

High-Resolution Photoemission Study of Fine Electronic Structure in the Vicinity of the Fermi Level of High-T_c Superconductors(Abstracts of Doctoral Dissertations, Annual Report(from April 2001 to March 2002))

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High-Resolution Photoemission Study of Fine Electronic Structure in the Vicinity of the Fermi Level of High- $T_{\rm c}$ Superconductors

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In order to understand the mechanism of superconductivity in high- $T_{\rm c}$ (transition temperature) superconductors (HTSCs), there are several important issues to be solved. Typical examples are to clarify 1) the topology of Fermi surface (FS) in Bi-family HTSCs whether it is electron-like or hole-like, 2) the universality of $d_{\rm x^2-y^2}$ -like superconducting gap in hole-doped HTSCs, 3) the origin of pseudogap, 4) whether the mechanism of the supercodnctivity is the same between electron- and hole-doped HTSCs. To address these issues in more detail, it is necessary to improve the resolution of photoemission spectrometer. The purpose of present study is, firstly, to construct a new photoemission spectrometer and achieve the high resolution, and, secondly, to investigate the fine electronic structure in the vicinity of the Fermi level ($E_{\rm F}$) of HTSCs by performing the photoemission spectroscopy (PES) so as to elucidate above issues.

I. Construction of High-Resolution Photoemission Spectrometer Hemispherical analyze Figure 1 shows schematic view of new high resolution photoemission spectrometer. The system consists mainly CCD Electron lens of electron analyzer, main chamber, preparation chamber, photon source, and pumping lines. Ultrahigh-resolution 2.1 meV has been Monochromator achieved (shown in the inset of Fig. 1) by Diamond file Evapolator \ ı-metal designing the \u03c4-metal shield chamber, degaussing the analyzer, Cryo pum Prep. chambe adjusting the voltage source Cryp pump cables, calibrating the settings of lens voltages as well as CCD $\Delta E=2.1 \text{ meV}$ camera. Other specifications of Au T = 13.5 Kthe spectrometer are listed below; He la (21.218 e • Angular resolution: 0.2 deg • Lowest temperature: 6.4 K

• Best pressure: $1.5 \times 10^{-11} \text{Torr}$ Fig. 1 Schematic view of high-resolution photoemission spectrometer, and PES spectrum of Au in the vicinity of E_{F} .

II. Topology of Fermi Surface in Bi-family HTSCs

Figure 2 shows angle-resolved PES (ARPES) spectral intensity at 25 K of Pb-substituted $Bi_2Sr_2CuO_6$ (Pb-Bi2201, T_c <4 K) in the vicinity of E_F . Bright area corresponds to the bands. We clearly find single band all over the measured momentum region, demonstrating that the present sample is actually superstructure free. At away from $(\pi, 0)$ (lines A-J), the intensity shows remarkable dispersion, while that around $(\pi, 0)$ (lines J-L) shows almost no dispersion. A finite intensity at $(\pi, 0)$ unambiguously indicates a hole-like FS centered at (π, π) in Pb-Bi2201.

It is also observed that the hole-like FS stands irrespective of the doping ratio or the energy of photons. Similar hole-like FS has been found to exist in $^{0.3\pi}$ Bi₂Sr₂CaCu₂O₈ (Bi2212).

III. Pseudogap in LSCO

High-resolution PES of La_{1.85}Sr_{0.15}CuO₄ $\stackrel{\text{E}}{=}$ 0.3 π - (LSCO, T_c =38 K) has been performed. $\stackrel{\text{O}}{=}$ 0.5 It is found that the pseudogap of 40 0.3 π - meV opens above T_c and it does not smoothly connect to superconducting gap (8 meV) below T_c . The energy scale of the pseudogap well corresponds to the characteristic temperature of magnetic susceptibility (400 K). This suggests that the observed pseudogap is

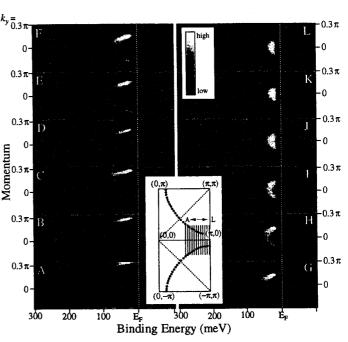


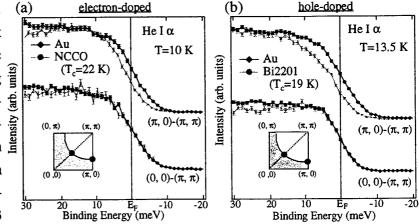
Fig. 2 ARPES intensity map of Pb-Bi2201(T_c<4 K) at 25 K measured along the lines in the Brillouin zone shown in inset.

formed by a development of antiferromagnetic correlation.

IV. Direct Observation of Superconducting Gap in Electron-Doped NCCO

Figure 3(a) shows ARPES spectra at 10 K of Nd_{1.85}Ce_{0.15}CuO₄ (NCCO, T_c =22 K) measured at two representative k-points. At Fermi vector (k_F) along (π , 0)-(π , π), we clearly find a leading edge shift in NCCO to higher binding energy with respect to the Au reference, showing that superconducting gap opens. On the other hand, leading edge along (0, 0)-(π , π) does not

show any remarkable shifts. These clearly indicate that NCCO has an anisotropic $d_{x^2-y^2}$ -like superconducting gap, as in the case of holedoped HTSCs. Furthermore, size of the gap in NCCO is found to be much smaller than that of holedoped Pb-Bi2201 (Fig.3 (b)), although the T_c 's are similar between the two.



(b)), although the T_c 's are Fig. 3 Comparison of ARPES spectra in the vicinity of E_F at the superconducting state between (a) $Nd_{1.85}Ce_{0.15}CuO_4(T_c=22 \text{ K})$ and (b) Pb-Bi2201($T_c=19 \text{ K}$). Spectra of Au are also shown as a reference.

V. Summary

Above experimental facts show that electron- and hole-doped HTSCs share common features in the topology of FS and the symmetry of superconducting gap. This indicates that essential framework of the superconductivity is the same between electron- and hole-doped HTSCs.