

Anisotropic Strain Enhanced Hydrogen Solubility in bcc Metals: The Independence on the Sign of Strain

Hong-Bo Zhou, Shuo Jin, Ying Zhang, and Guang-Hong Lu*
Department of Physics, Beihang University, Beijing 100191, China

Feng Liu

Department of Materials Science and Engineering, University of Utah, Salt Lake City, Utah 84112, USA
 (Received 28 March 2012; published 27 September 2012)

When an impurity is doped in a solid, it inevitably induces a local stress, tending to expand or contract the lattice. Consequently, strain can be applied to change the solubility of impurity in a solid. Generally, the solubility responds to strain “monotonically,” increasing (decreasing) with the tensile (compressive) strain if the impurity induces a compressive stress or vice versa. Using first-principles calculations, however, we discovered that the H solubility can be enhanced by anisotropic strain in some bcc metals, almost independent of the sign of strain. This anomalous behavior is found to be caused by a continuous change of H location induced by anisotropic strain. Our finding suggests a cascading effect of H bubble formation in bcc metals: the H solution leads to H bubble formation that induces anisotropic strain that in turn enhances H solubility to further facilitate bubble growth.

DOI: [10.1103/PhysRevLett.109.135502](https://doi.org/10.1103/PhysRevLett.109.135502)

PACS numbers: 62.20.-x

The solubility of impurity in solids is not only of scientific interest but also technological importance [1]. Strain engineering has been recognized to provide an effective way to enhance the solubility of impurity in solids, such as doping of impurity in semiconductors. In general, the impurity formation energy is a linear monotonic function of strain following elasticity theory, regardless of whether it is a hydrostatic or biaxial strain [2,3]. Thus, one can expect that, when a “large” impurity induces a compressive lattice stress, tending to expand the lattice, its solubility can only be enhanced by applying a tensile strain but decreased by a compressive strain, and the reverse is true for a “small” impurity. In this Letter, we report the discovery of an exceptional case that the H solubility in some bcc metals is found to be enhanced by applying both a tensile and compressive anisotropic strain, independent of the sign of strain. This surprising counterintuitive behavior is shown to be caused by an unusual strain induced continuous change of H minimum-energy location from one interstitial site to another.

Hydrogen is the most common impurity in metals. Understanding the H solubility in metals has broad and significant technological implications. Hydrogen retention assists vacancy formation in many metals that degrades their structural properties, a phenomenon known as H embrittlement [4–6]. Metal hydrides have been intensively studied for H storage [7,8]. In a fusion reactor, metals are used as the plasma facing material (PFM) that is exposed to extremely high fluxes of H isotope (deuterium-tritium) ions [9]. The accumulation of H in metals leads to formation of voids, bubbles, and blisters [10]. In this regard, our finding suggests a cascading effect of H bubble

formation in bcc metals, which has significant implications in using bcc metals as PFMs.

We have investigated the effect of strain on H solubility in bcc metals of W, Mo, Fe, and Cr, using first-principles calculation. Our calculations were performed using the pseudopotential plane-wave method as implemented in the VASP code [11,12] based on the density functional theory (DFT). We used the generalized gradient approximation of Perdew and Wang [13] and projected augmented wave potentials [14], with a plane-wave energy cutoff of 350 eV. The bcc supercell of 54 atoms was used, and its Brillouin zone was sampled with $(5 \times 5 \times 5)$ k points by the Monkhorst-Pack scheme [15]. The energy minimization is continued until the forces on all the atoms are converged to less than 10^{-3} eV \AA^{-1} . The H solution energy ($E_{\epsilon}^{\text{sol}}$) in the strained metal is calculated as $E_{\epsilon}^{\text{sol}} = E_{\epsilon, \text{H}} - E_{\epsilon} - \frac{1}{2}E_{\text{H}_2}$, where $E_{\epsilon, \text{H}}$ and E_{ϵ} are total energies with and without H in the supercell, respectively, ϵ is the applied strain, and E_{H_2} is the calculated total energy of an H_2 molecule (-6.76 eV). This corresponds to the binding energy of ~ 4.53 eV for an H_2 molecule, consistent with the experiment [16]. Since the H atom is a light-mass particle and therefore has a high vibrational energy compared to the relatively slowly moving heavy metal atoms, the zero-point energy of H has been taken into account by summing up the zero-point vibrational energies of H’s normal modes. For lattice relaxation, under hydrostatic strain, the lattice parameter is fixed at given strain values with atomic positions relaxed; under anisotropic biaxial strain, the in-plane x and y lattice parameters are fixed at given strain values, with the z lattice parameter relaxed and optimized with all atomic coordinates.

