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Study of $(e, e' \alpha)$ Reaction on ${}^9\text{Be}$

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The $(e, e' \alpha)$ cross section has been measured at energy transfers from 10.0 to 28.4 MeV and a momentum transfer of 99 MeV/c, using a 197 MeV continuous electron beam. The cross section rapidly increases with decreasing energy at angles smaller than 25° , while it appears flat at larger angles. The forward-peaked angular distribution was observed below 18 MeV, and the forward peak shrinks at higher energies. An amount of α particles from decay of ${}^5\text{He}$ is estimated.

§ 1. Introduction

There has been a growing interest in the study of neutron-halo and neutron-skin structure. Very recently, the RI Beam Factory project has started at RIKEN to explore this subject. The nucleus ${}^{11}\text{Be}$ is known as a neutron-halo nucleus [1], and there are several microscopic and macroscopic attempts at describing Be isotopes in a unified framework of two α particles and extra neutrons. The ${}^9\text{Be}$ nucleus is the simplest isotope, which has one extra neutron, and is known to have a typical $\alpha + \alpha + n$ cluster structure. Therefore the nucleus has already been investigated in various models.

The photo-disintegration of ${}^9\text{Be}$ is categorized in three separate mechanisms: (1) threshold to 5 MeV, where sharp resonances correspond to direct excitations of the unpaired neutron; (2) 5-18 MeV, where weak coupling of the unbound neutron to ${}^8\text{Be}$ dominates; (3) above 18 MeV, where a core is excited. The photo-nuclear reaction has been investigated in the (γ, n) [2-5], (γ, p) [6], (γ, d) [6], (γ, t) [6] and $(\gamma, {}^3\text{He})$ [6] channels. However, no (γ, α) cross section has been measured, presumably because the decayed α particles of ${}^5\text{He}$ that partly overlap on the spectrum make the analysis difficult. The threshold of the (γ, α) reaction is 2.53 MeV; this is the only charged-particle emission channel below 16 MeV. The cross section is expected to be large enough for measurement.

The residual nucleus ${}^5\text{He}$ of the ${}^9\text{Be}$ $(e, e' \alpha)$ reaction is same as the ${}^6\text{Li}$ $(e, e' p)$ reaction, which we investigated several years ago [7]. However, the missing energy spectra of both reactions may be different from each other, because they reflect their reaction mechanisms. We measured ${}^9\text{Be}$ $(e, e' \alpha)$ cross section at eleven angles at transferred energies between 10.0 and 28.4 MeV as a test experiment.

§ 2. Experimental procedure

The experiment was performed using a 197 MeV continuous electron beam from a stretcher-

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booster ring (STB). Electrons scattered with a 1.8 mg/cm^2 thick natural beryllium foil were analyzed with a magnetic spectrometer (LDM) at 30° . Missing energies were set at two energies: phase 1 ($10.0 \sim 21.0 \text{ MeV}$), phase 2 ($17.6 \sim 28.4 \text{ MeV}$). Corresponding momentum transfers are 98.60 MeV/c for phase 1 and 99.10 MeV/c for phase 2. Ejected α particles were measured with eleven SSD counter telescopes composed of two surface-barrier type SSD's (one $50 \mu\text{m}$ SSD + one 1 mm SSD at eight angles, and two 1 mm SSD's at three angles), out-of the scattering plane ($\phi_\alpha = 90^\circ$). The solid angle of the telescopes is 4.8 msr . Due to insufficient machine time, no measurement between 90° and 180° was made in this run.

§ 3. Results and Discussion

Figure 1 shows missing energy spectra at the momentum transfer direction; a peak at 2.5 MeV corresponds to α particle emission leaving the residual nucleus at the ground state of ^5He . A tail toward higher missing energies is observed in the spectrum for the lower energy transfer region; more structures can be seen in the spectrum for the higher energy transfer region. Considering a resolution of the experiment, events below $E_m = 7 \text{ MeV}$ were treated as α_0 events.

