

A Single Pulse Beam Emittance Measurement for the CERN Heavy Ion Linac.

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Abstract

A new device for transverse emittance measurement has been installed in the 4.2 MeV/u filter region of the CERN Heavy Ion Linac (Linac 3). It allows to obtain pulse-to-pulse (every 1.2 sec) visualisation of the Linac 3 beam parameters in order to tune the machine and to match the beam for injection into the first circular accelerator, the PS Booster.

The system is based on the "multi-slit" technique similar to the well-known "pepper pot" method.

A plate with a series of horizontal or vertical slits is placed in the beam, defining positions in the phase plane. Particles pass through the slits and drift to a scintillator screen where they produce light. The screen is looked at by an externally triggered high resolution CCD camera. For each slit position the light intensity distribution, in the limit of infinitesimal slit aperture, is proportional to the angle distribution of the particles and therefore, provides the angular distribution in the phase plane. The video signal from the camera is digitised and the result stored in a 512 x 512 x 8 frame store.

An application program visualising slit images (raw data) as well as the calculated emittance ellipses has been developed. This program also delivers numerical values for the rms-emittance and the Twiss parameters.

The paper discusses the method, the prototype realisation, the data treatment and first experimental results obtained during the running-in of Linac 3.

1. Introduction

The construction and assembly of a new injector for heavy ion physics at CERN has been completed during 1994 [1]. It consists of an ECR ion source producing a $Pb27^+$ beam. These ions are accelerated through an RFQ and an interdigital-H Linac. The beam is stripped and the charge state of $Pb53^+$ selected in a filter line[3] following the Linac (Fig.1).

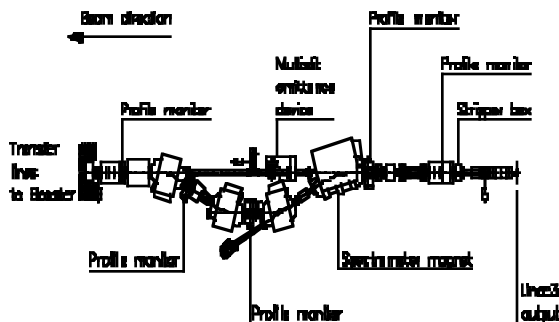


Fig.1 Beam measurement devices in filter line.

Theoretical values for accelerator output parameters are [4]: pulse width of 400 μ s with a repetition frequency of 1.2 s, nominal current of $Pb27^+$ 65 μ A, (20% of which is converted into $Pb53^+$), which gives 6×10^9 particles of $Pb27^+$ per pulse. Transverse emittances are about 0.8π mm mrad normalised.

In the context of this project, our task was to develop a device for measuring transverse emittances at each pulse, giving the possibility to check the pulse to pulse stability of the accelerator.

2. The multislit method

A plate with 10 slits is inserted in the beam at the measurement position, and particles passing through the slits are detected by a scintillator screen placed 600 mm after the plate. The image on the screen is read into a 512 x 512 pixel array by a single-shot CCD camera synchronised with the beam. Each pixel is digitised by an 8 bit ADC (256 intensity levels), and stored in a frame grabber. The plate and the scintillator screen are tilted at 45° to allow camera visualisation(Fig.2).

Image on scintillator screen

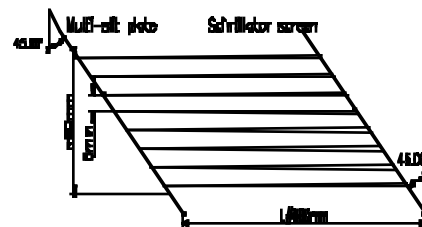


Fig.2 The method.

X or Y intensity profiles are calculated and displayed in a 512 channel projection graph (Fig.3). Phase plane reconstruction is based on these projections.

Each peak in this graph is associated with a slit position and therefore with a position in the phase plane. The width of these lines determines the angular extension (e.g. α_1 - α_2 in Fig.4) for the corresponding position.

Fig.3 Projections giving x and y intensity profiles. Different choices of the central peak do not influence emittance calculations, but care must be taken to avoid overlapping of light spots on the screen.

3. The prototype realisation

3.1 The multi-slit

The multislit is a stainless steel plate 5 mm thick, with a series of parallel slits. The number of slits can be 11, 17 or 21. The slit acceptance is ± 30 mrad. Each slit is 0.3 mm in width. The finite slit dimension limits the precision in position and gives a fixed over-estimation in angle of less than 0.5 mrad.

