

On the spectrum of gauge periodic point perturbations on the Lobachevsky plane

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Introduction

Let H be the periodic Schrödinger operator with a uniform magnetic field on the Euclidean plane \mathbb{R}^2 ; the spectral structure of H depends drastically on the flux η of the field through an elementary cell of the period lattice: If η is a rational number, then the spectrum of H has a band structure, whereas for irrational η regions with Cantor spectrum may appear [1]. The situation is different in the case of the Lobachevsky plane. Indeed, if the group of periods of H is the modular group $SL(2, \mathbb{Z})$, then the spectrum of H has band structure for any value of the flux η [2],[3]. This result is obtained under the condition that the periodic perturbation of the free magnetic Hamiltonian is the operator of multiplication by a periodic function. On the other hand, an interesting class of periodic Schrödinger operators is obtained by so-called point perturbations since these perturbations give a broad collection of explicitly solvable models [4], [5]. In particular, the point perturbations of the two-dimensional magnetic Schrödinger operator are widely used in theoretical physics to investigate the transport properties of two-dimensional systems [6], [7].

In the present paper, the results of the articles [2], [3] are extended to periodic point perturbations of magnetic Schrödinger operators on the Lobachevsky plane. It is proved that these operators have band spectrum, too, if the associated C^* -algebra has the Kadison property. This result seems to be relevant in studying how the geometry of a two-dimensional electron system influences its spectral and transport properties [8], [9].



1. The free Hamiltonian

We consider a two-dimensional complete Riemannian manifold X of negative curvature (the Lobachevsky plane). We suppose that X is realized as the Poincaré upper half-plane, \mathbb{H} ,

$$\mathbb{H} = \{z = x + iy \in \mathbf{C} : y > 0\}, \quad (1.1)$$

endowed with the metric

$$ds^2 = \frac{a^2}{y^2} (dx^2 + dy^2); \quad (1.2)$$

thus, the curvature of X is equal to $R = -1/a^2$. The geodesic distance between points $z = x + iy, z' = x' + iy' \in X$ has the form

$$d(z, z') = a \cosh^{-1} \left[1 + \frac{|z - z'|^2}{2yy'} \right], \quad (1.3)$$

and the volume form is given by

$$d\sigma = \frac{a^2}{y^2} dx \wedge dy. \quad (1.4)$$

By definition, a constant uniform magnetic field \mathbf{B} perpendicular to X is a 2-form

$$\mathbf{B} = B \frac{a^2}{y^2} dx \wedge dy, \quad B \in \mathbb{R}, \quad (1.5)$$

where B is the strength of the field. The form \mathbf{B} is exact, i.e. $\mathbf{B} = d\mathbf{A}$, where the 1-form \mathbf{A} is called a vector potential of \mathbf{B} . The vector potential \mathbf{A} is defined up to a gauge term; we shall use the so-called Landau gauge,

$$\mathbf{A} = \frac{Ba^2}{y} dx. \quad (1.6)$$

The Hamiltonian of a free quantum-mechanical particle (of mass m and charge e) moving in the plain X subjected to the external field \mathbf{B} has the form [10]

$$H^0 = \frac{y^2}{2ma^2} \left\{ \left(-i\hbar \frac{\partial}{\partial x} - \frac{e}{c} \frac{b}{y} \right)^2 - \hbar^2 \frac{\partial^2}{\partial y^2} \right\}, \quad (1.7)$$

where we write

$$b := Ba^2 \quad (1.8)$$

(as usual, c denotes the velocity of light and \hbar is the Planck constant). In what follows we use a system of units in which $e = c = \hbar = 1$ and $m = 1/2$. In this case, H^0 is a self-adjoint operator in $L^2(X)$ namely the closure of the symmetric operator τ^0 ,

$$\tau^0 = \frac{1}{a^2} \left\{ -y^2 \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) + 2iby \frac{\partial}{\partial x} + b^2 \right\}, \quad (1.9)$$

with domain $\mathcal{D}(\tau^0) = C_0^\infty(X)$. It is well known (see e.g. [11]) that $\mathcal{D}(H^0) \subset C(X)$. It is useful to note that with

$$Df := df + i\mathbf{A} \wedge f, \quad f \in C_0^\infty(X), \quad (1.10a)$$

we obtain

$$\tau^0 = D^*D. \quad (1.10b)$$

The spectrum of H^0 , $\text{spec } H^0$, consists of two parts. The first one is the pure point spectrum formed by the finitely many eigenvalues (the Landau levels)

$$\lambda_l = \frac{1}{a^2} \left[\frac{1}{4} + b^2 - \left(l + \frac{1}{2} - |b| \right)^2 \right], \quad 0 \leq l < |b| - \frac{1}{2}. \quad (1.11)$$

The second part of the spectrum is the absolutely continuous spectrum which fills out the whole semi-axis $\left[\frac{1}{a^2} \left(\frac{1}{4} + b^2 \right), \infty \right)$ [12].

The main role in this paper is played by the resolvent, $R^0(\zeta) = (H^0 - \zeta)^{-1}$, of H^0 . The integral kernel of $R^0(\zeta)$ (i.e., the Green's function $G^0(z; z'; \zeta)$ of H^0) is determined in [10]; it has the form

$$G^0(z, z'; \zeta) = \frac{\sigma^{-t}}{4\pi} \exp(ib\varphi) \frac{\Gamma(t+b)\Gamma(t-b)}{\Gamma(2t)} F\left(t+b, t-b; 2t; \frac{1}{\sigma}\right). \quad (1.12)$$

Here $F(\alpha, \beta; \gamma; z)$ is the Gauss hypergeometric function,

$$\sigma := \cosh^2 \frac{d(z, z')}{a}, \quad (1.13)$$

$$\varphi := 2 \arctg \frac{x - x'}{y + y'}, \quad (1.14)$$

and $t = t(\zeta), \zeta \in \mathbb{C} \setminus \text{spec } H^0$, is uniquely defined by the conditions

$$\zeta = \frac{t(1-t) + b^2}{a^2}, \quad \text{Re } t > 1/2. \quad (1.15)$$

In the following lemma, we collect the properties of G^0 which are needed below.

Lemma 1 *For any $z \in X$ there exists the limit*

$$\lim_{z' \rightarrow z} \left[G^0(z, z'; \zeta) - \frac{1}{2\pi} \log d(z, z') \right] =: q(\zeta). \quad (1.16)$$

This limit does not depend on z , in fact

$$q(\zeta) = -\frac{1}{4\pi} \left[\psi(t+b) + \psi(t-b) + 2C_E - \log 4a^2 \right]. \quad (1.17)$$

Here $\psi(z) = [\log \Gamma(z)]'$, and C_E is the Euler constant.

Proof According to [13; 2.3.1(2)] we have for $|1-z| < 1$

$$F(\alpha, \beta; \alpha + \beta; z) = \frac{\Gamma(\alpha + \beta)}{\Gamma(\alpha)\Gamma(\beta)} \sum_{n=0}^{\infty} \frac{(\alpha)_n (\beta)_n}{(n!)^2} [k_n - \log(1-z)] (1-z)^n, \quad (1.18)$$

where $(\alpha)_n = \alpha(\alpha+1)\dots(\alpha+n-1)$, and

$$k_n = 2\psi(n+1) - \psi(\alpha+n) - \psi(\beta+n). \quad (1.19)$$

In view of (1.12), we have to perform the limit $\sigma \rightarrow 1$. Substituting (1.17) and (1.18) in (1.11) and taking into account that $C_E = -\psi(1)$, we get (1.16). \square

Since $\text{Re } t(\zeta) \rightarrow -\infty$ as $\text{Re } \zeta \rightarrow -\infty$, we obtain the following assertion from the well known asymptotics of $\psi(z)$.

Lemma 2 $\text{Re } q(\zeta) \rightarrow -\infty$ as $\text{Re } \zeta \rightarrow -\infty$.

