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## Stability, Erosion, and Morphology Considerations for Sustainable Slope Design

A Dissertation Presented for the Doctor of Philosophy Degree The University of Tennessee, Knoxville

> Isaac Andres Jeldes Halty May 2014

Copyright © 2014 by Isaac A. Jeldes Halty All rights reserved. To my wife Daniela and my parents, whom I love more than anything.

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#### Abstract

The construction of more natural and sustainable earth slopes requires the consideration of erosion and runoff characteristics as an integral part of the design. These effects not only result in high costs for removal of sediment, but also a profound damage to the ecosystem. In this dissertation, innovative techniques are developed such that more natural appearing slopes can be designed to minimize sediment delivery, while meeting mechanical equilibrium requirements. This was accomplished by: a) examining the fundamental failure modes of slopes built with minimum compaction (FRA) to enhance quick establishment of forest, b) investigating the geomechanical and erosion stability of concave slopes, and c) developing design equations for a new type of inclined-face retaining structure, the Piling Framed Retaining Wall (PFRW), which in the limit is a confined slope. The analysis of several potential failures via Limit Equilibrium (LEM) and Finite Element (FEM) suggested that the governing failure of FRA slopes is shallow and well represented by infinite slope conditions, and laboratory and field data suggests that seasonal increase of stability due to matric suction is possible, while instability may occur under local seismicity. The investigation of the mechanical and erosion stability of concave slopes began with a mathematical definition of critical concave slopes at limiting equilibrium. Based on this, a mechanism to design concave slopes for a selected Factor of Safety (FS) was proposed. Results indicated that concave slopes can yield 15-40% less sediment than planar slopes of equal FS, and the stability is not compromised by errors in the construction. Concave slopes satisfying mechanical equilibrium are not necessarily in erosion equilibrium as observed in many natural landscapes. It was shown that when these two equilibrium conditions are met, the slopes become sustainable and a set of equations describing sustainable concave slopes was proposed. Finally, rational design equations

for the innovative PFRW were developed based on numerous FEM analyses for different soil and geometry conditions. The equations provided a good prediction of the soil stresses measured on a PFRW built in Knoxville, TN.

## Preface

This dissertation comprises the assembly of six manuscripts, which at the present time, are at different stages of review and publication in various peer review journals. Each chapter corresponds to a unique article, where methods, results, conclusions, references and in some cases appendices are individually addressed. The first five chapters (articles) deal with geoenvironmental slope stability issues, where the main objective is to investigate, and finally provide, mechanisms for more natural and sustainable design of earth slopes. This research was partially funded by the Office of Surface Mining Reclamation and Enforcement, and was motivated from the growing need for environmental-friendly techniques for landform construction and sustainable land management. In these first five chapters efforts have been made to include surficial water erosion as a key variable in the slope design, and attempts have been made to integrate water erosion and slope stability, which are traditionally treated as separated sciences. Through these chapters, the reader will find topics ranging from theoretical development of equations to illustrative design examples, with the overall goal of providing practical tools that engineers can use in design. The last article, on the other hand, addresses a different stability topic related to an innovative retaining wall system comprising an inclined wall face, which in a practical sense can be seen as a *confined slope*. This research was funded by the Tennessee Department of Transportation, with the objective of creating a rational design methodology for this new type of wall, eliminating the need to conduct extensive numerical analyses for each wall to be constructed.

## **Table of Contents**

Chapter 1. Introduction1
Dissertation Overview
Selected fundamental frameworks for mechanic and erosion soil modeling
The Low Compaction Grading Technique on steep slopes (chapters 2 and 3) 4
The mechanical and erosional stability of concave slopes (chapters 4, 5 and 6)
The Piling Framed Concrete Retaining Wall (chapter 7) 10
References:
Chapter 2. The Low Compaction Grading Technique on Steep Reclaimed Slopes: Soil
Characterization and Static Slope Stability14
Abstract
Introduction16
Methods18
Location of field sites
Site construction and reclamation process
Geotechnical characterization
Static long-term slope stability analyses
Results and Discussion
Geotechnical characterization of research sites
Static long-term slope stability analyses
Conclusions
Acknowledgments

References:	
Chapter 3. Partial Saturation and Seismicity on Steep Reclaimed Slopes	47
Abstract	
Introduction	49
Background	50
Unsaturated soil shear strength	50
Shear strength and the pseudo-static forces for seismic analyses	
Field and experimental methods	55
Study sites: construction and site characterization	55
Soil water characteristic curves for unsaturated stability analyses	59
Peak ground and spectral accelerations for seismic stability analyses	61
Analytical methods for stability analyses	61
Infinite slope equation for unsaturated soil strength	61
Infinite Slope equation for horizontal and vertical pseudo-static forces	65
Summary of the methodology	67
Results and Discussion	68
SWCCs and stability results for unsaturated FRA soils	68
Seismic stability analyses	74
Conclusions	
Acknowledgments	79
References:	81
Chapter 4. An Approximate Solution to the Sokolovskiĭ Concave Slope at Limi	iting
Equilibrium	
	ix

Abstract	90
Introduction	90
Background: the slip line field theory and the characteristic equations	92
Elaboration of Sokolovskiĭ solution for the critical slope in a weightless medium	96
Proposed solution for the critical slope in a medium with self-weight	99
Validation of the WMA solution	04
Geometry of the concave slopes	04
Critical FS and observed failure mechanisms	07
Effects of slope height on failure mechanisms and FS obtained from the proposed solution 1	10
Discussion and conclusions	13
References:	15
Appendix: The slip line field theory and the characteristic equations	18
Equations of internal equilibrium and principal stresses	18
Formulation of the equations of the characteristics	19
Chapter 5. Design of Stable Concave Slopes for Reduced Sediment Delivery	27
Abstract	28
Introduction12	28
Background	30
Mechanical stability of concave profiles13	30
Concave slopes and soil erosion	31
Methods13	34
Mechanical stability	34
Soil erosion and sediment yield	38
	Х

Construction and sensitivity	
Results and Discussion	
Concave slopes with pre-selected FS's for long-term ( $\phi > 0$ ) conditions	
Stability check for short-term ( $\phi = 0$ ) conditions	
Soil loss from the mechanically stable concave slopes	
Sensitivity to construction accuracy	150
Illustrative Examples	152
Finding the concave profile for long-term stability:	
Checking the short-term stability of the concave slope:	
Erosion analyses:	156
Conclusions	156
References:	159
Chapter 6. Sustainable Concave Slopes	
Chapter 6. Sustainable Concave Slopes	<b>164</b> 165
Chapter 6. Sustainable Concave Slopes Abstract Introduction	<b> 164</b> 165 166
Chapter 6. Sustainable Concave Slopes Abstract Introduction Background	164 165 166 167
Chapter 6. Sustainable Concave Slopes	<b> 164</b> 165 166 167 167
Chapter 6. Sustainable Concave Slopes	<b> 164</b> 165 166 167 167 170
Chapter 6. Sustainable Concave Slopes	<b> 164</b> 165 166 167 167 170 173
Chapter 6. Sustainable Concave Slopes	<b> 164</b> 165 166 167 167 170 173 173
Chapter 6. Sustainable Concave Slopes	164         165         166         167         167         170         173         173         178

Results and Discussion	
Concave slopes with constant rate of erosion	185
Sustainable slopes	
Limitations of this work	195
Conclusions	198
References:	199
Chapter 7. The Piling Framed Concrete Retaining Wall: Design Pressures and S	Stability
Evaluation	203
Abstract	
Introduction	205
Background	206
The PFRW concept and the construction sequence	206
Theoretical expressions for active, passive, and at-rest earth pressure coefficient	s 208
Methods	
Numerical Analyses	
Model stresses vs. field measured stresses: The SmartFix case	220
Results and Discussion	222
Normal stresses on battered wall	222
Destabilizing moments for overturning stability	227
Design equations to predict earth pressures and destabilizing moments	229
The SmartFix case: predicted vs. measured stresses on the inclined wall	
Conclusions	
Acknowledgments	
	xii

References:
Appendix
Active, passive, and at rest states of the soil mass
Active and passive lateral earth pressure coefficients: a literature review
At-rest lateral earth pressure coefficients: a literature review
Definition of boundary condition and distance to boundaries
Generalized FEM soil stresses in terms of the dimensionless quantities $\sigma_{_{Neq}}/c$ and $H\gamma/c$
Generalized FEM overturning moments in terms of the quantities $M\gamma/c$ and $H\gamma/c$ 256
Validation of simplified design equations
Determination of soil strength parameters
Chapter 8. Conclusions
The Low Compaction Grading Technique on steep slopes (chapters 2 and 3)
The mechanical and erosional stability of concave slopes (chapters 4, 5 and 6)
The Piling Framed Concrete Retaining Wall (chapter 7)
References:
Vita

## List of Tables

Table 2.1 Average slope length, width and inclination angle for the four plots at the Premium,
National and Mountainside sites
<b>Table 2.2</b> Mean values of Liquid Limit, Plastic Index, soil texture and soil classification (USCS)
Table 2.3 Means, standard deviations, 95% confidence intervals (C.I.), and 90% tolerance intervals (T.I) (80% coverage) for wet and dry unit weights for Premium, National and Mountainside sites         32
<b>Table 2.4</b> Summary of internal friction angle and cohesion for reclaimed mine materials
<b>Table 2.5</b> FS obtained for long-term static stability focused on the low strength surface layer 36
<b>Table 3.1</b> Summary of recommended pseudo-static values and expressions available in the literature (after Duncan and Wright, 2005)
<b>Table 3.2</b> Average values of slope length, inclination angle, unit weights, water content, andobserved angles of repose for Premium, National and Mountainside sites.60
<b>Table 3.3</b> Soil texture (USDA) and unit weights for ROSETTA model
<b>Table 3.4</b> Summary of $K_h$ , $K_v$ and obtained FS's for a 2 and 10% P.E. in 50 yr
<b>Table 4.1</b> Maximum values of $H_s(m)$ reported by Sokolovskii (1960) and used in stabilityanalyses (1960) solution from LEM105
<b>Table 4.2</b> Comparison of stability between proposed approximate solution and Sokolovskii (1960)       solution from LEM
Table 5.1 Soil composition, classification, erodibility and assumed internal friction angle of investigated soils
<b>Table 5.2</b> Reduction in soil loss $(1 - A_c / A_p)$ from concave contours reported in the literature 151

Table 5.3 Geometry, soil properties, and probable failure mechanisms of illustrative examples
<b>Table 5.4</b> Soil parameters and results from erosion analyses of illustrative examples
<b>Table 7.1</b> Wall configurations being investigated for every combination of $\phi$ and $c/\gamma$ 213
Table 7.2 Section Properties of Beam Elements       216
Table 7.3 Construction sequence as approximated in the FEM models         221

## List of Figures

Fig. 1.1. Illustration of the superior mechanical and erosion resistances of concave slopes [adapted
from Schor and Gray (2007)]7
<b>Fig. 2.1.</b> Location of field sites in northeastern Tennessee, referred to as Premium, National, and Mountainside
<b>Fig. 2.2.</b> National site after the FRA reclamation process and during the construction of the study plots
Fig. 2.3. Depiction of the reclamation process according to FRA
<b>Fig. 2.4.</b> Randomized systematic sampling technique; a) area of interest subdivided into small sub- areas (Sweigard et al. 2007c), where small circles represent the random measurement locations; b) application of this technique for NDG measurements at the National site
<b>Fig. 2.5.</b> Typical field determination of the angle of repose at National Site (White et al. 2009). Camera placed on a level and photo taken of material in loose state. Note the large number of oversize ( $> 0.3$ m) particles in the reclaimed material
<b>Fig. 2.6.</b> Shear strains and nodal displacements obtained from the FEM analysis assuming a very strong core. Upper left corner illustrates the geometric dimensions employed for LEM and FEM analyses
<b>Fig. 2.7.</b> Failure mechanisms and FS's obtained from FEM stability analysis for $\tan \phi_{core} / \tan \phi_{loose}$ = 1.3. The FS = 1.94 shown for the deeper failure mechanism was obtained when the strength reduction factor (SRF) search was restricted to be outside the zone where the shallow mechanism occurred. Shear strains showed for SRF = 2.02 to emphasize failure mode
<b>Fig. 2.8.</b> FEM analyses of the global stability for various values of $\tan \phi_{core} / \tan \phi_{loose}$ and the infinite slope equation
Fig. 3.1. Location of field sites in northeastern Tennessee, referred to as Premium, National, and Mountainside

Fig. 3.2. Depiction of the reclamation process according to FRA (Jeldes et al. 2013)
Fig. 3.3. Depiction of the 3 Spectral Accelerations for Premium, National and Mountainside (USGS 2012)
Fig. 3.4. The Infinite Slope method. Figure modified from Salgado (2008)
Fig. 3.5. Infinite Slope method with horizontal and vertical pseudo-static forces
<b>Fig. 3.6.</b> Soil water characteristic curves measured in laboratory experiments and the ROSETTA model (Schaap et al. 2001), for the Premium site
<b>Fig. 3.7.</b> Soil water characteristic curves measured in laboratory experiments and the ROSETTA model (Schaap et al. 2001), for the National site
<b>Fig. 3.8.</b> Soil water characteristic curves measured in laboratory experiments and the ROSETTA model (Schaap et al. 2001), for the Mountainside site
<b>Fig. 3.9.</b> Calculated monthly variations in stability (FS), volumetric water content and cumulative monthly rainfall from July 2009 to August 2010 at the Premium site
<b>Fig. 3.10.</b> Calculated monthly variations in stability (FS), volumetric water content and cumulative monthly rainfall from July 2009 to August 2010 at the National site
<b>Fig. 3.11.</b> Calculated monthly variations in stability (FS), volumetric water content and cumulative monthly rainfall from July 2009 to August 2010 at the Mountainside site
<b>Fig. 3.12.</b> FS chart for Mountainside site as a function of spectral accelerations and initial fundamental period
<b>Fig. 4.1.</b> Orientation of slip lines in a soil mass at limiting equilibrium [adapted from Sokolovskiĭ (1965)]
<b>Fig. 4.2.</b> Critical slope shape (planar) for a weightless medium. Inside zone AoC the principal directions are aligned with the reference frame. On the slope surface, the principal directions are aligned with the normal and tangential directions

Fig. 4.3. Slope surface in a medium possessing weight [after Sokolovskiĭ, (1965)]...... 100

Fig. 4.4. Weightless Medium Approximation: discretized weightless medium (e.g. rigid foam beads) supporting thin, external loads (e.g. frictionless lead sheets) separated by a finite interval: Fig. 4.5. Sokolovskii's critical slope vs the real (non-imaginary) component of the solution of Eq. Fig. 4.6. Comparison of Sokolovskii (1960) numerical solution and the proposed analytical Fig. 4.7. Comparison of Sokolovskii (1960) numerical solution and the proposed analytical Fig. 4.8. FEM results in terms of shear strains for the Sokolovskii (1960) critical concave slope. 108 Fig. 4.9. FEM results in terms of shear strains for the proposed critical concave slope. .......... 109 Fig. 4.10. Modes of failure depicted in the form of shear bands (maximum shear strains) as observed in concave slopes (FEM). Figure illustrate case for  $\phi = 20^{\circ}$  and  $c/\gamma = 2$  m, and two values of H<sub>s</sub>.....111 **Fig. 4.11.** Computed FS (LEM) on concave slopes with different vertical heights  $H_s$  and  $c/\gamma$  for  **Fig. 5.1.** Illustration of the reference frame and the components of the concave slope formulation. In practice the sharp bluff at the top of the slope may be rounded, as indicated by the dashed line.

Fig. 5.2	Long-term ( $\phi$	> 0) critical	contour vs	short-term ( $\phi$	= 0) critic	cal contours	defined by
differen	t S <sub>u</sub> values	••••••		•••••			

Fig. 5.3. Illustration of concave and equivalent planar slopes (same FS and slope height). ..... 139

Fig. 5.9. Decrease in FS due to low precision construction (200 mm vertical) of the concave slope.

Fig. 5.10. Example case for sandy soil: $\phi = 35^{\circ}$ , $c = 15 \text{ kN/m}^2$ , $\gamma = 19 \text{ kN/m}^3$ and $H_s = 15 \text{ m}$ .
Required $FS = 1.5$ . Shear strains showed for $SRF = 1.52$ to emphasize failure mode
Fig. 6.1. Conceptual model of slope morphology evolution by water erosion
Fig. 6.2. Illustration of the parallel retreat concept. The sharp edge at the top of the slope may be
naturally eroded (becoming convex) if there is runoff coming over the top edge of the slope. This effect is not included in the proposed model
Fig. 6.3. Erosion rates experienced by a planar profile as modeled by RUSLE2 perspective 175
<b>Fig. 6.4.</b> Illustration of a potential family of concave slopes with constant rate of erosion: a) high rate, b) intermediate rate, and c) low rate
Fig. 6.5. Discretization of the slope profile in small linear segments
<b>Fig. 6.6.</b> Illustration of steady-state or sustainable slope (profile 1) and a non-equilibrium slope with cyclic morphology changes (profile 2)
Fig. 6.7. Family of concave slopes with constant rate of erosion (equilibrium erosion profiles).
Fig. 6.8. The overall steepness (defined as the slope of the straight line that connects the top and
the bottom of the profile) increases at a much faster rate for $A_l^n \ge 1$ , as denoted by dashed line. 
<b>Fig. 6.9.</b> Erosion equilibrium profiles and a critical concave contour for $\phi = 20^{\circ}$ and $c / \gamma = 1$ m.
Concave slopes with $A_1^n \le 1.05$ are considered sustainable
<b>Fig. 6.10.</b> Erosion equilibrium profiles and a critical concave contour for $\phi = 30^{\circ}$ and $c/\gamma = 0.5$ m. Concave slopes with $A_1^n \le 1.039$ are considered sustainable
<b>Fig. 6.11.</b> Erosion equilibrium profiles and a critical concave contour for $\phi = 20^{\circ}$ , 25°, 30°, 35°, 40° and $c/\gamma = 0.5, 0.75, 1, 1.5, 2, 3$ m

<b>Fig. 6.12.</b> Limiting values of erosion $(A_L^n)$ defining sustainability
<b>Fig. 6.13.</b> Erosion equilibrium profiles and a concave contours with $FS = 1.25$ , for $\phi = 20^{\circ}$ , $25^{\circ}$ ,
30°, 35°, 40° and $c/\gamma = 0.5, 0.75, 1, 1.5, 2, 3$ m
<b>Fig. 6.14.</b> Erosion equilibrium profiles and a concave contours with $FS = 1.5$ , for $\phi = 20^{\circ}, 25^{\circ}, 30^{\circ}, 30$
35°, 40° and $c / \gamma = 0.5, 0.75, 1, 1.5, 2, 3$ m
<b>Fig. 6.15.</b> Erosion equilibrium profiles and a concave contours with $FS = 1.75$ , for $\phi = 20^{\circ}$ , $25^{\circ}$ ,
30°, 35°, 40° and $c/\gamma = 0.5, 0.75, 1, 1.5, 2, 3$ m
Fig. 6.16. Computed slope angles vs horizontal slope length
Fig. 7.1. The PFRWs built along the I-40/I-75 corridor (Knoxville, TN). a) The SmartFix wall and
b) The West Hills wall
Fig. 7.2. The piling frame system forming the basis of the PFRW concept
Fig. 7.3. Construction sequence for PFRW
Fig. 7.4. Illustration of the construction stages for the timber lagging, concrete face and concrete
cap. The black vertical material is a geosynthetic filter layer
Fig. 7.5. Definition of "effective" wall height, wall inclination and backfill slope 213
Fig. 7.6. Resulting wall deflection over the inclined face of the wall ( $\delta/L$ ). Triangle and square
dots are configurations of wall geometry and soil properties excluded from the result analyses due
to large wall deflection ( $\delta/L > 1/240$ )
Fig. 7.7. Typical FEM model for the PFRW: boundary conditions, distance to boundaries and
structural elements considered
Fig. 7.8. Earth pressure cells (white circles) positioned in Sections 1 and 2. Distances shown for
Section 1 are along the inclined plane and measured from the final grade

Fig. 7.10. Definition of equivalent normal distributed pressures  $\sigma_{Neq}$ : a) FEM and b) Coulomb and Fig. 7.11. Equivalent normal uniform stress  $\sigma_{Neq}$  vs. wall height H, for the  $\phi = 30^{\circ}$  and  $c/\gamma =$ Fig. 7.12. Illustration of a case when the theoretical Coulomb earth pressures are lower than FEM predictions, and therefore unconservative ( $\phi = 30^{\circ}$  and  $c/\gamma = 1$  m,  $\alpha = 20^{\circ}$ , and  $\beta = 50^{\circ}$ ). ... 226 Fig. 7.13. Dimensionless equivalent normal earth pressures vs. stability number as predicted by Fig. 7.14. Forces and moment arms employed for overturning moments. a) Numerical FEM and Fig. 7.15. Destabilizing moment M vs. wall height H, for the  $\phi = 30^{\circ}$  and  $c / \gamma = 1$  m soil case Fig. 7.16. Normalized overturning moment vs. stability number as predicted by FEM for  $\phi = 30^{\circ}$ . Fig. 7.18. Field measured stresses over time vs. theoretical and numerical earth pressures..... 232 Fig. 7.20. Two FEM models considered: a) the 0.5 m of the battered pile below the final grade is

Fig. 7.21. Difference in resulting normal stresses from the two battered wall end conditions
investigated (SmartFix case). The model comprising a pinned condition at the 0.5 m below final
grade resulted in higher overall stresses
Fig. 7.22. Dimensionless equivalent normal earth pressures vs stability number as predicted by
FEM for $\phi = 20^{\circ}$
<b>Fig. 7.23.</b> Dimensionless equivalent normal earth pressures vs stability number as predicted by
FEM for $\phi = 30^{\circ}$
Fig. 7.24. Dimensionless equivalent normal earth pressures vs stability number as predicted by
FEM for $\phi = 40^{\circ}$
<b>Fig. 7.25</b> . Normalized overturning moment vs. stability number as predicted by FEM, $\phi = 20^{\circ}$ .
<b>Fig. 7.26.</b> Normalized overturning moment vs. stability number as predicted by FEM, $\phi = 30^{\circ}$ .
<b>Fig. 7.27.</b> Normalized overturning moment vs. stability number as predicted by FEM $\phi = 40^{\circ}$ .
<b>Fig. 7.28.</b> Validation of design equations: a) $\sigma_{Neq}$ and b) <i>M</i> for $\phi = 20^{\circ}$
<b>Fig. 7.29.</b> Validation of design equations: a) $\sigma_{Neq}$ and b) <i>M</i> for $\phi = 30^{\circ}$
<b>Fig. 7.30.</b> Validation of design equations: a) $\sigma_{Neq}$ and b) <i>M</i> for $\phi = 40^{\circ}$
Fig. 7.31. Selected M-C strength parameters based on triaxial p-q diagram at 14% strain level.

# Chapter 1. Introduction

#### **Dissertation Overview**

The present dissertation is a collection of six manuscripts, each a single chapter, at different stages of review and publication in various peer review journals. The original contributions of this work are mostly centered on topics related to the broad area of slope stability, with efforts to integrate the traditionally separate areas of mechanical slope stability and surficial rain-driven water erosion. In this vein, *chapters 2* and *3* addresses stability issues related to a new reclamation technique for rapid reforestation and reduction of soil loss, while chapters 4, 5 and 6 focus on understanding the role of the topography of concave contours--similar to those existing in nature--on the mechanical and erosion stability of slopes. Through these first five chapters, topics ranging from theoretical development of equations to illustrative design examples are offered with the overall goal of providing mechanisms for the design of more natural and sustainable earth slopes. The last chapter (chapter 7), on the other hand, addresses a different stability topic related to an innovative retaining wall system comprising an inclined face, which in a practical sense can be seen as a *confined slope*. In this section, equations to predict earth pressures and overturning moments are developed to support the creation of a rational design methodology, eliminating the need to conduct extensive numerical analyses for each wall to be constructed.

#### Selected fundamental frameworks for mechanic and erosion soil modeling

In soil mechanics, numerous constitutive laws defining *failure* have been proposed such that strength and deformation characteristics of soils can be modeled for engineering purposes. For frictional materials like soils, the *Mohr-Coulomb* (M-C) yield criterion arises as the "best known" fundamental law (Smith and Griffiths 2004) and it has been the most widely used model for soil shear strength prediction (Griffiths and Lane 1999, Salgado 2008). This constitutive model does

not consider the effects of intermediate principal stresses (Desai and Siriwardane 1984), and as a result, it takes the form of an irregular hexagonal cone in the principal stress space, whose mathematical representation in terms of principal stresses is (Potts and Zdravković 1999):

$$f = \frac{1}{2} (\sigma_1 + \sigma_3) \sin \phi + c \cos \phi - \frac{1}{2} (\sigma_1 - \sigma_3)$$
(1.1)

where *f* is the failure function,  $\sigma_1$  and  $\sigma_3$  are the major and minor principal stresses,  $\phi$  is the internal friction angle, and *c* is the soil cohesion. From Eq. (1.1) it can be seen that the M-C model captures the behavior dependency on the mean normal stress that frictional materials such as soils exhibit, and it is based only on two laboratory determined parameters:  $\phi$  and *c*. Its simplicity and reliability makes the M-C constitutive model a suitable framework to study the mechanical behavior of earth slopes, and it has been used throughout the entire manuscript. A special emphasis on slopes with soils at critical equilibrium is given in *chapters 4, 5* and 6, which is the fundamental assumption in the development of *mechanically stable concave slopes* (Jeldes et al. 2013) In this special case, the soil body is taken into a limiting state (or at the verge of flowing plastically) by making the function *f* = 0. In other words, the largest difference between the shear stresses and the shear strength must satisfy the following mathematical relationship (Sokolovskiĭ 1960):

$$\max\left\{\left(\sigma_{1}-\sigma_{3}\right)\sin(2\lambda-\phi)-\sin\phi(\sigma_{1}+\sigma_{3}+2H)\right\}=0$$
(1.2)

Where  $\lambda$  is the angle between the direction of the normal to the plane and the major principal direction, and *H* is the tensile strength of the soil. The implications and limitations of the use of this constitutive framework are discussed throughout the manuscript when needed, always attempting to maintain a practical perspective in the discussion.

The Revised Universal Soil Loss Equation *RUSLE2* (USDA-ARS 2008) is one of the most widely used erosion equations and is known for its effectiveness and simplicity in accounting for the critical effects controlling erosion (Tiwari et al. 2000). RUSLE2 relies on concrete empirical soil loss data to model the important processes involving soil erosion:

$$A = R \cdot K \cdot LS \cdot C \cdot P \tag{1.3}$$

where the predicted soil loss A is directly proportional to: the rainfall erosivity R quantifying the rainfall's erosive potential; the soil erodibility K defining the soil's susceptibility to that erosivity; the topographic factor LS representing slope length and steepness effects; the surface cover factor C; and the conservation practices factor P. Its simplicity and accuracy in predicting erosion make RUSLE2 an ideal and reliable framework for soil loss investigations involving more complex (non-planar) slope topographies, and it was used here to investigate the effects of mechanically stable concave profiles on soil loss (*chapter 5*) and also to investigate the morphology of slopes in erosion equilibrium (*chapter 6*). The limitations of RUSLE2 in the context of this work were identified and discussed, and the assumptions made to overcome these limitations were properly justified throughout these two chapters.

### The Low Compaction Grading Technique on steep slopes (chapters 2 and 3)

Mine reclamation activities have traditionally incorporated compaction procedures to augment the strength of the reclaimed material and ensure stability of the restored slopes (Hoomehr et al. 2013, Jeldes et al. 2013). However, while compaction is important for strength and erosion resistance, it has negative impacts on tree survival and hampers reforestation efforts. The quick establishment of forest and ground cover is an important consideration for long-term erosion control, since vegetation absorbs raindrop impact, reduces flow energy by reducing runoff

velocity, increases soil infiltration, and increases soil resistance. A reclamation method that employs minimally compacted spoils to enhance native forest growth known as the Forest Reclamation Approach (FRA) is currently being promoted by the US Office of Surface Mining. The FRA method specifies the use of low compaction energy in the top 1.2 m to 1.5 m of the contour, which may be in conflict with general considerations for mechanical stability of steep slopes. In *chapters 2* and *3*, the stability of steep FRA slopes (steeper than 20 degrees) was investigated on three reclaimed coal mining sites in the Appalachian region of East Tennessee. These two chapters seek to answer the following research questions: a) what are the governing failure modes of steep reclaimed FRA slopes?, b) what are the key soil properties for proper FRA design?, c) considering that FRA comprises a plane of discontinuity parallel to the surface, are steep FRA slopes infinite slope candidates?, d) what are the implications of unsaturated mine spoils on the soil strength and how periodic changes in environmental conditions affect the stability of steep FRA slopes?, and e) what is the seismic response of these steep FRA slopes?.

In *chapter 2*, the geotechnical properties of low compacted spoils with a large number of oversize particles were investigated. Through the analysis of several potential modes of failure via Limit Equilibrium (LEM) and Finite Element Method (FEM) analyses, a rational method to design and evaluate the stability of steep FRA slopes is suggested. In *chapter 3*, a set of theoretical equations were developed to investigate the implications of unsaturated soils and seismicity on the stability of steep FRA slopes. These equations, coupled with field and laboratory analyses, demonstrated the seasonal variation of stability due to matric suction, while seismic analyses illustrated the conditions under which instability might occur. The likely field conditions would suggest that the FRA has no negative impact on slope stability, and the benefits of faster forest establishment in

terms of reduced erosion and sediment delivery make the FRA very attractive for future reclamation work.

#### The mechanical and erosional stability of concave slopes (chapters 4, 5 and 6)

Under the FRA concept, the quick establishment of ground cover and forest is important to reduce the high initial erosion and sediment delivery on reclaimed and constructed slopes. The efforts in chapters 4, 5 and 6 were then directed to understand the role of topography on erosion rates and to investigate the role of non-planar contour shapes--similar to those existing in nature--on the mechanical and erosion stability of slopes. While constructed slopes are traditionally designed to be planar in cross section, in nature curvilinear slopes with concave shapes are naturally formed as a result of evolutionary processes leading towards a state of erosion and sediment transport equilibrium. The superiority of concave slopes in terms of mechanical stability and erosion resistance can be explained by the conceptual model developed by Schor and Gray (2007). Assuming that a planar slope of height H and angle  $\beta$  can be discretized into a series of horizontal layers with equal thickness (Fig. 1.1), where the strength properties (internal friction angle  $\phi$  and cohesion c) and unit weight ( $\gamma$ ) remain constant for the entire slope, the mechanical stability of each layer will be dependent upon the particular stress-strength state and the angle on inclination of each layer. With the acceleration of gravity acting vertically downwards, the vertical stresses at each point within the soil mass increases proportional to the depth. To bring each layer into the same degree of equilibrium or Factor of Safety (FS), the inclination of each layer must be adjusted; steeper slopes will be allowed in the upper layers where less overburden mass exists, while gentler slopes will be necessary in the lower layers. The result of this process is a concave profile, which



**Fig. 1.1.** Illustration of the superior mechanical and erosion resistances of concave slopes [adapted from Schor and Gray (2007)].

provides a more uniform equilibrium state for the entire slope. On the other hand, surficial water erosion increases with slope length and slope angle. In a planar slope, water tractive forces will monotonically increase downslope, with higher energy in the lowest part of the contour. By gradually decreasing the steepness of each layer in the direction of the water flow, the increased erosion due to the increased slope length will be partially counteracted by the decreased erosion due to decreasing slope steepness, ultimately resulting in a concave slope with a minimized uniform erosion rate along the profile.

Realizing that not all concave shapes will be mechanically stable, it is desirable to have a description of concave slopes that provide mechanical stability for given set of soil properties. In these chapters answers for the following research questions are sought: a) what is the critical concave profile for mechanical stability considerations as defined by the Mohr-Coulomb constitutive model?, b) how can concave slopes be described such that they provide a desired degree of stability (or FS) for given soil properties?, c) how effective are these mechanically optimized concave slopes in reducing sediment delivery?, d) how does the accuracy of construction affects the mechanical stability of concave slopes?, e) what is the equilibrium concave profile for erosion considerations?, f) how do erosion equilibrium concave slopes compare with critical concave slopes defined for mechanical stability?, and g) can slopes be constructed to achieve both mechanical and erosional stability?

In *chapter 4*, an approximate analytical solution that defines the geometry of critical concave slopes (FS  $\approx$  1) for frictional soils with self-weight ( $\phi > 0$ , c > 0,  $\gamma > 0$ ) was developed, based on the slip line field method of Sokolovskiĭ (1960). The fundamentals behind Sokolovskiĭ's theory

were revisited, and the physical and mathematical derivations preceding the approximate solution are presented in detail. The approximate solution was validated using LEM and FEM analyses. *Chapter 5* builds on the previous chapter to provide a practical design approach such that geotechnical engineers will consider non-planar shapes, which provide a more natural appearance in addition to offering improved erosion resistance. This was accomplished by creating a mechanism to design slopes for a given FS, and suggesting a method for both long-term and shortterm stability investigations for concave slopes. The difference in soil loss (erosion) between planar and concave slopes satisfying the same degree of mechanical stability was also investigated, along with analyses to determine how sensitive concave slopes are to construction inaccuracies. An illustrative example was provided to demonstrate the design process and the benefits of meeting mass stability requirements while at the same time reducing surficial erosion and yielding a more natural looking slope.

While mechanically stable, the proposed concave slopes in *chapters 4* and 5 may not be in equilibrium from a water erosion perspective. Evidence exists that natural fluvial systems, seeking erosion and sediment transport equilibrium, may adjust the geometry in order to achieve a concave steady-state form that will be somewhat unchanged over time. In *chapter 6*, the concept of steady-state landforms was explored and a conceptual model of changes in slope morphology toward a concave erosion equilibrium shape was described. Based on the assumptions inherent to the RUSLE2 erosion model, concave profiles in water erosion equilibrium were identified and described, and an approach was proposed to discern between long-term mechanically stable and unstable erosion equilibrium shapes for any given combination of soil stresses and strength. A definition of the approximate limiting erosion rate at which equilibrium erosion shapes become

mechanically stable and thus "sustainable" is explored, and a mathematical expression to obtain this limiting erosion rate is offered as a function of the Mohr-Coulomb parameters.

