# Molecular Dynamics Simulation Studies of DNA and Proteins: Force Field Parameter Development of Small Ligands and Convergence Analysis for Simulations of Biomolecules 

Anne Loccisano

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# MOLECULAR DYNAMICS SIMULATION STUDIES OF DNA AND PROTEINS: FORCE FIELD PARAMETER DEVELOPMENT FOR SMALL LIGANDS AND CONVERGENCE ANALYSIS FOR SIMULATIONS OF BIOMOLECULES 

A Dissertation<br>Submitted to the Bayer School<br>of Natural and Environmental Sciences

Duquesne University

In partial fulfillment of the requirements for the degree of Doctor of Philosophy
by
Anne Elizabeth Loccisano

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## Anne Elizabeth Loccisano

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# MOLECULAR DYNAMICS SIMULATION STUDIES OF DNA AND PROTEINS: 

FORCE FIELD PARAMETER DEVELOPMENT FOR SMALL LIGANDS AND CONVERGENCE ANALYSIS FOR SIMULATIONS OF BIOMOLECULES

By<br>Anne Elizabeth Loccisano

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# ABSTRACT <br> MOLECULAR DYNAMICS SIMULATION STUDIES OF DNA AND PROTEINS: FORCE FIELD PARAMETER DEVELOPMENT FOR SMALL LIGANDS AND CONVERGENCE ANALYSIS FOR SIMULATIONS OF BIOMOLECULES 

By<br>Anne Elizabeth Loccisano

December 2007

Dissertation Supervised by: Dr. Jeffrey D. Evanseck
In the first part of this dissertation, CHARMM force field parameters for DNA minor groove-binding polyamides were developed. The parameterization involved the subdivision of the polyamides into model compounds, which were calibrated against MP2/6-31G(d) data. To test the new parameters, fourteen 10 ns molecular dynamics crystal simulations have been carried out on a DNA/polyamide complex at low (113K) and high ( 300 K ) temperatures. Of the 18 helical parameters examined, only one (stagger) is found to be statistically significant from the crystal structure with a t-test at the $95 \%$ confidence level. For the high temperature, stagger is non-significant at the $97 \%$ confidence level, which underscores the importance of running multiple trajectories. It is observed that when the simulations are run at 300 K , the DNA fragment begins to distort; however, better sampling is achieved. Competition between water and polyamides for hydrogen bonding to DNA is found to explain weak or unpredictable binding.

In the second part, force field parameters for retinoids were developed. The retinoids were divided into model compounds and calibrated against MP2/6-31G(d) data. To test the parameters, five molecular dynamics crystal simulations of reported x-ray structures of protein/retinoid complexes were performed. The structural and geometric analysis of these simulations compares well to experiment, and some dynamics that could be important to ligand binding were discovered. The new parameters can now be used in simulations of retinoid-binding proteins to better understand these systems and in drug design to make new retinoids with therapeutic and anticancer potential.

The last part explores the convergence of structural parameters in biomolecular systems. A simple statistical test was applied to the different parameters from a few long and many short simulations to observe which strategy is best. For the protein, both the long and short simulations gave similar results with respect to convergence. For the DNA, it was found that fraying effects penetrate four base pairs in from the ends of the helix. Structural parameters converge more quickly for the middle four bases than for all bases, and the long simulations yielded better results with respect to convergence than the short simulations.

To Ben and my family

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## Chapter 1

## Introduction

### 1.1 Historical Perspective on Simulation

The first molecular dynamics simulation of a biomolecular system was published in 1977. ${ }^{1}$ The simulation was of bovine pancreatic trypsin inhibitor (BPTI), and although the system was small ( 58 residues; 885 atoms), the simulation was performed in vacuum, the force field was primitive, and the simulation only lasted 9.2 ps , it showed that proteins exhibit substantial fluctuations on the picosecond timescale. ${ }^{1}$ The first molecular dynamics simulations of DNA followed in 1983 and were performed on a 12- and 24-base pair fragment in vacuum and without accounting for electrostatic forces. ${ }^{2}$ Since then, a wide range of protein and nucleic acid motions and structure features have been discovered and explored through molecular dynamics simulations. ${ }^{3-11}$ Advances in computer hardware, software, force fields, improvement in simulation methodology, and also the tremendous progress in experimental structure determination have allowed the field of MD simulation to grow at a rapid rate. Since the first BPTI simulation in 1977 in vacuum on a short timescale, these advances now enable large ( $10^{4}-10^{6}$ atoms) and complicated systems ${ }^{3,12}$ (such as a membrane-bound protein with explicit lipid and
water molecules) to be simulated on longer timescales (hundreds of nanoseconds to microseconds) with a more realistic representation of the solvent and surroundings.

Molecular dynamics simulations can provide the methodology to model, in ultimate detail, the internal atomic motions and fluctuations of a biomolecular system as a function of time. MD simulations allow for observation of phenomena and properties of systems that cannot be observed with experimental techniques. Although MD simulations have been widely used to give insight into the underlying natural dynamics, observe functionally important motions, make predictions, and facilitate experimental structure determination, the dynamics of many biomolecular systems are poorly or not understood. ${ }^{12-14}$ In addition, in order to compute and extract accurate information from the modeled system, the models used must be as realistic as possible (e.g., the force field used must be parameterized for the molecules of interest). Also, the various properties of interest that are calculated need to have reached satisfactory levels of convergence; the properties must be stabilized if analysis is to be performed on them. ${ }^{15,16}$

### 1.2 Force Field Development

In molecular dynamics and molecular mechanics, the intra- and intermolecular forces between the atoms are described by an empirical force field (the CHARMM force field, which is used in this dissertation, is discussed in Chapter 2). The mathematical terms in empirical force fields that are used to describe the intra- and intermolecular forces in biomolecules are relatively simple, and force fields assume that an atom is the smallest particle in the system of interest, rather than considering nuclei and electrons explicitly as in quantum mechanics. ${ }^{3,17}$ Both of these factors
allow for the computational demands of biomolecular simulation studies to be fulfilled. However, the force field parameters must be developed correctly in order to obtain accurate and reliable results and predictions.

The quality and accuracy obtained from a calculation performed with a particular force field is obviously based on the methods and target data used to optimize the parameters. Several considerations should be taken into account when parameterizing a force field for a new class of compounds. ${ }^{18}$ The first is the issue of transferability; this refers to the ability to take parameters optimized for a specific set of target data and then apply those parameters to molecules that are not included in the target data. In a transferable force field, the dihedral parameters that are optimized for rotation about the $\mathrm{C}-\mathrm{C}$ bond of ethane could then be used for butane. In a nontransferable force field, the dihedrals present in butane but not in ethane (C-C-C-C) must be optimized specifically by using target data for butane. However, how transferable parameters are depends on the similarity of the molecules, and what parameters are actually transferable and which are not is not well-defined. ${ }^{18}$ Any new compounds to be used with the force field must be tested in order to ensure that they are treated correctly and the required accuracy is obtained. Generally, the extent of transferability is considered to be minimal, and new parameters must be generated for each new class of molecules to be used with the force field. ${ }^{18}$

Since new parameters must usually be generated and optimized, a second consideration is the selection of target data that is to be used to optimize the new parameters. The form of the CHARMM force field is discussed in Chapter 2; therefore, the terms will not be described and shown here. The parameters needed to
extend the force field for a new class of compounds will fit into the different terms of the force field, and these are generally optimized using different sources. The internal parameters such as equilibrium bond lengths, equilibrium angle values, and dihedral multiplicity are often optimized to reproduce gas-phase geometric data obtained from quantum mechanical calculations, electron diffraction, or microwave experiments. The internal force constants, such as the bond and angle force constants, are usually optimized using vibrational spectra, which contain individual frequencies and their assignments. The external parameter optimization is usually more difficult than the internal parameter optimization; the amount of target data relative to the amount of needed parameters is decreased compared to the internal parameters. ${ }^{18}$ Upon the development of the OPLS force field, optimization of van der Waals terms improved. ${ }^{19}$ Heats of vaporization of pure solvents were computed through Monte Carlo simulations, and this data was then used to refine the van der Waals parameters. ${ }^{19}$ This same approach has also been applied to the CHARMM and AMBER force fields. ${ }^{20,21}$ Van der Waals parameters have also been optimized using the calculated heats of sublimation of crystals. ${ }^{22}$ The approaches for optimization of electrostatic parameters are mostly dominated by the reproduction of target data from QM calculations. ${ }^{18}$ One method is based on optimizing partial atomic charges to reproduce the electrostatic potential (ESP) around a molecule calculated with QM methods. A popular variation of this method is called the RESP method, ${ }^{23}$ where the charges on atoms that are minimally exposed to solvent are restrained. The goal of both methods is to produce the partial atomic charges that reproduce the electrostatic field created by the molecule. ${ }^{24}$

When parameterizing the force field for a new set of compounds, it is important to remember that the force field is "empirical". The choice of method and target data and even the form of the potential energy function is the decision of the developer of the force field. ${ }^{18}$ Even if the same form of the potential energy function is used, if developers have used different target data, the quality and results of the force field can differ. A great amount of correlation exists between all the different, individual parameters of the force field, and therefore, a number of different combinations of parameters can reproduce a given set of target data. ${ }^{18}$ Therefore, the optimization process is extremely important to the quality and results of the force field. Automatic parameterization procedures have been attempted; ${ }^{25-27}$ however, a significant amount of manual work is generally required. ${ }^{18}$

A third consideration is the procedure used in which to optimize the force field. When a force field is selected for use, the information that one wishes to extract must be considered. For example, if one is interested in examining the atomic details of water interactions with protein residues, the proper force field to use would be an all-atom force field (as opposed to an extended atom force field where hydrogens are not explicitly represented but treated as part of the atom to which they are bound) designed specifically for biomolecules that allows for explicit representation of water molecules. In this case, the AMBER or CHARMM force fields would be a suitable choice; however, they may not be of sufficient accuracy for a particular study. ${ }^{18}$ If the study mentioned above also involved different but structurally similar ligands being bound to the protein, the AMBER or CHARMM force field must be extended to treat the new class of compounds. The optimization
method used must be consistent with the method used for the original development of the force field. Chapters 3 and 4 of this dissertation contain the details of how the CHARMM force field was parameterized and tested for two new sets of compounds, polyamides and retinoids.

### 1.3 Project Overviews

This dissertation focuses on the three issues mentioned above: force field parameterization, gaining insight into dynamics of the systems using the new force field parameters, and analyzing biomolecular systems for convergence. The first part of the force field parameterization focuses on the polyamides, which are a class of DNA minor groove binders. ${ }^{28}$ The molecules have potential as DNA sequencespecific recognition agents; however, why some show poor binding affinity for their target sites and why some sites are more difficult than others to target is not well understood. Force field parameters for the polyamides were created, and through testing the parameters with molecular dynamics simulations in the crystal environment, a new conceptual idea as to why polyamides are weak binders emerged as well as an interesting phenomenon concerning the sampling of DNA helical parameters. The second part of the force field parameterization focuses on retinoids. Attempts to model $\pi$-conjugated systems have been difficult to implement ${ }^{29-31}$ or the compounds have not been explicitly parameterized. ${ }^{26,32}$ New parameters were made for conjugated systems with a focus on retinoids. Retinoids are an attractive class of compounds in anticancer therapy now; however, toxicity limits their use. ${ }^{33,34}$ Through testing the parameters with crystal simulations, interesting details were uncovered with respect to ligand motions that could be important in ligand binding
and dynamics. These motions enable better understanding of how the ligand behaves in the binding site of the protein, and this can aid in the future design of less toxic retinoids for use as anticancer agents. The last part of this dissertation focuses on the convergence of structural/geometric parameters in two biomolecular systems, DNA and a small protein. Multiple trajectories are commonly run in order to improve sampling of conformational space, ${ }^{35-37}$ and this approach was used in this study. A set of a few long simulations and a set of many short simulations were run, the parameters of interest were measured, and a statistical test was applied to the different parameters from the sets of simulations to observe which simulation strategy works best.

### 1.4 DNA Minor Groove-Binding Polyamides

The sequence specific recognition of DNA by small, organic molecules is an area of intense research due to the ability of these agents to interfere with critical biochemical processes involving DNA. Pioneering work by Dervan, ${ }^{38-57}$ Dickerson, ${ }^{58-63}$ and Lown ${ }^{64-70}$ on analogues of the natural products distamycin and netropsin (which are made up of three and two N-methylpyrrole (Py) rings, respectively) has led to the development of compounds that are capable of recognizing any base pair combination found in DNA. These compounds, termed polyamides or lexitropsins, have functional groups that are complementary to the distinct pattern and position of hydrogen bond donors and acceptors found in the minor groove of DNA. ${ }^{28,39,41,45,71-73}$ The polyamides are made up of five-membered (and recently, six-membered) heterocyclic amino acids linked by amide bonds that can be combined as antiparallel ring pairs in the minor groove of DNA to recognize
predetermined sequences of DNA, with affinities and specificities comparable to those of DNA-binding proteins (i.e., $\mathrm{K}_{\mathrm{d}}<1 \mathrm{nM}$ ). ${ }^{41}$

The heterocycle pairs used in polyamides are pieced together with the intervening amide bonds in order to recognize a predetermined DNA sequence based on hydrogen bond donors and acceptors, molecular shape, and electrostatic potential of the minor groove. The Dervan group has established a code describing this relationship. ${ }^{40,74-77}$ It was first discovered that a $\mathrm{Py} / \mathrm{Py}$ pair could distinguish an AT or TA pair from CG or GC pairs (AT/TA pairs are fairly symmetrical compared to the unsymmetrical nature of GC/CG pairs due to the exocyclic NH2 group of G); shortly after, it was found that an imidazole (Im)/Py pair could distinguish GC from the other three base pairs (due to a linear hydrogen bond between the Im nitrogen and the exocyclic NH2 group of G). It took more time to distinguish AT from TA; it was finally discovered that the Py/hydroxypyrrole (Hp) pair could distinguish AT from TA, due to the hydrogen bond formed between the hydroxyl group and the O 2 of thymine. Thus, the basic Dervan pairing code is as follows: Im/Py pair binds GC, a Py/Im pair binds CG, a Py/Py pair binds AT or TA, and an $\mathrm{Hp} / \mathrm{Py}$ pair and a $\mathrm{Py} / \mathrm{Hp}$ pair distinguish TA and AT, respectively. ${ }^{40,41,74}$ All polyamides contain a positively charged group at the tail end (dimethylaminopropyamine (Dp)), which interacts with the negatively charged phosphate backbone and increases affinity.

Polyamides can bind to DNA in either 1:1 or 2:1 stoichiometries; however, less is understood about polyamide recognition in the $1: 1$ mode. ${ }^{78}$ It is known that in the $1: 1$ binding mode, polyamides can form bifurcated hydrogen bonds between bases on either one of the two strands in DNA. Also, because DNA sequences containing a
run of purine bases has a particularly narrow minor groove ${ }^{79}$ and cannot accommodate two polyamides in the groove, the $1: 1$ motif shows promise in targeting these types of sequences. ${ }^{78,80}$

In the $2: 1$ binding mode, which is the most common binding mode, the heterocycles of two polyamide molecules are paired and stack on top of each other when bound to the minor groove of DNA. In this mode, one polyamide is capable of reading one strand of DNA, while the other polyamide recognizes the second strand. In addition to forming hydrogen bonds to the DNA, the interaction is stabilized by the DNA-polyamide van der Waals interactions and $\pi$-stacking interactions formed between the ring pairs of the polyamides. Figure 1.1 shows a 3-D representation of how two antiparallel polyamides bind in the minor groove of DNA (left) and a schematic diagram of how the paired heterocycles recognize the hydrogen bond donor-acceptor pattern of the minor groove (right).



Figure 1.1 On the left: 3-D representation of polyamides bound 2:1 in the minor groove of DNA ( PDB $^{81}$ id $1 \mathrm{CVY}^{73}$ ). DNA is shown as blue sticks, and the polyamides are shown in van der Waals representation (imidazoles in red, pyrroles in yellow, amide bonds in cyan, $\beta$-alanine linker in orange, and cationic tail in green). The structure contains the sequence 5'-CCAGATCTGG-3' complexed with two polyamides of sequence ImPyPyPy- $\beta$-Dp. On the right: Schematic diagram of how the paired heterocycles recognize the hydrogen bond donors and acceptor of the minor groove. The "ladder" in the middle is the DNA binding site, the dashed lines are the hydrogen bonds, the circles with the dots are the lone pairs on O 2 of T and C and N 3 of A and G , and circles with an H are the Hs of the exocyclic NH2 of G . The tail is made up of two parts: $\beta$-alanine, which targets $\mathrm{A}, \mathrm{T}$ bases and the cationic Dp , which interacts with the phosphate backbone (figure adapted from Dervan, et al., 2005 ${ }^{71}$ ).

In the $2: 1$ binding mode, the heterocycles are usually paired by the covalent linkage of two antiparallel polyamide strands; the covalent linkage results in increased affinity and specifity (100-fold higher affinity than their unlinked counterparts). ${ }^{28}$ The standard motif for polyamides is a hairpin containing eight heterocycles, where two strands containing four heterocycles each are linked by a $\gamma$ aminobutyric acid linker. Called a $\gamma$-turn, it shows preference for binding AT/TA pairs over GC/CG pairs (due to steric clashes with the exocyclic NH2 group of $\mathrm{G}^{82}$ ),
and it connects the carboxylic end of one polyamide to the amino end of the other polyamide (thus, the polyamides are antiparallel). The $\gamma$-turn stabilizes the register of the heterocycle pairings with the DNA bases, which is most likely the reason for the higher affinity and specificity. ${ }^{83}$ Hairpin polyamides align themselves with the DNA so that the amino end of the polyamide is oriented toward the 5' direction of the adjacent DNA strand, and the carboxylic end oriented toward the 3 ' direction of the adjacent DNA strand. ${ }^{84}$

Although the normal eight-ring hairpin motif binds with high affinity and specificity, this motif binds only six base pairs of DNA. In order for the polyamides to be useful in biology or medicine, specifically recognizing a very long stretch of DNA is important; longer sequences would be expected to occur less frequently in large genome. ${ }^{28,71}$ However, expanding the binding site size simply by adding extra heterocycles to the polyamides has been shown to decrease affinity and specificity because the curvature of the polyamides is tighter than that of the curvature of the DNA groove. ${ }^{85}$ Actually, beyond five heterocycles, the polyamide shape is no longer complementary to the DNA groove..$^{47}$ In order to allow more conformational freedom in the polyamides, a flexible $\beta$-alanine residue is used. The $\beta$-alanine residue is used as a replacement for Py , so $\beta / \mathrm{Py}$ and $\beta /$ Im pairs function as $\mathrm{Py} / \mathrm{Py}$ and $\mathrm{Py} / \mathrm{Im}$ pairs and simultaneously relaxe the curvature of the polyamide, restoring complementarity to the DNA groove. ${ }^{86}$ Polyamides with $\beta$-alanine have been used to target successfully up to 16 base pairs with high affinity and specificity. ${ }^{52}$

Although most polyamide research has focused on the five-membered heterocycles, other heterocycles have recently been shown to bind to DNA. ${ }^{28,71}$

These heterocycles are based on the benzimidazole (Bi) ring system, which consists of a $6-5$ bicyclic ring structure. Different substituents can be added to the sixmembered ring, and polyamides with the benzimidazole derivatives have a curvature that is more complementary to the DNA than the traditional five-membered ring polyamides. ${ }^{87,88}$ The benzimidazole derivatives (imidazopyridine (Ip) and hydroxybenzimidazole (Hz)) have been incorporated into the eight-ring hairpin polyamides and have been found to be as effective as the five-membered ring polyamides at recognizing the minor groove of DNA; the hydrogen bonding contacts between the DNA and polyamide are preserved. ${ }^{55}$ Of particular interest are the $\mathrm{Hz} / \mathrm{Py}$ and $\mathrm{Py} / \mathrm{Hz}$ pairing. This pair can replace $\mathrm{Hp} / \mathrm{Py}$ and $\mathrm{Py} / \mathrm{Hp}$ pairs, and because $\mathrm{Hp}-$ containing polyamides have been found to degrade over time and the Hz polyamides are chemically stable, the $\mathrm{Hz} / \mathrm{Py}$ pair is a candidate to replace the $\mathrm{Hp} / \mathrm{Py}$ pair. ${ }^{28}$ However, because the benzimidazole system is more hydrophobic and has a greater surface area, this may alter the DNA-polyamide van der Waals interactions and the $\pi$ stacking interactions between the polyamide ring pairs. ${ }^{28}$

Polyamides utilized in this dissertation are composed of three most common five-membered ring heterocycle types joined by an amide functional group. The NH group of the amide bonds is critical for binding as it forms specific hydrogen bonds with the bases when bound to the DNA. The three heterocycles, N-methylpyrrole (Py), N-methylimidazole (Im), and N-methyl-4-hydroxypyrrole (Hp) and the amide linkages are shown in Figure 1.2.


Figure 1.2 Schematic diagram of Im-Py-Hp polyamide fragment.
The Dervan pairing rules have been used to design literally hundreds of synthetic ligands that do bind their predetermined DNA target sequences with the high affinity and specificity comparable to DNA binding proteins. The polyamides have been shown to affect both transcriptional activation (by inhibiting repressor proteins and thus activating a particular gene) ${ }^{43,89-93}$ and repression (by inhibiting assembly of transcriptional proteins on DNA and thus downregulating a particular gene) in vitro and in cell cultures. ${ }^{39, ~ 42, ~ 69, ~ 94-98 ~ P o l y a m i d e s ~ l a b e l e d ~ w i t h ~ f l u o r e s c e n t ~}$ dyes have also been shown to be taken up by cells, although nuclear uptake is dependent on cell type. ${ }^{38,44,99}$ However, even though the polyamides bind to B-DNA and have been found to perform the functions mentioned above, there are limitations regarding targeted DNA sequences. Sequence dependent structural variations of the DNA minor groove affect the energetics of binding, a number of DNA sequences remain difficult to target (especially purine tracts), and binding affinity and specificity at some sequences is weak and/or unpredictable. Variation in the DNA minor groove width, curvature, flexibility, hydration, and the relative positions of hydrogen bond donors/acceptors could all possibly influence polyamide binding and affinity for a particular sequence. ${ }^{71}$ Because molecular dynamics simulations allow for an atomistic view and examination of the DNA minor groove, this is an ideal tool for better understanding of the interactions of the polyamides with the DNA.

Computational studies are important for the description on how distamycin A, netropsin, and related polyamides bind to DNA. ${ }^{72,76,78,83,86,100-121}$ For example, molecular dynamics simulations have been used to aid in the refinement of NMR structures of DNA/distamycin and DNA/polyamide complexes. ${ }^{\text {72, 78, 83, 100, 102-104, } 121}$ Free energy perturbation studies have been utilized to understand binding affinities of DNA/polyamide complexes, ${ }^{103,109}$ and unrestrained molecular dynamics have been employed to understand features of the complex such as hydrogen bonding, hydration, and structural distortions on the DNA induced by the polyamides. ${ }^{108,110-112,}$ ${ }^{114-120}$ Finally, quantum mechanical studies have been applied to understand the electronic and geometric properties of novel polyamides, and to interpret the results of experimental binding studies. ${ }^{76,101,106}$

The quality of the force fields implemented to treat polyamides varies greatly between studies. Close examination of the force field parameters utilized underscores the need for parameter development. For example, treatments of DNA complexes with an isopropyl-thiazole polyamide, ${ }^{100}$ lexitropsin, ${ }^{103} \mathrm{CPI}$-lexitropsin,,$^{102}$ distamycin $\mathrm{A},{ }^{104,}$, 107-109, 121 netropsin, ${ }^{109,} 114$, 117, 118, 122 2-imidazoledistamycin, ${ }^{107}$ carbocyclic analogs of netropsin and distamycin, ${ }^{110,111}$ carbocyclic derivatives of distamycin with a chlorambucil moiety, ${ }^{112}$ ImHpPyPy- $\beta$-Dp, ${ }^{115,}{ }^{116,} 119$ ImPyPy- $\gamma-$ PyPyDp, ${ }^{112}$ ImPy- $\beta$-Im- $\beta$-ImPy- $\beta$-Dp, ${ }^{78}$ Ac-ImPyPy- $\gamma-$ PyPyPy-Gly-Dp, ${ }^{72}$ and ImPyPy- $\gamma$-PyPyPy-Gly-Dp ${ }^{72}$ have used default force field parameters that are best guesses or extrapolations from similar functionality. Force field development based upon the recruitment of similar terms and parameters may be successful in some cases; ${ }^{18}$ however, in the case of polyamides made up of Im, Py, and Hp, serious errors
in the torsional terms, multiplicity, phase shift, and force constant magnitude have been found. The simulation results could lead to artificial behavior that could impact the interpretation of structure, dynamics, and properties of polyamides and their complexes.

Zhang and coworkers created polyamide parameters in the AMBER force field ${ }^{21}$ to determine the structure of a DNA/cyclic polyamide complex by NMR. ${ }^{105}$ Energy minimization and molecular dynamics with the AMBER force field were used to aid in structure refinement; their results for the structure indicate distortion of the DNA in the polyamide binding site and altered stacking of the ligand rings relative to noncyclic polyamides. ${ }^{105}$ A number of specific bonds, angles, and torsional terms were calibrated against B3LYP/6-31G(d) data. However, a majority of the terms were derived by comparison with other functionalities, and no comparison was made back to experimental structures. It is well known that density functional methods, such as B3LYP, do not account for the effects of dispersion. ${ }^{123-131}$ Dispersion effects are important in the development of polyamide force fields because polyamides are cresent-shaped, which complements the shape of the narrow minor groove and improves binding by promoting van der Waals interactions. ${ }^{132}$ Also, in the case of the 2:1 binders, the rings of the polyamides must stack on top of one another in the minor groove, which is stabilized mainly by nonbond interactions. ${ }^{132}$

Pang developed force field parameters for $A M B E R^{21}$ to study the formation of DNA/polyamide complexes. ${ }^{120}$ The results suggest that three and four ring polyamides have the same tendency for forming the same dimer conformations in water as they do in their ternary complexes with DNA and that the "pre-dimerization"
of polyamides in water assists in generating a tighter and more selective binding to DNA. ${ }^{120}$ The majority of bonds, angle, and dihedral terms were calibrated against HF/6-31G(d), and no comparison was made back to any experimental structures. In a separate study, Topham and Smith made parameters for the CHARMM force field for distamycin and netropsin. ${ }^{113}$ Netropsin and distamycin were divided into smaller model compounds, and potential energy surfaces for specific torsions were generated using HF/6-31G(d). The energy-minimized model of netropsin had an RMSD of 0.28 $\AA$ with the crystal structure. It is known that Hartree-Fock does not account for the effects of electron correlation. ${ }^{133}$

Wellenzohn, et al. performed simulation studies with the AMBER force field ${ }^{21}$ on DNA/netropsin ${ }^{114,118}$ and other DNA/polyamide complexes. ${ }^{21,115-117,119}$ The netropsin molecule and other polyamides were not explicitly parameterized for the force field. Only the partial charges on the atoms of the molecules were explicitly derived; the other parameters for these molecules were taken from analogous existing functionality and molecules in the force field. Generally, the extent of transferability of force field parameters is considered to be minimal, and new parameters must be generated for each new class of molecules to be used with the force field. ${ }^{18}$ No comparison back to experiment was reported.

Given the problem of interest in polyamide/DNA simulation and the lack of accurate force field parameters to perform molecular dynamics simulations, we have developed parameters for the polyamides for the CHARMM force field. ${ }^{20}$ Parameterization is accomplished by first dividing the polyamides into six model compounds and carrying out energy minimization using MP2/6-31G(d), which
includes the effects of electron correlation and dispersive interactions. The bond, angle, torsional, and improper terms were modified to match the quantum mechanical results. The parameters were tested by performing molecular dynamics crystal simulations of the reported x-ray crystal structures of a netropsin crystal, a 2:1 polyamide/DNA complex, a netropsin/DNA complex, a distamycin/DNA complex, and DNA without bound polyamides. The new parameters can be used for improved simulations of polyamides and their nucleic acid complexes.

### 1.5 Retinoids and $\pi$-Conjugated Systems

Retinoids make up a large group of naturally occurring and synthetic compounds related to retinol, otherwise known as vitamin A. ${ }^{134135136}$ Physiologically occuring retinoids consist of a $\beta$-ionene ring, a chain of conjugated double bonds (that includes the ring), and a polar functional group at the end of the chain. The figure below illustrates the general retinoid skeleton with two naturally occurring retinoids, retinol (1) and retinal (2), and two synthetic retinoids, axerophthene (3), and fenretinide (4).


1


3


2




Figure 1.3 Structures of four retinoids. 1, retinol; 2, retinal; 3, axerophthene; 4, fenretinide.

Retinoids are essential in a number of vertebrate physiological processes; they play vital roles in cell growth, ${ }^{137}$ reproduction (spermatogenesis, conception, and placental formation), ${ }^{138}$ embryogenesis, ${ }^{138}$ vision, ${ }^{136,139}$ learning and memory, ${ }^{140}$ and resistance to and recovery from infection. ${ }^{136}$ Retinol (also known as vitamin A) and its active metabolites (the most active is retinoic acid) also control the differentiation, proliferation, and apoptosis of many different types of cells from conception to death, ${ }^{138,141}$ including epithelial cells in the digestive tract, respiratory system, skin, bone, and the central nervous system. ${ }^{136,142,143}$ Retinol cannot be synthesized de novo by the human body, so it is an essential part of the diet. ${ }^{134}$ Dietary retinol is taken in as retinyl esters (obtained from animal products primarily as retinyl palmitate) or as $\beta$-carotene (obtained from vegetables). ${ }^{144}$ The retinyl esters and $\beta$ carotene are then metabolized into the various retinoids found in the body, including retinol, retinoic acid, and retinal. Retinyl palmitate is thought to be hydrolyzed to form retinol and a carboxylic acid in the intestine before being transported to the liver. In the liver, which is the main storage site of vitamin A, retinol is esterified again to maintain the storage levels. $\beta$-carotene is thought to be cleaved to form two retinal molecules in the intestine. The retinal can then be reduced by retinaldehyde reductase to yield retinoic acid or it can be esterified to retinyl esters for storage. ${ }^{33,138}$

The stored retinyl esters in the liver are hydrolyzed to all-trans retinol, which is the major circulating retinoid, for release into blood circulation. ${ }^{144}$ The carrier protein for all-trans retinol in the blood to various target tissues is binding protein (RBP), which controls the plasma levels of retinol. Retinoids have low solubility in their uncomplexed form in aqueous medium so they must be protected, solubilized,
and transported in the blood and in the cell by specific retinoid binding proteins. ${ }^{135}$, 145-147 In circulation, holo RBP is bound to transthyretin (TTR), which consists of four subunits. Formation of the complex is thought to prevent filtration of the relatively small RBP through the kidneys and is hypothesized to be the complex that interacts with cell surface receptors that mediate the uptake of retinol. ${ }^{148,149}$

Once inside the cells of target tissues, retinol can be metabolized by various enzymes to retinal and retinoic acid or converted back to the ester form for storage; all forms (alcohol, acid, aldehyde and ester) can be present in the cell and are interconvertible. ${ }^{150}$ Cytoplasmic retinol-binding proteins (CRBP I and CRBP II) and cytoplasmic retinoic acid-binding proteins (CRABP I and CRABP II) bind retinol and retinoic acid with high affinity, respectively, inside the cell. ${ }^{151}$ Other than the obvious binding of retinoids in the cell, the functions of these cytoplasmic retinoid-binding proteins remain to be defined; however, they are thought to modulate intracellular retinoid metabolism by influencing the amount of ligand available to nuclear receptors. ${ }^{135,151,152}$ The regulation of retinoid metabolism in the cell is crucial because a shortage or excess or retinoic acid can affect retinoid-signaling pathways and is known to result in developmental defects. ${ }^{33,150,153}$ However, it is known that retinol bound to CRBP I is the substrate of enzymes that synthesize retinyl esters or oxidize retinol to retinoic acid. Retinoic acid bound to CRABP I or II is then delivered to the nuclear receptors, which direct and regulate the expression of specific retinoic-acid inducible genes. ${ }^{143}$ CRBPs and CRABPs show a specific tissue distribution in adult and developing animals, which suggests that these proteins each serve different functions. ${ }^{151}$ Visual tissues contain other specific binding proteins for

11-cis retinal and 11-cis retinol (known to function only in vision), cellular retinaldehyde-binding protein (CRALBP) and interphotoreceptor retinol binding protein (IRBP). ${ }^{138,154}$

Two classes of nuclear receptors mediate the biological activities of retinoids, retinoic acid receptors (RARs) and retinoid X receptors (RXRs), both of which have three subtypes $(\alpha, \beta$, and $\gamma){ }^{34,134,140}$ RARs bind and are activated by all-trans retinoic acid and 9-cis retinoic acid with similar affinity but bind 13-cis retinoic acid only weakly; RXRs bind and are activated by only 9-cis retinoic acid. ${ }^{34}$ RARs and RXRs are transcription factors that bind to specific DNA sequence-retinoic acid response elements as RAR-RXR heterodimers and RXR homodimers, acting as ligand-dependent transcriptional regulators for retinoic acid responsive genes. ${ }^{34,155,156}$

Because retinoids play essential roles in cell differentiation, growth, proliferation, and apoptosis in all stages of life in vertebrates, their pharmaceutical and chemopreventative properties have been exploited. The many effects that retinoids have on cells have made this class of compounds an attractive candidate for use in cancer therapy and other diseases. ${ }^{34,155,157,158}$ Both naturally occurring and synthetic retinoids are used to treat various skin disorders (acne, psoriasis, and ageing), ${ }^{155}$ they have been in used to augment treatments of other diseases, such as type II diabetes, ${ }^{159}$ and they show promise for the treatment of several cancers, ${ }^{142,158,}$ ${ }^{160,161}$ HIV infection, ${ }^{162}$ and late-onset Alzheimer's disease. ${ }^{140}$ These agents have targeted the RARs and RXRs; however, the scope of retinoid therapy is limited due to high toxicity (only a few retinoids are used routinely for treatment of skin disorders $\left.{ }^{163-165}\right),{ }^{34}$ and the design of less toxic retinoids is necessary for further
clinical use. Because vitamin A and its derivatives play a multitude of roles in all stages of life through activation of widely expressed RARs and RXRs, administration of additional retinoids (natural or synthetic) causes toxic side effects that are characteristic of excess levels of vitamin A. ${ }^{33}$ Current efforts in retinoid development involve increasing the selectivity of synthetic retinoids to decrease the toxic side effects. ${ }^{155,166}$

Because of their biological importance and potential theraputic uses, it is of interest to understand the interactions and dynamics of retinoids with the various proteins that bind them at the molecular level. Molecular dynamics simulations are a powerful tool for elucidating the atomistic details of biomolecular systems; however, the empirical force field parameters used must be optimized so that they accurately treat the class of compounds that one is attempting to model. Previous attempts have been made to treat conjugated systems properly, which are fundamental to retinoids. The first attempts were made independently by Warshel and Karplus ${ }^{29}$ and Allinger and Sprague, ${ }^{30}$ where both involved a combined quantum mechanical/molecular mechanical approach.

The method of Warshel and Karplus was based on the formal separation of $\sigma$ and $\pi$ electrons, where the $\sigma$ electrons were treated using an empirical force field (the consistent force field, or $\mathrm{CFF}^{167}$ ), and $\pi$ electrons with Pariser-Parr-Pople semiempirical theory ${ }^{168,169}$ corrected for orbital overlap. The total energy of a molecule was then expressed as the sum of $\sigma$ and $\pi$ electron contributions. The goal of the method was to treat the equilibrium conformations and vibrational spectra of the ground and excited states of conjugated molecules. Experimental data from
ethylene, butadiene, benzene, and propylene were used to fit the $\sigma$ and $\pi$ parameters used in their method. ${ }^{29}$ The method was tested against the ground and excited states of several conjugated molecules, including 1,3-butadiene, 1,3-cyclohexadiene, and 1,3,5-hexatriene. The results were in good agreement with experimental geometries and vibrational frequencies; however, it was never generally used and thus never made widely available because at that time it required excessive computer resources, and it was not easily extended for large systems of interest.

Allinger also developed a combined quantum and molecular mechanical approach that separated the $\sigma$ and $\pi$ electrons. ${ }^{30}$ Allinger treated the $\sigma$ electrons with the MMP1 force field, ${ }^{170}$ and the $\pi$ electrons were treated with a variableelectronegativity self-consistent field (VESCF) calculation. The method was tested against the experimental geometries (x-ray crystallographic and electron diffraction structures) and heats of formation for a representative group of 65 compounds. The computed results yielded satisfactory agreement with experiment over the range of compounds with the largest difference between experiment and computation being $2.0 \mathrm{kcal} / \mathrm{mol}$ for the heats of formation. The program was never widely used by the scientific community due to the same issues limiting the Warshel and Karplus method.

A more recent method by Treboux for calculating geometries and vibrational frequencies of ground state and excited conjugated hydrocarbons also utilized the combined quantum mechanics and molecular mechanics methods. ${ }^{31}$ This method again partitioned the total energy into $\sigma$ and $\pi$ parts. In order to reduce the computational time involved with the iterative SCF algorithm utilized in the last two
methods described, this method makes use of a Heisenberg Hamiltonian for the $\pi$ electrons. ${ }^{31}$ The geometries and energy barriers calculated from various conjugated systems were in good agreement with experimental data; however, the authors made note of the fact that it could be practically applied to only 20 carbon atoms at the time the method was devised, which was in $1995 .{ }^{31}$

The most recent approach for treating a general class of large biomolecules, which include conjugated systems, was based completely on molecular mechanics. Both CHARMM and AMBER force fields were designed to be general purpose force fields to describe a wide variety of organic molecules, specifically, moieties that are commonly found in pharmaceuticals. ${ }^{20,26,32}$ Neither CHARMM nor AMBER were parameterized explicitly for conjugated systems.

The limitations of force fields to treat conjugated systems, such as retinoids, have been recognized. Numerous molecular dynamics studies involving retinalbinding proteins have been performed, primarily with the goal of understanding different aspects of the bacteriorhodopsin or rhodopsin photocycle. ${ }^{171-174} 175-185$ Some studies involved the parameterization of the retinal ligand using quantum mechanical calculations; however, the quantum mechanical calculations used to parameterize the retinal were performed using Hartree-Fock ${ }^{176,177,186}$ or B3LYP, ${ }^{180,183,187,188}$ which do not account for the effects of electron correlation and dispersion interactions, respectively. Many other studies used parameters for the retinal ligand that were simply adapted from other atoms or molecules in the force field being used in the study. ${ }^{171-175,178,179,181}$

In order to understand the interactions and dynamics of retinoid binding proteins and design new retinoid ligands with therapeutic potential using molecular dynamics simulations, this class of compounds has been explicitly parameterized for this dissertation. New force field parameters have been developed for the CHARMM force field for retinoids, which properly account for the alternating single and double bonds present in the conjugated systems of these compounds. The new force field has been evaluated against x-ray crystallographic data and quantum mechanical structures and energies. Through testing the parameters with crystal simulations, some ligand motions were also discovered that could be important in ligand binding and dynamics or interaction of the protein/ligand complex with other molecules. The observed motions aid in understanding more about how the ligands move internally within the protein, which can be used in design of new ligands. The parameters can now be used in simulations of retinoid binding proteins to better understand their structure, dynamics, and how they interact with other proteins and for the design of new retinoids that show promise as therapeutic and anticancer agents.

### 1.6 Convergence Analysis of Biomolecular Simulations

Understanding DNA and protein dynamics is critical not only for expanding our knowledge of DNA and protein structure in itself, but for understanding DNA and protein interactions and dynamics with other proteins, nucleic acids, and ligands. Therefore, it is crucial to build a knowledge base of DNA and protein dynamics, which is not currently available, especially for DNA. The energy landscapes of biomolecules are rough and rugged, containing many accessible local minima (corresponding to similar but slightly different conformations) separated by energy
barriers, and transitions between these minima take time. ${ }^{15,35,36,189,190}$ On the current time scale of MD simulations (hundreds of nanoseconds), the simulation can get stuck in the many minima on the potential energy surface, and thus not allow for complete sampling of all possible conformations of even a small (8-10 base pairs) DNA duplex or small protein (46 residues). ${ }^{191}$ Slow motions, such as large conformational changes, ion channel gating motions, or other functionally relevant motions, often occur on timescales that not easily accessible and thus are not wellcharacterized with simulations (microseconds to milliseconds). ${ }^{37,191}$ Furthermore, simulated biomolecular systems commonly contain a complex mixture of nucleic acids or proteins, water, and counterions (such as sodium, chlorine, or magnesium), and in some systems, lipid bilayers, which cause the convergence to occur more slowly, resulting from less efficient sampling. ${ }^{\text {.92-196 }}$

Sampling of minima is essential for the conformational flexibility of biomolecules, and therefore, to their function. ${ }^{35,197,198}$ Thus, if the biomolecular system of interest is not properly sampled, the properties measured from the simulation will not be representative of its true behavior, and predictions made could be incorrect. The results and predictions from simulation studies depend on proper sampling of the molecule's conformational space.

Simulations of biomolecular systems must address the issue of the ergodic hypothesis, which a problem for any system with multiple minima on its potential energy surface. ${ }^{35}$ The system may not cross all of the energy barriers sufficiently in the time that the simulation is run to sample properly all states accessible to it. The ergodic hypothesis states that average of a property over the entire simulation
trajectory is equal to an average over all states accessible to the system. ${ }^{15,35}$ The hypothesis is difficult to prove for any system; however, it is easier and possible to evaluate a necessary criterion for the ergodic hypothesis to be valid: At equilibrium, independent trajectories over an ergodic system must be self-averaging. ${ }^{15}$ This means that a property measured from two or more independent trajectories should be equal if the system dynamics are ergodic. In other words, if two independent simulations do not converge to equal values, the system is nonergodic on the timescale that the simulations were run. ${ }^{15}$ Obviously, different properties of the system will exhibit different convergence behavior (some will converge more quickly than others), so multiple properties should be examined in order to address ergodicity and convergence. ${ }^{35,199}$

Because of the complexity of biomolecular systems and the need for adequate sampling, one of the primary issues with MD simulations is convergence: How long does the simulation need to be run in order to extract reliable information? Acquiring reliable and accurate information on dynamics from MD simulations requires us to know how much sampling is required for the convergence of all of the properties of interest. Attaining convergence is obviously going depend on the system under study and the properties of interest, and it is a difficult question to answer. For effective sampling and therefore ergodicity, the simulations must correctly sample the regions of phase space that are accessible to the system under the given conditions (e.g., temperature and pressure). If only a single long simulation is run, the system can get stuck in one of the many minima and thus not sample effectively. Multiple shorter trajectories have been found to be more effective at sampling phase space than a
single long one, ${ }^{35-37,200-204}$ and this is the approach that is utilized in this dissertation for the convergence studies.

Rigid convergence of a system can never be proved, because there is no guarantee that the past dynamics of the system will predict the future dynamics, and in principle, new events or conformations of a system may be discovered with more extensive sampling. ${ }^{32,37,191,192}$ However, efforts can be made to ensure that the various properties calculated from simulations have reached a level of satisfactory convergence, and different approaches have been used to do this. Some methods are relatively simple and commonly used, such as monitoring the RMSD of the simulation structures over time with respect to the starting structure and calculating correlation times among different properties. ${ }^{32,196,205-208}$ However, the RMSD is not always a good indicator of convergence; ${ }^{191,209}$ we have found in this study that DNA simulations started from different atomic coordinates result in RMSDs that converge at different rates or fluctuate at different points when the simulation has run for a considerable amount of time ( $>100 \mathrm{~ns}$ ). In other words, the RMSD failed to settle down to a constant value; the system appears to be moving between various substates (see Chapter 5). This is not surprising for a flexible biomolecule like DNA, but it indicates that this method cannot determine when the simulation has converged. Also, if correlation times (time over which a particular property value takes to become independent from the previous values) longer than the overall trajectory length exist for the system, then the proper correlation time will not be observed; the correlation times will also depend on what part of the protein (a loop vs. a helix) or DNA (entire fragment vs. selected base pairs) is being examined. ${ }^{37,191,209}$ Another
commonly used method is block averaging. ${ }^{32,195,210}$ With block averaging, the total trajectory is broken up into independent, equally sized blocks, and the averages are then calculated and compared over the blocks. However, in order obtain independent blocks, the relaxation time of the property of interest must determined to figure out what length of block can be considered independent. ${ }^{32}$ This, as with correlation times, will present a problem if the trajectory has not been run sufficiently long enough for the property to relax.

Other more advanced methods, such as principal components analysis (PCA), ${ }^{198, ~ 206, ~ 211-216 ~}$ cluster population analysis, ${ }^{37,191, ~ 199, ~ 217-219}$ or energy-based methods ${ }^{189,220,221}$ have also been used to diagnose sampling and convergence. PCA, also called essential dynamics, ${ }^{211,222}$ has been to applied biomolecular simulations to visualize large-scale motions. ${ }^{198,206,211-216,223,224225,226}$ It is assumed that the dynamics of the large-scale motions analyzed by PCA are converged on the timescale of the simulation. ${ }^{37}$ However, this may not be the case; several studies have found that the motion visualized from PCA were not consistent among individual simulations. ${ }^{37,227-230}$ Although these studies showed that portions of a single trajectory or that individual simulations yield different results when PCA is applied, they do show that the simulation times examined are not long enough to sample properly the structure and fluctuations of the systems studied.

Cluster population analysis ${ }^{37,191,199,217-219}$ has been used to examine the relative populations of substates on the potential energy surface of the simulated system. This method aims to classify the structures generated from trajectories into clusters based on the RMSD between the trajectory structures and a reference
structure. In practice, sampling is limited so there are statistical variations among the cluster probability distributions; the magnitude of the variations is then considered a measure of the degree of convergence among the simulations. However, interpreting the cluster population variance can be difficult, ${ }^{37}$ and as a result, the whole procedure requires time and processing of data. Despite the time investment, this method has been used, and results suggest the same as those of PCA: the simulation times used were not long enough to produce ergodic behavior of the systems studied. ${ }^{37,191}$

The energy-base method (also called the "ergodic measure" ${ }^{189}$ and not limited to energy-based quantities) developed by Straub, Thirumalai, and Mountain ${ }^{189,220,221}$ examines the fluctuations of some quantity averaged over independent simulations, and the timescale required for that quantity to be ergodic is determined. However, a thorough sampling of the conformational space of the system's structures is required for convergence, and many structures are close in energy but not structurally similar, and energy-based methods will not be sensitive to this. ${ }^{191,209}$

In order to understand more about sampling and convergence in simulations of DNA and protein systems, we have approached the problem using several longtime simulations ( 4 simulations of 150 ns each for DNA; 10 simulations of 45 ns each for the protein crambin) and many short-time simulations ( 20 simulations of 2 ns each for both the DNA and crambin) and a simple statistical test. The statistical test used is called the potential scale reduction test, ${ }^{231}$ and it is a popular convergence diagnostic used by statisticians in order to monitor convergence in multiple Markov chain Monto Carlo simulations. ${ }^{232-235}$ Approximate convergence of a particular parameter measured from the simulations is assumed when the variance between the
different simulations for a particular parameter is no larger than the variance seen in that particular parameter from an individual simulation. ${ }^{235}$ This approach enables us to know which methodology is better (several long simulations or many short simulations) for achieving convergence of several geometric parameters. Eighteen helical parameters ${ }^{236}$ were measured from the DNA simulations and 16 parameters, including radius of gyration and hydrogen bonding distances, were measured from the crambin simulations. This approach will also give insight as to which parameters may take longer to stabilize over multiple simulations and how long they take to stabilize (effectively telling us how long we need to run the simulations in order to achieve convergence of a particular parameter thus obtain reliable results). To our knowledge, the potential scale reduction test has never been applied to biomolecular systems in order to assess convergence. Like other approaches, this test cannot diagnose absolute convergence, but the approach is very simple. All that is required are the values of the properties of interest over time; the statistical test (see Chapter 5 for details and formula) can easily be implemented in any statistical program.

In addition to gaining knowledge of the convergence of different properties, we have also discovered interesting and important information about which base pairs of the simulated DNA fragment should be included in the analysis of DNA simulations. For example, we have found that the effects of fraying (a well-known phenomenon where the end base pairs open and close rapidly $)^{237-240}$ penetrate beyond the end base pairs of the DNA. An understanding of sampling and convergence in biomolecular simulations will enable us to better model these systems and thus use
computational techniques to examine structural and dynamic features that cannot be observed by experiment.

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## Chapter 2

## Methods

This chapter discusses the background, methods, and procedures used for computations. The theoretical background of molecular dynamics is described in Section 2.1, the statistical tests used for analysis are described in Section 2.2, the theoretical background of quantum mechanical methods used is described in Section 2.3, and descriptions of the computations performed are in Section 2.4.

### 2.1 Molecular Dynamics

Molecular dynamics simulation provides a method for examination of molecular motion at the level of the atom and thus is a useful tool in chemistry and physics. ${ }^{1}$ MD is a classical simulation method, meaning that the motions of the particles in the system obey Newton's laws of motion and that electronic degrees of freedom are not explicitly described. Thus, one of the most important limitations of MD is that it cannot be used to describe phenomena involving bond breaking/formation events, excited states, or isomerizations. ${ }^{1}$ MD simulations enable the time dependence of properties of a molecular system (structural, dynamic, and thermodynamic) to be studied by numerically
solving Newton's equations of motion. A simple description of a simulation can be summarized by the following four steps: ${ }^{2}$

1. Initial coordinates are obtained (in the case of biomolecules, these are usually obtained from the Brookhaven Protein Data Bank ${ }^{3}$ or from the Rutgers Nucleic Acid Database ${ }^{4}$. Initial velocities are then assigned to the initial coordinates (usually assigned from some sort of distribution-Gaussian, Maxwell, etc.).
2. Compute the forces on acting on every particle in the system. This is the most time-consuming part of most simulations because the contribution of to the force on every particle due to all of its neighbors must be considered. If the system contains $N$ particles, and if no techniques are employed in order to speed up the evaluation of shortor long-range forces (use of nonbond cutoffs), the force evaluation scales as $N^{2}$.
3. Integrate Newton's equations of motion. After all forces are computed, Newton's equations of motion are integrated by breaking the calculation down into small time steps. Integration of Newton's equations is described in further detail below. Steps 2 and 3 make up the bulk of the simulation; they repeated until the time evolution of the system has been computed for the specified amount of time.
4. Determine desired averages and properties of the system. When the calculation is complete, the desired structural, dynamic, and thermodynamic properties can then be computed.
2.1.1 CHARMM Force Field. In molecular dynamics and molecular mechanics, the intra- and intermolecular forces between the atoms are described by an empirical force field. Empirical force fields use atomistic models, where atoms are the smallest particles in the system, rather than the electrons and nuclei used in quantum mechanics.

The empirical force field, as known as a potential energy function, allows for the potential energy $(V)$ of the system to be calculated as a function of the 3-D structure of the system $(R)$. The potential energy function consists of harmonic approximations of intramolecular (bond stretching, angle bending, and torsion interaction) terms and intermolecular (Coulombic and Van der Waals) terms (Eqn.2.1).

$$
V(R)_{\text {total }}=V(R)_{\text {intramolecular }}+V(R)_{\text {intermolecular }} \quad \text { Eqn. } 2.1
$$

The harmonic approximations used in the force field are sufficient for biomolecular simulations. Simulations of biomolecules are typically performed at temperatures where bond lengths and angles usually stay near their equilibrium values. Also, since there is no bond breakage/formation occurring in an MD simulation, a harmonic potential can accurately describe any bond and angle distortions.

The CHARMM force field, which is used in this dissertation, has the following form (Eqn.2.2).

$$
\begin{align*}
& V(R)=\sum_{\text {bonds }} K b\left(b-b_{0}\right)^{2}+\sum_{U B} K_{U B}\left(S-S_{0}\right)^{2}+\sum_{\text {angle }} K \theta\left(\theta-\theta_{0}\right)^{2}+ \\
& \sum_{\text {dihedrals }} K_{\chi}(1+\cos (n \chi-\delta))+\sum_{\text {impropers }} K_{\text {imp }}\left(\phi-\phi_{0}\right)^{2}+  \tag{Eqn. 2.2}\\
& \sum_{\text {nonbonded }} \varepsilon\left[\left(\frac{R \operatorname{minij}}{r_{i j}}\right)^{12}-\left(\frac{R \mathrm{mininij}^{6}}{r_{i j}}\right)^{6}\right]+\frac{q_{i} q_{j}}{\varepsilon_{1 r i j}}
\end{align*}
$$

The first five terms are the intramolecular terms, and the last two terms are the intermolecular terms. The parameters $K_{b}, K_{U B}, K_{\theta}, K_{\chi}$, and $K_{i m p}$ are the force constants and are usually determined experimentally or from quantum mechanical calculations. ${ }^{5}$ The terms in parentheses (e.g., $b-b_{0}$ ) describe the deviation $(b)$ of the parameter value from the equilibrium value $\left(b_{0}\right)$. For example, $\left(b-b_{0}\right)$ describes the deviation of the
current bond length from the equilibrium bond length (the other terms are analogous). The UB term is the Urey-Bradley term which describes 1,3-nonbonded interactions. In the dihedral term, $\chi$ is the torsion angle value, $\delta$ is the phase shift, and $n$ is the multiplicity. The first intermolecular term is the Lennard-Jones term, where $R_{\text {min,ij }}$ is the minimum interaction radius (this is dependent on the atoms interacting and is determined by experiment), $\varepsilon$ is the Lennard-Jones well depth, $1 / r^{6}$ describes the attractive interactions, and $1 / r^{12}$ describes the repulsive interactions. The second intermolecular term is the Coulombic term, where $q_{i}$ and $q_{j}$ are the partial charges on the interacting atoms, $r_{i j}$ is the distance between the interacting atoms, and $\varepsilon_{l}$ is the dielectric constant, which is generally treated as equal to 1 (permittivity of vacuum).

The 3-D structure (usually Cartesian coordinates obtained from the PDB or NDB in the case of biomolecules) and all parameters required in the force field are needed in order to begin the simulation. Once these are known, the energy of the system can then be calculated using Eqn.2.2. The first derivative of the energy with respect to the positions of the atoms yields the forces acting on the atoms, which is then used for the MD simulation (Eqn.2.3),

$$
\begin{equation*}
\frac{\partial V\left(x_{1}, x_{2}, \ldots x_{i} \ldots x_{N}\right)}{\partial x_{i}}=F\left(x_{1}, x_{2}, \ldots x_{i} \ldots x_{N}\right) \tag{Eqn. 2.3}
\end{equation*}
$$

where $x$ are the coordinates ( $x_{i}$ is the coordinate for the $i$ th particle and $N$ is the number of particles in the system).

There are currently a variety of existing force fields, and choice for which to use depends on the nature of the system. In this dissertation, the CHARMM force field is used for MD simulations, which is designed for use with biomolecular systems. ${ }^{6-8}$
2.1.2 Integration of Newton's equations. Once the 3-D coordinates are obtained, velocities are assigned to each particle in the system (these are usually assigned from a Gaussian or Maxwell distribution). The motion of the system is then simulated by integrating Newton's second law of motion,

$$
\begin{equation*}
F=m a \tag{Eqn. 2.4}
\end{equation*}
$$

where $F$ is the force acting on a particle (this is obtained from the potential energy function in Eqn.2.2), $m$ is the mass of the particle, and $a$ is its acceleration. The force can also be expressed as the gradient of the potential energy.

$$
\begin{equation*}
F=-\nabla V(R) \tag{Eqn. 2.5}
\end{equation*}
$$

Since acceleration is the second derivative of position with respect to time, this can be substituted in and Eqns. 2.4 and 2.5 can be combined to yield

$$
\begin{equation*}
-\frac{d V}{d R}=m \frac{d^{2} R}{d t^{2}} \tag{Eqn. 2.6}
\end{equation*}
$$

Newton's equation of motion can then relate the derivative of the potential energy to the changes in position as a function of time.

The potential energy is a function of all coordinates of all particles in the system. Numerical treatment is accomplished by using finite difference methods, which break the integration down into small timesteps ( $\Delta t$ ). The timestep for the simulation must be chosen so that the algorithm will accurately describe the fastest degrees of freedom (highest frequency motions) of the system. For biomolecular systems, the fastest motions are the stretching vibrations of hydrogens covalently bound to heavy atoms (this occurs on a timescale of $\sim 1 \mathrm{fs}$; therefore, the integration is only accurate for a timestep of $\sim 1-2$ fs). In order to use a timestep of 2 fs , the SHAKE algorithm ${ }^{9}$ is applied to covalent
hydrogen bonds in molecular dynamics simulations. SHAKE holds the covalent hydrogen bonds (e.g. C-H) constrained, eliminating the high frequency motions, and thus allows for a larger timestep. From the force, the accelerations of the particles can be calculated, and those are then combined with the positions and velocities at time $t$ to calculate the new positions and velocities at time $t+\Delta t$. The forces on the particles in the new positions are then calculated, which gives new positions and velocities at time $t+2 \Delta t$, etc.

There are several algorithms which are commonly used for integrating Newton's equations using finite difference methods. All of them assume that the positions, velocities, and accelerations of the particles in the system can be approximated by a Taylor series expansion. The most commonly used algorithm in simulations of biomolecules is the Verlet algorithm, which is based on two Taylor series expansions, a forward expansion $(t+\Delta t)$ and a backward $(t-\Delta t)$ expansion:

$$
\begin{align*}
& r_{n+1}=r_{n}+v_{n} \Delta t+\frac{1}{2}\left(\frac{F_{n}}{m}\right) \Delta t^{2}+O\left(\Delta t^{3}\right)  \tag{Eqn. 2.7}\\
& r_{n-1}=r_{n}+v_{n} \Delta t+\frac{1}{2}\left(\frac{F_{n}}{m}\right) \Delta t^{2}-O\left(\Delta t^{3}\right) \tag{Eqn. 2.8}
\end{align*}
$$

where $r_{n}$ indicates the position at step $n, n+1$ indicates the position at the next step, and $O\left(\Delta t^{n}\right)$ is the term of order $\Delta t^{n}$ or smaller. These two expansions can be summed together, which gives an algorithm for propagating the positions:

$$
\begin{equation*}
r_{n+1}=2 r_{n}-r_{n-1}+\left(\frac{F_{n}}{m}\right) \Delta t^{2}+O\left(\Delta t^{4}\right) \tag{Eqn. 2.9}
\end{equation*}
$$

The Verlet algorithm is carried out in two steps: use the current position $r_{n}$ to calculate the current force $F_{n}$ and then use the current and previous positions $r_{n}$ and $r_{n+1}$ together
with the current force $F_{n}$ to obtain the position in the next step, $r_{n+1}$. These two steps are repeated for every timestep for every particle in the system.

Because the velocities do not explicitly appear in the Verlet algorithm for computing the new positions, subtracting Eqn. 2.8 from Eqn. 2.7 yields an algorithm for propagating the velocities:

$$
\begin{equation*}
v_{n}=\frac{r_{n+1}-r_{n+1}}{2 \Delta t}+O\left(\Delta t^{2}\right) \tag{Eqn. 2.10}
\end{equation*}
$$

In this dissertation, the leap-frog integrator is used. The leap-frog integrator is a variation on the Verlet integrator; the modification was made to improve upon the large velocity errors of the velocity propagation of the Verlet algorithm. In the leap-frog integrator, velocities are determined at the mid-point of the position evaluation. The algorithm is written as:

$$
\begin{gather*}
r_{n+1}=r_{n}+v_{n+1 / 2} \Delta t  \tag{Eqn. 2.11}\\
v_{n+1 / 2}=v_{n-1 / 2}+\left(\frac{F_{n}}{m}\right) \Delta t \tag{Eqn. 2.12}
\end{gather*}
$$

where $v_{n \pm 1 / 2}$ is the velocity at the mid-step time, $t \pm(1 / 2) \Delta t$. The leap-frog integrator involves three steps: use the current position $r_{n}$ to calculate the current force $F_{n}$, use the current force and previous mid-step velocity $v_{n-1 / 2}$ to calculate the next mid-step velocity $v_{n+1 / 2}$, and finally use the current position and the next mid-step velocity to obtain the position in the next step $r_{n+1}$.
2.1.3 Nonbond Calculations. The expensive part of an MD simulation is the evaluation of nonbond forces. The number of bond, angle, and dihedral terms in the
force field are proportional to the number of particles in the system (the number of internal coordinates is $3 N-6$, which is linear in $N$ ); however, the number of nonbond terms increases on the order of $N^{2}$ (nonbond forces must be calculated between every pair of atoms in system). The non-bond attractive and repulsive terms of the Lennard-Jones potential decay quickly with distance; at $2.5 \sigma$ ( $\sigma$ is the distance at which the force between the two interacting particles is zero) the Lennard-Jones potential has only $1 \%$ of its value at $\sigma .{ }^{10}$ However, charge-charge interactions do not decay as rapidly with distance; the electrostatic energy is $\sim 1 / r$ for two point charges. In both instances, various algorithms are employed in order to reduce computational time., ${ }^{1,11}$
2.1.3.1. Use of Cutoffs. The most common way of reducing computational time and effort in the calculation of nonbond interactions is the use of nonbond cutoffs. These methods generally work by explicitly calculating the interaction energies between all pairs of particles within a specified distance (this is known as the cutoff distance and is usually set to $10-12 \AA$ for biomolecular systems) and approximating the contributions from the particles beyond the cutoff distance.

Abruptly truncating interactions at the cutoff distance causes discontinuities in the potential energy and in the force near the cutoff distance. ${ }^{10}$ This is obviously a problem, and methods have been developed to reduce the discontinuities. Two commonly used algorithms are known as switching and shifting functions. Both methods set the nonbond energies at zero beyond some distance value; however, the values for interactions at less than the cutoff distance, $r$, are treated differently. For a switching function $S(r), S(r)$ is a polynomial of $r$ that alters the nonbond energies smoothly and gradually over a buffer region $[a, b]$ so that the energy at $b=0$ and the energies at $r \leq a$ are not changed (the
buffer region is typically $1-2 \AA$ ). ${ }^{11}$ Shifting functions alter the nonbond energies more gradually over a larger region than switching function. This method underestimates short-range forces and alters the nonbond energies over the region $r \leq b$, rather than the smaller region that is affected in switching functions. ${ }^{11}$
2.1.3.2. Ewald Summation. Because long-range electrostatic interactions do not decay as rapidly as other nonbond interactions, long-range electrostatics can be problematic in MD. Until recently, cutoffs were applied to electrostatics (just like they were to other, more rapidly decaying nonbond interactions) in order to cut back on computational time. However, it was noticed that when cutoffs were used for electrostatic interactions, artifacts were introduced to the system, especially when the system contained many localized charges, as in the case of DNA. ${ }^{12,13}$ For example, it was found that even when explicit counterions were used in DNA simulations with periodic boundary conditions and cutoffs, substantial distortions were observed for the DNA structure. ${ }^{14}$ In simulations of polypeptides, the size of the cutoff distance was found to affect the stability of an $\alpha$-helix. ${ }^{15}$

Because of the artifacts caused by cutoffs, use of the Ewald summation, which is a more rigorous treatment of electrostatics, is required for highly charged systems, such as DNA. ${ }^{12,16}$ Studies have shown that use of the Ewald summation has led to more stable trajectories of DNA, RNA, and proteins. ${ }^{17-20}$ The Ewald summation is a technique for calculating the electrostatic energy of a system using periodic boundary conditions (simulated system is placed in a unit cell and is considered to have infinite images of itself in space). The method uses a combination of real and reciprocal space sums of interaction energies in order to evaluate long-range electrostatic energies. In
crystallography, the reciprocal lattice is an orthogonal system related to the orthogonal system associated with the atoms of the unit cell, which is the real space lattice. In the Ewald sum, two terms arise from the calculation of electrostatic energy; one term from pairs of atoms in the direct lattice (central simulation cell) and the other term from atom pairs corresponding to interactions with images of the central cell atoms. The Ewald sum is a periodic method; it includes a complete representation of electrostatic energies from an infinite number of images. Because it is a periodic method, the Ewald sum avoids the truncation effects in energy forces observed when using cutoff methods (and also requires use of periodic boundary conditions).

In the Ewald method, a particle interacts with all other particles in the simulation cell and with all of their images in the periodic system. There is thus a contribution to the total energy from interactions in the central simulation box along with the interactions between the central box and the periodic image boxes. The electrostatic contribution to the potential energy for all pairs of charges in the central simulation cell is written

$$
\begin{equation*}
V=\frac{1}{2} \sum_{i=1}^{N} \sum_{j=1}^{N} \frac{q_{i} q_{j}}{4 \pi \varepsilon_{0} r_{i j}} \tag{Eqn. 2.13}
\end{equation*}
$$

where $q_{i}$ is the partial charge on particle $i, \varepsilon_{0}$ is the dielectric constant, $N$ is the number of particles in the system, $r_{i j}$ is the minimum distance between particles $i$ and $j$, and $1 / 2$ is in front to avoid double-counting. The charge-charge contribution between the particles in the central cell and all the images of all the particles in the periodic system is written as

$$
\begin{equation*}
V=\frac{1}{2} \sum_{\mathbf{n}} \sum_{i=1}^{N} \sum_{j=1}^{N} \frac{q_{i} q_{j}}{4 \pi \varepsilon_{0}\left|r_{i j}+\mathbf{n}\right|} \tag{Eqn. 2.14}
\end{equation*}
$$

where $\mathbf{n}$ is a vector whose components are integral multiples of the length, $L$, of the central simulation box $\left(\mathbf{n}=n_{x} L, n_{y} L, n_{z} L\right)$. The sum in Eqn. 2.14 is very slow to converge, so the Ewald sum converts the summation to two series, both of which converge much more quickly. The split must take into account the slow decay of Coulomb interactions when $r$ is a large distance and also the variation Coulomb interactions when $r$ is a short distance. The split is of the form

$$
\frac{1}{r}=\frac{f(r)}{r}+\frac{1-f(r)}{r} \quad \text { Eqn. } 2.15
$$

where $f(r)$ must be a function to take both problems with Coulomb interaction into account. The first term is for real space and is short-ranged, and the second term is for reciprocal space and is a long-ranged term (this term can be Fourier-transformed). The Ewald sum uses the complementary error function $(\operatorname{erfc}(r))$ for $f(r)$.

$$
\begin{equation*}
\operatorname{erfc}(x)=\frac{2}{\sqrt{\pi}} \int_{x}^{\infty} \exp \left(-t^{2}\right) d t \tag{Eqn. 2.16}
\end{equation*}
$$

In the Ewald sum, each charge is considered to be surrounded by a neutralizing charge distribution which is of equal magnitude but opposite in sign to the charge. This is normally a Gaussian distribution of the form

$$
\begin{equation*}
\rho_{i}(r)=\frac{q_{i} \alpha^{3}}{\pi^{3 / 2}} \exp \left(-\alpha^{2} r^{2}\right) \tag{Eqn. 2.17}
\end{equation*}
$$

The sum is now the sum of the interactions between charges and the neutralizing Gaussians. This is the real space part of the Ewald sum.

$$
\begin{equation*}
V=\frac{1}{2} \sum_{i=1}^{N} \sum_{j=1}^{N} \sum_{\mid \mathbf{n}=0} \cdot \frac{q_{i} q_{j}}{4 \pi \varepsilon_{0}} \frac{\operatorname{erfc}\left(\alpha\left|r_{i j}+\mathbf{n}\right|\right.}{\left|r_{i j}+\mathbf{n}\right|} \tag{Eqn. 2.18}
\end{equation*}
$$

The summation with the prime indicates that the series does not include the interaction $i=j$ at $\mathbf{n}=0$. The new summation that uses the error function now converges quickly (and beyond some cutoff distance, the value is considered negligible). The rate of convergence depends on $\alpha$, the width of the neutralizing Gaussians, and the wider the Gaussian, the faster the sum converges.

Another charge distribution is then added to the system to counteract the first neutralizing distribution. This is the reciprocal space sum ( $\mathbf{k}$ are reciprocal vectors) and is a Fourier series.

$$
\begin{equation*}
V=\frac{1}{2} \sum_{k \neq 0} \sum_{i=1}^{N} \sum_{j=1}^{N} \frac{1}{\pi L^{3}} \frac{q_{i} q_{j}}{4 \pi \varepsilon_{0}} \frac{4 \pi^{2}}{k^{2}} \exp \left(-\frac{k^{2}}{4 \alpha^{2}}\right) \cos \left(\mathbf{k} \cdot \mathbf{r}_{i j}\right) \tag{Eqn. 2.19}
\end{equation*}
$$

The reciprocal space sum, like the real space sum, also converges much more quickly than the original sum. However, unlike the real space sum, the reciprocal space sum converges more quickly for a small $\alpha$ (narrower Gaussian). The number of terms that must be included in the reciprocal sum increases as $\alpha$ increases, which results in more time required for evaluation.

The sum of the neutralizing Gaussians in real space includes the interaction of each Gaussian with itself. Therefore, a term must be subtracted from the Ewald sum:

$$
\begin{equation*}
V=-\frac{\alpha}{\sqrt{\pi}} \sum_{k=1}^{N} \frac{q_{k}^{2}}{4 \pi \varepsilon_{0}} \tag{Eqn. 2.20}
\end{equation*}
$$

The sums described are convergent; however, they are convergent only for a system with a net charge of zero. If the system has a non-zero net charge, then an extra term must be added.

$$
\begin{equation*}
V=\frac{2 \pi}{3 L^{3}}\left|\sum_{i=1}^{N} \frac{q_{i}}{4 \pi \varepsilon_{0}} \mathbf{r}_{i}\right|^{2} \tag{Eqn. 2.21}
\end{equation*}
$$

The full Ewald summation is shown in Eqn. 2.22.

$$
V=\frac{1}{2} \sum_{i=1}^{N} \sum_{j=1}^{N}\left\{\begin{array}{l}
\sum_{\mathrm{n}=0 \mid}^{\infty} \frac{q_{i} q_{j}}{4 \pi \varepsilon_{0}} \frac{\operatorname{erfc}\left(\alpha\left|r_{i j}+\mathbf{n}\right|\right.}{\left|r_{i j}+\mathbf{n}\right|}  \tag{Eqn. 2.22}\\
+\sum_{k \neq 0} \frac{1}{\pi L^{3}} \frac{q_{i} q_{j}}{4 \pi \varepsilon_{0}} \frac{4 \pi^{2}}{k^{2}} \exp \left(-\frac{k^{2}}{4 \alpha^{2}}\right) \cos \left(\mathbf{k} \cdot \mathbf{r}_{i j}\right) \\
-\frac{\alpha}{\sqrt{\pi}} \sum_{k=1}^{N} \frac{q_{k}^{2}}{4 \pi \varepsilon_{0}}+\frac{2 \pi}{3 L^{3}}\left|\sum_{i=1}^{N} \frac{q_{i}}{4 \pi \varepsilon_{0}} \mathbf{r}_{k}\right|^{2}
\end{array}\right\}
$$

The Ewald method is computationally expensive to implement (it scales as $N^{2}$ ); however, methods have been devised to cut back on computational time. ${ }^{14,16}$ By optimizing $\alpha$, the relative rates of convergence of the real and reciprocal space sums can be adjusted to suit the system under study. The optimal balance between the sums will enable the Ewald sum to scale as $N^{3 / 2}$, which is still a considerable amount of computational time for large biomolecular systems. The particle mesh Ewald sum (PME) developed by Darden and coworkers reduces the Ewald sum to $N \log (N) .{ }^{16}$ In the PME method, the trigonometric function values of the Fourier series present in the reciprocal space sum are evaluated by a smooth approximation of the potential over a grid, or mesh. The smoothing function is known as the Euler spline, which expresses the value of the trigonometric function at the actual charge coordinates in terms of the charge value at
neighboring grid points. The resulting sums over the grid points are then evaluated efficiently by a fast Fourier transform. This dissertation uses the PME for all simulations; in the case of net-charged systems, a neutralizing background is used by default in the standard Ewald calculation. ${ }^{21}$

The Ewald sum was developed by P.P.Ewald in the $1920 \mathrm{~s},{ }^{22}$ long before the advent of computers, in order to compute the electrostatic energy of crystal. In the 1970s, it was applied to computer simulations of particles for precisely the same reason: to calculate efficiently the electrostatic energy of a periodic system (effectively, a crystal).
2.1.4. Boundaries. The treatment of boundaries and their effects is important to MD simulations because it allows macroscopic properties to be simulated using a relatively small number of particles. To be realistic, biomolecular simulations have to include some sort of description of a solvent environment. This modeling of bulk solvent is accomplished through the use of boundary conditions. There are various approaches to doing this, which include continuum boundary conditions, finite boundary conditions, and periodic boundary conditions.

With continuum boundary conditions, the solute (protein or DNA) is treated as a macroscopic object surrounded by a continuum which represents the solvent (water molecules are represented explicitly). With finite boundary conditions, the solute is surrounded by a layer of explicit water molecules while the bulk solvent molecules are modeled by some form of boundary potential at the solvent/vacuum border. ${ }^{1,12}$ These methods have the advantage of decreasing the number of particles that are explicitly represented and thus decrease computational expense; however, they have drawbacks. With continuum solvent methods, explicit interaction of the solute with water molecules
at the atomic level cannot be observed (e.g., interaction of specific water molecules with amino acids in a protein). ${ }^{12,23}$ With finite boundary conditions, water molecules cannot diffuse naturally due to the finite size of the system, and movements of water molecules can only be observed in the interior parts of the protein, such as binding cavities. ${ }^{12,23}$

Periodic boundary conditions allow simulations to be performed using a relatively small number of particles so that the particles in the system experience forces as if they were in bulk solvent. The simulated system (the solute and the solvent) is placed in a unit cell and is considered to have an infinite number of images in space. This replication of the unit cell forms an infinite lattice in three dimensions. In three dimensions, the atoms of the simulated system are contained within a unit cell analogous to the unit cell of a crystal. During the simulation, only the coordinates of the unit cell need are included. If an atom of the unit cell leaves the unit cell by crossing the boundary, then an image atom enters to replace it; therefore, the number of particles is conserved. In order to save computational time in evaluation of forces, the minimum image convention is customarily applied. With the minimum image convention, each atom only sees one image of every other atom in the system. The forces are calculated with only with closest periodic image.

The choice of geometry for the periodic cell is important. The cubic cell is the simplest periodic cell; however, it is appropriate to use a periodic cell that reflects the underlying geometry of the system. For example, a rectangular cell is probably best suited for a solvated DNA system, while a more spherical cell is better for a globular or spherical protein.

The simulations in this dissertation are all performed with periodic boundary conditions in the crystalline environment (crystal simulations allow for better comparison back to experiment when the experimental structure is a crystal structure ${ }^{24}$ ). The periodic cells and the dimensions used for the systems were those reported by the authors of the experimental crystal structures.

### 2.2 Analysis of Trajectories

2.2.1 RMSD. For a quantitative comparison of one structure to another (such as a simulation-generated structure to an experimental structure), some sort of similarity measurement is required. A commonly used method is to calculate the root-mean-square deviation (RMSD) between two structures or conformations. This is done by first superimposing the structures on top of one another so as to eliminate deviations caused by translation or rotation and then calculating the RMSD using Eqn. 2.23.

$$
\begin{equation*}
R M S D=\sqrt{\frac{1}{N} \sum_{k=1}^{N}\left(r_{k}^{i}-r_{k}^{j}\right)^{2}} \tag{Eqn. 2.23}
\end{equation*}
$$

where $N$ is the number of atoms in the structure, $k$ is the index over those atoms, and the $r_{k}$ 's are the coordinates of atom k in conformations $i$ and $j$. The summation in Eqn. 2.23 can include any set of atoms of the structure under consideration; it can include all the atoms in the structure or just a subset, such as the atoms of a ligand bound to a protein. In the analysis of biomolecular simulations, the summation usually includes only the heavy atoms (non-hydrogen atoms) and in the case of proteins and DNA, sometimes it only includes the alpha carbons of the protein backbone or the phosphate groups of the DNA backbone. RMSD is usually in the units of $\AA$.
2.2.2. T-tests. The t -test ${ }^{25}$ assess whether the means of two groups are statistically different from each other. The null hypothesis $\left(\mathrm{H}_{0}\right)$ of the test is that the means of the populations are equal. There are two forms of the $t$-test; one form assumes that the population standard deviation is the same for both sets of measurements being compared, and the other form assumes that the population standard deviations are different. In practice, both forms give about the same results (the numerical result of the test will fluctuate, but the overall significance of the test usually will not). In this dissertation, both forms of the t -test were used in order to compare the properties extracted from simulations to the experimental structure that was used as the starting point for the simulations. Eqns. 2.24 and 2.25 show both forms of the $t$-test.

$$
\begin{gather*}
t_{\text {calc }}=\frac{\overline{x_{1}}-\overline{x_{2}}}{s_{\text {pooled }}} \sqrt{\frac{n_{1} n_{2}}{n_{1}+n_{2}}}  \tag{Eqn. 2.24}\\
t_{\text {calc }}=\frac{\overline{x_{1}}-\overline{x_{2}}}{\sqrt{\left(s_{1}{ }^{2} / n_{1}\right)+\left(s_{2}{ }^{2} / n_{2}\right)}}
\end{gather*}
$$

Eqn. 2.25

Eqn. 2.24 is the form of the t-test that assumes that the population standard deviations are the same, and Eqn. 2.25 is the form that assumes they are different. In both equations, $s$ is the standard deviation for each population considered, $n$ is the number of samples used, and $x$ is the mean of each population. In Eqn. 2.24, $s_{\text {pooled }}$ is the pooled standard deviation (this makes use of both data sets) is calculated using Eqn. 2.26

$$
\begin{equation*}
s_{\text {pooled }}=\sqrt{\frac{s_{1}^{2}\left(n_{1}-1\right)+s_{2}^{2}\left(n_{2}-1\right)+\ldots+s_{k}^{2}\left(n_{k}-1\right)}{n_{1}+n_{2}+\ldots+n_{k}-k}} \tag{Eqn. 2.26}
\end{equation*}
$$

where $s$ is the standard deviation of each population considered, $n$ is the number of samples, and $k$ is the number of populations. When used in the t -test, $s_{\text {pooled }}$ (Eqn. 2.26)
involves two data sets (the t-test compares the means of two groups at a time, $k=2$ ); however, Eqn. 2.26 can be used to pool as many standard deviations as needed.

Once the $t_{\text {calc }}$ value is calculated, this number is compared with the $t_{\text {table }}$ number at the specified degrees of freedom and confidence level (the $95 \%$ confidence level is normally used). If $t_{\text {calc }}$ is greater than $t_{\text {table }}$ at the $95 \%$ confidence level, the results are considered to be statistically significantly different. The number of degrees of freedom is determined using Eqn. 2.27.

$$
d o f=n_{1}+n_{2}-2 \quad \text { Eqn. } 2.27
$$

In this dissertation, the dof was $\sim 100$ for the properties being analyzed. At the $95 \%$ confidence interval for 100 and greater $d o f$, the $t_{\text {table }}$ value is 1.96 , and this will be used throughout the document, unless otherwise noted. The specific details of what properties the $t$-test was applied to and how it was used (number of samples used, etc.) will be discussed with the simulation results.
2.2.3. Principal Components Analysis. Principal components analysis (PCA) ${ }^{26-29}$ is a multivariate analysis technique that is used to identify patterns in data and express the original data in such a way as to highlight similarities and differences. PCA transforms data to a new coordinate system such the greatest variance by any projection of the data comes to lie on the first coordinate (first principal component), the second variance lies on the second coordinate, and so on. PCA can be used to reduce the dimensionality of a data set by retaining those characteristics of the data that contribute most to the variance of the data set. The lower-order principal components are kept, while the higher-order ones are ignored because the lower order ones contribute most to the variance.

