

# The inverse problem of variational calculus and the problem of mixed endpoint conditions

Pedro Gonçalves Henriques

**Abstract.** P. A. Griffiths established the so-called mixed endpoint conditions for variational problems with non-holonomic constraints. We will present some results in this context and discuss the inverse problem of calculus of variations.

*Keywords:* Inverse problem of calculus of variations.

## 1. Introduction

The study of Calculus of Variations for multiple integrals was first developed by Caratheodory [1929], while Weil-De Donder [1936], [1935] advanced a different theory later. The two approaches were unified by Lepage [1936-1942], Dedecker [1953-1977] and Liesen [1967] in a framework using the  $n$ -Grassmannian manifold of a  $C^\infty$  manifold. Important contributions in the Calculus of Variations on smooth manifolds were made by R. Hermann [1966], H. Goldschmidt and S. Sternberg [1973] with their Hamilton-Cartan formalism, as well as by Ouzilou [1972], D. Krupka [1970-1975] and I. M. Anderson [1980]. The symplectic approach of P. L. Garcia and A. Pérez-Rendón [1969-1978], the multisymplectic version of Kijowski and Tulczyjew [1979] based on the theory of Dedecker, the polysymplectic approach of C. Günther [1987], Edelen [1961] and Rund [1966] are also important references in this field. Here we deal with the broader problem of finding extrema of a functional on a set of  $n$ -dimensional integral manifolds of a Pfaffian differential system.

In 1983, Griffiths proposed a new approach to variational problems based on techniques from the theory of exterior differential systems. His work dealt with the problem of finding extrema for a functional  $\phi$  defined on the set of one-dimensional integral manifolds of a differential system  $(I^*, L^*)$ .

---

Supported by the Center for Mathematical Analysis, Geometry, and Dynamical Systems.

This approach was established using intrinsic entities. In this work we present a general setting based on [25] (sections 2 to 8), and we deal with the inverse problem (section 9).

In 1887 Helmholtz addressed the following problem: given

$$P_i = P_i(x, u^j, u_x^j, u_{xx}^j),$$

is there a Lagrangian  $L(x, u^j, u_x^j)$  such that

$$E_i(L) = \partial L / \partial u^i - D_x \partial L / \partial u_x^i = P_i$$

where

$$D_x = \partial / \partial x + u_x^i \partial / \partial u^i + u_{xx}^i \partial / \partial u_x^i?$$

He found necessary conditions for  $P_i$  to form an Euler-Lagrange system of equations (see (9.1), (9.2) and (9.3)). Some years later, these conditions were proved to be locally sufficient. I. M. Anderson [1992], [1980], P. J. Olver [1986], F. Takens [1979], W. M. Tulczyjew [1980] and A. M. Vinogradov [1984] generalized Helmholtz's conditions for both higher order systems of partial differential equations and multiple integrals.

## 2. Integral manifolds of a differential system and valued differential systems

We assume that a Pfaffian differential system  $(I^*, L^*)$  is given on a real-manifold  $X$  by:

- i) a subbundle  $I^* \subset T^*X$ ,
- ii) another subbundle  $L^* \subset T^*X$  with  $I^* \subset L^* \subset T^*X$ ,

such that the rank  $(L^*/I^*) = n$  (with  $n$  being a natural number).

An integral manifold of  $(I^*, L^*)$  is given by an oriented connected compact  $n$ -dimensional smooth manifold  $N$  (possibly with a piecewise smooth boundary  $\partial N$ ) together with a smooth mapping

$$f : N \rightarrow X$$

satisfying

$$I_{f(x)}^* \perp = L_{f(x)}^* \perp + f_*(TN), \quad (2.1)$$

for all  $x \in N$ , where  $f_* : T_x N \rightarrow T_{f(x)} X$  is the differential of  $f$  at  $x$ .

We denote by  $V(I^*, L^*)$  the collection of integral manifolds  $f$  of  $(I^*, L^*)$ .

A valued differential system is a triple  $(I^*, L^*, \varphi)$ , where  $(I^*, L^*)$  is a Pfaffian differential system and  $\varphi$  is an  $n$ -form on  $X$ .

We define the functional  $\phi$  associated with  $(I^*, L^*, \varphi)$  in  $V(I^*, L^*)$  by:

$$\phi : V(I^*, L^*) \rightarrow R,$$

$$f \rightarrow \phi[f] = \int f^* \varphi. \tag{2.2}$$

### 3. Local embeddability

The following definition is a general setting for the study of problems in the Calculus of Variations. In [25] we proved that there exist locally defined mappings that induce  $(I^*, L^*)$  from the canonical system in  $J^1(R^n, R^s)$  possibly with some constraints, establishing a local coorespondence between these differential systems. Let us assume that  $d(C^\infty(X, L^*)) \subset C^\infty(X, L^* \wedge T^*X)$ , and let  $d' = \dim X$ ;  $s = \text{rank} I^*$  ( $d(C^\infty(X, L^*))$  is the set of images produced by the exterior derivative of  $C^\infty(X, L^*)$ ). Using the Frobenius theorem, we can set for every  $p \in X$  a chart coordinate system  $\{u^1, \dots, u^{s+n}, v^1, \dots, v^{d'-s-n}\}$  so that

i) 
$$L^* = \text{span}\{du^\alpha | 1 \leq \alpha \leq s + n\}, \tag{3.1}$$

ii) 
$$L^{*\perp} = \text{span}\left\{\frac{\partial}{\partial v^i} | 1 \leq i \leq d' - s - n\right\} \tag{3.2}$$

for an open subset  $U$  of  $X$  with  $p \in U$ .

**Definition 3.1.** Let  $(I^*, L^*)$  be a Pfaffian differential system with

$$d(C^\infty(X, L)) \subset C^\infty(X, L^* \wedge T^*X).$$

We say that  $(I^*, L^*)$  is locally embeddable if for every  $p \in X$  there exist an open neighborhood  $U$  of  $p$  and local coframes

$$CF = \{\theta_1, \dots, \theta_s\} \tag{3.3}$$

for  $I^*$  and

$$CF' = \{\theta_1, \dots, \theta_s, du''^s + 1, du''^s + n\} \tag{3.4}$$

for  $L^*_U$ , satisfying the following conditions:

(i) 
$$\delta(I^*_U \wedge \Omega) \subset T^* \wedge \Lambda^n(L^*_U) / (T^*U \wedge I^*_U \wedge \Lambda^{n-1}(L^*)) \tag{3.5}$$

(ii)  $\text{Ker } \delta$  is a constant rank subbundle of  $I^* \wedge \Omega$ ,

where  $\Omega = \text{span}\{du''^{s+1} \wedge \dots \wedge \widehat{du''^{s+\beta}} \wedge \dots \wedge du''^{s+n}\}$ ;  $\widehat{du''^{s+\beta}}$  -means deletion of the  $s + b$  factor (for  $n = 1, \widehat{du''^{s+1}} = 1$ ). We use  $u''$  since we may have to reorder these coordinates.

The map  $\delta : I^* \wedge \Omega \rightarrow \Lambda^{n+1}(T^*U)/I_u^* \wedge (\Lambda^n(T^*U))$  is induced by

$$d : C^\infty(U, I^* \wedge \Omega) \rightarrow C^\infty(U, \Lambda^{n+1}(T^*U))$$

on  $I^* \wedge \Omega$ .

This definition means that if  $I^*$  has no Cauchy characteristics, the structure equations are locally:

$$d\theta^i \equiv \pi_j^i \wedge du''^{s+j} + A_{i'\alpha}^{ij'} \pi_{j'}^{i'} \wedge \theta^\alpha + B_{\alpha\beta}^i \theta^\alpha \wedge du''^{s+\beta} \text{ mod } I \wedge I \quad (3.6)$$

$$1 \leq i, i', \alpha \leq s, 1 \leq j, j', \beta \leq n, I = C^\infty(X, I^*).$$

#### 4. The Cartan system of $\Psi$

Let  $(I^*, L^*, \varphi)$  be a valued differential system on  $X$ , and  $W$  be the total space of  $I^*$ . Let  $\chi$  be the canonical form on  $T^*X$ , and  $i$  the inclusion map  $W \xrightarrow{i} T^*X$ .

Let us assume that there exists a local  $n$ -form  $\omega$  inducing a nonzero section of  $\Lambda^n(L^*/I^*)$  and has the following form:

$$\omega = \omega^1 \wedge \dots \wedge \omega^n. \quad (4.1)$$

We define:

$$\omega_i = (-1)^{i-1} \omega^1 \wedge \dots \wedge \widehat{\omega^i} \wedge \dots \wedge \omega^n. \quad (4.2)$$

Let  $W^n$  be the  $n$ -Cartesian power of  $W$ , and  $Z$  be a subset of  $W^n$  defined by  $Z = \{z \in W^n : \pi'(z) \in \Delta X^n\}$ , where  $\pi'$  is the natural projection  $\pi' : W^n \rightarrow X^n$ , and  $\Delta X^n$  is the diagonal submanifold of  $X^n$ . The subset  $Z$  is a vector subbundle over  $X$  and  $\dim Z = d + sn$ . We define

$$\Psi = d\psi \quad (4.3)$$

where  $\psi$  is given by

$$\psi = \pi^* \varphi + (\pi^j \circ i')^* [i^*(\chi)] \wedge \pi^* \omega_j. \quad (4.4)$$

$\pi^j$  is the natural projection into the  $j^{\text{th}}$  component  $\pi^j : W^n \rightarrow W$ ,  $i'$  is the inclusion map  $Z \rightarrow W^n$  and  $\pi$  is the natural projection  $\pi : Z \rightarrow X$ .