#### The Piling Framed Concrete Retaining Wall (chapter 7)

This last chapter addresses a different stability topic related to an innovative retaining wall system called The Piling Framed Retaining Wall (PFRW). This wall system comprises an inclined wall face, which in a practical sense can be seen as a *confined slope* problem. The PFRW is ideal for applications where limited right-of-way (ROW) is available, or where adjacent structures and underground utilities limit the use of tie-back anchor systems. Since the soil pressures acting on the inclined wall face of this new system are not fully understood, a rational design method for this wall has not been developed. This chapter seeks answers to the following questions: a) what is the magnitude and distribution of the earth pressures on the face of the PFRW according to the associated mode of deformation?, b) how these earth pressures differ from theoretical predictions?, and c) how the earth pressures and overturning moments of PFRW's vary with respect to changes in wall geometry and soil properties?.

In this chapter, the soil pressures on the PFRW for various wall face inclinations, wall heights, and backfill slopes are investigated via FEM analyses, and results compared with theoretical expressions available in the literature. Simplified equations are developed from the FEM analyses to facilitate the design and stability calculations of PFRWs without the need to create geometricspecific finite element models. In addition, an approximate design approach based on the wellknown Coulomb earth pressure theory is demonstrated. The earth pressures predicted by the design equations and the approximate design approach were compared with field measured soil stresses on a PFRW built in Knoxville, TN.
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## Abstract

Since the Surface Mining and Control Reclamation Act of 1977, U.S. coal mining companies have been required by law to restore the approximate ground contours that existed prior to mining. To ensure mass stability and limit erosion, the reclaimed materials have traditionally been placed with significant compaction energy. The Forest Reclamation Approach (FRA) is a relatively new approach that has been successfully used to facilitate the fast establishment of native healthy forests. The FRA method specifies the use of low compaction energy in the top 1.2 m to 1.5 m of the contour, which may be in conflict with general considerations for mechanical slope stability. Although successful for reforestation, the stability of FRA slopes has not been fully investigated and a rational stability method has not been identified. Further, a mechanics-based analysis is limited due to the significant amount of oversize particles which makes the sampling and measurement of soil strength properties difficult. To investigate the stability of steep FRA slopes (steeper than 20 degrees), three reclaimed coal mining sites in the Appalachian region of East Tennessee were investigated. The stability was evaluated by several methods to identify the predominant failure modes. The infinite slope method, coupled with the estimation of the shear strength from field observations, was shown to provide a rational mean to evaluate the stability of FRA slopes. The analysis results suggest that the low compaction of the surface materials may not compromise the long-term stability for the sites and material properties investigated.

# Introduction

A reclamation method that employs minimally compacted spoils to enhance native forest growth, known as the Forest Reclamation Approach (FRA) is currently being promoted by the US Office of Surface Mining (OSM) (Angel et al. 2007, Sweigard et al. 2007). Since the Surface Mining and Control Reclamation Act of 1977 (SMCRA), coal companies in the U.S.A. have been required by law to restore the land to its pre-mined contours (USDoI 1977). Reclamation activities have traditionally incorporated compaction procedures to augment the strength of the reclaimed material and ensure stability of the restored slopes. However, while compaction is important for strength and erosion resistance, it diminishes soil porosity which restricts root penetration and reduces water infiltration with negative impacts on tree survival and grass reestablishment (Angel et al. 2007, Sweigard et al. 2007). FRA employs a low compaction effort in the uppermost 1.2 m to 1.5 m. The low-compaction grading technique has proven to be successful in encouraging tree growth, and demonstrates the potential for establishing healthy native forests on reclaimed mine lands (Angel et al. 2007, Barton et al. 2007). However, with the exception of Torbert and Burger (1994), most of these demonstrations were conducted on relatively flat lying terrain where stability issues were negligible. The stability of steep FRA slopes, defined as steeper than 20 degrees by the USDoI (2009), and the possible modes of failure have not been investigated and a rational stability analysis method has not been suggested.

Slope stability analysis requires the knowledge of soil properties in terms of density and strength; characteristics not easily determined for reclaimed mine spoil due to the significant amount of oversize material (> 0.3 m). The in situ density of soils consisting of large rock particles can be difficult to measure, which makes it difficult to quantify and awkward to provide proper construction quality control. Sweigard et al. (2007) have suggested correlations between dry bulk density and shovel penetration; though practical for reforestation efforts they are not appropriate for the evaluation of slope stability. Furthermore, because of the difficulties associated with sampling and testing due to the oversize particles, the shear strength properties are not typically measured in laboratory or field tests. For mine reclamation projects, the design is typically completed well in advance of mining activities and usually based on experience using assumed or traditional regional soil properties (Bell et al. 1989). Naturally, there is uncertainty associated with this practice, especially if low compaction is employed on steep reclaimed slopes. For example, in Kentucky the majority of slope failures in abandoned mine lands have occurred via translational and rotational failure mechanisms through the loose material placed prior to the SMCRA (Iannacchione and Vallejo 1995). The lack of proper compaction is a known cause of failure in constructed slopes, with the stability becoming worse under intense rainstorms (Chen et al. 2004). The failure of Sau Mau Ping slopes in Hong Kong is one dramatic example of the danger associated with poorly compacted slopes and lack of proper engineering design (Abramson 1996, Hong Kong Geotechnical Engineering Office 2007).

The objectives of this paper are to 1) characterize the geotechnical properties of low compacted spoils on steep slopes constructed according to the FRA, and 2) investigate the likely failure mechanisms associated to steep slopes reclaimed using the FRA. This is accomplished using three

reclaimed field sites at which the material characteristics are evaluated, and the results will be used to suggest a practical method to estimate the shear strength and evaluate the stability of slopes constructed using the low compaction grading technique.

## Methods

### **Location of field sites**

To investigate the potential effects on stability resulting from the implementation of the low compaction grading technique, three steep FRA slopes were studied. The three sites, referred to here by the name of the initial coal operator (Premium, National and Mountainside), are located in northeastern Tennessee, with Premium located in Anderson County, National in Campbell County and Mountainside located in Claiborne County (Fig. 2.1). Each of the mine operators played an instrumental role in the development of the study sites. Each site was divided into four different plots which while not discussed here, were instrumented in order to concurrently investigate the runoff hydrology and sediment yield on the FRA slopes (Hoomehr et al. 2013). Fig. 2.2 shows the National site during construction of the study plots.

### Site construction and reclamation process

At each of the three sites in this study, the construction procedure followed the contour *haulback method* (Sweigard and Kumar 2010), where a ramp is constructed on the contour bench and spoil is hauled up the ramp and dumped over the edge. The sequence of the construction process can be divided into four major steps (Sweigard et al. 2007) depicted schematically in Fig. 2.3: a) placement and compaction of the materials for the primary backfill core using traditional practices,



**Fig. 2.1.** Location of field sites in northeastern Tennessee, referred to as Premium, National, and Mountainside.



**Fig. 2.2.** National site after the FRA reclamation process and during the construction of the study plots.



Fig. 2.3. Depiction of the reclamation process according to FRA.

b) dumping of the soil that will constitute the loose surface layer (1.2 m-1.5 m thick), c) grading of the loose soil layer with the lightest equipment available using the fewest passes possible, and d) reforestation. The three research sites presented a very rough soil surface after the final grading, which is consistent with the FRA recommendations for successful reforestation (Sweigard et al. 2007). However, because the final layer at all three sites often included boulder-sized material, significant depressions and large rocks were left on the surface of the slope which is a deviation of Sweigard's recommendations for an ideal finished surface.

## **Geotechnical characterization**

The investigation proceeded with the characterization of the research sites and the analysis of their mechanical stability. The field characterization of the mine spoil included: a) determination of the site geometry; b) particle size analysis, index tests, and classification of the materials; c) determination of unit weight; and d) estimation of the Mohr-Coulomb (M-C) shear strength parameters.

## Geometry

The geometric characteristics of the research sites were determined via a series of Trimble<sup>TM</sup> total station topographical surveys. The purpose of these surveys was not only to obtain a more accurate estimation of the slope angles, but also to gather data to geo-reference the instrumented sites allowing an investigation of the spatial variation of material properties.

## Particle size analyses, index tests, and classification of reclaimed materials

Four soil samples of 0.02 m<sup>3</sup> each were randomly taken across the slope at each site for particle size analyses, index tests and classification. All samples were collected from a depth of at least 30 cm below the slope surface to a) avoid samples with fewer fines due to erosion armoring and b) avoid surficial soils affected by changes in fabric due to weathering. Particle size analysis (grain size distribution and hydrometer) and Atterberg limits tests were conducted in general accordance to ASTM D422-07 and ASTM D4318-10 respectively, and classification of the materials in general accordance to the Unified Soil Classification System (USCS), ASTM D2487-10. Visual inspection suggested that the materials may contain a large amount of agglomerated fines in the form of large particles, and therefore, traditional dry particle size analysis would indicate a larger amount of coarse material than really exists. This issue was investigated at all three sites by allowing soil samples to soak in water for 14 days. Very few aggregated fines were found at the Premium and National soils, but significant aggregated fines at Mountainside. For this reason, a wet preparation of the Mountainside samples was conducted in general accordance to ASTM D2217-04 before conducting the particle size analysis and index tests. The amount of oversize material (> 0.3 m) was estimated on a surface basis. This was accomplished by dividing the plot into multiple squares of 1 m side length; a photograph of each square was used to estimate the percentage of oversize particles per surface area.

#### Unit weight

An extensive data collection of the unit weight of the loose surface layer at each of the three sites was conducted using a Troxler 3411-B Nuclear Density Gage (NDG), in general accordance to the ASTM D6938-10. The measurements were obtained shortly after each slope was constructed.

22

During the data collection, periodic calibration of the NDG device was conducted at the beginning and middle of each work day, employing the calibration block provided by the manufacturer. A randomized systematic sampling technique (Sweigard et al. 2007) was used at one plot per site to reduce data tendency or bias in the measurements. The plot was divided into multiple squares of 3 m side length, where single measurements of bulk dry unit weight ( $\gamma_d$ ), wet unit weight ( $\gamma_T$ ) and moisture content (w) were taken at random locations inside the squares (Fig. 2.4). The soil surface was cautiously carved to obtain a planar surface before placing the nuclear gauge device, avoiding air interchange between the gauge base and the soil surface. Then, a hole perpendicular to the slope surface was driven inside each sub-area and the source rod inserted to obtain the readings at 300 mm (which is the maximum length of the source rod of the Nuclear Density Gauge). From previous observations of the materials it was concluded that the amount of hydrocarbons present at each site was very small and unlikely to affect the NDG readings.

#### Shear strength parameters

Slope stability analysis for long-term conditions assumes that positive excess pore water pressure dissipates during the construction or loading period, and thus, requires the estimation of the drained or effective shear strength parameters. The shear strength of soil is typically described using the Mohr-Coulomb (M-C) yield or failure criterion:

$$\tau_n = c + \sigma_n \tan \phi \tag{2.1}$$



**Fig. 2.4.** Randomized systematic sampling technique; a) area of interest subdivided into small subareas (Sweigard et al. 2007c), where small circles represent the random measurement locations; b) application of this technique for NDG measurements at the National site.

where  $\sigma_n$  and  $\tau_n$  are the normal and shear stresses acting on the failure plane within the soil body, and  $\phi$  and c are the internal friction angle and cohesion. Outside of mine reclamation, the shear strength material properties  $\phi$  and c are often measured in laboratory or obtained through correlations with in situ tests. However, due to the large particles present in mine spoils, traditional tests are very difficult to conduct. The current practice in mine reclamation is usually based on experience or assumed values of  $\phi$  and c. As discussed earlier, the FRA technique consists on having a 1.2 to 1.5 m of loose soil at the uppermost part of the contour, above a well compacted and stable core. Thus, this low density/low strength zone should be evaluated for stability. Since the angle of repose is the "steepest stable slope for loose packed granular material and represents the angle of internal friction at its loosest state" (Holtz and Kovacs 1981), it is suggested as a good representation of the internal friction angle of loose soil layer in a FRA slope. The use of observed angles of repose offers the additional advantage that the overall strength of the mass, including the contribution of the oversize particles, is captured. Angles of repose were obtained by observing the placement of spoils piles, and measuring the angle at which they hold in place. Because these tests were conducted during reclamation activities, the angle was measured via photographs of the fresh piles using a hand held level, for safety reasons. The level helped to ensure that the picture was taken parallel to the horizon. With a photo editing software the angles of the piles were measured (Fig. 2.5). Regarding the strength properties of the stronger core material, stability analysis will later show that the most critical condition for stability is insensitive to the selected strength values of the core.



**Fig. 2.5.** Typical field determination of the angle of repose at National Site (White et al. 2009). Camera placed on a level and photo taken of material in loose state. Note the large number of oversize (> 0.3 m) particles in the reclaimed material.

On the other hand, short-term stability analyses (undrained soil conditions) are often conducted to characterize the behavior of the soil during and immediately after construction, where the loading occurs much faster than the rate of dissipation of positive excess pore water pressure. However, it is assumed that the coarse mine spoil material will not develop significant excess positive pore water pressure under typical loadings (Duncan and Wright 2005); thus, a drained response is expected. This assumption is supported by the grain size distribution of the materials and relatively high void ratios obtained for each site which are reported in the results section. Furthermore, short-term stability is not considered to be important in the reclamation of mine slopes, since any short-term failure would have minimum consequences and would be repaired during construction or routine maintenance.

#### Static long-term slope stability analyses

The analyses of the mechanical stability focused on long-term analyses of the primary failure modes that are likely to be experienced by FRA slopes: a) shallow or local failure modes within the loose surface layer and b) global or deep rotational failure modes of the overall soil mass. Limit Equilibrium Methods (LEM) and the Finite Element Method (FEM) were employed in the analyses assuming 2-D plain strain conditions. LEM analyses were computed using Slide 6.0 (Rocscience Inc. 2011) with 10,000 critical surfaces analyzed. The FEM analyses were computed using Phase2 (Rocscience Inc. 2011) employing an elastic perfectly plastic stress-strain behavior and the M-C yield criterion with a non-associated flow rule (zero dilatancy angle) to avoid over-prediction of dilation and failure load for purely frictional materials (Griffiths and Lane 1999). Since the estimation of the factor of safety (FS) has been shown to be not significantly affected by the value of the elastic constants E (Young's modulus) and v (Poisson's ratio) used in the FEM solution (Cheng et al. 2007, Griffiths and Lane 1999), nominal values of  $E = 10^5$  kPa and v = 0.3 were assumed here for the surface and core materials. The model was created with 6-noded triangular elements, with a maximum of 500 iterations solved by Gaussian elimination. Because the geometry and material properties from all sites are reasonably similar, those from Mountainside ( $\beta = 28^{\circ}$ from the horizontal and H = 21 m height) will be used here to explore the various failure modes. These values are representative of all three sites and many steep slopes in the southern Appalachian coal fields. The thickness of the low-strength layer (z) was assumed to be 1.5 m.

## Shallow stability within the low strength surface layer

A shallow failure mode was investigated by assuming that the core was significantly stronger than the surface layer. The methods used were: a) LEM restricting the analyses to the shallow surface layer via the non-circular Janbu's method and a search block feature, b) FEM and the shear strength reduction method (Cheng et al. 2007, Griffiths and Lane 1999), and c) the infinite slope equation for cohesionless soils without seepage (FS =  $\tan \phi / \tan \beta$ ) and seepage (FS  $\approx 0.5 \tan \phi / \tan \beta$ ). The infinite slope equation idealizes the surface as an infinite plane with the failure mechanism running parallel to the surface (Skempton and Delory 1957) and it is appropriate when the ratio of depth to length of the sliding surface is small. The geometry of the reclaimed mine slopes constructed according to the FRA is ideally suited for investigation by the infinite slope method.

#### Deep rotational stability of the overall soil mass

Because the strength of the compacted core was not known, the deep rotational failure mode was investigated by performing a series of analyses where the strength of the loose surface material was held constant ( $\phi = 38^{\circ}$ , c = 0) while the internal friction angle of the core was increased. The analyses started with a homogeneous slope with the properties of the weak, loose surface layer analyzed via LEM (the circular Simplified Bishop's method) and FEM. Then, the shear strength of the core was increased such that the ratio  $\tan \phi_{core} / \tan \phi_{loose}$  was equal to 1.1, 1.2, 1.3, and 1.4 ( $\phi_{core}$  is the friction angle of the core material and  $\phi_{loose}$  is the friction angle of the loose surface layer.

## **Results and Discussion**

## Geotechnical characterization of research sites

### Geometry

The geometric information of the three sites obtained from the topographical survey is summarized in Table 2.1. Preliminary information of the slopes angles via a Suunto PM-5/360PC mechanical inclinometer reported angles between 26 and 30 degrees at Premium site, 20 and 22 degrees at National site, and 28 and 29 degrees at Mountainside site (White et al. 2009). These angles, collected shortly after the end of the reclamation process, coincide with the topographic information gathered 15 months later, and suggest no changes in the slope morphology and no slope failures during the study period.

## Particle size analysis, index tests, and classification of reclaimed materials

Results from the soil analyses are reported in Table 2.2. The grain size distribution was conducted on material smaller than 51 mm (2 in. sieve), while Atterberg limits were determined on material smaller than 0.42 mm (No. 40 sieve). According to the USCS, for all the research sites the material

Site	Average slope angle, β	Average slope length	Average slope width (m)		
	(degrees)	( <b>m</b> )	Тор	Bottom	
Premium	28	32.2	28.1	25.0	
National	20	48.4	22.4	25.4	
Mountainside	28	45.4	23.6	23.1	

Table 2.1 Average slope length, width and inclination angle for the four plots at the Premium,

National and Mountainside sites

Table 2.2 Mean values of Liquid Limit, Plastic Index, soil texture and soil classification (USCS)

Sites	Gravel particles 51 mm – 4.75 mm (%)	Sand Particles 4.75 mm - 0.075 mm (%)	Fines < 0.075 mm (%)	Clay particles < 2 µm (%)	Liquid Limit (LL)	Plastic Index (PI)	Soil classification (USCS)
Premium	59	28	13	6	29	13	GC to GP-GC
National	52	28	20	10	27	14	GC
Mountainside	37	22	41	19	32	15	GC

classify as clayey gravel (GC) with the exception of one plot at the Premium site that classifies as poorly graded clayey gravel (GP-GC) due to slightly less material finer than the number 200 sieve. Regarding oversize particles, it was estimated that material larger than 300 mm occupies 0 to 25% of 1 m<sup>2</sup> at Premium site, 0 to 10% of 1 m<sup>2</sup> at National site, and 5 to 40% of 1 m<sup>2</sup> at Mountainside site.

## Unit weight

Results of the unit weight measurements were as follows: the maximum and minimum measured  $\gamma_d$  were 18.8 and 13.0 kN/m<sup>3</sup> at Premium site, 21.4 and 14.6 kN/m<sup>3</sup> at National site, and 22.8 and 14.9 kN/m<sup>3</sup> at Mountainside site. Complementary results from statistical analyses are presented in Table 2.3. Field measures of density at a similar mine site in Kentucky (Sweigard et al. 2011) indicated similar variations as those found at Premium and National. The largest standard deviations (S.D.) were observed at Mountainside, which is consistent with the largest range of unit weights, and the largest amount of fines and observed number of oversize particles. At all sites, the spatial variation of unit weight reflects the large range of particle sizes in these reclaimed materials. A single unit weight measurement has a limited ability to represent the state of body forces acting on the complete FRA slope, and thus, mean unit weights with the probable upper and lower bounds (confidence and tolerance intervals) are desired for the mine material characterization. While confidence intervals (C.I.) provide an upper and lower bound of the true mean found at the constructed sites, tolerance intervals (T.I.) provide information of the probable future range of unit weights that each site will have on average. In any case, as discussed later, the determination of the unit weight for static long-term conditions may be of minor concern, but necessary for static unsaturated and seismic stability analyses.

Table 2.3 Mean	s, standard dev	iations, 95%	confidence	intervals (	C.I.), and	90% to	lerance
intervals (T.I)	(80% coverage	e) for wet and	l dry unit w	eights for F	Premium, I	Nation	al and

				95% C.I. for the Mean		90%/0.8 T.I. for the Mea		
Sites	Unit Weight	Mean kN/m <sup>3</sup>	S.D. kN/m <sup>3</sup>	Lower, kN/m <sup>3</sup>	Upper, kN/m <sup>3</sup>	Lower, kN/m <sup>3</sup>	Upper, kN/m <sup>3</sup>	
Premium	Dry	16.2	1.3	15.8	16.5	14.2	18.1	
	Wet	18.5	1.3	18.2	18.8	16.6	20.4	
National	Dry	18.5	1.0	18.3	18.7	17.2	19.9	
Inational	Wet	20.3	1.0	20.1	20.5	18.9	21.7	
Mountainside	Dry	18.6	2.2	18.1	19.1	15.5	21.7	
	Wet	20.4	2.2	19.9	20.9	17.2	23.6	

# Mountainside sites

Water and sand cone replacement tests were previously conducted to determine bulk unit weights at random locations on the four plots at each site (White et al. 2009). A comparison of the results indicates that the NDG device gives on average about 25% higher average unit weights than replacement methods at Premium site, 14% at National site and 21% at Mountainside. Since both replacement methods involved the removal of small samples, they did not take into account the effects of large rock fragments that are randomly embedded in the loose surface soil layer. On the other hand, the NDG calculates unit weights based on the velocity travel of gamma rays between the source and the detector, and any denser material that appears on the travel path will be counted in the measurement. In this regard, the collection of a sufficient amount of NDG readings will better represent the wide range of in-place density and provides a more representative average unit weight for stress analyses. Replacement methods may be preferable for the calculation of void ratio and soil porosity due to better representation of the soil matrix. The average void ratio of the loose surface layer calculated via replacement methods was 1.0 at Premium, 0.6 at National and 0.7 at Mountainside. The largest void ratio was calculated for the soils found at Premium, which is consistent with the lowest NDG unit weight measured. Overall, relatively large void ratios were obtained for all three sites, which is consistent with the FRA requirements for healthy tree growth.

## Shear strength parameters

The angle of repose of the looser soil layer was found to range between 37 and 39 degrees at Premium and Mountainside, and between 36 and 38 degrees at National site. Zero cohesion is usually employed for long-term analysis on coarse granular soils (Holtz and Kovacs 1981, Lambe and Whitman 1969) and normally consolidated fine soils (Skempton 1964), and would be appropriate for reclaimed materials receiving minimum compaction effort. While even a small amount of compaction will increase the density and strength of the soil, the angle of repose is a conservative estimate of the friction angle. Similar values of  $\phi$  and c for loose spoils in the Appalachian region were reported by Sweigard et al. (2011), while similar values for reclaimed spoils outside the Appalachian were found in the literature (Gutierrez et al. 2008, Kasmer and Ulusay 2006, Stormont and Farfan 2005, Sweigard et al. 2011, Ulusay et al. 1995) as summarized in Table 2.4.

#### Static long-term slope stability analyses

## Shallow stability within the low strength surface layer

Results from the limit equilibrium, finite element, and infinite slope analyses are summarized in Table 2.5. From a practical perspective, all the analyses yielded very similar FS's (approximately 1.47), implying that the shear strength along the most critical slip surface is about 47% greater than that required to maintain equilibrium in the long-term. For all cases, the most critical failure mechanism is shallow and is consistent with the assumed failure mechanism in the infinite slope method. Fig. 2.6 shows the FEM model of the shallow failure mode with a section of the slope enlarged (the 1.5 m thick surface layer is small with respect to the size of the model and may not be clearly distinguished in the full model). It also shows nodal displacement vectors. Larger strains are observed at the interface of the weak surface and core materials, with the displacement vectors acting parallel to the surface suggesting a planar failure mechanism. The obtained long-term FS's are valid for drained conditions in the absence of seepage forces due to transient flow. However, since the occurrence of downslope water flow through the complete thickness of the loose layer is highly unlikely, this condition represents a lower bound or worst case value for the stability of FRA, and would reduce the FS by a factor of 2.

Author	Origin of Material Tested	Type of Test	Sample Dimensions, mm	Internal friction angle $\phi$ , Degrees	Cohesion c , kN/m <sup>2</sup>
Ulusay et al. (1995)	Limestone, Claystone and Marl (Turkey)	In situ SPT Test	N/A	31-38	N/A
Ulusay et al. (1995).	Limestone, Claystone and Marl (Turkey)	Direct Shear Test	N/A	34 (peak) 33 (residual)	12 (peak) 9 (residual)
Ulusay et al. (1995)	Limestone, Claystone and Marl (Turkey)	Triaxial (CD) Test	Diameter= 191 Height = 382	23-35	0 - 10
Stormont and Farfan (2005)	N/A (San Juan, Colorado)	Direct Shear Test (Large Laboratory Box)	Length = 762 Width = 762 Height = 457	37	5
Gutierrez et al. (2008)	N/A (Northern New Mexico)	Direct Shear Test	Length = 51 Width = 51 Height = N/A	42-47 (peak) 37-41 (residual)	0
Kasmer and Ulusay (2006)	Limestone and marl (Turkey)	Direct Shear Test	N/A	31-34 (peak) 24-33(residual)	18-34 (peak) 6-10 (residual)
Sweigard et al. (2011)	Sandstone and Shale (Pike County, Kentucky)	Triaxial (CU) Test	N/A	37	0
FRA research sites (this study)	Sandstone and Shale (Northeast Tennessee)	Angle of repose	N/A	38	0

Table 2.4 Summary of internal friction angle and cohesion for reclaimed mine materials

slope stability analysis results for generic slope ( $\beta = 28^\circ$ , $H = 21$ m, $\phi = 38^\circ$ , $c = 0$ , $\gamma_T = 20.4$ kN/m <sup>3</sup> )						
Analysis Method	Assumptions	FS	Critical failure mode			
a) LEM	Rigid core and Search Block - non-linear Janbu's Method with 10,000 critical surfaces analyzed	1.48	Shallow planar failure surface			
b) FEM	Core much stronger than loose surface layer, Shear Strength Reduction Method to determine FS, with 500 iterations solved by Gaussian elimination	1.47	Shallow planar failure surface			
c) Analytical	Infinite slope equation (no seepage)	1.47	Shallow planar failure surface			

Table 2.5 FS obtained for long-term static stability focused on the low strength surface layer



**Fig. 2.6.** Shear strains and nodal displacements obtained from the FEM analysis assuming a very strong core. Upper left corner illustrates the geometric dimensions employed for LEM and FEM analyses.

## Deep rotational stability of the overall soil mass

The results from LEM and FEM analyses of the deep failure mode of the homogeneous slope yield a FS = 1.48, which is consistent with that obtained from the shallow stability analyses. Fig. 2.7 shows the results of a FEM analysis when the core strength was 30% stronger than the loose layer (i.e.  $\tan \phi_{core} / \tan \phi_{loose} = 1.3$ ). Here two possible failure mechanisms were observed in the form of shear bands; a deeper mechanism through the core material with a FS = 1.94, and a shallow mechanism with the lowest FS = 1.48 and highest shear strains at the interface of the materials. The FS of the shallow mechanism is equal to those obtained from the shallow analyses above. A similar trend is observed for cases when  $\tan \phi_{core} / \tan \phi_{loose} = 1.1$ , 1.2, and 1.4 (Fig. 2.8). As the strength of the core increases, the FS of the deeper mechanism increases; however, the lowest FS is found to be constant with a consistent shallow failure mode and dependent only on the strength level of the loose surface layer. Additional analyses at angles of inclination of 20 and 35 degrees yielded similar results and confirm that the critical failure mode is a surface failure.

These results are consistent with the observation that the failure mechanisms through the loose surface layer will govern and the determination of the strength parameters of the stronger dense core are not important for FRA slope design. Furthermore, since the infinite slope method adequately approximates the shallow failure mode, and accurately predicts the FS, it can be taken as a simple and reliable method to evaluate the performance of FRA slopes and more sophisticated computer analyses are not necessary for most applications. The use of the infinite slope method also simplifies the field characterization of materials and disregards the unit weight determination, because it only requires  $\phi_{loose}$  and  $\beta$  for long-term conditions. The simplicity of the method is



**Fig. 2.7.** Failure mechanisms and FS's obtained from FEM stability analysis for  $\tan \phi_{core} / \tan \phi_{loose}$ = 1.3. The FS = 1.94 shown for the deeper failure mechanism was obtained when the strength reduction factor (SRF) search was restricted to be outside the zone where the shallow mechanism occurred. Shear strains showed for SRF = 2.02 to emphasize failure mode.



**Fig. 2.8.** FEM analyses of the global stability for various values of  $\tan \phi_{core} / \tan \phi_{loose}$  and the infinite slope equation

appropriate for design of reclaimed mine slopes which are typically designed in advance of mineral extraction with assumed overburden properties. Accordingly, the lowest FS for drained or long-term conditions at each instrumented site using the infinite slope equation is approximately 1.47 for Premium, 2.07 for National and 1.47 for Mountainside.

# Conclusions

- Characterization and stability evaluation of three FRA slopes in northeastern Tennessee with inclinations as high as 28 degrees were conducted. A large number of oversize particles were found in the reclaimed materials. In general, the material finer than 51 mm classified as Clayey Gravels with the average Plasticity Index (PI) ranging from 13-15, suggesting that the physical characteristic of the soils are similar across the three research sites.
- Unit weights determined using a Nuclear Density Gage were found to be higher than those
  determined by replacement methods, yet vary significantly across the study plots. NDG
  measures are preferred for stability analyses because they better capture the effect of
  oversize particles on the in situ state of stresses of FRA slopes. It also allows more
  measurements to characterize the wide range of in-place density. Tolerance intervals were
  constructed to reflect the probable future range of unit weights that each site will have on
  average.
- The analysis of several potential modes of failure suggests that the governing failure mode is shallow and contained within the weak, loose surface layer. The determination of the strength parameters of the core is not important for FRA slope design.

- Because the infinite slope method adequately approximates the shallow failure mode and accurately predicts the FS, it may be an appropriate method to evaluate the performance of FRA slopes and more sophisticated analyses are not necessary for most applications. Since the unit weight of the material is not considered in the infinite slope expression, field measurements of the highly variable unit weight are not required for long-term analyses.
- The angle of repose was suggested to be a conservative estimate of the internal friction angle and it is consistent with the loose nature of the FRA material. This provides a means to quantify the friction angle of the mine spoil, which has been traditionally assumed based on experience.
- The shear strength along the most critical slip surface, for the typical FRA slope investigated, is at least 47% greater than that required to maintain static equilibrium in the long-term. In case that the entire loose surface zone becomes saturated with downslope seepage and no infiltration into the core, the FS is reduced by a factor of 2, suggesting that the slope would be unstable. However, these conditions are very unlikely and provide a lower limit to the factor of safety
- The likely conditions would suggest that the FRA has no negative impact on slope stability, and the benefits of faster forest establishment in terms of reduced erosion and sediment delivery make the FRA very attractive for future reclamation work.

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## Abstract

While traditional mine reclamation methods emphasize compaction to increase the strength of the materials and ensure stability of the restored slope, high compaction restricts the successful reforestation of reclaimed mine sites. The Forest Reclamation Approach (FRA), which uses low compaction in the uppermost 1.2 m - 1.5 m of the surface has been shown to facilitate the establishment of healthy native forests. Slope stability analyses of three steep FRA sites from the southern Appalachian region have shown that the long-term static stability is not compromised, and that the infinite slope method provides a rational method to evaluate the stability of steep FRA slopes. In this article, modifications of the infinite slope equation were utilized to a) include the effects of matric suction due to unsaturated soil conditions, and b) evaluate the seismic performance of FRA slopes based on spectral accelerations. Monthly variation of the water content at three research sites demonstrated the seasonal stability variation of FRA slopes due to matric suction, while seismic analyses illustrated the conditions under which instability may occur.

## Introduction

The Forest Reclamation Approach (FRA) (Angel et al. 2007, Sweigard et al. 2007) is a method of mine slope reclamation which maintains a loose surface layer (1.2 to 1.5 m thick) above a well compacted core to facilitate tree growth. This method has been shown to be successful in promoting healthy reforestation (Angel et al. 2007, Barton et al. 2007) and is currently being promoted by the US Office of Surface Mining (OSM). However, there has been limited experience with the FRA on steep slopes, which the OSM defines as those inclined at more than 20 degrees (USDoI 2009). Jeldes et al. (2013) conducted a series of static slope stability analyses (including effects of seepage water forces) via Finite Element Method (FEM) and Limit Equilibrium Methods (LEM), and concluded that: 1) the static long-term stability of FRA slopes with inclinations as much as 28 degrees is not compromised; 2) shallow failure modes inside the loose surface layer are the dominant failure mechanism, regardless of the strength level of the core; and 3) the infinite slope method using the angle of repose as a conservative estimate of the shear strength provides a rational method to evaluate the stability of steep FRA slopes.

The theoretical idealizations for the static long-term stability overlook two field conditions that affect the strength and the stress state of slopes: partial saturation of soils and seismicity. Specifically, the long-term stability involves the analysis of saturated soils under drained loading, which is not consistent with field conditions and the seasonal variations of soil water content. In reality, the solid and liquid soil phases cohabit with air and, in most cases of geotechnical interest, the water pore pressure can be negative with respect the air pore pressure, resulting in higher soil strengths (Fredlund and Rahardjo 1993). On the other hand, static analyses do not include seismic-induced forces that, if not accounted for, may trigger slope failures by increasing the shear stresses

and, sometimes, decreasing the strength (Abramson 1996). The Appalachian region is recognized as one of the most seismically active areas in the eastern North America (Chapman et al. 1997), with a concentration of events occurring in eastern Tennessee of magnitudes even larger than 4.5 (Bollinger et al. 1976, Chapman et al. 1997). Therefore, it becomes important to investigate the seismic effects on steep FRA slopes.

The objectives of this paper are to 1) investigate the implications of unsaturated soils on steep FRA slopes and 2) investigate the seismic response of steep FRA slopes. This is accomplished using three reclaimed field sites at which the geometric and material characteristics were determined (Jeldes et al. 2013). To evaluate the stability of unsaturated FRA slopes, an expansion of the infinite slope equation to include the effects of matric suction is proposed here. Similarly, to evaluate the seismic response of FRA slopes, an infinite slope equation which includes the spectral response is developed with solutions presented in terms of charts for design applications.

