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## Using DelPhi capabilities to mimic protein's conformational reorganization with amino acid specific dielectric constants

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### Abstract

Many molecular events are associated with small or large conformational changes occurring in the corresponding proteins. Modeling such changes is a challenge and requires significant amount of computing time. From point of view of electrostatics, these changes can be viewed as a reorganization of local charges and dipoles in response to the changes of the electrostatic field, if the cause is insertion or deletion of a charged amino acid. Here we report a large scale investigation of modeling the changes of the folding energy due to single mutations involving charged group. This allows the changes of the folding energy to be considered mostly electrostatics in origin and to be calculated with DelPhi assigning residue-specific value of the internal dielectric constant of protein. The predicted energy changes are benchmarked against experimentally measured changes of the folding energy on a set of 257 single mutations. The best fit between experimental values and predicted changes is used to find out the effective value of the internal dielectric constant for each type of amino acid. The predicted folding free energy changes with the optimal, amino acid specific, dielectric constants are within RMSD=0.86 kcal/mol from experimentally measured changes.

### Keywords

DelPhi; protein electrostatics; dielectric constant; Poisson-Boltzmann equation; protein flexibility; energy calculations; single point mutations

### Introduction

Electrostatic interactions play significant role in determining the structure, function and interactions of biomolecules.[1–4] Unlike the short-range force as van der Waals (vdW), electrostatic interactions are long ranging and highly sensitive to the surrounding environment such as solvent characteristics and ions concentration.[3] Being long range force, the electrostatics can affect the energetics and structural properties of distant objects ranging from macromolecular assemblages [5] to cluster of interacting amino acids. [6] Furthermore, the electrostatics may be the dominant energy contribution in describing variety of phenomena as pH and salt dependence of protein stability [7–10] and interactions [11, 12]. The last observation allows pH and salt dependent processes to be effectively studied by modeling the electrostatic component of the energy, while ignoring all other contributions.[10]

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The compiled folding energy changes, the corresponding structures and *in silico* made mutant structures, along with calculated energies are available from <http://compbio.clemson.edu/delphi> and clicking on the link “downloadable data”.

Due to the importance of electrostatic interactions in biomolecular systems, significant efforts were invested to develop methods for modeling the electrostatics and better understanding its contribution to macromolecular properties. Currently there are two distinctive methods for treating the electrostatics in biological macromolecules immersed in water: explicit and implicit methods [13, 14]. The explicit methods describe the water phase explicitly, i.e. as a sea of explicit water molecules [13, 15, 16] and calculate the electrostatic interactions and energies via Coulomb's law. However, these methods demand integration over countless solvent molecules degrees of freedom which requires significant computational time to obtain accurate results. On the other end of the spectrum are so termed implicit solvent models which treat the solvent implicitly as a homogeneous medium with specific dielectric properties and ion concentration [14, 17–19]. The implicit models greatly reduce the requirements of computational time while capturing most of the important electrostatics effects, especially for the bulk water phase [20–22]. However, simple implicit models may fail to accurately describe the electrostatic properties of the system of interest in the regions close to macromolecular surface and in cases of macromolecular systems experiencing conformational changes associated with the process being modeled. Because of that, hybrid methods were also developed treating explicitly the water molecules at the molecular surface while the bulk water is modeled as a continuum medium. [23–26] At the same time, fewer efforts were invested to account for the heterogeneity of macromolecules themselves and the effect of intrinsic flexibility on implicit electrostatic calculations. Our work, in conjunction with previously reported investigations [27, 28], is devoted to offer better representation of inhomogeneous dielectric properties of macromolecules in the framework of continuum electrostatics.

From point of view of continuum electrostatics, the protein molecules are highly polar objects, because they are made of amino acids carrying electrical charges and dipoles [29]. However, the distribution of charged and polar residues is not homogeneous. Polar and charged residues tend to appear on the surface of the molecules while hydrophobic residues are typically located in the core of the corresponding protein [29]. However, this may not hold in case of membrane proteins which frequently have charged and polar residues buried in their core [30]. This indicates that the polarity of biological macromolecules should vary from the core to the molecular surface or in general, will depend on the structure and amino acid composition of the macromolecule. The polarizability, on another hand, is more complicated quantity, directly related to the dielectric “constant” in continuum electrostatics. The polarizability reflects the presence of dipoles (and broadly speaking charges) which are capable of changing their orientations and positions in response to the local electrostatic field [31]. In terms of proteins, the polarizability critically depends on the ability of the corresponding structures to reorganize (backbone and side chain conformational changes) in response to the local electrostatic field.

While understanding that dielectric constant is a macroscopic quantity and cannot be used to describe microscopic systems, we argue that each amino acid can be viewed as space region of protein structure with specific dielectric properties, and thus having specific dielectric constant. [27, 28] Our motivation for such an approach is to investigate the dielectric properties from the view of continuum electrostatics. Each amino acid is considered to be an object made of the corresponding atoms and partial charges and described as a homogeneous medium with particular dielectric constant. The lowest dielectric constant is assigned to amino acids which do not carry much partial charge (hydrophobic residues) and are not very flexible (buried in the hydrophobic core). In contrast, charged and polar amino acids are typically on the surface (being able to make conformational changes) and carry large electrical dipole or even a net charge and therefore will have large dielectric response to the changes of the local electrostatic field.

