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I. 5 The Decay of a $T_z = -1/2$ Mirror Nucleus ^{57}Cu

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The neutron-deficient mirror nuclei of the $f_{7/2}p_{3/2}$ shell region are difficult to observe because of their small production cross sections and their short half-lives.¹⁻³⁾ In order to study such a nucleus far from the beta-stability line, we developed a high-speed Target Rotation system (TARO) mentioned before.⁴⁾ Using this system, we studied the decay of a new short-lived nucleus ^{57}Cu , which is very interesting because it consists of the ^{56}Ni core and one proton in the $p_{3/2}$ shell. From this experiment, the half-life and $E_{\beta^+}^{\text{max}}$ of the decay of ^{57}Cu were obtained.

^{57}Cu was produced in the $(p,2n)$ reaction on ^{58}Ni targets of 2-3 mg/cm² enriched to 99.76 %. The proton beam current and energy were typically 0.2-0.3 μA and 25-34 MeV, respectively. Sixteen ^{58}Ni targets were distributed on the rotating plate⁴⁾ in order to reduce the accumulation of long-lived isotopes. The proton beam was pulsed to irradiate these targets sequentially; the irradiated target was rotated in 60 ms to the shielded detector position and was measured.⁴⁾

The γ -rays and positrons were measured in two-parameter (energy and time) list modes with a 96 cm³ coaxial Ge(HP) and a 490 mm² \times 10 mm planer Ge(HP) detectors, respectively. We collected typically 10^7 events for positrons.

In the γ -ray energy spectrum we discovered the γ -ray of 1112 keV decaying with a half-life of $T_{1/2} = 233 \pm 16$ ms, which we assigned to depopulate the second-excited state ($I^\pi = 1/2^-$) of ^{57}Ni ⁵⁾ fed from the decay of ^{57}Cu . The γ -ray's energy was obtained to be 1111.7 ± 0.5 keV. Figure 1 shows a γ -ray spectrum and the decay of the 1111.7 keV γ -ray. This half-life value is consistent with a value of 0.2 s predicted from the gross theory of β -decay.⁶⁾ No other γ -rays from ^{57}Cu were found.

As an aid in the identification of ^{57}Cu , the γ -ray yield curves were compared with theoretical excitation functions calculated by ALICE.⁷⁾ The experimental curves were in good agreement with the theoretical ones as shown in Fig. 2.

The observed positron spectrum included the backgrounds due to ^{54}Co ($T_{1/2} = 196$ ms, $E_{\beta^+}^{\text{max}} = 7.2$ MeV⁵⁾) and ^{58}Cu ($T_{1/2} = 3.2$ s, $E_{\beta^+}^{\text{max}} = 7.5$ MeV⁵⁾). The $E_{\beta^+}^{\text{max}}$ of ^{57}Cu was predicted to be closed to $E_{\beta^+}^{\text{max}}$ of these background.⁸⁾ Therefore, we determined $E_{\beta^+}^{\text{max}}$ of ^{57}Cu using a "time-filtering" analysis. By taking advantage of the difference in half-life, we subtracted the background due to ^{58}Cu and others from the raw positron spectrum as shown in Fig. 3. This "time-filtered" spectrum should be composed of positrons from short-lived ($T_{1/2} \approx 200$ ms) nuclei. This spectrum showed a $E_{\beta^+}^{\text{max}}$ larger than that of ^{54}Co measured using the

$^{54}\text{Fe}(p,n)^{54}\text{Co}$ reaction. We deduced $E_{\beta^+}^{\text{max}}(^{57}\text{Cu})$ by the Kurie-plot analysis of the part of the "time-filtered" spectrum having $E_{\beta^+} > E_{\beta^+}^{\text{max}}(^{54}\text{Co})$ on the basis of the method of Rehfield.⁹⁾ Thus we obtained $E_{\beta^+}^{\text{max}}(^{57}\text{Cu}) = 7.72 \pm 0.13$ MeV. The branching ratio to the 1111.7 keV excited state of ^{57}Ni was found to be 3.7 ± 1.7 % from γ - and positron spectra.

