

## Scaling theory for double-gate SOI MOSFET's

著者	田中 徹
journal or publication title	IEEE Transactions on Electron Devices
volume	40
number	12
page range	2326-2329
year	1993
URL	<a href="http://hdl.handle.net/10097/47472">http://hdl.handle.net/10097/47472</a>

doi: 10.1109/16.249482

# Scaling Theory for Double-Gate SOI MOSFET's

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**Abstract**—We established a scaling theory for double-gate SOI MOSFET's, which gives a guidance for the device design (silicon thickness  $t_{si}$ ; gate oxide thickness  $t_{ox}$ ) so that maintaining a subthreshold factor for a given gate length  $L_G$ . According to our theory, a device can be designed with a gate length of less than  $0.1 \mu\text{m}$  while maintaining the ideal subthreshold factor, which is verified numerically with a two-dimensional device simulator.

## I. INTRODUCTION

ACCORDING to the Brew's scaling theory [1], the channel doping concentration in bulk MOSFET's should be increased to alleviate the short-channel effects, leading to more than  $10^{18} \text{ cm}^{-3}$  for a gate length of less than  $0.1 \mu\text{m}$  [2]. This high-doping concentration degrades device performance due to decreased mobility and increased junction capacitance.

A double-gate SOI MOSFET, in which the potential is controlled by front and back gates as shown in Fig. 1, is proposed to circumvent the scaling limitations of bulk MOSFET's. The characteristics of this device have been studied [2]–[11], revealing its ideal subthreshold factor, high transconductance, and short-channel effect immunity.

Since the potential distribution in double-gate SOI MOSFET's differs greatly from that in bulk and single-gate SOI MOSFET's because of symmetrical device structure with quite low channel doping concentration [4], [7], the scaling theory developed for bulk MOSFET's cannot be applied to double-gate SOI MOSFET's.

Yan *et al.* proposed a unique scaling theory for double-gate SOI MOSFET's [2]. According to their theory, the device should be designed maintaining

$$\alpha_1 = \frac{L_G}{2\lambda_1} \quad (1)$$

where  $\lambda_1$  is the so-called natural length which characterizes the short-channel effect and is given by

$$\lambda_1 = \sqrt{\frac{\epsilon_{si}}{2\epsilon_{ox}}} t_{si} t_{ox} \quad (2)$$

where  $\epsilon_{si}$  is the dielectric constant of silicon,  $\epsilon_{ox}$  is the dielectric constant of silicon dioxide,  $t_{si}$  is the SOI silicon

thickness, and  $t_{ox}$  is the gate oxide thickness. This natural length is an easy guide for choosing device parameters, and has simple physical meaning that a small natural length corresponds to superb short channel effect immunity. Although the authors showed that a small  $\alpha$  gives a degraded  $S$ -factor, they did not show whether the same  $\alpha$  with various device parameters gives the same  $S$ -factor.

Fig. 2(a) shows the dependence of  $S$ -factor on  $\alpha_1$  calculated using a two-dimensional device simulator [12]. Although a smaller  $\alpha_1$  gives a larger  $S$ -factor, the same  $\alpha_1$  with various combinations of device parameters does not give the same  $S$ -factor, that is, the theory does not provide the same guideline for different gate lengths.

The onset point where the  $S$ -factor degrades does not depend on the drain voltage,  $V_D$ , although the value of the  $S$ -factor after degradation strongly depends on  $V_D$ . Since the purpose of our analysis is to clarify the onset point, we restricted our analysis to  $V_D = 0.05 \text{ V}$  unless specified.

Yan *et al.* [2] assumed that the punchthrough current flows along the surface, which is invalid for a double-gate SOI MOSFET for the following reasons.

The maximum potential at the SOI center,  $\phi_c$ , is more sensitive to gate length than that at the surface,  $\phi_s$ , and, furthermore, the absolute value of  $\phi_c$  is smaller than that of  $\phi_s$  (Fig. 3), meaning that the punchthrough current dominantly flows at the SOI center.

In this paper, we derive a scaling theory relevant to  $\phi_c$ , and show that our scaling theory acts as a guide to design devices that hold proper  $S$ -factors.