Lemma 3 *For every $\varepsilon > 0$ and $\zeta \in \mathbb{C}$ with $\text{Re } \zeta < 0$, there exist constants $c_1(\varepsilon, \zeta) = c_1 > 0$ and $\tilde{c}_1(\varepsilon, \zeta) = \tilde{c}_1 > 0$ such that*

$$|G^0(z, z'; \zeta)| \leq c_1 \exp(-\tilde{c}_1 d(z, z')), \quad (1.20)$$

whenever $d(z, z') \geq \varepsilon$.

Moreover, if ε is fixed then $c_1(\zeta) = o(1)$ and $\tilde{c}_1(\zeta) \rightarrow \infty$ as $\text{Re } \zeta \rightarrow -\infty$.

Proof If $\operatorname{Re} \zeta < 0$, then from formula 2.12 (1) in [13] we get for $|z| < 1$

$$F(t+b, t-b; 2t; z) = \frac{\Gamma(2t)}{\Gamma(t+b)\Gamma(t-b)} \int_0^1 s^{t-b-1}(1-s)^{t+b-1}(1-zs)^{-t-b} ds. \quad (1.21)$$

The assumptions imply $\operatorname{Re} t > |b| + 1/2$ and $z \in \mathbb{R}$, $|z| < 1$. Then we estimate

$$\begin{aligned} & \left| \int_0^1 s^{t-b-1}(1-s)^{t+b-1}(1-zs)^{-t-b} ds \right| \\ &= \left| \int_0^1 s^{t-b-1/2} \left(\frac{1-s}{1-zs} \right)^{t+b-1/2} (1-zs)^{-1/2} (s(1-s))^{-1/2} ds \right| \\ &\leq (1-|z|)^{-1/2} \int_0^1 s^{\operatorname{Re} t - |b| - 1/2} (s(1-s))^{-1/2} ds, \end{aligned}$$

and the last integral converges to 0, by dominated convergence. Substituting this in (1.11), we conclude the proof. \square

Lemma 4 *Let K be a compact subset of X and z_0 a fixed point of X . Then for every $\varepsilon > 0$ and $\zeta \in \mathbb{C}$, $\operatorname{Re} \zeta < 0$, there exist constants $c_2(K, z_0, \varepsilon, \zeta) =: c_2 > 0$ and $\tilde{c}_2(\varepsilon, \zeta) =: \tilde{c}_2 > 0$ such that*

$$\sup\{|G^0(z, z'; \zeta)| : z' \in K\} \leq c_2 \exp(-\tilde{c}_2 d(z, z_0)), \quad (1.22)$$

whenever $d(z, K) \geq \varepsilon$.

Moreover, if K, z_0 , and ε remain fixed then $c_2(\zeta) = o(1)$ and $\tilde{c}_2(\zeta) \rightarrow \infty$ as $\operatorname{Re} \zeta \rightarrow -\infty$.

Proof Let $z, z' \in X$ such that $z' \in K, d(z, K) \geq \varepsilon$. Then $d(z, z') \geq \varepsilon$ and

$$d(z, z') \geq d(z, z_0) - d(z_0, z') \geq d(z, z_0) - k, \quad (1.23)$$

where $k = \sup\{d(z', z_0) : z' \in K\}$. Substituting (1.23) in (1.19) completes the proof. \square

We now recall that a bounded linear operator, L , in the space $L^2(X)$ is called a Carleman operator if L has an integral kernel $L(z, z')$ such that

$$\int_X |L(z, z')|^2 d\sigma(z') < \infty$$

for almost every $z \in X$ [14, Thm. 11.6].

Lemma 5 *For any $\zeta \in \mathbb{C} \setminus \text{spec } H^0$, the resolvent $R^0(\zeta)$ is a Carleman operator. Moreover, the integral*

$$\int_X |G^0(z, z'; \zeta)|^2 d\sigma(z') \tag{1.24}$$

does not depend on z .

Proof The last statement of the lemma is valid because, for a fixed ζ , the value $|G^0(z, z'; \zeta)|$ depends on $d(z, z')$ only. Fix $z_0 \in X$ and denote by $B_\varepsilon(z_0)$ the metric ball around z_0 in X of radius ε .

Taking into account Lemma 1 we obtain with $\varepsilon = \varepsilon(z_0, \zeta)$

$$\int_{B_\varepsilon} |G^0(z_0, z'; \zeta)|^2 d\sigma(z') < \infty, \tag{1.25}$$

and with Lemma 3,

$$\int_{B_1} |G^0(z_0, z'; \zeta)|^2 d\sigma(z') < \infty.$$

It is known that the area of the circle B_n is equal to $2\pi(\cosh \frac{n}{a} - 1)$. Hence, $\text{area}(B_{n+1} \setminus B_n) = O(e^n)$ as $n \rightarrow \infty$. According to Lemma 3 we can find a number ζ_0 , $\text{Re } \zeta_0 < 0$, such that

$$\int_{B_{n+1} \setminus B_n} |G^0(z_0, z'; \zeta_0)|^2 d\sigma(z') = O(e^{-2n}), \tag{1.26}$$

implying

$$\int_{X \setminus B_1} |G^0(z_0, z'; \zeta_0)|^2 d\sigma(z') < \infty.$$

Therefore, $R^0(\zeta_0)$ is a Carleman operator. It is known that the space of Carleman operators in $L^2(X)$ is a right ideal in the algebra $\mathcal{L}(L^2(X))$ of all bounded linear operators in $L^2(X)$ [14]. Using the Hilbert identity

$$R^0(\zeta) = R^0(\zeta_0) + (\zeta - \zeta_0)R^0(\zeta_0)R^0(\zeta) \quad (1.27)$$

we see that $R^0(\zeta)$ is a Carleman operator for any $\zeta \in \mathbb{C} \setminus \text{spec } H^0$. \square

Lemma 6 *Let K be a compact subset of X and z_0 any point in X . Then for any $\zeta \in \mathbb{C}, \text{Re } \zeta < 0$, there exist constants $c_3(K, z_0, \zeta) = c_3 > 0$ and $\tilde{c}_3(\zeta) = \tilde{c}_3 > 0$ such that*

$$\left[\int_K |G^0(z, z'; \zeta)|^2 d\sigma(z') \right]^{1/2} \leq c_3 \exp(-\tilde{c}_3 d(z, z_0)). \quad (1.28)$$

Moreover, $\tilde{c}_3(\zeta) \rightarrow \infty$ as $\text{Re } \zeta \rightarrow -\infty$. If K and z_0 are fixed, then $c_3(\zeta) = 0(1)$ as $\text{Re } \zeta \rightarrow -\infty$.

Proof This is an easy consequence of Lemmas 4 and 5. \square

2. Γ -Equivariance

Let Γ be a group of isometries of the plane X . The field \mathbf{B} is invariant with respect to Γ but the Hamiltonian H^0 is not. To obtain the invariance group of H^0 we must consider an extension of Γ , the so-called ‘‘magnetic translation group’’ [15], [3]. Let us recall the construction of this group.