The procedure by which PCA works is typically described using an example data set. Suppose that this data set consists of property with $p$ measurements over $n$ conformations for multiple MD simulations. This data set is matrix $\mathbf{X}$ with $p$ columns and $n$ rows. The purpose of performing PCA in this dissertation was to see the relationship of a particular property among the all conformations from different simulations and their relationship to an experimental structure that was used to initialize the simulations (i.e., which conformations among the simulations are similar or different to each other and the experimental structure). PCA plots were also used to select starting structures for additional simulations. The data set is multidimensional, and the whole point of PCA is to reduce dimensionality in order to examine the relationships of the different conformations from the simulations. First, for PCA to work properly, the mean of each dimension of the data set must be subtracted from each dimension (this produces a data set whose mean is zero). Next, a covariance matrix $\mathbf{C}$ is constructed. Covariance is measured between two dimensions, and it is used to find out how much the dimensions of the data vary from the mean with respect to each other (variance is the analogous measure in one dimension). The covariance between two dimensions $a$ and $b$, would be calculated by Eqn. 2.28.

$$
\begin{equation*}
\operatorname{cov}(a, b)=\frac{\sum_{i=1}^{n}\left(a_{i}-\bar{a}\right)\left(b_{i}-\bar{b}\right)}{(n-1)} \tag{Eqn. 2.28}
\end{equation*}
$$

If the data set contains more than two dimensions, obviously more than one covariance measurement can be calculated. For an $n$ dimensional data set, $\mathbf{C}$ will be an $n \times n$ matrix, and each element is the result of calculating the covariance between two separate
dimensions (Eqn. 2.28). The diagonal of $\mathbf{C}$ will contain the covariance between one of the dimensions and itself, which are just the variances for that dimension. Also, since $\operatorname{cov}(a, b)=\operatorname{cov}(b, a), \mathbf{C}$ is symmetric about the diagonal.

The next step is to diagonalize $\mathbf{C}$. By doing this, a matrix of eigenvectors $\mathbf{E}$ and their corresponding eigenvalues. For an $n \times n$ covariance matrix, $n$ eigenvectors will be obtained. All the eigenvectors of $\mathbf{C}$ are orthogonal to each other (as with any matrix), and this is important because it means that we will be able to express the original data in terms of these orthogonal eigenvectors rather in terms of the usual $x, y, a n d z$ axes (these will not show how each data point relates to the rest of the data). The eigenvalues that are associated with each eigenvector indicate the magnitude of the variances in the direction of their corresponding eigenvectors.

Once the eigenvectors and eigenvalues are found from $\mathbf{C}$, the next step is to sort them in order by eigenvalue, highest to lowest. The sorting can be visualized by using a Scree plot, which simply shows which components contribute most to the variance of the data set (the percentage that each component contributes to the variance of the data set is plotted versus the component number). Usually, only the first 1 or 2 components will contribute most to the variance, and the lower components can be ignored. Some information will be lost when components are ignored; however, if the eigenvalues are small, not much information is lost, and the dimensionality of the data set is reduced (if the data set has $n$ dimensions, and only the first $p$ eigenvectors are retained, the final data set will only contain $p$ dimensions). Another matrix $\mathbf{E}_{\mathbf{1}}$ is then formed from the chosen eigenvectors.

The final step in PCA is to project the original data set onto the new set given by the chosen eigenvectors. The new components are the principal components (PCs), and while the original data contained $n$ dimensions and may have been correlated, the new data set contains only $p$ dimensions (if p eigenvectors were retained), and the PCs are uncorrelated. The final data set (which is made up of the PCs) is computed by the following equation,

$$
M=E_{1}^{\mathrm{T}} \times X^{\mathrm{T}} \quad \text { Eqn. } 2.29
$$

where $\mathbf{M}$ is the final data matrix, $\mathbf{E}_{\mathbf{1}}{ }^{\mathbf{T}}$ is the transposed matrix of the chosen eigenvectors, and $\mathbf{X}^{\mathbf{T}}$ is the transpose of the original data matrix with the means subtracted out. To get the data in a plot-friendly format, it helps to transpose $\mathbf{M}$. The original data is now solely in terms of the chosen eigenvectors, and if not all eigenvectors were used, the dimensions of the data have been reduced. The PCs can now be plotted against one another in order to visualize the projections of the data onto the eigenvectors with the largest variances.

In this dissertation, PCA was applied to DNA/polyamide complex data sets. The S-plus program ${ }^{30}$ was used to perform PCA.

### 2.3 Quantum Mechanical Methods

2.3.1 Schrödinger Equation. The time-independent solution to the Schrödinger equation for a molecular system can be expressed by the following equation

$$
\begin{equation*}
H \Psi=E \Psi \tag{Eqn. 2.30}
\end{equation*}
$$

where $\Psi$ is the wavefunction, $E$ is the energy of the system, and $H$ is the Hamiltonian operator, which is equal to

$$
\begin{equation*}
H=\frac{-h^{2}}{8 \pi^{2} m} \nabla^{2}+V \tag{Eqn. 2.31}
\end{equation*}
$$

where $h$ is Planck's constant, $m$ is the mass of the particle, $\nabla$ is the Laplacian operator, and $V$ is the potential energy. The Hamiltonian is made up of the sum of the kinetic and potential energy contributions of the nuclei and electrons in the system. The kinetic energy is a summation of $\nabla^{2}$ over all the particles in the molecule (first part of Eqn. 2.31):

$$
\begin{equation*}
K E=\frac{-h^{2}}{8 \pi^{2}} \sum_{k} \frac{1}{m_{k}}\left(\frac{\partial^{2}}{\partial x_{k}^{2}}+\frac{\partial^{2}}{\partial y_{k}^{2}}+\frac{\partial^{2}}{\partial z_{k}^{2}}\right) \tag{Eqn. 2.32}
\end{equation*}
$$

The potential energy component is the Coulomb repulsion between each pair of charged particles:

$$
\begin{equation*}
V=\frac{1}{4 \pi \varepsilon_{0}} \sum_{j} \sum_{k<j} \frac{e_{j} e_{k}}{\Delta r_{j k}} \tag{Eqn. 2.33}
\end{equation*}
$$

$\Delta r_{j k}$ is the distance between the two particles $j$ and $k$, and $e_{j}$ and $e_{k}$ are the charges on the particles. For an electron, the charge is $-e$, and for a nucleus, the charge is $Z e(Z$ is the atomic number of the atom).

Therefore, we can write:

$$
\begin{equation*}
V=\frac{1}{4 \pi \varepsilon_{0}}\left(-\sum_{i} \sum_{I}\left(\frac{Z_{I} e^{2}}{\Delta r_{i I}}\right)+\sum_{i} \sum_{j<i}\left(\frac{e^{2}}{\Delta r_{i j}}\right)+\sum_{i} \sum_{J<I}\left(\frac{Z_{I} Z_{J} e^{2}}{\Delta R_{I J}}\right)\right) \tag{Eqn. 2.34}
\end{equation*}
$$

where the first term is for the nuclei-electron attraction, the second term is for electronelectron repulsion, and the third term is for nuclear-nuclear repulsion.

Under the Born-Oppenheimer approximation, ${ }^{31}$ nuclear and electronic motions can be separated, which simplifies the Schrödinger equation. This approximation is reasonable because the mass of a nucleus is thousands of times greater than that of an electron, and nuclei move very slowly with respect to the electrons (electrons move instantaneously in response to changes in nuclear position). Therefore, the nuclei are fixed relative to the electrons, and the motion of the electrons can be considered to occur in a fixed potential produced by the nuclei.

$$
\begin{equation*}
H=K E^{\text {elec }}(\stackrel{\rightharpoonup}{r})+K E^{\text {nuclear }}(\stackrel{\rightharpoonup}{R})+V^{\text {nuclear-elec }}(\stackrel{\rightharpoonup}{R}, \stackrel{\rightharpoonup}{r})+V^{\text {elec }}(\stackrel{\rightharpoonup}{r})+V^{\text {nuclear }}(\stackrel{\rightharpoonup}{R}) \tag{Eqn. 2.35}
\end{equation*}
$$

Where $\vec{r}$ and $\vec{R}$ designate the positions of the electrons and nuclei, respectively. The Born-Oppenheimer approximation says that the two parts (electronic and nuclear) can be solved independently, so a Hamiltonian that omits the nuclear kinetic energy term can be written:

$$
\begin{align*}
& H^{\text {elec }}=-\frac{1}{2} \sum_{i}\left(\frac{\partial^{2}}{\partial x_{i}^{2}}+\frac{\partial^{2}}{\partial y_{i}^{2}}+\frac{\partial^{2}}{\partial z_{i}^{2}}\right)-\sum_{i} \sum_{I}\left(\frac{Z_{I}}{\left|\vec{R}_{I}-\vec{r}_{i}\right|}\right)  \tag{Eqn. 2.36}\\
& +\sum_{i} \sum_{j<i}\left(\frac{1}{\left|\vec{r}_{i}-\overrightarrow{r_{j}}\right|}\right)+\sum_{I} \sum_{J<I}\left(\frac{Z_{I} Z_{J}}{\left|\overrightarrow{R_{I}}-\overrightarrow{R_{J}}\right|}\right)
\end{align*}
$$

Using this electronic Hamiltonian, a Schrödinger equation is obtained for describing the motion and energy of the electrons in a fixed field of nuclei, where $n$ is the number of electrons.

$$
\begin{equation*}
H^{e l e c} \Psi^{n}(\stackrel{\rightharpoonup}{r}, \stackrel{\rightharpoonup}{R})=E^{n}(\stackrel{\rightharpoonup}{R}) \Psi^{n}(\stackrel{\rightharpoonup}{r}, \stackrel{\rightharpoonup}{R}) \tag{Eqn. 2.37}
\end{equation*}
$$

The Schrödinger equation cannot be solved exactly with the exception of one-electron systems (such as the hydrogen atom). ${ }^{32,33}$ The electron-electron repulsion term in Eqn. 2.36 makes it impossible to find an analytical solution to the Schrödinger equation. Therefore, some sort of approximation method, such as perturbation theory, is necessary to obtain a solution.

The product of the wavefunction $\Psi$ with its complex conjugate $\left(\Psi^{*} \Psi\right)$ is interpreted the probability of finding the particle at some point in space. The square of $\Psi$ thus gives the electron density at any given point. Therefore, $\Psi$ must be normalized; if the probability of finding the particle is integrated over all space, the result must be 1 (the particle has to be somewhere).

$$
\begin{equation*}
\int \Psi * \Psi d \tau=1 \tag{Eqn. 2.38}
\end{equation*}
$$

$d \tau$ indicates integration over all space. Aside from the normality requirement, the wavefunction has the requirement of being antisymmetric, which means that $\Psi$ changes sign when two identical particles are interchanged (this follows from the fact that electrons are indistinguishable from one another, and there is a fundamental symmetry that the wavefunction must obey in describing the behavior of many electrons).

The solutions (the wavefunctions, $\Psi$ ) to the Schrödinger equation (Eqn. 2.37) describe some stationary state of the system (the solutions are called stationary-state wave functions because they are independent of time). The wavefunction depends on both the position of the particle ( $\mathrm{x}, \mathrm{y}, \mathrm{z}$ coordinates) and their spin states. Solving Eqn. 2.37 gives $\Psi^{\mathrm{n}}$ and $E^{n}$, where the $\Psi \mathrm{s}$ are the eigenfunctions of the Hamiltonian operator
and the $\mathrm{E}^{\mathrm{n}} \mathrm{s}$ are the eigenvalues (the $E^{n} \mathrm{~s}$ are the allowed energy values). $E_{0}{ }^{n}$ is the energy of the ground-state energy of the system, while all the $E^{n}$ s above that are the energies of the excited states.
2.3.2 Hartree-Fock Theory. Hartree-Fock (HF) is one of the most common approximations for determining the ground-state wavefunction and ground-state energies for a many-electron system. ${ }^{32-34}$ There is no correct solution for a many-electron system; there must be some way to determine if one wavefunction is better than another one. Hartree-Fock utilizes the variational principle, which says that for a time-independent Hamiltonian operator, any trial wavefunction $\phi$ will have an energy expectation value that is greater than or equal to the true ground state wavefunction corresponding to the given Hamiltonian.

$$
\begin{equation*}
\int \frac{\phi^{*} \hat{H} \phi d \tau}{\phi^{*} \phi d \tau} \geq E_{g s} \tag{Eqn. 2.39}
\end{equation*}
$$

Therefore, the Hartree-Fock energy is an upper bound to the true ground state energy of a given system.

The next approximation made by Hartree-Fock theory is that the trial wavefunctions to be used are approximated by a single Slater determinant. A Slater determinant is the simplest form of an orbital wavefunction that satisfies the asymmetry requirement. Because asymmetry must be considered, electron spin has to be taken into account. Electron spin can be conveniently treated by combining spatial orbitals with the spin functions $\alpha$ and $\beta$, resulting in a spin orbital that a function of location (spatial part) and spin (spin part). Eqn. 2.40 is an example of a Slater determinant

$$
\Psi=\frac{1}{\sqrt{N!}}\left(\begin{array}{lll}
\phi_{1}(1) \alpha & \phi_{1}(1) \beta \ldots & \phi_{n}(1) \beta  \tag{Eqn. 2.40}\\
\phi_{1}(2) \alpha & \phi_{1}(2) \beta \ldots & \phi_{n}(2) \beta \\
\ldots & \ldots & \ldots \\
\phi_{1}(n) \alpha & \phi_{1}(n) \beta \ldots & \phi_{n}(n) \beta
\end{array}\right)
$$

where $\frac{1}{\sqrt{N!}}$ is a factor that ensures that the wavefunction is normalized, $\phi_{i}(n)$ is the spatial orbital for the atom $n$, and $\alpha$ and $\beta$ are the spin functions.

The Hartree-Fock equation for each spatial orbital is

$$
\begin{equation*}
F_{i} \phi_{i}=\varepsilon_{i} \phi_{i} \tag{Eqn. 2.41}
\end{equation*}
$$

where $\phi_{i}$ is the wavefunction, $\varepsilon_{i}$ is the energy of the orbital, and $F_{i}$ is the Fock operator. The Fock operator has the form

$$
\begin{equation*}
F_{i}=-\frac{1}{2} \nabla_{i}^{2}-\sum_{A=1}^{M} \frac{Z_{A}}{r_{i a}}+v^{H F}(i) \tag{Eqn. 2.42}
\end{equation*}
$$

where $v^{H F}(i)$ is the average potential felt by the $i$ th electron due to the other electrons; this average potential is approximation to the third term of Eqn. 2.36.

When a Hartree-Fock calculation is started, neither $F_{i}$ nor $\phi_{i}$ are known. An initial guess is made for the $\phi_{i}$, which in turn allows calculation of $F_{i}$. By using the variational method, the $\phi_{i}$ and $F_{i}$ are optimized in order to approximate ground state energy for the system. This iterative procedure is known as the self-consistent field (SCF) method.

The largest drawback of Hartree-Fock theory is that it does not include effects of electron correlation. The electrons are considered to be moving in an average potential of the other electrons; thus, electrons are not specifically influenced by the presence of individual neighboring electrons. However, despite the lack of electron correlation, HF
does allow for quantum calculations to be applied to larger and more realistic systems, because it is less computationally demanding. HF is also a useful starting point for more advanced quantum mechanical methods.
2.3.3. Perturbation Theory. Because Hartree-Fock does not include effects of electron correlation, it will represent many properties incorrectly, such as bond formation/breakage. ${ }^{34,35}$ Methods that go beyond Hartree-Fock to include the effects of electron correlation are needed; these are known as post-Hartree-Fock or post-SCF methods.

One method of treating electron correlation is Møller-Plesset perturbation theory. This method estimates the electron correlation energy by splitting the Hamiltonian operator into two parts. The perturbation Hamiltonian is the difference between the exact and Hartree-Fock Hamiltonians. The true Hamiltonian $H$ is expressed as the sum of a zeroeth order Hamiltonian $H_{0}$ (for which a set of molecular orbitals can be obtained and therefore can be solved exactly) and a perturbation ( $\lambda V$ ), which is assumed to be small in comparison to $H_{0}$.

$$
\begin{equation*}
H=H_{0}+\lambda V \tag{Eqn. 2.43}
\end{equation*}
$$

The eigenfunctions of the true Hamiltonian operator are $\Psi_{i}$ with corresponding eigenvalues $E_{i}$. The eigenfunctions of the zeroeth order Hamiltonian are $\Psi_{\mathrm{i}}{ }^{0}$ with corresponding energies $E_{i}{ }^{0}$. The ground state wavefunction and energy are therefore $\Psi_{0}{ }^{0}$ and $E_{0}{ }^{0}$. The parameter $\lambda$ can vary from 0 to 1 ; when $\lambda=0$, the $H$ is equal to the zeroeth order

Hamiltonian, but when $\lambda=1, H$ is equal to its true value. The eigenfunctions $\Psi_{\mathrm{i}}$ and eigenvalues $E_{i}$ of $H$ can then be expressed as powers of $\lambda$ :

$$
\begin{align*}
& \Psi_{i}=\Psi_{i}^{0}+\lambda \Psi_{i}^{1}+\lambda^{2} \Psi_{i}^{2}+\lambda^{3} \Psi_{i}^{3}+\ldots  \tag{Eqn. 2.44}\\
& E_{i}=E_{i}^{0}+\lambda E_{i}^{1}+\lambda^{2} E_{i}^{2}+\lambda^{3} E_{i}^{3}+\ldots
\end{align*}
$$

Eqn. 2.45

The perturbed wavefunction and energy are then substituted into the full Schrodinger equation $\left(H \Psi_{i}=E_{i} \Psi_{i}\right)$ and like powers of $\lambda$ are equated to give the zeroeth (Eqn. 2.46), first (Eqn. 2.47), and second orders (Eqn. 2.48) of perturbation:

$$
\begin{gather*}
H_{0} \Psi_{i}^{0}=E_{i}^{0} \Psi_{i}^{0}  \tag{Eqn. 2.46}\\
\left(H_{0}-E_{i}^{0}\right) \Psi_{i}^{1}=\left(E_{i}^{1}-V\right) \Psi_{i}^{0}  \tag{Eqn. 2.47}\\
\left(H_{0}-E_{i}^{0}\right) \Psi_{i}^{2}=E_{i}^{2} \Psi_{i}^{0}+E_{i}^{1} \Psi_{i}^{1}-V \Psi_{i}^{1} \tag{Eqn. 2.48}
\end{gather*}
$$

Solving these equations yields the following for $E_{i}$ and $\Psi_{i}$ :

$$
\begin{align*}
& E_{i}=E_{i}^{0}+\lambda H_{i i}^{1}+\lambda^{2} \sum_{j \neq i} \frac{H_{i j}^{1} H_{j i}^{1}}{E_{i}^{0}-E_{j}^{0}}+\ldots  \tag{Eqn. 2.49}\\
& \Psi_{i}=\Psi_{i}^{0}+\lambda \sum_{j \neq 1} \frac{H_{j i}^{1}}{E_{i}^{0}-E_{j}^{0}} \Psi_{j}^{0}+\ldots \tag{Eqn. 2.50}
\end{align*}
$$

Where $H_{i j}^{1}=\int \Psi_{i}^{0} V \Psi_{j}^{0} d \tau$. The first-order correction to the energy can be calculated by using the unperturbed wavefunction and the perturbed Hamiltonian:

$$
\begin{equation*}
E_{i}^{1}=H_{i i}^{1}=\int \Psi_{i}^{0} V \Psi_{i}^{0} d \tau \tag{Eqn. 2.51}
\end{equation*}
$$

The higher-order terms are developed similarly; for example, the second-order correction is given by:

$$
\begin{equation*}
E_{i}^{2}=\sum_{j \neq i} \frac{\left(\int \Psi_{i}^{0} V \Psi_{j}^{0} d \tau\right)^{2}}{E_{i}^{0}-E_{j}^{0}} \tag{Eqn. 2.52}
\end{equation*}
$$

So far, these are the results of general perturbation theory. In the particular case of Møller-Plesset theory, the Hartree-Fock wavefunction is used as the zeroeth order approximation, and the difference between the true electron-electron repulsions and the averaged ones from HF is treated as the perturbation (the perturbation Hamiltonian is the difference between the exact and Hartree-Fock Hamiltonians). The zeroeth order Hamiltonian is a sum of Fock operators, and the zeroeth order energy is sum of orbital energies.

$$
\begin{align*}
& H_{0}=\sum_{i} F_{i}  \tag{Eqn. 2.53}\\
& E_{i}^{0}=\sum_{i} E_{i} \tag{Eqn. 2.54}
\end{align*}
$$

The first-order energy is given by Eqn. 2.54, and adding $E_{i}{ }^{0}$ to $E_{i}{ }^{l}$ gives the full HartreeFock energy. Therefore, to get any treatment of electron correlation, the second-order perturbation must be incorporated. The second-order energy and the first-order correction to the wavefunction are given by:

$$
\begin{align*}
& E_{i}^{2}=\sum_{j} \frac{\left|\int \Psi_{i}^{0} V \Psi_{j} d \tau\right|^{2}}{E_{i}^{0}-E_{j}}  \tag{Eqn. 2.55}\\
& \Psi_{i}^{1}=\sum_{j} \frac{\int \Psi_{i}^{0} V \Psi_{j}^{0} d \tau}{E_{i}^{0}-E_{j}} \Psi_{j} \tag{Eqn. 2.56}
\end{align*}
$$

where the sum is over all excited states, denoted as $\Psi_{j}$ and $E_{j}$. The $E_{i}^{2}$ numerator is always positive because it is squared, while the denominator is negative because $\Psi_{i}^{0}$ is the ground state and $\Psi_{j}$ is an excited state. Therefore, the second-order energy correction is always negative or zero. Higher order corrections may be either positive or negative.

The higher orders of Møller-Plesset theory perturbation are denoted as MPn, where $n$ is the order of perturbation. The MP1 energy is the same as the Hartree-Fock energy; MP2 includes the effects of double excitation and is the most practical treatment for electron correlation. ${ }^{34,35}$ Third, fourth, and higher orders of perturbation are derived similarly to second-order. The terms quickly become more mathematically complicated and computationally expensive. For $n$ basis functions, HF scales as $n^{4}$ and MP2 scales as $n^{5}$; for higher orders of perturbation, the calculations become more costly. MP2 calculations require more computational effort; however, MP2 yields a significant correction to geometries and energies over the corresponding HF calculations. ${ }^{34,35}$
2.3.4. Density Functional Theory. The central idea in density functional theory (DFT) is that a relationship exists between the total electronic energy and the overall electron density. DFT methods are based on the ideas of Hohenberg and Kohn, ${ }^{36}$ who stated that the ground state electronic energy of a molecule can be expressed exactly as a functional of the electron density of the molecule. The term functional refers to a function of a function, which in this case means that the total energy has a functional dependence on electron density, which in turn is dependent on the coordinates of the electrons of the system. The approximate functionals used in DFT methods use the following separation of the total electronic energy.

$$
\begin{equation*}
E=E^{T}+E^{V}+E^{J}+E^{X C} \tag{Eqn. 2.57}
\end{equation*}
$$

$E^{T}$ is the kinetic energy term arising from electron motion, $E^{V}$ is the potential energy term and includes nuclear-electron attraction and nuclear-nuclear repulsion, $E^{J}$ is electronelectron repulsion term (also described as the Coulomb self-interaction of the electron density), and $E^{X C}$ is the exchange-correlation term and includes the contributions due to electron exchange and correlation. All terms of Eqn. 2.57 with the exception the nuclearnuclear repulsion are functions of the electron density, $\rho$.

Eqn. 2.57 follows the ideas of Kohn and Sham, ${ }^{37}$ who suggested that the sum of the kinetic energy of the electrons and the contributions from electronic interactions should be expressed as a sum of three terms: the kinetic energy (which is defined as the kinetic energy of a system with non-interacting electrons), the electron-electron repulsion energy, and the electron exchange-correlation energy. In practice, Kohn-Sham DFT calculations are performed using an iterative approach, analogous to the SCF approach used in HF calculations. An initial guess is made for $\rho$, which allows for derivation of a set of orbitals that leads to an improved value for $\rho$. The improved $\rho$ value is then used in the second iteration and so on until convergence is reached. ${ }^{33}$

The exchange term in Eqn. 2.57 represents the effects of electron exchange and correlation on the total energy of the system, which is an advantage over HF, which does not treat effects of electron correlation. However, there is no known exact expression for $E^{X C}$ (Hohenberg and Kohn showed that $E^{X C}$ is dependent entirely on $\rho$; however, the theorem does not provide the form of this functional). ${ }^{33,36,37}$ A number of approximate expressions for $E^{X C}$ have been developed, which in turn lead to a variety of methods. In
practice, $E^{X C}$ is usually divided into separate parts, which are the exchange and correlation parts.

$$
\begin{equation*}
E^{X C}(\rho)=E^{X}(\rho)+E^{C}(\rho) \tag{Eqn. 2.58}
\end{equation*}
$$

All terms are again functionals of $\rho$. The functionals on the right side of the equation are the exchange and correlation functionals, respectively. Both parts can be one of two types, a local functional (depends only on $\rho$ ) or a gradient-corrected functional (depends on $\rho$ and the gradient of $\rho, \nabla \rho)$.

The local exchange functional is almost always defined as

$$
\begin{equation*}
E_{L D A}^{X}=-\frac{3}{2}\left(\frac{3}{4 \pi}\right)^{1 / 3} \int \rho^{4 / 3} d^{3} \vec{r} \tag{Eqn. 2.59}
\end{equation*}
$$

where $\vec{r}$ represents the coordinates of the electrons. This is known as the local density approximation (LDA), and it states that in regions of the system where charge density varies slowly, the exchange correlation energy can be considered the same as that of a uniform electron gas of the same charge density. ${ }^{35}$ The LDA approximation is generally not of high enough accuracy to be useful for determining structural properties or dissociation energies of molecules. ${ }^{33,35}$ However, the quality of the results can be improved by adding correction terms to Eqn. 2.59 that depend on the gradient of $\rho$.

The functional developed by Becke in $1988^{38,39}$ makes a correction to the LDA approximation that improves upon many of its deficiencies. The following gradientcorrected exchange functional is Becke's form, which is now commonly used,

$$
\begin{equation*}
E_{\text {Becke }}{ }^{x}=E_{L D A}{ }^{x}-\gamma \int \frac{\rho^{4 / 3} x^{2}}{\left(1+6 \gamma \sinh ^{-1} x\right)} d^{3}{ }^{3} \tag{Eqn. 2.60}
\end{equation*}
$$

where $x^{2}=\rho^{-4 / 3}|\nabla \rho|$, and $\gamma$ is a parameter chosen to fit the known exchange energies of inert gas atoms (Becke defines $\gamma=0.0042$ Hartrees). ${ }^{38}$

As for the exchange functional, corrections (local and gradient) exist for the correlation functional. A commonly used method is that of Lee, Yang, and Parr, ${ }^{40}$
$E_{L Y P}^{C}=-a \int \frac{1}{1+d \rho^{-1 / 3}}\left\{r+b \rho^{-2 / 3}\left[c_{F} \rho^{5 / 3}-2 t_{W}+\left(\frac{1}{9} t_{W}+\frac{1}{18} \nabla^{2} \rho\right) e^{-c r^{-1 / 3}}\right]\right\} d \vec{r}$
where $a, b, c$, and $d$ are constants and $t_{W}(\vec{r})=\sum_{i=1}^{N} \frac{\left|\nabla \rho_{i}(\vec{r})\right|^{2}}{\rho_{i}(\vec{r})}-\frac{1}{8} \nabla^{2} \rho$. Usually, DFT methods are formed by combining an exchange functional with a correlation functional. For example, the BLYP functional is formed by combining Becke's gradient-corrected exchange functional with Lee, Yang, and Parr's gradient-corrected correlation functional.

The key feature of DFT methods is the way in which electron exchange and correlation effects are directly incorporated; correlation effects are truly only considered in more complex, post-HF methods, such as MP2. Because DFT methods include both exchange and correlation effects, higher accuracy is achieved compared to Hartree-Fock, which does not include electron correlation effects. Despite the improvements in DFT in recent years, there are still difficulties in using density functional theory to describe properly certain intermolecular interactions, in particular, van der Waals (dispersion) forces. The poor treatment of dispersion by DFT renders this method unsuitable for the treatment of systems that are dominated by dispersion forces, such as noble gases. ${ }^{33}$
2.3.5. Basis Sets. A basis set is a set of functions used to create the molecular orbitals of a molecule. The functions that make up the basis set are known as basis
functions, and they are expanded as linear combinations to make up the basis set. ${ }^{33}$ The basis functions are usually centered on the atomic nuclei, so they represent atomic orbitals. The accuracy of quantum mechanical calculations depends largely on the basis set chosen. The basis set specifies which and how many basis functions are used to describe a molecule.

An individual molecular orbital can be defined as

$$
\begin{equation*}
\phi_{i}=\sum_{\mu=1}^{N} c_{\mu i} \chi_{\mu} \tag{Eqn. 2.62}
\end{equation*}
$$

where the coefficients $c_{\mu i}$ are known as the molecular orbital expansion coefficients, which are determined computationally and $\chi_{\mu}$ are the basis functions (these are chosen to be normalized). Gaussian $03,{ }^{41}$ which was used in this dissertation for quantum mechanical calculations, uses Gaussian-type atomic functions as basis functions. Gaussian functions have the general form

$$
\begin{equation*}
g(\alpha, \stackrel{\rightharpoonup}{r})=c x^{n} y^{m} z^{l} e^{-\alpha r^{2}} \tag{Eqn. 2.63}
\end{equation*}
$$

where $\vec{r}$ is made up of $x, y$, and $z$, and $\alpha$ is a constant that determines the size, or spread, of the function (a large $\alpha$ value gives a function that does not spread very far; a small $\alpha$ value will give a large spread). In a Gaussian function, $e^{-\alpha r^{2}}$ is multiplied by powers of $x$, $y, z$, and normalization constant, so that when the square of the Gaussian function is integrated over all space, the result is 1 . The parameter $c$ depends on $\alpha, l, m$, and $n$. Linear combinations of primitive Gaussian functions of the form in Eqn. 2.63 are used to form the basis functions. These are known as contracted Gaussians and have the form

$$
\begin{equation*}
\chi_{\mu}=\sum_{p} d_{\mu p} g_{p} \tag{Eqn. 2.64}
\end{equation*}
$$

where $d_{\mu p}$ are fixed constants within a given basis set. The basis functions are combined to yield the linear expansion that forms the molecular orbital (Eqn. 2.65).

$$
\begin{equation*}
\phi_{i}=\sum_{\mu} c_{\mu i} \chi_{\mu}=\sum_{\mu} c_{\mu i}\left(\sum_{p} d_{\mu p} g_{p}\right) \tag{Eqn. 2.65}
\end{equation*}
$$

In this dissertation, the $6-31 \mathrm{G}(\mathrm{d})$ basis set was used in all quantum mechanical calculations. This is a combination of split-valence and polarized basis sets. Valence electrons principally take part in chemical bonding; therefore, valence orbitals are represented by more than one basis function, which in turn are composed of a linear combination of Gaussian primitive functions. Split-valence basis sets have two or more sizes of basis function for each valence orbital. Split-valence basis sets allow orbitals to change size but not shape. Polarized basis sets overcome this by adding orbitals to each atom with angular momentum beyond what is required for the ground state (e.g., d functions are added to carbon, p functions can be added to hydrogen). In the $6-31 \mathrm{G}(\mathrm{d})$ basis set, one basis function comprised of six Gaussian primitives is applied to the core atomic orbitals, and the valence orbitals are split into two basis functions. The first one is composed of three Gaussian primitives and the second by one Gaussian primitive. In this specific basis set, d orbitals are added to heavy atoms.

### 2.4 Description of Calculations

2.4.1. General Methods. Calculations were conducted using the computational resources at the Center for Computational Sciences ${ }^{42,43}$ at Duquesne University and the

Pittsburgh Supercomputing Center. Electronic structure calculations were carried out with the Gaussian03 program ${ }^{41}$ with the second-order Møller-Plesset (MP2) method ${ }^{44}$ and the $6-31 \mathrm{G}(\mathrm{d})$ basis set. ${ }^{45,46}$ Molecular mechanics calculations were performed using the CHARMM program ${ }^{47}$ and the CHARMM force field for proteins ${ }^{8,48}$ and nucleic acids. ${ }^{6,7}$ The water model in all simulations was TIP3P. ${ }^{49}$ Crystal simulations employed the CRYSTAL ${ }^{50}$ facility in CHARMM. Molecular dynamics simulations were performed in the $\mathrm{NPT}^{51}$ ensemble using the Leap-Frog integrator and periodic boundary conditions with the $\mathrm{SHAKE}^{9}$ algorithm applied to covalent hydrogen bonds, which allowed for a 0.002 ps timestep. A dielectric constant of 1.0 was used in all simulations. Electrostatics were treated using the particle-mesh Ewald method. ${ }^{16}$ In the case of net-charged systems, a neutralizing background is used by default in the standard Ewald calculation. ${ }^{21}$ Ewald calculations were performed using the specified real space cutoffs with the LennardJones interactions truncated at the same distance. Kappa (screening parameter) was set to 0.36 ; the order parameter was set to $6 .{ }^{6}$ Nonbond pair lists were maintained to $14 \AA$, nonbond interactions were truncated at $12 \AA$, and the image cutoff was kept at $14.0 \AA$. Nonbond lists, hydrogen bond lists, and image lists were updated heuristically. The volumes of the systems were monitored to ensure that there were only minimal fluctuations in the volume over time. Snapshots of the systems were saved every ten picoseconds. Simulation details for each specific system are given in the following sections
2.4.2. Quantum Mechanical Calculations (polyamide model compounds). Each compound was energy minimized using MP2/6-31G(d). For comparison, the structures were also energy minimized using HF and B3LYP with the 6-31G(d) basis set.

Frequency calculations were also performed at the MP2/6-31G(d) level of theory. In order to obtain the potential energy surfaces (PES) for selected torsion angles, torsion energies were calculated by holding the selected torsion angle fixed at different increments ( $30^{\circ}$ increments from $-180^{\circ}$ to $180^{\circ}$ ) and geometry optimizing all other degrees of freedom at the three levels of theory mentioned above.
2.4.3. Netropsin Crystal Simulation. The starting structure for the simulation was an x-ray structure solved by Berman, et al. ${ }^{52}$ The crystal was of the $\mathrm{C} 2 / \mathrm{c}$ space group (monoclinic crystal; unit cell dimensions are $\mathrm{a}=21.7 \AA, \mathrm{~b}=6.37 \AA, \mathrm{c}=42.7 \AA ; \alpha=\gamma=$ $\left.90.0^{\circ}, \beta=107.0^{\circ}\right)$. The netropsin molecule was built using the CHARMM program, ${ }^{47}$ and the HBUILD algorithm ${ }^{53}$ was used to add hydrogens to netropsin. The structure also contained five crystallographic waters and one sulfate ion, which were retained for the simulation. No additional water or ions were added to the system. The system was then energy minimized using the steepest descent algorithm (SD) for 500 steps and for 100 steps using the ABNR algorithm to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$. Dynamics were then performed for 5 ns .
2.4.4. DNA with Bound Polyamides. The initial structure for the simulations was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{CVY} .{ }^{54}$ The crystal contained a DNA decamer (5'-CCAGATCTGG-3') with two polyamides bound antiparallel in the minor groove (polyamide sequence was ImPyPyPy- $\beta-\mathrm{Dp}$ ). The polyamides recognize the six base pair target site 5'-AGATCT-3'. In addition to the two DNA strands and two polyamides, the structure also contained 52 crystallographic waters. The structure was solved to $2.15 \AA$ resolution, and the crystal is of the $\mathrm{P} 2_{1} 2_{1} 2_{1}$ space group (orthorhombic crystal; unit cell dimensions are $\mathrm{a}=36.4 \AA, \mathrm{~b}=38.7 \AA, \mathrm{c}=47.2 \AA ; \alpha=\beta=\gamma=90.0^{\circ}$ ). The DNA-
polyamide complex was generated in CHARMM using the revised nucleic acids parameters. ${ }^{6,7}$ The HBUILD algorithm ${ }^{53}$ was used to add hydrogens to the DNA and polyamide molecules, and the hydrogens were then energy minimized for 50 steps using the steepest descent (SD) algorithm. The 52 crystallographic waters were also added to the system. The simulation was performed on the asymmetric unit of the crystal; this contained all the atoms of the DNA duplex and the two polyamides. Water was added to fill the vacuum spaces in the crystal by generating the primary and image atoms of the DNA, polyamide, and crystallographic waters and overlaying them with a waterbox that was larger than the dimensions of the asymmetric unit (waterbox dimensions were $59.9 \AA$ x $59.9 \AA$ x $63.8 \AA$ ). The waterbox was centered over the primary atoms, and any water molecules within $2.5 \AA$ of the DNA or polyamides were deleted. The waterbox was then overlayed three times with the original atoms and rotated $90^{\circ}$ each time in order to obtain the correct density for the crystal (the first time the waterbox was rotated $90^{\circ}$ degrees about the $x$-axis, the second time the waterbox was rotated $90^{\circ}$ about the $y$-axis, and the third time the waterbox was rotated $90^{\circ}$ about the z-axis). Any overlapping waters with the original atoms were deleted. An image update and image centering was then performed in order to delete the appropriate amount of waters for the size and symmetry of the unit cell. The waters that moved after the image update and centering were deleted. No ions were identified or reported in the published x-ray crystal structure, and none were added. The water, DNA, and polyamide were energy minimized for 500 steps of steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$. The water was then equilibrated for 20 ps using the NVT ensemble with the DNA and polyamide constrained (RMSD of waters converged to $\sim 5.6 \AA$ ). Following relaxation of water, the constraints
on the DNA and polyamide were released, and all atoms were allowed to relax for 20 ps in the NVT ensemble (RMSD of whole complex and waters converged to $\sim 5.8 \AA$; average RMSD of DNA/polyamide complex was $0.65 \AA$ ). All atoms were then subjected to dynamics for 10 ns at 300 K . Many studies involving DNA crystal simulations are performed at $300 \mathrm{~K} .{ }^{7,19,24,55,56}$ Two more simulations were also run at 300 K ; however, the systems were perturbed by rotating the waterbox $-90^{\circ}$ about the $y$-axis (rotated $90^{\circ}$ about the x - and z -axes) and $-90^{\circ}$ about the z -axis (rotated $90^{\circ}$ about the x - and y -axes) during the reoverlaying procedure (perturbing the systems in this way gave different initial coordinates for the solvent atoms and thus allowed for better sampling of conformational space $\left.{ }^{57-61}\right)$. The systems were prepared in the same way as described above, and data was then collected over 10 ns in the NPT ensemble. Because the x-ray crystal data was collected at $-160^{\circ} \mathrm{C}(113 \mathrm{~K})$, four simulations were run at 113 K . The systems were set up and run in the same way as the simulations run at 300 K (rotations of the waterbox, equilibration, and data collection were done in the same way as the 300 K crystal simulations).

In order to improve sampling of conformational space, four more simulations were run at 113 K and three more simulations were run at 300 K . Principal component analysis (PCA) was performed, and points (structures) that were close to the x-ray value on the PCA plots were used as the starting structures for the additional simulations. In order to perturb the systems, each simulation was started from different initial velocities.

A simulation was also performed with a hydroxypyrrole-containing polyamide. The initial structure for the simulations was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{CVX} .{ }^{54}$ The crystal contained the DNA decamer (5'-CCAGATCTGG-3') with two polyamides
bound antiparallel in the minor groove (polyamide sequence was $\operatorname{ImPyHpPy}-\beta-\mathrm{Dp}$ ). As with the other simulated Dervan polyamides, the polyamides recognize the six base pair target site 5'-AGATCT-3'. In addition to the two DNA strands and two polyamides, the structure also contained 30 crystallographic waters, which were retained for the simulation. The structure was solved to $2.27 \AA$ resolution, and the crystal is of the $\mathrm{P} 3_{1}$ space group (trigonal crystal; unit cell dimensions are $\mathrm{a}=31.2 \AA, \mathrm{~b}=31.2 \AA, \mathrm{c}=46.1 \AA$; $\left.\alpha=\beta=90.0^{\circ}, \gamma=120.0^{\circ}\right)$. The system was set up and generated similarly to the other DNA/polyamide complex simulations, and the simulation was run for 10 ns at 103 K .
2.4.5. DNA/Netropsin Complex. The initial structure for the simulation was the xray crystal structure $\mathrm{PDB}^{3}$ id 121D. ${ }^{62}$ The crystal contained a DNA dodecamer (5'-CGCAAATTTGCG-3') with the drug netropsin in the minor groove. Netropsin has a binding preference for AT sequences over GC sequences, and in this structure, netropsin covers five of the six AT base pairs of the DNA fragment. In addition to the two DNA strands and netropsin, the structure also contained 47 crystallographic waters, which were retained for the simulation. The structure was solved to $2.2 \AA$ resolution, and the crystal is of the $\mathrm{P} 2{ }_{1} 2_{1} 2_{1}$ space group (orthorhombic crystal; unit cell dimensions are $\mathrm{a}=25.65 \AA$, $\left.\mathrm{b}=42.03 \AA, \mathrm{c}=65.33 \AA ; \alpha=\beta=\gamma=90.0^{\circ}\right)$. The system was set up and generated similarly to the DNA/polyamide complex simulations, and the simulation was run for 10 ns at 300 K .
2.4.6. DNA/Distamycin Complex. The initial structure for the simulation was the x-ray crystal structure $\mathrm{PDB}^{3}$ id 121D. ${ }^{63}$ The crystal contained a DNA dodecamer (5'-CGCAAATTTGCG-3') with the drug netropsin in the minor groove. Like netropsin, distamycin has a binding preference for AT sequences over GC sequences, and in this
structure, distamycin covers five of the six AT base pairs of the DNA fragment. In addition to the two DNA strands and distamycin, the structure also contained 75 crystallographic waters, which were retained for the simulation. The structure was solved to $2.2 \AA$ resolution, and the crystal is of the $\mathrm{P} 2_{1} 2_{1} 2_{1}$ space group (orthorhombic crystal; unit cell dimensions are $\mathrm{a}=25.20 \AA, \mathrm{~b}=41.07 \AA, \mathrm{c}=64.65 \AA ; \alpha=\beta=\gamma=90.0^{\circ}$ ). The system was set up and generated similarly to the DNA/polyamide and DNA/netropsin complex simulations, and the simulation was run for 10 ns at 288 K .
2.4.7. Netropsin and DNA/Complex Simulations Using AMBER parameters. Two tests were performed using previously described parameters intended for use with the AMBER ${ }^{64}$ force field. ${ }^{65,66}$ The first test was a simulation of a netropsin crystal; the starting structure was the same used for the netropsin crystal simulation with the new CHARMM parameters, which is that by Berman, et al. ${ }^{52}$ The simulation was prepared and run for 5 ns in the same way as the netropsin simulation using the CHARMM parameters. The second test was a crystal simulation of a DNA/netropsin complex. The starting structure for the simulation was the same as that used for the new CHARMM parameters $\left(\mathrm{PDB}^{3}\right.$ id $\left.121 \mathrm{D}^{62}\right)$. The system was set up and run identically to the CHARMM DNA/netropsin simulation. The system contained the same amount of water and was started from the same initial velocities as the CHARMM simulation and run for 10 ns at 300 K .

### 2.4.8. Quantum Mechanical Calculations (retinoid model compounds).

 Optimization and frequency calculations were carried out on the model compounds at the MP2/6-31G(d) level of theory. In order to obtain the potential energy surfaces (PES) for selected torsion angles, torsion energies were calculated by holding the selected torsionangle fixed at different increments ( $20^{\circ}$ increments from $-180^{\circ}$ to $180^{\circ}$ ) and geometry optimizing all other degrees of freedom at the three levels of theory mentioned above.
2.4.9. Small Retinoid Crystal Simulations. For the retinoid crystal simulations, the unit cell was built using the CHARMM program. ${ }^{47}$ No water or ions were reported in the x-ray structures, and none were added. The systems were energy minimized using the steepest descent algorithm (SD) for 500 steps and for 1000 steps using the ABNR algorithm to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$. All simulations were carried out at the temperature in which the crystal structure was determined, unless noted otherwise.
2.4.9.1. All-Trans Retinal Crystal. The initial structure for the simulation was the x-ray crystal structure $\mathrm{CSD}^{67}$ id TRETAL01 ${ }^{68}$ (space group $\mathrm{P} 2{ }_{1}$ (monoclinic), unit cell dimensions were $\mathrm{a}=15.27 \AA, \mathrm{~b}=8.26 \AA, \mathrm{c}=14.94 \AA ; \alpha=\gamma=90.0^{\circ}, \beta=104.7^{\circ}$; four retinal molecules were reported for the unit cell, which was what was used for the simulation). Five molecular dynamics simulations started from different velocities were then performed for 1 ns at 273 K .
2.4.9.2. All-Trans Retinoic Acid Crystal (Monoclinic). The initial structure for the simulation was the x-ray crystal structure $\mathrm{CSD}^{67}$ id VITAAC01 ${ }^{69}$ (space group P21 (monoclinic), unit cell dimensions were $a=8.078 \AA, b=34.103 \AA, c=7.387 \AA \AA ; \alpha=\beta=$ $90.0^{\circ}, \gamma=118.76^{\circ}$ ). The retinoic acid molecule parameterized for the force field is the deprotonated form; however, the crystal structure of the ligand contained the protonated form. A patch analogous to the existing patches for neutral aspartic acid or neutral glutamic acid was made in order to protonate the retinoic acid. Five molecular dynamics simulations started from different velocities were then performed for 1 ns at 393 K .
2.4.9.3. All-Trans Retinoic Acid Crystal (Triclinic). The initial structure for the simulation was the x-ray crystal structure $\mathrm{CSD}^{67}$ id TRETAL10 ${ }^{70}$ (space group P1' (triclinic), unit cell dimensions were $\mathrm{a}=8.04 \AA, \mathrm{~b}=28.49 \AA, \mathrm{c}=5.996 \AA ; \alpha=50.58^{\circ}, \beta$ $=71.38^{\circ}, \gamma=95.7^{\circ}$; two molecules of retinoic acid were found in the unit cell, which was what was used for the simulation). Five molecular dynamics simulations started from different velocities were then performed for 1 ns at 123 K .

### 2.4.10. Protein/Retinoid Complex Crystal Simulations.

### 2.4.10.1. Retinoic Acid Complexed with cellular retinoic acid binding protein II

 (CRABP II). The initial structure for the simulation was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{CBS}^{71}$ (resolution $1.80 \AA$; space group $\mathrm{P} 2_{1} 2_{1} 2_{1}$ (orthorhombic), unit cell dimensions were $\mathrm{a}=45.65 \AA, \mathrm{~b}=47.56 \AA, \mathrm{c}=77.61 \AA ; \alpha=\beta=\gamma=90.0^{\circ}$ ). The retinoid-protein complex was generated in CHARMM using the revised protein parameters, ${ }^{48}$ the HBUILD algorithm ${ }^{53}$ in CHARMM was used to add hydrogens to the protein and ligand molecules, and the hydrogens were then energy minimized for 50 steps using the steepest descent (SD) algorithm (energy tolerance was $0.001 \mathrm{kcal} / \mathrm{mol}$ ). The 100 crystallographic waters present in the x-ray structure were added, and three sodium ions were added to neutralize the system (no ions were reported in the crystal structure). The sodium ions were placed randomly, and they were all placed at least $10 \AA$ away from the protein to avoid bad contacts. The crystal was then built using the CRYSTAL ${ }^{50}$ facility in CHARMM. Water was added to fill the vacuum spaces in the crystal by generating the primary and image atoms of the protein, ligand, and crystallographic waters and overlaying them with a waterbox that was larger than the dimensions of the asymmetric unit (waterbox dimensions were $59.9 \AA \times 59.9 \AA \times 63.8 \AA$ ). The waterbox was centeredover the primary atoms, and any water molecules within $2.5 \AA$ of the protein or ligand were deleted. The waterbox was then overlayed three times with the original atoms and rotated $90^{\circ}$ each time in order to obtain the correct density for the crystal (the first time the waterbox was rotated $90^{\circ}$ degrees about the x -axis, the second time the waterbox was rotated $90^{\circ}$ about the $y$-axis, and the third time the waterbox was rotated $90^{\circ}$ about the $z$ axis). Any overlapping waters with the original atoms were deleted (overlapping atoms were any atoms within $2.8 \AA$ of the original atoms). An image update and image centering was then performed in order to delete the appropriate amount of waters for the size and symmetry of the unit cell. The waters that moved after the image update and centering were deleted. The water, ions, protein, and ligand were energy minimized for 500 steps of steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$. The solvent was then equilibrated for 20 ps using the NVT ensemble with the protein and retinoic acid constrained. Following relaxation of water, the constraints on the protein/ligand complex were released, and all atoms were allowed to relax for 20 ps in the NVT ensemble (RMSD of whole complex and solvent converged to $\sim 14 \AA$; average RMSD of protein/retinoic acid complex after equilibration was $0.89 \AA$ ). All atoms were then subjected to dynamics for 2 ns at 277 K .

### 2.4.10.2. Retinol Complexed with Cellular Retinol Binding Protein (CRBP) from

 Zebrafish. The initial structure for the simulation was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{KQW}^{72}$ (resolution $1.38 \AA$; space group I4 (tetragonal), unit cell dimensions were $\mathrm{a}=$ $89.02 \AA, b=89.02 \AA, c=38.27 \AA ; \alpha=\beta=\gamma=90.0^{\circ}$ ). The retinoid-protein complex was generated in CHARMM in the same way as the retinoic acid/CRABP II system. The crystal structure was determined at pH 5.60 , so the His residues were protonated for thesimulation. The 194 crystallographic waters present in the x-ray structure were added, and three chlorine ions were added to neutralize the system (no ions were reported in the crystal structure). The chloride ions were placed randomly and at least $10 \AA$ away from the protein to avoid bad contacts. The crystal was then built, and water was added to fill the vacuum spaces in the crystal in the same way as for the retinoic acid/CRABP II system. Equilibration and dynamics were performed on the retinol/CRBP complex in the same way as for the retinoic acid/CRABP II complex. All atoms were then subjected to dynamics for 2 ns at 100 K .
2.4.10.3. Fenretinide Complexed with retinol binding protein ( $R B P$ ). The initial structure for the simulation was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{FEL}^{73}$ (resolution $1.80 \AA$; space group $\mathrm{P} 2_{1} 2_{1} 2_{1}$ (orthorhombic), unit cell dimensions were $\mathrm{a}=46.63 \AA, \mathrm{~b}=$ $49.26 \AA, \mathrm{c}=76.64 \AA ; \alpha=\beta=\gamma=90.0^{\circ}$ ). The last 8 residues $(176-183)$ are unresolved for this x-ray structure, so the missing residues were generated with homology modeling using the MOE homology program ${ }^{74,75}$ based on a segment matching procedure. ${ }^{76}$ The residues generated with the homology modeling were then energy minimized with steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$, while holding the rest of the protein/ligand complex fixed. The retinoid-protein complex was then generated in CHARMM in the same way as the retinoic acid/CRABP II and retinol/CRBP systems. Bovine RBP contains three disulfide bridges. The crystal structure was determined at pH 5.0-5.3, so the His residues were protonated for the simulation. The 163 crystallographic waters present in the x-ray structure were added, and one sodium ion was added to neutralize the system, which was placed randomly (no ions were reported in the crystal structure). The crystal was then built, and water was added to fill the vacuum
spaces in the crystal in the same way as for the retinoic acid/CRABP II and the retinol/CRBP systems. Equilibration and dynamics were performed on the fenretinide/RBP complex in the same way as for the other systems. All atoms were then subjected to dynamics for 3 ns at 275 K .
2.4.10.4. Retinal Complexed with Sensory Rhodopsin II (SRII). The initial structure for the simulation was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{H} 68^{77}$ (resolution $2.10 \AA$; space group $\mathrm{C} 222_{1}$ (orthorhombic), unit cell dimensions were $\mathrm{a}=84.83 \AA, \mathrm{~b}=$ $128.74 \AA, c=50.74 \AA ; \alpha=\beta=\gamma=90.0^{\circ}$ ). The asymmetric unit of the $x$-ray structure contains one chain, and this is how SRII is thought to exist biologically; however, the oligomeric state of SRII is unknown. The first residue and the last 20 residues (220-239) are unresolved for this x-ray structure, so the missing residues were generated with homology modeling using the MOE homology program. ${ }^{74,75}$ The residues generated with the homology modeling were then energy minimized with steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$ while holding the rest of the protein/ligand complex fixed. The retinoid-protein complex was then generated in CHARMM in the same way as the other protein/ligand systems. The crystal structure was determined at pH 4.6 so the His residues were protonated for the simulation. The 27 crystallographic waters and one chloride ion present in the x-ray structure were added, and three sodium ions were added to neutralize the system. The sodium ions were randomly placed at least $10 \AA$ away from the protein to avoid bad contacts. The crystal was then built, and water was added to fill the vacuum spaces in the crystal in the same way as for the other protein/ligand systems. Equilibration and dynamics were performed on the retinal/SRII complex similar to the other systems. The solvent was equilibrated for 30 ps in the NVT
ensemble, and the whole system was then equilibrated for 30 ps in the NVT ensemble (average RMSD of retinal/SRII complex after equilibration was $1.67 \AA$ ). All atoms were then subjected to dynamics for 2 ns at 100 K .
2.4.10.5. Retinal Complexed with Bacteriorhodopsin $(B R)$. The initial structure for the simulation was the x-ray crystal structure $\mathrm{PDB}^{3}$ id $1 \mathrm{PXR}^{78}$ (resolution $1.70 \AA$; space group $\mathrm{P} 12_{1} 1$ (monoclinic), unit cell dimensions were $\mathrm{a}=44.97 \AA, \mathrm{~b}=108.84 \AA$, c $\left.=55.98 \AA ; \alpha=\gamma=90.0^{\circ}, \beta=113.60^{\circ}\right)$. The asymmetric unit of this x-ray structure contains two chains; however, native BR exists in the purple membrane of Halobacterium salinarum as a trimeric structure. ${ }^{79}$ Also, this structure is a P50A mutant. The first five residues and the last 18 residues (232-249) are unresolved in both chains for this $x$-ray structure, so the missing residues were generated with homology modeling using the MOE homology program. ${ }^{74,75}$ The residues generated with the homology modeling were then energy minimized with steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$, while holding the rest of the protein/ligand complex fixed. The retinoidprotein complex was then generated in CHARMM in the same way as the other protein/ligand systems. The crystal structure was determined at pH 3.5 so the His residues were protonated for the simulation. Only solvent-exposed Glu and Asp residues were protonated. The 124 crystallographic waters present in the x-ray structure were added, and ten chloride ions were added to neutralize the system. The ions were randomly placed $\sim 10 \AA$ away from the protein to avoid any bad contacts. The crystal was then built, and water was added to fill the vacuum spaces in the crystal in the same way as for the other protein/ligand systems. Equilibration and dynamics were performed on the retinal/BR complex similarly to the other systems. The solvent was equilibrated
for 30 ps in the NVT ensemble, and the whole system was then equilibrated for 30 ps in the NVT ensemble. All atoms were then subjected to dynamics for 3.5 ns at 110 K .

### 2.4.11. Crambin Simulations (Convergence Analysis). Molecular dynamics

 simulations were performed using the CHARMM program ${ }^{47}$ and the CHARMM force field. ${ }^{6}$ Calculations were conducted using the computational resources at the Pittsburgh Supercomputing Center. The water model in all simulations was TIP3P. ${ }^{49}$ Crystal simulations employed the CRYSTAL ${ }^{50}$ facility in CHARMM. Molecular dynamics simulations were performed in the $\mathrm{NPT}^{51}$ ensemble using the Leap-Frog integrator and periodic boundary conditions with the $\mathrm{SHAKE}^{9}$ algorithm applied to covalent hydrogen bonds, which allowed for a 0.002 ps timestep. A dielectric constant of 1.0 was used in all simulations. Electrostatics were treated using the particle-mesh Ewald method. ${ }^{16}$ Ewald calculations were performed using the specified real space cutoffs with the LennardJones interactions truncated at the same distance. Kappa (screening parameter) was set to 0.36 ; the order parameter was set to 6 . Nonbond pair lists were maintained to $14 \AA$, nonbond interactions were truncated at $12 \AA$, and the image cutoff was kept at $14.0 \AA$. Nonbond lists, hydrogen bond lists, and image lists were updated heuristically. The volumes of the systems were monitored to ensure that there were only minimal fluctuations in the volume over time. Snapshots of the systems were saved every ten picoseconds.2.4.11.1. Long-time protein simulations. The initial protein (crambin) structure for all simulations was $\mathrm{PDB}^{3}$ id $1 \mathrm{EJG} .{ }^{80}$ The crystal contained only the 46 -residue protein crambin; no water or ions were reported in the structure. The structure was solved to 0.54 $\AA$ resolution, and the crystal was of the $\mathrm{P} 2_{1}$ space group (monoclinic crystal; unit cell
dimensions were $\left.\mathrm{a}=40.82 \AA, \mathrm{~b}=18.50 \AA, \mathrm{c}=22.37 \AA ; \alpha=\gamma=90.0^{\circ}, \beta=90.47^{\circ}\right)$. The protein was generated in CHARMM using the revised protein backbone parameters. ${ }^{48}$ The HBUILD algorithm ${ }^{53}$ was used to add hydrogens to the protein molecule, and the hydrogens were then energy minimized for 50 steps using the steepest descent (SD) algorithm. The simulation was performed on the asymmetric unit of the crystal, which contained only the protein. Water was added to fill the vacuum spaces in the crystal in the same way as was done for the DNA. No ions were added to the system, and the final system had a charge of 0 . The protein and water were energy minimized for 1000 steps of steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$. The water was then equilibrated for 20 ps using the NVT ensemble with the protein constrained. Following relaxation of water, the constraints on the protein were released, and all atoms were allowed to relax for 20 ps in the NVT ensemble. All atoms were then subjected to dynamics for 45 ns at 100 K , which was the temperature at which the x-ray crystal structure was determined. Nine more simulations were also run at 100 K ; however, the systems were perturbed by assigning different initial velocities to the atoms in each simulation. The systems were prepared in the same way as described above, and data was then collected over 45 ns in the NPT ensemble. Snapshots were saved every 10 ps . The total simulation time over all ten simulations was 450 ns .
2.4.11.2. Short-time protein simulations. The systems were set up in the same way as for the long-time simulations (same starting coordinates and protocol). Each simulation was started from different initial velocities in order to introduce a perturbation. Snapshots were collected every 10 ps . Ten 2 ns simulation were run. The first 2 ns of the long-time simulations were used along with the 10 short simulations for
data analysis of short-time simulations to give a total of 20 short simulations (total simulation time was 40 ns ).
2.4.12. DNA Simulations (Convergence Analysis). Molecular dynamics simulations were performed using the CHARMM program ${ }^{47}$ and the CHARMM force field. ${ }^{6}$ Calculations were conducted using the computational resources at the Pittsburgh Supercomputing Center. The water model in all simulations was TIP3P. ${ }^{49}$ Crystal simulations employed the CRYSTAL ${ }^{50}$ facility in CHARMM. Molecular dynamics simulations were performed in the $\mathrm{NPT}^{51}$ ensemble using the Leap-Frog integrator and periodic boundary conditions with the $\mathrm{SHAKE}^{9}$ algorithm applied to covalent hydrogen bonds, which allowed for a 0.002 ps timestep. A dielectric constant of 1.0 was used in all simulations. Electrostatics were treated using the particle-mesh Ewald method. ${ }^{16}$ In the case of charged systems, electrostatics were treated using the modified particle mesh Ewald method for net-charged systems. ${ }^{21}$ This modified Ewald summation is implemented in versions 25 and higher of the CHARMM program. Ewald calculations were performed using the specified real space cutoffs with the Lennard-Jones interactions truncated at the same distance. Kappa (screening parameter) was set to 0.36 ; the order parameter was set to 6 . Nonbond pair lists were maintained to $14 \AA$, nonbond interactions were truncated at $12 \AA$, and the image cutoff was kept at $14.0 \AA$. Nonbond lists, hydrogen bond lists, and image lists were updated heuristically. The volumes of the systems were monitored to ensure that there were only minimal fluctuations in the volume over time. Snapshots of the systems were saved every 10 ps .
2.4.12.1. Long-time DNA Simulations. The initial DNA structure for all simulations was the x-ray crystal structure $\mathrm{PDB}^{3}$ id 119D. ${ }^{81}$ The crystal contained a DNA
dodecamer ( $5^{\prime}$-CGTAGATCTACG-3') along with $1 \mathrm{Mg}^{+2}$ ion, which was retained for the simulations, and 137 crystallographic waters, which were also retained. The structure was solved to $2.25 \AA$ resolution, and the crystal was of the C2 space group (monoclinic crystal; unit cell dimensions were $\mathrm{a}=64.83 \AA, \mathrm{~b}=35.36 \AA, \mathrm{c}=25.35 \AA ; \alpha=$ $\gamma=90.0^{\circ}, \beta=92.24^{\circ}$ ). The DNA dodecamer was generated in CHARMM using the revised nucleic acids parameters. ${ }^{6,7}$ The HBUILD algorithm ${ }^{53}$ was used to add hydrogens to the DNA molecule, and the hydrogens were then energy minimized for 50 steps using the steepest descent (SD) algorithm. The 137 crystallographic waters were then added to the system. The simulation was performed on the asymmetric unit of the crystal; this contained all the atoms of the DNA duplex and the one $\mathrm{Mg}^{2+}$ ion. Water was added to fill the vacuum spaces in the crystal by generating the primary and image atoms of the DNA, $\mathrm{Mg}^{2+}$ ion, and crystallographic waters and then overlaying them with a waterbox that was larger than the dimensions of the asymmetric unit (waterbox dimensions were $59.9 \AA \times 59.9 \AA \times 63.8 \AA$ ). The waterbox was centered over the primary atoms, and any water molecules within $2.5 \AA$ of the DNA or magnesium ion were deleted. The waterbox was then overlayed three times with the original atoms and rotated $90^{\circ}$ each time in order to obtain the correct density for the crystal (the first time the waterbox was rotated $90^{\circ}$ degrees about the $x$-axis, the second time the waterbox was rotated $90^{\circ}$ about the $y$-axis, and the third time the waterbox was rotated $90^{\circ}$ about the $z$-axis). Any overlapping waters with the original atoms were deleted. An image update and image centering was then performed in order to delete the appropriate amount of waters for the size and symmetry of the unit cell. The waters that moved after the image update and centering were deleted. Only the $\mathrm{Mg}^{2+}$ ion that was identified and reported in the published x-ray crystal
structure was used in the system; no additional ions were added. The water, DNA, and magnesium ion were energy minimized for 500 steps of steepest descent to an energy convergence of $0.001 \mathrm{kcal} / \mathrm{mol}$. The water and $\mathrm{Mg}^{2+}$ ion were then equilibrated for 20 ps using the NVT ensemble with the DNA constrained (RMSD of waters converged to $\sim 5.6$ $\AA$ ). Following relaxation of water, the constraints on the DNA were released, and all atoms were allowed to relax for 20 ps in the NVT ensemble (RMSD of DNA and waters converged to $\sim 5.8 \AA$; average RMSD of DNA only was $0.65 \AA$ ). All atoms were then subjected to dynamics for 150 ns at 288 K , which was the temperature at which the x-ray crystal structure was determined. Many studies involving DNA crystal simulations are performed at $300 \mathrm{~K} .{ }^{7,19,24,55,56}$ Three more simulations were also run at 288 K ; however, the systems were perturbed by rotating the waterbox $-90^{\circ}$ about the $y$-axis (rotated $90^{\circ}$ about the x - and z -axes) and $-90^{\circ}$ about the z -axis (rotated $90^{\circ}$ about the x - and y -axes) during the reoverlaying procedure (perturbing the systems in this way gave different initial coordinates for the solvent atoms and thus allowed for better sampling of conformational space ${ }^{57,58,61}$ ). The systems were prepared in the same way as described above, and data was then collected over 150 ns in the NPT ensemble. The total simulation time over all four simulations was 600 ns .
2.4.12.2. Short-time DNA Simulations. The short DNA simulations were set up similarly to the long-time simulations (same starting DNA structure, same protocol used). Each system was starting from different initial solvent coordinates in order to introduce a perturbation in each simulation. This was achieved by rotating the waterbox in different degrees about the $\mathrm{x}-$, y -, and z -axes. A total of 20 simulations were run for 2 ns each
(total of 40 ns ). Snapshots were saved every 10 ps . Listed below are the rotations of the waterbox about each axis for the 20 simulations.

Table 2.1 Rotations about $\mathrm{x}, \mathrm{y}$, and z axes for short DNA simulations

| Simulation | $\mathbf{x}$ | $\mathbf{y}$ | $\mathbf{z}$ |
| :--- | :--- | :--- | :--- |
| 1 | 20 | 20 | -20 |
| 2 | 25 | 25 | -25 |
| 3 | 30 | 30 | -30 |
| 4 | 35 | 35 | -35 |
| 5 | 40 | 40 | -40 |
| 6 | 45 | 45 | -45 |
| 7 | 50 | 50 | -50 |
| 8 | 55 | 55 | -55 |
| 9 | 60 | 60 | -60 |
| 10 | 65 | 65 | -65 |
| 11 | 70 | 70 | -70 |
| 12 | 75 | 75 | -75 |
| 13 | 80 | 80 | -80 |
| 14 | 85 | 85 | -85 |
| 15 | 100 | 100 | -100 |
| 16 | 105 | 105 | -105 |
| 17 | 110 | 110 | -110 |
| 18 | 115 | 115 | -115 |
| 19 | 120 | 120 | -120 |
| 20 | 125 | 125 | -125 |

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## Chapter 3

## Force Field Parameters for DNA Minor Groove-Binding Pyrrole-Imidazole Polyamides and Significance of Multiple Simulations

This chapter describes the parameterization procedure for the DNA minor groovebinding polyamides, the results of the crystal simulations performed to test the parameters, and conclusions made from the study.

### 3.1 Parameterization of DNA minor groove-binding polyamides

Parameterization of the polyamides began by first dividing the polyamides into six model compounds and carrying out energy minimization using MP2/6-31G(d), which includes the effects of electron correlation and dispersive interactions. The bond, angle, torsional, and improper terms were modified to match the quantum mechanical results. The parameters were evaluated by performing molecular dynamics crystal simulations of the reported x-ray crystal structures of a netropsin crystal, 2:1 Dervan polyamide/DNA complexes, a netropsin/DNA complex, and a distamycin/DNA complex. The new parameters can be used for improved simulations of polyamides and their nucleic acid complexes.

Parameterization was performed by first dividing the molecule into components, otherwise known as model compounds. ${ }^{1-15}$ The models, shown in Figure 3.1, consist of
three commonly used heterocycles for DNA recognition (N-methylpyrrole, N methylimidazole, and $N$-methylhydroxypyrrole), $\beta$-alanine used for enhancing flexibility of the molecule, a cationic tail (dimethylaminopropylamide) used for phosphate binding, and $\gamma$-aminobutyric acid as a linker in hairpin polyamides. Since heterocycles and linkers are joined together in the polyamide by amide bonds, compounds that contain terminal amide groups to represent the linkage to another heterocycle or linker in the polyamide sequence were chosen.


1



2



3




Figure 3.1 Model compounds for polyamides (some hydrogens not shown for clarity).

Each model compound was energy minimized using MP2/6-31G(d). Model compound 1 was also minimized using CCSD/6-31G(d) to verify the MP2 results. Comparison of MP2 and CCSD energy minimized structures revealed that the error for bonds was $0.003 \pm 0.005 \AA, 0.2 \pm 0.1^{\circ}$ for angles, and $0.4 \pm 0.3^{\circ}$ for dihedral angles. The comparison indicated that the MP2/6-31G(d) level of theory is appropriate for force field parameterization, a result in agreement with others for parameterization of polymers. ${ }^{16,17}$ For comparison, the structures were also energy minimized using HF and B3LYP with the $6-31 \mathrm{G}(\mathrm{d})$ basis set. Initial topology and parameters were then created from existing
compounds in the CHARMM force field. ${ }^{18}$ Indole was the starting point for the pyrrolebased rings ( $\mathbf{1}$ and 3 ), histidine was the starting point for the imidazole ring (2), and $\mathrm{N}-$ methylacetamide was the starting point for the tail and linker compounds $(\mathbf{4}, \mathbf{5}, \mathbf{6})$.

The initial topology and parameters were then modified to match the quantum mechanical results. The minimized bond lengths and angle values obtained from the MP2 optimized structures were used as the equilibrium bond lengths and equilibrium angle values for the new CHARMM parameters for the associated bond and angle types. Each model compound was energy minimized with the CHARMM force field with steepest descent and Newton-Raphson (tolerance gradient was $10^{-6} \mathrm{kcal} / \mathrm{mol} / \AA$ for each minimization). The parameters where then adjusted to minimize the error with the MP2/6-31G(d) structure. The average error between CHARMM and MP2 is $0.02 \pm 0.02$ $\AA$ for bond distances and $2 \pm 2^{\circ}$ for angles.

In addition to bonds and angles, torsion terms were also parameterized. For each model compound, the potential energy surfaces (PES) for selected torsion angles (see Figures 3.2-3.7 for selected torsion angles; the bonds with the arrows were the bonds which were rotated about) were calculated by holding the selected torsion angle fixed at different increments ( $30^{\circ}$ increments from $-180^{\circ}$ to $180^{\circ}$ ) and geometry optimizing all other degrees of freedom. The torsion angles selected were for rotations about the bonds connecting the amide groups to the heterocycle, the bond connecting the methyl group to the heterocycle, and the carbon-carbon single bonds of the linker and tail model compounds. The torsion angles selected are necessary for the polyamide to optimize contacts with the bases of the minor groove of DNA. The CHARMM dihedral angle parameters force constant ( $K \chi$ ), multiplicity ( $n$ ), and phase shift $(\delta)$ were then modified to
match the quantum mechanical results from the potential energy surfaces. Each model compound was energy minimized in CHARMM (steepest descent and Newton-Raphson; tolerance gradient was $10^{-6} \mathrm{kcal} / \mathrm{mol} / \AA$ ) with the selected torsion angle held fixed, and the potential energy surface of the dihedral was calculated. If the surface did not match the MP2 results, the multiplicity, phase shift, or force constant were adjusted as necessary. After the results from the CHARMM compounds matched those obtained from the QM calculations, the model compounds were pieced together to make the entire polyamide molecule for use in molecular mechanics energy minimization or molecular dynamics simulations.

For N-methylpyrrole, (1) three dihedrals were selected for parameterization, which are the rotation of the methyl group and the rotation of the two amide groups on the ring. The most flexible points on the polyamide are most likely along the backbone where the heterocycles are connected to the amide groups. The polyamide will have to rotate around those bonds to optimize interactions with the DNA, and the methyl groups on the ring have free rotation, which will interact with the DNA through van der Waals contacts, so they will also have to rotate to optimize their interactions with the DNA. The heterocycles themselves are aromatic, so the rings stay planar. The initial CHARMM parameters for three important dihedrals borrowed from similar functionality did not agree with the MP2/6-31G(d) potential energy surface (Figure 3.2). For example, the methyl rotation $(\beta)$ is $180^{\circ}$ out of phase, both amide groups have energy barriers that are too high ( $\sim 3.4 \mathrm{kcal} / \mathrm{mol}$ too high for $\alpha$ and $\sim 4.1 \mathrm{kcal} / \mathrm{mol}$ too high for $\gamma$ ), and the second amide rotation $(\gamma)$ does not give the lowest energy at $0^{\circ}$. The energies obtained from HF/6-31G(d) and B3LYP/6-31G(d) do not necessarily match those obtained with MP2/6-
$31 \mathrm{G}(\mathrm{d})$. For the methyl rotation, B3LYP gives a barrier at $0^{\circ}$, whereas MP2 gives a minimum. HF shows the correct phase for methyl rotation with similar barriers to MP2 . As shown for the $\alpha$ torsional plot, HF barriers are $34 \%$ lower than MP2.

energy vs. $\alpha$

energy vs. $\beta$

energy vs. $\gamma$


Figure 3.2 Structure (top) and the three torsion energy plots for N -methylpyrrole. Blue diamond $(\bullet)$, initial CHARMM; red square $(+$ ), MP2/6-31G(d); green triangle ( $\boldsymbol{\Delta}$ ), final CHARMM; pink x (X), B3LYP/6-31G(d); purple star (*), HF/6-31G(d).

Similarly, for N -methylimidazole (2), three dihedrals were selected for
parameterization. The chosen dihedrals are again for the rotation of the amide groups
and the methyl group on the imidazole ring. The results for the imidazole torsions are shown in Figure 3.3. The initial CHARMM parameters for N-methylimidazole borrowed from similar functionality are also in poor agreement with the MP2/6-31G(d) results. The methyl rotation is out of phase by $60^{\circ}$, and the energy barriers are $\sim 1.5 \mathrm{kcal} / \mathrm{mol}$ too high. Both amide group rotations do not show the correct behavior. The first amide group rotation $(\alpha)$ should have a maximum at $180^{\circ}$ and $-180^{\circ}$; instead the maxima are at $90^{\circ}$ and $-90^{\circ}$. The second amide group rotation ( $\gamma$ ) has energy barriers that are $\sim 6.3$ $\mathrm{kcal} / \mathrm{mol}$ too high, and there are minima at $180^{\circ}$ and $-180^{\circ}$, which the MP2 does not show. The HF and B3LYP results agree well with MP2 for both the amide group rotations; however, for the methyl rotation, the B3LYP is $30^{\circ}$ out of phase.

energy vs. $\alpha$

energy vs. $\beta$



Figure 3.3 Structure (top) and the three torsion energy plots for N -methylimidazole. Blue diamond ( $\bullet$ ), initial CHARMM; red square ( + ), MP2/6-31G(d); green triangle ( $\boldsymbol{\Delta}$ ), final CHARMM; pink x (X), B3LYP/6-31G(d); purple star (*), HF/6-31G(d).

Four torsions were chosen for hydroxypyrrole (3), the two amide group rotations, the methyl group rotation, and the OH group rotation. The OH group rotation was chosen for this model compound because the H of the hydroxyl group hydrogen bonds with the O 2 of thymine, so this group will have to rotate to optimize interaction with the DNA. The results for hydroxypyrrole are shown in Figure 3.4. The initial CHARMM parameters are in poor agreement with the MP2/6-31G(d) results; all of the selected torsion angles are out of phase, the energy barriers are too high for the methyl rotation $(\sim 1.3 \mathrm{kcal} / \mathrm{mol})$ and the OH group rotation $(\sim 17.0 \mathrm{kcal} / \mathrm{mol})$, and all but the methyl rotation show the incorrect multiplicity. The final CHARMM parameters give correct methyl rotation barriers and show the right multiplicity, the rotation of the OH group on the pyrrole ring shows the correct maximum, the rotation of the first amide group $(\alpha)$ shows the correct minimum and multiplicity, and the rotation of the second amide group shows the correct barrier heights. The height of the small maximum at $0^{\circ}$ is $\sim 2.5$ $\mathrm{kcal} / \mathrm{mol}$ higher than that of the MP2 results; however, this is the best agreement that could be obtained in order to obtain simultaneous agreement for the other dihedrals of the hydroxypyrrole model compound. The second amide torsion plot $(\gamma)$ shows a minimum at $180^{\circ}$ and a small maximum at $0^{\circ}$, whereas the other heterocycles show the global minimum at $0^{\circ}$ for this torsion surface. The MP2 calculations are carried out in the gas phase, so the amide group rotates $180^{\circ}$ for the carbonyl oxygen to form a hydrogen bond with the H of the OH group, lowering the energy. The other heterocycles do not have the OH group; therefore, the minimum is at $0^{\circ}$ for the others. Both HF and B3LYP are in reasonable agreement with MP2 for $\delta$; they both give the correct phase, multiplicity, and a maximum at $0^{\circ}$. HF is $0.65 \mathrm{kcal} / \mathrm{mol}$ lower than MP2, and B3LYP is $0.52 \mathrm{kcal} / \mathrm{mol}$
lower than MP2. HF shows good agreement with MP2 for $\beta$ and $\gamma$. For $\beta$, HF shows the correct phase and multiplicity and is $0.28 \mathrm{kcal} / \mathrm{mol}$ higher than MP2. For $\gamma, \mathrm{HF}$ again shows the correct phase and multiplicity and is $0.39 \mathrm{kcal} / \mathrm{mol}$ higher than MP2. However, HF shows poor agreement with MP2 for $\alpha$; HF gives a maximum at $150^{\circ}$, and a minimum at $180^{\circ}$, while MP2 gives a maximum at $180^{\circ}$. B3LYP shows reasonable agreement with MP2 for $\alpha$ and $\beta$. For $\alpha$, B3LYP gives the correct phase and multiplicity, and the maximum at $180^{\circ}$ is $0.92 \mathrm{kcal} / \mathrm{mol}$ lower than MP2. For $\beta$, B3LYP shows the correct phase and multiplicity, and the barriers are $\sim 0.32 \mathrm{kcal} / \mathrm{mol}$ lower than MP2. For $\gamma$, B3LYP shows the correct phase and multiplicity; however, the barriers at $90^{\circ}$ and $-90^{\circ}$ are $3.41 \mathrm{kcal} / \mathrm{mol}$ higher than MP2.
energy vs. $\alpha$



energy vs. $\gamma$

energy vs. $\delta$


Figure 3.4 Structure (top left) and the four torsion energy plots for hydroxypyrrole. Blue diamond $(\bullet)$, initial CHARMM; red square $(+)$, MP2/6-31G(d); green triangle ( $\boldsymbol{\Delta}$ ), final CHARMM; pink x (X), B3LYP/6-31G(d); purple star (*), HF/6-31G(d).

Similar analyses were carried out for the cationic tail (4), the $\gamma$-aminobutyric acid linker, and the $\beta$-alanine linker. For the cationic tail, two dihedral angles were selected for rotation, which are for the rotation of the amide group and the rotation of the dimethyl amino group. The torsion plots are shown in Figure 3.5.

energy vs. $\alpha$

energy vs. $\beta$


Figure 3.5 Structure (top) and the two torsion energy plots for the dimethylaminopropyl cationic tail. Blue diamond ( $\downarrow$ ), initial CHARMM; red square ( + ), MP2/6-31G(d); green triangle ( $\mathbf{\Delta}$ ), final CHARMM; pink x (X), B3LYP/6-31G(d); purple star (*), HF/6$31 \mathrm{G}(\mathrm{d})$.

From the torsion energy plots for the cationic tail, it can be seen that the initial CHARMM parameters for this model compounds are in better agreement with the MP2
results than the initial heterocycle parameters. The MP2 calculations were performed with two dihedrals constrained at $180^{\circ}$, which are marked by the straight red arrows. The molecule was initially optimized with no constraints; however, in the gas phase, the molecule folded up on itself into a ball, so the constraints were necessary. The initial CHARMM parameters for the tail show the correct phase and number of minima and maxima; however, the energy barriers for both torsion angles are initially too low (1.34 $\mathrm{kcal} / \mathrm{mol}$ for $\alpha$, and $1.78 \mathrm{kcal} / \mathrm{mol}$ for $\beta$ ). Initially, constraints were placed on the same dihedrals in CHARMM when making the new parameters; however, the energy barriers were too high, and the constraints had to be removed. The height of the energy barriers in the final CHARMM parameters is in good agreement with the calculated MP2 energy barriers $(0.6 \mathrm{kcal} / \mathrm{mol}$ lower than MP2 for $\alpha$ and $0.43 \mathrm{kcal} / \mathrm{mol}$ lower for $\beta)$. The HF and B3LYP results are in reasonable agreement with the MP2 results for $\alpha$. HF is 0.5 $\mathrm{kcal} / \mathrm{mol}$ lower than MP2 and B3LYP is $0.74 \mathrm{kcal} / \mathrm{mol}$ lower than MP2. For $\beta$, the agreement is better with HF being $0.05 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 and B3LYP being 0.52 $\mathrm{kcal} / \mathrm{mol}$ lower than MP2.

