

# **Melting layer attenuation at Ka- and W-bands as derived from multi-frequency radar Doppler spectra observations**

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**Key Points:**

- Ka-/W-band melting layer attenuation is derived
- The melting layer attenuation is derived by analyzing X-, Ka- and W-band radar Doppler spectra
- Presented results are in good agreement with previously reported melting layer modeling results for rain rates smaller than 3 mm/h

## 1 **Abstract**

2 The melting layer of precipitation has a major impact on remote sensing and telecommunications.  
3 However, there is a shortage of observational studies to validate and constrain the melting layer  
4 models especially for high-frequency radar bands. In this paper, we report how multi-frequency  
5 radar Doppler spectra can be used to retrieve the melting layer attenuation at Ka- and W-bands.  
6 The presented analysis is based on identifying Rayleigh scattering regions in radar Doppler  
7 spectra measurements where dual-wavelength spectral ratios can be related to differential  
8 attenuation. We show that the estimated attenuation at Ka- and W-bands agrees reasonably well  
9 with previously reported studies, but there are indications of differences at higher rain rates. We  
10 advocate that this technique can be applied to long-term observations to advance our knowledge  
11 of the melting process. The parameterizations of melting layer attenuation as a function of rain  
12 rate and radar reflectivity are also presented.

## 13 **Plain Language Summary**

14 While the melting layer is a relatively narrow layer in precipitation systems, it has a significant  
15 impact on telecommunication and remote sensing applications. For spaceborne and ground-based  
16 radar measurements, unaccounted attenuation in the melting layer may cause significant errors in  
17 retrievals of rainfall rate and ice cloud properties, respectively. There are two approaches to  
18 quantify attenuation in the melting layer, namely through modeling or observations. Modeling  
19 studies have been carried out and are reported in the scientific literature, however, it is  
20 acknowledged that there is a shortage of observational studies that could help us to constrain the  
21 models. In this paper, we present a study that utilizes radar Doppler spectra recorded at X-, Ka-  
22 and W-bands to retrieve melting layer attenuation. We show that the estimated attenuation at Ka-

23 and W-bands agrees reasonably well with previous studies, but there are indications of  
24 differences at higher rain rates. With the utilization of this technique, long-term observations are  
25 expected to advance our knowledge of the melting process, as well as modeling of the melting  
26 layer attenuation for applications in millimeter-wavelength radars, passive microwave remote  
27 sensing, 5G/6G commercial links and space-Earth telecommunications.

## 28 **1 Introduction**

29 The melting layer of precipitation is an area where snowflakes melt into raindrops. In this layer  
30 the hydrometeors undergo rather complex microphysical processes, while changing from  
31 irregularly-shaped ice particles into spheroid-like water drops [Matsuo & Sasyo, 1981;  
32 Yokoyama & Tanaka 1984; Fabry & Zawadzki, 1995; Leinonen & von Lerber, 2018]. In  
33 traditional weather radar observations, the melting layer can be identified by the enhanced radar  
34 reflectivity factor, the so-called “bright band”, which is caused among other things by the change  
35 in the hydrometeor dielectric properties. Propagation of radio waves through the melting layer  
36 could lead to a detectable attenuation which is comparable to the path-integrated rain attenuation  
37 [Bellon et al., 1997; Kollias & Albrecht, 2005]. Although melting layer attenuation does not pose  
38 major problems for cm-wavelength radars if the propagation path through the layer is not too  
39 long [von Lerber et al. 2015], it is considerable at mm wavelengths and affects the interpretation  
40 of ground- and space-based cloud and precipitation radar observations [Illingworth et al., 2007;  
41 Mitrescu et al., 2010]. For passive satellite microwave remote sensing, the absence of an accurate  
42 estimate of excess extinction in the melting layer may also lead to significant errors [Bauer et al.,  
43 1999; Battaglia et al., 2003]. In addition, quantifying melting layer attenuation at millimeter  
44 wavelengths becomes increasingly important as the radio waves in commercial communications

45 [Siles et al., 2015] and Earth-satellite links [Panagopoulos et al., 2004] move towards higher  
46 frequency bands.

47 There are two main approaches for estimating the radar signal attenuation in the melting layer,  
48 namely model- or observation-based. In the literature, there are a number of modeling studies  
49 that somewhat differ in the parametrizations used in the microphysical and/or electromagnetic  
50 parts of the models [e.g., Zhang et al., 1994; Russchenberg & Ligthart, 1996; D'Amico et al.,  
51 1998; Fabry & Szyrmer, 1999; Olson et al., 2001; Skaropoulos & Russchenberg, 2003; Matrosov,  
52 2008; Liao et al., 2009; Planche et al., 2014; von Lerber et al. 2015]. Most of these models are  
53 capable of reproducing observed profiles of radar reflectivity factor [Russchenberg & Ligthart,  
54 1996; D'Amico et al., 1998; Fabry & Szyrmer, 1999; Matrosov, 2008; Liao et al., 2009; von  
55 Lerber et al. 2015], dual-polarized variables [Russchenberg & Ligthart, 1996] or Doppler  
56 velocity profiles [Fabry & Szyrmer, 1999; Skaropoulos & Russchenberg, 2003]. Using such a  
57 modeling approach, attenuation caused by the melting layer can be estimated [e.g., von Lerber et  
58 al., 2015]. At mm wavelengths, melting layer models are capable of reproducing radar signatures,  
59 such as the distinct “dark band” at the upper part of the melting layer [Kollias & Albrecht, 2005].  
60 However, due to the lack of knowledge on thermodynamical [Leinonen & von Lerber, 2018] and  
61 scattering processes [Botta et al., 2010; Johnson et al., 2016] that take place during melting, the  
62 computed melting layer attenuation values are expected to be uncertain. For example, two widely  
63 used studies [Matrosov, 2008; Haynes et al., 2009] differ in the computed W-band attenuation  
64 values. Therefore, it is important to validate such estimates by measurements.

65 Previously, such measurements were carried out by using cm-wavelength radars. Bellon et al.  
66 [1997] have used coinciding X-band and UHF profiling radar observations to retrieve melting  
67 layer attenuation at X-band. Since precipitation observations at both frequency bands fall within

68 the Rayleigh regime, the observed reflectivity value difference can be linked to the attenuation.  
69 Due to the presence of non-Rayleigh scattering and potential additional attenuation (e.g., from  
70 supercooled liquid water) this approach is not directly applicable at higher frequencies.  
71 Nakamura et al. [2018] have proposed to use two Ka-band radars located on a mountain slope to  
72 estimate melting layer attenuation at this frequency. The setup is difficult to replicate and there  
73 appears to be a lack of direct measurements of melting layer attenuation at mm wavelengths,  
74 especially at W-band. This study aims to address this topic.

75 Since several years, multi-frequency Doppler spectra observations are being recorded at a  
76 number of measurement sites. This research uses multi-frequency Doppler spectra observations  
77 collected by the US DOE Atmospheric Radiation Measurement (ARM) 2<sup>nd</sup> mobile facility  
78 (AMF2) during the Biogenic Aerosols Effects on Clouds and Climate (BAECC) field campaign  
79 in Finland [Petäjä et al., 2016]. At vertical incidence, radar Doppler spectra separate radar signals  
80 from hydrometeors according to their fall velocities [e.g., Kollias et al., 2002; Moisseev et al.,  
81 2006; Spek et al., 2008]. Since smaller particles usually fall slower, there may be a part of the  
82 Doppler spectrum where the scattering from observed hydrometeors can be described using the  
83 Rayleigh approximation even at W-band [Tridon et al., 2013; Kneifel et al., 2016].