Denote by U the standard unitary representation of Γ in $L^2(X)$:

$$U_\gamma f(z) = f(\gamma^{-1}z), \quad \gamma \in \Gamma, \quad f \in L^2(X). \quad (2.1)$$

If $\gamma^* \mathbf{A} \neq \mathbf{A}$ then $U_\gamma H^0 \neq H^0 U_\gamma$. Nevertheless, $d(\gamma^* \mathbf{A} - \mathbf{A}) = 0$ because $\gamma^* \mathbf{B} = \mathbf{B}$. Hence there exists a function $\omega_\gamma \in C^\infty(X)$ such that

$$d\omega_\gamma = \gamma^* \mathbf{A} - \mathbf{A}. \quad (2.2)$$

Fix for every $\gamma \in \Gamma$ such a function ω_γ ; for $\gamma = 1$ we put $\omega_1 = 0$. Introduce the unitary operator

$$W_\gamma f := \exp(i\omega_\gamma) f, \quad f \in L^2(X), \quad (2.3)$$

and define $T_\gamma^0 = W_\gamma U_\gamma$. Then $T_\gamma^0 H^0 = H^0 T_\gamma^0$ for each $\gamma \in \Gamma$, by (1.10). Unfortunately, the correspondence $\gamma \mapsto T_\gamma^0$ is not a unitary representation of Γ in $L^2(X)$ but only a projective representation in the sense that

$$T_\beta^0 T_\gamma^0 = \Theta(\beta, \gamma) T_{\beta\gamma}^0, \quad \beta, \gamma \in \Gamma, \quad (2.4)$$

where $\Theta(\beta, \gamma) \in \mathbb{C}$, $|\Theta(\beta, \gamma)| = 1$. The family $\Theta(\beta, \gamma)$ has the property

$$\Theta(\gamma_1, \gamma_2) \Theta(\gamma_1 \gamma_2, \gamma_3) = \Theta(\gamma_1, \gamma_2 \gamma_3) \Theta(\gamma_2, \gamma_3), \quad (2.5)$$

i.e. this family is a 2-cocycle of the group Γ with coefficients in $\mathbb{U}(1)$. This cocycle determines a group extension of Γ by $\mathbb{U}(1)$,

$$1 \rightarrow \mathbb{U}(1) \rightarrow M(\Gamma, \Theta) \rightarrow \Gamma \rightarrow 1; \quad (2.6)$$

the group $M(\Gamma, \Theta)$ is called the *magnetic translation group*. An explicit construction of $M(\Gamma, \Theta)$ is the following: $M(\Gamma, \Theta) = \Gamma \times \mathbb{U}(1)$ with multiplication defined by

$$(\gamma_1, \zeta_1)(\gamma_2, \zeta_2) = (\gamma_1 \gamma_2, \Theta(\gamma_1, \gamma_2) \zeta_1 \zeta_2). \quad (2.7)$$

Denote by $[\gamma, \zeta]$ the unitary operator ζT_γ^0 ; the correspondence $(\gamma, \zeta) \mapsto [\gamma, \zeta]$ is then a faithful unitary representation of the group $M(\Gamma, \Theta)$ in $L^2(X)$; we shall denote this representation by T . H^0 is invariant with respect to T ; we will refer to this fact as the *gauge-periodicity* of H^0 .

We need the following lemma.

Lemma 7 *Let L be a linear integral operator in $L^2(X)$ with kernel $L(z, z')$, $z, z' \in X$. The operator L is invariant with respect to T if and only if for any $\gamma \in \Gamma$ the following relation holds a.e.:*

$$\exp(i\omega_\gamma(z)) L(\gamma^{-1}z, z') = \exp(i\omega_\gamma(\gamma z')) L(z, \gamma z'), \quad (2.8)$$

or, equivalently,

$$L(z, z') = \exp[i(\omega_\gamma(z) - \omega_\gamma(\gamma z'))] L(\gamma^{-1}z, \gamma^{-1}z'). \quad (2.9)$$

Proof Let $f \in \mathcal{D}(L)$, $\gamma \in \Gamma$, then we have

$$[\gamma, 1] Lf(z) = \exp(i\omega_\gamma(z)) \int_X L(\gamma^{-1}z, z') f(z') d\sigma(z'), \quad (2.10)$$

$$\begin{aligned}
L[\gamma, 1]f(z) &= \int_X L(z, z') \exp(i\omega_\gamma(z')) f(\gamma^{-1}z') d\sigma(z') = \\
&= \int_X L(z, \gamma z') \exp(i\omega_\gamma(\gamma z')) f(z') d\sigma(z'). \quad (2.11)
\end{aligned}$$

Comparing (2.10) and (2.11), we get (2.8). \square

From now on, we impose the following requirements on the group Γ :

- ($\Gamma 1$) Γ acts properly discontinuously on X ,
- ($\Gamma 2$) the orbit space $\Gamma \backslash X$ is compact.

Fix once and for all a fundamental domain \mathcal{F} of Γ , i.e. a subset $\mathcal{F} \subset X$ such that: (a) $\overline{\mathcal{F}} = \overline{\text{Int}\mathcal{F}}$, (b) $\overline{\mathcal{F}}$ is a compact set, (c) the restriction to \mathcal{F} of the canonical projection $X \rightarrow \Gamma \backslash X$ is a bijective mapping.

To construct a gauge periodic point perturbation of H^0 we choose a finite subset $K \subset \mathcal{F}$ and denote by Λ the orbit of K : $\Lambda = \Gamma \cdot K$. The definition of Λ implies that each element $\lambda \in \Lambda$ has a unique representation of the form $\lambda = \gamma x$, where $\gamma \in \Gamma$ and $x \in K$. Define a unitary representation T^d of $M(\Gamma, \Theta)$ in the discrete space $l^2(\Lambda)$ by the rule

$$T_{(\gamma, \zeta)}^d \varphi(\lambda) = \zeta \exp(i\omega_\gamma(\lambda)) \varphi(\gamma^{-1}\lambda) \quad (2.12)$$

where $(\gamma, \zeta) \in M(\Gamma, \Theta)$, $\varphi \in l^2(\Lambda)$. We denote the operator $T_{(\gamma, \zeta)}^d$ by $[\gamma, \zeta]$, too.

The proof of the following lemma is similar to that of Lemma 7 and is omitted.

Lemma 8 *Let L be a densely defined closed linear operator in the space $l^2(\Lambda)$ having in the standard basis of this space the matrix $(L(\lambda, \mu))_{\lambda, \mu \in \Lambda}$. The operator L is T^d -invariant if and only if for any $\gamma \in \Gamma$ the following relation holds:*

$$\exp(i\omega_\gamma(\lambda)) L(\gamma^{-1}\lambda, \mu) = \exp(i\omega_\gamma(\gamma\mu)) L(\lambda, \gamma\mu), \quad (2.13)$$

or, equivalently,

$$L(\lambda, \mu) = \exp[i(\omega_\gamma(\lambda) - \omega_\gamma(\mu))] L(\gamma^{-1}\lambda, \gamma^{-1}\mu). \quad (2.14)$$

The following lemma is significant for the sequel.

Lemma 9 *There exist constants $c_\Lambda > 0$ and $\tilde{c}_\Lambda > 0$ such that for any $\lambda_0 \in \Lambda$ and $r \in \mathbb{R}, r > 0$, we have*

$$\#\{\lambda \in \Lambda : d(\lambda, \lambda_0) \leq r\} \leq c_\Lambda \exp(\tilde{c}_\Lambda r). \quad (2.15)$$

Proof Let d_Γ be the minimal word length metric with respect to a fixed finite set of generators in Γ . It is known (see [16]) that for some constants $k > 0, \tilde{k} > 0$ we have

$$\#\{\gamma \in \Gamma : d_\Gamma(\gamma, \gamma_0) \leq r\} \leq k \exp(\tilde{k}r), \quad (2.16)$$

where $\gamma_0 \in \Gamma$ is arbitrary. Moreover, there exists a constant $k_1 > 0$ such that

$$d_\Gamma(\gamma_1, \gamma_2) \leq k_1(\inf\{d(\gamma_1 z, \gamma_2 z') : z, z' \in \mathcal{F}\} + 1). \quad (2.17)$$

Let λ_0 be a point of Λ ; then $\lambda_0 = \gamma_0 x_0$, for some $\gamma_0 \in \Gamma, x_0 \in K$. Hence for $r > 0$

$$\begin{aligned} \# \{\lambda \in \Lambda : d(\lambda, \lambda_0) \leq r\} &= \#\{(\gamma, x) \in \Gamma \times K : d(\gamma x, \gamma_0 x_0) \leq r\} \\ &\leq \sum_{x \in K} \#\{\gamma \in \Gamma : d(\gamma x, \gamma_0 x_0) \leq r\}. \end{aligned} \quad (2.18)$$