Two dihedrals were selected for rotation for the $\gamma$-aminobutyric acid linker (5) and for the $\beta$-alanine linker (6). For the $\gamma$-aminobutyric acid linker, the two dihedrals selected are for rotation of the amide groups. The results are shown in Figure 3.6. The initial CHARMM parameters for the $\gamma$-linker are in slight agreement with MP2, so they did not require significant modification. As with the cationic tail model compound, the MP2 calculations were performed with two dihedrals constrained at $180^{\circ}$ (the constrained dihedrals are marked by the straight red arrows) for the same reason they were place on the cationic tail. Again, constraints were placed on the same dihedrals in CHARMM
initially, but the energy barriers were too high and could not be lowered until the constraints were removed. The plots show the correct phase and multiplicity. However, the one energy barrier is $\sim 1.15 \mathrm{kcal} / \mathrm{mol}$ too high $(\alpha)$, and the other $\sim 0.97 \mathrm{kcal} / \mathrm{mol}$ too low $(\beta)$. The final CHARMM gives energy barriers that are good agreement with MP2; $6.24 \mathrm{kcal} / \mathrm{mol}$ for $\alpha$, which is $0.19 \mathrm{kcal} / \mathrm{mol}$ higher than MP2, and $0.22 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 for $\beta$. The energies for $120^{\circ}$ and $150^{\circ}$ for $\beta$ do not follow the MP2 energies, but this is the best agreement that could be obtained in order to have both torsions agree well with MP2. The HF and B3LYP results are in reasonable agreement with the MP2 results for both torsions. For $\alpha$, HF is $0.73 \mathrm{kcal} / \mathrm{mol}$ higher than MP2, and B3LYP is 0.34 $\mathrm{kcal} / \mathrm{mol}$ higher than MP2. For $\beta$, the agreement is better; HF is $0.34 \mathrm{kcal} / \mathrm{mol}$ higher than MP2, and B3LYP is $0.09 \mathrm{kcal} / \mathrm{mol}$ higher than MP2.

energy vs. $\alpha$



Figure 3.6 Structure (top) and the two torsion energy plots for the $\gamma$-aminobutyric acid linker. Blue diamond ( $\bullet$ ), initial CHARMM; red square ( + ), MP2/6-31G(d); green triangle ( $\boldsymbol{\Delta}$ ), final CHARMM; pink $x(X)$, B3LYP/6-31G(d); purple star (*), HF/6$31 G(d)$.

As with the $\gamma$-aminobutyric acid linker, the two dihedrals selected for the $\beta$ alanine linker are for rotation of the amide groups. The results are shown in Figure 3.7. As for the tail and $\gamma$-linker model compounds, the torsion energy plots for the $\beta$-alanine linker show that the initial CHARMM parameters were in better agreement with MP2 than were the initial heterocycle parameters. One dihedral was constrained at $180^{\circ}$ (marked by the straight red arrow) for the MP2 calculations for the same reasons as the cationic tail and the $\gamma$-linker. When the structure was energy minimized with CHARMM,
no constraints were used because the energy barriers were too high. The energy barrier is $\sim 5.9 \mathrm{kcal} / \mathrm{mol}$ too high for $\alpha$ and the barrier is $\sim 0.94 \mathrm{kcal} / \mathrm{mol}$ too low for $\beta$; however, they are in phase and show the correct multiplicity. With adjustment of force constants, the final CHARMM parameters agree well with the MP2 results. The energy barrier for $\beta$ is $\sim 0.10 \mathrm{kcal} / \mathrm{mol}$ higher than MP2, and $\alpha$ is $2.85 \mathrm{kcal} / \mathrm{mol}$ higher than MP2; however, this better than the initial CHARMM parameters and is the best agreement that could be obtained and still have both torsions agree well with the MP2 results. HF and B3LYP are in good agreement with MP2 for both dihedrals. HF is $0.15 \mathrm{kcal} / \mathrm{mol}$ lower than MP2 and B3LYP is $0.49 \mathrm{kcal} / \mathrm{mol}$ lower than MP2 for $\alpha$. For $\beta$, HF is $0.41 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 and B3LYP is $0.39 \mathrm{kcal} / \mathrm{mol}$ lower than MP2.

energy vs. $\alpha$

energy vs. $\beta$


Figure 3.7 Structure (top) and the two torsion energy plots for the -alanine linker. Blue diamond $(\star)$, initial CHARMM; red square $(+)$, MP2/6-31G(d); green triangle ( $\boldsymbol{\Delta}$ ), final CHARMM; pink x (X), B3LYP/6-31G(d); purple star (*), HF/6-31G(d).

The same trend of systematic errors in multiplicity, phase shift, and force constant is observed for all six model compounds. Because HF and B3LYP do not necessarily show the same behavior as MP2, electron correlation and dispersion effects are probably important for this class of compounds and are necessary in polyamide parameter development. HF and B3LYP are therefore not appropriate methods to use for parameterization of these compounds.

Final CHARMM parameters for the model compounds produce structures and energies in good agreement with MP2/6-31G(d) results. Overall, the error between CHARMM and MP2 geometry optimized structures is $0.02 \pm 0.02 \AA$ for bonds, $2 \pm 2^{\circ}$ for angles, and $5 \pm 9^{\circ}$ for dihedrals. The error in barrier heights between CHARMM and MP2 is $0.38 \pm 0.40 \mathrm{kcal} / \mathrm{mol}$.

### 3.2 DNA/polyamide Simulations in the Crystal Environment

Simulations were performed in the crystal environment in order to validate the new force field parameters for polyamides. There are several high-resolution crystal structures of DNA with bound polyamides available, ${ }^{1-3,19}$ and performing the simulations in the crystal environment allows for a direct comparison back to experiment. Simulations were analyzed by comparing the calculated DNA helical parameters and DNA/polyamide distances back to the starting x-ray structure using the t -test.
3.2.1 Netropsin Crystal Simulation. As an initial test of the parameters, a 5 ns crystal simulation of netropsin $\left(\mathrm{CCSD}^{20}\right.$ id NETRSN ${ }^{21}$; R-factor 6.7\%) was performed. The system contained one netropsin molecule, five water molecules, and one sulfate ion. The average RMSD of the simulation was $0.42 \pm 0.09 \AA$, as shown in Figure 3.8 (top right). The low RMSD indicates that the parameters are in good agreement with the experimental crystal structure.



RMSD vs. time


Figure 3.8 Structure of netropsin (top left) and RMSD vs. time plot (top right) for the netropsin simulation (hydrogens not included in RMSD). The bottom plot is the RMSD vs. time for different parts of the netropsin molecule (rings in blue, amide backbone in red, and tails in green).

Exploration of the differences between the simulated and experimentally derived structures was done by examining the RMSDs of the pyrrole rings, amide backbone, and the cationic tail. It was found that the tail region moves more than any other region during the simulation (RMSD of $0.4 \pm 0.1 \AA$ ). The backbone has the next largest contribution with an average RMSD of $0.27 \pm 0.07 \AA$, and the rings contribute the least with an average RMSD of $0.22 \pm 0.08 \AA$ (Figure 3.8, bottom). This is reasonable because the rings are aromatic and planar so there will not be as much movement within
the rings, the backbone has a little more freedom to move, but not as much as the tail.
Rotation is allowed where the amide groups are attached to the rings, and the tails consist primarily of C-C single bonds so free rotation is allowed. The high RMSD of the tail with the amidino group at the end involves a rotation around the carbon-carbon bonds from $\sim 40^{\circ}$ to $80^{\circ}$ (Figure 3.9).


dihedral vs. time (NZ-CV-CC-CZ2)


Figure 3.9 Time series (bottom) for one dihedral of the amidino group tail (indicated by arrow on netropsin structure on the top). This end of the molecule appears to contribute most to the RMSD due to rotation about carbon-carbon bonds. The overlay of the x-ray (red) and average structures (green) shows that the simulated structure retains the same orientation of the amidino group.

The x-ray value for the dihedral shown is $65^{\circ}$; the simulation yields $54.7 \pm 9.7^{\circ}$. The overlay of the x-ray and average structures shows that even though the dihedrals of the tail show deviation, on average, the tail retains the same orientation as the x-ray structure.
3.2.2. Nucleic Acid Helical Parameters. The next tests of the developed parameters were crystal simulations of DNA/polyamide complexes. An important part of nucleic acid simulation analysis is the examination of helical parameters. Helical parameters describe how the bases and base pairs are oriented with respect to each other and/or the global helical axis of the DNA. ${ }^{22-25}$ Because they describe the relative orientation of the bases and base pairs, and a ligand that binds to a nucleic acid will affect these orientations, a successful force field for ligands that bind to nucleic acids should be able to reproduce these parameters accurately. ${ }^{26}$ As part of the validation of the new parameters, the helical parameters of the simulated DNA were compared the values of the starting x-ray crystal structure using a t-test at the $95 \%$ confidence level.

A complete set of helical parameters (definitions, names, and spatial arrangements of bases and base pairs) was established at a 1988 conference, ${ }^{25}$ a standard reference frame has also been adopted, ${ }^{23}$ which is used by the Nucleic Acids Database (NDB) ${ }^{27}$ for calculating helical parameters. Various computer programs ${ }^{22,28-35}$ have been developed to calculate nucleic acid helical parameters, and although the algorithms used are different, this has been found to have a limited effect on the computed parameters. ${ }^{24,36}$ However, the choice of reference frame was found to have a profound effect on the computed parameters. ${ }^{24}$ Figure 3.10 shows the standard reference frame and helical parameters that are calculated by most helical parameter analysis programs.


Figure 3.10 Standard reference frame (top left) and helical parameters. In the reference frame, the x -axis points away from the minor groove. Translations are shown in the upper half of the figure (x-displacment through rise) and rotations are shown in the bottom half of the figure (buckle through tip). Each small tile represents a DNA base; two of them together represent a base pair.

The helical parameters can be divided into translational motions and rotational motions; they can also be divided further into three families: base pair-axis parameters, intrabase pair parameters, and interbase pair parameters.

There are two different ways of describing nucleic acid helical parameters, which are termed global and local. ${ }^{37,38}$ In the global approach, the helical parameters calculated describe the overall arrangement of the bases and base pairs with respect to a global helical axis, which is an axis that runs vertically down a double-stranded helix and is usually taken to be linear. The difficulty with a global approach is that a linear helical axis is not an appropriate assumption for many irregular nucleic acid conformations (e.g., a structure with a bend). In the local approach, the helical parameters calculated describe the orientations of the bases/base pairs within the local framework of two successive base pairs along the nucleic acid fragment; a local helical axis is calculated for each base pair step rather than a global helical axis for the entire fragment. There has been no final decision on calculating global versus local helical parameters; both have their advantages and disadvantages. ${ }^{37}$ Local parameter algorithms have the advantage that they avoid the difficulty of defining a global helical axis, and they also yield parameters that depend only the conformation of the given base pair step. Global parameters depend on the conformation of the whole fragment and involve the determination of the global helical axis. However, global parameters have the advantage of distinguishing between the different helical conformations (e.g., A-DNA vs. B-DNA). If there is curvature in the fragment, the location and extent can be identified more easily with global parameters. ${ }^{37}$ The various available computer algorithms utilize one of the two approaches, but as mentioned, neither method has been found to be truly superior to another. ${ }^{37}$ In this
dissertation, the 3DNA program, ${ }^{28}$ which uses the local approach, was used for calculation of helical parameters.
3.2.3. DNA/Dervan polyamide complex simulations. Adequate sampling of conformational space is a difficult task, especially for large biomolecules such as DNA in which the potential energy surface contains many local minima. ${ }^{39-41}$ In this study, we employ multiple trajectories because other studies have reported that this method improves sampling of conformational space. ${ }^{42-45}$ For this system, four initial trajectories were performed at the temperature at which the starting x-ray structure was determined (113 K), and three initial trajectories were performed at 300 K . The theoretical melting temperature for this DNA fragment is $305 \mathrm{~K}\left(2^{\circ} \mathrm{C}\right.$ for AT pairs and $4^{\circ} \mathrm{C}$ for GC pairs $)$, which will be increased a few degrees by the polyamides in the minor groove, so the melting temperature should be above the simulation temperature.

RMSD. Shown in Figure 3.11 is the RMSD, relative to the starting x-ray crystal structure $\left(\mathrm{PDB}^{46}\right.$ id $\left.1 \mathrm{CVY}^{3}\right)$; left side of Figure 3.11), of the DNA/polyamide complexes from the seven initial simulations done at the low temperature (113 K) and high temperature ( 300 K ).


Figure 3.11 Starting structure for DNA/Dervan polyamide simulations ( $\mathrm{PDB}^{46}$ id $1 \mathrm{CVY}^{3}$; left). The DNA is in red, and the polyamides are in green. The structure of the polyamides bound in the minor groove is shown on the bottom right (there are two of them bound antiparallel in the minor groove). The plot on the top right is the RMSD vs. time for the 113 K simulations (bottom) and for the 300 K simulations (top) compared to the starting x-ray crystal structure.

The RMS deviation converges quickly to $0.8-1 \AA$ for all of the low temperature simulations, with the exception of one low temperature simulation (LT2), which converges with the rest of the low temperature simulations at $\sim 5 \mathrm{~ns}$. As expected for the 300 K simulations, the RMSD shows more fluctuation than for the lower temperature simulations. The RMSD stays below $3 \AA$, but there is great amount of fluctuation up to $\sim 6 \mathrm{~ns}$ of simulation time. After 6 ns , the RMS deviations start to converge to around 1.7 $\AA$, with the exception of one simulation (HT3), which shows deviations greater than 2.5

Å. Table 3.1 lists the average RMSDs of the complex, DNA, and polyamides separately.

Table 3.1 RMSDs of complexes, DNA, and polyamides

|  |  |  | RMSD $(\AA \mathbf{\AA})$ |  |
| :--- | :--- | :--- | :--- | :--- |
| simulation | Temp (K) | complex | DNA only | polyamides only |
| HT1 | $300.0 \pm 0.83$ | $1.64 \pm 0.13$ | $1.15 \pm 0.15$ | $2.52 \pm 0.33$ |
| HT2 | $300.0 \pm 0.84$ | $1.95 \pm 0.27$ | $1.87 \pm 0.28$ | $1.68 \pm 0.24$ |
| HT3 | $300.0 \pm 0.83$ | $1.71 \pm 0.24$ | $1.65 \pm 0.24$ | $1.42 \pm 0.25$ |
|  |  |  |  |  |
| LT1 | $113.0 \pm 0.29$ | $0.96 \pm 0.04$ | $0.89 \pm 0.04$ | $1.13 \pm 0.06$ |
| LT2 | $113.0 \pm 0.29$ | $0.86 \pm 0.04$ | $0.73 \pm 0.04$ | $1.01 \pm 0.05$ |
| LT3 | $113.0 \pm 0.30$ | $0.79 \pm 0.05$ | $0.58 \pm 0.04$ | $0.95 \pm 0.15$ |
| LT4 | $113.0 \pm 0.31$ | $0.90 \pm 0.05$ | $0.83 \pm 0.05$ | $1.00 \pm 0.05$ |
|  |  |  |  |  |

For all of the low temperature simulations and one high temperature simulation, the RMSD for the complex is influenced specifically from the tails of the polyamide, which contain $\mathrm{C}-\mathrm{C}$ single bonds with free rotation. In the remaining two high temperature simulations, the end base pairs of the DNA exhibit some fraying, which is the most likely cause of higher RMSD contribution of the DNA.

Statistical Analysis of Helical Parameters. Helical parameters were computed using the 3DNA program..$^{28}$ The helical parameters were computed and then averaged from snapshots from the simulations. A t-test was then performed between the computed and x-ray helical parameters at the $95 \%$ confidence level (see Chapter 2, Section 2.2.2 for a description and formulas for the $t$-test). The output from 3DNA gives a value for each base or base pair in the DNA fragment (depending on what parameter is being measured); therefore, there is an average and standard deviation associated with each parameter for each snapshot. Figure 3.12 shows a sample output from 3DNA.

|  | bp | Shear | Stretch | Stagger | Buckle | Propeller | Opening |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | C-G | 0.71 | -0.26 | -0.27 | 9.75 | -18.27 | -3.72 |
| 2 | C-G | 0.67 | -0.12 | -0.01 | 6.20 | -14.32 | 1.72 |
| 3 | A-T | -0.13 | 0.08 | -0. 50 | -15.35 | -25.94 | 26.90 |
| 4 | G-C | -0.21 | -0.08 | -0.19 | -21.65 | -8.44 | 2.00 |
| 5 | A-T | 0.03 | -0.23 | 0.30 | -5.25 | -4.01 | -3.27 |
| 6 | T-A | 0.11 | -0.18 | -0.08 | -5.73 | -10.23 | 0.31 |
| 7 | C-G | 0.10 | -0.12 | 0.05 | -15.78 | -3.56 | -0.07 |
| 8 | T-A | 0.17 | 0.53 | -0.27 | 14.11 | -18.42 | 24.29 |
| 9 | G-C | -0.41 | -0.32 | -0.35 | -6.05 | -5.31 | -2.27 |
| 10 | G-C | -0.64 | 0.16 | -0.67 | -10.19 | -7.20 | 5.08 |
|  | ave. | 0.04 0.42 | -0.05 0.25 | -0.20 0.28 | -4.99 11.73 | -11.57 7.43 | 5.10 11.14 |

Figure 3.12 Sample output from 3DNA for a DNA 10-mer for six helical parameters.

An average of each helical parameter was taken over all snapshots used for a simulation.
In order to take the standard deviation from each snapshot in account, a pooled standard deviation was performed for each helical parameter. For each 10 ns simulation, 100 snapshots, or samples, (one every 100 ps ) were evaluated using 3DNA, and an overall average and a pooled standard deviation (Eqn.3.1) was calculated for each helical parameter.
$s_{\text {pooled }}=\sqrt{\frac{s_{1}{ }^{2}\left(n_{1}-1\right)+s_{2}{ }^{2}\left(n_{2}-1\right)+\ldots+s_{k}^{2}\left(n_{k}-1\right)}{n_{1}+n_{2}+\ldots+n_{k}-k}}$
Eqn. 3.1
$n=10$ for each snapshot ( $n=9$ in the case of base pair step parameters), $s$ is the standard deviation for each snapshot taken from a simulation, and $k=100$ ( 100 snapshots used).

After the pooled standard deviation was calculated, the $t$-test comparing the simulation averages to the x-ray values was performed for each helical parameter.

$$
\begin{equation*}
t_{\text {calc }}=\frac{\overline{x_{1}}-\overline{x_{2}}}{\sqrt{\left(s_{1}{ }^{2} / n_{1}\right)+\left(s_{2}{ }^{2} / n_{2}\right)}} \tag{Eqn. 3.2}
\end{equation*}
$$

$x_{1}$ is the simulation average, $x_{2}$ is the x-ray average, $s_{1}$ is the simulation standard deviation (this number is the number obtained from the pooled standard deviation), $s_{2}$ is the x-ray standard deviation (obtained from the 3DNA output for the x-ray structure), $n_{l}$ is the number of samples for the simulation, and $n_{2}$ is the number of samples for the x-ray structure. For the simulations, $n_{1}=1000$ ( 100 snapshots * 10 values) for individual base pair parameters, or $n_{1}=900$ ( 100 snapshots $* 9$ values) for base pair step parameters (the DNA/Dervan polyamide systems contained 10 base pairs). For the DNA/netropsin complex, the DNA/distamycin complex, and DNA with no bound polyamides, $n_{1}=1200$ for individual base pairs, and $n_{1}=1100$ for base pair steps (these systems contained 12 base pairs). The number of samples $\left(n_{2}\right)$ for the x-ray structure was set to 10 (1 structure * 10 values); 9 was used for the helical parameters that involve base pair steps rather than individual base pairs (in the case of the DNA/netropsin complex, DNA/distamycin complex, and DNA with no bound polyamides, 12 and 11 were used).

T-test results for Helical Parameters. According to the $t$-test, there is good agreement between the helical parameters calculated from both the low and high temperature simulations and the x-ray structure, where 17 out of the 18 helical parameters measured are not statistically significant according to the t-test. The only helical parameter that was significantly different in all seven individual simulations is stagger at the $95 \%$ confidence level. Tables of the $t$-test results for the individual simulations can be found in Appendix A.

Because stagger was the only helical parameter that was found to be statistically significant in all seven individual simulations, t-tests were performed on the combined sets of simulations (one t-test performed on the low temperature simulations and one t-
test performed on the high temperature simulations) to see if the combined t-test result would be non-significant. The t-tests were performed using Eqn. 3.3,

$$
\begin{equation*}
t_{\text {calc }}=\frac{\overline{x_{1}}-\overline{x_{2}}}{s_{\text {pooled }}} \sqrt{\frac{n_{1} n_{2}}{n_{1}+n_{2}}} \tag{Eqn. 3.3}
\end{equation*}
$$

where $x_{1}$ is the simulation average, $x_{2}$ is the x-ray value, $s_{\text {pooled }}$ is the pooled standard deviation of the simulations and x-ray standard deviation, $n_{l}$ is the number of samples used from the simulations, and $n_{2}$ is the number of samples used for the x -ray structure. In order to see if the combined $t$-test results and how the $t$-values fluctuated over time, the t-test was performed in the following way. Rather than taking the average value of stagger over the ten bases of the DNA fragment for each snapshot, all ten values from the 100 snapshots were used for the $t$-test for each simulation. Because this $t$-test was performed over multiple simulations, the t -test for the high temperature simulations used 300 snapshots, and the t-test for the low temperature simulations used 400 snapshots. For comparison to the x-ray structure, the x-ray values of stagger for each base were compared to the corresponding bases from the simulations rather using the average over the ten bases. The snapshots from the simulations were ordered chronologically, the stagger value for each base at each snapshot was averaged over all simulations considered, and a t-test was then performed between each snapshot and the x-ray value. Because the stagger values were averaged over all simulations considered, there were 100 snapshots and thus a total of 100 t-tests performed. The resulting time series plots showed how the $t$-values for stagger evolve with time.


Figure 3.13 Time series plots of t -values over time for the stagger helical parameter. The time series for the low temperature simulations is on the top; the time series for the high temperature simulations is on the bottom. The red lines denote the $t$-table value (1.96) at the $95 \%$ confidence level.

Stagger is also observed to be statistically significant when all low temperature simulations are examined together and when all high temperature simulations are examined together. In the time series plot for the low temperature simulations, all 100 t tests performed yield a $t$-value greater than the $t$-table value (1.96) at the $95 \%$ confidence level. The average of the $t$-values is 3.16 ; stagger is statistically significant when all low temperature simulations are considered together. For the high temperature simulations, the average of the $t$-values is 2.51 , so stagger is still statistically significant when all high temperature simulations are considered together. For comparison, the t-test results of a helical parameter (buckle) that did not exhibit any statistical significance in the individual
simulations is shown in Figure 3.14. Over the four low temperature simulations, all tvalues are less than 1.96, and the average $t$-value is 0.93 . For the high temperature simulations, all t -values are again less then 1.96 , and the average t -value is 0.64 .


Figure 3.14 Time series plots of t-values over time for the buckle helical parameter. The time series for the low temperature simulations is on the top; the time series for the high temperature simulations is on the bottom. None of the individual or combined simulations exhibited significance with respect to buckle.

Principal Components Analysis (PCA). Neither the high nor the low temperature simulations appear to be sampling the x-ray value for the stagger parameter properly. In order to examine where the simulations are sampling with respect to the x-ray value of stagger and which simulations sample the x-ray value best (or not at all), PCA was performed on the stagger parameter for both the high and low temperature simulations (Figure 3.15).


Scree plot (stagger_113K)


Scree plot (stagger_300K)


Figure 3.15 Principal component analysis plots of stagger (low temperature simulations, top; high temperature simulations, middle). The different colors represent the different simulations. The red diamond on each PCA plot represents the x-ray value. The Scree plots associated with the PCA plots are shown at the bottom (low T, left; high T, right).

Upon examination of the PCA plots, the x-ray point falls outside of the regions where the simulations are sampling. Three of the low temperature simulations (green, light blue,
and most of the pink) are not even in the region of where the x-ray value falls; the dark blue simulation is sampling right on the edge of where the x-ray value falls. The high temperature simulations seem to better sample the region of conformational space where the x-ray value falls for stagger; however, the x-ray value lies on the edge of the sampling area. The PCA plots confirm that the simulations simply are not sampling enough conformational space to sample the x -ray value.

For comparison, PCA was also carried out for the buckle parameter, which did not show statistical significant using the $t$-test. For the low temperature simulations (Figure 3.16, top) the x-ray point falls in the middle of the region around where each of the simulations sample; the simulations average out to sample the x-ray value. For the high temperature simulations (Figure 3.16, middle), the x-ray point falls in the middle of the region where the simulations sample. The PCA plots show that x-ray value of buckle is sampled efficiently by both sets of simulations, unlike stagger.

## buckle (113 K)




Scree plot (buckle_113K)


Scree plot (buckle_300K)


Figure 3.16 Principal component analysis plots of buckle (low temperature simulations, top; high temperature simulations, middle). The different colors represent the different simulations. The red diamond on the PCA plots represents the x-ray value. The Scree plots associated with the PCA plots are shown at the bottom.

Additional Simulations. To help improve sampling of conformational space of the stagger parameter, additional simulations were performed (four new simulations at the lower temperature $(113 \mathrm{~K})$ and three new simulations at the higher temperature $(300 \mathrm{~K})$ ). The simulations were started from structures that were in regions not well sampled by the simulations on the PCA plots for stagger (each started from different initial velocities). Shown in Figure 3.17 is the RMSD vs. time for the DNA/polyamide complexes for the additional simulations.

## RMSD vs. tme (DNA/polyamide complex)



Figure 3.17 RMSD vs. time plots for the 113 K simulations (bottom) and for the 300 K simulations (top) compared to the starting structures for the simulations.

The RMS deviation converges quickly to $\sim 0.5 \AA$ for all of the low temperature simulations at $\sim 4 \mathrm{~ns}$. As expected for the 300 K simulations, the RMSD shows more fluctuation than for the lower temperature simulations. There is a great amount of fluctuation during the simulations, with HT3 having the most fluctuation. However, the RMSD does not rise above $\sim 2.5 \AA$ for any of the simulations, and HT3 has the lowest RMSD at $\sim 1-1.5 \AA$. Table 3.2 lists the average RMSDs of the complex, DNA, and polyamides separately.

Table 3.2 RMSDs of complexes, DNA, and polyamides

|  |  |  | RMSD (£̊) |  |
| :--- | :--- | :--- | :--- | :--- |
| simulation | Temp (K) | complex | DNA only | polyamides only |
| HT1 | $300.0 \pm 0.82$ | $1.80 \pm 0.24$ | $1.76 \pm 0.22$ | $1.56 \pm 0.57$ |
| HT2 | $300.0 \pm 0.85$ | $1.88 \pm 0.18$ | $1.81 \pm 0.18$ | $1.80 \pm 0.25$ |
| HT3 | $300.0 \pm 0.83$ | $1.26 \pm 0.23$ | $1.29 \pm 0.25$ | $0.86 \pm 0.16$ |
|  |  |  |  |  |
| LT1 | $113.0 \pm 0.30$ | $0.36 \pm 0.04$ | $0.42 \pm 0.05$ | $0.28 \pm 0.04$ |
| LT2 | $113.0 \pm 0.28$ | $0.38 \pm 0.04$ | $0.38 \pm 0.04$ | $0.30 \pm 0.03$ |
| LT3 | $113.0 \pm 0.32$ | $0.40 \pm 0.06$ | $0.42 \pm 0.07$ | $0.27 \pm 0.05$ |
| LT4 | $113.0 \pm 0.31$ | $0.31 \pm 0.03$ | $0.31 \pm 0.03$ | $0.25 \pm 0.03$ |

For two of the high temperature simulations and three of the low temperature simulations, the DNA contributes most to the RMSD of the complex. Some fraying of the end base pairs was observed, which is most likely the cause of the higher contribution of the DNA. In the remaining simulations, the polyamides contribute more to the RMSD which is due to the tails of the polyamides (free rotation is allowed about the $\mathrm{C}-\mathrm{C}$ single bonds of the tails).

Low-Temperature Simulations. A total of eight simulations were performed at 113 K (four original and four new). All individual simulations with the exception of one showed that stagger is statistically significantly different from the x-ray structure (all other helical parameters are non-significant; see Appendix A for tables of helical parameter and t-values). When all eight simulations were combined together, the stagger parameter is still statistically significant according to the t -test (Figure 3.18, top). All t values are greater than 1.96 ; the average $t$-value is 3.11 (this is a slight improvement over the four initial simulations, which had an average $t$-value of 3.16). A PCA plot of the stagger parameter shows that the four new simulations do improve sampling; more simulation points are in the region where the x-ray value (red diamond) lies. However,
some simulations are still not sampling even near where the x -ray value lies (e.g., the green, light blue, and pink simulations). Even when these three simulations were not considered in the analysis, the $t$-test still yielded statistical significance (average $t$-value was 3.04; plot not shown). In the case of the low temperature simulations, even the additional sampling is not sufficient because according to the $t$-test, stagger is still statistically significant from the x-ray structure at the $95 \%$ confidence level.

High-Temperature Simulations. A total of six simulations were performed at 300 K (three original and three new). In the original three individual simulations, stagger is the only statistically significant helical parameter; in the three new simulations, none of the helical parameters are statistically significant. When the six simulations were considered together, the t-test yields statistical significance at the $95 \%$ confidence level, but not at the $97 \%$ confidence level $\left(\mathrm{t}\right.$-table $\left.{ }_{97 \%}=2.24\right)$ (Figure 3.19 ; top). The average t value is 2.03 , which is just above 1.96 ; however, this value is less than 2.24 . The additional high temperature simulations did improve the average $t$-value (the three original simulations yielded an average $t$-value of 2.51 ), and the improvement observed is greater than that for the low temperature simulations. A PCA plot of stagger for the six simulations shows that the additional simulations improve the sampling of stagger (Figure 3.19, middle).

Because the six simulations together produced an average $t$-value that was close to the t -table value at the $95 \%$ confidence interval, only the three new high temperature simulations were then considered together for the t-test. With only the three new simulations, the t-test result for stagger is overall non-significance (Figure 3.20, top) at the $95 \%$ confidence level (average $t$-value is 1.64). Upon examination of a PCA plot of
stagger for only the three new simulations (Figure 3.20, middle), the simulations do sample the x-ray value more efficiently than the first three simulations and even when all six simulations are averaged together. For the three new simulations, all of the simulations appear to sample values that fall in the region of the x-ray value.



Figure 3.18 T-test for all eight low temperature simulations (top); PCA plot of all eight simulations (x-ray value is shown as a red diamond, middle); Scree plot corresponding to PCA plot (bottom).
stagger ( $\mathbf{3 0 0} \mathbf{K}$; 6 simulations)

stagger ( $\mathbf{3 0 0} \mathrm{K}$; 6 simulations)

Scree plot (stagger 300 K; 6 simulations)


Figure 3.19 T-test for all six high temperature simulations (top). The $t$-table value at the $95 \%$ CL (1.96) is marked by the red line, and the $t$-table value at the $97 \%$ CL (2.24) is marked by the blue line. PCA plot of all six simulations (x-ray value is shown as a red diamond, middle); Scree plot corresponding to PCA plot (bottom).


Figure 3.20 T-test for three new high temperature simulations (top). The t-table value at the $95 \%$ CL (1.96) is marked by the red line, and the t-table value at the $97 \%$ CL (2.24) is marked by the blue line. PCA plot of three new simulations (x-ray value is shown as a red diamond, middle); Scree plot corresponding to PCA plot (bottom).

The additional simulations improve the sampling in both the low and high temperature simulations. However, the low temperature average $t$-value is only improved upon
marginally ( 3.16 vs. 3.11 ), whereas in the high temperature simulations, a larger improvement is observed ( 2.51 vs. 2.03 ). From the average $t$-value and the PCA plot, the additional sampling with the low temperature simulations is not sufficient to obtain nonstatistical significance for the stagger parameter. The high temperature simulations do achieve non-significance. All six simulations together show statistical significance at the $95 \%$ confidence level but not at the $97 \%$ confidence level, which is already an improvement over the initial three simulations. When only the three new simulations are considered, stagger is not statistically significant at the $97 \%$ or $95 \%$ confidence level. As expected, the temperature at which the simulations are run influence the sampling of stagger for this particular system, and the higher temperature is the obvious choice for better sampling.

Although the high and low temperature simulations yielded a statistically significant t-value when considered by temperature, all 14 (8 low temperature and 6 high temperature) simulations gave a non-significant t -value for all helical parameters when considered together. The figure below shows the value of the $t$-test over time for all 14 simulations. The average $t$-value was 1.04 . None of the other helical parameters were statistically significant over all 14 high or low temperature simulations (see Appendix A for a table of t -values).


Figure 3.21 T-test for combined high and low temperature simulations. The t-table value at the $95 \%$ CL (1.96) is marked by the red line.

Although the $t$-test for stagger displays non-significant results when the 14 simulations are combined, the high temperature simulations are obviously contributing more to the non-significance. The additional high temperature simulations showed more improvement in the t-test results than the low temperature simulations. The low temperature simulations did not sample enough to get the x-ray value; however, the additional high temperature simulations do improve sampling. However, what is different about the high-temperature structures that lie close to the x-ray value and those that are far away from it on the PCA plot for stagger? Are there conformations/structural features of the DNA that cause stagger to deviate in the structures that are far away from the x-ray value? To answer these questions, structures that lay farthest from the x-ray value on the PCA stagger plot were examined by RMSD analysis, and the largest contribution to RMSD in structures was the DNA bases (not including backbone or sugars). An example of one of the simulated structures is shown below in Figure 3.22.


Figure 3.22 Comparison of x-ray structure and high temperature simulated structure. The x -ray structure is in red; polyamides are bound in the minor groove (also in red). One of the simulated structures that was far away from the x-ray point on the stagger PCA plot is shown with the DNA in green with the polyamides bound in the minor groove in colors. Gaps appear in the structures because only the bases of the DNA are shown (backbone of the DNA is not shown for clarity). The simulated structure has a wider minor groove and a more bent helix than the x-ray structure does.

The simulated structure is obviously distorted and has a wider minor groove (Figure 3.22, top left) and a more bent helix (Figure 3.22, bottom) than the x-ray structure does. The polyamides appear to have been "pushed" out the groove; the polyamides in the x-ray structure are deeper inside the groove than those in the simulated structures. The top
right picture shows the view from the major groove instead of the minor groove. This view is just to show that the polyamides really are not coming out of the groove (they look like like they are in the top left picture). The hydrogen bonding contacts are still where they should be; the polyamide tail parts that are flipped out of the groove appear to make hydrogen bonds with nearby waters. The distortion of the helix that is observed is why structures that look like this are farthest from the x-ray point on the stagger PCA plot.

The tilted/bent structures that are observed exhibit characteristics of A-DNA. 47-49 In A-DNA, the bases are more tilted with respect to the helical axis, which is seen in the structures. Also, the average inclination and x-displacement of the bases in those structures are closer to those of A-DNA, which are around $20^{\circ}$ and $-4 \AA$, respectively. A-DNA has a more shallow and a wider minor groove than B-DNA. The structures where the polyamide is getting "pushed out" of the minor groove are in line with A-form; the DNA adopts features of A-form, and the polyamide gets pushed out, because the groove has become narrower, and it also appears wider from the snapshots. However, the structures have not completely converged to the A-form; the rise ( $\sim 2.6 \AA$ in A-DNA and $3.4 \AA$ in B-DNA) and twist helical parameters ( $31^{\circ}$ in A-DNA and $36^{\circ}$ in B-DNA), the distance between the phosphates in each strand ( $\sim 5.0 \AA$ for A-DNA and $\sim 7.0 \AA$ for BDNA, and the width of the helix ( $26 \AA$ in A-DNA and $20 \AA$ in B-DNA) in these structures have B-form (or closer to B-form) values. An overstabilization of A-DNA was previously observed the CHARMM22 force field for nucleic acids; ${ }^{50}$ this has since been corrected with revised parameters, ${ }^{16,26}$ which were used in this study, so the force field parameters should not be an issue.

Structures from the low temperature simulations were also examined to see if they exhibited the same distortion of the DNA helix. The low temperature structures do not exhibit the widened minor groove or the distorted helix, and the polyamides stay where they are; they are not pushed out of the groove as in the high temperature structures.


Figure 3.23 Comparison of x-ray structure and low temperature simulated structure. The x-ray structure is in red; polyamides are bound in the minor groove (green). A representative example of a simulated low temperature structure is shown with the DNA in purple with the polyamides bound in the minor groove in colors. Gaps appear in the structures because only the bases of the DNA are shown (backbone of the DNA is not shown for clarity).

To further examine the high and low temperature structures, histograms were made based on the minor groove width of the structures. The histogram from the eight low temperature simulations (Figure 3.24) show that a few structures have a wider groove, but most have a narrower minor groove ( $\sim 13.0-13.6 \AA$ ), which is consistent with the representative structure shown in Figure 3.23 and the other structures examined from the low temperature simulations. The histogram from the six high temperature simulations (Figure 3.25) shows approximately a normal distribution of minor groove
widths, with most of them being wider ( $\sim 14.0-15.0 \AA$ ) than the low temperature simulations.


Figure 3.24 Probability density curve (top) and histogram bar plot (bottom) for minor groove width of low temperature simualtions.


Figure 3.25 Probability density curve (top) and histogram bar plot (bottom) for minor groove width of high temperature simualtions.

Because the low temperature simulations do not show the distortion of the helix, and most of the minor groove widths are narrower, the B-form is retained. However, in the
high temperature simulations, almost all structures have the wider minor groove that is characteristic of A-form DNA. Why are the distortion and the A-like features observed only in the high temperature structures? Also, because more distortion is observed in the high temperature simulations, why is more statistical significance not observed with the helical parameters (the high temperature simulations actually show better results for stagger than the low temperature)?

Force fields often do not get melting temperatures correct. ${ }^{51-56}$ This particular fragment of DNA has a theoretical melting temperature of 305 K ; the melting temperature of the DNA/polyamide complex is probably a few degrees higher. Polyamides have been found to increase the melting temperature ${ }^{57,58}$ of the DNA alone; the magnitude of the increase depends on the sequence of the polyamide, the DNA fragment, and how tightly the polyamide binds to the fragment. ${ }^{59}$ The simulation temperature for the high temperature simulations is 300 K . The simulation temperature is too close to that of the melting temperature, and because the force field does not account perfectly for the melting temperature, the DNA/polyamide complex is starting to melt in the high temperature simulations. The melting causes the distortion and thus the A-like features, which makes the minor groove wider and narrower, which in turn pushes the polyamides out. As far as the sampling is concerned, the higher temperature simulations do sample better and thus less statistical significance is observed for the stagger parameter. The low temperature simulations yield solid and undistorted structures, but they just do not sample well.

Polyamide/DNA interaction. Polyamides bind to the minor groove of DNA by hydrogen bonding, so these contacts were examined to ensure that they stayed in place
over time during the simulations and that the polyamide was not sliding around in the minor groove. For direct comparison back to the crystal structure, only heavy atom distances were measured. The distances that were measured are shown in the Figure 3.26 and denoted by the " $x$ " with a number above the distance lines. As with the helical parameters, a statistical comparison back to the starting x-ray crystal structure was also performed with the DNA/polyamide distances ( $95 \%$ confidence level) as a test of the parameters.


Figure 3.26 Polyamide distances measured (dashed lines) in the x-ray structure and from the low and high temperature crystal simulations. The "ladder" in the middle represents the DNA binding site of the polyamides. The circles in the middle represent the hydrogen bond donors or acceptors in the minor groove of the DNA.

The distances measured from the simulations are in good agreement with the x-ray structure (tables of measured distances and t -values can be found in Appendix A). Over the seven initial individual simulations, two distances were statistically different from the x-ray structure. The distance (x1a) is much larger than that in the crystal structure (x1a =
$4.99 \AA$ for the calculated average, and x1a $=3.3 \AA$ for the x-ray structure) in one low temperature simulation. Upon closer examination of the atoms of the polyamide and DNA that show this large distance, a water molecule was discovered above the minor groove. The x1a distance is from an NH group in the polyamide backbone to the O 2 of thymine. However, the NH group appears to be forming a strong hydrogen bond ( $\sim 1.8$ $\AA$ ) with the oxygen of the water molecule above the groove rather than forming a strong hydrogen bond with the O 2 of thymine. The other simulations do not have a water molecule in this position or anywhere near this position in order to form a hydrogen bond.

The distance from the H of the NH group to the water oxygen over the whole 10 ns of simulation 2 was plotted (shown in Figure 3.27, top). At the start of the simulation, the distance was $\sim 3.5 \AA$; also, waters were deleted within $2.5 \AA$ of the DNA/polyamide complex, but the water molecule moves right in to the polyamide, and the distance then stays small ( $\sim 1.9 \AA$ ). Upon closer examination of the starting structure, the water forming the hydrogen bond to the polyamide was actually in the minor groove between the two polyamides, but it was still at least $3 \AA$ from either of them and the DNA. It was trapped and moved closer to the NH of the polyamide backbone. The water and the DNA are therefore competing for hydrogen bonding to the polyamide.


Figure 3.27 Plot of distance from water to thymine O2 over the 10 ns time span of the simulation for low temperature simulation (top). The bottom plot shows the distance from polyamide to nearby water molecules. The water molecules exchange, but they still hydrogen bond to the polyamide. The straight red line on the bottom plot is at $3 \AA$, which is the average length of the hydrogen bond formed from the polyamide to the exchanging waters. There are water molecules hydrogen bonding to the polyamide at most times during the simulation.

One of the distances (x4) in one high temperature simulation is much larger ( $\sim 1$ $\AA$ ) than that of the crystal structure (calculated x4 = $3.95 \AA$; x-ray x4 = $2.95 \AA$ ). Upon closer examination of the simulation snapshots, a water molecule hydrogen was found to be hydrogen bonding to a carbonyl oxygen in the polyamide backbone. However, it was not the same water molecule hydrogen bonding to the polyamide for the duration of the simulation; the water molecules appear to be exchanging with each other, but nonetheless forming a strong hydrogen bond to a carbonyl oxygen in the polyamide backbone (Figure 3.27, bottom plot). The waters exchanging to form the hydrogen bond is probably a
consequence of the 300 K simulations having more kinetic energy than the 113 K simulations; the atoms move more and can thus exchange to form the hydrogen bond to the polyamide. Because good agreement between the calculated and x-ray distances was achieved with the initial seven simulations, distances were not measured for additional simulations (those were considered only for improvement of sampling).

Even though individual simulations exhibited statistical significance for DNA/polyamide distances due to intruding water molecules, when the distances were averaged over all simulations (low temperature and high temperature distances averaged separately), none of the distances showed statistical significance. The $t$-tests for the combined simulation distances were performed similarly to the combined t-tests for the helical parameters. For example, for distance x 1 for the low temperature simulations, 1000 values were taken from each of the four simulations, the values at each time were averaged, and a t-test was performed between the simulation values and x-ray value at each time ( 1000 t -tests total for each distance). A table of the average t -values obtained for each distance can be found in Appendix A.

### 3.2.4. DNA/Dervan hydroxypyrrole polyamide complex. A 10 ns crystal

 simulation of a different DNA/Dervan polyamide complex was also performed. In this simulation, one pyrrole heterocycle in each of the polyamides is replaced by the hydroxypyrrole heterocycle. The polyamides in the simulations discussed in Section 3.2.3 and in this section recognize the same bases in the minor groove of DNA (5'-AGATCT-3') and the DNA sequence is the same; however, the hydroxypyrrole moiety reduces the binding affinity due to elongation of the Watson-Crick hydrogen bonds of thetarget T-A base pairs. 3 The starting structure for the simulation was PDB46 id 1CVX. 3
The RMSD plot of the DNA/hydroxypyrrole polyamide is shown in Figure 3.28.


Figure 3.28 Starting structure for $\mathrm{DNA} /$ hydroxypyrrole polyamide simulation $\left(\mathrm{PDB}^{46} \mathrm{id}\right.$ $1 \mathrm{CVX}^{3}$; left). The DNA is in purple, and the polyamides are in silver in the minor groove. The structure of the polyamides bound in the minor grove is shown on the bottom right (two are bound antiparallel in the minor groove). The plot of RMSD vs. time of DNA/polyamide complex is on the top right.

The average RMSDs of the complex, DNA, and polyamides are $1.44 \pm 0.05 \AA, 1.44 \pm$ $0.05 \AA$, and $1.08 \pm 0.05 \AA$, respectively. The RMSD of the complex and the DNA stabilizes at $\sim 6 \mathrm{~ns}$ to $1.44 \AA$, while the RMSD of the polyamides is steady until about 7 ns. At 7 ns , the RMSD of the polyamides decreases to just above $1 \AA$. The DNA contributes more the RMSD than the polyamides, which is probably due to some fraying of the end bases. The contribution to the polyamide RMSD is most likely due to the polyamide tails, which contain C-C single bonds about which free rotation is allowed. This conformational freedom is reflected by the disorder in the starting x-ray structure,
because there were two possible conformations identified for the positions of the polyamide tails. ${ }^{3}$

The same analysis for the DNA helical parameters and DNA/polyamide distances was carried out for the hydroxypyrrole simulation. None of the 18 helical parameters examined were statistically significantly different from the x-ray structure values (a table of helical parameter and calculated t-values can be found in Appendix A). Figure 3.29 shows the DNA/polyamide distances measured (a table of measured distances is in Appendix A). Seven distances were measured for each polyamide to its respective DNA strand; only two distances out of the 14 measured were statistically significant according to the t -test. The distances showing statistical significance were the x 5 and x 1 a distances. The x-ray values for these distances are $3.16 \AA$ and $2.54 \AA$, respectively; whereas the average simulation distances are $4.93 \AA$ and $5.93 \AA$ (both are over an angstrom longer than the $x$-ray values).


Figure 3.29 Polyamide distances measured (dashed lines) in the x-ray structure and from the hydroxypyrrole polyamide crystal simulations. The "ladder" in the middle represents the DNA binding site of the polyamides. The circles in the middle represent the hydrogen bond donors or acceptors in the minor groove of the DNA.

These distances are at the ends of the polyamide and the DNA strand, where water molecules were identified in the minor groove in two of the other DNA/polyamide simulations. In this simulation, water molecules were also found to be hydrogen bonding to the DNA or polyamide in this region, causing the distances to be large. Not only does water get into the groove at the ends, but the identification of two possible orientations of the polyamide tails in the x-ray structure indicates that there is a lot of conformational freedom in the tails. In one orientation, the tail is inside the minor groove where it can hydrogen bond with the DNA, and in the other orientation, the tail is projected outside of the minor groove. The simulation does sample the latter conformation, which contributes to the large distances observed. new parameters was a 10 ns crystal simulation of a $\mathrm{DNA} /$ netropsin complex and a 10 ns crystal simulation of a DNA/distamycin complex. Netropsin and distamycin are both antibiotics that bind in the minor groove of DNA, preferentially to A-T rich sequences. Netropsin covers $\sim 4$ base pairs; distamycin covers $\sim 5$ base pairs. Like the Dervan polyamides, they contain pyrrole heterocycles linked by an amide backbone, and they recognize the minor groove by the hydrogen bonding pattern. Figure 3.30 shows the molecular structures of netropsin and distamycin and how each one binds in the minor groove of DNA.




Figure 3.30 Starting structure for DNA/netropsin simulation $\left(\mathrm{PDB}^{46} \mathrm{id} 121 \mathrm{D}^{60}\right.$; top left) and the molecular structure of netropsin (top right). The bottom panel shows the starting structure for the DNA/distamycin simulation ( $\mathrm{PDB}^{46} \mathrm{id} 2 \mathrm{DND}^{61}$; bottom left) and the molecular structure of distamycin (bottom right).

DNA/netropsin complex. The RMSD plot of the DNA and netropsin complex is shown below (Figure 3.31). The average RMSD of the complex was $1.85 \pm 0.29 \AA$. The DNA contributes the most to the RMSD (average RMSD was $1.83 \pm 0.31 \AA$ ), while the netropsin has a lower RMSD (average RMSD was $0.91 \pm 0.12 \AA$ ). The DNA/netropsin simulation was performed at 300 K , and the RMSD for the DNA/netropsin complex is in the same range as the DNA/polyamide high temperature simulations. The geometric parameters of the netropsin molecule were examined and compared back to the x-ray structure, and the simulated structure is in good agreement with experiment. The average bond error was $0.03 \pm 0.02 \AA$, the average angle error was $4 \pm 3^{\circ}$, and the average dihedral error was $8 \pm 12^{\circ}$. The large dihedral error stems from the rotation about the carbon-carbon bonds at the ends of the netropsin molecule. An overlay of the x-ray structure and the simulation average structure is shown in Figure 3.31. The tails of the molecule, especially the amidino group tail which contains two carbon-carbon single bonds, do not align with each other as the pyrrole rings in the middle. Free rotation is allowed about these carbon-carbon bonds, and this is responsible for the large dihedral angle error. The B-factors reported for the x-ray structure are slightly higher for the atoms of the tails of netropsin, ${ }^{60}$ indicating that there is some disorder in the crystal for the tails. This is most likely observed because the tails contain C-C single bonds, and free rotation is allowed about these bonds; thus, the tail could adopt different orientations, leading to uncertainty in the positions of the atoms.


Figure 3.31 RMSD vs. time of DNA/netropsin complex (left). Overlay of x-ray structure (red) of netropsin and simulation average structure (green) (right).

The helical parameters of the DNA were examined from the DNA/netropsin simulation. The analysis was performed in the same way as for the DNA/polyamide complex simulations, using the t -test at the $95 \%$ confidence level. None of the helical parameters were found to be significantly different from those of the x -ray structure (a table of the calculated helical parameters and $t$-values can be found in Appendix A). The hydrogen bond distances from netropsin to the DNA were also examined (Figure 3.32). As for the DNA/polyamide complex, a t-test was used to analyze the hydrogen bonding distances (table of calculated distance is in Appendix A). For the DNA/netropsin complex, three of the 11 distances were found to be significantly different from the values determined from the x-ray structure.


Figure 3.32 Hydrogen bond distances (numbered) from netropsin to DNA. The figure shows the middle eight base pairs of the DNA fragment and the hydrogen bonds formed from the netropsin N atoms to the atoms of the DNA bases (h-bonds represented by dashed lines). The netropsin molecule is represented by the numbered Ns in the middle of the figure. Red dashed lines show significant differences between experiment and simulation.

The three distances that were found to be significantly different from the x-ray structure are the G10 (N3) to netropsin (N1) distance, the T21 (O3') to netropsin (N9) distance, and the A6 (O3') to netropsin (N10) distance. These three distances involve nitrogens on the end of the netropsin molecule. It is observed that the netropsin moves most at the ends because free rotation is allowed at the single bonds. Therefore, it is not surprising that the length of these hydrogen bonds change more over time. Also, as in the simulations with DNA/polyamide complexes, there were water molecules found near the ends of the netropsin, which form hydrogen bonds with $\mathrm{N} 1, \mathrm{~N} 9$, and N 10 of netropsin. In the x-ray structure, no water molecules were identified as being hydrogen bonded to these three nitrogens of netropsin (they are hydrogen bonded to the DNA bases), but as the simulation progresses, those three nitrogens form hydrogen bonds with water molecules rather than the DNA. The water molecules appear to exchange during the simulation, but the netropsin is still hydrogen bonding to them (distances in the range of $\sim 2.6-3.3 \AA$ ).

DNA/distamycin complex. The RMSD plot of the complex, DNA, and distamycin is shown below (Figure 3.33). The average RMSD of the complex was $2.33 \pm 0.27 \AA$. The DNA contributes the most to the RMSD (average RMSD was $2.41 \pm 0.27 \AA$ ), while the distamycin has a lower RMSD (average RMSD was $0.75 \pm 0.10 \AA$ ). The cause of the high RMSD does not appear to be the distamycin; one base from the helix that does not interact with the distamycin flipped out slightly from the helix (THY7). The movement of this base did not appear to affect the DNA/distamycin distances or helical parameters greatly. The geometric parameters of the distamycin molecule were examined and compared back to the x-ray structure, and the simulated structure is in good agreement with the x-ray structure. The average bond error was $0.04 \pm 0.04 \AA$, the average angle error was $4 \pm 3^{\circ}$, and the average dihedral error was $12 \pm 24^{\circ}$. The dihedral error stems from the rotation about the carbon-carbon bonds at the ends of the distamycin molecule. An overlay of the x-ray structure and the simulation average structure is shown in Figure 3.33. The tails of the molecule, especially the formamide group tail where the carbonyl oxygens are turned in opposite directions, do not align with each other as the pyrrole rings in the middle.


Figure 3.33 RMSD vs. time of DNA/distamycin complex (left). Overlay of x-ray structure (red) of distamycin and simulation average structure (green) (right).

The helical parameters of the DNA were examined from the DNA/distamycin simulation. The analysis was performed in the same way as for the other simulations. One helical parameter was found to be significantly different than the x-ray value, which was the opening parameter. This statistical significance can be attributed to the flipping out of Thy7. The hydrogen bond distances from distamycin to the DNA were also examined (Figure 3.34). For the DNA/distamycin complex, one of the 7 distances were found to be significantly different from the values determined from the x-ray structure.


Figure 3.34 Hydrogen bond distances from distamycin to DNA (numbered). The figure shows the middle eight base pairs of the DNA fragment and the hydrogen bonds formed from the distamycin N atoms to the atoms of the DNA bases (h-bonds represented by dashed lines). The distamycin molecule is represented by the numbered Ns in the middle of the figure. The W indicates a water molecule found in the minor groove. Figure adapted from Coll, et al. ${ }^{61}$

The distance found to be significantly different from the x-ray structure was the A4 (N3) to distamycin (N9) distance. This distance involves a nitrogen on the end of distamycin. It is observed that the distamycin moves most at the ends because free rotation is allowed at the single bonds. Therefore, it is not surprising that this nitrogen may not maintain the hydrogen bond with DNA. Instead, there are water molecules in the groove near distamycin forming a hydrogen bond with N9. The water molecules appear to be exchanging, but they are still forming hydrogen bonds to N9 (distances were in the range
of 2.7-3.5 $\AA$ ). The water molecule identified in the minor groove that forms a hydrogen bond to Thy9 and Ade17 also exchanges with other waters during the simulation. Initially, the water molecule is $\sim 2.6 \AA$ from Thy9 (O2) and $3.1 \AA$ from Ade17 (N3). During the simulation, the water molecule exchanges with other waters but the water molecule always forms a hydrogen bond to the DNA bases. The distance between the water and the DNA bases changes slightly (distances range from 2.8 to $4.1 \AA$ ) and in some snapshots there are two water molecules (one forming a hydrogen bond to Thy9 and one forming the hydrogen bond to Ade17); however, the same trend is usually followed: the distance from Thy9 to the water is usually shorter than the distance from Ade 17 to the water.

### 3.2.6. DNA/Netropsin Crystal Simulation with AMBER Parameters. A 10 ns

 crystal simulation of the DNA/netropsin complex was also performed using the previously reported parameters for netropsin ${ }^{62}$ for use with the AMBER ${ }^{63}$ force field. The netropsin molecule was not explicitly parameterized for the force field; instead, analogous parameters in the force field were adopted for netropsin. The same starting conditions (same structure, same number of water molecules, same system preparation) were used as with the CHARMM simulation. The RMSD plot for the DNA/netropsin complex is in Figure 3.35. The RMSD of the complex was $2.17 \pm 0.21 \AA$, and the DNA contributes more to the $\operatorname{RMSD}(2.16 \pm 0.20 \AA)$ than the netropsin $(1.16 \pm 0.40 \AA)$. The RMSDs of all parts of the system are higher than those obtained from the CHARMM parameters. There is also a large jump in RMSD for the netropsin at $\sim 8 n s$. This is due to the large structural distortion of the netropsin relative to the starting x-ray structure (Figure 3.31).RMSD vs. time (DNA/netropsin complex--AMBER)



Figure 3.35 RMSD vs. time of DNA/netropsin complex using AMBER parameters (left). Comparison (right) of x-ray structure (red) of netropsin and simulation average structure (green).

The average simulation structure shows the failure of agreement of the AMBER parameters with the x-ray structure. The pyrrole ring on the left of the molecule is distorted, and the force constants for the connection between the pyrrole rings is obviously not large enough to keep the bond length correct (the $\mathrm{VMD}^{64}$ viewer did not even draw a bond). The bond length between the amide nitrogen and the carbon of the pyrrole ring should be $1.44 \AA$ (x-ray), but the average simulation structure yields $2.15 \AA$. The force constants are also flawed for the connection of the methyl groups to the nitrogen of the pyrrole rings. In the picture on the right, the green dots are the methyl groups that are supposed to be attached to the pyrrole rings (again, the viewer did not even draw a bond). The length in the x-ray structure for those bonds is $\sim 1.50 \AA$, but the average simulation structure yields an average value of $3.76 \AA$.

The helical parameters and DNA-netropsin distances were also examined using the t-test. Of the 18 helical parameters measured, nine of them are statistically significant from the x-ray value. The significant parameters are $x$-displacement, inclination, stretch, stagger, opening, slide, twist, helical rise, and helical twist. Despite the fact that the
netropsin molecule does not stay bonded together, it does stay in the minor groove of the DNA. For the DNA/netropsin distances, 3 out of the 11 distances measured were statistically significant. These three distances are also found to be statistically significant using the CHARMM parameters (tables of calculated helical parameters and distances can be found in Appendix A). Overall, the AMBER parameters do not allow for a good comparison back to experiment; therefore, an explicit parameterization of the compounds of interest must be performed rather than simply taking analogous parameters from existing molecules in the force field.

### 3.3 Conclusions

New force field parameters for the CHARMM force field for DNA minor groove binding pyrrole-imidazole polyamides have been developed. To test the parameters, crystal simulations of DNA with different polyamides bound in the minor groove have been performed.

The computed helical parameters were in excellent agreement with experiment (17 out of the 18 examined were not statistically significant). The stagger helical parameter was the only parameter that was statistically different from the x-ray structure in both sets of DNA/polyamide simulations (113K and 300 K ). However, upon running additional trajectories, the sampling of stagger was improved. Overall, the low temperature simulations still yield statistical significance with respect to stagger, and stagger is still significant for the high temperature simulations at the $95 \%$ confidence level but not at the $97 \%$ confidence level. When only the three additional high temperature simulations are considered, stagger is not statistically significant. Temperature obviously influences the sampling of stagger; the higher temperature allows
for better sampling and thus non-statistically significant results according to the t -test. The fact that the high and low simulations do not show non-significance overall for all helical parameters underscores the importance of running multiple trajectories.

The high temperature simulations yield structures that are distorted and exhibit ADNA features, and as a result, the polyamides are pushed out of the minor groove in these structures. The temperature of the simulations is too close to that of the melting temperature of the DNA fragment, causing the DNA to start melting, which gives rise to the distorted structures. However, the high temperature structures are still in good agreement with experiment; most helical parameters are non-significant, and we do not observe the rigid structure of the crystal. The authors of the x-ray structure most likely could not obtain crystals at 300 K either. The low temperature simulations yield solid and undistorted structures, much like the starting x-ray structure. Again, we get good agreement with experiment because we see the crystal at 113 K . Although the distortion is observed, the high temperature simulations sample better than the low temperature simulations, thus, stagger is sampled better at the high temperature and yields better t-test results.

The DNA/netropsin, DNA/distamycin, and DNA/hydroxypyrrole polyamide simulations show excellent agreement with the x-ray structure. The DNA/netropsin simulation and the DNA/hydroxypyrrole did not exhibit any statistically significant helical parameters. Opening was the only significantly different parameter in the DNA/distamycin simulation. This can be attributed to one base (Thy7) that flips out slightly out from the rest of the helix during the simulation. This slight flipping out did not affect the DNA/distamycin distances measured or the other helical parameters.

The distances measured from the atoms in the polyamides to their hydrogen bonding partners in the DNA show good agreement with the x-ray structure with the exception of two distances over all simulations. In one low temperature simulation, a water molecule was discovered to be hydrogen bonding to the H of an NH group in the polyamide backbone; as a result, the NH group was not hydrogen bonding to the DNA. In one high temperature simulation, a water molecule was again found to be hydrogen bonding to the polyamide; however, the water molecules appear to be exchanging with each other, nonetheless forming a hydrogen bond with the polyamide. The water molecules are therefore competing with the DNA for hydrogen bonding with the polyamide (which could happen with any molecule that relies on hydrogen bonding for recognition and binding to its target site), which is a possible factor for why some polyamides show reduced binding affinity for their target sequences. Even though individual simulations exhibited statistical significance for two of the DNA/polyamide distances, when the distances are averaged together over all low temperature and all high temperature simulations, none of the distances are significantly different. The distance measurements, like the helical parameters, show the utility of multiple trajectories; the behavior of the intruding water molecules may have gone unnoticed with only one simulation.

The DNA/distamycin, DNA/netropsin, and DNA/hydroxypyrrole polyamide simulations exhibited similar behavior with respect to water molecule hydrogen bonding to the ligand. The DNA/hydroxypyrrole simulation had two distances that were significantly different from the x-ray distances. These distances were at the ends of the DNA groove, which is accessible to water molecules. As in the other DNA/polyamide
simulations, water molecules were observed to be forming hydrogen bonds to the DNA or polyamide. The DNA/netropsin simulation had three distances that were significantly different from the x-ray structure and the DNA/distamycin simulation had one. These distances were from nitrogens on the ends of the ligands to DNA bases. Water molecules were observed to form hydrogen bonds to these nitrogens and although the water molecules appear to exchange during the simulation, they still form hydrogen bonds with the ligand. This does not seem unusual because the minor groove of DNA normally contains a spine of hydration, so water is present in the minor groove. The netropsin and distamycin ligands replace this spine of hydration in the middle of the groove, but water molecules were identified in the x-ray structures near the ends of the grooves closer to the ends of the ligands. ${ }^{60,61}$

The simulation performed using the $\mathrm{AMBER}{ }^{63}$ parameters emphasizes that explicit parameterization of the force field for new molecules is necessary. The netropsin molecule was not explicity parameterized for the AMBER force field ${ }^{63}$; instead, the parameters were guessed at by simply using analogous parameters for existing molecules in the force field. The simulation exhibited large fluctuations in the RMSD of the DNA/netropsin complex, and this was due to large structural distortions of the pyrrole rings of netropsin. The guessed bond force constants were incorrect because the bond lengths of the methyl groups to the pyrrole rings are much too long for a nitrogen-carbon bond. The AMBER parameters did not yield a satisfactory comparison back to experiment; thus, unique and explicit parameterization must be performed.

The simulations performed have validated the new force field parameters for the polyamides, and new simulation studies can now be performed to further understand the
interactions and dynamics between polyamides and DNA. From this knowledge, new ligands can be designed that have a higher affinity for or bind better to their target site (new molecules or different heterocycles or linkages may improve the current design of polyamides). Also, new ligands can be designed that use polyamides as a sequencespecific recognition element that incorporate another type of molecule that will distort the DNA in a way that transcriptional proteins cannot bind to their target site.

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## Chapter 4

## Force Field Parameters for the Simulation of Conjugated Dienes with a Focus on Retinoids

This chapter describes the parameterization procedure for the retinoid molecules, the results of the crystal simulations performed to test the parameters, and conclusions derived from the simulations.

### 4.1 Parameterization of Retinoids

CHARMM force field topology and parameters for $\pi$-conjugated systems were developed and tested against four retinoid compounds (retinoic acid, retinol, fenretinide, and retinal). The parameterization strategy involved the subdivision of key retinoid functionality into 17 small organic (model) compounds. In force field development, one possible approach is to divide a larger molecular system into smaller representative, or model compounds that contain the important functionalities. ${ }^{1}$ MacKerell and coworkers have used quantum mechanical data to parameterize various classes of molecules for the CHARMM force field. ${ }^{2-12}$ Quantum mechanical data has also been used to parameterize the Merck Molecular Force Field, ${ }^{13}$ the AMBER force field, ${ }^{14}$ the OPLS-AA force field,,${ }^{15,16}$ the GROMOS force field, ${ }^{17}$ and the CVFF force field. ${ }^{18}$ Typically, experimental data is critical for force field development. However, geometries derived from quantum mechanical data are particularly useful for force field parameterization, since the time or length scale of experimental studies could introduce significant error or prohibit the measurement process. ${ }^{13,19}$

Model compounds must be selected so that they capture the important bond and nonbond attributions of the molecular system, and be small enough to be computationally tractable. The model compounds are shown in Figure 4.1. Models $\mathbf{1 - 5}$ were selected because all retinoids contain a chain of alternating single and double bonds with methyl groups attached to various carbons on the chain. Structures $\mathbf{6}$ and $\mathbf{7}$ are for use in the fenretinide molecule, which contains amide and phenol groups at the tail end. The cyclohexene ring portion that is in all retinoid compounds was modeled by $\mathbf{8}$ and 9 . The tail ends were parameterized using $\mathbf{1 0 - 1 2}$. Model $\mathbf{1 0}$ is for retinoic acid, $\mathbf{1 1}$ is for retinol, and $\mathbf{1 2}$ is for retinal. Models $\mathbf{1 3}$ and $\mathbf{1 4}$ were parameterized to examine rotation/energetics about the double bonds in retinoids. Models 15-17 were parameterized for the Schiff base linkage for retinal (when retinal binds to rhodopsin or bacteriorhodopsin, it binds covalently to a lysine residue through a Schiff base linkage). Models $\mathbf{1 8}$ and 19 are used to create patches for making the covalent linkage from retinal to the lysine of the protein; the patches simply delete the hydrogens on the nitrogen of the side chain of lysine and the oxygen of the carbonyl group on the tail of retinal and make a double bond between the lysine nitrogen and the carbonyl carbon of retinal.

1

5

6

7


3

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12

13

8


14


17


18


19

Figure 4.1 Model compounds and patch residues for retinoids (some hydrogens not shown for clarity). Models 1-17 are the model compounds, and 18 and 19 are used to create patches for forming Schiff bases.

New atom types were created to allow for proper treatment of alternating single and double bonds. The previous CHARMM parameters used only two atom types for conjugated systems, CE1 for internal carbons and CE2 for terminal carbons. Previously, alternation between single and double bonds was not possible with only one atom type for the internal carbons (Figure 4.2). The previous parameters (when the structure was energy minimized) showed only a $0.002 \AA$ difference between the carbon-carbon single and double bonds, whereas microwave ${ }^{20}$ and infrared data ${ }^{21}$ resulted in a difference of $0.13 \AA$, Raman data gave a difference $0.14 \AA,{ }^{22}$ and electron diffraction data showed a difference between 0.12 to $0.15 \AA \AA^{23-26}$


Figure 4.2 Previous 1,3,5-hexatriene carbon atom types (top). New 1,3,5hexatriene atom types (bottom).

New atom types for $\pi$-conjugated systems have been created, which are CC 2 for terminal carbons and CC1A and CC1B for internal carbons (Figure 4.2). Each model compound was energy minimized at the MP2/6-31G(d) level of theory. ${ }^{13,27,28}$ The average MP2/6-31G(d) carbon-carbon bond lengths over all model compounds was 1.50 $\pm 0.03 \AA$ for single bonds, and $1.35 \pm 0.02 \AA$ for double bonds. The average bond lengths calculated from MP2 were used as the equilibrium bond lengths for the new CHARMM parameters for the associated bond type. The force constants were based on MP2/6-31G(d) frequency computations for $\mathbf{1}, \mathbf{1 3}, \mathbf{1 5}$, and 16 in Figure 4.1 as an initial guess. Each model compound was energy minimized with the CHARMM force field with steepest descent and Newton-Raphson (tolerance gradient was $10^{-6} \mathrm{kcal} / \mathrm{mol} / \AA$ for each minimization). The parameters where then adjusted to minimize error with the MP2/6-31G(d) structure. The average value for the CHARMM single bonds is $1.51 \pm$ $0.02 \AA$, and the average value for double bonds is $1.36 \pm 0.02 \AA$. The average error between MP2 and CHARMM for single bonds is $0.01 \pm 0.01 \AA$, and $0.005 \pm 0.006 \AA$ for double bonds. The average error between MP2 and CHARMM over all bonds (not just carbon-carbon bonds) in the model compounds is $0.012 \pm 0.009 \AA$.

New angle terms and angle modifications were also made. The minimized angle values obtained from the MP2 optimized structures were used as the equilibrium angle
values for the new CHARMM parameters. The force constants were based on MP2/6$31 \mathrm{G}(\mathrm{d})$ frequency computations for $\mathbf{1 , 1 3}, \mathbf{1 5}$, and $1 \mathbf{6}$ in Figure 4.1 as an initial guess. Each model compound was energy minimized with the CHARMM force field with steepest descent and Newton-Raphson (tolerance gradient was $10^{-6} \mathrm{kcal} / \mathrm{mol} / \AA \AA$ ). The parameters were adjusted to minimize the error across all model compounds. The average error between MP2 and CHARMM was calculated for angles as $2 \pm 2^{\circ}$.

Torsion angle terms were also reparameterized. The torsion angles chosen were selected because they involve low energy barriers that are important in describing the dynamics allowing the retinoids to interact with various proteins; the single bonds and any bonds linking the conjugated hydrocarbon chain of the retinoids to a functional group must be able to rotate in order for the retinoid molecule to maximize favorable interactions with amino acids in the binding site of the protein. For each model compound, the potential energy surfaces (PES) for selected torsion angles (see Figures 4.3-4.9 for selected torsion angles; the bond with the arrows were the bonds which were rotated about) were calculated by holding the selected torsion angle fixed at different increments $\left(20^{\circ}\right.$ increments from $-180^{\circ}$ to $180^{\circ}$ ) and geometry optimizing all other degrees of freedom. The torsion angles selected were for rotations about the single bond between two double bonds, rotation of a particular functional group that is necessary for optimizing ligand-protein contacts, or rotations necessary for ring puckering. The CHARMM dihedral angle parameters force constant $\left(K_{\chi}\right)$, multiplicity ( $n$ ), and phase shift ( $\delta$ ) were then modified to match the quantum mechanical results from the potential energy surfaces. Each model compound was energy minimized in CHARMM (steepest descent and Newton-Raphson; tolerance gradient was $10^{-6} \mathrm{kcal} / \mathrm{mol} / \AA$ ) with the selected
torsion angle held fixed, and the potential energy surface of the dihedral was calculated. If the surface did not match the MP2 results, then the multiplicity, phase shift, or force constant were adjusted as necessary.

For models 1-5 in Figure 4.1, the single bond between the two double bonds of the molecules was rotated from $-180^{\circ}$ to $180^{\circ}$ in $20^{\circ}$ increments. Figure 4.3 shows the torsion angle results for 1,3-butadiene and 1,3-pentadiene.

1,3-butadiene (1)



Figure 4.3 Torsion energy plots for 1,3-butadiene (top) and 1,3-pentadiene (bottom). Diamonds, final CHARMM; square, MP2/6-31G(d).

The initial CHARMM surface for 1,3-butadiene had barriers at $\sim 64 \mathrm{kcal} / \mathrm{mol}$ (Appendix B), which is $\sim 58 \mathrm{kcal} / \mathrm{mol}$ higher than the MP2 barriers; the initial CHARMM also did not show the correct multiplicity. This was corrected by changing the force constant ( $K_{\chi}$ )
to lower the barriers and the multiplicity $(n)$ to get the correct number of barriers. The final CHARMM results show the correct phase and multiplicity as the MP2 results. The CHARMM energy barrier at $0^{\circ}$, which corresponds to cis-1,3-butadiene, is $1.03 \mathrm{kcal} / \mathrm{mol}$ higher than the MP2 results, but overall, the results agree well.

1,3-Butadiene has been examined by several experiments, including electron diffraction, ${ }^{23-26} \mathrm{UV},{ }^{29} \mathrm{IR},{ }^{21,30,31}$ microwave, ${ }^{32}$ Raman, ${ }^{22,33}$ and computational studies. ${ }^{34-38}$ The trans form is the dominant conformation, and the cis $\left(0^{\circ}\right)$ or gauche $\left(\sim 40^{\circ}\right)$ conformation is the minor component. However, there is still debate on whether the cis or gauche conformer is the more stable minor conformer. ${ }^{39-44}$ Computational studies show that the cis form is a "transition structure" between the two gauche conformers, but generally, experimental studies say that the cis form is the minor conformer. ${ }^{31,45}$ This difference has been attributed to conditions under which the torsional potential was calculated or measured. Theoretical and vapor phase spectroscopy results favor the gauche form as the second stable conformer, while spectroscopic measurements performed in an argon matrix favor the cis form as the second stable conformer of 1,3butadiene (the interaction with the surrounding Ar matrix stabilizes the cis form and thus changes the torsional potential). ${ }^{46}$ Also, it is difficult to distinguish the cis and gauche forms from one another by an experimental direct measurement because only $\sim 1 \%$ of the minor conformer is present at room temperature, and 1,3-butadiene is reactive at higher temperatures. ${ }^{42,44}$

For 1,3-pentadiene (bottom of Figure 4.3), the initial CHARMM had energy barriers that are too high ( $\sim 92 \mathrm{kcal} / \mathrm{mol}$; see Appendix B), and the multiplicity was incorrect. The energy barriers were lowered by decreasing the value of the force constant
$\left(K_{\chi}\right)$ to $0.5600 \mathrm{kcal} / \mathrm{mol} / \AA^{2}$, and the correct number of barriers was obtained by changing the multiplicity $(n)$ to 2 . The final CHARMM surface shows the correct phase and multiplicity, and the energy barriers are in good agreement with the MP2 results. The CHARMM energy barriers at $100^{\circ}$ and $-100^{\circ}$ are $\sim 0.58 \mathrm{kcal} / \mathrm{mol}$ lower than the MP2, and the CHARMM energy barrier at $0^{\circ}$ is $\sim 0.34 \mathrm{kcal} / \mathrm{mol}$ lower than the MP2.

Figure 4.4 shows the torsion angle results for the conjugated methylated butanes and pentenes. As for 1,3-butadiene, all of the methylated dienes have two minima (gauche conformers) in addition to the global minimum (trans conformer) on their potential energy surfaces. Also, the initial CHARMM for the methylated dienes shows very high energy barriers and incorrect multiplicities (Appendix B). 2-Methyl-1,3butadiene shows a maximum at $116 \mathrm{kcal} / \mathrm{mol}$ for the initial CHARMM, 2-methyl-1,3pentadiene shows a maximum at $125 \mathrm{kcal} / \mathrm{mol}$, and 4-methyl-1,3-pentadiene has a maximum at $93 \mathrm{kcal} / \mathrm{mol}$. The final CHARMM parameters are in good agreement with the MP2 results. For 2-methyl-1,3-pentadiene, the CHARMM barrier height is 0.62 $\mathrm{kcal} / \mathrm{mol}$ lower than MP2 at $0^{\circ}$ and $0.57 \mathrm{kcal} / \mathrm{mol}$ lower than MP2 at $-100^{\circ}$ and $100^{\circ} .4-$ Methyl-1,3-pentadiene has a CHARMM surface that is $0.90 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 at $0^{\circ}$ and $0.76 \mathrm{kcal} / \mathrm{mol}$ higher at $100^{\circ}$. 2-Methyl-1,3-butadiene has good agreement between CHARMM and MP2 with CHARMM having barriers only $0.28 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 at the barrier at $0^{\circ}$.



Figure 4.4 Torsion energy plots for 2-methyl-1,3-butadiene (top left), 4-methyl-1,3pentadiene (top right), and 2-methyl-1,3-pentadiene (bottom). Diamonds-final CHARMM surfaces; square-MP2/6-31G(d) surface.

Model compounds $\mathbf{6}$ and $\mathbf{7}$ are for use in the fenretinide molecule, which contains and amide and phenol group at the tail end.

p-acetamide phenol (7)



Figure 4.5 Torsion energy plot for N-3-dimethyl-2-butenamide (top) and p-acetamide phenol (bottom). Diamonds; MP2/6-31G(d); squares; final CHARMM.

N-3-dimethyl-2-buteneamide shows high energy barriers for the initial CHARMM (Appendix B) and also shows the wrong phase and multiplicity. The final CHARMM shows much better agreement with MP2, with the CHARMM being $0.86 \mathrm{kcal} / \mathrm{mol}$ lower than MP2 at the maxima $\left(180^{\circ}\right.$ and $\left.-180^{\circ}\right)$.

The initial CHARMM energy barriers for $p$-acetamide phenol, 7, are too high ( $\sim 13 \mathrm{kcal} / \mathrm{mol}$; Appendix B), but the phase and multiplicity are correct. The final CHARMM maxima at $80^{\circ}$ are in better agreement with MP2, and the minimum at $0^{\circ}$ is $\sim 0.80 \mathrm{kcal} / \mathrm{mol}$ higher than MP2.

Models 8 and 9 were parameterized for the cyclohexene ring portion that is in all retinoid compounds.

## 1,6,6,-trimethyl-2-ethene-cyclohexene (8)



1,2,6,6,-tetramethylcyclcohexene (9)




Figure 4.6 Torsion energy surfaces for 1,6,6-trimethyl-2-ethene-cyclohexene (top) and 1,2,6,6-tetramethylcyclohexene (bottom). Diamonds, MP2/6-31G(d); squares, final CHARMM.

The initial CHARMM parameters for the 1,6,6-trimethyl-2-ethene cyclohexene behave like the others; the energy barriers are too high ( $\sim 120 \mathrm{kcal} / \mathrm{mol}$; Appendix B). The final CHARMM parameters are in better agreement, but they show an increase in energy at $150^{\circ}$ and $-150^{\circ}$, whereas MP2 does not. Also, CHARMM is $\sim 1.1 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 at the maximum at $0^{\circ}$. The dihedral surface for $1,2,6,6-$
tetramethylcyclcohexene was examined for the purpose of deciding which ring pucker was more favorable, C 2 -endo ( C 2 carbon above the plane of the ring) or C 3 -endo ( C 3 carbon above the plane of the ring). The side of the ring containing the double bond stays planar, but the other side of the ring can move and pucker with either the C 2 or C 3 above the ring plane. The dihedral examined was rotated from $-50^{\circ}$ to $50^{\circ}$ in $10^{\circ}$ increments (the dihedral is constrained being in the cyclohexene ring). When the dihedral is at $-10^{\circ}$, the ring is in the C 3 -endo conformation, and when the dihedral is at $10^{\circ}$, the ring is in the C 2 endo conformation. Since both conformations are at a minimum on the surface, each structure was geometry optimized to find which conformation was lowest in energy. The MP2 calculations showed that the both conformations were nearly the same in energy; the C2-endo structure was only $3.00 \times 10^{-4} \mathrm{kcal} / \mathrm{mol}$ lower in energy than the C3-endo structure. The initial and final CHARMM parameters are almost the same for this model compound (Appendix B), and after revision, the final CHARMM parameters are only 0.5 $\mathrm{kcal} / \mathrm{mol}$ higher than the initial parameters at $50^{\circ}$. MP2 is $\sim 1.3 \mathrm{kcal} / \mathrm{mol}$ higher than CHARMM at $10^{\circ}$ and $-10^{\circ}$.

Models 10, 11, and 12 were parameterized for the tail ends of the retinoids. Model 10 is for retinoic acid, Model 11 is for retinol, and Model 12 is for retinal. For Models 10, 11, and 12 , rotation to obtain the surfaces was about the $\mathrm{C}-\mathrm{C}$ single bond in the middle of the structure.
2-propenoic acid (10)







Figure 4.7 Torsion energy surfaces for 2-propenoic acid (top left), 2-propenol (top right), and 2-propenal (bottom). Diamonds, final CHARMM; squares, MP2/6-31G(d).

The initial CHARMM parameters for 2-propenoic acid, 2-propenol, and 2-propenal do not show the same magnitude of high energy barriers as the conjugated dienes.