In standard Poisson-Boltzmann calculations, the protein is typically treated as a homogeneous medium with dielectric constant between 2 (reflecting electronic polarizability [32]) and 4 (accounting for small backbone fluctuations [33]). At the same time, significant body of works exists claiming that the value of dielectric constant is underestimated [34] and that the dielectric constants should be position-dependent [30]. In terms of pKa calculations, it was shown that the use of static structures results in the best fit with experimental data if the dielectric constant is 20. [35] Other works, which were acquired from dipole fluctuation simulations in proteins, state various and even higher dielectric constants from 1–20 [36] up to 30–36 [37]. Also, the investigations of trypsin [38], lysozyme [39] and cytochrome *c* in water [33] have shown that the dielectric constant can vary from 2 in the center of protein to 25 or higher in the polar surface regions.

Alternative approach of accounting for inhomogeneous dielectric properties of proteins was introduced by Zhou and co-workers by using vdW surface instead of solvent exclusion (SE) surface in PB calculations [40–42]. Such an approach results in better solvation of the charges, both on the surface and in the interior of the proteins [43], and effectively raises the internal dielectric constant [44]. This method was applied to investigate the electrostatic contributions to the stability of several proteins. [40] The mutation sites were selected to be surface exposed to reduce the effect of other terms (not electrostatic energy terms) to the changes of the folding free energy. It was shown that the calculations performed with vdW surface provide better match to experimental data compared with those made with SE.

Previous numerical investigations of the local dielectric constant per amino acid type were reported. Karshikoff and co-worker developed an algorithm to estimate the dielectric constant per amino acid type [28]. It was done by assigning pair dipole moment  $\mu$  and polarizability  $\alpha$  (assumed to be fixed at this position) at every point inside of known molecular structure and assuming that all dipoles can freely rotate and their relaxation can be described by a Langevin function in the Debye approximation. It was found that the residue specific dielectric constant can be as large as 58 for Gln residue, while hydrophobic groups' dielectric constant stays with a single digit. In recent work [27], Song used molecular dynamics (MD) simulations to assess the conformational flexibility of three proteins and to deliver amino acid specific polarizabilities. The proteins were modeled as collection of polarizable dipoles in a cavity embedded inside water phase. It was shown that charged and polar residues consistently have large intrinsic polarizability (Table III in [27]), while small side chain and hydrophobic groups tend to have low polarizability. The calculated polarizabilities, however, were much smaller compared to those reported by Karshikoff and co-worker [28], perhaps, because of the use of more realistic modeling (MD) of protein structural fluctuations.

Experimental efforts to evaluate the inhomogeneous dielectric response in proteins were also reported. In series of works, Garcia-Moreno and co-workers used experimental and computational approaches to evaluate the optimal value of the effective dielectric constant in staphylococcus nuclease (SNase) with respect to pKa calculations. It was shown that the effective dielectric constant, as delivered to match the experimental interaction energies with Coulomb formula, should be different for amino acids situated on close distance from those at medium and large distances [45]. Furthermore, at different salt concentrations, different optimal effective dielectric constants were found [46], ranging from 36 up to more than 80.

In this work, by utilizing the capability of Delphi to assigning different dielectric constants inside biological macromolecule, we investigate the possibility of mimicking small structural changes occurring upon single point mutations involving charged amino acid. By calculations of protein folding energy with different dielectric constants for different groups of amino acids, especially polar and ionized amino acids, we obtain the optimal values of

dielectric constant for specific amino acids by optimizing the comparison with the experimental results of the changes of the folding free energy caused by the mutations.

## Methods

### Mutant Protein Database

Mutant proteins subjected to the study and the corresponding folding free energy changes were extracted from *ProTherm* Database[47], (<http://gibk26.bse.kyutech.ac.jp/jouhou/protherm/protherm.html>), which is a collection of experimental data of thermodynamic parameters. For our study, the database was purged to select only cases involving mutation of a charged amino acid (charged wild type to uncharged mutant or uncharged wild type to charged mutant). The experiments listed in *ProTherm* database to measure the thermodynamic parameters conclude Fluorescence spectroscopy, Circular Dichroism, Differential Scanning Calorimetry, NMR, etc. The experimental methods used to measure the unfolding free energy include Thermal, Urea and GdnHCl unfolding at various pH. In our investigation, we collected entries from *ProThem* database which unfolding free energy was measured in the pH range from 7 to 8. This resulted in an initial data set of 37 wild types and 397 mutants involving charged amino acid. After further investigation of structures, some structures were excluded due to large structural defects such as, for instance, large segments of missing polypeptide chains. Thus, our data set was reduced to 28 wild type proteins and 257 mutants (list is provided in supplementary material). The 3D structures of all corresponding wild type proteins were obtained from RCSB Protein Data Bank [48].

To avoid cases with large conformational changes for which the assumption that the electrostatics in the dominant component of the folding energy will not hold, we further purged the initial dataset with respect to the solvent accessible surface area (SASA). For this purpose we used the *msms* program [49] (<http://mgltools.scripps.edu/packages/MSMS/>) which calculates the surface area of each atom. Summing up the individual atomic SASA within an amino acid, we obtain the total SASA per residue and calculate the relative change of SASA ( $f_{ea}$ ) from the free state (isolated amino acid in solution) to the corresponding protein:

$$f_{ea} = \frac{SASA_{isolated\ state} - SASA_{in\ protein}}{SASA_{isolated\ state}} \quad (1)$$

where  $SASA_{in\ protein}$  is the calculated SASA of the residue in protein structure and  $SASA_{isolated\ state}$  is the SASA of the same amino acid extracted from the protein structure (labeled as “isolated state”). In our study we empirically select  $f_{ea} < 0.8$  as a measure for surface exposed amino acids. As mentioned above, this was done to avoid problems with large structural changes which may occur upon a mutation at buried site. Such plausible large structural changes, definitely will involve other, different from electrostatic, energy terms and the changes of the folding free energy cannot be predicted with electrostatics alone.