## Background

#### Unsaturated soil shear strength

The total shear strength of an unsaturated soil has been typically approached as an extension of traditional saturated soil mechanics, where matric suction and net normal stress are considered to be independent stress variables, and consequently the Mohr-Coulomb (M-C) failure criterion for unsaturated soils can be extended as (Fredlund et al. 1978):

$$\tau = c + (\sigma - u_a) \tan \phi + (u_a - u_w) \tan \phi^b$$
(3.1)

where c and  $\phi$  are the cohesion and the internal friction angle (effective strength parameters),  $u_a$ is the pore air pressure,  $u_w$  is the pore water pressure,  $(\sigma - u_a)$  is the net normal stress,  $(u_a - u_w)$ the soil matric suction, and  $\phi^b$  is the rate of increase of shear strength due to matric suction. Although this expression is widely used, it has been pointed out to be philosophically erroneous, since the matric suction is not a stress variable and transfer functions are required to upscale the negative pore pressure to a macroscopic level (Lu 2008). In this regard, an extension of the M-C criteria based on the classical effective stress expression  $\sigma' = (\sigma - u_a) + \chi(u_a - u_w)$  (Bishop 1959) may be more appropriate, since the effective stress parameter  $\chi$  is indeed a transfer function:

$$\tau = c + \left[ \left( \sigma - u_a \right) + \chi (u_a - u_w) \right] \tan \phi$$
(3.2)

Naturally,  $\chi$  is strongly related to the degree of saturation of the soil mass and it ranges from 0 (completely dry material) to 1 (complete saturation). Theoretical relationships to relate effective stresses and the degree of saturation have been developed for ideal arrangements of spherical particles (Cho and Santamarina 2001), which assume a matric suction only dependent on capillarity and valid for cases with degree of saturation less than 25% (Lu and Griffiths 2004). For higher degrees of saturation some empirical expressions have been proposed and reviewed (Vanapalli and Fredlund 2000) with the one based on the Vanapalli et al. (1996) approach being of particular interest since it allows the relationship between effective stresses and matric suction via the Soil Water Characteristic Curve (SWCC),

$$\chi = \left(\frac{\theta - \theta_r}{\theta_s - \theta_r}\right) \tag{3.3}$$

where  $\theta$  is the volumetric water content (volume of water over total volume of soil) at some level of matric suction,  $\theta_r$  is the residual volumetric water content, and  $\theta_s$  is the volumetric water 51 content at saturation. This expression [Eq. (3.3)] was shown to provide reasonable estimates of the unsaturated shear strength for suctions in the range of 0 – 1500 kPa (Sheng et al. 2011, Vanapalli and Fredlund 2000), but to be sensitive to small changes in residual water content (Sheng et al. 2011). On the other hand, the use of the SWCC to estimate the shear strength of unsaturated soils has been suggested as an adequate framework for stability problems in soil mechanics (Lu and Griffiths 2004) and also to overcome conceptual limitations regarding the stress suction (Lu 2008, Lu 2010, Lu and Likos 2006). Accordingly, Eqs. (3.2) and (3.3) will be employed here to investigate the stability of three steep FRA slopes under partially saturated conditions in the Appalachian region.

#### Shear strength and the pseudo-static forces for seismic analyses

For seismic analyses, the saturated-undrained shear strength is typically used since seismic excitation may be considered a short-term loading condition, and excess of positive pore water pressure may develop. However, it is assumed that the coarse, low density material resulting from the FRA compaction methods would drain appropriately, and the probability of significant development of excess positive pore water pressure during an earthquake event is low (Duncan and Wright 2005), especially during seasons of low water content where negative pore water pressure arises. Consequently, drained strength parameters will be employed here for dynamic slope stability analysis.

The pseudo-static approach remains as one of the most widely used methods for addressing seismic hazards in civil engineering practice, due to its simplicity and practicality (Duncan and Wright 2005). This method is based on d'Alembert's principle of mechanics, which assumes that

equilibrium may be achieved by adding a fictitious static inertial force equal to the maximum acceleration of the body, with the inertial force expressed as a fraction of the acceleration of gravity (g) (named pseudo-static coefficient K) times the mass of the structure (Paz 1997). However, the maximum horizontal acceleration experienced during an earthquake is usually not equivalent to the horizontal pseudo-static coefficient  $K_h$  (Towhata 2008), and therefore, many authors have proposed either arbitrary values of maximum acceleration for design or reduction factors for the maximum horizontal Peak Ground Acceleration  $PGA_h$  (Table 3.1). Regarding the vertical coefficient  $K_v$ , it has been usual practice to either assume a fraction of  $PGA_h$  or neglect its contribution. Examples of recorded seismic activity (e.g. Towhata, 2008) show that the maximum vertical acceleration can be assumed 0.5  $PGA_h$  with little error, while neglecting its effect may lead to unstable slopes.

More recently, a procedure to calculate  $K_h$  based on the 5% damped elastic spectral acceleration at the degraded fundamental period  $S_a$ , the maximum allowable displacement  $D_a$  (in cm), the initial fundamental period  $T_s$ , the earthquake magnitude M, and a random normal distributed variable  $\varepsilon$  was proposed (Bray and Travasarou 2009):

$$K_{h} = e^{\frac{-a + \sqrt{b}}{0.665}} \tag{3.4}$$

$$a = 2.83 - 0.566 \ln(S_a) \tag{3.5}$$

$$b = \begin{cases} a^{2} - 1.33 \left[ \ln(D_{a}) + 1.10 - 3.04 \ln S_{a} + 0.244 \ln^{2}(S_{a}) - 1.5T_{s} - 0.278(M - 7) - \varepsilon \right], T_{s} \ge 0.05 \\ a^{2} - 1.33 \left[ \ln(D_{a}) + 0.22 - 3.04 \ln S_{a} + 0.244 \ln^{2}(S_{a}) - 1.5T_{s} - 0.278(M - 7) - \varepsilon \right], T_{s} < 0.05 \end{cases}$$
(3.6)

	Seismic Coefficient K						
Method	Accel.	Accel. reduction	Cyclic degrade reduction	Min. FS	Max. displ. (m)	Requirements/ comments	
	0.10g	-	-	>1.00	-	-	
Terzaghi (1950)	0.25g	-	-	>1.00	-	-	
	0.50g	-	-	>1.00	-	-	
Noda et al. (1975)	$PGA_h$	$PGA_{h}^{1/3} / 3$	-	-	-	-	
Makdisi and Seed (1978)	0.20g	0.50	0.80	1.15	$\approx 1.0$	For M≈6.5	
	0.75g	0.20	0.80	1.15	$\approx 1.0$	For M≈8.25	
Seed (1979)	0.1-0.15g	-	-	-	-	Earth dams in several countries	
Hynes-Griffin and Franklin (1984)	$PGA_{\rm h}$	0.50	0.80	>1.00	1.0	-	
Kavazanjian et al. (1997) (*)	$PGA_h$	0.17	0.80	>1.00	1.0	Response analysis	
	$PGA_h$	0.50	0.80	>1.00	1.0	No response analysis	
Bray and Rathje (1998)	$PGA_h$	0.75	-	>1.00	0.15-0.30	Conservative strength	

**Table 3.1** Summary of recommended pseudo-static values and expressions available in theliterature (after Duncan and Wright, 2005)

(\*) Acceleration reduction factor reported for soil conditions. Strength reduction factor only for saturated or sensitive clays

The form of  $K_h$  in Eq. (3.4) was obtained from a semi-empirical probabilistic approach originally developed to determine permanent displacements (Bray and Travasarou 2007). The relationship between induced displacement and a single value of motion intensity was found to be optimally satisfied not by *PGA*, but by the 5% damped elastic spectral acceleration at the degraded fundamental period,  $S_a$ . The degraded period was found to be well represented by  $1.5 T_s$ , due to material non-linearity. Both, the traditional pseudo-static coefficients listed in Table 3.1 as well as Eqs. (3.4) – (3.6) will be employed to investigate the seismic stability of three steep FRA slopes constructed in the Appalachian region.

## Field and experimental methods

#### Study sites: construction and site characterization

### Location of field sites and construction process

To investigate the effects of the low-compaction grading technique on the mechanical stability of steep slopes (> 20 degrees), three reclaimed mine sites in northeastern Tennessee were constructed (Fig. 3.1), referred to here by the name of the initial coal operator (Premium, National and Mountainside). Each site was instrumented with weather stations to concurrently investigate the runoff hydrology and sediment erosion on the FRA slopes, as described elsewhere (Hoomehr et al. 2013).

At each of the three sites in this study, the construction procedure followed the contour *haulback method* (Sweigard and Kumar 2010), where a ramp is constructed on the contour bench and spoil is hauled up the ramp and dumped over the edge. The sequence of the construction process can be



**Fig. 3.1.** Location of field sites in northeastern Tennessee, referred to as Premium, National, and Mountainside.

divided into four major steps (Sweigard et al. 2007) depicted schematically in Fig. 3.2 (Jeldes et al. 2013): a) placement and compaction of the materials for the primary backfill core using traditional practices, b) dumping of the soil that will constitute the loose surface layer (1.2 -1.5 m thick), c) grading of the loose soil layer with the lightest equipment available using the fewest passes possible, and d) seeding of cover vegetation and reforestation.

#### Geometry, unit weight and saturated shear strength parameters

The slope lengths, widths, and angles of inclination of the research fields were determined using a total station instrument, while unit weights and gravimetric water contents were measured in situ through water and sand replacement methods and Nuclear Density Gauge (NDG) readings. A randomized systematic sampling technique was employed for the data collection, as described by Jeldes et al. (2013). Due to the extreme range of particles size (from clay particles to boulders), the determination of the shear strength parameters is difficult to conduct in mine spoil. Since by definition the angle of repose represents the friction angle at the soil's loosest state (Holtz and Kovacs 1981), it was chosen as a conservative estimate of the friction angle  $\phi$  for the surface layer which received minimum compaction effort. Because the design of these reclaimed slopes seldom includes laboratory testing, it is suggested that the use of the angle of repose also captures the overall strength of the mass, including the contribution of the oversize particles. A negligible value of cohesion was assumed. Jeldes et al. (2013) provide a comparative summary of the employed strength parameters and values reported in the literature. Similar values of  $\phi$  and c for loose spoils in the Appalachian region were reported by Sweigard et al. (2011), while similar values for reclaimed spoils outside the Appalachian are reported in the literature (Gutierrez et al. 2008, Kasmer and Ulusay 2006, Stormont and Farfan 2005, Sweigard et al. 2011, Ulusay et al. 1995).



Fig. 3.2. Depiction of the reclamation process according to FRA (Jeldes et al. 2013).

Table 3.2 summarizes the obtained slope lengths, slope angles, unit weights and angles of repose for the Premium, National and Mountainside sites.

#### Seasonal variation of volumetric water content

The field variation of volumetric water content  $\theta$  was obtained (Aubuchon 2010) for each of the three research sites at various depths in the loose layer, up to a maximum depth of 0.76 m using a capacitance sensor (Aqua Pro-Sensors LLC 2012) moisture probe. Since capacitance sensors can provide a good indication of seasonal variations in  $\theta$ , but may not yield good absolute values (Leib et al. 2003), the measured monthly variations of  $\theta$  taken from a depth of 0.15 m were corrected relative to the measured July 2009  $\theta$  values from the NDG and used to obtain monthly values of  $\theta$  for the range of depths. Corrected values of  $\theta$  from a depth of 0.76 m were then used to determine the variation in strength and stability for each FRA site during the first year of data collection.

## Soil water characteristic curves for unsaturated stability analyses

For each of the three sites, SWCCs were obtained via: a) suction table and pressure plate laboratory tests conducted in general accordance with Dane and Hopmans (2002) and the ASTM D6836-02; and b) the neural network ROSETTA model (Schaap et al. 2001).

Laboratory suction experiments were conducted on representative soil samples obtained from the field, which were reconstructed to densities of  $1.65 \text{ Mg/m}^3$  for Premium,  $1.66 \text{ Mg/m}^3$  for National, and  $1.57 \text{ Mg/m}^3$  for Mountainside. Two sets of suction table experiments, at pressures between 0.5 and 10 kPa, were conducted: a) with material smaller than 12.7 mm (1/2" sieve), and b) with

Sites	Slope length (m)	Slope angle (Degrees)	Dry unit           weight,           γ <sub>d</sub> (kN/m <sup>3</sup> )	Wet unit weight, $\gamma_T$ (kN/m <sup>3</sup> )	Gravimetric water content, w (%)	Volumetric water content, θ (%)	Angle of repose, φ (Degrees)
Premium	32	28	16.2	18.5	14.5	23.9	38
National	48	20	18.5	20.3	9.7	18.2	37
Mountainside	45	28	18.6	20.4	9.8	18.8	38

**Table 3.2** Average values of slope length, inclination angle, unit weights, water content, andobserved angles of repose for Premium, National and Mountainside sites.

material smaller than 3.36 mm (# 6 sieve). The sample size in both aforementioned cases was 102 mm in diameter and 64 mm height. Pressure plate experiments at pressures between 20 and 1100 kPa were conducted only on material smaller than 3.36 mm (# 6 sieve), with a sample size of 50 mm in diameter and 10 mm height. The pressure plate results were then adjusted to reflect the absence of gravel size particles, by subtracting the equivalent volume of water that would have been occupied by these oversize materials. In addition, the neural network ROSETTA model (Schaap et al. 2001) was used to determine SWCCs using transfer functions based on the van Genuchten (1980) equations. The input parameters used for the ROSETTA model are summarized in Table 3.3, with gravels included with the sand as the coarser fraction.

#### Peak ground and spectral accelerations for seismic stability analyses

Local spectral acceleration values as a function of the earthquake period (Fig. 3.3) and values of *PGA* were estimated from local hazard maps developed by the U.S. Geological Survey, USGS (<u>http://earthquake.usgs.gov/earthquakes</u>), for a 2 and 10% of probability of exceedance (P.E.) in 50 years. The three sites have small variations of *PGA* thus the maximum among them (0.2g for a 2% P.E. in 50 years and 0.07g for 10% P.E. in 50 years) will be used for stability analysis.

## Analytical methods for stability analyses

#### Infinite slope equation for unsaturated soil strength

The infinite slope method is a limit equilibrium analysis in which the failure surface is assumed to be parallel to the ground surface, at a depth that is small with respect to the length of the slope (Skempton and Delory 1957). The assumptions behind this method match well with the geometry of slopes constructed according to the FRA, because a surface layer of loose material runs

	Soil Fraction (%)				
	Premium	National	Mountainside		
Coarse (sand + gravel)	87.0	79.5	58.5		
Silt	7.3	11.0	22.6		
Clay	5.7	9.5	18.9		
Dry unit weight (Mg/m <sup>3</sup> )	1.65	1.66	1.57		

Table 3.3 Soil texture (USDA) and unit weights for ROSETTA model.



Fig. 3.3. Depiction of the 3 Spectral Accelerations for Premium, National and Mountainside (USGS 2012).

approximately parallel to the contour of the strong, dense core (Fig. 3.4). Also, the ratio of horizontal length to depth of the failure surface at each of the three research slopes is large enough so that they can be considered infinite (Griffiths et al. 2011, Milledge et al. 2012). FEM and LEM analyses conducted by Jeldes et al. (2013) demonstrated the appropriateness of the infinite slope method for static long-term analysis of steep FRA slopes, and therefore, it is extended here to include the effects of partially saturated soils. Fig. 3.4 shows a slip surface of thickness *z* below the ground surface inclined at  $\beta$  degrees from the horizontal. The weight of the slice is  $W = \gamma_T bz / cos\beta$ , where  $\gamma_T$  is the total unit weight of the soil, *b* is the horizontal width of the slice that can be expressed as  $b = l cos\beta$ , and *l* is the length of the corresponding slip segment. *P* and *T* are the corresponding normal and shear resultant forces at the bottom of the slice. Neglecting the effects of the side forces  $Q_L$  and  $Q_R$  and satisfying equilibrium of forces in a free diagram space, the factor of safety (FS) defined here as the ratio of the magnitude of resisting forces of the slide to the magnitude of the destabilizing forces becomes:

$$FS = \frac{cl + P'\tan\phi}{T} = \frac{c + (u_a - u_w) \left[ \tan\phi \left(\frac{\theta - \theta_r}{\theta_s - \theta_r}\right) \right]}{\gamma_T z \sin\beta} + \frac{\tan\phi}{\tan\beta}$$
(3.7)

where P' is the force corresponding to the effective stress written in terms of Bishop's formulation. In the presence of pure coarse material, Eq. (3.7) is only valid for  $\theta \ge \theta_r$ , since below this point capillarity decreases radically (Lu and Likos 2006). Similar expressions of the infinite slope equation have been proposed by Uchaipichat (2012) and Lu and Godt (2008). While the Uchaipichat (2012) equation incorporated the relationship between the effective stress parameter



Fig. 3.4. The Infinite Slope method. Figure modified from Salgado (2008).

and the matric suction proposed by Uchaipichat and Man-Koksung (2011), Lu and Godt (2008) assumed a form of stress suction that is a function of the infiltration rates, saturated hydraulic conductivity and empirical fitting parameters obtained from van Genuchten (1980). Eq. (3.7) provided here shares the same fundamental principles as the Lu and Godt (2008) approach, since the empirical model proposed by van Genuchten (1980) is directly related to  $(\theta - \theta_r)/(\theta_s - \theta_r)$ .

#### Infinite Slope equation for horizontal and vertical pseudo-static forces

Analogous to the formulation for the partially saturated soils, when horizontal and vertical pseudostatic forces ( $K_hW$  and  $K_vW$ ) are applied at the center of mass (Fig. 3.5), the following seismic infinite slope equation is obtained:

$$FS = \frac{c + \gamma_T z [(1 - K_v) \cos \beta - K_h \sin \beta] \tan \phi}{\gamma_T z [(1 - K_v) \sin \beta + K_h \cos \beta]}$$
(3.8)

Note that when the effects of the vertical ground acceleration are neglected, Eq. (3.8) simplifies to an equivalent form of the equation proposed by Duncan and Wright (2005). If  $K_h$  in Eq. (3.8) is forced to take the form defined by Eqs. (3.4) – (3.6), a single expression for the seismic infinite FRA slope based on spectral accelerations is obtained here,

$$FS = \frac{c + \gamma_T z \left\{ \left[ 1 - 0.5 \exp(\frac{-a + \sqrt{b}}{0.665}) \right] \cos\beta - \exp(\frac{-a + \sqrt{b}}{0.665}) \sin\beta \right\} \tan\phi}{\gamma_T z \left\{ \left[ 1 - 0.5 \exp(\frac{-a + \sqrt{b}}{0.665}) \right] \sin\beta + \exp(\frac{-a + \sqrt{b}}{0.665}) \cos\beta \right\}}$$
(3.9)

where a and b are defined by the Eqs. (3.5) and (3.6), and  $K_v$  is assumed to be 0.5  $K_h$ .



Fig. 3.5. Infinite Slope method with horizontal and vertical pseudo-static forces.

## Summary of the methodology

The stability of steep FRA slopes under partially saturated conditions will be accomplished by:

- Obtaining experimental and model based SWCCs for each research site.
- Using the measured field variations of  $\theta$ .
- Estimating the seasonal variations of FS via the proposed Eq. (3.7), for both sets of SWCCs and the measured  $\theta$ . The effective strength parameters, unit weights and geometry reported in Table 3.2 will be used, with z = 1.5 m.

The seismic stability of steep FRA slopes will be investigated by:

- Employing Eq. (3.8) and a variety of traditional pseudo-static coefficients (Table 3.1), for the highest *PGA* among the three sites and the soil and geometrical characteristic of Mountainside (Table 3.2) since they are the most critical among the 3 sites. Here, z = 1.5m.
- Comparing/validating results from Eq. (3.8) via the Simplified Bishop's Method of slices iterated with the software XSTABL (Interactive Software Designs INC 2008).
- Calculating site specific FS's based on the spectral accelerations of the field with the proposed Eq. (3.9) and presenting the results in the form of solution charts for design applications.

## **Results and Discussion**

#### SWCCs and stability results for unsaturated FRA soils

The resulting SWCCs from the ROSETTA model and laboratory suction are compared in Figs. 3.6, 3.7 and 3.8, for Premium, National and Mountainside respectively. The corrected SWCCs obtained from the pressure plate results on material < 3.36 mm, follow fairly well those obtained from the suction table with material smaller than 12.7 mm (1/2" sieve), and provide a complete SWCC for this case. Accordingly, SWCCs from the laboratory tests (material < 12.7 mm) and the ROSETTA model are used in the stability calculations.

Resulting seasonal variations of the FS [Eq. (3.7)], along with the monthly cumulative precipitation (obtained with the in situ weather stations) and measured water content are shown in Figs. 3.9, 3.10 and 3.11 for each of the sites. The figures illustrate the increase in FS that accompanies the dryer (higher suction) periods of the year, with the variation in FS depending upon the method used to obtain the SWCC. Also, smaller levels of stress suction may be found upon re-wetting of the loose surface layer (Likos and Lu 2004); however, these effects, related to the hysteresis of SWCC, are neglected here. The laboratory measured SWCCs yield larger variations in FS's than the ROSETTA model for National and Mountainside, suggesting that ROSETTA may yield more conservative estimates of suction for these soils. The maximum seasonal increase in stability at Premium and National sites occurred in March 2010, while the same occurred in June 2010 at Mountainside site. Saturation of the FRA material and a corresponding minimum FS were observed at least once during the year at each research site.



**Fig. 3.6.** Soil water characteristic curves measured in laboratory experiments and the ROSETTA model (Schaap et al. 2001), for the Premium site.



**Fig. 3.7.** Soil water characteristic curves measured in laboratory experiments and the ROSETTA model (Schaap et al. 2001), for the National site.



**Fig. 3.8.** Soil water characteristic curves measured in laboratory experiments and the ROSETTA model (Schaap et al. 2001), for the Mountainside site.





**Fig. 3.9.** Calculated monthly variations in stability (FS), volumetric water content and cumulative monthly rainfall from July 2009 to August 2010 at the Premium site.





**Fig. 3.10.** Calculated monthly variations in stability (FS), volumetric water content and cumulative monthly rainfall from July 2009 to August 2010 at the National site.





**Fig. 3.11.** Calculated monthly variations in stability (FS), volumetric water content and cumulative monthly rainfall from July 2009 to August 2010 at the Mountainside site.

Monthly cumulative rainfall data (Figs. 3.9, 3.10 and 3.11) shows that the sites were exposed to relatively large amounts of rainfall during late spring, summer, and early winter. Precipitation intensity was found to be significant during this period, with shorter and more intense rainfall events at National and Mountainside and longer but less intense events at premium site (Hoomehr et al. 2010). Rainfall events of this nature are capable of saturating the loose surface layer, and therefore, stability analyses employing traditional saturated strength parameters should always prevail in the design of FRA slopes.

#### Seismic stability analyses

#### Stability analysis with traditional pseudo-static coefficients

Results from the seismic stability analyses are shown in Table 3.4. All combinations of force orientations were examined, with the most critical being when the vertical pseudo static force is upward and the horizontal force is acting horizontally away from the slope.

Results show that the seismic infinite slope equation [Eq. (3.8)] is in good agreement with the Simplified Bishop's method, and it can be used with confidence. The Simplified Bishop's method with pseudo static forces predicts a shallow failure mechanism which is consistent with the assumptions of the infinite slope method. For a 10% P.E. in 50 yr. it is expected that the shear strength along the slip surface be 25-31% greater than the required to maintain equilibrium, whereas for a 2% P.E. in 50 yr. the shear strength is only 2-15% greater and the slope may be at the verge of failure. Since the static long-term FS was found to be 1.47 (Jeldes et al. 2013), the stability of these slopes is reduced up to 30% when seismic forces are introduced. Notice that the acceleration reduction proposed by Noda et al. (1975) for a 10% P.E. in 50 yr. yields a lower FS

	Seismic Coefficient			Computed Factor of Safety (FS)				
Method	K	2% P.E. in 50 yr. (PGA=0.2)	10% P.E. in 50 yr. ( <i>PGA</i> =0.07)	Infinite Slope Equation	Simplified Bishop's Method	Infinite Slope Equation	Simplified Bishop's Method	
				2% P.E. in 50 yr.		10% P.E. in 50 yr.		
Theoretical d'Alembert's Principle (no reduction)	$K_h$	0.20	0.07	0.91	0.92	1.24	1.25	
	K <sub>v</sub>	0.10	0.04					
Makdisi and Seed (1978)_(M≈6.5)	$K_h$	0.10	0.04	1.14	1.15	1.35	1.34	
	K <sub>v</sub>	0.10	0.04					
Hynes-Griffin and Franklin (1984)	$K_h$	0.10	0.04	1.14	1.15	1.35	1.34	
	K <sub>v</sub>	0.10	0.04					
Noda et al. (1975)	$K_h$	0.19	0.14	0.92	0.94	1.07	1.07	
	K <sub>v</sub>	0.10	0.04					
Bray and Rathje (1998)	$K_h$	0.15	0.05	1.02	1.03	1.29	1.31	
	K <sub>v</sub>	0.10	0.04					
Kavazanjian et al. (1997) (without response analysis)	K <sub>h</sub>	0.10	0.04	1.14	1.15	1.35	1.34	
	$K_{\nu}$	0.10	0.04					

# **Table 3.4** Summary of $K_h$ , $K_v$ and obtained FS's for a 2 and 10% P.E. in 50 yr.

than the case without reduction, which is an inconsistency. Noda et al. (1975) developed the equation based on limit-equilibrium analysis of structures (quay walls) that were already damaged by earthquakes, where the results and parameter selection probably have a high level of subjectivism (Towhata 2008). Regarding the seismic coefficient to be used, Hynes-Griffin and Franklin (1984), Bray and Rathje (1998) and Kavazanjian et al. (1997) are suggested here, because they express the seismic coefficient in terms of  $PGA_h$  rather than an arbitrary fraction of the acceleration of gravity. Current hazard maps have broad coverage and easier access, and a more site-based design is possible.

## Slope stability charts based on spectral accelerations and an Illustrative example

Graphical chart solutions of Eq. (3.9) for different earthquake magnitudes, maximum allowable displacements and soil properties were developed for each of the FRA reclaimed sites. Fig. 3.12 illustrates the solution chart for the conditions at Mountainside site when M = 5.0 (about the highest earthquake magnitude recorded in the area)  $D_a = 30$  cm for a 2% P.E. in 50 years. The random variable  $\varepsilon$  was selected to be 0.66 representing a 16% of allowable displacement exceedance (Bray and Travasarou 2009). Similar solution charts for other values of M,  $D_a$  and probability of exceedance can be created. Fig. 3.12 suggest that failure may occur at Mountainside when  $S_a > 0.4g$  and  $T_s > 1.5$  s.

To illustrate the use of the solution charts, the Mountainside slope is used with M = 5.0 and  $D_a = 30$  cm. The initial fundamental period of the slope may be calculated using  $T_s = 4H/Vs$ , where H is the depth of the sliding plane and Vs average shear wave velocity. Since Vs has not been



**Fig. 3.12.** FS chart for Mountainside site as a function of spectral accelerations and initial fundamental period.

measured, an average Vs = 200 m/s is assumed. For a shallow failure mode in the upper 1.5 m, the initial fundamental period becomes  $T_s = 0.03$  and the spectral acceleration at the degraded period  $S_a$  (T = 0.045)  $\approx 0.36$ g (Fig. 3.3). Then, for  $T_s = 0.03$  s and  $S_a = 0.36$ g the FS is found from Eq. (3.9) or Fig. 3.12 to be approximately 1.26.

## Conclusions

Static long-term stability analyses of the three constructed FRA slopes conducted by Jeldes et al. (2013) showed that shear strength is at least 47% greater than that required to balance the destabilizing forces, in the absence of downslope seepage. These results reflect the assumption of total saturation of the soil at the slip surface, which is not always consistent with the in situ FRA conditions, where the surface layer can be seen as an active unsaturated zone whose variations in water content will be directly affected by periodic changes in environmental conditions. Slope stability analyses for partially saturated mine soils showed that seasonal increments in the stability of steep FRA slopes are possible, and the static long-term stability is a lower bound of the real field performance. Naturally, the interpretation of results is restricted to the hydrological conditions experienced during the time period at which the data was collected. This period was characterized by relatively poor establishment of ground cover, and therefore, evaporation was probably the main mechanism inducing matric suction and the relative increase in stability. Different monthly stability patterns may be observed in the future depending on the amount of vegetation, precipitation and variations in temperature affecting evapotranspiration rates. In any case, temporal saturation of the loose surface layer is likely to occur and slope stability analysis employing traditional saturated strength parameters should prevail in the design of FRA slopes.

Under seismic loading conditions, the proposed infinite slope equation showed good agreement with the simplified Bishop's method of the slices. The selection of the pseudo-static coefficient requires an understanding of the assumptions behind the coefficient and practical knowledge of the local seismicity. Site specific results were obtained when the horizontal seismic coefficient proposed by Hynes-Griffin and Franklin (1984), Bray and Rathje (1998) and Kavazanjian et al. (1997) were employed, and suggest that FS can be as low as 1.02. While these methods based on a fraction of the probable PGA are most common, they do not consider the effects of frequency and duration (Bray 2007). A method based on spectral response may provide a better tool for evaluating slope stability, and a combination of Bray and Travasarou (2009) approach with the proposed modification of the infinite slope equation was suggested. Charts were developed for this method employing site-specific  $S_a$  values and local soil properties and slope geometry. Here, the FS is a function of parameters that have less selection subjectivity for the designer and are a function of the importance of the project, available seismic data, and accuracy of the site characterization. It is suggested that the use of spectral values for a 2% P.E. in 50 yr. or higher are sufficient, except when the designer has reasons for using lower local values.

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# Chapter 4. An Approximate Solution to the Sokolovskiĭ Concave Slope at Limiting Equilibrium

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# Abstract

The growth of precision auto-guidance systems on construction equipment suggests that nonplanar slopes and landforms can now be readily constructed. Slopes with concave cross sections not only appear more like natural slopes, but can also have superior stability and erosion resistance. Thus, it is desirable to have the description of concave slopes that provide mechanical stability for given set of soil properties. In this article, an approximate solution that defines the geometry of critical concave slopes (FS  $\approx$  1) in a frictional medium is developed, based on the slip line field method of Sokolovskii. The approximate solution is compared with the Sokolovskii's numerical results and validated via Limit Equilibrium and Finite Element Method analyses. The proposed solution is simple in form, and when implemented with precision construction equipment will allow the construction of embankments and reclaimed mine lands which appear more like those in nature, and yet are more erosion-resistant.

# Introduction

The growth of precision of auto-guidance construction equipment and 3D mapping technology allows more complex non-planar slopes and landforms to be readily constructed. This can lead to engineered slopes of concave cross sections which appear more like natural slopes and have

superior stability and erosion resistance (Schor and Gray 2007). In fact, experimental and numerical simulations have shown that concave slopes lead to less erosion than planar slopes (Meyer and Kramer 1969, Rieke-Zapp and Nearing 2005). The general concept of "form follows function" is observed in geomorphological evolutionary processes such us fluvial systems where concave slopes are suggested as the most probable shape in landform evolution (Leopold and Langbein 1962, Miyamoto et al. 2005). Furthermore, slopes subjected to physical weathering seem to evolve into steady-state concave-like forms to achieve erosional equilibrium (Nash 1980, Pelletier and Rasmussen 2009, Twidale 2007). The use of concave surfaces can become a powerful ecological building technique to reduce sediment yield in constructed embankments, reclaimed mine lands, and highway cut and fill sections, with shapes that appear more like those observed in nature. Thus, it is important to investigate and define the optimum concave shape relative to mechanical slope stability. While computational methods based on limit equilibrium and limit analysis techniques (Ahmed et al. 2012, Liu and Zhao 2012, Michalowski 2010) prevail in slope analyses and design, a rational mechanics approach based on the Sokolovskii's slip line field theory was employed here to define concave slope shapes at critical equilibrium.

Sokolovskii (1960, 1965) found that the slope surface at critical or limiting equilibrium has a concave shape, mathematically described as a function of the Mohr-Coulomb (M-C) shear strength parameters and the material unit weight. Since the solution for a non-frictional soil (i.e. undrained loading conditions) is mathematically straightforward, this work focuses on the  $c - \phi$  case ( $\phi$  is the soil internal friction angle and c is the soil cohesion), which has no analytic solution. For a soil with  $\phi > 0$ , c > 0, and unit weight  $\gamma > 0$ , Sokolovskii (1960) implemented a numeric solution

reporting the slope surface as a function of the ratio  $c/\gamma$ . This solution, however, has some practical design limitations. The Sokolovskiĭ geometric description of the contour is incomplete which limits the application to distinct values of  $\phi$  and slope heights  $y \cdot c/\gamma$ . In this article we: a) revisit and fully describe Sokolovskiĭ's formulation and solution based on slip line field theory for a weightless slope ( $\gamma = 0$ ), b) develop an approximate analytic solution which describes the concave slope at critical equilibrium for a medium with self-weight ( $\gamma > 0$ ), and c) validate the solution via the Limit Equilibrium (LEM) the and Finite Element Method (FEM). While the main objective of this piece of work is to provide a rational mathematic description of the critical concave slope shape (FS = 1), the solution can be easily extended to FS's > 1 as needed for design. This issue, along with those related to transient underground water flow and constructability are briefly addressed in the discussion.