**Definition 4.1.** Given the  $n+1$ -form  $\Psi$ , the Cartan system  $C(\Psi)$  is the ideal generated by the set of  $n$ -forms

$$\{v \lrcorner \Psi \text{ where } v \in C^\infty(Z, TZ)\}.$$

An integral manifold of  $(C(\Psi), \omega)$  is given by an oriented connected compact  $n$ -dimensional smooth manifold  $N$  (possibly with a piecewise smooth boundary  $\partial N$ ) together with a smooth mapping

$$f : N \rightarrow X$$

satisfying:

$$f^*\theta = 0 \quad \text{for every } \theta \in C(\Psi) \tag{4.5}$$

and

$$f^*(\omega) \neq 0. \tag{4.6}$$

A solution of  $(C(\Psi), \omega)$  projected in  $X$  will give an extremum of  $\phi$ .

**5. The momentum space, prolongation of  $(C(\Psi), \pi^*\omega)$  in the momentum space, non-degeneracy**

The momentum space is constructed in the following way. Suppose we are given on  $Z$  (see section 4):

- (i) a closed  $(n + 1)$ -form  $\Psi$  with the associated Cartan system  $C(\Psi)$ ,
- (ii)  $\pi'^*$  the pull back to  $Z$  of the  $\omega$   $n$ -form which induces a nonzero section on  $\Lambda^n(L^*/I^*)$ .

Integral elements of  $(C(\Psi), \pi'^*\omega)$  are defined in a similar way as the integral elements of  $(I^*, L^*)$ . The set of integral elements  $[x_0, E_0^n]$  gives a subset

$$V_n(C(\Psi), \pi'^*\omega) \subset G_n(Z) \quad (G_n(Z) \text{ is the } n\text{-Grassmanian}).$$

Denoting by  $\pi''$  the projection  $G_n(Z) \rightarrow Z$  and assuming regularity at each step, one inductively defines:

$$\begin{aligned} Z_1 &= \pi''(V_n(C(\Psi), \pi'^*\omega), V'_n(C(\Psi), \pi'^*\omega)) = \\ &= \{E \in V_n(C(\Psi), \pi'^*\omega) : E \text{ tangent to } Z_1\}, \end{aligned} \tag{5.1}$$

$$\begin{aligned} Z_2 &= \pi''(V'_n(C(\Psi), \pi'^*\omega), V''_n(C(\Psi), \pi'^*\omega)) = \\ &= \{E \in V'_n(C(\Psi), \pi'^*\omega) : E \text{ tangent to } Z_2\}. \end{aligned} \tag{5.2}$$

**Definition 5.1.** *Suppose  $(I^*, L^*, \varphi)$  is a valued differential system, with  $(I^*, L^*)$  being a locally embeddable differential system and  $\omega = \omega^1 \wedge \dots \wedge \omega^n$ . If there exists a  $k_0 \in N$  such that  $Z_{k_0} = Z_{k_0+1} = \dots = Z_{k_0+n'}$  ( $n' \in N$ ) in the above construction, with*

- (i)  $Z_{k_0}$  being a manifold of dimension  $(n + 1)m + n$  for  $m \in N$ , and
- (ii)  $(C(\Psi), \pi'^*\omega)_{Z_{k_0}}$  being a differential system in  $Z_{k_0}$  with  $r_n = 0$  (Cartan number in Cartan-Kähler Theorem) for all  $V_{n-1}(C(\Psi), \pi'^*\omega)$ ; (for  $n = 1$  we follow [23] and replace this condition by  $\psi \wedge \Psi^n \neq 0$  on  $Z_{k_0}$ ).

Then  $(I^*, L^*, \varphi)$  is a non-degenerate valued differential system, and  $Z = Y$  is called the momentum space.

We call  $(C(\Psi), \pi^*\omega)_Y$  the prolongation of  $(C(\Psi), \pi^*\omega)$  in the momentum space. By construction, the differential system  $(C(\Psi), \pi^*\omega)_Y$  satisfies:

- (i) the projection  $(C(\Psi), \pi^*\omega) \rightarrow Y$  is surjective,
- (ii) the integral manifolds of  $(C(\Psi), \pi^*\omega)$  on  $Z$  coincide with those of  $(C(\Psi), \pi^*\omega)$  on  $Y$ .

## 6. Well-posed valued differential systems

**Definition 6.1.**  $(I^*, L^*, \varphi, P^*, M^*)$  is a well-posed valued differential system, if the following conditions are satisfied:

- (i)  $(I^*, L^*, \varphi)$  is a non-degenerate valued differential system (with  $\dim Y = (n+1)m + n$ ) and  $\varphi = L\omega$  for a smooth function  $L$  on  $X$ ;
- (ii) there exists a subbundle  $P^*$  of  $I^*$  of rank  $m$  and a subbundle  $M^*$  of  $L^*$  of rank  $m+n$ , such that:
  - (a) 
$$\begin{array}{ccc} I^* & \subset & L^* \subset T^*X \\ \cup & & \cup \\ P^* & \subset & M^* \end{array}$$
  - (b) the locally given  $n$ -form  $\omega$  also induces a nonzero section on  $\Lambda^n(M^*/P^*)$ ,
  - (c)  $Y \subset (P^*)^n|_{\Delta X^n}$ , with  $Y$  a subbundle of  $(P^*)^n|_{\Delta X^n}$ ,
- (iii)  $\pi^{**}M^* = \text{span}\{\pi^*\theta | \theta \in C^\infty(X, M^*)\}$  is completely integrable on  $Y$ , where  $\pi^{**} = \pi \circ i$ . As before  $i$  denotes the inclusion mapping  $Y \rightarrow Z$  and  $\pi$  the projection  $Z \rightarrow X$ .

Let us assume that there exists a coframe  $CF = \{\theta^\alpha, du^{s+j}, \pi_j^{i'}, \pi_j^{i''} | 1 \leq \alpha \leq s, 1 \leq i' \leq s_l, j' \in L_{i'}, s_{l+1} \leq i'' \leq s, 1 \leq j \leq n\}$  for  $T^*X$  with  $L_{i'} \subset \{k \in N, 1 \leq k \leq n\}$  such that

$$(i) \quad I^* = \text{span}\{\theta^\alpha | 1 \leq \alpha \leq s\}; \quad (6.1)$$

$$(ii) \quad L^* = \text{span}\{\theta^\alpha, du^{s+j} | 1 \leq \alpha \leq s, 1 \leq j \leq n\}; \quad (6.2)$$

$$(iii) \quad T^*X = L^* \oplus R^* \quad (\oplus \text{ denotes a direct sum}) \text{ with } R^* = \text{span}\{\pi_j^{i'}, \pi_j^{i''} | 1 \leq i' \leq s_l, j' \in L_{i'}, s_{l+1} \leq i'' \leq s, 1 \leq j \leq n\};$$

$$(iv) \quad d\theta_j^{i''} \equiv 0 \pmod{I}, \text{ for } j'' \notin L_{i'}; \quad (6.3)$$

$$(v) \quad d\theta_j^{i'} \equiv \pi_j^{i'} \wedge \omega \pmod{I}, \text{ for } j' \in L_{i'}; \quad (6.4)$$

$$(vi) \quad d\theta_j^{i''} \equiv \pi_j^{i''} \wedge \omega \pmod{I}, \text{ when } 1 \leq j \leq n; \quad (6.5)$$

(vii)  $\pi_{j'}^{i'}, \pi_j^{i''}$  are linearly independent mod  $L$ .

We define  $\theta_j^\alpha \doteq \theta^\alpha \wedge \omega_j$ .

Let  $d\varphi \equiv L_{i''}^j \wedge \pi_j^{i''} + L_{i'}^j \wedge \pi_{j'}^{i'}$  mod  $I$  and  $dL_\nu^\alpha \equiv L_{\nu\nu'}^{\alpha\alpha'} \pi_{\alpha'}^{\nu'}$  mod  $\pi L^*$   $1 \leq \alpha, \alpha' \leq s$   $\nu \in L_\alpha$  and  $\nu' \in L_{\alpha'}$ .

**Quadratic form  $A$ :** Let  $(I^*, L^*, \varphi, P^*, M^*)$  be a well-posed valued differential system and  $A$  be a quadratic form defined in  $T^*X$  given by  $A(v, w) = L_{\nu\nu'}^{\alpha\alpha'} v_\nu^\alpha w_{\nu'}^{\alpha'}$ , where  $v = v_{\theta^\alpha} \partial / \partial \theta^\alpha + v_{\pi_\alpha^\nu} \partial / \partial \pi_\alpha^\nu$  and  $w = w_{\theta^\alpha} \partial / \partial \theta^\alpha + w_{\pi_\alpha^\nu} \partial / \partial \pi_\alpha^\nu$ . This quadratic form plays an important role in establishing necessary conditions for a local extremum.

**6.1. Generalized Lagrange Problem.** Let us describe the following problem:

**Generalized Lagrange Problem.** Let  $X = J^1(R^n, R^m)$  (the 1 jet manifold), with the canonical system  $I^*$  defined on  $X$  (i.e.  $I^* = \text{span}\{\theta^\alpha = dy^\alpha - y_{x^i}^\alpha dx^i\}$ ). Let  $\varphi = L\omega$  with  $\omega = dx^1 \wedge \dots \wedge dx^n$ . We choose  $x^1, \dots, x^n$  to be coordinates for  $R^n$ , and  $y^1, \dots, y^m$  to be coordinates for  $R^m$ .

We proved in [26] that a Lagrange problem for  $n = 1$  with  $L \det L_{\nu\nu'}^{\alpha\alpha'} \neq 0$ , and with constraints not involving more than one variable  $y$  in each equation of restriction is a well posed valued differential system.

## 7. The Euler-Lagrange differential system for a well-posed valued differential system

When we compute the first variation of  $\phi$ , we find an integral over  $N$  and another over the boundary  $\partial N$ . The volume integral will vanish for projections of integral manifolds of the Cartan system  $(C(\Psi), \pi^*\omega)$  into  $X$ . Choosing suitably the set of boundary conditions we can make the integral over the boundary to vanish as well, providing stationary integral manifolds for generalized Lagrange problems (see [25]).

### 7.1. The Euler-Lagrange differential system.

**Definition 7.1.** Let  $(I^*, L^*, \varphi)$  be a valued differential system. The Cartan system  $(C(\Psi), \pi^*\omega)$  is called the Euler-Lagrange differential system associated with  $(I^*, L^*, \varphi)$ .

