# Parametric Study of the Performance of a Biplane Joined at the Tips 

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# ABSTRACT <br> PARAMETRIC STUDY OF THE PERFORMANCE OF A BIPLANE JOINED AT THE TIPS 

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This thesis investigates the impact on performance of stagger and gap on a biplane wing joined at the tips by endplates. The effects due to these parameters on the aerodynamic performance of the biplane will be discussed in this thesis to better understand the interaction between the two wings. The experimental integrated force data were obtained from wind tunnel tests performed in the University of Dayton Low Speed Wind Tunnel (LSWT) at two different Reynolds numbers. Test conditions varied angle of attack from $-2^{\circ}$ to $25^{\circ}$ in $0.25^{\circ}$ increments. The computational study used the Athena Vortex Lattice (AVL) code based on the Vortex Lattice Method (VLM). Then, the computed results were verified by comparison with the experimental data and the two were found to compare well under some conditions, and not well under other conditions. The results show that both gap and stagger have positive effects on the coefficient of lift and aerodynamic efficiency. These effects become less important at higher values of gap, due to the smaller interactions between the two wings. The models with positive stagger show higher lift coefficient and higher aerodynamic efficiency
compared to the models with negative stagger. Curves of lift coefficient versus angle of attack show a change in the lift slope, $d C_{L} / d \alpha$, around the angle of attack of maximum aerodynamic efficiency. These interesting results will be discussed in this thesis as well as a more profound analysis of the aerodynamic performance using both experimental and theoretical results.

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## NOMENCLATURE

| AR | aspect ratio |
| :---: | :---: |
| $b$ | length of semi-span |
| $c$ | mean chord length |
| $C_{D}$ | coefficient of drag |
| $C_{\text {Do }}$ | drag coefficient at zero lift |
| $\mathrm{C}_{\mathrm{L}}$ | lift coefficient |
| $\mathrm{C}_{f}$ | skin friction coefficient |
| $D_{i}$ | induced drag |
| $D_{\text {interfernce }}$ | interference drag |
| $D_{\text {parsite }}$ | parasite drag |
| $D_{\text {skinfricion }}$ | skin friction drag |
| $D_{\text {toat }}$ | total Drag force |
| $e$ | span efficiency factor |
| g | gap |
| $L$ | lift force |
| L/D | lift over drag ratio |
| $\mathrm{k}_{\mathrm{i}}$ | Munk's Equivalent Span Factor |
| Re | Reynolds Number |
| $S$ | wing area |
| SVP | saturation vapor pressure for a given temperature T |
| s | stagger |
| $\mathrm{V}_{\infty}$ | freestream velocity |
| $\alpha$ | angle of attack |
| $\alpha_{i}$ | lift induced angle of attack |
| $\Gamma$ | circulation strength |
| $\Delta \mathrm{H}$ | manometer reading |
| $\mu$ | dynamic viscosity of the fluid |
| $\rho$ | density of the fluid |

coefficient of mutual influence
velocity potential (scalar)
vorticity vector in two dimensional flow about z -axis
curl
divergence
gradient
biplane's geometric parameter
biplane's geometric parameter
strength of a vortex $i$

# PARAMETRIC STUDY OF THE PERFORMANCE OF A BIPLANE JOINED AT THE TIPS 

## I. Introduction

### 1.1 Background

The science of aerodynamics is concerned with the motion of air and of bodies moving through air. The airflow around an aircraft is a highly complex phenomenon. To study it, in the present state of our knowledge, demands simplifying assumptions. These must be based, at least partly, on experimental observations. Making mathematical deductions and predictions belong to the theoretical side. To arrive at the knowledge of such mathematical laws, Galileo says, we must make use of sense and reason, by passing through three stages:

- The observation of the facts which fall within our experience;
- The elaboration of a mathematical hypothesis as a presumed explanation of the phenomena under observation;
- Verification of the hypothesis through new facts of experience. If the verification of experience agrees, the hypothesis becomes law.

The aim of this thesis is therefore to lay bare the behavior of airflow around a biplane joined at the tips by endplates, using the observation of the facts, elaborating hypotheses, and verifying these hypotheses using an experimental approach.

Airflow around a biplane is a complex study. It is important, as a first step, to understand which factors influence the performance of a biplane and
furthermore, a biplane joined at the tips by endplates. This report will take into account the influence of stagger and gap on the performance of a biplane joined at the tips by endplates.

### 1.2 Biplane

The first successful powered, piloted airplanes in history were biplane designs, with the Wright brothers' planes being the earliest example (Figure 1.1). The biplane is an aircraft with two sets of wings, an upper set and a lower set, separated by struts and wires and connected to the upper and lower parts of the fuselage. The reason for this was structural. Airplanes were initially quite fragile.


Figure 1.1: The beginning of the first flight, December 17, 1903. Credits - Library of Congress

The reason airplanes were so fragile had to do with the difficulty of flight in the early days. For the first decade or so of flight, because of low-powered engines and designs that created a lot of drag, airplanes could barely make it into the air at all. Planes therefore were constructed of the lightest materials, such as fabric and wood, and built in such a way that they used as little of the heaviest materials as possible. As a result, their wings were not capable of supporting much weight or handling wind gusts without crumpling in flight. Early airplanes had more than one wing because this reduced the wing loading, or amount of weight that the wing had to support in flight. With an upper and lower set of wings, the wings had to support less weight and the structural fittings between them, such as struts and wires, could reinforce them (Figure 1.2). While the biplane was the primary aircraft during the early days of flight, its design also led to its downfall. The primary shortcoming of the biplane design lies in the fact that the two wings interfere with each other from an aerodynamic standpoint. In short, the two sets of wings on a biplane produce more drag and less lift than does the lone set of wings on a monoplane. The triplane, a plane with three sets of wings, is even worse in this regard. Once designers were able to make stronger, thicker wings, removing the need for the bracing and stabilizers common to biplanes, monoplanes quickly became the norm. While biplanes are not often seen today, they are still used at air shows and for agricultural purposes.


Figure 1.2: Pitts S1s Aerobatic biplane in flight

Some parameters that determine the efficiency of a biplane configuration are the stagger of the wings, and the gap distance between the two biplane wings (Figure 1.3). The biplane has significant advantages over other configurations in terms of the weight of the finished aircraft. If spanwise, as well as vertical bracing is used between the wings, the combined pair can be treated much like a single beam with each individual wing carrying bending loads and the bracing carrying shear loads. If lateral inter-wing bracing is not used, smaller weight advantages can still be realized. The use of only inboard and outboard inter-wing struts/plates gives the aircraft wings much higher torsional stiffness compared to unbraced wings of comparable weight and design. This is because the wings and bracing now form an open box when viewed from the aircraft front. If the end plates/braces on the wings can carry bending moments at the wing/brace joint then some additional weight savings can be realized because now the braces carry shear loads between the two wings.

Figure 1.3 shows the section of a two-dimensional biplane by a plane perpendicular to the span, giving two profiles, 1 and 2, whose chords are inclined at the angle $\delta$ called the dècalage, which is assumed positive when the upper wing is at greater incidence than the lower, as in the diagram.


Figure 1.3: Biplane Parameters: Dècalage, Stagger, and Gap

The gap $g$ is defined as the vertical distance between the quarter chords of the two wings perpendicular to the freestream, and stagger $s$ is the distance between the quarter chords of the two wings parallel to the freestream. The stagger can be also defined by the angle $\sigma$ between the line perpendicular to the freestream and the line joining the quarter chords of the wings, and it is equal to:

$$
\begin{equation*}
\sigma=\tan ^{-1} \frac{s}{g} \tag{1.1}
\end{equation*}
$$

There is no defined convention for positive stagger. In this report, the stagger is assumed positive when the upper wing is fore of the lower wing. Both gap and stagger are referred to the chord length of the model. The effect of dècalage is not taken into account for this experimental study. However, it has minimal effect on overall lift efficiency [1].

### 1.3 Joined Wings

The joined wing concept was first introduced as Patent $4,365,773$ by Julian Wolkovitch of Stanford University in 1980. He states in his article [2],
"The joined-wing airplane may be defined as an airplane that incorporates tandem wings arranged to form diamond shapes in both plan and front views".

The joined wing has two wings, a front wing that is swept to the rear, and a back wing that is swept toward the front. The back wing is mounted at the top of the vertical tail and extends downward at a substantial anhedral angle (negative dihedral), meeting the attaching to the front wing. From the front this has a triangular shape, whereas from the top, a diamond shape. Area ruling is employed at the place of the attachment, where the back wing begins behind the maximum thickness point of the front wing. An example of a joined-wing aircraft can be seen in Figure 1.4


Figure 1.4: Joined Wing Demonstrator (wind tunnel model) NASA Langley Research Center

Compared to an aircraft with the same wing-span, an optimally configured joined-wing aircraft may have some advantages. Unlike normal biplanes, such an arrangement can have good transonic and aerodynamic characteristics. Additional tails are probably not required, and trailing edge surfaces on both wings provide pitch and roll and can provide direct lift and side force, if desired. The main benefit, though, is the substantial reduction in wing structural weight that is achievable, on the order of $30 \%$. On the negative side, it is difficult to get a trimmed maximum lift coefficient equal to a normal wing-tail configuration, and there can be excess wetted area and interference drag with so many component intersections [3].

### 1.4 Problem Statement

The root question that will be answered at the conclusion of this report is what is the best configuration in terms of aerodynamic efficiency across different models where each model has different values of stagger and gap. How does the interference between wings behave with changes in stagger and gap? Furthermore, can this interference between wings be used to optimize the aerodynamic efficiency?

## II. Literature Review

### 2.1 Chapter Overview

Before improvements can be made, the fundamentals of flight must first be understood. This work assumes prior knowledge of fluid dynamics in general, and aerodynamics in particular, referred to a single-wing system. Therefore this chapter includes the fundamentals of flight for a biplane that will be the basis of discussion throughout the paper.

### 2.2 Lift for a Biplane

By definition, lift is the component of the aerodynamic force that is perpendicular to the relative wind.


Figure 2.1: Resultant aerodynamic force and the components into which it splits (Reproduced from [4])

In, $V_{\infty}$ is the relative wind, defined as the flow velocity far ahead of the body, and is also called the freestream velocity, $c$ is the chord, defined as the linear
distance from the leading edge to the trailing edge of the body, $\alpha$ is the angle of attack defined as the angle between $c$ and $V_{\infty}, R$ is the resultant aerodynamic force which is sometimes split into components perpendicular and parallel to the chord. By definition, $N$ is the normal force, component of $R$ perpendicular to $c, A$ is the axial force, component of $R$ parallel to $c$. Lift is therefore equal to:

$$
\begin{equation*}
\text { Lift } \equiv L=N \cos \alpha-A \sin \alpha \tag{2.1}
\end{equation*}
$$

For an incompressible, potential flow (irrotational velocity field), the lift per unit span $L^{\prime}$ on the airfoil is given by the Kutta Joukowski Theorem:

$$
\begin{equation*}
L^{\prime}=\rho_{\infty} V_{\infty} \Gamma \tag{2.2}
\end{equation*}
$$

It states that the lift per unit span is directly proportional to circulation $\Gamma$. It allows us to determine the lift generated by a body in the flow without considering the physical sources of the forces on the body, which are pressure and shear stresses.

The Kutta-Joukowski Theorem is simply an alternative way of expressing the consequences of the surface pressure distribution. In the theory of incompressible, potential flow, it is generally easier to determine circulation around the body rather than calculate the detailed surface pressure distribution. Under these assumptions, the drag generated by a body in an inviscid, incompressible flow is zero, regardless of whether or not the flow has circulation about the body. The inviscid flow assumed in the Kutta-Joukowski Theorem does not model the proper physics for drag calculation. On the other hand, the prediction of lift via equation is quite realistic.

We suppose the biplane wings consist of thin airfoils of small camber, as shown in Figure 2.2. To a first approximation the effect of one wing on the other can be calculated by means of a substitution vortex of strength $\zeta_{1}$ for the upper wing and $\zeta_{2}$ for the lower wing, placed at the quarter point of the chord of each wing [5]. The chords of the airfoils are represented by $c_{1}, c_{2}$. The stagger is represented by $s$, the angle of attack by $\alpha$, the dècalage by $\delta$, the gap by $g$, and the freestream velocity by $V_{\infty}$.


Figure 2.2: thin airfoils with vortex locations (Reproduced from [5])

In this general case we get:

$$
\begin{align*}
& \zeta_{1}=\frac{1}{2} c_{1} V_{\infty}\left(\frac{\alpha-\xi_{2}(\alpha+\delta)}{1-\xi_{1} \xi_{2}}\right)  \tag{2.3}\\
& \zeta_{2}=\frac{1}{2} c_{2} V_{\infty}\left(\frac{\alpha+\delta-\xi_{1} \alpha}{1-\xi_{1} \xi_{2}}\right) \tag{2.4}
\end{align*}
$$

where, if we introduce the mean chord $c=\frac{1}{2}\left(c_{1}+c_{2}\right)$,

$$
\begin{align*}
& \xi_{1}=\frac{\frac{1}{2} c_{1}\left(\frac{1}{2} c-s-g \delta\right)}{g^{2}+\left(\frac{1}{2} c-s\right)^{2}-\frac{1}{2} c_{2} g \delta}  \tag{2.5}\\
& \xi_{2}=\frac{\frac{1}{2} c_{2}\left(\frac{1}{2} c+s\right)}{g^{2}+\left(\frac{1}{2} c+s\right)^{2}+\frac{1}{2} c_{2} g \delta} \tag{2.6}
\end{align*}
$$

and the lift on the biplane system is

$$
\begin{equation*}
L=2 \pi\left(\zeta_{1}+\zeta_{2}\right) \rho_{\infty} V_{\infty} \tag{2.7}
\end{equation*}
$$

In the case of a biplane without dècalage, and same chord length for both wings, we have:

$$
\begin{align*}
& \xi_{1}=\frac{\frac{1}{2} c\left(\frac{1}{2} c-s\right)}{g^{2}+\left(\frac{1}{2} c-s\right)^{2}}  \tag{2.8}\\
& \xi_{2}=\frac{\frac{1}{2} c\left(\frac{1}{2} c+s\right)}{g^{2}+\left(\frac{1}{2} c+s\right)^{2}} \tag{2.9}
\end{align*}
$$

and the strength of the vortices are equal to:

$$
\begin{align*}
& \zeta_{1}=\frac{1}{2} c \alpha V_{\infty}\left(\frac{1-\xi_{2}}{1-\xi_{1} \xi_{2}}\right)  \tag{2.10}\\
& \zeta_{2}=\frac{1}{2} c \alpha V_{\infty}\left(\frac{1-\xi_{1}}{1-\xi_{1} \xi_{2}}\right) \tag{2.11}
\end{align*}
$$

the parameters $\xi_{1}$ and $\xi_{2}$ depend only on geometric parameters such as chord length, stagger, and gap.

Substituting (2.10) and (2.11) into (2.7) we get the lift of the biplane with same chord length for both wings, without dècalage:

$$
\begin{equation*}
L=c \pi \alpha \rho_{\infty} V_{\infty}^{2}\left(\frac{2-\xi_{1}-\xi_{2}}{1-\xi_{1} \xi_{2}}\right)=c \pi \alpha \rho_{\infty} V_{\infty}^{2} \Lambda \tag{2.12}
\end{equation*}
$$

Where $\Lambda$ is a geometric parameter which depends on chord length, stagger and gap. Now consider that each plane is isolated. Allowing ( $g \rightarrow \infty$ ) in equations (2.8) and (2.9), we get:

$$
\begin{equation*}
L=2 c \pi \alpha \rho_{\infty} V_{\infty}^{2} \tag{2.13}
\end{equation*}
$$

we would get the same result allowing $(s \rightarrow \infty)$ in equations (2.8) and (2.9).

### 2.3 Drag for a Biplane

By definition, drag is the component of the aerodynamic force that is parallel to the relative wind and retards the forward motion of the aircraft (see ).

$$
\begin{equation*}
\text { Drag } \equiv D=N \sin \alpha+A \cos \alpha \tag{2.14}
\end{equation*}
$$

It is important to consider viscous flow to calculate the drag. In this case, at subsonic speeds, there are two kinds of drag: parasite drag and induced drag (see Figure 2.3):

$$
\begin{equation*}
D=D_{p}+D_{i} \tag{2.15}
\end{equation*}
$$

Parasite drag, $D_{p}$, is made up of three different components: skin friction drag ( $D_{\text {skinfricioin }}$ ), form drag $\left(D_{\text {form }}\right)$, and interference drag ( $D_{\text {intefeference }}$ ). Induced drag, $D_{i}$, also known as drag due to lift, is primarily caused by wing-tip vortices that form as the aircraft produces lift.


Fligh Speed, V
Figure 2.3: Shows the two components of Drag: Parasite Drag and Induced Drag as function of Flight Speed (Reproduced from [4])

Induced drag also includes the incremental change in pressure drag due to lift (due to the change in the angle of attack). Figure 2.4 shows the effect of induced drag on a finite wing.


Figure 2.4: Effect of Induced Drag on Finite Wing (Reproduced from [4])

The angle between the chord line and the direction of $V_{\infty}$ is the angle of attack $\alpha$. We now define $\alpha$ as the geometric angle of attack. In Figure 2.4 the local relative wing is inclined below the direction of $V_{\infty}$ by the angle $\alpha_{i}$ called the
induced angle of attack. This downward component of velocity is called downwash, denoted by the symbol $w$. The presence of downwash, and its effect on inclining the local relative wind in the downward direction, has two important effects on the local airfoil section:

1. The angle of attack actually seen by the local airfoil section is the angle between the chord line and the local relative wind. This angle is called effective angle of attack:

$$
\alpha_{e f f}=\alpha-\alpha_{i}
$$

2. The local lift vector is aligned perpendicular to the local relative wind. Consequently, there is a component of the local lift vector in the direction of $V_{\infty}$; that is, there is a drag created by the presence of downwash. This drag is defined as induced drag, denoted by $D_{i}$ in Figure 2.4.

Induced drag is a function of the wing spanload only, and is independent of the details of the particular airfoil used in the wing. The additional profile drag is associated with the airfoil used in the wing. At low lift coefficients this drag should be small, only becoming important as flow separation starts to develop on the airfoil section. The additional profile drag becomes large as wing stall is approached.

Because of biplane lift interference effects, Munk studied the induced drag of a biplane due to the interference of one wing on the other. These studies led to a main conclusive theory [6]:

> "The total induced drag of any multiplane lifting system is unaltered if any of the lifting elements are moved in the direction of the motion provided that the attitude of the elements is adjusted to maintain the same lift distribution of lift among them."

This theorem is known as "Munk's Stagger Theorem", and basically states three important results:

1. If constant section lift is maintained, the chordwise pressure distribution does not affect the induced drag.
2. If the spanwise lift distribution is constant, there will be no effect on the induced drag from a change in biplane stagger or wing sweep.
3. The sum of all the lifting surfaces can be made equivalent to a single lifting element, enabling easier calculation of the induced drag.

Furthermore, Munk stated that if the two wings of the biplane are parallel and unstaggered, the downwash of each wing induced by the other wing is equal.

Munk then concludes with a formula for determining the induced drag coefficient of a biplane compared to that of a monoplane, where $C_{D_{1}}$ and $C_{L}$ are the induced drag and lift coefficients of the monoplane, respectively. The subscripts denote terms relating to the monoplane and biplane, and $k$ denotes the "equivalent monoplane span" factor:

$$
\begin{equation*}
C_{D_{2}}=C_{D_{1}}-\frac{C_{L}^{2}}{\pi}\left[\frac{S_{1}}{b_{1}^{2} k_{1}^{2}}-\frac{S_{2}}{b_{2}^{2} k_{2}^{2}}\right] \tag{2.16}
\end{equation*}
$$

where $S$ is the entire area and $b$ the greatest span. $k_{1}$ and $k_{2}$ are factors which depend merely on the gap/span ratio of the biplane and assume the value $k=1$ for
a monoplane. If the two spans of a biplane are slightly different, an average span is to be substituted. The values of $k$ are determined by Munk empirically as described in a former paper [6], and are given in Table 2.1 and Figure 2.5.

Table 2.1: values of Munk's $k$ at different gap/span ratio

| gap/span | $\mathbf{k}$ |
| :---: | :---: |
| 0 | 1 |
| 0.05 | 1.02 |
| 0.1 | 1.05 |
| 0.15 | 1.09 |
| 0.2 | 1.15 |
| 0.3 | 1.21 |
| 0.4 | 1.25 |
| 0.5 | 1.27 |

The differences are not very great. In view of the fact that the comparison has been made with one wing section only, and that it is difficult to obtain exact values of $k$, these values are not very reliable and an average curve must be taken.


Figure 2.5: values of Munk's $k$ plotted against the gap/span ratio
The result of the calculation of the drag coefficient is practically unaffected by this small change of $k$. For a rough estimate it is even sufficient to take $k=1$ for all monoplanes and $k=1.1$ for all biplanes. Equation (2.16) basically states that, at the same lift coefficient, the induced drag of a biplane will be smaller than that of a
monoplane with the same span. Munk then found that, due to induction and interference between the upper and lower wing sections, a biplane will experience an induced angle of attack, causing a greater angle of attack than that of a monoplane with the same lift coefficient.

For the same reasons of induction and interference Munk concluded that the shift in the center of pressure CP due to a change in the lift coefficient for the monoplane and the biplane were about the same, and when the shift is due to a change in the angle of attack, the CP travels an even smaller distance. Even though the difference in travel of the center of pressure between the two configurations is small, the biplane chord is only about half the length of a monoplane chord having the same airfoil section and lift, thus experiencing only about half the overall CP travel compared to the monoplane. This fact proves very advantageous in determining the stability characteristics of the biplane.

The final equation developed by Munk [6] for obtaining the induced drag coefficient, and that will be used in the theoretical and experimental comparison, is Equation (2.17)

$$
\begin{equation*}
C_{D_{i E M}}=\frac{C_{L}^{2}}{\pi} \frac{S}{b^{2} k^{2}} \tag{2.17}
\end{equation*}
$$

where the subscript EM is referred to the equivalent monoplane factor $k$.
Prandtl, then collaborating with Munk on biplane theory, reaffirmed Munk's stagger theorem by stating that the sum of the induced downwash between the two wings will remain constant, given any longitudinal change in geometry, and at angles of attack such that the lift is constant [7]. Prandtl states
that in an unstaggered wing system, the drag $D_{12}$, induced by wing 1 on wing 2 , equals the drag $D_{21}$ induced by wing 2 on wing 1 . The drag $D_{12}$ is due to the fact that wing 1 produces a downward air current toward wing 2. This value of $D_{12}$ was found to be equal to:

$$
\begin{equation*}
D_{12}=D_{21}=\frac{\sigma}{\pi q} \frac{L_{1}}{b_{1}} \frac{L_{2}}{b_{2}} \tag{2.18}
\end{equation*}
$$

where $\sigma$ is the coefficient of mutual influence which depends upon gap/span ratio, and it is equal to:

$$
\begin{equation*}
\sigma \cong \frac{1-0.66 \mathrm{~g} / \mathrm{b}}{1.05+3.7 \mathrm{~g} / \mathrm{b}} \tag{2.19}
\end{equation*}
$$

Values of $\sigma$ are given in Figure 2.6 for different $g / b$ and different $r=b_{2} / b_{1}$.


Figure 2.6: values of $\sigma$ plotted against the gap/span ratio and different $r=b_{2} / b_{1}$ (Reproduced from [7])
The induced drag of the upper wing, for the unstaggered biplane, is

$$
\begin{equation*}
D_{1}=D_{11}+D_{12}=\frac{1}{\pi q}\left(\frac{L_{1}^{2}}{b_{1}^{2}}+\sigma \frac{L_{1}}{b_{1}} \frac{L_{2}}{b_{2}}\right) \tag{2.20}
\end{equation*}
$$

and that of the lower wing is

$$
\begin{equation*}
D_{2}=D_{21}+D_{22}=\frac{1}{\pi q}\left(\sigma \frac{L_{1}}{b_{1}} \frac{L_{2}}{b_{2}}+\frac{L_{2}^{2}}{b_{2}^{2}}\right) \tag{2.21}
\end{equation*}
$$

where there is a positive stagger, the drag of the upper wing is diminished by the upward air currents produced by the lower wing; but, on the other hand, the drag of the lower wing is increased, by exactly the same extent, by the downward air current produced by the upper wing, so that the total drag is the same as in the case of the unstaggered biplane and is equal to:

$$
\begin{equation*}
D=D_{1}+D_{2}=\frac{1}{\pi q}\left(\frac{L_{1}^{2}}{b_{1}^{2}}+2 \sigma \frac{L_{1}}{b_{1}} \frac{L_{2}}{b_{2}}+\frac{L_{2}^{2}}{b_{2}^{2}}\right) \tag{2.22}
\end{equation*}
$$

Also Prandtl [7] gives a final equation for obtaining the induced drag coefficient, and that will be used in the theoretical and experimental comparison, that is Equation (2.23)

$$
\begin{equation*}
C_{D_{i o s}}=\frac{C_{L}^{2}}{\pi} \frac{S}{b^{2}}(1+\sigma) \tag{2.23}
\end{equation*}
$$

where the subscript OM is referred to the orthogonal biplane theory developed by Prandtl.

### 2.4 Wing-Tip Vortices

A finite wing is a three dimensional body, and consequently the flow over the finite wing is three dimensional; that is, there is a component of flow in the spanwise direction. Figure 2.7 gives the top view of a finite wing.


Figure 2.7: Finite wing Streamline Curvature (Reproduced from [4])

The primary cause for the formation of wing-tip vortices is lift. Lift, generated by a wing, results when there is a net pressure difference between the upper and lower surfaces of the wing. When the higher pressure exists on the lower surface of the wing, then positive lift is created. As a by-product of this pressure imbalance, the flow near the tips tends to curl around the tips. This generates on the top surface of the wing, a spanwise component of flow from the tip toward the root. Similarly, on the bottom surface of the wing, there is a spanwise component of flow from the root toward the tip, as shown in Figure 2.7. This flow establishes a circulatory motion that trails downstream of the wing; that is, a trailing vortex is created at each wing tip. These wing-tip vortices are sketched in Figure 2.8.


Figure 2.8: Schematic of wing-tip vortices (Reproduced from [4])

The aerodynamic theory of a finite wing now becomes more complicated, due to the three-dimensionality of the flow. The explanation of this theory can be found in the next section where the concepts of vortex filament, the Biot-Savart law, and Helmholtz's Theorem will be the tools by which obtaining analytical results using a Vortex Lattice Code called AVL (Athena Vortex Lattice) is made possible.

### 2.5 Vortex Lattice Method

### 2.5.1 Overview

The vortex lattice method is an extension of Prandtl's classical lifting line representation of a wing. Prandtl's classical lifting line theory gives reasonable results for straight wings at moderate to high aspect ratio. However, for low-aspect-ratio straight wings, swept wings, and delta wings, classical lifting line theory is inappropriate. For such planforms a more sophisticated model must be used. This model is called the Vortex Lattice Method. In this approach, the wing is represented by a series of lifting lines on the plane of the wing, at different chordwise and spanwise stations. Each vortex loop is located in each station and consisting of a bound vortex, infinitely long trailing vortex lines and a control point. The circulation $\Gamma$ of the bound vortex generates a discontinuity in the tangential velocity component, corresponding to the velocity difference on the pressure and suction side of the wing. The lift force on the wing is calculated from
the strength of the bound vortex using the Kutta-Joukowski law, stated in Equation (2.2).

### 2.5.2 Introduction

The physical problem addressed was to find the aerodynamic forces acting on an aircraft flying at low subsonic speeds, below the stall limit. The governing equations used to solve the physical problem came from standard vortex lattice theory. The law of Biot-Savart was used to get the flowfield around a finite straight vortex line, one of the basic vortex segments needed for the lattice. These vortices induce a flow field in the air, and their strength was determined by the boundary conditions that no air should flow through the wings.

The forces acting on each vortex segment can be determined by employing the Kutta-Joukowski theorem. These forces may then be integrated to yield a composite force in three dimensions, which in turn may be used to compute aerodynamic coefficients. The computational problem is to create a good system for dealing with the mathematical results.

### 2.5.3 Physical Problem

As air flows around the airframe of an aircraft forces build up. These forces can be derived from pressure and friction acting on every free surface of the interface between the fluid (air) and the airframe (wings, body control surfaces etc.). The resulting force acting on the aircraft is given by integrating the distributed forces across the interface as shown in Figure 2.9


Figure 2.9: Showing aircraft with resulting force. Partial blowup showing pressure and tangential forces. (Reproduced from [23])

The force acting on an infinitesimal area of the interface can be divided into two components. One of them is perpendicular to the surface and the other is parallel. The physical interpretation being that the perpendicular force is a pressure force and the parallel force a shear, or friction, force.

The pressure forces are the forces acting along the normal of the interface. The total pressure in the fluid is constant and consisting of the static and the dynamic pressure. At rest, the total pressure is equal to the static pressure and constant at every point of the interface. Once in motion, the airframe creates a flow field in the air. As the interface is impenetrable the fluid must move to allow the passage of the airframe thus creating the flow field. The shape of this field is highly dependent of the shape of the interface. Thus parts of the airframe facing the wind experience a higher static pressure than areas parallel to the wind or on the offwind side. This accounts for the pressure forces on the interface. At the microscopic level, the smallest divisions of air located closest to the skin of the
airframe must be stationary relative the airframe. If we move along the normal of the surface the tangential speed picks up and at a certain distance the tangential speed becomes constant along the rest of the normal. This distance constitutes the thickness of the boundary layer as shown in Figure 2.10. The viscous motion of the boundary layer is the source of the friction forces. The thickness of this boundary layer is often small in comparison with other typical lengths in the interface.


A


Figure 2.10: The boundary layer, tangential speed and viscous movement. The arrows are velocity vectors of the air. A shows the profile and $B$ the local enlargement with visible boundary layer. (Reproduced from [23])

### 2.5.4 Mathematical Problem

## Potential Flow

When considering a vector field where the rotation along any closed path is zero, the conclusion arises that the line integral between two points in the field is independent of the path. Such a flow field is called conservative. The concept of conservative fields allows the definition of the potential:

$$
\begin{equation*}
A=\nabla \phi \tag{2.24}
\end{equation*}
$$

where $\phi$ equals the potential of $A$.
This mathematical theory can be, and is, applied to many physical problems. Among these are the theory of electric potential and the theory of velocity
potentials in flow fields. In the case of a fluid flow the field is defined as follows for an irrotational, inviscid, incompressible flow [4].

No mass is produced in the field so,

$$
\begin{equation*}
\nabla \cdot \vec{V}=0 \tag{2.25}
\end{equation*}
$$

Further, as the flow is irrotational,

$$
\begin{equation*}
\vec{V}=\nabla \phi \tag{2.26}
\end{equation*}
$$

Hence,

$$
\begin{equation*}
\nabla \cdot(\nabla \phi)=0 \tag{2.27}
\end{equation*}
$$

or

$$
\begin{equation*}
\nabla^{2} \phi=0 \tag{2.28}
\end{equation*}
$$

Equation (2.28) is Laplace's equation for which a number of solutions can be found.

Anderson [4] continues to evolve the matter and states that: "A complicated flow pattern for an irrotational, incompressible flow can be synthesized by adding together a number of elementary flows, which are also irrotational and incompressible." Such elementary flows may be the point source, the point sink, the doublet and the vortex line. These may be superpositioned in many ways including the formation of line sources, vortex sheets and so on. As one may use an arbitrary number of singularities the concept of using numerical methods is close at hand. Today, a wide variety of methods exist and one of them is the vortex lattice method (VLM).

## Vortex Theorem

In using vortex singularities to model lifting surfaces, let review some properties of vortices. The key properties are defined by the so-called vortex theorems. These theorems are associated with the names of Kelvin and Helmholtz. Three important results are:

1. Along a vortex line (tube) the circulation, $\Gamma$, is constant.
2. A vortex filament (or line) cannot begin or end abruptly in a fluid. The vortex line must i) be closed, ii) extend to infinity, or iii) end at a solid boundary. Furthermore, the circulation, $\Gamma$, about any section is the vortex strength.
3. An initially irrotational, inviscid flow will remain irrotational.

Related to these theorems, an important result can be stated:

- A sheet of vortices can support a jump in tangential velocity [i.e. a force], while the normal velocity is continuous. This means that a vortex sheet can be used to represent a lifting surface.


### 2.6 AVL

Athena Vortex Lattice [8] (AVL) is a code developed by Youngren and Drela that utilizes vortex-lattice theory for aerodynamic and dynamic stability analysis of a given aircraft geometry. By changing parameters ( $\mathrm{C}_{\mathrm{L}}, \mathrm{C}_{\mathrm{D}}, \mathrm{a}$, etc.), an AVL model can be used to predict aircraft performance in real-world scenarios. To create a model, AVL requires a text file that defines the aircraft geometry.

An example of model created using AVL is shown in Figure 2.11


Figure 2.11: Example of AVL Model

### 2.7 Previous Experiments Conducted on Biplanes

### 2.7.1 Overview

Although these early theories are very important, they lacked experimental confirmation of their accuracy and did not take into account the effect of streamline curvature, non-elliptical lift distribution, or any geometrical variables. For these reasons, some discrepancy between the theories and actual biplane performance were encountered.

### 2.7.2 Early Experiments

Soon after Munk's biplane theories were published, J.C. Hunsaker performed an experimental analysis on the inherent longitudinal stability of a "typical" biplane. The aircraft he tested did not vary any geometric parameters, only the aircraft angle of attack was varied to determine the lift, drag, and pitching moment characteristics and their effect on the stability of the aircraft. His results showed that the biplane was an inherently unstable aircraft configuration at low speeds and high angles of attack [9].

In 1918, F.H. Norton conducted an investigation similar to Hunsaker's work utilizing a three-dimensional, non-symmetric biplane model to determine the effect of staggering the wings. All other variables held constant, Norton found that the maximum efficiency and lift are achieved at the highest degree of stagger possible. Furthermore, the travel of the CP was greatly reduced with large positive stagger, which could ease in solving the dynamic stability problem [10].
H. Glauert of the Royal Aircraft Establishment then incorporated a new variable into Munk's angle of incidence formula to include an improved method of determining the effect of streamline curvature [11]. His results showed accurate correlation with experiment for positive stagger but did not give good results at negative stagger.

As flight research and the biplane became more popular, studies became more prominent throughout the United States and Europe. New concepts had been formulated, old ones improved and many experimental tests had begun. In 1929, a series of papers were written by Montgomery Knight and Richard W. Noyes [12] [13] [14]. They conducted wind tunnel tests on three dimensional asymmetrical biplane airfoils while varying the gap, stagger, decalage, dihedral, sweepback, overhang, and combinations thereof. Many useful results were obtained, some of which are as follows:

1. Increasing the gap or stagger in the positive direction tends to equalize the loads on the two wings, and also increases the maximum total lift coefficient of the biplane.
2. An increasing gap or lower sweep tends to decrease the travel of the center of pressure.
3. The deviation of decalage angle from zero tends to decrease the maximum lift coefficient when there is zero stagger.
4. With positive stagger, increasing the decalage angle tends to increase the maximum lift coefficient.

The last of these results may cause some confusion, but the reasoning is relatively straight-forward. To reach the maximum lift possible, the two wings of the biplane must stall at the same time. For this to occur, the wings must also be at the same effective angle of attack. When the wings are at positive stagger, there is an increased amount of downwash imposed by the upper wing onto the lower wing, causing a reduced effective angle of attack on the lower wing. Therefore, for the effective angles of attack to be equal, and hence stall at the same time, the lower wing's angle of incidence must be increased, creating a negative angle of decalage in the biplane wing configuration. These tests were all conducted at a low Reynolds number of 150,000.

Max Munk, still very interested in biplane theory, was influenced by the Bureau of Aeronautics of the Navy to conduct a series of tests on biplane and triplane models [15]. These tests were to determine the lift, drag, and pitching moment for different airfoils, systematically varying the gap and stagger, and then
to compare the results with the Army standards. The results of these tests were very interesting:

1. There was a general tendency of the upper wing to contribute more of the lift than the lower wing at positive stagger and less at negative stagger.
2. The gap/chord ratio had little effect on the relative lifts of the wings at high lift coefficients, but significant effects at low lift coefficients.
3. An increase in gap tends to equalize the lift of the wings over a wide range of angles of attack.
4. The gap/chord ratio had little effect on the positions of the center of pressures of the individual wings.
5. With an increase in positive stagger, the centers of pressure moved forward on the upper wing and aft on the lower wing, lying nearly together at zero stagger.

These results are very similar to the previous studies by Knight and Noyes, and would seem to verify the accuracy of both tests. Furthermore, Munk used two different airfoils, the RAF-15 and the USATS-5, to ensure that his own tests would give results unbiased to a given wing profile.

### 2.7.3 New Theories

While these extensive experimental tests were being performed, a few new theories on the biplane and its interference characteristics had been developed.

In 1930, Clark B. Millikan adopted a theory from Dr. Theodore Von Karman known as "Thin Airfoil Theory". In his paper, Millikan presented and used Von Karman's theory to develop a procedure for determining the characteristics of the individual wings of an arbitrary biplane configuration without sweepback or dihedral [16]. Although this process showed great success over current theories when compared with experimental data (at gap/chord ratios greater than $3 / 4$ ), the procedure was very tedious and cumbersome, and was therefore rarely incorporated into use by the designers of that day.

In 1933 Walter S. Diehl published a report on biplane theory [17]. His paper combined experimental and theoretical data by Fuchs and Hopf to obtain a series of curves from which the lift curves of the individual wings could be found. His results showed promise, but even Diehl agreed that:
"Millikan's treatment of the biplane theory...appear to give somewhat better agreement with test data...but it is very difficult for an engineer to follow the steps required in a typical calculation."

### 2.8 Comparison of Monoplane and Biplane Results

In all of the aforementioned studies, the only comparison of biplane and monoplane performance characteristics was theoretical. Actual experimental testing of the biplane compared to the monoplane had not been done, but in July of 1974, E. Carl Olson wrote his M.S. on the improved aerodynamic characteristics of
a biplane over that of a monoplane by comparing experimental data [18]. Olson's experiments consisted of a three-dimensional asymmetrical airfoil biplane configuration, and incorporated the results of Nenadovitch by varying the geometry about his optimum point: a gap of one chord length, a stagger of one chord length, and a decalage angle of negative six degrees. Furthermore, Olson also tested a monoplane system using the same area and similar aspect ratio of the biplane configuration.

The aspect ratio considered in his report is equal to:

$$
\begin{equation*}
A R=\frac{b^{2}}{S} \frac{\mu^{2}(1+r)^{2}}{\mu^{2}+2 \sigma \cdot r+r^{2}} \tag{2.29}
\end{equation*}
$$

where $S$ is the wing area, $b$ the span of the upper wing, $\sigma$ is the Prandtl interference factor defined in Equation (2.19), $\mu$ is the ratio of lower wing span over upper wing span, and $r$ is the ratio lower wing area over upper wing area.

His tests were also conducted with and without a fuselage. The results obtained are probably the most significant to date.

In the first phase of testing, Olson concluded that the best range of decalage angles to continue testing at were between $-5^{\circ}$ and $-7^{\circ}$, due to the high $L / D$ and low $C_{D}$ at these angles. Therefore, the remainder of his testing occurred between these angles.

In the second phase of his testing, a fuselage was incorporated into the configuration. In these tests, gap, stagger, and decalage angle were all varied, and
then the results of the biplane and monoplane configurations were compared and plotted against each other. There were many important conclusions obtained:

1. A substantial $C_{D}$ reduction with respect to the monoplane over a wide range of angles of attack was obtained for most biplane configurations, the most efficient showing a $25 \%$ decrease in the $C_{D}$ over the monoplane in a typical cruise condition.
2. A significant $L / D$ ratio increase for the biplane configuration was obtained over a wide range of lift conditions with respect to the monoplane. The largest increase was $31.2 \%$ at the maximum $L / D$, with $C_{D}$ being $21.4 \%$ lower than the monoplane at a $C_{L}$ of 0.175 .
3. The endurance of the biplane would be increased over a wide range of $C_{L}$ over that of the monoplane due to the $L^{3 / 2} / D$ curve.
4. While creating higher interference drag, the biplane realized a substantial increase in efficiency over the monoplane due to a decreased induced drag and/or altered pressure distribution over the wings, creating an overall reduction in the total drag.
5. The most efficient overall biplane configuration increased $L / D$ by $16.3 \%$, reduced the $C_{D}$ by $14.3 \%$ (at a $C_{L}$ of 0.175 ), but had a $10.6 \%$ decrease in $C_{L_{\max }}$ compared to the monoplane.
6. Increasing stagger tended to decrease the overall $C_{D}$.
7. Pitching moment characteristics of the biplane system were markedly improved over the monoplane system (a more negative moment curve slope).

Throughout all of the experimental testing and theoretical evaluations in the past, along with the recent studies, it is evident that a biplane wing configuration can hold many aerodynamic advantages over that of the monoplane.

### 2.9 Research Focus

The purpose of this report is to further investigate how gap and stagger influence the aerodynamic performance of a biplane joined at the tips by endplates. This parametric study has led to the selection of fourteen cases which have been test run first in the Athena Vortex Lattice (AVL) code, developed by Youngren and Drela, based on the Vortex Lattice Method (VLM). This method is based on solutions to Laplace's Equation and was used to estimate the inviscid aerodynamic coefficients and performance characteristics analytically. The fourteen different models selected for wind tunnel testing are represented in Table 2.2

Table 2.2: Model Configurations for wind tunnel testing

| Moclels |  |  |  |
| :---: | :---: | :---: | :---: |
| Gap | Stagger | $\boldsymbol{\sigma}\left(^{\circ}\right)$ | Chord Length (in) |
| 1 c | 0 c | 0 | 3.995 |
| 0.5 c | 0 c | 0 | 4.004 |
| 0.5 c | 0.5 c | 45 | 3.943 |
| 0.5 c | 1 c | 63.4 | 3.892 |
| 1 c | 0.5 c | 26.6 | 3.967 |
| 1 c | 1 c | 45 | 4.005 |
| 2 c | 1 c | 26.6 | 3.964 |
| 1 c | 1.5 c | 56.3 | 3.917 |
| 0.5 c | -0.5 c | -45 | 3.943 |
| 0.5 c | -1 c | -63.4 | 3.892 |
| 1 c | -0.5 c | -26.6 | 3.967 |
| 1 c | -1 c | -45 | 4.005 |
| 2 c | -1 c | -26.6 | 3.964 |
| 1 c | -1.5 c | -56.3 | 3.917 |

The wing is based on a flat plate airfoil section of thickness over chord ratio $t / \mathrm{c}=1.18 \%$. A wing of $12^{\prime \prime}$ semi-span in a vertical orientation was used to match the size and Reynolds number with a full-sized UAV. Asymmetric flow is prevented by mounting the model on the floor.

Such an arrangement, though obviously unsuited for yaw tests, yields accurate pitch, lift, and downwash data at the maximum Reynolds number [19]. The gap between the half wing and the ground was measured to be $0.04^{\prime \prime}$ and it was filled with petroleum jelly to prevent interaction between the flow and the model in such a gap.

In order to achieve aerodynamic analyses of the models, the entire test matrix was repeated three times in the University of Dayton/University of Dayton Research Institute Low Speed Wind Tunnel (LSWT).

## III. Methodology

### 3.1 Experimental Equipment

### 3.1.1 University of Dayton Low-Speed Wind Tunnel (LSWT)

The experimental integrated force data were obtained from wind-tunnel tests performed in the University of Dayton/University of Dayton Research Institute Low-Speed Wind Tunnel (LSWT) seen in Figure 3.1.


Figure 3.1: University of Dayton Low Speed Wind Tunnel (LSWT)

The DART CORPORATION constructed the tunnel in 1992. The LSWT is an Eiffel type with an 16:1 contraction ratio. The fan was designed and constructed by Hartzell and is driven by a 60 HP motor. The test section measures $30^{\prime \prime} \times 30^{\prime \prime} \times 90^{\prime \prime}$ ( $\sim 0.75 \times 0.75 \times 2.3 \mathrm{~m}$ ), and is easily exchanged with other test sections used for
demonstration and educational purposes to preserve the quality of the research test section. The LSWT has five anti-turbulence screens at the tunnel inlet.

The highest flow quality operable speed range of the LSWT is from $6.7 \mathrm{~m} / \mathrm{s}$ $(20 \mathrm{ft} / \mathrm{s})$ to $36.7 \mathrm{~m} / \mathrm{s}(120 \mathrm{ft} / \mathrm{s})$. The tunnel has a turbulence intensity in the freestream direction of less than $0.1 \%$ throughout the test section (measured by hot wire anemometer), and less than $0.05 \%$ throughout the center portion of the test section utilized for testing the different models.

The test models are mounted on a rotary stage which is located on the top of a force balance. The connections between the model and the rotary stage, the rotary stage and the balance, the balance and the support were designed by the author. The system balance - rotary stage is fixed to a support located underneath the tunnel as shown in Figure 3.2.


Figure 3.2: Rotary stage and balance

### 3.1.2 LSWT 25 lb Strain Gage Balance

The force balance used to determine the forces on the models is based on an ATI Gamma F/T Transducer. Figure 3.3 shows the balance.


Figure 3.3: ATI Gamma F/T Transducer (Credit - ATI Industrial Automation)

The transducer senses applied loading with six degrees of freedom (Fx, Fy, $\mathrm{Fz}, \mathrm{Tx}, \mathrm{Ty}$, and Tz ) and it has the interface board inside the transducer. The transducer cable is attached with a connector and it is specially designed for noise immunity. This durable cable protects the transducer signals from electrical fields and mechanical stress. The controller interfaces with the transducer to process the transducer data into usable force and torque data and to provide high-level functions. The controller is powered by standard AC power. This controller communicates over an RS-232 serial port and can also output loads via analog voltages.

The interface board electronics receive transducer gauge signals and convert them to readable DAQ card signals using noise immunity technology. Each interface board is calibrated. The interface board is mounted within the sensor. Since the wind tunnel installed transducer output is uncalibrated, an in-house
high-resolution calibration was performed by the author and a calibration matrix subsequently created.

Figure 3.4 shows the axes of the sensor.


Figure 3.4: Axes of the sensor

A list of the maximum allowable forces and moments is listed in Table 3.1. If forces or moments exceed the allowable range then the balance could be damaged and thus invalidate the calibration. Sensing ranges and resolution (typical for a 16bit data acquisition system) associated with the calibration matrix are listed in Table 3.2. The other dimensions of the balance can be seen in Appendix A.

Table 3.1: Max Allowable Forces and Moments for the Balance

| Component | Max Load |
| :--- | :--- |
| Lift $-\mathrm{F}_{X}$ | $\pm 235 \mathrm{lb}$ |
| Drag- $\mathrm{F}_{Y}$ | $\pm 235 \mathrm{lb}$ |
| Side $-\mathrm{F}_{Z}$ | $\pm 736 \mathrm{lb}$ |
| Yaw $-\mathrm{T}_{X}$ | $\pm 618 \mathrm{in}-\mathrm{lb}$ |
| Roll $-\mathrm{T}_{Y}$ | $\pm 618 \mathrm{in}-\mathrm{lb}$ |
| Pitch $-\mathrm{T}_{Z}$ | $\pm 727 \mathrm{in}-\mathrm{lb}$ |

Table 3.2: Sensing ranges and Resolution

| Component | Rated <br> Sensing <br> Ranges | Resolution |
| :--- | :--- | :--- |
| Lift $-\mathrm{F}_{\mathrm{X}}$ | $\pm 7.5 \mathrm{lb}$ | $1 / 2560 \mathrm{lb}$ |
| Drag- $\mathrm{F}_{\mathrm{Y}}$ | $\pm 7.5 \mathrm{lb}$ | $1 / 2560 \mathrm{lb}$ |
| Side $-\mathrm{F}_{\mathrm{Z}}$ | $\pm 25 \mathrm{lb}$ | $1 / 1280 \mathrm{lb}$ |
| Yaw $-\mathrm{T}_{\mathrm{X}}$ | $\pm 25 \mathrm{in}-\mathrm{lb}$ | $1 / 1280 \mathrm{in}-\mathrm{lb}$ |
| Roll $-\mathrm{T}_{\mathrm{Y}}$ | $\pm 25 \mathrm{in}-\mathrm{lb}$ | $1 / 1280 \mathrm{in}-\mathrm{lb}$ |
| Pitch $-\mathrm{T}_{\mathrm{Z}}$ | $\pm 25 \mathrm{in}-\mathrm{lb}$ | $1 / 1280 \mathrm{in}-\mathrm{lb}$ |

### 3.1.3 LSWT 275 lb Rotary Stage

Due to the sidewall mount orientation of the model in the test section, a rotary stage Model RM-5 manufactured by Newmark was used to change the angle of attack of the models. The rotary stage is designed with a maximum load of $\pm 275 \mathrm{lbs}$ and it uses a single axis controller. It provides a resolution of $0.36 \mathrm{arc}-\mathrm{sec}$ and an accuracy of 60 arc-sec. Figure 3.5 shows the rotary stage.


