Mirror Symmetry, Mirror Map and Applications to Complete Intersection Calabi-Yau Spaces

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Abstract:
We extend the discussion of mirror symmetry, Picard-Fuchs equations, instanton corrected Yukawa couplings and the topological one loop partition function to the case of complete intersections with higher dimensional moduli spaces. We will develop a new method of obtaining the instanton corrected Yukawa couplings through a study of the solutions of the Picard-Fuchs equations. This leads to closed formulas for the prepotential for the Kähler moduli fields induced from the ambient space for all complete intersections in nonsingular weighted projective spaces. As examples we treat part of the moduli space of the phenomenologically interesting three generation models which are found in this class. We also apply our method to solve the simplest model in which topology change was observed and discuss examples of complete intersections in singular ambient spaces.

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1. Introduction:

Complete intersection Calabi-Yau manifolds embedded in products of projective spaces ($CICY$) are the most prominent candidates for the compactification of the heterotic string. They have been intensively studied. Let us point out briefly the main results; a detailed account from a physical point of view can be found in [1]. Using the $c_1 = 0$ condition and curve identities it was recognized in [2] that all CICY's can be described by finitely many configurations of polynomials in products of projective spaces. Each configuration leads to a family of Calabi-Yau spaces whose generic member is smooth. By a computer classification 7868 configurations with Euler numbers between -200 and 0 were found in [3] [4]. In [5] all their Hodge numbers were calculated. There occur 265 different combinations. Application of a theorem of Wall [6] which states that the homotopy types of Calabi-Yau three-folds X can be classified by their Hodge numbers, their topological triple couplings $K_{ijk}^0 \equiv \int_X J_i \wedge J_j \wedge J_k$ and $c_2 \cdot J_i \equiv \int_X c_2 \wedge J_i$, reveals that there are at least 2590 topologically different models in this class [7]. Green and Hübsch have shown in [8] that all families of CICY's are connected by the process, already described in [3], of contracting a family to a nodal configuration and performing a small resolution of the latter. Certain quotients of them by discrete groups were constructed in [9] and [10] which have Euler number $\chi = -6$. Two of them, discussed below, have a nontrivial fundamental group ($\pi_1(X) = \mathbb{Z}_3$) and give rise to heterotic string compactifications with three generations and a natural option to break the E_6 gauge group by Wilson lines.

In this paper we want to extend the analysis of the Picard-Fuchs equations, the construction of the mirror map and the calculation of the instanton corrected Yukawa couplings, which was performed for one moduli hypersurfaces in a (weighted) projective space in [11] [12] [13] [14] and generalized to higher dimensional moduli spaces of general hypersurfaces in toric varieties in [15] [16] to the class of complete intersections in products of weighted projective spaces. It is natural to focus on the derivation of the prepotentials for the complex structure and the Kähler structure deformations, which encode all information of theses two (topological) subsectors of the theory at tree level. The knowledge of the prepotentials is quite relevant for the low energy phenomenology of the subsector of the moduli- and the $\mathbf{27}$, $\overline{\mathbf{27}}$ matter fields in the effective field theory.

We will show that there is an easily exactly solvable subsystem of the complex structure deformations, namely the one which is, by the mirror map, associated to the quantum cohomology of the subset of the elements in $H^{1,1}(X, \mathbb{Z})$ induced from the Kähler forms of the projective spaces. We focus on this subsystem and present results in a closed form for the instanton corrected cohomology of this system. In general the moduli space is enlarged due to exceptional divisors coming from the resolution of possible singularities. This situation we have treated quite generally for hypersurfaces in [16]. Here we will give two simple but illustrative examples how to deal with the exceptional divisors in the complete intersection case.

Even in the smooth case there are usually more elements in $H^{1,1}(X, \mathbb{Z})$ than the ones mentioned above. For the counting and the calculation of intersections of the latter see [1]. If they correspond under the mirror map to complex structure deformations, which can be represented as deformations of the polynomials by vector monomials, they can in principle be incorporated. This is e.g. the situation for the Tian-Yau manifold. If a

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Landau-Ginzburg prescription of the model is available, also perturbations which are not of this kind, can be represented as roots of monomial deformations, as was suggested in [17].

Our approach covers all cases treated in [11][12][13][14][18] [19] [20]. Due to threefold isomorphisms, it also covers the two two-moduli Calabi-Yau spaces which were treated in [15] and among others in [16], namely the degree 8 hypersurface in $\mathbb{P}^4[2,2,2,1,1]$ and the degree 12 hypersurface in $\mathbb{P}^4[6,2,2,1,1]$. We will see moreover that the described subsystems provide simple examples of higher dimensional Calabi-Yau moduli spaces, the simplest case being the subsystem of the Tian-Yau manifold. Apart from being phenomenologically interesting, there exists also a detailed mathematical study on the rational curves of this manifold [21].

We organize the material as follows. In section two we shortly review how to calculate classical intersections numbers. In the third section we explain our method of deriving the system of Picard-Fuchs equations. In section four we exhibit the general structure of the solutions of the Picard-Fuchs equations. This, when combined with the mirror hypothesis, leads to the main result of this section, namely a concise formula for the prepotential. This solves the problem of determining the moduli dependence of the instanton corrected Yukawa couplings and the Weil-Peterson metric for the states associated to the Kählerforms of the ambient space for the class of $(2,2)$ compactifications on CICY's in the large radius limit completely[1].

We demonstrate our method with some selected examples in section five. In section six we discuss the connection of certain complete intersection CY manifolds with rational superconformal field theories. This way we can also explain the occurrence of identical invariants for the rational and elliptic curves on some pairs of hypersurfaces and complete intersections by an identity of conformal field theories. Following the approach of [22] we extend in section seven the analysis to the topological one-loop partition function, which is a moduli dependent quantity describing the difference between the threshold corrections to the E_8 and E_6 gauge couplings [23] (see also [24]). The expansion of the one-loop partition function has a conjectural interpretation in terms of Gromov-Witten invariants for elliptic curves. As the geometrical understanding of these predictions is in a somewhat preliminary state, we found it useful to use our data to calculate them explicitly for various one, two and three moduli examples of different type. In the final section we discuss some open problems and possible avenues for future work.

2. Calculation of the classical topological data of CICY's

We consider in the following complete intersections of l hypersurfaces in products of k projective spaces. Since most formulas allow for an incorporation of weights we will state

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[1] Instead of presenting a lengthy list of examples, we distribute the Mathematica program INSTANTON which calculates the Yukawa couplings and counts the numbers of rational curves for any complete intersection Calabi-Yau manifold discussed in [3],[4] and other examples in singular ambient spaces. It is appended to the HEP-TH version of this paper.

them for the general case. Denote by $d_j^{(i)}$ the degree of the coordinates of $\mathbb{P}^{n_i}[\vec{w}^{(i)}]$ in the j-th polynomial p_j ($i = 1, \ldots, k$; $j = 1, \ldots, l$). We will use the notation

$$
\begin{pmatrix}
\mathbb{P}^{n_1}[w_1^{(1)}, \ldots, w_{n_1+1}^{(1)}] & d_1^{(1)} & , \ldots, & d_l^{(1)} \\
\vdots & \vdots & & \vdots \\
\mathbb{P}^{n_k}[w_1^{(k)}, \ldots, w_{n_k+1}^{(k)}] & d_1^{(k)} & , \ldots, & d_l^{(k)}
\end{pmatrix}_{\chi}^{h^{1,1}}
\tag{2.1}
$$

for a configuration. Let us associate to the Kähler forms induced from the i'th projective space the formal variable J_i and consider the map $\Pi : \mathbb{Z}[J_1, \ldots, J_k] \to \mathbb{Z}$ defined on the generators J_s as

$$
\Pi(J_s) = \left(\prod_{r=1}^{k} \frac{\partial_{J_r}^{n_r}}{n_r!} \right) \left(\frac{\prod_{i=1}^{k} \prod_{j=1}^{n_i+1} (1 + w_j^{(i)} J_i)}{\prod_{j=1}^{l} (1 + \sum_{i=1}^{k} d_j^{(i)} J_i)} \right) \left(\frac{\prod_{j=1}^{l} \sum_{i=1}^{k} d_j^{(i)} J_i}{\prod_{i=1}^{k} \prod_{j=1}^{n_i+1} w_j^{(i)}} \right) J_s \Bigg|_{J_1 = \ldots = J_k = 0} .
\tag{2.2}
$$

It follows from the adjunction formula that the term in the second bracket on the right-hand side yields, by formal expansion, the total Chern class

$$
c(X) = 1 + c_1 + c_2 + c_3 = 1 + c_1^a J_a + c_2^{ab} J_a J_b + c_3^{abc} J_a J_b J_c
\tag{2.3}
$$

of the Calabi-Yau manifold X. The coefficients are given by

$$
c_1^a = \sum_{i=1}^{n_a+1} w_i^{(a)} - \sum_{i=1}^{l} d_i^{(a)} \equiv 0
$$

$$
c_2^{ab} = \frac{1}{2} \left[-\delta^{ab} \sum_{i=1}^{n_a+1} \left(w_i^{(a)} \right)^2 + \sum_{i=1}^{l} d_i^{(a)} d_i^{(b)} \right]
\tag{2.4}
$$

$$
c_3^{abc} = \frac{1}{3} \left[\delta^{abc} \sum_{i=1}^{n_a+1} \left(w_i^{(a)} \right)^3 - \sum_{i=1}^{l} d_i^{(a)} d_i^{(b)} d_i^{(c)} \right] .
$$

Here we have enforced the vanishing of the first Chern class. The numerator of the third term in (2.2) is the top Chern class of the normal bundle of X and the denominator is a normalization of the volume of the weighted projective space. Applying Π to a monomial of the J_i's is equivalent to the integration of the wedge product of the corresponding $(1,1)$ forms, also denoted by J_i, wedged with the Chern class of dual form degree over X. We have therefore

$$
\chi = \int_X c_3 = \Pi(1), \quad \int_X c_2 \wedge J_m = \Pi(J_m), \quad K_{ijk}^0 = \int_X J_i \wedge J_j \wedge J_k = \Pi(J_i J_j J_k).
\tag{2.5}
$$

Note that these formulas, with the exception of the first line in (2.4), are valid only for the case in which X has no singularities. For the singular case the triple intersection gets modified to $\int_X h_i \wedge h_j \wedge h_k = \Pi(J_i J_j J_k) n_0^{(i)} n_0^{(j)} n_0^{(k)}$, where $n_0^{(i)}$ is the least common multiple

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of the orders of the isotropy groups of all fixed points of the manifold under $\mathbb{C}^*$-actions in the i-th weighted projective space. The modification of the first two integrals is more involved. E.g. for complete intersections in one weighted projective space one has [25]

$$\int_X c_2 \wedge J = \frac{12}{n_0!} \left(\frac{\partial}{\partial J}\right)^{n_0} \left. \frac{\prod_{i=1}^l (1 - J^{d_i})}{\prod_{i=1}^{n+1}(1 - J^{w_i})} \right|_{J=0} - 2 \int_X J \wedge J \wedge J. \tag{2.6}$$

In section four we will provide a more illuminative generalization of the formulas (2.4) to the case of hypersurfaces and complete intersections with desingularized quotient singularities.

For example, in a product of two (weighted) projective spaces we have a smooth configuration

$$\left(\begin{array}{c} \mathbb{P}^4[3,1,1,1,1] \\ \mathbb{P}^1 \end{array} \left\|\begin{array}{cc} 6 & 1 \\ 0 & 2 \end{array}\right)^2_{-252} \right. : \quad \begin{array}{l} K^0 = 4\,J_1^3 + 2\,J_1^2 J_2, \\[6pt] c_2 \cdot J_1 = 42, \; c_2 \cdot J_2 = 24, \end{array} \tag{2.7}$$

where we use the notation of [16] to display the intersection numbers.

3. Derivation of the Picard-Fuchs equations

We will begin this section by deriving some general results on the Picard-Fuchs equations as they apply to the models considered here, using the toric data of the manifolds. This approach gives one period by explicit integration and a holonomic system of linear differential operators, which are satisfied by all periods, but which allows for additional solutions. Among the finitely many solutions the periods can be singled out by the requirement that the monodromy acts irreducibly on them. Technically the problem of specifying them is solved here by factorizing the differential operators.

Let us first show how the system of Picard-Fuchs equations for a restricted set of k complex structure deformation parameters can be obtained from toric data of the (mirror) manifolds according to [20]. This system of Picard-Fuchs equations is equivalent to the first order Gauss-Main differential system which describes the variation of the Hodge structure of $H^3(X, \mathbb{Z})$, restricted to the holomorphic and antiholomorphic $(3,0)$ and $(0,3)$ forms and k $(2,1)$ and $(1,2)$ forms. By the mirror map we will identify the k complex structure deformations with the Kähler deformations in the restricted basis of Kähler forms specified in the previous section.

Here we consider only configurations in which the complete intersection does not intersect with singular loci of the ambient space $\mathbb{P}^{n_1}[\vec{w}^{(1)}] \times \cdots \times \mathbb{P}^{n_k}[\vec{w}^{(k)}]$. Without further restricting the generality we may choose $w_{n_i+1}^{(i)} = 1$ for all i. Each projective space $\mathbb{P}^{n_i}[\vec{w}^{(i)}]$ is a toric variety which can be described by a (reflexive) simplicial polyhedron Δ_i with integral vertices in $\mathbb{R}^{n_i}$ (see refs. [26] [16] for its determination).

Since the ambient space is a direct product of all $\mathbb{P}^{n_i}[\vec{w}^{(i)}]$, it is also a toric variety described by the reflexive polyhedra $\Delta = \Delta_1 \times \cdots \times \Delta_k$ in $\mathbb{R}^{n_1} \times \cdots \times \mathbb{R}^{n_k}$.

In refs. [20] [27] it was conjectured that the mirror manifold of the CICY of type (2.1) is given by the CICY constructed from the combinatorial data of the dual polyhedron

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Δ^* of Δ. Δ is reflexive and the corners of the dual Δ^* are the integral points $\nu_{i,1}^* = (1,0,\cdots,0),\cdots,\nu_{i,n_i}^* = (0,\cdots,0,1)$ and $\nu_{i,n_i+1}^* = (-w_1^{(1)},\cdots,-w_{n_i}^{(i)})$ in $\mathbb{R}^{n_i}$ of $\mathbb{R}^{n_1}\times\cdots\times\mathbb{R}^{n_k}$. These vertices satisfy the relations $\sum_{j=1}^{n_i+1} w_j^{(i)}\nu_{i,j}^* = 0$, $(i = 1,\cdots,k)$. We group the vertices $\nu_{i,j}^*$ into l (= number of defining polynomials) sets (a so-called as nef-partition)

$$\{\nu_{i,j}^*\}_{1\le i\le k, 1\le j\le n_i+1} = E_1 \cup \cdots \cup E_l \tag{3.1}$$

by defining E_m $(1 \le m \le l)$ so that it contains $d_m^{(i)}$ vertices from $\{\nu_{i,j}^*\}_{1\le j\le n_i+1}$ for each $i = 1,\cdots,k$. We extend each vertex $\nu_{i,j}^*$ of E_m to $\bar{\nu}_{i,j}^* = (\vec{e}^{(m)},\nu_{i,j})$ in $\mathbb{R}^l\times\mathbb{R}^{n_1}\times\cdots\times\mathbb{R}^{n_k}$ with $\vec{e}^{(m)}$ being the unit vector in the m'th direction of $\mathbb{R}^l$. After adding the additional vertices $\bar{\nu}_{0,p}^* = (\vec{e}^{(p)},\vec{0})$ one finds, as a consequence of the first relation in (2.4), k independent linear relations between the $\sum_{i=1}^l (n_i + 2)$ vertices $\bar{\nu}_{i,j}^*$ of the form

$$\sum l^{(s)}\bar{\nu}_{i,j}^* = \vec{0}. \tag{3.2}$$

The $l^{(s)}$ $(s = 1,\ldots,k, \ j = 1,\ldots,l)$ are given by

$$l^{(s)} = (-d_1^{(s)},\ldots,-d_l^{(s)};\ldots,0,w_1^{(s)},\ldots,w_{n_s+1}^{(s)},0,\ldots) \equiv (\{l_{0j}^{(s)}\};\{l_i^{(s)}\}) \tag{3.3}$$

The mirror manifold X^* can then be written conjectually as the complete intersection of the vanishing locus of the following Laurent polynomials in the variables $X_{m,n}$, $m = 1,\ldots,k$, $n = 1,\ldots,n_m$ (using the notation of ref.[16])

$$P_r = a_r - \sum_{\nu_{i,j}^* \in E_r} a_{i,j} X^{\nu_{i,j}^*} \qquad (r = 1,\ldots,l) \tag{3.4}$$

where the sum is over the (unextended) vertices in the m'th set E_m. The vanishing loci of (3.4) are considered in a toric variety $\mathbb{P}_{\Delta_{(1)}^* + \cdots + \Delta_{(l)}^*}$ with $\Delta_{(i)}^*$ being a convex hull of $\{0\}$ and the set E_i. In ref.[27] the combinatorial aspects of this construction has been nicely formulated.

Choosing a cycle Γ_0 determined by $|X_{m,n}| = 1 \ \forall m,n$, the corresponding period integral (see [26] and [16])

$$w_0(a) = \int_{\Gamma_0} \frac{a_1\cdots a_l}{P_1\cdots P_l} \prod_{m=1}^k \prod_{n=1}^{n_m} \frac{dX_{m,n}}{X_{m,n}}. \tag{3.5}$$

can be performed explicitly by expanding the integrand in a multiple power series in the $1/a_m$ and using the residue formula. After introducing the variables

$$z_s = \frac{a_{s,1}^{w_1^{(s)}}\cdots a_{s,n+1}^{w_{n_s+1}^{(s)}}}{a_1^{d_1^{(s)}}\cdots a_l^{d_l^{(s)}}} \equiv a^{l^{(s)}} \quad (s = 1,\cdots,k) \tag{3.6}$$

it can be easily verified that the period $w_0(a)$ is given by[2]

$$w_0(z) = \sum_{n_s \geq 0} c(n) z^n \tag{3.7}$$

with

$$c(n) = \frac{\prod_j \left(-\sum_{s=1}^{k} l_{0j}^{(s)} n_s\right)!}{\prod_i \left(\sum_s l_i^{(s)} n_s\right)!} = \frac{\prod_{j=1}^{l} \left(\sum_{i=1}^{k} n_i d_j^{(i)}\right)!}{\prod_{i=1}^{k} \prod_{j=1}^{n_i+1} (w_j^{(i)} n_i)!} \tag{3.8}$$

It satisfies the generalized hypergeometric system of Gelfand-Kapranov-Zelevinsky with the k linear differential operators

$$\mathcal{L}_s = \prod_{j=1}^{n_s+1} \left(w_j^{(s)} \theta_s\right) \left(w_j^{(s)} \theta_s - 1\right) \cdots \left(w_j^{(s)} \theta_s - w_j^{(s)} + 1\right)$$
$$- \prod_{j=1}^{l} \left(\sum_{i=1}^{k} d_j^{(i)} \theta_i\right) \cdots \left(\sum_{i=1}^{k} d_j^{(i)} \theta_i - d_j^{(s)} + 1\right) z_s \tag{3.9}$$

associated to the vector of the coefficients of the linear relations $l^{(s)}$, $s = 1, \ldots, k$ given in (3.3). Here the θ_i are logarithmic derivatives $\theta_i = z_i \frac{\partial}{\partial z_i}$. Similarly as in [16] one can show that these equations are satisfied for all periods w_j as they reflect the infinitesimal symmetries of (3.5), independent of the cycle chosen. This system is holonomic, which means that the left ideal I generated by (3.9) in the ring of linear partial differential operators D is of finite rank $\mathrm{rk}(I)$. This implies the existence of $\mathrm{rk}(I)$ independent solutions [28] where $\mathrm{rk}(I)$ is always larger than the expected number of periods.

We are interested in a subset of solutions of the system (3.9) which corresponds to period integrals over the $2k + 2$ cycles dual to the restricted basis of $H^3(X, \mathbb{Z})$. These solutions can be characterized by the requirement that the monodromy acts irreducibly on them [29]. Solving the Riemann-Hilbert problem for the reduced monodromy leads to a reduced system of (lower degree) differential operators $L_1, \ldots, L_h$. h denotes the number of Picard-Fuchs equations; cf. below. In many examples such a system can be specified by factorizing differential operators from (3.9) in the form $p_i(\theta) L_i := \sum_{j=1}^{k} q_j(\theta) \mathcal{L}_j$, where $p(\theta)$ and $q_j(\theta)$ are polynomials in θ. We remark that this is however not the generic situation (see [16]). The system $L_1, \ldots, L_h$ is again holonomic and generates a left ideal whose rank is $2(k + 1)$. In fact this is our criterium to check that a given system of PF equations is complete (cf. below).

Let us demonstrate the derivation of the system $L_1, \ldots, L_h$ for the following complete intersection Calabi-Yau manifold

$$\left(\begin{matrix} \mathbb{P}^3 \\ \mathbb{P}^2 \end{matrix} \middle\| \begin{matrix} 3 & 1 \\ 0 & 3 \end{matrix} \right)_{-54}^{8} . \tag{3.10}$$

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[2] Here and in the following we denote by θ, z, n and ρ the k-tuples $\theta_1, \ldots, \theta_k, z_1 \ldots z_k, n_1, \ldots, n_k$ and $\rho_1, \ldots, \rho_k$. We use obvious abbreviations such as $z^n := \prod_{s=1}^{k} z_s^{n_s}$, $n! := \prod_{s=1}^{k} n_s!$ etc.