Assuming that  $(I^*, L^*, \varphi)$  is non-degenerate, we now consider the restriction to  $Y$  of the Euler-Lagrange differential system associated with  $(I^*, L^*, \varphi)$ . The following proposition is easy to prove (see [25]):

**Proposition 7.1.** If  $g$  is an integral manifold of  $(C(\Psi), \pi^*\omega)$ , then  $\pi \circ g \in V(I^*, L^*)$ , where  $\pi$  is the natural projection  $\pi : Z \rightarrow X$ .

We denote by  $(V(C(\Psi), \pi^*\omega))$  the set of integral manifolds of  $(C(\Psi), \pi^*\omega)$ .

## 8. Examples

**Example 1.** *Strings* [41], [42]

Let  $X = J^1(N, R^m)$ ,  $N$  being a two-dimensional manifold. In this case  $I^* = \text{span}\{dx^\alpha - x'^\alpha d\sigma - \dot{x}^\alpha d\tau \mid 0 \leq \alpha \leq m-1, x^\alpha\}$  are coordinates in  $R^m$ , and  $\sigma, \tau$  are coordinates of  $N$ ,  $x'^\alpha = \frac{\partial x^\alpha}{\partial \sigma}$ ,  $\dot{x}^\alpha = \frac{\partial x^\alpha}{\partial \tau}$ . In  $R^m$  we take a metric defined in  $TR^m$  by  $g^{00} = -g^{11} = 1, 1 \leq i \leq m$  and  $g^{ij} = 0$  for  $i \neq j$ . The set  $X$  is given by:  $X = \{x \in X_0 \mid (\dot{x} \cdot \dot{x}) \geq 0 \text{ and } (x' \cdot x') \leq 0\}$  (where  $(\cdot)$  denotes the inner product with respect to the metric  $g$ ). The form  $\omega$  is  $\omega = d\sigma \wedge d\tau$ . We have

$$\varphi = L\omega = [(x' \cdot \dot{x})^2 - (\dot{x} \cdot \dot{x})(x' \cdot x')]^{1/2} d\sigma \wedge d\tau. \quad (8.1)$$

Note:  $L$  is a function of  $\dot{x}$  and  $x'$  only.

**First variation of  $\phi$ .** Let  $\phi = \int f^*(\varphi)$ , where  $f \in V(I^*, L^*)$ . Then

$$\delta\phi = \int f^*(v \lrcorner d\varphi + d(v \lrcorner \varphi)), \quad (8.2)$$

where  $v(\sigma, \tau) = F_*(\partial/\partial t)(t, \sigma, \tau)|_{t=0}, (\sigma, \tau) \in N, t \in [0, 1]$  and  $F$  is the one parameter variation of  $f$  i.e.  $F(t, \sigma, \tau)|_{t=t_1} \in V(I^*, L^*)$  for all  $0 \leq t_1 \leq 1$ . Hence the Lie derivative of  $dx^\alpha - x'^\alpha d\sigma - \dot{x}^\alpha d\tau$  by  $v$  along  $f(N)$  vanishes,  $(d(v \lrcorner (dx^\alpha - x'^\alpha d\sigma - \dot{x}^\alpha d\tau)) + (v \lrcorner (-dx'^\alpha \wedge d\sigma - d\dot{x}^\alpha \wedge d\tau)))|_{f(N)} = 0$ .

The form  $\Psi_Z$  is given by

$$\begin{aligned} \Psi_Z = & (L_{\dot{x}^\alpha} - \dot{\lambda}_\alpha) \pi^*(d\dot{x}^\alpha \wedge \omega) + (L_{x'^\alpha} - \lambda'_\alpha) \pi^*(dx'^\alpha \wedge \omega) + \\ & (d\dot{\lambda}_\alpha \wedge \pi^* d\sigma - d\lambda'_\alpha \wedge \pi^* d\tau) \wedge \pi^* dx^\alpha + (-\dot{x}^\alpha d\dot{\lambda}_\alpha - x'^\alpha d\lambda'_\alpha) \wedge \pi^* \omega \end{aligned} \quad (8.3)$$

The Cartan system in  $Z$  is:

$$(i) \quad \partial/\partial \dot{\lambda}_\alpha \lrcorner \Psi_Z = -\pi^*((dx^\alpha - \dot{x}^\alpha d\tau) \wedge \pi^* d\sigma) = 0, \quad (8.4)$$

$$(ii) \quad \partial/\partial \lambda'_\alpha \lrcorner \Psi_Z = -\pi^*((dx^\alpha - x'^\alpha d\tau) \wedge \pi^* d\sigma) = 0, \quad (8.5)$$

$$(iii) \quad \partial/\partial \dot{x}^\alpha \lrcorner \Psi_Z = -\pi^*(L_{\dot{x}^\alpha} - \dot{\lambda}_\alpha) \omega = 0, \quad (8.6)$$

$$(iv) \quad \partial/\partial x'^\alpha \lrcorner \Psi_Z = -\pi^*(L_{x'^\alpha} - \lambda'_\alpha) \omega = 0, \quad (8.7)$$

$$(v) \quad \partial/\partial x^\alpha \lrcorner \Psi_Z = -\pi^*(d\dot{\lambda}_\alpha \wedge \pi^* d\sigma - d\lambda'_\alpha \wedge \pi^* d\tau) = 0. \quad (8.8)$$



Hence

$$Z_1 = Z|L_{\dot{x}^\alpha - \dot{\lambda}_\alpha}, L_{x'^\alpha - \lambda'_\alpha}. \tag{8.9}$$

Note that from (i) and (ii) we have  $\theta^\alpha = 0$ ;

from (iii), (iv) and (v) we have  $E[L]\omega = (\partial L/\partial x^\alpha - D_\sigma \partial L/\partial x'^\alpha - D_\tau \partial L/\partial \dot{x}^\alpha)\omega = 0$  for  $D_\tau = \partial/\partial \tau + \dot{x}^\alpha \partial/\partial x^\alpha + \ddot{x}^\alpha \partial/\partial \dot{x}^\alpha$  and  $D_\sigma = \partial/\partial \sigma + x'^\alpha \partial/\partial x^\alpha + x''^\alpha \partial/\partial x'^\alpha$ .

The generalized momenta are given by

$$\dot{\lambda}_\alpha = \frac{x'^\alpha(x' \cdot \dot{x}) - (x' \cdot x')\dot{x}^\alpha}{[(x' \cdot \dot{x})^2 - (\dot{x} \cdot \dot{x})(x' \cdot x')]^{1/2}}, \tag{8.10}$$

$$\lambda'_\alpha = \frac{\dot{x}^\alpha(x' \cdot \dot{x}) - (\dot{x} \cdot \dot{x})x'^\alpha}{[(x' \cdot \dot{x})^2 - (\dot{x} \cdot \dot{x})(x' \cdot x')]^{1/2}}. \tag{8.11}$$

Let  $R^{2m}|(\dot{x} \cdot \dot{x}) \geq 0, (x' \cdot x') \leq 0 \xrightarrow{F'} R^{2m}$  be given by

$$F'(\dot{x}^\alpha, x'^\alpha) = (\lambda'_\alpha(\dot{x}^\alpha, x'^\alpha), \dot{\lambda}_\alpha(\dot{x}^\alpha, x'^\alpha)).$$

In this case  $F'$  has an inverse in  $R^{2m}|(\dot{x} \cdot \dot{x}) \geq 0, (x' \cdot x') \leq 0$  and  $F'^{-1}$  is given by:

$$\dot{x}^\alpha = \frac{\lambda'_\alpha(\lambda' \cdot \dot{\lambda}) - (\lambda' \cdot \lambda')\dot{\lambda}_\alpha}{[(\lambda' \cdot \dot{\lambda})^2 - (\dot{\lambda} \cdot \dot{\lambda})(\lambda' \cdot \lambda')]^{1/2}}, \tag{8.12}$$

$$x'^\alpha = \frac{\dot{\lambda}_\alpha(\lambda' \cdot \dot{\lambda}) - (\dot{\lambda} \cdot \dot{\lambda})\lambda'_\alpha}{[(\lambda' \cdot \dot{\lambda})^2 - (\dot{\lambda} \cdot \dot{\lambda})(\lambda' \cdot \lambda')]^{1/2}}. \tag{8.13}$$

The Cartan system in  $Z'_1 = Z_1|(\dot{\lambda} \cdot \dot{\lambda}) \geq 0, (\lambda' \cdot \lambda') \leq 0$  is given by (i), (ii), (iv) and (v) of the Cartan system in  $Z$ . Let  $Y = Z'_1$ . The prolongation of  $(C(\Psi), \pi^*\omega)$  ends at  $Z'_1$ . The dimension of  $Y$  is  $\dim Y = 3m+2$ . Every point in  $Y$  is a zero-dimensional integral element of  $(C(\Psi), \pi^*\omega)$ , and  $r_1 = 2m+1$ . The Cartan system is in involution at  $x$  if  $\det C(v)|_{X_0} \neq 0$ , and

$$C(v) = \begin{bmatrix} \langle v, d\tau \rangle I & \langle v, d\sigma \rangle I \\ m \times m & m \times m \\ A & B \\ m \times m & m \times m \end{bmatrix} \tag{8.14}$$

for every  $v \neq 0$  along  $E^1$ , with  $[x_0, E^1]$  being any integral element of  $(C(\Psi), \pi^*\omega)$ , where

$$A = \langle v, d\sigma \rangle L_{\dot{x}^\alpha \dot{x}^\beta} - \langle v, d\tau \rangle L_{x'^\alpha x'^\beta} \tag{8.15}$$

and

$$B = \langle v, d\sigma \rangle L_{\dot{x}^\alpha x'^\beta} - \langle v, d\tau \rangle L_{x'^\alpha x'^\beta}, \text{ with } 0 \leq \beta \leq m-1. \tag{8.16}$$

Let us define the energy momentum current  $P = (P^0, \dots, P^{m-1})$  on the surface  $\gamma = \{x^\alpha(\sigma, \tau), \sigma, \tau | 0 \leq \alpha \leq m - 1\}$  by

$$P^\alpha = \int \dot{P}^\alpha d\tau + P'^\alpha d\sigma \tag{8.17}$$

where  $\dot{P}^\alpha = -L_{\dot{x}^\alpha}, P'^\alpha = -L_{x'^\alpha}$ .