Nevertheless, the initial parameters do not agree with the MP2 results (see Appendix B).

The final CHARMM parameters for 2-propenoic acid are in good agreement; CHARMM shows the correct minima (at $0^{\circ}, 180^{\circ}$, and $-180^{\circ}$ ) and maxima; CHARMM is only 0.08 $\mathrm{kcal} / \mathrm{mol}$ higher than MP2 at the maxima. For 2-propenol, the final CHARMM parameters are $1.1 \mathrm{kcal} / \mathrm{mol}$ higher than MP2 at $180^{\circ}$ and $-180^{\circ}$, and they give the correct phase and multiplicity. The final CHARMM parameters for 2-propenal are in also in good agreement with MP2, showing the correct barrier heights; CHARMM is only 0.5 $\mathrm{kcal} / \mathrm{mol}$ higher than MP2 at the maxima.

In general, the initial CHARMM parameters delivered unrealistic torsional energetic barriers, which were parameterized by changing the force constants and multiplicities to match the MP2 data. For each case, the changes are given in Appendix B. The final CHARMM torsional surfaces agree well with those calculated from MP2. The largest difference observed between MP2 and CHARMM was $1.1 \mathrm{kcal} / \mathrm{mol}$.

Models 13 and 14 were parameterized to examine rotation/energetics about the double bonds in the retinoids. CHARMM follows MP2 at the higher values for the dihedral $\left(120^{\circ}-180^{\circ}\right)$, but as the compound approaches the lower values and gets closer to a cis conformation $\left(105^{\circ}\right.$ down to $\left.45^{\circ}\right)$, CHARMM deviates from MP2, and at $45^{\circ}$, there is a $49 \mathrm{kcal} / \mathrm{mol}$ difference. As for 2-butene, the CHARMM parameters for $1,3,5-$ hexatriene follow MP2 when the molecule has a trans geometry $\left(-180^{\circ}\right.$ to $\left.-120^{\circ}\right)$, but as the structure approaches a more cis conformation about the middle double bond, CHARMM starts to deviate from MP2. At $60^{\circ}$, CHARMM even starts to show a decrease in energy. This deviation is due to a conformational change of the end groups that are on either side of the double bond being rotated. The CHARMM structures do not show as much movement of the end groups.
2-butene (13)




Figure 4.8 Torsion energy surfaces for 2-butene (left) and 1,3,5-hexatriene (right). Diamonds, final CHARMM; squares, MP2/6-31G(d).

Models 15, 16, and 17 in Figure 4.1 were parameterized for the Schiff base linkage for retinal. When retinal binds to rhodopsin or bacteriorhodopsin, it binds covalently to a lysine residue through a Schiff base linkage. Model 15 is for a deprotonated Schiff base linkage, and 16 and 17 are for protonated Schiff base linkages. The CHARMM for the first two Schiff base compounds (Figure 4.9) follows the MP2 up to the last two points tested, and then it starts to deviate by about $9 \mathrm{kcal} / \mathrm{mol}$. The double bond was used to examine the rotation/energetics for the Schiff base linkage in retinal. The CHARMM parameters for the third Schiff base compound (Figure 4.9) are in excellent agreement with MP2; the multiplicity and phase are correct, and the barrier heights coincide almost exactly.


Schiff base--protonated (17)




Figure 4.9 Torsion energy plot for deprotonated Schiff base (top left) and protonated Schiff base (top right and bottom). Diamonds, final CHARMM; squares, MP2/6-31G(d).

The last two models in Figure 4.1 (18 and 19) are for patches making the covalent linkage from retinal to the lysine of the protein. The Schiff base linkages they contain were parameterized ( $\mathbf{1 5}, \mathbf{1 6}$, and $\mathbf{1 7}$ ); the patches simply delete the hydrogens on the nitrogen of the side chain of lysine and the oxygen of the carbonyl group on the tail of retinal and make a double bond between the lysine nitrogen and the carbonyl carbon of retinal.

The final CHARMM parameters for the model compounds are in good agreement with the MP2/6-31G(d) results. The initial CHARMM parameters not only treated the alternating single and double bonds incorrectly, but for many of the model compounds, the energy barriers for the torsion angles were too high (over $60 \mathrm{kcal} / \mathrm{mol}$ ), and some were out of phase. The final CHARMM parameters exhibit the correct barrier heights and phases, and largest difference between CHARMM and MP2 was $1.15 \mathrm{kcal} / \mathrm{mol}$ (CHARMM is higher than MP2 for 1,6,6-trimethyl-2-ethene-cyclohexene at the maxima at $-180^{\circ}$ and $\left.180^{\circ}\right)$. The average error between CHARMM and MP2 in bond distances is $0.012 \pm 0.009 \AA$, the average error for angles is $1 \pm 1^{\circ}$, and the average error for dihedrals is $2 \pm 4^{\circ}$.

### 4.2 Comparison to Available Experimental Data

Some of model compounds and the retinoids built from them were compared to available experimental data. Experimental data for all model compounds is not available, nor is experimental data for all four retinoids. The most experimental data was for 1,3butadiene, and this included mainly electron diffraction data; ${ }^{23-26}$ however, there was infrared, ${ }^{21}$ microwave, ${ }^{32}$ and Raman ${ }^{22}$ data as well. Electron diffraction data was also found for 2-methyl-1,3-butadiene ${ }^{47}$ and 1,3,5-hexatriene, ${ }^{25}$ and x-ray diffraction data was
found for $p$-acetamide phenol ${ }^{48}$ and 2-propenoic acid. ${ }^{49}$ For the final retinoids, x-ray structures of all-trans retinoic acid ${ }^{50,51}$ and all-trans retinal ${ }^{52}$ were compared to the final retinoids in CHARMM. The experimental comparisons with the new CHARMM parameters and MP2 results (in the case of the smaller model compounds) for 1,3butadiene are shown in Table 4.1 (CHARMM atom numbering referred to in Table 4.1 is shown above the table).


Table 0.1 Experimental, QM, and MM comparison for 1,3-butadiene.

| Method | C1-C2 | C2-C3 | C3-C4 | C-H (all) | C1-C2-C3 | C2-C3-C4 | C=C-C | C=C-H | C1-C2-C3-C4 |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| ED $^{\text {a }}$ | 1.35 | 1.46 | 1.35 | 1.06 | $124 \pm 2$ | $124 \pm 2$ |  |  | 180 |
| ED $^{\mathrm{b}}$ | 1.337 | 1.483 | 1.337 | 1.082 |  |  | 122.4 | 119.8 | 180 |
| ED $^{\mathrm{c}}$ | 1.344 | 1.467 | 1.344 | 1.094 |  |  | $122.8 \pm 0.5$ | $119.5 \pm 1$ | 180 |
| ED $^{\mathrm{d}}$ | 1.341 | 1.463 | 1.341 | 1.086 |  |  | $123.3 \pm 0.5$ | $121.8 \pm 1.2$ | 180 |
| IR $^{\mathrm{e}}$ | 1.338 | 1.464 | 1.338 | 1.086 |  |  | $123.2 \pm 0.2$ | $119.6 \pm 0.2$ | 180 |
| Microwave $^{\mathrm{f}}$ | 1.337 | 1.467 |  | 1.087 | 123.5 |  |  | $120.9 \pm 1.13$ | 180 |
| Raman $^{\text {g }}$ | 1.337 | 1.476 | 1.337 | 1.085 |  |  | $122.9 \pm 0.5$ | 120 | 180 |
| MP2/6-31G(d) | 1.344 | 1.458 | 1.344 | 1.087 | 123.7 | 123.7 | $123.7 \pm 0.2$ | $120.9 \pm 1.05$ | 180 |
| CHARMM | 1.346 | 1.471 | 1.346 | 1.1 | 123.7 | 123.7 | $123.7 \pm 0.3$ | $120.1 \pm 1.34$ | 180 |

Bonds in $\AA$; angles and dihedrals in degrees.
a see ref. 23
b see ref. 24 ; C=C bond error $0.005 \AA, \mathrm{C}-\mathrm{C}$ bond error $0.01 \AA, \mathrm{C}-\mathrm{H}$ bond error $0.01 \AA$.
c see ref. 25
d see ref. $26 ; \mathrm{C}=\mathrm{C}$ bond error $0.002 \AA, \mathrm{C}-\mathrm{C}$ bond error $0.003 \AA, \mathrm{C}-\mathrm{H}$ bond error $0.04 \AA$.
e see ref. 21 ; C-C bond error $0.003 \AA$.
f see ref. 32
g see ref. $22 \mathrm{C}-\mathrm{C}$ bond error $0.01 \AA$.

From Table 4.1, the MP2/6-31G(d) results and the final CHARMM results compare well to the experimental bonds, angles, and dihedral angle. MP2 and CHARMM both give the trans structure ( $180^{\circ}$ ) of 1,3-butadiene after energy
minimization; all of the experimental techniques find this as well. Also, the average error between MP2 and CHARMM compared with experimental bonds is $0.007 \pm 0.005 \AA$ and $0.009 \pm 0.008 \AA$, respectively, and the average error between MP2 and CHARMM compared with experimental angles is $0.7 \pm 0.5^{\circ}$ and $0.6 \pm 0.4^{\circ}$, respectively. When comparing the computed values to electron diffraction results, it should be kept in mind that the average carbon-carbon bond length can be determined with high accuracy when equivalent bonds are present, but if non-equivalent bonds are present, the individual values can be much less accurate. ${ }^{53}$

Table 4.2 shows the comparison of experimental data with the structures obtained from the final CHARMM parameters for isoprene (2-methyl-1,3-butadiene).


Table 0.2 Experimental, QM, and MM comparison for 2-methyl-1,3-butadiene.

|  |  | C=C | C2-C3 | C2-C5 Csp $^{2}-\mathrm{H} \mathrm{Csp}^{3}-\mathrm{H}$ | C1-C2-C3 | C2-C3-C4 | C1-C2-C5 | H12-C1-C2 | C2-C5-H1 |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ED $^{\mathrm{a}}$ | 1.34 | 1.46 | 1.51 | 1.08 | 1.11 | 121.4 | 127.3 | 121 | 124.3 | 109.1 |
| MP2/6-31G(d) | 1.35 | 1.47 | 1.51 | 1.09 | 1.09 | 122.6 | 122.6 | 119.8 | 121.7 | 111 |
| CHARMM | 1.35 | 1.49 | 1.51 | 1.1 | 1.1 | 121.8 | 128.4 | 119.4 | 121.3 | 114.3 |
| MP2/ED error | 0.01 | 0.01 | 0 | 0.01 | 0.02 | 1.2 | 4.7 | 1.2 | 2.6 | 1.9 |
| CH/ED error | 0.01 | 0.03 | 0 | 0.02 | 0.01 | 0.4 | 1.1 | 1.6 | 3 | 5.2 |
| MP2/CH error | 0 | 0.02 | 0 | 0.01 | 0.01 | 0.8 | 5.8 | 0.4 | 0.4 | 3.3 |

Bonds in $\AA$; angles in degrees.
a see ref. 47; no error reported
For isoprene, the MP2 and CHARMM results compare well to the ED data. The average error for bonds between MP2 and the ED data and CHARMM and the ED data is $0.01 \pm$ $0.01 \AA$ and $0.01 \pm 0.02 \AA$, respectively. The average error for angles compared to the ED
data is $2.5 \pm 1.4^{\circ}$ for MP2 and $3.1 \pm 2.7^{\circ}$ for CHARMM. The MP2 and CHARMM results are also in good agreement; the average error for bonds between MP2 and CHARMM is $0.008 \pm 0.009 \AA$, and the average error for angles is $2.4 \pm 2.2^{\circ}$.



Table 0.3 Experimental, QM, and CHARMM comparison of 2-propenoic acid.

|  | C2-CG | C1-C2 | CG-OD2 | CG-OD1 | OD1-OD2 | OD1-CG-C2 | OD1-CG-OD2 | C1-C2-CG-OD2 | C1-C2-CG-OD1 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| X-ray $^{\text {a }}$ | 1.47 | 1.3 | 1.28 | 1.26 | 2.66 | 116 | 122 | 0 | 180 |
| MP2/6-31G(d) | 1.48 | 1.34 | 1.27 | 1.26 | 2.29 | 113.9 | 130 | 0 | 180 |
| CHARMM | 1.5 | 1.34 | 1.26 | 1.26 | 2.22 | 116.2 | 123.6 | 0 | 180 |
| MP2/x-ray error | 0.01 | 0.04 | 0.01 | 0 | 0.37 | 2.1 | 8 | 0 | 0 |
| CH/x-ray error | 0.03 | 0.04 | 0.02 | 0 | 0.44 | 0.2 | 1.6 | 0 | 0 |
| MP2/CH error | 0.02 | 0 | 0.01 | 0 | 0.07 | 2.3 | 6.4 | 0 | 0 |

Bonds in $\AA$; angles and dihedrals in degrees. h see ref. 49

An x-ray structure has been reported for 2-propenoic acid. ${ }^{49}$ The average error for bonds between CHARMM and the x-ray data is $0.1 \pm 0.2 \AA$ (Table 4.3). This error is larger than the other comparisons when CHARMM is compared to experimental data; and this is probably due to the fact that CHARMM gives bond lengths that are slightly longer than the x-ray data for the carbon-carbon bonds, and one of the carbon-oxygen bonds is shorter than what was found in the x-ray structure. The error between MP2 and the x-ray data is smaller than that of CHARMM, with the error being $0.09 \pm 0.1 \AA$. The error between MP2 and CHARMM is smallest at $0.02 \pm 0.02 \AA$. Two of the MP2 angles are slightly less than the x-ray structure, and the angle between the two oxygens is larger
than the x-ray structure, giving an average error of $5.4 \pm 3.0^{\circ}$. The error between CHARMM and the x-ray structure is smaller for the angles at $1.2 \pm 0.8^{\circ}$. The average angle error between CHARMM and MP2 is $4.3 \pm 2.1^{\circ}$. An overlay of the x-ray structure (red) and the CHARMM minimized structure (blue) is shown above Table 4.3; the RMSD between the two structures is $0.0001 \AA$.

An electron diffraction structure has been reported for 1,3,5-hexatriene. ${ }^{25}$ All three structures give the trans conformation about the double bonds, and most of the bonds and angle from CHARMM and MP2 agree with those of the experimental structure (average bond errors are $0.01 \pm 0.03 \AA$ for MP2 and ED and $0.02 \pm 0.02 \AA$ for CHARMM and ED). However, CHARMM gives slightly longer carbon-carbon single bonds ( $\sim 0.02 \AA$ ) than the x-ray and MP2 structures (Table 4.4).


Table 0.4 Experimental, QM, and MM comparison of 1,3,5-hexatriene

|  | C1-C2 | C2-C3 | C3-C4 | C4-C5 | C5-C6 | C-H (all) | C=C-C | C=C-H | C-C-C-C |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| ED $^{a}$ | 1.345 | 1.45 | 1.345 | 1.45 | 1.345 | 1.103 | $124.3 \pm 1.5$ | 117.8 | 180 |
| MP2/6-31G(d) | 1.346 | 1.451 | 1.345 | 1.451 | 1.346 | 1.088 | $123.8 \pm 0.04$ | 120.4 | 180 |
| CHARMM | 1.346 | 1.475 | 1.346 | 1.476 | 1.345 | 1.1 | $123.8 \pm 0.26$ | 119.8 | 180 |
| MP2/ED error | 0.001 | 0.001 | 0 | 0.001 | 0.001 |  |  | 0 |  |
| CH/ED error | 0.001 | 0.025 | 0.001 | 0.026 | 0 |  |  | 0 |  |
| CH/MP2 error | 0 | 0.024 | 0.001 | 0.025 | 0.001 |  |  |  |  |
| BOM |  |  |  | 0 |  |  |  |  |  |

Bonds in $\AA$; angles and dihedrals in degrees.
a see ref. 25

An x-ray structure was reported for $p$-acetamide phenol, and the bonds, angles, and dihedrals from MP2 and CHARMM agree with those of the x-ray structure (Table
4.5). The average errors between MP2 and the x-ray structure are $0.01 \pm 0.01 \AA$ for bonds, $1.1 \pm 1.3^{\circ}$ for angles, and $4.1 \pm 5.8^{\circ}$ for dihedrals. The average errors between CHARMM and the x-ray structure are $0.02 \pm 0.01 \AA$ for bonds, $1.0 \pm 0.7^{\circ}$ for angles, and $1.7 \pm 1.1^{\circ}$ for dihedrals. CHARMM is in good agreement with the x-ray structure for dihedrals; the largest deviations for MP2 and the x-ray structure are the dihedrals C23-C22-N21-C15 and C27-C22-N21-C15, which are the two dihedrals connecting the amide group to the phenol ring. An overlay of the x-ray structure (red) and the CHARMM minimized structure (blue) is shown above Table 4.5; the RMSD between the two structures is $0.0010 \AA$. For clarity, the ring hydrogens and methyl group hydrogens on C14 are not shown.



Table 0.5 Experimental, QM, and MM comparison of $p$-acetamide phenol ${ }^{\text {b }}$

|  | C22-C27 | C22-N21 | C15-O29 | C27-C26 | C22-C27-C26-C25 | C25-C24-C23-C22 | O28-C25-C26-C27 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| X-ray $^{2}$ | 1.39 | 1.42 | 1.22 | 1.39 | 0.45 | 0.4 | -178.2 |
| MP2/6-31G(d) | 1.4 | 1.41 | 1.22 | 1.39 | 0.09 | 0.07 | -180 |
| CHARMM | 1.41 | 1.41 | 1.23 | 1.39 | 0.26 | 0.18 | -179.9 |
| MP2/x-ray error | 0.01 | 0.01 | 0.01 | 0 | 0.36 | 0.33 | 1.8 |
| CH/x-ray error | 0.02 | 0.01 | 0 | 0.01 | 0.19 | 0.22 | 1.7 |
| CH/MP2 error | 0.01 | 0.01 | 0.01 | 0.01 | 0.17 | 0.11 | 0.1 |

Bonds in $\AA$; dihedrals in degrees.
a see ref. 48
b Only some bonds and dihedrals are shown in the table. For a complete table, see Appendix B.

Two retinoids (retinal and retinoic acid) had available experimental data. The crystal structure of all-trans retinal was taken from the reference of Hamanaka, et al., ${ }^{52}$ (Cambridge Structural Database ${ }^{54}$ id TRETAL01), and the average bond error between x-ray and CHARMM was $0.02 \pm 0.02 \AA$, the average angle error was $1.93 \pm 1.73^{\circ}$, and the average dihedral error was $15.6 \pm 30.5^{\circ}$. The RMSD between CHARMM and x-ray is $1.92 \AA$. The cyclohexene ring of both structures has the same conformation (C2-endo), but several dihedrals in the chain show large differences between the structures; this is what causes the different orientations of the methyl groups (Figure 4.10). There are three dihedrals of the chain that show large differences between the CHARMM minimized and x-ray structures and two dihedrals of the ring that show large differences. These dihedrals are all rotations about single bonds of the chain, and because they are different in the two structures, the orientations of the methyl groups are different (see Appendix B for a table of bonds, angles, and dihedrals for CHARMM and the x-ray structure).



Figure 4.10 The left structure is that of retinal, and the right picture shows an overlay of the x-ray structure (red) and the CHARMM energy minimized structure (blue). The heavy atom RMSD between CHARMM and x-ray is $1.92 \AA$.

The triclinic crystal structure of all-trans retinoic acid was taken from the reference of Stam and MacGillavry ${ }^{51}$ (Cambridge Structural Database ${ }^{54}$ id VITAAC10). The average bond error between crystal and CHARMM is $0.04 \pm 0.09 \AA$, the average angle error is $2.4 \pm 1.7^{\circ}$, and the average dihedral error is $11.8 \pm 20.2^{\circ}$. The RMSD
between x-ray and CHARMM is $2.18 \AA$. The methyl groups of the chain point in the same direction in both structures; however, some dihedrals along the chain show large differences between the two structures, and the cyclohexene rings are in different conformations (Figure 4.11). The x-ray structure shows a C2-endo conformation for the ring, while CHARMM gives a C3-endo conformation (see Appendix B for a comparison of CHARMM to the x-ray structure bonds, angles, and dihedrals).


Figure 4.11 The structure on the right is that of retinoic acid, and the right picture is a comparison of the x-ray structure (red) and the CHARMM energy minimized structure (blue) for the triclinic form of retinoic acid.

The monoclinic crystal structure of all-trans retinoic acid was taken from the reference of Stam ${ }^{50}$ (Cambridge Structural Database ${ }^{54}$ id VITAAC01; search date $2 / 17 / 07$ ). The average bond error between crystal and CHARMM is $0.04 \pm 0.09 \AA$, the average angle error is $1.9 \pm 1.6^{\circ}$, and the average dihedral error is $13.6 \pm 26.0^{\circ}$. RMSD between x-ray and CHARMM is $2.72 \AA$. The cyclohexene rings have the same conformation; both are C3-endo. This is different than the triclinic crystal, where the chain sticks out from the ring, but the cyclohexene ring is in a different conformation ( C 2 -endo in the triclinic form and C 3 -endo in the monoclinic form). However, the chain off the ring is oriented differently and so are the methyl groups. This is the result of the chain dihedrals being different (Figure 4.12; for a table of bonds, angle, and dihedrals, see Appendix B).


Figure 4.12 Comparison of monoclinic crystal and CHARMM forms (red, x-ray; blue, CHARMM). RMSD between x-ray and CHARMM is $2.72 \AA$.

Overall, the CHARMM minimized structures of retinal and retinoic acid agree well with the x-ray structures in terms of bonds and angles (largest average bond error is $0.04 \pm 0.09 \AA$, and the largest angle error is $2.4 \pm 1.7^{\circ}$ ). However, there are some large differences in the dihedral angles, especially in the chain parts of the retinoids. In retinal, both the x-ray structure and CHARMM show the same C2-endo ring conformation for the cyclohexene ring, but three dihedrals of the chain cause the different orientations of the methyl groups. The largest dihedral difference in retinal arises because the last methyl group at the end of the chain (C20) is oriented differently in the CHARMM and x-ray structures. In retinoic acid, CHARMM shows differences with both the triclinic and monoclinic x-ray crystal forms. For the triclinic crystal, the cyclohexene rings are in different orientations (CHARMM gives C3-endo while the x-ray structure gives C2endo). As for retinal, there are several dihedrals of the chain that cause different orientations of the methyl groups and the end carboxyl group. Unlike retinal, however, the differences in the methyl group orientations are not as pronounced for the CHARMM and triclinic forms of retinoic acid. For the monoclinic crystal, the cyclohexene ring of the x-ray structure and the CHARMM structure are in the same C3-endo conformation; however, the chain attached to the ring is oriented differently (the chain from the x-ray structure seems to point up, and the chain from the CHARMM structure points out),
which gives rise to large differences for the dihedrals that connect the ring to the chain. The triclinic and monoclinic crystals show differences with each other and CHARMM. The triclinic crystal and CHARMM have different ring conformations, but the chain off the ring seems to stick out from the ring more, and it doesn't stick up. The monoclinic crystal and CHARMM show the same ring conformations, but the chain off the ring sticks up more in the monoclinic crystal, whereas the CHARMM chain sticks out and not up. The differences between the ring conformations and how the chain is oriented off the cyclohexene ring in the triclinic and monoclinic structures are probably due to different crystal packing arrangements. In all structures, none of the chains off the cyclohexene ring are completely planar, both in the x-ray and CHARMM structures. The reason for the large differences in the dihedrals is that a solid (x-ray) structure with neighboring molecules is being compared to a gas-phase minimized structure (CHARMM).

### 4.3 Retinoid Crystal Simulations

Crystal simulations were performed for retinal and both crystalline forms of retinoic acid. Five 1 ns simulations were performed for each ligand in the crystal environment, and geometric parameters (bonds, angles, and dihedrals) of the ligands were examined. For simulation details, see Chapter 2.
4.3.1 Retinoic acid (Triclinic Form). Two molecules of retinoic acid were reported for the unit cell of the triclinic form of retinoic acid, CSD id ${ }^{54}$ VITAAC10. ${ }^{51}$ The bonds, angles, and dihedrals for each retinoic acid molecule were measured from the simulations and compared back to the x-ray structure. The RMSDs over time (compared to starting x-ray structure) for both retinoic acid molecules in the crystal for the five simulations are in Table 4.6. The RMSD is under $0.5 \AA$ for all simulations, so the
simulated structure does not deviate much from the x-ray structure. The bond, angle, and dihedral error between the simulated and x-ray structure was also low (see Table 4.6 for average errors over all five 1 ns simulations). Unlike in the protein/retinoid simulations, which will be discussed later, the retinoic acid in the crystal simulation does not show large dihedral angle deviations about bonds of the chain. This may be because the simulation was done at a slightly lower temperature (123 K) than the protein/retinoic acid complex simulation ( 277 K ), so the system here does not have as much energy to rotate as much around the single bonds of the chain. In the crystal, the chain remains relatively planar throughout the simulations, and the chain does not move much relative to the cyclohexene ring. The cyclohexene rings in both ligands retain the C 2 '-endo conformation found in the x-ray structure. Figure 4.13 shows overlays of the simulation average structures (blue) and the x-ray structure (red).

Table 0.6 RMSD of ligands compared to x-ray structure for triclinic form of retinoic acid

|  | RTAC1 | RTAC2 |
| :--- | :--- | :--- |
| Simulation 1 $(1 \mathrm{~ns})$ | $0.33 \pm 0.08 \AA$ | $0.34 \pm 0.09 \AA$ |
| Simulation $2(1 \mathrm{~ns})$ | $0.35 \pm 0.08 \AA$ | $0.44 \pm 0.09 \AA$ |
| Simulation $3(1 \mathrm{~ns})$ | $0.38 \pm 0.12 \AA$ | $0.32 \pm 0.09 \AA$ |
| Simulation $4(1 \mathrm{~ns})$ | $0.32 \pm 0.09 \AA$ | $0.40 \pm 0.13 \AA$ |
| Simulation $5(1 \mathrm{~ns})$ | $0.33 \pm 0.09 \AA$ | $0.30 \pm 0.08 \AA$ |
| Average RMSD $(5 \mathrm{~ns})$ | $0.34 \pm 0.09 \AA$ | $0.36 \pm 0.11 \AA$ |
| Average bond error | $0.03 \pm 0.02 \AA$ | $0.03 \pm 0.03 \AA$ |
| Average angle error | $3.0 \pm 3.1^{\circ}$ | $3.0 \pm 3.0^{\circ}$ |
| Average dihedral error | $6.6 \pm 8.1^{\circ}$ | $6.3 \pm 6.6^{\circ}$ |



Figure 4.13 Examples of overlayed x-ray (red) and simulation average structures (blue) from Simulation 1.
4.3.2 Retinal (Monoclinic). Four molecules of retinal were reported for the unit cell of the monoclinic form of retinal, CSD id ${ }^{54}$ TRETAL01. ${ }^{52}$ The bonds, angles, and dihedrals for each retinal molecule were measured from the simulations and compared back to the x-ray structure. The RMSDs over time (compared to starting x-ray structure) for the four retinal molecules in the crystal for the five simulations are in Table 4.7. The RMSD is under $0.5 \AA$ for all simulations, so the simulated structure does not deviate much from the x-ray structure. The bond, angle, and dihedral error between the simulated and x-ray structure was also low (See Table 4.7 for average errors over the five 1 ns simulations). As for the retinoic acid, the retinal in the crystal simulation does not show large dihedral angle deviations about bonds of the chain. In the crystal, the chain remains relatively planar throughout the simulations, and the chain does not move much relative to the cyclohexene ring (the chain is slightly twisted relative to the ring in the x-ray structure, and the retinal molecules retain this twisted configuration during the simulation). The cyclohexene rings in all four molecules in the crystal retain the C2'endo conformation found in the x-ray structure. Figure 4.14 shows the overlays of simulated structures (blue) with the x-ray structure (red).

Table 0.7 RMSD of ligands compared to x-ray structure for retinal

|  | RTAL1 | RTAL2 | RTAL3 | RTAL4 |
| :--- | :--- | :--- | :--- | :--- |
| Simulation $1(1 \mathrm{~ns})$ | $0.30 \pm 0.06 \AA$ | $0.34 \pm 0.06 \AA$ | $0.38 \pm 0.09 \AA$ | $0.35 \pm 0.06 \AA$ |
| Simulation 2 ( 1 ns$)$ | $0.33 \pm 0.09 \AA$ | $0.37 \pm 0.07 \AA$ | $0.42 \pm 0.10 \AA$ | $0.39 \pm 0.09 \AA$ |
| Simulation $3(1 \mathrm{~ns})$ | $0.31 \pm 0.08 \AA$ | $0.37 \pm 0.08 \AA$ | $0.39 \pm 0.09 \AA$ | $0.41 \pm 0.09 \AA$ |
| Simulation $4(1 \mathrm{~ns})$ | $0.37 \pm 0.11 \AA$ | $0.39 \pm 0.08 \AA$ | $0.41 \pm 0.09 \AA$ | $0.42 \pm 0.10 \AA$ |
| Simulation $5(1 \mathrm{~ns})$ | $0.43 \pm 0.11 \AA$ | $0.41 \pm 0.08 \AA$ | $0.46 \pm 0.10 \AA$ | $0.46 \pm 0.09 \AA$ |
| Average RMSD $(5 \mathrm{~ns})$ | $0.35 \pm 0.09 \AA$ | $0.38 \pm 0.07 \AA$ | $0.41 \pm 0.09 \AA$ | $0.41 \pm 0.09 \AA$ |
| Average bond error | $0.02 \pm 0.01 \AA$ | $0.02 \pm 0.01 \AA$ | $0.02 \pm 0.02 \AA$ | $0.02 \pm 0.02 \AA$ |
| Average angle error | $3.0 \pm 4.5^{\circ}$ | $3.4 \pm 5.0^{\circ}$ | $4.3 \pm 5.9^{\circ}$ | $4.0 \pm 5.3^{\circ}$ |
| Average dihedral error | $6.6 \pm 9.6^{\circ}$ | $5.9 \pm 5.6^{\circ}$ | $7.4 \pm 6.4^{\circ}$ | $5.7 \pm 5.6^{\circ}$ |



Figure 4.14 Examples of overlayed x-ray (red) and simulation average structures (blue) (Simulation 2). Only ligands 1 (left) and 3 (right) of the x-ray structure are shown.
4.3.3 Retinoic acid (Monoclinic Form). Four molecules of retinoic acid were reported for the unit cell of the monoclinic form of retinoic acid, CSD id ${ }^{54}$ VITAAC01. ${ }^{50}$ The bonds, angles, and dihedrals for each retinoic acid molecule were measured from the simulations and compared back to the x-ray structure. The RMSDs over time (compared to starting x-ray structure) for the four retinoic acid molecules in the crystal for the five simulations are in Table 4.8.

Table 0.8 RMSD of ligands compared to x-ray structure for monoclinic form of retinoic acid

|  | RTAC | RTAC2 | RTAC3 | RTAC4 |
| :--- | :--- | :--- | :--- | :--- |
| Simulation 1 (1 ns) | $0.52 \pm 0.10 \AA$ | $0.51 \pm 0.11 \AA$ | $0.52 \pm 0.10 \AA$ | $0.52 \pm 0.11 \AA$ |
| Simulation 2 (1 ns) | $0.56 \pm 0.10 \AA$ | $0.50 \pm 0.10 \AA$ | $0.55 \pm 0.10 \AA$ | $0.55 \pm 0.09 \AA$ |
| Simulation 3 (1 ns) | $0.46 \pm 0.11 \AA$ | $0.49 \pm 0.11 \AA$ | $0.50 \pm 0.11 \AA$ | $0.45 \pm 0.11 \AA$ |
| Simulation 4 (1 ns) | $0.38 \pm 0.11 \AA$ | $0.40 \pm 0.09 \AA$ | $0.51 \pm 0.09 \AA$ | $0.50 \pm 0.10 \AA$ |
| Simulation 5 (1 ns) | $0.48 \pm 0.11 \AA$ | $0.49 \pm 0.09 \AA$ | $0.52 \pm 0.10 \AA$ | $0.51 \pm 0.08 \AA$ |
| Average RMSD (5 ns) | $0.48 \pm 0.11 \AA$ | $0.48 \pm 0.10 \AA$ | $0.52 \pm 0.10 \AA$ | $0.51 \pm 0.10 \AA$ |
| Average bond error | $0.02 \pm 0.02 \AA$ | $0.02 \pm 0.02 \AA$ | $0.03 \pm 0.04 \AA$ | $0.03 \pm 0.04 \AA$ |
| Average angle error | $3.6 \pm 3.3^{\circ}$ | $3.6 \pm 3.3^{\circ}$ | $3.4 \pm 3.3^{\circ}$ | $3.8 \pm 3.3^{\circ}$ |
| Average dihedral error | $21 \pm 33^{\circ}$ | $21 \pm 32^{\circ}$ | $19 \pm 32^{\circ}$ | $16 \pm 24^{\circ}$ |
|  |  |  |  |  |

Unlike the triclinic form of retinoic acid and retinal, the RMSD does go above $0.5 \AA$, but the RMSD does not go above $1 \AA$. The bond and angle error between the simulated and x-ray structure was also low (See Table 4.8 for average errors over the five 1 ns simulations). Unlike the triclinic form of retinoic acid and retinal, the monoclinic form of retinoic acid in the crystal simulation shows some large dihedral angle deviations about bonds of the chain, and the chain moves relative to the x-ray structure. The chain of
retinoic acid in the x-ray structure points upward from the ring, but during the simulation, the chain starts to point downward, and there is rotation about the single bonds of the chain. However, the cyclohexene rings in all four molecules in the crystal retain the C3'endo conformation found in the x-ray structure. The reason for the higher RMSDs and more movement about the single bonds of the chain could be that this set of simulations was performed at a high temperature ( 393 K ). Retinoic acid crystallizes in two forms (monoclinic and triclinic); however, the monoclinic form is metastable, and at $\sim 80^{\circ} \mathrm{C}$ and above, the monoclinic form converts irreversibly to the triclinic form. ${ }^{50,51}$ This is what appears to be happening in the high temperature simulations. The chain starts to point more outward than upward, which is observed in the triclinic form, and rotation occurs about the single bonds, probably because of the higher temperature. Figure 4.15 shows examples of simulation average structures (blue) overlayed with the starting x-ray structure (red).


Figure 4.15 Example of overlayed x-ray (red) and simulation average structures (blue) (Simulation 1). The chain on the cyclohexene ring appears to point more outward than upward.

### 4.4 Retinoid/Protein Complex Simulations

Five simulations of retinoid protein complexes in the crystal environment were performed in order to validate the new force field parameters for the retinoids. In each simulation, amino acid-ligand contacts, ligand-water contacts, amino-acid water contacts,
and any other distances or parameters of interest were examined and compared back to the starting x-ray crystal structure. Each simulation was run at the temperature at which the crystal structure was determined, and distances, angles, and torsions measured from the simulations are for heavy atoms only in order to allow for direct comparison back to the x-ray structure. The structures of the four retinoids, all-trans retinoic acid, all-trans retinol, all-trans retinal, and fenretinide, which were built from the model compounds, are shown in Figure 4.16. All retinoids considered differ only in the functional group on the end of the chain. For simulation details, see Chapter 2.





Figure 4.16 The four retinoids built from the model compounds for use in simulations. Retinoic acid, top left; retinol, top right; retinal, bottom left; fenretinide, bottom right.

### 4.4.1 Retinoic Acid Complexed with Celluar Retinoic Acid Binding Protein II.

Retinoic acid binds non-covalently to cellular retinoic acid binding protein II (CRABP
II). The simulation was started from the crystal structure of Kleywegt, et al., ${ }^{55}$ which is shown in Figure 4.17.


Figure 4.17 Retinoic acid complexed with CRABPII ( $\mathrm{PDB}^{56}$ id $1 \mathrm{CBS}^{55}$ ). Protein show in red; ligand shown in cyan with carboxyl group oxygens in red. Figure rendered using the VMD program. ${ }^{57}$

CRABPs (types I and II) are thought to modulate the amount of retinoic acid that is available to nuclear retinoic acid receptors. ${ }^{55}$ From the x-ray structure, it was observed that CRABPII has a one domain with two orthogonal $\beta$-sheets that form a barrel, which in turn are made up of $10 \beta$-strands. ${ }^{55}$ The middle of the first $\beta$-strand contains a bulge at residue 10 , which results in a directional change that allows this strand to form part of both $\beta$-sheets of the protein. This forms a cavity, which contains the retinoic acid ligand. The carboxyl group of the ligand points into the cavity, while the cyclohexene ring points out and is accessible to solvent. Shown below are the RMSD plots (compared to the starting x-ray crystal structure) for different parts of the complex.


Average RMSDs
Protein/ligand complex $1.18 \pm 0.12 \AA$
Protein only
Ligand only $1.44 \pm 0.25 \AA$ $0.74 \pm 0.01 \AA$

RMSD vs. time


Figure 4.18 RMSD plots of the retinoid acid/CRABPII complex (top left) and different parts of the system individually (bottom).

The RMSD of the whole complex shows initial fluctuation, but appears to stabilize at $\sim 700-800$ ps. Upon examination of the individual RMSDs of the protein and ligand, the ligand appears to contribute the most to the RMSD of the complex (average of rmsd of ligand is $\sim 0.29 \AA$ higher than the average RMSD of the protein alone). The ligand RMSD is higher than that of the protein, but it follows a similar pattern of initial fluctuation and then stabilizes around 800 ps .

The geometric parameters (bonds, angles, and dihedrals) of the ligand from the simulation were examined, and the error with respect to the x-ray structure was calculated. The error for bonds of the ligand was $0.02 \pm 0.02 \AA$, the error for angles of
the ligand was $5.5 \pm 3.0^{\circ}$, and the error for dihedrals of the ligand was $13 \pm 22^{\circ}$. The source of this large error is that the chain of the ligand turns with respect to the ring in the binding site of the protein. Also, there are three dihedrals of the chain that deviate from the orientation found in the x-ray structure that contribute to this large error. The chain rotates with respect to the cyclohexene ring during the simulation, which causes large deviations in the torsion angles that are involved around the connection of the ring to the chain. The other sources of large RMSD differences involve parts of the chain. These differences are seen during the simulation because these are single bonds, and the ligand will probably try to orient itself to optimize contacts with the protein side chains, and this does not appear to cause any major distortion of the surrounding protein residues (see Appendix B (Figure C.5) for figures of retinoic acid and large dihedral differences). The cyclohexene ring retains the $\mathrm{C} 2^{\prime}$-endo conformation during the simulation. Because of the large deviations of these ligand dihedrals, different parts of the ligand were examined more closely. Figure 4.19 shows the RMSD for different parts of the retinoic acid (denoted as RTAC) ligand.


Figure 4.19 RMSD over time of different parts of the retinoic acid ligand.

The ring (including methyl groups) by itself does not give a high RMSD, and the chain (excluding the carboxyl group on the end) gives an RMSD intermediate between the ring by itself and ring and chain together. However, the ring and the chain together (including methyl groups but not including the carboxyl group on the end) do give higher RMSDs. This appears to account for most of the RMSD of the ligand. Upon closer examination of the average structure from the simulation and several snapshots, it appears that the cyclohexene ring rotates around, while the ligand is in the binding site and that the chain rotates around two of the single bonds in the chain. Figure 4.20 shows some snapshots from the simulation.

From the top pictures in Figure 4.20, cyclohexene ring of the ligand appears to rotate in the binding site because the methyl groups are not aligned when the structures are superimposed. The bottom pictures show the starting structure from the simulation (bottom left) and the ending structure (bottom right). The ring starts out with the two methyl groups of the ring pointed to the right; however, during the simulation, the ring turns so that the two methyl groups are pointed in the opposite direction. Also, it can be seen from the overlays of the ligand structures that the chain appears to reorient itself in several places, probably to optimize contacts with protein sidechains. The rotation of the cyclohexene ring in the binding site and the reorientation of some of the chain dihedrals is the cause for the large RMSD of the ligand and for the large average error in the torsion angles between the simulated structure and the x-ray structure.


Figure 4.20 Overlay of the average structure of the ligand from the simulation (green) and the x-ray structure (red); top pictures. Starting structure (bottom left) and ending structure (bottom right) from the simulation.

Several ligand-protein and protein-protein contacts were mentioned by the authors of the x-ray structure. These distances were measured in the simulation, and they are shown in the table below and some of the plots are shown (for CHARMM atom names and remaining time series plots of these distances, see Appendix B).

Table 0.9 Protein-ligand and protein-protein distances (CRABPII/retinoic acid complex)

| Protein-ligand <br> contact | x-ray (£̊) | Computed (£) |
| :--- | :--- | :--- |
| Arg132(NH)-RTAC(O29) | 2.73 | $2.69 \pm 0.12$ |
| Tyr134(OH)--RTAC(O29) | 2.57 | $2.91 \pm 0.51$ |
| Ala32(CB)--RTAC(C8) | 3.97 | $3.78 \pm 0.30$ |
| Ala36(CB)--RTAC(C12) | 4.49 | $4.04 \pm 0.37$ |
| Leu28(CD1)--RTAC(C4) | 5.14 | $4.87 \pm 0.59$ |
| Ala35(CB)--RTAC(C19) | 4.45 | $3.90 \pm 0.39$ |
| Leu19(CD2)--RTAC(C16) | 8.63 | $4.69 \pm 1.32$ |
| Ile9(CD)--RTAC(C20) | 7.01 | $7.67 \pm 0.63$ |
| Phe15(CZ)--RTAC(C10) | 3.91 | $4.22 \pm 0.40$ |
| Pro39(CB)--RTAC(C20) | 3.80 | $3.90 \pm 0.38$ |
|  |  |  |
| Protein-protein |  |  |
| Arg59(NH)--Gln74(OE1) | 2.97 | $3.00 \pm 0.50$ |
| Glu73(OE2)--Gln97(NE2) | 2.97 | $3.09 \pm 0.29$ |
| Glu73(OE2)--Trp109(NE1) | 2.95 | $2.86 \pm 0.14$ |
| Ser4(O)--Arg136(NE) | 2.97 | $2.83 \pm 0.12$ |
| Asn2(O)--Arg136(NH2) | 2.97 | $3.18 \pm 0.36$ |
| Glu73(OE2)--Ser83(OG) | 2.81 | $2.71 \pm 0.13$ |
| Arg111(NH1)-Trp109(CZ) 3.19 | $3.41 \pm 0.20$ |  |

The carboxyl oxygens make contact with three residues of the protein, which are $\operatorname{Arg} 132$, Tyr134, and Arg111, which are monitored in Figure 4.21. The retinoic acid (RTAC) ligand starts out and stays close to $\operatorname{Arg} 132$, and the x-ray value falls within the average and standard deviation computed from the simulation. The RTAC ligand starts out near Tyr134, moves away during the $100-550 \mathrm{ps}$ time period, and then moves close to the Tyr residue again and stays there for the remainder of the simulation. The x-ray value falls into the overall average and standard deviation computed from the simulation.


Figure 4.21 Distance vs. time plots of $\operatorname{Arg} 132$ and $\operatorname{Tyr} 134$ to retinoic acid.

The contact between Arg111 and one of the carboxylic oxygens of RTAC is mediated by a water molecule that was found in the crystal structure (Figure 4.22). From the top plot in Figure 4.22, the water molecules exchange but they still hydrogen-bond to Arg111. The closest water in the x-ray structure was $3.02 \AA$ away (labeled as $\mathrm{H}_{2} \mathrm{O} 15$ ). However, there are several waters that exchange positions, because at different time periods of the simulation, different water molecules are closest to Arg111. In the second plot in Figure 4.22, it can be seen again that $\mathrm{H}_{2} \mathrm{O} 15$ and $\mathrm{H}_{2} \mathrm{O} 10$ exchange positions. At the beginning of the simulation, $\mathrm{H}_{2} \mathrm{O} 15$ is closest to the RTAC ligand, but as the simulation progresses, this water is replaced by $\mathrm{H}_{2} \mathrm{O} 10$. The third plot shows the distance over time from Arg111 to the RTAC ligand. The distance is greater than that observed in the x-ray structure (average computed distance is $1.76 \AA$ greater than x-ray distance); this
is probably because more water molecules are present in the simulation than in the x-ray structure, and thus more waters can get between the Arg111 and ligand, pushing them farther apart. The bottom plot in Figure 4.22 shows all the different water molecules that are close to the carboxyl oxygens of the ligand. The Arg111 distance is also shown for comparison, and water molecules appear to be closer to the ligand over most of the simulation time rather than Arg111.


RTAC (O29B) to water oxygens


Arg111(NH2) to RTAC (O29B)


RTAC--water/RTAC--Arg111


Figure 4.22 Distance from Arg111 to different water oxygens (top); distance from retinoic acid to different water oxygens (second from top), distance from Arg111 to RTAC (second from bottom), and ligand water distances (bottom).

Several other protein-ligand contacts from the simulation were also tracked. Most of these interactions were nonpolar residues with nonpolar parts of the ligand (most of the ligand is nonpolar, except the tail with the carboxyl group). Two of the contacts are shown in Figure 4.23.

Pro39(CB) to RTAC(C20)


Phe15 (CZ) to RTAC (C10)


Figure 4.23 Two of the nonpolar ligand-protein contacts (retinoic acid/CRABPII).

Pro39 and Phe15 are two of the nonpolar residues that were in contact with the ligand. For both residues, there is some initial fluctuation at the beginning of the simulation, then the movement of them stabilizes and fluctuates around the x-ray value for the remainder of the simulation. The other nonpolar ligand-protein contacts showed similar behavior.

Several protein-protein contacts were examined to make sure that the new parameters for the ligand were not distorting the protein. Two of these plots are shown in Figure 4.24.

Arg111 (NH1) to Trp109 (CZ)


Glu73 (OE2) to Ser83 (OG)


Figure 4.24 Two of the protein-protein distances (retinoic acid/CRABPII).

The Arg111-Trp109 contact is mediated through an "amido/aromatic ring hydrogen bond". ${ }^{55}$ Burley and Petsko explain that side-chain amino groups interact with aromatic side chains, and that the amino groups of amino acids such as Lys, Arg, etc. are preferentially located within $6 \AA$ of the ring centers of Phe, Tyr, and Trp. ${ }^{58}$ The distance from the amino group of Arg111 to one of the carbons of the aromatic ring of $\operatorname{Trp} 109$ was measured (top plot in Figure 4.24). The computed average is $0.02 \AA$ higher than the distance measured from the x-ray structure, which is in good agreement. The other protein-protein contact shown is for Glu73 and Ser83 (bottom plot in Figure 4.24), and this distance fluctuates around the x-ray value for the entire simulation.

Overall, the simulated retinoic acid/CRABPII complex agrees with the x-ray structure. The ligand gave a large RMSD over time and a large torsion angle error, which
can be attributed to the cyclohexene ring rotating in the binding site and the two chain dihedrals rotating. The ring is probably able to rotate freely because not many proteinligand contacts were identified for the ring portion of the ligand; most are for the chain and carboxyl end group. Not many protein residues are contacting that ring strongly and forcing it to stay in any one position. One carbon of the ring was identified as a close contact with the protein (Leu28), but even that residue is $5.14 \AA$ away from the ligand. The authors of the x-ray structure reported higher B-factors for the cyclohexene ring part of the ligand than for the chain part, ${ }^{55}$ which means that there is more uncertainty for the positions of the ring atoms. It is possible the ring is free to rotate; however, a single x-ray structure will capture only one orientation, and through the simulations, the rotation is observed. The chain dihedrals probably rotate in order to optimize contacts with protein sidechains. The turning motion of the cyclohexene ring in the binding site could be important in ligand binding and dynamics or interaction of the complex with other proteins.
4.4.2 Retinol Complexed with Cellular Retinol Binding Protein. Retinol binds non-covalently to cellular retinol binding proteins (CRBPs). The simulation was started from the crystal structure determined by Calderone, et al. ${ }^{59}$ The retinol ligand complexed with zebrafish CRBP is shown in Figure 4.25.


Figure 4.25 Retinol complexed with zebrafish CRBP ( $\mathrm{PDB}^{56}$ id $1 \mathrm{KQW}^{59}$ ). The ligand is in cyan with the OH oxygen in red, and the protein in yellow. Figure was rendered with VMD. ${ }^{57}$

CRBPs are found in a variety of tissues, are retinol-specific, and are suspected to aid in retinol storage, uptake, and metabolism. ${ }^{60,61}$ Four types of mammalian CRBPs have been identified and characterized, and zebrafish CRBP appears to have a similar amino acid sequence and structure to the mammalian types (sequence identity is highest with mammalian CRBPII—73\%). ${ }^{59}$ The CRBPs, intracellular lipid binding proteins (which include the CRABPs), and fatty acid binding proteins have been show to have low sequence identity, but they have a conserved structural fold, ${ }^{62}$ so the CRBPs look similar to the CRABPs. They have the same $\beta$-barrel structure formed from 10 antiparallel $\beta$ strands, which forms the cavity where the ligand binds, and the open end of the $\beta$-barrel is capped by an $\alpha$-helical region. In the x-ray structure of zebrafish CRBP complexed with retinol (as with rat CRBPI and CRBPII), the retinol is almost completely enclosed within the cavity of the protein. The hydroxyl group on the end is toward the inside of the cavity, while the cyclohexene ring is the near the $\alpha$-helical region near the top of the
cavity. ${ }^{63,64}$ Shown below in Figure 4.26 are the RMSD plots for different parts of the CRBP/retinol complex.

The RMSD of the protein/ligand complex stabilizes quickly to $0.6 \AA$ at the beginning of the simulation, and fluctuates around $0.6 \AA$ for the remainder of the simulation. The alpha carbons follow the same trend as the protein alone but at a lower RMSD (alpha carbons alone fluctuate at around $0.4 \AA$. As for the retinoic acid/CRABP complex, the retinol of the retinol/CRBP complex contributes the most to the RMSD of the complex (protein alone is $0.6 \AA$, and the ligand alone is $0.74 \AA$ ). The ligand and protein RMSDs follow different patterns; the protein RMSD stabilizes quickly at the beginning of the simulation, but the ligand RMSD is initially at $\sim 0.8 \AA$ and then falls to $\sim 0.7 \AA$ at 1 ns and stays around $0.7 \AA$ for the rest of the simulation.


RMSD vs. time


Figure 4.26 RMSD plots of the retinol/CRBP complex (top) and different parts of the system individually (bottom).

The geometric parameters (bonds, angles, and dihedrals) of the ligand from the simulation were examined, and the error with respect to the x-ray structure was calculated. The error for bonds was $0.02 \pm 0.01 \AA$, the error for angles was $2.9 \pm 3.0^{\circ}$, and the error for dihedrals was $21 \pm 25^{\circ}$. The origin of this large dihedral error is rotation about the $\mathrm{C} 12-\mathrm{C} 13$ bond of the chain of retinol.


Figure 4.27 Overlay of x-ray (red) and simulation average structure (blue) (top) and CHARMM atom names in retinol (bottom). The arrow indicates the dihedrals where the large deviations from the starting x-ray crystal structure occur.

The two dihedrals that show the largest deviations from the x-ray structure are rotations about the C12-C13 bonds. The deviation can be seen in the overlayed structures; the C11-C12 bond points in opposite directions in the x-ray structure and in the simulation average structure. As with the retinoic acid ligand in the CRAPBII simulation, the ligand probably tries to optimize contacts with protein side chains; thus, there is rotation about single bonds. The retinol ligand has only two dihedrals showing a major deviation from the x-ray structure along the chain, while the retinoic acid has several dihedrals with major deviations from the x-ray structure. However, unlike the retinoic acid/CRAPBII complex, the retinol does not show any rotation of the cyclohexene ring with respect to the chain of the ligand. This large deviation from the x-ray dihedral does not appear to cause any major distortion of the surrounding protein residues. C11 and C13 move closer to surrounding protein residues, while $\mathrm{C} 12, \mathrm{C} 14$, and C 20 remain relatively unaffected (see Appendix B for time series plots from these atoms to surrounding protein residues). The cyclohexene ring retains the $\mathrm{C}^{\prime}$ '-endo conformation during the simulation.

The reference for the x-ray structure listed the protein residues with at least one atom within $4.5 \AA$ of the ligand (some contacts were up to $6 \AA$ away), which was used for the protein-ligand contacts to measure from the simulation. Table 4.10 lists the calculated and x-ray distances. Some plots are shown; the rest can be found in Appendix B.

Table 0.10 Protein-ligand distances (CRBP/retinol complex)

| Protein-ligand |  |  |
| :--- | :--- | :--- |
| contact | x-ray (Å) | Computed (£) |
| Ile25(CD)--RTOL(C18) | 4.00 | $3.83 \pm 0.17$ |
| Thr29(CG2)--RTOL(C4) | 5.10 | $4.72 \pm 0.20$ |
| Gln38(CD)--RTOL(C17) | 3.88 | $4.32 \pm 0.33$ |
| Ile42(CD)--RTOL(C17) | 6.31 | $6.54 \pm 0.23$ |
| Thr53(OG1)--RTOL(C12) | 4.16 | $4.30 \pm 0.14$ |
| Phe57(CD2)--RTOL(C3) | 3.79 | $4.17 \pm 0.20$ |
| Val62(CG2)--RTOL(C14) | 7.36 | $7.78 \pm 0.31$ |
| Trp106(CZ2)--RTOL(C14) | 4.59 | $4.28 \pm 0.27$ |
| Leu117(CD2)--RTOL(C20) | 3.86 | $4.05 \pm 0.09$ |
| Ile119(CD)--RTOL(C19) | 3.95 | $4.35 \pm 0.26$ |
| Arg58(CB)--RTOL(C16) | 4.24 | $3.97 \pm 0.16$ |
| Thr53(CB)--RTOL(C11) | 5.10 | $4.04 \pm 0.16$ |
| Trp106(CZ2)--RTOL(C13) | 4.96 | $4.56 \pm 0.27$ |
| Met20(SD)--RTOL(C18) | 3.90 | $3.81 \pm 0.15$ |
| Ser55(OG)--RTOL(C17) | 4.04 | $3.90 \pm 0.17$ |
| Ala33(CB)--RTOL(C4) | 4.01 | $3.92 \pm 0.14$ |
| Leu36(CD1)--RTOL(CZ) | 4.25 | $4.00 \pm 0.26$ |
| Tyr60(CD2)--RTOL(C11) | 5.00 | $4.46 \pm 0.29$ |
| Thr51(OG1)-RTOL(OR) | 5.15 | $5.86 \pm 0.22$ |
| Gln108(NE2)-RTOL(OR) | 4.51 | $4.13 \pm 0.24$ |
| Gln108(NE2)-RTOL(OR) | 3.00 | $2.67 \pm 0.07$ |
| Val19(CG1)-RTOL(C19) | 6.41 | $7.30 \pm 0.41$ |

There were only two polar contacts mentioned, which are shown below (these would be hydrogen bonds, but since no hydrogens were identified in the x-ray structure, they are measured as heavy atom distances). The retinol ligand is designated as RTOL.


Figure 4.28 Polar protein/ligand contacts (retinol/CRBP).

The Thr51-RTOL distance from the simulation is larger than that measured from the x ray structure ( $\sim 0.71 \AA$ larger). However, there is a residue in closer contact with the hydroxyl group than Thr51, which is Gln108. The nitrogen of the sidechain of Gln108 is not only closer to the ligand in the x-ray structure ( $4.51 \AA$ vs. $5.15 \AA$ ), but the ligand stays near Gln 108 during the simulation. The distance between the ligand and oxygen of the sidechain of Gln108 fluctuates around $\sim 2.7 \AA$ during the whole simulation, which is about $0.3 \AA$ closer than that reported in the x-ray structure. The authors of the crystal structure refined two positions for the hydroxyl group of retinol so that in one position a hydrogen bond forms between the hydroxyl group and the oxygen of the sidechain of Gln 108, and in the other position, the hydrogen bond from the ligand is formed to Thr51. ${ }^{59}$ For the simulation, we chose the hydroxyl position in which the hydrogen bond
is formed with Gln 108 ; this mode of retinol binding was also observed in structures of rat CRBP I and rat CRBP II complexed with retinol. ${ }^{63,64}$ Therefore, even though the retinol ligand is not forming a hydrogen bond with Thr51, it is forming (keeping during the simulation) a hydrogen bond with sidechain oxygen of Gln108, and the amino group of Gln108 presents another possible hydrogen bond donor/acceptor for the hydroxyl group.

All of the other ligand-protein contacts listed by the paper were contacts of protein residues with the chain or ring of the ligand (not the hydroxyl group). Four of them are shown below.


Ser55 (OG) to RTOL (C17)


Figure 4.29 Two of the contacts from the retinol ligand chain to the protein (retinol/CRBP).

The amino acid residue contacts with the ligand chain measured from the simulation are in good agreement with the x-ray structure distances. For both distances shown, the
ligand chain and amino acid residues do not appear to move away from each other during the simulation. Met20 and the retinol ligand move $\sim 0.5 \AA$ closer to each other; Ser55 and retinol stay about the same distance apart during the whole simulation.


Figure 4.30 Two of the contacts of the retinol ligand ring to the protein (retinol/CRBP).

As with the amino acid residue contacts with the ligand chain, the amino acid contacts with the ligand ring measured from the simulation are in good agreement with the x-ray structure. The ring of the retinol ligand does appear to move away from either Ala33 or Leu36 during the simulation, and the distances fluctuate around the x-ray value. The other distances measured from the simulation are shown in Appendix B. Three of the average distances measured from the simulation became shorter over time than the distance measured from the x-ray structure; the Thr29 distance to the retinol ligand
decreased the most, becoming $0.38 \AA$ shorter than that of the x-ray structure. Several distances measured from the simulation increased over time (ligand and amino acid residues drifted apart), and the most increased distance was for Val19 to the ligand (plot is in Appendix B). This distance increased on average $0.89 \AA$ (initially increased to $\sim 8.0$ $\AA$ and then decreased to $\sim 7 \AA$ ); however, initially, this distance was $6.41 \AA$, which was not extremely close to the ligand in the x-ray structure. The distance increase of $0.89 \AA$ does not seem unusual, considering the starting distance and that water is present in the binding cavity, which could have pushed the residues farther apart.

The authors of the x-ray structure identified seven water molecules in the ligand binding cavity of zebrafish CRBP. The structure on the left of Figure 4.31 shows the seven water molecules in the binding cavity in the x-ray structure. The structure on the top right shows the average structure of the complex from the simulation. Six of the original crystallographic waters are still in the cavity; however, one of them appears to have moved out (that is the water molecule farthest to the right in the average structure of the simulation). Upon closer examination of the average structure, this water molecule appears to have moved outside of the binding cavity because it forms a strong hydrogen bond with Asn 13 .


Asn13 (HD22) to water262 (OH2)


Figure 4.31 Water in binding cavity of zebrafish CRBP. The protein is shown in blue, the ligand is shown in green, and the water oxygens are shown in red. The starting structure is shown on the top left, and the average structure from the simulation is on the top right. The plot on the bottom is the distance from the water molecule to Asn13.

When the water molecule moves toward the edge of the cavity, it forms a strong hydrogen bond with Asn13 $(2.00 \pm 0.18 \AA)$, and thus stays in that position for the remainder of the simulation.

Overall, the distances measured from the simulation agree well with the x-ray structure. Two of the distances from protein residues to the ligand decrease over time (by
$0.38 \AA$ at the most), and several of the distances increase. The largest increase was 0.89 $\AA$, and the distance between the amino acid and ligand was over $6 \AA$ in the starting structure, so movement of the protein and water in the binding site could have pushed the ligand and the residue apart. Six of the seven crystallographic waters stayed in the binding site during the simulation, and the one that did move to the edge was stabilized by a hydrogen bonding interaction with Asn13. Unlike the retinoic acid ligand complexed with CRABPII, the cyclohexene ring of the retinol ligand did not appear to rotate in the binding site of CRBP.


Figure 4.32 Starting structures for retinoic acid/CRABPII complex (left) and for retinol/CRBP complex (right).

Upon closer examination of the starting structures for both complexes, the retinol appears to be positioned deeper into the binding cavity of CRBP than retinoic acid is positioned into CRABPII (Figure 4.31). As a result, retinoic acid probably is not constrained as much by the protein residues around it (especially the ring at the top), so the ring has more space to rotate and move around. Because retinol is positioned farther down into
the binding site, it probably does not have as much room to move around due to the protein residues around it.
4.4.3 Fenretinide Complexed with Retinol Binding Protein (RBP). The starting xray crystal structure for the simulation was determined by Zanotti and coworkers. ${ }^{65}$ The structure of the fenretinide/RBP complex is shown in Figure 4.33.


Figure 4.33 Fenretinide/RBP complex $\left(\mathrm{PDB}^{56}\right.$ id $\left.1 \mathrm{FEL}^{65}\right)$. The ligand is shown in cyan with the nitrogen atom in blue and oxygens in red. The protein is in purple. Figure rendered with VMD. ${ }^{57}$

Retinol binding protein carries retinol through the blood and delivers retinol from storage sites to target cells. RBP contains one binding site for retinol, and in the blood, RBP is found bound to transthyretin (TTR); the formation of the complex is believed to prevent filtration of the small RBP through the kidneys. ${ }^{65}$ Like CRBPs and CRABPs, RBP binds the ligand noncovalently and contains a $\beta$-barrel which forms the binding cavity for the ligand. However, in the case of RBP, the cyclohexene ring binds inside the $\beta$-barrel, and the chain is extended out with the hydroxyl group reaching the surface (the ligand binds the opposite way in CRBPs and CRABPs). ${ }^{65}$ Studies have shown that modifications to
ligands in the vicinity of the cyclohexene ring render them unable to bind to RBP, however, retinoids modified in the area of the hydroxyl group will bind well to RBP. ${ }^{66-68}$

With the increasing interest in synthetic retinoids that show pharmacological and antitumor activity, some questions about how they interact with retinoid-binding proteins, other proteins, and nuclear retinoid receptors have been raised. Synthetic retinoids may interact with retinoid-binding proteins involved in retinoid metabolism, and it has been shown that some synthetic retinoids may affect plasma transport of retinol, which could be a result of their interaction with RBP. ${ }^{67}$ Fenretinide is a retinol analog in which the hydroxyl end group is replaced by a hydroxyphenyl amide group. This particular synthetic retinoid has been examined for use as an antitumor chemopreventative agent in humans. ${ }^{65}$ Studies have shown the fenretinide binds well to apo RBP, but the fenretinide/RBP complex does not show any affinity for TTR. ${ }^{69}$

In the x-ray structure of the fenretinide/RBP complex, the fenretinide ligand fits into the $\beta$-barrel cavity where the retinol ligand would fit. The cyclohexene ring and chain take the place of the corresponding parts of retinol (cyclohexene ring is inside the cavity with the chain protruding up), while the hydroxyphenyl amide group comes out toward the solvent, replacing the position of the hydroxyl group of retinol. Shown below are the RMSD plots for the whole complex and different parts of the complex.


Average RMSDs
Protein/ligand complex $1.78 \pm 0.16 \AA$ Ligand only $\quad 1.11 \pm 0.31 \AA$ Protein only $\quad 1.79 \pm 0.16 \AA$ Alpha carbons $\quad 1.34 \pm 0.13 \AA$


Figure 4.34 RMSD plots of the fenretinide/RBP complex (top) and different parts of the system individually (bottom).

The RMSD of the protein-ligand complex continues to increase until $\sim 1.5 \mathrm{~ns}$, after which the RMSD levels off and stays just under $2 \AA$. The RMSD for the alpha carbons of the protein follows the RMSD pattern of the whole protein, only the magnitude is $\sim 0.5 \AA$ lower for the alpha carbons. The ligand RMSD decreases from the beginning of the simulation to $\sim 1.5 \mathrm{~ns}$ into the simulation, then it jumps up to over $1.5 \AA$ and starts to decrease again. To examine what caused the large increase, the RMSDs for different parts of the ligand were plotted, which is shown at the top of Figure 4.35.


Figure 4.35 RMSDs for different parts of the ligand (top) and time series plot for rotation of the phenol ring of fenretinide (bottom).

The chain is the isoprene chain including the methyl groups on it, the cyclohexene ring includes the three methyl groups attached to it, the phenol/amide part is the phenol ring and the amide group, and the phenol/chain is the isoprene chain including methyl groups, the amide group, and the phenol ring. The cyclohexene ring contributes the least to the RMSD, while the phenol/amide and phenol/chain parts seem to contribute the most, so the phenol ring and amide are the parts that are moving around the most in the binding site. The time series of the dihedral for rotation of the phenol ring is shown in Figure 4.35 (bottom).

The time series shows a large increase in the dihedral value at $\sim 1.5 \mathrm{~ns}$, which is where the large increase is seen the RMSD plot for the ligand. From the plot, it appears that the phenol ring is turning in binding site during the simulation, and this will be supported further by distances measured from the ligand to nearby residues and from snapshots taken from the simulation.

The geometric parameters (bonds, angles, and dihedrals) of the ligand from the simulation were examined, and the error with respect to the x-ray structure was calculated. The error for bonds was $0.03 \pm 0.03 \AA$, the error for angles was $4.9 \pm 4.9^{\circ}$, and the error for dihedrals was $15 \pm 12^{\circ}$. The origin of this large dihedral error is the rotation of the phenol ring, which was mentioned above. This dihedral shows an approximate $30^{\circ}$ difference on average from the x -ray structure.

The authors of the x-ray structure mentioned few details about the ligand conformation and position; they explain that the resolution is not high enough and allows only a limited discussion of the bound ligand conformation. ${ }^{65}$ In the x-ray structure, the cyclohexene ring is in a half-chair conformation (C3-endo), and from the picture below, the average structure from the simulation also shows the C3-endo conformation.


Figure 4.36 Overlay of x-ray structure of ligand and simulation average structure of ligand. The x-ray structure is shown in red, and the simulation average structure is shown in green.

From the overlay of the two structures, the rotation/shift of the phenol ring can also be seen. The authors also mention that two water molecules were found near the hydroxyl group of the phenol ring (each water was $\sim 2.8 \AA$ away). The fenretinide ligand is abbreviated as FRET.

## FRET (O28) to various waters



Figure 4.37 Distance from hydroxy group of the phenol ring of the ligand to various water molecules.

Over time, the water molecules near the hydroxy group of the phenol ring appear to be exchanging, but are still close to the phenol ring.

Due to limited resolution, the paper did not make mention of any ligand-protein contacts. Amino acid-ligand distances that were within $\sim 6 \AA$ in the x-ray structure were measured from the simulation. The table below shows the x-ray and simulation averages of these distances (for CHARMM atom names and remaining distance plots, see Appendix B).

Table 0.11 Protein-ligand distances (RBP/fenretinide complex)

| Protein-ligand |  |  |
| :--- | :--- | :--- |
| contact | x-ray (£̊) | Computed (£) |
| Ile41(CG2)--FRET(C17) | 5.37 | $5.57 \pm 0.58$ |
| Ala55(CB)--FRET(C4) | 3.78 | $3.74 \pm 0.21$ |
| Val61(CG1)--FRET(O29) | 5.11 | $4.75 \pm 0.48$ |
| Phe137(CZ)--FRET(C17) | 4.61 | $4.07 \pm 0.33$ |
| Gln117(CD)--FRET(C16) | 4.64 | $4.28 \pm 0.28$ |
| Ala43(CB)--FRET(C2) | 4.55 | $3.81 \pm 0.24$ |
| His104(CE1)--FRET(C16) | 3.65 | $4.06 \pm 0.56$ |
| Phe36(O)--FRET(N21) | 3.00 | $3.58 \pm 1.15$ |
| Tyr90(CE1)--FRET(C19) | 5.46 | $4.57 \pm 0.70$ |
| Trp91(O)--FRET(C18) | 5.69 | $6.40 \pm 0.64$ |
| Gly75(N)--FRET(C18) | 4.27 | $3.90 \pm 0.32$ |
| Ala57(CB)--FRET(C5) | 3.69 | $3.93 \pm 0.30$ |
| Ala71(CB)--FRET(C12) | 5.64 | $6.13 \pm 1.59$ |
| Leu97(CD1)--FRET(C14) | 6.05 | $6.39 \pm 0.54$ |
| Leu97(CD2)--FRET(C14) | 4.23 | $4.88 \pm 0.54$ |
| Gln98(N)--FRET(O29) | 3.40 | $2.99 \pm 0.25$ |
| Leu97(CA)--FRET(O29) | 4.56 | $3.77 \pm 0.40$ |
| Leu64(CG)--FRET(O28) | 5.91 | $4.26 \pm 0.77$ |
| Leu64(CB)--FRET(O28) | 4.40 | $3.53 \pm 0.51$ |
| Leu37(CD2)--FRET(C12) | 4.27 | $4.05 \pm 0.38$ |
| Leu64(CD1)-FRRT(O28) | 6.66 | $4.59 \pm 1.01$ |
| Leu64(CD2)-FRET(O28) | 6.70 | $4.88 \pm 0.62$ |
|  |  |  |

The distances listed in Table 4.11 show good agreement between the x-ray structure and the simulation; the x-ray value is within the averages and standard deviations from the simulation. However, some distances increase (ligand and protein residue drift apart) and some distances decrease (ligand and protein residue drift closer together); the largest deviation from the x-ray structure is $2.07 \AA$. This large change in distance occurs for Leu64 (sidechain carbons) to fenretinide.

## Leu64 (CD1 and CD2) to FRET (O28)



Figure 4.38 Distance over time for sidechain carbons of Leu64 to ligand.

The sidechain carbons of Leu64 and the oxygen of the phenol ring of the ligand move closer together during the simulation and then move farther apart near the end of the simulation. However, the backbone part of Leu64 moves away from the ligand during the $1.5-2.0$ ns period of the simulation, which follows the pattern of the RMSD increase and the phenol ring rotation dihedral change. The backbone nitrogen of Leu64, the carbonyl oxygen of Arg62, and the carbonyl oxygen of Phe96 all show the same pattern of moving away from the ligand during 1.5-2.0 ns.


Figure 4.39 Distances over time for resides that move away from the ligand during the $1.5-2.0 \mathrm{~ns}$ time period.

The four distances shown in Figure 4.39 all exhibit the same pattern as the ligand RMSD and the dihedral for the phenol ring rotation (all show an increase during the $1.5-2.0 \mathrm{~ns}$ time period). Also, the ligand atoms in the plots of Figure 4.39 are all atoms of the phenol ring. An overlay of the ligand from different snapshots of the simulation shows that the fenretinide ligand does indeed move in the binding site.


Figure 4.40 Overlay of snapshots of the fenretinide ligand.

Figure 4.40 shows an overlay of snapshots of the fenretinide ligand at different times during the simulation. The green ( 500 ps ) and pink ( 1170 ps ) structures are from the first part of the simulation before the RMSD and distances increased. The red ( 1600 ps ) and yellow ( 1800 ps ) structures are from the time period where the RMSD and distances increased. The blue ( 3000 ps ) structure is from the last part of the simulation where the RMSD started to decrease again. The RMSD of the different parts of the ligand, the time series of dihedral for rotation of the phenol ring, and the increase in the protein-ligand distances for ligand atoms in or near the phenol ring are all evidence that the phenol ring at the end of the fenretinide ligand turns and shifts in the binding cavity, and that is what causes any large deviations in protein-ligand distances measured from the simulation.

A closer examination of some simulation snapshots reveals a possible explanation for the movement of the ligand. There are water molecules forming a "bridge" between the Leu35 oxygen to the hydroxyl group oxygen of the phenol ring of the ligand. The distance between Leu35 and the ligand decreases during the 1.5-2.0 ns; however, the distance is too large for a hydrogen bond--~4-6 $\AA$ (top plot in Figure 4.41). The water molecules start out in bulk solvent, but they move into the cavity with the ligand and form the "bridge" between Leu35 and the ligand.


Figure 4.41 Distance from Leu35 to ligand (top), distance from Leu35 to water molecule 414 (second plot), and distance from ligand to various water molecules (third plot), and distance from water 414 to different nearby water molecules (bottom).

Figure 4.41 shows the distance over time from the water molecule to Leu35 and the distance from the ligand to various water molecules. The water molecule (numbered 414 in the simulation) is initially far away from the ligand and Leu35, but it moves in during the 1.5-2.0 ns time period and forms a hydrogen bond with the oxygen of Leu35 during most of that time period (second plot from the top in Figure 4.41). Different water molecules move near the ligand during this period to form a hydrogen bond with the hydroxyl group of the ligand and with water molecule number 414 (second from the bottom and bottom plots in Figure 4.41). The hydroxyl group of the ligand forms a hydrogen bond to different waters, which in turn form a hydrogen bond to water 414, which then forms a hydrogen bond to the carbonyl oxygen of Leu35 (some distances appear large because the $\mathrm{O}-\mathrm{O}$ distance is shown in the plots. The $\mathrm{O}-\mathrm{H}$ distances measured were all under $3.2 \AA$ ).

This water "bridge" may not be what is actually causing the ligand to shift and turn in the binding site, but this is probably what is stabilizing the conformation(s) seen during the $1.5-2.0 \mathrm{~ns}$ time period. The ligand appears to go back to its initial position in the binding site (see Figure 4.40) because the position of the phenol ring at 3 ns looks similar to that at 500 and 1170 ps (at 3 ns , the ligand conformation looks more similar to the structure at 500 ps than the conformation at 1600 ps does). The protein does not appear to undergo any major conformation changes (Figure 4.42) from the overlay of the structure at 500 ps (yellow ligand/blue protein) and the structure at 1600 ps (green ligand/red protein).


Figure 4.42 Overlays of the ligand/protein complex at 500 ps (yellow ligand/blue protein; before large RMSD change) and at 1600 ps (green ligand/red protein; during large RMSD change).

The end of the ligand with the phenol group looks like it shifts/turns in the binding cavity, but this does not appear to cause major structural changes to the protein (at least not on the timescale of the simulation). As with retinoic acid and CRABP, there is probably enough space in the binding site so the ligand can move around with little
consequence for the protein. Shown below are the residues around the ligand at 500 ps and at 1600 ps (Figures 4.43 and 4.44).


Figure 4.43 Snapshot from 500 ps (before distance increase). The labeled residues are the ones that showed a large increase in distance from the ligand during the 1.52.0 ns time period.


Figure 4.44 Snapshot from 1600 ps (after distance increase). The labeled residues are the ones that showed a large increase in distance from the ligand during the 1.52.0 ns time period.

Figures 4.43 and 4.44 show the orientation of the ligand with respect to the surrounding residues (those that showed a large distance change). The turn and shift of the ligand is visible from the snapshots. Overall, the simulated fenretinide/RBP complex agrees well with the x-ray structure. Most of the distances measured were in good agreement, the bonds and angle errors are small, and the fenretinide is not floating out of the binding cavity or distorting the protein much even with turning and shifting of the phenol ring. From the pictures in Figure 4.42, it appears that the fenretinide just has some space to move around in the binding cavity without disrupting too much of the protein, especially near the top of the $\beta$-barrel near the phenol ring. Because the turning/shifting of the phenol ring is not affecting the protein much (except for the amino acids right in the vicinity of the ligand), this does not appear to be a problem (this is the cause of the large dihedral error with respect to the x-ray structure). These motions of the ligand could be important in binding and dynamics of the fenretinide or other synthetic retinoid ligands.
4.4.4 Retinal Complexed with Sensory Rhodopsin II (SRII). The starting x-ray crystal structure for the simulation was determined by Royant and coworkers. ${ }^{70}$ The structure of the SRII/retinal complex is shown in Figure 4.45.


Figure 4.45 Sensory rhodopsin II complexed with retinal. The retinal ligand along with Lys205 are shown in cyan, and the protein is in orange.

Sensory rhodopsins I and II belong to the family of archaeal rhodopsins, which contain seven transmembrane helices (A through G) and a retinal chromophore covalently bound to a Lys residue. ${ }^{70-72}$ SRI and SRII are coupled to the membrane-bound tranducer proteins, HtrI and HtrII. Excitation of SRI or SRII by light causes activation of these tranducer proteins, which in turn initiate phosphorylation cascades that result in attractant or repellent phototaxis, respectively (meaning the whole organism moves in response to the stimulus light). ${ }^{72,73}$ Analogous to bacteriorhodopsin (BR), SRI and SRII can work as light-driven proton pumps when they are not coupled to HtrI and HtrII, which indicates that the pump mechanism of BR is conserved in SRs. ${ }^{73}$ Stimulation by light induces tilting of the helices on the cytoplasmic side of the SRs, and the how this structural
change in SRII translates into signal tranduction is unclear. ${ }^{70,72}$ The x-ray structure was determined in order to understand more about the mechanism of action of SRII.

The RMSD of the protein/ligand complex stabilizes quickly and stays at $\sim 1.75 \AA$ for the remainder of the simulation. The RMSD of the alpha carbons follows the pattern of the protein RMSD; the alpha carbon RMSD is $\sim 0.5 \AA$ lower that that of the protein RMSD. The retinal ligand appears to contribute most to the RMSD at $\sim 1.77 \AA$; however, the ligand RMSD is stable and does not show any large increases or decreases.


Figure 4.46 RMSD plots of the protein-ligand complex (top) and RMSDs of protein, ligand and alpha carbons (bottom).

The geometric parameters (bonds, angles, and dihedrals) of the ligand from the simulation were examined, and the error with respect to the x-ray structure was
calculated. The error for bonds was $0.03 \pm 0.02 \AA$, the error for angles was $3.0 \pm 2.3^{\circ}$, and the error for dihedrals was $25 \pm 33^{\circ}$. The origin of this large dihedral error is the movement of the chain relative to the cyclohexene ring, and several rotations about bonds of the chain.



Figure 4.47 Overlay of the x-ray ligand (red) with the simulation average structure (blue) (top). The bottom picture shows the numbering of the carbon atoms of retinal and the dihedrals that contribute most to the dihedral error.

The chain of the ligand in the x-ray structure points upward, while the chain appears to move downwards during the simulation. The rotation about the $\mathrm{C} 6-\mathrm{C} 7$ bond contributes to the large dihedral error (C1-C6-C7-C8 and C5-C6-C7-C8). Three of the chain dihedrals also show large deviations from the x-ray dihedrals, and this can be seen in the overlayed structures in Figure 4.45 (C7-C8-C9-C19 and C11-C12-C13-C20). The rotation about the $\mathrm{C} 10-\mathrm{C} 11$ bond also shows deviation from the x -ray structure. This is similar to the retinoic acid ligand; rotations about the same single bonds of the chain show deviations from the x-ray structure. This movement probably occurs so that the ligand can optimize contacts with protein sidechains. The cyclohexene ring retains the

C2'-endo conformation. There does not appear to be much distortion of the surrounding protein residues, and most of the protein residue contacts stay in place over the time of the simulation. The largest deviation from the x-ray distances is $\sim 0.5 \AA$ (time series plots can be found in Appendix B).

Because the Schiff base linkage to the retinal ligand was parameterized, this distance was examined for the simulation to make sure it stayed in place. Figure 4.48 shows the Schiff base linkage over time, and the simulation shows excellent agreement with experiment (computed distance was $1.28 \pm 0.015 \AA$; x-ray was $1.28 \AA$ )

RTAL(C15) to Lys205(NZ) (Schiff base linkage)


Figure 4.48 Schiff base linkage over time.

The authors described the binding site of SRII and some protein-protein interactions in detail, and these distances and features were all examined in the simulation. Some of these distances are shown in Table 4.12, some of the time series plots are shown, and the remaining time series plots can be found in Appendix B.