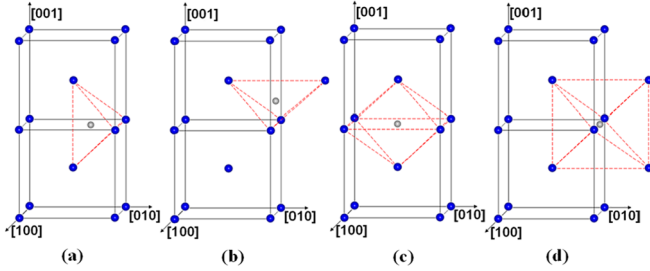


FIG. 1 (color online). The interstitial sites in a bcc lattice. (a) TIS-I, (b) TIS-II, (c) OIS-I, and (d) OIS-II. The larger blue spheres represent the metal atoms; the smaller gray spheres represent the H atoms.

When an H atom is introduced into a bcc metal lattice, it prefers to occupy an interstitial site for its small size relative to the host atom. This is confirmed by our calculations showing that in all the metals considered the H solution energy is much lower at the interstitial sites than at the substitutional site. Thus, only the interstitial sites will be considered below. There are two kinds of interstitial sites in a bcc lattice, the tetrahedral interstitial site (TIS) and the octahedral interstitial site (OIS), as shown in Fig. 1. The TIS has four nearest neighbors at $0.559a$, where a is the lattice constant, and the OIS has six nearest neighbors, with two of them located at $0.5a$ and four at $0.707a$. The calculated H solution energies in the equilibrium lattice are listed in Table I. In all four metals, H prefers TISs without strain, in good agreement with previous results [17–19].

A standard approach to understand the effect of strain on a point defect in a solid, such as H in metal, is by the concept of the force dipole tensor [2]. Here, we have calculated the H-induced lattice stress tensor, which is equivalent to the force dipole tensor, differing only by a volume normalization factor. Because of tetragonal lattice symmetry, at both the TIS and the OIS, the stress is anisotropic. However, the magnitude of stress anisotropy is small at the TIS but very large at the OIS, as shown in Table I.

The H behavior under strain is then dictated by the H-induced lattice stress tensor, i.e., the force dipole tensor, as has been demonstrated before [2,3]. The defect induced lattice stress generally has two contributions [2,3,20]: the atomic size effect (bond deformation) and the electronic band effect (Fermi level shift), termed as quantum electronic stress [20]. Typically, one assumes the stress tensor is independent of strain. Following linear elasticity theory, the H solution energy under strain can be calculated as

$$E_{\varepsilon}^{\text{sol}} = E_{\varepsilon=0}^{\text{sol}} + V(\sigma_{[100]}\varepsilon_{[100]} + \sigma_{[010]}\varepsilon_{[010]} + \sigma_{[001]}\varepsilon_{[001]}), \quad (1)$$

where $E_{\varepsilon=0}^{\text{sol}}$ is the H solution energy without strain, V is the volume of the supercell at equilibrium, σ is the H-induced lattice stress in the equilibrium lattice, and ε is the strain. Note that this equation applies to a special case of the force

TABLE I. H solution energies (in eV) and H-induced lattice stress (in GPa) in bcc metals.

		W	Mo	Fe	Cr	
Solution energy	TIS	1.08	0.78	0.34	0.79	
	OIS	1.48	1.03	0.44	0.92	
σ	TIS-I	$XX = ZZ$	-0.99	-0.97	-1.25	-1.42
		YY	-0.97	-0.86	-1.10	-1.18
	TIS-II	$XX = YY$	-0.99	-0.97	-1.25	-1.42
		ZZ	-0.97	-0.86	-1.10	-1.18
	OIS-I	$XX = YY$	-0.45	-0.39	-0.91	-0.57
		ZZ	-2.21	-2.06	-2.14	-2.88
OIS-II	XX	-2.21	-2.06	-2.14	-2.88	
	$YY = ZZ$	-0.45	-0.39	-0.91	-0.57	

