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著者	齊藤 伸
journal or	Journal of applied physics
publication title	
volume	93
number	10
page range	8182-8184
year	2003
URL	http://hdl.handle.net/10097/35355

doi: 10.1063/1.1558093

## Improvement of magnetocrystalline anisotropy of columnar structure for epitaxially grown CoCr-based perpendicular media

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(Presented on 14 November 2002)

To improve uniaxial magnetic anisotropy  $K_u^{\text{grain}}$  of *c*-plane-oriented columnar grains for CoCr-based perpendicular recording media, an epitaxial growth technique has been applied. The main results are as follows. (1) Using a Ru(3 nm) or  $\text{Co}_{60}\text{Cr}_{40}(20 \text{ nm})/\text{C}(1 \text{ nm})$  intermediate layer,  $K_u^{\text{grain}}$  for a  $\text{Co}_{69}\text{Cr}_{19}\text{Pt}_8\text{B}_4$  medium was enhanced from  $2.04 \times 10^6 \text{ erg/cm}^3$  to  $2.53 \times 10^6 \text{ erg/cm}^3$  or to  $2.76 \times 10^6 \text{ erg/cm}^3$ , respectively. (2) The enhancement of  $K_u^{\text{grain}}$  depends on the composition of the magnetic layer when using the  $\text{Co}_{60}\text{Cr}_{40}/\text{C}$  intermediate layer. (3) The epitaxial growth technique is also effective to increase  $K_u^{\text{grain}}V_{\text{act}}/kT$ . Especially, in the media with a high concentration of B or Cr, the thermal stability is improved mainly due to an increase of  $K_u^{\text{grain}}$  rather than an enlargement of  $V_{\text{act}}$ . (2) 2003 American Institute of Physics. [DOI: 10.1063/1.1558093]

To improve the thermal stability of high density recording for CoCr-based perpendicular media, it is essential to use materials with high uniaxial magnetocrystalline anisotropy energy,  $K_u^{\text{grain}}$ . We have reported that epitaxial growth of a  $\text{Co}_{69}\text{Cr}_{19}\text{Pt}_8\text{B}_4$  magnetic layer by utilizing a  $\text{Co}_{60}\text{Cr}_{40}/\text{C}$  intermediate layer is effective to improve  $K_u^{\text{grain}}$  of *c*-planeoriented columnar grains and eliminate the formation of an initial growth layer. In this study, the influence of magnetic layer and intermediate layer materials on  $K_u^{\text{grain}}$  of epitaxially grown media is further investigated.

CoCrPt and CoCrPtB perpendicular thin film media were fabricated by a dc magnetron sputtering method on 65 mm diam glass disk substrates using the so-called ultraclean sputtering system.<sup>1</sup> The substrate temperature was varied from 150 to 350 °C and was optimized such that the perpendicular coercivity of the media takes maximum. The film thickness of the media was varied from 5 to 50 nm. The underlayer and the protective layer were 25-nm-thick Ti, and 7-nm-thick C.

The crystal structure of the media was examined by the grazing incident angle x-ray diffraction (XRD) method (inplane XRD;  $2\theta_{\chi}$  scan) using Cu  $K\alpha$  radiation.<sup>2</sup> The incident angle was fixed at 0.4°, which corresponds to about a 20-nmthick penetration of the x ray from the incident surface. The saturation magnetization  $M_s$  and perpendicular magnetic anisotropy  $K_{u\perp}^{\exp}$  were evaluated by a vibrating sample magnetometer (VSM) and high sensitive torque magnetometer, respectively. The uniaxial magnetocrystalline anisotropy of columnar grains,  $K_u^{\text{grain}}$ , and the initial layer thickness,  $d_{\text{ini}}$ , were derived by perpendicular magnetic torque curve and determined from the  $K_{u\perp}^{\exp} \times d_{\text{mag}}$  versus  $d_{\text{mag}}$  plot.<sup>3,4</sup> Activation volume  $V_{\text{act}}$ , which corresponds with the coherent magnetic reversal volume in a medium, was determined by irreversible susceptibility measurement using the following equations:<sup>5</sup>  $V_{act} = (kT)/(M_sH_f)$ ,  $H_f = -[dH_r(t)/d \ln(t)]$ , where  $H_f$  is the fluctuation field, *T* is the absolute temperature, and *k* is the Boltzmann's constant.  $H_f$  was determined by linear extrapolation of the  $H_r$  versus  $\ln(t)$  plot.