Table 0.12 Protein-ligand and protein-protein distances (SRII/retinal complex)

| protein-ligand |  |  |
| :--- | :--- | :--- |
| contact | x-ray ( $\AA$ ) | Computed ( $\AA$ ) |
| Trp76(CE3)--RTAL(C10) | 4.18 | $4.10 \pm 0.19$ |
| Trp171(CZ2)--RTAL(C20) | 3.83 | $3.98 \pm 0.12$ |
| Tyr174(CZ)--RTAL(C12) | 3.60 | $3.85 \pm 0.15$ |
| Ile83(CD)--RTAL(C20) | 3.96 | $3.81 \pm 0.17$ |
| Ala111(CB)--RTAL(C3) | 5.69 | $5.62 \pm 0.32$ |
| Leu126(O)--RTAL(C3) | 5.38 | $4.98 \pm 0.16$ |
| Gly130(O)--RTAL(C4) | 4.24 | $4.01 \pm 0.15$ |
| Val108(CG1)--RTAL(C7) | 3.72 | $4.31 \pm 0.15$ |
| Pro175(CB)--RTAL(C4) | 3.87 | $4.17 \pm 0.17$ |
| Ile43(CG2)--RTAL(C15) | 5.75 | $6.10 \pm 0.21$ |
| Ala47(CB)--RTAL(C15) | 5.91 | $5.94 \pm 0.15$ |
| Asp75(OD2)--RTAL(C15) | 5.30 | $4.60 \pm 0.32$ |
| Thr79(OG1)--RTAL(C15) | 4.00 | $3.83 \pm 0.19$ |
| Met109(CE)--RTAL(C10) | 3.87 | $4.07 \pm 0.22$ |
| Trp171(CB)--RTAL(C19) | 3.72 | $4.15 \pm 0.22$ |
| Asp201(OD2)--RTAL(C14) | 3.44 | $3.36 \pm 0.11$ |
| Thr204(OG1)--RTAL(C15) | 3.74 | $3.99 \pm 0.27$ |
| Arg72(NH1)--Cl- ion | 3.56 | $4.28 \pm 0.18$ |
| Tyr73(OH)--Cl- ion | 3.51 | $3.47 \pm 0.28$ |
| Phe69(O)--Cl- ion | 3.35 | $3.51 \pm 0.12$ |
| Water406(OH2)--Cl- ion | $4.04,2.83$ | $3.28 \pm 0.12$ |
| Water407(OH2)--Cl- ion | 3.68 | $3.13 \pm 0.08$ |
| Lys205(NZ)--Asp75(OD1) | 4.07 | $3.58 \pm 0.16$ |
| Lys205(NZ)--Thr79(OG1) | 3.67 | $4.15 \pm 0.34$ |
| Thr79(OG1)--Asp75(OD1) | 2.81 | $2.81 \pm 0.14$ |
|  |  |  |

The authors identified one $\mathrm{Cl}^{-}$ion in the structure, and they explain that this ion is coordinated to $\operatorname{Arg} 72$, Tyr73, Phe69, and two water molecules. Figure 4.49 shows these distances over time during the simulation.
$\mathrm{Cl}^{-}$ion to Arg 72 (NH1)

$\mathrm{Cl}^{-}$ion to $\operatorname{Tyr} 73(\mathrm{OH})$

$\mathrm{Cl}^{-}$ion to Phe69 (N)

$\mathrm{Cl}^{-}$ion to water oxygens


Figure $4.49 \mathrm{Cl}^{-}$ion distances to nearby residues (retinal/SRII).

The $\mathrm{Cl}^{-}$ion moves away from $\operatorname{Arg} 72$ by $\sim 0.72 \AA$, but the ion stays close to $\operatorname{Tyr} 73$ and Phe69. Two positions were identified for water 406, and the $\mathrm{Cl}^{-}$ion stays in between the two positions. The ion and water 407 move $\sim 0.5 \AA$ closer together.

The protonated Schiff base is hydrogen bonded to a water molecule (numbered 402), which is also hydrogen bonded to two aspartic acid residues (Asp75 and Asp201). From the top plot of Figure 4.50, water 402 moves away from the Schiff base nitrogen (NZ) of Lys205, and an oxygen (OD1) of Asp201 appears to moves closer to the Schiff base nitrogen of Lys205, replacing the hydrogen bond that was formed between water 402 and Lys205. The second plot in Figure 4.50 shows the distance from Asp75 to water 402 and Lys205 over time. Water 402 moves away from the oxygen (OD2) of Asp75, but the Schiff base nitrogen of Lys75 stays close throughout the simulation. The contact between the Schiff base nitrogen and Asp75 (OD1) is part of the salt bridge that is hypothesized to keep SRII in its inactive conformation (the other part of the salt bridge is the Schiff base nitrogen to Asp75 (OD2). ${ }^{73}$ Even though water 402 moves away from Asp75, water 401 stays close during the entire simulation (third plot in Figure 4.50) and forms a hydrogen bond. The bottom plot of Figure 4.50 shows the distance over time from Asp201 to water 402, which stays about the same throughout the simulation.

Lys205 (NZ) to water402 (OH2)


- computed--lys to
water
_-ray--lys to water
_omputed--lys 205 to
asp201
- $x$-ray--lys to asp
Asp75 (OD2) to water402 (OH2)


Asp75 (OD2) to water401 (OH2)


Asp201 (OD2) to water402 (OH2)


Figure 4.50 Distances for Lys205, Asp75, Asp201, water 401, and water 402.

The chain of the retinal ligand is constrained at the Schiff base end by aromatic residues (Trp76, Trp171, and Tyr174).


Figure 4.51 Distances from retinal ligand to aromatic residues.

Figure 4.51 shows the distances from the Schiff base end of the ligand to nearby aromatic residues. All of the measured simulation distances are in good agreement with the x-ray values; the retinal ligand appears to stay in place at the Schiff base end.

The authors made note of several protein-protein and protein-water contacts in and near the binding site of the ligand.


Figure 4.52 Protein-protein contacts near the ligand (retinal/SRII).

Lys205 stays near Asp75 throughout the simulation (top plot in Figure 4.52). This distance is the other part of the salt bridge that keeps the receptor in its inactive conformation. During the simulation, these two residues move about $0.5 \AA$ closer
together. Lys205 moves away from Thr79 ( $\sim 0.5 \AA$ ) about half way through the simulation (second plot in Figure 4.50); however, Asp75 stays close enough to Thr79 to form a hydrogen bond from one its sidechain oxygens to the hydroxyl group of Thr79. Thr79 and Lys205 move away from each other at about the same time that water 402 moves away from Lys205 (see Figure 4.50, top plot); however, Asp201 moves in and forms a hydrogen bond to the Schiff base nitrogen of Lys205 (Figure 4.50). Figure 4.53 shows the protein-water distances mentioned by the authors of the x-ray structure. The waters that are present in the x -ray structure in the active site stay there during the course of the simulation and form hydrogen bonds to Asp201 and $\operatorname{Arg} 72$.


Figure 4.53 Protein water distances near the ligand (retinal/SRII).

The authors of the x-ray structure did not mention any other protein-ligand contacts, but several were measured from the simulation for comparison back to the x-ray structure. Three of them are shown here, and the rest can be found in Appendix B.

Ile83 (CD) to RTAL (C20)


Ala111(CB) to RTAL (C3)


Leu126 (O) to RTAL (C3)


Figure 4.54 Three of the protein-ligand contacts measured from the simulation (retinal/SRII).

The protein-ligand contacts shown in Figure 4.54 show that there is good agreement between the crystal structure and the simulation. The top plot is from Ile83 to the isoprene chain of retinal, and the bottom two plots are from protein residues to the cyclohexene ring of retinal. In Figure 4.54, the protein residues and the ligand fluctuate around the value from the x-ray structure with the exception of Leu126, in which case the
residue and ligand move $\sim 0.4 \AA$ closer together. Figure 4.51 also contains time series plots for distances from the isoprene chain of retinal to nearby protein residues.

The rest of the protein-ligand contacts measured from the simulation can be found in Appendix B. Some of the distances for the protein-ligand contacts deviate from the starting x-ray value over time; the largest deviation (increase in distance) from the x -ray structure for protein-ligand distances observed from the simulation was $\sim 0.5 \AA$ (Val108 to retinal (C7); Appendix B). Solvating the protein in water instead of using a lipid bilayer could be the cause of some of these deviations. Most other distances fluctuated around the x-ray value or did not deviate by more than $0.5 \AA$. The simulated retinal/SRII complex agrees well with the x-ray structure. The Schiff base linkage from Lys205 to retinal stays intact, and the retinal ligand stays inside the protein. Most of the proteinligand distances measured are in good agreement with the x-ray structure; in fact, this simulation had the most simulation distances that simply fluctuated around the x-ray structure values for the whole simulation time (did not show much deviation from x-ray). This is probably due to the fact that the retinal is covalently bound to SRII, rather than noncovalently, as in the other complexes. The average bond and angle error for the retinal is low; however, the average dihedral error is high. This is due to the chain of the ligand moving with respect to the cyclohexene ring and rotation around several of the single bonds of the chain. There does not seem to be any major distortion in the protein residue surrounding the ligand, so the ligand is probably moving simply to optimize contacts with the nearby sidechains.
$(B R)$. The starting x-ray crystal structure for the simulation was determined by Faham and coworkers. ${ }^{74}$


Figure 4.55 Retinal complexed with bacteriorhodopsin. The retinal ligand for each chain is shown in cyan, and the two protein chains of BR are shown in dark blue.

Bacteriorhodopsin ( BR ) is a light-driven proton pump (converts light energy into a proton gradient) present in the cell membrane of Halobacterium salinarum, an organism found in highly saline environments. BR contains seven transmembrane helices (A through G), which are linked by short loops on either side of the membrane. Each BR chain contains one molecule of retinal, which is covalently bound by a Schiff base linkage to Lys216 in helix G. Absorption of a photon by BR causes isomerization of all-trans retinal (protonated Schiff base linkage) to 13-cis retinal. The Schiff base is then depronated (which causes the retinal chromophore to straighten and push against helix F, causing it to tilt), and the proton is transferred to Asp85. Asp96 reprotonates the Schiff base, and a proton is taken in from the inner, cytoplasmic side to reprotonate Asp96 (helices F and G open a narrow channel in which a proton can come through). The proton on Asp85 is transferred through a network of hydrogen bonds and water molecules to the outside. The retinal then relaxes back to the all-trans configuration (ground state), and another
photocycle can then start. The proton pumping mechanism of BR has been studied extensively by a variety of experimental ${ }^{75-84}$ and computational methods, ${ }^{85-94}$ and the proton route through the membrane can now be followed in detail. The great interest in BR has continued for several reasons: BR is the most well-understood ion pump, it serves as a model for G-protein coupled receptors and other proteins which contain transmembrane helices, and it has become the paradigm of a membrane transporter. ${ }^{95}$

The simulation was performed on the asymmetric unit of the crystal structure, which was of a dimer of BR , rather than the native form, which is a trimer, or a monomer, which is commonly the result when the trimer form is solubilized by nonionic detergent. ${ }^{75,78}$ The RMSD and distance plots for protein-ligand or protein-protein distances were calculated for both chains, and only one x-ray distance is plotted because the x-ray distances were same for both chains of BR. Shown below are the RMSD plots for the whole complex, the protein only, the alpha carbons, and the ligands only.

The RMSD of the ligand/protein complex and the protein chains without the ligands both stabilize quickly to $\sim 3 \AA$ (both chains together and individually). The RMSD of the alpha carbons also stabilizes at the beginning of the simulation to just under $3 \AA$. The ligands together show an RMSD of $\sim 1.2 \AA$ until 2 ns into the simulation, where the RMSD then decreases to just above $1 \AA$. Ligand A follows the same pattern; a decrease in RMSD is observed at 2 ns . Ligand B shows an initial increase in RMSD, stabilizes at $\sim 1 \mathrm{~ns}$, and then converges to where Ligand A converges to $(\sim 0.7 \AA)$.


RMSD vs. time (ligands only)


| Average RMSDs | all | chain A | chain B |
| :--- | :---: | :---: | :---: |
| Protein/ligand complex | $3.04 \pm 0.16 \AA$ | $2.85 \pm 0.15 \AA$ | $2.92 \pm 0.16 \AA$ |
| Protein only | $3.05 \pm 0.16 \AA$ | $2.86 \pm 0.15 \AA$ | $2.93 \pm 0.16 \AA$ |
| Alpha carbons only | $2.92 \pm 0.16 \AA$ | $2.62 \pm 0.14 \AA$ | $2.79 \pm 0.15 \AA$ |
| Ligands only | $1.16 \pm 0.10 \AA$ | $0.71 \pm 0.06 \AA$ | $0.62 \pm 0.07 \AA$ |

Figure 4.56 RMSD plots for retinal/BR complex. Ligand/protein complex (top), protein chains and alpha carbons only (middle), and retinal ligands only (bottom).

The geometric parameters (bonds, angles, and dihedrals) of the ligands from the simulation were examined, and the error with respect to the x-ray structure was calculated. The error for bonds was $0.02 \pm 0.02 \AA$ for both ligands, the error for angles was $5 \pm 4^{\circ}$ for the first ligand and $5 \pm 5^{\circ}$ for the second ligand, and the error for dihedrals was $\sim 33 \pm 20^{\circ}$ for both ligands. The origin of these large dihedral errors is the rotation about single bonds of the chain, and movement of the chain relative to the cyclohexene ring.



Figure 4.57 Overlays of x-ray structure (red) and simulation average structure (blue). The bottom picture shows the rotations about the single bonds that deviate from the x-ray structure.

The largest differences in the dihedrals are seen for the attachment of the ring to the chain (rotation about the C6-C7 bond), and for rotation about several bonds of the chain (C8C9, C10-C11, and C12-C13). This is similar to the behavior seen the SRII/retinal and CRABPII/retinoic acid simulations; the deviations from the x-ray structure are rotations about the same bonds. The bonds that show rotation are single bonds, so the ligand
probably rotates about these to optimize contacts with protein sidechains. Both ligands retain the C3'-endo conformation for the cyclohexene ring found in the x-ray structure.

As for the retinal complexed with SRII, the Schiff base linkage was examined to make sure it stayed in place for both ligands. Figure 4.58 shows that the Schiff base linkage stays in place for both BR chains.


Figure 4.58 Schiff base linkages for retinal/BR complex.

Table 4.13 contains many of the protein-ligand distances measured from the simulation. Several time series plots are shown; additional properties measured can be found in Appendix B. The x-ray distances were measured for both chains, and the distance was found to be the same for both; thus, only one value is given for the x-ray distance in the table and on the time series plots.

Table 0.13 Protein-ligand and protein-protein distances (BR/retinal complex)

| contact | x-ray ( $\AA$ ) | computed chain $\mathbf{A}(\mathbf{\AA})$ | computed chain B ( $\AA$ ) |
| :---: | :---: | :---: | :---: |
| Ser141(OG)—RTAL(C2) | 3.61 | $3.54 \pm 0.12$ | $3.65 \pm 0.16$ |
| Tyr185(CD)-RTAL(C10) | 3.95 | $3.42 \pm 0.12$ | $3.78 \pm 0.19$ |
| Trp86(CD1)-RTAL(C14) | 3.41 | $3.42 \pm 0.10$ | $3.49 \pm 0.09$ |
| Met145(SD)-RTAL(C18) | 3.89 | $3.62 \pm 0.12$ | $3.68 \pm 0.15$ |
| Pro186(CB)-RTAL(C3) | 4.97 | $4.00 \pm 0.13$ | $3.80 \pm 0.14$ |
| Tyr185(CE1)-RTAL(C11) | 3.86 | $3.73 \pm 0.13$ | $4.19 \pm 0.22$ |
| Thr90(CG2)-RTAL(C12) | 3.77 | $3.81 \pm 0.16$ | $4.36 \pm 0.19$ |
| Asp85(CG)-RTAL(C15) | 5.36 | $4.33 \pm 0.20$ | $4.42 \pm 0.13$ |
| Met20(CE)-RTAL(C15) | 5.79 | $6.64 \pm 0.15$ | $6.87 \pm 0.42$ |
| Thr89(CB)-RTAL(C15) | 4.32 | $4.31 \pm 0.13$ | $4.72 \pm 0.16$ |
| Leu93(CD1)-RTAL(C20) | 3.55 | $4.20 \pm 0.37$ | $3.63 \pm 0.14$ |
| Lys41(NZ)-Asp38(OD1) | 4.14 | $4.58 \pm 0.10$ | $3.87 \pm 0.12$ |
| Lys41(NZ)—Asp38(OD2) | 3.78 | $2.59 \pm 0.07$ | $2.67 \pm 0.08$ |
| Thr46(OG1)-Asp96(OD2) | 2.88 | $2.62 \pm 0.05$ | $2.80 \pm 0.13$ |

Several protein-protein distances were mentioned by the authors of x-ray structure, and three of those are shown below. Lys41 (NZ) forms an intrahelical salt bridge with Asp38 (top plot in Figure 4.59), and Thr46 forms a hydrogen bond to Asp96 (second plot in Figure 4.59). The hydrogen bond from Thr 46 to Asp96 appears to stay intact over the entire simulation for both chains. The intrahelical salt bridge distance (Lys41 to Asp38) shortens for chain B but gets larger for chain A by $\sim 0.44 \AA$. However, the distance decreases for both chains when measured from Lys41 to the other oxygen of Asp38 (distances decreases from $3.78 \AA$ to $\sim 2.60 \AA$ ), so the intrahelical salt bridge contact stays in place, just to the other oxygen of Asp38. Tyr30 makes another intrahelical salt bridge with Lys30 (bottom of Figure 4.59). The original hydrogen bond is broken in both chains; however, new ones are formed. In chain A, Tyr43 forms a hydrogen bond to Leu224, and Lys30 forms a hydrogen bond to Tyr26. In chain B, Tyr43 forms hydrogen bond to Met32, and Lys30 forms a hydrogen bond to $\operatorname{Arg} 225$.


Thr46 (OG1) to Asp96 (OD2)



Tyr43 (OH) to Lys30 (NZ)


Figure 4.59 Protein-protein contacts measured from BR/retinal simulation.

Many contacts were measured from the retinal ligand to surrounding protein residues; however, only four of them will be shown here and the remaining ones are in Appendix B. The protein-ligand distances shown in Figure 4.60 show good agreement between the simulation and the x-ray structure. The distances measured from chain A and chain B appear to fluctuate around the x-ray value for the whole simulation time; however, the Tyr185 and Met145 distances to the ligand decrease $\sim 0.5 \AA$ from the x-ray value. The distances shown in Appendix B also show good agreement between the computed values and the x-ray structure. If any contacts appeared to be broken (i.e. ligand or protein distances increased), usually another protein residue near the ligand was found that replaced the original contact. This did not seem unusual; the ligand probably has some room to move around in the binding site even though it is covalently bound to the protein.

Ser141 (OG) to RTAL (C2)


Tyr185(CD1) to RTAL (C10)


Trp86 (CD1) to RTAL (C14)


Met145 (SD) to RTAL (C18)


Figure 4.60 Four of the ligand-protein distances measured from the BR/retinal simulation.

Studies on bacteriorhodopsin have found that water molecules near the retinal ligand are extensively involved in both the structure and function of the protein. ${ }^{79,82,83,96}$ Water molecules were identified in the active site of the x-ray structure, and some of the distances of the water molecules to the ligand are shown in Figure 4.61. The top two plots in Figure 4.61 are for the ligand of chain A, and the bottom two plots are for the ligand of chain B. In chain A, a crystallographic water (denoted WATX14) stays close to the ligand for the entire simulation. WATX6 and WATX114 drift away from the ligand but are replaced by other crystallographic waters or waters added for the simulation. In chain B, WATX 23 and the ligand move closer together as the simulation progresses, and a bulk water replaces the replaces the WATX32-retinal contact. Even though other crystallographic or bulk waters replace the original water-retinal contacts, the interactions appear to stay intact over the simulation time period.


Figure 4.61 Water-ligand distances measured from the BR/retinal simulation.

Overall, the $\mathrm{BR} /$ retinal simulation compares well to the x -ray structure. The Schiff base linkage stays in place and if any ligand-protein or ligand-water contacts were broken, usually another protein residue or water was found near the ligand to make a new contact. The RMSD is highest for the $\mathrm{BR} /$ retinal complex, and this could be due to the fact that the protein is membrane-bound but the simulation was performed in water in order to save computational time. The average bond and angle errors are low; however, the average dihedral error is high. This is due to the chain moving relative to the cyclohexene ring and to the rotation about several single bonds of the chain. The surrounding protein residues do not appear to be disturbed much by the movement of the ligands.

### 4.6 Conclusions

Presented are parameters for the CHARMM force field that allow for molecular modeling and simulation studies of systems with alternating single and double bonds, specifically, retinoids. The bonds, angles, and dihedral angles have been parameterized against MP2/6-31G(d) target data, and the error between CHARMM and the quantum mechanical data is in excellent agreement. The CHARMM model compounds and full retinoids obtained from the new parameters were compared against available experimental data (spectroscopic and x-ray crystal data), and the comparison yields good agreement with respect to geometric parameters. The retinoids made from new parameters were also compared against five reported x-ray structures of protein/retinoid complexes. The structural and geometric analysis of these simulations also compares well to the experimental structures.

The CHARMM torsion energy surfaces show excellent agreement with those obtained from MP2/6-31G(d). The largest difference over all model compounds between barrier heights is $\sim 1.15 \mathrm{kcal} / \mathrm{mol}$ (CHARMM is higher than MP2 for model 8). The bond, angle, and dihedral error for between CHARMM and MP2 is low, with bond error being $0.012 \pm 0.009 \AA$, angle error is $1 \pm 1^{\circ}$, and dihedral error is $2 \pm 4^{\circ}$.

Comparison of the new CHARMM parameters with available experimental data shows that CHARMM is in good agreement with experiment. The model compounds were compared to available experimental structures; the average bond error is $0.03 \pm 0.06$ $\AA$, the average angle error is $2 \pm 2^{\circ}$, and the average dihedral error is $3 \pm 3^{\circ}$. Crystal simulations of the new retinoids show that the new parameters yield the proper bond lengths for compounds with alternating single and double bonds. The simulated structures of retinal and retinoic acid do not deviate much from the starting x-ray crystal structures; the cyclohexene ring retains the proper conformation, the chain does move much relative to the cyclohexene ring, and the single bonds of the chain show minimal rotation (chain stays relatively planar). Another important observation from the crystal simulations is that the new CHARMM parameters show the proper crystal transformations at high temperature. Retinoic acid crystallizes in the monoclinic and triclinic forms; however, at high temperature, the monoclinic form transforms irreversibly to the triclinic form. The crystal simulations at high temperature of retinoic acid in the monoclinic form show the beginnings of this transformation. The chain starts to move outward from the cyclohexene ring during the simulation, which is what is seen in the triclinic form (the monoclinic structure has the chain pointing upward from the ring).

In the crystal simulations of the protein/retinoid complexes, some retinoids show rotation about the single bonds of the chain and rotation of the chain relative to the cyclohexene ring. This behavior is observed because the retinoid ligand will try to optimize contacts with nearby protein side chains. These movements of the retinoids do not appear to cause any major conformational changes or disruptions of the protein, and the retinoid ligands stay inside the binding cavities of the proteins during the simulations and retain contacts found in the x-ray structures. Some ligand motions were also discovered that could be important in ligand binding and dynamics or interaction of the protein/ligand complex with other molecules. In the case of the retinoic acid/CRABPII complex, the cyclohexene ring turns in the binding site of the protein. In the fenretinide/RBP complex, the phenol ring turns and shifts in the binding site of RBP.

The new parameters yield good agreement with experiment and now exhibit the proper bond lengths, angles, and torsion angles for compounds with alternating single and double bonds. The new parameters can now be used in simulations of proteins that bind retinoids to better understand the structure and dynamics of those systems and in drug design to make new retinoids that have therapeutic and anticancer potential.

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## Chapter 5

## Convergence Testing in Biomolecular Simulations

This chapter discusses the statistical tests performed to test for convergence in long and short simulations of DNA and proteins. The purpose of this study was to try to assess convergence in biomolecular systems, and this was done by performing a few very long simulations and many very short simulations. A simple statistical test (potential scale reduction test) was then applied to observe which simulation procedure exhibited better convergence among simulations.

### 5.1 Potential Scale Reduction Test

The potential scale reduction test, also termed the variance ratio method, was developed by Gelman and Rubin in $1992{ }^{1}$ and is a popular convergence diagnostic used by statisticians in order to monitor convergence in multiple Markov chain Monto Carlo simulations. ${ }^{2-4}$ The most obvious way to compare the parameters measured from each simulation (and see if they converge) is to look at time series plots of a particular parameter for each simulation. The potential scale reduction test (PSRT) is a quantitative way (does not rely on visual inspection of the simulations) to compare a particular parameter and is based on analysis of variance. Approximate convergence is assumed when the variance between the different simulations for a particular parameter is no
larger than the variance seen in that particular parameter from an individual simulation. ${ }^{5}$ This approach to monitoring convergence is based on whether or not the simulations have "forgotten" their starting points by comparing multiple simulations to see if they are indistinguishable. With these particular sets of simulations, different properties (18 helical parameters for DNA and 16 geometric parameters for crambin) were monitored for convergence. It is assumed that when the parameters measured have converged by the criteria of the PSRT, that they are sampling the same regions of conformational space at that point and the value of that parameter overlap. In other words, they have "forgotten" their starting points and have become indistinguishable because their values overlap.

For each parameter measured $\psi$, we have $\psi_{i j}$, where $j=1, \ldots, n$ and $i=1, \ldots, m$. The parameter $m$ is the number of simulations and $n$ is the number of samples used from each simulation. Two variances are calculated, the between-sequence variance (variance between different simulations), $B$, and the within-sequence variance (variance within an individual simulation), $W$ :

$$
\begin{gather*}
B=\frac{n}{m-1} \sum_{i=1}^{m}\left(\overline{\psi_{i .}}-\bar{\psi}\right)^{2}  \tag{Eqn. 5.1}\\
W=\frac{1}{m} \sum_{i=1}^{m} s_{i}^{2} \tag{Eqn. 5.2}
\end{gather*}
$$

where

$$
\begin{gather*}
\overline{\psi_{i}}=\frac{1}{n} \sum_{j=1}^{n} \psi_{i j}  \tag{Eqn. 5.3}\\
\bar{\psi}=\frac{1}{m} \sum_{i=1}^{m} \bar{\psi}_{i}  \tag{Eqn. 5.4}\\
s_{i}^{2}=\frac{1}{n-1} \sum_{j=1}^{n}\left(\psi_{i j}-\overline{\psi_{i}}\right)^{2} \tag{Eqn. 5.5}
\end{gather*}
$$

$B$ contains a factor of $n$ because it is based on the variance of the within-sequence means, $\overline{\psi_{i}}\left(\right.$ each of these is an average of $n$ values $\left.\psi_{i j}\right)$.

From $B$ and $W$, two estimates of variance of $\psi$ are calculated. The first, $\hat{v}(\psi)$, is an overestimate of the variance under the assumption that the starting points (in our case the seed for the random number generator for the starting velocities) of the simulations are drawn from a wide range of values (this is termed overdispersion).

$$
\begin{equation*}
\hat{v}(\psi)=\frac{n-1}{n} W+\frac{1}{n} B \tag{Eqn. 5.6}
\end{equation*}
$$

For any finite $n, W$ should underestimate the variance of $\psi$ because the individual simulations have not had time to range over all of the possible values and will thus have less variability. As $n$ approaches infinity, both $\hat{v}(\psi)$ and $W$ will approach the actual variance, $\operatorname{var}(\psi)$, but they do so from opposite directions.

The convergence of the simulations is monitored by estimating the factor by which the overestimate of the variance, $\hat{v}(\psi)$, can be reduced. In other words, it is the
ratio between the upper $(\hat{v}(\psi))$ and lower bounds $(W)$ for the standard deviation of $\psi$, which is termed the estimated potential scale reduction.

$$
\begin{equation*}
\hat{R}=\frac{\hat{v}(\psi)}{W} \tag{Eqn. 5.7}
\end{equation*}
$$

The potential scale reduction is designated as $\hat{R}$ instead $R$ because both the numerator and denominator are estimates of the upper and lower bounds of the variance.

As the simulations converge (and as $n$ goes to infinity), the within-sequence variance exceeds the between-sequence variance (there is more variability within an individual simulation than between the different simulations), and the potential scale reduction will approach 1 . The different simulations will now have overlapping values for the parameter being tested. If the potential scale reduction is high, then the value of the parameter being measured is not converging for all simulations, and the values sampled for a particular parameter are not overlapping. In practice, if $\hat{R}$ values are less than or equal 1.2 , then the particular property being measured is said to have converged. ${ }^{5}$

### 5.2 Protein Convergence

5.2.1. Long-Time Simulations. The PSRT was applied to 10 long ( 45 ns each) simulations of crambin in the crystal environment (see Chapter 2 for simulation methods). A total of 16 geometric parameters were examined, which were the radius of gyration, helix-helix distance (measured from the backbone nitrogen of Asn14 to the backbone nitrogen of Ala27), three beta turn distances (listed in Table 5.1; CA denotes the alpha carbon), and a total of 11 hydrogen bonds in the two $\alpha$-helices (helix 1 has seven hydrogen bonds and helix 2 has 4 hydrogen bonds; these distances are listed in Table 5.1). The hydrogen bonds were measured from the amide hydrogen (denoted HN)
to the backbone carbonyl oxygen (denote O). Table 5.1 lists the $\hat{R}$ values for each parameter. All snapshots saved were used for the analysis (total of 4500 snapshots for each simulation).

Table 5.1 PSRT results (long-time crambin simulations)
parameter $\hat{R}$ long
radius of gyration 3.3
helix-helix distance 1.4
beta turn 1 (Arg17 CA to Gly20 CA) 1.6
beta turn 2 (Pro41 CA to Tyr44 CA) 1.7
beta turn 3 (Gly42 CA to Ala45 CA) 2.6
helix 1
hbond 1 (Ser11 HN to Ile7 O) 1.3
hbond 2 (Asn12 HN to Val8 O) 2.1
hbond 3 (Phe13 HN to Ala9 O) 1.2
hbond 4 (Asn14 HN to Arg10 O) 1.4
hbond 5 (Val15 HN to Ser11 O) 1.5
hbond 6 (Cys16 HN to Asn12 O) 1.2
hbond 7 (Arg17 HN to Phe13 O) 1.2
helix 2
hbond 1 (Ala27 HN to Glu23 O) 1.2
hbond 2 (Thr 28 HN to Ala24 O) 1.1
hbond 3 (Tyr29 HN to Ile25 O) 1.2
hbond 4 (Thr30 HN to Cys26 O) 3.7

Only six of the measured parameters converge for all ten simulations, and all of these parameters are hydrogen bond distances. Two of the time series plots for the hydrogen bond distances that did converge are shown in Figure 5.1. All of the simulations overlap and thus sample approximately the same range of values.


Figure 5.1 Examples of time series plots for hydrogen bonds. Both of these parameters converged according to the PSRT. The different colors represent the different simulations.

The parameters that did not converge were also examined with time series plots to see which simulations were not sampling the same values as the others. Figure 5.2 shows the time series plots for two of the parameters that did not converge. The top plot shows the time series for the radius of gyration, and it is very obvious that the simulations do not sample the same values. For example, the red and dark blue simulations sample lower values than the light blue and green simulations. The bottom plot is for the helix-helix distance, which shows more overlap between the different simulations than the radius of gyration; however, the helix-helix distance still does not converge according to the PSRT.

Even though the simulations exhibit overlap for the helix-helix distance, the PSRT yields non-convergence, which shows the importance of using a quantitative and non-visual test. The other time series plots looked similar.


Figure 5.2 Time series plots for radius of gyration (top) and helix-helix distance (bottom). Neither of these parameters converged according to the PSRT. The different colors represent the different simulations.

The time series plots and Table 5.1 illustrate the importance of running multiple simulations. Some simulations sample different ranges of values than other simulations and thus sample different regions of conformational space. If only one simulation had been performed, none of these differences in behavior would have been discovered. The conclusions made from only one simulation would not apply to the rest of the simulations
and thus would not have been representative of the system. One simulation is not going to sample all possible values that a particular parameter can take on, and this would have gone unnoticed if only a single simulation had been run.

Because convergence could not be achieved for the radius of gyration by removing simulations (and because PCA has been used in other studies to assess convergence), a PCA plot was made to observe which simulations did not sample the same regions of conformational space as the others (Figure 5.3, top). None of the simulations appear to be outliers; they all appear to sample in the same regions. The Scree plot (Figure 5.3, bottom) does not show any large contributions to variance. The highest principal component only contributes $\sim 18 \%$ of the variance. The PCA plot clearly shows that the simulations sample in the same areas, but according the PSRT, the simulations do not converge. The ten simulations have averages for the radius of gyration in the range of 9.82 to 9.92 , which is only a tenth of an angstrom difference. PCA does not reflect this small difference, but the convergence test is sensitive enough that even a small difference is magnified.


Figure 5.3 PCA plot for radius of gyration for the long-time simulations (top) and Scree plot corresponding to PCA plot (bottom).
5.2.2. Short-time Simulations. A total of 20 short-time ( 2 ns each) simulations were performed of crambin in the crystal environment (see Chapter 2 for simulation methods). The same geometric parameters were tested in the short simulations, and all saved snapshots were used for analysis (200 snapshots for each simulation). Table 5.2
lists the $\hat{R}$ values for each parameter for the short simulations. As for the long-time simulations, only six parameters converge according to the PSRT, and they are all hydrogen bond distances. The converging hydrogen bond distances are not the same in the long and short simulations, but they have four in common, which are the third and seventh distances in helix 1 and the first and second distances in helix 2 . Two of the time series plots for the short-time simulations are in Figure 5.4. The top plot is for the second hydrogen in helix 2, and as for the long-time simulations, the short-time simulations overlap and sample the same range of values. The bottom plot is for the radius of gyration; some simulations sample higher values (the green and yellow simulations) while others (blue simulations) sample lower values.

Table 5.2 PSRT results (short-time crambin simulations)

| parameter | $\hat{R}$ short |
| :--- | :--- |
| radius of gyration | 2.6 |
| helix-helix distance | 2.1 |
| beta turn 1 (Arg17 CA to Gly20 CA) | 1.8 |
| beta turn 2 (Pro41 CA to Tyr44 CA) | 1.7 |
| beta turn 3 (Gly42 CA to Ala45 CA) | 2.1 |
|  |  |
| helix 1 |  |
| hbond 1 (Ser11 HN to Ile7 O) | 1.1 |
| hbond 2 (Asn12 HN to Val8 O) | 1.4 |
| hbond 3 (Phe13 HN to Ala9 O) | 1.2 |
| hbond 4 (Asn14 HN to Arg10 O) | 1.2 |
| hbond 5 (Val15 HN to Ser11 O) | 1.3 |
| hbond 6 (Cys16 HN to Asn12 O) | 1.3 |
| hbond 7 (Arg17 HN to Phe13 O) | 1.2 |
|  |  |
| helix 2 |  |
| hbond 1 (Ala27 HN to Glu23 O) | 1.1 |
| hbond 2 (Thr 28 HN to Ala24 O) | 1.1 |
| hbond 3 (Tyr29 HN to Ile25 O) | 1.3 |
| hbond 4 (Thr30 HN to Cys26 O) | 3.2 |

As for the long-time simulations, time series plots were examined for the short-time simulations to see which simulations were the outliers. Figure 5.4 shows the time series plots for hydrogen bond 2 of helix 2 (top) and the radius of gyration (bottom). For the hydrogen bond, the values of all simulations overlap and thus the parameter converges.

However, for the radius of gyration, the opposite is observed; the simulations do not have overlapping values, and the variance within individual simulations is greater than that among the different simulations.
hydrogen bond 2 (helix 2)

radius of gyration vs. time


Figure 5.4 Examples of time series plots for short-time simulations. The top plot is for the second hydrogen bond of helix 2 , which converged. The bottom plot is for the radius of gyration, which did not converge. The different colors represent the different simulations.

Table 5.3 shows the averages and standard deviations for the parameters for both the long and short sets of simulations.

Table 5.3 Parameter averages for long and short crambin simulations

| parameter | long average $(\mathbf{\AA})$ | short average $(\mathbf{\AA})$ <br> radius of gyration |
| :--- | :--- | :--- |
| $9.88 \pm 0.14$ | $9.89 \pm 0.19$ |  |
| helix-helix distance | $9.40 \pm 0.51$ | $8.37 \pm 0.84$ |
| beta turn 1 | $5.31 \pm 0.67$ | $5.30 \pm 1.35$ |
| beta turn 2 | $5.52 \pm 0.45$ | $5.51 \pm 0.65$ |
| beta turn 3 | $5.08 \pm 0.83$ | $5.08 \pm 1.06$ |
|  |  |  |
| helix 1 |  |  |
| hydrogen bond 1 | $1.87 \pm 0.27$ | $1.88 \pm 0.37$ |
| hydrogen bond 2 | $1.98 \pm 0.37$ | $1.96 \pm 0.43$ |
| hydrogen bond 3 | $2.25 \pm 0.51$ | $2.26 \pm 0.71$ |
| hydrogen bond 4 | $1.98 \pm 0.33$ | $1.94 \pm 0.40$ |
| hydrogen bond 5 | $2.25 \pm 0.51$ | $2.24 \pm 0.68$ |
| hydrogen bond 6 | $1.93 \pm 0.30$ | $1.94 \pm 0.44$ |
| hydrogen bond 7 | $2.05 \pm 0.33$ | $2.05 \pm 0.48$ |
|  |  |  |
| helix 2 |  |  |
| hydrogen bond 1 | $1.86 \pm 0.26$ | $1.87 \pm 0.36$ |
| hydrogen bond 2 | $1.96 \pm 0.30$ | $1.95 \pm 0.44$ |
| hydrogen bond 3 | $2.01 \pm 0.35$ | $1.97 \pm 0.48$ |
| hydrogen bond 4 | $2.21 \pm 0.86$ | $2.31 \pm 1.39$ |

Most of the average values calculated from the short and long simulations agree with one another or they are at least within the standard deviations of each other. Both sets of simulations appear to reach the same range of values, whether this happens in a long time or a short time. The long and short simulations appear to converge to the same values, even though the individual simulations within the sets do not converge according to the potential scale reduction test. Because they reach the approximately the same average values and show similar convergence behavior, the short simulations are probably the better choice for convergence when time and effort are considered. 5.3 DNA Convergence
5.3.1. Long-Time Simulations. The PSRT was applied to 4 long ( 150 ns each) simulations of DNA in the crystal environment (see Chapter 2 for simulation methods). The DNA fragment was a 12 base pair fragment of DNA (5'-CGTAGACTCAGC-3'; no ligands bound) and the same 18 helical parameters that were examined in DNA/polyamide study were analyzed with the PSRT. All 15,000 snapshots saved from each simulation were used for analysis. Table 5.4 lists the $\hat{R}$ values for each helical parameter for the four long simulations.

Table 5.4 PSRT results (long-time DNA simulations; 12 bases)

|  |  |
| :--- | :---: |
| parameter | $\hat{R}$ |
| shear | 1.2 |
| stretch | 2.9 |
| stagger | 2.6 |
| buckle | 1.2 |
| propeller | 3.1 |
| opening | 2.5 |
| shift | 2.0 |
| slide | 1.4 |
| rise | 2.2 |
| tilt | 2.4 |
| roll | 4.6 |
| twist | 2.6 |
| x-displacement | 3 |
| y-displacement | 1.8 |
| helical rise | 2.6 |
| inclination | 5.1 |
| tip | 2.8 |
| helical twist | 1.1 |

Almost all of the helical parameters with the exception of three (shear, buckle, and helical twist) do not converge according to the PSRT. Upon examination of time series plots of the parameters, it was obvious that the simulations had a wide range of values for some
parameters and that the simulations appear to sample very different ranges of values. Shown below are the time series plots for propeller and stretch.

stretch vs. time


Figure 5.5 Time series plots for propeller and stretch for long DNA simulations. The values plotted are averages over all 12 base pairs of the DNA fragment.

For propeller, the values sampled range from $\sim 5^{\circ}$ to $-25^{\circ}$, and Simulation 1 has higher values of propeller than the other simulations. Simulation 2 samples the widest range of values with a range of $\sim 20^{\circ}$. Simulation 2 also exhibits the widest range of values for $\operatorname{stretch}(0.5$ to $-1.5 \AA)$. Simulation 1 has a large range of value for stretch $(0.5$ to $-0.5 \AA)$, but not as large as Simulation 2. Simulations 3 and 4 appear to sample the same values for stretch and opening, and they do not appear to have as much variation.

Because fraying of the end base pairs is a commonly observed phenomenon in both experimental ${ }^{6-15}$ and computational ${ }^{16-23}$ studies of DNA, this is most likely occurring here. The terminal base pairs are not as constrained as internal bases, which are held in place by the stacking of the bases. ${ }^{24}$ The conformational flexibility possessed by the end base pairs contributes to non-hydrogen bonded base pairs in the Watson-Crick double helix. Fraying of the end bases causes the helical parameter values to sample different values than the bases in the middle of the helix, and when the averages are taken over the 12 base pairs, the large variation between the ends and the middle bases is reflected in the average. In order to see which bases in the four simulations contributed most to fraying, the RMSD of the DNA with different bases was plotted. The RMSD was plotted for all 12 bases, 10 bases (one base on each end removed), 8 bases (two bases on each end removed), 6 bases (three bases on each end removed), and 4 bases (central four bases).

RMSD vs. time (Simulation1)


RMSD vs. time (Simulation4)


Figure 5.6 RMSD plots for 12, 10, 8, 6, and 4 DNA bases.

From the RMSD plots, it can be seen that there are a lot of fluctuations and bumps (the system is most likely moving between various substates) in the RMSD, in particular for Simulations 1 and 2. Upon removal of one, two, and three base pairs, the RMSD over time is still bumpy. When four base pairs are removed from each end of the helix, most of the fluctuation is eliminated, which means that fraying effects penetrate four base pairs in from the ends of the helix (at least for this DNA fragment). From the RMSD plots, only the middle bases are left for analysis. Because most simulation studies on DNA use only small fragments (10-20 base pairs), many authors eliminate one base from each end for analysis or add an extra base pair "cap" to each end, which is then left out of the analysis. However, the RMSD plots obtained from the four 150 ns simulations described here show that fraying affects not only the end bases, but four bases in from the ends. This RMSD analysis was helpful in deciding which bases should be included in the convergence analysis. When the helical parameters were averaged over the 12 base pairs, the effects of fraying obviously affect the convergence analysis. The abnormal helical parameter values obtained from the end bases are reflected in the averages and thus skew the convergence test. Even though the whole fragment shows large changes in RMSD, the middle 4 pairs do not, and this shows that not only are the middle four bases not affected as much by fraying, but because they are not affected as much by the fraying, their dynamics appear to converge more quickly.

To see if the helical parameters of the middle really did exhibit better convergence than the whole fragment, the PSRT was performed on only the middle four base pairs of the DNA fragment ( $5^{\prime}$-GATC- $3^{\prime}$ ). The helical parameter values were averaged over the middle four bases. Table 5.5 gives the $\hat{R}$ values.

Table 5.5 PSRT results (long-time DNA simulations; 4 bases)

| parameter | $\hat{R}$ |
| :--- | :---: |
| shear | 1.1 |
| stretch | 1.0 |
| stagger | 1.0 |
| buckle | 1.1 |
| propeller | 1.0 |
| opening | 1.0 |
| shift | 1.2 |
| slide | 1.1 |
| rise | 1.1 |
| tilt | 1.1 |
| roll | 1.0 |
| twist | 1.0 |
| x-displacement | 1.2 |
| y-displacement | 1.3 |
| helical rise | 1.5 |
| inclination | 1.0 |
| tip | 1.0 |
| helical twist | 1.0 |

When only the middle four bases are used, only two helical parameters do not converge according the PSRT, which are y-displacement and helical rise. The fraying of the end bases do influence the convergence of the simulations, and even in 150 ns , the dynamics of the end bases do allow the helical parameters to converge. A total of 16 out of 18 the helical parameters for the middle four bases converge; the dynamics of the middle four bases obviously converge more quickly than those of the ends. Thus, the middle bases appear to be the better choice for analysis rather than the whole DNA fragment because they are not influenced by the fraying of the ends. The time series plots below are for the middle four base pairs for twist (which converged) and y-displacement (did not converge).

## twist vs. time



Figure 5.7 Time series plots for twist and y-displacement for the middle four bases of the DNA fragment.

Unlike the time series plots shown in Figure 5.5 for propeller and stretch for all 12 bases, twist does not show the large variation in values for the four simulations (other helical parameter plots were similar). All simulations appear to sample in the same small range of $\sim 34^{\circ}$ to $42^{\circ}$ and thus exhibit convergence. However, even for twist, which did converge, the blue simulation shows an increase up until 35 ns into the simulation. This simulation showed similar behavior for other parameters, so even up to 35 ns , the system structural properties have not relaxed. Simulations 1,3, and 4 have about the same values
for y-displacement; however, Simulation 2 has more variation in values, and this is most likely why y-displacement does not converge even for the middle four bases.
5.3.2. Short-Time Simulations. From the long-time DNA simulations, the middle four base pairs of this DNA fragment obviously converged more quickly and are the better choice for helical parameter analysis. However, due to the amount of computational time needed to run 150 ns of DNA dynamics, many short simulations were run to see if they yielded similar results with the PSRT. A total of 20 short ( 2 ns each) simulations were run with same DNA fragment, and the PSRT was applied to all 12 bases and the four middle bases. Table 5.6 gives the results of the PSRT for the short simulations using all 12 bases.

Table 5.6 PSRT results (short-time DNA simulations; 12 bases)

| parameter | $\hat{R}$ |
| :--- | :---: |
| shear | 3.3 |
| stretch | 3.1 |
| stagger | 1.8 |
| buckle | 1.7 |
| propellor | 1.7 |
| opening | 1.6 |
| shift | 2.9 |
| slide | 1.9 |
| rise | 2.0 |
| tilt | 1.4 |
| roll | 2.1 |
| twist | 1.5 |
| x-displacement | 1.5 |
| y-displacement | 1.4 |
| helical rise | 2.3 |
| inclination | 2.0 |
| tip | 1.6 |
| helical twist | 1.4 |

The convergence results for the short simulations show that none of the helical parameters converge when an average over all 12 base pairs is used. The time series plots below are for opening and shift for the 12 base pairs (other time series plots looked similar).


Figure 5.8 Time series plots for opening and shift for all 12 base pairs of the the 20 short DNA simulations.

As for the long simulations, the helical parameter values exhibit large variation (opening) and some simulations simply sample a different range of values than others (shift).

Again, this is due to fraying, and since the majority of the helical parameters did not
converge for the 150 ns simulations for all 12 base pairs, it is not surprising that none of the helical parameters converge in two 2 ns when all bases are considered. Even though more simulations were run for 2 ns than 150 ns , each of the 20 simulations will diverge, and the dynamics of the end bases in each of those will need time plenty of time to relax (from the long simulations, they obviously need $>150 \mathrm{~ns}$ ). From the plot of shift versus time in Figure 5.8, the system needs $\sim 500 \mathrm{ps}$ for the trajectories to exhibit divergence (for shift, the trajectories coincide up to 500 ps , and then start to diverge); therefore, in 2 ns , the systems will not have enough time for the helical parameters to converge for all bases.

As for the long simulations, the convergence of the helical parameters for the middle four bases was examined using the PSRT. Table 5.7 shows the results of the PSRT for the middle four bases of the short DNA simulations.

Table 5.7 PSRT results (short-time DNA simulations; 4 bases)

|  | $\hat{R}$ |
| :--- | :--- |
| parameter | 1.2 |
| shear | 1.2 |
| stretch | 1.1 |
| stagger | 1.2 |
| buckle | 1.3 |
| propellor | 1.4 |
| opening | 1.2 |
| shift | 1.7 |
| slide | 1.2 |
| rise | 1.2 |
| tilt | 1.5 |
| roll | 1.6 |
| twist | 1.3 |
| x-displacement | 1.4 |
| y-displacement | 1.2 |
| helical rise | 1.7 |
| inclination | 1.3 |
| tip | 1.4 |
| helical twist |  |

When only the middle four bases are used, more of the helical parameters converge than when all 12 bases are used, so there is some improvement. The plots below are for shift and opening (compare to Figure 5.8 for all 12 bases). Shift converged when only 4 bases are used; opening did not.

shift vs. time
opening vs. time

Figure 5.9 Time series plots for shift and opening for the middle four base pairs of the the 20 short DNA simulations.

Shift improved when the middle four bases were used; most of the simulations sample the same range of values for the middle four bases, while some trajectories clearly diverged when all 12 bases were used. For opening, some trajectories (most notably, the pink and blue simulations in Figure 5.9) sample a higher range of values than others even for the middle four bases and thus do not converge.

Even though half of the helical parameters converged when the middle four bases were used, nine parameters still did not converge. The long-time simulations exhibited convergence for all parameters with the exception of two for the middle four bases. The short-time simulations obviously do not show similar convergence behavior to the longtime simulations. A simulation of 2 ns is not enough for most of the helical parameters even for the middle four bases to converge. As with the 12 bases for 2 ns , some of the simulations diverge from the others, and it takes $\sim 500 \mathrm{ps}$ or longer for this to occur (for example, in Figure 5.9, it took 800 ps for opening to diverge). It takes $>500 \mathrm{ps}$ for divergence to be observed, and in 2 ns , the simulations do not have enough time to both diverge and then start to sample the same values again, and therefore, converge. The long simulations, in the case of the DNA, are the better choice. However, since half of the helical parameters did converge for the short simulations, the simulations probably do not need to be run for 150 ns nor do 20 of them need to be run. A smaller number of simulations and a short time period (for example, 10 simulations run for 10 ns each) may yield similar convergence results to the long-time simulations for the middle four bases.

Table 5.8 shows the averages and standard deviations for the parameters for both the long and short sets of DNA simulations.

Table 5.8 Parameter averages for long and short DNA simulations (12 and 4 bases)

|  | long simulations <br> $\mathbf{( 1 2 b p})$ | short simulations <br> $\mathbf{( 1 2 b p})$ | long simulations <br> $\mathbf{( 4 b p )}$ | short simulations <br> $\mathbf{( 4 b p )}$ |
| :--- | :--- | :--- | :--- | :--- |
| shear | $-0.04 \pm 0.36$ | $0.05 \pm 1.35$ | $-0.012 \pm 0.73$ | $0.05 \pm 0.88$ |
| stretch | $0.03 \pm 0.45$ | $-0.10 \pm 0.88$ | $-0.11 \pm 0.16$ | $-0.11 \pm 0.43$ |
| stagger | $0.04 \pm 1.13$ | $0.06 \pm 1.04$ | $0.03 \pm 0.40$ | $0.05 \pm 0.93$ |
| buckle | $2.33 \pm 8.83$ | $-0.04 \pm 23.5$ | $5.49 \pm 15.1$ | $4.40 \pm 25.4$ |
| propellor | $-8.75 \pm 11.32$ | $-8.75 \pm 16.6$ | $-20.5 \pm 13.7$ | $-18.7 \pm 18.8$ |
| opening | $2.48 \pm 6.27$ | $-0.31 \pm 18.4$ | $1.89 \pm 5.90$ | $2.43 \pm 18.6$ |
| shift | $-0.01 \pm 0.38$ | $0.28 \pm 1.14$ | $-0.03 \pm 0.88$ | $0.01 \pm 1.07$ |
| slide | $0.04 \pm 0.34$ | $0.37 \pm 1.10$ | $-0.07 \pm 0.48$ | $-0.11 \pm 1.27$ |
| rise | $3.31 \pm 0.31$ | $3.39 \pm 0.60$ | $3.51 \pm 0.32$ | $3.44 \pm 0.51$ |
| tilt | $-0.53 \pm 5.38$ | $-0.21 \pm 8.85$ | $0.37 \pm 4.59$ | $0.39 \pm 8.80$ |
| roll | $0.04 \pm 8.19$ | $0.88 \pm 10.6$ | $0.28 \pm 5.04$ | $1.34 \pm 10.1$ |
| twist | $33.6 \pm 1.91$ | $34.0 \pm 8.79$ | $38.5 \pm 7.56$ | $38.1 \pm 6.91$ |
| x-displacement | $-0.25 \pm 1.39$ | $0.28 \pm 2.68$ | $-0.26 \pm 1.01$ | $-0.59 \pm 2.12$ |
| y-displacement | $-0.16 \pm 1.01$ | $-0.52 \pm 2.67$ | $0.10 \pm 1.28$ | $0.07 \pm 2.17$ |
| helical rise | $3.08 \pm 0.45$ | $3.18 \pm 1.23$ | $3.47 \pm 0.30$ | $3.38 \pm 0.61$ |
| inclination | $0.46 \pm 15.0$ | $1.54 \pm 20.2$ | $0.75 \pm 7.62$ | $2.86 \pm 19.0$ |
| tip | $0.46 \pm 8.02$ | $-0.15 \pm 14.4$ | $-0.67 \pm 7.17$ | $-0.69 \pm 15.1$ |
| helical twist | $36.1 \pm 4.33$ | $36.4 \pm 12.2$ | $39.2 \pm 7.51$ | $38.8 \pm 6.41$ |

As for the protein simulations, most of the average values calculated from the short and long DNA simulations ( 12 and 4 bases) agree with one another or they are at least within the standard deviations of each other. Both sets of simulations appear to reach a similar range of values, whether this happens in a long time or a short time. The long and short simulations appear to sample and thus converge to similar values, even though the individual simulations within the sets do not converge according to the potential scale reduction test. The long and short simulations sample the same range of values for both 12 and 4 base pairs; however, they show different convergence behavior with the potential scale reduction test, especially when the middle four bases are used.

### 5.4 Conclusions

The potential scale reduction test (PSRT) was applied to multiple simulations of two biomolecular systems in order to assess convergence of several geometric properties. The systems examined were a small protein (crambin) and a 12-base pair fragment of DNA, both in the crystal environment. For both systems, a set of a few long simulations and a set of many short simulations were run in order to assess which set would exhibit better convergence with respect to the properties measured.

For the crambin system, 16 geometric parameters were measured for both sets of simulations (see Table 5.1). Ten long (45 ns each) and 20 short ( 2 ns each) simulations were run for crambin, and the PSRT was applied to the parameters measured for both sets of trajectories. For the long simulations, only six parameters measured converged for the ten simulations, all of which were hydrogen bond distances. The same was observed for the 20 short simulations; six parameters converged, and all were hydrogen bond distances. The distances that converged in the long and short simulations were not all the same distances, but they were all hydrogen bond distances. The other parameters measured (beta turn distances, radius of gyration, and helix-helix distance) are "looser" distances, and the hydrogen bond distances are "tighter"- the former are distances that can vary from simulation to simulation more easily than the hydrogen bond distances. The hydrogen bond distances measured were all in the $\alpha$-helices of the protein, and the hydrogen bonds in the helices do not deviate greatly unless the $\alpha$-helices break apart. The $\alpha$-helices stayed intact during the simulations, so most of the hydrogen bonds stayed around the same values for most simulations. A distance such as the helix-helix distance, however, can fluctuate more among simulations because the helices can drift apart during
the simulations (more or less in different simulations) without the protein unfolding, secondary structure elements breaking down, or any major conformational changes occurring. Because the hydrogen bond distances are "tighter", this is most likely why all the parameters that do converge are hydrogen bond distances; it is more likely that most simulations will sample the same values for hydrogen bonds rather than for a distance like the helix-helix distance. Even for the hydrogen bond distances that do not converge, the distances measured for all simulations are a normal value for a hydrogen bond ( $\sim 1.8$ $3.0 \AA$ ). Both the long and short simulations exhibit similar convergence behavior when examined with the PSRT. Six parameters converge for both sets, all of which are hydrogen bond distances.

When the averages over all long simulations and all short simulations are examined (Table 5.3), most of the average values calculated from the short and long simulations agree with one another or they are at least within the standard deviations of each other. Both the long and short sets of simulations appear to reach the same range of values, even though the individual simulations within the sets do not converge according to the potential scale reduction test (even in 45 ns ). Because they reach the approximately the same average values and show similar convergence behavior, the short simulations are probably the better choice for convergence when time and effort are considered.

For the DNA system, 18 helical parameters were examined. Four long (150 ns each) and 20 short ( 2 ns each) simulations were run for the DNA, and the PSRT was applied to the parameters measured for both sets of trajectories. The PSRT was applied to the helical parameter averages for all 12 bases of the DNA fragment and also to the
averages over the middle four bases of the fragment. When the PSRT was applied to the averages over all 12 bases of the long simulations, only three helical parameters converged. The lack of convergence is due to fraying of the end base pairs. Fraying causes the end bases to sample a different (and wider) range of values than the bases in the middle of the helix, and when the averages are taken over the 12 base pairs, the large variation between the ends and the middle bases is reflected in the average. The lack of convergence observed for the helical parameters over 12 bases shows agreement with experimentally observed results for DNA relaxation. Local structural relaxation of three different B-DNA 16-mers was measured, and complete relaxation and convergence was not observed even at 40 ns (the experiment did not go past 40 ns because reliability is lost beyond that). ${ }^{25}$ However, some relaxation times were only tens of picoseconds; the authors suggest that this broad timescale indicates a complex convergence over a large number of conformational substates. ${ }^{25}$ The simple convergence test used here shows that even at 150 ns , convergence of the helical parameters is not achieved for all 12 bases.

When the RMSD was plotted for different parts of the DNA fragment (eliminating 1, 2, 3 , and 4 base pairs from the ends), the fluctuations in the RMSD were removed only when four bases on each end were chopped off. The effects of fraying penetrate four base pairs in from the ends, and this left only the middle four bases for analysis. To see whether the helical parameter averages over the middle four bases exhibited better convergence than the helical parameter averages over all 12 bases, the PSRT was applied to the helical parameters obtained from the middle four bases. When only the middle four bases are used, the convergence was improved (only two helical parameters do not converge according the PSRT). The fraying of the end bases did
influence the convergence of the simulations, and even in 150 ns , the dynamics of the end bases do allow the helical parameters to converge. For a total of 16 out of 18 the helical parameters for the middle four bases converge; the dynamics of the middle four bases obviously converge more quickly than those of the ends. Thus, the middle bases are the better choice for analysis rather than the whole DNA fragment (at least for this particular piece of DNA). The RMSD analysis (chopping off successive bases from each end) should be performed on simulated DNA fragments to decide which bases should be used in further analysis of the simulations. The middle four bases are not affected as much by fraying, the values sampled for the helical parameters of these bases do not show as much variation, and thus, their dynamics converge more quickly for the simulations.

The same analysis was carried out on the short-time DNA simulations. When all 12 bases were used for analysis with the PSRT, none of the helical parameters converged. This is again due to fraying, and for some parameters, the system needs $\sim 500 \mathrm{ps}$ for the trajectories to even diverge. In the long simulations, the helical parameters did not even converge over 150 ns ; therefore, 2 ns is obviously not enough time for the trajectories to diverge and then start to sample the same values in order to converge. When only the middle four bases were used in the analysis of the short simulations, improvement was observed. Nine of the helical parameters converged, but nine still did not converge. The convergence behavior for the short simulations is obviously different than that of the long simulations, in which only two helical parameters did not converge.

Simulations of 2 ns is not enough for most of the helical parameters even for the middle four bases to converge. As with the 12 bases for 2 ns , some of the simulations diverge from the others, and it takes over 500 ps for divergence to be observed, and in 2
ns, the simulations do not have enough time to both diverge and then converge. Because the long and short simulations do not exhibit similar convergence behavior, the long simulations are the better choice in the case of the DNA.

Other computational studies have reported different results for convergence times for DNA helical parameters. From the results of 10 ns DNA (10 base pairs) simulations, Feig and Pettitt report that structural parameters of the DNA appear to relax around 4 ns into the simulation (for the middle eight bases) and suggest that a timescale of several nanoseconds is probably necessary to obtain convergence of structural parameters, although they did not specify the actual amount of time necessary. ${ }^{16}$ From the results of a 60 ns DNA simulation (12 base pairs), Ponomarev, et al. report that helical parameters relax at $\sim 500 \mathrm{ps}$ into the simulation; however, this was for only one base pair step in the middle of the helix and only one simulation was performed. ${ }^{26}$ From our multiple short simulations, 500 ps is not enough time for the helical parameters to relax for even the middle four base pairs. In some of the simulations for some parameters, 500 ps appears to be long enough for relaxation, but this does not occur in all simulations. Our results appear to be in better agreement with those of Feig and Pettitt; the structural parameters need a nanosecond timescale for relaxation (>150 ns for all 12 bases and at least 35 ns for relaxation for the middle four bases). Overall, for each system of interest, there is a balance between the number of simulations and the length of time that they need to be run. Detailed studies for each unique system are required in order to determine what that balance is.

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## Appendix A

This appendix contains tables of simulated helical parameter values and calculated t-values for the DNA/polyamide simulations. Tables A.1-A. 4 contain the helical parameter values and t -values for the DNA/Dervan polyamide simulations (low and high temperature), Tables A.5-A. 7 contain the DNA/Dervan polyamide distances, Tables A. 9 and A. 10 contain the DNA/hydroxpyrrole helical parameters and distances, Tables A. 11 and A. 12 contain the DNA/netropsin helical parameters and distances, Tables A. 13 and A. 14 contain the DNA/distamycin helical parameters and distances, Tables A. 15 and A. 16 contain the helical parameters and distances from the AMBER simulation of the DNA/netropsin complex, and Table A. 17 contains the helical parameter t -values for the low and high temperature simulations combined. The polyamide topology and parameters are at the end of this appendix.

Table A.1. Computed and x-ray averages and standard deviations of helical parameters for initial four low temperature ( 113 K ) simulations

| Simulation 1 <br> helical parameter $(\mathbf{\AA})$ <br> computed avg./sd | x-ray value/sd | t-test value | significant? |  |
| :--- | :--- | :--- | :--- | :--- |
| x-displacement | $0.29 \pm 1.69$ | $0.21 \pm 2.21$ | 0.10 | no |
| y-displacment | $0.017 \pm 1.14$ | $0.05 \pm 0.96$ | 0.10 | no |
| shear | $-0.038 \pm 0.67$ | $0.04 \pm 0.40$ | 0.58 | no |
| stretch | $-0.15 \pm 0.11$ | $-0.15 \pm 0.18$ | 0 | no |
| stagger | $-0.0095 \pm 0.25$ | $0.23 \pm 0.13$ | 5.47 | yes |
| shift | $0.12 \pm 0.87$ | $0.01 \pm 0.95$ | 0.36 | no |
| slide | $0.98 \pm 0.66$ | $0.82 \pm 0.73$ | 0.67 | no |
| rise | $3.33 \pm 0.22$ | $3.31 \pm 0.17$ | 0.42 | no |
| helical rise | $3.42 \pm 0.34$ | $3.29 \pm 0.18$ | 2.28 | yes |
| helical parameter $\left({ }^{\mathbf{o}}\right)$ |  |  |  |  |
| inclination | $12.27 \pm 9.71$ | $9.19 \pm 16.06$ | 0.57 | no |
| tip | $-0.09 \pm 7.12$ | $-0.57 \pm 6.57$ | 0.21 | no |
| buckle | $4.46 \pm 6.38$ | $-1.14 \pm 5.39$ | 3.11 | no |
| propellor | $-11.25 \pm 7.51$ | $-13.41 \pm 9.86$ | 0.66 | no |

## Table A. 1 (cont'd.)

| opening | $0.13 \pm 5.78$ | $-0.86 \pm 6.05$ | 0.49 | no |
| :--- | :--- | :--- | :--- | :--- |
| tilt | $0.10 \pm 4.24$ | $0.35 \pm 4.17$ | 0.18 | no |
| roll | $6.89 \pm 4.99$ | $4.63 \pm 8.70$ | 0.78 | no |
| twist | $34.4 \pm 7.02$ | $34.99 \pm 6.41$ | 0.26 | no |
| helical twist | $35.72 \pm 7.53$ | $36.62 \pm 4.96$ | 0.54 | no |

Simulation 2

| helical parameter $(\mathbf{A})$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x-displacement | $0.076 \pm 2.04$ | $0.21 \pm 2.21$ | 0.18 | no |
| y-displacment | $-0.053 \pm 1.12$ | $0.05 \pm 0.96$ | 0.32 | no |
| shear | $-0.029 \pm 0.46$ | $0.04 \pm 0.40$ | 0.51 | no |
| stretch | $-0.16 \pm 0.12$ | $-0.15 \pm 0.18$ | 0.13 | no |
| stagger | $-0.21 \pm 0.29$ | $0.23 \pm 0.13$ | 9.91 | yes |
| shift | $-0.51 \pm 0.75$ | $0.01 \pm 0.95$ | 0.19 | no |
| slide | $1.02 \pm 0.69$ | $0.82 \pm 0.73$ | 0.81 | no |
| rise | $3.29 \pm 0.35$ | $3.31 \pm 0.17$ | 0.37 | no |
| helical rise | $3.32 \pm 0.48$ | $3.29 \pm 0.18$ | 0.48 | no |
| helical parameter $\left(^{\circ}\right)$ |  |  |  |  |
| inclination | $14.18 \pm 12.87$ | $9.19 \pm 16.06$ | 0.93 | no |
| tip | $0.91 \pm 9.74$ | $-0.57 \pm 6.57$ | 0.67 | no |
| buckle | $0.31 \pm 9.23$ | $-1.14 \pm 5.39$ | 0.80 | no |
| propellor | $-14.52 \pm 4.96$ | $-13.41 \pm 9.86$ | 0.34 | no |
| opening | $-1.76 \pm 4.54$ | $-0.86 \pm 6.05$ | 0.45 | no |
| tilt | $-0.46 \pm 5.24$ | $0.35 \pm 4.17$ | 0.58 | no |
| roll | $7.45 \pm 6.29$ | $4.63 \pm 8.70$ | 0.97 | no |
| twist | $33.65 \pm 8.41$ | $34.99 \pm 6.41$ | 0.62 | no |
| helical twist | $35.55 \pm 7.09$ | $36.62 \pm 4.96$ | 0.64 | no |

Simulation 3
helical parameter $(\AA)$ computed avg./sd $x$-ray value/sd t-test value significant?

| x-displacement | $0.36 \pm 2.31$ | $0.21 \pm 2.21$ | 0.21 | no |
| :--- | :--- | :--- | :--- | :--- |
| y-displacment | $0.26 \pm 1.13$ | $0.05 \pm 0.96$ | 0.65 | no |
| shear | $-0.010 \pm 0.45$ | $0.04 \pm 0.40$ | 0.38 | no |
| stretch | $-0.16 \pm 0.10$ | $-0.15 \pm 0.18$ | 0.23 | no |
| stagger | $-0.02 \pm 0.37$ | $0.23 \pm 0.13$ | 5.55 | yes |
| shift | $-0.11 \pm 0.92$ | $0.01 \pm 0.95$ | 0.38 | no |
| slide | $1.22 \pm 0.95$ | $0.82 \pm 0.73$ | 1.63 | no |
| rise | $3.33 \pm 0.36$ | $3.31 \pm 0.17$ | 0.27 | no |
| helical rise | $3.39 \pm 0.65$ | $3.29 \pm 0.18$ | 1.60 | no |

helical parameter $\left(^{\circ}\right.$ )

| inclination | $13.17 \pm 16.59$ | $9.19 \pm 16.06$ | 0.74 | no |
| :--- | :--- | :--- | :--- | :--- |
| tip | $-0.17 \pm 8.83$ | $-0.57 \pm 6.57$ | 0.18 | no |
| buckle | $0.62 \pm 11.95$ | $-1.14 \pm 5.39$ | 0.96 | no |
| propellor | $-12.86 \pm 5.67$ | $-13.41 \pm 9.86$ | 0.17 | no |
| opening | $-1.27 \pm 3.52$ | $-0.86 \pm 6.05$ | 0.20 | no |
| tilt | $0.25 \pm 5.02$ | $0.35 \pm 4.17$ | 0.07 | no |

## Table A. 1 (cont'd.)

| roll | $6.98 \pm 9.09$ | $4.63 \pm 8.70$ | 0.81 | no |
| :--- | :--- | :--- | :--- | :--- |
| twist | $33.95 \pm 7.80$ | $34.99 \pm 6.41$ | 0.49 | no |
| helical twist | $36.23 \pm 6.33$ | $36.62 \pm 4.96$ | 0.24 | no |


| Simulation $\mathbf{4}$ |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- |
| helical parameter $(\mathbf{\AA})$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| x-displacement | $0.017 \pm 1.96$ | $0.21 \pm 2.21$ | 0.26 | no |
| y-displacment | $0.066 \pm 1.38$ | $0.05 \pm 0.96$ | 0.05 | no |
| shear | $0.081 \pm 0.51$ | $0.04 \pm 0.40$ | 0.32 | no |
| stretch | $-0.15 \pm 0.13$ | $-0.15 \pm 0.18$ | 0.09 | no |
| stagger | $-0.08 \pm 0.27$ | $0.23 \pm 0.13$ | 7.36 | yes |
| shift | $-0.042 \pm 0.93$ | $0.01 \pm 0.95$ | 0.16 | no |
| slide | $0.708 \pm 0.82$ | $0.82 \pm 0.73$ | 0.46 | no |
| rise | $3.35 \pm 0.30$ | $3.31 \pm 0.17$ | 0.68 | no |
| helical rise | $3.35 \pm 0.41$ | $3.29 \pm 0.18$ | 0.92 | no |
| helical parameter $\left({ }^{\circ}\right)$ |  |  |  |  |
| inclination | $10.12 \pm 9.38$ | $9.19 \pm 16.06$ | 0.17 | no |
| tip | $0.74 \pm 8.30$ | $-0.57 \pm 6.57$ | 0.59 | no |
| buckle | $1.39 \pm 7.90$ | $-1.14 \pm 5.39$ | 1.47 | no |
| propellor | $-11.99 \pm 8.18$ | $-13.41 \pm 9.86$ | 0.45 | no |
| opening | $-2.58 \pm 4.86$ | $-0.86 \pm 6.05$ | 0.90 | no |
| tilt | $-0.41 \pm 4.52$ | $0.35 \pm 4.17$ | 0.54 | no |
| roll | $5.36 \pm 4.83$ | $4.63 \pm 8.70$ | 0.25 | no |
| twist | $34.03 \pm 7.33$ | $34.99 \pm 6.41$ | 0.45 | no |
| helical twist | $35.12 \pm 6.67$ | $36.62 \pm 4.96$ | 0.90 | no |

*helical rise and helical twist are parameters which describe the regularity of the DNA helix, while rise and twist describe the geometry of a base pair step from a local perspective. For more details, see Lu, X.J. and Olson, W.K, Nuc. Acids Res. 2003, 31, 5108.

