

# External rays for a regular polynomial endomorphism of $\mathbb{C}^2$ associated with Chebyshev mappings

Shizuo Nakane\*

In this note, the dynamics of a regular polynomial map of  $\mathbb{C}^2$  is investigated. Especially, landing points of the external rays are completely characterized.

## 1 Introduction

In this note, we consider the external rays of the map  $F : \mathbb{C}^2 \rightarrow \mathbb{C}^2$  of the form :

$$F(x, y) = (x^2 - 2y, y^2 - 2x).$$

External rays were first defined for polynomial maps on  $\mathbb{C}$  to investigate the combinatorial properties of the dynamics on the Julia sets. Let  $P$  be a monic centered polynomial with degree  $d$  of one variable. Let  $\varphi = \varphi_P$  be its *Böttcher coordinate*, that is, a conformal map  $\varphi$  in a neighborhood of the point at  $\infty$  satisfying

$$\varphi(P(z)) = \varphi(z)^d, \quad \lim_{z \rightarrow \infty} \frac{\varphi(z)}{z} = 1.$$

By this functional equation, it can be continued analytically until it meets a critical point. Especially, if  $K(P)$  is connected, it extends to a conformal map  $\varphi : \mathbb{C} - K(P) \rightarrow \mathbb{C} - \overline{\mathbb{D}}$ . The *external ray*  $R_P(\theta)$  of *external angle*  $\theta$  is defined by the preimage of the ray  $\{re^{2\pi i\theta}; r > 1\}$  by  $\varphi$ . We say it lands at a point  $z \in J(P)$  if it is continued to  $r > 1$  and converges to  $z$  as  $r \rightarrow 1$ . Recently, Bedford and Jonsson [BJ] defined external rays for regular polynomial endomorphisms of  $\mathbb{C}^k$  and established a landing property with some additional assumptions. Although the map  $F$  does not satisfy the assumptions in [BJ], we can investigate the landing property from the explicit expression of its Böttcher coordinate.

The map  $F$  has dynamically distinguished properties. For example, it is *critically finite*, that is, the union of the forward orbit of the critical set forms an analytic subset of

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\*Professor, General Education and Research Center, Tokyo Polytechnic University  
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$\mathbb{C}^2$ . This is because it is closely related to the *Chebyshev maps* of two variables. A typical example of Chebyshev maps of one variable is the quadratic polynomial  $p(z) = z^2 - 2$ , which is critically finite, too. A natural extension of this Chebyshev map to the two variables case is  $f(z) = z^2 - 2\bar{z}$ . By virtue of the distinguished properties of Chebyshev maps, Uchimura [U1] has obtained many interesting results.

Here we show why Chebyshev maps are easily analyzed. Put  $z = \psi(t) = t + 1/t$ . Then

$$p(\psi(t)) = (t + 1/t)^2 - 2 = t^2 + 1/t^2 = \psi(t^2).$$

Hence a branch of its inverse  $\varphi = \psi^{-1}$  satisfies  $\varphi(p(z)) = \varphi(z)^2$  and it gives the Böttcher coordinate of  $p$ . Then the external ray  $R_p(\theta)$  is explicitly written by

$$z = \psi(re^{2\pi i\theta}) = re^{2\pi i\theta} + \frac{1}{r}e^{-2\pi i\theta}, \quad r > 1,$$

and it lands at the point

$$z_0 = e^{2\pi i\theta} + e^{-2\pi i\theta} = 2 \cos 2\pi\theta.$$

Consequently,  $J(p) = \{z = 2 \cos 2\pi\theta; \theta \in \mathbb{T}\}$ . In the sequel, we will apply this idea to the maps  $f$  and  $F$ .