To realize the epitaxial growth of the magnetic layer, the intermediate layer, which consists of *c*-plane-oriented hexagonal grains, Ru or  $\text{Co}_{60}\text{Cr}_{40}/\text{C}(1 \text{ nm})$ , was selected. Here, a 1 nm carbon seedlayer underlying the  $\text{Co}_{60}\text{Cr}_{40}$  intermediate layer plays a very important role to form Cr-deprived hcp grains with Cr segregation structure in the  $\text{Co}_{60}\text{Cr}_{40}$  layer without forming  $\sigma$ -phase grains.<sup>6,7</sup>

Figure 1 (left) shows the dependence of perpendicular magnetic properties  $H_c$  and S for 30-nm-thick  $\text{Co}_{69}\text{Cr}_{19}\text{Pt}_8\text{B}_4$  media on intermediate layer thickness,  $d_{\text{inter}}$ , of (A) Ru and

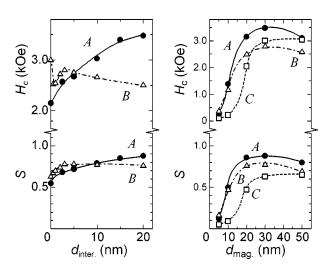


FIG. 1. Left: dependence of perpendicular magnetic properties  $H_c$  and S for 30-nm-thick Co<sub>69</sub>Cr<sub>19</sub>Pt<sub>8</sub>B<sub>4</sub> media on intermediate layer thickness,  $d_{\text{inter}}$ , of (A) Ru, and (B) Co<sub>60</sub>Cr<sub>40</sub>/C(1 nm), respectively. Right: dependence of  $H_c$  and S on magnetic film thickness,  $d_{\text{mag}}$ , for Co<sub>69</sub>Cr<sub>19</sub>Pt<sub>8</sub>B<sub>4</sub> media (A) with Ru(3 nm), (B) with Co<sub>60</sub>Cr<sub>40</sub>(20 nm)/C(1 nm), and (C) without an intermediate layer, respectively.

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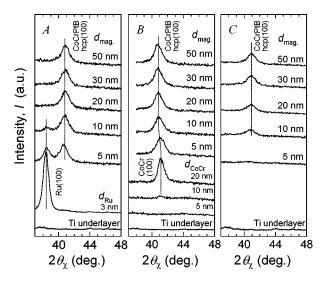


FIG. 2. In-plane XRD profiles for  $Co_{69}Cr_{19}Pt_8B_4$  media with various magnetic film thickness,  $d_{mag}$ , (A) with Ru (3 nm), (B) with a  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer, and (C) without an intermediate layer, respectively. In-plane XRD profiles for Ru films deposited on a Ti underlayer (in the left figure),  $Co_{60}Cr_{40}/C$  film deposited on a Ti underlayer (in the middle figure), and Ti film (in the right figure) are also shown.

(B)  $\operatorname{Co}_{60}\operatorname{Cr}_{40}/\operatorname{C}(1 \text{ nm})$ , respectively. For the medium A with an Ru intermediate layer, both  $H_c$  and S show the maximum values of 2.8 kOe and 0.78 at  $d_{\text{inter}}=3 \text{ nm}$ , respectively. On the other hand for the medium B with a  $\operatorname{Co}_{60}\operatorname{Cr}_{40}/\operatorname{C}$  intermediate layer, both  $H_c$  and S show the maximum values of 3.5 kOe and 0.88 at  $d_{\text{inter}}=20 \text{ nm}$ , respectively.