84 In this paper, we demonstrate the use of this “Rayleigh region” in multi-frequency radar Doppler  
85 spectra to derive melting layer attenuation. The proposed method can be used to correct for this  
86 attenuation in multi-frequency Doppler radar measurements. However, as discussed later in the  
87 article, matching the radar volumes is very important for the optimal performance of the method.  
88 That is why our focus is on deriving the parametrizations and validating results of previously  
89 reported melting layer attenuation studies, using carefully selected observations.

90 The topic presented in this work is of direct relevance to the cloud and precipitation remote  
91 sensing communities working with ground- and space-based observations. Furthermore, the  
92 study supports current and future cloud and precipitation satellite missions, namely the National  
93 Aeronautics and Space Administration (NASA) CloudSat, Global Precipitation Measurement  
94 mission (GPM), and the European Space Agency/Japan Aerospace Exploration Agency  
95 (ESA/JAXA) Earth Clouds, Aerosols and Radiation Explorer (EarthCARE).

96 This paper is organized as follows. Section 2 presents the measurement setup and associated data  
97 processing. Then, section 3 describes the procedure of this new technique and discusses the  
98 retrieval uncertainties. A “sanity” check, a case analysis as well as the retrieved melting layer  
99 attenuation are presented in section 4 and conclusions are summarized in section 5.

## 100 **2 Measurements**

101 The rain events analyzed in this study were collected during the BAecc field campaign, which  
102 was carried out at the University of Helsinki Hyytiälä Station from February to September 2014  
103 [Petäjä et al., 2016]. The radar data were recorded by the X/Ka-band scanning Atmospheric  
104 Radiation Measurement (ARM) cloud radar (X/Ka-SACR) as well as by the zenith-pointing Ka-  
105 band (KAZR) and W-band ARM cloud radars (MWACR) [Kollias et al., 2014; Kneifel et al.,  
106 2015; Falconi et al., 2018]. X- and Ka-SACR are dual-polarization radar systems that record the  
107 copolar correlation coefficient and the linear depolarization ratio, respectively. The range gate  
108 spacing for X/Ka-SACR and KAZR is 25 m, 5 m less than that of MWACR. All the radars were  
109 recording Doppler spectra every 2s using 512-sample Fast Fourier Transform (FFT) that resulted  
110 in the velocity resolution of about 2.3 cm/s for all the radars. The original spectra were averaged  
111 and the resulting temporal resolution is 14 s. The half-power beam widths of the X-SACR, Ka-  
112 SACR, KAZR and MWACR radars are  $1.27^\circ$ ,  $0.33^\circ$ ,  $0.33^\circ$  and  $0.38^\circ$  respectively.

113 The data used in this study were collected when all the radars were pointing vertically. To  
114 maximize the matching between measurements at various frequencies, X/Ka-SACR  
115 measurements were used for analysis of X/Ka-band spectra, while KAZR and MWACR data  
116 were used for the Ka/W-band spectra. The retrieval of W-band attenuation was done based on  
117 KAZR/MWACR combination, in which Ka-band attenuation was derived from Ka-SACR and  
118 X-SACR observations. This is done because the horizontal distance between X-SACR and  
119 MWACR was about 17 m [Petäjä et al., 2016], while the distance between KAZR and MWACR  
120 was just a couple of meters.

121 It should be noted that MWACR had a small antenna pointing error of about 0.5 to 1°, caused by  
122 freezing and thawing of the gravel under the radar container, which has led to some offset in  
123 measured vertical Doppler velocity [Kneifel et al., 2016]. Furthermore, due to the difference in  
124 the radar volumes, it was observed that in some cases X- and Ka- band spectra were not well  
125 matched. To minimize the associated errors, the MWACR (X-SACR) spectra were shifted until  
126 the best match with the KAZR (Ka-SACR) measured spectra was reached [Kneifel et al., 2016].  
127 It was found, however, that in the ice region the matching of single-peak spectra was very  
128 difficult to achieve. Given that there is no reliable indicator (e.g., presence of detectable small  
129 peaks or valleys) for identifying a good spectra match, the resulting uncertainty in dual-  
130 wavelength spectral ratios can be as large as the expected attenuation contribution. For multi-  
131 peak spectra this uncertainty can be mitigated. Because of this, in this study only cases with  
132 multi-modal spectra in the snow region were used. Additionally, the size of the Rayleigh region  
133 in W-band Doppler spectra can be rather small and is easily affected by turbulent broadening  
134 [Tridon & Battaglia, 2015]. The broadening of spectra due to turbulence and cross-wind [Doviak  
135 & Zrníc, 1984] results in power leakage from non-Rayleigh scattering spectra regions into the



136 Rayleigh region and therefore biases the needed DSR measurements. That is why Ka/W-band  
137 spectra analyses were only performed on cases with a well-defined supercooled water peak  
138 [Luke et al., 2010].

139 Prior to analyzing the radar data, gaseous attenuation was corrected by applying the millimeter-  
140 wave propagation model [Liebe, 1985] which was computed using the closest radio sounding  
141 data, which were carried out four times per day [Petäjä et al., 2016].

142 To verify X-SACR calibration and to compute rain specific attenuation at all the frequencies,  
143 radar variables were computed from 2D video disdrometer [Schuur et al., 2001] observations by  
144 using Python implementation of T-matrix scattering calculations (PyTmatrix) [Mishchenko &  
145 Travis, 1994; Wielaard et al., 1997; Leinonen, 2014]. The melting layer attenuation method is  
146 based on “dual-wavelength spectral ratio” measurements, as discussed in the next section, and  
147 therefore is independent of potential reflectivity calibration offsets. That is why no calibration is  
148 needed for Ka-SACR, KAZR or MWACR. The X-SACR reflectivity values, however, are used  
149 to derive parametrizations of melting layer attenuation as a function of radar reflectivity and rain  
150 rate, which are affected by the calibration bias as well as the radome attenuation. To mitigate this,  
151 the X-SACR calibration was done by matching the observed reflectivity at about 500 m to the  
152 reflectivity values computed from the 2D video disdrometer data. The time lag between the  
153 observed and simulated radar reflectivity values was found to be around 1 min. Note that we did  
154 not use the lowest range gate of X-SACR due to the near-field effects [Sekelsky, 2002; Falconi  
155 et al., 2018].

## 156 **3 Methods**

### 157 **3.1 General description**

158 Generally, radar signal attenuation through clouds and rain is greater at higher frequencies,  
159 resulting in so-called differential attenuation between radio waves of collocated radars operating  
160 at distinct frequencies [Tuttle & Rinehart, 1983]. If the radars observe weather echoes where  
161 hydrometeors are small enough to fall within the Rayleigh regime, the ratio of the observed  
162 reflectivity factors, the dual-wavelength ratio (DWR), is mainly caused by the differential  
163 attenuation. This measurement has been used to retrieve the liquid water content from dual-  
164 wavelength radar observations, since most liquid droplets are good Rayleigh scatterers at Ka-  
165 band [Ellis & Vivekanandan, 2011] as well as at W-band [Hogan et al., 2005]. By applying the  
166 DWR method to vertically pointing X-band and UHF profiling radars measurements, Bellon et al.  
167 [1997] have estimated the melting layer attenuation at X-band, as the sizes of most hydrometeors  
168 below and above the melting layer are much smaller than the X-band wavelength.

