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Analysis of the Effect of Soil Roughness in the Forward-Scattering Interference Pattern Using Second-Order Small Perturbation Method Simulations

Mariano Franco[®], Emanuel More, Esteban Roitberg[®], Francisco Grings[®], *Member, IEEE*, Estefanía Piegari, Vanesa Douna, and Pablo Perna

Abstract-Soil moisture (SM) is a key geophysical variable 1 that can be estimated at regional scales using remote sensing 2 techniques, by making use of the known relationship between 3 soil reflectivity and the dielectric constant in the microwave 4 regime. In this context, the exploitation of available illuminators 5 of opportunity that currently emit large amounts of power at microwave frequencies (compared to typical synthetic aperture radar systems) is promising. Some published techniques esti-8 mate SM by analyzing the interference pattern (IP) between 9 direct and reflected signal as measured by a single antenna 10 (i.e., IP technique). In this letter, a new approach to simulate 11 the IP is proposed, in which the soil roughness is modeled 12 straightforwardly using the second-order small perturbation 13 model. Results illustrate that the "notch" in the VV-polarization 14 15 IP (related to the Brewster angle) can only be directly observed for very low values of soil rms roughness (s < 0.5 cm). For typical 16 values of soil roughness ($s \sim 1.2$ cm), the notch disappears and 17 only a minimum in the IP is observed near the Brewster angle. 18

Index Terms—Electromagnetic and remote sensing, Global
 Navigation Satellite System data, microwave radiometry, surface
 and subsurface properties.

I. INTRODUCTION

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TYPICAL approach for estimating surface soil mois-23 A ture (SM) at regional scales and high spatial resolution is 24 based on the exploitation of the signal reflected from the soil 25 in the microwave regime, using the well-known relationship 26 between the dielectric constant and SM. However, the scattered 27 signal by the surface is not only determined by SM but also 28 surface properties as its rms height (s) and the geometry of the 29 incident wave. Then, in order to get a retrieval algorithm for 30 the SM, it is essential to count with a coherent scattering model 31 which physically relates the relevant parameters involved in 32 the scattering process. A well-established theoretical model 33 to relate the soil backscattering coefficient (σ^0) with the soil 34 dielectric constant (ϵ) and soil rms roughness s is the small 35

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perturbation model (SPM) [1], [2], which has been successfully studied under several different conditions for the surface or the incident wave [3]–[6]. However, to directly invert the second-order SPM (SPM-2) in order to estimate ϵ , σ^0 of the terrain at high resolution is required, a task that can only be accomplished by a synthetic aperture radar, a relatively expensive, and power-demanding instrument.

In this letter, illuminators of opportunity (IOO) present several operative advantages for estimating SM. First, IOO bistatic radar configurations do not require a dedicated transmitter, which significantly reduces implementation costs (in both power and overall mass). Second, at least at L-band $(\lambda \sim 25 \text{ cm})$, there is evidence that soil forward scattering presents a similar sensitivity to SM than the backscattered signal [7]. Typical SM retrieval techniques based on IOOs rely on the measurement of the soil forward-scattering coefficient, which is related to SM through the dielectric constant. The main benefit of this approach is that the ratio between reflected and transmitted fields is a direct proxy of SM. However, this approach has several drawbacks. One of the most significant drawbacks is the very good antenna isolation required to separate direct and reflected components (reflected component is usually -10 to -20 dB below direct component, see [8]–[11]). Therefore, a very low level of crosstalk between antennas can be tolerated.

As an alternative, for IOOs characterized by sufficiently long pulses, the interference pattern technique (IPT) was proposed [12]. This technique is based on measuring the vertically polarized component of the received signal at the antenna. This "vertically polarized interference pattern (IP)" is analyzed in order to find a minimum ("notch"), which according to theory should correspond to the Brewster angle, which is itself related to the surface dielectric constant and SM [12]. Therefore, this technique has the advantage of requiring only one antenna and thus it does not rely on signal separation, but on the signal coherent sum. This technique has been successfully implemented and validated using fieldwork in [8], [11], and [13].

However, the analysis presented in [13] and its subsequents [14]–[16] present a key limitation: they are based on a simplified model for soil specular scattering. This means that it is assumed that soil specular scattering can be modeled 77

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using a plane interface multiplied by a roughness term to 78 analyze the relationship between the angle in which the notch 79 is found and then the dielectric constant of the soil. Therefore, 80 the position and the presence of the notch itself are based on a 81 simplified scattering model, which does take into account soil 82 roughness but does not include multiple scattering effects on 83 the surface. Hence, this analysis could be improved by using 84 a more advanced scattering model. 85

In this letter, we implemented the SPM-2 [2]-[4] for simu-86 lating the vertically polarized IP expected in the antenna as a 87 function of both geometric and dielectric soil properties (soil 88 dielectric constant and roughness). We developed a model that 89 computes the coherent sum between the electric field scattered 90 by the soil and the direct component emitted by the IOO as a 91 function of surface characteristics, assuming incident circular 92 polarization and received vertical polarization. In the analysis, 93 we show that the SPM-2 predicts the notch for very smooth 94 soils. In this condition, the notch expected monotonic behavior 95 as a function of soil dielectric constant is also observed. 96 However, the existence of the notch depends strongly on soil 97 roughness due to multiple scattering processes that take place 98 in the illuminated surface at this frequency. 99