Figure 3.5: Newmark RM-5 Rotary Stage (Credit -Newmark Systems Inc.)

The specifications are listed in Table 3.3

Table 3.3: Rolary Stage Specifications

| Repeatability | $5 \mathrm{arc}-$ seconds |
| :---: | :---: |
| Resolution | $0.36 \mathrm{arc}-$ seconds |
| Accuracy | $72 \mathrm{arc}-$ seconds |
| Gear Ratio | $72: 1$ |
| Max. Load | 275 lbs |
| Moment | $260 \mathrm{in}-\mathrm{lb}$ |
| Max. Speed | 1200 RPM |
| Travel | $360^{\circ}$ Continuos |

The other dimensions of the rotary stage can be seen in Appendix B.

### 3.2 Collecting and Processing Data

### 3.2.1 Overview

Measurements can be taken by the balance once the wind tunnel reaches a desired velocity. Angle of attack can be accomplished by rotating the rotary stage and subsequent model using the angle control device. Data acquisition was performed on a Pentium IV PC running LabVIEW 8.0 and software written inhouse driving a PCI 6281 performing simultaneous sampling at 1000 Hz passed through an SCXI Chassis and SCXI 1140 hardware filter. Ten thousand samples were recorded in each data sweep, and these values were acquired four times for any given angle of attack and tunnel speed to provide 40,000 samples per data point represented.

All data files for each test run were stored on the hard drive of the acquisition system and were later retrieved for data reduction. The acquisition recorded the following values: $\mathfrak{a}$, tunnel speed, lift force, drag force, side force,
pitch moment, yaw moment, roll moment, and relative standard deviations. All forces and moments were measured about the balance center. The balance records the force data by comparing voltage measurements to the calibrated voltage measurements. This comparison allows the forces that act on the balance to be determined.

The balance - rotary stage system is outside the wind tunnel test section (see Figure 3.2) but the support structure rises $0.04^{\prime \prime}$ from the bottom wall of the test section. For this reason, before and after testing each model, an aerodynamic tare was run on the support structure alone to account for the drag of the support. Between each change of airspeed during testing, the wind tunnel velocity was brought back to zero to account for any balance drift.

### 3.2.2 Density, Viscosity, and Velocity

In the LSWT, an Angle of Attack from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on the models at two different Reynolds Numbers. Tunnel temperature, atmospheric pressure and air humidity were measured numerous times during execution of the experiments and were used in the correction of calculated velocity to real tunnel velocity.

Each model has a different chord length as listed in Table 2.2. The Reynolds number was based on a chord length equal to $\bar{c}=4^{\prime \prime}$ and is equal to:

$$
\begin{equation*}
\mathrm{Re}=\frac{V_{z} \cdot \bar{c} \cdot \rho_{a i t}}{\mu_{a i r}} \tag{3.1}
\end{equation*}
$$

The air density, $\rho_{\text {air }}$, is calculated using Equation (3.2):

$$
\begin{equation*}
\rho_{\text {air }}=1.2929\left(\frac{273.15}{T+273.15}\right)\left(\frac{p_{\text {atm }}-S V P \cdot H}{760}\right) \tag{3.2}
\end{equation*}
$$

Where $T\left({ }^{\circ} \mathrm{C}\right)$ is the tunnel temperature, $p_{\text {arm }}(m m H g)$ is the atmospheric pressure, SVP ( mmHg ) is the Saturation Vapor Pressure which depends on the temperature, and $H$ is the air humidity.

The air viscosity, $\mu_{\text {air }}$, is calculated using Sutherland's formula:

$$
\begin{equation*}
\mu_{\text {air }}=\mu_{0} \frac{T_{0}+C}{T+C}\left(\frac{T}{T_{0}}\right)^{3 / 2} \tag{3.3}
\end{equation*}
$$

For standard air, $\mu_{0}$ is the reference viscosity equal to $1.827 \cdot 10^{-5} \mathrm{~Pa} \cdot \mathrm{~s}$ at reference temperature $T_{0}(K)$ equal to $291.15 K, C$ is the Sutherland's constant equal to 120 , and $T(K)$ is the input temperature. Once the Reynolds Number was chosen, the freestream velocity was obtained.

### 3.3 Test Plan

The test were performed on all the models in Table 2.2 at two different Reynolds numbers 60,000 and 120,000. Each model was placed in the LSWT and mounted on the rotary stage. Data were taken for all configurations for 109 different angles of attack, $\alpha$, from $-2^{\circ}$ to $25^{\circ}$ by $0.25^{\circ}$ increments. Testing procedures were repeated for all models. The dates of the testing and the corresponding temperature, pressure, and humidity can be seen in Appendix $C$.

## IV. Results and Analysis

### 4.1 Chapter Overview

In this chapter, results for the wind tunnel tests conducted on the fourteen biplane configurations will be shown. The aerodynamic data for all fourteen models in Table 2.2 can be seen in Appendix D for Reynolds number 60,000 and in Appendix E for Reynolds number 120,000.

The results will be divided into five main sections. The first section will compare the results of models of similar gap at each speed. The second section will compare the results of models of similar stagger at each speed. The third section will compare the differences between positive stagger and negative stagger for each model having stagger at each speed. The fourth section will describe the change in lift slope at each speed. The fifth section will discuss the aerodynamic hysteresis. The sixth section will compare analytical and experimental results. The seventh section will discuss the method by which obtaining uncertainty analysis on aerodynamic efficiency and lift coefficient is made possible.

### 4.2 Wind Tunnel Balance Data - Similar Gap Models

In this section the effect of stagger on the performance of the biplane joined at the tips will be presented. For this analysis, data from models having similar gap will be presented to better understand the behaviour of changes in stagger. The aerodynamic performance analyzed will be essentially the lift coefficient, the drag polar, and the aerodynamic efficiency of the different models.

### 4.2.1 Gap 0.5c

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on the models having similar gap, $g=0.5 \mathrm{c}$, at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s}), \operatorname{Re}=120,000(\approx 18 \mathrm{~m} / \mathrm{s})$.

## Lift Coefficient

The Reynolds number does not affect the variation of lift coefficient for the different models. For this reason, the section concerning the variation of lift coefficient changing the stagger, will be presented with the data at Reynolds number 60,000 . At the end of the section, the percentage of variation of lift coefficient when changing the stagger will be presented for both Reynolds numbers.

Figure 4.1 and Figure 4.2 plot the lift coefficient, $C_{L}$, vs. $\alpha$ for models having gap 0.5 c at two different stagger conditions and Reynolds number 60,000. Figure 4.1 plots the lift coefficient for the models having positive stagger from $\mathrm{s}=0 \mathrm{c}$ to $\mathrm{s}=1 \mathrm{c}$, Figure 4.2 plots the lift coefficient the models having negative stagger from $s=-1 c$ to $s=0 c$. In both graphs, the curve of lift coefficient as function of angle of attack (lift curve) for the model having no stagger was plotted.


Figure 4.1: Comparison of Lift Curve with varying stagger in the positive direction and constant gap of 0.5 c at $\mathrm{Re}=60,000$


Figure 4.2: Comparison of Lift Curve with varying stagger in the negative direction and constant gap of 0.5 c at $\mathrm{Re}=60,000$

It can be seen that the lift coefficient increases when the stagger increases. Models with positive stagger do not show stall in the range of angle of attack
considered. They show a non linearity of the curve $C_{L}$ vs. $\alpha$ between $10^{\circ}$ and $15^{\circ}$ angle of attack. To compare the different models having gap $\mathrm{g}=0.5 \mathrm{c}$ and different stagger at the two different Reynolds numbers, the lift slope can be considered. Values of lift slope in the range of angle of attack between $-2^{\circ}$ and $8^{\circ}$ are shown in Table 4.1 and in Table 4.2 where the comparison between the different models is expressed in terms of a percent variation which is referenced to the model having no stagger.

Table 4.1: Lift Slope for models having $g=0.5 \mathrm{c}$ at Re 60,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | C La (-) |  |
|  | $-1 c$ | 60,000 | 0.0613 | 11.45 |
|  | -0.5 c | 60,000 | 0.0572 | 4.01 |
|  | 0 c | 60,000 | 0.0550 | 0.00 |
|  | 0.5 c | 60,000 | 0.0634 | 15.31 |
|  | 1 c | 60,000 | 0.0701 | 27.51 |

Table 4.2: Lift Slope for models having $\mathrm{g}=0.5 \mathrm{c}$ at Re 120,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | C $_{\text {La }}(-)$ |  |
|  | -1 c | 120,000 | 0.0598 | 8.22 |
|  | -0.5 c | 120,000 | 0.0569 | 2.89 |
|  | 0 c | 120,000 | 0.0553 | 0.00 |
|  | 0.5 c | 120,000 | 0.0625 | 13.04 |
|  | 1 c | 120,000 | 0.0693 | 25.34 |

It can be seen from the graphs in Figure 4.1 and in Figure 4.2 of lift coefficient vs. angle of attack at $\operatorname{Re} 60,000$ that the lift slope increases in both the positive and negative stagger direction. From Table 4.1 it can be seen that at $\operatorname{Re}$ 60,000 the lift slope increases by almost $5 \%$ for each 0.5 c stagger increment in the
negative direction and it increases by almost $15 \%$ for each 0.5 c stagger increment in the positive direction. Instead, from Table 4.2 at Re 120,000 the percentages of variation between different staggers is slightly different but the trend is the same.

Figure 4.3 shows the lift slope variation vs. stagger for the different models at different Reynolds number, where, again, the percentage is referenced to the model having no stagger.


Figure 4.3: Shows the Lift Slope Variation with varying stagger and constant gap of 0.5c at both Reynolds numbers of 60,000 and 120,000

## Aerodynamic Efficiency

Figure 4.4 and Figure 4.5 plot the aerodynamic efficiency, or lift over drag ratio, $L / D$, vs. $\alpha$ for the models having gap $g=0.5 \mathrm{c}$ at two different stagger conditions and Reynolds number 60,000. As before, Figure 4.4 plots the aerodynamic efficiency for the models having positive stagger from $s=0 c$ to $s=1 c$,

Figure 4.5 plots the aerodynamic efficiency for the models having negative stagger from $s=-1 c$ to $s=0 c$. In both graphs, the curve of aerodynamic efficiency as a function of angle of attack for the model having no stagger was plotted.


Angle of Attack, a $\left({ }^{\circ}\right)$
Figure 4.4: Comparison of Lift over Drag Ratio with varying stagger in the positive direction and constant gap of 0.5 c at $\mathrm{Re}=60,000$


Figure 4.5: Comparison of Lift over Drag Ratio with varying stagger in the negative direction and constant gap of 0.5 c at $\operatorname{Re}=60,000$

It can be seen from the graphs of aerodynamic efficiency vs. angle of attack at Re 60,000 that models with larger stagger experience higher lift over drag ratio. Due to noise in the data around the value of maximum aerodynamic efficiency, the comparison between different models having same gap will be conducted considering the area underneath the curve of L/D vs. angle of attack. This procedure will allow to obtain values of area and the model with highest aerodynamic efficiency will be the model having highest area underneath the same curve keeping the angle of attack range constant across comparisons. Figure 4.6 shows an example of area underneath the curve of $L / D$ vs. angle of attack that has been considered.


Figure 4.6: Shows an example of area underneath the curve L/D

This figure and the subsequent calculations were performed using the software Matlab and the calculation were carried out using the function $\operatorname{trapz}(x, y)$. This function computes an approximation of the integral of the function via the trapezoid rule. The error, if present, is declared to be equal to $1 / 1000$ of the step size of the function. Because the step size in this case is equal to the increment of angle of attack which is equal to $0.25^{\circ}$, the error using this function will be equal to $0.025 \%$. Also, in this section the comparison between areas, which means the largest total aerodynamic efficiency, of each model was performed by taking the percentage of variation referenced to the model with no stagger. This means that each calculation will be affected by the same error for each model and this error will not affect the final comparison between different models. Table 4.3 shows the variation of aerodynamic efficiency between different models at two different Reynolds numbers.

Table 4.3: Integrated Lift over Drag Ratio for modells having $g=0.5 \mathrm{c}$ at $\operatorname{Re} 60,000$

| Model |  | Speed | Integrated Lift <br> over Drag Ratio | Speed | Integrated Lift <br> over Drag Ratio |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | \% Variation | Re (-) | \% Variation |
| 0.5 | -1 | 60,000 | 14.90 | 120,000 | 8.24 |
|  | -0.5 | 60,000 | 11.44 | 120,000 | 5.45 |
|  | 0 | 60,000 | 0.00 | 120,000 | 0.00 |
|  | 0.5 | 60,000 | 10.24 | 120,000 | 28.07 |
|  | 1 | 60,000 | 21.18 | 120,000 | 34.48 |

From Table 4.3 it can be seen that at $\operatorname{Re} 60,000$ the integrated lift over drag ratio increases by almost $10 \%$ for the first 0.5 c stagger increment in the negative direction then it increases of another $5 \%$ for the second 0.5 c stagger increment in the negative direction. For the positive stagger increment it can be seen that the
maximum lift over drag ratio increases by almost $10 \%$ for each 0.5 c stagger increment in the positive direction. At $\operatorname{Re} 120,000$ the percentages of variation between different staggers is slightly different but the trend is the same. As has been done before for the lift slope, the variation of aerodynamic efficiency at different Reynolds number as function of stagger, in Figure 4.7 is plotted.


Figure 4.7: Shows the Integrated Lift over Drag Variation with varying stagger and constant gap of 0.5 c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and $\mathbf{1 2 0 , 0 0 0}$

Models with positive stagger at $\operatorname{Re} 120,000$ show large variation in aerodynamic efficiency when the stagger changes from $s=0 c$ to $s=1 c$. This variation at $\operatorname{Re} 120,000$ is lower when considering negative stagger configurations. At $\operatorname{Re} 60,000$ the models show almost the same variation in both negative and positive stagger configuration.

## Drag Polar

Figure 4.8 and Figure 4.9 plot the drag polar, $C_{D}$ vs. $C_{L}$, for the models with gap $g=0.5 \mathrm{c}$ at two different stagger conditions and Reynolds number 60,000.


Lift Coefficient, $\mathrm{C}_{\mathrm{L}}(-)$
Figure 4.8: Comparison of Drag Polar with varying stagger in the positive direction and constant gap of 0.5 c at $\operatorname{Re}=60,000$


Figure 4.9: Comparison of Drag Polar with varying stagger in the negative direction and constant gap of 0.5 c at $\mathrm{Re}=60,000$

As before, Figure 4.8 plots the drag polar for the models having positive stagger from $s=0 c$ to $s=1 c$, Figure 4.9 plots the drag polar for the models with negative stagger from $s=-1 c$ to $s=0 c$. In both graphs, the drag polar for the model having no stagger was plotted. From Figure 4.8 it can be seen that models with positive stagger have lower $\mathrm{C}_{\mathrm{Di}}$ compared to the models with negative stagger in Figure 4.9. This means that positive stagger models perform better than the models with negative stagger.

Table 4.4 and Table 4.5 compare the induced and parasite drag between all configurations at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$.

Table 4.4: Induced Drag over Total Drag Ratio for models having $g=0.5 \mathrm{c}$ at $\operatorname{Re} 60,000$

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
|  | -1 c | 60,000 | 0.051 | 0.019 | 0.031 | 61.69 |
|  | -0.5 c | 60,000 | 0.040 | 0.019 | 0.021 | 53.04 |
|  | 0 c | 60,000 | 0.041 | 0.022 | 0.019 | 46.03 |
|  | 0.5 c | 60,000 | 0.052 | 0.024 | 0.028 | 54.01 |
|  | 1 c | 60,000 | 0.042 | 0.018 | 0.024 | 57.38 |

Table 4.5: Induced Drag over Total Drag Ratio for models having g=0.5c at Re 120,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
|  | -1 c | 120,000 | 0.041 | 0.021 | 0.020 | 48.74 |
|  | -0.5 c | 120,000 | 0.033 | 0.022 | 0.012 | 35.58 |
|  | 0 c | 120,000 | 0.034 | 0.019 | 0.014 | 42.38 |
|  | 0.5 c | 120,000 | 0.031 | 0.021 | 0.010 | 32.33 |
|  | 1 c | 120,000 | 0.043 | 0.018 | 0.025 | 57.67 |

Figure 4.10 plots the induced drag percentage of total drag for both Reynolds numbers versus stagger.


Figure 4.10: Shows the Induced Drag over Total Drag Ratio with varying stagger and constant gap of 0.5 c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and 120,000

It can be seen that at $\operatorname{Re} 60,000$ the unstaggered model has the least amount of induced drag percentage of the total drag. At Re 120,000 the model with stagger $s=0.5 \mathrm{c}$ is the one with the least amount of induced drag percentage of the total drag.

### 4.2.2 Gap 1c

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on models with similar gap, $g=1 \mathrm{c}$, at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s})$, $\operatorname{Re}=120,000(\approx 18 \mathrm{~m} / \mathrm{s})$.

## Lift Coefficient

The Reynolds number does not affect the variation of lift coefficient for the different models. For this reason, the section concerning the variation of lift coefficient with change in stagger, will be presented with the data at Reynolds
number 60,000 . At the end of the section, the percent variation in lift coefficient with changing stagger will be presented for both Reynolds numbers.

Figure 4.11 and Figure 4.12 plot the lift coefficient, $C_{L}$, vs. $\alpha$ for models with gap $\mathrm{g}=1 \mathrm{c}$ at two different stagger conditions and Reynolds number 60,000. Figure 4.11 plots the lift coefficients for the models with positive stagger from $\mathrm{s}=$ $0 c$ to $s=1.5 \mathrm{c}$, Figure 4.12 plots the lift coefficient for the models with negative stagger from $s=-1.5 \mathrm{c}$ to $\mathrm{s}=0 \mathrm{c}$. In both graphs, the curve of lift coefficient as a function of angle of attack for the model with no stagger was plotted.


Figure 4.11: Comparison of Lift Curve with varying stagger in the positive direction and constant gap of 1 c at $\mathrm{Re}=60,000$


Figure 4.12: Comparison of Lift Curve with varying stagger in the negative direction and constant gap of 1 c at $\operatorname{Re}=60,000$

It can be seen from Figure 4.11 that the lift coefficient does not change when the stagger increases in the positive direction. Models with positive stagger do not show stall in the range of angle of attack considered. From Figure 4.12 it can be seen that models with negative stagger have different behavior than the models with positive stagger. In the negative stagger configurations each model presents a different lift coefficient vs. angle of attack curve. But nothing can be said about the models with gap $\mathrm{g}=1 \mathrm{c}$ because the trend is not clear as it was in the case of gap $g=0.5 c$. This could be due to the fact that higher gap means less interference between the two wings. This yields essentially the same lift coefficient vs. angle of attack behavior for all models having similar gap $g=1 \mathrm{c}$. Now compare the lift slope of the different models with gap $g=1 c$ and different stagger at the two different Reynolds numbers. Values of lift slope in the range of angle of attack
between $-2^{\circ}$ and $8^{\circ}$ are shown in Table 4.6 and Table 4.7 where the comparison between the different models is expressed in terms of a percentage referenced to the model having no stagger.

Table 4.6: Lift Slope for models having $g=1 c$ at Re 60,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | $C_{\mathrm{La}}(-)$ |  |
| 1 F | -1.5 c | 60,000 | 0.0697 | 0.56 |
|  | -1 c | 60,000 | 0.0693 | -0.10 |
|  | -0.5 c | 60,000 | 0.0737 | 6.32 |
|  | 0 c | 60,000 | 0.0693 | 0.00 |
|  | 0.5 c | 60,000 | 0.0748 | 7.83 |
|  | 1 c | 60,000 | 0.0756 | 9.12 |
|  | 1.5 c | 60,000 | 0.0803 | 15.89 |

Table 4.7: Lift Slope for models having $\mathrm{g}=1 \mathrm{c}$ at Re 120,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | C $_{\text {La }}(-)$ |  |
|  | $-1.5 c$ | 120,000 | 0.0725 | 2.36 |
|  | $-1 c$ | 120,000 | 0.0690 | -2.59 |
|  | -0.5 c | 120,000 | 0.0712 | 0.58 |
|  | 0 c | 120,000 | 0.0708 | 0.00 |
|  | 0.5 c | 120,000 | 0.0719 | 1.59 |
|  | 1 c | 120,000 | 0.0746 | 5.33 |
|  | 1.5 c | 120,000 | 0.0803 | 13.49 |

Figure 4.13 shows the lift slope variation vs. stagger for the different models at different Reynolds number, where the percentage is referenced to the model having no stagger.


Figure 4.13: Shows the Lift Slope Variation with varying stagger and constant gap of 1c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and 120,000

Looking at Figure 4.13 it can be seen that the lift slope of the models changes following a descendent trend going from stagger $s=-1.5 c$ to $s=-1 c$ at both $\operatorname{Re}$ 60,000 and $\operatorname{Re} 120,000$. After a value of stagger equal to $s=-1 c$, the lift slope starts to increase until it reaches its maximum value at the highest stagger configuration. Also, the only model that has lower lift slope compared to the model with no stagger is the model with stagger $s=-0.5$.

## Aerodynamic Efficiency

Figure 4.14 and Figure 4.15 plot the aerodynamic efficiency, or lift over drag ratio, $L / D$, vs. $\alpha$ for the models having gap $\mathrm{g}=1 \mathrm{c}$ at two different stagger conditions and Reynolds number 60,000.


Angle of Attack, $\alpha\left({ }^{\circ}\right)$
Figure 4.14: Comparison of Lift over Drag Ratio with varying stagger in the positive direction and constant gap of 1 c at $\mathrm{Re}=60,000$


Angle of Attack, $\alpha$ ( ${ }^{\circ}$ )
Figure 4.15: Comparison of Lift over Drag Ratio with varying stagger in the negative direction and constant gap of 1c at $\operatorname{Re}=\mathbf{6 0 , 0 0 0}$

As before, Figure 4.14 plots the aerodynamic efficiency for the models having positive stagger from $s=0 c$ to $s=1.5 \mathrm{c}$, Figure 4.15 plots the aerodynamic efficiency for the models having negative stagger from $s=-1.5 c$ to $s=0 c$. In both graphs, the curve of aerodynamic efficiency as a function of angle of attack for the model having no stagger was plotted.

To compare the different models in term of integrated efficiency the same method of computing the area beneath the curve of lift to drag ratio vs. angle of attack was used. This method, explained in the section concerning the models having similar gap $g=0.5 \mathrm{c}$, gives, for these models with similar gap $g=1 \mathrm{c}$, the following results represented in Table 4.8 where the variation of aerodynamic efficiency between different models is referenced to the model having no stagger at the two different Reynolds numbers.

Table 4.8: Integrated Lift over Drag Ratio for models having g=1c at $\operatorname{Re} 60,000$

| Model |  | Speed | Integrated Lift <br> over Drag Ratio | Speed | Integrated Lift <br> over Drag Ratio |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | L/D $_{\max }(-)$ | $\operatorname{Re}(-)$ | L/D $_{\max }(-)$ |
| 18 | $-1.5 c$ | 60,000 | 10.54 | 120,000 | -13.43 |
|  | $-1 c$ | 60,000 | 14.13 | 120,000 | -12.98 |
|  | $-0.5 c$ | 60,000 | 3.18 | 120,000 | -18.22 |
|  | $0 c$ | 60,000 | 0.00 | 120,000 | 0.00 |
|  | $0.5 c$ | 60,000 | -6.80 | 120,000 | -11.82 |
|  | $1 c$ | 60,000 | 13.83 | 120,000 | -0.48 |
|  | $1.5 c$ | 60,000 | 11.70 | 120,000 | -2.46 |

As before for the lift slope, the variation of aerodynamic efficiency at different Reynolds number as function of stagger is plotted in Figure 4.16.


Figure 4.16: Shows the Integrated Lift over Drag Variation with varying stagger and constant gap of 1c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and $\mathbf{1 2 0 , 0 0 0}$

From Table 4.8 and Figure 4.16 , it can be seen that at Re 60,000 the stagger has a positive effect on the aerodynamic efficiency, except in the case of stagger $\mathrm{s}=$ 0.5 c where the integrated lift over drag ratio decreases by almost $7 \%$ compared to the unstaggered model. Also, models with stagger $s= \pm 1 \mathrm{c}$ have higher lift over drag ratio than the models with stagger $s= \pm 1.5 \mathrm{c}$. It can also be seen that positive staggers $s=1 c$, and $s=1.5 c$ have the same aerodynamic efficiency as the negative cases $s=-1 c$, and $s=-1.5$ respectively. This means that at low Reynolds number the stagger could be positive or negative and it would affect the integrated aerodynamic efficiency in the same way. At Re 120,000 the stagger always has a negative effect on the aerodynamic efficiency. In this case, negative stagger has a
big deleterious effect on integrated aerodynamic efficiency compared to the positive stagger $s=1 c$ and $s=1.5 c$ where the effect of stagger is almost negligible. The model with stagger $s=0.5 \mathrm{c}$ does not follow this trend. But this model was also the only exception in the case of $\operatorname{Re} 60,000$.

## Drag Polar

Figure 4.17 and Figure 4.18 plot the drag polar, $C_{D}$ vs. $C_{L}$, for the models with gap $g=1 \mathrm{c}$ at two different stagger conditions and Reynolds number 60,000. As before, Figure 4.17 plots the drag polar for the models having positive stagger from $s=0 c$ to $s=1.5 c$, Figure 4.18 plots the drag polar for the models with negative stagger from $s=-1.5 c$ to $s=0 c$. In both graphs, the drag polar for the model with no stagger was plotted as a baseline reference.


Figure 4.17: Comparison of Drag Polar with varying stagger in the positive direction and constant gap of 1 c at $\mathrm{Re}=60,000$


Figure 4.18: Comparison of Drag Polar with varying stagger in the negative direction and constant gap of 1 c at $\operatorname{Re}=60,000$

Also in the case of the drag polar, positive stagger models follow the same trend. Instead, from Figure 4.18 it can be seen that models with negative stagger have different behavior than the models with positive stagger. In the negative
stagger configurations each model presents a different drag polar. But nothing can be said about the models with gap $g=1 \mathrm{c}$ because the trend is not clear as it was in the case of gap $g=0.5 \mathrm{c}$. And the reason is the same explained for the lift coefficient vs. angle of attack. The two wings interact more in the case of negative stagger than in the case of positive stagger. But still, high values of gap, which in this case means $g=1 c$, make the drag polar to be almost the same when changing the stagger. Table 4.9 and

Table 4.10 show the comparison between the induced and parasite drag of all the models at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$.

Table 4.9: Induced Drag over Total Drag Ratio for models having g=1c at Re 60,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{\text {d }}(-)$ | $\mathrm{C}_{\text {D } 0}(-)$ | $\mathrm{C}_{\text {Di }}(-)$ | $\% \mathrm{C}_{\mathrm{D}_{i} / C_{D}(-)}$ |
| 1c | -1.5c | 60,000 | 0.029 | 0.019 | 0.010 | 35.19 |
|  | -1c | 60,000 | 0.039 | 0.019 | 0.020 | 51.37 |
|  | -0.5c | 60,000 | 0.066 | 0.027 | 0.039 | 59.63 |
|  | 0c | 60,000 | 0.045 | 0.024 | 0.021 | 47.13 |
|  | 0.5c | 60,000 | 0.091 | 0.041 | 0.050 | 55.37 |
|  | 1c | 60,000 | 0.052 | 0.020 | 0.032 | 61.79 |
|  | 1.5 c | 60,000 | 0.053 | 0.022 | 0.030 | 57.87 |

Table 4.10: Induced Drag over Total Drag Ratio for models having g=1c at Re 120,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{i}} / C_{D}(-)$ |
|  | -1.5 c | 120,000 | 0.035 | 0.022 | 0.013 | 37.75 |
|  | -1 c | 120,000 | 0.039 | 0.018 | 0.021 | 53.68 |
|  | -0.5 c | 120,000 | 0.041 | 0.022 | 0.019 | 46.44 |
|  | 0 c | 120,000 | 0.028 | 0.013 | 0.015 | 54.34 |
|  | 0.5 c | 120,000 | 0.058 | 0.029 | 0.029 | 50.48 |
|  | 1 c | 120,000 | 0.029 | 0.019 | 0.009 | 33.18 |
|  | 1.5 c | 120,000 | 0.036 | 0.021 | 0.016 | 43.15 |

The induced drag percentage of total drag ratio for both Reynolds numbers versus stagger are plotted in Figure 4.19 providing:


Figure 4.19: Shows the Induced Drag over Total Drag Ratio with varying stagger and constant gap of 1c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and $\mathbf{1 2 0 , 0 0 0}$

It can be seen from Table 4.9 and Figure 4.19, that at Re 60,000 the model with stagger -1.5 c has the least amount of induced drag percentage of the total drag. At $\operatorname{Re} 120,000$ the model with stagger 1c has the least induced drag percentage of total drag.

### 4.2.3 Gap 2c

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on the models with similar gap, $g=2 c$, at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s})$, $\operatorname{Re}=120,000(\approx 18 \mathrm{~m} / \mathrm{s})$.

## Lift Coefficient

The Reynolds number does not affect the variation of lift coefficient for the different models. For this reason, the section concerning the variation of lift coefficient by varying the stagger, will be presented with the data at Reynolds number 60,000 . At the end of the section, the percentage of variation in lift coefficient changing the stagger will be presented for both Reynolds numbers.

Figure 4.20 plots the lift coefficient, $C_{L}$, vs. $\alpha$ for the models with gap g = 2c at Reynolds number 60,000. The two different models with gap $g=2 c$ have stagger $s= \pm 1 c$. There were just two models with gap $g=2 c$. This will result in a more superficial comparison between different models, but these results can potentially verify the considerations about the effect of stagger on models having similar gap.


Figure 4.20: Comparison of Lift Curve with varying stagger and constant gap of 2 c at $\mathrm{Re}=60,000$

It can be seen that high values of gap means little interference between the two wings. The difference between positive stagger, $s=1 c$, and negative stagger, $s=-1 c$, is negligible before an angle of attack equal of $10^{\circ}$. The two models also have the same lift slope, which is represented in Table 4.11 and Table 4.12 at both Reynolds numbers, and it is referenced to the model having negative stagger.

Table 4.11: Lift Slope for models having $g=1 \mathrm{c}$ at $\operatorname{Re} 60,000$

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | CLa $_{\text {( }}$ ) |  |
|  | -1c | 60,000 | 0.0881 | 0.00 |
|  | 1c | 60,000 | 0.0873 | -0.89 |

Table 4.12: Lift Slope for models having $\mathrm{g}=1 \mathrm{c}$ at Re 120,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{\text {La }}(-)$ |  |
| 2 c | -1 c | 120,000 | 0.0875 | 0.00 |
|  | 1 c | 120,000 | 0.0874 | -0.05 |

The variation of lift slope between the two models is almost negligible.

## Aerodynamic Efficiency

Figure 4.21 plots the aerodynamic efficiency, or lift over drag ratio, $L / D$, vs. $\alpha$ for the models with gap 2c at two different stagger conditions and Reynolds number 60,000.


Angle of Attack, $a\left({ }^{\circ}\right)$
Figure 4.21: Comparison of Lift over Drag Ratio with varying stagger and constant gap of 2 c at $\mathrm{Re}=60,000$

The aerodynamic efficiency does not vary too much between the two models, except in the region of angle of attack between $2^{\circ}$ and $6^{\circ}$ where the aerodynamic efficiency reaches the maximum value. Table 4.13 shows the variation of aerodynamic efficiency for the two models referred to the model having negative stagger at both Reynolds numbers.

Table 4.13: Integrated Lift over Drag Ratio for models having $\mathbf{g = 0 . 5 c}$ at $\operatorname{Re} \mathbf{6 0 , 0 0 0}$

| Model |  | Speed | Integrated Lift <br> over Drag Ratio | Speed | Integrated Lift <br> over Drag Ratio |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | L/D $\max ^{(-)}$ | $\operatorname{Re}(-)$ | L/D $_{\max }(-)$ |
| 2c | -1c | 60,000 | 0.00 | 120,000 | 0.00 |
|  | 1c | 60,000 | -12.12 | 120,000 | -8.92 |

Negative stagger $s=-1 c$ has a higher aerodynamic efficiency than positive stagger $s=1 \mathrm{c}$ at both Reynolds numbers. It is interesting, for this configuration with a gap of $g=2 c$, to take a look at the lift over drag ratio as function of angle of attack in the case of $\operatorname{Re} 120,000$. Figure 4.22 plots the aerodynamic efficiency, or lift over drag ratio, $L / D$, vs. $\alpha$ for the models with gap $g=2 c$ at two different stagger conditions and a Reynolds number of 120,000.


Figure 4.22: Comparison of Lift over Drag Ratio with varying stagger and constant gap of 2 c at $\mathrm{Re}=120,000$

In Figure 4.22, the two models show a interesting behavior beyond the angle of attack equal to $12^{\circ}$. It can be seen that this is the angle of stall for both models. This behavior will also be seen in the drag polar in the next section. This phenomenon is called stall flutter and it will be briefly explained in Appendix F.

## Drag Polar

Figure 4.23 and Figure 4.24 plot the drag polar, $C_{D}$ vs. $C_{L}$, for models with gap $g=2 c$ at two different Reynolds numbers 60,000 and 120,000.


Figure 4.23: Comparison of Drag Polar with varying stagger and constant gap of 2 c at $\mathrm{Re}=60,000$


Figure 4.24: Comparison of Drag Polar with varying stagger and constant gap of 2 c at $\mathrm{Re}=120,000$

It can be seen from Figure 4.23 and Figure 4.24, that the drag polar is almost the same for both models at two different Reynolds numbers until a lift coefficient $C_{L} \cong 0.8$ then $C_{D i}$ becomes lower for the positive stagger case, meaning that it performs better than the negative stagger case. Table 4.14 and Table 4.15 show the comparison between the induced and parasite drag of all the models at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$.

Table 4.14: Induced Drag over Total Drag Ratio for models having $g=2 \mathrm{c}$ at $\operatorname{Re} 60,000$

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
|  | -1 c | 60,000 | 0.036 | 0.020 | 0.015 | 42.44 |
|  | 1 c | 60,000 | 0.060 | 0.026 | 0.034 | 56.33 |

Table 4.15: Induced Drag over Total Drag Ratio for models having $\mathbf{g}=\mathbf{2 c}$ at Re 120,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
|  | -1 c | 120,000 | 0.031 | 0.020 | 0.011 | 36.44 |
|  | 1 C | 120,000 | 0.050 | 0.026 | 0.023 | 46.90 |

Plotting the ratio of induced drag over total drag for both Reynolds numbers versus stagger in Figure 4.25 we get:


Figure 4.25: Shows the Induced Drag over Total Drag Ratio with varying stagger and constant gap of $2 c$ at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and $\mathbf{1 2 0 , 0 0 0}$

It can be seen that the ratio induced drag over total drag of the positive stagger model is higher than the one with negative stagger at both Reynolds numbers.

### 4.3 Wind Tunnel Balance Data - Similar Stagger Models

In this section the effect of gap on the performance of the biplane joined at the tips will be presented. For this analysis, data from models having similar stagger will be presented to better understand the behaviour of such changes in gap. The aerodynamic performance analyzed will essentially be the lift coefficient, the drag polar, and the aerodynamic efficiency of the different models. In Table 2.2 it can be seen that there are five different stagger configurations, from $s=-1.5 \mathrm{c}$ to $s=1.5 \mathrm{c}$. In this section the effect of gap will be analyzed considering only two different stagger configurations: models having stagger 0 c and the models having
stagger -1c. At the end of each section, results referenced to all the other configurations will be presented.

### 4.3.1 Stagger 0c

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on the models having similar stagger, $s=0 c$, at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s}), \operatorname{Re}=120,000(\approx 18 \mathrm{~m} / \mathrm{s})$.

## Lift Coefficient

Figure 4.26 plots the lift coefficient, $C_{L}$ vs. $\alpha$ for models having stagger 0 c at Reynolds number 60,000.


Figure 4.26: Comparison of Lift Curve with varying gap and constant stagger of Oc at $\mathrm{Re}=60,000$

It can be seen that the lift coefficient increases mostly linearly until an angle of attack around $12^{\circ}$. The model with higher gap shows typical stall behavior, instead the model with lower gap does not show a clear stall break in the range of angle of attack considered. This means that low gap is a favourable condition for obtaining no stall. When the gap is higher, the interference between the two wings is less important and this yields a more common stall behavior. Taking a look now at the lift slope it can be seen that it increases when the gap increases. Values of lift slope in the range of angle of attack between $-2^{\circ}$ and $8^{\circ}$ are shown in Table 4.16 and Table 4.17 at two Reynolds numbers 60,000 and 120,000 . The comparison between the different models is expressed in terms of a percentage which is referenced to the model having a gap of 0.5c.

Table 4.16: Lift Slope for models having $s=0 \mathrm{c}$ at $\operatorname{Re} 60,000$

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | C $_{\text {La }}(-)$ |  |
| 0.5 c | 0 c | 60,000 | 0.0550 | 0.00 |
|  |  | 60,000 | 0.0693 | 26.06 |

Table 4.17: Lift Slope for models having $\mathrm{g}=1 \mathrm{c}$ at $\operatorname{Re} 120,000$

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | C $_{\text {La }}(-)$ |  |
| 0.5 c | 0 c | 120,000 | 0.0553 | 0.00 |
|  |  | 0.0708 | 28.07 |  |

The lift slope of the model having gap $\mathrm{g}=1 \mathrm{c}$ is always higher than the lift slope of the model having gap $g=0.5$ at both Reynolds numbers. The variation of lift slope is almost the same in both cases.

## Aerodynamic Efficiency

Figure 4.27 and Figure 4.28 plot the aerodynamic efficiency, or lift over drag ratio vs. $\alpha$ for models with stagger 0 c at Reynolds number 60,000 and $\operatorname{Re} 120,000$.


Angle of Attack, $\boldsymbol{\alpha}\left({ }^{\circ}\right)$
Figure 4.27: Comparison of Lift over Drag Ratio with varying gap and constant stagger of $0 c$ at $\operatorname{Re}=60,000$


Figure 4.28: Comparison of Lift over Drag Ratio with varying gap and constant stagger of $0 c$ at $R e=120,000$

It can be seen from the graphs of aerodynamic efficiency vs. angle of attack at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$ that models with larger gap experience higher lift over drag ratio. At higher Reynolds number this effect becomes important and it can be seen in Table 4.18 where the variation of aerodynamic efficiency is referenced to the model with stagger $s=0 \mathrm{c}$ :

Table 4.18: Integrated Lift over Drag Ratio for models having $s=0 \mathrm{c}$ at $\operatorname{Re} 120,000$

| Model |  | Speed | Integrated Lift <br> over Drag Ratio | Speed | Integrated Lift <br> over Drag Ratio |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | \% Variation | Re (-) | \% Variation |
| 0.5c | $0 c$ | 60,000 | 0.00 | 120,000 | 0.00 |
| 1c |  | 60,000 | 1.48 | 120,000 | 34.92 |

At $\operatorname{Re} 120,000$ the integrated lift over drag ratio of the model having gap $\mathrm{g}=1 \mathrm{c}$ is almost $35 \%$ higher than the integrated lift over drag ratio of the model with gap $g=0.5 \mathrm{c}$. This variation at $\operatorname{Re} 60,000$ is equal to just $1.5 \%$. This means that in terms of aerodynamic efficiency, the effect of changing the gap at higher Reynolds numbers is more important than the same effect at lower Reynolds numbers.

## Drag Polar

Figure 4.29 and Figure 4.30 plot the drag polar, $C_{D}$ vs. $C_{L}$, for the models having stagger $\mathrm{s}=0 \mathrm{c}$ at two different Reynolds number 60,000 and 120,000.


Figure 4.29: Comparison of Drag Polar with varying gap and constant stagger of 0 C at $\mathrm{Re}=\mathbf{6 0 , 0 0 0}$


Figure 4.30: Comparison of Drag Polar with varying gap and constant stagger of $0 c$ at $R e=120,000$

The drag polar is different at the two Reynolds numbers. This difference can be evaluated in terms of induced drag over total drag ratio. The values of this ratio are reported in Table 4.19 and Table 4.20 at two different Reynolds numbers.

Table 4.19: Induced Drag over Total Drag Ratio for models having s=0c at Re 60,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
| 0.5 c | $0 c$ | 60,000 | 0.041 | 0.022 | 0.019 | 46.03 |
|  |  | 60,000 | 0.045 | 0.024 | 0.021 | 47.13 |

Table 4.20: Induced Drag over Total Drag Ratio for models having s=0c at Re 120,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
| 0.5 c | 0 c | 120,000 | 0.034 | 0.019 | 0.014 | 42.38 |
|  |  | 120,000 | 0.029 | 0.013 | 0.016 | 56.25 |

It can be seen from Figure 4.29 and Figure 4.30, and from Table 4.19 and Table 4.20 that at Re 60,000 the gap has marginal effect on the induced drag. But at Re 120,000 it can be seen that model with gap $g=0.5$ experiences less induced drag than the model with gap $\mathrm{g}=1 \mathrm{c}$.

### 4.3.2 Stagger -1c

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on models with similar stagger, $s=-1 c$, at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s})$, $\operatorname{Re}=120,000(\approx 18 \mathrm{~m} / \mathrm{s})$.

## Lift Coefficient

Figure 4.31 plots the lift coefficient, $C_{L}$, vs. $\alpha$ for models with stagger $\mathrm{s}=-1 \mathrm{c}$ at Reynolds number 60,000.


Figure 4.31: Comparison of Lift Curve with varying gap and constant stagger of -1 c at $\mathrm{Re}=60,000$

It can be seen that the lift coefficient increases linearly until an angle of attack around $8^{\circ}$. The model with lower gap shows more typical stall behavior, instead the models with higher gap do not show stall in the range of angle of attack considered. Taking a look at the lift slope it can be seen that it increases when the gap increases. Values of lift slope in the range of angle of attack between $-2^{\circ}$ and $8^{\circ}$ are shown in Table 4.21 and Table 4.22 at two Reynolds numbers 60,000 and 120,000 . The comparison between the different models is expressed in terms of a percentage which is referenced to the model with gap $\mathrm{g}=0.5 \mathrm{c}$

Table 4.21: Lift Slope for models having $s=-1 c$ at Re 60,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\boldsymbol{R e}(-)$ | $\mathrm{C}_{\text {La }}(-)$ |  |
| 0.5c | -1c | 60,000 | 0.0613 | 0.00 |
| 1c |  | 60,000 | 0.0693 | 12.99 |
| 2 c |  | 60,000 | 0.0881 | 43.71 |

Table 4.22: Lift Slope for models having s=-1c at Re 120,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | Cla $^{(-)}$ |  |
| 0.5 c | $-1 \mathrm{c}$ | 120,000 | 0.0598 | 0.00 |
|  |  | 0.0690 | 15.27 |  |
| Ic |  | 120,000 | 0.0875 | 46.17 |

Figure 4.32 plots the variation of lift slope vs. gap at both Reynolds numbers.


Figure 4.32: Shows the Lift Slope Variation with varying gap and constant stagger of -1c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and 120,000

The lift slope is increasing by almost $15 \%$ with each 0.5 c gap increment. Higher gap models experience higher lift slope. This variation of lift slope is slightly higher at higher Reynolds numbers.

## Aerodynamic Efficiency

Figure 4.33 and Figure 4.34 plot the aerodynamic efficiency, or lift over drag ratio, $L / D$ vs. $\alpha$ for the models having stagger $\mathrm{s}=-1 \mathrm{c}$ at Reynolds number 60,000 and $\operatorname{Re} 120,000$.


Figure 4.33: Comparison of Lift over Drag Ratio with varying gap
and constant stagger of -1 c at $\mathrm{Re}=\mathbf{6 0 , 0 0 0}$


Angle of Attack, $\pi\left({ }^{\circ}\right)$

Figure 4.34: Comparison of Lift over Drag Ratio with varying gap and constant stagger of $\mathbf{- 1 c}$ at $\operatorname{Re}=\mathbf{1 2 0 , 0 0 0}$

It can be seen from the graphs of aerodynamic efficiency vs. angle of attack at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$ that models with larger gap experience higher lift over drag ratio. In Figure 4.34 it can be seen that the model with stagger $s=-1 \mathrm{c}$ and gap $\mathrm{g}=2 \mathrm{c}$ produces very noisy data. During the test large vibration were observed. It is believed that this may be stall flutter. At higher Reynolds number this effect becomes important and it can be seen in Table 4.23 where the variation of aerodynamic efficiency referenced to the model with stagger $s=-1 c$ is reported:

Table 4.23: Integrated Lift over Drag Ratio for models having $s=-1 c$ at Re 120,000

| Model |  | Speed | Integrated Lift over Drag Ratio | Speed | Integrated Lift over Drag Ratio |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\boldsymbol{R e}(-)$ | \% Variation | $\boldsymbol{R e}(-)$ | \% Variation |
| 0.5c | -1c | 60,000 | 0.00 | 120,000 | 0.00 |
| 1c |  | 60,000 | 0.80 | 120,000 | 8.47 |
| 2c |  | 60,000 | 11.70 | 120,000 | 38.78 |

At $\operatorname{Re} 120,000$ the integrated lift over drag ratio of the model with gap 1c is almost $9 \%$ higher than the integrated lift over drag ratio of the model with gap 0.5 c . The variation of integrated lift over drag ratio becomes equal to almost $40 \%$ when comparing the model with gap 2 c to the model with gap 0.5 c . This variation at Re 60,000 is equal to just $0.8 \%$ for the model with gap 1 c and almost $12 \%$ for the one with gap 2c. This means that in terms of aerodynamic efficiency, the effect of changing the gap at higher Reynolds numbers is more important than the same effect at lower Reynolds numbers. The same result was obtained when considering models with no stagger and different gap configurations. Figure 4.35 plots the variation of integrated aerodynamic efficiency for models with stagger -1c.


Figure 4.35: Shows the Integrated Lift over Drag Variation with varying gap and constant stagger of -1c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and $\mathbf{1 2 0 , 0 0 0}$

## Drag Polar

Figure 4.36 plots the drag polar, $C_{D}$ vs. $C_{L}$, for the models having stagger $s=-1 \mathrm{c}$ at Reynolds number 60,000.


Figure 4.36: Comparison of Drag Polar with varying gap and constant stagger of -1 c at $\mathrm{Re}=60,000$

It can be seen that models with higher gap perform better than the model with lower gap in terms of induced drag. Table 4.24 and Table 4.25 show the induced drag over total drag ratio at both Reynolds numbers for models with stagger $s=-1 c$.

Table 4.24: Induced Drag over Total Drag Ratio for models having s=-1c at $\operatorname{Re} 60,000$

| Models |  | $\begin{aligned} & \hline \text { Speed } \\ & \operatorname{Re}(-) \end{aligned}$ | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger |  | $\mathrm{C}_{\mathrm{D}}(-)$ | $\mathrm{C}_{\text {D0 }}(-)$ | $\mathrm{C}_{\mathrm{Di}}(-)$ | $\% \mathrm{C}_{\mathrm{Di}} / \mathrm{C}_{\mathrm{D}}(-)$ |
| 0.5c | -1c | 60,000 | 0.040 | 0.019 | 0.021 | 53.04 |
| 1 c |  | 60,000 | 0.039 | 0.019 | 0.020 | 51.37 |
| 2 c |  | 60,000 | 0.036 | 0.020 | 0.015 | 42.44 |

Table 4.25: Induced Drag over Total Drag Ratio for models having s=-1c at Re 120,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\boldsymbol{R e}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | $\mathrm{C}_{\text {D0 }}(-)$ | $\mathrm{C}_{\text {Di }}(-)$ | \% $\mathrm{C}_{\mathrm{Di}} / \mathrm{C}_{\mathrm{D}}(-)$ |
| 0.5c | -1c | 120,000 | 0.039 | 0.018 | 0.021 | 53.68 |
| 1c |  | 120,000 | 0.041 | 0.021 | 0.020 | 48.74 |
| 2c |  | 120,000 | 0.031 | 0.020 | 0.011 | 36.44 |

The ratio of induced drag over total drag has been plotted versus gap in Figure 4.37 at both Reynolds numbers.