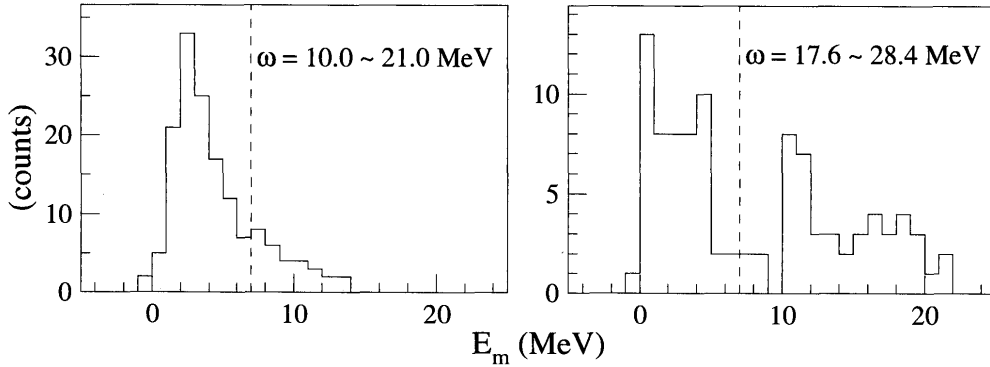


Fig.1. Missing energy spectra at 0° . A peak at 2.5 MeV corresponds to $(e, e' \alpha_0)$ reaction.

A transferred energy dependence of the differential $(e, e' \alpha_0)$ cross sections is shown in Fig.2. At forward angles, the cross section becomes larger as the energy decreases; while no such rise can be seen at backward angles. angular dependence of the differential cross section is shown in Fig.3. At low energies, a strong forward peaking is observed; data at backward angles are lacking because energies of the α particles are lower than the threshold energy of the detectors there.

In the $(e, e' \alpha)$ ^5He reaction, the residual nucleus ^5He decays into $n + \alpha$; this α particles partly overlap on the α particles of the $(e, e' \alpha)$ ^5He reaction. Simulated missing energy spectra of both α particles are demonstrated in Fig.4. In this calculation several assumptions were made:

- (1) The angular distribution of the ^9Be $(e, e' \alpha)$ ^5He reaction has a forward-backward symmetry in the center of mass system; Legendre parameters has been obtained by fitting the measured angular

distributions with even-order Legendre functions up to $l = 4$.

- (2) The angular distribution of the α particle due to $\alpha + \alpha + n$ decay of ${}^5\text{He}$ is isotropic in the center of mass system.

The lowest peak in Fig.4 corresponds to α particles from the ${}^9\text{Be}(e, e' \alpha) {}^5\text{He}$ reaction; the higher missing-energy part of two-peaked structure corresponds to decay products of ${}^5\text{He}$. Alpha-particles emitted toward the detector direction make the lower peak, while those to the opposite direction do the higher peak. The structure changes slowly at other angles. The two-peaked structure shifts toward higher missing-energies as the transferred energy ω increases. The missing energy spectra folded in the regime of the measurements are shown in Fig.5, where the structure is smeared in resolution (1 MeV FWHM). A tail in higher energy part for $\omega = 10 \sim 21$ MeV in Fig.1 is reproduced in this