This letter is organized as follows. In Section II, we intro-100 duce the scattering model and a solution for the proposed 101 geometry. In Section III, we present key simulation results 102 that show the expected vertically polarized IP for different soil 103 conditions and system configurations. In particular, the posi-104 tion and amplitude of the ITP minimum are studied. Finally, 105 in Section IV, we present some conclusions derived from the 106 analysis. 107

II. SCATTERING MODEL: SECOND-ORDER 109 SMALL PERTURBATION METHOD

The SPM is based on the hypothesis that soil surface has 110 a small rms height (s) with respect to the incident wave-111 length (λ), specifically that $2\pi s/\lambda \ll 0.3$ [1], [2]. Usually, 112 the SPM is used even in the limit $s \sim 0.05\lambda$, producing 113 satisfactory results (see [3]–[6]). For our case ($\lambda = 25$ cm), 114 the condition to be satisfied is s < 1.25 cm, which is a 115 reasonable assumption since for bare agricultural soils, typical 116 values of rms height are $s \sim 1$ cm [17]–[22]. This approach 117 is based on proposing that the scattered and transmitted fields 118 above and below the surface satisfy the boundary conditions 119 of the Maxwell equations. The scattered and transmitted fields 120 are written as a power series expansion in terms of the surface 121 height z, where each term accounts for different scattering 122 mechanisms on the surface. The zeroth-order term shows 123 the specular reflection, the first-order term gives a single 124

scattering behavior, and the second order represents a multiple-125 scattering process where the incident wave after the first 126 interaction with the surface goes to a second point, reflects 127 again and finally propagates to the free space. At this point, 128 it is important to remark that at the second order, the SPM 129 verifies energy conservation [23]. Thus, we are proposing a 130 physically based scattering theory that includes both surface 131 roughness and energy conservation effects. 132

We implemented the SPM up to the second order to compute 133 the IP observed in the antenna. We deal with circular polariza-134 tion (GPS systems); therefore, the incident wave and reflected 135 wave (RV) can be decomposed in terms of horizontal and 136 vertical polarizations. As the antenna only measures vertical 137 polarization, we must compute the vertical polarization of 138 the RV, which has two components: VV and HV. The first 139 one is due to a vertical incident polarization which remains 140 in the same polarization after it interacts with the surface; the 141 latter is due to the cross polarization effect, by which a wave 142 with horizontal polarization changes its state after it scatters 143 with the surface. It is well known that the cross polarization is 144 due to multiple scattering effects [2]-[5], and thus it requires 145 the development of the SPM up to the second order. 146

Based on SPM, the scattered field is expressed as [2]–[4] ¹⁴⁷

$$\mathbf{E}_{s} \approx \int \frac{d^{2}k}{(2\pi)^{2}} e^{i\mathbf{k}\cdot\mathbf{r}} \sum_{n} \left[\alpha^{(n)} \hat{h}_{s} + \beta^{(n)} \hat{v}_{s}\right]$$
(1) 148

being \hat{h}_s and \hat{v}_s the horizontal and vertical polarizations of the scattered field. Up to the second order, the amplitude of the vertical polarization mode is

$$\beta(\mathbf{k}) \approx \beta^{(0)} \,\delta(\mathbf{k} - \mathbf{k}_i) + \beta^{(1)} \,Z(\mathbf{k})$$

$$+ \int d^2 k' \,\beta^{(2)}(\mathbf{k}, \mathbf{k}') \,Z(\mathbf{k} - \mathbf{k}')$$
(2) 153

with a similar expression for the amplitude α (for a detailed development of these coefficients, see [2]–[4]). First- and second-order terms depend on the Fourier transform of the surface, $Z(\mathbf{k})$. We see that the second-order term shows a process of multiple scattering through the integration of an auxiliary mode in Z. We have a set of parameters { $\alpha^{(n)}, \beta^{(n)}$ } for the TE mode and a different set for the TM mode.

