair<double, double> TArray2D::Interpolate2D(double a, double b) const {
566:     const double EPS = 1e-5;
567:
568:     if (a < MIN[0]) { a = MIN[0]; } // Truncate before (possible) logarithm
569:     if (b < MIN[1]) { b = MIN[1]; } // Truncate before (possible) logarithm
570:
571:     // Logarithmic stepping or not
572:     if (islog[0]) { a = std::log(a); }
573:     if (islog[1]) { b = std::log(b); }
574:
575:     if (a > MAX[0] * (1 + EPS) || b > MAX[1] * (1 + EPS)) {
576:         printf(
577:             "Interpolate2D(%s,%s) Input out of grid domain: ",
578:             "a = %0.3F (%0.3E, %0.3E), b = %0.3F (%0.3E, %0.3E) \n",
579:             name[0].c_str(), name[1].c_str(), name[0].c_str(), a, MIN[0], MAX[0], name[1].c_str(), b,
580:             MIN[1], MAX[1]);
581:     }

```

```

./src/MSudakov.cc          8/8
582:
583: // Get indices
584: int i = std::floor((a - MIN[0]) / STEP[0]);
585: int j = std::floor((b - MIN[1]) / STEP[1]);
586:
587: // Lower boundary protection
588: if (i < 0) { i = 0; } // Int needed for this (not unsigned int)
589: if (j < 0) { j = 0; }
590:
591: // Upper boundary protection
592: if (i > (int)N[0]) { i = N[0] - 1; } // We got N+1 elements in F
593: if (j > (int)N[1]) { j = N[1] - 1; }
594:
595: // Aux variables for readability
596: const unsigned int X = 0;
597: const unsigned int Y = 1;
598:
599: const double x1 = F[i][j][X];
600: const double x2 = F[i + 1][j][X];
601: const double y1 = F[i][j][Y];
602: const double y2 = F[i][j + 1][Y];
603:
604: const double xstep = STEP[0];
605: const double ystep = STEP[1];
606: const double xval = a;
607: const double yval = b;
608:
609: double values[2] = {0.0};
610:
611: // 2 == Z, 3 == dz
612: for (std::size_t C = 2; C <= 3; ++C) {
613:     const double Q1 = F[i][j][C];
614:     const double Q2 = F[i + 1][j][C];
615:     const double Q3 = F[i + 1][j + 1][C];
616:     const double Q4 = F[i + 1][j + 1][C];
617:
618:     // Interpolated value
619:     values[C - 2] =
620:         ((y2 - yval) / ystep) * ((x2 - xval) / xstep * Q11 + (xval - x1) / xstep * Q21) +
621:         ((yval - y1) / ystep) * ((x2 - xval) / xstep * Q12 + (xval - x1) / xstep * Q22) +
622:         ((x2 - x1) / xstep) * ((y2 - yval) / ystep * Q31 + (yval - y1) / ystep * Q41) +
623:         ((y2 - y1) / ystep) * ((x2 - xval) / xstep * Q32 + (xval - x1) / xstep * Q42);
624:     return {values[0], values[1]};
625: }
626:
627: } // namespace gra
628:

```

```

./src/MForm.cc             1/11
1: // Form factors, structure functions, Regge trajectories etc. parametrizations
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <cstdlib>
9: #include <iostream>
10: #include <map>
11: #include <vector>
12:
13: // Own
14: #include "GranitLL/MForm.h"
15: #include "GranitLL/MGRanitLL.h"
16: #include "GranitLL/MELMonematics.h"
17: #include "GranitLL/MMath.h"
18: #include "GranitLL/MFProcess.h"
19: #include "GranitLL/MQSB.h"
20: #include "GranitLL/MSpin.h"
21:
22: // Libraries
23: #include "json.hpp"
24: #include "rang.hpp"
25:
26:
27: using gra::math::PI;
28: using gra::math::abs2;
29: using gra::math::imag2;
30: using gra::math::pow2;
31: using gra::math::pow3;
32: using gra::math::t2;
33:
34: using gra::PDD::mp;
35: using gra::PDD::mp2;
36:
37: namespace gra {
38:
39: // Structure functions and nuclear form factors
40: namespace PARAM_STRUCTURE {
41:
42:     std::string F2 = "DMF";
43:     std::string EM = "DIPOLE";
44:     std::string QED_alpha = "ZERO";
45:
46:     bool initialized = false;
47:
48:     // Read parameters from file
49:     void ReadParameters(const std::string& modelfile) {
50:         try {
51:             const std::string data = gra::aux::GetInputData(modelfile);
52:             nlohmann::json j = nlohmann::json::parse(data);
53:
54:             const std::string XID = "PARAM_STRUCTURE";
55:
56:             PARAM_STRUCTURE::F2 = j.at(XID).at("F2");
57:             PARAM_STRUCTURE::EM = j.at(XID).at("EM");
58:             PARAM_STRUCTURE::QED_alpha = j.at(XID).at("QED_alpha");
59:
60:             initialized = true;
61:         }
62:     } catch (...) {
63:         std::string str = "PARAM_STRUCTURE::ReadParameters: Error parsing " + modelfile +
64:             " (Check for extra/missing commas)";
65:         throw std::invalid_argument(str);
66:     }
67: }
68:
69: // namespace PARAM_STRUCTURE
70:
71: // Model parameters
72: namespace PARAM_SOFT {
73:     // Pomeron trajectory
74:     double DELTA_P = 0.0;
75:     double ALPHA_P = 0.0;
76:
77: // Couplings
78:     double gR_P = 0.0;
79:     double gP = 0.0;
80:     double gamma = 0.0;
81:
82: // Proton form factor
83:     double fcl = 0.0;

```

```

./src/MForm.cc             2/11
84:     double fc2 = 0.0;
85:
86: // Flux loop
87:     double fc3 = 0.0;
88:
89: // On/Off
90:     bool ODDERON_ON = false;
91:
92:     bool initialized = false;
93:
94:     std::string GetHashString() {
95:         std::string str = std::ito_string(PARAM_SOFT::DELTA_P) + std::ito_string(PARAM_SOFT::ALPHA_P) +
96:             std::ito_string(PARAM_SOFT::gR_P) + std::ito_string(PARAM_SOFT::gR_D) +
97:             std::ito_string(PARAM_SOFT::fc1) + std::ito_string(PARAM_SOFT::fc2) +
98:             std::ito_string(PARAM_SOFT::fc3) + std::ito_string(PARAM_SOFT::ODDERON_ON);
99:         return str;
100:     }
101:
102:     void PrintParam() {
103:         printf("PARAM_SOFT: Soft model parameters: \n\n");
104:         printf(" DELTA_P = %0.5F \n", DELTA_P);
105:         printf(" ALPHA_P = %0.5F [GeV^-2] \n", ALPHA_P);
106:         printf(" gR_P = %0.5F [GeV^-1] \n", gR_P);
107:         printf(" gR_D = %0.5F [GeV^-1] \n", gR_D);
108:         printf(" gP = %0.5F \n", gP / gR_P); // Convention
109:         printf(" gamma = %0.5F \n", gamma);
110:         printf(" fc1 = %0.5F [GeV^2] \n", fc1);
111:         printf(" fc2 = %0.5F [GeV^2] \n", fc2);
112:         printf(" fc3 = %0.5F [GeV^2] \n", fc3);
113:         printf(" ODDERON_ON = %i", ODDERON_ON);
114:         std::cout << "ODDERON_ON ? Error: " "false" << std::endl;
115:         std::cout << std::endl << std::endl;
116:     }
117:
118: // Read parameters from file
119: void ReadParameters(const std::string& modelfile) {
120:     try {
121:         const std::string data = gra::aux::GetInputData(modelfile);
122:         nlohmann::json j = nlohmann::json::parse(data);
123:
124:         const std::string XID = "PARAM_SOFT";
125:
126:         PARAM_SOFT::DELTA_P = j.at(XID).at("DELTA_P");
127:         PARAM_SOFT::ALPHA_P = j.at(XID).at("ALPHA_P");
128:         PARAM_SOFT::gR_P = j.at(XID).at("gR_P");
129:         PARAM_SOFT::gR_D = j.at(XID).at("gR_D");
130:
131:         double tripleSP = j.at(XID).at("g3P");
132:         PARAM_SOFT::g3P = tripleSP * PARAM_SOFT::gR_P; // Convention
133:         PARAM_SOFT::igamma = j.at(XID).at("gamma");
134:
135:         PARAM_SOFT::fc1 = j.at(XID).at("fc1");
136:         PARAM_SOFT::fc2 = j.at(XID).at("fc2");
137:         PARAM_SOFT::fc3 = j.at(XID).at("fc3");
138:
139:         PARAM_SOFT::ODDERON_ON = j.at(XID).at("ODDERON_ON");
140:
141:         initialized = true;
142:         PARAM_SOFT::PrintParam();
143:     }
144: } catch (...) {
145:     std::string str = "PARAM_SOFT::ReadParameters: Error parsing " + modelfile +
146:         " (Check for extra/missing commas)";
147:     throw std::invalid_argument(str);
148: }
149:
150:
151: // namespace PARAM_SOFT
152:
153: // Flat amplitude parameters
154: namespace PARAM_FLAT {
155:     int active = 0;
156:     double b = 0.0;
157:
158:     bool initialized = false;
159:
160:     // Read parameters from file
161:     void ReadParameters(const std::string& modelfile) {
162:         try {
163:             const std::string data = gra::aux::GetInputData(modelfile);
164:             nlohmann::json j = nlohmann::json::parse(data);
165:
166:             const std::string XID = "PARAM_FLAT";

```

```

./src/MForm.cc 3/11
167:
168:   PARAM_FLAT:b = j.at(XID).at("b");
169:
170:   initialized = true;
171: } catch (...) {
172:   std::string str = "PARAM_FLAT:ReadParameters: Error parsing " + modelfile +
173:                   " (Check for extra/missing commas)";
174:   throw std::invalid_argument(str);
175: }
176: }
177: } // namespace PARAM_FLAT
178:
179: // Forward proton excitation
180: namespace PARAM_NSTAR {
181: std::string fragment = "none";
182: std::vector<double> rc = {0.0, 0.0, 0.0};
183:
184: bool initialized = false;
185:
186: // Read parameters from file
187: void ReadParameters(const std::string &modelfile) {
188:   try {
189:     const std::string data = gra::aux::GetInputData(modelfile);
190:     nlohmann::json j = nlohmann::json::parse(data);
191:
192:     const std::string XID = "PARAM_NSTAR";
193:
194:     PARAM_NSTAR:fragment = j.at(XID).at("fragment");
195:     std::vector<double> rc = j.at(XID).at("rc");
196:     PARAM_NSTAR:rc = rc;
197:
198:     // Make sure they sum to one
199:     const double sum_rc = std::accumulate(rc.begin(), rc.end(), 0);
200:     for (std::size_t i = 0; i < PARAM_NSTAR:rc.size(); ++i) {PARAM_NSTAR:rc[i] /= sum_rc;}
201:
202:     initialized = true;
203:   } catch (...) {
204:     std::string str = "PARAM_NSTAR:ReadParameters: Error parsing " + modelfile +
205:                     " (Check for extra/missing commas)";
206:     throw std::invalid_argument(str);
207:   }
208: } // namespace PARAM_NSTAR
209:
210: namespace form {
211:
212: // Read resonance parameters
213: gra::PARAM_RES ReadResonance(const std::string &resparam_str, MRandom &rng) {
214: // =====
215: // Find global resonance parameters
216:
217: bool SPINDEC = false;
218: bool SPINDECC = false;
219: std::string FRAME = "null";
220: int JMAX = 0;
221:
222: // Read and parse
223: const std::string fullpath =
224:   gra::aux::GetBasePath(2) + "/modeldata/" + gra::MODELPARAM + "/GENERAL.json";
225:
226: const std::string data = gra::aux::GetInputData(fullpath);
227: nlohmann::json j;
228:
229: try {
230:   j = nlohmann::json::parse(data);
231: } catch (...) {
232:   std::string str =
233:     "form:ReadResonance: Error parsing " + fullpath + " (Check for extra/missing commas)";
234:   throw std::invalid_argument(str);
235: }
236:
237: SPINDEC = j.at("PARAM_SPIN")>at("SPINDEC");
238: SPINDECC = j.at("PARAM_SPIN")>at("SPINDECC");
239: FRAME = j.at("PARAM_SPIN")>at("FRAME");
240: JMAX = j.at("PARAM_SPIN")>at("JMAX");
241:
242: // =====
243:
244: std::cout << "gra:form:ReadResonance: Reading " + resparam_str + " ";
245:
246:
247:
248:
249: // Create a JSON object from file

```

```

./src/MForm.cc 4/11
250: std::string inputfile =
251:   gra::aux::GetBasePath(2) + "/modeldata/" + gra::MODELPARAM + "/" + resparam_str;
252:
253: // Read and parse
254: std::string data;
255: nlohmann::json j;
256:
257: try {
258:   data = gra::aux::GetInputData(inputfile);
259:   j = nlohmann::json::parse(data);
260: } catch (...) {
261:   throw std::invalid_argument("form:ReadResonance: Error parsing " + resparam_str + "");
262: }
263:
264: // Resonance parameters
265: gra::PARAM_RES res;
266:
267: try {
268: // =====
269: // Collect global variables
270: res.SPINDEC = SPINDEC;
271: res.SPINDECC = SPINDECC;
272: res.FRAME = FRAME;
273: res.JMAX = JMAX;
274: // =====
275:
276: // Complex coupling
277: double g_A = j.at("PARAM_RES")>at("g_A");
278: double g_phi = j.at("PARAM_RES")>at("g_phi");
279: res.g = g_A * std::exp(std::complex<double>(0, 1) * g_phi);
280:
281: // Form factor parameter
282: res.g_FF = j.at("PARAM_RES")>at("g_FF");
283:
284: // PDC code
285: res.p.pdc = j.at("PARAM_RES")>at("PDC");
286:
287: // Mass
288: res.p.mass = j.at("PARAM_RES")>at("M");
289: if (res.p.mass < 0) {
290:   std::string str = "Mux:ReadResonance: <" + resparam_str + "> Invalid M < 0 ";
291:   throw std::invalid_argument(str);
292: }
293:
294: // Width
295: res.p.width = j.at("PARAM_RES")>at("W");
296: if (res.p.width < 0) {
297:   std::string str = "Mux:ReadResonance: <" + resparam_str + "> Invalid W < 0 ";
298:   throw std::invalid_argument(str);
299: }
300:
301: // Spin
302: const double J = j.at("PARAM_RES")>at("J");
303: res.p.spin2 = J * 2;
304: if (res.p.spin2 < 0) {
305:   std::string str = "Mux:ReadResonance: <" + resparam_str + "> Invalid J < 0 ";
306:   throw std::invalid_argument(str);
307: }
308:
309: // Parity
310: res.p.p = j.at("PARAM_RES")>at("P");
311: if (1/(res.p.p == -1 || res.p.p == 1)) {
312:   std::string str = "Mux:ReadResonance: <" + resparam_str + "> Invalid P (not -1 or 1) ";
313:   throw std::invalid_argument(str);
314: }
315:
316: // =====
317: // Tensor Pomeron couplings vector
318: std::vector<double> g_Tensor = j.at("PARAM_RES")>at("g_Tensor");
319: res.g_Tensor = g_Tensor;
320:
321: if (J == 0 && res.g_Tensor.size() != 2) {
322:   throw std::invalid_argument("Mux:ReadResonance: <" + resparam_str +
323:     " Tensor Pomeron coupling array should be of size 2 for J = 0");
324: }
325: if (J == 1 && res.g_Tensor.size() != 2) {
326:   throw std::invalid_argument("Mux:ReadResonance: <" + resparam_str +
327:     " Tensor Pomeron coupling array should be of size 2 for J = 1");
328: }
329: if (J == 2 && res.g_Tensor.size() != 7) {
330:   throw std::invalid_argument("Mux:ReadResonance: <" + resparam_str +
331:     " Tensor Pomeron coupling array should be of size 7 for J = 2");
332: }

```

```

./src/MForm.cc 5/11
333: // =====
334: // Validity of these is taken care of in the functions
335: res.BW = j.at("PARAM_RES")>at("BW");
336:
337: const bool P_conservation = true;
338:
339: const int n = res.p.spin2 + 1; // n = 2J + 1
340: MMatrix<std::complex<double>> rho(s, pi);
341:
342: // If we have spin
343: if (res.p.spin2 != 0) {
344: // Draw Random Density matrices (until the number set by user)
345: if (j.at("PARAM_RES")>at("random_rho") > 0) {
346:   for (std::size_t k = 0; k < j.at("PARAM_RES")>at("random_rho"); ++k) {
347:     rho = gra::spin::RandomSho(res.p.spin2, 2, 0, P_conservation, rng);
348:   }
349: }
350:
351: // Construct spin density matrix from the input
352: } else {
353:   for (std::size_t a = 0; a < rho.size_row(); ++a) {
354:     for (std::size_t b = 0; b < rho.size_col(); ++b) {
355:       const double Re = j.at("PARAM_RES")>at("rho_real").at(a).at(b);
356:       const double Im = j.at("PARAM_RES")>at("rho_imag").at(a).at(b);
357:       rho(a)[b] = Re + xi * Im;
358:     }
359:   }
360: }
361: // Check positivity conditions
362: if (gra::spin::Positivity(rho, res.p.spin2 / 2.0) == false) {
363:   std::string str = "gra:form:ReadResonance: <" + resparam_str +
364:     " Input density matrix not positive definite";
365:   throw std::invalid_argument(str);
366: }
367: res.rho = rho;
368:
369: std::cout << rang::fg::green << "[DONE]" << rang::fg::reset << std::endl;
370: } catch (nlohmann::json::exception &e) {
371:   throw std::invalid_argument("form:ReadResonance: Missing parameter in " + resparam_str +
372:     " " + e.what());
373: }
374:
375: return res;
376: }
377:
378: // Regge signature factor, alpha_t is alpha(t), and signature sigma = +/-
379: // Remember that denominator hits pole at every integer Pi + alpha(t)
380: // Regge signature factor for linear trajectories at t -> 0
381: // Mandelstam t = alpha_0(t)
382: // ap = alpha'
383: // sigma = +/-
384: // Regge signature factor for linear trajectories at t -> 0
385: // Mandelstam t = alpha_0(t)
386: // ap = alpha'
387: // sigma = +/-
388: // Regge signature factor for linear trajectories at t -> 0
389: // Mandelstam t = alpha_0(t)
390: // ap = alpha'
391: // sigma = +/-
392: // Regge signature factor for linear trajectories at t -> 0
393: // Mandelstam t = alpha_0(t)
394: // ap = alpha'
395: // sigma = +/-
396: // Regge signature factor for linear trajectories at t -> 0
397: // Mandelstam t = alpha_0(t)
398: // ap = alpha'
399: // sigma = +/-
400: // Regge signature factor for linear trajectories at t -> 0
401: // Mandelstam t = alpha_0(t)
402: // ap = alpha'
403: // sigma = +/-
404: // Regge signature factor for linear trajectories at t -> 0
405: // Mandelstam t = alpha_0(t)
406: // ap = alpha'
407: // sigma = +/-
408: // Regge signature factor for linear trajectories at t -> 0
409: // Mandelstam t = alpha_0(t)
410: // ap = alpha'
411: // sigma = +/-
412: // Regge signature factor for linear trajectories at t -> 0
413: // Mandelstam t = alpha_0(t)
414: // ap = alpha'
415: // sigma = +/-

```

```

./src/MForm.cc 6/11
416: double SHPI(double tau, double t) {
417:   const double m = 1.0; // fixed scale (GeV)
418:   const double sqrtau = sqrt(t + tau);
419:
420: // Non-Pomeron form factor parametrization
421: const double F_2 = 1.0 / (1.0 - t / PARAM_SOFT:fcf);
422:
423: return (4.0 / tau) * pow(F_2, 3) *
424:   (2.0 * tau - std::pow(1.0 + tau, 3.0 / 2.0) * std::log(sqrtau + 1.0) / (sqrtau - 1.0)) +
425:   std::log(m * m) / (mpi * mpi);
426: }
427: // =====
428: // Elastic proton form factor parametrization
429: // <apply at amplitude level>
430: // [REFERENCE: Khose, Martin, Ryskin, arxiv.org/abs/hep-ph/0007359]
431: const double F_1(double t) {
432:   return (1.0 / (2 * t / PARAM_SOFT:fcf1)) * (1.0 / (1 - t / PARAM_SOFT:fcf2));
433: }
434: // =====
435: // Proton inelastic form factor / structure function
436: // parametrization for Pomeron processes (THIS FUNCTION IS ANSATZ - IMPROVE!)
437: // Motivated by arxiv.org/abs/hep-ph/9305319
438: // <apply at amplitude level>
439: double SHFIMEL(double t, double M2) {
440:   const double DELTA_P = 0.0808;
441:   const double A = 0.5616; // GeV^2
442:   double f = std::pow(std::abs(t) / M2 * (std::abs(t) + A), 0.5 * (1 + DELTA_P));
443:   return f;
444: }
445: // =====
446: // Proton inelastic structure function F2(x,Q^2) parametrization
447: // The basic idea is that at low-Q^2, a fully non-perturbative description
448: // (parametrization) is needed.
449: // At high Q^2, DGLAP evolution could be done in log(Q^2) starting from the
450: // input description.
451: // Now, some (very) classic ones have been implemented. Add new one here!
452: // [REFERENCE: Donnachie, Landshoff, arxiv.org/abs/hep-ph/9305319]
453: // [REFERENCE: Capella, Kaidalov, Merino, Tran Thanh Van, arxiv.org/abs/hep-ph/9405388v1]
454: // [REFERENCE: F2HQ2(double xBj, double Q2) {
455: //   if (PARAM_STRUCTURE:F2 == "DL") {
456: //     constexpr double A = 0.324;
457: //     constexpr double B = 0.098;
458: //     constexpr double DELTA_P = 0.0808;
459: //     constexpr double DELTA_R = 0.5475;
460: //     constexpr double b = 0.01133;
461: //     constexpr double F2 = A * std::ipow(xBj, -DELTA_P) * std::ipow(Q2 / (Q2 + A), 1 + DELTA_R) +
462: //       B * std::ipow(xBj, 1 - DELTA_R) * std::ipow(Q2 / (Q2 + B), DELTA_R);
463: //     return F2;
464: //   }
465: //   else if (PARAM_STRUCTURE:F2 == "CBMT") {
466: //     constexpr double A = 0.1502;
467: //     constexpr double B_u = 1.2064;
468: //     constexpr double B_d = 0.1798;
469: //     constexpr double alpha = 0.4302;
470: //     constexpr double DELTA_0 = 0.0800;
471: //     constexpr double B = 0.2631;
472: //     constexpr double b = 0.6452;
473: //     constexpr double c = 3.1468;
474: //     constexpr double d = 1.1170;
475: //     constexpr double n_Q2 = (3.0 / 2.0) * (1 + Q2 / (Q2 + c));
476: //     constexpr double DELTA_Q2 = DELTA_0 * (1 + (2 * Q2) / (Q2 + d));
477: //     constexpr double C1 = std::ipow(Q2 / (Q2 + A), 1.0 + DELTA_Q2);
478: //     constexpr double C2 = std::ipow(Q2 / (Q2 + B), alpha_R);
479: //     constexpr double F2 = A * std::ipow(xBj, -DELTA_Q2) * std::ipow(1 - xBj, n_Q2 + 4.0) * C1 +

```

```

./src/MForm.cc 7/11
499:      std::pow(xb_j, 1.0 - alpha_R) *
500:      (B_u * std::pow(1 - xb_j, n_Q2) + B_d * std::pow(1 - xb_j, n_Q2 + 1.0)) *
501:      c2;
502:
503:   return F2;
504: } else {
505:   throw std::invalid_argument("gra:form:F2x2: Unknown PARAM_STRUCTURE:F2 = " +
506:     PARAM_STRUCTURE:F2);
507: }
508: }
509:
510: // "Purely magnetic structure function"
511: //
512: // Callan-Gross relation for spin-1/2: F_2(x) = 2x_1(x) under Bjorken scaling
513: // For spin-0, F_1(x) = 0
514: //
515: // Longitudinal structure function (eq. G. QD2)
516: // F_L(xb_j, Q2) = (1 + 4*pow2(xb_jmp)/Q2) * F_2(xb_j, Q2) - 2xb_j * F_1(xb_j, Q2)
517: //
518: double F1xQ(double xb_j, double Q2) { return F2xQ(xb_j, Q2) / (2.0 * xb_j); }
519:
520: // QED: Proton magnetic moment in magneton units (mu_p / mu_N)
521: double mu_ratio() { return 2.792847337; }
522:
523: // =====
524: // Photon flux densities and form factors, input Q^2 as positive
525: //
526: // kT unintegrated coherent EPA photon flux as in:
527: //
528: // [REFERENCE: Luszczak, Schafer, Szcurek, arxiv.org/abs/1802.03244]
529: //
530: // Form factors:
531: // [REFERENCE: Funjabi et al., arxiv.org/abs/1503.01452v4]
532: //
533: // Proton electromagnetic form factors: Basic notions, present
534: // achievements and future perspectives, Physics Reports, 2015
535: // <www.sciencedirect.com/science/article/pii/S0370157314003184>
536: //
537: // [REFERENCE: Budnev, Ginzburg, Meledin, Serbo, EPA paper, Physics Reports, 1976]
538: // <www.sciencedirect.com/science/article/pii/S0370157379009395>
539: //
540: //
541: // Proton EM form factor F1 (Dirac)
542: double F1(double Q2) {
543:   const double tau = Q2 / pow2(2 * mp);
544:   return 1.0 / (tau + 1) * G_E(Q2) + tau / (tau + 1) * G_M(Q2);
545: }
546:
547: //
548: // Proton EM form factor F2 (Pauli)
549: double F2(double Q2) {
550:   const double tau = Q2 / pow2(2 * mp);
551:   return 1.0 / (tau + 1) * G_E(Q2) + 1.0 / (tau + 1) * G_M(Q2);
552: }
553:
554: // Rosenbluth separation:
555: // low-Q^2 dominated by G_E, high-Q^2 dominated by G_M
556: //
557: // Sachs Form Factor* goes as follows:
558: // G_E(0) = 1 for proton, 0 for neutron
559: // G_M(0) = mu_p for proton, mu_n for neutron
560: //
561: // <http://www.scholarpedia.org/article/Nucleon_Form_Factors>
562: //
563: double G_E(double Q2) {
564:   Q2 = std::abs(Q2); // For safety
565:   if (PARAM_STRUCTURE:EM == "DIPOLE") {
566:     return G_E_DIPOLE(Q2);
567:   } else if (PARAM_STRUCTURE:EM == "KELLY") {
568:     return G_E_KELLY(Q2);
569:   } else {
570:     throw std::invalid_argument("gra:form:G_E: Unknown proton EM-form factor chosen = " +
571:       PARAM_STRUCTURE:EM);
572:   }
573: }
574:
575: double G_M(double Q2) {
576:   Q2 = std::abs(Q2); // For safety
577: }
578:
579: double G_E(double Q2) {
580:   Q2 = std::abs(Q2); // For safety
581: }

```

```

./src/MForm.cc 8/11
582: if (PARAM_STRUCTURE:EM == "DIPOLE") {
583:   return G_E_DIPOLE(Q2);
584: } else if (PARAM_STRUCTURE:EM == "KELLY") {
585:   return G_M_KELLY(Q2);
586: } else {
587:   throw std::invalid_argument("gra:form:G_M: Unknown proton EM-form factor chosen = " +
588:     PARAM_STRUCTURE:EM);
589: }
590: }
591:
592: // The simplest possible: Dipole parametrization of nucleon EM-form factors
593: //
594: //
595: //
596: double G_E_DIPOLE(double Q2) { // Scaling assumption
597:   return G_M(Q2) / mu_ratio();
598: }
599:
600: double G_M_DIPOLE(double Q2) {
601:   const double lambda2 = 0.71; // Dipole parameter GeV^2
602:   return mu_ratio() / pow2(1.0 + Q2 / lambda2);
603: }
604: // Simple parametrization of nucleon EM-form factors
605: //
606: // [REFERENCE: Kelly, journals.aps.org/pr/pdf/10.1103/PhysRevC.70.068202]
607: double G_E_KELLY(double Q2) {
608:   static const std::vector<double> b = {1, -0.24};
609:   static const std::vector<double> b = {10.98, 12.82, 21.97};
610: }
611: const double tau = Q2 / pow2(2 * mp);
612: // Numerator
613: double num = 0.0; // 0
614: num += a[0]; // a_0 tau^0
615: num += a[1] * tau; // a_1 tau^1
616: // Denominator
617: double den = 1.0; // 1.0
618: den += b[0] * tau; // b_1 tau^1
619: den += b[1] * pow2(tau); // b_2 tau^2
620: den += b[2] * pow3(tau); // b_3 tau^3
621: return num / den;
622: }
623:
624: double G_M_KELLY(double Q2) {
625:   static const std::vector<double> a = {1, 0.12};
626:   static const std::vector<double> b = {10.97, 16.86, 6.55};
627:   const double tau = Q2 / pow2(2 * mp);
628:   // Numerator
629:   double num = 0.0; // 0
630:   num += a[0]; // a_0 tau^0
631:   num += a[1] * tau; // a_1 tau^1
632:   // Denominator
633:   double den = 1.0; // 1.0
634:   den += b[0] * tau; // b_1 tau^1
635:   den += b[1] * pow2(tau); // b_2 tau^2
636:   den += b[2] * pow3(tau); // b_3 tau^3
637:   return mu_ratio() * num / den;
638: }
639: // Coherent photon flux from proton
640: // xi - longitudinal momentum loss [0,1]
641: // p - Mandelstam s
642: // pt - proton transverse momentum
643: //
644: // p -----F----- p' with xi = 1 - p'^2/p_s
645: //
646: //
647: // Factors applied here, compatible with 2 -> N phase space sampling:
648: // 1/xi ["sub Moller flux]
649: // 1/pt2 ["kt-factorization] (cancels with pt2 from numerator)
650: // 16pi^2 ["kinematics volume factor]
651: //
652: double CohFlux(double xi, double t, double pt) {

```

```

./src/MForm.cc 9/11
665: const double pt2 = pow2(pt);
666: const double x12 = pow2(x1);
667: const double mp2 = pow2(mp);
668: const double Q2 = std::abs(Q2);
669:
670: const double PART1 =
671:   (4.0 * pow2(mp) * pow2(G_E(Q2)) + Q2 * pow2(G_M(Q2))) / (4.0 * pow2(mp) + Q2);
672: const double PART2 = pow2(G_M(Q2));
673: const double DELTA = pt2 / (pt2 + x12 * mp2); // Bjorken x
674: double f =
675:   qed:alpha_QED() / PI * ((1.0 - xi) * pow2(DELTA) * PART1 + (x12 / 4.0) * DELTA * PART2);
676:
677: // Factors
678: f *= x1;
679: f /= pt2;
680: f *= 16.0 * gra:math:PIPI;
681:
682: return f; // Use at cross section level
683: }
684:
685: // Incoherent photon flux from a dissociated proton with mass M.
686: // When M = mp, this reproduces CohFlux() if
687: // F2(N,Q^2) also reproduces the elastic limit (not all parametrizations do)
688: //
689: //
690: // p -----F2(N,Q^2)-----> p' with xi = 1 - p'^2 / p_s
691: //
692: //
693: //
694: //
695: //
696: //
697: // Factors applied as with CohFlux() above.
698: //
699: //
700: double IncohFlux(double x1, double t, double pt, double M2) {
701:   const exp: double mp2 = pow2(mp);
702: }
703: const double pt2 = pow2(pt);
704: const double x12 = pow2(x1);
705: const double Q2 = std::abs(Q2);
706: const double xb_j = Q2 / (Q2 + M2 - mp2); // Bjorken x
707: const double DELTA = pt2 / (pt2 + x12 * (M2 - mp2) + x12 * mp2);
708:
709: double f = qed:alpha_QED() / PI *
710:   ((1.0 - x1) * pow2(DELTA) * F2xQ(xb_j, Q2) * (Q2 + M2 - mp2) +
711:     (x12 / (4.0 * pow2(xb_j))) * DELTA * 2.0 * xb_j * F1xQ(xb_j, Q2) / (Q2 + M2 - mp2));
712:
713: // Factors
714: f *= x1;
715: f /= pt2;
716: f *= 16.0 * gra:math:PIPI;
717:
718: return f; // Use at cross section level
719: }
720:
721: // Drees-Zeppenfeld proton coherent gamma flux (collinear)
722: double DZFux(double x) {
723:   const double Q2min = (pow2(mp) * pow2(x)) / (1.0 - x);
724:   const double A = 1.0 + 0.71 / Q2min;
725: }
726: double f = qed:alpha_QED() / (2.0 * PI * x) * (1.0 + pow2(1.0 - x)) *
727:   (std::log(A) - 11.0 / 4.0 + 3.0 / A - 3.0 / (2.0 * pow2(A)) + 1.0 / (3 * pow3(A)));
728:
729: return f; // Use at cross section level
730: }
731:
732: // Breit-Wigner propagators / form factors
733: //
734: //
735: //
736: // Useful identity for normalization:
737: //
738: // int dm^2 1/(m^2 - M0^2)^2 + M0^2 Gamma^2) * logiv 1/(pi|M0 Gamma)
739: //
740: // based on squaring the complex propagator
741: // Dirac^2 = J / (m^2 - M0^2 + iM0 Gamma), and integrating.
742: //
743: std::complex<double> CBW(const gra:LOBNITS_SCALAR lts, const gra:PARAM_RES resonance) {
744:   switch (resonance.BW) {
745:     case 1:
746:       return CBW_FW(lts,m2, resonance.p.mass, resonance.p.width);
747:     case 2:

```

```

./src/MForm.cc 10/11
748: return CBW_RW(lts,m2, resonance.p.mass, resonance.p.width);
749: case 3:
750:   return CBW_BF(lts,m2, resonance.p.mass, resonance.p.width, resonance.p.spinX2 / 2.0,
751:     lts.decaytree[0].p4.M(), lts.decaytree[1].p4.M());
752: case 4:
753:   return CBW_JR(lts,m2, resonance.p.mass, resonance.p.width, resonance.p.spinX2 / 2.0);
754: default:
755:   throw std::invalid_argument("CBW: Unknown BW (Breit-Wigner) parameter: " +
756:     std::to_string(resonance.BW));
757: }
758: }
759:
760: // See e.g.
761: // [REFERENCE: TASI Lectures on propagators, users.lctip.it/~nmr244/tait-supplemental.pdf]
762: // [REFERENCE: Ciaccipaglia, Deandrea, Curtis, arxiv.org/abs/0906.3417v2]
763: // [REFERENCE:
764: // <www.t2.ucsd.edu/twiki/pub/UCSDT2/Physics2148spring2015/sjw-breit-wigner-cbw9-95.pdf>
765: //
766: //
767: // -----
768: // Delta function Delta(hat{s}) - M0^2) replacement function:
769: // int dhat{s} delta(hat{s}) - M0^2) = int dhat{s}
770: // delta2BW(hat{s},M0,Gamma)
771: //
772: // To be applied at cross section level
773: // double delta2BWase(double shat, double M0, double Gamma) {
774: //   return msqrt(delta2BWsect(shat, M0, Gamma));
775: // }
776: //
777: // To be applied at amplitude level
778: // double delta2BWamp(double shat, double M0, double Gamma) {
779: //   return msqrt(delta2BWsect(shat, M0, Gamma));
780: // }
781: // -----
782: // Breit-Wigner propagator parametrizations
783: //
784: //
785: //
786: // 1. Complex Fixed Width Relativistic Breit-Wigner
787: // std::complex<double> CBW_FW(double m2, double M0, double Gamma) {
788: //   return -1.0 / (m2 - M0 * M0 + z1 * M0 * Gamma);
789: // }
790: //
791: // 2. Complex Running Width Relativistic Breit-Wigner
792: // std::complex<double> CBW_RW(double m2, double M0, double Gamma) {
793: //   return -1.0 / (m2 - M0 * M0 + z1 * std::sqrt(m2) * Gamma);
794: // }
795: //
796: // 3. J = 0,1,2 Complex Relativistic Breit-Wigner
797: // with angular barrier effects type of Blatt-Weisskopf
798: // m_A and m_B are the masses of daughters (GeV)
799: // std::complex<double> CBW_BF(double m2, double M0, double Gamma, int J, double mA, double mB) {
800: //   const double s = msqrt(m2 - pow2(mA + mB)) * (m2 - pow2(mA - mB)) / (2 * msqrt(m2));
801: //   const double d = msqrt((M0 * M0 - pow2(mA + mB)) * (M0 * M0 - pow2(mA - mB))) / (2 * M0);
802: //   return -1.0 / (m2 - M0 * M0 + z1 * Gamma * M0 / msqrt(m2) * Bfactor);
803: // }
804: //
805: //
806: // 4. Spin dependent relativistic Breit-Wigner (should not be used blindly!)
807: //
808: // [REFERENCE: Alwall et al., arxiv.org/abs/1402.1178]
809: // std::complex<double> CBW_JR(double m2, double M0, double Gamma, double J) {
810: //   const std::complex<double> denom = (m2 - M0 * M0 + z1 * M0 * Gamma);
811: // }
812: // if (static_cast<int>(J) == 0) { // J = 0
813: //   return -1.0 / denom;
814: // } else if (static_cast<int>(J) == 1) { // J = 1/2
815: //   return -2.0 * msqrt(m2) / denom;
816: // } else if (static_cast<int>(J) == 1) { // J = 1
817: //   return -(1.0 - m2 / (M0 * M0)) / denom;
818: // } else if (static_cast<int>(J) == 3) { // J = 3/2
819: //   return -(2.0 / 3.0) * msqrt(m2) * (1.0 - m2 / (M0 * M0)) / denom;
820: // } else if (static_cast<int>(J) == 2) { // J = 2
821: //   return -(0.0 / 6.0 * (4.0 / 3.0) * (m2 / (M0 * M0) +
822: //     (2.0 / 3.0) * (m2 * m2) / (gra:math:pow4(M0))) /
823: //     denom;
824: // } else {
825: //   return -1.0 / denom; // Too high spin
826: // }
827: // }
828: //
829: // namespace form
830: // }

```

```

831:
832:
833:
834:
835:
836: // Coupling schemes
837:
838: // Dirac: alpha_g = g^2/(4pi)
839: // Beta-dirac alpha_g = (g*beta)^2 / (4pi)
840:
841: // [REFERENCE: Rajantie, physica_today.acitation.org/doi/pdf/10.1063/PT.3.3328]
842: // [REFERENCE: Dougall, Wick, arxiv.org/abs/0706.10470]
843: // [REFERENCE: Reis, Sauter, arxiv.org/abs/1707.04170v1]
844:
845: double MGamma::yyfbar(gra:LORENTZSCALAR kIts) {
846: // QED couplings
847: const double COUPL = 16.0 * pow(gra:math:PI * qed:alpha_QED()); // = e^4
848: const double mass = Its.DecayTree[0].p4.M(); // lepton, quark (or monopole) mass
849: const double mass2 = pow2(mass);
850: const bool MONOPOLE_MODE = (Its.DecayTree[0].p.pdg == PDG:PDG_monopole) ? true : false;
851:
852: if (PARAM_MONOPOLE:ign < 1) {
853: throw std::invalid_argument("MGamma::yyfbar: Parameter Dirac n less than 1");
854: }
855:
856: // Amplitude squared
857: double amp2 = 0.0;
858:
859: // Monopole-Antimonopole coupling
860: if (MONOPOLE_MODE) {
861: static const double g = 2.0 * math:PI * PARAM_MONOPOLE:sign / qed:alpha_QED();
862:
863: if (PARAM_MONOPOLE:icoupling == "beta-dirac") {
864: // Calculate beta (velocity)
865: // const: M4Vec p3 = pfinal[3] + pfinal[4];
866: // M4Vec p3 = pfinal[3];
867: // LorentzBoost(p3, p3.M(), p3, -1); // Boost to the
868: // CM frame of Mbar[M] pair
869: // double beta = std::sqrt( p3.Fx()*p3.Fx() +
870: // p3.Fy()*p3.Fy() + p3.Fz()*p3.Fz() ) / p3.E();
871:
872: // Faster way
873: const double beta = msqrt(1.0 - 4.0 * pow2(PARAM_MONOPOLE:M0) / Its.s_hat);
874: COUPL = pow4(g * beta);
875: } else if (PARAM_MONOPOLE:icoupling == "dirac") {
876: COUPL = pow4(g);
877: } else {
878: throw std::invalid_argument("MGamma::yyfbar: Unknown PARAM_MONOPOLE:icoupling " +
879: PARAM_MONOPOLE:icoupling);
880: }
881:
882: // QED tree level amplitude squared |M|^2, spin averaged and
883: // summed
884: amp2 = 2.0 * COUPL *
885: (Its.u_hat - mass2) / (Its.t_hat - mass2) + (Its.t_hat - mass2) / (Its.u_hat - mass2) +
886: 1.0
887: pow2(1.0 + (2.0 * mass2) / (Its.t_hat + mass2) + (2.0 * mass2) / (Its.u_hat + mass2));
888:
889: //
890: // FeynCalc result, less simplified, but exactly same result
891: amp2 = 2 * COUPL *
892: ( pow2(mass) * (3*pow2(Its.t_hat) +
893: 14*Its.t_hat*Its.u_hat + 3*pow2(Its.u_hat)
894: -mass2) * pow3(Its.t_hat) +
895: 7*pow2(Its.t_hat) * pow2(Its.u_hat) * Its.t_hat +
896: 4*Its.t_hat * Its.u_hat )
897: + 6*pow2(mass) *
898: Its.t_hat * Its.u_hat * (pow2(Its.t_hat) + pow2(Its.u_hat)) // (
899: pow2(Its.t_hat - mass2) * pow2(Its.u_hat - mass2) );
900:
901: //
902: // print("40.15f %0.15f in", amp2, amp2);
903:
904: //
905: // [MADGRAPH/HELAS, all helicity amplitudes individually -> for
906: // the screening loop
907: return AmpM5_yy_ll.CalcAmp2(Its, 0.0);
908: }
909:
910:
911:
912: // quark pair (charge 1/3 or 2/3), apply charge and color factors
913: if (std::abs(Its.DecayTree[0].p.pdg) <= 6) { // we have a quark
914: const double Q = Its.DecayTree[0].p.chargeX3 / 3.0;
915: const double NC = 3.0; // quarks come in three colors

```

```

1: // Gamma-Gamma Amplitudes
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT license <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <random>
9: #include <vector>
10:
11: // Own
12: #include "Granitti/MForm.h"
13: #include "Granitti/MGamma.h"
14: #include "Granitti/MCubola.h"
15: #include "Granitti/MKinematics.h"
16: #include "Granitti/Match.h"
17: #include "Granitti/MFDG.h"
18: #include "Granitti/MQED.h"
19: #include "Granitti/MQEDg.h"
20: #include "Granitti/MSpin.h"
21:
22:
23: using gra:math:PI;
24: using gra:math:abs;
25: using gra:math:msqrt;
26: using gra:math:pow2;
27: using gra:math:pow4;
28: using gra:math:pow4;
29: using gra:math:xi;
30:
31: using namespace gra:form;
32:
33: namespace gra {
34:
35: // Constructor
36: MGamma(MGamma(gra:LORENTZSCALAR kIts, const std::string& defFile) {
37: // @see MUST_INCREMENTING_LOCK_NEEDED FOR THE INITIALIZATION @
38: gra:ig_mutex.lock();
39:
40: // Monopolum process
41: if (PARAM_MONOPOLE:initialized) {
42: try {
43: PARAM_MONOPOLE:M0 = Its.PDG.FindByPDG(PDG:PDG_monopole).mass;
44: PARAM_MONOPOLE:ReadParameters(defFile);
45: } catch (...) {
46: gra:ig_mutex.unlock(); // need to release here, otherwise got infinite lock
47: throw;
48: }
49: }
50: gra:ig_mutex.unlock();
51:
52:
53: // =====
54: // yy -> fermion-antifermion pair
55: // yy -> e+e-, mu+mu-, tau+tau-, gbar or Monopole-Antimonopole (spin-1/2
56: // monopole)
57: // production
58: //
59: //
60: // This is the same amplitude as yy -> e+e- (two diagrams)
61: // obtained by crossing e+e- -> yy annihilation, see:
62: //
63: // http://theory.sinp.msu.ru/comphsp_old/tutorial/QED/node4.html
64: //
65: // beta = v/c of the monopole (or antimonopole) in their CM system
66: //
67: // =====
68: // Dirac quantization condition:
69: //
70: // g = 2pi lbarbar n / (lmu_0 e), where n = 1,2,3,...
71: // where alpha = e^2/(4*PI) is the running QED coupling
72: //
73: //
74: //
75: // Numerical values:
76: // alpha_g = g^2 / (4*PI) ^ 34 (when n = 1)
77: // alpha_em = e^2 / (4*PI) ^ 1/137
78: // =====
79: //
80: // From lepton pair to monopole pair:
81: // replace e -> g*beta
82: //
83: // 16 * pow2(PI * alpha_EM) -> pow(g*beta, 4)
84:
85:
86:
87:
88:
89:
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164:
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166:

```

```

85: //
86: // =====
87: //
88: // Dirac: alpha_g = g^2/(4pi)
89: // Beta-dirac alpha_g = (g*beta)^2 / (4pi)
90:
91: // [REFERENCE: Rajantie, physica_today.acitation.org/doi/pdf/10.1063/PT.3.3328]
92: // [REFERENCE: Dougall, Wick, arxiv.org/abs/0706.10470]
93: // [REFERENCE: Reis, Sauter, arxiv.org/abs/1707.04170v1]
94:
95: double MGamma::yyfbar(gra:LORENTZSCALAR kIts) {
96: // QED couplings
97: const double COUPL = 16.0 * pow(gra:math:PI * qed:alpha_QED()); // = e^4
98: const double mass = Its.DecayTree[0].p4.M(); // lepton, quark (or monopole) mass
99: const double mass2 = pow2(mass);
100: const bool MONOPOLE_MODE = (Its.DecayTree[0].p.pdg == PDG:PDG_monopole) ? true : false;
101:
102: if (PARAM_MONOPOLE:ign < 1) {
103: throw std::invalid_argument("MGamma::yyfbar: Parameter Dirac n less than 1");
104: }
105:
106: // Amplitude squared
107: double amp2 = 0.0;
108:
109: // Monopole-Antimonopole coupling
110: if (MONOPOLE_MODE) {
111: static const double g = 2.0 * math:PI * PARAM_MONOPOLE:sign / qed:alpha_QED();
112:
113: if (PARAM_MONOPOLE:icoupling == "beta-dirac") {
114: // Calculate beta (velocity)
115: // const: M4Vec p3 = pfinal[3] + pfinal[4];
116: // M4Vec p3 = pfinal[3];
117: // LorentzBoost(p3, p3.M(), p3, -1); // Boost to the
118: // CM frame of Mbar[M] pair
119: // double beta = std::sqrt( p3.Fx()*p3.Fx() +
120: // p3.Fy()*p3.Fy() + p3.Fz()*p3.Fz() ) / p3.E();
121:
122: // Faster way
123: const double beta = msqrt(1.0 - 4.0 * pow2(PARAM_MONOPOLE:M0) / Its.s_hat);
124: COUPL = pow4(g * beta);
125: } else if (PARAM_MONOPOLE:icoupling == "dirac") {
126: COUPL = pow4(g);
127: } else {
128: throw std::invalid_argument("MGamma::yyfbar: Unknown PARAM_MONOPOLE:icoupling " +
129: PARAM_MONOPOLE:icoupling);
130: }
131:
132: // QED tree level amplitude squared |M|^2, spin averaged and
133: // summed
134: amp2 = 2.0 * COUPL *
135: (Its.u_hat - mass2) / (Its.t_hat - mass2) + (Its.t_hat - mass2) / (Its.u_hat - mass2) +
136: 1.0
137: pow2(1.0 + (2.0 * mass2) / (Its.t_hat + mass2) + (2.0 * mass2) / (Its.u_hat + mass2));
138:
139: //
140: // FeynCalc result, less simplified, but exactly same result
141: amp2 = 2 * COUPL *
142: ( pow2(mass) * (3*pow2(Its.t_hat) +
143: 14*Its.t_hat*Its.u_hat + 3*pow2(Its.u_hat)
144: -mass2) * pow3(Its.t_hat) +
145: 7*pow2(Its.t_hat) * pow2(Its.u_hat) * Its.t_hat +
146: 4*Its.t_hat * Its.u_hat )
147: + 6*pow2(mass) *
148: Its.t_hat * Its.u_hat * (pow2(Its.t_hat) + pow2(Its.u_hat)) // (
149: pow2(Its.t_hat - mass2) * pow2(Its.u_hat - mass2) );
150:
151: //
152: // print("40.15f %0.15f in", amp2, amp2);
153:
154: //
155: // [MADGRAPH/HELAS, all helicity amplitudes individually -> for
156: // the screening loop
157: return AmpM5_yy_ll.CalcAmp2(Its, 0.0);
158: }
159:
160:
161:
162: // quark pair (charge 1/3 or 2/3), apply charge and color factors
163: if (std::abs(Its.DecayTree[0].p.pdg) <= 6) { // we have a quark
164: const double Q = Its.DecayTree[0].p.chargeX3 / 3.0;
165: const double NC = 3.0; // quarks come in three colors

```

```

167:
168: const double factor = pow4(Q * NC);
169: const double sqrt_factor = msqrt(factor); // sqrt to "amplitude level"
170:
171: // amplitude squared
172: amp2 = factor;
173: // helicity amplitudes
174: for (const auto& i : aux:indices(Its.hamp)) { Its.hamp[i] *= sqrt_factor; }
175:
176:
177: // =====
178: // For screening loop (approximation)
179: Its.hamp = (msqrt(amp2));
180:
181:
182: return amp2;
183:
184:
185: // =====
186: // A Monopolum (monopole-antimonopole) bound state process
187: //
188: // See e.g.
189: //
190: // [REFERENCE: Preskill,
191: // http://www.theory.caltech.edu/~preskill/pubs/preskill-1984-monopoles.pdf]
192: // [REFERENCE: Spele, Franchiotti, Garcia, Canal, Vento,
193: // https://arxiv.org/abs/hep-th/0701302v2]
194: // [REFERENCE: Barrio, Sugamoto, Yamashita, https://arxiv.org/abs/1607.03987v3]
195: // [REFERENCE: Franchiotti, Canal, Vento, https://arxiv.org/pdf/1703.06649.pdf]
196: // [REFERENCE: Reis, Sauter, https://arxiv.org/abs/1707.04170v1]
197: // =====
198: //
199: // Dirac quantization condition:
200: //
201: // g = 2pi lbarbar (lmu_0 e) n, where n = 1,2,3,...
202: // where alpha = e^2/(4*PI) is the running QED coupling
203: //
204: //
205: // Numerical values:
206: // alpha_g = g^2 / (4*PI) ^ 34 (when n = 1)
207: // alpha_em = e^2 / (4*PI) ^ 1/137
208: // =====
209: //
210: double MGamma::yymp(gra:LORENTZSCALAR kIts) const {
211: // Easy printing
212: gra:ig_mutex.lock();
213: PARAM_MONOPOLE:PrintParameters(Its.sqrt_s);
214: gra:ig_mutex.unlock();
215:
216: // Monopolum nominal mass and width parameters
217: static const double M = 2.0 * PARAM_MONOPOLE:M0 + PARAM_MONOPOLE:EMeryMP(PARAM_MONOPOLE:fin);
218: static const double Gamma_M = PARAM_MONOPOLE:GammaM0;
219:
220: if (M < 0) {
221: throw std::invalid_argument("MGamma::yymp: Increases ladder level parameter En. Monopolum "
222: "nominal mass " +
223: std::ito_string(M) + " < 0!");
224: }
225:
226: if (PARAM_MONOPOLE:ign < 1) {
227: throw std::invalid_argument("MGamma::yymp: Parameter Dirac n less than 1");
228: }
229:
230: // Two coupling scenarios:
231: static const double g = 2.0 * math:PI * PARAM_MONOPOLE:ign / qed:alpha_QED();
232: double beta = pow2(M / Its.t_hat);
233:
234: if (PARAM_MONOPOLE:icoupling == "beta-dirac") {
235: beta = msqrt(1.0 - pow2(M / Its.t_hat));
236: } else if (PARAM_MONOPOLE:icoupling == "dirac") {
237: beta = 1.0;
238: } else {
239: throw std::invalid_argument("MGamma::yymp: Unknown PARAM_MONOPOLE:icoupling " +
240: PARAM_MONOPOLE:icoupling);
241: }
242:
243: // Magnetic coupling
244: const double alpha_g = pow2(beta * g) / (4.0 * math:PI);
245:
246: // Running width
247: const double Gamma_E = PARAM_MONOPOLE:GammaMP(PARAM_MONOPOLE:fin, alpha_g);
248:
249: // print("alpha_g = 40.3E, Gamma_E = 40.3E, Gamma_M = 40.3E, Psi_MP = 40.3E

```



```

84: // * *
85: // *p*
86: // * M^2
87: // k2 * M^2
88: // a ===== a
89: //
90: // CD:
91: //
92: //
93: // a gW a a gW a
94: //
95: // i1 * j1
96: // *p*
97: // k * M^2
98: // *p*
99: // *p*
100: // i2 * j2
101: //
102: // a gW a a gW a
103: //
104: //
105: //
106: // [REFERENCE: Gribov, A Reggeon Diagram Technique, Soviet JETP, 1966,
107: // jetp.ac.ru/cgi-bin/dm/gw/G26_02_044.pdf]
108: //
109: // Strong coupling in the Pomeronchuk pole problem
110: // [REFERENCE: Gribov, Mgalia, Sov. Phys. JETP 28(4), 794-795 (1968)]
111: // [REFERENCE: Muller, 1972]
112: //
113: // For different forms of triple Pomeron coupling (weak/strong, scalar/vector):
114: // [REFERENCE: Luna, Khoze, Martin, Ryskin, arxiv.org/abs/1005.4864v1]
115: //
116: std::complex<double> MRegge(ME2(gra:LORRENTZSCALAR ilts, int mode) const {
117:   const double s0 = 1.0; // Energy scale GeV^2
118:
119:   // Pomeron trajectory intercept + 1
120:   const double alpha_0 = 1.0 + PARAM_SOFT:DELTA_T;
121:
122:   // Triple legs
123:   const double alpha_t_1 = S3PonAlpha(its.t);
124:   const double alpha_t_2 = alpha_t_1;
125:
126:   std::complex<double> A(0, 0);
127:
128:   if (mode == 1) { // Elastic
129:
130:     // Single Pomeron exchange
131:     // (this is the eikonal pomeron without eikonalization, just for test, use loop on)
132:     return MEikonal:SingleAmpElastic(its.s, its.t, 1);
133:
134:   } else if (mode == 2) { // SD triple Pomeron
135:
136:     // Which proton is excited
137:     const double MX_2 = (its.ss[1][1] > 1.0) ? its.ss[1][1] : its.ss[2][2];
138:
139:     const double amp2 = pow3(PARAM_SOFT:qgP) // Proton-Pomeron coupling ^ 3
140:       * pow2(S3P(its.t)) // Proton form factor ^ 2
141:       * PARAM_SOFT:g3P // Triple-Pomeron coupling
142:
143:       * std::pow(its.s / M0_2, // LEFT + RIGHT LEG
144:         alpha_t_1 + alpha_t_2); // LEFT + RIGHT LEG
145:
146:     * std::pow(MX_2 / s0, alpha_0); // DISCONTINUITY PART
147:
148:     // recover amplitude, we had the expression at cross section level
149:     A = transform(MRegge(alpha_t_1, 1) * msqrt(amp2));
150:
151:   } else if (mode == 3) { // SD triple Pomeron, note NO proton form
152:     // factor (both proton dissociate)
153:
154:     const double MX_2 = its.ss[1][1];
155:     const double MY_2 = its.ss[2][2];
156:
157:     const double amp2 = pow2(PARAM_SOFT:qgP) // Proton-Pomeron ^ 2
158:       * pow2(PARAM_SOFT:g3P) // Triple-Pomeron coupling ^ 2
159:
160:       * std::pow(its.s * s0 / (MX_2 * MY_2),
161:         alpha_t_1 + alpha_t_2); // LEFT + RIGHT LEG
162:
163:     * std::pow(MX_2 / s0, alpha_0) // DISCONTINUITY PART
164:     * std::pow(MY_2 / s0, alpha_0); // DISCONTINUITY PART
165:
166:     // recover amplitude, we had the expression at cross section level

```

```

250:   continue;
251: }
252: // Apply lower vertex helicity conservation
253: if (g_Vertex(its.t2, lambda_d[1][1], lambda_d[1][3]) !=
254:     std::pow(-1, lambda_d[1][1] - lambda_d[1][3]))
255:   continue;
256: }
257: }
258:
259: // Spin density matrix weight for this helicity
260: const int index = index + lambda_d[1][4] + resonance.p.spinX2; // Index the diagonal
261: const std::complex<double> rhosight =
262:   resonance.p.spinX2 != 0 ? resonance.rho[index][index] : 1.0;
263:
264: // Calculate amplitude
265: const std::complex<double> amp =
266:   rhosight * g_Vertex(its.t1, lambda_d[1][0], lambda_d[1][2]) *
267:   gik_Vertex(its.t1, its.t2, dphi, lambda_d[1][4], resonance.p.spinX2 / 2.0, resonance.p.P,
268:     common);
269: g_Vertex(its.t2, lambda_d[1][1], lambda_d[1][3]);
270:
271: // std::cout << amp << " : : << gra:math:abs2(amp) << std::endl;
272: its.hamp.push_back(amp);
273: }
274: }
275: // 1/4 (|M1|^2 + |M2|^2 + ... + |Mn|^2)
276: double amp2 = 0.0;
277: for (std::size_t i = 0; i < its.hamp.size(); ++i) { amp2 += gra:math:abs2(its.hamp[i]); }
278: amp2 *= 4; // Initial state average
279: return msqrt(amp2); // we expect amplitude
280:
281: // This function is used for parity conservation check:
282:
283: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
284: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
285: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
286: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
287: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
288: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
289: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
290: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
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294: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
295: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
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331: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h
332: // [ gamma_m1,m2 ] \ lambda_h = (-1) ^ lambda_h \lambda_i3 [ gamma_m1,-m2 ] / \ lambda_h

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167:   A = gra:form:ReggeEta(alpha_t_1, 1) * msqrt(amp2);
168: } else {
169:   throw std::invalid_argument("MRegge:ME2: Unknown mode parameter: " + std::ito_string(mode));
170: }
171:
172: // For screening loop
173: its.hamp *= (A);
174:
175: //
176: //
177: return A;
178: }
179:
180: // Helicity matrix element for Pomeron-Pomeron resonances
181: // [THIS IS UNDER CONSTRUCTION!]
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415: // 2. parity conservation

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416: // Particle-Particle-Pomeron coupling
417: const double gpp_P = PARAM_REGGE::c(t0) / PARAM_SOFT::gM_P;
418:
419:
420: const std::complex<double> A_t = PropOnly(its.s1[1][3], its.t1) * FF_A * (FF(its.t_hat, M2_M) *
421: gpp_P * prop(its.t_hat, M2_M) + (FF(its.t_hat, M2_M) * gpp_P *
422: PropOnly(its.s2[1][4], its.t2) * FF_B);
423:
424: // sign applied here
425: const std::complex<double> A_u =
426: sign * PropOnly(its.s1[1][4], its.t1) * FF_A * (FF(its.t_hat, M2_M) * gpp_P *
427: prop(its.t_hat, M2_M) + (FF(its.t_hat, M2_M) * gpp_P * PropOnly(its.s2[1][3], its.t2) * FF_B);
428:
429: // Total amplitude
430: const std::complex<double> A = A_t + A_u;
431:
432:
433: // For screening loop
434: its.hamp = (A);
435: // -----
436:
437: return A;
438: }
439:
440: // =====
441: // Regge matrix element ansatz for 2->2 Continuum spectrum
442: // -----
443: // =====
444: // =====
445: // +
446: // +
447: // zf----- a
448: // |
449: // |
450: // |
451: // +
452: // +
453: // zf----- c
454: // |
455: // |
456: // +
457: // =====
458: // =====
459: // For many similar amplitudes, see the literature of the era
460: // "pion-supercritical/generalized Veneziano amplitude". For example:
461: // [REFERENCE: Bardakci, Ruegg, journals.aps.org/pr/pdf/10.1103/PhysRev.181.1884]
462: // [REFERENCE: Hardaki, Ruegg, journals.aps.org/pr/pdf/10.1103/PhysRev.181.1884]
463: // -----
464: // -----
465: // Numerically, one may compare with:
466: // -----
467: // [REFERENCE: Kycia, Ledwiesko, Szczurek, Turnau, arxiv.org/abs/1702.07572]
468: // -----
469: // -----
470: std::complex<double> MRegge::ME6(gra:LORENTZSCALAR &its) const {
471: // Amplitude
472: std::complex<double> A(0, 0);
473:
474: // Offshell propagator masses^2 [for mixed states, such as pi+ pi- K+ K-]
475: const double M2_A = pow2(its.decaytree[0].p.mass);
476: const double M2_B = pow2(its.decaytree[1].p.mass);
477:
478: // Create function pointers
479: double (*f1)(double, double);
480: double (*f2)(double, double);
481:
482: const double sign = 1;
483:
484: if (sign > 0) { // positive sign amplitudes
485:
486: ff = &PARAM_REGGE::Meson_FF;
487: prop = &PARAM_REGGE::Meson_prop;
488: } else { // negative sign (alternative model spin-statistics) amplitude
489:
490: ff = &PARAM_REGGE::Baryon_FF;
491: prop = &PARAM_REGGE::Baryon_prop;
492: }
493:
494: // Proton / Dissociative system vertices
495: double FF_A = 0.0;
496: double FF_B = 0.0;
497: PomProtonVertex(its, FF_A, FF_B);
498:

```

```

582: ff = &PARAM_REGGE::Baryon_FF;
583: prop = &PARAM_REGGE::Baryon_prop;
584: }
585:
586: // Proton / Dissociative system vertices
587: double FF_A = 0.0;
588: double FF_B = 0.0;
589: PomProtonVertex(its, FF_A, FF_B);
590:
591: // Particle-Particle-Pomeron coupling
592: const double gpp_P = PARAM_REGGE::c(t0) / PARAM_SOFT::gM_P;
593:
594: // Loop over different permutations (max #289)
595: for (const auto i1 : indices(permutations6)) {
596: const unsigned int a = permutations6[1][0];
597: const unsigned int b = permutations6[1][1];
598: const unsigned int c = permutations6[1][2];
599: const unsigned int d = permutations6[1][3];
600: const unsigned int e = permutations6[1][4];
601: const unsigned int f = permutations6[1][5];
602:
603: // t-type Lorentz scalars [no need here, already calculated]
604: // const double tt_ab = (pbearm1.pfinal1 - pfinal1[0]).M2();
605: // const double tt_bc = (pbearm1.pfinal1 - pfinal1[1]).M2();
606: // const double tt_cd = (pbearm1.pfinal1 - pfinal1[2]).M2();
607: // const double tt_de = (pbearm1.pfinal1 - pfinal1[3]).M2();
608: // const double tt_ef = (pbearm1.pfinal1 - pfinal1[4]).M2();
609:
610: // collect scalars
611: const double tt_ab = its.t1[a];
612: const double tt_bc = its.t1_xy[a][b];
613: const double tt_cd = its.t1_xy[a][c];
614: const double tt_de = its.t1_xy[c][d];
615: const double tt_ef = its.t1[d];
616:
617: const std::complex<double> subamp =
618: PropOnly(its.s1[1][a], its.t1) * FF_A * (FF(tt_ab, M2_A) * gpp_P * prop(tt_ab, M2_A) *
619: (FF(tt_bc, M2_B) * gpp_P * PropOnly(its.s2[1][c], tt_bc) + (FF(tt_bc, M2_B) * gpp_P *
620: prop(tt_cd, M2_B) + (FF(tt_de, M2_B) * gpp_P * PropOnly(its.s2[1][d], tt_de) +
621: (FF(tt_de, M2_C) * gpp_P * prop(tt_ef, M2_C) + (FF(tt_de, M2_C) * gpp_P *
622: PropOnly(its.s2[1][f], its.t2) * FF_B);
623:
624: A += subamp;
625: }
626:
627: // -----
628: // For screening loop
629: its.hamp = (A);
630: // -----
631:
632: return A;
633: }
634:
635: // Proton-Proton-Pomeron elastic / inelastic vertex
636: // -----
637: void MRegge::PomProtonVertex(const gra:LORENTZSCALAR &its, double &FF_A, double &FF_B) const {
638: FF_A = its.excite1 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
639: : PARAM_SOFT::gM_P * gra:form:ISF(its.t1);
640: FF_B = its.excite2 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
641: : PARAM_SOFT::gM_P * gra:form:ISF(its.t2);
642:
643: }
644:
645:
646: // -----
647: // Generic resonance (sub)-cross section:
648: // -----
649: // |hat(s)|^2 = |pi |hat(s)|^2 = |W_f(hat(s)) W_f(hat(s))|^2
650: // |hat(s)|^2 = |hat(s)|^2 = |W_f(hat(s))|^2
651: // -----
652: // where M_{1,2} denotes effective initial (final) state factors
653: // Incorporating spin and color symmetry factors, couplings and the rest of the
654: // dynamics
655: // -----
656: // In general, the behavior can be different in domains:
657: // -----
658: // 1. hat(s) << M0^2, where M0 is the on-shell mass
659: // 2. hat(s) ~ M0^2
660: // 3. hat(s) >> M0^2
661: // -----
662: // and isolated resonances might break unitarity when |hat(s)| -> inf,
663: // but resonance + continuum amplitude does not.

```

```

499: // Particle-Particle-Pomeron coupling
500: const double gpp_P = PARAM_REGGE::c(t0) / PARAM_SOFT::gM_P;
501:
502: // Loop over different final state permutations (max #16)
503: for (const auto i1 : indices(permutations4)) {
504: const unsigned int a = permutations4[1][0];
505: const unsigned int b = permutations4[1][1];
506: const unsigned int c = permutations4[1][2];
507: const unsigned int d = permutations4[1][3];
508:
509: // Calculate t-type Lorentz scalars here [no need, done already]
510: // const double tt_ab = (pbearm1.pfinal1 - pfinal1[0]).M2();
511: // const double tt_bc = (pbearm1.pfinal1 - pfinal1[1]).M2();
512: // symmtry, same as (pbearm2 - pfinal1[2] - pfinal1[3] - pfinal1[4]).M2()
513: // const double tt_cd = (pbearm2.pfinal1 - pfinal1[2]).M2();
514:
515: // Get Lorentz scalars
516: const double tt_ab = its.t1[a];
517: const double tt_bc = its.t1_xy[a][b];
518: const double tt_cd = its.t1[d];
519:
520: const std::complex<double> subamp =
521: PropOnly(its.s1[1][a], its.t1) * FF_A * (FF(tt_ab, M2_A) * gpp_P * prop(tt_ab, M2_A) *
522: (FF(tt_bc, M2_B) * gpp_P * PropOnly(its.s2[1][c], tt_bc) + (FF(tt_bc, M2_B) * gpp_P *
523: prop(tt_cd, M2_B) + (FF(tt_cd, M2_B) * gpp_P * PropOnly(its.s2[1][d], its.t2) * FF_B);
524:
525: A += subamp;
526: }
527:
528: // -----
529: // For screening loop
530: its.hamp = (A);
531: // -----
532:
533: return A;
534: }
535:
536: // =====
537: // Regge matrix element ansatz for 2->8 continuum spectrum (generalization of
538: // 2->4)
539: // -----
540: // =====
541: // +
542: // +
543: // zf----- a
544: // |
545: // |
546: // |
547: // +
548: // +
549: // zf----- c
550: // |
551: // |
552: // +
553: // +
554: // zf----- e
555: // |
556: // |
557: // +
558: // =====
559: // =====
560: std::complex<double> MRegge::ME8(gra:LORENTZSCALAR &its) const {
561: // Amplitude
562: std::complex<double> A(0, 0);
563:
564: // Offshell propagator masses^2 [for mixed states, such as pi+ pi- K+ K- pi+
565: // pi-]
566: const double M2_A = pow2(its.decaytree[0].p.mass);
567: const double M2_B = pow2(its.decaytree[2].p.mass);
568: const double M2_C = pow2(its.decaytree[4].p.mass);
569:
570: // Create function pointers
571: double (*f1)(double, double);
572: double (*f2)(double, double);
573:
574: const double sign = 1;
575:
576: if (sign > 0) { // positive sign amplitudes
577:
578: ff = &PARAM_REGGE::Meson_FF;
579: prop = &PARAM_REGGE::Meson_prop;
580: } else { // negative sign (alternative model spin-statistics) amplitude
581:

```

```

666: // =====
667: // Simple matrix element ansatz for Pomeron-Pomeron resonances
668: // -----
669: // =====
670: // +
671: // +
672: // z-----
673: // |
674: // |
675: // +
676: // =====
677: // =====
678: std::complex<double> MRegge::ME3(gra:LORENTZSCALAR &its, gra:PARAM_RES &resonance) const {
679: // Proton / Dissociative system vertices
680: double FF_A = 0.0;
681: double FF_B = 0.0;
682: PomProtonVertex(its, FF_A, FF_B);
683:
684: // s-channel
685: const std::complex<double> A_sprod =
686: 2.0 * PropOnly(its.s1, its.t1) * FF_A * CBW(its, resonance) *
687: PARAM_REGGE::ResonanceFormFactor(its.m2, pow2(resonance.p.mass), resonance.g_FF) *
688: PropOnly(its.s2, its.t2) * FF_B;
689:
690: // Production and decay amplitude
691: const std::complex<double> A_spin =
692: spin:IProdAmp(its, resonance) * spin:IDecayAmp(its, resonance);
693:
694: // Flux
695: const double V = std::pow(1.0 / its.m2, PARAM_REGGE::iomega);
696:
697: // Full amplitude
698: const std::complex<double> A = A_sprod * A_spin * V;
699:
700: // -----
701: // For screening loop
702: its.hamp = (A);
703: // -----
704:
705: return A;
706: }
707:
708: // =====
709: // Simple matrix element ansatz for Odderon-Pomeron resonances
710: // -----
711: // -----
712: // O
713: // |
714: // |
715: // P
716: // |
717: // |
718: // P
719: // -----
720: // =====
721: std::complex<double> MRegge::ME3ODD(gra:LORENTZSCALAR &its, gra:PARAM_RES &resonance) const {
722: // Proton / Dissociative system vertices
723: double FF_A = its.excite1 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
724: : PARAM_SOFT::gM_P * gra:form:ISF(its.t1);
725: FF_B = its.excite2 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
726: : PARAM_SOFT::gM_P * gra:form:ISF(its.t2);
727:
728: const std::complex<double> A1 =
729: PropOnly(its.s1, its.t1) * FF_A * CBW(its, resonance) +
730: PARAM_REGGE::ResonanceFormFactor(its.m2, pow2(resonance.p.mass), resonance.g_FF) *
731: OdderonProp(its.s2, its.t2) * FF_B;
732:
733: // -----
734: // Proton / Dissociative system vertices
735: FF_A = its.excite1 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
736: : PARAM_SOFT::gM_P * gra:form:ISF(its.t1, its.pfinal1[1].M2());
737: FF_B = its.excite2 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
738: : PARAM_SOFT::gM_P * gra:form:ISF(its.t2, its.pfinal1[2].M2());
739: FF_C = its.excite2 ? msqr(PARAM_SOFT::g3P * PARAM_SOFT::gM_P) *
740: : PARAM_SOFT::gM_P * gra:form:ISF(its.t2, its.pfinal1[2].M2());
741:
742: const std::complex<double> A2 =
743: OdderonProp(its.s1, its.t1) * FF_A * CBW(its, resonance) +
744: PARAM_REGGE::ResonanceFormFactor(its.m2, pow2(resonance.p.mass), resonance.g_FF) *
745: PropOnly(its.s2, its.t2) * FF_B;
746:

```



```
./src/MRegge.cc 10/15
748: // -----
749:
750: // Production and Decay amplitude
751: const std::complex<double> A_spin =
752:   spin:ProdAmp(its, resonance) * spin:DecayAmp(its, resonance);
753:
754: // Flux
755: const double V = std::pow(1.0 / its.m2, PARAM_REGGE::omega);
756:
757: // Should sum here with negative sign if proton-antiproton initial state (anti-symmetric)
758: const std::complex<double> A_pprod = (its.beam_pdg == its.beam2_pdg ? (A1 + A2) : (A1 - A2));
759: const std::complex<double> A = A_pprod * A_spin * V;
760:
761: // -----
762: // For screening loop
763: its.hamp = (A);
764: // -----
765:
766: return A;
767: }
768:
769:
770: // -----
771: // Simple matrix element ansatz for Photoproduction of resonances
772: // -----
773: // -----
774: // -----
775: // -----
776: // -----
777: // -----
778: // -----
779: // -----
780: // -----
781: std::complex<double> MRegge::PhotM2(G3:1:LORENTZSCALAR &its, G3:PARAM_RES &resonance) const {
782: // Check spin
783: if (resonance.p.spinX2 != 2) {
784:   throw std::invalid_argument("MRegge::PhotM2(Resonance.p.spinX2 = " +
785:     resonance.p.spinX2 + " should be J = 1)!");
786: }
787: // Resonance part
788: const std::complex<double> common =
789:   CBW(its, resonance) *
790:   PARAM_REGGE::ResonanceFormFactor(its.m2, pow2(resonance.p.mass), resonance.g_FF);
791:
792: double gammaFlux1 = its.excite1 ? IncohFlux(its.x1, its.t1, its.q1, its.pfinal[1].M2(0))
793:   : CohFlux(its.x1, its.t1, its.q1);
794: double gammaFlux2 = its.excite2 ? IncohFlux(its.x2, its.t2, its.q2, its.pfinal[2].M2(0))
795:   : CohFlux(its.x2, its.t2, its.q2);
796:
797: // "To amplitude level"
798: gammaFlux1 = msqrt(gammaFlux1 / its.x1);
799: gammaFlux2 = msqrt(gammaFlux2 / its.x2);
800:
801: // Pomeron up (t1) & Photon down (t2)
802: const std::complex<double> A_1 =
803:   gammaFlux2 * common *
804:   PhotProp(its.t1, its.t1, pow2(resonance.p.mass), its.excite1, its.pfinal[1].M2(1));
805: // Photon up (t1) & Pomeron down (t2)
806: const std::complex<double> A_2 =
807:   gammaFlux1 * common *
808:   PhotProp(its.t2, its.t2, pow2(resonance.p.mass), its.excite2, its.pfinal[2].M2(1));
809:
810: // Should sum here with negative sign if proton-antiproton initial state (anti-symmetric)
811: const std::complex<double> A_pprod = (its.beam_pdg == its.beam2_pdg ? (A_1 + A_2) : (A_1 - A_2));
812:
813: // Production and Decay amplitude
814: const std::complex<double> A_spin =
815:   spin:ProdAmp(its, resonance) * spin:DecayAmp(its, resonance);
816:
817: // Full amplitude
818: const std::complex<double> A = A_pprod * A_spin;
819: // -----
820: // For screening loop
821: its.hamp = (A);
822: // -----
823: // -----
824: // -----
825: // -----
826: // -----
827: return A;
828: }
829:
830:
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./src/MH2.cc          3/6
167:
168: if (!FILLBUFF) { // Normal filling
169:     +fills;
170: }
171:
172: // Find out bins
173: const int xbin = GetIdx(xvalue, XMN, XMX, XBINS, LOGX);
174: const int ybin = GetIdx(yvalue, YMN, YMX, YBINS, LOGY);
175:
176: if (xbin == -3 || ybin == -3) { nanflow += 1; }
177:
178: if (xbin == -1) { underflow[0] += 1; }
179: if (ybin == -1) { underflow[1] += 1; }
180:
181: if (xbin == -2) { overflow[0] += 1; }
182: if (ybin == -2) { overflow[1] += 1; }
183:
184: if (ValidBin(xbin, ybin)) {
185:     weights[xbin|ybin] += weight;
186:     weights2[xbin|ybin] += weight * weight;
187:     counts[xbin|ybin] += 1;
188: }
189: } else { // Autorange initialization
190:     buff_values.push_back((xvalue, yvalue));
191:     buff_weights.push_back(weight);
192: }
193:
194: if (buff_values.size() > static_cast<unsigned int>(AUTOBUFFSIZE)) { FlushBuffer(); }
195: }
196: }
197:
198: void MH2::Clear() {
199:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
200:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) {
201:             weights[i|j] = 0;
202:             weights2[i|j] = 0;
203:             counts[i|j] = 0;
204:         }
205:     }
206:     fills = 0;
207:     underflow = {0, 0};
208:     overflow = {0, 0};
209:     nanflow = 0;
210: }
211:
212: // Automatic histogram range algorithm
213: void MH2::FlushBuffer() {
214:     if (FILLBUFF && buff_values.size() > 0) {
215:         FILLBUFF = false; // No more filling buffer
216:     }
217:     std::vector<double> min = {0.0, 0.0};
218:     std::vector<double> max = {0.0, 0.0};
219:     // Loop over dimensions
220:     for (std::size_t dim = 0; dim < 2; ++dim) {
221:         // Find out mean
222:         double sum = 0;
223:         double sumw = 0;
224:         for (std::size_t i = 0; i < buff_values.size(); ++i) {
225:             mu += buff_values[i][dim] * buff_weights[i];
226:             sumw += buff_weights[i];
227:         }
228:         if (sumw > 0) { mu /= sumw; }
229:     }
230:     // Variance
231:     double var = 0;
232:     for (std::size_t i = 0; i < buff_values.size(); ++i) {
233:         var += buff_weights[i] * std::pow(buff_values[i][dim] - mu, 2);
234:     }
235:     if (sumw > 0) { var /= sumw; }
236: }
237:
238: // Find minimum and maximum
239: double minval = 1e64;
240: double maxval = -1e64;
241: for (std::size_t i = 0; i < buff_values.size(); ++i) {
242:     if (buff_values[i][dim] < minval) { minval = buff_values[i][dim]; }
243:     if (buff_values[i][dim] > maxval) { maxval = buff_values[i][dim]; }
244: }
245: }
246:
247: // Set new histogram bounds
248: double std = std::sqrt(std::abs(var));
249: double xmin = mu - 2.5 * std;

```

```

./src/MH2.cc          4/6
250: double xmax = mu + 2.5 * std;
251:
252: // A numerical failure may happen with variance calculation, then use this
253: if (std::isnan(xmin) || std::isnan(xmax)) {
254:     xmin = minval;
255:     xmax = maxval;
256: }
257:
258: // If symmetric setup set by user
259: if (AUTOSYMMETRY[dim]) {
260:     double val = std::abs(xmin) + std::abs(xmax) / 2.0;
261:     xmin = -val;
262:     xmax = val;
263: }
264:
265: // We have only positive values, such as invariant mass
266: if (minval < 0.0) { xmin = std::max(0.0, xmin); }
267:
268: min[dim] = xmin;
269: max[dim] = xmax;
270: }
271:
272: // New histogram bounds
273: ResetBounds(YBINS, min[0], max[0], YBINS, min[1], max[1]);
274:
275: // Fill buffered events
276: for (std::size_t i = 0; i < buff_values.size(); ++i) {
277:     Fill(buff_values[i][1], buff_weights[i]);
278: }
279:
280: // Clear buffers
281: buff_values.clear();
282: buff_weights.clear();
283: }
284: }
285:
286: // Histogram mean (based on binned values)
287: double MH2::GetMeanX(int power) const {
288:     double sum = 0.0;
289:     double norm = 0.0; // Normalization
290:     const double binwidth = (XMAX - XMN) / XBINS;
291:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
292:         const double value = std::pow(binwidth * (i + 1) - binwidth / 2.0 + XMN, power);
293:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) {
294:             const double weight = GetBinWeight(i, j);
295:             sum += weight * value;
296:             norm += weight;
297:         }
298:     }
299:     if (norm > 0) {
300:         return sum / norm;
301:     } else {
302:         return 0.0;
303:     }
304: }
305:
306: double MH2::GetMeanY(int power) const {
307:     double sum = 0.0;
308:     double norm = 0.0; // Normalization
309:     const double binwidth = (YMAX - YMN) / YBINS;
310:     for (std::size_t i = 0; i < static_cast<unsigned int>(YBINS); ++i) {
311:         const double value = std::pow(binwidth * (i + 1) - binwidth / 2.0 + YMN, power);
312:         for (std::size_t j = 0; j < static_cast<unsigned int>(XBINS); ++j) {
313:             const double weight = GetBinWeight(i, j);
314:             sum += weight * value;
315:             norm += weight;
316:         }
317:     }
318:     if (norm > 0) {
319:         return sum / norm;
320:     } else {
321:         return 0.0;
322:     }
323: }
324:
325: double MH2::SumWeights() const {
326:     double sum = 0.0;
327:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
328:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) { sum += weights[i|j]; }
329:     }
330:     return sum;
331: }
332:

```