Now $d(\gamma x, \gamma_0 x_0) \leq r$ implies $d_\Gamma(\gamma, \gamma_0) \leq k_1(r + 1)$, by (2.17), so the proof follows from (2.16). \square

3. Gauge periodic point perturbations

We construct a point perturbation of the operator H^0 in the sense of [17]. Since $\mathcal{D}(H^0) \subset C(X)$ we may define the domain

$$\mathcal{D}(S) := \{f \in \mathcal{D}(H^0) : f(\lambda) = 0 \text{ for } \lambda \in \Lambda\}, \quad (3.1)$$

and the operator S as the restriction of H^0 to $\mathcal{D}(S)$; clearly, S is a symmetric operator in $L^2(X)$. A self-adjoint extension H of S is said to be a *point perturbation of H^0 supported on Λ* if $\mathcal{D}(H) \cap \mathcal{D}(H^0) = \mathcal{D}(S)$. It is an important fact that the point perturbations of H^0 can be described by means of the Krein resolvent formula [4, 5, 17]. To do so, we must find a Hilbert space \mathcal{G} isomorphic to each deficiency subspace of S , and two holomorphic functions

$$B : \mathbb{C} \setminus \text{spec } H^0 \rightarrow \mathcal{L}(\mathcal{G}, L^2(X)),$$

$$Q : \mathbb{C} \setminus \text{spec } H^0 \rightarrow \mathcal{L}(\mathcal{G}, \mathcal{G}),$$

satisfying some conditions which are called Krein's (Γ) - and (Q) - condition [17]; the functions B and Q are then called the Krein Γ - and Q -function, respectively. Fixing a Γ -function and a Q -function, we determine a one-to-one correspondence between point perturbations H of H^0 and (not necessarily bounded) self-adjoint operators A in \mathcal{G} . This correspondence is given by the *Krein resolvent formula* alluded to above:

$$(H - \zeta)^{-1} = (H^0 - \zeta)^{-1} - B(\zeta)[Q(\zeta) + A]^{-1}B^*(\bar{\zeta}). \quad (3.2)$$

We denote by H_A the point perturbation H that corresponds to A via (3.2); the resolvent of H_A will be denoted by $R_A(\zeta)$.

Now we give some explicit description of the Krein Γ - and Q -functions, using Theorem 4 and Proposition 4 from [18] (the proofs of these statements are given in [18] in the case where X is a domain in Euclidean space, but these proofs remain valid in the case of Riemannian manifolds X , too). Denote by \mathcal{G} the space $l^2(\Lambda)$ and by $Q(\zeta)$ the infinite matrix (cf. Lemma 1 for the notation)

$$Q(\lambda, \mu; \zeta) = \begin{cases} G^0(\lambda, \mu; \zeta), & \lambda, \mu \in \Lambda, \quad \lambda \neq \mu, \\ q(\zeta), & \lambda, \mu \in \Lambda, \quad \lambda = \mu. \end{cases} \quad (3.3)$$

Lemma 10 (1) *There exist constants $c_4(\zeta) = c_4 > 0$ and $\tilde{c}_4(\zeta) = \tilde{c}_4 > 0$ such that for $\text{Re } \zeta < 0$ we have*

$$|Q(\lambda, \mu; \zeta)| \leq c_4(\zeta) \exp(-\tilde{c}_4(\zeta)d(\lambda, \mu)), \quad (3.4)$$

whenever $\lambda \neq \mu$. Moreover, $c_4(\zeta) = o(1)$ and $\tilde{c}_4(\zeta) \rightarrow \infty$ as $\text{Re } \zeta \rightarrow -\infty$.

(2) $|Q(\lambda, \lambda; \zeta)| \rightarrow \infty$ as $\text{Re } \zeta \rightarrow -\infty$.

Proof This follows immediately from Lemmas 2 and 3. □

Lemma 11 *There exist a number $E_1 \in \mathbb{R}$ such that for any $\zeta \in \mathbb{C}$ with $\text{Re } \zeta < E_1$, the matrix $Q(\zeta)$ determines a bounded linear operator in $l^2(\Lambda)$ (this operator is denoted by $Q(\zeta)$ as well).*

Proof This is an immediate consequence of Lemmas 9, 10, and A2 (in the appendix). □

For every $\zeta \in \mathbb{C} \setminus \text{spec } H^0$ and $\lambda \in \Lambda$ we denote by $g_\lambda(\zeta)$ the function on X that takes each point $z \in X$ to $G^0(z, \lambda; \zeta)$. It follows from Lemma 5 that $g_\lambda(\zeta) \in L^2(X)$.

Lemma 12 *There exist $\zeta \in \mathbb{C} \setminus \text{spec } H^0$ such that the Gram matrix*

$$(\langle g_\lambda(\zeta) | g_\mu(\zeta) \rangle)_{\lambda, \mu \in \Lambda}$$

determines a bounded operator in $l^2(\Lambda)$.

Proof Let $\text{Im } \zeta \neq 0$ and $\text{Re } \zeta < E_1$, where E_1 is taken from Lemma 11. By the Hilbert resolvent identity, we have for $\lambda \neq \mu$

$$\langle g_\lambda(\zeta) | g_\mu(\zeta) \rangle = (\bar{\zeta} - \zeta)^{-1} [Q(\lambda, \mu; \bar{\zeta}) - Q(\lambda, \mu; \zeta)]. \quad (3.5)$$

Since the diagonal elements of the matrices $Q(\zeta)$ and $\langle g_\lambda(\zeta) | g_\mu(\zeta) \rangle$ are constants at any fixed ζ , the proof follows from Lemma 11. \square

Now we state the main result of this section.

Theorem 1 1. *For any $\zeta \in \mathbb{C} \setminus \text{spec } H^0$ the family $(g_\lambda(\zeta))_{\lambda \in \Lambda}$ is a Riesz basis for its own closed linear hull in $L^2(X)$.*

If $B(\zeta) : l^2(\Lambda) \rightarrow L^2(X)$ is defined by

$$B(\zeta)\varphi = \sum_{\lambda \in \Lambda} \varphi(\lambda)g_\lambda(\zeta), \quad \varphi \in l^2(\Lambda), \quad (3.6)$$

then $B(\zeta)$ is a Krein Γ -function of the pair (S, H^0) .

2. *There exists $E_0 \in \mathbb{R}$ such that for any $\zeta \in \mathbb{C}$ with $\text{Re } \zeta < E_0$ the matrix $Q(\zeta)$ determines a Krein Q -function of the pair (S, H^0) . Hence for any $f \in L^2(X)$ we have*

$$R_A(\zeta)f = R^0(\zeta)f - \sum_{\lambda \in \Lambda} \left(\sum_{\mu \in \Lambda} [Q(\zeta) + A]^{-1}(\lambda, \mu) \right) \langle g_\mu(\zeta) | f \rangle g_\lambda(\zeta). \quad (3.7)$$

Proof In view of Lemmas 11 and 12 the theorem follows immediately from Theorem 4 and Proposition 4 of [18]. \square

We are interested in T -invariant point perturbations, H_A , only. The following proposition provides a necessary and sufficient condition for T -invariance.

Proposition 1 *The operator H_A is T -invariant if and only if the operator A is T^d -invariant.*

To prove this proposition we need the following lemma.

Lemma 13 *For any $\zeta \in \mathbb{C} \setminus \text{spec } H^0$ and $\gamma \in \Gamma$*

$$[\gamma, 1]B(\zeta) = B(\zeta)[\gamma, 1]. \quad (3.8)$$

In other words, $B(\zeta)$ is an intertwining operator for the representations T and T^d .