## 2 External rays for the map $f$

Consider the map  $f$  studied in [U1] of the form :

$$f(z) = z^2 - 2\bar{z}.$$

The map  $f : \mathbb{C} \rightarrow \mathbb{C}$  is not holomorphic but is associated with the Chebyshev maps of two variables and its dynamics is completely determined. See [U1]. Since the jacobian of  $f$  is :

$$Jac(f) = |\partial f/\partial z|^2 - |\partial f/\partial \bar{z}|^2 = 4(|z|^2 - 1),$$

its *critical set*  $\mathcal{C}(f)$  is the unit circle  $|z| = 1$ . The *filled-in Julia set*  $K(f)$ , i.e., the set of points with bounded orbits, is parametrized (with some identification) as follows.

$$K(f) = \{z = z(\phi, \theta) = e^{2\pi i\phi} + e^{2\pi i\theta} + e^{-2\pi i(\phi+\theta)}; (\phi, \theta) \in \mathbb{T}^2\}. \tag{2.1}$$

And its boundary is the hypocycloid (see Figure 1) :

$$z = 2e^{2\pi i\theta} + e^{-4\pi i\theta}, \quad \theta \in \mathbb{T}.$$

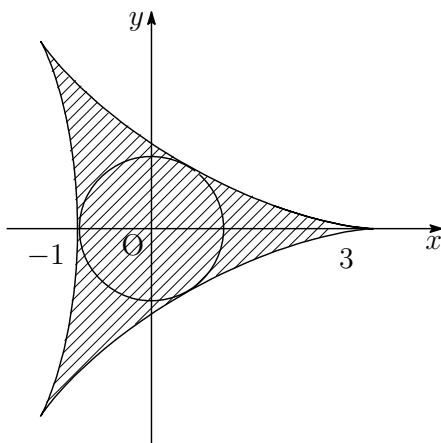


Figure 1: The dark region :  $K(f)$ , the circle :  $\mathcal{C}(f)$

Moreover, the dynamics of  $f$  on  $K(f)$  is expressed by  $f(z(\phi, \theta)) = z(2\phi, 2\theta)$ , which enables us to describe their dynamics by symbolic dynamics. Although this parametrization of  $K(f)$  seems a bit tricky, we will give a dynamical meaning of the parameters  $\phi, \theta$  above in the next section.

First we study the Böttcher coordinate of  $f$ . Put

$$z = \psi(t) = t + \frac{1}{t} + \frac{\bar{t}}{t}.$$

Then its jacobian  $Jac(\psi)$  satisfies

$$\begin{aligned} Jac(\psi) &= |z_t|^2 - |\bar{z}_t|^2 \\ &= \left|1 - \frac{\bar{t}}{t^2}\right|^2 - \left|\frac{1}{t} - \frac{1}{\bar{t}^2}\right|^2 \\ &= \left(1 - \frac{1}{|t|^2}\right) \left|1 - \frac{\bar{t}}{t^2}\right|^2. \end{aligned}$$

Thus  $\psi$  gives a diffeomorphism from  $\mathbb{C} - \overline{\mathbb{D}}$  onto  $\mathbb{C} - K(f)$  and it is easy to see

$$f(\psi(t)) = \psi(t^2), \quad \lim_{t \rightarrow \infty} \frac{\psi(t)}{t} = 1.$$

That is, the map  $\varphi = \psi^{-1}$  should be the Böttcher coordinate of  $f$  and we can define the external ray as follows :

$$R_f(\theta) = \psi(\{re^{2\pi i\theta}; r > 1\}).$$

Then we have the following.

**Theorem 2.1.** *The external ray  $R_f(\theta)$  is parametrized by*

$$z = re^{2\pi i\theta} + \frac{1}{r}e^{2\pi i\theta} + e^{-4\pi i\theta} \quad (r > 1)$$

*and it lands at the point*

$$z_0 = 2e^{2\pi i\theta} + e^{-4\pi i\theta} \in \partial K(f).$$

### 3 External rays for the map $F$

Now consider the following map in  $\mathbb{C}^2$ .

$$F(x, y) = (x^2 - 2y, y^2 - 2x),$$

which is closely related to the map  $f$  in the previous section. In fact, the map  $F$  restricted to  $H = \{(x, y) \in \mathbb{C}^2; y = \bar{x}\}$  is equivalent to  $f$ . Let  $\mathcal{C}(F)$  be the set of the critical points of  $F$ . By a direct calculation, it follows  $\mathcal{C}(F) = \{xy = 1\}$ .