```

./src/MH2.cc          5/6
333: double MH2::SumWeights2() const {
334:     double sum = 0.0;
335:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
336:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) { sum += weights2[i|j]; }
337:     }
338:     return sum;
339: }
340:
341: long long int MH2::SumBinCounts() const {
342:     long long int sum = 0;
343:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
344:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) { sum += counts[i|j]; }
345:     }
346:     return sum;
347: }
348:
349: double MH2::GetMaxWeight() const {
350:     double maxval = 0;
351:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
352:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) {
353:             if (weights[i|j] > maxval) { maxval = weights[i|j]; }
354:         }
355:     }
356:     return maxval;
357: }
358:
359: double MH2::GetMinWeight() const {
360:     double minval = 1e128;
361:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
362:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) {
363:             if (weights[i|j] < minval) { minval = weights[i|j]; }
364:         }
365:     }
366:     return minval;
367: }
368:
369: long long int MH2::GetBinCount(int xbin, int ybin) const {
370:     if (ValidBin(xbin, ybin)) {
371:         return counts[xbin|ybin];
372:     } else {
373:         return 0;
374:     }
375: }
376:
377: // Get weight of the bin
378: double MH2::GetBinWeight(int xbin, int ybin) const {
379:     if (ValidBin(xbin, ybin)) {
380:         return weights[xbin|ybin];
381:     } else {
382:         return 0;
383:     }
384: }
385:
386: // Get bin indices (i, j) corresponding to value (xvalue, yvalue)
387: void MH2::GetBin(double xvalue, double yvalue, int ixbin, int iybin) const {
388:     // Find out bins
389:     xbin = GetIdx(xvalue, XMN, XMX, XBINS, LOGX);
390:     ybin = GetIdx(yvalue, YMN, YMX, YBINS, LOGY);
391: }
392:
393: // Get table/histogram index for linearly or base-10 logarithmically spaced
394: // bins
395: // Gives exact uniform filling within bin boundaries.
396: //
397: // In the logarithmic case, MINVAL and MAXVAL > 0, naturally.
398: //
399: //
400: // Underflow returns -1
401: // Overflow returns -2
402: // nan/inf returns -3
403: //
404: int MH2::GetIdx(double value, double minval, double maxval, int nbins, bool logbins) const {
405:     if (std::isnan(value) || std::isinf(value)) { return -3; }
406:     if (value < minval) { return -1; } // Underflow
407:     if (value > maxval) { return -2; } // Overflow
408:     int idx = 0;
409:     // Logarithmic binning
410:     if (logbins) {
411:         // Check do we have a non-negative input
412:         idx = value > 0 ? std::floor(nbins * (std::log10(value) - std::log10(minval))) /
413:             : -1;

```

```

./src/MH2.cc          6/6
414:
415: // Linear binning
416: } else {
417:     const double BINWIDTH = (maxval - minval) / nbins;
418:     idx = std::floor((value - minval) / BINWIDTH);
419:     return idx;
420: }
421:
422: // Return Shannon entropy of the histogram in bits (base-2 log)
423: // (useful for validating statistical properties, for example)
424: double MH2::ShannonEntropy() const {
425:     double sum = 0.0;
426:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
427:         for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) { sum += weights[i|j]; }
428:     }
429:     if (sum > 0) {
430:         double S = 0;
431:         for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
432:             for (std::size_t j = 0; j < static_cast<unsigned int>(YBINS); ++j) {
433:                 if (weights[i|j] > 0) { S += weights[i|j] / sum * std::log2(weights[i|j] / sum); }
434:             }
435:         }
436:         return -S;
437:     } else {
438:         return 0.0;
439:     }
440: }
441:
442: // namespace gra
443:
444:
445:

```



```

333: // Build kinematics of 2->N
334: bool MContinuum::BNBuildKin(const std::vector<double>& kfs, const std::vector<double>& spins,
335:                             const std::vector<double>& kfs, double m1, double m2) {
336:     const unsigned int KF = Nc - 2; // Central system multiplicity
337:     static const M4Vec beamsum = M4Vec(0, 0, 0, 0);
338:     if (lts.decaortree.size() != KF) {
339:         std::string str =
340:             "MContinuum::BNBuildKin: Decaytree first level topology "
341:             " = "
342:             "std::to_string(lts.decaortree.size()) + " (should be " + std::to_string(KF) +
343:             " for this process!)";
344:         throw std::runtime_error(str);
345:     }
346:     // Forward protons pkpy
347:     M4Vec p1(pkf1 * std::cos(phi1), pkf1 * std::sin(phi1), 0, 0);
348:     M4Vec p2(pkf2 * std::cos(phi2), pkf2 * std::sin(phi2), 0, 0);
349:     // Auxiliary "difference momentum" q0 = p0 - p1 ...
350:     pkt.resize(Kf - 1);
351:     for (const auto k1 : indices(pkt_1)) {
352:         pkt_1[k1] = M4Vec(kf1 * std::cos(phi1[k1]), kf1 * std::sin(phi1[k1]), 0, 0);
353:     }
354:     // Apply linear system to get p
355:     std::vector<M4Vec> p(kf, M4Vec(0, 0, 0, 0));
356:     BLinearSystem(p, pkt_1, p1, p2);
357:     // Set pz and E for central final states
358:     M4Vec sumP(0, 0, 0, 0);
359:     for (const auto k1 : indices(p)) {
360:         const double m = lts.decaortree[k1].m_offshell; // Note offshell!
361:         const double mt = msqrt(pow2(m) + p1[p1].Pz());
362:         p1[1].SetPz(mt * std::sin(y[1]), mt * std::cos(y[1]));
363:         sumP += p[1];
364:     }
365:     // Check cruce energy overflow
366:     if (sumP.E() > lts.sqrts) { return false; }
367:     double p1z = gkfs.kinematics::solvePz(m1, m2, p1, p2, sumP.Pz(), sumP.E(), lts.s);
368:     double p2z = -sumP.Pz() + p1z; // By momentum conservation
369:     // Enforce scattering direction sp -> +p, -p -> -p (VERY RARE POLYNOMIAL
370:     // BRANCH FIZ)
371:     if (p1z < 0 || p2z > 0) { return false; }
372:     // ps and E of protons**
373:     p1.SetPz(p1z, msqrt(pow2(m1) + pow2(p1z)));
374:     p2.SetPz(p2z, msqrt(pow2(m2) + pow2(p2z)));
375:     // How boost if asymmetric beam
376:     if (std::abs(beamsum.Pz()) != 0) {
377:         constexpr int sign = 1; // positive -> boost to the lab
378:         kinematics::LorentzBoost(beamsum, lts.sqrts, p1, sign);
379:         kinematics::LorentzBoost(beamsum, lts.sqrts, p2, sign);
380:     }
381:     for (const auto k1 : indices(p)) { kinematics::LorentzBoost(beamsum, lts.sqrts, p[1], sign); }
382:     // First branch kinematics
383:     lts.pfinal[1] = p1; // Forward systems
384:     lts.pfinal[2] = p2;
385:     lts.pfinal[0].sumP = sumP; // Central system
386:     double sumE = 0;
387:     const unsigned int offset = 3;
388:     for (const auto k1 : indices(p)) {
389:         lts.decaortree[k1].p4 = p[k1];
390:         lts.pfinal[k1 + offset] = p[k1];
391:         sumE += p1.M();
392:     }
393:     // Kinematic checks
394:     // Check we are above mass threshold
395:     if (sumP.M() < sumE) { return false; }

```

```

499:     (5.0 / 6.0, 1.0 / 1.0, -2.0 / 3.0, -1.0 / 2.0, -1.0 / 3.0, -1.0 / 6.0),
500:     (1.0 / 3.0, 1.0 / 2.0, 1.0 / 3.0, -1.0 / 2.0, -1.0 / 3.0, -1.0 / 6.0),
501:     (-1.0 / 6.0, 1.0 / 3.0, 1.0 / 3.0, 1.0 / 2.0, -1.0 / 3.0, -1.0 / 6.0),
502:     (-2.0 / 3.0, -1.0 / 2.0, 1.0 / 3.0, 1.0 / 2.0, 2.0 / 3.0, -1.0 / 6.0),
503:     (7.0 / 6.0, -1.0 / 1.0, 1.0 / 3.0, 1.0 / 2.0, 2.0 / 3.0, 5.0 / 6.0),
504:     (11.0 / 7.0, 5.0 / 7.0, -5.0 / 7.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
505:     (11.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
506:     (4.0 / 7.0, 5.0 / 7.0, 2.0 / 7.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
507:     (1.0 / 14.0, 3.0 / 14.0, 2.0 / 7.0, 3.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
508:     (-3.0 / 7.0, -2.0 / 7.0, 2.0 / 7.0, 3.0 / 7.0, 4.0 / 7.0, 2.0 / 7.0, -1.0 / 7.0),
509:     (-13.0 / 14.0, -11.0 / 14.0, 2.0 / 7.0, 3.0 / 7.0, 4.0 / 7.0, 5.0 / 7.0, -1.0 / 7.0),
510:     (-10.0 / 7.0, -9.0 / 7.0, 2.0 / 7.0, 3.0 / 7.0, 4.0 / 7.0, 5.0 / 7.0, 6.0 / 7.0),
511:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
512:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
513:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
514:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
515:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
516:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
517:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
518:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
519:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
520:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
521:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
522:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
523:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
524:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
525:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
526:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
527:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
528:     (12.0 / 14.0, 17.0 / 14.0, -17.0 / 14.0, -4.0 / 7.0, -3.0 / 7.0, -2.0 / 7.0, -1.0 / 7.0),
529:     // Construct vector b
530:     const unsigned int KF = p.size(); // Number of central system particles
531:     std::vector<M4Vec> b(KF);
532:     const M4Vec p1p2sum = p1f + p2f;
533:     for (const auto k1 : indices(b)) {
534:         if (k1 == 0) {
535:             b[k1] = q[0] - p1p2sum;
536:         } else {
537:             b[k1] = q[k1 - 1] * (-1.0) - p1p2sum;
538:         }
539:     }
540:     // Apply linear system p = Ab to get pkpy components for each p1
541:     const unsigned int index = KF - 2; // -2 because of C++
542:     for (const auto k1 : indices(p)) {
543:         for (const auto k2 : indices(B)) {
544:             p[k1] = b[k2] * A[index][k1][k2]; // notice plus
545:         }
546:     }
547:     // -----
548:     // -----
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580:     // -----
581:     // -----

```

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416: // Total 4-momentum conservation
417: if (gkfs.math::CheckMC(beamsum - (lts.pfinal[1] + lts.pfinal[2] + lts.pfinal[0])) {
418:     return false;
419: }
420: // -----
421: // -----
422: // -----
423: // -----
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```

```
./src/MProcess.cc 1/23
1: // Abstract process class
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <algorithm>
8: #include <cmath>
9: #include <complex>
10: #include <future>
11: #include <iostream>
12: #include <iterator>
13: #include <map>
14: #include <random>
15: #include <regex>
16: #include <string>
17: #include <vector>
18:
19: // Own
20: #include "GravitLL/MBox.h"
21: #include "GravitLL/MFactorised.h"
22: #include "GravitLL/MForm.h"
23: #include "GravitLL/MFragments.h"
24: #include "GravitLL/MGlobin.h"
25: #include "GravitLL/MH1.h"
26: #include "GravitLL/MH2.h"
27: #include "GravitLL/MH3Mematics.h"
28: #include "GravitLL/MH4Mematics.h"
29: #include "GravitLL/MH5.h"
30: #include "GravitLL/MH6.h"
31: #include "GravitLL/MProcess.h"
32: #include "GravitLL/MSetup.h"
33: #include "GravitLL/MTimer.h"
34: #include "GravitLL/MUserOutputs.h"
35:
36: // RegMC3
37: #include "RegMC3/ForceVector.h"
38: #include "RegMC3/GenEvent.h"
39: #include "RegMC3/GenParticle.h"
40: #include "RegMC3/GenVertex.h"
41: #include "RegMC3/Print.h"
42:
43: // Libraries
44: #include "json.hpp"
45: #include "rang.hpp"
46:
47: using gra::aux::indices;
48: using gra::math::abs2;
49: using gra::math::msqrt;
50: using gra::math::pow2;
51: using gra::math::sqrt;
52:
53: namespace gra {
54:
55: void MProcess::PrintSetup() const {
56:     std::cout << std::endl;
57:     std::cout << rang::style::bold << "Process setup:" << rang::style::reset << std::endl;
58:     << std::endl;
59:     std::cout << "Random seed: " << random.GetSeed() << std::endl;
60:     std::cout << "Initial state: " << its.beam1.name << " << its.beam2.name << std::endl;
61:     print(" Beam energies: 10.0 if GeV vs, its.pbeam1.E(), its.pbeam2.E());
62:     print(" CMS energy: 20.0 if GeV vs, its.sqrt_s());
63:     std::cout << "M-Process: " << PROCESS << rang::fg::green << " [" <<
64:
65:     for (const auto &str : ProcPtr.GetProcessDescriptor(PROCESS)) { std::cout << str << " | " <<
66:     << rang::fg::reset << std::endl << std::endl;
67:
68:     // Subprocess
69:     std::cout << rang::style::bold << "Subprocess parameters:" << rang::style::reset << std::endl;
70:     << std::endl;
71:     std::cout << "Pomeron loop screening: " << std::boolalpha << SCREENING << std::endl;
72:
73:     // All other than inclusive processes
74:     if (ProcPtr.ISTATE != "K") {
75:         std::cout << "Final state: " << DECAVMODE << std::endl;
76:         std::cout << "Froton M* excitation: " << std::boolalpha << EXCITATION << std::endl;
77:
78:         if (EXCITATION == 0) {
79:             std::cout << rang::fg::green << " <elastic>" << rang::fg::reset << std::endl;
80:         }
81:         if (EXCITATION == 1) {
82:             std::cout << rang::fg::green << " <single>" << rang::fg::reset << std::endl;
83:         }
84:     }
85: }
```

```
./src/MProcess.cc 3/23
167:
168: if (its.hamp.size() == 0)
169:     const std::istring str =
170:     "MProcess::S3ScreenedAmplitude does not support screening, its.hamp.size() = 0";
171:     throw std::invalid_argument(str);
172: }
173:
174: // Save beam amplitudes
175: const std::vector<std::complex<double>> hamp_0 = its.hamp;
176:
177: // Proton pt vectors
178: const std::vector<double> pT = {its.pfinal[1].Pz(), its.pfinal[1].Py()};
179: const std::vector<double> pT = {its.pfinal[2].Pz(), its.pfinal[2].Py()};
180:
181: // Save old kinematics
182: its.pfinal_orig = its.pfinal;
183:
184: // ## C++11 handles multithreaded static initialization ##
185: // Discretization steps
186: const static double StepE =
187:     (Eikonal.Numerics.MaxLoopRT - Eikonal.Numerics.MinLoopRT) / Eikonal.Numerics.NumberLoopRT;
188: const static double MinPhi = 0.0;
189: const static double MaxPhi = 2.0 * gra::math::PI;
190: const static double StepPhi = (MaxPhi - MinPhi) / Eikonal.Numerics.NumberLoopPHI;
191:
192: // Init 2D-Simpson weight matrix (will be calculated only once, being
193: // static)
194: const static MMatrix<double> W Simpson =
195:     gra::math::Simpson3Weight2D(Eikonal.Numerics.NumberLoopPHI, Eikonal.Numerics.NumberLoopRT);
196:
197: // NOTE M + 1, init with zero!
198: std::vector<MMatrix<std::complex<double>>> f_hamp(
199:     its.hamp.size(), MMatrix<std::complex<double>>(Eikonal.Numerics.NumberLoopPHI + 1,
200:     Eikonal.Numerics.NumberLoopRT + 1, 0.0));
201:
202:
203: // 2D-Integral
204:
205: // Int d^2k A_sik(k;2) A(kt1,kt2)
206: // = Int d^2k A_sik(k;2) A(kt1,kt2) * A(A(kt1,kt2) * A(kt1,kt2)
207:
208: // +3
209: // +1
210: // +1
211: // +1
212: // +1
213: // +1
214: // +1
215: // +1
216:
217: for (std::size_t i = 0; i < Eikonal.Numerics.NumberLoopPHI + 1; ++i) {
218:     const double phi = MinPhi + i * StepPhi;
219:
220:     for (std::size_t j = 0; j < Eikonal.Numerics.NumberLoopRT + 1; ++j) {
221:         const double kt = Eikonal.Numerics.MinLoopRT + j * StepRT;
222:         const double kt2 = gra::math::pow2(kt);
223:
224:         // -----
225:
226:         // Get screening amplitude
227:         // (compiler will optimize this outside the loop)
228:         const std::complex<double> A_sik = Eikonal.MMA.InterpolateID(kt2);
229:
230:         // 0. Construct loop 2D kt-vector
231:         const std::vector<double> kt_kt = std::cos(phi), kt * std::sin(phi);
232:
233:         // 1. New proton pt vectors
234:         const std::vector<double> p1p = {pT[0] - kt_kt[0], pT[1] - kt_kt[1]};
235:         const std::vector<double> p2p = {pT[0] + kt_kt[0], pT[1] + kt_kt[1]};
236:
237:         // 2. Update kinematics
238:         if ((LoopKinematics(p1p, p2p)) <
239:             continue; // not valid kinematically
240:         }
241:
242:         // 3. Get new amplitudes to its.hamp
243:         ProcPtr.GetBeamAmplitude2(its);
244:         // -----
245:         // Loop over all helicity amplitudes
246:         for (const auto &h : indices(its.hamp)) {
247:             f_hamp[h][i][j] = A_sik * its.hamp[h] * kt; // note x kt (Jacobian)
248:         }
249:     }
250: }
```

```
./src/MProcess.cc 2/23
84: if (EXCITATION == 2) {
85:     std::cout << rang::fg::green << " <double>" << rang::fg::reset << std::endl;
86: }
87:
88: if (FLATAMP != 0) { std::cout << " - Flat amplitude mode: " << FLATAMP << std::endl; }
89:
90: std::cout << std::endl << std::endl;
91:
92:
93:
94: // Cascaded phase-space factor [VOLUME] x [MC INTEGRAL] / (2PI)
95: double MProcess::CascadedPS() const {
96:     // Cascade resonances phase-space
97:     double product = 1.0;
98:     double product2pi = 1.0;
99:     double volume = 1.0;
100:     int N_final = 0;
101:     for (const auto &i : indices(its.decautree)) {
102:         CalculatePhaseSpace(its.decautree[i], product, product2pi, volume, N_final);
103:     }
104:     return volume * product / product2pi;
105: }
106:
107: // Recursive function to calculate phase-space weights
108: void MProcess::CalculatePhaseSpace(const gra::MDecayBranch &branch, double &product,
109:     double &product2pi, double &volume, int &N_final) const {
110:     // There is decay information and this is activated by the amplitude
111:     if (branch.PS_active == true & branch.W.Integral() > 0) {
112:         product *= branch.W.Integral();
113:         product2pi *= (2 * math::PI);
114:
115:         // Flat mass^2 interval
116:         double M_sum = 0;
117:         for (const auto &i : indices(branch.legs)) {
118:             M_sum += branch.legs[i].p.mass;
119:         } // Daughters give lower limit
120:
121:         const double MIN = std::max(pow2(M_sum), pow2(branch.p.mass - branch.p.width * OFFSHELL));
122:         const double MAX = pow2(branch.p.mass + branch.p.width * OFFSHELL);
123:         volume *= (MAX - MIN);
124:
125:         for (const auto &i : indices(branch.legs)) {
126:             CalculatePhaseSpace(branch.legs[i], product, product2pi, volume, N_final);
127:         }
128:     }
129:     if (branch.legs.size() == 0) { ++N_final; } // Final state particle
130: }
131:
132: // Amplitude squared
133: double MProcess::GetAmp2() {
134:     double amp2 = (FLATAMP == 0) ? S3ScreenedAmp2() : GetFlatAmp2(its);
135: }
136:
137: // -----
138: // ** Custom sampling control initiated by amplitudes **
139: // if (its.FORCE_FLATMASS2 && !FLATMASS2_user) { FLATMASS2 = true; }
140: // if (its.FORCE_OFFSHELL == 0 && !OFFSHELL_user) { OFFSHELL = its.FORCE_OFFSHELL; }
141:
142: return amp2;
143: }
144:
145: // Pomeron Loop Screened Amplitude
146:
147: // =====
148: // *
149: // * -pT + kt = qT1
150: // *
151: // * +x=
152: // *
153: // * -pT - kt = qT2
154: // *
155: // =====
156:
157: double MProcess::S3ScreenedAmp2() {
158:     // Eikonal Loop Screening not on
159:     if (SCREENING == false) { return ProcPtr.GetBeamAmplitude2(its); }
160:
161:     // Elastic scattering, return directly eikonalized amplitude squared itself
162:     if (ProcPtr.CHANNEL == "EL") { return abs2(Eikonal.MMA.InterpolateID(-its.t)); }
163:
164:     // -----
165:     // First evaluate bare amplitudes to its.hamp
166:     ProcPtr.GetBeamAmplitude2(its);
167:
168:     // inner-loop
169:     // outer-loop
170:
171:     // Normalization
172:     const std::complex<double> norm = zi / (8.0 * gra::math::PIPI * its.s);
173:
174:     // Loop over amplitudes
175:     std::vector<std::complex<double>> hamp_loop(its.hamp.size(), 0.0);
176:
177:     for (const auto &h : indices(its.hamp)) {
178:         hamp_loop[h] = norm * gra::math::Simpson3Integral2D(f_hamp[h], W Simpson, StepPhi, StepRT);
179:     }
180:
181:     // Update back to tree level kinematics
182:     LoopKinematics(pT, pT);
183:
184:     // -----
185:     // Final amplitude (squared)
186:
187:     // Not Durban-QCD
188:     if (ProcPtr.ISTATE != "qq") {
189:         // Separate (non)resonant sum
190:         double amp2 = 0.0;
191:         for (const auto &h : indices(its.hamp)) { amp2 += abs2(hamp_0[h] + hamp_loop[h]); }
192:
193:         // Initial state helicity average, if we have all helicity amplitudes
194:         if (its.hamp.size() != 1) { amp2 /= 4; }
195:         return amp2;
196:     } else {
197:         // Coherent sum
198:         std::complex<double> A = 0.0;
199:         for (const auto &h : indices(its.hamp)) { A += hamp_0[h] + hamp_loop[h]; }
200:         // Helicity average already taken care of
201:         return abs2(A);
202:     }
203: }
204:
205: // Set CMS energy and beam particle 4-vectors
206: void MProcess::SetInitialState(const std::vector<istring> &beam,
207:     const std::vector<double> &energy) {
208:     if (beam.size() != 2) {
209:         throw std::invalid_argument("MProcess::SetInitialState: Input BEAM vector not dim 2!");
210:     }
211:     if (energy.size() != 2) {
212:         throw std::invalid_argument("MProcess::SetInitialState: Input ENERGY vector not dim 2!");
213:     }
214: }
215:
216: // Beam needs to be set before this!
217: void MProcess::SetBeamEnergies(double E1, double E2) {
218:     if (its.beam1.pdg == 0 || its.beam2.pdg == 0) {
219:         throw std::invalid_argument("MProcess::SetBeamEnergies: Beam PDG particles not set yet!");
220:     }
221: }
222:
223: // Beam 4-momentum (px,py,pz,E)
218: its.pbeam1 = M4Vec(0, 0, msqrt(pow2(E1) - pow2(its.beam1.mass)), E1); // positive z-axis
219: its.pbeam2 = M4Vec(0, 0, -msqrt(pow2(E2) - pow2(its.beam2.mass)), E2); // negative z-axis
220:
221: // Mandelstam s
222: its.s = (its.pbeam1 + its.pbeam2).M2();
223: its.sqrt_s = msqrt(its.s);
224:
225: if (its.sqrt_s < (its.beam1.mass + its.beam2.mass)) {
226:     std::istring str = "MProcess::SetBeamEnergies: Error with input CMS energy: " +
227:         std::string(its.sqrt_s) + " GeV - initial state masses!";
228:     throw std::invalid_argument(str);
229: }
230:
231: print("MProcess::SetBeamEnergies: beam: [%s, %s], energy = [%0.1f, %0.1f] \n",
232: }
```



```

./src/MProcess.cc          9/23
665: //
666: // In 2-body case, the decay phase space and decay amplitude squared
667: // factorize
668: //
669: bool TP_computed = false;
670: if (daughter_size() == 2) {
671:     // Calculate phase space part:
672:     // Gamma = PS * |A_decay|^2 then with |A_decay|^2 = 1 <<< Gamma = PS
673: }
674: const double amp2 = 1.0;
675: const double sym = 1.0;
676:
677: double PS = gra::kinematics::PFWbody(pow2(p.mass), pow2(daughter[0].mass),
678:                                     pow2(daughter[1].mass), amp2, sym);
679:
680: // -----
681: // Tensor Pomeron 2-body couplings for 2-body decays directly calculable from width and
682: // BR
683: if (hc.g_decay_tensor.size() == 0) { // Not set yet
684:     hc.g_decay_tensor =
685:         MTensorPomeron::IDecay(p.spinX2 / 2, p.mass, p.width, daughter[0].mass, hc.BR);
686:     if (std::fabs(PS) < ZERO_EPS) {
687:         // Try again with higher mother mass as a crude approximation, we might
688:         // be trying purely off-shell decay (such as f(980) -> K+K-)
689:         const unsigned int N_width = 3;
690:         for (std::size_t i = 1; i <= N_width; ++i)
691:             PS = gra::kinematics::PFWbody(pow2(p.mass + i * p.width), pow2(daughter[0].mass),
692:                                         pow2(daughter[1].mass), amp2, sym);
693:     }
694:     // -----
695:     // Tensor Pomeron 2-body couplings
696:     hc.g_decay_tensor = MTensorPomeron::IDecay(p.spinX2 / 2, p.mass + i * p.width, p.width,
697:                                               daughter[0].mass, hc.BR);
698:     if (PS > ZERO_EPS) { break; }
699: }
700: }
701: TP_computed = true;
702: }
703: }
704: }
705:
706: // BR [equiv Gamma/Gamma_tot = PS * |A_decay|^2] / Gamma_tot
707: // <= <= <= BR * Gamma_tot / Gamma_PS and sqrt to get 'amplitude'
708: // [A]^2 = BR * Gamma_tot / Gamma_PS and sqrt to get 'amplitude'
709: // level'
710: //
711: // ** GeV unit can be obtained by inspecting PFWbody function **
712: hc.g_decay = msqrt(hc.BR * p.width / PS);
713:
714: if (std::isnan(hc.g_decay) || std::isnan(hc.g_decay)) {
715:     std::string str =
716:         "MProcess::ProcessHelicityDecay: Kinematic coupling problem "
717:         "to: "
718:         "DECAYMODE + " with resonance PDG = " + p.name + "(" + std::to_string(p.pdg) + ")" +
719:         " (daughter too heavy?) (check RESONANCE, DECAYMODE
720:         "and BRANCHING tables)";
721:     throw std::invalid_argument(str);
722: }
723: }
724: }
725: }
726: // 3, 4, ... body cases [NOT TREATED YET, phase space and matrix
727: // element do not factorize there]
728: } else {
729:     hc.g_decay = 1.0; // For the rest, put 1.0
730: }
731: }
732: printf("Mass: Full width: %10.3E %10.3E GeV\n", p.mass, p.width);
733: printf("Branching ratio || Partial width: %10.3E || %10.3E GeV\n", hc.BR, hc.BR * p.width);
734: printf("=> Computed decay coupling: %10.3E\n", hc.g_decay);
735: // -----
736: // Tensor Pomeron 2-body couplings: 'J'
737: for (const auto &i : indices(hc.g_decay_tensor)) { printf("%0.3E ", hc.g_decay_tensor[i]); }
738: printf("\n", TP_computed ? " (computed from J, width and BR) : "");
739: }
740: }
741: }
742: }
743: }
744: }
745: }
746: aux::PrintBar("");
747: }

```

```

./src/MProcess.cc          11/23
831: const double a0 = 0.0;
832: double b0 = 2.0;
833: double avgM = a0 + b0 * (1 / QProb) * std::log(Q2_scale);
834: int outertrials = 0;
835:
836: const double M = nstar.M();
837:
838: while (true) { // OUTER
839:     // Draw fluctuating number of particles
840:     int N = 0;
841:     while (true) {
842:         N = random.PoissonNumber(avgM);
843:         // If N < 10 ( // low mass sampling treatment
844:         for (std::size_t i = 0; i < N; ++i) {
845:             N = PoissonRnd(N);
846:         }
847:     }
848:     if (nstar.M() > N * 0.14) break; // Minimal mass threshold
849:     }
850: }
851: // Boundary conditions
852: N = (N < 1) ? 1 : N; // At least 1
853: N = (N >= std::abs(B)) ? N : std::abs(B); // Baryon number
854: N = (N >= std::abs(Q)) ? N : std::abs(Q); // Charge conservation
855:
856: // -----
857: // Pick particles
858: std::vector<double> mass;
859: std::vector<int> pdgcode;
860: if (MFRAGMENT::PickParticles(M, N, B, Q, mass, pdgcode, lts.PDG, random)) {
861:     +outertrials;
862:     continue;
863: }
864: }
865: }
866: // -----
867: // Now decay >>
868: std::vector<M4Vec> p4;
869:
870: double g = std::pow(N, 0.11 / 3); // Powerlaw high-pt slope
871: double T = 0.065; // Temperature
872: if (etree.size() > 0) { // For MD soft-core, 'Temperature' rises ~ 1/impact parameter squared
873:     T = std::pow(1 / (bc * bc), 0.10);
874: }
875:
876: const double maxpt = 15.0; // Maximum Pt per particle
877: const double lambda = 5.0; // Get r(=2), for exponential pr^2
878: const double M = nstar.M();
879: const double MFRAGMENT::TubeFragment(nstar, M, mass, p4, g, T, lambda, maxpt, random, pt_distribution);
880:
881: if (N <= 0) {
882:     +outertrials;
883:     continue;
884: }
885: if (outertrials > OUTERTRIAL) {
886:     // std::cout << "MProcess::ExciteContinuum: Outer loop failure!" <<
887:     // std::endl;
888:     return false; // too many failures
889: } else {
890:     continue; // try again!
891: }
892: } else {
893:     // Re-check kinematics
894:     bool fail = false;
895:     for (const auto &i : indices(p4)) {
896:         if (std::isnan(p4[i].Rap()) ||
897:             // std::cout << "range: yellow <<
898:             // MFRAGMENT::ExciteContinuum: Kinematics failure!" <<
899:             // std::endl;
900:             // std::cout << "range: yellow <<
901:             // MFRAGMENT::ExciteContinuum: Kinematics failure!" <<
902:             // std::endl;
903:             // std::cout << "range: yellow <<
904:             // MFRAGMENT::ExciteContinuum: Outer
905:             // loop failure!" << std::endl;
906:             return false; // too many failures
907:         )
908:         fail = true;
909:         break;
910:     }
911:     if (fail) { continue; }
912: }
913: // Get corresponding PDG particles
914: std::vector<gra::MParticle> p(p4.size());
915: for (const auto &i : indices(p)) { p[i] = lts.PDG.FindByPDG(pdgcode[i]); }

```

```

./src/MProcess.cc          10/23
748: return hc;
749: }
750:
751: // Get intermediate off-shell mass
752: void MProcess::GetOffShellMass(const gra::MDecayBranch &branch, double &mass) {
753:     // If zero-width particle (electron, gamma, proton)
754:     if (branch.p.width < 1e-40) {
755:         mass = branch.p.mass;
756:         return;
757:     }
758:     const unsigned int OUTERMAXTRIAL = 1e4;
759:     const unsigned int INNERMAXTRIAL = 1e4;
760:     unsigned int outertrials = 0;
761:     while (true) {
762:         // We have decay daughters
763:         double daughter_masses = 0;
764:         if (branch.legs.size() != 0) {
765:             // Find random daughter offshell masses from BW
766:             for (const auto &i : indices(branch.legs)) {
767:                 const double auto_i1 : indices(branch.legs) {
768:                     0.0,
769:                     random.RelativeisticBRandom(branch.legs[i].p.width, OFFSHELL);
770:                 }
771:                 const double safe_margin = 1e-5;
772:                 const double auto_i2 : indices(branch.legs) {
773:                     0.0,
774:                     random.RelativeisticBRandom(branch.legs[i].p.width, OFFSHELL);
775:                 }
776:                 const double safe_margin = 1e-5;
777:                 const double auto_i3 : indices(branch.legs) {
778:                     0.0,
779:                     random.RelativeisticBRandom(branch.legs[i].p.width, OFFSHELL);
780:                 }
781:                 const double M = branch.p.mass;
782:                 const double W = branch.p.width;
783:                 if (!FLATMASE2) {
784:                     mass = std::max(daughter_masses, random.RelativeisticBRandom(M, W, OFFSHELL, daughter_masses));
785:                 } else {
786:                     mass = msqrt(random.Ustd::max(pow2(daughter_masses), pow2(M - OFFSHELL * W)),
787:                                 std::min((ts, pow2(M + OFFSHELL * W)));
788:                 }
789:                 +innertrials;
790:                 if (mass >= daughter_masses) {
791:                     return; // all done
792:                 }
793:                 if (innertrials > INNERMAXTRIAL) {
794:                     break; // try again with different daughter masses
795:                 }
796:             }
797:             +outertrials;
798:             if (outertrials > OUTERMAXTRIAL) {
799:                 // Impossible
800:                 std::string str =
801:                     "MProcess::GetOffShellMass: Kinematically impossible decay: " + branch.p.name + " + " +
802:                     str;
803:                 throw std::invalid_argument(str);
804:             }
805:         }
806:     }
807: }
808: }
809: }
810: }
811: }
812: }
813: }
814: }
815: // -----
816: // Proton continuum excitation
817: bool MProcess::ExciteContinuum(const M4Vec &nstar, gra::MDecayBranch &forward, double Q2_scale,
818:                               int B, int Q, const std::string &pt_distribution) {
819:     // Sanity check (2 x pion mass)
820:     if (msqrt(Q2_scale) < 0.3) {
821:         // return false; // not valid
822:     }
823:     const double QProb = 2.0 / 3.0; // Probability for a charged particle (isospin)
824:     const int OUTERTRIAL = 100;
825:     const int MULTIPlicity = 10;
826: }

```

```

./src/MProcess.cc          12/23
914: // Create decaytree
915: BranchForwardSystem(p4, p, nstar, forward);
916: break;
917: }
918: }
919: }
920: return true;
921: }
922: }
923: // Proton low mass (N* like) 2-body and 3-body excitation
924: // -----
925: // -----
926: bool MProcess::ExciteNstar(const M4Vec &nstar, gra::MDecayBranch &forward, const MParticle &pbear) {
927:     // Find nstar decaymode
928:     std::vector<int> pdgcodes;
929:     MFRAGMENT::NstarDecayTable(pbear.chargeX3 / 3, nstar.M(), pdgcodes, random);
930:     // Get corresponding PDG particles
931:     std::vector<MParticle> p(pdgcodes.size());
932:     std::vector<double> mass(pdgcodes.size(), 0.0);
933:     for (const auto &i : indices(pdgcodes)) {
934:         p[i] = lts.PDG.FindByPDG(pdgcodes[i]);
935:         mass[i] = p[i].mass;
936:     }
937:     // -----
938:     // Products 3-momenta
939:     std::vector<M4Vec> p4;
940:     // Do the 2 or 3-body isotropic decay
941:     if (mass.size() == 2) { gra::kinematics::TwoBodyPhaseSpace(nstar, nstar.M(), mass, p4, random); }
942:     if (mass.size() == 3) {
943:         const bool UNWEIGHT = true;
944:         gra::kinematics::ThreeBodyPhaseSpace(nstar, nstar.M(), mass, p4, UNWEIGHT, random);
945:     }
946:     // Create decaytree
947:     BranchForwardSystem(p4, p, nstar, forward);
948:     return true;
949: }
950: }
951: // This is used with forward excitation
952: void MProcess::BranchForwardSystem(const std::vector<M4Vec> &p4, const std::vector<MParticle> &p,
953:                                   const M4Vec &nstar, gra::MDecayBranch &forward) {
954:     // -----
955:     // -----
956:     std::vector<bool> isstable(N, false);
957:     MFRAGMENT::GetDecayStatus(pdgcode, isstable);
958:     // -----
959:     // -----
960:     // Construct decaytree
961:     forward = gra::MDecayBranch(); // Initialize!
962:     forward.p4 = nstar;
963:     // Daughters
964:     forward.legs.resize(p.size());
965:     forward.depth = 0;
966:     for (const auto &i : indices(p)) {
967:         // Decay particle
968:         gra::MDecayBranch branch;
969:         branch.p = p[i];
970:         branch.p4 = p4[i];
971:         // Treat pi0 -> gamma BR ~ 100%
972:         if (branch.p.pdg == PDG::PDG_pi0) {
973:             const std::vector<double> md = {0.0, 0.0};
974:             std::vector<double> pd;
975:             gra::kinematics::TwoBodyPhaseSpace(branch.p4.M(), md, pd, random);
976:             // Add gamma legs
977:             branch.legs.resize(2);
978:             for (std::size_t k = 0; k < 2; ++k) {
979:                 gra::MDecayBranch decaybranch;
980:                 decaybranch.p = lts.PDG.FindByPDG(PDG::PDG_gamma);
981:                 decaybranch.p4 = pd[k];
982:                 branch.legs[k] = decaybranch;
983:                 branch.legs[k].depth = 2;
984:             }
985:         }
986:     }
987: }

```



```

./src/MProcess.cc 13/23
997: // Treat rho0 -> pi+ pi- (BR ~ 100%)
998: else if (branch.p.pdg == PDG::PDG_rho0) {
999: // Decay
1000: const std::vector<double> md = {PDG::mpi, PDG::mpipi};
1001: const std::vector<int> pdg = {PDG::PDG_pi+, PDG::PDG_pi-};
1002: std::vector<MVeco> md_vec;
1003: gra::kinematics::TwoBodyPhaseSpace(branch.p4, branch.p4.M(), md, pd, random);
1004:
1005: // Add pion legs
1006: branch.legs.resize(2);
1007: for (std::size_t k = 0; k < 2; ++k) {
1008: gra::MDecayBranch decaybranch;
1009: decaybranch.p = lcs.PDG::PdgMVeco[branch.pdg[k]];
1010: decaybranch.m = md_vec[k];
1011: branch.legs[k].m = decaybranch.m;
1012: branch.legs[k].depth = 2;
1013: }
1014: }
1015: // Here, we could add other MAJOR resonant decays
1016: // ...
1017: // Add leg
1018: forward.legs[i] = branch;
1019: forward.legs[i].depth = 1;
1020: }
1021: }
1022: // Fragment forward systems in central production
1023: bool MProcess::CEPForwardFragment() {
1024: bool ok = true;
1025: const std::string pt = "wpp"; // Soft exponential pt
1026: if (lcs.excited() {
1027: const int Q = math::sign(lcs.beam.pdg); // Charge and baryon number
1028: const int B = Q;
1029: if (PARAM_NSTAR::fragment == "none") {
1030: BranchForwardSystem(1), lcs.pfinal[1], lcs.decayforward);
1031: } else if (PARAM_NSTAR::fragment == "teabody") {
1032: if (!lcs.isStar(lcs.pfinal[1], lcs.decayforward, lcs.beam)) (ok = false);
1033: } else if (PARAM_NSTAR::fragment == "cylinder") {
1034: if (!lcs.isContinuum(lcs.pfinal[1], lcs.decayforward, lcs.pfinal[1].M2(), B, Q, pt)) (ok = false);
1035: } else {
1036: // none
1037: }
1038: }
1039: }
1040: if (lcs.excited2) {
1041: const int Q = math::sign(lcs.beam2.pdg); // Charge and baryon number
1042: const int B = Q;
1043: if (PARAM_NSTAR::fragment == "none") {
1044: BranchForwardSystem(1), lcs.pfinal[2], lcs.decayforward2);
1045: } else if (PARAM_NSTAR::fragment == "teabody") {
1046: if (!lcs.isStar(lcs.pfinal[2], lcs.decayforward2, lcs.beam2)) (ok = false);
1047: } else if (PARAM_NSTAR::fragment == "cylinder") {
1048: if (!lcs.isContinuum(lcs.pfinal[2], lcs.decayforward2, lcs.pfinal[2].M2(), B, Q, pt)) (ok = false);
1049: } else {
1050: // none
1051: }
1052: }
1053: }
1054: return ok;
1055: }
1056: bool MProcess::VetoOuts() const {
1057: bool ok = true;
1058: // Veto cuts
1059: if (vetcuts.active) {
1060: // Forward system
1061: FindVetoCuts(lcs.decayforward, ok);
1062: FindVetoCuts(lcs.decayforward2, ok);
1063: }
1064: // Central system
1065: for (const auto i : indices(lcs.decaytree)) (FindVetoCuts(lcs.decaytree[i], ok);)
1066: }
1067: }
1068: }

```

```

./src/MProcess.cc 13/23
1163: if (branch.p4.Pt() > vetocuts.cuts[1].pt_min && branch.p4.Pt() <= vetocuts.cuts[1].pt_max &&
1164: branch.p4.Eta() >= vetocuts.cuts[1].eta_min &&
1165: branch.p4.Eta() <= vetocuts.cuts[1].eta_max) {
1166: ok = false; // VETO, does not pass
1167: } else {
1168: // does not trigger veto, do nothing
1169: }
1170: }
1171: }
1172: // ** RECURSION here **
1173: for (const auto i : indices(branch.legs)) (FindVetoCuts(branch.legs[i], ok);)
1174: }
1175: }
1176: // Recursive function to plot out the decay tree
1177: void MProcess::PrintDecayTree(const gra::MDecayBranch &branch) const {
1178: std::string spaces(branch.depth * 2, ' '); // Give empty space
1179: std::string lines = spaces + "1A2A/224A/224A/200A/224A/200B";
1180:
1181: if (branch.depth + 1 % 3 == 0) {
1182: std::cout << rang::fgblue << lines << rang::freset;
1183: } else if (branch.depth + 1 % 2 == 0) {
1184: std::cout << rang::fgyellow << lines << rang::freset;
1185: } else {
1186: std::cout << rang::fggreen << lines << rang::freset;
1187: }
1188: }
1189: print(
1190: "M = 10.2E, W = 10.2E (GeV) [ctau = 10.1E s] PDG = [std::to_string] | ka"
1191: " [Qs, Ds] | n",
1192: branch.p.mass, branch.p.width, branch.p.tau, branch.p.pdg > 0 ? " " : "", branch.p.pdg,
1193: branch.p.name_c_str(), gra::aux::Charge3toStr(branch.p.charge3).c_str(),
1194: gra::aux::Spin2toStr(branch.p.spin2).c_str());
1195: // ** RECURSION here **
1196: for (const auto i : indices(branch.legs)) (PrintDecayTree(branch.legs[i]);)
1197: }
1198: }
1199: // Recursive function to plot out the decay tree
1200: void MProcess::PrintPhaseSpace(const gra::MDecayBranch &branch, double &product, double &product2p,
1201: int &final) const {
1202: std::string spaces(branch.depth * 2, ' '); // Give empty space
1203: std::string lines = spaces + "1A2A/200A/224A/200B";
1204:
1205: if (branch.W.Integral() > 0) { // There is decay information
1206: print("%s \t [1->N] LIPS: 10.2E - 10.2E, lines.c_str(), branch.legs.size(),
1207: branch.W.Integral(), branch.W.IntegralError());
1208: }
1209: if (branch.PS_active) {
1210: std::cout << "[ " << rang::fggreen << "ACTIVE" << rang::freset;
1211: << " part of integral / (2PI) " << std::endl;
1212: } else {
1213: std::cout << "[ " << rang::fired << "INACTIVE" << rang::freset;
1214: << " part of integral << apply manual BR " << std::endl;
1215: }
1216: }
1217: }
1218: product = branch.W.Integral();
1219: product2p = (2 * math::PI);
1220: }
1221: for (const auto i : indices(branch.legs)) {
1222: PrintPhaseSpace(branch.legs[i], product, product2p, final);
1223: }
1224: }
1225: }
1226: if (branch.legs.size() == 0) { // Final state particle
1227: }
1228: }
1229: // Recursive decay tree kinematics (called event by event from inheriting
1230: // classes)
1231: bool MProcess::ConstructDecayKinematics(gra::MDecayBranch &branch) {
1232: // This leg has any daughters
1233: if (branch.legs.size() != 0) {
1234: // Generate decay product masses
1235: std::vector<double> m(branch.legs.size(), 0.0);
1236: while (true) {
1237: for (const auto i : indices(branch.legs)) {
1238: GetEffMass(branch.legs[i], branch.legs[i].m_offshell);
1239: m[i] = branch.legs[i].m_offshell;
1240: }
1241: // Check decay products masses are not over mother mass
1242: if (branch.p4.M() < std::accumulate(m.begin(), m.end(), 0.0)) {
1243: continue;
1244: } else {
1245: break;

```

```

./src/MProcess.cc 14/23
1080: }
1081: return ok;
1082: }
1083: // Common cuts for MResonance and MContinuum classes
1084: bool MProcess::CommonCuts() const {
1085: bool ok = true;
1086: // Fiducial cuts
1087: if (cuts.active == true) {
1088: // Check custom user cuts (do not substitute to kinematics_ok = UserCut...)
1089: if (UserCut(USERCuts, lcs)) {
1090: return false; // Not fine
1091: }
1092: // Check forward system variables
1093: if (CID != "P") { // Collinear class does not support these
1094: if (std::abs(lcs.t1) >= fcuts.forward_t_min && std::abs(lcs.t1) <= fcuts.forward_t_max &&
1095: std::abs(lcs.t2) >= fcuts.forward_t_min && std::abs(lcs.t2) <= fcuts.forward_t_max) {
1096: // fine
1097: } else {
1098: return false;
1099: }
1100: }
1101: if (lcs.excited) {
1102: if (lcs.pfinal[1].M() >= fcuts.forward_M_min && lcs.pfinal[1].M() <= fcuts.forward_M_max) {
1103: // fine
1104: } else {
1105: return false;
1106: }
1107: }
1108: if (lcs.excited2) {
1109: if (lcs.pfinal[2].M() >= fcuts.forward_M_min && lcs.pfinal[2].M() <= fcuts.forward_M_max) {
1110: // fine
1111: } else {
1112: return false;
1113: }
1114: }
1115: // Check system variables
1116: if (msqrt(lcs.m2) >= fcuts.M_min && msqrt(lcs.m2) <= fcuts.M_max && lcs.Y >= fcuts.Y_min &&
1117: lcs.Y <= fcuts.Y_max && lcs.Pt >= fcuts.Pt_min && lcs.Pt <= fcuts.Pt_max) {
1118: // fine, do not touch
1119: } else {
1120: return false; // Not fine
1121: }
1122: }
1123: // Check fiducial cuts of the central final state particles
1124: for (const auto i : indices(lcs.decaytree)) (FindDecayCuts(lcs.decaytree[i], ok);)
1125: }
1126: }
1127: // Recursive find out final daughters and check if they pass fiducial cuts
1128: void MProcess::FindDecayCuts(const gra::MDecayBranch &branch, bool &ok) const {
1129: // We are at the end -> must be a final state
1130: if (branch.legs.size() == 0) {
1131: // Check cuts
1132: if (branch.p4.Pt() >= fcuts.pt_min && branch.p4.Pt() <= fcuts.pt_max &&
1133: branch.p4.Eta() >= fcuts.eta_min && branch.p4.Eta() <= fcuts.eta_max &&
1134: branch.p4.Rap() >= fcuts.rap_min && branch.p4.Rap() <= fcuts.rap_max) {
1135: // Event passes cuts
1136: } else {
1137: ok = false; // does not pass
1138: }
1139: }
1140: // ** RECURSION here **
1141: for (const auto i : indices(branch.legs)) (FindDecayCuts(branch.legs[i], ok);)
1142: }
1143: // Recursive find out final daughters and check if they trigger veto cuts
1144: void MProcess::FindVetoCuts(const gra::MDecayBranch &branch, bool &ok) const {
1145: // We are at the end -> must be a final state
1146: if (branch.legs.size() == 0) {
1147: // Loop over veto domains
1148: for (const auto i : indices(vetcuts.cuts)) {
1149: // Check cuts
1150: }
1151: }
1152: }
1153: // For now, keep always UNWEIGHT = true;
1154: const bool UNWEIGHT = true;
1155: std::vector<MVeco> p;
1156: // 2-body
1157: gra::kinematics::MCW w;
1158: if (branch.legs.size() == 2) {
1159: w = gra::kinematics::TwoBodyPhaseSpace(branch.p4, branch.p4.M(), m, p, random);
1160: } // 3-body
1161: } else if (branch.legs.size() == 3) {
1162: w = gra::kinematics::ThreeBodyPhaseSpace(branch.p4, branch.p4.M(), m, p, UNWEIGHT, random);
1163: } else {
1164: w = gra::kinematics::NBodyPhaseSpace(branch.p4, branch.p4.M(), m, p, UNWEIGHT, random);
1165: }
1166: if (w.GetW() < 0) {
1167: std::string str =
1168: "MProcess::ConstructDecayKinematics: Fatal error: Weight < 0 (Check your decay tree)";
1169: std::cout << str << std::endl;
1170: return false;
1171: }
1172: // Collect weight
1173: branch.W += w;
1174: // Collect decay product 4-momenta
1175: for (const auto i : indices(branch.legs)) (branch.legs[i].p = p[i];)
1176: }
1177: // ** Now the IMNES recursion **
1178: for (const auto i : indices(branch.legs)) {
1179: if (!ConstructDecayKinematics(branch.legs[i])) (return false;);
1180: }
1181: }
1182: return true;
1183: }
1184: // Recursively add final states to the event structure
1185: void MProcess::WriteDecayKinematics(const gra::MDecayBranch &branch, const HepMC3::GenParticle &mother, HepMC3::GenVertex &evt) {
1186: // This particle has daughters
1187: if (branch.legs.size() > 0) {
1188: // Create new vertex with decay 4-position
1189: HepMC3::GenVertexPtr vertex =
1190: std::make_shared<HepMC3::GenVertex>(gra::aux::M4VecHepMC3(branch.decay_position));
1191: evt.add_vertex(vertex);
1192: // The decaying particle
1193: vertex->add_particle_in(mother);
1194: }
1195: // Add daughters
1196: for (const auto i : indices(branch.legs)) {
1197: const int STATE = (branch.legs[i].legs.size() > 0 ? PDG::PDG_DECAY : PDG::PDG_STABLE);
1198: // ADD HERE THE ctau > 1.0 cm definition for the status
1199: // code [TBD]
1200: HepMC3::GenParticlePtr particle = std::make_shared<HepMC3::GenParticle>(
1201: gra::aux::M4VecHepMC3(branch.legs[i].p4), branch.legs[i].p.pdg, STATE);
1202: vertex->add_particle_out(particle);
1203: }
1204: // ** RECURSION here **
1205: WriteDecayKinematics(branch.legs[i], particle, evt);
1206: }
1207: // Forward excitation mass sampling
1208: void MProcess::SampleForwardMasses(std::vector<double> &mwec, const std::vector<double> &randvec) {
1209: mwec = {lcs.beam.mass, lcs.beam2.mass};
1210: lcs.excited = false;
1211: lcs.excited2 = false;
1212: if (EXCITATION == 0) {
1213: M2_f_min = pow(1.05); // neutron + pi+ threshold
1214: M2_f_max = cuts.MI_max * lcs.s;
1215: log_M2_f_min = std::log(M2_f_min);
1216: log_M2_f_max = std::log(M2_f_max);

```



```

./src/MProcess.cc 18/23
1329:
1330: if (M2_f_max <= M2_f_min) {
1331:   throw std::invalid_argument("MProcess::ForwardVolume: EXCITATION variable in invalid state");
1332:   "MProcess::SampleForwardClasses: Forward leg XL : (max = " + std::ito_string(gcuts.XL_max) +
1333:   " gives mass = " + std::ito_string(sqrt(M2_f_max)) +
1334:   " GeV, below the invariant threshold. Increase the upper (max) bound.");
1335: }
1336:
1337: if (EXCITATION == 1) { // Single
1338:
1339:   // log-change of variable
1340:   const double u = log_M2_f_min + (log_M2_f_max - log_M2_f_min) * randvec[0];
1341:   const double r1 = std::exp(u);
1342:
1343:   if (random(0, 1) < 0.5) { // 50-50
1344:     mvec[0] = msqrt(r1);
1345:     lts.excite1 = true;
1346:   } else {
1347:     mvec[1] = msqrt(r1);
1348:     lts.excite2 = true;
1349:   }
1350: } else if (EXCITATION == 2) { // Double
1351:
1352:   // log-change of variable
1353:   const double u1 = log_M2_f_min + (log_M2_f_max - log_M2_f_min) * randvec[0];
1354:   const double r1 = std::exp(u1);
1355:
1356:   const double u2 = log_M2_f_min + (log_M2_f_max - log_M2_f_min) * randvec[1];
1357:   const double r2 = std::exp(u2);
1358:
1359:   mvec[0] = msqrt(r1);
1360:   mvec[1] = msqrt(r2);
1361:   lts.excite1 = true;
1362:   lts.excite2 = true;
1363: }
1364:
1365: }
1366:
1367: // This is called last by the initialization routines
1368: // as the last step before event generation.
1369: void MProcess::SetTechnicalBoundaries(gcuts:GENCUT gcuts, unsigned int EXCITATION) {
1370:   if (gcuts.forward_pt_min < 0.0) { // Not set yet by the USER
1371:     gcuts.forward_pt_min = 0.0;
1372:   }
1373:
1374:   if (gcuts.forward_pt_max < 0.0) { // Not set yet by the USER
1375:     if (EXCITATION == 0) { // Elastic forward protons, default values
1376:       gcuts.forward_pt_max = 2.5;
1377:     }
1378:   }
1379:
1380:   else if (EXCITATION == 1) { // Single excitation
1381:     gcuts.forward_pt_max = 50.0;
1382:   } else if (EXCITATION == 2) { // Double excitation
1383:     gcuts.forward_pt_max = 100.0;
1384:   }
1385: }
1386:
1387:
1388: // Forward leg integration volume
1389: double MProcess::ForwardVolume() const {
1390:   // Forward leg phi: phi12 volume gives (2pi)^2
1391:   const double Phi_vol = pow(2.0 * M_PI, 2);
1392:
1393:   // Forward leg pt volume with log-change of variables
1394:   const double J_pt = lts.pfinal[1].Pt() * lts.pfinal[2].Pt();
1395:   const double PT_vol =
1396:     J_pt * pow(std::log(gcuts.forward_pt_max) - std::log(gcuts.forward_pt_min + ZERO_EPS));
1397:
1398:   if (EXCITATION == 0) {
1399:     return Phi_vol * PT_vol;
1400:   }
1401:   else if (EXCITATION == 1) {
1402:     // log-change of variable Jacobian
1403:     // |det(J)| = |det( [ln(b1) ln(b2)] / [ln(b1) ln(b2)] )| * exp(ln du) where u = ln(M2)
1404:     const double JM2 = (lts.excite1 ? lts.pfinal[1].M2() : lts.pfinal[2].M2());
1405:     return Phi_vol * PT_vol * JM2 * (log_M2_f_max - log_M2_f_min);
1406:   }
1407:   else if (EXCITATION == 2) {
1408:     // log-change of variable Jacobian
1409:     const double JM2 = lts.pfinal[1].M2() * lts.pfinal[2].M2();
1410:     return Phi_vol * PT_vol * JM2 * pow(2.0 * M_PI, 2) * (log_M2_f_max - log_M2_f_min);
1411:   }
1412: }

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```

./src/MProcess.cc 18/23
1412: } else {
1413:   throw std::invalid_argument("MProcess::ForwardVolume: EXCITATION variable in invalid state");
1414: }
1415:
1416: // Print fiducial cuts set by user
1417: void MProcess::PrintFiducialCuts() const {
1418:   gra::aux::PrintBar("----");
1419:   std::cout << std::endl;
1420:   std::cout << std::endl;
1421:   std::cout << "Central final states" << std::endl;
1422:   std::cout << "Fiducial cuts: " << std::endl << std::endl;
1423:   std::cout << "range:style:reset";
1424:
1425:   if (fcuts.active == true) {
1426:     std::cout << "Central final states" << std::endl;
1427:     printf("-- eta [min, max] = [0.2E, 0.2E] \n", fcuts.eta_min, fcuts.eta_max);
1428:     printf("-- pt [min, max] = [0.2E, 0.2E] GeV \n", fcuts.pt_min, fcuts.pt_max);
1429:     printf("-- Et [min, max] = [0.2E, 0.2E] GeV \n", fcuts.Et_min, fcuts.Et_max);
1430:     printf("-- Rap [min, max] = [0.2E, 0.2E] \n", fcuts.rap_min, fcuts.rap_max);
1431:     std::cout << std::endl;
1432:   }
1433:
1434:   std::cout << "Central system" << std::endl;
1435:   printf("-- M [min, max] = [0.2E, 0.2E] GeV \n", fcuts.M_min, fcuts.M_max);
1436:   printf("-- s_min [min, max] = [0.2E, 0.2E] GeV \n", fcuts.s_min, fcuts.s_max);
1437:   std::cout << std::endl;
1438:
1439:   std::cout << "Forward kinematics" << std::endl;
1440:   printf("-- lts [min, max] = [0.2E, 0.2E] GeV^2 \n", fcuts.forward_M_min, fcuts.forward_M_max);
1441:   printf("-- lts [min, max] = [0.2E, 0.2E] GeV^2 \n", fcuts.forward_t_min, fcuts.forward_t_max);
1442:   } else {
1443:     fcuts.forward_t_max;
1444:   }
1445:   std::cout << "Not active" << std::endl;
1446:
1447:   gra::aux::PrintBar("----");
1448:   std::cout << std::endl;
1449:   std::cout << "range:style:tbody";
1450:   std::cout << "Extra custom fiducial cuts:" << std::endl << std::endl;
1451:   std::cout << "range:style:reset";
1452:   if (USERCUTS == 0) {
1453:     printf("-- Active ID = %d (see M0userCuts.cc) \n", USERCUTS);
1454:   } else {
1455:     std::cout << "Not active" << std::endl;
1456:   }
1457:   std::cout << "range:fg:reset";
1458:
1459:   gra::aux::PrintBar("----");
1460:   std::cout << std::endl;
1461:   std::cout << "range:style:tbody";
1462:   std::cout << "VETO cuts:" << std::endl << std::endl;
1463:   std::cout << "range:style:reset";
1464:
1465:   if (vetcuts.active == true) {
1466:     for (std::size_t i = 0; i < vetcuts.cuts.size(); ++i) {
1467:       std::cout << "VETO cut: " << i << std::endl;
1468:       printf("  eta [min, max] = [0.2E, 0.2E] \n", vetcuts.cuts[i].eta_min,
1469:             vetcuts.cuts[i].eta_max);
1470:       printf("  pt [min, max] = [0.2E, 0.2E] GeV \n", vetcuts.cuts[i].pt_min,
1471:             vetcuts.cuts[i].pt_max);
1472:       std::cout << "Not active" << std::endl;
1473:     }
1474:   } else {
1475:     std::cout << "Not active" << std::endl;
1476:   }
1477:
1478:   gra::aux::PrintBar("----");
1479:   std::cout << std::endl;
1480:   std::cout << "range:style:tbody";
1481:   std::cout << "Central system decay tree:" << std::endl << std::endl;
1482:   std::cout << "range:style:reset";
1483:   // Print out decaytree recursively
1484:   for (std::size_t i = 0; i < lts.decaytree.size(); ++i) { PrintDecayTree(lts.decaytree[i]); }
1485:
1486:   std::cout << std::endl;
1487:   printf("Final state symmetry factor (1/S = 1/0.0E) applied at cross section "
1488:         "level \n",
1489:         1_factor);
1490:
1491:   gra::aux::PrintBar("----");
1492:   std::cout << std::endl;
1493: }

```

```

./src/MProcess.cc 19/23
1495:
1496: // Build and check lorentz scalars
1497: // Input as the number of final states
1498: bool MProcess::GetLorentzScalars(unsigned int NF) {
1499:   // Example of 4-particle Lorentz scalars:
1500:   //
1501:   //
1502:   //
1503:   // -----> t1
1504:   // |
1505:   // | s13 ----->
1506:   // |
1507:   // | s hat, t hat, s hat (s34)
1508:   // |
1509:   // | -----> t2
1510:   // |
1511:   // | -----> t3
1512:   // |
1513:   // | -----> t4
1514:   // |
1515:   // s-type Lorentz scalars --> central system indexing starts from offset
1516:   const int offset = 3; // central system indexing starts from offset
1517:
1518:   // Upper right triangle
1519:   for (std::size_t i = 0; i <= NF; ++i) { // start at 0
1520:     for (std::size_t j = 0; j <= NF; ++j) {
1521:       if (i < j) { lts.s[i][j] = (lts.pfinal[i] + lts.pfinal[j]).M2(); }
1522:       if (i == j) { lts.s[i][i] = 0; // return false; }
1523:     }
1524:   }
1525:   // Copy the i<->j permuted to the left bottom triangle
1526:   // (faster than calculating twice)
1527:   for (std::size_t i = 0; i <= NF; ++i) { // start at 0
1528:     for (std::size_t j = i+1; j <= NF; ++j) {
1529:       if (i != j) { lts.s[j][i] = lts.s[i][j]; }
1530:     }
1531:   }
1532:
1533:   // Sub invariants w.r.t central system
1534:   lts.s1 = lts.s[1][0];
1535:   lts.s2 = lts.s[2][0];
1536:   if (lts.s1 < 0 || lts.s2 < 0) { return false; }
1537:
1538:   //
1539:   //
1540:   // t-type Lorentz scalars -->
1541:   // Propagator vectors
1542:   lts.q1 = lts.pbeam1 - lts.pfinal[1];
1543:   lts.q2 = lts.pbeam2 - lts.pfinal[2];
1544:
1545:   lts.t1 = lts.q1.M2();
1546:   lts.t2 = lts.q2.M2();
1547:
1548:   if (lts.t1 > 0 || lts.t2 > 0) { return false; }
1549:
1550:   // For 2-body central processes
1551:   if (lts.decaytree.size() == 2) {
1552:     lts.t_hat = (lts.q1 - lts.decaytree[0].p4).M2(); // note q1 on both
1553:     lts.u_hat = (lts.q1 - lts.decaytree[1].p4).M2(); // in t and u!
1554:   }
1555:
1556:   for (std::size_t i = offset; i <= NF; ++i) { lts.ttt[i] = (lts.q1 - lts.pfinal[i]).M2(); }
1557:   for (std::size_t i = offset; i <= NF; ++i) {
1558:     for (std::size_t j = offset; j <= NF; ++j) {
1559:       lts.ttt_xy[i][j] = (lts.q1 - lts.pfinal[i] - lts.pfinal[j]).M2();
1560:     }
1561:   }
1562:   for (std::size_t i = offset; i <= NF; ++i) { lts.ttt_2[i] = (lts.q1 - lts.pfinal[i]).M2(); }
1563:
1564:   // Fractional longitudinal momentum loss [0,1]
1565:   lts.x1 = (1 - lts.pfinal[1].Pz() / lts.pbeam1.Pz());
1566:   lts.x2 = (1 - lts.pfinal[2].Pz() / lts.pbeam2.Pz());
1567:
1568:   // Bjorken-x [0,1] (this Lorentz invariant expression
1569:   // gives i always for elastic central production forward leg)
1570:   lts.xb1 = lts.x1 / (1 - lts.pbeam1 - lts.q1);
1571:   lts.xb2 = lts.x2 / (1 - lts.pbeam2 - lts.q1);
1572:
1573:   // Propagator pt
1574:   lts.q1 = lts.q1.Pt();
1575:   lts.q2 = lts.q2.Pt();
1576:
1577:   // Often used system variables

```

```

./src/MProcess.cc 20/23
1578:   lts.m2 = lts.pfinal[0].M2();
1579:   lts.s_hat = lts.m2;
1580:   lts.Y = lts.pfinal[0].Rap();
1581:   lts.Pt = lts.pfinal[0].Pt();
1582:
1583:   return true;
1584: }
1585:
1586: // Event output recording (to HepMC containers)
1587:
1588: bool MProcess::CommonRecord(HepMC3::GenEvent evt) {
1589:   // Initial states (4-momentum, pdg-ID, status code)
1590:   HepMC3::GenParticlePtr gen_p1 = std::make_shared<HepMC3::GenParticle>(
1591:     gra::aux::MVec2HepMC3(lts.pbeam1), lts.beam1.pdg, PDG::PDG_BEAM);
1592:   HepMC3::GenParticlePtr gen_p2 = std::make_shared<HepMC3::GenParticle>(
1593:     gra::aux::MVec2HepMC3(lts.pbeam2), lts.beam2.pdg, PDG::PDG_BEAM);
1594:
1595:   // Final states (protons/excited system)
1596:   int PDG_ID1 = lts.beam1.pdg;
1597:   int PDG_ID2 = lts.beam2.pdg;
1598:
1599:   int PDG_status1 = PDG::PDG_STABLE;
1600:   int PDG_status2 = PDG::PDG_STABLE;
1601:
1602:   if (lts.excite1 == true) {
1603:     PDG_ID1 = std::abs(PDG::PDG_NSTAR) * math::sign(lts.beam1.pdg);
1604:     PDG_status1 = PDG::PDG_INTERMEDIATE;
1605:   }
1606:   if (lts.excite2 == true) {
1607:     PDG_ID2 = std::abs(PDG::PDG_NSTAR) * math::sign(lts.beam2.pdg);
1608:     PDG_status2 = PDG::PDG_INTERMEDIATE;
1609:   }
1610:
1611:   HepMC3::GenParticlePtr gen_p1f = std::make_shared<HepMC3::GenParticle>(
1612:     gra::aux::MVec2HepMC3(lts.pfinal[1]), PDG_ID1, PDG_status1);
1613:   HepMC3::GenParticlePtr gen_p2f = std::make_shared<HepMC3::GenParticle>(
1614:     gra::aux::MVec2HepMC3(lts.pfinal[2]), PDG_ID2, PDG_status2);
1615:
1616:   //
1617:   //
1618:   // Propagator 1 and 2
1619:   // It is ill-posed to try classify pomerons/gluons etc. here, thus,
1620:   // we tag only generic propagator ID in the record
1621:   HepMC3::GenParticlePtr gen_p1p = std::make_shared<HepMC3::GenParticle>(
1622:     gra::aux::MVec2HepMC3(lts.q1), PDG::PDG_propagator, PDG::PDG_INTERMEDIATE);
1623:   HepMC3::GenParticlePtr gen_p2p = std::make_shared<HepMC3::GenParticle>(
1624:     gra::aux::MVec2HepMC3(lts.q2), PDG::PDG_propagator, PDG::PDG_INTERMEDIATE);
1625:
1626:   //
1627:   //
1628:   // Central system / resonance
1629:   HepMC3::GenParticlePtr gen_g3 = std::make_shared<HepMC3::GenParticle>(
1630:     gra::aux::MVec2HepMC3(lts.pfinal[0]), PDG::PDG_system, PDG::PDG_INTERMEDIATE);
1631:
1632:   //
1633:   //
1634:   // Construct vertices
1635:   //
1636:   // Upper proton-propagator-proton
1637:   HepMC3::GenVertexPtr v1 = std::make_shared<HepMC3::GenVertex>(
1638:     v1->add_particle_in(gen_p1);
1639:     v1->add_particle_out(gen_p1f);
1640:     v1->add_particle_out(gen_p1p);
1641:   );
1642:   // Lower proton-propagator-proton
1643:   HepMC3::GenVertexPtr v2 = std::make_shared<HepMC3::GenVertex>(
1644:     v2->add_particle_in(gen_p2);
1645:     v2->add_particle_out(gen_p2f);
1646:     v2->add_particle_out(gen_p2p);
1647:   );
1648:   // Propagator-Propagator-System vertex
1649:   HepMC3::GenVertexPtr v3 = std::make_shared<HepMC3::GenVertex>(
1650:     v3->add_particle_in(gen_p1);
1651:     v3->add_particle_in(gen_p2);
1652:     v3->add_particle_out(gen_g3);
1653:   );
1654:   v1->add_vertex(v1);
1655:   v2->add_vertex(v2);
1656:   v3->add_vertex(v3);
1657:
1658:   // System->Decay products vertex
1659:   HepMC3::GenVertexPtr v4 = std::make_shared<HepMC3::GenVertex>(

```

```

./src/MProcess.cc                21/23
1661: evt.add_vertex(v4);
1662:
1663: // Add resonance in
1664: v4->add_particle_in(gen_p);
1665:
1666: // Add direct daughters
1667: for (const auto &i : indices(lts.decaytree)) {
1668:     const int STATE = (lts.decaytree[i].legs.size() > 0) ? PDG::PDG_DECAY : PDG::PDG_STABLE;
1669:
1670:     // TBD: ADD HERE THE ctau > 1.0 cm definition for the status code
1671:
1672:     if (STATE == PDG::PDG_DECAY) {
1673:         // -----
1674:         // Pick exponential lifetime in the rest frame
1675:         const double tau = random.ExpRandom(1.0 / lts.decaytree[i].p.tau);
1676:
1677:         // Get decay vertex coordinates in the lab
1678:         const double MM = 1E3; // meters to millimeters
1679:         lts.decaytree[i].decay_position = lts.decaytree[i].p.PropagatePosition(tau, MM);
1680:         // -----
1681:     }
1682:
1683:     HepMC3::GenParticlePtr particle = std::make_shared<HepMC3::GenParticle>(
1684:         gra::aux::IMVec2HepMC3(lts.decaytree[i].p), lts.decaytree[i].p.pdg, STATE);
1685:     v4->add_particle_out(particle);
1686:
1687:     WriteDecayKinematics(lts.decaytree[i], particle, evt);
1688: }
1689:
1690: // -----
1691: // Crucial, check do we need to fragment
1692: // With v4 source is inclusive, fragmentation done at last step here (for efficiency)
1693: if (!CEPForwardFragment()) { return false; }
1694:
1695: // Upper proton excitation
1696: if (lts.excite1 && lts.decayforward1.legs.size() != 0) {
1697:     SaveBranch(evt, lts.decayforward1.gen_p2f);
1698: }
1699:
1700: // Lower proton excitation
1701: if (lts.excite2 && lts.decayforward2.legs.size() != 0) {
1702:     SaveBranch(evt, lts.decayforward2.gen_p2f);
1703: }
1704:
1705: return true;
1706: }
1707:
1708: // Save branch to event with mother pX
1709: void MProcess::SaveBranch(HepMC3::GenEvent &evt, const gra::IMDecayBranch &branch,
1710:     const HepMC3::GenParticlePtr &px) {
1711:     // Create vertex
1712:     HepMC3::GenVertexPtr vx = std::make_shared<HepMC3::GenVertex>();
1713:     evt.add_vertex(vx);
1714:
1715:     // Add mother in
1716:     vx->add_particle_in(px);
1717:
1718:     // Add direct daughters
1719:     for (const auto &i : indices(branch.legs)) {
1720:         const int STATE = (branch.legs[i].legs.size() > 0) ? PDG::PDG_DECAY : PDG::PDG_STABLE;
1721:         // TBD: ADD HERE THE ctau > 1.0 cm definition for the status code
1722:
1723:         HepMC3::GenParticlePtr particle = std::make_shared<HepMC3::GenParticle>(
1724:             gra::aux::IMVec2HepMC3(branch.legs[i].p), branch.legs[i].p.pdg, STATE);
1725:         vx->add_particle_out(particle);
1726:
1727:         WriteDecayKinematics(branch.legs[i], particle, evt);
1728:     }
1729: }
1730:
1731: // Set process
1732: void MProcess::SetProcess(std::string &str, const std::vector<aux::OneCMD> &syntax) {
1733:     // SEP IT HERE!
1734:     PROCESS = str;
1735:
1736:     // Call this always first
1737:     lts.PDG.ReadParticleData(gra::aux::GetBasePath(2) + "/modeldata/mass_width_2018.mod");
1738:
1739:     // SYNTAX Read and set new PDG input
1740:     for (const auto &i : indices(syntax)) {
1741:         if (syntax[i].id == "PDG") {
1742:             // Take target string
1743:             std::string pdg_1

```

```

./src/MProcess.cc                23/23
1827: for (const auto &i : indices(str)) {
1828:     if (str[i] == '[' && !found) {
1829:         first = str[i];
1830:         continue;
1831:     } else if (str[i] == ']' && !found) {
1832:         found = true;
1833:         continue;
1834:     }
1835:     if (found && str[i] != ']') { second = str[i]; }
1836:     if (found && str[i] == ']') { break; }
1837: }
1838:
1839: // Third
1840: int mark1 = 0;
1841: int mark2 = 0;
1842: for (const auto &i : indices(str)) {
1843:     if (str[i] == '<' ) {
1844:         mark1 = i;
1845:         continue;
1846:     }
1847:     if (str[i] == '>' ) {
1848:         mark2 = i;
1849:         break; // First >, then break
1850:     }
1851: }
1852: third = str.substr(mark1 + 1, mark2 - mark1 - 1);
1853: }
1854:
1855: // namespace gra
1856: }
1857:

```

```

./src/MProcess.cc                22/23
1744:     if (syntax[i].target.size() == 1) {
1745:         pdg = syntax[i].target[0];
1746:     } else if (syntax[i].target.size() == 0) {
1747:         throw std::invalid_argument("Syntax error: invalid PDG[] without any target [!];");
1748:     } else if (syntax[i].target.size() > 1) {
1749:         throw std::invalid_argument("Syntax error: invalid PDG[] with multiple targets inside [!];");
1750:     }
1751:     // Conversion to Integer
1752:     int pdg = 0;
1753:
1754:     try {
1755:         pdg = std::stoi(pdg_);
1756:     } catch (...) {
1757:         throw std::invalid_argument(
1758:             "Syntax error: invalid PDG[] target number, string to int conversion fails!");
1759:     }
1760:
1761:     // Try to find the particle from PDG table, will throw exception if fails
1762:     MParticle p = lts.PDG.FindByPDG(pdg);
1763:
1764:     // Check if it has an anti-particle
1765:     MParticle p_anti;
1766:     bool found_anti = false;
1767:     try {
1768:         p_anti = lts.PDG.FindByPDG(-pdg);
1769:         found_anti = true;
1770:     } catch (...) {} // do nothing.
1771:
1772:     // Set new properties
1773:     for (const auto &k : syntax[i].arg) {
1774:         if (k.first == "M") {
1775:             p.mass = std::stod(k.second);
1776:             if (found_anti) { p_anti.mass = p.mass; }
1777:         }
1778:         if (k.first == "W") {
1779:             p.width = std::stod(k.second);
1780:             p.mass = PDG::ubar / p.width; // mean lifetime in the rest frame
1781:             if (found_anti) {
1782:                 p_anti.width = p.width;
1783:                 p_anti.tau = p.tau;
1784:             }
1785:         }
1786:     }
1787:     // Set new modified to the PDG table
1788:     lts.PDG.PDG_table[pg] = p;
1789:     if (found_anti) { lts.PDG.PDG_table[-pdg] = p_anti; }
1790:
1791:     std::cout << "rangif:read\n";
1792:     << "MProcess::SetProcess: New particle properties set "\n";
1793:     << "with PDGIDnumber[" <<keyval> "] syntax:"\n";
1794:     << "rangif:reset <& std::endl;";
1795:     p.print();
1796: }
1797: }
1798:
1799: // Remove whitespace
1800: std::string::iterator end_pos = std::remove(str.begin(), str.end(), ' ');
1801: str.erase(end_pos, str.end());
1802:
1803: // Parse commandline string
1804: std::string istate = "";
1805: std::string channel = "";
1806: std::string me = "";
1807: ParseCMD(str, istate, channel, me);
1808:
1809: // Setup subprocess
1810: ProcPtr.Initialize(istate, channel);
1811:
1812: // Check do we find the process
1813: if (ProcPtr.Processes.count(str)) {
1814:     // fine
1815: } else {
1816:     throw std::invalid_argument("MProcess::SetProcess: Unknown PROCESS: " + str);
1817: }
1818: }
1819:
1820: // "FR[RES]<C>" is an example of valid string to be parsed
1821: //
1822: void MProcess::ParseCMD(const std::string &str, std::string &istate, std::string &channel,
1823:     std::string &third, const {
1824:     // First and Second
1825:     bool found = false;
1826: }

```

```

./src/MQuasiElastic.cc          1/11
1: // QuasiElastic (EL,SD,DD) and soft ND (simplified) class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <iostream>
9: #include <random>
10: #include <vector>
11:
12: // HepMC3
13: #include "HepMC3/FourVector.h"
14: #include "HepMC3/GenEvent.h"
15: #include "HepMC3/GenParticle.h"
16: #include "HepMC3/GenVertex.h"
17:
18: // DGL
19: #include "Graniitti/NDux.h"
20: #include "Graniitti/MForm.h"
21: #include "Graniitti/MKinematics.h"
22: #include "Graniitti/MMath.h"
23: #include "Graniitti/MProcess.h"
24: #include "Graniitti/MQuasiElastic.h"
25: #include "Graniitti/MUserCuts.h"
26:
27: using gra::PDG::GeV2Barn;
28: using gra::PDG::Imp;
29: using gra::aux::Indices;
30: using gra::math::FJ;
31: using gra::math::abs2;
32: using gra::math::mag2;
33: using gra::math::pow2;
34: using gra::math::zi;
35:
36: namespace gra {
37: // This is needed by construction
38: MQuasiElastic::MQuasiElastic() { Initialize(); }
39:
40: // Constructor
41: MQuasiElastic::MQuasiElastic(std::string process, const std::vector<aux::OneCMD> &syntax) {
42:     Initialize();
43:     InitListParams();
44:     SetProcess(process, syntax);
45: }
46:
47: // Init final states
48: MAVec zeroVec(0, 0, 0, 0);
49: for (std::size_t i = 0; i < 10; ++i) { lts.pfnal.push_back(zeroVec); }
49:
50: // Random weights
51: MAXFORM = std::vector<double>(100, 0.0);
52: std::cout << "MQuasiElastic: [Constructor done]" << std::endl;
53: }
54:
55: void MQuasiElastic::Initialize() {
56:     const std::vector<std::string> supported = {"TK"},
57:         CID = "Q",
58:         ProcPtr = MSubProc(supported, CID);
59: }
60:
61: // Destructor
62: MQuasiElastic::~MQuasiElastic() {}
63:
64: // Initialize cut and process specific postsetup
65: void MQuasiElastic::post_Constructor() {
66:     // Set phase space dimension
67:     if (ProcPtr.CHANNEL == "EL") { ProcPtr.LIFSDIM = 1; }
68:     if (ProcPtr.CHANNEL == "SD") { ProcPtr.LIFSDIM = 2; }
69:     if (ProcPtr.CHANNEL == "DD") { ProcPtr.LIFSDIM = 3; }
70:     if (ProcPtr.CHANNEL == "ND") { ProcPtr.LIFSDIM = 1; } // Keep it 1
71:
72:     Eikonal.Numerics.MaxLoopMT = 3.0;
73: }
74:
75: // Fiducial user cuts
76: bool MQuasiElastic::FiducialCuts() const {
77:     if (fcuts.Active == true) {
78:         // EL cuts
79:         if (ProcPtr.CHANNEL == "EL") {
80:             if (fcuts.forward_min < std::abs(lts.t) && std::abs(lts.t) < fcuts.forward_t_max) {
81:                 // fine
82:             } else {
83:                 return false; // not fine

```

```

./src/MQuasiElastic.cc          2/11
84: }
85: }
86: }
87: // SD cuts
88: if (ProcPtr.CHANNEL == "SD") {
89:   if (lts.pfinal[1].M() > 1.0) { // this one is excited system
90:     if (fcuts.forward_m_min < lts.pfinal[1].M() && lts.pfinal[1].M() < fcuts.forward_m_max &&
91:         fcuts.forward_t_min < std::abs(lts.t) && std::abs(lts.t) < fcuts.forward_t_max) {
92:       // fine
93:     } else {
94:       return false; // not fine
95:     }
96:   } else {
97:     if (fcuts.forward_m_min < lts.pfinal[2].M() && lts.pfinal[2].M() < fcuts.forward_m_max &&
98:         fcuts.forward_t_min < std::abs(lts.t) && std::abs(lts.t) < fcuts.forward_t_max) {
99:       // fine
100:     } else {
101:       return false; // not fine
102:     }
103:   }
104: }
105: }
106: // DD cuts
107: if (ProcPtr.CHANNEL == "DD") {
108:   if (fcuts.forward_m_min < lts.pfinal[1].M() && lts.pfinal[1].M() < fcuts.forward_m_max &&
109:       fcuts.forward_m_min < lts.pfinal[2].M() && lts.pfinal[2].M() < fcuts.forward_m_max &&
110:       fcuts.forward_t_min < std::abs(lts.t) && std::abs(lts.t) < fcuts.forward_t_max) {
111:     // fine
112:   } else {
113:     return false; // not fine
114:   }
115: }
116: }
117: // Check user cuts (do not substitute to kinematics =
118: // UserCuts)
119: if (!UserCuts(USERCUTS, lts)) {
120:   return false; // not fine
121: }
122: }
123: }
124: }
125: }
126: }
127: }
128: }
129: }
130: }
131: }
132: }
133: }
134: }
135: }
136: }
137: }
138: }
139: }
140: }
141: }
142: }
143: }
144: }
145: }
146: }
147: }
148: }
149: }
150: }
151: }
152: }
153: }
154: }
155: }
156: }
157: }
158: }
159: }
160: }
161: }
162: }
163: }
164: }
165: }
166: }

```