**Case 1. Open strings.** Let  $N = [0, \pi] \times [t_1, t_2], (t_1, t_2) \in \mathbb{R}^2, t_1 < t_2$ . We will impose the following constraints on variations of  $f \in V(I^*, L^*)$ :

a) 
$$g^*(v \lrcorner \pi^* \omega)_{\partial N} = 0, \tag{8.18}$$

b) 
$$g^*(v \lrcorner \pi^*(dx^\alpha - \dot{x}^\alpha d\tau - x'^\alpha d\sigma))_B = 0 \tag{8.19}$$

where  $B = [0, \pi] \times t_1 \cup [0, \pi] \times t_2$ ,

c) 
$$\lambda'_\alpha = 0 \text{ on } g(A) \text{ where } A = N \setminus B. \tag{8.20}$$

In this case,  $G$  is any smooth lift of  $F$  to  $Y$  with  $G|_{t=0} = g, (\pi \circ g = f)$ , and  $v$  is a vector field defined along  $g$  with  $v = G_*(\partial/\partial t)|_{t=0}$ . The constraint c) forces the boundary term in the first variation of  $\phi(f)$  vanish.

**Case 2. Closed strings.** Let  $N = S_1 \times [t_1, t_2]$ , with  $S_1$  being the unit circle. Its coordinate  $\sigma \in [0, 2\pi]$ , and  $(t_1, t_2) \in \mathbb{R}^2, t_1 < t_2$ . We will replace the constraints on variations of  $f \in V(I^*, L^*)$  of the previous case with the following:

a) 
$$g^*(v \lrcorner \pi^* \omega)_{\partial N} = 0, \tag{8.21}$$

b) 
$$g^*(v \lrcorner \pi^*(dx^\alpha - \dot{x}^\alpha d\tau - x'^\alpha d\sigma))_B = 0 \tag{8.22}$$

where  $B = S_1 \times t_1 \cup [0, \pi] \times t_2$ .

**The quadratic form A..** The cone  $X' = X | (\dot{x} \cdot \dot{x}) \geq 0, (x' \cdot x') \leq 0$  is convex.  $F'$  has an inverse in  $X'$  with  $F' : X \xrightarrow{F^{-1}} \mathbb{R}^{2m}$  where  $X'' = \mathbb{R}^{2m} | (\dot{\lambda} \cdot \dot{\lambda}) \geq 0, (\lambda' \lambda') \leq 0$ . Hence the matrix

$$A' = \begin{bmatrix} L_{\dot{x}^\alpha \dot{x}^\beta} & L_{\dot{x}^\alpha x'^\beta} \\ L_{x'^\alpha \dot{x}^\beta} & L_{x'^\alpha x'^\beta} \end{bmatrix} \tag{8.23}$$

has an inverse. Therefore, the eigenvalues of  $A'$  do not vanish on  $X'$ . Thus, it suffices to know the eigenvalues of  $A'$  at an interior point of  $X'$  to determine the number of positive eigenvalues of  $A'$  in every point of  $X'$ .

Let

$$a = \{ \dot{x}^0 = 1, \dot{x}^i = 0, x'^1 = 1, x'^j = 0 \text{ with } 1 \leq i \leq m - 1, j = 0 \\ \text{or } 2 \leq j \leq m - 1 \}.$$

Then

$$L_{\dot{x}^0 x^1}(a) = -L_{\dot{x}^1 x^0}(a) = -L_{\dot{x}^i x^i}(a) = L_{x'^i x^i}(a) = 1, 2 \leq i \leq m - 1, \tag{8.24}$$

and all the other elements of  $A'$  are zero. We conclude that the matrix has  $m$ -positive eigenvalues and  $m$ -negative eigenvalues in  $X'$  and the quadratic form  $A$  is neither positive nor negative definite.

**Example 2.** Let  $X_0 = J^1(R^2, R^m), N \subset R^2$ , with  $N$  being a two-dimensional manifold with boundary. Let also

$I^* = \text{span}\{dx^\alpha - x'^\alpha d\sigma - \dot{x}^\alpha d\tau | 1 \leq \alpha \leq m\}$ ,  $x^\alpha$  are coordinates in  $R^m$  and  $x'^\alpha = \frac{\partial x^\alpha}{\partial \sigma}, \dot{x}^\alpha = \frac{\partial x^\alpha}{\partial \tau}$ . Moreover, let

$$\varphi = L\omega = \left[ \sum_{\alpha=1}^m (x'^\alpha)^2 + (\dot{x}^\alpha)^2 \right] d\sigma \wedge d\tau. \tag{8.25}$$

The Cartan system in  $Z$  is

(i) 
$$\partial/\partial \dot{\lambda}_\alpha \lrcorner \Psi_Z = -\pi^*((dx^\alpha - \dot{x}^\alpha d\tau) \wedge \pi^* d\sigma) = 0, \tag{8.26}$$

(ii) 
$$\partial/\partial \lambda'_\alpha \lrcorner \Psi_Z = -\pi^*((dx^\alpha - x'^\alpha d\tau) \wedge \pi^* d\sigma) = 0, \tag{8.27}$$

(iii) 
$$\partial/\partial \dot{x}^\alpha \lrcorner \Psi_Z = -\pi^*(2\dot{x}^\alpha - \dot{\lambda}_\alpha)\omega = 0, \tag{8.28}$$

(iv) 
$$\partial/\partial x'^\alpha \lrcorner \Psi_Z = -\pi^*(2x'^\alpha - \lambda'_\alpha)\omega = 0, \tag{8.29}$$

(v) 
$$\partial/\partial x^\alpha \lrcorner \Psi_Z = -\pi^*d\dot{\lambda}_\alpha \wedge \pi^*d\sigma - d\lambda'_\alpha \wedge \pi^*d\tau = 0. \tag{8.30}$$

Hence

$$Z_1 = Z | L_{\dot{x}^\alpha} = \dot{\lambda}_\alpha, L_{x'^\alpha} = \lambda'_\alpha. \tag{8.31}$$

The prolongation ends at  $Z_1$  with  $(C(\Psi), \pi^*\omega)$  on  $Z_1$  given by (8.26), (8.27) and (8.30). It is easy to prove that  $(C(\Psi), \pi^*\omega)$  in  $Y$  is in involution and  $(I^*, L^*, \varphi, I^*, L^*)$  is a well-posed valued differential system.

**Boundary conditions.** *The constraints on one-parameter variations  $F$  of  $f$  in  $V(I^*, L^*)$  are:*

$$\text{a)} \quad g^*(v \lrcorner \pi^* \omega)_{\partial N} = 0, \quad (8.32)$$

$$\text{b)} \quad g^*(v \lrcorner \pi^*(dx^\alpha - \dot{x}^\alpha d\tau - x'^\alpha d\sigma))_{\partial N} = 0. \quad (8.33)$$

*In this case, too,  $G$  is any smooth lift of  $F$  to  $Y$  with  $G|_{t=0} = g$ , ( $\pi \circ g = f$ ), and  $v$  is a vector field defined along  $g$  with  $v = G|_{t=0*}(\partial/\partial t)$ .*

**The quadratic form  $A$ .** *A simple computation yields*

$$L_{\dot{x}^\alpha \dot{x}^\beta} = 2\delta_{\alpha\beta}, L_{\dot{x}^\alpha x'^\beta} = 0, L_{x'^\alpha x'^\beta} = 2\delta_{\alpha\beta}. \quad (8.34)$$

*Thus, the quadratic form  $A$  is positive definite.*

**Example 3.** *Let  $X_0 = J^1(R^2, R^m)$ . We associate coordinates  $\sigma, \tau$  to  $R^2, x^i$ ,  $1 \leq i \leq m$  to  $R^m$ , and  $x'^i = \frac{\partial x^i}{\partial \sigma}, \dot{x}^i = \frac{\partial x^i}{\partial \tau}$ . Let  $X = X_0|_{g_1} = 0$ , where  $g_1(\dot{x}^1, x^2) = \dot{x}^1 - x^2 = 0$ . Let  $N = B_1$  be a ball with radius 1 centered at  $(0, 0)$ . Then*

$$x^1(t, b) - x^1(a, b) = \int_a^t \frac{\partial x^1}{\partial \tau} d\tau = \int_a^t x^2 d\tau, \quad (8.35)$$

*where  $a \leq 0$  and  $a^2 + b^2 = 1$ .*

**Boundary condition  $h_{A'}$ .** *We have the following system for  $v = F_*(\partial/\partial t)(t, x)|_{t=0}$  where  $F$  is a one-parameter variation of  $f$ :*

$$\frac{\partial v_{x^1}}{\partial \tau} - v_{\dot{x}^1} = 0, \quad (8.36)$$

$$\frac{\partial v_{x^1}}{\partial \sigma} - v_{x'^1} = 0, \quad (8.37)$$

$$\frac{\partial v_{x^1}}{\partial \tau} - v_{x^2} = 0, \quad (8.38)$$

$$\frac{\partial v_{x^1}}{\partial \sigma} - v_{x'^1} = 0. \quad (8.39)$$

*Let  $A' = \{(\tau, \sigma) \in R^2 | (\tau)^2 + (\sigma)^2 = 1 \text{ and } \tau \leq 0\}$ .  $A'$  is nowhere characteristic for (8.38) and the values of  $v_{x^1}$  at  $A'$  and  $v_{x^2}$  in  $N$  determine uniquely a solution in  $N$  for the system of equations. Let  $h_{A'}^1 : A' \rightarrow R$  and  $h_{\partial N}^j : \partial N \rightarrow R$  ( $2 \leq j \leq m$ ) be a smooth function. Assume  $f \in V(I^*, L^*)$ , and let  $I^*, L^*$  be as before. Then,  $f$  satisfies the boundary condition  $[h_{A'}]$  if*

$$x_{A'}^1 = h_{A'}^1 \quad \text{and} \quad x_{\partial N}^j = h_{\partial N}^j. \quad (8.40)$$