The distance between the slits limits the maximal beam angles that can be measured. In order to be ready to measure beams of different emittances, we prepared 3 different kinds of plates, with a distance between slits of 6, 4 or 3 mm. Most of the time we used the 6 mm plate that allows the measurement of beams of max. ± 5 mrad (Fig.4).

If the camera captures the whole scintillator screen on the CCD array, angular resolution is better than 0.2 mrad (and in this case the precision in angle is 4%).

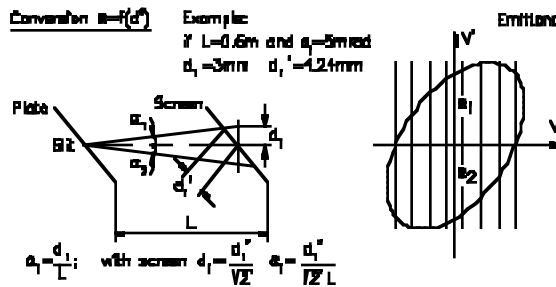


Fig.4 Some details of slits and phase plane reconstruction.

3.2 The scintillator screen

The main problem to overcome was to find a scintillator screen suitable for low intensity and low energy lead ions (compared to protons). Two different 60 x 60 x 1 mm fluorescent screens have been used: Cromox type 6 [5] (~700 nm emission wavelength, 3 ms fluorescence decay time) and the CsI thallium doped type (~500 nm emission wavelength, 0.5 μ s fluorescence decay time). Tests showed that the Cromox screen was more suitable for our purposes, since less subject to saturation at beam intensity levels encountered (Fig.5)

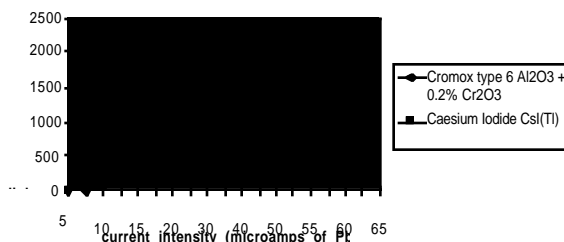


Fig.5 Light intensity output vs. beam intensity for

3.3 The CCD camera

Identical vertical and horizontal resolutions are needed, without rotating the camera. We chose a full-frame non-interlaced CCD camera, with 15 μ m x 15 μ m square pixels and analogue video output. The number of active pixels is 512 x 512 and the CCD array dimension is 7.7 mm x 7.7 mm.

An adjustable active low TTL pulse is needed to define the CCDs integration time (during which image acquisition is done). The camera accepts pulse lengths from a minimum of 150 μ s to a maximum of 1s (due to CCD thermal noise). The image readout takes 33 ms.

3.4 The frame grabber

We chose a commercial frame grabber with a modular structure, consisting of a mother board and two plug-in daughter modules. The mother board contains 3 memory blocks of 1024 x 512 bytes each for storage of the video frames. The two daughter modules are the acquisition module (8 bits ADC resolution), and the display module, regenerating the video signal from the digitised data. The board works on the VME bus, and can be configured to accept synchronisation pulses and data from non-standard video cameras. In particular, the image acquisition can be externally triggered.

In our configuration the camera is used as master and the board as slave. Three signals are needed to synchronise image acquisition: pixel clock, line blank and vertical blank.

3.5 Timing

A programmable Camac module gives precise timing to the system. It receives a trigger pulse from the accelerator timing system, the valid beam signal (VB0), and drives the camera strobe and the trigger to software real-time task. The camera strobe pulse can be delayed, to allow displacement along the beam pulse.

In most of our measurements we use short camera strobe times (max. 300 μ s) to avoid CCD saturation (Fig. 6), and a delay of 200 μ s with respect to VB0.

3 ms 150 μ s

Fig.6 CCD saturation with 3 ms strobe time (compared to 150 μ s).

3.6 Optical and mechanical aspects

The scintillator screen is mounted in a standard CERN/PS vacuum tank. Since these tanks are not explicitly designed for emittance measurements, the camera could not be placed perpendicular to the screen, resulting in problems of perspective and depth of field. Two MgF₂ mirrors facing each other inserted in the vacuum tank by means of a newly designed mechanical support solved the problems.

4. Controls

4.1 Hardware

The detector is remotely controlled through the

with strong support and useful suggestions during the critical debugging phase of the detector.

References:

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