dipole tensor (as well as the strain tensor) that is diagonal in the chosen cubic crystal system. In applying an isotropic (hydrostatic) strain, $\varepsilon_{[100]} = \varepsilon_{[010]} = \varepsilon_{[001]}$, one expects a linear dependence of the H solution on strain with a slope equals to the average stress, $\bar{\sigma} = 1/3(\sigma_{[100]} + \sigma_{[010]} + \sigma_{[001]})$. This is indeed confirmed by our DFT calculations, which show very good agreement between the prediction from Eq. (1) (solid lines), using the data of $E_{\varepsilon=0}^{\text{sol}}$ and σ in Table I, and the direct calculations of the H solution energy under different isotropic strains (data points), as shown in Fig. 2, except at large strains where nonlinear effects become important. Because the average stress is compressive (negative by convention; see Table I), the H solution energy decreases “monotonically” in all four metals with the increasing tensile strain but increases with the increasing compressive strain. Also, the TIS remains to be the preferred site under all strains. So, overall we may view the results under hydrostatic strain in Fig. 2 as the “normal behavior” in accordance with our common expectations [2,3].

Next, we consider the case of applying anisotropic (biaxial) strain. When biaxial strain is applied in the x - y plane along the $[100]$ and $[010]$ directions (see Fig. 1), the lattice parameter along the z direction ($[001]$) is fully relaxed according to the Poisson ratio (ν). So, the applied strain tensor is $\varepsilon_{[100]} = \varepsilon_{[010]} = -\varepsilon_{[001]}/\nu$, with lattice contraction (expansion) in the x - y plane and expansion (contraction) along the z direction. From our calculations, we obtain the Poisson ratios 0.29, 0.28, 0.31, and 0.21 for W, Mo, Fe, and Cr, respectively. Following Eq. (1), the H solution energy should still depend linearly on the strain with a slope of $\bar{\sigma} = 1/3(\sigma_{[100]} + \sigma_{[010]} - \nu\sigma_{[001]})$. The calculation results under biaxial strain are presented in Fig. 3. Because the biaxial strain breaks the lattice symmetry, causing a bcc to body-centered tetragonal (bct) structural transformation, here we have to consider two different TIS and OIS configurations, denoted as TIS-I, TIS-II, OIS-I, and OIS-II, respectively, in Fig. 1. The most important observation here is that the linear dependence is followed at all four locations only for very small strain but not for large strain. Also, we note that the results at the

