# Universal Pion Freeze-out Phase-Space Density

D. Ferenc, B. Tomášik, U. Heinz, Inst. f. Theor. Physik, Universität Regensburg

G. Bertsch has indicated [1] a possibility to measure the pion freeze-out phase-space density and thereby test the local thermal equilibrium in a pion source. In case of thermal equilibrium at temperature T, identical pions of energy Ewould follow the Bose-Einstein distribution

$$f = \frac{1}{e^{\frac{E}{T}} - 1}. (1)$$

$$\langle f \rangle (p_T, y) = \frac{\frac{\sqrt{\pi}}{2} \frac{\sqrt{\lambda_{\text{dir}}(p_T, y)}}{E_p T_{\text{eff}}^2(y)} \exp\left(-\frac{p_T}{T_{\text{eff}}(y)}\right) \frac{dn^-}{dy}(y)}{R_s(p_T, y) \sqrt{R_o^2(p_T, y) R_l^2(p_T, y) - R_{ol}^4(p_T, y)}}.$$
(2)

An average of this function over different phase-space regions is the quantity to be measured. When the  $p_T$ -spectrum is parameterized by an exponential with  $T_{\rm eff}(y)$  being the inverse slope parameter, averaging over the spatial coordinates yields  $\langle f \rangle (p_T,y) = \frac{\frac{\sqrt{\pi}}{2} \frac{\sqrt{\lambda_{\rm dir}(p_T,y)}}{E_p T_{\rm eff}^2(y)} \exp\left(-\frac{p_T}{T_{\rm eff}(y)}\right) \frac{dn^-}{dy}(y)}{R_s(p_T,y)\sqrt{R_o^2(p_T,y)R_l^2(p_T,y)-R_o^4(p_T,y)}}. \tag{2}$  This equation comprises information from essentially two different classes of experimental results: the single particle momentum spectra  $(dn^-/dy,T_{eff})$ , and the two pion Bose-Einstein correlations  $(R_s,R_o,R_l,R_{ol},\lambda_{dir})$ . We have calculated  $\langle f \rangle (p_T,y)$  for the S-S, S-Cu, S-Ag, S-Au, S-Pb and Pb-Pb data from the experiments NA35 [2], NA49 [3], NA44 [4], and for the  $\pi$ -p data from the NA22 experiment [5] at CERN-SPS. From the results for  $\langle f \rangle$  as a function of  $p_T$ , presented in Fig. 1, one may conclude:

1. universal phase-space density

All the nuclear collision data from the SPS in Fig. 1 are indistinguishably similar, in spite of a factor of  $\sim$  10 dif-

indistinguishably similar, in spite of a factor of  $\sim 10$  difference in multiplicity density.

#### 2. agreement with Bose-Einstein distribution

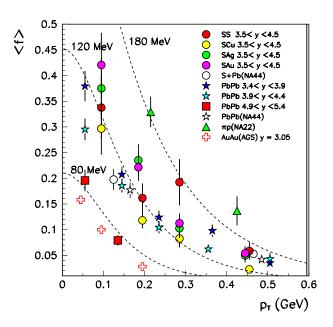
Using simultaneously pion spectra and pion correlations NA49 has disentangled thermal motion from collective expansion [3]. Taking the NA49 result for the local freezeout temperature T = 120 MeV as the **only** parameter in Eq. (1) one indeed finds good agreement with the data. This is consistent with the thermal nature of the pion source, and inconsistent with the presence of a hypothetic pion condensate.

### 3. radial flow

Looking in more detail, one finds that the data indicate a somewhat slower decrease with increasing  $p_T$  than the Bose-Einstein curve. This is most likely due to radial collective expansion which adds extra transverse momentum to particles, i.e. the local  $\langle f \rangle$  values appear in the measurement at a  $p_T$  that is higher than the local  $\langle p_T \rangle$  in the source reference frame. A detailed study is under way.

#### 4. rapidity dependence

A certain departure from the universal scaling is seen for the data at rapidities close to the projectile rapidity, both at AGS and SPS; moreover, the two results are consistent.



Phase-space density as a function of  $p_T$  for Figure 1: different data sets. Heavy-ion data from SPS are indistinguishably similar, although they span over an order of magnitude range in multiplicity. A Bose-Einstein function (Eq. (1)) is superimposed with the three choices of the local freeze-out temperature: 80 MeV, 120 MeV and 180 MeV.

# 5. high temperature decoupling in $\pi - p$ collisions In contrast to freeze-out in nuclear collisions which takes place in two steps (chemical at $T \simeq 170\text{-}180 \text{ MeV}$ and thermal at $T \simeq 120 \text{ MeV}$ ), pion production in $\pi$ -p collisions [5] is essentially immediate, without the second evolution stage, and therefore freeze-out temperatures of around 180 MeV should be expected. The data [5] are indeed consistent with this expectation, as seen in Fig. 1.

## References

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