Proof From (2.8) we have

$$[\gamma, 1]g_\lambda(\zeta) = \exp(i\omega_\gamma(\gamma\lambda))g_{\gamma\lambda}(\zeta), \quad (3.9)$$

hence for $\varphi \in l^2(\Lambda)$ we get

$$\begin{aligned} [\gamma, 1]B(\zeta)\varphi &= \sum_{\lambda \in \Lambda} \varphi(\lambda)[\gamma, 1]g_\lambda(\zeta) \\ &= \sum_{\lambda \in \Lambda} \varphi(\lambda) \exp(i\omega_\gamma(\gamma\lambda))g_{\gamma\lambda}(\zeta) \\ &= \sum_{\lambda \in \Lambda} \varphi(\gamma^{-1}\lambda) \exp(i\omega_\gamma(\lambda))g_\lambda(\zeta) = B(\zeta)[\gamma, 1]\varphi. \end{aligned}$$

□

Proof of Proposition 1: Take $E_0 \in \mathbb{R}$ from Theorem 1, then for $\zeta \in \mathbb{C}$ with $\text{Re } \zeta < E_0$ the operator $Q(\zeta)$ is T^d -invariant, by Lemmas 7 and 8. Consequently, $[Q(\zeta) + A]^{-1}$ is T^d -invariant if and only if A is T^d -invariant. Hence the proposition follows from Lemma 13 and from the fact that $R^0(\zeta)$ is T -invariant. □

In what follows we consider only self-adjoint extensions, H_A , that are invariant with respect to the representation T . For applications in physics, the most interesting case arises if A is a diagonal matrix in the standard basis of the space $l^2(\Lambda)$ [4], [6], [7]; only these operators appear as the limits of Hamiltonians with short-range potentials [4]. In this case, the invariance property

of A implies that there are only finitely many values among the diagonal elements of A . From now on we restrict ourselves to this class of operators. It follows from Lemma A2 and Lemma 8 that these operators are bounded. Moreover, Theorem A1 in the appendix implies the following assertion.

Theorem 2 *There is a number $E_A \in \mathbb{R}$ with the following properties:*

- (1) *for any $\zeta \in \mathbb{C}$ with $\operatorname{Re} \zeta < E_A$, the operator $Q(\zeta) + A$ has a bounded inverse.*
(2) *If $\operatorname{Re} \zeta < E_A$ then there are constants $c_5(\zeta) = c_5 > 0$ and $\tilde{c}_5(\zeta) = \tilde{c}_5 > 0$ such that for any $\lambda, \mu \in \Lambda$*

$$|[Q(\zeta) + A]^{-1}(\lambda, \mu)| \leq c_5 \exp(-\tilde{c}_5 d(\lambda, \mu)). \quad (3.10)$$

Moreover, $c_5(\zeta) = 0(1)$ and $\tilde{c}_5(\zeta) \rightarrow \infty$ as $\operatorname{Re} \zeta \rightarrow -\infty$.

Corollary 1 *The operator H_A is semibounded from below.*

4. The main result

First we recall the notion of the twisted group C^* -algebra $C^*(\Gamma, \Theta)$ of the pair (Γ, Θ) [19], [20], [21]. Let

$$C_0(\Gamma) = \{a : \Gamma \rightarrow \mathbb{C} : a \text{ has finite support}\}. \quad (4.1)$$

Define an associative multiplication in $C_0(\Gamma)$ by the rule

$$(a \cdot b)(\gamma) = \sum_{\beta \in \Gamma} \Theta(\gamma\beta^{-1}, \beta)^{-1} a(\gamma\beta)^{-1} b(\beta), \quad (4.2)$$

and a $*$ -operation by

$$a^*(\gamma) = \Theta(\gamma^{-1}, \gamma) \Theta(1, 1) \overline{a(\gamma^{-1})}. \quad (4.3)$$

There is an injective $*$ -homomorphism, I , of $C_0(\Gamma)$ into the operator algebra $L(l^2(\Gamma))$ that takes each $a \in C_0(\Gamma)$ to $\tilde{a} = Ia$, where

$$\tilde{a}(\varphi)(\gamma) = \sum_{\beta \in \Gamma} \Theta(\gamma\beta^{-1}, \beta)^{-1} a(\gamma\beta^{-1}) \varphi(\beta). \quad (4.4)$$

The twisted group C^* -algebra $C^*(\Gamma, \Theta)$ is defined as the completion of $C_0(\Gamma)$ with respect to the norm $\|a\| := \|Ia\|_{l^2(\Gamma)}$. The algebra $C^*(\Gamma, \Theta)$ has a standard trace, τ , defined by

$$\tau(a) = a(1). \quad (4.5)$$

Now denote by ρ_γ , for $\gamma \in \Gamma$, the operator in $l^2(\Gamma)$ defined by

$$(\rho_\gamma \varphi)(\beta) = \Theta(\beta, \gamma) \varphi(\beta\gamma). \quad (4.6)$$

As usual, let δ_γ denote the element of $l^2(\Gamma)$ with $\delta_\gamma(\beta) = \delta_{\beta\gamma}$.

Lemma 14 (1) *For any $\beta, \gamma \in \Gamma$ we have*

$$\rho_\beta \rho_\gamma = \Theta(\beta, \gamma) \rho_{\beta\gamma}, \quad (4.7)$$

i.e. ρ is a unitary projective representation of Γ in $l^2(\Gamma)$.

(2) *For any $\varphi \in l^2(\Gamma)$,*

$$\varphi(\gamma) = \Theta(1, \gamma)^{-1} (\rho_\gamma \varphi)(1). \quad (4.8)$$

(3)

$$\delta_\gamma = \Theta(\gamma, \gamma^{-1})^{-1} \rho_{\gamma^{-1}} \delta_1. \quad (4.9)$$

(4) *For any $a \in C_0(\Gamma)$ and $\gamma \in \Gamma$*

$$\rho_\gamma \tilde{a} = \tilde{a} \rho_\gamma. \quad (4.10)$$

Proof (1)

$$\begin{aligned} \rho_\beta \rho_\gamma \varphi(\alpha) &= \Theta(\alpha, \beta) \rho_\gamma \varphi(\alpha\beta) = \Theta(\alpha, \beta) \Theta(\alpha\beta, \gamma) \varphi(\alpha\beta\gamma) \\ &= \Theta(\beta, \gamma) \Theta(\alpha, \beta\gamma) \varphi(\alpha\beta\gamma) = \Theta(\beta, \gamma) \rho_{\beta\gamma} \varphi(\alpha). \end{aligned}$$

(2)

$$\rho_\gamma \varphi(1) = \Theta(1, \gamma) \varphi(\gamma).$$

(3)

$$\rho_{\gamma^{-1}} \delta_1(\beta) = \Theta(\beta, \gamma^{-1}) \delta_1(\beta\gamma^{-1}) = \Theta(\gamma, \gamma^{-1}) \delta_\gamma(\beta).$$

(4)

$$\begin{aligned} (\tilde{a} \rho_\gamma \varphi)(\alpha) &= \sum_{\beta} \theta(\alpha\beta^{-1}, \beta)^{-1} a(\alpha\beta^{-1}) (\rho_\gamma \varphi)(\beta) \\ &= \sum_{\beta} \Theta(\alpha\beta^{-1}, \beta)^{-1} \Theta(\beta, \gamma) a(\alpha\beta)^{-1} \varphi(\beta\gamma), \end{aligned} \quad (4.11)$$

$$\begin{aligned}
(\rho_\gamma \tilde{a}\varphi)(\alpha) &= \Theta(\alpha, \gamma) \tilde{a}\varphi(\alpha\gamma) \\
&= \Theta(\alpha, \gamma) \sum_{\beta} \Theta(\alpha\gamma\beta^{-1}, \beta)^{-1} a(\alpha\gamma\beta^{-1}) \varphi(\beta) \\
&= \Theta(\alpha, \gamma) \sum_{\beta} \Theta(\alpha\beta^{-1}, \beta\gamma)^{-1} a(\alpha\beta^{-1}) \varphi(\beta\gamma). \tag{4.12}
\end{aligned}$$

Using (2.5) we obtain

$$\Theta(\alpha\beta^{-1}, \beta) \Theta(\alpha, \gamma) = \Theta(\alpha\beta^{-1}, \beta\gamma) \Theta(\beta, \gamma), \tag{4.13}$$

and statement (4) is proved. \square

Now we define a “canonical” isomorphism

$$\Phi : L^2(X) \rightarrow l^2(\Gamma) \otimes L^2(F) = l^2(\Gamma, L^2(F)) \tag{4.14}$$

by the rule

$$(\Phi f)(\gamma) = r_F([\gamma, 1]f) = r_F(T_\gamma^0 f), \tag{4.15a}$$

where r_F denotes the restriction to F : $r_F f = f|_F$, as in [2].