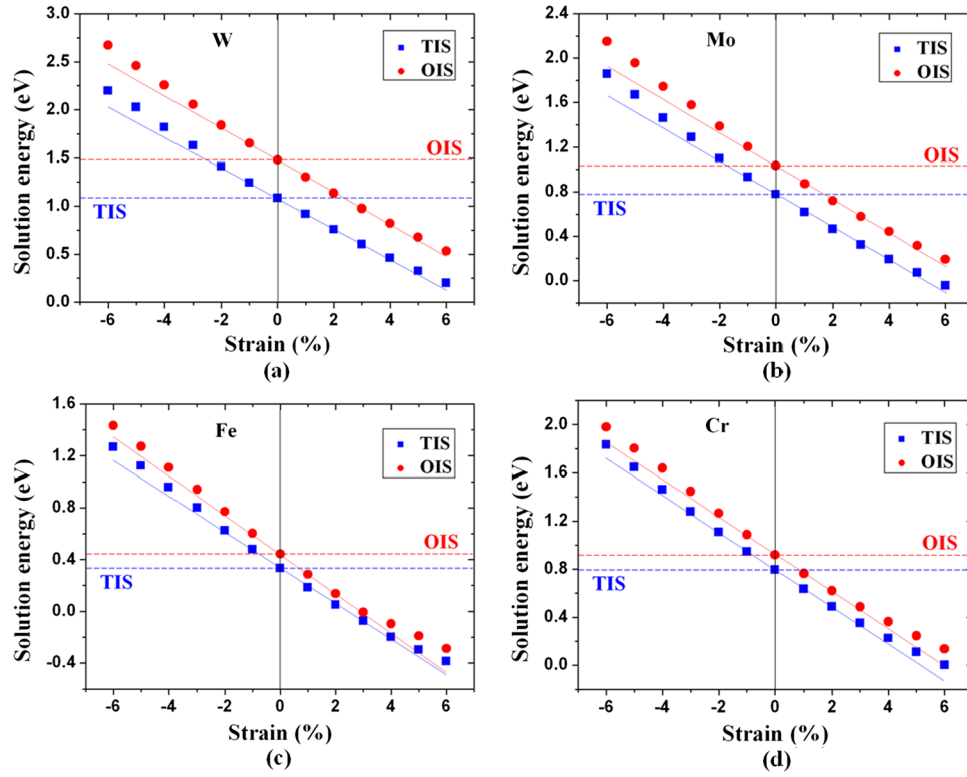


FIG. 2 (color online). The H solution energy at the TIS and the OIS as a function of applied hydrostatic strain. (a) W, (b) Mo, (c) Fe, and (d) Cr. The solid lines are the linear elasticity model predictions of Eq. (1) using the DFT data of solution energies and stresses without strain in Table I.

OIS-I have an opposite trend from those at the other three sites. This is caused by the large stress anisotropy and the Poisson effect. When compressive biaxial strains are applied in the x - y plane, a tensile uniaxial strain is applied in the z direction by the Poisson ratio. Note that the stress tensor components switch by symmetry between σ_{yy} and σ_{zz} at the TIS-I and the TIS-II and between σ_{xx} and σ_{zz} at the OIS-I and the OIS-II, respectively, as shown in Table I. Consequently, the z component stress is much larger than those of the x and y components at the OIS-I but not so at the OIS-II (see Table I), so that the z direction effect dominates at the OIS-I to show a net tensile strain effect when compressive strain is applied in the x - y plane. (The difference between TIS-I and TIS-II is too small to show an effect.)

Overall, the results at the TIS-II and the two OISs are relatively normal, showing a monotonic strain dependence of solution energy. Surprisingly, however, the TIS-I solution energy (black squares in Fig. 3) shows nonmonotonic dependence on strain, especially decreasing with the increase of both signs of strain. Thus, for the TIS-I site, the standard picture of Eq. (1) with the stress (force dipole) tensor [2,3] fails. Most importantly, since the TIS-I remains to be the preferred site with lower energy than the TIS-II and the two OISs under strain (see Fig. 3), this unusual TIS-I behavior may dominate the overall strain dependence of H solubility in bcc metals, which cannot be simply understood by the standard picture. Practically, it

implies that the H solubility can be enhanced by applying either a tensile or a compressive anisotropic strain.

To show the above more clearly, we plot in Fig. 4 the change of the overall H solution energy in the anisotropically strained metal relative to that in the strain-free metal ($E^{\text{dif}} = E_{\varepsilon, \text{H}}^{\text{sol}} - E_{\text{TIS, H}}^{\text{sol}}$). Except for a small increase (< 0.05 eV) in a limited range of small compressive strain, the H solution energy decreases over a wide range of strain signs and values. This means that the H solubility can be effectively enhanced in all the strained bcc metals, almost independent of the sign of strain.

We now discuss the physical origin underlying the anomalous sign independence of H solubility on strain in bcc metals. It is important to realize that the common expectation of a monotonic dependence of H solution energy on strain has a *prerequisite* condition that the stress tensor (or force dipole tensor) [2,3] is independent of strain to guarantee the applicability of Eq. (1). This implies that H stays at the same configuration under strain; i.e., neither the site occupation nor the location or the lattice displacement field induced by H change appreciably under strain. This condition is always satisfied by symmetry constraint when a hydrostatic strain is applied, as shown in Fig. 1. When a biaxial strain is applied, however, we found that this condition can break down due to the broken symmetry for H at the TIS-I. Figure 5 shows the continuous change of the H minimum-energy location in W, initially placed at the TIS-I, induced by anisotropic strain discovered from