169 However, the DWR method cannot be directly applied to retrieve the attenuation when the  
170 hydrometeors are not Rayleigh scatterers [Gaussiat et al., 2003]. Tridon et al. [2013] have  
171 proposed to use the slower falling part of the vertically pointing Doppler spectra, where the  
172 Rayleigh assumption is valid, to study the attenuation in rainfall.

173 Similar to rainfall, in the ice part of the precipitation the slower falling particles should be small  
174 enough to be Rayleigh scatterers. Spek et al. [2008] have reported that spectral dual-polarization  
175 radar signatures of these particles are similar to those of ice crystals. Kneifel et al. [2016] have  
176 shown that spectral multi-frequency observations in this Doppler spectrum part behave in  
177 accordance with the Rayleigh scattering approximation.

178 Figure 1 demonstrates multi-frequency Doppler spectra properties above and below the melting  
 179 layer as was discussed above. As shown in Figure 1 (a), the melting layer can be recognized by  
 180 the enhanced radar reflectivity. The melting layer boundaries, shown in the figure and used in  
 181 this study, were estimated using X-SACR measurements of the copolar correlation coefficient  
 182 [Ryzhkov & Zrnich, 1998; Matrosov et al., 2007]. The derivation of the differential attenuation in  
 183 Figure 1 (a) is highly dependent on the identification of the Rayleigh region, as illustrated in  
 184 panels (b) to (e), and will be discussed in the next section.

185 Assuming that multi-frequency Doppler radar spectra are well-matched, the Rayleigh parts of the  
 186 spectra can be recognized by the plateau in the observed spectral dual-wavelength spectral ratios  
 187 (DSR). The observed DSR values in the Rayleigh part of the Doppler spectra above the melting  
 188 layer will be caused by differential attenuation due to combined wet radome, rain, and melting  
 189 layer, as well as possible calibration offsets. The same factors, except for the melting layer  
 190 attenuation, contribute to DSR values in the Rayleigh region in rain. Therefore, by considering  
 191 the differences of the dielectric factors for hydrometeors below and above melting layer, the  
 192 differential melting layer attenuation can be estimated as:

$$A_{ML}(\lambda_1, \lambda_2) = \overline{DSR_{\text{Rayleigh, above ML}}(\lambda_1, \lambda_2)} - \overline{DSR_{\text{Rayleigh, below ML}}(\lambda_1, \lambda_2)} - 10 \log_{10} \left( \frac{|K_{\text{above ML}}(\lambda_1)|^2 |K_{\text{below ML}}(\lambda_2)|^2}{|K_{\text{below ML}}(\lambda_1)|^2 |K_{\text{above ML}}(\lambda_2)|^2} \right), \text{ [dB]} \quad (1)$$

193 where  $\lambda_1$  and  $\lambda_2$  are 30.9 mm and 8.5 mm (X- and Ka-SACR ) or 8.6 mm and 3.19 mm (KAZR  
 194 and MWACR) respectively.  $|K_{\text{above ML}}(\lambda)|^2$  is the dielectric constant of ice (X/Ka) or supercooled  
 195 water droplets (Ka/W) at  $\lambda$  and  $|K_{\text{below ML}}(\lambda)|^2$  is the dielectric constant of rain water at  $\lambda$  .

196  $\overline{DSR_{\text{Rayleigh, above ML}}(\lambda_1, \lambda_2)}$  stands for the average DSR in the Rayleigh part of the spectrum just

197 above the melting layer (ice for X/Ka, liquid water for Ka/W),  $\overline{DSR}_{\text{Rayleigh, below ML}}(\lambda_1, \lambda_2)$  is the  
 198 same but for below the melting layer (rain). As can be seen in Figure 1 (b, c), the supercooled  
 199 water DSR values are higher than those of ice. This difference can be explained by the fact that  
 200 the dielectric constants of ice at W-, Ka- and X-bands are almost the same, while the dielectric  
 201 constants of liquid water at these radar frequencies are different. Here  $\overline{DSR}_{\text{Rayleigh}}(\lambda_1, \lambda_2)$  is  
 202 defined as

$$\overline{DSR}_{\text{Rayleigh}}(\lambda_1, \lambda_2) = 10 \log_{10} \left( \frac{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_{\lambda_1}(v_i)}{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_{\lambda_2}(v_i)} \right), [\text{dB}] \quad (2)$$

203 where  $v_{\text{start}}$  and  $v_{\text{end}}$  are velocity boundaries of Rayleigh region and will be discussed later;  
 204  $S_{\lambda_1}(v)$  and  $S_{\lambda_2}(v)$  are spectral powers in linear scale at  $v$  observed by radars working at  $\lambda_1$  and  
 205  $\lambda_2$  respectively. As can be seen in Figure 1(a), the differential attenuation for Ka/W is about 3  
 206 times of that for X/Ka in rain. Their difference at 500 m is mainly due to the radome attenuation,  
 207 rain attenuation and other calibration issues. There is no significant increase of the differential  
 208 attenuation for either X/Ka or Ka/W in snow. The differential attenuation in the melting layer is  
 209 larger for Ka/W than for X/Ka, confirming that the W-band signal is attenuated more than the  
 210 Ka-band upon propagation through the melting layer.

211 The above-described procedure provides melting layer differential attenuation values for X/Ka  
 212 and Ka/W band radar observations. To retrieve absolute attenuation values, we use modeling  
 213 results of Matrosov [2008] that links melting layer attenuation to rain intensity for X-band radar  
 214 observations.

### 215 **3.2 Identification of Rayleigh scatterers in Doppler spectra**

216 The most important and somewhat challenging part of the presented method is the identification  
217 of the Rayleigh scatterers in the observed radar Doppler spectra. In this study we have used three  
218 distinct approaches for this, depending on the considered frequency bands and whether we are  
219 dealing with measurements in rain or snow.