### Structural Fixing, Protonating and Mutation Operations

In our data set, some protein structures had structural defects such as missing atoms or residues. Because of that, all structures were subjected to a protocol to fix these structural defects. We utilized *profix* program from Jackal package which is developed in Honig’s lab ([http://wiki.c2b2.columbia.edu/honiglab\\_public/index.php/Software:Jackal](http://wiki.c2b2.columbia.edu/honiglab_public/index.php/Software:Jackal)) to rebuild the missing atoms of both backbone and side chains. However, if the protein structure had fifteen missing side chains or more, the structure was removed from the dataset because

such large rebuilding would probably cause significant structural inaccuracies. Hereafter, we utilized *pdbxyz* and *xyzpdb* modules of TINKER software [50] to add missing hydrogen atoms with AMBER [51] force parameter. The single point mutations were done *in silico* with *scap* module from the Jackal package using default options with AMBER force field. Only mutations involving charged amino acids Asp(D), Glu(E), His(H), Lys(K) and Arg(R) were modeled.

### Folding Free Energy Difference Calculations with Delphi

The folding free energy change due to a mutation was calculated as the difference of the folding free energy of wild type and the folding free energy of the mutant:

$$\Delta\Delta G_{\text{folding\_mutant}} = \Delta G_{\text{folding\_wildtype}} - \Delta G_{\text{folding\_mutant}} \quad (2)$$

where  $\Delta G_{\text{folding}}$  is the difference of free energy of folded state and of unfolded state:

$$\Delta G_{\text{folding\_X}} = \Delta G_{\text{folded\_X}} - \Delta G_{\text{unfolded\_X}} \quad (3)$$

where X indicates either wild-type or mutant. Due to our lack of definition of unfolded state structure, we applied seven residue segment model to simulate the free energy of unfolded state (this is described in detail in our previous works [52–54]):

$$\Delta G_{\text{unfolded\_wildtype}} = G_{0(\text{unfolded\_wildtype})} - G_{7(\text{unfolded\_wildtype})} \quad (4)$$

where  $G_{7(\text{unfolding\_wildtype})}$  is free energy of the segments of mutation site with its neighboring six residues in unfolded state, and  $G_{0(\text{unfolded\_wildtype})}$  is the free energy of the rest of the unfolded state. For mutant, we obtain in similar manner:

$$\Delta G_{\text{unfolded\_mutant}} = G_{0(\text{unfolded\_mutant})} - G_{7(\text{unfolded\_mutant})} \quad (5)$$

Since the wild type protein and mutant were assumed to have similar structures in unfolded state except the mutation site with its six neighboring residues, free energy  $G_{0(\text{unfolded\_wildtype})}$  and  $G_{0(\text{unfolded\_mutant})}$  are assumed to be equal to each other. Thus,

$$\Delta\Delta G_{\text{folding\_mutant}} = G_{\text{folded\_wildtype}} - G_{7(\text{unfolded\_wildtype})} - G_{\text{folded\_mutant}} - G_{7(\text{unfolded\_mutant})} \quad (6)$$

For each protein, the wild type and the mutant, four electrostatic energy components were calculated according to Eq. (6). These energy terms were calculated with Delphi program[55, 56] with the following parameters: scale 4 grids/Å, external dielectric constant 80 and the convergence criterion RMSC was 0.0001 kT/e. The internal dielectric constant was varied as described below.

### Dielectric Constants Assignment

Two different scenarios of assigning the internal dielectric constant in Delphi calculations were explored: (a) Two dielectric model of protein. In this simple case, the protein was considered to have only two distinctive dielectric regions, the site of mutation and the rest of the protein. Since in this work the site of mutation is always charged amino acids (either in the wild type or in the mutant), the first region corresponds to a residue carrying a net charge; the second region is the rest of the protein and includes all types of amino acids. Obviously this is simplification since the second region includes charged groups as well. (b) Ten dielectric regions model of protein. Each protein is allowed to have ten different types of dielectric constants (Fig. 1). These local dielectric constants were assigned according to

the following rules: each ionizable group, Asp, Glu, His, Lys and Arg, was considered to have specific dielectric constant which was varied in our analysis from 15 to 28 (low values were ignored based on the analysis performed with two dielectric regions protocol); polar groups, Ser, Thr, Asn and Gln, were also assigned individual dielectric constant throughout the entire protein and the value was varied from 10 to 22, and the rest of amino acid types (including Tyr) were modeled with different dielectric constant which value was varied from 6 to 18. Thus, the entire protein was considered to be kind of mosaic object, each amino acid having specific dielectric constant.

In order to test the sensitivity of the results with respect to different force field parameters, three force fields were applied in our study for comparing and optimizing the results: Amber98 [51], Charmm22 [57] and OPLS [58].

### Computational and experimental results comparison

To seek the optimal combination of dielectric constants resulting in the best fit of the computational results from Delphi calculation and experimental data from *ProTherm* Database[47], we adopted RMSD (Root Mean Square Deviation) measure.

$$RMSD = \sqrt{\sum_{i=1}^n \frac{(\Delta\Delta G_i^{calc} - \Delta\Delta G_i^{exp})^2}{n}} \quad (7)$$

where  $\Delta\Delta G_i^{calc}$  is the calculated change of the folding free energy and  $\Delta\Delta G_i^{exp}$  is the experimental free energy change taken from *ProTherm* database,  $i$  indicates the corresponding mutant and  $n$  is the number of mutants considered in our work. The smallest RMSD indicates the best fit to the experimental data and was used to determine the optimal set of dielectric constants per residue. Below we show the graphs of the best fit between computational and experimental results and report the corresponding optimal dielectric constants which were used to obtain the best fit.