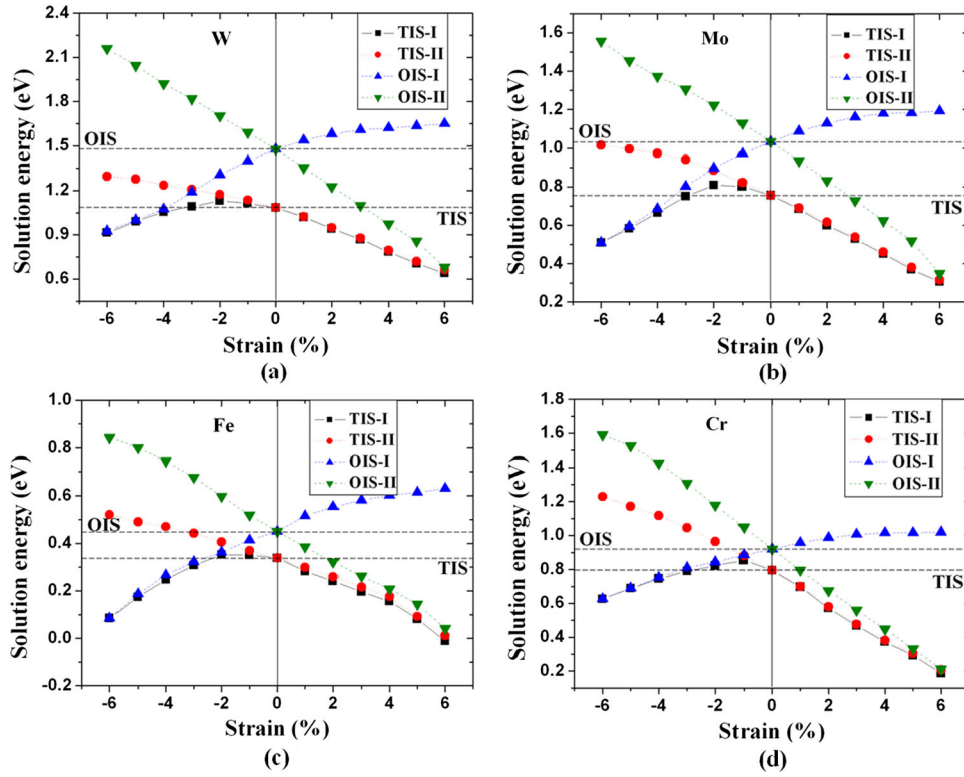


FIG. 3 (color online). The solution energy of H in bcc metals (W, Mo, Fe, and Cr) under the biaxial strain applied along the [100] and [010] directions.

our DFT calculations. Under compressive biaxial strain, the H moves towards the OIS-I site and reaches the OIS-I at about -4% strain, as shown in Fig. 3(a), beyond which the energy curves of the TIS-I and the OIS-I merge together. Under tensile biaxial strain, the H moves towards the OIS-II site and reaches the OIS-II at about $+6\%$ strain, as shown in Fig. 3(a), beyond which the energy curves of the TIS-I and the OIS-II merge together. Therefore, it is such an unusual strain induced change of the H position that is responsible for the observed anomalous nonmonotonic dependence of the H solution energy on strain, leading to the end result of energy decreasing under both signs of strain. Apparently, such a strain induced change of microscopic atomic configuration cannot be captured by the macroscopic continuum elasticity theory. Similar behaviors of H under biaxial strain are also found in other bcc metals of Mo, Fe, and Cr, which leads to anomalous nonmonotonic strain dependence of the H solution energy in all four bcc metals studied, as shown in Fig. 4.

We note that previous work has shown that H in bcc metals can switch from TIS to OIS as a function of isotopic mass [21]. Here, we discover a different mechanism for H switching from TIS to OIS as a function of anisotropic strain. The isotope effect changes the relative stability of the TIS versus the OIS and hence the relative H population between the two sites. Instead, the anisotropic strain moves the minimum-energy H position gradually from the TIS to the OIS with the increasing strain; at the intermediate strain, the minimum-energy position sits in between the

TIS and the OIS. A simple nearest-neighbor H-metal interaction model [22] may suggest the change of relative stability of the TIS vs the OIS under anisotropic strain but cannot resolve such a complex evolution of potential-energy surface versus strain.