Considering that we are interested in finding the scattered field in vertical polarization from an incident wave with circular polarization, we must compute the dot product between the expression (1) and \hat{v}_s . Thus, we will have the VV contribution given by the set of TE and the HV contribution given by the set of TE. Supplicity, we call them $\xi^{(n)}$. In Section II-A, we will compute the IP due to the incident and scattered waves.

$$\langle E_T E_T^* \rangle = \frac{1}{E_i^2} \langle (E_i + E_s) \cdot (E_i^* + E_s^*) \rangle$$

$$\langle E_T E_T^* \rangle \approx \left\langle \left\{ e^{\imath \mathbf{k}_i \cdot \mathbf{r}_A} + \int \frac{d^2 k}{(2\pi)^2} e^{\imath \mathbf{k} \cdot \mathbf{r}_A} \left[\xi^{(0)} \delta \left(\mathbf{k} - \mathbf{k}_i \right) + \xi^{(1)} Z(\mathbf{k}) + \int \frac{d^2 k'}{(2\pi)^2} \xi^{(2)}(\mathbf{k}, \mathbf{k}') Z(\mathbf{k} - \mathbf{k}') \right] \right\}$$

$$\times \left\{ e^{-\imath \mathbf{k}_i \cdot \mathbf{r}_A} + \int \frac{d^2 k}{(2\pi)^2} e^{-\imath \mathbf{k} \cdot \mathbf{r}_A} \left[\xi^{(0)} \delta \left(\mathbf{k} - \mathbf{k}_i \right) + \xi^{(1)} Z(\mathbf{k}) + \int \frac{d^2 k'}{(2\pi)^2} \xi^{(2)}(\mathbf{k}, \mathbf{k}') Z(\mathbf{k} - \mathbf{k}') \right]^* \right\} \right\rangle$$

$$(3)$$

168 A. Interference Pattern

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With the goal of computing the IP received by the antenna, 169 we must compute the mean value of the total field measured 170 by it. As we have a stochastic RV, due to the scattering from 171 a random rough surface, the IP will depend on statistical 172 properties of the surface (i.e., roughness spectrum $W(\mathbf{k})$, 173 rms height s, and correlation length l). Therefore, we must 174 compute the mean power of the total intensity normalized to 175 the incident field amplitude, which is shown in (3) shown at 176 the bottom of the previous page.

Here, $\mathbf{r}_A = (x_A, z_A)$ indicates the antenna position. To per-178 form the mean value, we only need to know that the sur-179 face can be modeled as a stationary and isotropic random 180 process, with a Gaussian distribution of heights and zero 181 mean. This kind of stochastic process has the following 182 properties: $\langle Z(\mathbf{k}) \rangle = 0$ and $\langle Z(\mathbf{k}) Z^*(\mathbf{k}') \rangle = W(\mathbf{k} - \mathbf{k}')$, 183 where the last quantity is the roughness spectrum of the 184 surface, which we consider Gaussian, namely, $W(\mathbf{k}) =$ 185 $(s^2 l^2/4\pi) \exp[-l^2 \mathbf{k}^2/4]$. By using these properties, the inten-186 sity pattern that emerges from a straightforward but tedious 187 calculation results in 188

$$\langle E_T \ E_T^* \rangle \approx 1 + 2 \,\xi^{(0)} \cos(2k_{sz} \, z_A) + [\xi^{(0)}]^2 + 2 \Re \left[e^{i \, 2k_{sz} \, z_A} \int d^2 k' \,\xi^{(2)} \, W(\mathbf{k}_i - \mathbf{k}') \right] + \left| \xi^{(1)}(\mathbf{k}_i) \right|^2 \int d^2 k' \, W(\mathbf{k}')$$

$$+ 2 \xi^{(0)} \int d^2 k' \Re[\xi^{(2)}(\mathbf{k}_i, \mathbf{k}') W(\mathbf{k}_i - \mathbf{k}')].$$

The first line shown in (4) is similar (but not exactly equal) to the expression derived in [13], where the power intensity is proportional to the Fresnel reflection coefficient (here noted by $\zeta^{(0)}$) multiplied by a Gaussian roughness factor and to the phase difference between the incident and RV (i.e., $\Delta \phi =$ $2k_{sz}z_A$)

$$\langle E_T \ E_T^* \rangle \approx 1 + 2 \, \xi^{(0)} \, \cos(2k_{sz} \, z_A) + [\xi^{(0)}]^2$$

$$= |1 + \xi^{(0)} e^{i 2 \, k_{sz} \, z_A}|^2.$$
(5)

As we are using a scattering method based on small 201 perturbations, the above-mentioned term does not include any 202 roughness factor, because it is the zeroth-order solution. The 203 surface roughness appears naturally in the following terms, 204 which are proportional to the roughness spectrum $W(\mathbf{k})$. 205 Moreover, for typical agricultural soils ($s \approx 1$ cm), both terms, 206 $\xi^{(1)}$ and $\xi^{(2)}$, are relevant because they take into account the 207 multiple scattering effects produced on the surface. 208

In Section III, we use (4) to simulate the behavior of the IP measured in the antenna.

III. SIMULATED RESULTS

In Section II, we obtained the theoretical expression of the Provide the incidence of the illumination geometry and surface parameters (geometric and dielectric). In Fig. 1, the behavior of the signal measured in the antenna is shown as a function of the incidence angle for several values of soil management [18], [19], [22]; other less typical management



Fig. 1. Simulated IP as a function of the incidence angle for different values of the surface rms height s. Results correspond to L-band, i.e., $\lambda = 25$ cm, correlation length l = 10 cm, dielectric constant $\epsilon = 20$ (which is related to a medium SM value), and the scattered angle is equal to the incidence angle $\theta_s = \theta_i$.