## Results

In this section we report the results of numerical calculations beginning with the data set originally compiled by Zhou and co-workers [40] and then extending our approach to much larger data set.

### (a) Benchmarking against Zhou and co-workers dataset [40]

The idea of assessing the changes of the folding free energy upon charged amino acid substitutions by computing the electrostatic component of the folding energy was first proposed by Zhou and co-workers [40]. For such a purpose they selected mutations which involve ionizable group and mutation sites are exposed to the solvent. This was done to reduce the effect of other (non electrostatic) energy terms to the changes of the folding free energy. The data set included 10 mutants which are K16E, R119E and R154E from T4 Lysozyme (PDB\_ID: 3LZM), D120N and D49N from Human Lysozyme (PDB\_ID: 1REX), D1K, D17K, D25K, E41K and E74K from Ribonuclease Sa (PDB\_ID: 1C54). The calculated by Zhou and co-workers values using vdW and SAS molecular surfaces are shown in Table 1 and in Fig. 2. Our calculations with DelPhi generally agree with these numbers, although small differences were accounted, due to different algorithms implemented in UHBD [40, 59], which originally was used by Zhou and coworkers. In parallel, we employed the two dielectric model (see method section) on the same dataset to calculate the folding energy changes. The best results, with respect to experimental data are shown in Table and drawn in Fig. 2. These best results were obtained with dielectric

constant of the mutation site equal to 21, while the rest of the corresponding protein was treated with dielectric constant of 7.

In Delphi calculations, we used solvent exposed (SE) surface and Amber98 force field, and varied the internal dielectric constants  $\epsilon_1$  from 2 to 30 and  $\epsilon_2$  from 2 to 30, in increments of 1. The results indicate that using two dielectrics model provides much better fit to the experimental data (see the diamonds in Fig. 2) than the single dielectric model with either vdW or SE molecular surfaces. In both cases, the outliers corresponding to over predicted energy charges are significantly improved and are much closer to experimental data when calculated with the two dielectrics model.

### (b) Two dielectrics model tested on the large data set

To further explore plausible advantages of the two dielectrics model, we tested it on much larger data set (see method section). To be consistent with the previous analysis, we used Amber force field parameters and the values of the dielectric constants  $\epsilon_1$  and  $\epsilon_2$  were varied in the same manner as above. The results were grouped into five categories according to the specific ionizable residue being mutated, Asp, Glu, Arg, His and Lys. For each group, the best set of dielectric constants was found and results shown in Fig. 3. For Asp (D) group, the best combination is  $\epsilon_1 = 20$  and  $\epsilon_2 = 5$ , for Glu (E) group,  $\epsilon_1 = 18$  and  $\epsilon_2 = 10$ , for Arg (R) group,  $\epsilon_1 = 28$  and  $\epsilon_2 = 6$ , for His (H) group,  $\epsilon_1 = 18$  and  $\epsilon_2 = 4$  and for Lys (K),  $\epsilon_1 = 25$  and  $\epsilon_2 = 5$ . And the corresponding RMSDs for these five sets are 0.57, 0.61, 0.87, 0.19 and 0.92.

Figure 3 shows that the calculated results are very close to experimental data, which can be seen from the corresponding RMSDs as well. Comparing the individual best values for the dielectric constants, it can be seen that the best  $\epsilon_1$  is always large number varying from 18 to 28, while the best value for  $\epsilon_2$  is much smaller ranging from 4 to 10. This observation will be used in our further analysis.

### (c) Ten dielectrics model

Two dielectrics model is overly simplified and does not account that different amino acid types may have different effective dielectric constants and that this can be attributed to any amino acid in the protein of interest, without being restricted to the mutation site only. To account for this, the ten dielectric constants model was applied for seeking more accurate description of dielectric properties of proteins. Each amino acid with charged side chain as Arg( $\epsilon_1$ ), His( $\epsilon_2$ ), Lys( $\epsilon_3$ ), Asp( $\epsilon_4$ ) and Glu( $\epsilon_5$ ) was considered to have distinctive dielectric constant. Since they are all charged residues and based on our previous analysis discussed above, the ranges for dielectric constants were set from 15 to 25. Another group was formed from amino acids with polar side chain as Ser ( $\epsilon_6$ ), Thr ( $\epsilon_7$ ), Asn ( $\epsilon_8$ ) and Gln ( $\epsilon_9$ ) and each of them was assigned distinctive dielectric constant. The value was varied from 10 to 20. To avoid combinatorial explosion, the rest of the amino acids, including Tyr, were modeled with a single dielectric constant which was ranged from 5 to 18. To assess the sensitivity of the results with regard to force field parameters, the calculations were done with three different force field parameters: Amber, Charmm and OPLS.

Two equally important investigations were pursued: (c1) The possibility of obtaining consistent (universal) values for amino acid specific dielectric constants and (c2) The possibility that such an approach will lower the overall RMSD of predicted changes of the folding free energy.

#### (c1) Obtaining consistent values for individual dielectric constants per amino acid types

Results are summarized in Table 2, where the optimal sets of dielectric constants for mutation involving all amino acids and specific amino acid type are provided as indicated in



column “X”. Two major observations can be made: (1) It can be seen that despite of using different force field parameters, the values are either unchanged or are quite similar. (2) The optimal dielectric constant per amino acid type remains the same or changes very little for different types of mutations. These observations speak in favor of the robustness of the finding and the applied protocol. Thus, the optimal dielectric constants for all types of amino acids modeled in this study are: Arg=18, Lys=20, Glu=19, Thr=12, Asn=13, Gln=13, His=21/22, Asp=17/18, Ser=10/11 and for other amino acids ranges from 14 to 18.