## II. THEORY

The Poisson equation of potential,  $\phi$ , is [13]

$$\frac{d^2\phi(x, y)}{dx^2} + \frac{d^2\phi(x, y)}{dy^2} = \frac{qN_A}{\epsilon_{si}} \quad (3)$$

where  $N_A$  is the channel doping concentration, and the  $y$ -axis is perpendicular and the  $x$ -axis is parallel to the channel (Fig. 1).

Using the same parabolic potential profile in the vertical direction as Young used [13] and applying the boundary condition of  $d\phi/dy = 0$  for  $y = t_{si}/2$ , we obtained the same  $\phi(x, y)$  as that in [2] given by

$$\begin{aligned} \phi(x, y) = & \phi_s(x) + \frac{\epsilon_{ox}}{\epsilon_{si}} \frac{\phi_s(x) - (V_G - V_{FB})}{t_{ox}} y \\ & - \frac{\epsilon_{ox}}{\epsilon_{si}} \frac{\phi_s(x) - (V_G - V_{FB})}{t_{ox} t_{si}} y^2 \end{aligned} \quad (4)$$

Manuscript received June 4, 1993; revised August 10, 1993. The review of this paper was arranged by Associate Editor J. R. Pfister.

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IEEE Log Number 9212783.

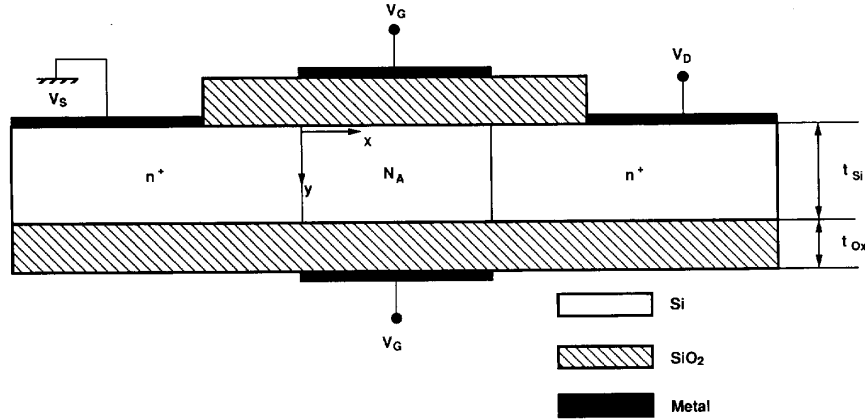


Fig. 1. Double-gate MOSFET.

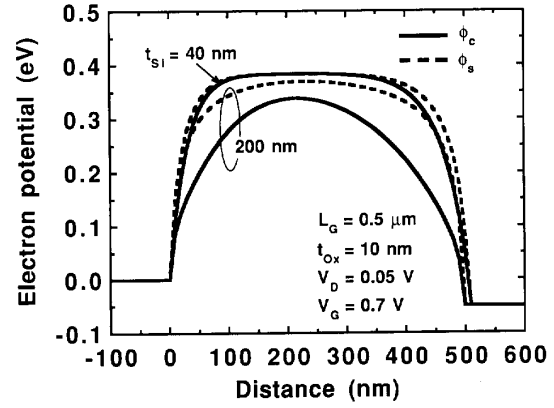
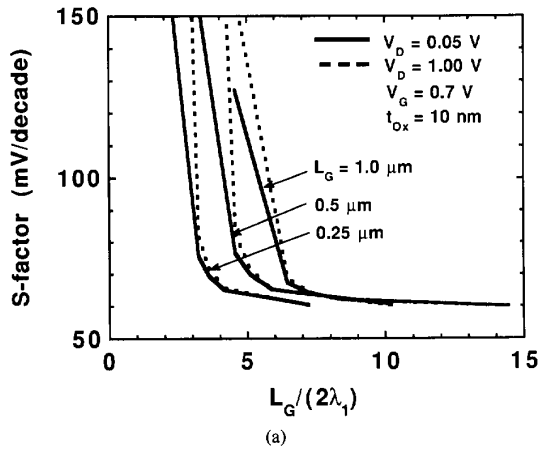
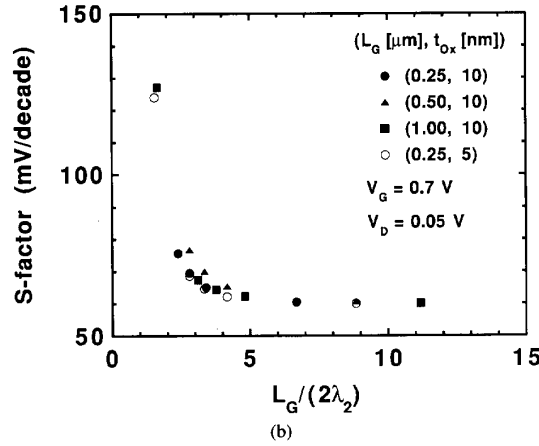
Fig. 3. Potential distribution along channel at the surface,  $\phi_s$ , and at the center of SOI,  $\phi_c$ .