In this case,

$$\phi[f] = \int f^* \varphi, \text{ where } f \in V(I^*, L^*, [h_{A'}^1]), \tag{8.41}$$

and

$$\varphi = Lw = [(x^1)^2 + \sum_j (\dot{x}^j)^2 + \sum_j (x'^j)^2] d\sigma \wedge d\tau. \tag{8.42}$$

**Momentum space.** The Cartan system in  $Z$  is:

$$(i) \quad \partial/\partial \dot{\lambda}_i \lrcorner \Psi_Z = -\pi^*((dx^i - \dot{x}^i d\tau) \wedge \pi^* d\sigma) = 0, \tag{8.43}$$

$$(ii) \quad \partial/\partial \lambda'_i \lrcorner \Psi_Z = -\pi^*((dx^i - x'^i d\tau) \wedge \pi^* d\sigma) = 0, \tag{8.44}$$

$$(iii) \quad \partial/\partial \dot{x}^j \lrcorner \Psi_Z = -\pi^*(2\dot{x}^j - \dot{\lambda}_j)\omega = 0, \tag{8.45}$$

$$(iv) \quad \partial/\partial x'^i \lrcorner \Psi_Z = -\pi^*(2x'^i - \lambda'_i)\omega = 0, \tag{8.46}$$

$$(v) \quad \partial/\partial x^i \lrcorner \Psi_Z = -\pi^* d\dot{\lambda}_i \wedge \pi^* d\sigma - d\lambda'_i \wedge \pi^* d\tau = 0. \tag{8.47}$$

From (8.46) and (8.47) we also have  $Y = Z_1 = Z|_{2\dot{x}^j=\dot{\lambda}_j, 2x'^i=\lambda'_i}$ . This Cartan system  $(C(\Psi), \pi^*\omega)$  is non-degenerate. Let us transfer the boundary condition to  $Q_i = Y|_{\pi^*L_i}$ , where  $L_i^* = \text{span}\{dx^i - \dot{x}^i d\tau - x'^i d\sigma, d\sigma, d\tau\}$ . Then,  $f \in V(I^*, L^*)$  satisfies the boundary condition  $h_{A'}$ , if for any lift  $g$  of  $f$  to  $Y$  we have:

$$(\omega'_1 \circ g)|_{A'} = h_{A'}^1 \quad \text{and} \quad (\omega'_j \circ g)|_{\partial N} = h_{\partial N}^j, \tag{8.48}$$

where  $h_{A'}^1 : A' \rightarrow Q_1$  with  $\pi_1 \circ h_{A'}^1 = h_{A'}^1$  and the projection

$\pi_i : Q_i \rightarrow R$  given by  $\pi_i(q) = x^i(q)$ .

Furthermore,  $g$  is a solution to the Euler-Lagrange system satisfying the mixed boundary condition  $[h_{A'}]$ , if  $g$  satisfies (8.43), (8.44) and (8.47), and

$$v \lrcorner (\dot{\lambda}_i \pi^*[dx^i - \dot{x}^i d\tau - x'^i d\sigma] \wedge d\tau + \lambda'_i \pi^*[dx^i - \dot{x}^i d\tau - x'^i d\sigma] \wedge d\sigma)_{g(\partial N \setminus A')} \equiv 0 \tag{8.49}$$

for any element,  $v = F_*(\partial/\partial t)(t, x)|_{t=0}$  where  $F$  is a one parameter variation of  $\pi \circ g$  satisfying  $v_{x^1}|_{A'=0}$  and  $v_{x^2}|_{N=0}$ .

Finally, the quadratic form  $A$  is positive definite.

## 9. Inverse problem for calculus of variations

**Example 4.** *In 1887, Helmholtz solved the following problem:*

*It is given  $P_i = P_i(x, u^j, u_x^j, u_{xx}^j)$ . Is there a Lagrangian  $L(x, u^j, u_x^j)$  such that  $E_i(L) = \partial L / \partial u^i - D_x \partial L / \partial u_x^i = P_i$ , where  $D_x = \partial / \partial x + u_x^i \partial / \partial u^i + u_{xx}^i \partial / \partial u_x^i$ ? He found the following necessary conditions for  $P_i$ :*

$$(i) \quad \partial P_i / \partial u_{xx}^j = \partial P_j / \partial u_{xx}^i, \quad (9.1)$$

$$(ii) \quad \partial P_i / \partial u_x^j = \partial P_j / \partial u_x^i + 2D_x \partial P_j / \partial u_{xx}^i, \quad (9.2)$$

$$(iii) \quad \partial P_i / \partial u^j = \partial P_j / \partial u^i - D_x \partial P_j / \partial u_x^i + D_{xx} \partial P_j / \partial u_{xx}^i. \quad (9.3)$$

*This problem led to the following studies ([2]):*

- (i) - *the derivation and analysis of Helmholtz conditions as necessary and (locally) sufficient conditions for a differential operator to coincide with the Euler-Lagrange operator for some Lagrangian;*
- ii) - *the characterization of the obstructions to the existence of global variational principles for different operators defined on manifolds;*
- iii) - *the invariant inverse problem for different operators with symmetry; and*
- (iv) - *the variational multiplier problem wherein variational principles are found, not for a given differential operator, but rather for the differential equations determined by that operator.*

*That is: find a matrix  $B = [B_i^j]$  such that  $B_i^j P_j = E_i(L)$  for some  $L$  with  $B$  being non-singular.*

*Let  $E \rightarrow M$  be a fibered manifold.  $J^\infty(E)$  is the infinite jet of  $E$ .*

*Let*

$$\theta^i = du^i - u_x^i dx \quad (9.4)$$

$$\theta_x^i = du_x^i - u_{xx}^i dx \quad (9.5)$$

*and*

$$\begin{aligned} \Omega_P = & P_i \theta^i \wedge dx + 1/2 [\partial P_i / \partial u_x^j - D_x \partial P_i / \partial u_{xx}^j] \theta^i \wedge \theta^j \\ & + 1/2 [\partial P_i / \partial u_{xx}^j + \partial P_j / \partial u_{xx}^i] \theta^i \wedge \theta_x^j. \end{aligned} \quad (9.6)$$

*If  $P$  satisfies the Helmholtz conditions, then  $d\Omega_P = 0$ . If the  $H^{n+1}(E) - n + 1$  de Rham cohomology group of  $E$  is trivial, then  $\Omega_P$  is exact. This fact implies that  $P_i$  is globally variational. If  $\theta_L = Ldx + \partial L / \partial u_x^i \theta^i$ ,*

then  $d\theta_L = \Omega_P$ . In 1913, Volterra showed that if  $L = \int_N u^i P_i(x, tu^j, tu^j_x, tu^j_{xx})dt$  where  $N = [0, 1]$ , then

$$E_i(L) = P_i. \tag{9.7}$$

Thus, we have a global solution for the inverse problem in the case of one independent variable and to equations  $P_i = 0$  of second order.

Vaingberg [1969] generalized this result; however his Lagrangian is usually of a much higher order than necessary.

In [2] we can find the following theorem.

**Theorem 9.1.** Let  $P_i$  be a differential operator of order  $2k$

$$P_i = P_i(x, u^j, u^j_1, \dots, u^j_{2k}). \tag{9.8}$$

Then  $P_i$  is the Euler-Lagrange operator of a  $k$ -th order Lagrangian  $L = L(x, u^j, u^j_1, \dots, u^j_k)$  if and only if the functions  $P_i$  satisfy the higher order Helmholtz conditions, and the functions

$$p_i(t) = P_i(x, u^j, u^j_1, \dots, u^j_k, tu^j_{k+1}, \dots, t^k u^j_{2k}) \tag{9.9}$$

are polynomials in  $t$  of degree less or equal to  $k$ .

**Example 5.** Let us now look to another example where we have a function of three independent variables  $x, y$  and  $z$ , with a single dependent variable  $u$ . Let  $T = T(x, y, z, u, u_x, u_y, u_z, u_{xx}, u_{xy}, \dots, u_{zz})$  be a second order operator.

$$E[L] = \partial L / \partial u - D_x \partial L / \partial u_x - D_y \partial L / \partial u_y - D_z \partial L / \partial u_z \tag{9.10}$$

Let  $v$  be a lift to the momentum space of an infinitesimal variation  $F_*(\partial/\partial t)$  of  $f = \pi \circ g$ , where  $g$  is a solution of  $(C(\Psi), \pi^*\omega)$ . The Lie-derivative of  $\psi = \pi^*L\omega + (\pi^j \circ i')^*[i^*(\chi)] \wedge \pi^*\omega_j$  by  $v$  is

$$\begin{aligned} v \lrcorner d\psi + d(v \lrcorner \psi) &= E[L](u)v^1 \pi^*(dx \wedge dy \wedge dz) \\ + d(\partial L / \partial u_x v^1 \pi^*(dy \wedge dz) - \partial L / \partial u_y v^1 \pi^*(dx \wedge dz) + \partial L / \partial u_z v^1 \pi^*(dx \wedge dy)). \end{aligned} \tag{9.11}$$

Suppose that for some vector  $w$  with  $\pi_*w \in T_fV(I^*, L^*, \varphi, [h])$  (i.e.  $w \lrcorner d\theta + d(w \lrcorner \theta)$  for  $\theta = du - u_x dx - u_y dy - u_z dz$  and  $w \lrcorner \theta|_{\partial N} = 0$ ) we have  $v \lrcorner d\psi + d(v \lrcorner \psi) =$

$$\begin{aligned} T[u]v^1 \pi^*(dx \wedge dy \wedge dz) + d(\partial L / \partial u_x w^1 \pi^*(dy \wedge dz) - \partial L / \partial u_y w^1 \pi^*(dx \wedge dz) \\ + \partial L / \partial u_z w^1 \pi^*(dx \wedge dy)). \end{aligned} \tag{9.12}$$

Then we have  $T[u] = E[L](u)$

If we identify  $e_1$  with  $dy \wedge dz$ ,  $e_2$  with  $dz \wedge dx$  and  $e_3$  with  $dx \wedge dy$  at each point of the integral manifold of  $(C(\Psi), \pi^*\omega)$ , we can write

$$d(\partial L / \partial u_x v^1 \pi^*(dy \wedge dz) - \partial L / \partial u_y v^1 \pi^*(dx \wedge dz)) \tag{9.13}$$

$$+\partial L/\partial u_z v^1 \pi^*(dx \wedge dy) = \text{Div}V[u] \pi^*(dx \wedge dy \wedge dz), \quad (9.14)$$

where

$$V[u] = \partial L/\partial u_x v^1 e_1 + \partial L/\partial u_y v^1 e_2 + \partial L/\partial u_z v^1 e_3. \quad (9.15)$$

The divergence operator is defined in terms of the total derivatives  $D_x, D_y$  and  $D_z$ .