#### **Background:** the slip line field theory and the characteristic equations

The slip line field theory in soil mechanics rests on the assumptions that the soil yield stress is independent of the strain level and that the strain-stress behavior can be described by either elastic perfectly plastic or rigid perfectly plastic idealizations. For frictional materials like soils, the Mohr-Coulomb shear strength expression has been traditionally used as the constitutive law defining the yield or maximum allowable stress. In this context, to bring the body into a limiting equilibrium state or at the verge of flowing plastically, the largest difference between the shear stresses and the shear strength must be zero (Sokolovskiĭ 1960). Expressed in terms of principal stresses:

$$\max\left\{\left(\sigma_{1}-\sigma_{3}\right)\sin(2\lambda-\phi)-\sin\phi(\sigma_{1}+\sigma_{3}+2H)\right\}=0$$
(4.1)

where  $\sigma_1$  is the major principal stress,  $\sigma_3$  is the minor principal stress (Fig. 4.1),  $H = c \cdot \cot \phi$  is



**Fig. 4.1.** Orientation of slip lines in a soil mass at limiting equilibrium [adapted from Sokolovskiĭ (1965)].

the tensile strength of the soil, and  $\lambda$  is the angle between the normal to the plane and the major principal direction. Approaching [Eq. (4.1)] as a mathematical optimization problem, the limiting condition is reached when

$$2\lambda = \phi + \frac{\pi}{2} \tag{4.2}$$

which implies that failure or slip line initiates at an angle  $\varepsilon$  (Fig. 4.1) inclined at  $|45^\circ - \phi/2|$  from the direction of the major principal stress and  $|45^\circ + \phi/2|$  from the direction of the minor principal stress. Note that the inclinations of the slip lines are independent of the strain level, which is not always consistent with experimental observations. Roscoe (1970) concluded that a strain based approach rather than a stress approach better predicts the inclination of the slip plane, and its magnitude can be approximated by  $45^\circ + \psi/2$  (measured from the minor principal direction), where  $\psi$  is the dilatancy angle at failure. Later studies on shear band formation in plane strain experiments confirmed Roscoe's findings (Alshibli and Sture 2000). Nevertheless, this approach is still valid in the context of associated plasticity ( $\phi = \psi$ ), which for relatively unconfined problems such as slope stability, does not significantly influence the prediction of the Factor of Safety (FS) (Griffiths and Lane 1999) nor the prediction of the critical slip or failure mechanism (except for cases of very low friction angle and high cohesion) (Cheng et al. 2007). By combining Eqs. (4.1) and (4.2), Sokolovskiĭ (1960) obtained the limiting equilibrium in terms of principal stresses as:

$$(\sigma_1 - \sigma_3) = \sin\phi(\sigma_1 + \sigma_3 + 2H) \tag{4.3}$$

If the Sokolovskii (1965) x-y reference frame is adopted (see Fig. 4.1) where the normal of an infinitesimal element located anywhere inside the soil mass is inclined at an angle  $\mu$  from the x-

axis, the slip lines will be inclined at an angle  $\xi \pm \varepsilon$  from the x-axis. Then, the equations of limiting stresses can be written in the directions of the adopted reference frame as:

$$\sigma_{xx} = \sigma(1 + \sin\phi\cos 2\xi) - H \tag{4.4}$$

$$\sigma_{vv} = \sigma(1 - \sin\phi \cos 2\xi) - H \tag{4.5}$$

$$\tau_{xy} = \sigma \sin\phi \sin 2\xi \tag{4.6}$$

and in the normal  $\hat{n}$  and tangential  $\hat{t}$  directions (Fig. 4.1)

$$\sigma_n = \sigma \left[ 1 + \sin\phi \cos 2(\xi - \mu) \right] - H \tag{4.7}$$

$$\sigma_{t} = \sigma \left[ 1 - \sin\phi \cos 2(\xi - \mu) \right] - H \tag{4.8}$$

where  $\sigma_{xx}$  and  $\sigma_{yy}$  are the normal limiting stresses acting in the x and y direction respectively,  $\tau_{xy}$  is the shear limiting stress acting on planes perpendicular to the x and y axes,  $\sigma_n$  and  $\sigma_t$  are the normal limiting stresses acting in the normal  $\hat{n}$  and tangential  $\hat{t}$  directions respectively, and  $\sigma = 1/2(\sigma_1 + \sigma_3) + H$  is the mean stress. When Eqs. (4.4), (4.5) and (4.6) are combined with those of a continuum solid at equilibrium (Malvern 1969), a set of two first order partial differential equations is obtained, which upon solution via the *method of characteristics* (Hill 1950) becomes a set of two differential equations that describe the direction of the slip lines in a soil body, called *characteristic equations*. Sokolovskiĭ (1965) obtained:

$$\frac{dy}{dx} = \tan(\xi \mp \varepsilon) \tag{4.9}$$

$$d\sigma \mp 2\sigma \tan \phi d\xi = \gamma \left[ dy \mp \tan \phi dx \right] \tag{4.10}$$

# Elaboration of Sokolovskiĭ solution for the critical slope in a weightless medium

By assuming a weightless medium ( $\gamma = 0$ ), the first-order hyperbolic Eq. (4.10) yields an analytic solution. Here, an alternative development of Sokolovskiĭ (1960) analytic solution of this problem is offered, aiming to provide clarity to the mathematical formulation.

Consider the boundary conditions under an external stress q depicted in Fig. 4.2, where the origin of the rectilinear system of coordinates is set at the intersection of the slope surface and the horizontal ground surface. At the top of the slope (along the horizontal ground surface), the stresses  $(\sigma_{xx} = 0, \sigma_{yy} = q, \tau_{xy} = 0)$  of the first boundary condition define the y-axis and the x-axis as the major and minor principal directions. Here,  $\xi = \pi/2$  and the mean normal stress  $\sigma^{surf}$  from Eq. (4.5) becomes:

$$\sigma^{surf} = \frac{q+H}{1+\sin\phi} , \quad \xi = \frac{\pi}{2}$$
(4.11)

Along the slope surface, the stresses redefine the major and minor principal directions as the tangential  $\hat{t}$  and normal  $\hat{n}$  components respectively. Therefore,  $\xi$  becomes the angle of the slope  $\beta$  measured from the x-axis. By defining the boundary normal stress  $\sigma_n = 0$  and  $\xi - \mu = \pi/2$ , Eq. (4.7) returns the magnitude of the mean normal stress along the slope surface  $\sigma^{cont}$ :

$$\sigma^{cont} = \frac{H}{1 - \sin\phi} , \quad \xi = \beta \tag{4.12}$$

Eqs. (4.11) and (4.12) establish the boundary conditions of the problem. For a weightless ( $\gamma = 0$ ) medium, Eq. (4.10) becomes:

$$\frac{d\sigma}{\sigma} = \pm 2\tan\phi d\xi \tag{4.13}$$



**Fig. 4.2.** Critical slope shape (planar) for a weightless medium. Inside zone AoC the principal directions are aligned with the reference frame. On the slope surface, the principal directions are aligned with the normal and tangential directions.

Integrating Eq. (4.13) and imposing the boundary condition along the slope surface from Eq. (4.12) yields

$$\frac{\cot\phi}{2}\ln\left[\frac{H}{1-\sin\phi}\right] + c1 = \pm\beta \tag{4.14}$$

where *c*1 is the constant of integration. Eq. (4.14) reveals the existence of two planes satisfying the stress condition; they have the same absolute inclination, but differ only in direction. Choosing  $-\beta$  and imposing the surface boundary condition from Eq. (4.11):

$$c1 = -\frac{\cot\phi}{2}\ln\left[\frac{q+H}{1+\sin\phi}\right] - \frac{\pi}{2}$$
(4.15)

Substituting Eq. (4.15) into Eq. (4.14)  $\beta$  is found to be constant revealing a planar slope,

$$\beta = \frac{\pi}{2} + \frac{\cot\phi}{2} \ln\left[\frac{q+H}{H}\frac{1-\sin\phi}{1+\sin\phi}\right]$$
(4.16)

which is equivalent to the Sokolovskii (1960) solution, with the discrepancy found only due to the difference in the employed coordinate system. The x-y coordinates can be found by solving the first equation of the characteristic system [Eq. (4.9)], with  $\varepsilon = 0$  and  $\xi = \beta$  at the slope surface:

$$\frac{dy}{dx} = \tan\left(\beta\right) \tag{4.17}$$

It is important to emphasize that for the case of a weightless medium, the slope is constant or planar. However, a medium with self-weight induces changing stress conditions which influences the principal directions, reducing the slope inclination downslope. Sokolovskii's general ( $\gamma \neq 0$ ) solution of Eq. (4.10) is the result of a numeric boundary value problem, which has no analytical counterpart. Unfortunately, Sokolovskii's solution is constrained geometrically (i.e. slope height), limited by the ratio  $c/\gamma$  and restricted to specific values of  $\phi$ . Furthermore, extending this

solution to other initial conditions requires the reconstruction of the sophisticated boundary value problem making its implementation in practical slope design unlikely.

#### Proposed solution for the critical slope in a medium with self-weight

An analytical approximation for the slope shape of a medium with self-weight is proposed here, based on a Weightless Medium Approximation (WMA). Our derivation begins with the initial conditions of the uppermost layer supporting the maximum vertical height of a continuous medium. This vertical height is limited by the height of the tension crack  $h_{cr}$  (Terzaghi 1943) and defines the upper portion of the slope (Fig. 4.3) as a tension zone lying above the x-axis.

$$h_{cr} = \frac{2c\cos\phi}{\gamma(1-\sin\phi)} \tag{4.18}$$

For the portion of the slope lying below the x-axis, the WMA assumes that the medium can be discretized as a series of finite weightless intervals with constant surface inclination, each supporting a line load representing the self-weight of the layer above. Fig. 4.4 depicts a model or analog for a slope composed of multiple layers of weightless media (rigid foam beads), with layers loaded by a line load (frictionless lead sheets). Each new layer adds a finite weight that changes the slope by a finite amount. As the layer thickness  $\Delta y \rightarrow 0$ , the x-y concave slope emerges. The weight of the material is therefore introduced as an external stress q including the weight of the tension zone [ $\gamma h_{cr}$ ] above the x-axis, and a linearly increasing component [ $\gamma y$ ] starting from the origin of the reference frame; thus:

$$q(y) = (h_{cr} + y)\gamma \tag{4.19}$$



Fig. 4.3. Slope surface in a medium possessing weight [after Sokolovskiĭ, (1965)].



**Fig. 4.4.** Weightless Medium Approximation: discretized weightless medium (e.g. rigid foam beads) supporting thin, external loads (e.g. frictionless lead sheets) separated by a finite interval: a) coarser discretization and b) finer discretization.

The analytic solution from Eq. (4.16) can be modified by replacing the constant load q with Eq. (4.19). Because q(y) is a function of the vertical direction, y is treated as the independent variable and the angular coordinates are transformed to measure the angle from the y-axis. This change of coordinates [ $\beta' \rightarrow \beta - \pi/2$ ] makes the slope of the first interval zero rather than infinite.

$$\beta'(y) = \frac{\cot\phi}{2} \ln\left[\frac{q(y) + H}{H} \frac{1 - \sin\phi}{1 + \sin\phi}\right]$$
(4.20)

Accordingly, the first equation of the characteristic [Eq. (4.9)] after transformation and integration over the closed interval [0, y], becomes:

$$x(y) = \int_{0}^{y} \tan\left[\beta'(y')\right] dy'$$
(4.21)

The exact solution of Eq. (4.21) involves a Gauss hypergeometric function  $_{2}F_{1}$  with complex (real and imaginary) arguments. If as a first approximation only the real component  $[x_{\mathbb{R}}(y)]$  of Eq. (4.21) is taken, it can be seen (Fig. 4.5) that a concave slope is depicted; however, this solution does not satisfy the critical equilibrium requirement. The discrepancy between Sokolovskii's critical slope and  $x_{\mathbb{R}}(y)$  is illustrated in Fig. 4.5 (with  $x_{\mathbb{R}}(y)$  plotted in the negative side of x-axis to match Sokolovskii). Notice that  $x_{\mathbb{R}}(y)$  soon diverges into a horizontal asymptote corresponding to an infinite slope  $(\Delta x / \Delta y \rightarrow \infty)$ . The resolution of this infinite-slope problem lies in a second approximation.

While our first WMA is a physical discretization and modification of Eq. (4.10), our second



**Fig. 4.5.** Sokolovskiĭ's critical slope vs the real (non-imaginary) component of the solution of Eq. (4.21).

approximation is a mathematical modification of Eq. (4.21), via examination of the Taylor series of the tangent of an angle  $\alpha$ 

$$\tan\left(\alpha\right) = \alpha + \frac{\alpha^3}{3} + \frac{2\alpha^5}{15} + O\left(\alpha^7\right)$$
(4.22)

with  $O(\alpha^7)$  representing truncation of the series from the 7<sup>th</sup> order. The higher order terms make the Taylor expansion diverge sooner near  $\alpha = \pm \pi/2$ . This effect is minimized by removing the high-order terms, replacing  $\tan(\alpha) \rightarrow \alpha$ . Applying the Taylor expansion without the high order terms reduces Eq. (4.21) to:

$$x(y) = \int_{0}^{y} \beta'(y') dy'$$
(4.23)

which has the following analytical solution:

$$x(y) = \begin{cases} 0 & , -h_{cr} \le y \le 0 \\ A\left[\sigma_{y} \left(B-1\right)\left(\cos e c \phi-1\right)+H B\left(\cos e c \phi+1\right)\right], & y > 0 \end{cases}$$
(4.24)

where:

$$A = \frac{\cos\phi}{2\gamma(1-\sin\phi)} \tag{4.25}$$

$$B = ln \left[ \frac{\sigma_y}{H} \left( \frac{1 - \sin \phi}{1 + \sin \phi} \right) + 1 \right] = ln \left[ \frac{\sigma_y}{H} K_a + 1 \right]$$
(4.26)

$$\sigma_{y} = \gamma y \tag{4.27}$$

$$H = c \cot \phi \tag{4.28}$$

Notice that the factor *B* is a function of the Rankine (1857) active coefficient of earth pressure  $K_a = (1 - \sin \phi) / (1 + \sin \phi).$ The proposed solution describes a critical slope surface in the

quadrant where x-axis and y-axis are positive, with the tension zone  $h_{cr}$  above the x-axis from the origin to the point of coordinates (0,-  $h_{cr}$ ).

# Validation of the WMA solution

Comparisons between the concave slope surface from the proposed solution [Eq. (4.24)] and the Sokolovskiĭ (1960) numerical approach were conducted in terms of geometry, critical FS, and failure mechanism. For comparison, the results from Eq. (4.24) will be plotted in the negative side of the x-axis to match Sokolovskiĭ's results. Concave slopes for the range of  $\phi$  and  $c/\gamma$  typically observed in soils are employed here. For fine grained soils the range c = 5 - 40 kN/m<sup>2</sup> covers the majority of the cases reported by Mesri and Abdelghaffar (1993), while  $\gamma = 10 - 23$  kN/m<sup>3</sup> covers most materials from clays to coarse granular soils (NAVFAC 1986). Thus,  $\phi = 20^{\circ}$ ,  $30^{\circ}$ , and  $40^{\circ}$ , and  $c/\gamma = 0.2$ , 1, 2, 3, and 4 m are employed here to compare Sokolovskiĭ (1960) with the proposed solution. The total vertical slope height ( $H_s$ ) employed was the maximum reported by Sokolovskiĭ (1960) for the selected range of  $c/\gamma$  (Table 4.1). The influence of  $H_s$  over the computed FS and failure mechanisms is discussed later.

#### Geometry of the concave slopes

The slope surfaces for  $\phi = 30^{\circ}$  and for a  $c/\gamma = 2$  m are compared in Fig. 4.6 for a maximum  $H_s$ = 47 m, while slopes for  $\phi = 20^{\circ}$ , 30°, and 40°, are compared in terms of non-dimensional coordinates  $(x \cdot \gamma/c, y \cdot \gamma/c)$  in Fig. 4.7. For all combinations of  $\phi$  and  $c/\gamma$ , a very close agreement at the upper portion of the slope is observed, while differences appear at the lower

	Internal friction angle $\phi$ (°)					
<i>c/y</i> (m)	20	30	40			
0.2	3.8	5.3	7.1			
1.0	19	26	35			
2.0	38	53	71			
3.0	57	79	106			
4.0	76	105	141			

Table 4.1 Maximum values of  $H_s$  (m) reported by Sokolovskii (1960) and used in stability



analyses (1960) solution from LEM

**Fig. 4.6.** Comparison of Sokolovskiĭ (1960) numerical solution and the proposed analytical solution for  $\phi = 30^{\circ}$  and  $c / \gamma = 2$  m.



Fig. 4.7. Comparison of Sokolovskii (1960) numerical solution and the proposed analytical solution in dimensionless coordinates for  $\phi = 20^{\circ}$ ,  $30^{\circ}$ , and  $40^{\circ}$ .

portion of the slope. The amount of variation seems to depend only on the value of  $\phi$ , with best geometrical agreement found for  $20^{\circ} \le \phi \le 30^{\circ}$ , which is the range found in many fine grained soils. Nevertheless, as discussed in the subsequent section, these geometrical variances only induce a minor difference in stability, because regardless of the value of  $\phi$  the critical failure mechanism develops in the upper zone of close agreement between the solutions.

#### Critical FS and observed failure mechanisms

Differences in stability between the two approaches were investigated via Finite Element Method (FEM) and Limit Equilibrium Method (LEM). FEM analyses were computed using the software *Phase2* (Rocscience Inc. 2011), while LEM analyses using the Software *Slide* (Rocscience Inc. 2011). In the FEM model, a non-associated flow rule with  $\psi = 0$  was employed to avoid effects of computation domain size (Cheng et al. 2007) and to limit the over-prediction of dilation (Griffiths and Lane 1999). Also, nominal values for Young's modulus ( $E = 2x10^4$  kPa) and Poisson's ratio (v = 0.3) were used, since they play a minor role in the calculation of the FS (Cheng et al. 2007, Griffiths and Lane 1999). In the LEM analyses, the Simplified Bishop's Method was used with 495,000 critical surfaces analyzed.

A comparison of the computed FS's via LEM are reported in Table 4.2, while the FEM results for Sokolovskiĭ's and the proposed solutions when  $\phi = 30^{\circ}$  and  $c/\gamma = 2$  m are shown in Figs. 4.8 and 4.9 respectively. Very similar results in terms of FS's were obtained from FEM and LEM approaches. For all combinations of  $\phi$  and  $c/\gamma$ , the computed critical FS's do not differ more than 4% between the proposed and Sokolovskiĭ's solutions. More importantly, all computed FS's

	Internal friction angle $\phi$ (°)								
	20		30		40				
<i>c/y</i> (m)	Proposed	Sokolovskiĭ	Proposed	Sokolovskiĭ	Proposed	Sokolovskiĭ			
	approximate	(1960)	approximate	(1960)	approximate	(1960)			
0.2	1.00	1.04	1.04	1.03	1.06	1.03			
1.0	1.00	1.04	1.04	1.03	1.06	1.03			
2.0	1.00	1.04	1.04	1.03	1.06	1.03			
3.0	1.00	1.04	1.04	1.03	1.06	1.03			
4.0	1.00	1.04	1.04	1.03	1.06	1.03			

 Table 4.2 Comparison of stability between proposed approximate solution and Sokolovskii (1960) solution from LEM



Fig. 4.8. FEM results in terms of shear strains for the Sokolovskii (1960) critical concave slope.



Fig. 4.9. FEM results in terms of shear strains for the proposed critical concave slope.

approximate 1, which is consistent with the limiting strength condition, and confirms the appropriateness of Eq. (4.24). For the investigated range of  $\phi$ , the proposed solution does not deviate more than 6% from the theoretical FS = 1, which can be considered a minor difference in terms of practical applications in geotechnical design.

# Effects of slope height on failure mechanisms and FS obtained from the proposed solution

To investigate the effects of the slope height  $H_s$  on the stability of the proposed concave slopes, LEM stability analyses were conducted for different values of  $\phi$ ,  $c/\gamma$ , and  $H_s$ .

Analyses of the results indicate the existence of a limiting slope height ( $h_L$ ) at which the failure mechanism of the proposed concave slope changes. Concave slopes with heights  $H_s < h_L$  develop toe-failure mechanisms, whereas those with  $H_s > h_L$  develop face-failure mechanisms (Fig. 4.10). The change in failure mechanism naturally alters the computed FS, as illustrated in Fig. 4.11 for the  $\phi = 20^\circ$  case. Concave slopes with  $H_s < h_L$  tend to have slightly greater FS's, which only in few cases exceeded the theoretical FS = 1 by more than 10%. As  $H_s$  approaches  $h_L$ , the FS decreases until it reaches a steady value (FS  $\approx$  1) that remains constant for all  $H_s > h_L$ . This steady equilibrium condition coincides with a somewhat steady position of the critical failure surface, and therefore,  $h_L$  defines the height of the failure slip for larger concave slopes. Similar analyses on concave slopes with  $\phi = 30^\circ$ , 38° and 45°, not only yielded equivalent results, but also showed that  $h_L$  can be approximated as a linear function of  $c / \gamma$  only (Fig. 4.12). Notice that over 99% of

110



**Fig. 4.10.** Modes of failure depicted in the form of shear bands (maximum shear strains) as observed in concave slopes (FEM). Figure illustrate case for  $\phi = 20^{\circ}$  and  $c/\gamma = 2$  m, and two values of  $H_s$ .



Fig. 4.11. Computed FS (LEM) on concave slopes with different vertical heights  $H_s$  and  $c/\gamma$  for  $\phi = 20^{\circ}$ .



**Fig. 4.12.** Relationship between the limiting height  $h_L$  and  $c / \gamma$ .

the variability in  $h_L$  can be explained by the linear relationship  $h_L = 15.8 c / \gamma + 0.3$ , which can be used to estimate the most likely failure mechanism of a concave slope.

# **Discussion and conclusions**

Concave slope shapes mimic the landforms observed in nature, and have been determined to be less prone to erosion and sediment production. A mathematical description of the concave slope surface at limiting equilibrium was developed by Sokolovskiĭ (1960), but the characteristic equations must be solved numerically for the most relevant cases, i.e. frictional soils with selfweight ( $\phi > 0$ ,  $\gamma > 0$ ). The lack of an analytic solution limits the application of concave slopes in practice.

In this article, an analytical mathematical approximation for the  $\phi > 0$  and  $\gamma > 0$  case was offered in order to obtain concave slopes at critical equilibrium, for any combination of  $\phi$ , c, and  $\gamma$ . This solution was based on the Sokolovskiĭ solution for a weightless ( $\gamma = 0$ ) medium. The approximation transformed a set of differential equations into a single algebraic expression, which can be more readily used in slope design with no need to reconstruct Sokolovskiĭ's sophisticated numerical boundary value problem for the  $\gamma > 0$  case. The solution offered here is simple, effective, and it was shown to provide good agreement with Sokolovskiī's numerical approach. Unlike Sokolovskiĭ's tabulated values, the proposed solution can be used for any desirable maximum value of  $H_s$  or any value of  $\phi$ . In theory, every point of the slope surface depicted by Eq. (4.24) should be in limiting equilibrium such that, independent of the size of the physical domain, the FS approximates 1. Nevertheless, small deviations arise due to the approximate nature of the solution. Though associated with the mode of failure, these deviations are small for geotechnical applications and indicate that the stability of the proposed critical concave slopes will range from FS = 1 - 1.1 for the majority of cases.

Although the solution for the concave slope was developed in terms of the critical slope shape (FS = 1), the proposed solution can be used as a base line for practical applications. For design, a given FS can be obtained if a strength reduction factor (Griffiths and Lane 1999, Spencer 1967) equal to the desired FS is used with our proposed solution. The effects of an external semi-infinite surcharge  $(q_e)$  acting on the ground surface can be included by adding  $q_e$  to the right side of Eq. (4.19) and modifying Eq. (4.18) as  $h_{er} = 2c \cos \phi / [\gamma (1-\sin \phi)] - q_e / \gamma$ . Once the desired concave shape is defined, the effects of transient ground water flow can be investigated using commercial slope stability software. While the successful construction of concave slopes will depend on the precision of the employed construction equipment, the increasing precision and availability of GPS based auto-guidance equipment (Koehrsen et al. 2001) makes the construction of concave slopes possible. Nevertheless, in actual slope construction the sharp bluff at the top (the tension crack zone) is not likely to be constructed since it may erode over time; instead this would be rounded for a more natural appearance, which would decrease the destabilizing forces and add conservatism to concave slope performance.

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# Appendix: The slip line field theory and the characteristic equations

#### Equations of internal equilibrium and principal stresses

From the assumption that soil medium can be treated as a continuum, the complete state of stresses acting on any plane can be specified if the resultant stresses in the directions of the axes of the reference frame are known (Malvern 1969). For the plane strain condition (e.g. the profile of a 2D slope), strains and shear stresses in the z direction (perpendicular to the paper) vanish and therefore, the equations of internal equilibrium in a xy space that not necessarily coincides with the vertical and horizontal can be written as:

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} = \gamma \sin \alpha \tag{4.29}$$

$$\frac{\partial \tau_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} = \gamma \cos \alpha \tag{4.30}$$

where  $\sigma_{xx}$  and  $\sigma_{yy}$  are the normal stresses acting in the x and y directions,  $\tau_{xy} = \tau_{yx}$  are the shear stresses,  $\gamma$  is the unit weight of the material, and  $\alpha$  is the angle between an horizontal plane and the x-axis. The complete state of stresses at a point, mathematically expressed as a second order tensor, can be described by the principal stresses and directions, which are independent of the selected reference frame. Any normal and tangential stress component on any plane can be determined if the principal stresses are known (Holtz and Kovacs 1981):

$$\sigma_n = \frac{1}{2} (\sigma_1 + \sigma_3) + \frac{1}{2} (\sigma_1 - \sigma_3) \cos 2\lambda$$
(4.31)

$$\tau_n = \frac{1}{2} (\sigma_1 - \sigma_3) \sin 2\lambda \tag{4.32}$$

118

where  $\sigma_1$  is the major principal stress,  $\sigma_3$  is the minor principal stress, and  $\lambda$  is the angle between the direction of the normal to the plane and the major principal direction.

#### Formulation of the equations of the characteristics

The slip line field theory in soil mechanics rests on the assumptions that the soil yield stress is independent on the strain level and that the strain-stress behavior can be described by either elastic perfectly plastic or rigid perfectly plastic idealizations. In this way, the formulation of the problem does not require stress and displacement boundary conditions, which combined impose a statically indeterminate condition (Scott 1963). For frictional materials like soils, Mohr-Coulomb has been traditionally used as the constitutive law defining the yield or maximum allowable stress. In this context, to bring the body into a limiting equilibrium state or at the verge of flowing plastically, the largest difference between the shear stresses and the shear strength must be zero (Sokolovskiĭ 1960):

$$\max\left\{\left|\tau_{n}\right| - (\sigma_{n} + H) \tan\phi\right\} = 0 \tag{4.33}$$

where  $\sigma_n$  and  $\tau_n$  are the normal and shear stresses acting on any plane within the soil body and H is the tensile strength of the soil. If the stresses in the soil body exceed this liming state, plastic deformation will occur. According to the slip line field theory, the plastic deformation will take place by the relative movement of two planes on which the inclination of the stresses is maximum (Scott 1963). These planes, called slip planes or shear bands, are oriented parallel to the straight lines that connect the pole of the Mohr circle and its contact with the failure condition in a  $\tau - \sigma$  stress space (Fig. 4.13). If  $\sigma_n$  and  $\tau_n$  are written in terms of principal stresses [Eqs. (4.31) and (4.32)], the limiting condition in Eq. (4.33) can be written as


Fig. 4.13. Illustration of slip planes on Mohr-Coulomb diagram.

$$\max\left\{\left(\sigma_{1}-\sigma_{3}\right)\sin(2\lambda-\phi)-\sin\phi(\sigma_{1}+\sigma_{3}+2H)\right\}=0$$
(4.34)

Approaching the condition of maximum difference as a mathematical optimization problem, it is found that the maximum difference between shear stresses and strength is reached when:

$$2\lambda = \phi + \frac{\pi}{2} \tag{4.35}$$

According to Sokolovskii (1960), at the limiting condition the absolute values of  $\lambda$  and  $\varepsilon$  (the angle between the slip planes and the direction of the major principal stress) are complementary, and therefore, the slip planes make an angle

$$\left|\varepsilon\right| = (45^\circ - \phi/2) \tag{4.36}$$

with the direction of the major principal stress and  $\pm (45^\circ + \phi/2)$  with the direction of the minor principal stress.

The solution in Eq. (4.36) is independent of the strain level, which is not always consistent with experimental observations. Roscoe (1970) concluded that a strain based approach rather than a stress approach better predicts the inclination of the slip plane, and its magnitude can be approximated by  $45^{\circ} + \psi/2$ , where  $\psi$  is the dilatancy angle at failure. Later studies on shear band formation in plane strain experiments confirmed Roscoe's findings (Alshibli and Sture 2000). Nevertheless, Eq. (4.36) is still valid in the context of the associated plasticity ( $\phi = \psi$ ), which for relatively unconfined problems such as slope stability, do not influence the prediction of the FS (Griffiths and Lane 1999) neither the prediction of the critical slip or failure mechanism (except for cases of very low friction angle and high cohesion) (Cheng et al. 2007). Accordingly, if Eqs.

(4.34) and (4.35) are combined, the equation for the limiting equilibrium in terms of principal stresses becomes (Sokolovskiĭ 1960):

$$(\sigma_1 - \sigma_3) = \sin\phi(\sigma_1 + \sigma_3 + 2H) \tag{4.37}$$

On the other hand, it can be shown that the slips on the physical plane (Fig. 4.14) are in general not straight, but curve, because shear boundary conditions may exist in real problems, changing the direction of the slips with depth (Scott 1963, Sokolovskiĭ 1960). This set of curve slips are called *slip lines*. If an x-y reference frame as shown in Fig. 4.14 is adopted, where the normal to the plane of an infinitesimal element is inclined at an angle  $\mu$  from the x-axis, the slip lines would be inclined at an angle  $\xi \pm \varepsilon$  from the x-axis. Consequently, stresses in x and y directions in terms of the principal stresses can be written as (Sokolovskiĭ 1960):

$$\sigma_{xx} = \frac{1}{2} \left( \sigma_1 + \sigma_3 \right) + \frac{1}{2} \left( \sigma_1 - \sigma_3 \right) \cos 2\xi$$
(4.38)

$$\sigma_{yy} = \frac{1}{2} \left( \sigma_1 + \sigma_3 \right) - \frac{1}{2} \left( \sigma_1 - \sigma_3 \right) \cos 2\xi$$
(4.39)

$$\tau_{xy} = \frac{1}{2} (\sigma_1 - \sigma_3) \sin 2\xi \tag{4.40}$$

Combining Eq. (4.37) and the Eqs. (4.38), (4.39) and (4.40), the equations of limiting stresses are obtained in the x and y directions:

$$\sigma_{xx} = \sigma(1 + \sin\phi\cos 2\xi) - H \tag{4.41}$$

$$\sigma_{yy} = \sigma(1 - \sin\phi\cos 2\xi) - H \tag{4.42}$$

$$\tau_{xy} = \sigma \sin\phi \sin 2\xi \tag{4.43}$$



Fig. 4.14. Orientation of slip lines (adapted from Sokolovskiĭ, 1965).

Analogously, the equations of limiting stresses in the normal  $\hat{n}$  and  $\hat{t}$  directions can be obtained (Sokolovskii 1960):

$$\sigma_n = \sigma \left[ 1 + \sin\phi \cos 2(\xi - \mu) \right] - H \tag{4.44}$$

$$\sigma_{t} = \sigma \left[ 1 - \sin\phi \cos 2(\xi - \mu) \right] - H \tag{4.45}$$

$$\tau_{nt} = \sigma \sin\phi \sin 2(\xi - \mu) \tag{4.46}$$

where  $\sigma = 1/2(\sigma_1 + \sigma_3) + H$ . When Eqs. (4.41), (4.42) and (4.43) are combined with the equations of equilibrium [Eqs. (4.29) and (4.30)], a set of two first order partial differential equations which govern the variation of the stresses in the x and y directions is obtained:

$$(1+\sin\phi\cos 2\xi)\frac{\partial\sigma}{\partial x}+\sin\phi\sin 2\xi\frac{\partial\sigma}{\partial y}-2\sigma\sin\phi\left(\sin 2\xi\frac{\partial\xi}{\partial x}-\cos 2\xi\frac{\partial\xi}{\partial y}\right)=\gamma\sin\alpha\qquad(4.47)$$

$$\sin\phi\sin 2\xi\frac{\partial\sigma}{\partial x} + (1-\sin\phi\cos 2\xi)\frac{\partial\sigma}{\partial y} + 2\sigma\sin\phi\left(\cos 2\xi\frac{\partial\xi}{\partial x} + \sin 2\xi\frac{\partial\xi}{\partial y}\right) = \gamma\cos\alpha \qquad (4.48)$$

Applying the *method of characteristics* (Hill 1950), the system of Eqs. (4.47) and (4.48), can be solved. Multiplying Eq. (4.47) by  $\sin(\xi \pm \varepsilon)$  and Eq. (4.48) by  $\cos(\xi \pm \varepsilon)$  and adding them, the following expression is obtained:

$$\left[\frac{\partial\sigma}{\partial x}\mp 2\sigma\tan\phi\frac{\partial\xi}{\partial x}-\gamma\frac{\sin(\alpha\mp\phi)}{\cos\phi}\right]\cos(\xi\mp\varepsilon)+\left[\frac{\partial\sigma}{\partial y}\mp 2\sigma\tan\phi\frac{\partial\xi}{\partial y}-\gamma\frac{\cos(\alpha\mp\phi)}{\cos\phi}\right]\sin(\xi\mp\varepsilon)=0$$
(4.49)

Using the auxiliary quantities *M* and *N*, as defined by:

$$M = \frac{\partial \sigma}{\partial x} \mp 2\sigma \tan \phi \frac{\partial \xi}{\partial x} - \gamma \frac{\sin(\alpha \mp \phi)}{\cos \phi}$$
(4.50)

$$N = \frac{\partial \sigma}{\partial y} \mp 2\sigma \tan \phi \frac{\partial \xi}{\partial y} - \gamma \frac{\cos(\alpha \mp \phi)}{\cos \phi}$$
(4.51)

Eq. (21) can be written as (Sokolovskiĭ 1960):

124

$$M + N \tan\left(\phi \mp \varepsilon\right) = 0 \tag{4.52}$$

Obtaining the characteristic differential equations requires the expression of the auxiliary quantities in ordinary derivative form:

$$d\sigma = \frac{\partial\sigma}{\partial x}dx + \frac{\partial\sigma}{\partial y}dy \tag{4.53}$$

$$d\xi = \frac{\partial\xi}{\partial x}dx + \frac{\partial\xi}{\partial y}dy$$
(4.54)

Introducing a new a new auxiliary quantity P,

$$P = \frac{d\sigma}{dx} \mp 2\sigma \tan\phi \frac{d\xi}{dx} - \frac{\gamma}{\cos\phi} \left[ \sin(\alpha \mp \phi) + \cos(\alpha \mp \phi) \frac{dy}{dx} \right]$$
(4.55)

*M* and *N*, can be rewritten as:

$$M = \frac{P dx \sin(\xi \mp \varepsilon)}{\sin(\xi \mp \varepsilon) dx - \cos(\xi \mp \varepsilon) dy}$$
(4.56)

$$N = \frac{P dx \cos(\xi \mp \varepsilon)}{\sin(\xi \mp \varepsilon) dx - \cos(\xi \mp \varepsilon) dy}$$
(4.57)

Note that the characteristics will be found when the numerator and denominator of the Eqs. (4.56) and (4.57) are simultaneously set to zero, and therefore the equations of characteristic are:

$$\frac{dy}{dx} = \tan(\xi \mp \varepsilon) \tag{4.58}$$

$$d\sigma \mp 2\sigma \tan \phi d\xi = \frac{\gamma}{\cos \phi} \left[ \sin \left( \alpha \mp \phi \right) dx + \cos \left( \alpha \mp \phi \right) dy \right]$$
(4.59)

Eq. (4.58) imposes the condition that the characteristics have to be inclined at  $(\xi \mp \varepsilon)^{\circ}$  from the x-axis. Therefore, they are indeed slip lines in the x-y space (Sokolovskii 1965). For cases when the body forces are aligned in the direction of vertical y-axis,  $\alpha = 0$  and Eq. (4.59) simplifies to:

$$d\sigma \mp 2\sigma \tan \phi d\xi = \frac{\gamma}{\cos \phi} \left[ \sin \left( \alpha \mp \phi \right) dx + \cos \left( \alpha \mp \phi \right) dy \right]$$
(4.60)

# Chapter 5. Design of Stable Concave Slopes for Reduced Sediment Delivery

This chapter was submitted as an original paper in the ASCE Journal of Geotechnical and Geoenvironmental Engineering, and has been revised according to the reviewers' suggestions. It is currently undergoing the second review. The co-authors of this work are Dr. Eric Drumm and Dr. Daniel Yoder. This article would be cited as:

Jeldes, I. A., Drumm, E. C., and Yoder, D. C. "Design of stable concave slopes for reduced sediment delivery." J. Geotech. Geoenviron. Eng., (in review).