220 For all frequency bands in rain we follow the approach of Tridon et al. [2013], who have  
221 demonstrated that the DSR in rain exhibit a well-defined plateau in the low-velocity region,  
222 which is defined by small raindrops that satisfy the Rayleigh scattering assumption. Analysis of  
223 the cases used in this study shows that the DSR plateau extends up to velocities of 2.5 ~ 3.5 m/s  
224 and 4 ~ 5 m/s for Ka/W and X/Ka spectral ratios, respectively, see Figure 1 (d, e) for an example.  
225 Tridon & Battaglia [2015] have shown that the spectral broadening, due to air motion, could  
226 affect the Rayleigh plateau and bias the DSR measurements. To mitigate this, the optimal  
227 estimation method has been utilized [Tridon & Battaglia, 2015; Tridon et al., 2017a; Tridon et al.,  
228 2017b]. After inspecting the cases analyzed in this paper (see the list in the supporting  
229 information), we found that the DSR in the Rayleigh region is rather flat and does not seem to be  
230 affected by the spectral broadening, which would skew the plateau [Tridon et al., 2013]. This  
231 applies to X/Ka and Ka/W DSR observations. Furthermore, in the majority of cases the Ka/W  
232 DSR exhibit distinct resonance scattering peaks in the larger raindrop region, see Figure 1 (e),  
233 which is also an indication of low turbulence conditions [Tridon et al., 2013]. Therefore, it is  
234 expected that the derived combined attenuation due to wet radome and rain, as well as possible  
235 calibration offsets can be estimated from the DSR Rayleigh plateaux in rain. To minimize the  
236 impact of noise, the summation limits in (2) were determined as follows:  $v_{\text{start}}$  is defined as  
237 Doppler velocity where the signal-to-noise ratio (SNR) reaches 10 dB and the width of the

238 Rayleigh plateau,  $v_{\text{end}} - v_{\text{start}}$ , is estimated to be 2 m/s and 1.5 m/s for X/Ka and Ka/W,  
239 respectively.

240 The identification of Rayleigh scatterers above the melting layer is more involved and depends  
241 on the frequency band. Because ice particle velocities do not change as much with size as  
242 velocities of raindrops, the Rayleigh plateau region is smaller in snow than in rain, as can be seen  
243 in Figure 1 (b, c). Analysis of our data shows that X/Ka DSR Rayleigh plateau extends over at  
244 least 0.5 m/s. Therefore, for X/Ka in snow, an SNR of 10 dB was adopted to determine  $v_{\text{start}}$  and  
245 the value of  $v_{\text{end}}$  depends on the extent of Rayleigh plateau. To avoid the potential impact of  
246 turbulence, cases with slanted Rayleigh region, namely not a flat plateau, due to significant  
247 spectral broadening were removed.

248 Locating the Rayleigh region in snow for W-band seems to be more difficult due to its relatively  
249 narrow span. However, liquid water droplets in cloud are small enough to satisfy Rayleigh  
250 scattering assumptions at Ka- and W- bands [e.g., Gaussiat et al., 2003; Hogan et al., 2005]. In  
251 radar Doppler spectra, the supercooled liquid water can be identified as a narrow well-defined  
252 peak close to zero Doppler velocity [Zawadzki et al., 2001; Luke et al., 2010; Shupe et al., 2004].  
253 That is why we have used supercooled water peak, visible and distinct in Figure 1 (c) just above  
254 the melting layer, for computing DSR and estimating the differential attenuation between Ka-  
255 band and W-band. Another benefit of adopting the supercooled water is that the total spectral  
256 power of this peak is almost not affected by turbulence. Observations show that such  
257 supercooled water peaks barely exceed the noise level by a few dBs, thus lower SNR (3 dB) was  
258 adopted to determine  $v_{\text{start}}$  and  $v_{\text{end}}$ . By contrast, the Rayleigh plateau in ice is used for X/Ka-band  
259 observations since the supercooled water peaks detected by Ka- and W-band radars are not  
260 always detectable by X-SACR.

261 **3.3 Uncertainty analysis**

262 There are two uncertainty types relevant to this study, namely statistical and methodological. The  
 263 estimation of the differential attenuation includes computations of DSR below and above melting  
 264 layer. Ignoring the much smaller contribution from the dielectric factor term in (1), the statistical  
 265 uncertainty of Ka-band melting layer attenuation is defined as

$$\Delta A_{\text{ML,Ka}} = \sqrt{\overline{\Delta DSR}_{\text{Rayleigh, above ML}}(X, Ka)^2 + \overline{\Delta DSR}_{\text{Rayleigh, below ML}}(X, Ka)^2} \quad (3)$$

266 Following the method utilized by Hogan et al. [2005],  $\overline{\Delta DSR}_{\text{Rayleigh}}(X, Ka)$  for high SNR may be  
 267 expressed as

$$\begin{aligned} \overline{\Delta DSR}_{\text{Rayleigh}}(X, Ka) &= \sqrt{\left( \frac{4.343 \sqrt{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} \left( \frac{S_X(v_i)}{\sqrt{M}} \right)^2}}{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_X(v_i)} \right)^2 + \left( \frac{4.343 \sqrt{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} \left( \frac{S_{\text{Ka}}(v_i)}{\sqrt{M}} \right)^2}}{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_{\text{Ka}}(v_i)} \right)^2}, \text{ [dB]} \quad (4) \\ &= \frac{4.343}{\sqrt{M}} \sqrt{\frac{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_X(v_i)^2}{\left( \sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_X(v_i) \right)^2} + \frac{\sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_{\text{Ka}}(v_i)^2}{\left( \sum_{v_i=v_{\text{start}}}^{v_{\text{end}}} S_{\text{Ka}}(v_i) \right)^2}} \end{aligned}$$

268 where  $M$  is the number of independent samples. The numbers of sample averaged in raw data  
 269 products are 3, 11 and 5 for X-SACR, Ka-SACR/KAZR and MWACR respectively. In this study,  
 270 we have used averaging over seven continuous Doppler spectra products, implying  $M = 21$  (X-  
 271 SACR), 77 (Ka-SACR/KAZR) and 35(MWACR).

272 For W-band attenuation the uncertainty can be expressed as

$$\Delta A_{ML,W} = \sqrt{\Delta A_{ML,Ka}^2 + \overline{\Delta DSR}_{\text{Rayleigh, above ML}}(Ka,W)^2 + \overline{\Delta DSR}_{\text{Rayleigh, below ML}}(Ka,W)^2} \quad (5)$$

273 We found that  $\Delta A_{ML,Ka}$  and  $\Delta A_{ML,W}$  of our estimates are 0.4 ~ 0.6 dB and 0.55 ~ 0.75 dB  
 274 respectively, mainly depending on the number of spectra averaged.

275 The methodological uncertainties include potential errors caused by imperfectly-matched  
 276 Doppler spectra at various wavelengths and computations of the absolute values of melting layer  
 277 attenuation from the differential attenuation. To derive the absolute values of the melting layer  
 278 attenuation, the relations linking X-band melting layer attenuation to rain rate proposed by  
 279 Matrosov [2008] are used. Von Lerber et al. [2015] have shown that the melting layer  
 280 attenuation depends on ice particle properties, i.e. whether they are rimed or unrimed. In the  
 281 cases we have collected, the maximum reflectivity (from X-SACR) at melting layer peak is  
 282 about 41 dB, which may correspond to around 1 dB X-band attenuation according to [Bellon et  
 283 al., 1997]. The corresponding value estimated by Matrosov [2008] is about 0.26 dB. Therefore,  
 284 Matrosov's relations could underestimate or overestimate the attenuation. In this study, most  
 285 melting layer cases are associated with rimed snowflakes as revealed by the supercooled water  
 286 peaks, thus we expect that the X-band attenuation should be on the smaller side according to  
 287 [von Lerber et al. 2015]. Assuming that the actual X-band attenuation is 0.2 to 5 times that of  
 288 Matrosov's relation, the expected error of the melting layer attenuation at X-band ranges from -  
 289 1.04 (underestimation) to 0.21 (overestimation) dB. This translates to 20 ~ 30% error at Ka-band  
 290 and < 10% at W-band in dB scale.