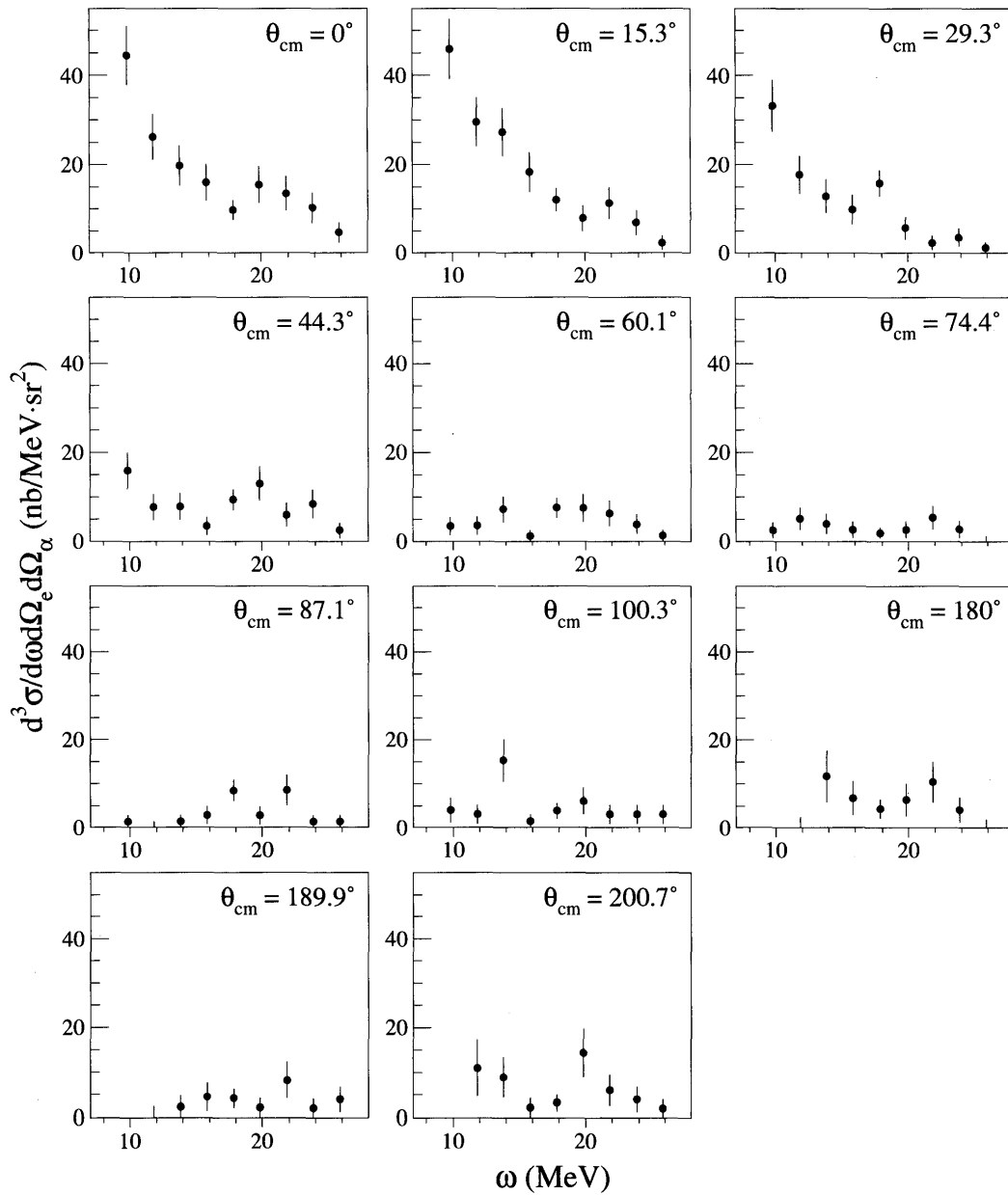


Fig.2. The ${}^9\text{Be}(e, e' \alpha)$ differential cross sections.

simulation. In case $\omega = 17.6 \sim 28.4$ MeV, the part $E_m > 10$ MeV cannot be reproduced in the simulation although the measured spectrum has large statistical errors. It may be a contribution of α particles from the ${}^9\text{Be}(e, e'n){}^8\text{Be}(16.6 \text{ MeV}) \rightarrow 2\alpha$.

We wish to thank the accelerator group and the computer group for their assistance during the measurements.