Using the f -value from Ref. 10), the $\log ft$ value for the ground state decay of ^{57}Cu was obtained to be $\log ft = 3.71 \pm 0.05$, and the GT matrix element was obtained¹¹⁾ to be $\langle \sigma \tau \rangle_{\text{exp}} = 0.37 \pm 0.11$. This value indicates a large reduction in comparison with the $p_{3/2}$ pure single-particle estimate $\langle \sigma \tau \rangle_{\text{s.p.}} = 1.291$.

The measured mass excess of ^{57}Cu is compared with several mass predictions¹²⁾ in Table 1. Agreement between experiment and theory is fairly good except the value estimated by Meyers with semiempirical droplet model. The Coulomb displacement energy was found to be $\Delta E_{\text{C}} = 9.524 \pm 0.13$ MeV from the present value. This value is in good agreement with the estimated value using the phenomenological formulae of Jänecke.¹³⁾

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Table 1. Comparison of mass predictions with the experimental value.

Mass prediction ¹²⁾	Mass excess (MeV)
Meyers	-51.47
Groote-Hilf-Takahashi	-47.73
Seeger-Howard	-47.7
Liran-Zeldes	-47.20
Beiner-Lombard-Mas	-44.9
Jänecke-Garvey-Kelson	-47.43
Experimental value	-47.36 (13)

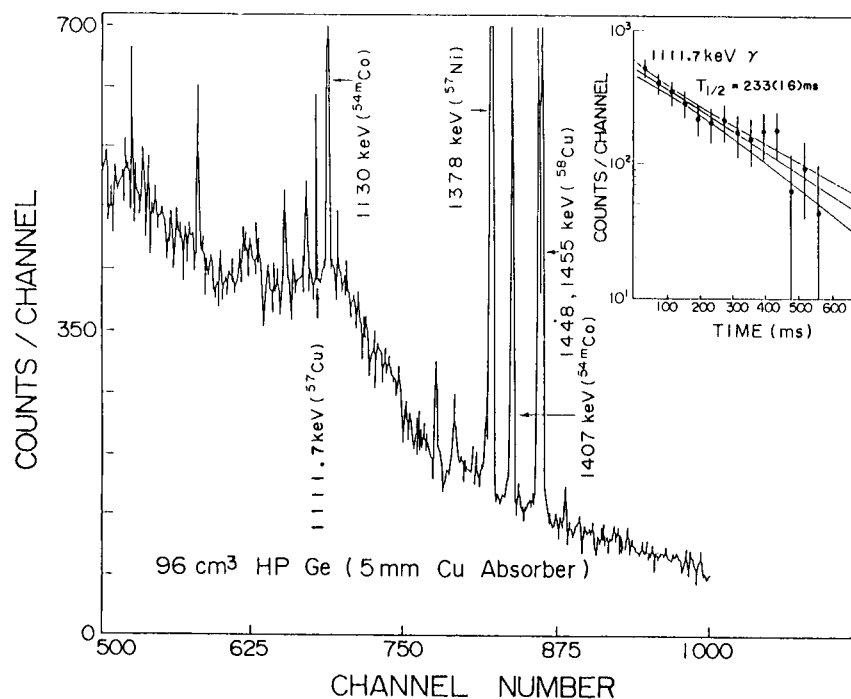


Fig. 1. The γ -ray spectrum and the half-life of 1111.7 keV γ -rays assigned to the decay of ^{57}Cu