Let  $f(z)$  be a polynomial endomorphism of  $\mathbb{C}^k$  of degree  $d$  and let  $f_h(z)$  be the degree  $d$  part of  $f(z)$ . It is *regular* if  $f_h^{-1}(0) = \{0\}$ . Note that regular polynomial maps extend to analytic maps of  $\mathbb{P}^k$ . Let  $\Pi$  denote the hyperplane at  $\infty$ , which is isomorphic to  $\mathbb{P}^{k-1}$ . In case  $k = 2$ ,  $\Pi$  is isomorphic to the Riemann sphere  $\overline{\mathbb{C}}$ . For a regular polynomial map  $f$ , we denote the filled-in Julia set also by  $K(f)$ . It is a compact subset of  $\mathbb{C}^k$ . And  $J(f)$  denotes the smallest Julia set of  $f$ , that is, the support of  $\mu = (dd^c G_f)^k$ . Here  $G_f$  is the Green function of  $f$ . And we put  $f_\Pi = f|_\Pi$ ,  $J_\Pi = J(f_\Pi)$ .

Let  $W^s(J_\Pi, f)$  be the stable set of  $J_\Pi$  :

$$W^s(J_\Pi, f) = \{z \in \mathbb{P}^k; \lim_{n \rightarrow \infty} \text{dist}(f^n(z), J_\Pi) = 0\}.$$

The inverse Böttcher coordinate  $\Psi$  is a homeomorphism  $W_{loc}^s(J_\Pi, f_h) \rightarrow W_{loc}^s(J_\Pi, f)$  conjugating  $f_h$  to  $f$ . It extends to  $W^s(J_\Pi, f_h)$  until it meets a critical point. Each local stable manifold  $W_{loc}^s(a)$  ( $a \in J_\Pi$ ) is a complex disk homeomorphic to  $\overline{\mathbb{C}} - \overline{\mathbb{D}_R}$  for some  $R > 1$ . External rays are the rays in  $W^s(J_\Pi, f)$  defined by the gradient lines of the Green function  $G_f$  restricted to  $W_{loc}^s(a)$ . Since the Böttcher coordinate transforms the Green function into a canonical form, external rays are the images of the actual rays by the inverse Böttcher coordinate, just as for polynomials of one variable.

Bedford and Jonsson [BJ] established the continuous landing property of external rays for regular polynomial endomorphisms of  $\mathbb{C}^2$ .

**Theorem 3.1.** (*[BJ], Theorem 10.2*)

*Let  $f$  be a regular polynomial endomorphism of  $\mathbb{C}^2$ . Assume*

- (1)  $f_{\Pi}$  is uniformly expanding on  $J_{\Pi}$ .  
 (2)  $f$  is uniformly expanding on  $J(f)$ .  
 (3) The non-wandering set of  $f$  in  $\partial K(f)$  consists of  $J(f)$  and a hyperbolic set  $\mathcal{S}_1$  of unstable index 1.  
 (4)  $W^s(\mathcal{S}_1) = \cup_{\hat{x} \in \mathcal{S}_1} W^s(\hat{x})$ .  
 (5)  $W^s(J_{\Pi}) \cap \mathcal{C}(f) = \emptyset$ .  
 Then all external rays land onto  $J(f)$  and landing points vary continuously.