Fig. 2. Simulated IP as a function of the incident angle for different values of soils dielectric constant ϵ . Results correspond to soil roughness s = 1 cm (typical of agricultural soils), correlation length l = 10 cm, and L-band operating frequency. The scattered angle is equal to the incidence angle $\theta_s = \theta_i$.

(e.g., harrow or roll) lead to values ~0.5 cm [20], [21]. ²¹⁹ In addition, in plowed fields, the soil rms height can reach values in the range of 2.5–4 cm [20], [21], for which notch detection will be more difficult, as we will show in Fig. 1. ²²² Finally, in order to plot the IP, the antenna is located at $z_A = 2 \text{ m}$ (or 8λ in the worst case, which satisfies the far-field condition). ²²⁵

As expected for a flat surface (s = 0 cm), a notch 226 corresponding to the Brewster angle is present for large values 227 of the incidence angle, which is in agreement with (5) and the 228 results presented in [13]. As s increases the notch starts to 229 vanish, almost disappearing for relatively large values of the 230 roughness ($s = 1.2 \text{ cm} \sim 0.05\lambda$). This is related to the fact that 231 for a very rough surface, the diffuse scattering (proportional to 232 the factors $\xi^{(1)}$ and $\xi^{(2)}$, which take into account the effects of 233 multiple scattering) becomes more relevant than the specular 234 reflection (just proportional to $\xi^{(0)}$). This result implies that 235



Fig. 3. (Top) Simulated IP as a function of the incidence angle according to two different models (modified Fresnel and SPM-2) for a low value of soil rms s = 0 cm (plane interface). (Bottom) Simulated IP as a function of the incidence angle according to two different models (modified Fresnel and SPM-2) for a typical value of soil rms s = 1.2 cm [20]. Results correspond to soil dielectric constant $\epsilon = 20$, correlation length l = 10 cm, and L-band operating frequency. The scattered angle is equal to the incidence angle $\theta_s = \theta_i$.

the SM retrieval based on simplified approaches (see [13]) will 236 have problems in detecting the notch of the IP, since it may 237 not be present. However, the simulations based on SPM-2 do 238 predict a reduction in the amplitude of the ITP (that, however, 239 could not correspond with the Brewster angle, see Fig. 1). 240 This is important since typical agricultural soils (related to 241 the "no tillage" practice) are characterized by mean values of 242 s of the order of 1 cm [17], a region in which (according 243 to our simulations) the notch will not be present and only a 244 minimum in the ITP will be observed. 245

Next, using typical values of bare soil rms roughness 246 s [17]–[20], the objective is now to test how the simulated 247 IP changes as a function of the dielectric constant of the 248 soil according to SPM-2 model. In Fig. 2, we simulated 249 the vertically polarized IP as a function of the incidence 250 angle for different values of ϵ , keeping the value of s fixed. 251 It can be seen that no notch is present, but a minimum in 252 the IP can be seen for all the range of dielectric constants. 253 Moreover, the angle in which this minimum occurs varies 254 with ϵ , as expected from theory. However, for this typical value 255 of agricultural soil roughness, the position of the minimum 256

will be difficult to locate with simple techniques—in particular, for large values of ϵ , in which the Brewster angle is known to saturate.

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Finally, it is important to evaluate how the results proposed 260 here differ from the standard approach in [13] and related 261 works. Reference [13, eqs. (4)-(9)] presents the scattering 262 model used to represent the observed signal in the antenna 263 [as discussed in Section II, the final expression is qualitatively 264 similar to the first term of the SPM2 developed here (5)]. 265 Combining simulations based on the expressions developed 266 in [13] and ours, in Fig. 3, the IPs for the two models are 267 shown for comparison, considering two extreme values of the 268 soil rms height s (s = 0 cm, s = 1.2 cm). 269

As seen, for low values of soil rms (s = 0 cm, a plane 270 interface), both models predict a notch which coincides with 271 the Brewster angle. For s = 1.2 cm, according to the modified 272 Fresnel model, the notch is still present and the overall effect 273 of soil roughness is to reduce the IP amplitude. However, 274 according to our simulations (SPM-2), the notch disappears, 275 but a minimum in the IP is still present. This differential 276 behavior between models is related to the fact that SPM-2 277 includes multiple scattering effects, which are neglected in the modified Fresnel approach. Nevertheless, if soil roughness has low values ($s \sim 0.5$ cm), the notch will be present and standard techniques can be used successfully.

