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May 2018

Modeling current flow in nanoparticle doped polymer film systems

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Seaman, David; Spradlin, Joshua; Leger, Janelle; and Rahmani, Armin, "Modeling current flow in nanoparticle doped polymer film systems" (2018). *Scholars Week*. 56.

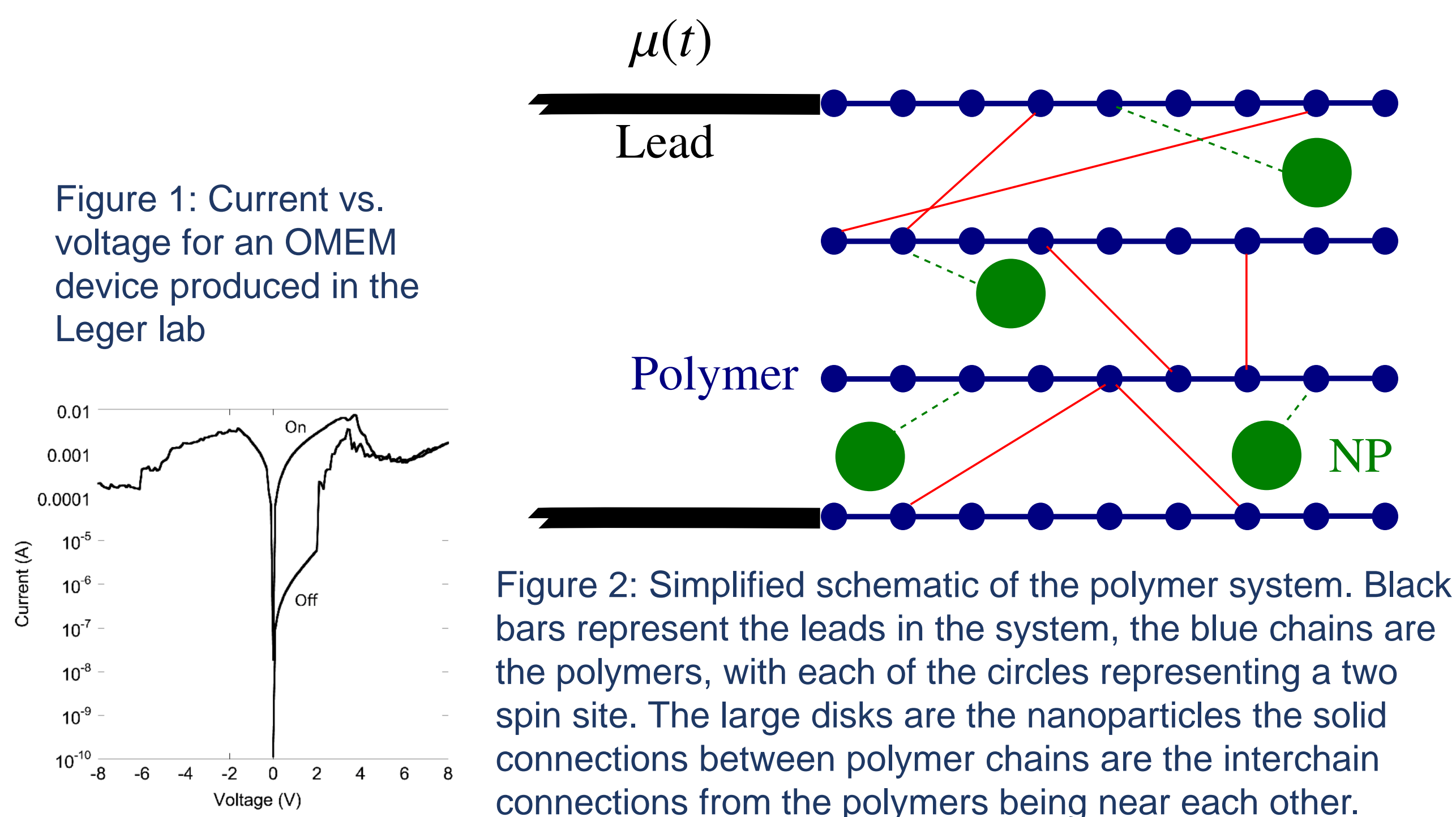
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Modeling current flow in nanoparticle doped polymer film systems

Introduction

- Nanoparticle doped polymer films are of interest for use in organic memory systems due to their ability to exhibit electrical bistabilities like that seen in figure 1 [1-3].
- Understanding of how current flows under these conditions is desired to better understand and predict how different films will behave.
- We attempt to find a computational model of the current through a simplified system with finite lead lengths as seen in figure 2