As a trivial example, we consider the map  $F_h(x, y) = (x^2, y^2)$ . Then

$$\begin{aligned} \mathcal{C}(F_h) &= \{x = 0\} \cup \{y = 0\}, & K(F_h) &= \{|x| \leq 1, |y| \leq 1\}, \\ J(F_h) &= \{|x| = |y| = 1\}, & F_{h, \Pi}(\zeta) &= \zeta^2, & J_{\Pi} &= \{|\zeta| = 1\}, \\ W^s(\zeta) &= \{y = \zeta x, |x| > 1\}, & W^s(J_{\Pi}, F_h) &= \{|x| = |y| > 1\}. \end{aligned}$$

And all the assumptions of the above theorem are satisfied. Then external rays for  $F_h$  are labelled by two angles  $(\phi, \theta) \in \mathbb{T}^2$ . Here  $\zeta = y/x = e^{2\pi i \phi} \in J_{\Pi}$  and  $\theta$  is the argument of the ray in the disk  $W^s(\zeta)$ . Hence the external ray  $R_{F_h}(\phi, \theta)$  is just the ray :

$$x = r e^{2\pi i \theta}, \quad y = \zeta x = r e^{2\pi i(\phi + \theta)}, \quad (r > 1),$$

which lands at  $(e^{2\pi i \theta}, e^{2\pi i(\phi + \theta)}) \in J(F_h)$ .

Our map  $F$  is regular but is not expanding on  $J(F)$  since  $J(F)$  contains critical points, as we will see later. Next lemma says that it satisfies the last condition (5). Its proof also implies that  $F$  is critically finite.

**Lemma 3.1.**  $W^s(J_{\Pi}, F) \cap \mathcal{C}(F) = \emptyset$ .

*proof.* Note that the critical set  $\mathcal{C}(F)$  is parametrized as  $x = t, y = 1/t$ . We calculate the critical orbits and by induction, we show

$$F^n(t, t^{-1}) = (t^{2^n} + 2t^{-2^{n-1}}, t^{-2^n} + 2t^{2^{n-1}}) \quad (n \geq 2).$$

In fact, it is true for  $n = 2$ . Suppose it is true for  $n = k$ . Then the first entry of  $F^{k+1}(t, t^{-1})$  is

$$\begin{aligned} (t^{2^k} + 2t^{-2^{k-1}})^2 - 2(t^{-2^k} + 2t^{2^{k-1}}) &= t^{2^{k+1}} + 4t^{2^{k-1}} + 4t^{-2^k} - 2t^{-2^k} - 4t^{2^{k-1}} \\ &= t^{2^{k+1}} + 2t^{-2^k}. \end{aligned}$$

The same holds for the second entry. Hence the case  $n = k + 1$  is also true.

Note that the map  $F$  has two super-attracting fixed points  $[1 : 0 : 0]$  and  $[0 : 1 : 0]$  in  $\Pi$  and  $W^s(J_{\Pi}, F)$  is contained in the common boundary of their basins. The above

calculation implies that the parts  $|t| > 1$  and  $|t| < 1$  are contained in the basins of the points  $[1 : 0 : 0]$  and  $[0 : 1 : 0]$  respectively and the part  $|t| = 1$  is contained in  $K(F)$ . Thus  $\mathcal{C}$  never intersects  $W^s(J_\Pi, F)$ . This completes the proof.  $\square$

Now we consider the external rays for  $F$ . Fortunately, we have an explicit expression of an inverse Böttcher coordinate of  $F$  and we can define them directly. Put

$$(x, y) = \Psi(u, v) = \left(u + \frac{1}{v} + \frac{v}{u}, v + \frac{1}{u} + \frac{u}{v}\right).$$

Then it satisfies the functional equation

$$F \circ \Psi(u, v) = \Psi(u^2, v^2) = \Psi \circ F_h(u, v).$$

The jacobian  $Jac(\Psi)$  is written by

$$Jac(\Psi)(u, v) = \left(1 - \frac{1}{uv}\right)\left(1 - \frac{u}{v^2}\right)\left(1 - \frac{v}{u^2}\right).$$

Hence it is invertible on  $W^s(J_\Pi, F_h)$ . The inverse  $\Phi = \Psi^{-1}$  is a Böttcher coordinate of  $F$ .