Figure 1 (right) shows the dependence of perpendicular magnetic properties  $H_c$  and S on magnetic film thickness,  $d_{\text{mag}}$ , for  $\text{Co}_{69}\text{Cr}_{19}\text{Pt}_8\text{B}_4$  media with an (A) Ru(3 nm) and (B)  $\text{Co}_{60}\text{Cr}_{40}(20 \text{ nm})/\text{C}(1 \text{ nm})$  intermediate layer, respectively. In the same figure,  $H_c$  and S for the medium (C) without an intermediate layer are also shown for comparison. For the media C with  $d_{\text{mag}} \leq 10 \text{ nm}$ ,  $H_c$  and S show low values of 0.1 kOe and 0.09, respectively. On the other hand, for the media A and B,  $H_c$  and S show relatively high values of about 1.4 kOe and 0.50, respectively, even when  $d_{\text{mag}} = 10 \text{ nm}$ . Especially for the medium B, the maximum values of  $H_c$  and S are 3.5 kOe and 0.88 at  $d_{\text{mag}} = 30 \text{ nm}$ , which are higher than those of the media A.

Figure 2 shows in-plane XRD profiles for  $Co_{69}Cr_{19}Pt_8B_4$ media with various magnetic film thickness,  $d_{mag}$ , (A) with Ru (3 nm), (B) with  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer, and (C) without intermediate layer, respectively. Inplane XRD profiles for Ru film deposited on a Ti underlayer

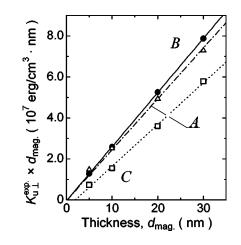


FIG. 3.  $K_{\mu\perp}^{exp} \times d_{mag}$  vs  $d_{mag}$  plots for  $Co_{69}Cr_{19}Pt_8B_4$  media (A) with Ru (3 nm) and (B) with a  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer, and (C) without an intermediate layer, respectively.

(in the left figure),  $Co_{60}Cr_{40}/C$  film deposited on a Ti underlayer (in the middle figure), and Ti film (in the right figure) are also shown. In the profiles for a Ru (3 nm) film and a  $Co_{60}Cr_{40}$  (20 nm) film, a diffracted line from the (100) plane from the hexagonal structure can be clearly seen. Therefore it was clarified that c-plane-oriented Ru grains or CoCr grains grow in the intermediate layer. Taking a look at the XRD patterns of the magnetic layer, it is found that for the medium C with  $d_{\text{mag}} = 5$  nm, a diffracted line can be hardly observed, and with increase  $d_{\text{mag}} \ge 10 \text{ nm}$ , a diffracted line from the (100) plane becomes observable. This fact indicates that for the medium C, c-plane-oriented columnar grains grow on the initial layer, which consist of nanocrystalline grains with random orientation. On the contrary, for the media A and B with  $d_{\text{mag}} = 5 \text{ nm}$ , a diffracted line from the (100) plane can be clearly observed. Therefore it is suggested that for the media A and B, c-plane-oriented columnar grains in the magnetic film grow from the initial growth stage.

Figure 3 shows a  $K_{u\perp}^{\exp} \times d_{mag}$  versus  $d_{mag}$  plot for  $Co_{69}Cr_{19}Pt_8B_4$  media (A) with Ru (3 nm), (B) with a  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer, and (C) without an intermediate layer, respectively. Judging from the intersection of the extended line of the linear portion of the  $K_{u\perp}^{\exp} \times d_{mag}$  curve with the  $d_{mag}$  axis, the medium C has the initial layer with the thickness of 2.0 nm. On the contrary, for the media A and B, the initial layer is completely eliminated. Based on the analyses of in-plane XRD in Fig. 2 and the  $K_{u\perp}^{\exp} \times d_{mag}$  plot in Fig. 3, the media A and B can be concluded to grow epitaxially. Furthermore, according to the

TABLE I.  $d_{\text{ini}}$ ,  $K_u^{\text{grain}}$ ,  $M_s$ ,  $H_c/H_k$ , and  $\alpha$  for Co<sub>69</sub>Cr<sub>19</sub>Pt<sub>8</sub>B<sub>4</sub> medium.  $H_c/H_k$  and  $\alpha$  are for the media with  $d_{\text{mag}} = 30$  nm.