Figure 4.37: Shows the Induced Drag over Total Drag Ratio with varying gap and constant stagger of -1c at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and 120,000

Changing Reynolds number, the induced drag over total drag ratio is almost the same and it decreases when the gap increases.

### 4.4 Wind Tunnel Balance Data - Positive Versus Negative Stagger

In this section the effect of changes from positive to negative stagger on the performance of the biplane joined at the tips will be presented. For this analysis, data from models having similar gap will be presented to better understand the behaviour of such changes in stagger from a positive to a negative value. The aerodynamic performance analyzed will essentially be the lift coefficient, the drag polar, and the aerodynamic efficiency of the different models. The model with the highest change in aerodynamic performance compared to the one with same gap but opposite stagger direction, should be the model with the lowest gap where the interferences of the two wings are higher, such that the positive stagger configuration should show different behavior than the negative stagger configuration. In this section the effect of changes in stagger will be analyzed considering only two different models having similar gap $\mathrm{g}=0.5 \mathrm{c}$ and two stagger configurations: the models having stagger $s=-1 c$ and the models having stagger $s=1 c$.

### 4.4.1 Gap 0.5c , Stagger $\pm 1 c$

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $25^{\circ}$ in steps of $0.25^{\circ}$ was performed on the models having similar gap, $g=0.5 \mathrm{c}$, and opposite stagger direction, $\mathrm{s}= \pm 1 \mathrm{c}$ at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s}), \operatorname{Re}=120,000(\approx 18 \mathrm{~m} / \mathrm{s})$.

## Lift Coefficient

Figure 4.38 plots the lift coefficient, $C_{L}$, vs. $\alpha$ for the models having stagger $s= \pm 1 \mathrm{c}$ at Reynolds number 60,000.


Figure 4.38: Comparison of Lift Curve positive vs. negative stagger and constant gap of 0.5 c at $\mathrm{Re}=60,000$

It can be seen that the model with positive stagger shows a higher lift coefficient in the entire range of angle of attack. The model with negative stagger shows the typical stall behavior, instead the model with positive stagger does not
show the stall in the range of angle of attack considered, and the lift coefficient increases linearly until an angle of attack around $12^{\circ}$. This means that positive stagger is a favourable condition for obtaining no stall. Taking a look at the lift slope it can be seen that it increases when the gap increases. Table 4.26 and Table 4.27 show the values of lift slope in the range of angle of attack between $-2^{\circ}$ and $8^{\circ}$ at two Reynolds numbers 60,000 and 120,000 . The comparison between the different models is expressed in terms of a percentage which is referenced to the model having negative stagger $s=-1 c$.

Table 4.26: Lift Slope for models having $g=0.5 c, s= \pm 1 \mathrm{c}$ at $\operatorname{Re} 60,000$

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | C $_{\text {La }}(-)$ |  |
| 0.5 c | -1 c | 60,000 | 0.0613 | 0.00 |
|  | 1c | 60,000 | 0.0701 | 14.41 |

Table 4.27: Lift Slope for models having $g=0.5 c, s= \pm 1 \mathrm{c}$ at Re 120,000

| Models |  | Speed | Lift Slope | \% Variation |
| :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | C La $^{(-)}$ |  |
| 0.5 c | -1 c | 120,000 | 0.0598 | 0.00 |
|  | 1 c | 120,000 | 0.0693 | 15.81 |

The lift slope of the model with positive stagger $s=1 c$ is always higher than the lift slope of the model with negative stagger $s=-1 \mathrm{c}$ at both Reynolds numbers, and the variation of lift slope is almost the same in both cases and equal to $\sim 15 \%$.

## Aerodynamic Efficiency

Figure 4.39 plots the aerodynamic efficiency, or lift over drag ratio, $L / D$, vs. $\alpha$ for the models having similar gap $g=0.5 \mathrm{c}$ and two different stagger configurations $s=-1 c$ and $s=1 c$ at Reynolds number 60,000.


Angle of Attack, $\alpha\left({ }^{\circ}\right)$
Figure 4.39: Comparison of Lift over Drag Ratio positive vs. negative stagger and constant gap of 0.5 c at $\operatorname{Re}=60,000$

It can be seen from the graph in Figure 4.33 of aerodynamic efficiency vs. angle of attack at $\operatorname{Re} 60,000$ that the model with positive stagger experiences a higher lift over drag ratio. Table 4.28 shows the variation of aerodynamic efficiency referenced to the model with negative stagger at both Reynolds numbers:

Table 4.28: Integrated Lift over Drag Ratio for models having $g=0.5 c, s= \pm 1 \mathrm{c}$

| Model |  | Speed | Integrated Lift <br> over Drag Ratio | Speed | Integrated Lift <br> over Drag Ratio |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | \% Variation | Re (-) | \% Variation |
| 0.5 c | -1 c | 60,000 | 0.00 | 120,000 | 0.00 |
|  | 1c | 60,000 | 5.47 | 120,000 | 18.32 |

At $\operatorname{Re} 120,000$ the integrated lift over drag ratio of the model having positive stagger is almost $18 \%$ higher than the integrated lift over drag ratio of the model with negative stagger. This variation at $\operatorname{Re} 60,000$ is equal to almost $5.5 \%$. This means that in terms of aerodynamic efficiency, models with positive stagger are
more efficient than models with negative stagger and that this behavior, and that this effect becomes more important at higher Reynolds numbers.

## Drag Polar

Figure 4.40 plots the drag polar, $C_{D}$ vs. $C_{L}$, for models having gap $g=0.5 \mathrm{c}$ and opposite stagger configuration at Reynolds number 60,000.


Figure 4.40: Comparison of Drag Polar positive vs. negative stagger and constant gap of 0.5 c at $\mathrm{Re}=60,000$

The drag polar in Figure 4.40 shows how models with positive stagger have higher drag compared to models with negative stagger. This difference can be evaluated in terms of induced drag over total drag ratio. The values of this ratio are reported in Table 4.29 and Table 4.30 at two different Reynolds numbers.

Table 4.29: Induced Drag over Total Drag Ratio for models having $g=0.5 \mathrm{c}, \mathrm{s}= \pm 1 \mathrm{c}$ at Re 60,000

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
|  | -1 c | 60,000 | 0.040 | 0.019 | 0.021 | 53.04 |
|  | 1 c | 60,000 | 0.042 | 0.018 | 0.024 | 57.38 |

Table 4.30: Induced Drag over Total Drag Ratio for models having $\mathrm{g}=0.5 \mathrm{c}, \mathrm{s}= \pm 1 \mathrm{c}$ at $\operatorname{Re} 120,000$

| Models |  | Speed | Drag Coefficients |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | $C_{D}(-)$ | $C_{D 0}(-)$ | $C_{D i}(-)$ | $\%_{C_{D i}} / C_{D}(-)$ |
|  | -1 c | 120,000 | 0.041 | 0.021 | 0.020 | 48.74 |
|  | 1 c | 120,000 | 0.043 | 0.018 | 0.025 | 57.67 |

It can be seen from Figure 4.40 and from Table 4.29 and Table 4.30 that models with positive stagger always have higher induced drag over total drag ratio. Also, it can be seen that models with positive stagger do not change their induced drag over total drag ratio at higher Reynolds numbers, meaning that the effect of Reynolds number is less important in models with positive stagger. The opposite behavior can be seen for models with negative stagger where this ratio decreases when the Reynolds number increases.

### 4.5 Wind Tunnel Balance Data - Change in Lift Slope

The plot of lift coefficient, $C_{L}$, vs. angle of attack, $a$, in the linear regime (angle of attack between $-2^{\circ}$ and $8^{\circ}$ ) shows a change in the slope of the curve around $5^{\circ}$ angle of attack for all the models tested in the LSWT at both Reynolds numbers. This section will provide a quantitative aspect to this analysis, showing different approaches followed for obtaining the variation in lift slope and how big this variation is compared to parameters such as stagger, gap, and angle of stagger. This section will also try to answer the question as to why this variation occurs and how the variation can be predicted and studied in future work.

## Summary

In the LSWT, an $\alpha$ sweep from $-2^{\circ}$ to $8^{\circ}$ in steps of $0.25^{\circ}$ was performed on the fourteen models at two different speeds: $\operatorname{Re}=60,000(\approx 9 \mathrm{~m} / \mathrm{s}), \operatorname{Re}=120,000$ $(\approx 18 \mathrm{~m} / \mathrm{s})$. Plotting lift coefficient as a function of angle of attack, a change in lift slope can be observed. The same change in lift slope was observed in previous experiments conducted on biplanes. In 1929, M. Knight and R. Noyes published a report on tests conducted on a series of biplanes, showing the effect of changes in stagger, and gap [12] [13] [14]. Using their data at Re 150,000 the change in lift slope observed was equal to almost $9 \%$ and it occurs around $0^{\circ}$ angle of attack.

A more recent work, conducted in 2007 by D. N. Killian [20] on a biplane joined at the tips by endplates, shows a change in lift slope at $\operatorname{Re} 60,000$ of almost $25 \%$ and it occurs around $2^{\circ}$ angle of attack.

In this section of the paper, the change in the lift slope will be quantified and an hypothesis will be elaborated, using the LSWT data. There are two methods used to compute the lift slope. The graphs will be shown in this paragraph and they will allow the reader to better understand the process by which obtaining the lift slope is made possible. The two methods have been applied to all fourteen models but in this paper the model that shows the greatest change in slope will be the only one considered. This model is the model with gap $\mathrm{g}=0.5 \mathrm{c}$ and stagger $\mathrm{s}=1 \mathrm{c}$ at Re 60,000 . Tables at the end of the paragraph will present the results for all fourteen models.

## First Method

The first method consists of plotting the entire curve of $C_{L}$ vs. $\alpha$ in a range of angle of attack between $-2^{\circ}$ and $8^{\circ}$. Then the slope has been computed across the entire range of angle of attack. Figure 4.41 shows $C_{L}$ vs. $a$ in the range of angle of attack between $-2^{\circ}$ and $8^{\circ}$ for the model having gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=1 \mathrm{c}$ at Re 60,000 and it also shows the linear curve of lift coefficient vs. angle of attack with the equation of the curve.


Figure 4.41: Lift Curve and linear curve in the entire region of angle of attack from $-2^{\circ}$ to $8^{\circ}$ for the model with gap 0.5 c and stagger 1 c at $\operatorname{Re} \mathbf{6 0 , 0 0 0}$

Table 4.31 shows the lift slope for the model having gap $\mathrm{g}=0.5 \mathrm{c}$ and stagger $s=1 c$ at both Reynolds numbers 60,000 and 120,000.

Table 4.31: Lift Slope for model having gap $\mathrm{g}=0.5 \mathrm{c}$ and stagger $\mathrm{s}=1 \mathrm{c}$, First Method

| Model |  | Speed | Lift Slope, $C_{L a}(\boldsymbol{-})$ |
| :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(\boldsymbol{-})$ | $\mathbf{- 2}^{\circ}<\mathbf{\alpha}<\mathbf{8}^{\circ}$ |
|  | 1 c | 60,000 | 0.0701 |
|  |  | 120,000 | 0.0693 |

## Second Method

The second method consists of plotting the entire curve of $C_{L}$ vs. $a$ in a range of angle of attack between $-2^{\circ}$ and $8^{\circ}$. From each of these graphs it is possible to recognize a region of transition, where the lift slope changes from one value to another. Eliminating this region of transition, it will be possible to determine a first slope in the region before the transition, and a second slope right after the same
transition. In this method, the slope of the first region will be called first slope and the slope of the second region, second slope.

Figure 4.42 plots lift coefficient vs. angle of attack for the model having gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=1 \mathrm{c}$ at $\operatorname{Re} 60,000$ in the range of angle of attack before the transition, and in the range of angle of attack after the transition. For this model, at $\operatorname{Re} 60,000$ the region of transition occurs between $3.22^{\circ}$ and $4.72^{\circ}$.


Figure 4.42: Lift Curve in the two different regions of angle of attack before and after the transition for the model with gap 0.5c and stagger 1c at Re 60,000

Then the percentage of change with respect to the first slope has been computed and it is reported in Table 4.32 for both Reynolds numbers.

Table 4.32: Lift Slope for model having gap g=0.5c and stagger $s=1 \mathrm{c}$, Second Method

| Model |  | Speed | Lift Slope, $C_{\text {La }}(-)$ | Lift Slope, $C_{L a}(-)$ | \% Variation |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | first slope | second slope |  |
| 0.5 c | 1 c | 60,000 | 0.0640 | 0.0895 | 39.87 |
|  |  | 120,000 | 0.0660 | 0.0796 | 20.57 |

Using the second method it can be seen that at both Reynolds numbers the lift slope increases after the region of transition. This variation at $\operatorname{Re} 60,000$ is equal to almost $40 \%$, instead at $\operatorname{Re} 120,000$ it is equal to almost $20.6 \%$.

## Results

Table 4.33 shows the results of the slopes at $\operatorname{Re} 60,000$ and 120,000 for all fourteen models using the first method.

Table 4.33: Lift Slope all models, First Method, Re 60,000-120,000

| Model |  | Speed | Lift Slope | Speed | Lift Slope |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\boldsymbol{R e}(-)$ | $\mathrm{C}_{\mathrm{La}}(-)$ | $\mathbf{R e}(-)$ | $\mathrm{C}_{\mathrm{La}}(-)$ |
| 1c | 0 c | 60,000 | 0.0708 | 120,000 | 0.0693 |
| 0.5c | 0c | 60,000 | 0.0553 | 120,000 | 0.0550 |
| 0.5c | 0.5c | 60,000 | 0.0625 | 120,000 | 0.0634 |
| 0.5c | 1c | 60,000 | 0.0693 | 120,000 | 0.0701 |
| 1c | 0.5c | 60,000 | 0.0719 | 120,000 | 0.0748 |
| 1 c | 1c | 60,000 | 0.0746 | 120,000 | 0.0756 |
| 2c | 1c | 60,000 | 0.0874 | 120,000 | 0.0873 |
| 1c | 1.5 c | 60,000 | 0.0786 | 120,000 | 0.0803 |
| 0.5c | -0.5c | 60,000 | 0.0569 | 120,000 | 0.0572 |
| 0.5c | -1c | 60,000 | 0.0598 | 120,000 | 0.0613 |
| 1 c | -0.5c | 60,000 | 0.0712 | 120,000 | 0.0737 |
| 1 c | -1c | 60,000 | 0.0690 | 120,000 | 0.0693 |
| 2c | -1c | 60,000 | 0.0875 | 120,000 | 0.0881 |
| 1 c | -1.5c | 60,000 | 0.0725 | 120,000 | 0.0697 |

The lift slope variation between different models having similar gap or similar stagger was discussed in each section, and reported in Figure 4.3, Figure 4.13, and Figure 4.32. Concerning the second method, Table 4.34 and Table 4.35 show the two lift slopes, first and second slopes, before and after the region of transition, at $\operatorname{Re} 60,000$ and 120,000 . The variation in each table represents the
percentage of change in lift slope referenced for each model to the first slope, meaning that positive variation means that the lift slope increases after the transition.

Table 4.34: Lift Slope all models, Second Method, Re 60,000

| Model |  | Speed | Lift Slope $\mathrm{C}_{\text {La }}(-)$ | Lift Slope C ${ }_{\text {La }}(-)$ | \% Variation |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\operatorname{Re}(-)$ | first region | second region |  |
| 1c | Oc | 60,000 | 0.0631 | 0.0844 | 33.76 |
| 0.5c | 0c | 60,000 | 0.0516 | 0.0641 | 24.10 |
| 0.5c | 0.5c | 60,000 | 0.0591 | 0.0724 | 22.43 |
| 0.5c | 1c | 60,000 | 0.0640 | 0.0895 | 39.87 |
| 1 c | 0.5c | 60,000 | 0.0735 | 0.0877 | 19.26 |
| 1 c | 1c | 60,000 | 0.0711 | 0.0910 | 27.97 |
| 2c | 1c | 60,000 | 0.0846 | 0.1054 | 24.65 |
| 1c | 1.5 c | 60,000 | 0.0777 | 0.0899 | 15.76 |
| 0.5c | -0.5c | 60,000 | 0.0576 | 0.0574 | -0.42 |
| 0.5c | -1c | 60,000 | 0.0639 | 0.0530 | -17.00 |
| 1 c | -0.5c | 60,000 | 0.0725 | 0.0812 | 12.08 |
| 1 c | -1c | 60,000 | 0.0697 | 0.0704 | 1.06 |
| 2c | -1c | 60,000 | 0.0845 | 0.0948 | 12.19 |
| 1 c | -1.5c | 60,000 | 0.0719 | 0.0708 | -1.51 |

Table 4.35: Lift Slope all models, Second Method, Re 120,000

| Model |  | Speed | Lift Slope $\mathrm{C}_{\mathbf{L a}}(-)$ | Lift Slope $\mathrm{C}_{\mathbf{L a}}(-)$ | \% Variation |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | Re (-) | first region | second region |  |
| 1c | 0c | 120,000 | 0.0659 | 0.0822 | 24.74 |
| 0.5 c | 0 c | 120,000 | 0.0521 | 0.0677 | 29.77 |
| 0.5 c | 0.5 c | 120,000 | 0.0583 | 0.0628 | 7.82 |
| 0.5 c | 1 c | 120,000 | 0.0660 | 0.0796 | 20.57 |
| 1c | 0.5 c | 120,000 | 0.0647 | 0.0824 | 27.35 |
| 1c | 1 c | 120,000 | 0.0690 | 0.0889 | 28.82 |
| 2c | 1 c | 120,000 | 0.0795 | 0.1010 | 27.07 |
| 1c | 1.5 c | 120,000 | 0.0733 | 0.0924 | 25.95 |
| 0.5 c | -0.5 c | 120,000 | 0.0580 | 0.0483 | -16.72 |
| 0.5 c | -1 c | 120,000 | 0.0648 | 0.0535 | -17.35 |
| 1c | -0.5 c | 120,000 | 0.0666 | 0.0823 | 23.46 |
| 1c | -1 c | 120,000 | 0.0689 | 0.0642 | -6.88 |
| 2c | -1 c | 120,000 | 0.0835 | 0.0989 | 18.43 |
| 1c | -1.5 c | 120,000 | 0.0716 | 0.0637 | -11.07 |

To better understand this kind of behavior, a series of different tables were generated, trying to find the causes of this kind of change in the slope. The best comparison found was between the percentage change of lift slope and the angle of stagger, $\sigma$, defined before as the angle between the vertical line and the line joining the leading edges of the wings.

Table 4.36 and Table 4.37 show this comparison at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$.

Table 4.36: Lift Slope all models with Angle of Stagger, Second Method, Re 60,000

| Model |  |  | Speed | Lift Slope $C_{L a}(-)$ | Lift Slope $C_{L a}(-)$ | \% Variation |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | $\boldsymbol{\sigma}\left({ }^{\circ}\right)$ | Re $(-)$ | first region | second region |  |
| 0.5 c | -1 c | -63.4 | 60,000 | 0.0639 | 0.0530 | -17.00 |
| 1c | -1.5 c | -56.3 | 60,000 | 0.0719 | 0.0708 | -1.51 |
| 0.5 c | -0.5 c | -45 | 60,000 | 0.0576 | 0.0574 | -0.42 |
| 1c | -1 c | -45 | 60,000 | 0.0697 | 0.0704 | 1.06 |
| 1c | -0.5 c | -26.6 | 60,000 | 0.0725 | 0.0812 | 12.08 |
| 2c | -1 c | -26.6 | 60,000 | 0.0845 | 0.0948 | 12.19 |
| 0.5 c | 0 c | 0 | 60,000 | 0.0516 | 0.0641 | 24.10 |
| 1c | 0 c | 0 | 60,000 | 0.0631 | 0.0844 | 33.76 |
| 1c | 0.5 c | 26.6 | 60,000 | 0.0735 | 0.0877 | 19.26 |
| 2c | 1 c | 26.6 | 60,000 | 0.0846 | 0.1054 | 24.65 |
| 0.5 c | 0.5 c | 45 | 60,000 | 0.0591 | 0.0724 | 22.43 |
| 1c | 1 c | 45 | 60,000 | 0.0711 | 0.0910 | 27.97 |
| 1c | 1.5 c | 56.3 | 60,000 | 0.0777 | 0.0899 | 15.76 |
| 0.5 c | 1 c | 63.4 | 60,000 | 0.0640 | 0.0895 | 39.87 |

Table 4.37: Lift Slope all models with Angle of Stagger, Second Method, Re 120,000

| Model |  |  | Speed | Lift Slope C ${ }_{\text {La }}(-)$ | Lift Slope C ${ }_{\text {La }}(-)$ | \% Variation |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gap | Stagger | \% $\left(^{\circ}\right.$ ) | $\operatorname{Re}(-)$ | first region | second region |  |
| 0.5c | -1c | -63.4 | 120,000 | 0.0648 | 0.0535 | -17.35 |
| 1c | -1.5c | -56.3 | 120,000 | 0.0716 | 0.0637 | -11.07 |
| 0.5c | -0.5c | -45 | 120,000 | 0.0580 | 0.0483 | -16.72 |
| 1c | -1c | -45 | 120,000 | 0.0689 | 0.0642 | -6.88 |
| 1 c | -0.5c | -26.6 | 120,000 | 0.0666 | 0.0823 | 23.46 |
| 2c | -1c | -26.6 | 120,000 | 0.0835 | 0.0989 | 18.43 |
| 0.5c | 0c | 0 | 120,000 | 0.0521 | 0.0677 | 29.77 |
| 1c | 0c | 0 | 120,000 | 0.0659 | 0.0822 | 24.74 |
| 1c | 0.5c | 26.6 | 120,000 | 0.0647 | 0.0824 | 27.35 |
| 2c | 1c | 26.6 | 120,000 | 0.0795 | 0.1010 | 27.07 |
| 0.5c | 0.5c | 45 | 120,000 | 0.0583 | 0.0628 | 7.82 |
| 1c | 1c | 45 | 120,000 | 0.0690 | 0.0889 | 28.82 |
| 1c | 1.5c | 56.3 | 120,000 | 0.0733 | 0.0924 | 25.95 |
| 0.5c | 1c | 63.4 | 120,000 | 0.0660 | 0.0796 | 20.57 |

Figure 4.43 shows the percentage change of lift slope by changing the angle of stagger at Re 60,000 and 120,000.


Figure 4.43: Shows the Change in Lift Slope between first and second region varying angle of stagger at both Reynolds numbers of $\mathbf{6 0 , 0 0 0}$ and $\mathbf{1 2 0 , 0 0 0}$

It seems that the lift slope decreases in the second region with respect to the first region, while the stagger moves in the negative direction. The largest change in slope is associated with the model having highest angle of stagger and the lowest change in slope is associated with the model having lowest angle of stagger. The only exception is given by the model with angle of stagger $\sigma=56.3^{\circ}$, that is the model with $\mathrm{g}=1 \mathrm{c}$ and $\mathrm{s}=1.5 \mathrm{c}$. This model does not follow this trend. But it is still reasonable to lay down an hypothesis concerning the change in the lift slope. It can be possible that with increasing the angle of stagger, which can be realized by decreasing the gap or increasing the stagger, the model experiences a positive change in lift slope. This can be translated into a better span efficiency factor after this change has occurred.

### 4.6 Wind Tunnel Balance Data - Aerodynamic Hysteresis

Aerodynamic hysteresis of an airfoil refers to airfoil aerodynamic characteristics as it becomes history dependent, i.e., dependent on the sense of change in the angle of attack, near the airfoil stall angle. The coefficients of lift, and drag, of the airfoil are found to be multiple-valued rather than single-valued functions of the angle of attack. Aerodynamic hysteresis is of practical importance because it produces widely different values of lift coefficient and lift over drag ratio for a given angle of attack. It could also affect the recovery from stall conditions [21].

This section reports data obtained in the LSWT on the model with gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=-1 \mathrm{c}$ at Reynolds numbers 60,000 and 120,000. The $\alpha$ sweep for this analysis is from $-2^{\circ}$ to $8^{\circ}$ in steps of $1^{\circ}$.

## Lift Coefficient

Figure 4.44 and Figure 4.45 show the lift coefficient vs. angle of attack for the model with gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=-1 \mathrm{c}$ at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$.


Figure 4.44: Hysteresis Analysis on Lift Curve for the model with gap 0.5 c and stagger -1c at Re 60,000


Figure 4.45: Hysteresis Analysis on Lift Curve for the model with gap 0.5c and stagger -1c at Re 120,000

From Figure 4.44 very mild hysteresis of lift coefficient at Re 60,000 can be observed for the angles of attack lying between $16^{\circ}$ and $23^{\circ}$. From Figure 4.45 hysteresis of lift coefficient at Re 120,000 can be observed for the angles of attack
lying between $4^{\circ}$ and $8^{\circ}$. The difference in hysteresis location can be explained considering the different kind of hysteresis loops that can occur when the flow becomes unsteady [22]. One of the most important parameters affecting the dynamic behavior of an airfoil under unsteady flow is the reduced frequency, defined as:

$$
k=\frac{\omega \cdot c}{V_{\infty}}
$$

where $\omega$ is the frequency of oscillation of the airfoil, c is the chord, and $V_{\infty}$ is the freestream velocity. Keeping chord length constant, small reduced frequency means high freestream velocity or small frequency of oscillation. At small angles of attack and relatively small reduced frequencies the airfoil behavior can be treated with a potential flow approach. Accordingly, the viscous effects can be neglected. The boundary layer remains attached and the airfoil generates a sinusoidal wake. The effect of the oscillating wake is to produce a time lag between the actual conditions and the state of the boundary layer. As a result the lift will be in delay, and the characteristic is a closed loop that is described clockwise. This loop can be observed clearly in Figure 4.45 for the angles of attack lying between $3^{\circ}$ and $6^{\circ}$. Figure 4.46 shows a hysteresis loop in the range of angle of attack between $3^{\circ}$ and $6^{\circ}$ for the model with gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=-1 \mathrm{c}$ at $\operatorname{Re} 120,000$. The hysteresis loop is clockwise and it occurs before stall.


Figure 4.46: Hysteresis Loop in the range of angle of attack between $3^{\circ}$ and $6^{\circ}$ for the model with gap 0.5 c and stagger -1 c at Re 120,000

If the oscillation occurs around a mean angle of attack close to $\mathrm{C}_{\text {Lmax }}$ (static stall) viscous effects become predominant. Starting from the point of minimum incidence, the dynamic lift follows the static lift, until the static lift curve deflects, due to increasing trailing edge separation. The dynamic lift, instead, keeps growing almost linearly until a breakdown occurs. At the breakdown point there is massive flow separation and the lift drops to levels far below those typical of the static curve. It will take some time to recover more regular behavior, but the lift will remain below the static lift for most of the remaining loop. The increase of the lift above the static $C_{\text {Lmax }}$ is attributed to the development of a leading edge vortex on the upper surface that grows and travels downstream. The breakdown is associated with the point when the leading edge vortex has travelled past the airfoil trailing edge. The loop is described clockwise.

From Figure 4.44 and Figure 4.45 it can be seen that there is no hysteresis around $C_{\text {Lmax }}$. The last kind of hysteresis loops that can occur is above $C_{\text {Lmax. }}$. They are characteristic of a loss of energy (due to viscous dissipation) that is proportional to the area enclosed by the loop and they usually occur at high reduced frequencies, which means low Reynolds numbers. This loop can be observed clearly in Figure 4.44 for the angles of attack lying between $18^{\circ}$ and $23^{\circ}$. Figure 4.47 shows a hysteresis loop in the range of angle of attack between $18^{\circ}$ and $23^{\circ}$ for the model with gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=-1 \mathrm{c}$ at $\operatorname{Re} 60,000$. The hysteresis loop is clockwise and it occurs after stall.


Figure 4.47: Hysteresis Loop in the range of angle of attack between $18^{\circ}$ and $23^{\circ}$ for the model with gap 0.5 c and stagger -1 c at $\operatorname{Re} \mathbf{6 0 , 0 0 0}$

### 4.7 Comparison Analytical - Experimental Results

## Overview

Computational analysis for the biplane configurations was carried out using AVL. This code utilizes a vortex-lattice representation of the lifting surfaces, and it assumes steady, irrotational, inviscid, incompressible attached flow. It predicts lift and induced drag coefficients only in the linear region before stall, and at low freestream velocities. In this section, analytical results of lift coefficients and induced drag coefficients, will be compared to the experimental results.

## Lift Coefficient

Once the value of lift coefficient obtained from the LSWT were analyzed, it was possible to compare the results with those obtained using AVL. This code, as can only generate lift and drag coefficients in the linear region of lift coefficient. Figure 4.48 plots the lift coefficient versus angle of attack in the linear region from $2^{\circ}$ to $10^{\circ}$ angle of attack, for models with similar gap $g=0.5 \mathrm{c}$, and both positive and negative stagger $s= \pm 0.5 \mathrm{c}$, at $\operatorname{Re} 60,000$. It also plots the data of lift coefficient versus angle of attack obtained using AVL. From Table 4.1 and Table 4.2, it can be seen that the difference at $\operatorname{Re} 60,000$ and 120,000 in lift slope in a range of angle of attack between $-2^{\circ}$ and $8^{\circ}$ for the two models with similar gap $g=0.5 \mathrm{c}$ and stagger s $= \pm 0.5 \mathrm{c}$ is really small, and equal to $1.4 \%$ for the positive stagger configuration and $0.5 \%$ for the negative stagger configuration. Due to the small difference in lift slope changing Reynolds number, for the comparison between analytical and
experimental results, the case $\operatorname{Re} 60,000$ is the only one considered. Also, it is important to notice that results from AVL show no difference between positive and negative stagger.


Figure 4.48: Comparison of Lift Curve between the AVL results and the LSWT data for the models with similar gap 0.5 c and stagger $\pm 0.5 \mathrm{C}$ at Re $\mathbf{6 0 , 0 0 0}$

The AVL results show good agreement with the LSWT data under the angle of attack range where the transition from one lift slope value to the other occurs. The change in lift slope, explained in the previous paragraph, is the only effect that causes AVL results to be different from the experimental one. All models tested show this agreement when comparing analytical and experimental results.

## Induced Drag Coefficient

The induced drag coefficient, $\mathrm{C}_{\mathrm{Di}}$, was predicted theoretically using both the equivalent monoplane method (EM) by Munk and orthogonal biplane method (OB)
by Prandtl. Figure 4.49 shows the comparison between analytical and experimental induced drag coefficient, $C_{\text {Di }}$, versus angle of attack, $\alpha$, for the models with similar gap $g=0.5 c$, and both positive and negative stagger $s= \pm 0.5 \mathrm{c}$ at $\operatorname{Re} 60,000$.


Figure 4.49: Comparison of Induced Drag Coefficient between the Orthogonal Biplane theory (OB), Equivalent Monoplane theory (EM) and the LSWT data for the models with similar gap 0.5 c and stagger $\pm 0.5 \mathrm{c}$ at $\operatorname{Re} 60,000$

The theory shows good agreement with the experimental results. And again, from Figure 4.49, it can be seen that the theory matches the experimental curves reasonably well under the angle of attack range where the transition from one lift slope value to the other occurs. All models tested show this agreement when comparing analytical and experimental results.

### 4.8 Uncertainty Analysis

### 4.8.1 Overview

Uncertainty analysis was performed on the lift over drag ratio and on the lift coefficient for the model with gap $\mathrm{g}=0.5 \mathrm{c}$ and stagger $\mathrm{s}=1 \mathrm{c}$.

### 4.8.2 Lift over Drag Ratio

In this case the quantity that we wish to determine is derived from measured quantities of lift and drag. To find the uncertainty in lift over drag ratio caused by the uncertainties in lift and drag, we consider the contribution due to the uncertainty in lift and the contribution due to the uncertainty in drag separately. Each contribution may be considered separately as long as the variables lift and drag are independent of each other. The total error is obtained by combining the individual contributions. The basic idea is to determine by how much the aerodynamic efficiency would change if lift (or drag) were changed by its uncertainty. This means we need to derive the expression of lift over drag ratio with respect to both lift and drag:

$$
\begin{gathered}
\frac{\partial(L / D)}{\partial L}=\frac{1}{D} \\
\frac{\partial(L / D)}{\partial D}=-\frac{L}{D^{2}}
\end{gathered}
$$

Since each uncertain variable will increase, not decrease the final uncertainty, we take the absolute value of these derivatives. We can consider two different scenarios. The worst scenario is the case where each possible error occurs in the
same direction. The realistic scenario is a geometric mean of the possible errors. The uncertainty in lift over drag ratio for these two scenarios is equal to:

$$
\begin{gather*}
\Delta(L / D)_{\text {worst }}=\left|\frac{\partial(L / D)}{\partial L} \Delta L\right|+\left|\frac{\partial(L / D)}{\partial D} \Delta D\right|=\left|\frac{1}{D} \Delta L\right|+\left|\frac{L}{D^{2}} \Delta D\right|  \tag{4.1}\\
\Delta(L / D)_{\text {realistic }}=\sqrt{\left(\frac{\partial(L / D)}{\partial L} \Delta L\right)^{2}+\left(\frac{\partial(L / D)}{\partial D} \Delta D\right)^{2}}=\sqrt{\left(\frac{1}{D} \Delta L\right)^{2}+\left(\frac{L}{D^{2}} \Delta D\right)^{2}} \tag{4.2}
\end{gather*}
$$

where $\Delta L$ and $\Delta D$ are the possible error in the lift force and drag force measurement.

We can now compute the uncertainty in the lift over drag ratio using the values specified by the manufacturer of the balance used for taking the forces in the LSWT for $\Delta L$ and $\Delta D$ in Equations (4.1) and (4.2). These values are listed in Table 3.2 and they are both equal to $1 / 2560 \mathrm{lb}$, and they are considered to be conservative for actual data acquisition. The standard deviation of the force measurements at each given angle of attack can also be determined. Multiple data points are logged for each angle, therefore the actual sample variance and hence the standard deviation can be computed. Each standard deviation will be the value used for $\Delta L$ and $\Delta D$ in Equations (4.1) and (4.2).

The uncertainty analysis performed by using the balance resolution will be called conservative uncertainty analysis. The uncertainty analysis performed by using the standard deviations will be called tested uncertainty analysis, and it gives a more realistic error value.

### 4.8.2.1 Conservative Uncertainty Analysis

Figure 4.50 and Figure 4.51 show the lift over drag ratio vs. angle of attack for both Reynolds numbers 60,000 and 120,000 with error bars calculated using the conservative method.


Figure 4.50: Shows the error bars of Lift over Drag Ratio for the worst case of the Conservative Uncertainty Analysis, for the model with gap 0.5c and stagger 1c at Re $60,000-120,000$


Angle of Attack, $\alpha\left({ }^{\circ}\right)$
Figure 4.51: Shows the error bars of Lift over Drag Ratio for the realistic case of the Conservative Uncertainty Analysis, for the model with gap 0.5c and stagger 1c at $\operatorname{Re} 60,000-120,000$

The results of this analysis can be seen in Table 4.38 and Table 4.39 where the percentage of error in lift over drag ratio is reported for the region where this percentage is higher. From Figure 4.50 and Figure 4.51 we can see that this region is where the lift over drag is higher which is in a range of angle of attack between $2^{\circ}$ and $6^{\circ}$.

Table 4.38: Conservative Uncertainty Analysis in L/D, $\mathrm{g}=\mathbf{0 . 5} \mathrm{c}$ and $\mathrm{s}=1 \mathrm{c}, \operatorname{Re} 60,000$

| $\boldsymbol{\alpha}(\boldsymbol{9})$ | L/D (-) | \% Worst Case Error | \% Realistic Case Error |
| :---: | :---: | :---: | :---: |
| 1.97 | 5.3214 | 14.0983 | 12.0758 |
| 2.22 | 6.6169 | 17.2616 | 15.1657 |
| 2.47 | 6.1775 | 14.6157 | 12.7431 |
| 2.72 | 6.2978 | 15.0791 | 13.1759 |
| 2.97 | 6.7500 | 16.2525 | 14.3099 |
| 3.22 | 7.4063 | 16.4045 | 14.5842 |
| 3.47 | 7.4063 | 16.4045 | 14.5842 |
| 3.72 | 7.9075 | 17.1446 | 15.3411 |
| 3.97 | 7.2115 | 12.6788 | 11.2413 |
| 4.22 | 7.7074 | 13.9027 | 12.4092 |
| 4.47 | 7.4193 | 12.3224 | 10.9570 |
| 4.72 | 7.8686 | 12.8462 | 11.4894 |
| 4.97 | 7.6472 | 11.3549 | 10.1273 |
| 5.22 | 8.1205 | 11.6498 | 10.4508 |
| 5.47 | 7.8313 | 10.0882 | 9.0185 |
| 5.72 | 7.5668 | 9.1884 | 8.1864 |
| 5.97 | 7.5533 | 8.8367 | 7.8716 |

Table 4.39: Conservative Uncertainty Analysis in $L / D, g=0.5 c$ and $s=1 c, \operatorname{Re} 120,000$

| $\alpha\left({ }^{9}\right)$ | L/D (-) | \% Worst Case Error | \% Realistic Case Error |
| :---: | :---: | :---: | :---: |
| 2.1225 | 6.1613 | 4.1622 | 3.6278 |
| 2.3725 | 7.2590 | 5.0730 | 4.5008 |
| 2.6225 | 7.2567 | 4.9419 | 4.3844 |
| 2.8725 | 7.4566 | 4.4377 | 3.9480 |
| 3.1225 | 7.7592 | 4.2040 | 3.7548 |
| 3.3725 | 7.7592 | 4.2040 | 3.7548 |
| 3.6225 | 7.5785 | 4.0481 | 3.6072 |
| 3.8725 | 7.5787 | 3.6665 | 3.2672 |
| 4.1225 | 7.4027 | 3.3187 | 2.9503 |
| 4.3725 | 7.7980 | 3.3523 | 2.9956 |
| 4.6225 | 7.7070 | 3.1308 | 2.7945 |
| 4.8725 | 8.0231 | 3.0790 | 2.7590 |
| 5.1225 | 8.1167 | 2.8481 | 2.5549 |
| 5.3725 | 7.8609 | 2.4913 | 2.2279 |
| 5.6225 | 7.7031 | 2.3348 | 2.0838 |
| 5.8725 | 7.5696 | 2.1283 | 1.8963 |
| 6.1225 | 7.4864 | 1.9981 | 1.7783 |

In general, for the entire range of angle of attack, at Reynolds number 60,000 , the worst case error ranges from $0.36 \%$ to $17.26 \%$, and the realistic case error ranges from $0.28 \%$ to $15.34 \%$. At Reynolds number 120,000 , the worst case error ranges from $0.09 \%$ to $5.07 \%$, and the realistic case error ranges from $0.07 \%$ to $4.50 \%$. The error values are decreasing as the Reynolds number is increased. This is due to higher velocities producing greater balance loading and therefore, increased sensitivity. The lower loadings are closer to the balance's minimum resolution which causes the error to increase.

### 4.8.2.2 Tested Uncertainty Analysis

Figure 4.52 and Figure 4.53 show the lift over drag ratio vs. angle of attack for both Reynolds numbers 60,000 and 120,000 with error bars calculated using the tested method.


Angle of Attack, $a\left({ }^{\circ}\right)$
Figure 4.52: Shows the error bars of Lift over Drag Ratio for the worst case of the Tested Uncertainty Analysis, for the model with gap 0.5c and stagger 1c at Re $60,000-120,000$


Figure 4.53: Shows the error bars of Lift over Drag Ratio for the realistic case of the Tested Uncertainty Analysis, for the model with gap 0.5c and stagger 1c at Re $60,000-120,000$

The results of this analysis can be seen in Table 4.40 where the percentage of error in lift over drag ratio is reported for the region where this percentage is higher. From Figure 4.52 and Figure 4.53 this region is where the lift over drag is higher which is in a range of angle of attack between $2^{\circ}$ and $6^{\circ}$.

Table 4.40: Tested Uncertainty Analysis in L/D, $\mathrm{g}=0.5 \mathrm{c}$ and $\mathrm{s}=1 \mathrm{c}, \operatorname{Re} 60,000$

| $\boldsymbol{\alpha}$ ( $\boldsymbol{9}$ | L/D (-) | \% Worst Case Error | \% Realistic Case Error |
| :---: | :---: | :---: | :---: |
| 2.1225 | 6.1613 | 4.1622 | 3.6278 |
| 2.3725 | 7.2590 | 5.0730 | 4.5008 |
| 2.6225 | 7.2567 | 4.9419 | 4.3844 |
| 2.8725 | 7.4566 | 4.4377 | 3.9480 |
| 3.1225 | 7.7592 | 4.2040 | 3.7548 |
| 3.3725 | 7.7592 | 4.2040 | 3.7548 |
| 3.6225 | 7.5785 | 4.0481 | 3.6072 |
| 3.8725 | 7.5787 | 3.6665 | 3.2672 |
| 4.1225 | 7.4027 | 3.3187 | 2.9503 |
| 4.3725 | 7.7980 | 3.3523 | 2.9956 |
| 4.6225 | 7.7070 | 3.1308 | 2.7945 |
| 4.8725 | 8.0231 | 3.0790 | 2.7590 |
| 5.1225 | 8.1167 | 2.8481 | 2.5549 |
| 5.3725 | 7.8609 | 2.4913 | 2.2279 |
| 5.6225 | 7.7031 | 2.3348 | 2.0838 |
| 5.8725 | 7.5696 | 2.1283 | 1.8963 |
| 6.1225 | 7.4864 | 1.9981 | 1.7783 |

Table 4.41: Tested Uncertainty Analysis in $L / D, g=0.5 \mathrm{c}$ and $s=1 c, \operatorname{Re} 120,000$

| $\boldsymbol{\alpha}(9)$ | L/D (-) | \% Worst Case Error | \% Realistic Case Error |
| :---: | :---: | :---: | :---: |
| 2.1225 | 6.1613 | 4.1622 | 3.6278 |
| 2.3725 | 7.2590 | 5.0730 | 4.5008 |
| 2.6225 | 7.2567 | 4.9419 | 4.3844 |
| 2.8725 | 7.4566 | 4.4377 | 3.9480 |
| 3.1225 | 7.7592 | 4.2040 | 3.7548 |
| 3.3725 | 7.7592 | 4.2040 | 3.7548 |
| 3.6225 | 7.5785 | 4.0481 | 3.6072 |
| 3.8725 | 7.5787 | 3.6665 | 3.2672 |
| 4.1225 | 7.4027 | 3.3187 | 2.9503 |
| 4.3725 | 7.7980 | 3.3523 | 2.9956 |
| 4.6225 | 7.7070 | 3.1308 | 2.7945 |
| 4.8725 | 8.0231 | 3.0790 | 2.7590 |
| 5.1225 | 8.1167 | 2.8481 | 2.5549 |
| 5.3725 | 7.8609 | 2.4913 | 2.2279 |
| 5.6225 | 7.7031 | 2.3348 | 2.0838 |
| 5.8725 | 7.5696 | 2.1283 | 1.8963 |
| 6.1225 | 7.4864 | 1.9981 | 1.7783 |

In general, for the entire range of angle of attack, at Reynolds number 60,000 , the worst case error ranges from $0.18 \%$ to $6.46 \%$, and the realistic case error ranges from $0.13 \%$ to $4.85 \%$. At Reynolds number 120,000 , the worst case error ranges from $0.10 \%$ to $4.43 \%$, and the realistic case error ranges from $0.08 \%$ to $3.96 \%$. The error values are decreasing as the Reynolds number is increased. This is due to higher velocities producing greater balance loading and therefore, increased sensitivity. The lower loadings are closer to the balance's minimum resolution which causes the error to increase.

As compared to the conservative error percentages, the tested error percentages are much smaller and thus reflect the actual precision of the balance for measurements.

### 4.8.3 Lift Coefficient

The lift coefficient is computed using Equation (4.3):

$$
\begin{equation*}
C_{L}=\frac{L}{\frac{1}{2} \rho_{a i r} S V_{\infty}^{2}} \tag{4.3}
\end{equation*}
$$

In this case the quantity that we wish to determine is derived from measured quantities of lift, air density, surface, and freestream velocity. As has been done before, the derivative of Equation (4.3) with respect to all the variables can be computed.

$$
\begin{gathered}
\frac{\partial C_{L}}{\partial L}=\frac{1}{\frac{1}{2} \rho_{\text {air }} S V_{-}^{2}} \\
\frac{\partial C_{l}}{\partial \rho_{\text {air }}}=-\frac{L}{\frac{1}{2} \rho_{\text {air }}^{2} S V_{\infty}^{2}} \\
\frac{\partial C_{L}}{\partial S}=-\frac{L}{\frac{1}{2} \rho_{a i r} S^{2} V_{\infty}^{2}} \\
\frac{\partial C_{L}}{\partial V_{\infty}}=-\frac{L}{\frac{1}{4} \rho_{a i r} S V_{\infty}^{3}}
\end{gathered}
$$

The total error is obtained by combining the individual contributions. Since each uncertain variable will increase, not decrease the final uncertainty, we take the absolute value of these derivatives. For uncertainty analysis on lift coefficient, only the realistic scenario is considered in this section, and the method used is the conservative uncertainty analysis.

$$
\begin{equation*}
\Delta\left(C_{L}\right)_{\text {realisitic }}=\sqrt{\left(\frac{\partial C_{L}}{\partial L} \Delta L\right)^{2}+\left(\frac{\partial C_{L}}{\partial \rho_{\text {air }}} \Delta \rho_{\text {air }}\right)^{2}+\left(\frac{\partial C_{L}}{\partial S} \Delta S\right)^{2}+\left(\frac{\partial C_{L}}{\partial V_{-}} \Delta V_{\infty}\right)^{2}} \tag{4.4}
\end{equation*}
$$

where, for the conservative method $\Delta L$ is equal to $1 / 2560 \mathrm{lb}$. The other possible errors $\Delta \rho_{\text {air }}, \Delta S$, and $\Delta V_{\infty}$ are functions of other variables, and they can be found by using the definition of air density, surface, and freestream velocity respectively.

Air density is calculated using Equation (3.2):

$$
\rho_{\text {air }}=1.2929\left(\frac{273.15}{r+273.15}\right)\left(\frac{p_{\text {amm }}-S V P \cdot H}{760}\right)
$$

Freestream velocity is calculated using Equation (3.1):

$$
\begin{aligned}
\operatorname{Re} & =\frac{V_{\infty} \cdot \bar{c} \cdot \rho_{\text {air }}}{\mu_{\text {air }}} \\
V_{\infty} & =\frac{\operatorname{Re} \cdot \mu_{\text {air }}}{\bar{c} \cdot \rho_{\text {air }}}
\end{aligned}
$$

where it can be seen that the freestream velocity is also function of other variables, so that the error due to these variables should be considered. The air viscosity, $\mu_{\text {air }}$, is calculated using Sutherland's formula in Equation (3.3):

$$
\mu_{\text {air }}=\mu_{0} \frac{T_{0}+C}{T+C}\left(\frac{T}{T_{0}}\right)^{3 / 2}
$$

Now, taking the derivative of air density with respect to temperature, atmospheric pressure, and humidity, the derivative of freestream velocity with respect to air viscosity, chord length, and air density, the derivative of air viscosity with respect to temperature, the possible errors associated with air density, and freestream velocity, can be computed. Using these values in Equation (4.3), the error on lift coefficient can be computed.

Figure 4.54 shows the lift coefficient vs. angle of attack for both Reynolds numbers 60,000 and 120,000 with error bars calculated using the conservative method, for the model with gap $g=0.5 \mathrm{c}$ and stagger $\mathrm{s}=1 \mathrm{c}$.