Our finding indicates that, inside a nonuniformly strained bcc metal, H is favored to dissolve into either

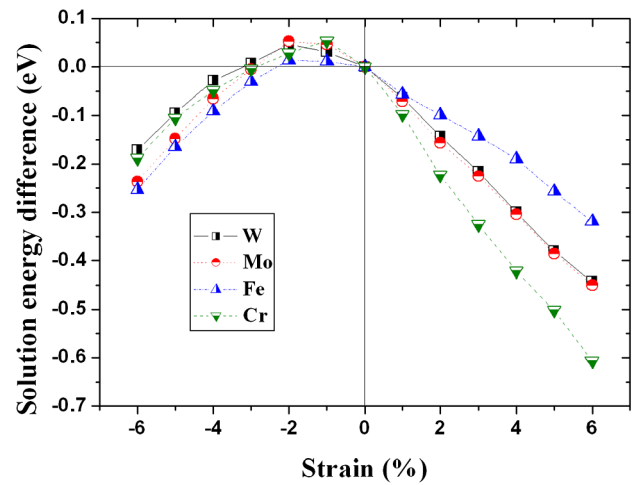


FIG. 4 (color online). Change of the H solution energy vs biaxial strain in the bcc metals studied. The biaxial strain is applied along the [100] and [010] directions. Negative values indicate enhanced H solubility in the strained bcc metals for given strain values.

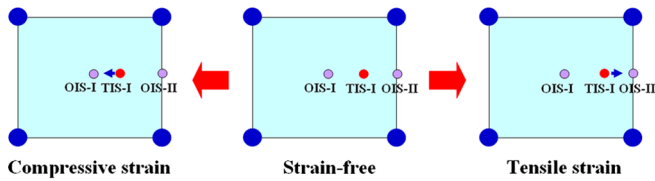


FIG. 5 (color online). The strain induced H position change starting at the TIS-I site. Compressive biaxial strain drives the H to the left towards the OIS-I site (left panel), and tensile biaxial strain drives the H to the right towards the OIS-II site (right panel).

compressive or tensile strained regions, except for a small range of compressive strain values. This has an important implication in using bcc metals as PFMs in a nuclear reactor. In particular, W has been considered as the most promising PFM, but H blistering in W PFM seriously influences the plasma stability and limits the lifetime of W PFM. H blistering is also a major concern in the fusion research. As such, many studies have been devoted to find ways to suppress H bubble formation in W [23–25]. The internal gas pressure in H bubbles is estimated up to a few tens of GPa [26,27], and thus H bubbles will exert a large strain to the surrounding W lattice, so that elastic strain is believed to be an important factor in affecting the H bubble formation [28]. However, the physical mechanism underlying the relationship between strain and H bubble formation remains poorly understood, and no microscopic insight has been given so far. Here, we propose that there exists a strain-triggered cascading effect on H bubble formation in bcc metals: the H solution leads to H bubble formation that induces anisotropic strain in the W lattice around the bubble, which in turn enhances local H solubility that facilitates further growth of H bubbles.

In particular, H bubbles are observed to form in the surface regions after W is exposed to H plasma irradiation [27,29] and to adopt different size and shape depending on the microstructure of the W crystal target and the irradiation temperature. Therefore, the strain in the W lattice surrounding the H bubble must be highly nonuniform and anisotropic due to the low-symmetry surface environment and the irregular size and shape of the H bubble. According to our study, such an anisotropic strain field will drive more H atoms to segregate into the vicinity of H bubbles, even independent of the sign of strain for most strain values, and the larger the anisotropic strain is, the lower the H solution energy will be. The increased H concentration will then facilitate further growth of the H bubble, and as the bubble grows bigger it induces even larger anisotropic strain around it, which in turn attracts more H to make the bubble grow even bigger, resulting in a cascading effect.

This research is supported by the National Magnetic Confinement Fusion Program with Grant No. 2009GB106003 and by the National Natural Science

Foundation of China (NSFC) with Grant No. 51061130558. F. Liu thanks the DOE-BES (Grant No. DE-FG02-04ER46148) and NSF-MWN (Grant No. DMR0909212) programs for support.