Calculating Current

- Working in the Heisenberg picture, we calculated current using the differential equation

$$\frac{d\hat{O}(T)}{dt} = i[H(T-t), \hat{O}(t)]$$

- Where $\hat{O}(T)$ is the Heisenberg operator at time T, which can be written as

$$\hat{O} = \Psi^\dagger O \Psi$$

- Which in the case we can write O as the following, where k is the hopping term and c the creation operators

$$\hat{O}(t=0) = ik(c_j^\dagger c_{j+1} - c_{j+1}^\dagger c_j)$$

- Where O is an LxL matrix representing the single particle Hamiltonian, and Ψ is the many particle wave function
- We can find for \hat{O} by solving for O(T) in:

$$\frac{dO(T)}{dt} = i[H(T-t), O(t)]$$

- To find the total current over time, we solve the differential equation for each T
- The expectation value of O(t) can be found with the initial state, |0>

$$\langle \hat{O}(T) \rangle = \langle 0 | \Psi^\dagger O(T) \Psi | 0 \rangle$$

- To find the above expectation value, we diagonalize H(0):

$$H(0) = V D V^\dagger$$

- V represents the eigenstates of the system, when we multiply both sides of D with

$$\Gamma = V^\dagger \Psi$$

- The expectation value for our operator at a given time is calculated using

$$\langle \hat{O}(T) \rangle = \langle 0 | \Gamma^\dagger V^\dagger O(T) V \Gamma | 0 \rangle$$

- When we substitute in $\tilde{O}(t) = V^\dagger O(t) V$ the calculation reduces to a summation along the diagonal of \tilde{O} or:

$$\langle \hat{O}(T) \rangle = \sum_{i \in \text{occupied}} \tilde{O}_{ii}(T)$$

Results: Non-Constant Potential

- When the duration of the non-constant potential phase lasted half time for $T_{total} = 1$, we saw that the current did not return to zero with the potential (figure 3).
- When the duration of non-constant potential was restricted to one tenth of the time $T_{total} = 10$, the current exhibits an oscillatory behavior without damping (figure 4).