We can conclude that  $v \lrcorner d\psi + d(v \lrcorner \psi) = (E[L](u)v + \text{Div}V[u]) \pi^*(dx \wedge dy \wedge dz)$ .

We have

$$E[L](u) = 0 \quad \text{whenever} \quad L[u] = \text{Div}W[u]. \quad (9.16)$$

Suppose  $T[u] = E[L](u)$ . Then the first variation formula is

$$v \lrcorner d\psi + d(v \lrcorner \psi) = (T[u]v^1 + \text{Div}W[u]) \pi^*(dx \wedge dy \wedge dz). \quad (9.17)$$

By applying the Euler-Lagrange operator (i.e.  $E[\alpha[u] \pi^*(dx \wedge dy \wedge dz)] \doteq E[\alpha[u] \pi^*(dx \wedge dy \wedge dz)]$ ), we obtain

$$E[v \lrcorner d\psi + d(v \lrcorner \psi)] = E[T[u]v] \pi^*(dx \wedge dy \wedge dz), \quad \text{since } E(\text{Div}W)(u) = 0. \quad (9.18)$$

We have

$$E[v \lrcorner d\psi + d(v \lrcorner \psi)] = (v \lrcorner dE[L](u) + d(v \lrcorner dE[L](u))) \pi^*(dx \wedge dy \wedge dz) \quad (9.19)$$

$$= (v \lrcorner dT + d(v \lrcorner dT)) \pi^*(dx \wedge dy \wedge dz). \quad (9.20)$$

Therefore

$$E[T[u]v] \pi^*(dx \wedge dy \wedge dz) = (v \lrcorner dT + d(v \lrcorner dT)) \pi^*(dx \wedge dy \wedge dz). \quad (9.21)$$

Let

$$\psi' = \pi^* T \omega + (\pi^j o i')^* [i^*(\chi)] \pi^* \omega_j, \quad (9.22)$$

and

$$v \lrcorner d\psi' + d(v \lrcorner \psi') = E[T[u]v] \pi^*(dx \wedge dy \wedge dz). \quad (9.23)$$

If we define

$$H(T)[v] \pi^*(dx \wedge dy \wedge dz) = v \lrcorner d\psi' + d(v \lrcorner \psi') - E[T(u)v] \pi^*(dx \wedge dy \wedge dz), \quad (9.24)$$

then  $H(T) = 0$  if  $T[u]$  is Euler-Lagrange. Helmholtz equations are:

$$(i) \quad \partial T/\partial u_x = D_x \partial T/\partial u_{xx} + 1/2 D_y \partial T/\partial u_{xy} + 1/2 D_z \partial T/\partial u_{xz}, \quad (9.25)$$

$$(ii) \quad \partial T/\partial u_y = D_y \partial T/\partial u_{yy} + 1/2 D_x \partial T/\partial u_{yx} + 1/2 D_z \partial T/\partial u_{yz}, \quad (9.26)$$

$$(iii) \quad \partial T/\partial u_z = D_z \partial T/\partial u_{zz} + 1/2 D_x \partial T/\partial u_{zx} + 1/2 D_y \partial T/\partial u_{zy}. \quad (9.27)$$



We have a sequences of spaces

$$\begin{array}{cccccccc}
 & & \text{Grad} & \text{Curl} & \text{Div} & E & H & \\
 0 & \rightarrow R & \rightarrow F[u] & \rightarrow V(u) & \rightarrow V(u) & \rightarrow F(u) & \rightarrow F(u) & \rightarrow V(u)
 \end{array} \tag{9.28}$$

that is cochain complex, the Euler-Lagrange complex. Since this complex is exact, the inverse problem is globally solved in this second example.

**9.1. Variational Bicomplex.** Let us introduce now a very important tool for a globalization of the inverse problem.

**Definition 9.1.** A  $p$  form  $\omega$  on  $J^\infty(E)$  is said to be of type  $(r, s)$ , where  $r + s = p$ , if at each point  $x$  of  $J^\infty(E)$

$$\omega(X_1, X_2, \dots, X_p) = 0, \tag{9.29}$$

whenever either

- (i) more than  $s$  of the vectors  $X_1, X_2, \dots, X_p$  are  $\pi_M^\infty$  vertical, or
- (ii) more than  $r$  of the vectors  $X_1, X_2, \dots, X_p$  annihilate all contact one forms.

Note:  $\Omega^{r,s}$  denotes the space of type  $(r, s)$  forms on  $J^\infty(E)$ .

- (i)  $\pi : E \rightarrow M$  be a fiber bundle.
- (ii) Let us assume that there exists a transformation group  $G$  acting on  $E$ , and
- (iii) that there exists a set of differential equations on sections of  $E$ .

$$d = d_H + d_V,$$

$$d_H : \Omega^{r,s}(J^\infty(E)) \rightarrow \Omega^{r+1,s}(J^\infty(E)), \tag{9.30}$$

$$d_V : \Omega^{r,s}(J^\infty(E)) \rightarrow \Omega^{r,s+1}(J^\infty(E)), \tag{9.31}$$

$$d_H^2 = 0, \quad d_H d_V = -d_V d_H, \quad d_V^2 = 0. \tag{9.32}$$

In local coordinates

$$d_H f = [\partial f / \partial x^i + u \alpha_i \partial f / \partial u^\alpha + u_{ij}^\alpha \partial f / \partial u_j^\alpha + \dots] dx^i \tag{9.33}$$

$$d_V f = \partial f / \partial u^\alpha \theta^\alpha + \partial f / \partial u_i^\alpha \theta_i^\alpha + \dots \tag{9.34}$$

The sequences of spaces

$$\begin{array}{cccccccc}
 & & & & & \uparrow d_V & I & \uparrow \delta_V & & \\
 0 & \rightarrow \Omega^{0,3} & \dots & & & \rightarrow \Omega^{n,3} & \rightarrow F^3 & \rightarrow 0 & & \\
 & \uparrow d_V & d_H \uparrow d_V & \dots & d_H \uparrow d_V & d_H \uparrow d_V & I & \uparrow \delta_V & & \\
 0 & \rightarrow \Omega^{0,2} & \rightarrow \Omega^{1,2} & \dots & \rightarrow \Omega^{n-1,2} & \rightarrow \Omega^{n,2} & \rightarrow F^2 & \rightarrow 0 & & \\
 & \uparrow d_V & d_H \uparrow d_V & \dots & d_H \uparrow d_V & d_H \uparrow d_V & I & \uparrow \delta_V & & \\
 0 & \rightarrow \Omega^{0,1} & \rightarrow \Omega^{1,1} & \dots & \rightarrow \Omega^{n-1,1} & \rightarrow \Omega^{n,1} & \rightarrow F^1 & \rightarrow 0 & & \\
 & \uparrow d_V & d_H \uparrow d_V & \dots & d_H \uparrow d_V & d_H \uparrow d_V & & & & \\
 0 & \rightarrow R & \rightarrow \Omega^{0,0} & \rightarrow \Omega^{1,0} & \dots & \rightarrow \Omega^{n-1,0} & \rightarrow \Omega^{n,0} & & & 
 \end{array}$$



$$-g_{u_{\sigma'}}^{\rho} a_{\sigma'}^{\sigma'} g_{u_x^{\mu}}^{\sigma} + g_{u_x^{\mu} u_x^{\sigma}}^{\rho} a_{\rho'}^{\sigma} (g_x^{\rho'} - g_{u_x^{\alpha}}^{\rho'} u_x^{\alpha}) + g_{u_x^{\mu} u_x^{\sigma}}^{\rho} a_{\rho'}^{\sigma} u_x^{\rho'} + g_{u_x^{\mu} u_x^{\nu}}^{\rho} u_x^{\nu}. \quad (9.42)$$

$$L_{\mu} = (\partial/\partial u_x^{\mu} - a_{\rho}^{\sigma} g_{u_x^{\mu}}^{\rho} \partial/\partial u_x^{\sigma}) L, \quad (9.43)$$

$$L_{\mu\nu} = (\partial/\partial u_x^{\mu} - a_{\rho}^{\sigma} g_{u_x^{\mu}}^{\rho} \partial/\partial u_x^{\sigma}) L_{\mu}, \quad (9.44)$$

and

$$A_{\mu\nu}(\lambda_1, \dots, \lambda_{m-l}) = L_{\mu\nu} + \lambda_{\rho} (g_{u_x^{\sigma} u_x^{\sigma'}}^{\rho} a_{\rho'}^{\sigma} g_{u_x^{\nu}}^{\rho'} a_{\rho''}^{\sigma'} g_{u_x^{\mu}}^{\rho''} - g_{u_x^{\sigma} u_x^{\mu}}^{\rho} a_{\rho'}^{\sigma} g_{u_x^{\nu}}^{\rho'} - g_{u_x^{\nu} u_x^{\sigma}}^{\rho} a_{\rho'}^{\sigma} g_{u_x^{\mu}}^{\rho'} + g_{u_x^{\nu}}^{\rho} u_x^{\mu}), \quad (9.45)$$