```

./src/MQuasiElastic.cc          4/11
250:   gra:aux:M4Vec2HepMC3(etree[1].q2), PDG:PDG_propagator, PDG:PDG_INTERMEDIATE);
251: }
252: // Virtual state
253: M4Vec system4vec(etree[1].k.px(), etree[1].k.py(), etree[1].k.pz(), etree[1].k.e());
254: }
255: HepMC3::GenParticlePtr gen_X = std::make_shared<HepMC3::GenParticle>(
256:   gra:aux:M4Vec2HepMC3(system4vec), PDG:PDG_system, PDG:PDG_INTERMEDIATE);
257: }
258: HepMC3::GenVertexPtr vx = std::make_shared<HepMC3::GenVertex>();
259: }
260: if (l != etree.size() - 1) {
261:   // Upper vertex
262:   HepMC3::GenVertexPtr vUX = std::make_shared<HepMC3::GenVertex>();
263:   vUX->add_particle_in(gen_p1);
264:   vUX->add_particle_out(gen_p2);
265:   vUX->add_particle_out(gen_p3);
266: }
267: // Lower vertex
268: HepMC3::GenVertexPtr vXD = std::make_shared<HepMC3::GenVertex>();
269: vXD->add_particle_in(gen_p2);
270: vXD->add_particle_out(gen_p2);
271: vXD->add_particle_out(gen_p2);
272: }
273: // Add vertices
274: evt.add_vertex(vUX);
275: evt.add_vertex(vXD);
276: }
277: // Pomeron-Pomeron-Virtual state vertex
278: vx->add_particle_in(gen_p1);
279: vx->add_particle_in(gen_p2);
280: vx->add_particle_out(gen_X);
281: }
282: } else { // Last MPI, fuse remnant + remnant -> system
283: }
284: // Proton-Proton-Virtual state vertex
285: vx = std::make_shared<HepMC3::GenVertex>();
286: vx->add_particle_in(gen_p1);
287: vx->add_particle_in(gen_p2);
288: vx->add_particle_out(gen_X);
289: }
290: }
291: // Add vertex to the event
292: evt.add_vertex(vX);
293: }
294: }
295: // Quantum numbers distributed here
296: int B = 0; // baryon number
297: int Q = 0; // charge
298: double Q2_scale = system4vec.M2();
299: }
300: // Beam fragments carry the initial state quantum numbers,
301: // we distribute them here for the second and the last MPI
302: if (l == etree.size() - 2) {
303:   B = math::sign(lts.beam2.pdg); // Proton(anti) proton beams only
304:   Q = B;
305: }
306: if (l == etree.size() - 1) {
307:   B = math::sign(lts.beam1.pdg);
308:   Q = B;
309: }
310: }
311: // Excite it
312: MDecayBranch branch;
313: if (!ExciteContinuum(system4vec, branch, Q2_scale, B, Q)) {
314:   std::cout << "HD failed with ExciteContinuum" << std::endl;
315:   return false; // failed
316: }
317: }
318: // Save it
319: SaveBranch(evt, branch, gen_X);
320: }
321: }
322: }
323: }
324: }
325: }
326: }
327: // Initial states (4-momentum, pdg-id, status code)
328: HepMC3::GenParticlePtr gen_p1 = std::make_shared<HepMC3::GenParticle>(
329:   gra:aux:M4Vec2HepMC3(lts.pbeam1), lts.beam1.pdg, PDG:PDG_BEAM);
330: HepMC3::GenParticlePtr gen_p2 = std::make_shared<HepMC3::GenParticle>(
331:   gra:aux:M4Vec2HepMC3(lts.pbeam2), lts.beam2.pdg, PDG:PDG_BEAM);
332: }

```

```

./src/MQuasiElastic.cc          3/11
167:   aux.vetocuts_ok = VetoCuts();
168: }
169: }
170: // ** EVENT WEIGHT **
171: const double LIPS = B3PhaseSpaceWeight(); // Phase-space weight
172: const double MatEQ = GetM2(); // Matrix element squared
173: }
174: // Total weight: phase-space x |M|^2 x barn units
175: W = LIPS * MatEQ; // Total weight:
176: }
177: }
178: aux.amplitude_ok = CheckInfNaN(W);
179: }
180: // Fill Histograms
181: // const double totalweight = W * aux.vegusweight;
182: // FillHistograms(totalweight, lts) // This seqfaults, because we do not
183: // have
184: // enough observables
185: } else { // Non-diffractive
186: }
187: aux.kinematic_ok = true; // This is always the case
188: aux.fiducial_ok = true;
189: aux.vetocuts_ok = true;
190: }
191: const double LIPS = B3PhaseSpaceWeight(); // Phase-space weight
192: const double MatEQ = abs(PolySoft(randvec)); // Matrix element squared
193: // W = LIPS * MatEQ; // Total weight: phase-space x |M|^2 x barn units
194: W = LIPS * MatEQ; // Total weight:
195: }
196: }
197: aux.amplitude_ok = CheckInfNaN(W);
198: }
199: // Trigger forced event generation (using last component of auxvar)
200: // (do not need for outside acceptance/rejection with this process)
201: if (W > 0) {
202:   aux.forced_accept = true;
203: } else {
204:   aux.forced_accept = false; // Failed event
205: }
206: }
207: }
208: }
209: }
210: }
211: // Record HepMC3 event
212: bool MQuasiElastic::EventRecord(HepMC3::GenEvent &evt) {
213: // Non-diffractive
214: if (ProcPtr.CHANNEL == "ND") {
215:   HepMC3::GenParticlePtr gen_p1;
216:   HepMC3::GenParticlePtr gen_p2;
217:   HepMC3::GenParticlePtr gen_p3;
218:   HepMC3::GenParticlePtr gen_p2f;
219: }
220: // Loop over multiple "cut pomérons"
221: for (const auto &l : indices(etrree)) {
222:   if (l == 0) { // First
223: }
224: // Initial state protons (4-momentum, pdg-id, status code)
225: gen_p1 = std::make_shared<HepMC3::GenParticle>(
226:   gra:aux:M4Vec2HepMC3(etrree[1].p1), (l == 0 ? lts.beam1.pdg : PDG:PDG_fragment,
227:   (l == 0 ? PDG:PDG_BEAM : PDG:PDG_INTERMEDIATE));
228: }
229: gen_p2 = std::make_shared<HepMC3::GenParticle>(
230:   gra:aux:M4Vec2HepMC3(etrree[1].p2), (l == 0 ? lts.beam2.pdg : PDG:PDG_fragment,
231:   (l == 0 ? PDG:PDG_BEAM : PDG:PDG_INTERMEDIATE));
232: }
233: } else { // Others recursively
234: }
235: gen_p1 = gen_p1f;
236: gen_p2 = gen_p2f;
237: }
238: }
239: // Final state (or intermediate in the chain) fragments
240: gen_p3 = std::make_shared<HepMC3::GenParticle>(gra:aux:M4Vec2HepMC3(etrree[1].p1f),
241:   PDG:PDG_fragment, PDG:PDG_INTERMEDIATE);
242: }
243: gen_p2f = std::make_shared<HepMC3::GenParticle>(gra:aux:M4Vec2HepMC3(etrree[1].p2f),
244:   PDG:PDG_fragment, PDG:PDG_INTERMEDIATE);
245: }
246: }
247: }
248: }
249: }

```

```

./src/MQuasiElastic.cc          5/11
333: // Propagator 4-vector and generator particle
334: M4Vec q1(lts.pbeam1 - lts.pfinal[1]);
335: HepMC3::GenParticlePtr gen_q1 = std::make_shared<HepMC3::GenParticle>(
336:   gra:aux:M4Vec2HepMC3(q1), PDG:PDG_propagator, PDG:PDG_INTERMEDIATE);
337: }
338: // Final state protons/excited systems
339: int PDG_ID1 = lts.beam1.pdg;
340: int PDG_ID2 = lts.beam2.pdg;
341: int PDG_status1 = PDG:PDG_STABLE;
342: int PDG_status2 = PDG:PDG_STABLE;
343: }
344: // EL
345: // already fine
346: }
347: // SD
348: if (ProcPtr.CHANNEL == "SD") {
349:   if (lts.excited) { // proton 1 excited
350:     PDG_ID1 = std::abs(PDG:PDG_NSTAR) * math::sign(lts.beam1.pdg);
351:     PDG_status1 = PDG:PDG_INTERMEDIATE;
352:   } else {
353:     PDG_ID1 = std::abs(PDG:PDG_NSTAR) * math::sign(lts.beam1.pdg);
354:     PDG_status2 = PDG:PDG_INTERMEDIATE;
355:   }
356: }
357: // DD
358: if (ProcPtr.CHANNEL == "DD") {
359:   PDG_ID1 = std::abs(PDG:PDG_NSTAR) * math::sign(lts.beam1.pdg);
360:   PDG_status1 = PDG:PDG_INTERMEDIATE;
361:   PDG_ID2 = std::abs(PDG:PDG_NSTAR) * math::sign(lts.beam2.pdg);
362:   PDG_status2 = PDG:PDG_INTERMEDIATE;
363: }
364: }
365: HepMC3::GenParticlePtr gen_p1f = std::make_shared<HepMC3::GenParticle>(
366:   gra:aux:M4Vec2HepMC3(lts.pfinal[1]), PDG_ID1, PDG_status1);
367: HepMC3::GenParticlePtr gen_p2f = std::make_shared<HepMC3::GenParticle>(
368:   gra:aux:M4Vec2HepMC3(lts.pfinal[2]), PDG_ID2, PDG_status2);
369: }
370: // Construct vertices
371: }
372: // Upper proton-pomeron-proton
373: HepMC3::GenVertexPtr v1 = std::make_shared<HepMC3::GenVertex>();
374: v1->add_particle_in(gen_p1);
375: v1->add_particle_out(gen_p1);
376: v1->add_particle_out(gen_q1);
377: }
378: // Lower proton-pomeron-proton
379: HepMC3::GenVertexPtr v2 = std::make_shared<HepMC3::GenVertex>();
380: v2->add_particle_in(gen_p2);
381: v2->add_particle_out(gen_p2);
382: v2->add_particle_in(gen_q1);
383: }
384: // Finally add all vertices
385: evt.add_vertex(v1);
386: evt.add_vertex(v2);
387: }
388: // Upper proton excitation
389: if (lts.excited) {
390:   const int B = math::sign(lts.beam1.pdg); // baryon number the same as charge
391:   const int Q = B;
392: }
393: if (!ExciteContinuum(lts.pfinal[1], lts.decayforward1, lts.pfinal[1].M2(), B, Q)) {
394:   return false; // failed
395: }
396: SaveBranch(evt, lts.decayforward1, gen_p1f);
397: }
398: // Lower proton excitation
399: if (lts.excited2) {
400:   const int B = math::sign(lts.beam2.pdg); // baryon number the same as charge
401:   const int Q = B;
402: }
403: if (!ExciteContinuum(lts.pfinal[2], lts.decayforward2, lts.pfinal[2].M2(), B, Q)) {
404:   return false; // failed
405: }
406: SaveBranch(evt, lts.decayforward2, gen_p2f);
407: }
408: }
409: }
410: }
411: }
412: // Print out setup
413: void MQuasiElastic::PrintInit(bool silent) const {
414:   if (!silent) {
415:     PrintSetup();
416:   }
417: }

```



```

./src/MQuasiElastic.cc          10/11
748:
749: // Calculate kinematically valid t-range
750: const double s1 = lts.pbeam1.M2();
751: const double s2 = lts.pbeam2.M2();
752:
753: // Mandelstam t-range calculation
754: gra:kkinematics:TwoToTwoLimit(lts.s, s1, s2, s3, s4, t_min, t_max);
755:
756: // Then limit the "diffraction cone" due to screening loop limit
757: t_min = std:min(std::max(-pow(Ekronal.Numerics.MaxLoopKT), t_min), t_max);
758:
759: // Logarithmic change of variable sampling
760: const double A = std:abs(t_max);
761: const double B = std:abs(t_min);
762:
763: const double r =
764: std:(log(A + ZERO_EPS) + std:(log(B + ZERO_EPS) - std:(log(A + ZERO_EPS))) * randvec(0);
765: const double t = -std:exp(r);
766:
767: return B3BuildKin(s3, s4, t);
768:
769:
770: // Build kinematics for elastic, single and double diffractive 2->2 quasielastic
771: bool MQuasiElastic:B3BuildKin(double s3, double s4, double t) {
772: static const double s1 = lts.pbeam1.M2();
773: static const double s2 = lts.pbeam2.M2();
774: static const M4Vec beamsum = lts.pbeam1 + lts.pbeam2;
775:
776: // Scattering angle based on invariants
777: double theta = std:acos(kinematics:CosThetaTtar(lts.s, t, s1, s2, s3, s4));
778:
779: // Forward/backward solution flip (skip these, rare)
780: if (std:cos(theta) < 0) { return false; }
781: // theta = std:cos(theta) < 0 ? gra:math:PI - theta : theta;
782:
783: // Outgoing 4-momentum by Kallen (triangle) function in the center-of-momentum
784: // frame
785: const double pnorm = kinematics:DecayMomentum(lts.sqrt_s, msqrt(s3), msqrt(s4));
786: M4Vec p3(0, 0, pnorm, 0.5 * (lts.s + s3 - s4) / lts.sqrt_s);
787: M4Vec p4(0, 0, -pnorm, 0.5 * (lts.s + s4 - s3) / lts.sqrt_s);
788:
789: // Transverse momentum by orienting with random rotation (theta, phi)
790: const double phi = random.U(0.0, 2.0 * gra:math:PI); // Flat phi
791: gra:kkinematics:Rotate(p3, theta, phi);
792: gra:kkinematics:Rotate(p4, theta, phi);
793:
794: // -----
795: // Now boost if asymmetric beams
796: if (std:abs(beamsum.Pz()) != 0) {
797: const rap ant sign = 1; // positive > boost to the Lab
798: kinematics:LorentzBoost(beamsum, lts.sqrt_s, p3, sign);
799: kinematics:LorentzBoost(beamsum, lts.sqrt_s, p4, sign);
800: }
801: // -----
802:
803: lts.pfinal[1] = p3;
804: lts.pfinal[2] = p4;
805:
806: // Check Energy-Momentum conservation
807: if (!gra:math:CheckEMC(beamsum - (lts.pfinal[1] + lts.pfinal[2]))) { return false; }
808:
809: return B3GetLorentzScalars();
810:
811:
812: // Build and check scalars
813: bool MQuasiElastic:B3GetLorentzScalars() {
814: // Calculate Lorentz scalars
815: lts.ss[1][1] = lts.pfinal[1].M2();
816: lts.ss[2][2] = lts.pfinal[2].M2();
817:
818: lts.t = (lts.pbeam1 - lts.pfinal[1]).M2();
819: lts.u = (lts.pbeam1 - lts.pfinal[2]).M2();
820:
821: // Test scalars
822: if (lts.ss[1][1] > lts.s) return false;
823: if (lts.ss[2][2] > lts.s) return false;
824: if (lts.t > 0) return false;
825: if (lts.u > 0) return false;
826:
827: printf("s = %E, s^1/2 = %E \n", lts.s, msqrt(lts.s));
828: printf("t = %E, u = %E \n", lts.t, lts.u);
829: printf("s = %E, sd = %E \n", lts.ss[1][1], lts.ss[2][2]);
830: printf("\n");
}

```

```

./src/MQuasiElastic.cc          11/11
831: //
832: return true;
833: }
834:
835: // 1/2/3-Dim Integral Volume [t] x [M^2] x [M^2]
836:
837: double MQuasiElastic:B3IntegralVolume() const {
838: // Mandelstam t, log change of variable, volume times Jacobian
839: const double A = std:abs(t_max);
840: const double B = std:abs(t_min);
841: const double t_VOL = (std:(log(A + ZERO_EPS) - std:(log(A + ZERO_EPS))) * std:abs(lts.t);
842:
843: if (ProcPtr.CHANNEL == "EL") {
844: return t_VOL;
845: }
846: } else if (ProcPtr.CHANNEL == "SD") {
847: // Jacobian from log-change of variable;
848: // \int_{t_min}^t f(M^2) dM^2 = \int_{t_min}^t f(u) du / (u ln(u)) f(u) du, where u = ln(M^2)
849: const double A = (lts.s - t_min) / lts.ss[1][1];
850: const double M2_VOL = (log(M2_f_max - log(M2_f_min)) * J;
851: return t_VOL + M2_VOL;
852: }
853: } else if (ProcPtr.CHANNEL == "DD") {
854: // Jacobian from log-change of variables
855: const double J = lts.ss[1][1] * lts.ss[2][2];
856: const double M2_VOL = (log(M2_f_max - log(M2_f_min)) * (log(M2_M2_max - log(M2_f_min)) * J;
857: return t_VOL + M2_VOL;
858: }
859: }
860: return 0;
861: }
862: }
863:
864: // Standard phase space for m1 + m2 -> m3 + m4
865:
866: double MQuasiElastic:B3PhaseSpaceWeight() const {
867: // expression -> id * ps * s^2 / (t * s * m3 * m4)
868: const double norm = 16.0 * gra:math:PI *
869: pow2(lts.s - gra:kkinematics:beta2(lts.s, lts.beam1.mass, lts.beam2.mass));
870:
871: if (ProcPtr.CHANNEL == "EL") {
872: return s.0 / norm;
873: } else if (ProcPtr.CHANNEL == "SD") {
874: return 2.0 / norm; // Factor of two in
875: // numerator from single
876: // diffraction left + right
877: } else if (ProcPtr.CHANNEL == "DD") {
878: return s.0 / norm;
879: } else if (ProcPtr.CHANNEL == "DD") {
880: return 1.0;
881: } else {
882: return 0;
883: }
884: }
885:
886: // namespace gra
887:

```

```

./src/MAux.cc                   1/10
1: // I/O aux functions
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <fstream>
9: #include <iostream>
10: #include <regex>
11: #include <vector>
12: // #include <experimental/filesystem>
13:
14: // C system functions
15: #include <fcntl.h>
16: #include <sys/resource.h>
17: #include <sys/stat.h>
18: #include <sys/statvfs.h>
19: #include <sys/time.h>
20: #include <sys/types.h>
21: #include <sys/utime.h>
22: #include <unistd.h>
23: #include <unistd.h>
24:
25: // Libraries
26: #include "HepMC3/FourVector.h"
27: #include "HepMC3/LHAPDF.h"
28:
29: // Own
30: #include "Granitt/M4Vec.h"
31: #include "Granitt/MAux.h"
32: #include "Granitt/M4Vec.h"
33: #include "Granitt/MSpin.h"
34: #include "Granitt/M5udakov.h"
35:
36: // Libraries
37: #include "rang.hpp"
38:
39: namespace gra {
40: namespace aux {
41:
42: // -----
43: // FIXED HERE manually
44:
45: double GetVersion() { return 1.05; }
46: std:string GetVersionType() { return "release"; }
47: std:string GetVersionDate() { return "26.02.2020"; }
48: std:string GetVersionDate() { return "performance, IO, new HepMC3 & fixes"; }
49:
50: // -----
51:
52: // Check do we have terminal output (true), or output to file (false)
53: static const bool IS_TERMINAL = isatty(fileno(stdout)) != 0;
54:
55: // Print input arguments
56: void PrintArgv(int argc, char *argv[]) {
57: std::cout << rang:fg:green << " " << "\n";
58: for (int i = 0; i < argc; ++i) {
59: const std:string s = std:string(argv[i]);
60: if (s.find(' ') != std:string::npos) { // e.g. "PP[COMP]< > pi+ pi-"
61: std::cout << "\"" << s << "\"\n";
62: } else {
63: std::cout << s << "\n";
64: }
65: }
66: }
67: std::cout << rang:fg:reset << std::endl;
68:
69: // -----
70: // Download LHAPDFset automatically
71: void AutoDownloadLHAPDF(const std:string pdfname) {
72: std::cout << rang:fg:red
73: << "aux: AutoDownloadLHAPDF: Trying automatic download." << rang:fg:reset
74: << std::endl;
75: std::cout << std::endl;
76:
77: // Get install path and remove "\n"
78: std:string INSTALLPATH = aux:ExecCommand("lhpadf-config --prefix");
79: INSTALLPATH.erase(std::remove(INSTALLPATH.begin(), INSTALLPATH.end(), '\\'), INSTALLPATH.end());
80:
81: if (INSTALLPATH.find("command not found") != std:string::npos) {
82: // there std:invalid_argument ("aux: AutoDownloadLHAPDF: Failure: lhpadf-config command missing");
83: }

```

```

./src/MAux.cc                   2/10
84:
85: // Download, untar
86: std:string cmd = "wget http://lhpadfsets.web.cern.ch/lhpadfsets/current/" + pdfname +
87: ".tar.gz -O - | " + " tar xz -C " + INSTALLPATH + "/share/LHAPDF";
88: std::cout << cmd << std::endl;
89: std:string OUTPUT = aux:ExecCommand(cmd);
90:
91:
92: // Run terminal command, get output to std:string
93: std:string ExecCommand(const std:string cmd) {
94: std::array<char, 128> buffer;
95: std:string result;
96: std::unique_ptr<FILE, decltype(&close)> pipe(popen(cmd.c_str(), "r"), &close);
97: if (!pipe) { throw std:runtime_error("aux: ExecCommand: popen() failed!"); }
98: while (!fgets(buffer.data(), buffer.size(), pipe.get()) != nullptr) { result += buffer.data(); }
99: return result;
100: }
101:
102: std:string GetExecutablePath() {
103: char buff[2048];
104: ssize_t len = readlink("/proc/self/exe", buff, sizeof(buff) - 1);
105: if (len != -1) {
106: buff[len] = '\0';
107: return std:string(buff);
108: }
109: return "";
110: }
111:
112: // level 0 returns same as GetExecutablePath;
113: // /home/user/granitt/bin/gz
114: // /
115: // level 1 returns
116: // /home/user/granitt/bin
117: // /
118: // etc..
119:
120: std:string GetBasePath(std::size_t level) {
121: std:string s = GetExecutablePath();
122: char sep = '/';
123:
124: for (std::size_t k = 0; k < level; ++k) {
125: // Search backwards from end of string
126: std::size_t i = s.rfind(sep);
127: if (i != std:string::npos) {
128: s = s.substr(0, i);
129: } else {
130: return s;
131: }
132: }
133: return s;
134: }
135:
136: // alias
137: namespace fs = std:experimental:filesystem;
138:
139: std:string GetCurrentPath() {
140: fs::path cwd = std:experimental:filesystem:current_path();
141: return cwd.string();
142: }
143:
144:
145: // Folder exists
146: bool FileExists(const fs::path p, fs::file_status s) {
147: if (fs::status_known(s) ? fs::exists(s) : fs::exists(p)) {
148: return true;
149: } else {
150: return false;
151: }
152: }
153:
154: //
155: // Get filesize in bytes
156: std::uintmax_t GetFileSize(const std:string filename) {
157: namespace fs = std:experimental:filesystem;
158: fs::path p = fs::current_path() / filename;
159: return fs::file_size(p);
160: }
161:
162: struct stat stat_buf;
163: use rc = stat(filename.c_str(), &stat_buf);
164: return rc == 0 ? stat_buf.st_size : 0;

```

```

167: }
168:
169: // Get Process Memory Usage (Linux/BSD/OSX) in bytes
170: void GetProcessMemory(double &peak_use, double &resident_use) {
171:     // Peak memory
172:     struct rusage usage;
173:     getrusage(RUSAGE_SELF, &usage);
174:     peak_use = usage.ru_maxrss * 1024L;
175:
176:     // Current memory
177:     long rss = 0;
178:     FILE *fp = NULL;
179:     if ((fp = fopen("/proc/self/statm", "r")) == NULL) {
180:         resident_use = 0.0;
181:         return;
182:     }
183:     if (fscanf(fp, "%ld", &rss) != 1) { // Reading problem
184:         resident_use = 0.0;
185:     } else { // fine
186:         resident_use = rss * (size_t)sysconf(_SC_PAGESIZE);
187:     }
188:     fclose(fp);
189: }
190:
191: // Get disk usage
192: void GetDiskUsage(const std::string &path, int64_t &size, int64_t &free, int64_t &used) {
193:     int64_t frsize;
194:     int64_t blocks;
195:     int64_t bfree;
196:
197:     struct statvfs buf;
198:     int ret = statvfs(path.c_str(), &buf);
199:
200:     if (!ret) {
201:         frsize = buf.f_frsize; // block size
202:         blocks = buf.f_blocks; // blocks
203:         bfree = buf.f_bfree; // free blocks
204:
205:         size = frsize * blocks;
206:         free = frsize * bfree;
207:         used = size - free;
208:     }
209: }
210:
211: // Get available memory in bytes
212: unsigned long long TotalSystemMemory() {
213:     long pages = sysconf(_SC_PHYS_PAGES);
214:     long page_size = sysconf(_SC_PAGE_SIZE);
215:     return pages * page_size;
216: }
217:
218: // System information
219: std::string SystemName() {
220:     struct utsname name;
221:     uname(&name);
222:
223:     std::string sysname(name.sysname);
224:     std::string nodename(name.nodename);
225:     std::string release(name.release);
226:     std::string version(name.version);
227:     std::string machine(name.machine);
228:
229:     std::string s = " ";
230:
231:     return sysname + s + release + s + version + s + machine;
232: }
233:
234: // Get system hostname
235: std::string HostName() {
236:     char hostname[2048];
237:     gethostname(hostname, 2048);
238:
239:     std::string str(hostname);
240:     return str;
241: }
242:
243: // Get current date/time, format is YYYY-MM-DD HH:mm:ss
244: // http://en.cppreference.com/w/cpp/chrono/c/strftime
245: // For more information about date/Time format
246: const std::string DateTime() {
247:     time_t now = time(0);
248:     struct tm tstruct;
249:     char buf[80];

```

```

333: // Compress all extra spaces (compress to a space)
334: // This also processes through the strings "", not problems for use,
335: // but (2009/08/25)
336: data = std::regex_replace(data, std::regex(R"(s+)", ""));
337:
338: // JSONS: <Objects may have a single trailing comma>
339: // Replace ",," or ",," with ","
340: data = std::regex_replace(data, std::regex(R"(,)", ""));
341:
342: // JSONS: <Arrays may have a single trailing comma>
343: // Replace ",," or ",," with ","
344: data = std::regex_replace(data, std::regex(R"(,)", ""));
345:
346: // std::cout << std::endl << data << std::endl << std::endl;
347: return data;
348: }
349:
350: // Check if it integer digits
351: bool IsIntegerDigits(const std::string &str) {
352:     return str.find_first_not_of("0123456789") == std::string::npos;
353: }
354:
355: // Return particle parity as a string
356: std::string ParityToString(int value) {
357:     if (value > 0) {
358:         return "+";
359:     } else if (value == 0) {
360:         return "=";
361:     } else {
362:         return "-";
363:     }
364: }
365:
366: // Return particle charge as a string
367: std::string Charge3XtoString(int q3) {
368:     const std::string sign = (q3 < 0) ? "-" : "+";
369:     const int absq3 = std::abs(q3);
370:
371:     if (absq3 == 6) {
372:         return sign + "2";
373:     } else if (absq3 == 3) {
374:         return sign + "1";
375:     } else if (absq3 == 2) {
376:         return sign + "2/3";
377:     } else if (absq3 == 1) {
378:         return sign + "1/3";
379:     } else {
380:         return "0";
381:     }
382: }
383:
384: // Return spin as a string
385: std::string Spin3toString(int J2) {
386:     if (J2 == 10) {
387:         return "5";
388:     } else if (J2 == 9) {
389:         return "9/2";
390:     } else if (J2 == 8) {
391:         return "4";
392:     } else if (J2 == 7) {
393:         return "7/2";
394:     } else if (J2 == 6) {
395:         return "3";
396:     } else if (J2 == 5) {
397:         return "5/2";
398:     } else if (J2 == 4) {
399:         return "2";
400:     } else if (J2 == 3) {
401:         return "3/2";
402:     } else if (J2 == 2) {
403:         return "1";
404:     } else if (J2 == 1) {
405:         return "1/2";
406:     } else {
407:         return "0";
408:     }
409: }
410: // Split a string to strings separated by delimiter
411: std::vector<std::string> SplitStr2Str(const std::string input, const char delim, bool trimextraspaces) {
412:     std::vector<std::string> output;
413:     std::stringstream ss(input);
414:
415:     // String by string

```

```

250: tstruct = "localtime(now);
251:
252: strftime(buf, sizeof(buf), "%Y-%m-%d %X", &tstruct);
253: return buf;
254: }
255:
256: // Print out progress bar visualization
257: void PrintProgress(double ratio) {
258:     if (ratio > 1.0) { ratio = 1.0; }
259:
260:     #define BAR "|||||
261:     const int WIDTH = 62;
262:     const int pos = static_cast<int>(ratio * 100);
263:     const int left = static_cast<int>(ratio * WIDTH);
264:     const int right = WIDTH - left;
265:
266:     if (getchar() != '\n') { // If no terminal, do not print!
267:         std::cout << "\r";
268:         print("\r[progress: %s%%]", pos, left, BAR, right, "");
269:         std::cout << "\r";
270:         std::cout << std::flush;
271:     }
272: }
273:
274: // Clear the line, and move cursor to the left
275: void ClearProgress() {
276:     std::cout << "\r";
277: }
278:
279:
280: // djb2 hash function (used for saving unique filenames)
281: unsigned long djb2hash(const std::string &s) {
282:     unsigned long hash = 5381; // Magic number
283:     for (auto c : s) {
284:         hash = (hash << 5) + hash + c; // same as: hash * 33 + c
285:     }
286:     return hash;
287: }
288:
289: // Read CSV file
290: void ReadCSV(const std::string &inputfile, std::vector<std::vector<std::string>> &output) {
291:     std::stringstream file;
292:     file.open(inputfile);
293:     std::string line;
294:
295:     // Read every line from the stream
296:     while (getline(file, line)) {
297:         std::stringstream stream(line);
298:         std::vector<std::string> columns;
299:         std::string element;
300:
301:         // Every line element separated by separator
302:         while (getline(stream, element, ",")) { columns.push_back(element); }
303:         output.push_back(columns);
304:     }
305:     file.close();
306: }
307:
308: // Get JSON file input
309: std::string GetInputData(const std::string &inputfile) {
310:     // Check if exists
311:     if (getchar() != '\n') {
312:         std::string str = "file: " + inputfile;
313:         throw std::invalid_argument(str);
314:     }
315:
316:     // Create a JSON object from file
317:     ifstream ifs(inputfile);
318:     std::string data((std::istreambuf_iterator<char>(ifs), (std::istreambuf_iterator<char>())));
319:
320:     // https://json.org/features:
321:     // We use C++11 raw string literals here R"(text)"
322:     // so it is easier to type in expressions without problems with slashes
323:
324:     // JSONS: <Single and multi-line comments are allowed>
325:     // First this. Remove C style block comments /* */
326:     data = std::regex_replace(data, std::regex(R"(\/\*\/)", ""));
327:     // Second this. Remove C++ style single line comments //
328:     data = std::regex_replace(data, std::regex(R"(\/\//)", ""));
329:
330:     // Remove "\n" or "\t" or "\r", needed for the following regex operations
331:     data = std::regex_replace(data, std::regex(R"(\\n|\\t|\\r)", ""));
332:
333:     // JSONS: <Additional white space characters are allowed>

```

```

416: while (ss.good()) {
417:     std::string substr;
418:     std::getline(ss, substr, delim);
419:
420:     if (trimextraspaces) { TrimExtraSpace(substr); }
421:     output.push_back(substr);
422: }
423: return output;
424: }
425:
426: // Split string to ints
427: std::vector<int> SplitStr2Int(const std::string &input, const char delim) {
428:     std::vector<int> output;
429:     std::stringstream ss(input);
430:
431:     // Get input files by comma
432:     while (ss.good()) {
433:         std::string substr;
434:         std::getline(ss, substr, delim);
435:
436:         TrimExtraSpace(substr);
437:         output.push_back(std::stoi(substr));
438:     }
439:     return output;
440: }
441:
442: // Trim leading, extra and trailing spaces
443: void TrimExtraSpace(std::string &value) {
444:     value = std::regex_replace(value, std::regex(R"(^ +| +)$"), "");
445: }
446:
447: void TrimLeadSpace(std::string &value) {
448:     value = std::regex_replace(value, std::regex(R"(^ +)"), "");
449: }
450:
451: void TrimTrailSpace(std::string &value) {
452:     value = std::regex_replace(value, std::regex(R"( +)$"), "");
453: }
454:
455: void TrimEmptySpace(std::string &value) {
456:     value = std::regex_replace(value, std::regex(R"(^ +| +)$"), "");
457: }
458:
459: void TrimAllSpace(std::string &value) {
460:     value = std::regex_replace(value, std::regex(R"(\\s)", ""));
461: }
462:
463:
464: // Extract words from a string
465: std::vector<std::string> Extract(const std::string &str) {
466:     std::vector<std::string> words;
467:     std::stringstream ss(str);
468:     std::string word;
469:     while (ss >> word) { words.push_back(word); }
470:     return words;
471: }
472:
473: // Check if file exists
474: bool FileExists(const std::string &name) {
475:     struct stat buffer;
476:     return (stat(name.c_str(), &buffer) == 0);
477: }
478:
479:
480: void PrintNotice() {
481:     std::cout << "\n";
482:     << "
483:     << "
484:     << "
485:     << "
486:     << "
487:     << "\n";
488: }
489:
490: void PrintWarning() {
491:     std::cout << "\n";
492:     << "
493:     << "
494:     << "
495:     << "
496:     << "\n";
497: }
498:

```



```

./src/MH1.cc 3/7
167: const double binwidth = (XMAX - XMIN) / XBINS;
168:
169: for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
170:     const double value = std::pow(binwidth * (i + 1) - binwidth / 2.0 * XMIN, power);
171:     const double weight = GetPositiveDefinite(i);
172:     sum += weight * value;
173:     norm += weight;
174: }
175: if (norm > 0) {
176:     return sum / norm;
177: } else {
178:     return 0.0;
179: }
180:
181:
182: template <class T>
183: bool MH1<T>::ValidBin(int idx) const {
184:     if (idx >= 0 && idx < XBINS) { return true; }
185:     return false;
186: }
187:
188: // Unweighted fill
189: template <class T>
190: void MH1<T>::FillAbs(double xvalue) {
191:     // Call the weighted with weight 1.0
192:     Fill(xvalue, 1.0);
193: }
194:
195: // Weighted fill
196: template <class T>
197: void MH1<T>::FillAbs(double xvalue, T weight) {
198:     if (!FILLBUFF) { // Normal filling
199:         fills += 1;
200:     }
201:     // Find out index
202:     const int idx = GetIdx(xvalue, XMIN, XMAX, XBINS, LOOK);
203:
204:     if (idx == -3) { (nanflow += 1); }
205:     if (idx == -1) { (underflow += 1); }
206:     if (idx == -2) { (overflow += 1); }
207:
208:     if (ValidBin(idx)) {
209:         weights[idx] += weight;
210:         weights2[idx] += std::abs(std::conj(weights) * weight);
211:         counts[idx] += 1;
212:     } else { // Autorange initialization
213:         buff_values.push_back(xvalue);
214:         buff_weights.push_back(weight);
215:     }
216:     if (buff_values.size() > static_cast<unsigned int>(AUTOBUFSIZE)) { FlushBuffer(); }
217: }
218:
219:
220: // Automatic histogram range algorithm
221: template <class T>
222: void MH1<T>::FlushBuffer() {
223:     if (FILLBUFF && buff_values.size() > 0) {
224:         FILLBUFF = false; // no more filling buffer
225:     }
226:     // Find out mean
227:     double mu = 0;
228:     double sum = 0;
229:     for (std::size_t i = 0; i < buff_values.size(); ++i) {
230:         mu += std::abs(buff_values[i]) * buff_weights[i];
231:         sum += std::abs(buff_weights[i]);
232:     }
233:     if (sum > 0) { mu /= sum; }
234:
235:     // Variance
236:     double var = 0;
237:     for (std::size_t i = 0; i < buff_values.size(); ++i) {
238:         var += std::abs(buff_weights[i]) * std::pow(buff_values[i] - mu, 2);
239:     }
240:     if (sum > 0) { var /= sum; }
241:
242:     // Minimum and maximum
243:     auto it1 = std::min_element(buff_values.begin(), buff_values.end());
244:     auto it2 = std::max_element(buff_values.begin(), buff_values.end());
245:     const double minval = *it1;
246:     const double maxval = *it2;
247: }
248:
249:

```

```

./src/MH1.cc 5/7
333: for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) { sum += weights2[i]; }
334: return sum;
335: }
336:
337: // Sum the number of bin counts (not the same as fills)
338: template <class T>
339: long long int MH1<T>::SumBinCounts() const {
340:     long long int sum = 0;
341:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) { sum += counts[i]; }
342:     return sum;
343: }
344:
345: // Get maximum histogram bin weight, for complex return |w|^2
346: template <class T>
347: double MH1<T>::GetMaxWeight() const {
348:     double maxval = 0;
349:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
350:         if (GetPositiveDefinite(i) > maxval) { maxval = GetPositiveDefinite(i); }
351:     }
352:     return maxval;
353: }
354:
355: // Get minimum histogram bin weight, for complex return |w|^2
356: template <class T>
357: double MH1<T>::GetMinWeight() const {
358:     double minval = 1e128;
359:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
360:         if (GetPositiveDefinite(i) < minval) { minval = GetPositiveDefinite(i); }
361:     }
362:     return minval;
363: }
364:
365: // Get number of event fills in the bin
366: template <class T>
367: long long int MH1<T>::GetBinCount(int idx) const {
368:     if (ValidBin(idx)) {
369:         return counts[idx];
370:     } else {
371:         return 0;
372:     }
373: }
374:
375: // Get weight of the bin, for complex return complex number
376: template <class T>
377: T MH1<T>::GetBinWeight(int idx) const {
378:     if (ValidBin(idx)) {
379:         return weights[idx];
380:     } else {
381:         return 0.0;
382:     }
383: }
384:
385: // Get weight of the bin, for complex return complex number
386: template <class T>
387: T MH1<T>::GetBinWeight2(int idx) const {
388:     if (ValidBin(idx)) {
389:         return weights2[idx];
390:     } else {
391:         return 0.0;
392:     }
393: }
394:
395: // Error estimate on bin
396: template <class T>
397: double MH1<T>::GetBinError(int idx) const {
398:     if (ValidBin(idx)) { return 0.0; }
399:     double err = 0;
400:     if (GetBinCount(idx) > 0) {
401:         if (isReal()) {
402:             err = std::sqrt(std::abs(GetBinWeight2(idx))); // [sum_i w_i^2]^1/2
403:         } else {
404:             err = std::abs(GetBinWeight2(idx)); // [sum_i w_i^2]^1/2
405:         }
406:     }
407:     return err;
408: }
409:
410: // Return weight (double) or |w|^2 (complex case)
411: template <class T>
412: double MH1<T>::GetPositiveDefinite(int idx) const {
413:     if (ValidBin(idx)) {
414:         if (isReal()) { return std::abs(weights[idx]); }
415:         return std::pow(std::abs(weights[idx]), 2);

```

```

./src/MH1.cc 4/7
250: // Set new histogram bounds
251: const double std = std::sqrt(std::abs(var));
252:
253: double xmin = mu - 2.5 * std;
254: double xmax = mu + 2.5 * std;
255:
256: // A numerical failure may happen with variance calculation, then use this
257: if (std::isnan(xmin) || std::isnan(xmax)) {
258:     xmin = minval;
259:     xmax = maxval;
260: }
261:
262: // If symmetric setup set by user
263: if (AUTOSYMMETRY) {
264:     double val = std::abs(xmin) + std::abs(xmax) / 2.0;
265:     xmin = -val;
266:     xmax = val;
267: }
268:
269: // We have only positive values, such as invariant mass
270: if (minval > 0.0) { xmin = std::max(0.0, xmin); }
271:
272: ResetBounds(XBINS, xmin, xmax);
273:
274: // Fill buffered events
275: for (std::size_t i = 0; i < buff_values.size(); ++i) { Fill(buff_values[i], buff_weights[i]); }
276:
277: // Clear buffers
278: buff_values.clear();
279: buff_weights.clear();
280: }
281:
282: // Reset histogram completely
283: template <class T>
284: void MH1<T>::ResetBounds(int nbins) {
285:     ResetBounds(nbins, 0.0, 0.0);
286:     FILLBUFF = true; // Autorange on, no explicit bounds
287: }
288:
289: // Reset histogram completely
290: template <class T>
291: void MH1<T>::ResetBounds(int nbins, double xmin, double xmax) {
292:     XMIN = xmin;
293:     XMAX = xmax;
294:     XBINS = nbins;
295:
296:     // Init
297:     std::vector<T> null(XBINS, 0.0);
298:     weights = null;
299:     weights2 = null;
300:     counts = std::vector<long long int>(XBINS, 0);
301:
302:     Clear(); // Call also this!
303:     FILLBUFF = false; // No autorange, explicit bounds provided
304: }
305:
306: // Clear the histogram data but keep the bounds
307: template <class T>
308: void MH1<T>::Clear() {
309:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) {
310:         weights[i] = 0.0;
311:         weights2[i] = 0.0;
312:         counts[i] = 0;
313:         fills = 0;
314:         underflow = 0;
315:         overflow = 0;
316:         nanflow = 0;
317:     }
318: }
319:
320: // Sum over all histogram bin weights
321: template <class T>
322: T MH1<T>::SumWeights() const {
323:     T sum = 0.0;
324:     for (std::size_t i = 0; i < static_cast<unsigned int>(XBINS); ++i) { sum += weights[i]; }
325:     return sum;
326: }
327:
328: // Sum over all histogram bin weights squared
329: template <class T>
330: T MH1<T>::SumWeights2() const {
331:     T sum = 0.0;
332: }
333:
334: } else {
335:     return 0.0;
336: }
337:
338: // Get bin index (idx) corresponding to value (xvalue)
339: template <class T>
340: void MH1<T>::GetBinIdx(double xvalue, int idx) {
341:     // Find out bins
342:     idx = GetIdx(xvalue, XMIN, XMAX, XBINS, LOOK);
343: }
344:
345: // Get bin value in units of X for a given bin index
346: // boundary = -1, 0, 1 (lower, center, upper)
347: template <class T>
348: double MH1<T>::GetBinValue(int idx, int boundary) const {
349:     if (idx > XBINS - 1) {
350:         throw std::invalid_argument("MH1::GetBinValue: idx = " + std::to_string(idx) +
351:                                     " > XBINS = " + std::to_string(XBINS) + "!");
352:     }
353:     if (LOOK) {
354:         const double log10step = (std::log10(XMAX) - std::log10(XMIN)) / XBINS;
355:         if (boundary == -1) {
356:             return std::pow(10, std::log10(XMIN) + idx * log10step);
357:         } else if (boundary == 0) {
358:             return std::pow(10, std::log10(XMIN) + idx * log10step +
359:                             std::pow(10, std::log10(XMIN) + (idx + 1) * log10step) / 2);
360:         } else if (boundary == 1) {
361:             return std::pow(10, std::log10(XMIN) + (idx + 1) * log10step);
362:         } else {
363:             throw std::invalid_argument("MH1::GetBinValue: Bin boundary not valid (-1,0,1)");
364:         }
365:     } else {
366:         const double binwidth = (XMAX - XMIN) / XBINS;
367:         const double value = XMIN + (idx + 1) * binwidth;
368:         if (boundary == -1) {
369:             return value - binwidth;
370:         } else if (boundary == 0) {
371:             return value - binwidth / 2.0;
372:         } else if (boundary == 1) {
373:             return value;
374:         } else {
375:             throw std::invalid_argument("MH1::GetBinValue: Bin boundary not valid (-1,0,1)");
376:         }
377:     }
378: }
379:
380: // Get table/histogram index for linearly or base-10 logarithmically spaced
381: // bins
382: // Gives exact uniform filling within bin boundaries.
383: // In the logarithmic case, MINVAL and MAXVAL > 0, naturally.
384: // Underflow returns -1
385: // Overflow returns -2
386: int MH1<T>::GetIdx(double value, double minval, double maxval, int nbins, bool logbins) const {
387:     if (std::isnan(value) || std::isinf(value)) { return -3; }
388:     if (value < minval) { return -1; } // underflow
389:     if (value > maxval) { return -2; } // overflow
390: }
391:
392: int idx = 0;
393: // Logarithmic binning
394: if (logbins) {
395:     // Check do we have non-negative input
396:     if (value > 0) { std::floor(nbins * (std::log10(value) - std::log10(minval)) /
397:                                     (std::log10(maxval) - std::log10(minval)))
398:     } else {
399:         // Linear binning
400:         const double BINWIDTH = (maxval - minval) / nbins;
401:         idx = std::floor((value - minval) / BINWIDTH);
402:     }
403: }
404:
405: // Instantiate (necessary for compilation)
406: template class MH1<double>;
407: template class MH1<std::complex<double>>;
408:
409: // namespace gra

```

```

499:
1: // KISS fragmentation class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT license <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <complex>
8: #include <random>
9: #include <valarray>
10: #include <vector>
11:
12: // Own
13: #include "Granitti/MFVec.h"
14: #include "Granitti/MBox.h"
15: #include "Granitti/MForm.h"
16: #include "Granitti/MFragment.h"
17: #include "Granitti/MKinematics.h"
18: #include "Granitti/MMath.h"
19: #include "Granitti/MMatrix.h"
20: #include "Granitti/MPDG.h"
21: #include "Granitti/MRandom.h"
22:
23: // Libraries
24: #include "rang.hpp"
25:
26: using gra::max;indices;
27: using gra::math::magr;
28: using gra::math::pow2;
29:
30: using gra::PDG::imp;
31: using gra::PDG::impi;
32:
33: using gra::PDG::PDG_gamma;
34: using gra::PDG::PDG_n;
35: using gra::PDG::PDG_p;
36: using gra::PDG::PDG_pi0;
37: using gra::PDG::PDG_pi_m;
38: using gra::PDG::PDG_pi_p;
39:
40: namespace gra {
41:
42: // Return decay status [RESERVATION!]
43: //
44: //
45: void MFragment::GetDecayStatus(const std::vector<int> spdgcode, std::vector<bool> isstable) {
46:     isstable.resize(spdgcode.size(), true);
47:     for (const auto i : indices(isstable)) {
48:         isstable[i] = (spdgcode[i] == 111 || spdgcode[i] == 311) ? false : true;
49:     }
50: }
51:
52: // Sample from Hagedorn type parametrization
53: //
54: // This function is useful to generate distribution with
55: // exponential at low pt, powerlaw at high pt.
56: // Does not generate jet like topologies,
57: // so this can be used as a strict null model (H0) for soft like topologies.
58: //
59: // [slow function, performance should be upgraded]
60: //
61: // Parameters <=>
62: // p0 = T/(q-1)
63: // n = 1/(q-1)
64: //
65: void MFragment::ExpPowWMD(double q, double T, double mact, const std::vector<double> fmass,
66:                            const unsigned int MAXTRIAL = 1e4, // Safety brack
67:                            const unsigned int Nbins = 1e4;
68:                            const double pstep = mact / Nbins;
69:                            //
70:                            // Calculate pt-bin values
71:                            std::vector<double> ptval(Nbins);
72:                            for (std::size_t i = 0; i < Nbins; ++i) ptval[i] = i * pstep; }
73: //
74: // Random integer from [0, Nbins-1]
75: // [C++], thread_local is also static
76: // Thread_local std::uniform_int_distribution<int> RANDI(0, Nbins - 1);
77: //
78: // Calculate for pion, kaon, proton masses
79: const std::vector<double> fmass = {PDG::mpi0, PDG::mpi, PDG::mk, PDG::mp};
80: std::vector<double> dsdpt(fmass.size(), std::vector<double>(Nbins, 0.0));
81: std::vector<double>

```

```

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30: using gra::PDG::imp;
31: using gra::PDG::impi;
32:
33: using gra::PDG::PDG_gamma;
34: using gra::PDG::PDG_n;
35: using gra::PDG::PDG_p;
36: using gra::PDG::PDG_pi0;
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47:     for (const auto i : indices(isstable)) {
48:         isstable[i] = (spdgcode[i] == 111 || spdgcode[i] == 311) ? false : true;
49:     }
50: }
51:
52: // Sample from Hagedorn type parametrization
53: //
54: // This function is useful to generate distribution with
55: // exponential at low pt, powerlaw at high pt.
56: // Does not generate jet like topologies,
57: // so this can be used as a strict null model (H0) for soft like topologies.
58: //
59: // [slow function, performance should be upgraded]
60: //
61: // Parameters <=>
62: // p0 = T/(q-1)
63: // n = 1/(q-1)
64: //
65: void MFragment::ExpPowWMD(double q, double T, double mact, const std::vector<double> fmass,
66:                            const unsigned int MAXTRIAL = 1e4, // Safety brack
67:                            const unsigned int Nbins = 1e4;
68:                            const double pstep = mact / Nbins;
69:                            //
70:                            // Calculate pt-bin values
71:                            std::vector<double> ptval(Nbins);
72:                            for (std::size_t i = 0; i < Nbins; ++i) ptval[i] = i * pstep; }
73: //
74: // Random integer from [0, Nbins-1]
75: // [C++], thread_local is also static
76: // Thread_local std::uniform_int_distribution<int> RANDI(0, Nbins - 1);
77: //
78: // Calculate for pion, kaon, proton masses
79: const std::vector<double> fmass = {PDG::mpi0, PDG::mpi, PDG::mk, PDG::mp};
80: std::vector<double> dsdpt(fmass.size(), std::vector<double>(Nbins, 0.0));
81: std::vector<double>

```

```

84: // -----
85: // Very slow
86:
87: // Pre-evaluate dsigma/dpt [L_y=0 (mid rapidity)]
88: for (const auto k : indices(fixmass)) {
89:     const double m2 = pow2(fixmass[k]);
90:     for (std::size_t i = 0; i < Nbins; ++i) {
91:         const double mval = msqrt(pow2(ptval[i]) + m2);
92:
93:         // PD
94:         dsdpt[k][i] = ptval[i] * mval * std::pow(1.0 + (q - 1.0) * mval / T, q / (1.0 - q));
95:
96:         // Save maximum
97:         maxval[k] = (dsdpt[k][i] > maxval[k]) ? dsdpt[k][i] : maxval[k];
98:     }
99: // -----
100:
101: // For each particle, generate pt
102: x.resize(mass.size(), 0.0);
103: for (const auto sp : indices(mass)) {
104:     // find the best distribution
105:     unsigned int best = 0;
106:     double minmist = 1e32;
107:     for (std::size_t k = 0; k < fixmass.size(); ++k) {
108:         const double dist = std::abs(fixmass[k] - mass[sp]);
109:         if (dist < minmist) {
110:             minmist = dist;
111:             best = k;
112:         }
113:     }
114:
115:     // Acceptance-Rejection, very slow
116:     unsigned int trials = 0;
117:     while (true) {
118:         const int BIN = RANDI(rng.rng);
119:         if (rng.D(0.0, 1.0) * maxval[best] < dsdpt[best][BIN]) {
120:             x[sp] = ptval[BIN];
121:             break;
122:         }
123:         ++trials;
124:         if (trials > MAXTRIAL) {
125:             break; // failed
126:         }
127:     }
128: }
129: }
130:
131: // N-body fragmentation with tube (cylinder) phase space
132: // (approximate target distribution is flat over rapidity, pt from a given
133: // distribution)
134: //
135: // [REFERENCE: Jadach, Computer Physics Communications, 9 (1975) 297-304]
136: // [REFERENCE: UG5, http://cds.cern.ch/record/17907/files/19870120.pdf]
137: //
138: // This function contains a mixture of dynamics and kinematics, i.e.,
139: // is not a pure phase space and is suitable for soft fragmentation studies.
140: //
141: // Return: 1.0 for a valid fragmentation and -1.0 for a kinematically impossible
142: //
143: double MFragment::TubeFragment(const MVec &mother, double M0, const std::vector<double> fmass,
144:                                double mact, MRandom &rng, const std::string &pt_distribution) {
145:
146:     // Number of criticals in a case of failing kinematics
147:     const unsigned int MAXTRIAL = 30;
148:     unsigned int trials = 0;
149:
150:     const unsigned int N = m.size();
151:     std::valarray<double> m2(N);
152:     for (const auto i : indices(m)) m2[i] = pow2(m[i]);
153:     // transverse momentum px, py
154:     std::valarray<double> px(N);
155:     std::valarray<double> py(N);
156:     std::valarray<double> pt2(N);
157:
158:     // rapidity and transverse mass
159:     std::valarray<double> y(N);
160:     std::valarray<double> mt(N);
161:
162:     while (true) {
163:         // Random sample pxpy and initial rapidity
164:         std::vector<double> pt(N);
165:
166:         if (pt_distribution == "powexp") {

```

```

167:         ExpPowWMD(q, T, mact, m, pt, rng);
168:     } else if (pt_distribution == "exp") {
169:         for (const auto i : indices(pt)) {
170:             pt2[i] = rng.ExpRandom(lambda);
171:             pt[i] = msqrt(pt2[i]);
172:         }
173:     } else {
174:         throw std::invalid_argument("MFragment::TubeFragment: Unknown pt-distribution parameter = " +
175:                                     pt_distribution);
176:     }
177:
178:     for (const auto i : indices(pt)) {
179:         // Sample angles, px, py
180:         const double phi = rng.D(0, 2.0 * gra::math::PI);
181:         px[i] = pt[i] * std::cos(phi);
182:         py[i] = pt[i] * std::sin(phi);
183:     }
184:
185:     // Rapidity
186:     y[i] = rng.D(0.0, 1.0);
187:     std::sort(begin(y), end(y)); // Rapidity ordering
188:
189:     // -----
190:     // Scale rapidities to span the full range [0,1] (min=0, max=1)
191:     y = (y - y[0]) / (y[N - 1] - y[0]);
192:
193:     // Set transverse momentum with zero sum
194:     px = px - px.sum() / N;
195:     py = py - py.sum() / N;
196:     pt2 = px * px + py * py;
197:
198:     // Transverse mass for all particles
199:     mt = sqrt(m2 + pt2);
200:
201:     // -----
202:     // Solve rapidity scale factor 'alpha'
203:     double alpha = 0.0;
204:     if (ISolveAlpha(alpha, M0, m, mt, y)) {
205:         ++trials;
206:         if (trials > MAXTRIAL) {
207:             return -1.0;
208:         }
209:         continue;
210:     }
211: }
212: // Scale all rapidities
213: y = alpha * y;
214:
215: // Set longitudinal momentum and energy using 'alpha'
216: // get boost factor to +- 0
217: const double Q = (m * exp(y)).sum();
218: y = y + std::log(M0 / Q);
219: const std::valarray<double> pz = m * sin(y);
220: const std::valarray<double> E = mt * cosh(y);
221:
222: // -----
223: // Set all particles 4-momenta and boost to the original frame
224: p.resize(N);
225: const int sign = 1; // positive
226: for (const auto i : indices(pt)) {
227:     p[i] = MVec(pz[i], py[i], px[i], E[i]);
228:     gra::Kinematics::torentBoost(mother, M0, p[i], sign);
229: }
230:
231: // Check EM-conservation
232: MVec p_sum(0, 0, 0, 0);
233: for (const auto n : p) p_sum += n;
234: bool valid = gra::math::CheckKMC(p_sum - mother);
235:
236: // Check rapidity (floating points can fail after boost in forward)
237: for (const auto i : indices(pt)) {
238:     if (std::isnan(p[i].Rap()) || std::isinf(p[i].Rap())) {
239:         valid = false;
240:         break;
241:     }
242: }
243: if (!valid) {
244:     ++trials;
245:     if (trials > MAXTRIAL) {
246:         return -1.0;
247:     } else {
248:         continue;
249:     }

```

```

./src/MFragment.cc 4/7
250: }
251: break; // Fragmentation successful!
252: }
253: return 1.0;
254: }
255: }
256: // -----
257: // Solve rapidly scale factor via Newton iteration
258: // (alternative strategies are viable, too, this does not converge
259: // always with non-gaussian pt-distributions)
260: //
261: // a_1 = a_0 - f(a_0)/f'(a_1), find root a such that f(a) = 0
262: // iteratively via
263: // a_{forj} = a_{j-1} - f(a_{j-1}) / f'(a_{j-1})
264: //
265: bool MFragment::SolveAlpha(double alpha, double MO, const std::vector<double>& em,
266: const std::valarray<double>& smt, const std::valarray<double>& ky) {
267: const double STOP_EPS = 1e-8;
268: const unsigned int MAXITER = 20;
269:
270: // Starting value
271: const int N = m.size();
272: const double C = std::log(MO * MO);
273: alpha = C - std::log(10) * m[N - 1];
274: std::vector<double> E = {0, 0, 0, 0};
275:
276: unsigned int iter = 0;
277: while (true) {
278: const std::valarray<double> x = exp(alpha * y);
279:
280: E[0] = (m * x).sum();
281: E[1] = (m * x).sum();
282: E[2] = (m * y * x).sum(); // Derivative d/dy
283: E[3] = (m * y * x).sum(); // Derivative d^2/dy^2
284:
285: // Iterate the solution, d/dx ln(x) = 1/x
286: const double DY = E[0] * E[1] * (C - std::log(E[0] * E[1])) / (E[0] * E[3] - E[1] * E[2]);
287: alpha = DY;
288:
289: if (std::abs(N * DY / alpha) < STOP_EPS) { return true; }
290: if (std::isnan(alpha) || (++iter) > MAXITER) { return false; }
291: }
292: }
293:
294: // N* decay table [set manually according to experimental data]
295: //
296: // 0 is the proton / antiproton charge (1,-1)
297: // MO is the M mass
298: //
299: void MFragment::MetaDecayTable(int Q, double MO, std::vector<int>& pdgcodes, MRandom &rng) {
300: int decaymode = 0;
301:
302: // Only 2-body decay possible, mass below 3-body threshold
303: if (MO < (PDG::mp + 2 * PDG::mpi)) {
304: decaymode = 0;
305: }
306: // 2- or 3-body decay possible
307: else {
308: // C++11, thread_local is also static
309: // 2->body / 3-body branching ratios from POC
310: thread_local std::discrete_distribution<> d1(0.60, 0.40);
311: decaymode = d1(rng, rng); // Draw random
312: }
313: }
314: // -----
315: // *(PT) = 1/2*pi Decay Parameters from POC
316: // Only major decays implemented
317:
318: std::vector<int> decays;
319:
320: // 2-body channel
321: if (decaymode == 0) {
322: // subchannels
323: // C++11, thread_local is also static
324: thread_local std::discrete_distribution<> subd1
325: { 2.0 / 3.0, 1.0 / 3.0 }; // Branching ratios from Clebsch-Gordan
326: const int channel = subd1(rng, rng);
327:
328: // PDG-ID of subchannels
329: const std::vector<std::vector<int>> ID = {{Q * PDG::n, Q * PDG::pip}, // (anti)neutron & pi(-)+
330: {Q * PDG::p, PDG::pi0}}; // (anti)proton & pi0
331:
332: // Choose the decay channel

```

```

./src/MFragment.cc 6/7
416: if (mX < (PDG::mp + PDG::mpi + safe_margin)) { // && mX < lts.sqrts_s()
417: mass = mX;
418: return;
419: }
420: }
421: }
422: }
423: // Simple statistical toy particle pick-up, nothing more
424: //
425: bool MFragment::PickParticles(MRandom M, unsigned int N, int B, int S, int Q,
426: const std::vector<double>& smas, std::vector<int>& pdgcodes, const MPOG &POG,
427: const unsigned int MAXTRIAL = 1e4,
428: MRandom &rng) {
429:
430: // C++11, thread_local is also static
431: thread_local std::uniform_int_distribution<int> RANDI3(0, 2);
432:
433: // Pion, Kaon, Proton ratios (MGAURUM)
434: const double denom3 = (1 + 0.1 + 0.06);
435: const std::vector<double> ratio3 = {1.0 / denom3, 0.1 / denom3, 0.06 / denom3};
436:
437: const std::vector<int> charged_pdg = {211, 321, 2212}; // pi+, K+, (anti)proton
438: const std::vector<int> charged_B = {0, 0, 1};
439: const std::vector<int> charged_S = {0, 1, 0};
440:
441: const std::vector<int> neutral_pdg = {111, 311, 2112}; // pi0, K0, neutron
442: const std::vector<int> neutral_B = {0, 0, 1};
443: const std::vector<int> neutral_S = {0, 1, 0};
444:
445: const double QProb = 2.0 / 3.0; // Charged probability
446: const double DOUBLE = 0.9; // Prob to pick particle pair
447: // -----
448: // Now resize!
449: mas.resize(N, 0);
450: pdgcode.resize(N, 0);
451: mas.resize(N, 0);
452: std::vector<int> Qcharges(N, 0);
453: std::vector<int> Bcharges(N, 0);
454: std::vector<int> Scharges(N, 0);
455: unsigned int trials = 0;
456:
457: while (true) {
458: unsigned int i = 0;
459:
460: do { // hadron picking
461:
462: // Charged
463: if (rng.U(0, 1) < QProb) {
464: // Charged pair
465: if (rng.U(0, 1) < DOUBLE && i <= N - 2) {
466: // Sample particle flavour
467: while (true) {
468: const int bin = RANDI3(rng, rng);
469:
470: if (rng.U(0, 1) < ratio3[bin]) {
471: Qcharges[i] = 1;
472: Bcharges[i] = charged_B[bin];
473: Scharges[i] = charged_S[bin];
474: pdgcode[i] = charged_pdg[bin];
475: ++i; // Next slot
476:
477: Qcharges[i] = -1;
478: Bcharges[i] = -charged_B[bin];
479: Scharges[i] = -charged_S[bin];
480: pdgcode[i] = -charged_pdg[bin];
481: ++i;
482: break;
483: }
484: }
485: // Single
486: } else {
487: while (true) {
488: const int bin = RANDI3(rng, rng);
489:
490: if (rng.U(0, 1) < ratio3[bin]) {
491: int sign = 0;
492: sign = (rng.U(0, 1) < 0.5) ? 1 : -1;
493:
494: Qcharges[i] = sign;
495: Bcharges[i] = sign * charged_B[bin];
496: Scharges[i] = sign * charged_S[bin];
497: pdgcode[i] = sign * charged_pdg[bin];

```

```

./src/MFragment.cc 5/7
333: pdgcodes = ID[channel];
334: }
335: }
336: // 3-body channel
337: if (decaymode == 1) {
338: // Subchannel
339: // C++11, thread_local is also static
340: thread_local std::discrete_distribution<> subd1
341: { 1.0 / 3.0, 1.0 / 3.0, 1.0 / 3.0 }; // Branching ratios ansatz
342: const int channel = subd1(rng, rng);
343:
344: // PDG-ID of subchannels
345: const std::vector<std::vector<int>> ID = {{Q * PDG::n, Q * PDG::pip, PDG::pi0},
346: {Q * PDG::p, PDG::pip, PDG::pi0},
347: {Q * PDG::p, PDG::pi0, PDG::pi0}};
348:
349: // Perhaps to add
350: // {Q*PDG::delat0, Q*PDG::pip, PDG::pi0}, // delat0 & pi+ & pi0
351: // {Q*PDG::delatp, PDG::pip, PDG::pi0}, // delat+ & pi+ & pi-
352: // {Q*PDG::delatp, PDG::pi0, PDG::pi0}, // delat+ & pi0 & pi0
353:
354: // Draw the decay channel
355: pdgcodes = ID[channel];
356: }
357: }
358:
359: // Return excited forward proton masses
360: void MFragment::GetForwardMass(double smas1, double smas2, bool lexcite1, bool excite2,
361: unsigned int excite, MRandom &rng) {
362: if (excite == 0) { // Fully elastic
363: mass1 = PDG::mp;
364: mass2 = PDG::mp;
365: excite1 = false;
366: excite2 = false;
367: } else if (excite == 1) { // Single excitation
368: if (rng.U(0, 1) < 0.5) {
369: mass1 = PDG::mp;
370: GetSingleForwardMass(mas2, rng);
371: excite1 = false;
372: excite2 = true;
373: } else {
374: GetSingleForwardMass(mas1, rng);
375: mass2 = PDG::mp;
376: excite1 = true;
377: excite2 = false;
378: }
379: } else if (excite == 2) { // Double excitation
380: GetSingleForwardMass(mas1, rng);
381: GetSingleForwardMass(mas2, rng);
382: excite1 = true;
383: excite2 = true;
384: }
385: }
386:
387: // Return excited forward system mass
388: void MFragment::GetSingleForwardMass(double smas, MRandom &rng) {
389: // Good-Walker resonance excitation probabilities
390: // C++11, thread_local is also static
391: thread_local std::discrete_distribution<> d
392: { PARAM_NSTAR::rc[0], PARAM_NSTAR::rc[1], PARAM_NSTAR::rc[2] };
393: const int state = d(rng, rng); // Draw random
394:
395: // Excited proton states
396: double MO = 0;
397: double width = 0.0;
398:
399: if (state == 0) {
400: MO = 1.440; // **
401: width = 0.325;
402: } else if (state == 1) {
403: MO = 1.680; // **
404: width = 0.140;
405: } else if (state == 2) {
406: MO = 2.190; // ***
407: width = 0.450;
408: }
409: }
410:
411: // Draw the excited state mass
412: const double safe_margin = 0.01;
413: double mX = 0;
414: while (true) {
415: mX = random.RelativisticEMRandom(MO, width, 1e6);

```

```

./src/MFragment.cc 7/7
499: ++i;
500: break;
501: }
502: }
503: }
504: }
505: // Neutral
506: } else {
507: while (true) {
508: // Sample particle flavour
509: const int bin = RANDI3(rng, rng);
510: if (rng.U(0, 1) < ratio3[bin]) {
511: Qcharges[i] = 0;
512: Bcharges[i] = neutral_B[bin];
513: Scharges[i] = neutral_S[bin];
514: pdgcode[i] = neutral_pdg[bin];
515: ++i;
516: break;
517: }
518: }
519: }
520: } while (i < N); // Picking loop
521:
522: // Sum charges
523: const int Q_sum = std::accumulate(Qcharges.begin(), Qcharges.end(), 0);
524: const int B_sum = std::accumulate(Bcharges.begin(), Bcharges.end(), 0);
525: const int S_sum = std::accumulate(Scharges.begin(), Scharges.end(), 0);
526:
527: // Get corresponding masses
528: for (std::size_t i = 0; i < N; ++i) {
529: // print("PDGcode[+id] = %d [Nsid] %s", i, pdgcode[i], N);
530: mas[i] = PDG::findByPDG(pdgcode[i], mass);
531:
532: const double M_sum = std::accumulate(mas.begin(), mas.end(), 0.0);
533:
534: // Check charge and mass threshold
535: bool Q_check = (Q_sum == 0) ? true : false;
536: bool B_check = (B_sum == 0) ? true : false;
537: bool S_check = (S_sum == 0) ? true : false;
538: bool M_check = (M_sum < MO) ? true : false;
539:
540: if (Q_check && B_check && S_check && M_check) {
541: break; // We are ok!
542: } else {
543: // print("PTICK Q_sum = %d, B_sum = %d, M_sum = %0.1f, N = %d
544: // %s", Q_sum, B_sum, M_sum, M, N);
545: ++trials;
546: if (trials > MAXTRIAL) { return false; }
547: }
548: } // Out loop
549:
550: return true;
551: }
552: }
553: // namespace gra

```

```

./src/MUserCuts.cc          1/3
1: // Custom user cuts which cannot be implemented directly via json steering files
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: //
6: // C++
7: #include <complex>
8: #include <iostream>
9: #include <vector>
10:
11: // Own
12: #include "Granititi/MForm.h"
13: #include "Granititi/MGLematics.h"
14: #include "Granititi/Mmath.h"
15: #include "Granititi/MUserCuts.h"
16:
17: namespace gra {
18: // USERCUTS (implement custom cuts here); cuts which cannot be implemented
19: // in .json steering files; label these by unique integers.
20:
21: bool UserCut(int id, const gra::LORNETSCALAR &its) {
22: // ** NO CUTS CASE, THIS SHOULD BE FIRST **
23: if (id == 0) {
24: return true;
25: }
26:
27: // -----
28: // "Spin-filter" cut ("Globeball" filter)
29:
30: // Forward proton |dpt| < 0.3 GeV
31: else if (id == -3) {
32: const double dpt = (its.pfinal[1] - its.pfinal[2]).Pt();
33: if (dpt < 0.3) {
34: // fine
35: } else {
36: return false; // did not pass
37: }
38: }
39:
40: // Forward proton |dpt| > 0.3 GeV
41: else if (id == 3) {
42: const double dpt = (its.pfinal[1] - its.pfinal[2]).Pt();
43: if (dpt > 0.3) {
44: // fine
45: } else {
46: return false; // did not pass
47: }
48: }
49:
50: // -----
51: // "Spin-filter" cut
52:
53: // Forward proton pT1 dot pT2 < 0
54: else if (id == -11) {
55: if (its.pfinal[1].DotPt(its.pfinal[2]) < 0) {
56: // fine
57: } else {
58: return false; // did not pass
59: }
60: }
61:
62: // Forward proton pT1 dot pT2 > 0
63: else if (id == 11) {
64: if (its.pfinal[1].DotPt(its.pfinal[2]) > 0) {
65: // fine
66: } else {
67: return false; // did not pass
68: }
69: }
70:
71: // -----
72: // "Spin-filter" cut
73:
74: // Forward proton |deltaphi| in (90, 180)
75: else if (id == 90180) {
76: const double deltaphiabs = its.pfinal[1].DeltaPhiAbs(its.pfinal[2]);
77: if (gra::math::DegRad(90) < deltaphiabs && deltaphiabs <= gra::math::DegRad(180)) {
78: // fine
79: } else {
80: return false; // did not pass
81: }
82: }
83: }

```

```

./src/MUserCuts.cc          3/3
167: }
168: } else if (30 <= M && M <= 70) { // GeV
169: if (its.decaytree[0].p4.Pt() > 10 && its.decaytree[1].p4.Pt() > 10) {
170: // fine
171: } else {
172: return false; // not passed
173: }
174: }
175: }
176:
177: // -----
178: // ATLAS pi pi - 13 TeV roman pot fiducial cuts
179: // (rather cuts needed in .json file)
180: // from R. Sikora, ATLAS pointer, Rad Homenf QCD School 2017
181: //
182: // (N.B. check the implementation
183: // c.f. cut |c| > 0.03 GeV^2 seems to give more physical results)
184:
185: else if (id == 1230123) {
186: // Forward protons |y1| and |phi|
187: std::vector<double> pyabs = {std::abs(its.pfinal[1].Py()), std::abs(its.pfinal[2].Py())};
188: std::vector<double> phiabs = {std::abs(its.pfinal[1].Phi()), std::abs(its.pfinal[2].Phi())};
189: for (std::size_t i = 0; i < 2; ++i) {
190: if ((0.17 < pyabs[i]) && (pyabs[i] < 0.5)) { // GeV
191: // fine
192: } else {
193: return false; // not passed
194: }
195: if ((gra::math::PI / 4 < phiabs[i]) && (phiabs[i] < 3.0 * gra::math::PI / 4)) {
196: // fine
197: } else {
198: return false; // not passed
199: }
200: }
201: }
202:
203: // Roman pot geometry
204: const double deltaphiabs = its.pfinal[1].DeltaPhiAbs(its.pfinal[2]);
205: if (deltaphiabs < gra::math::DegRad(40.0)) || (deltaphiabs > gra::math::DegRad(140.0)) {
206: // fine
207: } else {
208: return false; // not passed
209: }
210: } else { // Poor input
211: std::string str = "MUserCuts: Unknown cut ID: " + std::to_string(id);
212: throw std::invalid_argument(str);
213: }
214: return true;
215: }
216:
217: // namespace gra
218:

```

```

./src/MUserCuts.cc          2/3
84:
85: // Forward proton |deltaphi| in (0, 90)
86: else if (id == 90) {
87: const double deltaphiabs = its.pfinal[1].DeltaPhiAbs(its.pfinal[2]);
88: if (0 < deltaphiabs && deltaphiabs <= gra::math::DegRad(90)) {
89: // fine
90: } else {
91: return false; // did not pass
92: }
93: }
94:
95: // -----
96: // ATLAS/BKFT |logr(t)| = 200 GeV |pT| - to other cuts needed in .json file
97: // Indico.cern.ch/event/713101/contributions/310231/
98: // attachments/1705771/2748440/Difraction2018_RafalSikora.pdf
99:
100: else if (id == 280818) {
101: // Loop over forward protons
102: std::vector<int> indices = {1, 2};
103:
104: for (const auto &i : indices) {
105: if (gra::math::pow2(its.pfinal[i].Px() + 0.3) + gra::math::pow2(its.pfinal[i].Py()) <
106: 0.25) { // GeV^2
107: // fine
108: } else {
109: return false; // not passed
110: }
111:
112: if (0.2 < std::abs(its.pfinal[i].Py()) && std::abs(its.pfinal[i].Py()) < 0.4) { // GeV
113: // fine
114: } else {
115: return false; // not passed
116: }
117:
118: if (its.pfinal[i].Px() > -0.2) { // GeV
119: // fine
120: } else {
121: return false; // not passed
122: }
123: }
124:
125: // [arxiv.org/abs/1608.03765]
126:
127: } else if (id == 160803765) {
128: const double x1 = (its.pbeam.Pz() - its.pfinal[1].Pz()) / its.pbeam.Pz();
129: const double x12 = (its.pbeam.Pz() - its.pfinal[2].Pz()) / its.pbeam.Pz();
130: const double X1_MAX = 0.03;
131:
132: if (x1 < X1_MAX && x12 < X1_MAX) {
133: // fine
134: } else {
135: return false; // not passed
136: }
137: }
138: }
139:
140: // -----
141: // CDF exclusive dijets
142: // [arxiv.org/abs/0712.0604]
143:
144: else if (id == 7120604) {
145: // anti-proton longitudinal momentum loss fraction
146: const double xi_pbar = (its.pbeam.Pz() - its.pfinal[2].Pz()) / its.pbeam.Pz();
147:
148: if (0.03 < xi_pbar && xi_pbar < 0.08) {
149: // fine
150: } else {
151: return false; // not passed
152: }
153: }
154:
155: // -----
156: // ATLAS yy->mu+mu- 13 TeV fiducial cuts (rather cuts needed in .json file)
157: // [arxiv.org/abs/hep-ex/1706053]
158:
159: else if (id == 170805053) {
160: const double M = gra::math::sqrt(its.m2);
161:
162: if (12 <= M && M < 30) { // GeV
163: if (its.decaytree[0].p4.Pt() > 6 && its.decaytree[1].p4.Pt() > 6) {
164: // fine
165: } else {
166: return false; // not passed
167: }
168: }
169: }

```

```

./src/MSpherical.cc          1/8
1: // Functional methods for Spherical Harmonic Expansions
2: //
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: //
6: // C++
7: #include <complex>
8: #include <iostream>
9: #include <vector>
10:
11: // Own
12: #include "Granititi/MUser.h"
13: #include "Granititi/Mmath.h"
14: #include "Granititi/MMatrix.h"
15: #include "Granititi/MRandom.h"
16: #include "Granititi/MSpherical.h"
17:
18: // Libraries
19: #include "rang.hpp"
20:
21: using gra::max_indices;
22: using gra::math::PI;
23: using gra::math::sqrt;
24: using gra::math::pow2;
25: using gra::math::xi;
26:
27: namespace gra {
28: namespace spherical {
29: // Monte Carlo integral I = \int_{\Omega} f(x) dx, where f(x) = V <f(x)>
30: // True integral being I = \int_{\Omega} f(x) dx / V
31:
32: // This inner-product (overlap integral) matrix is identity if the phase space
33: // is flat (uniform). In a geometrically restricted phase space, these basis
34: // functions get mixed => lost orthogonality.
35:
36: MMatrix<double> GetGMixing(const std::vector<Omega> &events, const std::vector<std::size_t> &ind,
37: int l_max, const std::string &mode) {
38: const int NCOEF = (l_max + 1) * (l_max + 1);
39:
40: std::cout << "GetGMixing: mode = " << mode << std::endl;
41: std::cout << "Generated flat MC phase space events = " << ind.size() << std::endl;
42:
43: // Construct the efficiency coefficients EPSILON_LM with linear
44: // indexing
45: MMatrix<double> E(NCOEF, NCOEF, 0.0);
46: MMatrix<double> E2(NCOEF, NCOEF, 0.0); // for the uncertainty
47:
48: // Loop over GENERATED MC events and do the integral in effect via
49: // uniform MC sampling
50: // We evaluate the integral:
51: //
52: // \int_{\Omega} f(x) dx = \int_{\Omega} f(x) d\Omega / V
53: // ind: fiducial = 0
54: // ind: selected = 1
55:
56: const double VOL = 4.0 * PI; // [cos(theta) x phi] plane area
57:
58: for (const auto &k : ind) {
59: const bool fid = events[k].fiducial;
60: const bool sel = events[k].selected;
61:
62: // Flat phase space
63: if (mode == "flat") {
64: // all fine
65: }
66: // Geometric acceptance ("fiducial phase space")
67: else if (mode == "fid") {
68: if (fid) {
69: continue;
70: } else {
71: //fiducial;
72: }
73: }
74: // Geometric Efficiency ("detector level")
75: else if (mode == "det") {
76: if (fid) {
77: continue;
78: } else {
79: //fiducial;
80: }
81: if (sel) {
82: continue;
83: } else {

```



```

./src/MSpherical.cc          6/8
416: return sum;
417: }
418:
419: // TODO: add correlations in to the sum
420: double HamDotProdError() { return 0.0; }
421:
422: void PrintOutMoments(const std::vector<double> &ix, const std::vector<double> &ix_err,
423:                    const std::vector<bool> &ACTIVE, int IMAX) {
424:     for (int l = 0; l <= IMAX; ++l) {
425:         for (int m = -l; m <= l; ++m) {
426:             const unsigned int index = LinearInd(l, m);
427:             printf("M%d%d Vc = %8.1e +/- %8.1e ", l, m, ix[index], ix_err[index]);
428:
429:             if (ACTIVE[index]) {
430:                 std::cout << "[ " << rang::fg::green << "active" << rang::fg::reset << " ]" << std::endl;
431:             } else {
432:                 std::cout << "[ " << rang::fg::red << "inactive" << rang::fg::reset << " ]" << std::endl;
433:             }
434:         }
435:     }
436:     std::cout << std::endl;
437: }
438:
439: // Calculate indices for this interval
440:
441: std::vector<std::size_t> GetIndices(const std::vector<Omega> &events, const std::vector<double> &W,
442:                                  const std::vector<double> &PR, const std::vector<double> &X) {
443:     std::vector<std::size_t> ind;
444:
445:     for (const auto &i : indices(events)) {
446:         if (events[i].M > M[0] && events[i].M <= M[1] && events[i].Pt > Pt[0] && events[i].Pt <= Pt[1] &&
447:             events[i].Y > Y[0] && events[i].Y <= Y[1]) {
448:             ind.push_back(i);
449:         }
450:     }
451:
452:     gra::maxi::PrintBar(" ");
453:     std::cout << rang::fg::green;
454:     printf("MASS RANGE: [%0.3f, %0.3f] GeV, PT RANGE: [%0.3f, %0.3f] GeV, Y =\n"
455:           "RANGE: [%0.3f, %0.3f] =\n"
456:           "Events in this hypobasis %d/%d\n",
457:           M[0], M[1], Pt[0], Pt[1], Y[0], Y[1], ind.size(), events.size());
458:     std::cout << rang::fg::reset;
459:
460:     return ind;
461: }
462:
463: // Test integrals spherical harmonics
464: // This is a test function, all integrals should give 1 if the normalization is
465: // correct
466: // correct
467: void TestSphericalIntegrals(int IMAX) {
468:     WRandom random;
469:     const unsigned int NCOEF = (IMAX + 1) * (IMAX + 1);
470:     std::cout << "TestSphericalIntegrals: " << std::endl << std::endl;
471:
472:     // Construct the efficiency coefficients EPSILON_LM with linear indexing
473:     std::vector<double> E(NCOEF, 0.0);
474:     std::vector<double> E2(NCOEF, 0.0); // for the uncertainty
475:
476:     // We MC evaluate the integral:
477:     // \int [R_Y_LM(Omega)]^2 dOmega, Omega = (cos(theta), phi)
478:     const double N = 10000; // Number of samples
479:
480:     // Integral domain volume: |cos(theta) x phi| = [-1,1] x [0,2pi)
481:     const double V = 4.0 * gra::math::PI;
482:
483:     for (std::size_t k = 0; k < N; ++k) {
484:         const double costheta = random.U(-1.0, 1.0);
485:         const double phi = random.U(0.0, 2.0 * PI);
486:
487:         for (int l = 0; l <= IMAX; ++l) {
488:             for (int m = -l; m <= l; ++m) {
489:                 const std::complex<double> Y = gra::math::Y_complex_basis(costheta, phi, l, m);
490:
491:                 // Function value
492:                 double f = std::ipow(std::abs(gra::math::INRM(Y, l, m)), 2) * V;
493:
494:                 // Sum to the MC integral and its squared version (for uncertainty)
495:                 const int ind = LinearInd(l, m);
496:                 E[ind] += f;
497:                 E2[ind] += pow(f);
498:             }
499:         }
500:     }

```