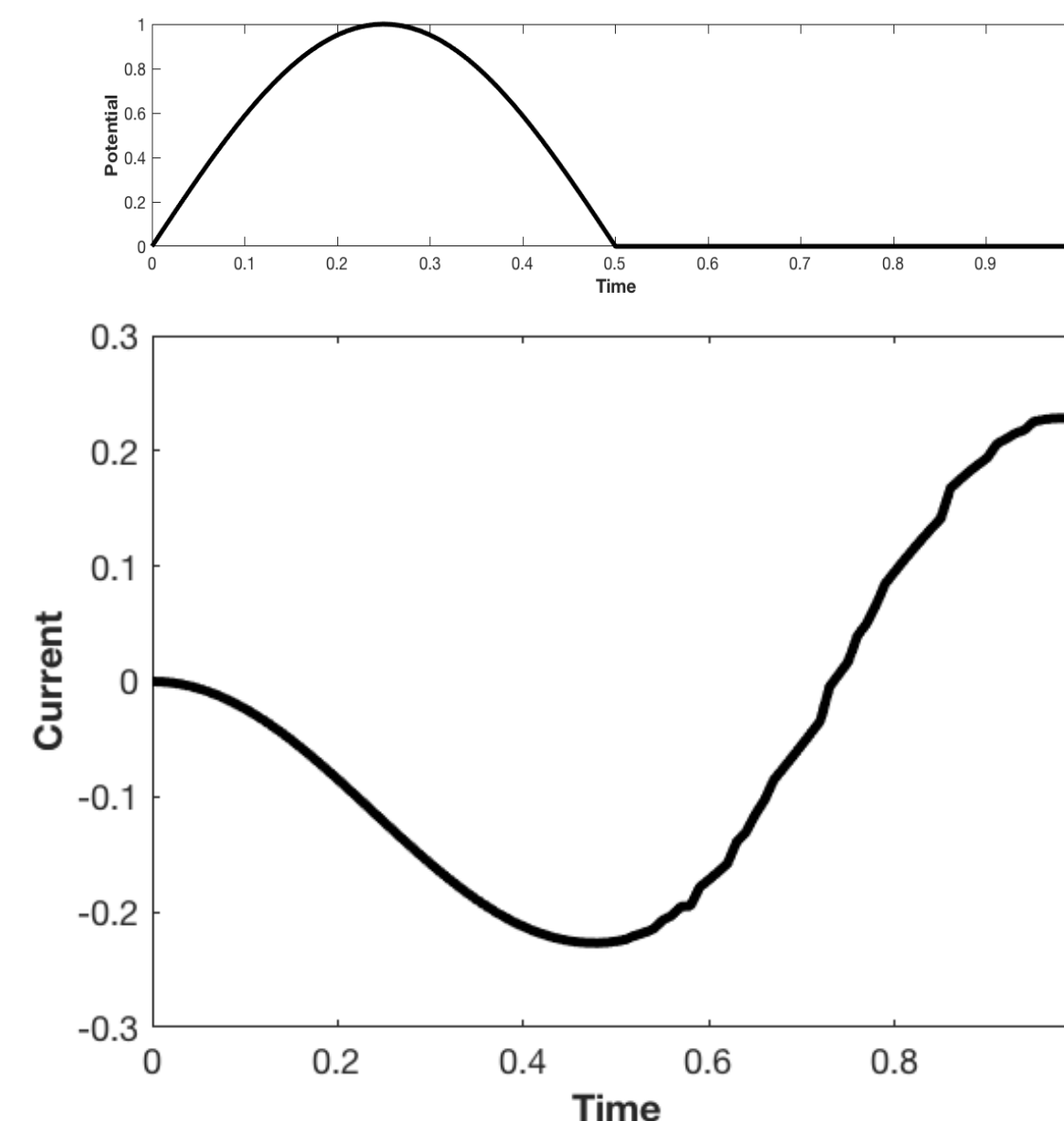


Figure 3: Corresponding current for the $T_{total} = 1$ case

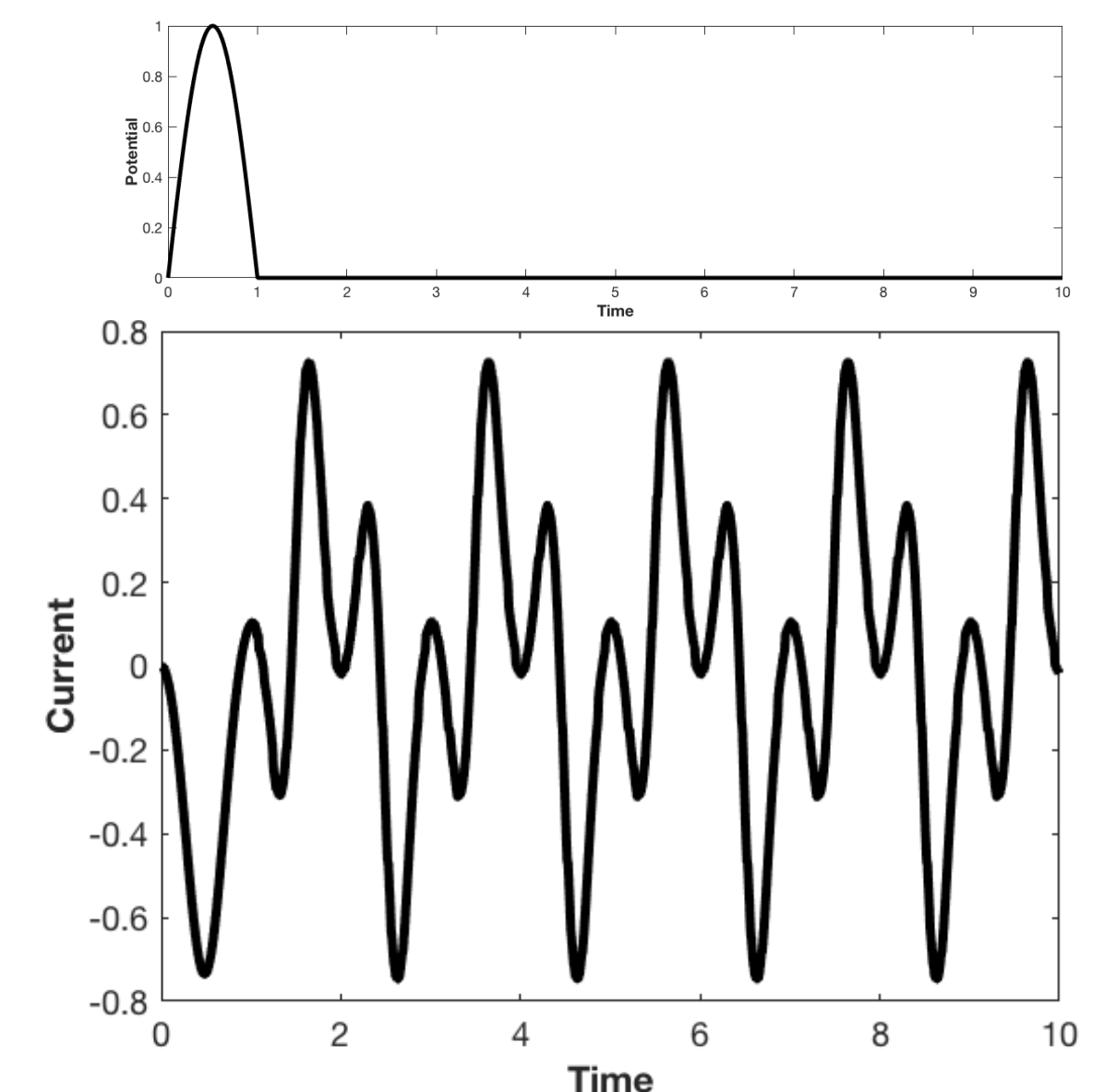


Figure 4: Corresponding current for the $T_{total} = 10$ case

Results: Constant Potential Varying Lead Lengths

- To simulate the infinite lead limit, we considered the impact of the length of the electrode leads