# Abstract

Constructed slopes have traditionally taken a planar form. However, natural slopes are more likely to be concave in cross section, and laboratory and computational studies have demonstrated that concave slopes yield less sediment than planar slopes. With the current auto-guided construction equipment, it is now possible to construct slopes with concave profiles and a more natural appearance, yet a simple method to describe such concave slopes for a given level of mechanical stability does not exist. This article begins with the examination of concave shapes satisfying a desired degree of stability and compare results with those from Finite Element and Limit Equilibrium methods. An erosion model is used to demonstrate that the concave slopes proposed here can yield 15-40% less sediment than planar slopes with same Factor of Safety. Finally, a sensitivity analysis suggests that reasonable construction deviations do not compromise the stability of typical concave slopes.

# Introduction

Constructed slopes designed with traditional planar cross sections are encountered in most land development, including highway cut and fill sections, constructed embankments, and reclaimed mines. Planar landscape profiles are seldom encountered in nature, however, where curvilinear slopes with concave shapes usually arise as the result of evolutionary processes in fluvial systems and hillslopes (Leopold and Langbein 1962, Miyamoto et al. 2005, Rodríguez-Iturbe et al. 1992, Twidale 2007, Yang and Song 1979). Landforming approaches such as the Geomorphic Reclamation of mine lands (Toy and Chuse 2005) include the construction of concave-like shapes in both the transverse (cross-slope) and longitudinal (downslope) directions to create natural selfsustainable ecosystems (Martín-Duque et al. 2010) with improved erosion resistance (Schor and Gray 2007). Nevertheless, not all concave shapes are mechanically stable, requiring a rational definition of those concave slopes that provide the desired degree of stability expressed in terms of a desired Factor of Safety (FS). For example, Howard et al. (2011) point out the risk associated with the growing practice of shaping slopes to reflect natural regional landforms without appropriate stability and erosion analyses and without accounting for the limited precision of the construction equipment employed to build concave profiles. Hence, the objectives of this work are to 1) describe concave shapes that provide a desired degree of stability (or FS) for given soil properties, 2) provide a quantitative measure of the difference in soil loss (erosion) between concave and planar slopes that satisfy the same degree of mechanical stability, and 3) investigate the precision to which concave forms can be constructed, and how this affects the desired slope stability. To accomplish these objectives this work focuses on slopes with concavity in the longitudinal direction, which have been shown to deliver less sediment than transversely-convex, transversely-concave, longitudinally-convex and planar slopes (Rieke-Zapp and Nearing 2005).

# Background

## Mechanical stability of concave profiles

Although a number of studies have shown that concave slopes produce less sediment than planar slopes, few studies have linked the benefits of concave slopes to mass stability. Utili and Nova (2007) demonstrated by means of the upper bound method that concave log-spiral slopes result in higher FS's than planar slopes of same height and width and than planar slopes of same soil mass. Based on the slip line field theory and the associated characteristic equations, Sokolovskii (1960) found a longitudinal concave critical slope contour (FS = 1) for plain strain conditions. Sokolovskii's solution is the result of a numeric boundary value problem, which for a medium with unit weight ( $\gamma$ ) > 0 and friction angle ( $\phi$ ) > 0 has no analytical counterpart. Recently, Jeldes et al. (2013) proposed an approximate analytical expression to obtain coordinates of the critical contour for any desired combination of  $\phi$ , c,  $\gamma$  and slope height ( $H_s$ ). This expression brings a system of differential equations into a single algebraic expression, which highly simplifies the manipulation and solution as:

$$x = \begin{cases} 0 & , -h_{cr} \le y \le 0\\ A \Big[ \sigma_y (B-1)(cosec\phi-1) + H B (cosec\phi+1) \Big], & y > 0 \end{cases}$$
(5.1)

where:

$$A = \frac{\cos\phi}{2\gamma(1-\sin\phi)} \tag{5.2}$$

$$B = ln \left[ \frac{\sigma_y}{H} \left( \frac{1 - \sin \phi}{1 + \sin \phi} \right) + 1 \right] = ln \left[ \frac{\sigma_y}{H} K_a + 1 \right]$$
(5.3)

$$\sigma_{y} = \gamma y \tag{5.4}$$

130

$$H = c \cot \phi \tag{5.5}$$

$$h_{cr} = \frac{2c\cos\phi}{\gamma(1-\sin\phi)} \tag{5.6}$$

Here, H is the tensile strength of the soil,  $K_a = (1 - \sin \phi) / (1 + \sin \phi)$  is the Rankine active coefficient of earth pressure,  $\sigma_y$  is the geo-static vertical stress, and  $h_{cr}$  is the height of the tension zone. The equation describes a slope contour in the quadrant with x-axis positive to the right and y-axis positive downward (Fig. 5.1), and with  $h_{cr}$  lying above the x-axis from the coordinates (0,0) to  $(0, -h_{cr})$ . A limiting slope height  $(h_L)$  was also found, defining where the failure mechanism changes. Slopes of heights  $H_s < h_L$  indicate a toe-failure mechanism with FS values slightly greater than 1, while those with  $H_s > h_L$  indicate a face-failure mechanism with a steady FS  $\approx 1$ . A linear correlation between  $h_L$  and  $c/\gamma$  was approximated by

$$h_{L} = 15.8 \ c \,/\, \gamma + 0.3 \tag{5.7}$$

which can be used as a preliminary estimate of the most likely failure mechanism for a concave slope (Jeldes et al. 2013). While in practice the sharp bluff at the top of the tension zone may be removed or rounded (as indicated by the dashed line in Fig. 5.1) adding conservatism, the full tension height was included in the development of the solution for mathematical completeness.

## **Concave slopes and soil erosion**

Experimental and computational studies have shown that slopes with longitudinally-concave shapes yield less sediment than slopes with planar or other curvilinear shapes. Thus, this work focuses on the difference in soil loss between longitudinally concave (referred to henceforth as concave) and planar slopes, which are the two shapes of interest in this study.



**Fig. 5.1.** Illustration of the reference frame and the components of the concave slope formulation. In practice the sharp bluff at the top of the slope may be rounded, as indicated by the dashed line.

Since erosion rates increase with slope length and steepness, for a concave slope the increased erosion due to the increased length is partially counteracted by the decreased erosion due to decreasing slope steepness. In examining this, Young and Mutchler (1969) used field-scale plots on a silt soil and measured 10% less sediment loss on concave slopes with the same average inclination as planar ones. D'Souza and Morgan (1976) used laboratory scale models with a sandy soil, and measured 20-25% less soil loss on concave slopes with the same average inclination as planar slopes. Rieke-Zapp and Nearing (2005) reported 75% less soil loss for concave slopes with similar surface area as planar slopes using experiments conducted on laboratory models with silt soil. Reductions in erosion for concave slopes can also be predicted via computational models. Meyer and Kramer (1969), using the Universal Soil Loss Equation (USLE), reported less soil loss for concave slopes of the same height and average steepness as planar slopes, and the difference in soil loss increased with horizontal slope distance. Williams and Nicks (1988), employing the field-scale model CREAMS, computed approximately 50% lower soil loss for a concave slope of the same average inclination as a planar slope; grass cover and filter strips did not alter the ratio of soil losses between concave and planar slopes. Hancock et al. (2003), using the landform evolution model SIBERIA, computed up to 80% less soil loss for concave slopes of the same horizontal length and average inclination as planar slopes, and showed that sediment loss decreases as the slope becomes more concave. Lee et al. (2004) simulated moderately intense storm conditions via the Hillslope Erosion Model and found that concave slopes yielded 17, 22, 24, and 28% less sediment loss than planar slopes when sandy clay, clay, silt, and silty clay soils were used in the simulation. Priyashantha et al. (2009), also using SIBERIA, computed approximately 50% reduction in soil loss between planar and concave slopes with same overall inclinations, and again found that the reduction occurs regardless of vegetation and climate type.

While all the aforementioned erosional studies reported significantly less soil loss for concave slopes, most compared the difference in soil loss between planar and concave slopes with similar average steepness, and none of them linked the concave geometry to mechanical stability considerations. Accordingly, the investigation begins with the examination of concave shapes satisfying a desired degree of stability, and then proceeds by investigating their ability to reduce sediment yield. This is accomplished by comparing the soil loss on concave and planar slopes with the same FS. Finally, the impact of construction precision on the concave slope stability is explored.

# Methods

## Mechanical stability

In this section concave slopes for a given design FS for long-term stability are defined and the undrained shear strength required such that short-term stability is unlikely to govern is determined. This approach is conservative in that it neglects any additional shear strength that may be present due to partial saturation and resulting soil suction.

## Stable concave slopes for $\phi > 0$ (long-term stability) with pre-selected FS

While Eqs. (5.1) - (5.6) generate concave slopes at limiting equilibrium, a design factor of safety  $(FS_D) > 1$  is desired for design and construction. The expression for the concave slope can be extended to any desired FS if a strength reduction factor (Griffiths and Lane 1999, Spencer 1967) equal to  $FS_D$  is used to reduce the strength parameters as follows:

$$\phi^* = \arctan\left[\frac{\tan\phi}{FS_D}\right] \tag{5.8}$$

$$c^* = \frac{c}{FS_D} \tag{5.9}$$

where  $\phi^*$  and  $c^*$  are reduced friction angle and cohesion employed in the design to achieve the desired FS. To obtain a concave slope satisfying a pre-selected  $FS_D$ , replace  $\phi$  and c in Eqs. (5.1) – (5.6) with  $\phi^*$  and  $c^*$  and determine the x-y coordinates of the concave slope. The resulting concave shape will satisfy a FS  $\approx FS_D$  under the original  $\phi$  and c strength conditions.

The proposed methodology was compared with results obtained via Limit Equilibrium Method (LEM) and Finite Element Method (FEM) analyses for soils with  $\phi = 20^{\circ}$ , 30°, and 40° and  $c/\gamma$  ranging from 0.2 to 5 m, where the range c = 5 - 40 kN/m<sup>3</sup> covers most of the cohesion values reported by Mesri and Abdelghaffar (1993), and the range  $\gamma = 10 - 23$  kN/m<sup>3</sup> covers most materials from clays to coarse granular soils (NAVFAC 1986). Each combination of  $\phi$  and  $c/\gamma$  describing a unique concave slope was analyzed for heights  $H_s$  ranging from 2 to 100 m. The FEM analyses were conducted using *Phase2* (Rocscience Inc. 2011), and the LEM analyses using *Slide* (Rocscience Inc. 2011). In the FEM analyses, an elastic perfectly plastic model with a Mohr-Coulomb yield criterion and a non-associated flow rule (zero dilatancy angle) was employed since it has been shown to provide reliable FS's (Griffiths and Lane 1999). Also, nominal values for Young's modulus ( $E = 2x10^4$  kPa) and Poisson's ratio (v = 0.3) were used, because they do not influence the calculation of the FS (Cheng et al. 2007, Griffiths and Lane 1999). In the LEM

analyses, the Morgenstern-Price, the Spencer and the Simplified Bishop Methods were used with over 30,000 critical surfaces analyzed for each set of variables.

# Stability check for short- term ( $\phi = 0$ ) loading conditions

Once the long-term ( $\phi > 0$ ) stable concave shapes are defined, it is necessary to investigate under what conditions short-term failure may occur, e.g., failure during the construction period where the stability is controlled by the undrained shear strength ( $S_u$ ), which is the soil shear strength for  $\phi = 0$  conditions. The question then becomes: what would be the minimum value of undrained shear strength ( $S_u^{\min}$ ) such that the concave slope defined for long-term stability is also stable for short-term conditions? This problem was approached by making use of the critical slope concept. A critical contour implies that slopes less steep than this contour are safe, and steeper ones are unsafe (under same soil properties and slope height). Thus, any critical contour defined for longterm stability that is less or equally steep than a critical contour defined for short-term ( $\phi = 0$ ) will be stable in the short-term for a given height *y*. Although the long-term critical contour can be obtained via Eqs. (5.1) – (5.6), the critical slope contour for the  $\phi = 0$  case can be obtained as a closed form solution of the characteristic equations (Sokolovskii 1965):

$$x_{u}(y) = -\frac{2S_{u}}{\gamma} \ln \left[ \cos(\frac{\gamma y}{2S_{u}}) \right], \quad y \ge 0$$
(5.10)

Here  $x_u$  defines the horizontal projection of the slope contour. In the tension zone  $(-h_{cru} \le y \le 0)$  $x_u = 0$ , and  $h_{cru} = 2S_u / \gamma$  (Fig. 5.2). This concept is illustrated in Fig. 5.2, where the long-term



**Fig. 5.2.** Long-term ( $\phi > 0$ ) critical contour vs short-term ( $\phi = 0$ ) critical contours defined by different  $S_u$  values.

critical contour is steeper than the short-term critical contours 1, 2 and 3 that have undrained shear strengths  $S_{u1}$ ,  $S_{u2}$ , and  $S_{u3}$  respectively, and therefore it is unsafe in the short-term. The opposite yet desirable case is shown for contour 4, where  $S_{u4}$  is sufficiently large that the global stability is governed by the long-term condition. The value of  $S_u^{\min}$  to satisfy this condition is found at the intersection of the long and short-term contours. To find  $S_u^{\min}$ , we must a) enforce that  $h_{cru} \ge h_{cr}$ , and b) numerically solve the transcendental equation  $x_u(y) = x_d(y + \Delta h_{cr})$ , where  $\Delta h_{cr} = h_{cru} - h_{cr}$ ,  $\Delta h_{cr} > 0$ . It should be noted that  $S_u^{\min}$  is a conservative estimate of the necessary undrained shear strength and can be taken as an upper limiting value of the exact solution.

#### Soil erosion and sediment yield

Soil loss for concave and planar slopes with same FS's was investigated via the widely recognized Revised Universal Soil Loss Equation RUSLE2 (USDA-ARS 2008):

$$A = R \cdot K \cdot LS \cdot C \cdot P \tag{5.11}$$

where the predicted soil loss *A* (units of Mg/ha/y) is directly proportional to: the rainfall erosivity *R* (units of  $MJ \cdot mm/ha \cdot h \cdot y$ ), quantifying the rainfall's erosive potential; the soil erodibility *K* (units of  $Mg \cdot ha \cdot h/ha \cdot MJ \cdot mm$ ) defining the soil's susceptibility to that erosivity; the topographic factor *LS* (dimensionless) representing slope length and steepness effects; the surface cover factor *C* (dimensionless); and the conservation practices factor *P* (dimensionless).

Through RUSLE2 pairs of planar and concave slopes [Eqs. (5.1) – (5.6)] that had the same FS under the same values of  $\phi$ , c,  $\gamma$  and  $H_s$  (Fig. 5.3) were investigated. For soil loss calculations,



Fig. 5.3. Illustration of concave and equivalent planar slopes (same FS and slope height).

the concave slopes were idealized using ten to twenty linear segments. Since RUSLE2 determines soil loss over the horizontal projection of the slope, for slope segments steeper than 45° (which are not defined in RUSLE2), the soil loss was calculated assuming a planar slope of 45° over the total horizontal length occupied by the steeper part of the contour. Three different soil textures were employed for the analyses: a silt loam (high erodibility), a clay (moderate erodibility), and a sandy loam (moderately low erodibility). A range of possible K values for each soil texture (Haan et al. 1994, USDA-ARS 2008) was used (Table 5.1), which conservatively encompasses a large variety of soils found in practice, even mine reclaimed soils (Hoomehr et al. 2014). All the simulations were conducted for two locations with very different climates: Dakota County, Minnesota (R =2,180 MJ·mm/ha·h·y) representing dry weather conditions and Monroe County, Florida (R =10,500 MJ·mm/ha·h·y) representing high rainfall conditions. For simplicity, a bare surface (C = 1) and no conservation practices (P = 1) were modeled. A representative value of friction angle was assumed for each soil texture (Budhu 2011, Hough 1957). Values of  $c/\gamma$  ranging from 0.2 to 3 m were employed for the clay soils and from 0.2 to 2 m for the silt and sandy loam soils. Table 5.1 summarizes the erosion and mechanical properties of the investigated soils.

#### **Construction and sensitivity**

The widespread use of high accuracy auto guidance GPS-based construction equipment for grading and earthwork operations suggests that the construction of more complex slope shapes can become more commonplace. The effects of construction precision on the stability were investigated by considering the worst case scenario: the vertical component of the contour is constructed deeper than designed, resulting in a steeper slope. Although the vertical accuracy of 3D grade control systems is commonly within 30 mm (Trimble 2013), an accuracy T = 200 mm

<b>RUSLE</b> soil	Co	omposition [USDA],	(%)	Equivalent	Range of	Assumed
texture	Clay	Silt	Sand	USCS	erodibility	$\phi(^{\circ})$
[USDA]	<0.002 mm	<0.05-0.002 mm	<2-0.05 mm	classification	$(K)^{**}$	Ψ()
Silt Loam	20	60	20	ML	0.037-0.057*	25
Clay	45	28	27	CL	0.032-0.042	20
Sandy Loam	10	25	65	SM	0.026-0.037*	35

Table 5.1 Soil composition, classification, erodibility and assumed internal friction angle of investigated soils

\* Upper bound reported for cases of low permeability and low organic matter content.

\*\*All K values in SI units of  $Mg \cdot ha \cdot h / ha \cdot MJ \cdot mm$ .

was used, which may be also achieved by conventional equipment. Because the horizontal accuracy of GPS receivers is usually up to 3 times higher than the vertical (Berber et al. 2012), the horizontal errors are expected to stay in the millimeter scale, with no significant effects on stability. The effect of construction accuracy on the stability was investigated by: a) lowering by 200 mm the vertical component of various critical concave contours (FS = 1) for  $\phi = 20^{\circ}$ , 30°, and 40° and  $c/\gamma = 0.2$ , 1, and 5 m [Eqs. (5.1) - (5.6)]; and b) conducting stability analyses on them via the Simplified Bishop's Method.

# **Results and Discussion**

# Concave slopes with pre-selected FS's for long-term ( $\phi > 0$ ) conditions

Results from LEM analyses are shown in Fig. 5.4 for the case  $\phi = 30^\circ$ , where the computed FS values are plotted against the dimensionless value  $H_s\gamma/c$ . Each point represents a stability analysis on a concave slope of height  $H_s$ , cohesion c, and unit weight  $\gamma$ , and each  $\phi^*$  reflects a different value of  $FS_D$ . Note that  $H_s\gamma/c$  is the inverse of Taylor's stability number N (Taylor 1948). A change in failure mechanism is observed to occur at similar values of 1/N regardless of  $\phi^*$ . The change in failure modes observed on the critical concave slopes (Jeldes et al. 2013) is not altered by the use of a strength reduction factor, and Eq. (5.7) can still be used to estimate the most probable mode of failure with  $c^*/\gamma$  instead of  $c/\gamma$ . Concave slopes with  $H_s > h_L$  show steady computed values of FS  $\approx FS_D$ , while those with  $H_s < h_L$  showed FS values higher (more conservative) than the  $H_s > h_L$  case, with values increasing as  $H_s\gamma/c$  decreases. Similar results were obtained for the  $\phi = 20^\circ$  and  $40^\circ$  cases.



**Fig. 5.4.** Computed FS vs  $H_s \gamma / c$  for  $\phi = 30^\circ$  and  $FS_D = 1.00, 1.25, 1.5$  and 2.00. Lines are results from Simplified Bishop's method, triangles are results from Morgenstern-Price's method and squares are results from Spencer's method.

It is important to note that the methodology proposed here is based on the slip line field theory, which requires a value of  $c \neq 0$  for a valid mathematical solution to exist. However, the method can still be used to simulate long-term conditions using a very small value of cohesion, as it is illustrated in the example problem. It must be also emphasized that Eqs. (5.1) - (5.6) constitute an approximate solution in analytical form, so even when  $H_s > h_L$  there may not be a perfect match between  $FS_D$  and the actual FS. As seen in Fig. 5.4, the best match occurs when  $\phi^* \approx 20^\circ$ . Concave slopes with  $\phi^* > 20^\circ$  yield FS values higher than the design  $FS_D$  (4% higher when  $\phi^* = 30^\circ$  and 6% higher when  $\phi^* = 40^\circ$ ), which is conservative for design. On the other hand, slopes with  $\phi^* <$ 20° will have FS <  $FS_{D}$ , which is not conservative. This is shown in Fig. 5.4 for the case  $\phi^*$  = 16.1°, where the resulting FS = 1.92 is 4% less than the desired  $FS_D = 2$ . Nevertheless, the deviations from the target  $FS_D$  obtained through the methodology do not compromise its applicability for design, considering that neither the geometry nor the shear strength parameters are usually determined within ±6% accuracy (Duncan 1996). Furthermore, since the FS values calculated by limit equilibrium methods that satisfy all conditions of static equilibrium can vary as much as 12% (Duncan 1996), an accuracy of  $\pm 6\%$  is certainly acceptable. Once the desired shape for the design FS has been obtained, verifications can be conducted employing commercial slope stability software. In the same way, the effects of transient ground water flow and/or external surcharges on the FS can be investigated, which extends the applications of this solution to other field conditions.

## Stability check for short-term ( $\phi = 0$ ) conditions

The minimum undrained shear strength  $S_u^{\min}$  that the soil must possess for concave slopes to have FS > *FS<sub>D</sub>* in the short term can be obtained from Fig. 5.5 as a function of the effective strength parameters  $\phi^*$  and  $c^*$ . Fig. 5.5 is a solution chart that assembles the numerical results from Eqs. (5.1) and (5.10) and is presented in terms of the dimensionless parameters  $y \cdot \gamma / c^*$  and  $S_u^{\min} / (c^* \cdot FS_D)$  for a range of  $\phi^*$ . Here, y is the height of the slope below the tension zone  $(y = H_s - |h_{cr}|, y > 0)$ . The inner chart in Fig. 5.5 provides higher resolution for low values of  $y \cdot \gamma / c^*$ . Note that Fig. 5.5 defines the conditions under which long-term stability will always govern, so the solutions are conservative. Also, this chart is based on an assumed failure mechanism, where the critical surface exists from or above the toe. Consequently, this solution provides a first check for the short-term stability of concave slopes and does not replace the necessity of conducting stability analyses once the undrained shear strength of the soil has been determined.

## Soil loss from the mechanically stable concave slopes

Soil loss results obtained from RUSLE2 for the silt loam are shown in Fig. 5.6. All the analyses were conducted on concave and planar slopes with FS =  $FS_D = 1$  ( $\phi = \phi^*$  and  $c = c^*$ ), since they provide generic results that can be extended to any FS. Figs. 5.6a and 5.6b compare the predicted total soil loss for concave and planar slopes of a silt loam with  $c^*/\gamma = 1$  under wet weather conditions in Monroe County, FL, and dry conditions in Dakota County, MN. Planar slopes



Fig. 5.5. Solution chart for short-term stability check.



**Fig. 5.6.** Erosion analyses for a silt loam soil with  $\phi = 25^{\circ}$ ; a) total soil loss for Monroe County, FL; b) total soil loss for Dakota County, MN; c) difference in soil loss for the Dakota (dry) and Monroe (wet) Counties; d)  $A_c / A_p$  vs  $H_s \gamma / c^*$ . Units: A (Mg/Ha/y),  $A_p - A_c$  (Mg/Ha/y), R (MJ·mm/ha·h·y), K (Mg·ha·h/ha·MJ·mm), and  $H_s$  (m).

yielded more sediment than equally stable concave slopes for all *K* and *H<sub>s</sub>* values investigated, regardless of climate. Soil loss increased with *K* and *H<sub>s</sub>* in both profiles, but the upper surfaces (soil loss from planar slopes) remain constantly above the lower surfaces (soil loss from concave slopes). Fig. 5.6c compares, for the two weather conditions, the variation in  $A_p - A_c$  ( $A_p$  = soil loss from planar slopes,  $A_c$  = soil loss from concave slopes), or the erosion savings incurred using equally stable concave slopes. This illustrates that concave slopes can reduce erosion to a much larger extent in climates with high precipitations and stronger rainfall intensities.

Generalized results for the silt loam soil are shown in Fig. 5.6d in terms of the ratio  $A_c / A_p$  and the stability number  $H_s \gamma / c^*$ . Since the same RUSLE2 *R* and *K* were used in both planar and concave slopes, they cancel out when the results are presented in terms of  $A_c / A_p$  and the only factor affecting differences in the soil loss is the shape of the slope profile. The use of  $H_s \gamma / c^*$ enables us to represent the results of different slopes with the same degree of stability for a given value of  $\phi^*$ . For the range of  $H_s \gamma / c^*$  values investigated,  $A_c / A_p$  ranges from 0.85 - 0.60 (Fig. 5.6d), indicating that concave slopes yield 15 - 40% less sediment than their planar counterparts. Similar results are shown in Figs. 5.7 and 5.8 in terms of  $A_p - A_c$  and the  $A_c / A_p$  ratio for  $\phi^* =$ 20° (clay soil) and  $\phi^* = 35^\circ$  (sandy loam soil) respectively. The computed  $A_c / A_p$  values show soil loss reduction of the same order and suggest that the concave slopes proposed in this article will reduce sediment delivery by 15-40% regardless of soil erodibility and climate. This finding



**Fig. 5.7.** Erosion analyses for the clay soil with  $\phi = 20^{\circ}$ : a) difference in soil loss for the Dakota (dry) and Monroe (wet) Counties; b)  $A_c / A_p$  vs  $H_s \gamma / c^*$ . Units:  $A_p - A_c$  (Mg/Ha/y), R (MJ·mm/ha·h·y), K (Mg·ha·h/ha·MJ·mm), and  $H_s$  (m).



**Fig. 5.8.** Erosion analyses for the sandy loam soil with  $\phi = 35^\circ$ : a) difference in soil loss for the Dakota (dry) and Monroe (wet) Counties; b)  $A_c / A_p$  vs  $H_s \gamma / c^*$ . Units:  $A_p - A_c$  (Mg/Ha/y), *R* (MJ·mm/ha·h·y), *K* (Mg·ha·h/ha·MJ·mm), and  $H_s$  (m).

is extendable to slopes having natural or artificial surface covers (C < 1) where the erosion is proportionally reduced for both concave and planar slopes.

In Table 5.2 a comparison of the relative reduction in soil loss  $(1 - A_c / A_p)$  obtained here with values reported in the literature is offered. Although a wide range of erosion reductions were found there, it should be noted that the concave and planar slopes in those studies likely do not satisfy the equal degree of mechanical stability (FS) purposely enforced here. Similar reductions were obtained among the studies except for those reported by Meyer and Kramer (1969), Hancock et al. (2003) and Rieke-Zapp and Nearing (2005), for which the reductions were significantly higher. Meyer and Kramer (1969), and Rieke-Zapp and Nearing (2005) observed that an important amount of sediment deposition occurred at the lowest part of the concave contours, which may explain the higher overall reduction in sediment yield for those studies. In contrast, sediment deposition was not indicated in the RUSLE2 prediction obtained here. Each segment of the concave slopes proposed in this article is the steepest possible to satisfy mass equilibrium requirements, and sedimentation on the profile will likely not occur, especially at large values of  $\phi^*$ .

## Sensitivity to construction accuracy

Fig. 5.9 shows results obtained from the construction sensitivity analyses for concave slopes with  $FS = FS_D = 1$  ( $\phi = \phi^*$  and  $c = c^*$ ) for a range of  $\phi^*$  and  $H_s \gamma / c^*$  values, with results expressed in terms of the percent decrease in FS due to the construction error. Slopes with  $H_s \leq h_L$  were excluded from the analyses, since it was previously shown that their performing FS's are higher than required (conservative), and they did not present stability variations below equilibrium.

150

Reference	Method	Details	$1 - A_c / A_p$ (%)
Young and Mutchler (1969)	Experimental	Field plots	10
Meyer and Kramer (1969)	Computational	USLE	0-85
D'Souza and Morgan (1976)	Experimental	Laboratory	20-25
Williams and Nicks (1988)	Computational	CREAMS	50-55
Hancock et al. (2003)	Computational	SIBERIA	50-80
Lee et al. (2004)	Computational	Hillslope Erosion	17-28
Rieke-Zapp and Nearing (2005)	Experimental	Laboratory	75
Priyashantha et al. (2009)	Computational	SIBERIA	50
This study	Computational	RUSLE2	15-40

**Table 5.2** Reduction in soil loss  $(1 - A_c / A_p)$  from concave contours reported in the literature



Fig. 5.9. Decrease in FS due to low precision construction (200 mm vertical) of the concave slope.

The results demonstrated that the FS's are not significantly influenced by improper construction within the 200 mm of vertical accuracy. The largest difference was found for slopes with  $c^* / \gamma = 0.2$  m, where 5 - 5.5% lower FS's were obtained. Considering that the analyses were conducted with a vertical accuracy only 1/6 that which can be achieved with typical auto-guidance equipment, it is concluded that the stability of concave slopes is not significantly compromised by the precision of GPS-guided equipment. In addition, practicality suggests that during construction, the top vertical section will likely not be built, but will likely be rounded back to an average 2:1 (H:V) or similar inclination. This in turn will increase the stability and add conservatism to the slope design.

# **Illustrative Examples**

The design of concave slopes for long-term mechanical and erosion stability is illustrated for two hypothetical soils in Monroe County, FL, indicating high-erosivity conditions. FEM analyses were conducted to verify the mechanical stability for each case. The short-term stability was checked, followed by erosion analyses for concave and planar slopes with equal FS. The two selected soils were sand (SM or SC) and a fine grained soil (ML or CL), with the assumed mechanical properties and design parameters for each case as shown in Table 5.3. A  $FS_D = 1.5$  was required for each slope.

## Finding the concave profile for long-term stability:

Table 5.3 summarizes the initial steps: i) obtain  $\phi^*$  and  $c^*$  with Eqs. (5.8) and (5.9); and ii) estimate the failure mechanism via  $h_L$  [Eq. (5.7)] with  $c^*$  instead of c. The x-coordinates of the concave profile were then obtained by introducing  $\phi^*$  and  $c^*$  in Eqs. (5.1) – (5.6) for different values of y.

						Step i Step ii			i		
Hypothetical soil	FS <sub>D</sub>	H <sub>s</sub> m	φ (°)	<i>c</i> kPa	γ kN/m <sup>3</sup>	φ* (°)	c <sup>*</sup> kPa	c* / γ m	h <sub>L</sub> m	Condition	Predicted failure
Sand (e.g. SM or SC)	1.5	15	35	15	19	25.02	10.00	0.53	8.6	$h_L < H_s$	Face failure FS $\approx FS_D$
Fine grained (e.g. CL or ML)	1.5	30	25	35	17	17.27	23.33	1.37	22	$h_L < H_s$	Face failure FS $\approx FS_D$

Table 5.3 Geometry, soil properties, and probable failure mechanisms of illustrative examples

For the sandy slope,  $h_L = 8.6 \text{ m} < 15 \text{ m}$  and the concave slope will have a face failure mechanism with a performing FS  $\approx FS_D = 1.5$ . Verification of this condition was provided via FEM analysis (Fig. 5.10), which yielded FS = 1.51. The equivalent planar slope with FS = 1.5 is depicted with dashed lines in Fig. 5.10. Similarly, the fine grained slope has a face failure ( $h_L = 22 \text{ m} < 30 \text{ m}$ ). In this case, since  $\phi^* < 20^\circ$  a small error was introduced and the FEM analysis indicated a FS = 1.45. Note that the horizontal extent of the concave and planar slopes is similar, but there is less material involved in the concave slope. Since an analytical expression for the concave slope is now available, this could be incorporated into standard cut and fill balance calculations. On projects with significant cut slopes, the additional excavation may increase costs. However, on fill slopes and embankments, and especially in mine reclamation where a shortage of material is common, the construction of concave slopes could be an advantage and provide the benefit of more natural looking slopes and lower sediment loads.

#### Checking the short-term stability of the concave slope:

A conservative verification can be made using Fig. 5.5. For the sandy slope ( $\phi = 35^{\circ}$  and  $\phi^* = 25^{\circ}$ ),  $y = H_s - |h_{cr}| = 13.4$  m and  $y \cdot \gamma / c^* = 25.4$ . From Fig. 5.5,  $S_u^{\min} / (c^* \cdot FS_D) = 5.5$ , and thus the long-term stability of the sandy slope will govern as long as  $S_u^{\min} \ge 83$  kPa. Similarly, for the fine grained soil y = 26.3,  $y \cdot \gamma / c^* = 19.2$ , and the long-term stability will govern provided that  $S_u^{\min} \ge 150$  kPa.



**Fig. 5.10.** Example case for sandy soil:  $\phi = 35^{\circ}$ ,  $c = 15 \text{ kN/m}^2$ ,  $\gamma = 19 \text{ kN/m}^3$  and  $H_s = 15 \text{ m}$ . Required FS = 1.5. Shear strains showed for SRF = 1.52 to emphasize failure mode.
#### **Erosion analyses:**

Total soil loss was computed for low and high erodibility values for each soil texture (Table 5.4) for the Monroe County, FL weather conditions. The results (Table 5.4) reveal that the sandy concave slope yields on average 24% less sediment than its planar counterpart. Similarly, the fine grained concave slope yields on average 35% less sediment than a planar slope with the same FS.