We also record the explicit form of the inverse:

$$\Phi^{-1} f = \sum_{\gamma \in \Gamma} \chi_\gamma T_\gamma^{0*} e_F(f(\gamma)), \tag{4.15b}$$

where $e_F : L^2(F) \rightarrow L^2(X)$ denotes extension by zero and χ_γ is the characteristic function of $\gamma^{-1}F$. We extend ρ to a projective unitary representation, $\tilde{\rho}$, in $l^2(\Gamma) \otimes L^2(F)$ by the formula

$$\tilde{\rho}_\gamma = \rho_\gamma \otimes 1. \tag{4.16}$$

Then the action of Γ on $l^2(\Gamma, L^2(F))$ is, again, given by (4.6).

Let \mathcal{K} be the algebra of compact operators in the space $L^2(F)$. We denote the tensor product $C^*(\Gamma, \Theta) \otimes \mathcal{K}$ by \mathcal{A} . \mathcal{K} has a natural trace, tr_F , which gives, with (4.5), the canonical trace

$$\tilde{\tau} := \tau \otimes \text{tr}_F$$

on \mathcal{A} . The isomorphism Φ defines a canonical embedding $I_{\mathcal{K}}$ of \mathcal{A} in the C^* -algebra $\mathcal{L}(L^2(X)) = \mathcal{L}(l^2(\Gamma) \otimes L^2(F))$. Denote by $\tilde{\mathcal{A}}$ the image $I_{\mathcal{K}}(\mathcal{A})$.

Lemma 15 (Cf. [2],[21]). *For any $\gamma \in \Gamma$ we have*

$$\Phi[\gamma, 1]\Phi^{-1} = \tilde{\rho}_\gamma. \quad (4.17)$$

Proof Let $f \in L^2(X)$, then

$$(\tilde{\rho}_\gamma \Phi f)(\beta) = \Theta(\beta, \gamma)\Phi f(\beta\gamma) = r_F(\Theta(\beta, \gamma)[\beta\gamma, 1]f), \quad (4.18)$$

$$(\Phi[\gamma, 1]f)(\beta) = r_F([\beta, 1][\gamma, 1]f).$$

Taking into account the multiplication rule (2.7), we get the result. \square

Now we denote by

$$\mathcal{M}(\Gamma, \Theta) := \{B \in \mathcal{L}(l^2(\Gamma) \otimes L^2(F)) : B\tilde{\rho}_\gamma = \tilde{\rho}_\gamma B \text{ for } \gamma \in \Gamma\} \quad (4.19)$$

the commutant of $(\tilde{\rho}_\gamma)_{\gamma \in \Gamma}$. From Lemmas 14, (4) and 15 we obtain

$$\tilde{\mathcal{A}} \subset \mathcal{M}(\Gamma, \Theta). \quad (4.20)$$

Besides, Lemma 15 implies that

$$R_A(\zeta) \in \mathcal{M}(\Gamma, \Theta), \quad \zeta \in \mathbb{C} \setminus \text{spec } H_A, \quad (4.21)$$

where R_A is given by (3.7) (recall that we consider only T -invariant operators H_A).

Now, following [2] we define the Fourier coefficients for $B \in \mathcal{M}(\Gamma, \Theta)$. For any $\gamma \in \Gamma$ the Fourier coefficient $\hat{B}(\gamma)$ is the operator in $L^2(F)$ given by

$$\hat{B}(\gamma)(u) = \tilde{\rho}_\gamma B(\delta_1 \otimes u)(1), \quad u \in L^2(F). \quad (4.22)$$

Lemma 16 (Cf. [2], [21]).

(1) *For any $B \in \mathcal{M}(\Gamma, \Theta)$ and $f \in l^2(\Gamma, L^2(F))$ we have*

$$Bf(\gamma) = \Theta(1, \gamma)^{-1} \sum_{\beta} \Theta(\beta, \beta^{-1})^{-1} \Theta(\gamma, \beta^{-1}) \hat{B}(\gamma\beta^{-1})(f(\beta)). \quad (4.23)$$

(2) *Let $\gamma \in \Gamma, B \in \mathcal{L}(L^2(F))$, then*

$$\widehat{(\tilde{\delta}_\gamma \otimes B)}(\beta) = \Theta(1, \beta)\Theta(\beta, 1)^{-1}(\delta_\gamma(\beta) \otimes B). \quad (4.24)$$

(3) *For $B \in \mathcal{M}(\Gamma, \Theta)$, we have*

$$\|B\| \leq \sum_{\gamma \in \Gamma} \|\hat{B}(\gamma)\|. \quad (4.25)$$

Proof (1) Let $f \in l^2(\Gamma, L^2(F))$. Then, from (4.9)

$$\begin{aligned} f &= \sum_{\beta \in \Gamma} \delta_\beta \otimes f(\beta) \\ &= \sum_{\beta \in \Gamma} \Theta(\beta, \beta^{-1})^{-1} \tilde{\rho}_{\beta^{-1}} (\delta_1 \otimes f(\beta)). \end{aligned} \quad (4.26)$$

Hence

$$Bf = \sum_{\beta} \Theta(\beta, \beta^{-1})^{-1} \tilde{\rho}_{\beta^{-1}} B(\delta_1 \otimes f(\beta)), \quad (4.27)$$

$$(Bf)(1) = \sum_{\beta} \Theta(\beta, \beta^{-1})^{-1} \tilde{\rho}_{\beta^{-1}} B(\delta_1 \otimes f(\beta))(1). \quad (4.28)$$

In view of Lemma 14, (2) and (1), we have

$$\begin{aligned} (Bf)(\gamma) &= \Theta(1, \gamma)^{-1} \sum_{\beta} \Theta(\beta, \beta^{-1})^{-1} \tilde{\rho}_\gamma \tilde{\rho}_{\beta^{-1}} B(\delta_1 \otimes f(\beta))(1) \\ &= \Theta(1, \gamma)^{-1} \sum_{\beta} \Theta(\beta, \beta^{-1})^{-1} \Theta(\gamma, \beta^{-1}) \tilde{\rho}_{\gamma\beta^{-1}} B(\delta_1 \otimes f(\beta))(1) \\ &= \Theta(1, \gamma)^{-1} \sum_{\beta} \Theta(\beta, \beta^{-1})^{-1} \Theta(\gamma, \beta^{-1}) \hat{B}(\gamma\beta^{-1})(f(\beta)), \end{aligned}$$

and (4.23) is proved.