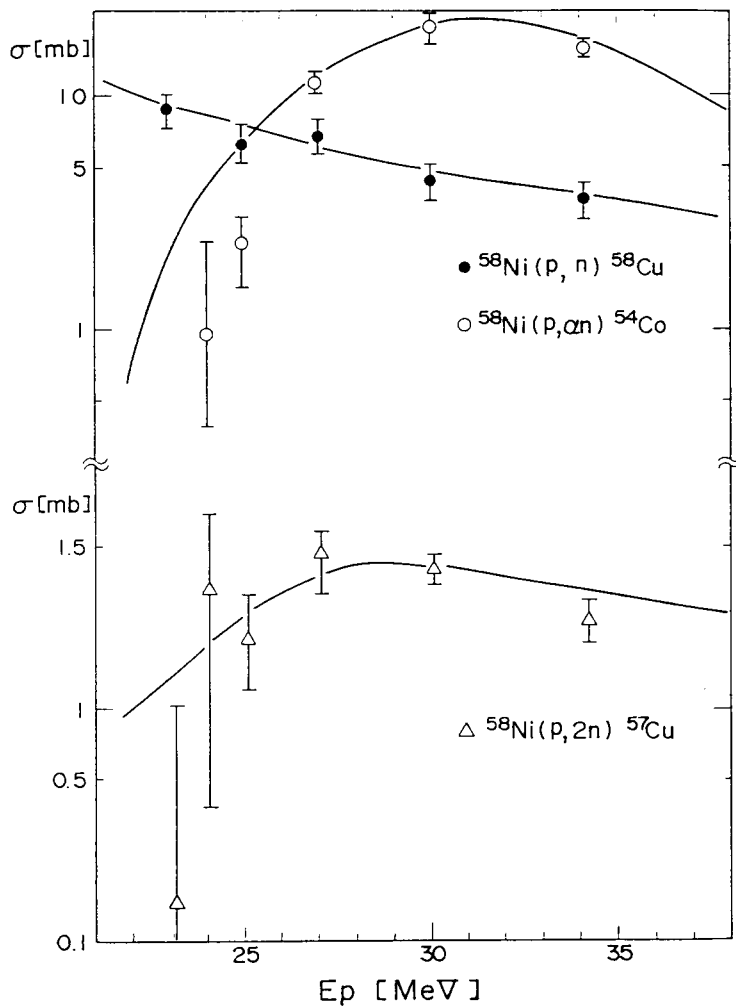


Fig. 2. Yield curves for ^{57}Cu , ^{54}Co and ^{58}Cu in the (p,2n), (p,an) and (p,n) reactions, respectively. The solid curves are theoretical ones calculated by the cross section code ALICE.⁷⁾ The experimental data are normalized at the bombarding energy of 30 MeV.

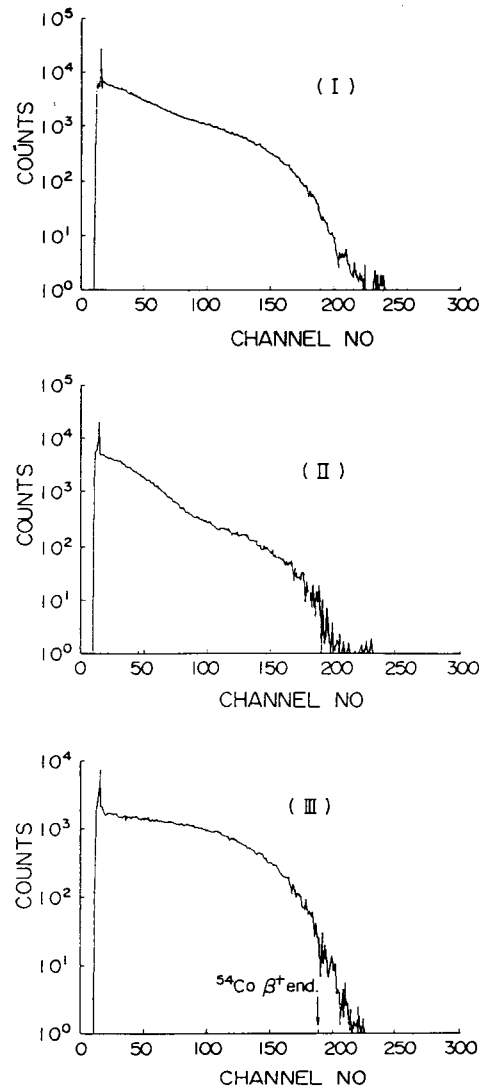


Fig. 3. Analysis of the positron spectrum; (I) is the raw spectrum and (II) and (III) are the "time filtered" spectra obtained by decomposing (I) by time-filtering of each channel. (II) is the long-lived part corresponding to ^{58}Cu , $^{54\text{m}}\text{Co}$ and others and (III) is the short-lived part corresponding to ^{54}Co and ^{57}Cu .