Figure 4.54: Shows the error bars of Lift Coefficient for the realistic case of the Conservative Uncertainty Analysis, for the model with gap 0.5c and stagger 1c at Re $60,000-120,000$

In general, for the entire range of angle of attack, at Reynolds number 60,000 the realistic case error ranges from $0.08 \%$ to $4.49 \%$. At Reynolds number 120,000 the realistic case error ranges from $0.02 \%$ to $0.72 \%$. The error values are decreasing as the Reynolds number is increased. This is due to higher velocities producing greater balance loading and therefore, increased sensitivity. The lower loadings are closer to the balance's minimum resolution which causes the error to increase.

## Bibliography

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## V. Conclusions and Recommendations

### 5.1 Conclusions of Research

The purpose of this study was to investigate the aerodynamic performance of a biplane joined at the tips by endplates using several biplane configurations having different stagger and gap. The models were based on rectangular flat plate profiles with the same reference planform area.

When comparing models with a gap of 0.5 c , at $\operatorname{Re} 60,000$ the lift slope increases by almost $5 \%$ for each 0.5 c stagger increment in the negative direction, and it increases by almost $15 \%$ for each 0.5 c stagger increment in the positive direction. At Re 120,000 the variation is slightly different but the trend is the same. It can be seen that at $\operatorname{Re} 120,000$ the lift slope increases by almost $4 \%$ for each 0.5 c stagger increment in the negative direction, and it increases by almost $13 \%$ for each 0.5 c stagger increment in the positive direction. The aerodynamic efficiency, or lift over drag ratio, increases by almost $10 \%$ for the first 0.5 c stagger increment in the negative direction then it increases by another $5 \%$ for the second 0.5 c stagger increment in the negative direction. For the positive stagger increment it can be seen that the maximum lift over drag ratio increases by almost $10 \%$ for each 0.5 c stagger increment in the positive direction. At Re 120,000 the percentages of variation between different staggers is slightly different but the trend is the same. It can be seen that at $\operatorname{Re} 120,000$ the aerodynamic efficiency increases by almost $5 \%$
for each 0.5 c stagger increment in the negative direction, and it increases by almost $20 \%$ for each 0.5 c stagger increment in the positive direction.

When comparing models with a gap of 1c, at $\operatorname{Re} 60,000$ and $\operatorname{Re} 120,000$ the variation of lift slope and aerodynamic efficiency is smaller than the case of models with a gap of 0.5 c due to the fact that higher gap means less interference between the two wings.

When comparing models with a gap of $2 c$, at both Reynolds numbers, the variation of lift slope and aerodynamic efficiency is almost negligible due to the higher gap between the two wings. Also, models with a gap of 2c experience stall flutter after an angle of attack equal to the angle of stall.

When comparing models with a stagger of $0 c$, the lift slope increases by almost $26 \%$ at $\operatorname{Re} 60,000$ and it increases by almost $28 \%$ at $\operatorname{Re} 120,000$ for each 0.5 c gap increment. The aerodynamic efficiency increases by almost $1.5 \%$ at $\operatorname{Re} 60,000$ and by almost $35 \%$ at $\operatorname{Re} 120,000$ for each 0.5 c gap increment.

When comparing models with a stagger of $-1 c$, the lift slope, at both Reynolds numbers, increases by almost $15 \%$ with each 0.5 gap increment. Higher gap models experience higher lift slope. At Re 120,000 the integrated lift over drag ratio of the model with a gap of 1 c is almost $9 \%$ higher than the integrated lift over drag ratio of the model with gap 0.5 c. The variation of integrated lift over drag ratio equals almost $40 \%$ when comparing the model with a gap of 2 c to the model with gap 0.5 c . This variation at $\operatorname{Re} 60,000$ is equal to just $0.8 \%$ for the model with a gap of 1 c and almost $12 \%$ for the one with a gap of 2 c . This means that in terms of
aerodynamic efficiency, the effect of changing the gap at higher Reynolds numbers is more important than the same effect at lower Reynolds numbers. The same result was obtained when considering models with no stagger and different gap configurations.

When comparing positive versus negative stagger, the lift slope of the models with positive stagger is always higher than those with negative stagger. At both Reynolds numbers this variation is equal to almost $15 \%$. At $\operatorname{Re} 120,000$ the integrated lift over drag ratio of the model having positive stagger is almost $18 \%$ higher than the integrated lift over drag ratio of the model with negative stagger. This variation at $\operatorname{Re} 60,000$ is equal to almost $5.5 \%$. This means that in terms of aerodynamic efficiency, models with positive stagger are more efficient than models with negative stagger and that this behavior, and that this effect becomes more important at higher Reynolds numbers.

A change in lift slope was observed for all the models tested in the LSWT. The largest change in lift slope is associated with the model with the highest angle of stagger and the lowest change in lift slope is associated with the model having lowest angle of stagger. It is possible that by increasing the angle of stagger, which can be realized by decreasing the gap or increasing the stagger, the model experiences a positive change in lift slope. This can be translated into a better span efficiency factor after this change has occurred. Also, this change in lift slope occurs around the same angle of attack as maximum aerodynamic efficiency.

Aerodynamic hysteresis analysis was conducted on the model with a gap of 0.5 c and a stagger of -1 c . Hysteresis on lift coefficient is very mild and at Reynolds number 60,000 it occurs after the angle of stall, instead at Reynolds number 120,000 hysteresis occurs before the angle of stall.

Comparison between analytical and experimental results was conducted on all the models tested in the LSWT by using for the analytical results the code AVL. The two were found to compare quite well under some conditions, and not well under other conditions. This is due to the fact that AVL, or any other code based on Laplace's Equation, utilizes a vortex-lattice representation of the lifting surfaces, and it assumes steady, irrotational, inviscid, incompressible attached flow.

It can be seen that both stagger and gap have positive effect on the aerodynamic performance of the biplane joined at the tips by endplates. Also, despite Munk's Stagger Theorem, it can be seen that positive and negative stagger do not show the same aerodynamic performance while keeping gap constant. Positive stagger configurations always perform better than negative stagger configurations in terms of aerodynamic performance. But the most important result obtained from this research is that all the models tested in the LSWT experience a change in lift slope at the angle of attack of maximum aerodynamic efficiency. Further investigation of the biplane joined at the tips will need to be conducted to determine the reason why this change in lift slope occurs, and if it is possible to take advantage of this change.

### 5.2 Recommendations for Future Research

It is recommended that there be continued testing of the biplane joined at the tips by endplates in order to better understand its aerodynamic performance, and in order to study the change in lift slope observed. A deeper analysis of lift slope could be conducted by using PIV tests or CFD analysis on the biplane with endplates, around an angle of attack where the maximum aerodynamic efficiency is experienced. Also, it would be interesting to run experimental wind tunnel tests on different monoplane configurations with same reference area, or same aspect ratio of the biplane configurations tested in the LSWT and presented in this thesis. It could be interesting to run tests at higher Reynolds numbers to see if the concept of a biplane joined at the tips by endplates could be used for higher speed and/or bigger airplanes.

Analytically, a deeper study on the interference of one wing on the other, and the effect of endplates on the aerodynamic performance should be conducted, considering viscous effects that have been neglected using a vortex lattice approach.

Appendix A: ATI Gamma F/T Transducer



Appendix B: Newmark RM-5 Rotary Stage


## Appendix C: Wind Tunnel Conditions

## Model \#1-Gap 1c, Stagger 0c

Table C.1: Model \#1 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: Model \#1 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Re (-) | V (m/s) | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.83 | 24 -jan-08 | 13 | 745.5 mmHg | $52 \%$ |
| $\mathbf{1 2 0 , 0 0 0}$ | 17.65 | $24-j a n-08$ | 13 | 745.5 mmHg | $52 \%$ |
| 60,000 | 9.07 | 28 -jan- 08 | 17 | 745.8 mmHg | $57 \%$ |
| 120,000 | 18.15 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 60,000 | 9.46 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.92 | $05-$ feb- 08 | 21 | 740.4 mmHg | $82 \%$ |

## Model \#2 - Gap 0.5c, Stagger 0c

Table C.2: Model \# 2 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: Model \#2 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Re (-) | V (m/s) | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.83 | 24 -jan- 08 | 13 | 745.5 mmHg | $52 \%$ |
| 120,000 | $\mathbf{1 7 . 6 5}$ | 24 -jan-08 | 13 | 745.5 mmHg | $52 \%$ |
| 60,000 | 9.07 | 28 -jan- 08 | 17 | 745.8 mmHg | $57 \%$ |
| 120,000 | 18.15 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 60,000 | 9.46 | $05-\mathrm{feb}-08$ | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.92 | 05 -feb-08 | 21 | 740.4 mmHg | $82 \%$ |

Model \#3 - Gap 0.5c, Stagger 0.5c
Table C.3: Model \#3 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: Model \#3 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\operatorname{Re}(-)$ | $\mathbf{V}(\mathrm{m} / \mathrm{s})$ | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.83 | 24 -jan-08 | 13 | 745.5 mmHg | $52 \%$ |
| 120,000 | 17.65 | 24 -jan-08 | 13 | 745.5 mmHg | $52 \%$ |
| 60,000 | 9.07 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 120,000 | 18.15 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 60,000 | 9.46 | $05-$ feb- 08 | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.92 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |

## Model \#4-Gap 0.5c, Stagger 1c

Table C.4: Model \#4 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Model \#4 |  |  |  |  |  |
| Re (-) | V (m/s) | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.82 | $25-$ jan-08 | 15 | 755 mmHg | $46 \%$ |
| 120,000 | 17.65 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |
| 60,000 | 9.07 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 120,000 | 18.15 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 60,000 | 9.46 | $05-\mathrm{feb}-08$ | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.92 | $05-\mathrm{feb}-08$ | 21 | 740.4 mmHg | $82 \%$ |

## Model \#5 - Gap 1c, Stagger 0.5c

Table C.5: Model \# 5 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model Configuration: | Model \#5 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\operatorname { R e }}(-)$ | V (m/s) | Date | Temperature $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric Pressure | Humidity |
| 60,000 | 8.82 | 25-ian-08 | 15 | 755 mmHg | 46\% |
| 120,000 | 17.65 | 25-ian-08 | 15 | 755 mmHg | 46\% |
| 60,000 | 9.07 | 28-ian-08 | 17 | 745.8 mmHg | 57\% |
| 120,000 | 18.15 | 28-ian-08 | 17 | 745.8 mmHg | 57\% |
| 60,000 | 9.46 | 05-feb-08 | 21 | 740.4 mmHg | 82\% |
| 120,000 | 18.92 | 05-feb-08 | 21 | 740.4 mmHg | 82\% |

Model \#6-Gap 1c, Stagger 1c
Table C.6: Model \# 6 Velocity, Test Dates, Temperature, Pressure, and Humidity


Model \# 7 - Gap 2c, Stagger 1c

Table C.7: Model \#7 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Model \#7 |  |  |  |  |  |
| Re (-) | $\mathbf{V ( m / s})$ | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.82 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |
| 120,000 | 17.65 | $25-j a n-08$ | 15 | 755 mmHg | $46 \%$ |
| 60,000 | 9.21 | 28 -jan-08 | 18 | 739.8 mmHg | $56 \%$ |
| 120,000 | 9.46 | $05-$ feb- 08 | 21 | 740.4 mmHg | $82 \%$ |
| 60,000 | 18.92 | 05 -feb-08 | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.15 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |

## Model \#8-Gap 1c, Stagger 1.5c

Table C.8: Model \#8 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Re (-) | Model \#8 (m/s) | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.82 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |
| 120,000 | 17.65 | $25-j a n-08$ | 15 | 755 mmHg | $46 \%$ |
| 60,000 | 9.21 | 28 -jan-08 | 18 | 739.8 mmHg | $56 \%$ |
| 120,000 | 18.42 | 28 -jan-08 | 18 | 739.8 mmHg | $56 \%$ |
| 60,000 | 9.49 | 05 -feb-08 | 21 | 738.1 mmHg | $81 \%$ |
| 120,000 | 18.97 | 05 -feb-08 | 21 | 738.1 mmHg | $81 \%$ |

## Model \#9 - Gap 0.5c, Stagger -0.5c

Table C.9: Model \#9 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: | Model \#9 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Re (-) | V (m/s) | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |  |
| 60,000 | 8.83 | 24 -jan- 08 | 13 | 745.5 mmHg | $52 \%$ |  |
| 120,000 | 17.65 | 24 -jan- 08 | 13 | 745.5 mmHg | $52 \%$ |  |
| 60,000 | 9.07 | 28 -jan- 08 | 17 | 745.8 mmHg | $57 \%$ |  |
| 120,000 | 18.15 | 28 -jan- 08 | 17 | 745.8 mmHg | $57 \%$ |  |
| 60,000 | 9.46 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |  |
| 120,000 | 18.92 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |  |

## Model \#10-Gap 0.5c, Stagger -1c

Table C.10: Model \#10 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: Model \#10 |  |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\operatorname{Re}(-)$ | $\mathbf{V}(\mathrm{m} / \mathrm{s})$ | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |  |  |  |  |  |  |  |
| 60,000 | 8.82 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |  |  |  |  |  |  |  |
| 120,000 | 17.65 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |  |  |  |  |  |  |  |
| 60,000 | 9.07 | 28 -ian- 08 | 17 | 745.8 mmHg | $57 \%$ |  |  |  |  |  |  |  |
| 120,000 | 18.15 | 28 -jan- 08 | 17 | 745.8 mmHg | $57 \%$ |  |  |  |  |  |  |  |
| 60,000 | 9.46 | 05 -feb-08 | 21 | 740.4 mmHg | $82 \%$ |  |  |  |  |  |  |  |
| 120.000 | 18.92 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |  |  |  |  |  |  |  |

## Model \#11-Gap 1c, Stagger -0.5c

Table C.11: Model \#11 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: Model \#11 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Re (-) | V (m/s) | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.82 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |
| 120,000 | 17.65 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |
| 60,000 | 9.07 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 120,000 | 18.15 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 60,000 | 9.46 | 05 -feb-08 | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.92 | 05 -feb-08 | 21 | 740.4 mmHg | $82 \%$ |

## Model \#12-Gap 1c, Stagger -1c

Table C.12: Model \#12 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model <br> Configuration: Model \#12 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Re (-) | $\mathbf{V}(\mathrm{m} / \mathrm{s})$ | Date | Temperature <br> $\left({ }^{\circ} \mathrm{C}\right)$ | Atmospheric <br> Pressure | Humidity |
| 60,000 | 8.82 | 25 -jan-08 | 15 | 755 mmHg | $46 \%$ |
| 120,000 | 17.65 | $25-j a n-08$ | 15 | 755 mmHg | $46 \%$ |
| 60,000 | 9.07 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 120,000 | 18.15 | 28 -jan-08 | 17 | 745.8 mmHg | $57 \%$ |
| 60,000 | 9.46 | 05 -feb-08 | 21 | 740.4 mmHg | $82 \%$ |
| 120,000 | 18.92 | 05 -feb- 08 | 21 | 740.4 mmHg | $82 \%$ |

Model \#13-Gap 2c, Stagger -1c
Table C.13: Model \#13 Velocity, Test Dates, Temperature, Pressure, and Humidity

| Model Configuration: | Model \#13 |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{R e}(-)$ | $\mathrm{V}(\mathrm{m} / \mathrm{s})$ | Date | Temperature ( ${ }^{\circ} \mathrm{C}$ ) | Atmospheric Pressure | Humidity |
| 60,000 | 8.82 | 25-jan-08 | 15 | 755 mmHg | 46\% |
| 120,000 | 17.65 | 25-jan-08 | 15 | 755 mmHg | 46\% |
| 60,000 | 9.21 | 28-ian-08 | 18 | 739.8 mmHg | 56\% |
| 120,000 | 18.42 | 28-ian-08 | 18 | 739.8 mmHg | 56\% |
| 60,000 | 9.46 | 05-feb-08 | 21 | 740.4 mmHg | 82\% |
| 120,000 | 18.92 | 05-feb-08 | 21 | 740.4 mmHg | 82\% |

Model \#14-Gap 1c, Stagger -1.5c
Table C.14: Model \#14 Velocity, Test Dates, Temperature, Pressure, and Humidity


## Appendix D: Aerodynamic Data at Reynolds Number 60,000

Table D.1: Model \#1, $g=1 c, s=0 c$, Aerodynamic Data at $\operatorname{Re} 60,000$

| $a\left({ }^{\circ}\right)$ | C. ${ }_{\text {( }}$ - | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | $\alpha\left({ }^{\circ}\right)$ | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | $\left.\alpha^{( }{ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | C ${ }_{\text {d }}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -2.13 | -0.1395 | 0.0262 | -5.3147 | 6.87 | 0.4964 | 0.0789 | 6.2901 | 15.87 | 0.7225 | 0.2245 | 3.2180 |
| -1.88 | -0.1230 | 0.0264 | -4.6535 | 7.12 | 0.5097 | 0.0841 | 6.0581 | 16.12 | 0.7229 | 02264 | 3.1933 |
| -1.63 | -0.1068 | 0.0272 | -3.9207 | 7.37 | 0.5448 | 0.0898 | 6.0681 | 16.37 | 0.7263 | 0.2301 | 3.1558 |
| -1.38 | -0.0866 | 0.0258 | -3.3618 | 7.62 | 0.5368 | 0.0897 | 5.9817 | 16.62 | 0.7197 | 0.2320 | 3.1021 |
| -1.13 | -0.0681 | 0.0257 | -2.6474 | 7.87 | 0.5666 | 0.0958 | 5.9168 | 16.87 | 0.7368 | 0.2403 | 3.0654 |
| -0.88 | -0.0508 | 0.0254 | -2.0024 | 8.12 | 0.5793 | 0.0995 | 5.8213 | 17.12 | 0.7378 | 0.2459 | 3.0010 |
| -0.63 | -0.0418 | 0.0239 | -1.7530 | 8.37 | 0.6054 | 0.1081 | 5.5997 | 17.37 | 0.7316 | 0.2455 | 2.9802 |
| -0.38 | -0.0219 | 0.0261 | -0.8382 | 8.62 | 0.6313 | 0.1171 | 5.3895 | 17.62 | 0.7450 | 0.2531 | 2.9439 |
| -0.13 | -0.0090 | 0.0242 | -0.3698 | 8.87 | 0.6505 | 0.1157 | 5.6230 | 17.87 | 0.7420 | 0.2561 | 2.8977 |
| 0.12 | -0.0063 | 0.0246 | -0.2562 | 9.12 | 0.6560 | 0.1254 | 5.2311 | 18.12 | 0.7144 | 0.2537 | 2.8164 |
| 0.37 | 0.0172 | 0.0265 | 0.6475 | 9.37 | 0.6784 | 0.1344 | 5.0490 | 18.37 | 0.7316 | 0.2613 | 2.8001 |
| 062 | 0.0326 | 0.0246 | 1.3239 | 9.62 | 0.6816 | 0.1278 | 5.3343 | 18.62 | 0.7237 | 0.2616 | 2.7663 |
| 0.87 | 0.0483 | 0.0291 | 1.6580 | 9.87 | 0.6959 | 0.1367 | 5.0923 | 18.87 | 0.7300 | 0.2691 | 2.7127 |
| 1.12 | 0.0671 | 0.0260 | 2.5769 | 10.12 | 0.7067 | 0.1412 | 5.0045 | 19.12 | 0.7363 | 0.2733 | 2.6939 |
| 1.37 | 0.0864 | 0.0306 | 2.8280 | 10.37 | 0.7156 | 0.1453 | 4.9246 | 19.37 | 0.7341 | 0.2753 | 2.6668 |
| 1.62 | 0.0973 | 0.0275 | 3.5333 | 10.62 | 0.7343 | 0.1539 | 4.7712 | 19.62 | 0.7395 | 0.2809 | 2.6326 |
| 1.87 | 0.1210 | 0.0298 | 4.0611 | 10.87 | 0.7471 | 0.1555 | 4.8035 | 19.87 | 0.7558 | 0.2929 | 2.5803 |
| 2.12 | 0.1350 | 0.0317 | 4.2613 | 11.12 | 0.7516 | 0.1639 | 4.5846 | 20.12 | 0.7494 | 0.2926 | 2.5616 |
| 2.37 | 0.1683 | 0.0347 | 4.8490 | 11.37 | 0.7338 | 0.1664 | 4.4112 | 20.37 | 0.7504 | 0.2978 | 2.5196 |
| 2.62 | 0.1844 | 0.0358 | 5.1530 | 11.62 | 0.7561 | 0.1775 | 4.2595 | 20.62 | 0.7518 | 0.2997 | 2.5083 |
| 2.87 | 0.1922 | 0.0335 | 5.7276 | 11.87 | 0.7543 | 0.1765 | 4.2749 | 20.87 | 0.7599 | 0.3077 | 2.4696 |
| 3.12 | 0.2031 | 0.0388 | 5.2335 | 12.12 | 0.7688 | 0.1833 | 4.1947 | 21.12 | 0.7641 | 0.3136 | 2.4369 |
| 3.37 | 0.2031 | 0.0388 | 5.2335 | 12.37 | 0.7508 | 0.1785 | 4.2063 | 21.37 | 0.7521 | 0.3128 | 2.4042 |
| 3.62 | 0.2163 | 0.0392 | 5.5180 | 12.62 | 0.7515 | 0.1833 | 4.0992 | 21.62 | 0.7426 | 0.3133 | 2.3702 |
| 3.87 | 0.2341 | 0.0403 | 5.8138 | 12.87 | 0.7348 | 0.1833 | 4.0079 | 21.87 | 0.7488 | 0.3159 | 2.3703 |
| 4.12 | 0.2410 | 0.0417 | 5.7729 | 13.12 | 0.7429 | 0.1909 | 3.8921 | 22.12 | 0.7584 | 0.3245 | 2.3370 |
| 4.37 | 0.2612 | 0.0433 | 60388 | 13.37 | 0.7349 | 0.1905 | 3.8572 | 22.37 | 0.7644 | 0.3323 | 2.3002 |
| 4.62 | 0.3043 | 0.0451 | 6.7434 | 13.62 | 0.7294 | 0.1966 | 3.7103 | 22.62 | 0.7657 | 0.3369 | 2.2725 |
| 4.87 | 0.3149 | 0.0485 | 6.4865 | 13.87 | 0.7321 | 0.1993 | 3.6740 | 22.87 | 0.7739 | 0.3435 | 2.2533 |
| 5.12 | 03504 | 0.0526 | 6.6590 | 14.12 | 0.7225 | 0.1978 | 3.6523 | 23.12 | 0.7818 | 0.3494 | 2.2377 |
| 5.37 | 0.3672 | 0.0576 | 6.3806 | 14.37 | 0.7190 | 0.2030 | 3.5414 | 23.37 | 0.7871 | 0.3592 | 2.1911 |
| 562 | 0.3759 | 0.0612 | 6.1425 | 14.62 | 0.7218 | 0.2054 | 3.5142 | 23.62 | 0.7825 | 0.3626 | 2.1579 |
| 5.87 | 0.4050 | 0.0608 | 6.6580 | 14.87 | 0.7276 | 0.2080 | 3.4979 | 23.87 | 0.7512 | 0.3530 | 2.1278 |
| 6.12 | 0.4184 | 0.0650 | 6.4337 | 15.12 | 0.7339 | 0.2155 | 3.4060 | 24.12 | 0.7568 | 0.3583 | 2.1126 |
| 6.37 | 0.4301 | 0.0710 | 6.0596 | 15.37 | 0.7251 | 0.2181 | 3.3246 | 24.37 | 0.7874 | 0.3730 | 2.1110 |
| 6.62 | 0.4638 | 0.0747 | 6.2076 | 15.62 | 0.7303 | 0.2208 | 3.3080 | 24.62 | 0.7724 | 0.3730 | 2.0706 |
|  |  |  |  |  |  |  |  | 24.87 | 0.7794 | 0.3801 | 2.0505 |

Table D.2: Model \#2, $g=0.5 \mathrm{c}, \mathrm{s}=0 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | - ${ }^{\circ}{ }^{\circ}$ ) | CL (-) | $\mathrm{Cu}^{(-)}$ | L/D (-) | a ( ${ }^{\circ}$ ) | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -2.06 | -0.1094 | 0.0269 | -4.0691 | 6.94 | 0.3911 | 0.0654 | 5.9772 | 15.94 | 0.7206 | 0.2254 | 3.1965 |
| -1.81 | -0.0937 | 0.0252 | -3.7188 | 7.19 | 0.4083 | 0.0718 | 5.6858 | 16.19 | 0.7042 | 0.2258 | 3.1187 |
| $-1.56$ | -0.0838 | 0.0237 | -3.5351 | 7.44 | 0.4206 | 0.0747 | 5.6278 | 16.44 | 0.6974 | 0.2268 | 3.0748 |
| -1.31 | -0.0661 | 0.0232 | -2.8468 | 7.69 | 0.4336 | 0.0760 | 5.7085 | 16.69 | 0.6982 | 0.2299 | 3.0367 |
| -1.06 | -0.0451 | 0.0232 | -1.9394 | 7.94 | 0.4459 | 0.0808 | 5.5198 | 16.94 | 0.7178 | 0.2411 | 2.9770 |
| -0.81 | -0.0393 | 0.0225 | -1.7498 | 8.19 | 0.4690 | 0.0848 | 5.5320 | 17.19 | 0.7001 | 0.2386 | 2.9346 |
| -0.56 | -0.0290 | 0.0222 | -1.3076 | 8.44 | 0.4782 | 0.0889 | 5.3764 | 17.44 | 0.7255 | 0.2496 | 2.9070 |
| -0.31 | -0.0123 | 0.0232 | -0.5308 | 8.69 | 0.4872 | 0.0932 | 5.2296 | 17.69 | 0.7214 | 0.2490 | 2.8968 |
| -0.06 | -0.0052 | 0.0235 | -0.2226 | 8.94 | 0.5089 | 0.0983 | 5.1765 | 17.94 | 0.7245 | 0.2571 | 2.8180 |
| 0.19 | 00108 | 0.0230 | 0.4711 | 9.19 | 0.5171 | 0.1013 | 5.1057 | 18.19 | 0.7280 | 0.2651 | 2.7465 |
| 0.44 | 0.0273 | 0.0233 | 1.1723 | 9.44 | 0.5335 | 0.1058 | 5.0405 | 18.44 | 0.7197 | 0.2631 | 2.7354 |
| 0.69 | 0.0382 | 0.0236 | 1.6231 | 9.69 | 0.5348 | 0.1084 | 4.9358 | 18.69 | 0.7297 | 0.2710 | 2.6928 |
| 0.94 | 0.0443 | 0.0228 | 1.9455 | 9.94 | 0.5502 | 0.1146 | 4.8027 | 18.94 | 0.7357 | 0.2749 | 2.6765 |
| 1.19 | 0.0641 | 0.0227 | 2.8189 | 10.19 | 0.5596 | 0.1189 | 4.7071 | 19.19 | 0.7513 | 0.2810 | 2.6732 |
| 1.44 | 0.0776 | 0.0235 | 3.3009 | 10.44 | 0.5663 | 0.1223 | 46291 | 19.44 | 0.7479 | 0.2866 | 2.6098 |
| 1.69 | 0.0873 | 0.0245 | 3.5674 | 10.69 | 0.5718 | 0.1260 | 4.5368 | 19.69 | 0.7394 | 0.2855 | 2.5898 |
| 1.94 | 0.0996 | 0.0257 | 38827 | 10.94 | 0.5854 | 0.1322 | 4.4286 | 19.94 | 0.7326 | 0.2841 | 2.5783 |
| 2.19 | 0.1125 | 0.0266 | 4.2201 | 11.19 | 0.5888 | 0.1353 | 4.3519 | 20.19 | 0.7381 | 0.2885 | 2.5588 |
| 244 | 0.1252 | 0.0281 | 4.4502 | 11.44 | 0.5913 | 0.1388 | 4.2606 | 20.44 | 0.7258 | 0.2974 | 2.4404 |
| 2.69 | 0.1407 | 0.0286 | 4.9205 | 11.69 | 0.6053 | 0.1436 | 4.2155 | 20.69 | 0.7315 | 0.2961 | 2.4706 |
| 2.94 | 0.1623 | 0.0306 | 5.3068 | 11.94 | 0.6060 | 0.1461 | 4.1482 | 20.94 | 0.7245 | 0.3047 | 2.3777 |
| 3.19 | 0.1682 | 0.0328 | 5.1226 | 12.19 | 0.6081 | 0.1502 | 4.0473 | 21.19 | 0.7390 | 0.3108 | 2.3779 |
| 3.44 | 0.1682 | 0.0328 | 5.1226 | 12.44 | 0.6117 | 0.1552 | 3.9408 | 21.44 | 0.7264 | 0.3054 | 2.3785 |
| 3.69 | 0.1898 | 0.0322 | 58860 | 12.69 | 0.6254 | 0.1612 | 3.8798 | 21.69 | 0.7334 | 0.3127 | 2.3453 |
| 3.94 | 0.1993 | 0.0343 | 5.8184 | 12.94 | 0.6253 | 0.1623 | 3.8533 | 21.94 | 0.7325 | 0.3171 | 2.3098 |
| 4.19 | 0.2103 | 0.0367 | 5.7220 | 13.19 | 0.6265 | 0.1672 | 3.7471 | 22.19 | 0.7329 | 0.3201 | 2.2897 |
| 4.44 | 0.2235 | 0.0386 | 5.7879 | 13.44 | 0.6349 | 0.1726 | 3.6795 | 22.44 | 0.7158 | 0.3147 | 2.2744 |
| 4.69 | 0.2448 | 0.0396 | 6.1733 | 13.69 | 0.6446 | 0.1775 | 3.6310 | 22.69 | 0.7402 | 0.3279 | 2.2576 |
| 4.94 | 0.2552 | 0.0411 | 6.2066 | 13.94 | 0.6467 | 0.1803 | 3.5861 | 22.94 | 0.7393 | 0.3437 | 2.1513 |
| 5.19 | 0.2874 | 0.0468 | 6.1416 | 14.19 | 0.6693 | 0.1885 | 3.5508 | 23.19 | 0.7289 | 0.3433 | 2.1235 |
| 5.44 | 0.2869 | 0.0487 | 5.8923 | 1444 | 0.6662 | 0.1925 | 3.4613 | 23.44 | 0.7248 | 0.3368 | 2.1520 |
| 5.69 | 0.3145 | 0.0512 | 6.1380 | 14.69 | 0.6731 | 0.1970 | 3.4166 | 23.69 | 0.7228 | 0.3464 | 2.0868 |
| 5.94 | 03314 | 0.0549 | 6.0395 | 14.94 | 0.6857 | 0.2023 | 3.3899 | 23.94 | 0.7223 | 0.3513 | 2.0558 |
| 6.19 | 0.3427 | 0.0559 | 6.1340 | 15.19 | 0.6884 | 0.2092 | 3.2911 | 24.19 | 0.7224 | 0.3422 | 2.1107 |
| 6.44 | 0.3586 | 0.0597 | 6.0053 | 15.44 | 0.7023 | 0.2175 | 3.2296 | 24.44 | 0.7279 | 0.3503 | 2.0779 |
| 6.69 | 0.3751 | 0.0640 | 5.8609 | 15.69 | 0.7125 | 0.2235 | 3.1870 | 24.69 | 0.7206 | 0.3558 | 2.0255 |
|  |  |  |  |  |  |  |  | 24.94 | 0.7254 | 0.3592 | 2.0196 |

Table D.3: Model \#3, $y_{s}=0.5 c_{y} s=0.5 \mathrm{c}$, Aerodynamic Data at Re 60,000

| - ${ }^{( }{ }^{\circ}$ | $\mathrm{CL}_{\mathrm{L}}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.93 | -0.1175 | 0.0257 | -4.5653 |
| -1.68 | -0.0892 | 0.0265 | -3.3603 |
| -1.43 | -0.0766 | 0.0269 | -2.8490 |
| -1.18 | -0.0624 | 0.0257 | -2.4278 |
| -0.93 | -0.0432 | 0.0269 | -1.6069 |
| -0.68 | -0.0420 | 0.0269 | -1.5621 |
| -0.43 | -0.0222 | 0.0238 | -0.9347 |
| -0.18 | -0.0139 | 0.0242 | -0.5736 |
| 0.07 | 0.0084 | 0.0258 | 0.3273 |
| 0.32 | 0.0338 | 0.0257 | 1.3141 |
| 0.57 | 0.0402 | 0.0254 | 1.5849 |
| 0.82 | 0.0513 | 0.0254 | 2.0226 |
| 1.07 | 0.0628 | 0.0250 | 2.5150 |
| 1.32 | 0.0776 | 0.0254 | 30528 |
| 1.57 | 0.0893 | 0.0254 | 3.5202 |
| 1.82 | 0.1116 | 0.0257 | 4.3389 |
| 2.07 | 0.1250 | 0.0269 | 4.6471 |
| 2.32 | 0.1500 | 0.0311 | 4.8166 |
| 2.57 | 0.1621 | 0.0311 | 5.2079 |
| 2.82 | 0.1696 | 0.0338 | 5.0201 |
| 3.07 | 0.1770 | 0.0315 | 5.6280 |
| 3.32 | 0.1956 | 0.0326 | 5.9963 |
| 3.57 | 0.1956 | 0.0326 | 5.9963 |
| 3.82 | 0.2214 | 0.0368 | 6.0185 |
| 4.07 | 0.2444 | 0.0369 | 6.6313 |
| 4.32 | 0.2588 | 0.0407 | 6.3641 |
| 4.57 | 0.2830 | 0.0426 | 6.6481 |
| 4.82 | 0.2784 | 0.0430 | 6.4786 |
| 5.07 | 0.3112 | 0.0468 | 6.6540 |
| 5.32 | 0.3479 | 0.0517 | 6.7218 |
| 5.57 | 0.3581 | 0.0545 | 6.5763 |
| 5.82 | 0.3529 | 0.0544 | 6.4893 |
| 6.07 | 0.3822 | 0.0605 | 6.3124 |
| 6.32 | 0.4117 | 0.0632 | 6.5181 |
| 6.57 | 0.4293 | 0.0697 | 6.1588 |
| 6.82 | 0.4453 | 0.0689 | 6.4601 |


| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.07 | 0.4686 | 0.0720 | 6.5051 |
| 7.32 | 0.4853 | 0.0800 | 6.0672 |
| 7.57 | 0.5005 | 0.0839 | 5.9685 |
| 7.82 | 0.5034 | 0.0862 | 5.8427 |
| 8.07 | 0.5206 | 0.0900 | 5.7875 |
| 8.32 | 0.5390 | 0.0942 | 5.7233 |
| 8.57 | 0.5426 | 0.0972 | 5.5821 |
| 8.82 | 0.5642 | 0.1030 | 5.4774 |
| 9.07 | 0.5840 | 0.1107 | 5.2768 |
| 9.32 | 0.5899 | 0.1117 | 5.2793 |
| 9.57 | 0.6010 | 0.1152 | 5.2159 |
| 9.82 | 0.6227 | 0.1232 | 5.0533 |
| 10.07 | 0.6390 | 0.1286 | 4.9681 |
| 1032 | 0.6481 | 0.1317 | 4.9205 |
| 10.57 | 0.6532 | 0.1344 | 4.8614 |
| 10.82 | 0.6701 | 0.1413 | 4.7437 |
| 11.07 | 0.6807 | 0.1428 | 4.7678 |
| 11.32 | 0.6840 | 0.1481 | 4.6183 |
| 11.57 | 0.7058 | 0.1546 | 4.5659 |
| 11.82 | 0.7148 | 0.1585 | 4.5102 |
| 12.07 | 0.7198 | 0.1635 | 4.4037 |
| 12.32 | 0.7253 | 0.1665 | 4.3569 |
| 12.57 | 0.7332 | 0.1703 | 4.3064 |
| 12.82 | 0.7335 | 0.1760 | 4.1673 |
| 13.07 | 0.7391 | 0.1786 | 4.1372 |
| 13.32 | 0.7499 | 0.1840 | 4.0755 |
| 13.57 | 0.7652 | 0.1928 | 3.9680 |
| 13.82 | 0.7708 | 0.1993 | 3.8671 |
| 14.07 | 0.7766 | 0.2005 | 3.8741 |
| 14.32 | 0.7880 | 0.2112 | 3.7316 |
| 14.57 | 0.7857 | 0.2111 | 3.7218 |
| 14.82 | 0.7823 | 0.2142 | 3.6516 |
| 15.07 | 0.7967 | 0.2207 | 3.6098 |
| 15.32 | 0.7890 | 0.2207 | 3.5752 |
| 15.57 | 0.7969 | 0.2276 | 3.5011 |
| 15.82 | 0.8172 | 0.2368 | 3.4508 |


| a $\left(^{\circ}\right.$ ) | $C_{L}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.07 | 0.8120 | 0.2409 | 3.3702 |
| 16.32 | 0.8025 | 0.2393 | 3.3529 |
| 16.57 | 0.8025 | 0.2416 | 3.3216 |
| 16.82 | 0.8053 | 0.2474 | 3.2553 |
| 17.07 | 0.8303 | 0.2574 | 3.2258 |
| 17.32 | 0.8260 | 0.2604 | 3.1724 |
| 17.57 | 0.8472 | 0.2720 | 3.1152 |
| 17.82 | 0.8308 | 0.2676 | 3.1044 |
| 18.07 | 0.8224 | 0.2689 | 3.0581 |
| 18.32 | 0.8170 | 0.2733 | 2.9891 |
| 18.57 | 0.8349 | 0.2814 | 2.9667 |
| 18.82 | 0.8594 | 0.2910 | 2.9528 |
| 19.07 | 0.8538 | 0.2956 | 2.8878 |
| 19.32 | 0.8480 | 0.2971 | 2.8541 |
| 19.57 | 0.8616 | 0.3063 | 2.8132 |
| 19.82 | 0.8685 | 0.3144 | 2.7621 |
| 20.07 | 0.8655 | 0.3160 | 2.7392 |
| 20.32 | 0.8552 | 0.3146 | 2.7185 |
| 20.57 | 0.8629 | 0.3234 | 2.6679 |
| 2082 | 0.8560 | 0.3261 | 2.6249 |
| 21.07 | 0.8600 | 0.3288 | 2.6154 |
| 21.32 | 0.8724 | 0.3372 | 2.5868 |
| 21.57 | 0.8883 | 0.3476 | 2.5551 |
| 21.82 | 0.8889 | 0.3522 | 2.5239 |
| 22.07 | 0.8899 | 0.3569 | 2.4936 |
| 22.32 | 0.8990 | 0.3634 | 2.4735 |
| 22.57 | 0.8767 | 0.3590 | 2.4424 |
| 22.82 | 0.8848 | 0.3668 | 2.4124 |
| 23.07 | 0.8778 | 0.3662 | 2.3972 |
| 23.32 | 0.8981 | 0.3791 | 2.3688 |
| 23.57 | 0.8820 | 0.3789 | 2.3278 |
| 23.82 | 0.9241 | 0.4008 | 2.3057 |
| 24.07 | 0.9212 | 0.4022 | 2.2905 |
| 24.32 | 0.8756 | 0.3858 | 2.2694 |
| 24.57 | 0.9103 | 0.4075 | 2.2341 |
| 24.82 | 0.8768 | 0.3969 | 2.2094 |
| 25.07 | 0.8826 | 0.4025 | 2.1927 |

Table D.4: Model \#4, $\mathrm{g}=0.5 \mathrm{c}, \mathrm{s}=1 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| $\alpha\left(^{\circ}\right)$ | $\mathrm{CL}_{\mathrm{L}}^{(-)}$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | a $\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -2.03 | -0.1223 | 0.0267 | -4.5880 | 6.97 | 0.4948 | 0.0707 | 6.9940 |
| -1.78 | -0.1164 | 0.0212 | -5.4945 | 7.22 | 0.5235 | 0.0793 | 6.5997 |
| -1.53 | -0.1034 | 0.0212 | -4.8713 | 7.47 | 0.5386 | 0.0812 | 6.6346 |
| -1.28 | -0.0844 | 0.0208 | -4.0489 | 7.72 | 0.5641 | 0.0867 | 6.5087 |
| -1.03 | -0.0611 | 0.0216 | -2.8240 | 7.97 | 0.5898 | 0.0955 | 6.1762 |
| -0.78 | -0.0590 | 0.0228 | -2.5897 | 8.22 | 0.6114 | 0.0960 | 6.3715 |
| -0.53 | -0.0414 | 0.0181 | -2.2882 | 847 | 0.6181 | 0.1002 | 6.1683 |
| -0.28 | -0.0204 | 0.0197 | -1.0325 | 8.72 | 0.6371 | 0.1040 | 6.1257 |
| -0.03 | -0.0164 | 0.0197 | -0.8337 | 8.97 | 0.6474 | 0.1079 | 6.0015 |
| 0.22 | 0.0111 | 0.0209 | 0.5307 | 9.22 | 0.6615 | 0.1129 | 5.8606 |
| 0.47 | 0.0187 | 0.0224 | 0.8358 | 9.47 | 0.6690 | 0.1148 | 5.8283 |
| 0.72 | 0.0328 | 0.0232 | 1.4142 | 9.72 | 0.6885 | 0.1217 | 5.6557 |
| 0.97 | 00482 | 0.0205 | 2.3547 | 9.97 | 0.7080 | 0.1295 | 5.4673 |
| 1.22 | 0.0759 | 0.0205 | 3.7079 | 10.22 | 0.7239 | 0.1310 | 5.5252 |
| 1.47 | 0.0922 | 0.0193 | 4.7760 | 10.47 | 0.7382 | 0.1396 | 5.2895 |
| 1.72 | 0.1013 | 0.0209 | 4.8575 | 10.72 | 0.7540 | 0.1445 | 5.2175 |
| 1.97 | 0.1296 | 0.0243 | 5.3232 | 10.97 | 0.7664 | 0.1511 | 5.0721 |
| 2.22 | 0.1590 | 0.0240 | 6.6257 | 11.22 | 0.7935 | 0.1569 | 5.0567 |
| 2.47 | 0.1649 | 0.0267 | 6.1814 | 11.47 | 0.8043 | 0.1650 | 4.8730 |
| 2.72 | 0.1655 | 0.0263 | 6.2963 | 11.72 | 0.8225 | 0.1701 | 4.8358 |
| 2.97 | 0.1747 | 0.0259 | 6.7523 | 11.97 | 0.8407 | 0.1767 | 4.7584 |
| 3.22 | 0.2060 | 0.0279 | 7.3962 | 12.22 | 0.8593 | 0.1844 | 4.6601 |
| 3.47 | 0.2060 | 0.0279 | 7.3962 | 12.47 | 0.8729 | 0.1929 | 4.5256 |
| 3.72 | 0.2231 | 0.0282 | 7.9136 | 12.72 | 0.8836 | 0.1952 | 4.5276 |
| 3.97 | 0.2537 | 0.0352 | 7.2112 | 12.97 | 0.8868 | 0.2021 | 4.3882 |
| 4.22 | 0.2622 | 0.0340 | 7.7047 | 13.22 | 0.9032 | 0.2106 | 4.2899 |
| 447 | 0.2750 | 0.0371 | 7.4161 | 13.47 | 0.9079 | 0.2176 | 4.1730 |
| 4.72 | 0.2948 | 0.0375 | 7.8636 | 13.72 | 0.9176 | 0.2175 | 4.2185 |
| 4.97 | 0.3160 | 0.0413 | 7.6449 | 13.97 | 0.9201 | 0.2264 | 4.0642 |
| 5.22 | 0.3450 | 0.0425 | 8.1246 | 14.22 | 0.9286 | 0.2303 | 4.0314 |
| 5.47 | 0.3722 | 0.0476 | 7.8227 | 14.47 | 0.9485 | 0.2403 | 3.9468 |
| 5.72 | 0.3829 | 0.0506 | 7.5611 | 14.72 | 0.9449 | 0.2422 | 3.9007 |
| 5.97 | 0.3968 | 0.0526 | 7.5450 | 14.97 | 0.9433 | 0.2473 | 3.8144 |
| 6.22 | 0.4368 | 0.0584 | 7.4848 | 15.22 | 0.9633 | 0.2566 | 3.7545 |
| 6.47 | 0.4513 | 0.0627 | 7.2029 | 15.47 | 0.9753 | 0.2604 | 3.7459 |
| 6.72 | 0.4775 | 0.0673 | 7.0958 | 15.72 | 0.9785 | 0.2700 | 3.6245 |


| $a\left({ }^{\circ}\right)$ | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.97 | 0.9851 | 0.2758 | 3.5712 |
| 16.22 | 0.9811 | 0.2774 | 3.5367 |
| 16.47 | 0.9886 | 0.2851 | 3.4676 |
| 16.72 | 0.9996 | 0.2890 | 3.4588 |
| 16.97 | 1.0075 | 0.2959 | 3.4045 |
| 17.22 | 1.0006 | 0.2998 | 3.3372 |
| 17.47 | 0.9989 | 0.3064 | 3.2603 |
| 17.72 | 1.0236 | 0.3172 | 3.2274 |
| 17.97 | 1.0183 | 0.3199 | 3.1832 |
| 18.22 | 1.0010 | 0.3211 | 3.1175 |
| 18.47 | 1.0014 | 0.3252 | 3.0795 |
| 18.72 | 1.0145 | 0.3330 | 3.0465 |
| 18.97 | 10379 | 0.3465 | 2.9954 |
| 19.22 | 1.0235 | 0.3457 | 2.9608 |
| 19.47 | 1.0275 | 0.3535 | 2.9067 |
| 19.72 | 1.0315 | 0.3578 | 2.8830 |
| 19.97 | 1.0276 | 0.3627 | 2.8329 |
| 20.22 | 1.0283 | 0.3650 | 2.8170 |
| 20.47 | 1.0217 | 0.3685 | 2.7726 |
| 20.72 | 1.0479 | 0.3812 | 2.7488 |
| 20.97 | 1.0444 | 0.3878 | 2.6929 |
| 21.22 | 1.0581 | 0.3983 | 2.6563 |
| 21.47 | 1.0560 | 0.4009 | 2.6339 |
| 21.72 | 1.0680 | 0.4106 | 2.6011 |
| 21.97 | 1.0645 | 0.4126 | 2.5797 |
| 22.22 | 1.0807 | 0.4248 | 2.5437 |
| 22.47 | 1.0689 | 0.4238 | 2.5223 |
| 22.72 | 1.0697 | 0.4287 | 24950 |
| 22.97 | 1.0682 | 0.4352 | 2.4546 |
| 23.22 | 1.0731 | 0.4431 | 2.4219 |
| 23.47 | 1.0583 | 0.4410 | 2.3997 |
| 23.72 | 1.0777 | 0.4500 | 2.3950 |
| 23.97 | 1.0699 | 0.4585 | 2.3336 |
| 24.22 | 1.0681 | 0.4612 | 2.3160 |
| 24.47 | 1.0711 | 0.4678 | 2.2895 |
| 24.72 | 1.0669 | 0.4716 | 2.2620 |
| 24.97 | 1.0831 | 0.4832 | 2.2417 |