291 The uncertainties associated with Doppler spectra matching are difficult to quantify. In rain, the  
 292 expected large Rayleigh plateau facilitates the matching. In snow, we have used multi-mode  
 293 spectra for matching.



## 294 **4 Results**

295 From all the rain cases recorded during BAECC, seven stratiform rain events were selected to  
296 estimate the Ka-band melting layer attenuation (~ 48 min in total). These cases have satisfied our  
297 criteria, namely presence of multimodal spectra in ice region. From these cases, only three events  
298 showed detectable supercooled liquid water peaks (~ 24 min in total) and were used to compute  
299 W-band attenuation (see supporting material for the list).

### 300 **4.1 “Sanity” check, comparison against cloud top DWR**

301 To verify whether our retrieval is meaningful, a comparison was made to the DWR at the cloud  
302 top ( $DWR_{CT}$ ) where the Rayleigh assumption is valid. Therefore, the retrievals based on the  
303  $DWR_{CT}$  and the DSR at melting layer top ( $DSR_{ML\_top}$ ) would be the same if there were no  
304 attenuation between the melting layer top and cloud top. However, the fact is that such  
305 attenuation due to snow and supercooled water above the melting layer could be significant,  
306 especially at W-band [Kneifel et al., 2015]. Therefore, we expect that the results of the presented  
307 technique ( $DSR_{ML\_top}$ ) should always be smaller than the estimates using the  $DWR_{CT}$ -based  
308 method.

309 To derive  $DWR_{CT}$ , the area of cloud top where Rayleigh’s assumption applies needs to be  
310 identified. Hogan et al. [2000] have observed that the radar reflectivity of Rayleigh scatterers in  
311 ice clouds at W-band is below -20 dBZ, but it seems that the threshold of -15 dBZ or higher can  
312 also be used [Stein et al., 2015]. After analysing the  $DWR_{CT}$  profiles, we found that the snow  
313 with reflectivity between -20 and -10 dBZ observed by X-SACR may also be treated as Rayleigh  
314 scatterers at W-band. The radar reflectivity where the Rayleigh assumption is satisfied seems to  
315 be dependent on snow type, which, however, is not the topic of this paper. In addition, the  
316  $DWR_{CT}$  values are not stable when SNR is low, since the radar signal at the highest detected

317 cloud top is affected by noise. In this comparison, the cloud top where X-band radar reflectivity  
318 ranges from -20 dBZ to -10 dBZ and all radars' SNR > 10 dB was used to calculate  $DWR_{CT}$ .

319 In Figure 2 the results of this comparison are shown. Note that the radome and rain attenuation  
320 was removed from  $DWR_{CT}$  and  $DSR_{ML\_top}$  values by subtracting DSR in the Rayleigh plateau  
321 derived just below the melting layer. As expected, the  $DWR_{CT}$ -based retrievals are larger than  
322  $DSR_{ML\_top}$ -based for X/Ka as well as for Ka/W, indicating the non-negligible attenuation above  
323 the melting layer top. Furthermore, the difference between the two approaches is larger for Ka/W  
324 than for X/Ka -bands. More specifically, the difference between the results from  $DWR_{CT}$  and  
325  $DSR_{ML\_top}$  can be as large as 1.5 dB for X/Ka and 4 dB for Ka/W in 11<sup>th</sup> May 2014, an event  
326 with significant supercooled water signatures in spectra observations. While this figure does not  
327 prove that the derived melting layer attenuation values are correct, it gives a level of confidence  
328 that our retrievals are reasonable.

#### 329 **4.2 Case study, May 11<sup>th</sup> 2014**

330 On May 11<sup>th</sup> 2014, a stratiform precipitation system passed over the Hyttiälä station. During  
331 BAECC the X-SACR and Ka-SACR were regularly performing range height indicator (RHI)  
332 scans and the Doppler spectra data are not continuous. Only a short period of about 22 min when  
333 all radars were pointing vertically is available for this event.

334 In Figure 3 (a), the X-SACR height time observations of reflectivity factor are shown. As can be  
335 seen the bright band is clearly visible and its height is around 1.2 km. There are three short  
336 periods of rain showers around 03:05, 03:15 and 03:20 UTC. The precipitation cloud top height  
337 is about 3 km. There is also a second cloud layer, which starts just below 4 km. The upper cloud  
338 is an ice cloud and particles in its upper part can be regarded as Rayleigh scatterers.

339 To compute the differential attenuation, both DSR and DWR observations were corrected for  
340 rain, radome attenuations and calibration offsets by subtracting estimated just below the melting  
341 layer. That is why  $DWR(X, Ka)$  values become negative in rain. This effect for  $DWR(Ka, W)$  is  
342 obscured by the non-Rayleigh scattering. The influence of differential attenuation on  
343 observations is clearly visible in the DWR measurements of the upper cloud, Figure 3 (b, d).

344 The comparison of the differential attenuation estimates using  $DSR_{ML\_top}$  and  $DWR_{CT}$  are shown  
345 in Figure 3 (c, e). As expected  $DWR_{CT}$ -based estimates are higher, because they also include  
346 attenuation due to snow and possibly supercooled liquid water. The melting layer differential  
347 attenuation between X-band and Ka-band (Figure 3c) shows dependence on precipitation  
348 intensity, namely the higher the precipitation intensity, the larger the melting layer attenuation.  
349 Such dependence can also be found in Figure 3 (e), however, with smaller relative variations.

350 The biggest difference between  $DWR_{CT}$  and  $DSR_{ML\_top}$  observations is found around 03:10 UTC  
351 and can be seen in Figure 3 (e). This larger difference between the  $DWR_{CT}$  and  $DSR_{ML\_top}$  values  
352 is most probably due to the attenuation above the melting layer, which is possibly caused by  
353 supercooled water that attenuates more at W-band.

354 Overall the comparison of  $DWR_{CT}$  and  $DSR_{ML\_top}$  based differential attenuation retrievals show  
355 that  $DWR_{CT}$  tends to overestimate the melting layer attenuation. Even though the attenuation due  
356 to snow could potentially be accounted for (e.g. Leinonen et al., [2011]), the attenuation from  
357 supercooled liquid water cannot be currently removed.

#### 358 **4.3 Parametrizations of melting layer attenuation**

359 It is customary to express the melting layer attenuation either as a function of precipitation rate  
360 or unattenuated reflectivity below the melting layer. For this we have used disdrometer corrected  
361 X-band reflectivity values and rain rates derived from these values. We have decided not to use

362 the surface rain rate measured by one of our ground-based precipitation sensors, because of the  
363 potential rain evaporation [Tridon et al., 2017b]. The disdrometer-corrected X-band reflectivity  
364 beneath the melting layer was converted to rain rates by the Z-R relations, which were derived  
365 from the hourly disdrometer data and not from the whole events (see the comparison in  
366 supporting material). This approach allows for adaptive Z-R conversions. It should be noted, that  
367 in order to compute the W-band attenuation, not only the differential attenuation between Ka-  
368 band and W-band but also the X-/Ka-band differential attenuation should be derived. Since the  
369 selection criteria for Ka/W and X/Ka spectra analysis cases are not the same, not in all cases  
370 Ka/W and X/Ka differential attenuation values can be estimated simultaneously. In such  
371 situations, the closest retrieved Ka-band attenuation value, within 1-min of Ka/W spectra  
372 observations, was used as the input in calculating the W-band attenuation.