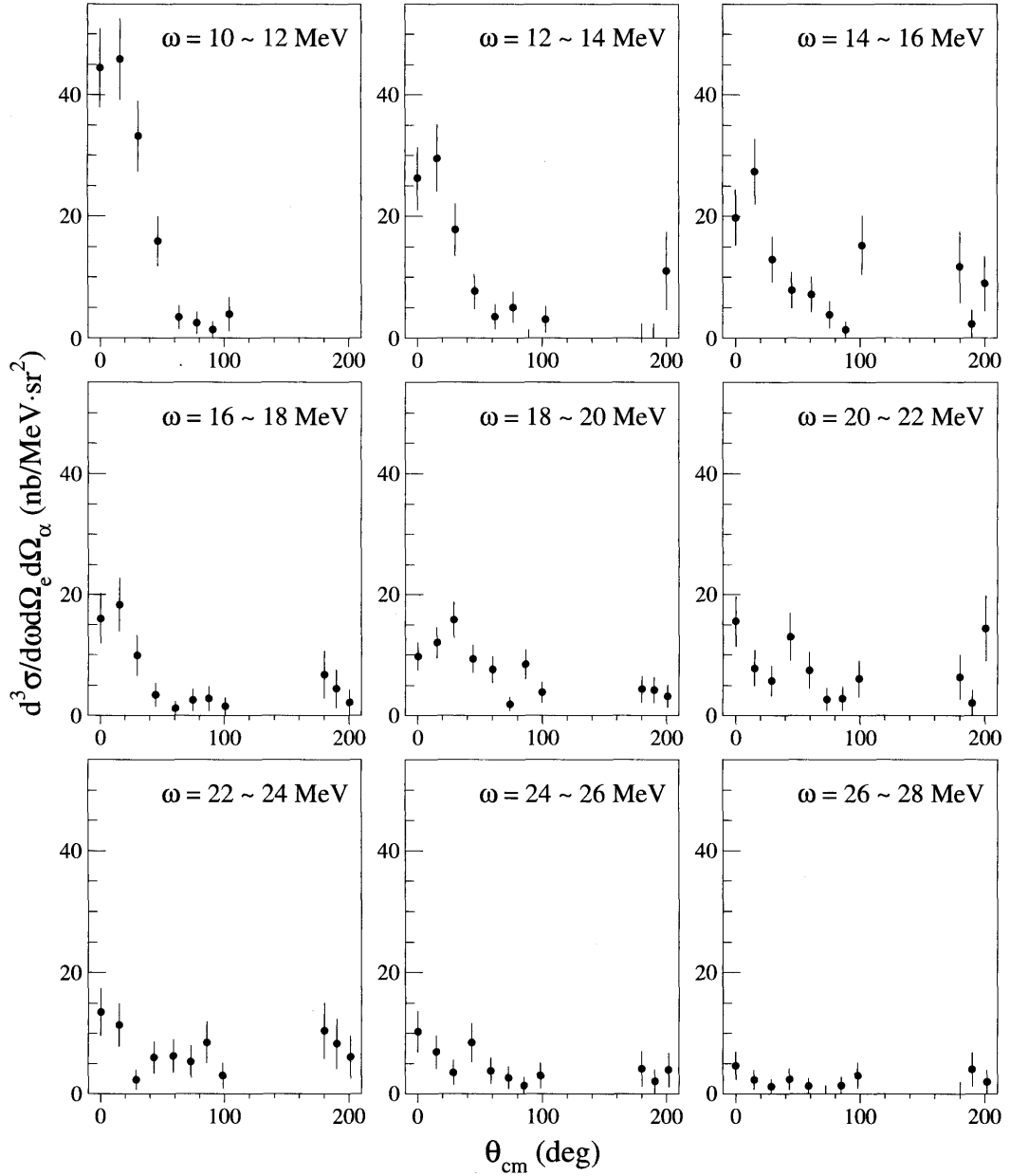


Fig.3. Angular distributions of the ${}^9\text{Be}(e, e' \alpha_0)$ reaction.