### (c2) Improving RMSD of predicted changes of the folding free energy

Here we report the predicted energy changes obtained with the optimal set of dielectric constants listed in Table 2 and calculated with specific force field parameters. Figure 4 shows the calculated changes of the folding free energy with Amber force field parameters benchmarked against experimental data. It can be seen that most of the data points lie on the main diagonal indicating that the calculated changes are very similar to the experimental results. However, there are several offset points, which are the main contributors to the RMSDs. With the exception of one case in Asp test (mutant D122A in 1RX4.pdb, where the folding free energy change calculated with Delphi is 1.73 Kcal/mol while experimental result is -1.6 Kcal/mol), the most difficult to predict are mutations involving Lys residue (low right panel in Fig. 4). In several cases, the numerical protocol underestimates the experimentally observed energy changes by much. These prominent cases are listed below: The folding energy change for the mutant K97V in 1IOB.pdb is predicted to be 0.17 Kcal/mol while the experimental energy change is 1.5 Kcal/mol resulting in 1.33 Kcal/mol difference. Another prominent case results in 2.38 Kcal/mol difference for 1IOB\_K97G (Delphi result is 0.28 Kcal/mol against experimental value -2.1 Kcal/mol). The third example is K46G mutant in 1ARR.pdb, for which Delphi calculates 0.52 Kcal/mol against the experimental value -1.6 Kcal/mol.

Figure 5 demonstrates the accuracy of the predictions made with Charmm force field parameters. It can be seen again that the most difficult to predict are changes of the folding free energy involving Lys mutation. Prominent examples are: 1STN\_K78A which has 3.3 Kcal/mol offset ( $\Delta\Delta G_{calc} = 2.73$  Kcal/mol,  $\Delta\Delta G_{exp} = -0.6$  Kcal/mol), 1STN\_K78G which has 3.83 Kcal/mol offset ( $\Delta\Delta G_{calc} = 2.72$  Kcal/mol,  $\Delta\Delta G_{exp} = -1.1$  Kcal/mol), 1IOB\_K97G which has 3.04 Kcal/mol offset ( $\Delta\Delta G_{calc} = 0.94$  Kcal/mol,  $\Delta\Delta G_{exp} = -2.1$  Kcal/mol). However, equally difficult to predict with Charmm force field parameters are mutations involving Arg residue as indicated by the next two cases: 1BTA\_R75L and 1STN\_R81G have 2.9 and 2.5 Kcal/mol offset respectively. At the same time, the offset point in case of mutation involving Asp residue calculated with Amber force field parameters (Fig. 4) is now predicted with better accuracy.

The last set of results was obtained with OPLS force field (Fig. 6). In this case, the most difficult to predict are the free energy changes upon mutation of Lys, Asp and Arg groups. Similar to the results with Charmm force field, the mutations K78A and K78G from wild type 1STN have 3.4 Kcal/mol and 3.6 Kcal/mol offset respectively. For the mutations involving Asp and Arg, 1HTI\_S105D ( $\Delta\Delta G_{calc} = 1.24$  Kcal/mol,  $\Delta\Delta G_{exp} = -1.3$  Kcal/mol), 1MSI\_D59N ( $\Delta\Delta G_{calc} = -2.60$  Kcal/mol,  $\Delta\Delta G_{exp} = -0.05$  Kcal/mol), 1STN\_R81G ( $\Delta\Delta G_{calc} = 0.59$  Kcal/mol,  $\Delta\Delta G_{exp} = -2.2$  Kcal/mol) and 1BTA\_R75L ( $\Delta\Delta G_{calc} = -1.07$  Kcal/mol,  $\Delta\Delta G_{exp} = 0.7$  Kcal/mol) have 2.5Kcal/mol, 2.6Kcal/mol, 2.8Kcal/mol and 1.8Kcal/mol offset respectively.

Comparing the list of prominent error cases listed above, it can be seen that the usage of different force field parameters results in different sets of outliers. Very seldom a given mutation is wrongly (more than 2kcal/mol difference from the experimental value) predicted by two different sets of force field parameters. This indicates that different force field

parameters have different issues associated with our approach and perhaps these issues are not directly related to the assumptions made.

Despite of the differences discussed above, the results obtained with ten dielectrics model, are better than those with two dielectrics. The results are summarized in Table 3. As it was indicated above, the results with Amber force are better than both Charmm and OPLS force fields according to the RMSDs comparison. For all dataset case, the RMSD with Amber force fields is 0.86 while the other two are higher than one. For all specific amino acids mutations, the RMSDs of the predictions with Amber force fields are less than 1.0 but for the other force fields are much higher and even to 1.54 with OPLS force fields (Table 3). Compared the RMSDs from two dielectric constants model, the RMSD for Asp improved from 0.57 to 0.31, for Arg improved from 0.87 to 0.64, for Asp, His and Lys also have slight improvements.

## Conclusions

This study demonstrated that small plausible structural rearrangements occurring upon mutations involving charged groups exposed to the water can be successfully modeled with PB formalism, as implemented in DelPhi, by assigning residue specific dielectric constant per amino acid type. The obtained optimal values of the internal dielectric constants for charged and polar groups are higher than those for less polar types of amino acids, in accordance with previous work ([27]). The robustness of the approach is demonstrated by the fact that the same optimal values of the residues specific dielectric constant were obtained with different force field parameters and with different residues involved in the mutation. The resulting RMSD from the benchmarking against experimental data of folding free energy changes are below 1 kcal/mol with Amber force field parameters, which is an excellent achievement from computational stand-point of view.