Table A.2. Computed and x-ray averages and standard deviations of helical parameters for initial three high temperature ( 300 K ) simulations

| Simulation 1 <br> helical parameter ( $\AA$ ) | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :---: | :---: | :---: | :---: | :---: |
| x-displacement | $0.0081 \pm 1.92$ | $0.21 \pm 2.21$ | 0.27 | no |
| y -displacment | $0.0041 \pm 1.46$ | $0.05 \pm 0.96$ | 0.14 | no |
| shear | $-0.033 \pm 0.60$ | $0.04 \pm 0.40$ | 0.54 | no |
| stretch | $-0.081 \pm 0.19$ | $-0.15 \pm 0.18$ | 1.15 | no |
| stagger | $-0.198 \pm 0.56$ | $0.23 \pm 0.13$ | 9.15 | yes |
| shift | $-0.0072 \pm 0.98$ | $0.01 \pm 0.95$ | 0.05 | no |
| slide | $0.84 \pm 0.79$ | $0.82 \pm 0.73$ | 0.08 | no |
| rise | $3.28 \pm 0.41$ | $3.31 \pm 0.17$ | 0.46 | no |
| helical rise | $3.29 \pm 0.46$ | $3.29 \pm 0.18$ | 0.02 | yes |
| Helical parameter ( ${ }^{\circ}$ ) inclination | $13.28 \pm 13.4$ | $9.19 \pm 16.06$ | 0.76 | no |

## Table A. 2 (cont'd).

| tip | $0.08 \pm 12.2$ | $-0.57 \pm 6.57$ | 0.22 | no |
| :--- | :--- | :--- | :--- | :--- |
| buckle | $0.77 \pm 17.8$ | $-1.14 \pm 5.39$ | 1.02 | no |
| propellor | $-12.17 \pm 9.36$ | $-13.41 \pm 9.86$ | 0.37 | no |
| opening | $-1.34 \pm 8.39$ | $-0.86 \pm 6.05$ | 0.24 | no |
| tilt | $0.29 \pm 6.72$ | $0.35 \pm 4.17$ | 0.04 | no |
| roll | $7.45 \pm 7.68$ | $4.63 \pm 8.70$ | 0.97 | no |
| twist | $33.8 \pm 8.07$ | $34.99 \pm 6.41$ | 0.54 | no |
| helical twist | $36.1 \pm 7.28$ | $36.62 \pm 4.96$ | 0.32 | no |


| Simulation 2 |  |
| :--- | :--- |
| helical parameter $(\mathbf{\AA})$ |  |
| computed avg./sd |  |
| x-displacement | $-0.571 \pm 2.45$ |
| y-displacment | $0.034 \pm 1.69$ |
| shear | $-0.0029 \pm 0.64$ |
| stretch | $-0.099 \pm 0.24$ |
| stagger | $-0.123 \pm 0.49$ |
| shift | $-0.063 \pm 1.18$ |
| slide | $0.69 \pm 0.89$ |
| rise | $3.33 \pm 0.37$ |
| helical rise | $3.27 \pm 0.53$ |

Helical parameter ( ${ }^{\circ}$ )

| inclination | $16.81 \pm 15.9$ |
| :--- | :--- |
| tip | $0.71 \pm 13.1$ |
| buckle | $1.37 \pm 12.3$ |
| propellor | $-7.85 \pm 10.5$ |
| opening | $1.83 \pm 9.76$ |
| tilt | $-0.29 \pm 7.33$ |
| roll | $9.33 \pm 8.46$ |
| twist | $31.9 \pm 7.46$ |
| helical twist | $35.4 \pm 6.28$ |

Simulation 3

| helical parameter $(\AA \mathbf{\AA})$ | computed avg./sd |
| :--- | :--- |
| x-displacement | $-0.507 \pm 2.04$ |
| y-displacment | $0.158 \pm 1.43$ |
| shear | $0.0064 \pm 0.59$ |
| stretch | $-0.107 \pm 0.19$ |
| stagger | $-0.164 \pm 0.46$ |
| shift | $-0.073 \pm 1.14$ |
| slide | $0.68 \pm 0.82$ |
| rise | $3.34 \pm 0.35$ |
| helical rise | $3.34 \pm 0.49$ |

Helical parameter ( ${ }^{\circ}$ )

| inclination | $15.88 \pm 13.3$ | $9.19 \pm 16.06$ | 1.24 | no |
| :--- | :--- | :--- | :--- | :--- |
| tip | $-0.33 \pm 11.9$ | $-0.57 \pm 6.57$ | 0.11 | no |
| buckle | $2.79 \pm 12.0$ | $-1.14 \pm 5.39$ | 2.25 | no |

## Table A. 2 (cont'd).

| propellor | $-9.87 \pm 9.11$ | $-13.41 \pm 9.86$ | 1.13 | no |
| :--- | :--- | :--- | :--- | :--- |
| opening | $0.148 \pm 6.79$ | $-0.86 \pm 6.05$ | 0.52 | no |
| tilt | $0.259 \pm 7.14$ | $0.35 \pm 4.17$ | 0.06 | no |
| roll | $9.16 \pm 7.42$ | $4.63 \pm 8.70$ | 1.56 | no |
| twist | $33.5 \pm 6.59$ | $34.99 \pm 6.41$ | 0.67 | no |
| helical twist | $36.2 \pm 5.58$ | $36.62 \pm 4.96$ | 0.23 | no |

Table A.3. Computed and x-ray averages and standard deviations of helical parameters for additional three high temperature ( 300 K ) simulations

| Simulation 1 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| helical parameter ( $\mathbf{A}$ ) | computed avg./sd | x-ray value/sd | t-test value | significant? |
| x-displacement | $-0.78 \pm 2.54$ | $0.21 \pm 2.21$ | 1.22 | no |
| y -displacment | $0.14 \pm 1.66$ | $0.05 \pm 0.96$ | 0.17 | no |
| shear | $0.02 \pm 0.70$ | $0.04 \pm 0.40$ | 0.08 | no |
| stretch | $-0.08 \pm 0.25$ | $-0.15 \pm 0.18$ | 0.91 | no |
| stagger | $-0.05 \pm 0.52$ | $0.23 \pm 0.13$ | 1.71 | no |
| shift | $-0.05 \pm 1.19$ | $0.01 \pm 0.95$ | 0.14 | no |
| slide | $0.63 \pm 0.84$ | $0.82 \pm 0.73$ | 0.70 | no |
| rise | $3.33 \pm 0.39$ | $3.31 \pm 0.17$ | 0.15 | no |
| helical rise | $3.24 \pm 0.53$ | $3.29 \pm 0.18$ | 0.27 | no |
| Helical parameter ( ${ }^{\circ}$ ) |  |  |  |  |
| inclination | $17.3 \pm 17.2$ | $9.19 \pm 16.06$ | 1.49 | no |
| tip | $-0.64 \pm 12.3$ | $-0.57 \pm 6.57$ | 0.02 | no |
| buckle | $3.61 \pm 15.8$ | $-1.14 \pm 5.39$ | 0.95 | no |
| propellor | $-8.43 \pm 10.2$ | $-13.41 \pm 9.86$ | 1.54 | no |
| opening | $2.64 \pm 9.85$ | $-0.86 \pm 6.05$ | 1.12 | no |
| tilt | $0.03 \pm 6.97$ | $0.35 \pm 4.17$ | 0.15 | no |
| roll | $9.21 \pm 8.91$ | $4.63 \pm 8.70$ | 1.62 | no |
| twist | $32.3 \pm 8.59$ | $34.99 \pm 6.41$ | 1.00 | no |
| helical twist | $35.5 \pm 7.08$ | $36.62 \pm 4.96$ | 0.51 | no |
| Simulation 2 |  |  |  |  |
| helical parameter ( $\AA$ ) | computed avg./sd | x-ray value/sd | t-test value | significant? |
| x -displacement | $-0.23 \pm 2.29$ | $0.21 \pm 2.21$ | 0.61 | no |
| y -displacment | $-0.11 \pm 1.57$ | $0.05 \pm 0.96$ | 0.33 | no |
| shear | $-0.003 \pm 0.64$ | $0.04 \pm 0.40$ | 0.21 | no |
| stretch | $-0.11 \pm 0.15$ | $-0.15 \pm 0.18$ | 0.80 | no |
| stagger | $-0.006 \pm 0.54$ | $0.23 \pm 0.13$ | 1.37 | no |
| shift | $0.08 \pm 1.13$ | $0.01 \pm 0.95$ | 0.21 | no |
| slide | $0.69 \pm 0.72$ | $0.82 \pm 0.73$ | 0.58 | no |
| rise | $3.32 \pm 0.32$ | $3.31 \pm 0.17$ | 0.12 | no |
| helical rise | $3.31 \pm 0.44$ | $3.29 \pm 0.18$ | 0.12 | no |
| Helical parameter ( ${ }^{\circ}$ ) |  |  |  |  |
| inclination | $11.7 \pm 14.4$ | $9.19 \pm 16.06$ | 0.54 | no |
| tip | $0.006 \pm 12.4$ | $-0.57 \pm 6.57$ | 0.15 | no |

## Table A. 3 (cont'd.)

| buckle | $1.34 \pm 14.2$ | $-1.14 \pm 5.39$ | 0.55 | no |
| :--- | :---: | :--- | :--- | :--- |
| propellor | $-10.4 \pm 9.22$ | $-13.41 \pm 9.86$ | 1.02 | no |
| opening | $0.03 \pm 6.17$ | $-0.86 \pm 6.05$ | 0.45 | no |
| tilt | $0.17 \pm 7.02$ | $0.35 \pm 4.17$ | 0.08 | no |
| roll | $5.98 \pm 7.07$ | $4.63 \pm 8.70$ | 0.60 | no |
| twist | $33.4 \pm 8.60$ | $34.99 \pm 6.41$ | 0.60 | no |
| helical twist | $35.4 \pm 7.39$ | $36.62 \pm 4.96$ | 0.52 | no |


| Simulation 3 |  |
| :--- | :--- |
| helical parameter $(\mathbf{\AA})$ computed avg./sd <br> x-displacement $-0.009 \pm 2.52$ <br> y-displacment $-0.07 \pm 2.08$ <br> shear $-0.007 \pm 0.66$ <br> stretch $-0.04 \pm 0.52$ <br> stagger $-0.01 \pm 0.53$ <br> shift $-0.009 \pm 1.39$ <br> slide $1.11 \pm 1.01$ <br> rise $3.36 \pm 0.38$ <br> helical rise $3.44 \pm 0.59$ |  |


| x-ray value/sd | t-test value | significant? |
| :--- | :--- | :--- |
| $0.21 \pm 2.21$ | 0.28 | no |
| $0.05 \pm 0.96$ | 0.18 | no |
| $0.04 \pm 0.40$ | 0.23 | no |
| $-0.15 \pm 0.18$ | 0.70 | no |
| $0.23 \pm 0.13$ | 1.46 | no |
| $0.01 \pm 0.95$ | 0.04 | no |
| $0.82 \pm 0.73$ | 0.91 | no |
| $3.31 \pm 0.17$ | 0.45 | no |
| $3.29 \pm 0.18$ | 0.82 | no |

Helical parameter ( ${ }^{\circ}$ )

| inclination | $17.6 \pm 17.9$ | $9.19 \pm 16.06$ | 1.48 | no |
| :--- | :---: | :--- | :--- | :--- |
| tip | $1.25 \pm 14.1$ | $-0.57 \pm 6.57$ | 0.41 | no |
| buckle | $-0.62 \pm 12.8$ | $-1.14 \pm 5.39$ | 0.13 | no |
| propellor | $-6.42 \pm 12.0$ | $-13.41 \pm 9.86$ | 1.83 | no |
| opening | $4.09 \pm 14.1$ | $-0.86 \pm 6.05$ | 1.10 | no |
| tilt | $-0.27 \pm 7.64$ | $0.35 \pm 4.17$ | 0.26 | no |
| roll | $9.46 \pm 9.52$ | $4.63 \pm 8.70$ | 1.59 | no |
| twist | $31.5 \pm 8.14$ | $34.99 \pm 6.41$ | 1.36 | no |
| helical twist | $35.1 \pm 6.43$ | $36.62 \pm 4.96$ | 0.73 | no |

Table A.4. Computed and x-ray averages and standard deviations of helical parameters for additional four low temperature ( 113 K ) simulations

## Simulation 1

| helical parameter $(\AA)$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x-displacement | $0.09 \pm 1.68$ | $0.21 \pm 2.21$ | 0.21 | no |
| y-displacment | $0.03 \pm 1.18$ | $0.05 \pm 0.96$ | 0.05 | no |
| shear | $-0.02 \pm 0.56$ | $0.04 \pm 0.40$ | 0.35 | no |
| stretch | $-0.18 \pm 0.11$ | $-0.15 \pm 0.18$ | 0.84 | no |
| stagger | $-0.02 \pm 0.26$ | $0.23 \pm 0.13$ | 3.07 | yes |
| shift | $0.005 \pm 0.90$ | $0.01 \pm 0.95$ | 0.02 | no |
| slide | $0.87 \pm 0.63$ | $0.82 \pm 0.73$ | 0.23 | no |
| rise | $3.36 \pm 0.23$ | $3.31 \pm 0.17$ | 0.63 | no |
| helical rise | $3.42 \pm 0.33$ | $3.29 \pm 0.18$ | 1.27 | no |

## Table A. 4 (cont'd.)

Helical parameter ( ${ }^{\circ}$ )

| inclination | $12.4 \pm 9.9$ | $9.19 \pm 16.06$ | 1.00 | no |
| :--- | :---: | :--- | :--- | :--- |
| tip | $-0.33 \pm 6.72$ | $-0.57 \pm 6.57$ | 0.11 | no |
| buckle | $3.21 \pm 6.37$ | $-1.14 \pm 5.39$ | 2.15 | yes |
| propellor | $-10.9 \pm 9.01$ | $-13.41 \pm 9.86$ | 0.88 | no |
| opening | $-0.95 \pm 4.11$ | $-0.86 \pm 6.05$ | 0.07 | no |
| tilt | $0.24 \pm 3.99$ | $0.35 \pm 4.17$ | 0.09 | no |
| roll | $7.08 \pm 5.41$ | $4.63 \pm 8.70$ | 1.42 | no |
| twist | $34.6 \pm 6.74$ | $34.99 \pm 6.41$ | 0.18 | no |
| helical twist | $35.9 \pm 6.16$ | $36.62 \pm 4.96$ | 0.34 | no |


| Simulation 2 |  |
| :--- | :---: |
| helical parameter $(\AA)$ | computed avg |
| x-displacement | $0.23 \pm 1.88$ |
| y-displacment | $-0.04 \pm 1.12$ |
| shear | $-0.07 \pm 0.71$ |
| stretch | $-0.15 \pm 0.45$ |
| stagger | $-0.002 \pm 0.44$ |
| shift | $0.06 \pm 0.90$ |
| slide | $0.94 \pm 0.72$ |
| rise | $3.33 \pm 0.26$ |
| helical rise | $3.39 \pm 0.34$ |

Helical parameter ( ${ }^{\circ}$ )

| inclination | $11.9 \pm 10.3$ | $9.19 \pm 16.06$ | 0.83 | no |
| :--- | :---: | :--- | :--- | :--- |
| tip | $-0.52 \pm 7.94$ | $-0.57 \pm 6.57$ | 0.02 | no |
| buckle | $6.53 \pm 6.84$ | $-1.14 \pm 5.39$ | 3.54 | yes |
| propellor | $-11.1 \pm 7.98$ | $-13.41 \pm 9.86$ | 0.93 | no |
| opening | $-1.31 \pm 9.72$ | $-0.86 \pm 6.05$ | 0.15 | no |
| tilt | $0.47 \pm 4.49$ | $0.35 \pm 4.17$ | 0.08 | no |
| roll | $6.56 \pm 6.82$ | $4.63 \pm 8.70$ | 0.89 | no |
| twist | $34.4 \pm 9.92$ | $34.99 \pm 6.41$ | 0.17 | no |
| helical twist | $35.7 \pm 9.58$ | $36.62 \pm 4.96$ | 0.30 | no |

Simulation 3

| helical parameter $(\mathbf{\AA})$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :--- | :---: | :--- | :--- | :--- |
| x-displacement | $0.33 \pm 1.95$ | $0.21 \pm 2.21$ | 0.19 | no |
| y-displacment | $0.13 \pm 1.09$ | $0.05 \pm 0.96$ | 0.23 | no |
| shear | $0.05 \pm 0.78$ | $0.04 \pm 0.40$ | 0.06 | no |
| stretch | $-0.14 \pm 0.09$ | $-0.15 \pm 0.18$ | 0.18 | no |
| stagger | $-0.03 \pm 0.25$ | $0.23 \pm 0.13$ | 3.35 | yes |
| shift | $-0.10 \pm 0.88$ | $0.01 \pm 0.95$ | 0.39 | no |
| slide | $1.03 \pm 0.71$ | $0.82 \pm 0.73$ | 0.93 | no |
| rise | $3.36 \pm 0.23$ | $3.31 \pm 0.17$ | 0.66 | no |
| helical rise | $3.44 \pm 0.34$ | $3.29 \pm 0.18$ | 1.36 | no |
| Helical parameter $\left({ }^{\circ}\right)$ |  |  |  |  |
| inclination | $12.4 \pm 11.3$ | $9.19 \pm 16.06$ | 0.88 | no |

## Table A. 4 (cont'd.)

| tip | $0.46 \pm 7.69$ | $-0.57 \pm 6.57$ | 0.42 | no |
| :---: | :---: | :---: | :---: | :---: |
| buckle | $6.67 \pm 6.97$ | $-1.14 \pm 5.39$ | 3.53 | yes |
| propellor | $-11.4 \pm 9.16$ | $-13.41 \pm 9.86$ | 0.70 | no |
| opening | $-0.004 \pm 7.21$ | $-0.86 \pm 6.05$ | 0.37 | no |
| tilt | $-0.16 \pm 4.86$ | $0.35 \pm 4.17$ | 0.33 | no |
| roll | $6.96 \pm 5.65$ | $4.63 \pm 8.70$ | 1.29 | no |
| twist | $34.8 \pm 7.35$ | $34.99 \pm 6.41$ | 0.08 | no |
| helical twist | $36.3 \pm 6.55$ | $36.62 \pm 4.96$ | 0.16 | no |
| Simulation 4 |  |  |  |  |
| helical parameter ( $\AA$ ) | computed avg./sd | x-ray value/sd | t-test value | significant? |
| x-displacement | $0.21 \pm 1.73$ | $0.21 \pm 2.21$ | 0.002 | no |
| $y$-displacment | $-0.03 \pm 1.06$ | $0.05 \pm 0.96$ | 0.25 | no |
| shear | $-0.08 \pm 0.56$ | $0.04 \pm 0.40$ | 0.69 | no |
| stretch | $-0.17 \pm 0.11$ | $-0.15 \pm 0.18$ | 0.46 | no |
| stagger | $-0.03 \pm 0.26$ | $0.23 \pm 0.13$ | 3.06 | yes |
| shift | $0.03 \pm 0.82$ | $0.01 \pm 0.95$ | 0.09 | no |
| slide | $0.89 \pm 0.71$ | $0.82 \pm 0.73$ | 0.35 | no |
| rise | $3.36 \pm 0.22$ | $3.31 \pm 0.17$ | 0.76 | no |
| helical rise | $3.43 \pm 0.35$ | $3.29 \pm 0.18$ | 1.27 | no |
| Helical parameter ( ${ }^{\circ}$ ) |  |  |  |  |
| inclination | $11.7 \pm 9.7$ | $9.19 \pm 16.06$ | 0.80 | no |
| tip | $0.005 \pm 6.91$ | $-0.57 \pm 6.57$ | 0.26 | no |
| buckle | $4.64 \pm 6.85$ | $-1.14 \pm 5.39$ | 2.66 | yes |
| propellor | $-11.1 \pm 7.54$ | $-13.41 \pm 9.86$ | 0.98 | no |
| opening | $-1.23 \pm 4.35$ | $-0.86 \pm 6.05$ | 0.27 | no |
| tilt | $0.03 \pm 3.90$ | $0.35 \pm 4.17$ | 0.25 | no |
| roll | $6.55 \pm 5.08$ | $4.63 \pm 8.70$ | 1.18 | no |
| twist | $34.3 \pm 6.82$ | $34.99 \pm 6.41$ | 0.29 | no |
| helical twist | $35.6 \pm 6.22$ | $36.62 \pm 4.96$ | 0.53 | no |

Table A.5. Polyamide distances (distances in $\AA$; standard deviation for the x-ray structure is the resolution, $2.15 \AA$ ) for three initial high temperature simulations $(300 \mathrm{~K})$

| Simulation 1 <br> distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x1 | $3.58 \pm 0.24$ | 3.08 | 0.74 | no |
| x2 | $3.55 \pm 0.27$ | 3.13 | 0.62 | no |
| x3 | $3.58 \pm 0.34$ | 2.93 | 0.96 | no |
| x4 | $3.95 \pm 0.36$ | 2.95 | 4.64 | yes |
| x5 | $2.94 \pm 0.22$ | 3.13 | 0.28 | no |
| x1a | $2.96 \pm 0.19$ | 3.3 | 0.50 | no |
| x2a | $3.83 \pm 0.36$ | 3.39 | 0.65 | no |
| x3a | $3.44 \pm 0.30$ | 3.12 | 0.47 | no |
| x4a | $3.46 \pm 0.24$ | 2.77 | 1.01 | no |

Table A. 5
(cont'd.)
$\begin{array}{lllll}x 5 a & 3.59 \pm 0.24 & 3.16 & 0.63 & \text { no }\end{array}$
Simulation 2

| distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x1 | $3.30 \pm 0.20$ | 3.08 | 0.32 | no |
| x2 | $3.27 \pm 0.21$ | 3.13 | 0.20 | no |
| x3 | $3.13 \pm 0.21$ | 2.93 | 0.29 | no |
| x4 | $3.12 \pm 0.22$ | 2.95 | 0.25 | no |
| x5 | $2.81 \pm 0.12$ | 3.13 | 0.47 | no |
| x1a | $2.93 \pm 0.20$ | 3.3 | 0.54 | no |
| x2a | $3.14 \pm 0.24$ | 3.39 | 0.36 | no |
| x3a | $3.12 \pm 0.22$ | 3.12 | 0 | no |
| x4a | $3.30 \pm 0.22$ | 2.77 | 0.77 | no |
| x5a | $3.29 \pm 0.20$ | 3.16 | 0.19 | no |

Simulation 3

| distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x1 | $3.22 \pm 0.18$ | 3.08 | 0.20 | no |
| x2 | $3.29 \pm 0.21$ | 3.13 | 0.23 | no |
| x3 | $3.12 \pm 0.21$ | 2.93 | 0.28 | no |
| x4 | $3.12 \pm 0.22$ | 2.95 | 0.25 | no |
| x5 | $2.82 \pm 0.12$ | 3.13 | 0.45 | no |
| x1a | $2.91 \pm 0.23$ | 3.3 | 0.57 | no |
| x2a | $3.22 \pm 0.27$ | 3.39 | 0.25 | no |
| x3a | $3.17 \pm 0.22$ | 3.12 | 0.07 | no |
| x4a | $3.32 \pm 0.22$ | 2.77 | 0.80 | no |
| x5a | $3.30 \pm 0.20$ | 3.16 | 0.20 | no |

Table A.6. Polyamide distances (distances in $\AA$; standard deviation for x-ray structure is $2.15 \AA$ ) for low temperature simulations ( 113 K )

| Simulation 1 <br> distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x1 | $3.15 \pm 0.091$ | 3.08 | 0.10 | no |
| x2 | $3.20 \pm 0.115$ | 3.13 | 0.10 | no |
| x3 | $3.34 \pm 0.15$ | 2.93 | 0.60 | no |
| x4 | $2.99 \pm 0.11$ | 2.95 | 0.47 | no |
| x5 | $3.27 \pm 0.19$ | 3.13 | 0.21 | no |
| x1a | $3.19 \pm 0.11$ | 3.3 | 0.66 | no |
| x2a | $3.12 \pm 0.13$ | 3.39 | 0.40 | no |
| x3a | $3.13 \pm 0.12$ | 3.12 | 0.01 | no |
| x4a | $2.85 \pm 0.085$ | 2.77 | 0.62 | no |
| x5a | $3.34 \pm 0.12$ | 3.16 | 0.26 | no |

## Table A. 6 <br> (cont'd.)

Simulation 2

| distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x1 | $3.47 \pm 0.12$ | 3.08 | 0.57 | no |
| x2 | $3.21 \pm 0.10$ | 3.13 | 0.12 | no |
| x3 | $3.07 \pm 0.11$ | 2.93 | 0.21 | no |
| x4 | $2.74 \pm 0.06$ | 2.95 | 0.21 | no |
| x5 | $3.09 \pm 0.12$ | 3.13 | 0.57 | no |
| x1a | $4.99 \pm 0.19$ | 3.3 | 2.49 | yes |
| x2a | $4.09 \pm 0.18$ | 3.39 | 1.03 | no |
| x3a | $3.22 \pm 0.14$ | 3.12 | 0.15 | no |
| x4a | $3.19 \pm 0.11$ | 2.77 | 0.62 | no |
| x5a | $3.19 \pm 0.10$ | 3.16 | 0.04 | no |
| Simulation 3 | computed average $\mathbf{x - r a y}$ | t-value | significant? |  |
| x1 | $3.20 \pm 0.10$ | 3.08 | 0.17 | no |
| x2 | $3.14 \pm 0.10$ | 3.13 | 0.01 | no |
| x3 | $3.33 \pm 0.14$ | 2.93 | 0.58 | no |
| x4 | $2.78 \pm 0.07$ | 2.95 | 0.13 | no |
| x5 | $3.04 \pm 0.11$ | 3.13 | 0.51 | no |
| x1a | $3.19 \pm 0.11$ | 3.3 | 0.54 | no |
| x2a | $3.63 \pm 0.21$ | 3.39 | 0.35 | no |
| x3a | $3.13 \pm 0.13$ | 3.12 | 0.01 | no |
| x4a | $2.93 \pm 0.11$ | 2.77 | 0.61 | no |
| x5a | $3.27 \pm 0.12$ | 3.16 | 0.16 | no |
| Simulation 4 | computed average | x-ray | t-value | significant? |
| x1 | $3.23 \pm 0.12$ | 3.08 | 0.22 | no |
| x2 | $3.19 \pm 0.11$ | 3.13 | 0.09 | no |
| x3 | $3.24 \pm 0.16$ | 2.93 | 0.46 | no |
| x4 | $2.90 \pm 0.09$ | 2.95 | 0.59 | no |
| x5 | $3.35 \pm 0.19$ | 3.13 | 0.34 | no |
| x1a | $3.16 \pm 0.10$ | 3.3 | 0.47 | no |
| x2a | $3.49 \pm 0.15$ | 3.39 | 0.15 | no |
| x3a | $3.01 \pm 0.10$ | 3.12 | 0.16 | no |
| x4a | $2.98 \pm 0.13$ | 2.77 | 0.57 | no |
| x5a | $3.40 \pm 0.12$ | 3.16 | 0.35 | no |
| xa |  |  |  |  |

Table A.7. T-test results for combined polyamide distances
Low temperature (113 K) simulations

| distance | \# significant tests | $\mathbf{\%}$ | avg. t-value significant? |  |
| :--- | :---: | :---: | :---: | :---: |
| x1 | 16 | 1.6 | 0.99 | no |
| x2 | 7 | 0.7 | 0.72 | no |
| x3 | 425 | 42.5 | 1.92 | no |
| x4 | 5 | 0.5 | 0.42 | no |
| x5 | 3 | 0.3 | 0.37 | no |
| x1a | 2 | 0.2 | 0.19 | no |
| x2a | 0 | 0 | 0.11 | no |
| x3a | 0 | 0 | 0.28 | no |
| x4a | 102 | 10.2 | 1.47 | no |
| x5a | 35 | 3.5 | 1.12 | no |

Table A.8. T-test results for combined polyamide distances

| Low temperature (300 K) simulations |  |  |  |  |
| :--- | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| distance | \# significant tests | \% | avg. t-value | significant? |
| x1 | 49 | 4.9 | 1.10 | no |
| x2 | 26 | 2.6 | 0.87 | no |
| x3 | 22 | 2.2 | 0.94 | no |
| x4 | 59 | 5.9 | 0.81 | no |
| x5 | 0 | 0 | 0.38 | no |
| x1a | 81 | 8.1 | 0.92 | no |
| x2a | 5 | 0.5 | 0.33 | no |
| x3a | 1 | 0.1 | 0.34 | no |
| x4a | 318 | 31.8 | 1.82 | no |
| x5a | 22 | 2.2 | 0.78 | no |

Table A.9. Computed and x-ray averages and standard deviations of helical parameters for hydroxypyrrole polyamide simulation

| helical parameter $(\mathbf{A})$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :--- | :---: | :---: | :--- | :--- |
| x-displacement | $0.59 \pm 1.76$ | $0.37 \pm 1.32$ | 0.41 | no |
| y-displacment | $-0.31 \pm 3.00$ | $-0.23 \pm 1.25$ | 0.08 | no |
| shear | $0.15 \pm 1.72$ | $-0.12 \pm 0.56$ | 0.49 | no |
| stretch | $-0.09 \pm 0.31$ | $-0.21 \pm 0.12$ | 1.22 | no |
| stagger | $0.01 \pm 0.56$ | $0.15 \pm 0.30$ | 0.79 | no |
| shift | $0.04 \pm 1.51$ | $0.01 \pm 0.84$ | 0.07 | no |
| slide | $0.76 \pm 1.13$ | $0.76 \pm 0.72$ | 0.006 | no |
| rise | $3.15 \pm 0.43$ | $3.29 \pm 0.26$ | 1.01 | no |
| helical rise | $3.20 \pm 0.81$ | $3.33 \pm 0.33$ | 0.50 | no |
| Helical parameter $\left(^{\circ}\right)$ |  |  |  |  |
| inclination | $7.44 \pm 14.1$ | $7.36 \pm 9.54$ | 0.02 | no |
| tip | $1.66 \pm 11.1$ | $0.49 \pm 8.79$ | 0.33 | no |

## Table A. 9 (cont'd.)

| buckle | $2.72 \pm 10.9$ | $-0.55 \pm 7.17$ | 0.94 | no |
| :--- | :---: | :---: | :---: | :---: |
| propellor | $-9.12 \pm 12.6$ | $-15.13 \pm 5.79$ | 1.51 | no |
| opening | $-3.39 \pm 6.04$ | $-0.53 \pm 5.60$ | 1.49 | no |
| tilt | $-0.36 \pm 5.70$ | $0.11 \pm 4.94$ | 0.26 | no |
| roll | $3.88 \pm 7.14$ | $4.24 \pm 4.75$ | 0.16 | no |
| twist | $32.6 \pm 9.00$ | $35.04 \pm 9.32$ | 0.84 | no |
| helical twist | $34.1 \pm 8.32$ | $35.95 \pm 8.88$ | 0.69 | no |

Table A.10. Polyamide distances (distances in $\AA$; standard deviation for x-ray structure is $2.27 \AA$ ) for hydroxypyrrole polyamide simulation

| Simulation 1 <br> distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| x1 | $3.43 \pm 0.13$ | 3.34 | 0.12 | no |
| x2 | $3.31 \pm 0.12$ | 2.92 | 0.55 | no |
| x3 | $3.13 \pm 0.14$ | 2.73 | 0.55 | no |
| x4 | $4.17 \pm 0.15$ | 3.23 | 1.31 | no |
| x5 | $4.93 \pm 0.17$ | 3.16 | 2.47 | yes |
| x6 | $3.28 \pm 0.15$ | 2.51 | 1.07 | no |
| x7 | $3.14 \pm 0.07$ | 2.99 | 0.21 | no |
| x1a | $5.53 \pm 0.28$ | 2.60 | 4.17 | yes |
| x2a | $4.36 \pm 0.14$ | 3.30 | 1.73 | no |
| x3a | $3.13 \pm 0.12$ | 2.67 | 0.64 | no |
| x4a | $3.06 \pm 0.7$ | 2.91 | 0.29 | no |
| x5a | $3.58 \pm 0.14$ | 3.25 | 0.47 | no |
| x6a | $2.98 \pm 0.09$ | 2.75 | 0.31 | no |
| x7a | $3.28 \pm 0.08$ | 3.00 | 0.94 | no |

Table A.11. Computed and x-ray averages and standard deviations of helical parameters for DNA/netropsin simulation

| helical parameter $(\mathbf{\AA})$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :--- | :--- | :---: | :--- | :--- |
| x-displacement | $-0.55 \pm 1.84$ | $0.24 \pm 1.31$ | 1.49 | no |
| y-displacment | $-0.05 \pm 0.20$ | $0.04 \pm 1.12$ | 1.41 | no |
| shear | $0.004 \pm 0.56$ | $0.09 \pm 0.55$ | 0.53 | no |
| stretch | $-0.13 \pm 0.17$ | $-0.22 \pm 0.24$ | 1.86 | no |
| stagger | $-0.09 \pm 0.47$ | $0.13 \pm 0.49$ | 1.65 | no |
| shift | $0.06 \pm 0.76$ | $-0.14 \pm 0.64$ | 0.91 | no |
| slide | $0.05 \pm 0.62$ | $0.27 \pm 0.61$ | 1.24 | no |
| rise | $3.29 \pm 0.33$ | $3.31 \pm 0.38$ | 0.18 | no |
| helical rise | $3.18 \pm 0.40$ | $3.27 \pm 0.40$ | 0.80 | no |
| Helical parameter $\left({ }^{\circ}\right)$ |  |  |  |  |
| inclination | $6.08 \pm 13.1$ | $0.82 \pm 8.32$ | 1.39 | no |

## Table A. 11 (cont'd.)

| tip | $-0.43 \pm 9.34$ | $0.87 \pm 8.87$ | 0.48 | no |
| :--- | ---: | ---: | ---: | :--- |
| buckle | $-1.62 \pm 12.1$ | $-3.23 \pm 10.6$ | 0.46 | no |
| propellor | $-10.7 \pm 10.7$ | $-11.81 \pm 10.5$ | 0.35 | no |
| opening | $1.30 \pm 6.18$ | $-0.09 \pm 7.21$ | 0.77 | no |
| tilt | $0.18 \pm 5.48$ | $-0.41 \pm 5.24$ | 0.37 | no |
| roll | $3.36 \pm 7.46$ | $0.07 \pm 5.14$ | 1.52 | no |
| twist | $34.8 \pm 6.12$ | $35.7 \pm 5.44$ | 0.54 | no |
| helical twist | $36.1 \pm 5.52$ | $36.4 \pm 5.28$ | 0.16 | no |

Table A.12. Polyamide distances (distances in $\AA \AA$; standard deviation for x-ray structure is $2.20 \AA$ (resolution)) for DNA/netropsin simulation

| Simulation 1 <br> distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| 1 | $3.59 \pm 0.66$ | 2.52 | 1.54 | no |
| 2 | $6.55 \pm 0.60$ | 3.24 | 4.76 | yes |
| 3 | $3.11 \pm 0.20$ | 3.40 | 0.42 | no |
| 4 | $3.25 \pm 0.18$ | 3.50 | 0.36 | no |
| 5 | $3.69 \pm 0.35$ | 2.70 | 1.42 | no |
| 6 | $3.38 \pm 0.26$ | 3.70 | 0.46 | no |
| 7 | $3.33 \pm 0.29$ | 2.70 | 0.91 | no |
| 8 | $3.35 \pm 0.18$ | 2.90 | 0.65 | no |
| 9 | $3.66 \pm 0.31$ | 3.30 | 0.52 | no |
| 10 | $5.43 \pm 0.63$ | 3.20 | 3.20 | yes |
| 11 | $5.87 \pm 0.74$ | 3.50 | 3.40 | yes |

Table A.13. Computed and x-ray averages and standard deviations of helical parameters for DNA/distamycin simulation

| helical parameter $(\AA) \mathbf{~ c o m p u t e d ~ a v g . / s d ~}$ | x-ray value/sd | t-test value | significant? |  |
| :--- | :---: | :---: | :--- | :--- |
| x-displacement | $0.06 \pm 2.08$ | $0.25 \pm 1.41$ | 0.44 | no |
| y-displacment | $-0.28 \pm 1.42$ | $0.16 \pm 0.79$ | 1.81 | no |
| shear | $0.04 \pm 0.85$ | $-0.07 \pm 0.28$ | 1.26 | no |
| stretch | $-0.15 \pm 0.38$ | $-0.11 \pm 0.08$ | 1.38 | no |
| stagger | $0.36 \pm 1.73$ | $0.14 \pm 0.36$ | 1.83 | no |
| shift | $0.20 \pm 1.01$ | $-0.08 \pm 0.47$ | 1.92 | no |
| slide | $0.35 \pm 1.05$ | $0.14 \pm 0.79$ | 0.88 | no |
| rise | $3.36 \pm 0.51$ | $3.37 \pm 0.38$ | 0.07 | no |
| helical rise | $3.46 \pm 1.05$ | $3.33 \pm 0.47$ | 0.89 | no |
| Helical parameter $\left({ }^{\circ}\right)$ |  |  |  |  |
| inclination | $3.01 \pm 14.8$ | $-0.49 \pm 9.87$ | 1.17 | no |
| tip | $1.29 \pm 10.6$ | $-0.02 \pm 7.50$ | 0.57 | no |
| buckle | $1.38 \pm 9.29$ | 0.28 | no |  |

Table A. 13 (cont'd).

| propellor | $-11.6 \pm 20.5$ | $-15.1 \pm 5.50$ | 1.91 | no |
| :--- | ---: | ---: | ---: | :--- |
| opening | $1.51 \pm 9.58$ | $-2.73 \pm 4.17$ | 3.29 | yes |
| tilt | $-0.89 \pm 7.06$ | $0.08 \pm 4.84$ | 0.66 | no |
| roll | $1.36 \pm 11.1$ | $-0.23 \pm 5.95$ | 0.87 | no |
| twist | $36.7 \pm 11.3$ | $36.5 \pm 4.64$ | 0.10 | no |
| helical twist | $38.5 \pm 12.1$ | $37.2 \pm 4.47$ | 0.88 | no |

Table A.14. Polyamide distances (distances in $\AA$; standard deviation for x-ray structure is $2.20 \AA$ ) for DNA/distamycin simulation

| Simulation 1 <br> distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| 1 | $4.65 \pm 0.86$ | 3.03 | 2.33 | yes |
| 2 | $3.26 \pm 0.21$ | 3.17 | 0.13 | no |
| 3 | $3.47 \pm 0.42$ | 2.93 | 0.77 | no |
| 4 | $3.54 \pm 0.29$ | 3.14 | 0.58 | no |
| 5 | $3.76 \pm 0.41$ | 3.25 | 0.74 | no |
| 6 | $4.18 \pm 0.97$ | 3.13 | 1.50 | no |
| 7 | $3.91 \pm 0.69$ | 3.24 | 0.96 | no |

Table A.15. Computed and x-ray averages and standard deviations of helical parameters for DNA/netropsin simulation with AMBER parameters

| helical parameter $(\mathbf{\AA})$ | computed avg./sd | x-ray value/sd | t-test value | significant? |
| :--- | :--- | :---: | :--- | :--- |
| x-displacement | $-0.78 \pm 2.98$ | $0.24 \pm 1.31$ | 2.53 | yes |
| y-displacment | $-0.25 \pm 3.27$ | $0.04 \pm 1.12$ | 0.82 | no |
| shear | $-0.22 \pm 0.97$ | $0.09 \pm 0.55$ | 1.86 | no |
| stretch | $0.12 \pm 0.61$ | $-0.22 \pm 0.24$ | 4.52 | yes |
| stagger | $-0.17 \pm 0.62$ | $0.13 \pm 0.49$ | 2.04 | yes |
| shift | $-0.38 \pm 1.43$ | $-0.14 \pm 0.64$ | 1.22 | no |
| slide | $-0.32 \pm 0.61$ | $0.27 \pm 0.61$ | 3.19 | yes |
| rise | $3.26 \pm 0.39$ | $3.31 \pm 0.38$ | 0.39 | no |
| helical rise | $2.95 \pm 0.88$ | $3.27 \pm 0.40$ | 2.57 | yes |
| Helical parameter $\left({ }^{\circ}\right)$ |  |  |  |  |
| inclination | $10.5 \pm 21.8$ | $0.82 \pm 8.32$ | 3.76 | yes |
| tip | $-0.61 \pm 11.2$ | $0.87 \pm 8.87$ | 0.55 | no |
| buckle | $-1.30 \pm 13.1$ | $-3.23 \pm 10.6$ | 0.61 | no |
| propellor | $-16.5 \pm 16.3$ | $-11.81 \pm 10.5$ | 1.47 | no |
| opening | $-9.22 \pm 42.8$ | $-0.09 \pm 7.21$ | 3.65 | yes |
| tilt | $0.57 \pm 6.83$ | $-0.41 \pm 5.24$ | 0.62 | no |
| roll | $0.65 \pm 11.7$ | $0.07 \pm 5.14$ | 0.36 | no |
| twist | $30.1 \pm 15.1$ | $35.7 \pm 5.44$ | 3.30 | yes |
| helical twist | $29.5 \pm 20.9$ | $36.4 \pm 5.28$ | 4.03 | yes |

Table A.16. Polyamide distances (distances in $\AA$; standard deviation for x-ray structure is $2.20 \AA$ (resolution)) for DNA/netropsin simulation with AMBER parameters

| Simulation 1 <br> distance | computed average | x-ray | t-value | significant? |
| :--- | :--- | :--- | :--- | :--- |
| 1 | $3.17 \pm 0.29$ | 2.52 | 0.94 | no |
| 2 | $5.78 \pm 0.41$ | 3.24 | 3.65 | yes |
| 3 | $2.85 \pm 0.12$ | 3.40 | 0.79 | no |
| 4 | $3.33 \pm 0.21$ | 3.50 | 0.25 | no |
| 5 | $3.06 \pm 0.22$ | 2.70 | 0.52 | no |
| 6 | $3.17 \pm 0.25$ | 3.70 | 0.77 | no |
| 7 | $3.06 \pm 0.20$ | 2.70 | 0.52 | no |
| 8 | $3.53 \pm 0.25$ | 2.90 | 0.91 | no |
| 9 | $3.40 \pm 0.27$ | 3.30 | 0.15 | no |
| 10 | $5.29 \pm 0.55$ | 3.20 | 3.01 | yes |
| 11 | $6.62 \pm 0.50$ | 3.50 | 4.49 | yes |

Table A.17. T-values for combined high and low temperature DNA/polyamide simulations

| Parameter ( $\AA$ ) | $\begin{aligned} & 300 \mathrm{~K}(6 \\ & \text { simulations) } \end{aligned}$ | $\begin{aligned} & 113 \mathrm{~K}(8 \\ & \text { simulations) } \end{aligned}$ | all (14 simulations) |
| :---: | :---: | :---: | :---: |
| shear | 0.391 | 0.228 | 0.271 |
| stretch | 0.752 | 0.278 | 0.35 |
| stagger | 2.24 | 3.11 | 1.4 |
| shift | 0.214 | 0.089 | 0.127 |
| slide | 0.508 | 0.121 | 0.25 |
| rise | 0.231 | 0.267 | 0.171 |
| x -displacement | 0.727 | 0.078 | 0.366 |
| y -displacement | 0.214 | 0.102 | 0.119 |
| helical rise | 0.391 | 0.505 | 0.331 |
| Parameter ( ${ }^{\circ}$ ) |  |  |  |
| buckle | 0.642 | 1.92 | 1.276 |
| propellor | 0.963 | 0.306 | 0.663 |
| opening | 0.443 | 0.303 | 0.198 |
| tilt | 0.222 | 0.178 | 0.113 |
| roll | 1.33 | 0.533 | 0.98 |
| twist | 0.668 | 0.291 | 0.478 |
| inclination | 1.11 | 0.676 | 0.949 |
| tip | 0.253 | 0.307 | 0.244 |
| helical twist | 0.349 | 0.425 | 0.397 |

## CHARMM Topology and Parameters for Polyamides


1


2

3



Figure A. 1 Model compounds for polyamides (some hydrogens not shown for clarity).

Topology for 1 (topology was added to the topology for small molecules, top_all22_model.inp and topology for DNA/proteins, top_all27_prot_na.inp)



Parameters for 1 (parameters were added to the parameters for DNA/proteins, par_all27_prot_na.inp)
BONDS
!n-methyl-pyrrole

| CPT | HP | 500.00 | 1.08 |
| :--- | :--- | :--- | :--- |
| CT3 | NY | 500.000 | 1.46 |
| C | CPT | 500.00 | 1.47 |
| NH1 | CY | 500.00 | 1.40 |
| C | H | 500.00 | 1.10 |

ANGLES

| ! N-methyl-pyrrole |  |  |  |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| CPT | CY | NH1 | 40.000 | 124.9 | CA | CA | NH1 | 40.000 | 127.2 |
| HP | CPT | CY | 40.000 | 126.9 |  |  |  |  |  |
| HP | CPT | CPT | 40.000 | 125.9 |  |  |  |  |  |
| CA | NY | CT3 | 40.000 | 123.3 |  |  |  |  |  |
| CPT | NY | CT3 | 40.000 | 127. |  |  |  |  |  |
| CPT | CPT | C | 40.000 | 130.2 |  |  |  |  |  |


| NY | CPT | C | 40.000 | 122.1 |
| :--- | :--- | :--- | :--- | :--- |
| NY | CT3 | HA | 40.000 | 109.5 |
| CPT | C | O | 40.000 | 123.5 |
| CPT | C | NH1 | 40.000 | 114.4 |
| H | NH1 | H | 40.000 | 115.7 |
| CY | NH1 | H | 40.000 | 117.2 |
| CY | NH1 | C | 40.00 | 125.2 |
| NH1 | C | H | 40.00 | 112.0 |
| CA | CY | NH1 | 40.000 | 127.2 |
| CA | CY | C | 40.00 | 100.3 |
| CPT | CY | C | 40.000 | 151.9 |
| CPT | CY | CPT | 40.000 | 36.3 |
| DIHEDRALS |  |  |  |  |
| !N-methyl-pyrrole |  |  |  |  |


| CA | NY | CT3 | HA | 1.60 | 1 | 0.00 | !methyl group |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CA | NY | CT3 | HA | 0.000 | 3 | -120.0 |  |
| CA | NY | CT3 | HA | 0.000 | 2 | -60.0 |  |
| $i$ |  |  |  |  |  | ---- |  |
| CPT | NY | CT3 | HA | 0.065 | 3 | 0.0 | !methyl group |
| CPT | NY | CT3 | HA | 0.00 | 1 | -120. |  |
| CPT | NY | CT3 | HA | 0.00 | 2 | -60.0 |  |
| HP | CPT | CY | CA | 2.0 | 1 | 180.0 |  |
| HP | CPT | CY | NH1 | 2.8 | 2 | 180.0 |  |
| CY | CA | NY | CT3 | 0.8 | 2 | 180.0 |  |
| CY | NH1 | C | $\bigcirc$ | 1.5 | 2 | 180.0 |  |
| CY | NH1 | C | H | 1.20 | 1 | 180.0 |  |
| CPT | CY | NH1 | H | 0.000 | 1 | 0.00 |  |
| CPT | CY | NH1 | C | 0.200 | 1 | 180.0 |  |
| CPT | CPT | NY | CT3 | 0.8 | 2 | 180.0 |  |
| CPT | CPT | C | $\bigcirc$ | 2.5 | 2 | 180.0 |  |
| CPT | CPT | C | NH1 | 1.5 | 2 | 180.0 |  |
| CA | NY | CPT | C | 2.00 | 1 | 180.0 |  |
| HP | CA | CY | NH1 | 1.00 | 2 | 180.0 |  |
| CA | CY | NH1 | H | 0.0 | 1 | 0.00 |  |
| CA | CY | NH1 | C | 2.0 | 2 | 180.0 |  |
| CA | CY | NH1 | C | 2.0 | 1 | 180.0 |  |
| HP | CA | NY | CT3 | 0.4000 | 2 | 180.0 |  |
| NY | CA | CY | NH1 | 2.0 | 2 | 180.0 |  |
| NY | CPT | C | 0 | 1.0 | 2 | 180.0 |  |
| NY | CPT | C | $\bigcirc$ | 3.0 | 1 | 180.0 |  |
| NY | CPT | C | NH1 | 1.0 | 1 | 0.0 |  |
| CPT | CPT | CY | NH1 | 3.0 | 2 | 180.0 |  |
| CT3 | NY | CPT | C | 0.8 | 2 | 180.0 |  |
| H | NH1 | C | H | 1.4 | 2 | 180.0 |  |
| CPT | C | NH1 | H | 2.5 | 2 | 180.0 |  |

Topology for 2 (topology was added to topology for small molecules, top_all22_model.inp and to DNA/proteins, top_all27_prot_na.inp) RESI IMID $0.00 \quad$ ! N -methyl-imidazole


| ic | 03 | CC | N3 | HC | 1.2233 | 122.84 | -180. | 119.23 | 0.9933 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ic | 03 | CC | N3 | HD | 1.2233 | 122.84 | 0.00 | 119.230 | 0.9933 |
| ic | HC | N3 | CC | CE1 | 0.9933 | 119.23 | 0 | 116.251 | 1.47 |
| ic | HD | N3 | CC | CE1 | 0.9933 | 119.23 | 180.0 | 116.251 | 1.47 |
| ic | N3 | CC | CE1 | ND1 | 1.3418 | 116.25 | 180. | 122.1 1 | 1.38 |
| ic | N3 | CC | CE1 | NE2 | 1.3418 | 116.25 | 0. | 130.31 | 1.40 |
| ic | ND1 | CE1 | CC | 03 | 1.38 | 122.0 | 0.00 | 123.51 | 1.2233 |
| ic | NE2 | CE1 | CC | 03 | 1.40 | 130.3 | 180.0 | 123.51 | 1.2233 |
| IC | CD2 | NE2 | CE1 | CC | 1.40 | 107.1 | -180.0 | 130.31 | 1.46 |
| ic | CG | ND1 | CE1 | CC | 1.37 | 109.8 | -180.0 | 122.1 | 1.46 |
| ic | C2 | ND1 | CE1 | CC | 1.46 | 127.0 | 0.00 | 122.1 1 | 1.46 |
| !ic's for amide group on CD2 |  |  |  |  |  |  |  |  |  |
| IMPR CD HY N4 O4 |  |  |  |  |  |  |  |  |  |
| ic | 04 | CD | N4 | HN | 1.2233 | 122.84 | 180. | 119.23 | 0.9933 |
| ic | HN | N4 | CD | HY | 0.9933 | 119.23 | 0. | 112.25 | 1.10 |
| ic | 04 | CD | N4 | CD2 | 1.223 | 3122.84 | 40. | 125.57 | 71.40 |
| ic | CD | N4 | CD2 | NE2 | 1.3418 | 8122.57 | $57-180$ | - 124.9 | 91.41 |
| ic | CD | N4 | CD2 | CG | 1.3418 | 122.57 | 0. | 127.2 | 1.40 |
| IC | N4 | CD2 | NE2 | CE1 | 1.40 | 124.9 | -180. | 107.1 | 1.40 |
| ic | N4 | CD2 | CG | ND1 | 1.40 | 127.2 | -180.0 | 107.6 | 1.37 |
| impr HN CG CD2 N4 |  |  |  |  |  |  |  |  |  |

Parameters for 2 (parameters were added to parameters for DNA/proteins, par_all27_prot_na.inp)
BONDS

| ! N-methyl-imidazole |  |  |
| :--- | :--- | :--- |
| NR1 CT3 | 500.00 |  |
| CPH2 C | 500.000 | 1.46 |
| CPH1 NH1 | 500.00 | 1.47 |
| C H | 500.00 | 1.10 |

ANGLES
! N-methyl-imidazole

| CPH1 | CPH1 | NH1 | 40.000 | 129.3 |
| :--- | :--- | :--- | :--- | :--- |
| NR2 | CPH1 | NH1 | 40.000 | 120.0 |
| CPH1 | NR1 | CT3 | 40.000 | 124.7 |
| CPH2 | NR1 | CT3 | 40.000 | 127.8 |
| NR1 | CPH2 | C | 40.000 | 123.8 |
| NR2 | CPH2 | C | 40.000 | 124.8 |
| NR1 | CT3 | HA | 40.000 | 107.4 |
| CPH2 | C | O | 40.000 | 122.6 |
| CPH2 | C | NH1 | 40.000 | 112.9 |
| H | NH1 | H | 40.000 | 121.1 |
| CPH1 | NH1 | H | 40.000 | 115.7 |
| CPH1 | NH1 | C | 40.000 | 124.5 |
| NH1 | C | H | 40.000 | 112.4 |

DIHEDRALS
! N-methyl-imidazole

| CPH1 NR1 | CT3 | HA | 1.60 | 1 | 0.00 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CPH1 NR1 | CT3 | HA | 0.000 | 3 | -120.0 | !methyl group |
| CPH1 NR1 | CT3 | HA | 0.000 | 2 | -60.0 |  |




Topology for 3 (added to topology for DNA/proteins, top_all27_prot_na.inp)



| ic C1 | NE1 | CE2 | CA | 1.46 | 127.0 | 0.00 | 122.1 | 1.46 |  |
| :--- | :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: | :---: |
| !ic's for amide | group | on CG |  |  |  |  |  |  |  |
| ic | O2 | CB | N2 | HM | 1.2233 | 122.84 | -180. | 119.23 | 0.9933 |
| ic | HM | N2 | CB | HZ | 0.9933 | 119.23 | 0. | 112.25 | 1.10 |
| ic | O2 | CB | N2 | CG | 1.2233 | 122.84 | 0. | 125.57 | 1.40 |
| ic | CB | N2 | CG | CD2 | 1.3418 | 122.57 | 180. | 124.9 | 1.41 |
| ic | CB | N2 | CG | CD1 | 1.3418 | 122.57 | 0. | 127.2 | 1.40 |
| IC | N2 | CG | CD2 | CE2 | 1.40 | 124.9 | 180. | 107.1 | 1.40 |
| ic | N2 | CG | CD1 | NE1 | 1.40 | 127.2 | -180.0 | 107.6 | 1.37 |

Parameters for 3 (added to parameters for DNA/proteins, par_all27_prot_na.inp) BONDS

| !hydroxy |  |  | pyrrole |
| :--- | :---: | :---: | :---: |
| CPT | OH1 | 500.000 | 1.37 |
| CT3 | NY | 500.000 | 1.46 |
| C | CPT | 500.00 | 1.45 |
| NH1 | CY | 500.000 | 1.40 |
| C | H | 500.000 | 1.10 |

ANGLES
! Hydroxy pyrrole

| CPT | CY | NH1 | 40.00 | 123.1 |
| :--- | :--- | :--- | :--- | :--- |
| CA | CY | NH1 | 40.000 | 129.7 |
| CY | CPT | OH1 | 40.00 | 123.1 |
| CPT | CPT | OH1 | 40.00 | 129.0 |
| CA | NY | CT3 | 40.000 | 124.3 |
| CPT | NY | CT3 | 40.000 | 125.9 |
| CPT | CPT | C | 40.000 | 128.3 |
| NY | CPT | C | 40.000 | 124.3 |
| CPT | OH1 | H | 40.000 | 106.4 |
| NY | CT3 | HA | 40.00 | 107.4 |
| CPT | C | O | 40.000 | 125.5 |
| CPT | C | NH1 | 40.00 | 112.5 |
| H | NH1 | H | 40.000 | 111.8 |
| CY | NH1 | H | 40.00 | 116.2 |
| CY | NH1 | C | 40.000 | 124.6 |
| NH1 | C | H | 40.000 | 112.3 |
| CPT | CPT | CC | 40.00 | 128.25394 |
| NY | CPT | CC | 40.00 | 124.31412 |
| CPT | CC | O | 40.00 | 125.47897 |
| CPT | CC | NH2 | 40.00 | 112.47125 |
| DIHEDRALS |  |  |  |  |
| !Hydroxy | Pyrrole |  |  |  |

!from pyrrole

$\begin{array}{lllll}\text { HP CPT CY CA } 2.0 & 180.0\end{array}$

| HP | CPT | CY | NH1 | 2.8 | 2 |  | 180.0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CY | CA | NY | CT3 | 0.8 | 2 |  | 180.0 |
| CY | NH1 | C | 0 | 2.0 | 2 |  | 180.0 |
| CY | NH1 | C | H | 0.3 | 1 |  | 180.0 |
| CPT | CY | NH1 | H | 0.000 | 1 |  | 0.00 |
| CPT | CY | NH1 | C | 0.5 | 1 |  | 180.0 |
| CPT | CPT | NY | CT3 | 0.8 | 2 |  | 180.0 |
| CPT | CPT | C | 0 | 2.5 | 2 |  | 180.0 |
| CPT | CPT | C | NH1 | 2.5 | 2 |  | 180.0 |
| CPT | CPT | C | NH1 | 2.0 | 1 |  | 0.00 |
| CA | NY | CPT | C | 2.30 | 2 |  | 180.0 |
| HP | CA | CY | NH1 | 1.0 | 2 |  | 180.0 |
| CA | CY | NH1 | H | 1.5 | 1 |  | 0.00 |
| CA | CY | NH1 | C | 1.5 | 2 |  | 180.0 |
| CA | CY | NH1 | C | 0.0 | 1 |  | 180.0 |
| HP | CA | NY | CT3 | 0.4000 | 2 |  | 180.0 |
| NY | CA | CY | NH1 | 1.0 | 2 |  | 180.0 |
| NY | CPT | C | $\bigcirc$ | 1.0 | 2 |  | 180.0 |
| NY | CPT | C | 0 | 5.5 | 1 |  | 180.0 |
| NY | CPT | C | NH1 | 1.0 | 1 |  | 180.0 |
| CPT | CPT | CY | NH1 | 4.0 | 2 |  | 180.0 |
| CT3 | NY | CPT | C | 0.8 | 2 |  | 180.0 |
| H | NH1 | C | H | 1.4 | 2 |  | 180.0 |
| CPT | C | NH1 | H | 4.0 | 2 |  | 180.0 |
| CY | CPT | OH1 |  | 0.0000 | 2 |  | 90.0 |
| CY | CPT | OH1 |  | 4.0 | 1 |  | 0.0 |
| CPT | CPT | OH1 H |  | 0.5 |  | 2 | 180.00 |
| CPT | CPT | OH1 H |  | 0.30 |  | 1 | 0.0 |
| OH1 | CPT | CY NH |  | 0.200 |  | 2 | 180.0 |
| CA | CY | CPT O | H1 | 3.1 |  | 2 | 180.0 |
| IMPROPER |  |  |  |  |  |  |  |
| !hydroxy-pyrrole |  |  |  |  |  |  |  |
| C | CPT | NH1 | 0 | 90.0 |  | 0 | 0.00 |
| C | NH1 | CPT | 0 | 90.0 |  | 0 | 0.00 |
| 0 | NY | CPT | C | 90.00 |  | 0 | 0.00 |
| CPT | CA | NY | CT3 | 90.000 |  | 0 | -180.0 |
| C | NH1 | $\bigcirc$ | CPT | 90.000 |  | 0 | 0.00 |
| NH1 | NY | CPT | C | 90.000 |  | 0 | 0.00 |
| 01 | CPT | CPT | C | 90.000 |  | 0 | 0.00 |


| H | CY | CPT | OH1 | 45.000 | 0 | -26.0 |
| :--- | :--- | :--- | :--- | :--- | :--- | :---: |
| OH1 | CY | CPT | CPT | 90.000 | 0 | -180.0 |
| H | CPT | CPT | OH1 | 90.000 | 0 | 14.0 |
| NH1 | CA | CY | CPT | 90.00 | 0 | -180.0 |

Topology for 4 (added to topology for DNA/proteins, top_all27_prot_na.inp)


| ATOM HR3 HA | $0.09!$ |  |
| :--- | :--- | ---: |
|  |  | $!$ |
| GROUP |  | $-0.11!$ |
| ATOM CS | CT2 | $0.09!$ |
| ATOM HS1 | HA | $0.09!$ |
| AtOM HS2 | HA | $0.11!$ |
| ATOM CT | CT2 | -0.1 |
| ATOM HT1 | HA | $0.09!$ |
| ATOM HT2 | HA | $0.09!$ |

GROUP

| ATOM N1 | NH3 | -0.30 |
| :--- | :--- | ---: |
| ATOM H1 | HC | 0.33 |
| ATOM CX | CT3 | 0.23 |
| ATOM HX1 | HA | 0.05 |
| ATOM HX2 | HA | 0.05 |
| ATOM HX3 | HA | 0.05 |
| ATOM CY | CT3 | 0.23 |
| ATOM HY1 | HA | 0.05 |
| ATOM HY2 | HA | 0.05 |
| ATOM HY3 | HA | 0.05 |

BOND CL C C N N CR N H
BOND CL HL1 CL HL2 CL HL3
BOND CR HR1 CR HR3 CR CS
DOUBLE C O
BOND CS HS1 CS HS2 CS CT
BOND CT HT1 CT HT2 CT N1
BOND N1 H1 N1 CX N1 CY
BOND CX HX1 CX HX2 CX HX3
BOND CY HY1 CY HY2 CY HY3

```
!-----------------------------------------------------------
```

!ic table
IMPR N C CR H
IMPR C CL N O
IC O C N H 1.2233 122.84 180. 119.23 0.9933


Parameters for 4 (added to parameters for DNA/proteins, top_all27_prot_na.inp) ANGLES

| !cationic tail |  |  |  |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: | :---: |
| CT2 | NH3 | CT3 | 40.00 | 111.4 |  |  |  |  |  |
| CT3 | NH3 | CT3 | 40.0 | 110.8 |  |  |  |  |  |

DIHEDRALS

| CT3 | C | NH1 | H | 2.5 | 2 | 180.0 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| CT3 | C | NH1 | CT2 | 1.6 | 1 | 0.00 |
| CT3 | C | NH1 | CT2 | 2.5 | 2 | 180.0 |
| HA | CT3 | C | O | 0.0 | 3 | 180.0 |
| HA | CT3 | C | NH1 | 0.0 | 3 | 0.00 |
|  |  |  |  |  |  |  |
| C | NH1 | CT2 | HA | 0.00 | 3 | 0.00 |
| C | NH1 | CT2 | CT2 | 1.0 | 2 | 0.00 |
| C | NH1 | CT2 | CT2 | 0.5 | 1 | 0.00 |
| O |  | C | NH1 | H | 2.5 | 2 |
| O | C | NH1 | CT2 | 2.5 | 2 | 180.0 |
| NH1 | CT2 | CT2 | HA | 0.1950 | 3 | 0.00 |
| NH1 | CT2 | CT2 | CT2 | 0.1950 | 3 | 0.00 |
| H | NH1 | CT2 | HA | 0.000 | 3 | 0.00 |
| H | NH1 | CT2 | CT2 | 0.000 | 1 | 0.00 |


| CT2 | CT2 | CT2 | HA | 0.1950 | 3 | 0.00 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| CT2 | CT2 | CT2 | NH3 | 0.1950 | 3 | 0.00 |
| H | CT2 | CT2 | HA | 0.1950 | 3 | 0.00 |
| CT2 | CT2 | NH3 | HC | 0.100 | 3 | 0.00 |
| CT2 | CT2 | NH3 | CT3 | 0.4000 | 3 | 0.000 |
|  |  |  |  |  |  |  |
| HA | CT2 | CT2 | NH3 | 0.1950 | 3 | 0.00 |
| CT2 | NH3 | CT3 | HA | 0.09000 | 3 | 0.00 |
|  |  |  |  |  |  |  |
| HA | CT2 | CT2 | HA | 0.1950 | 3 | 0.00 |
|  |  |  |  |  | 0.1000 | 3 |

Topology for 5 (added to topology for DNA/proteins, top_all27_prot_na.inp)

| RESI GAMM GROUP |  | 0.00 | !gamma aminobutyric acid linker for polyamides |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | ! |  |  |  |  |  |  |  |
| ATOM CL | CT3 | -0.27! |  |  |  |  |  |  |  |
| ATOM HL1 | HA | $0.09!$ | HL1 | H | HR1 | HS1 | HT1 | OM | H2 |
| ATOM HL2 | HA | 0.09! | $\backslash$ | \| | \| | \| | \| | \| | | / |
| ATOM HL3 | HA | $0.09!$ | HL2--CL | C---N | R | CS | - | CM- |  |
| ATOM C | C | 0.51 ! | / | \| | | \| | । | \| |  | $\backslash$ |
| ATOM O | $\bigcirc$ | -0.51! | HL3 | 0 | HR2 | HS2 | HT2 |  | H3 |
| ATOM N | NH1 | -0.47! |  |  |  |  |  |  |  |
| ATOM H | H | 0.47 ! |  |  |  |  |  |  |  |
| ATOM CR | CT2 | -0.18! |  |  |  |  |  |  |  |
| ATOM HR1 | HA | $0.09!$ |  |  |  |  |  |  |  |
| ATOM HR3 | HA | 0.09! |  |  |  |  |  |  |  |
|  |  | ! |  |  |  |  |  |  |  |
| GROUP |  | ! |  |  |  |  |  |  |  |
| ATOM CS | CT2 | -. 18 ! |  |  |  |  |  |  |  |
| ATOM HS1 | HA | 0.09 |  |  |  |  |  |  |  |
| ATOM HS2 | HA | 0.09 |  |  |  |  |  |  |  |
| ATOM CT | CT2 | -. 18 |  |  |  |  |  |  |  |
| ATOM HT1 | HA | 0.09 |  |  |  |  |  |  |  |
| ATOM HT2 | HA | 0.09 |  |  |  |  |  |  |  |
| GROUP |  |  |  |  |  |  |  |  |  |
| ATOM CM | CC | 0.51 |  |  |  |  |  |  |  |
| ATOM OM | $\bigcirc$ | -0.51 |  |  |  |  |  |  |  |
| ATOM N2 | NH2 | -0.47 |  |  |  |  |  |  |  |
| ATOM H2 | H | 0.235 |  |  |  |  |  |  |  |
| ATOM H3 | H | 0.235 |  |  |  |  |  |  |  |

```
BOND CL HL1 CL HL2 CL HL3
BOND CL C C N N H N CR
DOUBLE C O
BOND CR HR1 CR HR3 CR CS
BOND CS HS1 CS HS2 CS CT
BOND CT HT1 CT HT2 CT CM
DOUBLE CM OM
BOND CM N2 N2 H2 N2 H3
!ic table
IMPR N C CR H
IMPR C CL N O
IC O C N H 1.2233 122.84 180. 119.23 0.9933
IC H N C CL 0.9933 119.23 0. 116.25 1.5118
IC O C N CR 1.2233 122.84 0. 122.57 1.4488
IC N C CL HL1 1.3418 116.25 180. 109.3 1.109
IC N C CL HL2 1.3418 116.25 60. 109.3 1.109
IC N C CL HL3 1.3418 116.25 -60. 109.3 1.109
IC C N CR HR1 1.3418 122.57 -30. 110.7 1.113
IC C N CR CS 1.3418 122.57 90. 110.7 1.50
IC C N CR HR3 1.3418 122.57 -150. 110.7 1.113
IC N CR CS CT 1.47 109.5 -180.0 109.5 1.50
IC N CR CS HS1 1.47 109.5 -60.0 109.5 1.10
IC N CR CS HS2 1.47 109.5 60.0 109.5 1.10
IC CR CS CT HT1 1.50 109.5 60.0 109.5 1.10
IC CR CS CT HT2 1.50 109.5 -60.0 109.5 1.10
IC CR CS CT CM 1.50 109.5 180.0 109.5 1.50
IC CS CT CM OM 1.50 109.5 -30.00 123.1 1.22233
IC CS CT CM N2 1.50 109.5 150.0 115.0 1.37
IC CT CM N2 H2 1.50 115.0 -180.0 120.6 1.10
IC CT CM N2 H3 1.50 115.0 0.00 120.6 1.10
IMPR CM CT N2 OM
```

Parameters for 5 (added to parameters for DNA/proteins, par_all27_prot_na.inp)
ANGLES
!gamma linker
$\begin{array}{llll}\mathrm{H} & \text { NH1 H } & 25.00 & 117.0\end{array}$
DIHEDRALS

| CT3 | C | NH1 | H | 2.5 | 2 | 180.0 |
| :--- | :--- | :--- | :--- | :--- | :--- | ---: |
| CT3 | C | NH1 | CT2 | 1.6 | 1 | 0.00 |
| CT3 | C | NH1 | CT3 | 2.5 | 2 | 180.0 |
| HA | CT3 | C | O | 0.000 | 3 | 180.0 |
| HA | CT3 | C | NH1 | 0.000 | 3 | 0.00 |
|  |  |  |  |  |  |  |
| C | NH1 | CT2 | HA | 0.00 | 3 | 0.00 |
|  |  |  |  |  |  |  |
| C | NH1 | CT2 | CT2 | 0.5 | 2 | 0.00 |
| O | C | NH1 | H | 2.5 | 2 | 180.0 |
| O | C | NH1 | CT2 | 2.5 | 2 | 180.0 |


| NH1 | CT2 | CT2 | HA | 0.1950 | 3 |  | 0.00 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NH1 | CT2 | CT2 | CT2 | 0.1950 | 3 |  | 0.00 |
| H N | NH1 C | CT2 HA |  | 0.00 | 3 |  | 0.00 |
| H N | NH1 C | CT2 | CT2 | 0.00 | 1 |  | 0.00 |
| CT2 | CT2 | CT2 | HA | 0.1950 | 3 |  | 0.00 |
| CT2 | CT2 | CC | 0 | 0.050 | 6 |  | 0.00 |
| CT2 | CT2 | CC | NH2 | 0.020 | 6 |  | 180.00 |
| ! СT2 | 2 CT 2 | 2 CC | NH2 | 0.00 | 3 |  | 180.0 |
| HA | CT2 | CT2 | HA | 0.1950 |  | 3 | 0.00 |
| HA | CT2 | CT2 | CC | 0.1950 |  | 3 | 0.00 |
| CT2 | CC | NH2 | 2 H | 1.4 |  | 2 | 180.0 |
| HA | CT2 | CC | 0 | 0.00 |  | 3 | 0.00 |
| HA | CT2 | CC | NH2 | 0.000 |  | 3 | 0.00 |
| 0 | CC | NH2 | H | 1.4 |  | 2 | 180.0 |
| CT2 | CT2 | CT2 | CC | 0.1950 |  | 3 | 0.00 |
| IMPROPER |  |  |  |  |  |  |  |
| NH1 | C | CT2 | H | 20.00 | 0 |  | 0.00 |
| C | CT3 N | NH1 O | $\bigcirc$ | 120.0 | 0 |  | 0.00 |
| CC | CT2 | NH2 | 0 | 45.0 | 0 |  | 0.00 |

Topology for 6 (added to topology for DNA/proteins, top_all27_prot_na.inp)

| RESI BETA GROUP |  | 0.00 !Beta alanine linker for polyamides |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | I |  |  |  |  |  |  |
| ATOM CL | CT3 | -0.27! | HL2 | H | HR1 | HS1 |  | H1 |
| ATOM HL1 | HA | $0.09!$ | $\backslash$ | I | \| | 1 |  | / |
| ATOM HL2 | HA | 0.09! | HL1-CL | ---N- | R- | CS- | CT |  |
| ATOM HL3 | HA | 0.09! | / | \| | | \| | \| | \| | | $\backslash$ |
|  |  | ! | HL3 | 0 | HR3 | HS2 | OL | H2 |
| GROUP |  | ! |  |  |  |  |  |  |
| ATOM C | C | 0.51! |  |  |  |  |  |  |
| ATOM O | 0 | -0.51! |  |  |  |  |  |  |
| ATOM N | NH1 | -0.47! |  |  |  |  |  |  |
| ATOM H | H | 0.47! |  |  |  |  |  |  |
| ATOM CR | CT2 | -0.18! |  |  |  |  |  |  |
| ATOM HR1 | HA | 0.09! |  |  |  |  |  |  |
| ATOM HR3 | HA | 0.09 ! |  |  |  |  |  |  |
| GROUP |  |  |  |  |  |  |  |  |
| ATOM CS | CT2 | -0.18 |  |  |  |  |  |  |
| ATOM HS1 | HA | 0.09 |  |  |  |  |  |  |
| ATOM HS2 | HA | 0.09 |  |  |  |  |  |  |
| GROUP |  |  |  |  |  |  |  |  |
| ATOM CT | CC | 0.51 |  |  |  |  |  |  |
| ATOM OL | $\bigcirc$ | -0.51 |  |  |  |  |  |  |
| ATOM NS | NH2 | -0.47 |  |  |  |  |  |  |
| ATOM H1 | H | 0.235 |  |  |  |  |  |  |
| ATOM H2 | H | 0.235 |  |  |  |  |  |  |
| BOND HL1 CL HL2 CL HL3 CL |  |  |  |  |  |  |  |  |
| BOND CL | C | N CR |  |  |  |  |  |  |


| BOND N H |  |  |  |
| :---: | :---: | :---: | :---: |
| BOND HR1 CR HR3 | CR |  |  |
| DOUBLE C O CT OL |  |  |  |
| BOND CR CS CS CT CT NS |  |  |  |
| BOND NS H1 NS H2 CS HS1 CS HS2 |  |  |  |
| IMPR N C CR H |  |  |  |
| IMPR C CL N O |  |  |  |
| IC O CNH | 1.2233122 .84 | 180. 119.23 0.99 |  |
| IC H N C CL | 0.9933119 .23 | 0.116 .251 .51 |  |
| IC O C N CR | 1.2233122 .84 | 0.122 .571 .44 |  |
| IC N C CL HLI | 1.3418116 .25 | 150. 109.3 1.10 |  |
| IC N C CL HL2 | 1.3418116 .25 | 25. 109.3 1.1 |  |
| IC N C CL HL3 | 1.3418116 .25 | -95. 109.3 1.10 |  |
| IC C N CR CS | 1.3418122 .57 | 90. 110.71 .5 |  |
| IC C N CR HR1 | 1.3418122 .57 | -150. 110.7 |  |
| IC C N CR hR3 | 1.3418122 .57 | -30.0 110.7 |  |
| IC N CR CS HS1 | 1.4488110 .7 | 60.0109 .5 | 1.10 |
| IC N CR CS HS2 | 1.4488110 .7 | -60.0 109.5 | 1.10 |
| IC N CR CS CT | 1.4488110 .7 | -180.0 109.5 | 1.50 |
| IC CR CS CT NS | 1.50109 .5 | 150.0115 .0 | 1.40 |
| IC CR CS CT OL | 1.50109 .5 | -36.50 123.1 | 1.24 |
| IC OL CT NS H1 | 1.24121 .9 | 160.0121 .2 | 0.9933 |
| IC OL CT NS H2 | 1.24121 .9 | 0.00121 .2 | 0.9933 |
| IMPR CT CS NS OL |  |  |  |
| IMPR CT NS CS OL |  |  |  |
| IMPR NS CT H2 H1 |  |  |  |
| IMPR NS CT H1 H2 |  |  |  |

Parameters for 6 (added to parameters for DNA/proteins, par_all27_prot_na.inp) ANGLES !beta linker
H NH1 H 40.000 118.2 !A.E.L.

DIHEDRALS

| CT3 | C | NH1 | H | 2.5 | 2 | 180.0 |
| :--- | :--- | :--- | :--- | :--- | :--- | ---: |
| CT3 | C | NH1 | CT2 | 1.6 | 1 | 0.00 |
| CT3 | C | NH1 | CT2 | 2.5 | 2 | 180.0 |
|  |  |  |  |  |  |  |
| HA | CT3 | C | O | 0.00 | 3 | 180.0 |
| HA | CT3 | C | NH1 | 0.00 | 3 | 0.00 |
| C | NH1 | CT2 | HA | 0.00 | 3 | 0.00 |
|  |  |  |  |  |  |  |
| C | NH1 | CT2 | CT2 | 0.25 | 2 | 0.0 |
| C | NH1 | CT2 | CT2 | 0.3 | 1 | 0.0 |
|  |  |  |  |  |  |  |
| O | C | NH1 | H | 2.5 | 2 | 180.0 |
| O | C | NH1 | CT2 | 2.5 | 2 | 180.0 |
| NH1 | CT2 | CT2 | HA | 0.1950 | 3 | 0.00 |
| NH1 | CT2 | CT2 | CC | 0.1950 | 3 | 0.00 |


| H | NH1 | CT2 | HA | 0.00 | 3 | 0.00 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| H | NH1 | CT2 | CT2 | 0.000 | 1 | 0.00 |
|  |  |  |  |  |  |  |
| CT2 | CT2 | CC | O | 0.0 | 2 | 0.00 |
| CT2 | CT2 | CC | NH2 | 1.0 | 1 | 0.0 |
| CT2 | CT2 | CC | NH2 | 0.3 | 3 | 0.00 |
|  |  |  |  |  |  |  |
| !CT2 | CT2 | CC | NH2 | 1.0 | 1 | 0.0 |
|  |  |  |  |  |  |  |
| HA | CT2 | CT2 | HA | 0.1950 | 3 | 0.00 |
| HA | CT2 | CT2 | CC | 0.1950 | 3 | 0.00 |
| CT2 | CC | NH2 | H | 2.5 | 2 | 180.0 |
| HA | CT2 | CC | O | 0.00 | 3 | 180.0 |
| HA | CT2 | CC | NH2 | 0.00 | 3 | 0.00 |
| O | CC | NH2 | H | 2.5 | 2 | 180.0 |

Topology and parameters for netropsin and distamycin (streamed in; used with parameters for DNA/proteins, par_all27_prot_na.inp)

```
RESI NETR 2.00 !netropsin
!1st pyrrole ring
GROUP
ATOM HG HP 0.115
ATOM CG CY -0.145
ATOM CD2 CPT -0.020 !
ATOM CD1 CA 0.035
ATOM HD1 HP 0.115 !
ATOM NE1 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE2 CPT 0.130 !
GROUP
ATOM CH CT3 -0.27 !
ATOM H1 HA 0.09
ATOM H2 HA 0.09
ATOM H3 HA 0.09
GROUP
ATOM N2 NH1 -0.47
ATOM HM H 0.235
ATOM CB C 0.745
ATOM O2 O -0.51
!---------------------------------------------------
!2nd pyrrole ring
GROUP
ATOM HG1 HP 0.115 !
ATOM CG1 CY -0.145 !
ATOM CD3 CPT -0.020 !
ATOM CD4 CA 0.035 !
ATOM HD4 HP 0.115 !
ATOM NE2 NY -0.230 !
ATOM CE3 CPT 0.130 !
GROUP
ATOM CI CT3 -0.27
ATOM HI1 HA 0.09
ATOM HI2 HA 0.09
ATOM HI3 HA 0.09
```

| GROUP |  |  |
| :---: | :---: | :---: |
| ATOM NK | NH1 | -0.47 |
| ATOM HN | H | 0.235 |
| ATOM CK | C | 0.745 |
| ATOM OK | 0 | -0.51 |
| GROUP |  |  |
| ATOM NZ | NH1 | -0.47 |
| ATOM HZ | H | 0.235 |
| ATOM CZ | C | 0.745 |
| ATOM OZ | 0 | -0.51 |
| ! ATOM HZ | H | 0.235 |
| GROUP |  |  |
| ATOM CV | CT2 | -0.18 |
| ATOM HV1 | HA | 0.09 |
| ! ATOM HR2 | HA | 0.09 |
| ATOM HV3 | HA | 0.09 |
| GROUP |  |  |
| ATOM CC | CT2 | -0.18 |
| ATOM HC1 | HA | 0.09 |
| ATOM HC2 | HA | 0.09 |
| GROUP |  |  |
| ATOM CZ2 | C | 0.64 |
| ATOM NH1 | NC2 | -0.74 |
| ATOM HH11 | HC | 0.46 |
| ATOM HH12 | HC | 0.46 |
| ATOM NH2 | NC2 | -0.74 |
| ATOM HH22 | HC | 0.46 |
| ATOM HH21 | HC | 0.46 |
| GROUP |  |  |
| ATOM CC2 | CT2 | -0.18 |
| ATOM HC12 | HA | 0.09 |
| ATOM HC14 | HA | 0.09 |
| ATOM NC | NC2 | -0.47 |
| ATOM HC | HC | 0.47 |
| ATOM CZ3 | C | 0.64 |
| ATOM NH4 | NC2 | -0.74 |
| ATOM HH14 | HC | 0.46 |
| ATOM HH16 | HC | 0.46 |
| ATOM NH5 | NC2 | -0.74 |
| ATOM HH25 | HC | 0.46 |
| ATOM HH27 | HC | 0.46 |
| !bonds for 1st pyrrole ring |  |  |
| BONDBONDNE1lCE2 |  |  |
|  |  |  |
| BOND CD1 HD1 |  |  |
| DOUBLE CD1 CG CE2 CD2 |  |  |
| BOND CH NE1 |  |  |
| BOND CH H1 | CH | 2 CH H3 |
| BOND CE2 CK |  |  |
| DOUBLE CB O2 |  |  |
| BOND N2 CG N2 HM N2 CB |  |  |

```
!bonds for 2nd pyrrole ring
BOND CD3 HG1 CD3 CG1 NE2 CD4
BOND NE2 CE3
BOND CD4 HD4
DOUBLE CD4 CG1 CE3 CD3
BOND CI NE2
BOND CI HI1 CI HI2 CI HI3
BOND CZ CE3
DOUBLE CK OK
BOND NK CG1 NK HN NK CK
!----------------------------------------------
!1st tail amidino group
DOUBLE CZ OZ
BOND HV1 CV HV3 CV CV NZ CZ NZ
BOND CC CV CC HC1 CC HC2 NZ HZ
BOND CC CZ2 CZ2 NH2
DOUBLE CZ2 NH1
BOND NH1 HH12 NH1 HH11
BOND NH2 HH22 NH2 HH21
!----------------------------------------------
!2nd tail guanidino group
BOND CB CC2 CC2 HC12 CC2 HC14
BOND CC2 NC NC HC
BOND NC CZ3 CZ3 NH5
DOUBLE CZ3 NH4
BOND NH4 HH14 NH4 HH16
BOND NH5 HH25 NH5 HH27
\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|c|}
\hline \[
\begin{aligned}
& \text { !ic } \\
& \text { IC }
\end{aligned}
\] & tab & & NC & CC2 & CB & 1.47 & 105.82 & -180.0 & 109.5 & 1.54 \\
\hline ic & HM & N2 & CB & CC2 & 0.9933 & 119.23 & 0 . & 116.25 & 1.47 & \\
\hline ic & CG & N2 & CB & CC2 & 1.40 & 119.23 & -180.0 & 116.25 & 1.47 & \\
\hline ic & N2 & CB & CC2 & NC & 1.3418 & 116.25 & -180. & 109.5 & 1.47 & \\
\hline ic & CG & N2 & CB & O2 & 1.38 & 122.0 & 0.00 & 123.5 & 1.2233 & \\
\hline IC & CD2 & CG & N2 & CB & 1.40 & 126.4 & -180.0 & 130.3 & 1.46 & \\
\hline ic & CD1 & CG & N2 & CB & 1.37 & 126.4 & 0.0 & 122.1 & 1.46 & \\
\hline ic & N2 & CB & CC2 & HC12 & 1.36 & 115.0 & -60.0 & 109.5 & 1.07 & \\
\hline ic & N2 & CB & CC2 & HC14 & 1.36 & 115.0 & 60.0 & 109.5 & 1.07 & \\
\hline ic & 02 & CB & CC2 & HC12 & 1.22 & 123.0 & -114.0 & 109.5 & 1.07 & \\
\hline ic & 02 & CB & CC2 & HC14 & 1.22 & 123.0 & 115.0 & 109.5 & 1.07 & \\
\hline
\end{tabular}
!from arginine
\begin{tabular}{llllllllll} 
IC & NC & CB & *CC2 & HC14 & 1.5034 & 107.0900 & 120.6900 & 109.4100 & 1.1143 \\
IC & NC & CB & *CC2 & HC12 & 1.5034 & 107.0900 & -119.0400 & 111.5200 & 1.1150 \\
IC & CB & CC2 & NC & CZ3 & 1.5384 & 107.0900 & -180.0000 & 123.0500 & 1.3401 \\
IC & CZ3 & CC2 & *NC & HC & 1.3401 & 123.0500 & -180.0000 & 113.1400 & 1.0065 \\
IC & CC2 & NC & CZ3 & NH4 & 1.5034 & 123.0500 & 0.0000 & 118.0600 & 1.3311 \\
IC & NC & CZ3 & NH4 & HH14 & 1.3401 & 118.0600 & 10.2800 & 120.6100 & 0.9903 \\
IC & HH14 & CZ3 & *NH4 & HH16 & 0.9903 & 120.6100 & 180.1900 & 116.2900 & 1.0023 \\
IC & NH4 & NC & *CZ3 & NH5 & 1.3311 & 118.0600 & 160.6400 & 122.1400 & 1.3292 \\
IC & NC & CZ3 & NH5 & HH25 & 1.3401 & 122.1400 & 170.1400 & 119.9100 & 0.9899 \\
IC & HH25 & CZ3 & *NH5 & HH27 & 0.9899 & 119.9100 & 70.1600 & 116.8800 & 0.9914 \\
IC & CG & CD1 & NE1 & CE2 & 0.0000 & 0.0000 & 0.0000 & 0.0000 & 0.0000 \\
IC & CD2 & CB & *CG & CD1 & 0.0000 & 0.0000 & 180.0000 & 0.0000 & 0.0000
\end{tabular}
```

| IC | CD1 | CG | CD2 | CE2 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | CD2 | CG | CD1 | NE1 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| IC | CD1 | CE2 | *NE11 | CH | 0.0000 | 0.0000 |  |  |
| IC | CG | NE1 | *CD1 | HD1 | 0.0000 | 0.0000 | 180.0000 | 0.0000 |
| IC | CE2 | CG | CD2 | HG | 0.0000 | 0.0000 | -180.0000 | 126.9 |

!ic's for methyl group

| IC | CD1 | NE1 | CH H1 | H1 | 1.37 | 123.2 | 0.00 | 109.5 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | CD1 | NE1 | CH H | H2 | 1.37 | 123.2 | -120.0 | 109.5 |
| IC | CD1 | ND1 | CH H3 | H3 | 1.08 |  |  |  |
| iC | CE2 | NE1 | CH | H3 | 1.37 | 123.2 | 120. | 109.5 |


| ic | HN | NK | CK | CE2 | 0.9933 | 119.23 | 0. | 116.25 | 1.47 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :---: | :--- | :--- |
| ic | CG1 | NK | CK | CE2 | 1.40 | 119.23 | -180.0 | 116.25 | 1.47 |
| ic | NK | CK | CE2 | NE1 | 1.3418 | 116.25 | 180. | 122.1 | 1.38 |
| ic | NK | CK | CE2 | CD2 | 1.3418 | 116.25 | 0. | 130.3 | 1.40 |
| ic | NE1 | CE2 | CK | OK | 1.38 | 122.0 | 0.00 | 123.5 | 1.2233 |
| ic | CD2 | CE2 | CK | OK | 1.40 | 130.3 | 180.0 | 123.5 | 1.2233 |
| IC | CG | CD2 | CE2 | CK | 1.40 | 107.1 | -180.0 | 130.3 | 1.46 |
| ic | CD1 | NE1 | CE2 | CK | 1.37 | 109.8 | 180.0 | 122.1 | 1.46 |
| ic | CH | NE1 | CE2 | CK | 1.46 | 127.0 | 0.00 | 122.1 | 1.46 |

!ic's for amide group on CG

| ic | O2 | CB | N2 | HM | 1.2233 | 122.84 | 180. | 119.23 | 0.9933 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| ic | HM | N2 | CB | O2 | 0.9933 | 119.23 | 180. | 112.25 | 1.23 |
| ic | O2 | CB | N2 | CG | 1.2233 | 122.84 | 0. | 125.57 | 1.40 |
| ic | CB | N2 | CG | CD2 | 1.3418 | 122.57 | 180. | 124.9 | 1.41 |
| ic | CB | N2 | CG | CD1 | 1.3418 | 122.57 | 0. | 127.2 | 1.40 |
| ic | N2 | CG | CD2 | CE2 | 1.40 | 124.9 | 180. | 107.1 | 1.40 |
| ic | N2 | CG | CD1 | NE1 | 1.40 | 127.2 | -180.0 | 107.6 | 1.37 |
| !------------------------------------------ |  |  |  |  |  |  |  |  |  |
| !ic's for 2nd pyrrole ring |  |  |  |  |  |  |  |  |  |
| IC | CG1 | CD4 | NE2 | CE3 | 0.0000 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| IC | CD4 | CG1 | CD3 | CE3 | 0.0000 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| IC | CD3 | CG1 | CD4 | NE2 | 0.0000 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| IC | CD4 | CE3 | *NE2 | CI | 0.0000 | 0.0000 | -180.0000 | 127.0 | 1.46 |
| IC | CG1 | NE2 | *CD4 | HD4 | 0.0000 | 0.0000 | 180.0000 | 0.0000 | 0.0000 |
| IC | CE3 | CG1 | CD3 | HG1 | 0.0000 | 0.0000 | -180.0000 | 126.9 | 1.08 |