```

./src/MSpherical.cc          8/8
582: for (std::size_t m = 0; m < A.size_row(); ++m) {
583:     for (std::size_t n = 0; n < A.size_col(); ++n) | y[m] += pow2(A[m][n] * x[n]); }
584:     y[m] = sqrt(y[m]);
585: }
586: return y;
587: }
588:
589: // namespace spherical
590: | // namespace gra
591:
./src/MSpin.cc              1/16
1: // Spin polarization functions
2: |
3: | (Wigner-D functions, Clebsch-Gordan, Arbitrary Spin-Density matrix)
4: |
5: |
6: | In general, functions here use helicities for fermions which are
7: | not normalized by the spin vector norm.
8: | That is, functions use [-1/2, 1/2] not for example [-1, 1].
9: |
10: |
11: | TBD: Add permutation coherent spin chains with Bose-Einstein / Fermi-Dirac sign
12: | combinations, together with Breit-Wigner weights at amplitude level.
13: |
14: |
15: | (c) 2017-2020 Mikael Mieskolainen
16: | Licensed under the MIT license <http://opensource.org/licenses/MIT>.
17: |
18: | C++ standard
19: |#include <algorithm>
20: |#include <complex>
21: |#include <iostream>
22: |#include <random>
23: |#include <vector>
24: |
25: | Eigen
26: |#include <Eigen/Dense>
27: |
28: | Own
29: |#include "Granitti/MBus.h"
30: |#include "Granitti/MForm.h"
31: |#include "Granitti/MMath.h"
32: |#include "Granitti/MMatrix.h"
33: |#include "Granitti/MSpin.h"
34: |
35: using gra::math::msqrt;
36: using gra::math::pow2;
37: using gra::math::z1;
38:
39: namespace gra {
40: namespace spin {
41:
42:
43: // Used for checking eigenvalues
44: const double epsilon = 1e-6;
45:
46: // Initialize helicity decay amplitude matrix T (2s1 + 1) x (2s2 + 1),
47: // where s1, s2 are daughter spins (0, 1/2, 1, ...)
48: // T_j1 lambda_s1, lambda_s2
49: // T_j1 lambda_s1 | alpha_s1 |
50: // < lambda_s1 | alpha_s1 |
51: // < lambda_s2 | alpha_s2 |
52: // < lambda_s1 lambda_s2 | alpha_s1 alpha_s2 |
53:
54: void InitMatrix(gra::HELMatrix &hc, const gra::MParticle &sp, const gra::MParticle &sp1,
55:               const gra::MParticle &sp2) {
56:     const double J = p.spinX2 / 2.0;
57:     const int P = p.P;
58:
59:     const double s1 = p1.spinX2 / 2.0;
60:     const int P1 = p1.P;
61:
62:     const double s2 = p2.spinX2 / 2.0;
63:     const int P2 = p2.P;
64:
65: // Boson or Fermion
66: const bool is_boson1 = (int(s1)) ? true : false;
67: const bool is_boson2 = (int(s2)) ? true : false;
68: const bool identical = (p1.pdg == p2.pdg) ? true : false; // e.g. pi0 pi0 vs pi+ pi-
69:
70: // Helicity decay amplitude matrix
71:
72: hc.T = MMatrix<std::complex<double>>(static_cast<unsigned int>(2 * s1 + 1),
73:                                     static_cast<unsigned int>(2 * s2 + 1), 0.0);
74:
75: MMatrix<bool> need_to_set(20, 20, false);
76: MMatrix<bool> active(20, 20, false);
77:
78: bool tag_set = false;
79: unsigned int nonzero = 0;
80:
81: // Construct T-matrix (2s1 + 1) x (2s2 + 1) elements
82: std::cout << std::endl;

```

```

./src/MSpherical.cc          7/8
499: }
500: }
501:
502: for (int l = 0; l <= IMAX; ++l) {
503:     for (int m = -l; m <= l; ++m) {
504:         const int ind = LinearInd(l, m);
505:         E[ind] /= N;
506:         E2[ind] /= N;
507:
508:         // 1 sigma MC integration uncertainty
509:         double error = CalcError(E2[ind], E[ind], N);
510:         printf("T%d%d Vc = %8.1e +/- %8.1e (rel. error %9.3f percent) %n", l, m, E[ind], error,
511:               std::abs(error / E[ind]) * 100);
512:     }
513: }
514: std::cout << std::endl;
515: }
516:
517: // Find out linear index
518: int LinearInd(int l, int m) { return l * (l + 1) + m; }
519:
520: // Standard error
521: double CalcError(double f2, double f, double N) { return sqrt((f2 - std::pow(f, 2)) / N); }
522:
523: // Print matrix to a file
524: void PrintMatrix(FILE *fp, const std::vector<std::vector<double>> &A) {
525: // Error row coefficients
526: for (std::size_t i = 0; i < A.size(); ++i) {
527:     for (std::size_t j = 0; j < A[i].size(); ++j) { fprintf(fp, "%0.1e", A[i][j]); }
528:     fprintf(fp, "\n");
529: }
530: printf(fp, "\n");
531: }
532:
533: // Synthesize distributions with real SU-basis
534: // f(theta, phi) = \sum_{l=0}^N \sum_{m=-l}^l c_lm x Y_L^M(lm)(theta, phi)
535:
536: MMatrix<double> Y_real_synthesize(const std::vector<double> &c_lm, const std::vector<bool> &ACTIVE,
537:                                const int IMAX = msqrt(ACTIVE.size()) - 1);
538:
539: const int Nmax = math::linspace(-1.0, 1.0, N);
540: phi = math::linspace(-math::PI, math::PI, N);
541:
542: // Do the expansion
543: MMatrix<double> Z(M, N, 0.0);
544:
545: for (int l = 0; l <= IMAX; ++l) {
546:     for (int m = -l; m <= l; ++m) {
547:         const int index = gra::spherical::LinearInd(l, m);
548:         if (!ACTIVE[index]) | continue; // Not active
549:
550:         for (const std::size_t i : indices(costheta)) {
551:             for (const std::size_t j : indices(phi)) {
552:                 // Coeff
553:                 Z[i][j] += c_lm[index] * math::Y_real_basis(costheta[i], phi[j], l, m);
554:             }
555:         }
556:     }
557: }
558:
559: // Normalize elements
560: z = Z * (1.0 / max); // Normalize elements
561:
562:
563: if (normalized) {
564:     double max = 1e-32;
565:     for (std::size_t i = 0; i < Z.size_row(); ++i) {
566:         for (std::size_t j = 0; j < Z.size_col(); ++j) {
567:             if (Z[i][j] > max) { max = Z[i][j]; }
568:         }
569:     }
570:     z = Z * (1.0 / max); // Normalize elements
571: }
572:
573: return Z;
574: }
575:
576: // Vector Taylor expanded error propagation from x via A to y
577: // x contains standard deviations
578: // A is the system matrix
579: std::vector<double> ErrorProp(const MMatrix<double> &A, const std::vector<double> &x) {
580:     std::vector<double> y(A.size_row(), 0.0);
581: }

```

```

./src/MSpherical.cc          8/8
582: for (std::size_t m = 0; m < A.size_row(); ++m) {
583:     for (std::size_t n = 0; n < A.size_col(); ++n) | y[m] += pow2(A[m][n] * x[n]); }
584:     y[m] = sqrt(y[m]);
585: }
586: return y;
587: }
588:
589: // namespace spherical
590: | // namespace gra
591:
./src/MSpin.cc              1/16
1: // Spin polarization functions
2: |
3: | (Wigner-D functions, Clebsch-Gordan, Arbitrary Spin-Density matrix)
4: |
5: |
6: | In general, functions here use helicities for fermions which are
7: | not normalized by the spin vector norm.
8: | That is, functions use [-1/2, 1/2] not for example [-1, 1].
9: |
10: |
11: | TBD: Add permutation coherent spin chains with Bose-Einstein / Fermi-Dirac sign
12: | combinations, together with Breit-Wigner weights at amplitude level.
13: |
14: |
15: | (c) 2017-2020 Mikael Mieskolainen
16: | Licensed under the MIT license <http://opensource.org/licenses/MIT>.
17: |
18: | C++ standard
19: |#include <algorithm>
20: |#include <complex>
21: |#include <iostream>
22: |#include <random>
23: |#include <vector>
24: |
25: | Eigen
26: |#include <Eigen/Dense>
27: |
28: | Own
29: |#include "Granitti/MBus.h"
30: |#include "Granitti/MForm.h"
31: |#include "Granitti/MMath.h"
32: |#include "Granitti/MMatrix.h"
33: |#include "Granitti/MSpin.h"
34: |
35: using gra::math::msqrt;
36: using gra::math::pow2;
37: using gra::math::z1;
38:
39: namespace gra {
40: namespace spin {
41:
42:
43: // Used for checking eigenvalues
44: const double epsilon = 1e-6;
45:
46: // Initialize helicity decay amplitude matrix T (2s1 + 1) x (2s2 + 1),
47: // where s1, s2 are daughter spins (0, 1/2, 1, ...)
48: // T_j1 lambda_s1, lambda_s2
49: // T_j1 lambda_s1 | alpha_s1 |
50: // < lambda_s1 | alpha_s1 |
51: // < lambda_s2 | alpha_s2 |
52: // < lambda_s1 lambda_s2 | alpha_s1 alpha_s2 |
53:
54: void InitMatrix(gra::HELMatrix &hc, const gra::MParticle &sp, const gra::MParticle &sp1,
55:               const gra::MParticle &sp2) {
56:     const double J = p.spinX2 / 2.0;
57:     const int P = p.P;
58:
59:     const double s1 = p1.spinX2 / 2.0;
60:     const int P1 = p1.P;
61:
62:     const double s2 = p2.spinX2 / 2.0;
63:     const int P2 = p2.P;
64:
65: // Boson or Fermion
66: const bool is_boson1 = (int(s1)) ? true : false;
67: const bool is_boson2 = (int(s2)) ? true : false;
68: const bool identical = (p1.pdg == p2.pdg) ? true : false; // e.g. pi0 pi0 vs pi+ pi-
69:
70: // Helicity decay amplitude matrix
71:
72: hc.T = MMatrix<std::complex<double>>(static_cast<unsigned int>(2 * s1 + 1),
73:                                     static_cast<unsigned int>(2 * s2 + 1), 0.0);
74:
75: MMatrix<bool> need_to_set(20, 20, false);
76: MMatrix<bool> active(20, 20, false);
77:
78: bool tag_set = false;
79: unsigned int nonzero = 0;
80:
81: // Construct T-matrix (2s1 + 1) x (2s2 + 1) elements
82: std::cout << std::endl;

```

```

84:   std::cout << "gr:spin: "
85:   << " Particle 1: " << (is_boson1 ? "boson" : "fermion")
86:   << " Particle 2: " << (is_boson2 ? "boson" : "fermion") << std::endl;
87:
88:   if (hc.FConservation)
89:     std::cout << "Parity conservation: P = P1 x P2 x (-1)^{|P|} [P = '< P <', P1 = '< P1
90:     << ", P2 = '< P2 <']" << std::endl;
91:   }
92: }
93: std::cout << std::endl;
94: std::cout
95: << "gr:spin:InitMatrix: Calculating SU(2) decomposition [lambda = lambda1 - lambda2]: "
96: << std::endl;
97:
98: // Two loops, first check that sum |alpha_i|^2 = 1 for active ls values
99:
100: for (int MODE = 0; MODE < 2; ++MODE) {
101:   for (int s = 0; s <= static_cast<int>(s1 + s2); ++s) {
102:     for (int l = 0; l <= static_cast<int>(D + s); ++l) {
103:       for (int j = 0; j < static_cast<int>(D + s1 + 1); ++j) {
104:         for (int k = 0; k < static_cast<int>(D + s2 + 1); ++k) {
105:           const double lambda1 = s1 + static_cast<double>(l); // From negative to positive
106:           const double lambda2 = s2 + static_cast<double>(k); // From negative to positive
107:
108:           // Jacob-Wick
109:           const double lambda = lambda1 - lambda2;
110:
111:           // Check angular momentum conservation for z-axis if
112:           if (fabs(lambda1 - lambda2) <= J) ( continue; ) // Not conserved
113:
114:           // -----
115:           // Re-coupling coefficients << ... >>
116:
117:           // Normalization
118:           const double NORM = sqrt((2.0 * l + 1.0) / (2.0 * J + 1.0));
119:
120:           // Cj [lambda1 | s0 | lambda2]
121:           const double cgl = gra:spin:ClebschGordan(
122:             static_cast<double>(l), static_cast<double>(s), 0.0, lambda, J, lambda);
123:
124:           // Cj [lambda1 | s1 | lambda2, -lambda2]
125:           const double cg2 =
126:             gra:spin:ClebschGordan(s1, s2, lambda1, -lambda2, static_cast<double>(s), lambda);
127:
128:           // Angular momentum conserved (should be, by construction here)
129:           if (1*(std::abs(lambda) <= J) ( continue; ) // Not conserved
130:
131:           // If parity conserving decay
132:           if (hc.PConservation)
133:             // Parity of the final state (hold for both
134:             // bosons/fermions)
135:             const int P_tot = P1 * P2 * std::pow(-1, l);
136:             const int P_act = P1 * P2 * std::pow(-1, l);
137:             if (P_tot != P) ( continue; ) // Not conserved
138:           }
139:
140:           // Bose-Symmetry for identical Boson-Pairs (symmetric
141:           // wavefunction)
142:           if (is_boson1 && is_boson2 && identical {
143:             if (!BoseSymmetry(l, s)) ( continue; ) // Not right
144:           }
145:
146:           // Fermi-Symmetry for identical Fermion-Pairs (antisymmetric
147:           // wavefunction)
148:           if (!is_boson1 && !is_boson2 && identical {
149:             if (!FermiSymmetry(l, s)) ( continue; ) // Not right
150:           }
151:
152:           // CO-coupling product is non-zero
153:           if (std::abs(cgl * cg2) <= 1e-6) ( continue; )
154:
155:           // ** This coefficient is active **
156:           active[l][s] = true;
157:
158:           // INITIALIZATION MODE
159:           if (MODE == 0) {
160:             // Has not been set, throw error next
161:             if (alpha_set[l][s] == false) {
162:               need_to_set[l][s] = true;
163:               tag_set
164:             }
165:             // nonzero;
166:           }

```

```

250:   "missing from BRANCHING: (1,s) "
251:   " = ")
252:   throw std::invalid_argument(str0 + middle + str);
253: }
254: }
255:
256: // |l-s| <= J <= l + s, parity P = (-1)^{|l-s|}
257: // l is orbital angular momentum
258: // s is total spin
259: //
260: // Bose-Einstein statistics requires l-s to be even
261: // for identical Boson pairs.
262: bool BoseSymmetry(int l, int s) {
263:   const int number = l - s;
264:   if (number % 2 == 0) ( return true; )
265:   return false;
266: }
267:
268: // Fermi-Dirac statistics requires l-s to be even
269: // for identical Fermion pairs.
270: bool FermiSymmetry(int l, int s) {
271:   const int number = l + s;
272:   if (number % 2 == 0) ( return true; )
273:   return false;
274: }
275:
276: // Initial state spin treatment: Pomeron - Pomeron/Gamma -> Resonance X
277: //
278: // Apply couplings here
279:
280: std::complex<double> ProdAmp(const gra:ISORENTSCALAR &its, gra:PARAM_RES &res) {
281:   // Apply couplings here
282:   std::complex<double> A0 = res.g;
283:
284:   // Generation 2->1 spin correlations not active
285:   if (res.SPINGEN == false) ( return A0; )
286:
287:   // -----
288:   // Tensors J^P = 2+, 4+, 6+, ...
289:   // Pomeron effective spins
290:   double s1 = 0;
291:   double s2 = 0;
292:
293:   // SU(2) decomposition as given by InitMatrix
294:   //
295:   // only one alpha_(ls) = 1.0 needed
296:   MMatrix<std::complex<double>> T0 = {std::complex<double>(1.0)};
297:
298:   // -----
299:   // Scalar
300:   if (res.p.spinX2 == 0 && res.p.P == 1) ( return A0; )
301:
302:   // Pseudoscalar J^P = 0-, [e.g. eta(548), eta(958)]
303:   if (res.p.spinX2 == 0 && res.p.P == -1) {
304:     s1 = 1;
305:     s2 = 1;
306:     // SU(2) decomposition as given by InitMatrix
307:     //
308:     // only one alpha_(ls) = 1) = 1.0
309:     //
310:     // [std::complex<double>(0.707), std::complex<double>(0.000), std::complex<double>(0.000),
311:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
312:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
313:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
314:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
315:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
316:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
317:     // [std::complex<double>(0.000), std::complex<double>(0.000), std::complex<double>(0.000),
318:     // Forward proton pair deltaphi
319:     const double dphi = lts.pfinal[1].DeltaPhiAbs(lts.pfinal[2]);
320:
321:     // -----
322:     // M102 data (not fully spin) symmetric after MC, due to kinematics)
323:     //
324:     //
325:     //
326:     //
327:     // 0 ----- 180 deg
328:     //
329:     // Two cases m1=m2=1 and m1=m2=-1
330:     const double m1 = 1;
331:     const double m2 = 1;
332:

```

```

167:
168: // FINAL MODE
169: else if (MODE == 1) {
170:   // T matrix element
171:   hc.T[l][j] += hc.alpha[l][s] * NORM * cgl * cg2;
172:   print
173:   [ls0,s0] lambda2 = %4.1f, lambda2 = %4.1f : cgl*cg2 = %6.1f <alpha_ls = "
174:   "%0.5f %0.5f">";
175:   l, s, lambda1, lambda2, cgl * cg2, std::real(hc.alpha[l][s]),
176:   std::imag(hc.alpha[l][s]);
177: }
178: }
179: // lambda2
180: // lambda1
181: // l
182: // s
183:
184: // Check normalization
185: double sum_alphaq = 0;
186: if (MODE == 0) {
187:   for (std::size_t l = 0; l < hc.alpha.size_row(); ++l) {
188:     for (std::size_t s = 0; s < hc.alpha.size_col(); ++s) {
189:       if (active[l][s] == true) ( sum_alphaq += math::abs2(hc.alpha[l][s]); )
190:     }
191:   }
192: }
193:
194: // Now normalize the user alpha_ls input if needed such that: sum_(ls) |alpha_ls|^2 = 1
195: if (MODE == 0 && std::abs(sum_alphaq - 1.0) > 1e-6) {
196:   std::cout << "Normalizing input: sum_(ls) |alpha_ls|^2 = " << sum_alphaq << " = 1"
197:   << std::endl;
198:   // -----
199:   // Check normalization
200:   for (std::size_t l = 0; l < hc.alpha.size_row(); ++l) {
201:     for (std::size_t s = 0; s < hc.alpha.size_col(); ++s) {
202:       if (active[l][s] && hc.alpha[l][s] != sqrt(sum_alphaq); )
203:         hc.alpha[l][s] *= sqrt(sum_alphaq);
204:     }
205:   }
206: }
207: // MODE 0,1
208:
209: const std::string str0 = "gr:spin:InitMatrix: [" + gra:aux:Spin2toString(2 * J) + " + " +
210:   " + gra:aux:ParitytoString(P) + " - " + gra:aux:Spin2toString(2 * s1) +
211:   " + gra:aux:ParitytoString(P1) + " + " +
212:   " + gra:aux:Spin2toString(2 * s2) + " + " + gra:aux:ParitytoString(P2) + " +
213:   " + (identical == true ? std::string(Identical ? "true" : "false") : "") + " ]";
214:
215: if (nonzero == 0) {
216:   const std::string str = "Decay is impossible (spin-parity-statistics) [P_conservation = " +
217:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
218:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
219:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
220:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
221:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
222:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
223:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
224:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
225:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
226:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
227:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
228:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
229:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
230:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
231:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
232:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
233:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
234:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
235:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
236:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
237:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
238:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
239:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
240:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
241:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
242:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
243:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
244:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
245:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
246:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
247:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
248:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +
249:     hc.PConservation + "], (P_conservation = " + hc.PConservation + "], (P_conservation = " +

```

```

333: // See Kaidalov et al.
334: auto ReggeTheory = [(double t1, double t2, double m1, double m2) {
335:   return std::pow(std::abs(t1), std::abs(m1)) / (2.0 *
336:     std::pow(std::abs(t2), std::abs(m2)) / 2.0);
337: }];
338:
339: return A0 * ReggeTheory(lts.t1, lts.t2, m1, m2) * std::sin(dphi);
340: }
341:
342: // Axial vector J^P = 1- [e.g. f1(1260)]
343: if (res.p.spinX2 == 2 && res.p.P == 1) {
344:   s1 = 1;
345:   s2 = 1;
346:   // SU(2) decomposition as given by InitMatrix
347:   // only one alpha_(ls) = 2) = 1
348:   //
349:   T0 = {std::complex<double>(0.000), std::complex<double>(-0.500), std::complex<double>(0.000),
350:     std::complex<double>(0.500), std::complex<double>(0.000), std::complex<double>(-0.500),
351:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
352:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
353:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
354:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
355:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
356:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
357:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
358:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
359:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
360:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
361:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
362:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
363:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
364:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
365:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
366:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
367:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
368:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
369:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
370:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
371:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
372:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
373:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
374:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
375:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
376:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
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379:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
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387:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
388:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
389:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
390:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
391:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
392:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
393:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
394:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
395:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
396:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
397:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
398:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
399:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
400:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
401:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
402:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
403:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
404:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
405:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
406:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
407:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
408:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
409:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
410:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
411:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
412:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
413:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
414:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),
415:     std::complex<double>(0.000), std::complex<double>(0.500), std::complex<double>(0.000),

```



```

./src/MSpin.cc                               6/16
416: if (res.p.spinX2 == 0) {
417:     return res.hel.T[0][0] * A0; // (L,s) = (0,0)
418: }
419:
420: // Transition amplitude matrix f for X -> A + B in the X rest frame
421: // Daughter spins
422: const double s1 = its.decaytree[0].p.spinX2 / 2.0;
423: const double s2 = its.decaytree[1].p.spinX2 / 2.0;
424:
425: // Boost daughter to the A-rest frame
426: MVec boosted_daughter = its.decaytree[0].p4;
427: gra::kinematics:lorentzBoost(its.pfinal[0], its.pfinal[0].M(), boosted_daughter, -1);
428:
429: MMatrix<std::complex<double>> FX = MMatrix(res.hel.T, res.p.spinX2 / 2.0, s1, s2,
430:     boosted_daughter.Theta(), boosted_daughter.Phi());
431:
432: // Now go through the decay tree leaves recursively
433: TreeRecursion(its.decaytree[0]);
434: TreeRecursion(its.decaytree[1]);
435:
436: // Get tensor products back
437: MMatrix<std::complex<double>> fA;
438: TensorTree(its.decaytree[0], fA);
439:
440: MMatrix<std::complex<double>> fB;
441: TensorTree(its.decaytree[1], fB);
442:
443: // f matrices have dimensions: ((2s1 + 1)x(2s2 + 1)) x (2J + 1)
444:
445: // Total transition amplitude matrix as a tensor product
446: MMatrix<std::complex<double>> f;
447: if (its.decaytree[0].legs.size() == 0 && its.decaytree[1].legs.size() == 0) {
448:     f = FX; // No recursion
449: } else {
450:     f = gra::matoper:TensorProd(fA, fB) * FX; // Matrix x Matrix product
451: }
452:
453: MMatrix<std::complex<double>> fA_D = fA.Dagger();
454:
455: // -----
456: // Construct the D-matrix for an event-by-event spin space density operator
457: // rotation
458: double theta_R = 0.0;
459: double phi_R = 0.0;
460: GetRotation(its.res.FRAME, theta_R, phi_R);
461:
462: // -----
463: // Rotation does mixing of spin states. N.B. Eigenvalues do not change in
464: // rotation.
465: const MMatrix<std::complex<double>> D = gra::spin:DMatrix(res.p.spinX2 / 2.0, theta_R, phi_R);
466:
467: // rho_rot = D^\dagger rho D (keep this sandwich order!)
468: const MMatrix<std::complex<double>> rho_ROT = D.Dagger() * res.rho * D;
469:
470: // Weight (amplitude squared) of the event by the density matrix formalism:
471: // Tr(rho * f^\dagger dagger)
472: const MMatrix<std::complex<double>> wEvt = f * rho_ROT * f.Dagger();
473:
474: // (2s1+1) is normalization to match production coupling with spin-0
475: // (same cross section obtained as with spin-0 if density
476: // matrix = 1/(2s1+1) is totally unpolarized)
477: const double NORM = msqrt(res.p.spinX2 * 1.0);
478:
479: // Final weight at amplitude level
480: A0 = NORM * msqrt(std::real(Tr(fwd.Trace())));
481:
482: return A0;
483: }
484:
485: // Decay amplitude matrix X -> A + B
486:
487: MMatrix<std::complex<double>> CalculateMatrix(const MDecayBranch &branch) {
488:     const unsigned int A = 0;
489:     const unsigned int B = 1;
490:
491:     // Get daughter spins
492:     const double s1 = branch.legs[A].p.spinX2 / 2.0;
493:     const double s2 = branch.legs[B].p.spinX2 / 2.0;
494:
495:     // Rotate and boost daughters to the frame where z-spinned by X-flight direction (X-helicity rest
496:     // frame)
497:     std::vector<M4Vec> daughter = {branch.legs[A].p4, branch.legs[B].p4};

```

```

./src/MSpin.cc                               7/16
499: gra::kinematics:HXFrame(daughter, branch.p4);
500:
501: return MMatrix(branch.hel.T, branch.p.spinX2 / 2.0, s1, s2, daughter[A].Theta(),
502:     daughter[B].Phi());
503: }
504:
505: // Forward traverse the decay tree
506:
507: // -----
508: void TreeRecursion(MDecayBranch &branch) {
509:     if (branch.legs.size() == 2) {
510:         branch.f = CalculateMatrix(branch);
511:     } // Infinite recursion
512:     for (const auto &i : aux::indices(branch.legs) | TreeRecursion(branch.legs[i]); )
513:     }
514: }
515:
516: // -----
517: // Reverse traverse the sequential decay tree
518:
519: // -----
520: void TensorTree(const MDecayBranch &branch, MMatrix<std::complex<double>> fout) {
521:     if (branch.legs.size() == 2) {
522:         MMatrix<std::complex<double>> f0;
523:         MMatrix<std::complex<double>> f1;
524:         TensorTree(branch.legs[0], f0);
525:         TensorTree(branch.legs[1], f1);
526:
527:         // Now we are traversing backwards the binary tree at branch points
528:         if (branch.legs[0].legs.size() == 0 && branch.legs[1].legs.size() == 0) {
529:             out = branch.f;
530:         } else {
531:             out = matoper:TensorProd(f0, f1) * branch.f;
532:         }
533:     }
534: }
535:
536: // Last leg unit vector
537:
538: // -----
539: // Note the minus sign
540:
541: // -----
542: // Rotation angles to rotate the density matrix
543:
544: void GetRotation(const gra::LORENTZSCALAR &its, const std::string &IFRAME, double &theta_R,
545:     double &phi_R) {
546:     // Spin polarization/density matrix defined:
547:     // In direct non-rotated rest frame
548:     if (IFRAME == "CM") {
549:         theta_R = 0;
550:         phi_R = 0;
551:     }
552:     // In Helicity rest frame (quantization axis by beam orientation in the lab)
553:     else if (IFRAME == "HR") {
554:         theta_R = its.pfinal[0].Theta(); // +
555:         phi_R = its.pfinal[0].Phi(); // +
556:     }
557:     // In Collins-Soper rest frame (quantization axis by the beam bjector frame)
558:     else if (IFRAME == "CS") {
559:         MVec pib = its.pbeam1;
560:         MVec p2b = its.pbeam2;
561:         double theta = pib.Angle(p2b);
562:         double phi = p2b.Angle(pib);
563:         // Boost the beam particles
564:         gra::kinematics:lorentzBoost(its.pfinal[0], its.pfinal[0].M(), pib,
565:             -1); // Note the minus sign
566:         gra::kinematics:lorentzBoost(its.pfinal[0], its.pfinal[0].M(), p2b,
567:             -1); // Note the minus sign
568:     }
569:     // Now get the 3-momenta
570:     const std::vector<double> pibboost3 = pib.P3();
571:     const std::vector<double> p2bboost3 = p2b.P3();
572:
573:     // Collins-Soper bjector vector
574:     const std::vector<double> zaxis = gra::matoper:Unit(
575:         gra::matoper:MINUS(gra::matoper:Unit(pibboost3), gra::matoper:Unit(p2bboost3)));
576:
577:     MVec bjector;
578:     bjector.SetP3(zaxis);
579:
580:     // Now get the rotation angles
581:     theta_R = bjector.Theta(); // +

```

```

./src/MSpin.cc                               8/16
582:     phi_R = bjector.Phi(); // +
583:
584:     // In contridict-Jackson frame (quantization axis by the momentum transfer vector)
585:     else if (IFRAME == "DJ") {
586:         // Propagator
587:         MVec qlboost = its.q;
588:         gra::kinematics:lorentzBoost(its.pfinal[0], its.pfinal[0].M(), qlboost,
589:             -1); // Note the minus sign
590:
591:         theta_R = qlboost.Theta(); // +
592:         phi_R = qlboost.Phi(); // +
593:
594:     } else {
595:         // Throw exception
596:         const std::string str =
597:             "Spin specification: Unknown polarization Lorentz rest FRAME chosen: " + IFRAME +
598:             " (valid currently are: CM (no rotation), HK (helicity), CS (Collins-Soper));";
599:         throw std::invalid_argument(str);
600:     }
601: }
602:
603: // -----
604: // Create spin projections -J <= M <= J (number of 2s+1)
605: // negative to positive
606:
607: // Spin-1/2: [-1/2, 1/2]
608: // Spin-1: [-1, 0, 1]
609: // Spin-3/2: [-3/2, -1/2, 1/2, 3/2]
610: // Spin-2: [-2, -1, 0, 1, 2]
611:
612: std::vector<double> SpinProjections(double J) {
613:     std::vector<double> MMatrix<std::complex<double>> D(2 * J + 1, 0.0);
614:     double m = -J;
615:     for (std::size_t i = 0; i < M.size(); ++i) {
616:         M[i] = m;
617:         m += 1.0;
618:     }
619:     return M;
620: }
621:
622: // Returns transition amplitude matrix of size (2s1 + 1)x(2s2 + 1) x (2J + 1)
623:
624: // f_{(lambda_1)(lambda_2)M}(theta, phi) = D_{(lambda_2)M}^{(lambda_1)J}(theta, phi)
625: // T_{(lambda_1)(lambda_2)}
626:
627: // theta, phi defined in the decay rest frame
628:
629: // [REFERENCE: Jacob, Wick, On the general theory of particles with spin, 1959]
630: // [REFERENCE: Amisler, Bizot, Simulation of angular distributions, 1983]
631: MMatrix<std::complex<double>> MMatrix(const MMatrix<std::complex<double>> &T, double J, double s1,
632:     double s2, double theta, double phi) {
633:     if (T.size_row() != static_cast<unsigned int>(2 * s1 + 1) ||
634:         T.size_col() != static_cast<unsigned int>(2 * s2 + 1)) {
635:         throw std::invalid_argument("");
636:     }
637:     "gsa:spin:DMatrix:Input Matrix X
638:     (2s1+1)x(2s2+1) with wrong dimensions: " +
639:     std::to_string(T.size_row()) + " x " + std::to_string(T.size_col());
640:
641:     // Number of final state configurations
642:     const unsigned int N =
643:         static_cast<unsigned int>(2 * s1 + 1) * static_cast<unsigned int>(2 * s2 + 1);
644:
645:     // Initial state projections
646:     const std::vector<double> M = SpinProjections(J);
647:
648:     // Construct final state helicity configurations
649:
650:     // For indexing T matrix
651:     MMatrix<std::complex<double>> lambda_values(N, 2, 0.0);
652:     MMatrix<std::complex<double>> lambda_index(N, 2, 0);
653:
654:     unsigned int i = 0;
655:     for (int nu = 0; nu < static_cast<int>(2 * s1 + 1); ++nu) {
656:         for (int mu = 0; mu < static_cast<int>(2 * s2 + 1); ++mu) {
657:             lambda_values[i][0] = -s1 + static_cast<double>(mu);
658:             lambda_values[i][1] = -s2 + static_cast<double>(nu);
659:
660:             lambda_index[i][0] = mu;
661:             lambda_index[i][1] = nu;
662:
663:             ++i;
664:         }

```

```

./src/MSpin.cc                               9/16
665: // -----
666: // Construct transition amplitude matrix
667: MMatrix<std::complex<double>> f(N, 2 * J + 1, 0.0);
668:
669: // Rows = final state spin projections
670: for (std::size_t i = 0; i < N; ++i) {
671:     // Final state helicities
672:     const double lambda1 = lambda_values[i][0];
673:     const double lambda2 = lambda_values[i][1];
674:
675:     // Total lambda by definition
676:     const double lambda = lambda1 - lambda2;
677:
678:     // Columns = initial state polarizations
679:     for (std::size_t j = 0; j < 2 * J + 1; ++j) {
680:         // Wigner D * Helicity amplitude
681:         f[i][j] = WignerD(theta, phi, lambda, M[j], J) * T[lambda_index[i][0]][lambda_index[i][1]];
682:     }
683: }
684:
685: // -----
686: // std::cout << "f: " << std::endl;
687: // gra::matoper:PrintMatrixSeparate(f);
688: // std::endl;
689:
690: return f;
691: }
692:
693: // Construct Wigner D-matrix for a spin-space operator rotation
694: MMatrix<std::complex<double>> DMatrix(double J, double theta_R, double phi_R) {
695:     // Create spin projections
696:     const std::vector<double> M = SpinProjections(J);
697:
698:     const unsigned int n = static_cast<unsigned int>(2 * J + 1);
699:
700:     MMatrix<std::complex<double>> D(n, n, 0.0);
701:     for (std::size_t i = 0; i < n; ++i) {
702:         for (std::size_t j = 0; j < n; ++j) { D[i][j] = WignerD(theta_R, phi_R, M[i], M[j], J); }
703:     }
704:     return D;
705: }
706:
707: // SU(2) Clebsch-Gordan coefficients
708: double ClebschJordan(double j1, double j2, double m1, double m2, double j, double m) {
709:     // Selection rules
710:     if ((CGrules[j1][j2, m1, m2, j, m]) | (return 0.0); )
711:     }
712:     // Use Wigner-3j symbol to evaluate, note minus on m!
713:     return std::pow(-1.0, m + j1 - j2) * msqrt(2 * j + 1) * W3j(j1, j2, j, m1, m2, -m);
714: }
715:
716: // Check if half-integer
717: bool isHalfInteger(double x) {
718:     if (std::abs(2 * x - std::floor(2 * x)) < 1E-9) (return true; )
719:     return false;
720: }
721:
722: // Check if integer
723: bool isInt(double x) {
724:     if (std::abs(x - std::floor(x)) < 1E-9) (return true; )
725:     return false;
726: }
727:
728: // Floating point equality comparison
729: bool isEqual(double x, double y) {
730:     if (std::abs(x - y) < 1E-9) (return true; )
731:     return false;
732: }
733:
734: // CP-selection rules
735: double CGrules(double j1, double j2, double m1, double m2, double j, double m) {
736:     const bool PRINT_ON = false;
737:
738:     if (isHalfInt(j1) || isHalfInt(j2) || isHalfInt(j)) {
739:         if (PRINT_ON) print("CGrules: (j1, j2, j) not integer or half-integer!\n");
740:         return false;
741:     }
742:
743:     if (isInt(j1 + j2 + j)) { // Integer parameter rule
744:         if (PRINT_ON) print("CGrules: j1 + j2 + j not integer!\n");
745:         return false;
746:     }
747: }

```



```
./src/MSpin.cc 10/16
748: if (!chalfint(m1) || !chalfint(m2) || !chalfint(m)) {
749:   if (PRINT_ON) printf("Circles: (m1,m2,m) not integer or half-integer\n");
750:   return false;
751: }
752: if (!isequal(m1 + m2, m)) {
753:   if (PRINT_ON) printf("Circles: m1 + m2 != m\n");
754:   return false;
755: }
756: if (!isequal(j1 - m1, std::floor(j1 - m3))) {
757:   if (PRINT_ON) printf("Circles: parity of 2j_1 and 2m_1 does not match\n");
758:   return false;
759: }
760: if (!isequal(j2 - m2, std::floor(j2 - m3))) {
761:   if (PRINT_ON) printf("Circles: parity of 2j_2 and 2m_2 does not match\n");
762:   return false;
763: }
764: if (!isequal(j3 - m3, std::floor(j3 - m3))) {
765:   if (PRINT_ON) printf("Circles: parity of 2j_3 and 2m_3 does not match\n");
766:   return false;
767: }
768: if ((j1 + j2) < j3 || j3 < std::abs(j1 - j2)) {
769:   if (PRINT_ON) printf("Circles: |m| > j\n");
770:   return false;
771: }
772: if ((j1 + j2) < j3 || j3 < std::abs(j1 - j2)) {
773:   if (PRINT_ON) printf("Circles: |m| > j\n");
774:   return false;
775: }
776: if (std::abs(m1) > j1) {
777:   if (PRINT_ON) printf("Circles: |m1| > j1\n");
778:   return false;
779: }
780: if (std::abs(m2) > j2) {
781:   if (PRINT_ON) printf("Circles: |m2| > j2\n");
782:   return false;
783: }
784: if (std::abs(m3) > j3) {
785:   if (PRINT_ON) printf("Circles: |m3| > j3\n");
786:   return false;
787: }
788: if (std::abs(m) > j) {
789:   if (PRINT_ON) printf("Circles: |m| > j\n");
790:   return false;
791: }
792: return true;
793: }
794: // Wigner 3j symbol
795: //
796: // [1] J2 J3
797: // [m1 m1 m3]
798: // http://mathworld.wolfram.com/Wigner3j-Symbol.html
799: //
800: // Create coefficient structure
801: // Check selection rules
802: // (WJ)coeffs(j1, j2, j3, m1, m2, m3) (return 0.0)
803: //
804: // Boundaries for which factorials are non-negative
805: // const int upper = std::min(std::min(c1, c2), c3);
806: // double sum = 0;
807: // Evaluate sum term
808: for (int k = lower; k <= upper; ++k)
809:   sum += std::pow(-1, k) * (g1a:math::factorial(k) * g1b:math::factorial(c1 + k) *
810:     g2a:math::factorial(c2 + k) * g2b:math::factorial(c3 - k) *
811:     g3a:math::factorial(c4 - k) * g3b:math::factorial(c5 - k));
812: // Evaluate by Raab formula
813: // const double w = std::pow(-1, 0, j1 - j2 - m3) *
814: //   msqr(TriangCoeff(j1, j2, j3) * g1a:math::factorial(j1 + m1) *
815: //     g1b:math::factorial(j1 - m1) * g1c:math::factorial(j2 + m2) *
816: //     g1d:math::factorial(j2 - m2) * g1e:math::factorial(j3 + m3) *
817: //     g1f:math::factorial(j3 - m3));
```

```
./src/MSpin.cc 11/16
831: g1a:math::factorial(j3 - m3) *
832:   fsum;
833: return w;
834: }
835: // Triangle coefficient
836: // double TriangCoeff(double j1, double j2, double j3) {
837:   return g1a:math::factorial(j1 + j2 + j3) * g1b:math::factorial(j1 - j2 + j3) *
838:     g1c:math::factorial(-j1 + j2 + j3) / g1d:math::factorial(j1 + j2 + j3 + 1);
839: }
840: // Selection rules
841: // bool WJrules(double j1, double j2, double j3, double m1, double m2, double m3) {
842:   const bool PRINT_ON = false;
843:   if (rint(j1 + j2 + j3) != 0)
844:     if (PRINT_ON) std::cout << "WJ3R1: j1 + j2 + j3 not integer\n";
845:   if (rint(j1 - m1) != 0)
846:     if (PRINT_ON) std::cout << "WJ3R2: j1 - m1 not integer\n";
847:   if (rint(j2 - m2) != 0)
848:     if (PRINT_ON) std::cout << "WJ3R3: j2 - m2 not integer\n";
849:   if (rint(j3 - m3) != 0)
850:     if (PRINT_ON) std::cout << "WJ3R4: j3 - m3 not integer\n";
851:   if (rint(m1 + m2 + m3) != 0)
852:     if (PRINT_ON) std::cout << "WJ3R5: m1 + m2 + m3 != 0\n";
853:   if (rint(m1) != 0)
854:     if (PRINT_ON) std::cout << "WJ3R6: m1 != 0\n";
855:   if (rint(m2) != 0)
856:     if (PRINT_ON) std::cout << "WJ3R7: m2 != 0\n";
857:   if (rint(m3) != 0)
858:     if (PRINT_ON) std::cout << "WJ3R8: m3 != 0\n";
859:   if (rint(m) != 0)
860:     if (PRINT_ON) std::cout << "WJ3R9: m != 0\n";
861:   if (rint(j1) != 0)
862:     if (PRINT_ON) std::cout << "WJ3R10: j1 != 0\n";
863:   if (rint(j2) != 0)
864:     if (PRINT_ON) std::cout << "WJ3R11: j2 != 0\n";
865:   if (rint(j3) != 0)
866:     if (PRINT_ON) std::cout << "WJ3R12: j3 != 0\n";
867:   if (rint(j) != 0)
868:     if (PRINT_ON) std::cout << "WJ3R13: j != 0\n";
869:   if (rint(j1 + j2) < j3 || j3 < std::abs(j1 - j2))
870:     if (PRINT_ON) std::cout << "WJ3R14: j over valid domain\n";
871:   if (rint(j1) > j1)
872:     if (PRINT_ON) std::cout << "WJ3R15: |m1| > j1\n";
873:   if (rint(j2) > j2)
874:     if (PRINT_ON) std::cout << "WJ3R16: |m2| > j2\n";
875:   if (rint(j3) > j3)
876:     if (PRINT_ON) std::cout << "WJ3R17: |m3| > j3\n";
877:   if (rint(m) > j)
878:     if (PRINT_ON) std::cout << "WJ3R18: |m| > j\n";
879:   return true;
880: }
881: // Calculate quantum mechanical of Neumann entropy of a density matrix
882: // double NeumannEntropy(const MMatrix& rho) {
883:   const unsigned int N = rho.size_col();
884:   Eigen::MatrixXcd A(N, N);
885:   // Collect matrix elements
886:   for (std::size_t i = 0; i < N; ++i)
887:     for (std::size_t j = 0; j < N; ++j) A(i, j) = rho[i][j];
888:   // Solve eigenvalues
889:   Eigen::SelfAdjointEigenSolver<Eigen::MatrixXcd> es(A);
890:   if (es.info() != Eigen::Success) return -1.0;
891:   // Calculate von Neumann Entropy
892:   double entropy = 0.0;
893:   for (std::size_t i = 0; i < N; ++i)
894:     entropy += -es.eigenvalues()[i].real() * log(es.eigenvalues()[i].real());
895:   // Take real part for numerics
896:   return entropy;
```

```
./src/MSpin.cc 12/16
914: }
915: entropy = -entropy;
916: return entropy;
917: }
918: // Wigner small-d (real valued convention, z-y-z convention)
919: //
920: // [REFERENCE: en.wikipedia.org/wiki/Wigner-D-matrix#Wigner_small_d-matrix]
921: //
922: // m, mp, j all in integers (not half-integers)
923: //
924: // double WignerSmallD(double theta, double m, double mp, double j) {
925:   const double factor =
926:     g1a:math::msqr(g1a:math::factorial(c1 + m) * g1b:math::factorial(j - m) *
927:       g1c:math::factorial(j + mp) * g1d:math::factorial(j - mp));
928:   double s = 0.0;
929:   for (int k = 0; k <= j; ++k)
930:     for (int l = 0; l <= j - k; ++l)
931:       s += std::pow(-1, k + l) * g1e:math::factorial(j - k - l) * g1f:math::factorial(k + l) *
932:         g1g:math::factorial(j - m - k) * g1h:math::factorial(j + m - k) *
933:         g1i:math::factorial(k + mp - m) * g1j:math::factorial(k - mp);
934:   value = s * factor;
935:   return value;
936: }
937: // Fatal error
938: if (std::isnan(s))
939:   std::string str =
940:     "g1a:math::WignerSmallD: NaN value with (theta, m, mp, j) = " + std::to_string(theta) + " + " +
941:     std::to_string(m) + " + " + std::to_string(mp) + " + " + std::to_string(j);
942:   throw std::invalid_argument(str);
943: }
944: return factor * s;
945: }
946: // Wigner large-D
947: // D = (theta, phi, m, mp, j) exp(iLJ theta) exp(iLz phi)
948: //
949: // m, mp, j all in integers (not half-integers)
950: //
951: // const complexDouble WignerD(double theta, double phi, double m, double mp, double j) {
952:   return WignerSmallD(theta, m, mp, j) * std::exp(std::complexDouble(0, 1) * m * phi);
953: }
954: // Density matrix properties:
955: // 1. rho^2 = rho (projector)
956: // 2. rho^dagger = rho (hermiticity)
957: // 3. Tr(rho) = 1 (normalization)
958: // 4. rho == 0 (positivity, eigenvalues greater or equal to zero)
959: //
960: // Test positivity of the density matrix
961: // - Also checks the trivial fact that rho
962: // has right dimensions given J, that is (2J+1)
963: // - Also checks that the normalization is ok.
964: //
965: // TBD, write down generalization for any spin (e.g. done with Eigen library
966: // eigenvalues)
967: //
968: // bool Positivity(const MMatrix& rho) {
969:   const unsigned int N = rho.size_row();
970:   // Dimensions
971:   if ((N - 1) / 2 != J)
972:     return false;
973:   // Eigenvalues
974:   Eigen::MatrixXcd A(N, N);
975:   for (std::size_t i = 0; i < N; ++i)
976:     for (std::size_t j = 0; j < N; ++j) A(i, j) = rho[i][j];
977:   Eigen::SelfAdjointEigenSolver<Eigen::MatrixXcd> es(A);
978:   if (es.info() != Eigen::Success) return false;
979:   for (std::size_t i = 0; i < N; ++i)
980:     if (es.eigenvalues()[i].real() < 0) return false;
981:   return true;
982: }
```

```
./src/MSpin.cc 13/16
997: return false;
998: }
999: // Normalization, abs taken for comparison
1000: const double tracerho = std::abs(rho.Trace());
1001: if (std::abs(tracerho - 1.0) > 1e-2)
1002:   std::string str =
1003:     "g1a:math::Positivity: Density matrix is not properly "
1004:     "normalized (Tr(rho) = " +
1005:     std::to_string(tracerho) + ") (should be 1)";
1006:   g1a:math::PrintMatrixSeparate(rho);
1007:   throw std::invalid_argument(str);
1008: }
1009: return true;
1010: }
1011: // Negative for floating point reasons
1012: const double NEG_EPS = -1e6;
1013: // Spin-1
1014: if (m == 1)
1015:   // Component as:
1016:   // rho_1 = (0.5*(1-a), b+i*c, d,
1017:   //          b-i*c, a, -b+i*c,
1018:   //          d, -b-i*c, 0.5*(1+c));
1019:   // Test positivity of rho (ANALYTICAL)
1020:   const double a = std::real(rho[0][1]);
1021:   const double b = std::real(rho[0][2]);
1022:   const double c = std::real(rho[1][1]);
1023:   const double d = std::real(rho[0][2]);
1024:   const double pc1 = -4 * b * b - 4 * c * c - 2 * d * d + a * (3 * a + a) / 2.0 + 1 / 2.0;
1025:   const double pc2 =
1026:     -(3 * (2 * d - a + 1) * (2 * a * d - a + a * a + 4 * b * b + 4 * c * c)) / 2.0;
1027:   if (pc1 >= NEG_EPS && pc2 >= NEG_EPS)
1028:     return true;
1029:   else {
1030:     print("pc1 = " + std::to_string(pc1) + ", pc2 = " + std::to_string(pc2));
1031:     std::string str =
1032:       "g1a:math::Positivity: Eigenvalues of the density "
1033:       "matrix are not pos 0!";
1034:     g1a:math::PrintMatrixSeparate(rho);
1035:     throw std::invalid_argument(str);
1036:   }
1037: }
1038: // Spin-2
1039: if (m == 2) {
1040:   // Components as:
1041:   // rho_2 = [a, f+i*c, g+i*d, h+i*c, m,
1042:   //          f-i*c, b, j+i*k, l, -h-i*c,
1043:   //          g-i*d, j-i*k, j-2*(a+b), -j+i*k, g-i*d,
1044:   //          -h-i*c, b, -f-i*c,
1045:   //          m, -h-i*c, g+i*d, -f-i*c, a];
1046:   // Test positivity of rho (ANALYTICAL)
1047:   const double a = std::real(rho[0][0]);
1048:   const double b = std::real(rho[1][1]);
1049:   const double c = std::real(rho[0][1]);
1050:   const double d = std::real(rho[0][2]);
1051:   const double e = std::real(rho[0][3]);
1052:   const double f = std::real(rho[0][4]);
1053:   const double g = std::real(rho[0][5]);
1054:   const double h = std::real(rho[0][6]);
1055:   const double i = std::real(rho[0][7]);
1056:   const double j = std::real(rho[0][8]);
1057:   const double k = std::real(rho[1][2]);
1058:   const double l = std::real(rho[1][3]);
1059:   const double m = std::real(rho[0][4]);
1060:   const double n = std::real(rho[0][5]);
1061:   const double pc1 =
1062:     -6 * a * a - 8 * a * b + 4 * a - 6 * b * b + 4 * b - 4 * c * c - 4 * d * d -
1063:     2 * l * l - 2 * m * m;
1064:   const double pc2 =
1065:     -12 * a * a * a - 48 * a * a * b + 6 * a * a * d + 12 * a * a * e + 12 * a * a * f +
1066:     12 * a * a * g + 12 * a * a * h - 24 * a * a * j - 24 * a * a * k +
1067:     12 * a * m * m + 12 * b * b * d + 6 * b * b * e + 12 * b * b * f + 12 * b * b * g +
1068:     12 * b * d * d + 12 * b * e * e + 12 * b * f * f + 24 * b * g * g +
1069:     12 * b * h * h - 12 * b * j * j - 12 * b * k * k + 12 * b * l * l -
1070:     12 * c * c * c + 24 * c * c * d + 24 * c * c * e + 1 - 24 * c * e * m -
1071:     24 * c * g * k + 12 * d * d * m - 24 * d * e * j + 24 * d * f * k -
```

./src/MSpin.cc

14/16

```

1080:      24 * d * h * k - 12 * e * e + 24 * e * g * k - 12 * f * f + 24 * f * g * j +
1081:      24 * f * h - 12 * f * h * j - 12 * k * k - 12 * l * l - 6 * m * m;
1082:
1083:  const double pc3 =
1084:      48 * a * a * b - 168 * a * a * b * b - 96 * a * a * a * b + 96 * a * a * c * c +
1085:      96 * a * e * e + 96 * a * a * g * f + 96 * a * a * h * h - 48 * a * a * j * j -
1086:      48 * a * a * k * k + 72 * a * a * l * l - 96 * a * b * b * b + 48 * a * b * b * b *
1087:      144 * a * b * c * c - 96 * a * b * d * d + 144 * a * b * e * e + 144 * a * b * f * f -
1088:      96 * a * b * g * g + 144 * a * b * h * h - 96 * a * b * j * j - 96 * a * b * k * k +
1089:      96 * a * b * l * l + 96 * a * b * m * m - 48 * a * c * c * c + 96 * a * c * d * d * j -
1090:      96 * a * c * e * e + 192 * a * c * f * f - 96 * a * c * g * g * k - 96 * a * c * h * h * j +
1091:      96 * a * d * f * k - 96 * a * d * h * k - 48 * a * e * e * e + 96 * a * e * g * k +
1092:      48 * a * f * f * 96 * a * f * g * j - 96 * a * f * h * h * l + 192 * a * f * h * m -
1093:      96 * a * g * h * j - 48 * a * h * h * 96 * a * j * j * l - 96 * a * k * k * l -
1094:      48 * a * l * l * 96 * b * b * c * c - 48 * b * b * d * d * 96 * b * b * e * e * e +
1095:      96 * b * b * f * f * c - 48 * b * b * g * g * 96 * b * b * h * h + 72 * b * b * m * m -
1096:      48 * b * c * c * 96 * b * c * d * j - 192 * b * c * e * e + 196 * b * c * e * m -
1097:      96 * b * c * g * k + 96 * b * d * d * m - 96 * b * d * e * j - 96 * b * d * f * k -
1098:      96 * b * e * h * k - 48 * b * e * e * 96 * b * e * g * k - 48 * b * f * f * f +
1099:      96 * b * f * g * j - 192 * b * f * h * l + 96 * b * f * h * m + 96 * b * g * g * m -
1100:      96 * b * h * h * j - 48 * b * h * h * 48 * b * m * m + 24 * c * c * c * d +
1101:      48 * c * c * d * d - 48 * c * c * e * e - 48 * c * c * f * f + 48 * c * c * g * g +
1102:      48 * c * c * h * h + 48 * c * c * j * j + 48 * c * c * k * k + 48 * c * c * l * l * m +
1103:      96 * c * d * d * e + 96 * c * d * e * j - 96 * c * d * f * m - 192 * c * e * e * f * h +
1104:      96 * c * e * g * g + 96 * c * e * j * j + 96 * c * e * k * k + 96 * c * e * l * l -
1105:      96 * c * e * m - 96 * c * g * g * k * l + 96 * c * g * k * m + 48 * d * d * e * e +
1106:      96 * d * e * j * l + 96 * d * e * j * m - 96 * d * e * k * m - 196 * d * f * k * m -
1107:      96 * d * e * j * l + 96 * d * e * j * m - 96 * d * e * k * m - 196 * d * f * k * m -
1108:      96 * d * e * j * l + 96 * d * e * j * m - 96 * d * e * k * m - 196 * d * f * k * m +
1109:      48 * e * e * g * g + 48 * e * e * h * h + 48 * e * e * j * j + 48 * e * e * k * k +
1110:      48 * e * e * l * l + 96 * e * e * k * k - 196 * e * e * k * m - 24 * f * f * f * f +
1111:      48 * f * f * g * g - 48 * f * f * h * h + 48 * f * f * j * j + 48 * f * f * k * k +
1112:      48 * f * f * l * l + 96 * f * f * g * g + 96 * f * f * j * j - 196 * f * g * j * m +
1113:      96 * f * h * j + 96 * f * h * k * k + 96 * f * h * l * l - 96 * f * h * m +
1114:      48 * f * h * h * h + 48 * g * g * j * l - 96 * g * h * j * l - 96 * g * h * j * m +
1115:      24 * h * h * h * h + 48 * h * h * h * j + 48 * h * h * k * k + 48 * h * h * l * l * m +
1116:      48 * h * h * m * m + 48 * g * k * e * f * g * j + 48 * g * l * l * m - m;
1117:
1118:  const double pc4 =
1119:      120 *
1120:      (c * c + 2 * c * e + e * e + f * f + 2 * f * h + h * h - a * b - a * l + b * m + l * m) *
1121:      (a * l - 2 * c * e - a * b - 2 * f * h - b * m + l * m + 2 * a * b + b * h + 2 * a * b -
1122:      2 * b * e - 2 * b * f * c - 2 * a * h - 2 * a * b * g - 2 * b * h + h +
1123:      2 * a * j * j + 2 * a * k * k - 2 * a * l + 2 * b * m * m + 2 * d * d * l -
1124:      2 * g * g * l + 2 * j * m + 2 * k * k * m + c * c * e + f * f + h * h +
1125:      4 * a * c * e + 4 * b * c * d - 2 * a * b * l + 4 * a * f * h + 2 * a * b * m +
1126:      4 * b * h * h + 4 * c * d * j + 4 * a * d * e * j + 4 * a * d * e * k +
1127:      4 * g * h * k + 4 * e * g * k * e * f * g * j + 4 * g * l * l * m - a * l * m -
1128:      2 * b * l * m);
1129:
1130:  if (pc1 >= NEG_EPS && pc2 >= NEG_EPS && pc3 >= NEG_EPS && pc4 >= NEG_EPS) {
1131:      return true;
1132:  } else {
1133:      printf(
1134:          "pc1 = %0.5f, pc2 = %0.5f, pc3 = %0.5f, pc4 = "
1135:          "%0.5f\n",
1136:          pc1, pc2, pc3, pc4);
1137:      std::string str =
1138:          "grs:spin:Positivity: Eigenvalues of the density "
1139:          "matrix are not >= 0.";
1140:      gra:matoper::PrintMatrixSeparate(rho);
1141:      throw std::invalid_argument(str);
1142:      return false;
1143:  }
1144:  } else {
1145:      std::string str =
1146:          "grs:spin:Positivity: Not a valid spin, only J = 1 and 2 currently tested -- check "
1147:          "manually.";
1148:      gra:matoper::PrintMatrixSeparate(rho);
1149:      return true;
1150:  }
1151:  }
1152:
1153:  // Generate random density matrix for the resonance
1154:  // Create random (under Haar / Hilbert-Schmidt measure) spin-density matrices
1155:  // for any spin,
1156:  // and for spin 1 and 2 the possibility to enforce parity conservation
1157:  //
1158:  //
1159:  // TBD, write down the generalization of parity conservation constraint for any
1160:  // spin.
1161:  MMatrix<std::complex> RandomRho( unsigned int J, bool parity, MRandom rng) {
1162:      const std::complex> xi(0, 1); // imag unit

```

./src/MSpin.cc

16/16

```

1246:  } // namespace spin
1247:  } // namespace gra
1248:

```

./src/MSpin.cc

15/16

```

1163:  MMatrix<std::complex> rho;
1164:
1165:  // Draw random density matrices until valid found, takes a couple
1166:  // usually
1167:  while (true) {
1168:      // Create random complex matrix (n x n)
1169:      const unsigned int n = 2 * J + 1;
1170:      MMatrix<std::complex> r(n, n);
1171:
1172:      for (std::size_t i = 0; i < n; ++i) {
1173:          for (std::size_t j = 0; j < n; ++j) {
1174:              r[i][j] = rng.G0, 1) * z1 * rng.G0, 1); // real and imag part from
1175:              // normal distribution
1176:          }
1177:      }
1178:      // Take product (outer product per vector): r*r\dagger
1179:      rho = r * r.dagger();
1180:
1181:      // Normalize the trace to 1 (real just for conversion to double)
1182:      const double scale = 1.0 / std::real(rho.Trace());
1183:      rho = rho * scale;
1184:
1185:      // If parity conservation required, then symmetrize
1186:      if (parity) {
1187:          if (J == 1) {
1188:              // 1.1 Purely real elements by parity and
1189:              // hermiticity
1190:              rho[0][2] = std::real(rho[0][2]);
1191:
1192:              // 1.2 Parity constrained
1193:              rho[1][2] = -std::real(rho[0][1]) + xi * std::imag(rho[0][1]);
1194:
1195:              // 2. Lower triangle = hermiticity constrained
1196:              for (std::size_t i = 0; i < n; ++i) {
1197:                  for (std::size_t j = 0; j < i; ++j) { rho[i][j] = std::conj(rho[j][i]); }
1198:              }
1199:
1200:              // 3. Diagonal parity constrained
1201:              rho[2][2] = std::real(rho[0][0]);
1202:
1203:              if (J == 2) {
1204:                  // 1.1 Purely real elements by parity and
1205:                  // hermiticity
1206:                  rho[0][4] = std::real(rho[0][4]);
1207:                  rho[1][3] = std::real(rho[1][3]);
1208:
1209:                  // 1.2 Parity constrained
1210:                  rho[1][4] = -std::real(rho[0][3]) + xi * std::imag(rho[0][3]);
1211:                  rho[2][3] = -std::real(rho[1][2]) + xi * std::imag(rho[1][2]);
1212:                  rho[3][4] = std::conj(rho[0][2]);
1213:                  rho[3][4] = -std::real(rho[0][1]) + xi * std::imag(rho[0][1]);
1214:
1215:                  // 2. Lower triangle = hermiticity constrained
1216:                  for (std::size_t i = 0; i < n; ++i) {
1217:                      for (std::size_t j = 0; j < i; ++j) { rho[i][j] = std::conj(rho[j][i]); }
1218:                  }
1219:
1220:                  // 3. Diagonal parity constrained
1221:                  rho[3][3] = std::real(rho[0][0]);
1222:                  rho[4][4] = std::real(rho[0][0]);
1223:
1224:                  // Re-Normalize trace to 1 (real just for conversion to double)
1225:                  const double newscale = 1.0 / std::real(rho.Trace());
1226:                  rho = rho * newscale;
1227:              }
1228:
1229:              // printMatrix(rho);
1230:              // Check that the matrix is valid density matrix in terms of
1231:              // positivity requirement
1232:              try {
1233:                  if (Positivity(rho, J)) { break; }
1234:              } catch (...) {
1235:                  continue; // not valid density matrix
1236:              }
1237:          }
1238:      }
1239:
1240:      gra:matoper::PrintMatrixSeparate(rho);
1241:      return rho;
1242:  }
1243:
1244:
1245:

```

./src/MPDG.cc

1/8

```

1:  // PDG Class
2:
3:  // (c) 2017-2020 Mikael Mieskolainen
4:  // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6:  #include <map>
7:  #include <regex>
8:
9:  // Own
10: #include "Graniitti/MPDG.h"
11:
12: // Libraries
13: #include "rang.hpp"
14:
15: namespace gra {
16:
17: // Read PDG particle data in
18: // Input as the full path to PDG .mod file
19: void MPDG::ReadParticleData(const std::string& filepath) {
20:     std::cout << "MPDG:ReadParticleData: Reading in PDG tables: ";
21:     std::ifstream infile(filepath);
22:     std::string line;
23:
24: // Clear particle data
25:     PDG_table.clear();
26:
27:     while (std::getline(infile, line)) {
28:         std::istringstream iss(line);
29:
30:         // Check * (comment) lines
31:         std::string first;
32:         iss >> first;
33:         std::string str = first;
34:         str.erase(str.begin() + 1, str.end());
35:
36:         // Skip lines with *
37:         if (str.compare("") == 0) continue;
38:
39:         // ***** DEFINITION FROM PDG *****
40:         // 1 - 8 | Monte Carlo particle numbers as described in
41:         // the "Review of
42:         // 9 - 16 | Particle Physics". Charge states appear, as
43:         // appropriate,
44:         // 17 - 24 | from left-to-right in the order -, 0, +, ++.
45:         // 25 - 32
46:         // 33 | blank
47:         // 34 - 52 | central value of the mass (double precision)
48:         // 53 - 60 | positive error
49:         // 61 | blank
50:         // 62 - 69 | negative error
51:         // 70 | blank
52:         // 71 - 68 | central value of the width (double precision)
53:         // 69 | blank
54:         // 90 - 99 | positive error
55:         // 98 | blank
56:         // 99 - 106 | negative error
57:         // 107 | blank
58:         // 108 - 128 | particle name left-justified in the field and
59:         // charge states right-justified in the field.
60:         // This field is for ease of visual examination
61:         // of the file and
62:         // should not be taken as a standardized
63:         // presentation of
64:         // particle names.
65:
66:         const int L = 40;
67:         char cid[L] = {0};
68:         char cmass[L] = {0};
69:         char cwidth[L] = {0};
70:         char cname[L] = {0};
71:
72:         sscanf(line.c_str(), "%30s %30s %30s %21s", cid, cmass, cwidth, cname);
73:         // print(cid, cmass, cwidth, cname);
74:
75:         std::string sid(cid, L);
76:         std::string smass(cmass, L);
77:         std::string swidth(cwidth, L);
78:         std::string sname(cname, L);
79:
80:         std::string temp_str;
81:
82:
83: // -----

```

```

84: // ID
85: // First find out how many charges
86: std::vector<int> id;
87: std::stringstream ss(id);
88:
89: while (ss >> temp_str) {
90: // Convert string to int
91: int value = 0;
92: std::stringstream(temp_str) >> value;
93: if (value != 0) { id.push_back(value); }
94: }
95:
96: // -----
97: // MASS
98: // Mass and errors
99: std::vector<double> mass;
100: std::stringstream ss2(mass);
101:
102: while (ss2 >> temp_str) {
103: // Convert string to int
104: double value = 0;
105: std::stringstream(temp_str) >> value;
106: mass.push_back(value);
107: }
108:
109: // -----
110: // WIDTH
111: // Width and errors
112: std::vector<double> width;
113: std::stringstream ss3(width);
114:
115: // If the particle WIDTH columns were empty
116: // -> name column string will be empty by reading logic. This
117: // checks that.
118: bool empty = true;
119: if ((sname.find("0") != std::string::npos) || (sname.find("+") != std::string::npos) ||
120:     (sname.find("-") != std::string::npos) || (sname.find("*") != std::string::npos)) {
121:     empty = false;
122: }
123:
124: if (empty) {
125: while (ss3 >> temp_str) {
126: // Convert string to double
127: double value = 0;
128: std::stringstream(temp_str) >> value;
129: width.push_back(value);
130: }
131: }
132:
133: // -----
134: // NAME and Charges
135: std::stringstream ss4;
136: if (empty == true) {
137: ss4 = std::stringstream(swidth);
138: } else {
139: ss4 = std::stringstream(sname);
140: }
141: std::string name;
142: ss4 >> name;
143:
144: // Split by comma
145: std::vector<int> chargeX3;
146:
147: while (ss4.good()) {
148: std::string substr;
149: std::getline(ss4, substr, ',');
150:
151: // Remove extra whitespace
152: substr = std::regex_replace(substr, std::regex("^\s+|\s+$"), "");
153:
154: // Identify charge (ORDER is important here!)
155: if (substr.find("++") != std::string::npos) {
156: chargeX3.push_back(6);
157: } else if (substr.find("--") != std::string::npos) {
158: chargeX3.push_back(-6);
159: } else if (substr.find("-1/3") != std::string::npos) {
160: chargeX3.push_back(-1);
161: } else if (substr.find("+2/3") != std::string::npos) {
162: chargeX3.push_back(2);
163: } else if (substr.find("+") != std::string::npos) {
164: chargeX3.push_back(3);
165: } else if (substr.find("-") != std::string::npos) {
166: chargeX3.push_back(-3);

```

```

250: if (p.spinX2 == 0 || p.spinX2 == 2 || p.spinX2 == 4 || p.spinX2 == 8 ||
251:     p.spinX2 == 10) { // Is a boson
252:
253: if (std::to_string(p.pdg).length() == 3 || std::to_string(p.pdg).length() == 5 ||
254:     std::to_string(p.pdg).length() == 6 || std::to_string(p.pdg).length() == 7) {
255:     unsigned int N = std::to_string(p.pdg).length();
256:     unsigned int m = 0;
257:
258: if (N == 5) { m = 0; }
259: if (N == 6) { m = 1; }
260: if (N == 7) { m = 2; }
261:
262: // Code JPC L
263: // L = J+1, S = 1
264: // -----
265: // -----
266:
267: if (N == 3 || (std::to_string(p.pdg)[m] == '0' && std::to_string(p.pdg)[m+1] == '0')) {
268: // 00qgq 1--
269: if (std::to_string(p.pdg)[N-1] == '3') p.setPCL(-1, -1, 0);
270:
271: // 00qgq 2+-
272: if (std::to_string(p.pdg)[N-1] == '5') p.setPCL(1, 1, 1);
273:
274: // 00qgq 3--
275: if (std::to_string(p.pdg)[N-1] == '7') p.setPCL(-1, -1, 2);
276:
277: // 00qgq 4++
278: if (std::to_string(p.pdg)[N-1] == '9') p.setPCL(1, 1, 3);
279: }
280:
281: // L = J, S = 0
282: // -----
283: // -----
284: // 00qg1 0+-
285: if (N == 3 || (std::to_string(p.pdg)[m] == '0' && std::to_string(p.pdg)[m+1] == '0')) {
286: if (std::to_string(p.pdg)[N-1] == '1') p.setPCL(-1, 1, 0);
287: }
288:
289: if (N == 5) {
290: // 10qgq 1+-
291: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[N-1] == '3')
292: p.setPCL(-1, -1, 1);
293:
294: // 10qgq 2+-
295: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[N-1] == '5')
296: p.setPCL(-1, 1, 2);
297:
298: // 10qgq 3+-
299: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[N-1] == '7')
300: p.setPCL(+1, -1, 3);
301:
302: // 10qgq 4+-
303: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[N-1] == '9')
304: p.setPCL(-1, 1, 4);
305:
306: // L = J, S = 1
307: // -----
308: // -----
309: // 20qgq 1+-
310: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[N-1] == '3')
311: p.setPCL(1, 1, 1);
312:
313: // 20qgq 2--
314: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[N-1] == '5')
315: p.setPCL(-1, -1, 2);
316:
317: // 20qgq 3+-
318: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[N-1] == '7')
319: p.setPCL(1, 1, 3);
320:
321: // 20qgq 4--
322: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[N-1] == '9')
323: p.setPCL(-1, -1, 4);
324:
325: // L = J+1, S = 1
326: // -----
327: // -----
328: // 10qg1 0+-
329: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[N-1] == '1')
330: p.setPCL(1, 1, 1);
331:
332: // 30qgq 1--

```

```

167: } else if (substr.find("0") != std::string::npos) {
168: chargeX3.push_back(0);
169: }
170: }
171:
172: // Loop over different charge assignments
173: for (std::size_t k = 0; k < id.size(); ++k) {
174: // New particle
175: gpa:MParticle p;
176:
177: // Add properties
178: p.name = name;
179:
180: p.pdg = id[k];
181: if (mass.size() > 0) { p.mass = mass[0]; }
182: if (width.size() > 0) { // Unstable particles have width
183: p.width = width[0];
184: }
185: p.chargeX3 = chargeX3[k];
186: p.tau = PDG::hbar / p.width; // mean lifetime in the rest frame
187:
188: int lastdigit = p.pdg % 10; // Get last digit
189: p.spinX2 = (lastdigit - 1); // Get 2J
190: p.cut = 0; // Get shell mass (width) cut
191:
192: // Neutral mesons/baryons get 0 for their name
193: if (std::abs(p.chargeX3) == 0 && p.pdg > 100) {
194: p.name = p.name + "0";
195: } else if ((p.chargeX3 != 0) &&
196:           ((p.pdg % 11 && p.pdg % 6) || (p.pdg == 16))) { // quarks & neutrinos
197: p.pdg = 12 || p.pdg == 14 || p.pdg == 16;
198:
199: std::string signstr = p.chargeX3 > 0 ? std::string("abs(p.chargeX3 / 3), '+'") :
200: std::string("abs(p.chargeX3 / 3), '-'");
201: p.name = name + signstr;
202: }
203:
204: // -----
205: // SPECIAL CASES (not following spin numbering)
206:
207: // SM bosons
208: if (p.pdg == 21) // gluon
209: p.spinX2 = 2;
210: p.P = -1;
211: p.C = 0; // Not defined
212:
213: if (p.pdg == 22) // gamma
214: p.spinX2 = 2;
215: p.P = -1;
216: p.C = -1;
217:
218: if (std::abs(p.pdg) == 24) { // W-
219: p.spinX2 = 2;
220: p.P = 0; // Not defined
221: p.C = 0; // Not defined
222: }
223:
224: if (p.pdg == 23) // Z
225: p.spinX2 = 2;
226: p.P = 0; // Not defined
227: p.C = 0; // Not defined
228:
229: if (p.pdg == 25) // H
230: p.spinX2 = 0;
231: p.P = 1; // Scalar SM Higgs
232: p.C = 1; // Scalar SM Higgs
233: }
234:
235: // SM fermions
236: if (p.pdg >= 11 && p.pdg <= 16) { // leptons and neutrinos
237: p.spinX2 = 1;
238: p.P = 1;
239: p.C = 0; // Not defined
240:
241: if (p.pdg >= 14 && p.pdg <= 6) { // quarks
242: p.spinX2 = 1;
243: p.P = 1;
244: p.C = 0; // Not defined
245: }
246: }
247:
248: // Meson JPC (L) assignments
249:
333: if (std::to_string(p.pdg)[m] == '3' && std::to_string(p.pdg)[N-1] == '3')
334: p.setPCL(-1, -1, 2);
335:
336: // 30qg5 2+-
337: if (std::to_string(p.pdg)[m] == '3' && std::to_string(p.pdg)[N-1] == '5')
338: p.setPCL(1, 1, 3);
339:
340: // 30qg7 3--
341: if (std::to_string(p.pdg)[m] == '3' && std::to_string(p.pdg)[N-1] == '7')
342: p.setPCL(-1, -1, 4);
343:
344: // 30qg9 4+-
345: if (std::to_string(p.pdg)[m] == '3' && std::to_string(p.pdg)[N-1] == '9')
346: p.setPCL(1, 1, 5);
347: }
348:
349: // -----
350: // -----
351: // Baryon JPC (L) assignments
352:
353: if (p.spinX2 == 1 || p.spinX2 == 3 || p.spinX2 == 5 || p.spinX2 == 7 ||
354:     p.spinX2 == 9) { // Is a fermion
355:
356: if (std::to_string(p.pdg).length() == 4 || std::to_string(p.pdg).length() == 5 ||
357:     std::to_string(p.pdg).length() == 6 || std::to_string(p.pdg).length() == 7) {
358: // Code JPC L
359: unsigned int N = std::to_string(p.pdg).length();
360: unsigned int m = 0;
361:
362: if (N == 6) { m = 0; }
363: if (N == 7) { m = 1; }
364:
365: // 00qgq2
366: if (N == 4 || (std::to_string(p.pdg)[m] == '0' && std::to_string(p.pdg)[m+1] == '0')) {
367: if (std::to_string(p.pdg)[N-1] == '2') p.setPCL(1, 0, 0);
368: }
369:
370: // 20qgq2
371: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[m+1] == '0') {
372: if (std::to_string(p.pdg)[N-1] == '2') p.setPCL(1, 0, 0);
373: }
374:
375: // 21qgq2
376: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[m+1] == '1') {
377: if (std::to_string(p.pdg)[N-1] == '2') p.setPCL(1, 0, 0);
378: }
379:
380: // 10qgq2
381: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[m+1] == '0') {
382: if (std::to_string(p.pdg)[N-1] == '2') p.setPCL(-1, 0, 1);
383: }
384:
385: // 00qg4
386: if (N == 4 || (std::to_string(p.pdg)[m] == '0' && std::to_string(p.pdg)[m+1] == '0')) {
387: if (std::to_string(p.pdg)[N-1] == '4') p.setPCL(1, 0, 0);
388: }
389:
390: // 20qgq4
391: if (std::to_string(p.pdg)[m] == '2' && std::to_string(p.pdg)[m+1] == '0') {
392: if (std::to_string(p.pdg)[N-1] == '4') p.setPCL(1, 0, 0);
393: }
394:
395: // 11qgq4
396: if (std::to_string(p.pdg)[m] == '1' && std::to_string(p.pdg)[m+1] == '1') {
397: if (std::to_string(p.pdg)[N-1] == '4') p.setPCL(-1, 0, 1);
398: }
399:
400: // Check this case, PDG seems ambiguous here
401: if (N == 5) { p.setPCL(1, 0, 0); }
402: }
403:
404: // -----
405: // -----
406: // Finally, ADD TO THE TABLE
407: PDG::table.insert(std::make_pair(p.pdg, p));
408:
409: // ANTI-PARTICLE (id.size < 3 because 0,+-)
410: if ((!(std::abs(p.chargeX3) > 2.0) || (double)std::abs(p.chargeX3 / 2.0) < id.size() < 3)) ||
411:     (std::abs(p.chargeX3) > 0 && id.size() < 3)) {
412: gpa:MParticle anti(p);
413:
414: // antiquarks & antineutrinos get 1/ide
415: if (1 <= p.pdg && p.pdg <= 6 || (p.pdg == 12 || p.pdg == 14 || p.pdg == 16)) {

```



```

./src/MRandom.cc          2/3
84: if (Gamma <= 1e-40) { return m0; }
85:
86: const double m2 = math::pow2(m0);
87: const double m2max = gra::math::pow2(m0 + LIMIT * Gamma);
88: const double m2min = gra::math::pow2(m0 - LIMIT * Gamma);
89: std::max::max(0.0, math::pow2(M_MIN)), gra::math::pow2(m0 - LIMIT * Gamma));
90:
91: double m2val = 0.0;
92: while (true) {
93:   const double R = U(0, 1);
94:   m2val =
95:     m2 +
96:     m0 * Gamma *
97:     std::itan(std::atan2(m2min - m2, m0 * Gamma) +
98:       R * (std::atan2(m2max - m2, m0 * Gamma) - std::atan2(m2min - m2, m0 * Gamma)));
99:
100: // Note >= handles massless case, otherwise stuck with m = 0
101: if (m2val >= 0) { break; }
102: }
103: return gra::math::msqrt(m2val);
104:
105:
106: // Cauchy (non-relativistic Breit-Wigner) sampling
107: //
108: // Input as with RelativisticCBWRandom
109:
110: double MRandom::CauchyRandom(double m0, double Gamma, double LIMIT, double M_MIN) {
111:   if (Gamma <= 1e-40) { return m0; }
112:
113:   const double mmax = m0 + LIMIT * Gamma;
114:   const double mmin = std::max(std::max(0.0, M_MIN), m0 - LIMIT * Gamma);
115:
116:   double mval = 0.0;
117:   while (true) {
118:     const double R = U(0, 1);
119:     const double value = (Gamma * 0.5) * std::itan(gra::math::PI * (R - 0.5)) + m0;
120:
121:     if (value < mmax && value > mmin) {
122:       mval = value;
123:       break;
124:     }
125:   }
126:   return mval;
127: }
128:
129: // K-dimensional Dirichlet distribution with parameter vector alpha of length K
130: void MRandom::DirichletRandom(const std::vector<double> &alpha, std::vector<double> &ly) {
131:   unsigned int K = alpha.size();
132:
133:   // Draw K independent samples from Gamma(shape = alpha_i, scale = 1)
134:   y.resize(K, 0.0);
135:   for (std::size_t i = 0; i < K; ++i)
136:     std::igamma_distribution<double> gammand(alpha[i], 1.0);
137:   y[i] = gammand(rng); // Draw sample
138:
139:   // Take the sum
140:   double y_sum = 0.0;
141:   for (std::size_t i = 0; i < K; ++i) { y_sum += y[i]; }
142:   // Normalize
143:   for (std::size_t i = 0; i < K; ++i) { y[i] /= y_sum; }
144: }
145:
146: // Negative binomial distribution with parameters avgN and k
147: double MRandom::NBpdf(int n, double avgN, double k) {
148:   // return Chisom(n*k-1, n)
149:   return tgamma(n + 1) / (tgamma(n + 1) * tgamma(k) * std::pow(avgN / (k + avgN), n) *
150:     std::pow(k / (k + avgN), k));
151: }
152:
153: // Random sample from NBD distribution with parameters avgN and k
154: int MRandom::NBRandom(double avgN, double k, int maxvalue) {
155:   const unsigned int MAXTRIAL = 1e7;
156:
157:   // Random integer from [0, Mbins-1]
158:   std::uniform_int_distribution<int> RANDI(1, maxvalue);
159:
160:   // Acceptance-Rejection
161:   unsigned int trials = 0;
162:   while (true) {
163:     const int n = RANDI(rng);
164:     const double val = NBpdf(n, avgN, k);
165:     if (U(0, 1) < val) { return n; }
166:     ++trials;
167:   }
168: }

```

```

./src/MRandom.cc          3/3
167: if (trials > MAXTRIAL) {
168:   return avgN; // Return the mean value
169: }
170: }
171: }
172:
173: // Log-distribution with with parameter p
174: double MRandom::Logpdf(int k, double p) {
175:   return -1.0 / std::ilog1p(-p) * std::pow(p, k) / static_cast<double>(k);
176: }
177:
178: // Random sample from log-distribution with parameter p
179: int MRandom::LogRandom(double p, int maxvalue) {
180:   // Random integer from [0, Mbins-1]
181:   std::uniform_int_distribution<int> RANDI(0, maxvalue - 1);
182:
183:   // Acceptance-Rejection
184:   while (true) {
185:     const int n = RANDI(rng);
186:     const double val = Logpdf(n, p);
187:     if (U(0, 1) < val) { return n; }
188:   }
189: }
190:
191: // namespace gra
192:

```

```

./src/MTensorPomeron.cc 1/32
1: // Tensor Pomeron amplitudes
2: //
3: //
4: // [REFERENCE: Ewerz, Maniatis, Nachtmann, arxiv.org/abs/1309.3478]
5: // [REFERENCE: Labiodowicz, Nachtmann, Szczurek, arxiv.org/abs/1601.04537]
6: // [REFERENCE: LNS, arxiv.org/abs/1801.03902]
7: // [REFERENCE: LNS, arxiv.org/abs/1901.11490]
8: //
9: // (c) 2017-2020 Mikael Mieskolainen
10: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
11:
12: // C++
13: #include <complex>
14: #include <iostream>
15: #include <vector>
16:
17: // Own
18: #include "GraniTCL/M4Vec.h"
19: #include "GraniTCL/M4Vec.h"
20: #include "GraniTCL/MDIrac.h"
21: #include "GraniTCL/MFrac.h"
22: #include "GraniTCL/MHelMomematics.h"
23: #include "GraniTCL/MBach.h"
24: #include "GraniTCL/MWgts.h"
25: #include "GraniTCL/MWgts.h"
26: #include "GraniTCL/MTensorPomeron.h"
27:
28: // FTensor algebra
29: #include "FTensor.hpp"
30:
31: // LOOP MACROS
32: #define FOR_SACH_2(X) \
33:   for (const auto &u : X) { \
34:     for (const auto &v : X) { \
35:       \
36:     } \
37:   }
38:
39: #define FOR_SACH_3(X) \
40:   for (const auto &u : X) { \
41:     for (const auto &v : X) { \
42:       for (const auto &k : X) { \
43:         \
44:       } \
45:     } \
46:   }
47:
48: #define FOR_SACH_4(X) \
49:   for (const auto &u : X) { \
50:     for (const auto &v : X) { \
51:       for (const auto &k : X) { \
52:         for (const auto &l : X) { \
53:           \
54:         } \
55:       } \
56:     } \
57:   }
58:
59: #define FOR_SACH_5(X) \
60:   for (const auto &u : X) { \
61:     for (const auto &v : X) { \
62:       for (const auto &k : X) { \
63:         for (const auto &l : X) { \
64:           for (const auto &r : X) { \
65:             \
66:           } \
67:         } \
68:       } \
69:     } \
70:   }
71:
72: #define FOR_SACH_6(X) \
73:   for (const auto &u : X) { \
74:     for (const auto &v : X) { \
75:       for (const auto &k : X) { \
76:         for (const auto &l : X) { \
77:           for (const auto &r : X) { \
78:             for (const auto &s : X) { \
79:               \
80:             } \
81:           } \
82:         } \
83:       } \

```