The toric description of the mirror manifold was already given in [20]. One has two reflexive simplicial polyhedra Δ_1^*, Δ_2^* with vertices in $\mathbb{R}^3 \times \mathbb{R}^2$:

$$\nu_{1,1}^* = (1,0,0;0,0),\ \nu_{1,2}^* = (0,1,0;0,0),\ \nu_{1,3}^* = (0,0,1;0,0),\ \nu_{1,4}^* = (-1,-1,-1;0,0);$$

$$\nu_{2,1}^* = (0,0,0;1,0),\ \nu_{2,2}^* = (0,0,0;0,1),\ \nu_{2,3}^* = (0,0,0;-1,-1).$$

$$(3.11)$$

We now group the vertices into two sets $E_1 = \{\nu_{1,1}^*, \nu_{1,2}^*, \nu_{1,3}^*\}$ and $E_2 = \{\nu_{2,1}^*, \nu_{2,2}^*, \nu_{2,3}^*, \nu_{1,4}^*\}$ and define the extended vertices $\bar{\nu}^* = (\vec{e}_{1,2}; \nu^*)$ in $\mathbb{R}^7$ where we choose $\vec{e}_1^{(1)} = (1,0)$ and $\vec{e}^{(2)} = (0,1)$ for the vertices in the first and second sets, respectively. The results derived in the following are independent of how we group the vertices into two sets as long as the first set contains three vertices $\nu_{1,i}^*$ and the second set contains the remaining vertices. After adding the two vertices $\nu_{0,k}^* = (\vec{e}^{(k)}, \vec{0})$, $k = 1,2$ this leads by (3.4) to the following two Laurent polynomials:

$$P_1 = a_1 - a_{1,1}X_1 - a_{1,2}X_2 - a_{1,3}X_3$$
$$P_2 = a_2 - a_{2,1}Y_1 - a_{2,2}Y_2 - \frac{a_{2,3}}{Y_1 Y_2} - \frac{a_{1,4}}{X_1 X_2 X_3}$$

$$(3.12)$$

We now have independent linear relations[3] $\sum l^{(k)} \bar{\nu}_{i,j}^* = 0$, $k = 1,2$ between the vertices, namely

$$l^{(1)} = (-3, -1; 1, 1, 1, 1, 0, 0, 0) \quad \text{and} \quad l^{(2)} = (0, -3; 0, 0, 0, 0, 1, 1, 1) \qquad (3.13)$$

They are of the form (3.3) and define the variables $z_1 = \frac{a_{1,1} a_{1,2} a_{1,3} a_{1,4}}{a_1^3 a_2}$, $z_2 = \frac{a_{2,1} a_{2,2} a_{2,3}}{a_2^3}$ via (3.6). The parameters $a_{i,j}$ correspond to trivial automorphisms of (3.12) and can be set to one. Using (3.9) we can associate the following GKZ system to the $l^{(i)}$'s

$$\mathcal{L}_1 = \theta_1^4 - 3\,\theta_1\,(\theta_1 + 3\theta_2)\,(3\theta_1 - 1)\,(3\theta_1 - 2)\,z_1$$
$$\mathcal{L}_2 = \theta_2^3 - (\theta_1 + 3\theta_2)\,(\theta_1 + 3\theta_2 - 1)\,(\theta_1 + 3\theta_2 - 2)\,z_2 \,.$$

$$(3.14)$$

In addition to the power series solution the system has eleven logarithmic ones. To obtain an irreducible subsystem we factorize in the following way: $\mathcal{L}_1 =: \theta_1 L_1$ and $\mathcal{L}_1 + 27\mathcal{L}_2 =: \theta_1(\theta_1 + \theta_2)L_2$, which leads to the reduced system

$$L_1 = \theta_1^3 - 3\,(\theta_1 + 3\theta_2)\,(3\theta_1 - 1)\,(3\theta_1 - 2)\,z_1$$
$$L_2 = (\theta_1^2 - 3\,\theta_1\theta_2 + 9\,\theta_2^2) - 3\,(3\theta_1 - 1)(3\theta_1 - 2)\,z_1 - 27\,(\theta_1 + 3\theta_2 - 1)\,(\theta_1 + 3\theta_2 - 2)\,z_2.$$

$$(3.15)$$

The fact that the number of equations equals the number of moduli is special to the example here. In fact we will see in example (5 *vii*) that the numbers of linear differential operators describing the Picard-Fuchs system locally can vary in the different patches of the moduli space.

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[3] The components of the $l^{(k)}$ refer to vertices $\bar{\nu}_{i,j}^*$ with i, j in lexicographic order.

Let us finally comment on some generic features of the moduli space of the mirror manifold and its compactification, as far as we will need them for fixing the holomorphic ambiguity in section seven. It is always easy to find the invariance group of the Laurent polynomials, which acts by phase multiplication on the parameter a_i and the coordinates $X_i, Y_i, \ldots$. In our example one finds a $\mathbb{Z}_9$ invariance group of (3.12) acting by $(X_i, Y_j, a_1, a_2) \mapsto (\alpha^k X_i, \alpha^{6k} Y_i, \alpha^k a_1, \alpha^{6k} a_2)$, with $\alpha = \exp\frac{2\pi i}{9}$ and $k \in \mathbb{Z}$. Because of this invariance of the parameter space of (3.12), we have to define the moduli space of the mirror manifold as the quotient $\mathcal{M} = \tilde{\mathcal{M}}/\mathbb{Z}_9$, where $\tilde{\mathcal{M}}$ is parametrized by a_1, a_2. It is easy to see that the parameters (3.6) are in general invariant under the group action on the Laurent polynomials and hence well defined on the quotient[4]. The fact that (3.6) are invariant under the group action and hence well defined on the quotient can be seen in general[5].

The singularity of the moduli space of the Laurent polynomials at $a_i = 0$ due to the phase symmetries can now always be described by toric geometry. In general one has to consider as second step also the quotient with respect to invariances of the Laurent polynomials which are not acting simply by phase multiplication. In our example such symmetries are not present and we see that $\tilde{n}_1 = a_1^9$, $\tilde{n}_2 = a_1^3 a_2$ and $\tilde{n}_3 = a_2^3$ generate the multiplicative semigroup of invariant monomials under the $\mathbb{Z}_9$ but satisfy the relation $\tilde{n}_1 \tilde{n}_3 = \tilde{n}_2^3$, describing a rational A_2 double point. It is now straightforward to give a toric description of the moduli space using the secondary fan construction of ref. [30]. The secondary fan, whose dimension equals the number of Kähler moduli, contains the Kähler cone. Since it is a complete fan, it gives a toric description of the compactification of Kähler moduli space; for a review, see also [31]. If we define a matrix B whose rows are the generators $l^{(s)}$ of the Mori cone, then the columns of B generate the one-dimensional cones of the secondary fan. Their minimal generators for the model considered here are the vectors e_1, e_2, $-(e_1 + 3e_2)$ and $-e_1$, where e_1, e_2 generate a square lattice. The cone [6] $\langle e_1, e_2 \rangle$ is the Kähler cone and the cone $\langle -e_1, -(e_1 + 3e_2) \rangle$ describes the A_2 double point. This description of the A_2 double point is related to that given e.g. in [32] by a change of basis $e_1 \to -e_2$, $e_2 \to -e_1 + e_2$.

The general theory, due to Hironaka, tells us that the compactified moduli space can be resolved in such a way that all singularities of the Picard-Fuchs equation are regular divisors with normal crossings and part of the resolutions necessary to achieve this goal can be described by toric resolution if we start as above. Finding auch a resolution is necessary in order to analyse the full modular group. A thorough analysis of the modular groups for Calabi-Yau compactifications were given for one moduli examples in [11][13][14] and for two types of two moduli examples in [15].

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[4] This generalizes the reasoning, which leads for the quintic to the consideration of $z = 1/a^5$ [11]as the natural variable.

[5] This generalizes the reasoning, which leads for the quintic to the consideration of $z = 1/a^5$ [11] as the good variable.

[6] Here the conventions are as in [32].

4. Local behaviour of the solutions, mirror map and instanton corrected Yukawa couplings

In this section we calculate the singular locus of the Picard-Fuchs equations and discuss some essential parts of the local behaviour of their solutions. We explain how to introduce canonical coordinates and fix a canonical form of the period vector. This leads to simple expressions of the instanton corrected Yukawa couplings and the Kähler potential in terms of the solutions. We will see that all topological data of the Calabi-Yau manifold appear naturally in this period vector.

It was shown in [33] that general Picard-Fuchs systems have only regular singular points, i.e. locally the solutions are given by power series or series involving finite powers of logarithms in z_i.

Let us now describe the singular locus (cf. [28]). Denote the linear partial differential operators of degree m, defined in a neighborhood U of $z \in \mathcal{M}$ (the subspace of the complex structure moduli space to which our analysis applies), by $L_i = \sum_{|p| \le m} a_i^p(z) \left(\frac{d}{dz}\right)^p$. They define a left ideal I in the ring of partial differential operators on U. We now introduce the symbol of L_i: $\sigma(L_i) = \sum_{|p|=m} a_i^p(z)\xi_1^{p_1} \cdots \xi_k^{p_k}$, where $\xi_1, \ldots, \xi_k$ is a coordinate system in the fiber of the cotangent bundle T^*U at z. The ideal of symbols is defined by $\sigma(I) = \{\sigma(L)|L \in I\}$. The singular locus is $S(I) = \pi(\mathrm{Ch}(I) - U \times \{0\})$, where the characteristic variety $\mathrm{Ch}(I)$ is the subvariety in T^*U specified by the ideal of symbols and π is the projection along the fiber of T^*U. The fact that $\sigma(I)$ is generated by $\sigma(L_i)$ is a special property of Picard-Fuchs systems. This follows e.g. from the way the Picard-Fuchs equations are derived by the Griffith-Dwork-Katz reduction method and simplifies the calculation of $S(I)$. Let us demonstrate this for L_1, L_2 given in eq. (3.15). The symbols are

$$
\begin{aligned}
\sigma(L_1) &= z_1^3(\xi_1(1 - 27z_1) - 81z_2\xi_2)\xi_1^2 \\
\sigma(L_2) &= z_1{}^2(1 - 27z_1 - 27z_2)\xi_1^2 - 3z_1z_2(1 + 54z_2)\xi_1\xi_2 + 9z_2{}^2(1 - 27z_2)\xi_2^2.
\end{aligned}
\tag{4.1}
$$

Case by case analysis reveals that $\mathrm{Ch}(I)$ decomposes into the following components

$$
\begin{aligned}
\mathrm{Ch}(I) = &\{81z_2\xi_2 - (1 - 27z_1)\xi_1 = (1 - 27z_1)^3 - 27z_2 = 0\} \cup \{z_1 = z_2 = 0\} \\
&\cup \{\xi_1 = (1 - 27z_2) = 0\} \cup \{\xi_1 = z_2 = 0\} \cup \{\xi_2 = z_1 = 0\}.
\end{aligned}
\tag{4.2}
$$

Denoting the projections of the components on U by Δ_i we have $S(I) = \prod_{i=0}^3 \Delta_i$, with $\Delta_0 = (1 - 27z_1)^3 - 27z_2$, $\Delta_1 = (1 - 27z_2)$, $\Delta_2 = z_1$, $\Delta_3 = z_2$.

The singularities $\Delta_i = 0$ detected so far correspond to the discriminant locus in the moduli space on which the defining polynomials cease to be transverse, i.e. where $p_1 = \ldots = p_l = 0$ and $dp_1 \wedge \cdots \wedge dp_l = 0$ or, equivalently, where the Laurent polynomials fail to be Δ^* regular [26]. We calculate for comparison the discriminant of (3.12) at the end of section (5i). To study further singularities we have to compactify the moduli space. We can describe the compactification by the toric variety corresponding to the secondary fan constructed in the previous section. Then we find with the same method as above a singular locus of the Picard-Fuchs equation at the origin of the patch defined by σ_3, which

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is due to the $\mathbb{Z}_9$ identification of the parameter space discussed above. Such an analysis can be made for the general case and the singular locus determined this way is used in section seven in order to fix the holomorphic ambiguity for various examples.

From the definition of the (unnormalized) Yukawa-couplings (coefficients of the cubic form Ξ in [34]) and the variational property of the $(3,0)$-form [34] one finds (see also [35]) that the vanishing of the symbols at degree three gives relations for the Yukawa-couplings by the simple replacement $\xi_i \xi_j \xi_k \to K_{z_i z_j z_k}$; e.g. from $\sigma(L_1) = 0$ one has $K_{z_1 z_1 z_2} = \frac{(1-27z_1)}{81 z_2} K_{z_1 z_1 z_1}$. This determines the Yukawa couplings up to a gauge dependent overall function. The gauge in which the Picard-Fuchs equation are derived is defined by (3.5) , and the corresponding gauge dependent function can be calculated, up to a constant, by the methods outlined in [11][16].

These unnormalized Yukawa couplings have singularities on the discriminant of the Calabi-Yau manifold. From the derivation of the Yukawa couplings in [34][16] it follows that there is always one component which will appear in the denominator of all the Yukawa-couplings, which we call the general component Δ_0, following [36].

In the following we will demonstrate that the calculation of those Yukawa-couplings on X^* which are functions of the complex structure moduli is not necessary for the purpose of determining the number of rational and elliptic curves. I. e. one can directly compute the instanton corrected couplings on X.

Because of the fact that the symbols of I generate the ideal of symbols it is a rather simple algebraic problem to write down the associated first order Pfaffian system, equivalent to the Gauss-Manin connection. However we found that some local properties of the solutions are most easily obtained directly from $L_1, \ldots, L_h$. Most important is the analysis of the solutions around $z = 0$, which is obviously a singular point in every system derived by factorization from (3.9). For the analysis it is useful to introduce the ring $\mathcal{R}$, which we define as the polynomial ring $\mathbb{C}[\theta]$ modulo the ideal $\mathcal{I}$ generated by the principal parts I_s of the operators L_s

$$I_s(\theta) = \lim_{z \to 0} L_s(\theta, z), \tag{4.3}$$

i.e. $\mathcal{R} = \mathbb{C}[\theta]/\mathcal{I}$. First we note that the $I_s(\theta)$ are homogeneous polynomials and the solution to $I_s = 0$, $\forall s$ is $(2k + 2)$-times degenerate at $\theta = 0$. $\mathbb{C}[\theta]$ has a natural vector space structure with the monomials as orthogonal basis. We can choose representatives of $\mathcal{R}$ as homogeneous polynomials orthogonal to $\mathcal{I}$. The subspace R of $\mathbb{C}[\theta]$ spanned by them has in all cases dimensions $\{1, k, k, 1\}$ at degrees $\{0, 1, 2, 3\}$. As we will now show ,it can be identified with the vector space of solutions to $L_1, \ldots, L_h$. The grading translates to the fact that we have one pure power series solution, k solutions with a part linear in logarithms of z, k solutions with a part quadratic in the logarithms of z and one solution which has a part cubic in the logarithms. This identifies $z = 0$ as a point where all but one cycle in $H_3(X)$ are degenerate[37], which is also referred to as point of maximally unipotent monodromy[38] and provides precisely the structure of solutions one needs for the mirror map.

Extending the definition of $x! = \Gamma(x + 1)$ to $x \in \mathbb{R}$, we define the coefficient $c(n + \rho)$ for arbitrary values of the k parameters ρ_i and define the ρ-dependence of (3.7) as

$$w_0(z, \rho) = \sum_{n_i \geq 0} c(n + \rho) z^{n+\rho}. \tag{4.4}$$

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By the method of Frobenius, the logarithmic solutions are obtained by taking linear combinations of derivatives $D_\rho = \sum (b_n/n!)\partial_\rho^n$ of $w_0(z, \rho)$ evaluated at $\rho = 0$. Here we define $\partial_{\rho_i} := \frac{1}{2\pi i}\frac{\partial}{\partial_{\rho_i}}$. The factor $2\pi i$ will make the monodromy matrix around $z = 0$ integer. As $[L_s, \partial_{\rho_i}] = 0$, it is then sufficient to check whether

$$D_\rho\left(L_s w_0(z, \rho)\right)\big|_{\rho=0} = 0, \ \forall s \tag{4.5}$$

to establish $D_\rho w_0(z, \rho)|_{\rho=0}$ as a solution. By consideration of the explicit form of the series (4.4) one can show that the conditions for vanishing of the constant terms in (4.5)

$$D_\rho(I_s(\rho)c(0, \rho)z^\rho)\big|_{\rho=0} = 0, \ \forall s, \tag{4.6}$$

are in fact also sufficient.

We introduce an identification of the ring $\mathbb{C}[\theta]$ with the ring of the partial derivatives w.r.t. ρ

$$\varphi : \sum_{n_i \geq 0} b_n \theta^n \mapsto \sum_{n_i \geq 0} \frac{b_n}{n!}\partial_\rho^n, \tag{4.7}$$

which induces an isomorphism between the vector space R and the vector space of solutions to $L_1, \ldots, L_k$, as can be seen as follows. An element $r \in R$ is orthogonal to all polynomials of the form $m(\theta)I_s$, where $m(\theta)$ is monomial in θ. As r and I_s are homogeneous one has to check orthogonality only for monomials $m(\theta)$ of degree $\deg(r) - \deg(I_s)$. Suppose r is of lower degree, then I_s (4.6) will obviously hold for $D_\rho = \varphi(r)$, because all terms in $D_\rho(I_s(\rho)c(0, \rho)z^\rho)$ have positive degree in ρ and will vanish, after setting ρ to zero. If $\deg(r) = \deg(I_s) + n$, then $\deg(r) - n$ derivatives of $D_\rho = \varphi(r)$ have to act on I_s to give a nonvanishing term. By our choice of the factorials the vanishing of these contributions is equivalent to $r \perp m(\theta)I_s = 0$ for any $m(\theta)$ with $\deg m(\theta) = n$. Hence $r \in R$ iff $\varphi(r)w_0(z, \rho)|_{\rho=0}$ is a solution.

Besides the unique power-series solution (3.7) we choose for the k solutions linear in the logarithms the basis $\partial_{\rho_i} w(z, \rho)|_{\rho=0}$. They are given by

$$w_i(z) = \sum_{n_i \geq 0} d_i(n)z^n + w_0(z)\frac{\log z_i}{2\pi i} \tag{4.8}$$

with

$$d_i(n) = \frac{1}{2\pi i}\frac{\partial}{\partial_{\rho_i}}c(n + \rho)\bigg|_{\rho=0}$$

The top element of R is unique, up to a constant, which we will fix in a moment, and so is the solution cubic in the logarithms.

Defining now

$$D_i^{(1)} := \partial_{\rho_i}, \quad D_i^{(2)} := \frac{1}{2}\tilde{K}_{ijk}\partial_{\rho_j}\partial_{\rho_k} \ \text{and} \ D^{(3)} := -\frac{1}{6}\tilde{K}_{ijk}\partial_{\rho_i}\partial_{\rho_j}\partial_{\rho_k}, \tag{4.9}$$

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where summation is implicit, we have a natural basis for the period vector

$$\Pi(z) = \begin{pmatrix} w_0(z) \\ D_i^{(1)} w_0(z, \rho)|_{\rho=0} \\ D_i^{(2)} w_0(z, \rho)|_{\rho=0} \\ D^{(3)} w_0(z, \rho)|_{\rho=0} \end{pmatrix}. \tag{4.10}$$

When one actually performs the derivatives w.r.t. ρ_i, one has to be careful when treating the poles of the Gamma-function and its derivatives. We clear up these technical details in Appendix A.

So far we have only dealt with the complex structure deformation parameters z_i. To describe the mirror map between the subsectors of the theory depending only on the complex structure parameters and the complexified Kähler structure parameters[7], respectively, we will briefly review their common structure. They can be identified with topological field theories defined by the BRST operators $Q^C = G^+ + \bar{G}^-$ and $Q^K = G^+ + \bar{G}^+$, respectively. The $G^\pm, \bar{G}^\pm$ in Q^C and Q^K are related to the zero modes of the superpartners of the energy momentum tensor of the underlying $N = 2$ superconformal field theory by two different kinds of twist procedures which make either Q^C or Q^K to scalar operators [40][23]. At string tree level all relevant information is encoded in two prepotentials (free energy) F^C, F^K, which are sections of holomorphic line bundles L^2 over the complex and Kähler structure moduli spaces respectively. There exists a coordinate choice, so-called inhomogeneous special coordinates t^i, such that the Yukawa couplings (structure constants) are ordinary third derivatives ($\partial_i = \frac{\partial}{\partial t^i}$)

$$K_{ijk} = \partial_i \partial_j \partial_k F \tag{4.11}$$

of the prepotential. For general coordinate choices the derivatives have to be covariantized w.r.t. the connection of the line bundle L and the metric on the moduli space. Unlike for e.g. topological models on manifolds with $c_1 > 0$ [41] the prepotential F for Calabi-Yau threefolds cannot be derived from the associativity of the structure constants alone. For these cases one has on the other hand an additional geometrical structure known as special geometry, which was discovered first in the context of $N = 2$ supergravity theories [42] and derived for the Kähler and complex structure moduli spaces of Calabi-Yau and/or $N = 2$ string compactification in [43] [44] [45][39]. It implies that also the Kähler potential for the Weil-Peterson metric derives from the prepotential (hence the name) as

$$K = -\log \left((t^i - \bar{t}^i)(\partial_i F + \bar{\partial}_i \bar{F}) + 2\bar{F} - 2F \right). \tag{4.12}$$

On the complex structure side the Yukawa couplings are expressed for a choice of the holomorphic threeform Ω as [34] $K_{ijk} = \int \Omega \partial_i \partial_j \partial_k \Omega$ and the Kähler metric is given by [8]

---

[7] Recall that the moduli describing the variation of the Riccic flat Kähler metric and the antisymmetric tensor background field can be parametrized by $\delta g_{i\bar{j}} dz^i \wedge d\bar{z}^{\bar{j}} = \sum_{i=1}^k \delta \hat{t}^i J_i$ and $\delta B_{i\bar{j}} dz^i \wedge d\bar{z}^{\bar{j}} = \sum_{i=1}^k \delta \tilde{t}^i J_i$ where one can combine $t^i = \tilde{t}^i + i\hat{t}^i$ to the so-called complexified Kähler structure parameter [39].