# Conclusions

Concave slopes not only resemble natural contours, but also have superior erosion resistance. The design of concave slopes requires: a) the definition of concave shapes that provide a desired FS, b) a quantitative measure of the erosion/sediment delivery reduction, and c) determination of possible loss of mechanical stability due to improper construction. In this article:

- An approximate expression was described to provide the coordinates for concave slopes that satisfy a desired FS (long-term conditions) for any combination of φ, c, γ and H<sub>s</sub>. The methodology was shown to be conservative for φ<sup>\*</sup> > 20°, but provided slightly non-conservative values of design FS for cases with φ<sup>\*</sup> < 20°. The errors introduced by the methodology are less than ±6% for the majority of the cases, which is smaller than the accuracy to which the strength parameters are typically determined.</li>
- A simplified graphical solution was developed to estimate the required undrained shear strength  $S_u^{\min}$  such that the design long-term ( $\phi > 0$ ) stability is assured. The value of  $S_u^{\min}$  for concave slopes is obtained as a function of the effective strength parameters.
- Results from RUSLE2 analyses indicate that the concave slopes proposed here yield 15-40% less sediment than planar slopes of equal FS, regardless of soil erodibility and weather

Soil	φ* (°)	$\frac{1}{N^*} = \frac{H_s \gamma}{c^*}$	<i>R</i> (MJ·mm/ ha·h·y)	<i>K</i> (Mg·ha·h/ ha·MJ·mm)	A <sub>c</sub> (Mg/Ha/y)	$1 - \frac{A_c}{A_p}$	
						%	avg
Sand	25.02	28.5	10,500	0.026	557	22	24
				0.037	748	25	
Fine grained	17.27	21.8	10,500	0.037	1348	34	35
				0.057	2347	37	

 Table 5.4 Soil parameters and results from erosion analyses of illustrative examples

conditions. Although the analyses were conducted on slopes with  $FS_D = 1$ , the findings are valid for slopes with other FS's, as shown in the illustrative example.

• Results from the sensitivity analyses reveal that the stability of concave slopes is not significantly influenced by errors in the constructed profile of as great as 200 mm of vertical deviation. Therefore, the stability is not compromised when typical high accuracy GPS construction equipment is employed for concave slope construction.

The reduction of sediment delivery from constructed slopes is an important consideration in limiting environmental impact during land reclamation and site development. The proposed method suggests a rational design procedure which makes concave slopes an attractive ecotechnical alternative for more sustainable and natural appearing earthwork construction.

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# **Chapter 6. Sustainable Concave Slopes**

This chapter is soon to be submitted to the Journal of Soil and Water Conservation or Journal of Hydrological Processes. The co-authors of this work are Dr. Daniel Yoder and Dr. Eric Drumm. Excerpts from this chapter will be presented at the 2014 ASABE International Meeting in Montreal, QC, July 13-16, 2014

## Abstract

While manmade slopes are traditionally constructed with planar cross sections, natural stable slopes are usually curvilinear rather than planar. These curvilinear slope shapes with significant concave portions are obtained as a result of evolutionary processes where entropy is maximized through slope shape adjustments. Evidence suggests that this adjustment results in equilibrium between rain-driven erosion and sediment transport, yielding stable concave slopes with relatively constant shapes. This equilibrium shape is characterized by a parallel retreat with a uniform erosion rate over time. Concave slopes are more likely to maintain stable profiles than are planar slopes, where large amounts of soil are removed and delivered during the process of erosion equilibrium adjustment. Since a truly stable slope must be stable in terms of both mechanical and erosion processes, it becomes important to describe concave slopes in rainfall-erosion equilibrium and evaluate their mechanical stability. In this article concave profiles in rainfall erosion equilibrium are identified and described based on the well-known RUSLE2 model. Results indicate the existence of a family of potential slope shapes satisfying the condition of uniform erosion rate along the slope. Making use of the critical slope contour concept for a given mechanical soil strength, those steady concave shapes that also satisfy long-term stability are investigated. A true steady-state equilibrium slope is obtained when both erosion equilibrium and long-term mechanical stability are achieved. This suggests that concave slopes can be constructed to achieve 165

both minimal steady-state erosion equilibrium as well as mechanical stability, leading also to more "natural" landforms. Constructing slopes to reflect these stable shapes will provide more aesthetically pleasing results and minimize sediment delivery during initial slope adjustment.

## Introduction

Concave profiles have been presented as a way to optimize mechanical stability and water erosion resistance of slopes (Jeldes et al. 2014, Jeldes et al. 2013). In that work, efforts were concentrated on a) defining mechanically stable concave contours with a controlled degree of stability, and b) quantifying their effectiveness in reducing erosion and sediment delivery while meeting mass equilibrium requirements. This approach results in slopes that appear more natural, and constitutes an eco-technical solution to the design of slopes as a function of geotechnical parameters familiar to engineers. While effective in reducing erosion and providing mechanical stability, these concave profiles may not be in equilibrium from an erosion perspective. Evidence exists that natural fluvial systems seek erosion and sediment transport equilibrium, and in doing so may adjust their geometry in order to achieve a steady-state form (Yang and Song 1979) that will be somewhat unchangeable over time. This shape is usually referred as an equilibrium shape, with a concavelike form characterized by a parallel retreat (Penck 1953, Schumm 1956, Twidale 2007). In other words, man-made slopes (e.g., planar slopes) undergo continuous changes in slope morphology triggered by water erosion forces in order to ultimately achieve this equilibrium condition. During this process large amounts of soil are eroded.

Geomorphological evolution toward an erosion equilibrium state does not imply that the resulting slopes are mechanically stable. A true equilibrium state will be achieved only when both erosion 166

equilibrium and long-term mechanical stability are satisfied. The unchangeable shape and the minimization of sediment delivery to a relatively constant rate are attractive characteristics that can positively impact the way slopes are constructed, and can promote a more natural and "sustainable" design. The question becomes whether we can describe fully sustainable slopes, where both erosion and mechanical equilibrium requirements are met.

In this article the concept of steady-state landforms is explored, and from a conceptual perspective, the changes in slope morphology toward a concave erosion equilibrium shape are described. Based on the fundamentals of the RUSLE2 erosion model, concave profiles in water erosion equilibrium are identified and described, and a simple approach to discern between long-term mechanically stable and unstable erosion equilibrium shapes is proposed for any given combination of soil stresses and strength. A definition of the approximate limiting erosion rate at which equilibrium erosion shapes become mechanically stable and thus sustainable is explored, and a mathematical expression to obtain this limiting erosion rate is offered as a function of the Mohr-Coulomb parameters.

# Background

#### Concave profiles as erosion equilibrium shapes

In nature slope profiles are rarely uniform, but are usually curvilinear. This seems to be the result of a natural geomorphological evolution process of streams and slopes, where an optimum equilibrium profile is sought (Schor and Gray 2007). According to Leopold and Langbein (1962), the distribution of the energy in a fluvial system tends toward a state of maximum entropy, which is characterized by a uniform energy dissipation, where the same energy dissipation is experienced 167 at each point in the landscape (Miyamoto et al. 2005). Leopold and Langbein (1962) showed that the most probable longitudinal profile in a fluvial system is concave. Later studies have validated Leopold and Langbein's ideas (Molnár and Ramírez 1998, Rodríguez-Iturbe et al. 1992, Yang and Song 1979), but also have introduced new concepts such as minimum entropy production as a characteristic of steady equilibrium of a geomorphic system (Miyamoto et al. 2005), and the concept of minimum rate of energy dissipation experienced by equilibrium shapes (Yang and Song 1979). According to Yang and Song (1979), if fluvial systems deviate from an equilibrium state, the properties of slope inclination, roughness, channel geometry, and flow velocity will experience continuous adjustments in order to again achieve the steady equilibrium state. Other authors have also proposed concave profiles as the equilibrium or steady shape in fluvial systems (Goldrick and Bishop 2007, Larue 2008, Smith et al. 2000, Snyder et al. 2000).

Steady concave shapes have also been proposed as the result of hillslope evolution. According to Pelletier and Rasmussen (2009), slopes in more humid areas tend to be more complex in cross section (concave slope with convexity at the upper portion), while slopes in more arid regions tend to have concave forms. According to Schumm (1956), complex slope shapes are formed mainly by gravity driven forces, since in humid areas soils are usually covered by thick vegetation and the role of water erosion is reduced. On the other hand, when the slope surface is affected by rainfall driven erosion, the formation of concave equilibrium slopes is expected (Schumm 1956), and, at the point of equilibrium, parallel or lateral slope retreat occurs where the equilibrium concave shape is maintained over time (Nash 1980, Pelletier and Rasmussen 2009, Schumm 1956). The equilibrium shape is characterized by a uniform rate of erosion everywhere on the slope profile (Hack 1960, Montgomery 2001). Twidale and Milnes (1983) proposed that parallel retreat in

concave slopes occurs sequentially in a series of soil movement stages starting from the slope toe and evolving upward until the same slope concavity shape is formed again.

This geomorphology concept of a steady-state condition is not exempt from criticisms. Bracken and Wainwright (2008) reviewed several methods employed to measure geomorphological equilibrium, and conclude that all are flawed. They also argued that the lack of a precise definition of the spatial and temporal scales in geomorphology evolution makes many of the ideas related to steady forms untestable. Phillips (2011) states that even though steady-state conditions are sometime observed in nature, it is not a realistic representation of how geomorphologic evolution works in general. He also suggests that the idea of concave profiles as the equilibrium shape in fluvial systems is not well supported by empirical or theoretical evidence, casting doubt upon the Yang and Song (1979) concept of continuous system adjustments to reach equilibrium. Nevertheless, the fact that geomorphological studies cannot explain or include each and every single variable involved in this very complex behavior does not imply that they explain nothing. As Phillips (2011) recognized, steady-state assumptions have been useful in interpreting and modeling the evolutionary behavior of earth systems. Furthermore, there is a universal tendency for all natural processes--including geomorphological systems--to adjust and readjust themselves to dynamic equilibrium, i.e., through continuous adjustments between processes like erosion and soil resistance (Abrahams 1968, Ahnert 1994). Through these dynamic adjustments, steady-state landform shapes may be obtained (Abrahams 1968, Thorn and Welford 1994). Steady-state forms may indeed be dependent on the time scale and require unchanging boundary conditions (Thorn and Welford 1994), but if any external agent disturbs this steady equilibrium condition dynamic

adjustments will exist, so that the system will be naturally reshaped to re-obtain the steady-state form.

## Slope shape evolution and equilibrium profiles: A conceptual model

Considering that slope water erosion is a process involving water flow and sediment transport, the generalized principle of minimum rate of energy dissipation (Yang and Song 1979) may be used to explain the concept of steady-state equilibrium shapes and parallel retreat observed in slope evolution. Here, a simple conceptual model that attempts to explain the changes in morphology of slopes actively eroded by water is offered. For a slope in an initial un-balanced state (e.g., a planar slope), continuous changes in the slope geometry and water flow velocities are expected in order for the system to adjust and reach equilibrium. In this process, the overall erosion rate would gradually decrease until a constant rate of erosion along the profile is achieved. The evolution toward a constant rate of erosion implies a reduction in the difference in energy entering and leaving the system with respect to a previous point in time, eventually reaching a minimum. According to Hack (1960), the steady-state equilibrium of the landscape requires the existence of "opposing forces" counteracting and balancing each other, such that energy dissipation can exist. As it is explained in the following paragraph, these "opposing forces" can conceptually be the slope length and slope angle.

Erosion increases with slope length through the increase in cumulative runoff. The effectiveness of that runoff in detaching sediment or in transporting eroded sediment depends, however, not only on the amount of runoff, but also on its turbulent energy, which for a defined surface condition is determined by the slope steepness. This means that on a planar slope the detachment and transport 170

rates will increase (Fig. 6.1), as the steepness is constant but the slope length is increasing (Young and Mutchler 1969). In contrast, on a concave slope the steepness decreases with increased slope length, so the increase in detachment and transport potential due to the increased slope length may be counteracted by the reduction due to decreased steepness. This is consistent with Hack's (1960) idea of opposing effects balancing each other. This natural morphology evolution process toward a concave slope requires a combination of soil detachment along the contour and soil deposition at the slope base (Fig. 6.1) until an equilibrium state is reached, where the concave geometry will experience a constant detachment and transport rate at every point along the slope. At this stage, the slope profile will not suffer further adjustments and the slope will start eroding backwards in a parallel retreat fashion as illustrated in Fig. 6.2. If the slope deviates from this equilibrium state is reached again.

Notice that over time, the sharp edge at the top of the slope may be naturally shaped to a convex form (Fig. 6.2), yielding a more complex slope contour as seen in many natural landscapes in more humid areas. It is not clear if this convexity is formed due to high erodibility soil conditions triggered by runoff coming over the top edge that would infiltrate and saturate the soil, or due to a series of small mechanical failures given the steepness of the contour and the effects of transient flow through the top edge. Since the mechanisms behind the formation of convexity at the upper portion of the slope are not fully understood, this effect is not included in the development of the model in the subsequent sections.



Fig. 6.1. Conceptual model of slope morphology evolution by water erosion.



**Fig. 6.2.** Illustration of the parallel retreat concept. The sharp edge at the top of the slope may be naturally eroded (becoming convex) if there is runoff coming over the top edge of the slope. This effect is not included in the proposed model.

# Methodology

#### **Conceptual Development**

From the previous discussion it can be inferred that physical processes may shape slopes into steady concave profiles, and it will be assumed that concave profiles represent a topographic equilibrium condition that is more likely to remain steady over time than for planar profiles. Here, our efforts are not concentrated on modeling the time-scale geomorphological evolution of slopes, but rather on describing the equilibrium shape that will provide a parallel retreat. In this vein, once the concave equilibrium shape has been developed the main mechanism inducing parallel retreat is erosion, and therefore, soil deposition will not be included in this model.

As described above, the steady-state concave shape is characterized by a uniform water erosion rate. However, there are many potential concave profiles that can satisfy this condition, and the steady-state shape may not be unique. This may be explained using the widely recognized Revised Universal Soil Loss Equation RUSLE2 (USDA-ARS 2008):

$$A = R \cdot K \cdot LS \cdot C \cdot P \tag{6.1}$$

where the predicted soil loss *A* (units of Mg/ha/y) is directly proportional to: the rainfall erosivity *R* (units of  $MJ \cdot mm/ha \cdot h \cdot y$ ) quantifying the rainfall's erosive potential; the soil erodibility *K* (units of  $Mg \cdot ha \cdot h/ha \cdot MJ \cdot mm$ ) defining the soil's susceptibility to that erosivity; the topographic factor *LS* (dimensionless) representing slope length and steepness effects; the surface cover factor *C* (dimensionless); and the conservation practices factor *P* (dimensionless). To observe solely the influence of the landscape topography or slope shape on the erosion rate, let us assume that all the variables remain constant across the landscape except the topographic factor

LS, which corresponds to the combined action of the slope length factor L and the slope steepness factor S. The length factor is not a function of the length along the profile, but of the slope horizontal length  $\lambda$ . The steepness factor is a function of the angle of the slope with respect to the horizontal, and it varies in form depending on the relative values of steepness and horizontal slope length. The longer the horizontal slope length, the higher the observed erosion rate. Similarly, the steeper the slope, the higher the erosion rate. Since in a planar slope the steepness is constant, the erosion rate will uniformly increase downslope due to the increasing slope length. Fig. 6.3 illustrates the erosion rate experienced by four different points on a planar slope profile having bare soil surface conditions, assuming that all other inputs are equal to one in order to display the relative position differences. At the very top of the slope erosion is mainly controlled by raindrop detachment (sheet erosion), and the erosion rate will be proportional to the slope angle. At this point, little erosion is observed. As the runoff moves downslope it starts accumulating sufficient energy to detach soil particles, with that energy directly proportional to the runoff length and the slope angle. Thus, at point 1 on the profile, the experienced erosion rate  $A_{p1}$  will be equal to the product of the length factor  $L_{p1}$  (function of the horizontal slope length  $\lambda_{p1}$ ) and the steepness factor S (function of the slope angle  $\alpha$ ). Further down the slope runoff gains more energy due to greater cumulative flow rate, and therefore the erosion rate at point 2 will be greater than at point 1 ( $A_{p2} > A_{p1}$ ). Analogously, the erosion rate at point 3 will be greater than at point 2, and the result is an uninterrupted increasing erosion rate downslope ( $A_{p3} > A_{p2} > A_{p1} > A_{p0}$ ).



Fig. 6.3. Erosion rates experienced by a planar profile as modeled by RUSLE2 perspective.

On the other hand, what if any increment in slope length is compensated with a decrease in the slope angle in such a way that the erosion remains equal along the profile? This would be the case if the value of  $L_i S_i$  is constant for all segments. The resulting shape of this compensation process would be concave with a constant rate of erosion along the profile. Notice that the whole erosion phenomenon starts at the uppermost point of the slope, where only sheet erosion exists, and sheet erosion is a function only of the slope angle. Different initial slope angles will yield different initial rates of erosion, resulting in different subsequent downhill combinations of lengths and angles in order to achieve the same rate of erosion everywhere. Thus, many concave shapes may satisfy the condition of constant rate of erosion, each with a different level of constant erosion controlled by the assumed angle of the top segment. Fig. 6.4a shows a profile with an initial angle  $\alpha_{0h}$  at the top of the slope and an initial erosion rate  $A_{0h} = S(\alpha_{0h})$ . To maintain this erosion production constant along the profile  $(A_{0h} = A_{1h} = A_{2h} = A_{3h})$ , the angle of the slope  $\alpha_i$  must continuously decrease. Notice the erosion at  $\lambda_i$  needs to be a function of  $L_i$  and  $S_i$ , which in turn are functions of  $\lambda_i$ and  $\alpha_i$  respectively ( $f[L_i(\lambda_i) \cdot S_i(\alpha_i)]$ ). Fig. 6.4b shows a profile initially inclined at  $\alpha_{0m}$  at the top. Since  $\alpha_{0m} < \alpha_{0h}$  the initial erosion rate for this case will be  $A_{0m} < A_{0h}$ , and the adjustment process to achieve a constant erosion rate will yield a profile somewhat less steep. Similarly, the case illustrated in Fig. 6.4c corresponds to an even gentler concave profile, due to an even smaller initial rate of erosion  $A_{0/}$ . Notice that for the same soil strength an increase of mechanical stability from case a) to case c) is perceived due to the decrease in overall inclination.



**Fig. 6.4.** Illustration of a potential family of concave slopes with constant rate of erosion: a) high rate, b) intermediate rate, and c) low rate.

#### **RUSLE2** approximation to steady-state concave slopes

## Normalized rate of erosion and equations for LS

Following this conceptual idea, different levels of constant erosion will determine the concavity of the slope profile, starting from the initial angle at the top of the slope. To approach this problem, let us first modify Eq. (6.1) and define a normalized erosion rate  $A^n$  such that the amount of erosion is identical to the topographic factor *LS*:

$$A^{n} = \frac{A}{R \cdot K \cdot C \cdot P} = LS \tag{6.2}$$

Notice that in reality  $R \cdot K \cdot C \cdot P$  is not truly constant over time. For example, the cover factor *C* right after construction of a hillslope will be one (highly erodible and disturbed conditions), but as the soil starts consolidating and vegetation starts growing on the slope surface, the factor *C* starts decreasing. However, it is difficult to predict how *C* will change over time, and therefore, *C* is assumed to be constant. Also, notice that  $A^n$  is a dimensionless quantity of erosion and *LS* is a function of slope length and slope angle  $[LS(\lambda, \alpha)]$ . For horizontal slope lengths  $\lambda \leq 0.92$  m (3 ft) *LS* is defined as (USDA-ARS 2008):

$$LS = \left(\frac{4.6}{22.1}\right)^{m} \cdot \left(3 \cdot \sin(\alpha)^{0.8} + 0.56\right)$$
(6.3)

where  $\alpha$  is the slope angle measured in degrees. The factor *m* in Eq. (6.3) is called the slope length exponent, and can be calculated as follows (Foster et al. 1977):

$$m = \frac{\beta}{1+\beta} \tag{6.4}$$

where  $\beta$  is the ratio of rill to inter-rill erosion defined as (McCool et al. 1989):

$$\beta = \frac{\sin(\alpha)/0.0896}{3 \cdot \sin(\alpha)^{0.8} + 0.56}$$
(6.5)

178

For  $\lambda \ge 4.6$  m (15 ft) *LS* is computed as:

$$LS = \begin{cases} \left(\frac{\lambda}{22.1}\right)^{m} \cdot \left(10.8 \cdot \sin(\alpha) + 0.3\right), & \alpha < 5.1^{\circ} \\ \left(\frac{\lambda}{22.1}\right)^{m} \cdot \left(16.8 \cdot \sin(\alpha) - 0.5\right), & \alpha \ge 5.1^{\circ} \end{cases}$$

$$LS = \begin{cases} \left(\frac{\lambda}{72.6}\right)^{m} \cdot \left(10.8 \cdot \sin(\alpha) + 0.3\right), & \alpha < 5.1^{\circ} \\ \left(\frac{\lambda}{72.6}\right)^{m} \cdot \left(16.8 \cdot \sin(\alpha) - 0.5\right), & \alpha \ge 5.1^{\circ} \end{cases}$$
(6.6a)
$$(6.6b)$$

For slope lengths of 0.92- 4.6 m (3 - 15 ft), the algorithm employs a linear interpolation between the natural logarithm of Eq. (6.3) and the natural logarithm of Eq. (6.6) with  $\lambda = 4.6$  m (15 ft) to calculate intermediate values of *LS* (USDA-ARS 2008). Notice that Eq. (6.3) is not a function of the slope length  $\lambda$ , and the same erosion is computed at any point of the profile that is shorter than 0.92 m (3 ft). This is because only sheet erosion is observed for short flow lengths (Meyer and Harmon 1989).

## Discretization of the concave slope profile with linear segments of equal erosion rate

By assuming that a concave slope can be obtained from the assembly of a discrete number of small linear segments (Fig. 6.5) where the same erosion rate holds on every segment of the entire length, a selected initial erosion rate  $A_1^n$  can be employed to determine the angle of the first linear segment  $\alpha_1$  at the top of the slope. From this point forward, each subsequent segment must produce on average the same amount of erosion  $A_1^n$  on the segment's length:

$$A_i^n = \frac{p_i}{\Delta x_i} = A_1^n \tag{6.7}$$

179



Fig. 6.5. Discretization of the slope profile in small linear segments.

Here  $A_i^n$  and  $p_i$  are the erosion rate and the erosion production at the *i*-th segment respectively, and  $\Delta x_i$  is the horizontal length of the *i*-th segment. The erosion production is defined as the product of the total erosion and the total horizontal distance at a particular point in the profile. Therefore,  $p_i$  becomes the subtraction of the erosion produced at the initial and final points of the *i*-th segment (Fig. 6.5). Making use of Eq. (6.2),  $p_i$  can be expressed as:

$$p_i = LS(\lambda_i, \alpha_i) \cdot \lambda_i - LS(\lambda_{i-1}, \alpha_{i-1}) \cdot \lambda_{i-1}$$
(6.8)

Subsequently, the erosion rate at the *i*-th segment becomes:

$$A_i^n = \frac{LS(\lambda_i, \alpha_i) \cdot \lambda_i - LS(\lambda_{i-1}, \alpha_{i-1}) \cdot \lambda_{i-1}}{\Delta x_i} = A_1^n$$
(6.9)

Having a known initial erosion  $A_i^n$  and making the length of the first slope segment  $\Delta x_i = \lambda_i = 0.92$ m (since below this point erosion calculations are invariant),  $\alpha_i$  can be determined, and with it the coordinates of the end point of the first segment using simple trigonometry, with the initial point as the origin of the coordinate system (Fig. 6.5). To obtain the coordinates of the second segment, Eq. (6.9) can be employed. Here,  $\lambda_2 = \lambda_1 + \Delta x$  ( $\Delta x$  is the selected horizontal length of each subsequent linear segment) and  $A_2^n = A_1^n$ , so  $\alpha_2$  is the only remaining variable. Notice that  $\alpha_{i-1}$  is always known, and therefore  $\alpha_i$  is the only variable to be determined. Repeating this process, the angle of all the subsequent segments can be determined along with the associated (x, y) coordinates of each point, with  $\lambda$  corresponding to the x coordinate. Notice that in the selected reference frame the x-axis is positive to the right and y-axis is positive downward. Given the recursive nature of the proposed methodology and the non-linearity of the RUSLE2 *LS* equations, a numerical solution for  $\alpha_i$  was created. Since at the top of the slope ( $\lambda \le 0.92$  m) the erosion is dominated by raindrop detachment, the prediction of the slope contour results on a linear segment. However, here the overland flow path length is not important, and the linear shape of the first segment can be arbitrarily modified such that it preserves the overall slope shape predicted for  $\lambda > 0.92$  m (3 ft). A third-order polynomial of the form

$$V(\lambda) = a \cdot \lambda^3 + b \cdot \lambda^2 + c \cdot \lambda + d \tag{6.10}$$

was chosen to construct a cubic spline for  $\lambda \in [0, 0.92 \text{ m}]$ . The boundary conditions are:

$$V(\lambda) = 0, \text{ at } \lambda = 0 \tag{6.11}$$

$$V(\lambda) = 0.92 \cdot \tan(\alpha_1), \text{ at } \lambda = \Delta x_1 = 0.92 \text{ m} (3 \text{ ft})$$
 (6.12)

$$\frac{dV(\lambda)}{d\lambda} = \tan(\alpha_2), \text{ at } \lambda = \Delta x_1 = 0.92 \text{ m (3 ft)}$$
(6.13)

To satisfy a continuity condition,

$$\frac{d^2 V(\lambda)}{d\lambda^2} = 0$$
, at  $\lambda = \Delta x_1 = 0.92$  m (3 ft) (6.14)

After imposing these boundary conditions on Eq. (6.10), the problem becomes a linear system of equations that will provide the resulting values of *a*, *b*, *c*, and *d*, defining the polynomial function. Different functions are obtained for different initial erosion rates  $A_1^n$ .

#### Sustainable slopes: mechanical and erosion stability

The approach described above produces a family of concave profiles that satisfy the water erosion equilibrium condition. However, a truly stable slope must be stable in terms of both mechanical and water erosion processes. Employing the critical slope contour approximation proposed by Jeldes et al. (2013) those steady concave shapes that also have long-term stability for a given

mechanical soil strength defined by the Mohr-Coulomb parameters of friction angle  $\phi$ , cohesion c, and unit weight  $\gamma$  were investigated, using the relationship

$$x = \begin{cases} 0 & ,-h_{cr} \leq y \leq 0 \\ A \left[ \sigma_{y} \left( B - 1 \right) \left( cosec\phi - 1 \right) + H B \left( cosec\phi + 1 \right) \right], & y > 0 \end{cases}$$
(6.15)

where:

$$A = \frac{\cos\phi}{2\gamma(1-\sin\phi)} \tag{6.16}$$

$$B = ln \left[ \frac{\sigma_y}{H} \left( \frac{1 - \sin \phi}{1 + \sin \phi} \right) + 1 \right] = ln \left[ \frac{\sigma_y}{H} K_a + 1 \right]$$
(6.17)

$$\sigma_{y} = \gamma y \tag{6.18}$$

$$H = c \cot \phi \tag{6.19}$$

$$h_{cr} = \frac{2c\cos\phi}{\gamma(1-\sin\phi)} \tag{6.20}$$

In these relationships *H* is the tensile strength of the soil,  $K_a = (1 - \sin \phi)/(1 + \sin \phi)$  is the Rankine active coefficient of earth pressure,  $\sigma_y$  is the geo-static vertical stress, and  $h_{cr}$  is the height of the tension zone. The equation describes a slope contour in the quadrant with *x*-axis positive to the right and *y*-axis positive downward, with  $h_{cr}$  lying above the *x*-axis (Jeldes et al. 2013). In many instances the sharp cusp at the origin of the coordinate system (Fig. 6.6) would not exist, but it is maintained to be consistent with the theoretical mathematical solution.

Slopes in erosion equilibrium (for given values of  $A_1^n$ ) that are less steep than the critical slope Eqs. (6.15) - (6.20) are mechanically stable, while steeper ones are unstable. This idea is further



**Fig. 6.6.** Illustration of steady-state or sustainable slope (profile 1) and a non-equilibrium slope with cyclic morphology changes (profile 2).

illustrated in Fig. 6.6, where two different slopes in erosion equilibrium are shown for a soil defined by a unique combination of  $\phi$ , *c*, and  $\gamma$ . Here, slope profile 1 satisfies both the erosion and mechanical equilibrium requirements, and therefore becomes a sustainable slope with a steady-state shape over time. Profile 2, on the other hand, only satisfies erosion equilibrium and mechanical failures are expected, which would alter the erosion equilibrium shape. However, natural adjustments may bring the slope to a new erosion equilibrium state, which if not mechanically stable, will repeat the failure and natural re-shaping process, creating cyclic changes in morphology over time.

## **Results and Discussion**

#### Concave slopes with constant rate of erosion

A family of concave slopes with constant rates of erosion ( $A_1^n = 0.506$  to 1.068) along the profile is shown in Fig. 6.7. These concave slopes are the result of a numerical solution of Eqs. (6.9) – (6.14), with  $\Delta x = 0.03$  m (0.1 ft). The chosen discretization length was selected to be small enough such that the ratio  $\alpha_i / \alpha_{i-1}$  remained within the 0.9 – 0.99 range for the majority of slope length. The computations stopped once  $\alpha_i$  fell below 3°. In this family of profiles there is a clear increase in the overall slope steepness as  $A_1^n$  increases. The overall steepness, defined as the slope of the straight line that connects the top and the bottom of the slope, increases at a much faster rate for  $A_1^n \ge 1$  (Fig. 6.8). This means that the predicted equilibrium concave profiles are more sensitive to values of  $LS \ge 1$ . Although the entire family of profiles is in erosion equilibrium (equal erosion rate along the profile), these concave slopes may or may not be mechanically stable, so their stability must be analyzed based on the critical concave slope concept.



Fig. 6.7. Family of concave slopes with constant rate of erosion (equilibrium erosion profiles).



**Fig. 6.8.** The overall steepness (defined as the slope of the straight line that connects the top and the bottom of the profile) increases at a much faster rate for  $A_1^n \ge 1$ , as denoted by dashed line.

Notice that in the analyses the constant erosion is implied for the whole extent of the concave slope. In reality, when the concave slope encounters a relatively horizontal surface (e.g., at the bottom of the slope), part of the detached soil is deposited. For erosion to occur along a slope, the transport capacity of the flow  $(T_c)$  must be greater than its sediment load (q). Transport capacity, which is a measure of how much sediment the runoff can carry, is a function of the kinetic energy of the flow, which decreases as the slope becomes less steep. When the flow reaches a relatively flat area,  $T_c$  becomes less than q, and the flow is not capable of detaching new soil nor transporting the eroded soil, leading to deposition. Even though the employed equations do not capture this effect, the eroding portion of the profile will not be affected by deposition at the bottom end of the slope, so, the predicted profile with constant erosion rate would remain as predicted.

#### Sustainable slopes

Fig. 6.9 illustrates the sustainable slope concept for a medium soft clay soil with  $\phi = 20^{\circ}$  and  $c/\gamma$ = 1 m. Using Eqs. (6.15) - (6.20), the critical concave slope (FS = 1) for this specific soil was determined and plotted together with the erosion equilibrium shapes. For the purpose of this comparison, the vertical tension zone of height  $|h_{cr}|$  is not displayed. The results suggest that, for this specific soil erosion equilibrium slopes with  $A_1^n \le 1.05$  would be mechanically stable, and therefore would be sustainable with steady shapes over time. On the other hand, equilibrium concave shapes with  $A_1^n > 1.05$  (shown with dashed lines in Fig. 6.9) are likely to experience mechanical failures, and changes in morphology may be observed. A second case is illustrated in Fig. 6.10 for a sandy soil with  $\phi = 30^{\circ}$  and  $c/\gamma = 0.5$  m. This time erosion equilibrium shapes



**Fig. 6.9.** Erosion equilibrium profiles and a critical concave contour for  $\phi = 20^{\circ}$  and  $c / \gamma = 1$  m. Concave slopes with  $A_1^n \le 1.05$  are considered sustainable.



Fig. 6.10. Erosion equilibrium profiles and a critical concave contour for  $\phi = 30^{\circ}$  and  $c/\gamma = 0.5$  m. Concave slopes with  $A_1^n \le 1.039$  are considered sustainable.

with  $A_{l}^{n} \leq 1.039$  will be sustainable, while concave shapes with  $A_{l}^{n} > 1.039$  will not. Note that  $A_{l}^{n} = 1.05$  and 1.039 are approximated limiting values  $(A_{L}^{n})$  that separate sustainable and nonsustainable concave slopes for the given soil mechanical properties. Similar chart results are shown in Fig. 6.11 for a wide range of mechanical soil properties ( $\phi = 20^{\circ}, 25^{\circ}, 30^{\circ}, 35^{\circ}$  and  $40^{\circ}$ ;  $c/\gamma =$ 0.5, 0.75, 1, 1.5, 2 and 3 m). Generalized values of  $A_{L}^{n}$  are shown in Fig. 6.12 as a function  $\phi$  and  $c/\gamma$ . From this chart, the following approximation for  $A_{L}^{n}$  is proposed: for the range of  $A_{l}^{n}$ ,  $\phi$ , and  $c/\gamma$  investigated here, concave erosion equilibrium slopes will be sustainable if:

$$A_{\rm l}^n \le 0.02 \frac{\gamma}{c} \cdot \ln[\tan(\phi)] + 0.01 \frac{c}{\gamma} + 1.06 \tag{6.21}$$

In practice, slopes are rarely designed for mechanical stability with FS = 1, and some additional margin of safety is desired. Depending on the conditions and the impact of potential failure, slopes are designed for a minimum FS always greater than 1. The concept of sustainable slopes is extended in Figs. 6.13, 6.14, and 6.15 by introducing critical concave contours with FS > 1 (Jeldes et al. 2014). Fig. 6.13 shows erosion equilibrium slopes and a concave contours satisfying FS = 1.25, for  $\phi = 20^{\circ}$ ,  $25^{\circ}$ ,  $30^{\circ}$ ,  $35^{\circ}$  and  $40^{\circ}$ ; and  $c/\gamma = 0.5$ , 0.75, 1, 1.5, 2, and 3 m. Erosion equilibrium slopes less steep than the critical contour for a given set of  $\phi$  and  $c/\gamma$  will not only be sustainable, but also their shear strength along the most critical failure surface would be at least 25% greater than that required to maintain mechanical equilibrium in the long-term. Similar charts for FS = 1.5 and 1.75 are shown in Figs. 6.14 and 6.15, respectively.



Fig. 6.11. Erosion equilibrium profiles and a critical concave contour for  $\phi = 20^{\circ}$ , 25°, 30°, 35°, 40° and  $c/\gamma = 0.5, 0.75, 1, 1.5, 2, 3$  m.