Then each stable manifold  $W_F^s(\zeta)$  of  $\zeta \in J_\Pi$  for  $F$  is the image of  $W_{F_h}^s(\zeta) = \{(t, \zeta t); |t| > 1\} \cong \mathbb{C} - \overline{\mathbb{D}}$  by  $\Psi$ . This coordinate gives the Böttcher coordinate of the restriction of  $F$  on  $W_F^s(\zeta)$ . Hence the external ray  $R_F(\phi, \theta)$  is the image of  $R_{F_h}(\phi, \theta)$  by  $\Psi$ .

**Theorem 3.2.** *The external ray  $R_F(\phi, \theta)$  is expressed by*

$$\begin{aligned} x &= r e^{2\pi i \theta} + \frac{1}{r} e^{-2\pi i(\phi+\theta)} + e^{2\pi i \phi} \\ y &= r e^{2\pi i(\phi+\theta)} + \frac{1}{r} e^{-2\pi i \theta} + e^{-2\pi i \phi} \quad (r > 1). \end{aligned}$$

*Its landing point depends continuously on  $(\phi, \theta) \in \mathbb{T}^2$  :*

$$\begin{aligned} x_0 &= e^{2\pi i \theta} + e^{-2\pi i(\phi+\theta)} + e^{2\pi i \phi} \\ y_0 &= e^{2\pi i(\phi+\theta)} + e^{-2\pi i \theta} + e^{-2\pi i \phi} = \overline{x_0}. \end{aligned}$$

Thus  $(x_0, y_0) \in H$ . Recall that this parametrization of  $x_0$  coincides with that of  $K(f)$  described in (2.1) in the previous section.

**Lemma 3.2.**  $K(F) = \{(x, \bar{x}) \in H; x \in K(f)\}$ .

*proof.* Note that the numbers  $u, \frac{1}{v}, \frac{v}{u}$  (resp.  $v, \frac{1}{u}, \frac{u}{v}$ ) in the definition of  $\Psi$  are the roots of the cubic equation  $t^3 - xt^2 + yt - 1 = 0$ , (resp.  $t^3 - yt^2 + xt - 1 = 0$ .) Thus the map  $\Psi : (\mathbb{C} - \{0\})^2 \rightarrow \mathbb{C}^2$  is surjective. Hence, for any  $(x, y) \in \mathbb{C}^2$ , there exists a point

$(u, v) \in (\mathbb{C} - \{0\})^2$  such that  $(x, y) = \Psi(u, v)$ . Then we have  $F^n(x, y) = \Psi \circ F_h^n(u, v)$  and it easily follows that  $F^n(x, y) \rightarrow \infty$  if and only if  $F_h^n(u, v) \rightarrow \infty$ . Since  $\Psi(\frac{1}{v}, \frac{1}{u}) = \Psi(u, v)$ , it is easy to see that  $(x, y) \in K(F)$  if and only if  $|u| = |v| = 1$ . This completes the proof.  $\square$

**Lemma 3.3.**  $J(F) = K(F)$ .

*proof.* Note that  $J(F) \subset K(F)$ . Since the critical value set of  $\Psi$  intersects  $K(F)$  only at the boundary of  $K(f)$ ,  $\Psi$  is locally invertible in the interior of  $K(F)$  in  $H$ . Let  $\Phi_j, 0 \leq j \leq 5$  be the branches of  $\Psi^{-1}$  there. Then, we have

$$(dd^c G)^2 = \frac{1}{3} \sum_{j=0}^5 \Phi_j^*(dd^c G_h)^2.$$

Hence,  $J(F) = \text{supp}(dd^c G)^2$  contains the image of  $J(F_h) = \{|u| = |v| = 1\}$  by  $\Psi$ . Thus  $K(F) \subset J(F)$ . This completes the proof.  $\square$

Now the parameters  $\phi$  and  $\theta$  turn out to be the external angles for  $F$ . Note that  $\mathcal{C}(F) \cap H = \{(x, \bar{x}); |x| = 1\}$  coincides with  $\mathcal{C}(f)$  and is contained in  $J(F)$ . See Figure 1. Thus  $F$  is not expanding on  $J(F)$ .