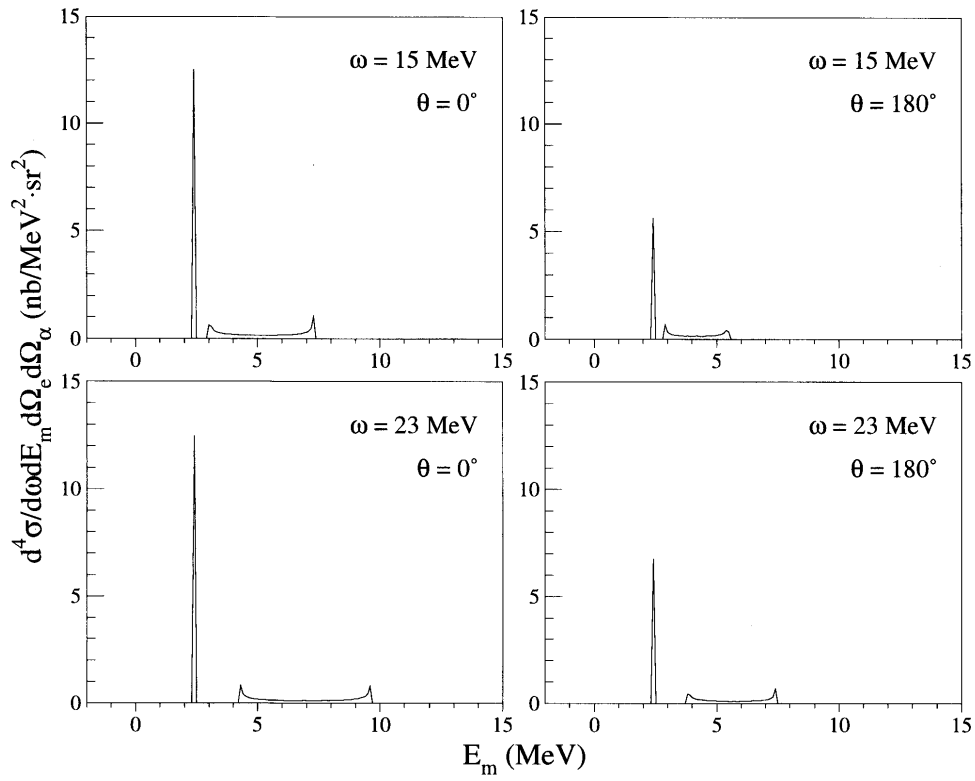


Fig.4. Simulated missing energy spectra of the ${}^9\text{Be}(e, e' \alpha_0)$ reaction (a peak at $E_m = 2.5$ MeV) and the α emission from ${}^9\text{He}$ (two-peaked part at higher missing energies).

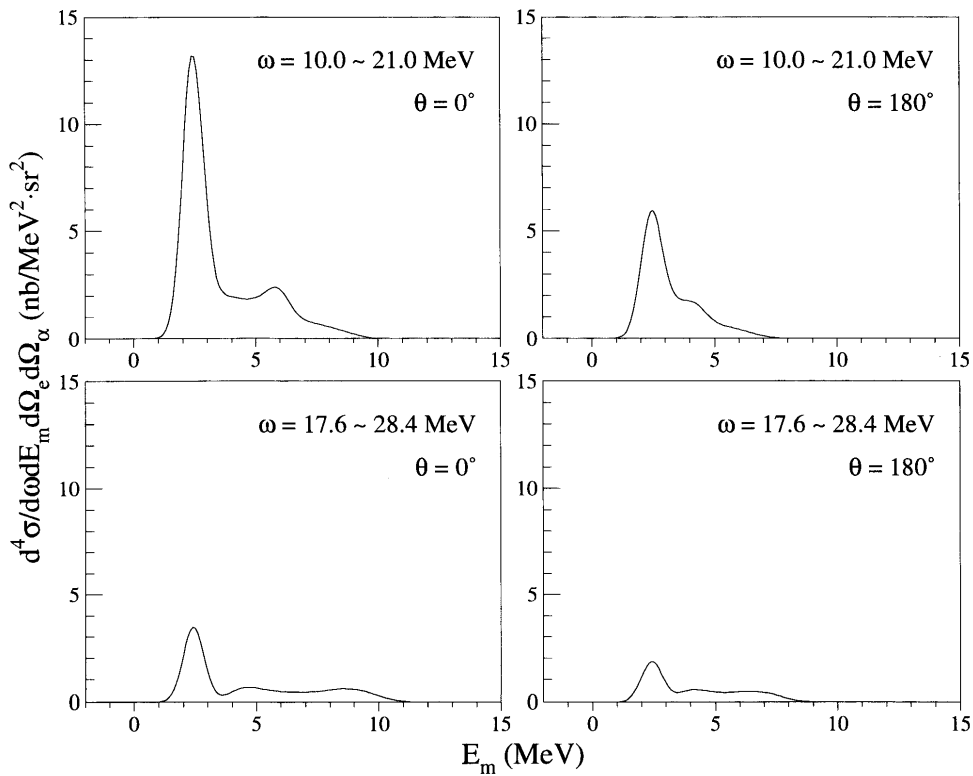


Fig.5. Simulated missing energy spectra folded in the regime of the measurements, which are smeared with a resolution of 1 MeV FWHM.

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