**Fig. 6.12.** Limiting values of erosion  $(A_L^n)$  defining sustainability.


Fig. 6.13. Erosion equilibrium profiles and a concave contours with FS = 1.25, for  $\phi = 20^{\circ}$ ,  $25^{\circ}$ ,  $30^{\circ}$ ,  $35^{\circ}$ ,  $40^{\circ}$  and  $c/\gamma = 0.5$ , 0.75, 1, 1.5, 2, 3 m.



Fig. 6.14. Erosion equilibrium profiles and a concave contours with FS = 1.5, for  $\phi = 20^{\circ}$ ,  $25^{\circ}$ ,  $30^{\circ}$ ,  $35^{\circ}$ ,  $40^{\circ}$  and  $c / \gamma = 0.5$ , 0.75, 1, 1.5, 2, 3 m.



Fig. 6.15. Erosion equilibrium profiles and a concave contours with FS = 1.75, for  $\phi = 20^{\circ}$ , 25°, 30°, 35°, 40° and  $c/\gamma = 0.5$ , 0.75, 1, 1.5, 2, 3 m.

In this article, the argument that slopes in nature are seldom planar, and concave contours are naturally formed as a result of evolutionary processes that seek erosion equilibrium has been presented. Therefore, why must constructed slopes be planar? The results presented here suggest that concave slopes can be constructed to achieve both minimal steady-state erosion equilibrium as well as mechanical stability. The desired slope profile could be selected by balancing erosion rates and desired mechanical stability for specific site conditions and regulation requirements such as discharge into sensitive waterbodies, highly erodible soils, and space limitations. Then, Figs. 6.13, 6.14, and 6.15 can be used as design tools for the construction of slopes reflecting more natural landforms, while minimizing sediment delivery during initial slope shape adjustments. If this is done, erosion and sediment control activities to handle the eroded sediment can be less rigorous. The increasing precision of GPS-based construction equipment suggests that such complex landforms can be now achieved with high accuracy, while deviations in the constructed slope profile due to typical construction inaccuracies do not compromise the stability of concave slopes (Jeldes et al. 2014). With the current construction technology, sustainable concave slopes can become more than a theoretical exercise.

## Limitations of this work

The methodology presented here has limitations that must be acknowledged. First, the topographic equations of RUSLE2 provide the fundamental prediction of erosion equilibrium slopes. RUSLE2, however, is defined for slope angles equal to or less than 45°, while in this analysis initial erosion rates sometimes requiring an initial angle (angle of the first segment  $\Delta x_1 = \lambda_1 = 0.92$  m) exceeding this definition were used. For example, the erosion equilibrium slope defined for  $A_1^n = 0.927$ 

requires an initial angle of 50°, while the one defined for  $A_1^n = 1.068$  requires an initial angle of 80°. In this sense, the limits of the empirical basis on which RUSLE2 was developed have been stretched. However, as the flow path crosses the  $\lambda = 0.92$  m limit, the angles rapidly decrease and fall below the 45° definition (Fig. 6.16). Notice that the angles of the concave slope with the highest erosion rate ( $A_1^n = 1.068$ ) reach 45° at x = 1.3 m, and therefore the assumption that RUSLE2 holds for  $\alpha > 45^\circ$  only affects a small portion of the slope. Future empirical work on slopes steeper than 45° with short flow paths is required to validate the solution for such steep initial segments.

In addition, natural slopes often include a convex portion at the top, which may be the result of runoff coming over the top edge of the slope, increasing the soil detachment on a section that otherwise is only affected by sheet erosion. This effect was not included in this model, and future work is needed to understand and model this complex mechanism.

Finally, this approach does not consider the effects of transient subsurface flow, nor the short-term stability of sustainable slopes. However, once the desired sustainable concave shape has been selected, the effects of transient ground water flow can be investigated using commercial slope stability software. Similarly, the short-term stability could be evaluated once the site specific undrained shear strength has been determined. The use of sustainable concave slopes in practice will not only provide more aesthetically pleasing results, but also contribute to the growing need for environmental-friendly techniques for landform construction and sustainable land management for agricultural development.

![](_page_221_Figure_0.jpeg)

Fig. 6.16. Computed slope angles vs horizontal slope length.

## Conclusions

Natural landforms adjust their shape over time in order to seek erosion and sediment transport equilibrium. The resulting equilibrium shape is usually concave with a uniform erosion rate along the profile. In theory, this equilibrium shape would not suffer further adjustments, so the total erosion production is minimized. In this article, a conceptual model was offered attempting to explain how concave equilibrium shapes slopes are formed in nature. Also from a conceptual perspective, it was shown that the equilibrium shape may not be unique, and a family of potential slope shapes can exist according to the level of constant erosion the slope experiences.

Using the fundamentals of the widely recognized RUSLE2 model, a set of equations to describe concave slopes in erosion equilibrium at different levels of constant erosion was developed. The resulting concave profiles were compared with critical concave slopes (FS = 1) and concave slopes having FS = 1.25, 1.5, and 1.75, identifying those shapes that satisfy both erosion equilibrium and the desired degree of mechanical stability. A mathematical expression defining the limiting erosion at which equilibrium erosion shapes become sustainable was provided as a function of the Mohr-Coulomb parameters, and the limitations of this work were discussed. The obtained results suggest that concave slopes can be constructed to achieve both minimal steady-state erosion equilibrium as well as mechanical stability, and that the slope profile can be chosen such that erosion rates and desired mechanical stability are balanced to satisfy site conditions and regulatory requirements. Sustainable concave slopes not only offer superior erosion control, but also a more natural appearing landform, being an attractive alternative to the traditional planar slopes that are typically constructed today.

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# Chapter 7. The Piling Framed Concrete Retaining Wall: Design Pressures and Stability Evaluation

This chapter was submitted as an original paper to the ASCE Practice Periodical on Structural Design and Construction, and it is currently undergoing the review process. The co-authors of this work are Dr. Eric Drumm, Dr. Richard Bennett, and Nikola Zisi. This article would be cited as:

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## Abstract

The Piling Framed Retaining Wall (PFRW) is an innovative earth retention system applicable for soils underlain by rock, which is ideal for applications where only limited right-of-way is available or adjacent structures limit the use of tie-back anchors. Two PFRW's were successfully built along the I-40/I-75 corridor in Knoxville, TN, with significant cost savings over traditional retaining wall designs. Although the walls were designed using conventional earth pressure theories, the soil pressures and forces acting on the wall face are not fully understood, and a rational design method has not been fully developed. Traditional theories of lateral earth pressure assume rigid translations or rotations as the fundamental deformation mode, when in reality more complex mechanisms of deformation and earth pressure distributions may exist. A series of Finite Elements analyses were used to evaluate the soils stresses on the face of the wall for various configurations of wall geometry, backfill slopes and soil properties. From the results, simplified design equations were developed to predict the earth pressures on the wall face and the overturning moments for stability analyses. The proposed design equations were validated against traditional expressions, and compared with earth pressures measured on the wall over a three year period. The measured stresses and the numerical results suggest that the typical earth pressure distribution of the PFRW is neither linear nor monotonically increasing, and the proposed design equations yield conservative results for practical combinations of geometry and soil properties. The proposed design methods offer a reliable way to predict wall pressures and overturning moments, and eliminate the need to conduct extensive numerical analyses for each wall to be constructed.

## Introduction

The Piling Framed Retaining Wall (PFRW) is an innovative earth retaining system, which is ideal for applications where the overburden soil is underlain by competent rock and where only limited right-of-way (ROW) is available or adjacent structures and underground utilities limit the use of tie-back anchor systems. Traditional retaining structures, such as concrete cantilever walls, usually require large excavations behind the wall, and thus may not be suitable when ROW is limited. Similarly, the need for underground easements behind the wall to avoid buried utilities or adjacent building foundations may constrain the use of tie-back walls. To overcome these limitations, the Tennessee Department of Transportation (TDOT) developed this innovative PFRW concept (Pate and Haddad 2007), and successfully implemented it in two walls along the I-40/I-75 corridor in Knoxville, TN (Fig. 7.1). This wall concept was found to be more economical than traditional retaining wall systems under the given geologic conditions. Despite its successful application, a rational design method for this wall has not been fully developed, and the soil pressures and forces acting on the structural elements are not fully understood. While traditional theories of lateral earth pressure assume that the wall behaves like a rigid body undergoing translations and/or rotations, in reality more complex mechanisms of deformation may exist, which in turn create earth pressure distributions different than those predicted by classical theoretical expressions.

In this investigation, the Finite Element Method (FEM) was used to create numerical models to investigate the soil pressures on the PFRW for various wall face inclinations, wall heights, and backfill slopes. The results were compared with theoretical expressions available in the literature for various configurations of wall geometry and soil properties. Simplified equations were developed from the FEM analyses to facilitate the design and stability calculations of PFRWs without the need to create geometric-specific finite element models. In addition, an approximate design approach based on the well-known Coulomb earth pressure theory is demonstrated. Finally, the earth pressures predicted by the design equations and the approximate design approach were compared with soil stresses measured on one of the walls built in Knoxville, TN, referred to here as the SmartFix wall (Fig. 7.1).

## Background

#### The PFRW concept and the construction sequence

The PFRW configuration consists of vertical and battered driven H-piles supported on rock to create the structural frame. The frames are spaced 3 m (10 ft) apart and are connected near the top by a horizontal waler (Fig. 7.2). After the piles are driven, vertical tie-down anchors are installed through the waler to provide stability against overturning (Fig. 7.2). The face of the wall is aligned with the battered piles and consists of timber lagging that is placed sequentially from top to bottom as the soil is being excavated. The face is finished with cast-in-place reinforced concrete placed over the lagging and acting composite with the battered piles.

![](_page_231_Picture_0.jpeg)

**Fig. 7.1.** The PFRWs built along the I-40/I-75 corridor (Knoxville, TN). a) The SmartFix wall and b) The West Hills wall.

![](_page_231_Figure_2.jpeg)

Fig. 7.2. The piling frame system forming the basis of the PFRW concept.

A large cast-in-place concrete cap or parapet provides mass for additional stability against overturning. Pate and Haddad (2007) outline the details of construction of the first PFRW built in Knoxville, Tennessee in 2005. The construction sequence can be summarized as follows: a) Initial dry excavation and driving of battered piles (Fig. 7.3a), where the ground surface is excavated to a depth slightly below the connection of the battered pile, the vertical pile and waler, and then the battered piles are driven to refusal. b) Installation of walers and driving of vertical piles (Fig. 7.3b), where the tops of the battered piles are coped to fit inside the waler which is installed with the strong axis in the horizontal direction and welded to the top of the battered piles; the vertical piles are driven to refusal using the waler as a template. c) Installation of tie-down anchors (Fig. 7.3c), where holes are cut in the web of the waler, through which the tie-down anchors are drilled and anchored to rock and a preload applied to resist overturning. d) Excavation of soil and installation of timber lagging (Fig. 7.3d), where the soil in front of the wall is excavated and timber lagging installed between the battered piles as the excavation progresses (Fig. 7.4). e) Construction of structural concrete facing (Fig. 7.3e), where shear studs are installed on the top flange of the battered pile, reinforcing steel placed, and the structural concrete facing poured; a decorative ashlar stone finish can be obtained with a reusable form liner (Fig. 7.4). f) Pouring of concrete cap or parapet wall (Fig. 7.3f) of 2.6 m x 0.9 m (8 ft x 3 ft) (Fig. 7.4), which provides a finished appearance and is intended to introduce additional overturning resistance. The walls are then backfilled and graded.

#### Theoretical expressions for active, passive, and at-rest earth pressure coefficients

Traditionally, lateral earth pressures for rigid walls have been estimated using simplifications or assumptions based on elasticity and plasticity theory. The most rigorous expression for the at-rest

![](_page_233_Figure_0.jpeg)

Fig. 7.3. Construction sequence for PFRW.

![](_page_234_Picture_0.jpeg)

**Fig. 7.4.** Illustration of the construction stages for the timber lagging, concrete face and concrete cap. The black vertical material is a geosynthetic filter layer.

lateral earth pressure coefficient is based on the theory of elasticity [ $K_0 = \upsilon/(1-\upsilon)$ ], where  $\upsilon$  is the Poisson's ratio, while empirical correlations with strength properties such as the friction angle have also been proposed (Jaky 1944):

$$K_0 = 1 - \sin\phi \tag{7.1}$$

Methods that assume a plastic soil behavior employ the active and passive states as limiting conditions. These limits are determined via a) static approaches, such as the *Rankine state of stresses* (Rankine 1856) [Eq. (7.2)], where the active lateral earth pressure coefficient  $K_a$  is determined based on the equilibrium equations and the Mohr-Coulomb yield criterion:

$$K_a = \frac{1 - \sin\phi}{1 + \sin\phi} \tag{7.2}$$

or b) kinematics methods, such as the *Coulomb's wedge method* (Coulomb 1776) [Eq. (7.3)], where wedge failures are assumed and the resultant earth pressure is calculated by means of force equilibrium or by means of external rate of work and internal rate of energy dissipation:

$$K_{a} = \frac{\sin^{2}(\beta + \phi)}{\sin^{2}(\beta)\sin(\beta - \delta)\left[1 + \sqrt{\frac{\sin(\phi + \delta)\sin(\phi - \alpha)}{\sin(\beta - \delta)\sin(\alpha + \beta)}}\right]^{2}}$$
(7.3)

Numerical solutions for  $K_a$  from the static *Slip-line Network* (Sokolovskii 1960) and the kinematics *Upper Bound in Limit Analysis* (Chen and Rosenfarb 1973) and the *Non-linear wedge analysis* (Caquot and Kérisel 1948) are also available in the literature. All these aforementioned methods, however, have been developed for rigid walls under translation, rotation about the top, or rotation about the toe as the fundamental deformation mode. In reality, walls may exhibit more complex mechanisms of deformation, dictated by the structural rigidity and the presence or absence of anchorage, which in turn creates more complex failure mechanisms than those assumed 211

by theoretical expressions. The distribution, magnitude and direction of the earth pressures acting on the face of the PFRW remain unknown.

## Methods

#### **Numerical Analyses**

When designing a PFRW the number of H pile sections available is relatively limited. For modeling purposes sizes that were successfully used in the construction of the Tennessee walls were employed. Then, the two important variables become a) the earth pressures on the wall face that will dictate the bending stresses on the wall facing, which in turn will determine the thickness and reinforcement of the concrete, and b) the overturning moments that will allow designers to determine how much anchorage force is needed to provide overturning stability.

To understand how the magnitude and distribution of the earth pressures behave on the face of a PFRW and to determine the magnitude of the overturning moments, a series of two-dimensional plane strain models using the Finite Element Method (FEM) were created with the software *Phase2* (Rocscience Inc. 2011) for different backfill slopes ( $\alpha$ ), wall face angles ( $\beta$ ), and effective wall heights (*H*) (Fig. 7.5). Table 7.1 summarizes the wall configurations investigated; note that  $\beta$ =71.6° corresponds to a 1:3 (H:V) slope that was used on the Tennessee walls. This inclination is commonly used in driven battered piles for bridge foundations and therefore easily accommodated with typical pile driving equipment.

![](_page_237_Figure_0.jpeg)

Fig. 7.5. Definition of "effective" wall height, wall inclination and backfill slope.

	H = 5  m		H = 7.62  m		H = 10  m		<i>H</i> = 15 m	
	(16 ft)		(25 ft)		(33 ft)		(49 ft)	
Wall Inclination	$\alpha = 0^{\circ}$	$\alpha = 20^{\circ}$						
$\beta = 50^{\circ}$	1	1	1	1	1	1	1	1
$\beta = 60^{\circ}$	1	1	1	1	1	1	<b>I</b>	1
$\beta = 71.6^{\circ}$	<b>√</b>	$\checkmark$	1	1	<b>\</b>	<b>√</b>	<b>√</b>	1
$\beta = 80^{\circ}$		$\checkmark$			✓			\[         \]     \[         \[         \]     \[         \[         \]     \[         \[         \]     \[         \[         \]     \[         \[         \[         \[

**Table 7.1** Wall configurations being investigated for every combination of  $\phi$  and  $c/\gamma$ 

### Soil properties for FEM modeling

To simulate the field conditions where the PFRW is appropriate, all the considered FEM models consisted of soil overlaying a very strong and stiff rock. An elastic perfectly plastic stress-strain behavior with a Mohr-Coulomb yield criterion was used to model the soil. For every combination of wall geometry parameters  $\alpha$ ,  $\beta$ , and H (Table 7.1) values of friction angle  $\phi = 20^{\circ}$ ,  $30^{\circ}$ , and 40° and cohesion over unit weight ( $c/\gamma$ ) ranging from 0.3 to 3 m (1 to 10 ft), were considered, where the range c = 5 - 40 kPa (104 – 835 psf) covers most of the cohesion values reported by Mesri and Abdelghaffar (1993), and the range  $\gamma = 10 - 23 \text{ kN/m}^3$  (64 – 146 pcf) would cover most materials from clays to coarse granular soils (NAVFAC 1986). A total of 288 FEM models were created. A different value of Young's modulus (E) was input for each value of  $\phi$  employed. An  $E = 5 \times 10^4$  kPa (1,044 ksf) was used for soils with  $\phi = 20^{\circ}$  to represent a medium stiff clay, an  $E = 1 \times 10^5$  kPa (2,089 ksf) was used for  $\phi = 30^{\circ}$  to represent a medium dense sand, and an E = $1.5 \times 10^5$  kPa (3,133 ksf) was used for  $\phi = 40^{\circ}$  to represent dense sand or medium dense gravel (Bardet 1997). A single Poisson's ratio v = 0.3 was used for all the models, since it has a small range of variation among soils (Bardet 1997). A non-associated flow rule (dilatancy angle  $\psi = 0$ ) was implemented. The effects of using an associated flow rule ( $\psi = \phi$ ) in the model were also investigated for the SmartFix case.

#### Structural and Geometric Properties for FEM modeling

Since there is only a limited number of H pile sections suitable for these conditions, the analysis was limited to those employed in the construction of the Tennessee walls (e.g. the SmartFix wall) for all the models. The battered pile was a HP 10x42 and the vertical pile was a HP 12x53. The

concrete face was 0.3 m (1 ft) thick and the concrete cap was 0.9 m (3 ft) wide and 2.4 m (8 ft) high. All the structural members were modeled on a 3 m (10 ft) basis. Standard linear elastic Timoshenko beam elements were employed to model all piles. A Young's modulus of 200 GPa (29,000 ksi) and a Poisson's ratio of 0.29 were assigned to the beams. The moment of inertia  $I_g$  with respect to the strong axis and the cross sectional area  $A_g$  of the steel H-piles were obtained from the AISC Steel Construction Manual and divided by the spacing between piles (Table 7.2). The wall elements comprised of steel H-piles and concrete (inclined wall face and cap) were treated as composite sections given that mechanical shear connectors provide a sufficient bond between the materials. The dimensions of the concrete elements were transformed to an equivalent steel section and the section properties of the composite sections calculated as if they were entirely comprised of steel (Table 7.2). Interface elements were employed to allow slip and separation between the soil and the battered composite wall. Due to a rough soil-lumber contact, a Mohr-Coulomb slip criterion with a soil-wall interface friction angle equal to the soil friction angle  $\phi$  was chosen.

To restrict the parametric study to practical combinations of soil properties and geometry that would likely be constructed, the geometric/soil combinations that resulted in excessive wall deformations, assumed here to be deflection over the length of the inclined face of the wall  $\delta/L$  > 1/240 (Council 2011), were excluded from the summary analyses. The combinations investigated are illustrated in Fig. 7.6. The triangles and squares are combinations with deflections exceeding 1/240 (the square markers are wall configurations that exceeded the more severe  $\delta/L$  > 1/120 limit). The zones drawn in Fig. 7.6 capture the potential combinations that exceeded

Cross Section	Moment of Inertia, $I_g$				Area, $A_g$				
	m <sup>4</sup>	in <sup>4</sup>	m <sup>4</sup> /m	in <sup>4</sup> /ft	m <sup>2</sup>	in <sup>2</sup>	m²/m	in²/ft	
Batter Pile HP 10x42	8.75E-05	210	2.87E-05	21.0	7.98E-03	12.4	2.62E-03	1.24	
Vertical Pile HP 12x53	1.63E-04	393	5.37E-05	39.3	1.00E-02	15.5	3.28E-03	1.55	
Batter Composite (wall face)	1.43E-03	3436	4.70E-04	343.6	1.11E-01	172.1	3.65E-02	17.21	
Vertical Pile Top Encased	2.96E-03	7111	9.73E-04	711.1	1.54E-01	238.7	5.08E-02	23.87	

Table 7.2 Section Properties of Beam Elements

![](_page_240_Figure_2.jpeg)

Fig. 7.6. Resulting wall deflection over the inclined face of the wall  $(\delta/L)$ . Triangle and square dots are configurations of wall geometry and soil properties excluded from the result analyses due to large wall deflection  $(\delta/L > 1/240)$ .

the 1/240 and 1/120 limits, defined by the ratio  $\beta/\phi$  and the stability number  $H\gamma/c$ . Circular dots outside those zones represent the cases without excessive wall deflection. Circular dots inside the shaded high deflection zones indicate that there were cases without excessive deformation in the shaded areas and suggest that the presented criteria will be conservative if used to evaluate serviceability of the PFRW.

#### Modeling of tie-down anchors

The vertical tie-down anchors were modeled using a linear elastic bar element with 3 m (10 ft) of bonded length inside the limestone rock. A Young's modulus of 200 GPa (29,000 ksi) was used. The pre-tensioning force input to each FEM model was estimated based on stability calculations against overturning using the Coulomb earth pressure theory. In reality, more than one anchor may be required and they would be evenly distributed between pile frames (as shown in Fig. 7.2). However, for plane strain modeling purposes a single anchor with the total force and the necessary cross-sectional area was implemented on a per unit width basis. If the predicted horizontal translation of the wall was much larger than that required to mobilize an active condition (0.001 - 0.01) H (Salgado 2008) for the specific soil being analyzed, the pre-tensioning force was increased until this condition was satisfied. For simplicity, the anchor element was aligned with the centerline of the vertical pile. In reality, the anchor is installed on the centerline of the walers, which creates an eccentric force and moment. This moment was neglected for modeling purposes, because it was insignificant for the prediction of the earth pressures.

#### Definition of boundary condition and distance to boundaries

The vertical boundaries of the soil mass were modeled such that only vertical translation and rotation were allowed. In this way, the soil was free to experience settlement/heave at each construction stage. The boundary conditions at the bottom of the soil mass were chosen to allow rotation only. While the H-piles are deeply embedded into the soil which creates a soil-pile connection somewhere between a fully fixed and truly pinned condition, they were modeled as pinned at the contact with the rock. To reflect the additional horizontal restraint offered by the pavement system, a second set of models were created where the battered concrete wall was pinned at a depth of at least 0.5 m (1.6 ft) below final grade, which also corresponds to the depth of embedment of the concrete wall face. Results indicated that this additional restraint resulted in slightly higher normal stresses on the battered wall (Figs. 7.20 and 7.21 in Appendix), and thus it was chosen for being more conservative.

The extent of the model boundaries was investigated via a series of FEM analyses. The analyses indicated that boundary effects are negligible if the horizontal distance to the vertical boundaries are not less than 6H for the  $\alpha = 0^{\circ}$  cases and no less than 12H for the  $\alpha = 20^{\circ}$  including 4H of sloped backfield as shown in Fig. 7.7. A soil depth below the excavation level equal to H was implemented for each model. Even though this value is somewhat arbitrary, the pinned condition of the wall at the final grade level helps to reduce the influence of the embedment depth of the battered pile on the numerical computations, making the predicted stresses more insensitive to the selected soil depth in the models.

![](_page_243_Figure_0.jpeg)

**Fig. 7.7.** Typical FEM model for the PFRW: boundary conditions, distance to boundaries and structural elements considered.

#### Construction and excavation sequence

Table 7.3 outlines the construction and excavation sequence as approximated in the FEM model. The concrete cap and the fill were simultaneously introduced at the last construction step, even though they are constructed sequentially.

#### Model stresses vs. field measured stresses: The SmartFix case

The SmartFix wall was constructed as part of the 2003-2009 reconstruction of I-40 in downtown Knoxville. The geometry of the wall at the tallest section was: H = 7.62 m (25 ft),  $\alpha = 20^{\circ}$ , and  $\beta = 71.6^{\circ}$ . The soil stratigraphy in the area is typical of the Chapman Ridge formation, i.e. silty-sandy clay residual soils over limestone and sandstone rock, and thus it is ideal for the PFRW system. The geotechnical exploration indicated a diversity of soil types from clays (CL, CH) and silts (MH, ML) to silty sands (SM) and clayey sands (SC) (Wilbur Smith and Associates 2006). Results obtained from 14 consolidated undrained (CU) triaxial tests and in situ borehole shear tests indicated that the effective Mohr-Coulomb shear strength parameters could be represented by  $\phi = 32^{\circ}$  and c = 11.7 kPa (244 psf). The total unit weight was taken as  $\gamma = 19$  kN/m<sup>3</sup> (121 pcf).

During the wall construction, an instrumentation system to measure the soil pressures on the battered wall was installed and continuously monitored from April 2008 to June 2013. The soil pressures were measured with Geokon Model 4800 Earth Pressure Cells, which were installed at the bottom half of the wall (approximately at 1/6, 1/3 and 1/2 of the total height of the wall) between the timber lagging and the concrete face. Two different sections were instrumented, with three pressure cells installed at the center and three at the edge of each section (Fig. 7.8).

Step	Description/Load
1	Geostatic stresses (self-weight of soil).
2	Dry Excavation.
3	Driving of steel piles including self-weight.
4	Installation and pre-tensioning of anchors.
5	Removal of soil in layers with thickness of $1-1.5 \text{ m} (3.3 - 5 \text{ ft})$ . This
	step involves 4 - 10 construction stages.
6	Installation of concrete face, concrete cap and backfill soil. Section
	properties of battered element are increased to reflect the increased
	rigidity of the composite section (steel + concrete) and the self-weight.

Table 7.3 Construction sequence as approximated in the FEM models

![](_page_245_Picture_2.jpeg)

**Fig. 7.8.** Earth pressure cells (white circles) positioned in Sections 1 and 2. Distances shown for Section 1 are along the inclined plane and measured from the final grade.

The field measured earth pressures at different time points were compared with those predicted by a FEM model with the geometry and material properties corresponding to the SmartFix wall. The implemented elastic constants for the soil were  $E = 1 \times 10^5$  kPa (2,089 ksf) and v = 0.3.

#### **Results and Discussion**

#### Normal stresses on battered wall

The typical distribution of the normal stresses on the battered wall obtained with FEM analyses are shown in Fig. 7.9 for a wall of effective height H = 7.62 m (25 ft) (SmartFix height) constructed in a sandy soil with  $\phi = 30^{\circ}$  and  $c/\gamma = 0.3$  m (1 ft) [e.g. c = 5.7 kPa (119 psf) and  $\gamma = 19$  kN/m<sup>3</sup> (121 pcf)]. The theoretical Rankine active stress on a vertical plane for this soil is also plotted in dashed lines for comparison purposes. The FEM results are consistent with the literature (Coulomb 1776, Sokolovskiĩ 1960) where for inclined walls the earth pressures decreases as the wall face inclination  $\beta$  approaches the friction angle  $\phi$ . Higher earth pressures are obtained when the backfill slope  $\alpha$  increases. The classical Rankine stresses are, in general, much higher than those predicted by the FEM, and would produce over-conservative force resultants and overturning moments for stability calculations. The FEM results for all the geometric combinations show a sharp increase of stress as the depth approaches the final grade, which is a result of a higher horizontal restraint at that point. This is particularly evident for the  $\beta = 50^{\circ}$  case where the large inclination and pinned pile tip induce a more at-rest state of lateral soil pressure. Similar earth pressure distributions were obtained for all the investigated cases.

![](_page_247_Figure_0.jpeg)

Fig. 7.9. Typical FEM results for normal stresses on battered wall for *H*=7.62 m (SmartFix wall height). Results shown for a sandy soil with  $\phi = 30^{\circ}$  and  $c / \gamma = 0.3$  m.

To provide a simplified means to design the structural members of the wall, equivalent uniform pressures normal to the wall face ( $\sigma_{Neq}$ ) were calculated for all the cases (Fig. 7.10a). The equivalent uniform pressure is that pressure which will produce the same force resultant as the original FEM pressure distribution, with no significant changes in the design moments and shear forces. The resulting  $\sigma_{Neq}$  vs. *H* are shown in Fig. 7.11 for the  $\phi = 30^{\circ}$  and  $c/\gamma = 1$  m (3.3 ft) soil case, for a wall with  $\beta = 71.6^{\circ}$  (SmartFix wall inclination) for  $\alpha = 0^{\circ}$  and 20°. Also shown in Fig. 7.11 are the equivalent uniform pressures obtained via the theoretical Coulomb and the Coulomb with tension correction methods (Bowles 1968) (Fig. 7.10b). The Coulomb theory accounts for the effects of sloped backfills and inclined wall faces, and the tension correction removes the small artificial tensile stress that is calculated due to the soil cohesion. The Coulomb uniform pressures with tension correction resulted in a conservative prediction of the normal stresses for the most detrimental conditions (high and steep walls with sloped backfills and soils with low cohesion); however, the Coulomb predicted soil pressures become lower than the FEM prediction (unconservative) as the soil cohesion increases and the wall inclination decreases as illustrated in Fig. 7.12 for the  $\beta = 50^{\circ}$  and  $\alpha = 20^{\circ}$  case (same soil properties). The Coulomb uniform pressure with tension correction may constitute a conservative simplified design approach if the predicted stresses are higher than those predicted by the FEM. This matter is illustrated later for the SmartFix field case study.

Generalized FEM results in terms of the dimensionless quantities  $\sigma_{Neq}/c$  and  $H\gamma/c$  are shown in Fig. 7.13 for the  $\phi = 30^{\circ}$  case. This figure shows the potential range of normal stresses over the battered wall, represented here by the two bounds of wall inclinations investigated

![](_page_249_Figure_0.jpeg)

**Fig. 7.10.** Definition of equivalent normal distributed pressures  $\sigma_{Neq}$ : a) FEM and b) Coulomb and Coulomb with tension correction.

![](_page_249_Figure_2.jpeg)

Fig. 7.11. Equivalent normal uniform stress  $\sigma_{Neq}$  vs. wall height *H*, for the  $\phi = 30^{\circ}$  and  $c/\gamma = 1$  m soil case and the SmartFix wall inclination  $\beta = 71.6^{\circ}$ . a)  $\alpha = 0^{\circ}$  and b)  $\alpha = 20^{\circ}$ .

![](_page_250_Figure_0.jpeg)

Fig. 7.12. Illustration of a case when the theoretical Coulomb earth pressures are lower than FEM predictions, and therefore unconservative ( $\phi = 30^{\circ}$  and  $c/\gamma = 1$  m,  $\alpha = 20^{\circ}$ , and  $\beta = 50^{\circ}$ ).

![](_page_250_Figure_2.jpeg)

Fig. 7.13. Dimensionless equivalent normal earth pressures vs. stability number as predicted by FEM for  $\phi = 30^{\circ}$ . a)  $\alpha = 0^{\circ}$  and b)  $\alpha = 20^{\circ}$ .

 $(\beta = 50^{\circ} \text{ and } 80^{\circ})$ , and can be used to obtain the expected equivalent uniform pressure given the geometrical wall configuration ( $\alpha$ ,  $\beta$ , and H) and the Mohr-Coulomb soil properties ( $\phi$  and  $c/\gamma$ ). Similar charts for the  $\phi = 20^{\circ}$  and  $40^{\circ}$  were created (Figs. 7.22, 7.23 and 7.24 in Appendix). For each set of data, a linear regression was fit. Within each value of  $\phi$ , the slope for  $\beta = 50^{\circ}$  is somewhat similar for  $\alpha = 0^{\circ}$  and  $20^{\circ}$ , indicating that the stress  $\sigma_{Neq}/c$  is not highly affected by the backfill inclination at low values of  $\beta$ . The role of the backfill, however, becomes significant as the wall becomes steeper.

#### Destabilizing moments for overturning stability

The forces and moment arms used to calculate the wall overturning moments (*M*) are illustrated in Fig. 7.14. FEM results are compared with theoretical values from Coulomb pressures acting on the inclined wall and at-rest pressures acting on vertical cap. The FEM shear forces along the vertical edge of the cap were neglected, since they create a moment that resists the overturning. The FEM and Coulomb shear components on the battered wall act through the point of rotation and do not play any role in the overturning moment. Fig. 7.15 shows results for the  $\phi = 30^{\circ}$  and  $c/\gamma = 1$  m (3.3 ft) soil case for walls with  $\beta = 71.6^{\circ}$  (SmartFix wall inclination) and  $\alpha = 0^{\circ}$  and 20°. Similarly to the normal stress cases, the moments calculated via Coulomb's theory are a conservative prediction for the most detrimental wall conditions, but it under predicts the moments for the cases with low  $\beta$ . Notice that FEM and Coulomb moments follow a similar trend as *H* increases.


**Fig. 7.14.** Forces and moment arms employed for overturning moments. a) Numerical FEM and b) Theoretical Coulomb with tension correction on wall + at-rest condition on cap.



Fig. 7.15. Destabilizing moment *M* vs. wall height *H*, for the  $\phi = 30^{\circ}$  and  $c / \gamma = 1$  m soil case and the SmartFix wall inclination  $\beta = 71.6^{\circ}$ .

FEM results in terms of a normalized moment  $M_{eq} = M\gamma/c$  and the dimensionless number  $H\gamma/c$ are shown in Fig. 7.16 for the  $\phi = 30^{\circ}$  case. The normalized moments follow a linear trend. Note that the results in Fig. 7.16 include all the investigated values of  $\beta$ , and therefore the regression equation is independent of the wall inclination. For all the investigated cases  $R^2 \ge 0.85$  and indicates that a linear trend is appropriate. Similar resulting linear trends were also obtained for the  $\phi = 20^{\circ}$ and  $40^{\circ}$  cases (Figs. 7.25, 7.26 and 7.27 in Appendix).