	Layer structure	d <sub>ini</sub> (nm)	$K_u^{\text{grain}}$ (erg/cm <sup>3</sup> )	$M_s$ (emu/cm <sup>3</sup> )	$\frac{H_c/H_k}{(30 \text{ nm})}$	α (30 nm)
А	Co <sub>69</sub> Cr <sub>19</sub> Pt <sub>8</sub> B <sub>4</sub> /Ru (3 nm) /Ti (25 nm)	0.2	2.53×10 <sup>6</sup>	488	0.27	1.45
В	$Co_{69}Cr_{19}Pt_8B_4/Co_{60}Cr_{40}$ (20 nm) /C(1 mm)/Ti (25 nm)	0	$2.76 \times 10^{6}$	477	0.30	1.35
С	Co <sub>69</sub> Cr <sub>19</sub> Pt <sub>8</sub> B <sub>4</sub> /Ti (25 nm)	2.0	$2.04 \times 10^{6}$	433	0.32	1.15

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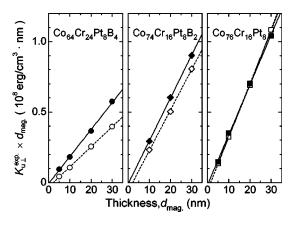
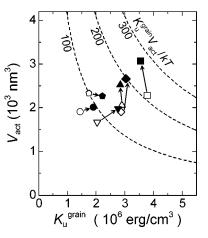


FIG. 4.  $K_{u\perp}^{exp} \times d_{mag}$  vs  $d_{mag}$  plots for the media with and without the  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer for (left)  $Co_{64}Cr_{24}Pt_8B_4$ , (center)  $Co_{74}Cr_{16}Pt_8B_2$ , and (right)  $Co_{76}Cr_{16}Pt_8$  media are shown in separate graphs. A closed or an open symbol corresponds to the medium with or without the intermediate layer, respectively.

evaluation of the gradient of the linear portion of  $K_{u\perp}^{exp}$ × $d_{mag}$  curves, the values of  $K_u^{grain}$  for the media A (2.53 × 10<sup>6</sup> erg/cm<sup>3</sup>) and B (2.76×10<sup>6</sup> erg/cm<sup>3</sup>) are remarkably larger than that for the medium C (2.04×10<sup>6</sup> erg/cm<sup>3</sup>). Therefore for the Co<sub>69</sub>Cr<sub>19</sub>Pt<sub>8</sub>B<sub>4</sub> medium, the epitaxial growth using *c*-plane-oriented Ru and a Co<sub>60</sub>Cr<sub>40</sub>/C intermediate layer enables one to increase  $K_u^{grain}$  by 24%–35%. The evaluated properties for the media A, B, and C are summarized in Table I. Comparing the medium A with B,  $H_c/H_k^{grain}$ for the medium B is larger and  $\alpha$  is smaller than those for the medium A. This fact means that magnetic grains in the medium B is well isolated compared with that in the medium A. Hence it is suggested that the promotion of Cr segregation to the grain boundary in epitaxially grown media depends on the material of the intermediate layer.

We have further applied the epitaxial growth technique utilizing the  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer to enhance the  $K_{\mu}^{\text{grain}}$  of several perpendicular media. In Fig. 4,  $K_{u\perp}^{\exp} \times d_{\max}$  versus  $d_{\max}$  plots are shown with and without the  $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$ intermediate layer for  $Co_{64}Cr_{24}Pt_8B_4$ ,  $Co_{74}Cr_{16}Pt_8B_2$ , and  $Co_{76}Cr_{16}Pt_8$  media, separately. A closed or an open symbol corresponds to the medium with or without the intermediate layer, respectively. In the medium without a intermediate layer, a several-nmthick initial layer exists, on the other hand for the medium using the intermediate layer, the initial layer is confirmed to be almost eliminated, hence the media with the intermediate layers are suggested to grow epitaxially. However, judging from the gradient of the linear portion of  $K_{u\perp}^{exp} \times d_{mag}$  curves, the enhancement of  $K_u^{\text{grain}}$  depends on the composition of the magnetic layer.