```

./src/MTensorPomeron.cc 2/32
84: } \
85: } \
86: }
87: #define FOR_PP_HELICITY \
88:   for (const auto &ha : {0, 1}) { \
89:     for (const auto &hb : {0, 1}) { \
90:       for (const auto &hl : {0, 1}) { \
91:         for (const auto &h2 : {0, 1}) { \
92:           #define FOR_PP_HELICITY_END \
93:           \
94:         } \
95:       } \
96:     } \
97:   }
98: using gra::aux::indices;
99: using gra::math::PI;
100: using gra::math::msqrt;
101: using gra::math::pow2;
102: using gra::math::sin;
103:
104: using FTensor::Tensor0;
105: using FTensor::Tensor1;
106: using FTensor::Tensor2;
107: using FTensor::Tensor3;
108: using FTensor::Tensor4;
109:
110: namespace gra {
111:
112: // Constructor
113: MTensorPomeron::MTensorPomeron(gra::LORENTZSCALAR &Its, const std::string &modelFile) {
114:   CalcTensor(); // Pre-Calculate tensors
115:
116: // ## MULTITHREADING LOCK NEEDED FOR THE INITIALIZATION ##
117:   gra::ig_mutex.lock();
118:
119:   if (!Param.initialized) {
120:     try {
121:       Param.ReadParameters(modelFile);
122:       PARAM_REGIOE::ReadParameters(0, modelFile);
123:
124:     } catch (...) {
125:       gra::ig_mutex.unlock(); // need to release here, otherwise get infinite lock
126:       throw;
127:     }
128:   }
129:   gra::ig_mutex.unlock();
130: }
131:
132: // Return decay coupling constant for resonance (M_Gamma) with decay daughter mass m
133: // BR being the branching ratio BR = Width_partial / Width_Total
134: // Decay: Mother (spin = 0/1/2) -> scalar or pseudoscalar daughters
135: //
136: double MTensorPomeron::GDecay(int J, double M, double Gamma, double BR) {
137:   const double SO = 1.0; // Should be set as the same scale as in decay amplitudes [G]
138:   const double partialWidth = Gamma * BR;
139:
140:   const double P = sqrt(1 - 4 * pow2(mf / M));
141:
142:   if (J == 0) {
143:     return sqrt(partialWidth / (1.0 / (16 * PI * M) * pow2(SO * P)));
144:   } else if (J == 1) {
145:     return sqrt(partialWidth / (M / (192 * PI * M) * math::ipow3(P)));
146:   } else if (J == 2) {
147:     return sqrt(partialWidth / (M / (480 * PI) * pow2(M / SO) * math::ipow5(P)));
148:   } else if (J == 4) {
149:     return 1.0; // Future
150:   } else {
151:     throw std::invalid_argument("MTensorPomeron::GDecay: Unknown input spin J = " +
152:       std::ito_string(J));
153:   }
154: }
155:
156: // 2 -> 3 amplitudes
157: //
158: // return value: matrix element complex helicities summed and averaged over
159: // Its.hamp: individual complex helicity amplitudes
160: //
161: double MTensorPomeron::M3(gra::LORENTZSCALAR &Its) const {
162:   // Two tensor indices (second parameter denotes the range of index)
163:   FTensor::Index<'a', 4> mu;

```



```

499:         ID.push_back(temp);
500:     }
501: }
502: } else {
503:     throw std::invalid_argument("MTensorPomeron:ME3: Unknown decay mode for tensor resonance");
504: }
505: }
506: // Over all central helicity combinations
507: Tensor<std::complex<double>, 4, 4, 4, 4> temp;
508: std::vector<Tensor<std::complex<double>, 4, 4, 4, 4>> cvtx;
509: for (const auto kind : indices(ID)) {
510:     // Construct central vertex
511:     FOR_EACH_4(1);
512:     temp[0][0][0][0] = 0.0;
513:     for (auto const k1ho : I1) {
514:         for (auto const sigma1 : I1) {
515:             temp[0][v, k, l] += GPPF2((u, v, k, l, rho, sigma)) * ID[ind](rho, sigma);
516:         }
517:     }
518:     FOR_EACH_4_END;
519:     cvtx.push_back(temp);
520: }
521: }
522: // Two helicity states for incoming and outgoing protons
523: std::size_t index = 0;
524: FOR_PP_HELICITY;
525: }
526: // Apply proton leg helicity conservation / No helicity flip (high energy limit)
527: if (ha != hb || hb != hc) continue;
528: }
529: // Full proton-Pomeron-proton spinor structure (upper and lower vertex)
530: const Tensor<std::complex<double>, 4, 4> iG_L1 = iG_Ppp(p1, pa, ubar_1[h1], u_a[h1]);
531: const Tensor<std::complex<double>, 4, 4> iG_L2 = iG_Ppp(p2, pb, ubar_2[h2], u_b[h2]);
532: }
533: // s-channel amplitude
534: for (const auto kind : indices(cvtx)) {
535:     const std::complex<double> A = (-z1) * iG_L1[mu1, nu1] * iD_P[mu, nu, alpha, beta] *
536:         iD_P_2(alpha, beta, mu2, nu2) * iG_L2[mu2, nu2];
537:     // Add it
538:     its.hamp[index] += A; // Note ==
539:     ++index;
540: }
541: }
542: FOR_PP_HELICITY_END;
543: maxindex = index > maxindex ? index : maxindex;
544: } else {
545:     throw std::invalid_argument("MTensorPomeron:ME3: Unknown spin-parity input");
546: }
547: }
548: }
549: // =====
550: }
551: // Loop over resonances
552: // =====
553: // Remove empty
554: its.hamp.resize(maxindex);
555: }
556: // Get total amplitude squared 1/4 |sum_h[A_A]|^2
557: double SumAmp2 = 0.0;
558: for (const auto k1 : indices(its.hamp)) { SumAmp2 += gra:math:abs2(its.hamp[l1]); }
559: SumAmp2 /= 4; // Initial state helicity average
560: }
561: }
562: return SumAmp2; // Amplitude squared
563: }
564: }
565: // 2 -> 4 amplitudes
566: }
567: // return value: matrix element squared with helicities summed over
568: // its.hamp: individual helicity amplitudes
569: }
570: }
571: // =====
572: double MTensorPomeron::ME4(gra:LORENTZSCALAR fits) const {
573:     // Kinematics
574:     const MVec pa = its.pbeam1;
575:     const MVec pb = its.pbeam2;
576:     const MVec p1 = its.pfinal[1];
577:     const MVec p2 = its.pfinal[2];
578:     const MVec p3 = its.pdecay[0].p4;
579:     const MVec p4 = its.pdecay[1].p4;
580:     }
581: }

```

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665: // Lepton or quark pair via two photon fusion
666: if (SPINMODE == "2xS") {
667:     // Full proton-gamma-proton spinor structure (upper and lower vertex)
668:     const Tensor<std::complex<double>, 4, 4> iG_L1 = iG_Ppp(p1, pa, ubar_1[h1], u_a[h1]);
669:     const Tensor<std::complex<double>, 4, 4> iG_L2 = iG_Ppp(p2, pb, ubar_2[h2], u_b[h2]);
670:     // Photon propagators
671:     const Tensor<std::complex<double>, 4, 4> iD_L1 = iD_P[its.t1];
672:     const Tensor<std::complex<double>, 4, 4> iD_L2 = iD_P[its.t2];
673:     // Fermion propagator
674:     const MMatrix<std::complex<double>, iSF_L = iD_P(p1, M_L);
675:     const MMatrix<std::complex<double>, iSF_U = iD_P(pu, M_U);
676:     // Central spinors (2 helicities)
677:     const std::array<std::vector<std::complex<double>, 2>> v_3 = SpinorStates(p3, "u");
678:     const std::array<std::vector<std::complex<double>, 2>> ubar_4 = SpinorStates(p4, "ubar");
679:     // =====
680:     double FACTOR = 1.0;
681:     // quark pair (charge 1/3 or 2/3), apply charge and color factors
682:     if (std::abs(its.decaytree[0].p.pdg) == 6) { // we have a quark
683:         const double Q = its.decaytree[0].p.chargeX3 / 3.0;
684:         const double NC = 3.0; // quarks come in three colors
685:         FACTOR = sqrt(math::ipow4(Q * NC)); // sqrt to "amplitude level"
686:     }
687:     // =====
688:     for (const auto k1 : indices(v_3)) {
689:         for (const auto k2 : indices(ubar_4)) {
690:             // Vertex functions
691:             const Tensor<std::complex<double>, 4, 4> iYF_U = iG_Yeebar(ubar_4[h4], iSF_U, v_3[h3]);
692:             const Tensor<std::complex<double>, 4, 4> iYF_U = iG_Yeebar(ubar_4[h4], iSF_U, v_3[h3]);
693:             // Full amplitude: t-channel + u-channel
694:             std::complex<double> M_U = (-z1) * iG_L1[mu1, nu1] * iD_P[mu, nu, alpha, beta] *
695:                 iYF_U[mu2, nu2] + iYF_U[mu1, nu2] * iD_P_2(mu2, nu2) *
696:                 iG_L2[mu2, nu2];
697:             M_U *= FACTOR;
698:             its.hamp.push_back(M_U);
699:         }
700:     }
701:     // =====
702:     continue; // skip parts below, only for Pomeron amplitudes
703: }
704: // =====
705: // Full proton-Pomeron-proton spinor structure (upper and lower vertex)
706: const Tensor<std::complex<double>, 4, 4> iG_L1 = iG_Ppp(p1, pa, ubar_1[h1], u_a[h1]);
707: const Tensor<std::complex<double>, 4, 4> iG_L2 = iG_Ppp(p2, pb, ubar_2[h2], u_b[h2]);
708: // =====
709: // 2 x pseudoscalar (pion pair, kaon pair ...)
710: if (SPINMODE == "2xS") {
711:     double gP = 0.0;
712:     // Pion or Kaon coupling
713:     gP = (its.decaytree[0].p.mass < 0.2) ? Param.gtpi1 : Param.gPKK;
714:     // t-channel blocks
715:     const Tensor<std::complex<double>, 4, 4> iG_T = iG_Ppp(p1, -p3, gP);
716:     const std::array<std::vector<std::complex<double>, 2>> iDMS_U = iD_MES(pu, M_U);
717:     const Tensor<std::complex<double>, 4, 4> iG_TB = iG_Ppp(p4, pt, gP);
718:     // u-channel blocks
719:     const Tensor<std::complex<double>, 4, 4> iG_U = iG_Ppp(p4, pu, gP);
720:     const std::array<std::vector<std::complex<double>, 2>> iDMS_U = iD_MES(pu, M_U);
721:     const Tensor<std::complex<double>, 4, 4> iG_UB = iG_Ppp(pu, -p3, gP);
722:     // =====
723:     std::complex<double> M_U;
724:     std::complex<double> M_U;
725:     // t-channel
726:     M_U = iG_L1[mu1, nu1] * iD_P_13(mu1, nu1, alpha, beta) * iG_T(alpha, beta, mu2, nu2) *
727:         iD_P_2(alpha, beta, mu2, nu2) * iG_UB(mu2, nu2);
728:     // u-channel
729:     M_U = iG_L1[mu1, nu1] * iD_P_13(mu1, nu1, alpha, beta) * iG_U(alpha, beta, mu2, nu2) *
730:         iD_P_2(alpha, beta, mu2, nu2) * iG_TB(mu2, nu2);
731:     // =====
732:     double g1 = 0.0;
733:     double g2 = 0.0;
734:     // Couplings (rho770 meson)
735:     if (pdg == 113) {
736:         g1 = Param.gFrho[0];
737:         g2 = Param.gFrho[1];
738:     }
739:     // Couplings (phi1020 meson)
740:     else if (pdg == 333) {
741:         g1 = Param.gPhi[0];
742:         g2 = Param.gPhi[1];
743:     }
744:     // =====
745:     FTensor::Index<'t', 4> rho4;
746:     FTensor::Index<'u', 4> rho4;
747:     // t-channel blocks
748:     const Tensor<std::complex<double>, 4, 4, 4, 4> iD_Pvt = iD_Pvt(p1, -p3, g1, g2);
749:     const Tensor<std::complex<double>, 4, 4, 4, 4> iD_Uvt = iD_Uvt(pu, M_U, its.pfinal[0].M2(1));

```

```

582: // Intermediate boson/fermion mass
583: const double M_ = its.decaytree[0].p.mass;
584: }
585: // Momentum convention of sub-diagrams
586: // =====
587: // anti-particle p4
588: // \hat{t} (arrow down)
589: // ----- particle p4
590: // ----- particle p4
591: // \hat{u} (arrow up)
592: // ----- anti-particle p3
593: // =====
594: const MVec pa = p1 - p1 - p3; // => q1 = pt + p3, q2 = pt - p4
595: const MVec pa = p4 - pa + p1; // => q2 = pu + p1, q1 = pu - p3
596: // =====
597: // Incoming and outgoing proton spinors (2 helicities)
598: const std::array<std::vector<std::complex<double>, 2>> u_a = SpinorStates(pa, "u");
599: const std::array<std::vector<std::complex<double>, 2>> u_b = SpinorStates(pb, "u");
600: const std::array<std::vector<std::complex<double>, 2>> ubar_1 = SpinorStates(p1, "ubar");
601: const std::array<std::vector<std::complex<double>, 2>> ubar_2 = SpinorStates(p2, "ubar");
602: // =====
603: // t-channel Pomeron propagators
604: const Tensor<std::complex<double>, 4, 4, 4, 4> iD_P_13 = iD_P[its.s1][13], its.t1;
605: const Tensor<std::complex<double>, 4, 4, 4, 4> iD_P_24 = iD_P[its.s2][14], its.t2;
606: // =====
607: // u-channel Pomeron propagators
608: const Tensor<std::complex<double>, 4, 4, 4, 4> iD_P_14 = iD_P[its.s1][14], its.t1;
609: const Tensor<std::complex<double>, 4, 4, 4, 4> iD_P_23 = iD_P[its.s2][13], its.t2;
610: // =====
611: // High Energy Limit proton-Pomeron-proton spinor structure
612: const Tensor<std::complex<double>, 4, 4> iG_L1 = iG_PppHE(p1, pa);
613: const Tensor<std::complex<double>, 4, 4> iG_L2 = iG_PppHE(p2, pb);
614: // =====
615: // Reset
616: its.hamp.clear();
617: // =====
618: // Free Lorentz indices [second parameter denotes the range of index]
619: FTensor::Index<'b', 4> nu1;
620: FTensor::Index<'c', 4> rho1;
621: FTensor::Index<'d', 4> rho2;
622: FTensor::Index<'g', 4> alpha;
623: FTensor::Index<'h', 4> beta;
624: FTensor::Index<'i', 4> alpha2;
625: FTensor::Index<'j', 4> beta2;
626: FTensor::Index<'k', 4> mu2;
627: FTensor::Index<'l', 4> nu2;
628: // =====
629: // DEUCE subprocesses
630: std::string SPINMODE;
631: // =====
632: if (its.decaytree[0].p.spinX2 == 0 && its.decaytree[1].p.spinX2 == 0) {
633:     SPINMODE = "2xS";
634: } else if (its.decaytree[0].p.spinX2 == 1 && its.decaytree[1].p.spinX2 == 1) {
635:     SPINMODE = "2xP";
636: } else if (its.decaytree[0].p.spinX2 == 2 && its.decaytree[1].p.spinX2 == 2) {
637:     SPINMODE = "2xV";
638: } else {
639:     throw std::invalid_argument("
640:         MTensorPomeron:ME4: Invalid daughter spin (J = 0, 1/2, 1 pairs supported)");
641: }
642: // =====
643: // Two helicity states for incoming and outgoing protons
644: FOR_PP_HELICITY;
645: // Apply proton leg helicity conservation / No helicity flip
646: // (high energy limit)
647: // This gives at least 4 x speed improvement
648: if (ha != hb || hb != hc) continue;
649: // =====
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./src/MTensorPomeron.cc 19/32
1495: // Form FACTOR
1496: const double LAMBDA = 1.0; // GeV
1497: const double F_rilde = FM(q1.M2(), 1.0) * FM(q2.M2(), 1.0) * F((q1 + q2).M2(), M0, LAMBDA);
1498:
1499: // Tensor structure
1500: Tensor<std::complex<double>, 4, 4, 4, 4> T;
1501:
1502: if (mode == "scalar") {
1503:     static const Tensor<std::complex<double>, 4, 4, 4, 4> IGO = MTensorPomeron::iG_PPFS_0();
1504:     const Tensor<std::complex<double>, 4, 4, 4, 4> IGI =
1505:         MTensorPomeron::iG_PPFS_1(q1, q2, g_PPFS(1));
1506:
1507:     FOR_EACH_4(L1);
1508:     T(u, v, k, l) = (g_PPFS[0] * IGO(u, v, k, l) + IGI(u, v, k, l)) * F_rilde;
1509:     FOR_EACH_4_END;
1510:
1511: } else if (mode == "pseudoscalar") {
1512:     const Tensor<std::complex<double>, 4, 4, 4, 4> IGO =
1513:         MTensorPomeron::iG_PPFS_0(q1, q2, g_PPFS(1));
1514:     const Tensor<std::complex<double>, 4, 4, 4, 4> IGI =
1515:         MTensorPomeron::iG_PPFS_1(q1, q2, g_PPFS(1));
1516:
1517:     FOR_EACH_4(L1);
1518:     T(u, v, k, l) = (IGO(u, v, k, l) + IGI(u, v, k, l)) * F_rilde;
1519:     FOR_EACH_4_END;
1520:
1521: } else {
1522:     throw std::invalid_argument(
1523:         "MTensorPomeron::iG_PPFS_Total: Unknown mode (should be scalar or pseudoscalar)");
1524: }
1525:
1526: return T;
1527: }
1528:
1529: // Pomeron-Pomeron-Scalar coupling structure #0
1530: // iG_1(u mu nu kappa lambda) ~ (1,s) = (0,0)
1531: //
1532: // Input as contravariant (upper index) 4-vectors
1533: //
1534: // Apply coupling g_PPFS[0] outside this!
1535: //
1536: Tensor<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iG_PPFS_0(const
1537:     const double S0 = 1.0; // Mass scale (GeV)
1538:     const std::complex<double> FACTOR = z1 * g_PPFS / (2 * S0);
1539:
1540:     Tensor<std::complex<double>, 4, 4, 4, 4> T;
1541:     FOR_EACH_4(L1);
1542:     T(u, v, k, l) = FACTOR * (g[u][k] * g[v][l] + g[u][l] * g[v][k]) - 0.5 * g[u][v] * g[k][l];
1543:     FOR_EACH_4_END;
1544:
1545:     return T;
1546: }
1547:
1548: // Pomeron-Pomeron-Scalar coupling structure #1
1549: // iG_1(u mu nu kappa lambda) ~ (1,s) = (2,2)
1550: //
1551: // Input as contravariant (upper index) 4-vectors
1552: //
1553: Tensor<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iG_PPFS_1(const M4Vec q1, const M4Vec q2,
1554:     const double S0 = 1.0; // Mass scale (GeV)
1555:     const double q1q2 = q1 * q2;
1556:     const std::complex<double> FACTOR = z1 * g_PPFS / (2 * S0);
1557:
1558:     Tensor<std::complex<double>, 4, 4, 4, 4> T;
1559:     FOR_EACH_4(L1);
1560:     T(u, v, k, l) = FACTOR * ((q1[k] * (q2[l] * g[u][l] + (q1[k] * (q2[l] * g[u][l] +
1561:         (q1[l] * (q2[k] * g[u][k] + (q1[l] * (q2[k] * g[u][k] -
1562:         2.0 * g[u][v] * g[k][l]))))))) * g[u][l] + g[v][k] * g[u][l]);
1563:     FOR_EACH_4_END;
1564:
1565:     return T;
1566: }
1567:
1568:
1569: // Pomeron-Pomeron-Pseudoscalar coupling structure #0
1570: // iG_1(u mu nu kappa lambda) ~ (1,s) = (1,1)
1571: //
1572: // Input as contravariant (upper index) 4-vectors
1573: //
1574: Tensor<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iG_PPFS_0(const M4Vec q1,
1575:     const M4Vec q2,
1576:     const double S0 = 1.0; // Mass scale (GeV)
1577:
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1594: // Pomeron-Pomeron-Tensor coupling structure #0
1595: // iG_1(u mu nu kappa lambda rho sigma) ~ (1,s) = (0,2)
1596: //
1597: // This is kinematics independent, can be pre-calculated.
1598: //
1599: // Remember to apply g_PPT (coupling constant) outside this
1600: //
1601: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_00(const
1602:     const double S0 = 1.0; // Mass scale (GeV)
1603:     const std::complex<double> FACTOR = z1 * S0 * S0;
1604:
1605:     // Init with zeros!
1606:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1607:
1608:     // Contract total vertex by summing up the tensors
1609:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1610:
1611:     FOR_EACH_6(L1);
1612:
1613:     // Indexing
1614:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1615:
1616:     // For number 10, we apply coupling here
1617:     T(ind) = IGO(ind) * g_PPT[0];
1618:
1619:     // For number 1-6, couplings have been already applied, just loop over
1620:     for (const auto i1 : indices(IGM)) T(ind) += IGM[i1(ind)];
1621:
1622:     // Apply form factors
1623:     T(ind) = F_rilde;
1624:     FOR_EACH_6_END;
1625:
1626:     return T;
1627: }
1628:
1629:
1630: // Pomeron-Pomeron-Tensor coupling structure #1
1631: // iG_1(u mu nu kappa lambda rho sigma) ~ (1,s) = (2,2)
1632: //
1633: // This is kinematics independent, can be pre-calculated.
1634: //
1635: // Remember to apply g_PPT (coupling constant) outside this
1636: //
1637: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_10(const M4Vec q1, const M4Vec q2,
1638:     const double S0 = 1.0; // Mass scale (GeV)
1639:     const std::complex<double> FACTOR = z1 * S0 * S0;
1640:
1641:     // Init with zeros!
1642:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1643:
1644:     // Contract total vertex by summing up the tensors
1645:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1646:
1647:     FOR_EACH_6(L1);
1648:
1649:     // Indexing
1650:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1651:
1652:     for (auto const i1 : I1) {
1653:         for (auto const i2 : I2) {
1654:             for (auto const i3 : I3) {
1655:                 for (auto const i4 : I4) {
1656:                     for (auto const i5 : I5) {
1657:                         for (auto const i6 : I6) {
1658:                             T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1659:                         }
1660:                     }
1661:                 }
1662:             }
1663:         }
1664:     }
1665:
1666:     // Apply form factors
1667:     T(ind) = F_rilde;
1668:     FOR_EACH_6_END;
1669:
1670:     return T;
1671: }
1672:
1673: // Pomeron-Pomeron-Tensor coupling structures #1 and #2
1674: // iG_1(u mu nu kappa lambda rho sigma)
1675: // ~ (1,s) = (2,0) + (2,2)
1676: // ~ (1,s) = (0,0) + (2,2)
1677: //
1678: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_12(const M4Vec q1, const M4Vec q2,
1679:     const double S0 = 1.0; // Mass scale (GeV)
1680:     const std::complex<double> FACTOR = z1 * S0 * S0;
1681:
1682:     // Init with zeros!
1683:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1684:
1685:     // Contract total vertex by summing up the tensors
1686:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1687:
1688:     FOR_EACH_6(L1);
1689:
1690:     // Indexing
1691:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1692:
1693:     for (auto const i1 : I1) {
1694:         for (auto const i2 : I2) {
1695:             for (auto const i3 : I3) {
1696:                 for (auto const i4 : I4) {
1697:                     for (auto const i5 : I5) {
1698:                         for (auto const i6 : I6) {
1699:                             T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1700:                         }
1701:                     }
1702:                 }
1703:             }
1704:         }
1705:     }
1706:
1707:     // Apply form factors
1708:     T(ind) = F_rilde;
1709:     FOR_EACH_6_END;
1710:
1711:     return T;
1712: }
1713:
1714: // Pomeron-Pomeron-Tensor coupling structure #3
1715: // iG_1(u mu nu kappa lambda rho sigma) ~ (1,s) = (2,4)
1716: //
1717: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_03(const M4Vec q1, const M4Vec q2,
1718:     const double S0 = 1.0; // Mass scale (GeV)
1719:     const std::complex<double> FACTOR = z1 * S0 * S0;
1720:
1721:     // Init with zeros!
1722:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1723:
1724:     // Contract total vertex by summing up the tensors
1725:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1726:
1727:     FOR_EACH_6(L1);
1728:
1729:     // Indexing
1730:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1731:
1732:     for (auto const i1 : I1) {
1733:         for (auto const i2 : I2) {
1734:             for (auto const i3 : I3) {
1735:                 for (auto const i4 : I4) {
1736:                     for (auto const i5 : I5) {
1737:                         for (auto const i6 : I6) {
1738:                             T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1739:                         }
1740:                     }
1741:                 }
1742:             }
1743:         }
1744:     }
1745:
1746:     // Apply form factors
1747:     T(ind) = F_rilde;
1748:     FOR_EACH_6_END;
1749:
1750:     return T;
1751: }
1752:
1753: // Final rank-6 tensor
1754: MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1755:
1756: FOR_EACH_6(L1);
1757:
1758: const std::vector<size_t> ind = {u, v, k, l, r, s};
1759:
1760: for (auto const i1 : I1) {
1761:     for (auto const i2 : I2) {
1762:         for (auto const i3 : I3) {
1763:             for (auto const i4 : I4) {
1764:                 for (auto const i5 : I5) {
1765:                     for (auto const i6 : I6) {
1766:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1767:                     }
1768:                 }
1769:             }
1770:         }
1771:     }
1772: }
1773:
1774: // Apply form factors
1775: T(ind) = F_rilde;
1776: FOR_EACH_6_END;
1777:
1778: return T;
1779: }
1780:
1781: // NAIVE LOOP
1782: for (auto const i1 : I1) {
1783:     for (auto const i2 : I2) {
1784:         for (auto const i3 : I3) {
1785:             for (auto const i4 : I4) {
1786:                 for (auto const i5 : I5) {
1787:                     for (auto const i6 : I6) {
1788:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1789:                     }
1790:                 }
1791:             }
1792:         }
1793:     }
1794: }
1795:
1796: // Apply form factors
1797: T(ind) = F_rilde;
1798: FOR_EACH_6_END;
1799:
1800: return T;
1801: }

```

```

./src/MTensorPomeron.cc 20/32
1578: const M4Vec q1_q2 = q1 - q2;
1579: const M4Vec q1 = q1 + q2;
1580:
1581: const std::complex<double> FACTOR = z1 * g_PPFS / (2.0 * S0);
1582: Tensor<std::complex<double>, 4, 4, 4, 4> T;
1583:
1584: FOR_EACH_4(L1);
1585: T(u, v, k, l) = 0.0;
1586:
1587: // Contract r and s indices
1588: for (const auto i1 : I1) {
1589:     for (const auto i2 : I2) {
1590:         for (const auto i3 : I3) {
1591:             for (const auto i4 : I4) {
1592:                 T(u, v, k, l) = FACTOR *
1593:                     (g[u][l] * eps_lo1(u, l, r, s) + g[v][k] * eps_lo1(u, l, r, s) +
1594:                     g[u][l] * eps_lo1(v, k, r, s) + g[v][l] * eps_lo1(u, k, r, s)) *
1595:                     q1_q2[i1] * q[i2];
1596:             }
1597:         }
1598:     }
1599: }
1600: return T;
1601:
1602: // Pomeron-Pomeron-Pseudoscalar coupling structure #1
1603: // iG_1(u mu nu kappa lambda) ~ (1,s) = (2,2)
1604: //
1605: // Input as contravariant (upper index) 4-vectors
1606: //
1607: Tensor<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iG_PPFS_1(const M4Vec q1, const M4Vec q2,
1608:     const double S0 = 1.0; // Mass scale (GeV)
1609:     const double q1q2 = q1 * q2;
1610:     const std::complex<double> FACTOR = z1 * g_PPFS / gra:math:pow(3);
1611:
1612:     Tensor<std::complex<double>, 4, 4, 4, 4> T;
1613:     FOR_EACH_4(L1);
1614:     T(u, v, k, l) = 0.0;
1615:
1616:     // Contract r and s indices
1617:     for (const auto i1 : I1) {
1618:         for (const auto i2 : I2) {
1619:             for (const auto i3 : I3) {
1620:                 for (const auto i4 : I4) {
1621:                     T(u, v, k, l) = FACTOR *
1622:                         (eps_lo1(u, l, r, s) * ((q1[k] * (q2[l] * g[u][l] + (q1[k] * (q2[l] * g[u][l] +
1623:                             (q1[l] * (q2[k] * g[u][k] + (q1[l] * (q2[k] * g[u][k] -
1624:                             2.0 * g[u][v] * g[k][l]))))))) * g[u][l] + g[v][k] * g[u][l]) +
1625:                         eps_lo1(v, k, r, s) * ((q1[l] * (q2[k] * g[u][k] + (q1[l] * (q2[k] * g[u][k] -
1626:                             2.0 * g[u][v] * g[k][l])))) * q1_q2[i1] * q[i2];
1627:                     }
1628:                 }
1629:             }
1630:         }
1631:     }
1632:
1633:     return T;
1634: }
1635:
1636: // Pomeron-Pomeron-Tensor coupling structure #3
1637: // iG_1(u mu nu kappa lambda rho sigma) ~ (1,s) = (2,4)
1638: //
1639: // Remember to apply g_PPT (coupling constant) outside this
1640: //
1641: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_03(const M4Vec q1, const M4Vec q2,
1642:     const double S0 = 1.0; // Mass scale (GeV)
1643:     const std::complex<double> FACTOR = z1 * S0 * S0;
1644:
1645:     // Init with zeros!
1646:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1647:
1648:     // Contract total vertex by summing up the tensors
1649:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1650:
1651:     FOR_EACH_6(L1);
1652:
1653:     // Indexing
1654:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1655:
1656:     for (auto const i1 : I1) {
1657:         for (auto const i2 : I2) {
1658:             for (auto const i3 : I3) {
1659:                 for (auto const i4 : I4) {
1660:                     for (auto const i5 : I5) {
1661:                         for (auto const i6 : I6) {
1662:                             T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1663:                         }
1664:                     }
1665:                 }
1666:             }
1667:         }
1668:     }
1669:
1670:     // Apply form factors
1671:     T(ind) = F_rilde;
1672:     FOR_EACH_6_END;
1673:
1674:     return T;
1675: }
1676:
1677: // Final rank-6 tensor
1678: MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1679:
1680: FOR_EACH_6(L1);
1681:
1682: const std::vector<size_t> ind = {u, v, k, l, r, s};
1683:
1684: for (auto const i1 : I1) {
1685:     for (auto const i2 : I2) {
1686:         for (auto const i3 : I3) {
1687:             for (auto const i4 : I4) {
1688:                 for (auto const i5 : I5) {
1689:                     for (auto const i6 : I6) {
1690:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1691:                     }
1692:                 }
1693:             }
1694:         }
1695:     }
1696: }
1697:
1698: // Apply form factors
1699: T(ind) = F_rilde;
1700: FOR_EACH_6_END;
1701:
1702: return T;
1703: }
1704:
1705: // NAIVE LOOP
1706: for (auto const i1 : I1) {
1707:     for (auto const i2 : I2) {
1708:         for (auto const i3 : I3) {
1709:             for (auto const i4 : I4) {
1710:                 for (auto const i5 : I5) {
1711:                     for (auto const i6 : I6) {
1712:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1713:                     }
1714:                 }
1715:             }
1716:         }
1717:     }
1718: }
1719:
1720: // Apply form factors
1721: T(ind) = F_rilde;
1722: FOR_EACH_6_END;
1723:
1724: return T;
1725: }

```

```

./src/MTensorPomeron.cc 21/32
1661: if (std::abs(g_PP_T[1]) > EPSILON) IGM.push_back(iG_PP_T_12(q1, q2, g_PPT[1], 1));
1662:
1663: if (std::abs(g_PP_T[2]) > EPSILON) IGM.push_back(iG_PP_T_12(q1, q2, g_PPT[2], 2));
1664:
1665: if (std::abs(g_PP_T[3]) > EPSILON) IGM.push_back(iG_PP_T_03(q1, q2, g_PPT[3]));
1666:
1667: if (std::abs(g_PP_T[4]) > EPSILON) IGM.push_back(iG_PP_T_04(q1, q2, g_PPT[4]));
1668:
1669: if (std::abs(g_PP_T[5]) > EPSILON) IGM.push_back(iG_PP_T_05(q1, q2, g_PPT[5]));
1670:
1671: if (std::abs(g_PP_T[6]) > EPSILON) IGM.push_back(iG_PP_T_06(q1, q2, g_PPT[6]));
1672:
1673: // Contract total vertex by summing up the tensors
1674: MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1675:
1676: FOR_EACH_6(L1);
1677:
1678: // Indexing
1679: const std::vector<size_t> ind = {u, v, k, l, r, s};
1680:
1681: // For number 10, we apply coupling here
1682: T(ind) = IGO(ind) * g_PPT[0];
1683:
1684: // For number 1-6, couplings have been already applied, just loop over
1685: for (const auto i1 : indices(IGM)) T(ind) += IGM[i1(ind)];
1686:
1687: // Apply form factors
1688: T(ind) = F_rilde;
1689: FOR_EACH_6_END;
1690:
1691: return T;
1692: }
1693:
1694: // Pomeron-Pomeron-Tensor coupling structure #0
1695: // iG_1(u mu nu kappa lambda rho sigma) ~ (1,s) = (0,2)
1696: //
1697: // This is kinematics independent, can be pre-calculated.
1698: //
1699: // Remember to apply g_PPT (coupling constant) outside this
1700: //
1701: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_00(const
1702:     const double S0 = 1.0; // Mass scale (GeV)
1703:     const std::complex<double> FACTOR = z1 * S0 * S0;
1704:
1705:     // Init with zeros!
1706:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1707:
1708:     // Contract total vertex by summing up the tensors
1709:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1710:
1711:     FOR_EACH_6(L1);
1712:
1713:     // Indexing
1714:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1715:
1716:     for (auto const i1 : I1) {
1717:         for (auto const i2 : I2) {
1718:             for (auto const i3 : I3) {
1719:                 for (auto const i4 : I4) {
1720:                     for (auto const i5 : I5) {
1721:                         for (auto const i6 : I6) {
1722:                             T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1723:                         }
1724:                     }
1725:                 }
1726:             }
1727:         }
1728:     }
1729:
1730:     // Apply form factors
1731:     T(ind) = F_rilde;
1732:     FOR_EACH_6_END;
1733:
1734:     return T;
1735: }
1736:
1737: // Pomeron-Pomeron-Tensor coupling structures #1 and #2
1738: // iG_1(u mu nu kappa lambda rho sigma)
1739: // ~ (1,s) = (2,0) + (2,2)
1740: // ~ (1,s) = (0,0) + (2,2)
1741: //
1742: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_12(const M4Vec q1, const M4Vec q2,
1743:     const double S0 = 1.0; // Mass scale (GeV)
1744:     const std::complex<double> FACTOR = z1 * S0 * S0;
1745:
1746:     // Init with zeros!
1747:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1748:
1749:     // Contract total vertex by summing up the tensors
1750:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1751:
1752:     FOR_EACH_6(L1);
1753:
1754:     // Indexing
1755:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1756:
1757:     for (auto const i1 : I1) {
1758:         for (auto const i2 : I2) {
1759:             for (auto const i3 : I3) {
1760:                 for (auto const i4 : I4) {
1761:                     for (auto const i5 : I5) {
1762:                         for (auto const i6 : I6) {
1763:                             T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1764:                         }
1765:                     }
1766:                 }
1767:             }
1768:         }
1769:     }
1770:
1771:     // Apply form factors
1772:     T(ind) = F_rilde;
1773:     FOR_EACH_6_END;
1774:
1775:     return T;
1776: }
1777:
1778: // Final rank-6 tensor
1779: MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1780:
1781: FOR_EACH_6(L1);
1782:
1783: const std::vector<size_t> ind = {u, v, k, l, r, s};
1784:
1785: for (auto const i1 : I1) {
1786:     for (auto const i2 : I2) {
1787:         for (auto const i3 : I3) {
1788:             for (auto const i4 : I4) {
1789:                 for (auto const i5 : I5) {
1790:                     for (auto const i6 : I6) {
1791:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1792:                     }
1793:                 }
1794:             }
1795:         }
1796:     }
1797: }
1798:
1799: // Apply form factors
1800: T(ind) = F_rilde;
1801: FOR_EACH_6_END;
1802:
1803: return T;
1804: }
1805:
1806: // NAIVE LOOP
1807: for (auto const i1 : I1) {
1808:     for (auto const i2 : I2) {
1809:         for (auto const i3 : I3) {
1810:             for (auto const i4 : I4) {
1811:                 for (auto const i5 : I5) {
1812:                     for (auto const i6 : I6) {
1813:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1814:                     }
1815:                 }
1816:             }
1817:         }
1818:     }
1819: }
1820:
1821: // Apply form factors
1822: T(ind) = F_rilde;
1823: FOR_EACH_6_END;
1824:
1825: return T;
1826: }

```

```

./src/MTensorPomeron.cc 22/32
1744: const double q1q2 = q1 * q2;
1745: const std::complex<double> FACTOR = -2.0 * z1 / S0 * g_PPT;
1746:
1747: // Coupling structure 1 or 2
1748: const double sign = (mode == '1' ? -1.0 : 1.0);
1749:
1750: // Init with zeros!
1751: MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1752: FOR_EACH_6(L1);
1753: const std::vector<size_t> ind = {u, v, k, l, r, s};
1754:
1755: // Pre-calculated tensor contraction structures T1,T2,T3 used here
1756: T(ind) += q1q2 * T1(ind);
1757:
1758: for (auto const i1 : I1) {
1759:     for (auto const i2 : I2) {
1760:         // Note ==
1761:         T(ind) += sign * q[i1] * (q1[i2] * T2(u, v, k, l, r, s, u, i1));
1762:     }
1763: }
1764:
1765: for (auto const i1 : I1) {
1766:     for (auto const i2 : I2) {
1767:         // Note ==
1768:         T(ind) += sign * q1[i1] * (q2[i2] * T3(u, v, k, l, r, s, u, i1));
1769:     }
1770: }
1771:
1772: for (auto const i1 : I1) {
1773:     for (auto const i2 : I2) {
1774:         T(ind) += (q1[i2] * (q2[i1] * R_RDDDD(u, v, k, l) * R_RDDDU(i, r, s, i1));
1775:     }
1776: }
1777:
1778: T(ind) = FACTOR;
1779: FOR_EACH_6_END;
1780:
1781: return T;
1782: }
1783:
1784: // Pomeron-Pomeron-Tensor coupling structure #3
1785: // iG_1(u mu nu kappa lambda rho sigma) ~ (1,s) = (2,4)
1786: //
1787: MTensor<std::complex<double>, MTensorPomeron::iG_PP_T_03(const M4Vec q1, const M4Vec q2,
1788:     const double S0 = 1.0; // Mass scale (GeV)
1789:     const std::complex<double> FACTOR = z1 / S0 * g_PPT;
1790:
1791:     Tensor<double, 4, 4, 4, 4> A;
1792:     Tensor<double, 4, 4, 4, 4> B;
1793:     Tensor<double, 4, 4, 4, 4> C;
1794:     Tensor<double, 4, 4, 4, 4> D;
1795:
1796:     // Manual contractions
1797:     FOR_EACH_3(L1);
1798:     A(u, v, k) = 0.0;
1799:     B(u, v, k) = 0.0;
1800:     for (auto const i1 : I1) {
1801:         // Note ==
1802:         A(u, v, k) += q2[i1] * R_RDDDD(u, v, k, k);
1803:         B(u, v, k) += q1[i1] * R_RDDDD(u, v, k, k);
1804:     }
1805:     FOR_EACH_3_END;
1806:
1807:     // Final rank-6 tensor
1808:     MTensor<std::complex<double>, T = MTensor((4, 4, 4, 4, 4, 4), std::complex<double>(0.0));
1809:
1810:     FOR_EACH_6(L1);
1811:
1812:     const std::vector<size_t> ind = {u, v, k, l, r, s};
1813:
1814:     for (auto const i1 : I1) {
1815:         // Note ==
1816:         T(ind) += (A(u, v, v1) * B(k, l, l1) + A(u, v, l1) * B(k, l, v1)) * R_RUDD(v1, l1, r, s);
1817:     }
1818:
1819:     T(ind) = FACTOR;
1820:
1821:     return T;
1822: }
1823:
1824: // NAIVE LOOP
1825: for (auto const i1 : I1) {
1826:     for (auto const i2 : I2) {
1827:         for (auto const i3 : I3) {
1828:             for (auto const i4 : I4) {
1829:                 for (auto const i5 : I5) {
1830:                     for (auto const i6 : I6) {
1831:                         T(ind) += IGM[i1][i2][i3][i4][i5][i6];
1832:                     }
1833:                 }
1834:             }
1835:         }
1836:     }
1837: }
1838:
1839: // Apply form factors
1840: T(ind) = F_rilde;
1841: FOR_EACH_6_END;
1842:
1843: return T;
1844: }

```

```

./src/MTensorPomeron.cc 23/32
1827: T(iind) += FACTOR * (q2[u1] * R_DDDD(u,v,vi) * q1[a1] * R_DDDD(k,l,al,1)) +
1828: R_UUDD(v1,l1,r,s);
1829:
1830: }
1831: }
1832: }
1833: }
1834: }
1835: FOR_EACH_6_END;
1836:
1837: return T;
1838: }
1839:
1840: // Pomeron-Pomeron-Tensor coupling structure #4
1841: // iG_lmu lnu kappas lambda rho sigmas ^ (1,s) = (4,2)
1842: //
1843: MTensor<std::complex>> MTensorPomeron:iG_PPT_04(const M4Vec iqt, const M4Vec iqt2,
1844: const double S0 = 1.0; // Mass scale (GeV)
1845: const std::complex>> FACTOR = -2.0 * zi / math::pow(S0) * g_PPT;
1846: const double q1q2 = q1 * q2;
1847:
1848: Tensor<double> q1_D = (q1 & 0, q1 & 1, q1 & 2, q1 & 3);
1849: Tensor<double> q2_D = (q2 & 0, q2 & 1, q2 & 2, q2 & 3);
1850:
1851:
1852: Tensor<double> 4, 4, 4, 4, 4; A;
1853: Tensor<double> 4, 4, 4, 4, 4; B;
1854: Tensor<double> 4, 4, 4; C;
1855:
1856: // Manual contractions (not possible with autocontraction)
1857: FOR_EACH_4(I1);
1858: A(u, v, k, l) = 0.0;
1859: B(u, v, k, l) = 0.0;
1860:
1861: for (auto const kul : LI) {
1862: for (auto const svi : LI) {
1863: for (auto const kul : LI) {
1864: // Note ==
1865: A(u, v, k, l) += q2[u1] * R_DDDD(u, v, vi, al) * q1[u1] * R_DDDD(k, l, ul, al);
1866: B(u, v, k, l) += q2[u1] * R_DDDD(u, v, vi, al) * q1[v1] * R_DDDD(k, l, vl, al);
1867: }
1868: }
1869: }
1870: FOR_EACH_4_END;
1871:
1872: // Autocontractions
1873: {
1874: FTensor::Index<'a', 4> r;
1875: FTensor::Index<'b', 4> s;
1876: FTensor::Index<'c', 4> al;
1877: FTensor::Index<'d', 4> ll;
1878:
1879: C(r, s) = q1_D(al) * q2_D(ll) * R_UUDD(al, ll, r, s);
1880: }
1881:
1882: // Final rank-6 expression
1883: MTensor<std::complex>> T = MTensor<(4, 4, 4, 4, 4, 4), std::complex>>(0.0);
1884: FOR_EACH_6(I1);
1885: T(lu, v, k, l, r, s) =
1886: FACTOR * (A(u, v, k, l) + B(u, v, k, l) - 2.0 * q1q2 * R_DDDD(u, v, k, l)) * C(r, s);
1887:
1888: }
1889: // NAIVE LOOP
1890: for (auto const svi : LI) {
1891: for (auto const kul : LI) {
1892: for (auto const kul : LI) {
1893: for (auto const kul : LI) {
1894: // Note ==
1895: T(iind) += FACTOR * (q2[u1] * R_DDDD(u,v,vi,al) * q1[u1] * R_DDDD(k,l,ul,al)
1896: + q1[u1] * R_DDDD(u,v,vi,al) * q1[v1] * R_DDDD(k,l,vl,al)
1897: - 2.0 * q1q2 * R_DDDD(u,v,k,l)) * (q1 & al) * (q2 & ll) *
1898: R_UUDD(al, ll, r, s);
1899: }
1900: }
1901: }
1902: }
1903: }
1904: }
1905: FOR_EACH_6_END;
1906: return T;
1907: }
1908:
1909: // Pomeron-Pomeron-Tensor coupling structure #05

```

```

./src/MTensorPomeron.cc 25/32
1993: FTensor::Index<'b', 4> b;
1994: FTensor::Index<'c', 4> c;
1995: FTensor::Index<'d', 4> d;
1996:
1997: Tensor<double> q1_U = (q1[0], q1[1], q1[2], q1[3]);
1998: Tensor<double> q2_U = (q2[0], q2[1], q2[2], q2[3]);
1999:
2000: Tensor<double> 4, 4, 4; A;
2001: Tensor<double> 4, 4, 4; B;
2002: Tensor<double> 4, 4, 4; C;
2003:
2004: // Autocontractions
2005: A(a, b) = q2_U(c) * q2_U(d) * R_DDDD(a, b, c, d);
2006: B(a, b) = q1_U(c) * q1_U(d) * R_DDDD(a, b, c, d);
2007: C(a, b) = q1_U(c) * q2_U(d) * R_DDDD(a, b, c, d);
2008:
2009: MTensor<std::complex>> T = MTensor<(4, 4, 4, 4, 4, 4), std::complex>>(0.0);
2010: FOR_EACH_6(I1);
2011: T(lu, v, k, l, r, s) = FACTOR * A(u, v) * B(k, l) * C(r, s);
2012: FOR_EACH_6_END;
2013: return T;
2014: }
2015: }
2016: }
2017: }
2018: // =====
2019:
2020: // f2 (Scalar) - (Pseudo)Scalar - (Pseudo)Scalar vertex function
2021: // i/Gamma
2022: //
2023: // Input as contravariant (upper index) 4-vectors, meson on-shell mass M0
2024: // Relations: p3_lmu iG^lmu lnu = 0, p4_lnu iG^lmu lnu = 0
2025: //
2026: std::complex>> MTensorPomeron:iG_f0s(const M4Vec iqt, const M4Vec iqt2, double M0,
2027: const double S0 = 1.0; // Mass scale (GeV)
2028: // Form FACTOR
2029: const double lambda4 = 1.0; // GeV^4 (normalization scale)
2030: auto F_f0s = [k](double q2) { return std::exp(-pow(q2 - M0 * M0) / lambda4); };
2031:
2032: return zi * q1 * S0 * F_f0s((p3 + p4).M2());
2033: }
2034: }
2035: }
2036: }
2037: }
2038:
2039: // f2 (Scalar) - Vector (Massive) - Vector (Massive) vertex function
2040: // i/Gamma_lmu lnu
2041: //
2042: // Input as contravariant (upper index) 4-vectors, meson on-shell mass M0
2043: // Relations: p3_lmu iG^lmu lnu = 0, p4_lnu iG^lmu lnu = 0
2044: //
2045: Tensor<std::complex>> 4, 4, 4; MTensorPomeron:iG_f0vv(const M4Vec iqt, const M4Vec iqt2,
2046: const double S0 = 1.0; // Mass scale (GeV)
2047: const double LAMBDA = 1.0; // GeV
2048: const double p3p4 = p3 * p4;
2049: const double p3sq = p3.M2();
2050: const double p4sq = p4.M2();
2051:
2052: const double F_v = FM(p3sq, 0.5) * FM(p4sq, 0.5) * F((p3 + p4).M2(), M0, LAMBDA);
2053: const std::complex>> FACTOR = zi * q1 * 2.0 / math::pow(S0) * F_v;
2054: const std::complex>> FACTOR2 = zi * q2 * 2.0 / S0 * F_v;
2055:
2056: Tensor<std::complex>> 4, 4, 4; T;
2057: FOR_EACH_2(I1);
2058: T(u, v) = FACTOR1 * (p3sq * p4sq * q1[u][v] - p4sq * (p3 & u) * (p3 & v) -
2059: p3sq * (p4 & u) * (p4 & v) + p3p4 * (p3 & u) * (p4 & v)) +
2060: FACTOR2 * ((p4 & u) * (p3 & v) - p3p4 * q1[u][v]);
2061:
2062: return T;
2063: }
2064: }
2065: }
2066: }
2067: }
2068:
2069: // Vector (Massive) - Pseudoscalar - Pseudoscalar vertex function
2070: // i/Gamma_lmu(k1,k2)
2071: //
2072: // Input as contravariant (upper index) 4-vectors
2073: //
2074: Tensor<std::complex>> 4, 4; MTensorPomeron:iG_Vps(const M4Vec iqt, const M4Vec iqt2,
2075: const double M0, double g1) const {

```

```

./src/MTensorPomeron.cc 24/32
1910: // iG_lmu lnu kappas lambda rho sigmas ^ (1,s) = (4,4)
1911: //
1912: MTensor<std::complex>> MTensorPomeron:iG_PPT_05(const M4Vec iqt, const M4Vec iqt2,
1913: const double S0 = 1.0; // Mass scale (GeV)
1914: const std::complex>> FACTOR = zi / math::pow(S0) * g_PPT;
1915:
1916: Tensor<double> 4, 4; q1_U = (q1[0], q1[1], q1[2], q1[3]);
1917: Tensor<double> 4, 4; q2_U = (q2[0], q2[1], q2[2], q2[3]);
1918: Tensor<double> 4, 4; q1_D = (q1 & 0, q1 & 1, q1 & 2, q1 & 3);
1919: Tensor<double> 4, 4; q2_D = (q2 & 0, q2 & 1, q2 & 2, q2 & 3);
1920: Tensor<double> 4, 4, 4, 4, 4, 4; A;
1921: Tensor<double> 4, 4, 4, 4, 4, 4; B;
1922: Tensor<double> 4, 4, 4, 4, 4, 4; C;
1923: Tensor<double> 4, 4, 4, 4, 4, 4; D;
1924: Tensor<double> 4, 4, 4, 4, 4, 4; E;
1925:
1926: // Manual contractions (not possible with autocontraction)
1927: MTensor<std::complex>> T = MTensor<(4, 4, 4, 4, 4, 4), std::complex>>(0.0);
1928: FOR_EACH_6(I1);
1929: A(u, v, r, s) = 0.0;
1930: C(k, l, r, s) = 0.0;
1931:
1932: for (auto const kul : LI) {
1933: for (auto const svi : LI) {
1934: for (auto const kul : LI) {
1935: // Note ==
1936: A(u, v, r, s) += q2_U(u1) * R_DDDD(u, v, vi, v1) * q2_D(r1) * R_UUDD(v1, r1, r, s);
1937: C(k, l, r, s) += q1_U(u1) * R_DDDD(k, l, ul, vl) * q1_D(r1) * R_UUDD(v1, r1, r, s);
1938: }
1939: }
1940: }
1941: }
1942: FOR_EACH_6_END;
1943:
1944: // Autocontractions
1945: {
1946: FTensor::Index<'a', 4> u;
1947: FTensor::Index<'b', 4> v;
1948: FTensor::Index<'c', 4> r;
1949: FTensor::Index<'d', 4> s;
1950: FTensor::Index<'e', 4> al;
1951: FTensor::Index<'f', 4> ll;
1952:
1953: B(k, l) = q1_D(al) * q1_U(ll) * R_DDDD(k, l, al, ll);
1954: D(u, v) = q2_D(al) * q2_U(ll) * R_DDDD(u, v, al, ll);
1955: }
1956:
1957: // Final rank-6 expression
1958: FOR_EACH_6(I1);
1959: T(lu, v, k, l, r, s) = FACTOR * (A(u, v, r, s) * B(k, l) + C(k, l, r, s) * D(u, v));
1960:
1961: }
1962: // NAIVE LOOP
1963: for (auto const svi : LI) {
1964: for (auto const kul : LI) {
1965: for (auto const kul : LI) {
1966: for (auto const kul : LI) {
1967: // Note ==
1968: T(iind) += FACTOR * (q2[u1] * (q2 & r1) * R_DDDD(u,v,vi,vi) * R_UUDD(v1,r1,r,s) * q1[al]
1969: * q1[l1] * R_DDDD(k,l,al,al)
1970: + q1[u1] * R_DDDD(u,v,vi,vi) * q1[v1] * (q1 & r1) * R_DDDD(k,l,vi,vi) * R_UUDD(v1,r1,r,s) * q2[al]
1971: * q2[l1] * R_DDDD(u,v,al,al));
1972: }
1973: }
1974: }
1975: }
1976: }
1977: }
1978: }
1979: }
1980: }
1981: return T;
1982: }
1983: }
1984: // Pomeron-Pomeron-Tensor coupling structure #6
1985: // iG_lmu lnu kappas lambda rho sigmas ^ (1,s) = (6,4)
1986: //
1987: MTensor<std::complex>> MTensorPomeron:iG_PPT_06(const M4Vec iqt, const M4Vec iqt2,
1988: const double S0 = 1.0; // Mass scale (GeV)
1989: const std::complex>> FACTOR = -2.0 * zi / math::pow(S0) * g_PPT;
1990:
1991: FTensor::Index<'a', 4> a;
1992:
1993: const M4Vec p = k1 - k2;
1994:
1995: // Form FACTOR
1996: const double LAMBDA = 1.0; // GeV
1997: const double F_LM = FM(k1.M2(), 0.5) * FM(k2.M2(), 0.5) * F((k1 + k2).M2(), M0, LAMBDA);
1998: const std::complex>> FACTOR2 = -0.5 * zi * q1 * g;
1999:
2000: Tensor<std::complex>> 4, 4; T;
2001: for (const auto kmu : LI) { T[mu] = FACTOR * (p & kmu); }
2002: return T;
2003: }
2004:
2005: // pseudoscalar - vector - vector vertex function
2006: // i/Gamma_lmu(kappa1,k1,k2)
2007: //
2008: // Input as contravariant (upper index) 4-vectors, M0 is the pseudoscalar on-shell mass
2009: Tensor<std::complex>> 4, 4; MTensorPomeron:iG_Vpv(const M4Vec iqt, const M4Vec iqt2,
2010: const double M0, double g1) const {
2011: // Form FACTOR
2012: const double LAMBDA = 1.0; // GeV
2013: const double S0 = 1.0; // Mass scale (GeV)
2014: const std::complex>> FACTOR = -zi * g1 / (2 * S0) * F((p3 + p4).M2(), M0, LAMBDA);
2015:
2016: Tensor<std::complex>> 4, 4; p3 = (p3[0], p3[1], p3[2], p3[3]);
2017: Tensor<std::complex>> 4, 4; p4 = (p4[0], p4[1], p4[2], p4[3]);
2018:
2019: FTensor::Index<'a', 4> mu;
2020: FTensor::Index<'b', 4> nu;
2021: FTensor::Index<'c', 4> rho;
2022: FTensor::Index<'d', 4> sigma;
2023:
2024: // Contract
2025: Tensor<std::complex>> 4, 4, 4; T;
2026: T(mu, nu) = FACTOR * p3[rho] * p4[sigma] * eps_l(mu, nu, rho, sigma);
2027:
2028: return T;
2029: }
2030: }
2031: }
2032: // f2 - pseudoscalar - pseudoscalar vertex function
2033: // i/Gamma_lmu(kappa1,k1,k2)
2034: //
2035: // Input as contravariant (upper index) 4-vectors, M0 is the f2 meson on-shell mass
2036: //
2037: Tensor<std::complex>> 4, 4; MTensorPomeron:iG_f2ps(const M4Vec iqt, const M4Vec iqt2,
2038: const double M0, double g1) const {
2039: // Form FACTOR
2040: const double LAMBDA = 1.0; // GeV
2041: const double S0 = 1.0; // Mass scale (GeV)
2042: const M4Vec k1_k2 = k1 - k2;
2043: const std::complex>> FACTOR = -zi * g1 / (2 * S0) * F((k1 + k2).M2(), M0, LAMBDA);
2044:
2045: const double k1_k2_sq = k1_k2.M2();
2046:
2047: Tensor<std::complex>> 4, 4; T;
2048: FOR_EACH_2(I1);
2049: T(u, v) = FACTOR * ((k1_k2 & u) * (k1_k2 & v) - 0.25 * q1[u][v] * k1_k2_sq);
2050:
2051: return T;
2052: }
2053: }
2054: }
2055: }
2056: // f2 - Vector (massive) - Vector (Massive) vertex function
2057: // i/Gamma_lmu lnu kappas lambda(k1,k2)
2058: //
2059: // Input as contravariant (upper index) 4-vectors, M0 is the f2 meson on-shell mass
2060: //
2061: Tensor<std::complex>> 4, 4, 4, 4; MTensorPomeron:iG_f2vv(const M4Vec iqt, const M4Vec iqt2,
2062: const double M0, double g1,
2063: const double M02, double g2) const {
2064: // Form FACTOR
2065: const double S0 = 1.0; // Mass scale (GeV)
2066: const double LAMBDA = 1.0; // GeV
2067: const Tensor<std::complex>> 4, 4, 4; G0 = Gamma0(k1, k2);
2068: const Tensor<std::complex>> 4, 4, 4, 4; G2 = Gamma2(k1, k2);
2069:
2070: const std::complex>> F_v =
2071: FM(k1.M2(), 0.5) * FM(k2.M2(), 0.5) * F((k1 + k2).M2(), M0, LAMBDA);
2072: const std::complex>> FACTOR1 = zi * 2.0 / math::pow(S0) * g1 * F_v;
2073: const std::complex>> FACTOR2 = -zi * 1.0 / S0 * g2 * F_v;
2074:
2075: Tensor<std::complex>> 4, 4, 4, 4; T;

```

```

./src/MTensorPomeron.cc          27/32
2150: FOR_EACH_4(L1);
2151: T(u, v, k, l) = FACTOR1 * G0(u, v, k, l) + FACTOR2 * G2(u, v, k, l);
2152: FOR_EACH_4_END;
2153: return T;
2154: }
2155:
2156: // f2 - gamma - gamma vertex function
2157: // iGamma_{\mu\nu}(kappa\lambda\lambda\lambda)(k1,k2)
2158:
2159: // Input as contravariant (upper index) 4-vectors, M0 is the f2 meson on-shell mass
2160: // Couplings g1,g2
2161:
2162: // Example couplings:
2163: // const double e = msqrt(qed:\alpha_QED() * 4.0 * PI); // ~ 0.3, no running
2164: // const double a_f2yy = pow(e) / (4 * PI) + 1.45; // gV^*V(-3)
2165: // const double a_f2yy = pow(e) / (4 * PI) + 2.49; // gV^*V(-1)
2166:
2167: Tensor4<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iG_f2yy(const M4Vec k1, const M4Vec k2,
2168: double M0, double g1,
2169: double g2) const {
2170: // Form FACTOR
2171: const double LAMBDA = 1.0; // gV
2172:
2173: const Tensor4<std::complex<double>, 4, 4, 4, 4> G0 = Gamma(k1, k2);
2174: const Tensor4<std::complex<double>, 4, 4, 4, 4> G2 = Gamma(k1, k2);
2175:
2176: const std::complex<double> FACTOR =
2177: z1 * FM(k1.M2(), 0.5) * FM(k2.M2(), 0.5) * F((k1+k2).M2(), M0, LAMBDA);
2178:
2179: Tensor4<std::complex<double>, 4, 4, 4, 4> T;
2180: FOR_EACH_4(L1);
2181: T(u, v, k, l) = FACTOR * (2.0 * g1 * G0(u, v, k, l) - g2 * G2(u, v, k, l));
2182: FOR_EACH_4_END;
2183: return T;
2184: }
2185:
2186: // Pomeron-Pseudoscalar-Pseudoscalar vertex function
2187: // iGamma_{\mu\nu}(p',p)
2188:
2189: // Input as contravariant (upper index) 4-vectors
2190: // Couplings g1,g2
2191:
2192: Tensor2<std::complex<double>, 4, 4> MTensorPomeron::iD_Pppp(const M4Vec pprime, const M4Vec sp,
2193: double M0, double g1,
2194: double g2) const {
2195: const M4Vec psum = pprime + sp;
2196: const double M2 = psum.M2();
2197: const std::complex<double> FACTOR = -zi * 2.0 * g1 * FM(pprime - p.M2(), 0.5);
2198:
2199: Tensor2<std::complex<double>, 4, 4> T;
2200: FOR_EACH_2(L1);
2201: T(u, v) = FACTOR * (psum.k u) * (psum.k v) - 0.25 * g1[u]v * M2;
2202: FOR_EACH_2_END;
2203: return T;
2204: }
2205:
2206: // Pomeron-Vector(massive)-Vector(massive) vertex function
2207: // iGamma_{\alpha\beta}(gamma\delta)(p',p)
2208:
2209: // Input M0 is the vector meson on-shell mass
2210: // Couplings g1,g2
2211:
2212: // Input as contravariant (upper index) 4-vectors
2213:
2214: Tensor4<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iD_Pvv(const M4Vec pprime, const M4Vec sp,
2215: double M0, double g1,
2216: double g2) const {
2217: const Tensor4<std::complex<double>, 4, 4, 4, 4> G0 = Gamma(pprime, -p);
2218: const Tensor4<std::complex<double>, 4, 4, 4, 4> G2 = Gamma(pprime, -p);
2219:
2220: const std::complex<double> FACTOR = zi * FM(pprime - p.M2(), 0.5);
2221:
2222: Tensor4<std::complex<double>, 4, 4, 4, 4> T;
2223: FOR_EACH_4(L1);
2224: T(u, v, k, l) = FACTOR * (2.0 * g1 * G0(u, v, k, l) - g2 * G2(u, v, k, l));
2225: FOR_EACH_4_END;
2226: return T;
2227: }
2228:
2229: // Gamma-Vector meson transition vertex
2230: // iGamma_{\mu\nu}(nu)
2231:
2232: Tensor2<std::complex<double>, 4, 4, 4> T;
2233: FOR_EACH_2(L1);
2234: T(u, v) = FACTOR * (2.0 * g1 * G0(u, v, k, l) - g2 * G2(u, v, k, l));
2235: FOR_EACH_2_END;
2236: return T;
2237: }
2238:
2239:
2240:
2241:

```

```

./src/MTensorPomeron.cc          28/32
2242:
2243: // pdg = 113 / "rho0", 223 / "omega0", 333 / "phi0"
2244:
2245: Tensor2<std::complex<double>, 4, 4> MTensorPomeron::iG_VV(double g2, int pdg) const {
2246: const double e = msqrt(qed:\alpha_QED() * 4.0 * PI); // ~ 0.3, no running
2247: double gammav = 0.0;
2248: double mv = 0.0;
2249:
2250: // Vector Meson Dominance (VMD) type parameters
2251: // If pdg = 113 ( / "rho
2252: gammav = msqrt(4 * PI / 0.496);
2253: mv = 0.770;
2254: } else if (pdg == 223 ( / "rho
2255: gammav = msqrt(4 * PI / 0.042);
2256: mv = 0.785;
2257: } else if (pdg == 333 ( / "phi
2258: gammav = (-1.0) * msqrt(4 * PI / 0.0716);
2259: mv = 1.020;
2260: } else
2261: throw std::invalid_argument("MTensorPomeron::iG_VV: Unknown vector meson pdg = " +
2262: std::ito_string(pdg));
2263: }
2264:
2265: Tensor2<std::complex<double>, 4, 4> T;
2266: FOR_EACH_2(L1);
2267: T(u, v) = -zi * e * pow2(mv) / gammav * g[u]v;
2268: FOR_EACH_2_END;
2269: return T;
2270: }
2271:
2272:
2273:
2274: // Propagators
2275:
2276: // Tensor Pomeron propagator (covariant = contravariant)
2277:
2278: // iDelta_{\mu\nu}(kappa\lambda\lambda\lambda)(s,t) = iDelta_{\nu\mu}(kappa\lambda\lambda\lambda)(s,t)
2279:
2280: // Input as (sub)-Mandelstam invariants: s,t
2281:
2282: // Symmetry relations:
2283: // iDelta_{\mu\nu}(kappa\lambda\lambda\lambda) = iDelta_{\nu\mu}(kappa\lambda\lambda\lambda)
2284: // iDelta_{\mu\nu}(kappa\lambda\lambda\lambda) = iDelta_{\nu\mu}(kappa\lambda\lambda\lambda)
2285:
2286: // Contraction with g^{\mu\nu} or g^{\rho\sigma}(kappa\lambda\lambda\lambda) gives 0
2287:
2288: Tensor2<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iD_P(double s, double t) const {
2289: const std::complex<double> FACTOR =
2290: 1.0 / (4.0 * s) * std::pow(-zi * s * Param.ap.P, alpha_P(t) - 1.0);
2291:
2292: Tensor4<std::complex<double>, 4, 4, 4, 4> T;
2293: FOR_EACH_4(L1);
2294: T(u, v, k, l) = FACTOR * g[u]k * g[v]l + g[u]l * g[v]k - 0.5 * g[u]v * g[k]l;
2295: FOR_EACH_4_END;
2296: return T;
2297: }
2298:
2299:
2300: // Tensor Reggeon propagator ID (f2-a2-reggeon)
2301:
2302: // Input as (sub)-Mandelstam invariants: s,t
2303:
2304: Tensor2<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::iD_2R(double s, double t) const {
2305: const std::complex<double> FACTOR =
2306: 1.0 / (4.0 * s) * std::pow(-zi * s * Param.ap.2R, alpha_2R(t) - 1.0);
2307:
2308: Tensor4<std::complex<double>, 4, 4, 4, 4> T;
2309: FOR_EACH_4(L1);
2310: T(u, v, k, l) = FACTOR * g[u]k * g[v]l + g[u]l * g[v]k - 0.5 * g[u]v * g[k]l;
2311: FOR_EACH_4_END;
2312: return T;
2313: }
2314:
2315:
2316: // Vector Reggeon propagator ID (rho-omega-reggeon)
2317:
2318: // Input as (sub)-Mandelstam invariants: s,t
2319:
2320: Tensor2<std::complex<double>, 4, 4> MTensorPomeron::iD_1R(double s, double t) const {
2321: const std::complex<double> FACTOR =
2322: zi / pow2(Param.M1R) * std::pow(-zi * s * Param.ap.1R, alpha_1R(t) - 1.0);
2323:
2324: Tensor2<std::complex<double>, 4, 4, 4> T;
2325: }
2326:
2327:
2328:
2329:
2330:
2331:
2332:
2333:
2334: // Input as (sub)-Mandelstam invariants s,t
2335:
2336: Tensor2<std::complex<double>, 4, 4> MTensorPomeron::iD_0(double s, double t) const {
2337: const std::complex<double> FACTOR =
2338: -zi * Param.sta.0 / pow2(Param.M0) * std::pow(-zi * s * Param.ap.0, alpha_0(t) - 1.0);
2339:
2340: Tensor2<std::complex<double>, 4, 4> T;
2341: FOR_EACH_2(L1);
2342: T(u, v) = FACTOR * g[u]v;
2343: FOR_EACH_2_END;
2344: return T;
2345: }
2346:
2347:
2348:
2349: // Vector propagator ID_{\mu\nu}(nu) = ID^{\nu\mu}(nu) with (naive) Reggeization
2350:
2351: // Input as contravariant 4-vector and the particle mass M0
2352:
2353: Tensor2<std::complex<double>, 4, 4> MTensorPomeron::iD_V(const M4Vec sp, double M0,
2354: double s34) const {
2355: const double m2 = p.M2();
2356: const double alpha = 0.5 + 0.9 * m2; // Typical Reggeon trajectory
2357: const double s_thresh = 4 * pow2(M0);
2358: const double phi = PI / 2.0 * std::exp((ln(s_thresh - s34) / s_thresh) - PI / 2.0);
2359:
2360: // Reggeization
2361: std::complex<double> REGGIEZE = std::pow(std::exp(zi * phi) * s34 / s_thresh, alpha - 1.0);
2362:
2363: Tensor2<std::complex<double>, 4, 4> T;
2364: FOR_EACH_2(L1);
2365: T(u, v) = g[u]v / (m2 - pow2(M0)) * REGGIEZE;
2366: FOR_EACH_2_END;
2367: return T;
2368: }
2369:
2370:
2371: // Scalar propagator ID with zero width
2372:
2373: // Input as 4-vector and the particle mass M0
2374:
2375: std::complex<double> MTensorPomeron::iD_M0(const M4Vec sp, double M0) const {
2376: return zi / (p.M2() - pow2(M0));
2377: }
2378:
2379: // Scalar propagator ID with finite Breit-Wigner width
2380:
2381: // Input as 4-vector and the particle mass M0 and full width
2382:
2383: std::complex<double> MTensorPomeron::iD_MS(const M4Vec sp, double M0, double Gamma) const {
2384: const double Delta_L = 1.0 / (p.M2() - pow2(M0) + zi * M0 * Gamma);
2385:
2386: return zi * Delta;
2387: }
2388:
2389: // Massive vector propagator ID_{\mu\nu}(nu) or ID^{\nu\mu}(nu) with INDEX_UP == true
2390:
2391: // Input as contravariant 4-vector and the particle mass M0 and full width Gamma
2392:
2393: Tensor2<std::complex<double>, 4, 4> MTensorPomeron::iD_VMS(const M4Vec sp, double M0, double Gamma,
2394: bool INDEX_UP) const {
2395: // Transverse part
2396: const double m2 = p.M2();
2397: const std::complex<double> Delta_L = 1.0 / (m2 - pow2(M0) + zi * M0 * Gamma);
2398:
2399: // Longitudinal part (does not enter here)
2400: const double Delta_L = 0.0;
2401:
2402: Tensor2<std::complex<double>, 4, 4> T;
2403: IF (INDEX_UP) {
2404: FOR_EACH_2(L1);
2405: T(u, v) =
2406: zi * ( -g[u]v + (p[u] * (p[v] / m2) * Delta_T - zi * (p[u] * (p[v] / m2 * Delta_L,
2407: FOR_EACH_2_END;

```

```

./src/MTensorPomeron.cc          31/32
2491: |
2492: |
2493: |-----
2494: | Tensor functions
2495: |
2496: | Output:  $\langle \Gamma_{\text{gamma}} \rangle (0)_{\dots} / (\mu \text{nu} \kappa \lambda \text{ambda})$ 
2497: |
2498: | Input must be contravariant (upper index) 4-vectors
2499: |
2500: | Tensor4<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::Gamma0(const MVec k1,
2501: | const MVec k2) const {
2502: |     const double k1k2 = k1 * k2; // 4-dot product
2503: |     Tensor4<std::complex<double>, 4, 4, 4, 4> T;
2504: |     FOR_EACH_4(L1) {
2505: |         T(u, v, w, l) = (k1k2 * g[u][v] - (k2 * l) * (k1 * w) + (k1 * l) * (k2 * w) -
2506: |             ((k1 * k) * (k2 * l) + (k2 * k) * (k1 * l)) - 0.5 * k1k2 * g[k][l]);
2507: |     }
2508: |     return T;
2509: | }
2510: |
2511: |
2512: | Output:  $\langle \Gamma_{\text{gamma}} \rangle (2)_{\dots} / (\mu \text{nu} \kappa \lambda \text{ambda})$ 
2513: |
2514: | Input must be contravariant (upper index) 4-vectors
2515: |
2516: | Tensor4<std::complex<double>, 4, 4, 4, 4> MTensorPomeron::Gamma2(const MVec k1,
2517: | const MVec k2) const {
2518: |     const double k1k2 = k1 * k2; // 4-dot product
2519: |     Tensor4<std::complex<double>, 4, 4, 4, 4> T;
2520: |     FOR_EACH_4(L1) {
2521: |         T(u, v, w, l) = k1k2 * (g[u][k] * g[v][l] + g[u][l] * g[v][k]) +
2522: |             g[u][v] * ((k1 * k) * (k2 * l) + (k2 * k) * (k1 * l)) -
2523: |             (k1 * v) * (k2 * l) + g[u][k] - (k1 * v) * (k2 * k) * g[u][l] -
2524: |             (k2 * u) * (k1 * l) + g[v][k] - (k2 * u) * (k1 * k) * g[v][l] -
2525: |             (k1k2 * g[u][v] - (k2 * u) * (k1 * w) * g[k][l]);
2526: |     }
2527: |     return T;
2528: | }
2529: |
2530: |
2531: | Pre-calculate tensors for speed
2532: |
2533: | 1. Minkowski metric tensor
2534: | 2. Auxiliary (hep) tensor
2535: | 1/2 g_{\mu\nu} / (\mu \nu \kappa \lambda \text{ambda}) + 1/2 g_{\mu\lambda} g_{\nu\kappa} / (\nu \kappa \lambda \text{ambda}) - 1/4
2536: | g_{\mu\nu} g_{\kappa\lambda} / (\kappa \lambda \text{ambda})
2537: |
2538: | void MTensorPomeron::CalcTensor() {
2539: |     // Minkowski metric tensor (+1,-1,-1,-1)
2540: |     FOR_EACH_2(L1) {
2541: |         gT(u, v) = g[u][v];
2542: |     }
2543: |     FOR_EACH_2_RND;
2544: | }
2545: |
2546: |-----
2547: | Covariant (lower index) 4D-epsilon tensor
2548: |
2549: | const MTensor<int> epsilon = math::EpsTensor(4);
2550: |
2551: | FOR_EACH_4(L1) {
2552: |     eps_lo(u, v, w, l) = static_cast<double>(epsilon(ind));
2553: | }
2554: |
2555: | Free Lorentz indices (second parameter denotes the range of index)
2556: | FTensor::Index<'a', 4> a;
2557: | FTensor::Index<'b', 4> b;
2558: | FTensor::Index<'c', 4> c;
2559: | FTensor::Index<'d', 4> d;
2560: | FTensor::Index<'g', 4> g;
2561: | FTensor::Index<'h', 4> h;
2562: |
2563: | Contravariant version (make it in two steps)
2564: | eps_hi(a, b, c, d) = eps_lo(alfa, beta, c, d) * gT(a, alfa) * gT(b, beta);
2565: | eps_hi(a, b, c, d) = eps_hi(a, b, alfa, beta) * gT(c, alfa) * gT(d, beta);
2566: |
2567: |
2568: | Aux tensor R
2569: | FTensor::Tensor<double, 4, 4, 4, 4> R;
2570: |
2571: | R(a, b, c, d) =
2572: |     0.3 * gT(a, c) * gT(b, d) + 0.5 * gT(a, d) * gT(b, c) + 0.25 * gT(a, b) * gT(c, d);
2573: |

```

```

./src/MTensorPomeron.cc          32/32
2574: | Different mixed index covariant/contravariant versions
2575: | By contraction with g_{\mu\nu}
2576: | R_DDDDD = R;
2577: | R_DDDDU(a, b, c, d) = R(a, b, c, alfa) * gT(d, alfa);
2578: | R_DDDDU(a, b, c, d) = R(a, b, alfa, beta) * gT(c, alfa) * gT(d, beta);
2579: | R_DDDDU(a, b, c, d) = R(alfa, beta, c, d) * gT(a, alfa) * gT(b, beta);
2580: |
2581: |-----
2582: | Pre-calculated tensor contractions for tensor resonances
2583: |
2584: | auto FIX1 = {4}();
2585: | MTensor<double> A = MTensor<{4, 4, 4, 4, 4, 4}, double(0.0); // Init with zeros!
2586: | FOR_EACH_6(L1) {
2587: |     const std::vector<size_t> ind = {u, v, k, l, r, s};
2588: |     for (auto const k1 : L1) {
2589: |         for (auto const k2 : L1) {
2590: |             A(ind) = R_DDDDU(u, v, r, s, al) * R_DDDDU(k, l, s, al) * R_DDDDU(r, s, al, s);
2591: |         }
2592: |     }
2593: | }
2594: |
2595: | FOR_EACH_6_RND;
2596: | return A;
2597: | }
2598: |
2599: | auto FIX2 = {4}();
2600: | MTensor<double> A = MTensor<{4, 4, 4, 4, 4, 4, 4, 4}, double(0.0); // Init with zeros!
2601: | FOR_EACH_8(L1) {
2602: |     for (auto const k1 : L1) {
2603: |         for (auto const k2 : L1) {
2604: |             for (auto const k3 : L1) {
2605: |                 A({u, v, k, l, r, s, ut, r1}) +=
2606: |                     R_DDDDU(u, v, ut, al) * R_DDDDU(k, l, s, al) * R_DDDDU(r, s, al, s);
2607: |             }
2608: |         }
2609: |     }
2610: | }
2611: |
2612: | FOR_EACH_8_RND;
2613: | return A;
2614: | }
2615: |
2616: | auto FIX3 = {4}();
2617: | MTensor<double> A = MTensor<{4, 4, 4, 4, 4, 4, 4, 4, 4}, double(0.0); // Init with zeros!
2618: | FOR_EACH_10(L1) {
2619: |     for (auto const k1 : L1) {
2620: |         for (auto const k2 : L1) {
2621: |             for (auto const k3 : L1) {
2622: |                 for (auto const k4 : L1) {
2623: |                     A({u, v, w, l, r, s, ut, ut, al}) +=
2624: |                         R_DDDDU(u, v, r, al) * R_DDDDU(k, l, s, al) * R_DDDDU(r, s, al, s);
2625: |                 }
2626: |             }
2627: |         }
2628: |     }
2629: | }
2630: | FOR_EACH_10_RND;
2631: | return A;
2632: | }
2633: | T1 = FIX1();
2634: | T2 = FIX2();
2635: | T3 = FIX3();
2636: |
2637: | }
2638: | } // namespace gra

```

```

./src/MFlux.cc                   1/2
1: | Photon and other fluxes
2: |
3: | (c) 2017-2020 Mikael Mieskolainen
4: | Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: |
6: | C++
7: | #include <complex>
8: | #include <random>
9: | #include <vector>
10: |
11: | Own
12: | #include "Grannitt1/MFlux.h"
13: |
14: | namespace gra {
15: | namespace flux {
16: |
17: | using math::abs2;
18: | using math::msqrt;
19: |
20: | Apply non-collinear EPA fluxes at cross section level
21: | Use with full 2 - 2 kinematics
22: | double ApplyEPAFluxes(double amp2, gra::LORENTZSCALARs lts) {
23: |     const double gammaflux1 = lts.excite1;
24: |     gTiform:form::IncohFlux(lts.x1, lts.t1, lts.q1, lts.pfinal(1),M2())
25: |         f form:CohFlux(lts.x1, lts.t1, lts.q1);
26: |     const double gammaflux2 = lts.excite2;
27: |     gTiform:form::IncohFlux(lts.x2, lts.t2, lts.q2, lts.pfinal(2),M2())
28: |         f form:CohFlux(lts.x2, lts.t2, lts.q2);
29: |
30: |     const double phasespace = lts.s / lts.s_hat; // This gives problems at low masses
31: |     const double phasespace = 1.0 / (lts.x1 * lts.x2); // Consistent with kt-factorization
32: |
33: |     // Combine all
34: |     const double fluxes = gammaflux1 * gammaflux2 * phasespace;
35: |     const double tot = fluxes * amp2;
36: |
37: | }
38: |
39: | Apply fluxes to helicity amplitudes
40: | const double sqrt_fluxes = msqrt(fluxes); // to "amplitude level"
41: | for (const auto& i : aux::indices(lts.hamp1) | lts.hamp1 | sqrt_fluxes);
42: |
43: | }
44: |
45: |
46: | Apply Gamma-Gamma collinear Drees-Zeppenfeld (coherent flux) at cross section level
47: | Use with collinear kinematics
48: | double ApplyGammaFluxes(double amp2, gra::LORENTZSCALARs lts) {
49: |     // Evaluate gamma pdfs
50: |     const double f1 = form:IDFflux(lts.x1);
51: |     const double f2 = form:IDFflux(lts.x2);
52: |     const double phasespace = 1.0 / (lts.x1 * lts.x2);
53: |     const double tot = f1 * f2 * amp2 * phasespace;
54: |     return tot;
55: | }
56: |
57: |
58: | Apply Gamma-Gamma LIX-pdf (use at |mu| > 10 GeV) at cross section level
59: | Use with collinear kinematics
60: | double ApplyLIXFluxes(double amp2, gra::LORENTZSCALARs lts) {
61: |     // #pragma omp lock nequn for initialization ##
62: |     gra::mutex.lock();
63: |
64: |     // Not created yet
65: |     if (lts.GlobalPdfPtr == nullptr) {
66: |         std::string pdfname = lts.LHAPDFSET;
67: |
68: |         retry:
69: |         try {
70: |             lts.GlobalPdfPtr = LHAPDF::mkPDF(pdfname, 0);
71: |             lts.pdf_trials = 0; // fine
72: |         } catch (...) {
73: |             ++lts.pdf_trials;
74: |             std::string str = "ApplyLIXFluxes: Problem with reading LHAPDF '" + pdfname + "'";
75: |             aux::AutoDownloadLHAPDF(pdfname); // Try autodownload
76: |             gra::mutex.unlock(); // Remember before throw, otherwise deadlock
77: |             if (lts.pdf_trials >= 2) { // too many failures
78: |                 throw std::invalid_argument(str);
79: |             } else {
80: |                 goto retry;
81: |             }
82: |         }
83: |     }