!ic's for methyl group

| IC | CD4 | NE2 | CI HI1 | 1.37 | 123.2 | 0.00 | 109.5 | 1.08 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | CD4 | NE2 | CI | HI2 | 1.37 | 123.2 | -120.0 | 109.5 |
| IC | CD4 | NC1 | CI | HI3 | 1.37 | 123.2 | 120. | 109.5 |
| iC | CE3 | NE2 | CI | HI3 | 1.37 | 123.2 | -60.0 | 109.5 |

ic HZ NZ CZ CE3 $0.9933119 .23 \quad 0.116 .25 \quad 1.47$
ic CG2 NZ CZ CE3 1.40 119.23 -180.0 116.25 1.47
ic NZ CZ CE3 NE2 1.3418 116.25 180. 122.1 1.38
ic NZ CZ CE3 CD3 1.3418 116.25 0. 130.3 1.40
ic NE2 CE3 CZ OZ $1.38 \quad 122.0 \quad 0.00 \quad 123.5 \quad 1.2233$
ic CD3 CE3 CZ OZ 1.40 130.3 180.0 123.5 1.2233
IC CG1 CD3 CE3 CZ 1.40 107.1 -180.0 130.31 .46
ic CD4 NE2 CE3 CZ $1.37 \quad 109.8 \quad 180.0122 .1 \quad 1.46$
ic CI NE2 CE3 CZ $1.46 \quad 127.0 \quad 0.00 \quad 122.1 \quad 1.46$
ic OK CK NK HN 1.2233 122.84 -180. 119.23 0.9933
ic HN NK CK CE2 0.9933 119.23 0. 112.25 1.47 !!


```
ATOM HG1 HP 0.115 !
ATOM CG1 CY -0.145 !
ATOM CD3 CPT -0.020 !
ATOM CD4 CA 0.035 !
ATOM HD4 HP 0.115 !
ATOM NE2 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE3 CPT 0.130 !
GROUP
ATOM CI CT3 -0.27 !
ATOM HI1 HA 0.09
ATOM HI2 HA 0.09
ATOM HI3 HA 0.09
GROUP
ATOM NZ NH1 -0.47
ATOM HZ H 0.235
ATOM CZ C 0.745
ATOM OZ O -0.51
!3rd pyrrole ring
GROUP
ATOM HG2 HP 0.115 !
ATOM CG2 CY -0.145 !
ATOM CD5 CPT -0.020 !
ATOM CD6 CA 0.035 !
ATOM HD6 HP 0.115 !
ATOM NE3 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE4 CPT 0.130 !
GROUP
ATOM CF CT3 -0.27 !
ATOM HF1 HA 0.09
ATOM HF2 HA 0.09
ATOM HF3 HA 0.09
ATOM CP C 0.51
ATOM OP O -0.51
ATOM NP NH1 -0.47
ATOM HP H 0.47
ATOM CR CT2 -0.18
ATOM HR1 HA 0.09
ATOM HR3 HA 0.09
!GROUP
ATOM CS CT2 -0.18
ATOM HS1 HA 0.09
ATOM HS2 HA 0.09
GROUP
ATOM CZ2 C 0.64
ATOM NH1 NC2 -0.74
ATOM HH11 HC 0.46
ATOM HH12 HC 0.46
ATOM NH2 NC2 -0.74
ATOM HH22 HC 0.46
ATOM HH21 HC 0.46
!bonds
BOND CB N2 N2 HM CB H N2 CG
DOUBLE CB O2
!bonds for lst pyrrole ring
```

```
BOND CD2 HG CD2 CG NE1 CD1
BOND NE1 CE2
BOND CD1 HD1
DOUBLE CD1 CG CE2 CD2
BOND CH NE1
BOND CH H1 CH H2 CH H3
BOND CE2 CK
!bonds for 2nd pyrrole ring
BOND CD3 HG1 CD3 CG1 NE2 CD4
BOND NE2 CE3
BOND CD4 HD4
DOUBLE CD4 CG1 CE3 CD3
BOND CI NE2
BOND CI HI1 CI HI2 CI HI3
BOND CZ CE3
DOUBLE CK OK
BOND NK CG1 NK HN NK CK
!bonds for 3rd pyrrole ring
BOND CD5 HG2 CD5 CG2 NE3 CD6
BOND NE3 CE4
BOND CD6 HD6
DOUBLE CD6 CG2 CE4 CD5
BOND CF NE3
BOND CF HF1 CF HF2 CF HF3
BOND CP CE4
DOUBLE CZ OZ
BOND NZ CG2 NZ HZ NZ CZ
!bonds for tail
BOND CP NP NP CR NP HP
BOND CR HR1 CR HR3 CR CS
DOUBLE CP OP
BOND CS HS1 CS HS2
BOND CS CZ2 CZ2 NH2
DOUBLE CZ2 NH1
BOND NH1 HH12 NH1 HH11
BOND NH2 HH22 NH2 HH21
!ic table
IMPR CB H N2 O2 CB N2 H O2
IMPR N2 CB CG HM N2 CB HM CG
IC O2 CB N2 CG 1.2012 124.70 -18.0 118.63 1.40 !
IC O2 CB N2 HM 1.2012 124.70 -162. 120.92 0.9930 !
IC CG N2 CB H 1.40 118.63 163. 115.65 1.34 !
IC HM N2 CB H N 1.0015 118.63 19.7 115.54 1.34
!ic's for 1st pyrrole ring
IC CG CD1 NE1 CE2 0.0000 0.0000 0.0000 0.0000 0.0000
IC CD2 CB *CG CD1 0.0000 0.0000 180.0000 0.0000 0.0000
IC CD1 CG CD2 CE2 0.0000 0.0000 0.0000 0.0000 0.0000
IC CD2 CG CD1 NE1 0.0000 0.0000 0.0000 0.0000 0.0000
IC CD1 CE2 *NE1 CH 0.0000 0.0000 -180.0000 127.0 1.46
IC CG NE1 *CD1 HD1 0.0000 0.0000 180.0000 0.0000 0.0000
IC CE2 CG CD2 HG 0.0000 0.0000 -180.0000 126.9 1.08
```

| !ic's for methyl group |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC CD1 NE1 CH H1 | 1.37 | 123.2 | 0.00 | 109.5 | 1.08 |
| IC CD1 NE1 CH H2 | 1.37 | 123.2 | -120.0 | 109.5 | 1.08 |
| IC CD1 ND1 CH H3 | 1.37 | 123.2 | 120. | 109.5 | 1.08 |
| iC CE2 NE1 CH H3 | 1.37 | 123.2 | -60.0 | 109.5 | 1.08 |

!ic's for amide group on CE2

| !ic | O1 | CA | N1 | HA |  | 1.2233 | 122.84 | 180. | 119.23 | 0.9933 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| !ic | O1 | CA | N1 | HB |  | 1.2233 | 122.84 | 0.00 | 119.23 | 0.9933 |
| ic | HN | NK | CK | CE2 | 0.9933 |  |  |  | 119.23 | 0. |
| ic | CG1 | NK | CK | CE2 |  | 116.25 | 1.47 |  |  |  |
| ic | NK | CK | CE2 | NE1 | 119.23 | 1.3418 | 116.25 | 180.0 | 116.25 | 1.47 |
| ic | NK | CK | CE2 | CD2 | 1.3418 | 116.25 | 0. | 122.1 | 1.38 |  |
| ic | NE1 | CE2 | CK | OK | 1.38 | 122.0 | 0.00 | 123.5 | 1.40 |  |
| ic | CD2 | CE2 | CK | OK | 1.40 | 130.3 | 180.0 | 123.5 | 1.22333 |  |
| IC | CG | CD2 | CE2 | CK | 1.40 | 107.1 | -180.0 | 130.3 | 1.46 |  |
| ic | CD1 | NE1 | CE2 | CK | 1.37 | 109.8 | 180.0 | 122.1 | 1.46 |  |
| ic | CH | NE1 | CE2 | CK | 1.46 | 127.0 | 0.00 | 122.1 | 1.46 |  |

!ic's for amide group on CG

| ic | O2 | CB | N2 | HM | 1.2233 | 122.84 | -180. | 119.23 | 0.9933 |
| :--- | :--- | :--- | :--- | :--- | :---: | ---: | :---: | :---: | :---: |
| ic | HM | N2 | CB | H | 0.9933 | 119.23 | 0. | 112.25 | 1.10 |
| ic | O2 | CB | N2 | CG | 1.2233 | 122.84 | 0. | 125.57 | 1.40 |
| ic | CB | N2 | CG | CD2 | 1.3418 | 122.57 | 180. | 124.9 | 1.41 |
| ic | CB | N2 | CG | CD1 | 1.3418 | 122.57 | 0. | 127.2 | 1.40 |
| IC | N2 | CG | CD2 | CE2 | 1.40 | 124.9 | 180. | 107.1 | 1.40 |
| ic | N2 | CG | CD1 | NE1 | 1.40 | 127.2 | -180.0 | 107.6 | 1.37 |


| !ic's for 2nd pyrrole ring |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC CG1 | CD4 | NE2 | CE3 | 0.0000 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| ! IC CD2 | CB | * CG | CD1 | 0.0000 | 0.0000 | 180.0000 | 0.0000 | 0.0000 |
| IC CD4 | CG1 | CD3 | CE3 | 0.0000 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| IC CD3 | CG1 | CD4 | NE2 | 0.0000 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| IC CD4 | CE3 | *NE2 | CI | 0.0000 | $0.0000-1$ | 180.0000 | 127.0 | 1.46 |
| IC CG1 | NE2 | *CD4 | HD4 | 0.0000 | 0.0000 | 180.0000 | 0.0000 | 0.0000 |
| IC CE3 | CG1 | CD3 | HG1 | 0.0000 | 0.0 .0000 | -180.00 | 0126. | 91.08 |


| !ic's for methyl group |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | CD4 NE2 CI HI1 | 1.37 | 123.2 | 0.00 | 109.5 | 1.08 |
| IC | CD4 NE2 CI HI2 | 1.37 | 123.2 | -120.0 | 109.5 | 1.08 |
| IC | CD4 NC1 CI HI3 | 1.37 | 123.2 | 120. | 109.5 | 1.08 |
| ic | CE3 NE2 CI HI3 | 1.37 | 123.2 | -60.0 | 109.5 | 1.08 |

!ic's for amide group on CE3

| !ic | OJ | CJ | NA | HC | 1.2233 | 3122.84 | 180. | 119.23 | 0.9933 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| !ic | OJ | CJ | NA | HD | 1.2233 | 3122.84 | 0.00 | 119.23 | 0.9933 |
| ic | HZ | NZ | CZ | CE3 | 0.9933 | 119.23 | 0. | 116.25 | 1.47 |
| ic | CG2 | NZ | CZ | CE3 | 1.401 | 119.23-180 | 80.0 | 116.25 | 1.47 |
| ic | NZ | CZ | CE3 | NE2 | 1.3418 | 116.25 | 180. | 122.1 | 1.38 |
| ic | NZ | CZ | CE3 | CD3 | 1.3418 | 116.25 | 0 | 130.3 | 1.40 |
| ic | NE2 | CE3 | CZ | OZ | 1.38 | 122.0 | 0.00 | 123.5 | 1.2233 |
| ic | CD4 | CE3 | CZ | OZ | 1.40 | 130.3 | 180.0 | 123.5 | 1.2233 |
| IC | CG1 | CD3 | CE3 | CZ | 1.40 | 107.1 | -180.0 | 0130.3 | 1.46 |
| ic | CD4 | NE2 | CE3 | CZ | 1.3710 | 109.8180 | 80.0 | 122.11 | 1.46 |
| ic | CI | NE2 | CE3 | CZ | 1.46 | 127.0 | 0.00 | 122.1 | 1.46 |


| !ic's for amide group on CG1 |  |  |  |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| ic | OK | CK | NK | HN | 1.2233 | 122.84 | -180. | 119.23 | 0.9933 |
| ic | HN | NK | CK | CE2 | 0.9933 | 119.23 | 0. | 112.25 | 1.10 | !!


| !ic's for methyl group |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC CD6 NE3 CF HF1 | 1.37 | 123.2 | 0.00 | 109.5 | 1.08 |  |
| IC CD6 NE3 CF HF2 | 1.37 | 123.2 | -120.0 | 109.5 | 1.08 |  |
| IC | CD6 NF1 CF HF3 | 1.37 | 123.2 | 120. | 109.5 | 1.08 |
| ic CE4 NE3 CF HF3 | 1.37 | 123.2 | -60.0 | 109.5 | 1.08 |  |

!ic's for amide group on CE4

| !ic | OL | CL | NB | HE | 1.2233 | 122.84 | 180. | 119.23 | 0.9933 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| !ic | OL | CL | NB | HF | 1.2233 | 122.84 | 0.00 | 119.23 | 0.9933 |
| ic | HP | NP | CP | CE4 | 0.9933 | 119.23 | 0. | 116.251. | . 47 |
| ic | CR | NP | CP | CE4 | 1.40119 | $19.23-180$ | 80.0 | 116.251 .47 |  |
| ic | NP | CP | CE4 | NE3 | 1.3418 | 116.25 | 180 | 122.11 .38 |  |
| ic | NP | CP | CE 4 | CD5 | 1.3418 | 116.25 | 0 | 130.31 .40 |  |
| ic | NE3 | CE5 | CP | OP | 1.38 | 122.0 | 0.00 | 123.51 .2 | 2233 |
| ic | CD6 | CE4 | CP | OP | 1.40 | 130.3 | 180.0 | 123.51 .22 | 2233 |
| IC | CG2 | CD5 | CE4 | CP | 1.40 | 107.1 -1 | 180.0 | 130.31 .46 |  |
| ic | CD6 | NE3 | CE4 | CP | 1.3710 | 09.8 18 | 80.0 | 122.11 .46 |  |
| ic | CF | NE3 | CE4 | CP | 1.46 | 127.0 | 0.00 | 122.11 .4 |  |
| !ic's for amide group on CG2 |  |  |  |  |  |  |  |  |  |
| ic | OZ | CZ | NZ | HZ | 1.2233 | 122.84 | -180 | . 119.23 | 0.9933 |
| ic | HZ | NZ | CZ | CE3 | 0.9933 | 33119.23 | 0 | 112.25 | 1.10 ! ! |
| ic | OZ | CZ | NZ | CG2 | 1.2233 | 33 122.84 | 0 | 125.57 | 1.40 |
| ic | CZ | NZ | CG2 | CD5 | 1.3418 | 8122.57 | 180 | . 124.9 | 1.41 |
| ic | CZ | NZ | CG2 | CD6 | 1.3418 | 8122.57 | 0 | . 127.2 | 1.40 |
| IC | NZ | CG2 | CD5 | CE 4 | 1.40 | 124.9 | 180 | . 107.1 | 1.40 |
| ic | NZ | CG2 | CD6 | NE3 | 1.40 | 127.2 | -180 | . 0107.6 | . 37 |

!ic's for cationic tail
IMPR NP CP CR HP
IMPR CP CE4 NP OP

| IC | OP | CP | NP | HP | 1.2233 | 122.84 | -166. | 119.23 | 0.9933 |  |
| :--- | :--- | :---: | :--- | :--- | :---: | :--- | :---: | :--- | :--- | :--- |
| IC | HP | NP | CP | CE4 | 0.9933119 .23 | 14. | 116.25 | 1.40 |  |  |
| IC | OP | CP | NP | CR | 1.2233 | 122.84 | -15. | 117.57 | 1.4488 |  |
| IC | NP | CP | CE4 | CD5 | 1.3418 | 116.25 | 0. | 130.0 | 1.39 |  |
| IC | NP | CP | CE4 | NE3 | 1.3418 | 116.25 | 180. | 121.3 | 1.39 |  |
| !IC | NP | CP | CE4 4 | HL3 | 1.3418 | 116.25 | -87. | 109.3 | 1.109 |  |
| IC | CP | NP | CR | HR1 | 1.3418117 .00 | 60. | 110.7 | 1.113 | $1-180$. |  |







| CA | CD | NH2 | CT2 | 1.6 | 1 | 0.00 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CA | $C D$ | NH2 | CT2 | 2.5 | 2 | 180.0 |  |
| $C D$ | NH2 | CT2 | HA | 0.2000 | 1 | 60.0 |  |
| $C D$ | NH2 | CT2 | CT2 | 0.200 | 1 | 180.0 |  |
| OB | CD | NH2 | CT2 | 1.5 | 2 | 0.0 |  |
| CY | CA | $C D$ | NH2 | 2.32 |  | -180.00 !A.E.L. |  |
| NY | CA | $C D$ | NH2 | 2.32 |  | -180.0 ! A |  |
| CY | CA | $C D$ | OB | 2.32 |  | 180.00 ! A |  |
| NY | CA | $C D$ | OB | 2.32 |  | 180.00 !A.E.L. |  |
| !dihedrals for ipy drug !N-methyl-pyrrole |  |  |  |  |  |  |  |
| CA | NY | CT3 | HA | 0.00 |  | 10.00 | !methyl group |
| CA | NY | CT3 | HA | 0.000 |  | $3-120.0$ |  |
| CA | NY | CT3 | HA | 0.000 |  | $2-60.0$ |  |
| CPT | NY | CT3 | HA | 0.055 | 3 | 0.0 | !methyl group |
| CPT | NY | CT3 | HA | 0.00 | 1 | -120. |  |
| CPT | NY | CT3 | HA | 0.00 | 2 | -60.0 |  |
| HP | CPT | CY | CA | 2.0 | 1 | 180.0 |  |
| HP | CPT | CY | NH1 | 2.8 | 2 | 180.0 |  |
| CY | CA | NY | CT3 | 0.8 | 2 | 180.0 |  |
| CY | NH1 | C | 0 | 1.5 | 2 | 180.0 |  |
| CY | NH1 | C | H | 1.20 | 1 | 180.0 |  |
| CPT | CY | NH1 | H | 0.000 | 1 | 0.00 |  |
| CPT | CY | NH1 | C | 0.200 | 1 | 180.0 |  |
| CPT | CPT | NY | CT3 | 0.8 | 2 | 180.0 |  |
| CPT | CPT | C | 0 | 2.5 | 2 | 180.0 |  |
| CPT | CPT | C | NH1 | 1.5 | 2 | 180.0 |  |
| CA | NY | CPT | C | 2.00 | 1 | 180.0 |  |
| HP | CA | CY | NH1 | 1.00 | 2 | 180.0 |  |
| CA | CY | NH1 | H | 0.0 | 1 | 0.00 |  |
| CA | CY | NH1 | C | 2.5 | 2 | 180.0 |  |
| CA | CY | NH1 | C | 3.0 | 1 | 180.0 |  |
| HP | CA | NY | CT3 | 0.4000 | 2 | 180.0 |  |
| NY | CA | CY | NH1 | 2.0 | 2 | 180.0 |  |



| CPH1 | CPH1 | 1 NR1 | CT3 | 0.2000 | 1 |  | -180.0 |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\bigcirc \mathrm{C}$ | NH1 | 1 H |  | 2.5000 | 2 | 2 | 180.0 |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |
| CPH1 | NR1 | CPH2 | C | 0.2000 | 1 | 1 | -180.0 | ! |  |  |
| HR3 | CPH1 | 1 NR1 | CT3 | 0.2000 | 1 | , | 0.00 | ! |  |  |
| CPH1 | NR2 | CPH2 | C | 0.2000 | 1 | 1 | -180.0 |  |  |  |
| NR1 | CPH2 | 2 C |  | 5.000 |  | 2 | 180.0 | $!5.5$ ! | ! 5.0 |  |
| NR1 | CPH2 | 2 C |  | 4.500 |  | 1 | 180.0 | $!3.00$ |  |  |
| NR1 | CPH2 | 2 C N |  | 1.000 |  | 2 | 0.0 ! |  |  |  |
| NR2 | CPH2 | C 0 |  | 1.0000 |  | 1 | 0.0 | ! |  |  |
| NR2 | CPH2 | C NH1 |  | 0.0000 |  | 2 | 0.00 | ! |  |  |
| CPH2 | C | NH1 H |  | 0.2000 |  | 1 | 0.00 |  |  |  |
| CPH2 | C | NH1 H |  | 0.2000 |  | 2 | 180.0 |  |  |  |
| NR2 | CPH2 | NR1 C |  | 0.2000 |  | 1 | 180.0 |  |  |  |
| CT3 | NR1 | CPH2 |  | 0.2000 |  | 1 | 0.00 |  |  |  |
| CPH2 | C N | NH1 CY |  | 0.2000 | 1 | 1 | 0.00 |  |  |  |
| CPH2 | C N | NH1 CY |  | 0.2000 | 2 | 2 | 180.0 |  |  |  |
| ! netropsin |  |  |  |  |  |  |  |  |  |  |
| ! CY | NH1 | C |  | 0.2000 |  | 1 | 0.00 |  |  |  |
| CY | NH1 C | C CT2 |  | 1.000 | 1 |  | -180.0 |  |  |  |
| $\bigcirc$ | C | CT2 NC2 |  | 0.5000 | 1 | 1 | 0.000 |  |  |  |
| CT2 | CT2 C | C NC2 |  | 0.000 | 1 | 1 | 0.00 |  |  |  |
| NH1 | C | CT2 NC2 |  | 0.6000 | 1 |  | 0.00 |  |  |  |
| HA | CT2 C | C NC2 |  | 0.0000 | 3 | 3 | 0.00 |  |  |  |
| !distamycin |  |  |  |  |  |  |  |  |  |  |
| CC | NH2 | CY CP |  | 0.2000 | 1 |  | 180.0 |  |  |  |
|  | NH2 | CY CA |  | 2.5 | 2 | 2 | 180.0 |  |  |  |
| CC | NH2 | CY CA |  | 3.7 | 1 | 1 | 180.0 |  |  |  |
| NH2 | CY | CPT HP |  | 2.8 | 2 |  | 180.0 |  |  |  |
| NH2 | CY | CA HP |  | 1.00 | 1 | 1 | 180.0 |  |  |  |
| H | NH2 | CC H |  | 1.4 | 2 |  | 180.0 |  |  |  |
| H | NH2 | CY CP |  | 0.050 | 1 |  | 0.00 |  |  |  |
| H | NH2 | CY CA |  | 0.05 | 1 | , | 0.00 |  |  |  |
| 0 | CC | NH2 CY |  | 1.50 | 2 | 2 | 180.0 |  |  |  |
| H | CC | NH2 CY |  | 1.20 | 1 |  | 180.0 |  |  |  |
| IMPROPER |  |  |  |  |  |  |  |  |  |  |
| !ato | m typ | pes |  | Kpsi |  |  | psi0 |  |  |  |
| !polyamide impropers |  |  |  |  |  |  |  |  |  |  |
| CA | CA | CA | NY | 90.000 |  |  | 0 | -180.00 | 0 ! A | A.E.L. |
| CPT | CA | NY | CA | 90.000 |  |  | 0 | 180.00 | 0 ! A | A.E.L. |
| CY | CPT | CPT | NH2 | 90.000 |  |  | 0 | 180.00 | 0 ! A | A.E.L. |
| CA | CA | CA | CT1 | 90.000 |  |  | 0 | 180.00 | 0 ! A | A.E.L. |
| OH1 | CA | CD | OB | 90.000 |  |  | 0 | 180.0 | ! A | A.E.L. |


| CD | NY | CA | CY | 90.000 | 0 | 180.0 | !A.E.L. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CT3 | NY | CPT | CPT | 90.000 | 0 | -180.0 | !A.E.L. |
| CA | CPT | NY | CT3 | 90.000 | 0 | -180.0 | !A.E.L. |
| CA | CPT | CY | NH2 | 90.000 | 0 | 180.00 | !A.E.L. |
| CPT | CY | CPT | NH2 | 90.000 | 0 | -180.0 | !A.E.L. |
| CPT | NY | CPT | F HP | 90.000 | 0 | 180.0 | !A.E.L. |
| OB | CA | $C D$ | OH1 | 90.00 | 0 | -180.0 | !A.E.L. |
| OH1 | OB | $C D$ | CA | 90.00 | 0 | -180.0 | !A.E.L. |
| OB | OH 1 | 1 CD | CA | 90.00 | 0 | 180.0 | !A.E.L. |
| CY | CD | OH1 | H | 90.000 | 0 | -180.0 | !A.E.L. |
| HP | CA | CD | OH1 | 90.00 | 0 | 0.00 | !A.E.L. |
| OH1 | CA | CY | HP | 90.000 | 0 | 0.00 | !A.E.L. |
| OB | CA | NY | CT3 | 90.000 | 0 | 0.00 | !A.E.L. |
| NH2 | CT2 | 2 CT 2 | 2 HA | 90.000 | 0 | -180.0 | !A.E.L. |
| OB | NH2 | 2 CT 2 | 2 CT 2 | 90.000 | 0 | 180.0 | !A.E.L. |
| CA | OB | $C D$ | OH1 | 90.000 | 0 | 180.0 | !A.E.L. |
| NY | $C D$ | OH1 | H | 90.000 | 0 | 180.0 | !A.E.L. |
| CY | $C D$ | CA | NY | 90.000 | 0 | -180.0 | !A.E.L. |
| H | $C D$ | CA | CA | 90.000 | 0 | 0.000 | 0!A.E.L. |
| H | $C D$ | CA | NY | 90.000 | 0 | -180.0 | O A.E.L. |
| CD | CY | CA | NY | 90.000 | 0 | 180.0 | 0 !A.E.L. |
| CD | CA | CPT | CPT | 90.000 | 0 | 0.00 | !A.E.L. |
| CA | CA | NY | CPT | 90.000 | 0 | -180.0 | 0 ! A.E.L. |
| CA | CA | CA | HP | 90.000 | 0 | 180.0 | O ! A.E.L. |
| CT1 | CA | CA | HP | 90.000 | 0 | -180.0 | O ! A.E.L. |
| HC | HC | NH2 C | CPT | 90.000 | 0 | -120 | O.0 ! A.E.L. |
| HC | CPT | NH2 H |  | 90.000 | 0 | 120 | 0.0 ! A.E.L. |
| CA | CPT | NH2 H |  | 90.000 | 0 |  | 20.33 !A.E.L. |
| NY | CD N | NH2 CP |  | 90.00 | 0 | 180.0 | 0 ! A.E.L. |
| NH2 | CD | CPT | HC | 20.000 | 0 | 0.00 | 0 !A.E.L. |
| CD | CA | NH2 | OB | 45.000 | 0 | 0.00 | 0 !A.E.L. |
| CA | CPT | T NH2 | HC | 90.000 | 0 | 0.00 | 0 !A.E.L. |
| CC | OS | NH2 | 0 | 45.000 | 0 | 0.00 ! | !A.E.L. |
| CC | NH2 | OS | 0 | 45.000 | 0 | 0.00 ! | !A.E.L. |
| OS | NH2 | CC | 0 | 90.000 | 0 | -180.0 ! | !A.E.L. |
| OS | 0 | CC | NH2 | 90.00 | 0 | 180.0 ! | !A.E.L. |
| CT3 | OS | CT1 | CT3 | 90.000 | 0 | -120.0! | !A.E.L. |
| NH2 | CC | CPT | HC | 20.00 | 0 | 0.00 ! | ! A.e.L. |
| NH2 | CD | CT2 | HC | 20.000 | 0 | 0.00 ! | !A.E.L. |
| CPT | CA | CY H | HP | 100.000 | 0 | 180.0 |  |
| NY | CPT | CPT H | HP | 100.000 | 0 | -180.0 |  |
| HP | CPT | CY C | CA | 100.000 | 0 | 180.0 |  |
| HP | CA | CY C | CPT | 100.000 | 0 | -180.0 |  |
| CD | CA | NH2 | OB | 45.000 | 0 | 0.00 |  |
| CD | NH2 | CA | OB | 45.000 | 0 | 0.000 |  |
| !impropers for ipy drug !Boc Binder |  |  |  |  |  |  |  |
| NH1 | C | H | CY | 4.00 | 0 | 0.00 |  |


| !imid |  |  |  |
| :---: | :---: | :---: | :---: |
| CT3 CPH1 NR1 CPH2 | 90.000 | 0 | 180.0 |
| NR1 CPH1 CPH2 CT3 | 90.000 | 0 | 0.00 |
| NR1 CPH2 CPH1 CT3 | 90.000 | 0 | 0.00 |
| CPH2 NR1 NR2 C | 90.000 | 0 | 0.00 |
| CPH2 NR2 NR1 C | 90.000 | 0 | 0.00 |
| O CPT C NH1 | 90.00 | 0 | -180.0 |
| CPH1 CPH2 NR1 CT3 | 90.000 | 0 | 180.0 |
| O CT2 C NH1 | 90.000 | 0 | -180.0 |
| O NH1 C CT2 | 90.000 | 0 | 180.0 |
| CY C NH1 H | 90.00 | 0 | 180.0 |
| C CA CY NH1 | 90.000 | 0 | 0.00 |
| H C NH1 CT2 | 90.000 | 0 | 180.0 |
| NC2 O C CT2 | 90.000 | 0 | 0.00 |
| C NC2 CT2 HA | 90.000 | 0 | 120.0 |
| NH1 C CT2 CT2 | 90.0 | 0 | 0.00 |
| C NC2 NC2 CT2 | 90.00 | 0 | 0.00 |
| !distamycin |  |  |  |
| NH2 CC H CY | 4.0 | 0 | 0.000 |
| NH1 CPT C O | 90.0 | 0 | 180.0 |
| O NH1 C CPT | 90.0 | 0 | 180.0 |
| CT2 C NH1 H | 90.0 | 0 | 180.0 |
| CC H NH2 O | 120.0 | 0 | 0.00 |
| CC NH2 H O | 120.0 | 0 | 0.00 |

NONBONDED nbxmod 5 atom cdiel shift vatom vdistance vswitch cutnb 14.0 ctofnb 12.0 ctonnb 10.0 eps 1.0 e14fac 1.0 wmin 1.5 !adm jr., 5/08/91, suggested cutoff scheme !
!atom ignored epsilon Rmin/2 ignored eps,1-4 Rmin/2,1-4 !
end

Topology and parameters for Dervan polyamides; IPY is the imidazole-pyrrole polyamide and HP2 is the hydroxypyrrole polyamide (streamed in; used with parameters for DNA/proteins, par_all27_prot_na.inp).


```
ATOM HI1 HA 0.09
ATOM HI2 HA 0.09
ATOM HI3 HA 0.09
GROUP
ATOM NK NH1 -0.47
ATOM HN H 0.235
ATOM CK C 0.745
ATOM OK O -0.51
!ATOM HZ H 0.235
!-----------------------------------------------
!3rd pyrrole ring
GROUP
ATOM HG2 HP 0.115 !
ATOM CG2 CY -0.145 !
ATOM CD5 CPT -0.020 !
ATOM CD6 CA 0.035 !
ATOM HD6 HP 0.115 !
ATOM NE3 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE4 CPT 0.130 !
GROUP
ATOM CF CT3 -0.27 !change C1 to CH
ATOM HF1 HA 0.09
ATOM HF2 HA 0.09
ATOM HF3 HA 0.09
GROUP
ATOM CL C 0.51
ATOM OL O -0.51
GROUP
ATOM NZ NH1 -0.47
ATOM HZ H 0.235
ATOM CZ C 0.745
ATOM OZ O -0.51
!ATOM HZ H 0.235
!---------------------------------------------
!beta-alanine
GROUP
ATOM NB NH1 -0.47
ATOM HB H 0.47
ATOM CV CT2 -0.18
ATOM HV1 HA 0.09
!ATOM HR2 HA 0.09
ATOM HV3 HA 0.09
GROUP
ATOM CC CT2 -0.18
ATOM HC1 HA 0.09
ATOM HC2 HA 0.09
GROUP
ATOM CT C 0.51
ATOM OT O -0.51
!ATOM NS NH1 -0.47
!ATOM HS1 H 0.235
!ATOM HS2 H 0.235
```

```
!-------------
GROUP
ATOM NC NH1 -0.47
ATOM HC H 0.47
ATOM CE CT2 -0.11
ATOM HM1 HA 0.09
ATOM HM3 HA 0.09
GROUP
ATOM CJ CT2 -0.11
ATOM HJ1 HA 0.09
AtOM HJ2 HA 0.09
ATOM CM CT2 -0.11
ATOM HN1 HA 0.09
ATOM HN2 HA 0.09
GROUP
ATOM NN NH3 -0.30
ATOM HH HC 0.33
ATOM CX CT3 0.23
ATOM HX1 HA 0.05
ATOM HX2 HA 0.05
ATOM HX3 HA 0.05
ATOM CY CT3 0.23
ATOM HY1 HA 0.05
ATOM HY2 HA 0.05
ATOM HY3 HA 0.05
!--------------------------------------------
!bonds for imidazole
BOND NZ2 CZ2 NZ1 CZ1 CZ3 NZ1
BOND CZ1 HE1 CZ2 HE2
DOUBLE NZ2 CZ3 CZ2 CZ1
BOND NZ1 CZ4 CZ4 HZ1 CZ4 HZ2 CZ4 HZ3
BOND CZ3 CB
!------------------------------------------
!bonds for 1st pyrrole ring
BOND CD2 HG CD2 CG NE1 CD1
BOND NE1 CE2
BOND CD1 HD1
DOUBLE CD1 CG CE2 CD2
BOND CH NE1
BOND CH H1 CH H2 CH H3
BOND CE2 CK
DOUBLE CB O2
BOND N2 CG N2 HM N2 CB
!-------------------------------------------------
!bonds for 2nd pyrrole ring
BOND CD3 HG1 CD3 CG1 NE2 CD4
BOND NE2 CE3
BOND CD4 HD4
DOUBLE CD4 CG1 CE3 CD3
BOND CI NE2
BOND CI HI1 CI HI2 CI HI3
BOND CZ CE3 !NA CJ NA HC NA HD
DOUBLE CK OK
```

BOND NK CG1 NK HN NK CK ! CB HZ

```
!----------------------------------------------
!bonds for 3rd pyrrole ring
BOND CD5 HG2 CD5 CG2 NE3 CD6
BOND NE3 CE4
BOND CD6 HD6
DOUBLE CD6 CG2 CE4 CD5
BOND CF NE3
BOND CF HF1 CF HF2 CF HF3
BOND CL CE4 CL NB !NB HE NB HF
DOUBLE CZ OZ CL OL
BOND NZ CG2 NZ HZ NZ CZ !CB HZ
!--------------------------------------------------
!bonds for beta-alanine
BOND NB CV
BOND NB HB
BOND HV1 CV HV3 CV
DOUBLE CT OT
BOND CV CC CC CT !CT NS
BOND CT NC !NS HS1 NS HS2
BOND CC HC1 CC HC2
!---------------------------------------------
!bonds for cationic tail
BOND NC HC NC CE
!BOND CL HL1 CL HL2 CL HL3
BOND CE HM1 CE HM3 CE CJ
!DOUBLE C O
BOND CJ HJ1 CJ HJ2 CJ CM
BOND CM HN1 CM HN2 CM NN
BOND NN HH NN CX NN CY
BOND CX HX1 CX HX2 CX HX3
BOND CY HY1 CY HY2 CY HY3
```

| !ic's for imidazole ring |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ! KEEPS HYDROGENS IN RING PLANE |  |  |  |  |  |  |  |  |
| ! IMPH NZ1 | CZ1 | CZ3 C | Z 4 | NZ1 CZ3 | 3 CZ1 | CZ4 |  |  |
| IMPH CZ2 | CZ1 | NZ2 | E2 | CZ2 NZ2 | $2 \mathrm{CZ1}$ | HE2 |  |  |
| IMPH CZ3 | NZ1 | NZ2 CB |  | CZ3 NZ2 | NZ1 CB |  |  |  |
| IMPH CZ1 | CZ2 | NZ1 H | E1 | CZ1 N | NZ1 CZ2 | HE1 |  |  |
| IC HE1 | CZ1 | NZ1 | CZ3 | 1.07 | 122.67 | -173.67 | 109.79 |  |
| 1.2987 |  |  |  |  |  |  |  |  |
| IC CZ1 | NZ1 | CZ3 | NZ2 | 1.2854 | 109.79 | 0.21 | 110.31 |  |
| 1.3071 |  |  |  |  |  |  |  |  |
| IC NZ1 | CZ3 | NZ2 | CZ2 | 1.2987 | 110.31 | 0.03 | 105.82 |  |
| 1.3165 |  |  |  |  |  |  |  |  |
| IC CZ3 | NZ2 | CZ2 | CZ1 | 1.3071 | 105.82 | -0.23 | 108.68 | 1.3758 |
| IC NZ2 | CZ2 | CZ1 | NZ1 | 1.3165 | 108.68 | 0.35 | 105.39 | 1.2854 |
| IC NZ2 | CZ2 | CZ1 | HE1 | 1.3165 | 108.68 | 172.86 | 131.52 | 1.07 |
| IC CZ2 | CZ1 | NZ1 | CZ3 | 1.3758 | 105.39 | -0.34 | 109.79 | 1.2987 |
| IC CZ2 | NZ2 | CZ3 | CB | 1.3165 | 105.82 | -180.0 | 124.7 | 1.46 |
| IC NZ2 | CZ3 | NZ1 | CZ 4 | 1.34 | 110.31 | 180.0 | 127.8 | 1.46 |
| IC HE1 | CZ1 | CZ2 | HE2 | 1.07 | 131.52 | 0.00 | 129.3 | 1.07 |

!ic's for methyl group
$\begin{array}{llllllll}\text { IC CZ1 NZ1 CZ4 HZ1 } & 1.37 & 124.7 & 0.00 & 109.5 & 1.08\end{array}$
$\begin{array}{lllllll}\text { IC CZ1 NZ1 CZ4 HZ2 } & 1.37 & 124.7 & -120.0 & 109.5 & 1.08\end{array}$



!ic's for amide group on CE4

| !!ic | OL | CL | NB | HE | 1.2233 | 122.84 | 180. | 119.23 | 0.9933 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| !!ic | OL | CL | NB | HF | 1.2233 | 122.84 | 0.00 | 119.23 | 0.9933 |




```
ATOM NE1 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE2 CPT 0.130 !
GROUP
ATOM CH CT3 -0.27
ATOM H1 HA 0.09
ATOM H2 HA 0.09
ATOM H3 HA 0.09
GROUP
ATOM N2 NH1 -0.47
ATOM HM H 0.235
ATOM CB C 0.745
ATOM O2 O -0.51
!ATOM HZ H 0.235
!---------------------------------------------
!2nd pyrrole ring--hydroxypyrrole
GROUP
!ATOM HG1 HP 0.115 !
ATOM OG1 OH1 -0.326
ATOM HG1 H 0.397
ATOM CG1 CY -0.140 !
ATOM CD3 CPT 0.020 !
ATOM CD4 CA 0.014 !
ATOM HD4 HP 0.115 !
ATOM NE2 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE3 CPT 0.150 !
GROUP
ATOM CI CT3 -0.27 !change C1 to CH
ATOM HI1 HA 0.09
ATOM HI2 HA 0.09
ATOM HI3 HA 0.09
GROUP
ATOM NK NH1 -0.47
ATOM HN H 0.235
ATOM CK C 0.745
ATOM OK O -0.51
!ATOM HZ H 0.235
!-----------------------------------------------
!3rd pyrrole ring
GROUP
ATOM HG2 HP 0.115 !
ATOM CG2 CY -0.145 !
ATOM CD5 CPT -0.020 !
ATOM CD6 CA 0.035 !
ATOM HD6 HP 0.115 !
ATOM NE3 NY -0.230 !
!ATOM HE1 H 0.380 !
ATOM CE4 CPT 0.130 !
GROUP
ATOM CF CT3 -0.27 !change C1 to CH
ATOM HF1 HA 0.09
ATOM HF2 HA 0.09
ATOM HF3 HA 0.09
```

| GROUP |  |  |
| :---: | :---: | :---: |
| ATOM CL | C | 0.51 |
| ATOM OL | 0 | -0.51 |
| ! ATOM NB | NH1 | -0.62 |
| ! ATOM HE | H | 0.31 |
| ! ATOM HF | H | 0.31 |
| GROUP |  |  |
| ATOM NZ | NH1 | -0.47 |
| ATOM HZ | H | 0.235 |
| ATOM CZ | C | 0.745 |
| ATOM OZ | 0 | -0.51 |
| ! ATOM HZ | H | 0.235 |
| ! beta-alanine |  |  |
| GROUP |  |  |
| ATOM NB | NH1 | -0.47 |
| ATOM HB | H | 0.47 |
| ATOM CV | CT2 | -0.18 |
| ATOM HV1 | HA | 0.09 |
| !ATOM HR2 | HA | 0.09 |
| ATOM HV3 | HA | 0.09 |
| GROUP |  |  |
| ATOM CC | CT2 | -0.18 |
| ATOM HC1 | HA | 0.09 |
| ATOM HC2 | HA | 0.09 |
| GROUP |  |  |
| ATOM CT | C | 0.51 |
| ATOM OT | 0 | -0.51 |
| ! ATOM NS | NH1 | -0.47 |
| !ATOM HS1 | H | 0.235 |
| !ATOM HS2 | H | 0.235 |
| !cationic tail |  |  |
| GROUP |  |  |
| ATOM NC | NH1 | -0.47 |
| ATOM HC | H | 0.47 |
| ATOM CE | CT2 | -0.11 |
| ATOM HM1 | HA | 0.09 |
| ATOM HM3 | HA | 0.09 |
| GROUP |  |  |
| ATOM CJ | CT2 | -0.11 |
| ATOM HJ1 | HA | 0.09 |
| AtOM HJ2 | HA | 0.09 |
| ATOM CM | CT2 | -0.11 |
| ATOM HN1 | HA | 0.09 |
| ATOM HN2 | HA | 0.09 |
| GROUP |  |  |
| ATOM NN | NH3 | -0.30 |
| ATOM HH | HC | 0.33 |
| ATOM CX | CT3 | 0.23 |
| ATOM HX1 | HA | 0.05 |
| ATOM HX2 | HA | 0.05 |
| ATOM HX3 | HA | 0.05 |
| ATOM CY | CT3 | 0.23 |
| ATOM HY1 | HA | 0.05 |
| ATOM HY2 | HA | 0.05 |

```
ATOM HY3 HA 0.05
```

!-----------------------------------------------
!bonds for imidazole
BOND NZ2 CZ2 NZ1 CZ1 CZ3 NZ1
BOND CZ1 HE1 CZ2 HE2
DOUBLE NZ2 CZ3 CZ2 CZ1
BOND NZ1 CZ4 CZ4 HZ1 CZ4 HZ2 CZ4 HZ3
BOND CZ3 CB
!----------------------------------------
!bonds for 1st pyrrole ring
BOND CD2 HG CD2 CG NE1 CD1
BOND NE1 CE2
BOND CD1 HD1
DOUBLE CD1 CG CE2 CD2
BOND CH NE1
BOND CH H1 CH H2 CH H3
BOND CE2 CK
DOUBLE CB O2
BOND N2 CG N2 HM N2 CB
!---------------------------------------------
!bonds for 2nd pyrrole ring
BOND CD3 OG1 OG1 HG1 CD3 CG1 NE2 CD4
BOND NE2 CE3
BOND CD4 HD4
DOUBLE CD4 CG1 CE3 CD3
BOND CI NE2
BOND CI HI1 CI HI2 CI HI3
BOND CZ CE3 !NA CJ NA HC NA HD
DOUBLE CK OK
BOND NK CG1 NK HN NK CK ! CB HZ
!------------------------------------------------
!bonds for 3rd pyrrole ring
BOND CD5 HG2 CD5 CG2 NE3 CD6
BOND NE3 CE4
BOND CD6 HD6
DOUBLE CD6 CG2 CE4 CD5
BOND CF NE3
BOND CF HF1 CF HF2 CF HF3
BOND CL CE4 CL NB ! NB HE NB HF
DOUBLE CZ OZ CL OL
BOND NZ CG2 NZ HZ NZ CZ ! CB HZ
!------------------------------------------------
!bonds for beta-alanine
BOND NB CV
BOND NB HB
BOND HV1 CV HV3 CV
DOUBLE CT OT
BOND CV CC CC CT ! CT NS
BOND CT NC !NS HS1 NS HS2
BOND CC HC1 CC HC2
!---------------------------------------------------
!bonds for cationic tail
BOND NC HC NC CE
! BOND CL HL1 CL HL2 CL HL3
BOND CE HM1 CE HM3 CE CJ

```
!DOUBLE C O
BOND CJ HJ1 CJ HJ2 CJ CM
BOND CM HN1 CM HN2 CM NN
BOND NN HH NN CX NN CY
BOND CX HX1 CX HX2 CX HX3
BOND CY HY1 CY HY2 CY HY3
```

| !ic's for imidazole ring |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ! KEEPS HYDROGENS IN RING PLANE |  |  |  |  |  |  |  |  |
| ! IMPH NZ1 | CZ1 | CZ3 C | 4 | NZ1 CZ3 | 3 CZ1 | CZ4 |  |  |
| IMPH CZ2 | CZ1 | NZ2 H | 2 | CZ2 NZ2 | 2 CZ1 | HE2 |  |  |
| IMPH CZ3 | NZ1 | NZ2 CB |  | CZ3 NZ2 | NZ1 CB |  |  |  |
| IMPH CZ1 | CZ2 | NZ1 H |  | CZ1 N | NZ1 CZ2 | HE1 |  |  |
| IC HE1 | CZ1 | NZ1 | CZ3 | 1.07 | 122.67 | -173.67 | 109.79 |  |
| 1.2987 |  |  |  |  |  |  |  |  |
| IC CZ1 | NZ1 | CZ3 | NZ2 | 1.2854 | 109.79 | 0.21 | 110.31 |  |
| 1.3071 |  |  |  |  |  |  |  |  |
| IC NZ1 | CZ3 | NZ2 | CZ2 | 1.2987 | 110.31 | 0.03 | 105.82 |  |
| 1.3165 |  |  |  |  |  |  |  |  |
| IC CZ3 | NZ2 | CZ2 | CZ1 | 1.3071 | 105.82 | -0.23 | 108.68 | 1.3758 |
| IC NZ2 | CZ2 | CZ1 | NZ1 | 1.3165 | 108.68 | 0.35 | 105.39 | 1.2854 |
| IC NZ2 | CZ2 | CZ1 | HE1 | 1.3165 | 108.68 | 172.86 | 131.52 | 1.07 |
| IC CZ2 | CZ1 | NZ1 | CZ3 | 1.3758 | 105.39 | -0.34 | 109.79 | 1.2987 |
| IC CZ2 | NZ2 | CZ3 | CB | 1.3165 | 105.82 | -180.0 | 124.7 | 1.46 |
| IC NZ2 | CZ3 | NZ1 | CZ4 | 1.34 | 110.31 | 180.0 | 127.8 | 1.46 |
| IC HE1 | CZ1 | CZ2 | HE2 | 1.07 | 131.52 | 0.00 | 129.3 | 1.07 |

!ic's for methyl group

| IC | CZ1 | NZ1 | CZ4 | HZ1 | 1.37 | 124.7 | 0.00 | 109.5 |
| :--- | :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: |
| IC | CZ1 | NZ1 | CZ4 | HZ2 | 1.087 | 124.7 | -120.0 | 109.5 |
| IC | CZ1 | NZ1 | CZ4 | HZ3 | 1.37 | 124.7 | 120. | 109.5 |
| iC | CZ3 | NZ1 | CZ4 | HZ3 | 1.37 | 124.7 | -60.0 | 109.5 |

impr CZ4 CZ1 NZ1 CZ3
!ic's for amide group on CE1
IMPR CB CZ3 N2 O2


| IC | CD2 | CG | CD1 | NE1 | 0.0000 | 0.0000 | 0.0000 | 0.0000 |
| :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: | :---: |
| IC | CD1 | CE2 | *NE1 | CH | 0.0000 | 0.0000 | -180.0000 | 127.0 |
| IC | CG | NE1 | *CD1 | HD1 | 0.0000 | 0.0000 | 180.0000 | 0.0000 |
| IC | CE2 | CG | CD2 | HG | 0.0000 | 0.0000 | -180.0000 | 126.9 |


| !ic's for methyl group |  |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | CD1 NE1 CH H1 | 1.37 | 123.2 | 0.00 | 109.5 | 1.08 |
| IC | CD1 NE1 CH H2 | 1.37 | 123.2 | -120.0 | 109.5 | 1.08 |
| IC | CD1 ND1 CH H3 | 1.37 | 123.2 | 120. | 109.5 | 1.08 |
| ic CE2 NE1 CH H3 | 1.37 | 123.2 | -60.0 | 109.5 | 1.08 |  |


!ic's for amide group on CG

| ic | O2 | CB | N2 | HM |  | 1.2233 | 122.84 | -180. |  | 9. 23 | 0.9 | 993 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ic | HM | N2 | CB | O 2 |  | 0.9933 | 119.23 | 180. |  | 12.25 |  | . 23 |  |
| ic | 02 | CB | N2 | CG |  | 1.2233 | 122.84 | 0. |  | . 57 | 1.4 |  |  |
| ic | CB | N2 | CG | CD2 |  | 1.3418 | 122.57 | 180. |  | 4. 9 | 1.41 |  |  |
| ic | CB | N2 | CG | CD1 |  | 1.3418 | 122.57 | 0 . |  | 7. 2 | 1.40 |  |  |
| IC | N2 | CG | CD2 | 2 CE 2 |  | 1.40 | 124.9 | 180. |  | 7.1 | 1.40 |  |  |
| ic | N2 | CG | CD1 | 1 NE1 |  | 1.40 | 127.2 | -180.0 |  | 7. 6 | 1.37 |  |  |
| !ic's for 2nd pyrrole ring |  |  |  |  |  |  |  |  |  |  |  |  |  |
| IC | CG1 | CD4 |  | NE2 CE3 |  | 0.0000 | $0.0000 \quad 0.00$ |  | 0.0 .0000 |  |  | 0.0000 |  |
| ! IC | CD2 | CB |  | *CG | CD1 | 0.0000 | 0.000 | 0 180.00 | 000 | 0.0 | 000 |  | 0.0000 |
| IC | CD4 | CG1 CD |  | CD3 | CE3 | 0.0000 | 0.0000 | 0.000 |  | 0.00 |  |  | . 0000 |
| IC | CD3 | CG1 |  | CD4 | NE2 | 0.0000 | 0.000 | 0.00 |  | 0.00 |  |  | . 0000 |
| IC | CD4 | CE3 * |  | +NE2 CI |  | $0.00000 .0000-180.0000$ |  |  |  | 127. | 0 | , | 46 |
| IC | CG1 | NE2 |  | *CD4 HD4 |  | 0.0000 | 0.000 | 00180.0 | 000 | 0.0 | 000 |  | 0.0000 |
| IC | CE3 | CG1 |  | CD3 | OG1 | 0.000 | 00.00 | 000-180 | . 000 | 001 | 26.9 |  | 1.08 |
| IC | CE3 | CD3 |  | OG1 | HG1 | 0.000 | 0.00 | 0.00 |  | 0.0 |  |  | 0.00 |

!ic's for methyl group

| IC | CD4 | NE2 | CI HI1 | 1.37 | 123.2 | 0.00 | 109.5 | 1.08 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | CD4 | NE2 | CI | HI2 | 1.37 | 123.2 | -120.0 | 109.5 |
| IC | CD4 | NC1 | CI | HI3 | 1.37 | 123.2 | 120. | 109.5 |
| iC | CE3 | NE2 | CI HI3 | 1.37 | 123.2 | -60.0 | 109.5 | 1.08 |

!ic's for amide group on CE3

| !!ic | OJ | C | NA | HC | 1.2233 | 3122.84 | 4180. | 119.23 | 30.9933 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| !!ic | OJ | CJ | NA | HD | 1.2233 | 3122.84 | 40.00 | 119.23 | 30.9933 |
| ic | HZ | NZ | CZ | CE3 | 0.99331 | 119.23 | 0. | 116.25 | 1.47 |
| ic | CG2 | NZ | CZ | CE3 | 1.40 | 119.23 | -180.0 | 116.25 | 1.47 |
| ic | NZ | CZ | CE3 | NE2 | 1.34181 | 116.25 | 180. | 122.11 | 1.38 |
| ic | NZ | CZ | CE3 | CD3 | 1.34181 | 116.25 | 0 . | 130.31 | 1.40 |



| !ic's for amide group on CE4 |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| !!ic | c OL | CL | NB | B HE | 1.2233 | 33122.84 | 4180. | 119.23 | 0.993 |
| ! !ic | c OL | CL | NB | B HF | 1.2233 | 33 122.84 | 40.00 | 119.23 | 0.993 |
| ic | HB | NB | CL | CE4 | 0.99331 | 119.23 | 0. | 116.251 | . 47 |
| !ic | HF | NB | CL | CE 4 | 0.9933 | 119.23 | -180. | 116.25 | 1.47 |
| ic | CV | NB | CL | CE4 | $1.40 \quad 1$ | $119.23-18$ | -180.0 | 116.251 | . 47 |
| ic | NB | CL | CE4 | NE3 | 1.34181 | 116.25 | 180. | 122.11 .38 |  |
| ic | NB | CL | CE4 | CD5 | 1.34181 | 116.25 | 0. | 130.31 .4 |  |
| ic | NE3 | CE4 | CL | OL | 1.381 | 122.0 | 0.00 | 123.51 .2 | 2233 |
| ic | CD5 | CE4 | CL | OL | 1.401 | 130.3 | 180.0 | 123.51 .22 | 2233 |
| IC | CG2 | CD5 | CE4 | CL | 1.401 | 107.1 - | -180.0 | 130.31 .46 |  |
| ic | CD6 | NE3 | CE4 | CL | 1.3710 | 109.8180 | 180.012 | 22.11 .46 |  |
| ic | CF | NE3 | CE4 | CL | 1.461 | 127.0 | 0.00 | 122.11 .4 |  |
| !ic's for amide group on CG2 |  |  |  |  |  |  |  |  |  |
| ic | OZ | CZ | NZ | HZ | 1.2233 | 122.84 | -180. | 119.23 | 0.993 |
| ic | HZ | NZ | CZ | CE3 | 0.9933 | 33119.23 | 30. | 112.25 | 1.47 |
| ic | OZ | CZ | NZ | CG2 | 1.2233 | 33 122.84 | 40. | 125.57 | 1.40 |
| ic | CZ | NZ | CG2 | CD5 | 1.3418 | 18122.57 | 780. | 124.9 | 1.41 |
| ic | CZ | NZ | CG2 | CD6 | 1.3418 | 18122.57 | 70. | 127.2 | 1.40 |
| IC | NZ | CG2 | CD5 | CE4 | 1.40 | 124.9 | 180. | 107.1 | 1.40 |
| ic | NZ | CG2 | CD6 | NE3 | 1.40 | 127.2 | -180. | . 107.6 | 1.37 |
| !ic's for beta-alanine |  |  |  |  |  |  |  |  |  |
| IMPR NB CL CV HB |  |  |  |  |  |  |  |  |  |
| IMPR CL CE4 NB OL |  |  |  |  |  |  |  |  |  |
| IC | OL C | CL NB | HB |  | 1.2233122 | $22.84-180$ | 180. 11 | 9.230 .9933 |  |
| IC | HB N | JB CL | CE4 |  | 0.9933119 | 19.23 | 0.116 | . 251.47 |  |
| IC | OL C | CL NB | CV |  | 1.2233122. | 2.840 | 0.122. | . 571.4488 |  |


| IC | NB CL CE4 CD5 | 1.3418 | 116.25 | 0.109 .31. | . 40 |
| :---: | :---: | :---: | :---: | :---: | :---: |
| IC | NB CL CE4 NE3 | 1.3418 | 116.25 | -180. 109.3 | 1.38 |
| IC | NB CL CE4 HL3 | 1.3418 | 116.25 | -60. 109.3 1 | 1.109 |
| IC | CL NB CV CC | 1.3418 | 122.57 | 180. 110.7 1 | 1.50 |
| IC | CL NB CV HV1 | 1.3418 | 122.57 | -60. 110.7 1 | 1.50 |
| IC | CL NB CV HV3 | 1.3418 | 122.57 | 60. 110.7 1 | 1.50 |
| IC | NB CV CC HC1 | 1.4488 | 110.7 | $60.0 \quad 109.5$ | 1.10 |
| IC | NB CV CC HC2 | 1.4488 | 110.7 | -60.0 109.5 | 1.10 |
| IC | NB CV CC CT | 1.4488 | 110.7 | 180.0109 .5 | 1.50 |
| ! IC | CV CC CT NS | 1.50 | 109.5 | 180.0115 .0 | 1.40 |
| IC | CV CC CT NC | 1.50 | 109.5 | -180.0 115.0 | 1.40 |
| IC | CV CC CT OT | 1.50 | 109.5 | 0.00123 .1 | 1.24 |
| IC | OT CT NC HC | 1.24 | 121.9 | 180.0121 .2 | 0.9933 |
| IC | OT CT NC CE | 1.24 | 121.9 | 0.00121 .2 | 1.47 |
| ! IC | OT CT NS HSI | 1.24 | 121.9 | 180.0121 .2 | 20.9933 |
| ! IC | OT CT NS HS2 | 1.24 | 121.9 | 0.00121 .2 | 21.47 |

!IMPR CT CC NS OT
IMPR CT CC NC OT

!ic's for tail
IMPR NC CT CE HC
IMPR CT CC NC OT

| IC | OT | CT | T NC | HC | 1.2233122 .84180 |  |  | 0. 119.23 | . 9933 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC | HC | NC | C CT | CL | 0.9933119 .23 |  |  | . 116.251. | . 5118 |
| IC | OT | CT | T NC | CE | 1.2233122 .8 |  | . 840. | . 122.571. | . 4488 |
| IC | NC | CT | T CC | HC1 | 1.3418116 .25 |  | $25-60$. | . 109.31. | . 109 |
| IC | NC | CT | T CC | HC2 | 1.3418116 .25 |  |  | . 109.3 1. | . 109 |
| IC | NC | CT | T CC | CV | 1.3418116 .25 |  | $5-180$. | 109.31 .50 | 50 |
| IC | CT | NC | C CE | HM1 | 1.3418122 .57 |  | 57 60. | 110.71 .11 | 113 |
| IC | CT | NC | C CE | CJ | 1.3418122 .57 |  | 57180. | 110.71 .50 |  |
| IC | CT | NC | C CE | HM3 | 1.3418122 .57 |  | $57-60$. | 110.71. | . 113 |
| IC | NC | CE | CJ CM |  | 1.47109 .5 |  | 180.0 | 109.51. | . 50 |
| IC | NC | CE | CJ HJ1 |  | 1.47109 .5 |  | -60.0 | 109.51. | . 10 |
| IC | NC | CE | CJ HJ2 |  | 1.47109 .5 |  | 60.0 | 109.51. | . 10 |
| IC | CE | CJ | CM HN1 |  | 1.50109 .5 |  | 60.0 | 109.51 .10 |  |
| IC | CE | CJ | CM HN2 |  | 1.50 | 109.5 | -60.0 | 109.51 .10 | 10 |
| IC | CE | CJ | CM NN |  | 1.50 | 109.5 | 180.0 | 109.51 .40 | 40 |
| IC | CJ | CM N | NN HH |  | 1.50 | 109.5 | 60.00 | 109.5 | 1.10 |
| IC | CJ | CM N | NN CX |  | 1.50 | 109.5 | -60.0 1 | 109.51 .47 |  |
| IC | CJ | CM N | NN CY |  | 1.50 | 109.5 | 180.0 | 109.51 .4 |  |
| IC | CM | NN C | CX HX1 |  | $1.47$ | 109.5 | -60.0 | 109.51 .10 | 10 |
| IC | CM | NN CX | CX HX2 |  | 1.47 | 109.5 | 60.0 | 109.51 .10 | 10 |
| IC | CM | NN C | CX HX3 |  | 1.47 | 109.5 | -180. | 109.51 .10 |  |
| IC | CM | NN C | CY HY1 |  | 1.47 | 109.5 | 60.0 | 109.51 .10 |  |
| IC | CM | N1 C | CY HY2 |  | 1.47 | 109.5 | 180.0 | 109.51 .10 |  |
| IC | CM | N1 C | CY HY3 |  | 1.47 | 109.5 | -60.0 | 109.51 .10 |  |
| ! IC | H1 | N1 | CX HX1 |  | 1.10 | 109.5 | -180. | . 0109.5 | 1.10 |
| IC | HY2 | CY | NN HH |  | 1.10 | 109.5 | -60.0 | 0.109 .5 | 1.10 |
| IC | HY3 | CY | NN HH |  | 1.10 | 109.5 | 60.0 | O 109.5 | 1.10 |

```
impr OT CC CT NC
impr OT CV CC CT
impr OT NC CT CC
end
read parameter card append
BONDS
!
!atom types Kb b0
!
!sulfate bond
OC S 400.000 1.43
!polyamide bonds
NH2 CPT 500.000 1.41 !A.E.L.
CD CA 500.000 1.45 !A.E.L.
CT3 NY 500.000 1.46 !A.E.L.
CPT HP 500.000 1.083 !A.E.L.
NC2 CPT 500.000 1.41 !A.E.L.
CD NH2 500.00 1.36 !A.E.L.
OS CC 400.000 1.35 !a.E.L.
OS CT1 400.00 1.46 !A.E.L.
!NH2 CT2 500.000 1.46 !A.E.L.
NC2 H 500.000 1.00 !A.E.L.
CC H 500.00 1.10
NH2 CY 500.00 1.40
!-----------------------------------------------
!bonds for ipy drug
!N-methyl-pyrrole
\begin{tabular}{llccl} 
CPT & HP & 500.00 & 1.08 & !A.E.L. \\
CT3 & NY & 500.000 & 1.46 & \\
C & CPT & 500.00 & 1.47 & \\
NH1 & CY & 500.00 & 1.40 & \\
C & H & 500.00 & 1.10 & \\
OS & CT1 & 400.00 & 1.46 & !A.E.L. \\
OS & C & 400.000 & 1.35 & !a.E.L. \\
CPT & OH1 & 500.000 & 1.37 & !ael; for hydroxypyrrole
\end{tabular}
!imidazole
NR1 CT3 500.00 1.46 !A.E.L.
CPH2 C 500.000 1.47
CPH1 NH1 500.00 1.40
\begin{tabular}{llll}
C & H & 500.00 & 1.10
\end{tabular}
ANGLES
!
!atom types Ktheta Theta0 Kub S0
!
```



| CT2 | NC2 | CT3 | 40.000 | 109.5 | !A.E.L. |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CT2 | NC2 | H | 40.000 | 109.5 | !a.E.L. |  |
| CT3 | NC2 | H | 40.000 | 109.5 | !A.E.L. |  |
| CD | NH2 | CT2 | 40.000 | 120.6 | !A.E.L. |  |
| CT2 | C | NC2 | 40.00 | 120.0 | !A.E.L. |  |
| C | CT2 | NC2 | 40.00 | 35.00 | !A.E.L. |  |
| !angles for ipy drug |  |  |  |  |  |  |
| !Boc group |  |  |  |  |  |  |
| NH1 | C O |  | 50.00 | 116.50 | 50.000 | 2.45000 |
| 0 | C O | OS | 15.00 | 121.00 | 50.00 | 2.44000 |
| C | OS C | CT1 | 33.00 | 109.50 | 30.000 | 2.16300 |
| OS | CT1 | CT3 | 34.500 | 110.10 | 22.53 | 2.17900 |
| ! N-methyl-pyrrole |  |  |  |  |  |  |
| CPT | CY | NH1 | 40.000 | 124.9 | !A.E.L. |  |
| CA | CA | NH1 | 40.000 | 127.2 |  |  |
| HP | CPT | CY | 40.000 | 126.9 |  |  |
| HP | CPT | CPT | 40.000 | 125.9 |  |  |
| CA | NY | CT3 | 40.000 | 123.3 |  |  |
| CPT | NY | CT3 | 40.000 | 127. |  |  |
| CPT | CPT | C | 40.000 | 130.2 |  |  |
| NY | CPT | C | 40.000 | 122.1 |  |  |
| NY | CT3 | HA | 40.000 | 109.5 |  |  |
| CPT | C | 0 | 40.000 | 123.5 |  |  |
| CPT | C | NH1 | 40.000 | 114.4 |  |  |
| H | NH1 | H | 40.000 | 115.7 |  |  |
| CY | NH1 | H | 40.000 | 117.2 |  |  |
| CY | NH1 | C | 40.00 | 125.2 |  |  |
| NH1 | C | H | 40.00 | 112.0 |  |  |
| CA | CY | NH1 | 40.000 | 127.2 |  |  |
| CA | CY | C | 40.00 | 100.3 |  |  |
| CPT | CY | C | 40.000 | 151.9 |  |  |
| CPT | CY | CPT | 40.000 | 36.3 |  |  |
| !imidazole |  |  |  |  |  |  |
| CPH1 | CPH1 | 1 NH1 | 40.000 | 129.3 | !A.E.L. |  |
| NR2 | CPH1 | 1 NH1 | 40.000 | 120.0 |  |  |
| CPH1 | NR1 | CT3 | 40.000 | 124.7 |  |  |
| CPH2 | NR1 | CT3 | 40.000 | 127.8 |  |  |
| NR1 | CPH2 | 2 C | 40.000 | 123.8 |  |  |
| NR2 | CPH2 | 2 C | 40.000 | 124.8 |  |  |
| NR1 | CT3 | HA | 40.000 | 107.4 |  |  |
| CPH2 | C | $\bigcirc$ | 40.000 | 122.6 |  |  |
| CPH2 | C | NH1 | 40.000 | 112.9 |  |  |
| H | NH1 | H | 40.000 | 121.1 |  |  |
| CPH1 | NH1 | H | 40.000 | 115.7 |  |  |
| CPH1 | NH1 | C | 40.000 | 124.5 |  |  |
| NH1 | C | H | 40.000 | 112.4 | . 4 |  |
| !cationic tail |  |  |  |  |  |  |
| CT2 | NH3 C | CT3 | 40.000 | 109.5 | ! A.E.l. |  |
| CT3 | NH3 C | CT3 | 40.00 | 110.8 |  |  |
| CA | CY C | CC | 40.0 | 100.3 |  |  |
| CPT | CY C | CC | 40.0 | 151.9 |  |  |




| CY | CA | CD | OB | 2.32 |  | 180.00 !A.E.L. |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NY | CA | $C D$ | OB | 2.32 |  |  | 180.00 ! A . | .L. |
| ! dihedrals for ipy drug |  |  |  |  |  |  |  |  |
| CA | NY | CT3 | HA | 0.00 | 1 | 1 | 0.00 | !methyl group |
| CA | NY | CT3 | HA | 0.000 | 3 | 3 | -120.0 |  |
| CA | NY | CT3 | HA | 0.000 | 2 | 2 | -60.0 |  |
| CPT | NY | CT3 | HA | 0.055 | 3 |  | 0.0 | !methyl group |
| CPT | NY | CT3 | HA | 0.00 | 1 |  | -120. |  |
| CPT | NY | CT3 | HA | 0.00 | 2 |  | -60.0 |  |
| HP | CPT | CY | CA | 2.0 | 1 |  | 180.0 |  |
| HP | CPT | CY | NH1 | 2.8 | 2 |  | 180.0 |  |
| CY | CA | NY | CT3 | 0.8 | 2 |  | 180.0 |  |
| CY | NH1 | C | 0 | 1.5 | 2 |  | 180.0 |  |
| CY | NH1 | C | H | 1.20 | 1 |  | 180.0 |  |
| CPT | CY | NH1 | H | 0.000 | 1 |  | 0.00 |  |
| CPT | CY | NH1 | C | 0.200 | 1 |  | 180.0 |  |
| CPT | CPT | NY | CT3 | 0.8 | 2 |  | 180.0 |  |
| CPT | CPT | C | 0 | 2.5 | 2 |  | 180.0 |  |
| CPT | CPT | C | NH1 | 1.5 | 2 |  | 180.0 |  |
| CA | NY | CPT | C | 2.00 | 1 |  | 180.0 |  |
| HP | CA | CY | NH1 | 1.00 | 2 |  | 180.0 |  |
| CA | CY | NH1 | H | 0.0 | 1 |  | 0.00 |  |
| CA | CY | NH1 | C | 2.5 | 2 |  | 180.0 |  |
| CA | CY | NH1 | C | 3.0 | 1 |  | 180.0 |  |
| HP | CA | NY | CT3 | 0.4000 | 2 |  | 180.0 |  |
| NY | CA | CY | NH1 | 2.0 | 2 |  | 180.0 |  |
| NY | CPT | C | 0 | 1.3 | 2 |  | 180.0 |  |
| NY | CPT | C | 0 | 3.7 | 1 |  | 180.0 |  |
| NY | CPT | C | NH1 | 1.0 | 1 |  | 0.0 |  |
| CPT | CPT | CY | NH1 | 3.0 | 2 |  | 180.0 |  |
| CT3 | NY | CPT | C | 0.8 | 2 |  | 180.0 |  |
| H | NH1 | C | H | 1.4 | 2 |  | 180.0 |  |
| CPT | C | NH1 | H | 2.5 |  | 2 | 180.0 |  |



| CPH2 | 2 C | NH1 C |  | 0.2000 | 1 | 0.00 |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CPH2 | 2 C | NH1 CY |  | 0.2000 | 2 | 180.0 |  |  |
| !netropsin |  |  |  |  |  |  |  |  |
| ! CY | NH1 | C | CT2 | 0.2000 | 1 | 0.00 |  |  |
| CY | NH1 | C | CT2 | 1.000 | 1 | -180.0 |  |  |
| 0 | C | CT2 N | NC2 | 0.5000 | 1 | 0.000 |  |  |
| CT2 | CT2 | C N | NC2 | 0.000 | 1 | 0.00 |  |  |
| NH1 | C | CT2 N | NC2 | 0.6000 | 1 | 0.00 |  |  |
| HA | CT2 | C N | NC2 | 0.0000 | 3 | 0.00 |  |  |
| !distamycin |  |  |  |  |  |  |  |  |
| CC | NH2 | CY | CPT | 0.2000 | 1 | 180.0 |  |  |
| CC | NH2 | CY | CA | 2.5 | 2 | 180.0 |  |  |
|  | NH2 | CY | CA | 3.7 | 1 | 180.0 |  |  |
| NH2 | CY | CPT | HP | 2.8 | 2 | 180.0 |  |  |
| NH2 | CY | CA | HP | 1.00 | 1 | 180.0 |  |  |
| H | NH2 | CC | H | 1.4 | 2 | 180.0 |  |  |
| H | NH2 | CY | CPT | 0.050 | 1 | 0.00 |  |  |
|  | NH2 | CY | CA | 0.05 | 1 | 0.00 |  |  |
| 0 | CC | NH2 |  | 1.50 | 2 | 180.0 |  |  |
| H | CC | NH2 | CY | 1.20 | 1 | 180.0 |  |  |
| !hydroxypyrrole |  |  |  |  |  |  |  |  |
| OH1 | CPT | CY | CA | 3.1 | 2 | 180.0 |  |  |
| OH1 | CPT | CY | NH1 | 0.20 | 2 | 180.0 |  |  |
| H | OH1 | CPT | CY | 0.00 | 2 | 90.0 |  |  |
| H | OH1 | CPT | CY | 4.0 | 1 | 0.00 |  |  |
| H | OH1 | CPT | CPT | 0.50 | 2 | 180.0 |  |  |
| H | OH 1 | CPT | CPT | 0.30 | 1 | 0.00 |  |  |
| IMPROPER |  |  |  |  |  |  |  |  |
| !atom types |  |  |  | Kpsi |  | psi0 |  |  |
| !polyamide impropers |  |  |  |  |  |  |  |  |
| CA | CA | CA | NY | 90.000 |  | 0 | -180.00 | !A.E.L. |
| CPT | CA | NY | CA | 90.000 |  | 0 | 180.00 | !A.E.L. |
| CY | CPT | CPT | T NH2 | 90.000 |  | 0 | 180.00 | !A.E.L. |
| CA | CA | CA | CT1 | 90.000 |  | 0 | 180.00 | !A.E.L. |
| OH1 | CA | CD | OB | 90.000 |  | 0 | 180.0 | !A.E.L. |
| CD | NY | CA | CY | 90.000 |  | 0 | 180.0 | !A.E.L. |
| CT3 | NY | CPT | T CPT | 90.000 |  | 0 | -180.0 | !A.E.L. |
| CA | CPT | NY | CT3 | 90.000 |  | 0 | -180.0 | !A.E.L. |
| CA | CPT | CY | NH2 | 90.000 |  | 0 | 180.00 | !A.E.L. |
| CPT | CY | CPT | T NH2 | 90.000 |  | 0 | -180.0 | !A.E.L. |
| CPT | NY | CPT | T HP | 90.000 |  | 0 | 180.0 | !A.E.L. |
| OB | CA | $C D$ | OH1 | 90.00 |  | 0 | -180.0 | !A.E.L. |
| OH1 | OB | $C D$ | CA | 90.00 |  | 0 | -180.0 | !A.E.L. |
| OB | OH1 | CD | CA | 90.00 |  | 0 | 180.0 | !A.E.L. |


| CY | $C D$ | OH1 | 1 H | 90.000 | 0 | -180.0 | !A.E.L. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| HP | CA | $C D$ | OH1 | 90.00 | 0 | 0.00 | !A.E.L. |
| OH1 | CA | CY | HP | 90.000 | 0 | 0.00 | !A.E.L. |
| OB | CA | NY | CT3 | 90.000 | 0 | 0.00 | !A.E.L. |
| NH2 | CT2 | 2 CT 2 | 2 HA | 90.000 | 0 | -180.0 | !A.E.L. |
| OB | NH2 | 2 CT 2 | 2 CT 2 | 90.000 | 0 | 180.0 | !A.E.L. |
| CA | OB | $C D$ | OH1 | 90.000 | 0 | 180.0 | !A.E.L. |
| NY | CD | OH1 | H | 90.000 | 0 | 180.0 | !A.E.L. |
| CY | $C D$ | CA | NY | 90.000 | 0 | -180.0 | !A.E.L. |
| H | $C D$ | CA | CA | 90.000 | 0 | 0.000 | 0!A.E.L. |
| H | $C D$ | CA | NY | 90.000 | 0 | -180.0 | O ! A.E.L. |
| CD | CY | CA | NY | 90.000 | 0 | 180.0 | 0 ! A.E.L. |
| CD | CA | CPT | CPT | 90.000 | 0 | 0.00 | !A.E.L. |
| CA | CA | NY | CPT | 90.000 | 0 | -180.0 | 0 ! A.E.L. |
| CA | CA | CA | HP | 90.000 | 0 | 180.0 | 0 ! A.E.L. |
| CT1 | CA | CA | HP | 90.000 | 0 | -180.0 | 0 ! A.E.L. |
| HC | HC | NH2 C | CPT | 90.000 | 0 | -120.0 | 0 ! A.E.L. |
| HC | CPT | NH2 H |  | 90.000 | 0 | 120.0 | 0 ! A.E.L. |
| CA | CPT | NH2 H |  | 90.000 | 0 | -20.33 | 3 !A.E.L. |
|  | CD N | NH2 CP |  | 90.00 | 0 | 180.0 | 0 !A.E.L. |
| NH2 | CD | CPT | HC | 20.000 |  | 0.00 | 0 !A.E.L. |
| CD | CA | NH2 |  | 45.000 |  | 0.00 | 0 !A.E.L. |
| CA | CPT | T NH2 |  | 90.000 |  | 0.00 | 0 !A.E.L. |
| CC | OS | NH2 | 0 | 45.000 | 0 | 0.00 | !A.E.L. |
| CC | NH2 | OS | 0 | 45.000 | 0 | 0.00 ! | !A.E.L. |
| OS | NH2 | CC | 0 | 90.000 | 0 | -180.0 ! | !A.E.L. |
| OS | 0 | CC | NH2 | 90.00 | 0 | 180.0 ! | !A.E.L. |
| CT3 | OS | CT1 | CT3 | 90.000 | 0 | -120.0! | !A.E.L. |
| NH2 | CC | CPT | HC | 20.00 | 0 | 0.00 | ! A.e.L. |
| NH2 | $C D$ | CT2 | HC | 20.000 | 0 | 0.00 | !A.E.L. |
| CPT | CA | CY HP | HP | 100.000 | 0 | 180.0 |  |
| NY | CPT | CPT HP | HP | 100.000 | 0 | -180.0 |  |
| HP | CPT | CY C | CA | 100.000 | 0 | 180.0 |  |
| HP | CA | CY C | CPT | 100.000 | 0 | -180.0 |  |
| CD | CA | NH2 | OB | 45.000 | 0 | 0.00 |  |
| CD | NH2 | CA | OB | 45.000 | 0 | 0.000 |  |
| !impropers for ipy drug ! |  |  |  |  |  |  |  |
| NH1 | C | H | CY | 4.00 | 0 | 0.00 |  |
| !imid |  |  |  |  |  |  |  |
| CT3 CPH1 NR1 CPH2 |  |  |  | 90.000 | 0 | 180.0 |  |
| NR1 CPH1 CPH2 CT3 |  |  |  | 90.000 | 0 | 0.00 |  |
| NR1 CPH2 CPH1 CT3 |  |  |  | 90.000 | 0 | 0.00 |  |
| CPH2 NR1 NR2 C |  |  |  | 90.000 | 0 | 0.00 |  |
| CPH2 NR2 NR1 C |  |  |  | 90.000 | 0 | 0.00 |  |
| O CPT C NH1 |  |  |  | 90.00 | 0 | -180.0 |  |
| CPH1 CPH2 NR1 CT3 |  |  |  | 90.000 | 0 | 180.0 |  |
| $0 \quad \text { CT2 C NH1 }$ |  |  |  | 90.000 | 0 | -180.0 |  |


| $\bigcirc \mathrm{N}$ | NH1 | C | CT2 | 90.000 | 0 |  | 180.0 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CY C | C N | NH1 | H | 90.00 | 0 |  | 180.0 |
| C CA | A CY | Y NH |  | 90.000 | 0 |  | 0.00 |
| H C | NH | H1 CT |  | 90.000 | 0 |  | 180.0 |
| NC2 0 | $\bigcirc \mathrm{C}$ | CT |  | 90.000 | 0 |  | 0.00 |
| C NC2 | 2 CT 2 | 2 HA |  | 90.000 | 0 |  | 120.0 |
| NH1 C | C CT2 | 2 CT |  | 90.0 | 0 |  | 0.00 |
| C NC2 NC2 CT2 |  |  |  | 90.00 | 0 |  | 0.00 |
| !distamycin |  |  |  |  |  |  |  |
| NH2 | CC | H | CY | 4.0 |  | 0 | 0.000 |
| NH1 | CPT | C | $\bigcirc$ | 90.0 |  | 0 | 180.0 |
| $\bigcirc$ | NH1 | C | CPT | 90.0 |  | 0 | 180.0 |
| CT2 | C | NH1 | H | 90.0 |  | 0 | 180.0 |
| CC | H | NH2 | 0 | 120.0 |  | 0 | 0.00 |
| CC | NH2 | H |  | 120.0 |  | 0 | 0.0 |

NONBONDED nbxmod 5 atom cdiel shift vatom vdistance vswitch cutnb 14.0 ctofnb 12.0 ctonnb 10.0 eps 1.0 e14fac 1.0 wmin 1.5 !adm jr., 5/08/91, suggested cutoff scheme
!
!atom ignored epsilon Rmin/2 ignored eps,1-4
Rmin/2,1-4
!
!Butadiene
CC1A 0.0 -0.0680 2.0900 !
CC1B $0.0-0.0680 \quad 2.0900$ !
CC2 $0.0 \quad-0.0640 \quad 2.0800$ !
!sulfate
! OC $\quad 2.1400 \quad-0.6469 \quad 1.6000$
!S $\quad 0.3400 \quad-0.0430 \quad 1.890$
!SCH1, SCH2
NS1 $0.000000-0.200000 \quad 1.850000$ ! N for deprotonated Schiff's
base
NS2 $0.000000-0.200000 \quad 1.850000$ ! N for protonated Schiff's base
end

## Appendix B

This appendix contains the initial CHARMM parameters for the retinoid model compounds, CHARMM atom names and numbering for the retinoids used in the protein/retinoid simulations, and distances measured from the protein/retinoid simulations. The retinoid topology and parameters are at the end of this appendix.