Fig. 2. Dependence of subthreshold factor on the scaling parameter: (a) former model; (b) proposed model.

where  $V_G$  is the gate voltage and  $V_{FB}$  is the flatband voltage.

Since  $\phi_c$  should be relevant to the punchthrough current, we obtained the relation between  $\phi_s$  and  $\phi_c$  from (4)

by substituting  $y = t_{si}/2$  as

$$\phi_s(x) = \frac{1}{1 + \frac{\epsilon_{ox} t_{si}}{4\epsilon_{si} t_{ox}}} \left[ \phi_c(x) + \frac{\epsilon_{ox} t_{si}}{4\epsilon_{si} t_{ox}} (V_G - V_{FB}) \right], \quad (5)$$

and then expressed  $\phi(x, y)$  using  $\phi_c$  as

$$\phi(x, y) = \left[ 1 + \frac{\epsilon_{ox} y}{\epsilon_{si} t_{ox}} - \frac{\epsilon_{ox} y^2}{\epsilon_{si} t_{ox} t_{si}} \right] \frac{\phi_c(x) + \frac{\epsilon_{ox} t_{si}}{4\epsilon_{si} t_{ox}} (V_G - V_{FB})}{1 + \frac{\epsilon_{ox} t_{si}}{4\epsilon_{si} t_{ox}}} - \frac{\epsilon_{ox} y}{\epsilon_{si} t_{ox}} (V_G - V_{FB}) + \frac{\epsilon_{ox} y^2}{\epsilon_{si} t_{ox} t_{si}} (V_G - V_{FB}). \quad (6)$$

Substituting (6) into (3), we obtain

$$\frac{d^2 \phi_c(x)}{dx^2} + \frac{V_G - V_{FB} - \phi_c(x)}{\lambda_2^2} = \frac{qN_A}{\epsilon_{si}} \quad (7)$$

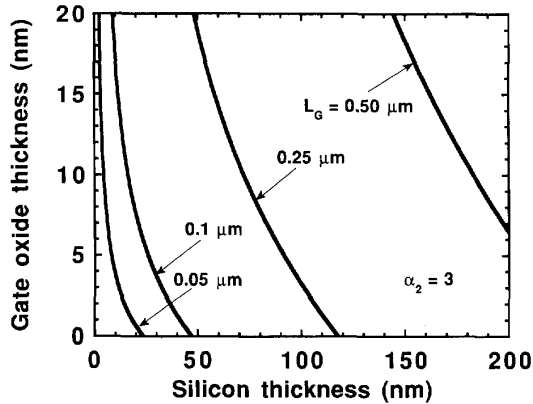


Fig. 4. Relationship between silicon thickness and gate oxide thickness with various gate lengths alleviating short channel effect.

where  $\lambda_2$  is the natural length having the same physical meaning as  $\lambda_1$  and is

$$\lambda_2 = \sqrt{\frac{\epsilon_{si}}{2\epsilon_{ox}} \left( 1 + \frac{\epsilon_{ox} t_{si}}{4\epsilon_{si} t_{ox}} \right) t_{si} t_{ox}}. \quad (8)$$

Note that  $\lambda_2$  is always larger than  $\lambda_1$ , that is, the double-gate SOI MOSFET suffers from the short-channel effect more than predicted by Yan's model [2].

Exactly the same analysis as in [2] is then followed. The key conclusion from this analysis is that the same  $\alpha_2$  leads to the same  $S$ -factor, where

$$\alpha_2 = \frac{L_G}{2\lambda_2}. \quad (9)$$

### III. RESULTS AND DISCUSSION

Throughout this analysis, we fixed  $N_A$  to  $10^{15} \text{ cm}^{-3}$  and  $V_{FB}$  to 0.29 V which corresponds to  $p^+$  polysilicon gate. The threshold voltage is about 1.0 V independent of device parameters and the  $S$ -factor is also independent of  $V_G$  in the subthreshold region when it has almost ideal value [14]. We evaluated  $S$ -factor for  $V_G = 0.7$  V.