Table D.5: Model \# 5, $\mathrm{g}=1 \mathrm{c}, \mathrm{s}=0.5 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{CD}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.89 | -0.1382 | 0.0449 | -3.0740 |
| -1.64 | -0.1296 | 0.0462 | -2.8064 |
| -1.39 | -0.1107 | 0.0452 | -2.4491 |
| -1.14 | -0.0928 | 0.0440 | -2.1062 |
| -0.89 | $-0.0720$ | 0.0449 | -1.6051 |
| -0.64 | -0.0464 | 0.0438 | -1.0596 |
| -0.39 | -0.0343 | 0.0413 | -0.8294 |
| -0.14 | -0.0174 | 0.0412 | -0.4217 |
| 0.11 | 0.0046 | 0.0418 | 0.1103 |
| 0.36 | 0.0233 | 0.0420 | 0.5550 |
| 0.61 | 0.0441 | 0.0435 | 1.0133 |
| 0.86 | 0.0673 | 00406 | 1.6566 |
| 1.11 | 0.0730 | 0.0426 | 1.7152 |
| 1.36 | 0.0933 | 0.0448 | 2.0814 |
| 1.61 | 0.1144 | 0.0478 | 2.3919 |
| 1.86 | 0.1374 | 0.0476 | 2.8841 |
| 2.11 | 0.1313 | 0.0431 | 3.0449 |
| 2.36 | 0.1645 | 0.0464 | 3.5468 |
| 2.61 | 0.1916 | 0.0502 | 3.8148 |
| 2.86 | 0.2095 | 0.0509 | 4.1146 |
| 3.11 | 0.2259 | 0.0516 | 4.3785 |
| 3.36 | 0.2404 | 0.0547 | 4.3922 |
| 3.61 | 0.2404 | 0.0547 | 4.3922 |
| 3.86 | 0.2575 | 0.0564 | 4.5672 |
| 4.11 | 0.2748 | 0.0562 | 4.8896 |
| 4.36 | 0.2935 | 0.0603 | 4.8654 |
| 4.61 | 0.3275 | 0.0627 | 5.2265 |
| 4.86 | 0.3343 | 0.0643 | 5.1974 |
| 5.11 | 0.3495 | 0.0646 | 5.4093 |
| 5.36 | 0.3874 | 0.0710 | 5.4594 |
| 5.61 | 0.4144 | 0.0749 | 5.5310 |
| 5.86 | 0.4395 | 0.0784 | 5.6058 |
| 6.11 | 0.4611 | 0.0820 | 5.6217 |
| 6.36 | 0.4761 | 0.0848 | 5.6155 |
| 6.61 | 0.4962 | 0.0888 | 5.5858 |
| 6.86 | 0.5247 | 0.0910 | 5.7676 |


| a $\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.11 | 0.5442 | 0.0966 | 5.6314 |
| 7.36 | 0.5524 | 00988 | 5.5891 |
| 7.61 | 0.5740 | 0.1041 | 5.5140 |
| 7.86 | 0.6079 | 0.1110 | 5.4784 |
| 8.11 | 0.6323 | 0.1152 | 5.4868 |
| 8.36 | 0.6345 | 0.1171 | 5.4190 |
| 8.61 | 0.6591 | 0.1247 | 5.2843 |
| 8.86 | 0.6647 | 0.1284 | 5.1771 |
| 9.11 | 0.6900 | 0.1351 | 5.1060 |
| 9.36 | 0.7033 | 0.1378 | 5.1046 |
| 9.61 | 0.7212 | 0.1410 | 5.1142 |
| 9.86 | 0.7473 | 0.1520 | 4.9173 |
| 10.11 | 0.7651 | 0.1584 | 4.8307 |
| 10.36 | 0.7689 | 0.1597 | 4.8147 |
| 10.61 | 0.7782 | 0.1669 | 4.6634 |
| 10.86 | 0.7938 | 0.1687 | 4.7054 |
| 11.11 | 0.8056 | 0.1759 | 4.5811 |
| 11.36 | 0.8213 | 0.1837 | 4.4707 |
| 11.61 | 0.8308 | 0.1879 | 4.4203 |
| 11.86 | 0.8276 | 0.1926 | 4.2968 |
| 12.11 | 0.8421 | 0.1985 | 4.2420 |
| 12.36 | 0.8498 | 0.2037 | 4.1715 |
| 12.61 | 0.8624 | 0.2076 | 4.1533 |
| 12.86 | 0.8632 | 0.2100 | 4.1103 |
| 13.11 | 0.8665 | 0.2167 | 3.9980 |
| 13.36 | 0.8709 | 0.2201 | 3.9566 |
| 13.61 | 0.8828 | 0.2268 | 3.8923 |
| 13.86 | 0.8923 | 0.2342 | 3.8104 |
| 14.11 | 0.8882 | 0.2368 | 3.7514 |
| 14.36 | 0.8983 | 0.2446 | 3.6720 |
| 14.61 | 0.8992 | 0.2479 | 3.6265 |
| 14.86 | 0.8990 | 0.2508 | 3.5845 |
| 15.11 | 0.9058 | 0.2557 | 3.5416 |
| 15.36 | 0.9026 | 0.2594 | 3.4794 |
| 15.61 | 0.9028 | 0.2643 | 3.4155 |
| 15.86 | 0.9003 | 0.2690 | 3.3465 |


| a ( ${ }^{\circ}$ ) | CL (-) | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.11 | 0.9069 | 0.2721 | 3.3327 |
| 16.36 | 0.9046 | 0.2786 | 3.2472 |
| 16.61 | 0.9039 | 0.2806 | 3.2214 |
| 16.86 | 0.9058 | 0.2840 | 3.1899 |
| 17.11 | 0.9150 | 0.2923 | 3.1306 |
| 17.36 | 0.9130 | 0.2969 | 3.0752 |
| 17.61 | 0.9096 | 0.2966 | 3.0671 |
| 17.86 | 0.9196 | 0.3058 | 3.0073 |
| 18.11 | 0.9180 | 0.3069 | 2.9912 |
| 18.36 | 0.9125 | 0.3094 | 2.9489 |
| 18.61 | 0.9300 | 0.3203 | 2.9036 |
| 18.86 | 0.9217 | 0.3224 | 2.8587 |
| 19.11 | 0.9184 | 0.3204 | 2.8664 |
| 19.36 | 0.9177 | 0.3275 | 28020 |
| 19.61 | 0.9295 | 0.3357 | 2.7692 |
| 19.86 | 0.9294 | 0.3386 | 2.7453 |
| 20.11 | 0.9348 | 0.3482 | 2.6849 |
| 20.36 | 0.9310 | 0.3470 | 2.6826 |
| 20.61 | 0.9459 | 0.3579 | 2.6427 |
| 20.86 | 0.9315 | 0.3619 | 2.5742 |
| 21.11 | 0.9345 | 0.3629 | 2.5755 |
| 21.36 | 0.9431 | 0.3696 | 2.5519 |
| 21.61 | 0.9321 | 0.3761 | 2.4784 |
| 21.86 | 0.9503 | 0.3847 | 2.4700 |
| 22.11 | 0.9450 | 0.3837 | 2.4628 |
| 22.36 | 0.9620 | 0.3949 | 2.4361 |
| 22.61 | 0.9476 | 0.3960 | 2.3931 |
| 22.86 | 0.9524 | 0.3975 | 2.3959 |
| 23.11 | 0.9560 | 0.4074 | 2.3464 |
| 23.36 | 0.9594 | 0.4112 | 2.3335 |
| 23.61 | 0.9542 | 0.4134 | 2.3081 |
| 23.86 | 0.9659 | 0.4225 | 2.2861 |
| 24.11 | 0.9714 | 0.4303 | 2.2577 |
| 24.36 | 0.9809 | 0.4409 | 2.2247 |
| 24.61 | 0.9728 | 0.4387 | 2.2173 |
| 24.86 | 0.9709 | 0.4471 | 2.1718 |
| 25.11 | 0.9809 | 0.4526 | 2.1674 |

Table D.6: Model \#6, $\mathrm{g}=1 \mathrm{c}, \mathrm{s}=1 \mathrm{c}$, Aerodynamic Data at Re 60,000

| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{Cu}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.03 | -0.1580 | 0.0246 | -6.4227 |
| -1.78 | -0.1316 | 0.0276 | -4.7698 |
| -1.53 | -0.1272 | 0.0250 | -5.0968 |
| -1.28 | -0.1054 | 0.0250 | -4.2244 |
| -1.03 | -0.1055 | 0.0250 | -4.2166 |
| -0.78 | -0.0673 | 0.0220 | -3.0624 |
| -0.53 | -0.0414 | 0.0262 | -1.5837 |
| -0.28 | -0.0261 | 0.0223 | -1.1712 |
| -0.03 | -0.0034 | 0.0197 | -0.1738 |
| 0.22 | 0.0211 | 0.0212 | 0.9949 |
| 0.47 | 0.0303 | 0.0209 | 1.4536 |
| 0.72 | 0.0439 | 0.0197 | 2.2251 |
| 0.97 | 0.0684 | 0.0220 | 3.1070 |
| 1.22 | 0.0609 | 0.0208 | 2.9291 |
| 1.47 | 0.0857 | 0.0216 | 3.9684 |
| 1.72 | 0.1087 | 0.0220 | 4.9419 |
| 1.97 | 0.1274 | 0.0227 | 5.6107 |
| 2.22 | 0.1425 | 0.0280 | 5.0966 |
| 2.47 | 0.1552 | 0.0296 | 5.2496 |
| 2.72 | 0.1810 | 0.0299 | 6.0592 |
| 2.97 | 0.2085 | 0.0336 | 6.2020 |
| 3.22 | 0.2255 | 0.0340 | 6.6296 |
| 3.47 | 0.2255 | 0.0340 | 6.6296 |
| 3.72 | 0.2420 | 0.0367 | 6.5993 |
| 3.97 | 0.2663 | 0.0378 | 7.0456 |
| 4.22 | 0.2784 | 0.0404 | 6.8889 |
| 4.47 | 0.2967 | 0.0404 | 7.3460 |
| 4.72 | 0.3170 | 0.0434 | 7.3043 |
| 4.97 | 0.3205 | 0.0467 | 6.8603 |
| 5.22 | 0.3686 | 0.0512 | 7.1952 |
| 5.47 | 0.3752 | 0.0524 | 7.1626 |
| 5.72 | 0.4046 | 0.0516 | 7.8393 |
| 5.97 | 0.4347 | 0.0622 | 6.9911 |
| 6.22 | 0.4564 | 0.0651 | 7.0109 |
| 6.47 | 0.4794 | 0.0659 | 7.2730 |
| 6.72 | 0.4971 | 0.0685 | 7.2533 |


| $\square\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{0}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.97 | 0.5180 | 0.0756 | 6.8499 |
| 7.22 | 0.5411 | 0.0779 | 6.9430 |
| 7.47 | 0.5708 | 0.0843 | 6.7710 |
| 7.72 | 0.5912 | 0.0892 | 6.6306 |
| 7.97 | 0.6002 | 0.0967 | 6.2076 |
| 8.22 | 0.6200 | 0.1004 | 6.1759 |
| 8.47 | 0.6320 | 0.1041 | 6.0699 |
| 8.72 | 0.6380 | 0.1094 | 5.8309 |
| 8.97 | 0.6744 | 0.1143 | 5.8984 |
| 9.22 | 0.7072 | 0.1233 | 5.7356 |
| 9.47 | 0.7113 | 01286 | 5.5317 |
| 9.72 | 0.7360 | 0.1315 | 5.5960 |
| 9.97 | 0.7393 | 0.1382 | 5.3481 |
| 10.22 | 0.7511 | 0.1402 | 5.3580 |
| 10.47 | 0.7677 | 0.1466 | 5.2363 |
| 10.72 | 0.7909 | 0.1518 | 5.2099 |
| 10.97 | 0.7996 | 0.1593 | 5.0180 |
| 11.22 | 0.8105 | 0.1668 | 4.8587 |
| 11.47 | 0.8231 | 0.1709 | 4.8153 |
| 11.72 | 0.8252 | 0.1752 | 4.7112 |
| 11.97 | 0.8294 | 0.1792 | 4.6281 |
| 12.22 | 0.8295 | 0.1830 | 4.5335 |
| 12.47 | 0.8476 | 0.1875 | 4.5207 |
| 12.72 | 0.8480 | 0.1924 | 4.4073 |
| 12.97 | 0.8517 | 0.1976 | 4.3104 |
| 13.22 | 0.8452 | 0.2040 | 4.1426 |
| 13.47 | 0.8628 | 0.2093 | 4.1228 |
| 13.72 | 0.8717 | 0.2126 | 4.0998 |
| 13.97 | 0.8723 | 0.2206 | 3.9550 |
| 14.22 | 0.8703 | 0.2239 | 3.8864 |
| 1447 | 0.8730 | 0.2314 | 3.7725 |
| 14.72 | 0.8776 | 0.2313 | 3.7942 |
| 14.97 | 0.8818 | 0.2393 | 3.6843 |
| 15.22 | 0.8850 | 0.2438 | 3.6296 |
| 15.47 | 0.8863 | 0.2463 | 3.5980 |
| 15.72 | 0.8990 | 0.2516 | 3.5728 |


| $\alpha\left(^{\circ}\right.$ ) | $\mathrm{C}_{2}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.97 | 0.8851 | 0.2546 | 3.4763 |
| 16.22 | 0.8882 | 0.2550 | 3.4829 |
| 16.47 | 0.8878 | 0.2659 | 3.3390 |
| 16.72 | 0.8901 | 0.2715 | 3.2783 |
| 16.97 | 0.8953 | 0.2726 | 3.2844 |
| 17.22 | 0.8828 | 0.2741 | 3.2208 |
| 17.47 | 0.8816 | 0.2801 | 3.1470 |
| 17.72 | 0.8811 | 0.2821 | 3.1235 |
| 17.97 | 0.8898 | 0.2918 | 3.0492 |
| 18.22 | 0.9104 | 0.3008 | 3.0265 |
| 18.47 | 0.9177 | 0.3045 | 3.0136 |
| 18.72 | 0.9019 | 0.3083 | 2.9252 |
| 18.97 | 0.9211 | 0.3188 | 2.8893 |
| 19.22 | 0.8960 | 0.3151 | 2.8436 |
| 19.47 | 0.8974 | 0.3187 | 2.8155 |
| 19.72 | 0.9224 | 0.3305 | 2.7911 |
| 19.97 | 0.9210 | 0.3314 | 2.7787 |
| 20.22 | 0.9291 | 0.3383 | 2.7466 |
| 20.47 | 0.9306 | 0.3428 | 2.7146 |
| 20.72 | 0.9147 | 0.3388 | 2.6997 |
| 20.97 | 0.9272 | 0.3485 | 2.6607 |
| 21.22 | 0.9258 | 0.3545 | 2.6117 |
| 21.47 | 0.9395 | 0.3582 | 2.6228 |
| 21.72 | 0.9341 | 0.3646 | 2.5618 |
| 21.97 | 0.9274 | 0.3657 | 2.5358 |
| 22.22 | 0.9375 | 0.3765 | 2.4901 |
| 22.47 | 0.9283 | 0.3758 | 2.4702 |
| 22.72 | 0.9362 | 0.3864 | 2.4231 |
| 22.97 | 0.9418 | 0.3882 | 2.4263 |
| 23.22 | 0.9382 | 0.3984 | 2.3547 |
| 23.47 | 0.9360 | 0.4032 | 2.3212 |
| 23.72 | 0.9386 | 0.4069 | 2.3066 |
| 23.97 | 0.9442 | 0.4129 | 2.2869 |
| 24.22 | 0.9553 | 0.4221 | 2.2635 |
| 24.47 | 0.9569 | 0.4253 | 2.2502 |
| 24.72 | 0.9420 | 0.4257 | 2.2127 |
| 24.97 | 0.9491 | 0.4310 | 2.2023 |

Table D.7: Model \#7, $g=2 c, s=1 c$, Aerodynamic Data at Re 60,000

| a ( ${ }^{\circ}$ ) | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.14 | -0.1938 | 0.0341 | -5.6899 |
| -1.89 | -0.1679 | 0.0333 | -5.0394 |
| -1.64 | -0.1388 | 0.0315 | -4.4012 |
| -1.39 | -0.1295 | 0.0321 | $-4.0370$ |
| -1.14 | -0.1079 | 0.0283 | -3.8101 |
| -0.89 | -0.0877 | 0.0260 | -3.3690 |
| -0.64 | -0.0698 | 0.0273 | -2.5569 |
| -0.39 | -0.0443 | 0.0303 | -1.4627 |
| -0.14 | -0.0204 | 0.0271 | -0.7552 |
| 0.11 | 0.0061 | 0.0305 | 0.2008 |
| 0.36 | 0.0256 | 0.0282 | 0.9075 |
| 0.61 | 0.0390 | 0.0285 | 1.3666 |
| 0.86 | 00658 | 0.0283 | 2.3258 |
| 111 | 0.0903 | 0.0266 | 3.4004 |
| 1.36 | 0.1070 | 0.0266 | 4.0273 |
| 1.61 | 0.1226 | 0.0292 | 4.1955 |
| 186 | 0.1379 | 0.0331 | 4.1710 |
| 2.11 | 0.1695 | 0.0315 | 5.3831 |
| 2.36 | 0.1929 | 0.0305 | 6.3149 |
| 2.61 | 0.2153 | 0.0387 | 5.5570 |
| 2.86 | 0.2306 | 0.0395 | 5.8433 |
| 3.11 | 0.2451 | 0.0438 | 5.6013 |
| 3.36 | 0.2451 | 0.0438 | 5.6013 |
| 3.61 | 0.2675 | 0.0428 | 6.2574 |
| 3.86 | 0.2927 | 0.0429 | 6.8212 |
| 4.11 | 0.3164 | 0.0452 | 6.9928 |
| 4.36 | 0.3443 | 0.0467 | 7.3717 |
| 4.61 | 0.3530 | 0.0514 | 6.8735 |
| 4.86 | 0.3719 | 0.0534 | 6.9664 |
| 5.11 | 0.4222 | 0.0593 | 7.1199 |
| 5.36 | 0.4413 | 0.0596 | 7.4076 |
| 5.61 | 0.4760 | 0.0658 | 7.2363 |
| 5.86 | 0.5039 | 0.0681 | 7.4006 |
| 6.11 | 0.5260 | 0.0743 | 7.0832 |
| 6.36 | 0.5486 | 0.0780 | 7.0353 |
| 661 | 0.5776 | 0.0809 | 7.1415 |


| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.86 | 0.5993 | 0.0859 | 6.9757 |
| 7.11 | 0.6277 | 0.0899 | 6.9860 |
| 7.36 | 0.6563 | 0.0954 | 6.8810 |
| 761 | 0.6817 | 0.1003 | 6.7947 |
| 7.86 | 06955 | 0.1030 | 6.7557 |
| 8.11 | 0.7078 | 0.1128 | 6.2732 |
| 8.36 | 0.7299 | 0.1104 | 6.6127 |
| 8.61 | 0.7420 | 0.1239 | 5.9889 |
| 8.86 | 0.7555 | 0.1255 | 6.0202 |
| 9.11 | 0.7724 | 0.1293 | 5.9736 |
| 9.36 | 0.7792 | 0.1350 | 5.7717 |
| 9.61 | 0.7776 | 0.1404 | 5.5376 |
| 9.86 | 0.7950 | 0.1445 | 5.5006 |
| 10.11 | 0.7984 | 0.1498 | 5.3291 |
| 10.36 | 0.8133 | 0.1566 | 5.1950 |
| 10.61 | 0.8028 | 0.1627 | 4.9333 |
| 10.86 | 0.8193 | 0.1665 | 4.9220 |
| 11.11 | 0.8247 | 0.1684 | 4.8966 |
| 11.36 | 0.8156 | 0.1668 | 4.8898 |
| 11.61 | 0.8319 | 0.1737 | 4.7885 |
| 11.86 | 0.8414 | 0.1778 | 4.7316 |
| 12.11 | 0.8279 | 0.1790 | 4.6254 |
| 12.36 | 0.8322 | 0.1861 | 4.4707 |
| 12.61 | 0.8468 | 0.1893 | 4.4734 |
| 12.86 | 0.8553 | 0.1965 | 4.3528 |
| 13.11 | 0.8337 | 0.1966 | 4.2400 |
| 13.36 | 08427 | 0.2063 | 4.0853 |
| 13.61 | 0.8421 | 0.2130 | 3.9530 |
| 13.86 | 0.8402 | 0.2134 | 3.9368 |
| 14.11 | 0.8293 | 0.2122 | 3.9072 |
| 14.36 | 0.8406 | 0.2197 | 3.8265 |
| 14.61 | 0.8475 | 0.2227 | 3.8052 |
| 14.86 | 0.8430 | 0.2315 | 3.6416 |
| 15.11 | 0.8450 | 0.2270 | 3.7224 |
| 15.36 | 0.8565 | 0.2382 | 3.5956 |
| 15.61 | 0.8733 | 0.2425 | 3.6014 |


| $0\left(^{\circ}\right.$ ) | CL (-) | $\mathrm{Cb}_{\text {( }}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.86 | 0.8677 | 0.2433 | 3.5667 |
| 16.11 | 0.8756 | 0.2478 | 3.5343 |
| 16.36 | 0.8609 | 0.2519 | 3.4177 |
| 16.61 | 0.8841 | 0.2601 | 3.3993 |
| 16.86 | 0.8689 | 0.2631 | 3.3018 |
| 17.11 | 0.8669 | 0.2615 | 3.3156 |
| 17.36 | 0.8693 | 0.2646 | 3.2854 |
| 17.61 | 0.8804 | 0.2736 | 3.2176 |
| 17.86 | 0.8968 | 0.2791 | 3.2133 |
| 18.11 | 0.8827 | 0.2823 | 3.1270 |
| 18.36 | 0.8721 | 0.2872 | 3.0366 |
| 18.61 | 0.9020 | 0.2969 | 3.0376 |
| 18.86 | 0.8956 | 0.2940 | 3.0461 |
| 19.11 | 0.8765 | 0.2997 | 2.9248 |
| 19.36 | 0.8798 | 0.3034 | 2.8996 |
| 19.61 | 08838 | 0.3118 | 2.8349 |
| 19.86 | 0.8708 | 0.3108 | 2.8013 |
| 20.11 | 0.8800 | 0.3182 | 2.7653 |
| 20.36 | 0.8812 | 0.3196 | 2.7572 |
| 20.61 | 0.8974 | 0.3237 | 2.7723 |
| 20.86 | 0.8909 | 0.3274 | 2.7209 |
| 21.11 | 0.8930 | 0.3351 | 2.6645 |
| 21.36 | 0.8897 | 0.3364 | 2.6448 |
| 21.61 | 0.8922 | 0.3406 | 2.6192 |
| 21.86 | 0.8964 | 0.3514 | 2.5507 |
| 22.11 | 0.9205 | 0.3548 | 2.5944 |
| 22.36 | 0.8921 | 0.3561 | 2.5055 |
| 22.61 | 0.9112 | 0.3614 | 2.5211 |
| 22.86 | 0.8893 | 0.3633 | 2.4479 |
| 23.11 | 0.8972 | 0.3709 | 2.4187 |
| 23.36 | 0.8857 | 0.3720 | 2.3810 |
| 23.61 | 0.8992 | 0.3736 | 2.4068 |
| 23.86 | 0.9113 | 0.3826 | 2.3817 |
| 24.11 | 0.8971 | 0.3854 | 2.3275 |
| 24.36 | 0.9030 | 0.3918 | 2.3046 |
| 24.61 | 0.9162 | 0.4020 | 2.2792 |
| 24.86 | 0.9384 | 0.4050 | 2.3172 |

Table D.8: Model \#8, $\mathrm{g}=1 \mathrm{c}$, $\mathrm{s}=1.5 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| - ${ }^{(0}$ ) | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | a ( ${ }^{\circ}$ ) | C $\mathrm{L}_{\text {( }}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -2.42 | -0.1897 | 0.0290 | -6.5323 | 6.58 | 0.5194 | 0.0737 | 7.0522 |
| -2.17 | -0.1676 | 0.0308 | -5.4385 | 6.83 | 0.5527 | 0.0795 | 6.9525 |
| -1.92 | -0.1492 | 0.0288 | -5.1827 | 7.08 | 0.5860 | 0.0837 | 7.0014 |
| -1.67 | -0.1334 | 0.0280 | -4.7600 | 7.33 | 0.6106 | 0.0912 | 6.6968 |
| -1.42 | -0.1004 | 0.0245 | -4.1048 | 7.58 | 0.6267 | 0.0921 | 6.8087 |
| -1.17 | $-0.0853$ | 0.0257 | -3.3141 | 7.83 | 0.6297 | 0.0995 | 6.3286 |
| -0.92 | -0.0787 | 0.0252 | -3.1213 | 8.08 | 0.6583 | 0.1064 | 6.1885 |
| -0.67 | -0.0529 | 0.0224 | -2.3614 | 8.33 | 0.6797 | 0.1119 | 6.0748 |
| -0.42 | -0.0337 | 0.0222 | -1.5225 | 8.58 | 0.6996 | 0.1134 | 6.1716 |
| -0.17 | -0.0116 | 0.0222 | -0.5235 | 8.83 | 0.7095 | 0.1229 | 5.7749 |
| 0.08 | 0.0070 | 0.0258 | 0.2715 | 9.08 | 0.7334 | 0.1271 | 5.7685 |
| 0.33 | 0.0247 | 0.0242 | 1.0223 | 9.33 | 0.7511 | 0.1309 | 5.7392 |
| 0.58 | 0.0426 | 0.0256 | 1.6616 | 9.58 | 0.7551 | 0.1338 | 5.6426 |
| 0.83 | 0.0583 | 0.0249 | 2.3401 | 9.83 | 0.7753 | 0.1412 | 5.4911 |
| 1.08 | 0.0917 | 0.0249 | 3.6770 | 10.08 | 0.7749 | 0.1448 | 5.3522 |
| 1.33 | 0.1006 | 0.0259 | 3.8804 | 10.33 | 0.7990 | 0.1518 | 5.2648 |
| 1.58 | 0.1182 | 0.0274 | 4.3087 | 10.58 | 0.8103 | 0.1588 | 5.1031 |
| 1.83 | 0.1433 | 0.0241 | 5.9460 | 10.83 | 0.8237 | 0.1639 | 5.0268 |
| 2.08 | 0.1597 | 0.0285 | 5.6118 | 11.08 | 0.8264 | 0.1649 | 5.0126 |
| 2.33 | 0.1879 | 0.0301 | 6.2405 | 11.33 | 0.8264 | 0.1751 | 4.7182 |
| 2.58 | 0.2064 | 0.0313 | 6.6033 | 11.58 | 0.8499 | 0.1778 | 4.7793 |
| 2.83 | 0.2257 | 00307 | 7.3451 | 11.83 | 0.8539 | 0.1859 | 4.5923 |
| 3.08 | 0.2257 | 0.0307 | 7.3451 | 12.08 | 0.8619 | 0.1882 | 4.5804 |
| 3.33 | 0.2482 | 0.0383 | 6.4852 | 12.33 | 0.8505 | 0.1905 | 4.4645 |
| 3.58 | 0.2585 | 0.0362 | 7.1326 | 12.58 | 0.8547 | 0.1980 | 4.3173 |
| 3.83 | 0.2893 | 0.0390 | 7.4123 | 12.83 | 0.8654 | 0.1992 | 4.3443 |
| 4.08 | 0.3017 | 0.0401 | 7.5248 | 13.08 | 08605 | 0.2041 | 4.2162 |
| 4.33 | 0.3289 | 0.0454 | 7.2506 | 13.33 | 0.8775 | 0.2123 | 4.1326 |
| 4.58 | 0.3497 | 0.0458 | 7.6350 | 13.58 | 0.8892 | 0.2171 | 4.0952 |
| 4.83 | 0.3882 | 0.0504 | 7.6954 | 13.83 | 0.8734 | 0.2207 | 3.9579 |
| 5.08 | 0.4030 | 0.0553 | 7.2868 | 14.08 | 0.8804 | 0.2254 | 3.9054 |
| 5.33 | 0.4235 | 0.0526 | 8.0554 | 14.33 | 0.8779 | 0.2275 | 3.8591 |
| 5.58 | 0.4549 | 0.0608 | 7.4786 | 14.58 | 0.8654 | 0.2327 | 3.7186 |
| 5.83 | 0.4751 | 0.0644 | 7.3814 | 14.83 | 0.8827 | 0.2355 | 3.7486 |
| 6.08 | 0.5036 | 0.0694 | 7.2526 | 15.08 | 0.8859 | 0.2387 | 3.7109 |
| 6.33 | 0.5057 | 0.0711 | 7.1130 | 15.33 | 0.8806 | 0.2438 | 3.6124 |


| a ( ${ }^{\circ}$ ) | $\mathrm{CL}_{\text {L ( }}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.58 | 0.8889 | 0.2477 | 3.5881 |
| 15.83 | 0.8875 | 0.2545 | 3.4876 |
| 16.08 | 0.9026 | 0.2558 | 3.5285 |
| 16.33 | 0.9090 | 0.2651 | 34285 |
| 16.58 | 0.8958 | 0.2658 | 33704 |
| 16.83 | 0.9094 | 0.2733 | 3.3274 |
| 17.08 | 0.8887 | 0.2735 | 3.2497 |
| 17.33 | 0.9004 | 0.2764 | 3.2581 |
| 17.58 | 0.9092 | 0.2889 | 3.1468 |
| 17.83 | 0.9060 | 0.2929 | 3.0930 |
| 18.08 | 0.9062 | 0.2922 | 3.1009 |
| 18.33 | 0.9099 | 0.2985 | 3.0483 |
| 18.58 | 0.9188 | 0.3073 | 2.9900 |
| 18.83 | 0.9185 | 0.3076 | 2.9860 |
| 19.08 | 0.9151 | 0.3153 | 2.9027 |
| 19.33 | 0.9092 | 0.3196 | 2.8447 |
| 19.58 | 0.9109 | 0.3249 | 2.8038 |
| 19.83 | 0.9111 | 0.3246 | 2.8072 |
| 20.08 | 0.9361 | 0.3333 | 2.8090 |
| 20.33 | 0.9380 | 0.3445 | 2.7230 |
| 20.58 | 0.9258 | 0.3452 | 2.6819 |
| 20.83 | 0.9393 | 0.3537 | 2.6558 |
| 21.08 | 0.9374 | 0.3558 | 2.6346 |
| 21.33 | 0.9395 | 0.3578 | 2.6261 |
| 21.58 | 0.9333 | 0.3656 | 2.5528 |
| 21.83 | 0.9419 | 0.3702 | 2.5444 |
| 22.08 | 0.9340 | 0.3720 | 2.5109 |
| 22.33 | 0.9354 | 0.3771 | 2.4804 |
| 22.58 | 0.9449 | 0.3840 | 2.4606 |
| 22.83 | 0.9334 | 0.3850 | 2.4246 |
| 23.08 | 0.9454 | 0.3919 | 2.4125 |
| 23.33 | 0.9399 | 0.3966 | 2.3699 |
| 23.58 | 0.9358 | 0.4005 | 2.3363 |
| 23.83 | 0.9478 | 0.4085 | 2.3201 |
| 24.08 | 0.9494 | 0.4126 | 2.3008 |
| 24.33 | 0.9486 | 0.4170 | 2.2750 |
| 24.58 | 0.9594 | 0.4263 | 2.2506 |

Table D.9: Model \#9, $\mathrm{g}=0.5 \mathrm{c}, \mathrm{s}=-0.5 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| ¢ ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{Cu}_{0}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.93 | -0.1136 | 0.0249 | -4.5684 |
| -1.68 | -0.0957 | 0.0245 | -3.9096 |
| -1.43 | -0.0828 | 0.0234 | -3.5436 |
| -1.18 | -0.0654 | 0.0226 | -2.8904 |
| -0.93 | -0.0526 | 0.0205 | -2.5671 |
| -0.68 | -0.0384 | 0.0227 | -1.6935 |
| $-0.43$ | -0.0183 | 0.0220 | -0.8323 |
| -0.18 | -0.0105 | 0.0194 | -0.5403 |
| 0.07 | -0.0007 | 0.0201 | -0.0333 |
| 0.32 | 0.0175 | 0.0226 | 0.7730 |
| 0.57 | 0.0248 | 0.0216 | 1.1476 |
| 0.82 | 0.0327 | 0.0208 | 1.5714 |
| 1.07 | 0.0605 | 0.0212 | 2.8515 |
| 1.32 | 0.0805 | 0.0226 | 3.5546 |
| 1.57 | 0.0901 | 0.0230 | 3.9209 |
| 1.82 | 0.0990 | 0.0219 | 4.5232 |
| 2.07 | 0.1125 | 0.0226 | 4.9750 |
| 2.32 | 0.1315 | 0.0263 | 5.0056 |
| 2.57 | 0.1489 | 0.0266 | 5.5999 |
| 2.82 | 0.1701 | 0.0291 | 5.8373 |
| 3.07 | 0.1814 | 0.0281 | 6.4589 |
| 3.32 | 0.1955 | 0.0306 | 6.3894 |
| 3.57 | 0.1955 | 0.0306 | 6.3894 |
| 3.82 | 0.2053 | 0.0310 | 6.6261 |
| 4.07 | 0.2241 | 0.0331 | 6.7642 |
| 4.32 | 0.2322 | 0.0343 | 6.7738 |
| 4.57 | 0.2464 | 0.0368 | 6.6960 |
| 4.82 | 0.2601 | 0.0372 | 6.9879 |
| 5.07 | 0.2838 | 0.0426 | 6.6539 |
| 5.32 | 0.3134 | 0.0437 | 7.1708 |
| 5.57 | 0.3322 | 0.0474 | 7.0084 |
| 5.82 | 0.3498 | 0.0518 | 6.7575 |
| 6.07 | 0.3634 | 0.0506 | 7.1798 |
| 6.32 | 0.3670 | 0.0561 | 6.5420 |
| 6.57 | 0.3808 | 0.0586 | 6.4947 |
| 6.82 | 0.3984 | 0.0630 | 6.3264 |


| $a\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | $\mathrm{L} / \mathrm{D}(-)$ |
| :---: | :---: | :---: | :---: |
| 7.07 | 0.4076 | 0.0652 | 6.2514 |
| 7.32 | 0.4131 | 0.0689 | 5.9996 |
| 7.57 | 0.4289 | 0.0732 | 5.8591 |
| 7.82 | 0.4376 | 0.0769 | 5.6916 |
| 8.07 | 0.4431 | 0.0769 | 5.7637 |
| 8.32 | 0.4604 | 0.0802 | 5.7433 |
| 8.57 | 0.4691 | 0.0838 | 5.5991 |
| 8.82 | 0.4829 | 0.0878 | 5.5010 |
| 9.07 | 0.4877 | 0.0929 | 5.2472 |
| 9.32 | 0.4914 | 0.0929 | 5.2898 |
| 9.57 | 0.4968 | 0.0976 | 5.0883 |
| 9.82 | 0.5068 | 0.1009 | 5.0235 |
| 10.07 | 0.5210 | 0.1039 | 5.0137 |
| 10.32 | 0.5228 | 0.1082 | 4.8302 |
| 10.57 | 0.5298 | 0.1112 | 4.7650 |
| 1082 | 0.5405 | 0.1170 | 4.6213 |
| 11.07 | 0.5436 | 0.1200 | 4.5314 |
| 14.57 | 0.5988 | 0.1703 | 3.5167 |
| 14.82 | 0.5886 | 0.1692 | 3.4786 |
| 15.07 | 0.5936 | 0.1739 | 3.4135 |
| 15.32 | 0.5936 | 0.1765 | 3.3628 |
| 11.57 | 0.5880 | 0.1805 | 3.2584 |
| 11.57 | 0.5482 | 0.1218 | 4.5018 |
| 11.82 | 0.5586 | 0.6091 | 0.1874 | 33.2496


| c $\left(^{\circ}\right.$ ) | $\mathrm{C}_{1}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.07 | 0.5991 | 0.1905 | 3.1442 |
| 16.32 | 0.5949 | 0.1900 | 3.1315 |
| 16.57 | 0.5910 | 0.1947 | 3.0360 |
| 16.82 | 0.5988 | 0.1957 | 3.0600 |
| 17.07 | 0.6015 | 0.2012 | 2.9903 |
| 17.32 | 0.5939 | 0.1997 | 2.9735 |
| 17.57 | 0.6036 | 0.2064 | 2.9241 |
| 17.82 | 0.5989 | 0.2067 | 2.8970 |
| 18.07 | 0.5942 | 0.2075 | 2.8633 |
| 18.32 | 0.5994 | 0.2169 | 2.7634 |
| 18.57 | 0.6096 | 0.2185 | 2.7897 |
| 18.82 | 0.5882 | 0.2158 | 2.7256 |
| 19.07 | 0.5953 | 0.2198 | 2.7088 |
| 19.32 | 0.5971 | 0.2238 | 2.6687 |
| 19.57 | 0.5907 | 0.2247 | 2.6284 |
| 19.82 | 0.6041 | 0.2315 | 2.6095 |
| 20.07 | 0.5838 | 0.2312 | 2.5251 |
| 20.32 | 0.5940 | 0.2382 | 2.4942 |
| 20.57 | 0.5937 | 0.2374 | 2.5011 |
| 20.82 | 0.5833 | 0.2396 | 2.4346 |
| 21.07 | 0.5706 | 0.2363 | 2.4146 |
| 21.32 | 0.5681 | 0.2367 | 2.3998 |
| 21.57 | 0.5738 | 0.2433 | 2.3589 |
| 21.82 | 0.5708 | 0.2465 | 2.3155 |
| 22.07 | 0.5757 | 0.2494 | 2.3086 |
| 22.32 | 0.5619 | 0.2476 | 2.2694 |
| 22.57 | 0.5555 | 0.2458 | 2.2599 |
| 22.82 | 0.5695 | 0.2571 | 2.2151 |
| 23.07 | 0.5591 | 0.2540 | 2.2010 |
| 23.32 | 0.5582 | 0.2597 | 2.1494 |
| 23.57 | 0.5717 | 0.2669 | 2.1424 |
| 23.82 | 0.5584 | 0.2616 | 2.1348 |
| 24.07 | 0.5468 | 0.2612 | 2.0935 |
| 24.32 | 0.5504 | 0.2670 | 2.0619 |
| 24.57 | 0.5426 | 0.2603 | 2.0847 |
| 24.82 | 0.5358 | 0.2618 | 2.0469 |
| 25.07 | 0.5333 | 0.2652 | 2.0111 |

Table D.10: Model \#10, $g=0.5 \mathrm{c}, \mathrm{s}=-1 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| a $\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{0}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.71 | -0.1190 | 0.0251 | -4.7427 |
| -1.46 | -0.0950 | 0.0216 | $-4.4050$ |
| -1.21 | -0.0729 | 0.0205 | -3.5509 |
| -0.96 | -0.0591 | 0.0205 | -2.8746 |
| -0.71 | -0.0432 | 0.0202 | -2.1361 |
| -0.46 | -0.0323 | 0.0215 | -1.5029 |
| -0.21 | -0.0087 | 0.0190 | -0.4585 |
| 0.04 | 0.0086 | 0.0215 | 0.3972 |
| 0.29 | 0.0099 | 0.0210 | 0.4691 |
| 0.54 | 0.0356 | 0.0210 | 1.6912 |
| 0.79 | 0.0544 | 0.0210 | 2.5891 |
| 1.04 | 0.0690 | 0.0212 | 3.2513 |
| 1.29 | 0.0827 | 0.0220 | 3.7670 |
| 1.54 | 0.0963 | 0.0219 | 4.3919 |
| 1.79 | 0.1090 | 0.0250 | 4.3621 |
| 2.04 | 0.1258 | 0.0260 | 4.8375 |
| 2.29 | 0.1484 | 0.0270 | 5.5044 |
| 2.54 | 0.1643 | 0.0279 | 5.8877 |
| 2.79 | 0.1809 | 0.0292 | 6.1884 |
| 3.04 | 0.2039 | 0.0304 | 6.6980 |
| 3.29 | 0.2146 | 0.0304 | 7.0563 |
| 3.54 | 0.2376 | 0.0337 | 7.0477 |
| 3.79 | 0.2376 | 0.0337 | 7.0477 |
| 4.04 | 0.2553 | 0.0354 | 7.2027 |
| 4.29 | 0.2710 | 0.0380 | 7.1393 |
| 4.54 | 0.2922 | 0.0402 | 7.2690 |
| 4.79 | 0.2998 | 0.0404 | 7.4151 |
| 5.04 | 0.3106 | 0.0448 | 6.9385 |
| 5.29 | 0.3369 | 0.0480 | 7.0196 |
| 5.54 | 0.3548 | 0.0533 | 6.6572 |
| 5.79 | 0.3752 | 0.0568 | 6.6097 |
| 604 | 0.3802 | 0.0601 | 6.3305 |
| 6.29 | 0.4046 | 0.0613 | 6.5992 |
| 6.54 | 0.4116 | 0.0646 | 6.3750 |
| 6.79 | 0.4255 | 0.0676 | 6.2960 |
| 7.04 | 0.4387 | 0.0715 | 6.1360 |


| - ( ${ }^{\circ}$ ) | C $\mathrm{l}^{(-)}$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.29 | 0.4449 | 0.0726 | 6.1324 |
| 7.54 | 0.4535 | 0.0798 | 5.6814 |
| 7.79 | 0.4636 | 0.0816 | 5.6838 |
| 804 | 0.4691 | 0.0831 | 5.6441 |
| 8.29 | 0.4780 | 0.0863 | 5.5391 |
| 8.54 | 0.4831 | 0.0888 | 5.4399 |
| 8.79 | 0.4933 | 0.0956 | 5.1602 |
| 9.04 | 0.5040 | 0.0979 | 5.1453 |
| 9.29 | 0.5026 | 0.0982 | 5.1201 |
| 9.54 | 0.5044 | 0.1044 | 4.8294 |
| 9.79 | 0.5030 | 0.1064 | 4.7261 |
| 1004 | 0.5073 | 0.1102 | 4.6037 |
| 10.29 | 0.5093 | 0.1082 | 4.7045 |
| 10.54 | 0.5138 | 0.1136 | 4.5246 |
| 10.79 | 0.5139 | 0.1190 | 4.3180 |
| 11.04 | 0.5270 | 0.1228 | 4.2909 |
| 11.29 | 0.5173 | 0.1213 | 4.2656 |
| 11.54 | 0.5222 | 0.1223 | 4.2704 |
| 11.79 | 0.5251 | 0.1286 | 4.0842 |
| 12.04 | 0.5250 | 0.1269 | 4.1354 |
| 12.29 | 0.5278 | 0.1325 | 3.9839 |
| 12.54 | 0.5192 | 0.1320 | 3.9347 |
| 12.79 | 0.5260 | 0.1353 | 3.8875 |
| 13.04 | 0.5258 | 0.1406 | 3.7391 |
| 13.29 | 0.5089 | 0.1409 | 3.6123 |
| 13.54 | 0.5083 | 0.1429 | 3.5570 |
| 13.79 | 0.5206 | 0.1467 | 3.5488 |
| 14.04 | 0.5129 | 0.1485 | 3.4533 |
| 14.29 | 0.5098 | 0.1474 | 3.4576 |
| 14.54 | 0.4929 | 0.1453 | 3.3917 |
| 14.79 | 0.4985 | 0.1502 | 3.3183 |
| 15.04 | 0.4894 | 0.1484 | 3.2987 |
| 15.29 | 0.4842 | 0.1506 | 3.2142 |
| 15.54 | 0.4817 | 0.1519 | 3.1716 |
| 15.79 | 0.4762 | 0.1515 | 3.1443 |
| 16.04 | 0.4710 | 0.1525 | 3.0883 |


| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.29 | 0.4650 | 0.1523 | 3.0541 |
| 16.54 | 0.4632 | 0.1550 | 2.9877 |
| 16.79 | 0.4640 | 0.1563 | 2.9687 |
| 17.04 | 0.4501 | 0.1548 | 2.9083 |
| 17.29 | 0.4430 | 0.1553 | 2.8522 |
| 17.54 | 0.4447 | 0.1555 | 2.8591 |
| 17.79 | 0.4393 | 0.1572 | 2.7949 |
| 18.04 | 0.4258 | 0.1529 | 2.7843 |
| 18.29 | 0.4203 | 0.1545 | 2.7199 |
| 18.54 | 0.4149 | 0.1522 | 2.7258 |
| 18.79 | 0.4086 | 0.1553 | 2.6315 |
| 19.04 | 0.3955 | 0.1487 | 2.6601 |
| 19.29 | 0.4086 | 0.1614 | 2.5310 |
| 19.54 | 0.4019 | 0.1560 | 2.5765 |
| 19.79 | 0.3990 | 0.1590 | 2.5102 |
| 20.04 | 0.3835 | 0.1559 | 2.4596 |
| 20.29 | 0.3913 | 0.1604 | 2.4392 |
| 20.54 | 0.3922 | 0.1574 | 2.4922 |
| 20.79 | 0.3814 | 0.1593 | 2.3945 |
| 21.04 | 0.3791 | 0.1587 | 23894 |
| 21.29 | 0.3751 | 0.1622 | 2.3122 |
| 21.54 | 0.3771 | 0.1626 | 2.3187 |
| 21.79 | 0.3708 | 0.1645 | 2.2538 |
| 22.04 | 0.3624 | 0.1640 | 2.2094 |
| 22.29 | 0.3657 | 0.1674 | 2.1850 |
| 22.54 | 0.3696 | 0.1680 | 2.2002 |
| 22.79 | 0.3677 | 0.1716 | 2.1425 |
| 23.04 | 0.3611 | 0.1691 | 2.1358 |
| 23.29 | 0.3613 | 0.1698 | 2.1282 |
| 23.54 | 0.3595 | 0.1691 | 2.1258 |
| 23.79 | 0.3615 | 0.1730 | 2.0896 |
| 24.04 | 0.3593 | 0.1758 | 2.0435 |
| 24.29 | 0.3628 | 0.1769 | 2.0508 |
| 24.54 | 0.3600 | 0.1763 | 2.0419 |
| 24.79 | 0.3529 | 0.1760 | 2.0054 |
| 25.04 | 0.3595 | 0.1792 | 2.0062 |
| 25.29 | 0.3641 | 0.1838 | 1.9808 |