Now Lemma 3.1 says  $W^s(J_\Pi) \cap \mathcal{C}(F) = \emptyset$ . Then it follows from Theorem 7.4 in [BJ] that  $\Psi$  extends to a homeomorphism from  $W^s(J_\Pi, F_h)$  onto  $W^s(J_\Pi, F)$  conjugating  $F_h$  to  $F$ . In our case, this is trivial and we have a global parametrization of  $W^s(J_\Pi, F)$  as the union of the stable manifolds  $W_F^s(\zeta)$  with  $\zeta = e^{2\pi i \phi}$ .

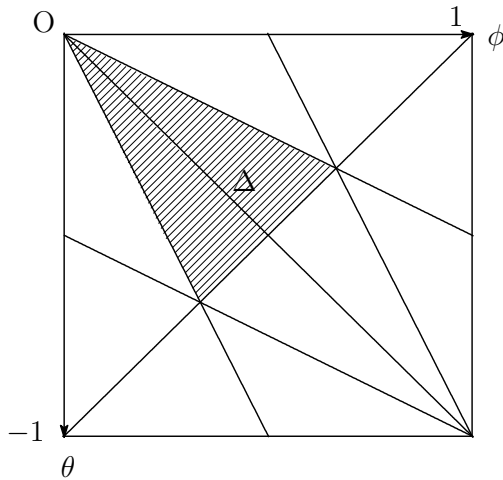


Figure 2: Equivalence on  $\mathbb{T}^2$  and the fundamental region  $\Delta$

Note that  $(\phi, \theta)$  and the parameters

$$\begin{aligned} \rho_1(\phi, \theta) &= (1 + \theta, \phi - 1) \\ \rho_2(\phi, \theta) &= (\phi, -\phi - \theta) \\ \rho_3(\phi, \theta) &= (-\phi - \theta, \theta) \end{aligned}$$

give a same landing point  $(x_0, y_0)$ . That is, several rays land at a same point. We will investigate this in details. We remark that  $\rho_1, \rho_2$  and  $\rho_3$  are the reflections with respect to the lines  $\phi = \theta + 1, \phi = -2\theta$  and  $\theta = -2\phi$  respectively. These reflections give an equivalence relation in  $\mathbb{T}^2$ . The fundamental region is the closed triangular region  $\Delta$  surrounded by the three lines :

$$\phi = \theta + 1, \quad \phi = -2\theta, \quad \theta = -2\phi.$$

Figure 2 shows the torus  $\mathbb{T}^2$ , where the dark region indicates the fundamental region  $\Delta$ . Each triangle is equivalent to one of the two halves of  $\Delta$ . Now the next lemma is easy to see.

**Lemma 3.4.** *The equivalence class of a point in the interior of  $\Delta$  consists of 6 points, while that of a point on one of the three edges of  $\partial\Delta$  consists of 3 points and that of a vertex of  $\partial J(F)$  consists of a single point itself.*

Since  $\Delta$  and  $\partial\Delta$  correspond respectively to  $J(F)$  and  $\partial J(F)$ , we have the following.

**Theorem 3.3.** *Each point  $z = (x, y)$  in  $J(F)$  is the landing point of exactly one, 3 or 6 external rays if  $z$  is a cusp point on  $\partial J(F)$ ,  $z$  is a non-cusp point on  $\partial J(F)$  or  $z \in \text{int } J(F)$ , respectively.*

Finally note that the restriction of  $\Psi$  to  $H$  is

$$\Psi(t, \bar{t}) = \left(t + \frac{1}{\bar{t}} + \frac{\bar{t}}{t}, \bar{t} + \frac{1}{t} + \frac{t}{\bar{t}}\right) = (\psi(t), \overline{\psi(t)}),$$

where  $\psi$  is the Böttcher coordinate of  $f$ . Thus the external rays for the map  $f$  studied in the previous section are just the restriction of the rays for the map  $F$  to  $H$ .

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