373 In addition to total attenuation, specific melting layer attenuation was also computed. The  
374 specific attenuation is defined as the total attenuation divided by the melting layer thickness.  
375 Additionally, the rain specific attenuation values, for the studied cases, were computed from the  
376 disdrometer data by using PyTmatrix.

377 In Figure 4, the derived parametrizations are presented and compared to the attenuation-rain rate  
378 (A-R) expressions of Matrosov, [2008]. In Figures 4 (a<sub>1</sub>, b<sub>1</sub>), the rain rate in Matrosov [2008]  
379 relations was converted to the radar reflectivity factor by applying the Marshall and Palmer  
380 relation, i.e.  $Z = 200R^{1.6}$  [Marshall and Palmer 1948; Marshall et al., 1955]. As can be seen, the  
381 melting layer attenuation for Ka- and W-bands depends on the reflectivity factor, rain rate, as  
382 well as the melting layer thickness. The melting layer thickness increases as the reflectivity  
383 factor rises, which is in agreement with previous studies [e.g., Fabry & Zawadzki, 1995;  
384 Wolfensberger et al., 2016]. Similar to Bellon et al. [1997], we found that the melting layer peak

385 reflectivity at X-band seems to be better correlated to melting layer attenuation than the rain rate  
386 (see supporting material).

387 As shown in Figure 4 (a, b), the observed results both for Ka- and W- bands agree reasonably  
388 well with Matrosov's relations, at least in cases where the rain rate is less than 3 mm/h ( $Z < 30$   
389 dBZ). Such good agreement is noteworthy, given the recent discussion of the applicability of the  
390 Effective-Medium Approximation (EMA) for modeling scattering properties of snowflakes and  
391 melting ice particles [Johnson et al., 2016]. It could possibly be explained by the fact that the  
392 forward scattering computations do not necessarily need complex particle models, as shown by  
393 scattering computations of snowflakes [e.g., Tyynelä et al., 2013, Hogan and Westbrook, 2014;  
394 Hogan et al., 2017; Leinonen et al., 2018]. However, these studies focused on dry snowflakes  
395 and no similar computations for melting particles were done.

396 For higher rain rates ( $Z > 30$  dBZ), the observations show smaller attenuation values than the  
397 model despite the fact that a very limited number of cases were analyzed. It is interesting to note  
398 that this difference between the model and observations is similar for Ka- and W-bands. This  
399 attenuation overestimation in modeling results might be due to the assumptions made about snow  
400 physical properties. The model assumes that the snowflakes are unrimed/lightly rimed [Matrosov,  
401 2008], whilst most of our observation cases, especially when  $Z > 30$  dBZ, exhibit well-defined  
402 supercooled water signatures which could lead to higher degrees of riming. Von Lerber et al.  
403 [2015] have shown that under similar conditions, i.e. rain intensity, the rimed snowflakes would  
404 lead to lower melting layer attenuation. It should also be noted that, in addition to density,  
405 snowflake shape is also affected by riming [Li et al., 2018], which may have a significant impact  
406 on scattering properties of ice particles especially at mm wavelengths.

407 The fitted relations for melting layer attenuation (dB) at Ka-band are

$$A_{Ka} = 0.97R^{0.61} \text{ or } A_{Ka} = 0.13Z_{lin}^{0.38} \text{ (} Z_{lin} \text{ is reflectivity in mm}^6\text{/m}^3\text{)}, \quad (6)$$

408 and at W-band,

$$A_W = 2.9R^{0.42} \text{ or } A_W = 0.67Z_{lin}^{0.27} . \quad (7)$$

409 As shown in Figure 4 (c) and (d), the melting layer specific attenuation at Ka- and W-bands  
 410 increases with reflectivity factor and rain rate as well as melting layer width. Compared to rain,  
 411 the specific attenuation in the melting layer is significantly larger. The derived relations for  
 412 melting layer and rain one-way specific attenuations at Ka-band are as follows,

$$k_{ML\_Ka} = 1.2R^{0.42} \text{ or } k_{ML\_Ka} = 0.29Z_{lin}^{0.27} , \quad (8)$$

$$k_{rain\_Ka} = 0.18R^{1.24} \text{ or } k_{rain\_Ka} = 0.0074Z_{lin}^{0.67} , \quad (9)$$

413 and at W-band,

$$k_{ML\_W} = 3.4R^{0.3} \text{ or } k_{ML\_W} = 1.2Z_{lin}^{0.2} , \quad (10)$$

$$k_{rain\_W} = 1.12R^{0.93} \text{ or } k_{rain\_W} = 0.14Z_{lin}^{0.47} . \quad (11)$$

414 As shown in Figure 4 (c<sub>2</sub>, d<sub>2</sub>), the fit for rain agrees rather well with that derived from field  
 415 campaign observations [Matrosov, 2005; Matrosov, 2007].

416 Using these fits, we calculated the ratio of the specific attenuation of melting layer and rain, and  
 417 found that it decreases from 4.3 to 1.5 and from 2.4 to 1.2 for Ka- and W-bands, respectively, as  
 418 the reflectivity rises from 23 to 36 dBZ. A similar decreasing behavior of specific attenuation  
 419 ratio was observed for X-band by Bellon et al. [1997] (see the comparison in supporting  
 420 material).

## 421 **5 Conclusions**

422 In this paper, we present a technique for deriving the melting layer attenuation at Ka- and W-  
423 bands using X-, Ka- and W-band vertically pointing radar Doppler spectra. The retrieval is based  
424 on the measurements of differential attenuation, which avoids potential radar calibration and wet  
425 radome issues. Using data collected during US DOE ARM funded BAECC experiment, melting  
426 layer attenuation was derived and parametrized. The presented “sanity” check and case study  
427 show the necessity of utilizing the  $DSR_{ML\_top}$  instead of simply using the reflectivity difference at  
428 the cloud top. Despite the limited observations (~ 48 min and ~24 min for Ka- and W-bands,  
429 respectively), the results show similar trends as previously found at X-band radar observations  
430 and agree reasonably well with the previously reported modeling studies. This agreement  
431 indicates a possibility of adopting simpler melting snowflake models for forward scattering  
432 computations. It is, however, found that the modeling-based attenuation parametrization seems to  
433 overestimate the attenuation for moderate to heavy rainfall. This difference could possibly be  
434 explained by the assumed snow microphysical properties that are used as input into the melting  
435 layer model. This stresses the need for more comprehensive modeling- and observation-based  
436 studies of melting layer properties.

437 Despite the focus of this study on deriving melting layer attenuation parametrizations, the  
438 proposed technique can be applied to mitigate the melting layer attenuation in ground-based  
439 radar measurements. This, however, requires multi-frequency radar observations with well-  
440 matched radar volumes. The application of this method to space- and aircraft- based observations  
441 is further complicated by the requirements of high-fidelity Doppler spectra measurements.

442 Mainly due to the expensive data storage in the past, spectra observations are still under-  
443 exploited. Long-term records of multi-frequency radar Doppler spectra are therefore expected to  
444 advance our knowledge on the melting layer, and to develop more realistic models for  
445 applications in commercial/satellite communications, ground/space borne radar remote sensing,  
446 quantitative precipitation estimation as well as radar data assimilation.

