

## Photoinduced ir-Active Vibrations in *trans*-(CD)<sub>x</sub>: A Three-Mode System

In a recent Letter Blanchet *et al.*<sup>1</sup> have extended photoinduced absorption measurements of *trans*-polyacetylene to lower energy than obtained before<sup>2</sup> for both isotopes, (CH)<sub>x</sub> and (CD)<sub>x</sub>. While in (CH)<sub>x</sub> three strong ir-active vibrations (IRAV) are photoinduced, only two IRAV were observed in (CD)<sub>x</sub>, at  $\omega_1 = 400 \text{ cm}^{-1}$  and at  $\omega_2 = 1045 \text{ cm}^{-1}$ . This appeared strange, since from doping-induced IRAV and from resonant Raman scattering (RRS) it is known that (CD)<sub>x</sub> is also a three-mode system.<sup>3,4</sup> Because of the asymmetry of the mode at  $1045 \text{ cm}^{-1}$ , it was suggested that two vibrations may be contained in this line, but high-resolution examination of the line failed to resolve it.<sup>1</sup>

We report a high-resolution photoinduced absorption measurement of (CD)<sub>x</sub> at 80 K obtained with a double integration technique using a setup described before,<sup>2</sup> in which we have observed the third IRAV in (CD)<sub>x</sub>. As shown in Fig. 1  $\omega_3$  peaks at  $1225 \text{ cm}^{-1}$  and its relative intensity  $I_3$  is more than an order of magnitude smaller than  $I_2$  ( $I_2/I_3 \approx 14$ ). Its location and small intensity may explain why this mode was not observed before.<sup>1</sup> On the other hand both its location and intensity were recently predicted with use of the amplitude-mode formalism<sup>3</sup> to describe the strongly coupled phonons in polyacetylene.<sup>4</sup>

When charges are added to the chain the induced absorption  $\Delta\alpha(\omega)$  is given by<sup>3</sup>

$$\Delta\alpha(\omega) \sim \omega D_0(\omega) / [1 + (1 - \alpha_p) D_0(\omega)], \quad (1)$$

where

$$D_0(\omega) = \sum_{n=1}^3 [(\omega/\omega_n^0)^2 - 1]^{-1} \lambda_n / \lambda.$$

In Eq. (1),  $\omega_n^0$  and  $\lambda_n$  are the bare phonon frequencies and their relative  $e$ - $p$  coupling and  $\alpha_p$  is a pinning parameter.<sup>3</sup> The IRAV are the poles of Eq. (1) and satisfy the equation  $D_0(\omega) = -(1 - \alpha_p)^{-1}$ .  $D_0(\omega)$  for (CD)<sub>x</sub> was inferred from RRS measurements<sup>4</sup> and is plotted in Fig. 1. The intersections of  $D_0(\omega)$  with the horizontal line drawn at  $-(1 - \alpha_p)^{-1}$  with  $\alpha_p = 0.055$  give accurately the experimental IRAV frequencies. These include the two modes  $\omega_2$  and  $\omega_3$  (observed at  $1045$  and  $1225 \text{ cm}^{-1}$ ) as well as the lowest mode  $\omega_1$  ("pinned mode") which, however, could not be reached with our equipment, but it is reported at  $\omega_1 \approx 400 \text{ cm}^{-1}$  [Ref. (1)]. The dashed line de-

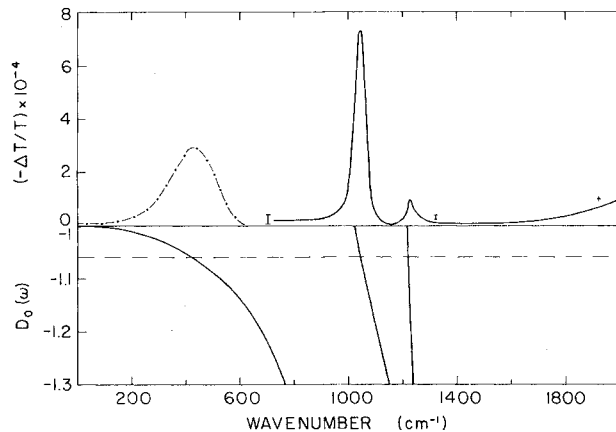


FIG. 1. Photoinduced absorption spectrum of *trans*-(CD)<sub>x</sub> at 80 K. The dot-dashed line is theoretical.  $D_0(\omega)$  is displayed and the horizontal line intersections give the IRAV frequencies.

scribing the "pinned mode" is a normalized theoretical<sup>3</sup> prediction for its line shape based on a Gaussian distribution for  $\alpha_p$  centered around 0.055 and with a width of 0.025.

In this model the IRAV relative intensities  $I_n$  are proportional to  $\omega$  times the residues [in Eq. (1)] at each pole  $\omega_n$ . This residue is inversely proportional to the derivative  $D' = \partial D_0(\omega) / \partial \omega$  at  $\omega = \omega_n$ . This explains the relative intensities in Fig. 1:  $I_1 > I_2 \gg I_3$  while the order is reversed in  $D'$ .

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<sup>1</sup>Graciela B. Blanchet, C. R. Fincher, T. C. Chung, and A. J. Heeger, *Phys. Rev. Lett.* **50**, 1938 (1983).

<sup>2</sup>Z. Vardeny, J. Orenstein, and G. L. Baker, *Phys. Rev. Lett.* **50**, 2032 (1983).

<sup>3</sup>B. Horovitz, *Solid State Commun.* **41**, 729 (1982).

<sup>4</sup>Z. Vardeny, E. Ehrenfreund, O. Brafman, and B. Horovitz, to be published.