[8] $\Sigma = \begin{pmatrix} 0 & -\sigma \\ \sigma & 0 \end{pmatrix}, \sigma = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$K = -\log(\Pi^\dagger \Sigma \Pi)$ where Π is the period vector for a symplectic basis of $H_3(X)$. As the holomorphic three-form Ω is defined only up to gauge transformations $\Omega \to f(z)\Omega$ ($f(z)$ holomorphic) it takes values in a holomorphic line bundle L over the moduli space. The transformation properties of the quantities introduced are: $\Omega \in L$, $F \in L^2$, the Yukawas K_{ijk} transform as components of elements in $L^2 \otimes \mathrm{Sym}((T^*_{\mathcal{M}})^{\otimes 3})$ and $e^{-K} \in L \otimes \bar{L}$. $e^{2K} g^{i\bar\imath} g^{j\bar\jmath} g^{k\bar k} K_{ijk} \bar{K}_{\bar\imath\bar\jmath\bar k}$ is then invariant under Kähler and coordinate transformations. The physical Yukawa couplings appearing in the effective Lagrangian are those for canonically normalized matter fields.

To relate the solution of the Picard-Fuchs equations directly to F we may use the form of the Picard-Fuchs differential operators in special coordinates in the Kähler gauge [46],

$$\sum_{i=1}^{k} \partial_j \partial_p (K_r^{-1})^{li} \partial_r \partial_i. \tag{4.13}$$

The period vector is then trivially expressed in terms of F: $\Pi(t) = (1, t^i, \partial_i F, 2F - t^i \partial_i F)$ (here and below we again use summation convention). The relation of the coordinates in Kähler gauge t^i to the homogeneous special coordinates (X^0, X^i), in which the period vector reads $(X^0, X^i, (\partial \mathcal{F}/\partial X^i), (\partial \mathcal{F}/\partial X^0))$, is given by $t^i = X^i/X^0$ with $\mathcal{F} = (X^0)^2 F$.

Guided by the mirror hypothesis we should have the same structure for the Kähler side and therefore, for a formal large radius expansion of F^K

$$F^K = \frac{1}{6} K^0_{ijk} t^i t^j t^k + \frac{1}{2} a_{ij} t^i t^j + b_i t^i + \frac{1}{2} c + F_{\text{inst.}} \tag{4.14}$$

around $\mathrm{Im}(t^i) \to \infty$, the "period vector"

$$X^0 \begin{pmatrix} 1 \\ t^i \\ \frac{1}{2} K^0_{ijk} t^j t^k + a_{ij} t^j + b_i + \partial_{t_i} F_{\text{inst.}} \\ -\frac{1}{6} K^0_{ijk} t^i t^j t^k + b_i t^i + c + (2F_{\text{inst.}} - t^i \partial_{t_i} F_{\text{inst.}}) \end{pmatrix} = X^0 \Pi(t) \tag{4.15}$$

Here K^0_{ijk} are the classical intersections calculated in section two and $F_{\text{inst.}}$ is a power series in $q_i = e^{2\pi i t^i}$, which encodes the instanton corrections.

Starting from (4.10) in z_i coordinates we have now to make a coordinate transformation to suitable inhomogenous coordinates, which are defined by ratios of solutions of the Picard-Fuchs equations as

$$t^i(z) = \frac{w_i(z)}{w_0(z)}. \tag{4.16}$$

The reason for picking this quotient is that the invariance of the theory under the monodromies around $z_i = 0$ can be identified with the invariance of the topological sigma model under real shifts $t^i \to t^i + 1$ of the antisymmetric background field or equivalently of the effective field theory under discrete Peccei-Quinn symmetries. As will become clear below, this choice identifies the form of the prepotential for the complex structure moduli space on X^* with the one expected for the Kähler structure moduli space on X. It strongly resembles the elliptic curve case where the ratios of two arbitrary solutions of the period equation

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and its inverse satisfy a third order non-linear differential equations introduced by Schwarz. Inserting the solutions $z(t)$ in the J invariant of the elliptic curve gives the wellknown expansion $J(\frac{at+b}{ct+d}) = \frac{1}{1728}(q^{-2} + 144 + \ldots)$ of the J function in terms of $q = e^{2\pi i(\frac{at+b}{ct+d})}$. The choice of the logarithmic solution for w_1 and the power series solution for w_0 at the point of maximal unipotent monodromy corresponds then to $a = d = 1, b = c = 0$. For Calabi-Yau spaces one likewise finds that the mirror maps $z_i(q_1, \ldots, q_h)$ always have an integral expansion, which is a necessary condition for obtaining integer instanton numbers. One therefore expects that the mirror map plays an important rôle for the theory of modular forms on the moduli space of Calabi-Yau space, which however seems to be much more intrigate as the modular group acts in general non-arithmetically on the geometrical parameters.

Inserting the inverse map of (4.16) into (4.10) we obtain the transformed period vector

$$
\tilde{\Pi}(t) = w_0(t) \begin{pmatrix} 1 \\ t^i \\ \frac{1}{2}\tilde{K}^0_{ijk}t^j t^k + \hat{b}_i + \partial_{t_i}\tilde{F}_{\text{inst.}} \\ -\frac{1}{6}\tilde{K}^0_{ijk}t^i t^j t^k - \hat{b}_i t^i + \hat{c} + (2\tilde{F}_{\text{inst.}} - t^i\partial_{t_i}\tilde{F}_{\text{inst.}}) \end{pmatrix}. \tag{4.17}
$$

By comparison with (4.15) we can view all quantities marked with a tilde, up to one overall normalisation, as predictions of mirror symmetry. Especially we have to identify K^0_{ijk} with $\tilde{K}^0_{ijk}$. In fact, for all complete intersection cases the top element of the ring R encodes the intersection numbers in the integral basis of divisors coming from the ambient spaces. This turns out to be true also for the hypersurfaces discussed in [16], for the basis of divisors which generate $H_2(X, \mathbb{Z})$.

After fixing the normalisation, we see that the instanton corrected Yukawa coupling K_{ijk} can be uniquely expressed by the solutions of the Picard-Fuchs equations as

$$
K_{ijk}(t) = \partial_{t_i}\partial_{t_j}\frac{D^{(2)}_k w_0|_{\rho=0}}{w_0}(t) \tag{4.18}
$$

or, equivalently, by the derivatives of the prepotential

$$
F(t) = \frac{1}{2}\left(\frac{1}{w_0}\right)^2 \left\{ w_0 D^{(3)} w_0 + D^{(1)}_l w_0 D^{(2)}_l w_0 \right\}(t) \,|_{\rho=0}. \tag{4.19}
$$

We will use formula (4.18),(4.19), which does not require the evaluation of the Yukawa couplings on the mirror manifold, to compute a convergent expansion for the instanton corrected Yukawas and the prepotential in the large radius limit and to predict the number of rational curves for various manifolds.

So far we have seen that the natural choice for the period vector (4.10) matches the leading terms in t of the components of (4.15) and leads to a prediction of the instanton corrections. Let us now compute also the lower order terms in t i.e. the constants $\hat{b}_i, \hat{c}$.

For this purpose one has to take the derivatives w.r.t. ρ explicitly

$$\frac{\partial}{\partial \rho_i} c(0) = -\left(\sum_\alpha l_{0\,\alpha}^{(i)} + \sum_\alpha l_\alpha^{(i)} \right) \Gamma'(1) =: c_1^i \Gamma'(1) = 0$$

$$\frac{\partial}{\partial \rho_i} \frac{\partial}{\partial \rho_j} c(0) = \frac{\pi^2}{6} \left(\sum_\alpha l_{0\,\alpha}^{(i)} l_{0\,\alpha}^{(j)} - \sum_\alpha l_\alpha^{(i)} l_\alpha^{(j)} \right) =: \frac{\pi^2}{3} c_2^{ij} \qquad (4.20)$$

$$\frac{\partial}{\partial \rho_i} \frac{\partial}{\partial \rho_j} \frac{\partial}{\partial \rho_k} c(0) = 2 \left(\sum_\alpha l_{0\,\alpha}^{(i)} l_{0\,\alpha}^{(j)} l_{0\,\alpha}^{(k)} + \sum_\alpha l_\alpha^{(i)} l_\alpha^{(j)} l_\alpha^{(k)} \right) \zeta(3) =: 6 c_3^{ijk} \zeta(3)$$

where the coefficients c_n coincide for complete intersections in products of nonsingular weighted projective spaces with the ones given in (2.4). Using the Gauss-Bonnet formula we get from this the following remarkable identities

$$\int_X c_2 J_i = \int_X c_2^{jk} J_j J_k J_i = \frac{3}{\pi^2} \int_X J_j J_k J_i \frac{\partial}{\partial \rho_j} \frac{\partial}{\partial \rho_k} c(0) = -24 D_i^{(2)} c(0)$$

$$\chi = \int_X c_3^{ijk} J_i J_j J_k = \frac{1}{6\zeta(3)} \int_X J_i J_j J_k \frac{\partial}{\partial \rho_i} \frac{\partial}{\partial \rho_j} \frac{\partial}{\partial \rho_k} c(0) = i \frac{(2\pi)^3}{\zeta(3)} D^{(3)} c(0),$$

$$(4.21)$$

with $D_i^{(2)}$ and $D^{(3)}$ defined in (4.9), which express the Euler number and $\int_M c_2 J_i$ in terms of the intersection numbers and the generators of the Mori cone $l^{(i)}$. These identities hold also for the canonically resolved manifolds, if the complete intersection had canonical quotient singularities. The linear terms in the third row of (4.17) are thus $\hat{b}_i = -\frac{1}{24} \int_M c_2 J_i$ and in the last row $\hat{c} = -i \frac{\zeta(3)}{(2\pi)^3} \chi$.

The imaginary part of the constant $\hat{c}$ can be identified with a σ-model loop contribution and is proportional to the Euler number of the manifold. The constant of proportionality seems to be universal and has been calculated explicitly for the quintic hypersurface in $\mathbb{P}^4$ and other one moduli cases [14][13] as $-i \frac{\zeta(3)}{(2\pi)^3}$. $\mathrm{Im}(a_{ij}) = \mathrm{Im}(b_i) = 0$ is a necessary condition for having a continuous Peccei-Quinn symmetry $t_j \to t_j + \alpha_j$, α_j real, which is broken by instanton corrections to discrete shifts $t^i \to t^i + 1$ as mentioned before.

While the real constants b_i, a_{ij} in (4.15) are irrelevant for physical quantities, they are important for fixing an integral symplectic basis for the period vector. In fact we find that the constant $\hat{b}_i$ for the one moduli cases [13][14] is also correctly reproduced in the third line. The constants in front of the first subleading terms in (4.17) do not correspond to a choice of an integral basis as can be seen by comparison with (4.13). Especially the constants a_{ij} do not seem to be directly related to topological numbers; in fact there are no further at our disposal. They have to be fixed by analysing the monodromy matrices themselves. The monodromy operation $T_i : t^i \to t^i + 1$ on (4.15) is obvious. The requirement that it is integral and symplectic yields restrictions on a_{ij}, e.g. for the one moduli cases $a = (\frac{K^o}{2} \bmod \mathbb{Z})$ (and $2b = ((1 - \frac{K^o}{6}) \bmod \mathbb{Z}))$. The integer part is irrelevant as it can be absorbed by an $SP(4, \mathbb{Z})$ transformation, hence the basis can be specified. Although we have no general proof, it is tempting to conjecture that occurence of the topological numbers K_{ijk}^0, χ and, up to $SP(2h + 2, \mathbb{Z})$ transformations, also $\int_X c_2 J_i$ in

F at the point of maximal unipotent monodromy in a integral sympletic basis is a general feature.

The instanton part of the Yukawa couplings comes from stationary points of the classical string action, which correspond to holomorphic mappings from $\mathbb{P}^1$ to the CY manifold. In the coordinates t^i defined in (4.16) it enjoys the following expansion

$$
\begin{aligned}
K_{klm} &= \int_X J_k \wedge J_l \wedge J_m + \sum_d \int_{\mathcal{M}_{C_d}} d\mu \; \frac{e^{2\pi i \int_{C_d} K(X)}}{1 - e^{2\pi i \int_{C_d} K(X)}} \\
&= K^0_{klm} + \sum_{d_1,\ldots,d_k} \frac{n^r_{d_1,\ldots,d_k} d_k \, d_l \, d_m}{1 - \prod_{i=1}^k q_i^{d_i}} \prod_{i=1}^k q_i^{d_i}
\end{aligned}
\tag{4.22}
$$

Here we have defined the degree of the curve as $d_i = \int_C J_i$, which is an integer for a basis $J_i \in H^{1,1}(\hat{X}, \mathbb{Z})$. The denominator $(1 - e^{\int_{C_d} K(x)})$ gives the correct combinatorical contributions from multi-covers of the curves s.t. the integral $d\mu$ of the Euler class over the compactified moduli space $\mathcal{M}_{C_d}$ of holomorphic maps of multi-degree d from $\mathbb{P}^1$ [47] can be taken over single cover curves only. The resulting invariants $n^r_{d_1,\ldots,d_k}$ are expected to be integers, for isolated curves they just count their numbers. In [16] examples of negative invariants n^r were found. They admit only the interpretation that there are non-isolated singular curves at the corresponding degree. The occurence of the terms $d_k \, d_l \, d_m$ in the second line is due to the integral of the part of the moduli space describing the reparametrisation of $\mathbb{P}^1$, as was explained for isolated curves in [48].

5. Selected examples

(i) As our first example we will calculate the topological invariants n^r for the manifold (3.10). The ideal $\mathcal{I}$ is generated by $I_1 = \theta_1^2 - 3\,\theta_1\theta_2 + 9\,\theta_2^2$ and $I_2 = \theta_1^3$, so R is spanned by

$$
\begin{aligned}
&1 \\
&\theta_1, \;\; \theta_2 \\
&9\,\theta_1\theta_2 + 3\,\theta_2^2, \;\; 9\,\theta_1^2 + 3\,\theta_1\theta_2 \\
&9\,\theta_1^2\theta_2 + 3\,\theta_1\theta_2^2 .
\end{aligned}
\tag{5.1}
$$

By (4.7) this translates to a basis of solutions for L_1, L_2 as one may verify. The fact that the top element of R coincides, up to scaling, with $K^0 = 9\,J_1^2 J_2 + 3\,J_1 J_2^2$, calculated by (2.5), is a nontrivial check on the mirror hypothesis. Using (4.18) one obtains concise formulas for the instanton corrected intersection numbers

$$
K_{i,j,1}(t) = \partial_{t_i} \partial_{t_j} \frac{(9\,\partial_{\rho_1}\partial_{\rho_2} + \frac{3}{2}\partial_{\rho_2}^2)w_0|_{\rho=0}}{w_0}(t) \quad K_{i,j,2}(t) = \partial_{t_i} \partial_{t_j} \frac{(\frac{9}{2}\partial_{\rho_1}^2 + 3\,\partial_{\rho_1}\partial_{\rho_2})w_0|_{\rho=0}}{w_0}(t)
\tag{5.2}
$$

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from which the topological invariants n^r follow by comparison with (4.22). Below we display them for multi-degrees (d_{J_1}, d_{J_2}), $d_{J_1} + d_{J_2} \leq 6$.

(1,0)	243	(0,1)	63
(2,0)	243	(0,2)	63
(3,0)	54	(0,3)	54
(4,0)	243	(0,4)	63
(5,0)	243	(0,5)	63
(6,0)	54	(0,6)	54

(1,1)	972
(2,2)	156249
(3,3)	60018786
(2,1)	15309
(4,2)	111401163
(3,1)	179901
(4,1)	1558845
(5,1)	11558295

(1,2)	3402
(2,4)	4803867
(3,2)	4830597
(1,3)	9720
(2,3)	977589
(1,4)	25515
(1,5)	61236

In fact, J_i is a basis of a subspace of $H^2(X, \mathbb{Z})$ in which the Kähler cone is simply given by

$$\sigma(K) = \{ \sum_i t_i J_i | t_i > 0 \} \tag{5.3}$$

In [10] a three generation model ($\chi = -6$) was constructed, by dividing the manifold (3.10) by a $G = \mathbb{Z}_3 \times \mathbb{Z}_3$ symmetry and resolving the singular quotient. The group G is generated, on the homogeneous coordinates of $\mathbb{P}^3 \times \mathbb{P}^2$, by $(x_0, x_1, x_2, x_3, y_1, y_2, y_3) \mapsto (x_0, x_1, x_2, x_3, \alpha y_1, y_2, \alpha^{-1} y_3)$ ($\alpha = \exp\frac{2\pi i}{3}$), with three fixed tori $T_i = \{ x \in \mathbb{P}^3 | x_i = 0, p_1 = 0 \}$, $i = 1, 2, 3$, and a freely acting cyclic permutation $(x_0, x_1, x_2, x_3, y_1, y_2, y_3) \mapsto (x_0, x_2, x_3, x_1, y_2, y_3, y_1)$. This action is only compatible with a subspace of the moduli space of complex structure deformations, namely with the following perturbations

$$p_1 = \sum_{i=1}^{3} x_i y_i^3 - \tilde{a}_1 x_0 y_1 y_2 y_3$$

$$p_2 = \sum_{i=0}^{3} x_i^3 - \tilde{a}_2 x_1 x_2 x_3. \tag{5.4}$$

To relate this description to (3.4), one relates the variables $X_{1,i}$ ($i = 1, \cdots, 4$) and $X_{2,j}$ ($j = 1, 2, 3$) in (3.4) to the following Laurent monomials of the homogeneous coordinates of $\mathbb{P}^3 \times \mathbb{P}^2$,

$$X_{1,1} = \frac{y_1^3 x_1}{x_0 y_1 y_2 y_3} \ , \ X_{1,2} = \frac{y_2^3 x_2}{x_0 y_1 y_2 y_3} \ , \ X_{1,3} = \frac{y_3^3 x_3}{x_0 y_1 y_2 y_3} \ ,$$

$$X_{1,4} = \frac{x_0^3}{x_1 x_2 x_3} \ , \ X_{2,1} = \frac{x_1^3}{x_1 x_2 x_3} \ , \ X_{2,2} = \frac{x_2^3}{x_1 x_2 x_3} \ , \ X_{2,3} = \frac{x_1^3}{x_1 x_2 x_3} \ , \tag{5.5}$$

by which we can identify (3.12) with (5.4) and the parameter a_i with $\tilde{a}_i$ ($a_{ij} = 1$). Furthermore we see that (3.12) is invariant under G, which implies that the subsystem of

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the moduli space we are considering is in the invariant sector of G and part of the moduli space of the three generation model.

The mirror manifold can also be constructed as the quotient of (5.4) w.r.t. the group of order 27 generated by $(x_0, x_1, x_2, x_3, y_1, y_2, y_3; a_1, a_2) \mapsto (x_0, \beta^{3(m+n)} x_1, \beta^{6m} x_2, \beta^{6n} x_3, \beta^{-m-n} y_1, \beta^{-8m} y_2, \beta^{-8n} y_3; a_1, a_2)$ with $\beta = e^{\frac{2\pi i}{9}}$ $m, n \in \mathbb{Z}$. The full invariance group of (5.4) is generated by the above transformations and $(x_0, x_1, x_2, x_3, y_1, y_2, y_3; a_1, a_2) \mapsto (x_0, \beta^{3k} x_1, x_2, x_3, \beta^{-k} y_1, y_2, y_3; \beta^k a_1, \beta^{6k} a_2)$, $k \in \mathbb{Z}$, i.e. we have the same identification of the parameter space of the mirror manifold as we found for (3.12).

We can always find a map analogous to (5.5) which maps the Laurent polynomials $P_1, \ldots, P_l$ to quasi-homogeneous polynomials $p_1, \ldots, p_l$. In terms of those the fundamental period w_0 can be obtained as the integral [17]

$$w_0 = \prod_{i=1}^{l} (-a_i) \int \frac{1}{p_1 \cdots p_l} dx_{1,1}, \ldots dx_{k,n_k+1} \tag{5.6}$$

with integration contours $|x_{i,j}| = 1$. While we can always define a map such that w_0 is given by this expression, in most of the cases however this map cannot be chosen such that the polynomials $p_1, \ldots, p_l$ define a transverse configuration in $\mathbb{P}^{n_1}[\vec{w}^{(1)}] \times \cdots \times \mathbb{P}^{n_k}[\vec{w}^{(k)}]$. An example of such a case will be given at the end of section six.

Let us demonstrate the direct calculation of the discriminant locus for the manifold defined in (3.12) by examining the conditions of Δ-regularity which reads that for all faces Θ_1^* and Θ_2^* of Δ_1 and Δ_2^* we need $P_{\Theta_1^*} = P_{\Theta_2^*} = 0$ and that all 2×2 sub-determinants of the matrix whose two rows are $X_{i,j} \frac{\partial P_{\Theta_1^*}}{\partial X_{i,j}}$ and $X_{i,j} \frac{\partial P_{\Theta_2^*}}{\partial X_{i,j}} = 0$ (no sum) vanish simultaneously.

For convenience we rename the parameters a_i, $a_{i,j}$ occuring in P_1 and P_2 as $a_0, \ldots, a_4$ and $b_0, \ldots, b_3$. If we introduce variables $U_i \equiv U_{1,i}$ and $V_j \equiv U_{2,j}$ such that $X_{i,j} = \frac{U_{ij}}{U_{i0}}$ the conditions for non-regularity become for $\Theta_1^* = \Delta_1^*$ and $\Theta_2^* = \Delta_2^*$

$$\sum_{i=0}^{3} a_i U_i = 0, \quad \sum_{j=0}^{3} b_j V_j + a_4 U_4 = 0$$

$$\frac{a_4 U_4}{V_0 U_0}(a_i U_i - a_j U_j) = 0, \qquad \frac{a_i U_i}{U_0 V_0}(b_3 V_3 - b_k V_k) = 0, \quad i,j = 1,2,3, \quad k = 1,2 \tag{5.7}$$

Note that $\frac{U_1 U_2 U_3 U_4}{U_0^3 V_0} = \frac{V_1 V_2 V_3}{V_0^3} = 1$. There are several ways to satisfy the second set of equations. E.g. if we choose to satisfy them by equating the expressions in parantheses to zero, we find, using the variables $z_1 = \frac{a_1 a_2 a_3 a_4}{a_0^3 b_0}$, $z_2 = \frac{b_1 b_2 b_3}{b_0^3}$ and introducing $A = a_4 U_4$ and $B = b_0 V_0$, $z_1 = -\frac{1}{27} \frac{A}{B}$ and $z_2 = \frac{1}{27} \frac{(A+B)^3}{B^3}$, which satisfy $(1 - 27 z_1)^3 - 27 z_2 = 0$. Another possible choice is to set $a_4 U_4 = 0$ and $(b_3 V_3 - b_k V_k) = 0$. This leads to $1 - 27 z_2 = 0$. Other components of the discriminant can be obtained by considering other (lower-dimensional) faces of Δ_1^* and Δ_2^*.