- For a constant potential, the current should also remain constant, following Ohm's law
- However, we see that our model predicts a non-constant current flow
- To isolate the cause, we examine the effect of the leads has on the oscillations
- The lead length has no noticeable impact on the current flow on the time scale we considered (figure 5)

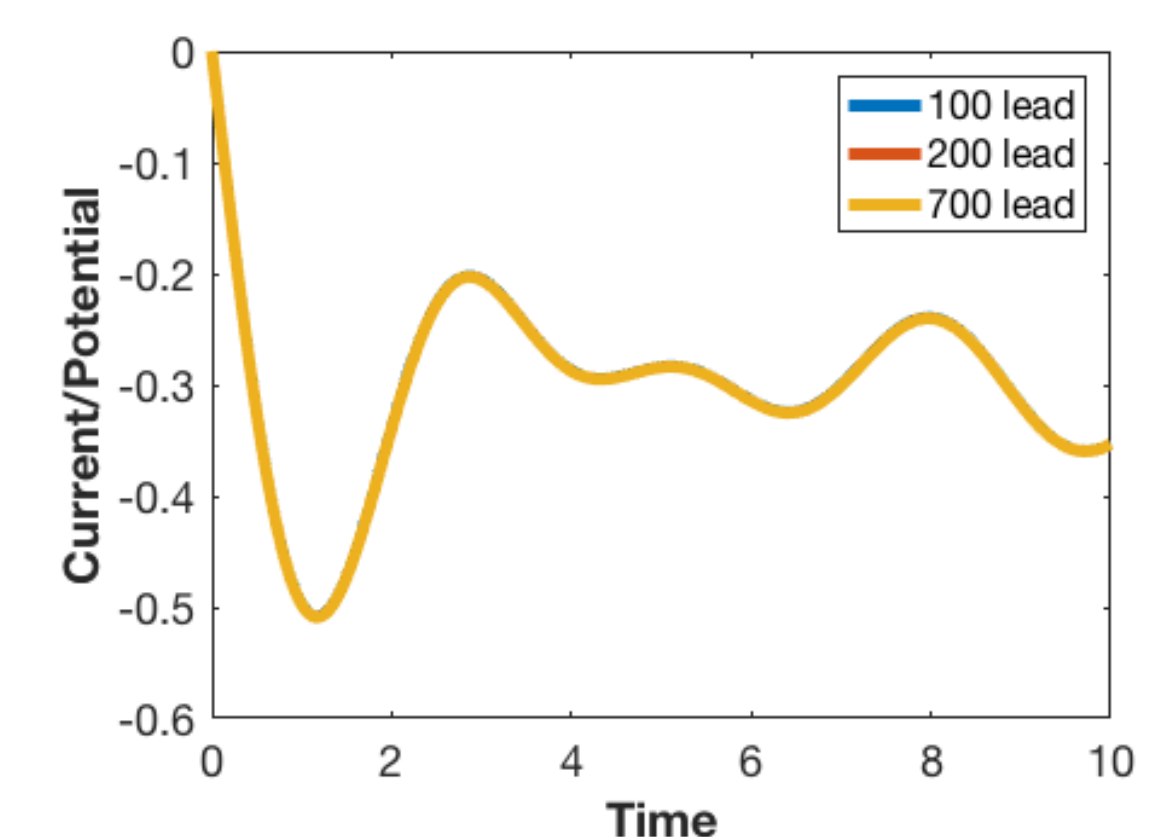


Figure 5: Current flow at constant potential ($V=0.01$).