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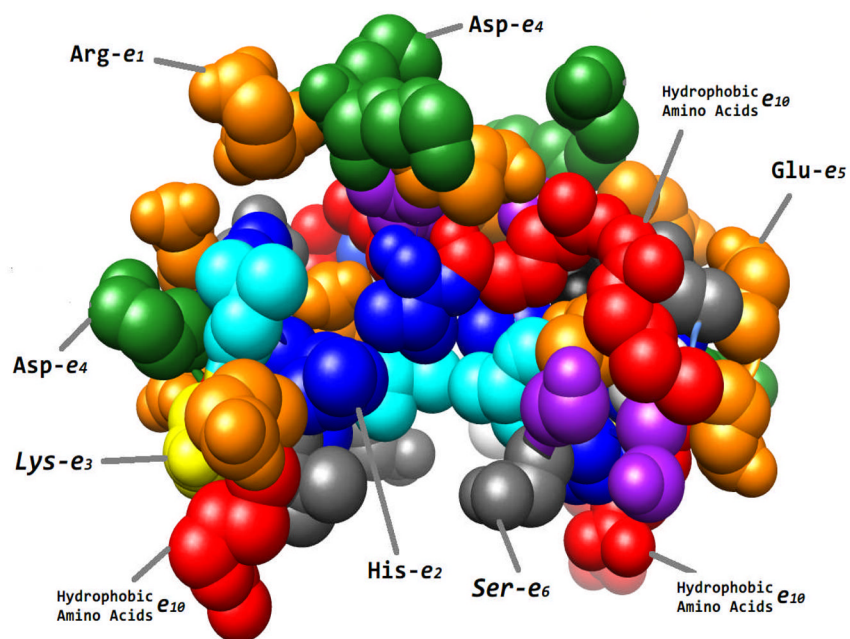
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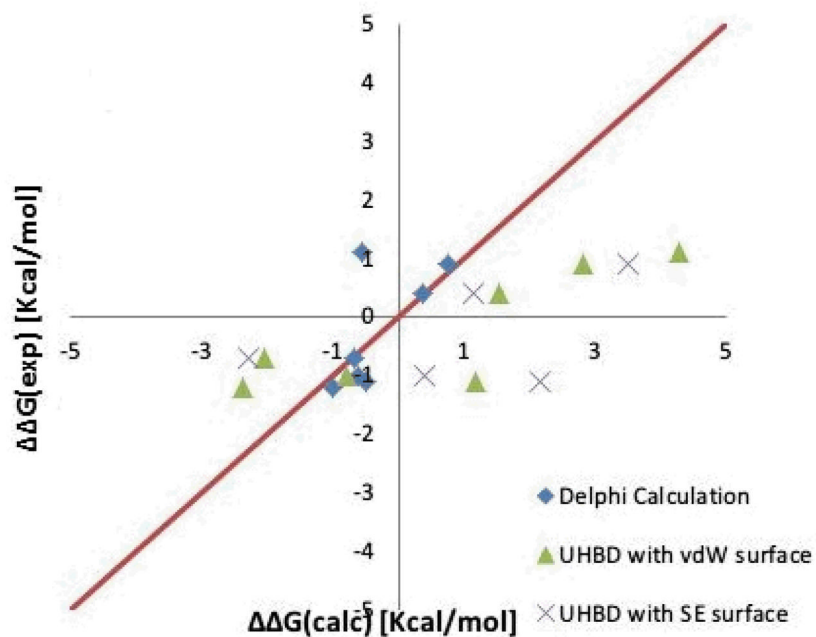
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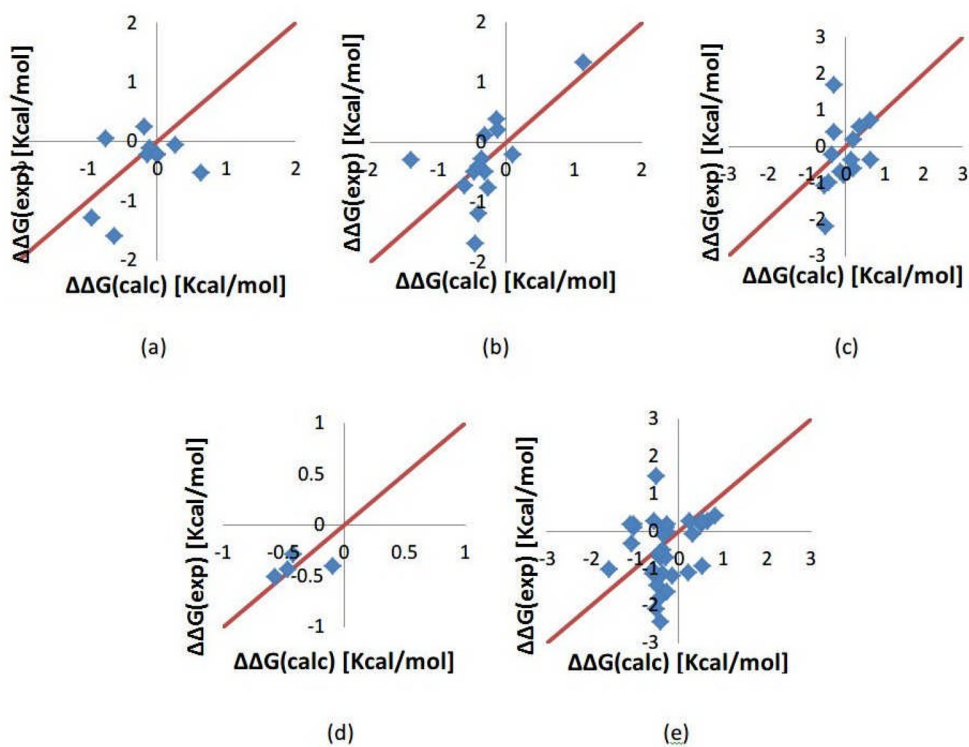
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**Figure 1.** Cartoon presentation of the protocol of assigning different dielectric constants per amino acid type