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(*ii*) Our next example is the Tian-Yau manifold given by the configuration

$$\begin{pmatrix} \mathbb{P}^3 & \Big\| & 3 & 0 & 1 \\ \mathbb{P}^3 & \Big\| & 0 & 3 & 1 \end{pmatrix}^{14}_{-18} \quad : \quad \begin{array}{l} K^0 = 9\,J_1^2 J_2 + 9\,J_1 J_2^2\,, \\ c_2 J_1 = c_2 J_2 = 36\,. \end{array} \tag{5.8}$$

The quotient of this manifold w.r.t. the $\mathbb{Z}_3$ group acting by $(x_0, x_1, x_2, x_3, y_0, y_1, y_2, y_3) \mapsto (x_0, x_1, x_2, \alpha x_3, y_0, y_1, y_2, \alpha^2 y_3)$ $(\alpha = e^{\frac{2\pi i}{3}})$ on the homogeneous coordinates of $\mathbb{P}^3 \times \mathbb{P}^3$ yields a simple realisation of a three generation compactification, which is diffeomorphically equivalent to the one discussed in (*5 i*). A preliminary phenomenological discussion was given in [49] and [50]. For the above configuration all complex structure deformations can be represented as monomials. Hence the full moduli dependence of all Yukawa couplings and the Kähler potential for 27's and $\overline{27}$'s can be in principle calculated by a straightforward, though very tedious, application of the methods described in [16] and [15].

The vertices of the dual polyhedra are now $\nu^*_{1,1} = (1, 0, 0; 0, 0, 0), \ldots,$ $\nu^*_{2,4} = (0, 0, 0; -1, -1, -1)$ which we group into three sets as $E_1 = \{\nu^*_{1,1}, \nu^*_{1,2}, \nu^*_{1,3}\}$, $E_2 = \{\nu^*_{2,1}, \nu^*_{2,2}, \nu^*_{2,3}\}$ and $E_3 = \{\nu^*_{1,4}, \nu^*_{2,4}\}$ [9]. This corresponds to the three Laurent polynomials

$$P_1 = a_0 - a_1 X_1 - a_2 X_2 - a_3 X_3\,, \quad P_2 = b_0 - b_1 Y_1 - b_2 Y_2 - b_3 Y_3\,,$$
$$P_3 = c_0 - \frac{c_1}{X_1 X_2 X_3} - \frac{c_2}{Y_1 Y_2 Y_3}\,. \tag{5.9}$$

The period (3.7) follows straightforwardly by performing the integral (see (3.5))

$$w_0(z_1, z_2) = \int_{|X_i| = |Y_j| = 1} \frac{a_0 \, b_0 \, c_0}{P_1 P_2 P_3} \prod \frac{dY_i}{Y_i} \frac{dX_i}{X_i} \tag{5.10}$$

which gives $w_0(z_1, z_2)$ as in eq.(3.7) with $z_1 = \frac{a_1 a_2 a_3 c_1}{a_0^3 c_0}$ and $z_2 = \frac{b_1 b_2 b_3 c_2}{b_0^3 c_0}$. This corresponds to introducing extended vertices $\bar{\nu}^* = (\vec{e_i}, \nu^*)$, where $\vec{e_1} = (1, 0, 0)$, $\vec{e_2} = (0, 1, 0)$ and $\vec{e_3} = (0, 0, 1)$ for the vertices in the three sets. The linear relations between the $\bar{\nu}^*$ are then in accordance with the general formula (3.9) and read

$$l^{(1)} = (-3, 0, -1; 1, 1, 1, 1, 0, 0, 0, 0) \quad l^{(2)} = (0, -3, -1; 0, 0, 0, 0, 1, 1, 1, 1). \tag{5.11}$$

The associated differential operators

$$\mathcal{L}_1 = \theta_1^3 - 3\,(\theta_1 + \theta_2)\,(3\,\theta_1 - 1)\,(3\,\theta_1 - 2)\,z_1$$
$$\mathcal{L}_2 = \theta_2^3 - 3\,(\theta_1 + \theta_2)\,(3\,\theta_2 - 1)\,(3\,\theta_2 - 2)\,z_2 \tag{5.12}$$

can be factorized $\mathcal{L}_1 + \mathcal{L}_2 \equiv (\theta_1 + \theta_2)L_2$ with the second order operator $(L_1 = \mathcal{L}_1)$

$$L_2 = (\theta_1^2 - \theta_1\theta_2 + \theta_2^2) - 3\,(3\,\theta_1 - 1)(3\,\theta_1 - 2)\,u - 3\,(3\,\theta_2 - 1)(3\,\theta_2 - 2)\,z_2\,. \tag{5.13}$$

---

[9] The results are of course again independent of the way we group the vertices into three sets, as long as the first and second sets contain three vertices pertaining to the first and the second $\mathbb{P}^3$, respectively.

The Yukawa couplings are

$$\tilde{K}^{(3,0)} = \frac{1}{z_2^2 \Delta_0 \Delta_1^2}, \quad \tilde{K}^{(2,1)} = \frac{1}{27 \, z_1^2 z_2 \Delta_0 \Delta_1,} \tag{5.14}$$

where the general component of the discriminant surface is $\Delta_0 = 1 - 27\,z_1 - 27\,z_2$ and a second component reads $\Delta_1 = (1 - 27z_1)$. For symmetry reasons $\tilde{K}^{(0,3)}$, $\tilde{K}^{(1,2)}$ are given by the above expressions but with z_1 and z_2 exchanged. Due to this symmetry we list below the invariants n_{d_1,d_2}^r only for $d_1 \leq d_2$.

(0,1)	81
(0,2)	81
(0,3)	18
(0,4)	81
(0,5)	81
(0,6)	18

(1,1)	729
(2,2)	33534
(3,3)	5433399
(1,2)	2187
(2,4)	1708047

(1,3)	6885
(1,4)	18954
(1,5)	45927
(2,3)	300348

The number of curves up to degree three agree with those calculated by algebraic counting methods in [21]. This calculation confirms the lines $(0,1)$ and the degree $(0,2)$, $(2,1)$ curves on this CICY. According to [21] there are no $(0,3)$ curves and 567 degree $(1,1)$ curves, leaving aside the possibility of nodal cubics and degenerate rational curves respectively. More recently, all entries up to degree three have been confirmed [51]. The invariant $n_{0,3}^r = 18$ was calculated as the Euler number of the tangent bundle of the moduli space of a family of nodal cubics.

In Appendix B we present general formulas for the predicted numbers of lines for all possible manifolds of type $(5\ i)$ in $\mathbb{P}^3 \times \mathbb{P}^2$ and of type $(5\ iii)$ in $\mathbb{P}^3 \times \mathbb{P}^3$. In deriving these formulas, we utilized formula (4.18) for the instanton corrected Yukawa couplings.

Next we will treat two examples, which involve a somewhat more complicated factorisation procedure, but skip in the following most of the details.

(iii) For the manifold defined by the configuration matrix

$$\left(\begin{array}{c} \mathbb{P}^3 \\ \mathbb{P}^3 \end{array} \middle\| \begin{array}{ccc} 3 & 1 & 0 \\ 1 & 1 & 2 \end{array} \right)^{12}_{-120}$$

the linear differential operators (3.9) are, after trivial factorization, of fourth and third order, respectively. We get a second order operator by factorising $27\mathcal{L}_1 + (40\theta_1 + 13\theta_2)\mathcal{L}_2 = (\theta_1 + \theta_2)(3\theta_1 + \theta_2)L_1$ with

$$L_1 = (9\theta_1 - 12\theta_1\theta_2 + 13\theta_2^2) - 27(3\theta_1 + \theta_2 - 1)(3\theta_1 + \theta_2 - 2)\,z_1 - 2(2\theta_2 - 1)(40\theta_1 + 13\theta_2)\,z_2. \tag{5.15}$$

Due to their complexity, we refrain from giving the Yukawa couplings explicitly. The expansion of the prepotential by (4.19) does not require their knowlededge. From the first

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few term in the expansion of the prepotential

$$F = \frac{2}{6}(t^1)^3 + \frac{6}{2}t^1(t^2)^2 + \frac{8}{2}(t^1)^2 t^2 + \frac{44}{24}t^1 + \frac{48}{24}t^2 + 120\frac{i\zeta(3)}{(2\pi)^3} + 180q_1 + \frac{405}{2}q_1^2 + \frac{380}{3}q_1^3$$

$$+ 48q_2 + 876q_1q_2 + 9772q_1^2 q_2 + 46q_2^2 + 3536q_1q_2^2 + \frac{16}{9}q_2^3 + O(q^4)$$

$$(5.16)$$

we can read off the number of lines of degrees $(1,0)$ and $(0,1)$ as 180 and 48, respectively. Other curves of low degree come with multiplicity $n_{2,0}^r = 180$, $n_{0,2}^r = 40$, $n_{1,1}^r = 876$, $n_{1,2}^r = 3536$ and $n_{2,1}^r = 9672$.

(iv) As the last two moduli example we choose a model with five bilinear constraints

$$\left(\begin{array}{c} \mathbb{P}^4 \\ \mathbb{P}^4 \end{array} \left\| \begin{array}{ccccc} 1 & 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 & 1 \end{array} \right. \right)^{12}_{-32}$$

Starting from

$$\mathcal{L}_1 = \theta_1^5 - (\theta_1 + \theta_2)^5 z_1, \qquad \mathcal{L}_2 = \theta_2^5 - (\theta_1 + \theta_2)^5 z_2, \qquad (5.17)$$

we factorize in five steps, namely (1) $\mathcal{L}_1 + \mathcal{L}_2 = (\theta_1 + \theta_2)\mathcal{L}_3$, (2) $\theta_2\mathcal{L}_3 - 5\mathcal{L}_2 = (\theta_1 + \theta_2)\mathcal{L}_4$, (3) $2\mathcal{L}_3 + \mathcal{L}_4 = (\theta_1 + \theta_2)\mathcal{L}_5$, (4) $2\mathcal{L}_3 - \theta_1\mathcal{L}_5 = (\theta_1 + \theta_2)\mathcal{L}_6$, (5) $\mathcal{L}_5 + 2\mathcal{L}_6 = (\theta_1 + \theta_2)\mathcal{L}_7$ and choose eg. $\mathcal{L}_5$ and $\mathcal{L}_7$ as the third and second order differential operators L_1 and L_2, respectively. The prepotential as derived from (4.19) reads

$$F = \frac{5}{6}\sum_i (t^i)^3 + \frac{10}{2}\sum_{i,j}' t^i(t^j)^2 + \frac{20}{24}\sum_i t^i + 32\frac{i\zeta(3)}{(2\pi)^3} + \sum_i \left(50q_i + \frac{25}{4}q_i^2 + \frac{25}{32}q_i^4 \right)$$

$$+ \sum_{i,j}' \left(\frac{650}{2}q_iq_j + 1475q_i^2 q_i + 650q_i^3 q_j + \frac{117725}{8}q_i^2 q_j^2 \right) + O(q^5),$$

$$(5.18)$$

here the summation indicated with $'$ is over distinct indices only.

Hence we have in total 100 lines and the non-zero invariants of curves up to degree four are $n_{1,1}^r = n_{1,3}^r = 650$ $n_{1,2}^r = 1475$ and $n_{2,2}^r = 29350$.

(v) As a three moduli example we consider

$$\left(\begin{array}{c} \mathbb{P}^3 \\ \mathbb{P}^1 \\ \mathbb{P}^1 \end{array} \left\| \begin{array}{cc} 3 & 1 \\ 0 & 2 \\ 0 & 2 \end{array} \right. \right)^9_{-48} : \quad \begin{array}{l} K^0 = 6\, J_1^2 J_2 + 6\, J_1^2 J_3 + 3 J_1 J_2 J_3, \\ c_2 J_1 = 36, \ c_2 J_2 = c_2 J_3 = 24. \end{array}$$

for which we have from (3.9)

$$\mathcal{L}_1 = \theta_1^3 - 3(3\theta_1 - 2)(3\theta_1 - 1)(\theta_1 + 2\theta_2 + 2\theta_3)z_1$$

$$\mathcal{L}_2 = \theta_2^2 - (\theta_1 + 2\theta_2 + 2\theta_3 - 1)(\theta_1 + 2\theta_2 + 2\theta_3)z_2 \qquad (5.19)$$

$$\mathcal{L}_3 = \theta_3^2 - (\theta_1 + 2\theta_2 + 2\theta_3 - 1)(\theta_1 + 2\theta_2 + 2\theta_3)z_3.$$

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An independent second order differential operator L_4 ($L_i = \mathcal{L}_i$, $i = 1, 2, 3$) can be factorized from the system in the following way

$$\mathcal{L}_1 + (16\theta_3 - 4\theta_1)\mathcal{L}_2 + (16\theta_2 - 4\theta_1)\mathcal{L}_3 = (\theta_1 + 2\theta_2 + 2\theta_3)L_4. \qquad (5.20)$$

Due to their symmetry, we have $n^r_{d_1,d_2,d_3} = n^r_{d_1,d_3,d_2}$ and will thus list the non-zero invariants only for $d_2 \leq d_3$ in the range $d_1 + d_2 + d_3 \leq 6$.

(0,0,1)	18
(0,1,1)	60
(0,2,2)	48
(0,3,3)	60
(0,1,2)	18
(0,2,3)	18
(0,3,4)	18

(1,0,0)	216
(2,0,0)	216
(3,0,0)	48
(4,0,0)	216
(5,0,0)	216
(6,0,0)	48
(1,0,1)	216
(2,0,2)	216
(3,0,3)	48

(1,1,1)	1512
(2,2,2)	621000
(1,1,2)	1512
(1,1,3)	216
(1,2,2)	7128
(1,2,3)	7128
(2,0,1)	2106
(4,0,2)	414720
(2,1,1)	28232

(2,1,2)	85212
(2,1,3)	28232
(3,0,1)	17856
(3,0,2)	17856
(3,1,1)	656952
(3,1,2)	2984904
(4,0,1)	95094
(4,1,1)	8757828
(5,0,1)	414720

In section (7) we discuss also the invariants associated to the elliptic curves of this three moduli example.

The closed formulas (4.18) and (4.19) can be easily evaluated for general toric varities with higher dimensional moduli space provided the generators of the Mori cone $l^{(i)}$ (cf. [16]) and the intersection numbers in the corresponding basis are known. By (4.22) they become a very useful tool for enumerative geometry.

We will demonstrate this in the following for a complete intersection with a six dimensional Kähler moduli space.

(vi) For the six moduli complete intersection

$$\begin{pmatrix} \mathbb{P}^1 & \Big\| & 0 & 1 & 1 \\ \mathbb{P}^1 & \Big\| & 0 & 1 & 1 \\ \mathbb{P}^1 & \Big\| & 1 & 1 & 0 \\ \mathbb{P}^1 & \Big\| & 1 & 1 & 0 \\ \mathbb{P}^1 & \Big\| & 1 & 0 & 1 \\ \mathbb{P}^1 & \Big\| & 1 & 0 & 1 \end{pmatrix}^6_{-48}$$

straightforward evaluation of (4.19) gives immediatly the prepotential, which reads up to order four in the q_i

$$F = \frac{1}{3}\sideset{}{'}\sum_{i,j,k} t^i t^j t^k + \sum_i t^i + 48\frac{i\zeta(3)}{(2\pi)^3} + \sum_i \left(16\, q_i + 2\, q_i^2 + \frac{16}{27}\, q_i^3 + \frac{1}{4}\, q_i^4 \right)$$
$$+ \sum_{i,j} \frac{1}{2}\left(8\, q_i q_j + q_i^2 q_j^2 \right) + \frac{8}{3}\sideset{}{'}\sum_{i,j,k} q_i q_j q_k + \frac{10}{3}\sideset{}{'}\sum_{i,j,k,l} q_i q_j q_k q_l + O(q^5); \qquad (5.21)$$

here again the primed sums are to be taken over distinct indices only. From this expression we can readily predict the number of lines $n^r_{1,0,0,0,0,0} = 8$, curves of multidegree $n^r_{1,1,0,0,0,0} = 8$, $n^r_{1,1,1,0,0,0} = 16$ and $n^r_{1,1,1,1,0,0} = 80$. Other invariants follow by permutation. There are no curves of degree $(2,0,0,0,0,0)$, $(3,0,0,0,0,0)$ and $(4,0,0,0,0,0)$. The expressions for the Yukwa couplings and the Weil-Peterson metric in the large radius expansions follow from (4.11) and (4.12).

We use our next examples to demonstrate that the analysis of the large complex respectively Kähler structure limit in the previous section and formulas (4.18),(4.19) have an application also to embbedings in toric varieties with Gorenstein singularities.

(vii) We consider first the hypersurface of degree 18 in the weighted projective space $\mathbb{P}^4[6,6,3,2,1]$, which reads in the shorthand notation (2.1) as ($\mathbb{P}^4[6,6,3,2,1] \parallel 18$)$^7_{-144}$. This model was discussed qualitatively in the context of topology changes by flops in [31], but it was not solved.

Generally, for quasi-smooth hypersurfaces and complete intersections in weighted $\mathbb{P}^n$ the singularities of the ambient space intersect the Calabi-Yau hypersurface X in sets of codimension two and three, see eg.[32][25].

In our example we have a singular curve C_1 in X with a $\mathbb{Z}_2$-action on its normal bundle, which lies inside the singular stratum where the third and the fifth homogeneous coordinate of the weighted projective space are set to zero. The resolution introduces a $\mathbb{P}^1$ fibration over that curve, which gives rise to one exceptional divisor D_1. Similarly, the singular stratum of the weighted projective space where $z_4 = z_5 = 0$ intersects X in a curve C_2 with a $\mathbb{Z}_3$-action on the normal bundle, whose resolution leads to a fibration with an $\mathbb{P}^1 \wedge \mathbb{P}^1$ sphere-tree over that curve, hence two irreducible divisors D_2, D_3 on the resolved space. The singular curves meet in three $\mathbb{Z}_6$−singular points inside the stratum $z_3 = z_4 = z_5 = 0$, whose resolutions support one exceptional divisor for each point, $\tilde{D}_4, \tilde{D}_5, \tilde{D}_6$. Hence we have, including the divisor D_0 from the Picard group of the singular space, a seven moduli case.

The toric describtion of the mirror pair in terms of reflexive polyhedra was given in [26], and reviewed in [16][31]. Associated to the manifold X and its mirror X^* are the simplicial polyhedra Δ and Δ^*, which are defined as the convex hull of the following points

$$\Delta = \text{conv} \begin{pmatrix} \nu_1 & = (& 2,-1,-1,-1) \\ \nu_2 & = (-1, & 2,-1,-1) \\ \nu_3 & = (-1,-1, & 5,-1) \\ \nu_4 & = (-1,-1,-1, & 8) \\ \nu_5 & = (-1,-1,-1,-1) \end{pmatrix}, \quad \Delta^* = \text{conv} \begin{pmatrix} \nu_1^* & = (& 1, & 0, & 0, & 0) \\ \nu_2^* & = (& 0, & 1, & 0, & 0) \\ \nu_3^* & = (& 0, & 0, & 1, & 0) \\ \nu_4^* & = (& 0, & 0, & 0, & 1) \\ \nu_5^* & = (-6, -6, -3, -2) \end{pmatrix} \quad (5.22)$$

Five points in the dual polyhedron Δ^*, namely the point in its interior; $\nu_0^* = (0,0,0,0)$ and the points on edges of codimension three: $\nu_6^* = (-3,-3,-1,-1)$, $\nu_7^* = (-2,-2,-1,0)$, $\nu_8^* = (-4,-4,-2,-1)$ and on the face of codimension two: $\nu_9^* = (-1,-1,0,0)$ can be identified with the (exceptional) divisors of the Calabi-Yau hypersurface X. Only for these divisors the toric describtion of the Kähler cone and its dual, the Mori cone, can be applied[10]. We prefer to work with the Mori cone as its generators are directly related

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[10] This technical complication could have been avoided if we had instead considered the model

to the expressions for the PF equations and the definition of the local coordinates for the large complex structure limit.