$$[a_{\rho}^{\sigma}] = [g_{u_x^{\rho}}^{\sigma}]^{-1} \text{ with } 1 \leq \rho, \rho', \rho'', \sigma, \sigma' \leq m-l \text{ and } m-l+1 \leq \mu, \nu \leq m. \quad (9.46)$$

$$\begin{aligned} \psi \equiv & (L_{\mu} - \lambda_{\mu}) \pi^* (du_x^{\mu} \wedge dx) + (d\lambda_{\mu} - (A_{\mu} + \lambda_{\rho} B_{\mu}^{\rho}) \pi^* dx + \lambda_{\rho} A_{\mu\nu}^{\rho} \pi^* du_x^{\nu}) \wedge \pi^* \theta^{\mu} \\ & + (d\lambda_{\sigma} - (A_{\sigma} + \lambda_{\rho} B_{\sigma}^{\rho}) \pi^* dx + \lambda_{\rho} A_{\mu\sigma}^{\rho} \pi^* du_x^{\mu}) \wedge \pi^* \theta^{\sigma} \\ & \text{mod}\{\pi^*(\theta^{\alpha} \wedge \theta^{\alpha'}) | 1 \leq \alpha, \alpha' \leq m\}, \end{aligned} \quad (9.47)$$

with

$$A_{\mu} = L_{u^{\mu}} - L_{u_x^{\sigma'}} a_{\rho'}^{\sigma'} g_{u_x^{\mu}}^{\rho'} + L_{u_x^{\sigma'}} a_{\rho'}^{\sigma'} g_{u_x^{\sigma}}^{\rho'} a_{\rho''}^{\sigma''} g_{u_x^{\mu}}^{\rho''} - L_{u^{\rho}} a_{\sigma}^{\rho} g_{u_x^{\mu}}^{\sigma} \quad (9.48)$$

$$A_{\sigma} = L_{u^{\rho}} a_{\sigma}^{\rho} - L_{u_x^{\sigma'}} a_{\rho'}^{\sigma'} g_{u_x^{\rho}}^{\rho'} a_{\sigma}^{\rho}. \quad (9.49)$$

The Cartan system is

$$\pi^* \theta^{\alpha} \quad (1 \leq \alpha \leq m), \quad (9.50)$$

$$(L_{\mu} - \lambda_{\mu}) \pi^* dx \quad (m-l+1 \leq \mu \leq m), \quad (9.51)$$

$$(d\lambda_{\mu} - (A_{\mu} + \lambda_{\rho} B_{\mu}^{\rho}) \pi^* dx + \lambda_{\rho} A_{\mu\nu}^{\rho} \pi^* du_x^{\nu}) \quad (m-l+1 \leq \mu \leq m), \quad (9.52)$$

$$(d\lambda_{\sigma} - (A_{\sigma} + \lambda_{\rho} B_{\sigma}^{\rho}) \pi^* dx + \lambda_{\rho} A_{\mu\sigma}^{\rho} \pi^* du_x^{\mu}) \quad (1 \leq \sigma \leq m-l). \quad (9.53)$$

**Proposition 9.2.** *Let  $(I^*, L^*)$  be a locally embeddable differential system defined in  $X = J^1(R, R^m)|g^\rho(x, u^j, u_x^j) = 0$ ,  $\text{rank}[\partial g^\rho/\partial u_x^j] = m - l$ ,  $1 \leq j \leq m$  and  $1 \leq \rho \leq m - l$ ,  $l \geq 0$ , where  $I^* = \text{span} \{\theta^\alpha | 1 \leq \alpha \leq m\}$  and  $L^* = \text{span} \{\theta^\alpha, dx | 1 \leq \alpha \leq m\}$ ,*

$$\theta^\rho = g_{u_x^\sigma}^\rho (du^\sigma - u_x^\sigma dx) + g_{u_x^\mu}^\rho (du^\mu - u_x^\mu dx) \quad 1 \leq \sigma, \rho \leq m - l, \quad (9.54)$$

$$\theta^\mu = du^\mu - u_x^\mu dx \quad m - l + 1 \leq \mu, \nu \leq m. \quad (9.55)$$

Let  $Q_i = Q_i(x, u^j, u_x^j, u_{xx}^\mu, \lambda_\rho \lambda_{\rho x})$ ,  $1 \leq i \leq m$ , with  $Q_i(x, u^j, u_x^j, tu_{xx}^\mu, \lambda_\rho \lambda_{\rho x})$  being polynomial in  $t$  of degree less or equal to 1, and

$$P_\mu = Q_\mu + \lambda_\rho B_\mu^\rho - \lambda_\rho A_{\mu\nu}^\rho \frac{du_x^\nu}{dx}, \quad (9.56)$$

$$R_\sigma = Q_\sigma - \lambda_{\sigma x} + \lambda_\rho B_\sigma^\rho - \lambda_\rho A_{\mu\sigma}^\rho \frac{du_x^\mu}{dx}, \quad (9.57)$$

and

$$R_\mu = P_\mu + D_x(\partial P_\mu/\partial u_{xx}^\mu). \quad (9.58)$$

Furthermore, let us assume that the functions  $P_\mu$  satisfy the Helmholtz conditions, that the functions  $R_\alpha$  do not depend on  $\lambda_\rho$  and  $(\lambda_\rho)_x$  coordinates, and the 1-form  $\Theta = R_\alpha(x, u^j, u_x^\mu, u_{xx}^\mu)\theta^\alpha$  is closed mod  $R$ , where  $R = C^\infty(Z, R^*)$ ,  $Z = J^2(R, R^m)|g^\rho(x, u^j, u_x^j) = 0$  with coordinates  $\{x, u^j, u_x^\mu, u_{xx}^\mu\}$  and  $R^* = \text{span} \{dx, du_x^\mu, du_{xx}^\mu\}$ . Then,  $Q_i$  is locally a Euler-Lagrange operator for a Lagrangian  $L(x, u^j, u_x^\mu)$ .

Proof: From Theorem 9.1 we know that a function  $F(x, u^j, u_x^j)$  can be found that does not depend on  $u_{xx}^\nu$ , such that  $E_\mu(F) = \partial F/\partial u^\mu - D_x \partial F/\partial u_x^\mu = P_\mu$  (note that if  $R_\mu$  does not depend on  $\lambda_\rho$ , then neither does  $P_\mu$ ).

Therefore,

$$\partial P_\mu/\partial u_{xx}^\nu = F_{\mu\nu}, \quad (9.59)$$

where

$$F_{\mu\nu} = (\partial/\partial u_x^\mu - a_\rho^\sigma g_{u_x^\mu}^\rho \partial/\partial u_x^\sigma) F_\nu, \quad (9.60)$$

and

$$F_\mu = (\partial/\partial u_x^\mu - a_\rho^\sigma g_{u_x^\mu}^\rho \partial/\partial u_x^\sigma) F. \quad (9.61)$$

The  $R_\mu$  functions satisfy

$$R_\mu = (\partial/\partial u^\mu - a_\rho^\sigma g_{u^\mu}^\rho \partial/\partial u_x^\sigma) F. \quad (9.62)$$

Hence, if the  $\Theta$ -form is closed mod  $R$ , then locally

$$R_\sigma = (\partial/\partial u^\sigma - a_\rho^{\sigma'} g_{u^\sigma}^\rho \partial/\partial u_x^{\sigma'}) F. \quad (9.63)$$

Finally, we make  $F = L$ .

In addition, if the domain of the  $R_\alpha$  functions is simply connected and

$$\begin{aligned} \Omega_P &= P_\mu \theta^\mu \wedge dx + 1/2[\partial P_\mu / \partial u_x^j - D_x \partial P_\mu / \partial u_{xx}^j] \theta^\mu \wedge \theta^j \\ &\quad + 1/2[\partial P_\mu / \partial u_{xx}^j + \partial P_j / \partial u_{xx}^\mu] \theta^\mu \wedge \theta_x^j. \end{aligned} \tag{9.64}$$

is exact, then we have a global solution for the inverse problem.

**Example 6.** Let  $X$  be the  $J^1(R, R^3)|g(v, y, z, v_x, y_x, z_x) = 0$ , where

$$g(v, y, z, v_x, y_x, z_x) = mvv_x - mgz_x + R\sqrt{1 + (y_x)^2 + (z_x)^2}. \tag{9.65}$$

Let

$$Q_1 = -\lambda_{\rho_x} - \frac{\sqrt{1 + (y_x)^2 + (z_x)^2}}{mv^3} = 0, \tag{9.66}$$

and

$$\begin{aligned} Q_2 &= -\frac{Ry_x}{mv^3} - \frac{v(1 + z_x^2)y_{xx} - y_x z_x z_{xx} - v_x y_x \sqrt{1 + (y_x)^2 + (z_x)^2}}{v^2(\sqrt{1 + (y_x)^2 + (z_x)^2})^3} \\ &\quad - \lambda_1 \left( \frac{R(1 + z_x^2)y_{xx}}{(\sqrt{1 + (y_x)^2 + (z_x)^2})^3} + \frac{Rz_x y_x z_{xx}}{(\sqrt{1 + (y_x)^2 + (z_x)^2})^3} \right) = 0, \end{aligned} \tag{9.67}$$

$$\begin{aligned} Q_3 &= -\frac{\sqrt{1 + (y_x)^2 + (z_x)^2}}{mv^3} \left( mg - \frac{Rz_x}{\sqrt{1 + (y_x)^2 + (z_x)^2}} \right) \\ &\quad - \frac{v(1 + y_x^2)z_{xx} - y_x z_x y_{xx} - v_x z_x \sqrt{1 + (y_x)^2 + (z_x)^2}}{v^2(\sqrt{1 + (y_x)^2 + (z_x)^2})^3} \\ &\quad - \lambda_1 \left( \frac{R(1 + y_x^2)z_{xx}}{(\sqrt{1 + (y_x)^2 + (z_x)^2})^3} + \frac{Rz_x y_x y_{xx}}{(\sqrt{1 + (y_x)^2 + (z_x)^2})^3} \right) = 0. \end{aligned} \tag{9.68}$$