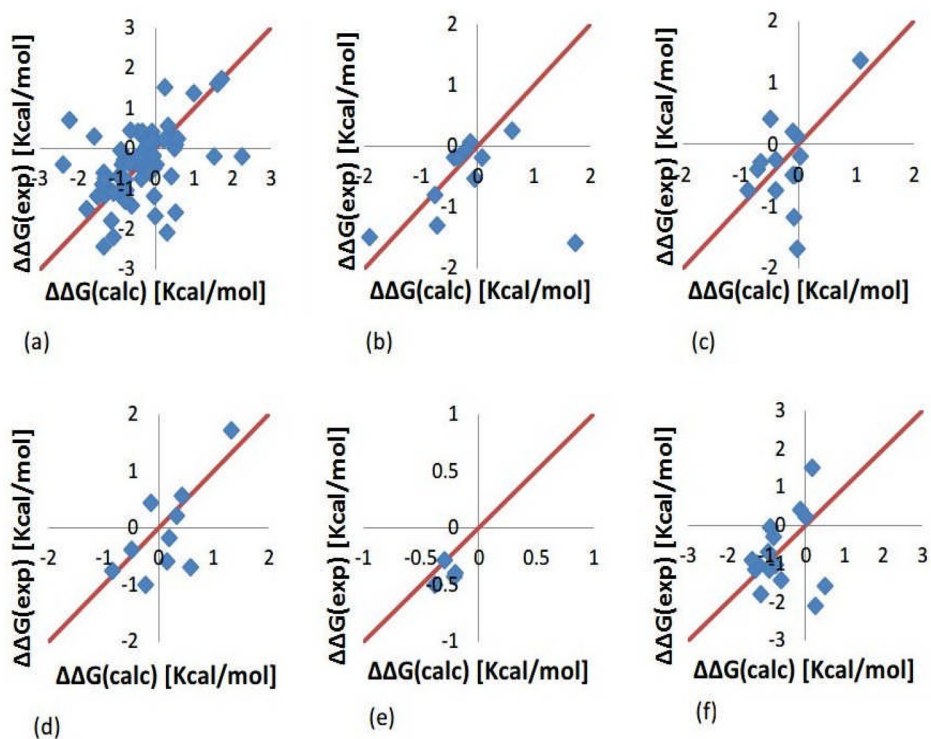
**Figure 2.** Comparison of Delphi results with optimal set of dielectric constants  $\epsilon_1 = 21$   $\epsilon_2 = 7$  (diamond), UHBD[30, 42] results with vdW surface (triangle) and SE (cross) with the experimental results using Zhou and coworkers dataset [40]



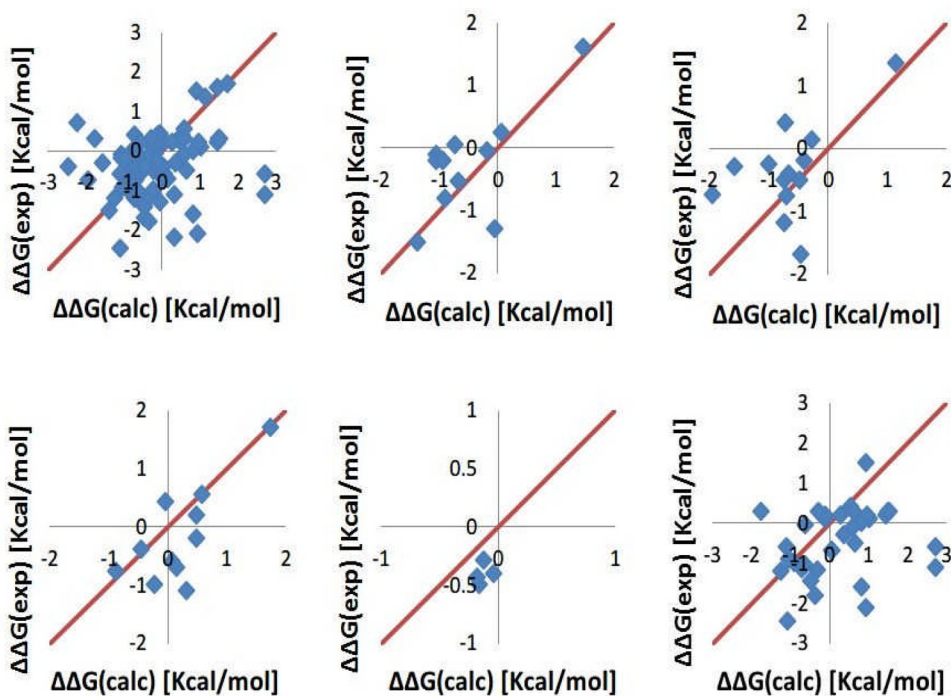
**Figure 3.**

Calculated folding free energy changes compared with the experimental data from *ProTherm* database for five types of mutations. The corresponding optimal dielectric constants are provided as well. (a) Mutation involving Asp residue; optimal dielectric constants  $\epsilon_1 = 20$  and  $\epsilon_2 = 5$  (b) Glu;  $\epsilon_1 = 18$  and  $\epsilon_2 = 10$  (c) Arg;  $\epsilon_1 = 28$  and  $\epsilon_2 = 6$  (d) His;  $\epsilon_1 = 18$  and  $\epsilon_2 = 4$  (e) Lys;  $\epsilon_1 = 25$  and  $\epsilon_2 = 5$ .