The $\mathbb{Z}_6$ singularity of type $\mathbb{C}^3/G$ with $G = \mathrm{diag}(\alpha, \alpha^2, \alpha^3)$ $(\alpha = \exp(\frac{2\pi i}{6}))$ is described in this construction by the three dimensional cone Σ, spanned from the origin ν_0^* by $\{\vec{\nu}_3^*, \vec{\nu}_4^*, \vec{\nu}_5^*\}$. The cone Σ is not basic, i.e. the three vectors $\{\vec{\nu}_3^*, \vec{\nu}_4^*, \vec{\nu}_5^*\}$ do not generate the lattice $\Sigma \cap M$, with $M = N = \{(n_1, n_2, n_3, n_4)|n_i \in \mathbb{Z}\}$. The volume of $\sigma \cap \Delta^* = \mathrm{conv}(\nu_0^*, \nu_3^*, \nu_4^*, \nu_5^*)$ is six, which is the order of the defining group for this singularity. $\Sigma \cap \Delta^*$ contains the point ν_6^* on the edge $\{\nu_3^*, \nu_5^*\}$, which can be identified with the exceptional divisor D_1 in the resolution over the curve C_1, as well as the points ν_7^*, ν_8^* on the edge $\{\nu_4^*, \nu_5^*\}$, which correspond to D_2, D_3. On the other hand, the point ν_9^* inside the triangle with corners $\{\nu_3^*, \nu_4^*, \nu_5^*\}$ corresponds to only one divisor D_4 in the resolution of three $\mathbb{Z}_6$ singular points. In the formula for the dimension of $H^{1,1}(X)$ [26] the multiplicity three is taken into account by an additional term, which multiplies the interior points in the three dimensional cone Σ with the number of points on the interior of its two dimensional dual, which in this case is spanned by the origin in Δ and $\{\vec{\nu}_1, \vec{\nu}_2\}$. So it contains two points $\delta_1 = (1, 0, -1, -1)$ and $\delta_2 = (0, 1, 1, -1)$.

There are five different canonical resolutions of the $\mathbb{C}^3/\mathbb{Z}_6$ singularity, corresponding to the five possible subdivison of the cone Σ into six basic cones of volume one, which are spanned from ν_0^* by

$$A: \quad \{\vec{\nu}_3^*, \vec{\nu}_6^*, \vec{\nu}_9^*\}, \{\vec{\nu}_6^*, \vec{\nu}_5^*, \vec{\nu}_8^*\}, \{\vec{\nu}_8^*, \vec{\nu}_9^*, \vec{\nu}_6^*\}, \{\vec{\nu}_8^*, \vec{\nu}_7^*, \vec{\nu}_9^*\}, \{\vec{\nu}_7^*, \vec{\nu}_4^*, \vec{\nu}_9^*\}, \{\vec{\nu}_4^*, \vec{\nu}_3^*, \vec{\nu}_9^*\}$$

$$B: \quad \{\vec{\nu}_3^*, \vec{\nu}_6^*, \vec{\nu}_8^*\}, \{\vec{\nu}_6^*, \vec{\nu}_5^*, \vec{\nu}_8^*\}, \{\vec{\nu}_8^*, \vec{\nu}_9^*, \vec{\nu}_3^*\}, \{\vec{\nu}_8^*, \vec{\nu}_7^*, \vec{\nu}_9^*\}, \{\vec{\nu}_7^*, \vec{\nu}_4^*, \vec{\nu}_9^*\}, \{\vec{\nu}_4^*, \vec{\nu}_3^*, \vec{\nu}_9^*\}$$

$$C: \quad \{\vec{\nu}_3^*, \vec{\nu}_6^*, \vec{\nu}_9^*\}, \{\vec{\nu}_6^*, \vec{\nu}_5^*, \vec{\nu}_9^*\}, \{\vec{\nu}_5^*, \vec{\nu}_8^*, \vec{\nu}_9^*\}, \{\vec{\nu}_8^*, \vec{\nu}_7^*, \vec{\nu}_9^*\}, \{\vec{\nu}_7^*, \vec{\nu}_4^*, \vec{\nu}_9^*\}, \{\vec{\nu}_4^*, \vec{\nu}_3^*, \vec{\nu}_9^*\}$$

$$D: \quad \{\vec{\nu}_3^*, \vec{\nu}_6^*, \vec{\nu}_9^*\}, \{\vec{\nu}_6^*, \vec{\nu}_5^*, \vec{\nu}_8^*\}, \{\vec{\nu}_8^*, \vec{\nu}_7^*, \vec{\nu}_6^*\}, \{\vec{\nu}_7^*, \vec{\nu}_9^*, \vec{\nu}_6^*\}, \{\vec{\nu}_7^*, \vec{\nu}_4^*, \vec{\nu}_9^*\}, \{\vec{\nu}_4^*, \vec{\nu}_3^*, \vec{\nu}_9^*\}$$

$$E: \quad \{\vec{\nu}_3^*, \vec{\nu}_6^*, \vec{\nu}_9^*\}, \{\vec{\nu}_6^*, \vec{\nu}_5^*, \vec{\nu}_8^*\}, \{\vec{\nu}_8^*, \vec{\nu}_7^*, \vec{\nu}_6^*\}, \{\vec{\nu}_7^*, \vec{\nu}_4^*, \vec{\nu}_6^*\}, \{\vec{\nu}_4^*, \vec{\nu}_9^*, \vec{\nu}_6^*\}, \{\vec{\nu}_4^*, \vec{\nu}_3^*, \vec{\nu}_9^*\}.$$

$$(5.23)$$

for the different resolutions. The generators of the Mori cone $l^{(i)}$ and the Picard-Fuchs system in the corresponding large complex structure coordinates z_i can be obtained for each subdivision by the methods described in [16]. They will of course depend on the subdivision.

The subdivions are connected by flops, e.g. subdivision B is obtained from subdivison A by the flop, which blows down the $\mathbb{P}^1$ represented by the subcone spanned from ν_0^* by $\{\vec{\nu}_6^*, \vec{\nu}_9^*\}$ in subdivision A and subsequently blows up the $\mathbb{P}^1$ associated with the new subcone spanned from ν_0^* by $\{\vec{\nu}_3^*, \vec{\nu}_8^*\}$ in subdivision B. We therefore expect simple relations of the generators of the Mori cones and large complex structure coordinates for all subdivisions to those of subdivision A. After extending the vectors ν_i^* to $\bar{\nu}_i^* = (1, \nu_i^*)$

$(\mathbb{P}^4[1, 2, 3, 12, 18] \| 36)^5_{-360}$, which has same type of singularities, but all divisors can be described by toric geometry in this case.

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the generators of the Mori cone can be constructed as described in [30],[16] to yield

$$l_A^{(1)} = (-3; 1, 1, 0, 0, -1, \quad 1, \quad 0, \quad 1, \quad 0), \; l_A^{(2)} = (\quad 0; 0, 0, 1, 0, 0, -1, \quad 0, \quad 1, -1)$$

$$l_A^{(3)} = (\quad 0; 0, 0, 0, 0, \quad 1, -1, \quad 0, -1, \quad 1), \; l_A^{(4)} = (\quad 0; 0, 0, 0, 0, 0, \quad 1, \quad 1, -1, -1). \quad (5.24)$$

$$l_A^{(5)} = (\quad 0; 0, 0, 0, 1, \quad 0, \quad 0, -2, \quad 1, \quad 0)$$

They define the large complex structure variables as $z_1^{(A)} = -a^{l_A^{(1)}}$ and $z_i^{(A)} = a^{l_A^{(i)}}$ for $i = 2, 3, 4, 5$. As next step we set up the $\bar{\Delta}^*$-hypergeometric system (see e.g. [26][16] for details). Every linear relation among the $\bar{\nu}_i^*$, not just the the ones encoded in the $l_A^{(i)}$ (comp. (3.2)), defines a linear differential operator of this system, which is satisfied by all periods. It should be mentioned again that this system is not equivalent to the Picard-Fuchs system, the nontrivial task will be to find, possibly after factorisation, a minimal systems of differential operators which uniquely determine the $2k + 2$-dimensional period vector in the domain containing the points $z^{(A, \dots, E)} = 0$. These systems are equivalent to the first order Gauss-Manin systems in these domains and all other differential operators are elements of the local left ideal associated to it (comp. section four). As we will see, the minimal numbers of differential operators generating this ideal can vary from domain to domain. The consideration of the rings $R^{(A, \dots, E)}$ will be essential in order find these generators.

From the fact that the free polynomial ring $\mathbb{C}[\theta_z^{(A)}]$ has to be truncated at degree two from fifteen to five elements by the principal part of the linear differential operators, we know that there must be ten second order differential operators. Indeed we find that the linear relations between the sites of the points in $\bar{\Delta}^*$ which are expressed by the following vectors $l_A^{(2)}$, $l_A^{(3)}$, $l_A^{(4)}$, $l_A^{(5)}$, $l_A^{(2)} + l_A^{(3)}$, $l_A^{(2)} + l_A^{(4)}$, $l_A^{(3)} + l_A^{(4)}$, $l_A^{(4)} + l_A^{(5)}$, $l_A^{(3)} + l_A^{(4)} + l_A^{(5)}$ define directly nine second order differential operators of the $\bar{\Delta}^*$-hypergeometric system (cf.[26][16])

$$L_1^{(A)} = \theta_{a_3}\theta_{a_8} - a^{l_A^{(2)}}\theta_{a_6}\theta_{a_9}, \qquad\qquad L_2^{(A)} = \theta_{a_5}\theta_{a_9} - a^{l_A^{(3)}}\theta_{a_6}\theta_{a_8},$$

$$L_3^{(A)} = \theta_{a_6}\theta_{a_7} - a^{l_A^{(4)}}\theta_{a_8}\theta_{a_9}, \qquad\qquad L_4^{(A)} = \theta_{a_4}\theta_{a_8} - a^{l_A^{(5)}}(\theta_{a_7} - 1)\theta_{a_7},$$

$$L_5^{(A)} = \theta_{a_3}\theta_{a_5} - a^{l_A^{(2)}+l_A^{(3)}}(\theta_{a_6} - 1)\theta_{a_6}, \quad L_6^{(A)} = \theta_{a_3}\theta_{a_7} - a^{l_A^{(2)}+l_A^{(4)}}(\theta_{a_9} - 1)\theta_{a_9}, \quad (5.25)$$

$$L_7^{(A)} = \theta_{a_5}\theta_{a_7} - a^{l_A^{(3)}+l_A^{(4)}}(\theta_{a_8} - 1)\theta_{a_8}, \quad L_8^{(A)} = \theta_{a_4}\theta_{a_6} - a^{l_A^{(4)}+l_A^{(5)}}\theta_{a_7}\theta_{a_9},$$

$$L_9^{(A)} = \theta_{a_4}\theta_{a_5} - a^{l_A^{(3)}+l_A^{(4)}+l_A^{(5)}}\theta_{a_7}\theta_{a_8}.$$

The logarithmic derivatives $\theta_{a_i} := a_i \frac{\partial}{\partial a_i}$ can be readily translated to logarithmic derivatives in the z_i by the identity $\theta_{a_i} = \sum_{k=1}^{h} l_i^{(k)}\theta_k$. Via this relation we also see that we can factor a θ_1 operator from the differential operator for $l_A^{(1)} + l_A^{(3)}$ to yield the tenth second order differential operator $L_{10}^{(A)}$

$$\theta_1 L_{10}^{(A)} = \theta_{a_1}\theta_{a_2}\theta_{a_9} - a^{l_A^{(1)}+l_A^{(3)}}(\theta_{a_0} - 3)(\theta_{a_0} - 2)(\theta_{a_0} - 1)$$

$$= \theta_1 \big(\theta_1(\theta_2 - \theta_3 + \theta_4) + 3z_1 z_2(3\theta_1 + 2)(3\theta_1 + 1) \big). \qquad (5.26)$$

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While $\mathbb{C}[\theta_z]$ is truncated by the ideal $I_s(\theta_z)$, $s = 1, \ldots, 10$ obtained from (5.25),(5.26) to five elements at degree two, one can easily check that we have still two elements at degree three. So we need one further independent 3rd order differential operator. We find e.g. that $l_A^{(1)}$ leads, after factorization of θ_1, to the operator

$$L_{11}^{(A)} = \theta_1(\theta_1 - \theta_2 - \theta_3 + \theta_4)(\theta_1 + \theta_2 - \theta_3 - \theta_4 + \theta_5) + 3z_1(3\theta_1 + 2)(3\theta_1 + 1)(\theta_1 - \theta_3),$$
$$(5.27)$$

which truncates $\mathbb{C}[\theta_z]$ it to one element at degree three and no element at any higher order. Now we can construct, as in chapter four, via the ring $R^{(A)}$ the period vector and calculate the prepotential explicitly. The classical intersection part and the corresponding expression for the deformed prepotential can be found most simply from (4.18) to be, up to terms of order q^4,

$$\begin{aligned}
F^{(A)} = {} & \frac{18}{3!}(t^1)^3 + \frac{9}{2}(t^1)^2 t^2 + \frac{3}{2}t^1(t^2)^2 + \frac{21}{2}(t^1)^2 t^3 + 9t^1 t^2 t^3 + \frac{3}{2}(t^2)^2 t^3 + \frac{21}{2}t^1(t^3)^2 \\
& + \frac{9}{2}t^2(t^3)^2 + \frac{21}{3!}(t^3)^3 + \frac{12}{2}(t^1)^2 t^4 + 6t^1 t^2 t^4 + 12t^1 t^3 t^4 + 6t^2 t^3 t^4 + \frac{12}{2}(t^3)^2 t^4 + \frac{6}{2}t^1(t^4)^2 \\
& + \frac{6}{2}t^3(t^4)^2 + \frac{6}{2}(t^1)^2 t_5 + 3t^1 t^2 t^5 + 6t^1 t^3 t^5 + 3t^2 t^3 t^5 + \frac{6}{2}(t^3)^2 t^5 + 3t^1 t^4 t^5 + 3t^3 t^4 t^5 \\
& + \frac{1}{24}(72t^1 + 36t^2 + 78t^3 + 48t^4 + 24t^5) + 144\frac{i\zeta(3)}{(2\pi)^3} + (27\,q_1 + 3\,q_2 + 3\,q_3 + 3\,q_4) \\
& + \left(\frac{3\,q_2{}^2}{8} - \frac{405\,q_1{}^2}{8} + 108\,q_1\,q_3 + \frac{3\,q_3{}^2}{8} - 6\,q_2\,q_4 + \frac{3\,q_4{}^2}{8} + 3\,q_4\,q_5\right) \\
& + \left(244q_1{}^3 + \frac{q_2{}^3}{9} + 81q_1{}^2 q_3 + 27q_1 q_2 q_3 + \frac{q_3{}^3}{9} + 27q_1 q_3 q_4 + 3q_2 q_3 q_4 + \frac{q_4{}^3}{9} - 6q_2 q_4 q_5\right).
\end{aligned}$$
$$(5.28)$$

We can next use the mirror hypothesis to interpret the coefficients of the cubic terms in t^i of (5.28) as triple intersection numbers of $H_2(X, \mathbb{Z})$. Note that all triple intersection numbers are positive, as expected in the bases J_i which generates the Kähler cone. The linear and the quadratic terms in the t^i obtained by (4.20), (4.21) can be compared with the classical calculation of these topological numbers performed e.g. in [25]. These formulas relate all topological data and the $l^{(i)}$ and provide an excellent check for our calculation. As it stands, the classical part of $F^{(A)}$ does not refer to the basis of divisors D_i for which these intersection numbers can be directly obtained by the formulas for resolved complete intersections in toric varieties summarized e.g. in [16][25]. To make the comparison we can easily transform to that basis by passing to the variables t_{D_i}[11] via $t^i = \delta^{i,1}t_{D_0} + \sum_{j=1}^{k-1} l_{A,j+5}^{(i)} t_{D_j}$, i.e. from the $z_i^{(A)}$ variables to the complex structure deformation variables which occur explicitly in the Laurent polynomial, namely $\frac{1}{a_0^3}, a_6, a_7, a_8, a_9$. This

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[11] Note that this reparametrisation is quite different from the transformations into different domains, discussed below, because we still keep $z_i = 0$ as the point around which we expand our solutions of the PF-equation.

leads to the prediction

$$F_{\rm cl}^{(A)} = \frac{18}{3!}t_{D_0}^3 - \frac{6}{2}t_{D_0}t_{D_1}^2 - \frac{3}{2}t_{D_1}^3 - \frac{6}{2}t_{D_0}t_{D_2}^2 - \frac{3}{2}t_{D_1}^2t_{D_3} + 3t_{D_0}t_{D_2}t_{D_3} - \frac{6}{3}t_{D_0}t_{D_3}^2 -$$
$$\frac{3}{2}t_{D_1}t_{D_3}^2 - \frac{3}{3!}t_{D_3}^3 - \frac{3}{2}t_{D_1}^2t_{D_4} - \frac{6}{2}t_{D_2}^2t_{D_4} + 3t_{D_1}t_{D_3}t_{D_4} + 3t_{D_2}t_{D_3}t_{D_4} - \frac{3}{2}t_{D_3}^2t_{D_4} -$$
$$\frac{3}{2}t_{D_1}t_{D_4}^2 - \frac{3}{2}t_{D_3}t_{D_4}^2 + \frac{21}{3!}t_{D_4}^3 + \frac{1}{24}(72t_{D_0} + 6t_{D_1} + 6t_{D_3} + 6t_{D_4}) + 144\frac{i\zeta(3)}{(2\pi)^3},$$

$$(5.29)$$

which is in agreement with the formulas of the classical intersections in the basis of the D_i. Note however that all intersection numbers among the divisors on the triangle with corners $\{\nu_3^*, \nu_4^*, \nu_5^*\}$ are multiplied by three in $F_{\rm cl}^{(A)}$. This is due to the fact that the $\mathbb{Z}_6$-fixed point has multiplicity three and the toric description of this singularity by Σ refers to the symmetric combination $D_4 = \tilde{D}_4 + \tilde{D}_5 + \tilde{D}_6$ of the exeptional divisors over the three singular points on X, with $K^0_{\tilde{D}_i\tilde{D}_j\tilde{D}_k} = \begin{cases} 6 & \text{if } i = j = k \\ 0 & \text{otherwise} \end{cases}$.

To investigate the theory in different domains of the moduli space we note the simple transformation property of the θ_i operators in the Picard-Fuchs system, i.e. for $z_i^{(B)} = \prod_j(z_j^{(A)})^{m_{i,j}} \leftrightarrow l_B^{(i)} = \sum_j m_{i,j}l_A^{(i)}$ the θ_i transform as $\theta_i^{(A)} = \sum_j m_{j,i}\theta_j^{(B)}$. We are interested in transformations which lead outside the Mori cone, i.e. in which not all entries of the matrix m are positive. A quick look at (5.25) and (5.26) reveals that the possibilities for such transformations are rather restricted if we insist in a completely degenerate large complex structure limit in the new domain B, i.e. a ring structure $R^{(B)}$ with the properties discussed in chapter four. For instance, we cannot just invert z_5 (which would correspond to the replacement $l_A^{(5)} \to -l_A^{(5)}$) without generating inhomogeneous (in θ) terms in $I(\theta_z)$. which is incompatible with the required ring structure of R. It is easy to see that the only possibilities are to invert z_2, z_3, z_4 and z_4z_5, accompanigned by transformations of the other variables. These transformations correspond to the flops leading to the coordinate patches described by the Mori cone of subdivisons B, C, D and E. They form part of the secondary fan as described in [30].

We start with the inversion of z_2, which leads to subdivision B for which the generators of the Mori cone read $l_B^{(1)} = l_A^{(1)}$, $l_B^{(2)} = -l_A^{(2)}$, $l_B^{(3)} = l_A^{(2)} + l_A^{(3)}$, $l_B^{(4)} = l_A^{(2)} + l_A^{(4)}$ and $l_B^{(5)} = l_A^{(5)}$. Note that the matrix m for a single flop squares to unity and the principal parts of the transformed system truncates $\mathbb{C}[\theta_z]$ at degree two to five elements but two elements remain at degree three. We can remedy the situation by adding the operator associated to $l_A^{(1)} + l_A^{(2)} + l_A^{(3)}$ to our system (5.25)-(5.27)

$$\theta_1 L_{12}^{(A)} = \theta_{a_1}\theta_{a_2}\theta_{a_3}\theta_{a_8} - a^{l_A^{(1)}+l_A^{(2)}+l_A^{(3)}}(\theta_{a_0} - 3)(\theta_{a_0} - 2)(\theta_{a_0} - 1)\theta_{a_6}, \qquad (5.30)$$

which factorizes to a third order operator $L_{12}^{(A)}$. For subdivision C: $l_C^{(1)} = l_A^{(1)} + l_A^{(3)}$, $l_C^{(2)} = l_A^{(2)} + l_A^{(3)}$, $l_C^{(3)} = -l_A^{(3)}$, $l_C^{(4)} = l_A^{(3)} + l_A^{(4)}$ and $l_C^{(5)} = l_A^{(5)}$ no further operator has to be added. To complete our system w.r.t. subdivison D: $l_D^{(1)} = l_A^{(1)}$, $l_D^{(2)} = l_A^{(2)} + l_A^{(4)}$, $l_D^{(3)} = l_A^{(3)} + l_A^{(4)}$, $l_D^{(4)} = -l_A^{(4)}$ $l_D^{(5)} = l_A^{(4)} + l_A^{(5)}$ and E: $l_E^{(1)} = l_A^{(1)}$, $l_E^{(2)} = 2l_4^{(1)} + l_A^{(2)} + l_A^{(5)}$,

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$l_E^{(3)} = l_A^{(3)} + l_A^{(4)}$, $l_E^{(4)} = l_A^{(5)}$ $l_E^{(5)} = -l_A^{(4)} - l_A^{(5)}$ we have to add three third order operators $L_{13}^{(A)}, L_{14}^{(A)}, L_{15}^{(A)}$:

$$\theta_1 L_{13}^{(A)} = \theta_{a_1}\theta_{a_2}\theta_{a_6}\theta_{a_7} - a^{l_A^{(1)}+l_A^{(2)}+l_A^{(4)}}(\theta_{a_0}-3)(\theta_{a_0}-2)(\theta_{a_0}-1)\theta_{a_8}$$

$$\theta_1 L_{14}^{(A)} = \theta_{a_1}\theta_{a_2}\theta_{a_4}\theta_{a_6} - a^{l_A^{(1)}+l_A^{(3)}+l_A^{(4)}+l_A^{(5)}}(\theta_{a_0}-3)(\theta_{a_0}-2)(\theta_{a_0}-1)\theta_{a_7} \qquad (5.31)$$

$$L_{15}^{(A)} = \theta_{a_3}\theta_{a_4}\theta_{a_5} - a^{l_A^{(2)}+l_A^{(3)}+2l_A^{(4)}+l_A^{(5)}}\theta_{a_8}(\theta_{a_9}-1)\theta_{a_9}$$

The system $L_1^{(A)},\ldots,L_{15}^{(A)}$ contains the information, which is necessary to extract the prepotential in all large complex structure regions. The basis of solutions given in (4.10) can now be obtain explicitly in all regions. As they are all solutions to the same system of Picard-Fuchs equations, expressed in different patches of the moduli space, they are analytic continuations of each other with trivial monodromy. Especially the prepotentials $F^{(A,B,C,D,E)}$ encoded in the $2k + 2$ component of (4.10) is the same analytic function whose expansion in the domains corresponding to the different resolutions can be evaluated with the attached program. Also, for $F^{(B,C,D,E)}$ we obtain, after transformation via $t^i = \delta^{i,1}t_{D_0} + \sum_{j=1}^{k-1} l_{j+5}^{(i)} t_{D_j}$ to the D_i basis, the classical intersection numbers as calculated e.g. by the formulas of [16], with the enhancement factor three for intersections among divisors on the triangle.