```

```

./src/MFlux.cc                   2/2
84: |     gra::mutex.unlock();
85: |     ## MULTITHREADING UNLOCK ##
86: |
87: |     // pdf factorization scale
88: |     const double Q2 = lts.s_hat / 4.0;
89: |
90: |     // Evaluate gamma pdfs
91: |     double f1 = 0.0;
92: |     double f2 = 0.0;
93: |     try {
94: |         // Divide x out
95: |         f1 = lts.GlobalPdfPtr->xfxQ2(PDG:PDG_gamma, lts.x1, Q2) / lts.x1;
96: |         f2 = lts.GlobalPdfPtr->xfxQ2(PDG:PDG_gamma, lts.x2, Q2) / lts.x2;
97: |     } catch (...) { throw std::invalid_argument("ApplyGammaFluxes: Failed evaluating LHAPDF"); }
98: |
99: |     const double phasespace = 1.0 / (lts.x1 * lts.x2);
100: |     const double tot = f1 * f2 * amp2 * phasespace;
101: |     return tot;
102: | }
103: |
104: |
105: | } // namespace flux
106: | } // namespace gra

```

```

./src/MGraniitti.cc 1/24
1: // GRANIITTI Monte Carlo main class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <algorithm>
8: #include <complex>
9: #include <cstdlib>
10: #include <fstream>
11: #include <iostream>
12: #include <memory>
13: #include <mutex>
14: #include <random>
15: #include <regex>
16: #include <string>
17: #include <string_view>
18: #include <string>
19: #include <threads>
20: #include <vector>
21:
22: // Own
23: #include "Graniitti/Muax.h"
24: #include "Graniitti/MContinuum.h"
25: #include "Graniitti/MFactorized.h"
26: #include "Graniitti/MGraniitti.h"
27: #include "Graniitti/MKinematics.h"
28: #include "Graniitti/MMath.h"
29: #include "Graniitti/MNeuroChemobian.h"
30: #include "Graniitti/MParaton.h"
31: #include "Graniitti/MProcess.h"
32: #include "Graniitti/MQuasiElastic.h"
33: #include "Graniitti/MTimer.h"
34:
35: // HepMC3
36: #include "HepMC3/GenEvent.h"
37: #include "HepMC3/WriterAsciiHepMC2.h"
38: #include "HepMC3/WriterHEPEVT.h"
39:
40: // Libraries
41: #include "json.hpp"
42: #include "rang.hpp"
43:
44: namespace gra {
45:
46: // .....
47: // GLOBALS from MGlobal.h initialized by MGraniitti()
48:
49: // Model tune
50: std::string MODELPARAM;
51:
52: // Multithreading
53: std::mutex g_mutex;
54: std::exception_ptr globalExceptionPtr;
55:
56: // .....
57:
58: using json = nlohmann::json;
59:
60: using gra::aux::indices;
61: using gra::math::P;
62: using gra::math::msqrt;
63: using gra::math::pow;
64: using gra::math::pow2;
65: using gra::math::r;
66:
67: // Constructors
68: MGraniitti::MGraniitti() {
69: // .....
70: // Init program globals
71:
72: gra::MODELPARAM = "null";
73: gra::globalExceptionPtr = nullptr;
74: // .....
75:
76: // Print general layout
77: PrintInit();
78:
79: }
80:
81: // Destructor
82: MGraniitti::~MGraniitti() {
83: // Destroy processes

```

```

./src/MGraniitti.cc 2/24
84: for (std::size_t i = 0; i < pvec.size(); ++i) { delete pvec[i]; }
85:
86: std::cout << "MGraniitti [DONE]" << std::endl;
87:
88:
89: void MGraniitti::PrintHistograms() {
90: HistogramFusion();
91: proc->PrintHistograms();
92: }
93:
94: // Unify histogram bounds across threads, so the histogram fusion is possible
95: void MGraniitti::UnifyHistogramBounds() {
96: // Loop over all 2D histograms
97: for (auto const &point : proc->h1) {
98: int xbins = 0;
99: double xmin = 0;
100: double xmax = 0;
101:
102: // Fuse buffer values [start index]
103: for (std::size_t p = 1; p < pvec.size(); ++p) {
104: proc->h1[xpoint.first].FuseBuffer(pvec[p]->h1[xpoint.first]);
105: }
106:
107: proc->h1[xpoint.first].FlushBuffer();
108: proc->h1[xpoint.first].GetBounds(xbins, xmin, xmax);
109: proc->h1[xpoint.first].ResetBounds(xbins, xmin, xmax); // New start
110:
111: // Loop over processes and set histogram bounds [start index = 1]
112: for (std::size_t p = 1; p < pvec.size(); ++p) {
113: pvec[p]->h1[xpoint.first].ResetBounds(xbins, xmin, xmax);
114: }
115:
116: // Loop over all 2D histograms
117: for (auto const &point : proc->h2) {
118: int ybins = 0;
119: double ymin = 0;
120: double ymax = 0;
121:
122: int xbins = 0;
123: double xmin = 0;
124: double xmax = 0;
125:
126: // Fuse buffer values [start index]
127: for (std::size_t p = 1; p < pvec.size(); ++p) {
128: proc->h2[xpoint.first].FuseBuffer(pvec[p]->h2[xpoint.first]);
129: }
130:
131: proc->h2[xpoint.first].FlushBuffer();
132: proc->h2[xpoint.first].GetBounds(xbins, xmin, xmax, ybins, ymin, ymax);
133: proc->h2[xpoint.first].ResetBounds(xbins, xmin, xmax, ybins, ymin, ymax); // New start
134:
135: // Loop over processes and set histogram bounds [start index = 1]
136: for (std::size_t p = 1; p < pvec.size(); ++p) {
137: pvec[p]->h2[xpoint.first].ResetBounds(xbins, xmin, xmax, ybins, ymin, ymax);
138: }
139: }
140:
141:
142: // Fuse histograms for N-fold statistics
143: void MGraniitti::HistogramFusion() {
144: if (hist_fusion_done = false) {
145: // START with process index 1, because 0 is the base
146: for (std::size_t i = 1; i < pvec.size(); ++i) {
147: // Loop over all 2D histograms
148: for (auto const &point : proc->h1) {
149: proc->h1[xpoint.first] = proc->h1[xpoint.first] + pvec[i]->h1[xpoint.first];
150: }
151: // Loop over all 2D histograms
152: for (auto const &point : proc->h2) {
153: proc->h2[xpoint.first] = proc->h2[xpoint.first] + pvec[i]->h2[xpoint.first];
154: }
155: }
156: hist_fusion_done = true;
157: } else {
158: std::cout << "MGraniitti::HistogramFusion: Multithreaded histograms have been "
159: << "fused already"
160: << std::endl;
161: }
162: }
163:
164: // Generate events
165: void MGraniitti::Generate() {
166: if (NEVENTS > 0) { // Generate events

```

```

./src/MGraniitti.cc 3/24
167:
168: // Just before event generation
169: // (in order not to generate unnecessary empty files
170: // if file name / output type is changed)
171: InitFileOutput();
172: CallIntegrator(NEVENTS);
173: }
174:
175:
176: // This is called just before event generation
177: // Check again multiplier is done below, because if the output is already set
178: // externally,
179: // this function is not used.
180: void MGraniitti::InitFileOutput() {
181: if (NEVENTS > 0) {
182: if (OUTPUT == "") { // OUTPUT must be set
183: throw std::invalid_argument("MGraniitti::InitFileOutput: OUTPUT filename not set!");
184: }
185:
186: FULL_OUTPUT_STR = gra::aux::GetBasePath(2) + "output/" + OUTPUT + ". " + FORMAT;
187:
188: // .....
189: // Generator info
190: runInfo = std::make_shared<HepMC3::GenRunInfo>();
191:
192: struct HepMC3::GenRunInfo::ToolInfo generator = {
193: std::string("MGraniitti (" + gra::MODELPARAM + ")"),
194: std::string(aux::GetVersion()).substr(0, 5), std::string("Generator");
195: runInfo->tools().push_back(generator);
196:
197: struct HepMC3::GenRunInfo::ToolInfo config = {FULL_INPUT_STR, "1.0",
198: std::string("Steering card");
199: runInfo->tools().push_back(config);
200:
201: // .....
202:
203: if (FORMAT == "hepmc3" && outputHepMC3 = nullptr) {
204: outputHepMC3 = std::make_shared<HepMC3::WriterAscii>(FULL_OUTPUT_STR, runInfo);
205: } else if (FORMAT == "hepmc2" && outputHepMC2 = nullptr) {
206: outputHepMC2 = std::make_shared<HepMC3::WriterAsciiHepMC2>(FULL_OUTPUT_STR, runInfo);
207: } else if (FORMAT == "hepevt" && outputHEPEVT = nullptr) {
208: outputHEPEVT = std::make_shared<HepMC3::WriterHEPEVT>(FULL_OUTPUT_STR);
209: }
210: }
211:
212: // Return process numbers available
213: std::vector<std::string> MGraniitti::GetProcessNumbers() const {
214: std::cout << rang::style::bold;
215: std::cout << "Available processes: ";
216: std::cout << rang::style::reset << std::endl << std::endl;
217:
218: std::cout << rang::style::bold << "2-3" << "(N-2) LIPS: " << rang::style::reset << std::endl;
219:
220: std::cout << rang::style::bold << "2-3" << "1-3" << "(N-2) LIPS: " << rang::style::reset << std::endl;
221: std::vector<std::string> list1 = proc_F.ProcPtr.PrintProcesses();
222: std::cout << std::endl;
223:
224: std::cout << rang::style::bold << "2-3" << "N" LIPS: " << rang::style::reset << std::endl;
225: std::vector<std::string> list2 = proc_C.ProcPtr.PrintProcesses();
226: std::cout << std::endl;
227:
228: std::cout << rang::style::bold << "Quasielastic: " << rang::style::reset << std::endl;
229: std::vector<std::string> list3 = proc_Q.ProcPtr.PrintProcesses();
230: std::cout << std::endl;
231:
232: std::cout << rang::style::bold << "2-3" << "1" << "(N-1) collinear: " << rang::style::reset << std::endl;
233: std::vector<std::string> list4 = proc_P.ProcPtr.PrintProcesses();
234: std::cout << std::endl;
235:
236: // Concatenate all
237: list1.insert(list1.end(), list2.begin(), list2.end());
238: list1.insert(list1.end(), list3.begin(), list3.end());
239: list1.insert(list1.end(), list4.begin(), list4.end());
240:
241: return list1;
242: }
243:
244: // (Re-)assign the pointers to the local memory space
245: void MGraniitti::InitProcessMemory(std::string process, unsigned int seed) {
246: // These must be here!
247: PROCESS = process;
248:
249: // <<> process

```

```

./src/MGraniitti.cc 4/24
250: if (proc_Q.ProcPtr.ProcessExist(process)) {
251: proc_Q = MQuasiElastic(process, syntax);
252: proc = &proc_Q;
253: }
254: // <<> processes
255: } else if (proc_F.ProcPtr.ProcessExist(process)) {
256: proc_F = MFactorized(process, syntax);
257: proc = &proc_F;
258: }
259: // <<> processes
260: } else if (proc_C.ProcPtr.ProcessExist(process)) {
261: proc_C = MContinuum(process, syntax);
262: proc = &proc_C;
263: }
264: // <<> processes
265: } else if (proc_P.ProcPtr.ProcessExist(process)) {
266: proc_P = MParaton(process, syntax);
267: proc = &proc_P;
268: } else {
269: std::string str = "MGraniitti::InitProcessMemory: Unknown PROCESS: " + process;
270: GetProcessNumbers(); // This is done by main program
271: throw std::invalid_argument(str);
272: }
273:
274: // Set random seed last!
275: proc->Random::SetSeed(seed);
276:
277: void MGraniitti::InitMemory() {
278: // Init multithreading memory by making copies of the process **
279: for (std::size_t i = 0; i < pvec.size(); ++i) { delete pvec[i]; }
280: pvec.resize(CORES, nullptr);
281:
282: // Create new process objects for each thread
283: for (int i = 0; i < CORES; ++i) {
284: if (proc_Q.ProcPtr.ProcessExist(PROCESS)) {
285: pvec[i] = new MQuasiElastic(proc_Q);
286: } else if (proc_F.ProcPtr.ProcessExist(PROCESS)) {
287: pvec[i] = new MFactorized(proc_F);
288: } else if (proc_C.ProcPtr.ProcessExist(PROCESS)) {
289: pvec[i] = new MContinuum(proc_C);
290: } else if (proc_P.ProcPtr.ProcessExist(PROCESS)) {
291: pvec[i] = new MParaton(proc_P);
292: }
293: }
294:
295: // RANDOM SEED PER THREAD (IMPORTANT!)
296: for (int i = 0; i < CORES; ++i) {
297: const unsigned int tid = i + 1;
298: // Deterministic seed
299: const unsigned int thread_seed =
300: static_cast<unsigned int>(std::max(0, (int)pvec[i]->random.GetSeed()) * tid);
301: pvec[i]->random.SetSeed(thread_seed);
302: }
303:
304: // SET main control pointer to the first one
305: // (needed for printing etc.)
306: proc = pvec[0];
307: proc->PrintInit(HILJAA);
308:
309: // Set simple MC parameters
310: void MGraniitti::SetMCParam(MCPARAM &in) { mcparam = in; }
311:
312: // Set VEGAS parameters
313: void MGraniitti::SetVegasParam(const VEGASPARAM &in) { vparam = in; }
314:
315: // Read parameters from a single JSON file
316: void MGraniitti::ReadInput(const std::string &inputfile, const std::string &cmd_PROCESS) {
317: // Save if for later use
318: FULL_INPUT_STR = inputfile;
319:
320: std::cout << rang::fg::green << "MGraniitti::ReadInput: " + inputfile << rang::fg::reset
321: << std::endl;
322: << std::endl;
323:
324: ReadGeneralParam(inputfile);
325: ReadProcessParam(inputfile, cmd_PROCESS);
326: ReadModelParam(gra::MODELPARAM);
327:
328: // General parameter initialization
329: void MGraniitti::ReadGeneralParam(const std::string &inputfile) {

```



```

./src/MGraniitti.cc          5/24
333: // Read and parse
334: const std::string data = gra::aux::GetInputData(inputfile);
335: json j;
336: try {
337:     j = json::parse(data);
338: } catch (...) {
339:     std::string str = "MGraniitti:ReadGeneralParam: Error parsing " + inputfile +
340:                     " (Check for extra/missing commas)";
341:     throw std::invalid_argument(str);
342: }
343:
344: // JSON block identifier
345: const std::string XID = "GENERALPARAM";
346:
347: // Setup parameters (order is important)
348: SetNumberOfProcessors(j.at(XID).at("PROCESSORS"));
349: SetOutput(j.at(XID).at("OUTPUT"));
350: SetFormat(j.at(XID).at("FORMAT"));
351: SetWeights(j.at(XID).at("WEIGHTS"));
352: SetIntegrator(j.at(XID).at("INTEGRATOR"));
353: SetCores(j.at(XID).at("CORES"));
354:
355: // Save for later use
356: gra::MODELPARAM = j.at(XID).at("MODELPARAM");
357:
358: // General model parameters initialized from .json file
359:
360: void MGraniitti::ReadModelParam(const std::string inputfile) const {
361:     const std::string fullpath =
362:         gra::aux::GetBasePath(2) + "/modeldata/" + inputfile + "/GENERAL.json";
363:
364:     // Read generic blocks
365:     PARAM_SOFT: ReadParameters(fullpath);
366:     PARAM_STRUCTURE: ReadParameters(fullpath);
367:     PARAM_FLAT: ReadParameters(fullpath);
368:     PARAM_NSTRM: ReadParameters(fullpath);
369:
370:
371: // The rest are handled by specific amplitude classes
372:
373:
374: // Process parameter initialization, call proc-syopt constructor() after this
375: void MGraniitti::ReadProcessParam(const std::string inputfile, const std::string cmd_PROCESS) {
376:     // Read and parse
377:     const std::string data = gra::aux::GetInputData(inputfile);
378:     json j;
379:     try {
380:         j = json::parse(data);
381:     } catch (...) {
382:         std::string str = "MGraniitti:ReadProcessParam: Error parsing " + inputfile +
383:                         " (Check for extra/missing commas)";
384:         throw std::invalid_argument(str);
385:     }
386:     const std::string XID = "PROCESSPARAM";
387:
388:     // -----
389:     // Initialize process
390:
391:     std::string fullstring;
392:     if (cmd_PROCESS == "null") {
393:         fullstring = j.at(XID).at("PROCESS");
394:     } else { // commandline override
395:         fullstring = cmd_PROCESS;
396:     }
397:
398:     // -----
399:     // First separate possible extra arguments by ...
400:     std::vector<std::string> markerspos = aux::FindCommand(fullstring, "#");
401:
402:     if (markerspos.size() != 0) {
403:         syntax = aux::SplitCommands(fullstring.substr(markerspos[0]),
404:                                     fullstring.substr(0, markerspos[0] - 1));
405:     }
406:     // -----
407:
408:     // Now separate process and decay parts by "=="
409:     std::string PROCESS_PART = "";
410:     std::string DECAY_PART = "";
411:     std::size_t pos = 0;
412:     std::size_t pos1 = fullstring.find("=="); // ISOLATED Phase-Space
413:     std::size_t pos2 = fullstring.find("#");
414:
415:     pos = (pos1 != std::string::npos) ? pos1 : std::string::npos; // Try to find ==

```

```

./src/MGraniitti.cc          6/24
416:     if (pos == std::string::npos) {
417:         pos = (pos2 != std::string::npos) ? pos2 : std::string::npos; // Try to find &#
418:     }
419:
420:     if (pos != std::string::npos) {
421:         PROCESS_PART = fullstring.substr(0, pos - 1); // beginning to the pos-1
422:         DECAY_PART = fullstring.substr(pos + 2); // from pos+2 to the end
423:     } else {
424:         PROCESS_PART = fullstring; // No decay defined
425:     }
426:
427: // Trim extra spaces away
428: gra::aux::TrimExtraSpace(PROCESS_PART);
429: gra::aux::TrimExtraSpace(DECAY_PART);
430:
431: InitProcessMemory(PROCESS_PART, j.at(XID).at("MEMBER"));
432:
433: // -----
434: // SETUP process: process memory needs to be initialized before!
435:
436: proc->SetLambda(j.at(XID).at("LAMBDA"));
437: proc->SetScreening(j.at(XID).at("SCREENING"));
438: proc->SetExcitation(j.at(XID).at("EXCITATION"));
439: proc->SetIntegrator(j.at(XID).at("INTEGRATOR"));
440:
441: if (pos2 != std::string::npos) {
442:     if (PROCESS_PART.find("#") != std::string::npos &&
443:         PROCESS_PART.find("&#") != std::string::npos) {
444:         throw std::invalid_argument("
445:         "MGraniitti:ReadProcessParam: Phase space "
446:         "isolation used with &# or &# character");
447:     }
448:     proc->SetISOLATE(true);
449: }
450:
451:
452: // -----
453: // Decaymode setup
454: proc->SetDecayMode(DECAY_PART);
455:
456: // Read free resonance parameters (only if needed)
457: std::map<std::string, gra::PARAM_RES> RESONANCES;
458: if (PROCESS_PART.find("#") != std::string::npos) {
459:     // From .json input
460:     std::vector<std::string> RES = j.at(XID).at("RES");
461:
462:     // Command syntax override
463:     for (const auto &i : indices(syntax)) {
464:         if (syntax[i].id == "RES") {
465:             std::vector<std::string> temp;
466:             for (const auto &ix : syntax[i].arg) {
467:                 if (ix.second == "true" || ix.second == "1") {
468:                     temp.push_back(ix.first); // CG_960, ...
469:                 }
470:             }
471:             RES = temp;
472:         }
473:     }
474:
475:     // Read resonance data
476:     for (std::size_t i = 0; i < RES.size(); ++i) {
477:         const std::string str = "RES_" + RES[i] + ".json";
478:         RESONANCES[RES[i]] = gra::form::ReadResonance(str, proc->random);
479:     }
480:
481:     // Command syntax override "on-the-flight parameters"
482:     for (const auto &i : indices(syntax)) {
483:         if (syntax[i].id == "R") {
484:             // Take target string
485:             RESNAME = syntax[i].arg;
486:             if (syntax[i].target.size() == 1) {
487:                 RESNAME = syntax[i].target[0];
488:             } else if (syntax[i].target.size() == 0) {
489:                 throw std::invalid_argument("Syntax error: invalid R[] without any target []");
490:             } else if (syntax[i].target.size() > 1) {
491:                 throw std::invalid_argument("Syntax error: invalid R[] with multiple targets inside []");
492:             }
493:
494:             // Do we find target
495:             if (RESONANCES.find(RESNAME) == RESONANCES.end()) {
496:                 throw std::invalid_argument("Syntax error: invalid R[" + RESNAME + "] not found");
497:             }
498:             bool couplings_touched = false;

```

```

./src/MGraniitti.cc          7/24
499:     bool density_touched = false;
500:
501:     MMatrix<std::complex<double>> newrho(1, 1, 0.);
502:     if (RESONANCES[RESNAME].p.spin2 != 0) {
503:         const int N = RESONANCES[RESNAME].rho.size_row();
504:         newrho = MMatrix<std::complex<double>>(N, N, 0.);
505:     }
506:
507:     // Loop over keyval arguments
508:     for (const auto &ix : syntax[i].arg) {
509:         if (ix.first == "R") {
510:             RESONANCES[RESNAME].p.mass = std:: stod(ix.second);
511:             std::cout << "rang:fg:reson << "R[" << RESNAME
512:                 << " | new mass: << RESONANCES[RESNAME].p.mass << "rang:fg:reset
513:                 << std::endl;
514:         }
515:         if (ix.first == "W") {
516:             RESONANCES[RESNAME].p.width = std:: stod(ix.second);
517:             std::cout << "rang:fg:reson << "R[" << RESNAME
518:                 << " | new width: << RESONANCES[RESNAME].p.width << "rang:fg:reset
519:                 << std::endl;
520:         }
521:     }
522:
523:     // Set couplings
524:     for (std::size_t n = 0; n < RESONANCES[RESNAME].g_Tensor.size(); ++n) {
525:         if (ix.first == "c" + std::to_string(n)) {
526:             RESONANCES[RESNAME].g_Tensor[n] = std:: stod(ix.second);
527:             couplings_touched = true;
528:         }
529:     }
530:
531:     // Set diagonal spin density elements
532:     for (std::size_t n = 0; n <= (newrho.size_row() - 1) / 2; ++n) {
533:         const int J = (newrho.size_row() - 1) / 2;
534:         if (ix.first == "D" + std::to_string(n)) {
535:             const double value = std:: stod(ix.second);
536:             if (value < 0) {
537:                 throw std::invalid_argument("MGraniitti:ReadProcessParam: [R] JZ value < 0");
538:             }
539:
540:             newrho(newrho.size_row() - 1 - J - n, newrho.size_row() - 1 - J - n) = value;
541:             newrho(newrho.size_row() - 1 - J + n, newrho.size_row() - 1 - J + n) = value;
542:             density_touched = true;
543:         }
544:     }
545:
546: }
547:
548: // Normalize the trace and save it
549: if (density_touched) {
550:     const std::complex<double> trace = newrho.Trace();
551:     if (std::abs(trace) > 0) {
552:         newrho = newrho * (1.0 / trace);
553:     } else {
554:         throw std::invalid_argument("
555:         "MGraniitti:ReadProcessParam: [R] Spin density error with <= 0 trace");
556:     }
557:     RESONANCES[RESNAME].rho = newrho;
558:
559:     std::cout << "rang:fg:reson << "R[" << RESNAME
560:         << " | new spin density set: << std::endl;
561:     RESONANCES[RESNAME].rho.Print();
562:     std::cout << "rang:fg:reset << std::endl;
563: }
564:
565: // Print new coupling array
566: if (couplings_touched) {
567:     std::cout << "rang:fg:reson << "R[" << RESNAME
568:         << " | new production couplings set: g_Tensor = [";
569:     for (std::size_t k = 0; k < RESONANCES[RESNAME].g_Tensor.size(); ++k) {
570:         printf("%0.3E", RESONANCES[RESNAME].g_Tensor[k]);
571:     }
572:     if (k < RESONANCES[RESNAME].g_Tensor.size() - 1) { std::cout << ", ";
573:     std::cout << "]" << "rang:fg:reset << std::endl;
574: }
575:
576: }
577: proc->SetResonances(RESONANCES);
578:
579: // Setup resonance branching (final step)
580: proc->SetBranching();
581:

```

```

./src/MGraniitti.cc          8/24
582: // -----
583:
584: // Command syntax parameters
585: for (const auto &i : indices(syntax)) {
586:     if (syntax[i].id == "FLATMAP") { proc->SetFLATMAP(std::stoi(syntax[i].arg["SINGLETON"])); }
587:     if (syntax[i].id == "FLATBASE") { proc->SetFLATBASE(std::stoi(syntax[i].arg["SINGLETON"])); }
588:     if (syntax[i].id == "OFFSHIELD") { proc->SetOFFSHIELD(std::stoi(syntax[i].arg["SINGLETON"])); }
589:     if (syntax[i].id == "SPINDEC") { proc->SetSPINDEC(std::stoi(syntax[i].arg["SINGLETON"])); }
590:     if (syntax[i].id == "FRAME") { proc->SetFRAME(syntax[i].arg["SINGLETON"]); }
591:     if (syntax[i].id == "JMAX") { proc->SetJMAX(std::stoi(syntax[i].arg["SINGLETON"])); }
592: }
593:
594: // Always last (we need PDG tables)
595: std::vector<std::string> beam = j.at(XID).at("BEAM");
596: std::vector<double> energy = j.at(XID).at("ENERGY");
597: proc->SetInitialState(beam, energy);
598:
599: // Now rest of the parameters
600: ReadIntegralParam(inputfile);
601: ReadDecays(inputfile);
602: ReadFCuts(inputfile);
603: ReadVetoCuts(inputfile);
604:
605: // MC integrator parameter initialization
606: void MGraniitti::ReadIntegralParam(const std::string inputfile) {
607:     using namespace gra::aux;
608:
609:     // Read and parse
610:     const std::string data = gra::aux::GetInputData(inputfile);
611:     json j;
612:     try {
613:         j = json::parse(data);
614:     } catch (...) {
615:         std::string str = "MGraniitti:ReadIntegralParam: Error parsing " + inputfile +
616:                         " (Check for extra/missing commas)";
617:         throw std::invalid_argument(str);
618:     }
619:     const std::string XID = "INTEGRALPARAM";
620:
621:     // Numerical loop integration
622:     const int ND = j.at(XID).at("POMLOOP").at("ND");
623:     AssertRange(ND, -10, 10, "POMLOOP: ND", true);
624:     proc->Eikonal.Numerics.SetLoopDiscretization(ND);
625:
626:     // FLAT (naive) MC parameters
627:     mparam.PRECISION = j.at(XID).at("FLAT").at("PRECISION");
628:     mparam.PRECISION = j.at(XID).at("FLAT").at("PRECISION");
629:     AssertRange(mparam.PRECISION, 0.0, 1.0, "FLAT: PRECISION", true);
630:     mparam.MIN_EVENTS = j.at(XID).at("FLAT").at("MIN_EVENTS");
631:     AssertRange(mparam.MIN_EVENTS, 10, (unsigned int)1e9, "FLAT: MIN_EVENTS", true);
632:     SetMCParam(mparam);
633:
634:     try {
635:         j.at(XID).at("VEGAS").at("BINS");
636:     } catch (...) {
637:         std::cout << "MGraniitti:ReadIntegralParam: Did not found VEGAS parameter block "
638:             "from json input, using default";
639:     }
640:     return; // Did not find VEGAS parameter block at all, use default
641:
642:
643: // VEGAS parameters
644: VEGASPARAM vparam;
645: vparam.BINS = j.at(XID).at("VEGAS").at("BINS");
646: AssertRange(vparam.BINS, 0, (unsigned int)1e9, "VEGAS: BINS", true);
647: if (vparam.BINS % 2 != 0) {
648:     throw std::invalid_argument("VEGAS: BINS = + std::to_string(vparam.BINS) +
649:         " should be even number");
650: }
651:
652: vparam.LAMBDA = j.at(XID).at("VEGAS").at("LAMBDA");
653: AssertRange(vparam.LAMBDA, 1.0, 10.0, "VEGAS: LAMBDA", true);
654: vparam.NCALL = j.at(XID).at("VEGAS").at("NCALL");
655: AssertRange(vparam.NCALL, 10, (unsigned int)1e9, "VEGAS: NCALL", true);
656: vparam.ITER = j.at(XID).at("VEGAS").at("ITER");
657: AssertRange(vparam.ITER, 1, (unsigned int)1e9, "VEGAS: ITER", true);
658: vparam.CHIZMAX = j.at(XID).at("VEGAS").at("CHIZMAX");
659: AssertRange(vparam.CHIZMAX, 0.0, 1e3, "VEGAS: CHIZMAX", true);
660: vparam.PRECISION = j.at(XID).at("VEGAS").at("PRECISION");

```

./src/MGranitti.cc 9/24

```

665: AssertRange(vpm.PRECISION, (0.0, 1.0), "VEGAS:PRECISION", true);
666: vpm.DEBUG = j.at(XID).at("VEGAS").at("DEBUG");
667: AssertSet(vpm.DEBUG, {-1, 0, 1}, "VEGAS:DEBUG", true);
668:
669: SetVegasParam(vpm);
670:
671:
672: // Generator cuts
673: void MGranitti::ReadGenCuts(const std::string& inputfile) {
674: // Read and parse
675: const std::string data = gra::aux::GetInputData(inputfile);
676: json j;
677: try {
678: j = json::parse(data);
679: } catch (...) {
680: std::string str =
681: "MGranitti::ReadGenCuts: Error parsing " + inputfile + " (Check for extra/missing commas)";
682: throw std::invalid_argument(str);
683: }
684: const std::string XID = "GENCUTS";
685:
686: gra::GENCUT gcuts;
687:
688: // Continuum phase space class
689: if (PROCESS.find("<C>") != std::string::npos) {
690: // Daughter rapidity
691: std::vector<double> rap;
692: try {
693: std::vector<double> temp = j.at(XID).at("<C>").at("Rap");
694: rap = temp;
695: } catch (...) {
696: throw std::invalid_argument(
697: "MGranitti::ReadGenCuts: <C> phase space class requires from user: "
698: "\"GENCUTS\": { \"<C>\": { \"Rap\": [min, max] }}");
699: }
700: gcuts.rap_min = rap[0];
701: gcuts.rap_max = rap[1];
702: gra::aux::AssertCut(rap, "GENCUTS:<C>:Rap", true);
703:
704: // This is optional, intermediate kt
705: std::vector<double> kt;
706: try {
707: std::vector<double> temp = j.at(XID).at("<C>").at("kt");
708: kt = temp;
709: } catch (...) {
710: // Do nothing
711: }
712: if (kt.size() != 0) {
713: gcuts.kt_min = kt[0];
714: gcuts.kt_max = kt[1];
715: gra::aux::AssertCut(kt, (0.0, 1e32), "GENCUTS:<C>:kt", true);
716: }
717:
718: // This is optional, forward leg pt
719: std::vector<double> pt;
720: try {
721: std::vector<double> temp = j.at(XID).at("<C>").at("Pt");
722: pt = temp;
723: } catch (...) {
724: // Do nothing
725: }
726: if (pt.size() != 0) {
727: gcuts.forward_pt_min = pt[0];
728: gcuts.forward_pt_max = pt[1];
729: gra::aux::AssertCutRange(pt, (0.0, 1e32), "GENCUTS:<C>:Pt", true);
730: }
731:
732: // This is optional, forward leg excitation
733: std::vector<double> Xi;
734: try {
735: std::vector<double> temp = j.at(XID).at("<C>").at("Xi");
736: Xi = temp;
737: } catch (...) {
738: // Do nothing
739: }
740: if (Xi.size() != 0) {
741: gcuts.Xi_min = Xi[0];
742: gcuts.Xi_max = Xi[1];
743: gra::aux::AssertCutRange(Xi, (0.0, 1.0), "GENCUTS:<C>:Xi", true);
744: }
745:
746:
747: // Factorized phase space class

```

./src/MGranitti.cc 11/24

```

831: j = json::parse(data);
832: } catch (...) {
833: std::string str =
834: "MGranitti::ReadFidCuts: Error parsing " + inputfile + " (Check for extra/missing commas)";
835: throw std::invalid_argument(str);
836: }
837:
838: // Fiducial cuts
839: gra::FIDUCUT fcuts;
840: const std::string XID = "FIDUCUTS";
841:
842: try {
843: fcuts.active = j.at(XID).at("active");
844: } catch (...) {
845: return; // Did not find cut block at all
846: }
847:
848: // Particles
849: {
850: std::vector<double> eta = j.at(XID).at("PARTICLE").at("Eta");
851: fcuts.eta_min = eta[0];
852: fcuts.eta_max = eta[1];
853: gra::aux::AssertCut(eta, "FIDUCUTS:PARTICLE:Eta", true);
854:
855: std::vector<double> rap = j.at(XID).at("PARTICLE").at("Rap");
856: fcuts.rap_min = rap[0];
857: fcuts.rap_max = rap[1];
858: gra::aux::AssertCut(rap, "FIDUCUTS:PARTICLE:Rap", true);
859:
860: std::vector<double> pt = j.at(XID).at("PARTICLE").at("Pt");
861: fcuts.pt_min = pt[0];
862: fcuts.pt_max = pt[1];
863: gra::aux::AssertCut(pt, "FIDUCUTS:PARTICLE:Pt", true);
864:
865: std::vector<double> Et = j.at(XID).at("PARTICLE").at("Et");
866: fcuts.Et_min = Et[0];
867: fcuts.Et_max = Et[1];
868: gra::aux::AssertCut(Et, "FIDUCUTS:PARTICLE:Et", true);
869: }
870:
871: // System
872: {
873: std::vector<double> M = j.at(XID).at("SYSTEM").at("M");
874: fcuts.M_min = M[0];
875: fcuts.M_max = M[1];
876: gra::aux::AssertCut(M, "FIDUCUTS:SYSTEM:M", true);
877:
878: std::vector<double> Y = j.at(XID).at("SYSTEM").at("Rap");
879: fcuts.Y_min = Y[0];
880: fcuts.Y_max = Y[1];
881: gra::aux::AssertCut(Y, "FIDUCUTS:SYSTEM:Rap", true);
882:
883: std::vector<double> Pt = j.at(XID).at("SYSTEM").at("Pt");
884: fcuts.Pt_min = Pt[0];
885: fcuts.Pt_max = Pt[1];
886: gra::aux::AssertCut(Pt, "FIDUCUTS:SYSTEM:Pt", true);
887: }
888:
889: // Forward
890: {
891: std::vector<double> t = j.at(XID).at("FORWARD").at("t");
892: fcuts.forward_t_min = t[0];
893: fcuts.forward_t_max = t[1];
894: gra::aux::AssertCut(t, "FIDUCUTS:FORWARD:t", true);
895:
896: std::vector<double> M = j.at(XID).at("FORWARD").at("M");
897: fcuts.forward_M_min = M[0];
898: fcuts.forward_M_max = M[1];
899: gra::aux::AssertCut(M, "FIDUCUTS:FORWARD:M", true);
900: }
901:
902: // Set fiducial cuts
903: proc->SetFidCuts(fcuts);
904:
905: // Set user cuts
906: proc->SetUserCuts(j.at(XID).at("USERCUTS"));
907:
908:
909: // Fiducial cuts
910: void MGranitti::ReadVetoCuts(const std::string& inputfile) {
911: // Read and parse
912: const std::string data = gra::aux::GetInputData(inputfile);
913: json j;

```

./src/MGranitti.cc 10/24

```

748: if (PROCESS.find("<F>") != std::string::npos) {
749: // System rapidity
750: std::vector<double> Y;
751: try {
752: std::vector<double> temp = j.at(XID).at("<F>").at("Rap");
753: Y = temp;
754: } catch (...) {
755: throw std::invalid_argument(
756: "MGranitti::ReadGenCuts: <F> phase space class requires from user: "
757: "\"GENCUTS\": { \"<F>\": { \"Rap\": [min, max] }}");
758: }
759: gcuts.Y_min = Y[0];
760: gcuts.Y_max = Y[1];
761: gra::aux::AssertCut(Y, "GENCUTS:<F>:Rap", true);
762:
763: // System mass
764: std::vector<double> M;
765: try {
766: std::vector<double> temp = j.at(XID).at("<F>").at("M");
767: M = temp;
768: } catch (...) {
769: throw std::invalid_argument(
770: "MGranitti::ReadGenCuts: <F> phase space class requires from user: "
771: "\"GENCUTS\": { \"<F>\": { \"M\": [min, max] }}");
772: }
773: gcuts.M_min = M[0];
774: gcuts.M_max = M[1];
775: gra::aux::AssertCutRange(M, (0.0, 1e32), "GENCUTS:<F>:M", true);
776:
777: // This is optional, forward leg pt
778: std::vector<double> pt;
779: try {
780: std::vector<double> temp = j.at(XID).at("<F>").at("Pt");
781: pt = temp;
782: } catch (...) {
783: // Do nothing
784: }
785: if (pt.size() != 0) {
786: gcuts.forward_pt_min = pt[0];
787: gcuts.forward_pt_max = pt[1];
788: gra::aux::AssertCutRange(pt, (0.0, 1e32), "GENCUTS:<F>:Pt", true);
789: }
790:
791: // This is optional, forward leg excitation
792: std::vector<double> Xi;
793: try {
794: std::vector<double> temp = j.at(XID).at("<F>").at("Xi");
795: Xi = temp;
796: } catch (...) {
797: // Do nothing
798: }
799: if (Xi.size() != 0) {
800: gcuts.Xi_min = Xi[0];
801: gcuts.Xi_max = Xi[1];
802: gra::aux::AssertCutRange(Xi, (0.0, 1.0), "GENCUTS:<F>:Xi", true);
803: }
804:
805:
806: // Quasielastic phase space class
807: if (PROCESS.find("<Q>") != std::string::npos) {
808: std::vector<double> Xi;
809: try {
810: std::vector<double> temp = j.at(XID).at("<Q>").at("Xi");
811: Xi = temp;
812: } catch (...) {
813: throw std::invalid_argument(
814: "MGranitti::ReadGenCuts: <Q> phase space class requires from user: "
815: "\"GENCUTS\": { \"<Q>\": { \"Xi\": [min, max] }}");
816: }
817: gcuts.Xi_min = Xi[0];
818: gcuts.Xi_max = Xi[1];
819: gra::aux::AssertCutRange(Xi, (0.0, 1.0), "GENCUTS:<Q>:Xi", true);
820: }
821:
822: proc->SetGenCuts(gcuts);
823: }
824:
825: // Fiducial cuts
826: void MGranitti::ReadFidCuts(const std::string& inputfile) {
827: // Read and parse
828: const std::string data = gra::aux::GetInputData(inputfile);
829: json j;
830: try {
831:
832:
833:
834:
835:
836:
837:
838:
839:
840:
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842:
843:
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```



```

./src/MGranittt.cc          13/24
997:   gra::aux::PrintVersion();
998:   gra::aux::PrintBar("-");
999:
1000:   const double GB = pow3(1024.0);
1001:   printf("Running on %d CORES / %0.2f GB RAM at %s\n", gra::aux::HostName().c_str(),
1002:         std::thread::hardware_concurrency(), gra::aux::TotalSystemMemory() / GB);
1003:   gra::aux::DateAndTime().c_str();
1004:   int64_t size = 0;
1005:   int64_t free = 0;
1006:   gra::aux::GetUsage("MB", size, free, used);
1007:   printf("Path %s / size %d used %d free %0.1f %0.1f GB\n", size / GB, used / GB,
1008:         free / GB);
1009:   std::cout << "Program path: " << gra::aux::GetBasePath(2) << std::endl;
1010:   std::cout << gra::aux::SystemName() << std::endl;
1011:   gra::aux::PrintBar("-");
1012:   }
1013:   }
1014:   }
1015:
1016: void MGranittt::Initialize() {
1017:   // ** ALWAYS HERE ONLY AS LAST! **
1018:   proc->post_Constructor();
1019:
1020:   // Print out basic information
1021:   std::cout << std::endl;
1022:   std::cout << rang::style::bold << "General setup:" << rang::style::reset << std::endl;
1023:   std::cout << "Output file: " << OUTPUT << std::endl;
1024:   std::cout << "Output format: " << FORMAT << std::endl;
1025:   std::cout << "Multithreading: " << CORES << std::endl;
1026:   std::cout << "Integrator: " << INTEGRATOR << std::endl;
1027:   std::cout << "Number of events: " << NEUTRONS << std::endl;
1028:   std::cout << "Parameter setup: " << gra::MODELPARAM << std::endl;
1029:
1030:   std::string str = "WEIGHTED == true ? 'weighted' : 'unweighted';";
1031:   std::cout << rang::fg::green << "Generation mode: " << str << rang::fg::reset << std::endl;
1032:   std::cout << std::endl;
1033:
1034:   // If Eikonal is not yet initialized and pomeron screening loop is on
1035:   if (proc->Eikonal::IsInitialized() == false && proc->GetScreening() == true) {
1036:     proc->Eikonal::SConstructor(proc->GetMandelstam_s(), proc->GetInitialState(), false);
1037:   }
1038:
1039:   // Integrate
1040:   CallIntegrator(0);
1041:
1042:   // Initialize with external Eikonal
1043:   void MGranittt::Initialize(const MKikonal eikonal_in) {
1044:     // ** ALWAYS HERE ONLY AS LAST! **
1045:     proc->post_Constructor();
1046:
1047:     // Set input eikonal
1048:     proc->SetKikonal(eikonal_in);
1049:
1050:     // Integrate
1051:     CallIntegrator(0);
1052:   }
1053:
1054:   void MGranittt::CallIntegrator(unsigned int N) {
1055:     // Initialize global cores
1056:     if (N == 0) { global_tictoc = MTimer(true); }
1057:
1058:     // Sample the phase space
1059:     if (INTEGRATOR == "VEGAS") {
1060:       SampleVegas(N);
1061:     }
1062:     else if (INTEGRATOR == "FLAT") {
1063:       SampleFlat(N);
1064:     }
1065:     else if (INTEGRATOR == "NEURO") {
1066:       SampleNeuro(N);
1067:     }
1068:     else {
1069:       throw std::invalid_argument("MGranittt::CallIntegrator: Unknown INTEGRATOR = " + INTEGRATOR +
1070:                                   " (use VEGAS, FLAT, NEURO)");
1071:     }
1072:   }
1073:
1074:   // Initialize and generate events using VEGAS MC
1075:   void MGranittt::SampleVegas(unsigned int N) {
1076:     if (N == 0) {
1077:       InitMIMemory();
1078:       MODE = 0; // Pure integration
1079:     }
1080:   }

```

```

./src/MGranittt.cc          14/24
1080:   MODE = 1; // Event generation
1081:
1082:   // *****
1083:   // *****
1084:   // *****
1085:   // *****
1086:   // *****
1087:   // *****
1088:   // *****
1089:   // *****
1090:   // Pure integration mode
1091:   if (MODE == 0) {
1092:     const double MINTIME = 0.1; // Seconds
1093:     unsigned int BURIN_ITER = 3; // BUR-IN iterations (default!)
1094:
1095:     // Initialize GRID
1096:     do {
1097:       unsigned int init = 0;
1098:       int factor = 0;
1099:
1100:       do { // Loop until stable
1101:         factor = VEGAS(init, vparam.NCALL, BURIN_ITER, N);
1102:         vparam.NCALL = vparam.NCALL * factor;
1103:         if (factor == 1) { break; } // We are done
1104:       } while (true);
1105:
1106:       // Increase CALLS and re-run, if too fast
1107:       if (istertime < MINTIME) {
1108:         const double time_per_iter = istertime / vparam.NCALL;
1109:
1110:         // Max, because this is only minimum condition
1111:         vparam.NCALL =
1112:           std::max(unsigned int)vparam.NCALL, (unsigned int)(MINTIME / time_per_iter);
1113:         VEGAS(init, vparam.NCALL, BURIN_ITER, N);
1114:       }
1115:
1116:       // Now re-calculate by skipping burn-in (init = 0) iterations
1117:       // because they deteriorate the integral value by bad grid
1118:       init = 1;
1119:       VEGAS(init, vparam.NCALL, vparam.ITER, N);
1120:
1121:       if (factor == 1) {
1122:         break;
1123:       } else {
1124:         BURIN_ITER = 2 * BURIN_ITER;
1125:       }
1126:     } while (true);
1127:   }
1128:
1129:   // Event generation mode
1130:   if (MODE == 1) {
1131:     const unsigned int init = 2;
1132:     const unsigned int itermn = 1E3;
1133:     VEGAS(init, vparam.NCALL * 10, itermn, N);
1134:   }
1135:
1136:   // Multithreaded VEGAS integrator
1137:   // [close to optimal importance sampling iff
1138:   // integrated factorizes dimension by dimension]
1139:   // Original algorithm from:
1140:   // [REFERENCE: Lepage, G.F. Journal of Computational Physics, 1978]
1141:   // [http://en.wikipedia.org/wiki/VEGAS_algorithm]
1142:
1143:   int MGranittt::VEGAS(unsigned int init, unsigned int itermn, unsigned int N) {
1144:     // First the initialization
1145:     VD.Init(init, vparam);
1146:
1147:     if (init == 1 && HILJAA) {
1148:       gra::aux::ClearProgress();
1149:       std::cout << rang::style::bold;
1150:       printf("VEGAS basic iterations: %d\n", N);
1151:       std::cout << rang::style::reset;
1152:     }
1153:
1154:     if (init == 1 && HILJAA) {
1155:       if (init == 1 && HILJAA) {
1156:         gra::aux::ClearProgress();
1157:         gra::aux::PrintBar("-");
1158:         std::cout << rang::style::bold;
1159:         std::cout << "VEGAS importance sampling:" << std::endl << std::endl;
1160:         std::cout << rang::style::reset;
1161:       }
1162:
1163:       stat.sigma_err2 = pow2(stat.sigma_err);
1164:
1165:       // Ch12
1166:       double ch12this = (VD.sumch12 - pow2(VD.sumdata / VD.sweight) / (iter + 1E-5));
1167:       ch12this = (ch12this < 0 ? 0.0 : ch12this);
1168:       stat.ch12 = ch12this;
1169:       // -----
1170:       // Got enough events generated
1171:       if (MODE == 1 && stat.generated >= N) { goto stop; }
1172:
1173:       // Fatal error and convergence restart treatment
1174:       if (MODE == 0 && iter > 0) {
1175:         if (std::isnan(stat.sigma) || std::isnan(stat.sigma)) {
1176:           gra::aux::ClearProgress();
1177:           throw std::invalid_argument(
1178:             "VEGAS: Integral inf/nan; FATAL ERROR, cuts or parameters =
1179:             "need fixing. Try <P> phase space class if not in use.");
1180:         }
1181:
1182:         if (stat.sigma < 1e-60) {
1183:           gra::aux::ClearProgress();
1184:           throw std::invalid_argument(
1185:             "VEGAS: Integral < 1e-60; Check VEGAS parameters, decaymode =
1186:             "nanity, generation and fiducial cut! Try <P> phase space class if not in use.");
1187:         }
1188:
1189:         if (stat.ch12 > 50) {
1190:           gra::aux::ClearProgress();
1191:           printf("VEGAS: ch12 = %0.1f > 50; Convergence problem, increasing 10 x calls. %d",
1192:                 stat.ch12);
1193:           printf("Try <P> phase space class if not in use. %d",
1194:                 return 10; // 10 times more calls
1195:           );
1196:         }
1197:
1198:         if (vparam.DEBUG >= 0) {
1199:           gra::aux::ClearProgress();
1200:           printf(
1201:             "VEGAS: local iter = %d; integral = %0.5E +- std =
1202:             \"%0.5E %t (global integral = %0.5E +- std = %0.5E) %t =
1203:             \"%0.5E %t, %t_this, gra::math::msqrt(I2_this), stat.sigma, stat.sigma_err, ch12this);
1204:
1205:           if (vparam.DEBUG > 0) {
1206:             for (std::size_t j = 0; j < VD.FDIM; ++j)
1207:               printf("VEGAS: data for dimension %d = %d\n", j, VD.FDIM);
1208:
1209:             for (std::size_t i = 0; i < vparam.BINS; ++i)
1210:               printf(
1211:                 "mat[%3lu][j] = %0.5E, fmat[%3lu][j] = %0.5E %t",
1212:                 i, VD.xmat[i][j], VD.fmat[i][j]);
1213:             );
1214:           }
1215:
1216:           // We are not below the ch12 or precision condition -> increase iterations
1217:           if (init > 0) { do not consider in burn-in (init = 0) mode }
1218:           if (iter == (iterm - 1) && stat.ch12 > vparam.CH12MAX) ||
1219:               (iter == (iterm - 1) && stat.sigma_err / stat.sigma > vparam.PRECISION) {
1220:             ++iterm;
1221:           }
1222:
1223:           // Status and grid optimization
1224:           if (MODE == 0) {
1225:             if (sttime.ElapsedSec() > 2.0) { init == 0 }
1226:             PrintStatus(stat.evaluations, M, local_tictoc, -1.0);
1227:             stime.Reset();
1228:
1229:             if (sttime.ElapsedSec() < 0.01) {
1230:               gra::aux::PrintProgress((iter + 1) / static_cast<double>(iterm));
1231:               atime.Reset();
1232:             }
1233:
1234:             // ***** Update VEGAS grid (not during event generation) *****
1235:             VD.OptimizeGrid(vparam);
1236:           }
1237:
1238:           // Unify histogram boundaries after burn-in across different threads
1239:           // (due to adaptive histogramming)

```

```

./src/MGranittt.cc          15/24
1163:   printf("Multithreading CORES = %d\n", CORES);
1164:
1165:   printf(" dimension = %u\n", VD.FDIM);
1166:   printf(" %d\n", N);
1167:   printf(" itermn = %d\n", itermn);
1168:   printf(" calls = %d\n", calls);
1169:   printf(" BINS = %u\n", vparam.BINS);
1170:   printf(" PRECISION = %0.4f\n", vparam.PRECISION);
1171:   printf(" CH12MAX = %0.4f\n", vparam.CH12MAX);
1172:   printf(" LAMBDA = %0.4f\n", vparam.LAMBDA);
1173:   printf(" DEBUD = %d\n", vparam.DEBUD);
1174:   gra::aux::PrintBar("-");
1175:
1176:   // Reset local timers
1177:   local_tictoc = MTimer(true);
1178:
1179:   MTimer stime = MTimer(true); // For statusprint
1180:   atime = MTimer(true); // For programmer
1181:
1182:   // Create number of calls per thread, their sum = calls
1183:   std::vector<unsigned int> LOCALcalls = VD.GetLocalCalls(calls, CORES);
1184:
1185:   MTimer gridtic;
1186:
1187:   // VEGAS grid iterations
1188:   for (std::size_t iter = 0; iter < itermn; ++iter) {
1189:     if (init == 0 && iter == 2) { // Save time for one iteration
1190:       itertime = gridtic.ElapsedSec() / 3; // Average over 3 iter
1191:     }
1192:
1193:     for (std::size_t j = 0; j < VD.FDIM; ++j) {
1194:       for (std::size_t i = 0; i < vparam.BINS; ++i) {
1195:         VD.fmat[i][j] = 0.0;
1196:         VD.f2mat[i][j] = 0.0;
1197:       }
1198:     }
1199:
1200:     // Init here outside parallel processing
1201:     VD.fsum = 0.0;
1202:     VD.f2sum = 0.0;
1203:
1204:     // -----
1205:     // SPAWN PARALLEL PROCESSING HERE
1206:     // -----
1207:     std::vector<std::thread> threads;
1208:     for (int tid = 0; tid < CORES; ++tid) {
1209:       threads.push_back(std::thread(+1) { VEGASMultiThread(N, tid, init, LOCALcalls[tid]); });
1210:     }
1211:
1212:     // Loop again to join the threads
1213:     for (auto it : threads) { it.join(); }
1214:
1215:     if (gra::global_exception) { // Exception handling of threads
1216:       std::rethrow_exception(gra::global_exception);
1217:     }
1218:
1219:     // Estimates based on this iteration
1220:     VD.fsum *= 1.0 / static_cast<double>(calls);
1221:     VD.f2sum = pow2(1.0 / static_cast<double>(calls));
1222:
1223:     // Local integral estimate
1224:     const double I_this = VD.fsum;
1225:
1226:     VD.f2sum = msqrt(VD.f2sum * calls);
1227:     VD.f2sum = (VD.f2sum - VD.fsum) * (VD.f2sum + VD.fsum); // - +
1228:     if (VD.f2sum <= 0.0) { VD.f2sum = vparam.EPS; }
1229:
1230:     // Local error squared estimate
1231:     const double I2_this = VD.f2sum * 1.0 / static_cast<double>(calls);
1232:
1233:     // Update global estimate
1234:     const double alpha = 1.0 / I2_this;
1235:
1236:     VD.sumdata += alpha * I_this;
1237:     VD.sumch12 += alpha * pow2(I_this);
1238:     VD.sweight += alpha;
1239:
1240:     // Integral and its error estimate
1241:     stat.sigma = VD.sumdata / VD.sweight;
1242:     stat.sigma_err = msqrt(1.0 / VD.sweight);

```

```

./src/MGranittt.cc          16/24
1246:   stat.sigma_err2 = pow2(stat.sigma_err);
1247:
1248:   // Ch12
1249:   double ch12this = (VD.sumch12 - pow2(VD.sumdata / VD.sweight) / (iter + 1E-5));
1250:   ch12this = (ch12this < 0 ? 0.0 : ch12this);
1251:   stat.ch12 = ch12this;
1252:   // -----
1253:   // Got enough events generated
1254:   if (MODE == 1 && stat.generated >= N) { goto stop; }
1255:
1256:   // Fatal error and convergence restart treatment
1257:   if (MODE == 0 && iter > 0) {
1258:     if (std::isnan(stat.sigma) || std::isnan(stat.sigma)) {
1259:       gra::aux::ClearProgress();
1260:       throw std::invalid_argument(
1261:         "VEGAS: Integral inf/nan; FATAL ERROR, cuts or parameters =
1262:         "need fixing. Try <P> phase space class if not in use.");
1263:     }
1264:
1265:     if (stat.sigma < 1e-60) {
1266:       gra::aux::ClearProgress();
1267:       throw std::invalid_argument(
1268:         "VEGAS: Integral < 1e-60; Check VEGAS parameters, decaymode =
1269:         "nanity, generation and fiducial cut! Try <P> phase space class if not in use.");
1270:     }
1271:
1272:     if (stat.ch12 > 50) {
1273:       gra::aux::ClearProgress();
1274:       printf("VEGAS: ch12 = %0.1f > 50; Convergence problem, increasing 10 x calls. %d",
1275:             stat.ch12);
1276:       printf("Try <P> phase space class if not in use. %d",
1277:             return 10; // 10 times more calls
1278:       );
1279:
1280:     if (vparam.DEBUG >= 0) {
1281:       gra::aux::ClearProgress();
1282:       printf(
1283:         "VEGAS: local iter = %d; integral = %0.5E +- std =
1284:         \"%0.5E %t (global integral = %0.5E +- std = %0.5E) %t =
1285:         \"%0.5E %t, %t_this, gra::math::msqrt(I2_this), stat.sigma, stat.sigma_err, ch12this);
1286:
1287:       if (vparam.DEBUG > 0) {
1288:         for (std::size_t j = 0; j < VD.FDIM; ++j)
1289:           printf("VEGAS: data for dimension %d = %d\n", j, VD.FDIM);
1290:
1291:         for (std::size_t i = 0; i < vparam.BINS; ++i)
1292:           printf(
1293:             "mat[%3lu][j] = %0.5E, fmat[%3lu][j] = %0.5E %t",
1294:             i, VD.xmat[i][j], VD.fmat[i][j]);
1295:           );
1296:       }
1297:
1298:       // We are not below the ch12 or precision condition -> increase iterations
1299:       if (init > 0) { do not consider in burn-in (init = 0) mode }
1300:       if (iter == (iterm - 1) && stat.ch12 > vparam.CH12MAX) ||
1301:           (iter == (iterm - 1) && stat.sigma_err / stat.sigma > vparam.PRECISION) {
1302:         ++iterm;
1303:       }
1304:
1305:       // Status and grid optimization
1306:       if (MODE == 0) {
1307:         if (sttime.ElapsedSec() > 2.0) { init == 0 }
1308:         PrintStatus(stat.evaluations, M, local_tictoc, -1.0);
1309:         stime.Reset();
1310:
1311:         if (sttime.ElapsedSec() < 0.01) {
1312:           gra::aux::PrintProgress((iter + 1) / static_cast<double>(iterm));
1313:           atime.Reset();
1314:         }
1315:
1316:         // ***** Update VEGAS grid (not during event generation) *****
1317:         VD.OptimizeGrid(vparam);
1318:       }
1319:
1320:       // Unify histogram boundaries after burn-in across different threads
1321:       // (due to adaptive histogramming)

```

```

./src/MGranitti.cc 17/24
1329: if (init == 0 && iter == itermin - 1) { UnitHistogramBounds(); }
1330: } // Main grid iteration loop
1331:
1332: stop; // We jump here once finished (GOTO point)
1333:
1334: // Final statistics
1335: if (init > 0) { PrintStatistics(N); }
1336:
1337: return 1; // Return 1 for good
1338: }
1339:
1340: // This is called once for every VEGAS grid iteration
1341: void MGranitti::VEGASMultithread(unsigned int N, unsigned int THREAD_ID, unsigned int init,
1342: double z0 = 0.0; unsigned int LOCALcalls) {
1343: double z0 = 0.0; unsigned int LOCALcalls) {
1344: double ac = 0.0;
1345: }
1346:
1347: try {
1348: for (std::size_t k = 0; k < LOCALcalls; ++k) { // ** LOCALcalls = calls / CORES **
1349:
1350: double vegasweight = 1;
1351:
1352: // Phase space point vector
1353: std::vector<double> xpoint(pvec[THREAD_ID] ->getdIPSDim(), 0.0);
1354:
1355: // Loop over dimensions and construct random vector xpoint
1356: std::vector<double> indvec(VD.FDIM, 0);
1357: for (std::size_t j = 0; j < VD.FDIM; ++j) {
1358: // Draw random number
1359: int proc[THREAD_ID] = random.U(0, 1) * vparam.BINS + 1.0;
1360: indvec[j] =
1361: std::max((unsigned int)1, std::min((unsigned int)zn, (unsigned int)vparam.BINS));
1362:
1363: if (indvec[j] > 1) {
1364: z0 = VD.xmat[indvec[j] - 1][j] - VD.xmat[indvec[j] - 2][j];
1365: ac = VD.xmat[indvec[j] - 2][j] + (zn - indvec[j]) * z0;
1366: } else {
1367: z0 = VD.xmat[indvec[j] - 1][j];
1368: ac = (zn - indvec[j]) * z0;
1369: }
1370:
1371: // Multidim space vector component
1372: xpoint[j] = VD.invec[j] + ac * VD.dvvec[j];
1373: vegasweight *= z0 * vparam.BINS;
1374:
1375: } // VD.FDIM loop
1376:
1377: // ***** Call the process under integration to get the weight *****
1378:
1379: gra:AuxInData aux;
1380: aux.vegasweight = vegasweight;
1381: aux.burn_in_mode = (init == 0) ? true : false;
1382:
1383: const double W = pvec[THREAD_ID] ->EventWeight(xpoint, aux);
1384:
1385: // *****
1386:
1387: // ** Multithreading lock (FAST SECTION) **
1388: gra:mutex.lock();
1389:
1390: // Increase statistics
1391: stat.Accumulate(aux);
1392:
1393: // ** Importance weighting **
1394: const double f = W * vegasweight;
1395: const double f2 = pow2(f);
1396:
1397: VD.fsุม += f;
1398: VD.f2sum += f2;
1399:
1400: // Loop over dimensions and add importance weighted results
1401: for (std::size_t j = 0; j < VD.FDIM; ++j) {
1402: VD.fmat[indvec[j] - 1][j] += f;
1403: VD.f2mat[indvec[j] - 1][j] += f2;
1404: }
1405:
1406: // -----
1407: // Initialization (integration) mode
1408: if (GMODE == 0) {
1409: // Do not consider burn-in phase weights (unstable)
1410: if (init != 0) {
1411:

```

```

./src/MGranitti.cc 19/24
1496: PrintStatus(stat. evaluations, N, local.tictoc, 10.0);
1497:
1498: if ( (stat.sigma_err / stat.sigma) < mparam.PRECISION &&
1499: stat. evaluations > mparam.MIN_EVENTS ) {
1500: break;
1501: }
1502: // Progressbar
1503: if (atime.ElapsedSec() > 0.1) {
1504: gra:aux:PrintProgress(stat. evaluations / static_cast<double>(mparam.MIN_EVENTS));
1505: atime.Reset();
1506: }
1507:
1508: // Event generation mode
1509: if (GMODE == 1) {
1510: SaveEvent(proc, M, stat.maxW, aux);
1511: PrintStatus(stat.generated, N, local.tictoc, 10.0);
1512: if (stat.generated >= N) { break; }
1513: }
1514: // Progressbar
1515: if (atime.ElapsedSec() > 0.1) {
1516: gra:aux:PrintProgress(stat.generated / static_cast<double>(N));
1517: atime.Reset();
1518: }
1519:
1520: }
1521:
1522: PrintStatus(stat.generated, N, local.tictoc, -1.0);
1523: PrintStatistics(N);
1524: }
1525:
1526: // Neural net Monte Carlo (small scale PROTOTYPE)
1527: void MGranitti::SampleNeuro(unsigned int N) {
1528: // Integration mode
1529: if (N == 0) {
1530: GMODE = 0;
1531: proc->PrintInit(HILJAA);
1532: }
1533: // Event generation mode
1534: if (N > 0) { GMODE = 1; }
1535:
1536: // Get dimension of the phase space
1537: const unsigned int D = proc->getdIPSDim();
1538:
1539: // Reset timers
1540: local.tictoc = MTimer(true);
1541: atime = MTimer(true);
1542:
1543: // Reservation (test)
1544: // Using namespace std::placeholders; // ..._1, _2, ... come from here
1545: // std::function<double>(const std::vector<double>&, int, double,
1546: // std::vector<double>&,
1547: // boost::shared_ptr<MQuasiElastic::EventWeight, sproc_0, _1, _2, _3, _4, _5>)
1548: // std::bind(4MQuasiElastic::EventWeight, sproc_0, _1, _2, _3, _4, _5);
1549:
1550: if (N == 0) {
1551: // -----
1552: // Bind the object and call it
1553: gra:neurojac:procptr = proc; // First set address
1554: // -----
1555:
1556: gra:neurojac:MNeuroJacobian neurojac;
1557: gra:neurojac:MATCHSIEE = 100;
1558:
1559: // Set network layer dimensions [first, ... output]
1560: gra:neurojac:par.D = D; // Integrand dimension
1561:
1562: gra:neurojac:par.L.push_back(gra:neurojac:layer(2, D)); // Input
1563: gra:neurojac:par.L.push_back(gra:neurojac:layer(2, 2));
1564: gra:neurojac:par.L.push_back(gra:neurojac:layer(2, 2));
1565: gra:neurojac:par.L.push_back(gra:neurojac:layer(2, 2)); // Output
1566:
1567: neurojac.Optimize();
1568:
1569: // -----
1570: // How to the event generation
1571:
1572: // Lambda capture
1573: std::vector<double> u(D);
1574:
1575: auto NeuroSample = [&](gra:AuxInData aux) {
1576:
1577:

```

```

./src/MGranitti.cc 18/24
1412: // Maximum raw weight (for general information, not used
1413: // here)
1414: if (W > stat.maxW) stat.maxW = W;
1415: // Maximum total VEGAS importance weighted
1416: if (f > stat.maxf) stat.maxf = f;
1417: }
1418:
1419: gra:mutex.unlock();
1420: // ** Multithreading unlock (FAST SECTION) **
1421:
1422: // -----
1423: // Event generation mode
1424: if (GMODE == 1) {
1425: // Enough events
1426: if (stat.generated == (unsigned int)getNumberOfEvents()) { break; }
1427:
1428: // Event trial
1429: SaveEvent(pvec[THREAD_ID], f, stat.maxf, aux);
1430:
1431: if (THREAD_ID == 0 && atime.ElapsedSec() > 0.5) {
1432: PrintStatus(stat.generated, N, local.tictoc, 10.0);
1433: gra:aux:PrintProgress(stat.generated / static_cast<double>(N));
1434: atime.Reset();
1435: }
1436:
1437: // calls loop
1438: } catch (...) {
1439: // Set the global exception pointer if exception arises
1440: // This is because of multithreading
1441: gra:globalExceptionPtr = std::current_exception();
1442: }
1443:
1444: // Generate events using plain simple MC (for reference/DEBUG purposes)
1445: void MGranitti::SampleFlat(unsigned int N) {
1446: // Integration mode
1447: if (N == 0) {
1448: GMODE = 0;
1449: proc->PrintInit(HILJAA);
1450:
1451: // Event generation mode
1452: if (N > 0) { GMODE = 1; }
1453:
1454: // Get dimension of the phase space
1455: const unsigned int dim = proc->getdIPSDim();
1456: std::vector<double> randvec(dim, 0.0);
1457:
1458: // Reset local timer
1459: local.tictoc = MTimer(true);
1460:
1461: // Progressbar
1462: atime = MTimer(true);
1463:
1464: // Reservation (test)
1465: // Using namespace std::placeholders; // ..._1, _2, ... come from here
1466: // std::function<double>(const std::vector<double>&, int, double,
1467: // std::vector<double>&,
1468: // boost::shared_ptr<MQuasiElastic::EventWeight, sproc_0, _1, _2, _3, _4, _5>)
1469: // std::bind(4MQuasiElastic::EventWeight, sproc_0, _1, _2, _3, _4, _5);
1470:
1471: // Event loop
1472: while (true) {
1473: // Generate new random numbers **
1474: for (const auto &i: indices.randvec) { randvec[i] = proc->random.U(0, 1); }
1475:
1476: // Generate event
1477: // Used for in-out control of the process
1478: gra:AuxInData aux;
1479: aux.vegasweight = 1.0;
1480: aux.burn_in_mode = false;
1481: const double W = proc->EventWeight(randvec, aux);
1482:
1483: // Increase statistics
1484: stat.Accumulate(aux);
1485: stat.Msum += W;
1486: stat.W2sum += pow2(W);
1487:
1488: // Update cross section estimate
1489: stat.CalculateCrossSection();
1490:
1491: // Initialization
1492: if (GMODE == 1) {
1493: stat.maxW = (W > stat.maxW) ? W : stat.maxW; // Update maximum weight

```

```

./src/MGranitti.cc 20/24
1578: // Prior p(x) distribution sampling
1579: VectorXdXu(z,u.size());
1580: for (std::size_t i = 0; i < D; ++i) { z[i] = proc->random.G(0, 1); }
1581: const double p = val(gra:neurojac:igaussprob(z, 0, 1));
1582:
1583: // Evaluate network map
1584: VectorXdXu_u = gra:neurojac:u_net(z, u);
1585:
1586: // Evaluate the Jacobian matrix du/dz
1587: MatrixXd dudz = jacobian(gra:neurojac:u_net, u, z);
1588:
1589: // Abs Jacobian determinant and inverse prior
1590: const double jacweight = abs(dudz.determinant()) / p;
1591:
1592: for (std::size_t i = 0; i < D; ++i) { u[i] = val(u[i]); }
1593:
1594: // Evaluate event weight
1595: aux.vegasweight = jacweight;
1596: aux.burn_in_mode = false;
1597: const double weight = proc->EventWeight(u, aux) * jacweight;
1598:
1599: return weight;
1600: }
1601:
1602: // Event loop
1603: while (true) {
1604: // aux used for in-out control of the process
1605: gra:AuxInData aux;
1606: const double W = NeuroSample(aux);
1607:
1608: // Increase statistics
1609: stat.Accumulate(aux);
1610: stat.Msum += W;
1611: stat.W2sum += pow2(W);
1612:
1613: // Update cross section estimate
1614: stat.CalculateCrossSection();
1615:
1616: // Initialization
1617: if (GMODE == 0) {
1618: stat.maxW = (W > stat.maxW) ? W : stat.maxW; // Update maximum weight
1619: PrintStatus(stat. evaluations, N, local.tictoc, 10.0);
1620:
1621: if ( (stat.sigma_err / stat.sigma) < mparam.PRECISION &&
1622: stat. evaluations > mparam.MIN_EVENTS ) {
1623: break;
1624: }
1625: }
1626: // Progressbar
1627: if (atime.ElapsedSec() > 0.1) {
1628: gra:aux:PrintProgress(stat. evaluations / static_cast<double>(mparam.MIN_EVENTS));
1629: atime.Reset();
1630: }
1631:
1632: // Event generation mode
1633: if (GMODE == 1) {
1634: SaveEvent(proc, W, stat.maxW, aux);
1635: PrintStatus(stat.generated, N, local.tictoc, 10.0);
1636: if (stat.generated >= N) { break; }
1637: }
1638: // Progressbar
1639: if (atime.ElapsedSec() > 0.1) {
1640: gra:aux:PrintProgress(stat.generated / static_cast<double>(N));
1641: atime.Reset();
1642: }
1643:
1644: }
1645:
1646: PrintStatus(stat.generated, N, local.tictoc, -1.0);
1647: PrintStatistics(N);
1648: }
1649:
1650: // Save unweighted or weighted event
1651: int MGranitti::SaveEvent(MProc sproc, double weight, double MAXHEIGHT,
1652: const gra:AuxInData aux) {
1653: gra:mutex.lock();
1654: stat.totals += 1; // This is one trial more
1655:
1656: // 0. Hit-Miss
1657: bool hit_in = proc->random.U(0, 1) < (weight / MAXHEIGHT);
1658:
1659: // 2. We see weight larger than maxweight
1660: if (1WEIGHTED && weight > MAXHEIGHT) {

```



```
./src/MDirac.cc 1/11
1: // "Brute Force" Dirac Gamma Algebra for amplitudes
2: //
3: // Next step, optimize this class using more efficient spinor-products.
4: // For those, see e.g.:
5: //
6: // www.phys.nthu.edu.tw/~class/helicity/VY9A_Note_01.pdf
7: // www.physics.wyome.edu/~abechman/main/Research_files/shm.pdf
8: // R. Kleiss, W.G. Stirling, Spinor Techniques, Nucl.Phys. B262 (1985) 235-262
9: // F. Richardson, arxiv.org/abs/hep-ph/0110108
10: //
11: //
12: // (c) 2017-2020 Mikael Mieskolainen
13: // Licensed under the MIT license <http://opensource.org/licenses/MIT>.
14: //
15: // C++
16: #include <complex>
17: #include <iostream>
18: #include <vector>
19:
20: // Tensor algebra
21: #include "FTensor.hpp"
22:
23: // Own
24: #include "GraniTLL/M4Vec.h"
25: #include "GraniTLL/M4u.h"
26: #include "GraniTLL/MDirac.h"
27: #include "GraniTLL/MKinematics.h"
28: #include "GraniTLL/MMath.h"
29: #include "GraniTLL/MSpin.h"
30:
31: using FTensor::Tensor;
32: using FTensor::Tensor2;
33:
34: using gra::aux::indices;
35: using gra::math::magqr;
36: using gra::math::pow2;
37: using gra::math::zi;
38:
39: namespace gra {
40: MDirac::MDirac() { InitGammaMatrices("DIRAC"); }
41: MDirac::MDirac(const std::string &basis) { InitGammaMatrices(basis); }
42:
43: void MDirac::InitGammaMatrices(const std::string &basis) {
44:     // Use gamma basis
45:     if (basis == "DIRAC") {
46:         BASIS = "D";
47:     } else if (basis == "CHIRAL") {
48:         BASIS = "C";
49:     } else {
50:         throw std::invalid_argument("MDirac::InitGammaMatrices: Unknown gamma basis chosen: " + basis);
51:     }
52: }
53:
54: // PARITY OPERATOR = \gamma^0
55: // Intrinsic parity of fermions +1, anti-fermion -1
56:
57: const MMatrix<std::complex, double> > y0_chiral {
58:     std::vector<std::complex, double>{1.0, 0.0, 0.0, 1.0, 0.0},
59:     std::vector<std::complex, double>{0.0, 0.0, 0.0, 1.0},
60:     std::vector<std::complex, double>{1.0, 0.0, 0.0, 0.0},
61:     std::vector<std::complex, double>{0.0, 0.0, -1.0, 0.0},
62: };
63:
64: const MMatrix<std::complex, double> > y0_dirac {
65:     std::vector<std::complex, double>{1.0, 0.0, 0.0, 0.0},
66:     std::vector<std::complex, double>{0.0, 1.0, 0.0, 0.0},
67:     std::vector<std::complex, double>{0.0, 0.0, 1.0, 0.0},
68:     std::vector<std::complex, double>{0.0, 0.0, 0.0, -1.0},
69: };
70:
71: // \equiv i \gamma^{\mu\nu} \gamma^{\mu\nu} \gamma^{\mu\nu} \gamma^{\mu\nu}
72: const MMatrix<std::complex, double> > y5_chiral {
73:     std::vector<std::complex, double>{-1.0, 0.0, 0.0, 0.0},
74:     std::vector<std::complex, double>{0.0, -1.0, 0.0, 0.0},
75:     std::vector<std::complex, double>{0.0, 0.0, 1.0, 0.0},
76:     std::vector<std::complex, double>{0.0, 0.0, 0.0, 1.0},
77: };
78:
79: // \equiv i \gamma^{\mu\nu} \gamma^{\mu\nu} \gamma^{\mu\nu} \gamma^{\mu\nu}
80: const MMatrix<std::complex, double> > y5_dirac {
81:     std::vector<std::complex, double>{0.0, 0.0, 1.0, 0.0},
82:     std::vector<std::complex, double>{0.0, 0.0, 0.0, 1.0},
83:     std::vector<std::complex, double>{1.0, 0.0, 0.0, 0.0},
84:     std::vector<std::complex, double>{0.0, 1.0, 0.0, 0.0},
85: };
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```

./src/MDirac.cc          9/11
665:
666: const std::vector<int> lambda = {-1, 0, 1};
667:
668: for (const auto &m : indices(lambda)) { // loop over helicities
669:     eps[m] = MDirac::EpsMassiveSpin(p, lambda[m]);
670: }
671: if (type == "conj") { // Take complex conjugate per element
672:     for (const auto &mu : LI) { eps[m](mu) = std::conj(eps[m](mu)); }
673: }
674: if (!INDEX_UP) { // Return covariant (lower index) version
675:     for (const auto &mu : {1, 2, 3}) { eps[m](mu) = -eps[m](mu); }
676: }
677: return eps;
678: }
679: }
680:
681: //-----
682: // Test functions
683:
684: // Test gamma matrix anticommutation relation:
685: // (\gamma^\mu \gamma^\nu + \gamma^\nu \gamma^\mu) = 2g^{\mu\nu} I_4
686: //
687: double MDirac::TestGammaAntiCommutation() const {
688:     double diffsum = 0.0;
689:
690:     for (const auto &mu : LI) {
691:         for (const auto &nu : LI) {
692:             const MMatrix<std::complex<double>> AC_lo =
693:                 gamma_lo[mu] * gamma_lo[nu] + gamma_lo[nu] * gamma_lo[mu];
694:             std::cout << "gamma_lo: mu=" << mu << " nu=" << nu << std::endl;
695:             AC_lo.Print();
696:         }
697:     }
698:     const MMatrix<std::complex<double>> AC_up =
699:         gamma_up[mu] * gamma_up[nu] + gamma_up[nu] * gamma_up[mu];
700:     std::cout << "gamma_up: mu=" << mu << " nu=" << nu << std::endl;
701:     AC_up.Print();
702: }
703: }
704: }
705: return diffsum;
706: }
707: }
708: // Test slash(p)slash(p) = p^2 I_4 (identity matrix being I4)
709: //
710: double MDirac::TestSlashSlash(const MVec &p) const {
711:     const MMatrix<std::complex<double>> A = FSlash(p) * FSlash(p);
712:     const MMatrix<std::complex<double>> B = I4 * p.M(1);
713:     const double norm = (A - B).Frobnorm();
714:     std::cout << "MDirac::TestSlashSlash: Frobenius norm |A-B|=" << norm << std::endl;
715:     return norm;
716: }
717: }
718:
719: // Check Dirac u- and v-spinor normalization and completeness relations
720: //
721: // 1. Normalization of the Lorentz invariant inner product:
722: //
723: // \bar{u}_s(p) u_s(p) = 2m \delta_{s's'}
724: // \bar{v}_s(p) v_s(p) = -2m \delta_{s's'}
725: // \bar{u}_s(p) v_s(p) = 0
726: // \bar{v}_s(p) u_s(p) = 0
727: //
728: // Note that \bar{v}_s(p) u_s(p) = 2E \delta_{s's'}
729: // (standard relativistic normalization of a wavefunction)
730: //
731: // 2. Completeness relation (Projection operators):
732: //
733: // \sum_{s'} u_{s'}(p) \bar{u}_{s'}(p) = \slashed{p} + m
734: // \sum_{s'} v_{s'}(p) \bar{v}_{s'}(p) = \slashed{p} - m
735: //
736: //
737: // Set type = "u" or "v"
738: //
739: // Normalization test (add to unit tests!)
740: //
741: //
742: // SPINOR    CHIRAL    DIRAC
743: //-----
744: // uDirac    -         OK
745: // vDirac    -         OK
746: // uSubDirac -         OK
747: // vSubDirac -         OK

```

```

./src/MDirac.cc          11/11
831: //
832: // \sum_{\lambda} (-1, 0, 1) eps^\mu(k, \lambda) eps^\nu(k, \lambda) = \eta^{\mu\nu} k / \lambda
833: // = \eta^{\mu\nu} k / \lambda + k^\mu k^\nu / M^2
834: //
835: double MDirac::TestMassiveSpinComplete(const MVec &k) const {
836:     double absdiffsum = 0.0;
837:
838:     for (const auto &mu : LI) {
839:         for (const auto &nu : LI) {
840:             // Massive spin-1 helicities
841:             std::complex<double> sum = 0.;
842:             for (const int &lambda : {-1, 0, 1}) {
843:                 const Tensor<std::complex<double>> &eps = EpsMassiveSpin(k, lambda);
844:                 sum += eps(mu) * std::conj(eps(nu));
845:             }
846:             const double lhs = std::real(sum);
847:
848:             // Right hand side
849:             const double rhs = -g[mu][nu] + k[mu] * k[nu] / k.M(1);
850:             print("MDirac::TestMassiveSpinComplete: lhs = 0.2E, rhs = 0.2E \n", lhs, rhs);
851:             absdiffsum += std::abs(lhs - rhs);
852:         }
853:     }
854: }
855: return absdiffsum;
856: }
857: }
858: }
859: // namespace gra
860: }

```

```

./src/MDirac.cc          10/11
748: // uHel    | OK    | -
749: // vHel    | OK    | -
750: //
751: // mode = "helicity" (default), "gauge", "spin"
752: //
753: double MDirac::TestSpinComplete(const MVec &p, const std::string &type,
754:                                 const std::string &mode) const {
755:     std::cout << "MDirac::TestSpinComplete: Type=" << type << std::endl;
756:     // InitGammaMatrices(basis);
757:     MMatrix<std::complex<double>> lhs(4, 4, 0.0); // Init with zero!
758:     const double SIGN = (type == "u") ? 1.0 : -1.0;
759:
760:     for (const auto &lambda : SPINORSTATE) {
761:         std::vector<std::complex<double>> spinor;
762:
763:         if (type == "u") {
764:             if (BASIS == "D") {
765:                 spinor = uHelDirac(p, lambda);
766:             }
767:             if (mode == "spin") { spinor = uDirac(p, lambda); }
768:         }
769:         if (BASIS == "C") { spinor = uHelChiral(p, lambda); }
770:         if (mode == "gauge") { spinor = uGauge(p, lambda); }
771:     }
772:     else if (type == "v") {
773:         if (BASIS == "D") {
774:             spinor = vHelDirac(p, lambda);
775:         }
776:         if (mode == "spin") { spinor = vDirac(p, lambda); }
777:     }
778:     if (BASIS == "C") { spinor = vHelChiral(p, lambda); }
779:     if (mode == "gauge") { spinor = vGauge(p, lambda); }
780: }
781: }
782: }
783: }
784: }
785: // Adjoint
786: const std::vector<std::complex<double>> spinorbar = Bar(spinor);
787: //
788: // Check normalization
789: const double nrhs =
790:     std::real(gra::matoper::VecVecMultiply(spinorbar, spinor)); // Take real to cast to double
791: const double nrhs = SIGN * 2. * p.M();
792: print("s = %d: Normalization = ", lambda);
793: if (gra::aux::AssertRatio(nrhs, nrhs, 3E-3)) {
794:     std::cout << "rang::fg:reset << nrhs << " OK" << " rang::fg:reset
795:     << std::endl;
796: }
797: }
798: }
799: }
800: }
801: // Take outerproduct, sum
802: lhs = gra::matoper::OuterProd(spinor, spinorbar);
803: }
804: }
805: std::cout << std::endl;
806: //
807: // Completeness relation
808: const MMatrix<std::complex<double>> rhs = FSlash(p) + I4 * p.M() * SIGN;
809: //
810: // Compare
811: lhs.Print("sum_s |sum_s(p) \bar{spinor}_s(p)|");
812: std::cout << std::endl;
813: rhs.Print("slash(p) + i*m*SIGN");
814: std::cout << std::endl;
815: //
816: const double frobnorm = (lhs - rhs).Frobnorm();
817: if (frobnorm < 0.1) {
818:     std::cout << "rang::fg:reset << "Completeness relation OK, Frobenius norm = " << frobnorm
819:     << " rang::fg:reset << std::endl;
820: }
821: }
822: }
823: }
824: std::cout << std::endl;
825: //
826: return frobnorm;
827: }
828: }
829: }
830: // Check massive spin-1 polarization sum (completeness relation)

```