#### Design equations to predict earth pressures and destabilizing moments

From the regression analyses based on the FEM results previously described, analytical approximations were developed to predict the equivalent uniform pressures  $\sigma_{Neq}$  and the overturning moments M as a function of the wall geometry and the soil properties. By assuming linear behavior of the quantities between  $\alpha = 0^{\circ}$  and  $20^{\circ}$ , and between  $\phi = 20^{\circ}$  and  $30^{\circ}$ , and  $\phi = 30^{\circ}$  and  $40^{\circ}$ , the following equations for  $\sigma_{Neq}$  and M were obtained:

$$\sigma_{Neq} = 0.018 H \gamma \frac{\sqrt{\tan \beta}}{\tan \phi} (1 + \tan \alpha)$$
(7.4)

$$M = 45H(1+1.5\tan\alpha), \text{ with } H \text{ in m (SI units)}$$

$$M = 3.1H(1+1.5\tan\alpha), \text{ with } H \text{ in ft (U.S. customary units)}$$
(7.5)

Note that M has SI units of kN-m/m, since the constant "45" has units of kN/m. Likewise, M has U.S. customary units of kip-ft/ft. A validation of these simplified models is offered in Fig. 7.17 for all the cases comprising  $\phi = 30^{\circ}$ . Here,  $\sigma_{Neq}$  (Fig. 7.17a) and M (Fig. 7.17b) predicted by Eqs. (7.4) and (7.5) respectively, are plotted against the FEM predictions. Notice that the distribution of the points tends to follow the 1:1 (45°) line, indicating correlation between the



Fig. 7.16. Normalized overturning moment vs. stability number as predicted by FEM for  $\phi = 30^{\circ}$ .



Fig. 7.17. Validation of design equations: a)  $\sigma_{Neq}$  (kPa) and b) M (kN-m/m).

simplified proposed equations and the FEM results. The ratio between the design equation and the FEM prediction for  $\sigma_{Neq}$  is 1.2 and for M is 1.5, suggesting that on average the design equations predict 20% higher values of  $\sigma_{Neq}$  and 50% higher values of M, making the design equations more conservative than the FEM prediction. Similar ratios were obtained for the  $\phi = 20^{\circ}$  and 40° cases. The coefficients of variation of these ratios are less than or equal to 40% for all cases. It is important to keep in mind that the range of applicability of the design equations is constrained to values of  $\alpha = 0 - 20^{\circ}$ ,  $\beta = 50 - 80^{\circ}$ , H = 5 - 15 m (16 – 49 ft),  $\phi = 20 - 40^{\circ}$  and  $c/\gamma = 0.3 - 3$  m (1 – 10 ft), and values outside the range of applicability are not defined by the Eqs. (7.4) and (7.5).

#### The SmartFix case: predicted vs. measured stresses on the inclined wall

Measured earth pressures over a period of 3 years (2010 – 2013) are shown in Fig. 7.18, along with the FEM predicted distribution of the earth pressures for the SmartFix model configuration  $[H = 7.62 \text{ m} (25 \text{ ft}), \alpha = 20^\circ, \beta = 71.6^\circ, \phi = 32^\circ, c = 11.7 \text{ kPa} (244 \text{ psf}), \gamma = 19 \text{ kN/m}^3 (121 \text{ pcf})]$  using both, a non-associated ( $\psi = 0^\circ$ ) and an associated ( $\psi = \phi$ ) flow rule, where  $\psi$  is the dilation angle. For comparison, two theoretical earth pressures that include the effects of  $\alpha$ ,  $\beta$ , and the soil-wall interface friction angle are also shown: Wedge Equilibrium (Coulomb 1776) and Upper Bound Limit Analyses (Chen and Rosenfarb 1973). The vertical line represents the corresponding value of  $\sigma_{Neq}$  obtained from the proposed design equation [Eq. (7.4)]. Also shown is the equivalent uniform stress obtained from the simplified Coulomb design approach that includes tension correction.



Fig. 7.18. Field measured stresses over time vs. theoretical and numerical earth pressures.

The results in Fig. 7.18 indicate that the pressures predicted by the numerical models are in reasonable agreement with the field measurements, with the exception of the larger predicted stresses near the bottom of the wall. The FEM predicted stresses comprising an associated soil  $(\psi = \phi)$  are on average 11% higher than those obtained with the non-associated model; similar behavior has been previously described in the literature (Benmeddour et al. 2012). While dilation effects are observed in soils, they rarely follow an associated flow rule, and dilation angles are usually smaller than  $\phi$ . It has been also shown that associated models tend to over predict failure loads in the context of "confined" problems. The PFRW can be considered as a partially confined system, and thus the degree of "over-prediction" is moderated, which makes the model more insensitive to the selected value of  $\psi$ . The results in Fig. 7.18 suggest that the equivalent uniform pressure Eq. (7.4) (developed from non-associated FEM results) predicts the measured earth pressures with a moderate degree of conservatism, and it is a better representation than the Coulomb and Chen approaches for this type of wall. The simplified desing approach based on uniform Coulomb stresses with tension correction provides, for this particular wall configuration, a prediction of stresses higher than those obtained from Eq. (7.4), and therefore can also be used for design. It should be kept in mind that the classical solutions assume linearly increasing pressure distributions, typical of walls with rigid translations as the deformation mode. It has been demonstrated that the typical earth pressure distribution of the PFRW is neither linear nor monotonicly increasing, and classical solution methods are not always conservative for this type of wall. A PFRW specific method such as the design equations based on FEM analyses [Eqs. (7.4) and (7.5)] and the uniform Coulomb stresses with tension correction (when predicted values are higher than design equations), provides a more reasonable design prediction for the potential geometrical configurations of PFRW systems covered here.

## Conclusions

The Piling Framed Retaining Wall is a novel system which allows top-down construction and requires minimal right of way. This wall concept was found to be significantly less expensive than traditional retaining wall systems for two walls of this type successfully built along the I-40/I-75 corridor in Knoxville, TN. For their design, the magnitude and distribution of the earth pressures and overturning moments were estimated based on traditional earth pressure theories. However, a rational design procedure for this wall system has not been fully developed.

To develop a simple design approach, a series of FEM analyses were conducted to compute the earth pressures on the wall face and the overall destabilizing moments for a range of wall geometries and different soil properties. The numerical results were used to develop simplified equations for design. It was found that the typical earth pressure distribution on the face of the PFRW is neither linear nor monotonically increasing. The design equations yield conservative results, relative to the FEM models, for most practical geometries and soil properties. Field measured earth pressures on the SmartFix prototype were well predicted by both, the design equations and the design method based on Coulomb stresses with tension correction. The proposed methods can be used to design future configurations of PFRW systems with no need to create advanced numerical models.

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## Appendix

#### Active, passive, and at rest states of the soil mass

For a soil medium stressed only by its own weight, any point on a horizontal plane at depth z will experience the same level vertical effective stress  $\gamma z$  ( $\gamma$  is the unit weight of soil), with negligible shear stresses. This vertical direction becomes therefore a principal stress direction  $\sigma_1$ . From the assumption that soil is an isotropic continuous medium, the other principal direction must be orthogonal and act in the horizontal direction  $\sigma_3$ . In reality, a third principal directions exist, but for simplification of this model to 2-D, let us assume that plane strain conditions are valid and the principal stress acting perpendicular to the paper remains intermediate between the two principal directions described.

In geo-mechanics it is a usual practice to consider that the magnitude of the horizontal stress inside the soil medium is proportional to the vertical stress by a constant K. In cases when the soil has remained unaltered since formation, this constant is considered to be at-rest conditions and is usually referred as  $K_0$ . In technical terms,  $K_0$  is the earth pressure coefficient when no lateral deformation is allowed to occur and corresponds to the in situ state of most geo-materials. Thus, when no external stresses act over the soil, the horizontal principal stress becomes  $K_0\gamma z$ . For most normally consolidated soils  $K_0 < 1$  and the horizontal direction usually becomes the minor principal direction. If, by some mechanism, the soil is permitted to yield laterally by gradually decreasing the horizontal stress level and keeping the vertical stress constant, it will be observed a progressive lateral expansion of the soil until a point when plastic deformation will start taking place along slip lines, ideally, at constant stress values (for a perfectly plastic behavior). This point corresponds to the active state in the soil medium. Since the vertical stress  $\sigma_1 = \gamma z$  remained constant, the horizontal stress has to be proportional to smaller value of K. This constant of proportionality is called  $K_a$ , the active earth pressure coefficient, and thus the lateral earth pressure at an active state is  $\sigma_{3a} = K_a \gamma z$ , where  $K_a < K_0$ . On the other hand, if the soil is permitted to yield by gradually increasing the level of lateral stress and keeping the vertical stress constant, it will be observed a lateral compression until a point where the plastic deformation will start taking place. At this point the soil has reached a passive condition and, analogously to the active case, the lateral stress will be  $\sigma_{3p} = K_p \gamma z$ , being  $K_p$  the passive earth pressure coefficient, where  $K_p > K_0$ . From this discussion  $\sigma_{3a} < \sigma_3 < \sigma_{3p}$  and therefore,  $K_a < K_0 < K_p$ . Fig. 7.19 illustrates this concept employing the Mohr-Coulomb yield criteria. These lateral extension and compression mechanisms in the soil are observed in retaining wall problems. Wherever the structural rigidity of the system is such as the wall will deform moving "away" from the soil a sufficient amount, an active condition may be achieved. The opposite is true for the passive condition. The level of soil strain required to achieve the active and passive condition is highly conditioned to the rigidity of the structural system and its mode of deformation.

#### Active and passive lateral earth pressure coefficients: a literature review

Plasticity methods for the active and passive lateral earth pressure coefficient can be grouped in two major categories: a) static based methods or b) kinematics based methods.



Fig. 7.19. Conceptual model of the different stress states on the retaining wall problem.

The *static* based approaches usually assume that portion of the material behind the wall is at the verge of flowing plastically, and only the equilibrium equations and yield conditions are satisfied. Some of the methods based on statics are listed below.

*Slip-line network*. This method attempts to determine the family of slip-line curves that will develop inside the soil mass by combining the equations of equilibrium of a continuous medium with the Mohr-Coulomb yield criterion. The results of the combination of these equations are the differential equations of the characteristics that describe the planes where the plastic deformation will take place. This formulation allows the implementation of more complex stress boundary conditions; however, for materials possessing weight, the characteristic equations become complex and exact mathematical solutions are available for a limited number of simple cases. Sokolovskii (1960), provided solutions for the lateral earth pressure problem by numerically integrating the characteristic equations. Sokolovskiĭ reported values of active and passive lateral and tangential earth pressures on a vertical wall considering: a) non-linearity of the slip-lines due to soil-wall friction interaction, b) sloped backfields and c) surcharges. The shortcomings of the slip-line method are in general related to the validity of the theoretical assumptions. Slip-line methods only take into account stress boundary conditions, when in many practical problems deformation conditions exist (Chen 2007); e.g. Sokolovskii's solution do not consider the amount of deformation that must exist to bring the soil to a limiting condition. Also, in this technique, it is assumed that only the zone inside the slip-network is in plastic equilibrium and outside this zone the stress distribution is unknown (Chen 2007).

**Rankine states of the soil mass.** From the active and passive states of the soil defined by Rankine (1856), expressions to calculate the active  $K_a$  and passive  $K_p$  earth pressure coefficients were developed. If it is assumed that no friction exist in the soil-wall interface, any rigid rotation or translation of the wall will develop a stress field composed of straight slip-lines that will intersect the horizontal at  $45^\circ + \phi/2$  for the active condition and  $45^\circ - \phi/2$  for the passive case. For these conditions the vertical and horizontal planes coincide with the principal planes and, therefore, Rankine's formulation [Eq. (7.6)] is an exact closed form solution of the slip-field problem.

$$K_a = \frac{1}{K_p} = \frac{1 - \sin\phi}{1 + \sin\phi} \tag{7.6}$$

The more general expression includes the case of a vertical wall with a sloping backfield at  $\alpha^{\circ}$ . degrees from the horizontal, as expressed in Eq. (7.7) (upper sign denotes the active condition and lower sign the passive)

$$\frac{K_a}{K_p} = \cos\alpha \, \frac{\cos\alpha \mp \sqrt{\cos^2\alpha - \cos^2\phi}}{\cos\alpha \pm \sqrt{\cos^2\alpha - \cos^2\phi}}$$
(7.7)

The limitations of Rankine's formulation are related to: a) the incapacity to include the soil-wall friction interaction and b) the inability to obtain lateral earth pressures in planes other than vertical ones.

*Lower Bound in Limit Analysis.* Limit analysis is a plasticity based theory that is rigorous within the context of its assumptions. It defines two extremes, the lower and the upper bound, that correspond to the lowest and highest collapse load that can be possibly found, whereas the true collapse load is somewhere in between these two bounds. Similar to the slip-line method, the lower bound theorem only requires the satisfaction of the equilibrium equations and yield conditions; 245

however, here a complete stress distribution satisfying equilibrium and yield condition everywhere needs to be established. Lancellotta (2002), based on the lower bound theorem, developed an analytical expression for the lateral active and passive earth pressures on vertical walls [Eqs. (7.8) and (7.9)]. The solution includes the effects of the soil-wall friction ( $\delta$ ) and it simplifies to the Rankine's expression for frictionless soil-wall interfaces:

$$\frac{K_a}{K_p} = \left[\frac{\cos\delta}{1\pm\sin\phi} \left(\cos\delta\mp\sqrt{\sin^2\phi - \sin^2\delta}\right)\right] \exp\left(\mp 2\theta\tan\phi\right)$$
(7.8)

$$\theta = \frac{1}{2} \left[ \arcsin\left(\frac{\sin\delta}{\sin\phi}\right) \mp \delta \right]$$
(7.9)

Though rigorous, Lacellotta's approach does not allow the evaluation of cases with a) sloped backfield and b) non-vertical walls.

*Kinematics* based methods, on the other hand, usually involves a variety of assumed slip surfaces that define different soil wedges. The resultant earth pressure is calculated by means of force equilibrium or by means of external rate of work and internal rate of energy dissipation. Some of the methods based on kinematics are listed below.

*Coulomb's wedge method.* Based on the concept of limit equilibrium, Coulomb (1776) proposed the existence of one critical slip plane where the failure of a soil mass will take place, and will produce the maximum mobilized lateral stresses over the wall. In reality, many wedges of planar failure surfaces can be constructed, and an optimization process to find most critical wedge is necessary. Coulomb solved this problem in terms of force equilibrium, considering the soil weight

and soil-wall friction force, for a horizontal backfield. The maximum lateral force that resulted from the equilibrium and optimization is:

$$P_a = \frac{1}{2} K_a \gamma H^2 \tag{7.10}$$

$$K_{a} = \tan^{2} \left( 45 - \frac{\phi}{2} \right) = \frac{1 - \sin \phi}{1 + \sin \phi}$$
(7.11)

where  $\gamma$  is the soil unit weight, *H* is the wall height. Note  $K_a$  coincides with the coefficient of earth pressure obtained from the Rankine's stress analysis. Subsequent work conducted by Poncelet and Woodbury (1854), Rebhann (1871) and Müller-Breslau (1906) extended Coulomb's wedge analysis to include the effects of a) soil-wall friction  $\delta$ , b) sloped backfield at angle  $\alpha$  to the horizontal and c) wall inclination at angle  $\beta$  from to the vertical. The obtained active and passive coefficients are (Salgado 2008):

$$\frac{K_a}{K_p} = \frac{\sin^2(\beta \pm \phi)}{\sin^2(\beta) \sin(\beta \mp \delta) \left[1 \pm \sqrt{\frac{\sin(\phi + \delta)\sin(\phi \mp \alpha)}{\sin(\beta \mp \delta)\sin(\alpha + \beta)}}\right]^2}$$
(7.12)

In Eq. (7.12) it is assumed that the slip failure is planar, when in reality the soil-friction along the vertical boundary modifies the shape of the slip-lines to be curved. Also, this procedure does not predict the distribution of the vertical and lateral stresses. Perhaps, those are the two major shortcomings of the Coulomb's formulation.

*Non-linear wedge analysis*. Recognizing the inconsistency of assuming a planar slip failure when friction exists between soil and wall, Caquot and Kérisel (1948) provided an advanced mathematical formulation for the earth pressure problem assuming a log-spiral failure mechanism.

Caquot and Kérisel reported tables of active and passive earth pressure coefficients as a function of a) internal friction angle, b) soil-wall friction, c) angle of the sloped backfield and d) angle of wall face inclination.

*Upper Bound in Limit Analysis.* This theorem states that an external load can only be resisted by the soil mass if the external rate of work of an assumed deformation mode is less than the internal rate of energy dissipation. If the velocity boundary conditions and the strain compatibility equations are satisfied by any load resulting by equating the external and internal energy rates, then the load is higher than the failure load. This theorem, assumes a perfectly plastic soil behavior with an associated flow rule. Chen and Rosenfarb (1973) reported upper bound solutions for the active and passive earth pressure coefficients from the optimization of a variety of assumed failures. More recently, Soubra and Macuh (2002) proposed different active and passive earth pressure coefficients for gravity loads, surcharge and soil cohesion, assuming a log-spiral failure and the effects of soil-wall friction and sloped backfield. The combination of three coefficients constitutes the lateral force acting on the wall.

Other kinematic based approaches for the passive earth pressure are a) *the method of triangular slices* (Zhu and Qian 2000) and b) *the upper-bond numerical analysis* (Antão et al. 2011). In general terms, all the methods previously described provide similar results; e.g. the lower and upper bounds of the active earth pressure coefficients are close, and Sokolovskiĭ's active coefficients falls somewhere in between these two bounds (Loukidis and Salgado 2011). Also, in one way or the other, all these methods assume a perfectly plastic soil behavior and they do not account for progressive failure experienced in soils. In reality more than one active state can occur

after the peak or yielding has occurred (Loukidis and Salgado 2011). All methods, except those based on Limit Analysis, do not include the effects of dilitancy and none of them include the effects of volumetric strains in shear zones. Despite of the extensive research that has been conducted on earth pressure coefficients, theoretical solutions still diverge from experimental results, perhaps due to the complex soil deformational behavior (Niedostatkiewicz et al. 2011).

#### At-rest lateral earth pressure coefficients: a literature review

Among all the equations that have been proposed to calculate  $K_0$ , the expression based on the Elasticity Theory seems to be the most theoretically rigorous. This expression is a function of the Poisson's ratio of the material:

$$K_0 = \frac{\upsilon}{1 - \upsilon} \tag{7.13}$$

For cohesionless normally consolidated materials, Jaky (1944) offered and empirical equation based on the internal friction angle:

$$K_0 = 1 - \sin\phi \tag{7.14}$$

Eq. (7.14) was validated for cases when the backfield is in a loose state and improved for dense compacted backfield cases (Sherif et al. 1984). This expression [Eq. (7.15)] is a function of  $\phi$ , the in situ dry unit weight  $\gamma_d$  and the dry unit weight of soil at the loosest state  $\gamma_{dmin}$ 

$$K_0 = 1 - \sin\phi + \left(\frac{\gamma_d}{\gamma_{dmin}} - 1\right) \cdot 5.5$$
(7.15)

To consider the effects of the soil stress history, Jaky's equation can be modified to include the over consolidation ratio  $OCR = \sigma'_{vp} / \sigma'_{v}$ , where  $\sigma'_{vp}$  is the pre-consolidation effective pressure and  $\sigma'_{v}$  is the effective vertical stress being experienced by the soil (Salgado 2008):

$$K_{0} = K_{0(NC)} * \sqrt{OCR} = (1 - \sin \phi) \sqrt{OCR}$$
(7.16)

Mayne and Kulhawy (1982) also provided an empirical expression to account for the effects of *OCR* on Jaky's equation.

Hendron (1963) and Hendron et al. (1963) formulated a theoretical expression for  $K_0$  for a body constituted of perfect uniform spheres as a function of  $\phi$  and the friction coefficient f:

$$K_0 = \frac{1}{2} \left( \frac{1 - f}{1 + f} \right)$$
(7.17)

$$f = 3\sqrt{6/8}\sin\phi - \sqrt{6/8} \tag{7.18}$$

For normally consolidated cohesive soils, Brooker and Ireland (1965) suggested that Jaky's equation tends to over predict  $K_0$  values and proposed a new expression that better fit this condition:

$$K_0 = 0.95 - \sin\phi \tag{7.19}$$

Brooker and Ireland (1965) also studied the effect of OCR and plasticity index *PI* over  $K_o$ , and reported their findings in terms of a solution chart. For over consolidated clays, Fioravante et al. (1998) suggested the following empirical expression based on in situ seismic shear waves experiments:

$$K_0 = 0.54 * OCR^{0.57} \tag{7.20}$$

250

Alpan (1967) exploring empirical correlations for  $K_0$  suggested that the following Eq. (7.21) is reliable for normally consolidated clays and provides a good agreement with Brooker and Ireland's formulation [Eq. (7.19)]

$$K_0 = 0.19 + 0.233\log(PI), PI \text{ in }\%$$
 (7.21)

Other expressions available in the literature consider the effects of the secondary compression during loading (Mesri and Castro 1987) and the non-linear behavior of  $K_0$  at high vertical pressures (Tian et al. 2009, Zhao et al. 2010). An empirical equation based on CPTU test data was proposed by (Kulhawy et al. 1990), and expressions based on non-destructive SCTPU and Cross Hole tests have been developed by Sully and Campanella (1995) and Fioravante et al. (1998), and validated by Guojun et al. (2011).

#### Definition of boundary condition and distance to boundaries

The vertical boundaries of the soil mass were selected such that only vertical translation and rotation were allowed. In this way, the soil was free to experience settlements/heave at each construction stage. The horizontal boundary at the bottom had restricted translation and was only allowed to rotate. Regarding the structural members, the H-piles are driven to refusal onto rock, which creates a soil-pile connection in between a fully fixed and truly pinned condition. It was assumed for modeling purposes that the piles were truly pinned at the bottom. The battered concrete wall, on the other hand, is extended to a depth of at least 0.5 m below final grade, which also creates a boundary in between an unrestricted translation and fully pinned condition. Two different models were considered: a) the 0.5 m below final grade is truly pinned and b) the 0.5 m below final grade is not pinned, and the translation is restricted only by a passive soil pressure (Fig. 7.20).



**Fig. 7.20.** Two FEM models considered: a) the 0.5 m of the battered pile below the final grade is truly pinned and b) the battered pile is truly pinned at the contact with the basal rock.



**Fig. 7.21.** Difference in resulting normal stresses from the two battered wall end conditions investigated (SmartFix case). The model comprising a pinned condition at the 0.5 m below final grade resulted in higher overall stresses.

Preliminary FEM results indicated the first model (Fig. 7.20a) predicted in general higher normal stresses on the battered wall (Fig. 7.21), and therefore, it was chosen for being more conservative.

Generalized FEM soil stresses in terms of the dimensionless quantities  $\sigma_{_{Neq}}/c$  and  $H\gamma/c$ 



Fig. 7.22. Dimensionless equivalent normal earth pressures vs stability number as predicted by FEM for  $\phi = 20^{\circ}$ .



Fig. 7.23. Dimensionless equivalent normal earth pressures vs stability number as predicted by FEM for  $\phi = 30^{\circ}$ .



Fig. 7.24. Dimensionless equivalent normal earth pressures vs stability number as predicted by FEM for  $\phi = 40^{\circ}$ .





Fig. 7.25. Normalized overturning moment vs. stability number as predicted by FEM,  $\phi = 20^{\circ}$ .



Fig. 7.26. Normalized overturning moment vs. stability number as predicted by FEM,  $\phi = 30^{\circ}$ .



Fig. 7.27. Normalized overturning moment vs. stability number as predicted by FEM  $\phi = 40^{\circ}$ .

## Validation of simplified design equations



**Fig. 7.28.** Validation of design equations: a)  $\sigma_{Neq}$  and b) *M* for  $\phi = 20^{\circ}$ .



**Fig. 7.29.** Validation of design equations: a)  $\sigma_{Neq}$  and b) *M* for  $\phi = 30^{\circ}$ .



**Fig. 7.30.** Validation of design equations: a)  $\sigma_{Neq}$  and b) M for  $\phi = 40^{\circ}$ .

# **Determination of soil strength parameters**



Fig. 7.31. Selected M-C strength parameters based on triaxial p-q diagram at 14% strain level.

# **Chapter 8. Conclusions**

## The Low Compaction Grading Technique on steep slopes (chapters 2 and 3)

The analysis of several potential modes of failures via LEM and FEM analyses suggest that for FRA slopes the governing failure mode is shallow and contained within the weak loose surface layer. The determination of the strength parameters of the core is not important for FRA slope design. Because the infinite slope method adequately approximates the shallow failure mode and accurately predicts the FS, it may be an appropriate method to evaluate the performance of FRA slopes, so more sophisticated analyses are not necessary for most applications. Since the unit weight of the material is not considered in the infinite slope expression, field measurements of the highly variable unit weight are not required for long-term analyses, but necessary for partially saturated conditions and seismic analyses. The angle of repose was suggested to be a conservative estimate of the internal friction angle and it is consistent with the loose nature of the FRA material. This provides a means to quantify the friction angle of mine spoil having large number of oversize particles, which has been traditionally assumed based on experience.

Slope stability analyses for partially saturated mine soils showed that seasonal increments in the stability of steep FRA slopes due to matric suction are possible, and the static long-term stability is a lower bound of the real field performance. However, it was shown that temporal saturation of the loose surface layer is likely to occur and slope stability analysis employing traditional saturated strength parameters should prevail in the design of FRA slopes. The interpretation of results is restricted to the hydrological conditions experienced during the time period at which the data was collected. This period was characterized by relatively poor establishment of ground cover, and therefore, evaporation was probably the main mechanism inducing matric suction and the relative increase in stability. Under seismic loading conditions, the proposed infinite slope equation 262

showed good agreement with the simplified Bishop's method of the slices. While the selection of the pseudo-static coefficient based on a fraction of the probable *PGA* is common practice, the effects of frequency and duration are not included. A method based on spectral response may provide a better tool for evaluating slope stability, and a combination of the Bray and Travasarou (2009) approach with the proposed modification of the infinite slope equation was suggested. Charts were developed for this method employing site-specific  $S_a$  values and local soil properties and slope geometry. Here, the Factor of Safety (FS) is a function of parameters that have less selection subjectivity for the designer and are a function of the importance of the project, available seismic data, and accuracy of the site characterization.

The results presented in these two chapters suggest that the FRA has no negative impact on slope stability. It is important to keep in mind that since the material receives minimum compaction effort, as shearing takes place this material will initially densify, becoming stronger over time. Also, further densification will occur through root reinforcement after reforestation, suggesting that the most detrimental conditions for the stability of FRA slopes are those experienced during the first years after construction. In this sense, the benefits of quick and proper establishment of vegetation go beyond the reduction of soil loss by water erosion, influencing also the mechanical stability of steep FRA slopes in the long-term. Future research work on steep FRA slopes could be directed to understand the role of vegetation on the stability of this type of reclamation process. Research questions to be addressed include the following: To what extent do plant roots positively affect the soil strength and mechanical stability of steep FRA slopes over time? Is there any arching soil effect that develops in the mine spoils due to this natural "anchor reinforcement" system? If so, how would this arching effect be quantified in stability calculations for FRA slopes? Would 263
the fundamental failure mode of these slopes be changed? How will root transpiration rates affect the unsaturated stability patterns observed on these sites? Would overloading and overturning of woody vegetation during high velocity wind events counteract the positive effects of root reinforcement on steep FRA slopes?

The answers to these research questions would help to understand the behavior of steep FRA slopes once the vegetation has been successfully established, providing a broader picture of the field performance in the very long-term.

## The mechanical and erosional stability of concave slopes (chapters 4, 5 and 6)

Concave slopes not only resemble natural contours, but also have superior water erosion resistance. The design of concave slopes requires: a) the definition of concave shapes that provide a desired FS, b) a quantitative measure of the erosion/sediment delivery reduction, and c) determination of possible loss of mechanical stability due to improper construction. A mathematical description of the concave slope surface at limiting equilibrium (FS = 1) for the  $\phi > 0$  and  $\gamma > 0$  case was developed, based on the slip line field theory of Sokolovskii (1960). The approximation transformed a set of differential equations into a single algebraic expression, which was validated using LEM and FEM analyses. Employing the developed equation, a mechanism to describe the coordinates of concave slopes that satisfy a desired FS (long-term conditions) for any combination of  $\phi$ , c,  $\gamma$  and slope height was offered. A simplified graphical solution was developed to estimate the required undrained shear strength such that the design long-term stability is assured. Results from RUSLE2 analyses indicate that these concave slopes yield 15-40% less sediment than

planar slopes of equal FS, regardless of soil erodibility and weather conditions. Sensitivity analyses reveal that the stability of concave slopes is not significantly influenced by errors in the constructed profile of as great as 200 mm of vertical deviation. Therefore, the stability is not compromised when typical high accuracy GPS construction equipment is employed for concave slope construction.

Similarly, natural landforms seek to achieve erosion and sediment transport equilibrium, and in doing so evolve into concave forms. In theory, this equilibrium shape would not suffer further adjustments so the total erosion production is minimized. From a conceptual perspective, it was shown that the equilibrium shape may not be unique, and a family of potential slope shapes exist according to the level of constant erosion the slope experiences. From the fundamentals of the widely recognized RUSLE2 model, a set of equations to describe concave slopes in erosion equilibrium at different levels of constant erosion was developed. The resulting concave profiles were compared with critical concave slopes (FS = 1) and with concave slopes having FS = 1.25, 1.5, and 1.75, identifying those shapes that satisfy both erosion equilibrium and the desired degree of mechanical stability. A mathematical expression defining the limiting erosion at which equilibrium erosion shapes become sustainable was provided as a function of the Mohr-Coulomb parameters.

The findings presented here suggest that concave slopes can be constructed to achieve both minimal erosion as well as mechanical stability, and that the slope profile can be chosen such that erosion rates and desired mechanical stability are balanced to satisfy site conditions and regulatory requirements. Concave slopes not only offer superior erosion control, but also a more natural

appearing landform, being an attractive alternative to the traditional planar slopes that are typically constructed today.

The results from the illustrative example in *chapter 5* suggest that concave and planar slopes having the same FS can potentially have a similar lateral extent, but that concave slopes comprise significantly less amount of soil. Since an analytical expression for the concave slope is now available, this could be incorporated into standard cut and fill balance calculations. On projects with significant cut slopes, the additional excavation may be costly. However, on fill slopes and embankments, and especially in mine reclamation where a shortage of material is common, the construction of concave slopes could be an advantage. Notice that under the FRA reclamation idea, as discussed in *chapters 2* and *3*, the quick establishment of vegetation is crucial for surficial water erosion control. The results obtained in here indicate that erosion reduction could be significantly enhanced if the FRA slopes are constructed with concave profiles. This requires, however, that not only the loose surface, but also the strong core must be concave. Future research on mine reclamation could be focused on understanding the erosion and mechanical slope stability behaviors of concave FRA slopes. Full-scale concave FRA slopes could be constructed and instrumented to investigate hydrology and sediment yield, and the results compared to those reported by Hoomehr et al. (2012) and Hoomehr et al. (2013) on steep planar FRA slopes. The quantitative benefit in terms of reduced soil loss when FRA slopes are built to be concave needs to be determined. Since the plane of discontinuity is not planar anymore, it will be necessary to investigate the fundamental failure modes of concave FRA slopes. Would the plane of discontinuity be the governing failure condition instead of the toe and face failure mechanism inherent of concave slopes? If yes, it will be necessary to investigate a new rational design

methodology for concave FRA slopes. On the other hand, what would be the implications of concave FRA slopes on the construction methodology? Would the contour haulback method typically used for FRA construction still be a valid construction methodology? The answer to these research questions would guide the development of an improved FRA methodology (FRA 2.0) that will be beneficial for erosion reduction not only post reforestation, but also during initial slope adjustments before vegetation is well established.

## The Piling Framed Concrete Retaining Wall (chapter 7)

The Piling Framed Retaining Wall or confined slope system is a novel concept allowing top-down construction and requires minimal right of way. This wall concept was found to be significantly less expensive than traditional retaining wall systems under these conditions. To develop a rational design approach for this type of wall, a series of FEM analyses were conducted to compute the earth pressures on the wall face and the overall destabilizing moments for a range of wall geometries and different soil properties. The numerical results were used to develop simplified equations for design. It was found that the typical earth pressure distribution on the face of the PFRW is neither linear nor monotonically increasing. The design equations yield conservative results--relative to the FEM models--for most practical geometries and soil properties. Field measured earth pressures on the SmartFix prototype were well predicted by both the design equations and a proposed design method based on Coulomb stresses with tension correction. The proposed methods can be used to design future configurations of PFRW systems with no need to create advanced numerical models.

The PFRW concept could be extended to applications beyond highway development. Smoother integration of nature and structural reinforcing elements could be achieved thanks to the inclined wall face. For example, in urban land development, there are cases where a combination of natural slopes and retaining structures are required to avoid extensive cut or fill slopes and to preserve the natural slope contour. In this sense, the PFRW can be an attractive alternative to the traditional vertical or terrace walls, since the wall face inclination can be adjusted. Not only that, but the decorative ashlar stone finish could be replaced with a combination of precast concrete blocks and compacted soil installed in a staggered arrangement such that vegetation can grow over the wall. This same idea would apply when concave slopes are used for more natural urban land development, where additional reinforcement is needed at intermediate points on the landscape (e.g., immediately below a roadway). In fact, the mode of deformation associated to PFRW suggests that the reinforcement would not only benefit slopes below (or after) the wall, but also it would prevent any potential development of progressive failure on excavated concave slopes above (or before) wall construction.

In summary, this dissertation examined slopes and the many factors affecting them that are typically overlooked during design and construction. It examined both compaction based slopes like those constructed during mine reclamation, and excavated slopes like the confined slopes with the piling framed retaining wall. The superior performance of concave slope contours in terms of the mechanical and water erosion stability was demonstrated, as well as the interaction between these two stability conditions. This dissertation examined novel ways of addressing some of the many issues that govern the design and sustainability of slopes, and contributed to the continuing development of our understanding of the behavior of soil slopes and to the growing need for environmental friendly design and construction techniques.

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## Vita

Isaac A. Jeldes Halty is a civil engineer whose passion is geotechnical engineering. He graduated with a Bachelor of Science in Civil Engineering from the University of Santiago of Chile and immediately after joined the R&V geotechnical consultants, where he conducted a variety of geotechnical projects for two and half years. During this period of time he started to have fundamental questions about some of the solutions that civil engineers employ in design, and felt the necessity to increase his theoretical understanding. Since August 2008, Isaac has been working in his doctoral degree under the guidance of Dr. Eric Drumm at the University of Tennessee. In the course of his program he earned a Master's Degree in Civil Engineering. Isaac's research has been primarily focused on: a) slope stability and restoration of mine reclaimed sites; b) the design of eco-friendly slopes to reflect natural shapes with reduced erosion rates and optimum mechanical stability; c) numerical modeling and stability analysis of a new type of retaining wall to support the creation of a rational design methodology.