Finally, we will discuss the thermal stability. In Fig. 5, the change in  $K_u^{\text{grain}}$  and  $V_{\text{act}}$  by epitaxial growth using a



Relationship between activation volume  $V_{act}$  and uniaxial magnetocrystalline anisotropy of the columnar structure,  $K_{\mu}^{grain}$  for 30-nm-thick media. In the figure, equivalue lines of thermal stability factor,  $K_{\mu}^{grain} V_{act}/kT$ , are also drawn as a reference. An open or closed symbol corresponds to the medium with or without the  $Co_{60}C_{20}(20 \text{ nm})/C$  (1 nm) intermediate layer, respectively, and (•), (•), (•), (•), (•), (•), (a), and (•) are represented properties for  $Co_{64}Cr_{24}Pt_8B_4$ ,  $Co_{60}Cr_{20}Pt_1_6B_4$ ,  $Co_{69}Cr_{19}Pt_8B_4$ ,  $Co_{74}Cr_{16}Pt_8B_2$ ,  $Co_{64}Cr_{20}Pt_{16}$ , and  $Co_{76}Cr_{16}Pt_8$  media, respectively. In the figure, equivalue lines of the thermal stability factor,  $K_{\mu}^{grain} V_{act}/kT$ , are also drawn.

 $Co_{60}Cr_{40}(20 \text{ nm})/C(1 \text{ nm})$  intermediate layer is shown for CoCrPt(B) media with various compositions. A closed or open symbol corresponds to the medium with or without the intermediate layer, respectively, while  $(\bullet)$ ,  $(\bullet)$ ,  $(\bullet)$ ,  $(\bullet)$ ,  $(\bullet)$ ,  $(\bullet)$ , and  $(\blacksquare)$  represent the properties for  $Co_{64}Cr_{24}Pt_8B_4$ ,  $Co_{60}Cr_{20}Pt_{16}B_4$ ,  $Co_{69}Cr_{19}Pt_8B_4$ ,  $Co_{74}Cr_{16}Pt_8B_2$ ,  $Co_{64}Cr_{20}Pt_{16}$ , and  $Co_{76}Cr_{16}Pt_8$  media, respectively. In the figure, equivalue lines of the thermal stability factor,  $K_u^{\text{grain}} V_{\text{act}} / kT$ , are also drawn as a reference. For all the media, utilizing an epitaxial growth technique is effective to increase  $K_u^{\text{grain}} V_{\text{act}} / kT$ . In the case of  $Co_{64}Cr_{24}Pt_8B_4$ ,  $Co_{60}Cr_{20}Pt_{16}B_4$ , and  $Co_{69}Cr_{19}Pt_8B_4$  media,  $K_u^{\text{grain}}$  substantially improves while an only a slight change of V<sub>act</sub> was observed. On the contrary, in the case of  $\text{Co}_{74}\text{Cr}_{16}\text{Pt}_8\text{B}_2$ ,  $\text{Co}_{64}\text{Cr}_{20}\text{Pt}_{16}$ , and  $\text{Co}_{76}\text{Cr}_{16}\text{Pt}_8$  media,  $K_u^{\text{grain}}$ changes slightly and  $V_{act}$  drastically increases. Therefore in the media with a high concentration of B or Cr, the thermal stability is improved mainly due to an increase of  $K_u^{\text{grain}}$ rather than an enlargement of  $V_{act}$ . To clarify the origin of this phenomenon from the viewpoint of the microstructure, further study including quantity analysis concerning the existence of stacking fault and composition of the columnar grains is necessary.

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