We also checked, up to order eight, that the expansions of $F^{(A,B,C,D,E)}$ are compatible with an integer expansion for n^r in (4.22). E.g. from $F^{(A)}$ we obtain, up to degree eight, instanton contributions $n_{i,0,0,0,0}^r$ with alternating sign

i	1	2	3	4	5	6	7	8
n^r	27	-54	243	-1728	15255	-153576	1696086	-20053440

We can read from the prepotential that

$$n_{0,1,0,0,0}^r = n_{0,0,1,0,0}^r = n_{0,0,0,1,0}^r = n_{0,0,0,1,1}^r = 3.$$

There are no curves of degree $(0, i, 0, 0, 0)$, $(0, 0, i, 0, 0)$, $(0, 0, 0, i, 0)$ and $(0, 0, 0, i, i)$ for $i > 1$ and no curve $(0, 0, 0, 0, i)$ for all i. By comparing with the expansions in the different regions we find that the three rational curves with degrees $(0, 1, 0, 0, 0)$, $(0, 0, 1, 0, 0)$, $(0, 0, 0, 1, 0)$ and $(0, 0, 0, 1, 1)$, respectively, are those which are shrunk to zero volume and whose corresponding invariant changes sign under the process of the four possible flop operations interrelating them, starting from resolution A (cf. the discussion in [31]).

For certain directions in the Kählercone e.g. $(i, 0, i, 0, 0)$ one has a periodicity in the invariants $n_r = 108, 108, 144, \ldots$ as in the cases (5 i,ii,iii,v) before. It is remarkable that in all cases, where we obtain periodicity, always one of the numbers, the third in the scheme $a, a, b \ldots$, is the negative of the Euler number of the manifold X.

(viii) So far we have considered complete intersections in nonsingular ambient spaces and in the last example as well as in [16], hypersurfaces in ambient spaces with Gorenstein

singularities. Let us investigate in the following the more general situation of complete intersections of codimension $n - 3$ in an n dimensional singular ambient space and take the following hypersurface and complete intersection two moduli cases as examples:

$$A : (\, \mathbb{P}^4[2,2,2,1,1] \, \| \, 8 \,)^2_{-168}$$

$$B : (\, \mathbb{P}^5[2,2,2,2,1,1] \, \| \, 4,6 \,)^2_{-132} \qquad (5.32)$$

$$C : (\, \mathbb{P}^6[2,2,2,2,2,1,1] \, \| \, 4,4,4 \,)^2_{-112} \, .$$

They all have polynomial constraints of Fermat-type and exhibit the simplest singular locus, namely a singular curve with an $\mathbb{Z}_2$-action on the normal bundle induced from the ambient space, whose resolution gives as exceptional divisor a $\mathbb{P}^1$ bundle over that curve. For all complete intersections of Fermat type in a weighted space $\mathbb{P}^n[\vec{w}]$ with $w_{n+1} = 1$ we can define, in generalization of (5.22), an n-dimensional pair of simplicial reflexive polyhedra Δ, Δ^* as the convex hulls

$$\Delta = \mathrm{conv} \begin{pmatrix} \nu_1 & = (\frac{\sum_i d_i}{w_1} - 1, -1, \ldots, -1) \\ \vdots & \vdots \\ \nu_n & = (-1, \ldots, -1, \frac{\sum_i d_i}{w_n} - 1) \\ \nu_{n+1} & = (-1, \ldots\ldots\ldots, -1) \end{pmatrix}, \quad \Delta^* = \mathrm{conv} \begin{pmatrix} \nu_1^* & = (1, 0 \ldots, 0) \\ \vdots & \vdots \\ \nu_n^* & = (0, \ldots, 0, 1) \\ \nu_{n+1}^* & = (-w_1, \ldots, -w_n) \end{pmatrix},$$

$$(5.33)$$

in an n dimensional lattice. For all three examples we have, beside the orign $\nu_0^* = (0, \ldots, 0)$, exactly one additional point in Δ^*, namely $\nu_{n+2}^* = (-1, \ldots, -1, 0)$. Extending the lattice by $n - 3$ dimensions, as described in section three, we find the linear relations among the extended lattice sites $\bar{\nu}_i^*$, which are summarized in the l-vectors

$$A : l^{(1)} = (-4; 1, 1, 1, 0, 0, 1), \quad l^{(2)} = (0; 0, 0, 0, 1, 1, -2)$$

$$B : l^{(1)} = (-2, -3; 1, 1, 1, 1, 0, 0, 1), \quad l^{(2)} = (0, 0; 0, 0, 0, 0, 1, 1, -2) \qquad (5.34)$$

$$C : l^{(1)} = (-2, -2, -2; 1, 1, 1, 1, 1, 0, 0, 1), \quad l^{(2)} = (0, 0, 0; 0, 0, 0, 0, 0, 1, 1, -2)$$

(The above choices for the generators $l_{B,C}$ do not uniquely define the nef partition of the lattice points of Δ^*.) The associated GKZ-system factorize in all cases to a third and second order Picard-Fuchs equation, where the latter has the form

$$L_2 = \theta_2^2 - z_2(2\theta_2 - \theta_1)(2\theta_2 - \theta_1 + 1). \qquad (5.35)$$

Also the principal part of the third order operator is universal. We find

$$A : L_1 = \theta_1^2(2\theta_2 - \theta_1) - 4z_1(4\theta_1 + 3)(4\theta_1 + 2)(4\theta_1 + 1)$$

$$B : L_1 = \theta_1^2(2\theta_2 - \theta_1) - 6z_1(2\theta_1 + 1)(3\theta_1 + 2)(3\theta_1 + 1) \qquad (5.36)$$

$$C : L_1 = \theta_1^2(2\theta_2 - \theta_1) - 8z_1(2\theta_1 + 1)^3,$$

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while the terms proportional to z_1 signal a different structure for the expansions of the solutions arround the singularity at $\frac{1}{z_1}, \frac{1}{z_2} \to 0$, namely six pure power series solutions for A and B and solutions involving logarithms for C. The Yukawa couplings

$$K_{111} = \frac{1}{z_1^3 \Delta_0}, \ K_{112} = \frac{1 - \kappa z_1}{z_1^2 z_2 \Delta_0}, \ K_{122} = -\frac{1 - 2\kappa z_1}{z_1 z_2 \Delta_0 \Delta_1}, \ K_{222} = \frac{1 - \kappa z_1 + 4z_2 - 12\kappa z_1 z_2}{2 z_2^2 \Delta_0 \Delta_1^2}$$
$$(5.37)$$

and the components of the discriminant $\Delta_0 = (1 - \kappa z_1)^2 - 4\kappa^2 z_1^2 z_2$, $\Delta_1 = (1 - 4z_2)$ can be parametrized by $\kappa = 256, 108, 64$ for the three cases in turn.

Using the model specific data (5.34)-(5.36) in our general formulas, we get the following predictions of the topological data and the invariants of the rational curves listed here up to degree three

Model	χ	$\int c_2 J_1$	$\int c_2 J_2$	K_{111}^0	K_{112}^0	$n_{1,0}^r$	$n_{2,0}^r$	$n_{2,1}^r$	$n_{3,0}^r$	$n_{3,1}^r$	$n_{0,1}^r$
$A: -168$	56	24	8	4	640	10032	72224	288384	7539200	4
$B: -132$	60	24	12	6	360	2682	17064	35472	770280	6
$C: -112$	64	24	16	8	256	1248	7232	10496	197632	8

In all cases we observe that $n_{0,1}^r$ is the only nonvanishing invariant for $n_{0,i}^r$, generally $n_{i,j} = 0$ for $j > i$ and similar as in [15] $n_{i,j}^r = n_{i,i-j}^r \ \forall i > 0, \ j \leq [i/2]$. Furthermore we obtain that $n_i^r = \sum_j n_{i,j}^r$, where n_i^r are the invariants for the rationals curves of the models $(\ \mathbb{P}^5 \ \| \ 4,2\ )^1_{-176}$, $(\ \mathbb{P}^6 \ \| \ 3,2,2\ )^1_{-144}$ and $(\ \mathbb{P}^7 \ \| \ 2,2,2,2\ )^1_{-120}$ and $n_{i,j}^r$ are the ones for A, B and C respectively. The invariants of the elliptic curves will be evaluated in section 7. As before, we have checked that the topological numbers coincide after the change of basis with the ones calculated in [25].

In the general case the Picard-Fuchs equation will not follow as easy as above by factorisation of the GKZ system. Rather the analysis of additional symmetries of the period will be necessary similar as it is described for hypersurfaces in [16]. On the other hand the examples indicate that our description of the instanton corrected Yukawa couplings also apply to the rich class of complete intersections with Gorenstein singularities.

The higher degree invariants for all complete intersection in products of weighted projective spaces and for all other examples discussed in this section can be evaluated by the programm INSTANTON.

6. Connection with rational superconformal theories.

In this section we like to comment on different realisations of equivalent manifolds and their relation to exactly solvable superconformal theories. The sigma model on a Calabi-Yau manifold can be identified with a $(2, 2)$ superconformal two dimensional field theory, whose partition function and correlation functions are sometimes known exactly, at least at a special point in moduli space. Although more general identifications should exist, in the known examples the SCFT is a GSO projected tensor product of minimal $(2, 2)$ superconformal field theories [52]. The classification of the latter follows an *ADE*

pattern and there is a one-to-one correspondence to the classification of modality zero *ADE* singularities. The defining equation of the latter can be viewed as Landau-Ginzburg potentials for two dimensional $(2,2)$ supersymmetric field theory having the SCFT as its infrared limit. The partition functions [53] of the *ADE* superconformal models at level k, as well as their coupings are explicitly known [54]. The identification of their LG potentials is as follows:

$$
\begin{aligned}
A_k &\sim z^{k+2}, && k \in \mathbb{Z}, \\
D_k &\sim z^{\frac{k+2}{2}} + zy^2, && k \in 2\mathbb{Z} \\
E_6 &\sim y^3 + z^4, && k = 10, \\
E_7 &\sim z^3 + zy^3, && k = 16, \\
E_8 &\sim y^3 + z^5, && k = 28
\end{aligned}
\tag{6.1}
$$

and for tensor product models the LG potential is simply the sum of the corresponding LG potential terms. The central charge $c = \sum_{i=1}^{n} \frac{3k_i}{k_i+2}$ is the sum of the central charges of the factor theories and has to be nine to cancel the conformal anomaly.

In [55] a large number of identifications between GSO projected partition functions was found, among them $G_1 \equiv (A_2, A_2, A_2, A_6, A_6) \simeq G_2 \equiv (A_2, A_2, A_2, D_6, D_6)$. We will argue that this implies an identification of the full string theory at a special point in the moduli space of the hypersurface $X_1 = (\,\mathbb{P}^4[2,2,2,1,1]\,\|\,8\,)^2_{-168}$ and of the complete intersection $X_2 = \begin{pmatrix} \mathbb{P}^4 \ \Big\| \ 4 & 1 \\ \mathbb{P}^1 \ \Big\| \ 0 & 2 \end{pmatrix}^2_{-168}$.

To see this one has to perform a geometrical analog of the GSO projection in the tensor product theory on the LG model. One can either apply an heuristic path integral argument due to ref. [56][55], or gauge the LG model as proposed in [57]. Both operations involve, similar to the GSO projection in the tensor product model, an orbifoldisation and one has to be careful to end up with the same symmetry group as in the SCFT. Before GSO projection one has in the SCFT a $\mathbb{Z}_{k+2}$ symmetry in each factor model for the A_k theories and for the D_k theories if $k \in 4\mathbb{Z}+2$. These symmetries are readily identified with the symmetries generated by $z \mapsto \exp\frac{2\pi i}{k+2}z$ and $z \mapsto \exp\frac{4\pi i}{k+2}z$, $y \mapsto \exp\frac{-2\pi i}{k+2}y$ on the LG fields of the A_k and D_k models, respectively. The GSO projection on the tensor product theory is implemented by orbifoldisation with respect to the diagonal subgroup which in the above case is a $\mathbb{Z}_{\mathrm{lcm}\{k_i+2\}}$. The symmetry group of the GSO projected theory will therefore be $\mathcal{S} \times \prod_{i=1}^{5} \mathbb{Z}_{k_i+2}/\mathbb{Z}_{\mathrm{lcm}\{k_i+2\}}$, where $\mathcal{S}$ is permutation of identical factors. In the first argument [56] the orbifoldisation is replaced by a map of the variables $z_i \to \xi_i$ with constant Jacobian such that the LG potential becomes linear in one or more of the ξ_i. They can be viewed as Lagrange multipliers and integrating them out restricts the field configuration to an affine patch of a product of weighted projective spaces. For the SCFT of type G_1 we have $z_1 = \xi_1^{\frac{1}{k_1+2}}$, $z_i = \xi_i \xi^{\frac{1}{k_i+2}}$, $i = 2,\ldots,5$ with $|\frac{\partial \xi}{\partial x}| = $ const. precisely because $c = 9$ implies $\sum_{i=1}^{5} \frac{1}{k_i+2} = 1$. Integrating out ξ_1 and going back to homogenous variables yields manifolds of type X_1, i.e. hypersurfaces in $\mathbb{P}^4[\vec{w}]$. The diagonal subgroup of the phase symmetries on the z_i is now trivial in $\mathbb{P}^4[\vec{w}]$, so that we end up with the same

symmetry group as in the SCFT. Similarly, for the second type (G_2) we have $z_1 = \xi_1^{\frac{1}{k_1+2}}$, $z_i = \xi_i \xi_1^{\frac{1}{k_i+2}}$, $i = 2, 3, 4, 6$, $z_5 = \xi_5^{\frac{1}{2}}/\xi_1^{\frac{1}{2(k_1+2)}}$, $z_7 = \xi_7 \xi_5^{\frac{1}{2}}/\xi^{\frac{1}{2(k_1+2)}}$ and integrating out ξ_1, ξ_5 yields manifolds[12] of type X_2. The identification in $\mathbb{P}^4$ and $\mathbb{P}^1$ trivialize a $\mathbb{Z}_8$, s.t. the remaining symmetry is again as in the GSO projected SCFT. By the same combinatorics it is possible to introduce one or two gauge group operations respectively, which leave the superpotential invariant and lead, by the argument of [57], to the same geometrical interpretation.

Using the basis of divisors J, D for the singular hypersurface described in [16] and the explicit formulas given there we can calculate the intersection numbers in this basis. The evaluation of the second Chern form on J is given in formula (2.6). These data and their analogs for the complete intersection, calculated by (2.5), are displayed below.

	X_1			X_2
K_{JJJ}	8		$K_{J_1 J_1 J_1}$	8
K_{JJD}	0		$K_{J_1 J_1 J_2}$	4
K_{JDD}	-8		$K_{J_1 J_2 J_2}$	0
K_{DDD}	-16		$K_{J_2 J_2 J_2}$	0
$\int c_2 h_J$	56		$\int c_2 h_{J_1}$	56
$\int c_2 h_D$	8		$\int c_2 h_{J_2}$	24

The theorem of Wall, applicable for manifolds without torsion, states that X_1 and X_2 are homotopy equivalent if these topological numbers coincide, up to a linear transformation of the basis. Identifying $J_1 = J$ and $J_2 = \frac{1}{2}(J - D)$ we see that this is in fact the case. Model X_1 has been treated in great detail in refs.[16] and [15]. In fact, one can prove that X_1 and X_2 are diffeomorphically equivalent, by realising both as singular fiber spaces where the generic fiber is a K_3 over $\mathbb{P}^1$ with pairs of points identified on the latter. Similary we can find for the complete intersections B and C in one singular projective space, which were discussed in (5 $viii$), diffeomorphic realisations in products of nonsigular projective spaces. We have the following equivalences

$$A : \left(\mathbb{P}^4[2,2,2,1,1] \,\|\, 8 \right)^2_{-168} \simeq \left(\begin{matrix} \mathbb{P}^4 \\ \mathbb{P}^1 \end{matrix} \middle\| \begin{matrix} 4 & 1 \\ 0 & 2 \end{matrix} \right)^2_{-168}$$

$$B : \left(\mathbb{P}^5[2,2,2,2,1,1] \,\|\, 4,6 \right)^2_{-132} \simeq \left(\begin{matrix} \mathbb{P}^5 \\ \mathbb{P}^1 \end{matrix} \middle\| \begin{matrix} 2 & 3 & 1 \\ 0 & 0 & 2 \end{matrix} \right)^2_{-132} \qquad (6.2)$$

$$C : \left(\mathbb{P}^6[2,2,2,2,2,1,1] \,\|\, 4,4,4 \right)^2_{-112} \simeq \left(\begin{matrix} \mathbb{P}^6 \\ \mathbb{P}^1 \end{matrix} \middle\| \begin{matrix} 2 & 2 & 2 & 1 \\ 0 & 0 & 0 & 2 \end{matrix} \right)^2_{-112}.$$

---

[12] We can interprete the LG potential $W = \sum_{i=1}^5 x_i^4 + x_4 x_6^2 + x_5 x_7^2 = 0$ also as four dimensional hypersurface in the five dimensional ambient space $\mathbb{P}^5(2,2,2,2,3,3)$ with $c_1 = 8$ and consider as in [58] a restricted cohomology of this space to define the six periods, which leads to the same prepotential and hence the same physical theory.

Let us discuss the solution for the model X_2. Here we have the vertices of the dual polyhedra $\nu_{1,1}^* = (1,0,0,0;0)$, $\nu_{1,2}^* = (0,1,0,0;0)$, $\nu_{1,3}^* = (0,0,1,0;0)$, $\nu_{1,4}^* = (0,0,0,1;0)$, $\nu_{1,5}^* = (-1,-1,-1,-1;0)$ and $\nu_{2,1}^* = (0,0,0,0;1)$, $\nu_{2,2}^* = (0,0,0,0;-1)$. We group them into two sets: $E_1 = \{\nu_{1,1}^*, \nu_{1,2}^*, \nu_{1,3}^*, \nu_{1,4}^*\}$ and $E_2 = \{\nu_{1,5}^*, \nu_{2,1}^*, \nu_{2,2}^*\}$ and proceed as described in section three. This leads to

$$l^{(1)} = (-4,-1;1,1,1,1,1,0,0) \quad , \qquad l^{(2)} = (0,-2;0,0,0,0,0,1,1) \qquad (6.3)$$

We can also write down the Laurent polynomials. They are

$$P_1 = a_0 - a_1 X_1 - a_2 X_2 - a_3 X_3 - a_4 X_4$$
$$P_2 = b_0 - \frac{b_1}{X_1 X_2 X_3 X_4} - b_2 Y_1 - \frac{b_3}{Y_1} \qquad (6.4)$$

Note that due to the freedom to rescale all variables and each polynomial, there are only two relevant parameters in P_1 and P_2, corresponding to the two complex structure moduli on X_2^*. The period (3.7) follows straightforwardly by performing the integral (3.5). One also finds $z_1 = \frac{a_1 a_2 a_2 a_4 b_1}{a_0^4 b_0}$ and $z_2 = \frac{b_2 b_3}{b_0^2}$. $l^{(1)}$ and $l^{(2)}$ lead, after trivial factorization, to differential operators of orders four ($\mathcal{L}_1$) and two ($\mathcal{L}_2 = L_2$), respectively. A third order operator L_1 can be obtained via $\mathcal{L}_1 - 4\theta_1^2 \mathcal{L}_2 = (2\theta_2 + \theta_1) L_1$. The Yukawa couplings for the model X_2 are found to be

$$K^{(3,0)} = \frac{1}{z_1^3 \Delta_0} , \qquad K^{(2,1)} = \frac{1 - 256\, z_1 - 4\, z_2}{2\, z_1^2 z_2\, \Delta_0\, \Delta_1} ,$$
$$K^{(1,2)} = \frac{3 - 512\, z_1 + 4\, z_2 z_1 z_2}{\Delta_0\, \Delta_1^2} , \qquad K^{(0,3)} = \frac{1 - 256\, z_1 + 24\, z_2 - 3072\, z_1\, z_2 + 16\, z_2^2}{2\, z_2^2\, \Delta_0\, \Delta_1^3}$$
$$(6.5)$$

with

$$\Delta_0 = (1 - 256\, z_1)^2 - 4 z_2, \quad \Delta_1 = (1 - 4 z_2)$$

The resulting topological invariants for the rational and elliptic curves for the case A, B and C in (6.2) are of course the same for both realisations. Model X_1 was solved in [16],[15]. In fact, in [15] the degrees of the curves are given with respect to the basis appropriate for X_2. We can find a map analogous to the one in (5.5) and hence a configuration $\left(\begin{matrix} \mathbb{P}^4 \\ \mathbb{P}^1 \end{matrix} \Big\| \begin{matrix} 4 & 1 \\ 0 & 2 \end{matrix} \right)^2_{-168}$, which reproduces, via the integral (5.6), the expression for w_0. However, this configuration cannot be chosen to be transverse for generic points in the two dimensional subspace of the moduli space under consideration. It is also interesting to note that the coordinates used here and the one used in [16],[15] are connected by a transcendental function.

