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Abstract

A transducer swing mechanism is introduced to accomplish a new scanning acoustic microscope system in the interference mode, which is applied to the quantitative measurement of acoustic velocities in thin materials. In this system, the distance between a pair of aligned focusing transducers along the beam axis is varied periodically, as the specimen is synchronously scanned in the X direction. Measurements of acoustic velocities are performed with this system for two kinds of polymer films, polyester and PVF₂ films, and a dog's cardiac infarcted tissue at a frequency of 150 MHz.

1. Introduction

The quantitative measurement of the acoustic properties of materials in microscopic scale is very important in the fields of polymer science, biological science, medical ultrasonics, nondestructive evaluation, etc.. An acoustic microscope is perhaps an only tool which is available for these purposes. Recently, a scanning laser acoustic microscope (SLAM) has been operated in the interference mode to detect the acoustic velocities in small area of the specimen, and the quantitative measurements of the acoustic velocities in biological tissues have been demonstrated by Goss et al. [1],[2]. In the scanning acoustic microscope (SAM) system [3], the combination system of plane-wave and focusing-wave transducers has been also developed to accomplish the interference mode operation by tilting the plane-wave transducer, so that the acoustic velocity distributions can be measured in the similar way [4], [5]. In the acoustic microscope system, a coherent plane-wave illumination is not always accepted favorably due to the following problems: (i) the uniform plane wave cannot easily generated, (#) the effect of diffraction influences on the received signals much more than the confocal SAM system, (iii) the radiation power of the transmitted transducer cannot be effectively used.

In this paper, we describe a mechanism, newly introduced into a conventional scanning acoustic microscope, which vibrates the focusing transducer along the beam axis to vary periodically the distance between a pair of aligned focusing transducers. This mechanism and its operation are applicable to both the transmission mode and reflection mode. If used in the transmission mode, an interference microscope of a new type becomes available without the plane-wave illumination. The construction and the performance of the new SAM for interference-mode operation in a frequency range around 150 MHz are described. Experiments are performed with this SAM at 150 MHz to measure acoustic velocities for two kinds of polymer films, a polyester film and a PVF₂ film, and a dog's cardiac infarcted tissue.

2. Operation Principle

Figure 1 shows the configuration of acoustical parts in the new interference microscope. A pair of focusing transducers are confocally aligned along the beam axis and a specimen is placed around the focal plane in a coupling liquid. To produce the interferogram in this system, the distance between a transmitting and receiving transducers is varied periodically, and the specimen is syncronously scanned in the X direction. The output signal of the receiving transducer is electrically mixed with the reference signal and then interferogram can be shown on the CRT display by using the conventional electronic techniques, as the specimen is slowly displaced in the Y direction.



Fig. 1. Basic configuration of the new interference microscope.

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Figure 2 schematically depicts the relative positions of the receiving transducer to the specimen and their relation in the interference pattern displayed on the CRT. The spacing of the fringe lines depends on the acoustic velocity of the coupling liquid, the acoustic frequency and the ratio B/A, where A is the Z-vibrating width and B is the X-scanning width. When the specimen is placed on a propagation path of an acoustic beam, the fringe lines will shift to the right or left depending on the acoustic velocity of the specimen. The fringe lines in Fig.2 show the shift toward right for the specimen with a velocity higher than that of coupling liquid. The spacing of the fringe lines D is given as follows:

$$D = (B/A)\lambda_{7} = (B/A) \cdot (v_{7}/f), \qquad (1)$$

where λ_{χ} and v_{χ} are the acoustic wavelength and the acoustic velocity, respectively, in the coupling liquid, and f is the acoustic frequency.

The fringe shift d depends on the thickness h of a specimen, the ratio B/A and the acoustic velocity ratio v_l/v_g , where v_g is the acoustic velocity in the specimen. That is, d is given by follows:

$$d=h(B/A)\cdot(1-v_l/v_s), \qquad (2)$$

From the equations (1) and (2), the acoustic velocity $v_{\rm g}$ in the specimen is,

$$v_{s} = v_{l} / (1 - d\lambda_{l} / Dh) = v_{l} / (1 - N\lambda_{l} / h), \qquad (3)$$

where N is the normalized lateral fringe shift that is defined as the ratio of the fringe shift d and the fringe spacing D. Therefore, if the acoustic velocity of the coupling liquid and the thickness of specimen are known and the fringe shift is measured on the CRT as explained in Fig.2, the acoustic velocity in the specimen v_g can be calculated by the equation (3).

3. Construction and Performance of the System

Figure 3 shows the block diagram of our system of interference microscope. In the system, the differentiator is a newly introduced electrical circuit for differentiating the video signals to obtain the clear fringe lines that make the measurement of the fringe shift easy. The phase detector is composed of a limiter amplifier and a double balanced mixer, and the accuracy of phase detector is measured to be better than 1 %, when the input dynamic range is kept within 30 dB.



Fig. 3. Block diagram of the new acoustic interference microscope system.

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Fig. 4. Photograph of the mechanical part of the system.