Table D.11: Model \#11, $g=1 \mathrm{c}, \mathrm{s}=-\mathbf{0 . 5 c}$, Aerodynamic Data at $\operatorname{Re} 60,000$

| व ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{CD}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.07 | -0.1402 | 00305 | -4.5939 |
| -1.82 | -0.1258 | 0.0314 | -4.0036 |
| $\cdot 1.57$ | -0.1037 | 0.0312 | -3.3253 |
| -1.32 | -0.0907 | 0.0302 | -3.0064 |
| -1.07 | -0.0702 | 0.0319 | -2.1996 |
| -0.82 | -0.0373 | 0.0285 | -1.3099 |
| -0.57 | -0.0329 | 0.0332 | -0.9903 |
| -0.32 | -0.0213 | 0.0311 | -0.6858 |
| -0.07 | 0.0026 | 0.0267 | 0.0980 |
| 0.18 | 0.0108 | 0.0288 | 0.3749 |
| 0.43 | 0.0324 | 0.0303 | 1.0689 |
| 0.68 | 0.0597 | 0.0285 | 2.0952 |
| 0.93 | 0.0738 | 0.0287 | 2.5725 |
| 1.18 | 0.0873 | 0.0277 | 3.1567 |
| 1.43 | 0.1191 | 0.0288 | 4.1343 |
| 1.68 | 0.1290 | 0.0303 | 4.2626 |
| 1.93 | 0.1462 | 0.0328 | 4.4543 |
| 2.18 | 0.1614 | 0.0341 | 4.7273 |
| 2.43 | 0.1823 | 0.0342 | 5.3225 |
| 2.68 | 0.2067 | 0.0376 | 5.5006 |
| 2.93 | 0.2248 | 0.0405 | 5.5545 |
| 3.18 | 0.2442 | 0.0416 | 5.8662 |
| 3.43 | 0.2442 | 0.0416 | 5.8662 |
| 3.68 | 0.2717 | 0.0433 | 6.2794 |
| 3.93 | 0.2789 | 0.0435 | 6.4163 |
| 4.18 | 0.2894 | 0.0496 | 5.8352 |
| 4.43 | 0.3135 | 0.0507 | 6.1820 |
| 4.68 | 0.3301 | 0.0524 | 6.2968 |
| 4.93 | 0.3549 | 0.0574 | 6.1867 |
| 5.18 | 0.3872 | 0.0621 | 6.2371 |
| 5.43 | 0.4145 | 0.0657 | 6.3076 |
| 5.68 | 0.4405 | 0.0661 | 6.6611 |
| 5.93 | 0.4510 | 0.0701 | 6.4301 |
| 6.18 | 0.4698 | 0.0752 | 6.2439 |
| 6.43 | 0.4737 | 0.0813 | 5.8280 |
| 6.68 | 0.5041 | 0.0855 | 5.8991 |


| $\square\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | Co (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.93 | 0.5332 | 0.0889 | 6.0002 |
| 7.18 | 0.5430 | 0.0924 | 5.8748 |
| 7.43 | 0.5449 | 0.0941 | 5.7883 |
| 7.68 | 0.5714 | 0.0997 | 5.7300 |
| 7.93 | 0.6103 | 0.1095 | 5.5749 |
| 8.18 | 0.6191 | 0.1123 | 5.5114 |
| 8.43 | 0.6251 | 0.1174 | 5.3229 |
| 8.68 | 0.6472 | 0.1244 | 5.2005 |
| 8.93 | 0.6655 | 0.1297 | 5.1322 |
| 9.18 | 0.6748 | 0.1315 | 5.1301 |
| 9.43 | 0.6823 | 0.1347 | 5.0636 |
| 9.68 | 0.6817 | 0.1389 | 4.9071 |
| 9.93 | 0.6930 | 0.1423 | 4.8705 |
| 10.18 | 0.6944 | 01462 | 4.7494 |
| 10.43 | 0.6986 | 0.1525 | 4.5820 |
| 1068 | 0.7171 | 0.1609 | 4.4571 |
| 10.93 | 0.7115 | 0.1654 | 4.3012 |
| 11.18 | 0.7111 | 0.1654 | 4.2999 |
| 1143 | 0.7165 | 0.1728 | 4.1461 |
| 11.68 | 0.7149 | 0.1768 | 4.0431 |
| 11.93 | 0.7115 | 0.1810 | 3.9311 |
| 12.18 | 0.7142 | 0.1802 | 3.9646 |
| 12.43 | 0.7151 | 0.1867 | 3.8293 |
| 12.68 | 0.7087 | 0.1856 | 3.8181 |
| 12.93 | 0.7141 | 0.1930 | 3.6997 |
| 13.18 | 0.7085 | 0.1954 | 3.6259 |
| 13.43 | 0.6989 | 0.1963 | 3.5602 |
| 13.68 | 0.7011 | 0.1966 | 3.5658 |
| 13.93 | 0.7138 | 0.2047 | 3.4876 |
| 14.18 | 0.7171 | 0.2126 | 3.3724 |
| 14.43 | 0.6995 | 0.2100 | 3.3307 |
| 14.68 | 0.7060 | 0.2161 | 3.2670 |
| 14.93 | 0.7069 | 0.2193 | 3.2241 |
| 15.18 | 0.7154 | 0.2259 | 3.1666 |
| 15.43 | 0.7088 | 0.2257 | 3.1410 |
| 15.68 | 0.7022 | 0.2256 | 3.1121 |


| $\alpha\left({ }^{\circ}\right)$ | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.93 | 0.7206 | 0.2360 | 3.0533 |
| 16.18 | 0.7169 | 0.2401 | 2.9864 |
| 16.43 | 0.7146 | 0.2424 | 2.9478 |
| 16.68 | 0.7147 | 0.2461 | 2.9042 |
| 16.93 | 0.7294 | 0.2540 | 2.8719 |
| 17.18 | 0.7112 | 0.2522 | 2.8201 |
| 17.43 | 0.7159 | 0.2608 | 2.7449 |
| 17.68 | 0.7065 | 0.2586 | 2.7325 |
| 17.93 | 0.7121 | 0.2630 | 2.7074 |
| 18.18 | 0.7154 | 0.2706 | 2.6432 |
| 18.43 | 0.7329 | 0.2782 | 2.6346 |
| 18.68 | 0.7248 | 0.2803 | 2.5856 |
| 18.93 | 0.7168 | 0.2791 | 2.5684 |
| 19.18 | 0.7184 | 0.2865 | 2.5075 |
| 19.43 | 0.7290 | 0.2891 | 2.5215 |
| 19.68 | 0.7115 | 0.2912 | 2.4433 |
| 19.93 | 0.7212 | 0.2979 | 2.4211 |
| 20.18 | 0.7243 | 0.3024 | 2.3954 |
| 20.43 | 0.7338 | 0.3095 | 2.3711 |
| 20.68 | 0.7362 | 0.3112 | 2.3658 |
| 20.93 | 0.7217 | 0.3097 | 2.3306 |
| 21.18 | 0.7201 | 0.3167 | 2.2738 |
| 21.43 | 0.7225 | 0.3192 | 2.2632 |
| 21.68 | 0.7269 | 0.3229 | 2.2507 |
| 21.93 | 0.7280 | 0.3289 | 2.2134 |
| 22.18 | 0.7320 | 0.3332 | 2.1970 |
| 22.43 | 0.7179 | 0.3360 | 2.1362 |
| 22.68 | 0.7184 | 0.3351 | 2.1442 |
| 22.93 | 0.7169 | 0.3391 | 2.1144 |
| 23.18 | 0.7226 | 0.3475 | 2.0794 |
| 23.43 | 0.7217 | 0.3534 | 2.0424 |
| 23.68 | 0.7201 | 0.3517 | 2.0473 |
| 23.93 | 0.7201 | 0.3566 | 2.0195 |
| 24.18 | 0.7151 | 0.3568 | 2.0041 |
| 24.43 | 0.7081 | 0.3598 | 1.9683 |
| 24.68 | 0.7124 | 0.3659 | 1.9470 |
| 24.93 | 0.7128 | 0.3668 | 1.9434 |

Table D.12: Model \#12, $g=1 \mathrm{c}, \mathrm{s}=-1 \mathrm{c}$, Aerodynamic Data at Re 60,000

| a ( ${ }^{\circ}$ ) | CL. (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.89 | -0.1364 | 0.0262 | -5.2140 |
| -1.64 | -0.1192 | 0.0247 | -4.8274 |
| -1.39 | -0.0960 | 0.0242 | -3.9613 |
| -1.14 | -0.0847 | 0.0224 | -3.7864 |
| -0.89 | -0.0761 | 0.0228 | -3.3402 |
| -0.64 | -0,0572 | 0.0201 | -2.8419 |
| -0.39 | -0.0240 | 0.0217 | -1.1073 |
| -0.14 | -0.0050 | 0.0209 | -0.2418 |
| 0.11 | 0.0123 | 0.0224 | 0.5498 |
| 0.36 | 0.0204 | 0.0220 | 0.9281 |
| 0.61 | 0.0453 | 0.0191 | 2.3735 |
| 0.86 | 0.0530 | 0.0250 | 2.1179 |
| 1.11 | 0.0748 | 0.0243 | 3.0797 |
| 1.36 | 0.0995 | 0.0269 | 3.6953 |
| 1.61 | 0.1175 | 0.0261 | 4.5011 |
| 1.86 | 0.1340 | 0.0258 | 5.1896 |
| 2.11 | 0.1542 | 0.0265 | 5.8112 |
| 2.36 | 0.1627 | 0.0258 | 6.3127 |
| 2.61 | 0.1724 | 0.0299 | 5.7634 |
| 2.86 | 0.1885 | 0.0288 | 6.5409 |
| 3.11 | 0.2030 | 0.0284 | 7.1538 |
| 3.36 | 0.2298 | 0.0322 | 7.1392 |
| 3.61 | 0.2298 | 0.0322 | 7.1392 |
| 3.86 | 0.2465 | 0.0337 | 7.3213 |
| 4.11 | 0.2643 | 0.0366 | 7.2188 |
| 4.36 | 0.3020 | 0.0419 | 7.2046 |
| 4.61 | 0.3078 | 0.0392 | 7.8496 |
| 4.86 | 0.3262 | 0.0457 | 7.1438 |
| 5.11 | 0.3405 | 0.0490 | 6.9434 |
| 5.36 | 0.3641 | 0.0535 | 6.8001 |
| 5.61 | 0.3965 | 0.0536 | 7.4007 |
| 5.86 | 0.4125 | 0.0591 | 6.9737 |
| 6.11 | 0.4395 | 0.0659 | 6.6668 |
| 6.36 | 0.4577 | 0.0671 | 6.8241 |
| 6.61 | 0.4584 | 0.0712 | 6.4423 |
| 6.86 | 0.4737 | 0.0734 | 6.4512 |


| a $\left(^{\circ}\right.$ ) | $\mathrm{Cl}_{1 .}(-)$ | $C_{\text {d }}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.11 | 0.4906 | 0.0746 | 6.5804 |
| 7.36 | 0.5068 | 0.0821 | 6.1744 |
| 7.61 | 0.5231 | 0.0817 | 64046 |
| 7.86 | 0.5289 | 0.0896 | 5.9045 |
| 8.11 | 0.5523 | 0.0911 | 6.0629 |
| 8.36 | 0.5654 | 0.1001 | 5.6498 |
| 8.61 | 0.5692 | 0.1023 | 5.5660 |
| 8.86 | 0.5761 | 0.1049 | 5.4901 |
| 9.11 | 0.5855 | 0.1091 | 5.3679 |
| 9.36 | 0.5954 | 0.1136 | 5.2399 |
| 9.61 | 0.5978 | 0.1177 | 5.0784 |
| 9.86 | 0.6080 | 0.1226 | 4.9588 |
| 10.11 | 0.6107 | 0.1267 | 4.8190 |
| 10.36 | 0.6363 | 0.1350 | 4.7136 |
| 10.61 | 0.6389 | 0.1357 | 4.7073 |
| 10.86 | 0.6321 | 0.1339 | 4.7209 |
| 11.11 | 0.6291 | 0.1380 | 4.5606 |
| 11.36 | 0.6381 | 0.1436 | 4.4435 |
| 11.61 | 0.6418 | 0.1455 | 4.4114 |
| 11.86 | 0.6361 | 0.1477 | 4.3074 |
| 1211 | 0.6431 | 0.1511 | 4.2546 |
| 12.36 | 0.6554 | 0.1593 | 4.1146 |
| 12.61 | 0.6730 | 0.1676 | 4.0152 |
| 12.86 | 0.6598 | 0.1654 | 3.9899 |
| 13.11 | 0.6640 | 0.1718 | 3.8661 |
| 13.36 | 0.6670 | 0.1736 | 3.8419 |
| 13.61 | 0.6610 | 0.1755 | 3.7669 |
| 13.86 | 0.6595 | 0.1751 | 3.7658 |
| 14.11 | 0.6680 | 0.1826 | 3.6581 |
| 14.36 | 0.6766 | 0.1900 | 3.5603 |
| 14.61 | 0.6700 | 0.1935 | 3.4632 |
| 14.86 | 0.6693 | 0.1973 | 3.3925 |
| 15.11 | 0.6623 | 0.1974 | 3.3559 |
| 15.36 | 0.6873 | 0.2080 | 3,3037 |
| 15.61 | 0.6829 | 0.2134 | 3.1997 |
| 15.86 | 0.6686 | 0.2081 | 3.2138 |


| a ( ${ }^{\circ}$ ) | CL (-) | $C_{0}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.11 | 0.6639 | 0.2077 | 3.1958 |
| 16.36 | 0.6744 | 0.2152 | 3.1335 |
| 16.61 | 0.6674 | 0.2201 | 3.0318 |
| 16.86 | 0.6860 | 0.2302 | 2.9802 |
| 17.11 | 0.6750 | 0.2269 | 2.9746 |
| 17.36 | 0.6724 | 0.2307 | 2.9142 |
| 17.61 | 0.6751 | 0.2362 | 2.8581 |
| 17.86 | 0.6685 | 0.2348 | 2.8473 |
| 18.11 | 0.6858 | 0.2453 | 2.7959 |
| 18.36 | 0.6862 | 0.2437 | 2.8160 |
| 18.61 | 0.6646 | 0.2438 | 2.7254 |
| 18.86 | 06464 | 0.2401 | 2.6925 |
| 19.11 | 0.6636 | 0.2508 | 2.6459 |
| 19.36 | 0.6674 | 0.2610 | 2.5571 |
| 19.61 | 0.6689 | 0.2618 | 2.5553 |
| 19.86 | 0.6502 | 0.2588 | 2.5122 |
| 20.11 | 0.6708 | 02671 | 2.5114 |
| 20.36 | 0.6678 | 0.2723 | 2.4525 |
| 20.61 | 0.6661 | 0.2704 | 2.4632 |
| 20.86 | 0.6736 | 0.2832 | 2.3786 |
| 21.11 | 0.6441 | 0.2740 | 2.3502 |
| 21.36 | 0.6905 | 0.2925 | 2.3607 |
| 21.61 | 0.6681 | 0.2888 | 2.3135 |
| 21.86 | 0.6651 | 0.2969 | 2.2401 |
| 22.11 | 0.6734 | 0.2973 | 2.2649 |
| 22.36 | 0.6591 | 0.2937 | 2.2440 |
| 22.61 | 0.6774 | 0.3025 | 2.2390 |
| 22.86 | 0.6889 | 0.3174 | 2.1708 |
| 23.11 | 0.6717 | 0.3042 | 2.2080 |
| 23.36 | 0.7080 | 0.3281 | 2.1582 |
| 23.61 | 0.7172 | 0.3354 | 2.1384 |
| 23.86 | 0.7415 | 0.3474 | 2.1346 |
| 24.11 | 0.7660 | 0.3636 | 2.1063 |
| 24.36 | 0.7565 | 0.3651 | 2.0721 |
| 24.61 | 0.7174 | 0.3464 | 2.0713 |
| 24.86 | 0.7292 | 0.3573 | 2.0407 |
| 25.11 | 0.7180 | 0.3616 | 1.9857 |

Table D.13: Model \#13, $g=2 c, s=-1 c$, Aerodynamic Data at Re 60,000

| $\alpha\left(^{\circ}\right.$ | $\mathrm{C}_{\mathrm{L}}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.83 | -0.1629 | 0.0261 | -6.2491 |
| -1.58 | -0.1486 | 0.0265 | -5.6049 |
| -1.33 | -0.1172 | 0.0250 | -4.6899 |
| -1.08 | -0.0850 | 0.0250 | -3.3943 |
| -0.83 | -0.0723 | 0.0205 | -3.5302 |
| -0.58 | -0.0557 | 0.0247 | -2.2539 |
| -0.33 | -0.0252 | 0.0232 | $-1.0878$ |
| -0.08 | -0.0073 | 0.0254 | -0.2873 |
| 0.17 | 0.0117 | 0.0250 | 0.4688 |
| 0.42 | 0.0321 | 0.0227 | 1.4151 |
| 0.67 | 0.0460 | 0.0239 | 1.9267 |
| 0.92 | 0.0736 | 0.0216 | 3.4100 |
| 1.17 | 0.1030 | 0.0250 | 4.1117 |
| 1.42 | 0.1255 | 0.0254 | 4.9454 |
| 1.67 | 0.1446 | 0.0273 | 5.2902 |
| 1.92 | 0.1563 | 0.0261 | 5.9809 |
| 2.17 | 0.1770 | 0.0285 | 6.2154 |
| 2.42 | 0.1975 | 0.0302 | 6.5291 |
| 2.67 | 0.2132 | 0.0314 | 6.7953 |
| 2.92 | 0.2425 | 0.0337 | 7.1943 |
| 3.17 | 0.2675 | 0.0340 | 7.8595 |
| 3.42 | 0.2999 | 0.0397 | 7.5491 |
| 3.67 | 0.2999 | 0.0397 | 7.5491 |
| 3.92 | 0.3023 | 0.0356 | 8.5015 |
| 4.17 | 0.3341 | 0.0442 | 7.5523 |
| 4.42 | 0.3375 | 0.0419 | 8.0541 |
| 4.67 | 0.3763 | 0.0465 | 8.0996 |
| 492 | 0.4024 | 0.0492 | 8.1837 |
| 5.17 | 0.4354 | 0.0547 | 7.9546 |
| 5.42 | 0.4674 | 0.0585 | 7.9895 |
| 5.67 | 0.4928 | 0.0660 | 7.4626 |
| 5.92 | 0.5294 | 0.0732 | 7.2312 |
| 6.17 | 0.5627 | 0.0743 | 7.5684 |
| 642 | 0.5620 | 0.0777 | 7.2326 |
| 6.67 | 0.5950 | 0.0826 | 7.2035 |
| 6.92 | 0.6178 | 0.0856 | 7.2157 |


| a ( ${ }^{\circ}$ ) | CL (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.17 | 0.6475 | 0.0924 | 7.0054 |
| 7.42 | 0.6515 | 0.0979 | 6.6522 |
| 7.67 | 0.6767 | 0.1002 | 6.7572 |
| 7.92 | 0.7008 | 0.1047 | 6.6923 |
| 8.17 | 0.7165 | 0.1137 | 6.3002 |
| 8.42 | 0.7232 | 0.1111 | 6.5099 |
| 8.67 | 0.7494 | 0.1153 | 6.5019 |
| 8.92 | 0.7344 | 0.1202 | 6.1086 |
| 9.17 | 0.7631 | 0.1243 | 6.1382 |
| 9.42 | 0.7710 | 0.1342 | 5.7435 |
| 9.67 | 0.7585 | 0.1339 | 5.6630 |
| 9.92 | 0.7790 | 0.1461 | 5.3323 |
| 10.17 | 0.7870 | 0.1495 | 5.2646 |
| 10.42 | 0.7835 | 0.1503 | 5.2132 |
| 10.67 | 0.7975 | 0.1568 | 5.0848 |
| 10.92 | 0.7888 | 0.1555 | 5.0722 |
| 11.17 | 0.7724 | 0.1641 | 4.7074 |
| 11.42 | 0.7870 | 0.1679 | 4.6878 |
| 11.67 | 0.7898 | 0.1716 | 4.6033 |
| 11.92 | 0.7827 | 0.1671 | 4.6829 |
| 12.17 | 0.7846 | 0.1738 | 4.5135 |
| 12.42 | 0.7916 | 0.1773 | 4.4652 |
| 12.67 | 0.7837 | 0.1736 | 4.5145 |
| 12.92 | 0.7875 | 0.1827 | 4.3114 |
| 13.17 | 0.7715 | 0.1880 | 4.1033 |
| 13.42 | 0.7892 | 0.1910 | 4.1324 |
| 13.67 | 0.7940 | 0.1921 | 4.1340 |
| 13.92 | 0.7739 | 0.1914 | 4.0424 |
| 14.17 | 0.7859 | 0.2043 | 3.8459 |
| 1442 | 0.7998 | 0.2159 | 3.7039 |
| 14.67 | 0.8010 | 0.2169 | 3.6923 |
| 14.92 | 0.7827 | 0.2137 | 3.6624 |
| 15.17 | 0.7958 | 0.2183 | 3.6449 |
| 15.42 | 0.7852 | 0.2249 | 3.4919 |
| 15.67 | 0.7933 | 0.2261 | 3.5086 |
| 15.92 | 0.7967 | 0.2277 | 3.4995 |


| $\mathrm{a}\left({ }^{\circ}{ }^{\text {) }}\right.$ | C. (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.17 | 0.7934 | 0.2283 | 3.4753 |
| 16.42 | 0.8168 | 0.2379 | 3.4335 |
| 16.67 | 0.8054 | 0.2404 | 3.3503 |
| 16.92 | 0.7981 | 0.2476 | 3.2233 |
| 17.17 | 0.8220 | 0.2611 | 3.1487 |
| 17.42 | 0.7917 | 0.2533 | 3.1250 |
| 17.67 | 0.7999 | 0.2683 | 2.9814 |
| 17.92 | 0.8105 | 0.2653 | 3.0547 |
| 18.17 | 0.8058 | 0.2616 | 3.0808 |
| 18.42 | 0.7966 | 0.2728 | 2.9205 |
| 18.67 | 0.8108 | 0.2735 | 2.9646 |
| 18.92 | 0.8042 | 0.2734 | 2.9418 |
| 19.17 | 0.8140 | 0.2826 | 28807 |
| 19.42 | 0.8027 | 0.2871 | 2.7956 |
| 19.67 | 0.8021 | 0.2930 | 2.7373 |
| 19.92 | 0.8114 | 0.3000 | 2.7051 |
| 20.17 | 0.8066 | 0.2944 | 2.7396 |
| 20.42 | 0.8019 | 0.3041 | 2.6371 |
| 20.67 | 0.8019 | 0.3105 | 2.5821 |
| 20.92 | 0.8175 | 0.3094 | 2.6418 |
| 21.17 | 0.8113 | 0.3153 | 2.5732 |
| 21.42 | 0.8039 | 0.3167 | 2.5385 |
| 21.67 | 0.8151 | 0.3301 | 2.4691 |
| 21.92 | 0.8243 | 0.3204 | 2.5727 |
| 22.17 | 0.7931 | 0.3135 | 2.5303 |
| 22.42 | 0.8083 | 0.3324 | 2.4318 |
| 22.67 | 0.7960 | 0.3356 | 2.3721 |
| 22.92 | 0.8062 | 0.3400 | 2.3713 |
| 23.17 | 0.8053 | 0.3482 | 2.3128 |
| 23.42 | 0.8083 | 0.3548 | 2.2783 |
| 23.67 | 08004 | 0.3501 | 2.2860 |
| 23.92 | 0.7921 | 0.3461 | 2.2889 |
| 24.17 | 0.7948 | 0.3579 | 2.2210 |
| 24.42 | 0.8313 | 0.3745 | 2.2198 |
| 24.67 | 0.7944 | 0.3626 | 2.1907 |
| 24.92 | 0.8103 | 0.3728 | 2.1734 |
| 25.17 | 0.8063 | 0.3756 | 2.1465 |

Table D.14: Model \#14, $g=1 c, s=-1.5 c$, Aerodynamic Data at $\operatorname{Re} 60,000$

| a ( ${ }^{\circ}$ ) | $\mathrm{Cl}_{\mathrm{L}}(-)$ | $\mathrm{CD}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.88 | -0.1382 | 0.0257 | $-5.3788$ |
| -1.63 | -0.1116 | 0.0247 | -4.5153 |
| -1.38 | -0.1037 | 0.0217 | -4.7879 |
| -1.13 | -0.0770 | 0.0230 | -3.3519 |
| -0.88 | -0.0575 | 0.0235 | $-2.4497$ |
| -0.63 | -0.0404 | 0.0203 | -1.9897 |
| -0.38 | -0.0239 | 0.0190 | -1.2595 |
| -0.13 | -0.0046 | 0.0234 | -0.1946 |
| 0.12 | 0.0107 | 0.0231 | 0.4643 |
| 0.37 | 0.0191 | 0.0224 | 0.8521 |
| 0.62 | 0.0430 | 0.0264 | 1.6276 |
| 0.87 | 0.0649 | 0.0231 | 2.8043 |
| 1.12 | 0.0813 | 0.0242 | 3.3527 |
| 1.37 | 0.0985 | 0.0244 | 4.0402 |
| 1.62 | 0.1177 | 0.0277 | 4.2405 |
| 1.87 | 0.1392 | 0.0276 | 5.0421 |
| 2.12 | 0.1550 | 0.0275 | 5.6302 |
| 2.37 | 0.1754 | 0.0271 | 6.4749 |
| 2.62 | 0.1838 | 0.0301 | 6.1164 |
| 2.87 | 0.2090 | 0.0360 | 5.8075 |
| 3.12 | 0.2181 | 0.0293 | 7.4525 |
| 3.37 | 0.2401 | 0.0377 | 6.3743 |
| 3.62 | 0.2401 | 0.0377 | 6.3743 |
| 3.87 | 0.2617 | 0.0380 | 6.8797 |
| 4.12 | 0.2836 | 0.0400 | 7.0966 |
| 4.37 | 0.3010 | 0.0430 | 7.0065 |
| 4.62 | 0.3057 | 0.0427 | 7.1536 |
| 4.87 | 0.3370 | 0.0481 | 7.0017 |
| 5.12 | 0.3492 | 0.0504 | 6.9325 |
| 5.37 | 0.3766 | 0.0573 | 6.5708 |
| 5.62 | 0.3982 | 0.0600 | 6.6332 |
| 5.87 | 0.4178 | 0.0628 | 6.6501 |
| 6.12 | 0.4424 | 0.0670 | 6.6069 |
| 6.37 | 0.4546 | 0.0685 | 6.6410 |
| 6.62 | 0.4699 | 0.0746 | 6.3023 |
| 6.87 | 0.4788 | 0.0753 | 6.3572 |


| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.12 | 0.4966 | 0.0830 | 5.9828 |
| 7.37 | 0.5153 | 0.0881 | 5.8520 |
| 7.62 | 0.5219 | 0.0883 | 5.9113 |
| 7.87 | 0.5475 | 0.0956 | 5.7265 |
| 8.12 | 0.5648 | 0.1038 | 5.4390 |
| 8.37 | 0.5652 | 0.1001 | 5.6465 |
| 8.62 | 0.5858 | 0.1053 | 5.5618 |
| 8.87 | 0.5986 | 0.1145 | 5.2270 |
| 9.12 | 0.6069 | 0.1132 | 5.3601 |
| 9.37 | 0.6199 | 0.1208 | 5.1315 |
| 9.62 | 0.6204 | 0.1257 | 4.9350 |
| 9.87 | 0.6246 | 0.1257 | 4.9709 |
| 10.12 | 0.6309 | 0.1329 | 4.7475 |
| 10.37 | 0.6395 | 0.1406 | 4.5494 |
| 1062 | 0.6532 | 0.1447 | 4.5147 |
| 10.87 | 06490 | 0.1459 | 4.4477 |
| 11.12 | 0.6511 | 0.1455 | 4.4736 |
| 11.37 | 0.6530 | 0.1537 | 4.2484 |
| 11.62 | 0.6576 | 0.1554 | 4.2319 |
| 11.87 | 0.6571 | 0.1573 | 4.1786 |
| 12.12 | 0.6763 | 0.1646 | 4.1083 |
| 12.37 | 0.6661 | 0.1641 | 4.0599 |
| 12.62 | 0.6838 | 0.1782 | 3.8378 |
| 12.87 | 0.6864 | 0.1756 | 3.9092 |
| 13.12 | 0.6725 | 0.1771 | 3.7964 |
| 13.37 | 0.6919 | 0.1857 | 3.7258 |
| 13.62 | 0.6909 | 0.1883 | 3.6688 |
| 13.87 | 0.6897 | 0.1921 | 3.5910 |
| 14.12 | 0.6991 | 0.1964 | 3.5593 |
| 14.37 | 0.6921 | 0.2007 | 3.4487 |
| 14.62 | 0.6914 | 0.2009 | 3.4411 |
| 14.87 | 0.6881 | 0.2027 | 3.3941 |
| 15.12 | 0.6967 | 0.2078 | 3.3528 |
| 15.37 | 0.6873 | 0.2099 | 3.2753 |
| 15.62 | 0.6920 | 0.2158 | 3.2058 |
| 15.87 | 0.6819 | 0.2218 | 3.0745 |


| a ( ${ }^{\circ}$ ) | CL (-) | $C_{D}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.12 | 0.7036 | 0.2241 | 3.1398 |
| 16.37 | 0.7040 | 0.2342 | 3.0059 |
| 16.62 | 0.7176 | 0.2428 | 2.9550 |
| 16.87 | 0.6965 | 0.2335 | 2.9836 |
| 17.12 | 0.7015 | 0.2396 | 2.9283 |
| 17.37 | 0.7128 | 0.2407 | 2.9618 |
| 17.62 | 0.7151 | 0.2523 | 2.8342 |
| 17.87 | 0.6881 | 0.2434 | 2.8272 |
| 18.12 | 0.6760 | 0.2391 | 2.8273 |
| 18.37 | 0.6782 | 0.2428 | 2.7929 |
| 18.62 | 0.6757 | 0.2461 | 2.7455 |
| 18.87 | 0.6556 | 0.2463 | 2.6616 |
| 19.12 | 0.6735 | 0.2561 | 2.6304 |
| 19.37 | 0.6786 | 0.2636 | 2.5747 |
| 19.62 | 0.6789 | 0.2697 | 2.5175 |
| 19.87 | 0.6544 | 0.2521 | 2.5963 |
| 20.12 | 0.6548 | 0.2642 | 2.4782 |
| 20.37 | 0.6522 | 0.2625 | 2.4844 |
| 20.62 | 0.6650 | 0.2745 | 2.4227 |
| 20.87 | 0.6579 | 0.2796 | 2.3528 |
| 21.12 | 0.6479 | 0.2761 | 2.3465 |
| 21.37 | 0.6552 | 02840 | 2.3071 |
| 21.62 | 0.6372 | 0.2755 | 2.3127 |
| 21.87 | 0.6285 | 0.2759 | 2.2782 |
| 22.12 | 0.6767 | 0.2958 | 2.2878 |
| 22.37 | 0.6719 | 0.2983 | 2.2523 |
| 22.62 | 0.6689 | 03076 | 2.1748 |
| 22.87 | 0.6713 | 0.3036 | 2.2110 |
| 23.12 | 0.6610 | 0.3037 | 2.1767 |
| 23.37 | 0.6237 | 0.2979 | 2.0934 |
| 23.62 | 0.6749 | 0.3190 | 2.1154 |
| 23.87 | 0.6743 | 0.3239 | 2.0816 |
| 24.12 | 0.6801 | 0.3279 | 2.0739 |
| 24.37 | 0.6318 | 0.3065 | 2.0616 |
| 24.62 | 0.6186 | 0.3078 | 2.0096 |
| 24.87 | 0.6440 | 0.3270 | 1.9695 |
| 25.12 | 0.5852 | 02949 | 1.9843 |

Appendix E: Aerodynamic Data at Reynolds Number 120,000
Table E.1: Model \#1, $\mathrm{g}=1 \mathrm{c}, \mathrm{s}=0 \mathrm{c}$, Aerodynamic Data at Re 120,000

| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $C_{\text {D }}(-)$ | L/D (-) | a ( ${ }^{\circ}$ ) | CL (-) | $C_{0}(-)$ | L/D (-) | a $\left(^{\circ}\right.$ ) | $\mathrm{C}_{L}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.70 | -0.1018 | 0.0195 | -5.2293 | 7.30 | 0.5220 | 0.0812 | 6.4272 | 16.30 | 0.7840 | 0.2541 | 3.0849 |
| -1.45 | -0.0931 | 0.0158 | -5.8821 | 7.55 | 0.5484 | 0.0874 | 6.2758 | 16.55 | 0.7833 | 0.2578 | 3.0383 |
| -1.20 | -0.0720 | 0.0127 | -5.6646 | 7.80 | 0.5722 | 0.0938 | 6.1026 | 16.80 | 0.7743 | 0.2582 | 2.9985 |
| -0.95 | -0.0646 | 0.0139 | -4.6380 | 8.05 | 0.5875 | 0.0981 | 5.9874 | 17.05 | 0.7855 | 0.2673 | 2.9390 |
| -0.70 | -0.0570 | 0.0138 | -4.1208 | 8.30 | 0.5988 | 0.1023 | 5.8511 | 17.30 | 0.7852 | 0.2707 | 2.9010 |
| -0.45 | -0.0303 | 0.0136 | -2.2236 | 8.55 | 0.6191 | 0.1064 | 5.8198 | 17.55 | 0.7825 | 0.2750 | 2.8454 |
| -0.20 | -0.0148 | 0.0137 | -1.0840 | 8.80 | 0.6503 | 0.1156 | 5.6245 | 17.80 | 0.7801 | 0.2772 | 2.8140 |
| 0.05 | 0.0019 | 0.0132 | 0.1452 | 9.05 | 06449 | 0.1190 | 5.4200 | 18.05 | 0.7720 | 0.2780 | 2.7767 |
| 0.30 | 0.0261 | 0.0159 | 1.6375 | 9.30 | 0.6742 | 0.1282 | 5.2572 | 18.30 | 0.7748 | 0.2844 | 2.7238 |
| 0.55 | 0.0382 | 0.0166 | 2.2965 | 9.55 | 0.6942 | 0.1347 | 5.1545 | 18.55 | 0.7691 | 0.2859 | 2.6901 |
| 0.80 | 0.0565 | 0.0173 | 3.2734 | 9.80 | 0.7131 | 0.1418 | 5.0290 | 18.80 | 0.7653 | 0.2896 | 2.6427 |
| 1.05 | 0.0639 | 0.0175 | 3.6456 | 10.05 | 0.7310 | 0.1469 | 4.9750 | 19.05 | 0.7603 | 0.2892 | 2.6287 |
| 1.30 | 0.0862 | 0.0149 | 5.7888 | 10.30 | 0.7422 | 0.1534 | 4.8375 | 19.30 | 0.7501 | 0.2893 | 2.5928 |
| 1.55 | 0.1026 | 0.0146 | 7.0062 | 10.55 | 0.7449 | 0.1527 | 4.8788 | 19.55 | 0.7511 | 02938 | 2.5568 |
| 1.80 | 0.1230 | 0.0170 | 7.2382 | 10.80 | 0.7532 | 0.1619 | 4.6525 | 19.80 | 0.7431 | 0.2970 | 2.5020 |
| 2.05 | 0.1444 | 0.0191 | 7.5749 | 11.05 | 0.7636 | 0.1677 | 4.5545 | 20.05 | 0.7474 | 0.2999 | 2.4921 |
| 2.30 | 0.1517 | 0.0201 | 7.5332 | 11.30 | 0.7763 | 0.1743 | 4.4528 | 20.30 | 0.7377 | 0.3032 | 2.4334 |
| 2.55 | 0.1687 | 0.0219 | 7.7136 | 11.55 | 0.7869 | 0.1848 | 4.2575 | 20.55 | 0.7323 | 0.3027 | 2.4197 |
| 280 | 0.1827 | 0.0236 | 7.7522 | 11.80 | 0.7914 | 0.1819 | 4.3515 | 2080 | 0.7472 | 0.3126 | 2.3906 |
| 3.05 | 0.2134 | 0.0262 | 8.1348 | 12.05 | 0.8017 | 0.1930 | 4.1539 | 21.05 | 0.7350 | 0.3121 | 2.3555 |
| 3.30 | 0.2276 | 0.0278 | 8.1728 | 12.30 | 0.8074 | 0.1961 | 4.1177 | 21.30 | 0.7500 | 0.3220 | 2.3288 |
| 3.55 | 0.2461 | 0.0291 | 8.4637 | 12.55 | 0.8142 | 0.2035 | 4.0000 | 21.55 | 0.7463 | 0.3232 | 2.3092 |
| 3.80 | 0.2461 | 0.0291 | 8.4637 | 12.80 | 0.8182 | 0.2094 | 3.9078 | 21.80 | 0.7479 | 0.3269 | 2.2874 |
| 4.05 | 0.2618 | 0.0279 | 9.3987 | 13.05 | 0.8199 | 0.2118 | 3.8719 | 22.05 | 0.7355 | 0.3263 | 2.2543 |
| 430 | 0.2697 | 0.0325 | 8.3028 | 13.30 | 0.8197 | 0.2143 | 3.8259 | 22.30 | 0.7477 | 0.3361 | 2.2245 |
| 4.55 | 0.2998 | 0.0369 | 8.1215 | 13.55 | 0.8180 | 0.2201 | 3.7169 | 22.55 | 0.7526 | 0.3409 | 2.2078 |
| 4.80 | 0.3170 | 0.0399 | 7.9408 | 13.80 | 0.8252 | 0.2267 | 3.6397 | 2280 | 0.7489 | 0.3430 | 2.1832 |
| 5.05 | 0.3306 | 0.0408 | 8.1115 | 14.05 | 0.8171 | 0.2261 | 3.6140 | 23.05 | 0.7407 | 0.3441 | 2.1526 |
| 5.30 | 0.3446 | 0.0422 | 8.1726 | 14.30 | 0.8174 | 0.2341 | 3.4909 | 23.30 | 0.7518 | 0.3507 | 2.1435 |
| 5.55 | 0.4012 | 0.0496 | 8.0916 | 14.55 | 0.8141 | 0.2358 | 3.4521 | 23.55 | 0.7556 | 0.3593 | 2.1030 |
| 5.80 | 0.4159 | 0.0562 | 7.4060 | 14.80 | 0.8152 | 0.2390 | 3.4100 | 23.80 | 0.7600 | 0.3633 | 2.0919 |
| 6.05 | 0.4216 | 0.0589 | 7.1577 | 15.05 | 0.8073 | 0.2423 | 3.3316 | 24.05 | 0.7496 | 0.3625 | 2.0680 |
| 630 | 0.4475 | 0.0649 | 6.8948 | 15.30 | 0.8030 | 0.2453 | 3.2742 | 24.30 | 0.7579 | 0.3705 | 2.0453 |
| 6.55 | 0.4611 | 0.0661 | 6.9710 | 15.55 | 0.7936 | 0.2472 | 3.2099 | 24.55 | 0.7517 | 0.3726 | 2.0176 |
| 680 | 0.4879 | 0.0717 | 6.8042 | 15.80 | 0.7931 | 0.2483 | 3.1941 | 24.80 | 0.7607 | 0.3798 | 2.0028 |
| 7.05 | 0.4992 | 0.0772 | 6.4623 | 16.05 | 0.7826 | 0.2494 | 3.1374 | 2505 | 0.7467 | 0.3776 | 1.9776 |
|  |  |  |  |  |  |  |  | 25.30 | 0.7545 | 0.3860 | 1.9549 |

Table E.2: Model \#2, $g=0.5 c, s=0 c_{\text {, }}$ Aerodynamic Data at $\operatorname{Re} \mathbf{1 2 0 , 0 0 0}$

| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{CD}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.24 | -0.1139 | 0.0224 | -5.0781 |
| -1.99 | -0.1037 | 0.0218 | -4.7506 |
| -1.74 | -0.0829 | 0.0209 | -3.9744 |
| -1.49 | -0.0911 | 0.0226 | $-4.0280$ |
| . 1.24 | -0.0677 | 0.0205 | -3.3088 |
| -0.99 | -0.0512 | 0.0193 | -2.6524 |
| -0.74 | -0.0405 | 0.0211 | -1.9238 |
| -0.49 | -0.0290 | 0.0211 | -1.3788 |
| -0.24 | -0.0135 | 0.0207 | -0.6518 |
| 0.01 | 0.0005 | 0.0209 | 0.0256 |
| 0.26 | 0.0237 | 0.0209 | 1.1352 |
| 0.51 | 0.0229 | 0.0199 | 1.1516 |
| 0.76 | 0.0357 | 0.0209 | 1.7133 |
| 1.01 | 0.0507 | 0.0213 | 2.3867 |
| 1.26 | 0.0612 | 0.0224 | 2.7310 |
| 1.51 | 0.0815 | 0.0224 | 3.6334 |
| 1.76 | 0.0910 | 0.0236 | 3.8577 |
| 2.01 | 0.1071 | 0.0257 | 4.1645 |
| 2.26 | 0.1145 | 0.0251 | 4.5551 |
| 2.51 | 0.1313 | 0.0240 | 54746 |
| 2.76 | 0.1315 | 0.0263 | 4.9963 |
| 3.01 | 0.1630 | 0.0281 | 5.8076 |
| 3.26 | 0.1630 | 0.0281 | 5.8076 |
| 3.51 | 0.1682 | 0.0296 | 5.6798 |
| 3.76 | 0.1931 | 0.0308 | 6.2733 |
| 4.01 | 0.1873 | 0.0308 | 6.0837 |
| 4.26 | 0.2143 | 0.0335 | 6.3962 |
| 4.51 | 0.2338 | 0.0374 | 6.2510 |
| 4.76 | 0.2464 | 00401 | 6.1418 |
| 5.01 | 0.2669 | 0.0428 | 6.2281 |
| 5.26 | 0.2715 | 0.0444 | 6.1150 |
| 5.51 | 0.3029 | 0.0487 | 6.2208 |
| 5.76 | 0.3007 | 0.0483 | 6.2266 |
| 6.01 | 0.3377 | 0.0553 | 6.1064 |
| 6.26 | 0.3482 | 0.0563 | 6.1876 |
| 6.51 | 0.3752 | 0.0627 | 5.9853 |


| a ${ }^{\circ}{ }^{\circ}$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $C_{\text {D }}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.76 | 0.3877 | 0.0677 | 5.7222 |
| 7.01 | 0.4042 | 0.0691 | 5.8487 |
| 7.26 | 0.4110 | 0.0728 | 5.6453 |
| 7.51 | 0.4307 | 0.0777 | 5.5448 |
| 7.76 | 0.4394 | 0.0806 | 5.4526 |
| 8.01 | 0.4509 | 0.0839 | 5.3745 |
| 8.26 | 04554 | 0.0847 | 5.3779 |
| 8.51 | 0.4766 | 0.0911 | 5.2317 |
| 8.76 | 0.4980 | 0.0991 | 5.0265 |
| 9.01 | 0.4997 | 0.1000 | 4.9951 |
| 9.26 | 0.5266 | 0.1067 | 4.9371 |
| 9.51 | 0.5235 | 0.1072 | 4.8812 |
| 9.76 | 0.5418 | 0.1154 | 4.6941 |
| 10.01 | 0.5398 | 0.1170 | 4.6150 |
| 10.26 | 0.5585 | 0.1222 | 4.5694 |
| 10.51 | 0.5560 | 0.1228 | 4.5271 |
| 10.76 | 0.5610 | 0.1294 | 4.3348 |
| 11.01 | 0.5620 | 0.1319 | 4.2591 |
| 11.26 | 0.5742 | 0.1368 | 4.1973 |
| 11.51 | 0.5735 | 0.1395 | 4.1098 |
| 11.76 | 0.5791 | 0.1428 | 4.0541 |
| 12.01 | 0.5937 | 0.1495 | 3.9723 |
| 12.26 | 0.5902 | 0.1522 | 38782 |
| 12.51 | 0.6032 | 0.1570 | 3.8411 |
| 12.76 | 0.5978 | 0.1590 | 3.7599 |
| 13.01 | 0.6165 | 0.1666 | 3.7007 |
| 13.26 | 0.5993 | 0.1650 | 3.6319 |
| 13.51 | 0.6155 | 0.1732 | 3.5538 |
| 13.76 | 0.6268 | 0.1767 | 3.5472 |
| 14.01 | 0.6211 | 0.1790 | 3.4694 |
| 14.26 | 0.6299 | 0.1847 | 3.4109 |
| 14.51 | 0.6381 | 0.1905 | 3.3492 |
| 14.76 | 0.6462 | 0.1960 | 3.2978 |
| 15.01 | 0.6281 | 0.1915 | 32804 |
| 15.26 | 0.6394 | 0.1975 | 3.2374 |
| 15.51 | 0.6338 | 0.2012 | 3.1499 |


| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.76 | 0.6303 | 0.2059 | 3.0614 |
| 1601 | 0.6421 | 0.2103 | 3.0527 |
| 16.26 | 0.6406 | 0.2129 | 3.0091 |
| 16.51 | 0.6419 | 0.2148 | 2.9882 |
| 16.76 | 0.6474 | 0.2220 | 2.9159 |
| 17.01 | 0.6493 | 0.2267 | 2.8644 |
| 17.26 | 0.6476 | 0.2271 | 2.8518 |
| 17.51 | 0.6536 | 0.2323 | 2.8133 |
| 17.76 | 0.6538 | 0.2353 | 2.7792 |
| 18.01 | 0.6643 | 0.2456 | 2.7053 |
| 18.26 | 0.6591 | 0.2454 | 2.6860 |
| 18.51 | 0.6612 | 0.2496 | 2.6485 |
| 18.76 | 06637 | 0.2510 | 2.6443 |
| 19.01 | 0.6676 | 0.2586 | 2.5817 |
| 19.26 | 0.6655 | 0.2617 | 2.5428 |
| 19.51 | 0.6847 | 0.2699 | 2.5372 |
| 19.76 | 0.6678 | 0.2697 | 2.4763 |
| 20.01 | 0.6672 | 0.2710 | 2.4617 |
| 20.26 | 0.6735 | 0.2773 | 2.4289 |
| 20.51 | 06781 | 0.2804 | 2.4186 |
| 20.76 | 0.6682 | 0.2810 | 2.3782 |
| 21.01 | 06801 | 0.2880 | 2.3616 |
| 21.26 | 06744 | 0.2899 | 2.3263 |
| 21.51 | 0.6949 | 0.3020 | 2.3010 |
| 21.76 | 0.6791 | 0.2996 | 2.2664 |
| 22.01 | 0.6834 | 0.3008 | 2.2718 |
| 22.26 | 0.6717 | 0.3031 | 2.2158 |
| 22.51 | 0.6836 | 0.3123 | 2.1889 |
| 22.76 | 0.6865 | 0.3168 | 2.1672 |
| 23.01 | 0.6997 | 0.3273 | 2.1381 |
| 23.26 | 0.6791 | 0.3212 | 2.1140 |
| 2351 | 0.6684 | 0.3216 | 2.0782 |
| 23.76 | 0.6785 | 0.3284 | 2.0659 |
| 24.01 | 06935 | 0.3374 | 2.0555 |
| 24.26 | 0.6857 | 0.3374 | 2.0324 |
| 24.51 | 0.6882 | 0.3409 | 2.0190 |
| 24.76 | 0.6818 | 0.3434 | 1.9854 |