IV. CONCLUSION

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In this letter, a new approach to simulate the IP based on 283 SPM-2 was proposed. We showed simulations that confirm 284 the basic behavior of the observed signal (in particular of the 285 "notch"), and its expected monotonic behavior as a function of 286 soil dielectric constant. However, according to our simulations, 287 the notch amplitude and position depend strongly on soil 288 roughness, a behavior that, according to our knowledge, is not 289 present in published models. 290

The IPT is promising since by measuring in the forward 291 direction, it has low sensitivity requirements for the receptor. 292 Moreover, by measuring the IP, no narrow angular pattern 293 of the receptor antenna is required. A drawback to this 294 technique is that it necessarily operates at low elevation angles 295 (very large scattering angles) for which the Brewster angle 296 is present. Moreover, it saturates quickly for large values of 297 the dielectric constant. In addition to these known constraints, 298 according to our simulations, soil roughness needs to be taken 299 into account in the SM retrieval scheme. Finally, the results 300 presented here are only valid for bare soils (a very atypical 301 condition of agricultural soils). In future work, we expect to 302 include the contribution of the vegetation to both signal scat-303 tering and attenuation in the computation of the interference 304 pattern. 305

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Analysis of the Effect of Soil Roughness in the Forward-Scattering Interference Pattern Using Second-Order Small Perturbation Method Simulations

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TYPICAL approach for estimating surface soil mois-23 A ture (SM) at regional scales and high spatial resolution is 24 based on the exploitation of the signal reflected from the soil 25 in the microwave regime, using the well-known relationship 26 between the dielectric constant and SM. However, the scattered 27 signal by the surface is not only determined by SM but also 28 surface properties as its rms height (s) and the geometry of the 29 incident wave. Then, in order to get a retrieval algorithm for 30 the SM, it is essential to count with a coherent scattering model 31 which physically relates the relevant parameters involved in 32 the scattering process. A well-established theoretical model 33 to relate the soil backscattering coefficient (σ^0) with the soil 34 dielectric constant (ϵ) and soil rms roughness s is the small 35

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perturbation model (SPM) [1], [2], which has been successfully studied under several different conditions for the surface or the incident wave [3]–[6]. However, to directly invert the second-order SPM (SPM-2) in order to estimate ϵ , σ^0 of the terrain at high resolution is required, a task that can only be accomplished by a synthetic aperture radar, a relatively expensive, and power-demanding instrument.

In this letter, illuminators of opportunity (IOO) present several operative advantages for estimating SM. First, IOO bistatic radar configurations do not require a dedicated transmitter, which significantly reduces implementation costs (in both power and overall mass). Second, at least at L-band $(\lambda \sim 25 \text{ cm})$, there is evidence that soil forward scattering presents a similar sensitivity to SM than the backscattered signal [7]. Typical SM retrieval techniques based on IOOs rely on the measurement of the soil forward-scattering coefficient, which is related to SM through the dielectric constant. The main benefit of this approach is that the ratio between reflected and transmitted fields is a direct proxy of SM. However, this approach has several drawbacks. One of the most significant drawbacks is the very good antenna isolation required to separate direct and reflected components (reflected component is usually -10 to -20 dB below direct component, see [8]–[11]). Therefore, a very low level of crosstalk between antennas can be tolerated.

As an alternative, for IOOs characterized by sufficiently long pulses, the interference pattern technique (IPT) was proposed [12]. This technique is based on measuring the vertically polarized component of the received signal at the antenna. This "vertically polarized interference pattern (IP)" is analyzed in order to find a minimum ("notch"), which according to theory should correspond to the Brewster angle, which is itself related to the surface dielectric constant and SM [12]. Therefore, this technique has the advantage of requiring only one antenna and thus it does not rely on signal separation, but on the signal coherent sum. This technique has been successfully implemented and validated using fieldwork in [8], [11], and [13].

However, the analysis presented in [13] and its subsequents [14]–[16] present a key limitation: they are based on a simplified model for soil specular scattering. This means that it is assumed that soil specular scattering can be modeled 77

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using a plane interface multiplied by a roughness term to 78 analyze the relationship between the angle in which the notch 79 is found and then the dielectric constant of the soil. Therefore, 80 the position and the presence of the notch itself are based on a 81 simplified scattering model, which does take into account soil 82 roughness but does not include multiple scattering effects on 83 the surface. Hence, this analysis could be improved by using 84 a more advanced scattering model. 85

In this letter, we implemented the SPM-2 [2]-[4] for simu-86 lating the vertically polarized IP expected in the antenna as a 87 function of both geometric and dielectric soil properties (soil 88 dielectric constant and roughness). We developed a model that 89 computes the coherent sum between the electric field scattered 90 by the soil and the direct component emitted by the IOO as a 91 function of surface characteristics, assuming incident circular 92 polarization and received vertical polarization. In the analysis, 93 we show that the SPM-2 predicts the notch for very smooth 94 soils. In this condition, the notch expected monotonic behavior 95 as a function of soil dielectric constant is also observed. 96 However, the existence of the notch depends strongly on soil 97 roughness due to multiple scattering processes that take place 98 in the illuminated surface at this frequency. 99

This letter is organized as follows. In Section II, we intro-100 duce the scattering model and a solution for the proposed 101 geometry. In Section III, we present key simulation results 102 that show the expected vertically polarized IP for different soil 103 conditions and system configurations. In particular, the posi-104 tion and amplitude of the ITP minimum are studied. Finally, 105 in Section IV, we present some conclusions derived from the 106 analysis. 107