A similar comparison can be made between the two models ($\mathbb{P}^4[6,2,2,1,1] \| 12)^2_{-252}$. and $\left(\begin{matrix} \mathbb{P}^4[3,1,1,1,1] \\ \mathbb{P}^1 \end{matrix} \Big\| \begin{matrix} 6 & 1 \\ 0 & 2 \end{matrix} \right)^2_{-252}$. At special points in moduli space they correspond to the Gepner models $(A_4, A_4, A_{10}, A_{10})$ and $(A_4, A_4, D_{10}, D_{10})$, respectively. One finds again the relations $J_1 = J$ and $J_2 = \frac{1}{2}(J - D)$ between the divisors and identical topological invariants.

7. Topological one-loop partition function and the number of elliptic curves

Knowledge of the periods and the canonical coordinates allows, up to the difficulty of fixing an holomorphic ambiguity at each step, to recursively calculate higher loop topological partition funcions as was shown in [23]. To fix the ambiguity we need as gobal properties of the moduli space $\mathcal{M}$ the singularities discussed in sec. 4.

We will focus on the one loop case and calculate the expression defined in the $N = 2$ SCFT on the torus as

$$F_1 = \frac{1}{2} \int \frac{d\tau}{\tau_2} \mathrm{Tr} \, (-1)^F F_L F_R q^{L_0} \bar{q}^{\bar{L}_0}, \tag{7.1}$$

where the trace is to be taken over the left- and right moving Ramond sectors. As shown in [23] and [24] this quantity is, for the heterotic string with canonical embedding of the spin connection into the gauge group, related to the difference of the threshold corrections to the gauge couplings of the E_6 and E_8, namely $12F_1 = \Delta(E_6) - \Delta(E_8)$. Using the holomorphic anomaly equation it was shown in [22] that it can be written as

$$F_1 = \log \left[M(z, \bar{z}) |f(z)|^2 \right], \tag{7.2}$$

where the holomorphic-antiholomorphic mixing $M(z, \bar{z})$ is given by

$$\log M = \sum_{p,q} (-1)^{p+q} \frac{p+q}{2} \mathrm{Tr}_{p,q} \log\det(g) - \frac{1}{12} \mathrm{Tr}(-1)^F; \tag{7.3}$$

here g is the $t, \bar{t}$ metric introduced in [59]. It is related to the Weil-Peterson metric $G_{i\bar{j}}$ by $g_{i\bar{j}}/g_{0\bar{0}} = g_{i\bar{j}} \exp(K) = G_{i\bar{j}}$ with $G_{i\bar{j}} = \partial_i \partial_{\bar{j}} K$. For σ-models on Calabi-Yau spaces (7.2) can be rewritten as

$$F_1 = \log \left[\exp \left((3 + k - \frac{\chi}{12}) K \right) \det[G_{i\bar{j}}]^{-1} |f(z)|^2 \right]. \tag{7.4}$$

In our application we will finally understand F_1 as a function of the Kähler moduli $t, \bar{t}$ which are related to $a, \bar{a}$ or $z, \bar{z}$ by the mirror map. To fix the holomorphic ambiguity $f(z)$ one considers the large volume limit $t, \bar{t} \to \infty$ for which one has the asymptotic behaviour

$$\lim_{t, \bar{t} \to \infty} F_1 = -\frac{2\pi i}{12} \sum_{i=1}^{k} (t_i + \bar{t}_i) \int c_2 J_i. \tag{7.5}$$

It was conjectured in [22] that $F_1^{top} \equiv \lim_{\bar{t} \to \infty} F_1$ has the following expansion

$$F_1^{top} = \text{const.} - \frac{2\pi i}{12} \sum_{i=1}^{k} t_i \int c_2 J_i - \sum_{n_l} \left[2n_{d_1,\ldots,d_k}^e \log(\eta(\prod_{i=1}^{k} q_i^{d_i})) + \frac{1}{6} n_{d_1 \ldots d_k}^r \log(1 - \prod_{i=1}^{k} q_i^{d_i}) \right] \tag{7.6}$$

in terms of the Euler numbers $n_{d_1,\ldots,d_h}^r$ and $n_{d_1,\ldots,d_k}^e$ of the tangentbundle over the appropriate compactified moduli space of the mappings from $\mathbb{P}^1$ and T^2, respectively, to the Calabi-Yau space. In the case of isolated curves they count the number of rational curves

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and elliptic curves respectively. Using (4.15),(4.12) one gets in the general case as the $\bar{t} \to \infty$ limit of (7.4)

$$F_1^{top} = \log\left[\left(\frac{a_1 \cdots a_l}{\omega_0}\right)^{3+k-\chi/12} \frac{\partial(a_1 \ldots a_k)}{\partial(t_1 \ldots t_k)} f(z)\right] + \text{const}. \qquad (7.7)$$

The factor $(a_1 \cdots a_l)/\omega_0$ corresponds to the gauge choice with the fundamental period normalized as in (3.5)[13] The holomorphic anomaly is determined by the requirement that F_1^{top} has to be a regular function everywhere in moduli space except possibly at the components of the discriminante surface determined by $\Delta_i = 0$, which can be directly determined from the Laurent polynomial (cf. (5i)). Besides the components $\Delta_i = 0$ there appear also other singular loci $\delta_i = 0$ in systems of Picard-Fuchs equations, which can be understood as identification singularities of the parameters space of the Laurent polynomial as it was discussed at the end of section three. In the cases we discuss we have $\delta_i = a_i$. While it is evident from (3.5) that $\left(\frac{a_1\ldots a_k}{\omega_0}\right)$ is regular at $\delta_i = 0$ the Jacobian of the mirror map (5.5) might have singularities at δ_i. We therefore make the following general ansatz for $f(z)$

$$f(z) = \Delta_0^{r_0} \ldots \Delta_m^{r_m} \delta_1^{s_1} \ldots \delta_k^{s_k}. \qquad (7.8)$$

Of course the $\Delta_i = z_i$ singularities are included in this ansatz. Inserting this ansatz in (7.7) and comparing the leading term with (7.5) yields equations for the r_i and s_i. If the manifold happens to be transverse at $a = 0$ the powers of the δ_i, which in this case only have to compensate possible singularities of the Jacobian, can in principle be determined by analytic continuation of the periods and hence the mirror map to the point $a = 0$. In the general case we use the values the numbers $n_{d_1,\ldots,d_h}^e$ of a few elliptic curves of low polydegree w.r.t. an integral basis of divisors, typically the fact they have to vanish, to fix all parameters and predict the other numbers.

Let us first discuss as examples various one moduli cases realized as hypersurfaces X_{k_1} of degree k_1 or complete intersections $X_{k_1,\ldots,k_n}$ of multidegree $k_1,\ldots,k_n$ in a single weighted projective space. From [13][18][19]we have one component of the singular locus for all cases at $\Delta_0 = (1 - a^{\sum k_i})$. We start with a short review of the one moduli hypersurfaces in a weighted $\mathbb{P}^4$, denoted by $X_k(\vec{w})$. In this case Δ_0 is the only component of the singular locus and it was observed in [23] that $r_0 = -1/6$ for all cases yielding the following invariants for the elliptic curves:

Model	χ	$\int c_2 J$	s_1	n_1^e	n_2^e	n_3^e	n_4^e
X_5	-200	50	0	0	0	609250	37214316625
X_6	-204	42	0	0	7884	145114704	1773044315001
X_8	-296	44	1	0	41312	21464350592	1805292092664544
X_{10}	-288	34	1	280	207680680	161279120326560	103038403740690105440

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[13] F_1^{top} depends of course on k parameter a_i $i = 1,\ldots,k$. By the $\mathbb{C}^*$ symmetries, acting on the parameters of the Laurent polynomials, we can set $l - k$ of the a_i appearing in (3.5) to one, if the number of polynomial constraints l exceeds k.

Similarly for the one moduli complete intersections in ordinary projective spaces (cf. [18]) we found $r_0 = -1/6$ and the s_1 value indicated in the following table. This result was obtained by requiring $n_1^e = 0$, which holds for intersections in ordinary projective spaces, and imposing (7.5). Note that these manifolds are not transverse at $\Delta_2 = a = 0$. In order to compare with the results of section (5 $viii$) we list the rational curves for these models

Model	χ	$\int c_2 J$	s_1		$j=1$	$j=2$	$j=3$	$j=4$	$j=5$
$X_{3,3}$	-144	54	11	n_j^r	9	1053	52812	6424326	11394483834
				n_j^e	0	0	3402	5520393	482074484
$X_{4,2}$	-176	56	$\frac{28}{3}$	n_j^r	8	1280	92288	15655168	3883902528
				n_j^e	0	0	2560	17407072	24834612736
$X_{3,2,2}$	-144	60	$\frac{115}{6}$	n_j^r	12	720	22428	1611504	168199200
				n_j^e	0	0	64	265113	198087264
$X_{2,2,2,2}$	-128	64	$\frac{85}{3}$	n_j^r	16	512	9728	416256	25703936
				n_j^e	0	0	0	14752	8782848

For the complete intesections in weighted projective spaces (cf. [19]) we found that $r_0 = -1/6$ likewise and s_1 from (7.5) gives the following integral expansion for the n_d^e

Model	χ	$\int c_2 J$	s_1	n_1^e	n_2^e	n_3^e	n_4^e
$X_{4,3}$	-156	48	$\frac{67}{6}$	0	27	16124238	38170438
$X_{6,2}$	-256	52	$\frac{29}{3}$	0	-504	1228032	79275665304
$X_{4,4}$	-144	40	11	0	1408	6953728	2684185380
$X_{6,4}$	-156	32	$\frac{31}{3}$	8	258336	5966034464	1267294361302800
$X_{6,6}$	-120	22	9	360	40691736	4956204918240	616199133057629184

All invariants of the $X_{6,2}$ model are consistent with the indentification of this model with the one parameter subspace of $(\, \mathbb{P}^4[6,2,2,1,1] \,\|\, 12\,)^2_{-252}$ [15], which implies $n_i^e = \sum_j n_{i,j}^e$. Especially $n_2^e = -504$, which is the only negative invariant for a one parameter family, is in agreement with $n_{2,0}^e = n_{2,2}^e = -492$, $n_{2,1}^e = 480$ for the two parameter hypersurface (comp. [15]). We also checked that the identification of $X_{6,4}$ with the one parameter subspace of $(\, \mathbb{P}^4[4,3,2,2,1] \,\|\, 12\,)^2$ [16] holds at the one loop level. E.g. the lowest invariants of elliptic curves for the latter model are $n_{1,0}^e = -2$, $n_{1,1}^e = 6$ and $n_{2,0}^e = 762$, $n_{2,1}^e = -3060$, $n_{2,2}^e = 18918$, $n_{2,3}^e = 225096$ with the general symmetry $n_{i,j} = n_{i,3i-j}$ and $n_{i,j} = 0$ for $j > 3i$; they reproduce the first two entries for the $X_{6,4}$ model above.

In summary the topological one loop partition function for the one parameter models is given by

$$F_1^{top} = \log \left[\frac{a^{\frac{\sum_i k_i}{12} \int c_2 J}}{w_0^{4 - \frac{\chi+4}{12}}} \left(\frac{1}{a} \frac{\partial a}{\partial t} \right)^{1/2} K_{ttt}^{\frac{1}{6}} \right] . \tag{7.9}$$

Next we treat a hypersurface $X_{14}(7,2,2,2,1)$ with two moduli. The singular locus is, in addition to the lines $z_1 = 0$, $z_2 = 0$, given by

$$\Delta_0 = 1 + 27z_1 - 63z_1z_2 + 56z_1z_2^2 - 112z_1z_2^3 - (7 - 4z_2)^4 z_1^2 z_2^3$$
$$\Delta_1 = 4z_2 - 7 \tag{7.10}$$

Equation (3.6) and the analog of (3.3) from [16] $l^{(1)} = (-7; 0, 1, 1, 1, -3, -7)$ and $l^{(2)} = (0; 1, 0, 0, 0, 1, -2)$ define[14] the relation $z_1 = a_2^7/a_1^7$ and $z_2 = 7/a_2^2$ by which we transform the expressions of singular components (7.10) to the a_i variables. Beside this components we have to care about the sets $\delta_1 = a_1 = 0$ and $\delta_2 = a_2 = 0$. For the further calculation it turns out to be advantegous to get rid of the denominators in the transformed expressions (7.10) by rescaling $\Delta_1 \to a^{14}\Delta_1$ and $\Delta_2 \to b^2\Delta_2$.

The Euler number is $\chi = -240$ the Hodge numbers are $h^{1,1} = 2, h^{2,1} = 122$. We calculate $\int c_2 \wedge J_1 = 44$ and $\int c_2 \wedge J_2 = 126$, where J_1, J_2 are the basis which generates $H^2(X, \mathbb{Z})$. It is connected with the $(1,1)$ forms dual to natural basis of the divisors in the polyheder construction J, D used in [16], by $J_1 = J$ and $J_2 = \frac{1}{2}(7J - D)$. From (7.5) we get two equations

$$r_0 = -\frac{s_1}{14} - \frac{1}{42}, \quad r_1 = -\frac{s_2}{2} - \frac{17}{6}. \tag{7.11}$$

The vanishing of the numbers of curves $n^e_{0,1} = s_2$ and $n^e_{1,0} = \frac{27}{28}(2 - s_1)$ enforces $s_1 = 2$ and $s_2 = 0$. Other numbers of elliptic curves $n^e_{i,j}$ with $i < 3$ are then given in the following table, where we list for convenience also the number of rational curves $n^r_{i,j}$

$n_{i,j}$	$j = 0$	$j = 1$	$j = 2$	$j = 3$	$j = 4$	$j = 5$	$j = 6$	$j = 7$
$n^r_{0,j}$	0	28	0	$\forall j > 1$				
$n^e_{0,j}$	0	$\forall j$						
$n^r_{1,j}$	3	-56	378	14427	14427	378	-56	3
$n^e_{1,j}$	0	$\forall j$						
$n^r_{2,j}$	-6	140	-1512	9828	-69804	500724	29683962	68588248
$n^e_{2,j}$	0	0	0	0	0	378	6496	27564

Here we have a symmetry $n_{i,j} = n_{i,7i-j}$ and $n_{i,j} = 0$ for $j > 7i$. The identification of the one parameter subspace of this model via $n_i = \sum_j n_{i,j}$ with the hypersurface $(\mathbb{P}^4[1,1,1,1,4] \,\|\, 8\,)^1_{-256}$ observed in [16] can also be checked at one-loop level.

Next we treat a hypersurface in a product of two projective spaces

$$\left(\begin{matrix} \mathbb{P}^2 \,\|\, 3 \\ \mathbb{P}^2 \,\|\, 3 \end{matrix} \right)^2_{-162}.$$

The number $n^r_{d_1,d_2}$ of rational curves of bidegree (d_1, d_2) was obtained in [16]. The evaluation of the second Chern class on J_1, J_2 can be read off from (2.5); $\int c_2 \wedge J_1 = \int c_2 \wedge J_2 = 36$.

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[14] Note the scale factor 7 introduced in [16] in order to simplify (7.10).

Because of the symmetry we have $s_1 = s_2$. In this case we have to consider only one component

$$\Delta_0 = 1 - (1 - z_1)^3 + (1 - z_2)^3 + 3z_1 z_2 (z_1 + z_2 + 7) \qquad (7.12)$$

of the discriminant of the complete intersection, where the connection with the parameters a_1, a_2 is encoded in the $l^{(i)}$ given by (3.3), $z_1 = 3^3/a_1^3$ and $z_2 = 3^3/a_2^3$. Again after clearing the denominator, the comparison with the large radius limit gives us the relation $r_0 = -\frac{7}{6} - \frac{s_1}{9}$, leaving us with one unknown constant s_1, which is determined by requiring $n_{0,1}^e = \frac{9}{2}(9 + s_1) = 0$. The following table contains the numbers of rational and elliptic curves up to bidegree $d_1 + d_2 \leq 6$. Because of the exchange symmetry we list only $n^i(d_1, d_2)$ with $d_1 \leq d_2$.

$(0,1)$	n^r	189
	n^e	0
$(0,2)$	n^r	189
	n^e	0
$(0,3)$	n^r	162
	n^e	3
$(0,4)$	n^r	189
	n^e	0
$(0,5)$	n^r	189
	n^e	0
$(0,6)$	n^r	162
	n^e	0

$(1,1)$	n^r	8262
	n^e	0
$(1,2)$	n^r	142884
	n^e	0
$(1,3)$	n^r	1492290
	n^e	-378
$(1,4)$	n^r	11375073
	n^e	-16524
$(1,5)$	n^r	69962130
	n^e	-285768

$(2,2)$	n^r	1310892
	n^e	8262
$(2,3)$	n^r	516953097
	n^e	1519434
$(2,4)$	n^r	12289326723
	n^e	71809416
$(3,3)$	n^r	55962304650
	n^e	818388234

Here we observe $n_{1,3}^e = -2\, n_{0,1}^r$, $n_{2,2}^e = n_{1,1}^r$ and $n_{1,d}^e = -2\, n_{1,d_2-3}^r$ for $d_2 \geq 4$.

For the three generation complete intersection case (5 i) the relevant components of the discriminant surface were obtained by different methods above: $\Delta_0 = (1 - 27 z_1)^3 - 27 z_2$, $\Delta_1 = (1 - 27 z_2)$. Proceeding as before we get $r_0 = -\frac{1}{6} - \frac{s_2}{9}$, $r_1 = 3 - s_1 + s_2$ and the vanishing of $n_{1,0}^e = \frac{9}{2} s_2$ requires $s_2 = 0$ and then from $n_{1,1}^e = \frac{3}{2}(9 s_1 - 27) = 0$ we get $s_1 = 3$. The corresponding predictions for n_{d_1,d_2}^e are

(1,0)	0
(2,0)	0
(3,0)	3
(4,0)	0
(5,0)	0
(6,0)	0

(0,1)	0
(0,2)	-27
(0,3)	81
(0,4)	-324
(0,5)	1728
(0,6)	-8955

(1,1)	0
(2,2)	-16028
(3,3)	-124719
(2,1)	0
(4,2)	-924372
(3,1)	-126
(4,1)	-1944
(5,1)	-30618

(1,2)	972
(2,4)	-426222
(3,2)	159678
(1,3)	-486
(2,3)	27945
(1,4)	22356
(1,5)	-72900

For the Tian-Yau manifold example (5 ii) the discriminant can be read off from the Yukawa couplings $\Delta_0 = 1 - 27\,z_1 - 27\,z_2$, $\Delta_1 = (1 - 27\,z_1)$, $\Delta_2 = (1 - 27\,z_2)$. Obviously $r_1 = r_2$ and $s_1 = s_2$ and (7.5) yields $r_0 = \frac{1}{2} + \frac{s_1}{3} - r_1$. The following predictions are obtained by imposing $n^e_{0,1} = \frac{9}{2}(s_1 - 2) = 0$ and $n^3_{1,1} = \frac{243}{2}(3r_1 - 4) = 0$

(0,1)	0						
(0,2)	-27	(1,1)	0				
(0,3)	81	(2,2)	-2916	(1,3)	1458		
(0,4)	-324	(3,3)	108180	(1,4)	4374		
(0,5)	1728	(1,2)	324	(1,5)	15066		
(0,6)	-8955	(2,4)	-31104	(2,3)	-13176		

Note that invariants $n^e_{0,i}$ of the example (5 i) coincide with $n^e_{i,0} = n^e_{0,i}$ of the present one.

The three moduli example in section (5 v) has the general discriminant

$$\Delta_0 = 1 - 108z_1 + 4374z_1^2 - 78732z_1^3 + 531441z_1^4 - 8z_2 + 432z_1 z_2$$
$$- 5832z_1^2 z_2 + 16z_2^2 - 8z_3 + 432z_1 z_3 - 5832z_1^2 z_3 - 23z_2 z_3 + 16z_3^2 \tag{7.13}$$

and a second component of the discriminant locus

$$\Delta_1 = 1 - 8z_2 + 16z_2^2 - 8z_3 - 32z_2 z_3 + 16z_3^2. \tag{7.14}$$

As a slight technical simplification we replace $\left(\frac{a_1 \dots a_k}{w_0}\right)$ by $\frac{1}{w_0}$ in (7.7) , set $\delta_i = z_i$ in (7.8) and work throughout with the large complex structure parameters z_i, which is possible as the z_i are the good coordinates on $\mathcal{M}$ (cf. sect. 3).