```

./src/MFactorized.cc     1/6
1: // Factorized type phase space class
2: //
3: // (c) 2019-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5: //
6: // C++
7: #include <algorithm>
8: #include <complex>
9: #include <iostream>
10: #include <random>
11: #include <vector>
12: //
13: // Dep
14: #include "Granitti/MDux.h"
15: #include "Granitti/MFactorized.h"
16: #include "Granitti/MForm.h"
17: #include "Granitti/MFragment.h"
18: #include "Granitti/MKinematics.h"
19: #include "Granitti/MMatch.h"
20: #include "Granitti/MMatrix.h"
21: #include "Granitti/MProcess.h"
22: #include "Granitti/MSpin.h"
23: #include "Granitti/MUserCuts.h"
24: //
25: // Libraries
26: #include "rang.hpp"
27: //
28: using gra::aux::indices;
29: //
30: using gra::math::CheckEMC;
31: using gra::math::FPI;
32: using gra::math::abs;
33: using gra::math::msqrt;
34: using gra::math::pow;
35: using gra::math::pow2;
36: using gra::math::pow4;
37: using gra::math::pow8;
38: using gra::math::r4;
39: //
40: using gra::PDG::GeV2barn;
41: //
42: namespace gra {
43: //
44: // This is needed by construction
45: MFactorized::MFactorized() { Initialize(); }
46: //
47: // Constructor
48: MFactorized::MFactorized(std::string process, const std::vector<aux::OneCMD &syntax) {
49:     Initialize();
50:     InitHistograms();
51:     SetProcess(process, syntax);
52: }
53: //
54: // Init final states
55: MVec zerovec(0, 0, 0, 0);
56: for (std::size_t i = 0; i < 10; ++i) { lhs.pfinal.push_back(zerovec); }
57: std::cout << "MFactorized: (Constructor done) << std::endl;
58: //
59: void MFactorized::Initialize() {
60:     const std::vector<std::string> supported = {"pp", "p\bar{p}", "pp", "pp", "pp"};
61:     CID = MSUBPROC(supported, CID);
62:     ProcPtr = MSUBPROC(supported, CID);
63: }
64: //
65: // Destructor
66: MFactorized::~MFactorized() {}
67: //
68: // Initialize cut and process specific poststep
69: void MFactorized::post_Constructor() {
70:     // Set sampling boundaries
71:     SetTechnicalBoundaries(gouts, EXCITATION);
72: }
73: //
74: // Initialize phase space dimension
75: ProcPtr::LIPSDIM = 5 + 1; // All processes, +1 from central system mass
76: //
77: if (EXCITATION == 1) ProcPtr::LIPSDIM = 1; }
78: if (EXCITATION == 2) ProcPtr::LIPSDIM = 2; }
79: //
80: // Update kinematics (screening Kt loop calls this)
81: // Exact 4-momentum conservation at loop vertices, that is, by using this one
82: // does not assume varying external momenta in the screening loop calculation.
83: //

```



```

416: exact = 0;
417: // Massive exact closed form for 2-body case
418: if (lts.decaytree.size() == 2) {
419:   exact = gra::kinematics::FS2Massive(lts.m2, pow2(lts.decaytree[0].p4.M()),
420:   pow2(lts.decaytree[1].p4.M()));
421: }
422: // Massless case
423: if (lts.decaytree.size() > 2) {
424:   exact = gra::kinematics::FSMassless(lts.m2, lts.decaytree.size());
425: }
426: }
427:
428: // 5-Dim Integral Volume [int_{M^2_MIN}^{M^2_MAX} { dM^2 } [pb1] x [pb2] x
429: // [pt1] x [pt2] x [y]
430: // Integral over central mass^2 is separate [phase-space factorization], but
431: // encapsulated here (in dimension)
432: double MFactorized::B5IntegralVolume() const {
433: // Forward leg integration
434: const double forwardVolume = ForwardVolume();
435:
436: return (pow2(M_MAX) - pow2(M_MIN)) * (gcuts.Y_max - gcuts.Y_min) * forwardVolume;
437: }
438:
439: // 5-Dim phase space weight
440: double MFactorized::B5PhaseSpaceWeight() const {
441: const double J =
442: 1.0 / std::abs(lts.pfinal[1].Pz() / lts.pfinal[1].E() -
443: lts.pfinal[2].Pz() / lts.pfinal[2].E()); // Jacobian, close to 0.5
444:
445: const double factor = (1.0 / 2.0) * (1.0 / pow2(2.0 * gra::math::PI)) *
446: (lts.pfinal[1].PE() / (2.0 * lts.pfinal[1].E())) *
447: (lts.pfinal[2].PE() / (2.0 * lts.pfinal[2].E())) * J;
448:
449: return factor;
450: }
451:
452: // For high mass limit kinematics, see e.g. [arxiv.org/pdf/hep-ph/9903279.pdf]
453:
454: } // namespace gra
455:

```

```

1: // "Metaclass" for Spherical Harmonics Expansion - contains
2: // dataprocessing/Minuit/ROOT plotting functions.
3:
4: // All mathematical spherical harmonic processing is separated to MSpherical namespace.
5:
6:
7: // (c) 2017-2020 Mikael Hieska
8: // Licensed under the MIT license <http://opensource.org/licenses/MIT>.
9:
10: // C++
11: #include <complex>
12: #include <iostream>
13: #include <random>
14: #include <regex>
15: #include <string>
16: #include <vector>
17:
18: // ROOT
19: #include "Canvas.h"
20: #include "Color.h"
21: #include "TGraphErrors.h"
22: #include "TH.h"
23: #include "TH2.h"
24: #include "TList.h"
25: #include "TLegend.h"
26: #include "TLine.h"
27: #include "TMatrixVector.h"
28: #include "TMath.h"
29: #include "TMathLog.h"
30: #include "TProfile.h"
31: #include "TStyle.h"
32:
33: // Own
34: #include "Gravitl/Analysis/MHarmonic.h"
35: #include "Gravitl/Analysis/MROOT.h"
36: #include "Gravitl/MKVec.h"
37: #include "Gravitl/MBox.h"
38: #include "Gravitl/MKinematics.h"
39: #include "Gravitl/MMesh.h"
40: #include "Gravitl/MROD.h"
41: #include "Gravitl/MSpherical.h"
42:
43: // Eigen
44: #include <Eigen/Dense>
45:
46: using gra::aux::indices;
47: using gra::math::PI;
48: using gra::math::INTEGR;
49: using gra::math::pow2;
50: using gra::math::z;
51:
52: namespace gra {
53: // Constructor
54: MHarmonic::MHarmonic() {}
55:
56: // Initialize
57: void MHarmonic::Init(const HPARAM shp) {
58:   param = shp;
59:   NCOEF = (param.LMAX + 1) * (param.LMAX + 1);
60:
61: // -----
62: // SETUP parameters of the fit
63: ACTIVE.resize(NCOEF); // Active moments
64: ACTIVEDEF = 0;
65:
66: for (int l = 0; l <= param.LMAX; ++l) {
67:   for (int m = -l; m <= l; ++m) {
68:     const int index = gra::spherical::lLinearInd(l, m);
69:     ACTIVE[index] = true;
70:
71: // FIX ODD MOMENTS TO ZERO
72: if (param.REMOVEODD && (l & 2) != 0) { ACTIVE[index] = false; }
73: // FIX NEGATIVE N TO ZERO
74: if (param.REMOVENEGATIVEM && (m < 0)) { ACTIVE[index] = false; }
75: if (ACTIVE[index]) { ++ACTIVEDEF; }
76: }
77: }
78:
79: // -----
80: std::cout << std::endl;
81: param.Print();
82: std::cout << std::endl;
83: // -----

```

```

84:
85: // Init arrays used by Minuit function
86: t_lm = std::vector<double>(NCOEF, 0.0);
87: t_lm_error = std::vector<double>(NCOEF, 0.0);
88: ermat = MMatrix<double>(NCOEF, NCOEF, 0.0);
89: covmat = MMatrix<double>(NCOEF, NCOEF, 0.0);
90:
91: // Test functions
92: gra::spherical::TestSphericalIntegrals(param.LMAX);
93:
94: std::cout << rangi::fyellow
95: << "Spherical Harmonic Based (cos(theta), phi) r, F."
96: << "Decomposition and Efficiency Inversions"
97: << std::endl;
98: << std::endl;
99: std::cout << "REMOVED ODD" << rangi::freset << std::endl;
100: std::cout << " (G) Generated == Events in the angular flat phase space (no "
101: "cuts on final "
102: "states, only on the system)"
103: << std::endl;
104: std::cout << " (F) Fiducial == Events within the strict fiducial "
105: " (geometric-kinematic)"
106: "Final state phase space (cuts on final states)"
107: << std::endl;
108: std::cout << " (D) Detector == Events after the detector efficiency losses, selection AND "
109: "fiducial cuts"
110: << std::endl;
111: std::cout << " (E) subset of (F) subset of (G) (this strict hierarchy might "
112: "be violated in "
113: "some special cases)"
114: << std::endl;
115: std::cout << std::endl;
116: std::cout << " The basic idea is to define (G) such that minimal extrapolation " << std::endl;
117: std::cout << " is required from the strict fiducial (geometric) phase space (F). " << std::endl;
118: std::cout << " Flatness requirement of (G) is strictly required to represent "
119: "moments in "
120: "an unimbed basis (non-flat phase space <=> geometric moment mixing)."
121: << std::endl;
122: std::cout << std::endl;
123: std::cout << rangi::fyellow << "EXAMPLE OF A FORMALLY VALID DEFINITION: " << rangi::freset
124: << std::endl;
125: std::cout << " G = { |y(system)| < 0.9 } " << std::endl;
126: std::cout << " F = { |eta(p1)| * 0.9 && pt(p1) > 0.1 GeV } " << std::endl;
127: std::cout << " D = { |eta(p1)| * 0.9 && pt(p1) > 0.1 GeV } " << std::endl;
128: std::cout << " Note also the rotation between inactive coefficients and "
129: "active one, due to "
130: "moment mixing "
131: << std::endl;
132: << std::endl;
133: std::cout << std::endl << std::endl;
134:
135: // pause (?)
136: }
137:
138: // Destructor
139: MHarmonic::~MHarmonic() {}
140:
141: bool MHarmonic::PrintLog(const std::string outpath) const {
142: // Add here code to print output to files ...
143:
144: return true;
145: }
146:
147: // Plot all figures
148: void MHarmonic::PlotAll(const std::string outpath) const {
149: for (const auto OBSERVABLE :
150: // *** EFFICIENCY DECOMPOSITION ***
151: PlotDEfficiency(OBSERVABLE, outpath);
152: // *** ALGEBRAIC INVERSE MOMENTS ***
153: PlotFigures(det, OBSERVABLE, "h_Moments[MPP]<det>", 33, outpath);
154: PlotFigures(fid, OBSERVABLE, "h_Moments[MPP]<fid>", 33, outpath);
155: PlotFigures(fia, OBSERVABLE, "h_Moments[MPP]<fia>", 33, outpath);
156: PlotFigures(fia, OBSERVABLE, "h_Moments[MPP]<fia>", 33, outpath);
157: // *** EXTENDED MAXIMUM LIKELIHOOD INVERSE MOMENTS ***
158: if (param.EML) {
159: PlotFigures(det, OBSERVABLE, "h_Moments[EM]<det>", 33, outpath);
160: PlotFigures(fid, OBSERVABLE, "h_Moments[EM]<fid>", 33, outpath);
161: PlotFigures(fia, OBSERVABLE, "h_Moments[EM]<fia>", 33, outpath);
162: }
163: }
164: }
165:
166: // 2D

```

```

167: std::vector<std::vector<int>> OBSERVABLES = {{0, 1}, {0, 2}, {1, 2}}; // Different pairs
168:
169: for (const auto OBSERVABLE : OBSERVABLES) {
170: // *** EFFICIENCY DECOMPOSITION ***
171: PlotFigures2D(fia, OBSERVABLE, "(Response)[FLAT_REFERENCE]<fia>", 17, outpath);
172: PlotFigures2D(fid, OBSERVABLE, "(Response)[FIDUCIAL_ACCEPTANCE]<fid>", 33, outpath);
173: PlotFigures2D(det, OBSERVABLE, "(Response)[ACCEPTANCE_EFFICIENCY]<det>", 29, outpath);
174:
175: // *** ALGEBRAIC INVERSE MOMENTS ***
176: PlotFigures2D(fia, OBSERVABLE, "(Moments)[MPP]<fia>", 17, outpath);
177: PlotFigures2D(fid, OBSERVABLE, "(Moments)[MPP]<fid>", 33, outpath);
178: PlotFigures2D(det, OBSERVABLE, "(Moments)[MPP]<det>", 29, outpath);
179:
180: // *** EXTENDED MAXIMUM LIKELIHOOD INVERSE MOMENTS ***
181: if (param.EML) {
182: PlotFigures2D(fia, OBSERVABLE, "(Moments)[EM]<fia>", 17, outpath);
183: PlotFigures2D(fid, OBSERVABLE, "(Moments)[EM]<fid>", 33, outpath);
184: PlotFigures2D(det, OBSERVABLE, "(Moments)[EM]<det>", 29, outpath);
185: }
186: }
187:
188: const int OBSERVABLE = 0;
189: Plot2DExpansion(fid, OBSERVABLE, "h_Moments[MPP]<fid>", 33, outpath);
190: Plot2DExpansion(fia, OBSERVABLE, "h_Moments[MPP]<fia>", 33, outpath);
191: if (param.EML) {
192: Plot2DExpansion(fid, OBSERVABLE, "h_Moments[EM]<fid>", 33, outpath);
193: Plot2DExpansion(fia, OBSERVABLE, "h_Moments[EM]<fia>", 33, outpath);
194: }
195: }
196:
197: // Synthesized (cos(theta), phi) plots
198: //
199: void MHarmonic::Plot2DExpansion(
200: const std::map<gra::spherical::Meta, MTensor<gra::spherical::SD>> &tensor,
201: unsigned int OBSERVABLE, const std::string TYPESTRING, int barcolor,
202: const std::string outpath) const {
203: // -----
204: // Extract name strings
205:
206: // Find (string)
207: std::smatch sma;
208: std::regex_search(TYPESTRING, sma, std::regex(R"(\.(.*)\{)")); // R"()" for raw string literals
209: std::string DATATYPE = sma[0];
210: DATATYPE = DATATYPE.substr(1, DATATYPE.size() - 2);
211:
212: // Find <string>
213: std::smatch smb;
214: std::regex_search(TYPESTRING, smb, std::regex(R"(\.(.*)\{)")); // R"()" for raw string literals
215: std::string SPACE = smb[0];
216: SPACE = SPACE.substr(1, SPACE.size() - 2);
217:
218: // Find (string)
219: std::smatch smc;
220: std::regex_search(TYPESTRING, smc, std::regex(R"(\.(.*)\{)")); // R"()" for raw string literals
221: std::string ALGO = smc[0];
222: ALGO = ALGO.substr(1, ALGO.size() - 2);
223:
224: // -----
225: // z(cos(theta), phi) M synthesis
226:
227: std::size_t N = 50;
228:
229: // Cos(theta) and phi
230: std::vector<double> costheta;
231: std::vector<double> phi;
232:
233: std::size_t BINS = 0;
234:
235: // TH2 for each
236: std::vector<std::vector<TH2D *>> h2;
237:
238: // Loop over fia
239: std::size_t source_ind = 0;
240: for (const auto source : tensor) {
241: // legends[find] = source.FIRST.LEGEND;
242: BINS = source.second.size(OBSERVABLE);
243:
244: // Add histograms
245: h2.push_back(std::vector<TH2D *>(BINS, NULL));
246:
247: // Loop over observable
248:
249: for (std::size_t bin = 0; bin < BINS; ++bin) {

```



```

250: // Get x-axis point
251: // const double value = grid(OBSERVABLE)[bin].center();
252:
253: // Set indices {0,0,0, ..., 0}
254: std::vector<int> size_t_cell(grid.size(), 0);
255: cell(OBSERVABLE) = bin;
256:
257: // Synthesize distribution
258: MMatrix<double> Z;
259:
260: if (ALGO == "MPP") {
261:     Z = spherical::Y_real_synthesize(source.second(cell).t_lm_MPP, ACTIVE, N, costheta, phi, true);
262: }
263: } else if (ALGO == "EML") {
264:     Z = spherical::Y_real_synthesize(source.second(cell).t_lm_EML, ACTIVE, N, costheta, phi, true);
265: }
266:
267: // Create histogram
268: const double EPS = 1e-3;
269: const std::string name = "h2_" + std::to_string(source_ind) + "_" + std::to_string(bin);
270: h2[source_ind][bin] = new TH2D(name.c_str(), "c1", 200 + theta, phi, (rad), N, -(1 + EPS), 1 + EPS, N, -(math::PI + EPS), math::PI + EPS);
271:
272: for (std::size_t i = 0; i < N; ++i) {
273:     for (std::size_t j = 0; j < N; ++j) {
274:         h2[source_ind][bin] += fillObservable(i, phi[i], Z[i][j]);
275:     }
276: }
277:
278: }
279:
280: ++source_ind;
281: // Loop over sources
282:
283: // Loop over bins
284: for (std::size_t bin = 0; bin < BINS; ++bin) {
285:     TCanvas* c1 = new TCanvas("c1", "c1", 200 + tensor.size() * 400, 500); // horizontal, vertical
286:     c1->Divide(tensor.size(), 1, 0.002, 0.001);
287:
288:     std::size_t source_ind = 0;
289:     std::string FRAME;
290:     for (const auto &source : tensor) {
291:         c1->cd(source_ind + 1);
292:         c1->cd(source_ind + 1)->SetRightMargin(0.13);
293:         //
294:         FRAME = source.first.FRAME;
295:         h2[source_ind][bin]->SetTitle(source.first.LEGEND.c_str());
296:         h2[source_ind][bin]->GetXaxis()->CenterTitle();
297:         h2[source_ind][bin]->GetYaxis()->SetRangeLower(0.1, 1.0);
298:         h2[source_ind][bin]->Draw("COLZ");
299:         //
300:         // Draw Lorentz Frame string on left
301:         TText* t1 = new TText("0.45, 0.85, Form(\"s = {0.2E, {0.2E} %s\", FRAME.c_str(), grid(OBSERVABLE)[bin].min, grid(OBSERVABLE)[bin].max, xlabel(OBSERVABLE).c_str());");
302:         t1->SetNDC();
303:         t1->SetTextAlign(22);
304:         t1->SetTextColor(8);
305:         t1->SetTextFont(43);
306:         t1->SetTextSize(18);
307:         t1->SetTextStyle(45);
308:         t1->Draw("same");
309:         //
310:         ++source_ind;
311:     }
312: }
313:
314: // CreateDirectory("figs");
315: // CreateDirectory("figs/harmonicfit");
316: // CreateDirectory("figs/harmonicfit/" + outputpath + "/synthesis");
317:
318: const std::string subpath = "SPACER_" + std::to_string(OBSERVABLE) + "_" + FRAME + "_" + ALGO;
319: const std::string fullpath = "figs/harmonicfit/" + outputpath + "/synthesis" + subpath;
320: // CreateDirectory(fullpath);
321: // Print Form("s/1001a.pdf", fullpath.c_str(), bin);
322: delete c1;
323:
324: }

```

```

417: std::string yaxis_label = "Events / bin";
418: for (const auto &source : tensor) {
419:     legendstrs[ind] = source.first.LEGEND;
420:     BINS = source.second.size(OBSERVABLE);
421:
422:     // Loop over moments
423:     int k = 0;
424:
425:     double SCALE = source.first.SCALE;
426:
427:     if (SCALE < 0 && source.first.Y_AXIS == "") { yaxis_label = "Normalized to 1"; }
428:     if (source.first.Y_AXIS != "") { yaxis_label = source.first.Y_AXIS; }
429:
430:     for (int l = 0; l <= param.LMAX; ++l) {
431:         for (int m = -l; m <= l; ++m) {
432:             const int index = gra::spherical::LinearInd(l, m);
433:             if (!ACTIVE(index)) { continue; } // Not active
434:
435:             // Set canvas position
436:             c1->cd(k + 1);
437:
438:             // Loop over observable
439:             double x[BINS] = {0.0};
440:             double y[BINS] = {0.0};
441:             double w_err[BINS] = {0.0};
442:             double y_err[BINS] = {0.0};
443:
444:             for (std::size_t bin = 0; bin < BINS; ++bin) {
445:                 // Get x-axis point
446:                 x[bin] = grid(OBSERVABLE)[bin].center();
447:
448:                 // Set indices {0,0,0, ..., 0}
449:                 std::vector<int> size_t_cell(grid.size(), 0);
450:                 cell(OBSERVABLE) = bin;
451:
452:                 // CHOOSE DATAMODE
453:                 if (ALGO == "MPP") {
454:                     y[bin] = source.second(cell).t_lm_MPP[index];
455:                     y_err[bin] = source.second(cell).t_lm_MPP_error[index];
456:                 } else if (ALGO == "EML") {
457:                     y[bin] = source.second(cell).t_lm_EML[index];
458:                     y_err[bin] = source.second(cell).t_lm_EML_error[index];
459:                 } else {
460:                     throw std::invalid_argument("MHarmmonic::PlotFigures: Unknown input: " + DATAMODE + " = ALGO = " + ALGO);
461:                 }
462:
463:                 // Scaling (e.g. luminosity)
464:                 // Normalize lm = 0 to sum cp 1 -> then apply to all
465:                 if (l == 0 && m == 0 && SCALE < 0.0) {
466:                     double sum = 0.0;
467:                     for (std::size_t bin = 0; bin < BINS; ++bin) { sum += y[bin]; }
468:                     if (sum > 0) { SCALE = 1.0 / sum; }
469:                 }
470:
471:                 // Apply scale by user
472:                 for (std::size_t bin = 0; bin < BINS; ++bin) {
473:                     y[bin] *= SCALE;
474:                     y_err[bin] *= SCALE;
475:                 }
476:
477:                 // Save maximum for visualization
478:                 for (std::size_t bin = 0; bin < BINS; ++bin) {
479:                     MAXVAL[k] = y[bin] > MAXVAL[k] ? y[bin] : MAXVAL[k];
480:                     MINVAL[k] = y[bin] < MINVAL[k] ? y[bin] : MINVAL[k];
481:                 }
482:
483:                 // Data displayed using TGraphErrors
484:                 g[ind][k] = new TGraphErrors(BINS, x, y, w_err, y_err);
485:
486:                 // Colors
487:                 g[ind][k]->SetMarkerColor(colors[ind]);
488:                 g[ind][k]->SetLineColor(colors[ind]);
489:                 g[ind][k]->SetFillColor(colors[ind]);
490:                 g[ind][k]->SetLineWidth(2.0);
491:                 g[ind][k]->SetMarkerStyle(23); // square
492:                 g[ind][k]->SetMarkerSize(0.3);

```

```

333: // Delete all histograms
334: for (std::size_t i = 0; i < h2.size(); ++i) {
335:     for (std::size_t j = 0; j < h2[i].size(); ++j) { delete h2[i][j]; }
336: }
337:
338: void GetLegendPosition(unsigned int N, double &x1, double &x2, double &y1, double &y2, const std::string legendposition) {
339:     // North-East
340:     x1 = 0.63;
341:     x2 = x1 + 0.18;
342:     y1 = 0.75 - 0.01 * N;
343:
344:     // South-East
345:     if (legendposition.compare("southeast") == 0) { y1 = 0.10 - 0.01 * N; }
346:     y2 = y1 + 0.05 * N; // Scale by the number of histograms
347: }
348:
349: // LM-Expansion plots as a function of observables
350: void MHarmmonic::PlotFigures(
351:     const std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>& tensor,
352:     unsigned int OBSERVABLE, const std::string &TYPESTRING, int barcolor,
353:     const std::string &outputpath) const {
354:     //
355:     const std::map<gra::spherical::Meta, MTensor<gra::spherical::SH>& tensor,
356:     unsigned int OBSERVABLE, const std::string &TYPESTRING, int barcolor,
357:     const std::string &outputpath) const {
358:         if (OBSERVABLE > xlabel.size() - 1) {
359:             throw std::invalid_argument("MHarmmonic::PlotFigures: Unknown observable " +
360:                 std::to_string(OBSERVABLE));
361:         }
362:         const std::string xlabel = xlabel(OBSERVABLE);
363:         // Extract name strings
364:         //
365:         // Find string
366:         std::string search, sma, std::regex("^(\\.[^\\.]*)"); // R "/" for raw string literals
367:         std::string DATAMODE = sma[0];
368:         DATAMODE = DATAMODE.substr(1, DATAMODE.size() - 2);
369:         // Find <string>
370:         std::string smb;
371:         std::string search, smc, std::regex("^(\\.[^\\.]*)"); // R "/" for raw string literals
372:         std::string SPACE = SPACE.substr(1, SPACE.size() - 2);
373:         // Find <string>
374:         std::string smc;
375:         std::string search, smc, std::regex("^(\\.[^\\.]*)"); // R "/" for raw string literals
376:         std::string ALGO = smc[0];
377:         ALGO = ALGO.substr(1, ALGO.size() - 2);
378:         //
379:         // Loop over data sources
380:         std::size_t BINS = 0;
381:         int ind = 0;
382:         // Graphs for each data source & [lm-moment]
383:         std::vector<std::vector<TGraphErrors >> g(tensor.size(), std::vector<TGraphErrors >>(ACTIVEDEF, NULL));
384:         // Legend titles
385:         std::vector<std::string> legendstrs(tensor.size());
386:         // Minimum and maximum y-values for each plot
387:         std::vector<double> MINVAL(ACTIVEDEF, 1e32);
388:         std::vector<double> MAXVAL(ACTIVEDEF, -1e32);
389:         // Turn of horizontal errors
390:         gStyle->SetErrorX(0);
391:         //
392:         std::vector<TMultiGraph > mg(ACTIVEDEF, NULL);
393:         for (std::size_t k = 0; k < ACTIVEDEF; ++k) { mg[k] = new TMultiGraph(); }
394:         //
395:         std::string FRAME;
396:         std::vector<std::string> TITLES;
397:         // Y-axis title
398:         FRAME = source.first.FRAME;
399:         TITLES = source.first.TITLES;
400:         // Add to the multigraph
401:         mg[k]->Add(g[ind][k]);
402:         ++k;
403:         // over m
404:         // over l
405:         // Loop over sources
406:         // Draw multigraph
407:         unsigned int k = 0;
408:         // Aux variables
409:         std::shared_ptr<TPad> tpad;
410:         std::shared_ptr<TText> t1;
411:         std::shared_ptr<TText> t2;
412:         for (int l = 0; l <= param.LMAX; ++l) {
413:             for (int m = -l; m <= l; ++m) {
414:                 const int index = gra::spherical::LinearInd(l, m);
415:                 if (!ACTIVE(index)) { continue; } // Not active
416:
417:                 // Set canvas position
418:                 c1->cd(k + 1);
419:                 // First draw, then setup (otherwise crash)
420:                 mg[k]->Draw("AC");
421:                 // Title and y-axis
422:                 if (k == 0) {
423:                     // Title
424:                     if (SPACE == "det") {
425:                         mg[k]->SetTitle(Form("s {l}m {m} @<{l}m>", TITLES[0].c_str(), l, m));
426:                     }
427:                     if (SPACE == "fid") {
428:                         mg[k]->SetTitle(Form("s {l}m {m} @<{l}m>", TITLES[1].c_str(), l, m));
429:                     }
430:                     if (SPACE == "fia") {
431:                         mg[k]->SetTitle(Form("s {l}m {m} @<{l}m>", TITLES[2].c_str(), l, m));
432:                     }
433:                 }
434:                 // Y-axis (for some ROOT reason, this needs to be always after SetTitle)
435:                 mg[k]->GetYaxis()->SetTitle(yaxis_label.c_str());
436:                 gPad->SetLeftMargin(0.15); // 15 per cent of pad for left margin, default is 10%
437:                 mg[k]->GetYaxis()->SetTitleOffset(1.25);
438:             }
439:             //
440:             // Set x-axis
441:             mg[k]->GetXaxis()->SetTitle(xlabel.c_str());
442:             gStyle->SetBarWidth(0.5);
443:             // mg[k]->SetFillColor(0);
444:             mg[k]->GetXaxis()->SetTitleSize(0.05);
445:             mg[k]->GetYaxis()->SetTitleSize(0.05);
446:             mg[k]->GetYaxis()->SetTitleSize(0.05);
447:             // Set y-axis
448:             mg[k]->GetYaxis()->SetTitle("Intensity");
449:             gStyle->SetTitleFontSize(0.08);
450:             //
451:             if (k == 0) {
452:                 gStyle->SetTitleW(0.95); // width percentage
453:             }
454:             // Y-axis range
455:             if (k == 0) {
456:                 mg[k]->GetHistogram()->SetMaximum(MAXVAL[k] * 1.1);
457:                 mg[k]->GetHistogram()->SetMinimum(0.0);
458:             } else {
459:                 // Skip the first element (lm=0)

```

```

./src/Analysis/MHarmomic.cc      8/19
582:     const double maxval = "std::max_element(std::begin(MAXVAL) + 1, std::end(MAXVAL));
583:     const double minval = "std::min_element(std::begin(MINVAL) + 1, std::end(MINVAL));
584:     const double bound = std::max(std::abs(minval), std::abs(maxval));
585:
586:     mg[k]->GetHistogram()->SetMaximum(bound * 1.1);
587:     mg[k]->GetHistogram()->SetMinimum(-bound * 1.1);
588: }
589:
590: //
591: // Draw Lorentz FRAME string on left
592: TText *t2 = new TText(0.225, 0.825, FRAME_c_str());
593: t2->SetTextAlign(22);
594: t2->SetTextAlign(22);
595: t2->SetTextColor(kRed + 2);
596: t2->SetTextFont(43);
597: t2->SetTextSize(std::ceil(1.0 / sqrt(ACTIVEVDF)) * 16);
598: // t2->SetTextAngle(45);
599: t2->Draw("same");
600: //
601: //
602: // Draw horizontal line
603: if (k != 0) {
604:     TLine *line = new TLine(grid(OBSERVABLE)[0].min, 0.0,
605:                             grid(OBSERVABLE)[grid(OBSERVABLE).size() - 1].max, 0.0);
606:     line->SetLineColor(kBlack);
607:     line->SetLineWidth(1.0);
608:     line->Draw("same");
609: }
610: //
611: //
612: // Who made it
613: if (k == ACTIVEVDF - 1) {
614:     // New pad on top of all
615:     c1->cd(); // Important!
616:     gStyle->SetLineStyle(1);
617:     gStyle->SetLineStyle(1);
618:     gStyle->SetLineStyle(1);
619: }
620: const double xpos = 0.99;
621: gStyle->SetLineStyle(1);
622: //
623: //
624: //
625: //
626: //
627: //
628: // Draw legend to the upper left most
629: c1->cd(1);
630: //
631: //
632: //
633: // Create legend
634: double x1, x2, y1, y2 = 0.0;
635: const std::string legendposition = "northeast";
636: GetLegendPosition(2, x1, x2, y1, y2, legendposition);
637: TLegend *legend = new TLegend(x1, x2, y1, y2);
638: legend->SetFillColor(0); // White background
639: legend->SetBorderSize(0); // No box
640: legend->SetTextSize(0.03);
641: //
642: // Add legend entries
643: for (const auto &i : indices(gpr)) { legend->AddEntry(gpr[i][0], legendstrs[i].c_str()); }
644: //
645: // Draw legend
646: legend->Draw("same");
647: //
648: //
649: // Require that we have data
650: if (BINS > 1) {
651:     aux::CreateDirectory("./figs");
652:     aux::CreateDirectory("./figs/harmonicfit");
653:     aux::CreateDirectory("./figs/harmonicfit/" + outputpath);
654:     const std::string subpath = "OBS_" + std::to_string(OBSERVABLE) + "_" + FRAME;
655:     const std::string fullpath = "./figs/harmonicfit/" + outputpath + "/" + subpath;
656:     aux::CreateDirectory(fullpath);
657:     c1->Print(Form("%s/%s.pdf", fullpath.c_str(), TYPETESTING.c_str()));
658: }
659: // Merge pdfs using Ghostscript (gs)
660: const std::string cmd = "gs -q -dNOCACHE -dNOPAUSE -q -dSDFTC=pdfwrite -dOutputFile=" + fullpath +
661:     " -sDEVICE=pdfwrite -sOutputFile=" + fullpath + ".h*.pdf";
662: if (system(cmd.c_str()) == -1) {
663:     throw std::invalid_argument("Error: Problem executing Ghostscript merge on pdfs!");
664: }

```

```

./src/Analysis/MHarmomic.cc      9/19
665: }
666: //
667: //
668: for (std::size_t i = 0; i < gpr.size(); ++i) {
669:     delete gpr[i];
670: }
671: //
672: //
673: //
674: // LN-Efficiency Figures
675: //
676: void MHarmomic::PlotINEfficiency(unsigned int OBSERVABLE, const std::string outputpath) const {
677: //
678: std::shared_ptr<TCanvas> c1;
679: gStyle->SetLineStyle(1);
680: //
681: //
682: if (OBSERVABLE > labels.size() - 1) {
683:     throw std::invalid_argument("MHarmomic::PlotFigures: Unknown observable " +
684:                                 std::to_string(OBSERVABLE));
685: }
686: const std::string xlabel = labels(OBSERVABLE);
687: //
688: // Graphs for each efficiency [level] x [l-moment]
689: std::vector<std::vector<TGraphErrors>> gr3, std::vector<TGraphErrors>> (ACTIVEVDF, NULL);
690: //
691: // Legend titles
692: std::vector<std::string> legendstrs(3);
693: //
694: // Minimum and maximum y-values for each plot
695: std::vector<double> MINVAL(ACTIVEVDF, 1e32);
696: std::vector<double> MAXVAL(ACTIVEVDF, -1e32);
697: //
698: // Turn off horizontal errors
699: gStyle->SetErrorX(0);
700: //
701: std::vector<MultiGraph> mg(ACTIVEVDF, NULL);
702: for (std::size_t k = 0; k < ACTIVEVDF; ++k) { mg[k] = new TMultiGraph(); }
703: //
704: //
705: //
706: //
707: //
708: //
709: //
710: //
711: //
712: //
713: //
714: //
715: //
716: for (std::size_t level = 0; level < 3; ++level) {
717:     // Loop over moments
718:     int k = 0;
719:     for (int l = 0; l < param.LMAX; ++l) {
720:         for (int m = -1; m <= l; ++m) {
721:             const int index = gra::spherical::linearrnd(l, m);
722:             if (!ACTIVE[index] || !continue; ) // Not active
723:                 continue;
724:             // Set canvas position
725:             c1->cd(k + 1);
726:             //
727:             // Loop over mass
728:             double x[BINS] = {0.0};
729:             double y[BINS] = {0.0};
730:             double x_err[BINS] = {0.0};
731:             double y_err[BINS] = {0.0};
732:             //
733:             for (std::size_t bin = 0; bin < BINS; ++bin) {
734:                 // Set x-axis point
735:                 x[bin] = grid(OBSERVABLE)[bin].center();
736:                 //
737:                 //
738:                 //
739:                 //
740:                 //
741:                 //
742:                 //
743:                 //
744:                 //
745:                 //
746:                 //
747:                 //
748:                 //
749:                 //
750:                 //
751:                 //
752:                 //
753:                 //
754:                 //
755:                 //
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```

```

./src/Analysis/MHarmomic.cc      10/19
748:     y_err[bin] = fid_DET(cell).E_lm_error[index];
749:     } else if (level == 2) {
750:         y[bin] = fia_DET(cell).E_lm[index];
751:         y_err[bin] = fia_DET(cell).E_lm_error[index];
752:     }
753: }
754: // Save maximum for visualization
755: MAXVAL[k] = y[bin] > MAXVAL[k] ? y[bin] : MAXVAL[k];
756: MINVAL[k] = y[bin] < MINVAL[k] ? y[bin] : MINVAL[k];
757: //
758: //
759: // Data displayed using TGraphErrors
760: gr[level][k] = new TGraphErrors(BINS, x, y, x_err, y_err);
761: //
762: // Colors
763: gr[level][k]->SetMarkerColor(colors[level]);
764: gr[level][k]->SetLineColor(colors[level]);
765: gr[level][k]->SetFillColor(colors[level]);
766: gr[level][k]->SetLineStyle(1);
767: gr[level][k]->SetMarkerStyle(21); // square
768: gr[level][k]->SetMarkerSize(0.3);
769: //
770: // Add to the multigraph
771: mg[k]->Add(gr[level][k]);
772: //
773: //
774: //
775: //
776: //
777: //
778: //
779: //
780: //
781: unsigned int k = 0;
782: //
783: std::shared_ptr<TPad> tpad;
784: std::shared_ptr<TLabel> l1;
785: std::shared_ptr<TLabel> l2;
786: //
787: for (int l = 0; l <= param.LMAX; ++l) {
788:     for (int m = -1; m <= l; ++m) {
789:         const int index = gra::spherical::linearrnd(l, m);
790:         if (!ACTIVE[index] || !continue; ) // Not active
791:             continue;
792:         // Set canvas position
793:         c1->cd(k + 1);
794:         //
795:         // First draw, then setup (otherwise crash)
796:         mg[k]->Draw("ac");
797:         //
798:         //
799:         //
800:         //
801:         //
802:         //
803:         //
804:         //
805:         //
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```

./src/Analysis/MHarmomic.cc      11/19
831:     const double maxval = "std::max_element(std::begin(MAXVAL) + 1, std::end(MAXVAL));
832:     const double minval = "std::min_element(std::begin(MINVAL) + 1, std::end(MINVAL));
833:     const double bound = std::max(std::abs(minval), std::abs(maxval));
834:
835:     mg[k]->GetHistogram()->SetMaximum(bound * 1.1);
836:     mg[k]->GetHistogram()->SetMinimum(-bound * 1.1);
837: }
838:
839: // Draw Lorentz FRAME string on left
840: TText *t2 = new TText(0.225, 0.825, FRAME_c_str());
841: t2->SetAlign(22);
842: t2->SetTextAlign(22);
843: t2->SetTextColor(kRed + 2);
844: t2->SetTextFont(43);
845: t2->SetTextSize(std::ceil(1.0 / sqrt(ACTIVEVDF)) * 16);
846: // t2->SetTextAngle(45);
847: t2->Draw("same");
848: // Draw horizontal line
849: if (k != 0) {
850:     TLine *line = new TLine(grid(OBSERVABLE)[0].min, 0.0,
851:                             grid(OBSERVABLE)[grid(OBSERVABLE).size() - 1].max, 0.0);
852:     line->SetLineColor(kBlack);
853:     line->SetLineWidth(1.0);
854:     line->Draw("same");
855: }
856: //
857: //
858: // Who made it
859: if (k == ACTIVEVDF - 1) {
860:     // New pad on top of all
861:     c1->cd(); // Important!
862:     gStyle->SetLineStyle(1);
863:     gStyle->SetLineStyle(1);
864:     gStyle->SetLineStyle(1);
865: }
866: const double xpos = 0.99;
867: gStyle->SetLineStyle(1);
868: //
869: //
870: //
871: //
872: //
873: //
874: //
875: //
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912: //
913: //

```

```

./src/Analysis/MHarmomic.cc          12/19
914: void MHarmomic::PlotFigures2D(
915: const std::map<gras:spherical:Meta, MTensor>:gras:spherical:SSH> &tensor,
916: const std::vector<gras:SO3SEWV3E2E2> &const std::string sTYPSTRNG, int barcolor,
917: const std::string &outpath) const {
918: /* Implement here ... */
919: }
920:
921: // Loop over system mass, pt, rapidity
922:
923: void MHarmomic::HyperLoop(void (*fifunc)(int &, double *, double *, double *, int),
924: const double M_STEP = (hp.PT[2] - hp.PT[1]) / hp.PT[0],
925: const std::vector<gras:spherical:Omega> &MC,
926: // Initialize detector expansion tensor
927: fild_DET = MTensor>:gras:spherical:SH_DET =
928: (unsigned int)hp.M[0], (unsigned int)hp.PT[0], (unsigned int)hp.Y[0]);
929: fild_DET = fild_DET;
930: det_DET = fild_DET;
931:
932: // -----
933: // Create grid discretization
934:
935: grid.resize(3);
936: grid[0].resize(hp.M[0]);
937: grid[1].resize(hp.PT[0]);
938: grid[2].resize(hp.Y[0]);
939:
940: // Mass steps
941: const double M_STEP = (hp.M[2] - hp.M[1]) / hp.M[0];
942: const double PT_STEP = (hp.PT[2] - hp.PT[1]) / hp.PT[0];
943: const double Y_STEP = (hp.Y[2] - hp.Y[1]) / hp.Y[0];
944:
945: for (std::size_t i = 0; i < hp.M[0]; ++i)
946:   grid[0][i].min = 1 * M_STEP + hp.M[1];
947:   grid[0][i].max = grid[0][i].min + M_STEP;
948:
949: for (std::size_t j = 0; j < hp.PT[0]; ++j)
950:   grid[1][j].min = 1 * PT_STEP + hp.PT[1];
951:   grid[1][j].max = grid[1][j].min + PT_STEP;
952:
953: for (std::size_t k = 0; k < hp.Y[0]; ++k)
954:   grid[2][k].min = 1 * Y_STEP + hp.Y[1];
955:   grid[2][k].max = grid[2][k].min + Y_STEP;
956:
957: // -----
958: // Expand the detector transfer function
959:
960: for (const auto &i : indices(grid[0])) {
961:   for (const auto &j : indices(grid[1])) {
962:     for (const auto &k : indices(grid[2])) {
963:       const std::vector<std::size_t> MC_ind = gras:spherical:GetIndices(
964:         MC, {grid[0][i].min, grid[0][i].max}, {grid[1][j].min, grid[1][j].max},
965:         {grid[2][k].min, grid[2][k].max});
966:
967: // Acceptance mixing matrices
968: fild_DET((i, j, k), k).MIXM = gras:spherical:GetMixing(MC, MC_ind, param.LMAX, "fid");
969: fild_DET((i, j, k), k).MIXM = gras:spherical:GetMixing(MC, MC_ind, param.LMAX, "fid");
970: det_DET((i, j, k), k).MIXM = gras:spherical:GetMixing(MC, MC_ind, param.LMAX, "det");
971:
972: // -----
973: // det efficiency decomposition for this interval
974: std::pair<std::vector<double>, std::vector<double>> EO =
975: gras:spherical:GeELEM(MC, MC_ind, param.LMAX, "fid");
976: std::pair<std::vector<double>, std::vector<double>> EI =
977: gras:spherical:GeELEM(MC, MC_ind, param.LMAX, "fid");
978: std::pair<std::vector<double>, std::vector<double>> E2 =
979: gras:spherical:GeELEM(MC, MC_ind, param.LMAX, "det");
980:
981: fild_DET((i, j, k), k).E1m = EO.first;
982: fild_DET((i, j, k), k).E1m_err = EO.second;
983:
984: fild_DET((i, j, k), k).E2m = EI.first;
985: fild_DET((i, j, k), k).E2m_err = EI.second;
986:
987: det_DET((i, j, k), k).E1m = E2.first;
988: det_DET((i, j, k), k).E1m_err = E2.second;
989:
990: }
991: }
992:
993: // -----
994: // Loop over data sources
995: for (const auto &ind : indices(DATA)) {

```

```

./src/Analysis/MHarmomic.cc          13/19
997: //
998: // Pre-Calculate once Spherical Harmonics for the MINUIT fit
999: DATA_events = DATA[ind].EVENTS;
1000: Y_min = gras:spherical:YLM(DATA_events, param.LMAX);
1001: // -----
1002:
1003: double ch2 = 0.0;
1004:
1005: // Data source identifier string
1006: const gras:spherical:Meta META = DATA[ind].META;
1007:
1008: // Initialize tensor arrays
1009: MTensor>:gras:spherical:SSH = temp;
1010: (unsigned int)hp.M[0], (unsigned int)hp.PT[0], (unsigned int)hp.Y[0]);
1011: fild(META) = temp;
1012: fild(META) = temp;
1013: det(META) = temp;
1014:
1015: // Expand Data
1016: for (const auto &i : indices(grid[0])) {
1017:   for (const auto &j : indices(grid[1])) {
1018:     for (const auto &k : indices(grid[2])) {
1019:       // Data indices
1020:       DATA_ind = gras:spherical:GetIndices(DATA_events, {grid[0][i].min, grid[0][i].max},
1021:         {grid[1][j].min, grid[1][j].max},
1022:         {grid[2][k].min, grid[2][k].max});
1023:
1024:       const unsigned int MINEVENTS = 75;
1025:       std::cout << "rangif:read << WARNING: Less than " << MINEVENTS << " in the call!"
1026:         << rangif:resd << std::endl;
1027:     }
1028:   }
1029: }
1030: // =====
1031: // ALGORITHM 1: DIRECT / OBSERV / ALGEBRAIC decomposition
1032:
1033: fild(META)((i, j, k), k).t_LMPP =
1034: gras:spherical:SphericalMoments(DATA_events, DATA_ind, param.LMAX, "fid");
1035: det(META)((i, j, k), k).t_LMPP =
1036: gras:spherical:SphericalMoments(DATA_events, DATA_ind, param.LMAX, "det");
1037:
1038: // Forward matrix
1039: Eigen::MatrixX M = gras:aux:Matrix2Eigen(det_DET((i, j, k), k).MIXM);
1040:
1041: // Matrix pseudo-inverse
1042: Eigen::VectorX b = gras:aux:Vector2Eigen(det_META)((i, j, k), k).t_LMPP;
1043: Eigen::VectorX x = gras:aux:MatrixInverse(M, param.SVOREG) * b;
1044:
1045: // Collect the inversion result
1046: fild(META)((i, j, k), k).t_LMPP_err = gras:aux:EigenVector(x);
1047:
1048: // =====
1049: // Update these to proper errors (TBD. FOR NAIVE POISSON
1050: // COUNTING FOR NOW)
1051: fild(META)((i, j, k), k).t_LMPP_error =
1052: std::vector<double>(&fild(META)((i, j, k), k).t_LMPP_err.size(), 0.0);
1053: for (const auto &e : indices(fild(META)((i, j, k), k).t_LMPP_err)) {
1054:   fild(META)((i, j, k), k).t_LMPP_error[e] =
1055:     msqrt(std::abs(fild(META)((i, j, k), k).t_LMPP_err[e]));
1056: }
1057:
1058: // Propagate errors assuming no error on mixing matrix
1059: // (infinite reference MC statistics limit)
1060:
1061: fild(META)((i, j, k), k).t_LMPP_error =
1062: gras:spherical:ErrorProp(fild_DET((i, j, k), k).MIXM, fild(META)((i, j, k), k).t_LMPP_error);
1063: det(META)((i, j, k), k).t_LMPP_error =
1064: gras:spherical:ErrorProp(det_DET((i, j, k), k).MIXM, fild(META)((i, j, k), k).t_LMPP_error);
1065:
1066: // (unmixed) moments in the (angular flat) -
1067: // "reference phase space."
1068:
1069: std::cout << "Algebraic Moore-Penrose/SVD inverted "
1070:   << std::endl;
1071: gras:spherical:PrintOutMoments(fild(META)((i, j, k), k).t_LMPP,
1072:   fild(META)((i, j, k), k).t_LMPP_error, ACTIVE, param.LMAX);
1073:
1074: std::cout << "Algebraic (mixed) moments in the fiducial "
1075:   << "phase space."
1076:   << std::endl;
1077: gras:spherical:PrintOutMoments(fild(META)((i, j, k), k).t_LMPP,
1078:   fild(META)((i, j, k), k).t_LMPP_error, ACTIVE, param.LMAX);
1079:
1080: std::cout << "Algebraic (mixed) moments in the detector space." << std::endl;

```

```

./src/Analysis/MHarmomic.cc          14/19
1080: gras:spherical:PrintOutMoments(det(META)((i, j, k), k).t_LMPP,
1081:   det(META)((i, j, k), k).t_LMPP_error, ACTIVE, param.LMAX);
1082:
1083: // =====
1084: // ALGORITHM 2: Extended Maximum Likelihood Fit
1085:
1086: if (param.EMG) {
1087:   // ** det EMG-based fit decomposition **
1088:   MomentFit(META, (i, j, k), fitfunc);
1089:
1090:   // Save result
1091:   fild(META)((i, j, k), k).t_LMML = t_Lm;
1092:   fild(META)((i, j, k), k).t_LMML_err = EM;
1093:   for (const auto &ind : indices(fild(META)((i, j, k), k).t_LMML_err)) {
1094:     // Vector
1095:     det(META)((i, j, k), k).t_LMML_err = EM;
1096:     // Moment forward rotation: Matrix *
1097:     det_DET((i, j, k), k).MIXM * t_Lm;
1098:   }
1099:   // Vectors
1100:   fild(META)((i, j, k), k).t_LMML_error = t_Lm_error;
1101:
1102: // Propagate errors assuming no error on mixing
1103: // matrix (infinite reference MC statistics limit)
1104: fild(META)((i, j, k), k).t_LMML_error =
1105: gras:spherical:ErrorProp(fild_DET((i, j, k), k).MIXM, t_Lm_error);
1106: det(META)((i, j, k), k).t_LMML_error =
1107: gras:spherical:ErrorProp(det_DET((i, j, k), k).MIXM, t_Lm_error);
1108:
1109: // =====
1110: std::cout << "Extended Maximum-Likelihood inverse "
1111:   << "fitted (unmixed) moments in the "
1112:   << "(angular flat) phase space."
1113:   << std::endl;
1114: gras:spherical:PrintOutMoments(fild(META)((i, j, k), k).t_LMML,
1115:   fild(META)((i, j, k), k).t_LMML_err, ACTIVE,
1116:   param.LMAX);
1117:
1118: std::cout << "Extended Maximum-Likelihood -
1119:   "Back-Projected (mixed) moments in -
1120:   "the fiducial phase space."
1121:   << std::endl;
1122: gras:spherical:PrintOutMoments(fild(META)((i, j, k), k).t_LMML,
1123:   fild(META)((i, j, k), k).t_LMML_err, ACTIVE,
1124:   param.LMAX);
1125:
1126: std::cout << "Extended Maximum-Likelihood -
1127:   "Back-Projected (mixed) moments in -
1128:   "the detector space."
1129:   << std::endl;
1130: gras:spherical:PrintOutMoments(det(META)((i, j, k), k).t_LMML,
1131:   det(META)((i, j, k), k).t_LMML_err, ACTIVE,
1132:   param.LMAX);
1133:
1134: }
1135:
1136: // Print results
1137: ch2 += PrintOutHyperCall(META, (i, j, k));
1138:
1139: // -----
1140: // Make comparison of synthetic MC data vs estimate
1141:
1142: if (DATA[ind].META.MODE == "MC") {
1143:   int fiducial = 0;
1144:   int selected = 0;
1145:   for (const auto &i : DATA[ind] {
1146:     if (DATA_events[i].fiducial) { ++fiducial; }
1147:     if (DATA_events[i].fiducial && DATA_events[i].selected) { ++selected; }
1148:   }
1149:   std::cout << std::endl;
1150:   std::cout << rangif:yellow << "MC GROUND TRUTH: " << rangif:reset << std::endl;
1151:   print(
1152:     " MC synthetic data events generated = "
1153:     " %u\n",
1154:     (unsigned int)DATA_ind.size());
1155:   print(
1156:     " MC synthetic data events fiducial = "
1157:     " %u (acceptance %0.1f percent)\n",
1158:     fiducial, fiducial / (double)DATA_ind.size() * 100);
1159:   print(
1160:     " MC synthetic data events fiducial and = "
1161:     " selected = %0.1f (efficiency %0.1f percent)\n",
1162:

```

```

./src/Analysis/MHarmomic.cc          15/19
1163:   selected, selected / (double)fiducial * 100);
1164:   std::cout << std::endl;
1165: }
1166:
1167: }
1168:
1169: if (param.EMG) {
1170:   gras:aux:PrintBar("");
1171:   double reducedch2 = ch2 / (double)ACTIVEINDF * hp.M[0] * hp.PT[0] * hp.Y[0];
1172:   if (reducedch2 < 3) {
1173:     std::cout << rangif:green;
1174:   } else {
1175:     std::cout << rangif:red;
1176:   }
1177:   print(
1178:     "Total ch2 (NPF - EM) / (ACTIVEINDF * BIN) = %0.2E / (%d * %d) = "
1179:     "%0.2E\n",
1180:     ch2, ACTIVEINDF, (int)hp.M[0] * hp.PT[0], hp.Y[0], reducedch2);
1181:
1182: std::cout << rangif:reset;
1183: gras:aux:PrintBar("");
1184: std::cout << std::endl << std::endl;
1185: }
1186:
1187: // Loop over data sources
1188:
1189: // Print out results for the hypercell
1190:
1191: double MHarmomic::PrintOutHyperCell(const gras:spherical:Meta & META,
1192:   const std::vector<std::size_t> &cell) {
1193:   double ch2 = 0;
1194:
1195:   // Print information
1196:   META.Print();
1197:
1198:   // Extended Maximum Likelihood
1199:   if (param.EMG == true) {
1200:     // Loop over moments
1201:     for (int l = 0; l <= param.LMAX; ++l) {
1202:       for (int m = -l; m <= l; ++m) {
1203:         const int index = gras:spherical:LinearInd(l, m);
1204:         const double obs = det(META)(cell, t_LMML[index]);
1205:         const double fit = det(META)(cell, t_LMML[index]);
1206:
1207:         // Moment active
1208:         if (ACTIVE(index) && (ch2 += pow2(obs - fit) / pow2(obs);))
1209:           ;
1210:       }
1211:     }
1212:     std::cout << std::endl;
1213:
1214:     const double reducedch2 = ch2 / (double)ACTIVEINDF;
1215:     if (reducedch2 < 3) {
1216:       std::cout << rangif:green;
1217:     } else {
1218:       std::cout << rangif:red;
1219:     }
1220:     print(
1221:       "ch2 (NPF - EM) / (NDF * %0.3E / %d * %0.3E) = %0.2E, ch2, ACTIVEINDF, reducedch2);
1222:     std::cout << rangif:reset << std::endl;
1223:
1224:     const double sum_fla = fild(META)(cell, t_LMML[gras:spherical:LinearInd(0, 0)]);
1225:     print(
1226:       "EM: Estimate of events in this hyperbin in the flat phase space = "
1227:       "%0.1E +/- %0.1E\n",
1228:       sum_fla, msqrt(sum_fla)); // Poisson error
1229:
1230:     const double sum_FID = gras:spherical:HarrodProd(fild_DET(cell), E_Lm, fild(META)(cell), t_LMML,
1231:       ACTIVE, param.LMAX);
1232:     print(
1233:       "EM: Estimate of events in this hyperbin in the fiducial phase "
1234:       "space = %0.1E\n",
1235:       sum_FID, msqrt(sum_FID)); // Poisson error
1236:
1237:     const double sum_DET = gras:spherical:HarrodProd(det_DET(cell), E_Lm, fild(META)(cell), t_LMML,
1238:       ACTIVE, param.LMAX);
1239:     print(
1240:       "EM: Estimate of events in this hyperbin in the detector "
1241:       "space = %0.1E\n",
1242:

```

```

./src/Analysis/MHarmmonic.cc      16/19
1246:   sum_DET, msqrt(sum_DET)); // Poisson error
1247: }
1248: std::cout << std::endl;
1249:
1250: // Algebraic Inverse
1251: const double sum_fla = fla[META](cell).t_lm_MPP(gra:spherical:LinearInd(0, 0));
1252: printf(
1253: "MPP: Estimate of events in this hyperbin in the flat phase space = %0.1f\n",
1254: sum_fla, msqrt(sum_fla)); // Poisson error
1255:
1256:
1257: const double sum_FID =
1258: gra:spherical:HarDotProd(fid_DET[cell]).E_lm, fla[META](cell).t_lm_MPP, ACTIVE, param.LMAX);
1259: printf(
1260: "MPP: Estimate of events in this hyperbin in the fiducial phase =
1261: "space = %0.1f\n",
1262: sum_FID, msqrt(sum_FID)); // Poisson error
1263:
1264:
1265: const double sum_DET =
1266: gra:spherical:HarDotProd(det_DET[cell]).E_lm, fla[META](cell).t_lm_MPP, ACTIVE, param.LMAX);
1267: printf(
1268: "MPP: Estimate of events in this hyperbin in the detector
1269: "space = %0.1f\n",
1270: sum_DET, msqrt(sum_DET)); // Poisson error
1271:
1272:
1273: return chi2;
1274: }
1275:
1276: // MINUIT based fit routine for the Extended Maximum Likelihood formalism
1277: //
1278: void MHarmmonic::MomentFit(const gra:spherical:Meta META, const std::vector<std::size_t> &cell,
1279: std::cout << "MomentFit: Starting ..." << std::endl << std::endl;
1280: // **** This must be set for the loss function ****
1281:
1282: activecell = cell;
1283:
1284: // Init TMinuit
1285: TMinuit *gMinuit = new TMinuit(NCOEFF); // initialize TMinuit with a maximum of N params
1286: gMinuit->SetFCN(fitfunc);
1287:
1288: // Set Print Level
1289: // -1 no output
1290: // 1 standard output
1291: gMinuit->SetPrintLevel(-1);
1292:
1293: // Set error Definition
1294: // 1 for Chi square
1295: // 0.5 for negative log likelihood
1296: gMinuit->SetErrorDef(0.5);
1297:
1298: double arglist[1];
1299: int iarglist = 0;
1300: arglist[0] = 0.5; // 0.5 <= We use negative log likelihood cost function
1301: gMinuit->mexccom("SET ERR", arglist, 1, iarglist);
1302:
1303: // double fminbest = 1e32;
1304: // Try different initial values to find out true minimum
1305: const std::size_t TRIALMAX = 1;
1306:
1307: for (std::size_t trials = 0; trials < TRIALMAX; ++trials) {
1308:   for (int l = 0; l <= param.LMAX; ++l)
1309:     for (int m = -1; m <= l; ++m) {
1310:       const int index = gra:spherical:LinearInd(l, m);
1311:
1312:       const std::string str = "t_" + std::ito_string(l) + std::ito_string(m);
1313:
1314:       const double max = 1e9; // Physically bounded ultimately by the number of events
1315:       const double min = -1e5;
1316:
1317:       // *****
1318:       // *** Use the algebraic inverse solution as a good starting value
1319:       // *****
1320:       const double start_value = fla[META](activecell).t_lm_MPP[index];
1321:       // *****
1322:       const double step_value = 0.1; // in units of Events
1323:
1324:       gMinuit->mnparm(index, str, start_value, step_value, min, max, ierflag);
1325:
1326:       // After first trial, fix t_00 (error are not estimated with
1327:       // constant parameters)
1328:

```

```

./src/Analysis/MHarmmonic.cc      18/19
1412:   std::vector<double> T(NCOEFF, 0.0);
1413:
1414:   for (const auto &i : indices(T)) {
1415:     T[i] = par[i];
1416:     // printf("T[%d] = %0.5f\n", i);
1417:   }
1418:   // This is the number of events in the current phase space point at detector
1419:   // Level
1420:   const int METHOD = 2;
1421:   double nhat = 0.0;
1422:
1423:   // Equivalent estimator 1
1424:   if (METHOD == 1) {
1425:     const std::vector<double> t_lm_det = det_DET(activecell).MIXlm * T; // Matrix * Vector
1426:     nhat = t_lm_det[0];
1427:   }
1428:   // Equivalent estimator 2
1429:   if (METHOD == 2) {
1430:     nhat = gra:spherical:HarDotProd(det_DET(activecell)).E_lm, T, ACTIVE, param.LMAX);
1431:   }
1432:
1433:   // For each event, calculate |sum_lm| t_lm
1434:   // Re[Y_lm](cos(theta), phi_k), k is the event index
1435:   std::vector<double> Y;
1436:   const double V = msqrt(4.0 * P1); // Normalization volume
1437:
1438:   for (const auto &k : DATA_ind) {
1439:     // Event is accepted
1440:     if (DATA_events[k].fiducial && DATA_events[k].selected) {
1441:       // fine
1442:     } else {
1443:       continue;
1444:     }
1445:
1446:     // Loop over (l,m) terms
1447:     double sum = 0.0;
1448:     for (int l = 0; l <= param.LMAX; ++l) {
1449:       for (int m = -1; m <= l; ++m) {
1450:         const int index = gra:spherical:LinearInd(l, m);
1451:
1452:         if (ACTIVE[index]) {
1453:           // Calculate here
1454:           // const std::complex<double> Y =
1455:           // gra:math:Y_complex_basis(DATA_events[k].costheta, DATA_events[k].phi,
1456:           // l, m);
1457:           // const double ReY = gra:math:RMReY(Y, l, m);
1458:
1459:           // Pre-calculated for speed
1460:           const double ReY = Y_lm[k][index];
1461:
1462:           // Add
1463:           sum += T[index] * ReY;
1464:         }
1465:       }
1466:     }
1467:     10.push_back(V * sum);
1468:   }
1469:   // Fidelity term
1470:   double logL = 0;
1471:
1472:   for (const auto &k : indices(T)) {
1473:     if (T[k] > 0) {
1474:       logL += std::log(T[k]);
1475:       // Positivity is not satisfied, make strong penalty
1476:     } else {
1477:       logL -= 1e32;
1478:       // The Re[Y_lm] (purely real) formulation here does
1479:       // not, explicitly, enforce non-negativity.
1480:       // Where as |h|^2 formulation would enforce it by
1481:       // construction,
1482:       // but that would have -1 <= m <= l parameters, i.e.,
1483:       // over redundant representation problem
1484:       // on the other hand.
1485:     }
1486:   }
1487:
1488:   // Note the minus sign on logL term
1489:   f = -logL + nhat;
1490:
1491:   // High penalty, total event count estimate has gone out of physical
1492:   if (nhat < 0) { f = 1e32; }
1493:
1494:   // L1-norm regularization (Laplace prior), -> + ln(P_Laplace)

```

```

./src/Analysis/MHarmmonic.cc      17/19
1329: // if (trials > 0) {
1330: //   gMinuit->mnparm(0, "t_00", t_lm[0], 0, 0, 0, ierflag);
1331: //   gMinuit->FixParameter(0);
1332: // }
1333:
1334: // FIX ONE MOMENTS TO ZERO
1335: if (param.REMOVEZEROS && ((l+1) > 1) &= 0) {
1336:   gMinuit->mnparm(index, str, 0, 0, 0, 0, ierflag);
1337:   gMinuit->FixParameter(index);
1338: }
1339:
1340: // FIX NEGATIVE N TO ZERO
1341: if (param.REMOVENEGATIVES && (m < 0)) {
1342:   gMinuit->mnparm(index, str, 0, 0, 0, 0, ierflag);
1343:   gMinuit->FixParameter(index);
1344: }
1345:
1346: // Scan main parameter
1347: arglist[0] = 10000; // Minimum number of function calls
1348: // Minimum tolerance
1349:
1350: // First simplex to find approximate answers
1351: gMinuit->mexccom("SIMPLEX", arglist, 2, ierflag);
1352:
1353: // Numerical Hessian (2nd derivatives matrix), inverse of this -> covariance
1354: // matrix
1355: gMinuit->mexccom("MIGRAD", arglist, 2, ierflag);
1356:
1357: // Confidence intervals based on the profile likelihood ratio
1358: gMinuit->mexccom("MINOS", arglist, 2, ierflag);
1359:
1360: // Calculate error covariance matrix
1361: gMinuit->memat(ermat[0][0], NCOEFF);
1362:
1363: for (int i = 0; i < NCOEFF; ++i) {
1364:   for (int j = 0; j < NCOEFF; ++j) {
1365:     covmat[i][j] = sqrt(ermat[i][i] * ermat[j][j]);
1366:     if (covmat[i][j] > 1E-80) {
1367:       covmat[i][j] = ermat[i][j] / covmat[i][j];
1368:     } else {
1369:       covmat[i][j] = 0.0;
1370:     }
1371:   }
1372:   ermat.Print("Error matrix");
1373:   covmat.Print("Covariance matrix");
1374:
1375:   // Print results
1376:   double fmin, fedm, errdef = 0.0;
1377:   int nvpars, nparx, istat = 0;
1378:   gMinuit->mstat(fmin, fedm, errdef, nvpars, nparx, istat);
1379:
1380:   // Collect fit result
1381:   for (int l = 0; l <= param.LMAX; ++l) {
1382:     for (int m = -1; m <= l; ++m) {
1383:       const int index = gra:spherical:LinearInd(l, m);
1384:
1385:       // Set results into t_lm_err[]
1386:       gMinuit->GetParameter(index, t_lm[index], t_lm_error[index]);
1387:     }
1388:   }
1389: } // TRIALS LOOP END
1390:
1391: // Sum of N random variables Y = X_1 + X_2 + ... X_N
1392: // sigma_Y^2 = sum_i^N sigma_i^2 + 2 * sum_i sum_j < j
1393: // covariance (X_i, X_j)
1394:
1395: // Print out results (see MINUIT manual for these parameters)
1396: double fmin, fedm, errdef = 0.0;
1397: int nvpars, nparx, istat = 0;
1398: gMinuit->mstat(fmin, fedm, errdef, nvpars, nparx, istat);
1399:
1400: std::cout << "MINUIT (MIGRAD+MINOS) done." << std::endl << std::endl;
1401:
1402: delete gMinuit;
1403:
1404:
1405:
1406: // Unbinned Extended Maximum Likelihood function
1407: // Extended means that the number of events (itself) is a Poisson distributed
1408: // random variable and that is incorporated to the fit.
1409: void MHarmmonic::logFunctionEpar, double *gin, double sr, double *par, int iflag) const {
1410:   // Collect fit t_lm coefficients from MINUIT

```

```

./src/Analysis/MHarmmonic.cc      19/19
1495: double lterm = 0.0;
1496: const unsigned int START = 1;
1497: for (std::size_t i = START; i < T.size(); ++i) { lterm += std::tanh(T[i]); }
1498: f += lterm * param.L1REG * nhat; // nhat for scale normalization
1499: // printf("MHarmmonic: cost-functional = %0.4f, nhat = %0.1f\n", f, nhat);
1500:
1501: } // namespace gra
1502:
1503: } // namespace gra
1504:

```

```

1: // Fast MC analysis class
2:
3: // (c) 2017-2020 Mikael Mieskolainen
4: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6: // C++
7: #include <algorithm>
8: #include <iostream>
9: #include <memory>
10: #include <vector>
11:
12: // C-file processing
13: #include <sys/atomic.h>
14: #include <sys/types.h>
15: #include <unistd.h>
16:
17: // ROOT
18: #include "TBranch.h"
19: #include "TCanvas.h"
20: #include "TColor.h"
21: #include "TFile.h"
22: #include "TFile.h"
23: #include "TFile.h"
24: #include "TFile.h"
25: #include "TFile.h"
26: #include "TFile.h"
27: #include "TFile.h"
28: #include "TFile.h"
29: #include "TFile.h"
30: #include "TFile.h"
31: #include "TFile.h"
32: #include "TFile.h"
33:
34: // HepMC3
35: #include "HepMC3/FourVector.h"
36: #include "HepMC3/GenCrossSection.h"
37: #include "HepMC3/GenEvent.h"
38: #include "HepMC3/GenParticle.h"
39: #include "HepMC3/GenVertex.h"
40: #include "HepMC3/ReaderAscii.h"
41: #include "HepMC3/ReaderAscii.h"
42: #include "HepMC3/Selector.h"
43: #include "HepMC3/Selector.h"
44:
45: // Own
46: #include "Granniti/Analysis/MNalyzer.h"
47: #include "Granniti/Analysis/ResultPlot.h"
48: #include "Granniti/M4Vec.h"
49: #include "Granniti/M4Vec.h"
50: #include "Granniti/M4Vec.h"
51: #include "Granniti/M4Vec.h"
52:
53: const bool DEBUG = false;
54:
55: using gra::aux::indices;
56: using gra::math::magr;
57:
58: namespace gra {
59:
60: // Constructor with unique ID string for ROOT bookkeeping reasons
61: MNalyzer(MNalyzer(const std::string& id) {
62: // Initialize histograms
63: const int NBINS = 150;
64:
65: // Energy
66: hE_Pions = std::make_shared<TH1D>(Form("hE_p", "Energy ppi (GeV)", ID.c_str()),
67: "Energy (GeV)", NBINS, 0, sqrts / 2.0);
68: hE_Gamma = std::make_shared<TH1D>(Form("hE_gamma", "Energy gamma (GeV)", ID.c_str()),
69: "Energy (GeV)", NBINS, 0, sqrts / 2.0);
70: hE_Neutron = std::make_shared<TH1D>(Form("hE_n", "Energy n (GeV)", ID.c_str()),
71: "Energy (GeV)", NBINS, 0, sqrts / 2.0);
72: hE_GammaNeutron = std::make_shared<TH1D>(Form("hE_gn", "Energy gn (GeV)", ID.c_str()),
73: "Energy (GeV)", NBINS, 0, sqrts / 2.0);
74:
75: // Feynman-x
76: hXF_Pions = std::make_shared<TH1D>(Form("hXF_p", "hXF p", ID.c_str()), "Feynman-x/Events",
77: NBINS, -1.0, 1.0);
78: hXF_Gamma = std::make_shared<TH1D>(Form("hXF_gamma", "hXF gamma", ID.c_str()), "Feynman-x/Events",
79: NBINS, -1.0, 1.0);
80: hXF_Neutron = std::make_shared<TH1D>(Form("hXF_n", "hXF n", ID.c_str()), "Feynman-x/Events",
81: NBINS, -1.0, 1.0);
82:
83: // Forward systems

```

```

167: // Read event from input file
168: HepMC3::GenEvent evt(HepMC3::Units::IGeV, HepMC3::Units::IMm);
169: input_file.read_event(evt);
170:
171: // Reading failed
172: if (input_file.failed()) {
173: if (events_read == 0) {
174: throw std::invalid_argument("MNalyzer::HepMC3Read: File " + totalpath + " is empty!");
175: } else {
176: break;
177: }
178: }
179: if (events_read == 0) {
180: HepMC3::Print::listing(evt);
181: HepMC3::Print::content(evt);
182: }
183: //event_read;
184:
185: // ** Get generator cross section (in picobarns by HepMC3 convention) **
186: std::shared_ptr<HepMC3::GenCrossSection> cs =
187: evt.attribute<HepMC3::GenCrossSection>("GenCrossSection");
188: if (cs) {
189: cross_section = 1E-12 * cs->xsec(0); // turn into barns
190: } else {
191: std::cout << "Problem accessing 'GenCrossSection' attribute" << std::endl;
192: }
193: // -----
194: // ** Get event weight (always in barn units) **
195: double W = 1.0;
196: if (evt.weights().size() != 0) // check do we have weights saved
197: W = evt.weights()[0]; // take the first one
198:
199: totalW += W;
200: // -----
201:
202: // Central particles
203: std::vector<M4Vec> pip;
204: std::vector<M4Vec> pim;
205:
206: for (HepMC3::ConstGenParticlePtr p1 :
207: HepMC3::applyFilter(HepMC3::StandardSelector::PDG_ID == finalPDG, evt.particles()) {
208: M4Vec pvec = gra::aux::HepMC2M4Vec(p1->momentum());
209:
210: // Check that ancestor is a central system
211: std::vector<HepMC3::ConstGenParticlePtr> results =
212: HepMC3::applyFilter("abs(HepMC3::StandardSelector::PDG_ID) == PDG::PDG_system,
213: HepMC3::Relatives::ANCESTORS(p1));
214:
215: if (results.size() != 0) { pip.push_back(pvec); }
216: }
217: for (HepMC3::ConstGenParticlePtr p1 :
218: HepMC3::applyFilter(HepMC3::StandardSelector::PDG_ID == NEGfinalPDG, evt.particles()) {
219: M4Vec pvec = gra::aux::HepMC2M4Vec(p1->momentum());
220:
221: // Check that ancestor is a central system
222: std::vector<HepMC3::ConstGenParticlePtr> results = HepMC3::applyFilter(
223: HepMC3::StandardSelector::PDG_ID == PDG::PDG_system, HepMC3::Relatives::ANCESTORS(p1));
224: if (results.size() != 0) { pim.push_back(pvec); }
225: }
226:
227: // CHECK CONDITION
228: if (pip.size() + pim.size() != (unsigned int) multiplicity) {
229: print(
230: "MNalyzer::ReadHepMC3: Multiplicity condition not filled [" + W +
231: "]=[" + W + "], multiplicity);
232: pip.clear();
233: continue; // skip event
234: }
235:
236: // -----
237: // CENTRAL SYSTEM pions
238: M4Vec system;
239: for (const auto& k : pip) { system += k; }
240: for (const auto& k : pim) { system += k; }
241:
242: std::vector<HepMC3::GenParticlePtr> beam_protons =
243: HepMC3::applyFilter(HepMC3::StandardSelector::STATUS == PDG::PDG_BEAM &&
244: HepMC3::StandardSelector::PDG_ID == PDG::PDG_p,
245: evt.particles());
246:
247: std::vector<HepMC3::GenParticlePtr> final_protons =
248: HepMC3::applyFilter(HepMC3::StandardSelector::STATUS == PDG::PDG_STABLE &&
249: HepMC3::StandardSelector::PDG_ID == PDG::PDG_p,

```

```

84: hEta_Pions =
85: std::make_shared<TH1D>(Form("hEta_p", "hEta p", ID.c_str()), "hEta/Events", NBINS, -12, 12);
86: hEta_Gamma =
87: std::make_shared<TH1D>(Form("hEta_gamma", "hEta gamma", ID.c_str()), "hEta/Events", NBINS, -12, 12);
88: hEta_Neutron =
89: std::make_shared<TH1D>(Form("hEta_n", "hEta n", ID.c_str()), "hEta/Events", NBINS, -12, 12);
90: hM4STAR =
91: std::make_shared<TH1D>(Form("hM4STAR", "M (GeV)", ID.c_str()), "M (GeV)/Events", NBINS, 0, 10);
92:
93: // Legendre polynomials, DO NOT CHANGE THE Y-RANGE [-1,1]
94: for (std::size_t i = 0; i < 8; ++i)
95: hP[i] =
96: std::make_shared<TF1>(Form("hP%i", "hP%i", ID.c_str()), "", 100, 0.0, 4.0, -1, 1);
97: hP[i] = >SetXTitle(Form("System M (GeV)"));
98: hP[i] = >SetYTitle(Form("Legendre P%i", ID.c_str()));
99:
100: // Costheta correlations between different frames
101: for (std::size_t i = 0; i < analyzer::FRAMES.size(); ++i) {
102: h2CosTheta[i][j] = std::make_shared<TH2D>(
103: Form("h2CosTheta%i%i", "h2CosTheta%i%i", analyzer::FRAMES[i].c_str(),
104: analyzer::FRAMES[j].c_str()), ID.c_str());
105:
106: Form("h2CosTheta%i%i", "h2CosTheta%i%i", analyzer::FRAMES[i].c_str(),
107: analyzer::FRAMES[j].c_str()), ID.c_str());
108: Form("h2CosTheta%i%i", "h2CosTheta%i%i", analyzer::FRAMES[i].c_str(),
109: analyzer::FRAMES[j].c_str()), ID.c_str());
110: NBINS, -1, 1, NBINS, -1, 1);
111: }
112:
113: // Phi correlations between different frames
114: for (std::size_t i = 0; i < analyzer::FRAMES.size(); ++i) {
115: h2Phi[i][j] = std::make_shared<TH2D>(
116: Form("h2Phi%i%i", "h2Phi%i%i", analyzer::FRAMES[i].c_str(),
117: analyzer::FRAMES[j].c_str()), ID.c_str());
118: Form("h2Phi%i%i", "h2Phi%i%i", analyzer::FRAMES[i].c_str(),
119: analyzer::FRAMES[j].c_str()), ID.c_str());
120: Form("h2Phi%i%i", "h2Phi%i%i", analyzer::FRAMES[i].c_str(),
121: analyzer::FRAMES[j].c_str()), ID.c_str());
122: NBINS, -grai::math::PI, grai::math::PI, NBINS, -grai::math::PI, grai::math::PI);
123: }
124: }
125:
126: // Destructor
127: MNalyzer::~MNalyzer() {}
128:
129: // Oracle histogram filler:
130: // Oracle here means that in this function we may use event tree information,
131: // not just pure fiducial final state information based on purely physical observables.
132: double MNalyzer::HepMC3_OracleFill(const std::string input, unsigned int multiplicity,
133: std::string finalPDG, unsigned int MAXOVERLAP,
134: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h1,
135: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h2,
136: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h3,
137: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h4,
138: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h5,
139: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h6,
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355: std::map<std::string, std::shared_ptr<HepMC3::GenParticle>> & h222,
356: std::map
```

```

./src/Analysis/MAnalyzer.cc          5/11
333: // const double t2 = -(p_beam_minus - p_final_minus).M(1);
334:
335: // Deltapl1
336: deltapl1_pp = p_final_minus.DeltaPhiAbs(p_final_minus);
337: M4Vec pp_diff = p_final_minus - p_initial_minus;
338: const double pp_dpt = pp_diff.P();
339:
340: h1["hl_pp_dphi"]->h[SID]>F(11, deltapl1_pp, W);
341: h1["hl_pp_ct"]->h[SID]>F(11, ct, W);
342: h1["hl_pp_dpt"]->h[SID]>F(11, pp_dpt, W);
343:
344: h2["h2_s_m_dphi"]->h[SID]>F(11, M, deltapl1_pp, W);
345: h2["h2_s_m_dpt"]->h[SID]>F(11, M, pp_dpt, W);
346: h2["h2_s_m_ct"]->h[SID]>F(11, M, ct, abs(ct), W);
347:
348:
349: // 2D
350: h2["h2_s_m_pt"]->h[SID]>F(11, M, Pt, W);
351: h2["h2_s_m_pt"]->h[SID]>F(11, M, a.Pt(), W);
352:
353: // 2-Body only
354: if (multiplicity == 2) {
355:   h1["hl_2b_m_dphi"]->h[SID]>F(11, M, a.DeltaPhi(), W);
356:   h1["hl_2b_m_dphi"]->h[SID]>F(11, M, a.DeltaPhi(), W);
357:   h1["hl_2b_diffrap"]->h[SID]>F(11, b.Rap(), a.Rap(), W);
358:   h2["h2_2b_m_dphi"]->h[SID]>F(11, M, a.DeltaPhi(), W);
359:   h2["h2_2b_m_dpt"]->h[SID]>F(11, M, a.Pt(), W);
360:
361: // Frame transform
362: const int direction = j;
363: const M4Vec X = a, b;
364:
365: std::vector<M4Vec> CM = {a, b};
366: gra::Kinematics::CMFrame(CM, X);
367:
368: std::vector<M4Vec> HX = {a, b};
369: gra::Kinematics::HXFrame(HX, X);
370:
371: std::vector<M4Vec> CS = {a, b};
372: gra::Kinematics::CSFrame(CS, X, p_beam_plus, p_beam_minus);
373:
374: std::vector<M4Vec> GJ = {a, b};
375: gra::Kinematics::GJFrame(GJ, X, direction, p_beam_plus - p_final_plus,
376:   p_beam_minus - p_final_minus);
377:
378: std::vector<M4Vec> PG = {a, b};
379: gra::Kinematics::PFrame(PG, X, direction, p_beam_plus, p_beam_minus);
380:
381:
382: h1["hl_costheta_CM"]->h[SID]>F(11, CM[0].CosTheta(), W);
383: h1["hl_costheta_HX"]->h[SID]>F(11, HX[0].CosTheta(), W);
384: h1["hl_costheta_CS"]->h[SID]>F(11, CS[0].CosTheta(), W);
385: h1["hl_costheta_GJ"]->h[SID]>F(11, GJ[0].CosTheta(), W);
386: h1["hl_costheta_PG"]->h[SID]>F(11, PG[0].CosTheta(), W);
387: h1["hl_costheta_P"]->h[SID]>F(11, P.CosTheta(), W);
388:
389:
390: h1["hl_phi_CM"]->h[SID]>F(11, CM[0].Phi(), W);
391: h1["hl_phi_HX"]->h[SID]>F(11, HX[0].Phi(), W);
392: h1["hl_phi_CS"]->h[SID]>F(11, CS[0].Phi(), W);
393: h1["hl_phi_GJ"]->h[SID]>F(11, GJ[0].Phi(), W);
394: h1["hl_phi_PG"]->h[SID]>F(11, PG[0].Phi(), W);
395: h1["hl_phi_P"]->h[SID]>F(11, P.Phi(), W);
396:
397: h1["hl_phi_LAB"]->h[SID]>F(11, a.Phi(), W);
398:
399: h2["h2_2b_costheta_phi_CM"]->h[SID]>F(11, CM[0].CosTheta(), CM[0].Phi(), W);
400: h2["h2_2b_costheta_phi_HX"]->h[SID]>F(11, HX[0].CosTheta(), HX[0].Phi(), W);
401: h2["h2_2b_costheta_phi_CS"]->h[SID]>F(11, CS[0].CosTheta(), CS[0].Phi(), W);
402: h2["h2_2b_costheta_phi_GJ"]->h[SID]>F(11, GJ[0].CosTheta(), GJ[0].Phi(), W);
403: h2["h2_2b_costheta_phi_PG"]->h[SID]>F(11, PG[0].CosTheta(), PG[0].Phi(), W);
404: h2["h2_2b_costheta_phi_LAB"]->h[SID]>F(11, a.CosTheta(), a.Phi(), W);
405:
406:
407: h2["h2_2b_m_costheta_CM"]->h[SID]>F(11, M, CM[0].CosTheta(), W);
408: h2["h2_2b_m_costheta_HX"]->h[SID]>F(11, M, HX[0].CosTheta(), W);
409: h2["h2_2b_m_costheta_CS"]->h[SID]>F(11, M, CS[0].CosTheta(), W);
410: h2["h2_2b_m_costheta_GJ"]->h[SID]>F(11, M, GJ[0].CosTheta(), W);
411: h2["h2_2b_m_costheta_PG"]->h[SID]>F(11, M, PG[0].CosTheta(), W);
412: h2["h2_2b_m_costheta_LAB"]->h[SID]>F(11, M, a.CosTheta(), W);
413:
414:
415: h2["h2_2b_m_phi_CM"]->h[SID]>F(11, M, CM[0].Phi(), W);