II. SCATTERING MODEL: SECOND-ORDER 109 SMALL PERTURBATION METHOD

The SPM is based on the hypothesis that soil surface has 110 a small rms height (s) with respect to the incident wave-111 length (λ), specifically that $2\pi s/\lambda \ll 0.3$ [1], [2]. Usually, 112 the SPM is used even in the limit $s \sim 0.05\lambda$, producing 113 satisfactory results (see [3]–[6]). For our case ($\lambda = 25$ cm), 114 the condition to be satisfied is s < 1.25 cm, which is a 115 reasonable assumption since for bare agricultural soils, typical 116 values of rms height are $s \sim 1$ cm [17]–[22]. This approach 117 is based on proposing that the scattered and transmitted fields 118 above and below the surface satisfy the boundary conditions 119 of the Maxwell equations. The scattered and transmitted fields 120 are written as a power series expansion in terms of the surface 121 height z, where each term accounts for different scattering 122 mechanisms on the surface. The zeroth-order term shows 123 the specular reflection, the first-order term gives a single 124

scattering behavior, and the second order represents a multiple-125 scattering process where the incident wave after the first 126 interaction with the surface goes to a second point, reflects 127 again and finally propagates to the free space. At this point, 128 it is important to remark that at the second order, the SPM 129 verifies energy conservation [23]. Thus, we are proposing a 130 physically based scattering theory that includes both surface 131 roughness and energy conservation effects. 132

We implemented the SPM up to the second order to compute 133 the IP observed in the antenna. We deal with circular polariza-134 tion (GPS systems); therefore, the incident wave and reflected 135 wave (RV) can be decomposed in terms of horizontal and 136 vertical polarizations. As the antenna only measures vertical 137 polarization, we must compute the vertical polarization of 138 the RV, which has two components: VV and HV. The first 139 one is due to a vertical incident polarization which remains 140 in the same polarization after it interacts with the surface; the 141 latter is due to the cross polarization effect, by which a wave 142 with horizontal polarization changes its state after it scatters 143 with the surface. It is well known that the cross polarization is 144 due to multiple scattering effects [2]-[5], and thus it requires 145 the development of the SPM up to the second order. 146

Based on SPM, the scattered field is expressed as [2]–[4] ¹⁴⁷

$$\mathbf{E}_{s} \approx \int \frac{d^{2}k}{(2\pi)^{2}} e^{i\mathbf{k}\cdot\mathbf{r}} \sum_{n} \left[\alpha^{(n)} \hat{h}_{s} + \beta^{(n)} \hat{v}_{s}\right]$$
(1) 148

being \hat{h}_s and \hat{v}_s the horizontal and vertical polarizations of the scattered field. Up to the second order, the amplitude of the vertical polarization mode is

$$\beta(\mathbf{k}) \approx \beta^{(0)} \,\delta(\mathbf{k} - \mathbf{k}_i) + \beta^{(1)} \,Z(\mathbf{k})$$

$$+ \int d^2 k' \,\beta^{(2)}(\mathbf{k}, \mathbf{k}') \,Z(\mathbf{k} - \mathbf{k}')$$
(2) 153

with a similar expression for the amplitude α (for a detailed development of these coefficients, see [2]–[4]). First- and second-order terms depend on the Fourier transform of the surface, $Z(\mathbf{k})$. We see that the second-order term shows a process of multiple scattering through the integration of an auxiliary mode in Z. We have a set of parameters { $\alpha^{(n)}, \beta^{(n)}$ } for the TE mode and a different set for the TM mode.

Considering that we are interested in finding the scattered field in vertical polarization from an incident wave with circular polarization, we must compute the dot product between the expression (1) and \hat{v}_s . Thus, we will have the VV contribution given by the set of TE and the HV contribution given by the set of TE. Supplicity, we call them $\xi^{(n)}$. In Section II-A, we will compute the IP due to the incident and scattered waves.

$$\langle E_T E_T^* \rangle = \frac{1}{E_i^2} \langle (E_i + E_s) \cdot (E_i^* + E_s^*) \rangle$$

$$\langle E_T E_T^* \rangle \approx \left\langle \left\{ e^{\imath \mathbf{k}_i \cdot \mathbf{r}_A} + \int \frac{d^2 k}{(2\pi)^2} e^{\imath \mathbf{k} \cdot \mathbf{r}_A} \left[\xi^{(0)} \delta \left(\mathbf{k} - \mathbf{k}_i \right) + \xi^{(1)} Z(\mathbf{k}) + \int \frac{d^2 k'}{(2\pi)^2} \xi^{(2)}(\mathbf{k}, \mathbf{k}') Z(\mathbf{k} - \mathbf{k}') \right] \right\}$$