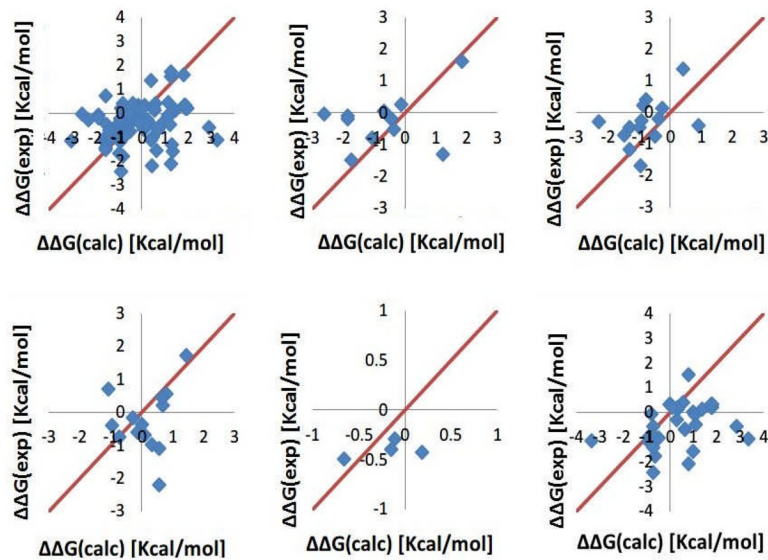




**Figure 4.** The folding free energy changes calculated with Amber force field against experimental data shown with diamonds. (a) All mutations. (b) Mutation involving Asp. (c) Mutation involving Glu. (d) Mutation involving Arg. (e) Mutation involving His. (f) Mutation involving Lys.



**Figure 5.** The folding energy changes calculated with Charmm force field against the experimental data shown with diamonds. All mutations (upper-left), mutation Asp (upper-middle), mutation Glu (upper-right), mutation Arg (lower-left), mutation His (lower-middle), and mutation Lys (lower-right)



**Figure 6.**

The folding energy changes calculated with OPLS force field parameters against the experimental data shown with diamonds. All mutations (upper-left), mutation Asp (upper-middle), mutation Glu (upper-right), mutation Arg (lower-left), mutation His (lower-middle), and mutation Lys (lower-right)

Calculated folding energy changes with three different methods and experimental results for 10 mutations (Zhou and coworkers dataset [40]) involving charged residues. In the third column are the results obtained with DelPhi using SE with probe radius of 1.4Å, the fourth and fifth columns are results taken from Ref [40] with vdW and SE surface, respectively. The last column is the experimental data taken from Ref. [40]. All energies are in [Kcal/mol]

**Table 1**

	Mutations	$\Delta\Delta G(\text{delphi}1=21 \text{ e}2=7)$	$\Delta\Delta G(\text{vdW surface})$	$\Delta\Delta G(\text{SE surface})$	$\Delta\Delta G(\text{exp})$
T4 Lysozyme	K16E	1.94	0.72	1.03	0.5
	R119E	0.09	1.3	2.58	0
	R154E	-1.96	-1.64	-1.79	-1.1
Human Lysozyme	D49N	-0.61	-0.8	0.41	-1
	D120N	-0.69	-2.06	-2.28	-0.7
	D1K	0.35	1.53	1.14	0.4
Ribonuclease Sa	D17K	-0.51	1.18	2.15	-1.1
	D25K	0.75	2.81	3.49	0.9
	E41K	-1.01	-2.37	-7.2	-1.2
	E74K	-0.57	4.28	6.2	1.1

**Table 2**

Optimal sets of dielectric constants reported for the entire dataset and for mutations involving specific amino acid type calculated with Amber, Charm, and OPLSforce fields. The “X” stands for type of mutation, for example X=Asp indicates a mutation involving Asp.

	X	Arg $\epsilon_1$	His $\epsilon_2$	Lys $\epsilon_3$	Asp $\epsilon_4$	Glu $\epsilon_5$	Ser $\epsilon_6$	Thr $\epsilon_7$	Asn $\epsilon_8$	Gln $\epsilon_9$	Other $\epsilon_{10}$
Amber	All	18	21	20	18	19	10	12	13	13	14
	Asp	18	22	20	17	19	11	12	12	13	14
	Glu	18	21	20	17	19	11	12	13	13	14
	Arg	18	22	20	18	20	11	12	13	13	16
	His	18	21	20	18	19	10	12	13	13	18
Charmm	Lys	18	22	20	18	19	10	12	13	14	14
	All	18	22	20	17	19	11	12	13	13	14
	Asp	18	22	20	17	19	11	12	13	13	14
	Glu	18	22	20	17	19	11	12	13	13	14
	Arg	18	22	20	17	19	11	12	13	13	14
OPLS	His	18	21	20	17	19	11	12	13	13	18
	Lys	18	21	20	17	19	11	12	13	13	14
	All	18	21	20	17	19	11	12	13	13	16
	Asp	18	21	20	17	19	11	12	13	13	18
	Glu	18	21	20	18	19	11	12	13	13	16
	Arg	18	21	20	17	19	11	12	13	13	14
	His	18	21	20	17	19	11	12	13	13	18
	Lys	18	22	20	18	19	11	12	13	13	18

**Table 3**

The RMSDs chart for all mutations and specific mutations under Amber, Charmm and OPLS force fields with 10 dielectric constants model

Force Field	All Mutations	Asp (D)	Glu (E)	Arg (R)	His (H)	Lys (K)
Amber	0.86	0.31	0.64	0.63	0.16	0.90
Charmm	1.12	0.65	0.73	1.31	0.35	1.36
OPLS	1.35	1.28	0.91	1.11	0.45	1.54