Initial CHARMM surfaces for retinoid model compounds

1

2


4

5

6

7

8

9

10

11

12

13

14

15

16

17


18


19

Figure B. 1 Model compounds and patch residues for retinoids (some hydrogens not shown for clarity). Models 1-17 are the model compounds, and 18 and 19 are used to create patches for forming Schiff bases.
energy vs. dihedral (1)

energy vs. dihedral (2)

energy vs. dihedral (3)

energy vs. dihedral (4)


Figure B. 2 Initial CHARMM surfaces for model compounds 1-4.
energy vs. dihedral (5)

$\rightarrow-$ final CHARMM
$\rightarrow-$ MP2/6-31G(d)
$\rightarrow$ intial CHARMM
energy vs. dihedral (6)

energy vs. dihedral (7)

energy vs. dihedral (8)


Figure B. 3 Initial CHARMM surfaces for model compounds 5-8.
energy vs. dihedral (9)

energy vs. dihedral (10)

energy vs. dihedral (11)



Figure B. 4 Initial CHARMM surfaces for model compounds 9-12.

Additional tables for comparison of CHARMM energy minimized structures and experimental data

Table B. 1 Experimental, QM, and MM comparison of $p$-acetamide phenol (7) ${ }^{\text {a }}$

| Bond (£) | x-ray | MP2 | CHARMM | Error (MP2 \& x-ray) | Error <br> (CHARMM <br> \& x-ray) | $\begin{aligned} & \text { Error (MP2 } \\ & \& \\ & \text { CHARMM) } \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C22-C27 |  |  |  |  |  |  |
| C27-C26 | 1.39 | 1.39 | 1.40 | 0 | 0.01 | 0.01 |
| C26-C25 | 1.38 | 1.40 | 1.40 | 0.02 | 0.02 | 0 |
| C25-C24 | 1.38 | 1.40 | 1.40 | 0.02 | 0.02 | 0 |
| C24-C23 | 1.39 | 1.40 | 1.41 | 0.01 | 0.02 | 0.01 |
| C25-O28 | 1.38 | 1.38 | 1.41 | 0 | 0.03 | 0.03 |
| C22-N21 | 1.42 | 1.41 | 1.42 | 0.01 | 0 | 0.01 |
| N21-C15 | 1.34 | 1.38 | 1.34 | 0.04 | 0 | 0.04 |
| C15-O29 | 1.22 | 1.23 | 1.22 | 0.01 | 0 | 0.01 |
| C15-C14 | 1.51 | 1.52 | 1.48 | 0.01 | 0.03 | 0.04 |
| Angle ( ${ }^{\mathbf{0}}$ ) |  |  |  |  |  |  |
| C22-C27-C26 | 119.4 | 119.4 | 120.5 | 0 | 1.1 | 1.1 |
| C27-C26-C25 | 120.7 | 121 | 119.9 | 0.3 | 0.8 | 1.1 |
| C26-C25-C24 | 120.2 | 119.6 | 120 | 0.6 | 0.2 | 0.4 |
| C25-C24-C23 | 119.5 | 119.7 | 120 | 0.2 | 0.5 | 0.3 |
| C24-C23-C22 | 120.4 | 120.9 | 120.6 | 0.5 | 0.2 | 0.3 |
| C23-C22-C27 | 119.8 | 119.3 | 119 | 0.5 | 0.8 | 0.3 |
| C26-C25-O28 | 118.9 | 123.3 | 120.4 | 4.4 | 1.5 | 2.9 |
| C24-C25-O28 | 120.8 | 117.1 | 119.6 | 3.7 | 1.2 | 2.5 |
| C23-C22-N21 | 116.5 | 117.2 | 115.9 | 0.7 | 0.6 | 1.3 |
| C27-C22-N21 | 123.8 | 123.5 | 125.1 | 0.3 | 1.3 | 1.6 |
| C22-N21-C15 | 129.8 | 128.7 | 129.2 | 1.1 | 0.6 | 0.5 |
| N21-C15-O29 | 122.7 | 124.2 | 124.8 | 1.5 | 2.1 | 0.6 |
| N21-C15-C14 | 114.6 | 113.8 | 115.2 | 0.8 | 0.6 | 1.4 |
| O29-C15-C14 | 122.7 | 122 | 119.9 | 0.7 | 2.8 | 2.1 |
| Dihedral ( ${ }^{\mathbf{0}}$ ) |  |  |  |  |  |  |
| C22-C27-C26-C25 | 0.45 | 0.09 | 0.26 | 0.36 | 0.19 | 0.17 |
| C27-C26-C5-C24 |  |  |  |  |  |  |
|  | 2.23 | 0.13 | 0 | 2.1 | 2.23 | 0.13 |
| C26-C25-C24-C23 | -2.65 | -0.11 | -0.04 | 2.54 | 2.61 | 0.07 |
| C25-C24-C23-C22 | 0.4 | 0.07 | 0.18 | 0.33 | 0.22 | 0.11 |
| C24-C23-C22-C27 | 2.27 | 0.03 | 0.43 | 2.24 | 1.84 | 0.4 |
| C23-C22-C27-C26 | -2.68 | -0.05 | -0.48 | 2.63 | 2.2 | 0.43 |
| O28-C25-C26-C27 | -178.2 | -180 | -179.9 | 1.8 | 1.7 | 0.1 |
| O28-C25-C24-C23 | 177.8 | 180 | 179.9 | 2.2 | 2.1 | 0.1 |
| C23-C22-N21-C15 | -164.3 | -179.7 | -164.4 | 15.4 | 0.1 | 15.3 |
| C27-C22-N21-C15 | 17.7 | 0.2 | 14.2 | 17.5 | 3.5 | 14 |
| C22-N21-C15-O29 | 0.57 | 0.9 | 3.39 | 0.33 | 2.82 | 2.49 |
| C22-N21-C15-C14 | -178.8 | -177 | -179.4 | 1.8 | 0.6 | 2.4 |
| ${ }^{\text {a }}$ Table C. 1 is continued from Table 4.5 in Chapter 4. |  |  |  |  |  |  |

Table B. 2 Comparison of CHARMM (energy-minimized structure) and x-ray data for retinal

| Bond ( $\AA$ ) | x-ray | CHARMM |  |
| :--- | ---: | ---: | ---: |
| C1-C2 | 1.55 | 1.55 | 0 |
| C1-C6 | 1.54 | 1.53 | 0.01 |
| C1-C17 | 1.51 | 1.55 | 0.04 |
| C1-C18 | 1.54 | 1.55 | 0.01 |
| C2-C3 | 1.42 | 1.53 | 0.11 |
| C3-C4 | 1.49 | 1.53 | 0.04 |
| C4-C5 | 1.51 | 1.51 | 0 |
| C5-C6 | 1.33 | 1.37 | 0.04 |
| C5-C16 | 1.51 | 1.51 | 0 |
| C6-C7 | 1.48 | 1.5 | 0.02 |
| C7-C8 | 1.32 | 1.35 | 0.03 |
| C8-C9 | 1.47 | 1.48 | 0.01 |
| C9-C10 | 1.35 | 1.35 | 0 |
| C9-C19 | 1.49 | 1.51 | 0.02 |
| C10-C11 | 1.44 | 1.48 | 0.04 |
| C11-C12 | 1.34 | 1.35 | 0.01 |
| C12-C13 | 1.45 | 1.48 | 0.03 |
| C13-C14 | 1.34 | 1.35 | 0.01 |
| C13-C20 | 1.5 | 1.51 | 0.01 |
| C14-C15 | 1.46 | 1.48 | 0.02 |
| C15-O | 1.2 | 1.2 | 0 |

Angle ( ${ }^{\circ}$ )

| C2-C1-C6 | 110.2 | 112 | 1.8 |
| :--- | ---: | ---: | ---: |
| C1-C6-C5 | 122.8 | 121 | 1.8 |
| C6-C5-C4 | 122.6 | 122.8 | 0.2 |
| C5-C4-C3 | 115.3 | 114.4 | 0.9 |
| C4-C3-C2 | 115.5 | 109.9 | 5.6 |
| C3-C2-C1 | 115.8 | 111.9 | 3.9 |
| C4-C5-C16 | 113.1 | 113.9 | 0.8 |
| C6-C5-C16 | 124.3 | 123.2 | 1.1 |
| C2-C1-C18 | 104.7 | 107.1 | 2.4 |
| C2-C1-C17 | 112 | 109.2 | 2.8 |
| C6-C1-C18 | 110.3 | 110.8 | 0.5 |
| C6-C1-C17 | 110.2 | 1109.4 | 0.8 |
| C1-C6-C7 | 122.2 | 116.1 | 1.1 |
| C5-C6-C7 | 124.5 | 122 | 0.2 |
| C6-C7-C8 | 126.5 | 127.2 | 2.7 |
| C7-C8-C9 | 118.3 | 125.7 | 0.8 |
| C8-C9-C10 | 118.4 | 121.6 | 3.3 |
| C8-C9-C19 | 123.3 | 115 | 3.4 |
| C10-C9-C19 | 127.2 | 123.3 | 0 |
| C9-C10-C11 | 125.7 | 1.5 |  |

## Table B. 2 (cont'd)

| C10-C11-C12 | 123.6 | 123.8 | 0.2 |
| :---: | :---: | :---: | :---: |
| C11-C12-C13 | 126 | 124.9 | 1.1 |
| C12-C13-C14 | 118.3 | 123.3 | 5 |
| C12-C13-C20 | 118.3 | 112.7 | 5.6 |
| C14-C13-C20 | 123.4 | 123.5 | 0.1 |
| C13-C14-C15 | 125.9 | 121.9 | 4 |
| C14-C15-O | 123.2 | 122.8 | 0.4 |
| Dihedral ( ${ }^{\circ}$ ) |  |  |  |
| C1-C2-C3-C4 | 50 | 59.9 | 9.9 |
| C2-C3-C4-C5 | -30.8 | -45.4 | 14.6 |
| C3-C4-C5-C6 | 6.83 | 21.3 | 14.47 |
| C4-C5-C6-C1 | -2.15 | -9.54 | 7.39 |
| C5-C6-C1-C2 | 18.7 | 22.6 | 3.9 |
| C6-C1-C2-C3 | -42.6 | -48.2 | 5.6 |
| C16-C5-C4-C3 | -172.3 | -162.8 | 9.5 |
| C16-C5-C6-C1 | 176.9 | 175 | 1.9 |
| C16-C5-C6-C7 | -2.37 | -6.5 | 4.13 |
| C18-C1-C6-C5 | 133.9 | 142 | 8.1 |
| C18-C1-C6-C7 | -46.8 | -48.9 | 2.1 |
| C18-C1-C2-C3 | -161.2 | -169.8 | 8.6 |
| C17-C1-C6-C5 | -105.5 | -98.7 | 6.8 |
| C17-C1-C6-C7 | 73.9 | 70.5 | 3.4 |
| C17-C1-C2-C3 | 80.5 | 73.2 | 7.3 |
| C5-C6-C7-C8 | -58.3 | -58.7 | 0.4 |
| C1-C6-C7-C8 | 122.4 | 132.3 | 9.9 |
| C6-C7-C8-C9 | -179.3 | -179.7 | 0.4 |
| C7-C8-C9-C10 | 175.8 | 134.2 | 41.6 |
| C7-C8-C9-C19 | -4.47 | -43.1 | 38.63 |
| C8-C9-C10-C11 | 179.1 | 178.5 | 0.6 |
| C9-C10-C11-C12 | -179.1 | -177 | 2.1 |
| C19-C9-C10-C11 | -0.64 | -1.46 | 0.82 |
| C10-C11-C12-C13 | 175.5 | 174.6 | 0.9 |
| C11-C12-C13-C14 | -178.4 | -63.4 | 115 |
| C11-C12-C13-C20 | -0.63 | -124.9 | 124.27 |
| C12-C13-C14-C15 | 177.7 | 179.5 | 1.8 |
| C20-C13-C14-C15 | 0 | -8.71 | 8.71 |
| C13-C14-C15-O | 179.3 | 178.4 | 0.9 |

Table B. 3 Comparison of CHARMM (energy-minimized structure) and x-ray data for retinoic acid (triclinic form)

| Bond ( $\AA$ ) | x-ray | CHARMM | error |
| :---: | :---: | :---: | :---: |
| C1-C2 | 1.54 | 1.55 | 0.01 |
| C1-C6 | 1.54 | 1.53 | 0.01 |
| C1-C16 | 1.53 | 1.54 | 0.01 |
| C1-C17 | 1.51 | 1.54 | 0.03 |
| C2-C3 | 1.43 | 1.53 | 0.1 |
| C3-C4 | 1.5 | 1.53 | 0.03 |
| C4-C5 | 1.51 | 1.51 | 0 |
| C5-C6 | 1.34 | 1.37 | 0.03 |
| C5-C18 | 1.51 | 1.51 | 0 |
| C6-C7 | 1.47 | 1.49 | 0.02 |
| C7-C8 | 1.34 | 1.35 | 0.01 |
| C8-C9 | 1.45 | 1.48 | 0.03 |
| C9-C10 | 1.35 | 1.35 | 0 |
| C9-C19 | 1.51 | 1.51 | 0 |
| C10-C11 | 1.44 | 1.48 | 0.04 |
| C11-C12 | 1.34 | 1.35 | 0.01 |
| C12-C13 | 1.46 | 1.49 | 0.03 |
| C13-C14 | 1.34 | 1.35 | 0.01 |
| C13-C20 | 1.49 | 1.51 | 0.02 |
| C14-C15 | 1.47 | 1.49 | 0.02 |
| C15-O29 | 1.31 | 1.26 | 0.05 |
| C15-O30 | 1.21 | 1.26 | 0.05 |
| O29-030 | 2.66 | 2.22 | 0.44 |
| Angle ( ${ }^{\circ}$ ) |  |  |  |
| C2-C1-C6 | 110.2 | 112.5 | 2.3 |
| C2-C1-C16 | 105.4 | 106.9 | 1.5 |
| C2-C1-C17 | 110.9 | 109.1 | 1.8 |
| C6-C1-C16 | 111.5 | 111.9 | 0.4 |
| C6-C1-C17 | 109.1 | 108 | 1.1 |
| C16-C1-C17 | 109.8 | 108.3 | 1.5 |
| C1-C2-C3 | 116.3 | 112.4 | 3.9 |
| C2-C3-C4 | 112.5 | 108.9 | 3.6 |
| C3-C4-C5 | 113.6 | 113.5 | 0.1 |
| C4-C5-C6 | 123 | 122.8 | 0.2 |
| C4-C5-C18 | 112.3 | 114.5 | 2.2 |
| C6-C5-C18 | 124.6 | 122.6 | 2 |
| C1-C6-C5 | 122.1 | 121.9 | 0.2 |
| C1-C6-C7 | 115 | 117.1 | 2.1 |
| C5-C6-C7 | 122.9 | 121 | 1.9 |
| C6-C7-C8 | 127.7 | 125.9 | 1.8 |
| C7-C8-C9 | 125.6 | 124.8 | 0.8 |
| C8-C9-C10 | 118.6 | 123.5 | 4.9 |
| C8-C9-C19 | 118.3 | 112.7 | 5.6 |

## Table B. 3 (cont'd).

| C10-C9-C19 | 123.2 | 123.7 | 0.5 |
| :--- | ---: | ---: | ---: |
| C9-C10-C11 | 126.2 | 125.7 | 0.5 |
| C10-C11-C12 | 124.5 | 123.9 | 0.6 |
| C11-C12-C13 | 126.1 | 126.3 | 0.2 |
| C12-C13-C14 | 118.4 | 121.2 | 2.8 |
| C12-C13-C20 | 117.9 | 115.5 | 2.4 |
| C14-C13-C20 | 123.7 | 123.4 | 0.3 |
| C13-C14-C15 | 126.5 | 129.9 | 3.4 |
| C14-C15-O30 | 112.3 | 114.8 | 2.5 |
| C14-C15-O30 | 126 | 121.7 | 4.3 |
| O29-C15-O30 | 121.7 | 123 | 1.3 |

Dihedral ( ${ }^{\circ}$ )

| C1-C2-C3-C4 | -60 | -60.8 | 0.8 |
| :--- | ---: | ---: | ---: |
| C2-C3-C4-C5 | -45 | 48.7 | 93.7 |
| C3-C4-C5-C6 | -11.5 | -20.3 | 8.8 |
| C4-C5-C6-C1 | 9.74 | 1.45 | 8.29 |
| C5-C6-C1-C2 | -11 | -12.3 | 1.3 |
| C6-C1-C2-C3 | 43.4 | 42.3 | 1.1 |
| C18-C5-C4-C3 | 175 | 177.9 | 2.9 |
| C18-C5-C6-C1 | 165.4 | 163 | 2.4 |
| C18-C5-C6-C7 | -4 | -1.69 | 2.31 |
| C16-C1-C6-C5 | 102.5 | 108.2 | 5.7 |
| C16-C1-C6-C7 | -108 | -72.2 | 35.8 |
| C16-C1-C2-C3 | -77 | -77.5 | 0.5 |
| C17-C1-C6-C5 | -126 | -132.7 | 6.7 |
| C17-C1-C6-C7 | 45.4 | 46.9 | 1.5 |
| C17-C1-C2-C3 | 168.1 | 165.6 | 2.5 |
| C1-C6-C7-C8 | 132.9 | 115.5 | 17.4 |
| C5-C6-C7-C8 | -63.7 | -64.7 | 1 |
| C6-C7-C8-C9 | -180 | -170.3 | 9.7 |
| C7-C8-C9-C19 | 0 | -38.3 | 38.3 |
| C7-C8-C9-C10 | 178.5 | 140.3 | 38.2 |
| C8-C9-C10-C11 | -177.7 | -178.1 | 0.4 |
| C19-C9-C10-C11 | 0.66 | -1.85 | 2.51 |
| C9-C10-C11-C12 | 179.2 | 178.3 | 0.9 |
| C10-C11-C12-C13 | -179.5 | -176.6 | 2.9 |
| C11-C12-C13-C20 | 1.2 | -33.8 | 35 |
| C11-C12-C13-C14 | -180 | -146.4 | 33.6 |
| C12-C13-C14-C15 | -178 | -179.4 | 1.4 |
| C20-C13-C14-C15 | 0.68 | 0.3 | 0.38 |
| C13-C14-C15-O29 | -7.61 | -11.7 | 4.09 |
| C13-C14-C15-O30 | 176.4 | 175.5 | 0.9 |
|  |  |  |  |

Table B. 4 Comparison of CHARMM (energy-minimized structure) and x-ray data for retinoic acid (monoclinic form)

| Bond (£) | x-ray |
| :--- | ---: |
| C1-C2 | 1.55 |
| C1-C6 | 1.53 |
| C1-C16 | 1.54 |
| C1-C17 | 1.53 |
| C2-C3 | 1.5 |
| C3-C4 | 1.51 |
| C4-C5 | 1.5 |
| C5-C6 | 1.36 |
| C5-C18 | 1.52 |
| C6-C7 | 1.46 |
| C7-C8 | 1.34 |
| C8-C9 | 1.45 |
| C9-C10 | 1.35 |
| C9-C19 | 1.5 |
| C10-C11 | 1.44 |
| C11-C12 | 1.34 |
| C12-C13 | 1.45 |
| C13-C14 | 1.34 |
| C13-C20 | 1.5 |
| C14-C15 | 1.47 |
| C15-O29 | 1.32 |
| C15-O30 | 1.22 |
| O29-O30 | 2.67 |

Angle ( ${ }^{\circ}$ )

| C2-C1-C6 | 110.1 | 112.5 | 2.4 |
| :--- | ---: | ---: | ---: |
| C2-C1-C16 | 104.5 | 106.9 | 2.4 |
| C2-C1-C17 | 110.1 | 109.1 | 1 |
| C6-C1-C16 | 111.1 | 111.9 | 0.8 |
| C6-C1-C17 | 111.3 | 108 | 3.3 |
| C16-C1-C17 | 109.6 | 108.3 | 1.3 |
| C1-C2-C3 | 112.2 | 112.4 | 0.2 |
| C2-C3-C4 | 109.5 | 108.9 | 0.6 |
| C3-C4-C5 | 113.6 | 113.5 | 0.1 |
| C4-C5-C6 | 123.7 | 122.8 | 0.9 |
| C4-C5-C18 | 112.5 | 114.5 | 2 |
| C6-C5-C18 | 123.7 | 122.6 | 1.1 |
| C1-C6-C5 | 121.6 | 121.9 | 0.3 |
| C1-C6-C7 | 120.7 | 117.1 | 3.6 |
| C5-C6-C7 | 117.7 | 121 | 3.3 |
| C6-C7-C8 | 131.5 | 125.9 | 5.6 |
| C7-C8-C9 | 125 | 124.8 | 0.2 |
| C8-C9-C10 | 119.3 | 123.5 | 4.2 |
| C8-C9-C19 | 118.4 | 112.7 | 5.7 |

## Table B. 4 (cont'd)

| C10-C9-C19 | 122.3 | 123.7 | 1.4 |
| :---: | :---: | :---: | :---: |
| C9-C10-C11 | 127.3 | 125.7 | 1.6 |
| C10-C11-C12 | 123.4 | 123.9 | 0.5 |
| C11-C12-C13 | 126.6 | 126.3 | 0.3 |
| C12-C13-C14 | 117.4 | 121.2 | 3.8 |
| C12-C13-C20 | 118.5 | 115.5 | 3 |
| C14-C13-C20 | 124.1 | 123.4 | 0.7 |
| C13-C14-C15 | 128.8 | 129.9 | 1.1 |
| Dihedral ( ${ }^{\circ}$ ) |  |  |  |
| C14-C15-O29 | 112.1 | 114.8 | 2.7 |
| C14-C15-O30 | 125.9 | 121.7 | 4.2 |
| O29-C15-O30 | 122 | 123 | 1 |
| C1-C2-C3-C4 | -63.3 | -60.8 | 2.5 |
| C2-C3-C4-C5 | 43.2 | 48.7 | 5.5 |
| C3-C4-C5-C6 | -10.4 | -20.3 | 9.9 |
| C4-C5-C6-C1 | -4.8 | 1.45 | 6.25 |
| C5-C6-C1-C2 | -13.4 | -12.3 | 1.1 |
| C6-C1-C2-C3 | 47.5 | 42.3 | 5.2 |
| C18-C5-C6-C1 | 175.1 | 177.9 | 2.8 |
| C18-C5-C4-C3 | 169.6 | 163 | 6.6 |
| C18-C5-C6-C7 | -3.1 | -1.69 | 1.41 |
| C16-C1-C6-C5 | 108.9 | 108.2 | 0.7 |
| C16-C1-C6-C7 | -72.9 | -72.2 | 0.7 |
| C16-C1-C2-C3 | -75.5 | -77.5 | 2 |
| C17-C1-C6-C5 | -128.7 | -132.7 | 4 |
| C17-C1-C6-C7 | 49.5 | 46.9 | 2.6 |
| C17-C1-C2-C3 | 167 | 165.6 | 1.4 |
| C1-C6-C7-C8 | 15.9 | 115.5 | 99.6 |
| C5-C6-C7-C8 | -165.8 | -64.7 | 101.1 |
| C6-C7-C8-C9 | 179.6 | 179.6 | 0 |
| C7-C8-C9-C10 | 176.2 | 140.3 | 35.9 |
| C7-C8-C9-C19 | -1.9 | -38.1 | 36.2 |
| C8-C9-C10-C11 | -178.2 | -178.1 | 0.1 |
| C19-C9-C10-C11 | -0.13 | -1.85 | 1.72 |
| C9-C10-C11-C12 | 172 | 178.3 | 6.3 |
| C10-C11-C12-C13 | -177.8 | -176.6 | 1.2 |
| C11-C12-C13-C14 | 179.4 | 146 | 33.4 |
| C11-C12-C13-C20 | 1.9 | 33.8 | 31.9 |
| C12-C13-C14-C15 | -179.8 | -179.4 | 0.4 |
| C15-C14-C13-C20 | -2.5 | 0.3 | 2.8 |
| O30-C15-C14-C13 | -6.9 | -11.7 | 4.8 |
| O29-C15-C14-C13 | 173.7 | 175.5 | 1.8 |




Figure B. 5 Overlay of cyclohexene rings (left) and overlay of chains (right). X-ray structure is shown in red; average structure from the simulation is shown in blue. Bottom picture shows the CHARMM atom names of the carbons in the retinoic acid ligand, and it shows the dihedrals that contribute to the large average dihedral error.

The chain rotates with respect to the cyclohexene ring during the simulation, which causes large deviations in the torsion angles that are involved around the connection of the ring to the chain (rotation about the C6-C7 bond). The other sources of large RMSD differences involve parts of the chain (rotation about $\mathrm{C} 8-\mathrm{C} 9, \mathrm{C} 10-\mathrm{C} 11$, and $\mathrm{C} 12-\mathrm{C} 13$ ). These differences are seen during the simulation because these are single bonds, and the ligand will probably try to orient itself to optimize contacts with the protein side chains, and this does not appear to cause any major distortion of the surrounding protein residues. The cyclohexene ring retains the $\mathrm{C}^{\prime}$ '-endo conformation during the simulation.

The protein-ligand contact time series plots not shown in Chapter 4 are shown below in Figures B. 6 and B.7.

Ala32 (CB) to RTAC (C8)

distance vs. time (leu28--rtac)


Ala35 (CB) to RTAC (C19)


Figure B. 6 Retinoic acid-protein distances.

## Leu19 (CD2) to RTAC (C16)



Ile9 (CD) to RTAC (C20)


Figure B. 7 Retinoic acid-protein distances.

The protein-protein contacts not shown in Chapter 4 are shown in Figure B.8.

Glu73 (OE2) to Gln97 (NE2)


Ser4 (O) to Arg136 (NE)


Asn2 (O) to Arg136 (NH2)


Figure B. 8 Protein-protein distances from CRABPII/retinoic acid simulation.

## Additional information and plots for CRBP/retinol simulation

The protein-ligand distance plots not shown in Chapter 4 are shown in Figures
B.9-B. 12 .


Figure B. 9 Protein-ligand distances from CRBP/retinol simulation.


Figure B.10 Protein-ligand distances from CRBP/retinol simulation.


Figure B. 11 Protein-ligand distances from CRBP/retinol simulation.


Figure B. 12 Protein-ligand distances from CRBP/retinol simulation.

## Additional information and plots for RBP/fenretinide simulation

The plots shown are for the ligand-protein distances not shown in Chapter 4 (B.13-B.17).


Ala55 (CB) to FRET (C4)


Val61 (CG1) to FRET (O29)


Figure B. 13 Protein-ligand distances from RBP/fenretinide simulation.


Figure B. 14 Protein-ligand distances from RBP/fenretinide simulation.


Figure B. 15 Protein-ligand distances from RBP/fenretinide simulation.


Figure B. 16 Protein-ligand distances from RBP/fenretinide simulation.


Figure B. 17 Protein-ligand distances from RBP/fenretinide simulation.

Additional information and plots for SRII/retinal simulation
The plots below are the ligand-protein distances not shown in Chapter 4.

Gly130 ( C) to RTAL (C4)


Val108 (CG1) to RTAL (C7)


Pro175 (CB) to RTAL (C4)


Figure B. 18 Protein-ligand distances from SRII/retinal simulation.


Figure B. 19 Protein-ligand distances from SRII/retinal simulation.


Figure B. 20 Protein-ligand distances from SRII/retinal simulation.

Thr204 (OG1) to RTAL (C15)


Phe134 (CG) to RTAL (C18)


Leu170 (O) to RTAL (C19)


Figure B. 21 Protein-ligand distances from SRII/retinal simulation.

## Additional information and plots for BR/retinal simulation

The plots below are the ligand-protein distances not shown in Chapter 4 (B.22-

## B.23).



Figure B. 22 Protein-ligand distances from BR/retinal simulation.


Figure B. 23 Protein-ligand distances from $\mathrm{BR} /$ retinal simulation.

## CHARMM Topology and Parameters for Retinoids and Retinoid Model Compounds

Topology and parameters for retinoids were streamed into topology and parameters for proteins (top_all27_prot_na.inp and par_all27_prot_na.inp with CMAP backbone dihedral correction). The CHARMM atom names and numbering used in plots and tables can also be found in the topology.

Topology for 1

```
RESI 13DB 0.00 ! 1,3-dibutene
!
GROUP
ATOM C1 CC2 -0.42 ! H11 H21
ATOM H11 HE2 0.21 ! \ / 
ATOM H12 HE2 0.21 ! C1=C2 H41
GROUP
ATOM C2 CC1A -0.15 ! H12 C3=C4
ATOM H21 HE1 0.15 ! / \
ATOM C3 CC1B -0.15 ! H31 H42
ATOM H31 HE1 0.15
GROUP
ATOM C4 CC2 -0.42
ATOM H41 HE2 0.21 !
ATOM H42 HE2 0.21 !
BOND C1 C2 C2 C3 C3 C4
BOND C1 H11 C1 H12 C2 H21 C3 H31
BOND C4 H41 C4 H42
IC C1 C2 C3 C4 0.00 0.00 180.0 0.00 0.00
IC C2 C3 C4 H41 0.00 0.00 180.0 0.00 0.00
IC C3 C2 C1 H11 0.00 0.00 180.0 0.00 0.00
IC C2 H11 *C1 H12 0.00 0.00 180.0 0.00 0.00
IC C3 H41 *C4 H42 0.00 0.00 180.0 0.00 0.00
IC C1 C3 *C2 H21 0.00 0.00 180.0 0.00 0.00
IC C2 C4 *C3 H31 0.00 0.00 180.0 0.00 0.00
PATCH FIRST NONE LAST NONE
```


## Topology for 2



```
GROUP
```



## Topology for 3

```
RESI DMB1 0.00 ! 2-methyl-1,3-butadiene
!
GROUP
ATOM C1 CC2 -0.42 !
ATOM H11 HE2 0.21 ! H2M1 H2M2
ATOM H12 HE2 0.21 ! \ /
GROUP ! H11 C2M-H2M3
ATOM C2 CC1A 0.00 ! \ C1=C2 /
ATOM H2M1 HA rro.09 ! / \ / N41
ATOM H2M2 HA 0.09 ! H12 C3=C4
ATOM H2M3 HA 0.09 ! / \
GROUP ! H31 H42
ATOM C3 CC1B -0.15 !
ATOM H31 HE1 0.15
GROUP
ATOM C4 CC2 -0.42 !
ATOM H41 HE2 0.21 !
ATOM H42 HE2 0.21 !
\begin{tabular}{lllllllllll} 
BOND & C1 & C2 & C2 & C3 & C3 & C4 & C2 & C2M & \\
BOND & C1 & H11 & C1 & H12 & C3 & H31 & C4 & H41 & C4 \\
BOND & C2M & H2M1 & C2M & H2M2 & C2M & H2M3 & & & \\
IC & C1 & C2 & C3 & C4 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C2 & C3 & C4 & H41 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C3 & C2 & C1 & H11 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C2 & H11 & *C1 & H12 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C1 & C3 & *C2 & C2M & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C2 & C4 & *C3 & H31 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C3 & H41 & *C4 & H42 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C1 & C2 & C2M & H2M1 & 0.00 & 0.00 & 180.0 & 0.00 & 0.00 \\
IC & C2 & H2M1 & *C2M & H2M2 & 0.00 & 0.00 & 120.0 & 0.00 & 0.00 \\
IC & C2 & H2M1 & *C2M & H2M3 & 0.00 & 0.00 & -120.0 & 0.00 & 0.00
\end{tabular}
```

PATCH FIRST NONE LAST NONE
Topology for 4


## Topology for 5

```
RESI DMP2 0.00 ! 2-methyl-1,3-pentadiene
!
llllll
ATOM H31 HE1 0.15 !
GROUP
ATOM C4 CC1B -0.15 !
ATOM H41 HE1 0.15 !
ATOM C5 CT3 -0.27
ATOM H51 HA 0.09 !
ATOM H52 HA 0.09 !
ATOM H53 HA 0.09 !
BOND C1 C2 C2 C3 C3 C4 C4 C5 C5 C2 C2M
BOND C1 H11 C1 H12 C3 H31 C4 H41
BOND C2M H2M1 C2M H2M2 C2M H2M3
BOND C5 H51 C5 H52 C5 H53
IC C1 C2 C3 C4 0.00 0.00 180.0 0.00 0.00
IC C2 C3 C4 C5 0.00 0.00 180.0 0.00 0.00
IC C3 C4 C5 H51 0.00 0.00 180.0 0.00 0.00
IC C3 C2 C1 H11 0.00 0.00 180.0 0.00 0.00
IC C2 H11 *C1 H12 0.00 0.00 180.0 0.00 0.00
IC C1 C3 *C2 C2M 0.00 0.00 180.0 0.00 0.00
IC C2 C4 *C3 H31 0.00 0.00 180.0 0.00 0.00
IC C3 C5 *C4 H41 0.00 0.00 180.0 0.00 0.00
IC C1 C2 C2M H2M1 0.00 0.00 180.0 0.00 0.00
IC C2 H2M1 *C2M H2M2 0.00 0.00 120.0 0.00 0.00
IC C2 H2M1 *C2M H2M3 0.00 0.00 -120.0 0.00 0.00
IC C4 H51 *C5 H52 0.00 0.00 120.0 0.00 0.00
IC C4 H51 *C5 H53 0.00 0.00 -120.0 0.00 0.00
PATCH FIRST NONE LAST NONE
```


## Topology for 6



PATCH FIRST NONE LAST NONE

Topology for 7


PATCH FIRST NONE LAST NONE
Topology for 8


| IC | C1 | C5 | *C6 | C7 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | ---: | :--- | :--- | :--- |
| IC | C5 | C6 | C7 | C8 | 0.00 | 0.00 | 60.0 | 0.00 | 0.00 |
| IC | C6 | C8 | *C7 | H71 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C6 | C7 | C8 | H81 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C7 | H81 | *C8 | H82 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C6 | C2 | *C1 | C16 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C6 | C2 | *C1 | C17 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C2 | C1 | C16 | H161 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C1 | H161 | *C16 | H162 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C1 | H161 | *C16 | H163 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C2 | C1 | C17 | H171 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C1 | H171 | *C17 | H172 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C1 | H171 | *C17 | H173 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C4 | C6 | *C5 | C18 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C4 | C5 | C18 | H181 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C5 | H181 | *C18 | H182 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C5 | H181 | *C18 | H183 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| PATCH | FIRST | NONE | LAST | NONE |  |  |  |  |  |

## Topology for 9



| BOND | C1 | C2 | C2 | C3 | C3 | C4 | C4 | C5 | C5 | C6 | C6 | C1 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| BOND | C6 | C7 | C1 | C16 | C1 | C17 | C5 | C18 |  |  |  |  |
| BOND | C2 | H21 | C2 | H22 | C3 | H31 | C3 | H32 | C4 | H41 | C4 | H42 |
| BOND | C16 | H161 | C16 | H162 | C16 | H163 | C17 | H171 | C17 | H172 | C17 | H173 |
| BOND | C7 | H71 | C7 | H72 | C7 | H73 | C18 | H181 | C18 | H182 | C18 | H183 |

!initial geometry is planar

| IC | C1 | C2 | C3 | C4 | 0.00 | 0.00 | -63.0 | 0.00 | 0.00 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC | C2 | C3 | C4 | C5 | 0.00 | 0.00 | 46.0 | 0.00 | 0.00 |
| IC | C3 | C4 | C5 | C6 | 0.00 | 0.00 | -13.0 | 0.00 | 0.00 |
| IC | C1 | C3 | ${ }^{*} \mathrm{C} 2$ | H21 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C1 | C3 | ${ }^{*} \mathrm{C} 2$ | H22 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C2 | C4 | ${ }^{*} \mathrm{C} 3$ | H31 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C2 | C4 | ${ }^{*} \mathrm{C} 3$ | H32 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C3 | C5 | ${ }^{*} \mathrm{C} 4$ | H41 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C3 | C5 | * C 4 | H42 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C1 | C5 | * C 6 | C7 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C5 | C6 | C7 | H71 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C6 | H71 | * C 7 | H72 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C6 | H71 | * C 7 | H73 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C6 | C2 | ${ }^{+} \mathrm{C} 1$ | C16 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C6 | C2 | ${ }^{+} \mathrm{C} 1$ | C17 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C2 | C1 | C16 | H161 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C1 | H161 | *C16 | H162 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C1 | H161 | *C16 | H163 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C2 | C1 | C17 | H171 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C1 | H171 | *C17 | H172 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C1 | H171 | *C17 | H173 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |
| IC | C4 | C6 | * C 5 | C18 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C4 | C5 | C18 | H181 | 0.00 | 0.00 | 180.0 | 0.00 | 0.00 |
| IC | C5 | H181 | *C18 | H182 | 0.00 | 0.00 | 120.0 | 0.00 | 0.00 |
| IC | C5 | H181 | *C18 | H183 | 0.00 | 0.00 | -120.0 | 0.00 | 0.00 |

PATCH FIRST NONE LAST NONE

## Topology for 10

| RESI PRAC -1.00 |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| GROUP |  |  |  |  |  |  |  |  |
| ATOM C1 | CC2 | -0.42 | ! |  |  |  |  |  |
| ATOM H11 | HE2 | 0.21 | ! |  |  |  |  |  |
| ATOM H12 | HE2 | 0.21 | ! H11 | 1 OD1 |  |  |  |  |
| GROUP |  | ! | $\backslash$ | \| | |  |  |  |  |
| ATOM C2 | CC1A | -0.25 | $!\mathrm{C} 1$ | 1 CG |  |  |  |  |
| ATOM H21 | HE1 | 0.15 | ! / |  |  |  |  |  |
| / \} |  |  |  |  |  |  |  |  |
| ATOM CG | CC | 0.62 | ! H12 | C2 OD2 | (-) |  |  |  |
| ATOM OD1 | OC | -0.76 | ! | \| |  |  |  |  |
| ATOM OD2 | OC | -0.76 | ! | H21 |  |  |  |  |
| BOND C1 | H11 | C1 H12 | C1 | C2 C2 | H21 |  |  |  |
| BOND C2 | CG | CG OD1 | CG | OD2 |  |  |  |  |
| IC C1 | C2 | CG | OD1 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC OD1 | C2 | * CG | OD2 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H11 | C1 | C2 | CG | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H11 | C2 | * C 1 | H12 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC CG | C1 | ${ }^{*} \mathrm{C} 2$ | H21 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |

PATCH FIRST NONE LAST NONE

## Topology for 11



## Topology for 12

| RESI PRAL |  | 0.00 ! 1-propenal |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| GROUP |  |  |  |  |  |  |  |  |  |  |
| ATOM C1 | CC2 | -0.42 | ! |  |  |  |  |  |  |  |
| ATOM H11 | HE2 | 0.21 | ! H1 | OD1 |  |  |  |  |  |  |
| ATOM H12 | HE2 | 0.21 | ! $\$ & \| | & & & &  \hline ATOM C2 & CC1A & -0.01 & ! C & CG & & & &  \hline ATOM H21 & HE1 & 0.15 & ! / & / \} & & & &  \hline ATOM CG & $C D$ | 0.19 | ! H12 | C2 H |  |  |  |  |
| ATOM OD1 | $\bigcirc$ | -0.43 | ! | \| |  |  |  |  |  |  |
| ATOM H | HR1 | 0.10 | ! | H21 |  |  |  |  |  |  |
| BOND C1 | H11 | C1 H12 | 2 C 1 | C2 C2 | H21 |  |  |  |  |  |
| BOND C2 | CG | CG OD1 | 1 CG | H |  |  |  |  |  |  |
| IMPR CG | C2 | OD1 H |  |  |  |  |  |  |  |  |
| IC C1 | C2 | CG | OD1 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC H11 | C1 | C2 | CG | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC H11 | C2 | * C 1 | H12 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC CG | C1 | * C 2 | H21 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC OD1 | C2 | *CG | H | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |

## Topology for 13



PATCH FIRS NONE LAST NONE

## Topology for 14



GROUP
$\begin{array}{llr}\text { ATOM C6 } & \text { CC2 } & -0.42 \\ \text { ATOM H61 } & \text { HE2 } & 0.21 \\ \text { ATOM H62 } & \text { HE2 } & 0.21\end{array}$

| BOND | C1 | C2 | C2 | C3 | C3 | C4 | C4 | C5 | C5 | C6 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| BOND | C1 | H11 | C1 | H12 | C2 | H21 | C3 | H31 |  |  |
| BOND | C4 | H41 | C5 | H51 | C6 | H61 | C6 | H62 |  |  |


| IC | C1 | C2 | C3 | C4 | 0.0000 | 0.00 | 180.00 | 0.00 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | C2 | C3 | C4 | C5 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | C3 | C4 | C5 | C6 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | H11 | C2 | *C1 | H12 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | H11 | C1 | C2 | C3 | 0.0000 | 0.00 | 180.00 | 0.000 |
| IC | C3 | C1 | *C2 | H21 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | C4 | C2 | *C3 | H31 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | C5 | C3 | *C4 | H41 | 0.0000 | 0.0000 |  |  |
| IC C6 | C4 | *C5 | H51 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C4 | C5 | C6 | H61 | 0.0000 | 180.00 | 0.00 | 0.0000 |
| IC | H61 | C5 | *C6 | H62 | 0.0000 | 0.00 | 180.00 | 0.00 |

PATCH FIRST NONE LAST NONE

## Topology for 15



## Topology for 16

| RESI SCH2 ! |  | 1.00 | ! Schiff's base model |  |  | compoun | protonated |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| GROUP |  |  |  |  |  |  |  |  |
| ATOM C1 | CT3 | 0.18 | ! H | H11 H21 | H31 |  |  |  |
| ATOM H11 | HA | 0.09 | . | $\backslash$ | / |  |  |  |
| ATOM H12 | HA | 0.09 | ! H1 | $2--\mathrm{C} 1-\mathrm{N} 2$ | C3 H |  |  |  |
| ATOM H13 | HA | 0.09 | ! | / | $\backslash /$ |  |  |  |
| ATOM N2 | NS2 | -0.40 | ! | 13 | C4-- | H2 |  |  |
| ATOM H21 | HC | 0.38 | ! |  | $\backslash$ |  |  |  |
| ATOM C3 | CC1B | 0.37 | ! |  |  | 3 |  |  |
| ATOM H31 | HR1 | 0.20 | ! hyd | rogen wi | interm | ediate | radi |  |
| GROUP |  | $!$ | ! |  |  |  |  |  |
| ATOM C4 | CT3 | -0.27 | . |  |  |  |  |  |
| ATOM H41 | HA | 0.09 | ! |  |  |  |  |  |
| ATOM H42 | HA | 0.09 |  |  |  |  |  |  |
| ATOM H43 | HA | 0.09 |  |  |  |  |  |  |
| BOND C1 | H11 | C1 H12 | C1 | H13 |  |  |  |  |
| BOND C1 | N2 | N2 H21 | N2 | C3 C3 | H31 |  |  |  |
| BOND C3 | C4 | C4 H41 | C4 | H42 C4 | H43 |  |  |  |
| IC C1 | N2 | C3 | C4 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC N2 | C3 | C4 | H41 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C1 | C3 | *N2 | H21 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C3 | N2 | C1 | H11 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| IC H11 | N2 | * C 1 | H12 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H11 | N2 | * C 1 | H13 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C4 | N2 | * C 3 | H31 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H41 | C3 | * C 4 | H42 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H41 | C3 | *C4 | H43 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |

## Topology for 17

```
RESI SCH3 1.00 ! Schiff's base model compound 3, protonated
! ! new atom type, NS2
GROUP
```



```
ATOM H31 HR1 0.20 ! hydrogen with intermediate VDW radius
GROUP !
ATOM C4 CC1A -0.15
ATOM H41 HE1 0.15
ATOM C5 CC2 -0.42
ATOM H51 HE2 0.21
ATOM H52 HE2 0.21
```

| BOND C1 | H11 | C1 | H12 | C1 | H13 |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| BOND C1 | N2 | N2 | H21 | N2 | C3 | C3 | H31 |  |  |  |
| BOND C3 | C4 | C4 | H41 |  |  |  |  |  |  |  |
| BOND C4 | C5 | C5 | H51 | C5 | H52 |  |  |  |  |  |
| IC C1 | N2 | C3 | C4 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC N2 | C3 | C4 | C5 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC H11 | C1 | N2 | C3 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC H11 | N2 | *C1 | H12 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |  |  |
| IC H11 | N2 | *C1 | H13 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |  |  |
| IC C3 | C1 | *N2 | H21 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC C4 | N2 | *C3 | H31 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC C5 | C3 | *C4 | H41 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |  |
| IC | C3 | C4 | C5 | H51 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |
| IC | H51 | C4 | *C5 | H52 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |  |

## Topology for 18

```
PRES SCK0 0.00 ! patch to link lysine with retinal to form a
    ! deprotonated schiff's base
    ! follow with AUTOgenerate ANGLes DIHEdrals
command
! residue 1: lysine
! residue 2: retinal
!
DELETE ATOM 1HZ1
DELETE ATOM 1HZ2
DELETE ATOM 1HZ3
DELETE ATOM 2015
\begin{tabular}{|c|c|c|c|c|c|}
\hline GROUP & & & & HE1 HE2 & \\
\hline ATOM 1CE & CT2 & 0.04 & ! & \ / & \\
\hline ATOM 1HE1 & HA & 0.09 & ! & --1CE1 & 2H15 \\
\hline ATOM 1HE2 & HA & 0.09 & ! & \(\backslash\) & / \\
\hline ATOM 1NZ & NS1 & -0.60 & ! & & \\
\hline ATOM 2C15 & CC1B & 0.23 & ! & / & \(\backslash\) \\
\hline ATOM 2H15 & HE1 & 0.15 & ! & 2HZ1 & 2 C \\
\hline GROUP & & & & & / \\
\hline ATOM 2C14 & CC1A & -0.15 & & & E14 \\
\hline
\end{tabular}
ATOM 2H141 HE1 0.15
```

BOND 1NZ 2C15

| IC | 1 CD | 1 CE | 1 NZ | 2 C 15 | 0.0000 | 0.00 | 180.00 | 0.00 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | 1CE | 1NZ | 2C15 | 2C14 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | 1NZ | 2C15 | 2C14 | 2C13 | 0.0000 | 0.00 | 180.00 | 0.00 |
| IC | 2C14 | 1NZ | *2C15 | 2H15 | 0.0000 | 0.0000 |  |  |
| IC | 2C5 | 2C6 | 2C1 | 2C2 | 0.0000 | 0.00 | 180.00 | 0.00 |
| 1 | 0.000 | 0.00 | 0.0000 |  |  |  |  |  |

! required to build retinal

## Topology for 19

```
PRES SCK1 1.00 ! patch to link lysine with retinal to form a
    ! protonated schiff's base
    ! follow with AUTOgenerate ANGLes DIHEdrals
command
! residue 1: lysine
! residue 2: retinal
!
DELETE ATOM 1HZ2
DELETE ATOM 1HZ3
DELETE ATOM 2O15
GROUP ! HE1 HE2
ATOM 1CE CT2 0.27 ! \ /
ATOM 1HE1 HA 0.09 ! --1CE1 2H15
ATOM 1HE2 HA 0.09 ! \
ATOM 1NZ NS2 -0.40! 1NZ=2C15
ATOM HZ1 HC 0.38 ! / \
ATOM 2C15 CC1B 0.37 ! 2HZ1 2C14==
ATOM 2H15 HR1 0.20 ! /
GROUP ! 2HE14
ATOM 2C14 CC1A -0.15
ATOM 2H141 HE1 0.15
BOND 1NZ 2C15
\begin{tabular}{|c|c|c|c|c|c|c|c|c|}
\hline IC 1CD & 1 CE & 1NZ & 2 C 15 & 0.0000 & 0.00 & 180.00 & 0.00 & 0.0000 \\
\hline IC 1CE & 1NZ & 2 C 15 & 2 C 14 & 0.0000 & 0.00 & 180.00 & 0.00 & 0.0000 \\
\hline IC 1NZ & 2C15 & 2 C 14 & 2 C 13 & 0.0000 & 0.00 & 180.00 & 0.00 & 0.0000 \\
\hline IC 2C15 & 1 CE & *1NZ & 1HZ1 & 0.0000 & 0.00 & 180.00 & 0.00 & 0.0000 \\
\hline IC 2C14 & 1 NZ & *2C15 & 2H15 & 0.0000 & 0.00 & 180.00 & 0.00 & 0.0000 \\
\hline IC 2C5 & 2C6 & 2C1 & 2 C 2 & 0.0000 & 0.00 & 0.00 & 0.00 & 0.0000 \\
\hline
\end{tabular}
! required to build retinal
end
```


## Topology for fenretinide

```
RESI FRET 0.00 ! fenretinide, nomenclature consistent with
retinol
!
!
GROUP
ATOM C1 CT 0.00
ATOM C2 CT2 -0.18
ATOM H21 HA 0.09
ATOM H22 HA 0.09
ATOM C3 CT2 -0.18
ATOM H31 HA 0.09
ATOM H32 HA 0.09
ATOM C4 CT2 -0.18
ATOM H41 HA 0.09
```

```
ATOM H42 HA 0.09
ATOM C5 CC1A 0.00
ATOM C6 CC1A 0.00
GROUP
ATOM C7 CC1B -0.15
ATOM H71 HE1 0.15
ATOM C8 CC1B -0.15
ATOM H81 HE1 0.15
ATOM C9 CC1A 0.00
ATOM C10 CC1A -0.15
ATOM H101 HE1 0.15
GROUP
ATOM C11 CC1B -0.15 !H162 H163 H171 H172
ATOM H111 HE1 0.15 ! \ | | /
ATOM C12 CC1B -0.15 !H161-C16 C17-H173 H191 H192
ATOM H121 HE1 0.15 ! \ / \ /
ATOM C13 CC1A 0.00 ! H21 C1 H71 H81 C19
ATOM C14 CC1A -0.15 ! \ / \ | | |
ATOM H141 HE1 0.15 ! H22-C2 C6------C7======C8----C9 H193
GROUP ! | || | |
ATOM C16 CT3 -0.27 ! H31-C3 C5 H181 C10-H101
ATOM H161 HA 0.09 ! / \ / \ /
```



```
    C11-H111
```



```
| |
ATOM C17 CT3 -0.27
ATOM H171 HA 0.09
ATOM H172 HA 0.09
ATOM H173 HA 0.09
GROUP !
ATOM C18 CT3 -0.27 !
ATOM H181 HA 0.09
ATOM H182 HA 0.09
ATOM H183 HA 0.09
GROUP
ATOM C19 CT3 -0.27
ATOM H191 HA 0.09
ATOM H192 HA 0.09
ATOM H193 HA 0.09
GROUP
!
ATOM C20 CT3 -0.27
ATOM H201 HA 0.09 !
ATOM H202 HA 0.09
ATOM H203 HA 0.09
GROUP !
ATOM C15 C 0.51 ! charges from CROT
ATOM O29 O -0.51 !
GROUP
ATOM N21 NH1 -0.47 !
ATOM H211 H 0.33 ! charges from PACP
ATOM C22 CA 0.14 !
GROUP
ATOM C23 CA -0.115
ATOM H231 HP 0.115
GROUP
ATOM C24 CA -0.115
ATOM H241 HP 0.115
```

GROUP
ATOM C26 CA -0.115
ATOM H261 HP 0.115
GROUP
ATOM C27 CA -0.115
ATOM H271 HP 0.115
GROUP
ATOM C25 CA 0.11
ATOM O28 OH1 -0.54
ATOM H281 H 0.43
BOND C1 C2 C2 C3 C3 C4 C4 C5 C5 C6 C6 C1
BOND C6 C7 C7 C8 C8 C9 C9 C10 C10 C11 C11 C12

$\begin{array}{lllllllllllll}\text { BOND } & \text { C23 } & \text { C24 } & \text { C24 } & \text { C25 } & \text { C25 } & \text { C26 } & \text { C26 } & \text { C27 } & \text { C27 } & \text { C22 } & & \\ \text { BOND } & \text { C1 } & \text { C16 } & \text { C1 } & \text { C17 } & \text { C5 } & \text { C18 } & \text { C9 } & \text { C19 } & \text { C13 } & \text { C20 } & \text { C15 } & \text { O29 }\end{array}$
BOND C25 O28 O28 H281 N21 H211
BOND C2 H21 C2 H22 C3 H31 C3 H32 C4 H41 C4 H42
BOND C7 H71 C8 H81 C10 H101 C11 H111 C12 H121 C14 H141
BOND C16 H161 C16 H162 C16 H163 C17 H171 C17 H172 C17 H173
BOND C18 H181 C18 H182 C18 H183 C19 H191 C19 H192 C19 H193
BOND C20 H201 C20 H202 C20 H203
BOND C23 H231 C24 H241 C26 H261 C27 H271

| IC C1 | C2 | C3 | C4 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC C2 | C3 | C4 | C5 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C2 | C3 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| IC C1 | C6 | C7 | C8 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C6 | C7 | C8 | C9 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C7 | C8 | C9 | C10 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C2 | C6 | * C 1 | C16 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C2 | C6 | * C 1 | C17 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C3 | C1 | * C 2 | H21 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C3 | C1 | * C 2 | H22 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C4 | C2 | * C 3 | H31 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C4 | C2 | * C 3 | H32 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C5 | C3 | ${ }^{*} \mathrm{C} 4$ | H41 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C5 | C3 | * C 4 | H42 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C5 | C1 | * C 6 | C7 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C8 | C6 | * C 7 | H71 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C9 | C7 | * C 8 | H81 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C10 | C8 | *C9 | C19 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C8 | C9 | C10 | C11 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C11 | C9 | *C10 | H101 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C9 | C10 | C11 | C12 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C12 | C10 | *C11 | H111 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C10 | C11 | C12 | C13 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C13 | C11 | * C12 | H121 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C11 | C12 | C13 | C14 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C14 | C12 | * C13 | C20 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C12 | C13 | C14 | C15 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C15 | C13 | * C14 | H141 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C16 | H161 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H161 | C1 | * C16 | H162 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H161 | C1 | *C16 | H163 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C17 | H171 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H171 | C1 | * C17 | H172 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H171 | C1 | * C17 | H173 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C6 | C4 | * C 5 | C18 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |


| IC | C4 | C5 | C18 | H181 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | H181 | C5 | *C18 | H182 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC | H181 | C5 | *C18 | H183 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC | C8 | C9 | C19 | H191 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | H191 | C9 | *C19 | H192 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC | H191 | C9 | *C19 | H193 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC | C12 | C13 | C20 | H201 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | H201 | C13 | *C20 | H202 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC | H201 | C13 | *C20 | H203 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC | C13 | C14 | C15 | N21 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | N21 | C14 | *C15 | O29 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C14 | C15 | N21 | C22 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C22 | C15 | *N21 | H211 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C15 | N21 | C22 | C27 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C27 | N21 | *C22 | C23 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | N21 | C22 | C23 | C24 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C24 | C22 | *C23 | H231 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C22 | C23 | C24 | C25 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| IC | C25 | C23 | *C24 | H241 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C23 | C24 | C25 | O28 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | O28 | C24 | *C25 | C26 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C27 | C25 | *C26 | H261 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C26 | C22 | *C27 | H271 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C24 | C25 | O28 | H281 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |

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## Topology for retinol




| IC C2 | C6 | ${ }^{*} \mathrm{C} 1$ | C16 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC C2 | C6 | ${ }^{*} \mathrm{C} 1$ | C17 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C3 | C1 | * C 2 | H21 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C3 | C1 | ${ }^{*} \mathrm{C} 2$ | H22 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C4 | C2 | * C 3 | H31 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C4 | C2 | * C 3 | H32 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C5 | C3 | * C 4 | H41 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C5 | C3 | ${ }^{*} \mathrm{C} 4$ | H42 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C5 | C1 | * C 6 | C7 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C8 | C6 | ${ }^{*} \mathrm{C} 7$ | H71 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C9 | C7 | * C 8 | H81 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C10 | C8 | * C 9 | C19 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C8 | C9 | C10 | C11 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C11 | C9 | *C10 | H101 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C9 | C10 | C11 | C12 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C12 | C10 | * C11 | H111 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C10 | C11 | C12 | C13 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C13 | C11 | *C12 | H121 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C11 | C12 | C13 | C14 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C14 | C12 | *C13 | C20 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C12 | C13 | C14 | C15 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C15 | C13 | *C14 | H141 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C13 | C14 | C15 | OR | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC OR | C14 | *C15 | H151 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC OR | C14 | *C15 | H152 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C14 | C15 | OR | HR | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C16 | H161 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H161 | C1 | *C16 | H162 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H161 | C1 | *C16 | H163 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C17 | H171 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H171 | C1 | * C17 | H172 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H171 | C1 | * C17 | H173 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C6 | C4 | * C 5 | C18 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C4 | C5 | C18 | H181 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H181 | C5 | * C 18 | H182 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H181 | C5 | * C 18 | H183 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C8 | C9 | C19 | H191 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H191 | C9 | *C19 | H192 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H191 | C9 | *C19 | H193 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C12 | C13 | C20 | H201 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H201 | C13 | * C 20 | H2 02 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H201 | C13 | ${ }^{*} \mathrm{C} 20$ | H203 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |

PATCH FIRST NONE LAST NONE

## Topology for retinal

| ATOM H193 HA | 0.09 |  |  |
| :--- | :--- | ---: | :--- |
| GROUP |  |  | $!$ |
| ATOM C20 CT3 | -0.27 | ! |  |
| ATOM H201 HA | 0.09 | ! |  |
| ATOM H202 HA | 0.09 | ! |  |
| ATOM H203 HA | 0.09 |  |  |


| BOND | C1 | C2 | C2 | C3 | C3 | C4 | C4 | C5 | C5 | C6 | C6 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| C1 |  |  |  |  |  |  |  |  |  |  |  |
| BOND | C6 | C7 | C7 | C8 | C8 | C9 | C9 | C10 | C10 | C11 | C11 |
| C12 |  |  |  |  |  |  |  |  |  |  |  |
| BOND | C12 | C13 | C13 | C14 | C14 | C15 | C15 | O15 | C15 | H15 |  |
| BOND C1 | C16 | C1 | C17 | C5 | C18 | C9 | C19 | C13 | C20 |  |  |
| BOND | C2 | H21 | C2 | H22 | C3 | H31 | C3 | H32 | C4 | H41 | C4 |
| H42 |  |  |  |  |  |  |  |  |  |  |  |
| BOND C7 | H71 | C8 | H81 | C10 | H101 | C11 | H111 | C12 | H121 | C14 | H141 |
| BOND | C16 | H161 | C16 | H162 | C16 | H163 | C17 | H171 | C17 | H172 | C17 |
| H173 |  |  |  |  |  |  |  |  |  |  |  |
| BOND | C18 | H181 | C18 | H182 | C18 | H183 | C19 | H191 | C19 | H192 | C19 |
| H193 |  |  |  |  |  |  |  |  |  |  |  |

BOND C20 H201 C20 H202 C20 H203

| 1 Cl | C2 | C3 | C4 | 0. | 0.0 | 0.00 | 0. | 0.0000 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| IC C2 | C3 | C4 | C5 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C2 | C3 | 0.0000 | 0.00 | 0.00 | 0.00 | 0.0000 |
| IC C1 | C6 | C7 | C8 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C6 | C7 | C8 | C9 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C7 | C8 | C9 | C10 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C8 | C9 | C10 | C11 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C2 | C6 | ${ }^{*} \mathrm{C} 1$ | C16 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C2 | C6 | ${ }^{+} \mathrm{C} 1$ | C17 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C3 | C1 | ${ }^{*} \mathrm{C} 2$ | H21 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C3 | C1 | ${ }^{*} \mathrm{C} 2$ | H22 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C4 | C2 | ${ }^{*} \mathrm{C} 3$ | H31 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C4 | C2 | * C 3 | H32 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C5 | C3 | ${ }^{*} \mathrm{C} 4$ | H41 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC C5 | C3 | * C 4 | H42 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C5 | C1 | * C 6 | C7 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C8 | C6 | * C 7 | H71 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C9 | C7 | ${ }^{*} \mathrm{C} 8$ | H81 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C10 | C8 | * C 9 | C19 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C11 | C9 | ${ }^{*} \mathrm{C} 10$ | H101 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C9 | C10 | C11 | C12 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C12 | C10 | *C11 | H111 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C10 | C11 | C12 | C13 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C13 | C11 | *C12 | H121 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C11 | C12 | C13 | C14 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C14 | C12 | *C13 | C20 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C12 | C13 | C14 | C15 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C15 | C13 | *C14 | H141 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C13 | C14 | C15 | 015 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC 015 | C14 | * C15 | H15 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C16 | H161 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H161 | C1 | *C16 | H162 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H161 | C1 | *C16 | H163 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C6 | C1 | C17 | H171 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H171 | C1 | * C 17 | H172 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H171 | C1 | *C17 | H173 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C6 | C4 | * C 5 | C18 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC C4 | C5 | C18 | H181 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC H181 | C5 | * C 18 | H182 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC H181 | C5 | *C18 | H183 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC C8 | C9 | C19 | H191 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |


| IC | H191 | C9 | ${ }^{*}$ C19 | H192 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | ---: | :--- | :--- |
| IC | H191 | C9 | ${ }^{*}$ C19 | H193 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC | C12 | C13 | C20 | H201 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | H201 | C13 | *C20 | H202 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC | H201 | C13 | ${ }^{*}$ C20 | H203 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |

PATCH FIRST NONE LAST NONE

## Topology for retinoic acid

```
RESI RTAC -1.00 ! all-trans-retinoic acid, nomenclature from PDB
!
! retinoic acid is depronated form which exists at physiological pH
GROUP
ATOM C1 CT 0.00
ATOM C16 CT3 -0.27
ATOM H161 HA 0.09
ATOM H162 HA 0.09
ATOM H163 HA 0.09
ATOM C17 CT3 -0.27
ATOM H171 HA 0.09
ATOM H172 HA 0.09
ATOM H173 HA 0.09
GROUP
ATOM C2 CT2 -0.18
ATOM H21 HA 0.09
ATOM H22 HA 0.09
ATOM C3 CT2 -0.18
ATOM H31 HA 0.09
ATOM H32 HA 0.09
ATOM C4 CT2 -0.18
ATOM H41 HA 0.09
ATOM H42 HA 0.09
GROUP
ATOM C5 CC1A 0.00
ATOM C6 CC1A 0.00
GROUP
ATOM C18 CT3 -0.27
ATOM H181 HA 0.09
ATOM H182 HA 0.09
ATOM H183 HA 0.09
GROUP
ATOM C7 CC1B -0.15
ATOM H71 HE1 0.15
ATOM C8 CC1B -0.15
ATOM H81 HE1 0.15
GROUP
ATOM C9 CC1A 0.00
ATOM C19 CT3 -0.27 !H162 H163 H171 H172
ATOM H191 HA 0.09 ! \ | | /
ATOM H192 HA 0.09 !H161-C16 C17-H173 H191 H192
ATOM H193 HA 0.09 ! \ / \ /
ATOM C10 CC1A -0.15 ! H21 C1 H71 H81 C19
```



```
GROUP C11 CC1B -0.15 ! H22-C2 C6------C7=====C8----C9 H193
```

| ATOM H111 | HE1 | 0.15 | ! | H31-C3 | C5 | H181 | C10-H101 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ATOM C12 | CC1B | -0.15 | ! | 1 \} | 1 \} | / | \| |
| ATOM H121 | HE1 | 0.15 | . | H32 |  | C18-H182 | C11-H111 |
| GROUP |  |  | . | / | $\backslash$ | $\backslash$ | \| | |
| ATOM C13 | CC1A | 0.00 | ! | H41 | H42 | H183 | C12-H121 H201 |
| ATOM C20 | CT3 | -0.27 | ! |  |  |  | / |
| ATOM H201 | HA | 0.09 | ! |  |  |  | C13---C20-H2O2 |
| ATOM H202 | HA | 0.09 | ! |  |  |  | 11 \} |
| ATOM H203 | HA | 0.09 | ! |  |  |  | C14-H141 H203 |
| GROUP |  |  | ! |  |  |  | \| |
| ATOM C14 | CC1A | -0.25 | ! |  |  |  | $\mathrm{C} 15=029$ |
| ATOM H141 | HE1 | 0.15 | ! |  |  |  | \| (-) |
| ATOM C15 | CC | 0.62 | , |  |  |  | O29B |
| ATOM O29 | OC | -0.76 | ! |  |  |  |  |


| BOND C1 | C2 | C2 | C3 | C3 | C4 | C4 | C5 | C5 | C6 | C6 | C1 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| BOND C6 | C7 | C7 | C8 | C8 | C9 | C9 | C10 | C10 | C11 | C11 | C12 |
| BOND C12 | C13 | C13 | C14 | C14 | C15 | C15 | 029 | C15 | O29B |  |  |
| BOND C1 | C16 | C1 | C17 | C5 | C18 | C9 | C19 | C13 | C20 |  |  |
| BOND C2 | H21 | C2 | H22 | C3 | H31 | C3 | H32 | C4 | H41 | C4 | H42 |
| BOND C7 | H71 | C8 | H81 | C10 | H101 | C11 | H111 | 1 C 12 | H121 | C14 | H141 |
| BOND C16 | H161 | C16 | H162 | C16 | H163 | C17 | H171 | 1 C 17 | H172 | C17 | H173 |
| BOND C18 | H181 | C18 | H182 | C18 | H183 | C19 | H191 | 1 C 19 | H192 | C19 | H193 |
| BOND C20 | H201 | C20 | H202 | C20 | H203 |  |  |  |  |  |  |
| IC C1 | C2 | C3 | C4 |  | 0.0000 |  | 0.00 | 0.00 |  | 0.00 | 0.0000 |
| IC C2 | C3 | C4 | C5 |  | 0.0000 |  | 0.00 | 0.00 |  | 0.00 | 0.0000 |
| IC C6 | C1 | C2 | C3 |  | 0.0000 |  | 0.00 | 0.00 |  | 0.00 | 0.0000 |
| IC C1 | C6 | C7 | C8 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C6 | C7 | C8 | C9 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C7 | C8 | C9 | C10 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C2 | C6 | ${ }^{*} \mathrm{C} 1$ | C16 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC C2 | C6 | * C 1 | C17 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C6 | C1 | C16 | H161 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC H161 | C1 | * C16 | H162 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC H161 | C1 | *C16 | H163 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C6 | C1 | C17 | H171 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC H171 | C1 | *C17 | H172 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC H171 | C1 | *C17 | H173 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C3 | C1 | * C 2 | H21 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC C3 | C1 | ${ }^{*} \mathrm{C} 2$ | H22 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C4 | C2 | ${ }^{*} \mathrm{C} 3$ | H31 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC C4 | C2 | ${ }^{*} \mathrm{C} 3$ | H32 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C5 | C3 | ${ }^{*} \mathrm{C} 4$ | H41 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC C5 | C3 | ${ }^{*} \mathrm{C} 4$ | H42 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C6 | C4 | * C 5 | C18 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C4 | C5 | C18 | H181 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC H181 | C5 | * C 18 | H182 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC H181 | C5 | * C 18 | H183 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C5 | C1 | * C 6 | C7 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C8 | C6 | ${ }^{*} \mathrm{C} 7$ | H71 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C9 | C7 | * C 8 | H81 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C10 | C8 | * C 9 | C19 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC C8 | C9 | C19 | H191 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |
| IC H191 | C9 | *C19 | H192 |  | 0.0000 |  | 0.00 | 120.00 |  | 0.00 | 0.0000 |
| IC H191 | C9 | *C19 | H193 |  | 0.0000 |  | $0.00-$ | -120.00 |  | 0.00 | 0.0000 |
| IC C8 | C9 | C10 | C11 |  | 0.0000 |  | 0.00 | 180.00 |  | 0.00 | 0.0000 |


| IC | C11 | C9 | ${ }^{*}$ C10 | H101 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| IC | C9 | C10 | C11 | C12 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C12 | C10 | *C11 | H111 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C10 | C11 | C12 | C13 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C13 | C11 | *C12 | H121 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C11 | C12 | C13 | C14 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C14 | C12 | *C13 | C20 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C12 | C13 | C20 | H201 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | H201 | C13 | *C20 | H202 | 0.0000 | 0.00 | 120.00 | 0.00 | 0.0000 |
| IC | H201 | C13 | *C20 | H203 | 0.0000 | 0.00 | -120.00 | 0.00 | 0.0000 |
| IC | C12 | C13 | C14 | C15 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C15 | C13 | *C14 | H141 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | C13 | C14 | C15 | O29 | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |
| IC | O29 | C14 | *C15 | O29B | 0.0000 | 0.00 | 180.00 | 0.00 | 0.0000 |

PATCH FIRST NONE LAST NONE

## Parameters for retinoids

```
read parameter card append
BONDS
!
!atom types Kb b0
!
!BTE2, 2-butene
HE1 CC1A 360.500 1.100
HE1 CC1B 360.500 1.100
HE2 CC2 365.000 1.100
CC1A CC1A 440.000 1.340
CC1B CC1B 440.000 1.340
!13DB, Butadiene
CC1A CC2 500.000 1.342
CC1B CC2 500.000 1.342
CC1A CC1B 300.000 1.470
!13DP, 1,3-Pentadiene
CC1B CT3 383.000 1.504
!MECH
CC1A CT2 365.000 1.502
CC1A CT 365.000 1.502
CC1A CT3 383.000 1.504
!TMCH/MECH
CT CT2 222.500 1.538
CT CT3 222.500 1.538
!PACP
CA NH1 305.000 1.414
!PRAL
CC1A CD 300.000 1.4798
CD O 720.000 1.205
CD HR1 330.000 1.110
!PRAC
CC1A CC 440.000 1.489
!CROT
CC1A C 440.000 1.489
!SCH1, Schiff's base model compound 1, deprotonated
```



| CT | CT3 | HA | 33.43 | 110.10 | 22.53 | 2.179 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CT | CC1A | CT3 | 48.00 | 123.50 |  |  |
| ! PACP |  |  |  |  |  |  |
| CA | NH1 | C | 50.00 | 120.00 |  |  |
| H | NH1 | CA | 34.00 | 117.00 |  |  |
| NH1 | CA | CA | 40.00 | 120.00 | 35.00 | 2.4162 |
| ! CROT |  |  |  |  |  |  |
| CT3 | CC1A | CT3 | 47.00 | 125.20 |  |  |
| CC1A | CC1A | C | 48.00 | 123.50 |  |  |
| HE1 | CC1A | C | 52.00 | 119.50 |  |  |
| $\bigcirc$ | C | CC1A | 80.00 | 122.50 |  |  |
| NH1 | C | CC1A | 80.00 | 116.50 |  |  |
| ! PROL |  |  |  |  |  |  |
| HE1 | CC1A | CT2 | 40.00 | 116.00 |  |  |
| OH1 | CT2 | CC1A | 75.70 | 110.10 |  |  |
| CT2 | CC1A | CC2 | 28.00 | 126.00 |  |  |
| ! PRAL |  |  |  |  |  |  |
| CC2 | CC1A | $C D$ | 60.00 | 120.00 |  |  |
| CC1A | CC1A | $C D$ | 60.00 | 120.00 |  |  |
| HE1 | CC1A | $C D$ | 32.00 | 122.00 | ! sum=2 |  |
| CC1A | CD | $\bigcirc$ | 75.00 | 124.00 |  |  |
| CC1A | $C D$ | HR1 | 15.00 | 115.00 |  |  |
| $\bigcirc$ | $C D$ | HR1 | 75.00 | 121.00 | ! fixed |  |
| ! PRAC |  |  |  |  |  |  |
| HE1 | CC1A | CC | 52.00 | 119.50 | $!$ |  |
| CC2 | CC1A | CC | 40.00 | 119.00 | 35.00 | 2.5267 |
| CC1A | CC | OC | 40.00 | 118.00 | 50.00 | 2.3880 |
| CC1A | CC1A | CC | 48.00 | 123.50 | ! |  |
| !SCH1, Schiff's base model compound 1, deprotonated |  |  |  |  |  |  |
| CT3 | NS1 | CC1B | 67.00 | 111.00 |  |  |
| NS1 | CC1B | CT3 | 52.00 | 123.00 | ! sum=24 |  |
| NS 1 | CC1B | HE1 | 38.00 | 119.50 | ! |  |
| HA | CT3 | NS1 | 42.00 | 113.50 |  |  |
| !SCH2, Schiff's base model compound 2, protonated |  |  |  |  |  |  |
| CT3 | NS 2 | CC1B | 67.00 | 123.60 |  |  |
| CT3 | NS2 | HC | 38.00 | 117.40 |  |  |
| CC1B | NS2 | HC | 38.00 | 118.80 |  |  |
| NS2 | CC1B | CT3 | 40.00 | 125.60 |  |  |
| NS2 | CC1B | HR1 | 38.00 | 114.00 |  |  |
| CT3 | CC1B | HR1 | 42.00 | 120.40 |  |  |
| HA | CT3 | NS2 | 42.00 | 110.10 |  |  |
| !SCH3, Schiff's base model compound 3, protonated |  |  |  |  |  |  |
| NS2 | CC1B | CC1A | 40.00 | 125.60 |  |  |
| HR1 | CC1B | CC1A | 42.00 | 120.40 |  |  |
| !SCK0, deprotonated Schiff's base, lysine retinal patch |  |  |  |  |  |  |
| CT2 | CT2 | NS1 | 67.70 | 110.00 | !from 1 | ine |
| HA | CT2 | NS1 | 42.00 | 113.50 |  |  |
| CT2 | NS1 | CC1B | 67.00 | 111.00 |  |  |
| NS1 | CC1B | CC1A | 40.00 | 123.00 |  |  |
| !SCK1, protonated Schiff's base, lysine retinal patch |  |  |  |  |  |  |
| CT2 | NS2 | CC1B | 67.00 | 123.60 |  |  |
| CT2 | NS2 | HC | 38.00 | 117.40 |  |  |
| CT2 | CT2 | NS2 | 67.70 | 110.00 | !from 1 | ine |
| HA | CT2 | NS2 | 42.00 | 110.10 |  |  |



| CT | CC1A | CC1A | CT2 | 10.0000 | 2 | 180.00 | ! double bond |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CT | CC1A | CC1A | CT3 | 10.0000 | 2 | 180.00 | ! double bond |
| CC1A | CC1A | CT | CT2 | 0.5000 | 2 | 0.00 | ! c2-c1-c6=c5 |
| CC1A | CC1A | CT | CT2 | 0.3000 | 3 | 0.00 | ! c2-c1-c6=c5 |
| CC1A | CC1A | CT | CT3 | 0.5000 | 2 | 0.00 | ! c16/c17-c1-c6=c5 |
| CC1A | CC1A | CT | CT3 | 0.4000 | 3 | 0.00 | ! c16/c17-c1-c6=c5 |
| CC1A | CC1A | CT2 | CT2 | 0.5000 | 2 | 0.00 | ! c3-c4-c5=c6 |
| CC1A | CC1A | CT2 | CT2 | 0.3000 | 3 | 0.00 | ! c3-c4-c5=c6 |
| CT2 | CT2 | CC1A | CT3 | 0.1900 | 3 | 0.00 | ! c3-c4-c5-c18 |
| CT2 | CT | CC1A | CT3 | 0.4000 | 3 | 0.00 | ! c2-c1-c6-c7 |
| CT3 | CT | CC1A | CT3 | 0.4000 | 3 | 0.00 | ! c16/c17-c1-c6-c7 |
| CC1A | CT2 | CT2 | CT2 | 0.1900 | 3 | 0.00 | ! |
| CT2 | CT2 | CT | CC1A | 0.2000 | 3 | 0.00 | ! |
| CT2 | CT2 | CT | CT3 | 0.2000 | 3 | 0.00 | ! |
| CT3 | CC1A | CT2 | HA | 0.1900 | 3 | 0.00 |  |
| CC1A | CT | CT2 | HA | 0.1900 | 3 | 0.00 |  |
| CT3 | CT | CT2 | HA | 0.1900 | 3 | 0.00 |  |
| CC1A | CT2 | CT2 | HA | 0.1900 | 3 | 0.00 |  |
| CT2 | CC1A | CT3 | HA | 0.1600 | 3 | 0.00 |  |
| CT | CC1A | CT3 | HA | 0.1600 | 3 | 0.00 |  |
| CT2 | CT | CT3 | HA | 0.1600 | 3 | 0.00 |  |
| CC1A | CT | CT3 | HA | 0.1600 | 3 | 0.00 |  |
| CT3 | CT | CT3 | HA | 0.1600 | 3 | 0.00 |  |
| ! MECH |  |  |  |  |  |  |  |
| CT | CC1A | CC1B | CC2 | 0.9000 | 1 | 0.00 |  |
| CT | CC1A | CC1B | CC2 | 2.1000 | 2 | 180.00 |  |
| CT | CC1A | CC1B | CC2 | 0.2200 | 3 | 0.00 |  |
| CT | CC1A | CC1B | CC2 | 0.2500 | 5 | 180.00 |  |
| CT | CC1A | CC1B | CC2 | 0.1000 | 6 | 0.00 |  |
| CT | CC1A | CC1B | CC1B | 0.9000 | 1 | 0.00 |  |
| CT | CC1A | CC1B | CC1B | 2.1000 | 2 | 180.00 |  |
| CT | CC1A | CC1B | CC1B | 0.2200 | 3 | 0.00 |  |
| CT | CC1A | CC1B | CC1B | 0.2500 | 5 | 180.00 |  |
| CT | CC1A | CC1B | CC1B | 0.1000 | 6 | 0.00 |  |
| CC1B | CC1A | CT | CT2 | 0.3000 | 3 | 0.00 |  |
| CC1B | CC1A | CT | CT3 | 0.3000 | 3 | 0.00 |  |
| CC1B | CC1A | CC1A | CT2 | 0.5600 | 1 | 180.00 | ! double bond |
| CC1B | CC1A | CC1A | CT2 | 7.0000 | 2 | 180.00 | ! double bond |
| CC1B | CC1A | CC1A | CT3 | 0.5600 | 1 | 180.00 | ! double bond |
| CC1B | CC1A | CC1A | CT3 | 7.0000 | 2 | 180.00 | ! double bond |
| CT | CC1A | CC1B | HE1 | 1.0000 | 2 | 180.00 |  |
| CC1A | CC1A | CC1B | HE1 | 1.0000 | 2 | 180.00 |  |
| ! PACP |  |  |  |  |  |  |  |
| $\bigcirc$ | C | NH1 | CA | 2.5000 | 2 | 180.00 | ! |
| HP | CA | CA | NH1 | 4.2000 | 2 | 180.00 | ! |
| CA | CA | CA | NH1 | 3.1000 | 2 | 180.00 | ! |
| C | NH1 | CA | CA | 1.2000 | 2 | 180.00 | ! |
| C | NH1 | CA | CA | 1.0000 | 3 | 180.00 | ! |
| H | NH1 | CA | CA | 0.5000 | 2 | 180.00 | ! |
| CA | NH1 | C | CT3 | 2.5000 | 2 | 180.00 | ! |
| ! CROT |  |  |  |  |  |  |  |
| CT3 | CC1A | CT3 | HA | 0.3000 | 3 | 0.00 | ! |
| CC1B | CC1A | CC1A | C | 0.5600 | 1 | 180.00 | !double bond |
| CC1B | CC1A | CC1A | C | 7.0000 | 2 | 180.00 | ! double bond |
| 0 | C | CC1A | CC1A | 0.7000 | 1 | 180.00 | . |
| 0 | C | CC1A | CC1A | 1.2000 | 2 | 180.00 | ! |
| 0 | C | CC1A | CC1A | 0.1000 | 3 | 180.00 | ! |



!SCH1, SCH2

| NS1 | 0.000000 | -0.200000 | 1.850000 | $!N$ for deprotonated Schiff's |
| :--- | :--- | :--- | :--- | :--- |
| base | 0.000000 | -0.200000 | 1.850000 | $!N$ for protonated Schiff's base |
| NS2 | 0.000 |  |  |  |