(2) By (4.22),

$$\begin{aligned} \widehat{(\tilde{\delta}_\gamma \otimes B)}(\beta)(u) &= \tilde{\rho}_\beta(\tilde{\delta}_\gamma \otimes B)(\delta_1 \otimes u)(1) \\ &= \Theta(1, \beta)(\tilde{\delta}_\gamma \otimes B)(\delta_1 \otimes u)(\beta) \\ &= \Theta(1, \beta) \tilde{\delta}_\gamma(\delta_1)(\beta) \otimes B(u). \end{aligned} \quad (4.29)$$

On the other hand, by (4.4),

$$\begin{aligned} \tilde{\delta}_\gamma(\delta_1)(\beta) &= \sum_{\alpha \in \Gamma} \Theta(\beta\alpha^{-1}, \alpha)^{-1} \delta_\gamma(\beta\alpha^{-1}) \delta_1(\alpha) \\ &= \Theta(\beta, 1)^{-1} \delta_\gamma(\beta). \end{aligned} \quad (4.30)$$

Substituting (4.30) in (4.29), we obtain (4.24).

(3) The proof is similar to the proof of Lemma 3 in [2]. \square

Lemma 17 [2] *Let $B \in \mathcal{M}(\Gamma, \Theta)$. If for all $\gamma \in \Gamma$, the operator $\hat{B}(\gamma)$ is compact and $\sum_{\gamma \in \Gamma} \|\hat{B}(\gamma)\| < \infty$, then $B \in \tilde{\mathcal{A}}$.*

Proof Cf. the proof of the Corollary of Lemma 3 in [2]. \square

Now let us state the main result of the paper.

Theorem 3 *Let A be a T^d -invariant self-adjoint operator in $l^2(\Lambda)$ with diagonal matrix. Then for any $\zeta \in \mathbb{C} \setminus \text{spec } H_A$ the resolvent $R_A(\zeta)$ belongs to $\tilde{\mathcal{A}}$.*

Proof First note that it is enough to prove that $R_A(E) \in \tilde{\mathcal{A}}$ for E ranging over some semi-axis $(-\infty, E_0)$. Indeed, let this property be satisfied but $R_A(\zeta) \notin \tilde{\mathcal{A}}$ for some $\zeta_0 \in \mathbb{C} \setminus \text{spec } H_A$. By the Hahn-Banach Theorem, there exists a continuous linear functional $\Psi \in \mathcal{L}(L^2(X))'$ such that $\Psi(B) = 0$ for each $B \in \tilde{\mathcal{A}}$ and $\Psi(R_A(\zeta_0)) \neq 0$. But this contradicts the analyticity of the function $\zeta \mapsto \Psi(R_A(\zeta))$.

Now note that $R^0(\zeta) \in \tilde{\mathcal{A}}$ for any $\zeta \in \mathbb{C} \setminus \text{spec } H^0$. In fact, the proof given in [2], [3] for the fact that $e^{-tH_0} \in \tilde{\mathcal{A}}$ for each $t \geq 0$ carries over to the case at hand without major changes. Using the Laplace transform, we show that $R^0(E) \in \tilde{\mathcal{A}}$ for any $E < 0$ and hence for each $\zeta \in \mathbb{C} \setminus \text{spec } H^0$.

Thus, it remains to prove that $V(E) := R^0(E) - R_A(E) \in \tilde{\mathcal{A}}$ if E ranges over some semi-axis $(-\infty, E_0)$.

By Theorem 1,2) we can find a number $E_0 \in \mathbb{R}$ such that for any $f \in L^2(X)$ we have

$$V(\zeta)f = \sum_{\lambda \in \Lambda} \left(\sum_{\mu \in \Lambda} M(\lambda, \mu; \zeta) \langle g_\mu(\bar{\zeta}) | f \rangle \right) g_\lambda(\zeta), \quad (4.31)$$

whenever $\text{Re } \zeta < E_0$. Here we have written

$$M(\lambda, \mu; \zeta) = [Q(\zeta) + A]^{-1}(\lambda, \mu). \quad (4.32)$$

Moreover, by (3.10) there are constants c_0 and $\tilde{c}_0(\zeta)$ such that

$$|M(\lambda, \mu; \zeta)| \leq c_0 \exp(-\tilde{c}_0(\zeta)d(\lambda, \mu)); \quad (4.33)$$

c_0 is independent of ζ , and $\tilde{c}_0(\zeta) \rightarrow \infty$ as $\text{Re } \zeta \rightarrow -\infty$. So we can suppose that

$$\tilde{c}_0(\zeta) > 3\tilde{c}_\Lambda, \quad (4.34)$$

where \tilde{c}_Λ is the constant from Lemma 9.

To see that $V(\zeta)$ is in $\tilde{\mathcal{A}}$ it is enough, in view of the definition and Lemma 17, to show that

$$\widehat{V(\zeta)}(\gamma) \text{ is compact in } L^2(F), \quad (4.35a)$$

and

$$\sum_{\gamma \in \Gamma} \|\widehat{V(\zeta)}(\gamma)\| < \infty. \quad (4.35b)$$

We compute with (4.22), (4.31), (4.16), and (4.6), for $v \in L^2(F)$,

$$\widehat{V(\zeta)}(\gamma) = \Theta(1, \gamma) \exp(i\omega\gamma) \sum_{\lambda, \mu \in \Lambda} M(\lambda, \mu; \bar{\zeta}) \langle g_\mu(\bar{\zeta}), \Phi^{-1}(\delta_1^v) \rangle r_F(g_\lambda(\zeta) \circ \gamma^{-1}). \quad (4.36)$$

Using (4.15b) and Lemma 6 we obtain

$$\begin{aligned} |\langle g_\mu(\bar{\zeta}), \Phi^{-1}(\delta_1^v) \rangle| &= |\langle g_\mu(\bar{\zeta}), e_F(v) \rangle| \\ &= \left| \int_F G^0(z', \lambda; \bar{\zeta}) v(z') d\sigma(z') \right| \\ &\leq c_3 e^{-\tilde{c}_3(\zeta) d(\lambda, \kappa_0)} \|v\|_{L^2(F)}, \end{aligned} \quad (4.37a)$$

where we may, again, assume that

$$\tilde{c}_3(\zeta) > 3\tilde{c}_\Lambda. \quad (4.38)$$

Finally, we see with Lemmas 6 and 7 that

$$\begin{aligned} \|r_F(g_\lambda(\zeta) \circ \gamma^{-1})\|_{L^2(F)} &\leq \left[\int_F |G^0(\lambda, \gamma^{-1}(z'); \zeta)|^2 d\sigma(z') \right]^{1/2} \\ &\leq c_3 e^{-\tilde{c}_3(\zeta) d(\gamma\lambda, \kappa_0)}. \end{aligned} \quad (4.39)$$

Now we write $\lambda = \alpha\kappa, \mu = \beta\kappa'$ with $\kappa, \kappa' \in K$ and $\alpha, \beta \in \Gamma$, and further $\beta =: \alpha\beta', \gamma' := \gamma\alpha$ and find

$$\begin{aligned} &\sum_{\kappa, \kappa' \in K} \sum_{\alpha \in \Gamma} \sum_{\beta', \gamma' \in \Gamma} e^{-3\tilde{c}_\Lambda(d(\alpha\kappa, \alpha\beta'\kappa') + d(\alpha\kappa, \kappa_0) + d(\gamma\alpha\kappa, \kappa_0))} \\ &= \sum_{\substack{\kappa, \kappa' \in K \\ \alpha \in \Gamma}} e^{-3\tilde{c}_\Lambda d(\alpha\kappa, \kappa_0)} \sum_{\beta', \gamma' \in \Gamma} e^{-3\tilde{c}_\Lambda(d(\beta'\kappa', \kappa) + d(\gamma'\kappa, \kappa_0))} \\ &\leq c_4 < \infty, \end{aligned} \quad (4.40)$$

by Lemma 9 and A1. Now we use (4.38) and (4.39) in (4.36) to see, similarly, that the sum is norm convergent; since all summands are operators of rank one, the compactness of $\widehat{V(\zeta)}(\gamma)$ follows. The summation (4.40) proves (4.35b) and the theorem is proved. \square

Corollary 2 *Let $E_1, E_2 \in \mathbb{R} \setminus \text{spec } H_A$ and $E_1 \leq E_2$. Then the spectral projector $P_{[E_1, E_2]}$ of the operator H_A belongs to $\tilde{\mathcal{A}}$.*

Proof Indeed, we may write for $E \notin \text{spec } H_A$: $P_{[E_1, E_2]} = \psi(R_A(E))$ where ψ is a continuous function with compact support. \square

Fix now a number $E' \in \mathbb{R}$ such that $E' < \inf \text{spec } H_A$, and consider the function

$$N(E) = \begin{cases} \tau P_{[E', E]}, & E \geq E' \\ 0, & E < E'. \end{cases}$$

Then this function is independent of the choice of E' . The values of $N(E)$ are constant on each gap of the spectrum of H_A . Therefore these values label in a natural way the gaps of H_A [22].