Table E.3: Model \#3, $\mathrm{g}=\mathbf{0 . 5} \mathrm{c}$, $\mathrm{s}=0.5 \mathrm{c}$, Aerodynamic Data at $\operatorname{Re} 120,000$

| a $\left(^{\circ}\right.$ ) | C. $(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.61 | -0.0919 | 0.0251 | -3.6626 |
| -1.36 | -0.0718 | 0.0233 | -3.0793 |
| -1.11 | -0.0607 | 0.0233 | -2.6040 |
| -0.86 | -0.0491 | 0.0234 | -2.0960 |
| -0.61 | -0.0394 | 0.0224 | -1.7603 |
| -0.36 | -0.0191 | 0.0224 | -0.8503 |
| -0.11 | -0.0068 | 0.0222 | -0.3051 |
| 0.14 | 0.0086 | 0.0223 | 0.3876 |
| 0.39 | 0.0273 | 0.0214 | 1.2765 |
| 0.64 | 0.0393 | 0.0211 | 1.8609 |
| 0.89 | 0.0604 | 0.0223 | 2.7125 |
| 1.14 | 0.0729 | 0.0222 | 3.2785 |
| 1.39 | 0.0796 | 0.0210 | 3.7980 |
| 1.64 | 0.0922 | 0.0226 | 4.0851 |
| 1.89 | 0.1127 | 0.0228 | 4.9319 |
| 2.14 | 0.1329 | 0.0229 | 5.8034 |
| 2.39 | 0.1409 | 0.0240 | 5.8784 |
| 2.64 | 0.1558 | 0.0259 | 6.0153 |
| 2.89 | 0.1691 | 0.0251 | 6.7461 |
| 3.14 | 0.1856 | 0.0257 | 7.2245 |
| 3.39 | 0.2038 | 0.0276 | 7.3712 |
| 3.64 | 0.2149 | 0.0294 | 7.3005 |
| 3.89 | 0.2149 | 0.0294 | 7.3005 |
| 4.14 | 0.2447 | 0.0310 | 7.9006 |
| 4.39 | 0.2499 | 0.0325 | 7.6981 |
| 4.64 | 0.2706 | 0.0353 | 7.6649 |
| 4.89 | 0.2807 | 0.0387 | 7.2488 |
| 5.14 | 0.2977 | 0.0399 | 7.4530 |
| 5.39 | 0.3110 | 0.0406 | 7.6602 |
| 5.64 | 0.3602 | 0.0498 | 7.2275 |
| 5.89 | 0.3710 | 0.0505 | 7.3413 |
| 6.14 | 0.3943 | 0.0534 | 7.3867 |
| 6.39 | 0.4052 | 0.0570 | 7.1134 |
| 6.64 | 0.4272 | 0.0597 | 7.1520 |
| 6.89 | 0.4383 | 0.0618 | 7.0919 |
| 7.14 | 0.4500 | 0.0656 | 6.8621 |


| $0\left({ }^{\circ}\right)$ | C. (-) | $\mathrm{CD}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.39 | 0.4693 | 0.0691 | 6.7867 |
| 7.64 | 0.4862 | 0.0738 | 6.5899 |
| 7.89 | 0.5052 | 00781 | 6.4658 |
| 8.14 | 0.5146 | 0.0827 | 6.2231 |
| 8.39 | 0.5311 | 0.0847 | 6.2682 |
| 8.64 | 0.5406 | 0.0890 | 6.0744 |
| 8.89 | 0.5643 | 0.0955 | 5.9116 |
| 9.14 | 0.5735 | 0.0978 | 5.8650 |
| 9.39 | 0.5963 | 0.1037 | 5.7496 |
| 9.64 | 0.6071 | 0.1078 | 5.6323 |
| 9.89 | 0.6163 | 0.1120 | 5.5030 |
| 10.14 | 0.6267 | 0.1149 | 5.4519 |
| 10.39 | 0.6427 | 0.1230 | 5.2271 |
| 10.64 | 0.6532 | 0.1261 | 5.1788 |
| 10.89 | 06534 | 0.1273 | 5.1315 |
| 11.14 | 0.6697 | 0.1327 | 5.0477 |
| 11.39 | 0.6752 | 0.1374 | 4.9128 |
| 11.64 | 0.6913 | 0.1440 | 4.8003 |
| 11.89 | 0.7006 | 0.1490 | 4.7009 |
| 12.14 | 0.7019 | 0.1511 | 4.6462 |
| 12.39 | 0.7176 | 0.1591 | 4.5096 |
| 12.64 | 0.7309 | 0.1645 | 4.4426 |
| 12.89 | 0.7336 | 0.1690 | 4.3416 |
| 13.14 | 0.7474 | 0.1737 | 4.3037 |
| 13.39 | 0.7459 | 0.1778 | 4.1951 |
| 13.64 | 0.7508 | 0.1793 | 4.1884 |
| 13.89 | 0.7487 | 0.1823 | 4.1064 |
| 14.14 | 0.7689 | 0.1928 | 3.9874 |
| 14.39 | 0.7700 | 0.1945 | 3.9597 |
| 14.64 | 0.7851 | 0.2009 | 3.9072 |
| 14.89 | 0.7762 | 0.2037 | 3.8098 |
| 15.14 | 0.7869 | 0.2083 | 3.7776 |
| 15.39 | 0.7927 | 0.2135 | 3.7134 |
| 15.64 | 0.7977 | 0.2198 | 3.6289 |
| 15.89 | 0.8041 | 0.2254 | 3.5673 |
| 16.14 | 0.8017 | 0.2284 | 3.5106 |


| $\alpha\left({ }^{\circ}\right)$ | C. ${ }^{(-)}$ | $\mathrm{Co}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.39 | 0.8123 | 0.2332 | 3.4837 |
| 16.64 | 0.8229 | 0.2383 | 3.4529 |
| 16.89 | 0.8181 | 0.2422 | 3.3780 |
| 17.14 | 0.8197 | 0.2463 | 3.3275 |
| 17.39 | 0.8286 | 0.2519 | 3.2889 |
| 17.64 | 0.8271 | 0.2557 | 3.2349 |
| 17.89 | 0.8357 | 0.2637 | 3.1687 |
| 18.14 | 0.8381 | 0.2683 | 3.1234 |
| 18.39 | 0.8296 | 0.2665 | 3.1133 |
| 18.64 | 08497 | 0.2775 | 3.0626 |
| 18.89 | 0.8499 | 0.2822 | 3.0118 |
| 19.14 | 0.8416 | 0.2841 | 2.9624 |
| 19.39 | 0.8353 | 0.2819 | 2.9636 |
| 19.64 | 0.8635 | 0.2973 | 2.9045 |
| 19.89 | 0.8562 | 0.2989 | 2.8643 |
| 20.14 | 0.8609 | 0.3045 | 2.8276 |
| 20.39 | 0.8826 | 0.3180 | 2.7753 |
| 20.64 | 0.8731 | 0.3173 | 2.7521 |
| 20.89 | 0.8509 | 0.3125 | 2.7230 |
| 21.14 | 0.8838 | 0.3267 | 2.7055 |
| 21.39 | 0.8821 | 0.3304 | 2.6701 |
| 21.64 | 08696 | 0.3308 | 2.6291 |
| 21.89 | 0.8772 | 0.3364 | 2.6076 |
| 22.14 | 0.8928 | 0.3473 | 2.5707 |
| 22.39 | 0.8887 | 0.3492 | 2.5448 |
| 22.64 | 0.8797 | 0.3492 | 2.5190 |
| 22.89 | 08879 | 0.3577 | 2.4820 |
| 23.14 | 0.8918 | 0.3617 | 2.4656 |
| 23.39 | 0.8848 | 0.3650 | 2.4243 |
| 23.64 | 0.8810 | 0.3662 | 2.4059 |
| 23.89 | 0.8791 | 0.3704 | 2.3737 |
| 24.14 | 0.8875 | 0.3758 | 2.3615 |
| 24.39 | 0.9131 | 0.3909 | 2.3361 |
| 24.64 | 0.8963 | 0.3873 | 2.3143 |
| 24.89 | 0.9226 | 0.4046 | 2.2799 |
| 25.14 | 0.9126 | 0.4050 | 2.2531 |
| 25.39 | 0.9232 | 0.4153 | 2.2230 |

Table E.4: Model \#4, $\mathbf{f = 0 . 5 c}$, $s=1 \mathrm{c}$, Aerodynamic Data at Re 120,000

| - ${ }^{\circ}$ ) | $\mathrm{CLL}_{(-)}$ | $\mathrm{CDD}^{\text {( }}$ - | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.13 | -0.1312 | 0.0235 | -5.5901 |
| -1.88 | -0.1204 | 0.0226 | -5.3193 |
| -1.63 | -0.1015 | 0.0212 | -4.7967 |
| -1.38 | -0.0907 | 0.0202 | -4.4922 |
| -1.13 | -0.0750 | 0.0192 | -3.9032 |
| -0.88 | -0.0602 | 0.0187 | -3.2240 |
| -0.63 | -0.0449 | 0.0199 | -2.2556 |
| -0.38 | -0.0280 | 0.0192 | -1.4538 |
| -0.13 | 00022 | 0.0184 | 0.1198 |
| 0.12 | 0.0101 | 0.0198 | 0.5087 |
| 0.37 | 0.0279 | 0.0190 | 1.4708 |
| 0.62 | 0.0398 | 0.0203 | 1.9631 |
| 0.87 | 0.0589 | 0.0208 | 2.8332 |
| 1.12 | 0.0724 | 0.0184 | 3.9466 |
| 1.37 | 0.0929 | 0.0206 | 4.5116 |
| 1.62 | 0.1026 | 0.0194 | 5.2750 |
| 1.87 | 0.1287 | 0.0221 | 5.8140 |
| 2.12 | 0.1440 | 0.0233 | 6.1716 |
| 2.37 | 0.1605 | 0.0221 | 7.2744 |
| 2.62 | 0.1646 | 0.0226 | 7.2697 |
| 2.87 | 0.1928 | 0.0258 | 7.4603 |
| 3.12 | 0.2194 | 0.0283 | 7.7660 |
| 3.37 | 0.2194 | 0.0283 | 7.7660 |
| 3.62 | 0.2179 | 0.0287 | 7.5803 |
| 3.87 | 0.2406 | 0.0317 | 7.5841 |
| 4.12 | 0.2542 | 0.0343 | 7.4153 |
| 4.37 | 0.2776 | 0.0355 | 7.8103 |
| 4.62 | 0.2907 | 0.0377 | 7.7155 |
| 4.87 | 0.3191 | 0.0397 | 8.0311 |
| 5.12 | 0.3523 | 0.0434 | 8.1237 |
| 5.37 | 0.3792 | 0.0482 | 7.8680 |
| 5.62 | 0.3895 | 0.0506 | 7.7012 |
| 5.87 | 0.4132 | 0.0546 | 7.5719 |
| 6.12 | 0.4313 | 0.0576 | 7.4865 |
| 6.37 | 0.4497 | 0.0629 | 7.1518 |
| 6.62 | 0.4686 | 0.0639 | 7.3380 |


| - $\left(^{\circ}\right.$ ) | $\mathrm{C}_{1}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.87 | 0.4928 | 0.0697 | 7.0661 |
| 7.12 | 0.5078 | 0.0736 | 6.9002 |
| 7.37 | 0.5259 | 0.0762 | 6.9001 |
| 7.62 | 0.5399 | 0.0809 | 6.6718 |
| 7.87 | 0.5565 | 0.0836 | 6.6601 |
| 8.12 | 0.5778 | 0.0932 | 6.1972 |
| 8.37 | 0.5894 | 0.0953 | 6.1840 |
| 8.62 | 0.6171 | 0.1012 | 6.0949 |
| 8.87 | 0.6263 | 0.1046 | 5.9866 |
| 9.12 | 0.6453 | 0.1104 | 5.8453 |
| 9.37 | 0.6570 | 0.1144 | 57448 |
| 9.62 | 0.6758 | 0.1208 | 5.5931 |
| 9.87 | 0.7063 | 0.1256 | 5.6235 |
| 10.12 | 0.7194 | 0.1336 | 5.3861 |
| 10.37 | 0.7353 | 0.1384 | 5.3124 |
| 10.62 | 0.7357 | 0.1433 | 5.1337 |
| 1087 | 0.7518 | 0.1490 | 5.0474 |
| 11.12 | 0.7684 | 0.1547 | 4.9671 |
| 11.37 | 0.7908 | 0.1625 | 4.8651 |
| 11.62 | 0.7961 | 0.1667 | 4.7762 |
| 11.87 | 0.8200 | 0.1724 | 4.7556 |
| 12.12 | 0.8336 | 0.1813 | 4.5979 |
| 12.37 | 0.8506 | 0.1876 | 4.5336 |
| 12.62 | 0.8606 | 0.1920 | 4.4818 |
| 12.87 | 0.8680 | 0.1977 | 4.3908 |
| 13.12 | 0.8781 | 0.2058 | 4.2676 |
| 13.37 | 0.8801 | 0.2104 | 4.1837 |
| 13.62 | 0.8912 | 0.2144 | 4.1566 |
| 13.87 | 0.8977 | 0.2211 | 4.0596 |
| 14.12 | 0.9121 | 0.2254 | 4.0469 |
| 14.37 | 0.9219 | 0.2347 | 3.9276 |
| 14.62 | 0.9261 | 0.2396 | 3.8647 |
| 14.87 | 0.9399 | 0.2477 | 3.7939 |
| 15.12 | 0.9468 | 0.2540 | 3.7284 |
| 15.37 | 0.9473 | 0.2539 | 3.7319 |
| 15.62 | 0.9424 | 0.2588 | 36418 |


| $0\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.87 | 0.9503 | 0.2668 | 3.5614 |
| 16.12 | 0.9538 | 0.2709 | 3.5204 |
| 16.37 | 0.9648 | 0.2780 | 3.4701 |
| 16.62 | 0.9619 | 0.2850 | 3.3748 |
| 16.87 | 0.9718 | 0.2856 | 3.4033 |
| 17.12 | 0.9733 | 0.2922 | 3.3309 |
| 17.37 | 0.9757 | 0.3000 | 3.2526 |
| 17.62 | 0.9697 | 0.3038 | 3.1923 |
| 17.87 | 0.9823 | 0.3042 | 3.2292 |
| 18.12 | 0.9854 | 0.3120 | 3.1579 |
| 18.37 | 0.9929 | 0.3183 | 3.1193 |
| 18.62 | 0.9867 | 0.3233 | 3.0517 |
| 18.87 | 0.9865 | 0.3259 | 3.0267 |
| 19.12 | 09980 | 0.3381 | 2.9523 |
| 19.37 | 0.9921 | 0.3392 | 2.9251 |
| 19.62 | 1.0070 | 0.3492 | 2.8833 |
| 19.87 | 0.9898 | 0.3474 | 2.8494 |
| 20.12 | 1.0010 | 0.3545 | 2.8237 |
| 20.37 | 0.9997 | 0.3601 | 2.7764 |
| 20.62 | 0.9965 | 0.3635 | 2.7415 |
| 20.87 | 0.9962 | 0.3702 | 2.6906 |
| 21,12 | 1.0016 | 0.3742 | 2.6765 |
| 21.37 | 0.9959 | 0.3783 | 2.6326 |
| 21.62 | 1.0156 | 0.3886 | 2.6137 |
| 21.87 | 1.0073 | 0.3913 | 2.5743 |
| 22.12 | 1.0046 | 0.3964 | 2.5346 |
| 22.37 | 10056 | 0.3995 | 2.5171 |
| 22.62 | 1.0155 | 0.4066 | 2.4978 |
| 22.87 | 1.0223 | 0.4166 | 2.4541 |
| 23.12 | 1.0333 | 0.4233 | 2.4410 |
| 23.37 | 1.0303 | 0.4300 | 2.3960 |
| 23.62 | 10237 | 0.4359 | 2.3487 |
| 23.87 | 1.0340 | 0.4420 | 2.3394 |
| 24.12 | 1.0249 | 0.4446 | 2.3050 |
| 24.37 | 1.0339 | 0.4508 | 2.2933 |
| 24.62 | 1.0367 | 0.4584 | 2.2614 |
| 24.87 | 1.0422 | 0.4666 | 2.2337 |

Table E.5: Model \#5, $\|=1 \mathrm{c}, \mathrm{s}=\mathbf{0 . 5 c}$, Aerodynamic Data at Re $\mathbf{1 2 0 , 0 0 0}$

| a $\left({ }^{\circ}\right)$ | CL (-) | $\mathrm{CD}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.76 | -0.0987 | 0.0331 | -2.9805 |
| -1.51 | -0.1015 | 0.0329 | -3.0836 |
| -1.26 | -0.0795 | 0.0318 | -2.5023 |
| -1.01 | -0.0623 | 0.0299 | -2.0820 |
| -0.76 | -0.0482 | 0.0297 | -1.6222 |
| -0.51 | -0.0294 | 0.0305 | -0.9653 |
| -0.26 | -0.0152 | 0.0288 | -0.5257 |
| -0.01 | 0.0032 | 0.0298 | 0.1069 |
| 0.25 | 0.0155 | 0.0286 | 0.5422 |
| 0.50 | 0.0352 | 0.0301 | 1.1707 |
| 0.75 | 0.0559 | 0.0315 | 1.7726 |
| 1.00 | 0.0805 | 0.0302 | 2.6684 |
| 1.25 | 0.0726 | 0.0304 | 2.3896 |
| 1.50 | 00946 | 0.0315 | 2.9988 |
| 1.75 | 0.1070 | 0.0308 | 3.4685 |
| 2.00 | 0.1313 | 0.0316 | 4.1512 |
| 2.25 | 0.1536 | 0.0326 | 4.7147 |
| 2.50 | 0.1661 | 0.0333 | 4.9809 |
| 2.75 | 0.1819 | 0.0355 | 5.1262 |
| 3.00 | 0.1969 | 0.0358 | 5.5006 |
| 3.25 | 0.2151 | 0.0389 | 5.5264 |
| 3.50 | 0.2404 | 0.0396 | 6.0699 |
| 3.75 | 0.2404 | 0.0396 | 6.0699 |
| 4.00 | 0.2570 | 0.0408 | 6.2961 |
| 4.25 | 0.2728 | 0.0439 | 6.2092 |
| 4.50 | 0.2886 | 0.0451 | 6.3996 |
| 4.75 | 0.3075 | 0.0480 | 6.4049 |
| 500 | 0.3348 | 0.0494 | 6.7798 |
| 5.25 | 0.3556 | 0.0530 | 6.7051 |
| 5.50 | 0.4001 | 0.0577 | 6.9390 |
| 5.75 | 0.4121 | 0.0628 | 6.5622 |
| 6.00 | 0.4330 | 0.0632 | 6.8464 |
| 6.25 | 0.4454 | 0.0647 | 6.8881 |
| 6.50 | 0.4729 | 0.0690 | 6.8569 |
| 6.75 | 0.4965 | 0.0760 | 6.5351 |
| 7.00 | 0.5137 | 0.0802 | 6.4047 |


| a ( ${ }^{\circ}$ ) | Cli (-) | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.25 | 0.5318 | 0.0859 | 6.1910 |
| 7.50 | 0.5465 | 0.0862 | 6.3395 |
| 7.75 | 0.5787 | 0.0919 | 6.2975 |
| 8.00 | 0.5956 | 0.0998 | 59665 |
| 8.25 | 0.6138 | 0.1022 | 6.0064 |
| 8.50 | 0.6220 | 0.1050 | 5.9231 |
| 8.75 | 0.6547 | 0.1120 | 5.8454 |
| 9.00 | 0.6615 | 0.1155 | 5.7275 |
| 9.25 | 0.6974 | 0.1247 | 5.5944 |
| 9.50 | 0.6955 | 0.1286 | 5.4073 |
| 9.75 | 0.7331 | 0.1348 | 5.4387 |
| 10.00 | 0.7267 | 0.1371 | 5.3016 |
| 10.25 | 0.7585 | 0.1467 | 5.1697 |
| 10.50 | 0.7668 | 0.1496 | 5.1269 |
| 10.75 | 0.7903 | 0.1590 | 4.9705 |
| 11.00 | 0.7943 | 0.1608 | 4.9397 |
| 11.25 | 0.8138 | 0.1667 | 4.8822 |
| 11.50 | 0.8161 | 0.1718 | 4.7503 |
| 11.75 | 0.8220 | 0.1759 | 46732 |
| 12.00 | 0.8319 | 0.1823 | 4.5635 |
| 12.25 | 0.8369 | 0.1909 | 4.3844 |
| 12.50 | 0.8328 | 0.1894 | 4.3977 |
| 12.75 | 0.8389 | 0.1953 | 4.2944 |
| 13.00 | 0.8346 | 0.1985 | 4.2053 |
| 13.25 | 0.8375 | 0.2040 | 4.1045 |
| 13.50 | 0.8463 | 0.2084 | 4.0610 |
| 13.75 | 0.8434 | 0.2107 | 4.0024 |
| 14.00 | 0.8447 | 0.2168 | 3.8959 |
| 14.25 | 0.8435 | 0.2183 | 3.8648 |
| 14.50 | 0.8384 | 0.2221 | 3.7751 |
| 14.75 | 0.8311 | 0.2225 | 3.7361 |
| 15.00 | 0.8499 | 0.2298 | 3.6990 |
| 15.25 | 0.8389 | 0.2321 | 3.6138 |
| 15.50 | 0.8465 | 0.2384 | 3.5505 |
| 15.75 | 0.8413 | 02409 | 3.4924 |
| 16.00 | 0.8455 | 0.2457 | 3.4417 |


| $\mathrm{a}\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.25 | 0.8578 | 0.2515 | 3.4107 |
| 16.50 | 0.8456 | 0.2541 | 3.3277 |
| 16.75 | 0.8563 | 0.2596 | 3.2986 |
| 17.00 | 0.8574 | 0.2636 | 3.2533 |
| $\$ 7.25$ | 0.8615 | 0.2700 | 3.1911 |
| 17.50 | 0.8561 | 0.2708 | 3.1617 |
| 17.75 | 0.8472 | 0.2711 | 3.1254 |
| 1800 | 0.8587 | 0.2809 | 3.0572 |
| 18.25 | 0.8522 | 0.2802 | 3.0410 |
| 18.50 | 0.8565 | 0.2856 | 2.9992 |
| 18.75 | 0.8582 | 0.2905 | 2.9545 |
| 19.00 | 0.8681 | 0.2976 | 2.9170 |
| 19.25 | 0.8533 | 0.2959 | 2.8842 |
| 19.50 | 0.8615 | 0.3050 | 2.8246 |
| 19.75 | 0.8515 | 0.3054 | 2.7881 |
| 20.00 | 0.8579 | 0.3114 | 2.7555 |
| 20.25 | 0.8588 | 0.3175 | 2.7052 |
| 20.50 | 0.8582 | 0.3191 | 2.6900 |
| 20.75 | 0.8633 | 0.3234 | 2.6696 |
| 21.00 | 0.8800 | 0.3346 | 2.6303 |
| 21.25 | 0.8712 | 0.3346 | 2.6042 |
| 21.50 | 0.8820 | 0.3422 | 2.5774 |
| 21.75 | 0.8776 | 0.3457 | 2.5387 |
| 22.00 | 0.8705 | 0.3448 | 2.5244 |
| 22.25 | 0.8858 | 0.3553 | 2.4934 |
| 22.50 | 0.8893 | 0.3602 | 2.4687 |
| 22.75 | 0.8954 | 0.3697 | 2.4222 |
| 23.00 | 0.8928 | 0.3710 | 2.4065 |
| 23.25 | 0.9071 | 0.3791 | 2.3930 |
| 23.50 | 0.8872 | 0.3760 | 2.3599 |
| 23.75 | 0.8838 | 0.3812 | 2.3186 |
| 24.00 | 0.8890 | 0.3863 | 2.3011 |
| 24.25 | 0.8948 | 0.3921 | 2.2821 |
| 24.50 | 0.9081 | 0.4026 | 2.2558 |
| 24.75 | 0.8954 | 0.4005 | 2.2357 |
| 25.00 | 0.9028 | 0.4089 | 2.2080 |
| 25.25 | 0.9029 | 0.4141 | 2.1803 |

Table E.6: Model \#6, $g=1 c, s=1 c$, Aerodynamic Data at Re 120,000

| a ( ${ }^{\circ}$ ) | CL (-) | CD ( - ) | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.90 | -0. 1435 | 0.0248 | -5.7852 |
| -1.65 | -0.1193 | 0.0227 | -5.2446 |
| -1.40 | -0.1061 | 0.0230 | -4.6042 |
| -1.15 | -0.0901 | 0.0221 | -4.0785 |
| -0.90 | -0.0731 | 0.0215 | -3.4038 |
| -0.65 | -0.0523 | 0.0208 | -2.5157 |
| -0.40 | -0.0331 | 0.0204 | -1.6220 |
| -0.15 | -0.0074 | 0.0197 | -0.3745 |
| 0.10 | 0.0019 | 0.0209 | 0.0924 |
| 0.35 | 0.0248 | 0.0191 | 1.2999 |
| 0.60 | 0.0378 | 0.0272 | 1.7043 |
| 0.85 | 0.0476 | 0.0199 | 2.3870 |
| 1.10 | 0.0773 | 0.0216 | 3.5820 |
| 1.35 | 0.0898 | 0.0208 | 4.3095 |
| 1.60 | 0.1006 | 0.0217 | 4.6330 |
| 1.85 | 0.1274 | 0.0222 | 5.7282 |
| 2.10 | 0.1350 | 0.0232 | 5.8096 |
| 2.35 | 0.1550 | 0.0242 | 6.4194 |
| 2.60 | 0.1631 | 0.0253 | 6.4526 |
| 2.85 | 0.1832 | 0.0270 | 6.7775 |
| 3.10 | 0.2116 | 0.0271 | 7.8045 |
| 3.35 | 0.2227 | 0.0293 | 7.6122 |
| 360 | 0.2227 | 0.0293 | 7.6122 |
| 3.85 | 0.2409 | 0.0286 | 8.4358 |
| 4.10 | 0.2704 | 0.0323 | 8.3758 |
| 4.35 | 0.2880 | 0.0358 | 8.0506 |
| 4.60 | 0.3017 | 0.0359 | 8.4076 |
| 4.85 | 0.3225 | 0.0394 | 8.1850 |
| 5.10 | 0.3479 | 0.0427 | 8.1438 |
| 5.35 | 0.3726 | 0.0455 | 8.1889 |
| 5.60 | 0.4034 | 0.0499 | 8.0811 |
| 5.85 | 0.4245 | 0.0532 | 7.9793 |
| 6.10 | 0.4451 | 0.0581 | 7.6580 |
| 6.35 | 0.4768 | 0.0617 | 7.7264 |
| 6.60 | 0.4906 | 0.0650 | 7.5460 |
| 6.85 | 0.5197 | 0.0689 | 7.5381 |


| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $C_{D}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.10 | 0.5381 | 0.0711 | 7.5691 |
| 7.35 | 0.5451 | 0.0768 | 7.0952 |
| 7.60 | 0.5735 | 0.0811 | 7.0694 |
| 7.85 | 0.5905 | 0.0918 | 6.4300 |
| 8.10 | 0.6289 | 0.0951 | 6.6106 |
| 8.35 | 0.6562 | 0.1041 | 63040 |
| 8.60 | 0.6622 | 0.1042 | 6.3547 |
| 885 | 0.6782 | 0.1104 | 6.1403 |
| 9.10 | 0.7037 | 0.1109 | 6.3467 |
| 9.35 | 0.7226 | 0.1196 | 6.0438 |
| 9.60 | 0.7359 | 0.1308 | 5.6255 |
| 9.85 | 0.7462 | 0.1301 | 5.7362 |
| 10.10 | 0.7762 | 0.1413 | 5.4944 |
| 10.35 | 0.8042 | 0.1445 | 5.5665 |
| 10.60 | 0.8146 | 0.1500 | 5.4296 |
| 10.85 | 0.8319 | 0.1579 | 5.2690 |
| 11.10 | 0.8451 | 0.1633 | 5.1763 |
| 11.35 | 0.8510 | 0.1674 | 5.0830 |
| 11.60 | 0.8579 | 0.1741 | 4.9281 |
| \$1.85 | 0.8748 | 0.1821 | 4.8033 |
| $\pm 2.10$ | 0.8678 | 0.1826 | 4.7532 |
| 12.35 | 0.8650 | 0.1906 | 4.5394 |
| 12.60 | 0.8712 | 0.1920 | 4.5384 |
| 12.85 | 0.8790 | 0.1995 | 4.4063 |
| 13.10 | 0.8838 | 0.2044 | 4.3237 |
| 13.35 | 0.8835 | 0.2106 | 4.1946 |
| 13.60 | 0.8755 | 0.2113 | 4.1434 |
| 13.85 | 0.8778 | 0.2148 | 4.0875 |
| 14.10 | 0.8731 | 0.2158 | 4.0465 |
| 14.35 | 0.8661 | 0.2177 | 3.9785 |
| 14.60 | 0.8710 | 0.2259 | 3.8553 |
| 14.85 | 0.8681 | 0.2274 | 3.8177 |
| 15.10 | 0.8736 | 0.2315 | 3.7739 |
| 15.35 | 0.8605 | 0.2347 | 3.6659 |
| 15.60 | 0.8842 | 0.2430 | 3.6387 |
| 15.85 | 0.8769 | 0.2431 | 3.6070 |


| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.10 | 0.8870 | 0.2517 | 3.5244 |
| 16.35 | 0.8769 | 0.2523 | 3.4751 |
| 16.60 | 0.8760 | 0.2553 | 3.4312 |
| 16.85 | 0.8756 | 0.2606 | 3.3603 |
| 17.10 | 0.8853 | 0.2665 | 3.3222 |
| 17.35 | 0.8856 | 0.2717 | 3.2598 |
| 17.60 | 0.8870 | 0.2745 | 3.2313 |
| 17.85 | 0.8818 | 0.2785 | 3.1666 |
| 18.10 | 0.8809 | 0.2827 | 3.1162 |
| 18.35 | 0.8848 | 0.2883 | 3.0693 |
| 18.60 | 0.8847 | 0.2908 | 3.0420 |
| 18.85 | 0.8857 | 0.2971 | 2.9812 |
| 19.10 | 0.8932 | 0.3009 | 2.9683 |
| 19.35 | 08838 | 0.3056 | 2.8916 |
| 19.60 | 0.8869 | 0.3074 | 2.8849 |
| 19.85 | 0.8890 | 0.3152 | 2.8202 |
| 20.10 | 0.8929 | 0.3207 | 2.7841 |
| 20.35 | 0.8961 | 0.3231 | 2.7732 |
| 20.60 | 0.9005 | 0.3291 | 2.7360 |
| 20.85 | 0.8991 | 0.3322 | 2.7064 |
| 21.10 | 0.9056 | 0.3382 | 2.6777 |
| 21.35 | 0.8919 | 0.3393 | 2.6290 |
| 21.60 | 0.9161 | 0.3479 | 2.6337 |
| 21.85 | 0.9184 | 0.3537 | 2.5965 |
| 22.10 | 0.8975 | 0.3520 | 2.5496 |
| 22.35 | 0.8930 | 0.3547 | 2.5178 |
| 22.60 | 09041 | 0.3634 | 2.4877 |
| 22.85 | 0.9045 | 0.3688 | 2.4527 |
| 23.10 | 0.9185 | 0.3783 | 2.4276 |
| 23.35 | 0.9227 | 0.3853 | 2.3948 |
| 23.60 | 0.9052 | 0.3816 | 2.3717 |
| 23.85 | 0.9270 | 0.3974 | 2.3327 |
| 24.10 | 0.9335 | 0.4006 | 2.3303 |
| 24.35 | 0.9315 | 0.4046 | 2.3021 |
| 24.60 | 0.9295 | 0.4093 | 2.2710 |
| 24.85 | 0.9222 | 0.4095 | 2.2521 |
| 25.10 | 0.9281 | 0.4184 | 2.2183 |

Table E.7: Model \#7, $\mathrm{g}=\mathbf{2 c} \mathrm{c}, \mathrm{s}=1 \mathrm{c}$, Aerodynamic Data at Re 120,000

| $\alpha\left({ }^{\circ}\right)$ | $C_{L}(-)$ | $C_{D}(-)$ | L/D $(-)$ |
| :---: | :---: | :---: | :---: |
| -2.08 | -0.1552 | 0.0281 | -5.5224 |
| -1.83 | -0.1349 | 0.0288 | -4.6813 |
| -1.58 | -0.1187 | 0.0288 | -4.1197 |
| -1.33 | -0.1010 | 0.0285 | -3.5473 |
| -1.08 | -0.0738 | 0.0281 | -2.6282 |
| -0.83 | -0.0571 | 0.0278 | -2.0542 |
| -0.58 | -0.0401 | 0.0263 | -1.5222 |
| -0.33 | -0.0151 | 0.0280 | -0.5416 |
| -0.08 | -0.0025 | 0.0284 | -0.0867 |
| 0.17 | 0.0081 | 0.0286 | 0.2825 |
| 0.42 | 0.0562 | 0.0282 | 1.9923 |
| 0.67 | 0.0607 | 0.0275 | 2.2110 |
| 0.92 | 0.0835 | 0.0276 | 3.0274 |
| 1.17 | 0.1082 | 0.0287 | 3.7687 |
| 1.42 | 0.1192 | 0.0284 | 4.1923 |
| 1.67 | 0.1496 | 0.0273 | 5.4835 |
| 1.92 | 0.1495 | 0.0295 | 5.0616 |
| 2.17 | 0.1870 | 0.0303 | 6.1719 |
| 2.42 | 0.2034 | 0.0326 | 6.2334 |
| 2.67 | 0.2206 | 0.0320 | 6.9023 |
| 2.92 | 0.2351 | 0.0340 | 6.9079 |
| 3.17 | 0.2627 | 0.0368 | 7.1408 |
| 3.42 | 0.2627 | 0.0368 | 7.1408 |
| 3.67 | 0.2954 | 0.0383 | 7.7220 |
| 3.92 | 0.3200 | 0.0409 | 7.8221 |
| 4.17 | 0.3329 | 0.0415 | 8.0164 |
| 4.42 | 0.3540 | 0.0446 | 7.9329 |
| 4.67 | 0.3948 | 0.0488 | 8.0913 |
| 4.92 | 0.4056 | 0.0496 | 8.1741 |
| 5.17 | 0.4501 | 0.0568 | 7.9215 |
| 5.42 | 0.4703 | 0.0606 | 7.7654 |
| 5.67 | 0.5080 | 0.0644 | 7.8900 |
| 5.92 | 0.5340 | 0.0700 | 7.6318 |
| 6.17 | 0.5629 | 0.0730 | 7.7123 |
| 6.42 | 0.5735 | 0.0760 | 7.5490 |
| 6.67 | 0.6064 | 0.0821 | 7.3882 |


| $\sigma\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | $\mathrm{L} / \mathrm{D}(-)$ |
| :---: | :---: | :---: | :---: |
| 6.92 | 0.6395 | 0.0874 | 7.3151 |
| 7.17 | 0.6555 | 0.0923 | 7.1017 |
| 7.42 | 0.6731 | 0.0952 | 7.0702 |
| 7.67 | 0.6937 | 0.1006 | 6.8954 |
| 7.92 | 0.7171 | 0.1083 | 6.6212 |
| 8.17 | 0.7276 | 0.1120 | 6.4962 |
| 8.42 | 0.7484 | 0.1186 | 6.3111 |
| 8.67 | 0.7589 | 0.1245 | 6.0959 |
| 8.92 | 0.7675 | 0.1290 | 5.9507 |
| 9.17 | 0.7817 | 0.1339 | 5.8370 |
| 9.42 | 0.7926 | 0.1405 | 5.6394 |
| 9.67 | 0.7908 | 0.1412 | 5.6023 |
| 9.92 | 0.8000 | 0.1472 | 5.4335 |
| 10.17 | 0.8073 | 0.1521 | 5.3064 |
| 1042 | 0.8120 | 0.1565 | 5.1891 |
| 10.67 | 0.8139 | 0.1601 | 5.0828 |
| 10.92 | 0.8246 | 0.1645 | 5.0124 |
| 11.17 | 0.8124 | 0.1679 | 4.8373 |
| 11.42 | 0.8282 | 0.1737 | 4.7678 |
| 11.67 | 0.8260 | 0.1737 | 4.7540 |
| 11.92 | 0.8278 | 0.1741 | 4.7554 |
| 12.17 | 0.8196 | 0.1662 | 4.9316 |
| 12.42 | 0.8313 | 0.1677 | 4.9562 |
| 12.67 | 0.8305 | 0.1682 | 4.9371 |
| 12.92 | 0.8300 | 0.1765 | 4.7013 |
| 13.17 | 0.8302 | 0.1675 | 4.9569 |
| 13.42 | 0.8440 | 0.1796 | 4.6985 |
| 13.67 | 0.8167 | 0.1724 | 4.7381 |
| 13.92 | 0.8275 | 0.1749 | 4.7311 |
| 14.17 | 0.8432 | 0.1778 | 4.7426 |
| 14.42 | 0.8298 | 0.1751 | 4.7402 |
| 14.67 | 0.8437 | 0.1857 | 4.5442 |
| 14.92 | 0.8530 | 0.1946 | 4.3833 |
| 15.17 | 0.8259 | 0.1884 | 4.3829 |
| 15.42 | 0.8371 | 0.1925 | 4.3491 |
| 15.67 | 0.8390 | 0.2081 | 4.0326 |
|  |  |  |  |


| a $\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.92 | 0.8499 | 0.2111 | 4.0251 |
| 16.17 | 0.8636 | 0.2090 | 4.1321 |
| 16.42 | 0.8599 | 0.2095 | 4.1036 |
| 16.67 | 0.8509 | 0.2181 | 3.9021 |
| 16.92 | 0.8665 | 0.2303 | 3.7619 |
| 17.17 | 0.8753 | 0.2438 | 3.5897 |
| 17.42 | 0.8554 | 0.2416 | 3.5405 |
| 17.67 | 0.8702 | 0.2348 | 3.7070 |
| 17.92 | 0.8849 | 0.2445 | 3.6188 |
| 18.17 | 0.8635 | 0.2556 | 3.3781 |
| 18.42 | 0.8514 | 0.2568 | 3.3158 |
| 18.67 | 0.8824 | 0.2695 | 3.2739 |
| 18.92 | 0.8791 | 0.2674 | 3.2877 |
| 19.17 | 0.8621 | 0.2596 | 3.3207 |
| 19.42 | 0.8819 | 0.2707 | 3.2574 |
| 19.67 | 0.8765 | 0.2886 | 3.0376 |
| 19.92 | 0.8589 | 0.2697 | 3.1841 |
| 20.17 | 0.8571 | 0.2722 | 3.1494 |
| 20.42 | 0.8804 | 0.2841 | 3.0986 |
| 20.67 | 0.8632 | 0.2870 | 3.0080 |
| 20.92 | 0.8588 | 0.2827 | 3.0384 |
| 21.17 | 0.8694 | 0.2879 | 3.0195 |
| 21.42 | 0.8682 | 0.2945 | 2.9476 |
| 21.67 | 0.8456 | 0.2971 | 2.8465 |
| 21.92 | 0.8888 | 0.3006 | 2.9565 |
| 22.17 | 0.8564 | 0.3052 | 2.8065 |
| 22.42 | 0.8651 | 0.3146 | 2.7496 |
| 22.67 | 0.8723 | 0.3100 | 2.8138 |
| 22.92 | 0.8767 | 0.3297 | 2.6593 |
| 23.17 | 0.8509 | 0.3246 | 2.6218 |
| 23.42 | 0.8867 | 0.3301 | 2.6864 |
| 23.67 | 0.8909 | 0.3205 | 2.7799 |
| 23.92 | 0.9021 | 0.3362 | 2.6833 |
| 24.17 | 0.9071 | 0.3416 | 2.6550 |
| 24.42 | 0.9090 | 0.3577 | 2.5415 |
| 24.67 | 0.8861 | 0.3476 | 2.5493 |
| 24.92 | 0.8881 | 0.3533 | 2.5135 |

Table E.8: Model \#8, $\mathrm{g}=1 \mathrm{c}$, $\mathrm{s}=1.5 \mathrm{c}$, Aerodynamic Data at Re 120,000

| a ( ${ }^{\circ}$ ) | $C_{L}(-)$ | $\mathrm{CD}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.07 | -0.1469 | 0.0257 | -5.7056 |
| -1.82 | -0.1309 | 0.0258 | -5.0727 |
| -1.57 | -0.1126 | 0.0246 | -4.5790 |
| -1.32 | -0.0991 | 0.0240 | -4.1361 |
| -1.07 | -0.0804 | 0.0226 | -3.5593 |
| -0.82 | -0.0552 | 0.0210 | -2.6265 |
| -0.57 | -0.0583 | 0.0207 | -2.8194 |
| -0.32 | -0.0394 | 0.0225 | -1.7524 |
| -0.07 | -0.0093 | 0.0227 | -0.4116 |
| 0.19 | 0.0174 | 0.0226 | 0.7704 |
| 0.44 | 0.0320 | 0.0232 | 1.3756 |
| 0.69 | 0.0517 | 0.0224 | 2.3075 |
| 0.94 | 0.0572 | 0.0225 | 2.5363 |
| 1.19 | 0.0833 | 0.0217 | 3.8402 |
| 1.44 | 0.0995 | 0.0241 | 4.1352 |
| 1.69 | 0.1248 | 0.0231 | 5.4102 |
| 1.94 | 0.1384 | 0.0252 | 5.4841 |
| 2.19 | 0.1647 | 0.0250 | 65881 |
| 2.44 | 0.1788 | 0.0255 | 7.0114 |
| 269 | 0.1978 | 0.0267 | 7.4096 |
| 2.94 | 0.2122 | 0.0275 | 7.7073 |
| 3.19 | 0.2384 | 0.0304 | 7.8404 |
| 3.44 | 0.2384 | 0.0304 | 7.8404 |
| 369 | 0.2679 | 0.0322 | 8.3289 |
| 3.94 | 0.2835 | 0.0336 | 8.4321 |
| 4.19 | 0.3081 | 0.0363 | 8.4768 |
| 4.44 | 0.3186 | 0.0392 | 8.1364 |
| 4.69 | 0.3305 | 0.0401 | 8.2429 |
| 4.94 | 0.3563 | 0.0447 | 7.9649 |
| 5.19 | 0.4001 | 0.0491 | 8.1470 |
| 5.44 | 0.4201 | 0.0515 | 8.1605 |
| 5.69 | 0.4450 | 0.0568 | 7.8319 |
| 5.94 | 0.4695 | 0.0599 | 7.8340 |
| 6.19 | 0.4802 | 0.0628 | 7.6525 |
| 6.44 | 0.5031 | 0.0659 | 7.6372 |
| 6.69 | 0.5456 | 0.0744 | 7.3367 |


| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(\cdot)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.94 | 0.5573 | 0.0762 | 7.3181 |
| 7.19 | 0.5862 | 0.0826 | 7.0924 |
| 7.44 | 0.6036 | 0.0874 | 6.9087 |
| 7.69 | 0.6005 | 0.0880 | 6.8278 |
| 7.94 | 0.6392 | 0.0959 | 6.6652 |
| 8.19 | 0.6514 | 0.1000 | 6.5141 |
| 8.44 | 0.6679 | 0.1061 | 6.2949 |
| 8.69 | 0.6986 | 0.1137 | 6.1433 |
| 8.94 | 0.7090 | 0.1170 | 6.0611 |
| 9.19 | 0.7489 | 0.1258 | 5.9508 |
| 9.44 | 0.7624 | 0.1329 | 5.7374 |
| 9.69 | 0.7757 | 0.1367 | 5.6758 |
| 9.94 | 0.8069 | 0.1446 | 5.5813 |
| 10.19 | 0.8168 | 0.1492 | 5.4740 |
| 10.44 | 0.8418 | 0.1565 | 5.3792 |
| 10.69 | 0.8508 | 0.1603 | 5.3068 |
| 10.94 | 0.8597 | 0.1691 | 5.0848 |
| 1119 | 0.8810 | 0.1766 | 4.9898 |
| 1144 | 0.8941 | 0.1807 | 4.9474 |
| 11.69 | 0.9027 | 0.1850 | 4.8808 |
| 11.94 | 0.9169 | 0.1937 | 4.7322 |
| 12.19 | 0.9074 | 0.1955 | 4.6424 |
| 12.44 | 0.9166 | 0.2022 | 4.5340 |
| 12.69 | 0.9228 | 0.2067 | 4.4646 |
| 12.94 | 0.9251 | 0.2103 | 4.3992 |
| 13.19 | 0.9252 | 0.2179 | 4.2469 |
| 13.44 | 0.9179 | 0.2195 | 4.1813 |
| 13.69 | 0.9118 | 0.2212 | 4.1227 |
| 13.94 | 0.9226 | 0.2250 | 4.1003 |
| 14.19 | 0.8999 | 0.2275 | 3.9551 |
| 14.44 | 0.9091 | 0.2327 | 3.9064 |
| 14.69 | 0.9037 | 0.2359 | 3.8305 |
| 14.94 | 0.8979 | 0.2381 | 3.7715 |
| 1519 | 0.9015 | 0.2440 | 3.6942 |
| 1544 | 0.9025 | 0.2453 | 3.6795 |
| 15.69 | 0.8997 | 0.2490 | 3.6128 |


| a $\left(^{\circ}\right.$ ) | C $\mathrm{L}^{(-)}$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.94 | 0.8963 | 0.2531 | 3.5405 |
| 16.19 | 0.8954 | 0.2597 | 3.4476 |
| 16.44 | 0.8793 | 0.2590 | 3.3947 |
| 16.69 | 0.8919 | 0.2663 | 3.3489 |
| 16.94 | 0.8997 | 0.2720 | 3.3076 |
| 17.19 | 0.9089 | 0.2743 | 3.3137 |
| 17.44 | 0.9032 | 0.2817 | 3.2059 |
| 17.69 | 0.9197 | 0.2875 | 3.1991 |
| 17.94 | 0.9078 | 0.2913 | 3.1166 |
| 18.19 | 0.9113 | 0.2968 | 3.0709 |
| 18.44 | 0.9033 | 0.2961 | 3.0510 |
| 18.69 | 0.8966 | 0.2989 | 2.9994 |
| 18.94 | 0.9103 | 0.3066 | 2.9686 |
| 19.19 | 0.9059 | 0.3074 | 2.9468 |
| 19.44 | 0.9204 | 0.3150 | 2.9219 |
| 19.69 | 0.9072 | 0.3177 | 2.8552 |
| 19.94 | 0.8863 | 0.3175 | 2.7915 |
| 20.19 | 0.8914 | 0.3223 | 2.7660 |
| 20.44 | 0.9199 | 0.3354 | 2.7425 |
| 20.69 | 0.8995 | 0.3301 | 2.7249 |
| 20.94 | 0.9184 | 0.3415 | 2.6889 |
| 21,19 | 0.9156 | 03468 | 2.6400 |
| 21.44 | 0.9091 | 0.3451 | 2.6346 |
| 21.69 | 0.9186 | 0.3544 | 2.5920 |
| 21.94 | 0.9269 | 0.3591 | 2.5807 |
| 22.19 | 0.9387 | 0.3666 | 2.5607 |
| 22.44 | 0.9239 | 0.3680 | 2.5105 |
| 22.69 | 0.9067 | 0.3706 | 2.4462 |
| 22.94 | 0.9142 | 0.3736 | 2.4467 |
| 23.19 | 0.9164 | 0.3823 | 2.3970 |
| 23.44 | 0.9244 | 03880 | 2.3824 |
| 23.69 | 0.9222 | 0.3915 | 2.3555 |
| 23.94 | 0.9174 | 0.3922 | 2.3392 |
| 24.19 | 0.9503 | 0.4089 | 2.3242 |
| 24.44 | 0.9156 | 0.4046 | 2.2631 |
| 24.69 | 0.9472 | 0.4161 | 2.2765 |
| 24.94 | 0.9361 | 0.4187 | 2.2359 |

Table E.9: Model \#9, $g=0.5 c, s=-0.5 c$, Aerodynamic Data at Re 120,000

| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.48 | -0.1435 | 0.0246 | -5.8345 |
| -2.23 | -0.1212 | 0.0242 | -5.0143 |
| -1.98 | -0.1106 | 0.0230 | -4.8021 |
| -1.73 | -0.0935 | 0.0228 | -4.1019 |
| -1.48 | $-0.0831$ | 0.0225 | -3.6954 |
| -1.23 | -0.0751 | 0.0221 | -3.4013 |
| -0.98 | -0.0592 | 0.0222 | -2.6670 |
| -0.73 | -0.0374 | 0.0217 | -1.7254 |
| -0.48 | -0.0336 | 0.0229 | -1.4685 |
| -0.23 | $-0.0107$ | 0.0215 | -0.4986 |
| 0.03 | 0.0033 | 0.0229 | 0.1426 |
| 0.28 | 0.0169 | 0.0216 | 0.7858 |
| 0.53 | 0.0360 | 0.0233 | 1.5463 |
| 0.78 | 0.0501 | 0.0230 | 2.1768 |
| 1.03 | 0.0580 | 0.0243 | 2.3900 |
| 1.28 | 0.0736 | 0.0238 | 3.0956 |
| 1.53 | 0.0927 | 0.0246 | 3.7696 |
| 1.78 | 0.1115 | 0.0258 | 4.3157 |
| 2.03 | 0.1221 | 0.0270 | 4.5262 |
| 2.28 | 0.1269 | 0.0278 | 4.5613 |
| 2.53 | 0.1503 | 0.0281 | 5.3420 |
| 2.78 | 0.1609 | 0.0304 | 5.2889 |
| 3.03 | 0.1609 | 0.0304 | 5.2889 |
| 3.28 | 0.1637 | 0.0305 | 5.3731 |
| 3.53 | 0.1874 | 0.0334 | 5.6048 |
| 3.78 | 0.2007 | 0.0340 | 5.9002 |
| 4.03 | 0.2167 | 0.0359 | 6.0274 |
| 4.28 | 0.2280 | 0.0375 | 6.0885 |
| 4.53 | 0.2468 | 0.0399 | 6.1817 |
| 4.78 | 0.2729 | 0.0432 | 6.3236 |
| 5.03 | 0.3014 | 0.0469 | 64319 |
| 5.28 | 0.3122 | 0.0500 | 6.2487 |
| 5.53 | 0.3284 | 0.0520 | 6.3111 |
| 5.78 | 0.3352 | 0.0551 | 6.0884 |
| 6.03 | 0.3525 | 0.0585 | 6.0309 |
| 6.28 | 0.3667 | 0.0606 | 6.0555 |


| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{CLL}^{(-)}$ | $C_{D}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.53 | 0.3729 | 0.0644 | 5.7886 |
| 6.78 | 0.3846 | 0.0683 | 5.6289 |
| 7.03 | 0.4048 | 0.0714 | 5.6677 |
| 7.28 | 0.4080 | 0.0731 | 5.5846 |
| 7.53 | 0.4206 | 0.0760 | 5.5309 |
| 7.78 | 0.4289 | 0.0806 | 5.3191 |
| 8.03 | 0.4423 | 0.0841 | 5.2604 |
| 8.28 | 0.4547 | 0.0870 | 5.2255 |
| 8.53 | 0.4570 | 0.0904 | 5.0564 |
| 8.78 | 0.4680 | 0.0941 | 4.9740 |
| 9.03 | 0.4752 | 0.0970 | 4.8985 |
| 9.28 | 0.4858 | 0.1022 | 4.7549 |
| 9.53 | 0.4886 | 0.1046 | 4.6711 |
| 9.78 | 0.4929 | 0.1074 | 4.5892 |
| 10.03 | 0.5025 | 0.1106 | 4.5458 |
| 10.28 | 0.5104 | 0.1155 | 4.4212 |
| 10.53 | 0.5193 | 0.1206 | 4.3050 |
| 10.78 | 0.5213 | 0.1220 | 4.2722 |
| 11.03 | 0.5236 | 0.1251 | 4.1857 |
| 11.28 | 0.5314 | 0.1294 | 4.1063 |
| 11.53 | 0.5358 | 0.1325 | 4.0450 |
| 11.78 | 0.5348 | 0.1355 | 3.9473 |
| 12.03 | 0.5434 | 0.1399 | 3.8852 |
| 12.28 | 0.5479 | 0.1431 | 3.8288 |
| 12.53 | 0.5524 | 0.1471 | 3.7561 |
| 12.78 | 0.5585 | 0.1505 | 3.7104 |
| 13.03 | 0.5623 | 0.1541 | 3.6482 |
| 13.28 | 0.5637 | 0.1589 | 3.5471 |
| 13.53 | 0.5638 | 0.1606 | 3.5098 |
| 13.78 | 0.5563 | 0.1627 | 3.4200 |
| 14.03 | 0.5600 | 0.1648 | 3.3977 |
| 14.28 | 0.5632 | 0.1698 | 3.3166 |
| 14.53 | 0.5675 | 0.1709 | 3.3211 |
| 14.78 | 0.5693 | 0.1758 | 3.2385 |
| 15.03 | 0.5678 | 0.1789 | 3.1748 |
| 15.28 | 0.5642 | 0.1794 | 3.1447 |


| $\mathrm{a}\left({ }^{\circ}\right)$ | CL (-) | Co (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.53 | 0.5702 | 0.1854 | 3.0758 |
| 15.78 | 0.5790 | 0.1918 | 3.0189 |
| 16.03 | 0.5808 | 0.1931 | 3.0077 |
| 16.28 | 0.5645 | 0.1920 | 2.9405 |
| 16.53 | 0.5624 | 0.1928 | 2.9163 |
| 16.78 | 0.5693 | 0.1995 | 2.8530 |
| 17.03 | 0.5637 | 0.1994 | 2.8278 |
| 17.28 | 0.5721 | 0.2042 | 2.8013 |
| 17.53 | 0.5828 | 0.2114 | 2.7565 |
| 17.78 | 0.5797 | 0.2147 | 2.7001 |
| 18.03 | 0.5774 | 0.2158 | 2.6748 |
| 18.28 | 0.5741 | 0.2165 | 2.6516 |
| 18.53 | 0.5794 | 0.2226 | 2.6033 |
| 18.78 | 0.5698 | 0.2207 | 2.5820 |
| 19.03 | 0.5607 | 0.2201 | 2.5470 |
| 19.28 | 0.5770 | 0.2312 | 2.4952 |
| 19.53 | 0.5700 | 0.2314 | 2.4635 |
| 19.78 | 0.5665 | 0.2305 | 2.4578 |
| 20.03 | 0.5764 | 0.2375 | 2.4265 |
| 20.28 | 0.5695 | 0.2350 | 2.4228 |
| 20.53 | 0.5650 | 0.2411 | 2.3431 |
| 20.78 | 0.5722 | 0.2464 | 2.3221 |
| 21.03 | 0.5702 | 0.2488 | 2.2920 |
| 21.28 | 0.5643 | 0.2474 | 2.2808 |
| 21.53 | 0.5689 | 0.2520 | 2.2574 |
| 21.78 | 0.5660 | 0.2586 | 2.1890 |
| 22.03 | 0.5722 | 0.2596 | 2.2038 |
| 22.28 | 0.5669 | 0.2641 | 2.1463 |
| 22.53 | 0.5542 | 0.2591 | 2.1387 |
| 22.78 | 0.5564 | 0.2610 | 2.1313 |
| 23.03 | 0.5487 | 0.2632 | 2.0844 |
| 23.28 | 0.5420 | 0.2626 | 2.0639 |
| 23.53 | 0.5465 | 0.2666 | 2.0498 |
| 23.78 | 0.5415 | 0.2679 | 2.0214 |
| 24.03 | 0.5365 | 0.2684 | 1.9990 |
| 2428 | 0.5366 | 0.2707 | 1.9819 |
| 24.53 | 0.5353 | 0.2763 | 1.9375 |