Hence,

$$P_2 = -\frac{Ry_x}{mv^3} - \frac{v(1 + z_x^2)y_{xx} - y_x z_x z_{xx} - v_x y_x \sqrt{1 + (y_x)^2 + (z_x)^2}}{v^2(\sqrt{1 + (y_x)^2 + (z_x)^2})^3}, \tag{9.69}$$

$$\begin{aligned} P_3 &= -\frac{\sqrt{1 + (y_x)^2 + (z_x)^2}}{mv^3} \left( mg - \frac{Rz_x}{\sqrt{1 + (y_x)^2 + (z_x)^2}} \right) \\ &\quad + \frac{v(1 + y_x^2)z_{xx} - y_x z_x y_{xx} - v_x z_x \sqrt{1 + (y_x)^2 + (z_x)^2}}{v^2(\sqrt{1 + (y_x)^2 + (z_x)^2})^3}, \end{aligned} \tag{9.70}$$

and

$$R_1 = -\frac{\sqrt{1 + (y_x)^2 + (z_x)^2}}{mv^3}, \quad (9.71)$$

$$R_2 = -\frac{Ry_x}{mv^3\sqrt{1 + (y_x)^2 + (z_x)^2}}, \quad (9.72)$$

$$R_3 = -\frac{\sqrt{1 + (y_x)^2 + (z_x)^2}}{mv^3}\left(mg - \frac{Rz_x}{\sqrt{1 + (y_x)^2 + (z_x)^2}}\right). \quad (9.73)$$

It is easy to verify that  $P_2$  and  $P_3$  satisfy Helmholtz conditions, and that the 1-form  $\Theta = R_1\theta^1 + R_2\theta^2 + R_3\theta^3$  is closed mod  $R$ , with  $R^* = \text{span}\{dx, dy_x, dz_x\}$  and  $R = C^\infty(X, R^*)$ . The Lagrangian for this example is  $L = \frac{\sqrt{1+(y_x)^2+(z_x)^2}}{v}$ .

## References

- [1] I. M. Anderson, Natural variational principles on Riemannian manifolds, Ann. of Math., 120 (1984), 329-370.
- [2] —, Introduction to the variational bicomplex, Contemporary Mathematics, 132 (1992), 51-73.
- [3] —, On the existence of global variational principles, Amer. J. Math., 102 (1980), 781-868
- [4] R. Bryant, S. S. Chern, R. Gardner and P. Griffiths, Essays on exterior differential systems, Springer-Verlag, New York, 1990.
- [5] C. Caratheodory, Variationsrechnung bei mehrfachen Integralen, Acta Szeged 4 (1929).
- [6] E. Cartan, Les systèmes différentielles extérieures et leurs applications géométriques, Herman, Paris, (1945).
- [7] P. Dedecker, Calcul des variations, formes différentielles et champs géodésiques. Colloques. Internat. du C.N.R.S, Strasbourg, (1953).
- [8] —, Calcul des variations et topologie algébrique Mem. Soc. Roy. Sc. Liège 4-e série, XIX, fase I (1957).
- [9] —, On the generalization of symplectic geometry to multiple integrals in the calculus of variations, Lecture Notes in Math., Vol. 570, Springer, Berlin and New York, 1977.
- [10] Th. De Donder, Théorie invariante du calcul de variations, Gauthier-Villars, Paris, 1935.
- [11] D. G. B. Edelen, The invariance group for Hamiltonian systems of partial differential equations, Arch.Rational Mech. Anal., 5(1961), 95-176.
- [12] —, Nonlocal variations and local invariance of fields, American Elsevier. New York, 1969.
- [13] H. I. Eliasson, Variational integrals in fiber bundles, Proc. Sympos. Pure Math., Vol. 16, Amer. Math. Soc., Providence, RI (1970), 67-89.
- [14] G. B. Folland, Introduction to partial differential equations, Mathematical Notes 17, Princeton University Press, Princeton, 1976.

- [15] P. L. Garcia, The Poincaré-Cartan invariant in the calculus of variations, *Sympos. Math.*, 14 (1974), 219-227.
- [16] —, Gauge algebras, curvature and symplectic structure. *J. Differential Geometry*, 12 (1977), 209-246.
- [17] —, Critical principal connections and gauge invariance, *Rep. Math. Phys.*, 13 (1978), 337-344.
- [18] —, Tangent structure of Yang-Mills equations and Hodge Theory, *Lecture Notes in Math.*, Vol. 775, Springer, Berlin and New York, 1980.
- [19] P. L. Garcia, A. Pérez-Rendón, Symplectic approach to the theory of quantized fields I, *Comm. Math Phys.*, 13 (1969), 24-44.
- [20] —, Symplectic approach to the theory of quantized fields. II, *Arch. Rational Mech. Anal.*, 43 (1971), 101-124.
- [21] —, Reducibility of the symplectic structure of minimal interactions, *Lecture Notes in Math.*, Vol. 676, Springer, Berlin and New York, 1978, 409-433.
- [22] H. Goldschmidt, S. Sternberg, The Hamilton-Cartan formalism in the calculus of variations, *Ann. Inst. Fourier (Grenoble)*, 23 (1973), 203-267.
- [23] P. Griffiths, *Exterior differential systems and the calculus of variations*, Birkhäuser Boston, Basel, Stuttgart, 1983.
- [24] C. Günther, The polysymplectic Hamiltonian formalism in the field theory and calculus of variations I: The local case, *J. Differential Geometry*, 25 (1987), 23-53.
- [25] P. G. Henriques, *Calculus of variations in the context of exterior differential systems*, *Differential Geometry and its Applications (North-Holland)*, 3 (1993), 331-372
- [26] —, Well-posed variational problem with mixed endpoint conditions, *Differential Geometry and its Applications (North-Holland)*, 3 (1993), 373-393.
- [27] —, The Noether theorem and the reduction procedure for the variational calculus in the context of differential systems, *C. R. Acad. Sci. Paris Sér. I Math.*, 317 (1993), 987-992.
- [28] R. Hermann, *Differential geometry and the calculus of variations*, Academic Press, New York, 1968.
- [29] J. Kijowski, W. M. Tulczyjew, *A symplectic framework for field theories*, *Lecture Notes in Math.*, Vol 107, Springer, Berlin and New York, 1979.
- [30] D. Krupka, Lagrange theory in fibered manifolds, *Rep. Math. Phys.* 2 (1970), 121-133.
- [31] —, A geometric theory of ordinary first order variational problems in fibered manifolds I: Critical sections, *J. Math Anal. Appl.*, 49 (1975), 180-206.
- [32] —, A geometric theory of ordinary first order variational problems in fibered manifolds II: Invariance, *J. Math Anal. Appl.* 49, (1975), 469-476.
- [33] T. Lepage, Sur les champs géodésiques des intégrales multiples, *Bull. Acad. Roy. Belg, Cl. Sc. , 5.ème série* 22 (1936).
- [34] —, Sur les champs géodésiques des intégrales multiples, *Bull. Acad. Roy. Belg, Cl. Sc. , 5.ème série* 27 (1941).
- [35] —, Champs stationnaires, champs géodésiques et formes intégrables I, II. *Bull. Acad. Roy. Belg, Cl. Sc. , 5.ème série* 28 (1942).
- [36] A. Liesen, Feldtheorie in der Variatonrechnung mehfacder Integrale, *Math Annalen* I-171 (1967), 194-218, II-171 (1967), 273-292.
- [37] P. J. Olver, Euler operators and conservation laws of BBM equation, *Math. Proc. Cam. Phil. Soc.*, 85 (1979), 143-160.
- [38] —, *Applications of Lie groups to differential equations*, Springer-Verlag, New York, 1986.

- [39] R. Ouzilou, Expression symplectic des problèmes variationnels, *Sympos. Math* 14 (1972), 85-98.
- [40] H. Rund, *The Hamilton-Jacobi theory in the calculus of variations*, Van Nostrand, Princeton, NJ, 1966.
- [41] J. Scherk, An introduction to the theory of dual models and strings, *Rev Mod. Physics* 47 (1975) 123-164.
- [42] J. H. Schwarz, *Superstring theory*, Physics Report 89 (North Holland Publishing Company), 3 (1982) 223-322.
- [43] I. Stakgold, *Green's functions and boundary value problems*, A. Wiley, 1979.
- [44] F. Takens, Symmetries, conservation laws and variational principles, *Lectures Notes in Mathematics*, Vol. 597, Springer-Verlag, New York, (1977), 581-603.
- [45] —, A global version of the inverse problem to the calculus of variations, *J. Differential Geometry*, 14 (1979), 543-562.
- [46] W. M. Tulczyjew, The Euler-Lagrange resolution, *Lecture Notes in Mathematics*, Vol. 836, Springer-Verlag, New York, (1980), 22-48.
- [47] —, Cohomology of the Lagrange complex, *Ann. Scuola. Norm. Sup. Pisa* (1988), 217-227.
- [48] A. M. Vinogradov, the C-spectral sequence, Lagrangian formalism and conservation laws I, II, *J. Math Anal. Appl* 100 (1984), 1-129.
- [49] —, Symmetries and conservation laws of partial differential equations: basic notions and results, *Acta Appl. Math.*, 15 (1989), 3-22.
- [50] —, Scalar differential invariants, diffeities and characteristic classes, *Mechanics, Analysis and Geometry: 200 Years after Lagrange*, M. Francavglia(ed), Elsevier Amsterdam, (1991), 379-416.
- [51] H. Weyl, Geodesic fields in the calculus of variations for multiple integrals, *Annals of Math.*, 36 (1935), 607-629.