$$\times \left\{ e^{-\imath \mathbf{k}_i \cdot \mathbf{r}_A} + \int \frac{d^2 k}{(2\pi)^2} e^{-\imath \mathbf{k} \cdot \mathbf{r}_A} \left[\xi^{(0)} \delta \left(\mathbf{k} - \mathbf{k}_i \right) + \xi^{(1)} Z(\mathbf{k}) + \int \frac{d^2 k'}{(2\pi)^2} \xi^{(2)}(\mathbf{k}, \mathbf{k}') Z(\mathbf{k} - \mathbf{k}') \right]^* \right\} \right\rangle$$

$$(3)$$

168 A. Interference Pattern

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With the goal of computing the IP received by the antenna, 169 we must compute the mean value of the total field measured 170 by it. As we have a stochastic RV, due to the scattering from 171 a random rough surface, the IP will depend on statistical 172 properties of the surface (i.e., roughness spectrum $W(\mathbf{k})$, 173 rms height s, and correlation length l). Therefore, we must 174 compute the mean power of the total intensity normalized to 175 the incident field amplitude, which is shown in (3) shown at 176 the bottom of the previous page.

Here, $\mathbf{r}_A = (x_A, z_A)$ indicates the antenna position. To per-178 form the mean value, we only need to know that the sur-179 face can be modeled as a stationary and isotropic random 180 process, with a Gaussian distribution of heights and zero 181 mean. This kind of stochastic process has the following 182 properties: $\langle Z(\mathbf{k}) \rangle = 0$ and $\langle Z(\mathbf{k}) Z^*(\mathbf{k}') \rangle = W(\mathbf{k} - \mathbf{k}')$, 183 where the last quantity is the roughness spectrum of the 184 surface, which we consider Gaussian, namely, $W(\mathbf{k}) =$ 185 $(s^2 l^2/4\pi) \exp[-l^2 \mathbf{k}^2/4]$. By using these properties, the inten-186 sity pattern that emerges from a straightforward but tedious 187 calculation results in 188

$$\langle E_T \ E_T^* \rangle \approx 1 + 2 \,\xi^{(0)} \cos(2k_{sz} \, z_A) + [\xi^{(0)}]^2 + 2 \Re \left[e^{i \, 2k_{sz} \, z_A} \int d^2 k' \,\xi^{(2)} \, W(\mathbf{k}_i - \mathbf{k}') \right] + \left| \xi^{(1)}(\mathbf{k}_i) \right|^2 \int d^2 k' \, W(\mathbf{k}')$$

$$+ 2\,\xi^{(0)} \int d^2k' \Re[\xi^{(2)}(\mathbf{k}_i,\mathbf{k}') W(\mathbf{k}_i-\mathbf{k}')]$$

The first line shown in (4) is similar (but not exactly equal) to the expression derived in [13], where the power intensity is proportional to the Fresnel reflection coefficient (here noted by $\xi^{(0)}$) multiplied by a Gaussian roughness factor and to the phase difference between the incident and RV (i.e., $\Delta \phi =$ $2k_{sz}z_A$)

$$\langle E_T \ E_T^* \rangle \approx 1 + 2 \, \xi^{(0)} \, \cos(2k_{sz} \, z_A) + [\xi^{(0)}]^2$$

$$= |1 + \xi^{(0)} \ e^{i 2 \, k_{sz} \, z_A}|^2.$$
(5)

As we are using a scattering method based on small 201 perturbations, the above-mentioned term does not include any 202 roughness factor, because it is the zeroth-order solution. The 203 surface roughness appears naturally in the following terms, 204 which are proportional to the roughness spectrum $W(\mathbf{k})$. 205 Moreover, for typical agricultural soils ($s \approx 1$ cm), both terms, 206 $\xi^{(1)}$ and $\xi^{(2)}$, are relevant because they take into account the 207 multiple scattering effects produced on the surface. 208

In Section III, we use (4) to simulate the behavior of the IP measured in the antenna.

III. SIMULATED RESULTS

In Section II, we obtained the theoretical expression of the Provide the incidence of the illumination geometry and surface parameters (geometric and dielectric). In Fig. 1, the behavior of the signal measured in the antenna is shown as a function of the incidence angle for several values of soil management [18], [19], [22]; other less typical management



Fig. 1. Simulated IP as a function of the incidence angle for different values of the surface rms height s. Results correspond to L-band, i.e., $\lambda = 25$ cm, correlation length l = 10 cm, dielectric constant $\epsilon = 20$ (which is related to a medium SM value), and the scattered angle is equal to the incidence angle $\theta_s = \theta_i$.



Fig. 2. Simulated IP as a function of the incident angle for different values of soils dielectric constant ϵ . Results correspond to soil roughness s = 1 cm (typical of agricultural soils), correlation length l = 10 cm, and L-band operating frequency. The scattered angle is equal to the incidence angle $\theta_s = \theta_i$.