*LGH@buaa.edu.cn

- [1] G. Leibfried and N. Breuer, in *Springer Tracts in Modern Physics* (Springer-Verlag, New York, 1978), Vol. 81.
- [2] C. Elsässer, M. Fähnle, L. Schimmele, C. T. Chan, and K. M. Ho, *Phys. Rev. B* **50**, 5155 (1994).
- [3] J. Y. Zhu, F. Liu, G. B. Stringfellow, and S. H. Wei, *Phys. Rev. Lett.* **105**, 195503 (2010).
- [4] S. M. Myers, M. I. Baskes, H. K. Birnbaum, J. W. Corbett, G. G. DeLeo, S. K. Estreicher, E. E. Haller, P. Jena, N. M. Johnson, R. Kirchheim, S. J. Pearton, and M. J. Stavola, *Rev. Mod. Phys.* **64**, 559 (1992).
- [5] Y. Fukai and N. Okuma, *Phys. Rev. Lett.* **73**, 1640 (1994).
- [6] Y. Fukai, *Phys. Scr.* **T103**, 11 (2003).
- [7] P. Chen, Z. T. Xiong, J. Z. Luo, J. Y. Lin, and K. L. Tan, *Nature (London)* **420**, 302 (2002).
- [8] S. Li, P. Jena, and R. Ahuja, *Phys. Rev. B* **73**, 214107 (2006).
- [9] A. R. Troiano, *Trans. Am. Soc. Met.* **52**, 54 (1960).
- [10] J. B. Condon and T. J. Schober, *J. Nucl. Mater.* **207**, 1 (1993).
- [11] G. Kresse and J. Hafner, *Phys. Rev. B* **47**, 558 (1993).
- [12] G. Kresse and J. Furthmüller, *Phys. Rev. B* **54**, 11 169 (1996).
- [13] J. P. Perdew and Y. Wang, *Phys. Rev. B* **45**, 13 244 (1992).
- [14] P. E. Blöchl, *Phys. Rev. B* **50**, 17 953 (1994).
- [15] H. J. Monkhorst and J. D. Pack, *Phys. Rev. B* **13**, 5188 (1976).
- [16] K. P. Huber and G. Hertzberg, *Molecular Spectra and Molecular Structure IV: Constants of Diatomic Molecules* (Van Nostrand Reinhold, New York, 1979).
- [17] Y. L. Liu, Y. Zhang, H. B. Zhou, G. H. Lu, F. Liu, and G. N. Luo, *Phys. Rev. B* **79**, 172103 (2009).
- [18] C. Duan, Y. L. Liu, H. B. Zhou, Y. Zhang, S. Jin, G. H. Lu, and G.-N. Luo, *J. Nucl. Mater.* **404**, 109 (2010).
- [19] Y. A. Du, L. Ismer, J. Rogal, T. Hicel, J. Neugebauer, and R. Drautz, *Phys. Rev. B* **84**, 144121 (2011).
- [20] H. Hu, M. Liu, Z. F. Wang, J. Zhu, D. Wu, H. Ding, Z. Liu, and F. Liu, *Phys. Rev. Lett.* **109**, 055501 (2012).
- [21] H. Krimmel, L. Schimmele, C. Elsässer, and M. Fähnle, *J. Phys. Condens. Matter* **6**, 7705 (1994).
- [22] R. Griessen, *Phys. Rev. B* **38**, 3690 (1988).
- [23] H. B. Zhou, Y. L. Liu, J. Shuo, Y. Zhang, G. N. Luo, and G. H. Lu, *Nucl. Fusion* **50**, 115010 (2010).
- [24] O. V. Ogorodnikova, T. Schwarz-Selinger, K. Sugiyama, and V. K. Alimov, *J. Appl. Phys.* **109**, 013309 (2011).
- [25] M. J. Baldwin, R. P. Doerner, W. R. Wampler, D. Nishijima, T. Lynch, and M. Miyamoto, *Nucl. Fusion* **51**, 103021 (2011).
- [26] A. V. Veen, H. A. Filius, J. D. Vries, K. R. Bijkerk, G. J. Rozing, and D. Segers, *J. Nucl. Mater.* **155**, 1113 (1988).
- [27] W. M. Shu, E. Wakai, and T. Yamanishi, *Nucl. Fusion* **47**, 201 (2007).
- [28] M. Balden, S. Lindig, A. Manhard, and J. H. You, *J. Nucl. Mater.* **414**, 69 (2011).
- [29] R. A. Causey and T. J. Venhaus, *Phys. Scr.* **T94**, 9 (2001).