Corollary 3 (*Gap Labelling Theorem*). *The value of the function $N(E)$ on a gap of $\text{spec } H_A$ belongs to $\tau^*(K_0 C^*(\Gamma, \Theta))$, a countable set of real numbers (here $K_0 \mathcal{B}$ denotes the K_0 -group of a C^* -algebra \mathcal{B}).*

Recall that the pair (Γ, Θ) is said to have the *Kadison property* if there exist a constant $c_K > 0$ such that $\tau(P) \geq c_K$ for every nonzero self-adjoint projection P in $C^*(\Gamma, \Theta) \otimes \mathcal{K}$. It now follows as in [2], [16], [21]:

Corollary 4 *If the pair (Γ, Θ) has the Kadison property then the spectrum of H_A has band structure.*

Appendix

In this appendix, we provide some general results concerning a discrete metric space, Λ , with metric d . We suppose that the following condition on the “volume growth” of metric balls is fulfilled (cf. (2.15)):

There are constants $c_\Lambda > 0$ and $\tilde{c}_\Lambda > 0$ such that for any $\lambda_0 \in \Lambda$ and any $r \in (0, \infty)$ we have

$$\#\{\lambda \in \Lambda : d(\lambda, \lambda_0) \leq r\} \leq c_\Lambda \exp(\tilde{c}_\Lambda r). \quad (\text{A1})$$

Lemma A1 *Let $\phi : \Lambda \rightarrow \mathbb{C}$ be a function such that*

$$|\phi(\lambda)| \leq c \exp[-(1 + \delta)\tilde{c}_\Lambda d(\lambda, \mu)], \quad (\text{A2})$$

where c and δ are positive constants and μ is any fixed element of Λ . Then

$$\sum_{\lambda \in \Lambda} |\phi(\lambda)| \leq c \cdot c_\Lambda \cdot \delta^{-1}. \quad (\text{A3})$$

Proof See [23]. □

Lemma A2 (*Schur's test*). Let $(L(\lambda, \mu))_{\lambda, \mu \in \Lambda}$ be an infinite matrix such that for some $c' > 0$ we have

$$\sup_{\mu \in \Lambda} \sum_{\lambda \in \Lambda} |L(\lambda, \mu)| \leq c', \quad \sup_{\lambda \in \Lambda} \sum_{\mu \in \Lambda} |L(\lambda, \mu)| \leq c'. \quad (\text{A4})$$

Then the matrix $L(\lambda, \mu)$ determines a bounded linear operator L in the space $l^2(\Lambda)$ and

$$\|L\| \leq c'. \quad (\text{A5})$$

Proof See [23], [14]. □

Theorem A1 Let $(K_n)_{n \geq 0}$ be a sequence of bounded linear operators in the space $l^2(\Lambda)$, having the matrices $(K_n(\lambda, \mu))_{\lambda, \mu \in \Lambda}$ with respect to the standard basis $l^2(\Lambda)$. Suppose that the following conditions are satisfied:

(1) if $\lambda \neq \mu$, then

$$|K_n(\lambda, \mu)| \leq a \exp(-b_n d(\lambda, \mu)), \quad (\text{A6})$$

where a is independent of n and $\lim_{n \rightarrow \infty} b_n = \infty$;

(2)

$$\inf_{\lambda \in \Lambda} |K_n(\lambda, \lambda)| \rightarrow \infty \text{ as } n \rightarrow \infty. \quad (\text{A7})$$

Then, for any $\alpha \in (0, 1)$ there is $n_0 \in \mathbb{N}$ such that for every $n \geq n_0$ the operator K_n has a bounded inverse, $L_n := K_n^{-1}$. The matrix $(L_n(\lambda, \mu))_{\lambda, \mu \in \Lambda}$ of this operator admits the estimate

$$|L_n(\lambda, \mu)| \leq 2c_n \exp(-\alpha b_n d(\lambda, \mu)), \quad (\text{A8})$$

where

$$c_n = \left(\inf_{\lambda \in \Lambda} |K_n(\lambda, \lambda)| \right)^{-1}. \quad (\text{A9})$$

Proof We introduce operators D_n, S_n by

$$D_n(\lambda, \mu) := K_n(\lambda, \mu)\delta_{\lambda\mu}, S_n(\lambda, \mu) := K_n(\lambda, \mu) - D_n(\lambda, \mu).$$

Moreover, we fix $\alpha \in (0, 1)$ and determine $n_0 \in \mathbb{N}$ such that for $n \geq n_0$

$$\inf_{\lambda \in \Lambda} |K_n(\lambda, \lambda)| \geq 1 \quad \text{and} \quad (1 - \alpha)b_n \geq 2\tilde{c}_\Lambda. \quad (\text{A10})$$

Then D_n^{-1} and S_n are bounded in $l^2(\Lambda)$ in view of Lemmas A1 and A2 and, clearly,

$$\begin{aligned} K_n &= D_n(1 + D_n^{-1}S_n), \\ L_n &= K_n^{-1} = (1 + D_n^{-1}S_n)^{-1}D_n^{-1} =: T_n D_n^{-1}. \end{aligned}$$

Thus, the theorem follows if we prove the estimate

$$|T_n(\lambda, \mu)| \leq 2 \exp(-\alpha b_n d(\lambda, \mu)). \quad (\text{A11})$$

Now

$$T_n(\lambda, \mu) = \sum_{j \geq 0} (-1)^j (D_n^{-1}S_n)^j(\lambda, \mu),$$

and it is enough to show that, with some constant $A \geq 1$,

$$|(D_n^{-1}S_n)^j(\lambda, \mu)| \leq (aAc_n)^j e^{-\alpha b_n d(\lambda, \mu)}, \quad (\text{A12})$$

since $c_n \rightarrow 0$ as $n \rightarrow \infty$. This estimate is obvious for $j = 0, 1$; inductively, we find with (A5) and (A7)

$$\begin{aligned} |(D_n^{-1}S_n)^{j+1}(\lambda, \mu)| &= \left| \sum_{\nu \in \Lambda} K_n(\lambda, \nu)^{-1} S_n(\lambda, \nu) (D_n^{-1}S_n)^j(\nu, \mu) \right| \\ &\leq \sum_{\nu \in \Lambda} c_n a e^{-b_n d(\lambda, \nu)} (Aac_n)^j e^{-\alpha b_n d(\nu, \mu)} \\ &\leq A^j (c_n a)^{j+1} e^{-\alpha b_n d(\lambda, \mu)} \sum_{\nu \in \Lambda} e^{-(1-\alpha)b_n d(\lambda, \nu)}. \end{aligned}$$

In view of (A8) and Lemma A1, the last sum has the bound c_Λ . The assertion (A10) for $j + 1$ follows if we put $A := \max\{1, c_\Lambda\}$. \square

Remark Theorem A1 strengthens Theorem 2.1 from [23]. The estimate there is insufficient for proving our results.

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