(e.g., harrow or roll) lead to values ~0.5 cm [20], [21]. ²¹⁹ In addition, in plowed fields, the soil rms height can reach values in the range of 2.5–4 cm [20], [21], for which notch detection will be more difficult, as we will show in Fig. 1. ²²² Finally, in order to plot the IP, the antenna is located at $z_A = 2 \text{ m}$ (or 8λ in the worst case, which satisfies the far-field condition). ²²⁵

As expected for a flat surface (s = 0 cm), a notch 226 corresponding to the Brewster angle is present for large values 227 of the incidence angle, which is in agreement with (5) and the 228 results presented in [13]. As s increases the notch starts to 229 vanish, almost disappearing for relatively large values of the 230 roughness ($s = 1.2 \text{ cm} \sim 0.05\lambda$). This is related to the fact that 231 for a very rough surface, the diffuse scattering (proportional to 232 the factors $\xi^{(1)}$ and $\xi^{(2)}$, which take into account the effects of 233 multiple scattering) becomes more relevant than the specular 234 reflection (just proportional to $\xi^{(0)}$). This result implies that 235



Fig. 3. (Top) Simulated IP as a function of the incidence angle according to two different models (modified Fresnel and SPM-2) for a low value of soil rms s = 0 cm (plane interface). (Bottom) Simulated IP as a function of the incidence angle according to two different models (modified Fresnel and SPM-2) for a typical value of soil rms s = 1.2 cm [20]. Results correspond to soil dielectric constant $\epsilon = 20$, correlation length l = 10 cm, and L-band operating frequency. The scattered angle is equal to the incidence angle $\theta_s = \theta_i$.

the SM retrieval based on simplified approaches (see [13]) will 236 have problems in detecting the notch of the IP, since it may 237 not be present. However, the simulations based on SPM-2 do 238 predict a reduction in the amplitude of the ITP (that, however, 239 could not correspond with the Brewster angle, see Fig. 1). 240 This is important since typical agricultural soils (related to 241 the "no tillage" practice) are characterized by mean values of 242 s of the order of 1 cm [17], a region in which (according 243 to our simulations) the notch will not be present and only a 244 minimum in the ITP will be observed. 245

Next, using typical values of bare soil rms roughness 246 s [17]–[20], the objective is now to test how the simulated 247 IP changes as a function of the dielectric constant of the 248 soil according to SPM-2 model. In Fig. 2, we simulated 249 the vertically polarized IP as a function of the incidence 250 angle for different values of ϵ , keeping the value of s fixed. 251 It can be seen that no notch is present, but a minimum in 252 the IP can be seen for all the range of dielectric constants. 253 Moreover, the angle in which this minimum occurs varies 254 with ϵ , as expected from theory. However, for this typical value 255 of agricultural soil roughness, the position of the minimum 256

will be difficult to locate with simple techniques—in particular, for large values of ϵ , in which the Brewster angle is known to saturate.

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Finally, it is important to evaluate how the results proposed 260 here differ from the standard approach in [13] and related 261 works. Reference [13, eqs. (4)-(9)] presents the scattering 262 model used to represent the observed signal in the antenna 263 [as discussed in Section II, the final expression is qualitatively 264 similar to the first term of the SPM2 developed here (5)]. 265 Combining simulations based on the expressions developed 266 in [13] and ours, in Fig. 3, the IPs for the two models are 267 shown for comparison, considering two extreme values of the 268 soil rms height s (s = 0 cm, s = 1.2 cm). 269

As seen, for low values of soil rms (s = 0 cm, a plane 270 interface), both models predict a notch which coincides with 271 the Brewster angle. For s = 1.2 cm, according to the modified 272 Fresnel model, the notch is still present and the overall effect 273 of soil roughness is to reduce the IP amplitude. However, 274 according to our simulations (SPM-2), the notch disappears, 275 but a minimum in the IP is still present. This differential 276 behavior between models is related to the fact that SPM-2 277 includes multiple scattering effects, which are neglected in the modified Fresnel approach. Nevertheless, if soil roughness has low values ($s \sim 0.5$ cm), the notch will be present and standard techniques can be used successfully.

IV. CONCLUSION

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In this letter, a new approach to simulate the IP based on 283 SPM-2 was proposed. We showed simulations that confirm 284 the basic behavior of the observed signal (in particular of the 285 "notch"), and its expected monotonic behavior as a function of 286 soil dielectric constant. However, according to our simulations, 287 the notch amplitude and position depend strongly on soil 288 roughness, a behavior that, according to our knowledge, is not 289 present in published models. 290

The IPT is promising since by measuring in the forward 291 direction, it has low sensitivity requirements for the receptor. 292 Moreover, by measuring the IP, no narrow angular pattern 293 of the receptor antenna is required. A drawback to this 294 technique is that it necessarily operates at low elevation angles 295 (very large scattering angles) for which the Brewster angle 296 is present. Moreover, it saturates quickly for large values of 297 the dielectric constant. In addition to these known constraints, 298 according to our simulations, soil roughness needs to be taken 299 into account in the SM retrieval scheme. Finally, the results 300 presented here are only valid for bare soils (a very atypical 301 condition of agricultural soils). In future work, we expect to 302 include the contribution of the vegetation to both signal scat-303 tering and attenuation in the computation of the interference 304 pattern. 305

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