By (7.5) we can fix in this case $s_1 = -4$, $s_2 = s_3 = -3$ and from $n^e_{0,0,1} = -9(1+6r_0) = 0$ we have again $r_0 = -1/6$. Enforcing also $n^e_{1,0,0} = -4(r_0 + r_1) = 0$ we get the following predictions for the non zero invariants of the elliptic curves up to bidegree seven

(0,1,1)	-12	(2,1,3)	-3024		
(0,2,2)	15	(2,2,3)	-2052	(4,0,1)	-432
(1,1,1)	288	(3,0,0)	4	(4,0,2)	-4212
(2,2,2)	-13284	(3,0,1)	-36	(4,0,3)	-432
(1,1,2)	2888	(3,0,2)	-36	(4,1,1)	-464296
(1,2,2)	2160	(3,0,3)	85212	(4,1,2)	-330300
(1,2,3)	2160	(3,1,1)	17580	(5,0,1)	-4212
(1,3,3)	11664	(3,1,2)	48024	(5,0,2)	-191484
(2,1,1)	-3024	(3,1,3)	48024	(5,1,1)	-46008
(2,1,2)	-4320	(3,2,2)	991536	(6,0,1)	-35820

We finally evaluate the invariants of the elliptic curves for the hypersurface and the complete intersections in one singular projective space discussed in (5 *viii*). The topological data can be found in the Table of section (5 *viii*) and the discriminant may be read of from (5.37)

Model	r_i	s_i	$n_{i,j}^e$	$j = 0$	$j = 1$	$j = 2$
$A:$	$r_0 = -\frac{1}{6}$	$s_1 = -\frac{17}{6}$		$n_{i,j}^e = 0,\ i < 3$		
	$r_1 = -\frac{5}{6}$	$s_2 = -3$	$i = 3$	-1280	2560	2560
			$i = 4$	-317864	1047280	15948240
$B:$	$r_0 = -\frac{1}{6}$	$s_1 = -6$		$n_{i,j}^e = 0,\ i < 2$		
	$r_1 = -1$	$s_2 = -3$	$i = 2$	-16	48	-16
			$i = 3$	-5364	18972	237897
$C:$	$r_0 = -\frac{1}{6}$	$s_1 = -\frac{19}{6}$		$n_{i,j}^e = 0,\ i < 4$		
	$r_1 = -\frac{7}{6}$	$s_2 = -3$	$i = 4$	-280	1120	13072
			$i = 5$	-20992	119808	429608

The invariants for case A are in accordance with [15]. Also for the complete intersections we obtain for the invariants $n_{i,j}^e = 0$ for $j > i$ as well as the symmetry $n_{i,j}^e = n_{i,j-i}^e$. Comparison with the models $(\mathbb{P}^5 \| 4, 2)_{-176}^1$, $(\mathbb{P}^6 \| 3, 2, 2)_{-144}^1$ and $(\mathbb{P}^7 \| 2, 2, 2, 2)_{-120}^1$ reveals the expected relation $n_i^e = \sum_j n_{i,j}^e$, which is an affirmative consistency check of the calculations at one loop level.

In all our examples we observed that the holomorphic anomaly is of index $r_0 = -1/6$ at the general component of the discriminante. This also holds true for examples with higher dimensional moduli spaces and for models with a somewhat different type of singularity such as $(\mathbb{P}^4[9, 6, 1, 1, 1] \| 18)_{-540}^2$ or $(\mathbb{P}^4[12, 8, 2, 1, 1] \| 24)_{-480}^3$. This is related to the fact that the manifolds approach always a nodal configuration along this component of the discriminant. The exponent of the holomorphic ambiguity seems to be universal for this type of singularity.

8. Discussion

To extend the discussion of mirror symmetry to CICY manifolds with higher dimensional moduli spaces, we have described how to set up the Picard-Fuchs equations and specified the point of maximal unipotent monodromy. We have developed a convenient way to construct all its solutions around this point by showing the equilvalence of the solutions with the elements of special representatitves of the ring $\mathcal{R}$. In fact, the top element of $\mathcal{R}$ corresponds to the cubic monomials of intersection numbers for the generating elements of $H^2(\hat{X}, \mathbb{Z})$. We have found very simple formulas for the instanton corrected intersection numbers from which the number of rational curves can be obtained, always assuming that mirror symmetry is correct. In this paper we have focussed on the region of the moduli space of large Kähler and complex structure. The extension to the whole moduli space requires analytical continuation of the periods. This is in principle straightforward using an integral representation of the solutions of the Picard-Fuchs equations.

Technically this is however rather involved and was performed so far only for one moduli cases [11][13][14] and for two types of two moduli models [15].

The form of the topological partition function is fixed by the holomorphic anomaly equation up to a holomorphic function. To specify the latter we had to analyse their asymptotic behaviour at the singular locus of the Picard-Fuchs system. In general we had to use the vanishing of the elliptic curves of low degree to provide this information.

From a more technical point of view the method described here requires, for the prediction of the instanton expansion in the large Kähler structure limit, only the generators of the Mori cone $l^{(i)}$ and the associated intersection numbers. Given these data the expansion for the corrected Yukawa couplings and the prepotential can be simply obtained via (4.18) (4.19). As it bypasses the evaluation of the Yukawa couplings on the complex structure side it is applicable to higher dimensional moduli spaces where the evaluation of these complicated algebraic expressions is extremely tedious. It should be clear by example (5 *vii*) and (5 *viii*) that the data mentioned above can be provided more generally for the moduli associated to algebraic deformations in the general class of Calabi-Yau manifolds representable as hypersurfaces or complete intersections in toric varieties[15]. It would be interesting to see if these data can also be obtained for the twisted sectors of the models in the above class and for those of the 3345 Landau-Ginzburg models with more then five fields constructed in [62], which can not be reduced to CY-threefolds as the examples in section 6 and others in [55], but have only an interpretation as higher dimensional manifolds with $c_1 > 0$.

The calculation of the prepotential was discussed. This problem can also be considered from the point of view of topological field theory. In fact in many cases the prepotential can be obtained from the axioms of topological field theory contained in the Witten-Dijkgraaf-Verlinde-Verlinde equations [63]. However for this approach the Calabi-Yau manifolds are a critical case because here the operator algebra of topological field theory is nilpotent, at least if we restrict ourself to the massless perturbations. For the threefolds the information from the WDVV equations just defines special geometry, but do not give further information on the prepotential.

From the identification of the two hypersurfaces $(\,\mathbb{P}^4[2,2,2,1,1]\,\|\,8\,)^2_{-168}$ and $(\,\mathbb{P}^4[6,2,2,1,1]\,\|\,12\,)^2_{-252}$ with complete intersections, which was also confirmed by comparison with the Gepner models, we are taught to view the system $L_1,\ldots,L_k$, which contains (at the point $z = 0$) the information about the ring $\mathcal{R}$, the object of primary interest.

Going one step further back one may see as the basic input the data of a Riemann-Hilbert problem with sympletic integral representations of the monodromy and singular points, where the solutions are characterised by a graded ring whose homogeneous subspaces are of suitable type. In fact these data seem to encode all topological data required for the classification of the homotopy type of families of Calabi-Yau manifolds by the theorem of Wall and might lead to a refined classification of $N = 2$ vacua of string theory and

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[15] In fact, the methods can be easily extended to evaluate threepoint functions on higher dimensional Calabi-Yau spaces with higher dimensional moduli spaces [60]. Some results for one parameter families were obtained in [61]

Calabi-Yau threefolds.

Appendix A. The pole structure in the coefficients of the logarithmic solutions to the Picard-Fuchs equation

In this appendix we exhibit the pole structures of the logarithmic solutions of the Picard-Fuchs equations.

The starting point are the generators of the Mori cone $\vec{l}^{(\alpha)}$ which are all of the form [16]

$$\vec{l}^{(\alpha)} = (-\{l_{0j}^{(\alpha)}\}, \{l_i^{(\alpha)}\}) \quad j = 1, \ldots, \text{number of polynomials}$$

$$\alpha = 1, \ldots, h^{1,1}$$

where

$$-\sum_j l_{0j}^{(\alpha)} + \sum_i l_i^{(\alpha)} = 0$$

$$l_i^{(\alpha)} \in \mathbf{Z}$$

$$l_{0j}^{(\alpha)} \in \mathbf{Z}_{\geq}$$

The fundamental period is

$$w_0(z) = \sum_n c(n) z^n$$

where the sum is over all non-negative integers n_α and the expansion coefficients are

$$c(n) = \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!}{\prod_i (\sum_\alpha l_i^{(\alpha)} n_\alpha)!}$$

We then define

$$w_0(z, \rho) = \sum_n c(n, \rho) z^{n+\rho}$$

with

$$c(n, \rho) = \frac{\prod_j \Gamma(\sum_\alpha l_{0j}^{(\alpha)}(n_\alpha + \rho_\alpha) + 1)}{\prod_i \Gamma(\sum_\alpha l_i^{(\alpha)}(n_\alpha + \rho_\alpha) + 1)}$$

The logarithmic solutions contain the coefficients

$$\text{(I)} \quad \partial_\beta c(n) \equiv \left. \frac{\partial}{\partial_\beta} c(n + \rho) \right|_{\rho = 0}$$

$$\text{(II)} \quad \partial_\beta \partial_\gamma c(n) \qquad \text{and} \qquad \text{(III)} \quad \partial_\beta \partial_\gamma \partial_\delta c(n)$$

The following the definitions will become useful:

$$A^{(\beta)}(n) = \sum_j l_{0j}^{(\beta)} \psi(\sum_\alpha l_{0j}^{(\alpha)} n_\alpha + 1) - \sum_i l_i^{(\beta)} \psi(\sum_\alpha l_i^{(\alpha)} n_\alpha + 1)$$

$$B^{(\beta\gamma)}(n) = \sum_j l_{0j}^{(\beta)} l_{0j}^{(\gamma)} \psi'(\sum_\alpha l_{0j}^{(\alpha)} n_\alpha + 1) - \sum_i l_i^{(\beta)} l_i^{(\gamma)} \psi'(\sum_\alpha l_i^{(\alpha)} n_\alpha + 1)$$

$$C^{(\beta\gamma\delta)}(n) = \sum_j l_{0j}^{(\beta)} l_{0j}^{(\gamma)} l_{0j}^{(\delta)} \psi''(\sum_\alpha l_{0j}^{(\alpha)} n_\alpha + 1) - \sum_i l_i^{(\beta)} l_i^{(\gamma)} l_i^{(\gamma)} \psi''(\sum_\alpha l_i^{(\alpha)} n_\alpha + 1)$$

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[16] We have changed the sign of the components $l_{0j}^{(\alpha)}$ as compared to (3.3) for notational convenience.

where ψ is the logarithmic derivative of the Gamma function. We then get the expressions

$$\partial_\beta c(n) = c(n) A^{(\beta)}(n) \tag{1}$$

$$\partial_\beta \partial_\gamma c(n) = c(n) \left\{ A^{(\beta)}(n) A^{(\gamma)}(n) + B^{(\beta\gamma)}(n) \right\} \tag{2}$$

$$\partial_\beta \partial_\gamma \partial_\delta c(n) = c(n) \Big\{ A^{(\beta)}(n) A^{(\gamma)}(n) A^{(\delta)}(n) + A^{(\beta)}(n) B^{(\gamma\delta)}(n) + A^{(\gamma)}(n) B^{(\delta\beta)}(n)$$
$$ + A^{(\delta)}(n) B^{(\beta\gamma)}(n) + C^{(\beta\gamma\delta)}(n) \Big\} \tag{3}$$

whose pole structures we have to examine. Before doing this we define

$$A^{(\beta)}_{k_1,\ldots,k_p}(n) = \sum_j l^{(\beta)}_{0j} \psi\Big(\sum_\alpha l^{(\alpha)}_{0j} n_\alpha + 1\Big) - \sum_{i \neq k_1,\ldots,k_p} l^{(\beta)}_i \psi\Big(\sum_\alpha l^{(\alpha)}_i n_\alpha + 1\Big)$$

and likewise for $B^{(\beta\gamma)}_{k_1,\ldots,k_p}(n)$ and $C^{(\beta\gamma\delta)}_{k_1,\ldots,k_p}(n)$. For $\sum_\alpha l^{(\alpha)}_k n_\alpha < 0$ we set

$$\sum_\alpha l^{(\alpha)}_k n_\alpha \equiv -m_k$$

and

$$\tilde{A}^{(\beta)}_{k_1,\ldots,k_p} = A^{(\beta)}_{k_1,\ldots,k_p} - \sum_{i=1}^p l^{(\beta)}_{k_i} \psi(m_{k_i})$$

$$\tilde{B}^{(\beta\gamma)}_{k_1,\ldots,k_p} = B^{(\beta\gamma)}_{k_1,\ldots,k_p} - \sum_{i=1}^p l^{(\beta)}_{k_i} l^{(\gamma)}_{k_i} \big(\pi^2 - \psi'(m_{k_i})\big)$$

The last ingredients we need are the pole structures of Γ and ψ for $m \in \mathbf{Z}_>$:

$$\Gamma(1-m) = \frac{1}{\Gamma(m)} \frac{\pi}{\sin(\pi m)}$$

$$\psi(1-m) = \psi(m) + \pi \cot(\pi m)$$

$$\psi'(1-m) = -\psi'(m) + \pi^2(1 + \cot^2(\pi m))$$

$$\psi''(1-m) = \psi''(m) + 2\pi^3 \cot(\pi m) + 2\pi^3 \cot^3(\pi m)$$

We now consider the cases (I), (II) and (III) in turn:

(I) If we have more than one $-m_k < 0$, then $c(n)$ has a double zero. $A^{(\beta)}$ however only has a simple pole, i.e. we have to consider only the case where $-m_k < 0$ for one k only. Then

$$c(n) = \frac{\prod_j (\sum_\alpha l^{(\alpha)}_{0j} n_\alpha)!}{\prod_{i \neq k}(\sum_\alpha l^{(\alpha)}_i n_\alpha)!}(m_k - 1)! \frac{1}{\pi} \sin(\pi m_k)$$

$$A^{(\beta)}(n) = -l^{(\beta)}_k \psi(1 - m_k) + \text{finite}$$

$$ = -\pi l^{(\beta)}_k \cot(\pi m_k) + \text{finite}$$

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and we get

$$\partial_\beta c(n) = -(-)^{m_k} l_k^{(\beta)} \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!(m_k - 1)!}{\prod_{i \neq k}(\sum_\alpha l_i^{(\alpha)} n_\alpha)!} \qquad m_k \in \mathbf{Z}_>$$

(II) Since A^2 and B have double poles, we have to distinguish two possibilities:

(i) $-m_k < 0$ for one k only

(ii) $-m_k < 0$ for k_1 and k_2

For the two cases we find:

(i)

$$\partial_\beta \partial_\gamma c(n) = -(-)^{m_k} \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!(m_k - 1)!}{\prod_{i \neq k}(\sum_\alpha l_i^{(\alpha)} n_\alpha)!} (\tilde{A}_k^{(\beta)} l_k^{(\gamma)} + \tilde{A}_k^{(\gamma)} l_k^{(\beta)})$$

(ii)

$$\partial_\beta \partial_\gamma c(n) = (-)^{m_{k_1} + m_{k_2}} \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!(m_{k_1} - 1)!(m_{k_2} - 1)!}{\prod_{i \neq k_1, k_2}(\sum_\alpha l_i^{(\alpha)} n_\alpha)!} (l_{k_1}^{(\beta)} l_{k_2}^{(\gamma)} + l_{k_1}^{(\gamma)} l_{k_2}^{(\beta)})$$

(III) We now have to distinguish three cases:

(i) $-m_k < 0$ for one k only

(ii) $-m_k < 0$ for k_1 and k_2

(iii) $-m_k < 0$ for k_1, k_2 and k_3

For the three cases we find:

(i)

$$\partial_\beta \partial_\gamma \partial_\delta c(n) = -(-)^{m_k} \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!(m_k - 1)!}{\prod_{i \neq k}(\sum_\alpha l_i^{(\alpha)} n_\alpha)!}$$
$$\times \left\{ 2\pi^2 l_k^{(\beta)} l_k^{(\gamma)} l_k^{(\delta)} + [\tilde{B}_k^{(\beta\gamma)} l_k^{(\delta)} + \tilde{B}_k^{(\gamma\delta)} l_k^{(\gamma)} + \tilde{B}_k^{(\delta\beta)} l_k^{(\gamma)}] \right.$$
$$\left. + [\tilde{A}_k^{(\beta)} \tilde{A}_k^{(\gamma)} l_k^{(\delta)} + \tilde{A}_k^{(\gamma)} \tilde{A}_k^{(\delta)} l_k^{(\beta)} + \tilde{A}_k^{(\delta)} \tilde{A}_k^{(\beta)} l_k^{(\gamma)}] \right\}$$

(ii)

$$\partial_\beta \partial_\gamma \partial_\delta c(n) = (-)^{m_{k_1} + m_{k_2}} \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!(m_{k_1} - 1)!(m_{k_2} - 1)!}{\prod_{i \neq k_1, k_2}(\sum_\alpha l_i^{(\alpha)} n_\alpha)!}$$
$$\times \left\{ \tilde{A}_{k_1, k_2}^{(\beta)} (l_{k_1}^{(\gamma)} l_{k_2}^{(\delta)} + l_{k_1}^{(\delta)} l_{k_2}^{(\gamma)}) + \tilde{A}_{k_1, k_2}^{(\gamma)} (l_{k_1}^{(\delta)} l_{k_2}^{(\beta)} + l_{k_1}^{(\beta)} l_{k_2}^{(\delta)}) \right.$$
$$\left. + \tilde{A}_{k_1, k_2}^{(\delta)} (l_{k_1}^{(\beta)} l_{k_2}^{(\gamma)} + l_{k_1}^{(\gamma)} l_{k_2}^{(\beta)}) \right\}$$

(iii)

$$\partial_\beta \partial_\gamma \partial_\delta c(n) = -(-)^{m_{k_1} + m_{k_2} + m_{k_3}} \frac{\prod_j (\sum_\alpha l_{0j}^{(\alpha)} n_\alpha)!(m_{k_1} - 1)!(m_{k_2} - 1)!(m_{k_3} - 1)!}{\prod_{i \neq k_1, k_2, k_3}(\sum_\alpha l_i^{(\alpha)} n_\alpha)!}$$
$$\times \left\{ l_{k_1}^{(\beta)} l_{k_2}^{(\gamma)} l_{k_3}^{(\delta)} + l_{k_1}^{(\gamma)} l_{k_2}^{(\delta)} l_{k_3}^{(\beta)} + l_{k_1}^{(\delta)} l_{k_2}^{(\beta)} l_{k_3}^{(\gamma)} \right.$$
$$\left. + l_{k_1}^{(\gamma)} l_{k_2}^{(\beta)} l_{k_3}^{(\delta)} + l_{k_1}^{(\beta)} l_{k_2}^{(\delta)} l_{k_3}^{(\gamma)} + l_{k_1}^{(\delta)} l_{k_2}^{(\gamma)} l_{k_3}^{(\beta)} \right\}$$

Appendix B. Predicted numbers of lines for complete intersections in $\mathbb{P}^3 \times \mathbb{P}^3$ and $\mathbb{P}^3 \times \mathbb{P}^2$

For complete intersections in $\mathbb{P}^3 \times \mathbb{P}^3$,

$$\begin{pmatrix} \mathbb{P}^3 \\ \mathbb{P}^3 \end{pmatrix} \left\Vert \begin{matrix} s_1 & s_2 & s_3 \\ t_1 & t_2 & t_3 \end{matrix} \right) \, , \tag{B.1}$$

there are eight possible (non-trivial) configurations: $(s_1, s_2, s_3 \,|\, t_1, t_2, t_3) = (3, 0, 1 | 0, 3, 1)$, $(0, 2, 2 | 2, 2, 0)$, $(2, 1, 1 | 2, 1, 1)$, $(2, 1, 1 | 1, 2, 1)$, $(2, 1, 1 | 0, 3, 1)$, $(2, 1, 1 | 0, 2, 2)$, $(3, 1, 0 | 0, 2, 2)$, $(2, 2, 0 | 0, 0, 4)$. For all of these, the predicted numbers of lines with bi-degree (n_1, n_2) with respect to J_1 and J_2 from $\mathbb{P}^3$ and $\mathbb{P}^3$, respectively, are given generally by

$$N(1, 0) = 10 t_1 t_2 t_3 s_1! s_2! s_3!$$

$$- s_1! s_2! s_3! \left\{ \frac{1}{2} \left(s_2 s_3 t_1{}^3 + 3 s_1 s_3 t_1{}^2 t_2 + 3 s_1 s_2 t_1{}^2 t_3 + 3 s_1{}^2 t_1 t_2 t_3 \right) \left(\sum_{r=1}^{s_1} \frac{1}{r^2} \right) \right.$$

$$+ 4 \left(s_3 t_1{}^2 t_2 + s_2 t_1{}^2 t_3 + 2 s_1 t_1 t_2 t_3 \right) \left(\sum_{r=1}^{s_1} \frac{1}{r} \right)$$

$$- \frac{1}{2} \left(s_2 s_3 t_1{}^3 + 3 s_1 s_3 t_1{}^2 t_2 + 3 s_1 s_2 t_1{}^2 t_3 + 3 s_1{}^2 t_1 t_2 t_3 \right) \left(\sum_{r=1}^{s_1} \frac{1}{r} \right)^2$$

$$- \left(2 s_2 s_3 t_1{}^2 t_2 + 2 s_1 s_3 t_1 t_2{}^2 + s_2{}^2 t_1{}^2 t_3 + 4 s_1 s_2 t_1 t_2 t_3 + s_1{}^2 t_2{}^2 t_3 \right) \left(\sum_{r=1}^{s_1} \frac{1}{r} \right) \left(\sum_{r=1}^{s_2} \frac{1}{r} \right)$$

$$+ \left(\text{cyclic permutations} : \; ((s_1, t_1) \to (s_2, t_2) \to (s_3, t_3)) \; \right) \left. \right\} \, ,$$

$$\tag{B.2}$$

and by $N(0, 1) = N(1, 0) \big|_{(\, s_i \leftrightarrow t_i \,)}$.

The predictions for seven possible non-trivial examples in $\mathbb{P}^3 \times \mathbb{P}^2$, $(s_1, s_2 \,|\, t_1, t_2) = (2, 2 | 0, 3)$, $(3, 1 | 0, 3)$, $(2, 2 | 1, 2)$, $(4, 0 | 1, 2)$, $(1, 3 | 2, 1)$, $(3, 1 | 2, 1)$, $(0, 4 | 3, 0)$, can be extracted from the above general formulas via the manifold identity,

$$\begin{pmatrix} \mathbb{P}^3 \\ \mathbb{P}^2 \end{pmatrix} \left\Vert \begin{matrix} s_1 & s_2 \\ t_1 & t_2 \end{matrix} \right) \cong \begin{pmatrix} \mathbb{P}^3 \\ \mathbb{P}^3 \end{pmatrix} \left\Vert \begin{matrix} s_1 & s_2 & 0 \\ t_1 & t_2 & 1 \end{matrix} \right) \, , \tag{B.3}$$

Further reduction simply reproduces the result for the bi-cubic model in $\mathbb{P}^2 \times \mathbb{P}^2$ and $(4|2)$ in $\mathbb{P}^3 \times \mathbb{P}^1$ which are treated in sect.7.

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