Figure 4 is a photograph of the mechanical part of the system. A loudspeaker is employed to scan the specimen in the X direction, and a commercially obtained vibrator is employed to vibrate the receiving transducer in the Z direction at a frequency of 50 Hz, where the displacement of the vibrator is detected by an accelerometer. The driving signals for them are supplied in common from an AF oscillator, and the phase of the Z-vibrating signal is adjusted by the variable phase shifter, to vibrate the receiving transducer syncronously in the Z direction as the specimen is scanned in the X direction. The spacing of interference fringe lines can be varied by changing the vibrating width in the Z direction.



Fig. 5. Frequency response of insertion loss at the confocal arrangement (including the loss of water ; about 15 dB at 150 MHz).

The focusing transducers employ ZnO pizoelectric films formed on the flat end of acoustic sapphire lenses, of which the diameter is 1 mm and the center frequency is about 150 MHz, and SiO₂ films with the thickness of 9.96 μ m are deposited onto the concave surface of the lenses as an acoustic antireflection coating (AARC) [6]. The radius of curvature of the lens is 1.5 mm and the half aperture angle of Θ =20° is determined to be smaller than that of the conventional SAM system, so that the longer depth of focus can be obtained. This could be possible to get a wide swing range in the Z direction.

Figure 5 shows a frequency response of the insertion loss between the sending and receiving transducers including the loss of water as a coupling liquid, when both the transducers are The minimum insertion loss of arranged confocally. 47 dB is obtained around 150 MHz. To examine the allowable swing range in the Z direction, the output signal of the receiving transducer is measured while one of two transducers is relatively moved off along the beam axis from the confocally arranged position at a frequency of 150 MHz. In Fig.6, the variation of the output signal is plotted with respect to the Z axis. The -3 dB width of the curve is observed to be 350 µm along the Z axis. So, our system can take at least the swing range of 350 µm for the vibration of the receiving transducer to obtain the clear interference patterns.



Fig. 6. Variation of receiving output with respect to Z axis, when the receiving transducer was moved from the confocally arranged position at 150 MHz,

Next, we have observed the received output signal when a stainless steel wire of $35 \ \mu m$ in diameter has been scanned in the X direction on the focal plane in water, in order to examine an influence to the acoustic wave diffraction. The variation of the output voltage along the X axis is shown in Fig.7 (a). This is compared to that obtained in the combination system of plane-wave and focusingwave transducers shown in Fig.7 (b). It can be easily seen that the diffraction effect is almost removed in the case of confocal system.

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Fig. 7. Variation of receiving output voltage along X axis, when a stainless steel wire of 35 µm in diameter has been scanned in the X direction on the focal plane in water (a), and that obtained by the combination system of plane-wave and focusing-wave transducers for comparison (b).

4. Experiments and Discussion

Various experiments for acoustic velocity measurements have been performed with some polymer films and biological tissues at a frequency of 150 MHz in water as a coupling liquid.

The interferograms observed for two kinds of polymer films are shown in Fig.8. They show the interferograms observed with a polyester film of 38 µm in thickness at the temperature of 24 °C in (a) and with a polyvinylidene fluoride (PVF₂) film of 30 μm in thickness at the temperature of 26 °C in (b), respectively. The fringe lines shift toward right in both figures at the water-specimen boundary, since the specimens have velocities higher than the velocity of water. The average acoustic velocity of the polyester film is calculated to be 2310 m/s and that of PVF2 film is calculated to be 2280 m/s by using Eq.(3) as the normalized fringe shifts N being 1.35 and 1.03, respectively. In the calculation, the data of velocity of the coupling water reported by Greenspan et al. in Ref.7 are used.

As an example of biological tissue, an interferogram observed at the temperature of 23 °C with a dog's cardiac infarcted tissue of 50 μ m in average thickness is shown in Fig.9 (a) and its optical image is shown in Fig.9 (b). The specimen has been fixed by formalin at first, embeded by paraffin to slice and the experiment is carried out after removing the paraffin completely. The average fringe shift observed at the boundary between water and the tissue is measured to be 0.56 although the fringe shift is slightly meandered depending on the tissue structure. So, the average velocity is determined to be 1680 m/s.



Fig. 8. Acoustic interferograms obtained at 150 MHz for polyester film of 38 µm thickness at 24 °C (a) and PVF₂ film of 30 µm thickness at 26 °C (b).



(a) INTERFEROGRAM



(b) OPTICAL IMAGE

Fig. 9. Acoustic interferogram obtained at 150 MHz for dog's cardiac infarcted tissue of 50 µm thickness at 23 °C (a) and its optical image (b).

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5. Conclusions

A new acoustic interference microscope has been proposed in this paper to measure acoustic velocities of thin materials quantitatively, in which one of the aligned focusing transducer is vibrated periodically along the beam axis. The construction and the performance of the SAM have been described. It has been shown that diffraction effects observed in the combination system of plane and focusing transducers were preferably eliminated by using the focusing transducer system, presented here. Experiments of velocity measurement have been further demonstrated with specimens such as polymer films and a biological tissue at a frequency of 150 MHz. This interference microscope system composed of focusing transducers is expected to be used widely as a very useful tool for tissue characterization, material characterization and nondestructive evaluation, because the system can be easily constructed by a slight modification of the conventional SAM system,

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