Fig. 2(b) shows the dependence of the  $S$ -factor on  $\alpha_2$  calculated using a two-dimensional device simulator [12]. The same  $\alpha_2$  leads to the same value of  $S$ -factor with various combinations of device parameters,  $t_{ox}$ ,  $t_{si}$ , and  $L_G$ . We found that an  $\alpha_2$  of less than 3 is needed to alleviate the punchthrough, that is, short-channel effect.

Once  $\alpha_2$  is determined from Fig. 2 (b), the relationship between  $t_{ox}$  and  $t_{si}$  is expressed from (9) as

$$t_{ox} = \frac{\epsilon_{ox} L_G^2}{2\alpha_2^2 \epsilon_{si} t_{si}} - \frac{\epsilon_{ox}}{4\epsilon_{si}} t_{si}. \quad (10)$$

$t_{ox}$  and  $t_{si}$  at a given  $L_G$  should be designed in the lower region of the corresponding  $L_G$  curve in Fig. 4. As  $L_G$  decreases, the allowable region decreases. When  $L_G = 0.1 \mu\text{m}$  with  $t_{ox} = 5$  nm,  $t_{si}$  should be less than 25 nm, and when  $L_G = 0.05 \mu\text{m}$  with  $t_{ox} = 3$  nm,  $t_{si}$  should be less than 10 nm.

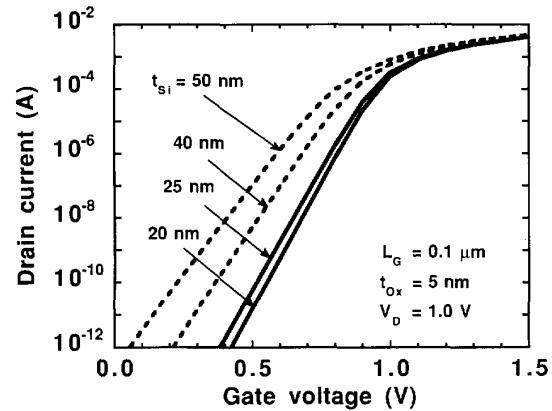


Fig. 5. Subthreshold characteristics with various silicon thickness. Solid lines correspond to the devices adhering to the scaling rule and dashed lines to those not doing so.

Fig. 5 shows the subthreshold characteristics of the devices for  $L_G = 0.1 \mu\text{m}$ , where solid lines correspond to  $\alpha_2 > 3$  and dashed lines to  $\alpha_2 < 3$  and  $V_D = 1.0$  V which may be the supply voltage for  $L_G = 0.1 \mu\text{m}$ . We can expect almost an ideal  $S$ -factor with the device if we design the devices so that  $\alpha_2$  is more than 3. We also verified the same numerical results with the device for  $L_G = 0.05 \mu\text{m}$ .

### IV. SUMMARY

We derived a natural length relevant to the scaling theory for double-gate SOI MOSFET's, and described how to design  $t_{ox}$  and  $t_{si}$  for a given gate length maintaining a proper  $S$ -factor. According to the theory, almost the ideal  $S$ -factor value can be expected even with  $L_G$  of less than  $0.1 \mu\text{m}$ . The key factor that determines the shortest gate length device we can design is the thinnest  $t_{si}$  we can produce. A  $t_{si}$  of less than 10 nm is needed at  $L_G = 0.05 \mu\text{m}$  with  $t_{ox} = 3$  nm.

### ACKNOWLEDGMENT

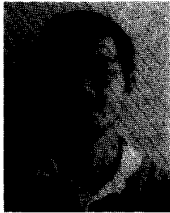
We thank S. Satoh and H. Tashiro for helping with the numerical calculations and N. Nakayama and Dr. T. Ito for their encouragement.

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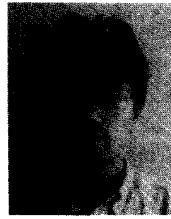
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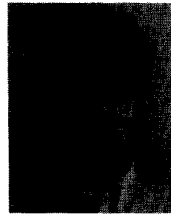
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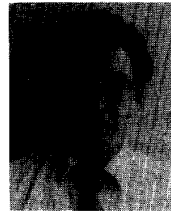
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