```

```

./src/Analysis/MAnalyzer.cc          6/11
416: h2["h2_2b_m_phi_HX"]->h[SID]>F(11, M, HX[0].Phi(), W);
417: h2["h2_2b_m_phi_CS"]->h[SID]>F(11, M, CS[0].Phi(), W);
418: h2["h2_2b_m_phi_GJ"]->h[SID]>F(11, M, GJ[0].Phi(), W);
419: h2["h2_2b_m_phi_PG"]->h[SID]>F(11, M, PG[0].Phi(), W);
420: h2["h2_2b_m_phi_LAB"]->h[SID]>F(11, M, a.Phi(), W);
421:
422: // -----
423: NP["np_s_m_p12_CM"]->h[SID]>F(11, M, math::LegendreP(2, CM[0].CosTheta()), W);
424: NP["np_s_m_p12_CM"]->h[SID]>F(11, M, math::LegendreP(2, CM[0].CosTheta()), W);
425: NP["np_s_m_p12_CM"]->h[SID]>F(11, M, math::LegendreP(4, CM[0].CosTheta()), W);
426: NP["np_s_m_p12_CM"]->h[SID]>F(11, M, math::LegendreP(4, CM[0].CosTheta()), W);
427:
428: // -----
429:
430: // 4-body only
431: if (multiplicity == 4) {
432:   // ...
433:
434:   auto catch (...) {
435:     throw std::invalid_argument("MAnalyzer::HepMC3Read: Problem filling histogram!");
436:   }
437:
438: // << SUPERPLOTTER
439: // -----
440:
441: if (events_read == MAXEVENTS) {
442:   std::cout << "MAnalyzer::HepMC3Read: Maximum event count "<< MAXEVENTS << " reached!";
443:   break; // Enough events
444: }
445:
446: if (events_read % 10000 == 0) {
447:   std::cout << std::endl << "Events processed: " << events_read << std::endl;
448: }
449: // [THIS AS LAST]: Sum selected event weights
450: selectW = W;
451:
452: std::cout << std::endl;
453: std::cout << "MAnalyzer::HepMC3Read: Events processed in total: " << events_read << std::endl;
454:
455: // Close HepMC3 file
456: input_file.close();
457:
458: if (selectW == 0.0) {
459:   throw std::invalid_argument("MAnalyzer::HepMC3Read: Valid events in " << totalpath << " + "
460:     + " out of " << events_read << " events.");
461: }
462: // Take into account extra fiducial cut efficiency here
463: double efficiency = selectW / totalW;
464: printf("MAnalyzer::HepMC3Read: Fiducial cut efficiency: %0.3f %\n", efficiency);
465: std::cout << std::endl;
466:
467: return cross_section * efficiency;
468:
469: // Sanity check
470: double MAnalyzer::CheckEnergyMomentum(HepMC3::GenEvent evt) const {
471:   std::vector<HepMC3::GenParticlePtr> all_init = HepMC3::applyFilter(
472:     HepMC3::StandardSelector::STATUS == PDG::PDG_BSM, evt.particles()); // Beam
473:   std::vector<HepMC3::GenParticlePtr> all_final = HepMC3::applyFilter(
474:     HepMC3::StandardSelector::STATUS == PDG::PDG_STABLE, evt.particles()); // Final state
475:
476:   M4Vec beam(0, 0, 0, 0);
477:   for (const HepMC3::GenParticlePtr p : all_init)
478:     beam += gra::aux::HepMC2M4Vec(p->momentum());
479:
480:   M4Vec final(0, 0, 0, 0);
481:   for (const HepMC3::GenParticlePtr p : all_final)
482:     final += gra::aux::HepMC2M4Vec(p->momentum());
483:
484:   if (fabs(smash::CheckMC(beam - final)) > 1e-10) {
485:     if (grs::smash::CheckMC(beam - final)) {
486:       grs::aux::PrintWarning();
487:       std::cout << rang::fg::red << "Energy-Momentum not conserved!" << rang::fg::reset << std::endl;
488:       (beam - final).Print();
489:       HepMC3::Print::listing(evt);
490:       HepMC3::Print::icount(evt);
491:     }
492:   }
493:   return beam.M();
494: }
495:
496: // 2-body angular observables
497: void MAnalyzer::FrameObservables(double W, HepMC3::GenEvent evt, const M4Vec ep_beam_plus,
498:   const M4Vec ep_beam_minus, const M4Vec ep_final_plus,

```

```

./src/Analysis/MAnalyzer.cc          7/11
499: // Find index
500: const auto ind = k;
501: const auto ind = k;
502: const auto ind = k;
503: const auto ind = k;
504: const auto ind = k;
505: const auto ind = k;
506: const auto ind = k;
507: const auto ind = k;
508: const auto ind = k;
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571: const auto ind = k;
572: const auto ind = k;
573: const auto ind = k;
574: const auto ind = k;
575: const auto ind = k;
576: const auto ind = k;
577: const auto ind = k;
578: const auto ind = k;
579: const auto ind = k;
580: const auto ind = k;
581: const auto ind = k;

```

```

./src/Analysis/MAnalyzer.cc          8/11
582: bool excited_minus = false;
583: for (const HepMC3::GenParticlePtr p : search_nstar) {
584:   M4Vec pvec = gra::aux::HepMC2M4Vec(p->momentum());
585:   hX_PIONS->Fill(pvec.M(), W);
586:
587:   if (pvec.P() > 0) { excited_plus = true; }
588:   if (pvec.P() < 0) { excited_minus = true; }
589: }
590: // Excited system found
591: if (excited_plus || excited_minus) { h_STAR_ON = true; }
592:
593: // * system decay products
594:
595: // Gammas
596: double gamma_e_plus = 0;
597: double gamma_e_minus = 0;
598:
599: // Gammas
600: for (const HepMC3::GenParticlePtr p : search_gamma) {
601:   M4Vec pvec = gra::aux::HepMC2M4Vec(p->momentum());
602:
603:   // Check that ancestor is excited forward system
604:   std::vector<HepMC3::GenParticlePtr> ancestor =
605:     HepMC3::applyFilter(HepMC3::StandardSelector::PDG_ID == PDG::PDG_NSTAR ||
606:       HepMC3::StandardSelector::PDG_ID == -PDG::PDG_NSTAR,
607:       HepMC3::Relatives::ANCESTORS(p));
608:
609:   if (ancestor.size() != 0) {
610:     hE_GAMMA->Fill(pvec.E(), W);
611:     hE_GAMMA->Fill(pvec.E(), W);
612:     hX_PIONS->Fill(pvec.P() / (sqrt(2)), W);
613:
614:     if (excited_plus && pvec.Rap() > 0) { gamma_e_plus += pvec.E(); }
615:     if (excited_minus && pvec.Rap() < 0) { gamma_e_minus += pvec.E(); }
616:   }
617: }
618:
619: // P+
620: for (const HepMC3::GenParticlePtr p : search_pi) {
621:   // HepMC3::Print::listing(p);
622:
623:   // Check that parent is the excited system
624:   std::vector<HepMC3::GenParticlePtr> parents = p->parents();
625:   bool found = false;
626:   for (const auto k : indices(parents)) {
627:     if (std::abs(parents[k]-p[0]) == PDG::PDG_NSTAR) { found = true; }
628:   }
629:   if (found) {
630:     M4Vec pvec = gra::aux::HepMC2M4Vec(p->momentum());
631:     hE_PIONS->Fill(pvec.E(), W);
632:     hE_PIONS->Fill(pvec.E(), W);
633:     hX_PIONS->Fill(pvec.P() / (sqrt(2)), W);
634:   }
635: }
636:
637: // P-
638: for (const HepMC3::GenParticlePtr p : search_pi) {
639:   // HepMC3::Print::listing(p);
640:
641:   // Check that parent is the excited system
642:   std::vector<HepMC3::GenParticlePtr> parents = p->parents();
643:   bool found = false;
644:   for (const auto k : indices(parents)) {
645:     if (std::abs(parents[k]-p[0]) == PDG::PDG_NSTAR) { found = true; }
646:   }
647:   if (found) {
648:     M4Vec pvec = gra::aux::HepMC2M4Vec(p->momentum());
649:     hE_PIONS->Fill(pvec.E(), W);
650:     hE_PIONS->Fill(pvec.E(), W);
651:     hX_PIONS->Fill(pvec.P() / (sqrt(2)), W);
652:   }
653: }
654:
655: // Neutrons
656: double neutron_e_plus = 0;
657: double neutron_e_minus = 0;
658:
659: for (const HepMC3::GenParticlePtr p : search_neutrons) {
660:   M4Vec pvec = gra::aux::HepMC2M4Vec(p->momentum());
661:
662:   // Check that parent is the excited system
663:   std::vector<HepMC3::GenParticlePtr> parents = p->parents();
664:   bool found = false;

```

```

./src/Analysis/MAnalyzer.cc          9/11
665:   for (const auto &k : indices(parents)) {
666:       if (std::abs(parents[k]->pid()) == PDG::PDG_NSTAR) { found = true;
667:       }
668:       if (found) {
669:           hEta_Neutron->Fill(pvec.Eta(), W);
670:           hE_Neutron->Fill(pvec.E(), W);
671:           hXF_Neutron->Fill(pvec.Pz() / (sqrtz / 2), W);
672:       }
673:       if (excited_plus && pvec.Rap() > 0) { neutron_n_plus += pvec.E(); }
674:       if (excited_minus && pvec.Rap() < 0) { neutron_n_minus += pvec.E(); }
675:   }
676: }
677:
678: // Gamma/Neutron energy histogram
679: if (excited_plus) { hE_GammaNeutron->Fill(gamma_n_plus + neutron_n_plus, W); }
680: if (excited_minus) { hE_GammaNeutron->Fill(gamma_n_minus + neutron_n_minus, W); }
681: }
682:
683: double powerlaw(double *x, double *par) {
684:     return par[0] / std::pow(1.0 + std::pow(x[0], 2) / (std::pow(par[1], 2) * par[2]), par[2]);
685: }
686:
687: double exponential(double *x, double *par) { return par[0] * exp(par[1] * x[0]); }
688:
689: // Custom plotter
690: void MAnalyzer::PlotAll(const std::string &titlestr) {
691:     // Create output directory if it does not exist
692:     const std::string FOLDER = gra::aux::GetBasePath(2) + "/" + titlestr + " + inputfile;
693:     aux::CreateDirectory(FOLDER);
694:
695:     // FIT FUNCTIONS
696:     std::shared_ptr<TF1> fD = std::make_shared<TF1>("exp_fit", exponential, 0.05,
697:     0.5, 2);
698:     fD->SetParameter(0, 10.0); // A
699:     fD->SetParameter(1, -8.0); // b
700:
701:     std::shared_ptr<TF1> fa = std::make_shared<TF1>("pow_fit", powerlaw, 0.5, 3.0,
702:     3);
703:     fa->SetParameter(0, 10.0); // A
704:     fa->SetParameter(1, 0.15); // T
705:     fa->SetParameter(2, 0.1); // n
706:
707:     //
708:     hF2->Fit("exp_fit", "R"); // "R" for range
709:     hF2->Fit("pow_fit", "R"); // "R" for range
710:
711:     // ***** FORWARD EXCITATION *****
712:     if (!N_STAR_ON) {
713:         // Draw histograms
714:         TCanvas c1("c", "c", 800, 600);
715:         c1.Divide(2, 2, 0.0002, 0.0002);
716:
717:         c1.cd(1);
718:         // gPad->SetLog();
719:         hEta_Gamma->SetLineColor(2);
720:         hEta_Pions->SetLineColor(4);
721:         hEta_Neutron->SetLineColor(8);
722:
723:         hEta_Gamma->Draw("same");
724:         hEta_Neutron->Draw("same");
725:         hEta_Pions->Draw("same");
726:
727:         // Set x-axis range
728:         const double eta_min = 3;
729:         const double eta_max = 15;
730:         hEta_Pions->SetAxisRange(eta_min, eta_max, "X");
731:         hEta_Gamma->SetAxisRange(eta_min, eta_max, "X");
732:         hEta_Neutron->SetAxisRange(eta_min, eta_max, "X");
733:
734:         c1.cd(2);
735:         // gPad->SetLog();
736:         hXF_Gamma->SetLineColor(2);
737:         hXF_Pions->SetLineColor(4);
738:         hXF_Neutron->SetLineColor(8);
739:
740:         hXF_Gamma->Draw("same");
741:         hXF_Pions->Draw("same");
742:         hXF_Neutron->Draw("same");
743:
744:         c1.cd(3);
745:         gPad->SetLog();
746:         // gra->SetLogz();
747:         hM_NSTAR->Draw();

```

```

./src/Analysis/MAnalyzer.cc          11/11
831:   hP1[1]->SetMinimum(-0.4); // Y-axis minimum
832:   hP1[1]->SetMaximum(0.4); // Y-axis maximum
833:
834:   // Adjust legend
835:   leg1[1]->SetFillColor(0); // White background
836:   leg1[1]->SetBorderSize(0); // No box
837:   leg1[1]->AddEntry(hP1[1].get(), Form("1 = %lu", 1 + 1), "l");
838:   leg1[1]->Draw("SAME");
839:
840:   // Horizontal line
841:   line[1] = new TLine(xBOUND[0], 0, xBOUND[1], 0);
842:   line[1]->Draw("SAME");
843:
844:   // Titledr
845:   if (l == 0 || l == 4) { hP1[1]->SetTitle(titlestr.c_str()); }
846: }
847:
848: c15->SaveAs(
849:     Form("%s/figs/figs/hP1_t04.pdf", gra::aux::GetBasePath(2).c_str(), inputfile.c_str()));
850: c16->SaveAs(
851:     Form("%s/figs/figs/hP1_t08.pdf", gra::aux::GetBasePath(2).c_str(), inputfile.c_str()));
852:
853: for (std::size_t i = 0; i < 8; ++i) {
854:     delete leg1[i];
855: }
856: delete c15;
857: delete c16;
858: }
859:
860: // namespace gra
861:

```

```

./src/Analysis/MAnalyzer.cc          10/11
748:   c1.cd(4);
749:   gPad->SetLog();
750:   hXF_Gamma->SetLineColor(2);
751:   hXF_Pions->SetLineColor(4);
752:   hXF_Neutron->SetLineColor(8);
753:
754:   hXF_Gamma->Draw("same");
755:   hXF_Pions->Draw("same");
756:   hXF_Neutron->Draw("same");
757:
758:   c1.SaveAs(Form("%s/figs/figs/forward.pdf", gra::aux::GetBasePath(2).c_str(), inputfile.c_str()));
759:
760:   // *****
761:   // *****
762:   // *****
763:   // *****
764:   // *****
765:   // *****
766:   TCanvas c2("c", "c", 800, 800);
767:   c2.Divide(analyzer::FRAMES.size(), analyzer::FRAMES.size(), 0.0001, 0.0002);
768:
769:   int k = 1;
770:   for (std::size_t i = 0; i < analyzer::FRAMES.size(); ++i) {
771:       for (std::size_t j = 0; j < analyzer::FRAMES.size(); ++j) {
772:           c2.cd(k);
773:           ++k;
774:           if (j >= 1) { h2CosTheta[4][j]->Draw("COLZ"); }
775:
776:           // Titledr
777:           if ((i == 0) && (j == 0)) { h2CosTheta[4][j]->SetTitle(titlestr.c_str()); }
778:       }
779:   }
780:   c2.SaveAs(Form("%s/figs/figs/h2_frame_correlations_ootheta.pdf", gra::aux::GetBasePath(2).c_str(),
781:   inputfile.c_str()));
782:
783:   // *****
784:   // *****
785:   TCanvas c3("c", "c", 800, 800);
786:   c3.Divide(analyzer::FRAMES.size(), analyzer::FRAMES.size(), 0.0001, 0.0002);
787:
788:   int k = 1;
789:   for (std::size_t i = 0; i < analyzer::FRAMES.size(); ++i) {
790:       for (std::size_t j = 0; j < analyzer::FRAMES.size(); ++j) {
791:           c3.cd(k);
792:           ++k;
793:           if (j >= 1) { h2Phi[4][j]->Draw("COLZ"); }
794:
795:           // Titledr
796:           if ((i == 0) && (j == 0)) { h2Phi[4][j]->SetTitle(titlestr.c_str()); }
797:       }
798:   }
799:   c3.SaveAs(Form("%s/figs/figs/h2_frame_correlations_phi.pdf", gra::aux::GetBasePath(2).c_str(),
800:   inputfile.c_str()));
801:
802:   // *****
803:   // *****
804:   // Legendre polynomials in the Rest Frame (non-rotated one)
805:
806:   const int colors[8] = {48, 53, 98, 32, 48, 53, 98, 32};
807:   const double XBOUND[2] = {0.0, 4.0};
808:
809:   // ...
810:   TCanvas *c15 = new TCanvas("c15", "Legendre polynomials 1-4", 600, 400);
811:   TCanvas *c16 = new TCanvas("c16", "Legendre polynomials 5-8", 600, 400);
812:
813:   TLegend *leg1[8];
814:   TLine *line[8];
815:   c15->Divide(2, 2, 0.001, 0.001);
816:   c16->Divide(2, 2, 0.001, 0.001);
817:
818:   for (std::size_t i = 0; i < 8; ++i) {
819:       if (i < 4) {
820:           c15->cd(i + 1);
821:           *line = new TLine(XBOUND[0], 0, XBOUND[1], 0);
822:           c15->cd(i - 3);
823:       }
824:
825:       // Legend
826:       leg1[i] = new TLegend(0.15, 0.75, 0.4, 0.85); // x1,y1,x2,y2
827:
828:       // Data
829:       hP1[i]->SetLineColor(colors[i]);
830:       hP1[i]->Draw("SAME");

```

```

./src/Analysis/MMultiplot.cc        1/8
1:   (Histo, Histo, Histo2, ..., Histo N-1) Multiplot ROOT histograms
2:
3:   (c) 2017-2020 Mikael Mieskolainen
4:   Licensed under the MIT License <http://opensource.org/licenses/MIT>.
5:
6:   C++
7:   #include <string>
8:   #include <tuple>
9:
10:  // ROOT
11:  #include "TColor.h"
12:  #include "TAxis.h"
13:  #include "TH1.h"
14:  #include "TH2.h"
15:  #include "TMath.h"
16:  #include "TLegend.h"
17:  #include "TText.h"
18:  #include "TProfile.h"
19:
20:  // Own
21:  #include "Graniiti/Analysis/Multiplot.h"
22:  #include "Graniiti/Analysis/MROOT.h"
23:  #include "Graniiti/MAux.h"
24:
25:  using gra::aux::indices;
26:
27:  namespace gra {
28:
29:  void GetLegendPosition(unsigned int N, double x1, double x2, double y1, double y2,
30:   const std::string &legendposition) {
31:      // North-East
32:      x1 = 0.70;
33:      x2 = x1 + 0.18;
34:      y1 = 0.75 - 0.02 * N;
35:
36:      // South-East
37:      if (legendposition.compare("southeast") == 0) { y1 = 0.10 - 0.01 * N; }
38:      y2 = y1 + 0.04 * N; // Scale by the number of histograms
39:  }
40:
41:  hMultiplot::hMultiplot(const std::string &name, const std::string &labeltext, int N,
42:   double minval, double maxval, const std::vector<std::string> &legendtext) {
43:      N_ = N;
44:      name_ = name;
45:      minval_ = minval;
46:      maxval_ = maxval;
47:      legendtext_ = legendtext;
48:      legendposition_ = "southeast";
49:
50:      // Initialize histogram vector container size
51:      h.resize(legendtext.size());
52:
53:      for (const auto &i : indices(h)) {
54:          h[i] = new TH1D(Form("%s%lu", name.c_str(), i), labeltext.c_str(), N, minval, maxval);
55:          h[i]->Sumw2(); // Error saving on
56:      }
57:
58:
59:  // Histogram normalization
60:  void hMultiplot::NormalizeAll(const std::vector<double> &cross_section,
61:   const std::vector<double> &multiplier) {
62:      for (const auto &i : indices(h)) {
63:          double scale = 1.0;
64:
65:          // Takes into account weighted and unweight (weight=1) filling
66:          // We take into account also under/overflow here with macro: (0, n bins+1)
67:          const double integral = h[i]->Integral(0, h[i]->GetNbinsX() + 1);
68:          if (integral > 0) { scale /= integral; }
69:
70:          // Binwidth
71:          const int bin = 1;
72:          const double binwidth = h[i]->GetXaxis()->GetBinWidth(bin);
73:          if (binwidth > 0) { scale /= binwidth; }
74:
75:          // Cross Section
76:          if (cross_section[i] > 0) { scale *= cross_section[i]; }
77:
78:          // Additional scale factor
79:          scale *= multiplier[i];
80:
81:          h[i]->Scale(scale);
82:      }
83:  }

```



```

./src/Analysis/Multiplet.cc      2/8
84:
85: std::vector<double> hMultiplet::SaveFig(const std::string fullpath, bool RATIOPLOTT const (
86: // New color scheme
87: std::vector<int> color;
88: std::vector<int> shared_ptr<Color> rootcolor;
89: rootstyle::CreateColorMap(color, rootcolor);
90:
91: // -----
92: // Apply the chi2 test and retrieve the residuals
93: gra::aux::PrintBar("");
94: std::vector<double> chi2ndf(h.size(), 0.0);
95:
96: for (const auto k1 : indices(h) {
97:     double res[N];
98:     printf("%s [%lu] : %n", legendtext_[1].c_str(), k1);
99:     double c2ndf = h[0] >= chi2Test(h[1], "M = CH2/ND", res);
100:     chi2ndf[k1] = c2ndf;
101:
102:     printf("chi2/ndf = %0.3f %s\n", c2ndf);
103: }
104: gra::aux::PrintBar("");
105:
106: TCanvas c0("c", "c", 750, 800);
107:
108: // Upper plot will be in pad1
109: std::shared_ptr<TPad> pad1;
110:
111: if (RATIOPLOTT) {
112:     pad1 = std::make_shared<TPad>("pad1", "pad1", 0, 0.3, 1.0, 1.0);
113:     pad1->SetBottomMargin(0.015); // Upper and lower plot are joined
114: } else {
115:     pad1 = std::make_shared<TPad>("pad1", "pad1", 0, 0.0, 1.0, 1.0);
116: }
117: pad1->SetGrid(); // Vertical grid
118: pad1->Draw(); // Draw the upper pad: pad1
119: pad1->cd(); // pad1 becomes the current pad
120: h[0]->SetStats(0); // No statistics on upper plot
121:
122: // Find maximum value for y-range limits
123: double MAXVAL = 0.0;
124: for (const auto k1 : indices(h) {
125:     const double max = h[1]->GetMaximum();
126:     if (max > MAXVAL) MAXVAL = max;
127: }
128:
129: // Find minimum (non-zero) value for y-range limits
130: double MINVAL = 1e32;
131: for (const auto k1 : indices(h) {
132:     for (int k = 0; k < h[1]->GetBinWidth(k); ++k) {
133:         double value = h[1]->GetBinContent(k);
134:         if (value < MINVAL && value > 0) MINVAL = value;
135:     }
136: }
137:
138: // Loop over histograms
139: for (const auto k1 : indices(h) {
140:     h[1]->SetLineColor(color[k1]);
141:     h[1]->SetLineStyle(2);
142:     h[1]->SetMarkerStyle(20 + k1);
143:     h[1]->SetMarkerColor(color[k1]);
144:     h[1]->SetMarkerSize(20 + k1);
145:     h[1]->SetMarkerSize(0.73);
146:
147:     std::cout << MINVAL << std::endl;
148:     h[1]->GetYaxis()->SetRangeUser(0, MAXVAL * 1.35);
149:
150:     // h[1]->Draw("x2 same"); // Needs this combo to draw box-lines
151:     h[1]->Draw("hist same");
152:
153:     // h[1]->SetFillStyle(4050);
154:     // h[1]->SetFillColor(color[k1]);
155:     // h[1]->SetFillColorAlpha(color[k1], 0.1);
156:     // h[1]->Draw("x hist same"); // Filled
157: }
158:
159: pad1->RedrawAxis(); // Fix overlapping histogram lines on borders
160:
161: // -----
162: // Create dynamically sized legend (depending on number of legend entries)
163:
164: double x1, x2, y1, y2 = 0.0;
165: legendposition(h.size(), x1, x2, y1, y2, legendposition);
166: std::shared_ptr<TLegend> legend = std::make_shared<TLegend>(x1, y1, x2, y2);

```

```

./src/Analysis/Multiplet.cc      4/8
250: // Save logscale pdf
251: if (MINVAL > 0) {
252:     pad1->cd(); // Sety axis() // pad2 becomes the current pad
253:     for (const auto k1 : indices(h) {
254:         h[1]->GetYaxis()->SetRangeUser(std::max(MINVAL, MAXVAL * 1e-8), MAXVAL * 5);
255:     }
256:     fullfile = fullpath + name_ + ".logy" + ".pdf";
257:     c0.SaveAs(fullfile.c_str());
258: }
259:
260: // Remove histograms
261: if (RATIOPLOTT) {
262:     for (const auto k1 : indices(hR)) { delete hR[k1]; }
263: }
264:
265: return chi2ndf;
266: }
267:
268: h2Multiplet::h2Multiplet(const std::string &name, const std::string &labeltext, int N1,
269:     double minval1, double maxval1, int N2, double minval2, double maxval2,
270:     const std::vector<std::string> &legendtext) {
271:     name_ = name;
272:     N1_ = N1;
273:     N2_ = N2;
274:     minval1_ = minval1;
275:     maxval1_ = maxval1;
276:     minval2_ = minval2;
277:     maxval2_ = maxval2;
278:
279:     legendtext_ = legendtext;
280:
281:     // Initialize histogram vector container size
282:     h.resize(legendtext_.size());
283:
284:     for (const auto k1 : indices(h)) {
285:         h[k1] = new TH2D(Form("%s_%lu", name.c_str(), k1), labeltext.c_str(), N1, minval1, maxval1, N2,
286:             minval2, maxval2);
287:     }
288:     h[1]->Sum2(); // Error saving on
289: }
290:
291:
292:
293: void h2Multiplet::NormalizeAll(const std::vector<double> &cross_section,
294:     const std::vector<double> &multiplier) {
295:     for (const auto k1 : indices(h)) {
296:         double scale = 1.0;
297:
298:         // Takes into account weighted and unweight (weight=1) filling
299:         // We take into account also under/overflow here with macro: (0, hbins)
300:         const double integral = h[1]->Integral(0, h[1]->GetBinWidth(k1) + 1,
301:             if (integral > 0) (scale /= integral);
302:
303:         // Binwidth
304:         const int bin = 1;
305:         const double binwidth = h[1]->GetBinWidth(bin) * h[1]->GetXaxis()->GetBinWidth(bin);
306:         if (binwidth > 0) (scale /= binwidth);
307:
308:         // Cross Section
309:         if (cross_section[k1] > 0) (scale *= cross_section[k1]);
310:
311:         // Additional scale factor
312:         scale *= multiplier[k1];
313:
314:         h[1]->Scale(scale);
315:     }
316: }
317:
318: double h2Multiplet::SaveFig(const std::string fullpath, bool RATIOPLOTT const (
319:     const double scale = (RATIOPLOTT == true) ? 2.0 : 1.0;
320:
321: TCanvas c0("c", "c", h.size() == 1 ? 750 : 750 / 3.0 * h.size(),
322:     260 * scale); // scale canvas according to number of sources
323: c0.Divide(h.size(), h.size(), 1, 2, 1, scale, 0.01, 0.02); // [columns] x [rows]
324:
325: // Normalize by the first histogram
326: const double ZMAX = h[0]->GetMaximum();
327:
328: // Histograms on TOP ROW
329: for (const auto k1 : indices(h)) {
330:     c0.cd(1 + k1); // choose position
331:     c0.cd(1) >= SetRightMargin(0.13);
332:

```

```

./src/Analysis/Multiplet.cc      3/8
167: legend->SetFillColor(0); // White background
168: // legend->SetBorderSize(0); // No box
169:
170: // Add legend entries
171: for (const auto k1 : indices(h)) { legend->AddEntry(h[k1], legendtext_[1].c_str()); }
172:
173:
174: // -----
175: // Create Labels
176: std::shared_ptr<TLabel> l1;
177: std::shared_ptr<TLabel> l2;
178: gra::rootstyle::MakeInFinland(l1, l2, 0.935, (RATIOPLOTT == true ? 0.03 : (0.16, 0.58)));
179:
180:
181: // Ratio plots
182: std::vector<TH2D *> hR(h.size(), nullptr);
183: std::shared_ptr<TPad> pad2;
184: std::shared_ptr<TLine> line;
185:
186: if (RATIOPLOTT) {
187:     c0.cd();
188:     pad2 = std::make_shared<TPad>("pad2", "pad2", 0, 0.05, 1, 0.3);
189:     pad2->SetTopMargin(0.025);
190:     pad2->SetBottomMargin(0.025);
191:     pad2->SetGrid(); // vertical grid
192:     pad2->Draw();
193:     pad2->cd(); // pad2 becomes the current pad
194:
195:     for (const auto k1 : indices(h)) {
196:         hR[k1] = static_cast<TH2D *>(h[1]->Clone(Form("ratio_%lu", k1)));
197:         hR[k1]->Divide(h[0]);
198:
199:         hR[k1]->SetMinimum(0.0); // y-axis range
200:         hR[k1]->SetMaximum(2.0); //
201:         hR[k1]->SetStats(0); // statistics box off
202:         hR[k1]->Draw("same"); // ratio plot
203:
204:         // Ratio plot (h3) settings
205:         hR[k1]->SetTitle(""); // Remove the ratio title
206:
207:         // Y axis ratio plot settings
208:         hR[k1]->GetYaxis()->SetTitle("Ratio");
209:         hR[k1]->GetYaxis()->GetCenterTitle();
210:         hR[k1]->GetYaxis()->SetMajorDivisions(505);
211:         hR[k1]->GetYaxis()->SetLabelFont(43);
212:         hR[k1]->GetYaxis()->SetLabelFont(43);
213:         hR[k1]->GetYaxis()->SetLabelOffset(1.55);
214:         hR[k1]->GetYaxis()->SetLabelFont(43); // Absolute font size in pixel (precision 3)
215:         hR[k1]->GetYaxis()->SetLabelSize(15);
216:
217:         // X axis ratio plot settings
218:         hR[k1]->GetXaxis()->SetLabelFont(43);
219:         hR[k1]->GetXaxis()->SetLabelFont(43);
220:         hR[k1]->GetXaxis()->SetLabelOffset(4.);
221:         hR[k1]->GetXaxis()->SetLabelFont(43); // Absolute font size in pixel (precision 3)
222:         hR[k1]->GetXaxis()->SetLabelSize(15);
223:     }
224:
225:     // Draw horizontal line
226:     const double ymax = 1.0;
227:     line = std::make_shared<TLine>(minval1, ymax, maxval1, ymax);
228:     line->SetLineColor(15);
229:     line->SetLineWidth(2.0);
230:     line->Draw();
231:
232: // -----
233: // Remove x-axis of UPPER PLOT
234: h[0]->GetXaxis()->SetLabelOffset(999);
235: h[0]->GetXaxis()->SetLabelSize(0);
236: // -----
237:
238: } // RATIOPLOTT
239:
240: // Give y-axis title some offset to avoid overlapping with numbers
241: h[0]->GetYaxis()->SetTitleOffset(1.45);
242:
243: // Create output directory if it does not exist
244: aux::CreateDirectory(fullpath);
245:
246: // Save pdf
247: std::string fullfile = fullpath + name_ + ".pdf";
248: c0.SaveAs(fullfile.c_str());
249:
333: h[1]->SetStats(0);
334: h[1]->Draw("COLS");
335: h[1]->GetYaxis()->SetTitleOffset(1.3);
336: // h[1]->GetXaxis()->SetRangeUser(0.0, ZMAX);
337: h[1]->SetTitle(legendtext_[1].c_str());
338: }
339:
340: // Ratio histograms on BOTTOM ROW
341: std::vector<TH2D *> hR(h.size(), nullptr);
342:
343: if (RATIOPLOTT) {
344:     if (h.size() > 1) { // Only if we have more than 1
345:         for (const auto k1 : indices(h)) {
346:             c0.cd(1 + k1 * h.size()); // Choose position
347:             c0.cd(1 + k1 * h.size()) >= SetRightMargin(0.13);
348:
349:             hR[k1] = static_cast<TH2D *>(h[1]->Clone(Form("h2R_%lu", k1)));
350:
351:             hR[k1]->Divide(h[0]); // Divide by 0-th histogram
352:             hR[k1]->SetLabelOffset(1.3);
353:             hR[k1]->SetStats(0); // No statistics on upper plot
354:             hR[k1]->Draw("COLS");
355:             hR[k1]->GetXaxis()->SetRangeUser(0.0, 2.0);
356:             hR[k1]->SetTitle(Form("Ratio = %s / %s", legendtext_[1].c_str(), legendtext_[0].c_str()));
357:         }
358:     }
359: }
360:
361: // -----
362: // Create Labels
363: c0.cd(); // Important!
364: std::shared_ptr<TPad> tpad;
365: gra::rootstyle::TransparentPad(tpad);
366:
367: std::shared_ptr<TLabel> l1;
368: std::shared_ptr<TLabel> l2;
369: gra::rootstyle::MakeInFinland(l1, l2, 0.999);
370:
371: // -----
372:
373: // Create output directory if it does not exist
374: aux::CreateDirectory(fullpath);
375:
376: // Save pdf
377: std::string fullfile = fullpath + name_ + ".pdf";
378: c0.SaveAs(fullfile.c_str());
379:
380: if (RATIOPLOTT) {
381:     // Delete ratio histograms from memory
382:     if (h.size() > 1) {
383:         for (const auto k1 : indices(hR)) { delete hR[k1]; }
384:     }
385: }
386: return 0.0;
387:
388:
389:
390: hProcMultiplet::hProcMultiplet(const std::string &name, const std::string &labeltext, int N,
391:     double minval1, double maxval1, double minval2, double maxval2,
392:     const std::vector<std::string> &legendtext) {
393:     name_ = name;
394:     N_ = N;
395:     minval1_ = minval1;
396:     maxval1_ = maxval1;
397:     minval2_ = minval2;
398:     maxval2_ = maxval2;
399:
400:     legendtext_ = legendtext;
401:     legendposition_ = "northeast";
402:
403:     // Initialize histogram vector container size
404:     h.resize(legendtext_.size());
405:
406:     for (const auto k1 : indices(h)) {
407:         h[k1] = new TProfile(Form("%s_%lu", name.c_str(), k1), labeltext.c_str(), N, minval1, maxval1,
408:             minval2, maxval2);
409:         h[1]->Sum2(); // Error saving on
410:     }
411: }
412:
413:
414:

```



```

./src/Analysis/MMultiplet.cc      6/8
416: double hProcMultiplet::SaveFig(const std::string &fullpath, bool RATIOPLOTT) const {
417: // New colorscheme
418: std::vector<int> colors;
419: std::vector<std::shared_ptr<TCOLOR>> rootcolor;
420: rootstyle::CreateColorMap(colors, rootcolor);
421:
422: TCanvas c0("c", "c", 750, 800);
423:
424: // Upper plot will be in pad1
425: std::shared_ptr<TPad> pad1;
426:
427: if (RATIOPLOTT) {
428: pad1 = std::make_shared<TPad>("pad1", "pad1", 0, 0.3, 1.0, 1.0);
429: pad1->SetBottomMargin(0.015); // Upper and lower plot are joined
430: } else {
431: pad1 = std::make_shared<TPad>("pad1", "pad1", 0, 0.0, 1.0, 1.0);
432: }
433: // pad1->SetDirx(1); // Vertical grid
434: pad1->Draw(); // Draw the upper pad; pad1
435: pad1->cd(); // pad1 becomes the current pad
436: h[0]->SetStats(0); // No statistics on upper plot
437:
438: // Find maximum value for y-range limits
439: double MAXVAL = 0.;
440: double MINVAL = 1e32;
441: for (const auto &i : indices(h)) {
442: if (h[i]->GetMaximum() > MAXVAL) { MAXVAL = h[i]->GetMaximum(); }
443: if (h[i]->GetMinimum() < MINVAL) { MINVAL = h[i]->GetMinimum(); }
444: }
445:
446: // Loop over histograms
447: for (const auto &i : indices(h)) {
448: h[i]->SetLineColor(color[i]);
449: h[i]->SetMarkerColor(color[i]);
450: h[i]->SetMarkerStyle(20);
451: h[i]->SetMarkerSize(0.5);
452: h[i]->GetXaxis()->SetRangeUser(MINVAL, MAXVAL * 1.25);
453:
454: h[i]->Draw("L, SAME");
455: }
456:
457: // LEGEND
458: // North-East
459: double x1, x2, y1, y2 = 0.0;
460: GetLegendPosition(h.size(), x1, x2, y1, y2, legendposition_);
461: std::shared_ptr<TLegend> legend = std::make_shared<TLegend>(x1, x2, y1, y2);
462: legend->SetFillColor(0); // White background
463: legend->SetSizeStyle(0); // No box
464:
465: // Add legend entries
466: for (const auto &i : indices(h)) { legend->AddEntry(h[i], legendtext_[i].c_str(), "l"); }
467: legend->Draw();
468:
469: // GRANITTTI text
470:
471: // -----
472: std::shared_ptr<TText> l1;
473: std::shared_ptr<TText> l2;
474: gra::rootstyle::MakeFinland(l1, l2, 0.935, (RATIOPLOTT == true ? 0.03 : 0.16), 0.58);
475:
476: // -----
477: // Ratio plots
478: std::vector<TProfile> * hP(h.size(), nullptr);
479: std::vector<TH1D> * hR(h.size(), nullptr);
480: TH1D * h0 = nullptr;
481: std::shared_ptr<TPad> pad2;
482: std::shared_ptr<TLine> line;
483:
484: if (RATIOPLOTT) {
485: c0.cd();
486: pad2 = std::make_shared<TPad>("pad2", "pad2", 0, 0.05, 1, 0.3);
487: pad2->SetTopMargin(0.025);
488: pad2->SetBottomMargin(0.25);
489: pad2->SetDirx(1); // Vertical grid
490: pad2->Draw();
491: pad2->cd(); // pad2 becomes the current pad
492:
493: for (const auto &i : indices(h)) {
494: hP[i] = static_cast<TProfile> *(h[i]->Clone(Form("ratio_%.1f", i)));
495: }
496: }

```

```

./src/Analysis/MMultiplet.cc      8/8
582: }
583: return 0.;
584: }
585:
586: // namespace gra
587:
588:

```

```

./src/Analysis/MMultiplet.cc      7/8
499: // To get proper errors, we need to use TH1 instead of TProfile
500: hR[i] = hP[i]->ProjectXonX();
501:
502: if (i == 0) { // this is needed as the denominator
503: h0 = static_cast<TH1D> *(hP[i]->Clone("denominator"));
504: }
505:
506: // Set same colors as above
507: hR[i]->SetLineColor(color[i]);
508: hR[i]->SetMarkerColor(color[i]);
509: hR[i]->SetMarkerStyle(20);
510: hR[i]->SetMarkerSize(0.5);
511:
512: // -----
513: // Get ratio to 0-th histogram
514: hR[i]->Divide(h0);
515:
516: hR[i]->SetMinimum(0.0); // y-range
517: hR[i]->SetMaximum(2.0); //
518: hR[i]->SetStats(0); // no statistics box
519: hR[i]->Draw("same"); // ratio plot
520:
521: // Ratio plot (h3) settings
522: hR[i]->SetTitle(""); // remove title
523:
524: // Y axis ratio plot settings
525: hR[i]->GetYaxis()->SetTitle("Ratio");
526: hR[i]->GetYaxis()->CenterTitle();
527: hR[i]->GetYaxis()->SetNDivisions(505);
528: hR[i]->GetYaxis()->SetTitleSize(20);
529: hR[i]->GetYaxis()->SetTitleFont(43);
530: hR[i]->GetYaxis()->SetTitleOffset(1.55);
531: hR[i]->GetYaxis()->SetLabelFont(43); // in pixel (precision 3)
532: hR[i]->GetYaxis()->SetLabelSize(15);
533:
534: // X axis ratio plot settings
535: hR[i]->GetXaxis()->SetTitleSize(20);
536: hR[i]->GetXaxis()->SetTitleFont(43);
537: hR[i]->GetXaxis()->SetTitleOffset(4.);
538: hR[i]->GetXaxis()->SetLabelFont(43); // in pixel (precision 3)
539: hR[i]->GetXaxis()->SetLabelSize(15);
540:
541: }
542:
543: // Draw horizontal line
544: const double ymax = 1.0;
545: line = SetLineColor(15);
546: line->SetLineStyle(2,0);
547: line->Draw();
548:
549: // -----
550: // Remove x-axis of UPPER PLOT
551: h[0]->GetXaxis()->SetLabelOffset(999);
552: h[0]->GetXaxis()->SetLabelSize(0);
553:
554: // -----
555: } // RATIOPLOTT
556:
557: // Give y-axis title some offset to avoid overlapping with numbers
558: h[0]->GetYaxis()->SetTitleOffset(1.45);
559:
560: // Create output directory if it does not exist
561: aux::CreateDirectory(fullpath);
562:
563: // Save pdf
564: std::string fullfile = fullpath + name_ + ".pdf";
565: c0.SaveAs(fullfile.c_str());
566:
567: // Save logscale pdf
568: if (MINVAL > 0) {
569: pad1->cd()->SetLogy(); // pad2 becomes the current pad
570: fullfile = fullpath + name_ + "_logy" + ".pdf";
571: c0.SaveAs(fullfile.c_str());
572: }
573:
574: if (RATIOPLOTT) {
575: // Remove histograms
576: for (const auto &i : indices(hR)) {
577: delete hP[i];
578: delete hR[i];
579: }
580: delete h0;
581: }

```

```

./src/Analysis/MEikonal.cc        1/7
1: // Eikonal density and screening class
2: //
3: //
4: // (c) 2017-2020 Mikael Mieskolainen
5: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
6:
7: // C++
8: #include <math>
9: #include <fstream>
10: #include <future>
11: #include <iomanip>
12: #include <iostream>
13: #include <vector>
14:
15: // Own
16: #include "Granitti/MBux.h"
17: #include "Granitti/MEikonal.h"
18: #include "Granitti/MForm.h"
19: #include "Granitti/MStat.h"
20: #include "Granitti/MPDG.h"
21: #include "Granitti/MTime.h"
22:
23: // Libraries
24: #include "math.h"
25: #include "json.hpp"
26: #include "rang.hpp"
27:
28: using gra::aux::indices;
29: using gra::math::magrt;
30: using gra::math::pow;
31: using gra::math::pi;
32:
33: namespace gra {
34:
35: MEikonal::MEikonal() {}
36: MEikonal::~MEikonal() {}
37:
38: // Return total, elastic, inelastic cross sections
39: void MEikonal::GetTotXS(double &tot, double &el, double &in) const {
40: tot = sigma_tot;
41: el = sigma_el;
42: in = sigma_inel;
43: }
44:
45: // Construct density and amplitude
46: void MEikonal::B3Constructor(double s_in, const std::vector<gra::MParticle> &initialstate_in,
47: bool onlyDensity, int NumberBT, int NumberRT2) {
48: std::cout << "MEikonal::B3Constructor" << std::endl;
49:
50: // This first
51: Numerics.ReadParameters();
52:
53: // Override
54: if (NumberBT > 0) { Numerics.NumberBT = NumberBT; }
55: if (NumberRT2 > 0) { Numerics.NumberRT2 = NumberRT2; }
56:
57: // Mandelstam s and initial state
58: s = s_in;
59: INITIALSTATE = initialstate_in;
60:
61: // -----
62: // Calculate hash based on all free variables -> if something changed,
63: // calculate new densities
64:
65: // Proton density
66: {
67: std::cout << "Initializing eikonal density array" << std::endl;
68: MBT.mgts = magrt(s); // FIRST THIS
69: MBT.mts = Numerics.MIBT, Numerics.MasBT, Numerics.MasRT, Numerics.NumberBT, Numerics.logBT;
70: MBT.InitArray(); // Initialize (call last!)
71:
72: const std::string hstr = std::to_string(s) + std::to_string(INITIALSTATE[0].pdg) + "_" +
73: std::to_string(INITIALSTATE[1].pdg) + PARAM_SOFT::GetHashString() +
74: Numerics.GetHashString();
75: const unsigned long hash = gra::aux::djb2hash(hstr);
76: const std::string filename = gra::aux::GetBasePath(C) + "/eikonal/" + "MBT_" +
77: std::to_string(INITIALSTATE[0].pdg) + "-" +
78: std::to_string(INITIALSTATE[1].pdg) + "-" +
79: gra::aux::Tostring(magrt(s), 0) + "_" + std::to_string(hash);
80:
81: // Try to read pre-calculated
82: bool ok = MBT.ReadArray(filename);
83:

```



```

./src/MEIkonal.cc          6/7
416:
417: MTimer timer(true);
418: std::cout << "ArrayID:WriteArray: ";
419: unsigned int line_number = 0;
420:
421: try {
422:     for (std::size_t i = 0; i < F.size_row(); ++i) {
423:         // Write to file
424:         file << std::setprecision(15) << std::real(F[i][X]) << ", " << std::real(F[i][Y]) << ", "
425:         << std::imag(F[i][Y]) << std::endl;
426:         ++line_number;
427:     }
428: } catch (...) {
429:     throw std::invalid_argument("ArrayID:WriteArray: Error in file " + filename + " at line " +
430:     std::ito_string(line_number));
431: }
432:
433: printf("Time elapsed %0.1f sec \n", timer.ElapsedSec());
434: file.close();
435: return true;
436: }
437:
438: // Read the array from a file
439: bool ArrayID::ReadArray(const std::string filename) {
440:     std::ifstream file;
441:     file.open(filename);
442:     if (!file.is_open()) {
443:         std::string str = "ArrayID:ReadArray: Fatal IO-error with: " + filename;
444:         return false;
445:     }
446:     std::string line;
447:     unsigned int fills = 0;
448:     std::cout << "ArrayID:ReadArray: ";
449:     unsigned line_number = 0;
450:
451:     try {
452:         for (std::size_t i = 0; i < F.size_row(); ++i) {
453:             // Read every line from the stream
454:             getline(file, line);
455:
456:             std::stringstream stream(line);
457:             std::vector<double> columns(3, 0.0);
458:             std::string element;
459:
460:             // Get every line element (3 of them) separated by separator
461:             int k = 0;
462:             while (getline(stream, element, ',')) {
463:                 columns[k] = std::atof(element); // string to double
464:                 ++k;
465:             }
466:             ++fills;
467:             F[i][X] = columns[0];
468:             F[i][Y] = std::complex<double>(columns[1], columns[2]);
469:
470:             ++line_number;
471:         }
472:     } catch (...) {
473:         throw std::invalid_argument("ArrayID:ReadArray: Error in file " + filename + " at line " +
474:         std::ito_string(line_number));
475:     }
476:     file.close();
477:
478:     if (fills != 3 * (H + 1)) {
479:         std::string str = "Corrupted file: " + filename;
480:         std::cout << str << std::endl;
481:         return false;
482:     }
483:     std::cout << rang::fgy:green << "[DONE]" << rang::fgy:reset << std::endl;
484:     return true;
485: }
486:
487: // Standard 1D-linear interpolation
488:
489: std::complex<double> ArrayID::InterpolatD(double a) const {
490:     const double EPS = 1e-5;
491:     if (a < MIN) { a = MIN; } // Truncate before (possible) logarithm
492:
493:     // Logarithmic stepping or not
494:     if (islog) { a = std::ilog(a); }
495:
496:     if (a > MAX * (1 + EPS)) {
497:         printf(
498:             "ArrayID::InterpolatD(a) Input out of grid domain: "

```

```

./src/MEIkonal.cc          7/7
499:         "a = %0.3f [%0.3f, %0.3f] \n",
500:         name_c_str(), name_c_str(), a, MIN, MAX);
501:
502:     throw std::invalid_argument("InterpolatD: Out of grid domain");
503: }
504:
505: int i = std::floor((a - MIN) / STEP);
506:
507: // Boundary protection
508: if (i < 0) { i = 0; } // Int needed for this, instead of unsigned int
509: if (i == (int)N { i = N - 1; } // We got N+1 elements in F
510: // y = y0 + (x - x0) * (y1 - y0) / (x1 - x0)
511: return F[i][Y] + (a - F[i][X]) * (F[i+1][Y] - F[i][Y]) / (F[i+1][X] - F[i][X]);
512: }
513:
514: // Calculate the number of out soft Pomerons for the inelastic
515: //
516: void MEIkonal::S3NinOutPomerons() {
517:     std::cout << "MEIkonal::S3NinOutPomerons: [PROTEST]" << std::endl;
518:
519:     // Numerical integral loop over impact parameter (b,t) space
520:     const double STEP = (Numerics.MaxBT - Numerics.MinBT) / Numerics.NumberBT;
521:     P_array = std::vector<std::vector<double>>(MCUT, std::vector<double>(Numerics.NumberBT + 1, 0.0));
522:
523:     for (std::size_t j = 0; j < Numerics.NumberBT + 1; ++j) {
524:         const double bt = Numerics.MinBT + j * STEP;
525:         const double xi = std::imag(MBT.InterpolatD(bt));
526:
527:         // Poisson probabilities P_m(bt)
528:         for (std::size_t m = 1; m < MCUT; ++m) {
529:             // Poisson ansatz
530:             double P_m = std::pow(2 * XI, m) / gamma::math::factorial(m * std::exp(-2 * XI));
531:             P_array[m][j] = P_m * bt; // *bt from Jacobian (int d^2b ...
532:         }
533:     }
534:
535:     // Impact parameter <b> average probabilities
536:     P_cut.resize(MCUT, 0.0);
537:     for (std::size_t m = 0; m < MCUT; ++m) {
538:         P_cut[m] = gamma::math::Integral(P_array[m], STEP) / (Numerics.MaxBT - Numerics.MinBT);
539:         printf("P_cut[%d]=%21u = %0.5f \n", m, P_cut[m]);
540:     }
541:     std::cout << "-----" << std::endl;
542:     printf("P_cut[SUM] = %0.5f \n", std::accumulate(P_cut.begin(), P_cut.end(), 0.0));
543:
544:     // Calculate zero-truncated average
545:     double avg = 0;
546:     for (std::size_t m = 1; m < P_cut.size(); ++m) { avg += m * P_cut[m]; }
547:     printf("<b>P_cut[avg]= %0.2f \n \n", avg);
548: }
549:
550: // namespace gra
551:

```

```

./src/MParton.cc          1/4
1: // Collinear parton model 2->2 or (2->N) type phase space class
2: // for calculations without proton (remnant) considerations
3: //
4: // (c) 2017-2020 Mikael Mieskolainen
5: // Licensed under the MIT License <http://opensource.org/licenses/MIT>.
6:
7: // C++
8: #include <algorithm>
9: #include <complex>
10: #include <ostream>
11: #include <random>
12: #include <vector>
13:
14: // Own
15: #include "GraniTTL/Mkux.h"
16: #include "GraniTTL/MFkonal.h"
17: #include "GraniTTL/MFfragment.h"
18: #include "GraniTTL/MFkinematics.h"
19: #include "GraniTTL/MFmesh.h"
20: #include "GraniTTL/MFmatrix.h"
21: #include "GraniTTL/MFparton.h"
22: #include "GraniTTL/MFprocess.h"
23: #include "GraniTTL/MFspin.h"
24: #include "GraniTTL/MFscut.h"
25:
26: // Libraries
27: #include "rang.hpp"
28:
29: using gra::aux::indices;
30:
31: using gra::math::checkMC;
32: using gra::math::Pi;
33: using gra::math::abs2;
34: using gra::math::mag2;
35: using gra::math::pow2;
36: using gra::math::pow3;
37: using gra::math::pow4;
38: using gra::math::pow5;
39: using gra::math::zi;
40:
41: using gra::PDG::GeVZbar;
42:
43: namespace gra {
44: // This is needed by construction
45: MParton::MParton() { Initialize(); }
46:
47: // Constructor
48: MParton::MParton(std::string process, const std::vector<aux::OneCDM> &syntax) {
49:     Initialize();
50:     InitHistograms();
51:     SetProcess(process, syntax);
52:
53:     // Init final states
54:     MVec zerovec(0, 0, 0, 0);
55:     for (std::size_t i = 0; i < 10; ++i) { its.final_push_back(zerovec); }
56:     std::cout << "MParton: Constructor done" << std::endl;
57: }
58:
59: void MParton::Initialize() {
60:     const std::vector<std::string> supported = {"yy_LUX", "yy_DZ"};
61:     CID = " ";
62:     ProcPtr = MSUBProc(supported, CID);
63: }
64:
65: // Destructor
66: MParton::~MParton() {}
67:
68: // Initialize cut and process specific postsetup
69: void MParton::post_Constructor() {
70:     if (ProcPtr.CHANNEL == "RES") {
71:         // Here we support only single resonances
72:         if (its.RESONANCES.size() != 1) {
73:             std::string str =
74:                 "MParton: post_Constructor: only single resonance supported for this "
75:                 "process (RESPARAM.size() != 1)";
76:             throw std::invalid_argument(str);
77:         }
78:     }
79: }
80: // Set sampling boundaries
81: SetTechnicalBoundaries(gcuts, EXCITATION);
82:
83: // Initialize phase space dimension

```

```

./src/MParton.cc          2/4
84: ProcPtr.LIPSDIM = 2; // All processes
85:
86: // Not applicable here
87: // if (EXCITATION == 1) ProcPtr.LIPSDIM += 1; }
88: // if (EXCITATION == 2) ProcPtr.LIPSDIM += 2; }
89: }
90:
91: // No screening loop kinematics considered here
92: bool MParton::LoopKinematics(const std::vector<double> &pip, const std::vector<double> &pp2) {
93:     return false;
94: }
95:
96: // Return Monte Carlo integrand weight
97: double MParton::EventWeight(const std::vector<double> &randvec, AuxInData &aux) {
98:     double W = 0.0;
99:
100:     // Kinematics and cuts
101:     aux.kinematics_ok = B2RandomKin(randvec);
102:     aux.fiducial_ok = FiducialCut();
103:     aux.vetocuts_ok = VetoCuts();
104:
105:     if (aux.Valid()) {
106:         // Matrix element squared
107:         const double MatEQ = Getamp2();
108:
109:         // Calculate central system Phase Space volume
110:         double exact = 0.0;
111:         DecayWidthPS(exact);
112:         its.DW_sum_exact.Add(gca::kinematics:MCW(exact, pow2(exact), 1), aux.vegasweight);
113:
114:         // Add to the integral sum and take into account the VEGAS weight
115:         its.DW_sum.Add(its.DW, aux.vegasweight);
116:
117:         double C_space = 1.0;
118:         // We have some legs in the central system
119:         if (its.decaytree.size() != 0 && its.PS_active) {
120:             C_space = its.DW.Integral();
121:
122:             // -----
123:             // Cascade resonances phase-space
124:             C_space *= CascadePS();
125:             // -----
126:         }
127:
128:         // ** EVENT WEIGHT **
129:         W = C_space * (1.0 / S_factor) * MatEQ * B2IntegralVolume() * B2PhaseSpaceWeight() * GeVZbar /
130:         MollerFlux();
131:
132:         aux.amplitude_ok = CheckInfMan(W);
133:
134:         // As the last STEP: Histograms
135:         // If (aux.burn_in_mode) {
136:         const double totalweight = W * aux.vegasweight;
137:         FillHistograms(totalweight, its);
138:         // }
139:         return W;
140:     }
141: }
142:
143:
144: // Fiducial cuts
145: bool MParton::FiducialCuts() const { return CommonCuts(); }
146:
147: // Record event
148: bool MParton::EventRecord(BgMC3::GeVEvent &evt) { return CommonRecord(evt); }
149:
150: void MParton::PrintInit(bool silent) const {
151:     if (!silent) {
152:         PrintSetup();
153:
154:         // Construct preprint diagram
155:         std::string proton1 = "-----pF----->";
156:         std::string proton2 = "-----pF----->";
157:
158:         /*
159:         if (EXCITATION == 1) {
160:             proton1 = "-----F2-xxxxxxx*";
161:         }
162:         if (EXCITATION == 2) {
163:             proton1 = "-----F2-xxxxxxx*";
164:             proton2 = "-----F2-xxxxxxx*";
165:         }
166:         */

```


./src/MDurham.cc

4/10

```

250:   { 1, 1, 1, 1, -1, 1, 14
251:   { 1, 1, 1, 1, 1, 1, 15
252:   }
253:
254: // Alternative (semi-ud) notation for the scale choice
255: inline void MDurham::DBScaleChose(double q2, double q2_2, double Q1_2_scale,
256:   double Q2_2_scale) const {
257:   if (Param_PDF_scale == "MIN") {
258:     Q1_2_scale = std::min(q2, q1_2);
259:     Q2_2_scale = std::min(q2, q2_2);
260:   } else if (Param_PDF_scale == "MAX") {
261:     Q1_2_scale = std::max(q2, q1_2);
262:     Q2_2_scale = std::max(q2, q2_2);
263:   } else if (Param_PDF_scale == "T") {
264:     Q1_2_scale = q1_2;
265:     Q2_2_scale = q2_2;
266:   } else if (Param_PDF_scale == "EX") {
267:     Q1_2_scale = q2;
268:     Q2_2_scale = q2;
269:   } else if (Param_PDF_scale == "AVG") {
270:     Q1_2_scale = (q2 + q1_2) / 2.0;
271:     Q2_2_scale = (q2 + q2_2) / 2.0;
272:   } else {
273:     throw std::invalid_argument("MDurham::DBScaleChose: Unknown 'Durham::PDF_scale' option!");
274:   }
275: }
276:
277: // [REFERENCE: Khose, Martin, Ryskin, Journals.aps.org/prd/pdf/10.1103/PhysRevD.56.5867]
278: // [REFERENCE: Harland-Lang, Khose, Martin, Ryskin, Stirling, arxiv.org/abs/1405.0018v2]
279: //
280: // See also:
281: // [REFERENCE: Lonnblad, Ziech, arxiv.org/abs/1608.03745]
282: //
283: //
284: // Durham loop integral amplitude:
285: // A = pi^2 \int_{\text{triangle}} \frac{d^2 Q_\perp}{(2\pi)^2} M(gg \to X) [(Q_\perp^2 - (p_\perp - it)^2)^{-2} (Q_\perp^2 - p_\perp^2)^{-2}]^2
286: // * \int_{\text{triangle}} \frac{d^2 K_\perp}{(2\pi)^2} \frac{d^2 L_\perp}{(2\pi)^2} \text{Im}(2t_1 - 1) X
287: // L_\perp^2 = K_\perp^2 - Q_\perp^2 - 2 Q_\perp \cdot t_2
288: //
289: double MDurham::DQLoop(gra:LORIENTSCALAR & lts,
290:   std::vector<std::vector<std::complex<double>>> amp) {
291: // Forward proton system pt-vectors
292: const std::vector<double> pt1 = {lts.pfinal[1].Px(), lts.pfinal[1].Py()};
293: const std::vector<double> pt2 = {lts.pfinal[2].Px(), lts.pfinal[2].Py()};
294: //
295: // ** Process scale (GeV) / Sudakov suppression factor integral upper bound **
296: const double msqr = msqr(lts.s_hat / Param.alpha_s_scale);
297: //
298: //
299: //
300: // Init 2D-Simpson weight matrix (will be calculated only once, being
301: // static), C++11 handles multithreaded static initialization
302: const static Matrix<double> WSimpson = math::SimpsonWeight2D(Param.Mt_q, Param.Mphi);
303: //
304: // NOTE N + 1, init with zero!
305:
306: // Amp.size = 4 for example for the gg -> gg process [initial states summed coherently]
307: std::vector<std::vector<std::complex<double>>> f1;
308: Amp.size(0), MMatrix<std::complex<double>>(Param.Mt_q + 1, Param.Mphi + 1, 0.0);
309: //
310: // Spin-Parity
311: std::vector<std::complex<double>> JzF(4, 0.0);
312: //
313: // 2D-loop integral
314: // \int d^2 q^2 \text{vec}(qt) [...] = \int d^2 q^2 \int d^2 q^2 qt [...]
315: //
316: //
317: // Linearly discretized qt-loop, Nt!
318: for (std::size_t i = 0; i < Param.Nt_q + 1; ++i) {
319:   const double qt = Param.qt_MIN + i * Param.qt_STEP;
320:   const double q2 = pow2(qt);
321:   //
322:   // Linearly discretized phi in [0, pi], Nt!
323:   for (std::size_t j = 0; j < Param.Nphi + 1; ++j) {
324:     const double qphi = j * Param.phi_STEP;
325:     //
326:     //
327:     //
328:     // Loop vector
329:     const std::vector<double> qt = {qt * std::cos(qphi), qt * std::sin(qphi)};
330:     //
331:     // Fusing gluon pt-vectors
332:     const std::vector<double> q1 = {qt[0] - pt1[0], qt[1] - pt1[1]};

```

./src/MDurham.cc

6/10

```

416: //
417: void MDurham::DggZhi0(const gra:LORIENTSCALAR & lts,
418:   std::vector<std::vector<std::complex<double>>> amp) {
419:   const std::vector<double> q1t, const std::vector<double> q2t const {
420:     const double alpha_s = lts.GlobalSudakovPr->alpha_s_Q2(lts.s_hat / Param.alpha_s_scale);
421:   };
422:   const double gs2 = 4.0 * PI * alpha_s // coupling
423:   const double K_NLO = 1.69 // NLO correction
424:   const double C0 = 3.41474; // ch1_c(0) mass (GeV)
425:   const double NC = 0.1098; // ch1_c(0) width (GeV)
426:   const double NC = 0.0; // factors
427: //
428: // Gluon width \Gamma_{\text{gluon}}(ch1_c(0) -> gg), see references
429: std::complex<double> A = K_NLO * 4.0 * math::i * q2 * NO * msqr(0.075) / msqr(PI * M * NC);
430: //
431: // Apply Breit-Wigner propagator shape delta-function
432: A *= gra:form:deltaWamp(lts.s_hat, MO, MO);
433: //
434: // Initial state gluon helicity combinations,
435: // ** and - give contribution for 0+ state (see DheProj function)
436: std::vector<std::complex<double>> is(4, 0.0);
437: is[0] = A;
438: is[1] = 0;
439: is[2] = 0;
440: is[3] = A;
441: //
442: // No polarization for the final state to loop over, only one index
443: Amp[0] = is;
444: //
445: //
446: // =====
447: // gg -> gg tree-level helicity amplitudes
448: //
449: //
450: // Basic result: d\hat{\sigma}(s)/ds = 9/4 |pi| alpha_s^2 / E_T^4
451: //
452: // In pure gluon amplitudes: when gluon helicities are the same,
453: // or at most one is different from the rest, vanish for any n > 4,
454: // where n is the total number of gluons (inout)
455: //
456: // [REFERENCE: Dixon, arxiv.org/abs/1310.5353v1]
457: //
458: // Note that incoming gluons (s, b) are in a color singlet state,
459: // i.e., a = b or 'delta'(ab) Tr [ color algebra. ] has been applied
460: //
461: void MDurham::DggZg(const gra:LORIENTSCALAR & lts,
462:   std::vector<std::vector<std::complex<double>>> amp) {
463:   const double alpha_s = lts.GlobalSudakovPr->alpha_s_Q2(lts.s_hat / Param.alpha_s_scale);
464: //
465: // Vertex factor coupling, g^2 = 4 pi alpha_s
466: double norm = 4.0 * PI * alpha_s;
467: //
468: // Color part
469: const double NC = 3.0;
470: norm *= NC / msqr(NC * NC - 1);
471: norm *= 2.0;
472: //
473: // Helicity phase
474: const double phi = lts.decaytree[0].p4.Phi();
475: const std::complex<double> negphase = std::exp(-i * phi);
476: const std::complex<double> posphase = std::exp(i * phi);
477: //
478: const double aux0 = pow2(lts.u_hat) / (lts.u_hat + lts.t_hat);
479: const double aux2 = lts.u_hat / (lts.t_hat);
480: //
481: // Use natural binary order / Madgraph order here
482: std::vector<std::vector<std::complex<double>>> dualamp(16, 0.0);
483: //
484: // Loop over final state gluon pair helicity combinations
485: for (std::size_t h = 0; h < f2z_size(); ++h) {
486:   std::vector<std::complex<double>> f2z(4, 0.0);
487:   if (f2z[h] == PP || f2z[h] == MM) {
488:     is[MM] = (f2z[h] == PP) ? 0 : norm * aux0; // == ++ or --
489:     is[PP] = 0; // == ++ or --
490:     is[PM] = 0; // == ++ or --
491:     is[MP] = (f2z[h] == PP) ? norm * aux0 : 0; // == ++ or --
492:   } else if (f2z[h] == FF || f2z[h] == MM) {
493:     is[MM] = 0; // == ++ or --
494:     is[PM] = negphase * norm * ((f2z[h] == PM) ? 1 : aux2); // == ++ or --
495:     is[MP] = posphase * norm * ((f2z[h] == PM) ? 1 : aux2); // == ++ or --
496:     is[PP] = 0; // == ++ or --
497:   }
498:   Amp[h] = is;

```

./src/MDurham.cc

5/10

```

333: const std::vector<double> q2 = {qt[0] + pt2[0], qt[1] + pt2[1]};
334:
335: const double q1_2 = math::pow2(q1);
336: const double q2_2 = math::pow2(q2);
337:
338: // Get fusing gluon spin-parity (L,S,P) components
339: // [q1, q2] -> [0^+, 0^+, -2^+, -2^-]
340: DHelicity(q1, q2, JzF);
341:
342: // ** Durham scale choice **
343: double Q1_2_scale = 0.0;
344: double Q2_2_scale = 0.0;
345: DBScaleChose(q2, q1_2, q2_2, Q1_2_scale, Q2_2_scale);
346:
347: // Minimum scale cutoff
348: if (Q1_2_scale < Param.qt2_MIN || Q2_2_scale < Param.qt2_MIN) { continue; }
349:
350: //
351: // Get amplitude level pdfs
352: const double fg_1 = lts.GlobalSudakovPr->fg_xQ2M(lts.x1, Q1_2_scale, MO);
353: const double fg_2 = lts.GlobalSudakovPr->fg_xQ2M(lts.x2, Q2_2_scale, MO);
354:
355: // Amplitude weight:
356: // * pi^2 * see original KMR papers: [alpha_s CF -> pi x f_g] for fg_1 and fg_2
357: // * 2 * factor from initial state boson-statistics, check it!
358: // * q : jacobian of 2q -> dphi dgt q
359: std::complex<double> weight = fg_1 * fg_2 / (q2 * q1_2 * q2_2);
360: weight *= math::PIPI * 2.0 * qt;
361:
362: // Loop over (outgoing) helicity combinations.
363: // Amp[h] contains initial state gluon helicity combinations --, +, -, +, +
364: // Here we sum coherently
365: for (std::size_t h = 0; h < f.size(); ++h) { f[h][1][j] = weight * DHeProj[Amp[h], JzF]; }
366:
367: // phi-loop
368: // |q,t|-loop
369: //
370: // Evaluate the total numerical integral for each helicity amplitude
371: std::vector<std::complex<double>> sum(f.size(), 0.0);
372: for (std::size_t h = 0; h < f.size(); ++h) {
373:   sum[h] = math::SimpsonIntegrand2D(f[h], WSimpson, Param.qt_STEP, Param.phi_STEP);
374: }
375:
376: // ****Make sure it is of right size****
377: lts.hamp.resize(f.size());
378:
379: // Outgoing helicity combinations
380: double A2 = 0.0;
381: for (std::size_t h = 0; h < f.size(); ++h) {
382:   lts.hamp[h] = sum[h];
383: }
384: // Apply proton form factors
385: lts.hamp[h] *= lts.excite1 ? gra:form:SBFFIN(lts.l1, lts.pfinal[1].M2()) : gra:form:SBF(lts.l1);
386: lts.hamp[h] *= lts.excite2 ? gra:form:SBFFIN(lts.l2, lts.pfinal[2].M2()) : gra:form:SBF(lts.l2);
387:
388: // Apply phase space factors
389: lts.hamp[h] *= msqr(16.0 * math::PIPI);
390: lts.hamp[h] *= msqr(16.0 * math::PIPI);
391: lts.hamp[h] *= msqr(lts.s / lts.s_hat);
392:
393: // For the total amplitude squared (for no eikonal screening applied)
394: A2 += math::abs2(lts.hamp[h]);
395: }
396: //
397: // Initial state helicity average 1/4 not needed here
398:
399: //
400: // Amplitude cutoff (hard perturbative limit)
401: if (gra:math:msqr(lts.m2) < 2.0 * lts.A2) {
402:   A2 = 0.0;
403:   lts.hamp = std::vector<std::complex<double>>(f.size(), 0.0);
404: }
405: //
406: //
407: return A2;
408: }
409: //
410: //
411: // =====
412: // Exclusive \chi_{1,c}(0) sub-amplitude
413: //
414: // [REFERENCE: Khose, Martin, Ryskin, Stirling, arxiv.org/abs/hp/040328]
415: // [REFERENCE: Pasechnik, Sczurek, Teruya, arxiv.org/abs/0709.0937v1]

```

./src/MDurham.cc

7/10

```

499: for (std::size_t k = 0; k < 4; ++k) { dualamp[4 * k + h] = is[k]; }
500:
501: //
502: //
503: //
504: const bool DEBUG = false;
505: //
506: // TEST amplitude squared WITH MADGRAPH (UNDER IMPLEMENTATION)
507: double madsum = 0.0;
508: double thinum = 0.0;
509: for (std::size_t i = 0; i < 16; ++i) {
510:   print("i=%d: |mad|^2 = %0.3f, |amp|^2 = %0.3e \n", i, math::abs2(lts.hamp[i]),
511:     math::abs2(dualamp[i]));
512:   madsum += math::abs2(lts.hamp[i]);
513:   thinum += math::abs2(dualamp[i]);
514: }
515: Asum += thinum / madsum;
516: print("i=0: Asum = %1.0f\n", Asum);
517: print("i=0: thinum/madsum = %0.10f \n", thinum / madsum);
518: //
519: //
520: // =====
521: //
522: // gg -> meson pair (gg -> q\bar{q} q\bar{q})
523: //
524: // [REFERENCE: Harland-Lang, Khose, Ryskin, Stirling, arxiv.org/pdf/1105.626.pdf]
525: //
526: // Meson wave-function (without scale Q^2-dependence)
527: // x the longitudinal momentum fraction of a parton within the meson
528: // m the meson decay constant
529: //
530: // Normalization: \int_0^1 dx \phi(x) / \int_0^1 dx \phi(x) = 1
531: //
532: //
533: double MDurham::phi_C0(double x, double FM const {
534:   return 5.0 * std::sqrt(3.0) * FM * x * (1.0 - x) * pow2(1.0 - 2.0 * x);
535: // return 6.0*std::sqrt(3.0)*FM * x*(1.0 - x);
536: //
537: //
538: // For tabulating the meson wave function
539: //
540: // [REFERENCE: Takizawa,
541: // http://www.nuclth.phys.sci.osaka-u.ac.jp/~jp-usw/takizawa.pdf]
542: //
543: // Here, one could perhaps update the phenomenology (check more recent papers).
544: //
545: // Three charge neutral state are in the SU(3)_F quark model nonet
546: // (octet+singlet):
547: pi0, eta8, eta0
548: //
549: // Rotation (mixing) to physical basis via:
550: //
551: // [eta] = [cos theta -sin theta] [eta8]
552: // [eta'] = [sin theta cos theta] [eta0]
553: //
554: //
555: // Decay constants:
556: //
557: // f_\pi(eta8) = f_\pi * cos theta
558: // f_\pi(eta') = f_\pi * sin theta
559: //
560: // f_\pi(eta0) = -f_\pi * sin theta
561: // f_\pi(eta') = f_\pi * cos theta
562: //
563: std::vector<double> MDurham::EvalPhi(int N, int pdg) const {
564: // Meson decay constants
565:   double FM = 0.0;
566:   static const std::vector<double> supported = {111, 211, 321, 311}; // pi0, pi+, K+, K0
567:   if (std::find(supported.begin(), supported.end(), pdg) == supported.end()) {
568:     FM = PDG::FM_meson.at(pdg);
569:   } else {
570:     FM = PDG::FM_meson.at(111); // else, take pi0
571:   }
572: //
573: //
574: //
575: // Mixing angle
576: static const double THETA_eta = math::Deg2Rad(-15.4);
577: const double cosh = std::cos(THETA_eta);
578: const double sinh = std::sin(THETA_eta);
579: //
580: // Evaluate
581: std::vector<double> f(N + 1);

```

```

./src/MDurham.cc                               8/10
582: const double STEP = 1.0 / N;
583: for (const auto & i : aux::indices(f)) {
584:     const double s = 1 + STEP;
585:     if (pdg == 221) { // eta via mixing
586:         f[i] = cosh * phi_CZ(x, fM_eta8) - sinh * phi_CZ(x, fM_eta0);
587:     } else if (pdg == 331) { // eta prime via mixing
588:         f[i] = sinh * phi_CZ(x, fM_eta8) + cosh * phi_CZ(x, fM_eta0);
589:     } else {
590:         f[i] = phi_CZ(x, fM);
591:     }
592: }
593: return f;
594: }
595:
596: // Meson pair amplitude
597: //
598: void MDurham::Dgg2Mbar(const gra::LORENTZSCALAR & lts,
599:                        std::vector<std::vector<std::complex<double>>> &amp;mp) {
600:     // Wave function x-discretization
601:     // [Easy speed & accuracy improvement: Check lambda-functions below,
602:     // they could be pre-calculated and interpolated as a function of cosheta]
603:     const int Nx = 96;
604:     const double STEPx = 1.0 / Nx;
605:
606:     // Init 2D-Simpson weight matrix (will be calculated only once, being
607:     // static), C++11 handles multithreaded static initialization
608:     const static MMatrix<double> WSimpson = math::Simpson3Weight2D(Nx, Nx);
609:
610:     const double delta_AB = 8; // Sum over |delta|_AB [gluons in with the same color]
611:     const double NC = 3; // Three colors
612:     const double CF = (pow2(NC) - 1.0) / (NC * 2.0); // SU(3) algebra
613:
614:     const double alpha_s = lts.GlobalSudakovPtr->alpha_s_Q2(lts.s_hat / Param.alphas_scale);
615:
616:     const double shat = lts.s_hat;
617:     const double m = lts.decaytree[0].p4.M();
618:     const double beta = Kinematic::Beta(pow2(m), shat);
619:     const double cosheta = (1.0 + 2.0 * lts.t_hat / lts.s_hat - 2.0 * pow2(m) / lts.s_hat) / beta;
620:     const double cosheta2 = pow2(cosheta);
621:
622:     // -----
623:     // ** Hard angular cut-off **
624:     // some sub-amplitudes are singular when |cosheta| -> 1
625:
626:     if (std::abs(cosheta) > Param.MAXCOS) {
627:         Amp[0] = std::vector<std::complex<double>>(4, 0.0); // Return zero
628:         return;
629:     }
630:
631:     // Helicity phase
632:     const double phi = lts.decaytree[0].p4.Phi();
633:     const std::complex<double> posphase = std::exp(2.0 * i * phi);
634:     const std::complex<double> negphase = std::exp(-2.0 * i * phi);
635:
636:     // -----
637:     // Mesons scalar flavor octet (non-singlet) amplitude:
638:     // |psi0>|psi0>, |psi+>|psi->, |K+>>, |S0>|S0>
639:
640:     auto T_SFO_PM = [4](double x, double y) {
641:         const double a = (1.0 - x) * (1.0 - y) + x * y; // +
642:         const double b = (1.0 - x) * (1.0 - y) - x * y; // -
643:
644:         return 1.0 / (x * y * (1.0 - x) * (1.0 - y)) * (x * (1.0 - x) + y * (1.0 - y)) /
645:             (pow2(a) - pow2(b) * cosheta2) * (NC / 2.0) * (cosheta2 - 2.0 * CF / NC * a);
646:     };
647:
648:     // -----
649:     // SU(3)_F scalar Flavor-singlet amplitude:
650:     // |eta'>|eta'>, |eta>|eta'>, |eta>|eta'>, |eta>|eta'>
651:     // -----
652:     auto T_SFS_PP = [4](double x, double y) {
653:         return 1.0 / (x * y * (1.0 - x) * (1.0 - y)) * (1.0 + cosheta2) / pow2(1.0 - cosheta2);
654:     };
655:
656:     auto T_SFS_PM = [4](double x, double y) {
657:         // T_+ = T_+
658:         // T_- = T_-
659:         return 1.0 / (x * y * (1.0 - x) * (1.0 - y)) * (1.0 + cosheta2) / pow2(1.0 - cosheta2);
660:     };
661:
662:     // T_+ = T_+
663:     auto T_SFS_PM = [4](double x, double y) {

```

```

./src/MDurham.cc                               10/10
748: }
749:
750: // namespace gra
751:

```

```

./src/MDurham.cc                               9/10
665:
666:     return 1.0 / (x * y * (1.0 - x) * (1.0 - y)) * (1.0 + 3.0 * cosheta2) /
667:         (2.0 * pow2(1.0 - cosheta2));
668: };
669: // -----
670: // p10, p1+, K+, K0
671: static const std::vector<int> SFS_PDG = {111, 211, 321, 311};
672:
673: // eta, eta'
674: static const std::vector<int> SFS_PDG = {221, 331};
675:
676: static const int pdg0 = std::abs(lts.decaytree[0].p.pdg);
677: static const int pdg1 = std::abs(lts.decaytree[1].p.pdg);
678:
679: // Evaluate only once the meson wave functions
680: static const std::vector<double> wfphi0 = EvalPhi(Nx, pdg0);
681: static const std::vector<double> wfphi1 = EvalPhi(Nx, pdg1);
682:
683: // -----
684: const double CUTOFF = 1e-15; // To avoid singularity at x = 0, x = 1
685: static const std::vector<double> xval = math::linspace(CUTOFF, 1.0 - CUTOFF, Nx + 1);
686:
687: // -----
688: // M(int_0^1 dx dy |psi_N(x) |psi_Lbar(M)(y) T.(lambda/lambda)
689: // |N_y, lbar(s), theta)
690:
691: // Normalization factor
692: const double norm = (delta_AB / NC) * (64.0 * math::PIPI * pow2(alpha_s)) / shat;
693:
694: // Scalar Flavor octet
695: if ((std::find(SFS_PDG.begin(), SFS_PDG.end(), pdg0) != SFS_PDG.end()) ||
696:     std::vector<std::complex<double>> is(4)) {
697:     MMatrix<double> f_PP(Nx + 1, Nx + 1, 0.0);
698:     for (std::size_t i = 0; i < Nx + 1; ++i) {
699:         for (std::size_t j = 0; j < Nx + 1; ++j) {
700:             f_PP[i][j] = wfphi0[i] * wfphi1[j] * T_SFO_PM(xval[i], xval[j]);
701:         }
702:     }
703:
704:     is[PP] = 0.0;
705:     is[MM] = 0.0;
706:     is[PM] = norm * math::Simpson3Integral2D(f_PP, WSimpson, STEPx, STEPx) * posphase;
707:     is[MP] = is[PM] * negphase;
708:
709:     Amp[0] = is;
710: }
711:
712: // Scalar flavor singlet
713: else if ((std::find(SFS_PDG.begin(), SFS_PDG.end(), pdg0) != SFS_PDG.end()) ||
714:         std::vector<std::complex<double>> is(4)) {
715:     MMatrix<double> f_PP(Nx + 1, Nx + 1, 0.0);
716:     MMatrix<double> f_PM(Nx + 1, Nx + 1, 0.0);
717:
718:     for (std::size_t i = 0; i < Nx + 1; ++i) {
719:         for (std::size_t j = 0; j < Nx + 1; ++j) {
720:             f_PP[i][j] = wfphi0[i] * wfphi1[j] * T_SFS_PP(xval[i], xval[j]);
721:             f_PM[i][j] = wfphi0[i] * wfphi1[j] * T_SFS_PM(xval[i], xval[j]);
722:         }
723:     }
724:
725:     is[PP] = norm * math::Simpson3Integral2D(f_PP, WSimpson, STEPx, STEPx);
726:     is[PM] = is[PP];
727:     is[MP] = norm * math::Simpson3Integral2D(f_PM, WSimpson, STEPx, STEPx) * posphase;
728:     is[MP] = is[PM] * negphase;
729:
730:     Amp[0] = is;
731: } else {
732:     throw std::invalid_argument("MDurham::Dgg2Mbar: Unsupported meson: " +
733:                                 std::to_string(lts.decaytree[0].p.pdg));
734: }
735:
736: // -----
737: // gg -> qqbar tree-level helicity amplitudes
738: // -----
739: // -----
740: // -----
741: // -----
742: // -----
743: // -----
744: // -----
745: void MDurham::Dgg2qbar(const gra::LORENTZSCALAR & lts,
746:                        std::vector<std::vector<std::complex<double>>> &amp;mp) {
747:     throw std::invalid_argument("MDurham::DusHamQCD: qqbar amplitude in the next version");

```


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