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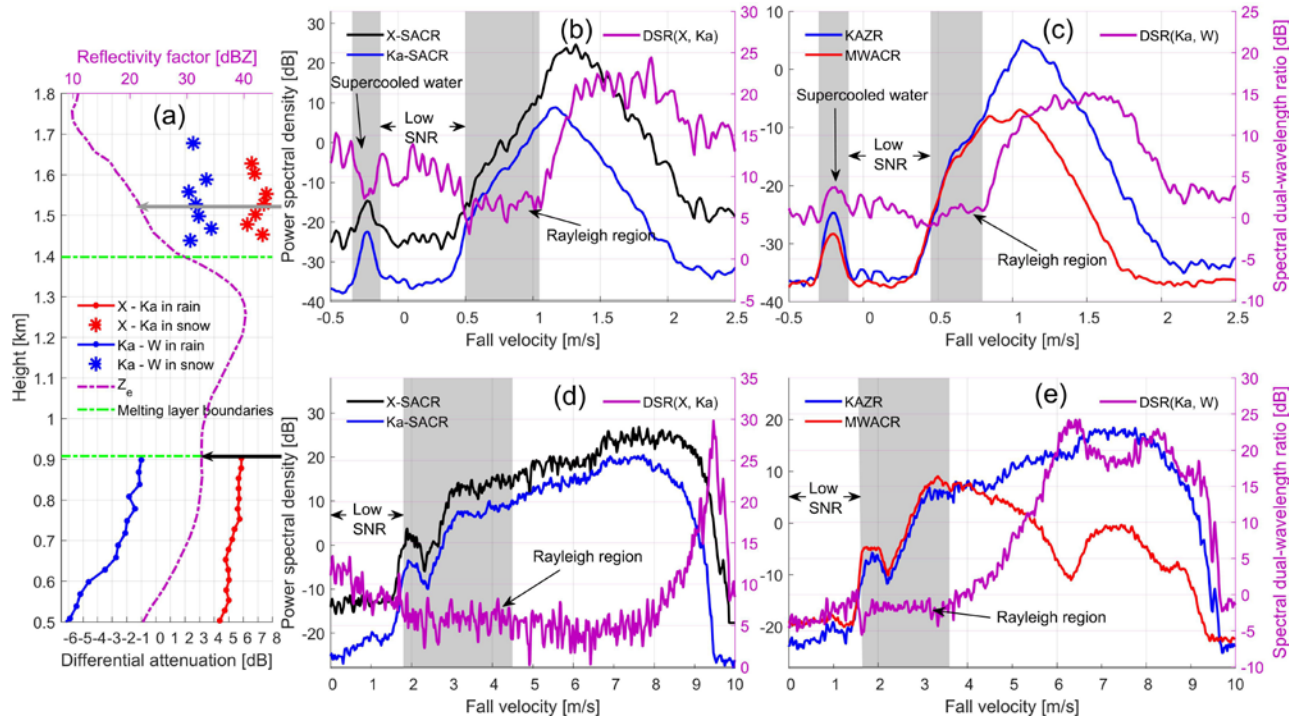


Figure 1. Differential attenuation and X-band observed reflectivity factor (a). Spectra observations in snow and rain areas indicated by the grey and black arrows are shown in panels (b, c) and (d, e) respectively. The grey shading identifies regions of spectra where hydrometeors are Rayleigh scatterers.

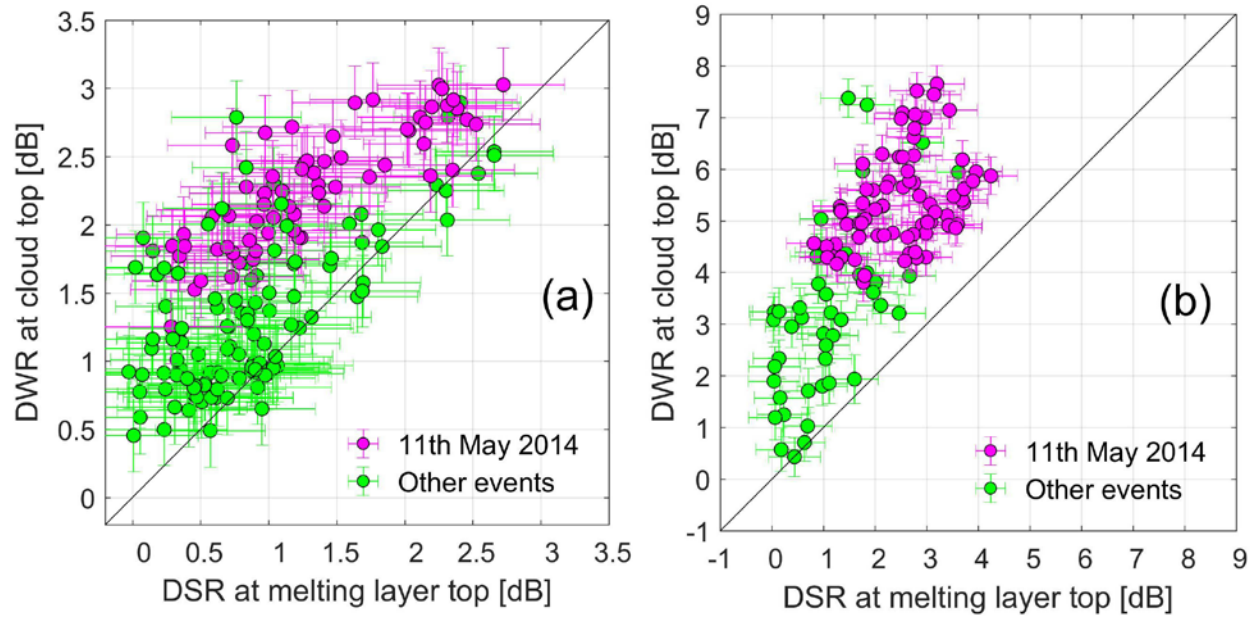


Figure 2. Comparison of calculated differential attenuations using DWR at cloud top and DSR at the melting layer top for X-SACR/Ka-SACR (a) and KAZR/MWACR (b). Note that dielectric factor for supercooled water in (b) was scaled to ice.