Table E.10: Model \#10, $g=0.5 \mathrm{c}, \mathrm{s}=-1 \mathrm{c}$, Aerodynamic Data at Re 120,000

| $a\left({ }^{\circ}\right)$ | $\mathrm{Cl}_{\mathrm{L}}(-)$ | $\mathrm{CD}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -1.90 | $-0.1271$ | 0.0227 | -5.6082 |
| -1.65 | -0.1094 | 0.0242 | -4.5131 |
| -1.40 | -0.0940 | 0.0230 | -4.0880 |
| -1.15 | -0.0718 | 0.0226 | -3.1764 |
| -0.90 | -0.0460 | 0.0213 | -2.1596 |
| -0.65 | -0.0506 | 0.0211 | -2.4011 |
| $-0.40$ | -0.0300 | 0.0232 | -1.2948 |
| -0.15 | -0.0077 | 0.0216 | -0.3569 |
| 0.10 | 0.0098 | 0.0230 | 0.4277 |
| 0.35 | 0.0121 | 0.0214 | 0.5669 |
| 0.60 | 0.0417 | 0.0215 | 1.9382 |
| 0.85 | 0.0568 | 0.0241 | 2.3609 |
| 1.10 | 0.0659 | 0.0238 | 2.7665 |
| 1.35 | 0.0906 | 0.0237 | 3.8285 |
| 1.60 | 0.1022 | 0.0247 | 4.1387 |
| 1.85 | 0.1170 | 0.0254 | 4.6107 |
| 2.10 | 0.1319 | 0.0259 | 50943 |
| 2.35 | 0.1530 | 0.0278 | 5.5090 |
| 2.60 | 0.1690 | 0.0297 | 5.6941 |
| 2.85 | 0.1859 | 0.0288 | 6.4569 |
| 3.10 | 0.1973 | 0.0314 | 6.2747 |
| 3.35 | 0.2113 | 0.0338 | 6.2488 |
| 3.60 | 0.2113 | 0.0338 | 6.2488 |
| 3.85 | 0.2299 | 0.0359 | 6.3964 |
| 4.10 | 0.2527 | 0.0374 | 6.7514 |
| 4.35 | 0.2575 | 0.0385 | 6.6827 |
| 4.60 | 0.2801 | 0.0411 | 6.8080 |
| 4.85 | 0.2846 | 0.0439 | 6.4791 |
| 5.10 | 0.3035 | 0.0473 | 6.4093 |
| 5.35 | 0.3344 | 0.0534 | 6.2629 |
| 5.60 | 0.3568 | 0.0545 | 6.5447 |
| 585 | 0.3613 | 0.0585 | 6.1804 |
| 6.10 | 0.3776 | 0.0623 | 6.0638 |
| 6.35 | 0.3924 | 0.0656 | 5.9778 |
| 6.60 | 0.4035 | 00692 | 5.8310 |
| 6.85 | 0.4184 | 0.0709 | 5.9003 |


| $\mathrm{a}\left({ }^{\circ}\right)$ | CI (-) | $\mathrm{Cob}^{(-)}$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 7.10 | 0.4223 | 0.0748 | 5.6466 |
| 7.35 | 0.4318 | 0.0773 | 5.5871 |
| 7.60 | 0.4366 | 0.0783 | 5.5764 |
| 7.85 | 0.4492 | 0.0843 | 5.3296 |
| 8.10 | 0.4515 | 0.0859 | 5.2543 |
| 8.35 | 0.4628 | 0.0900 | 5.1402 |
| 8.60 | 0.4634 | 0.0916 | 5.0613 |
| 8.85 | 0.4707 | 0.0946 | 4.9741 |
| 9.10 | 0.4772 | 0.0994 | 4.7985 |
| 9.35 | 0.4833 | 0.1013 | 4.7735 |
| 9.60 | 0.4865 | 0.1070 | 4.5486 |
| 985 | 0.4951 | 0.1090 | 4.5417 |
| 10.10 | 0.4941 | 0.1104 | 4.4767 |
| 10.35 | 0.4972 | 0.1122 | 4.4306 |
| 10.60 | 0.5004 | 0.1169 | 4.2808 |
| 10.85 | 0.4941 | 0.1189 | 4.1556 |
| 11.10 | 0.5017 | 0.1223 | 4.1023 |
| 11.35 | 0.5063 | 0.1254 | 4.0379 |
| 11.60 | 0.5090 | 0.1273 | 3.9991 |
| 11.85 | 0.5099 | 0.1281 | 3.9794 |
| 12.10 | 0.5040 | 0.1316 | 3.8300 |
| 12.35 | 0.5061 | 0.1320 | 3.8347 |
| 12.60 | 0.5054 | 0.1343 | 3.7616 |
| 12.85 | 0.5059 | 0.1388 | 3.6441 |
| 13.10 | 0.5021 | 0.1411 | 3.5593 |
| 13.35 | 0.5032 | 0.1386 | 3.6320 |
| 13.60 | 0.4991 | 0.1401 | 3.5618 |
| 13.85 | 0.5026 | 0.1469 | 3.4202 |
| 14.10 | 0.4941 | 0.1499 | 3.2961 |
| 14.35 | 0.4986 | 0.1464 | 3.4058 |
| 14.60 | 0.4890 | 0.1483 | 3.2974. |
| 14.85 | 0.4847 | 0.1508 | 3.2131 |
| 15.10 | 0.4843 | 0.1499 | 3.2303 |
| 15.35 | 0.4778 | 0.1534 | 3.1146 |
| 15.60 | 0.4748 | 0.1534 | 3.0948 |
| 15.85 | 0.4668 | 0.1537 | 3.0365 |


| $\square\left(^{\circ}\right.$ ) | CL (-) | $\mathrm{Cb}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 16.10 | 0.4617 | 0.1522 | 3.0345 |
| 16.35 | 0.4565 | 0.1588 | 2.8743 |
| 16.60 | 0.4525 | 0.1542 | 2.9352 |
| 16.85 | 0.4524 | 0.1565 | 2.8900 |
| 17.10 | 0.4361 | 0.1586 | 2.7502 |
| 17.35 | 0.4329 | 0.1466 | 2.9534 |
| 17.60 | 0.4317 | 0.1452 | 2.9728 |
| 17.85 | 0.4228 | 0.1614 | 2.6197 |
| 18.10 | 0.4126 | 0.1502 | 2.7468 |
| 18.35 | 0.4115 | 0.1459 | 2.8210 |
| 18.60 | 0.4045 | 0.1527 | 26486 |
| 18.85 | 0.4080 | 0.1552 | 2.6290 |
| 19.10 | 0.3942 | 0.1570 | 2.5103 |
| 19.35 | 0.3893 | 0.1521 | 2.5591 |
| 19.60 | 0.3871 | 0.1497 | 2.5853 |
| 19.85 | 0.3831 | 0.1540 | 2.4879 |
| 20.10 | 0.3775 | 0.1575 | 2.3965 |
| 20.35 | 0.3775 | 0.1584 | 2.3823 |
| 20.60 | 0.3767 | 0.1549 | 2.4313 |
| 20.85 | 0.3664 | 0.1561 | 2.3473 |
| 21.10 | 0.3706 | 0.1575 | 2.3536 |
| 21.35 | 0.3605 | 0.1582 | 2.2785 |
| 21.60 | 0.3602 | 0.1541 | 2.3376 |
| 21.85 | 0.3508 | 0.1572 | 2.2315 |
| 22.10 | 0.3533 | 0.1584 | 2.2299 |
| 22.35 | 0.3528 | 0.1625 | 2.1712 |
| 22.60 | 0.3461 | 0.1572 | 2.2016 |
| 22.85 | 0.3484 | 0.1629 | 2.1381 |
| 23.10 | 0.3429 | 0.1603 | 2.1389 |
| 23.35 | 0.3473 | 0.1657 | 2.0963 |
| 23.60 | 0.3398 | 0.1634 | 2.0797 |
| 23.85 | 0.3529 | 0.1680 | 2.1005 |
| 24.10 | 0.3433 | 0.1682 | 2.0413 |
| 24.35 | 0.3454 | 0.1688 | 2.0457 |
| 24.60 | 0.3465 | 0.1747 | 1.9829 |
| 24.85 | 0.3482 | 0.1743 | 1.9977 |
| 25.10 | 0.3429 | 0.1753 | 1.9560 |

Table E.11: Model \#11, $g=1 \mathrm{c}, \mathrm{s}=-\mathbf{0 . 5 c}$, Aerodynamic Data at Re 120,000

| $\mathrm{a}\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{CD}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.03 | -0.1404 | 0.0256 | -5.4801 |
| -1.78 | -0.1268 | 0.0237 | -5.3408 |
| -1.53 | -0.1108 | 0.0237 | -4.6730 |
| -1.28 | -0.0973 | 0.0249 | -3.9159 |
| -1.03 | -0.0687 | 0.0245 | $-2.8021$ |
| -0.78 | -0.0566 | 0.0227 | -2.4940 |
| -0.53 | -0.0377 | 0.0225 | -1.6751 |
| -0.28 | -0.0225 | 0.0242 | -0.9289 |
| -0.03 | -0.0064 | 0.0223 | -0.2867 |
| 0.22 | 0.0046 | 0.0232 | 0.1988 |
| 0.47 | 0.0203 | 0.0225 | 0.9017 |
| 0.72 | 0.0319 | 0.0225 | 1.4153 |
| 0.97 | 0.0607 | 0.0218 | 2.7854 |
| 1.22 | 0.0660 | 0.0240 | 2.7492 |
| 1.47 | 0.0918 | 00239 | 3.8455 |
| 1.72 | 0.1028 | 0.0263 | 3.9069 |
| 1.97 | 0.1237 | 0.0258 | 4.7887 |
| 2.22 | 0.1392 | 0.0258 | 5.3977 |
| 2.47 | 0.1584 | 0.0276 | 5.7335 |
| 2.72 | 0.1817 | 0.0305 | 5.9635 |
| 2.97 | 0.1940 | 0.0318 | 6.0984 |
| 3.22 | 0.2089 | 0.0331 | 6.3054 |
| 3.47 | 0.2089 | 0.0331 | 6.3054 |
| 3.72 | 0.2314 | 0.0354 | 6.5374 |
| 3.97 | 0.2419 | 0.0373 | 6.4774 |
| 4.22 | 0.2576 | 0.0380 | 6.7838 |
| 4.47 | 0.2938 | 0.0407 | 7.2191 |
| 4.72 | 0.3153 | 0.0451 | 6.9845 |
| 4.97 | 0.3200 | 0.0471 | 6.7957 |
| 5.22 | 0.3574 | 0.0535 | 6.6805 |
| 5.47 | 0.3853 | 0.0568 | 6.7837 |
| 5.72 | 0.4019 | 0.0590 | 6.8075 |
| 5.97 | 0.4218 | 0.0630 | 6.6943 |
| 6.22 | 0.4433 | 0.0671 | 6.6110 |
| 6.47 | 0.4622 | 0.0725 | 6.3789 |
| 6.72 | 0.4799 | 0.0776 | 6.1823 |


| $\square\left(^{\circ}\right.$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.97 | 0.4984 | 0.0796 | 6.2653 |
| 7.22 | 0.5058 | 0.0840 | 6.0222 |
| 7.47 | 0.5306 | 0.0889 | 5.9662 |
| 7.72 | 0.5591 | 0.0942 | 5.9343 |
| 7.97 | 0.5798 | 0.0995 | 5.8291 |
| 8.22 | 0.5845 | 0.1033 | 5.6578 |
| 8.47 | 0.6111 | 0.1096 | 5.5735 |
| 8.72 | 0.6225 | 0.1160 | 5.3680 |
| 8.97 | 0.6439 | 0.1204 | 5.3473 |
| 9.22 | 0.6678 | 0.1283 | 5.2042 |
| 9.47 | 0.6750 | 0.1319 | 5.1180 |
| 9.72 | 0.6931 | 0.1387 | 4.9962 |
| 9.97 | 0.7032 | 0.1429 | 4.9195 |
| 10.22 | 0.7019 | 0.1457 | 4.8186 |
| 10.47 | 0.7211 | 0.1514 | 4.7642 |
| 10.72 | 0.7298 | 0.1583 | 4.6100 |
| 10.97 | 0.7389 | 0.1627 | 4.5426 |
| 11.22 | 0.7445 | 0.1671 | 44553 |
| 11.47 | 0.7505 | 0.1736 | 4.3239 |
| 11.72 | 0.7540 | 0.1770 | 4.2595 |
| 11.97 | 0.7547 | 0.1805 | 4.1815 |
| 12.22 | 0.7572 | 0.1860 | 4.0706 |
| 12.47 | 0.7525 | 0.1892 | 3.9766 |
| 12.72 | 0.7536 | 0.1906 | 3.9543 |
| 12.97 | 0.7539 | 0.1947 | 3.8725 |
| 13.22 | 0.7569 | 0.1995 | 3.7938 |
| 13.47 | 0.7499 | 0.2024 | 3.7045 |
| 13.72 | 0.7452 | 0.2040 | 3.6529 |
| 13.97 | 0.7388 | 0.2070 | 3.5683 |
| 14.22 | 0.7296 | 0.2068 | 3.5288 |
| 14.47 | 0.7294 | 0.2100 | 3.4739 |
| 14.72 | 0.7360 | 0.2160 | 34079 |
| 14.97 | 0.7259 | 0.2161 | 3.3594 |
| 15.22 | 0.7200 | 0.2191 | 3.2858 |
| 15.47 | 0.7106 | 0.2216 | 3.2068 |
| 15.72 | 0.7050 | 0.2212 | 3.1874 |


| a ( ${ }^{\circ}$ ) | C ${ }_{\text {L }}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.97 | 0.6941 | 0.2229 | 3.1142 |
| 16.22 | 0.7024 | 0.2278 | 3.0842 |
| 16.47 | 0.6848 | 0.2268 | 3.0190 |
| 16.72 | 0.6859 | 0.2298 | 2.9847 |
| 16.97 | 0.6754 | 0.2307 | 2.9278 |
| 17.22 | 0.6762 | 0.2339 | 2.8906 |
| 17.47 | 0.6663 | 0.2357 | 2.8270 |
| 17.72 | 0.6824 | 0.2445 | 2.7909 |
| 17.97 | 0.6797 | 0.2454 | 2.7696 |
| 18.22 | 0.6879 | 0.2531 | 2.7175 |
| 18.47 | 0.6792 | 0.2540 | 2.6735 |
| 18.72 | 0.6830 | 0.2569 | 2.6583 |
| 18.97 | 0.6789 | 0.2582 | 2.6293 |
| 19.22 | 0.6797 | 0.2623 | 2.5909 |
| 19.47 | 0.6888 | 0.2687 | 2.5640 |
| 19.72 | 0.6892 | 0.2711 | 2.5425 |
| 19.97 | 0.6871 | 0.2780 | 2.4711 |
| 20.22 | 0.6966 | 0.2831 | 2.4607 |
| 20.47 | 0.6879 | 0.2827 | 2.4333 |
| 20.72 | 0.6847 | 0.2859 | 2.3951 |
| 20.97 | 0.6820 | 0.2896 | 2.3551 |
| 21.22 | 0.6800 | 0.2895 | 2.3485 |
| 21.47 | 0.6857 | 0.2957 | 2.3189 |
| 21.72 | 0.6812 | 0.2994 | 2.2750 |
| 21.97 | 0.6742 | 0.3003 | 2.2450 |
| 22.22 | 0.6780 | 0.3051 | 2.2219 |
| 22.47 | 0.6797 | 0.3105 | 2.1893 |
| 22.72 | 0.6743 | 0.3092 | 2.1811 |
| 22.97 | 0.6586 | 0.3056 | 2.1549 |
| 2322 | 0.6687 | 0.3143 | 2.1277 |
| 23.47 | 0.6683 | 0.3185 | 2.0983 |
| 23.72 | 0.6671 | 0.3190 | 2.0911 |
| 23.97 | 0.6815 | 0.3308 | 2.0604 |
| 24.22 | 0.6767 | 0.3295 | 2.0536 |
| 24.47 | 0.6807 | 0.3370 | 2.0199 |
| 24.72 | 0.6826 | 0.3414 | 1.9995 |
| 24.97 | 0.6743 | 0.3401 | 1.9828 |

Table E.12: Model \#12, $\mathrm{g}=1 \mathrm{c}, \mathrm{s}=-1 \mathrm{c}$, Aerodynamic Data at Re 120,000

| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | $\mathrm{a}\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(\cdot)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) | a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -2.03 | -0.1370 | 0.0226 | -6.0702 | 6.97 | 0.4804 | 0.0750 | 6.4079 | 15.97 | 0.6338 | 0.2032 | 3.1186 |
| -1.78 | -0.1167 | 0.0217 | -5.3775 | 7.22 | 0.4979 | 0.0811 | 6.1429 | 16.22 | 0.6498 | 0.2102 | 3.0915 |
| -1.53 | -0.0990 | 0.0225 | -4.4084 | 7.47 | 0.5132 | 0.0848 | 6.0527 | 16.47 | 0.6470 | 0.2149 | 3.0103 |
| -1.28 | -0.0899 | 0.0198 | -4.5519 | 7.72 | 0.5272 | 0.0919 | 5.7345 | 16.72 | 0.6465 | 0.2205 | 2.9320 |
| $-1.03$ | -0.0681 | 0.0204 | -3.3444 | 7.97 | 0.5510 | 0.0938 | 5.8756 | 16.97 | 0.6545 | 0.2248 | 2.9115 |
| -0.78 | -0.0450 | 0.0201 | -2.2349 | 8.22 | 0.5628 | 0.0977 | 5.7585 | \$7.22 | 0.6532 | 0.2247 | 2.9075 |
| -0.53 | -0.0429 | 0.0214 | -2.0000 | 8.47 | 0.5768 | 0.1035 | 5.5705 | 77.47 | 0.6529 | 0.2273 | 2.8722 |
| -0.28 | -0.0256 | 0.0211 | -1.2130 | 8.72 | 0.5873 | 0.1077 | 5.4545 | 17.72 | 0.6596 | 0.2337 | 2.8231 |
| -0.03 | -0.0051 | 0.0181 | -0.2785 | 8.97 | 0.6044 | 0.1121 | 5.3916 | 17.97 | 0.6553 | 0.2370 | 2.7652 |
| 0.22 | 0.0102 | 0.0194 | 0.5259 | 9.22 | 0.6100 | 0.1161 | 5.2531 | 18.22 | 0.6556 | 0.2400 | 2.7316 |
| 0.47 | 0.0253 | 0.0205 | 1.2360 | 9.47 | 0.6229 | 0.1189 | 5.2393 | 18.47 | 0.6384 | 0.2373 | 2.6907 |
| 0.72 | 0.0504 | 0.0215 | 2.3403 | 9.72 | 0.6171 | 0.1248 | 4.9433 | 18.72 | 0.6477 | 0.2408 | 2.6903 |
| 0.97 | 0.0624 | 0.0202 | 3.0875 | 9.97 | 0.6299 | 0.1287 | 4.8922 | 18.97 | 0.6475 | 0.2463 | 2.6287 |
| 1.22 | 0.0762 | 0.0223 | 3.4182 | 10.22 | 0.6329 | 0.1317 | 4.8042 | 19.22 | 0.6524 | 0.2508 | 2.6015 |
| 1.47 | 0.0917 | 0.0210 | 4.3732 | 10.47 | 0.6381 | 0.1361 | 4.6872 | 19.47 | 0.6438 | 0.2486 | 2.5895 |
| 1.72 | 0.1124 | 0.0233 | 4.8183 | 10.72 | 0.6418 | 0.1390 | 4.6185 | 19.72 | 0.6327 | 0.2513 | 2.5181 |
| 1.97 | 0.1392 | 0.0237 | 5.8795 | 10.97 | 0.6457 | 0.1442 | 4.4780 | 19.97 | 0.6476 | 0.2603 | 2.4877 |
| 2.22 | 0.1575 | 0.0253 | 6.2376 | 11.22 | 0.6508 | 0.1460 | 4.4582 | 20.22 | 0.6402 | 0.2592 | 2.4700 |
| 2.47 | 0.1728 | 0.0269 | 6.4312 | 11.47 | 0.6505 | 0.1505 | 4.3217 | 20.47 | 0.6421 | 0.2642 | 2.4304 |
| 2.72 | 0.1987 | 0.0262 | 7.5815 | 11.72 | 0.6503 | 0.1520 | 4.2782 | 20.72 | 0.6417 | 0.2666 | 2.4069 |
| 2.97 | 0.2102 | 0.0311 | 6.7629 | 11.97 | 0.6486 | 0.1562 | 4.1529 | 20.97 | 0.6475 | 0.2751 | 2.3540 |
| 3.22 | 0.2225 | 0.0314 | 7.0761 | 12.22 | 0.6526 | 0.1592 | 4.0980 | 21.22 | 0.6700 | 0.2858 | 2.3441 |
| 3.47 | 0.2225 | 0.0314 | 7.0761 | I2.47 | 0.6481 | 0.1621 | 3.9970 | 2147 | 0.6559 | 0.2834 | 2.3142 |
| 3.72 | 0.2525 | 0.0351 | 7.1938 | 12.72 | 0.6459 | 0.1651 | 3.9126 | 21.72 | 0.6760 | 0.2897 | 2.3339 |
| 3.97 | 0.2481 | 0.0341 | 7.2799 | $\$ 2.97$ | 06462 | 0.1694 | 3.8133 | 21.97 | 0.6853 | 0.3002 | 2.2830 |
| 4.22 | 0.2631 | 0.0358 | 7.3434 | 13.22 | 0.6551 | 0.1752 | 3.7390 | 22.22 | 0.6878 | 0.3054 | 2.2523 |
| 4.47 | 0.3012 | 0.0392 | 7.6888 | 13.47 | 0.6510 | 0.1790 | 3.6371 | 22.47 | 0.6645 | 0.2996 | 2.2179 |
| 4.72 | 0.3020 | 0.0420 | 7.1817 | 13.72 | 0.6531 | 0.1805 | 3.6194 | 22.72 | 0.6984 | 0.3149 | 2.2183 |
| 4.97 | 0.3377 | 0.0444 | 7.6087 | 13.97 | 0.6553 | 0.1837 | 3.5679 | 22.97 | 0.6811 | 0.3123 | 2.1809 |
| 5.22 | 0.3685 | 0.0496 | 7.4319 | 14.22 | 0.6437 | 0.1831 | 3.5152 | 23.22 | 0.7075 | 0.3315 | 2.1343 |
| 5.47 | 0.3800 | 0.0536 | 7.0924 | 14.47 | 0.6481 | 0.1877 | 3.4528 | 23.47 | 0.7205 | 0.3384 | 2.1288 |
| 5.72 | 0.4117 | 0.0589 | 6.9920 | 14.72 | 0.6655 | 0.1969 | 3.3797 | 23.72 | 0.6992 | 0.3297 | 2.1208 |
| 5.97 | 0.4155 | 0.0624 | 6.6565 | 14.97 | 0.6522 | 0.1989 | 3.2792 | 23.97 | 0.7231 | 0.3466 | 2.0865 |
| 6.22 | 0.4400 | 0.0662 | 6.6502 | 15.22 | 0.6586 | 0.2010 | 3.2770 | 24.22 | 0.7176 | 0.3442 | 2.0850 |
| 6.47 | 0.4532 | 0.0702 | 6.4563 | 15.47 | 0.6647 | 0.2063 | 3.2224 | 24.47 | 0.7347 | 0.3612 | 2.0338 |
| 6.72 | 0.4659 | 0.0716 | 6.5096 | 15.72 | 0.6360 | 0.2021 | 3.1465 | 24.72 | 0.7048 | 0.3431 | 2.0544 |
|  |  |  |  |  |  |  |  | 24.97 | 0.7138 | 0.3554 | 2.0084 |

Table E.13: Model \#13, $g=2 c, s=-1 c$, Aerodynamic Data at Re 120,000

| a ( ${ }^{\circ}$ ) | $\mathrm{C}_{\mathrm{L}}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.13 | -0.1713 | 0.0271 | -6.3167 |
| -1.88 | -0.1590 | 0.0258 | -6.1559 |
| -1.63 | -0.1226 | 0.0263 | -4.6577 |
| -1.38 | -0.1057 | 0.0246 | -4.2929 |
| -1.13 | -0.0903 | 0.0259 | -3.4942 |
| -0.88 | -0.0757 | 0.0240 | -3.1497 |
| -0.63 | -0.0530 | 0.0247 | -2.1449 |
| -0.38 | -0.0302 | 0.0238 | -1.2697 |
| -0.13 | -0.0150 | 0.0227 | -0.6618 |
| 0.12 | 0.0177 | 0.0222 | 0.7973 |
| 0.37 | 0.0310 | 0.0216 | 1.4395 |
| 0.62 | 0.0519 | 0.0216 | 2.4084 |
| 0.87 | 0.0819 | 0.0200 | 4.0984 |
| 1.12 | 0.0914 | 0.0216 | 4.2303 |
| 1.37 | 0.1187 | 0.0238 | 4.9956 |
| 1.62 | 0.1346 | 0.0241 | 5.5778 |
| 1.87 | 0.1669 | 0.0249 | 6.7079 |
| 2.12 | 0.1830 | 0.0244 | 7.4912 |
| 2.37 | 0.1990 | 0.0249 | 7.9923 |
| 2.62 | 0.2262 | 0.0269 | 84047 |
| 2.87 | 0.2508 | 0.0287 | 8.7471 |
| 3.12 | 0.2591 | 0.0294 | 8.8238 |
| 3.37 | 0.2591 | 0.0294 | 8.8238 |
| 3.62 | 0.2921 | 0.0306 | 9.5388 |
| 3.87 | 0.3071 | 0.0314 | 9.7692 |
| 4.12 | 0.3208 | 0.0338 | 9.4856 |
| 4.37 | 0.3477 | 0.0369 | 9.4219 |
| 4.62 | 0.3766 | 0.0390 | 9.6649 |
| 4.87 | 0.3977 | 0.0423 | 9.3949 |
| 5.12 | 0.4369 | 0.0470 | 9.3016 |
| 5.37 | 0.4765 | 0.0515 | 9.2468 |
| 5.62 | 0.4956 | 0.0566 | 8.7547 |
| 5.87 | 0.5294 | 0.0605 | 8.7458 |
| 6.12 | 0.5469 | 0.0638 | 8.5685 |
| 6.37 | 0.5774 | 0.0672 | 8.5911 |
| 6.62 | 0.6096 | 0.0726 | 8.3924 |


| a $\left(^{\circ}\right.$ ) | $\mathrm{Cl}_{\mathrm{L}}(-)$ | Co (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.87 | 0.6167 | 0.0753 | 8.1942 |
| 7.12 | 0.6380 | 0.0797 | 8.0038 |
| 7.37 | 0.6630 | 0.0851 | 7.7919 |
| 7.62 | 06801 | 0.0890 | 7.6427 |
| 7.87 | 0.6923 | 0.0902 | 7.6778 |
| 8.12 | 0.7283 | 0.0974 | 7.4740 |
| 8.37 | 0.7351 | 0.1027 | 7.1548 |
| 8.62 | 0.7559 | 0.1086 | 6.9620 |
| 8.87 | 0.7803 | 0.1151 | 6.7804 |
| 912 | 0.7881 | 0.1202 | 6.5589 |
| 9.37 | 0.8062 | 0.1248 | 6.4582 |
| 9.62 | 0.8031 | 0.1288 | 6.2373 |
| 9.87 | 0.8067 | 0.1316 | 6.1323 |
| 10.12 | 0.8139 | 0.1361 | 5.9823 |
| 10.37 | 0.8112 | 0.1401 | 5.7912 |
| 10.62 | 0.8198 | 0.1371 | 5.9808 |
| 10.87 | 0.8237 | 0.1472 | 5.5969 |
| 11.12 | 0.8352 | 0.1534 | 5.4433 |
| 11.37 | 0.8278 | 0.1542 | 5.3684 |
| 11.62 | 0.8243 | 0.1539 | 5.3573 |
| 11.87 | 0.8170 | 0.1530 | 5.3412 |
| 12.12 | 0.8149 | 0.1573 | 5.1796 |
| 12.37 | 0.8109 | 0.1593 | 50889 |
| 12.62 | 0.7975 | 0.1539 | 5.1812 |
| 12.87 | 0.7976 | 0.1291 | 6.1790 |
| 13.12 | 0.8052 | 0.1486 | 5.4196 |
| 13.37 | 0.8031 | 0.1265 | 6.3465 |
| 13.62 | 0.8072 | 0.1253 | 6.4438 |
| 13.87 | 0.8025 | 0.1293 | 6.2041 |
| 14.12 | 0.8065 | 0.1026 | 7.8589 |
| 14.37 | 0.7944 | 0.1204 | 6.5956 |
| 14.62 | 0.7698 | 0.1088 | 7.0754 |
| 14.87 | 0.7944 | 0.1210 | 6.5657 |
| 15.12 | 0.8142 | 0.1214 | 6.7070 |
| 15.37 | 0.7940 | 0.1139 | 6.9693 |
| 15.62 | 0.7820 | 0.1526 | 5.1229 |


| $\mathrm{a}\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.87 | 0.8103 | 0.1271 | 6.3728 |
| 16.12 | 0.7907 | 0.1489 | 5.3103 |
| 16.37 | 0.7987 | 0.1256 | 6.3602 |
| 16.62 | 0.7947 | 0.1502 | 5.2895 |
| 16.87 | 0.7868 | 0.1667 | 4.7214 |
| 17.12 | 0.7944 | 0.1360 | 5.8413 |
| 17.37 | 0.7706 | 0.1548 | 4.9769 |
| 17.62 | 0.7947 | 0.1615 | 4.9214 |
| 17.87 | 0.7832 | 0.1533 | 5.1105 |
| 18.12 | 0.7933 | 0.1504 | 5.2734 |
| 18.37 | 0.8249 | 0.1507 | 5.4748 |
| 18.62 | 0.8018 | 0.1666 | 4.8135 |
| 18.87 | 0.8096 | 0.1605 | 5.0439 |
| 19.12 | 0.8015 | 0.1658 | 4.8355 |
| 19.37 | 0.7965 | 0.1783 | 4.4679 |
| 19.62 | 0.7802 | 0.1913 | 4.0775 |
| 19.87 | 0.8226 | 0.1719 | 4.7853 |
| 20.12 | 0.8277 | 0.1813 | 4.5652 |
| 20.37 | 0.8139 | 0.1872 | 4.3468 |
| 20.62 | 0.8096 | 0.1942 | 4.1687 |
| 20.87 | 0.8372 | 0.1979 | 4.2316 |
| 21.12 | 0.8597 | 0.2047 | 4.1989 |
| 21.37 | 0.8004 | 0.2242 | 3.5695 |
| 21.62 | 0.8380 | 0.2204 | 3.8026 |
| 21.87 | 0.7989 | 0.2242 | 3.5632 |
| 22.12 | 0.8354 | 0.2146 | 3.8918 |
| 22.37 | 0.8619 | 0.2179 | 3.9565 |
| 22.62 | 0.8039 | 0.2196 | 3.6602 |
| 22.87 | 0.8070 | 0.2046 | 3.9442 |
| 23.12 | 0.8287 | 0.2300 | 3.6030 |
| 23.37 | 0.7747 | 0.2456 | 3.1540 |
| 23.62 | 0.7882 | 0.2380 | 3.3117 |
| 23.87 | 0.7435 | 0.2552 | 2.9129 |
| 24.12 | 0.8365 | 0.2438 | 3.4304 |
| 24.37 | 0.7909 | 0.2454 | 3.2226 |
| 24.62 | 0.8222 | 0.2523 | 3.2584 |
| 24.87 | 0.7860 | 0.3081 | 2.5513 |

Table E.14: Model \#14, $g=1 c, s=-1.5 c$, Aerodynamic Data at Re 120,000

| $\alpha\left({ }^{\circ}\right)$ | $\mathrm{C}_{\mathrm{L}}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| -2.06 | -0.1498 | 0.0255 | -5.8804 |
| -1.81 | -0.1305 | 0.0236 | -5.5397 |
| -1.56 | -0.1096 | 0.0235 | -4.6586 |
| -1.31 | -0.1014 | 0.0230 | -4.4065 |
| -1.06 | -0.0712 | 0.0232 | -3.0744 |
| -0.81 | -0.0540 | 0.0241 | -2.2366 |
| -0.56 | -0.0375 | 0.0229 | -1.6392 |
| -0.31 | -0.0232 | 0.0231 | -1.0074 |
| -0.06 | -0.0004 | 0.0220 | -0.0160 |
| 0.19 | 0.0092 | 0.0240 | 0.3830 |
| 0.44 | 0.0300 | 0.0220 | 1.3636 |
| 0.69 | 0.0473 | 0.0232 | 2.0365 |
| 0.94 | 0.0739 | 0.0237 | 3.1141 |
| 1.19 | 0.0865 | 0.0245 | 3.5313 |
| 1.44 | 0.1040 | 0.0234 | 4.4371 |
| 1.69 | 0.1156 | 0.0257 | 4.4922 |
| 1.94 | 0.1393 | 0.0257 | 5.4236 |
| 2.19 | 0.1546 | 0.0267 | 5.7879 |
| 2.44 | 0.1805 | 0.0289 | 6.2464 |
| 2.69 | 0.1837 | 0.0296 | 6.1983 |
| 2.94 | 0.2108 | 0.0303 | 6.9580 |
| 3.19 | 0.2275 | 0.0330 | 6.8843 |
| 3.44 | 0.2275 | 0.0330 | 6.8843 |
| 3.69 | 0.2354 | 0.0335 | 7.0242 |
| 3.94 | 0.2708 | 0.0353 | 7.6638 |
| 4.19 | 0.2833 | 0.0396 | 7.1551 |
| 4.44 | 0.2978 | 0.0420 | 7.0867 |
| 4.69 | 0.3233 | 0.0438 | 7.3759 |
| 4.94 | 0.3351 | 0.0465 | 7.2070 |
| 5.19 | 0.3792 | 0.0509 | 7.4514 |
| 5.44 | 0.4107 | 0.0563 | 7.2952 |
| 5.69 | 0.4233 | 0.0602 | 7.0298 |
| 5.94 | 0.4400 | 0.0621 | 7.0799 |
| 6.19 | 0.4580 | 0.0680 | 6.7339 |
| 6.44 | 0.4819 | 0.0729 | 6.6112 |
| 6.69 | 0.4963 | 0.0758 | 6.5483 |


| $\mathrm{d}\left({ }^{\circ}\right.$ ) | $\mathrm{C}_{1}(-)$ | $\mathrm{C}_{\mathrm{D}}(-)$ | L/D (-) |
| :---: | :---: | :---: | :---: |
| 6.94 | 0.5104 | 0.0802 | 6.3652 |
| 7.19 | 0.5239 | 0.0843 | 6.2150 |
| 7.44 | 0.5382 | 0.0881 | 6.1121 |
| 7.69 | 0.5571 | 0.0922 | 6.0408 |
| 7.94 | 0.5632 | 0.0946 | 5.9559 |
| 8.19 | 0.5843 | 0.1014 | 5.7613 |
| 8.44 | 0.5908 | 0.1042 | 5.6710 |
| 869 | 0.6083 | 0.1119 | 5.4374 |
| 8.94 | 0.6287 | 0.1179 | 5.3319 |
| 9.19 | 0.6354 | 0.1217 | 5.2193 |
| 944 | 0.6574 | 0.1279 | 5.1383 |
| 969 | 0.6527 | 0.1284 | 5.0836 |
| 9.94 | 0.6519 | 0.1320 | 4.9369 |
| 10.19 | 0.6850 | 0.1416 | 4.8374 |
| 10.44 | 0.6836 | 0.1449 | 4.7179 |
| 1069 | 0.6946 | 0.1487 | 4.6699 |
| 10.94 | 0.6905 | 0.1525 | 4.5274 |
| 11.19 | 0.7125 | 0.1604 | 4.4408 |
| 11.44 | 0.7223 | 0.1666 | 4.3348 |
| 11.69 | 0.7245 | 0.1679 | 4.3143 |
| 11.94 | 0.7201 | 0.1711 | 4.2093 |
| 12.19 | 0.7275 | 0.1780 | 4.0862 |
| 12.44 | 0.7170 | 0.1770 | 4.0513 |
| 12.69 | 0.7390 | 0.1871 | 3.9497 |
| 12.94 | 0.7336 | 0.1889 | 3.8836 |
| 13.19 | 0.7244 | 0.1899 | 3.8141 |
| 13.44 | 0.7180 | 0.1916 | 3.7482 |
| 13.69 | 0.7183 | 0.7953 | 3.6782 |
| 13.94 | 0.7219 | 0.1981 | 3.6434 |
| 14.19 | 0.7421 | 0.2103 | 3.5287 |
| 14.44 | 0.6976 | 0.2005 | 3.4793 |
| 14.69 | 0.7070 | 0.2064 | 3.4258 |
| 14.94 | 0.7050 | 0.2082 | 3.3856 |
| 15.19 | 0.7085 | 0.2168 | 3.2674 |
| 15.44 | 0.6951 | 0.2118 | 3.2819 |
| 15.69 | 0.7042 | 0.2220 | 3.1729 |


| $\mathrm{a}\left({ }^{\circ}\right)$ | $\mathrm{C}_{2}(-)$ | CD (-) | L/D (-) |
| :---: | :---: | :---: | :---: |
| 15.94 | 0.6915 | 0.2218 | 3.1172 |
| 16.19 | 0.7152 | 0.2311 | 3.0947 |
| 16.44 | 0.7008 | 0.2319 | 3.0221 |
| 16.69 | 0.6821 | 0.2278 | 2.9947 |
| 16.94 | 0.6864 | 0.2316 | 2.9640 |
| 17.19 | 0.6858 | 0.2351 | 2.9168 |
| \$7.44 | 0.7131 | 0.2485 | 2.8691 |
| \$7.69 | 0.7002 | 0.2467 | 2.8383 |
| $\pm 7.94$ | 0.7127 | 0.2524 | 2.8238 |
| 18.19 | 0.6879 | 0.2475 | 2.7797 |
| 1844 | 0.6789 | 0.2470 | 2.7484 |
| 18.69 | 0.6868 | 0.2561 | 2.6819 |
| 18.94 | 0.6679 | 0.2528 | 2.6418 |
| 19.19 | 0.6821 | 0.2610 | 2.6132 |
| 19.44 | 0.6698 | 0.2560 | 2.6168 |
| 19.69 | 0.6667 | 0.2606 | 2.5588 |
| 19.94 | 0.6744 | 0.2715 | 2.4843 |
| 20.19 | 0.6966 | 0.2874 | 2.4240 |
| 20.44 | 0.6948 | 0.2796 | 2.4847 |
| 20.69 | 0.7303 | 0.2965 | 2.4630 |
| 20.94 | 0.6863 | 0.2819 | 2.4348 |
| 21.19 | 0.6692 | 0.2805 | 2.3859 |
| 21.44 | 0.6834 | 0.2933 | 2.3304 |
| 21.69 | 0.7463 | 0.3059 | 2.4402 |
| 21.94 | 0.7189 | 0.3106 | 2.3148 |
| 22.19 | 0.7581 | 0.3204 | 2.3663 |
| 22.44 | 0.6982 | 0.3077 | 2.2693 |
| 22.69 | 0.6806 | 0.3087 | 2.2048 |
| 22.94 | 0.6939 | 0.3256 | 2.1314 |
| 23.19 | 0.6577 | 0.2959 | 2.2226 |
| 23.44 | 0.6323 | 0.2879 | 2.1963 |
| 23.69 | 0.5896 | 0.2843 | 2.0739 |
| 23.94 | 0.6200 | 0.3054 | 2.0298 |
| 24.19 | 0.5813 | 0.2828 | 2.0552 |
| 24.44 | 0.5855 | 0.2841 | 2.0611 |
| 24.69 | 0.5911 | 0.2748 | 2.1506 |
| 24.94 | 0.5915 | 0.3132 | 1.8883 |

## Appendix F: Stall Flutter

## Overview

This section is intended to clarify the behavior seen in Figure 4.22 and Figure 4.34. Most of this section will discuss how flutter occurs and the effect that it can have on the structure, such as wings and airplanes.

## Introduction

Flutter is a dangerous phenomenon encountered in flexible structures subjected to aerodynamic forces. This includes aircraft, buildings, telegraph wires, stop signs, and bridges. Flutter occurs as a result of interactions between aerodynamics, stiffness, and inertial forces on a structure. In an aircraft, as the speed of the wind increases, there may be a point at which the structural damping is insufficient to damp out the motions which are increasing due to aerodynamic energy being added to the structure. This vibration can cause structural failure and therefore considering flutter characteristics is an essential part of designing an aircraft.

## Flutter Motion

The basic type of flutter seen in aircraft wings is described here. Flutter may be initiated by a rotation of the airfoil (see $t=0$ in Figure F.1). As the increased force causes the airfoil to rise, the torsional stiffness of the structure returns the airfoil to zero rotation ( $\mathrm{t}=\mathrm{T} / 4$ in Figure F.1). The bending stiffness of the structure tries to
return the airfoil to the neutral position, but now the airfoil rotates in a nose-down position ( $t=\mathrm{T} / 2$ in Figure F.1). Again the increased force causes the airfoil to plunge and the torsional stiffness returns the airfoil to zero rotation ( $t=3 T / 4$ ). The cycle is completed when the airfoil returns to the neutral position with a nose-up rotation. Notice that the maximum rotation leads the maximum rise or plunge by 90 degrees (T/4). As time increases, the plunge motion tends to damp out, but the rotational motion diverges. If the motion is allowed to continue, the forces due to rotation will cause the structure to fail.


Figure F.1: Rotation and Plunge Motion for an Airfoil Exhibiting Flutter

This flutter is caused by the coalescence of two structural modes - pitch and plunge (or wing-bending) motion. This example wing has two basic degrees of freedom or natural modes of vibration: pitch and plunge (bending). The pitch mode is rotational and the bending mode is a vertical up and down motion at the wing tip. As the airfoil flies at increasing speed, the frequencies of these modes coalesce to create one mode at the flutter frequency and flutter condition. This is
flutter resonance. Stall flutter is a torsional mode of flutter that occurs on wings at high loading conditions near the stall speed. Because the airflow separates during stall, this single degree-of-freedom flutter cannot be explained by classical flutter theory. Figure 4.22 and Figure 4.34 show the aerodynamic efficiency of models tested in the LSWT. It can be seen that the region where the data of lift over drag ratio begin to oscillate is the region where the same model reaches the highest $C_{L}$, that is the region of stall.

## Classical Flutter Definition

The following example of a simple two degree-of-freedom model is fundamental to understanding flutter behavior. Aerodynamic forces excite the structural spring/mass system. The plunge spring represents the bending stiffness of the structure and the rotation spring represents the torsional stiffness. The shape of the airfoil determines the aerodynamic center. The center of gravity is determined by the mass distribution of the cross-section (that is, how the airfoil is constructed).


Figure F.2: Aerodynamic Lag

In Figure F.2, at time $t=t_{0}$ no angle of attack means symmetrical pressures on the upper and lower surfaces. At $t=t_{1}$ the airfoil has an angle of attack. No circulation nor pressures have built up. At $t=t_{1}+\Delta t$ aero pressures have fully developed. When this last situation occurs, the airfoil, at certain velocities, can pitch and plunge. Figure F. 3 and Figure F. 4 represent two modes, plunge and pitch.


Figure F.3: Mode 1 - Plunge
Figure F.4: Mode 2 - Pitch

As we said before, the flutter is caused by the coalescence of these structural modes - pitch and plunge. The frequencies of these modes can coalesce and generate flutter resonance.

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## VITA

Nicola Genco was born in Altamura (BA), Italy, on the 20th of February, 1984. He graduated from Galileo Galilei High School in chemistry, in Altamura (BA), Italy, in 2003. He went on to earn a Bachelors degree in Aerospace Engineering from the University of Rome "La Sapienza", Italy from 2003 to 2006, with a thesis in Propulsion on "The Hollow Cathode". In January 2007 he began attending the University of Dayton, Ohio for completion of his Masters Degree in Aerospace Engineering. During graduation, he presented his research project on biplane joined at the tips by endplates at DCASS (Dayton-Cincinnati Aerospace Sciences Symposium) and at two regional student conferences hosted by AIAA. He won the first place in the 2008 Regional Student Paper Competition for the Graduate division, and he will participate in the AIAA Foundation International Student Conference, part of the AIAA Aerospace Sciences Meeting, held in Orlando, Florida, 5 January 2009.


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