Results: Flat Potentials

- Plotted current for different constant potentials to investigate the impact of the potential
- To determine if the oscillations are arising from setting the potential too high
- We consider a variety of constant potentials varying from $V = 0.00001$ to $V = 0.1$
- The potential has no effect on the time scale we examine in the system (figure 6)

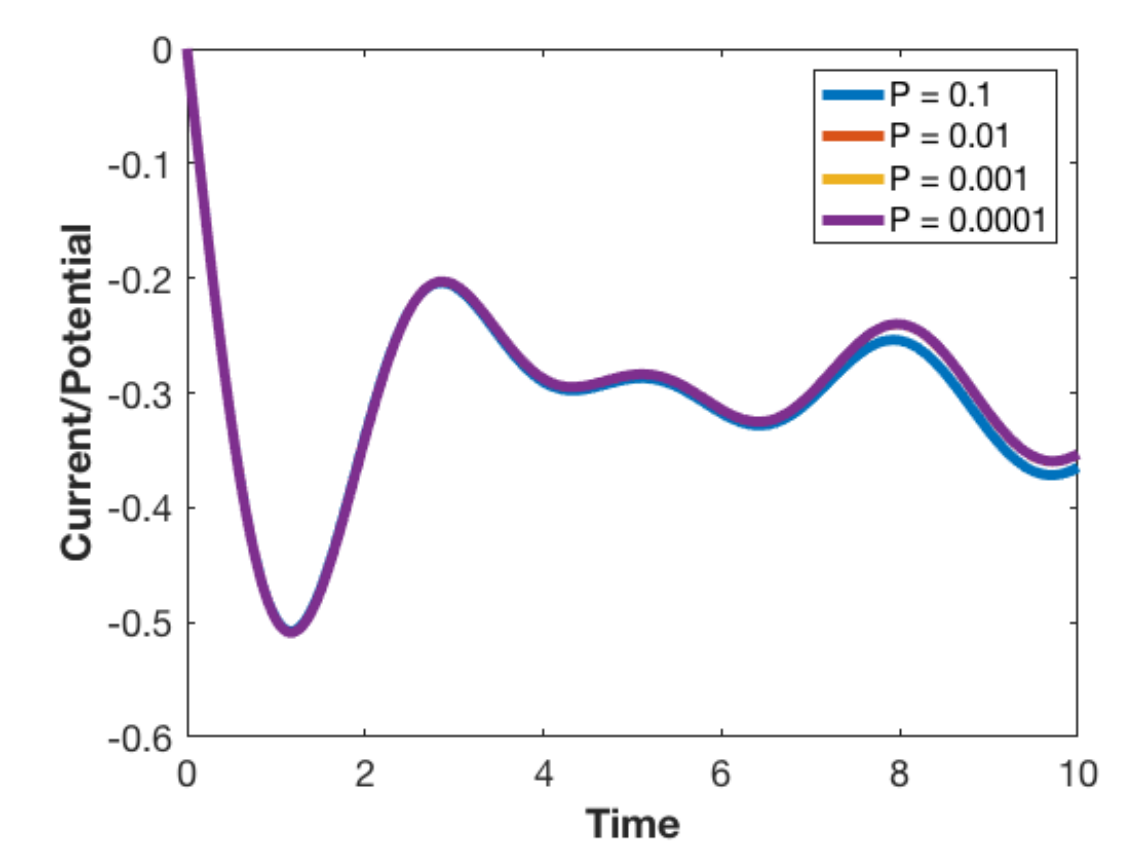


Figure 6: Current with constant leads of 600 sites

Conclusions

- Our results are preliminary without structure change.
- Size of the leads and current values below $V = 0.1$ has no effect on the current behavior in our model, with both exhibiting similar oscillatory behavior
- Next step is verifying our data with linear response by comparing to a model that treats complex current junctions as a scattering problem utilizing the Landauer method as seen in Wu et. al. [4]

Acknowledgments and Funding

- Leger Group
 - Joshua Spradlin
- WWU AMSEC SEED Grant funding

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