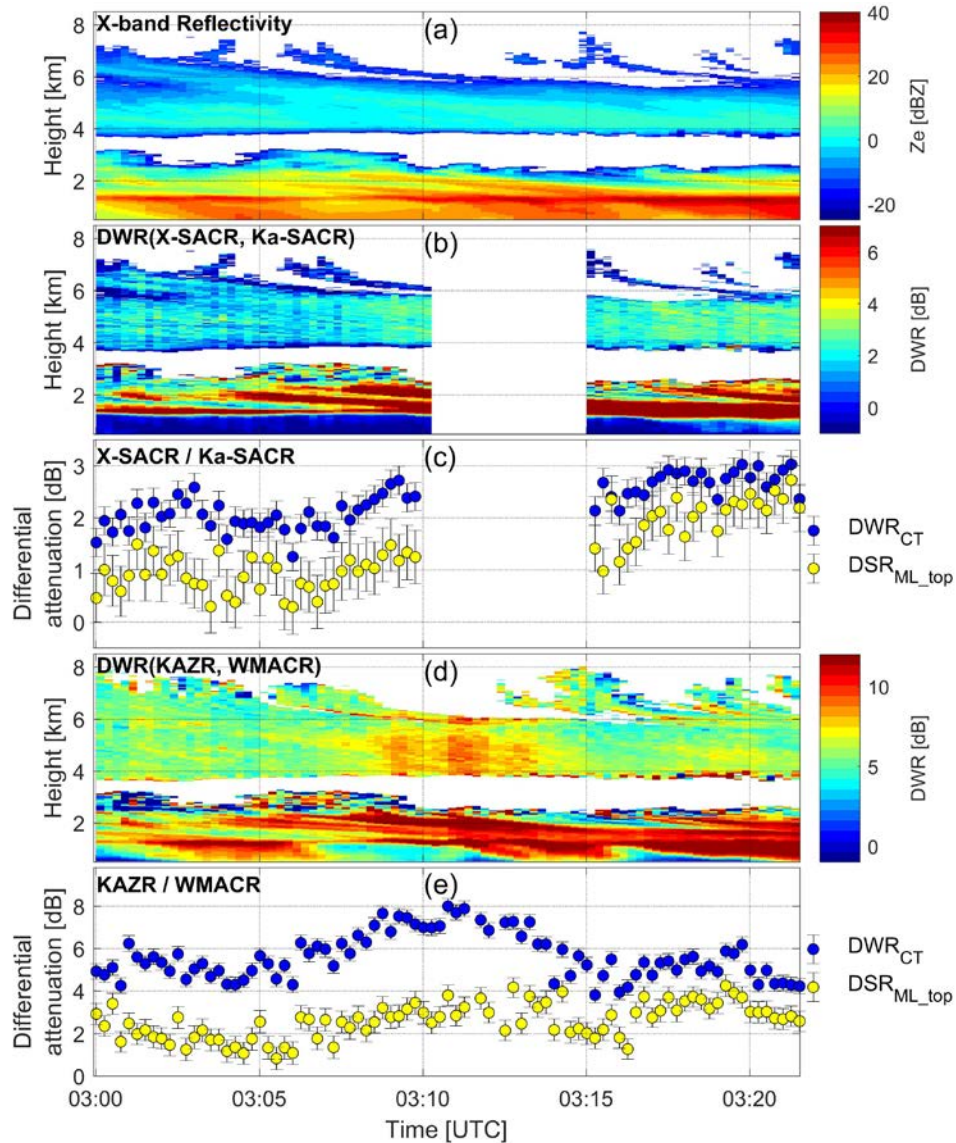


Figure 3. Differential attenuation in the melting layer for May 11<sup>th</sup> 2014 precipitation event. Radar reflectivity profiles observed by X-SACR are given in panel (a), DWR profiles for X-/Ka-SACR in (b) and KAZR/WMACR in (d). The retrieved melting layer differential attenuation from X-/Ka-SACR observations is shown in panel (c) and KAZR/WMACR in (e). Similar to Figure 2, panels (c) and (e) compare the DWR values calculated at cloud top and DSR values derived from the Rayleigh parts of spectra observations at melting layer top, respectively. The gaps during 03:10 ~ 03:15 in (b) and (c) are due to missing of X- and Ka-SACR spectra data.

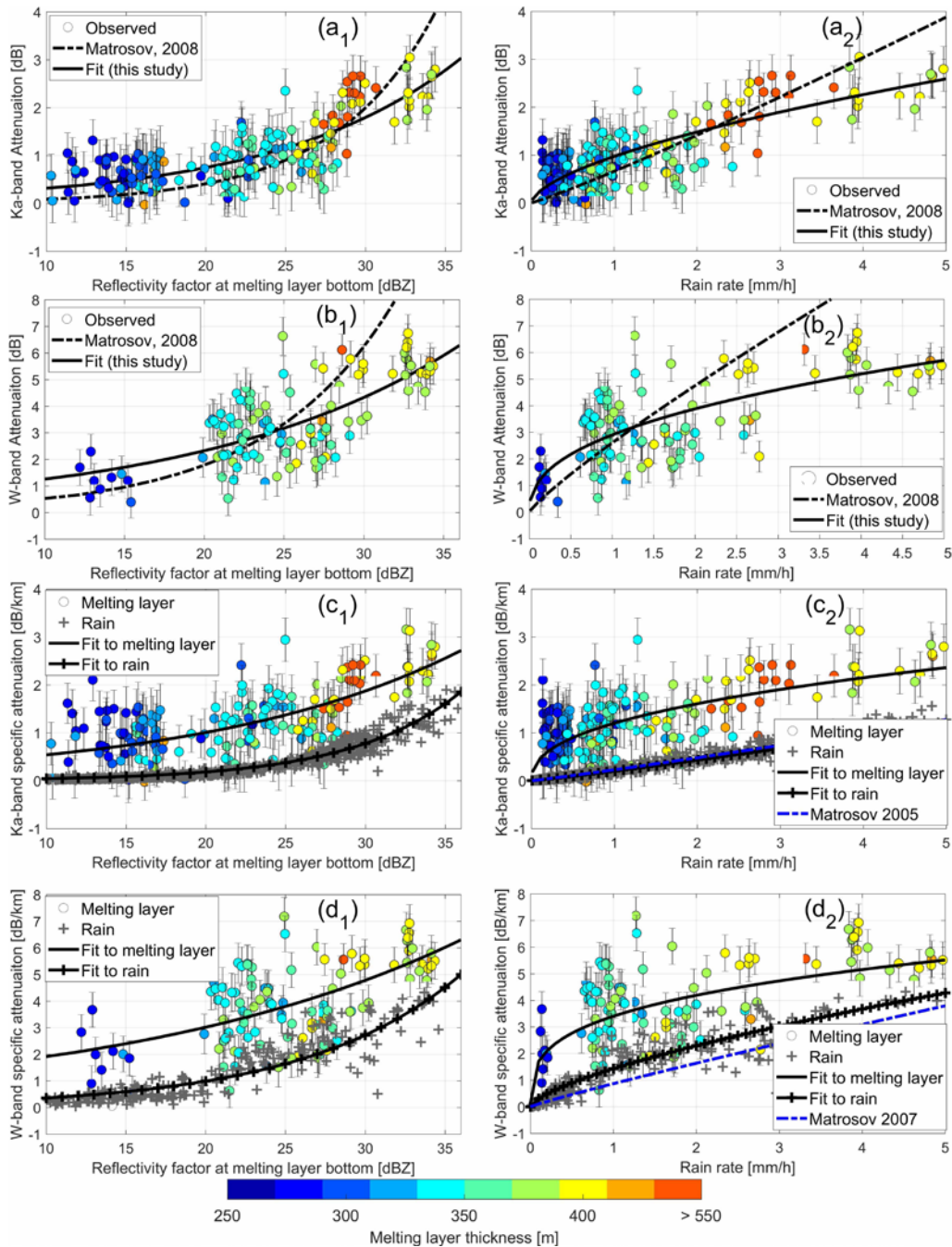


Figure 4. The derived melting-layer attenuation and specific attenuation for Ka-band (a and c) and W-band (b and d). The subscripts 1 and 2 indicate the functions of radar reflectivity factor and rain rate, respectively.